# **Modeling of JPA**

S. Pagano (+ collaborators in Salerno) Dip. Fisica Univ. Salerno INFN G.C. Salerno

Aim: modelling of a JPA without simplifications of the nonlinear components

#### Model by Claudio Gatti



The Lagrangian density of the circuit in figure 1 is obtained by the sum of the capacitive energies minus the inductive energies:

$$\mathcal{L} = \left(\frac{1}{2}C_L \dot{\phi_L}^2 - \frac{1}{2L_L} \phi_L'^2\right) \sigma(-x) + \left(\frac{1}{2}C_R \dot{\phi_R}^2 - \frac{1}{2L_R} \phi_R'^2\right) [\sigma(x) - \sigma(x-L)](1) \\ + \left(\frac{1}{2}C_c (\dot{\phi}_L - \dot{\phi}_R)\right) \delta(x) + \left[\frac{C_J}{2} \dot{\phi}_J^2 + E_J (\phi_{ext}) \cos\left(2\pi \frac{\phi_J}{\phi_0}\right)\right] \delta(x-L)$$

where  $\phi_{L,R}(x)$  is the flux variable defined as  $\phi(x) = \int^t dt' V(t')$  in the left (right) transmission line and  $C_{L,R}$  and  $L_{L,R}$  are their capacitances and inductances per unit length, and  $\phi_J = \phi_R(x = L)$ .  $\sigma(x)$  is the Heaviside (step) function. The (symmetric) SQUID Josephson energy is:

$$E_J(\phi_{ext}) = 2E_{0J}\cos\left(\pi\frac{\phi_{ext}}{\phi_0}\right) \tag{2}$$

where  $E_{0J}$  is the Josephson energy of the single junction.

Wave equations for  $\phi_L$  and  $\phi_R$ 

$$\sigma(-x) \left[ C_L \ddot{\phi}_L - \frac{1}{L_L} \phi_L'' \right] + \delta(x) \left[ C_c (\ddot{\phi}_L - \ddot{\phi}_R) + \frac{1}{L_L} \phi_L' \right] = 0(4)$$
$$[\sigma(x) - \sigma(x - L)] \left[ C_R \ddot{\phi}_R - \frac{1}{L_R} \phi_R'' \right] + \delta(x) \left[ -C_c (\ddot{\phi}_L - \ddot{\phi}_R) - \frac{1}{L_R} \phi_R' \right] + \delta(x - L) \left[ C_J \ddot{\phi}_R + \frac{1}{L_R} \phi_R' + \frac{2\pi}{\phi_0} E_J(\phi_{ext}) \sin\left(2\pi \frac{\phi_R}{\phi_0}\right) \right] = 0$$

#### The solution is the sum of an outgoing and an incoming wave

 $\phi(x,t) = \phi^{in}(t - x/v_p) + \phi^{out}(t + x/v_p)$  and  $v_p = 1/\sqrt{LC}$ . The other equations give the boundary conditions:

$$-\frac{1}{L_L}\phi'_L = C_c(\ddot{\phi}_L - \ddot{\phi}_R) \qquad (x=0)$$

•

$$\phi' = \frac{1}{v_p} \left( \dot{\phi}^{out} - \dot{\phi}^{in} \right)_{in \text{ and } out \text{ fields.}} - \frac{1}{L_R} \phi'_R = C_c (\ddot{\phi}_L - \ddot{\phi}_R) \qquad (x = 0)$$
  
$$\dot{\phi} = \dot{\phi}^{in} + \dot{\phi}^{out} \qquad C_J \ddot{\phi}_R + \frac{1}{L_R} \phi'_R + \frac{2\pi}{\phi_0} E_J(\phi_{ext}) \sin\left(2\pi \frac{\phi_R}{\phi_0}\right) = 0 \qquad (x = L)$$

$$2\ddot{\phi}_{R}^{in} + \frac{1}{CcZ_{0}}\dot{\phi}_{R}^{in} = \ddot{\phi}_{L}^{in} + \frac{1}{CcZ_{0}}\dot{\phi}_{R}^{out} \quad (x=0)$$

$$C_{J}\ddot{\phi}_{J} + \frac{1}{Z_{0}}\dot{\phi}_{J} + \frac{2\pi}{\phi_{0}}E_{J}(\phi_{ext})\sin\left(2\pi\frac{\phi_{J}}{\phi_{0}}\right) = \frac{2}{Z_{0}}\dot{\phi}_{R}^{in} \quad (x=L) \quad \phi_{R}^{in}(t,x=L) = \phi_{R}^{in}(t-L/v_{p},x=0)$$

$$\phi_R^{out}(t, x = L) = \phi_R^{out}(t + L/v_p, x = 0) = \phi_J(t) - \phi_R^{in}(t, x = L)$$
  
