

MEV-GEV DIRECT DETECTION OVERVIEW



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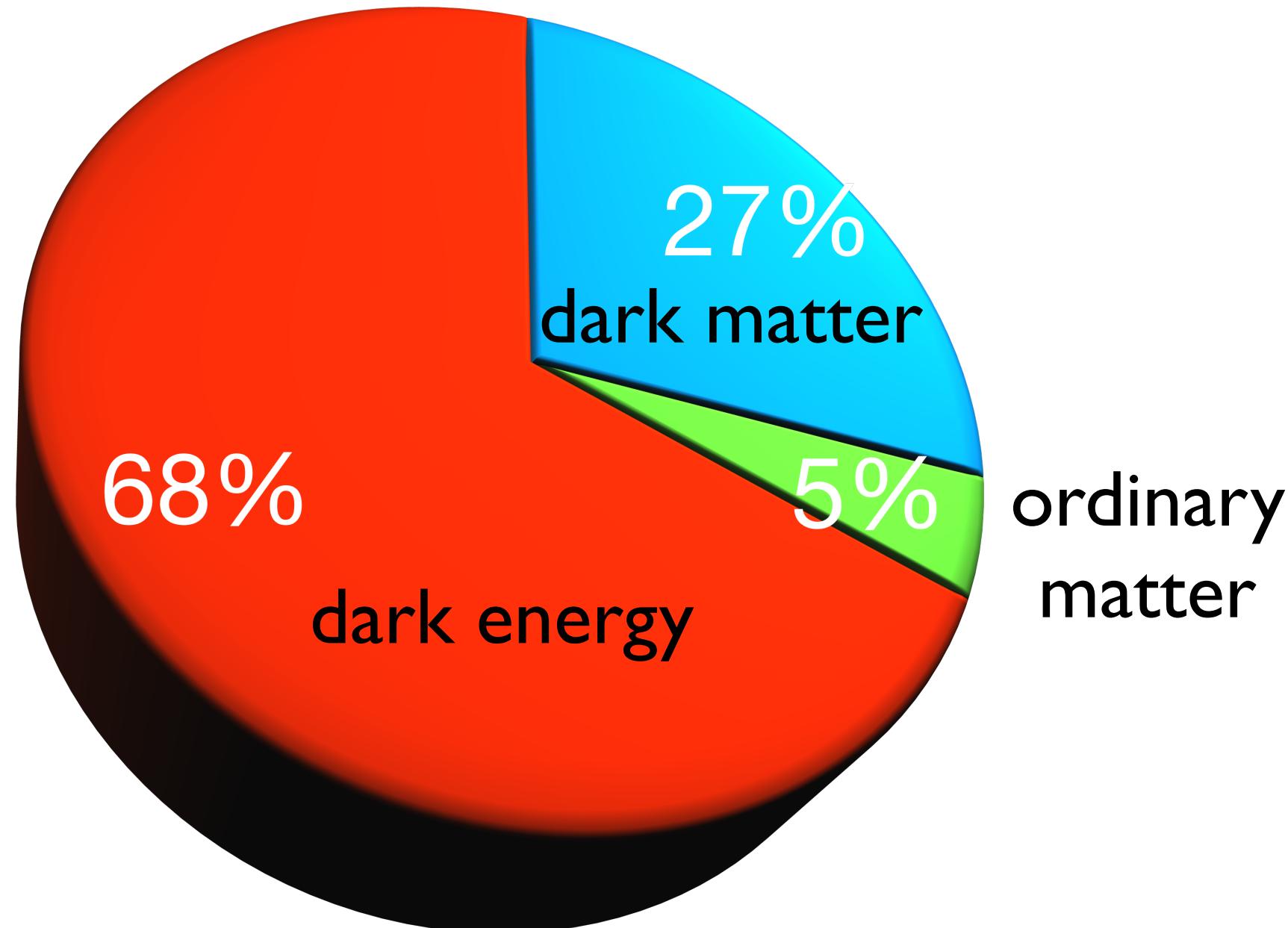
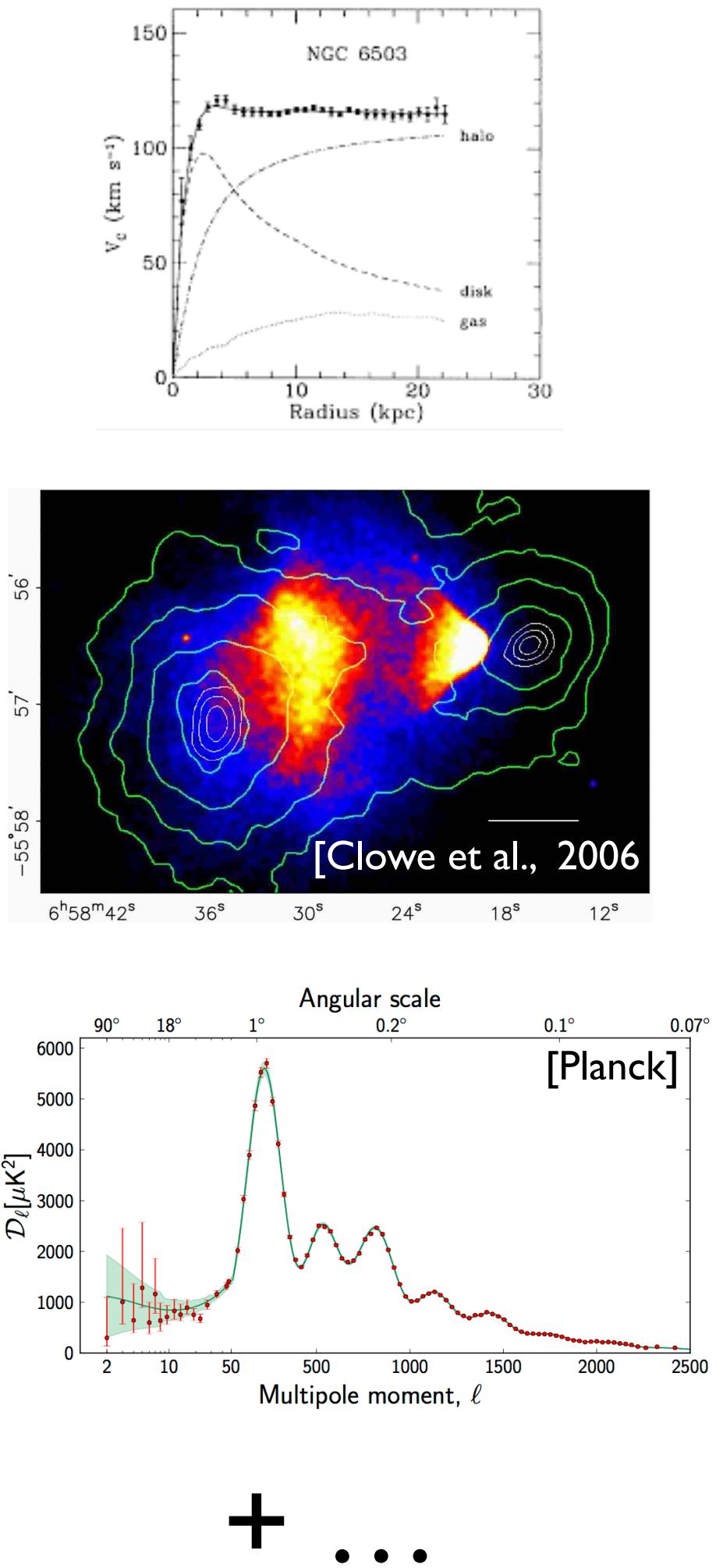
Outline

- **Introduction**
 - DM hints
 - (Theory) DM challenges
- **Direct detection phenomenology**
 - Migdal effect
 - Cosmic ray boosted dark matter
 - Blazar boosted dark matter
- **Conclusions**

Introduction

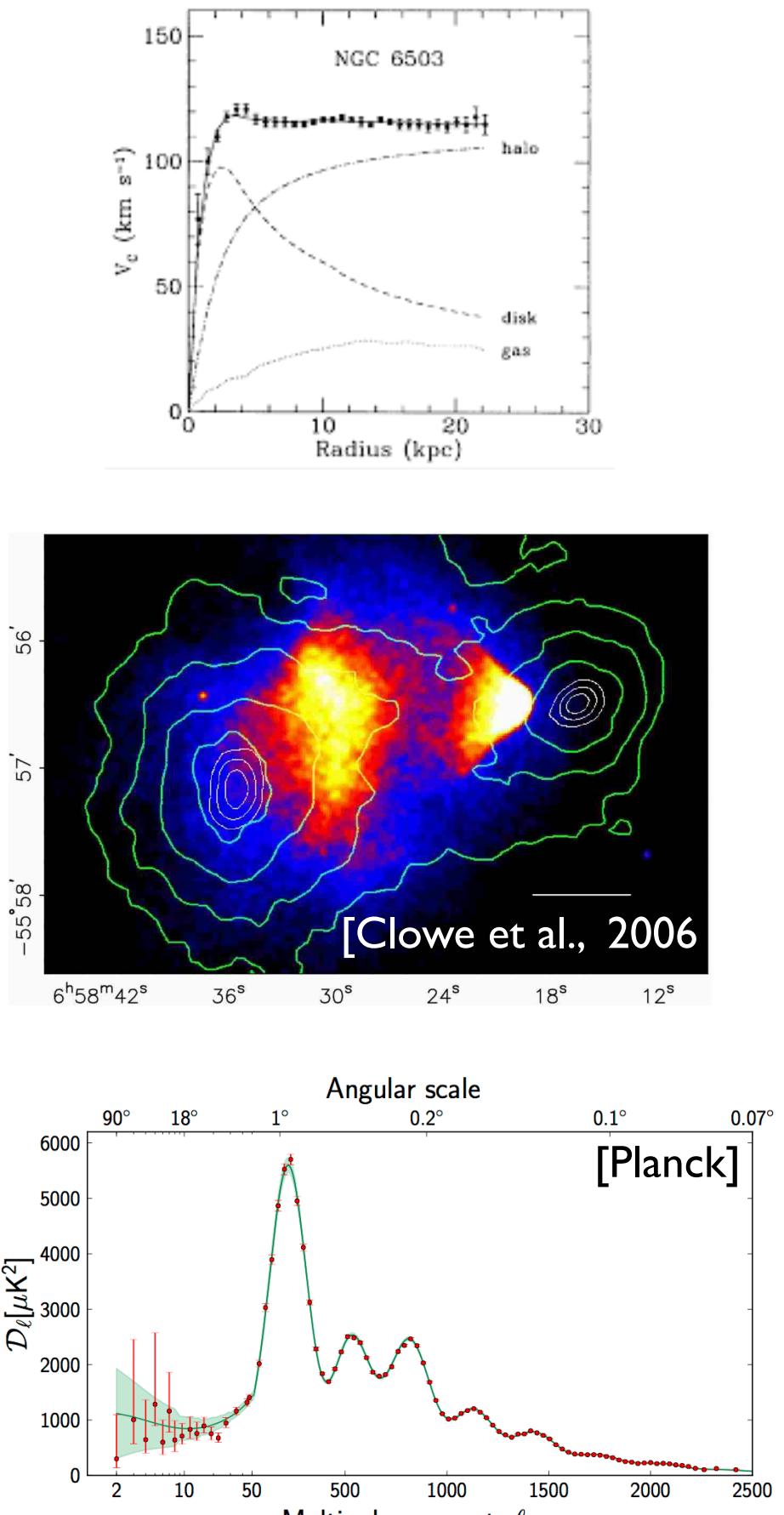
What is the universe made of?

Observations

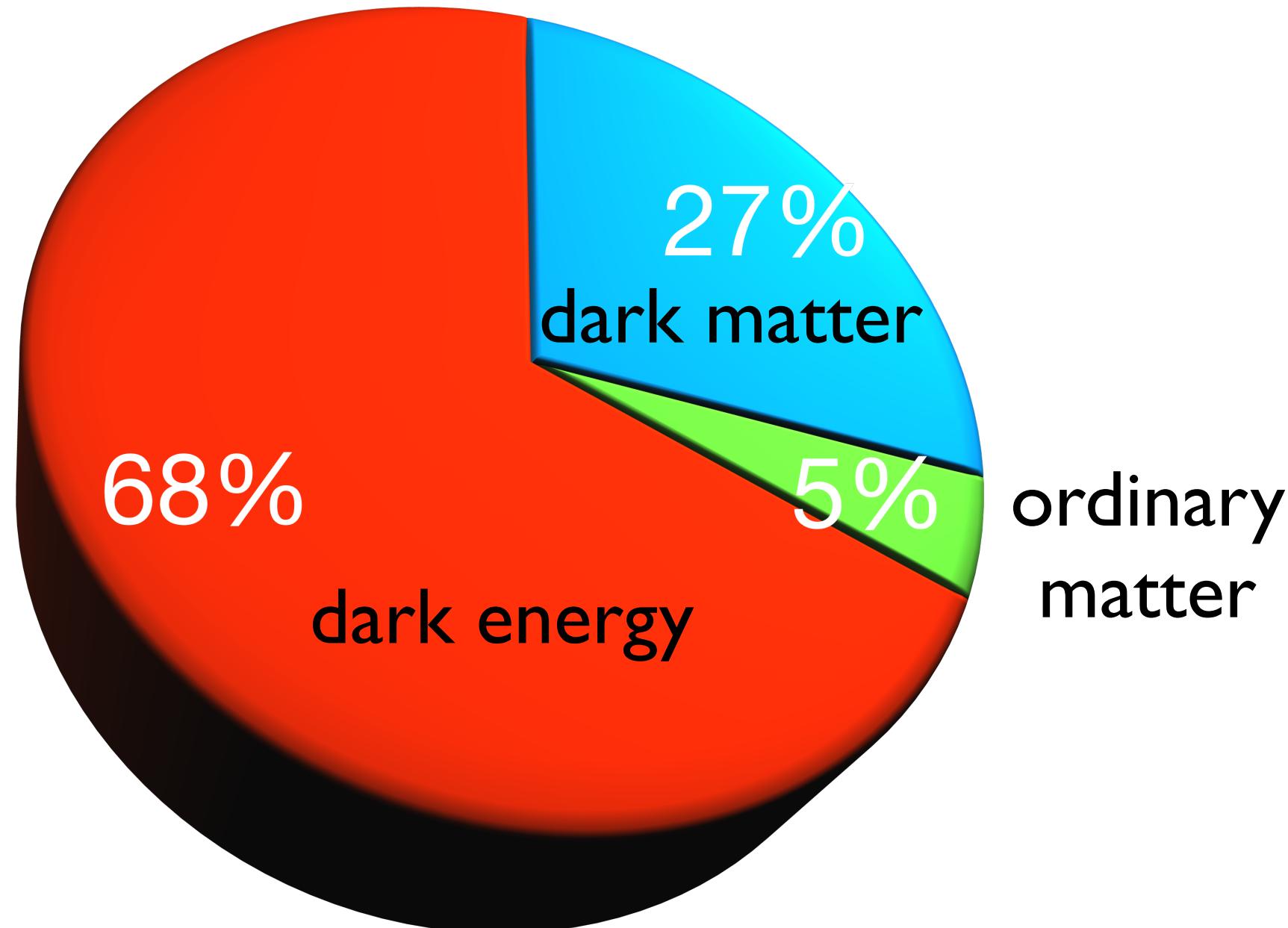


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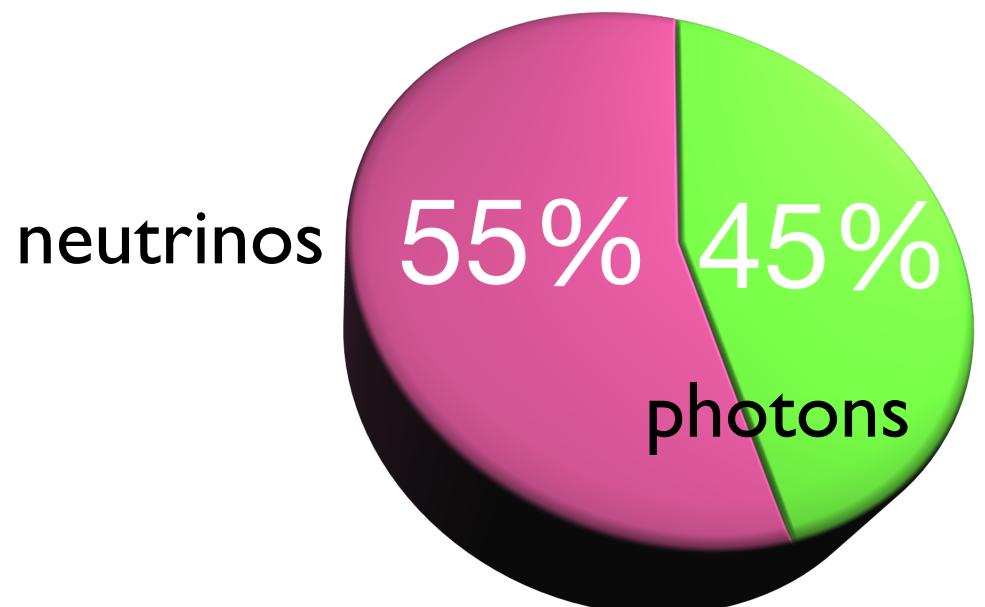
+ ...



