

The Δ resonance at different physical parameters

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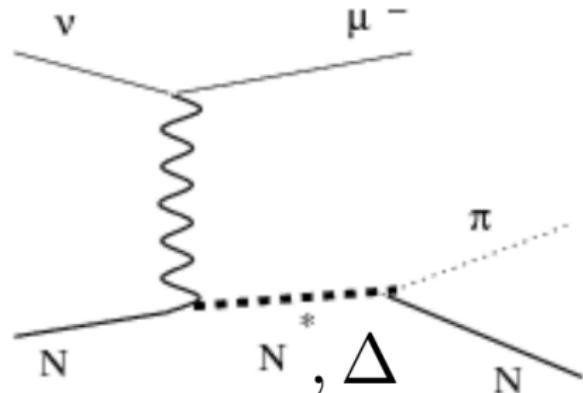
EINN2023

Collaborators

- Stefan Meinel (U Arizona)
- Gumaro Rendon (BNL)
- John W. Negele, Andrew Pochinsky (MIT)
- Luka Leskovec (Jozef Stefan Institute)
- Sergey Syritsyn (RIKEN BNL & Stony Brook U)
- Constantia Alexandrou (Cyl & U Cyprus)
- Li Yan (U Cyprus)
- Simone Bacchio (Cyl)
- Giannis Koutsou (Cyl)
- Marcus Petschlies (U Bonn)
- Srijit Paul (U. Edinburgh)
- Theodoros Leontiou (Frederick University)

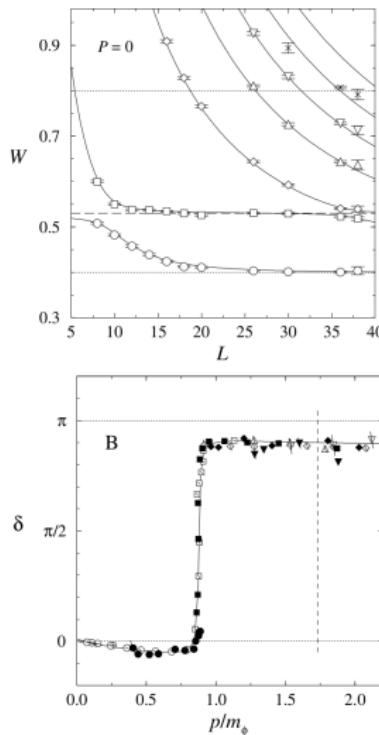
Delta resonance

- $\Delta(1232)$
- Lightest baryonic resonance
- Dominant in the p wave $N\pi$ scattering
- Simplest resonance: 3 quark and $N\pi$ contribution
- Resonances are not eigenstates of the QCD Hamiltonian



- They decay via strong interactions
- Finite volume influences the two-hadron spectrum

Lüscher Method (Nucl.Phys.B 1991)



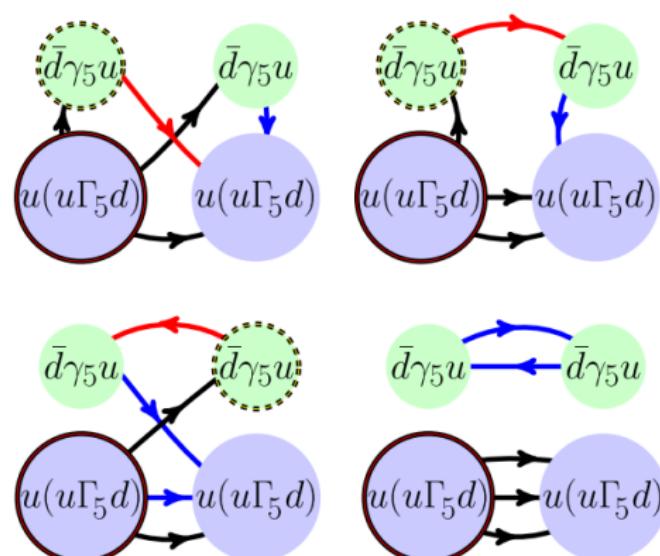
- Turn the finite volume to an advantage
- $\det(F^{-1}(E_{\text{cm}}) - 8\pi iT(E_{\text{cm}})) = 0$
- F is the box function (L), T is scattering amplitude
- E_{cm} can be determined
- We have two Different L -s at two different pion masses
- Physical point: threshold very low

Diagrams: $N - N, N - \pi N, \pi N - \pi N$

2pt correlator

- $C(t) = \langle 0 | O_{\mathcal{X}}(t) \bar{O}_{\mathcal{X}'}(0) | 0 \rangle \quad \mathcal{X}, \mathcal{X}' \in \{N, N\pi\}$

- Simplest case NN is already expensive
- Three point-to-all propagators have to be multiplied. Each has dimensions $\sim \mathcal{O}(10^9)$



- Per lattice site:
- $\epsilon_{a,b,c} \epsilon_{l,m,n} S1_{\alpha\alpha_0}^{c,l} \Gamma_{f\alpha_0,\alpha_1} S2_{\beta_0,\alpha_1}^{b,m} \Gamma_{f\beta_0,\beta_1} S3_{\beta_1,\beta}^{a,n}$
- Ideal task for a GPU-kernel
<https://github.com/cylqcd/PLEGMA>

