LIGHT DARK MATTER DETECTION

Angelo Esposito



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CYGNO Collaboration meeting, December 2022



• Searching for keV-MeV dark matter



- Searching for keV-MeV dark matter
- Down to the MeV: Migdal effect



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- <u>Down to the MeV</u>: Migdal effect
- <u>Down to the keV</u>: collective excitations



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- Outlook





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- However... huge possible mass range
 →
 detection techniques
 vary widely depending on the dark matter mass

































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- Need new materials or observables





• For an elastic scattering, it must be

$$E_T = \frac{m_{\chi}/m_T}{\left(1 + m_{\chi}/m_T\right)^2} E_{\chi} \gtrsim E_{threshold}$$



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- For $m_{\chi} \lesssim 1$ GeV elastic scattering off nuclei is very inefficient
- Two possibilities:
 - I. Look into lighter scattering targets
 - 2. Look into inelastic processes









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Need to account for the complicated many-body physics (correlations, strong coupling, ...)







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phys.

 (0^{3})

puys.

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phys.

• Need to find theoretical tools that allow to solve or bypass these problems (measured correlation functions, EFTs, ...)



Down to the MeV Migdal effect in semiconductors









For sub-GeV dark matter nuclear recoil signals become challenging
 sensitivity can be lowered by looking for inelastic processes
 [e.g., Essig, Mardon, Volansky - PRD 2012, 1108.5383; Kouvaris, Pradler - PRL 2017, 1607.01789]



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Less likely... but lower threshold! -> sensitivity down to • O(100 MeV) masses (see Giovanni's talk) [e.g., Ibe, Nakano, Shoji, Suzuki - JHEP 2018, 1707.07258; DarkSide - 2207.11967]





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 How to describe nucleus-nucleus and nucleus-electron interactions in a strongly correlated system?


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3. Harmonic approximation \rightarrow valid only for $m_{\chi} \simeq 1 - 10$ MeV

[Mo, Zheng, Zhang - PRD 2022, 2205.03395]



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 $q \,[\text{keV}]$ 0.550250 $\sqrt{2m_{\rm N}\omega_{\rm g}}$ EFT $(\Delta E_{\lambda} \ll \omega)$ 1/a $\sqrt{2m_{\rm N}\langle E_{\rm ph}\rangle}$ free ion impulse incoherent harmonic $m_{\chi} \; [\text{MeV}]$ 0.550250

10/24

[Knapen, Kozaczuk, Lin - PRL 2021, 2011.09496]



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$$H_{eff} = \frac{1}{m_N \omega^2} \overrightarrow{\nabla} H_{\chi L} \cdot \overrightarrow{\nabla} H_{eL} + \mathcal{O}\left(\frac{1}{\omega^3}\right) \quad s$$

[Berghaus, **AE**, Essig, Sholapurkar — JHEP 2020, 2210.06490]



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Now simple to determine the rate for Migdal emission

$$\frac{d^{2}\Gamma}{d\omega dE_{ph}} \propto \sum_{\mathbf{k}} \sum_{\mathbf{K},\mathbf{Q}} \frac{\mathbf{q} \cdot (\mathbf{k} + \mathbf{K}) \mathbf{q} \cdot (\mathbf{k} + \mathbf{K})}{|\mathbf{k} + \mathbf{K}| |\mathbf{k} + \mathbf{Q}|} \operatorname{Im} \left(-\epsilon_{\mathbf{K}\mathbf{Q}}^{-1}(\mathbf{k},\omega)\right) S(\mathbf{q} - \mathbf{k} - \mathbf{K}, E_{ph})$$



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 Energy loss function is already well studied

[e.g., Knapen, Kozaczuk, Lin - PRD 2021, 2101.08275; Hochberg et al. - PRL 2021, 2101.08263; Knapen, Kozaczuk, Lin - PRD 2022, 2104.12786]



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• Structure factor should be measured from neutron scattering data

$$\frac{d^2\sigma_n}{d\Omega dE} = \frac{\sigma_n}{4\pi} \frac{k_f}{k_i} S(q, E)$$





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• No data yet in the range of interest ($q \simeq 10 \text{ keV} - 100 \text{ keV}$)



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 If only interested in electron energy, the rate is independent on the details of the crystal lattice

$$\int_{0}^{\infty} dE S(q, E) = 1 \implies \frac{dR}{d\omega} \propto \int d^{3}\mathbf{k}_{e} \sum_{\mathbf{K}} \operatorname{Im}\left(-\epsilon_{\mathbf{KK}}^{-1}(\mathbf{k}_{e}, \omega)\right)$$



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[Berghaus, AE, Essig, Sholapurkar - JHEP 2020, 2210.06490]

• The description of Migdal effect in semiconductor is now complete





Down to the keV collective excitations









• For $m_{\gamma} \leq \mathcal{O}(\text{MeV})$, dark matter scattering can transfer a momentum