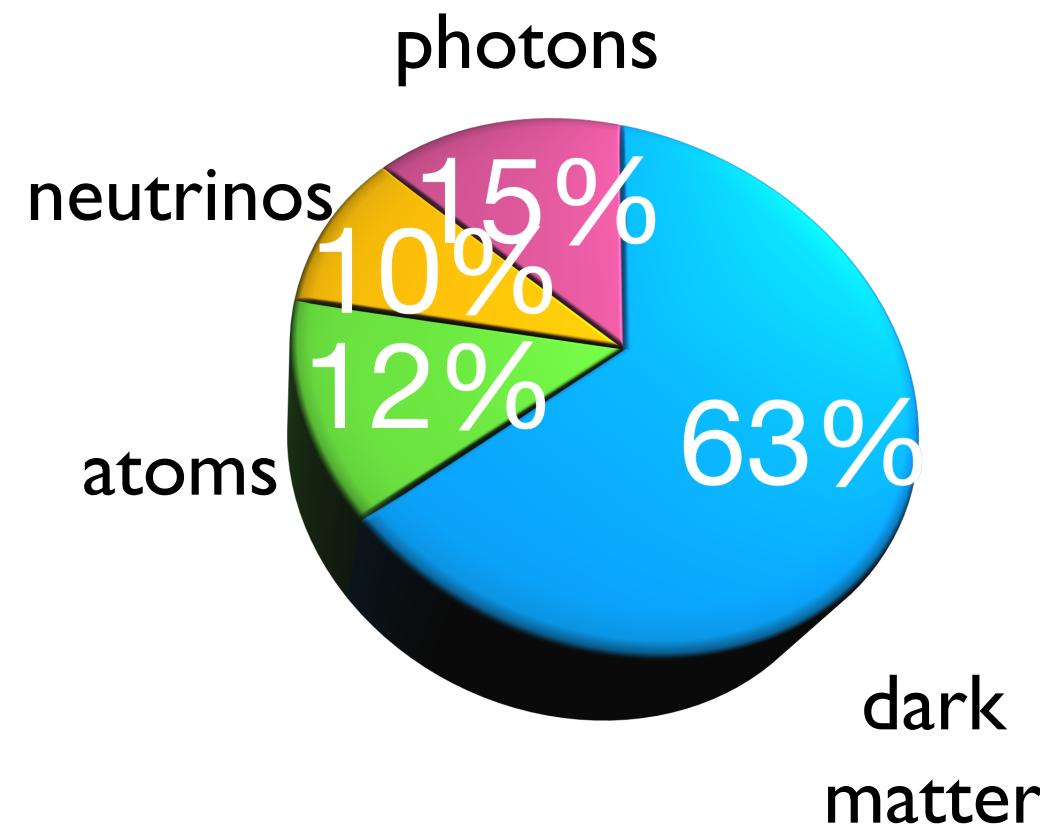
valid now and at large scales

What was the universe made of?

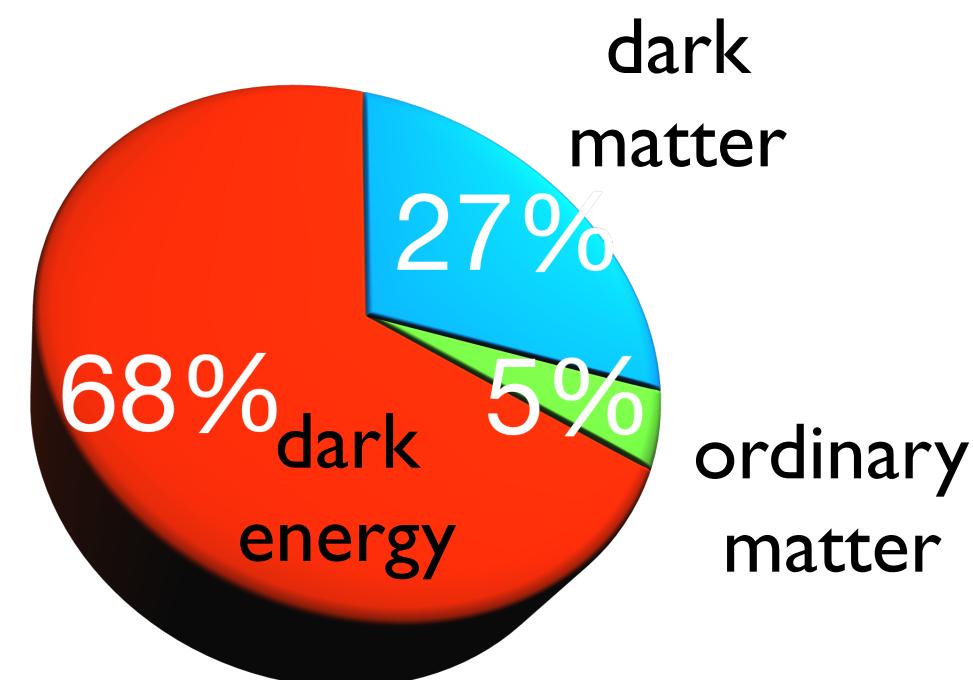
BBN



Recombination



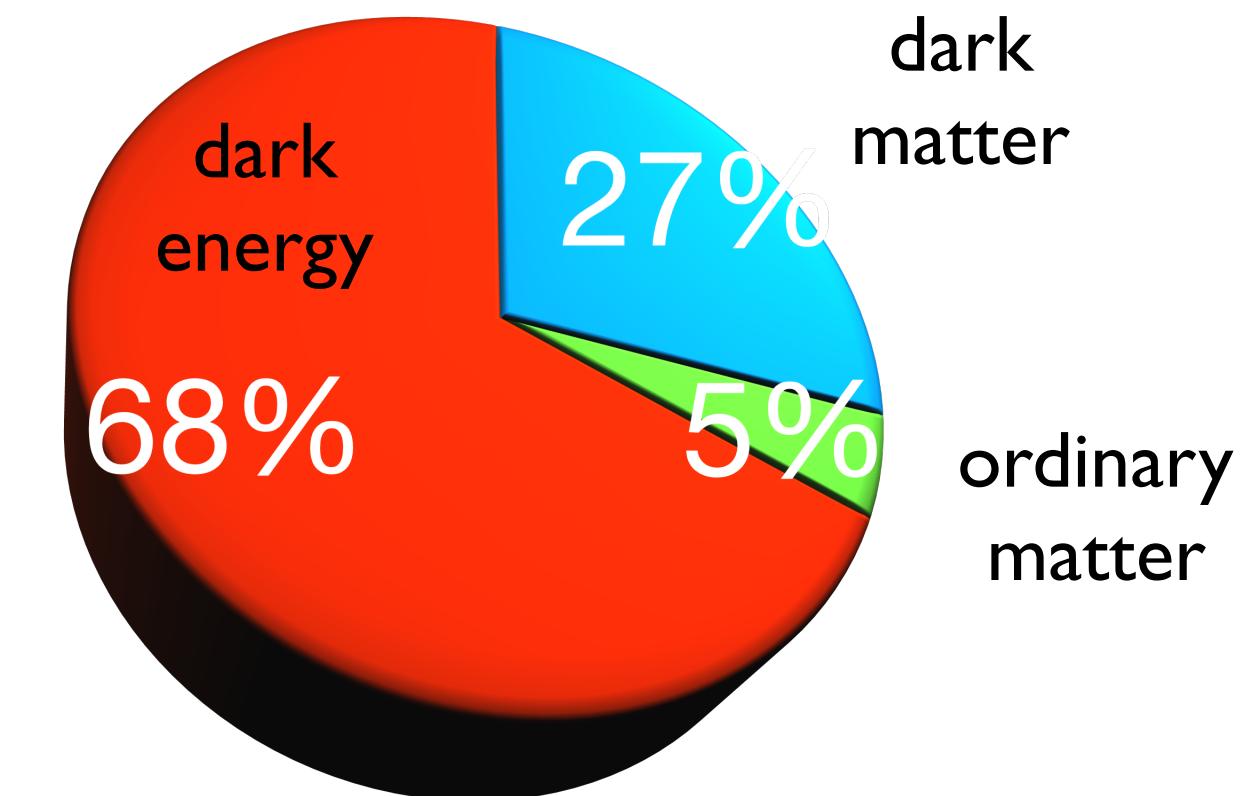
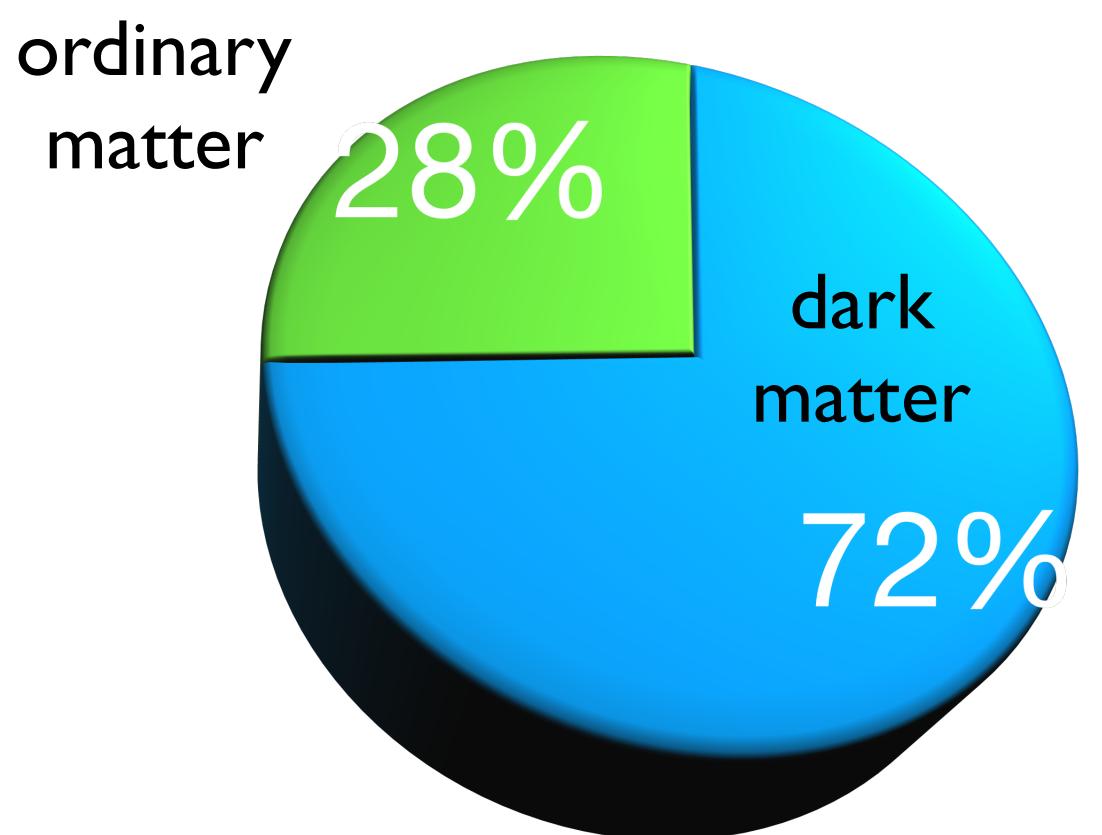
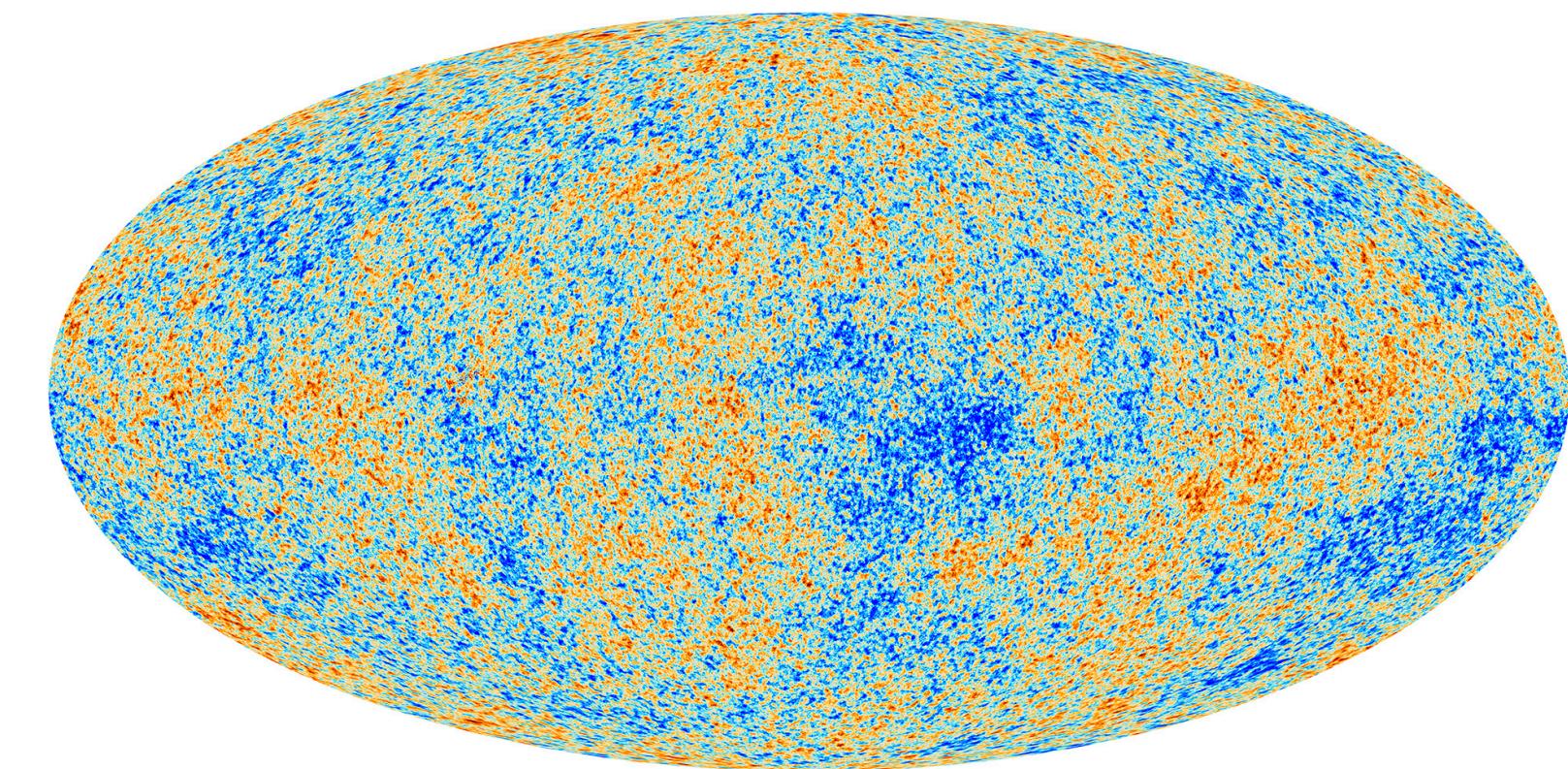
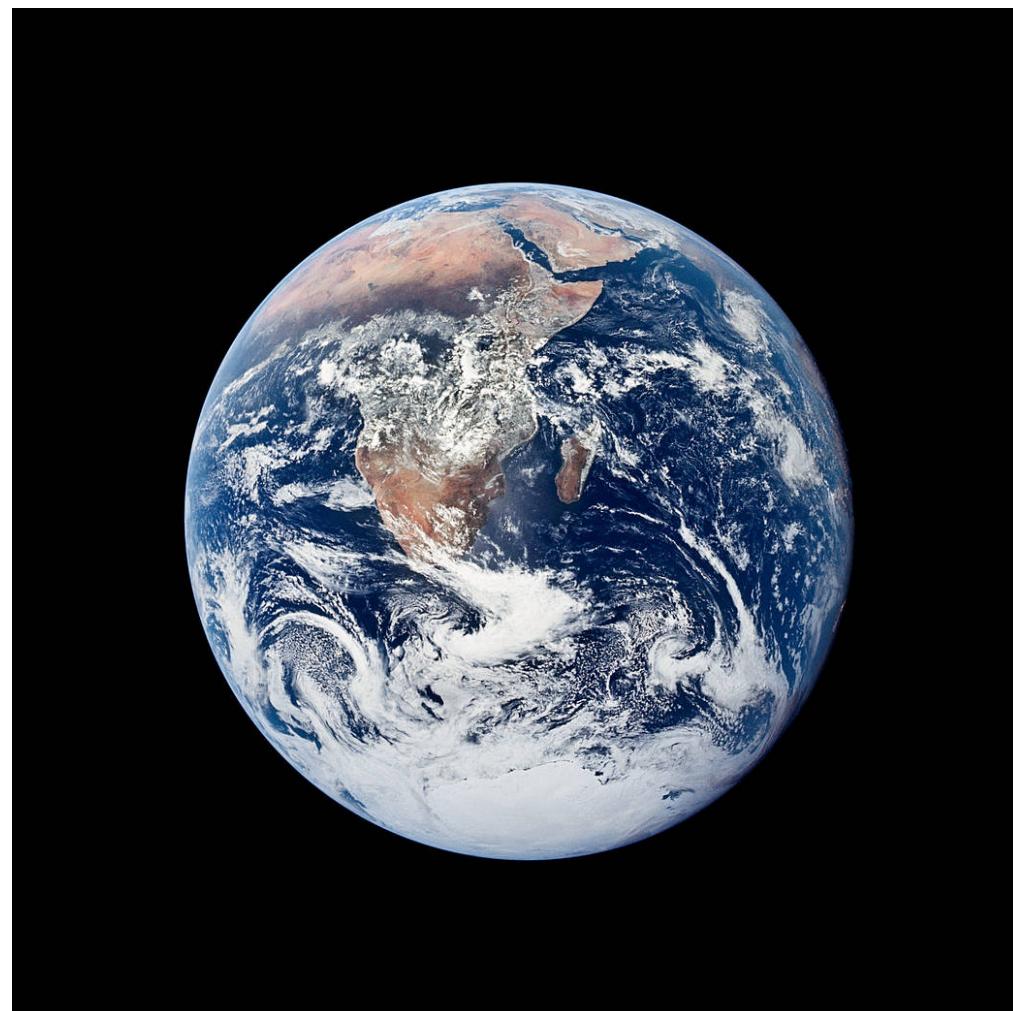
Now



