# **QCD** phase diagram from lattice simulations

#### A. Yu. Kotov







 $T \sim \sqrt[3]{n} \sim \Lambda$ 



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Quarks & Gluons

Color superconductor ?







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SKETCH!

Precise data: (mainly) lattice QCD in the blue region



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What is covered in this talk?

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Overview of the <u>lattice QCD</u> results on:

Color superconductor ?

 $\mu_B$ 



SKETCH!

Precise data: (mainly) lattice QCD in the blue region



What is covered in this talk?

Overview of the lattice QCD results on:

- Phase diagram of QCD in  $(T, \mu)$  plane
- Equation of State (EoS) and conserved charge fluctuations

**Color superconductor ?** 

 $\mu_B$ 





## Thermodynamics and lattice QCD **Grand-canonical partition function**

All field configurations

 $Z = \operatorname{tr} e^{-(H_{\text{QCD}} - \mu_B B - \mu_Q Q - \mu_S S)/T} \to \int DA_{\mu} D\bar{\psi} D\psi e^{-S_{\text{QCD}}} \to \operatorname{Computer}$ 



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 $T \neq 0$ : Compact time  $A_{\mu}(t + L_4, x, y, z) = A_{\mu}(t, x, y, z), \quad L_4 = 1/T$ 



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 $\rightarrow \int DA_{\mu}D\bar{\psi}D\psi e^{-S_{\rm QCD}} \rightarrow \text{Computer}$ 





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• First-principle method



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  - Infinite volume  $V \to \infty$ , continuum limit  $a \to 0$



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  - Solutions used in real simulations:
    - Taylor expansion in  $\mu_R$
    - Analytic continuation:  $\mu_B \rightarrow i \mu_B$





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Only work for small  $\mu_B/T \lesssim 2.5 - 3$ 





# Zero $\mu$ , nonzero T





## **Zero** $\mu$ , **nonzero** TNature of the chiral phase transition and (pseudo)-critical temperature $T_c$







# Zero $\mu$ , nonzero T







# [TWEXT, 2021]





<u>New Chiral-Spin</u>  $SU(2)_{CS}$ <u>symmetry</u>:  $T_{\rm ch} \lesssim T \lesssim 3T_{\rm ch}$ 





[Rohrhofer et al., 2020]



- IR phase:  $T \leq 250$  MeV [A.Alexandru and I. Horvath, 2019]

• Monopole condensation:  $T \leq 275$  MeV [M.Cardinali, M.D'Elia, A. Pasqui, 2021]

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# Zero $\mu$ , nonzero T



Possible transitions or thresholds:  $T \sim 2 - 3 T_c$ No consensus or full understanding of the physics

#### • Several (lattice) indications that regions $T \gtrsim T_c$ and $T \gg T_c$ are different



### **Small** $\mu$ , **nonzero** TTaylor expansion

$$T_{c}(\mu_{B}) = T_{0} \left( 1 - \kappa_{2} \left( \frac{\mu_{B}}{T_{0}} \right)^{2} - \kappa_{4} \left( \frac{\mu_{B}}{T_{0}} \right)^{4} \right), \quad T_{0} \equiv T_{c}(\mu_{B} = 0)$$



#### **Heavy Ion Collisions:**

 $S = 0, \quad Q \approx 0.4B \longrightarrow \mu_{Q,S} \equiv \mu_{Q,S}(\mu_B)$ 



### Small $\mu$ , nonzero T**Taylor expansion**

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[WB, 2020]

## Small $\mu$ , nonzero TEquation of State

• Pressure  $p(T, \mu_B)$ :

• Baryon density  $n_B = \frac{\partial p}{\partial \mu_B}$ • Entropy density  $s = \frac{\partial p}{\partial T}$ 

• Energy density:  $\epsilon = Ts - p + \mu_B B$ 



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EoS @  $\mu = 0$ : [WB, 2014] [HotQCD, 2014]









## Small $\mu$ , nonzero TEquation of State

4.5

3.5

1.5

0.5

120

(E) 2.5 t/d 2

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EoS @  $\mu = 0$ : [WB, 2014] [HotQCD, 2014]

#### Equation of State: • $\mu_B/T \lesssim 2.5$ for $T \lesssim 200$ MeV • $\mu_B/T \lesssim 3$ for $T \gtrsim 200$ MeV









# Small $\mu \rightarrow 0$ , nonzero T













• Taylor expansion: convergence radius in  $\mu_B$ 

$$\chi_2^B(T,\mu_B) = \sum_{n=0}^{\infty} \frac{1}{(2n)!} \chi_{2n+2}^B(\mu_B/T)^{2n} \qquad r_{2n}^{\chi} = \left| \frac{2n(2n-1)\chi_{2n}^B}{\chi_{2n+2}^B} \right|^{\frac{1}{2n}}$$



