
Nucleon Axial, Scalar, and Tensor Charges From Lattice QCD Simulations at the Physical Point

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Nucleon Structure with Physical Point Ensembles

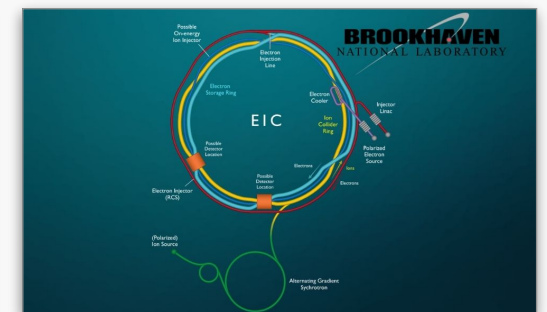
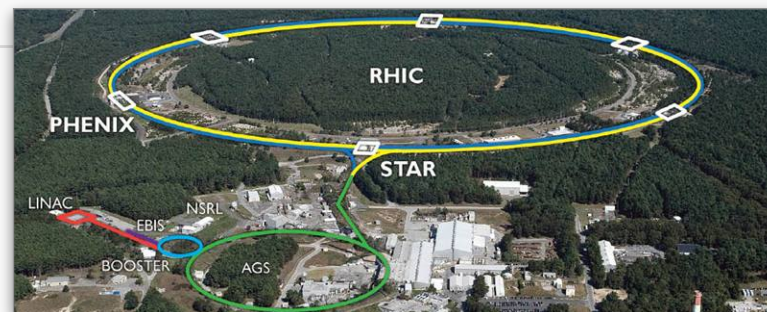
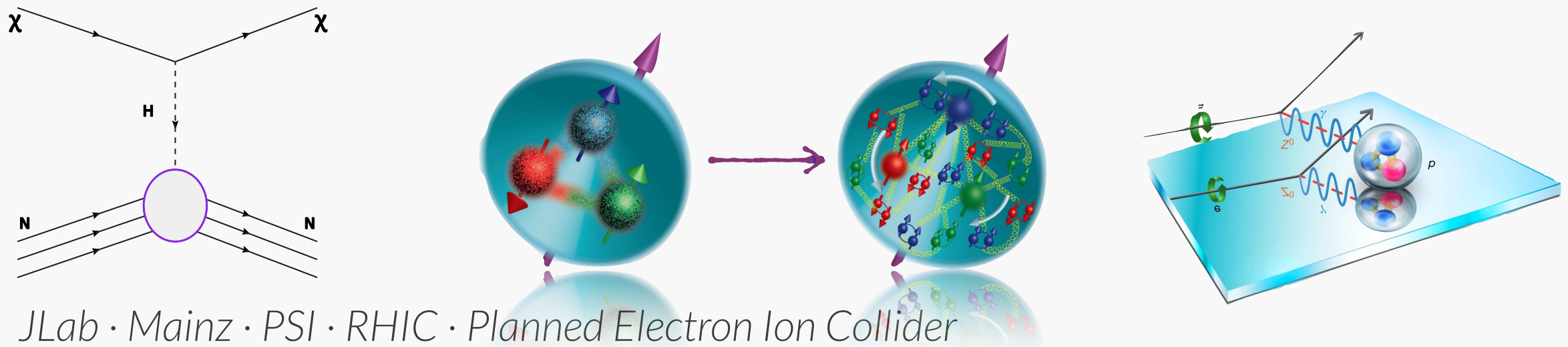
Outline

- Matrix elements on the lattice
- Lattice methods
- $N_f=2+1+1$ ensembles at physical point (twisted mass + clover)
- Statistics and uncertainties
- Results



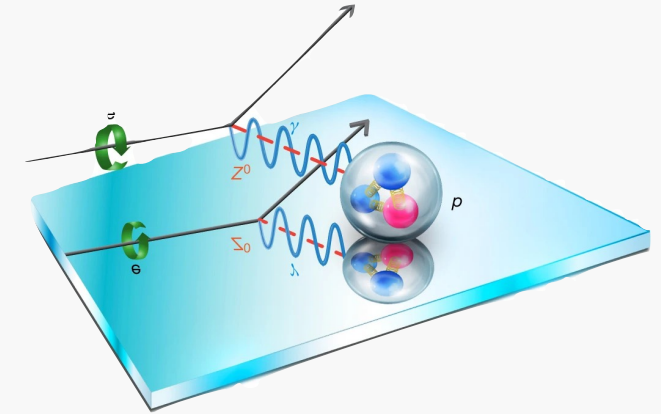
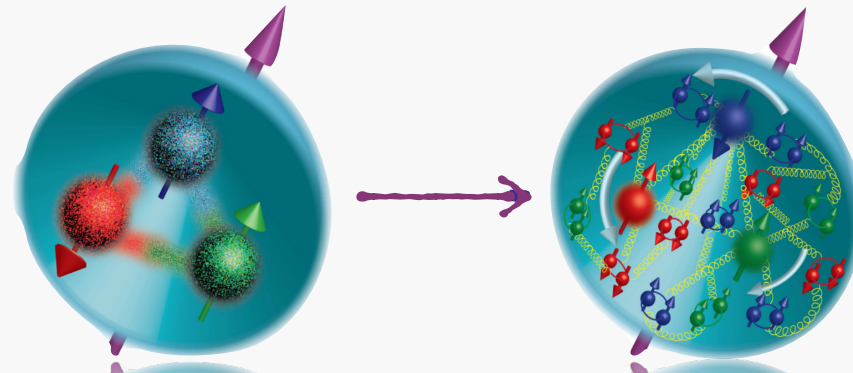
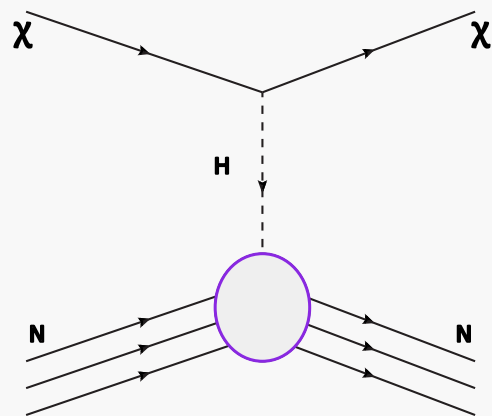
Nucleon Matrix Elements

- Scalar and tensor charges → novel interactions/dark matter searches
- Axial matrix elements → origin of nucleon spin
- σ -terms → mass decomposition of nucleon
- Electromagnetic form factors → radii and moments well known experimentally
- Axial form factors → PCAC and pion pole dominance relations
- Strange form factors → connect to weak charges and constraints on new physics
- Momentum fraction, moments of PDFs and GPDs, ...

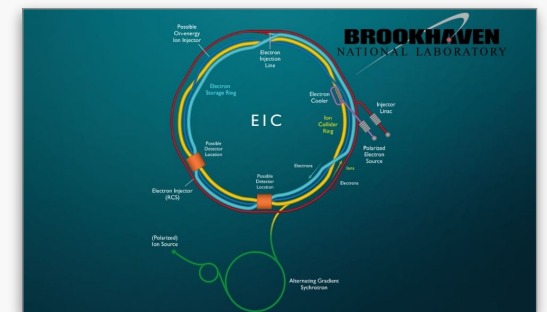
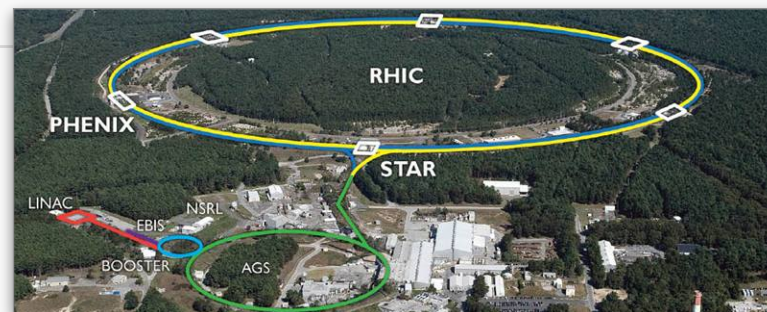


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JLab · Mainz · PSI · RHIC · Planned Electron Ion Collider



Lattice QCD

$$\langle \mathcal{O} \rangle = \frac{1}{Z} \int \mathcal{D}[\mathbf{U}] \mathcal{O}(\mathbf{D}_f^{-1}[\mathbf{U}], \mathbf{U}) \left(\prod_{f=u,d,s,c} \text{Det}(\mathbf{D}_f[\mathbf{U}]) \right) e^{-S_{\text{QCD}}[\mathbf{U}]}$$

1



2



3



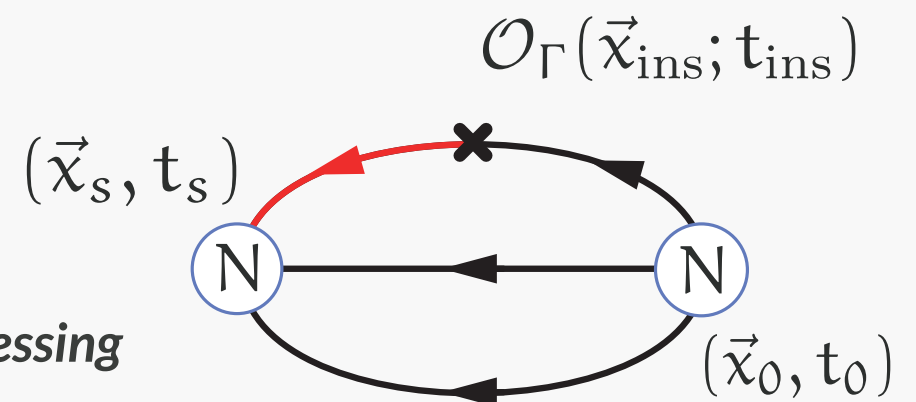
Simulation

- Markov chain Monte Carlo to generate ensembles of gluon-field configurations $\{\mathbf{U}\}$

$$P[\mathbf{U}] = \frac{1}{Z} \left(\prod_{f=u,d,s,c} \text{Det}(\mathbf{D}_f[\mathbf{U}]) \right) e^{-S_{\text{QCD}}[\mathbf{U}]}$$

Analysis

- Construction of hadron correlation functions on background field configurations