... in a (distant)
future



What is the universe made of?



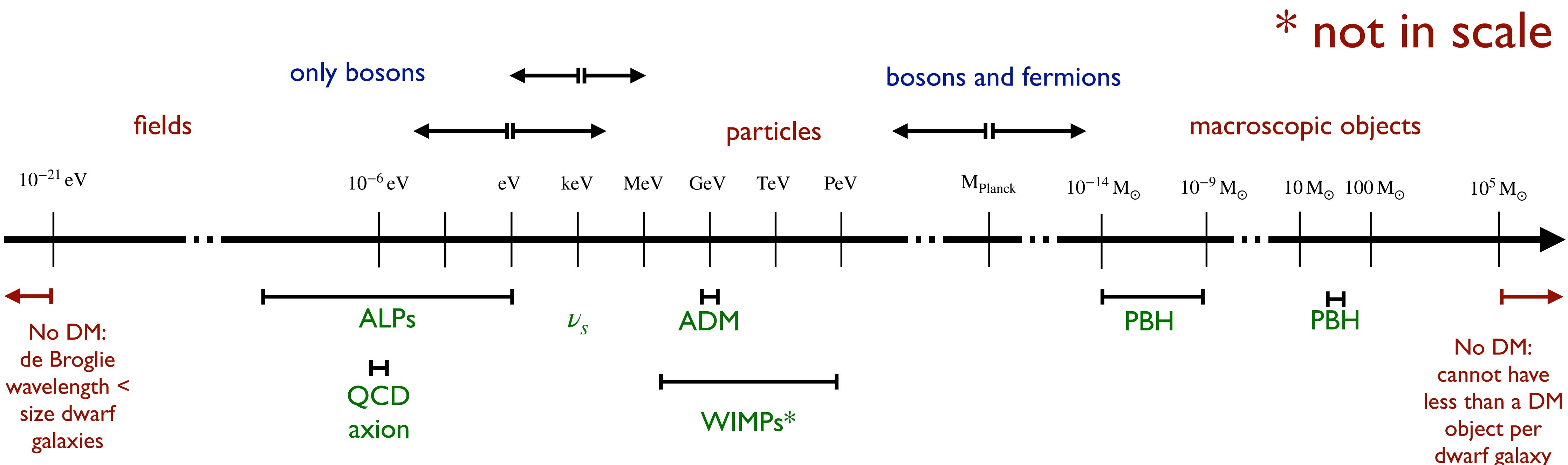
Can X be the DM?

A DM candidate must:

1. match the measured **relic density**;
2. be **cold**;
3. be **neutral**;
4. be **consistent with BBN**;
5. leave **stellar evolution unchanged**.

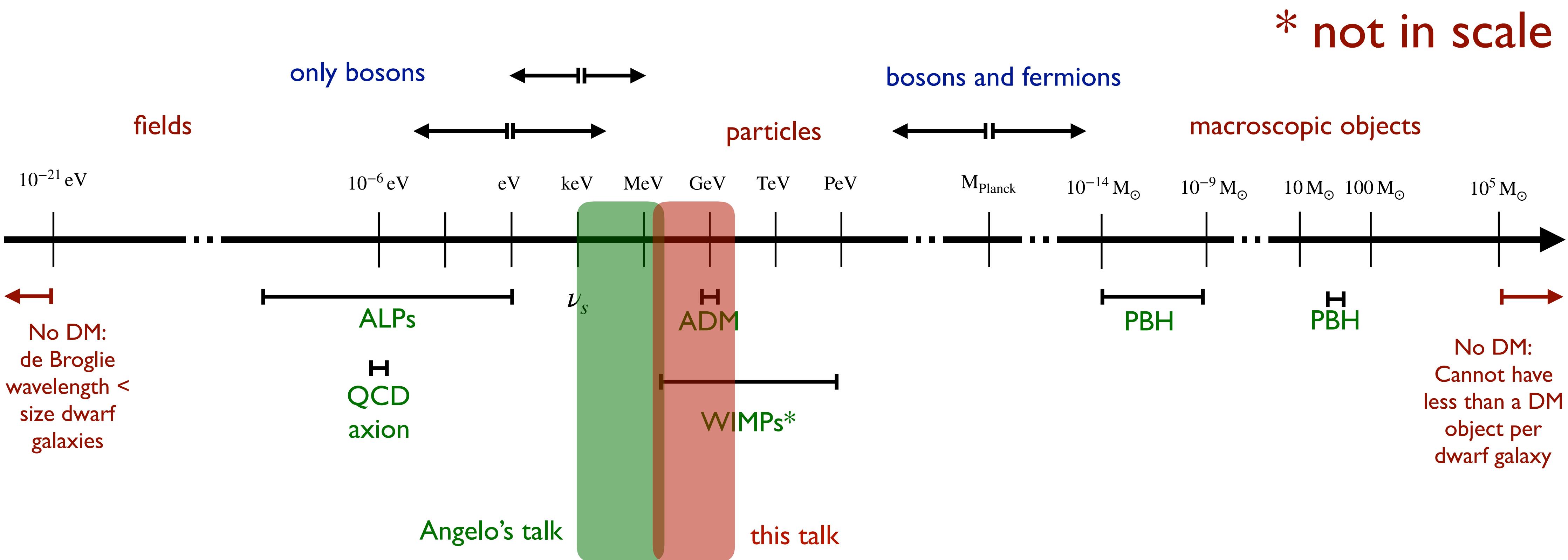
Challenge

We do not know what to search for



Challenges

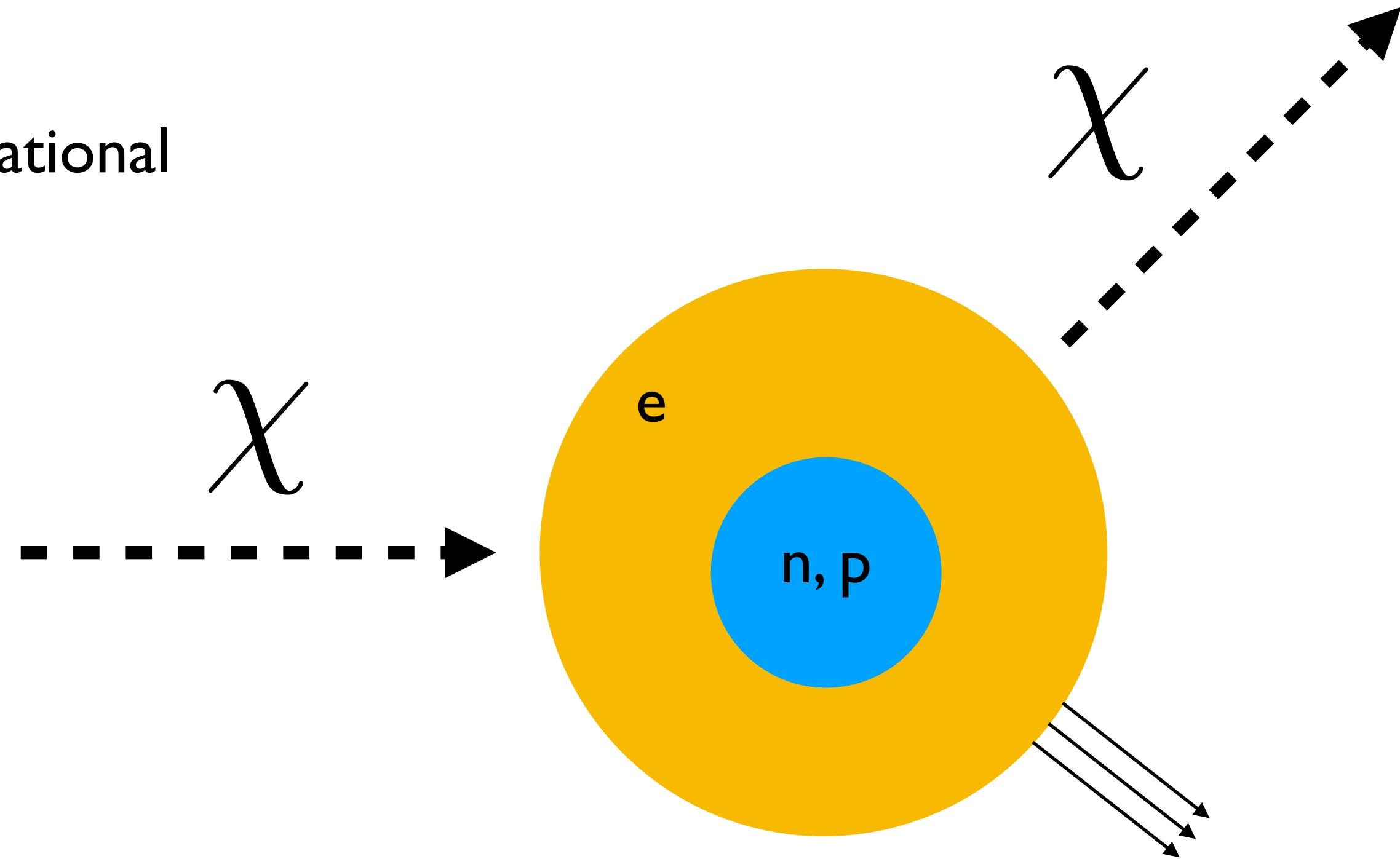
We do not know what to search for



Direct detection phenomenology

Direct detection phenomenology

*assuming non-gravitational interaction exists



$$\frac{d^2R_{\text{nr}}}{dE_R d\Omega_R} = \frac{1}{2\pi} \frac{\rho_0 M_{\text{det}}}{m_N m_\chi} \int \frac{d\sigma}{dE_R} v^2 f(\vec{v}, t) \delta(\vec{v} \cdot \hat{q} - v_{\min}) d^3v$$

Direct detection phenomenology

$$\frac{d^2 R_{\text{nr}}}{dE_R d\Omega_R} = \frac{1}{2\pi} \frac{\rho_0 M_{\text{det}}}{m_N m_\chi} \int \frac{d\sigma}{dE_R} v^2 f(\vec{v}, t) \delta(\vec{v} \cdot \hat{q} - v_{\min}) d^3 v$$

Detector dependence

m_N affects also the cross section

Direct detection phenomenology

Astrophysical inputs

$$\frac{d^2 R_{\text{nr}}}{dE_R d\Omega_R} = \frac{1}{2\pi} \frac{\rho_0 M_{\text{det}}}{m_N m_\chi} \int \frac{d\sigma}{dE_R} v^2 f(\vec{v}, t) \delta(\vec{v} \cdot \hat{q} - v_{\min}) d^3 v$$

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Direct detection phenomenology

