

Axion Stars and How to Find Them

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Energy conservation and axion back-reaction in an external magnetic field

Srimoyee Sen and Lars Sivertsen, JHEP 03 (2023) 097, arxiv:
2210.01149

Srimoyee Sen and Lars Sivertsen, JHEP 05 (2022) 192, arxiv:
2111.08728

Outline of the Talk

- ▶ Axions and axion-like particles
- ▶ Axion Condensates (Axion Stars)
- ▶ The modified Maxwell's Equations
- ▶ Electromagnetic Radiation from Axion Condensates
- ▶ Decay of axion stars
- ▶ Take Home Message

- ▶ The full QCD Lagrangian has a CP violating term leading to a non-zero magnetic dipole moment for the neutron
- ▶ Axions suggested as a solution to the strong CP problem
- ▶ To solve the strong CP problem, need $m_a f_a \approx m_\pi f_\pi \approx \Lambda_{\text{QCD}}^2$
- ▶ Axions turns out to be extremely light (order eV or smaller), and interact very weakly with ordinary matter
- ▶ Excellent candidate for dark matter as well!

Axions and Axion-like Particles

- ▶ Particles with axion like properties are predicted by string theory
- ▶ Multiple different axions spanning many orders of magnitude
- ▶ f_a determined by string compactifications
- ▶ This motivates the search for axion-like particles with $m_a f_a \neq \Lambda_{\text{QCD}}^2$

Axion Condensates

- ▶ Axions can form spherically symmetric, coherently oscillating lumps of Bose-Einstein condensates (axion stars)

$$\phi(\mathbf{x}, t) \approx \phi_0 \operatorname{sech}(r/R) \cos(\omega t).$$

- ▶ Can be dense ($m_a R \sim 1$)

- Dominated by self interactions

- $\omega \lesssim m_a$

- $\phi(\mathbf{x} = 0, t) \sim f_a$

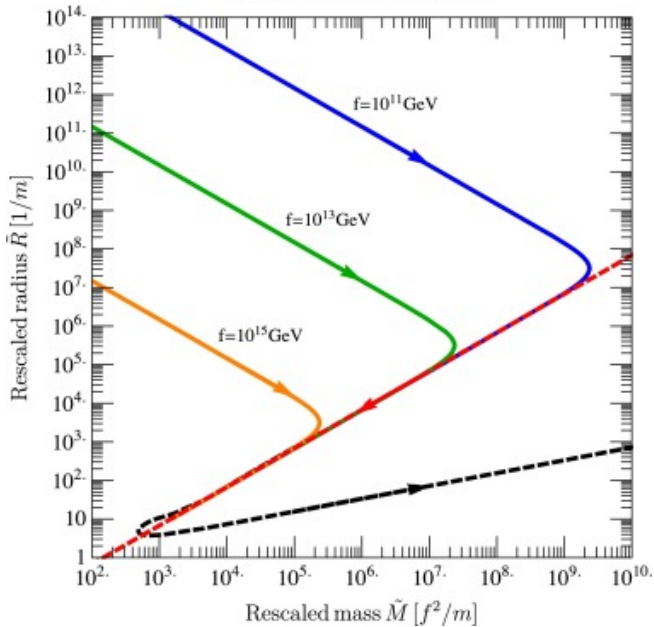
- ▶ Or dilute ($m_a R \gg 1$)

- Dominated by gravitational interactions

- $\omega \approx m_a$

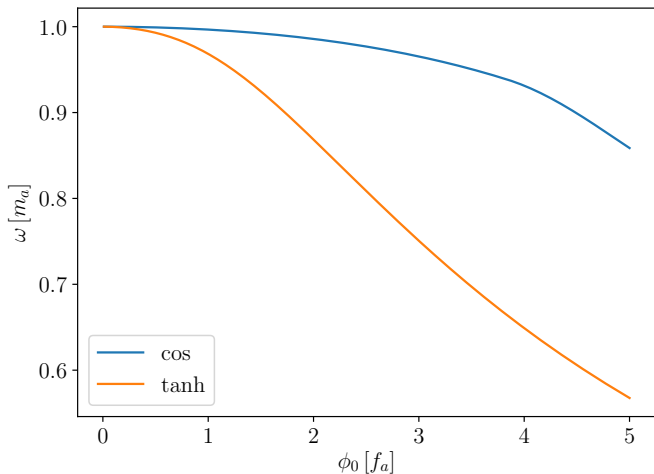
- $\phi(\mathbf{x} = 0, t) \ll f_a$

Axion star radius vs mass



Axion Condensates

- ▶ The frequency of the condensate is dependent on the central amplitude



Electromagnetic Radiation from Axion Stars in an External Magnetic Field

► Axion-photon Lagrangian

$$\mathcal{L} = -\frac{1}{4}F^{\mu\nu}F_{\mu\nu} + J_m^\mu A_\mu + \frac{C\beta}{4\pi f_a}\phi\epsilon^{\mu\nu\lambda\rho}F_{\mu\nu}F_{\lambda\rho} + \frac{1}{2}(\partial_\mu\phi)(\partial^\mu\phi) - V(\phi) \quad (1)$$