=  $\phi_J(t) - \phi_R^{in}(t - L/v_p, x = 0)$ 

 $\phi_L^{out} = \phi_R^{out} - \phi_R^{in} + \phi_L^{in} \qquad (x=0)$ 

#### Transient Analysis of Lossless Transmission Lines

Manuscript received July 7, 1967; revised September 8, 1967.







The method of characteristics for solving these equations is based on a transformation in the x-t plane which accomplishes the conversion of (1) and (2) into a pair of ordinary differential equations. Each of these two ordinary differential equations holds true along a different family of *characteristic curves* in the x-t plane, one family corresponding to the forward or incident wave and the other to the backward or reflected wave. These two families of characteristic curves are, in effect, the basis for a new coordinate system instead of the canonical coordinates x = constantand t = constant.

$$dx/dt = 1/\sqrt{LC}$$
 and  $dx/dt = -1/\sqrt{LC}$ 

$$\sqrt{\frac{L}{C}}di + \left(Ri + \sqrt{\frac{L}{C}}Ge\right)dx + de = 0$$
(3)

which holds true only along the *forward* characteristics, defined as  $dx/dt = 1/\sqrt{L/C}$ , and

$$-\sqrt{\frac{L}{C}}di + \left(Ri - \sqrt{\frac{L}{C}}Ge\right)dx + de = 0$$
<sup>(4)</sup>

which holds true only along the *backward* characteristics, defined as  $dx/dt = -1/\sqrt{LC}$ .

#### Transient Analysis of Lossless Transmission Lines

Manuscript received July 7, 1967; revised September 8, 1967.



F. H. BRANIN, JR. Systems Development Div. Lab. IBM Corporation Kingston, N. Y.

$$v = v(x,t)$$
  $i = i(x,t)$   
 $\tau = \sqrt{LC} l$  line delay

 $Z_0 = \sqrt{\frac{L}{c}}$  line impedance

$$\begin{aligned} v(l,t) &= -Z_0 i(l,t) + [v(0,t-\tau) + Z_0 i(0,t-\tau)] \\ v(0,t) &= +Z_0 i(0,t) + [v(l,t-\tau) - Z_0 i(l,t-\tau)] \\ v_2(t) &= -Z_0 i_2(t) + [v_1(t-\tau) + Z_0 i_1(t-\tau)] \\ v_1(t) &= +Z_0 i_1(t) + [v_2(t-\tau) - Z_0 i_2(t-\tau)] \end{aligned}$$

$$v_2(t) = -Z_0 i_2(t) - E_i(t - \tau)$$
  
$$v_1(t) = +Z_0 i_1(t) - E_r(t - \tau)$$

$$E_i(t) = - [2v_1(t) + E_r(t - \tau)]$$
  

$$E_r(t) = -[2v_2(t) + E_i(t - \tau)]$$



$$\begin{split} \dot{v}_1 &= -\frac{1}{2 C_c Z_0} v_1 + \frac{1}{2} \dot{v}_a(t) - \frac{1}{2} \dot{E}_r(t-\tau) - \frac{1}{2 C_c Z_0} E_r(t-\tau) \\ \dot{v}_2 &= -\left(\frac{1}{C_j Z_0} + \frac{1}{C_j R_j}\right) v_2 - \frac{I_j}{C_j} \cos\left(\pi \frac{\phi_{ext}}{\phi_0}\right) \sin \varphi - \frac{1}{C_j Z_0} E_i(t-\tau) \\ \dot{\phi} &= \frac{2 e}{\hbar} v_2 \\ \text{where} \\ v_s(t) &= v_{rf} \sin \omega t \\ \dot{v}_s(t) &= v_{rf} \omega \cos \omega t \\ \phi_{ext}(t) &= \phi_{dc} + \phi_{ac} \cos \omega_p t \\ Z_0 &= 50 \Omega; \quad \tau = \frac{T}{4} = \frac{\pi}{2 \omega_{ris}} ; \\ I_j &= 2I_{j0}, C_j &= 2C_{j0} \text{ and } R_j = \frac{1}{2}R_{j0} \text{ are the SQUID critical current capacitance and resistance} \\ \text{We start with } Ei \text{ and } Er \text{ set to zero} \\ \text{The values of } Ei \text{ and } Er \text{ are updated during the integration according to:} \\ E_i(t) &= -[2v_1(t) + E_r(t-\tau)] \\ E_r(t) &= -[2v_2(t) + E_i(t-\tau)] \end{split}$$

