

Neutrino Theory



Adapted from Artwork by Sandbox Studio, Chicago with Ana Kova

Elvio Lisi
INFN, Bari

Neutrino Theory

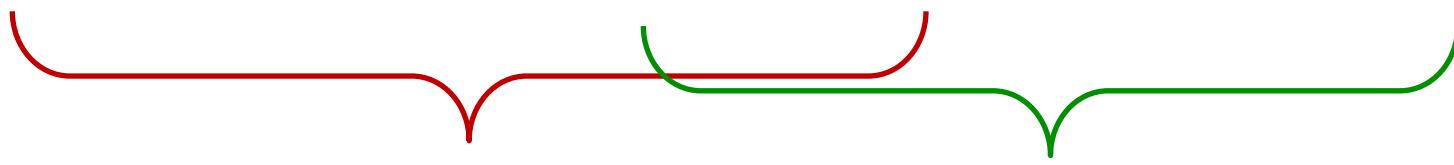
**Organizers' request: give a theory/pheno introduction to
 ν physics + ν -related experiments in this School, such as:**

LEGEND, CUORE/CUPID, PTOLEMY, RES-NOVA, AUGER ...

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Learning about neutrinos
(ν properties)

Learning from neutrinos
(ν sources)

For further exploration, browse the very complete
and updated neutrino website: www.nu.to.infn.it

Outline:

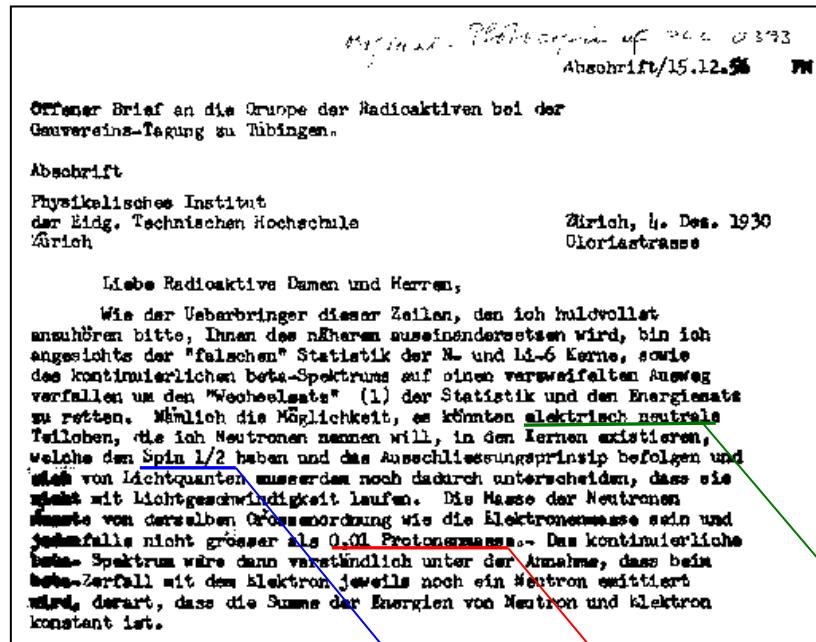
- Prologue: ν basic properties
- (Un)knowns in neutrino oscillations
- Non-oscillation ν observables
- Learning from neutrinos
- Epilogue: What's in a name?

Outline:

- **Prologue: ν basic properties**
- (Un)knowns in neutrino oscillations
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1930: ν hypothesis and first kinematical properties

Solving β -decay paradoxes: A famous letter by Wolfgang Pauli



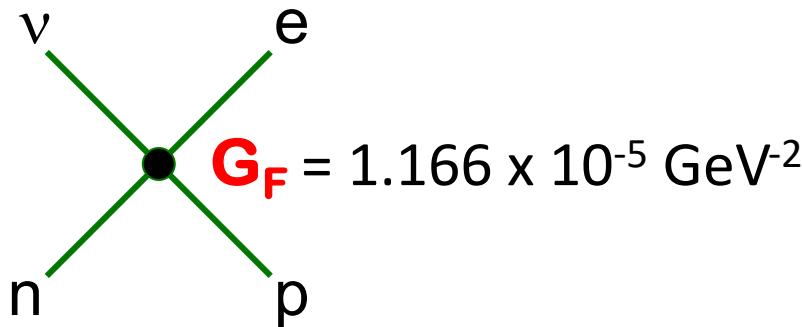
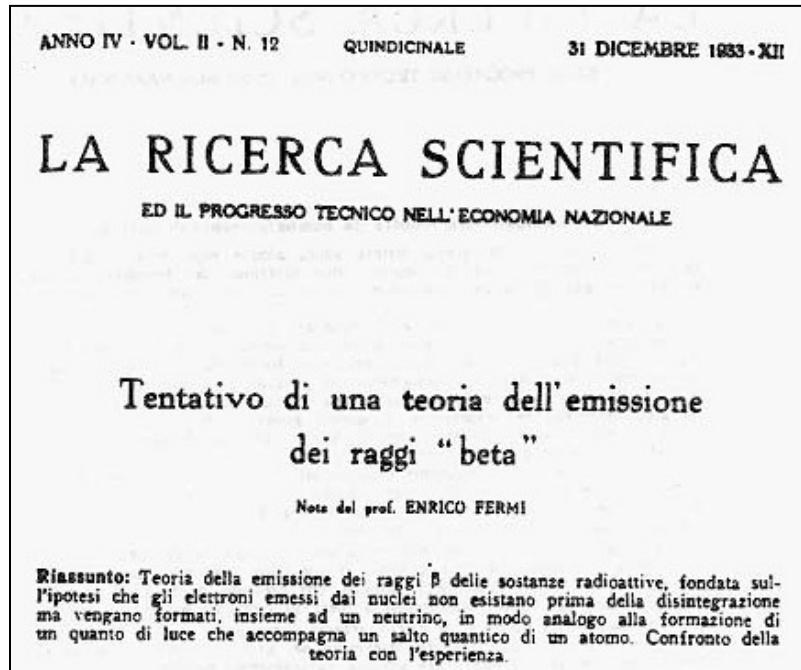
© CERN, Geneva

spin 1/2, tiny mass, zero electric charge

1930: $m_\nu < 0.01 \text{ GeV}$
Today: $m_\nu < 0.1 - 1 \text{ eV}$

Three years later: ν name and first dynamical properties

Neutrinos & QFT: A famous paper by Enrico Fermi

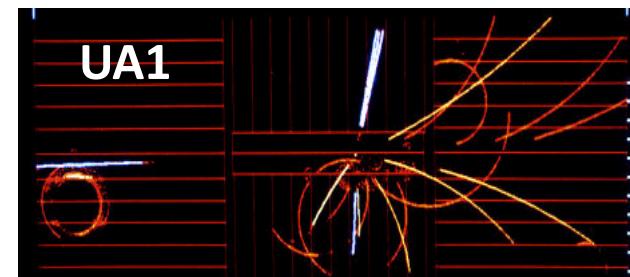
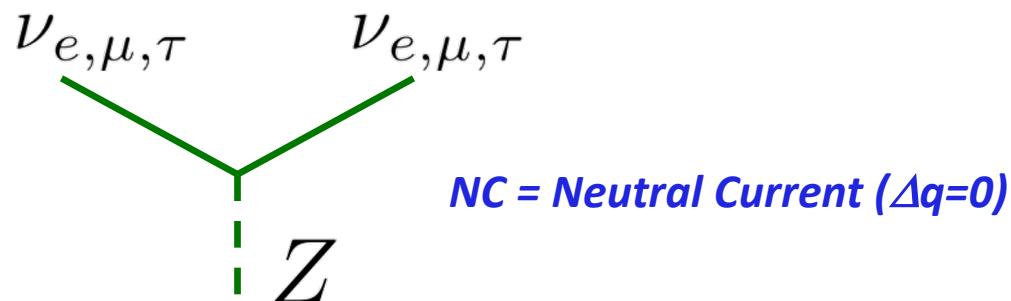
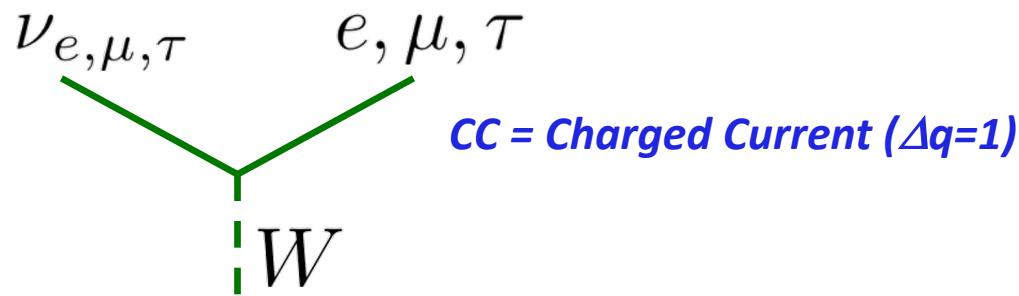


→ Sets the energy scale
 $\sqrt{(1/G_F)} \sim O(\text{few } 10^2) \text{ GeV}$
of weak interactions

Many decades of research have revealed other ν properties: There are
 3 different ν “flavors” $e \mu \tau$

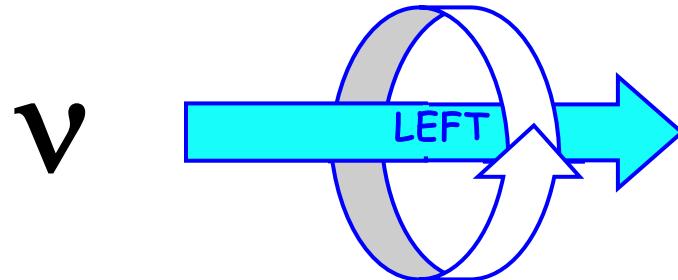
$$\begin{pmatrix} \nu_e \\ e \end{pmatrix} \begin{pmatrix} \nu_\mu \\ \mu \end{pmatrix} \begin{pmatrix} \nu_\tau \\ \tau \end{pmatrix} \quad \leftarrow \quad q = 0 \quad \leftarrow \quad q = -1 \quad (\Delta q = 1)$$

and their Fermi **interactions** are mediated by a charged **vector boson W**,
 with a neutral counterpart, the **Z boson**

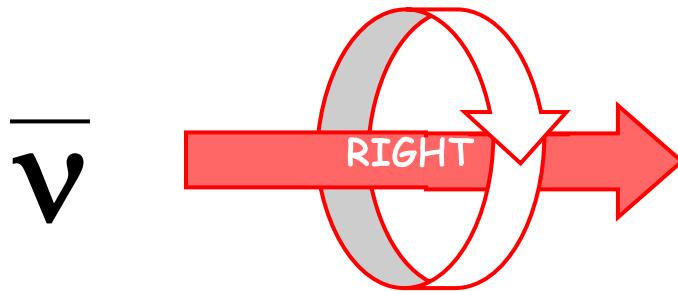


Such interactions are chiral (= not mirror-symmetric):

Neutrinos are created in
a left-handed (LH) state



Anti-nus are created in
a right-handed (RH) state

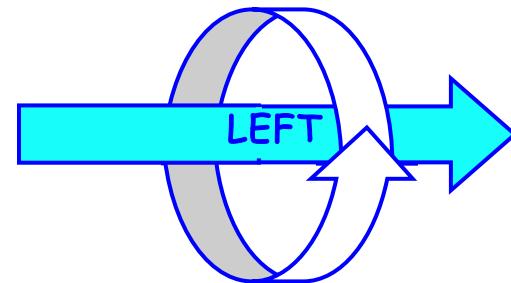


P parity symmetry is violated. *Nature is fundamentally chiral!*

Such interactions are chiral (= not mirror-symmetric):

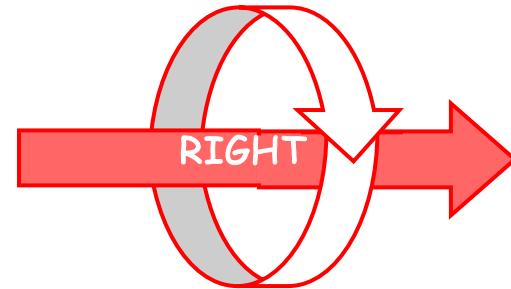
Neutrinos are created in a left-handed (LH) state

ν



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$\bar{\nu}$



Parity symmetry is violated. *Nature is fundamentally chiral!*

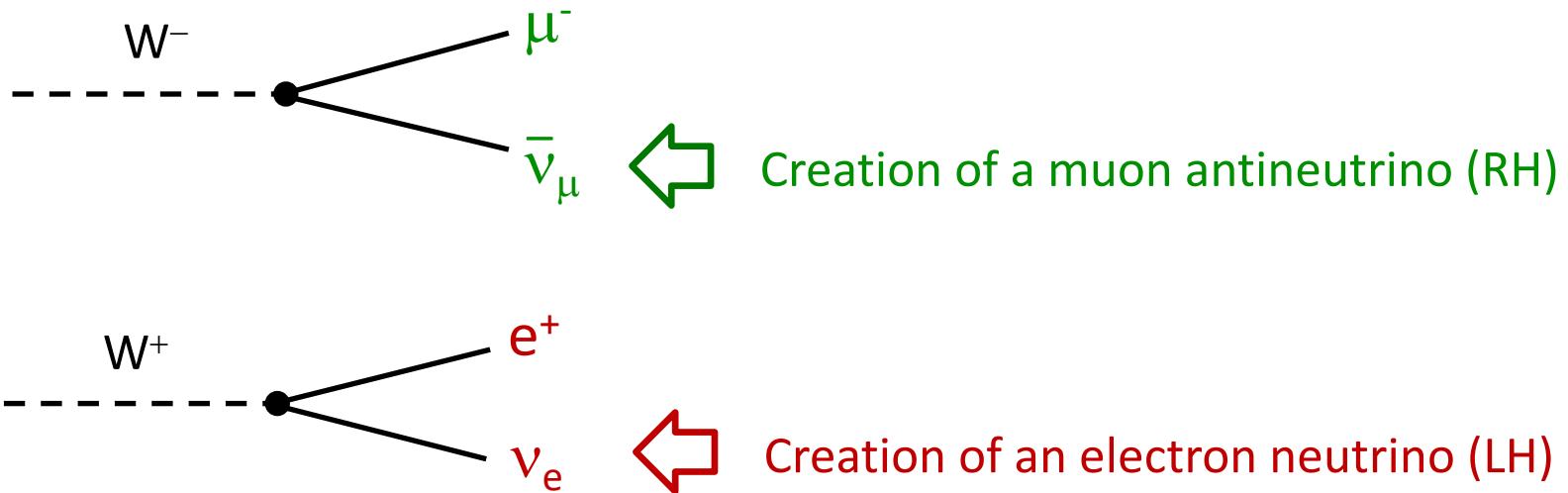
Other discrete symmetries in particle physics:

C charge conjugation (particle-antiparticle exchange)

T time reversal (change arrow of time)

Combined **CPT** conserved in QFT. **CP** violation required for baryogenesis.

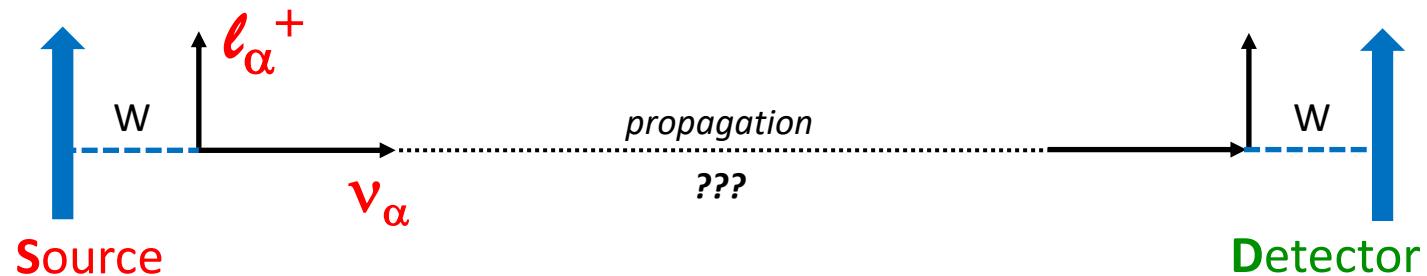
CC processes at **production** provide an operative definition of ν flavor, via the corresponding charged (anti)lepton. E.g., in leptonic decays:



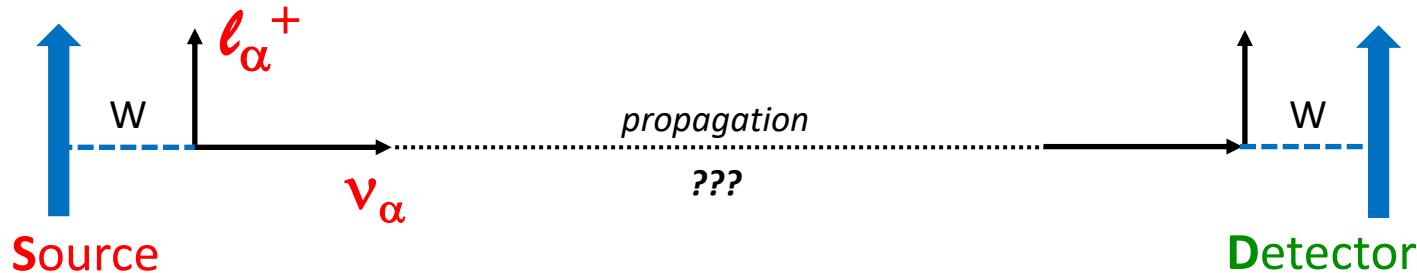
Similarly at **detection**, e.g.:



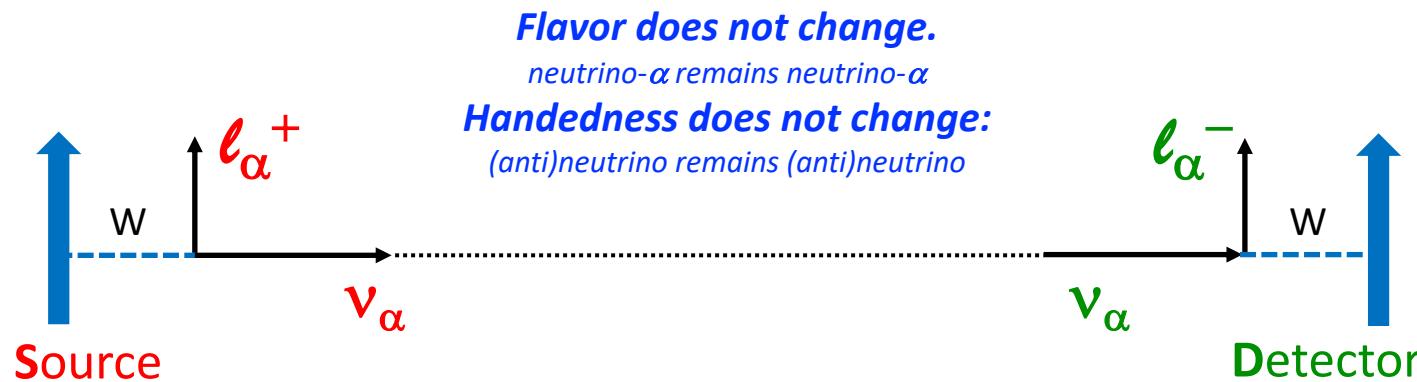
But... what happens in between?



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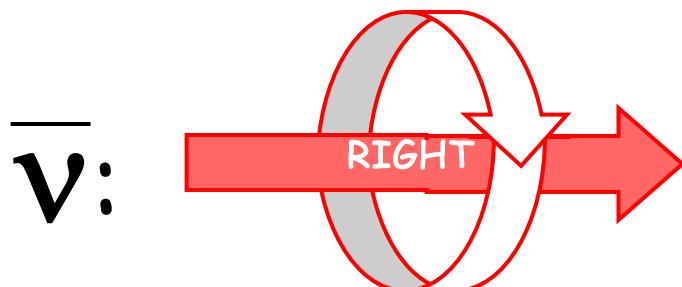
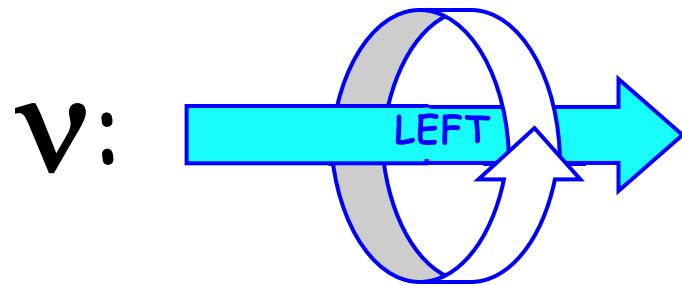
If neutrinos are massless: $v=c \rightarrow$ “clock” is frozen \rightarrow no change!



However, If v have mass, interesting things may happen to handedness and flavor...

Handedness: is a constant of motion for massless neutrinos

[You would see handedness reversal if you could travel faster... but you can't ($v=c$)!]

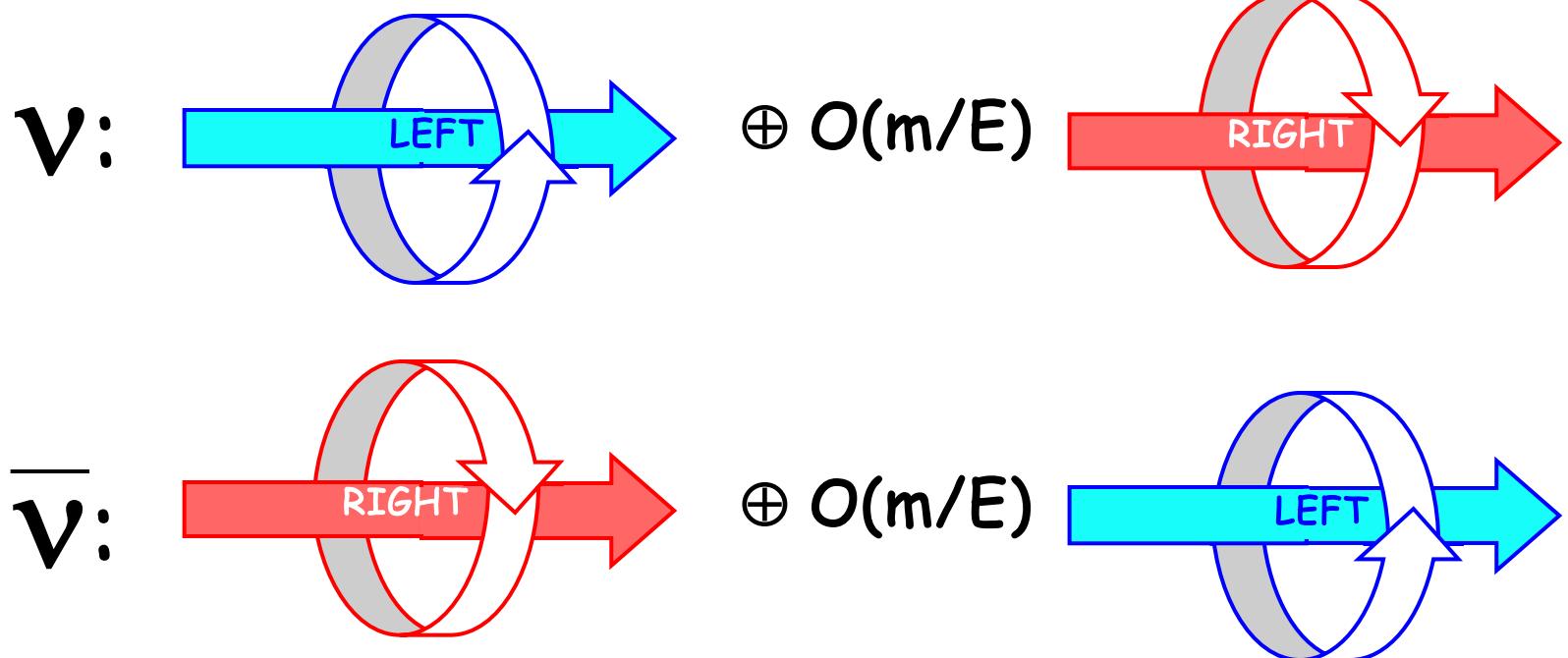


This is a massless “Weyl” two-spinor with 2 independent d.o.f

[And this was also the theoretical prejudice in the construction of the Standard Model]

Massive ν can develop the “other” handedness at $O(m/E)$

E = neutrino energy; the Dirac equation couples RH and LH states for $m \neq 0$



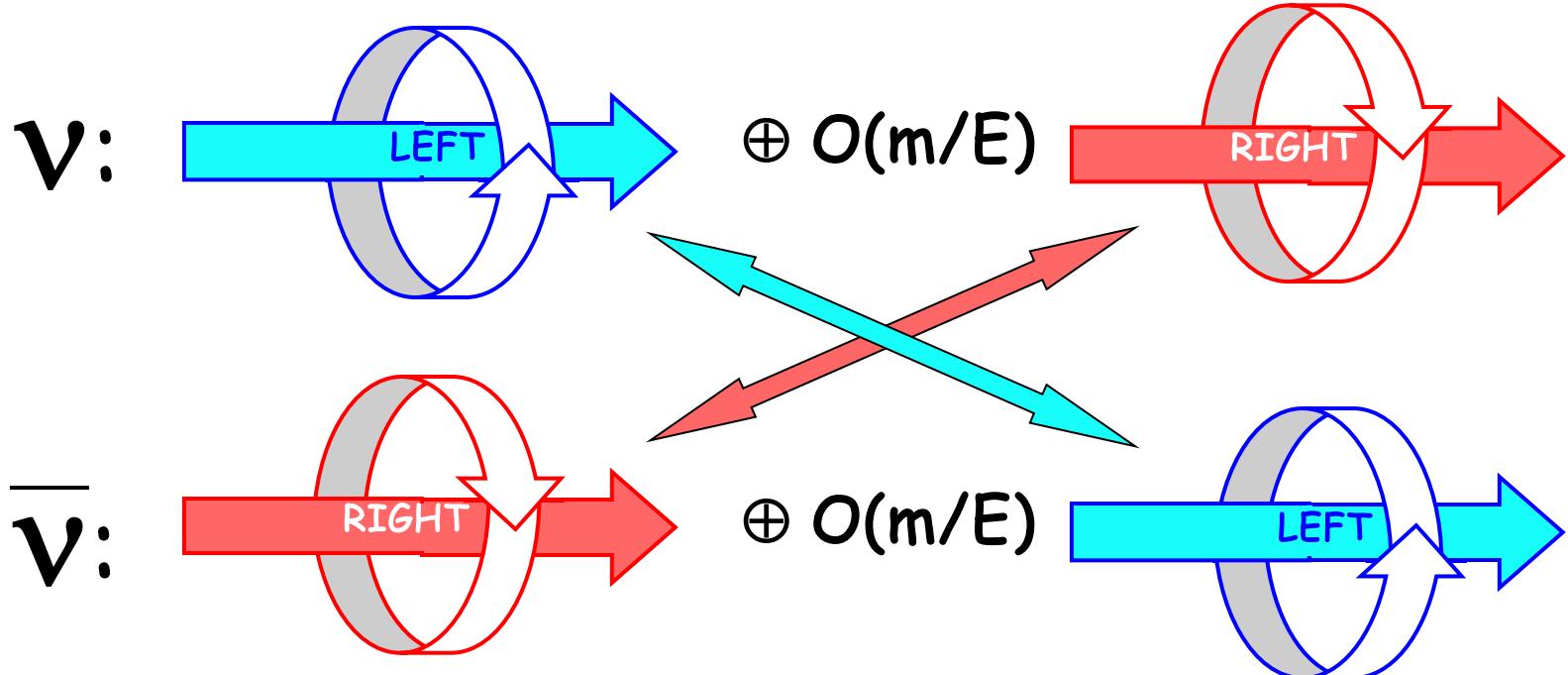
If these 4 d.o.f. are independent: massive “Dirac” four-spinor

Nu and anti-nu are different, just as the charged fermions are different.

