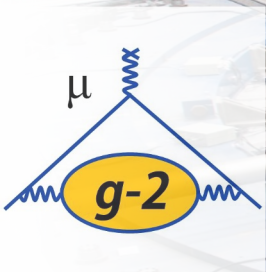


Recent measurement and experimental prospects of the Muon $g-2$ experiment at Fermilab

E. Bottalico

EINN 2023

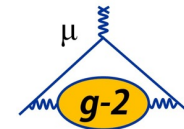
31 October 2023



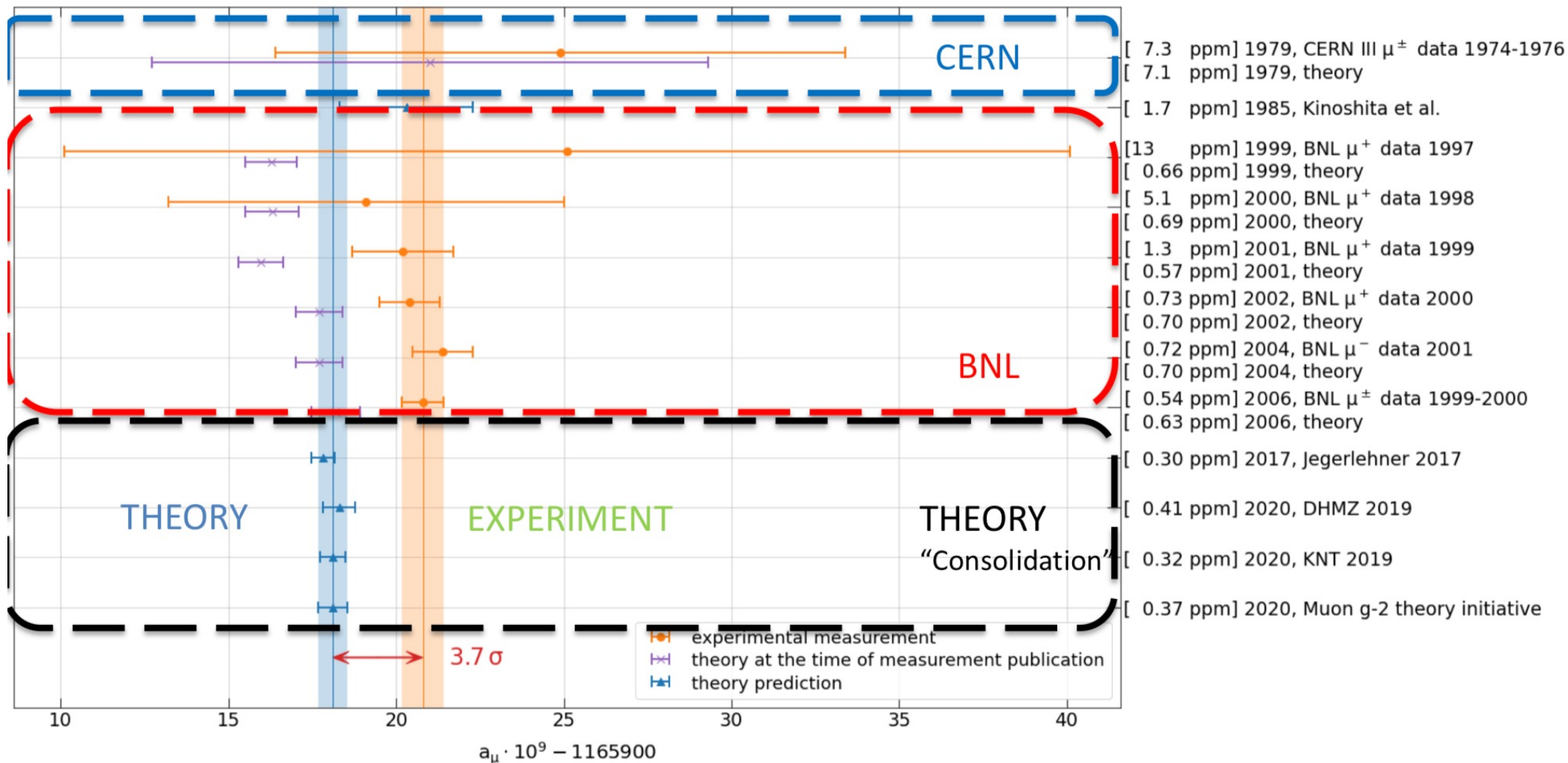
EINN2023



The history of Muon g-2

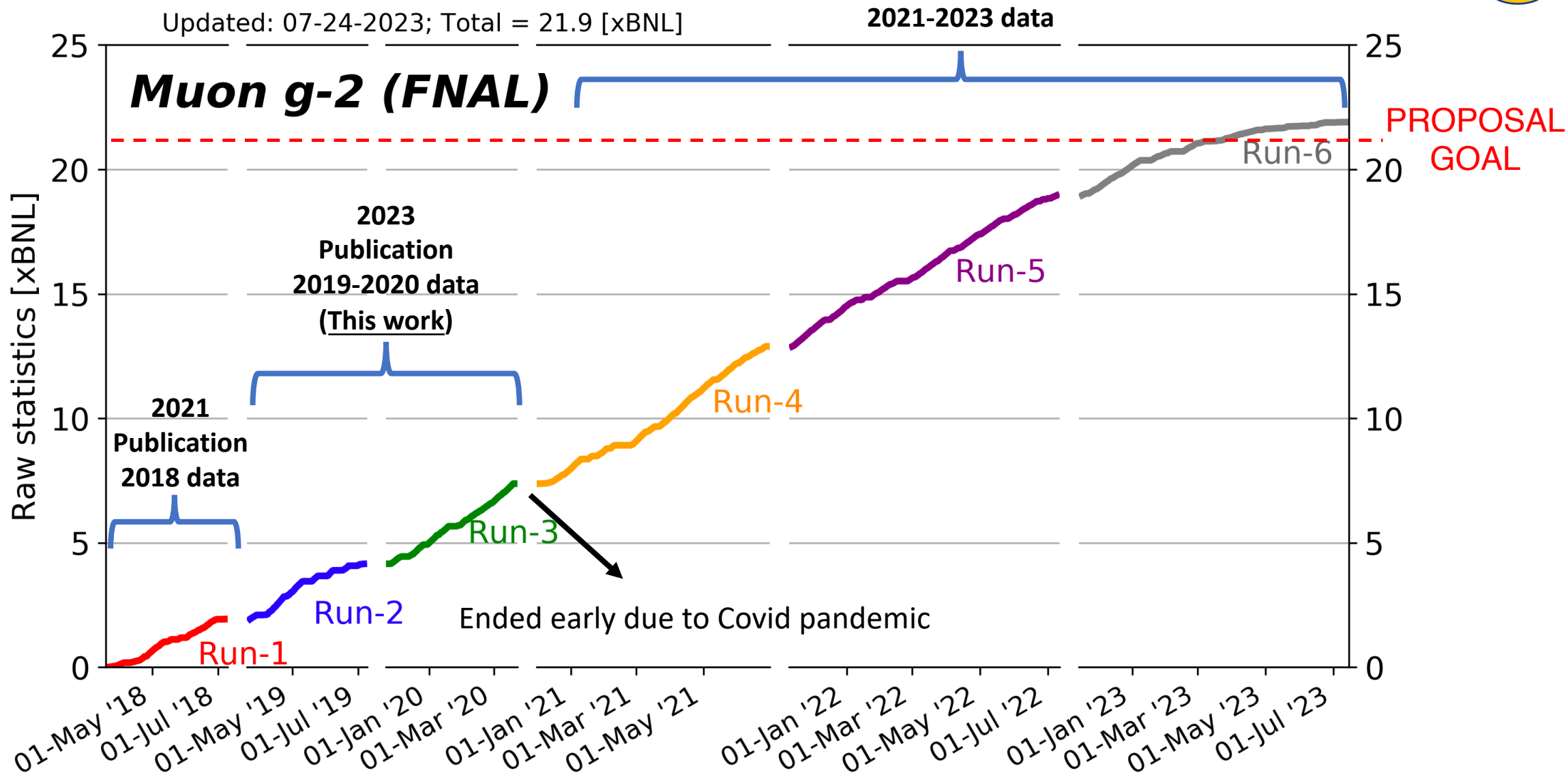
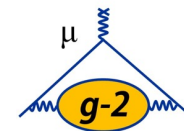


The history of the Muon g-2 experiments finds its roots in the series of experiment at CERN





Muon g-2 today: Data collection





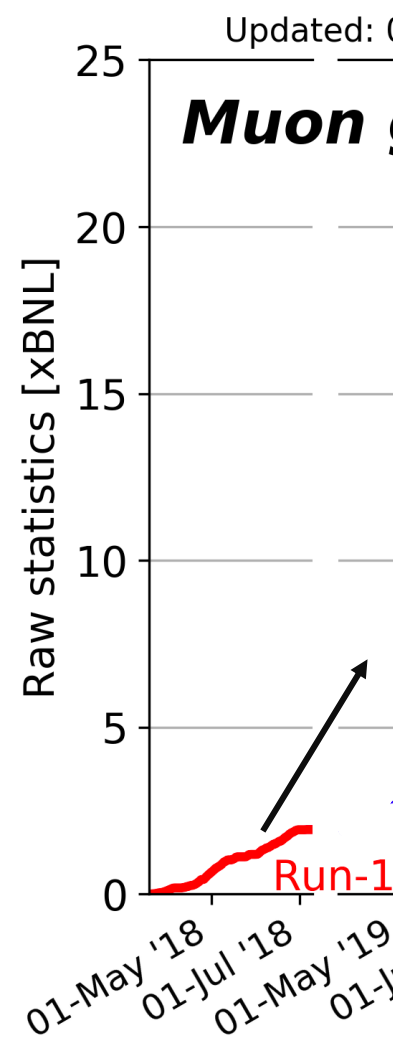
Run-1 Publication



Published on PRL – 7th April 2021

Measurement of the Positive Muon Anomalous Magnetic Moment to 0.46 ppm

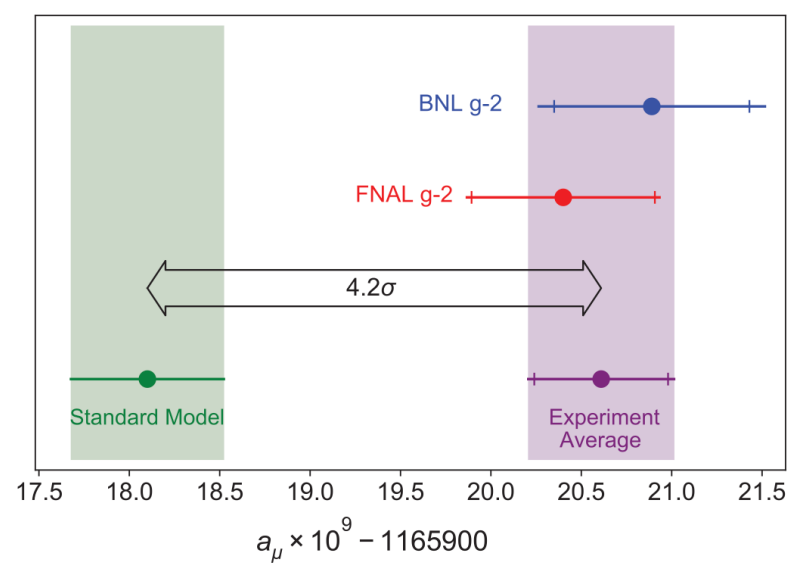
B. Abi,⁴⁴ T. Albahri,³⁹ S. Al-Kilani,³⁶ D. Allspach,⁷ L. P. Alonzi,⁴⁸ A. Anastasi,^{11,a} A. Anisenkov,^{4,b} F. Azfar,⁴⁴ K. Badgley,⁷ S. Baeßler,^{47,c} I. Bailey,^{19,d} V. A. Baranov,¹⁷ E. Barlas-Yucel,³⁷ T. Barrett,⁶ E. Barzi,⁷ A. Basti,^{11,32} F. Bedeschi,¹¹ A. Behnke,²² M. Berz,²⁰ M. Bhattacharya,⁴³ H. P. Binney,⁴⁸ R. Bjorkquist,⁶ P. Bloom,²¹ J. Bono,⁷ E. Bottalico,^{11,32} T. Bowcock,³⁹ D. Boyden,²² G. Cantatore,^{13,34} R. M. Carey,² J. Carroll,³⁹ B. C. K. Casey,⁷ D. Cauz,^{35,8} S. Ceravolo,⁹ R. Chakraborty,³⁸ S. P. Chang,^{18,5} A. Chapelain,⁶ S. Chappa,⁷ S. Charity,⁷ R. Chislett,³⁶ J. Choi,⁵ Z. Chu,^{26,e} T. E. Chupp,⁴² M. E. Convery,⁷ A. Conway,⁴¹ G. Corradi,⁹ S. Corrodi,¹ L. Cotrozzi,^{11,32} J. D. Crnkovic,^{3,37,43} S. Dabagov,^{9,f} P. M. De Lurgio,¹ P. T. Debevec,³⁷ S. Di Falco,¹¹ P. Di Meo,¹⁰ G. Di Sciascio,¹² R. Di Stefano,^{10,30} B. Drendel,⁷ A. Driutti,^{35,13,38} V. N. Duginov,¹⁷ M. Eads,²² N. Eggert,⁶ A. Epps,²² J. Esquivel,⁷ M. Farooq,⁴² R. Fatemi,³⁸ C. Ferrari,^{11,14} M. Fertl,^{48,16} A. Fiedler,²² A. T. Fienberg,⁴⁸ A. Fioretti,^{11,14} D. Flay,⁴¹ S. B. Foster,² H. Friedsam,⁷ E. Frlež,⁴⁷ N. S. Froemming,^{48,22} J. Fry,⁴⁷ C. Fu,^{26,e} C. Gabbanini,^{11,14} M. D. Galati,^{11,32} S. Ganguly,^{37,7} A. Garcia,⁴⁸ D. E. Gastler,² J. George,⁴¹ L. K. Gibbons,⁶ A. Gioiosa,^{29,11} K. L. Giovanetti,¹⁵ P. Girotti,^{11,32} W. Gohn,³⁸ T. Gorringer,³⁸ J. Grange,^{1,42} S. Grant,³⁶ F. Gray,²⁴ S. Haciomeroglu,⁵ D. Hahn,⁷ T. Halewood-Leagas,³⁹ D. Hampai,⁹ F. Han,³⁸ E. Hazen,² J. Hempstead,⁴⁸ S. Henry,⁴⁴ A. T. Herrod,^{39,d} D. W. Hertzog,⁴⁸ G. Hesketh,³⁶ A. Hibbert,³⁹ Z. Hodge,⁴⁸ J. L. Holzbauer,⁴³ K. W. Hong,⁴⁷ R. Hong,^{1,38} M. Iacovacci,^{10,31} M. Incagli,¹¹ C. Johnstone,⁷ J. A. Johnstone,⁷ P. Kammel,⁴⁸ M. Kargiantoulakis,⁷ M. Karuza,^{13,45} J. Kaspar,⁴⁸ D. Kawall,⁴¹ L. Kelton,³⁸ A. Keshavarzi,⁴⁰ D. Kessler,⁴¹ K. S. Khaw,^{27,26,48,e} Z. Khechadorian,⁶ N. V. Khomutov,¹⁷ B. Kiburg,⁷ M. Kiburg,^{7,21} O. Kim,^{18,5} S. C. Kim,⁶ Y. I. Kim,⁵ B. King,^{39,a} N. Kinnaird,² M. Korostelev,^{19,d} I. Kourbanis,⁷ E. Kraegeloh,⁴² V. A. Krylov,¹⁷ A. Kuchibhotla,³⁷ N. A. Kuchinskiy,¹⁷ K. R. Labe,⁶ J. LaBounty,⁴⁸ M. Lancaster,⁴⁰ M. J. Lee,⁵ S. Lee,⁵ S. Leo,³⁷ B. Li,^{26,1,e} D. Li,^{26,g} L. Li,^{26,e} I. Logashenko,^{4,b} A. Lorente Campos,³⁸ A. Lucà,⁷ G. Lukicov,³⁶ G. Luo,²² A. Lusiani,^{11,25} A. L. Lyon,⁷ B. MacCoy,⁴⁸ R. Madrak,⁷ K. Makino,²⁰ F. Marignetti,^{10,30} S. Mastroianni,¹⁰ S. Maxfield,³⁹ M. McEvoy,²² W. Merritt,⁷ A. A. Mikhailichenko,^{6,a} J. P. Miller,² S. Miozzi,¹² J. P. Morgan,⁷ W. M. Morse,³ J. Mott,^{2,7} E. Motuk,³⁶ A. Nath,^{10,31} D. Newton,^{39,a,d} H. Nguyen,⁷ M. Oberling,¹ R. Osofsky,⁴⁸ J.-F. Ostiguy,⁷ S. Park,⁵ G. Pauletta,^{35,8} G. M. Piacentino,^{29,12} R. N. Pilato,^{11,32} K. T. Pitts,³⁷ B. Plaster,³⁸ D. Počanić,⁴⁷ N. Pohlman,²² C. C. Polly,⁷ M. Popovic,⁷ J. Price,³⁹ B. Quinn,⁴³ N. Raha,¹¹ S. Ramachandran,¹ E. Ramberg,⁷ N. T. Rider,⁶ J. L. Ritchie,⁴⁶ B. L. Roberts,² D. L. Rubin,⁶ L. Santi,^{35,8} D. Sathyan,² H. Schellman,^{23,h} C. Schlesier,³⁷ A. Schreckenberger,^{46,2,37} Y. K. Semertzidis,^{5,18} Y. M. Shatunov,⁴ D. Shemyakin,^{4,b} M. Shenk,²² D. Sim,³⁹ M. W. Smith,^{48,11} A. Smith,³⁹ A. K. Soha,⁷ M. Sorbara,^{12,33} D. Stöckinger,²⁸ J. Stapleton,⁷ D. Still,⁷ C. Stoughton,⁷ D. Stratakis,⁷ C. Strohmaier,⁶ T. Stuttard,³⁶ H. E. Swanson,⁴⁸ G. Sweetmore,⁴⁰ D. A. Sweigart,⁶ M. J. Syphers,^{22,7} D. A. Tarazona,²⁰ T. Teubner,³⁹ A. E. Tewsley-Booth,⁴² K. Thomson,³⁹ V. Tishchenko,³ N. H. Tran,² W. Turner,³⁹ E. Valetov,^{20,19,27,d} D. Vasilkova,³⁶ G. Venanzoni,¹¹ V. P. Volnykh,¹⁷ T. Walton,⁷ M. Warren,³⁶ A. Weisskopf,²⁰ L. Welty-Rieger,⁷ M. Whitley,³⁹ P. Winter,¹ A. Wolski,^{39,d} M. Wormald,³⁹ W. Wu,⁴³ and C. Yoshikawa⁷



31/10/23

(Muon g-2 Collaboration)

Quantity	Correction terms (ppb)	Uncertainty (ppb)
ω_a^m (statistical)	...	434
ω_a^m (systematic)	...	56
C_e	489	53
C_p	180	13
C_{ml}	-11	5
C_{pa}	-158	75
$f_{\text{calib}} \langle \omega_p(x, y, \phi) \times M(x, y, \phi) \rangle$...	56
B_k	-27	37
B_q	-17	92
$\mu'_p(34.7^\circ)/\mu_e$...	10
m_μ/m_e	...	22
$g_e/2$...	0
Total systematic	...	157
Total fundamental factors	...	25
Totals	544	462





What is the Muon g-2?

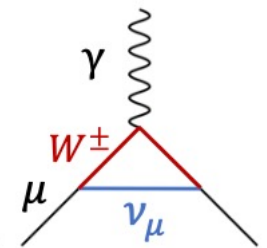
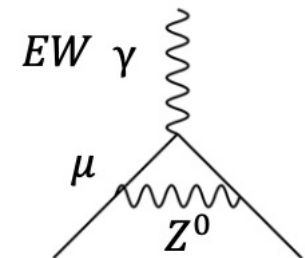
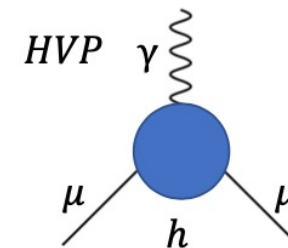
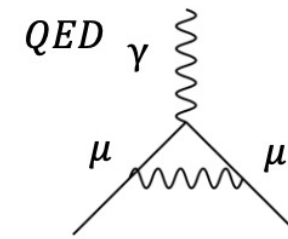


The intrinsic magnetic moment of a particle with spin is:

$$\vec{\mu} = g \frac{q}{2m} \vec{S}$$

The g-factor (gyromagnetic) defines the coupling between the spin and the magnetic field:

- $g = 1$ classic theory;
- $g = 2$ Dirac quantum theory;
- $g = 2.00233\dots$ quantum field theory.



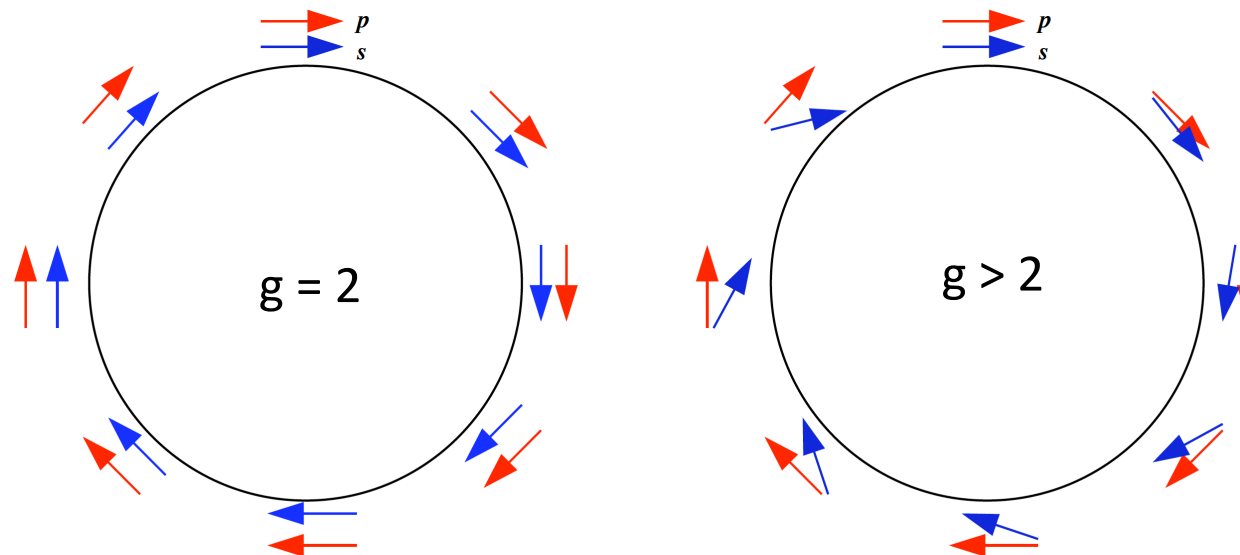


How do we measure a_μ ?



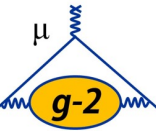
- Muons are injected in a magnetic field of **1.45 T**
- When a muon moves in a uniform magnetic field its spin rotates faster than its momentum (if **$g > 2$**) and the difference between the spin and momentum angular velocity is the anomalous precession velocity $\vec{\omega}_a$:

$$\vec{\omega}_a = \vec{\omega}_s - \vec{\omega}_c = g \frac{e\vec{B}}{2m} - \frac{e\vec{B}}{m} = \left(\frac{g-2}{2} \right) \cdot \frac{e\vec{B}}{m} = a_\mu \frac{e\vec{B}}{m}$$





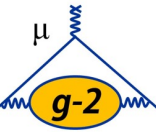
How such precision is possible?



- 4 nature gifts allow to reach this very high precision:



How such precision is possible?



- 4 nature gifts allow to reach this very high precision:

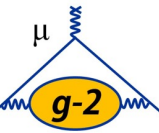
1. *Muons strongly polarized (95%):*

- It is possible thanks to the weak pion decay





How such precision is possible?