For Vout we have to decide how to simulate the circulator: Assuming that the reflected wave do not see the signal source, its component at the signal frequency will see an RC circuit  $\omega_s Z_0 C_c$ 



Alternatively, we can take the voltage at that the input of the coupling capacitor:



## Simulation results

As parameter values the experimental data from LNF are used:

Resonator Length 2.62mm Substrate silicon eps\_eff=6.346 Bare resonator frequency **11.3 GHz** Impedance **Z0=50 Ohm** Coupling capacitor **14-18 fF** S=12 μm e w=7 μm

SQUID Parameters: Junction area 0.3-0.6 μm2 Junction **Ic=0.47 μA** (quella misurata) Junction **Cj<0.5 pF** Junction freq fjj>> 10 GHz DC SQUID Area 4x4 μm2 M=70 fH (misurata)

The effects of the SQUID load depend strongly on the critical current



Scan of pump frequency with fpump = 2 fsignal

The effects of the SQUID load depend strongly on the critical current

A resonance occurs for fsignal near 8 Ghz

fpump=2 fsignal

Zi = 50 Ohm Cc = 0.015 pF Z0 = 50 Ohm **I0 = 2.0 uA** Cj = 0.5 pF Rj = 100000 Ohm Fplasma = 110.24538380186623 GHz

Lambda/4 Resonator frequency = 11.3 GHz Time delay = 22.12 ps

input RF signal = 0.1 uV = -130.0 dBm

Fstart = 7.0 GHz Fend = 9.0 GHz Fstep = 100.0 MHz

Magnetic flux amplitude **DC = 0.29** Phi0 **RF = 0.15** Phi0 Frequency = 10.2 GHz

Integration time = 20.0 ns time step = 0.1 ps integrating 200000 time steps with 221 lag steps

increase of dc flux bias





Effect of dc flux bias

Ic=2uA

fpump=2 fsignal

Zi = 50 Ohm Cc = 0.015 pF Z0 = 50 Ohm **I0 = 2.0 uA** Cj = 0.5 pF Rj = 100000 Ohm

Fplasma = 110.24 GHz

Lambda/4 Resonator frequency = 11.3 GHz Time delay = 22.12 ps

input RF signal = 0.1 uV = -130.0 dBm

Fstart = 7.0 GHz Fend = 9.0 GHz Fstep = 100.0 MHz

Magnetic flux amplitude **DC = 0.31** Phi0 RF = 0.15 Phi0 Frequency = 10.2 GHz



Now we reduce the SQUID critical current: Ic=1uA

It is necessary more pump power

The resonance moves to lower frequency

fpump=2 fsignal

Zi = 50 Ohm Cc = 0.015 pF Z0 = 50 Ohm **I0 = 1.0 uA** Cj = 0.5 pF Rj = 100000 Ohm Fplasma = 77.95 GHz