Can define a conserved charge, the “lepton number” (+1 for leptons, -1 for antileptons)

But, for neutral fermions, two components might be identical !

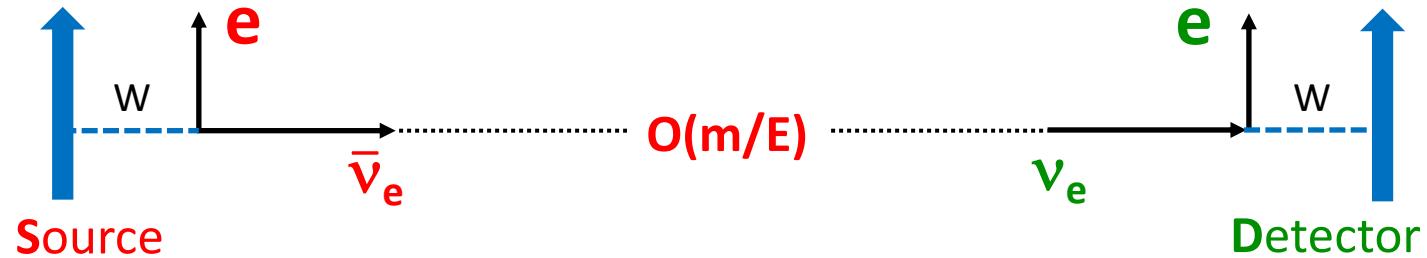
[Pairing components between electron and positron is forbidden: violates electric charge.]



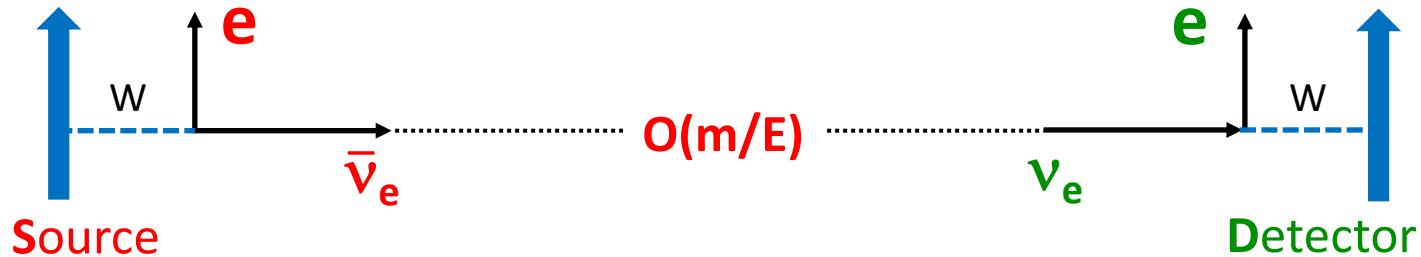
Only 2 independent d.o.f.: Massive “Majorana” four-spinor

No fundamental distinction between nu / antinu, up to a possible “Majorana phase”: A *very* neutral particle with no charge at all: **no electric charge, no lepton number ...**

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But, we have never seen anything like that so far!

E.g., the inverse beta-decay (IBD) reaction has been observed using reactor anti- ν_e ($E \sim \text{few MeV}$) with a statistics of $O(10^7)$ events:

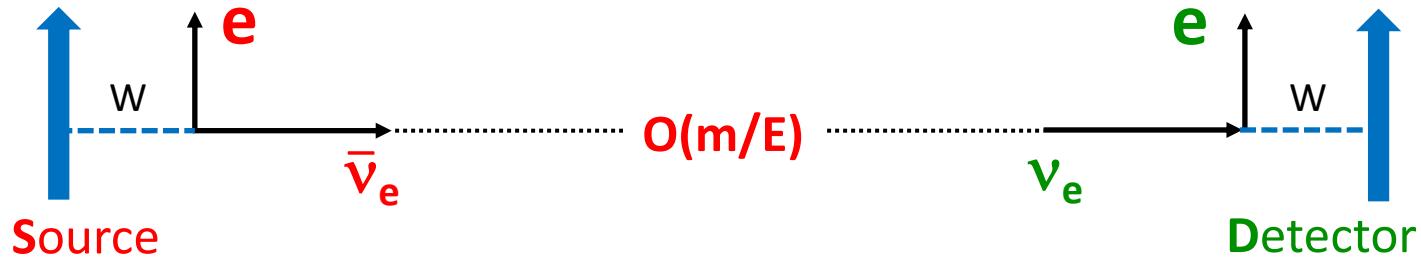


... but no event has ever been observed for the reaction:



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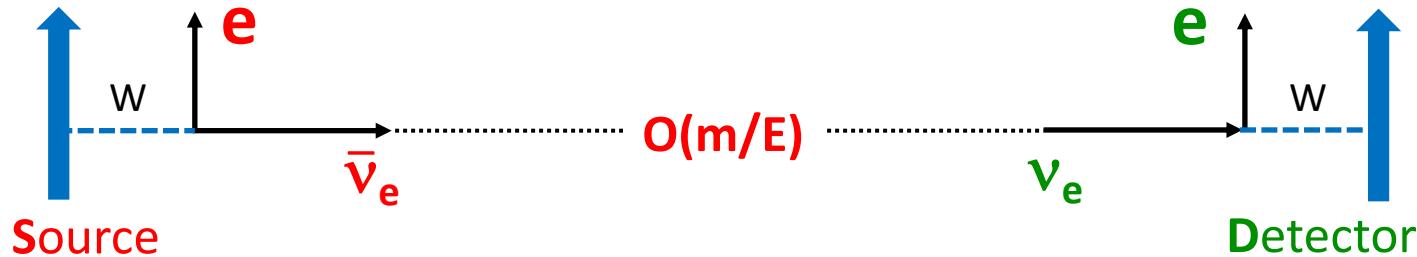
... but no event has ever been observed for the reaction:



Does this imply that neutrinos are not Majorana? NO! Still two valid options:

- 1) ν 's are Dirac: the 2nd reaction is strictly forbidden
- 2) ν 's are Majorana: the 2nd reaction is allowed... but suppressed by $m/E < 10^{-7}$!

If neutrinos are Majorana, expect $\bar{\nu} \rightarrow \nu$ transition:



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E.g., the inverse beta-decay (IBD) reaction has been observed using reactor anti- ν_e ($E \sim$ few MeV) with a statistics of $O(10^7)$ events:



... but no event has ever been observed for the reaction:



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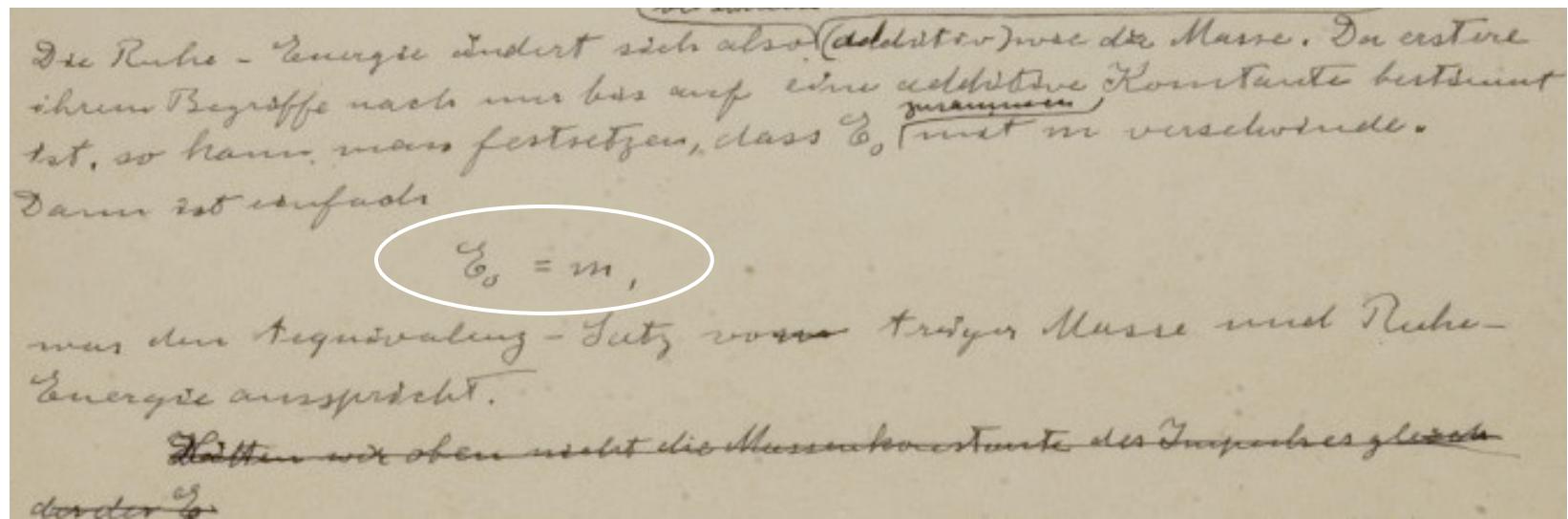
- 1) ν 's are Dirac: the 2nd reaction is strictly forbidden
- 2) ν 's are Majorana: the 2nd reaction is allowed... but suppressed by $m/E < 10^{-7}$!

Lesson: to search for Majorana ν , better compare $O(1)$ event to ~ 0 than to 10^N !

Neutrino flavor oscillations

Proof that ν are massive: not found (yet) via chirality flips but via **flavor change**, involving macroscopic distances $x=L$ and observable when $O(m^2 L/E) \sim O(1)$:

Let's start from a celebrated equation, already handwritten in natural units:



... namely, for $p \neq 0$:

$$E = \sqrt{m^2 + p^2}$$

Expand at small p/m or $m/p \rightarrow$

Our ordinary experience takes place in the limit: $p \ll m$

$$E \simeq m + \frac{p^2}{2m}$$

mass kinetic
energy

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... while ν often experience the opposite limit: $p \gg m$

$$E \simeq p + \frac{m^2}{2p}$$

Energy difference between two neutrinos ν_i e ν_j with mass m_i e m_j in the same beam ($p_i = p_j \simeq E$):

$$\Delta E \simeq \frac{\Delta m_{ij}^2}{2E}$$

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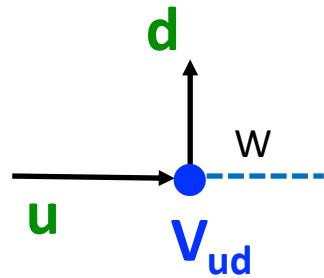
$$\Delta E \simeq \frac{\Delta m_{ij}^2}{2E}$$

Tiny $\Delta E \rightarrow$ Probed at large (macroscopic) $x=L \sim \Delta t$ from uncertainty relation:

$$1 \sim \Delta E \Delta t \simeq \frac{m_i^2 - m_j^2}{2E} L$$

Besides (different) neutrino masses, a second important ingredient of neutrino oscillations is **mixing**. In the Standard Model, mixing matrices arise in CC interaction vertices involving massive fermions:

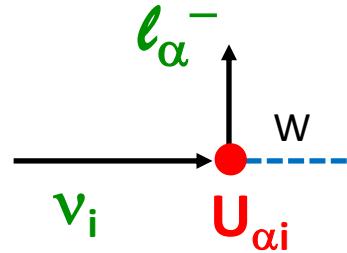
Quarks:



- = CC strength
 $\propto V$ element
 with $VV^\dagger=1$

CKM = Cabibbo-Kobayashi-Maskawa

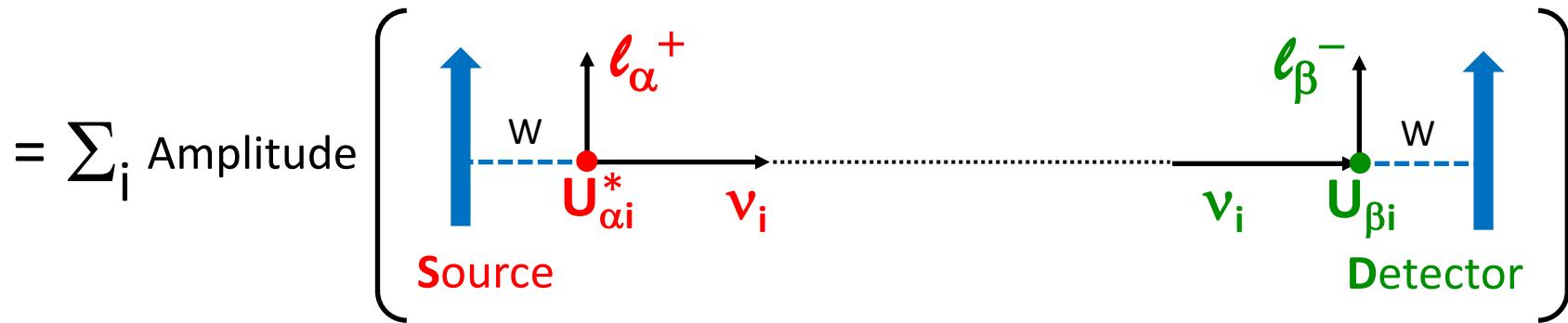
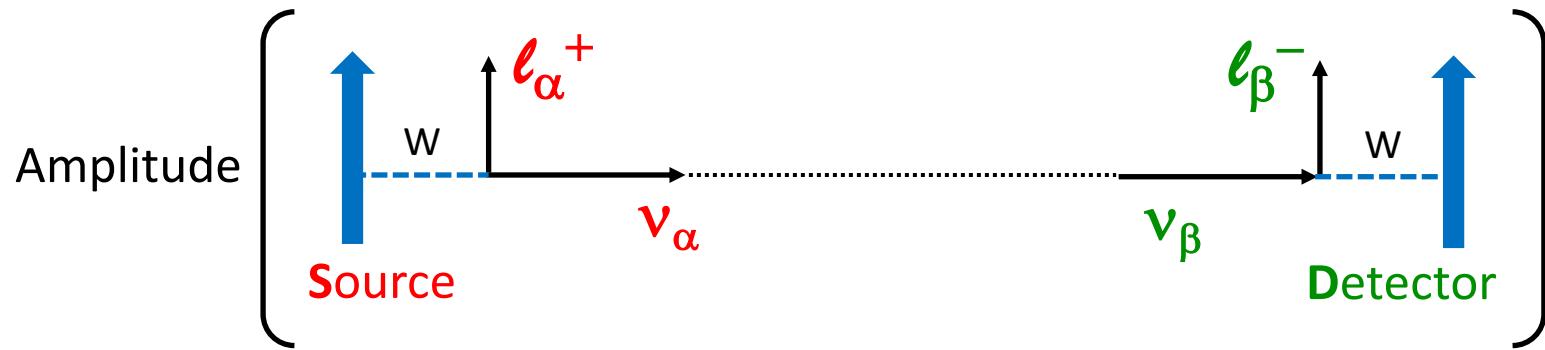
Leptons:



- = CC strength
 $\propto U$ element
 with $UU^\dagger=1$

PMNS = Pontecorvo-Maki-Nakagawa-Sakata

With both ingredients... flavor may change from α (production) to β (detection)!



$$v_\beta = U_{\beta i} v_i$$

$$v_i = U_{\alpha i}^* v_\alpha$$

Oscill. probability = $| \text{Amplitude} |^2$

Note: for \bar{v}

$$\ell_\alpha^\pm \rightarrow \ell_\alpha^\mp$$

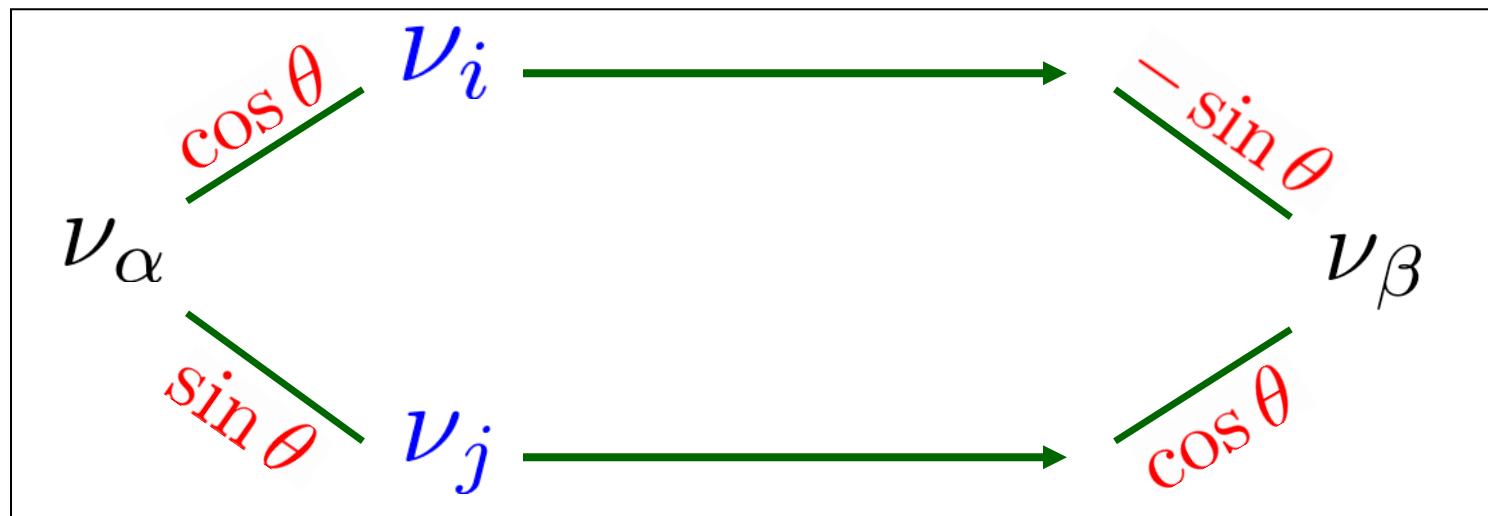
$$U \rightarrow U^*$$

The simplest oscillations: two neutrinos evolving in vacuum

(flavors α and β , masses m_i and m_j , $U = \theta$ -rotation)

$$\begin{bmatrix} \nu_\alpha \\ \nu_\beta \end{bmatrix} = \begin{bmatrix} \cos \theta & \sin \theta \\ -\sin \theta & \cos \theta \end{bmatrix} \begin{bmatrix} \nu_i \\ \nu_j \end{bmatrix}$$

Analogy with a double-slit interference experiment:

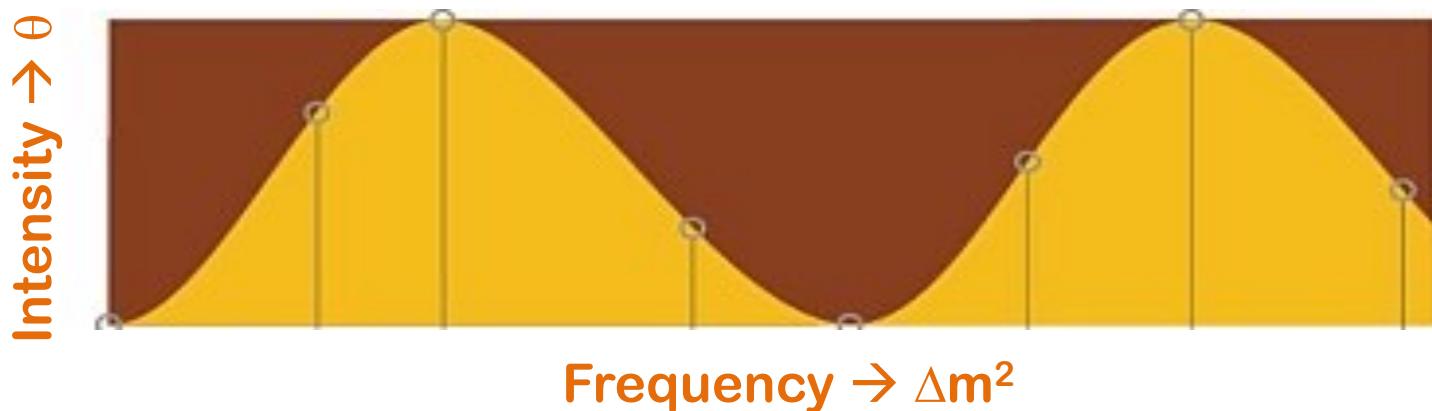


$$\Delta(\text{phase factor}) = e^{-i\Delta E \cdot x} = \exp\left(-i\frac{\Delta m_{ij}^2 x}{2E}\right)$$

Pontecorvo's 2ν oscillation probability:

$$P_{\alpha\beta} = \sin^2(2\theta) \sin^2 \left(\frac{\Delta m^2 x}{2E} \right)$$

Mixing angles govern osc. intensity, squared mass differences govern osc. frequency



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Mixing angles govern osc. intensity, squared mass differences govern osc. frequency

Change from natural units:

$$\frac{\Delta m^2 x}{4E} = 1.267 \left(\frac{\Delta m^2}{\text{eV}^2} \right) \left(\frac{x}{\text{m}} \right) \left(\frac{\text{MeV}}{E} \right)$$

In many textbooks: $1.267 \simeq 1.27$, no longer adequate in subpercent precision expts!

Brief recap

If neutrinos are massless: “clock” is frozen, no change in propagation

*Flavor does not change.
Handedness does not change.*



If neutrinos are massive, the other handedness develops at $O(m/E) \ll 1$

Iff neutrinos are Majorana:



If v are massive and mixed, other flavors develop at $O(\Delta m^2 L/E)$

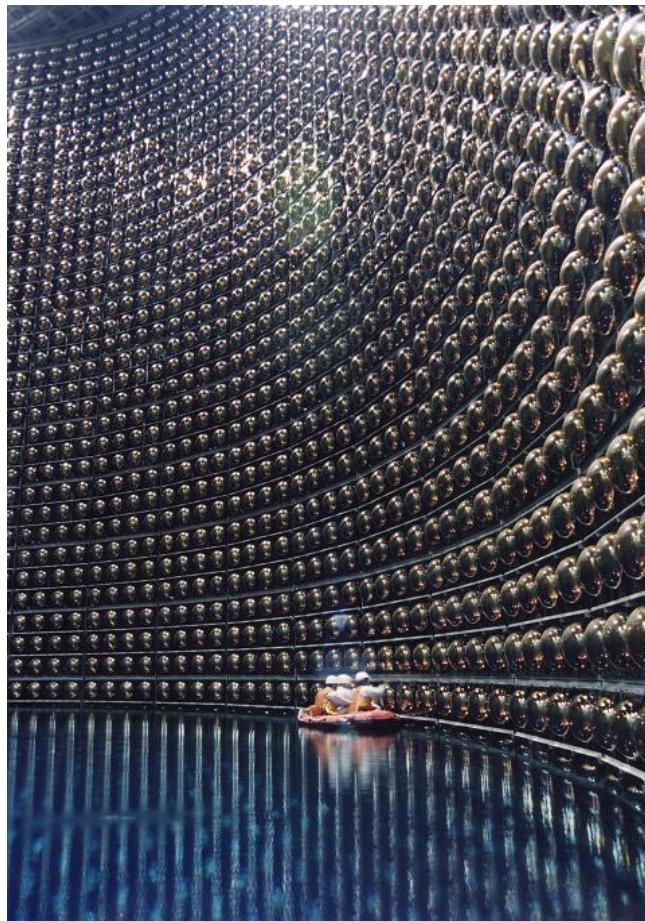
Flavor oscillations / transitions:



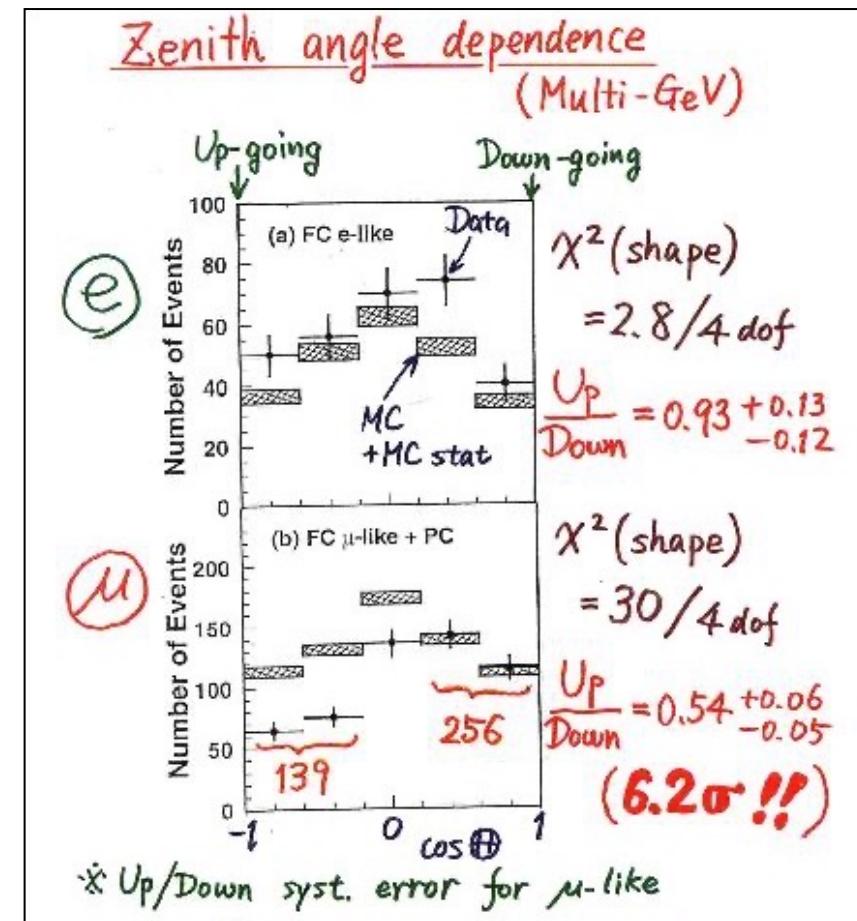
Outline:

- Prologue: ν basic properties
- **(Un)knowns in neutrino oscillations**
- Non-oscillation ν observables
- Learning from neutrinos
- Epilogue: What's in a name?

First evidence of oscillations in 1998 boosted neutrino physics...