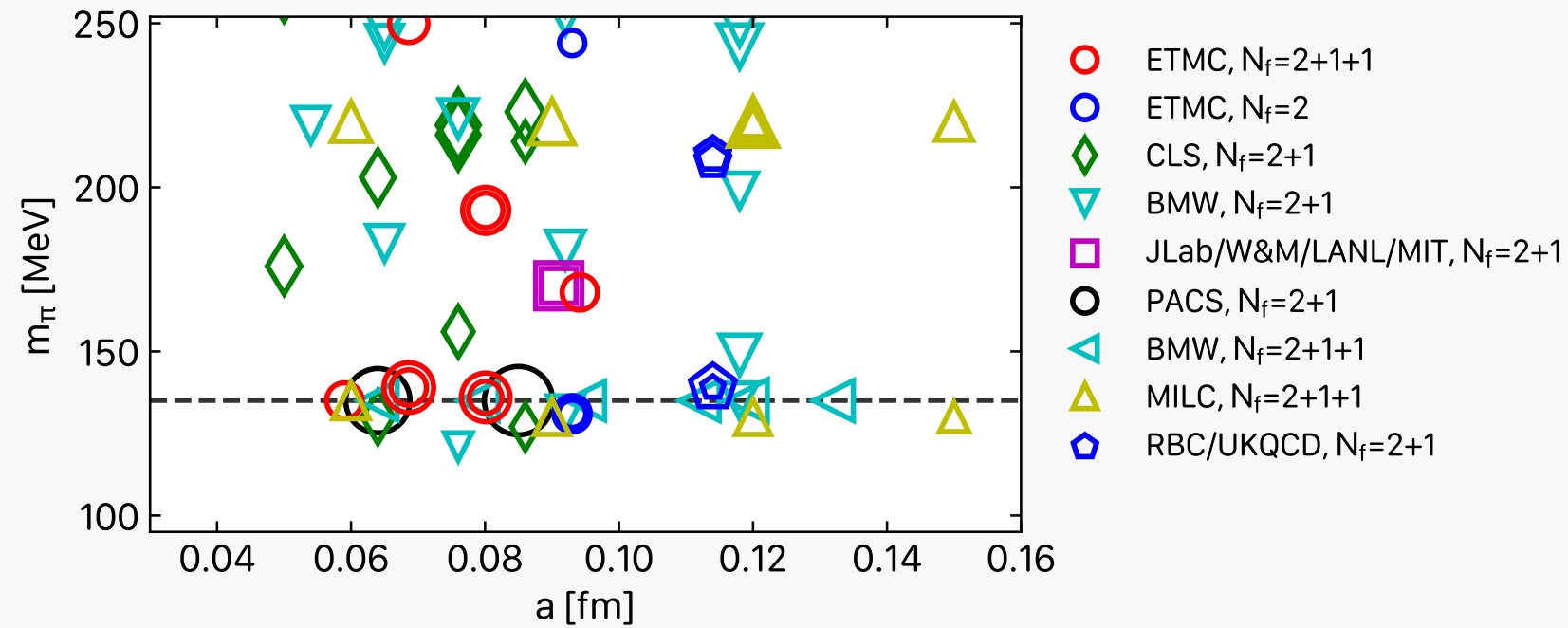


Data analysis - post-processing

- Statistical analysis, resampling
- Statistical and stochastic errors
- Continuum and infinite volume extrapolation

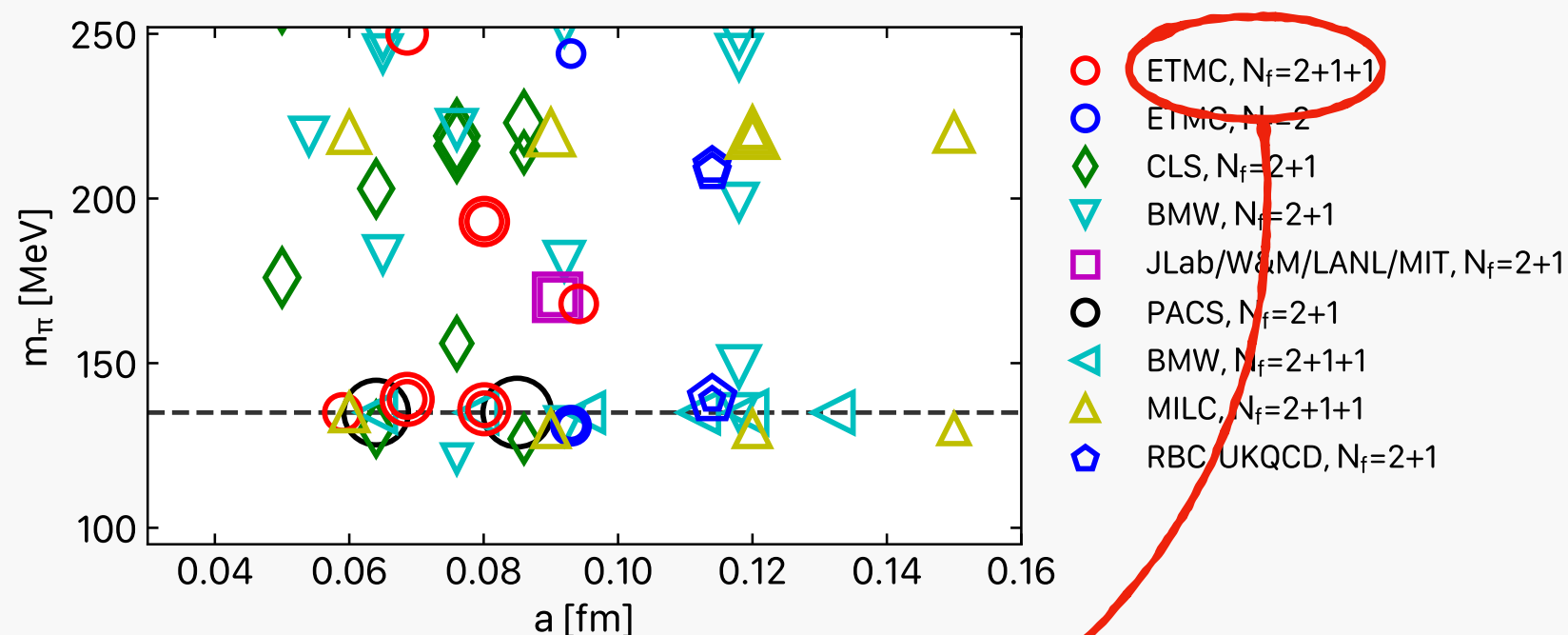
Ensembles

Landscape of ensembles used for nucleon structure



Ensembles

Landscape of ensembles used for nucleon structure



ETMC: three $N_f=2+1+1$ ensembles at physical pion mass

Ens. ID (abbrv.)	Vol.	a [fm]
cB211.072.64 (cB64)	64×128	0.080
cC211.060.80 (cC80)	80×160	0.068
cD211.054.96 (cD96)	96×192	0.057

- Three lattice spacings at physical point
- Ongoing generation of finer ensembles and larger volumes
- ***This talk:*** 3 ensembles with:
 $a = 0.057 - 0.068$ fm

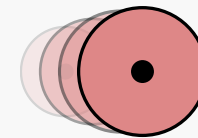
Local matrix elements

$$G_{\Gamma}(P; \vec{q}; t_s, t_{\text{ins}}) = \sum_{\vec{x}_s, \vec{x}_{\text{ins}}} e^{-i\vec{q} \cdot \vec{x}_{\text{ins}}} p^{\alpha\beta} \langle \bar{\chi}_{\text{N}}^{\beta}(\vec{x}_s; t_s) | \mathcal{O}^{\Gamma}(\vec{x}_{\text{ins}}; t_{\text{ins}}) | \chi_{\text{N}}^{\alpha}(\vec{0}; 0) \rangle$$

Unpolarised

$$\mathcal{O}_{\text{V}}^{\mu\mu_1\mu_2\dots\mu_n} = \bar{\psi} \gamma^{\mu} iD^{\mu_1} iD^{\mu_2} \dots iD^{\mu_n} \psi$$

$$\langle 1 \rangle_{u-d} = g_V, \quad \langle x \rangle_{u-d}, \quad \dots$$



Helicity

$$\mathcal{O}_{\text{A}}^{\mu\mu_1\mu_2\dots\mu_n} = \bar{\psi} \gamma_5 \gamma^{\mu} iD^{\mu_1} iD^{\mu_2} \dots iD^{\mu_n} \psi$$

$$\langle 1 \rangle_{\Delta u - \Delta d} = g_A, \quad \langle x \rangle_{\Delta u - \Delta d}, \quad \dots$$



Transverse

$$\mathcal{O}_{\text{T}}^{\nu\mu\mu_1\mu_2\dots\mu_n} = \bar{\psi} \sigma^{\nu} iD^{\mu_1} iD^{\mu_2} \dots iD^{\mu_n} \psi$$

