

# EVOLVED PULSAR WIND NEBULAE

BARBARA OLMI  
INAF - OAPa/OAA



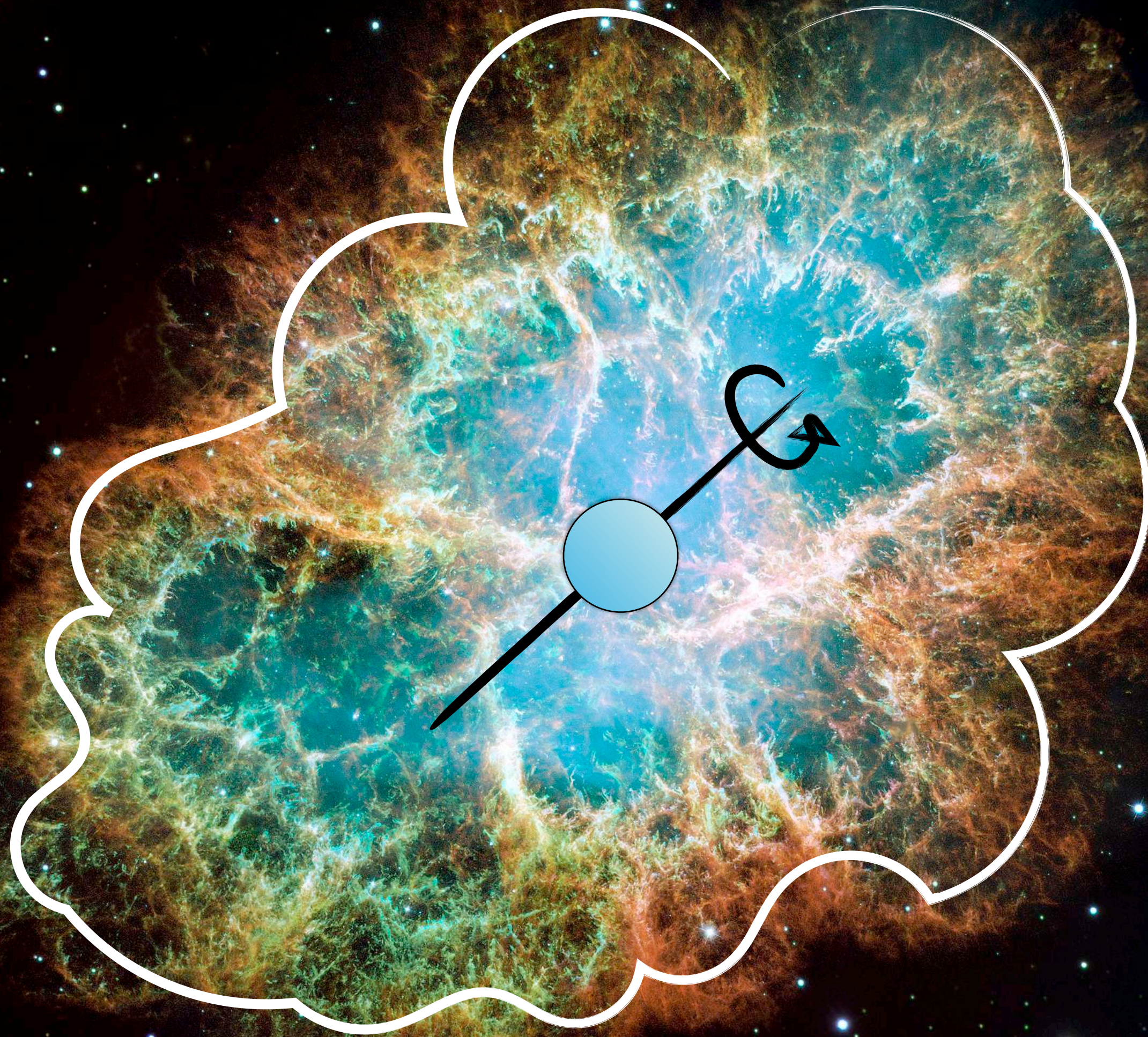
SEPTEMBER 11-15 2023  
NAPOLI



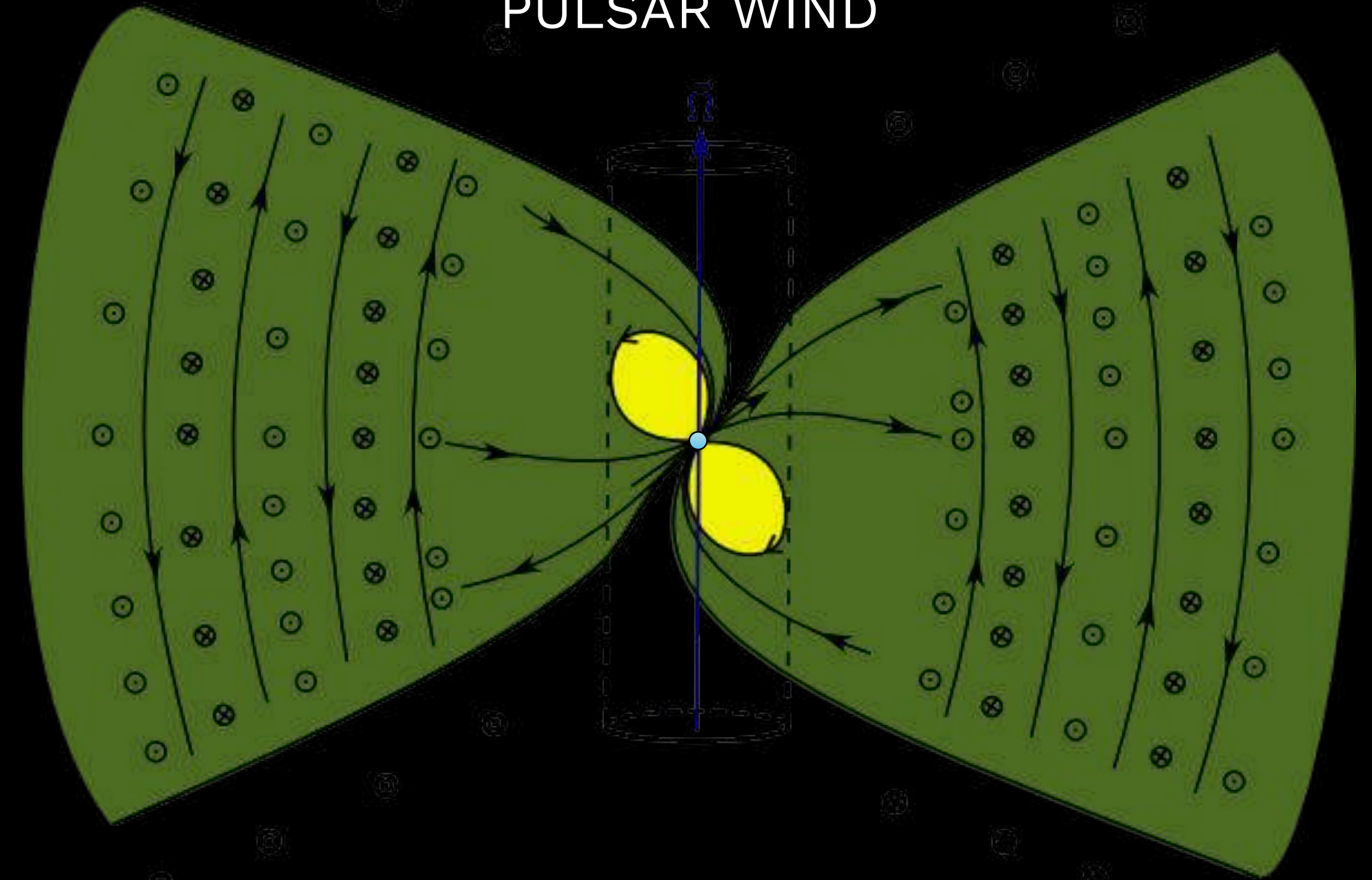
# INTRO: WHAT IS A PWN?

THE DEBRIS OF THE SUPERNOVA EXPLOSION OF A  
MASSIVE STAR ( $M \gtrsim 8 M_{\odot}$ )

MAIN INGREDIENTS:  
PULSAR  
PULSAR WIND



[NOT IN SCALE]

