New Experimental Results on GPDs and Perspectives at JLab with a Positron Beam

GPDs @ Ce+BAF

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(i)

(ii)

(v)





Generalized parton distributions Deeply virtual Compton scattering (iii) Future measurements (iv) Perspectives Ce⁺BAF

EINN 2023 ~ XVth European Research Conference on Electromagnetic Interactions with Nucleons and Nuclei





Parton Promenade

D. Müller, D. Robaschik, B. Geyer, F.M. Dittes, J. Horejsi, FP 42 (1994) 101 X. Ji, PRD 55 (1997) 7114 A. Radyushkin, PRD 56 (1997) 5524

 Generalized Parton Distributions (GPDs) encode the correlations between partons and contain information about the internal dynamics of hadrons which express in properties like the angular momentum or the distribution of the forces experienced by quarks and gluons inside hadrons.







Nucleon Femtography

M. Burkardt PRD 62 (2000) 071503 M. Diehl EPJC 25 (2002) 223 J.P. Ralston, B. Pire PRD 66 (2002) 111501 A.V. Belitsky, D. Müller, NPA 711 (2002) 118c

$$\rho_{H}^{q}(x, \boldsymbol{b}_{\perp}) = \int \frac{d^{2} \boldsymbol{\Delta}_{\perp}}{(2\pi)^{2}} e^{i\boldsymbol{b}_{\perp} \cdot \boldsymbol{\Delta}_{\perp}} \left[H^{q}(x, 0, -\Delta_{\perp}^{2}) + H^{q}(-x, 0, -\Delta_{\perp}^{2}) \right]$$



- The transverse densities of partons in nucleons and nuclei are related to the transverse momentum transfer $(-\Delta_{\perp}^2)$ dependence of GPDs at zero-skewness.
- The experimental knowledge of the ξ-dependence of GPDs at fixed longitudinal momentum fraction *x* and transverse momentum transfer *t* is required to constrain the zero-skewness extrapolation.





Nucleon Spín X. Ji, PRL 78 (1997) 610



$$\lim_{t \to 0} \int_{-1}^{1} x \left[H^{q}(x,\xi,t) + E^{q}(x,\xi,t) \right] dx = J^{q}$$

- The total angular momentum of partons inside the nucleon can be inferred from the Ji sum rule which involves the forward limit of the second Mellin moment of partons helicity conserving GPDs.
- The *E^q* GPD is a new piece of information which does not have a Deep Inelastic Scattering equivalent.
- The experimental knowledge of the *x* and *t*-dependence of GPDs at fixed skewness ξ is required to access the orbital angular momentum carried by partons inside the nucleon.

M. Mazouz et. al. PRL 99 (2007) 242501 A. Airapetian et al. JHEP 06 (2008) 066





Nucleon Forces

X. Ji, PRL 78 (1997) 610 M.V. Polyakov, PLB 555 (2003) 57 M.V. Polyakov, P. Schweitzer, IJMP A 33 (2018) 1830025



$$\int_{-1}^{1} x \sum_{q} H^{q}(x,\xi,t) \ dx = M_{2}(t) + \frac{4}{5}\xi^{2}d_{1}(t)$$

 The skewness dependence of the second Mellin moment of the GPD H accesses to the Gravitational Form Factors (GFFs) of the energy momentum tensor of the nucleon.

$$\mathfrak{N}e[\mathcal{H}(\xi,t)] \stackrel{\text{\tiny LO}}{=} D(t) + \mathcal{P}\left\{\int_{-1}^{1} \left[\frac{1}{\xi-x} - \frac{1}{\xi+x}\right] \mathfrak{I}m[\mathcal{H}(x,t)] dx\right\}$$

- Dispersion relations provide another access to the D-term GFF via the real and imaginary parts of the Compton Form Factors (CFFS).
- The experimental knowledge of the x-, ξ and t-dependence is required to access the GFFs of the nucleon.

V. Burkert, L. Elouadrhiri, F.-X. Girod, Nat. 557 (2018) 396; arXiv:nucl-ex:2104.02031; arXiv:hep-ph:2310.11568 K. Kumerički, Nature 570 (2019) E1





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Experimental Access

 Several GPDs parameterize the non-perturbative structure of the nucleon, but they are universal and can be accessed in different processes, as long as the factorization of the reaction amplitude is preserved.