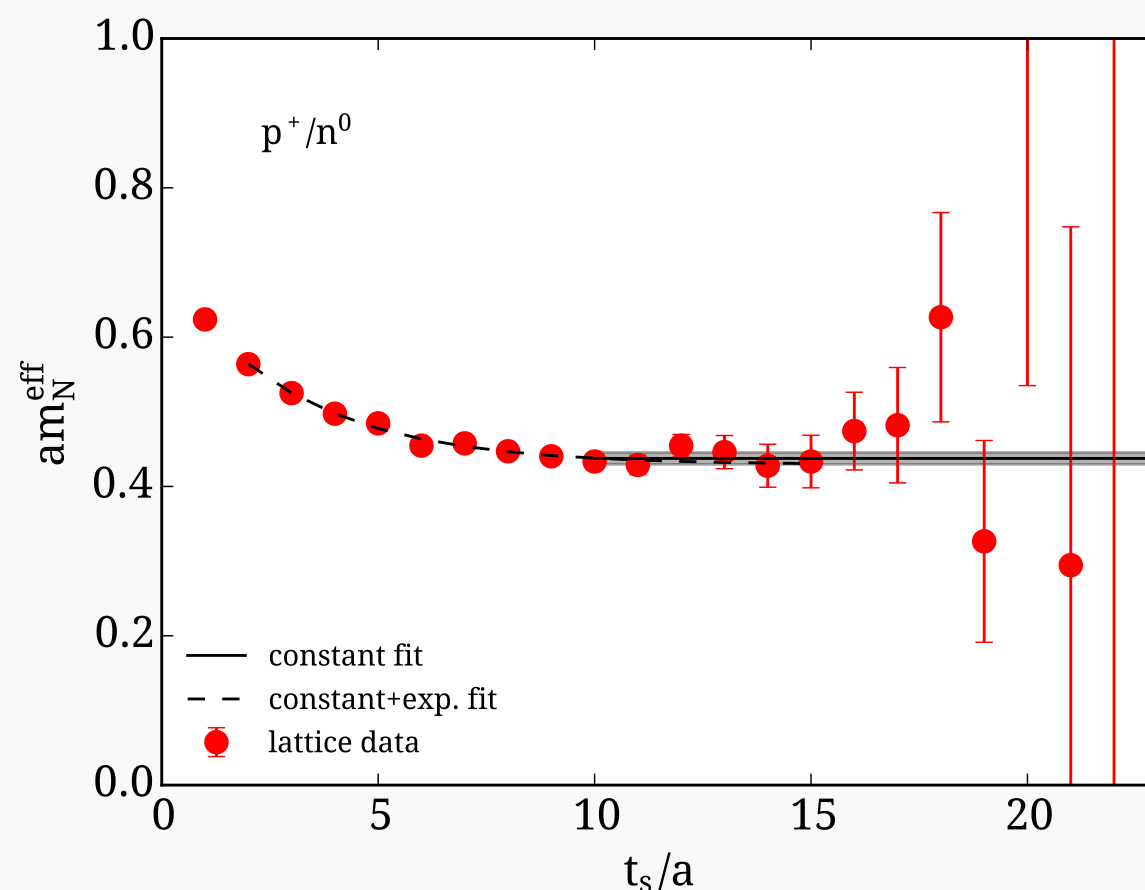
$$\langle 1 \rangle_{\delta u - \delta d} = g_T, \quad \langle x \rangle_{\delta u - \delta d}, \quad \dots$$



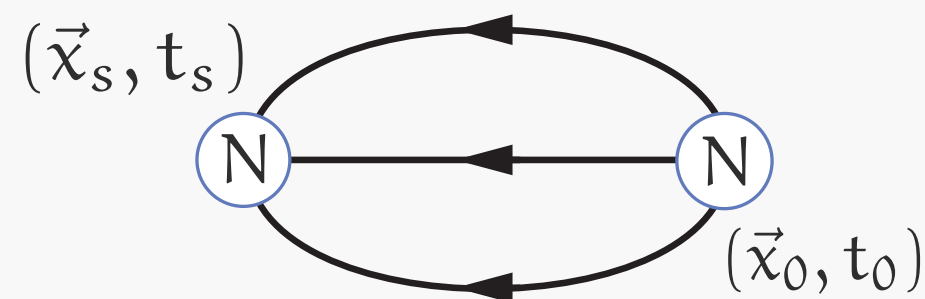
Nucleon structure on the lattice

Two-point correlation functions

- **Statistical error:** $N^{-1/2}$ with Monte Carlo samples
- Correlation functions exponentially decay with time-separation
- Contamination from higher energy states



$$\sum_{\vec{x}_s} \Gamma^{\alpha\beta} \langle \bar{\chi}_N^\beta(x_s) | \chi_N^\alpha(0) \rangle = c_0 e^{-E_0 t_s} + c_1 e^{-E_1 t_s} + \dots$$



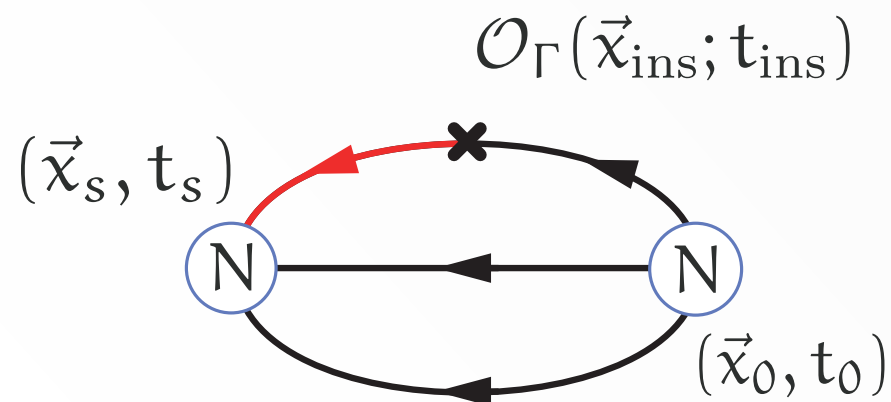
Matrix elements on the Lattice

General three-point function:

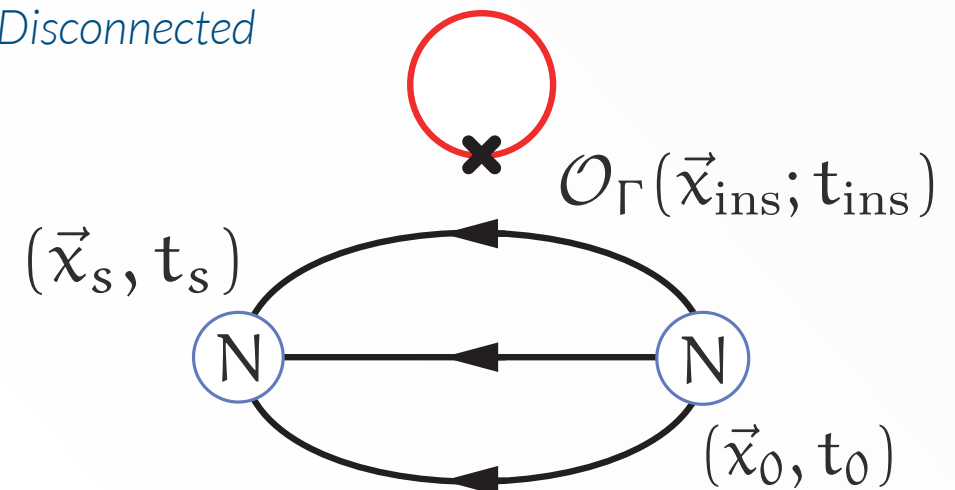
$$G_{\Gamma}(P; \vec{q}; t_s, t_{\text{ins}}) = \sum_{\vec{x}_s, \vec{x}_{\text{ins}}} e^{-i\vec{q} \cdot \vec{x}_{\text{ins}}} p^{\alpha\beta} \langle \bar{\chi}_{\text{N}}^{\beta}(\vec{x}_s; t_s) | \mathcal{O}_{\Gamma}(\vec{x}_{\text{ins}}; t_{\text{ins}}) | \chi_{\text{N}}^{\alpha}(\vec{0}; 0) \rangle$$

At quark level gives rise to both so-called connected and disconnected contributions

Connected



Disconnected



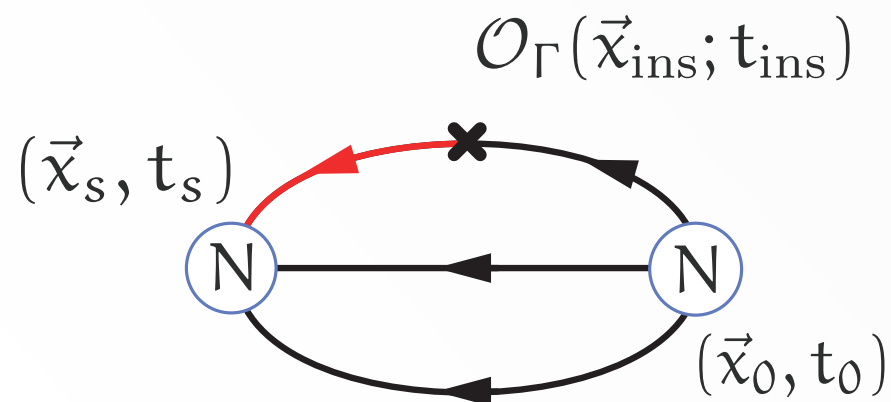
Matrix elements on the Lattice

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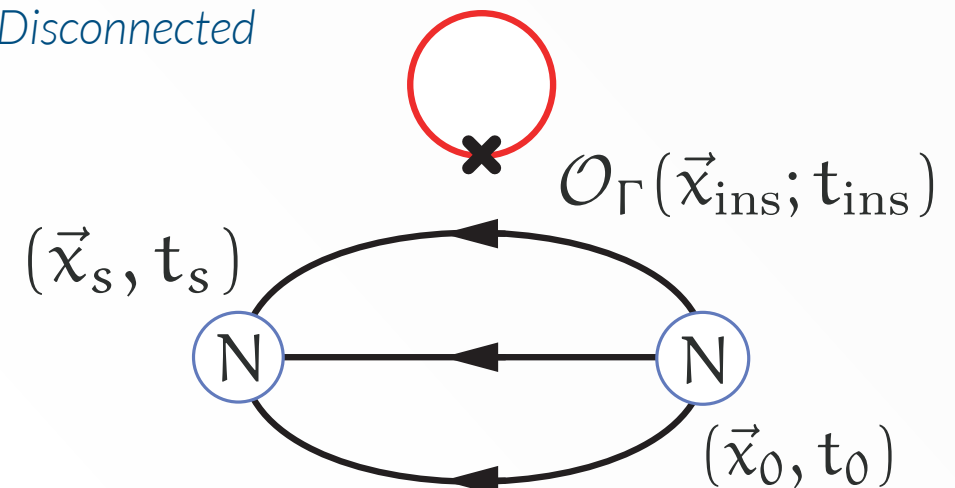
$$G_{\Gamma}(P; \vec{q}; t_s, t_{\text{ins}}) = \sum_{\vec{x}_s, \vec{x}_{\text{ins}}} e^{-i\vec{q} \cdot \vec{x}_{\text{ins}}} p^{\alpha\beta} \langle \bar{\chi}_{\text{N}}^{\beta}(\vec{x}_s; t_s) | \mathcal{O}_{\Gamma}(\vec{x}_{\text{ins}}; t_{\text{ins}}) | \chi_{\text{N}}^{\alpha}(\vec{0}; 0) \rangle$$

At quark level gives rise to both so-called connected and disconnected contributions