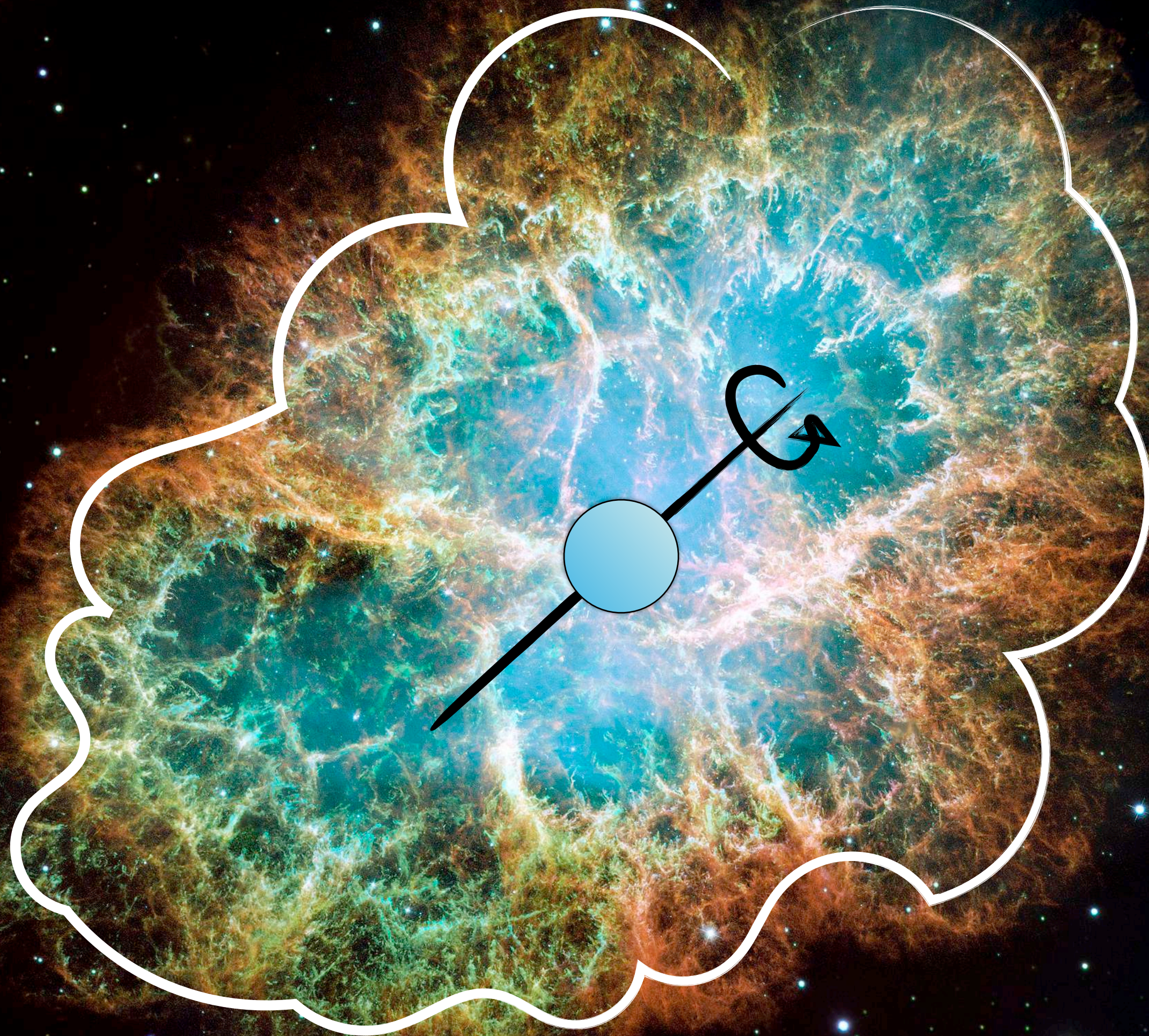




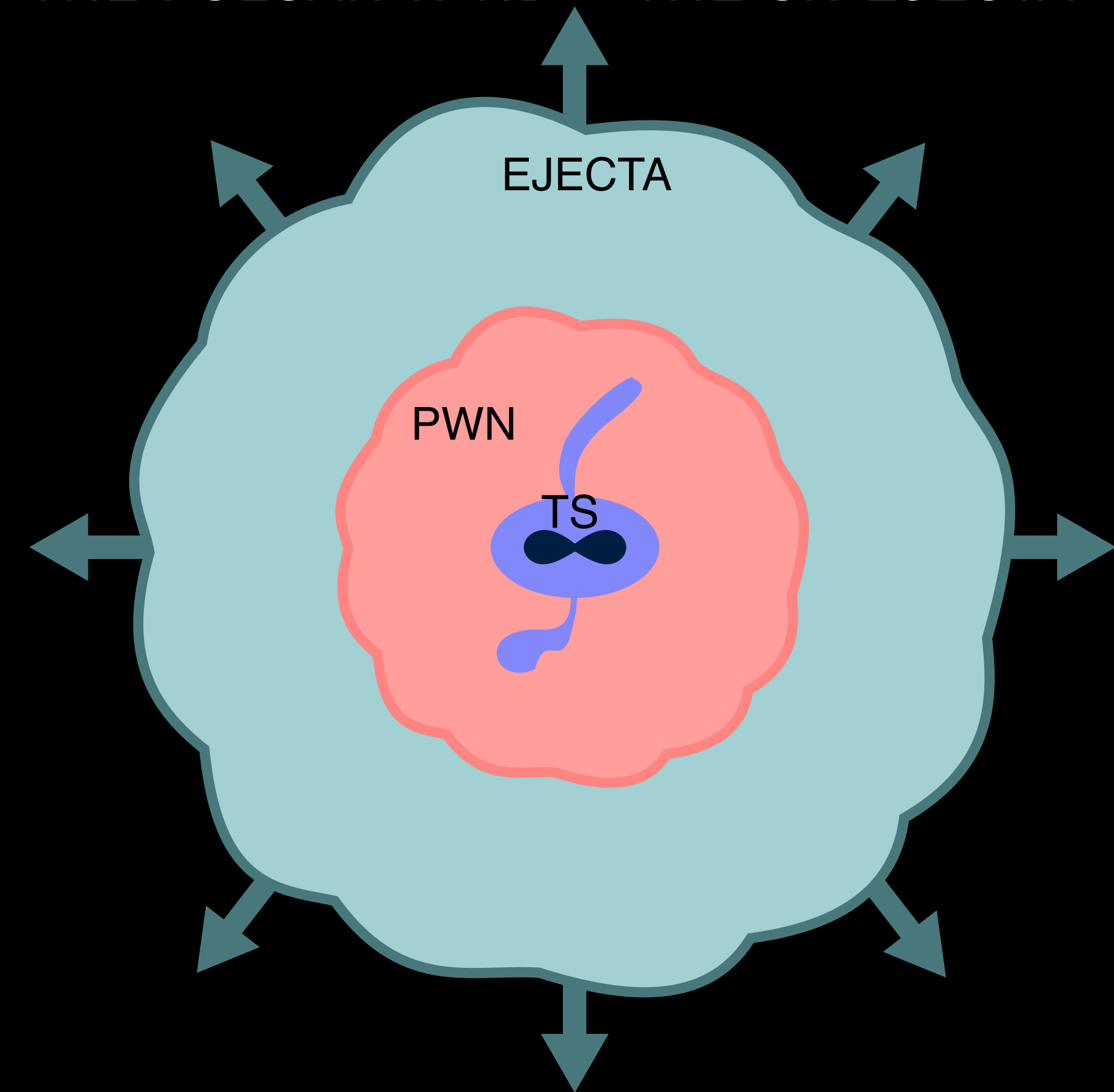
# INTRO: WHAT IS A PWN?

THE DEBRIS OF THE SUPERNOVA EXPLOSION OF A  
MASSIVE STAR ( $M \gtrsim 8 M_{\odot}$ )

MAIN INGREDIENTS:  
THE PULSAR WIND + THE SN EJECTA



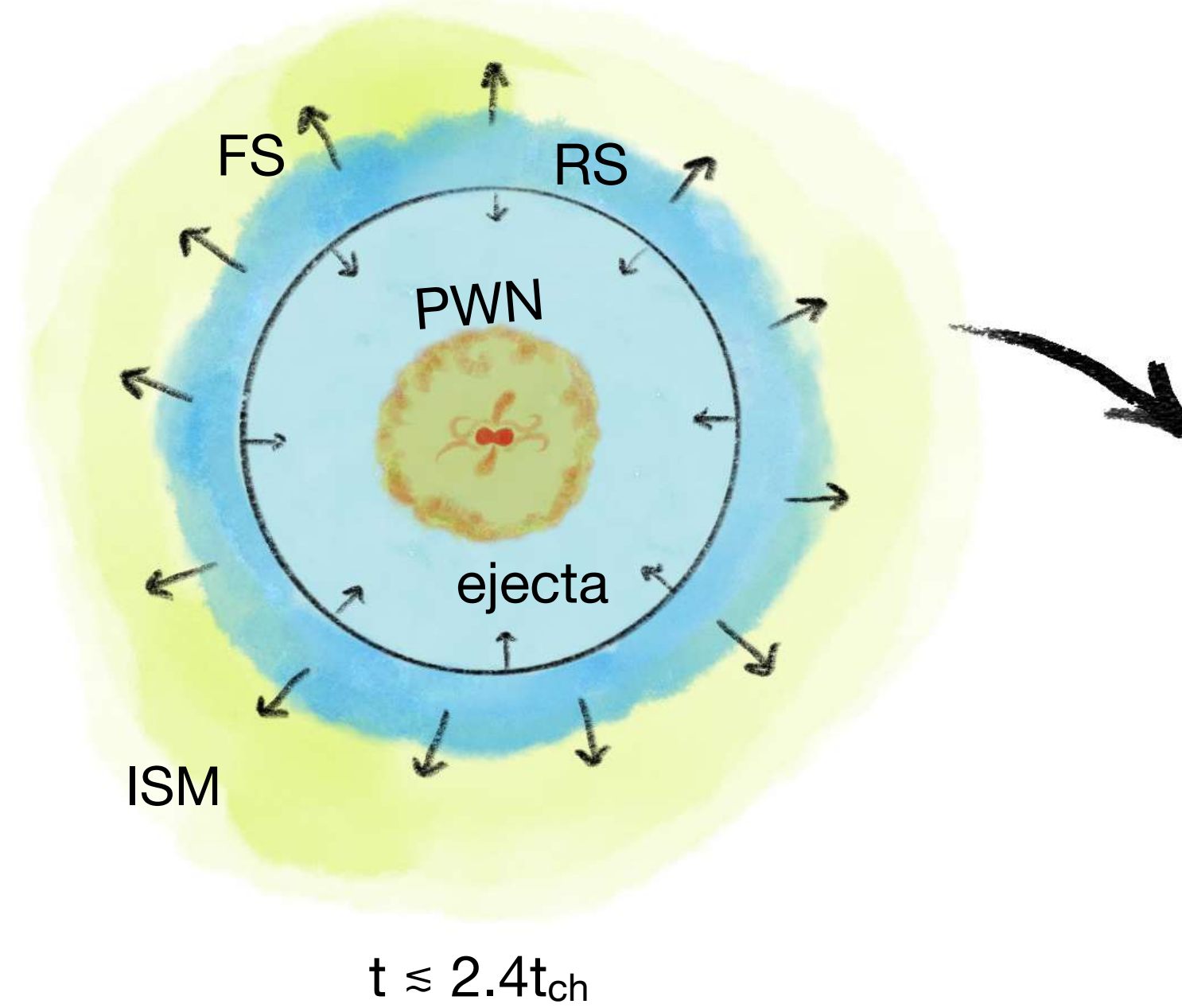
[NOT IN SCALE]