Finite volume lattice: projections

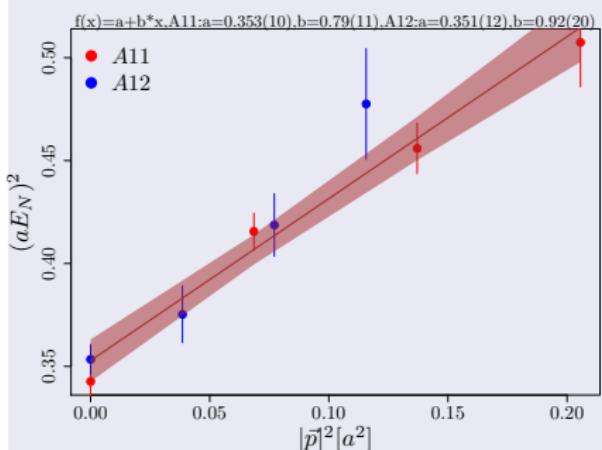
- We use single and two hadron interpolating operators with $I = 3/2, I_3 = 3/2$
- Finite volume we no longer have continuous rotational symmetry
- Symmetry group in the centre-off-mass frame is the double cover of the octahedral group $2O_h$
- Finite number of irreducible representations

| $\frac{L}{2\pi} \vec{P}$ | (0, 0, 0) | (0, 0, 1) | (0, 1, 1) | (1, 1, 1) |
|--|---|--|---|---|
| Group LG | $O_h^{(D)}$ | $C_{4v}^{(D)}$ | $C_{2v}^{(D)}$ | $C_{3v}^{(D)}$ |
| Axis and planes of symmetry | | | | |
| g_{LG} | 96 | 16 | 8 | 12 |
| $\Lambda(J^P) : \pi(0^-)$ | $A_{1u}(0^-, 4^-, \dots)$ | $A_2(0, 1, \dots)$ | $A_2(0, 1, \dots)$ | $A_2(0, 1, \dots)$ |
| $\Lambda(J^P) : N(\frac{1}{2}^+)$ | $G_{1g}(\frac{1}{2}^+, \frac{7}{2}^+, \dots)$ | $G_1(\frac{1}{2}, \frac{3}{2}, \dots)$ | $G(\frac{1}{2}, \frac{3}{2}, \dots)$ | $G(\frac{1}{2}, \frac{3}{2}, \dots)$ |
| $\Lambda(J^P) : \Delta(\frac{3}{2}^+)$ | $H_g(\frac{3}{2}^+, \frac{5}{2}^+, \dots)$ | $G_1(\frac{1}{2}, \frac{3}{2}, \dots) \oplus G_2(\frac{3}{2}, \frac{5}{2}, \dots)$ | $(2)G(\frac{1}{2}, \frac{3}{2}, \dots)$ | $G(\frac{1}{2}, \frac{3}{2}, \dots) \oplus F_1(\frac{3}{2}, \frac{5}{2}, \dots)$ $\oplus F_2(\frac{3}{2}, \frac{5}{2}, \dots)$ |

Simulation details

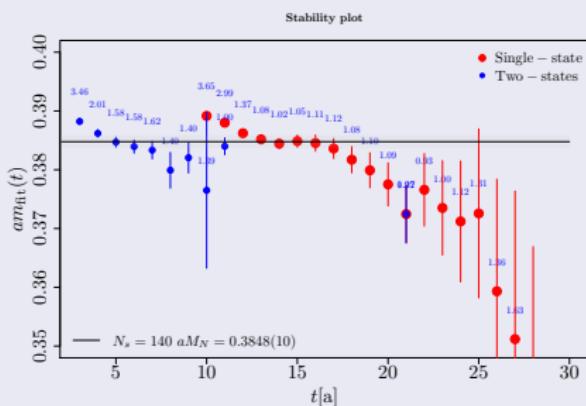
$N_f = 2 + 1$ Clover, $a = 0.1163\text{fm}$

- A11,A12:
 $M_\pi = 200\text{MeV}, L = 2.8 - 3.7\text{fm}$
- A7,A8: $M_\pi = 250\text{MeV}, L = 2.8 - 3.7\text{fm}$
- A15: $M_\pi = 137\text{MeV}, L = 5.5\text{fm}$



$N_f = 2 + 1 + 1$ Twisted-Clover
 $a = 0.08\text{fm}$

- $M_\pi = 139\text{MeV}, L = 5.12\text{fm}$

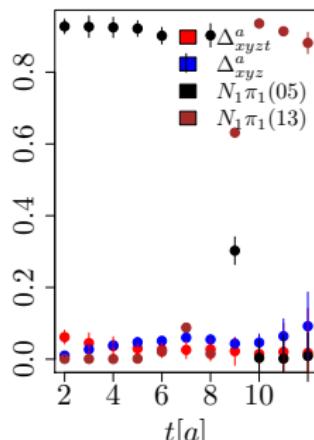
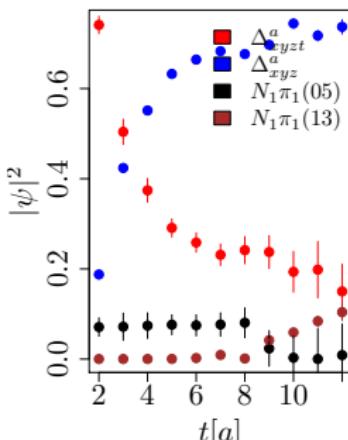