Connected

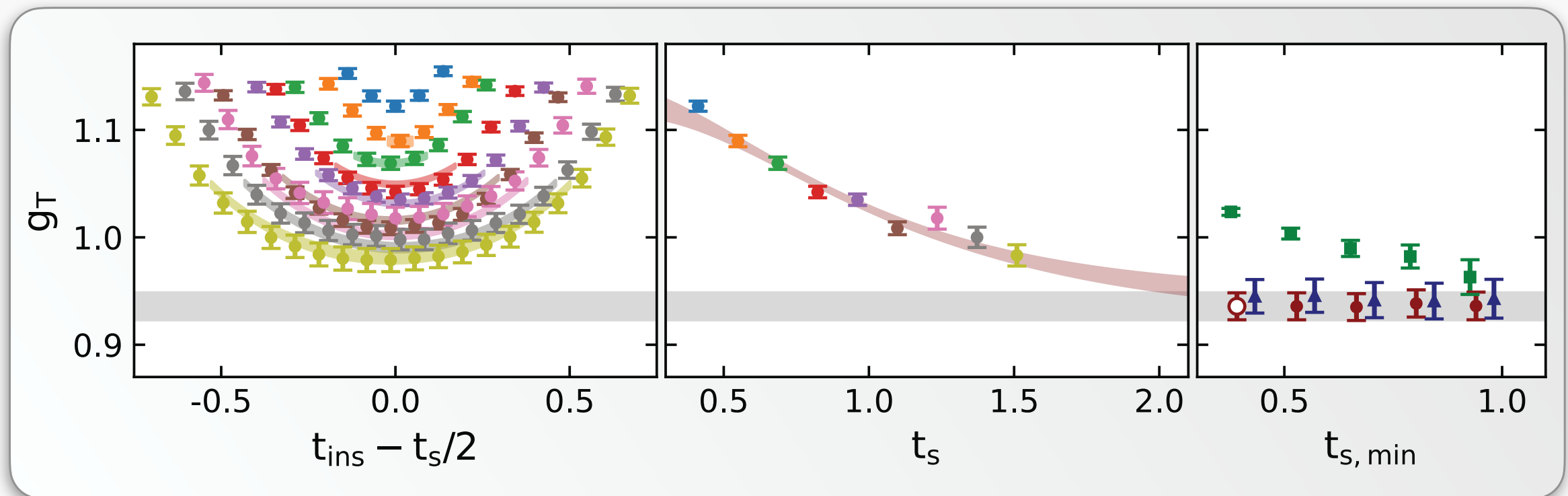


Disconnected



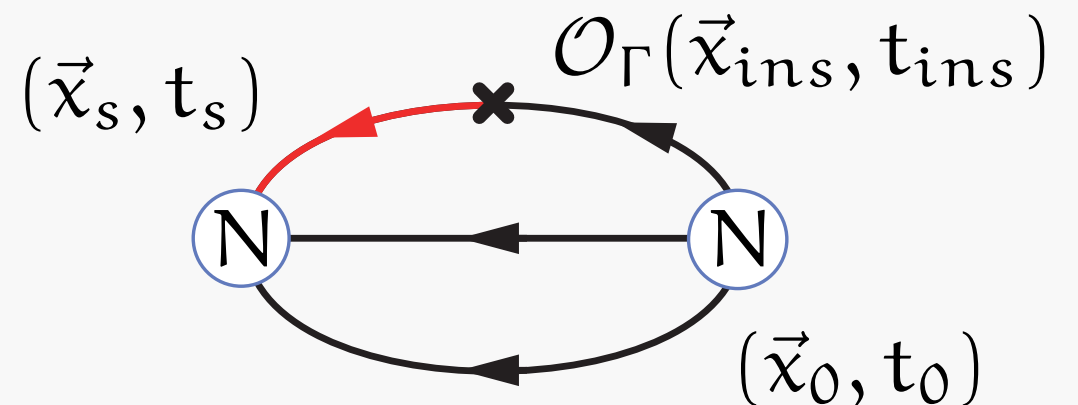
Stochastic evaluation of loop – stochastic error in addition to statistical

Treatment of excited states

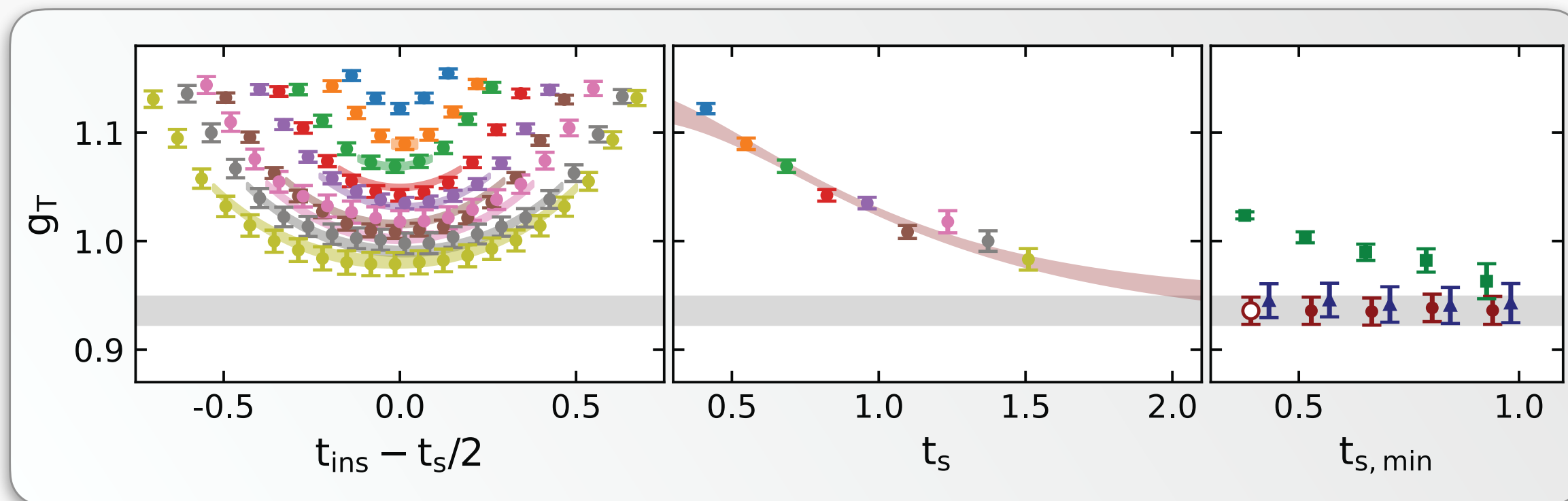


Example from intermediate α

- Isovector tensor charge (only connected)
- Increasing statistics with separation t_s
- Summation, two- and three-state fits



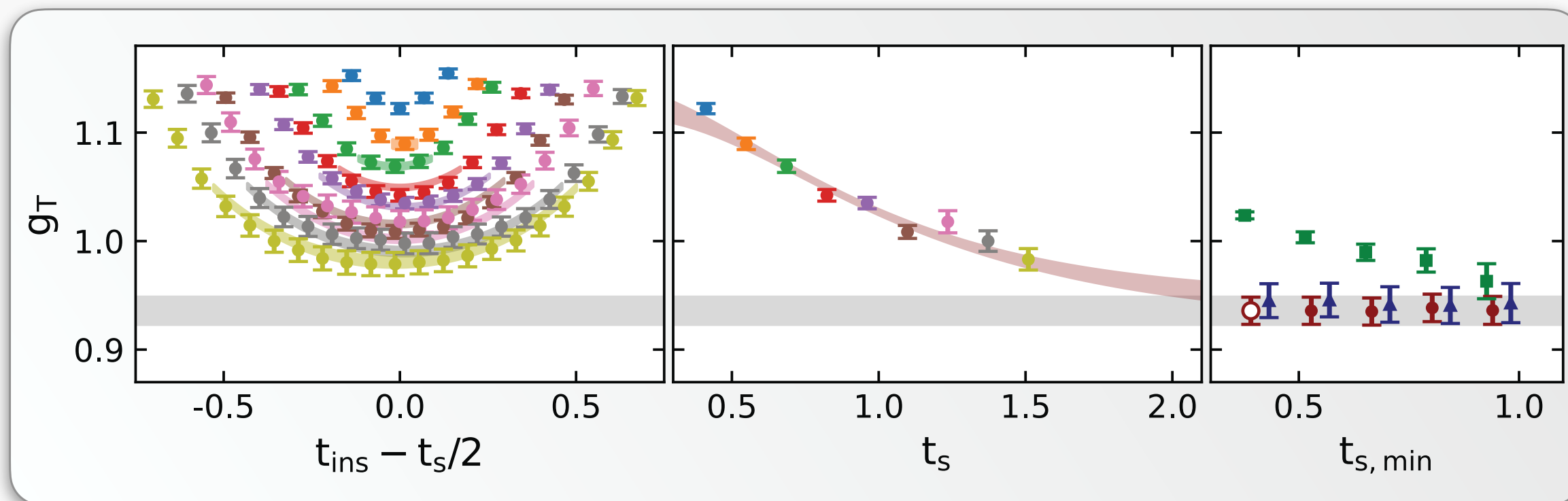
Treatment of excited states



Summation method

$$S_{\Gamma}(\vec{q}; t_s) = \sum_{t_{\text{ins}}=\tau}^{t_s-\tau} R_{\Gamma}(\vec{q}; t_s; t_{\text{ins}}) \rightarrow \mathcal{M}t_s + C$$

Treatment of excited states



Two-state fit

$$G_{\Gamma}(\vec{q}; t_s, t_{\text{ins}}) = \sum_{i=0}^1 \sum_{j=0}^1 c_{ij} e^{-E_i(0)(t_s - t_{\text{ins}})} e^{-E_j(\vec{q})t_{\text{ins}}}$$

$$\mathcal{M} = \frac{c_{00}}{\sqrt{a_0(\vec{0})a_0(\vec{q})}}$$

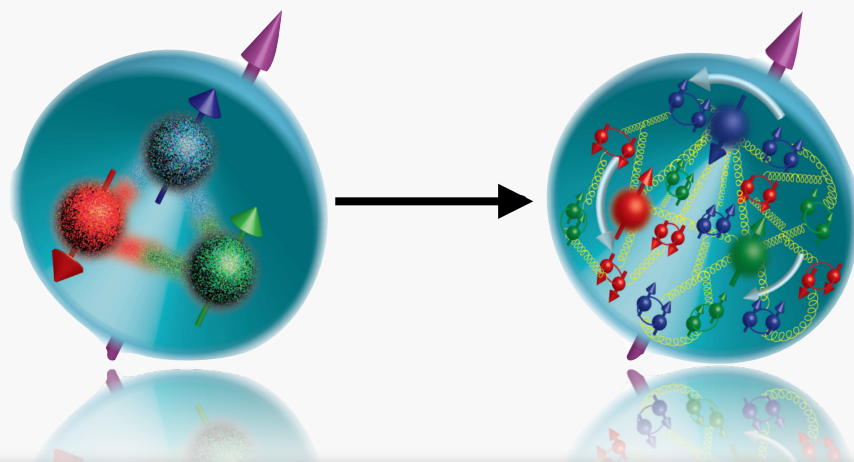
$$G(\vec{q}; t_s) = a_0(\vec{q})e^{-\varepsilon_0(\vec{q})t_s} + a_1(\vec{q})e^{-\varepsilon_1(\vec{q})t_s}$$