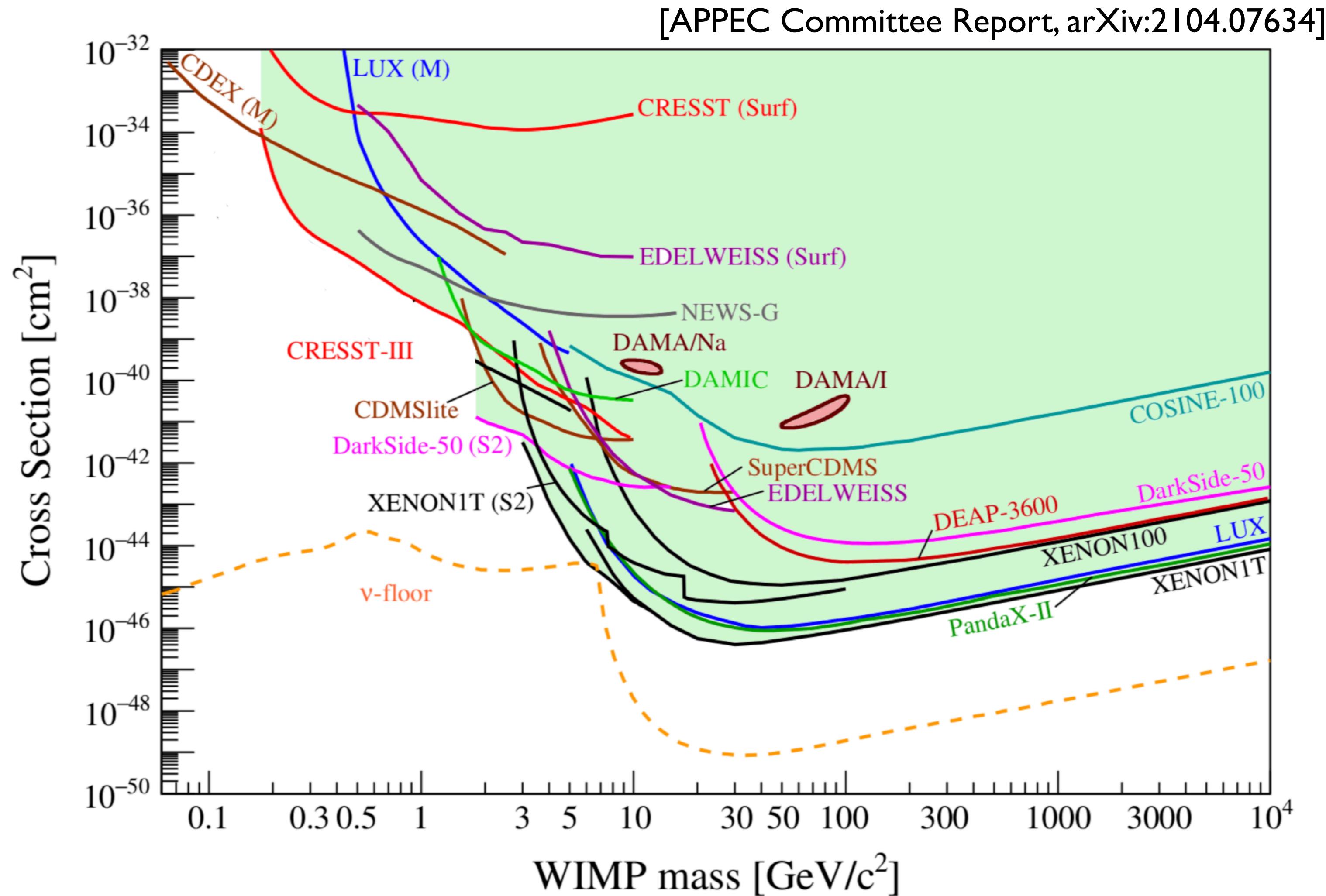
particle
physics Astrophysical inputs

$$\frac{d^2 R_{\text{nr}}}{dE_R d\Omega_R} = \frac{1}{2\pi} \frac{\rho_0 M_{\text{det}}}{m_N m_\chi} \int \frac{d\sigma}{dE_R} v^2 f(\vec{v}, t) \delta(\vec{v} \cdot \hat{q} - v_{\min}) d^3 v$$

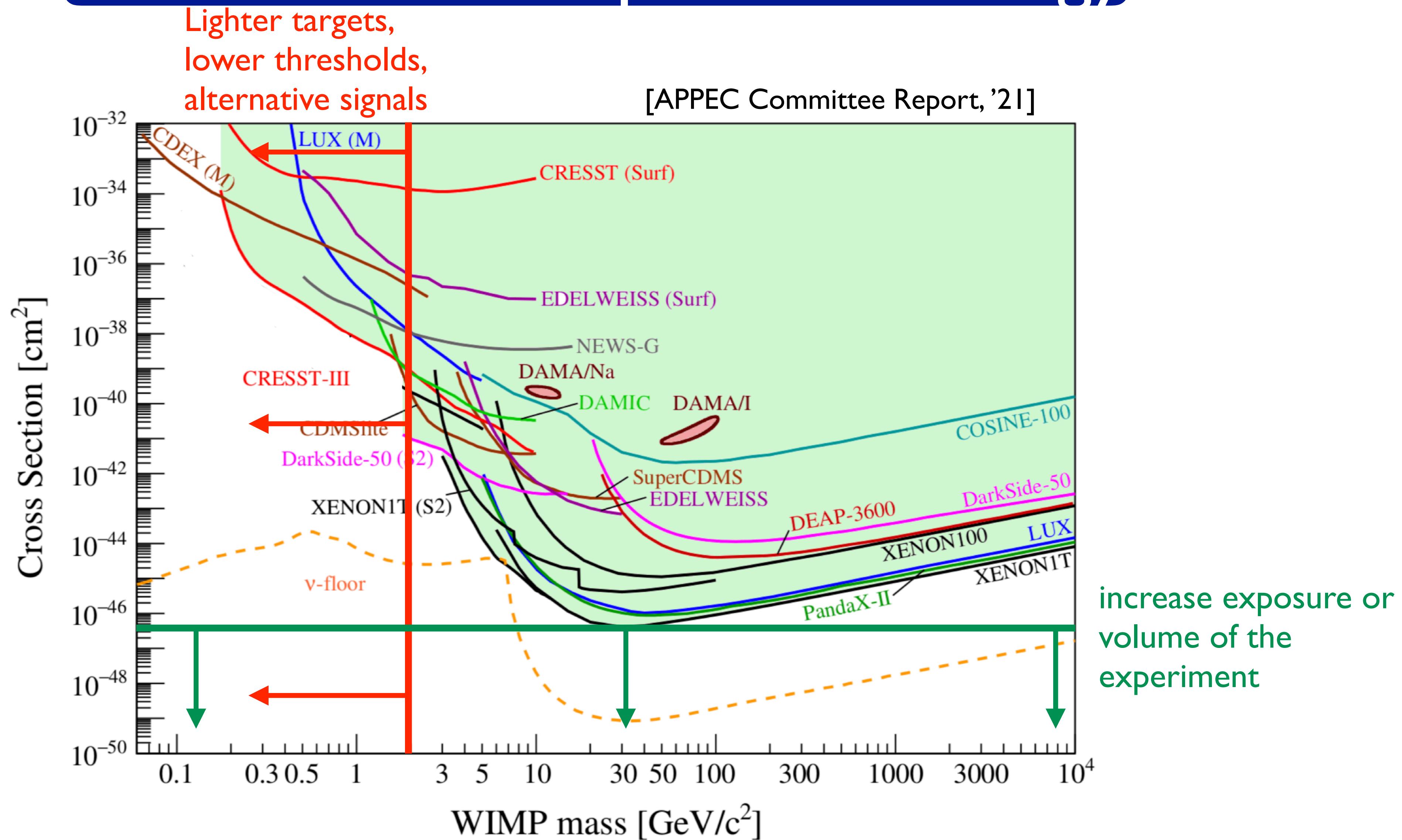
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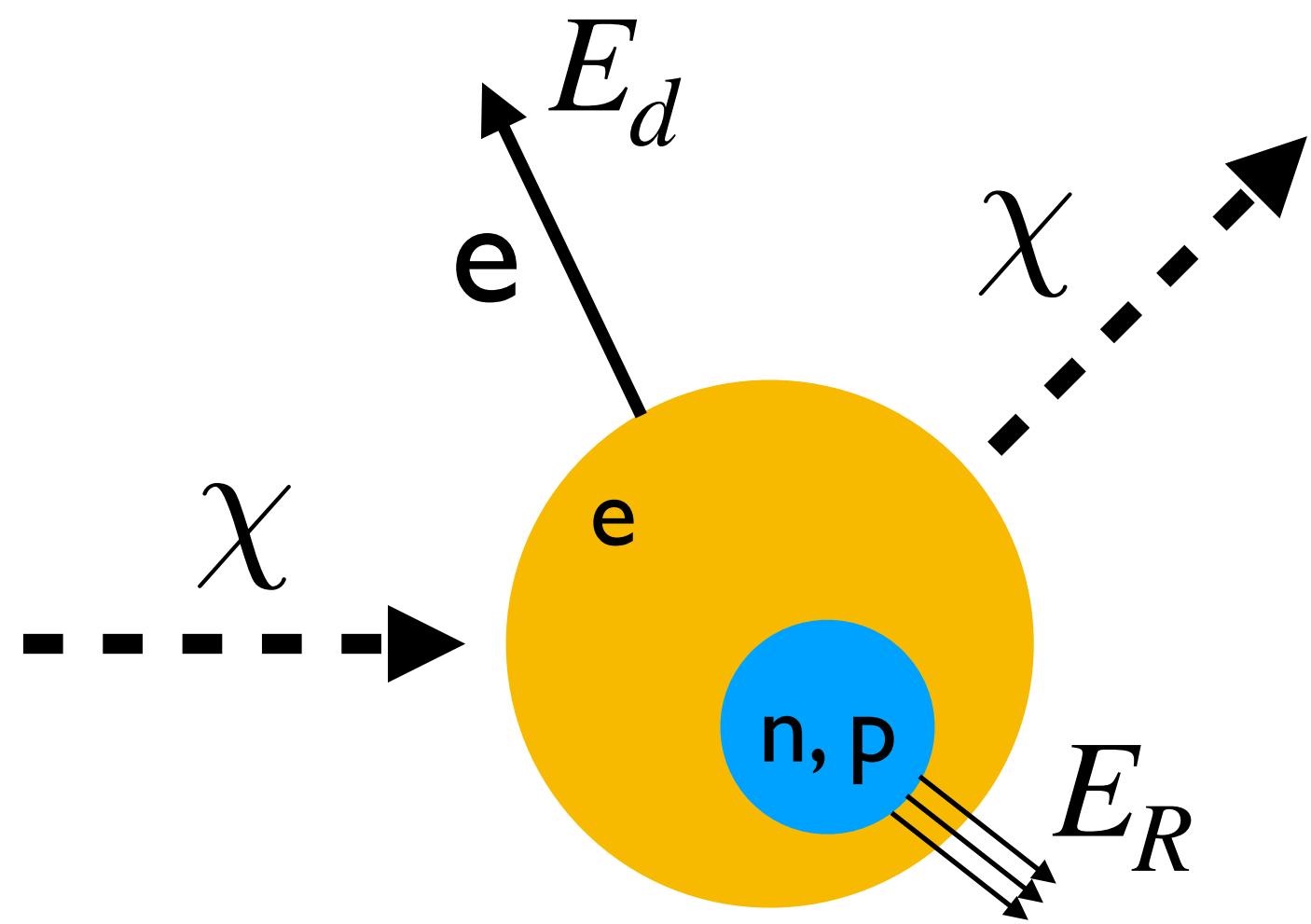