- 4 nature gifts allow to reach this very high precision:

- Muons strongly polarized (95%):*

- It is possible thanks to the weak pion decay



- Precession frequency proportional to (g-2)*

- $$\vec{\omega}_a = \vec{\omega}_S - \vec{\omega}_c = \left(\frac{g-2}{2}\right) \cdot \frac{e\vec{B}}{m}$$



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2. Precession frequency proportional to (g-2)

- $$\vec{\omega}_a = \vec{\omega}_S - \vec{\omega}_c = \left(\frac{g-2}{2}\right) \cdot \frac{e\vec{B}}{m}$$

3. Magic momentum $P_\mu = 3.094 \text{ GeV}/c$:

- $$\vec{\omega}_a = \frac{e}{m} \left[a_\mu \vec{B} - \left(a_\mu - \frac{1}{\gamma^2 - 1} \right) (\vec{\beta} \times \vec{E}) \right]$$



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- Magic momentum $P_\mu = 3.094 \text{ GeV}/c$: $\gamma = \sqrt{1 + \frac{1}{a_\mu}} \sim 29.3$*

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How such precision is possible?



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➤ It is possible thanks to the weak pion decay



- 2. Precession frequency proportional to (g-2)**

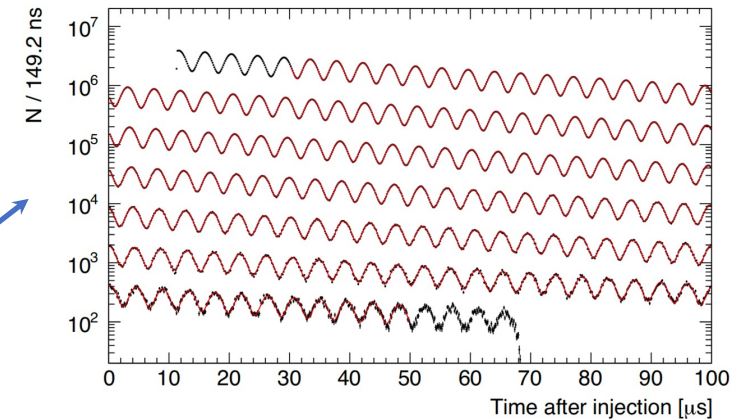
➤
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$$\vec{\omega}_a = \frac{e}{m} \left[a_\mu \vec{B} - \left(a_\mu - \frac{1}{\gamma^2 - 1} \right) (\vec{\beta} \times \vec{E}) \right]$$

- 4. Positron emitted preferably in direction of the muon spin**





a_μ Extraction (simplified)



- The magnetic anomaly a_μ is obtained by measuring many quantities:

$$a_\mu = \frac{\omega_a}{\tilde{\omega}'_p} \frac{\mu_p m_\mu g_e}{\mu_e m_e 2}$$

Ratio measured at Muon g-2 experiment:

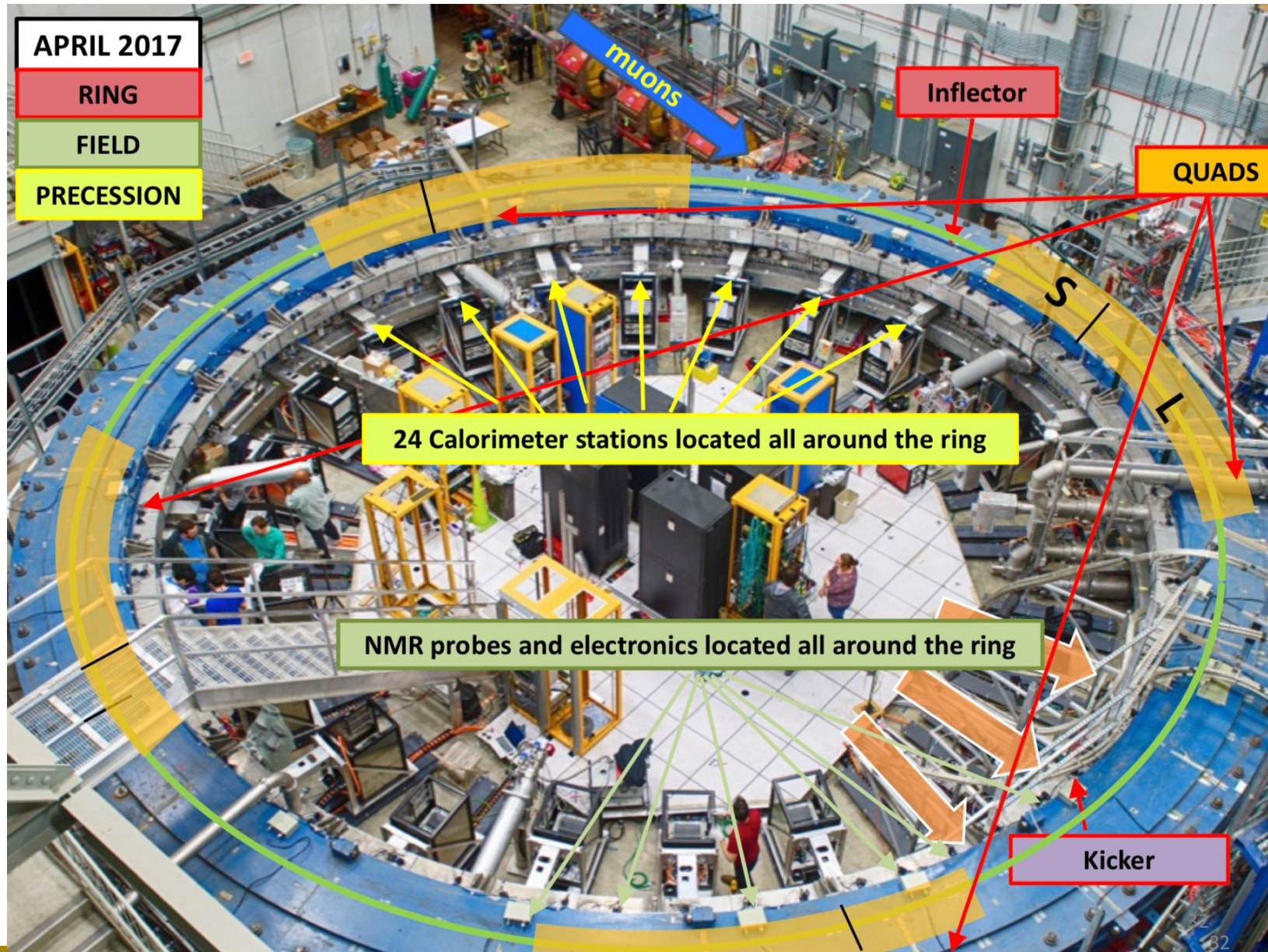
$$R' = \frac{\omega_a}{\tilde{\omega}'_p}$$

Quantities measured in other experiments

$\tilde{\omega}'_p$ represents the magnetic field as proton Larmor frequency (in a water sample) weighted with the beam distribution inside the ring.

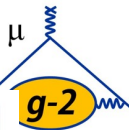


$g-2$ short recap: The Ring

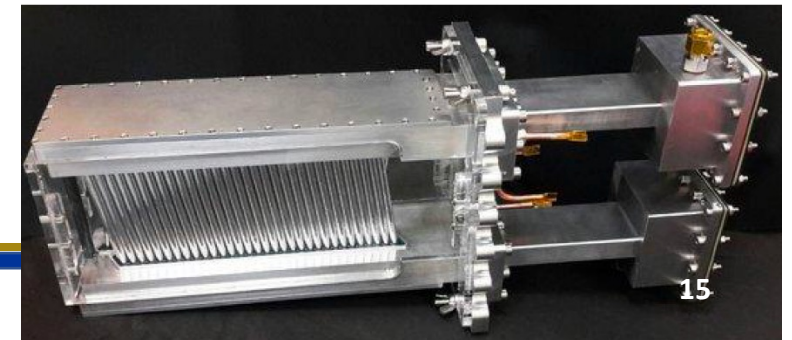
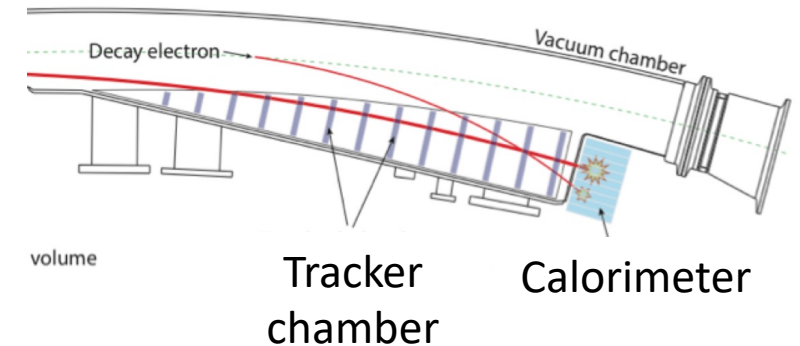
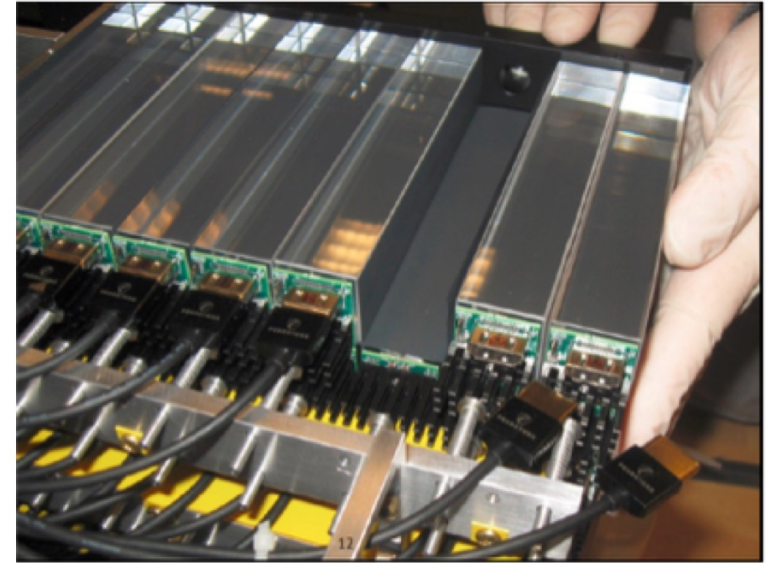




Detectors

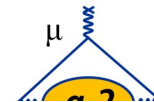


- **24 electromagnetic calorimeters:**
 - 54 PbF2 crystals read by 54 SIPMs.
 - Crystal length 14 cm, $15 X_0$.
 - Cherenkov light faster than showers (signal width \sim nanoseconds).
 - **Laser calibration system**, allows the energy and time calibration of the calorimeters
- **Two straw tubes trackers.**
 - 32 planes of drift tubes filled with a 50:50 mixture of Ar/Ethane.





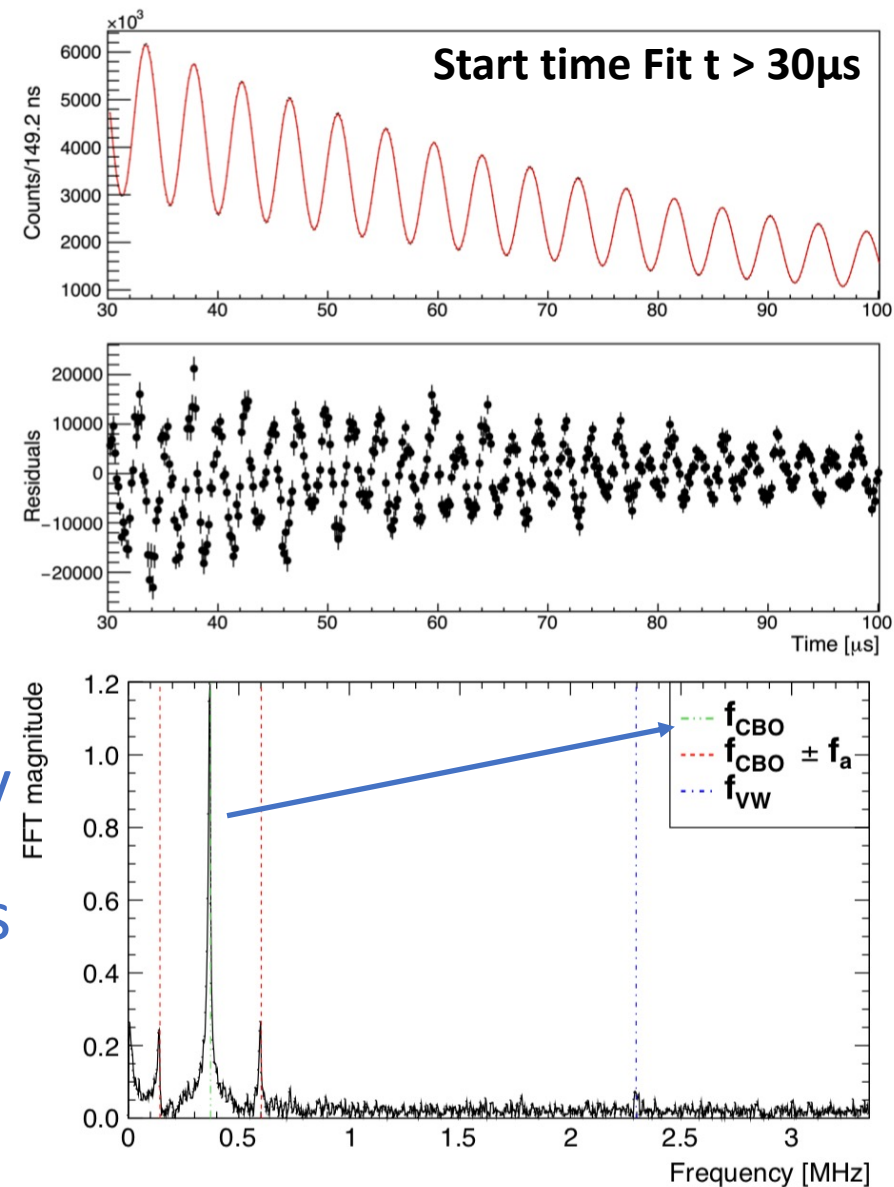
ω_a measurement



- The simplest function which describes the number of emitted positron from muon decay (so called “wiggle plot”) is:

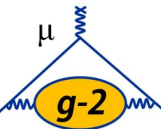
$$N(t) = N_0 \cdot e^{-\frac{t}{\tau}} \cdot (1 + A \cdot \cos(\omega_a \cdot t + \varphi))$$

- The **FFT** of the fit’s residual shows many frequency peaks due to beam dynamics effects that are not modeled by the previous function.

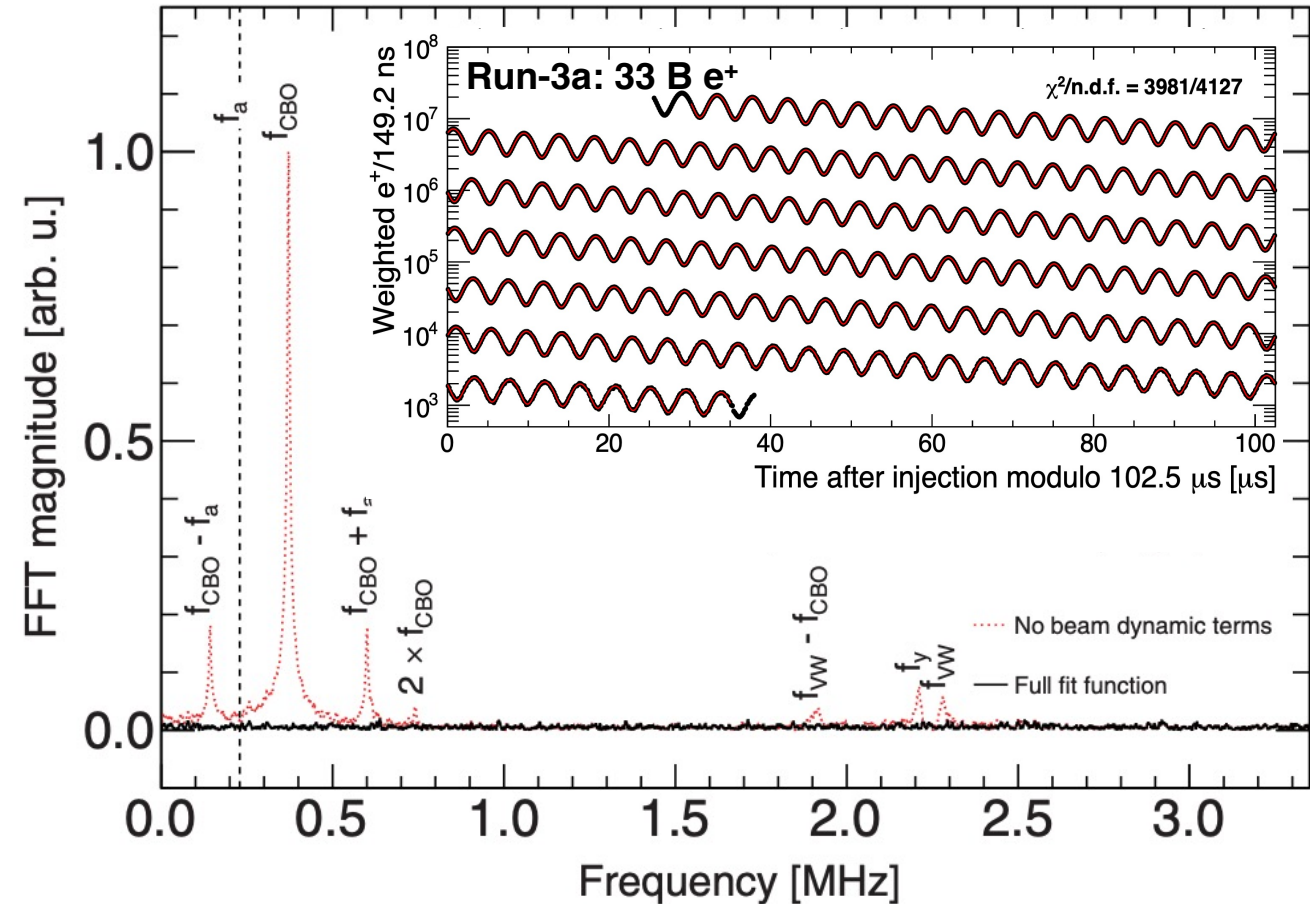
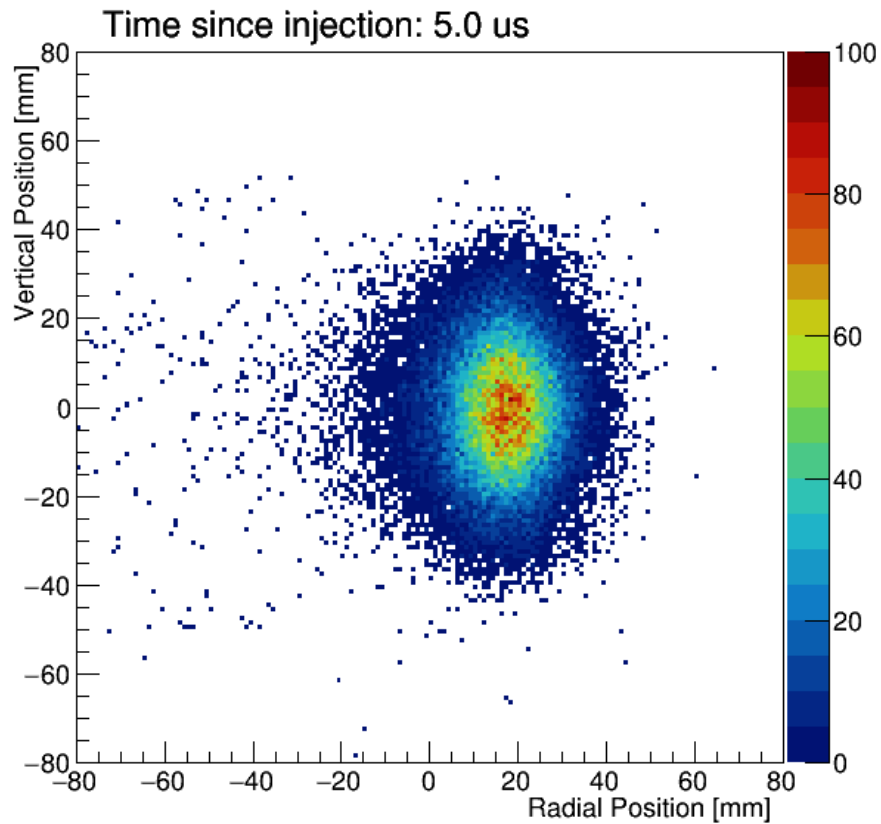




ω_a measurement

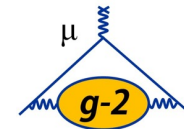


- Taking into account for the beam motion, the fit function gets more complicated up to contain up to 27 parameters.

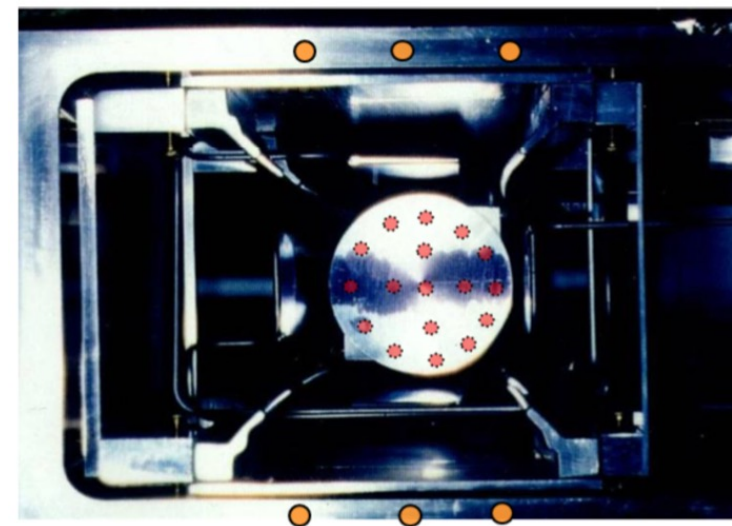




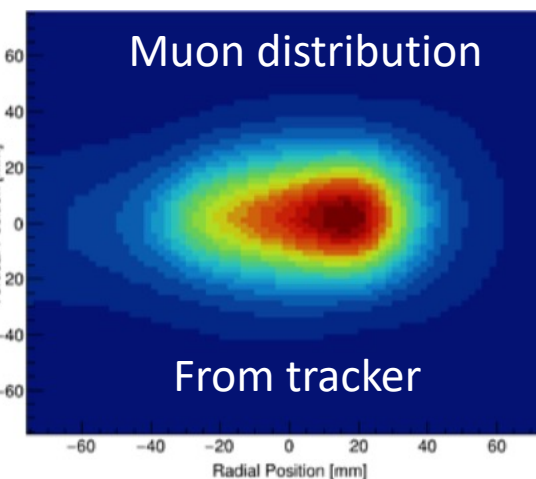
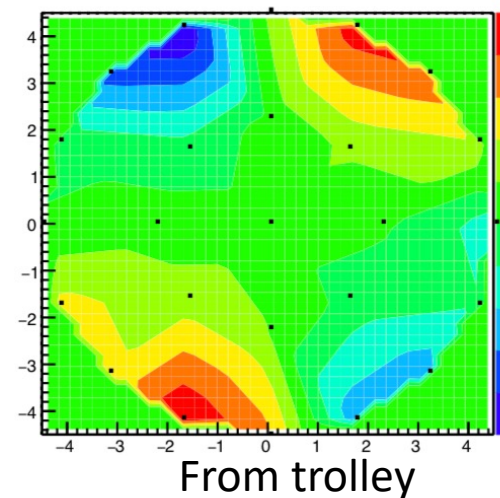
$\tilde{\omega}'_p$ Measurement



- The magnetic field is measured by:
 - 378 fixed probes around the ring;
 - 17 NMR probes moved around the ring via a trolley.
- The tracker measures the muon distribution around the ring.
- The magnetic field map is weighted with the muon distribution to obtain the effective field experienced by muons.