Nucleon axial charge

Matrix element of the axial current

Isovector case well known from β -decay: $\langle p | \bar{u} \gamma_5 \gamma_k d | n \rangle$

Flavor-separated contributions to axial charge relate to quark intrinsic spin contributions to nucleon spin

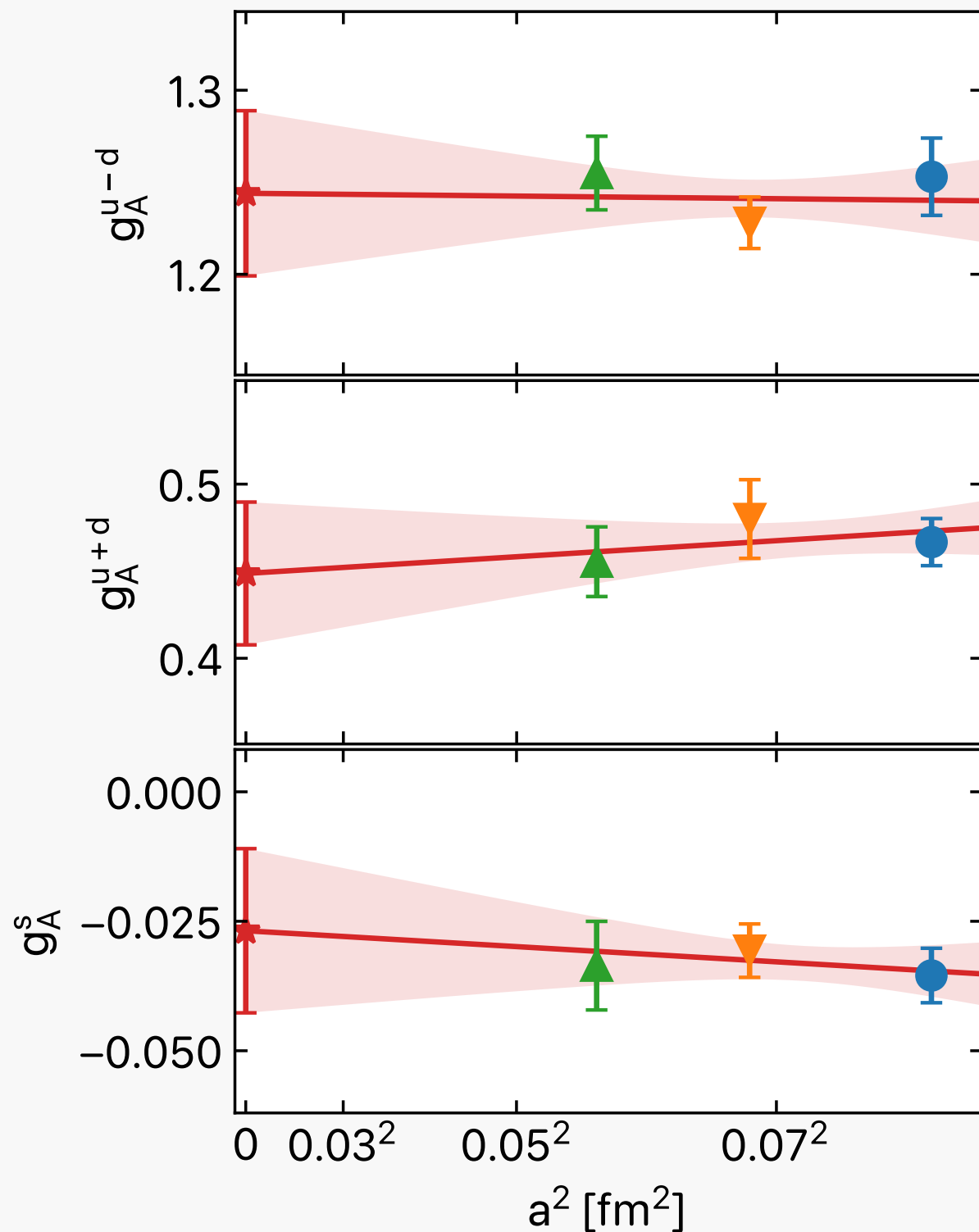


$$\frac{1}{2} \Delta \Sigma = \frac{1}{2} \sum_{q=u,d,s,\dots} g_A^q$$

Quark intrinsic spin contributions to nucleon spin

- Need linear combination of isovector ($u-d$) and isoscalar ($u+d$) contributions for individual up- and down-quarks
- Strange quark contribution is sea-quark contribution only (disconnected diagrams)

Nucleon axial charge



See arXiv:2309.05774 [hep-lat];

Includes isovector axial form-factors

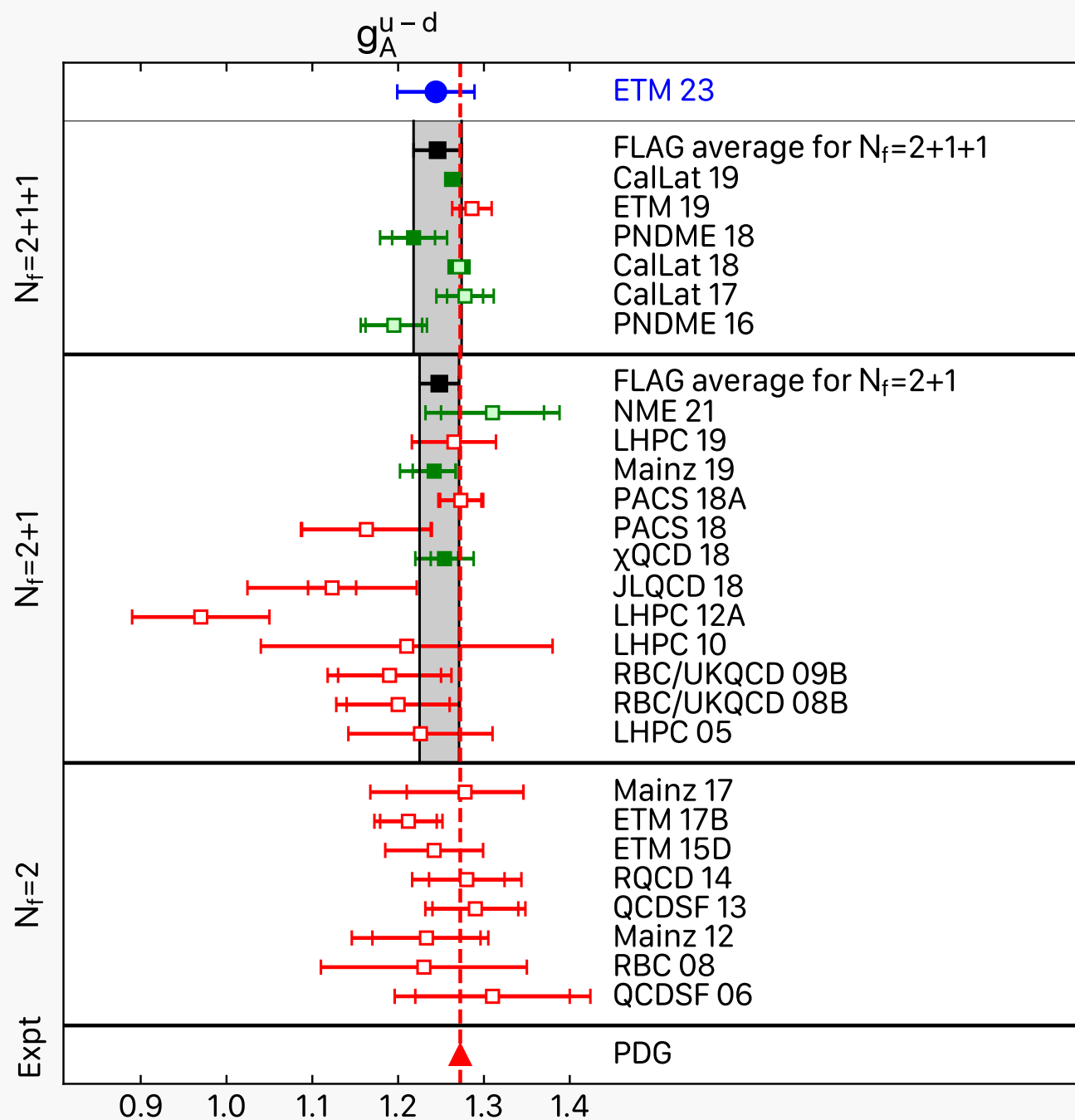
● Errors for each ensemble include *statistical* and *systematic* due to excited state contamination

● Model averaged based on AIC (see e.g. arXiv:2208.14983)

Preliminary!

Complete analysis for flavour separated charges ongoing

Nucleon axial charge



Latest FLAG21 values

- ETM23 consistent with FLAG average
- Only result with three physical point ensembles
- Agreement for g_A means confidence for more challenging quantities
- E.g.