Factorization is **fully proven** for **DVCS/TCS/DDVCS** ($-t \ll Q^2, Q^2 \gg M^2$) and **partially proven** for **DVMP** at leading order for longitudinally polarized photons.





Experimental Access

 Several GPDs parameterize the non-perturbative structure of the nucleon, but they are universal and can be accessed in different processes, as long as the factorization of the reaction amplitude is preserved.



- GPDs in **blue** reduce to usual PDFs in the forward limit.
- GPDs in **red** vanish if there is no quark orbital angular momentum.
- **Transversity** GPDs are chiral odd and other GPDs are chiral even.







M. Diehl at the CLAS12 European Workshop, Genova, February 25-28, 2009



$$d^{5}\sigma_{PS}^{e} = d^{5}\sigma_{P0}^{e} + S\left[P d^{5}\Delta\sigma_{BH} + (Pd^{5}\Delta\sigma_{DVCS} + d^{5}\Delta\tilde{\sigma}_{DVCS}) - e(Pd^{5}\Delta\sigma_{INT} + d^{5}\Delta\tilde{\sigma}_{INT})\right]$$

Polarized electrons and positrons allow to separate the unknown amplitudes of the cross section for electro-production of photons.





Deeply Vírtual Compton Scatteríng

M. Diehl at the CLAS12 European Workshop, Genova, February 25-28, 2009



Polarized electrons and positrons allow to separate the unknown amplitudes of the cross section for electro-production of photons.





Compton Form Factors



GPDs enter the epγ cross section via the Compton Form Factors (CFFs) representing an integral over the intermediate quark longitudinal momentum.

$$d^{5}\sigma \propto \int_{-1}^{+1} dx \ \frac{\text{GPD}(x,\xi,t)}{x \pm \xi \mp i\epsilon} = \mathcal{P} \int_{-1}^{+1} dx \ \frac{\text{GPD}(x,\xi,t)}{x \pm \xi} \pm i\pi \ \text{GPD}(x = \pm \xi,\xi,t)$$
$$\tilde{\sigma}_{INT,DVCS}$$

> At twist-2 and leading α_{QCD} -order, the epy reaction accesses the four chiral even and parton helicity conserving GPDs $\{H, \tilde{H}, E, \tilde{E}\}$ of the proton via the CFFs $\{\mathcal{H}, \mathcal{H}, \mathcal{E}, \tilde{\mathcal{E}}\}$.

$$\mathcal{C}^{DVCS} = 4(1-x_B) \left[\frac{\mathcal{H}\mathcal{H}^*}{\mathcal{H}^*} + \tilde{\mathcal{H}}\tilde{\mathcal{H}}^* \right] - x_B^2 \left[\mathcal{H}\mathcal{E}^* + \mathcal{E}\mathcal{H}^* + \tilde{\mathcal{H}}\tilde{\mathcal{E}}^* + \tilde{\mathcal{E}}\tilde{\mathcal{H}}^* \right] - \left(x_B^2 + (2-x_B)^2 \frac{t}{4M^2} \right) \mathcal{E}\mathcal{E}^* - x_B^2 \frac{t}{4M^2} \tilde{\mathcal{E}}\tilde{\mathcal{E}}^*$$

$$\mathcal{C}^{INT} = F_1 \mathcal{H} - \xi [F_1 + F_2] \widetilde{\mathcal{H}} - \frac{t}{4M^2} \mathcal{E}$$

Importance of the **separation** of the **DVCS** and **INT** reaction amplitudes for the **determination** of **CFFs**.





---- {++,0+,-+} CFFs fit

{++} CFFs fit KM15

DVCS Cross Section

(Hall A Collaboration) F. Georges et al. PRL 128 (2022) 252002

- First 12 GeV era results exploring DVCS at high $Q^2(2.7-8.4 \text{ GeV}^2)$ and $x_B(0.36-0.60)$.
- Analysis of the cross sections in the NLO+HT scenario using 3 different beam energies.
- Determination of the 4 helicity-conserving CFFs.