Super-Kamiokande

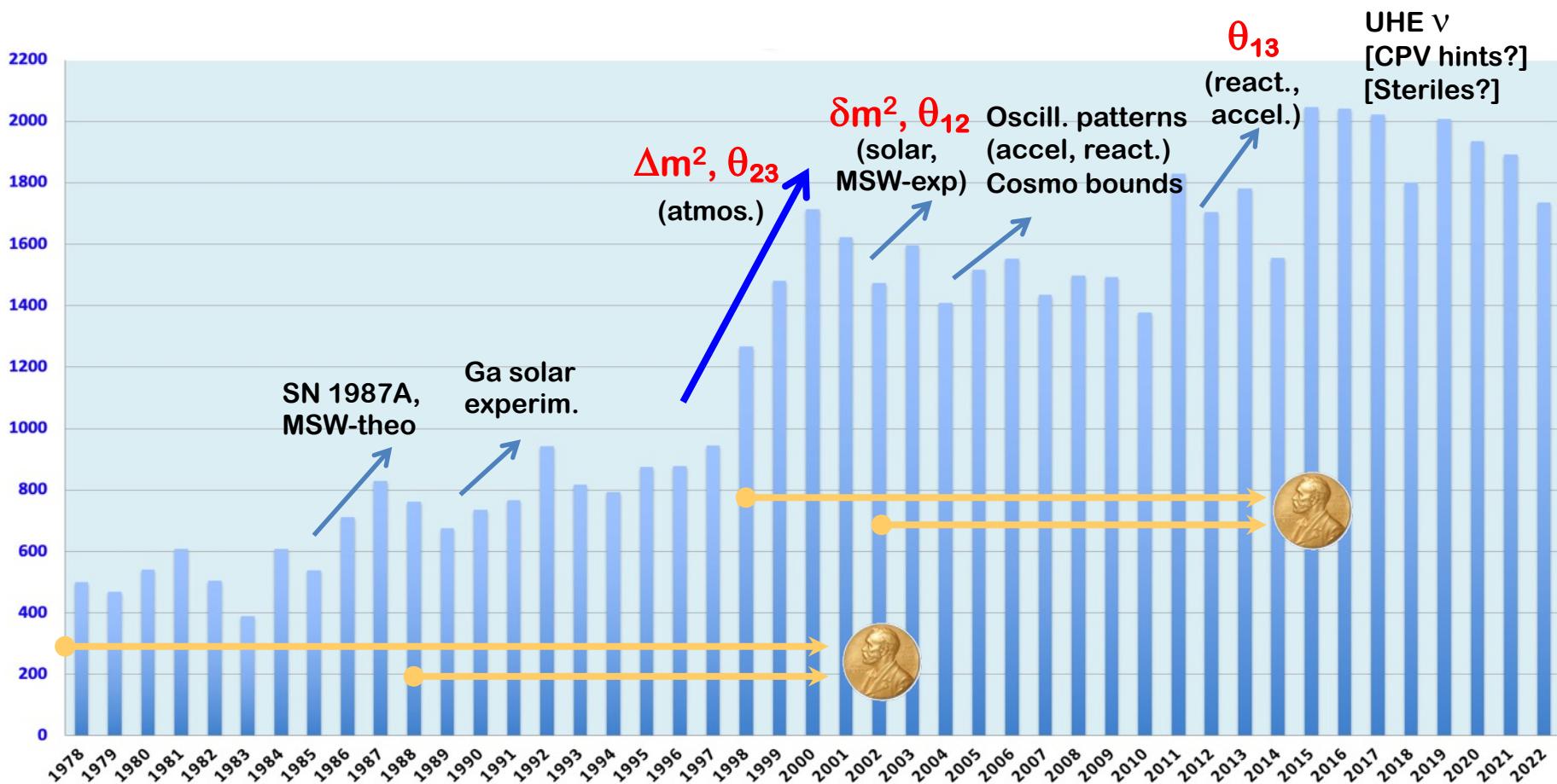


(T. Kajita at Neutrino' 98, Takayama)

Also confirmed by the MACRO atmospheric ν expt at LNGS

First evidence of oscillations in 1998 boosted neutrino physics...

Papers with *neutrino* in the title, yearly trend from INSPIRE



... and further results established a “standard 3ν framework”

The standard 3ν framework & mass-mixing parameters

Mixing matrix: CKM → PMNS [Pontecorvo-Maki-Nakagawa-Sakata]

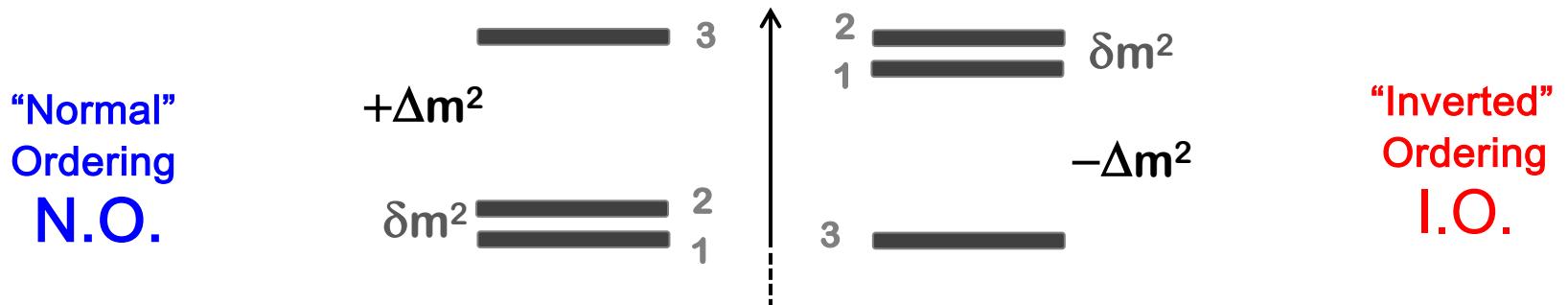
$$U_{\alpha i} = \begin{bmatrix} 1 & 0 & 0 \\ 0 & c_{23} & s_{23} \\ 0 & -s_{23} & c_{23} \end{bmatrix} \begin{bmatrix} c_{13} & 0 & s_{13}e^{-i\delta} \\ 0 & 1 & 0 \\ -s_{13}e^{i\delta} & 0 & c_{13} \end{bmatrix} \begin{bmatrix} c_{12} & s_{12} & 0 \\ -s_{12} & c_{12} & 0 \\ 0 & 0 & 1 \end{bmatrix} \begin{bmatrix} 1 & 0 & 0 \\ 0 & e^{i\alpha/2} & 0 \\ 0 & 0 & e^{i\beta/2} \end{bmatrix}$$

2-3 rotation 1-3 rotation
+ CPV “Dirac” phase
 $U(v) \rightarrow U^*(\bar{v})$

$c=\cos, s=\sin$

1-2 rotation Extra CPV phases
[if Majorana]
not tested in oscillat.

Mass [squared] spectrum)

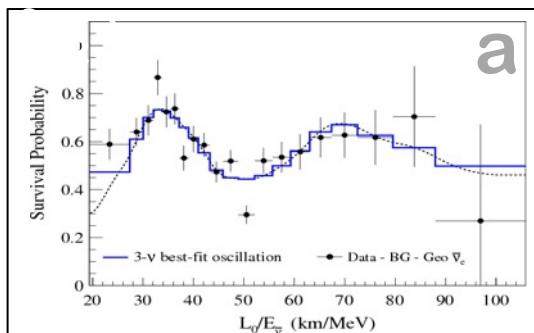


$$\delta m^2 = \Delta m_{21}^2, \quad \Delta m^2 = (\Delta m_{32}^2 + \Delta m_{31}^2)/2$$

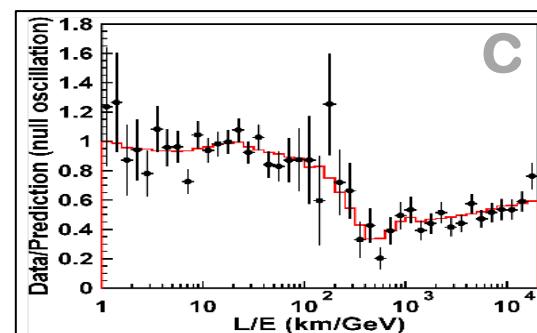
[Absolute ν mass scale not tested in oscillations]

3ν oscillations probed by many experiments in different flavor channels...

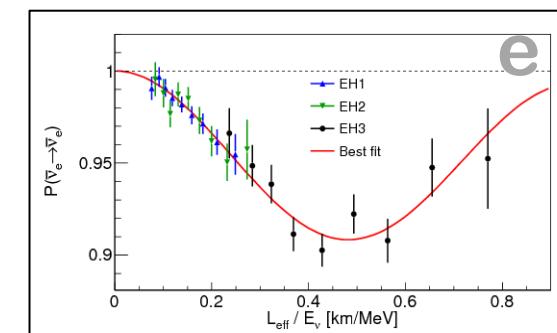
$e \rightarrow e$ (KamLAND, KL)



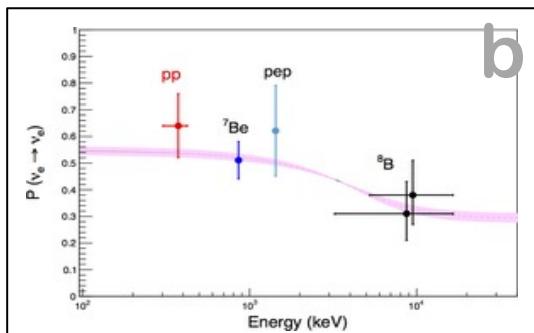
$\mu \rightarrow \mu$ (Atmospheric)



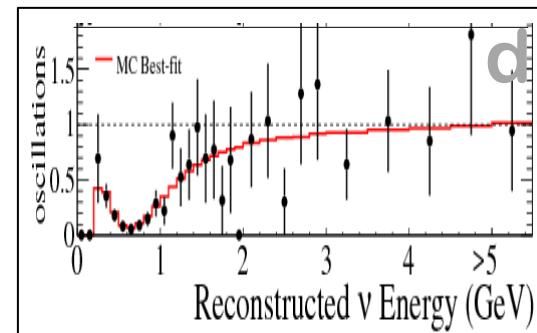
$e \rightarrow e$ (SBL Reac.)



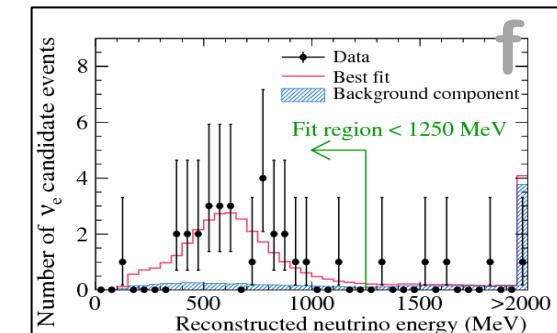
$e \rightarrow e$ (Solar)



$\mu \rightarrow \mu$ (LBL Accel.)



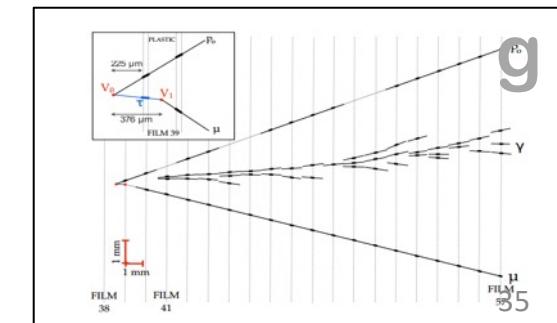
$\mu \rightarrow e$ (LBL Accel.)



LBL = Long baseline (few x 100 km); SBL = short baseline (~1 km)

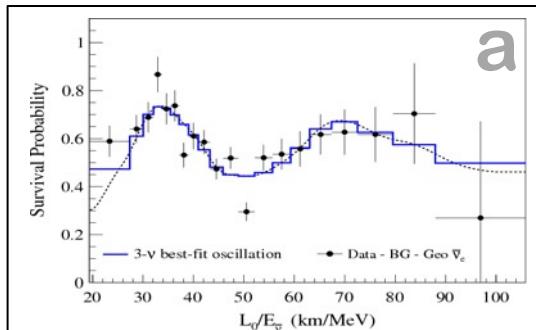
(a) KamLAND reactor [plot]; (b) Borexino [plot], Homestake, Super-K, SAGE, GALLEX/GNO, SNO; (c) Super-K atmosph. [plot], DeepCore, MACRO, MINOS etc.; (d) T2K (plot), NOvA, MINOS, K2K LBL accel.; (e) Daya Bay [plot], RENO, Double Chooz SBL reactor; (f) T2K [plot], MINOS, NOvA LBL accel.; (g) OPERA [plot] LBL accel., Super-K and IC-CD atmospheric.

$\mu \rightarrow \tau$ (OPERA, SK, DC)

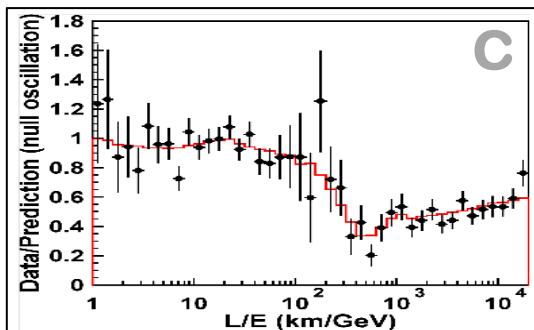


... with amplitude and frequency governed by 2 (or 3) leading parameters

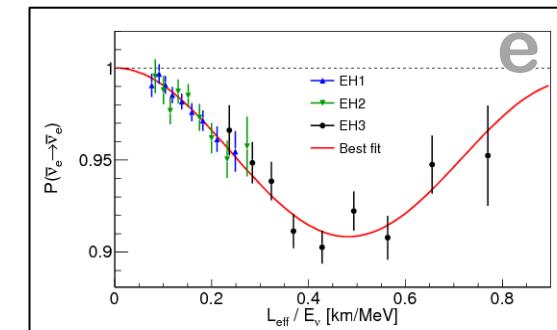
$e \rightarrow e$ ($\delta m^2, \theta_{12}$)



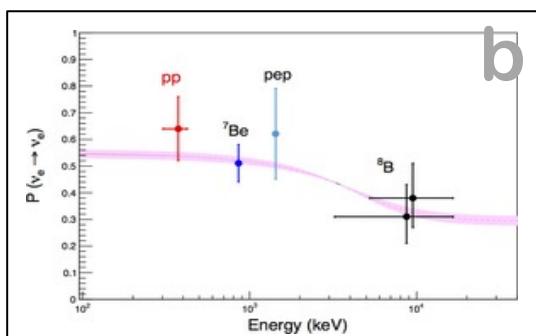
$\mu \rightarrow \mu$ ($\Delta m^2, \theta_{23}$)



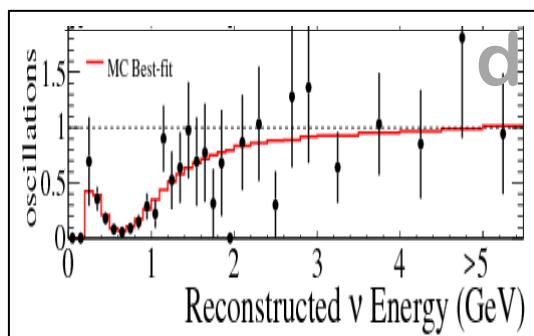
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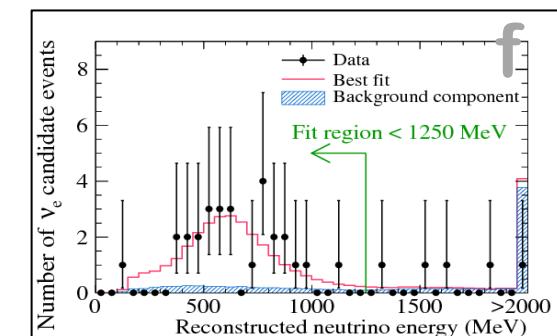
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$\mu \rightarrow \mu$ ($\Delta m^2, \theta_{23}$)



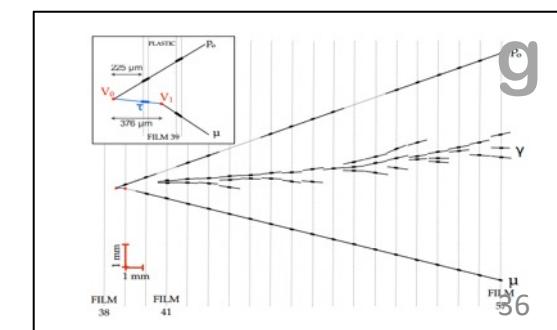
$\mu \rightarrow e$ ($\Delta m^2, \theta_{13}, \theta_{23}$)



5 param.'s known & (over)constrained \rightarrow consistency

Currently: focus on unknown par. & subleading effects,
especially CPV via $\nu_\mu \rightarrow \nu_e$ in LBL accel. and atmos. expts
and NO/IO mass spectrum via reactor + accel + atmos.

$\mu \rightarrow \tau$ ($\Delta m^2, \theta_{23}$)



Sketchy 3ν overview (see pdg.lbl.gov review for details)

5 knowns:

$$\begin{aligned}\delta m^2 &\sim 8 \times 10^{-5} \text{ eV}^2 \\ |\Delta m^2| &\sim 2 \times 10^{-3} \text{ eV}^2 \\ \sin^2 \theta_{12} &\sim 0.3 \\ \sin^2 \theta_{23} &\sim 0.5 \\ \sin^2 \theta_{13} &\sim 0.02\end{aligned}$$

5 unknowns:

- Oscillations*
- δ CPV Dirac phase
 - $\text{sign}(\Delta m^2) \rightarrow \text{NO/IO}$
 - θ_{23} octant ($>$ or $<$ $\pi/4$?)
- Non-oscillat.*
- absolute mass scale
 - Dirac/Majorana nature



How do $\nu_\mu \rightarrow \nu_e$ oscillation searches probe CP (T) violation?



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TIME REVERSAL VIOLATION IN NEUTRINO OSCILLATION

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*Laboratoire de Physique Théorique et Hautes Energies, Paris, France***

Received 11 October 1977

We discuss the possibility of CP or T violation in neutrino oscillation. CP requires $\nu_\mu \leftrightarrow \nu_e$ and $\bar{\nu}_\mu \leftrightarrow \bar{\nu}_e$ oscillations to be equal. Time reversal invariance requires the oscillation probability to be an even function of time. Both conditions can be violated, even drastically, if more than two neutrinos exist.

For two neutrinos, no CPV:

$$\stackrel{(-)}{\nu}_e = \cos\theta_{12} \nu_1 + \sin\theta_{12} \nu_2$$

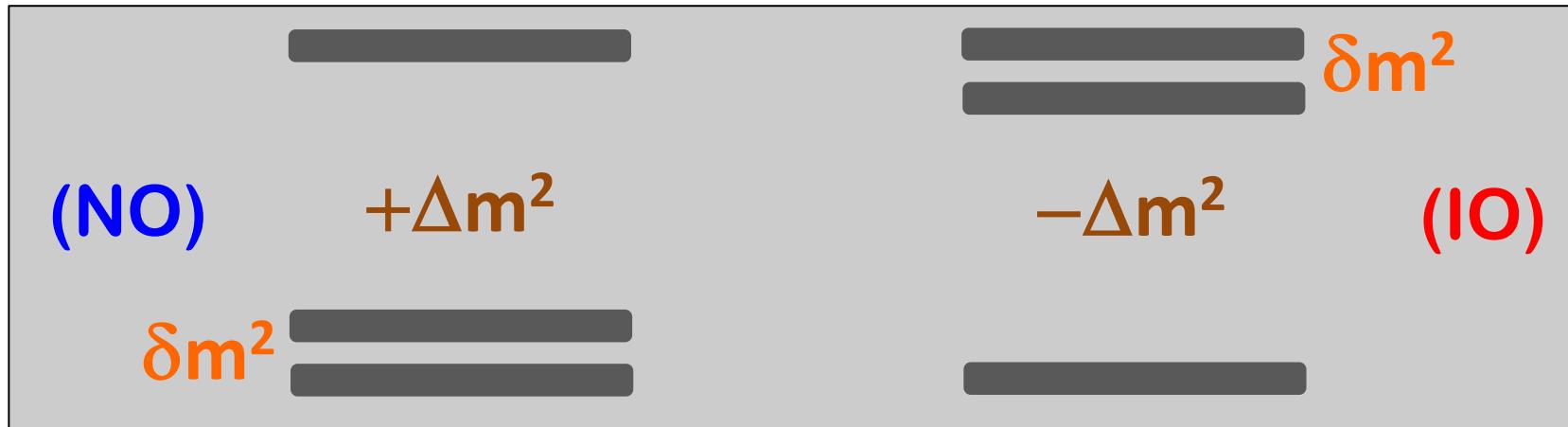
For three neutrinos: possible **CPV phase δ** , tested via ν versus $\bar{\nu}$

$$\stackrel{(-)}{\nu}_e = \cos\theta_{13} (\cos\theta_{12} \nu_1 + \sin\theta_{12} \nu_2) + e^{\pm i\delta} \sin\theta_{13} \nu_3$$

Question: $\nu_\mu \rightarrow \nu_e \neq \bar{\nu}_\mu \rightarrow \bar{\nu}_e ?$

Currently: hints at $\sim 2\sigma$ from T2K (NOvA?)

How do oscillation searches probe mass ordering?



Observe **interference effects** of oscill. driven by $\pm \Delta m^2$ with oscill. driven by another quantity Q with known sign. Options:

$$Q \sim \delta m^2$$

medium-baseline reactors in vacuum (JUNO)

$$Q \sim e^- \text{ density}$$

e^- background effects on oscillations (MSW) →

$$Q \sim \nu \text{ density}$$

ν background effects at high density (self-interac.)

Currently: global hints for NO at $\sim 2.5\sigma$ level

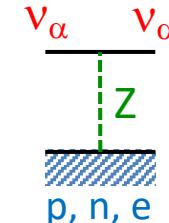
How do ν oscillate in background matter? The MSW effect

(MSW = Mykheev-Smirnov-Wolfenstein)

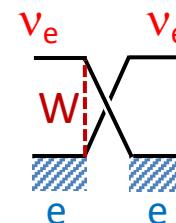
Neutrinos propagating in matter (Earth, Sun, SN...) feel background fermions:
nonzero amplitude that they interact without changing direction

Coherent forward amplitude of $O(G_F)$ - not an absorption process of $O(G_F^2)$!

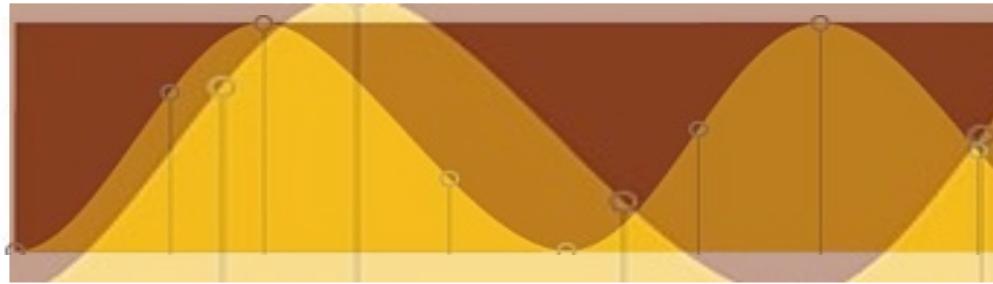
All ν flavors feel all bkgd fermions via NC...



... but only ν_e feels bkgd electrons via CC...

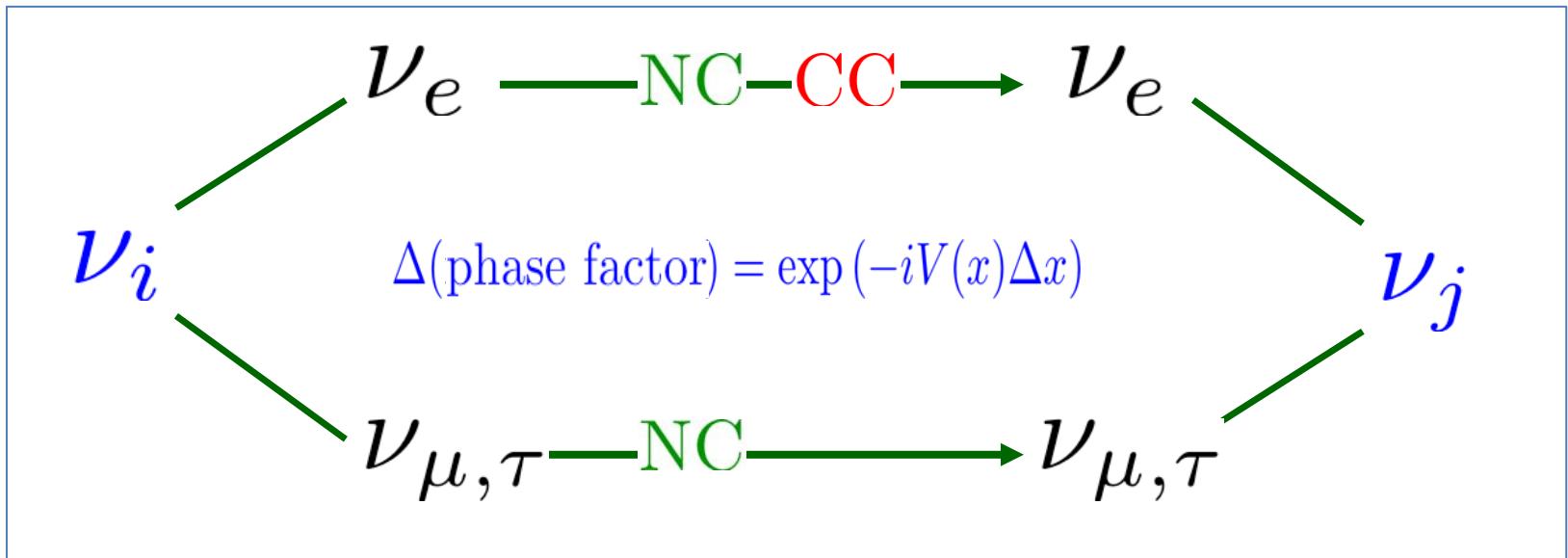


Tiny extra ν_e “interaction energy” or “potential”:
 $V \propto G_F \cdot E \cdot (e^- \text{ density})$



→ density-dependent change in frequency & intensity

Two-slit analogy: one interferometer arm feels a different refraction index

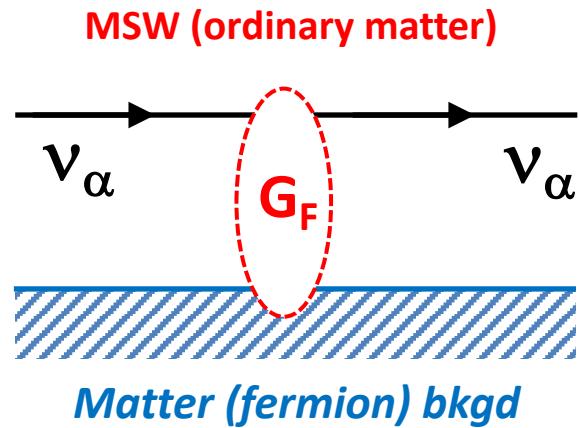


Probability of flavor change may change considerably in matter.

Not necessarily periodic: flavor oscillations → flavor transitions

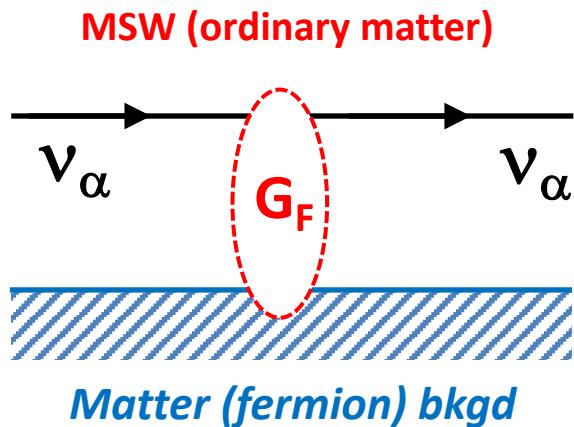
Seen in solar neutrino experiments as Borexino at LNGS

MSW effect: ν 's are the beam;
(n, p, e) are the background



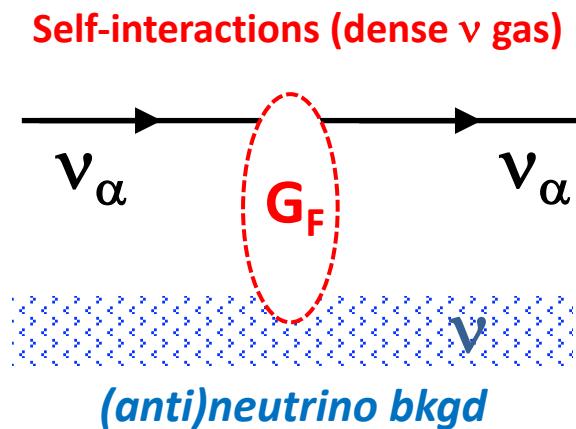
*Relevant for propagation
e.g. in the **Earth**, the **Sun**,
Supernovae (SN), ...*

MSW effect: ν 's are the beam;
(n, p, e) are the background



*Relevant for propagation
e.g. in the **Earth**, the **Sun**,
Supernovae (SN), ...*

"Self-interaction" effects: neutrinos
are both the beam and the bkgd!