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- Transition width from  $\mu_R^2 < 0$  to  $\mu_R^2 > 0$







### Lar CEP

• Taylc

• Lee-

• Trans



[WB, 2020]





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- Understanding properties (phases, thresholds) in QCD at finite T•
- Higher order fluctuations  $\chi_{B,O,S}^n$
- Improving CEP estimations
- Approach to solve or overcome the sign problem
  - Reweighting technique [WB, 2020]
  - Complex Langevin
  - Lefschetz thimbles
  - Density of states lacksquare
  - Machine Learning

- [M. Cristoforetti et al., 2012]
- [K. Langfeld et al., 2012]



[S. Lawrence, Y. Yamauchi et al., 2012]



## Conclusions **QCD** phase diagram on the lattice

- Precise data on the crossover at zero and small  $\mu_R$
- Non-trivial properties (phases) of finite-T QCD





• EoS of QCD at  $\mu_R/T \lesssim (2.5 - 3)$ , conserved charge fluctuations at  $\mu_R = 0$ 



# BACKUP

<u>Spectrum of the Dirac operator:</u>

$$D | \lambda \rangle = \lambda | \lambda \rangle$$

Density  $\rho(\lambda)$  contains a lot of information about the system and its symmetries

**Banks-Casher relation:** 

$$\frac{1}{V} \langle \bar{\psi} \psi \rangle \to \int_0^\infty \rho(\lambda) \frac{m}{\lambda^2 + m^2} \to \pi \rho(0)$$

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$$\frac{n \rho n \alpha \sigma \sigma}{\rho(\lambda) \sim \lambda^{-1+\delta}}$$

IR phase

Can be described by topological fluctuations ()







## **Complex Langevin method and Lefschetz thimbles**

- Complexify:  $x \to z \in \mathbb{C}$
- Complex Langevin:  $\dot{z} = -\partial_{z}S(z) + \eta$
- lacksquarecontour to get milder sign problem (Machine Learning)

$$Z = \int dx e^{-S(x)} dx e^{-S(x)}$$

Lefschetz thimbles: instead of integral along real axis  $z = x \in \mathbb{R}$ , deform a



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# **Chiral-spin symmetry**

• 
$$\psi \to \exp\left(i\frac{\epsilon^n \Sigma^n}{2}\right)\psi$$
  $\Sigma_n = \{\gamma_k, -i\gamma_5\gamma_k, \gamma_5\}$ 

• Symmetry of the chromo-electric color charge

$$Q^a = \int d^3x q^{\dagger}(x)$$

Larger symmetry then in the free theory(!) 

 $T^a q(x)$ 

$$(0,0) \qquad \overline{\Psi} \begin{pmatrix} f_{1}(0,1^{++}) \\ \mathbb{I}_{F} \otimes \gamma^{5} \gamma^{i} \end{pmatrix} \Psi \qquad \overline{\Psi} \begin{pmatrix} 0,1^{--} \\ \mathbb{I}_{F} \otimes \gamma^{i} \end{pmatrix} \Psi$$

$$(1/2,1/2)_{a} \qquad \qquad \overline{\Psi} \begin{pmatrix} f_{1}(1,1^{+-}) \\ \mathbb{V}(1)_{A} \end{pmatrix} \stackrel{b_{1}(1,1^{+-}) \\ \mathbb{V}(\tau^{a} \otimes \gamma^{5} \gamma^{4} \gamma^{i}) \Psi & \longrightarrow \overline{\Psi} \begin{pmatrix} 0,1^{--} \\ \mathbb{I}_{F} \otimes \gamma^{4} \gamma^{i} \end{pmatrix} \Psi \qquad \qquad \overline{\Psi} \begin{pmatrix} 0,1^{--} \\ \mathbb{V}(1_{F} \otimes \gamma^{5} \gamma^{4} \gamma^{i}) \Psi \\ \mathbb{V}(1)_{A} \end{pmatrix} \stackrel{c}{\Psi} \begin{pmatrix} \rho(1,1^{--}) \\ \mathbb{V}(\tau^{a} \otimes \gamma^{i}) \Psi & \longrightarrow \overline{\Psi} \begin{pmatrix} 0,1^{--} \\ \mathbb{V}(\tau^{a} \otimes \gamma^{5} \gamma^{i}) \Psi \\ \mathbb{V}(1_{F} \otimes \gamma^{5} \gamma^{i}) \Psi \end{pmatrix} \qquad \qquad U(1)$$