DVCS Beam Spin Asymmetry

(CLAS Collaboration) G. Christiaens et al. PRL 130 (2023) 211902

- A new set of BSAs extending over 64 bins in (Q^2, x_B, t) and exploring an unmeasured large phase space.
- More statistics to come as these results corresponds to 25% of the allocated beam-time.



K. Kumerički, D. Müller, EPJ Web Conf. 112 (2016) 01012 B. Berthou et al. EPJ C 78 (2018) 478 H. Moutarde, P. Sznadjer, J. Wagner, EPJ C 79 (2019) 614





GPD

Tímelíke Compton Scatteríng

(CLAS Collaboration) P. Chatagnon et al. PRL 127 (2021) 262501

- First ever measurement of this channel conjugate to DVCS and testing GPDs universality.
- The asymmetry with respect to the orientation of circular polarization of the photon accesses the imaginary part of the CFFs, similarly to the BSA.
- The forward-backward asymmetry accesses the real part of the CFFs, sensitive to the D-term.







Transversity GPDs

(Hall A Collaboration) M. Dlamini et al. PRL 127 (2021) 152301 (CLAS Collaboration) A. Kim et al. arXiv:2307.07874

$$\frac{d\sigma}{dt} = \frac{d\sigma_T}{dt} + \varepsilon \, \frac{d\sigma_L}{dt} + \sqrt{2\varepsilon(1+\epsilon)} \, \frac{d\sigma_{LT}}{dt} \cos(\phi) + \varepsilon \frac{d\sigma_{TT}}{dt} \cos(2\phi) + h \sqrt{2\varepsilon(1-\epsilon)} \frac{d\sigma_{LT'}}{dt} \sin(\phi)$$

 $d\sigma_{IT}$

0.3

0.25

1 1.2 1.4 1.6 1.8 -t [GeV²]

 $d\sigma_{LT'}$

GK calculations

0.2

JML

S.V. Goloskokov, P. Kroll, EPJ A 47 (2011) 112



 π^{0} -electroproduction accesses the partonic content of the nucleon through the coupling of twist-3 distribution amplitudes with transversity GPDs.

J.-M. Laget PLB 695 (2011) 199

$$\begin{split} d\sigma_T &\propto (1-\xi^2) |\mathcal{H}_T|^2 - \frac{t'}{8m^2} \left| 2\widetilde{\mathcal{H}}_T + \mathcal{E}_T \right|^2 \\ d\sigma_{TT} &\propto \left| 2\widetilde{\mathcal{H}}_T + \mathcal{E}_T \right|^2 \\ d\sigma_{LT'} &\propto \Im m \big[\big(2\widetilde{\mathcal{H}}_T^* + \mathcal{E}_T^* \big) \widetilde{\mathcal{H}} + \mathcal{H}_T^* \widetilde{\mathcal{E}} \big] \end{split}$$

- π^{0} -electroproduction is confirmed to be almost purely transverse.
- Experimental data are fairly described by a GPD based model (GK).





Transition GPDs

(CLAS Collaboration) S. Diehl et al. PRL 131 (2023) 021901

- o Investigation of a novel DVCS channel experimentally unexplored.
- The partonic content of the p→∆ transition is parametrerized by 8 helicity conserving transition GPDs and 8 helicity-flip transition GPDs.
- Access to transition GPDs via the BSA.
- On-going theoretical work.



P. Kroll, K. Passek-Kumerički, PRD 107 (2023) 054009 K.M. Semenov-Tian-Shansky, M. Vanderhaeghen PRD 108 (2023) 034021







More Experiments

Analyses ín progress

- · Beam spin asymmetries off the neutron
- Beam spin asymmetries off the bound proton in deuterium
- Longitudinal target asymmetries off proton and neutron
- DVCS cross sections off the proton from Hall B
- New DVCS and π^0, ρ^0 -electroproduction data from COMPASS

Short & Míd term Data Takíng







- o DVCS cross sections off the proton and the neutron with the Neutral Particle Spectrometer (NPS) in Hall C
- $\circ~$ Coherent and incoherent DVCS beam spin asymmetries off helium with ALERT in Hall B

Talks of A. Hobart, S. Niccolai, C. Riedl





LoI12-23-003/LoI12-23-014 M. Defurne et al.