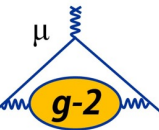


Field map





a_μ Extraction – Real world



For the measurement of a_μ the measured ω_a and ω_p need to be corrected by:

Beam dynamics corrections

$$R'_\mu \approx \frac{f_{clock} \omega_a^m (1 + C_e + C_p + C_{ml} + C_{pa} + C_{dd})}{f_{calib} \langle \omega'_p(x, y, \phi) \times M(x, y, \phi) \rangle (1 + B_k + B_q)}$$

Transient field corrections

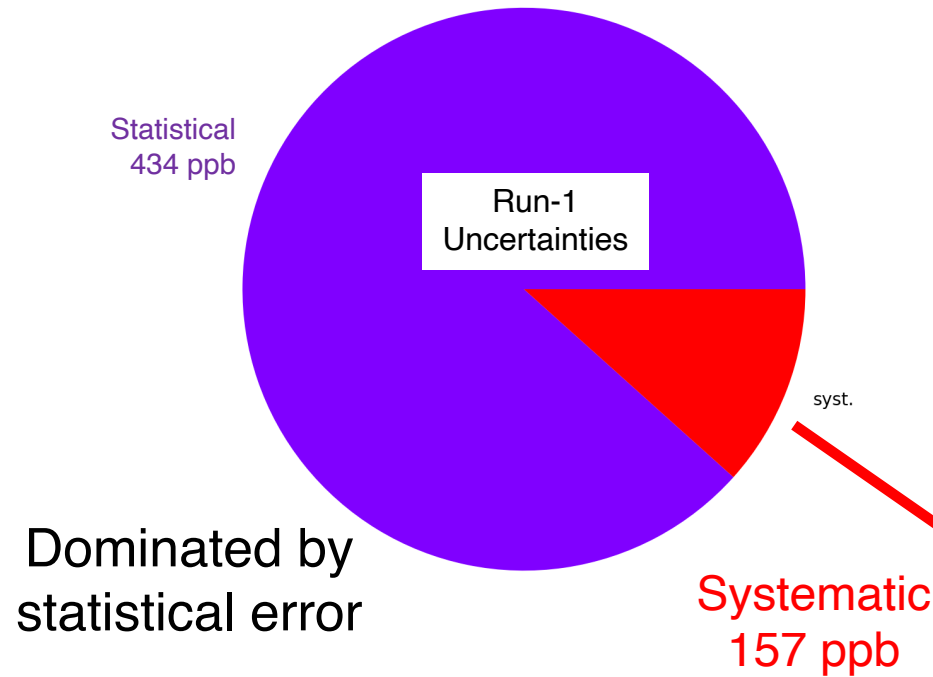
The total size of the corrections in **Run-2/3** is **622 ppb** dominated by C_e and C_p .
Corrections are small, but dominated Run-1 systematics...



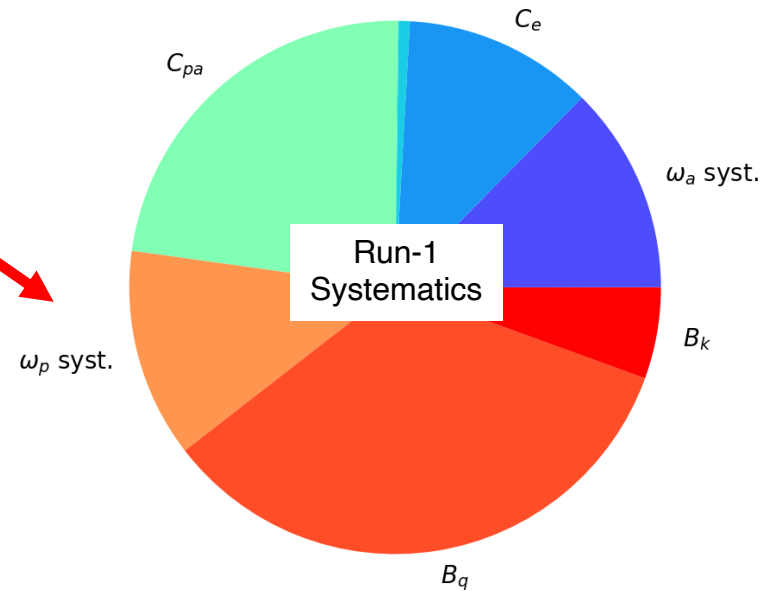
Run-1 systematic: Where can we improve?



Total Uncertainty: 462 ppb



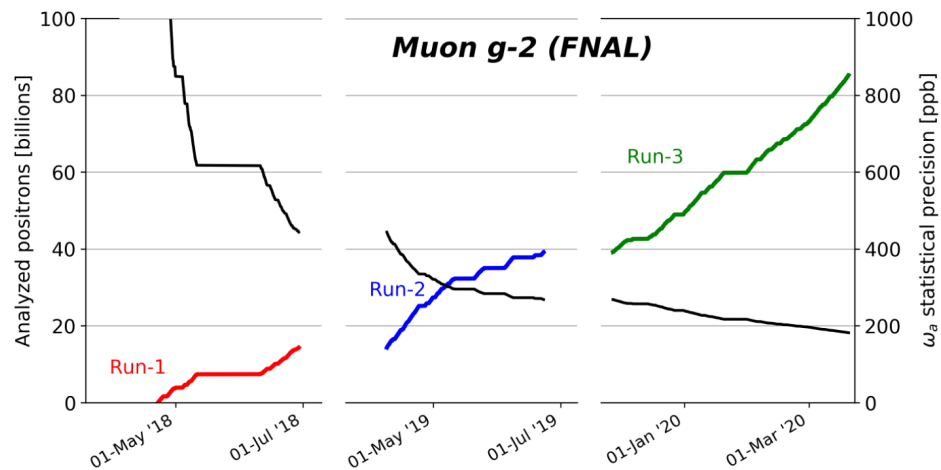
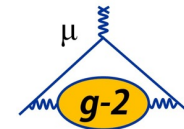
Total Systematic: 157 ppb



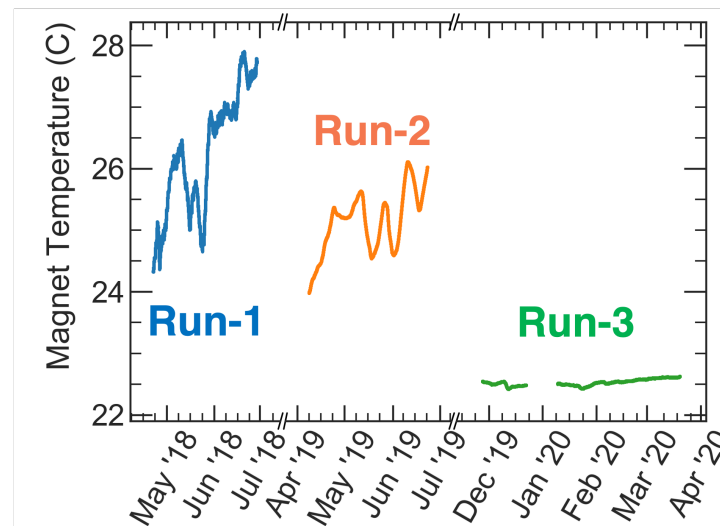
Target: quadrupole transient field B_q
and phase-acceptance effect C_{pa}



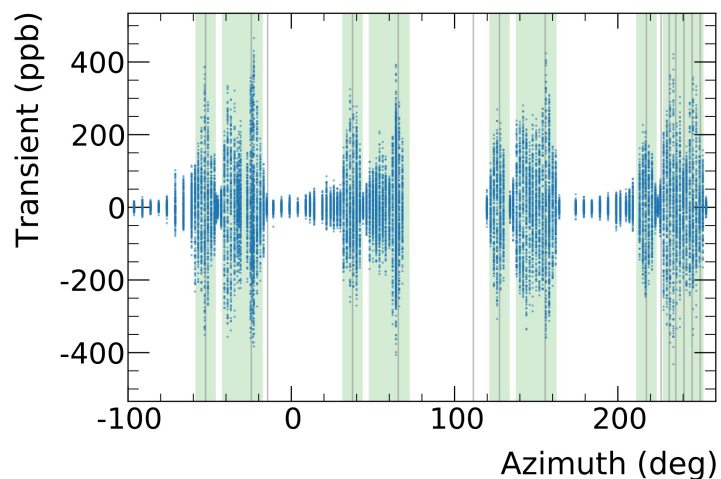
Run-2/3 - Improvements



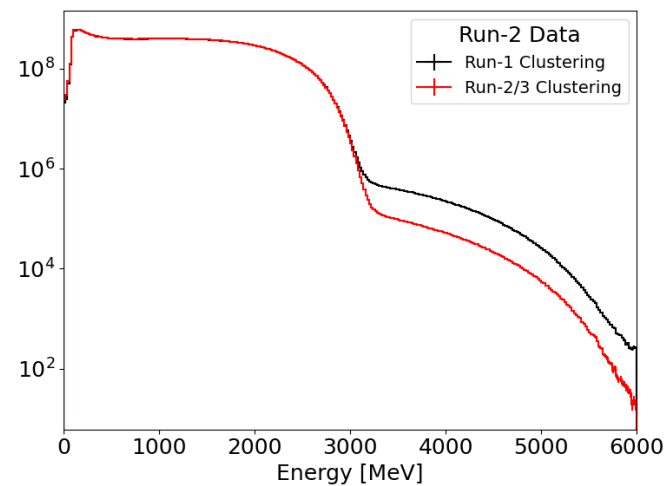
Statistics



Running Conditions



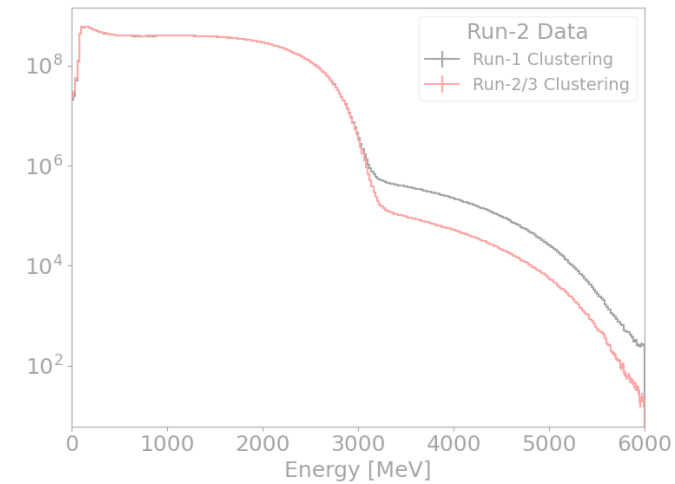
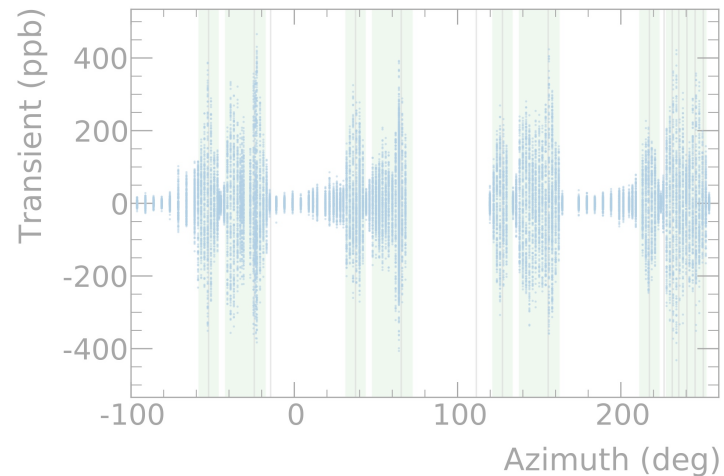
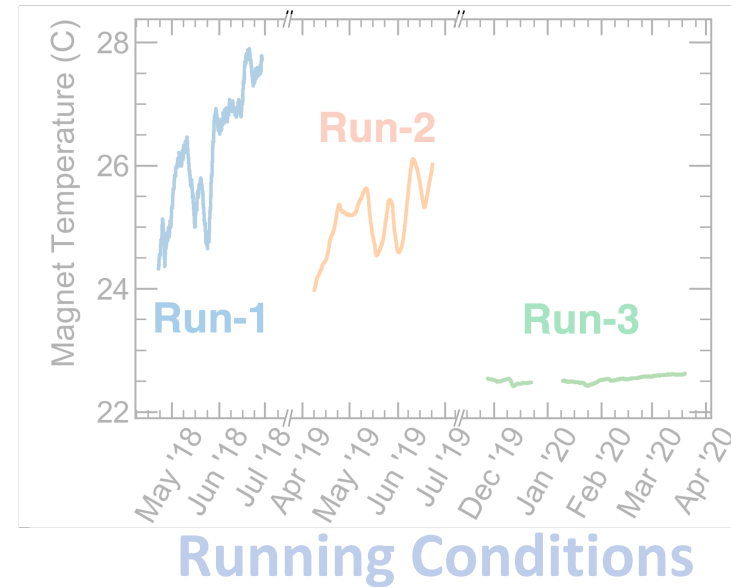
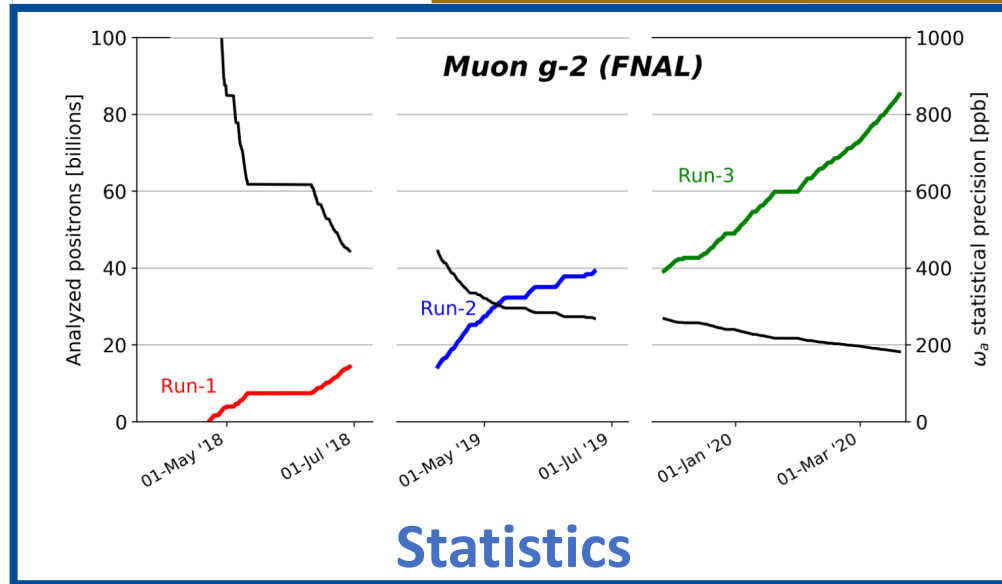
Systematic Measurements & Studies



Analysis Improvements



Run-2/3 - Improvements



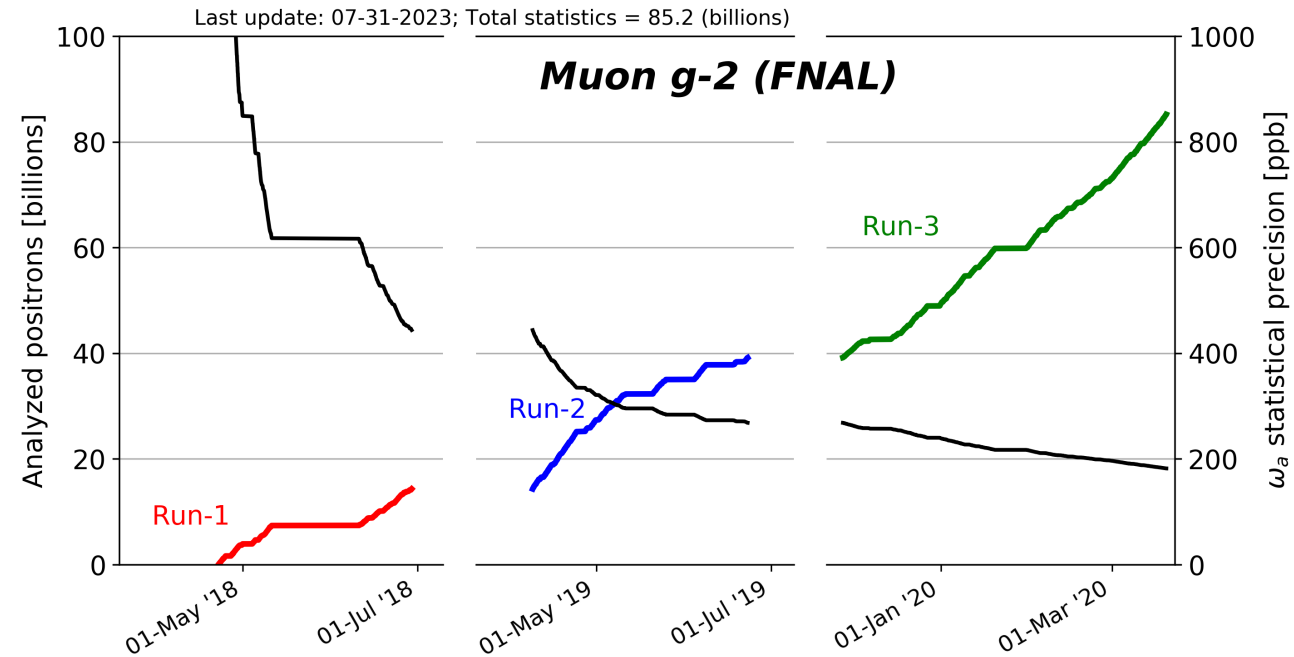


Run-2/3 – Improvements: Statistics



Weighted e^+ in our final fit after quality control

$E > 1 \text{ GeV}$
 $t > 30 \text{ us}$

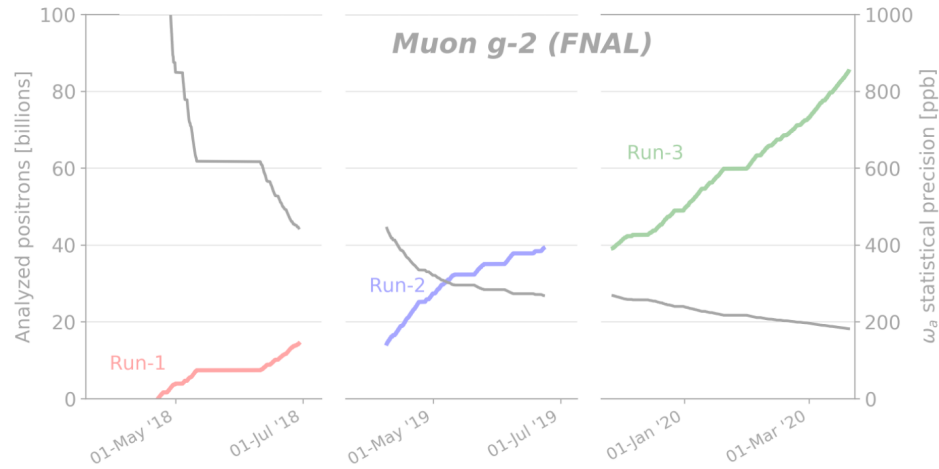


- Factor 4.7 more data in Run-2/3 than Run-1

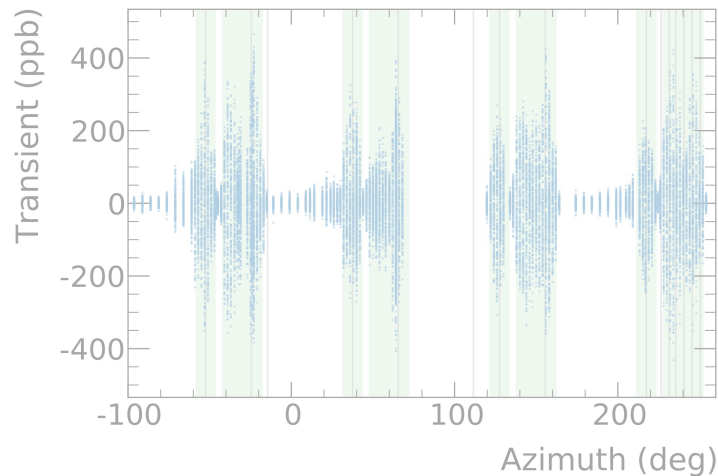
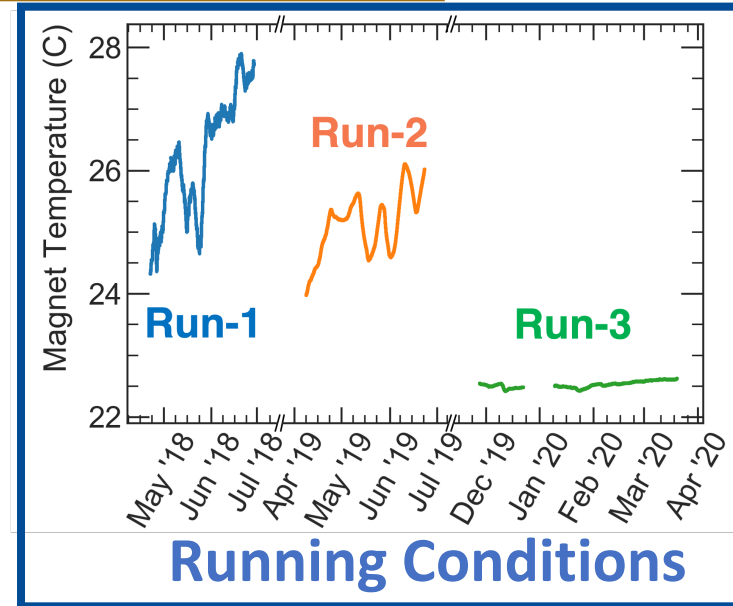
Dataset	Statistical Error [ppb]
Run-1	434
Run-2/3	201
Run-1 + Run-2/3	185



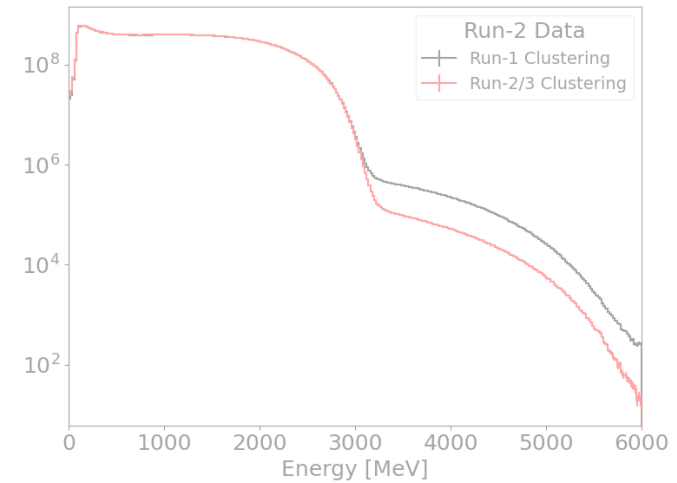
Run-2/3 - Improvements



Statistics



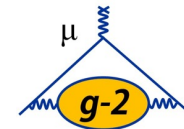
Systematic Measurements & Studies



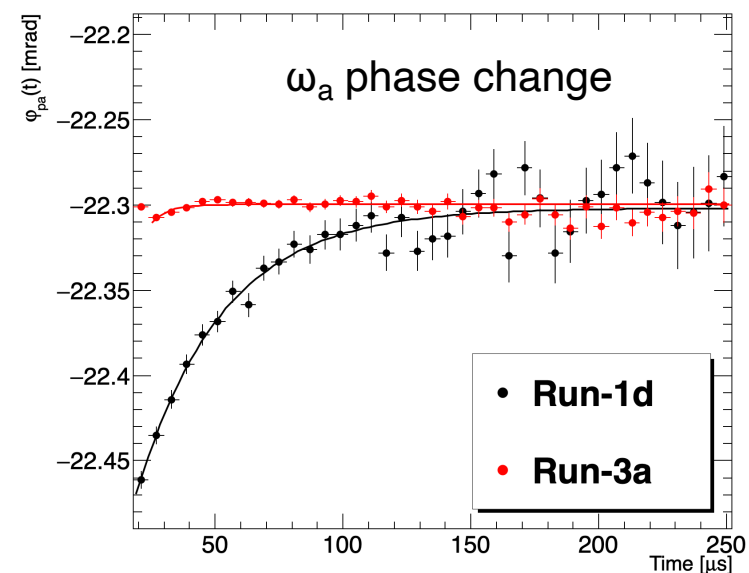
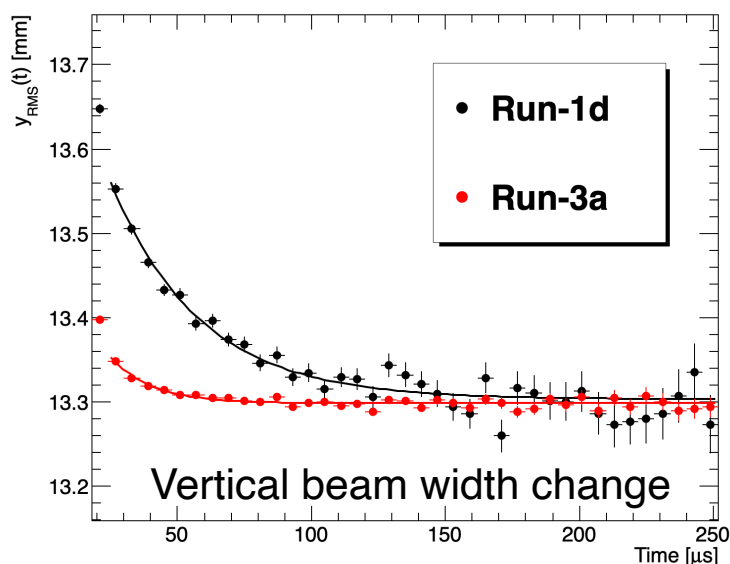
Analysis Improvements