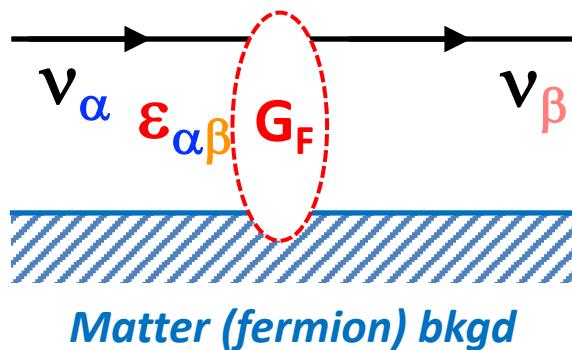


*Relevant for propagation
in very **dense neutrino gas**,
e.g. soon after **SN** explosion,
compact object mergers...*

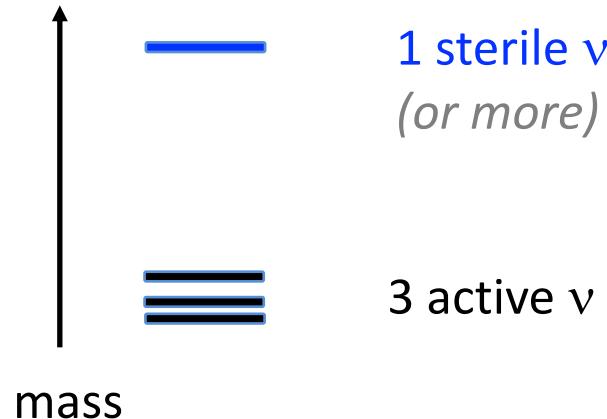
*Also called collective effects,
with highly nonlinear evolution.
May be sensitive to NO/IO.
Important but difficult topic!*

... Beyond standard oscill.? New interactions? More than 3 ν ?

Nonstandard interactions (e.g., FCNC)



Mixing with noninteracting (sterile) ν



Possibly suggested by some anomalous (but controversial) results.

In particular, active-sterile oscillations at O(eV) mass scale remain under investigation via muon-electron neutrino flavor oscillations.

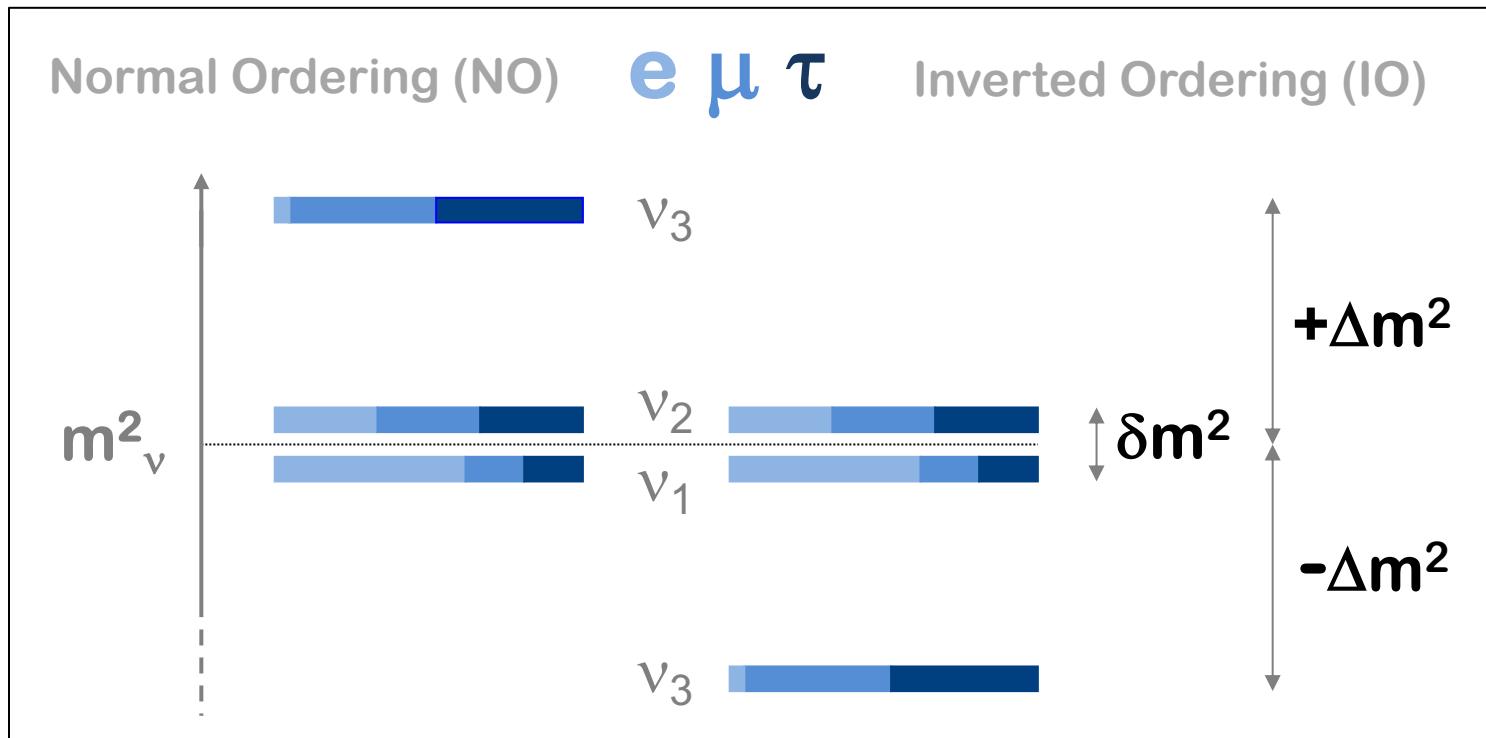
Standard 3ν framework: Recap

5 knowns:

$$\begin{array}{ll} \delta m^2 & \sim 8 \times 10^{-5} \text{ eV}^2 \\ |\Delta m^2| & \sim 2 \times 10^{-3} \text{ eV}^2 \\ \sin^2 \theta_{12} & \sim 0.3 \\ \sin^2 \theta_{23} & \sim 0.5 \\ \sin^2 \theta_{13} & \sim 0.02 \end{array}$$

5 unknowns:

- Oscillations*
- δ CPV Dirac phase
 - $\text{sign}(\Delta m^2) \rightarrow \text{NO/IO}$
 - θ_{23} octant ($>$ or $<$ $\pi/4$?)
- Non-oscillat.*
- absolute mass scale
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Standard 3ν framework: Recap

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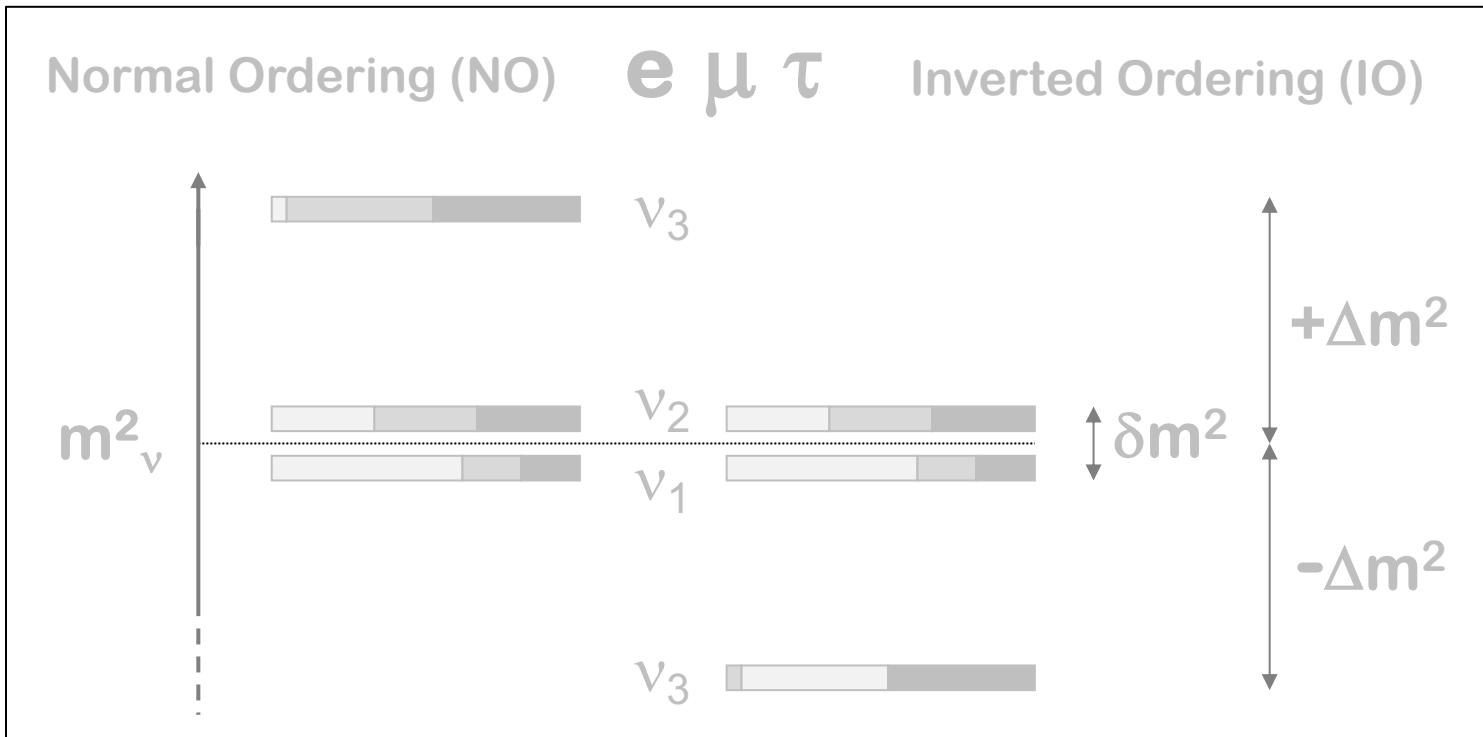
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Oscillations

Non-oscillat.

5 unknowns:

- δ CPV Dirac phase
- $\text{sign}(\Delta m^2) \rightarrow \text{NO/IO}$
- θ_{23} octant ($>$ or $<$ $\pi/4$?)
- absolute mass scale**
- Dirac/Majorana nature**



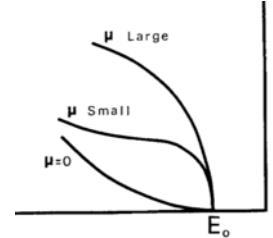
Outline:

- Prologue: ν basic properties
- (Un)knowns in neutrino oscillations
- **Non-oscillation ν observables**
- Learning from neutrinos
- Epilogue: What's in a name?

Absolute neutrino mass observables: (m_β , $m_{\beta\beta}$, Σ)

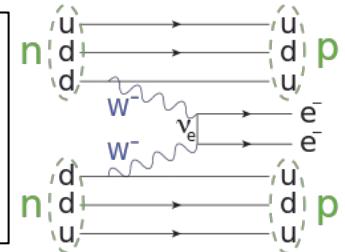
β decay, sensitive to the “effective electron neutrino mass”:

$$m_\beta = [c_{13}^2 c_{12}^2 m_1^2 + c_{13}^2 s_{12}^2 m_2^2 + s_{13}^2 m_3^2]^{\frac{1}{2}}$$



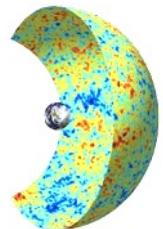
$0\nu\beta\beta$ decay: only if Majorana. “Effective Majorana mass” (+phases):

$$m_{\beta\beta} = |c_{13}^2 c_{12}^2 m_1 + c_{13}^2 s_{12}^2 m_2 e^{i\phi_2} + s_{13}^2 m_3 e^{i\phi_3}|$$



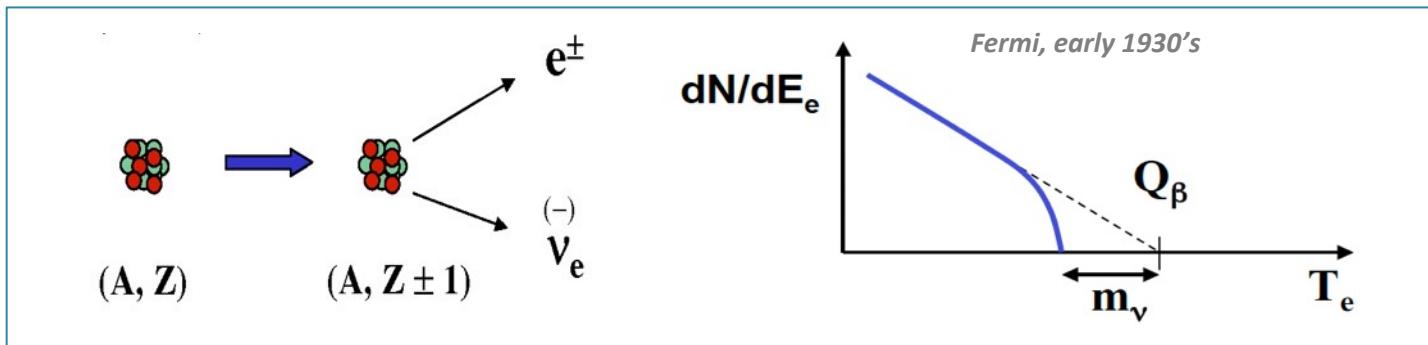
Cosmology: Dominantly sensitive to sum of neutrino masses:

$$\Sigma = m_1 + m_2 + m_3$$



Sensitive to absolute neutrino masses in different ways
May provide additional handles to distinguish NO vs IO

Beta decay: Classic kinematic search for neutrino mass

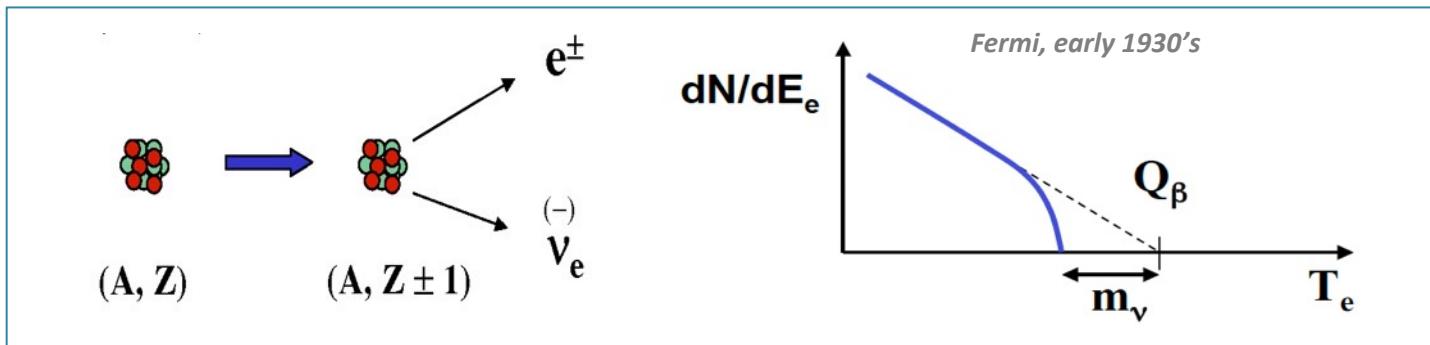


For **three** families ν_i , and for individual masses m_i experimentally unresolved in beta decay: sensitivity to the sum of $m^2(\nu_i)$, weighted by squared mixings $|U_{ei}|^2$ with the electron neutrino. Observable kink parametrized by:

$$m_\beta = [c_{13}^2 c_{12}^2 m_1^2 + c_{13}^2 s_{12}^2 m_2^2 + s_{13}^2 m_3^2]^{\frac{1}{2}}$$

(so-called “effective electron neutrino mass”)

Beta decay: Classic kinematic search for neutrino mass



Tritium (${}^3\text{H}$): low Q and fast decays (Hanna & Pontecorvo 1949)

Upper limits on m_β

1949: $m_\beta < 500$ eV

2023: $m_\beta < 0.8$ eV KATRIN with $O(10^{-4})$ g gaseous tritium

Frontiers:

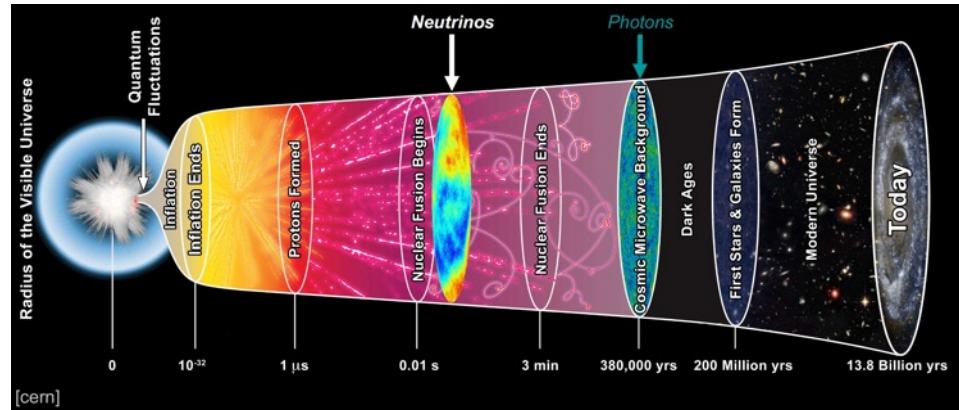
202X: $m_\beta < 0.2$ eV KATRIN final goal

20XY: $m_\beta \sim O(10^{-2})$ eV Needed to fully explore 3ν expectations...

A long-term project aiming at $O(10^2)$ g (!) solid-state Tritium →

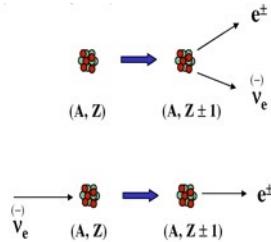


P on-
T ecorvo
O bservatory for
L ight,
E arly-universe,
M assive-neutrino
Y ield



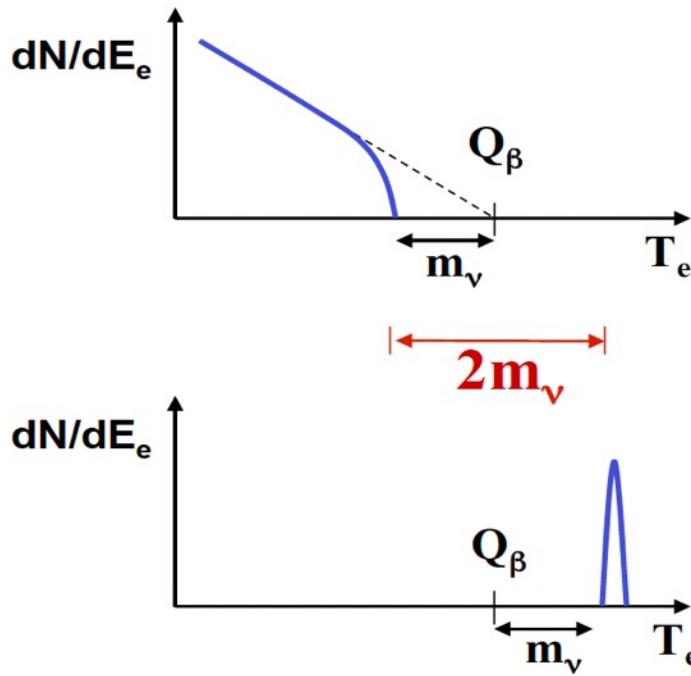
PTOLEMY dream: not just β -decay, but capture big-bang relic ν as “cold” as $T_\nu \sim O(10^{-4})$ eV

Nuclear Beta decay



Neutrino Capture on a Beta Decaying Nucleus

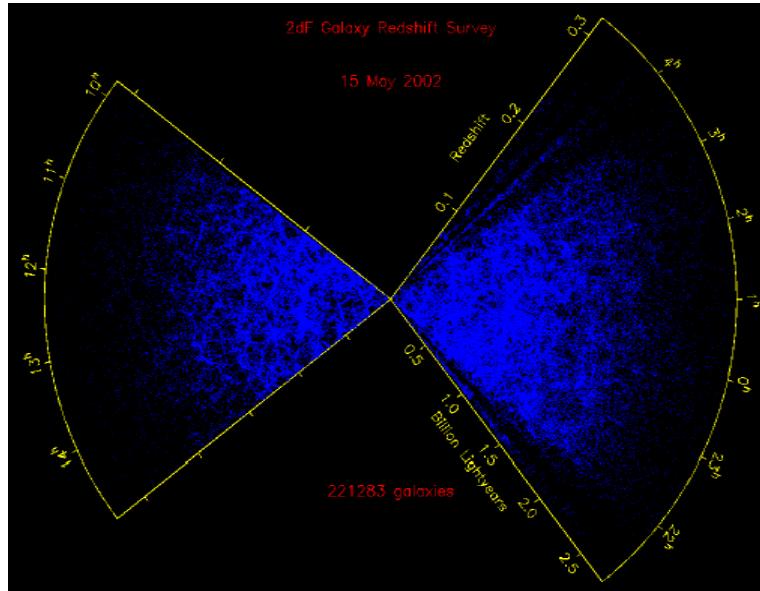
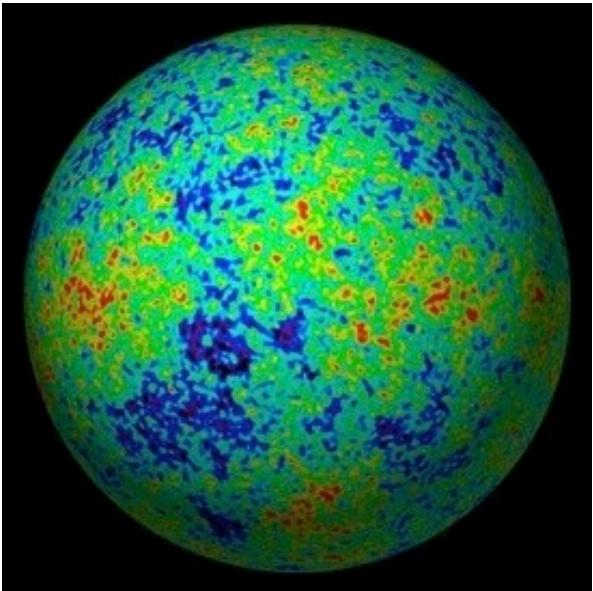
No-threshold reaction!



Cosmology: Dynamical effects of ν masses in evolution of the Universe

We can feel relic ν 's by their gravit. charge: $\Sigma = m_1 + m_2 + m_3$

Σ slightly affects CMB anisotropies + large scale structures (LSS)



From cosmol. data: $\Sigma < 0.1 - 0.2$ eV

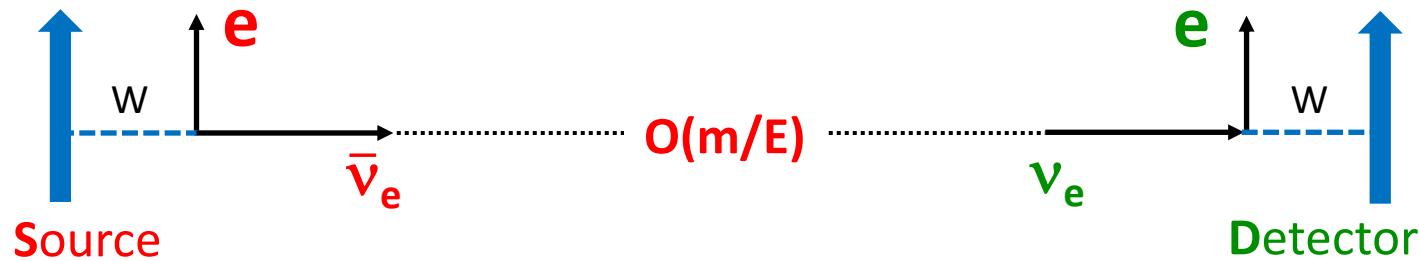
From 3ν oscillations: $\Sigma > 0.06$ eV (NO), > 0.1 eV (IO)

Close to a cosmological discovery of absolute ν mass?

Neutrinoless Double Beta Decay

Reminder:

If neutrinos are Majorana, expect $\bar{\nu} \rightarrow \nu$ transition:

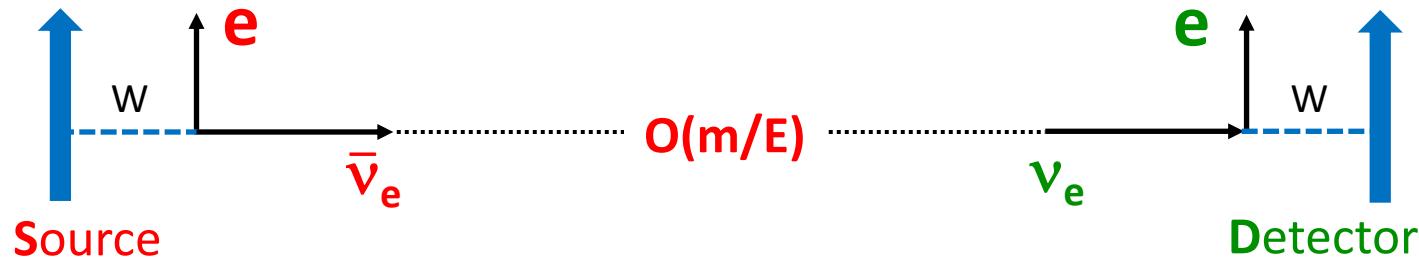


Not found so far, even in the highest-statistics oscill. experiments with 10^7 events
Lesson: to search for Majorana ν , better compare $O(1)$ event to ~ 0 than to 10^N !

