

Neutrino emissions associated with late-time emissions of GRB

Riki Matsui, Tohoku University

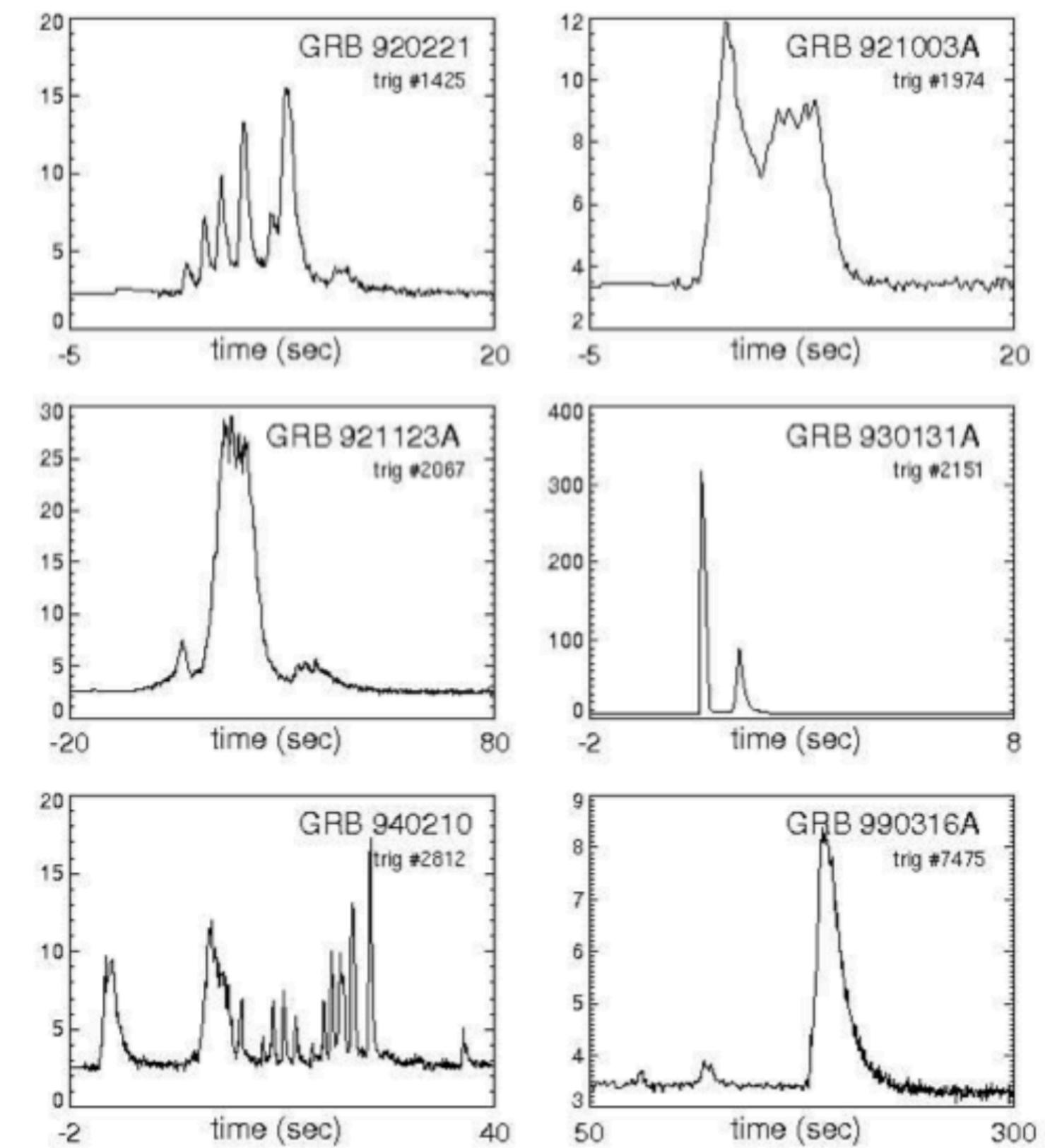
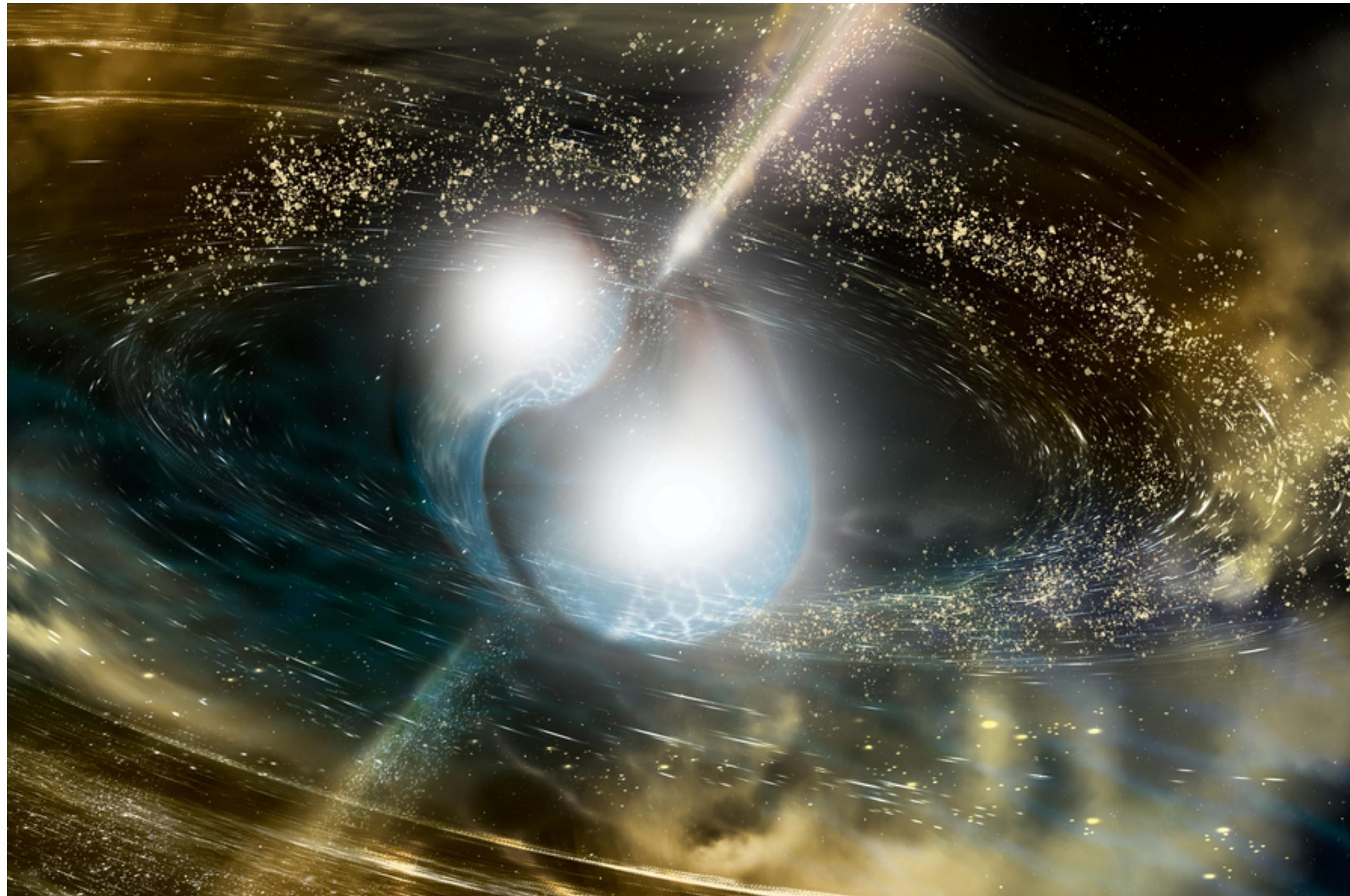
Co-auther : Shigeo S. Kimura

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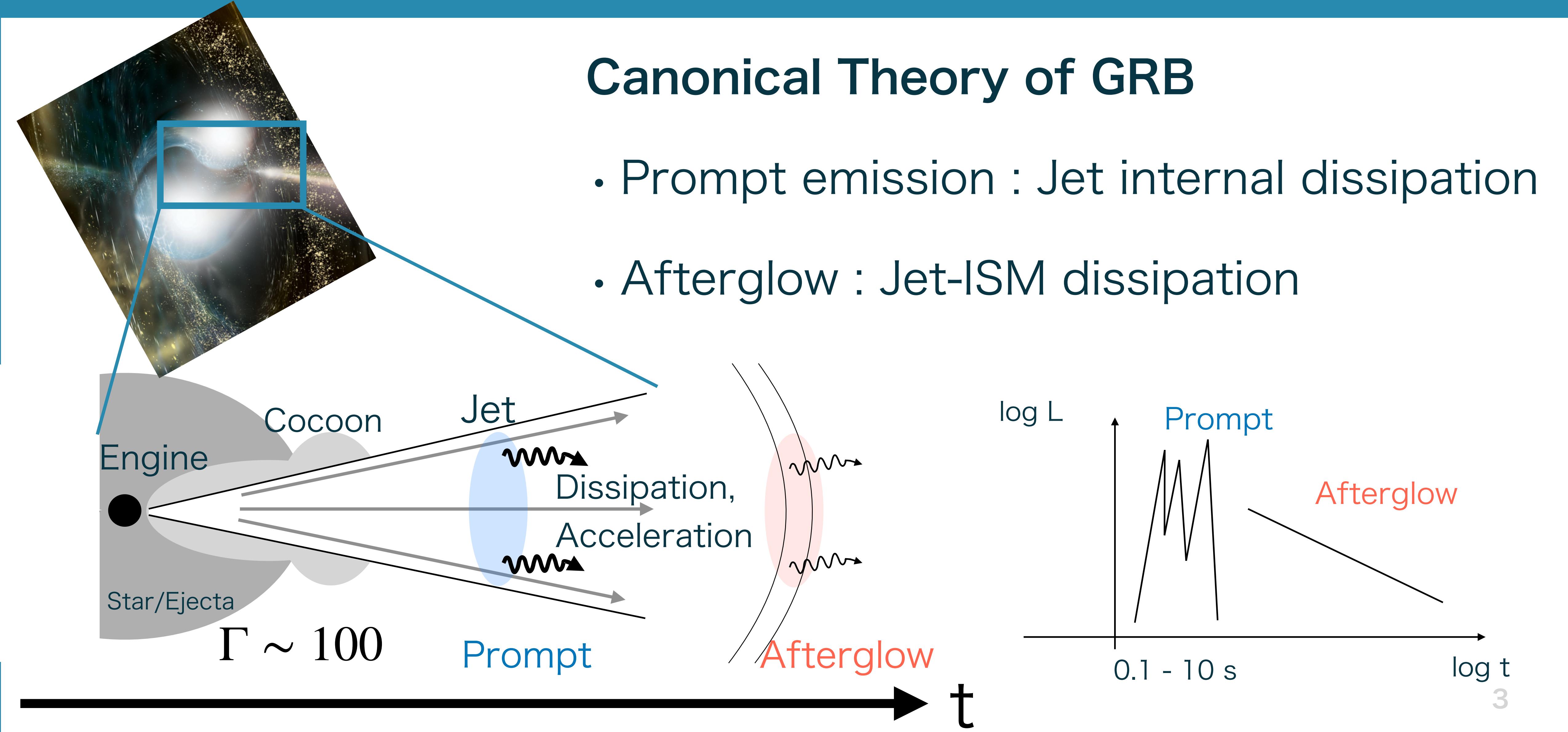
Gamma-Ray Burst (GRB)

The most luminous transient (in EM wave).

Mechanism ? Central object ? (Progenitor ?)