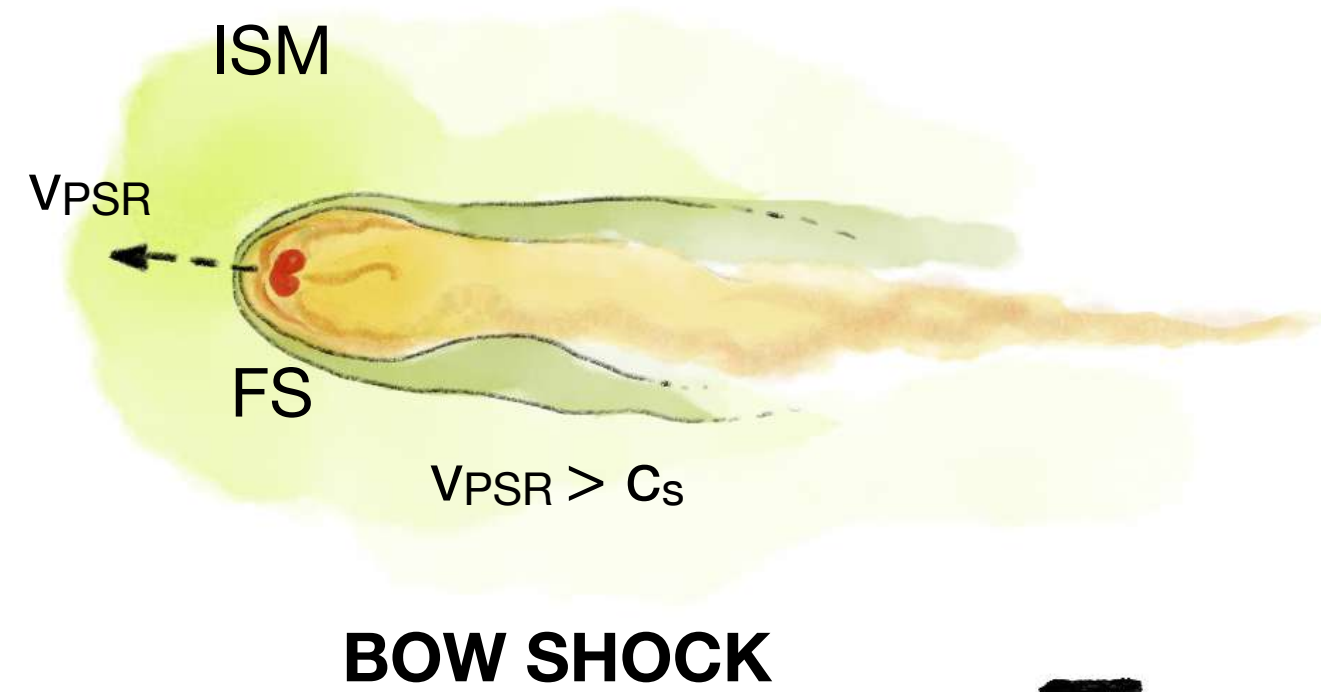
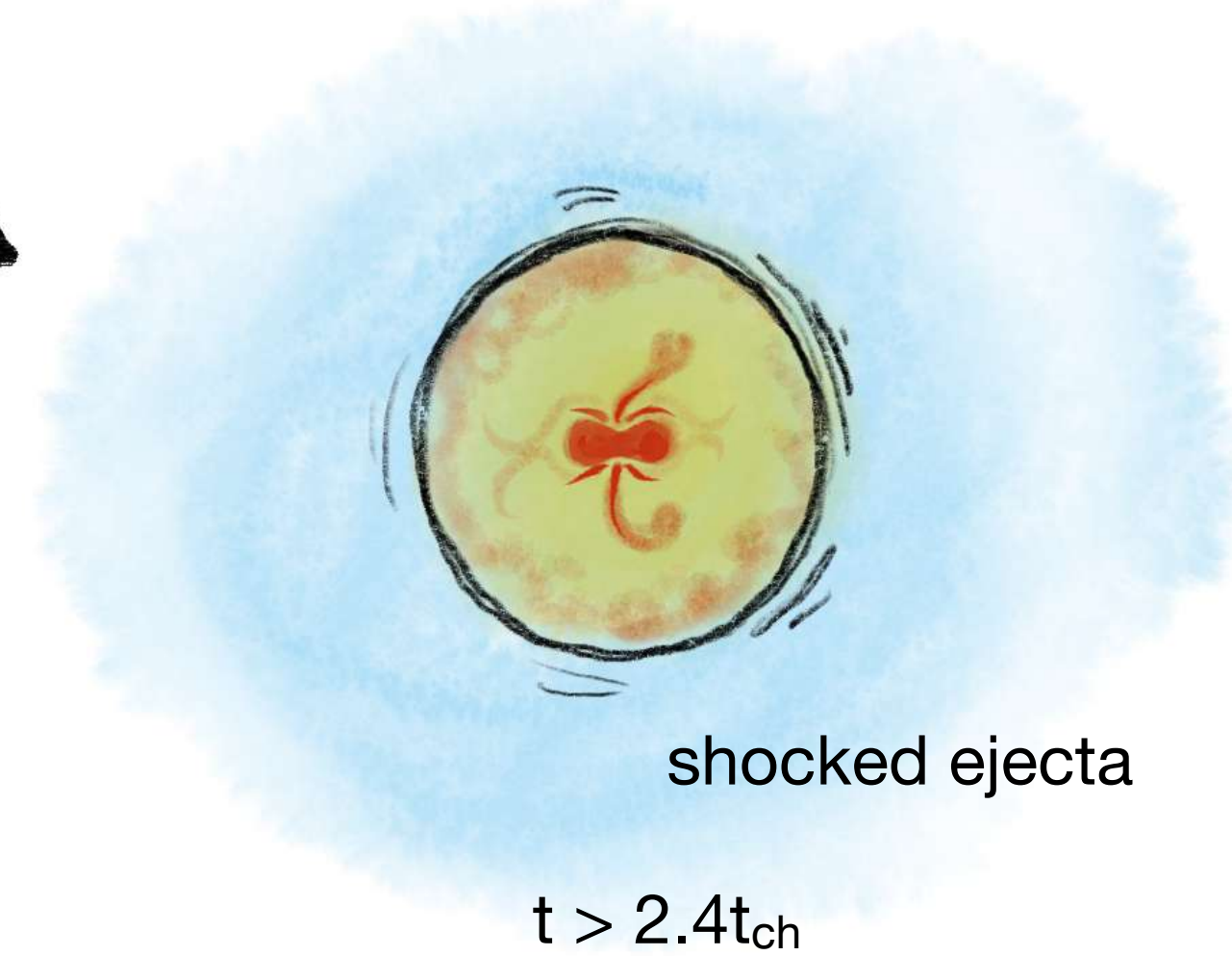