– Scalar ME, σ -terms

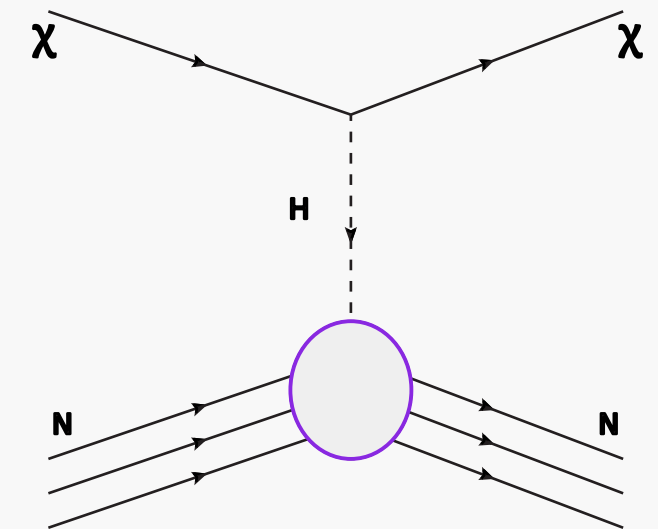
$$\sigma_{\pi N} = m_{ud} \langle N | \bar{u}u + \bar{d}d | N \rangle$$

– Tensor ME

$$g_T = \langle 1 \rangle_{\delta u - \delta d} \leftarrow \langle N | \bar{u} \sigma_{\mu\nu} u + \bar{d} \sigma_{\mu\nu} d | N \rangle$$

Scalar charge – σ -terms

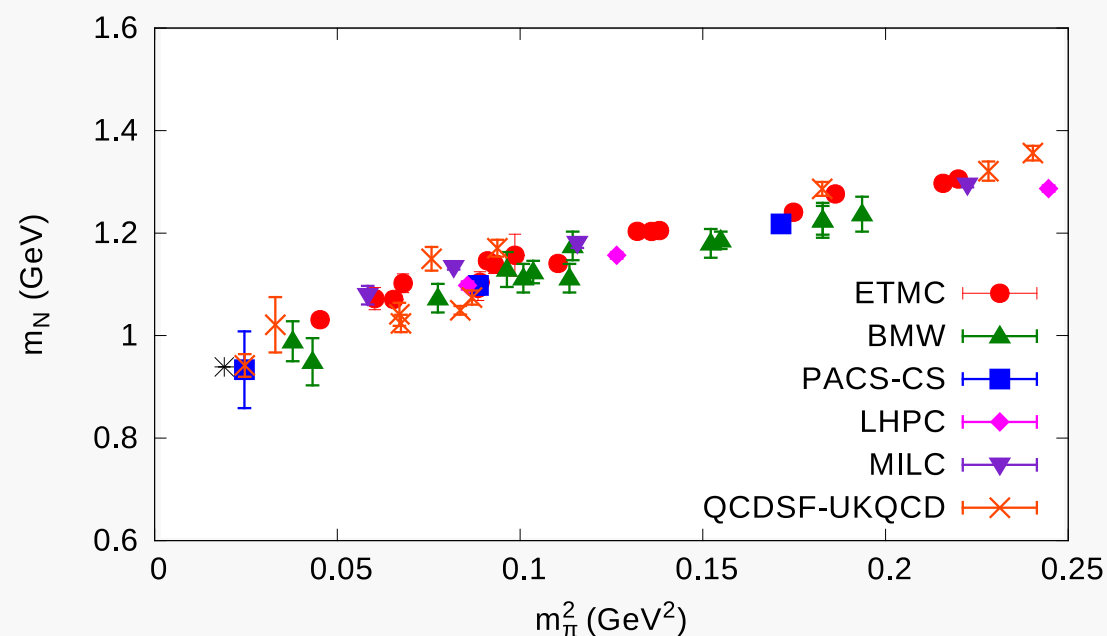
- Pion nucleon σ -term: $\sigma_{\pi N} = m_{ud} \langle N | \bar{u}u + \bar{d}d | N \rangle$
- Strange σ -term: $\sigma_s = m_s \langle N | \bar{s}s | N \rangle$
- Enter super-symmetric candidate particle scattering cross sections with nucleon (e.g. neutralino through Higgs)



1. Direct calculation of matrix elements

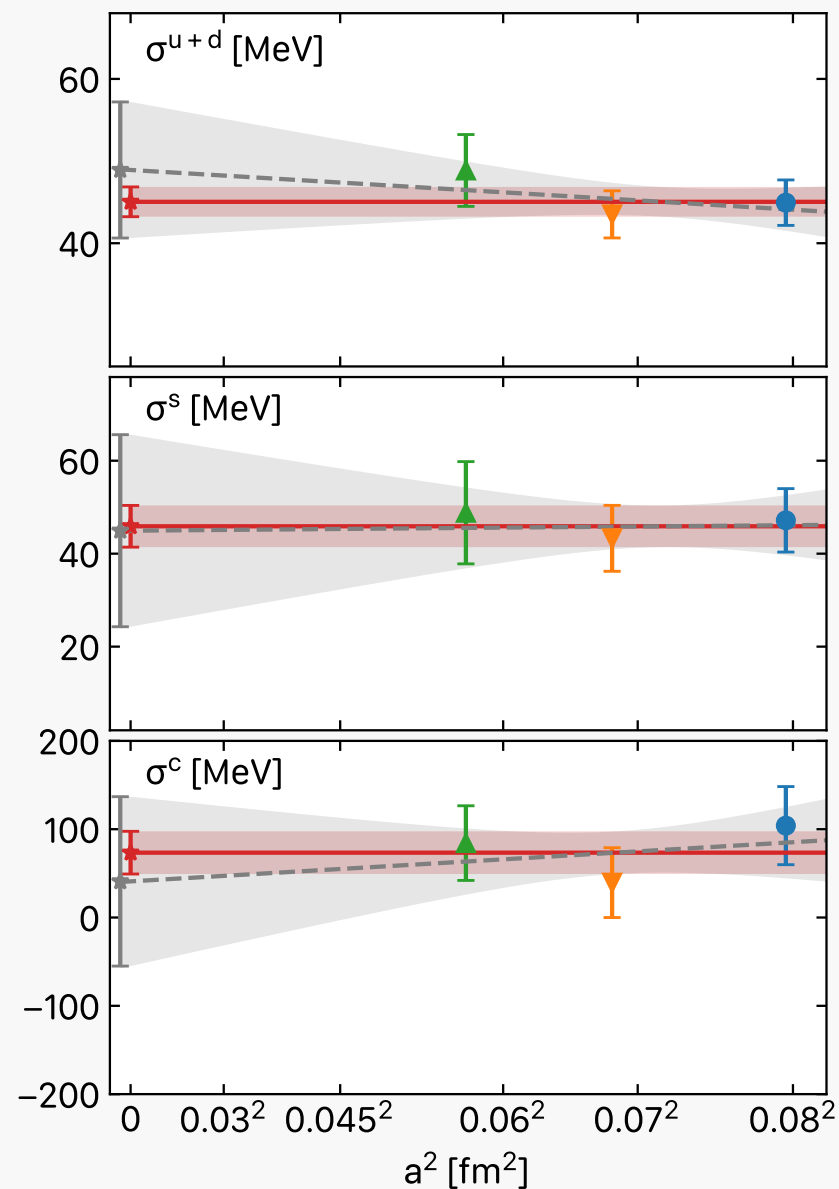
Involves disconnected contributions

2. Through Feynman - Hellmann theorem: $\sigma_{\pi N} = m_{ud} \frac{\partial m_N}{\partial m_{ud}}$ $\sigma_s = m_s \frac{\partial m_N}{\partial m_s}$