Direct detection phenomenology

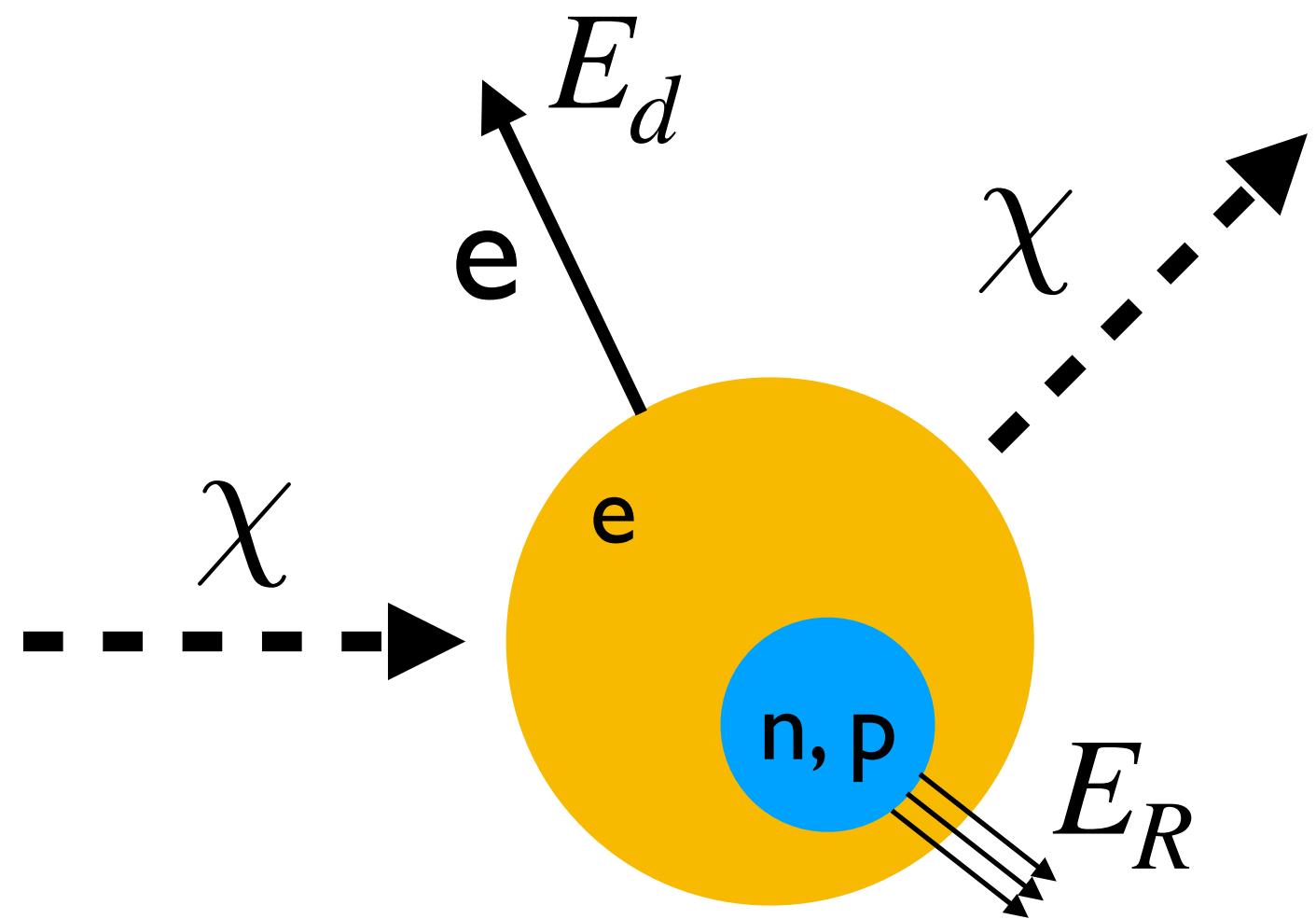


The Migdal effect

The Migdal effect



The Migdal effect



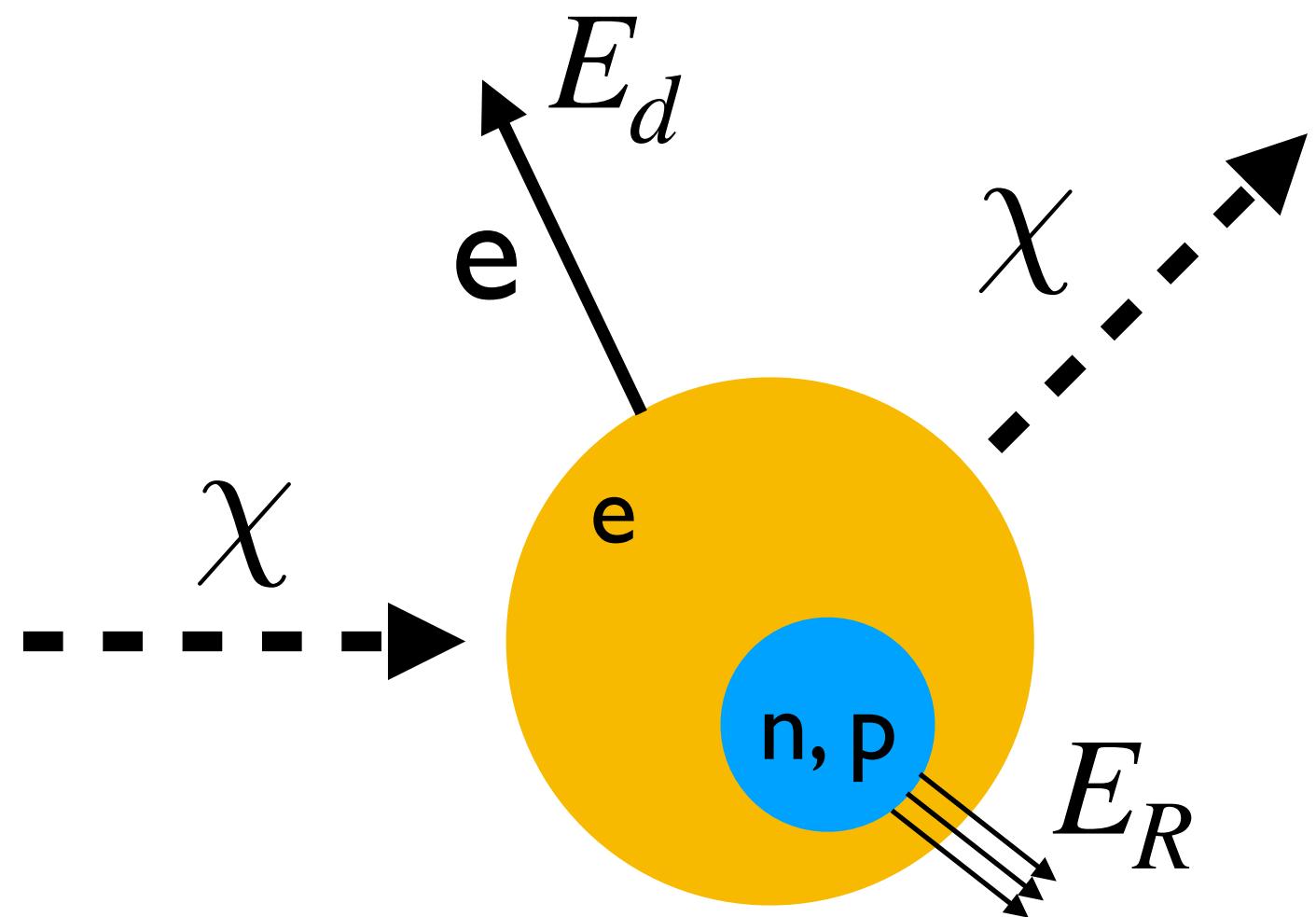
Kinematics

$$v_{\min} = \sqrt{\frac{m_N E_R}{2\mu_N^2}} + \frac{E_d}{\sqrt{2m_N E_R}}$$

$$E_R^{\max} = \frac{2\mu_N^2 v_{\max}^2}{m_N}$$

$$E_d^{\max} = \frac{\mu_N v_{\max}^2}{2}$$

The Migdal effect



Kinematics

Nuclear recoil energy

$$v_{\min} = \sqrt{\frac{m_N E_R}{2 \mu_N^2}} + \frac{E_d}{\sqrt{2 m_N E_R}}$$

Electron detected energy

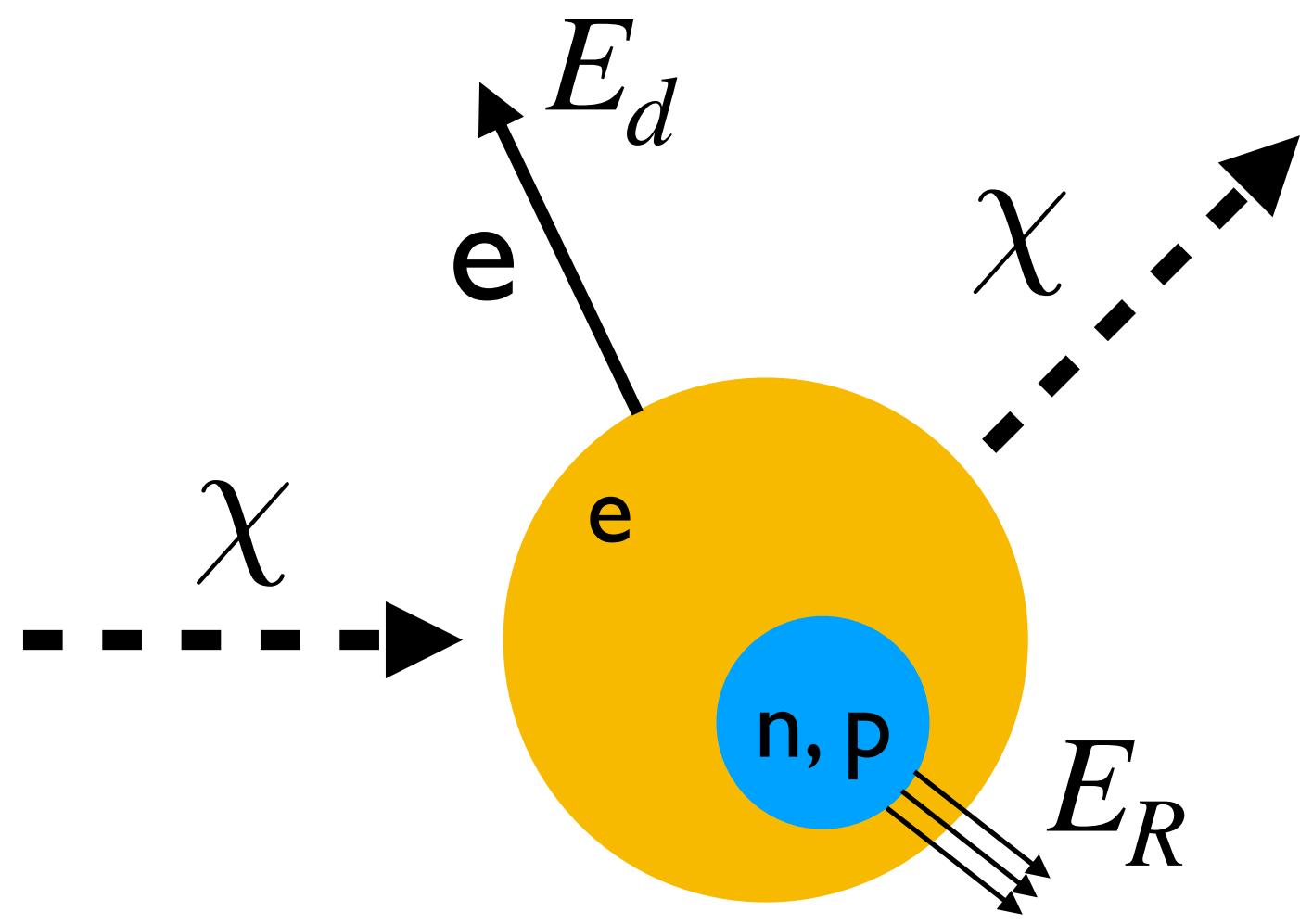
DM-nucleus reduced mass

nucleus mass

$$E_R^{\max} = \frac{2 \mu_N^2 v_{\max}^2}{m_N}$$

$$E_d^{\max} = \frac{\mu_N v_{\max}^2}{2}$$

The Migdal effect

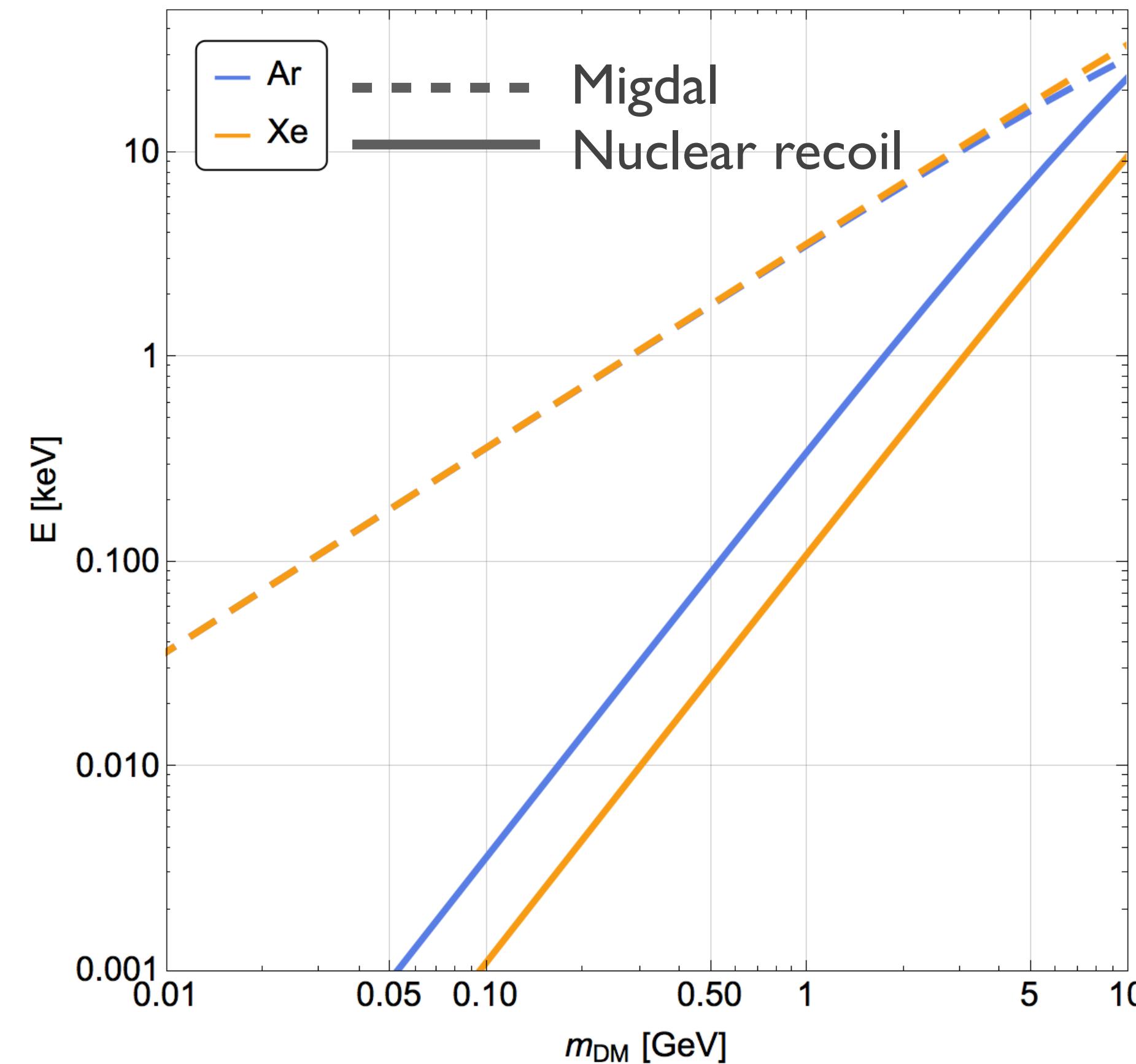


1. Threshold $E \lesssim 1 \text{ keV}$
2. Sensitivity loss for $m_{DM} \lesssim 2 \text{ GeV}$
3. $E_d^{\max} > E_R^{\max}$ for $m_{DM} \ll m_N$
4. The Migdal effect is sensitive to sub-GeV masses

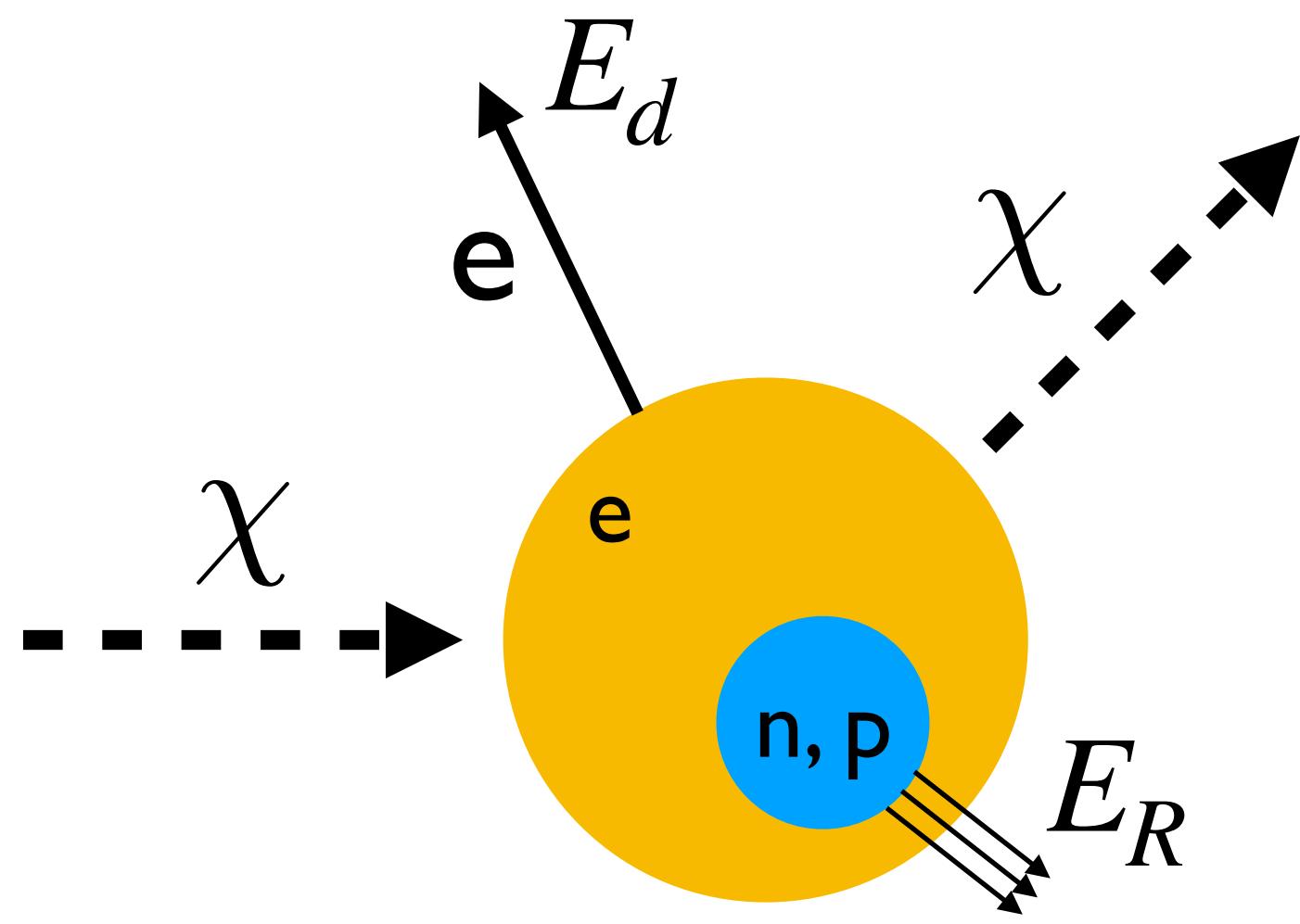
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The Migdal effect