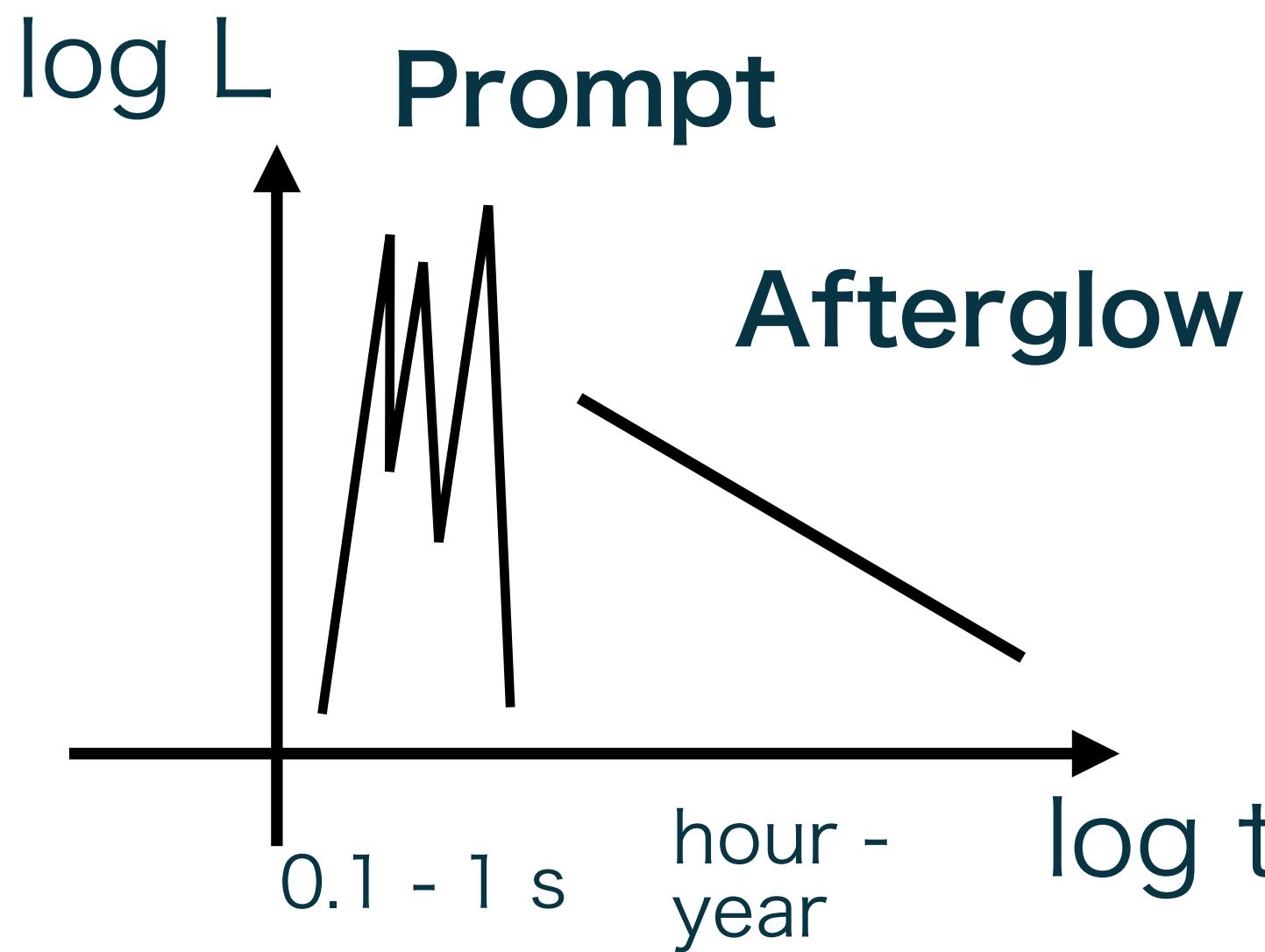


Relativistic Jets with $\nu \sim c, \Gamma \sim 100$

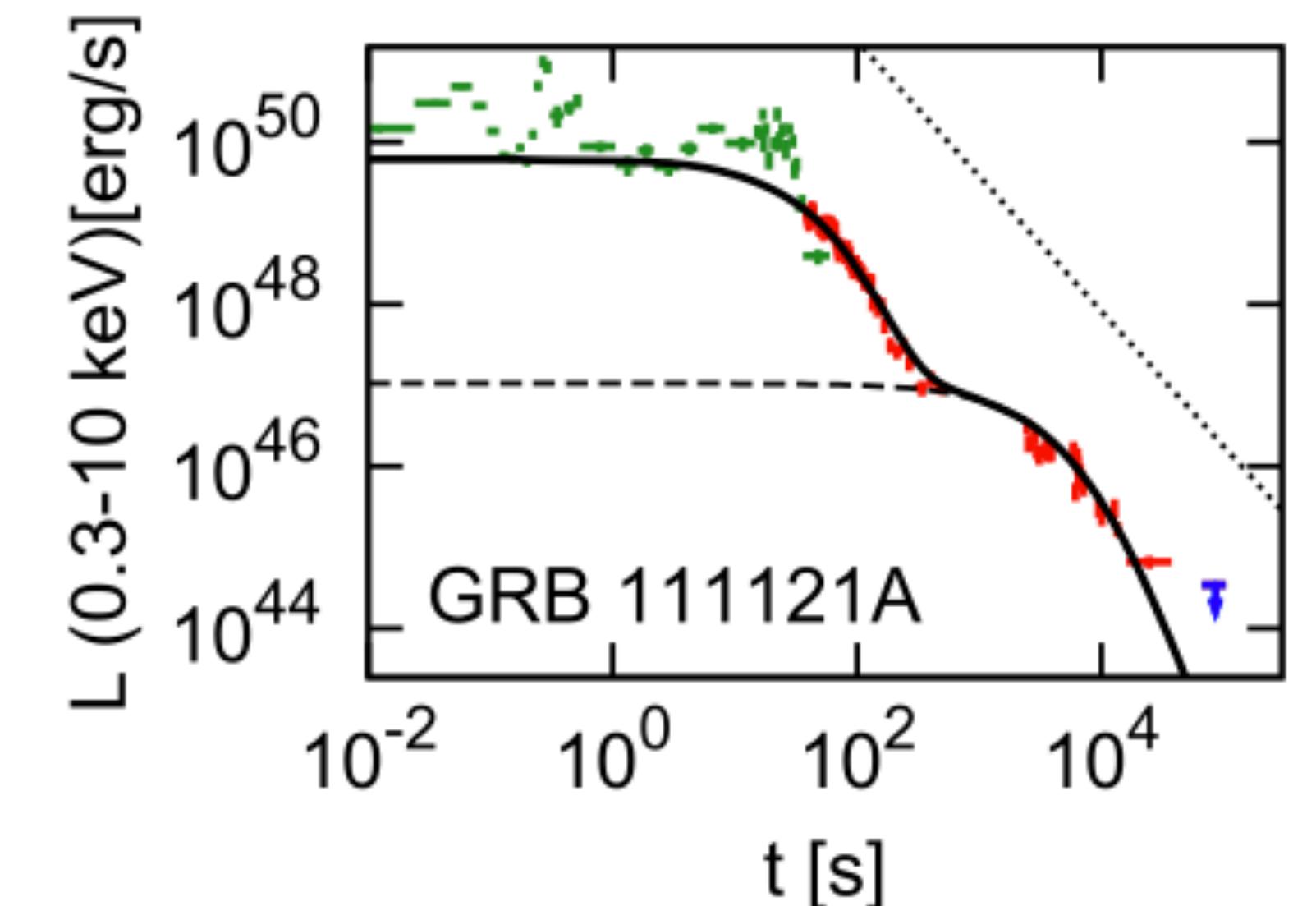


Late-time Emission of Gamma-Ray Burst

Theory of GRB

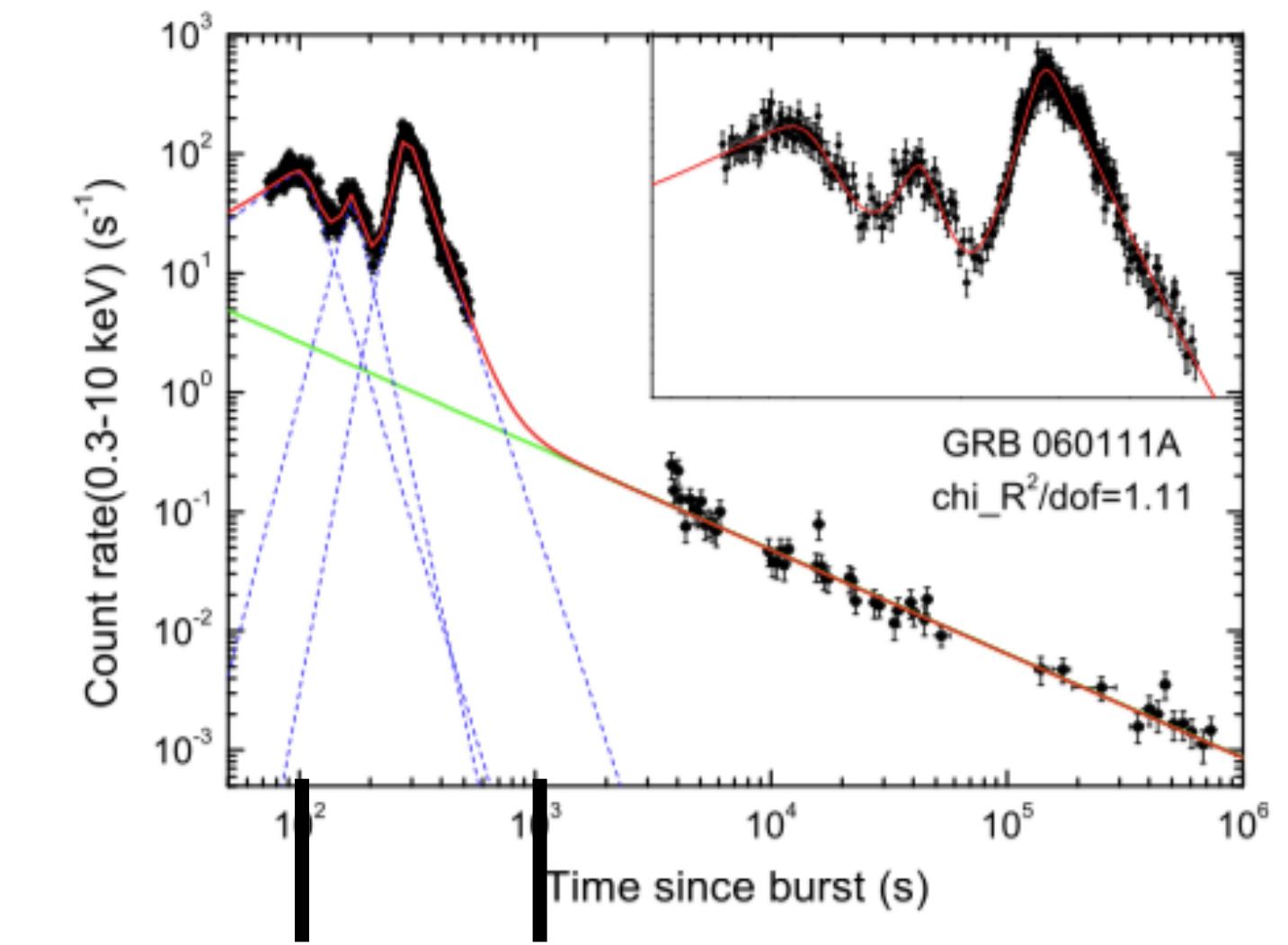


Short GRB



“Extended Emission”
(~ 2/3 of SGRBs)

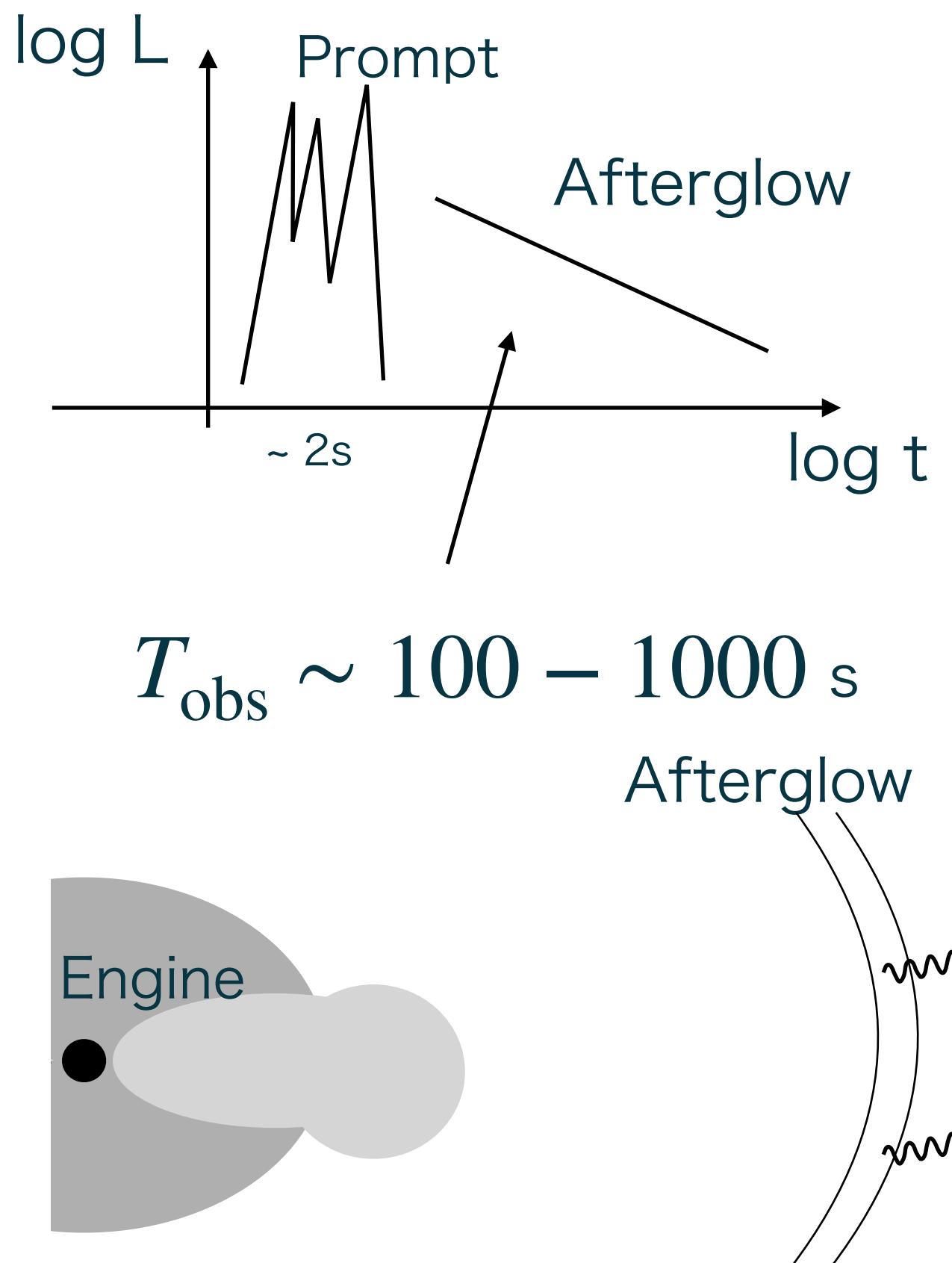
Long GRB



“X-ray Flare”
(~ 1/3 of LGRBs)