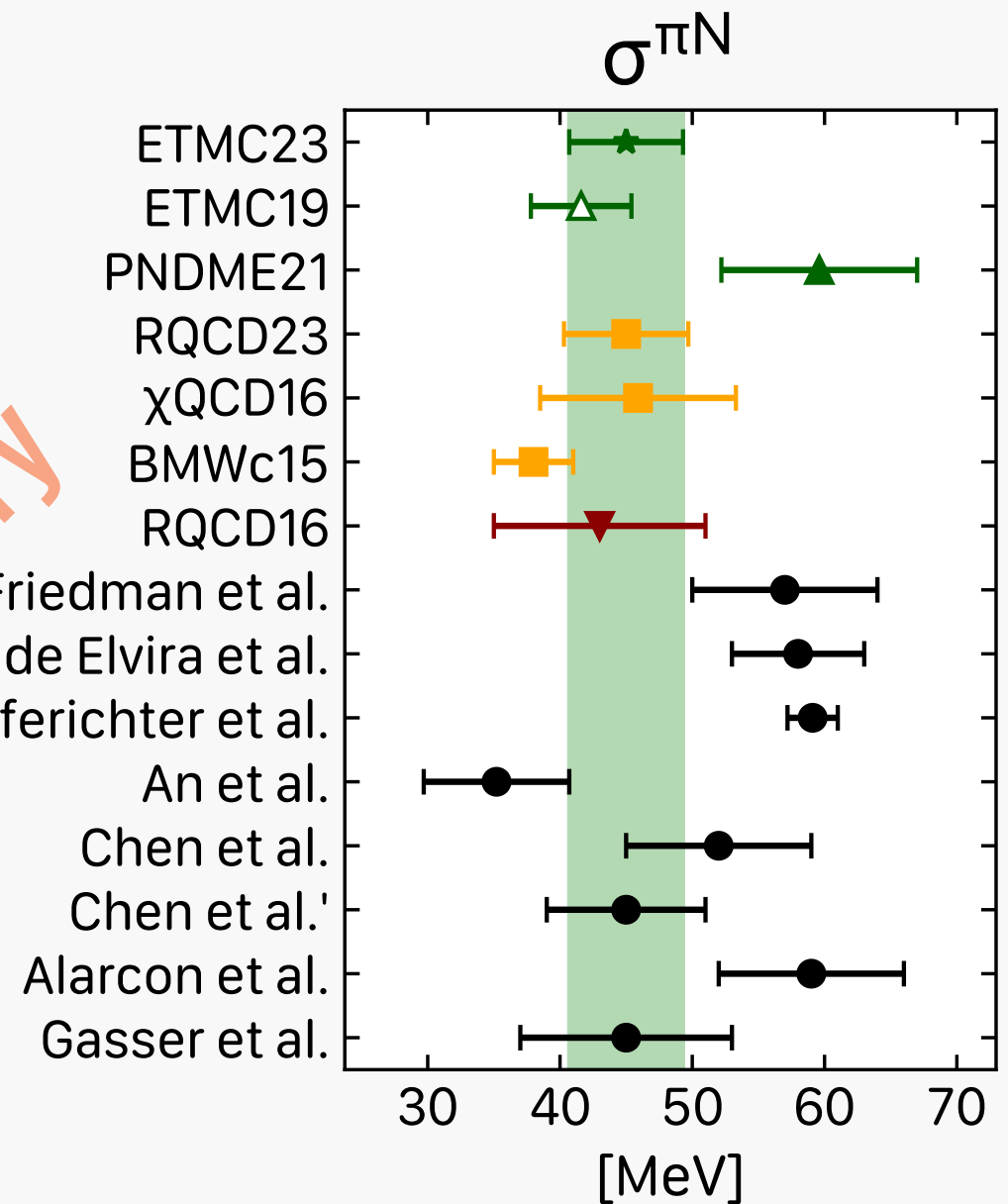


- Reliance on effective theories for dependence on m_π
- Weak dependence on m_s

Scalar charge – σ -terms

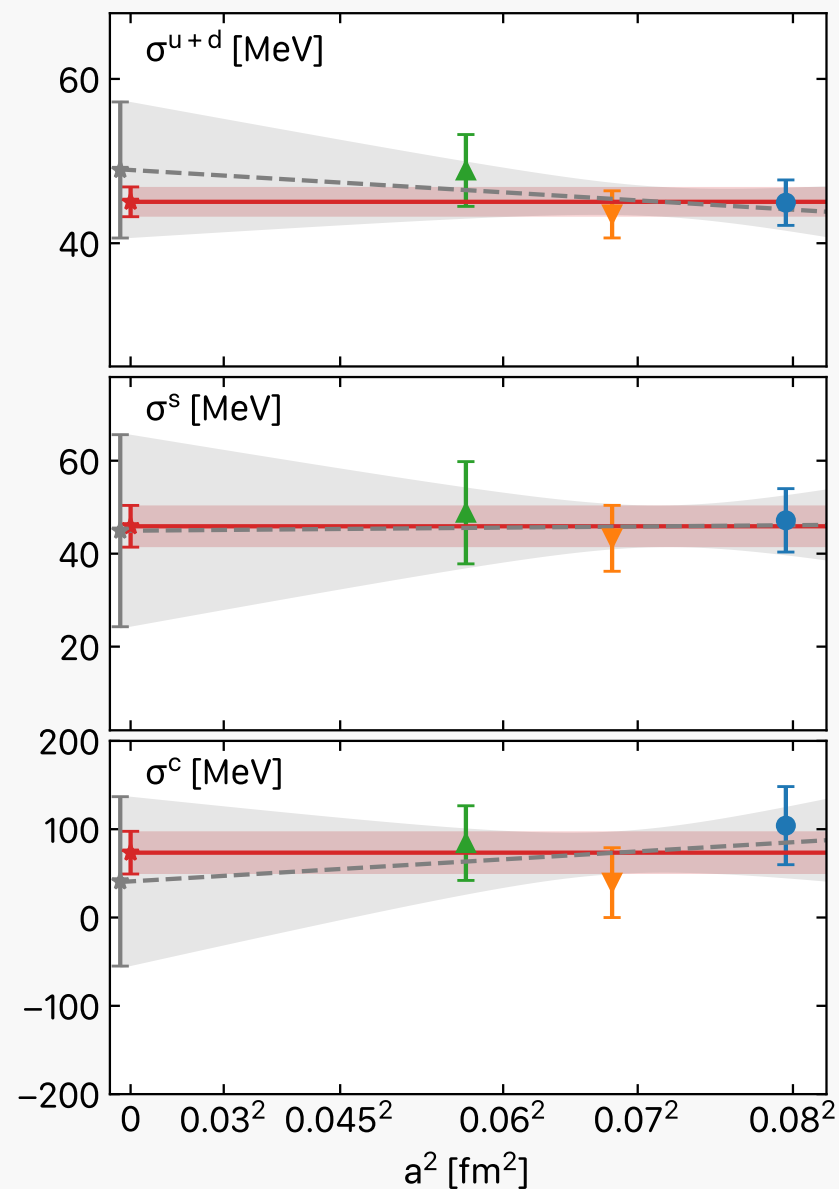


Preliminary

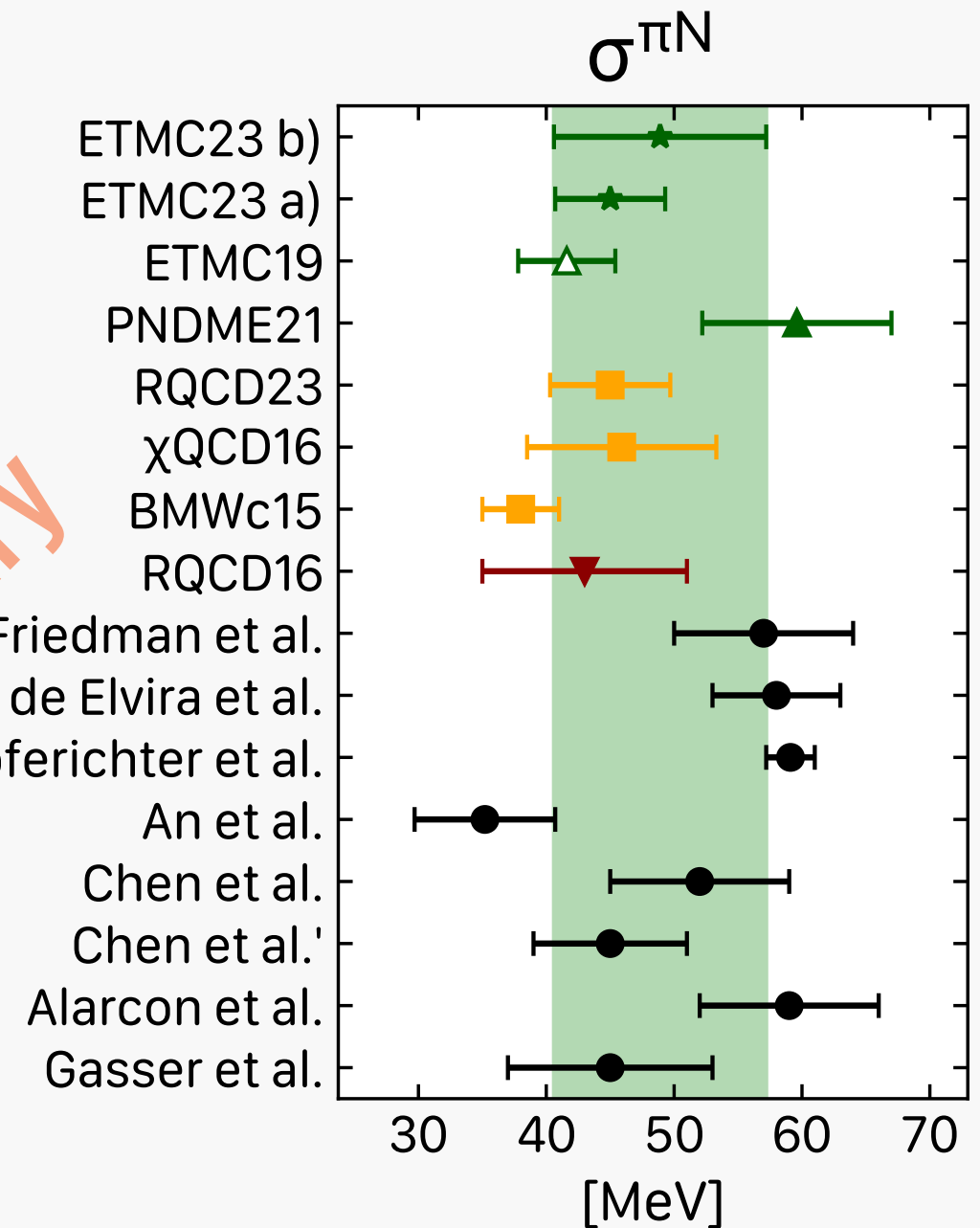


- One result using FH method (BMW)
- Los Alamos group: 59.6(7.4) MeV, when explicitly including πN energy as prior

Scalar charge – σ -terms

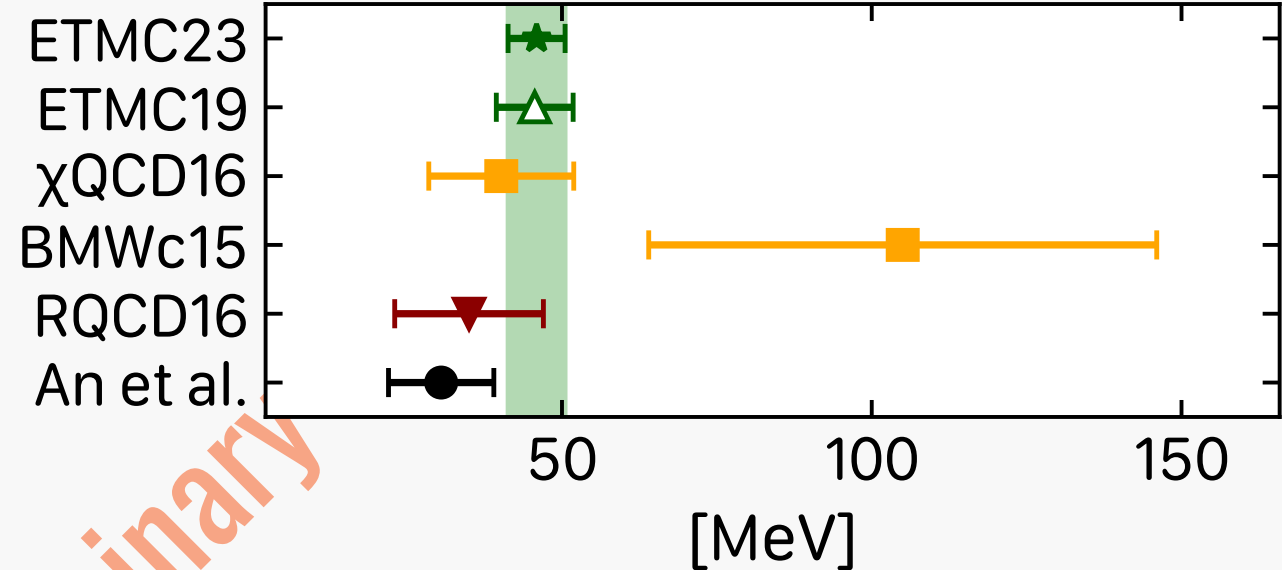
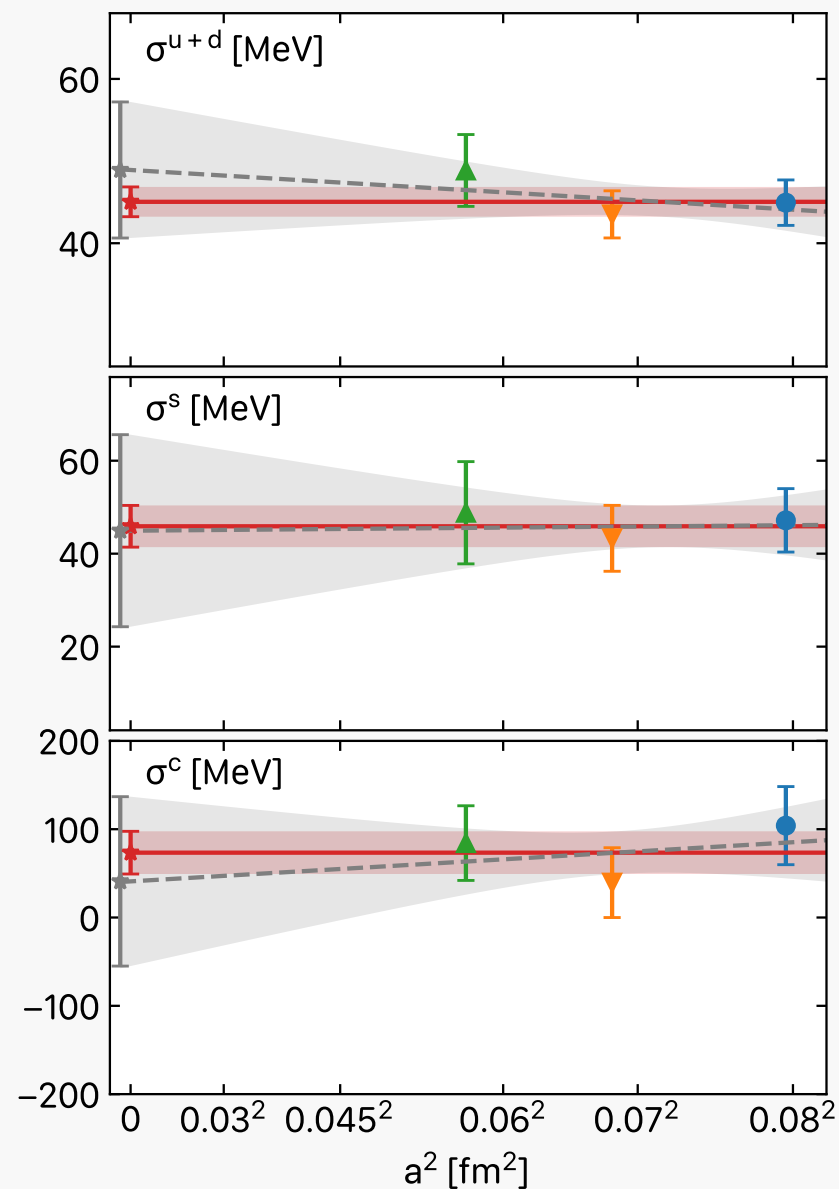


