

Minimum “length scale”: gravity as a regulator redux

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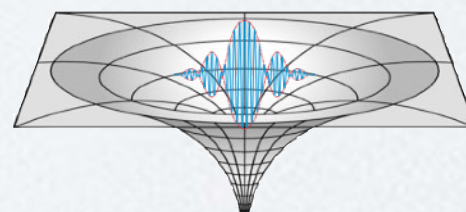
FLaG Meeting

Bologna

21/12/2022



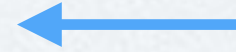
Theory and Phenomenology
of Fundamental Interactions
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Gravity as a regulator

Planck scale:

$$m = m_{\text{p}} \equiv \sqrt{\frac{c \hbar}{G_{\text{N}}}}$$
$$\lambda_{\text{C}} = \ell_{\text{p}} \equiv \sqrt{\frac{\hbar G_{\text{N}}}{c^3}}$$



$$\frac{\hbar}{m c} \equiv \lambda_{\text{C}} \sim R_{\text{H}} \equiv \frac{2 G_{\text{N}} m}{c^2}$$

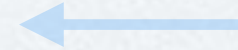


Gravity as a regulator of (mathematical) divergences [1]

Gravity as a regulator

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Gravity as a regulator of (mathematical) divergences [1]:

1. Planck **length** (mass) as a **geometric minimum length** (maximum energy) $\sim$ spacetime crystal / lattice
2. Planck **scale** as an emergent / dynamical regulator (**from non-linearity and quantum**) [2,3]

[2] R.C., I. Kuntz, *Revisiting the minimum length in the Schwinger-Keldysh formalism*, EPJ C 80 (2020) 958 [arXiv:2006.08450]

[3] R.C., W. Feng, I. Kuntz, F. Scardigli, *Minimum length (scale) in QFT, GUP and the non-renormalisability of gravity* [arXiv:2210.12801]

EFT metric corrections

EFT of gravity [4]:

Field redefinition $\mapsto$ effective massive(less) GW

$$S = \frac{1}{8\pi G_N} \int d^4x \sqrt{-g} \left(R + c_1 R^2 + c_2 R_{\mu\nu} R^{\mu\nu} + \dots \right)$$
$$\mapsto \int d^4x \sqrt{-g} \left[\frac{R}{8\pi G_N} - \frac{1}{2} (\nabla \chi_i)^2 - m_i^2 \chi_i^2 + \dots \right]$$

Background field approx:

$$g_{\mu\nu}(x) = \bar{g}_{\mu\rho}(x) \left(e^{\sqrt{\frac{32\pi\ell_p}{m_p}} h(x)} \right)_\nu^\rho \simeq \eta_{\mu\nu} + \sqrt{\frac{32\pi\ell_p}{m_p}} h_{\mu\nu}(x) + \frac{16\pi\ell_p}{m_p} \lim_{x \rightarrow y} h_{\mu\rho}(x) h^\rho_\nu(y) + \dots$$

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Geometric length [2,5]:

$$\ell_{\text{in-in}}^2 = \langle 0_{\text{in}} | \hat{g}_{\mu\nu} | 0_{\text{in}} \rangle dx^\mu dx^\nu = \ell^2 + \frac{16 \pi \ell_p}{m_p} \lim_{x \rightarrow y} \langle 0_{\text{in}} | \hat{h}_{\mu\rho}(x) \hat{h}^\rho_\nu(y) | 0_{\text{in}} \rangle dx^\mu dx^\nu$$

$$\equiv \ell^2 + \frac{16 \pi \ell_p}{m_p} \lim_{x \rightarrow y} G^{\text{R}}_{\mu\rho}{}^\rho{}_\nu(x, y) dx^\mu dx^\nu$$

Length scale [6]:

$$\ell_{\text{in-out}}^2 = \langle 0_{\text{out}} | \hat{g}_{\mu\nu} | 0_{\text{in}} \rangle dx^\mu dx^\nu = \ell^2 + \frac{16 \pi \ell_p}{m_p} \lim_{x \rightarrow y} \langle 0_{\text{out}} | \hat{h}_{\mu\rho}(x) \hat{h}^\rho_\nu(y) | 0_{\text{in}} \rangle dx^\mu dx^\nu$$

$$\equiv \ell^2 + \frac{16 \pi \ell_p}{m_p} \lim_{x \rightarrow y} G^{\text{F}}_{\mu\rho}{}^\rho{}_\nu(x, y) dx^\mu dx^\nu$$