Generalized eigenvalue problem (GEVP)

$$C_{ik}^{\Lambda, \vec{P}}(t) u_k^n(t, t_0) = \lambda^n(t, t_0) C_{ij}(t_0) u_j^n(t)$$

$$\Lambda^n(t, t_0) \propto e^{-E_n^{\Lambda, \vec{P}}(t-t_0)}$$

- Key point: Selecting a basis
- Aim: Robustness, "good" signal quality, eigenvectors



$$\begin{aligned}\langle GS | GS \rangle &= \\ |\psi|_{\Delta_{xyzt}}^2 &+ \\ |\psi|_{\Delta_{xyz}}^2 &+ \\ |\psi|_{N\pi(05)}^2 &+ \\ |\psi|_{N\pi(13)}^2 &\equiv 1\end{aligned}$$

Four different methods

Single state fits

- For each principal correlators of the GEVP $\Lambda^n(t, t_0) \propto e^{-E_n^{\Lambda, \vec{P}}(t-t_0)}$

Hankel

(Fischer et.al. Eur.Phys.J.A(2020))

- For each principal correlators of the GEVP
- $H_{ij}^0(t) = C^0(t + i\Delta + j\Delta)$
- $\sum_{k=0}^{n-1} e^{-E_k t} e^{-E_k i\Delta} e^{-E_k j\Delta} c_k$

AMIAS (Finkenrath et.al. PoS LATTICE2016)

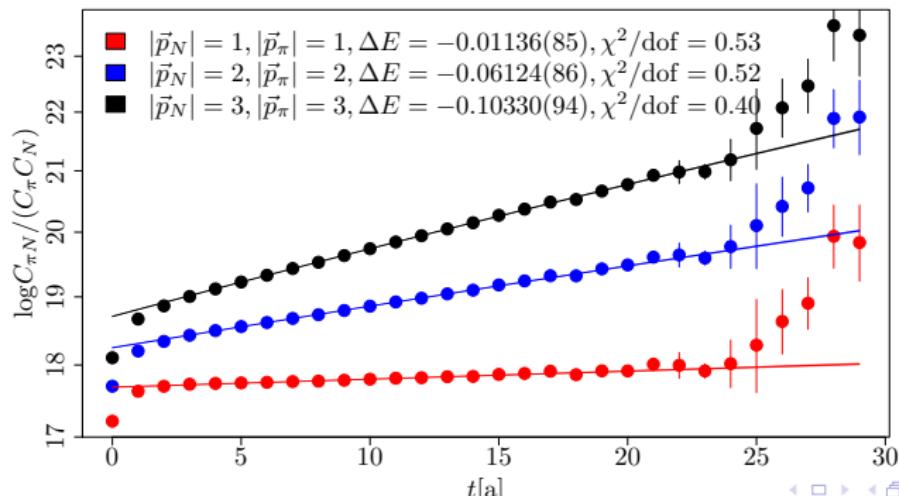
- Statistically sampling the space of fit parameters according to the χ^2 value of the fit function