Run-2/3 – Improvements: Running conditions



- Run-1 had **damaged resistors** in 2/32 quad plates leading to **unstable beam storage**
- Resistors re-designed & replaced before Run-2



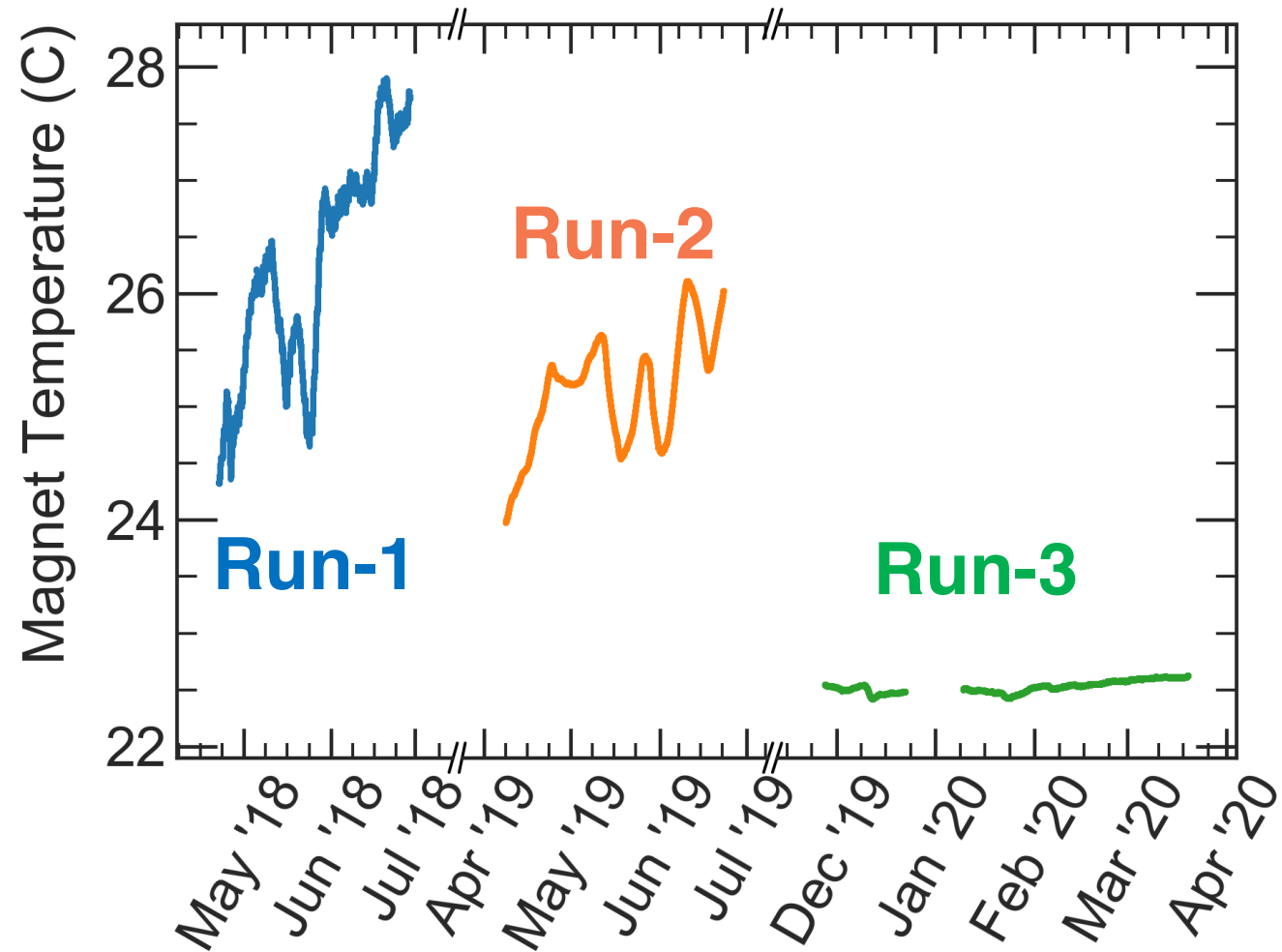
- C_{pa} uncertainty is reduced (75 ppb \rightarrow 13 ppb)
- Beam oscillation frequencies are also more stable



Run-2/3 – Improvements: Running conditions



- Temperature stability makes magnetic field less variable

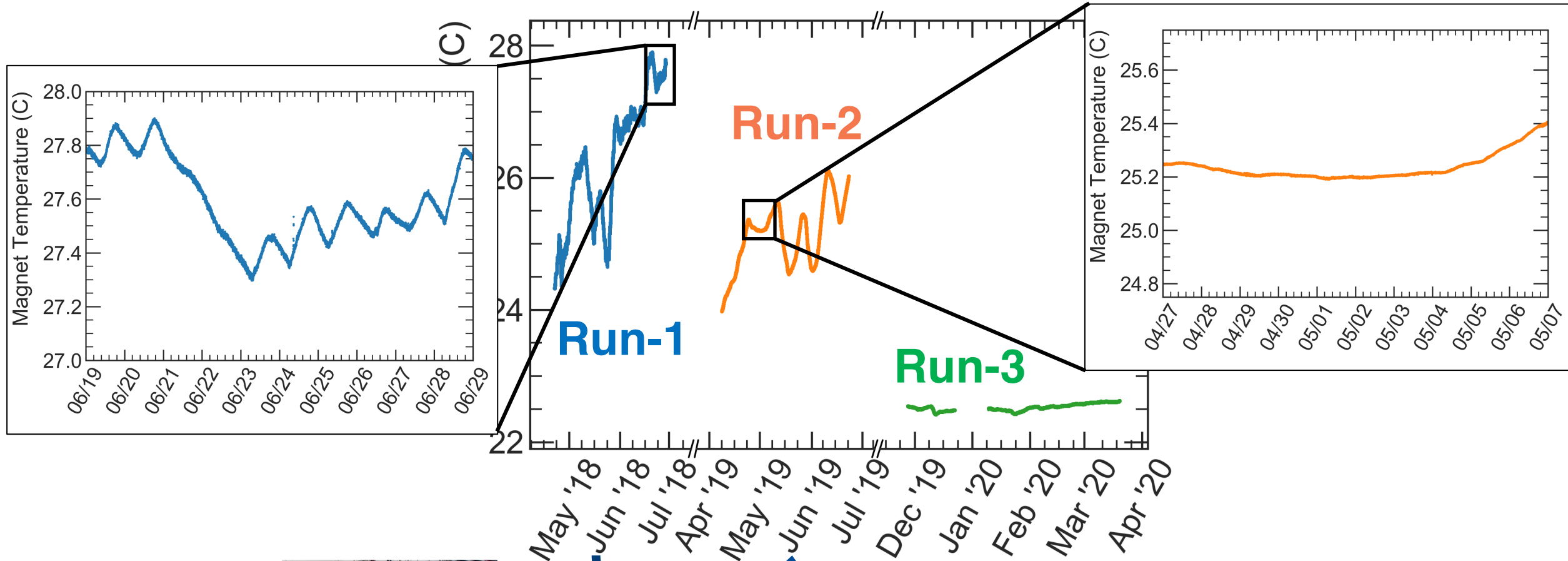




Run-2/3 – Improvements: Running conditions

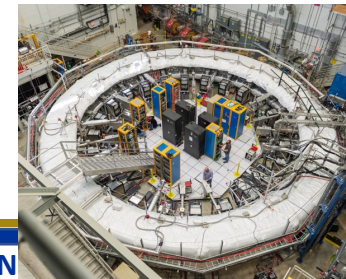
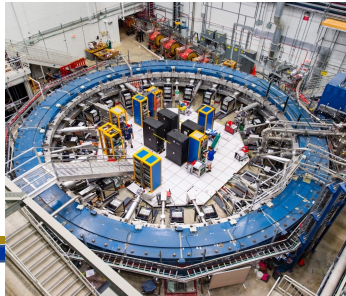


- Temperature stability makes magnetic field less variable



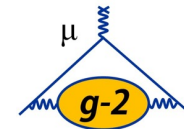
May '18 Jun '18 Jul '18 Apr '19 May '19 Jun '19 Jul '19 Dec '19 Jan '20 Feb '20 Mar '20 Apr '20

Added Insulation

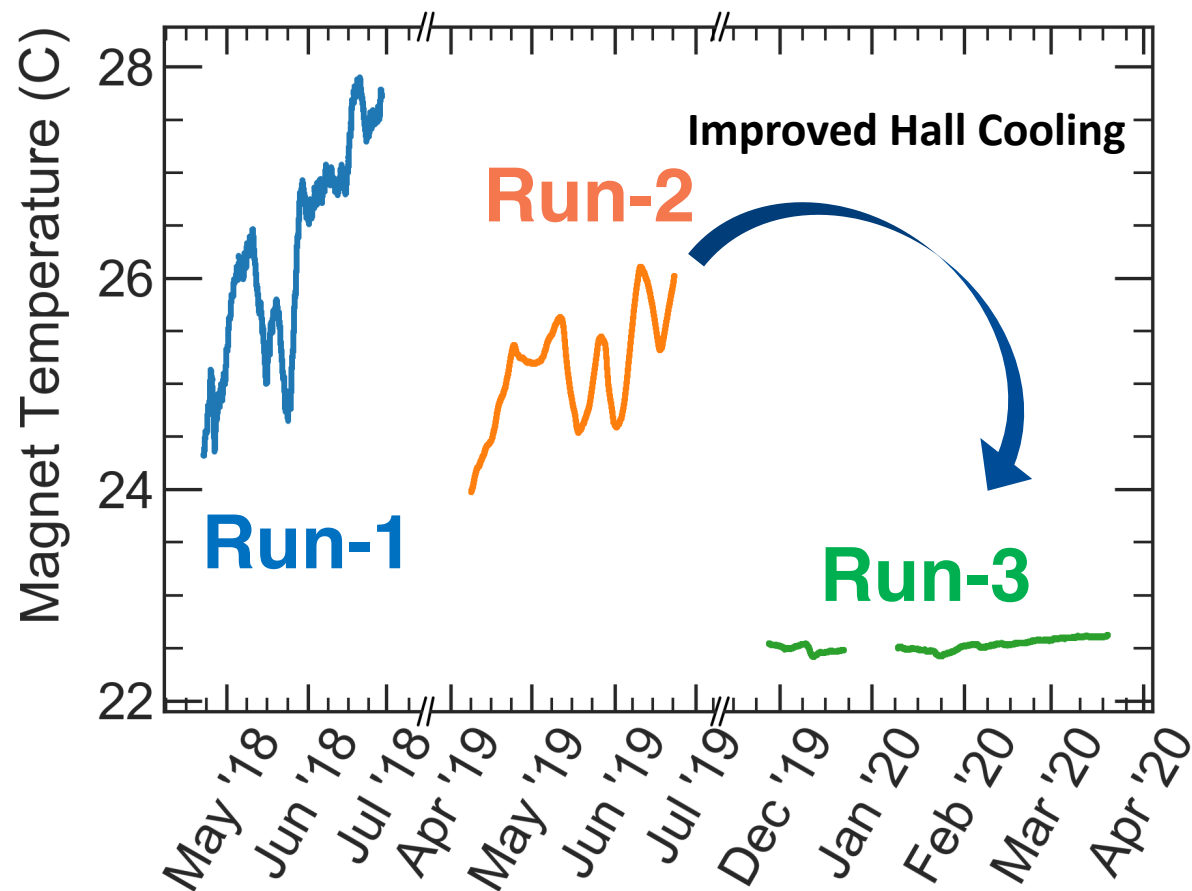




Run-2/3 – Improvements: Running conditions

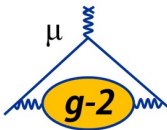


- Temperature stability makes magnetic field less variable

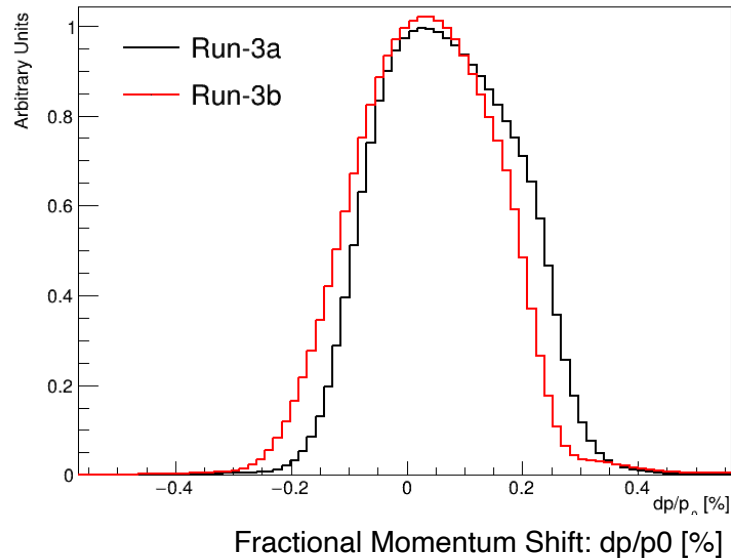




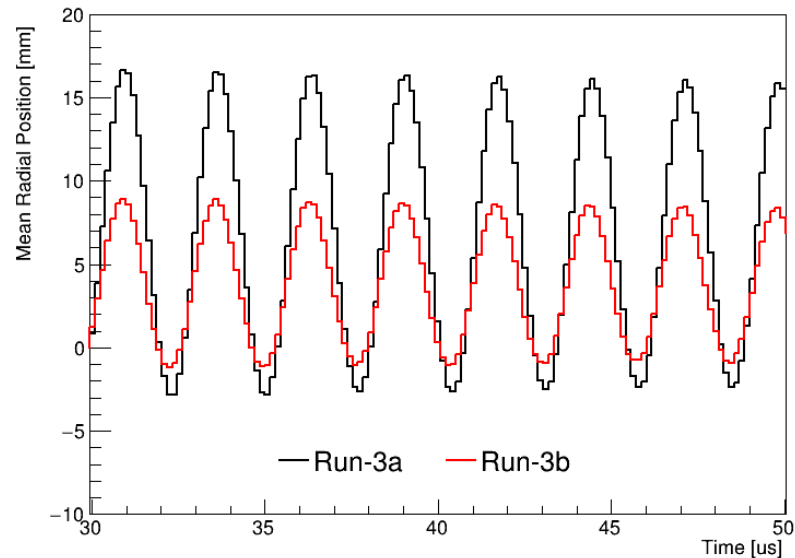
Run-2/3 – Improvements: Running conditions



- Last 18% of Run-2/3 has **upgraded, stronger kicker**



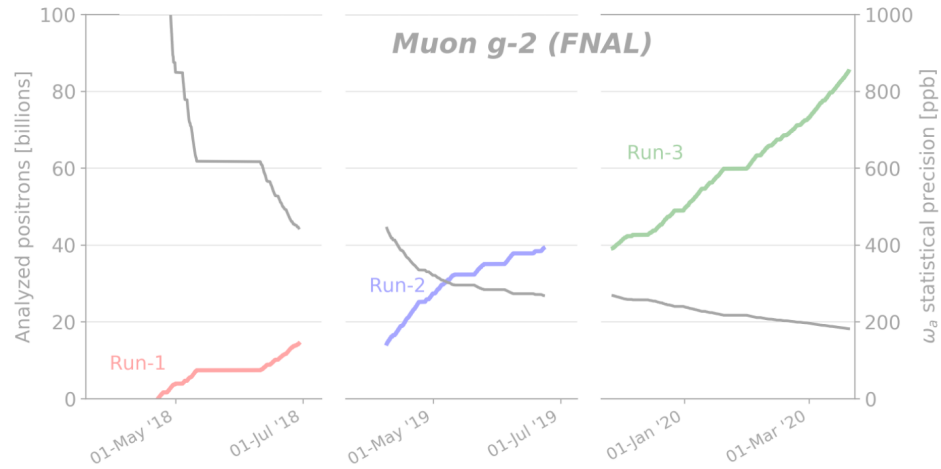
- Momentum distribution more centered
- **Lower E-field correction C_e**



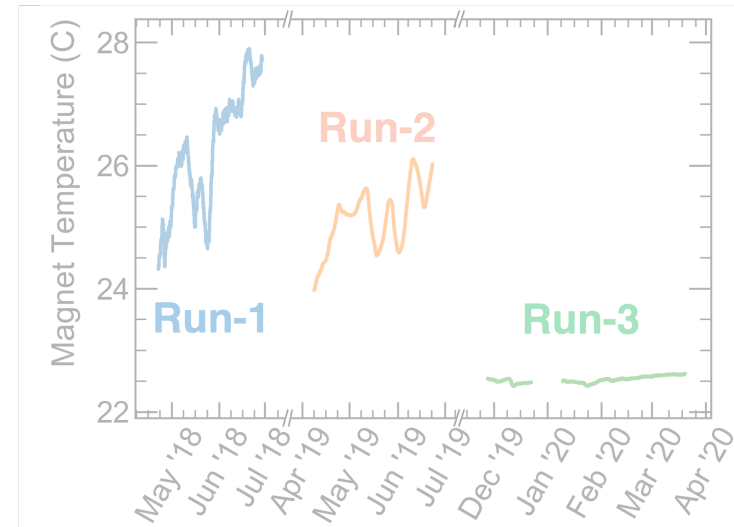
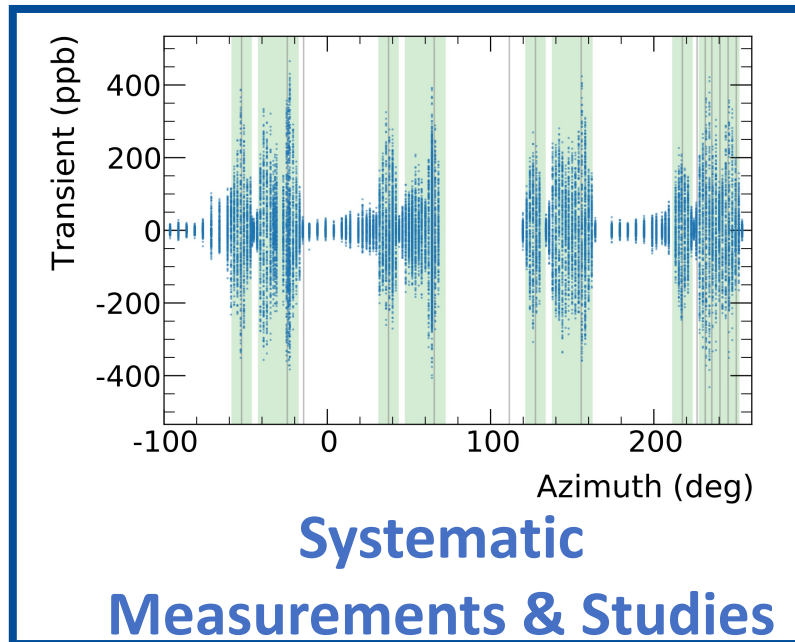
- Phase space matching improved
- **Smaller beam oscillations**



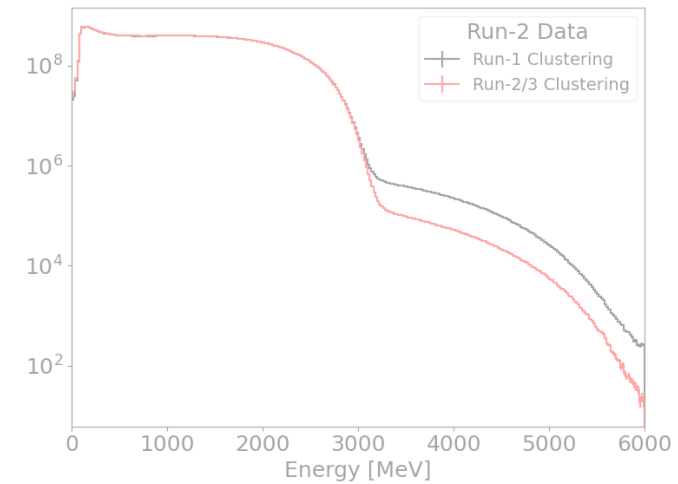
Run-2/3 - Improvements



Statistics



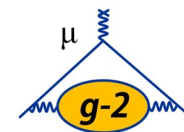
Running Conditions



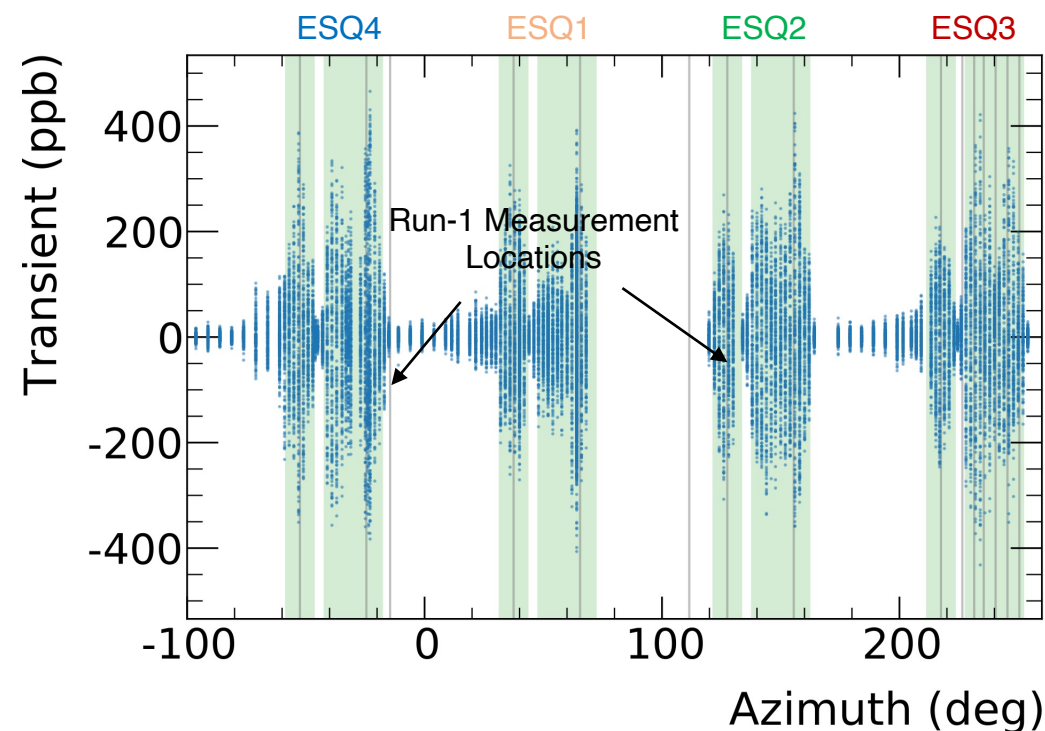
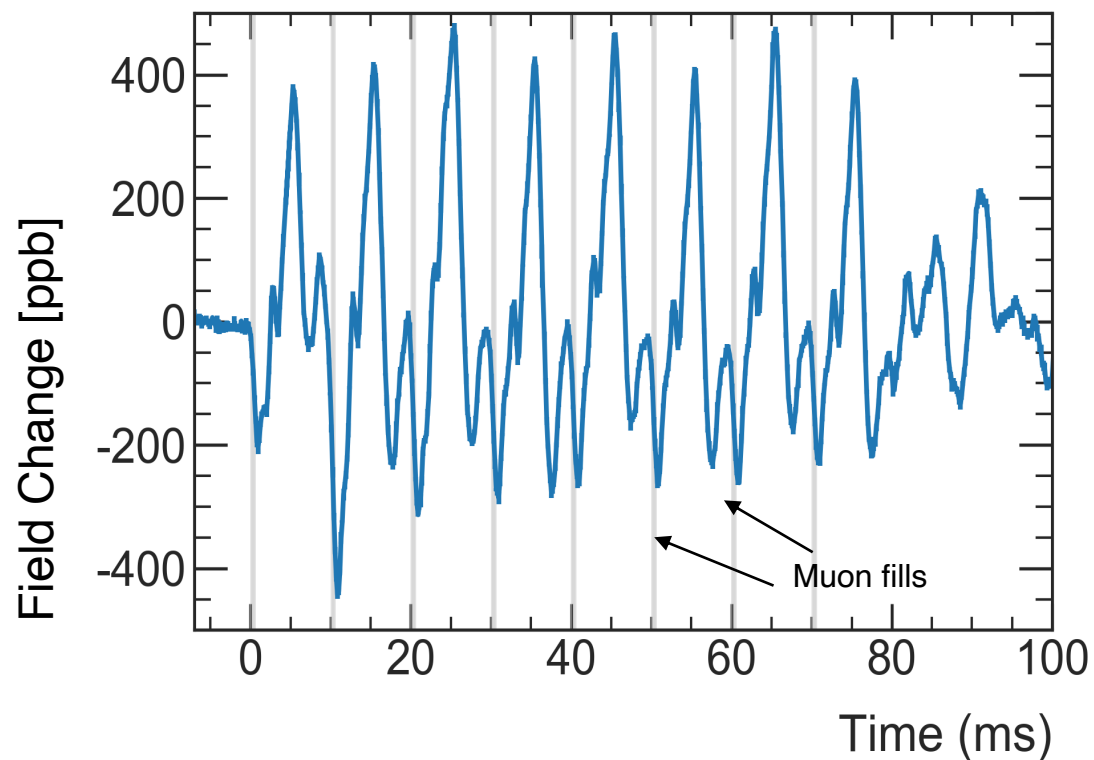
Analysis Improvements