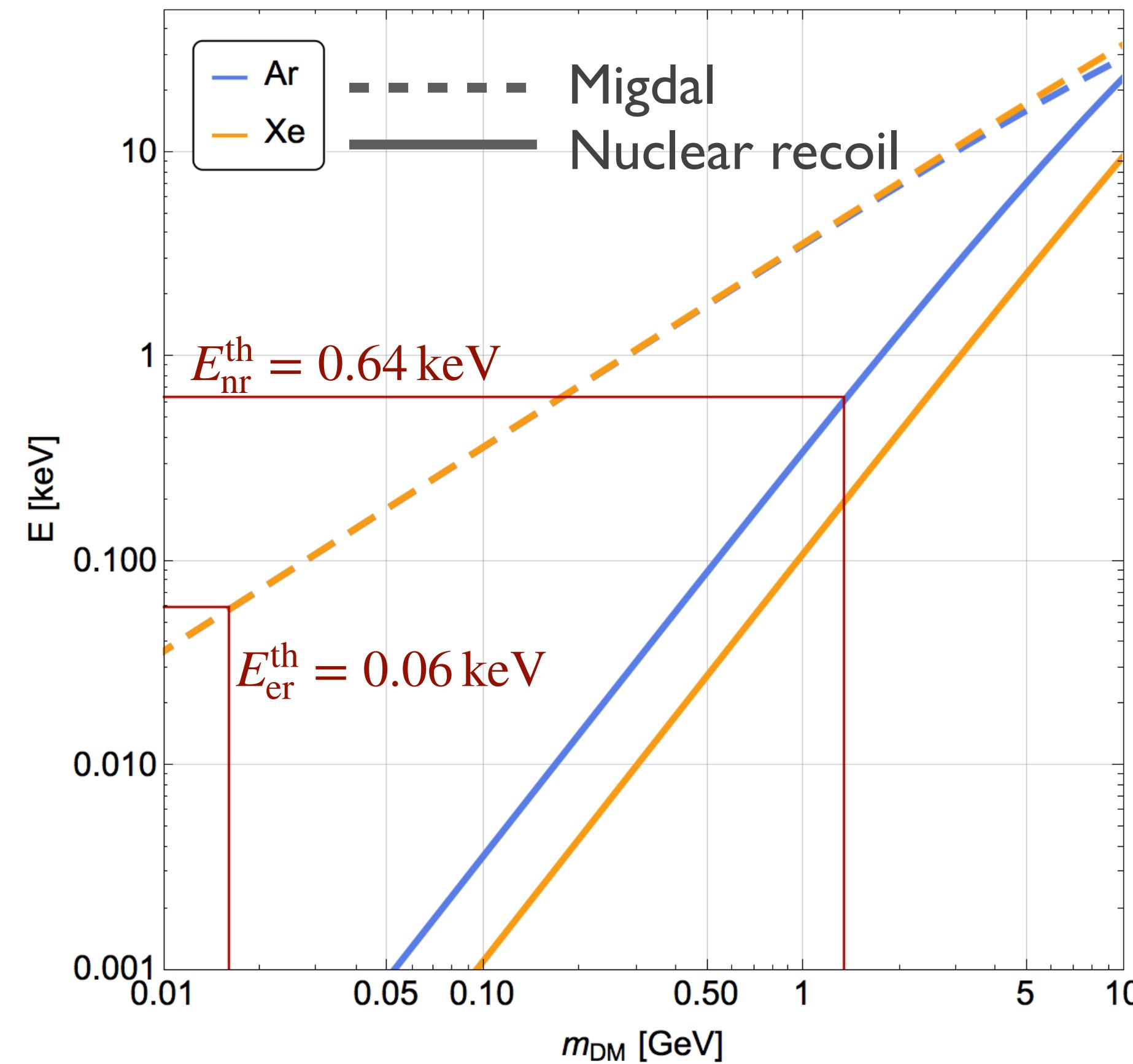


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The Migdal effect

$$\frac{d^3 R_{\text{mig}}}{dE_R dE_e dv} = \frac{d^2 R_{\text{nr}}}{dE_R dv} |Z(E_R, E_e)|^2$$

The Migdal effect

$$\frac{d^3 R_{\text{mig}}}{dE_R dE_e dv} = \boxed{\frac{d^2 R_{\text{nr}}}{dE_R dv}} |Z(E_R, E_e)|^2$$

DM nuclear recoil

The Migdal effect

$$\frac{d^3 R_{\text{mig}}}{dE_R dE_e dv} = \frac{d^2 R_{\text{nr}}}{dE_R dv} |Z(E_R, E_e)|^2$$

NR
de-excitation:
negligible

ionization
rate

DM nuclear recoil

$|Z(E_R, E_e)|^2 \simeq 1 + |Z_{\text{de}}|^2 + |Z_{\text{ion}}|^2$

The Migdal effect

$$\frac{d^3 R_{\text{mig}}}{dE_R dE_e dv} = \frac{d^2 R_{\text{nr}}}{dE_R dv} |Z(E_R, E_e)|^2$$

NR
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DM nuclear recoil

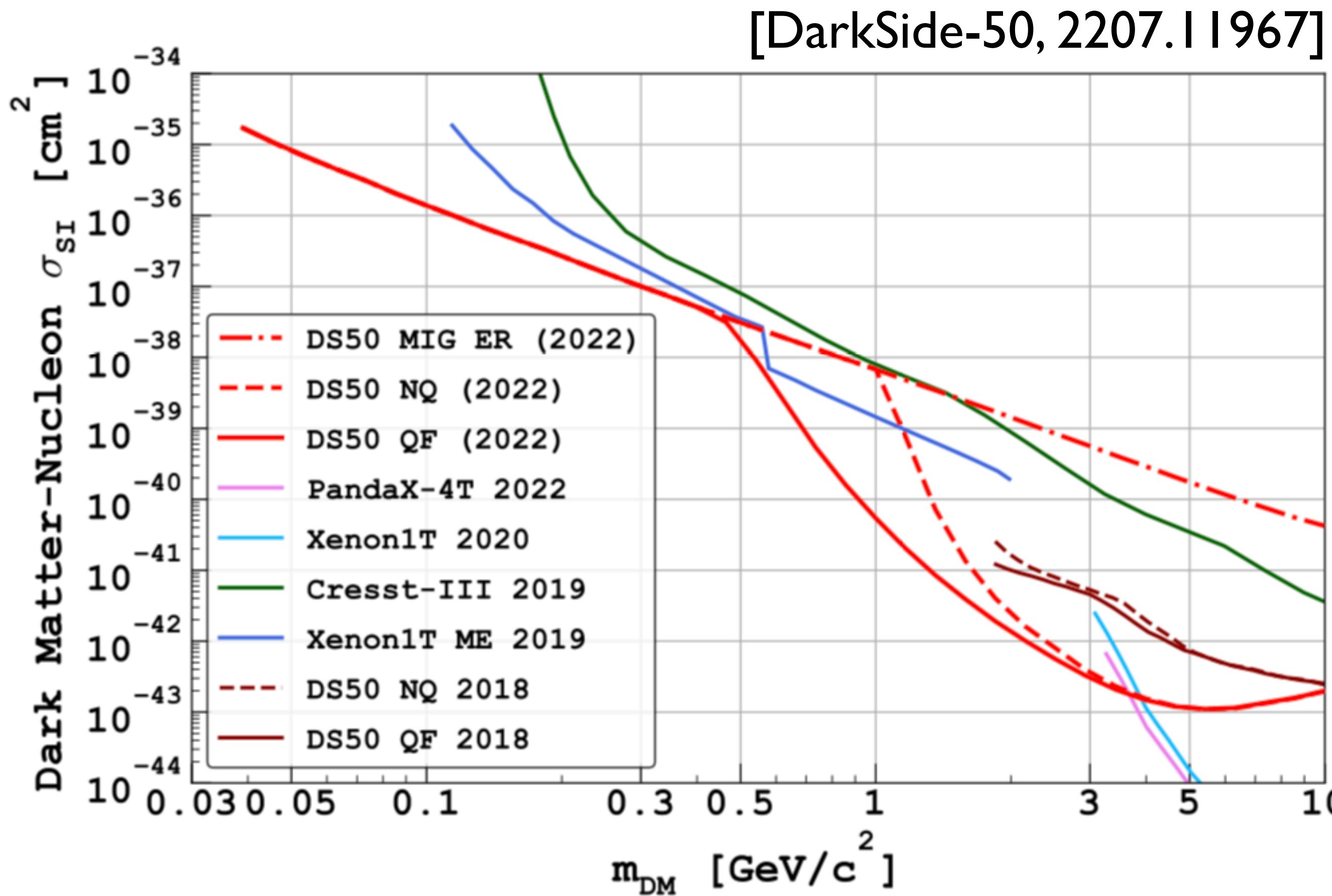
$|Z(E_R, E_e)|^2 \simeq 1 + |Z_{\text{de}}|^2 + |Z_{\text{ion}}|^2$

$$|Z_{\text{ion}}(E_R, E_e)|^2 = \frac{1}{2\pi} \sum_{n,\ell} \int dE_e \frac{dp_{q_e}^c(n\ell \rightarrow E_e)}{dE_e}$$

Migdal

Computed in [Ibe et al. JHEP03(2018)194] for C, F, Ar, Xe, ..., see also [Cox et al., 2208.12222]

The Migdal effect

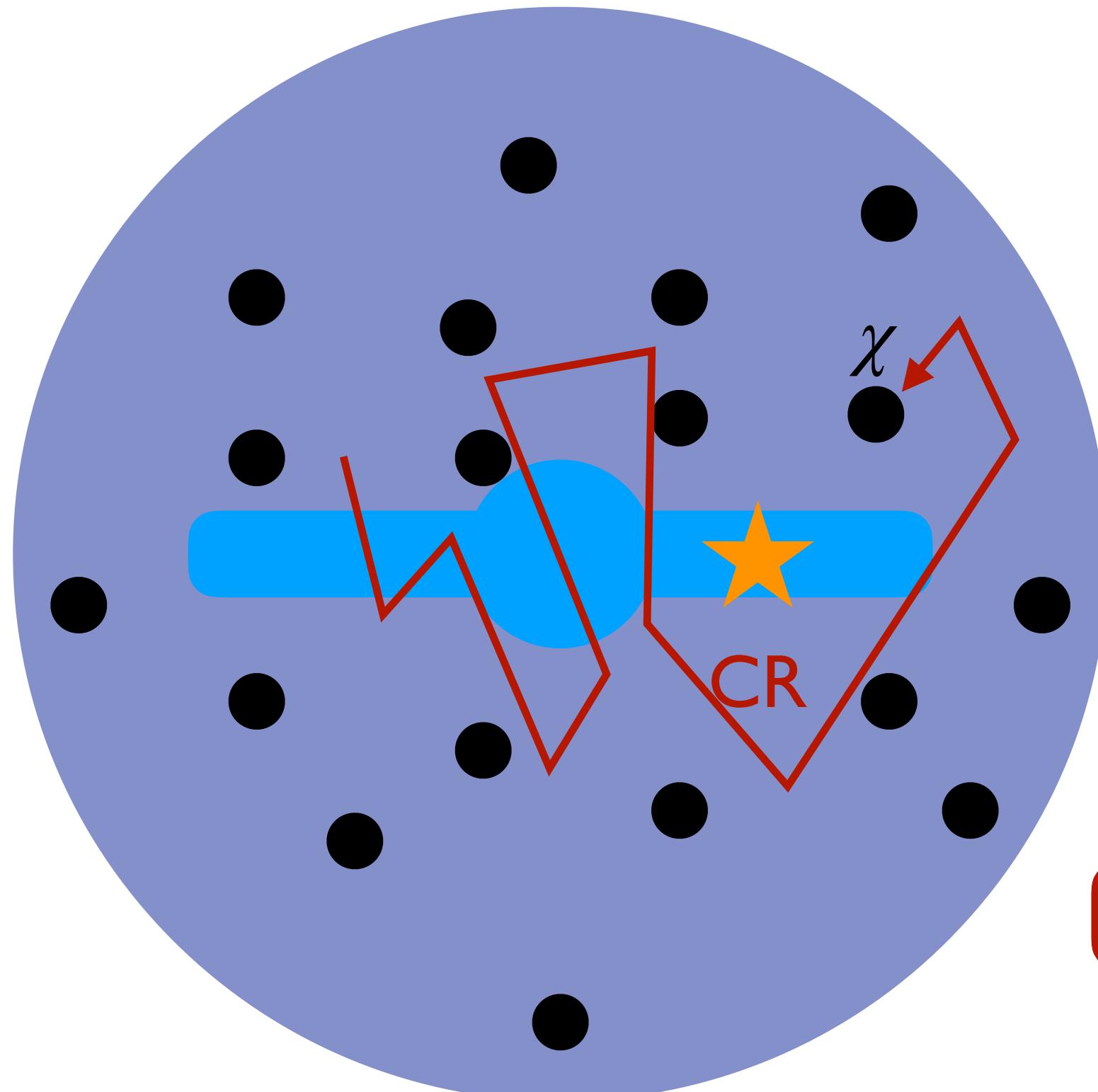


Boosted Dark Matter

Cosmic Ray boosted DM

[Bringmann and Pospelov, 1810.10543]

New isotropic component of DM extending to higher energies



DM density as a function of the distance to the galactic centre

CR flux as a function of the distance to the galactic centre

$$\frac{d\Phi_\chi}{dT_i} = \int \frac{d\Omega}{4\pi} \int_{l.o.s.} d\ell \frac{\sigma_{\chi i}}{m_\chi} \rho \frac{d\Phi_i}{dT_i}$$

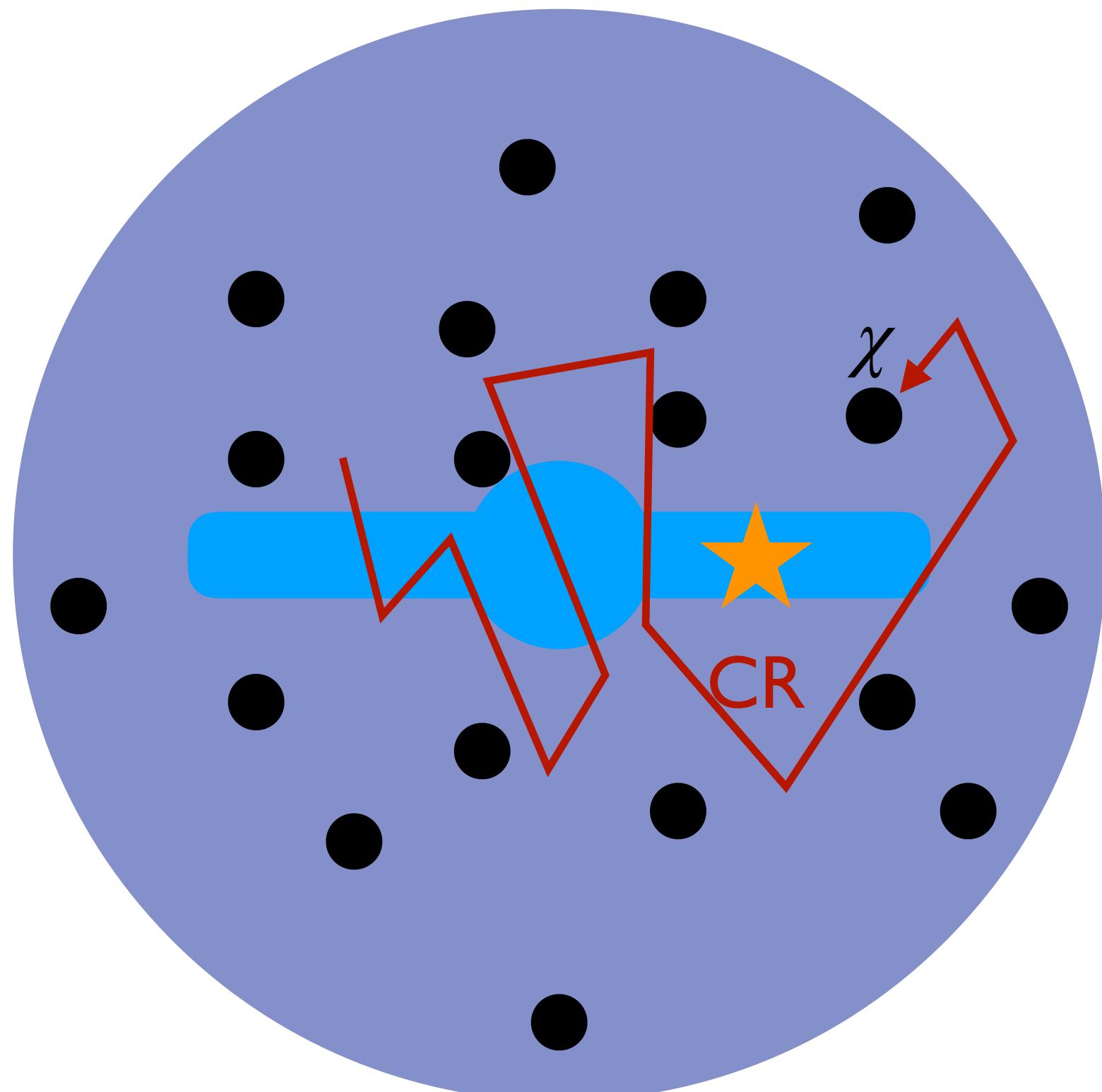
Denotes the CR

CR-DM cross section

Cosmic Ray boosted DM

[Bringmann and Pospelov, 1810.10543]