► To first order in $\frac{C\beta}{\pi f_a}\phi$ ($\phi \lesssim f_a$)

$$\square \mathbf{A}_r(\mathbf{x}, t) = -\frac{C\beta}{\pi f_a}(\partial_t\phi(\mathbf{x}, t))\mathbf{B}_0 \equiv \mathbf{J}_a(\mathbf{x}, t) \quad (2)$$

$$\square \Phi_r(\mathbf{x}, t) = \frac{C\beta}{\pi f_a}\nabla\phi(\mathbf{x}, t) \cdot \mathbf{B}_0 \equiv \rho_a(\mathbf{x}, t) \quad (3)$$

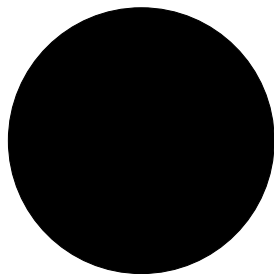
$$(\square + m_a^2)\phi(\mathbf{x}, t) + \partial_\phi V(\phi) = -\frac{C\beta}{\pi f_a}\mathbf{E}_r(\mathbf{x}, t) \cdot \mathbf{B}_0(\mathbf{x}, t) \quad (4)$$

Electromagnetic Radiation from Axion Stars in an External Magnetic Field

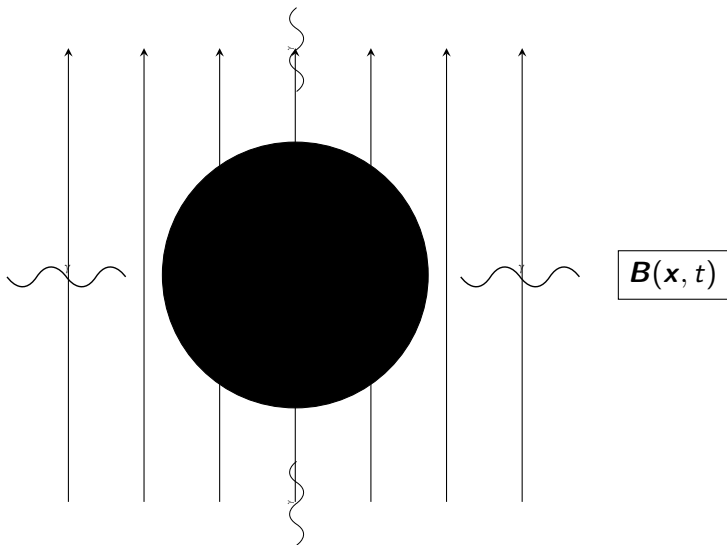
- ▶ Axions interact extremely weakly with matter
- ▶ However, axion stars have high occupation numbers, resonance effects can occur
- ▶ Axion stars in the presence of a (strong) magnetic field will radiate, slowly changing the condensate frequency.
- ▶ Presence of a plasma can enhance radiation

Electromagnetic Radiation from Axion Condensates in External Magnetic Field

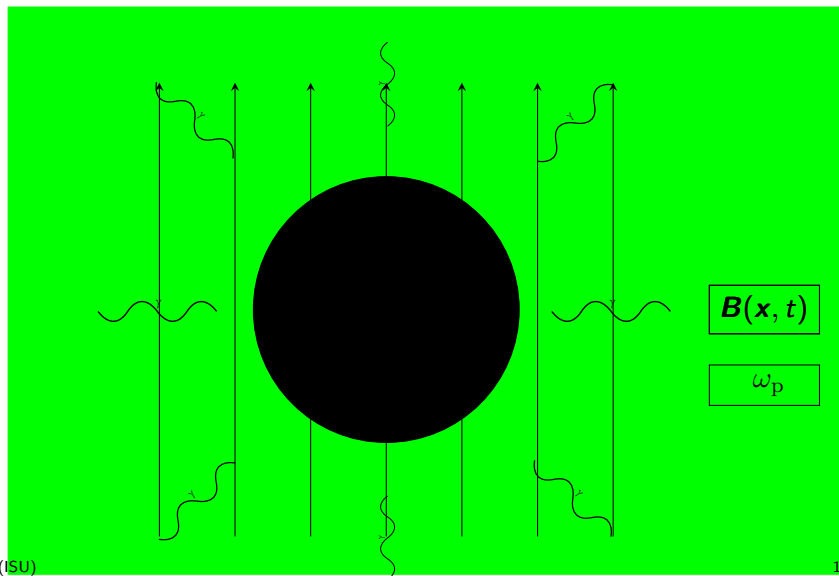
Electromagnetic Radiation from Axion Condensates in External Magnetic Field