# EVOLUTIONARY PHASES OF PWNe

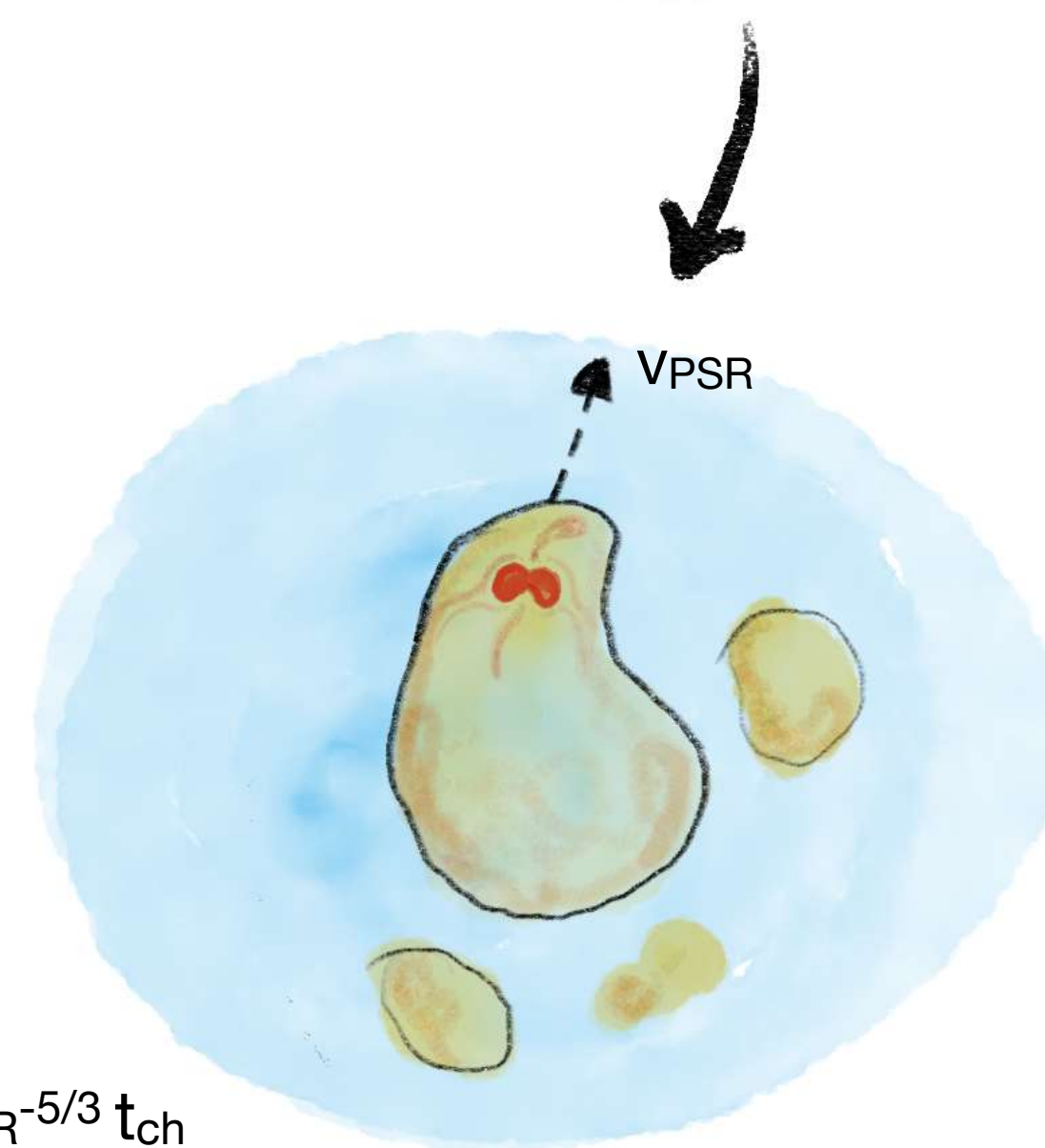
## FREE-EXPANSION



## REVERBERATION



## BOW SHOCK

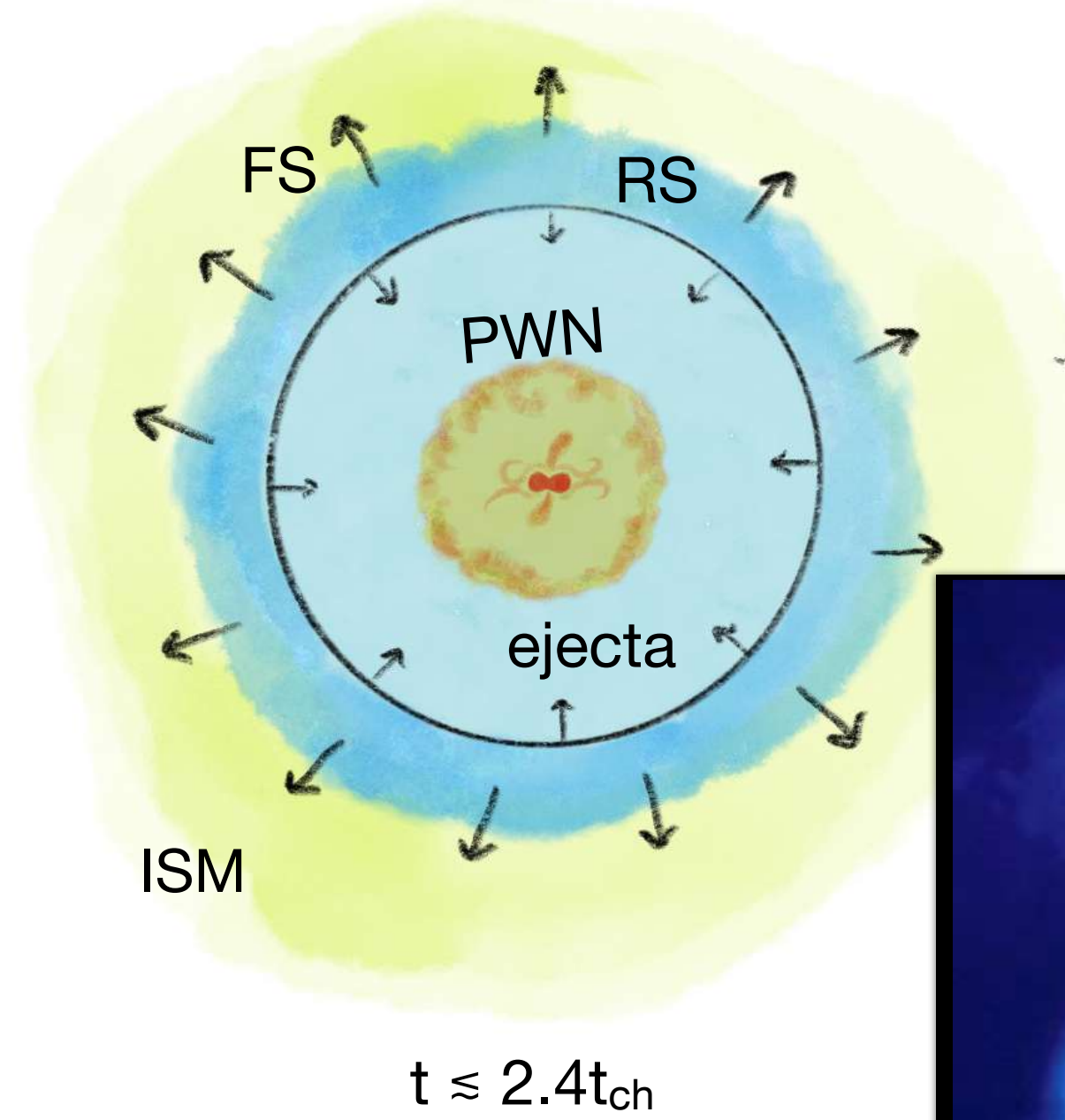


## POST-REVERBERATION

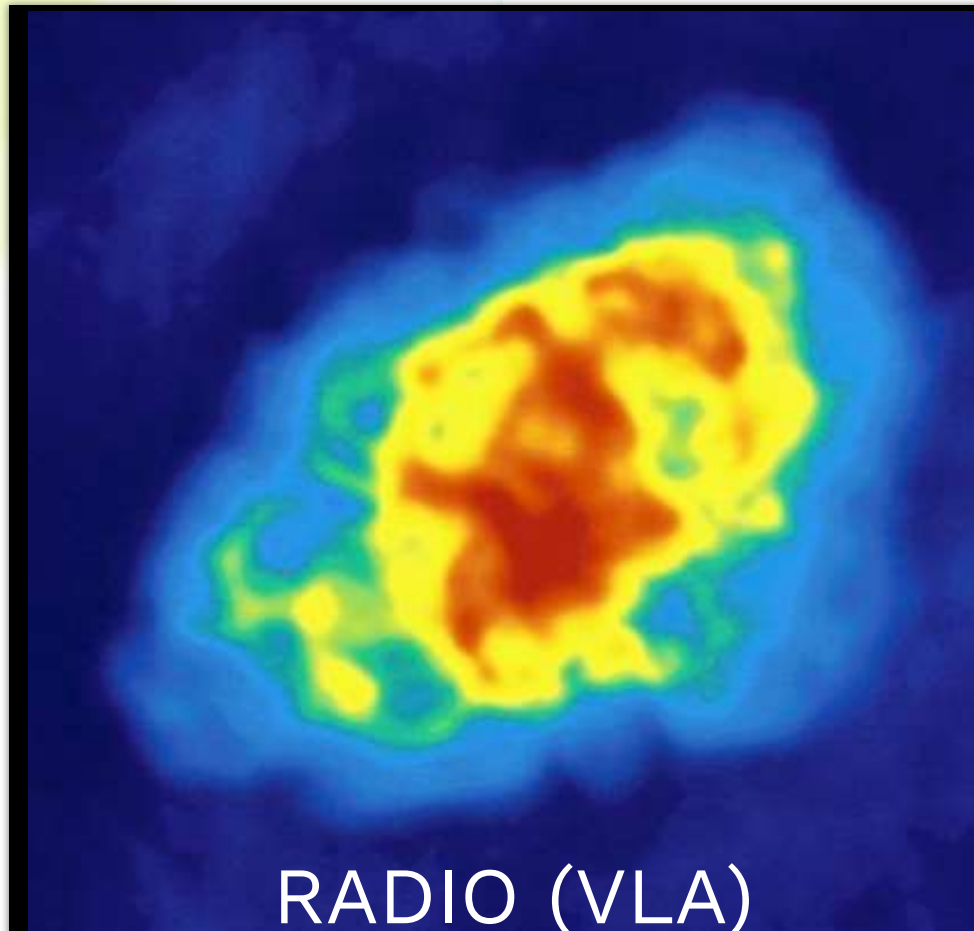


# EVOLUTIONARY PHASES OF PWNe

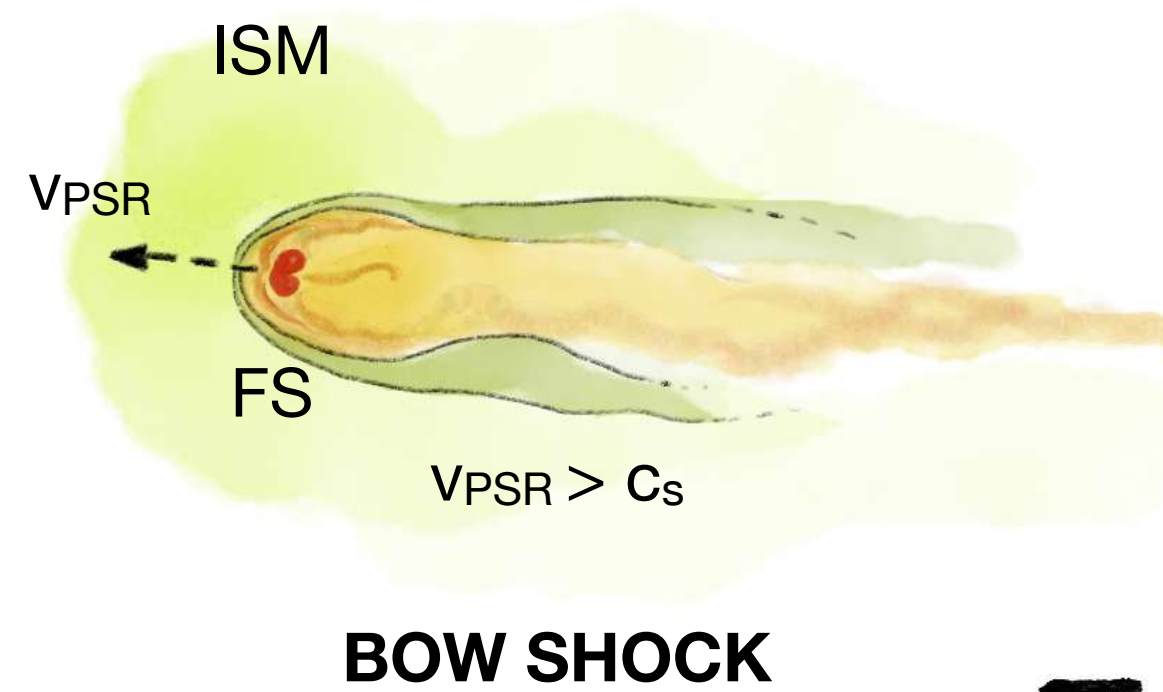
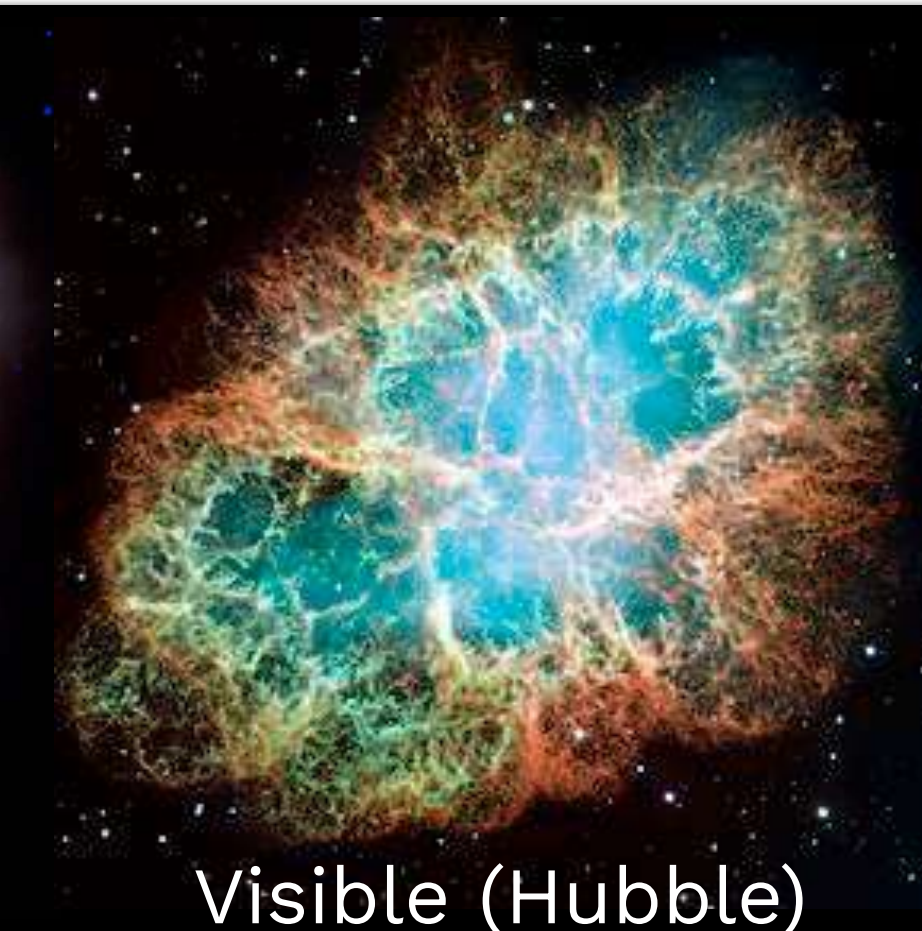
## FREE-EXPANSION