Ratio method

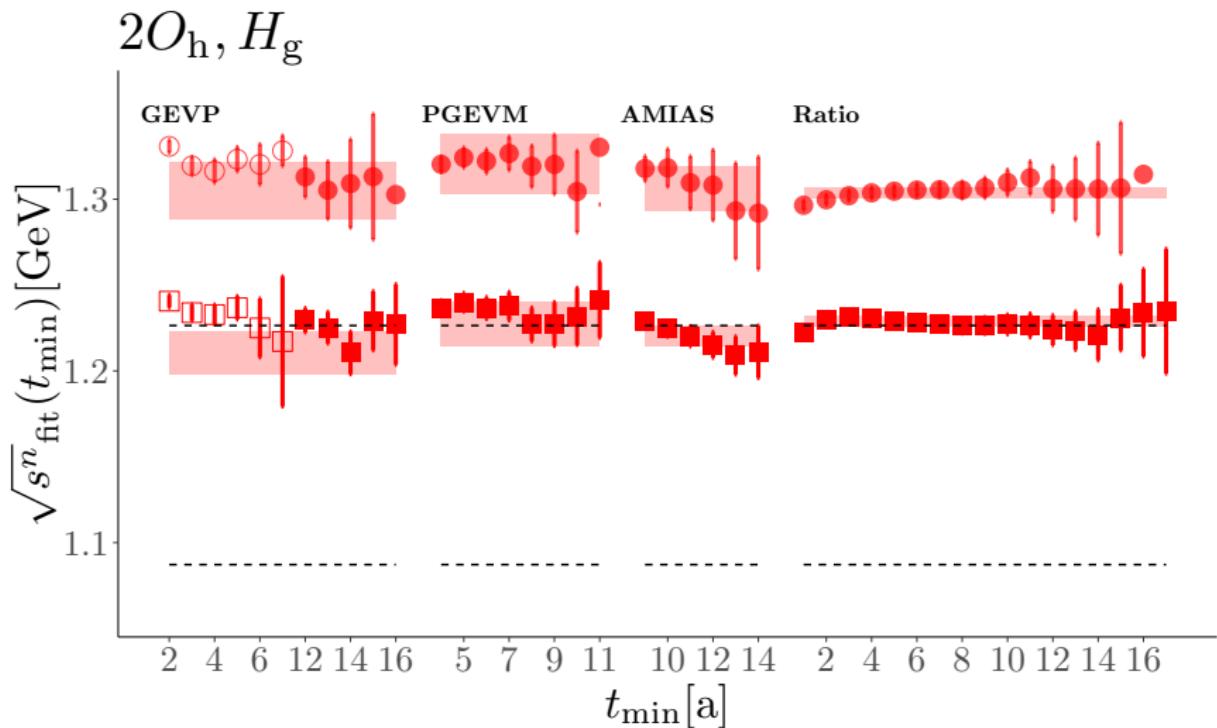
- We fit the energy shift directly

Ratio method

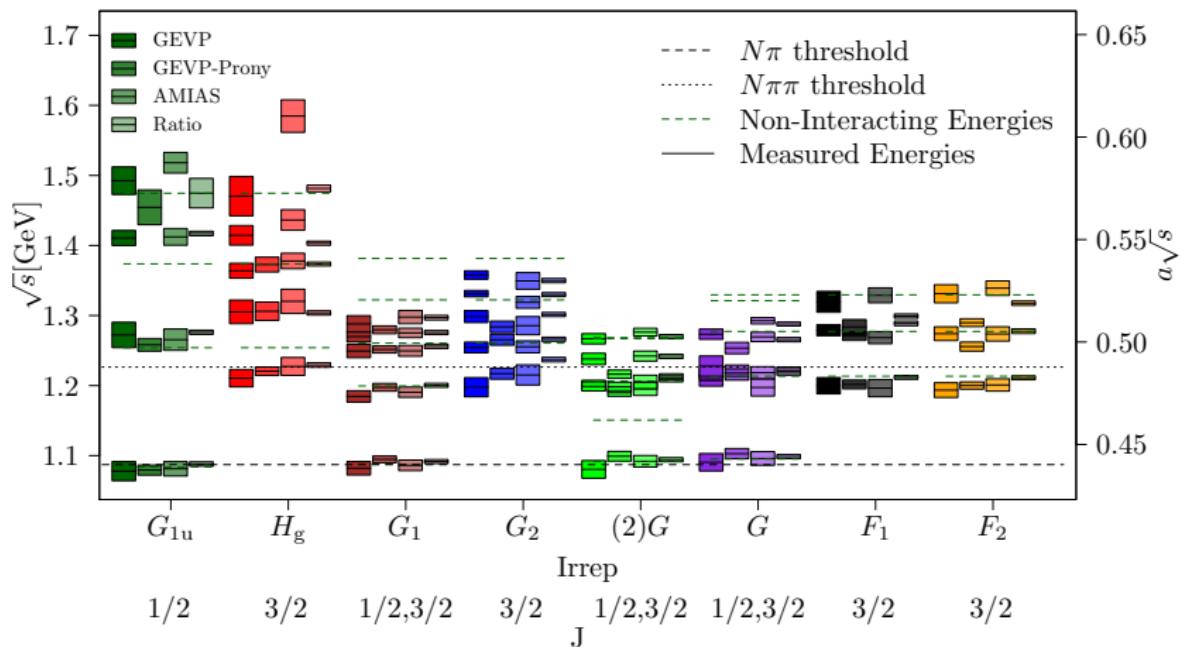
- Single hadron 2pt and two hadron 2pt functions are correlated
- Take the log of the ratio of $C_{\pi N}(t)/(C_N(t)C_\pi(t))$
- We can measure the shift relative to different non-interacting levels



Comparison of different methods

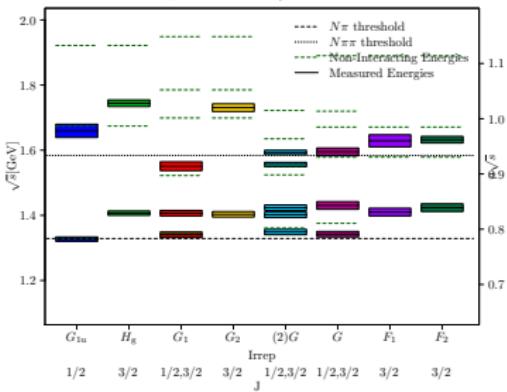


Spectrum summary

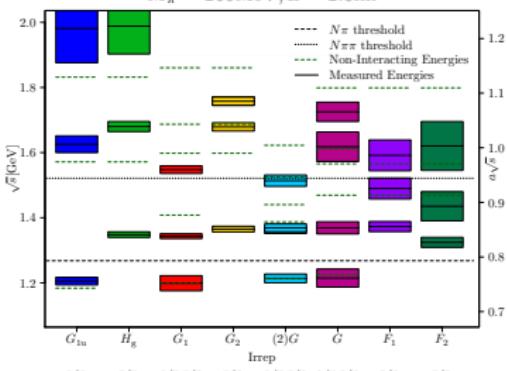


Spectrum summary Clover ensembles

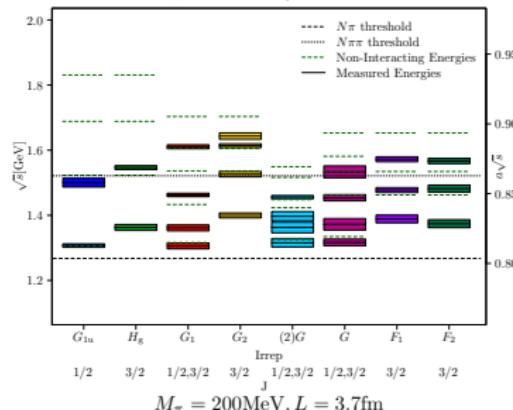
$M_\pi = 250\text{MeV}$, $L = 2.8\text{fm}$