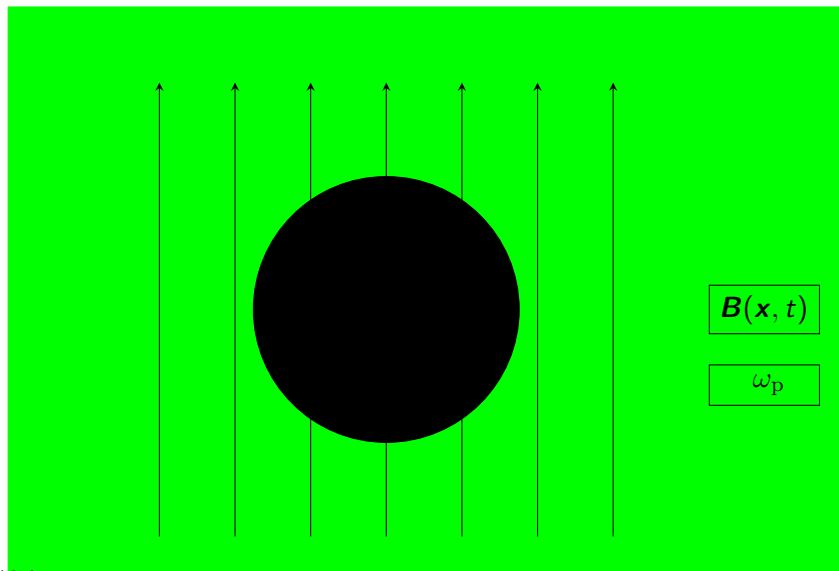
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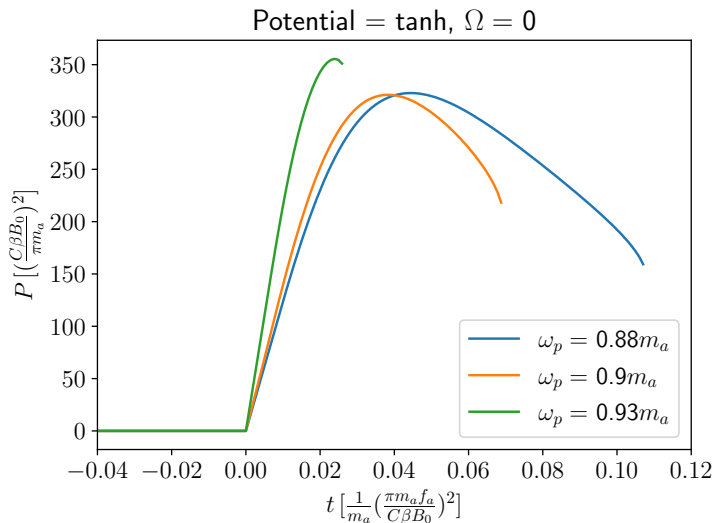
Electromagnetic Radiation from Axion Condensates in External Magnetic Field



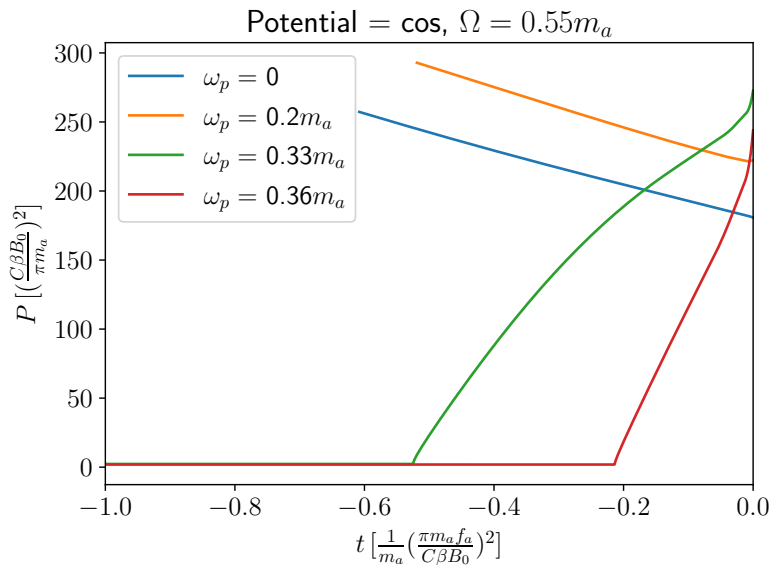
Electromagnetic Radiation from Axion Condensates in External Magnetic Field



Decay of axion stars



Decay of Axion Stars



Take Away Message

- ▶ Axions can form spherically symmetric axion condensates, axion stars
- ▶ Axion BEC start to give off electromagnetic radiation when subject to an external magnetic field
- ▶ For dense condensates, the frequency becomes time dependent when back-reaction is included
- ▶ Resonances can occur as time passes

Thanks for listening!