O. Bessidskaia Bylund, M. Defurne, P.AM. Guichon PRD 107 (2023) 014020

- The polarization of the final state particles in the DVCS process provides a novel access to GPDs:
 - the recoil proton polarization component P_y normal to the scattering plane is sensitive to **E**
 - the recoil proton polarization component P_x transverse to the proton momentum is sensitive to \tilde{H}
 - the linear polarization of the photon P_l is sensitive to the gluon transversity GPDs.





Polarimeters ask for high luminosities which can only be accessed at JLab.





Double Deeply Vírtual Compton Scattering

M. Guidal, M. Vanderhaeghen, PRL 90 (2003) 012001 A.V. Belitsky, D. Müller PRL 90 (2003) 022001; PRD 68 (2003) 116005

• Because of the virtuality of the final photon, DDVCS allows a direct access to GPDs at $x \neq \pm \xi$, which is of importance for their modeling and for the investigation of nuclear dynamics.

$$\mathcal{F}(\xi',\xi,t) = \mathcal{P} \int_{-1}^{1} dx \, F_{+}(x,\xi,t) \left[\frac{1}{x-\xi'} \pm \frac{1}{x+\xi'} \right] - i\pi F_{+}(\xi',\xi,t) \qquad F_{+}(x,\xi,t) = \sum_{q} \left(\frac{e_{q}}{e} \right)^{2} \left[F^{q}(x,\xi,t) \mp F^{q}(-x,\xi,t) \right]$$

$$\mathcal{F}_{+}(x,\xi,t) = \sum_{q} \left(\frac{e_{q}}{e} \right)^{2} \left[F^{q}(x,\xi,t) \mp F^{q}(-x,\xi,t) \right]$$

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sign, providing a testing ground of **GPDs universality**.

I.V. Anikin et. al. APPB 49 (2018) 741

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LoI12-16-004 S. Stepanyan et al.

A modified CLAS12 spectrometer featuring a new tracker, a new calorimeter and a tungsten shield in the 0 HTCC volume region would allow to increase CLAS12 luminosity up to 10³⁷cm⁻²·s⁻¹.



- The new detection system serves electron identification
- The CLAS12 forward detector acts as a muon detector
- This reconfigured CLAS12 spectrometer would measure DDVCS and J/Ψ electroproduction over a 100 days run.

350





LoI12-23-012 M. Boer, A. Camsonne, E. Voutier, Z. Zhao et al.

S. Zhao et al. EPJ A 57 (2021) 240

The SoLID apparatus completed with a muon detector at forward and large angles, enables DDVCS measurements with both polarized electron and positron beams at 11 GeV.





- The initial LoI discussed electron BSA measurements over a 50 days run parasitic to the J/ Ψ approved experiment.
- Completing this program with a **50 days positron beam** run would provide **unpolarized BCA** data.





Impact of Posítron Measurements

H. Dutrieux, V. Bertone, H. Moutarde, P. Sznajder, EPJ A 57 (2021) 300

- The existing DVCS world data set (H1, ZEUS, HERMES, JLab 6 GeV, COMPASS) is analyzed within a global fit based on an Artificial Neural Network procedure within PARTONS to extract CFFs.
- The impact of projected CLAS12 BCA data on the proton is evaluated from a Bayesian reweighting analysis of CFFs.







PR12+23-002 E. Voutier, V. Burkert, S. Niccolai, R. Paremuzyan et al.

V. Burkert et al. EPJ A 57 (2021) 186

• Measurements of beam charge asymmetries with CLAS12 will provide a full set of new GPD observables:

- the unpolarized beam charge asymmetry A_{UU}^{C} , sensitive to the CFF real part;
- the polarized beam charge asymmetry A_{LU}^{C} , sensitive to the CFF imaginary part;
- the charge averaged beam spin asymmetry A_{LU}^0 , signature of higher twist effects.







PR12+23-006 C. Muñoz Camacho, M. Mazouz et al.

A. Afanasev et al. EPJ A 57 (2021) 300

Combining the HMS and the NPS spectrometers, precise cross section measurements with unpolarized 0 positron beam are proposed at selected kinematics where electron beam data will soon be accumulated.