New isotropic component of DM extending to higher energies



$$\frac{d\Phi_\chi}{dT_i} = \int \frac{d\Omega}{4\pi} \int_{l.o.s.} d\ell \sigma_{\chi i} \frac{\rho}{m_\chi} \frac{d\Phi_i}{dT_i} \equiv \sigma_{\chi i}^{\text{local}} \frac{\rho_\chi^{\text{local}}}{m_\chi} \frac{d\Phi_i^{LIS}}{dT_i} D_{\text{eff}}$$

Denotes the CR

CR-DM cross section

Effective distance parametrizing uncertainty from l.o.s. integration

Cosmic Ray boosted DM

Production

Form factor

$$\frac{d\Phi_\chi}{dT_\chi} = D_{\text{eff}} \frac{\rho_\chi^{\text{local}}}{m_\chi} \sum_i \sigma_\chi i F_i^2(2m_\chi T_\chi) \int_{T_i^{\min}}^{\infty} \frac{dT_i}{T_\chi^{\max}(T_i)} \frac{d\Phi_i^{LIS}}{dT_i}$$

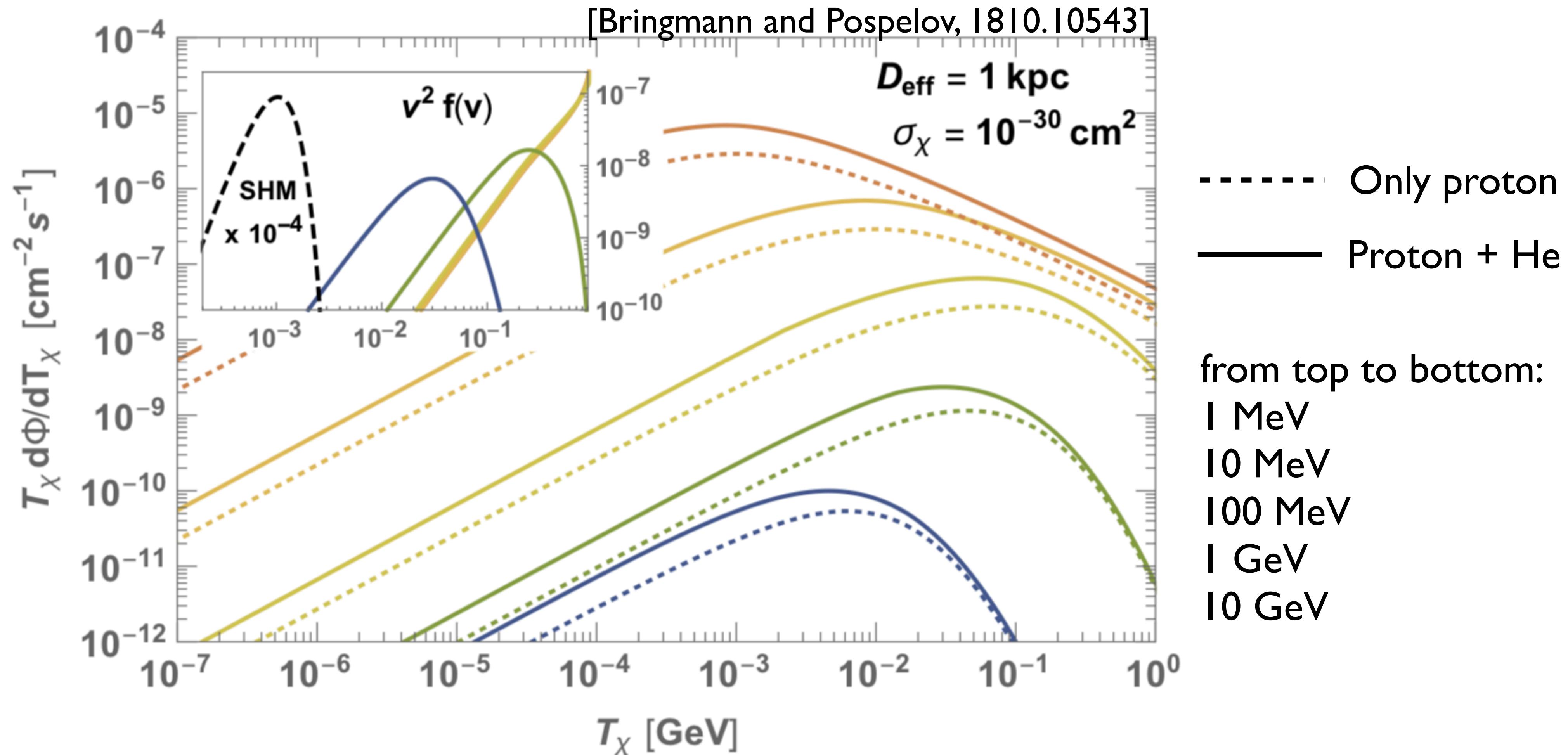
$$i = \{\text{p, He, ...}\}$$

$$T_\chi^{\max} = \frac{T_i^2 + 2m_i T_i}{T_i + (m_i + m_\chi)^2 / (2m_\chi)}$$

additional terms compared to the non-relativistic expression

Cosmic Ray boosted DM

Production



Cosmic Ray boosted DM

Earth attenuation

Solve energy loss equation

Easily solvable if it does not depend on the form factor

$$\frac{d\sigma_{\chi N}}{dT_r} = \frac{\sigma_{\chi N}}{T_r^{\max}} F^2(T_r)$$

$$\frac{dT_\chi}{dz} = - \sum_N n_N \int_0^{T_r^{\max}} \frac{d\sigma_{\chi N}}{dT_r} T_r dT_r$$

to obtain flux at the detector depth

$$\frac{d\Phi_\chi}{dT_\chi^z} = \left(\frac{dT_\chi}{dT_\chi^z} \right) \frac{d\Phi_\chi}{dT_\chi} = \frac{4m_\chi^2 e^{z/\ell}}{(2m_\chi + T_\chi^z - T_\chi^z e^{z/\ell})^2} \frac{d\Phi_\chi}{dT_\chi}$$

where the mean free path is

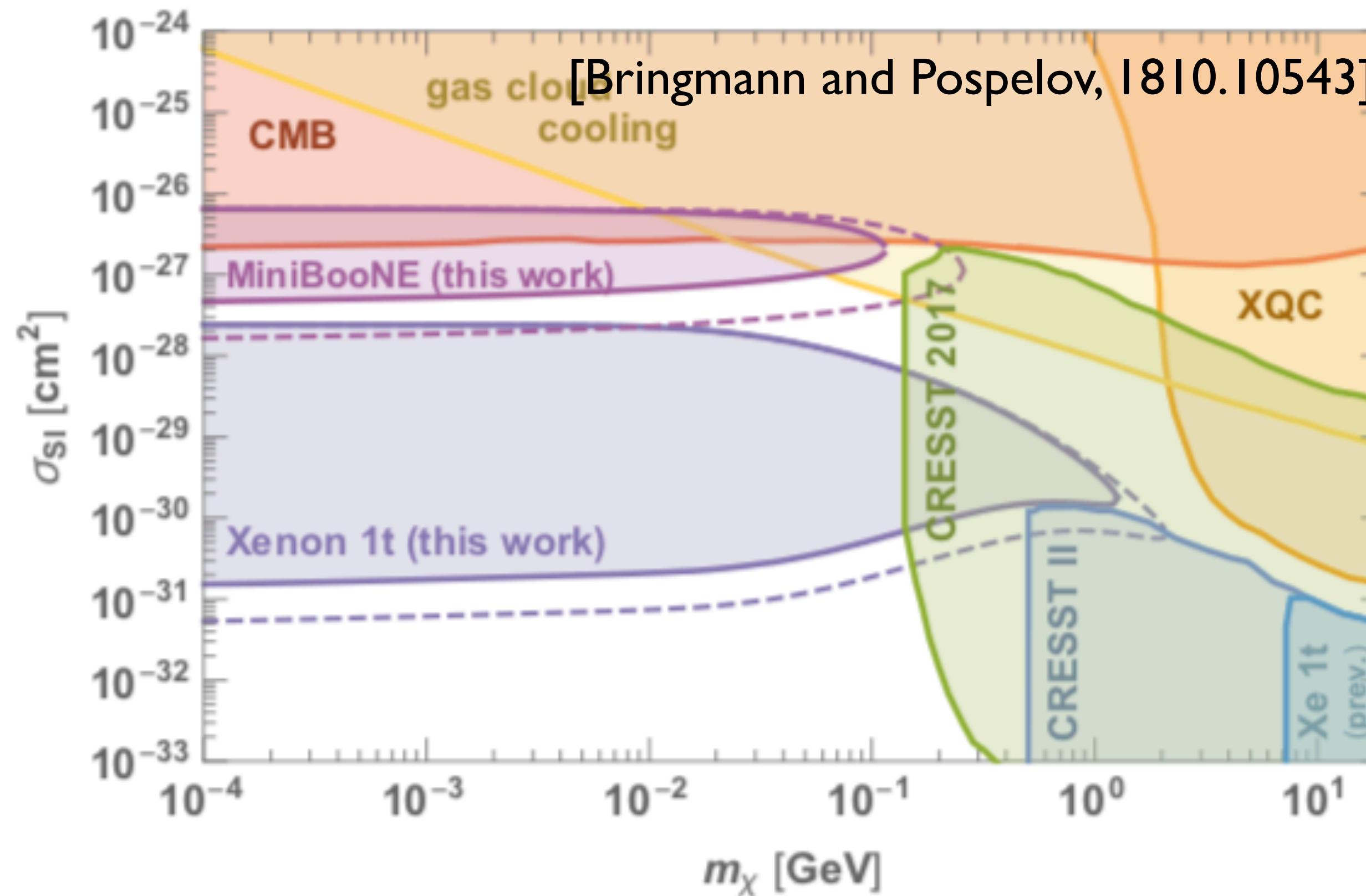
$$\ell^{-1} \equiv \sum_N n_N \sigma_{\chi N} \frac{2 m_N m_\chi}{(m_N + m_\chi)^2}$$

Cosmic Ray boosted DM

Detection

$$\frac{d\Gamma_N}{dT_N} = \int_{T_\chi^{\min}}^{\infty} dT_\chi \frac{d\sigma_{\chi N}}{dT_N} \frac{d\Phi_\chi}{dT_\chi} = \sigma_{\chi N} F_N^2(2m_N T_N) \int_{T_\chi(T_\chi^{\text{z, min}})}^{\infty} \frac{dT_\chi}{T_N^{\max}} \frac{d\Phi_\chi}{dT_\chi}$$

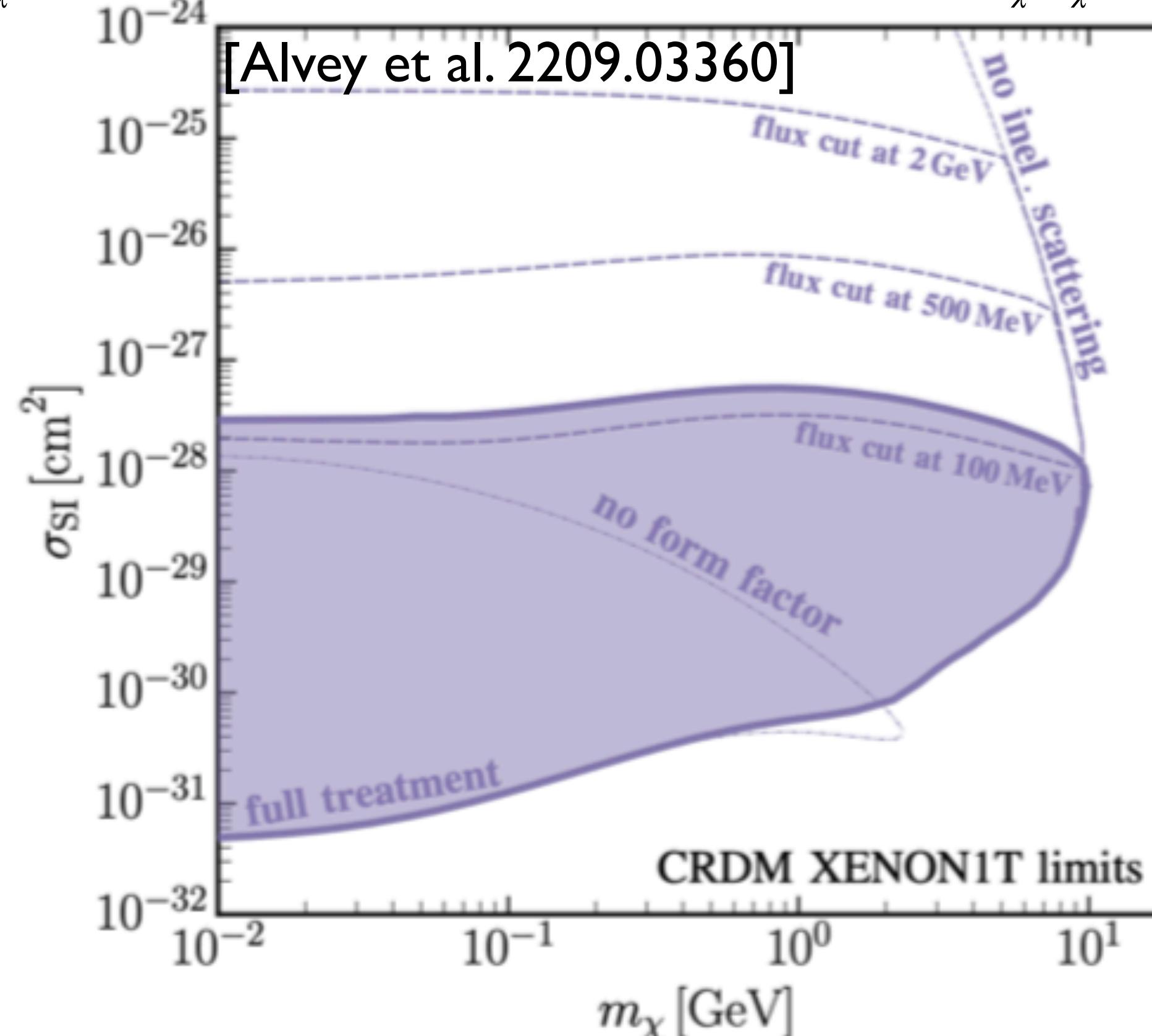
Assuming no form factor in
the Earth attenuation



Cosmic Ray boosted DM

Detection

$$\frac{d\Gamma_N}{dT_N} = \int_{T_\chi^{\min}}^{\infty} dT_\chi \frac{d\sigma_{\chi N}}{dT_N} \frac{d\Phi_\chi}{dT_\chi} = \sigma_{\chi N} F_N^2(2m_N T_N) \int_{T_\chi(T_\chi^{\text{z, min}})}^{\infty} \frac{dT_\chi}{T_N^{\max}} \frac{d\Phi_\chi}{dT_\chi}$$