The next slides are extra

► Modified Maxwell equations

$$\begin{aligned}\nabla \times \mathbf{B}(\mathbf{x}, t) - \partial_t \mathbf{E}(\mathbf{x}, t) - \mathbf{J}_m(\mathbf{x}, t) \\ = -\frac{C\beta}{\pi f_a} \left[(\partial_t \phi(\mathbf{x}, t)) \mathbf{B}(\mathbf{x}, t) + \nabla \phi(\mathbf{x}, t) \times \mathbf{E}(\mathbf{x}, t) \right] \quad (5)\end{aligned}$$

$$\nabla \times \mathbf{E}(\mathbf{x}, t) = -\partial_t \mathbf{B}(\mathbf{x}, t) \quad (6)$$

$$\nabla \cdot \mathbf{E}(\mathbf{x}, t) = \rho_m(\mathbf{x}, t) + \frac{C\beta}{\pi f_a} \nabla \phi(\mathbf{x}, t) \cdot \mathbf{B}(\mathbf{x}, t) \quad (7)$$

$$\nabla \cdot \mathbf{B}(\mathbf{x}, t) = 0 \quad (8)$$

► Axion Equation of motion

$$(\square + m_a^2)\phi(\mathbf{x}, t) + \partial_\phi V(\phi) = -\frac{C\beta}{\pi f_a} \mathbf{E}(\mathbf{x}, t) \cdot \mathbf{B}(\mathbf{x}, t) \quad (9)$$

Estimation of Energy Radiated

- ▶ Energy stored in condensates $R \sim m_a^{-1}$

$$E_\phi \sim m_a^2 \phi_0^2 R^3 \sim m_a^2 f_a^2 R^3 \sim \frac{f_a^2}{m_a} \quad (10)$$

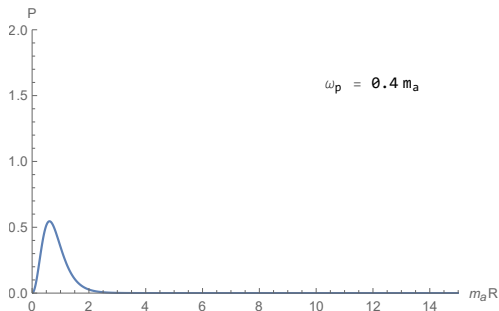
- ▶ A few orders of magnitude above or below a solar mass depending on the axion mass
- ▶ Energy stored in condensates $R \gg m_a^{-1}$

$$E_\phi \sim m_a^2 \phi_0^2 R^3 \sim \frac{m_{\text{P}}^2}{m_a} \frac{1}{(m_a R)} \quad (11)$$

Static External Magnetic Field

- ▶ Already found by Amin et al. to be ($k_\omega = \sqrt{\omega^2 - \omega_p^2}$)

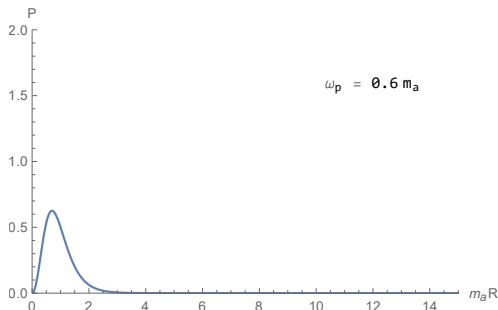
$$\langle P(t) \rangle_T = \left(\frac{C\beta}{\pi f_a} \right)^2 \left(\frac{\phi_0^2 B_0^2 \omega^3 R^4 \pi^5}{12 k_\omega} \right) \left(\frac{\tanh(\pi k_\omega R/2)}{\cosh(\pi k_\omega R/2)} \right)^2 \quad (12)$$