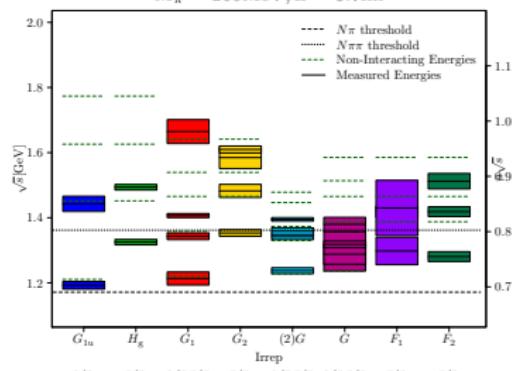
$M_\pi = 200\text{MeV}$, $L = 2.8\text{fm}$



$M_\pi = 250\text{MeV}$, $L = 3.7\text{fm}$



$M_\pi = 200\text{MeV}$, $L = 3.7\text{fm}$



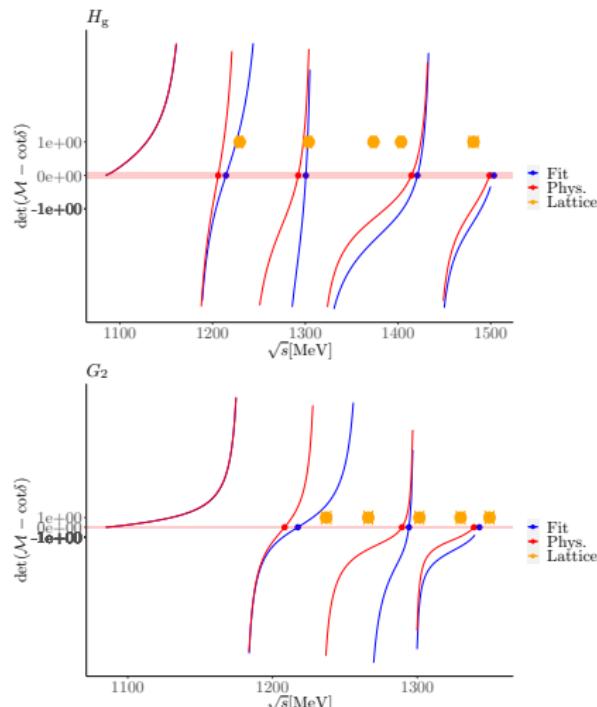
Getting the phase shift

Lüscher-method

- Two particle energy levels in a finite box with size L
- Volume dependence of the energy shift related to scattering observables at $L = \infty$

$$\det \left(\mathcal{M}_{J\ell\mu, J'\ell'\mu'}^P - \delta_{JJ'} \delta_{\ell\ell'} \delta_{\mu\mu'} \cot \delta_{J\ell} \right) = 0$$

- Determinant is taken in angular momentum space
- Important: For $\ell = 1$ dominant irreps there is a one-to-one correspondence between phase-shift and finite volume energy levels (ignoring contributions from higher partial waves)



Parametrization of the resonance

Possible mixing of partial waves

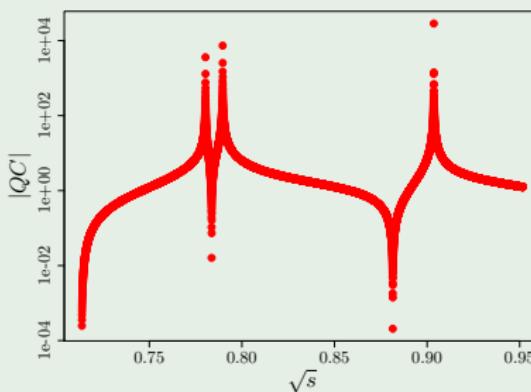
Quantization conditions (QC) Göckeler et. al PRD 2012

- Phase shift parametrization:

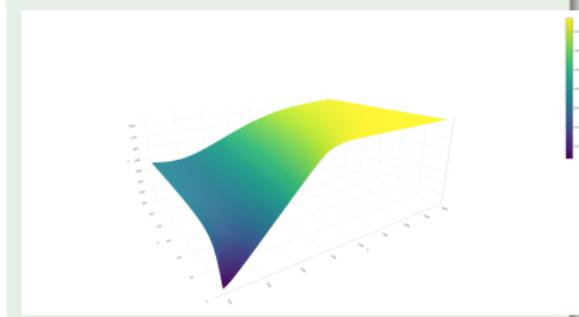
- $\ell = 0 \rightarrow \cot\delta_{\ell=0} = a_0 q_{\text{cmf}}, \quad \ell = 1 \rightarrow \tan\delta_{\ell=1} = \frac{\sqrt{s}\Gamma(R,s)}{M_R^2 - s}$