Neutrinoless Double Beta Decay

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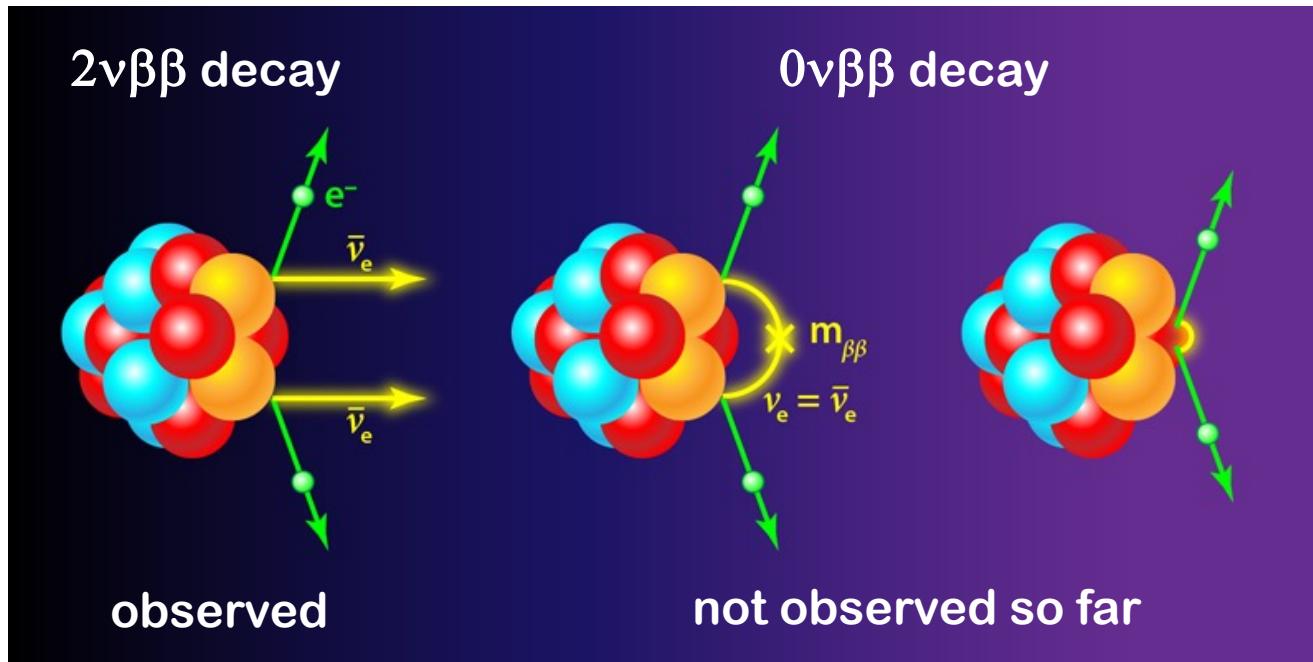
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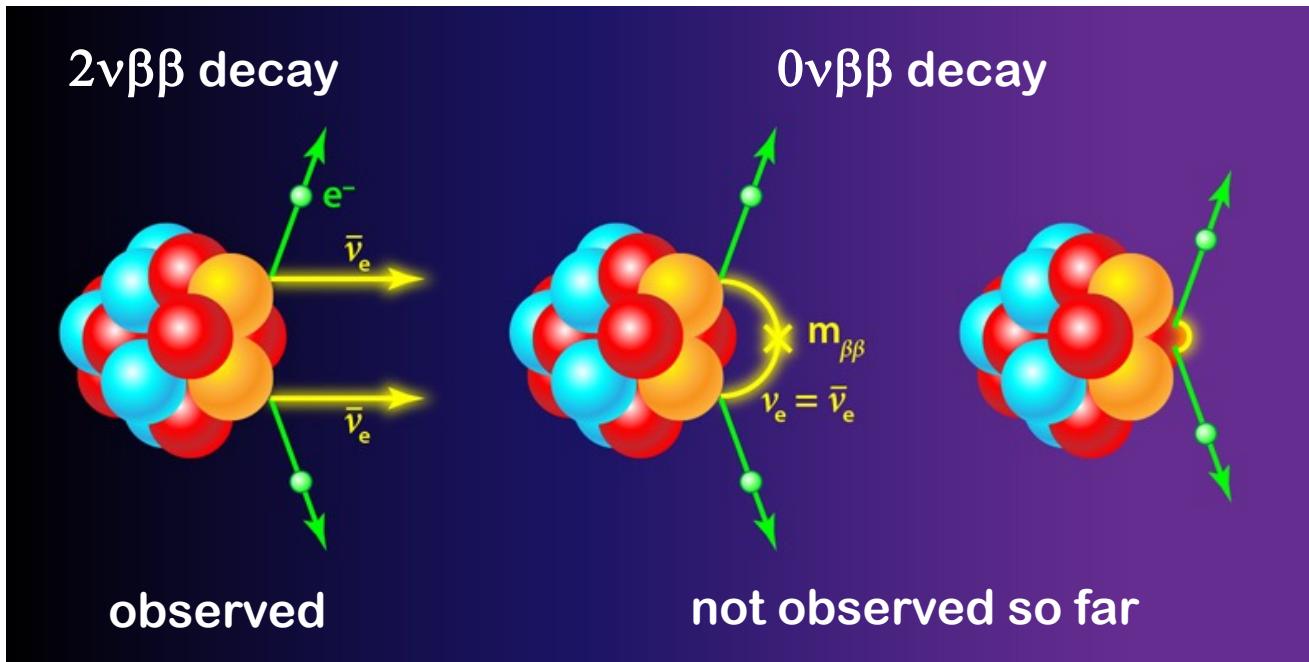
→ Look for rare $\nu \rightarrow \bar{\nu}$ transitions induced by Majorana ν , on top of small or zero bkgd

Only known realistic process of this kind: Neutrinoless double beta decay ($0\nu\beta\beta$)



Very rare: “weak squared” and suppressed by m/E
Possible if and only if neutrinos are Majorana.

Only known realistic process of this kind: Neutrinoless double beta decay ($0\nu\beta\beta$)

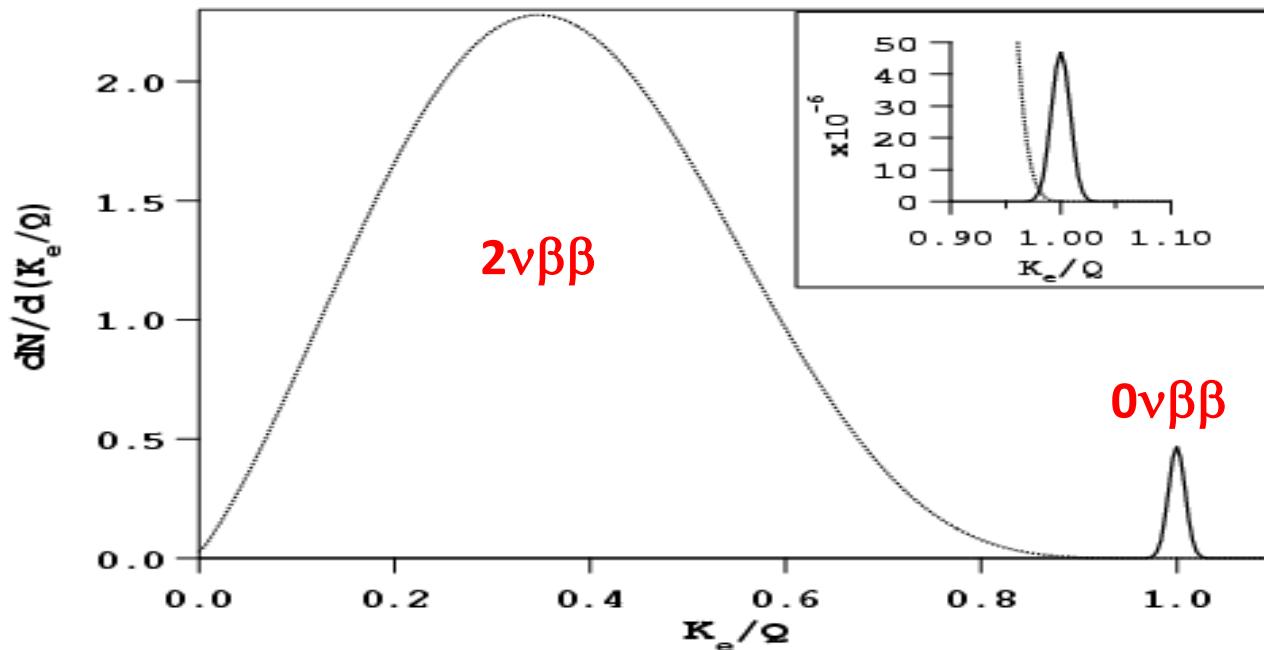


Very rare: “weak squared” and suppressed by m/E
Possible if and only if neutrinos are Majorana.

Discovery of $0\nu\beta\beta$ and Majorana neutrino masses would have profound theoret. implications.
Standard Model (SM): Higgs mechanisms → Dirac fermions. Majorana neutrinos → Beyond SM!
 $0\nu\beta\beta$ “creates” charged leptons: theory links with leptogenesis and new (HE) physics scales

Experimentally: Look at sum energy of both electrons

Need to see the $0\nu\beta\beta$ line emerge above bkgd, at endpoint spectrum of “conventional” $2\nu\beta\beta$ decay.

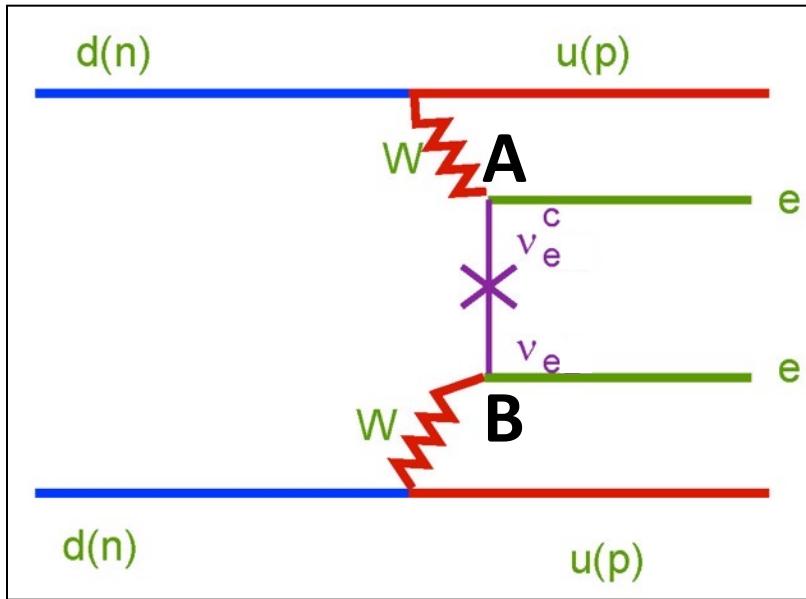


$0\nu\beta\beta$ decay rate \propto inverse half life ($1/T$)

Current expt generation probing $T \sim O(10^{26})$ y.

Ton-scale projects (LEGEND, CUPID, nEXO) aiming at: $T \sim O(10^{28})$ y

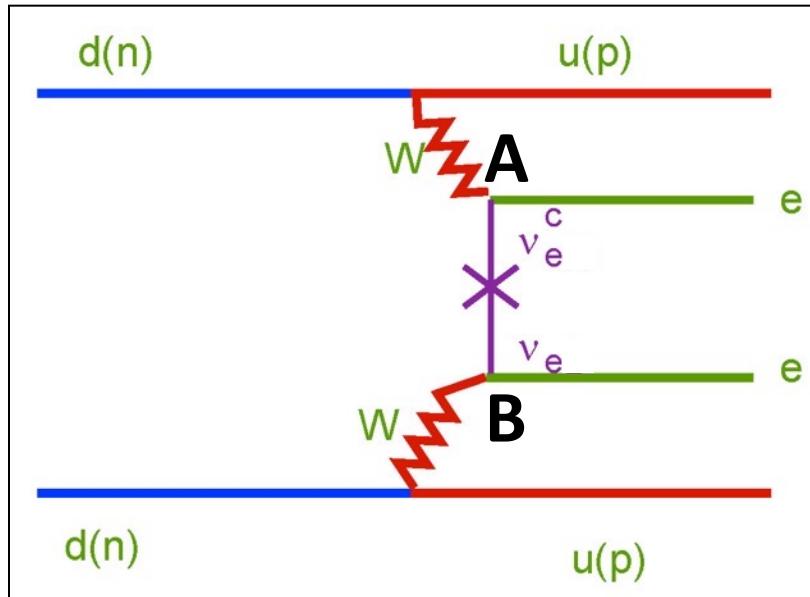
Neutrinoless double beta decay at elementary particle level



Can occur only for Majorana neutrinos. Intuitive picture:

- 1) A RH antineutrino is emitted at point "A" together with an electron
- 2) If it is massive, at $O(m/E)$ it develops a LH component (not possible if Weyl)
- 3) If neutrino=antineutrino, this component is a LH neutrino (not possible if Dirac)
- 4) The LH (Majorana) neutrino is absorbed at "B" where a 2nd electron is emitted

Neutrinoless double beta decay at elementary particle level



← mixing of ν_e with ν_i

← mass of ν_i [$O(m/E)$]

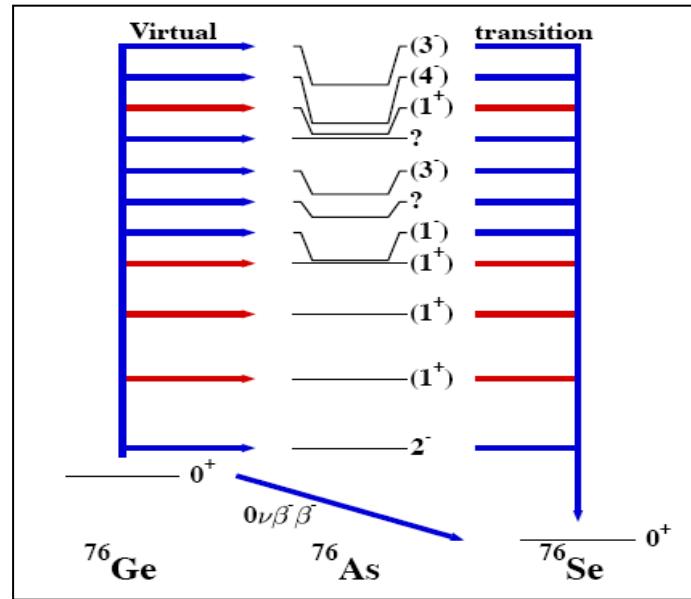
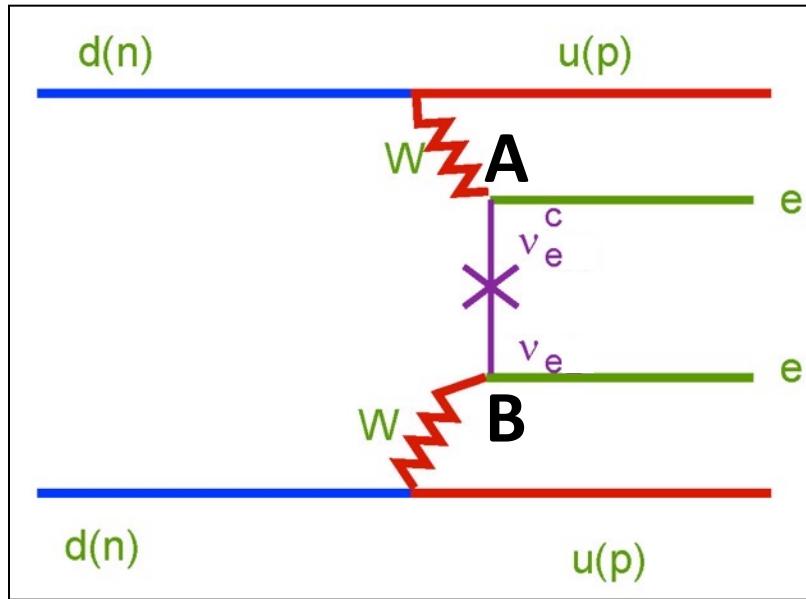
← mixing of ν_i with ν_e

(times an unknown ν_i phase)

Summing up for three massive neutrinos: Amplitude \propto “effective Majorana mass”

$$m_{\beta\beta} = |c_{13}^2 c_{12}^2 m_1 + c_{13}^2 s_{12}^2 m_2 e^{i\phi_2} + s_{13}^2 m_3 e^{i\phi_3}|$$

Neutrinoless double beta decay at nuclear level: complicated...



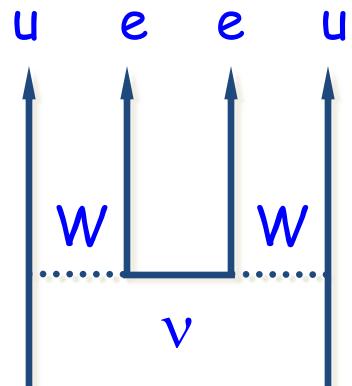
In case of positive decay signal, a major concern is the accuracy of the **nuclear matrix element $|M|$** , rather than the expt. error on the decay half life:

$$T_i^{-1} = G_i |M'_i|^2 m_{\beta\beta}^2$$

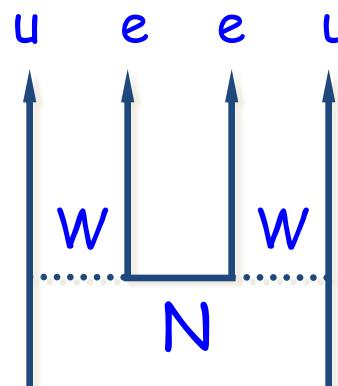
↑ ↑ ↓

Half-life Phase space Matrix element

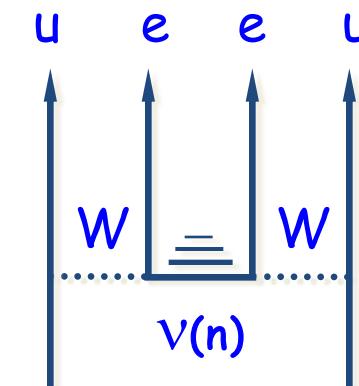
Note: $0\nu\beta\beta$ decays might also be induced by nonstandard physics



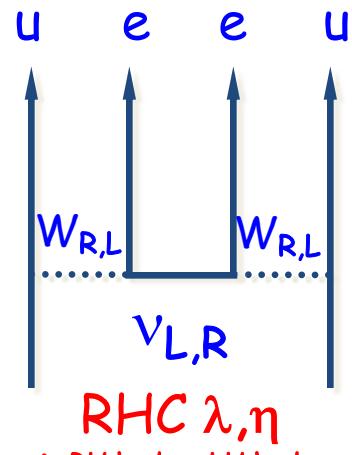
Standard



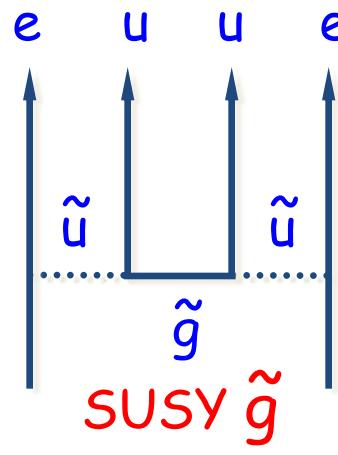
Heavy ν



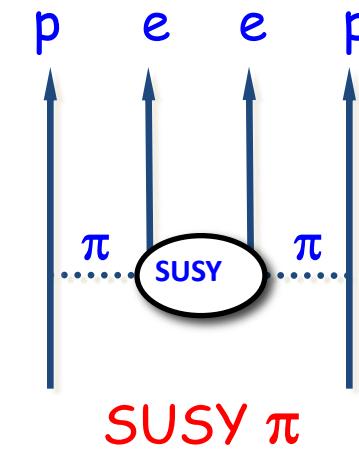
Kaluza-Klein
(KK ± 1 Brane: $a = 10^{\pm 1}/\text{GeV}$)



RHC λ, η
 $\lambda = \text{RH had}, \eta = \text{LH had}$

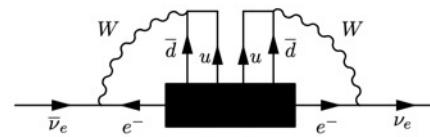


SUSY \tilde{g}



SUSY π

In any case, $0\nu\beta\beta$ decay implies Majorana ν :



Combining 3ν constraints from:

Oscillation data

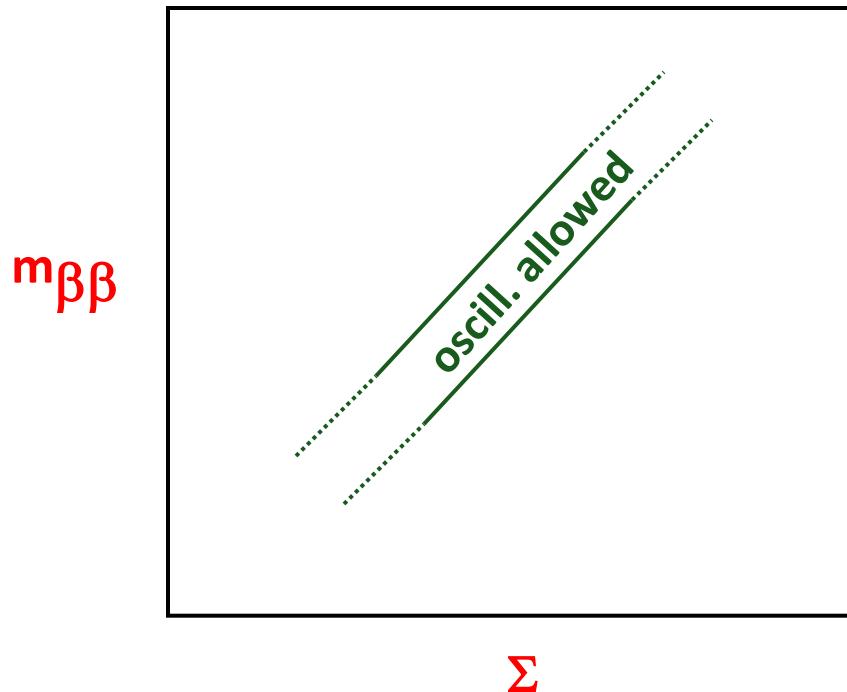
+

$$m_\beta = [c_{13}^2 c_{12}^2 m_1^2 + c_{13}^2 s_{12}^2 m_2^2 + s_{13}^2 m_3^2]^{\frac{1}{2}}$$

$$m_{\beta\beta} = |c_{13}^2 c_{12}^2 m_1 + c_{13}^2 s_{12}^2 m_2 e^{i\phi_2} + s_{13}^2 m_3 e^{i\phi_3}|$$

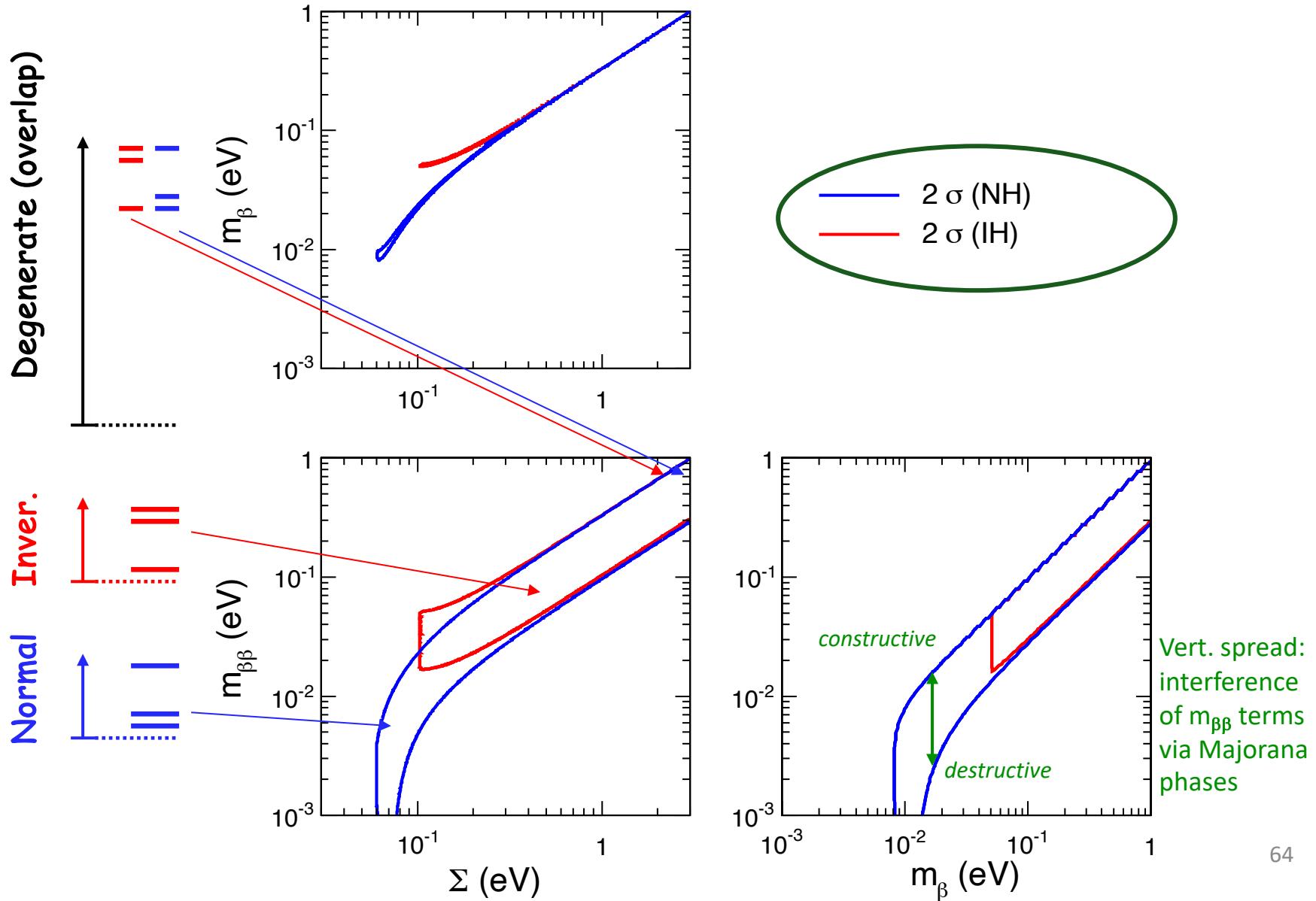
$$\Sigma = m_1 + m_2 + m_3$$

Interplay: **Oscillations** fix the **mass² splittings**, and thus induce **positive correlations** between any pair of the three observables **(m_β , $m_{\beta\beta}$, Σ)**, e.g.:

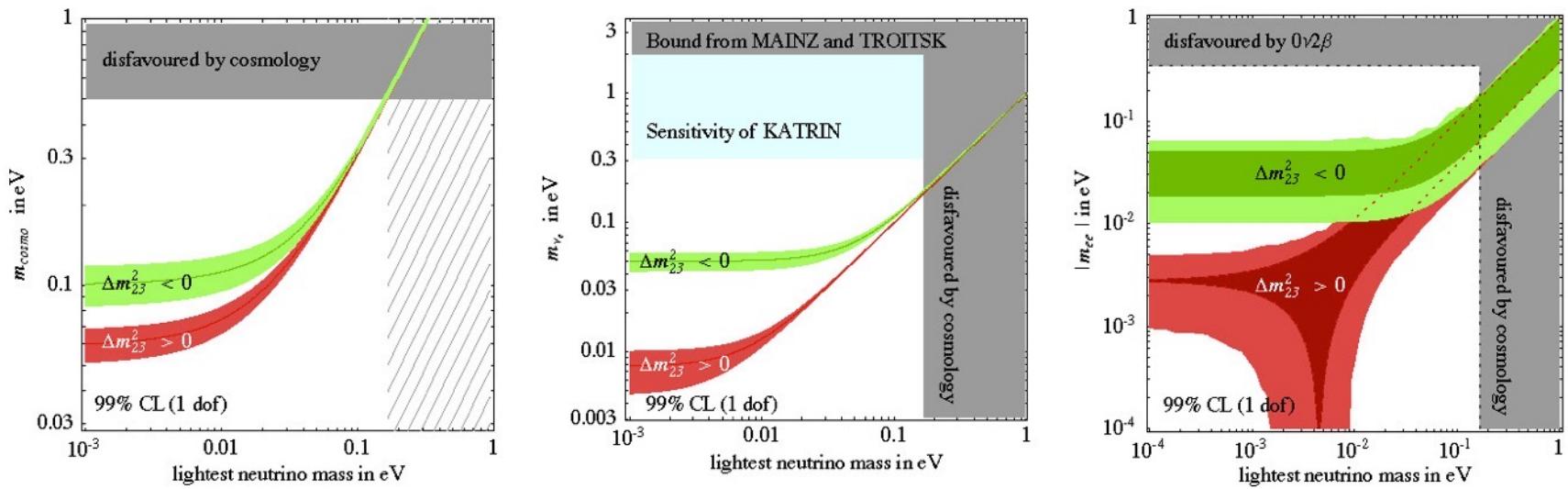


i.e., if one observable increases, the other one (typically) must increase to match mass splitting

Oscillation data constrain the $(m_\beta, m_{\beta\beta}, \Sigma)$ parameters within two bands:

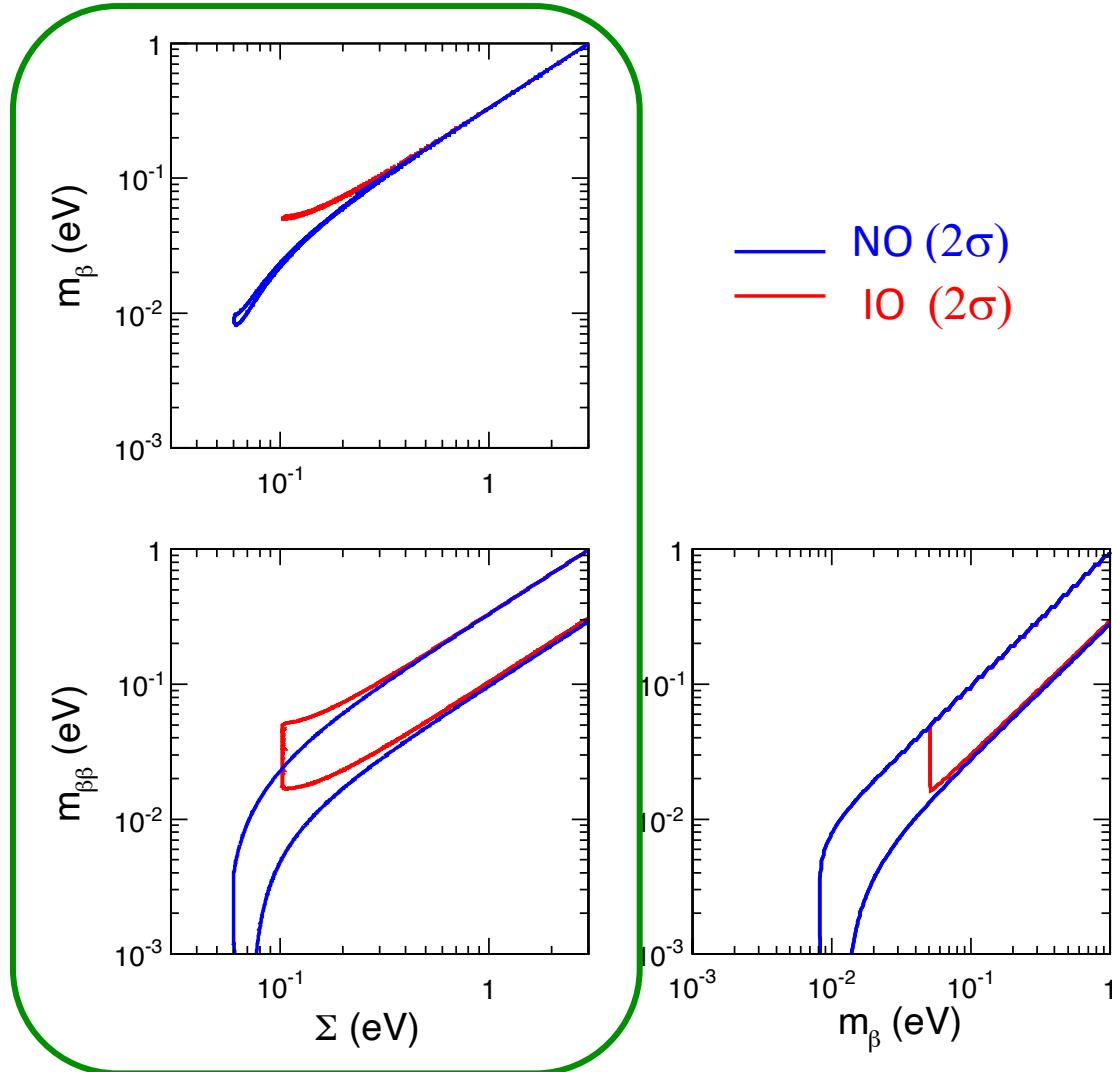


Previous plots project away the “unobservable” lightest neutrino mass from popular graphs like:

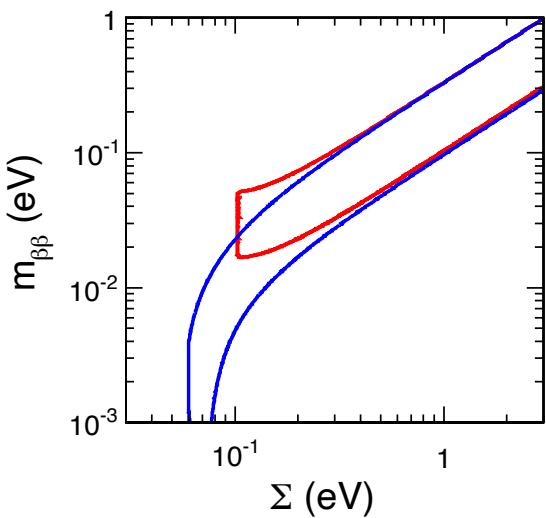
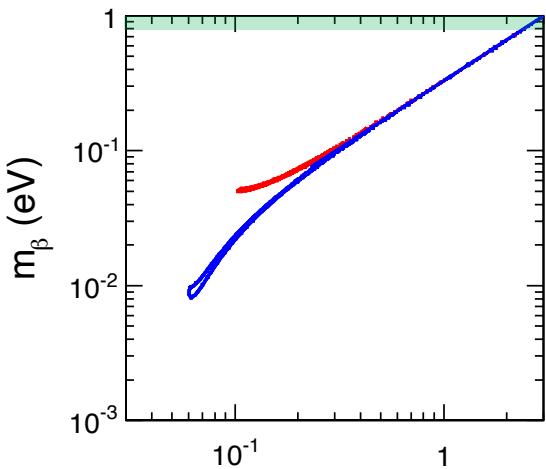


Taken from Strumia and Vissani, 2006

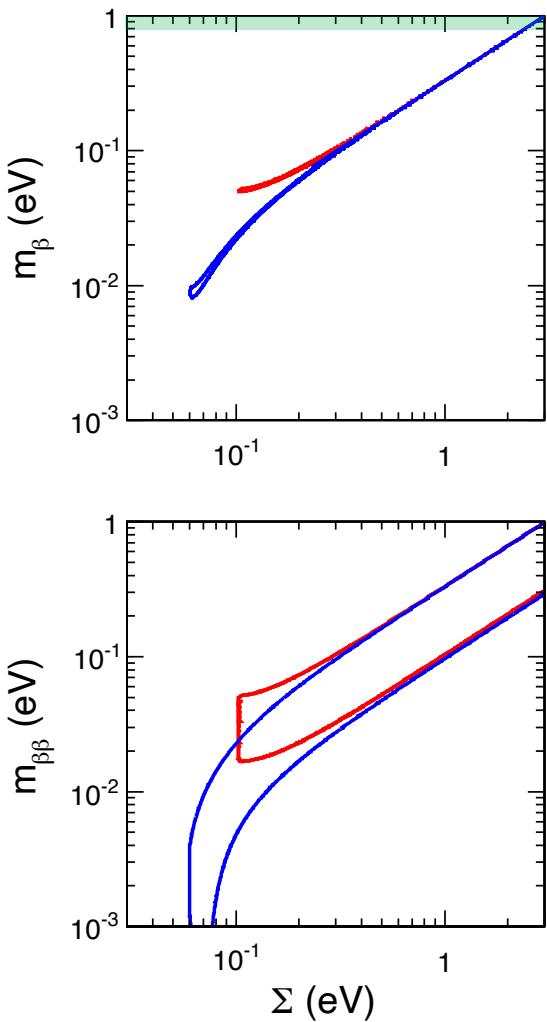
Absolute mass observables: currently, only upper bounds...



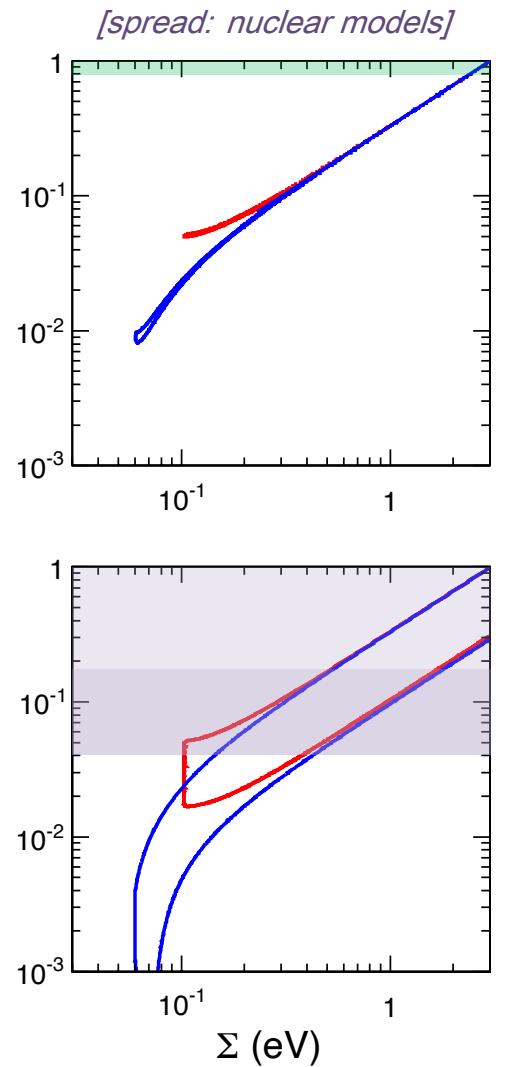
Focus on these planes



β : KATRIN

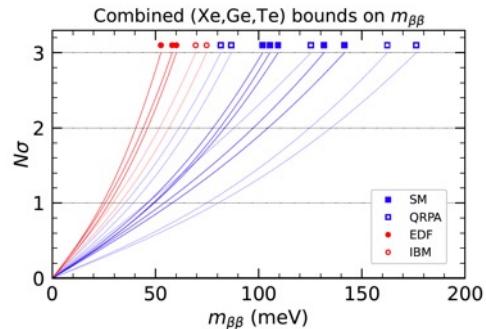


$0\nu\beta\beta$: KL-Zen, Exo,
GERDA, Cuore...

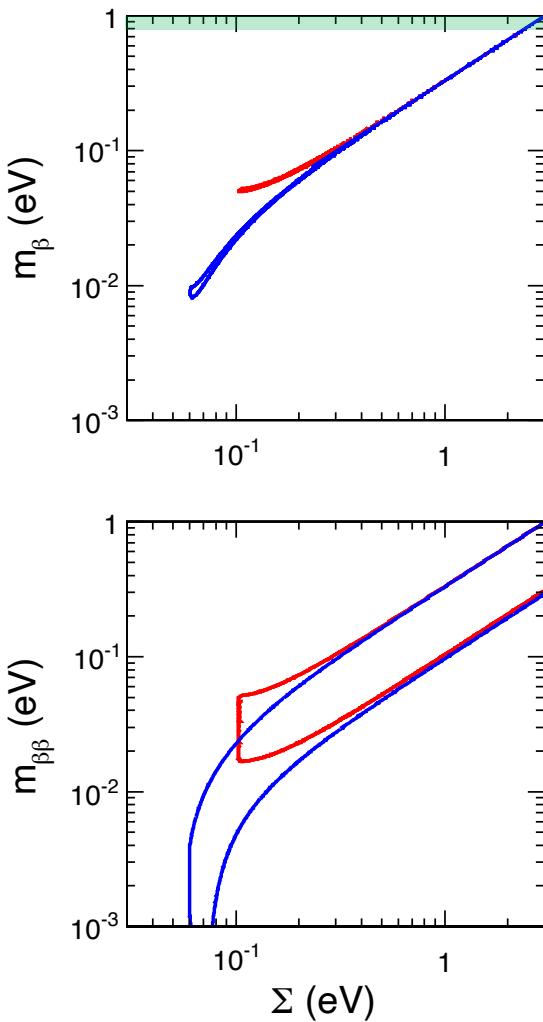


[spread: nuclear models]

E.g., spread of upper bounds
from Xe+Ge+Te data by using
15 nuclear matrix elements
from 4 classes of nucl. models.
e-print 2204.09569

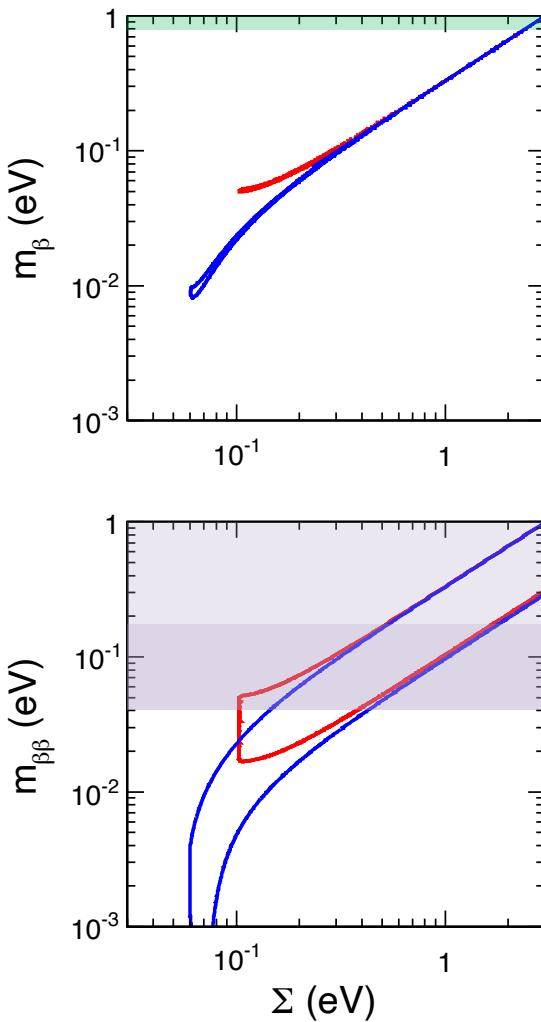


β : KATRIN

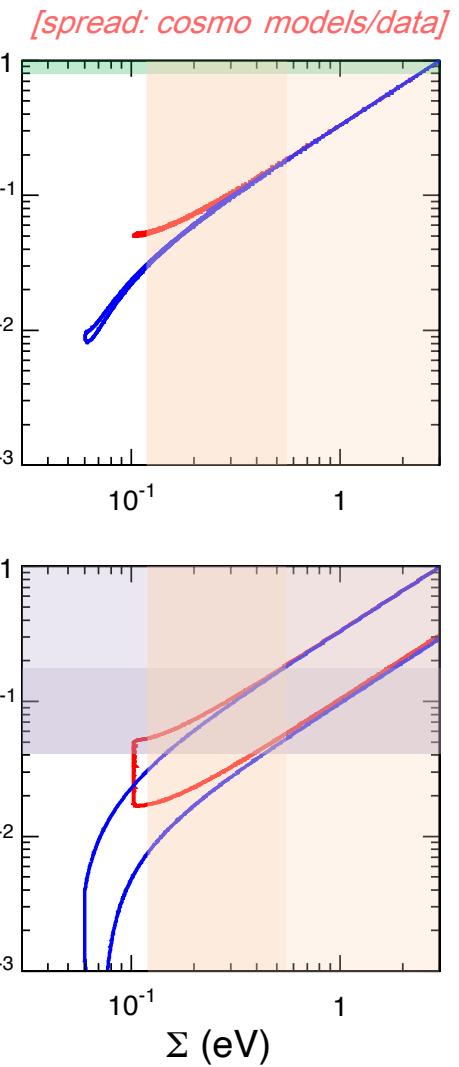


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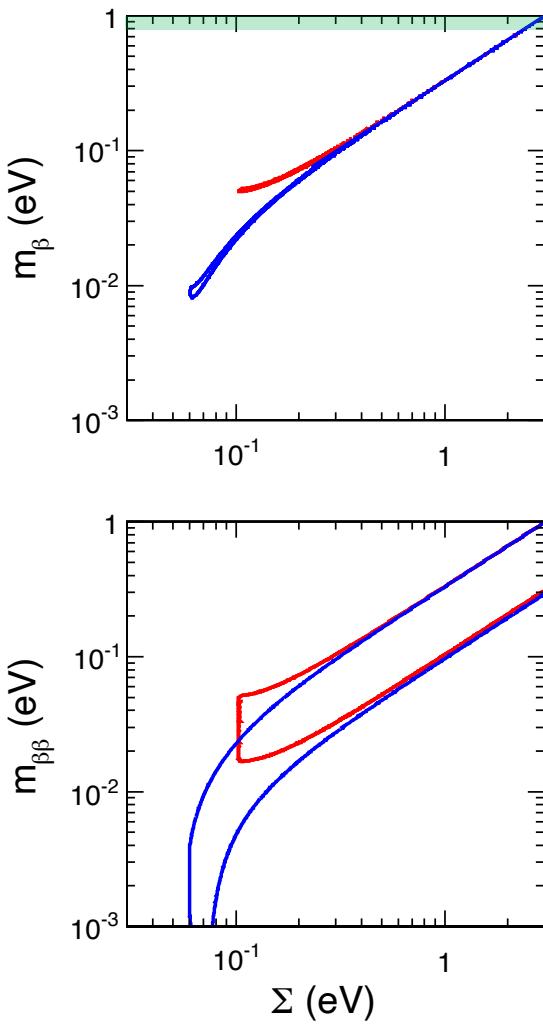
[spread: nuclear models]



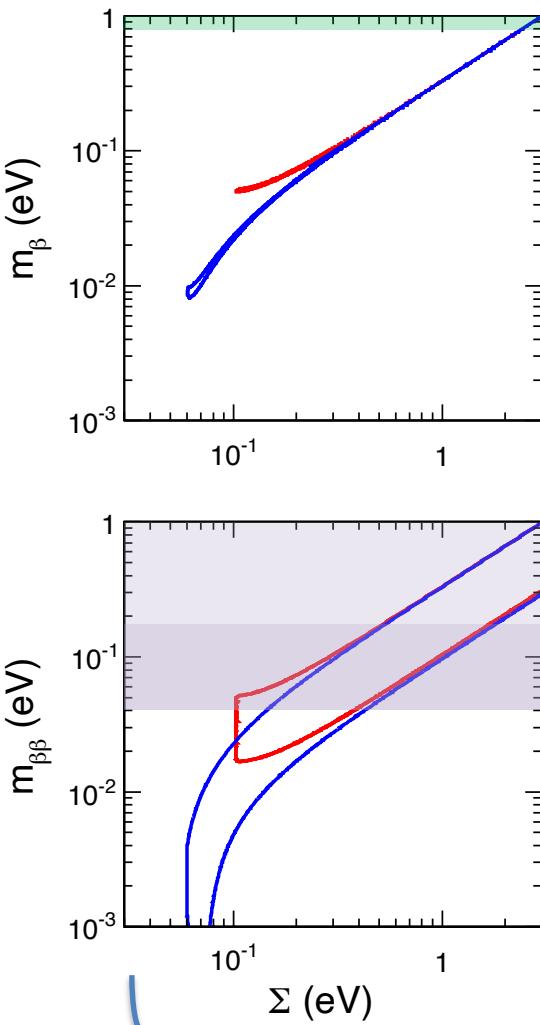
Σ : Planck, BAO,
lensing ...



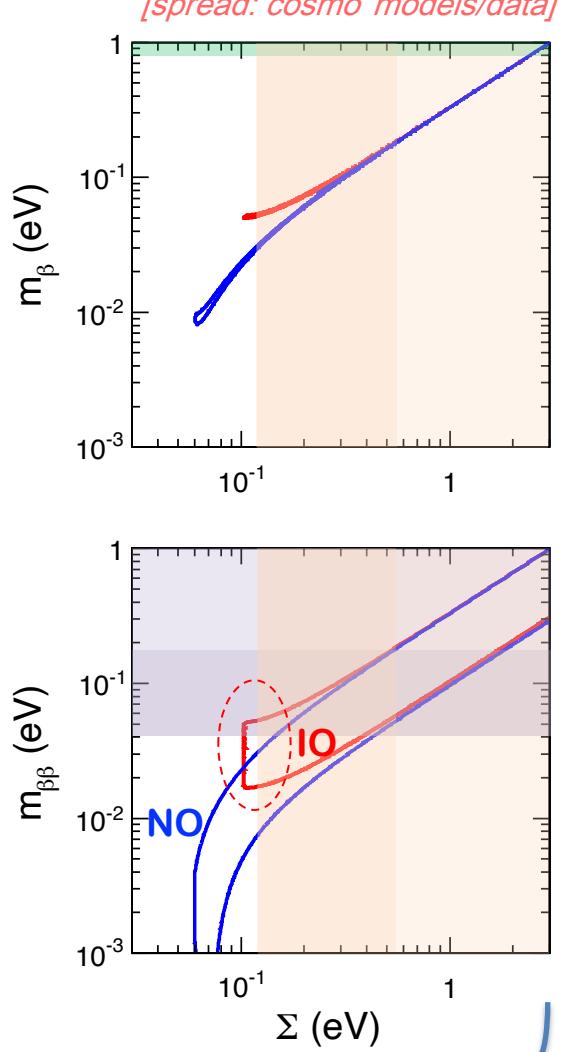
β : KATRIN



$0\nu\beta\beta$: KL-Zen, Exo,
GERDA, Cuore...

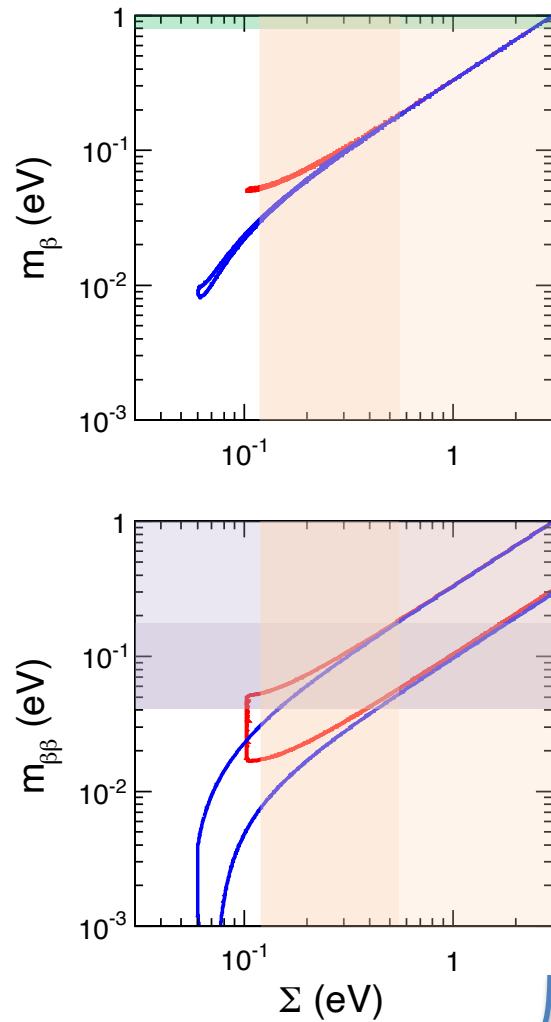
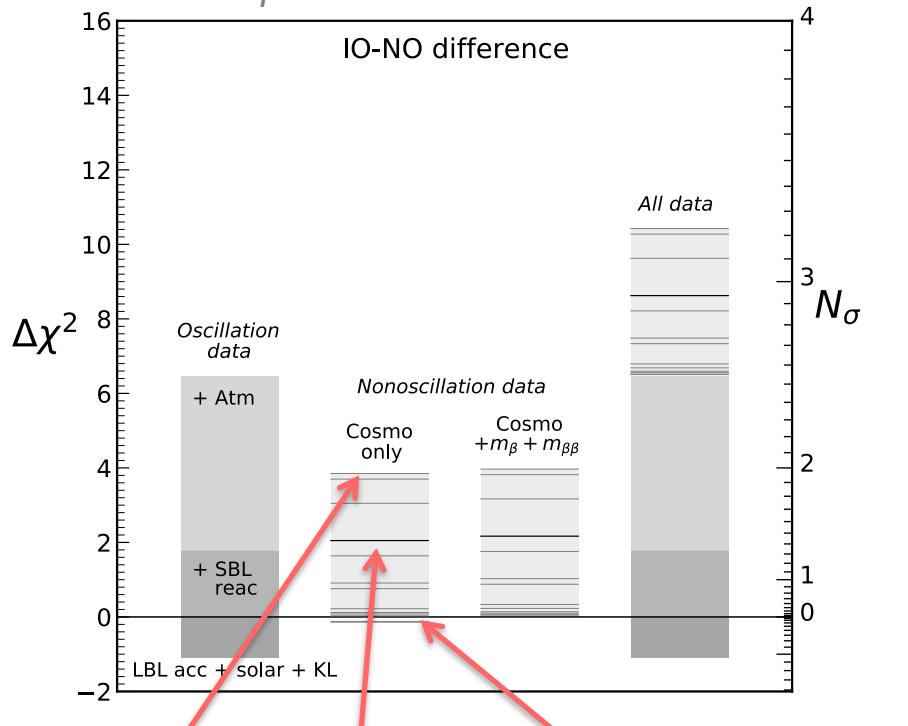


Σ : Planck, BAO,
lensing ...



IO “under pressure” but not excluded yet
 $0\nu\beta\beta$ expts need to probe IO fully (+ part of NO)

From Capozzi+ 2107.00532

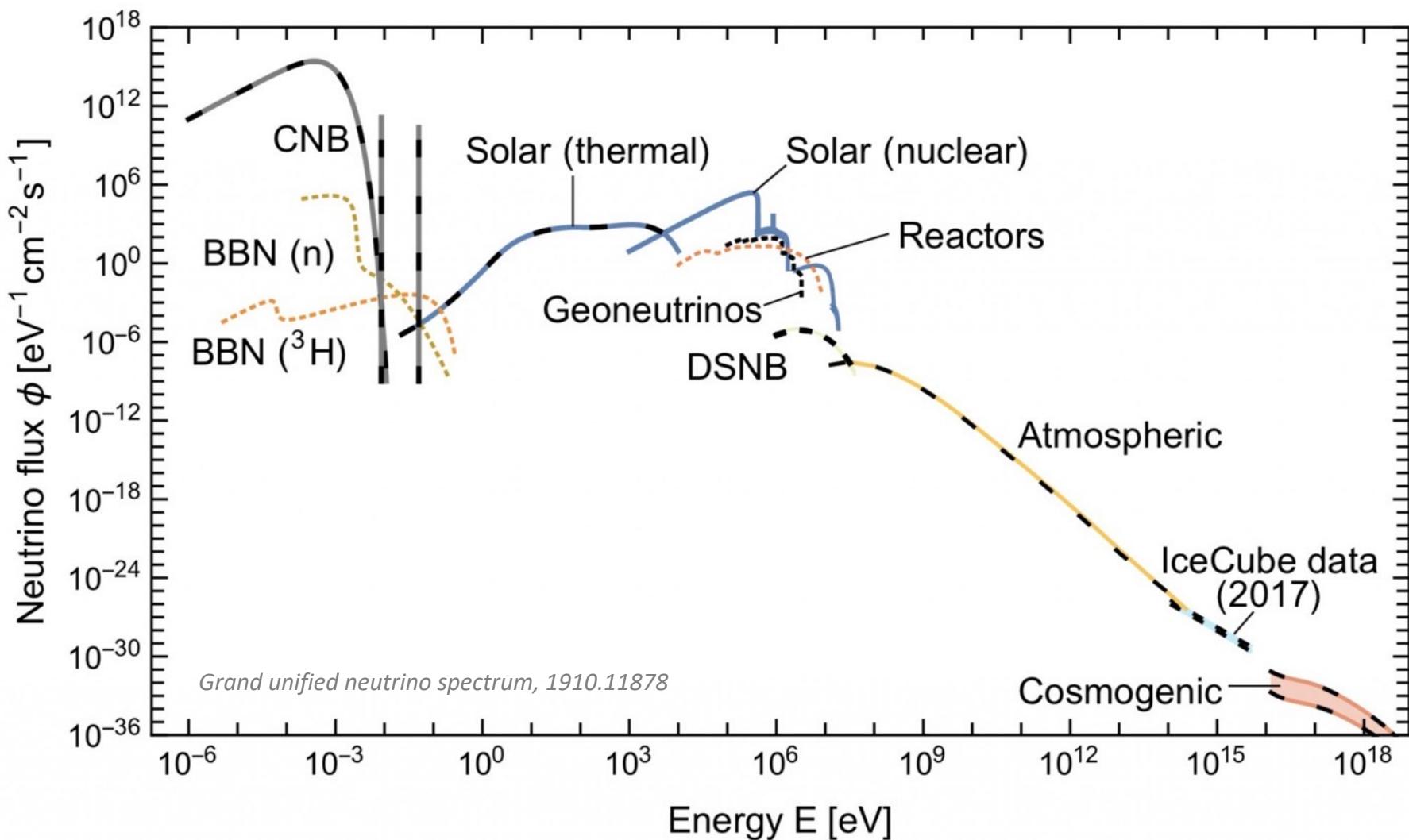


IO currently disfavored at $\sim 3\sigma$ by combining oscillation + nonoscillation data

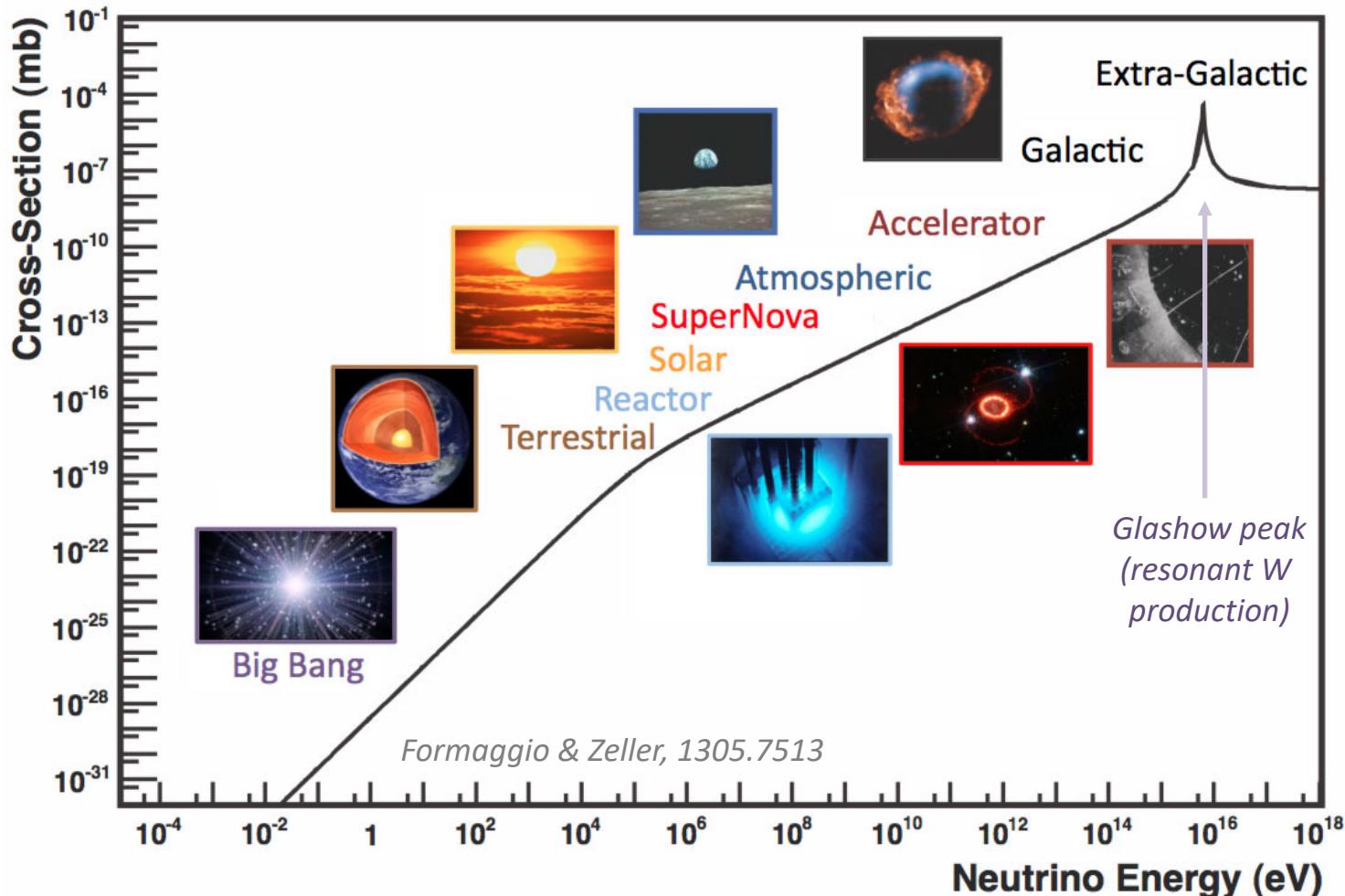
Outline:

- Prologue: ν basic properties
- (Un)knowns in neutrino oscillations
- Non-oscillation ν observables
- **Learning from neutrinos**
- Epilogue: What's in a name?