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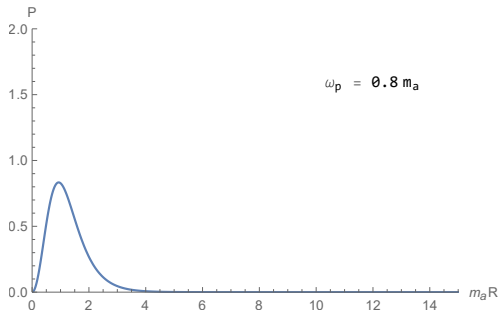
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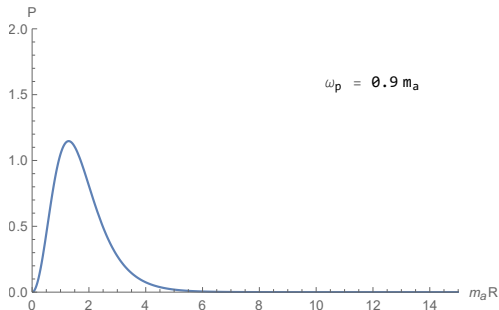
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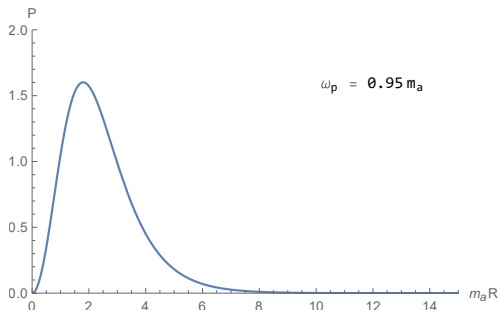
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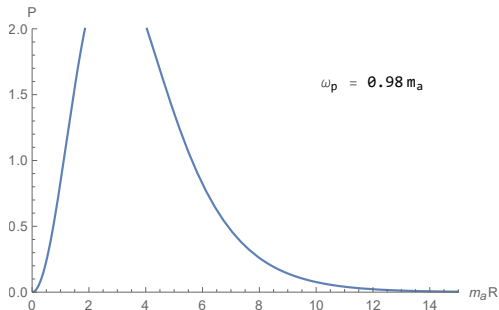
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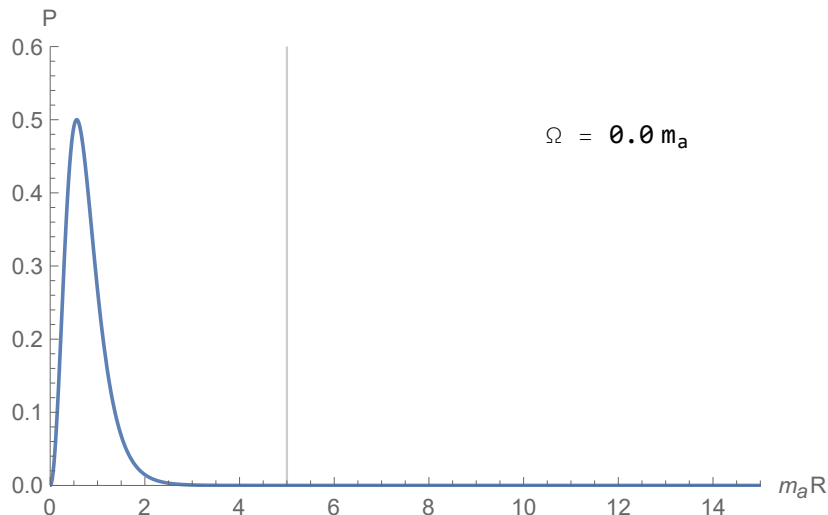


Static External Magnetic Field, takeaways

- ▶ Radiated power vanishes exponentially with system size ωR
- ▶ Tuning plasma frequency ω_p allows larger condensates to radiate efficiently
- ▶ Want to see if we can have similar effects by making the external magnetic field oscillate

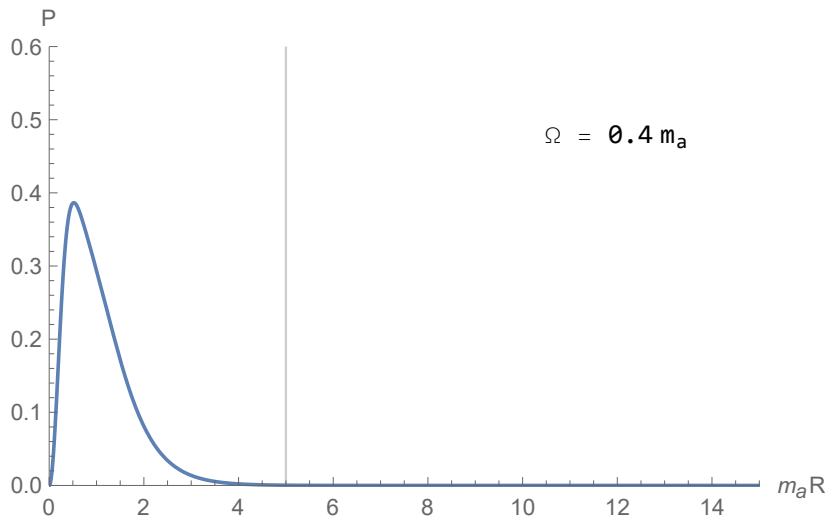
Oscillating External Magnetic Field

- ▶ When $\mathbf{B}_0 \rightarrow \mathbf{B}_0 \cos(\Omega t)$, the effective frequency splits in two $\omega \rightarrow \omega \pm \Omega$, wavenumber $k_{\pm} = |\Omega \pm \omega|$