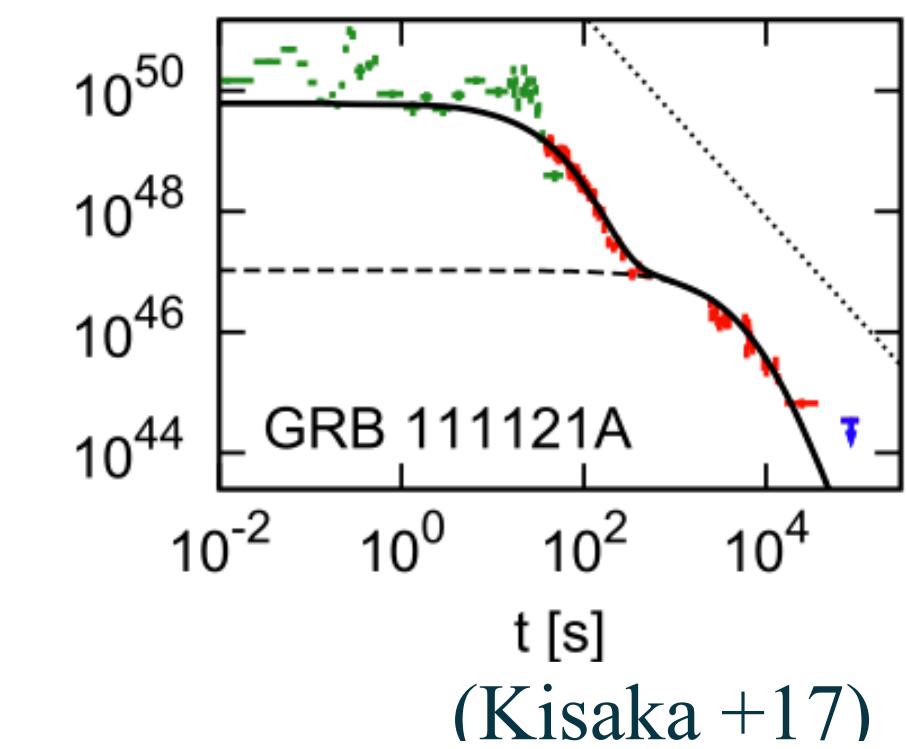
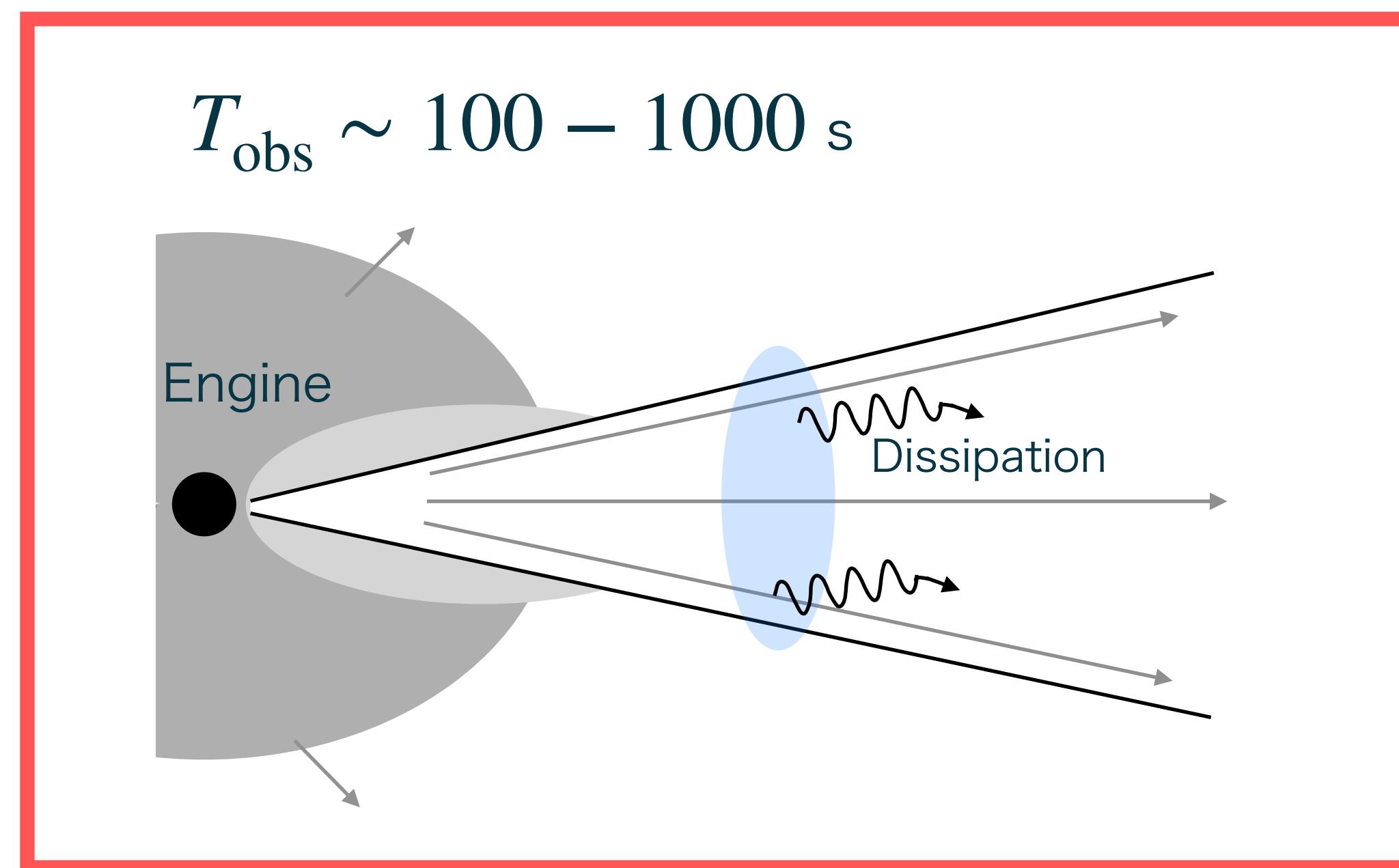
Prolonged central engine activity

Theory of GRB



Standard Scenario for late-time emission : Prolonged engine activity & Jet

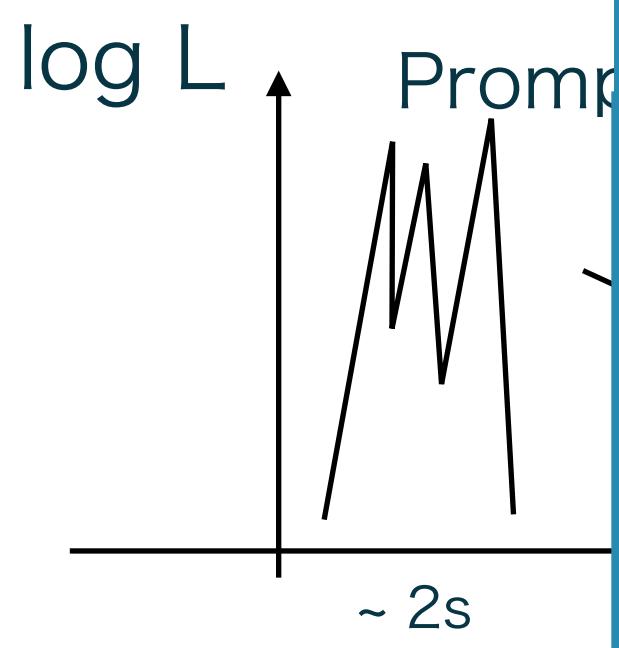
(e.g. Ioka et al. 2005, Perna et al. 2006)



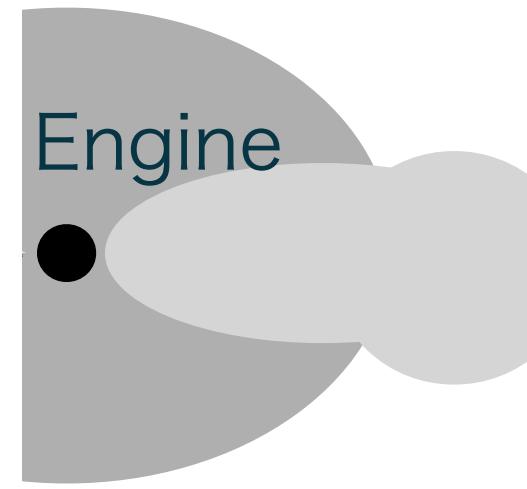
Problem : Physical quantities are unknown.

Prolonged central engine activity

Theory



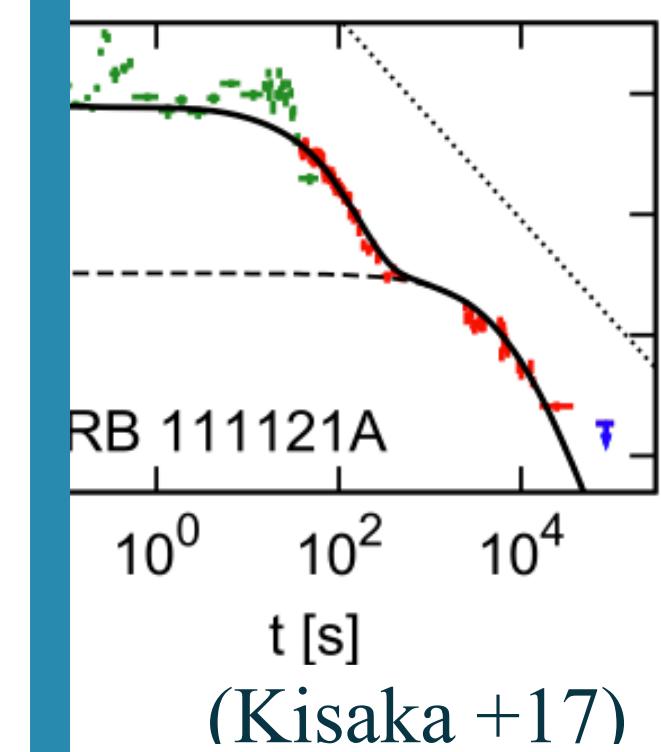
$T_{\text{obs}} \sim$



Question

Standard Scenario for late-time emission :
Prolonged engine activity & Jet

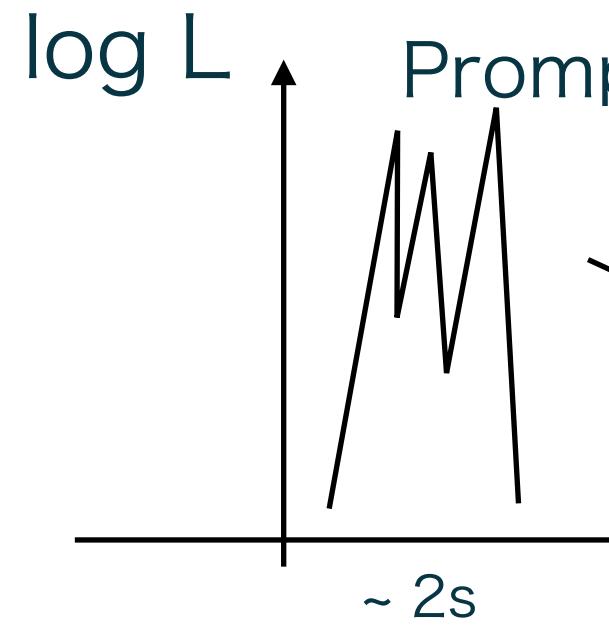
Can we characterize it by
other observations ?



// Problem : Physical quantities are unknown.