**YOUNG SYSTEMS**  
FILL CENTERED (NON-THERMAL) EMISSION  
ACROSS THE EM SPECTRUM

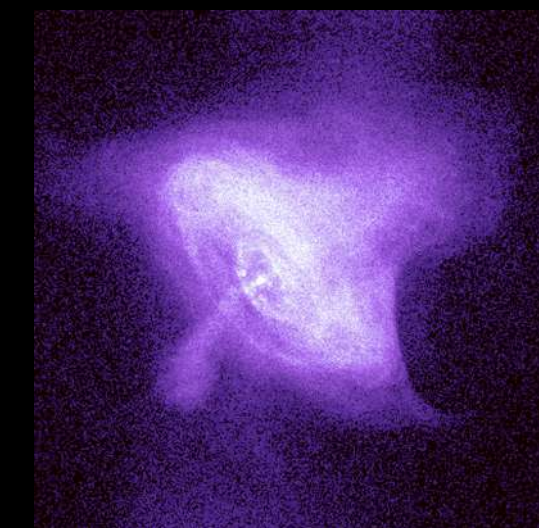
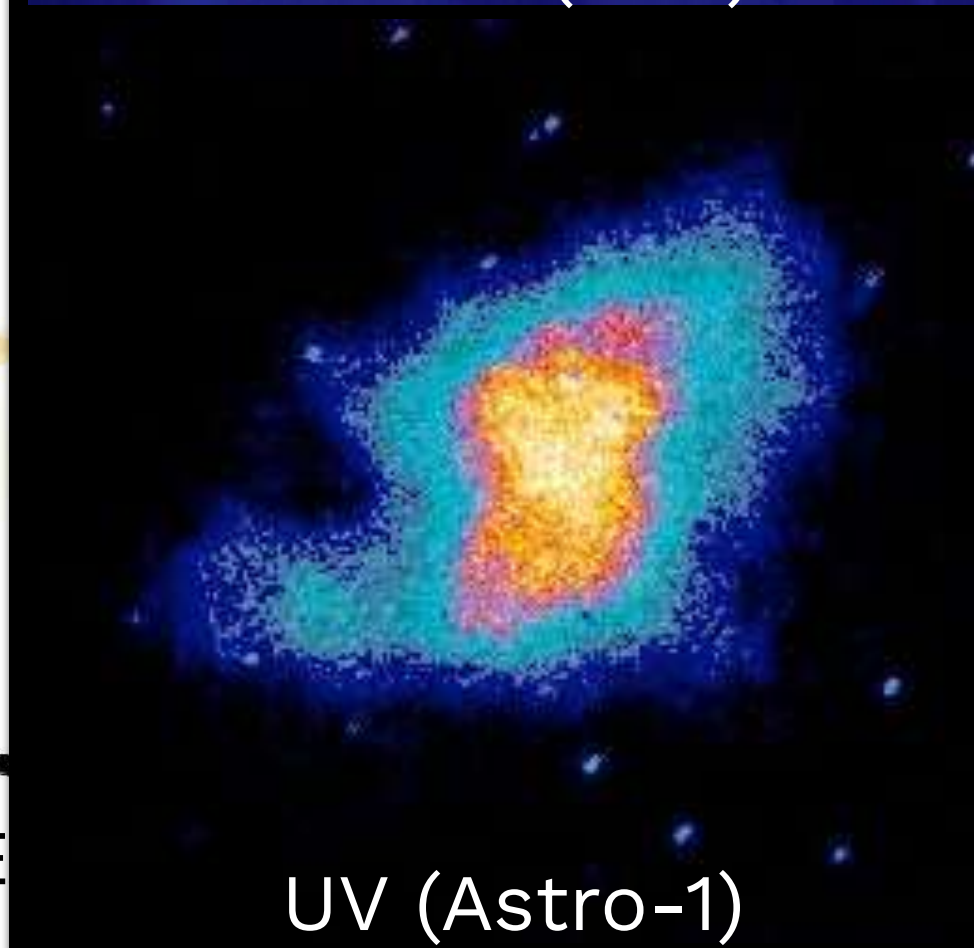