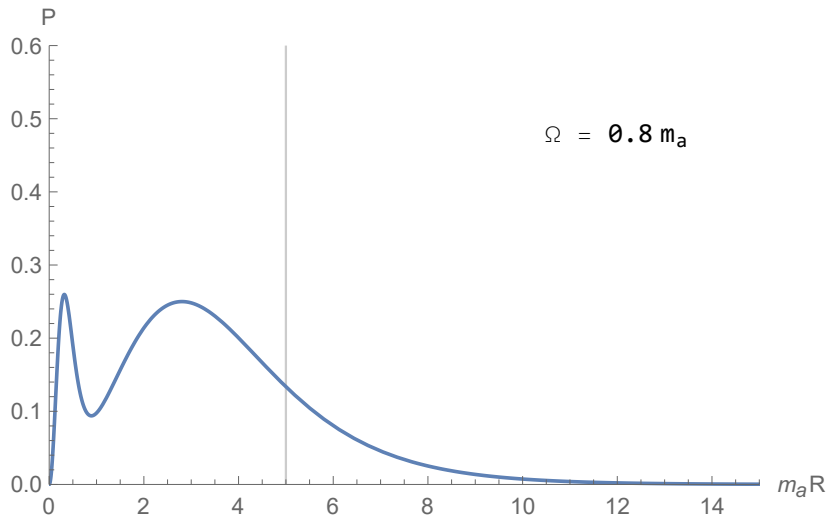
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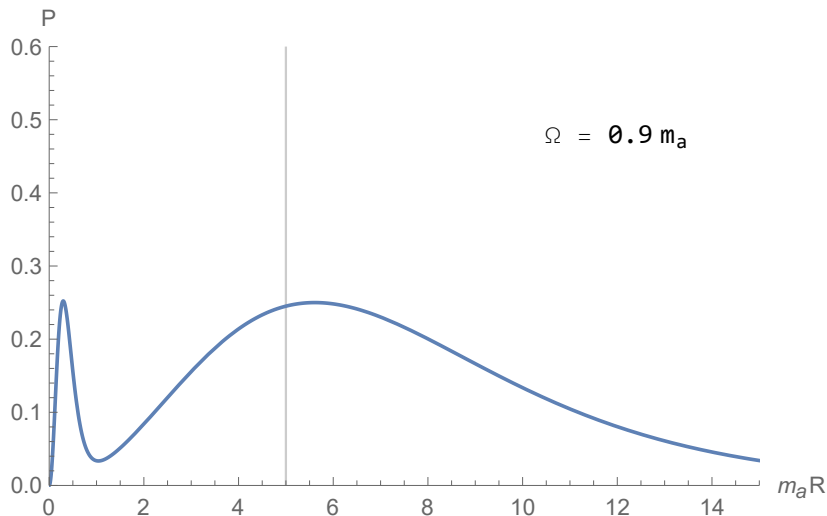
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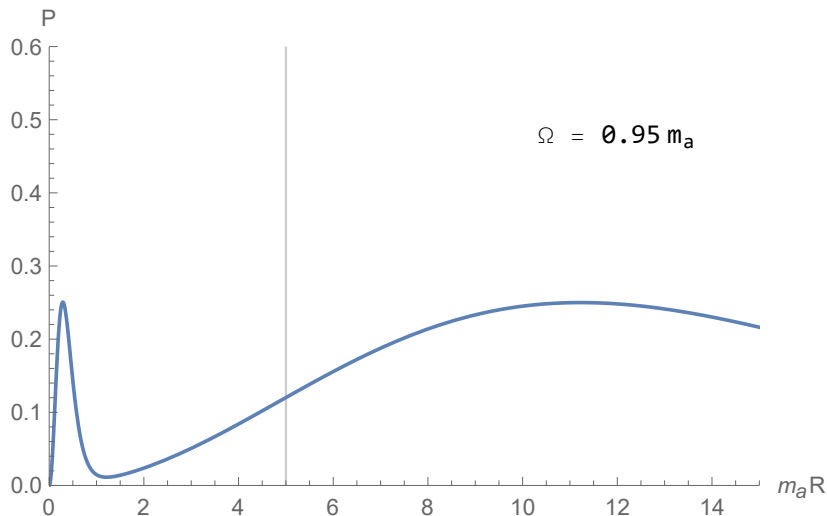
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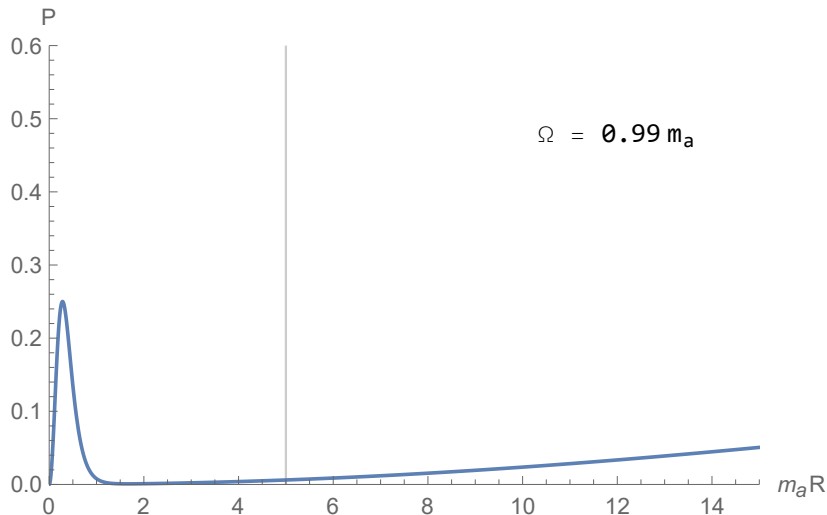