$$(1/2,1/2)_{b} \qquad \qquad \overline{\Psi} \begin{pmatrix} f_{1}(0,1^{++}) \\ \mathbb{V}(\tau^{a} \otimes \gamma^{5} \gamma^{i}) \Psi \\ \mathbb{V} \begin{pmatrix} 0,0 \end{pmatrix} & \overline{\Psi} \begin{pmatrix} f_{1}(0,1^{++}) \\ \mathbb{V}(\tau^{a} \otimes \gamma^{5} \gamma^{i}) \Psi \\ \mathbb{V} \begin{pmatrix} 0,0 \end{pmatrix} & \overline{\Psi} \begin{pmatrix} f_{1}(0,1^{++}) \\ \mathbb{V}(\tau^{a} \otimes \gamma^{5} \gamma^{i}) \Psi \\ \mathbb{V} \begin{pmatrix} 0,0 \end{pmatrix} & \overline{\Psi} \begin{pmatrix} f_{1}(1,1^{+-}) \\ \mathbb{V}(\tau^{a} \otimes \gamma^{5} \gamma^{i} \gamma^{i}) \Psi \\ \mathbb{V} \begin{pmatrix} 0,0 \end{pmatrix} & \overline{\Psi} \begin{pmatrix} f_{1}(1,1^{+-}) \\ \mathbb{V}(\tau^{a} \otimes \gamma^{4} \gamma^{i}) \Psi \\ \mathbb{V} \begin{pmatrix} 0,0 \end{pmatrix} & \overline{\Psi} \begin{pmatrix} 0,1^{--} \\ \mathbb{V} \begin{pmatrix} 0,0 \end{pmatrix} & \overline{\Psi} \begin{pmatrix} 0,1^{--} \\ \mathbb{V} \begin{pmatrix} 0,0 \end{pmatrix} & \overline{\Psi} \begin{pmatrix} 0,1^{--} \\ \mathbb{V} \begin{pmatrix} 0,0 \end{pmatrix} & \overline{\Psi} \begin{pmatrix} 0,1^{--} \\ \mathbb{V} \begin{pmatrix} 0,0 \end{pmatrix} & \overline{\Psi} \begin{pmatrix} 0,1^{--} \\ \mathbb{V} \begin{pmatrix} 0,0 \end{pmatrix} & \overline{\Psi} \begin{pmatrix} 0,1^{--} \\ \mathbb{V} \begin{pmatrix} 0,0 \end{pmatrix} & \overline{\Psi} \begin{pmatrix} 0,1^{--} \\ \mathbb{V} \begin{pmatrix} 0,0 \end{pmatrix} & \overline{\Psi} \begin{pmatrix} 0,1^{--} \\ \mathbb{V} \begin{pmatrix} 0,0 \end{pmatrix} & \overline{\Psi} \begin{pmatrix} 0,1^{--} \\ \mathbb{V} \begin{pmatrix} 0,0 \end{pmatrix} & \overline{\Psi} \begin{pmatrix} 0,1^{--} \\ \mathbb{V} \begin{pmatrix} 0,0 \end{pmatrix} & \overline{\Psi} \end{pmatrix} \end{pmatrix}$$

[L. Glozman, 2022]





# Lee-Yang zeros

- $Z \equiv Z(V, T, \mu_R)$  is real for  $\mu_R \in \mathbb{R}$
- $Z \equiv Z(V, T, \mu_R)$  can have complex  $\mu_R^i$  roots
- Lee-Yang theorem: phase transition when  $\mu_B^i \to \mu_B^c \in \mathbb{R}$ [Lee, Yang, 1952]
- Lattice data:  $B(\mu_B) \sim \frac{\partial \log Z}{\partial \mu_B}$  at  $-i\mu_B^0 \in \mathbb{R}$  or  $\chi_n^B$  at  $\mu_B = 0$
- Reconstruct Pade approximation of  $B(\mu_B)$  for all  $\mu_B$
- Study its zeros and when they approach real axis



[Bielefeld-Parma, 2023]





# Reweighting for the sign problem

• 
$$p \sim e^{-S_G} \det D = e^{-S_G} |\det D| e^{i\theta}$$
  
•  $p \sim e^{-S_G} \operatorname{Re} \det D = e^{-S_G} |\operatorname{Re} \det D| \operatorname{sign}(\operatorname{Re} \det D)$ 

Weight

• 
$$\int dUp[U]O[U] \to \int dUp_w[U]p_O[U]$$

Issues with staggered fermions 

#### Observable

[U]O[U]



[WB, 2021]