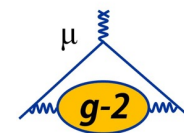
Run-2/3 – Improvements: Improved measurement



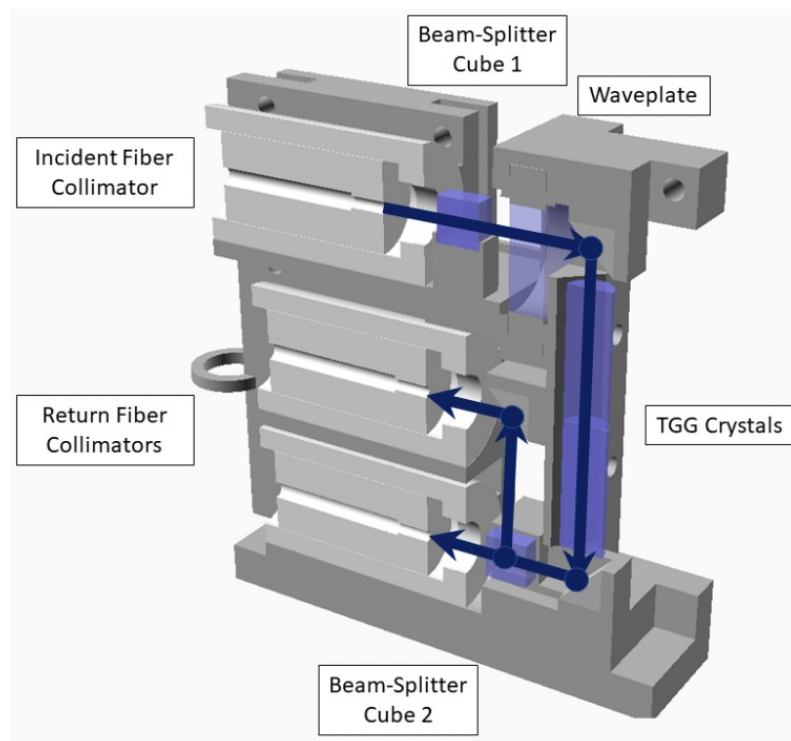
- Pulsing quads vibrate → oscillating magnetic field
- Measured with a new NMR probe



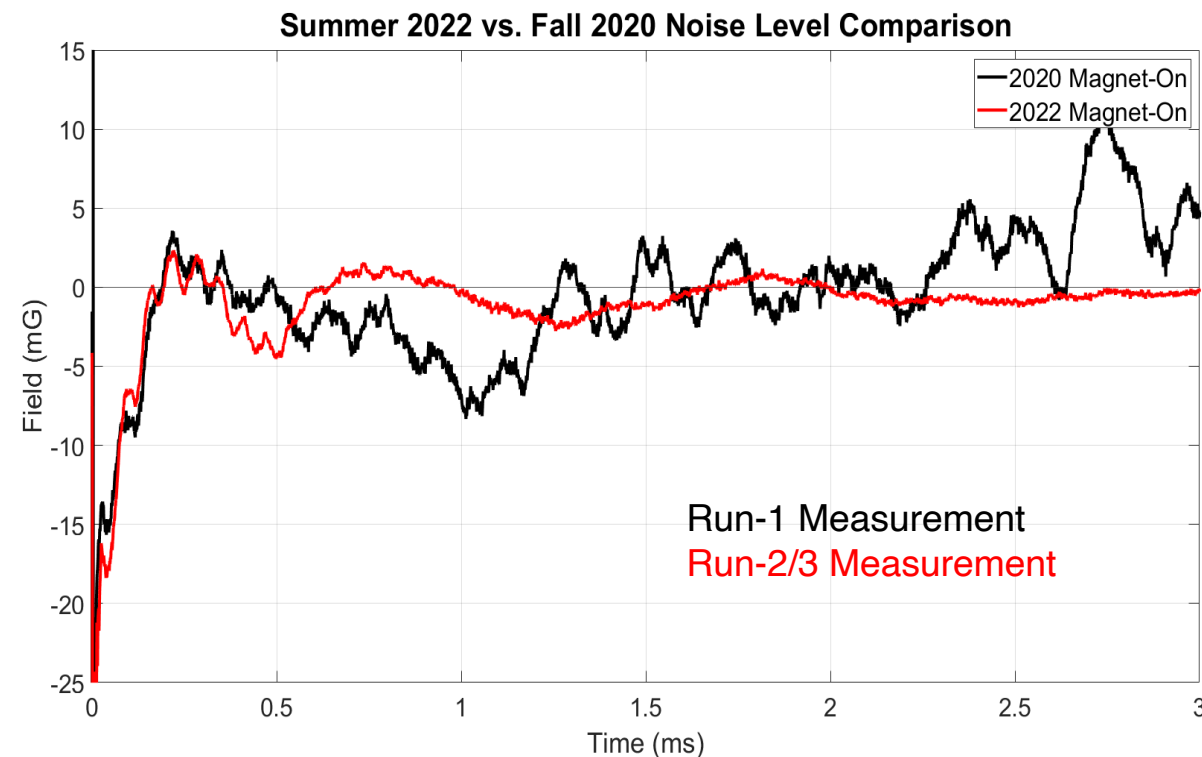
For Run-1 analysis, we had **limited measurement positions**
Largest Run-1 systematic: **92 ppb** reduced to **20 ppb**



- Kicker creates eddy currents → transient magnetic field



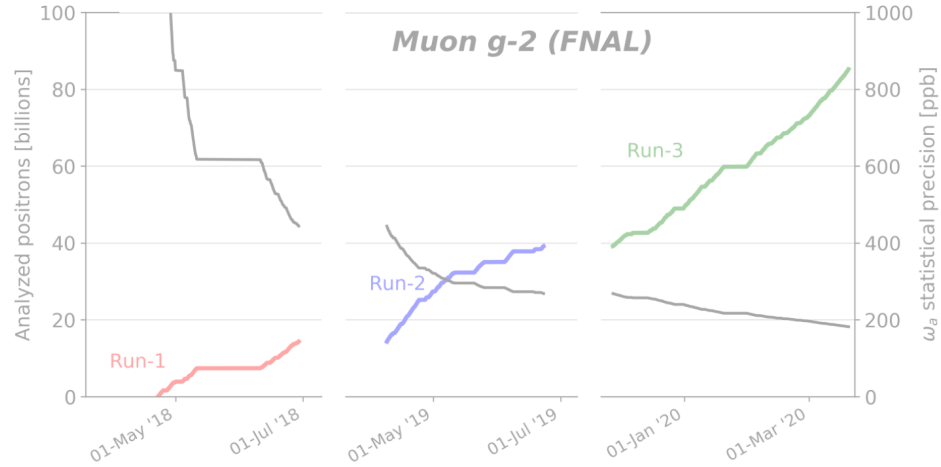
Faraday magnetometer



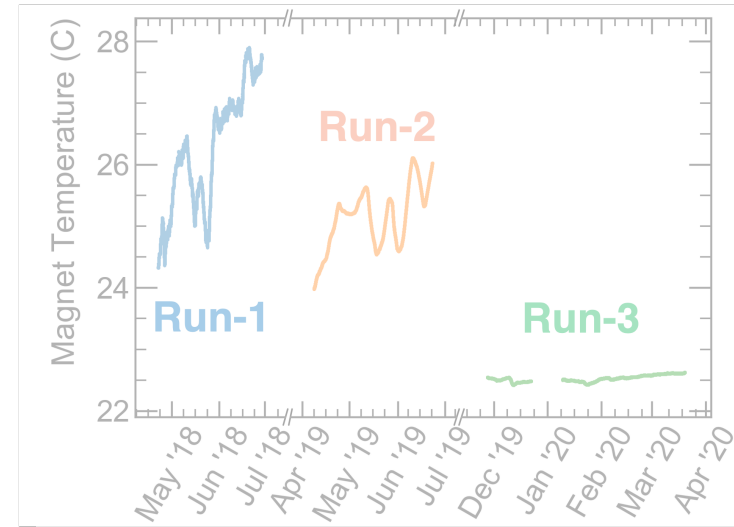
- Run-2/3 has **lower vibration noise vs. Run-1**
- Uncertainty reduced from **37 ppb to 13 ppb**



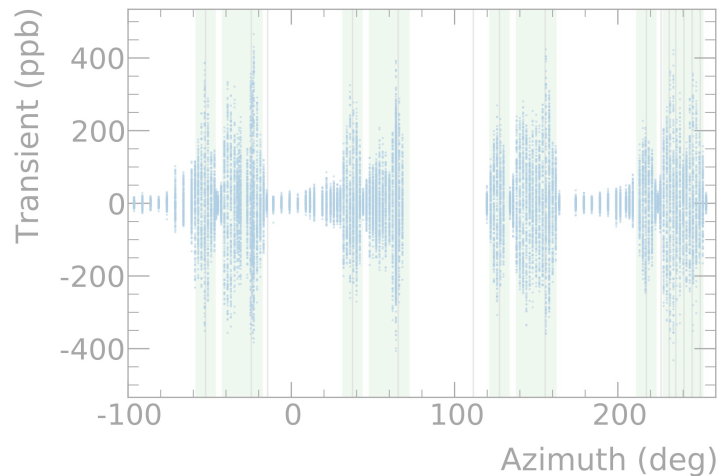
Run-2/3 - Improvements



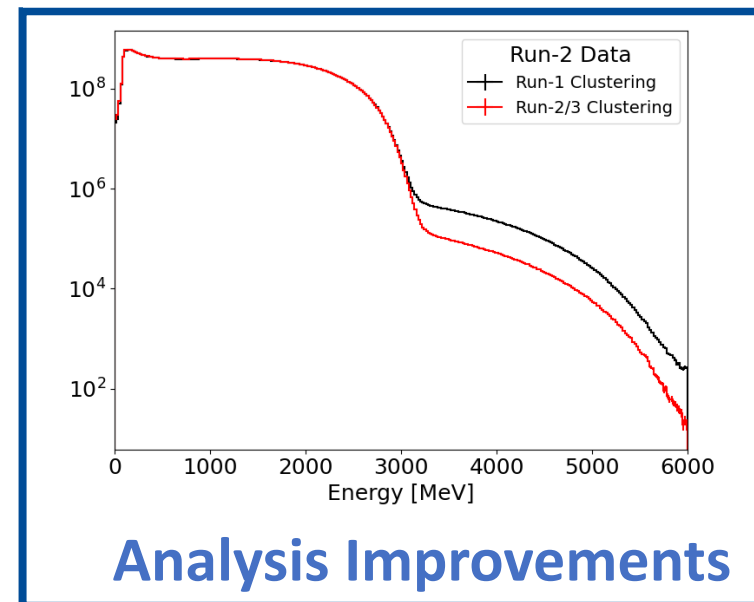
Statistics



Running Conditions



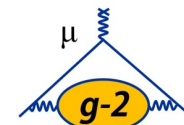
Systematic Measurements & Studies



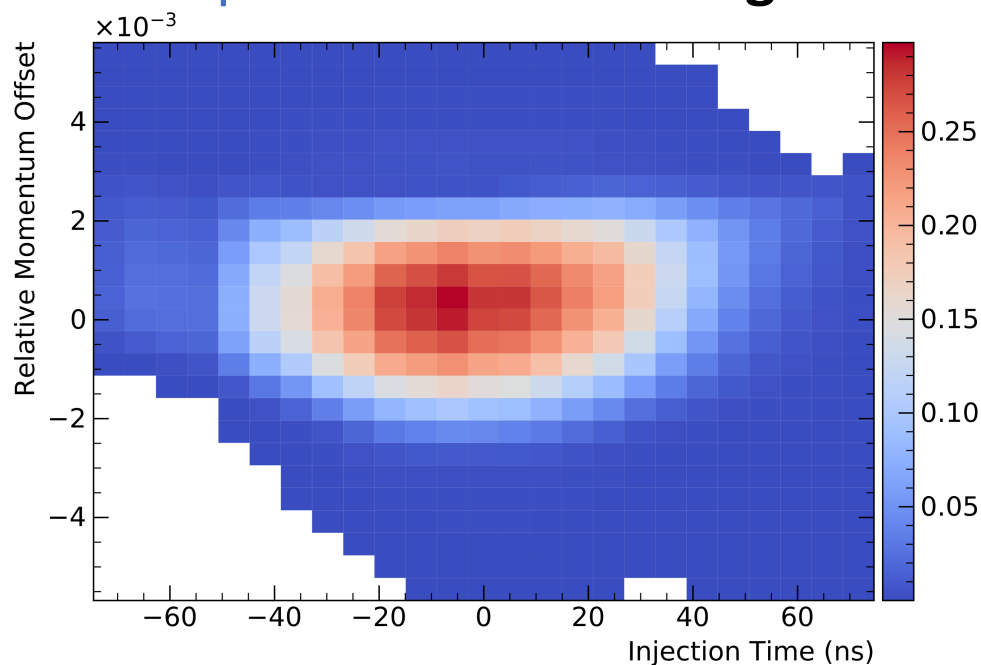
Analysis Improvements



Run-2/3 – Improvements: Analysis improvements

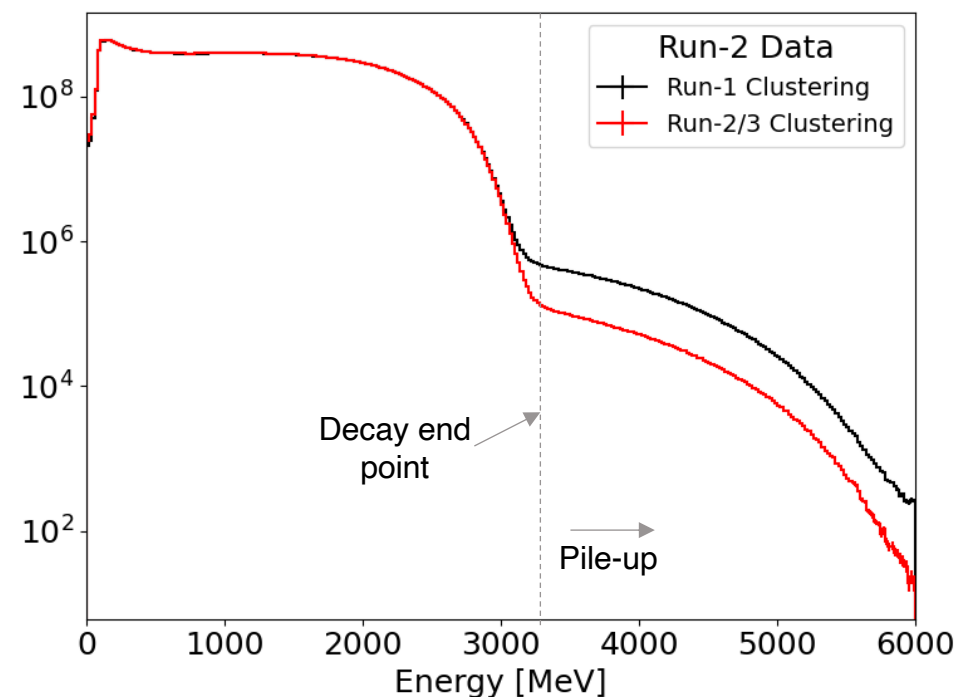


- 2 e^+ arriving at same time can be mistaken for 1: can bias ω_a
- Reduced uncertainty by:
 - Improved reconstruction
 - Improved correction algorithms



Run-1: 53 ppb to Run-2/3: 32 ppb

Run-1: 35 ppb to Run-2/3: 7 ppb



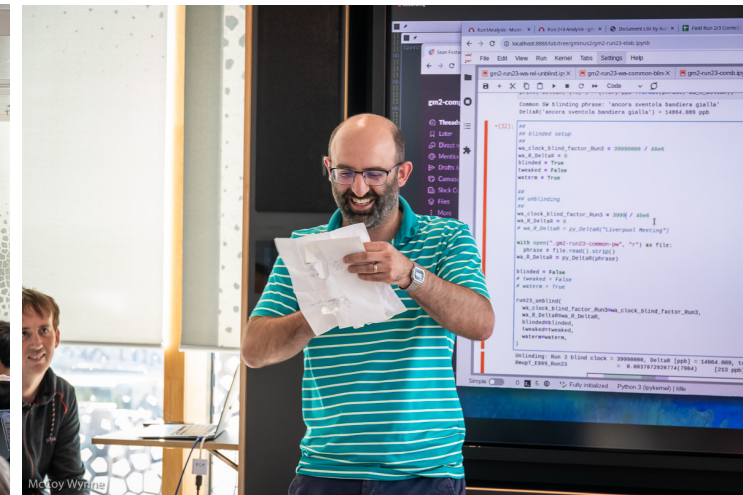
- E-field correction depends on muon momentum distribution
- Now include correlations between momentum & time of injection



Blind analysis and unblinding



- The analysis is performed with a double blinding one software and one hardware.
- Hardware blinding comes from altering our clock frequency randomly **$\pm 25\text{ppm}$ shift**
- On **24th July** the secret numbers have been unveiled, revealing the hidden frequency

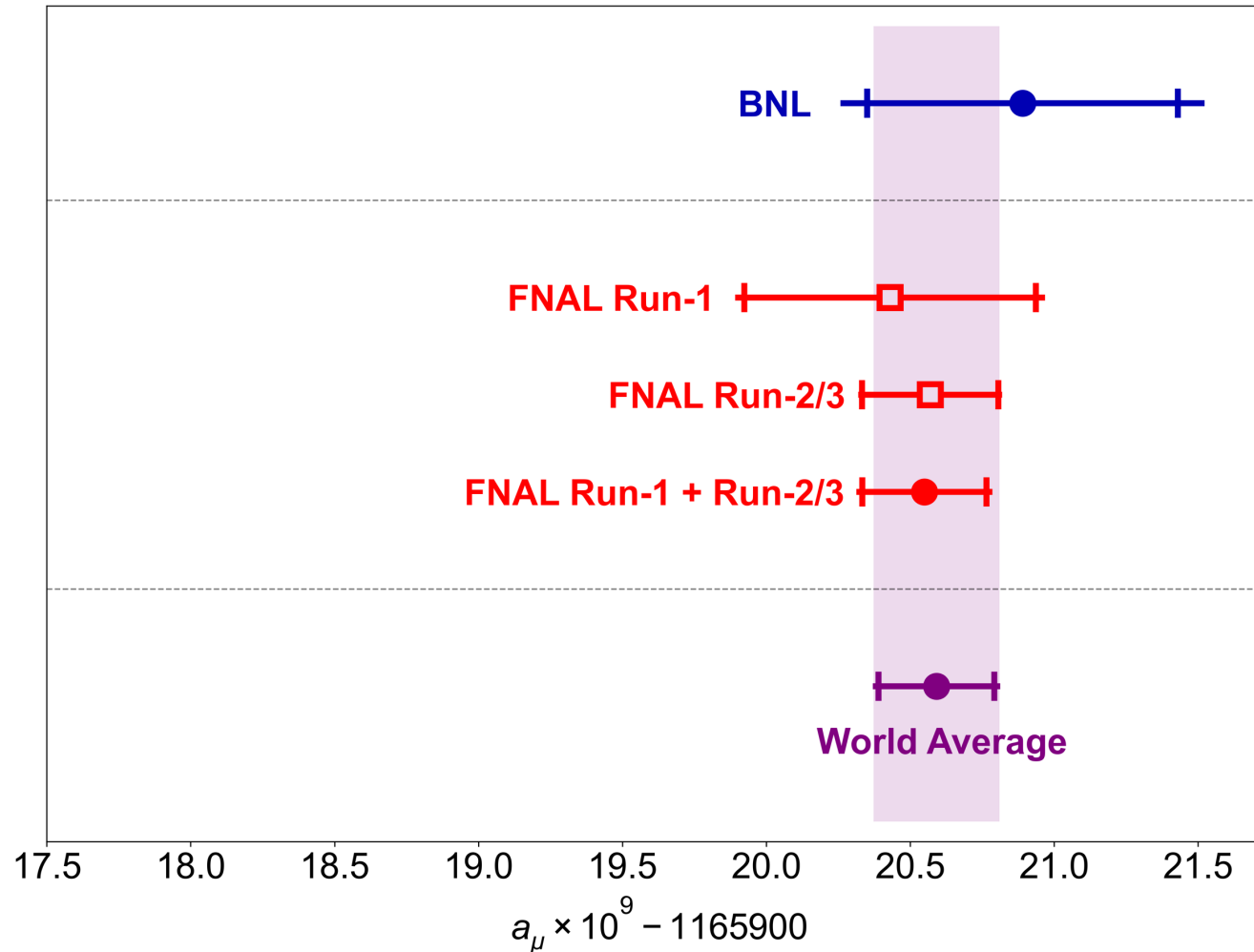




Run-2/3 Result



$$a_\mu(\text{FNAL}) = 0.00\ 116\ 592\ 055(24) [203\ \text{ppb}]$$



$$a_\mu(\text{Exp}) = 0.00\ 116\ 592\ 059(22) [190\ \text{ppb}]$$

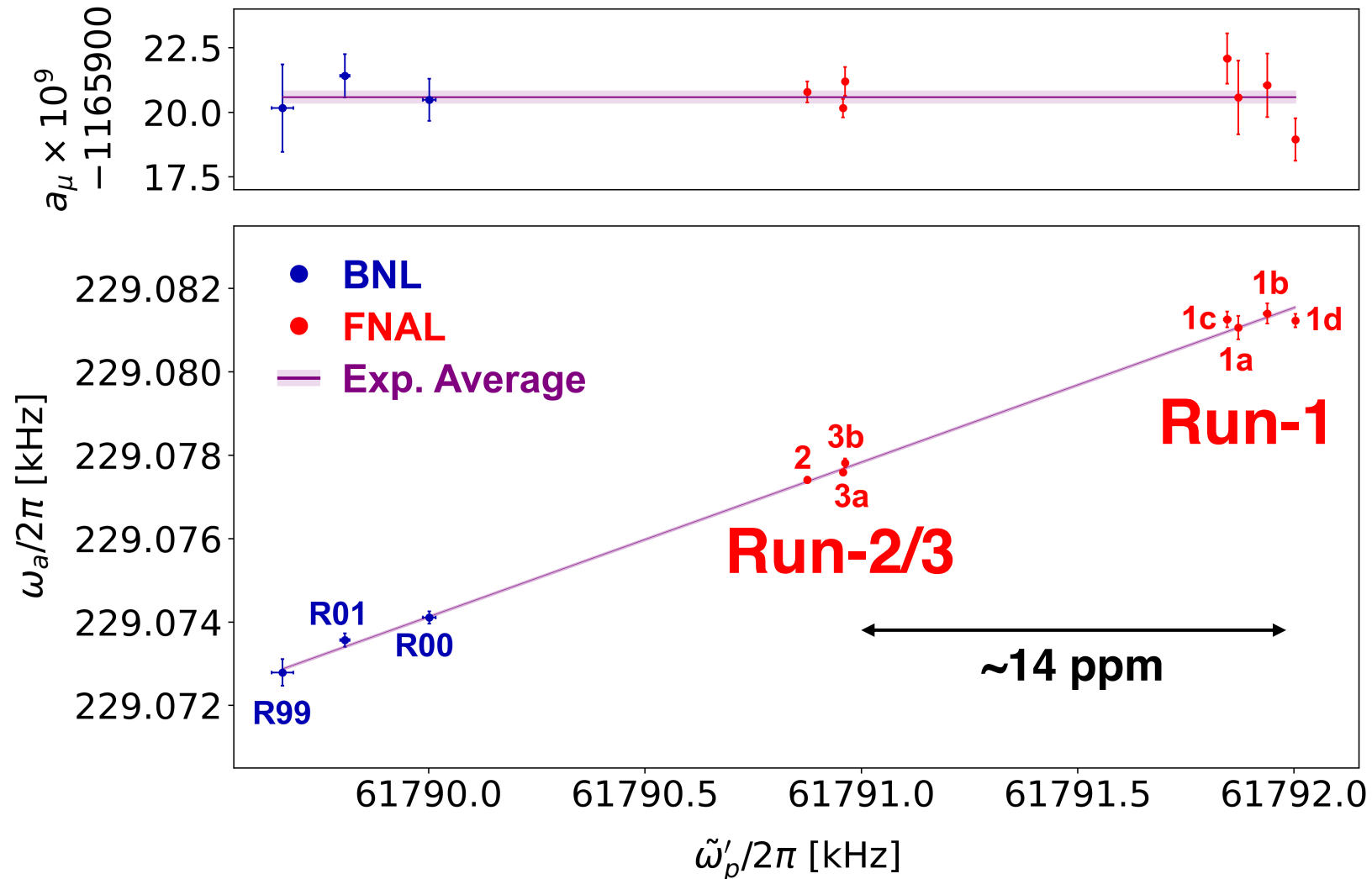
- Both FNAL and BNL dominated by statistical error
- Combined world average **dominated by FNAL** values.