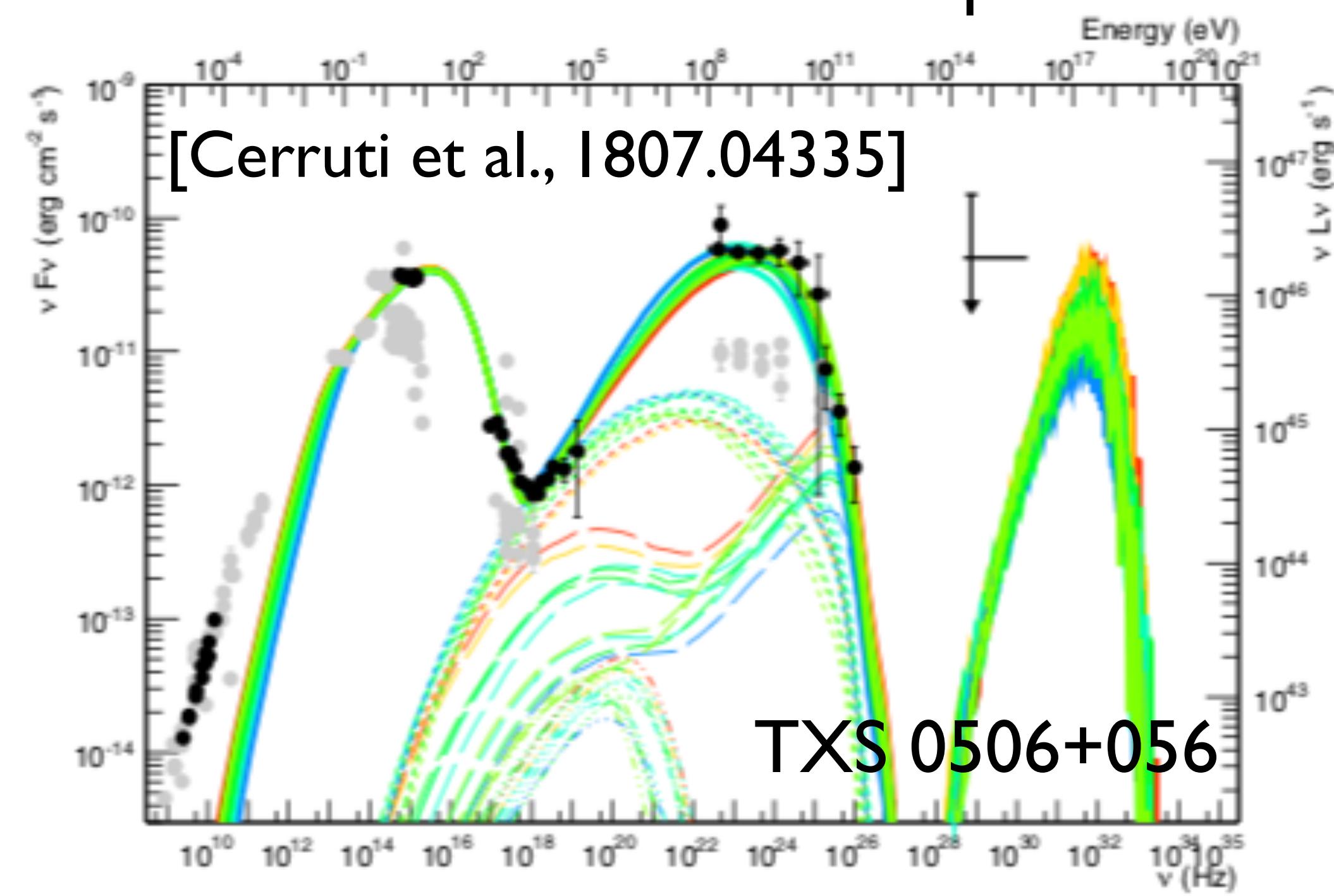
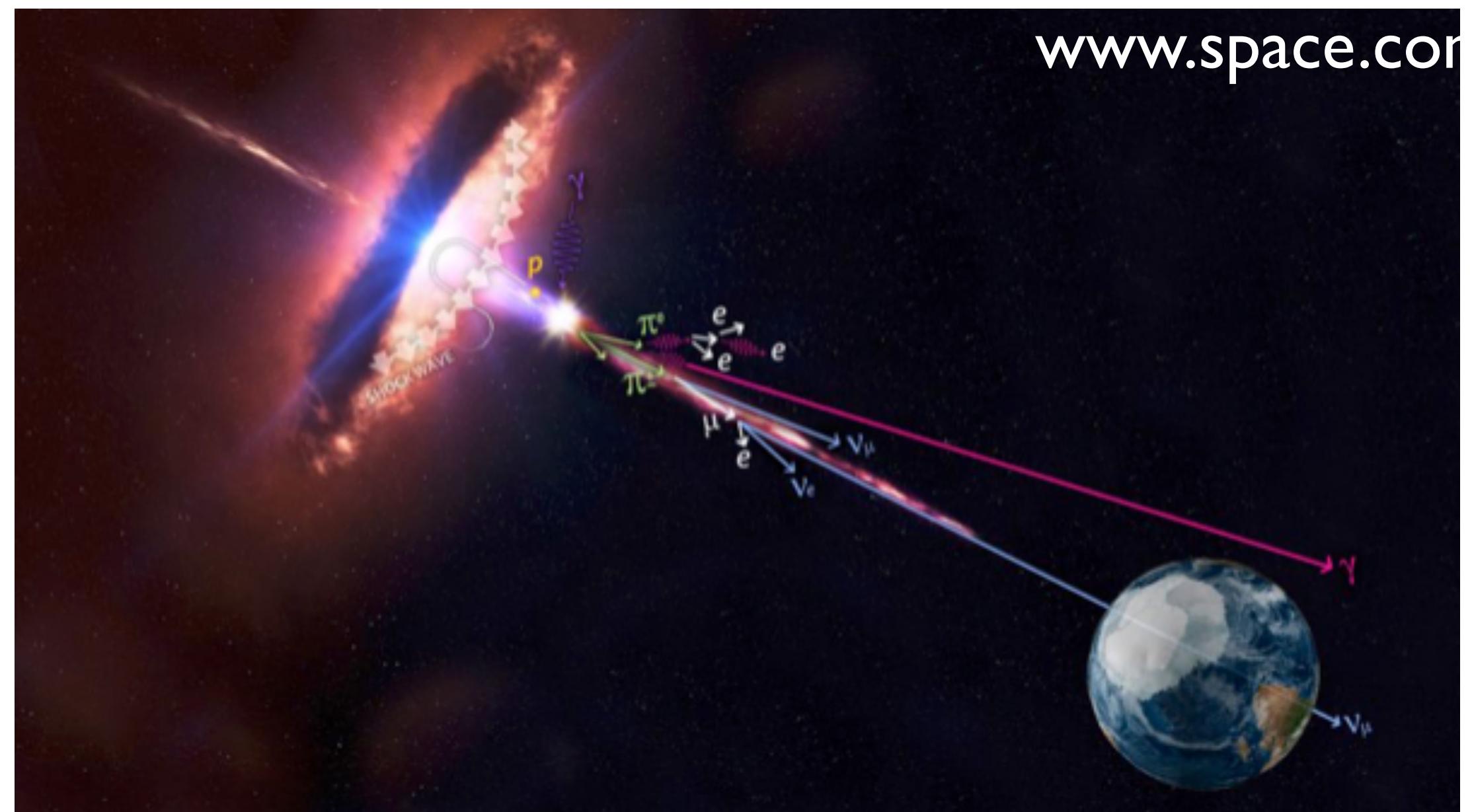


Assuming form factors in
the Earth attenuation and taking into
account inelastic scattering at high
recoil energies.

Blazar boosted DM

[Granelli et al., 2111.13644]

Blazars are AGNs pointing to Earth.
The BH at their center accrete
ordinary matter and focus DM.
They emit a back to back paired jet
of relativistic electrons and protons.



Blazar boosted DM

Blazar boosted dark matter at Earth:

$$\Sigma_{\text{DM}} \equiv \int_{r_{\min}}^r \rho_{\text{DM}}(r') dr'$$

Differential cross section

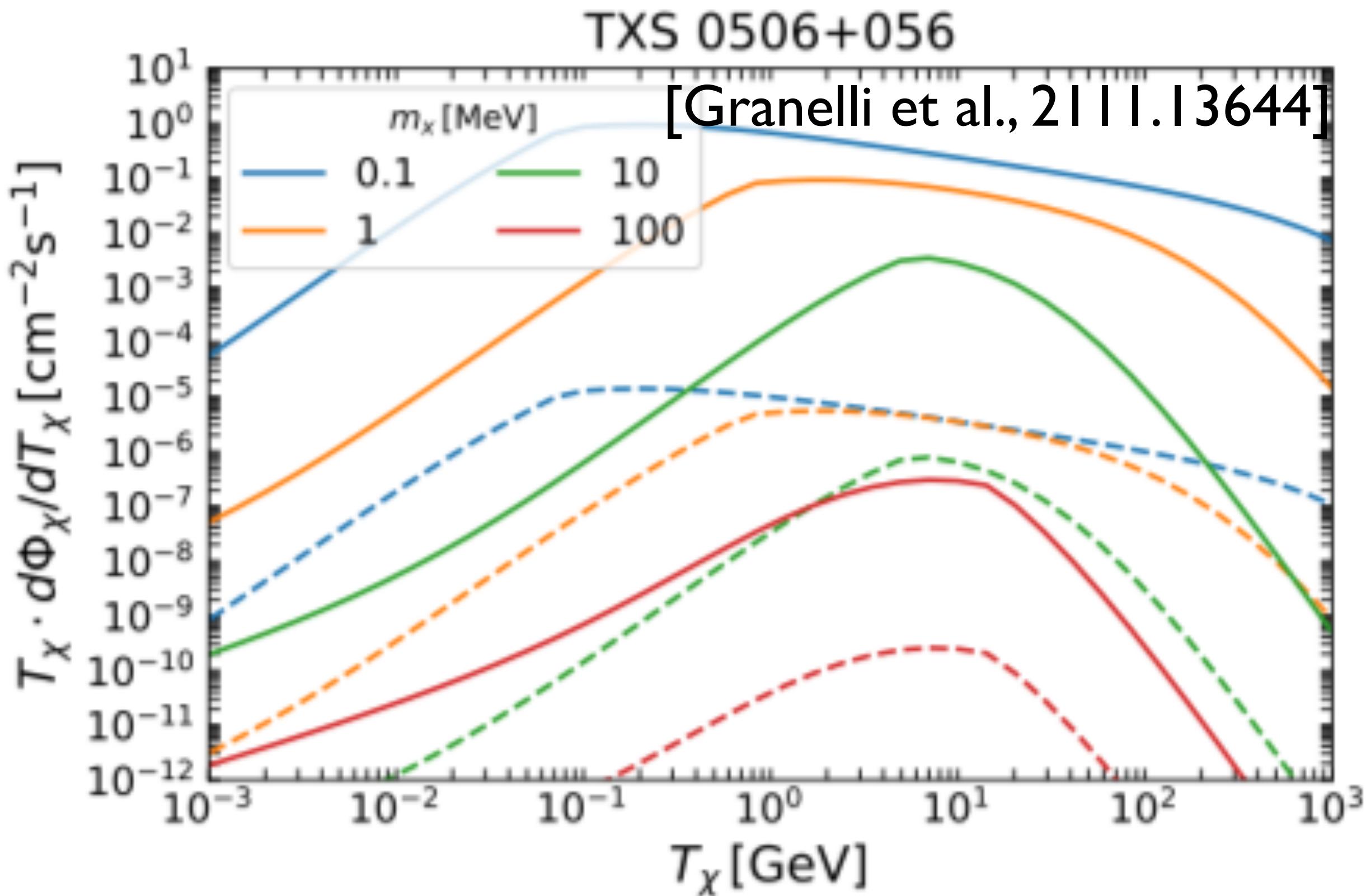
$$\frac{d\Phi_\chi}{dT_\chi} = \frac{\Sigma_{\text{DM}}^{\text{tot}}}{m_\chi d_L^2} \int_{T_j^{\min}}^{T_j^{\max}}$$

$$dT_j \frac{d\sigma_{\chi p}}{dT_\chi} \frac{d\Gamma_j}{dT_j d\Omega}$$

proton spectrum

Luminosity distance

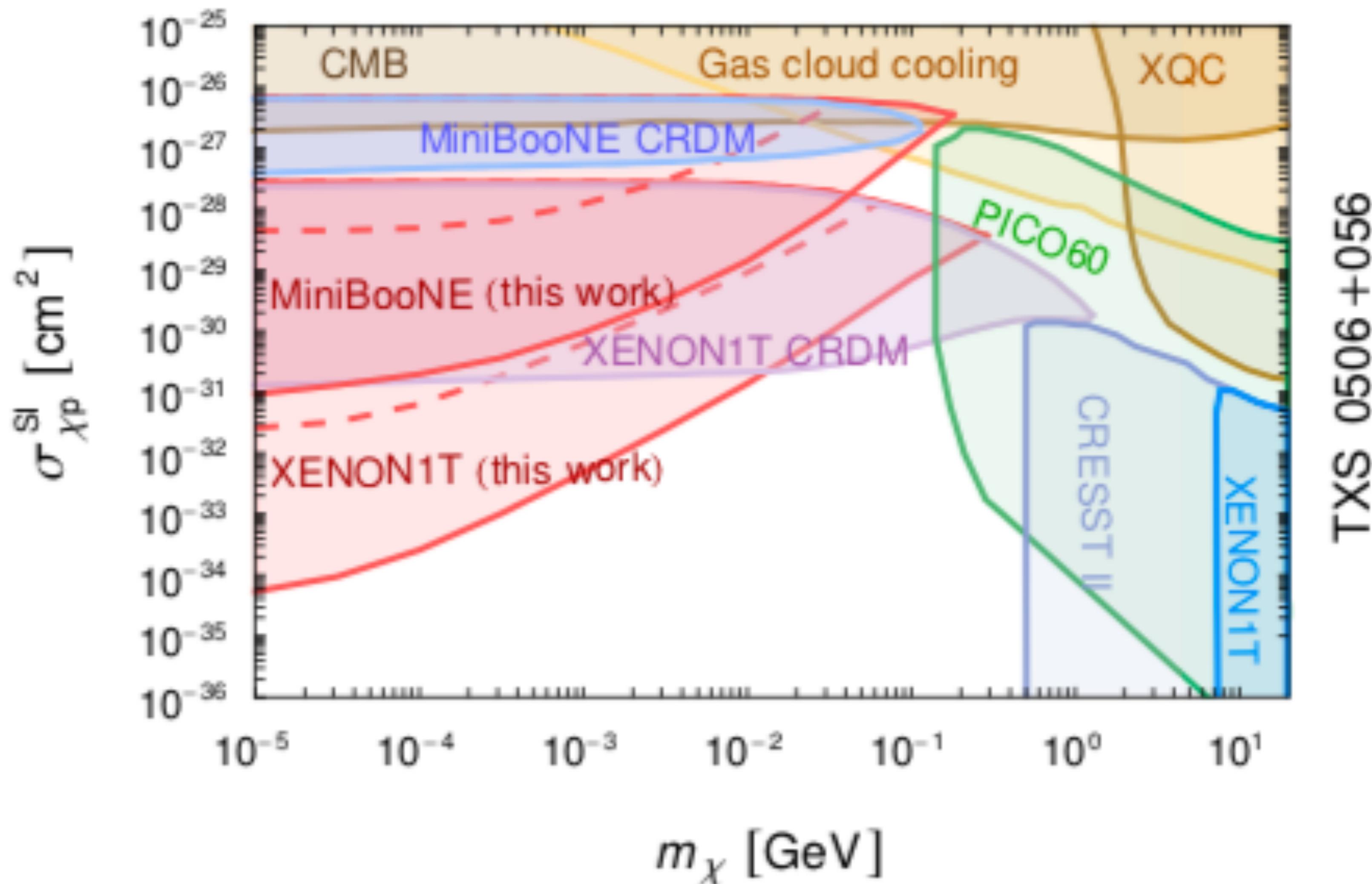
Blazar boosted DM



Different assumptions on the DM profile:

- $\langle \sigma v \rangle_0 = 0$
- $\langle \sigma v \rangle_0 = 10^{-28} \text{ cm}^3 \text{ s}^{-1}$

Blazar boosted DM



Conclusions

Conclusions

- Many hints of the DM existence but where to look for it? We should look everywhere.
- Below the GeV standard direct detection searches are insensitive.
- For the MeV- GeV range one needs to rely on alternative signals: Migdal effect and boosted DM.
- What about \lesssim MeV range? See Angelo's talk.