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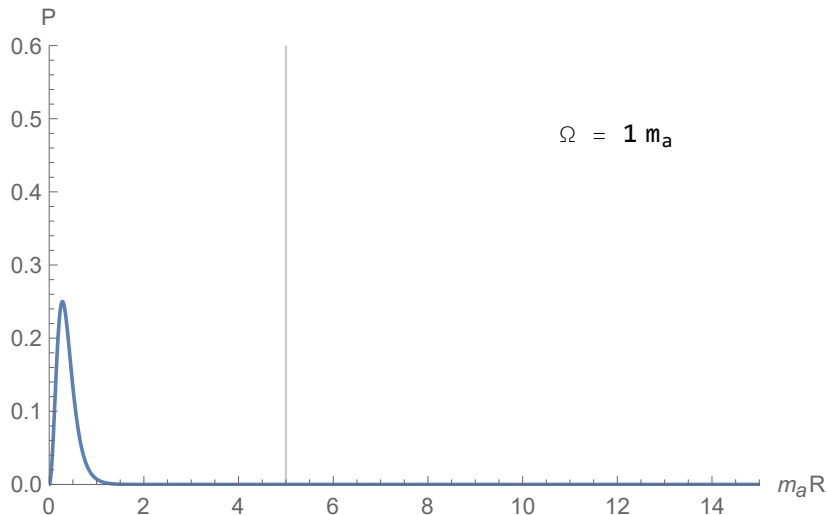
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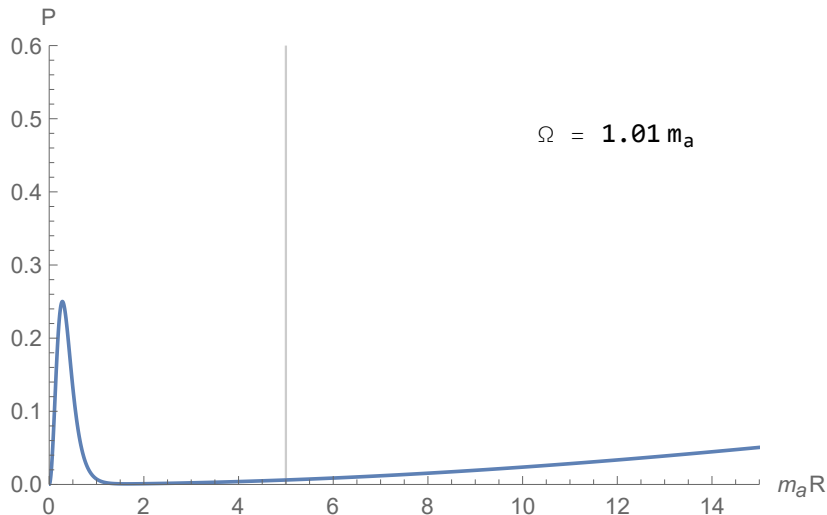
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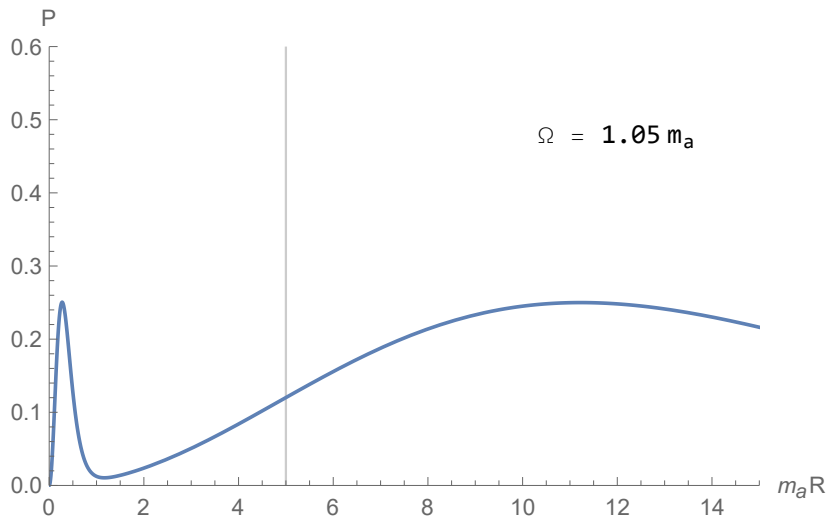
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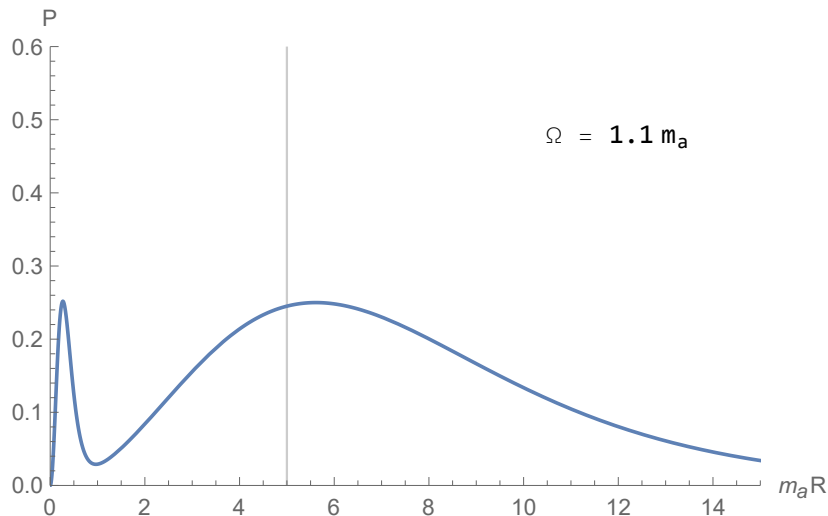