 $x_B = 0.36$ $Q^2 = 4.0 \text{ GeV}^2$





L0I12-18-004 Jefferson Lab Positron Working Group

S. Niccolai, P. Chatagnon, M. Hoballah, D. Marchand, C. Muñoz Camacho, E. Voutier, EPJ A 57 (2021) 226

 Contrary to H, the GPD E flips the spin of the nucleon and is consequently not constrained by Deep Inelastic Scattering data.



X. Ji, PRL 78 (1997) 610
$$\int_{-1}^{1} x \left[H(x,\xi,t \to 0) + E(x,\xi,t \to 0) \right] dx = J$$

$$A_{UU}^{C} \propto \frac{1}{F_{2}} \Re e \left[\xi \widetilde{H}_{n} - \frac{t}{4M^{2}} E_{n} \right]$$

The **BCA** on the neutron accesses the **real part** of the CFF **E**, and is sensitive to \tilde{H} at some kinematics.





Coherent He-DVCS

M. Rinaldi, S. Scopetta, PRC 85 (2012) 062201; PRC 87 (2013) 035208 S. Fucini et al. EPJ A 57 (2021) 273

 The association of the ALERT recoil detector and the CLAS12 spectrometer together with high-energy electron and positron beams offer a new tool for the investigation of the nuclear force.

$$\Re e[\mathcal{H}_{A}(\xi,t)] = \mathcal{P}\int_{0}^{1} \left[\frac{1}{\xi+x} + \frac{1}{x-\xi}\right] H_{A}(x,\xi,t) \, dx = \frac{1}{\pi} \mathcal{P}\int_{0}^{1} \left[\frac{1}{\xi+x} + \frac{1}{x-\xi}\right] \Im m[H_{A}(x,x,t)] \, dx - \delta(t)$$









Posítron Working Group



- The JLab Positron Working Group (PWG) developed the perspectives of an experimental program with positron beams at CEBAF in a topical EPJ A issue.
- This document constitutes the final JLab Positron White Paper, gathering 19 single contributions and a summary article, all peer-reviewed.

JLab PWG = \sim 250 Physicists from 75 Institutions and 16 countries

(Jefferson Lab Positron Working Group) A. Accardi et al. EPJ A 57 (2021) 261

Subscribe to the JLab Positron Working Group mailing list pwg@jlab.org





Jefferson Lab PAC51

 The Positron Experimental Program at JLab has formally started with the C1 approval of 5 positron proposals at the PAC meeting of Juy 2023, constituting 3 calendar years of single hall running.

NUMBER	TITLE	PHYSICS THEME	CONTACT PERSON	HALL	DAYS AWARDED	SCIENTIFIC RATING	PAC DECISION
PR12+23-002	Beam Charge Asymmetries for Deeply Virtual Compton Scattering on the Proton at CLAS12	GPDs	Eric Voutier	В	100	A -	C1
PR12+23-003	Measurement of Deep Inelastic Scattering from Nuclei with Electron and Positron Beams to Constrain the Impact of Coulomb Corrections in DIS	TPE	Dave Gaskell	С	9.3	A-	C1
PR12+23-006	Deeply Virtual Compton Scattering using a positron beam in Hall C	GPDs	Carlos Muñoz Camacho	С	137	A -	C1
PR12+23-008	A Direct Measurement of Hard Two-Photon Exchange with Electrons and Positrons at CLAS12	TPE	Axel Schmidt	В	55	Α	C1
PR12+23-012	A measurement of two-photon exchange in unpolarized elastic positron–proton and electron–proton scattering	TPE	Michael Nycz	С	56	A -	C1

C1 = Conditionally Approved with Technical Review by the Lab





e+@LERF

(Ce⁺BAF Working Group) J. Grames *et al.* JACoW IPAC2023 (2023) MOPL152; arXiv2309.15581

Taking advantage of the existing infrastructure (electric and cryogenic power supplies, shielding...), a new positron injector based on the PEPPo (Polarized Electrons for Polarized Positrons) concept will be installed at the LERF.







e⁺@CEBAJ

(Ce⁺BAF Working Group) J. Grames *et al.* JACoW IPAC2023 (2023) MOPL152; arXiv2309.15581

 A new beam transport line attached to the ceiling of the existing tunnel will guide the 123 MeV e⁺ beam till the injection point at the entrance of the North LinAc.