Prolonged central engine activity

Theory

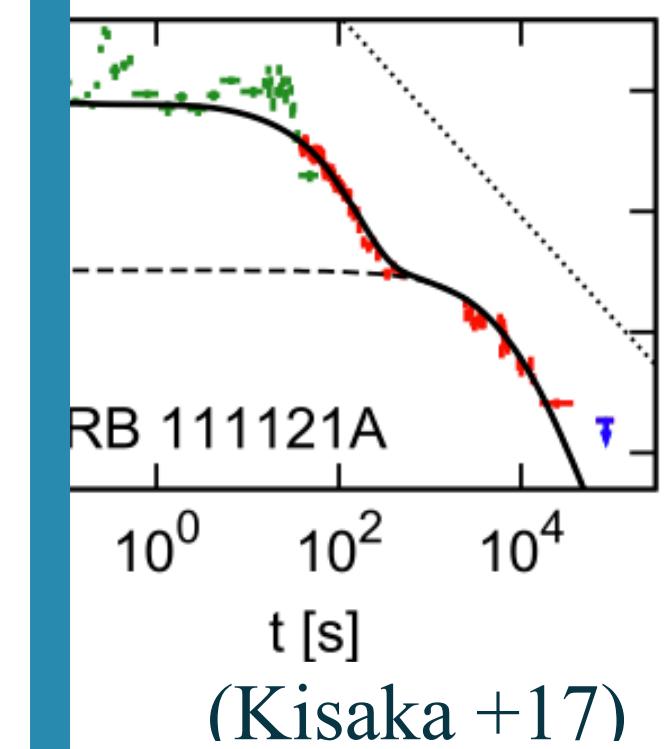


Question

Standard Scenario for late-time emission :
Prolonged engine activity & Jet

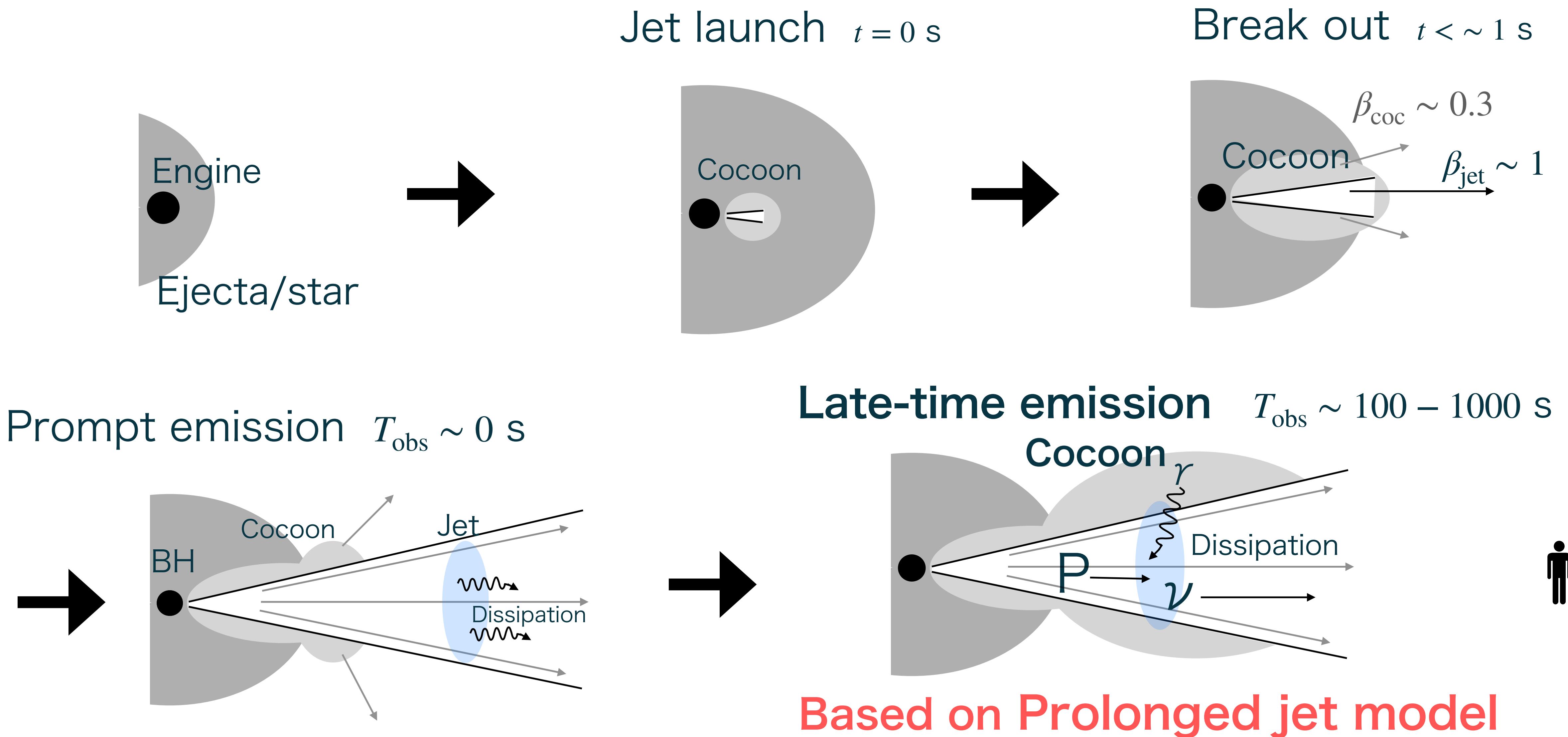
Can we characterize it by
other observations ?

A. Yes. (Neutrino)

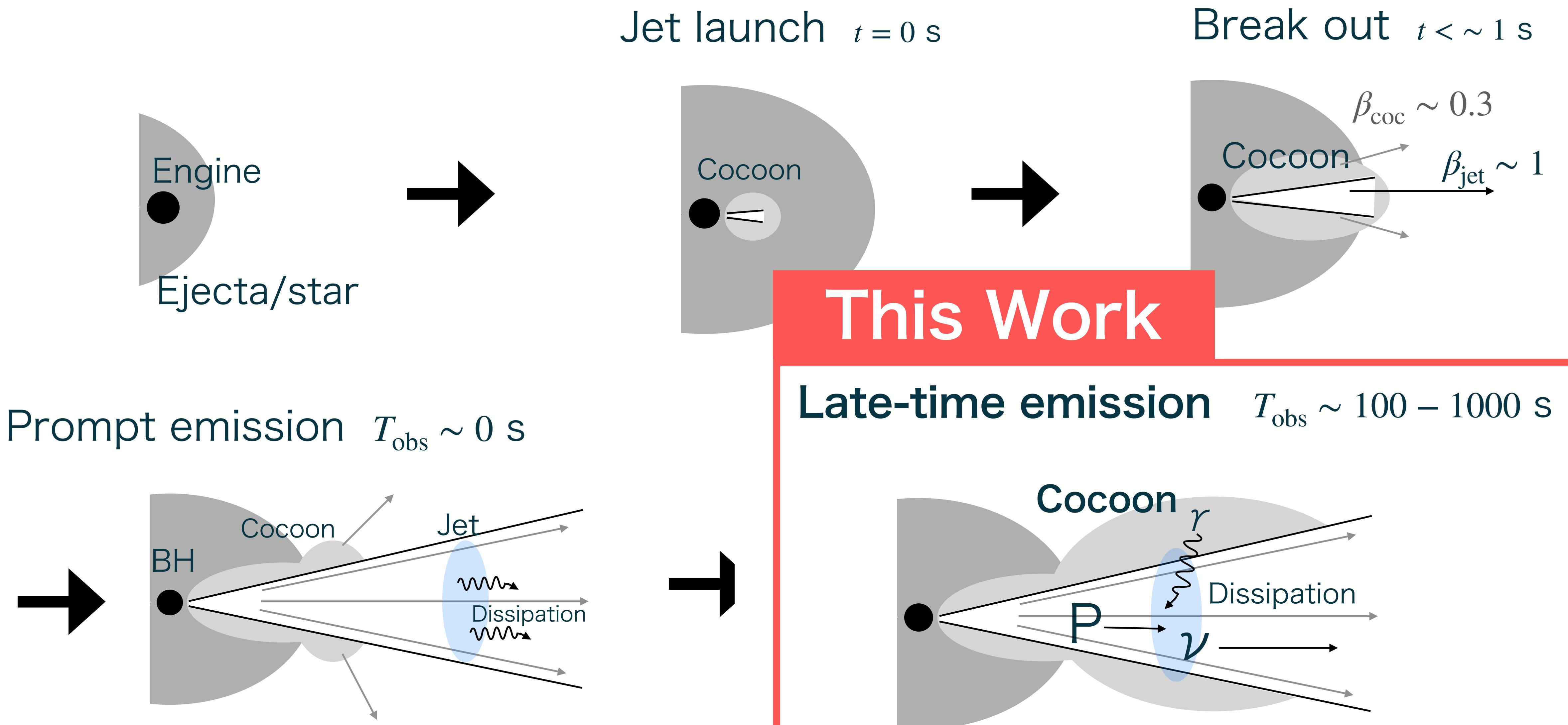


//
Problem : Physical quantities are unknown.