Preliminary



- One result using FH method (BMW)
- Los Alamos group: $59.6(7.4)$ MeV, when explicitly including πN energy as prior
- Our result with linear continuum extrapolation (Grey band in left plot)

Scalar charge – σ -terms



Preliminary

Strange content of the nucleon

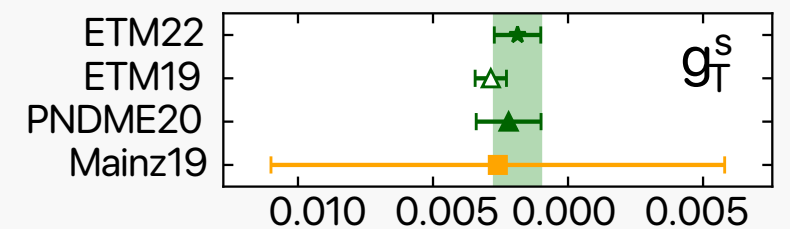
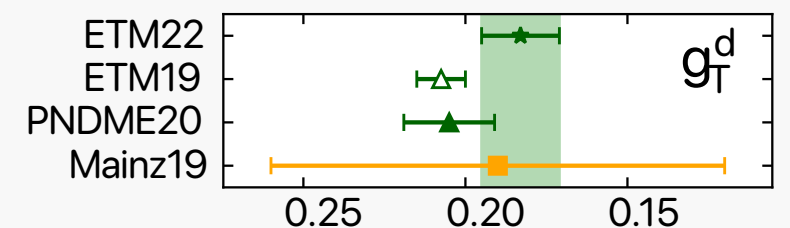
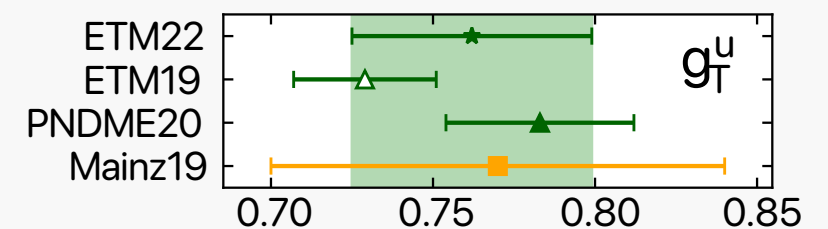
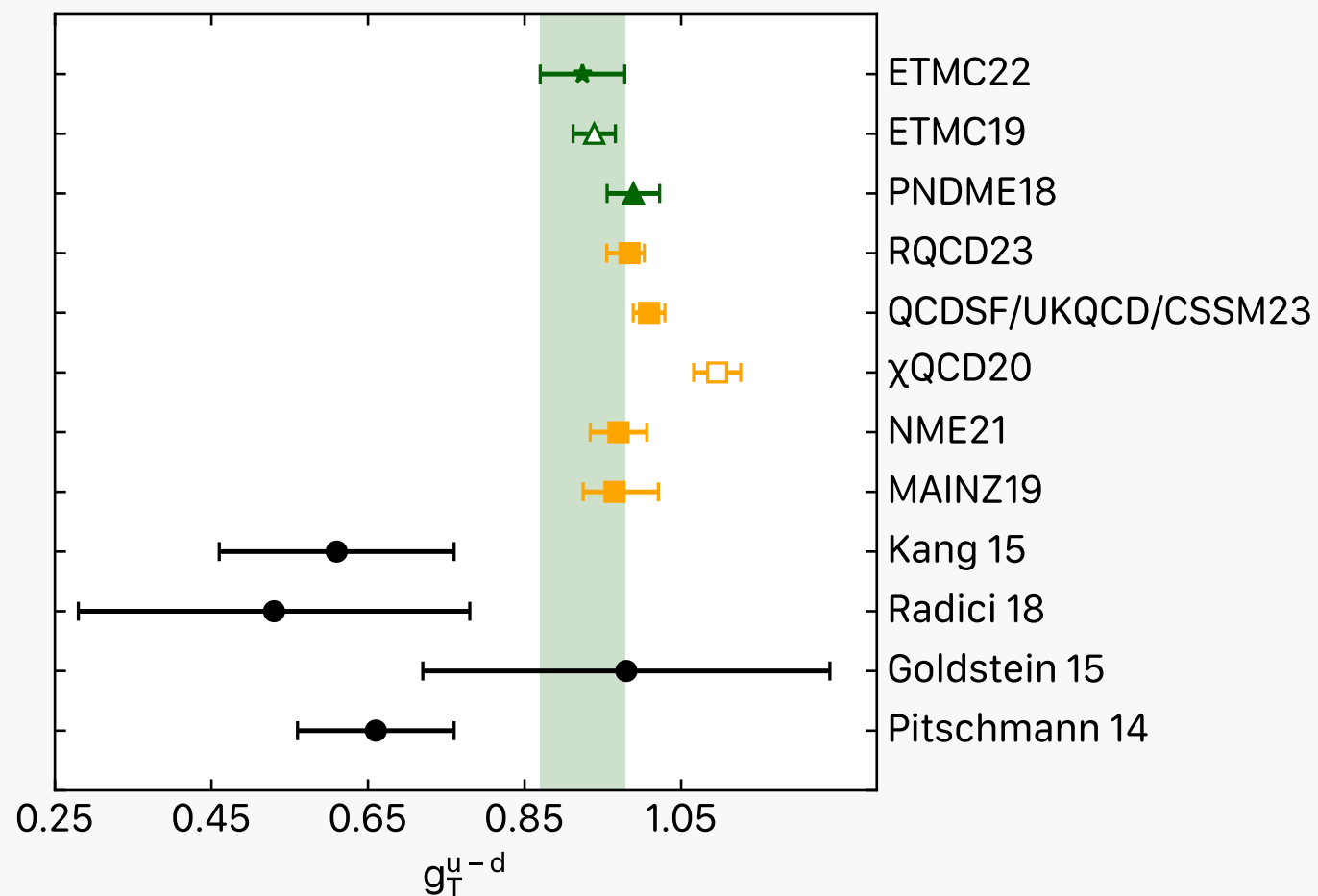
- Weaker dependence on lattice spacing
- Overall general agreement between lattice formulations

Tensor charge

- Tensor matrix element

$$g_T = \langle 1 \rangle_{\delta u - \delta d} \leftarrow \langle N | \bar{u} \sigma_{\mu\nu} u + \bar{d} \sigma_{\mu\nu} d | N \rangle$$

- Moment of transversity PDF; Can be used to constrain experimental analyses, e.g. JAM in [arXiv:2205.00999](https://arxiv.org/abs/2205.00999)



Preliminary

See: "First moments of the nucleon transverse quark spin densities using lattice QCD" [arXiv:2202.09871](https://arxiv.org/abs/2202.09871)

Summary & Outlook

- Lattice QCD with physical point ensembles at multiple lattice spacings
 - Here, three lattice spacings → Continuum limit directly at physical point
- Reproduction of well-known nucleon structure quantities, e.g. axial charge
- Requires thorough study of systematic uncertainties
 - Main systematic is excited state effects
- Agreement for axial charge leads to confidence in other less well known quantities
 - Flavor-separated axial charges
 - σ -terms
 - Tensor charges
- Analysis ongoing
 - For full systematic due to excited state effects in all charges and form factors
 - Additional lattice spacing at smaller values of a e.g. $a \simeq 0.05$ fm
 - Moments of PDFs, e.g. $\langle \chi \rangle_{u-d}$, $\langle \chi \rangle_{\Delta u - \Delta d}$, $\langle \chi \rangle_{\delta u - \delta d}$

Acknowledgements



Με τη συγχρηματοδότηση
της Ευρωπαϊκής Ένωσης



EXCELLENCE/0421/0043 EXCELLENCE/0421/0195