Run-2/3 Result

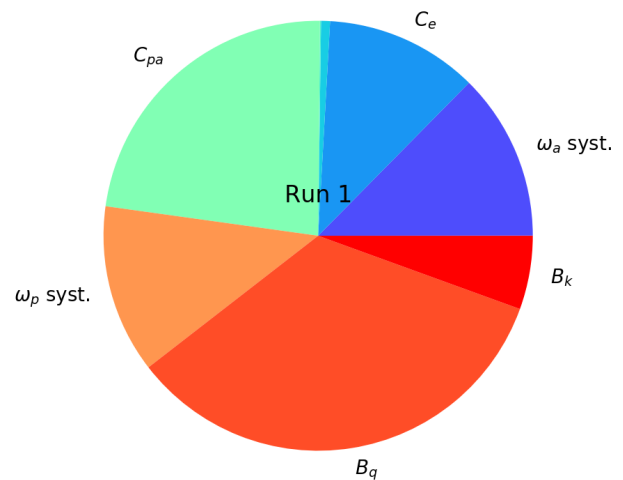
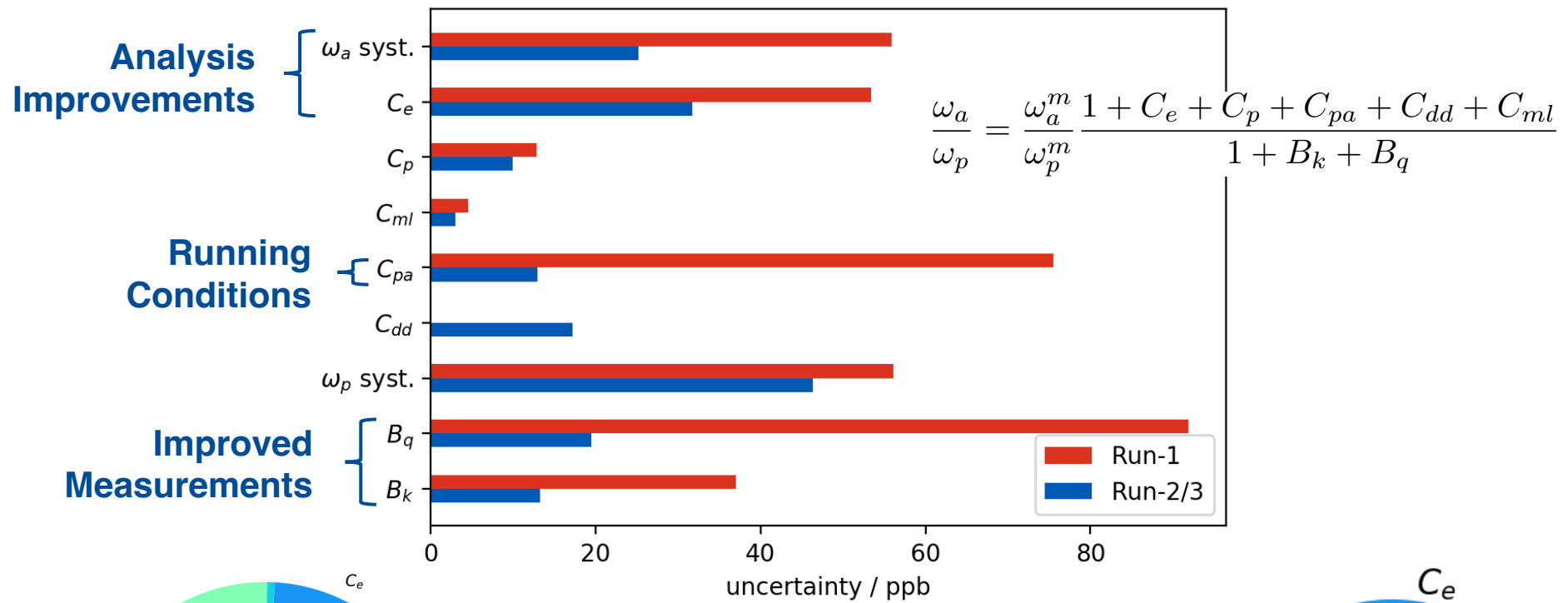


- Datasets were taken at slightly different field settings
- Allows a cross check with one of the most basic “handles”





Run-2/3 Uncertainties summary



157 ppb to 70 ppb





Run-2/3 Publication

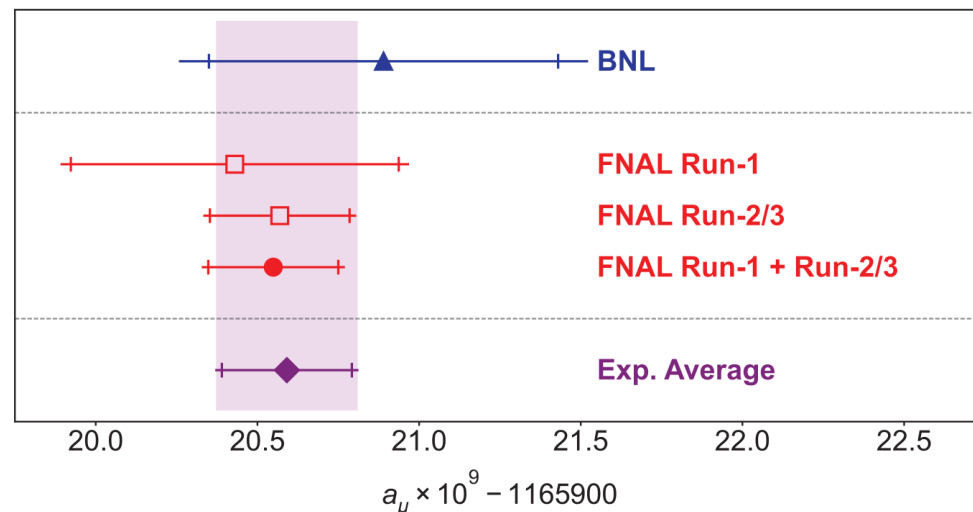


Published on PRL – 17th October 2023!

Measurement of the Positive Muon Anomalous Magnetic Moment to 0.20 ppm

D. P. Aguillard³³, T. Albahri³⁰, D. Allspach⁷, A. Anisenkov^{4,a}, K. Badgley⁷, S. Baeßler^{35,b}, I. Bailey^{17,c},
 L. Bailey²⁷, V. A. Baranov^{15,d}, E. Barlas-Yucel²⁸, T. Barrett⁶, E. Barzi⁷, F. Bedeschi¹⁰, M. Berz¹⁸,
 M. Bhattacharya⁷, H. P. Binney³⁶, P. Bloom¹⁹, J. Bono⁷, E. Bottalico³⁰, T. Bowcock³⁰, S. Braun³⁶,
 M. Bressler³², G. Cantatore^{12,e}, R. M. Carey², B. C. K. Casey⁷, D. Cauz^{26,f}, R. Chakraborty²⁹, A. Chapelain⁶,
 S. Chappa⁷, S. Charity³⁰, C. Chen^{23,22}, M. Cheng²⁸, R. Chislett²⁷, Z. Chu^{22,g}, T. E. Chupp³³, C. Claessens³⁶,
 M. E. Convery⁷, S. Corrodi¹, L. Cotrozzi^{10,h}, J. D. Crnkovic⁷, S. Dabagov^{8,i}, P. T. Debevec²⁸, S. Di Falco¹⁰,
 G. Di Sciascio¹¹, B. Drendel⁷, A. Driutti^{10,h}, V. N. Duginov^{15,d}, M. Eads²⁰, A. Edmonds², J. Esquivel⁷,
 M. Farooq³³, R. Fatemi²⁹, C. Ferrari^{10,j}, M. Fertl¹⁴, A. T. Fienberg³⁶, A. Fioretti^{10,j}, D. Flay³², S. B. Foster²,
 H. Friedrichs⁷, N. S. Froemming²⁰, C. Gabbanini^{10,j}, I. Gaines⁷, M. D. Galati^{10,h}, S. Ganguly⁷, A. Garcia³⁶,
 J. George^{32,k}, L. K. Gibbons⁶, A. Gioiosa^{25,l}, K. L. Giovanetti¹³, P. Girotti¹⁰, W. Gohn²⁹, L. Goodenough⁷,
 T. Gorringer²⁹, J. Grange³³, S. Grant^{1,27}, F. Gray²¹, S. Haciomeroglu^{5,m}, T. Halewood-Leagas³⁰, D. Hampai⁸,
 F. Han²⁹, J. Hempstead³⁶, D. W. Hertzog³⁶, G. Hesketh²⁷, E. Hess¹⁰, A. Hibbert³⁰, Z. Hodge³⁶, K. W. Hong³⁵,
 R. Hong^{29,1}, T. Hu^{23,22}, Y. Hu^{22,g}, M. Iacovacci^{9,n}, M. Incagli¹⁰, P. Kammel³⁶, M. Kargiantoulakis⁷,
 M. Karuza^{12,o}, J. Kaspar³⁶, D. Kawall³², L. Kelton²⁹, A. Keshavarzi³¹, D. S. Kessler³², K. S. Khaw^{23,22},
 Z. Khechadorian⁶, N. V. Khomutov¹⁵, B. Kiburg⁷, M. Kiburg^{7,19}, O. Kim³⁴, N. Kinnaird², E. Kraegeloh³³,
 V. A. Krylov¹⁵, N. A. Kuchinskiy¹⁵, K. R. Labe⁶, J. LaBounty³⁶, M. Lancaster³¹, S. Lee⁵, B. Li^{22,1,p}, D. Li^{22,q},
 L. Li^{22,g}, I. Logashenko^{4,a}, A. Lorente Campos²⁹, Z. Lu^{22,g}, A. Lucà⁷, G. Lukicov²⁷, A. Lusiani^{10,r},
 A. L. Lyon⁷, B. MacCoy³⁶, R. Madrak⁷, K. Makino¹⁸, S. Mastroianni⁹, J. P. Miller², S. Miozzi¹¹, B. Mitra³⁴,
 J. P. Morgan⁷, W. M. Morse³, J. Mott^{7,2}, A. Nath^{9,n}, J. K. Ng^{23,22}, H. Nguyen⁷, Y. Oksuzian¹, Z. Omarov^{16,5},
 R. Osofsky³⁶, S. Park⁵, G. Pauletta^{26,s}, G. M. Piacentino^{25,t}, R. N. Pilato³⁰, K. T. Pitts^{28,u}, B. Plaster²⁹,
 D. Počanić³⁵, N. Pohlman²⁰, C. C. Polly⁷, J. Price³⁰, B. Quinn³⁴, M. U. H. Qureshi¹⁴, S. Ramachandran^{1,k},
 E. Ramberg⁷, R. Reimann¹⁴, B. L. Roberts², D. L. Rubin⁶, L. Santi^{26,f}, C. Schlesier^{28,v}, A. Schreckenberger⁷,
 Y. K. Semertzidis^{5,16}, D. Shemyakin^{4,a}, M. Sorbara^{11,w}, D. Stöckinger²⁴, J. Stapleton⁷, D. Still⁷, C. Stoughton⁷,
 D. Stratakis⁷, H. E. Swanson³⁶, G. Sweetmore³¹, D. A. Sweigart⁶, M. J. Syphers²⁰, D. A. Tarazona^{6,30,18},
 T. Teubner³⁰, A. E. Tewsley-Booth^{29,33}, V. Tishchenko³, N. H. Tran^{2,x}, W. Turner³⁰, E. Valetov¹⁸,
 D. Vasilkova^{27,30}, G. Venanzoni^{30,1}, V. P. Volnykh¹⁵, T. Walton⁷, A. Weisskopf¹⁸, L. Welyt-Rieger⁷, P. Winter¹,
 Y. Wu¹, B. Yu³⁴, M. Yucel⁷, Y. Zeng^{23,22} and C. Zhang³⁰

Quantity	Correction (ppb)	Uncertainty (ppb)
ω_a^m (statistical)	...	201
ω_a^m (systematic)	...	25
C_e	451	32
C_p	170	10
C_{pa}	-27	13
C_{dd}	-15	17
C_{ml}	0	3
$f_{\text{calib}} \cdot \langle \omega'_p(\vec{r}) \times M(\vec{r}) \rangle$...	46
B_k	-21	13
B_q	-21	20
$\mu'_p(34.7^\circ)/\mu_e$...	11
m_μ/m_e	...	22
$g_e/2$...	0
Total systematic for \mathcal{R}'_μ	...	70
Total external parameters	...	25
Total for a_μ	622	215



PhysRevLett.131.161802

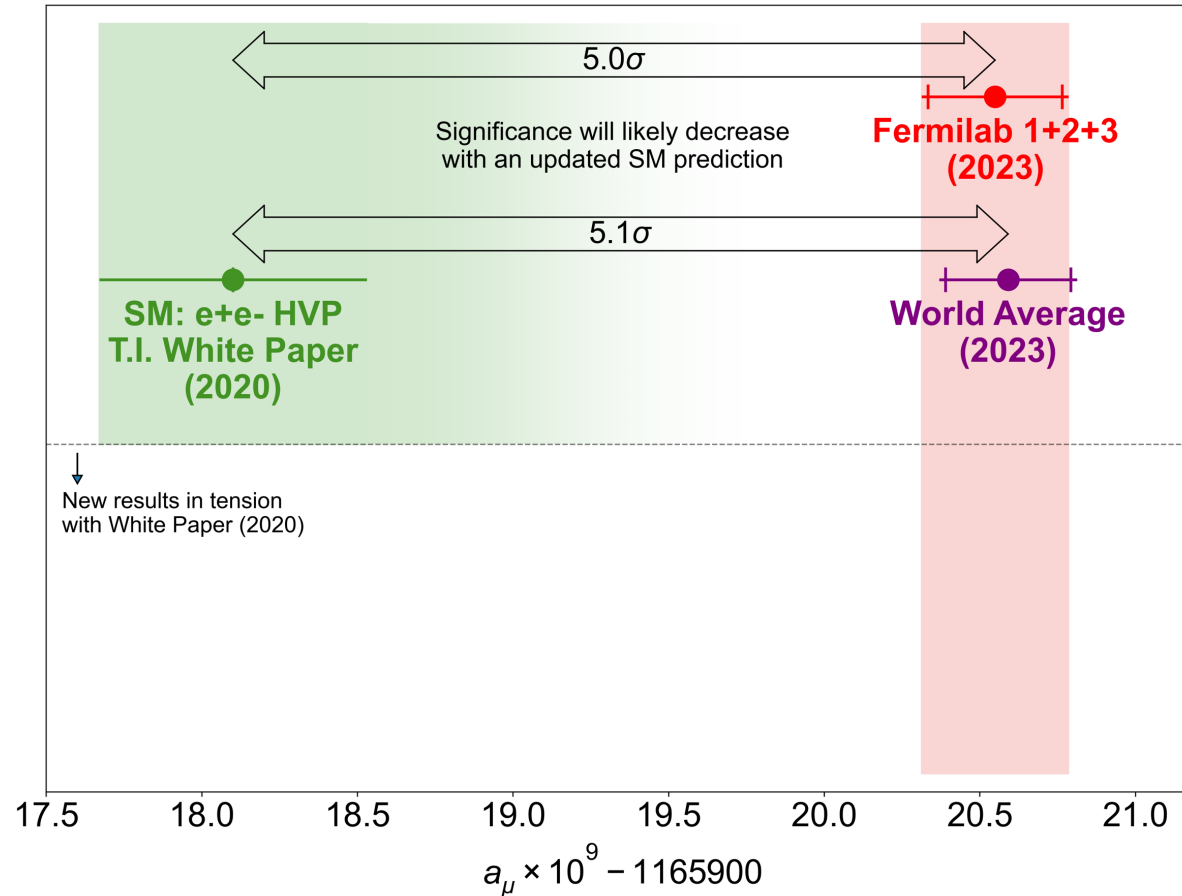
(The Muon $g-2$ Collaboration)



Theory comparison – The puzzle

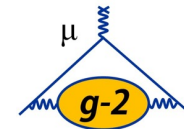


- Theory prediction is less clear now.
- Comparing with WP (2020) a large discrepancy is found 5.0σ

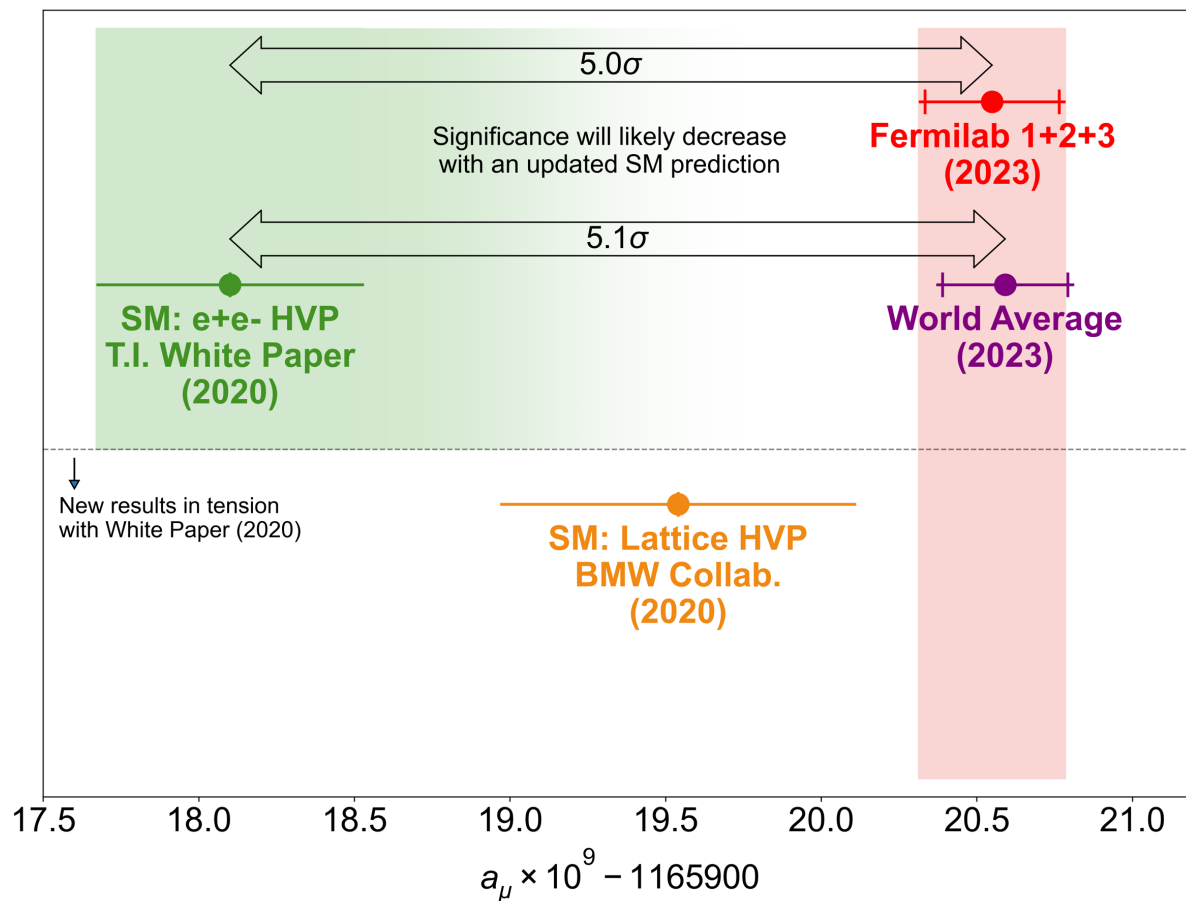




Theory comparison – The puzzle in the puzzle



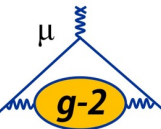
- Theory prediction is less clear now.
- Comparing with WP (2020) a large discrepancy is found 5.0σ



- Considering **BMW** result for the HVP term, instead of WP the discrepancy is reduced.



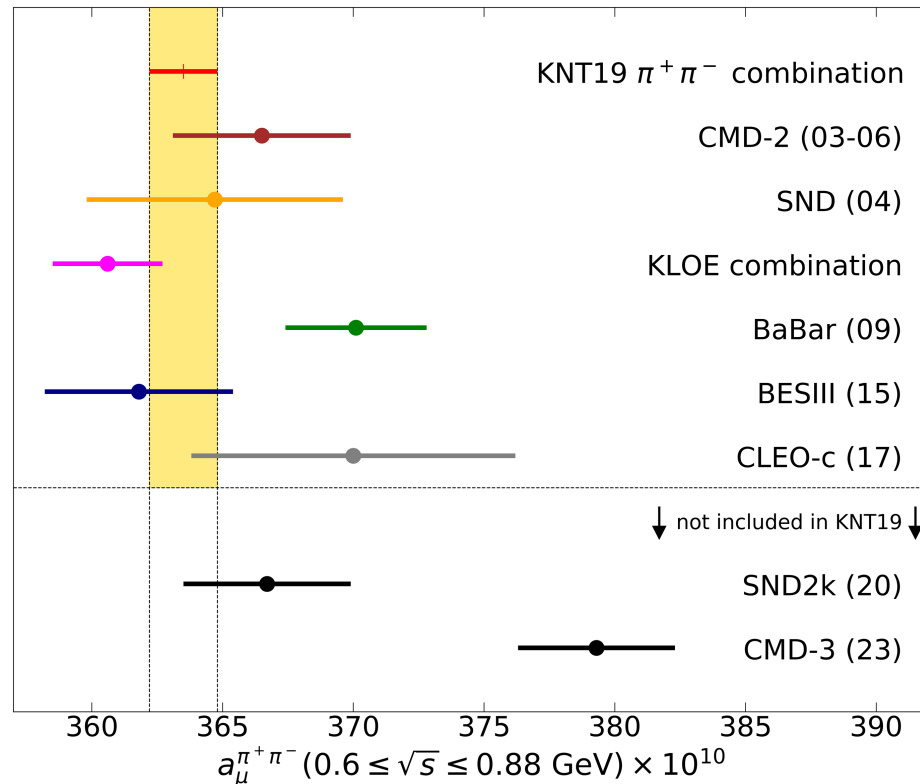
Theory comparison – The puzzle³



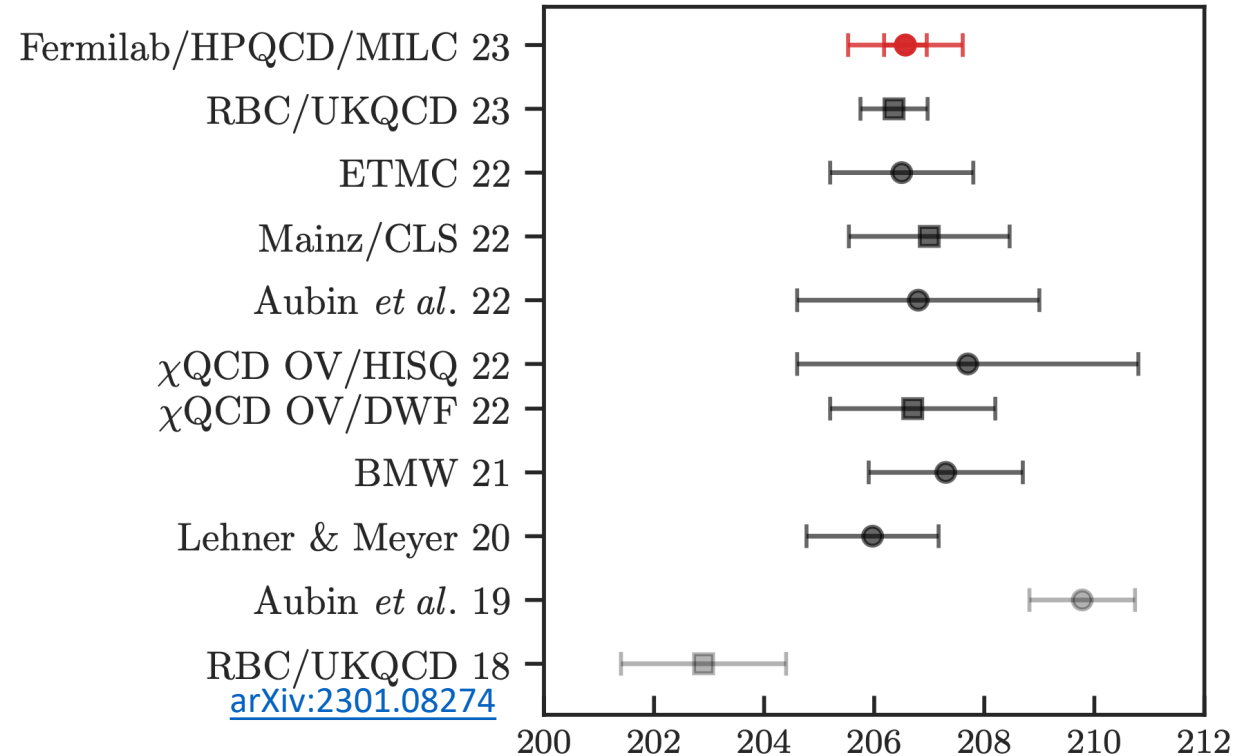
- New results from **SND2k** and **CMD-3** since White Paper, **CMD-3** shows a discrepancy.
- From lattice side, other groups are improving the precision of the calculation, converging on a common value.