- We restrict ourselves to $\ell = 0, 1$ and check for $\ell \geq 2$

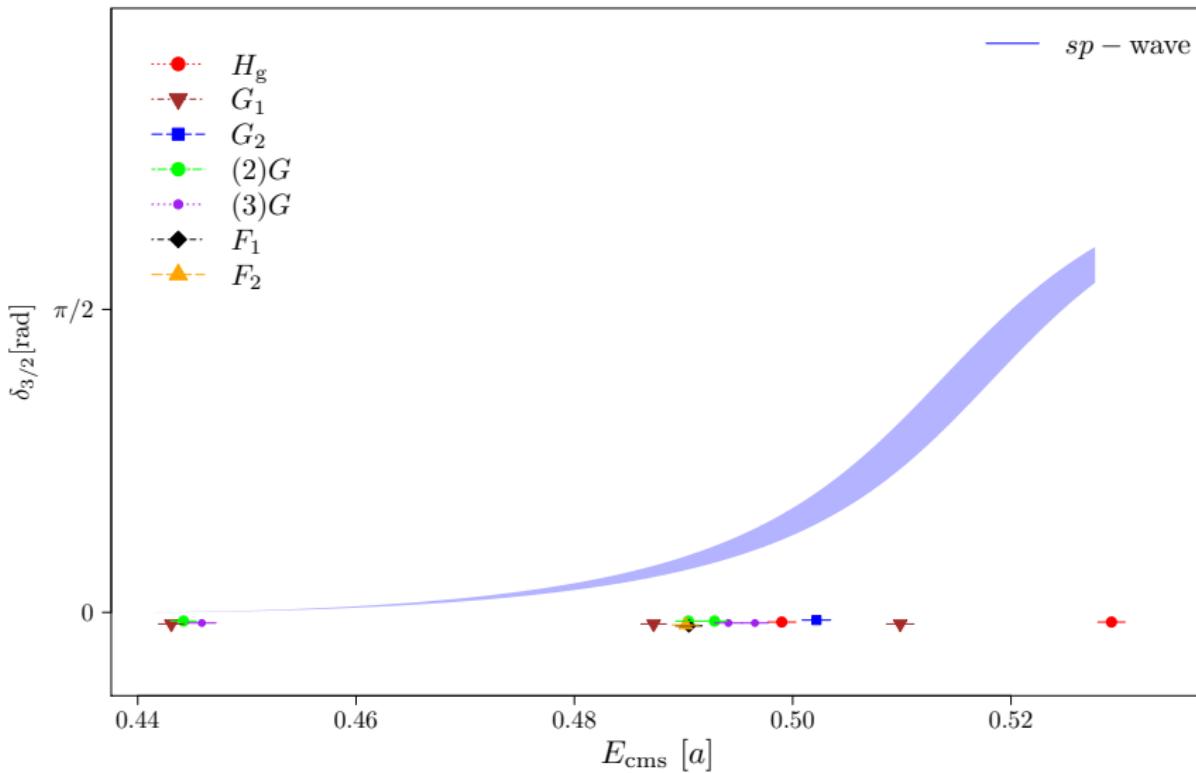
Example



2d Spline interpolation



LQC fit



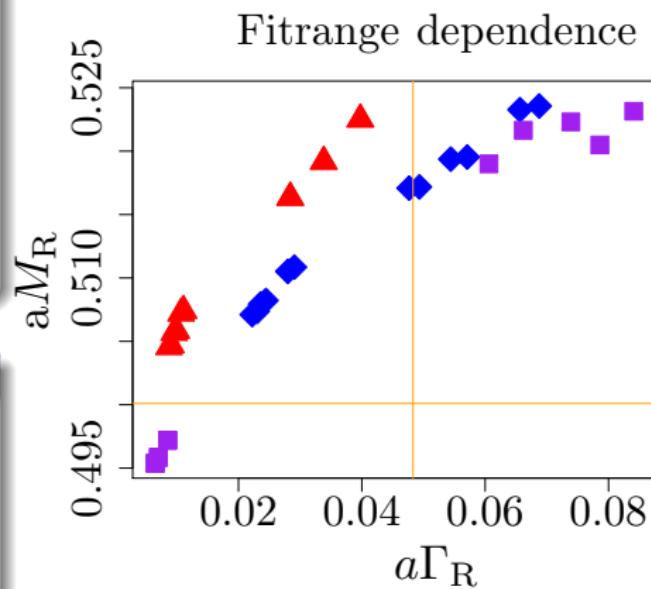
Luescher fits on the physical point ensemble

Performing several fits:

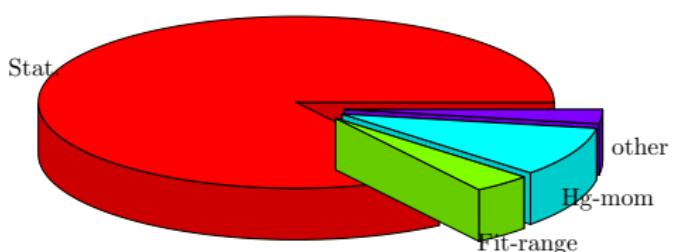
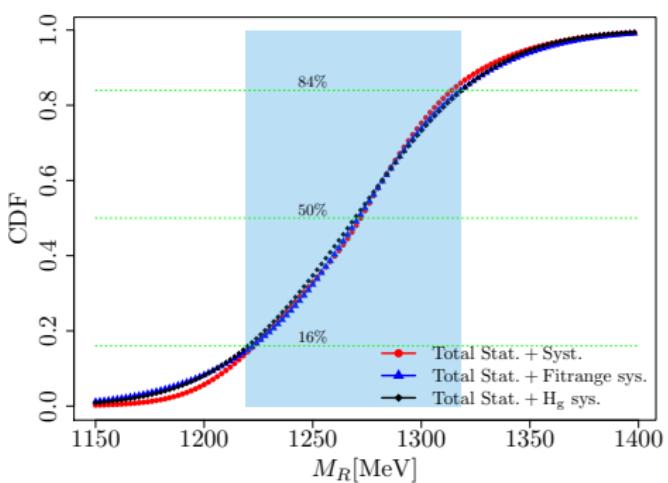
- p -wave only
- sp -wave
- below- $N\pi\pi$
- Different fitranges
- Different extraction