Lambda/4 Resonator frequency = 11.3 GHz Time delay = 22.12 ps

input RF signal = 0.1 uV = -130.0 dBm

Fstart = 5.0 GHz Fend = 7.0 GHz Fstep = 100.0 MHz

Magnetic flux amplitude **DC = 0.33 Phi0 RF = 0.25 Phi0** Frequency = 10.2 GHz



Reducing further the SQUID critical current: Ic=0.5uA

More pump power is necessary

The resonance frequency is further reduced fpump=2 fsignal

Zi = 50 Ohm Cc = 0.015 pF Z0 = 50 Ohm I0 = 0.5 uA Cj = 0.5 pF Rj = 100000 Ohm

Fplasma = 55.12 GHz

Lambda/4 Resonator frequency = 11.3 GHz Time delay = 22.12 ps

input RF signal = 0.1 uV = -130.0 dBm

Fstart = 2.5 GHz Fend = 4.5 GHz Fstep = 100.0 MHz

Magnetic flux **amplitude DC = 0.33 Phi0 RF = 0.3 Phi0** Frequency = 10.2 GHz



#### Back to Ic=1uA

Fix the pump frequency to 11.8 GHz and sweep the signal frequency

After 20 ns the system is not stationary

Zi = 50 Ohm Cc = 0.015 pF Z0 = 50 Ohm I0 = 1.0 uA Cj = 0.5 pF Rj = 100000 Ohm Fplasma = 77.95 GHz

Lambda/4 Resonator frequency = 11.3 GHz Time delay = 22.12 ps input RF signal = 0.1 uV = -130.0 dBm

Fstart = 5.0 GHz Fend = 7.0 GHz Fstep = 100.0 MHz

Magnetic flux amplitude DC = 0.33 Phi0 RF = 0.25 Phi0 Frequency = 11.8 GHz



#### Increase to 50 ns

Zi = 50 Ohm Cc = 0.015 pF Z0 = 50 Ohm I0 = 1.0 uA Cj = 0.5 pF Rj = 100000 Ohm

Fplasma = 77.95 GHz

Lambda/4 Resonator frequency = 11.3 GHz Time delay = 22.12 ps input RF signal = 0.1 uV = -130.0 dBm

Fstart = 5.0 GHz Fend = 7.0 GHz Fstep = 100.0 MHz

Magnetic flux amplitude DC = 0.33 Phi0 RF = 0.25 Phi0 Frequency = 11.8 GHz



Lowering the signal does not change the result

The system is self oscillating, triggered by the signal

Zi = 50 Ohm Cc = 0.015 pF Z0 = 50 OhmI0 = 1.0 uA Cj = 0.5 pF Rj = 100000 Ohm

Fplasma = 77.95 GHz Lambda/4 Resonator frequency = 11.3 GHz Time delay = 22.12 ps input RF signal = 0.01 uV = -150.0 dBm

Fstart = 5.0 GHz Fend = 7.0 GHz Fstep = 100.0 MHz

Magnetic flux amplitude DC = 0.33 Phi0 RF = 0.25 Phi0 Frequency = 11.8 GHz



### Scan of different dc and rf amplitudes of magnetic flux



Effect of increasing the dc magnetic flux

$$Fdc = 0.30$$

Zi = 50 Ohm Cc = 0.015 pF Z0 = 50 Ohm I0 = 1.0 uA Cj = 0.5 pF Rj = 100000 Ohm Fplasma = 77.95525848081317 GHz Lambda/4 Resonator frequency = 11.3 GHz Time delay = 22.123893805309738 ps input RF signal = 0.01 uV = -150.0 dBm Fstart = 2.0 GHz Fend = 10.0 GHz Fstep = 100.0 MHz **Magnetic flux amplitude DC = 0.30 Phi0 RF = 0.25 Phi0 Frequency = 11.3 GHz** Integration time = 50.0 ns time step = 0.1 ps

integrating 500000 time steps with 221 lag steps



Effect of increasing the dc magnetic flux



Effect of increasing the dc magnetic flux

Zi = 50 Ohm Cc = 0.015 pF Z0 = 50 Ohm I0 = 1.0 uA Cj = 0.5 pF Rj = 100000 Ohm

Lambda/4 Resonator frequency = 11.3 GHz Time delay = 22.123893805309738 ps

input RF signal = 0.01 uV = -150.0 dBm

Fstart = 2.0 GHz Fend = 10.0 GHz Fstep = 100.0 MHz

Magnetic flux amplitude DC = 0.32 Phi0 RF = 0.25 Phi0 Frequency = 11.3 GHz

Integration time = 50.0 ns time step = 0.1 ps

integrating 500000 time steps with 221 lag steps



# Scan of signal frequency for fixed pump



# Detail of dynamics



### Conclusions (for now)

1) the numerical model seems to work

2) Extension to a real SQUID is in progress

3) The value of experimental parameters are important (Cc and Io)

4) The JPA simulation show expected responses but not amplifications

5) Self oscillations are triggered by input signal

6) More tests are necessary to fully validate the model

7) The role of circulator needs to be clarified