Dispersive Approach

Keshavarzi, Nomura, Teubner: Priv. Comm.

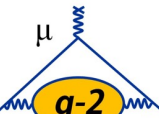


Lattice QCD Approach



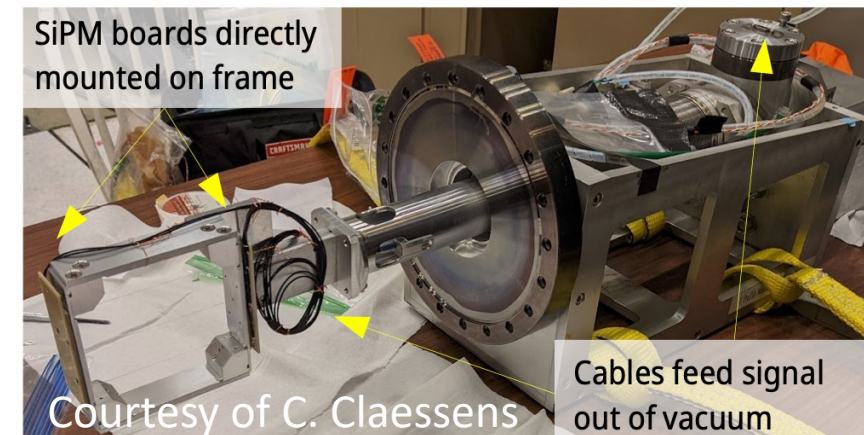
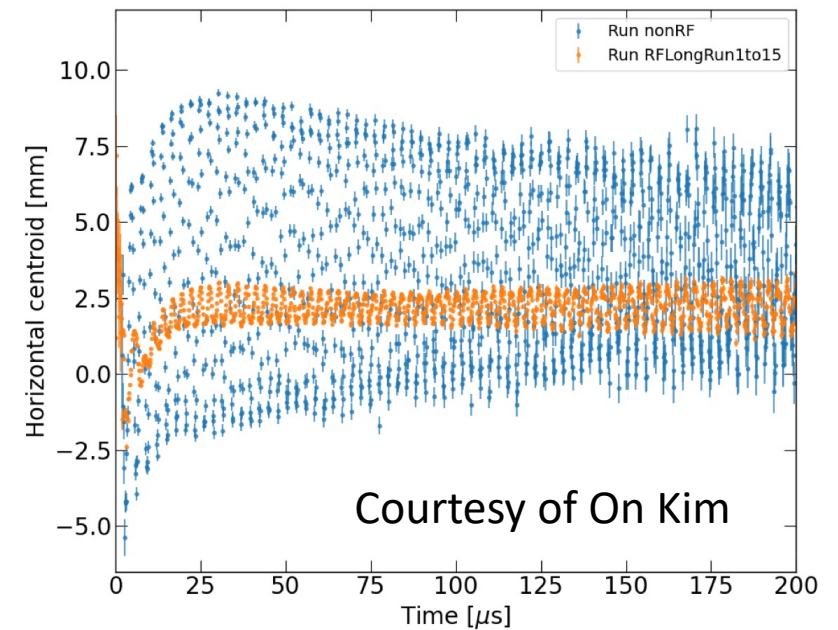


Muon $g-2$ - Run-4/5/6 further improvements



CBO amplitude: factor of 8.6
5.3 \rightarrow 0.6 mm (for 30-50 μ s fit)

- During Run-4, the quadruple **RF dampening system** has been installed and tuned.
- In Run-5 the RF has been used for production data, in order to **reduce radial oscillation of the beam**
- During Run-5/6 has been installed a **Minimally intrusive Scintillating Fiber detector (MiniSciFi)** useful to measure the time momentum correlation for E-field correction and the CBO oscillation.

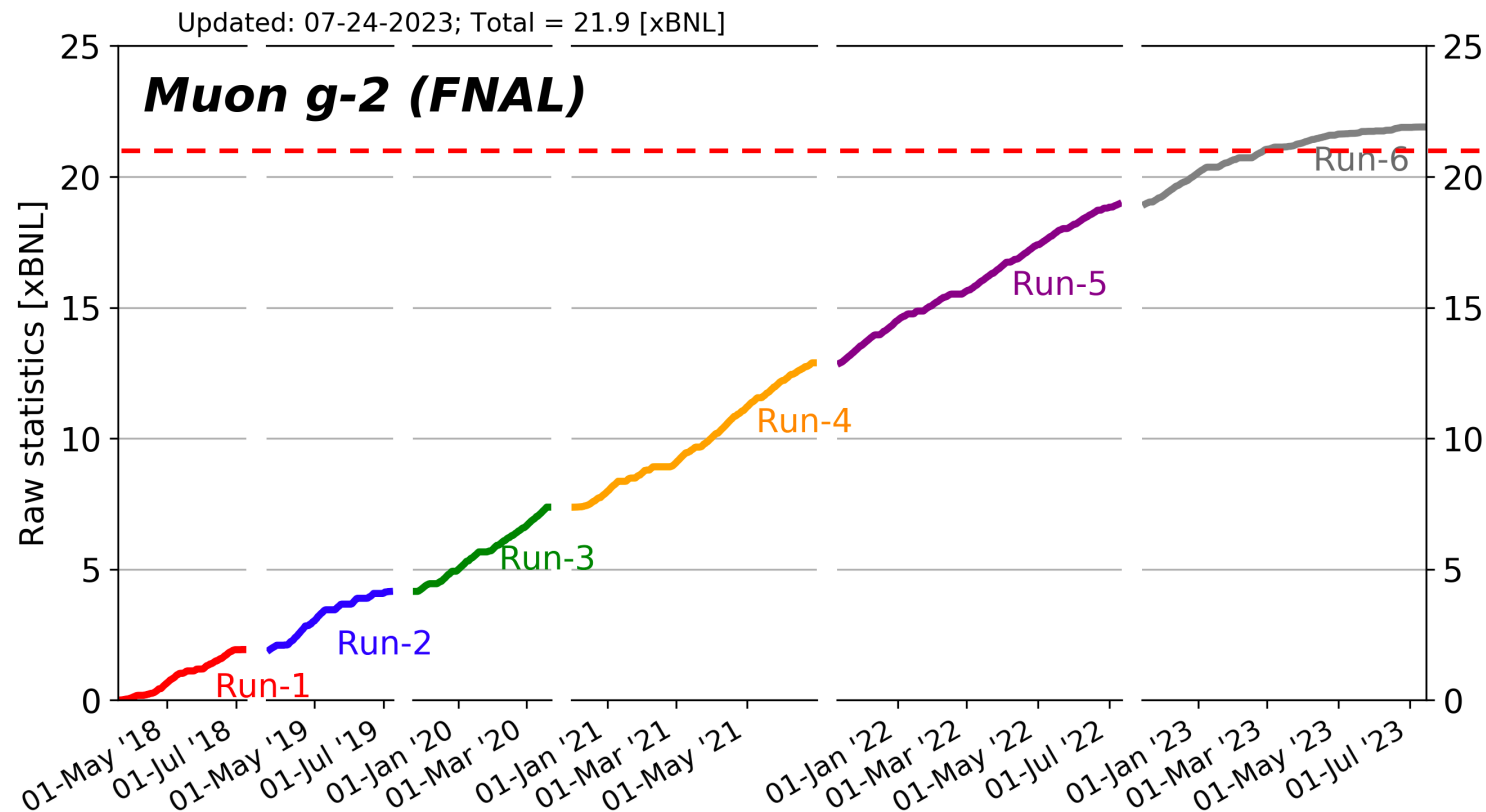




Muon $g-2$ Outlook



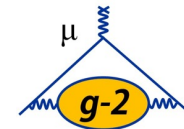
- The data acquisition is finished on 9th July 2023.
- We reached the TDR goal of 21 BNL on 27th February 2023.
- Run-4/5/6 result is expected in ~2025.



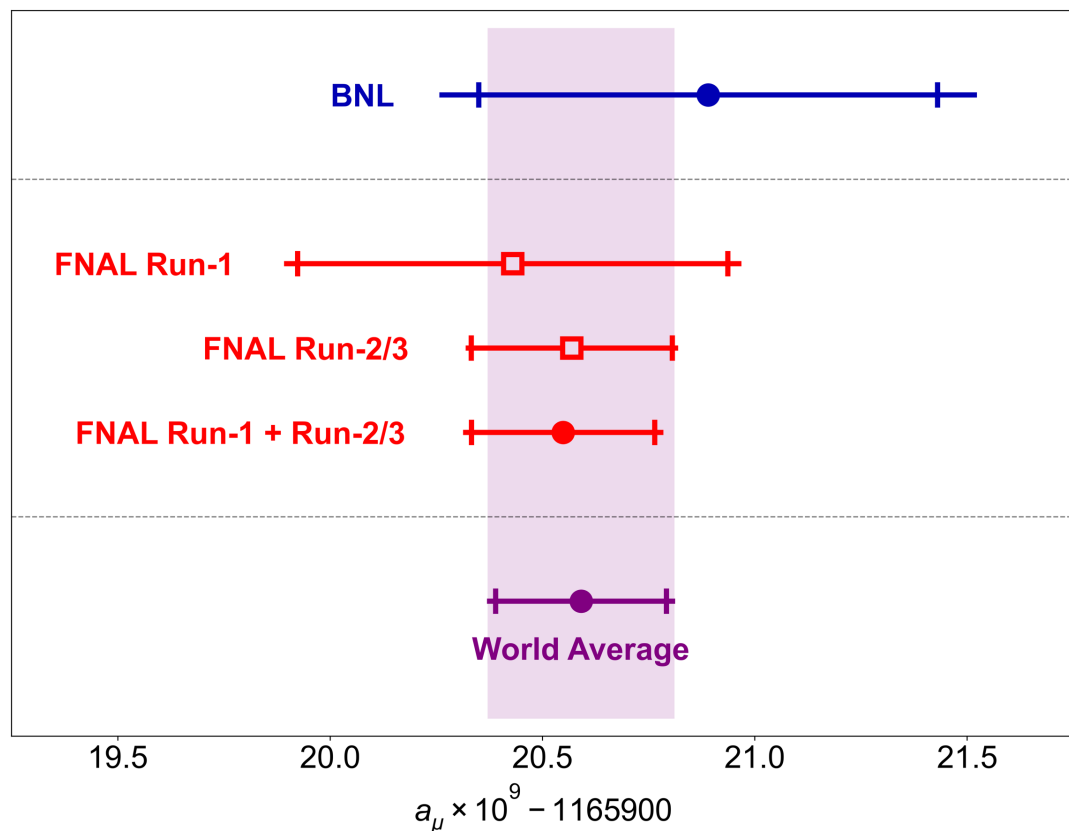
PROPOSAL GOAL



Conclusion



- We've determined a_μ to an unprecedented **203 ppb** precision

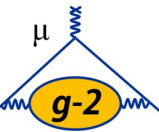


New result is in **excellent agreement** with **Run-1 & BNL**

More than **halved the total uncertainty** from Run-1

Improved our design goal with systematic uncertainty of **70 ppb**.

- There is **more data** to analyze and we'll squeeze uncertainty down further in our future results!

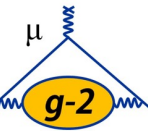


“The closer you look the more there is to see”

F. Jegerlehner

Thank you!!!

- For any question or just to have a chat – elia.bottalico@liverpool.ac.uk

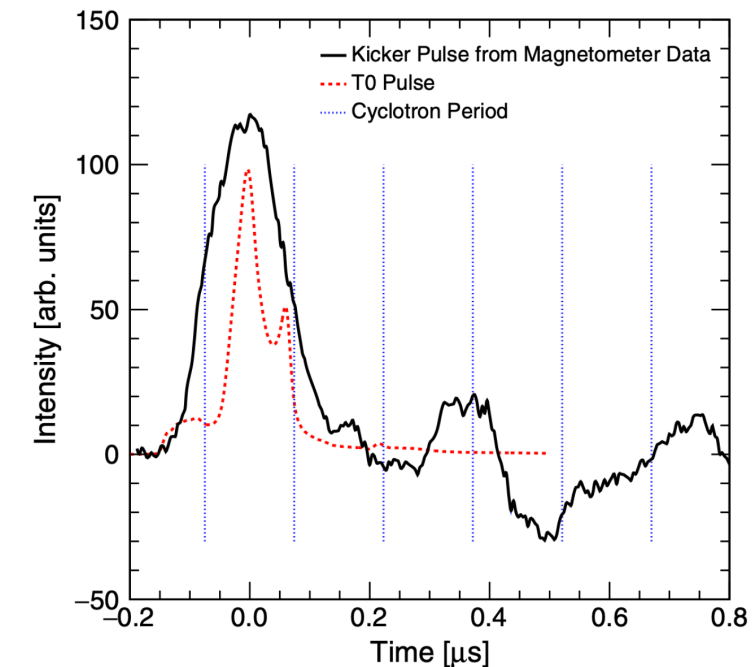
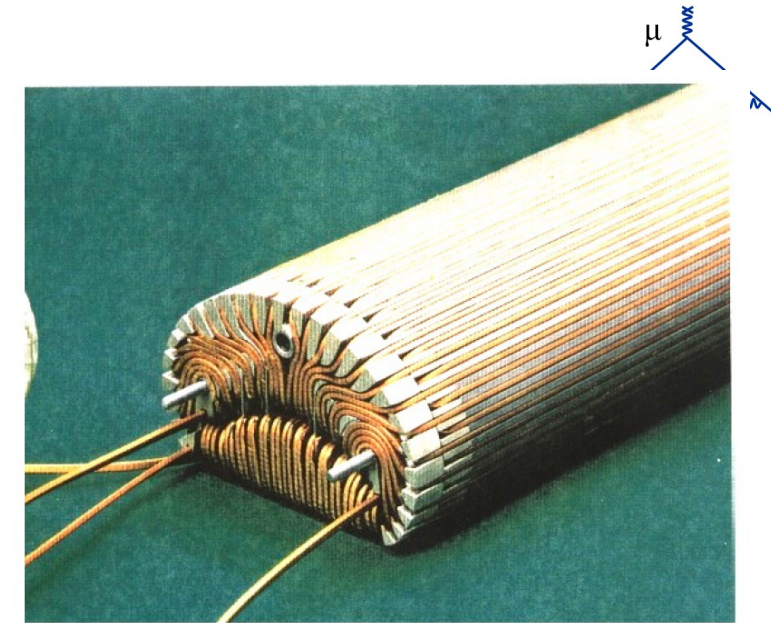


BACK-UP



Kickers and Inflector

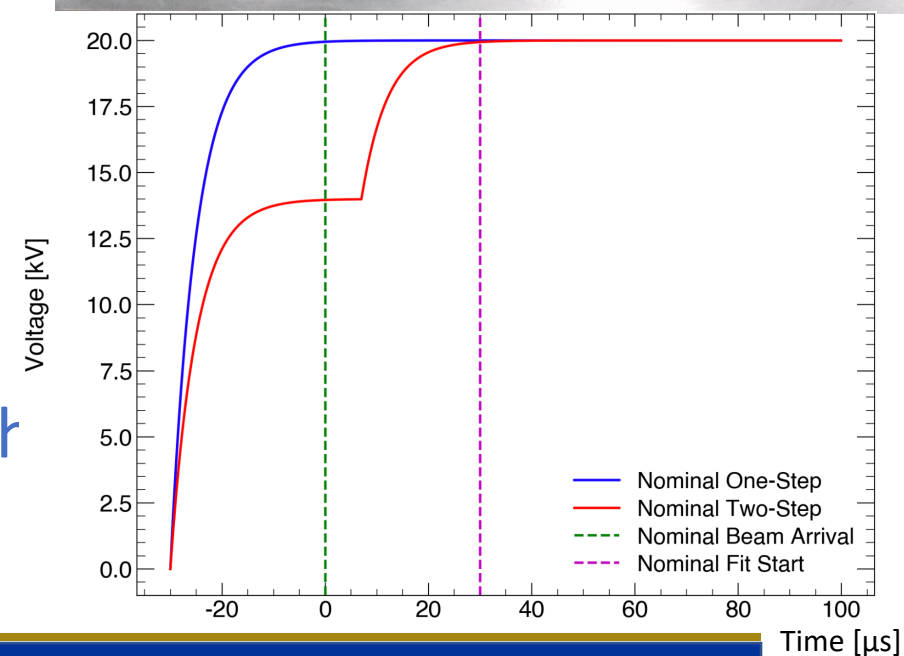
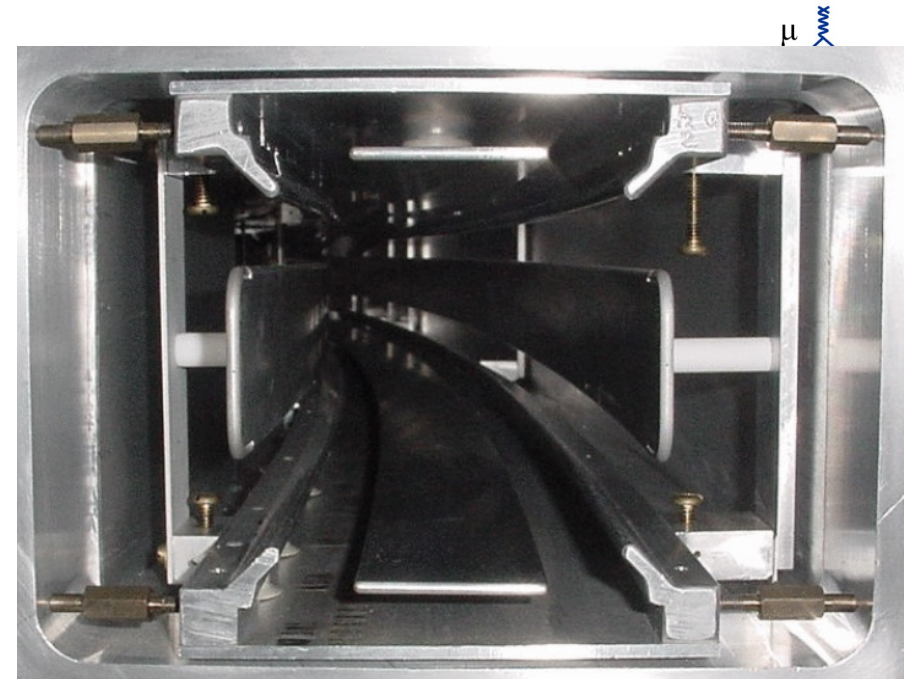
- The **inflector** cancels the storage ring field such that the muons are not deflected by the main **1.45 T** field.
- Superconducting, operational current ~ 2.6 kA.
- **3 Kickers** are necessary to inject magic momentum muons along the magic radius (7.11 m) with a required kick at order of 10 mrad.
- 4 kA current in 200 ns pulse.
- Design kick strength has been reached in Run-3 (~ 160 kV).





Quadrupoles

- The Electrostatic Quadrupoles (ESQ) system allows to strongly focus the beam vertically, four ESQ stations are symmetrically placed around the ring.
- The plates are raised from ground to operating voltage prior to each *fill* with RC charging time constants of $\sim 5 \mu\text{s}$.
- This procedure, known as **scraping**, initially displaces the beam vertically and horizontally with respect to the central closed orbit.