Stat.+Syst. error

- Model averaging
- Each fit is weighted
- $\sim e^{-0.5(\chi^2 - 2N_{\text{data}} + 2N_{\text{param}})}$



Result at the physical point twisted clover



Conclusion, outlook

This work (Details [arXiv:2307.12846])

| Ensemble | m_π [MeV] | L | m_Δ [MeV] | $g_{\Delta-\pi N}$ |
|----------------|---------------|---------|------------------|--------------------|
| Twisted-Clover | 139 MeV | 5.12 fm | 1267(46) MeV | |
| Nf2+1 Clover | 200 MeV | 3.7 fm | 1320(10) MeV | 17.6(2.7) |
| Nf2+1 Clover | 250 MeV | 2.8 fm | 1380(7) MeV | 13.6(5) |
| Nf2+1 Clover | 250 MeV | 3.7 fm | 1373(6) MeV | 10.3(1.6) |

| Collaboration | m_π [MeV] | Methodology | m_Δ [MeV] | $g_{\Delta-\pi N}$ |
|--------------------------|---------------|------------------------------|--------------------------|-------------------------|
| Verduci(2014) | 266 | Distillation, Lüscher | 1396(19) | 19.9(8) |
| Alexandrou et.al. (2013) | 360 | LO pert., Michael & McNeile | 1535(25) | 26.7(1.5) |
| Alexandrou et.al. (2015) | 180 | LO pert., Michael & McNeile | 1350(50) | 23.7(1.3) |
| Andersen et.al. (2017) | 280 | Stoch. distillation, Lüscher | 1344(20) | 37.1(9.2) |
| Morningstar et.al.(2022) | 200 | Stoch. distillation, Lüscher | 1290(7) | 14.41(53) _{BW} |
| Silvi et.al. (2021) | 255 | Smeared sources, Lüscher | 1380(7)(9) _{BW} | 13.6(5) _{BW} |

Summary

- Perform analysis on all the ensembles
- Perform chiral extrapolations



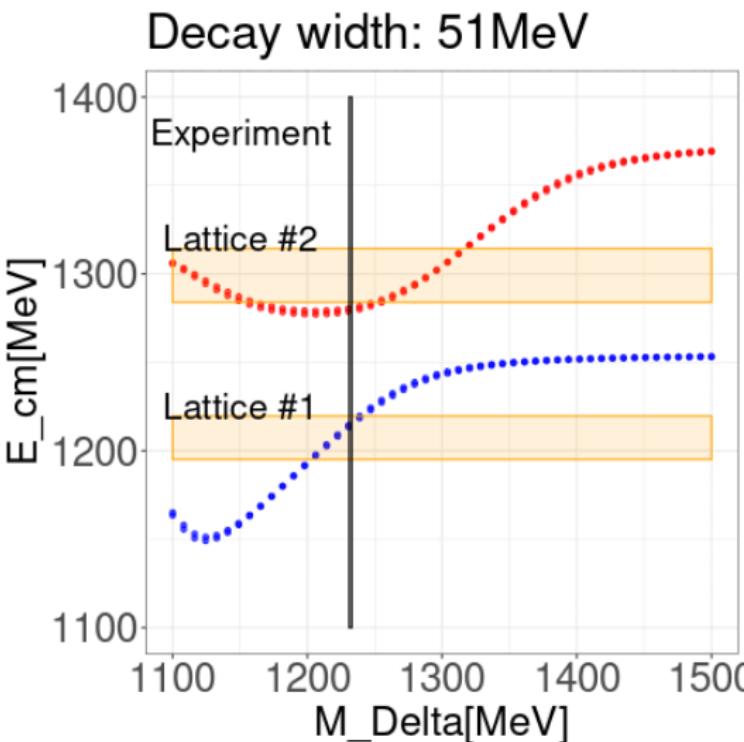
Acknowledgement

Thank you very much for your attention

The project acknowledge support

- Nice Quarks project
- Pizdaint supercomputer
- Juwelsbooster supercomputer
- NERSC supercomputer