Prolonged Engine and Cocoon Scenario

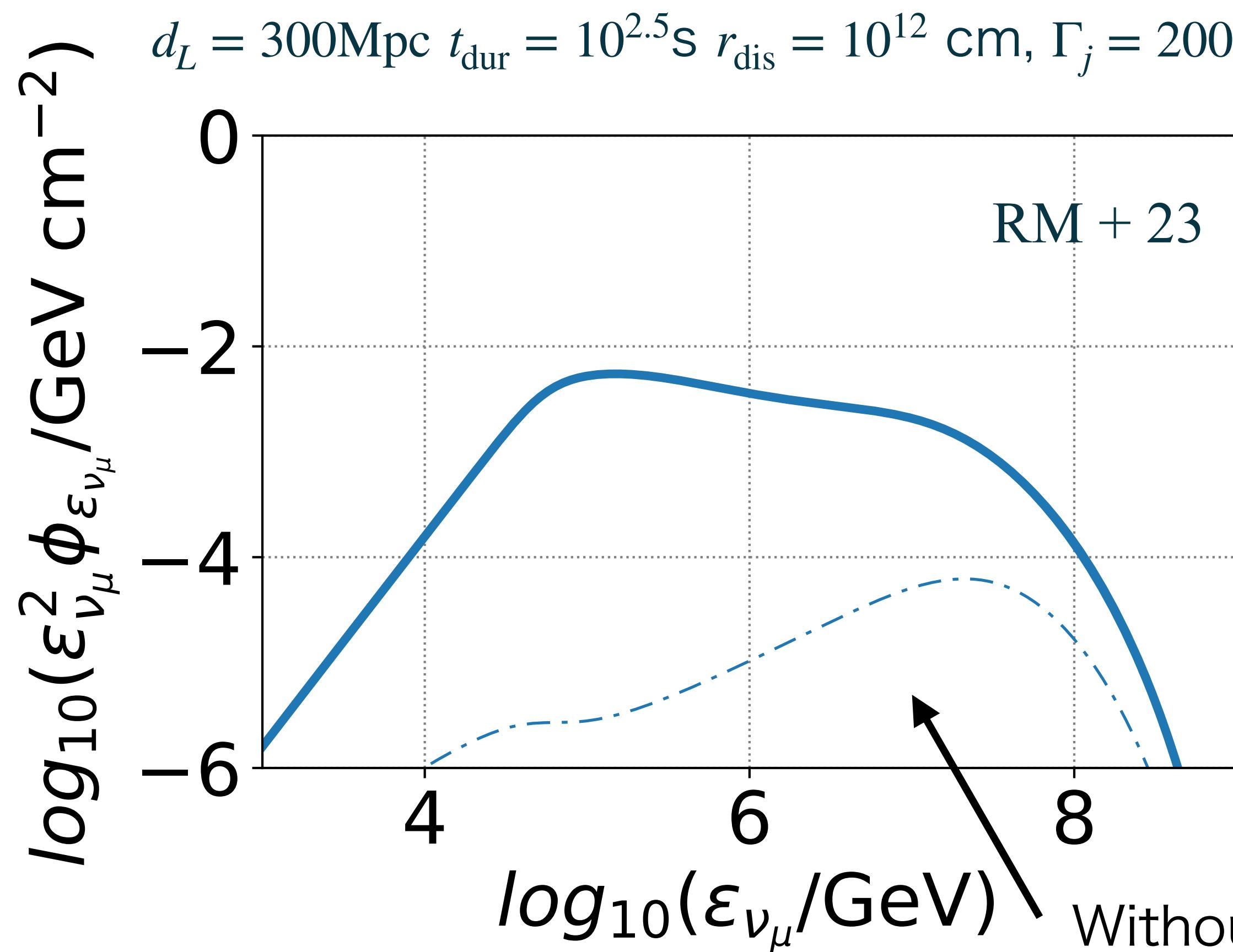


Prolonged Engine and Cocoon Scenario

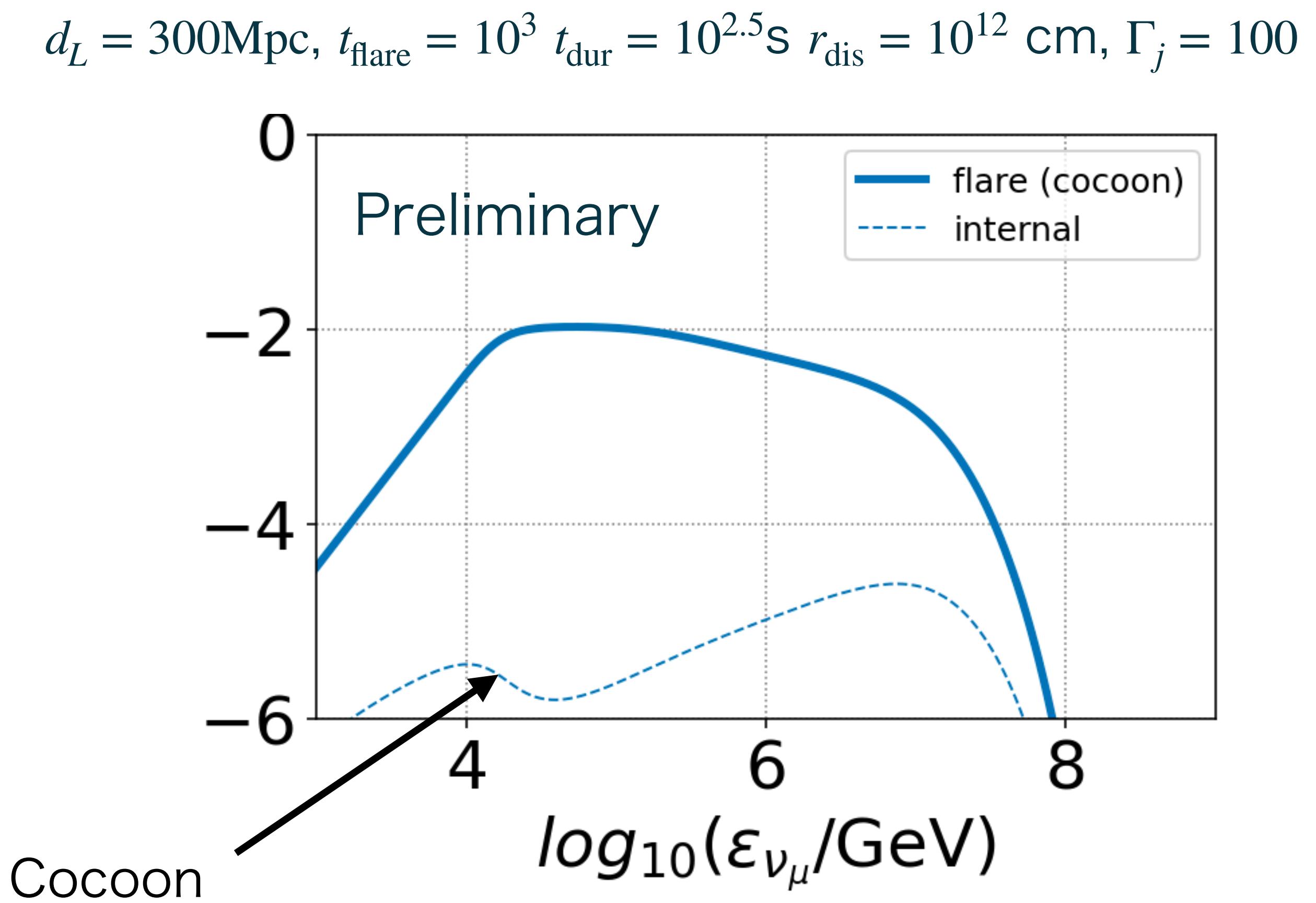


Results 1 : Neutrino Spectra

Extended Emission



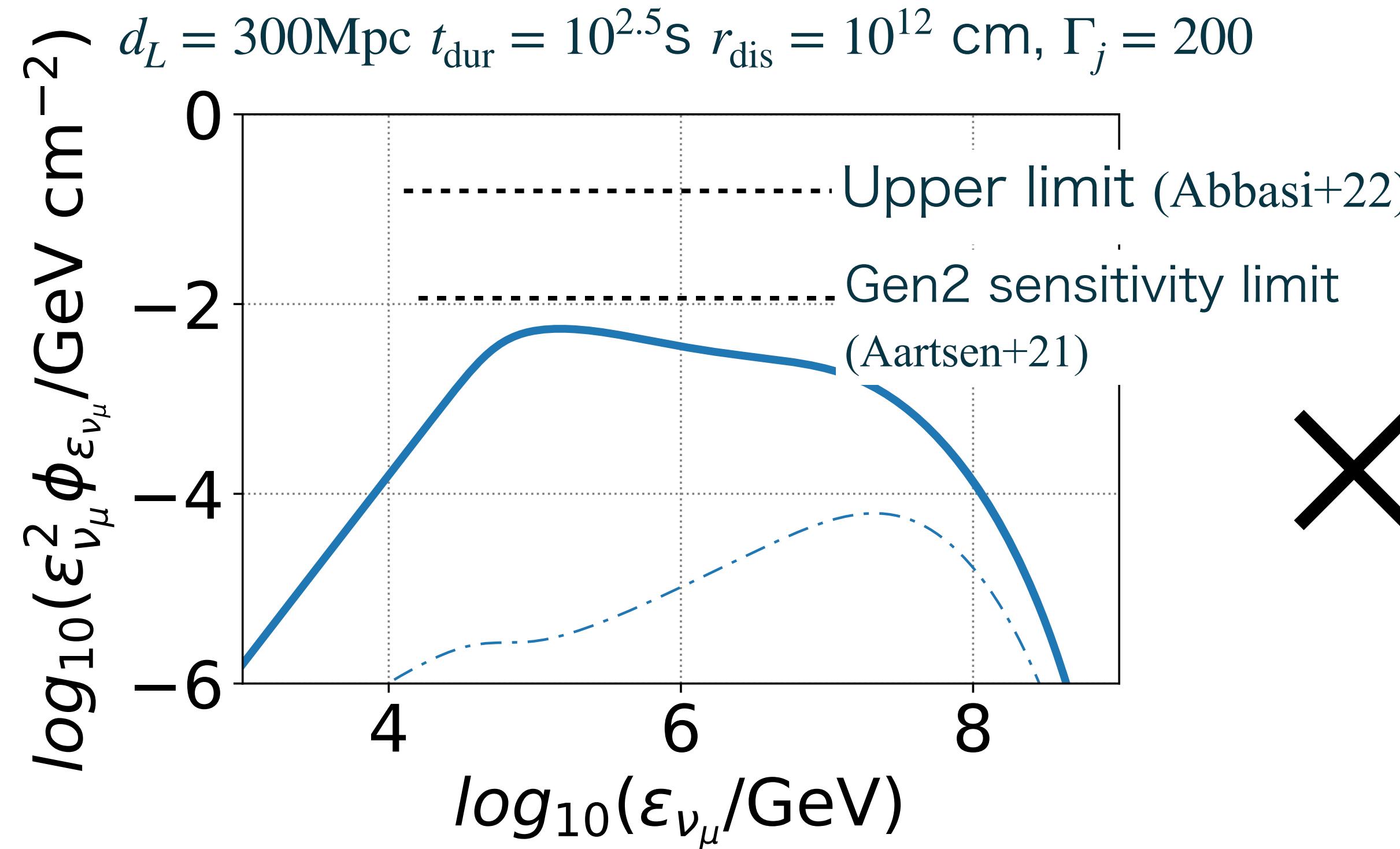
X-ray Flare



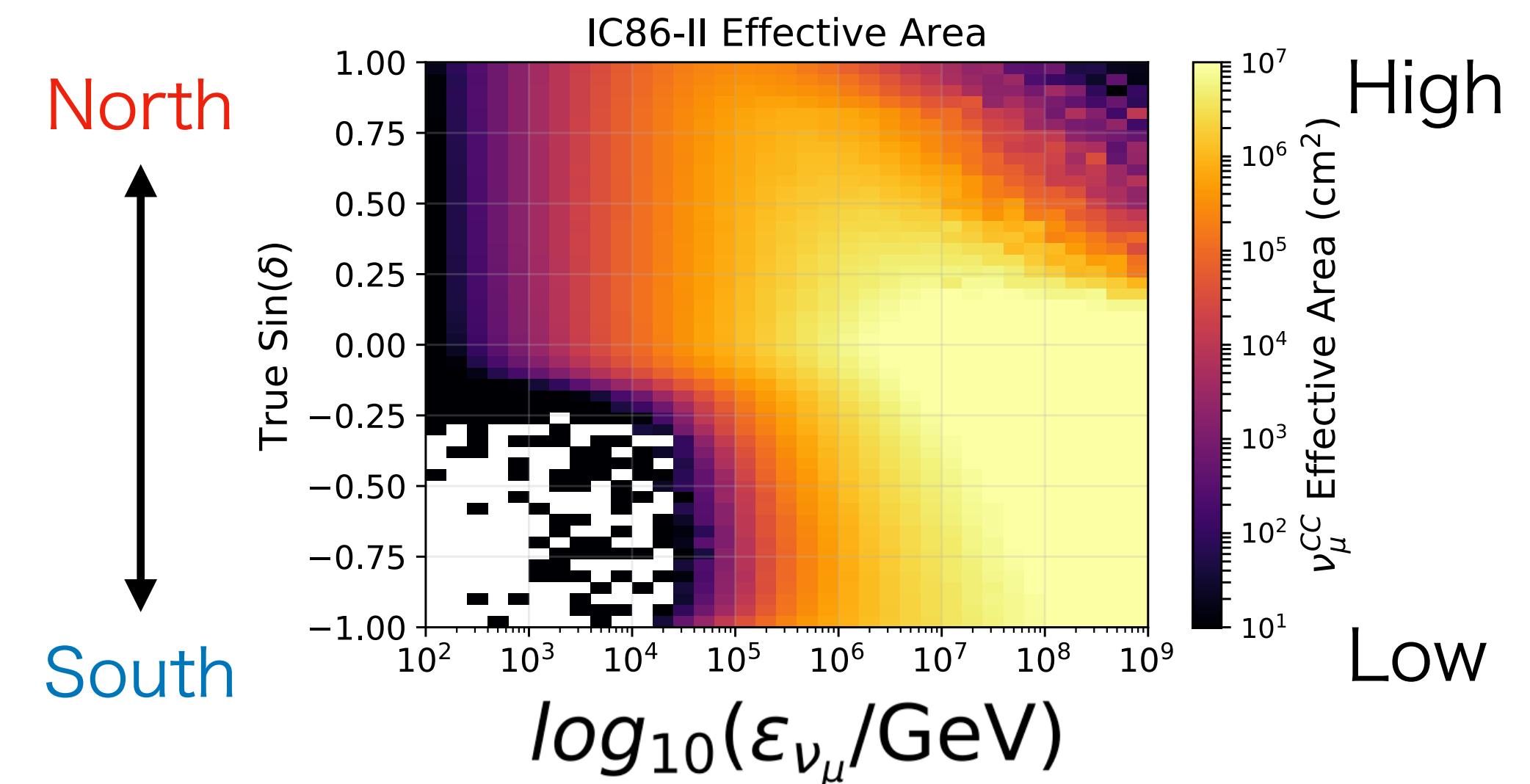
Cocoon increase the ν flux by 1000 times at ~PeV.

Results 2 : Expected Number of Detection

Neutrino Spectra



IceCube Effective Area



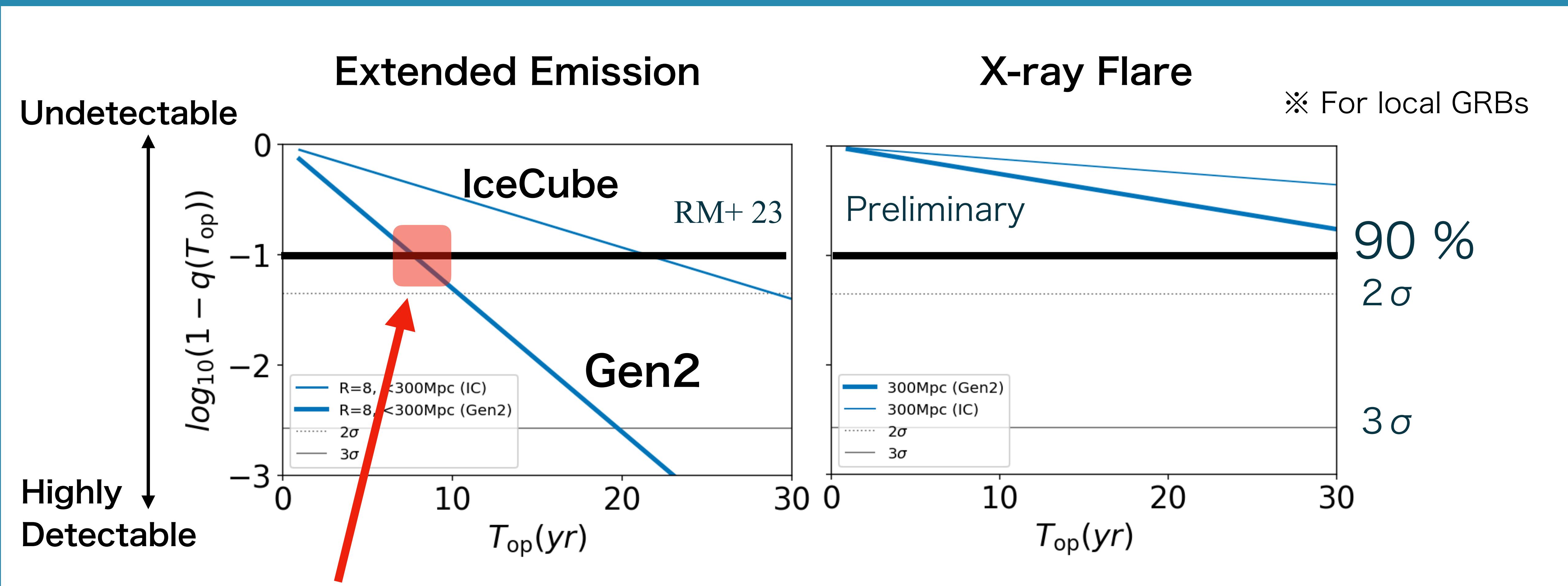
$$\bar{N}_{\nu_\mu} = \int d\epsilon_{\nu_\mu} \phi_{\nu_\mu + \bar{\nu}_\mu}(\epsilon_{\nu_\mu}) A_{\text{eff}}(\delta, \epsilon_{\nu_\mu})$$

Number of the Detection

1 ν / 10 GRBs by IceCube \otimes Fixed $d_L = 300 \text{ Mpc}$, $t_{\text{dur}} \sim 300 \text{ s}$



Results 3 : Operation time



Neutrinos are 90% detectable in 10 years (Gen2, E.E of local sGRB).

※ X-ray flares take more than 30 years.

Model Parameters

Table 1. fiducial parameters

Parameters	Γ_j	t_{dur} (s)	L_X (erg/s)	r_{dis} (cm)	$\varepsilon_{\gamma,\text{pk}}$ (keV)
Extended	200	$10^{2.5}$	10^{48}	10^{12}	10
Plateau	100	10^4	10^{46}	10^{13}	1
either case	p_{inj}	ξ_p	ξ_B	d_L (Mpc)	Energy band (keV)
	2.0	10	0.33	300	0.3 – 10 (XRT)