Grand view of the neutrino landscape: flux vs energy

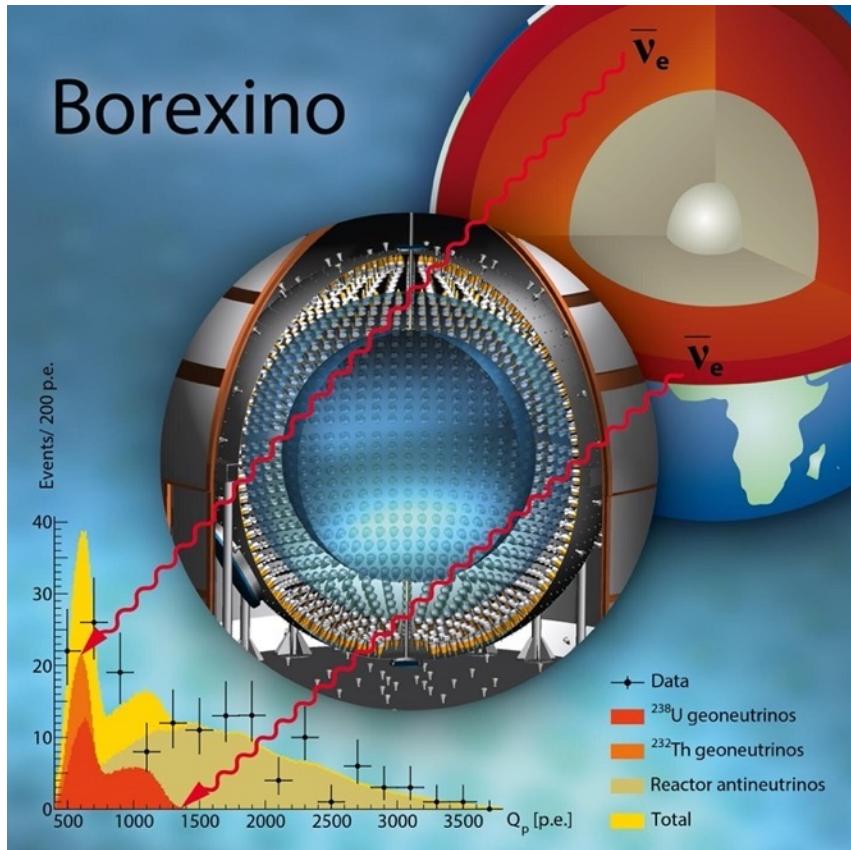


Typical neutrino cross section ($\bar{\nu}_e e^- \rightarrow \bar{\nu}_e e^-$) vs energy



NOTE: ν absorption in Earth (!) relevant >100 TeV & Glashow event @ ~ 6 PeV: seen in IceCube

Learning from ν about the Earth's interior and heat budget...



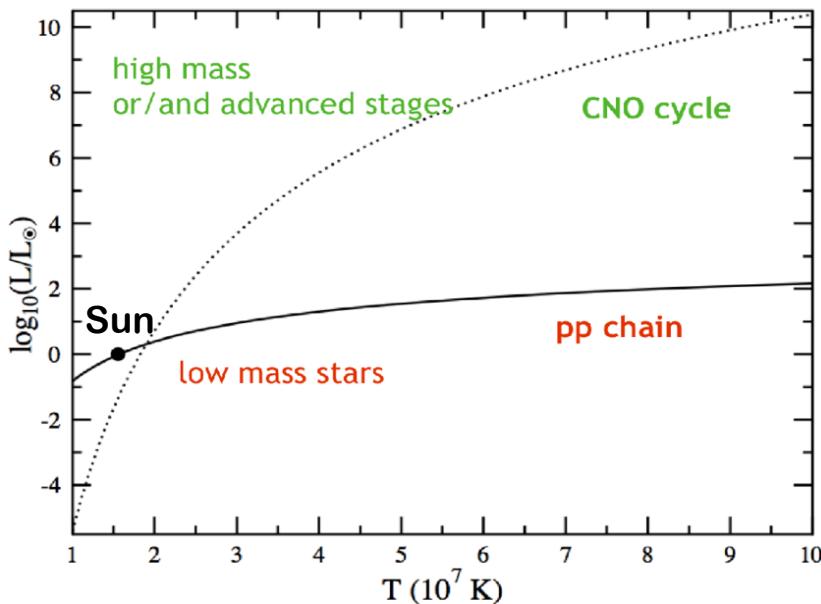
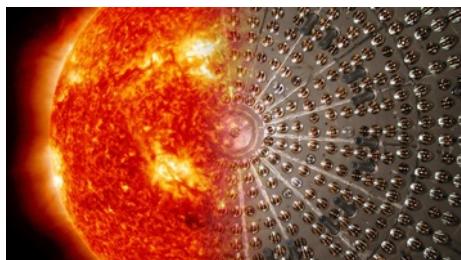
Geoneutrinos from
U, Th decays in the
Earth observed in
KamLAND & Borexino

Test global
geo-chemical
models

Frontiers:
JUNO and other
large-volume,
low-energy detectors

Co-messengers of elastic waves (seismology)

Learning from ν about the Solar nuclear reaction chains ...



CNO neutrinos recently observed in Borexino

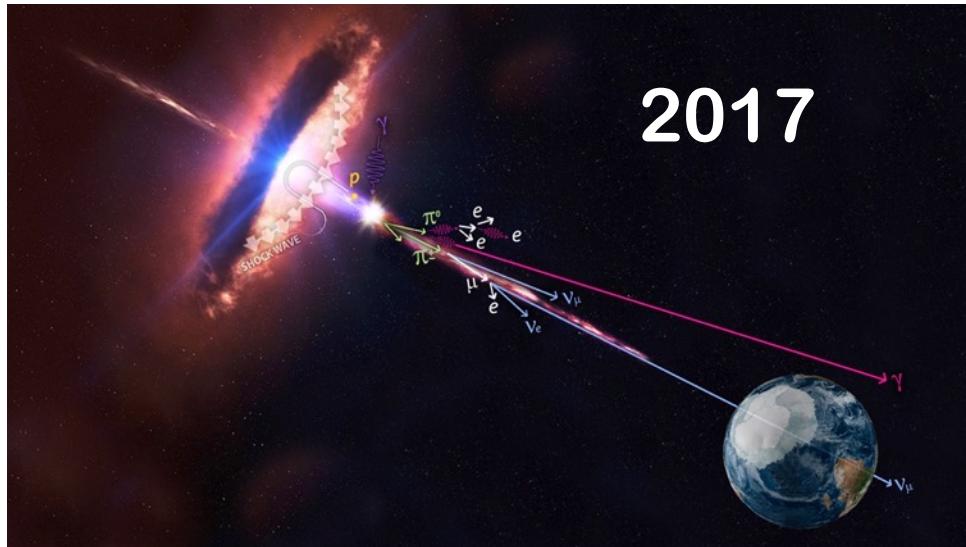
Test solar and stellar evolution models*

**Also w/ nuclear astrophysics inputs, e.g., from LUNA expt.*

Frontiers:
JUNO, DUNE, HK
and other large-volume,
low-energy detectors

Co-messengers of elastic waves (helioseismology)
Co-messengers of photons

Learning from ν about violent “particle accelerator” sources...



RESEARCH ARTICLE | NEUTRINO ASTROPHYSICS

Multimessenger observations of a flaring blazar coincident with high-energy neutrino IceCube-170922A

The IceCube Collaboration, *Fermi-LAT*, *MAGIC*, *AGILE*, *ASAS-SN*, *HAWC*, *H.E.S.S.*, *INTEGRAL*, *Kanata*, *Kiso*, *Kapteyn*, *Liverpool Telescope*, *Subaru*, *Swift/NuSTAR*, *VERITAS*, *VLA/17B-403* teams*,†

✉†Email: analysis@icecube.wisc.edu

* The full lists of participating members for each team and their affiliations are provided in the supplementary materials.
- Hide authors and affiliations

Science 12 Jul 2018:
eaat1378
DOI: 10.1126/science.aat1378

Flourishing field
of multimessenger
astronomy with
neutrino telescopes
(IceCube, KM3Net...)

New sources...

Co-messengers of photons

Co-messengers of gravitational waves?... hopefully in the future!

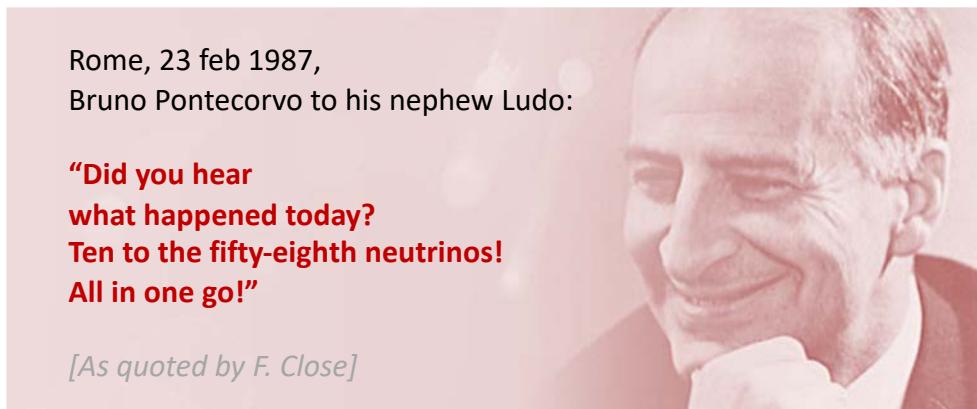
Learning from ν about the death of stars (core-collapse SN)



Rome, 23 feb 1987,
Bruno Pontecorvo to his nephew Ludo:

"Did you hear
what happened today?
Ten to the fifty-eighth neutrinos!
All in one go!"

[As quoted by F. Close]



Frontiers:

- Observe diffuse SN ν bkgd flux
- Understand collective ν effects
- Be ready for next near SN explosion!

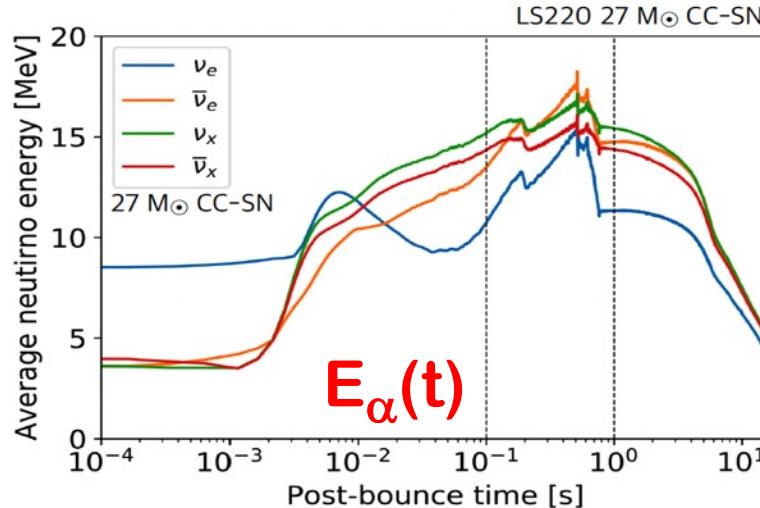
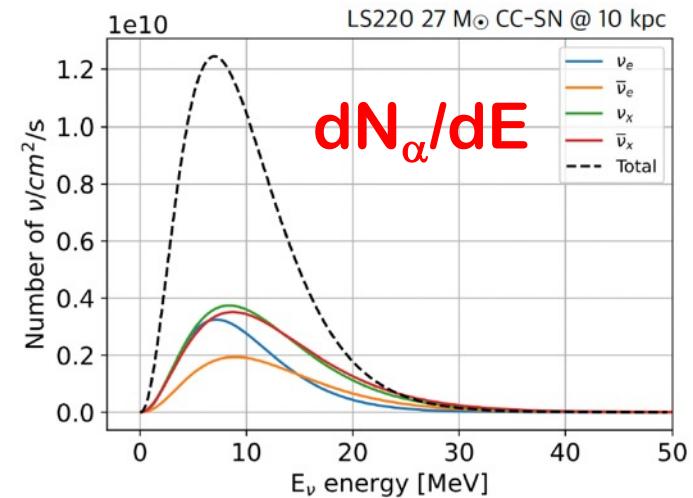
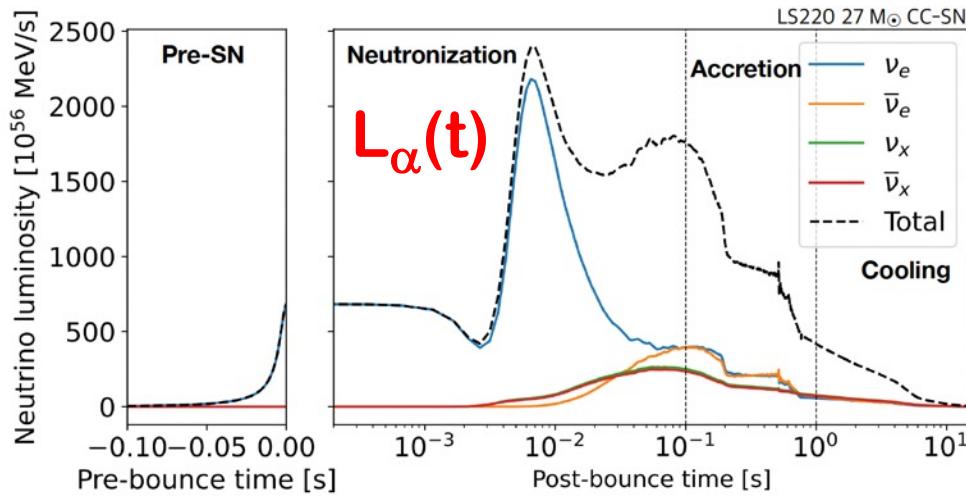
Co-messengers of photons

...as anticipated by B. Pontecorvo, JETP 36, 1625 (1959)

Co-messengers of gravitational waves?... hopefully in the future!

Towards the next SN event: frontiers in ν detection

Typical (simulated, unoscillated) SN neutrino signals for different flavors α :



Flavor transformations
re-weight such spectra
(and may even alter the
SN evolution ← feedback)

X-sections differ with $\alpha \rightarrow$
CC: $\alpha = e$ NC: $\alpha = e, \mu, \tau$

Neutrino interactions in the ~10 MeV range

Slide from
Kate Scholberg

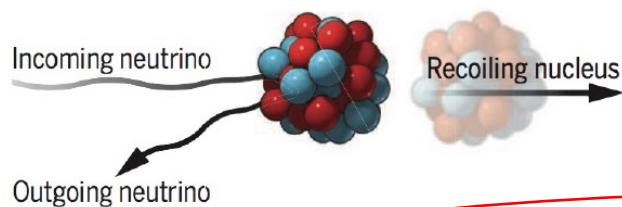
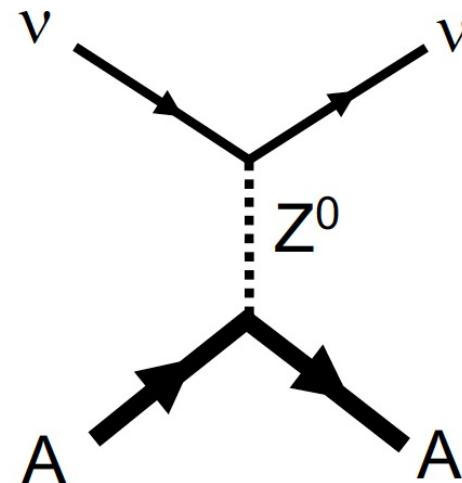
	Electrons	Protons	Nuclei
Charged current	<p>Elastic scattering</p> $\nu + e^- \rightarrow \nu + e^-$	<p>Inverse beta decay</p> $\bar{\nu}_e + p \rightarrow e^+ + n$	$\nu_e + (N, Z) \rightarrow e^- + (N - 1, Z + 1)$ $\bar{\nu}_e + (N, Z) \rightarrow e^+ + (N + 1, Z - 1)$
Neutral current	<p>Useful for pointing</p>	<p>Elastic scattering</p> $\nu \rightarrow \nu$ <p>very low energy recoils</p>	$\nu + A \rightarrow \nu + A^*$ $\nu + A \rightarrow \nu + A$ <p>Coherent elastic (CEvNS)</p>

Coherent Elastic neutrino Nucleus Scattering = CEvNS
(read: "sevens"): Predicted in 1974, observed in 2017

The gentlest (coherent) interaction of a neutrino with a nucleus...

$$\nu + A \rightarrow \nu + A$$

A neutrino smacks a nucleus via exchange of a Z , and the nucleus recoils as a whole; **coherent** up to $E_\nu \sim 50$ MeV



Nucleon wavefunctions in the target nucleus are **in phase with each other** at low momentum transfer

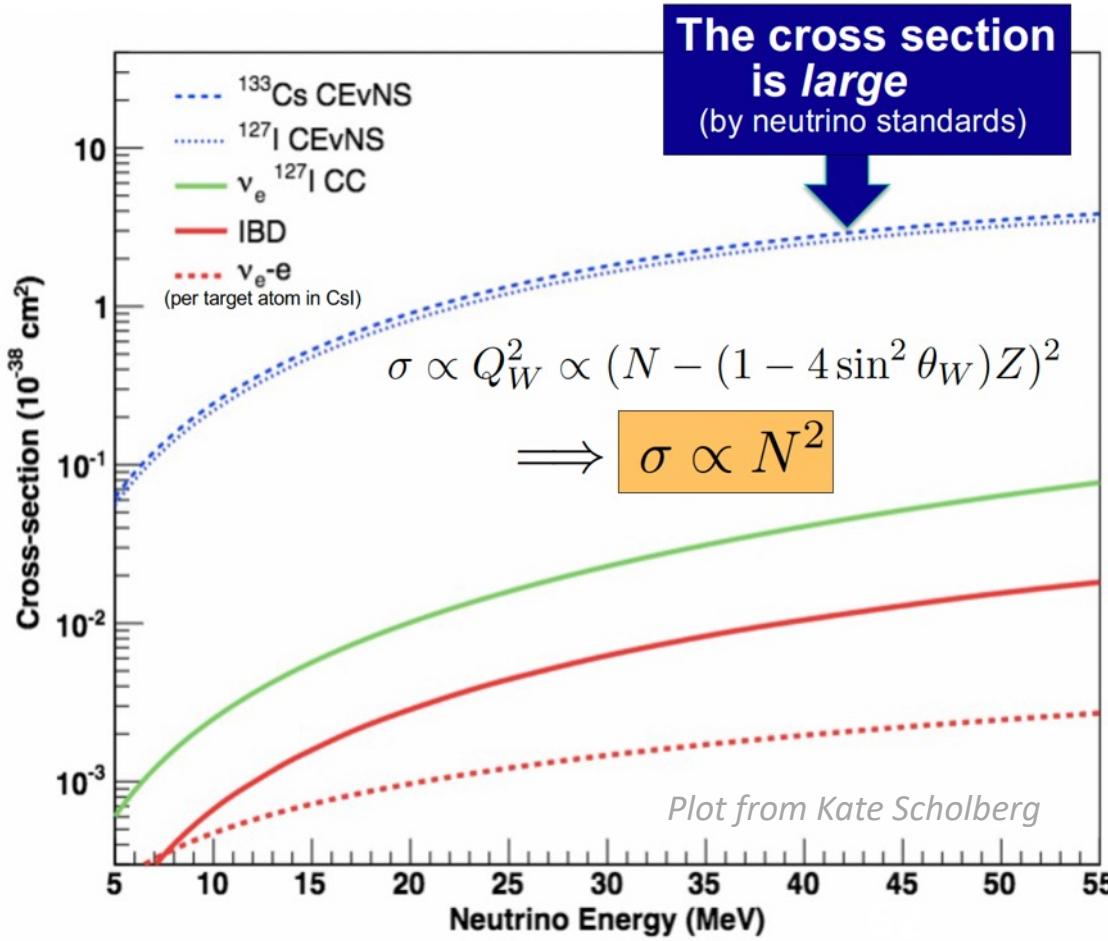
For $QR \ll 1$, [total xscn] $\sim A^2 * [\text{single constituent xscn}]$

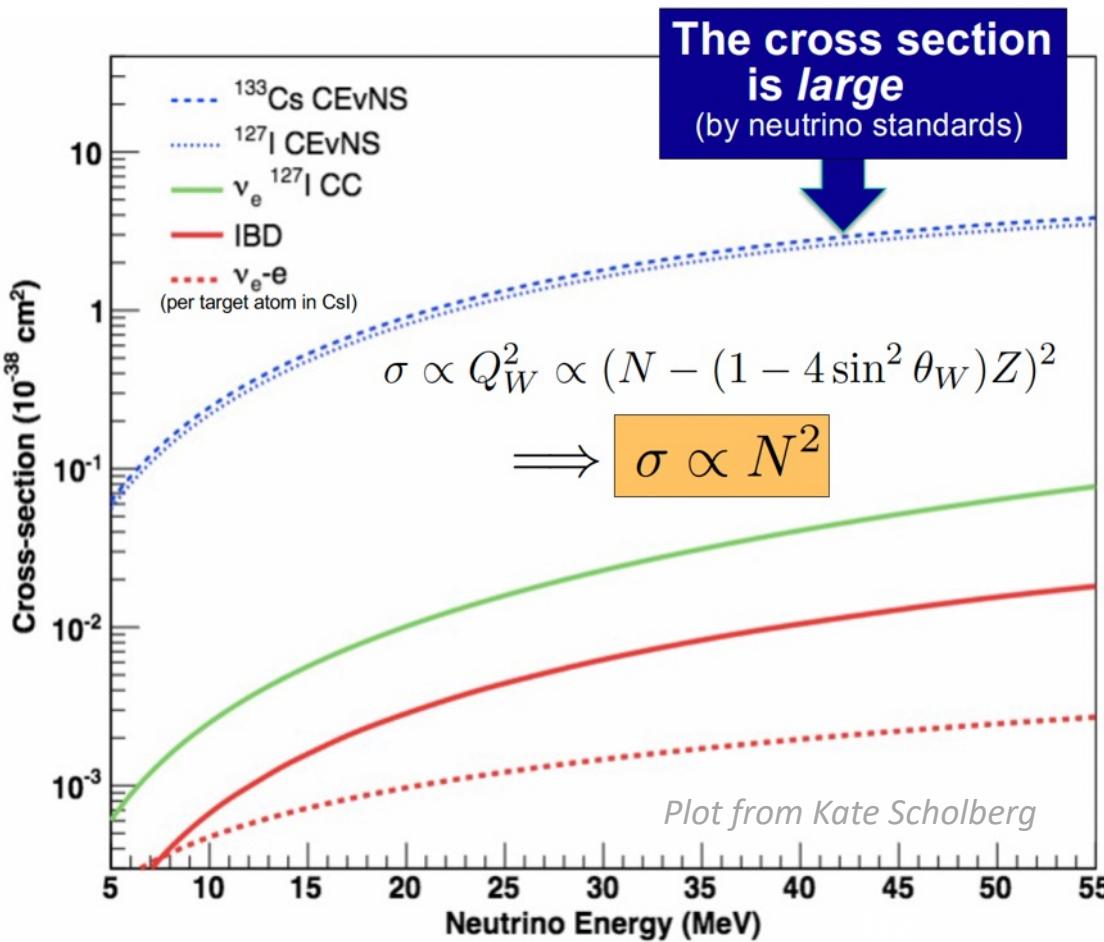
Image: J. Link Science Perspectives

A: no. of constituents

Slide from Kate Scholberg

... but with the highest cross section, proportional to N^2 !