Conclusion

- Thirty years after their theoretical formulation, the experimental landscape of GPDs in the valence region is rich and improving rapidly.
- New promising channels and observables have been identified and will become available in a short future.
- Polarized and unpolarized positron beams at Jefferson Lab would uniquely enlarge the physics reach of GPDs at CEBAF.
- A strong accelerator R&D effort is progressing towards the final design and implementation of positron beams at CEBAF.

Merci de Votre Attention





Rosenbluth-like DVCS separation

(Hall A Collaboration) M. Defurne et al. Nat. Com. 8 (2017) 1408



 The DVCS and INT contributions to the unpolarized epγ cross section are separated following on the one hand a scenario including Next-to-Leading Order (NLO) contributions, and on the other hand Higher-Twists effects (HT).



The **DVCS** and **INT** amplitudes signifcantly differ in each scenario.





PR12+23-006 C. Muñoz Camacho, M. Mazouz et al.

A. Afanasev et al. EPJ A 57 (2021) 300

 The combination of electron and positron measurements strongly reduce correlations between CFFs, from 94% without to 39% at the considered kinematics.











- Add a **new fast-tracking detector** in between HTCC and DC regions
- Remove HTCC and install a Møller cone in tungsten
- Add a new PbWO4 calorimeter to cover the 7°-30° region together with a thick tungsten shield protecting the full FD region
- Install a **new Micro Pattern Gas Detector** in front of the calorimeter

October 31st-November 4th, 2023





Posítron Injector

S. Habet, Y. Roblin et al. JACoW IPAC (2022) 457 S. Habet, Doctorate, Université Paris-Saclay, 12/2023

• The design of the JLab positron source evolved towards the today's latest concept :



High duty cycle, intensity, and polarization distinguish JLab positron beam from any past or existing others.





Ce⁺BAF Working Group

J. Benesch, A. Bogacz, L. Cardman, J. Conway, S. Covrig, <u>J. Grames</u>, J. Gubeli, C. Gulliford, S. Habet, C. Hernandez-Garcia, D. Higinbotham, A. Hofler, R. Kazimi, V. Kostroun, F. Lin, V. Lizarraga-Rubio, M. Poelker, Y. Roblin, A. Seryi, K. Smolenski, M. Spata, R. Suleiman, A. Sy, D. Turner, A. Ushakov, C. Valerio, E. Voutier, S. Zhang, Y. Zhang

(Ce⁺BAF Working Group) J. Grames et al. JACoW IPAC2023 (2023) MOPL152

S. Habet *et al.* JACoW IPAC2022 (2022) 457 R. Kazimi *et al.* JACoW IPAC2023 (2023) WEPA035 A. Sy *et al.* JACoW IPAC2023 (2023) MOPM081 A. Ushakov *et al.* JACoW IPAC2023 (2023) WEPM120







Rough Tímelíne

D. Dean at the International Workshop on CLAS12 Physics and Future Perspectives at JLab, Paris, March 21-24, 2023



- Phase 1 includes building a positron source and the tunnel & beamline connecting to CEBAF
- Phase 2 includes new permanent magnets to allow 22 GeV within current CEBAF footprint





Long Range Plan



 The 2023 NP LRP recommends priorities to US funding agencies for the next 7-8 year cycle.

RECOMMENDATION 1

The highest priority of the nuclear science community is to capitalize on the extraordinary opportunities for scientific discovery made possible by the substantial and sustained investments of the United States. We must draw on the talents of all in the nation to achieve this goal.

RECOMMENDATION 3

We recommend the expeditious completion of the EIC as the highest priority for facility construction.

RECOMMENDATION 2

As the highest priority for new experiment construction, we recommend that the United States lead an international consortium that will undertake a neutrinoless double beta decay campaign, featuring the expeditious construction of ton-scale experiments, using different isotopes and complementary techniques.

RECOMMENDATION 4

We recommend capitalizing on the unique ways in which nuclear physics can advance discovery science and applications for society by investing in additional projects and new strategic opportunities.

Recommendations 1&4 strengthen current CEBAF operations and future upgrades.