Proton number

γ number density

$$\theta_j = 5^\circ$$

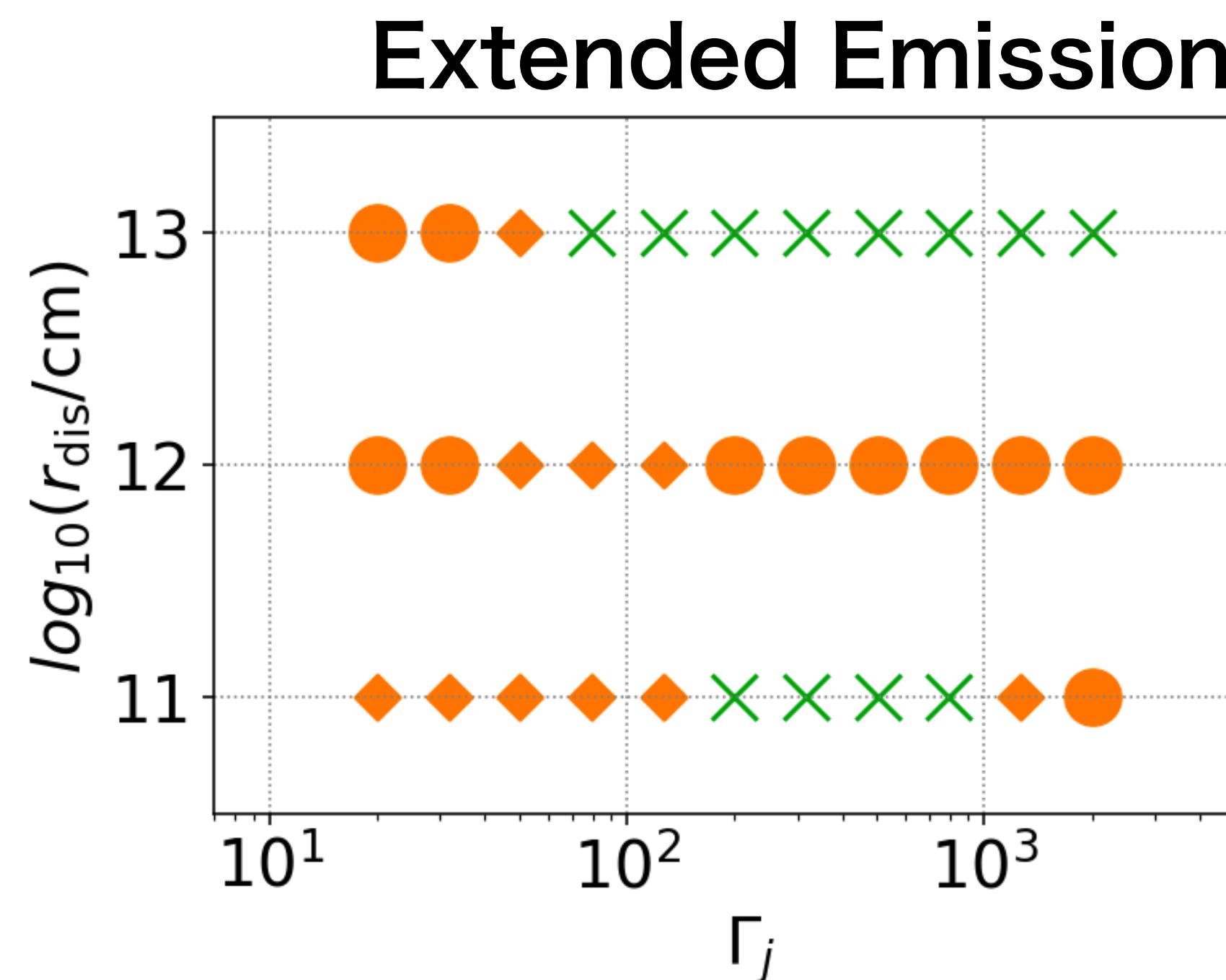
$$\xi_p = U_p/U_\gamma,$$

$$\xi_B = U_B/U_\gamma$$

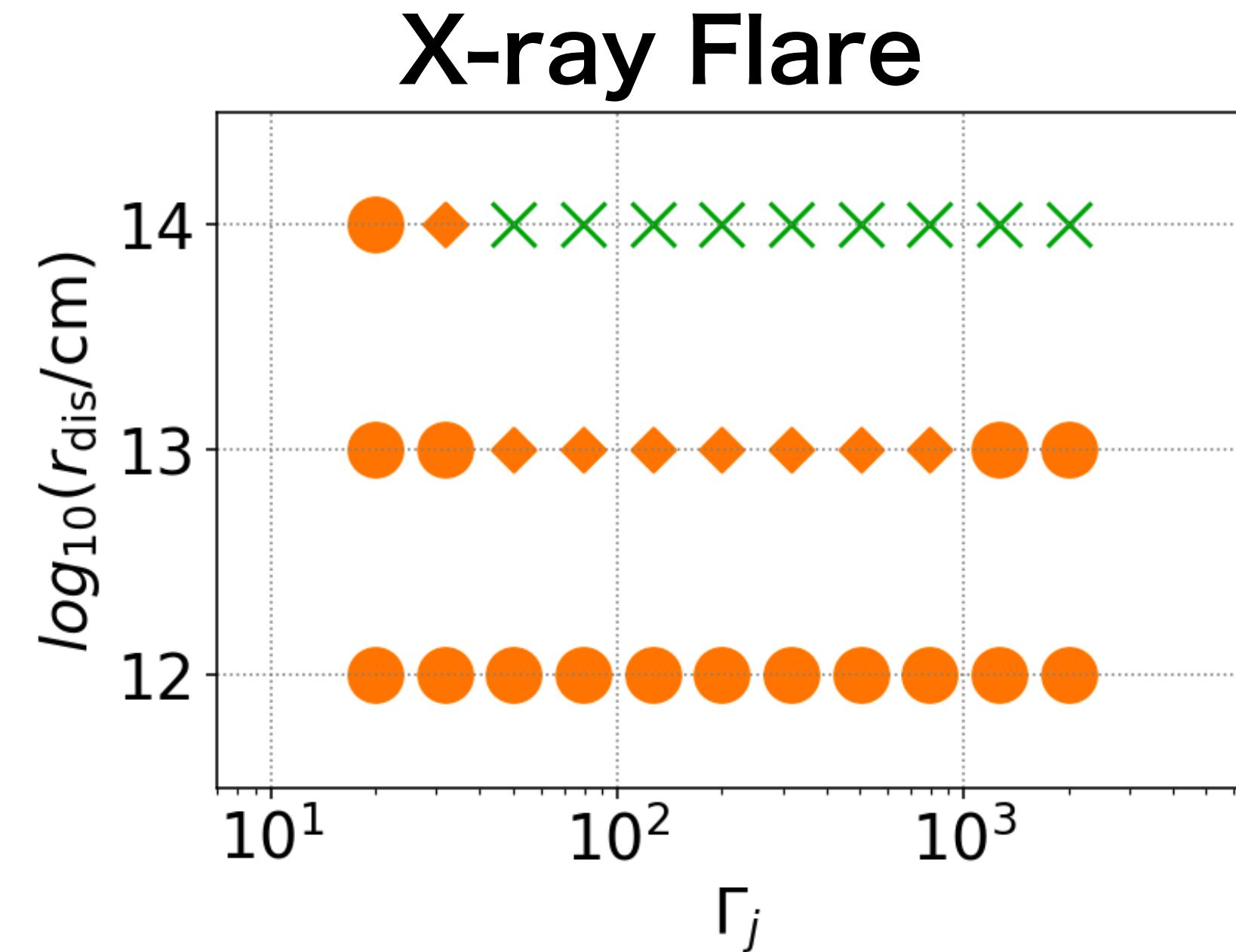
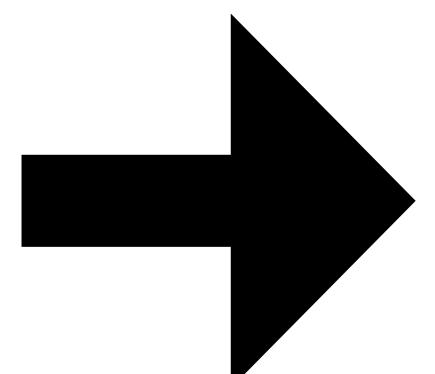
$$\epsilon_p = U_p/U_k \sim 0.3,$$

$$\epsilon_B = U_B/U_k \sim 0.01$$

Discussion : Constraints on physics of sGRB



>> 1 event/10 yr
or
<< 1 event/10yr



Under Debate

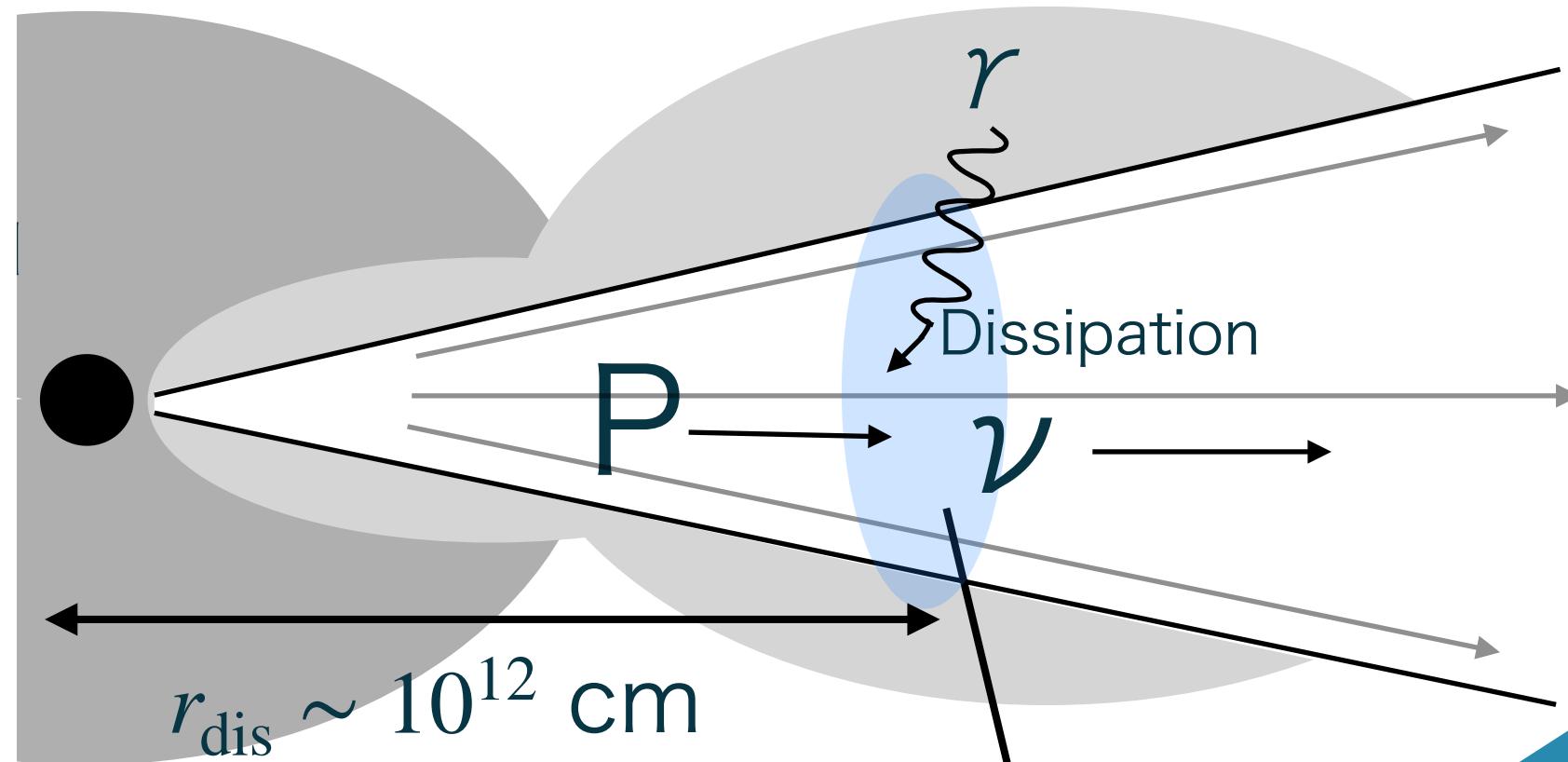
- How jets launched ?
- How jets dissipated ?
- How many CRs ?

Summary

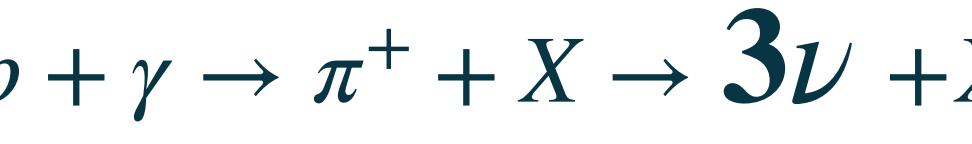
- We calculated the neutrino emission associated with the late-time emission of Gamma-Ray Bursts, considering the prolonged engine activity model and cocoon photons.
- 1 neutrino will be detected with 10 local GRBs ($d_L \sim 300$ Mpc)
- Neutrinos should be detected from GRBs in 10 years.
(Gen2 : 90%)
- The dissipation radius and the Lorentz factor of late-time jet will be constrained by the high energy observation in the future.

Back Up

Method 1 : Neutrino emission



$t'_{p\gamma}$



+ **Adiabatic** cooling

+ Bethe Heitler process

+ Synchrotron cooling

Neutrino Spectrum (One-Zone approx.)

$$\varepsilon_{\nu_\mu}^2 \frac{dN_{\nu_\mu}}{d\varepsilon_{\nu_\mu}} \approx \frac{1}{8} f_{p\gamma} \varepsilon_p^2 \frac{dN_p}{d\varepsilon_p}, \quad \varepsilon_p = \varepsilon_{\nu_\mu}/0.05$$

where $f_{p\gamma} = t'_{\text{cool}}/t'_{p\gamma}$,

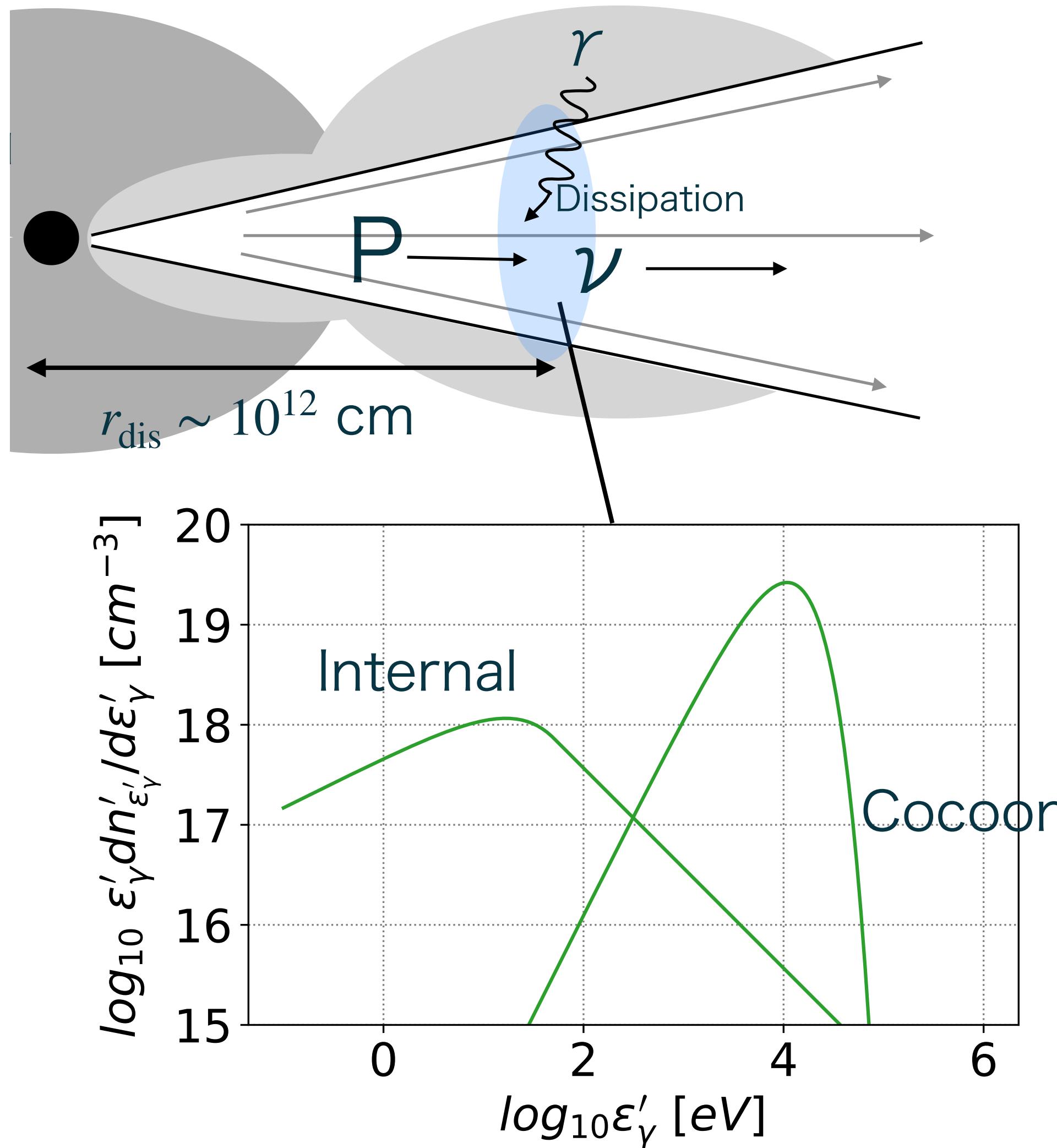
$$t'^{-1}_{p\gamma} = \left(\frac{c}{2\gamma_p'^2} \int_{\bar{\varepsilon}_{th}}^{\infty} d\bar{\varepsilon}_\gamma \sigma_{p\gamma} \kappa_{p\gamma} \bar{\varepsilon}_\gamma \int_{\bar{\varepsilon}_\gamma/2\gamma_p}^{\infty} d\varepsilon'_\gamma \varepsilon'^{-2} \frac{dn'_\gamma}{d\varepsilon'_\gamma} \right) \sim n_\gamma \sigma c$$

+ diffusive acceleration

t'_{acc}

Spectra of Photon and Proton are required.