 $(m_{\chi}v_{\chi})^{-1} \gtrsim \mathcal{O}(1\text{ Å}) \sim \text{inter-atomic distance}$





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 Typically, no more single particle final states
 signatures

 involve collective excitation
 [see e.g., Trickle et al. - JHEP 2020, 1910.08092; Griffin et al. - PRD 2020, 1910.10716; Coskuner et al. - PRD 2022, 2102.09567]

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 - Superconductors → breaking of Cooper pairs and production of collective excitations

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2. Solid crystals (GaAs, SiO, ...) → multi-phonon

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COLLECTIVE EXCITATIONS

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- 5. and more... [for a review, Kahn, Lin Rept.Prog.Phys. 2022, 2108.03239]
- I will focus on those I know best...



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 Superfluid ⁴He is an interesting target to probe dark matter with spin-independent interactions



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 - I. $E_{ion} \simeq 25 \text{ eV} \rightarrow \text{low electronic background}$
 - 2. High radiopurity
 - 3. Multi-phonon processes allow to probe down to $m_{\gamma} \sim O(\text{keV})$
- Idea: look for events where the dark matter produces more than one phonon

[Guo, McKinsey - PRD 2013, 1302.0534; Schutz, Zurek - PRL 2016, 1604.08206; Knapen, Lin, Zurek - PRD 2017, 1611.06228; Acanfora, AE, Polosa - EPJC 2019, 1902.02361; Caputo, AE, Polosa - PRD 2019, 1907.10635; Baym et al. - PRD 2021, 2005.08824; Caputo, AE, Piccini, Polosa, Rossi - PRD 2021, 2012.01432; Matchev et al. - JHEP 2022, 2108.07275; You et al. - 2208.14474]







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I. HeRALD

[Guo, McKinsey - PRD 2013, 1302.0534; Hertel et al. -PRD 2019, 1810.06283; TESSERACT - SnowMass]







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3. DELight

[Krosigk et al. - 2209.10950]







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 —> phonon emission
 rate is hard



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- Standard approach
 → start from
 atomic Hamiltonian
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 → bypass the atomic Hamiltonian



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$$\mathscr{L}_{EFT} \sim \dot{\pi}^2 - c_s^2 (\nabla \pi)^2 + \lambda_1 \dot{\pi}^3 + \lambda_2 \dot{\pi} (\nabla \pi)^2 + \dots$$



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Observables are computed in the standard particle physics way



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• How about dark matter with spin-dependent interactions?



- How about dark matter with spin-dependent interactions?
- One could use (anti-)ferromagnets





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 Ways to detect few magnons have been proposed and under work (TES, MKID, quantum sensors) [Trickle, Zhang, Zurek - PRL 2020, 1905.13744; Lachance-Quirion et al. - Science Advances 2017; Lachance-Quirion et al. - Science 2020]

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• First proposed to use ferromagnets

[Trickle, Zhang, Zurek - PRL 2020, 1905.13744; Mitridate et al. - PRD 2020, 2005.10256; Chigus, Moroi, Nakayama - PRD 2020, 2001.10666; Trickle, Zhang, Zurek - PRD 2022, 2009.13534]



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Traditional approach: start from Heisenberg model

$$H = \frac{1}{2} \sum_{\ell,\ell'}^{N} \sum_{j,j'}^{n} J_{\ell\ell'jj'} \mathbf{S}_{\ell j} \cdot \mathbf{S}_{\ell'j'} \to \sum_{\nu=1}^{n} \sum_{\mathbf{q}\in 1BZ} \omega_{\nu,\mathbf{q}} b_{\nu,\mathbf{q}}^{\dagger} b_{\nu,\mathbf{q}}$$



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•

Only single-magnon emission \rightarrow probes $m_{\gamma} \gtrsim \mathcal{O}(10 \text{ keV})$



[Trickle, Zhang, Zurek - PRL 2020, 1905.13744]



22/24



- A better class of materials turns out to be anti-ferromagnets
 - I. Also multi-magnon emission \rightarrow probe down to $m_{\gamma} \sim \mathcal{O}(1 \text{ keV})$
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- Gapless magnon = Goldstone \rightarrow bypass several problems $\mathscr{L}_{EFT} = \frac{c_1}{2}\dot{\mathbf{n}}^2 - \frac{c_2}{2}(\partial_i \mathbf{n})^2$


ANTI-FERROMAGNETS

- A better class of materials turns out to be anti-ferromagnets
 - I. Also multi-magnon emission \rightarrow probe down to $m_{\chi} \sim O(1 \text{ keV})$
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[**AE**, Pavaskar - 2210.13516]



CYGNO, December 2022

Angelo Esposito



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Thanks for your attention!