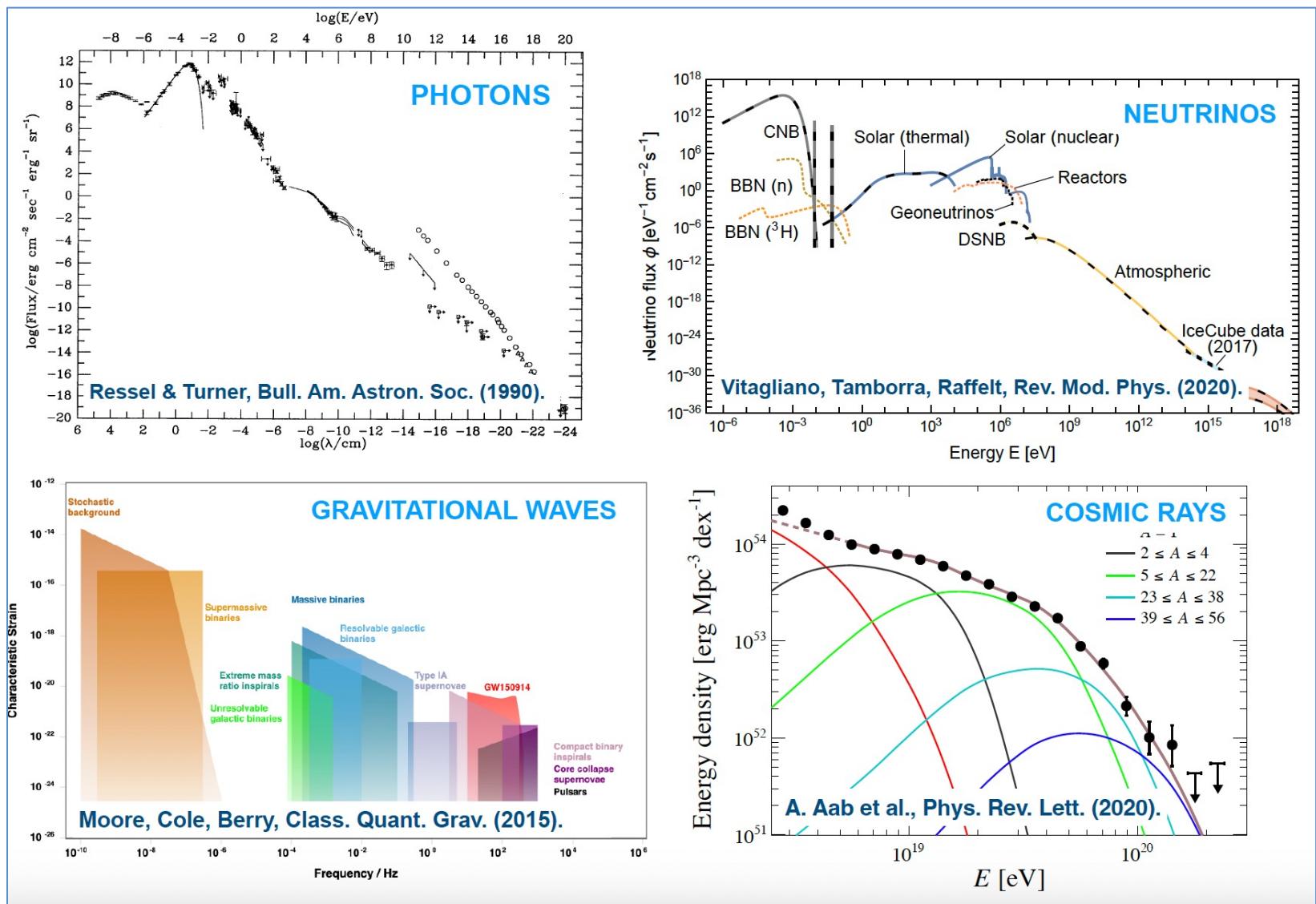
...but recoil energy is O(1-100) keV: small & difficult to measure!

Detection requires:
Low threshold
Low background
High E resolution

RES-NOVA approach:
Archaeological Pb & cryogenic detector

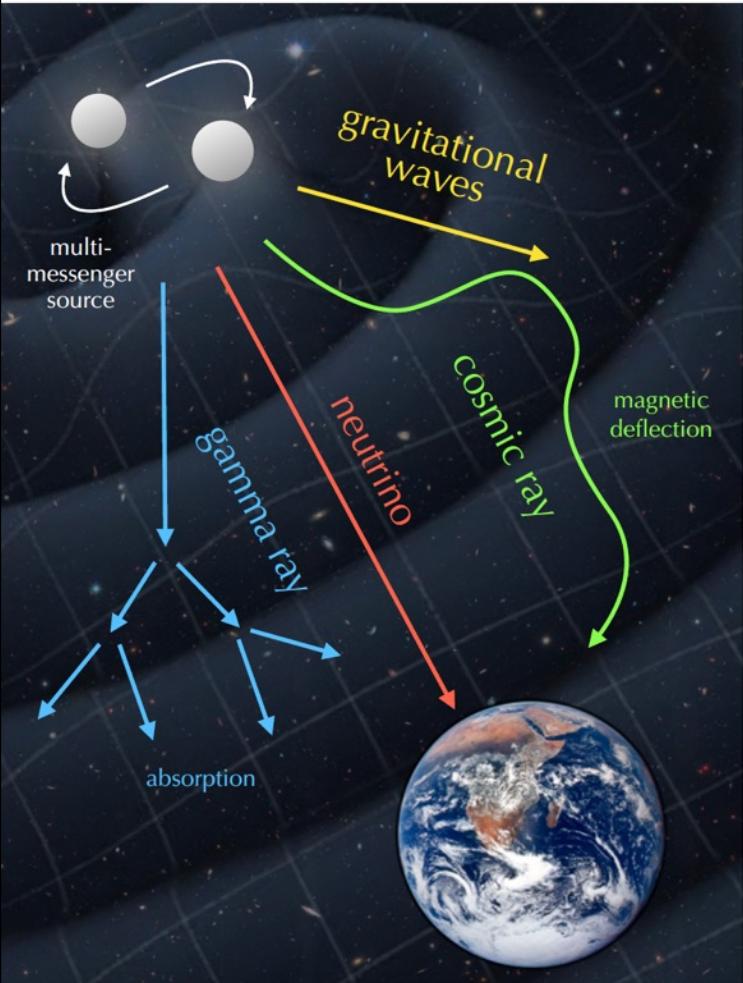


Back to neutrinos as (co)messengers of natural sources...

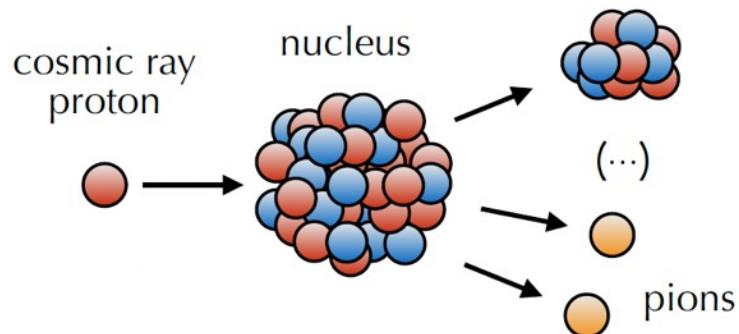


In general, expect (model-dep.) correlation of messengers from a single source...

Multimessenger Astronomy



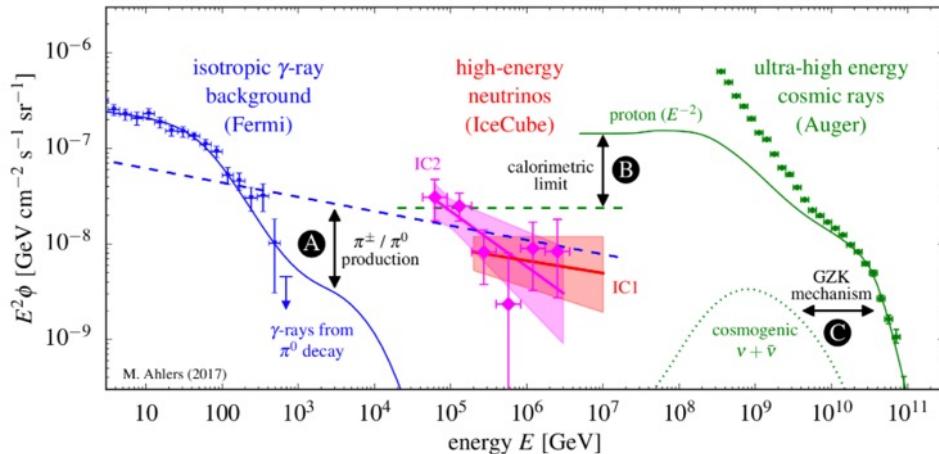
Acceleration of charged nuclei (**cosmic rays**) - especially in the aftermath of cataclysmic events, sometimes visible in **gravitational waves**.



Secondary **neutrinos** and **gamma-rays** from pion decays:

$$\begin{aligned}\pi^+ &\rightarrow \mu^+ + \nu_\mu & \pi^0 &\rightarrow \gamma + \gamma \\ && \downarrow & e^+ + \nu_e + \bar{\nu}_\mu\end{aligned}$$

... and also at a “statistical” level from diffuse sources



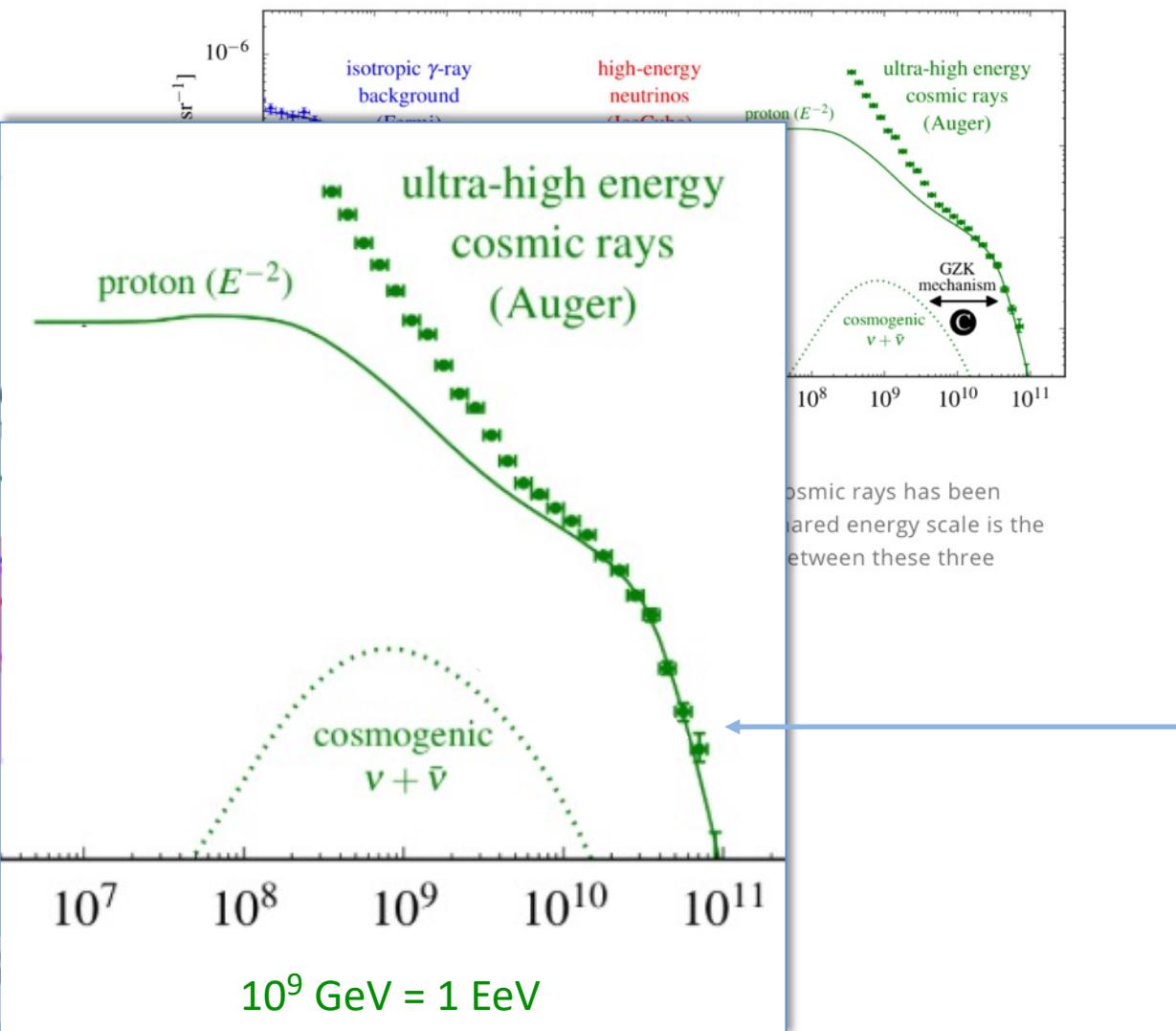
The observed diffuse flux of neutrinos, gamma rays and cosmic rays has been shown to be of comparable magnitude. This apparently shared energy scale is the main motivation behind the search for a common origin between these three messengers.

A: The joined production of charged pions and neutral pions in cosmic-ray interactions leads to the emission of neutrinos (dashed blue) and gamma-rays (solid blue), respectively. What does gamma-ray observations tell us about the neutrino emission?

B: Cosmic ray emission models (solid green) of the most energetic cosmic rays imply a maximal flux (calorimetric limit) of neutrinos from the same sources (green dashed). Is there a deep connection between these unknown CR sources and the recent IceCube observation?

C: The same cosmic ray model predicts the emission of cosmogenic neutrinos from the collision with cosmic background photons (GZK mechanism). What are the requirements for future neutrino observatories to reach this low flux of cosmogenic to test the GZK mechanism?

... and also at a “statistical” level from diffuse sources

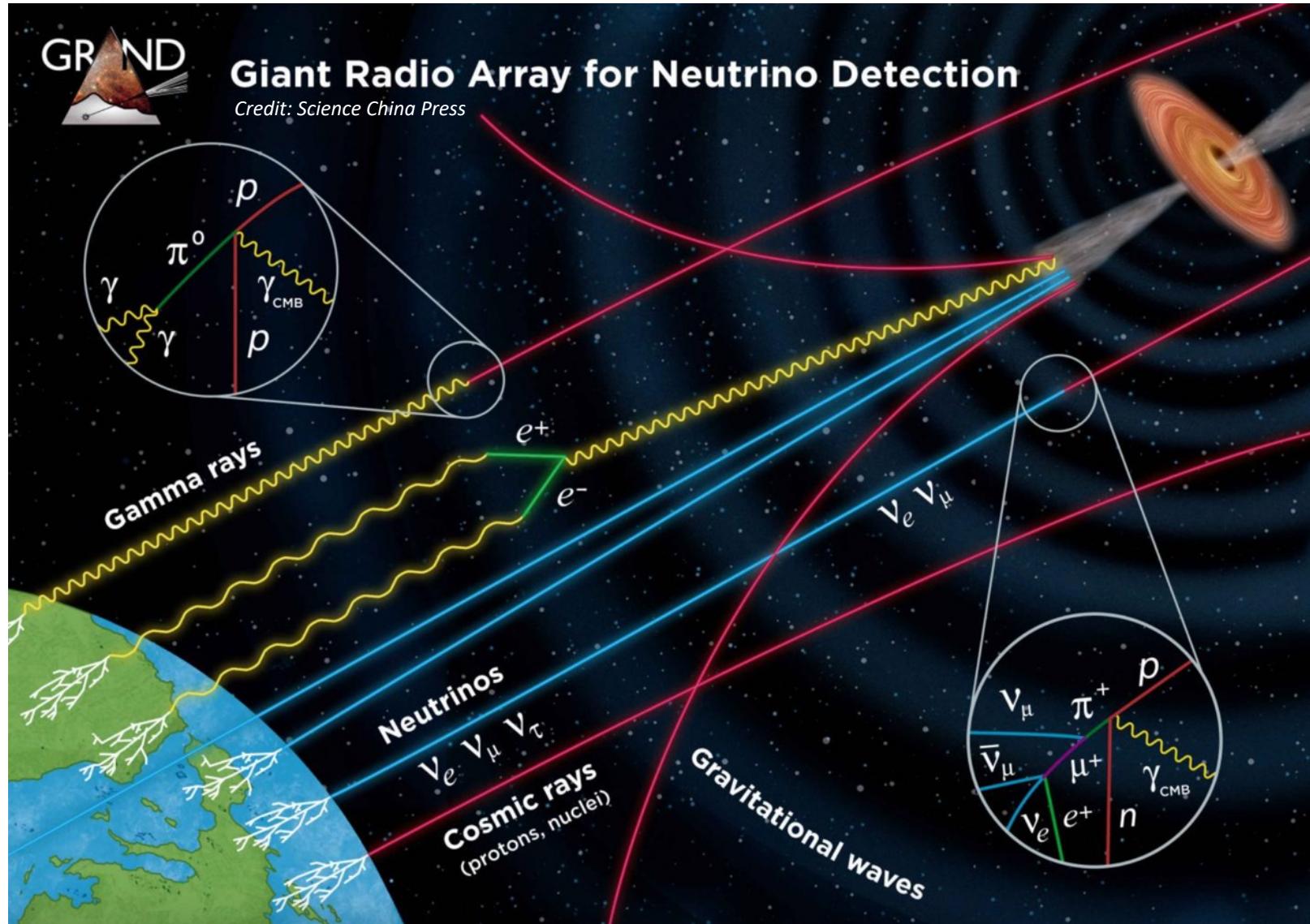


Of interest for this school (AUGER):

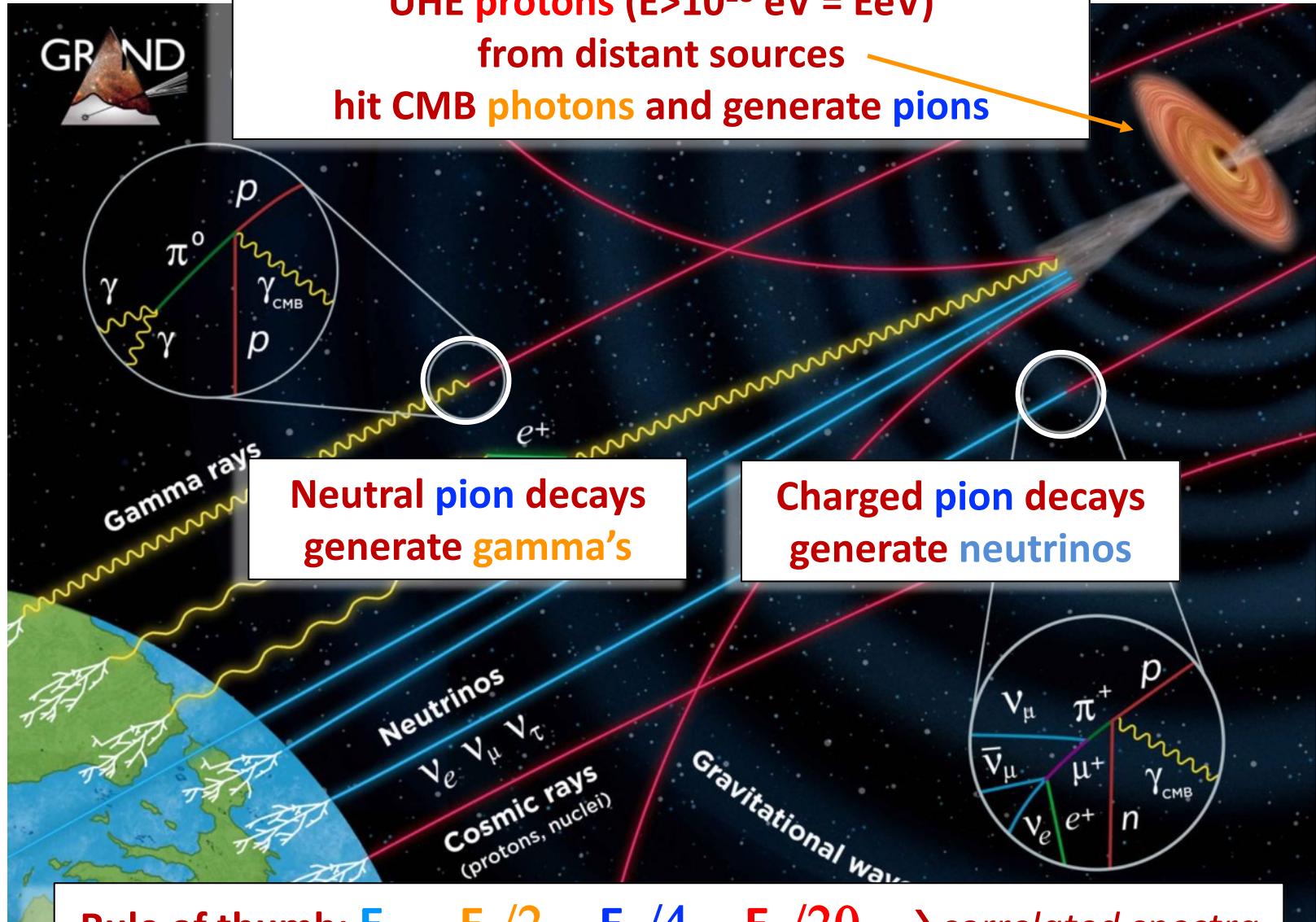
Linking the high-energy tails of CR and ν spectra

(UHECR & cosmogenic ν)

Creation and propagation of ultra-high energy particles in the Universe



(Above: "standard" processes. Possible UHECR from new physics: heavy dark matter, topological defects, etc.)





Complications:

Sources: Require assumptions

UHECR: Not only protons (also nuclei)

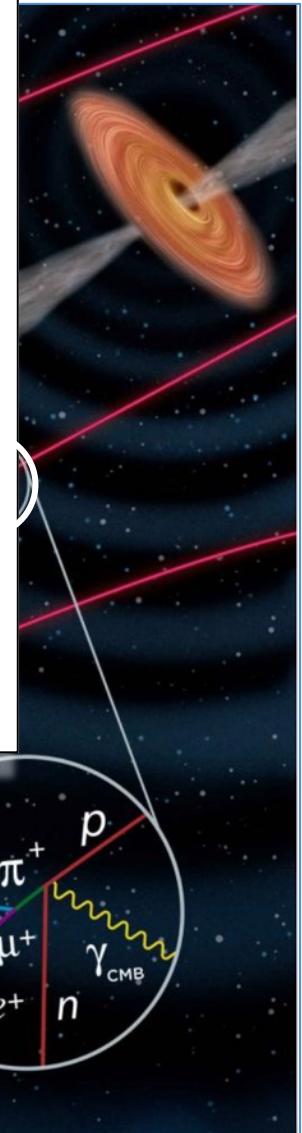
Bkgd photons: Not only CMB

Interactions: Not only pion production

Univ. expansion: Source history, E redshift

...

Need simulations for CR evolution with various source+sink terms & their variants to characterize observable UHE spectra

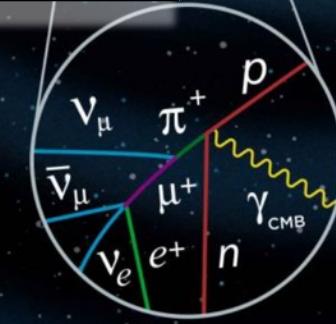


Neutrinos

$\nu_e \bar{\nu}_e \nu_\mu \bar{\nu}_\mu \nu_\tau \bar{\nu}_\tau$

Cosmic rays
(protons, nuclei)

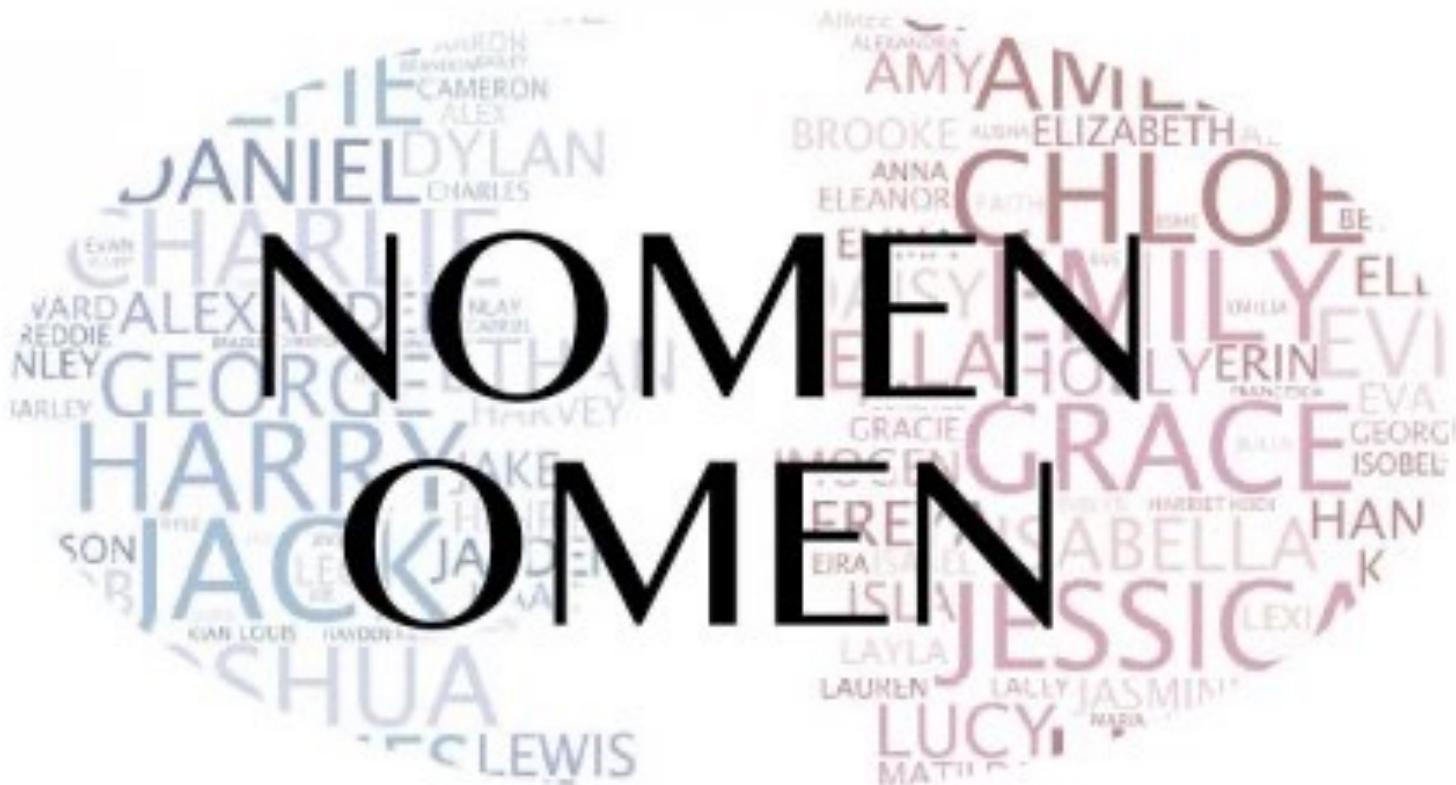
Gravitational waves



Outline:

- Prologue: ν basic properties
- (Un)knowns in neutrino oscillations
- Non-oscillation ν observables
- Learning from neutrinos
- **Epilogue: What's in a name?**

An old latin dictum:



A name, a destiny...
Neutrino name: which destiny?

Language	Word tree	<i>...Some branches</i>	Meaning
Physics (Fermi 1934)	NEUTR-INO		Little neutral one
Italian	NEUTRO		Neutral
Latin	NE-UTER		Not either; neutral
Latin	UTER		Either
Greek		OUDETEROS	<i>Neutral</i>
Old High German		HWEDAR	<i>Which of two; whether</i>
Phonetic change/loss	[K]UOTER[US]		Which of the two?
Ionic Greek	KOTEROS		Which of the two?
Sanskrit	KATARAS		Which of the two?

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Sanskrit	KATARAS		Which of the two?
Latin		QUANTUS	How much?
Sanskrit		KATAMAS	Which out of many?
Sanskrit		KATHA	How?
Sanskrit		KAS	Who?
Indo-European root	KA or KWA		Interrogative base

The root of the name [neutrino] ... is a [kwa]stion

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Sanskrit		KATHA	<i>How?</i>
Sanskrit		KAS	<i>Who?</i>
Indo-European root	KA or KWA		Interrogative base

Nomen Omen...



The destiny
of neutrinos is....
to raise new questions!