Method 2 : Photons in Dissipation Region



Photon Spectrum

- Internal photons (Kimura +17)

$$\frac{d n'^{\text{in}}}{d \varepsilon'_\gamma} \propto \varepsilon'_\gamma^{-0.5} \quad (\varepsilon'_\gamma \geq \varepsilon'_{\gamma, \text{pk}})$$

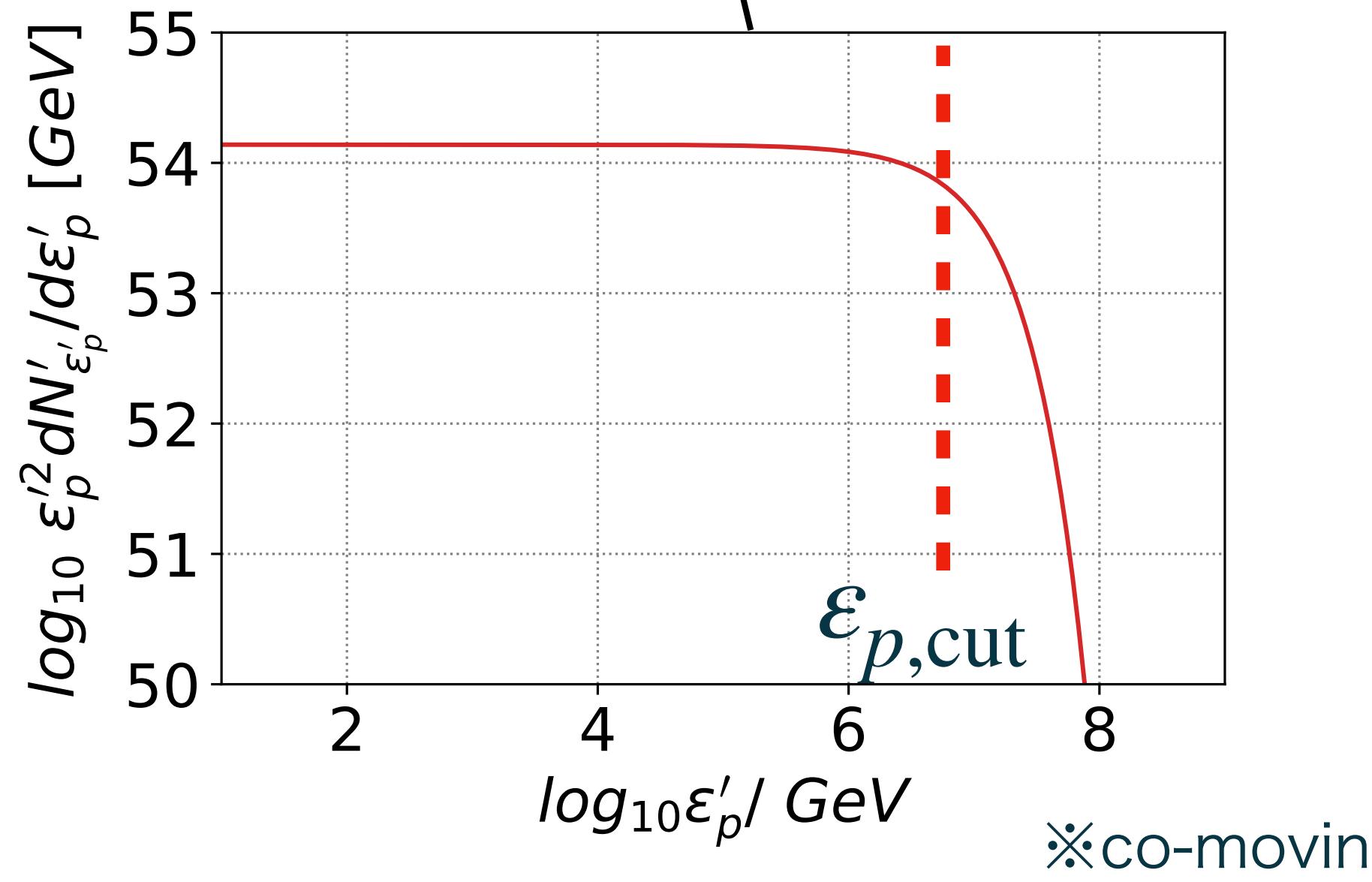
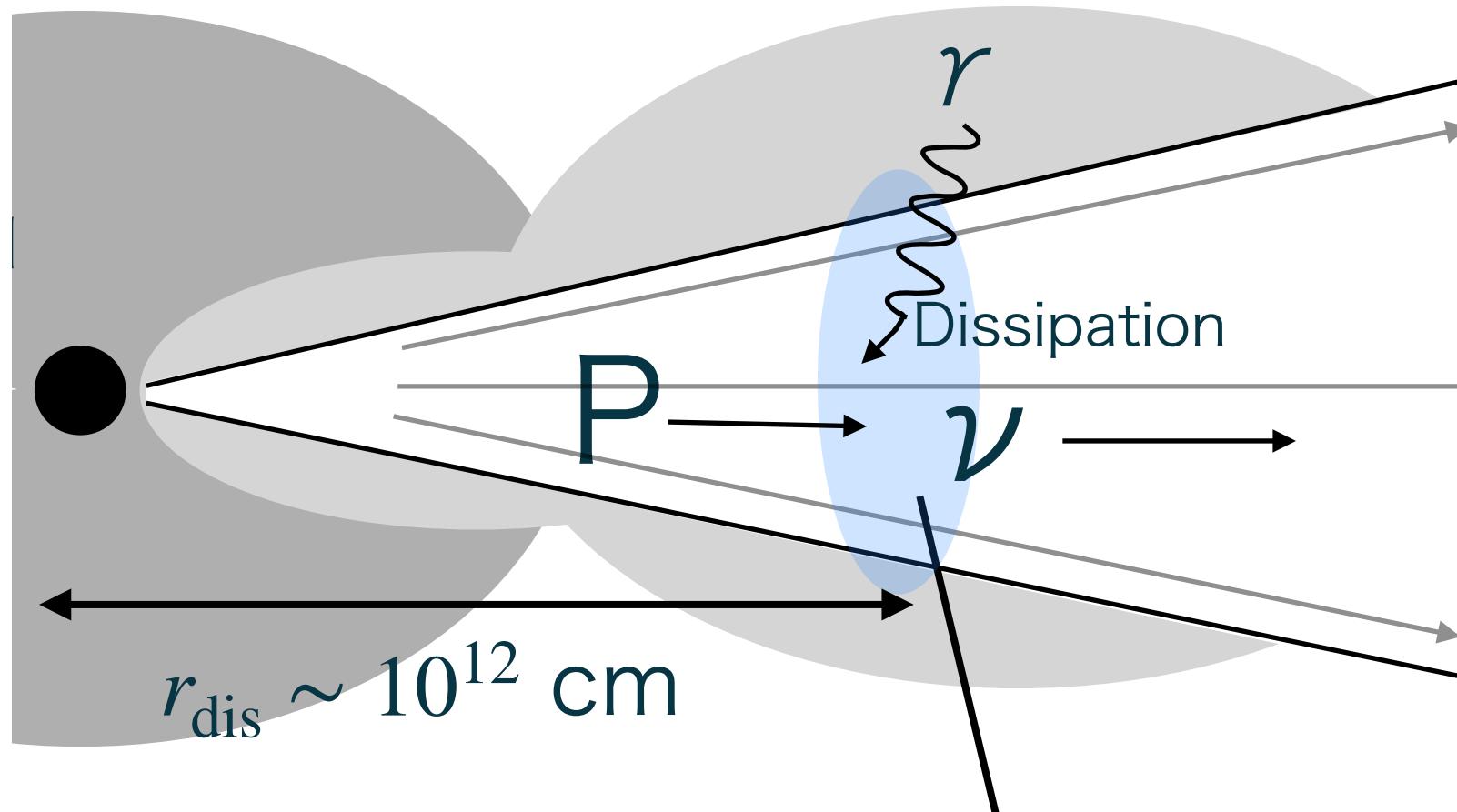
$$\propto \varepsilon'_\gamma^{-2} \quad (\varepsilon'_\gamma < \varepsilon'_{\gamma, \text{pk}})$$

- Cocoon photons (\leftarrow New)

$$\varepsilon'_\gamma \frac{d n'^{\text{ex}}}{d \varepsilon'_\gamma} = \Gamma_j \frac{8\pi(\varepsilon'_\gamma/\Gamma_j)^3}{h^3 c^3} \frac{1}{\exp(\varepsilon'_\gamma/\Gamma_j k_B T_{\text{coc}}) - 1}$$

$$k_B T_{\text{coc}} \sim 20 \text{ eV}$$

Method 3 : Protons in Dissipation Region



Proton (Cosmic-Ray) Spectrum

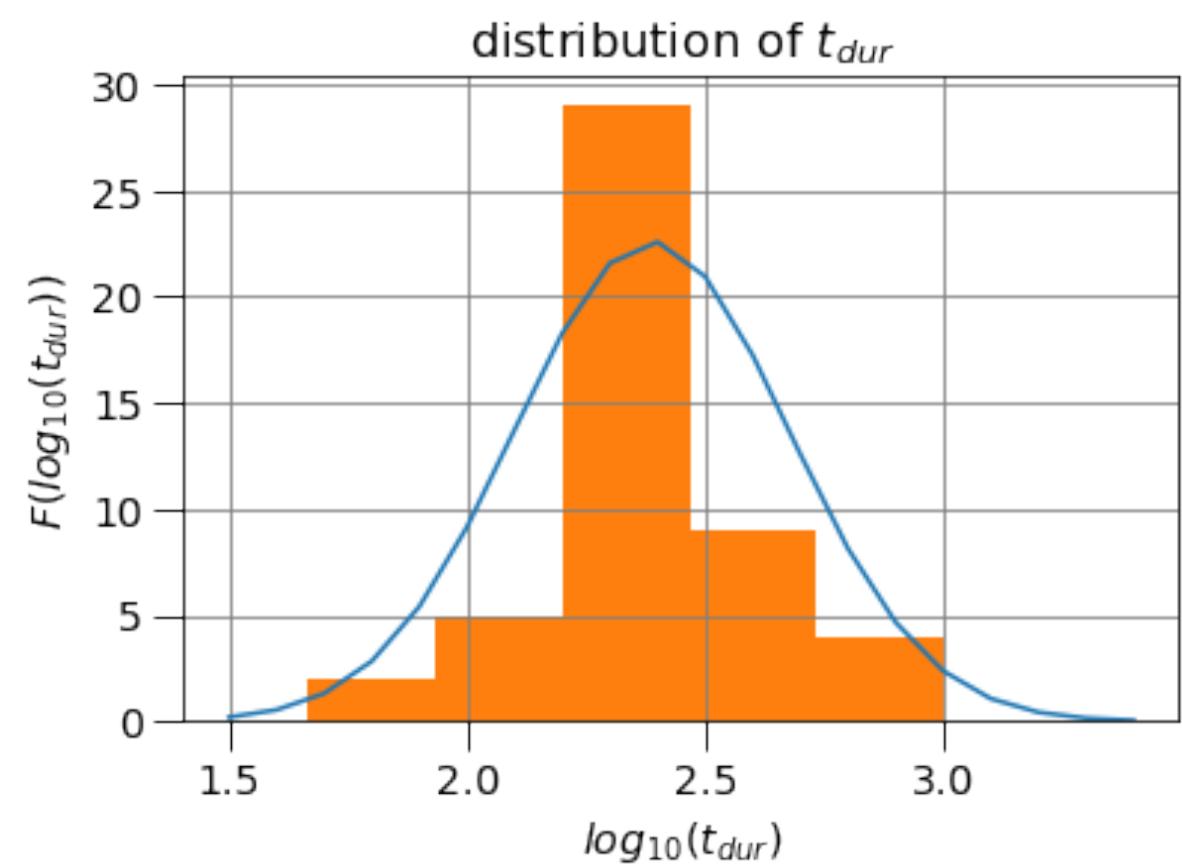
$$\frac{dN_p}{d\varepsilon_p} = N_{\varepsilon_p,\text{nor}} \left(\frac{\varepsilon_p}{\varepsilon_{p,\text{cut}}} \right)^{-2} \exp \left(-\frac{\varepsilon_p}{\varepsilon_{p,\text{cut}}} \right),$$