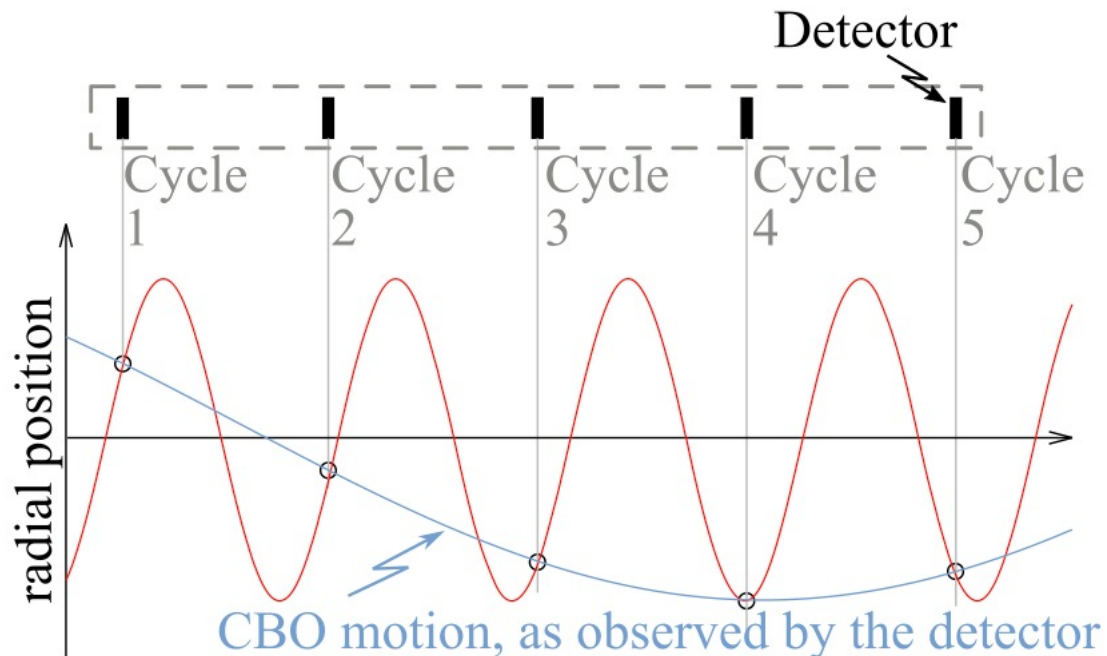




ω_a measurement – CBO oscillation



- Given the restoring force by radial magnetic field, the beam oscillates radially (vertically too) as the betatron frequency: $\omega_{BO} = \omega_C \sqrt{1 - n}$, where n is the field-index.
- The beam is measured by detectors, calorimeters and trackers.
- The $\omega_{BO} < \omega_C$, so calorimeters see a different phase at each turn, measuring an oscillation called **Coherent Betatron Oscillation (CBO)**, given by $\omega_{CBO} = \omega_C - \omega_{BO}$



$$2\pi f_{CBO} = \omega_C - \omega_{BO} = \omega_C(1 - \sqrt{1 - n})$$

$$\omega_{CBO} = 2.34 \text{ rad}/\mu\text{s}$$

Where $T_C \sim 0.149 \text{ ns}$ and $n \sim 0.108$



Beam dynamics correction to ω_a : C_e



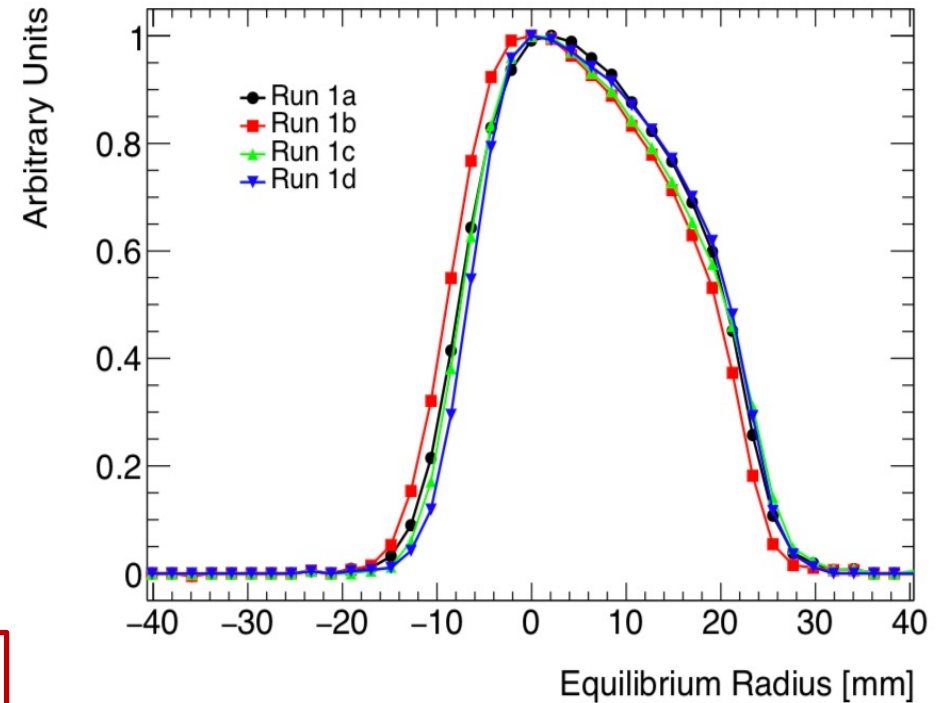
Considering the extended expression of the spin precession frequency in a magnetic field:

$$\vec{\omega}_a = \frac{e}{m} \left[a_\mu \vec{B} - \left(a_\mu - \frac{1}{\gamma^2 - 1} \right) (\vec{\beta} \times \vec{E}) - a_\mu \left(\frac{\gamma}{\gamma + 1} \right) (\vec{\beta} \cdot \vec{B}) \vec{\beta} \right]$$

This term introduces a bias on ω_a that needs to be corrected by Electric Field correction:

$C_e = 2n(1 - n)\beta^2 \frac{\langle x_e^2 \rangle}{R_0^2}$ is proportional to the equilibrium radius distribution x_e .

$$C_e \sim 489 \text{ ppb}$$





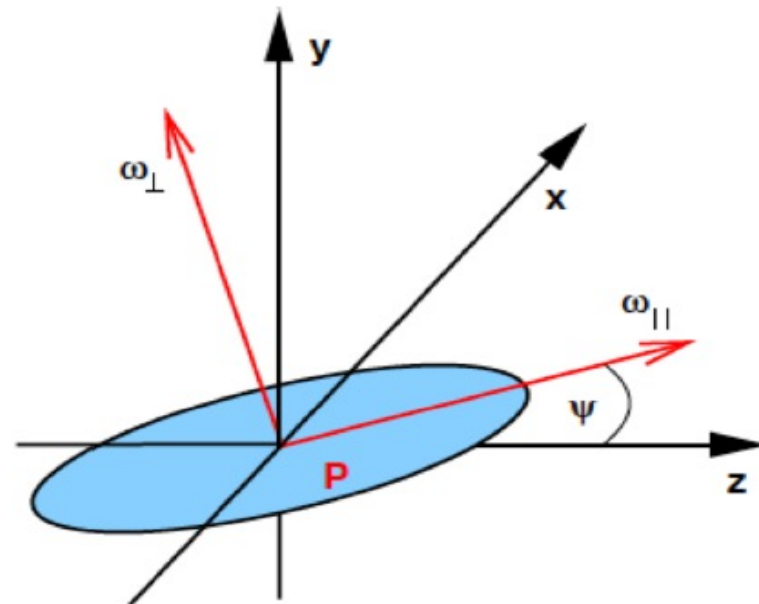
Beam dynamics correction to ω_a : C_p



Considering the extended expression of the spin precession frequency in a magnetic field:

$$\vec{\omega}_a = \frac{e}{m} \left[a_\mu \vec{B} - \left(a_\mu - \frac{1}{\gamma^2 - 1} \right) (\vec{\beta} \times \vec{E}) - a_\mu \left(\frac{\gamma}{\gamma + 1} \right) (\vec{\beta} \cdot \vec{B}) \vec{\beta} \right]$$

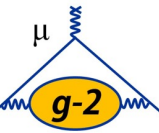
$$C_p \sim 200 \text{ ppb}$$



C_p : the pitch correction $C_p = n \langle A_y^2 \rangle / 4R_0^2$ depends on amplitude vertical oscillation (A_y).



a_μ systematic sources



Many systematics come from effects that change the phase of the detected positrons over time and introduce a bias on ω_a :

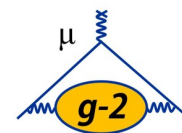
$$\begin{aligned}\cos(\omega_a t + \phi(t)) &= \cos(\omega_a t + \phi_0 + \phi' t + \dots) \\ &= \cos((\omega_a + \phi')t + \phi_0 + \dots)\end{aligned}$$

In general, anything that changes from early-to-late within each muon fill can be a cause of systematic error, as:

- Beam distortion
- Muon losses
- Varying lifetime
- Rate dependent reconstruction



Beam dynamics correction to ω_a : C_{lm}



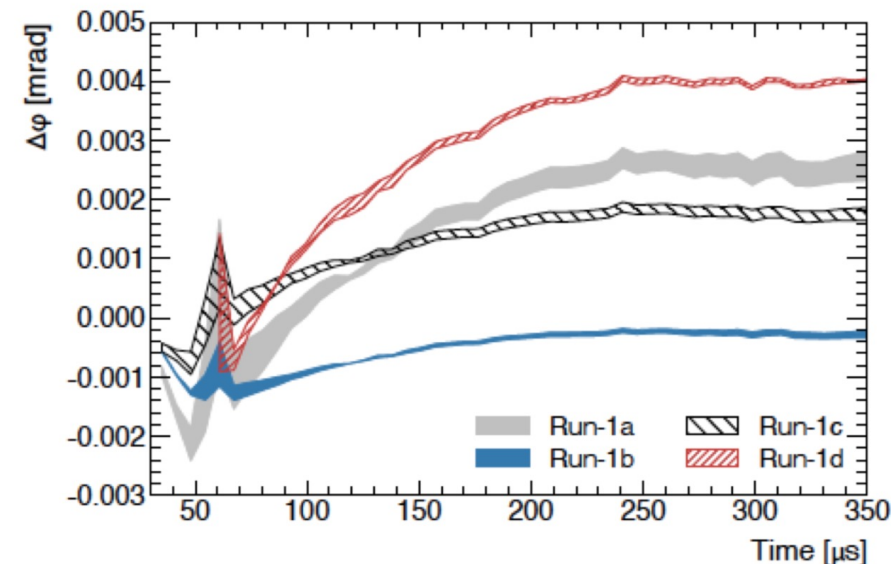
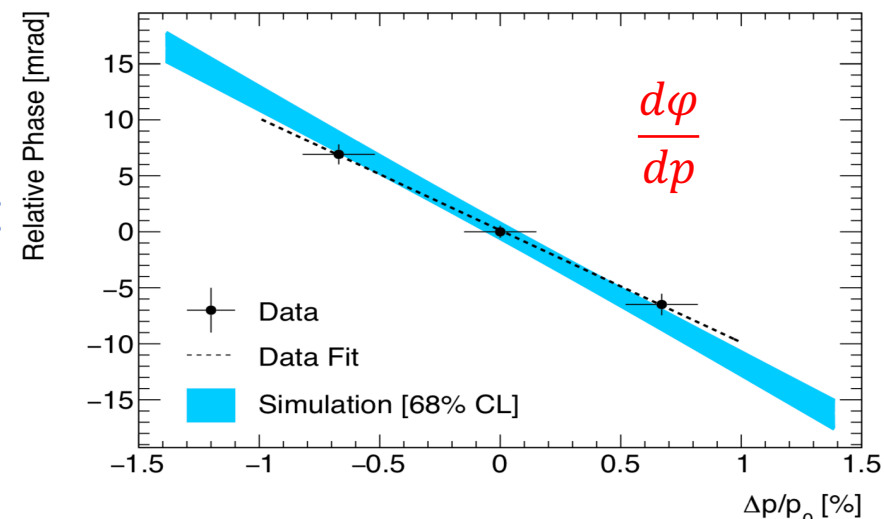
C_{lm} : describes the motion introduced on ω_a phase due to the loss of muon during the *fill*. It's explained by:

1. Muons with different **momentum** have different **phase**;
2. The number of loss muon change as function of momentum.

$$\Delta\omega_a = \frac{d\varphi}{dt} = \frac{d\varphi}{dp} \cdot \frac{dp}{dt}$$

$$C_{lm} < 20 \text{ ppb}$$

$$d\varphi_0/dp = (-10.0 \pm 1.6) \text{ mrad}/(\% \Delta p/p)$$

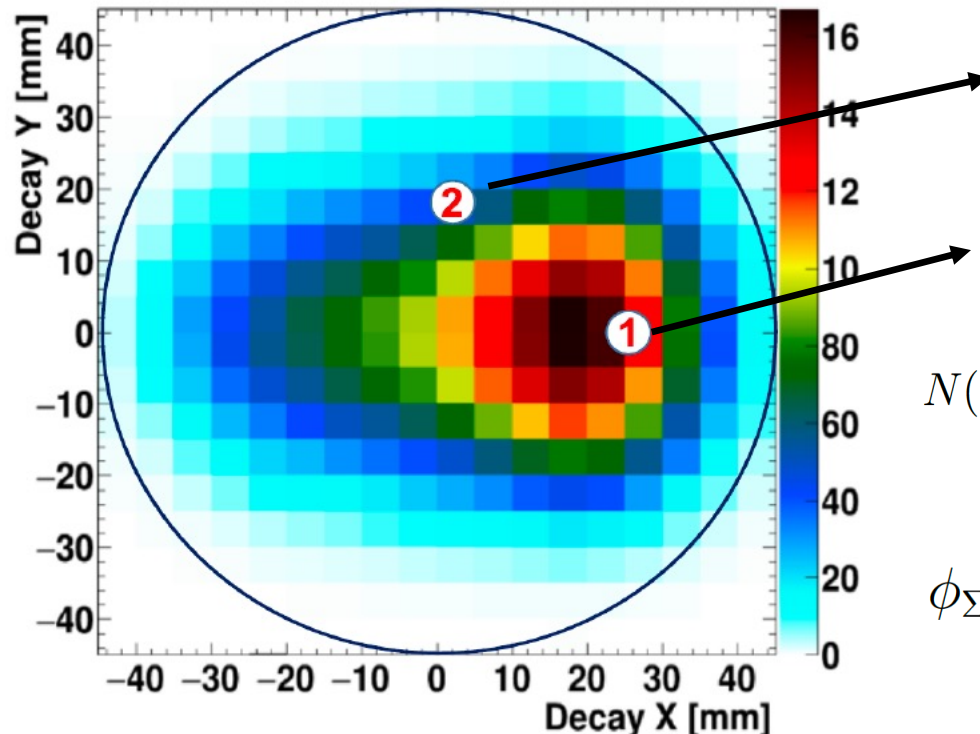




Beam dynamics correction to ω_a : C_{pa}



- The measured $g-2$ phase of the muon is decay vertex position dependent.
- It is obtained as weighted average of the phases measured by each (x,y) pair position.



$$N_2(t) = N_{02}e^{-t/\tau} [1 + A_2 \cos(\omega_a t + \phi_2)]$$

$$N_1(t) = N_{01}e^{-t/\tau} [1 + A_1 \cos(\omega_a t + \phi_1)]$$

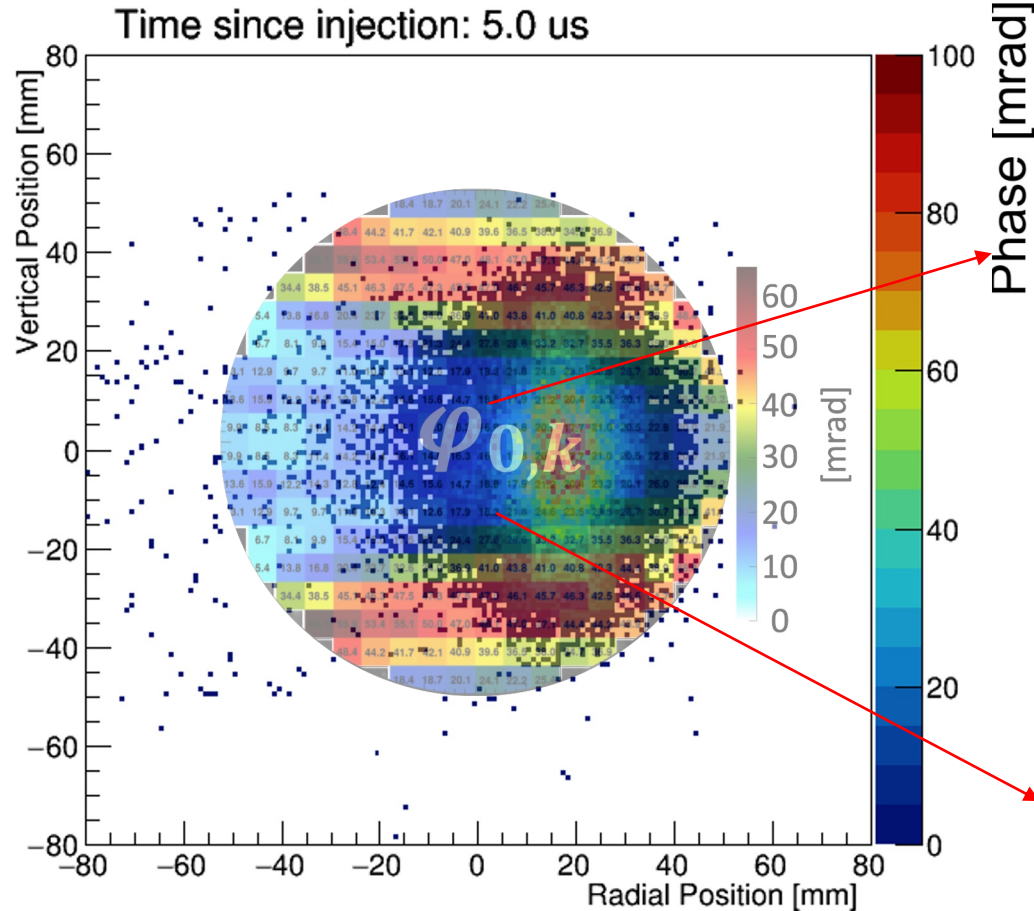
$$N(t) = N_1(t) + N_2(t) = N_{\Sigma}e^{-t/\tau} [1 + A_{\Sigma} \cos(\omega_a t + \phi_{\Sigma})]$$

$$\phi_{\Sigma} = \arctan \frac{N_{01}A_1 \sin(\phi_1) + N_{02}A_2 \sin(\phi_2)}{N_{01}A_1 \cos(\phi_1) + N_{02}A_2 \cos(\phi_2)}$$

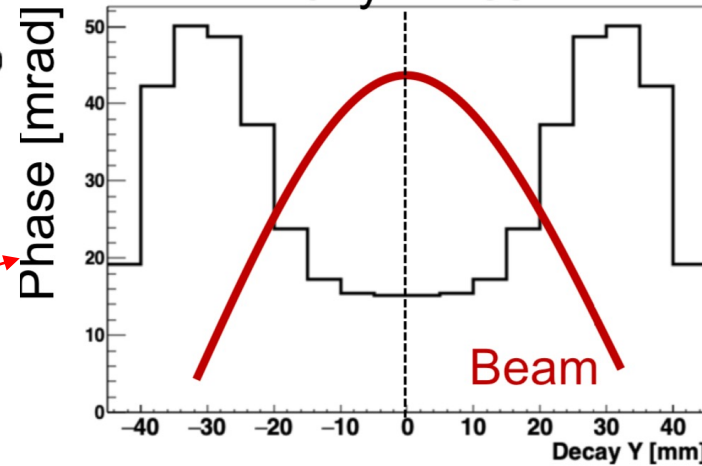
Phase acceptance: Beam Motion Effects



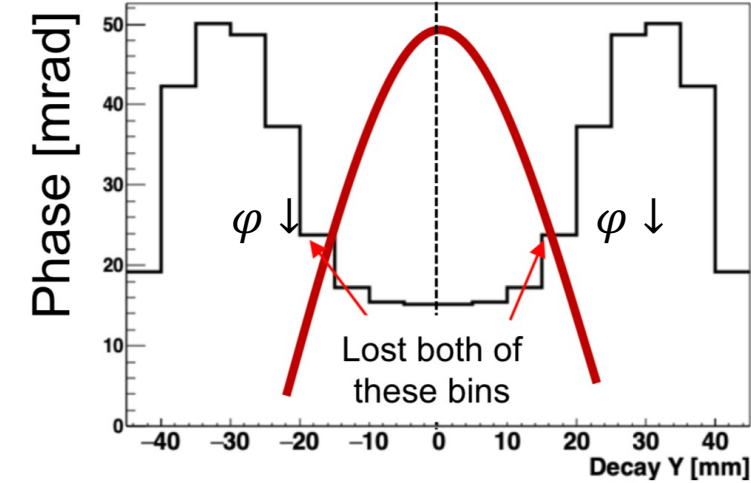
VERTICAL WIDTH VARIATION



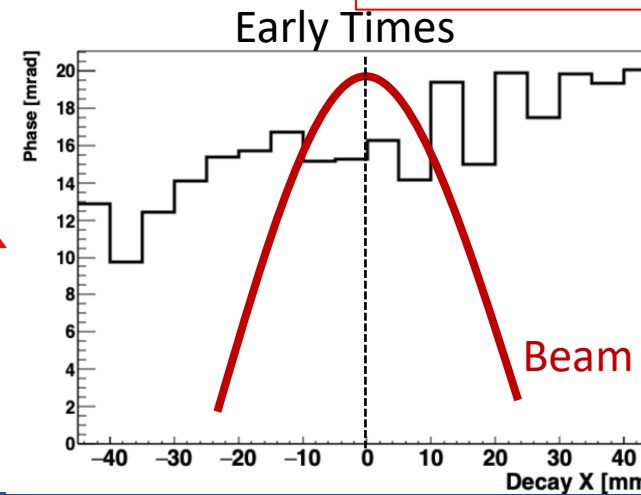
Early Times



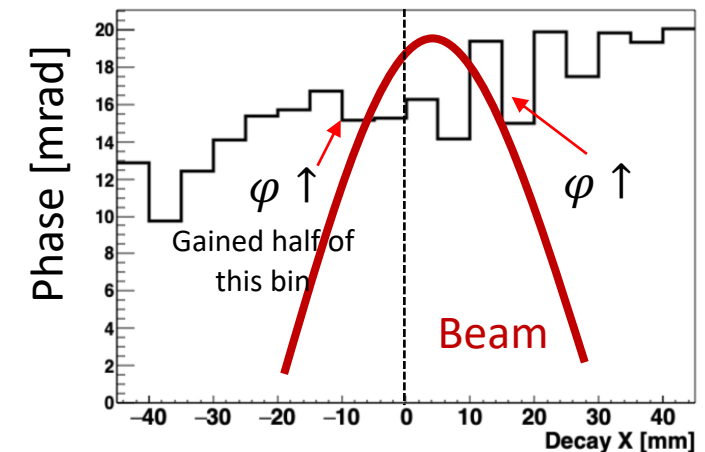
Late Times



RADIAL MEAN VARIATION



Late Times





Beam dynamics correction to ω_a : C_{pa}



C_{pa} : it is a Phase Acceptance effect. It is due to:

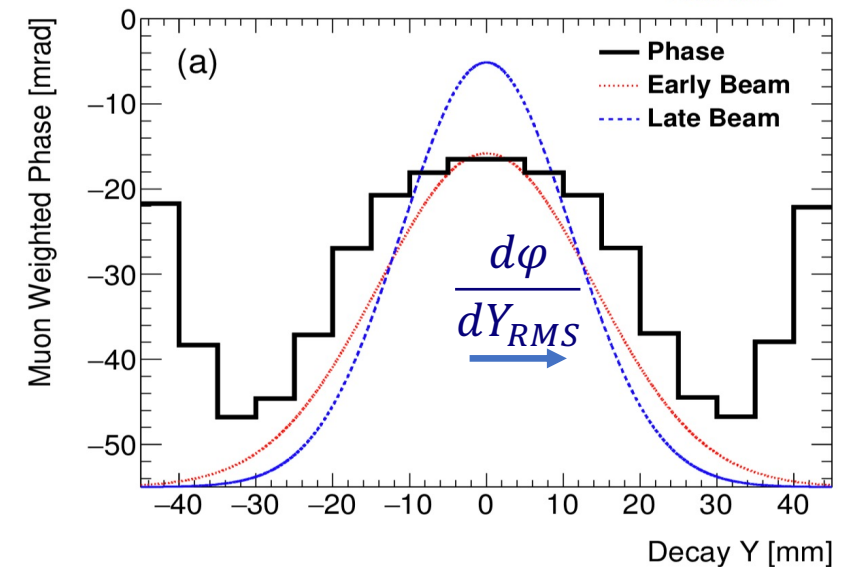
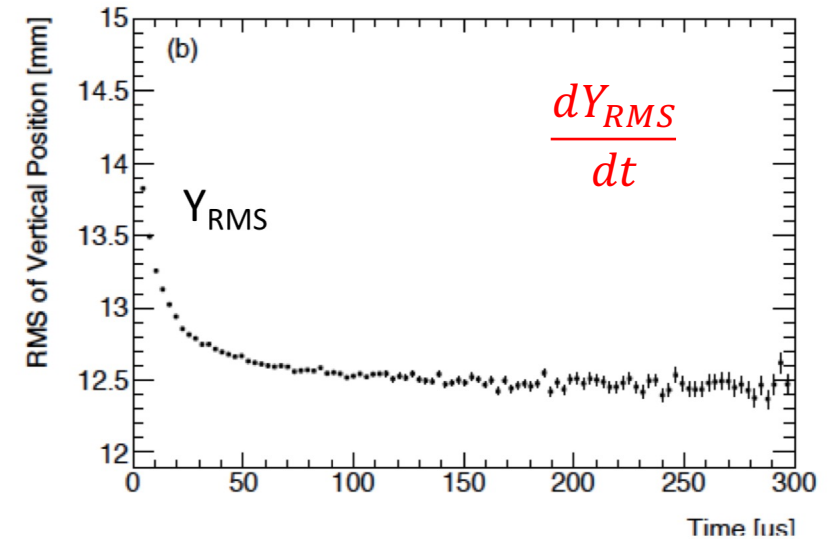
1. Beam variation during the *fill*;
2. Phase measured as function of the decay position.

$$\Delta\omega_a = \frac{d\varphi}{dt} = \frac{dY_{RMS}}{dt} \cdot \frac{d\varphi}{dY_{RMS}}$$

The effect was large in Run1 due to *broken resistors*

$$C_{pa} \sim 180 \text{ ppb}$$

We expect a reduction in Run2/3 ($\sim 50\text{ppb}$ / $\sim 20\text{ppb}$)





Beam dynamics correction to ω_a



- These are the results for the BD corrections from Run-1, the phase acceptance (C_{pa}) correction was one of the topic I addressed during my PhD.
- Now analysis is ongoing to finalize the Run-2/3 beam dynamics corrections, stay tuned!

	Correction Factor [ppb]	Uncertainty [ppb]
ω_a (stat.)	—	434
ω_a (syst.)	—	56
f_b/f_0	—	2
C_e	489	53
C_p	180	13
C_{ml}	-11	5
C_{pa}	-158	75
$f_{calib} \langle \omega'_p(x, y, \phi) \cdot M(x, y, \phi) \rangle$	—	56
B_q	-17	92
B_k	-27	37
$\mu'_p(34.7^\circ C)/\mu_e$ [PCK77]	—	10
m_μ/m_e [LAMPF-99; CD-2018]	—	22
$g_e/2$ [HFG08]	—	0
Total Systematic	—	157
Total Fundamental Factors	—	25
Total	544	461