Backup: The decay width

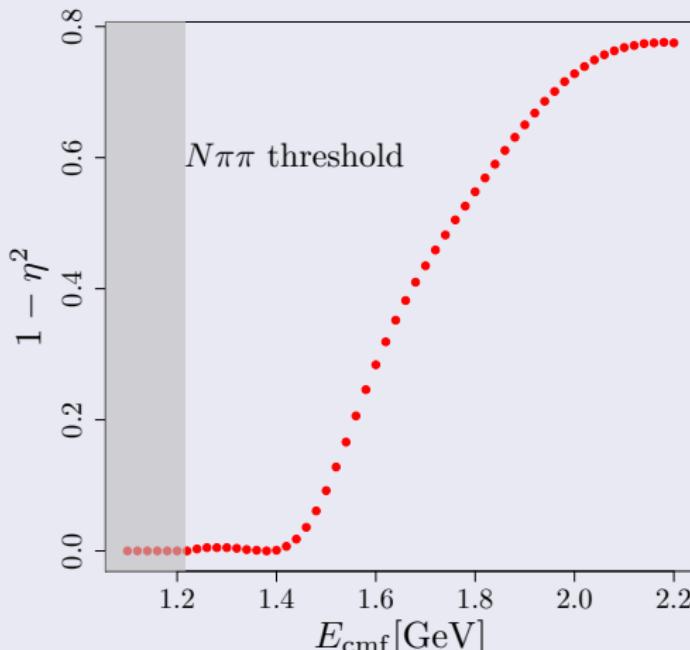


Backup: Scattering at the physical point

Data from <http://gwdac.phys.gwu.edu>

Challenge: $N\pi\pi$ threshold is very low

At the physical point $m_N + m_\pi < m_\Delta \rightarrow \Delta$ is unstable

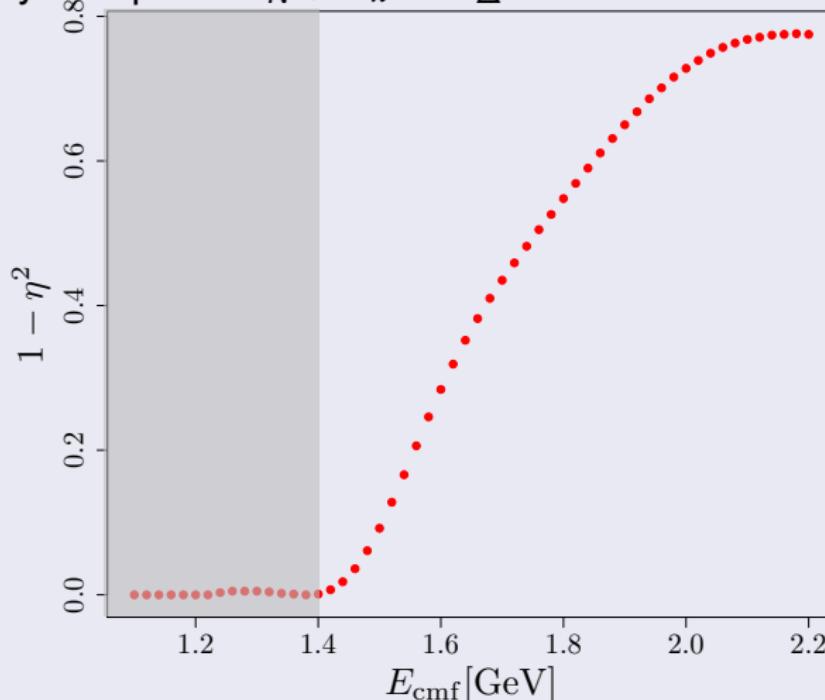


Backup: Scattering at the physical point

Data from <http://gwdac.phys.gwu.edu>

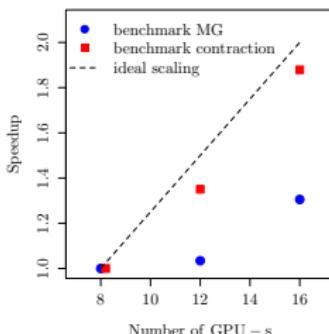
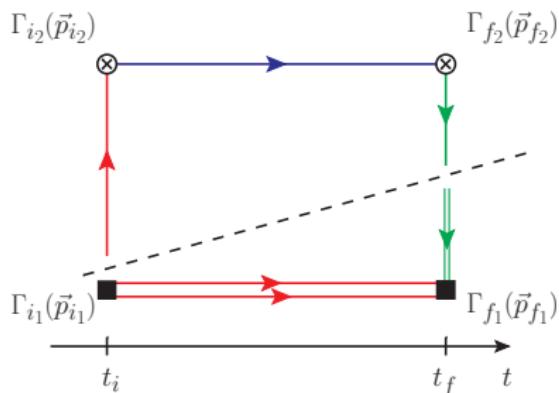
Challenge: $N\pi\pi$ threshold is very low

At the physical point $m_N + m_\pi < m_\Delta \rightarrow \Delta$ is unstable



Backup Code: Reductions on the GPU

B diagram: $U(x_{f1}, x_{i1})(\Gamma_{f1} D(x_{f1}, x_{f2}) \Gamma_{f2} U(x_{f2}, x_{i2}) \Gamma_{i2} D(x_{i2}, x_{i1}) \Gamma_{i1})^t U(x_{f1}, x_{i1})$



- Two correlated spatial sum (pion(f_2), nucleon(f_1))
- The problematic is the green line (sink-to-sink)
- Estimate it stochastically $D(x_{f1}, x_{f2}) = \sum_r \xi_r(f_1) \phi_r^\dagger(f_2)$
- Cut the diagram into factors
- Factors be combined to diagrams
- Many different diagrams share the same factors

Backup: Gramm-Schmidt irreducible representations

- irrep, irrep row(μ), # occurrences, # combinations of momenta
- As an example we have a 12×12 correlation matrix for the delta
- In the process of projection this matrix will be block diagonalized GS transformations
- Pion nucleon correlation matrix

| \vec{p}_{tot} , irrep name | N_{dim} |
|-------------------------------------|------------------|
| $\vec{p} = (0, 0, 0), G_{1u}$ | 8x8 |
| $\vec{p} = (0, 0, 0), H_g$ | 9x9 |
| $\vec{p} = (0, 0, 1), G_1$ | 24x24 |
| $\vec{p} = (0, 0, 1), G_2$ | 18x18 |
| $\vec{p} = (1, 1, 0), (2)G$ | 30x30 |
| $\vec{p} = (1, 1, 1), (3)G$ | 16x16 |
| $\vec{p} = (1, 1, 1), F_1$ | 6x6 |
| $\vec{p} = (1, 1, 1), F_2$ | 6x6 |