(assuming Fermi acceleration)

where $t_{\text{acc}}'^{-1}(\varepsilon_p') = t_{\text{cool}}'^{-1}(\varepsilon_p') \rightarrow \varepsilon_{p,\text{cut}}$

Empirical relation for Extended emission

- Duration, $\log t_{\text{dur}}$: Gaussian by Observational data
- $L_X \propto t_{\text{dur}}^{-2.5}$ (Empirical Low , Kisaka +17)
- Location : homogenous, $d_L < 300\text{Mpc}$
- Event rate : $R_{\text{sGRB}} \sim 8 \text{ Gpc}^{-3}\text{yr}^{-1}$



(Kisaka +17)

Flare cocoon model

$$\text{Internal energy density : } aT_{\text{coc}}^4 = \frac{E_{i,\text{coc,br}}}{4\pi r_{\text{diss}}^3/3} \frac{r_*}{ct_{\text{flare}}}$$

$$\text{Internal E. at break out : } E_{i,\text{coc,br}} = L_j t_b \times (1 - \beta_h) \quad (\text{Hamidani \& Ioka 2022})$$

$$\text{Jet (kinetic) luminosity : } L_j = L_{k,\text{iso}} \theta_j^2 / 4$$

$$\text{Jet isotropic-equivalent luminosity : } L_{k,\text{iso}} = 30 L_{\text{prompt,iso}}$$

$$\text{break out time : } t_b = 0.9 (r_*^2 \theta_j^4 M_*/L_j)^{1/3} \quad (\text{Hamidani \& Ioka 2022})$$

$$\text{Jet velocity before break out : } \beta_h = r_*/ct_b$$

where Prompt emission Luminosity : $L_{\text{prompt,iso}}$, Dissipation radis : r_{diss} , Jet opening angle θ_j ,
Steller mass, radius : M_* , r_* ,

Empirical relation and distribution for X-ray flare

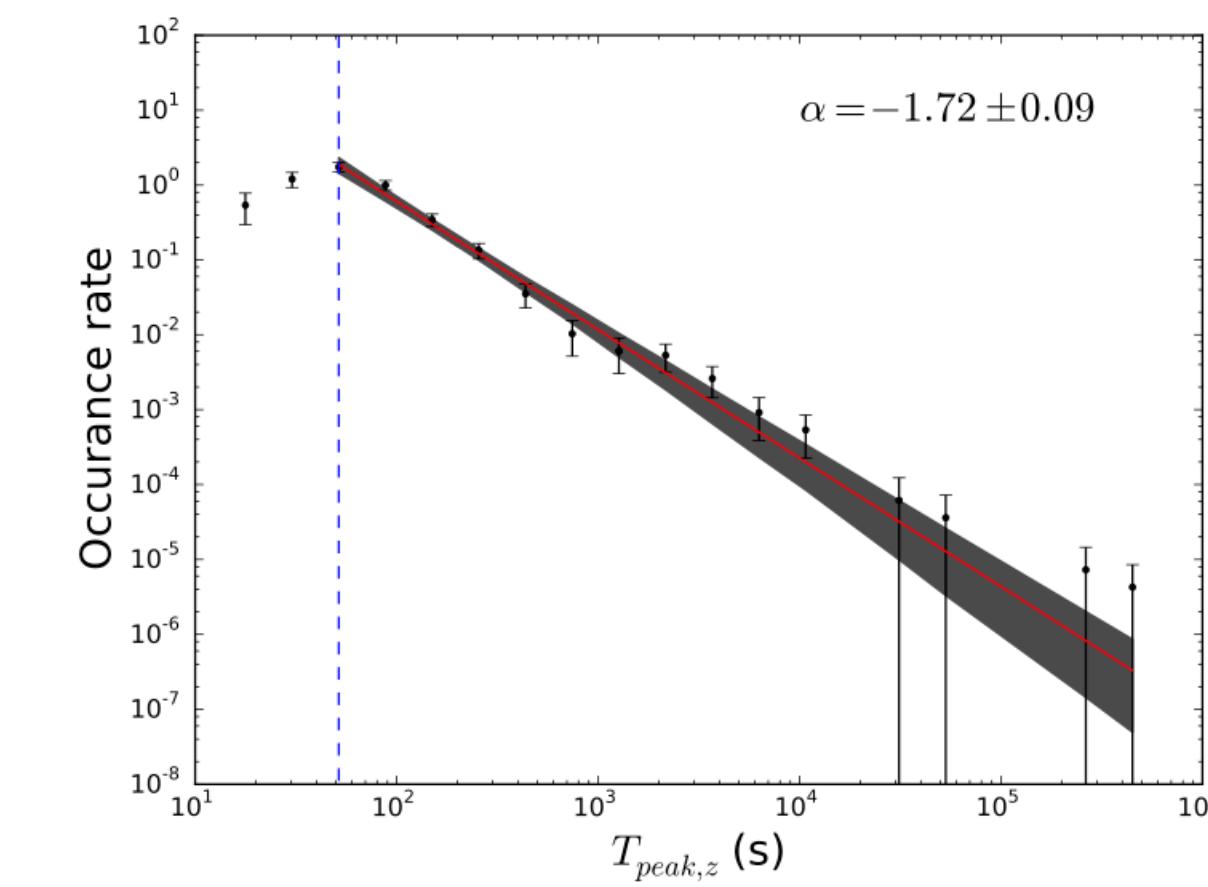
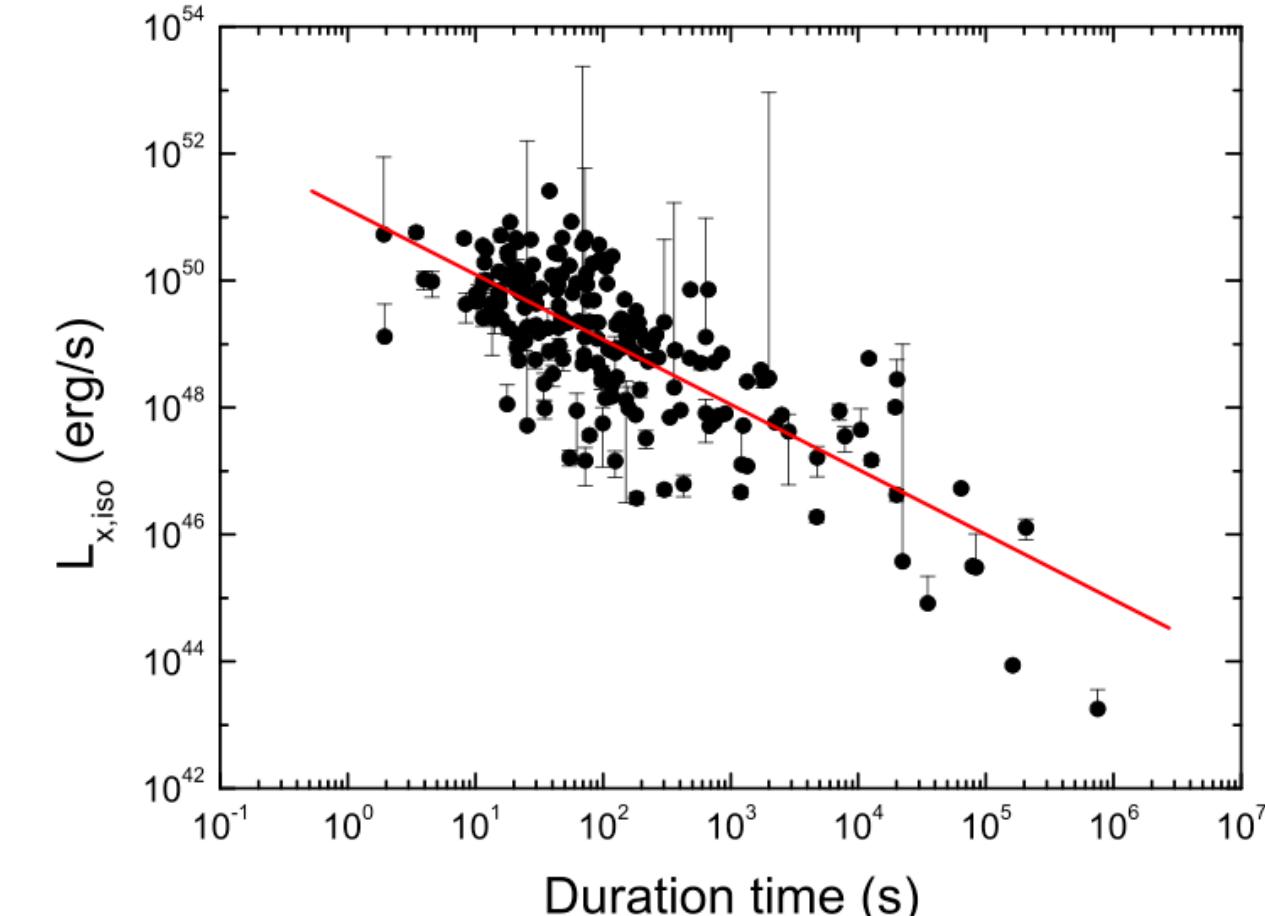
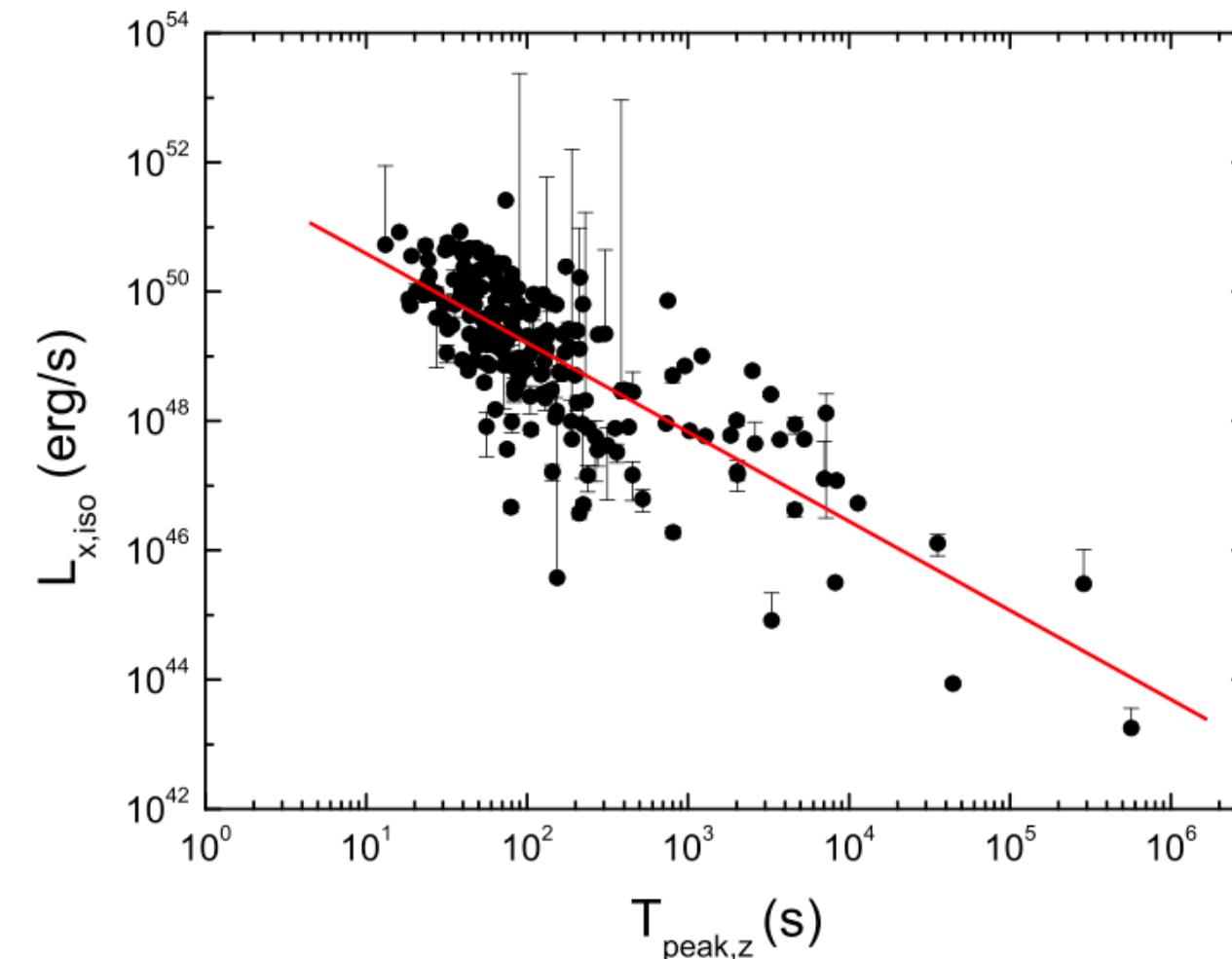
Prompt emission (radiation) luminosity : $L_{\text{prompt,iso}} \sim 30 L_{x,\text{flare}}$ (Liu & Mao 19)

X-ray flare luminosity : $L_{x,\text{iso}} = 1.3 \times 10^{48} \times (t_{\text{flare},3} (1+z))^{-1.02} \text{erg/s}$ (Yi+16)

Duration of X-ray flare : $t_{\text{dur}} = 450 \times (t_{\text{flare},3} (1+z))^{1.12} \text{s}$ (Yi+16)

Distribution of t_{flare} : $F(t_{\text{flare}}) dt_{\text{flare}} \propto t_{\text{flare}}^{-1.72} dt_{\text{flare}}$ (Yi+16)

where X-ray flare peak time : t_{flare} , $t_{\text{flare},3} = t_{\text{flare}}/(1000 \text{ s})$



Yi+16