## CRAB NEBULA



## BOW SHOCK

$t_{esc} \propto (E)$

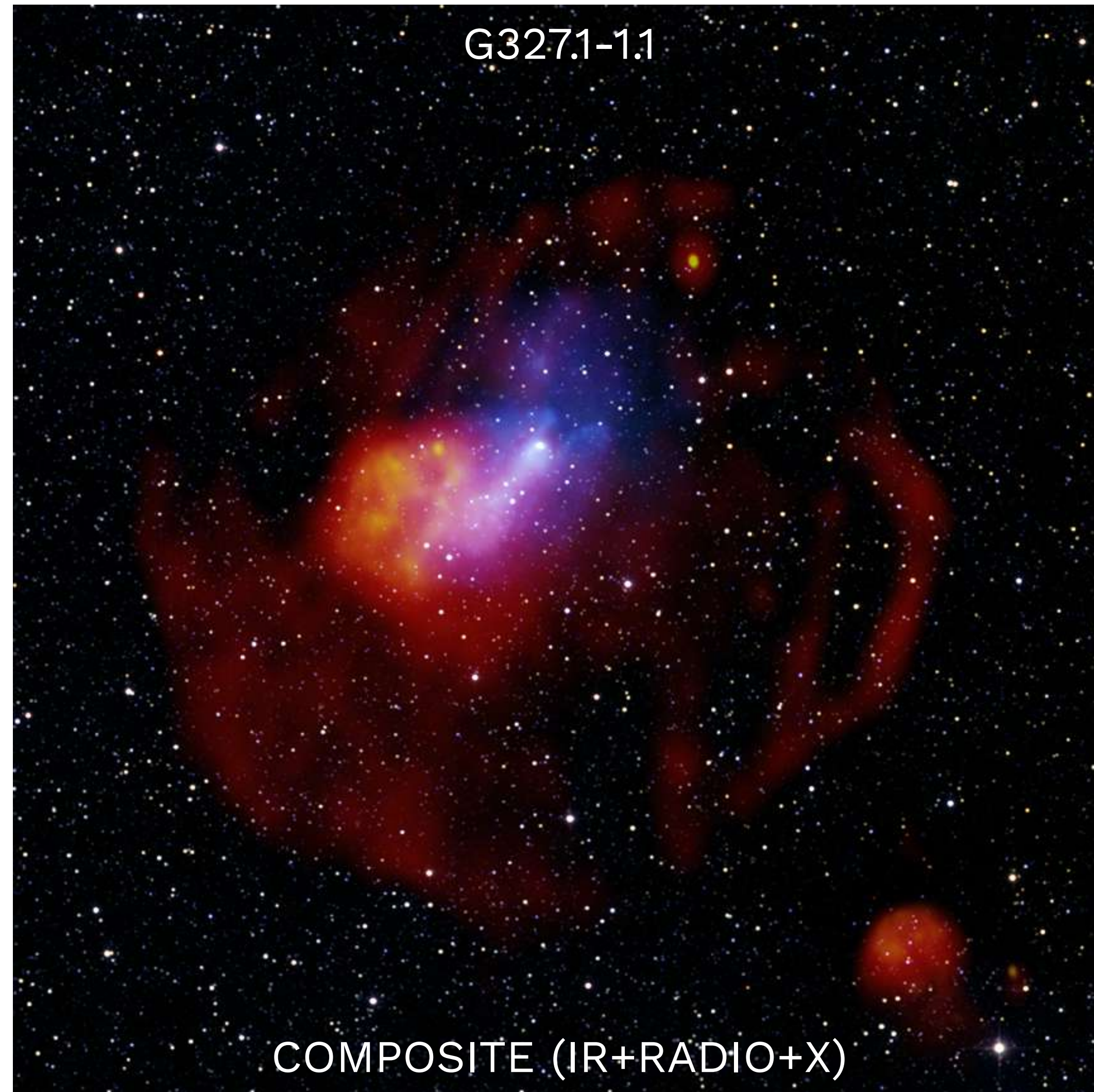


Pixel Size  
← →

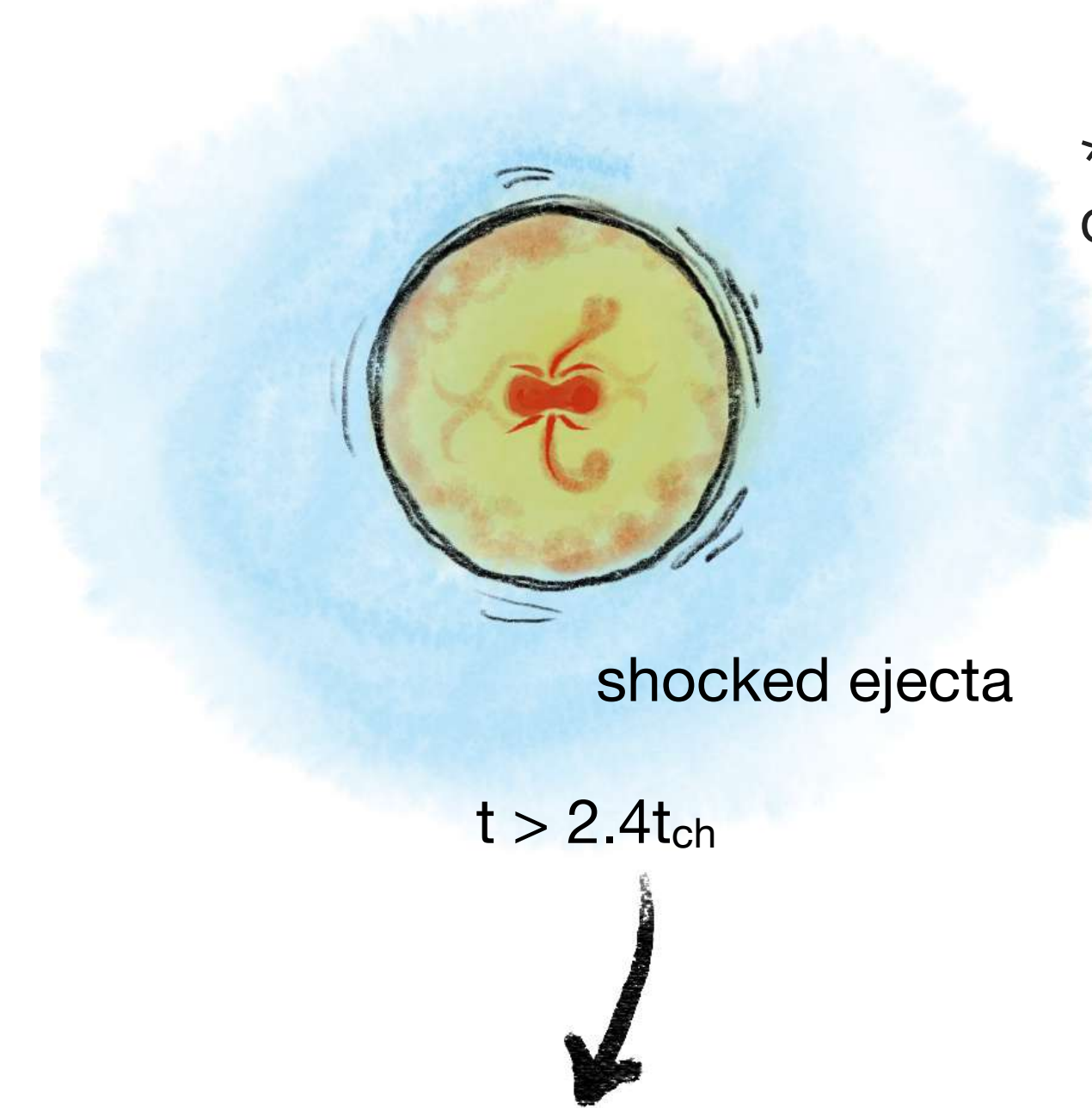
Hard X-Ray (HEFT)



# EVOLUTIONARY PHASES OF PWNe



## REVERBERATION

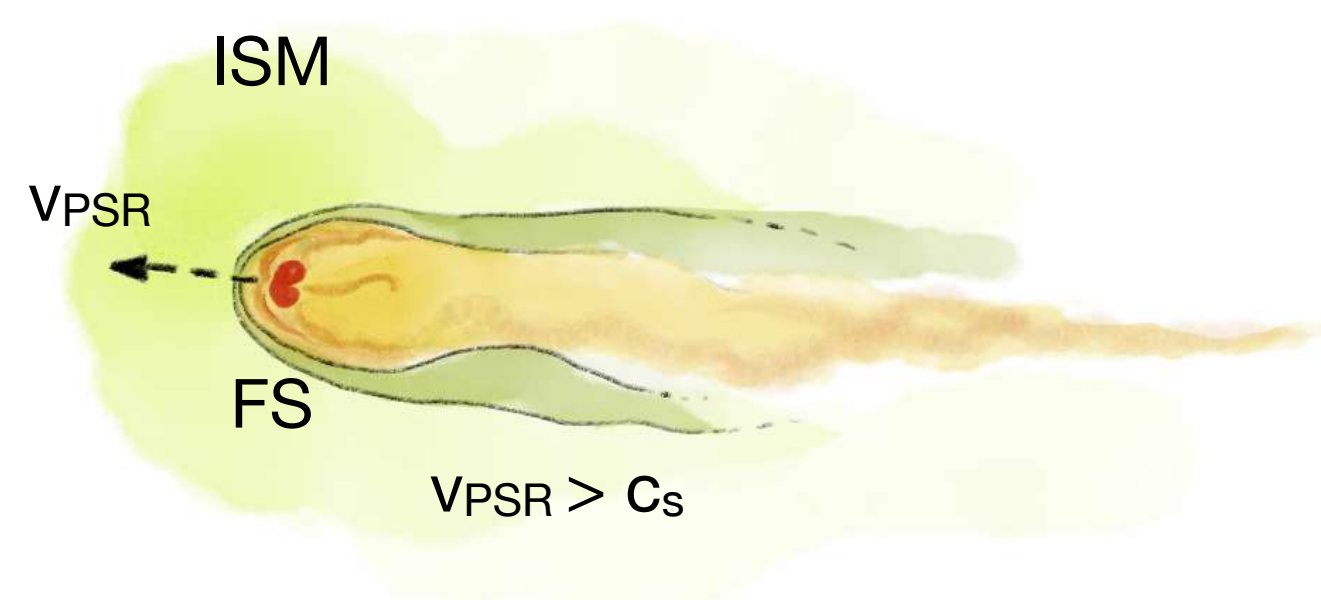
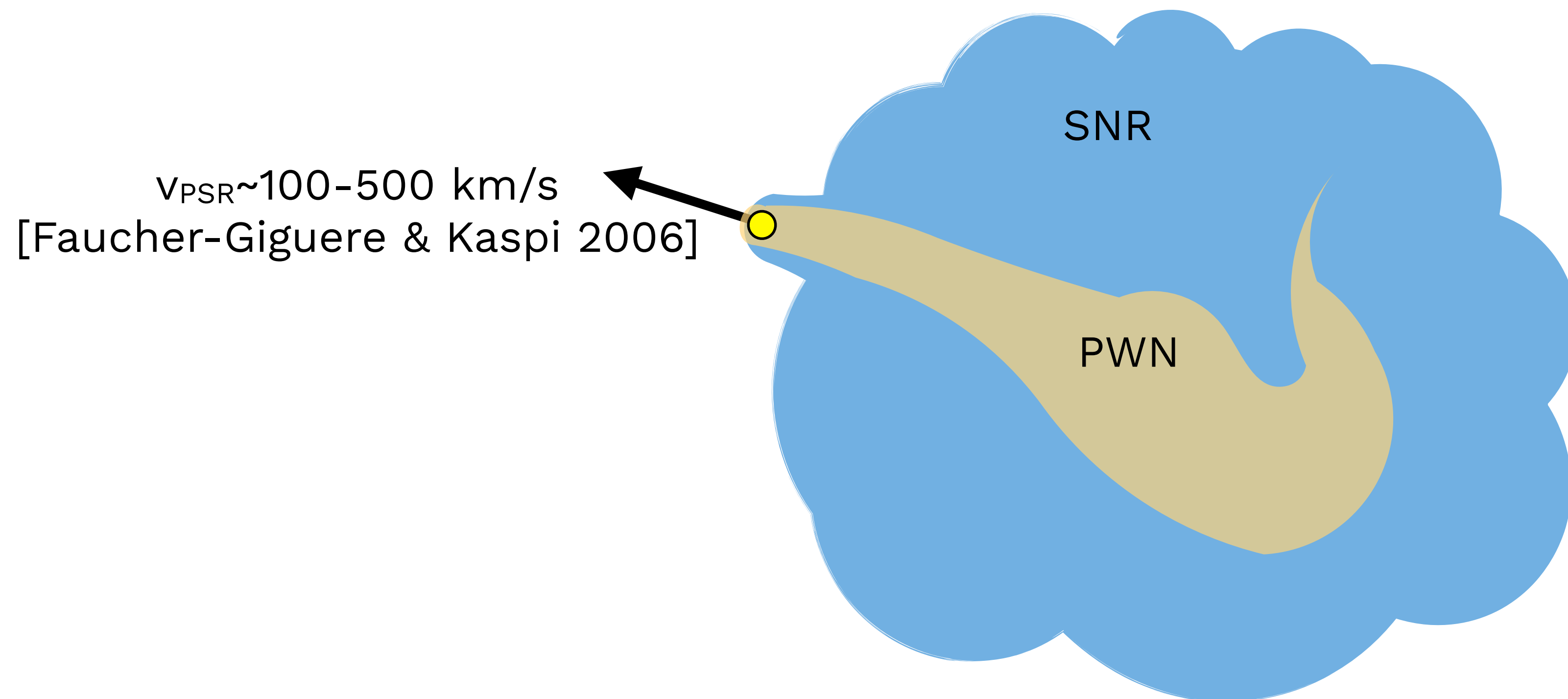


## MIDDLE AGED - REVERBERATING SYSTEMS

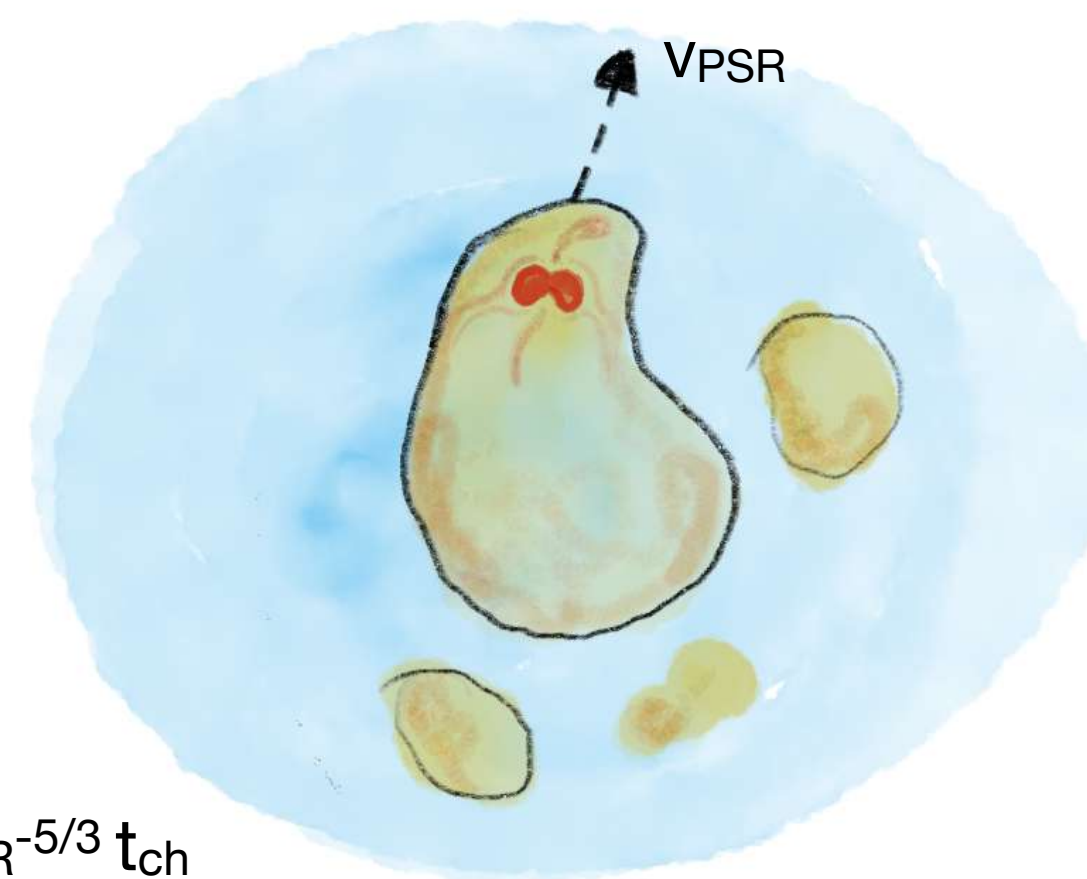
ONLY A FEW IDENTIFIED  
COMPLEX MORPHOLOGY DUE TO  
INTERACTION WITH SNR