[5] L.V. Keldysh, Zh. Eksp. Teor. Fiz. 47 (1964) 1515

[6] T. Padmanabhan, *Planck length as the lower bound to all physical length scales*, GRG 17 (1985) 215

EFT metric corrections

Geometric length ($\ell_{\text{in-in}} \in \mathbb{R}$): metric $\sim$ geodesics $\sim$ (free) retarded (GW) propagator

$$G_{\mu\nu\rho\sigma}^{\text{R}}(x, y) = \sum_i \left[-\frac{\theta(x^0 - y^0)}{2\pi} \delta(\ell^2) + \theta(x^0 - y^0) \theta(\ell^2) \frac{m_i J_1(m_i \ell)}{4\pi \ell} \right] \hbar P_{\mu\nu\rho\sigma}^i \sim \ell^0 + \ell^{-1}$$

Length scale: metric $\sim$ graviton scatterings $\sim$ Feynman propagator

$$G_{\mu\nu\rho\sigma}^{\text{F}}(x, y) = \sum_i \frac{\hbar P_{\mu\nu\rho\sigma}^i}{4\pi^2 (x - y)^2} + \mathcal{O}(|x - y|) \sim \ell^{-2}$$

$i = 1, \dots$ effective massive(less) GW: $P_{\mu\nu\rho\sigma}^i = \alpha_i \eta_{\mu\rho} \eta_{\nu\sigma} + \beta_i \eta_{\mu\sigma} \eta_{\nu\rho} + \gamma_i \eta_{\mu\nu} \eta_{\rho\sigma}$

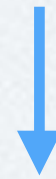
Background length: $\ell^2 = \eta_{\mu\nu} dx^\mu dx^\nu$

EFT minimum lengths

Coincidence limit $dx^\mu \sim \ell \rightarrow 0$

Minimum geometric length:

$$\ell_{\text{in-in}}^2 \sim (\ell^0 + \ell^{-1}) \ell^2 \rightarrow 0$$



Free propagation not affected by ℓ_p [2]

Minimum length scale:

$$\ell_{\text{in-out}}^2 \sim \begin{cases} e^{-\frac{\ell_p^2}{\ell^2} |\Sigma|} \ell^2 \rightarrow 0 & \Sigma \equiv \sum_i (\alpha_i + 4\beta_i + \gamma_i) \leq 0 \\ \ell_p^2 \frac{\Sigma}{\ell^2} \ell^2 \rightarrow \ell_p^2 |\Sigma| & \Sigma \equiv \sum_i (\alpha_i + 4\beta_i + \gamma_i) > 0 \end{cases}$$



Scattering processes (measurements) cannot probe below ℓ_p [2]

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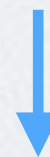
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Heisenberg microscope [7] / uncertainty relation modified to GUP [8]:

$$\Delta x \Delta p \geq \frac{\hbar}{2}$$

$$\Delta x \Delta p \geq \frac{\hbar}{2} \left(1 + \delta_0 \frac{\Delta p^2}{m_p^2} \right) \longrightarrow \ell = \ell_p \sqrt{\delta_0} = \ell_p \sqrt{\Sigma}$$



$$\delta_0 = \frac{2}{\pi} \sum_i (\alpha_i + 4\beta_i + \gamma_i)$$

N.B. (much) larger minimum length scales still possible in many-body systems [9]

[7] W. Heisenberg, *The physical principles of the quantum theory* (University of Chicago Press, 1930)

[8] M.P. Bronstein, Phys. Zeithchr. Der Sowjetunion 9 (1936) 140

[9] R.C. A quantum bound on the compactness, EPJ C 82 (2022) 10 [arXiv:21-3.14582] ~ talk at FLaG in TN 2022