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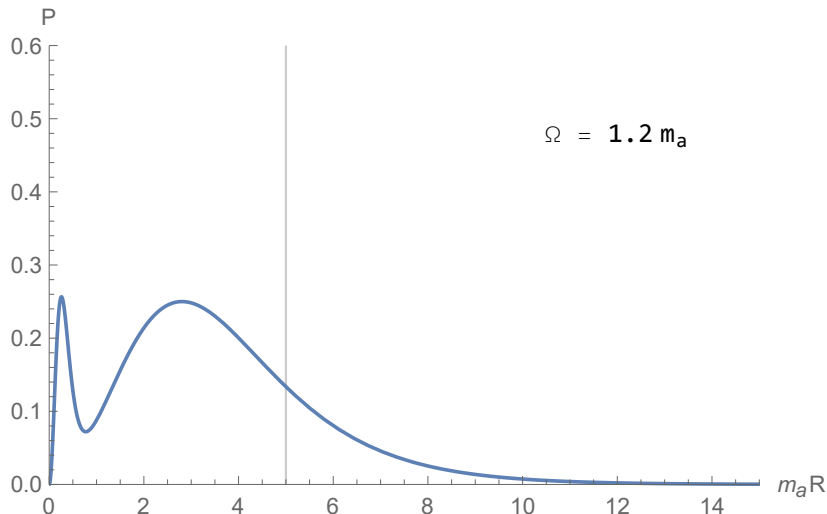
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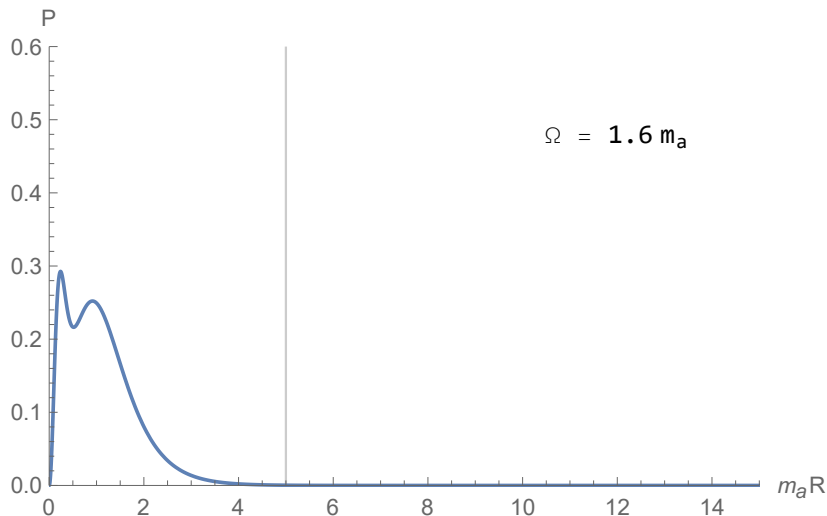
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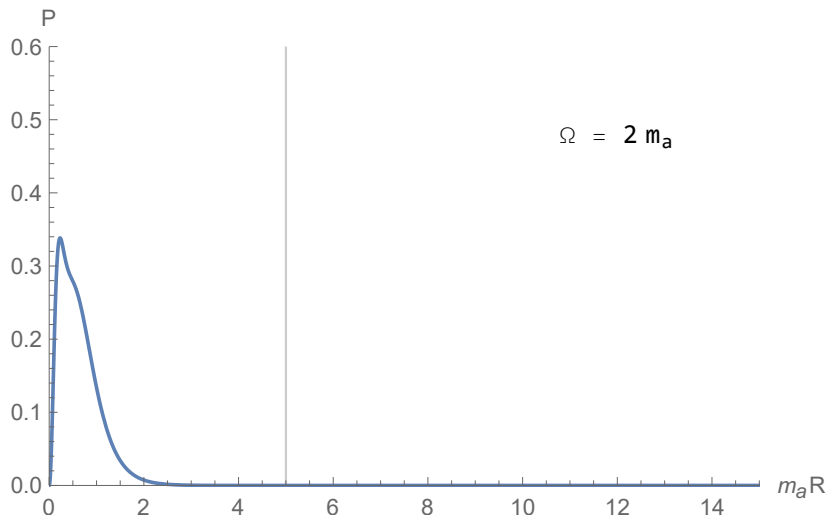
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