# EVOLUTIONARY PHASES OF PWNe



**BOW SHOCK**



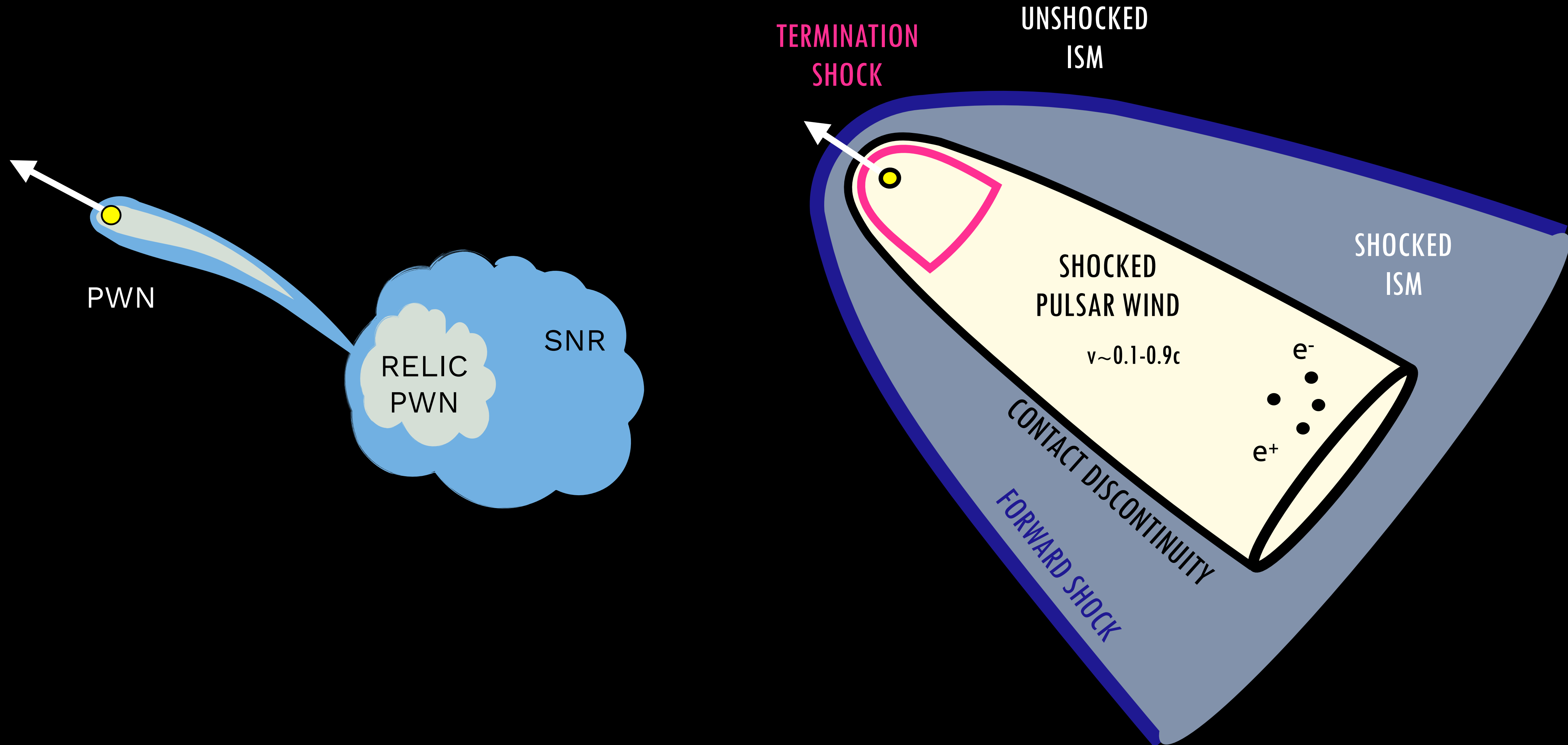
**POST-REVERBERATION**

$t_{esc} \propto (E_{sn}/M_{ej})^{5/6} V_{PSR}^{-5/3} t_{ch}$

# BOW SHOCK PWNe

THE PWN IS RESHAPED BY THE SUPERSONIC MOTION IN THE ISM

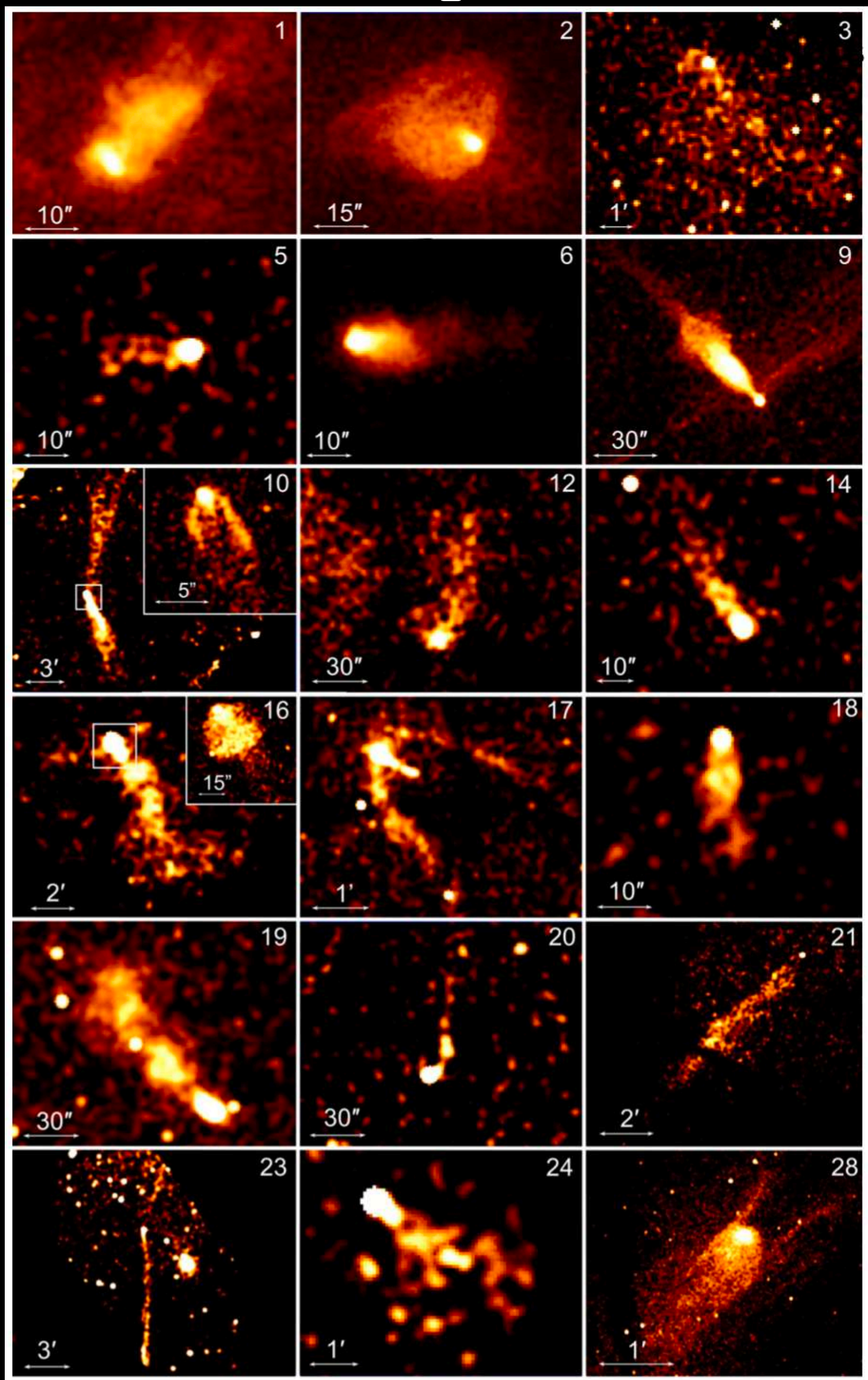
$$v_{\text{PSR}} \sim 100\text{--}500 \text{ km/s} > c_s$$