Gravity as a regulator

Planck **scale** as an emergent / dynamical regulator in **non-renormalisable** theories [3]:

$$\delta_0^{\text{GR}} = \frac{8}{\pi}$$



Minimum scale makes theory finite (not just effective!)



$$\delta_0^{\text{Stelle}} = 0 \quad [10]$$

$$\delta_0^{\text{Spacetime crystal}} \leq 0 \quad [11]$$



Possibly fundamental (“complete”) theories with no minimum scale

[3] R.C., W. Feng, I. Kuntz, F. Scardigli, [arXiv:2210.12801]

[10] K.S. Stelle, PRD 16 (1977) 953

[11] P. Jizba, H. Kleinert, F. Scardigli, PRD 81 (2010) 084030

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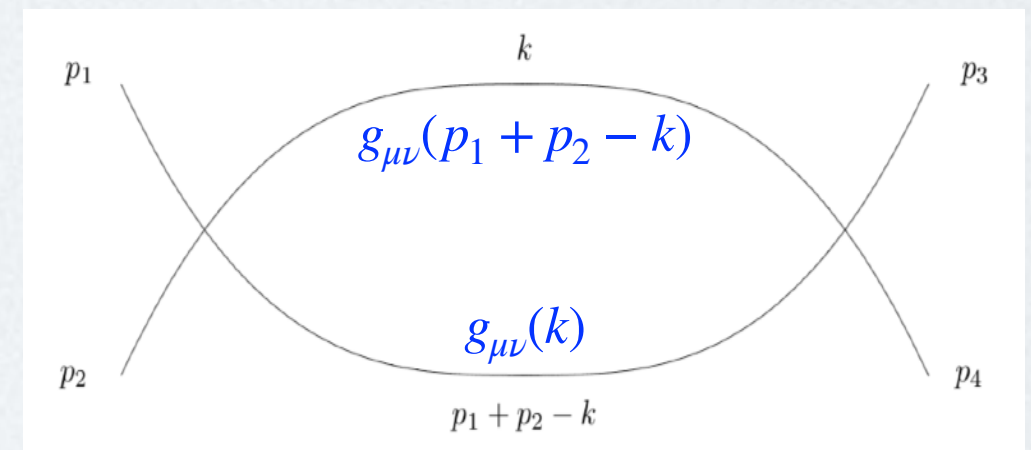
Non-linearity ~ beyond background field approx



Momentum-dependent propagators in loops [12]:

$$G^{\text{F}}(k) \implies G_{p_1+p_2-k}^{\text{F}}(k) \sim \Omega(p_1 + p_2 - k) G^{\text{F}}(k)$$

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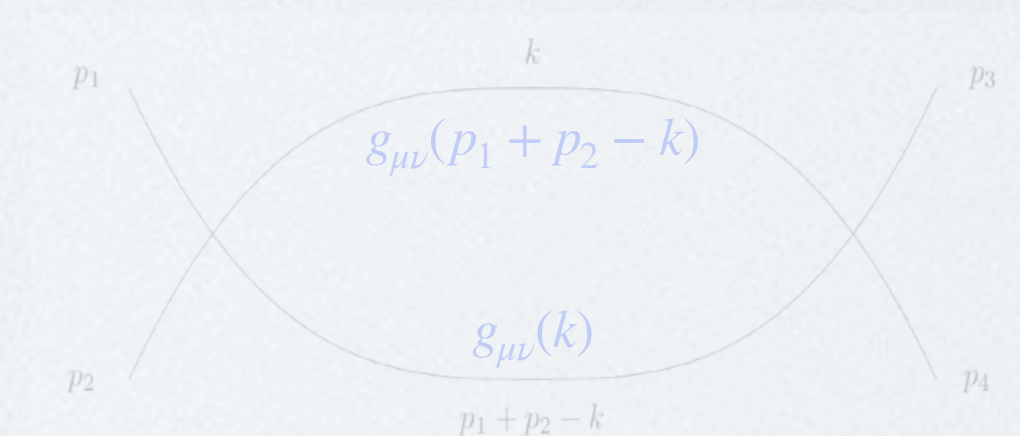
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