# HOW DO THEY DO LOOK LIKE

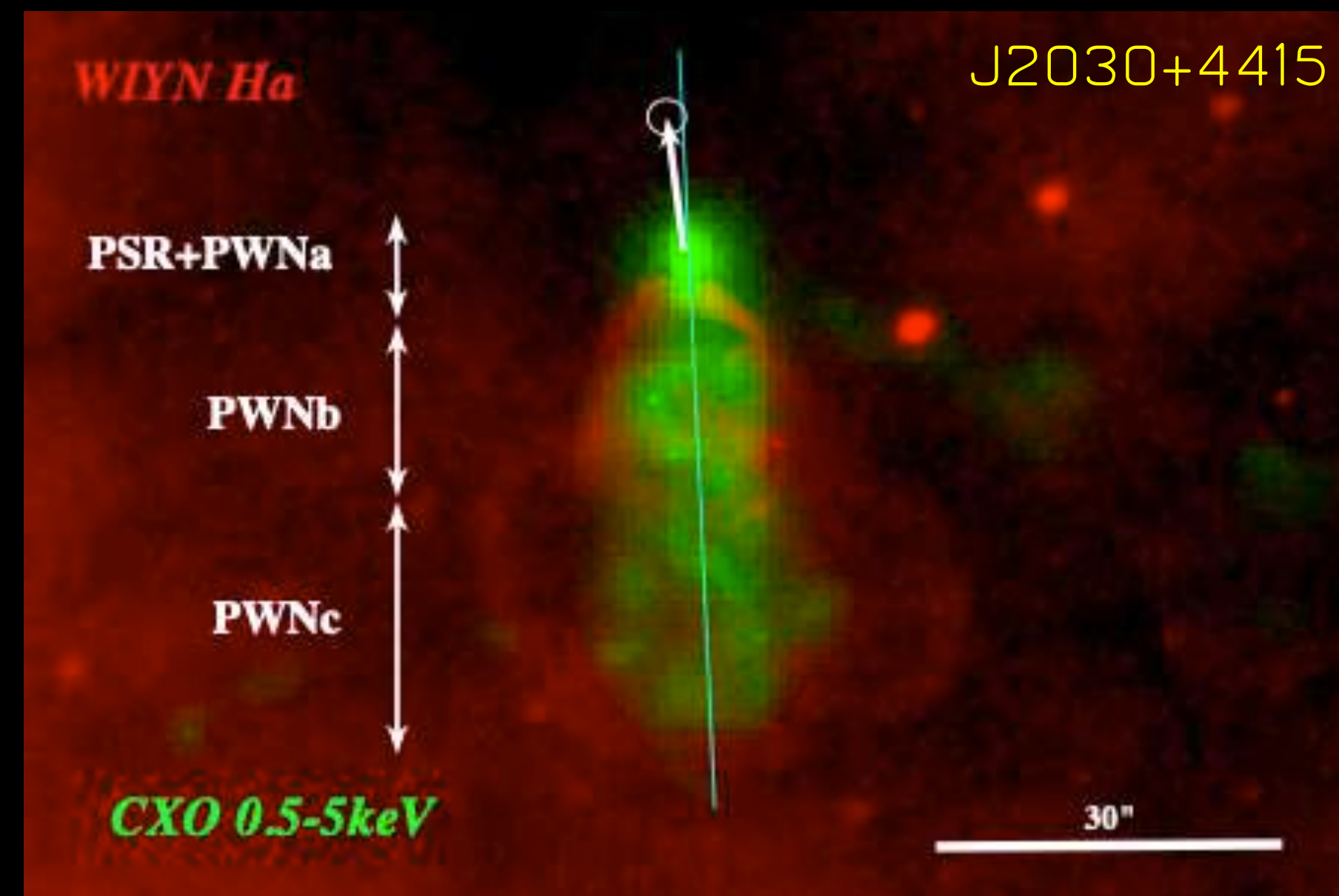
X-rays



[Kargaltsev et al. 2017]

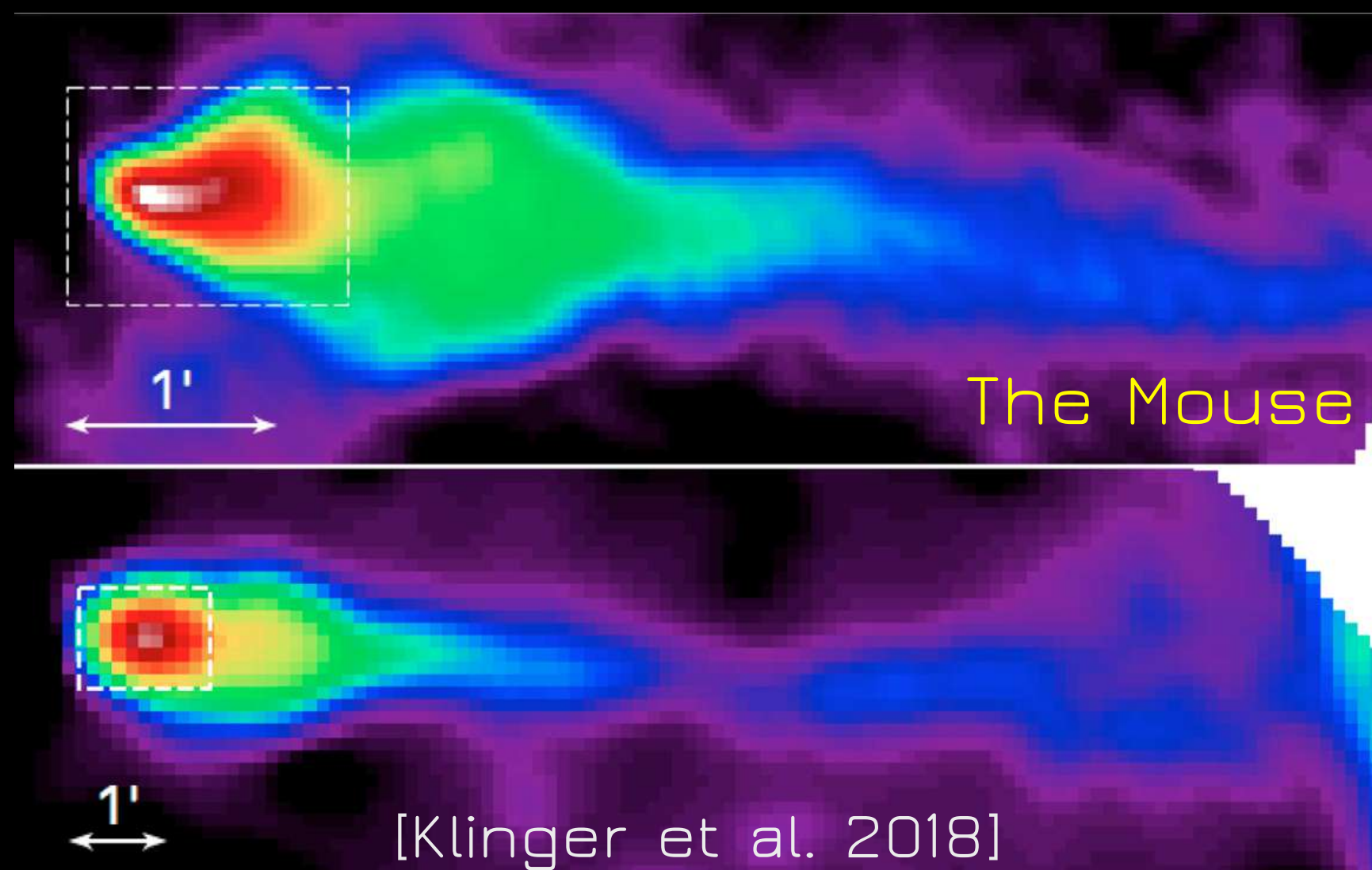
COMETARY SHAPED:  
BRIGHT HEADS AND  
EXTENDED TAILS  
BEHIND THE PULSAR

Combined Ha + X-rays



[De Vrie & Romani 2020]

Radio (VLA)

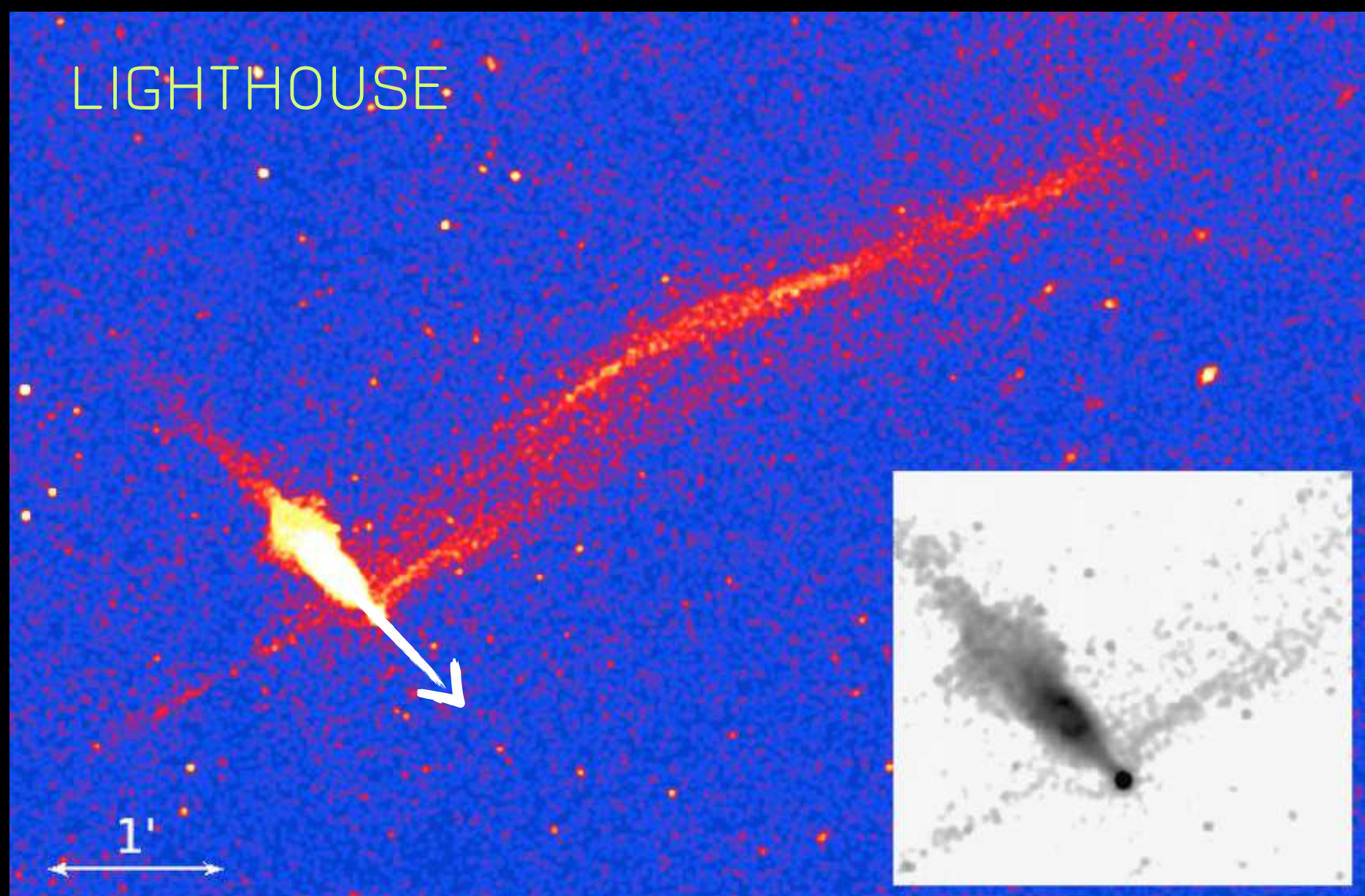


[Klinger et al. 2018]

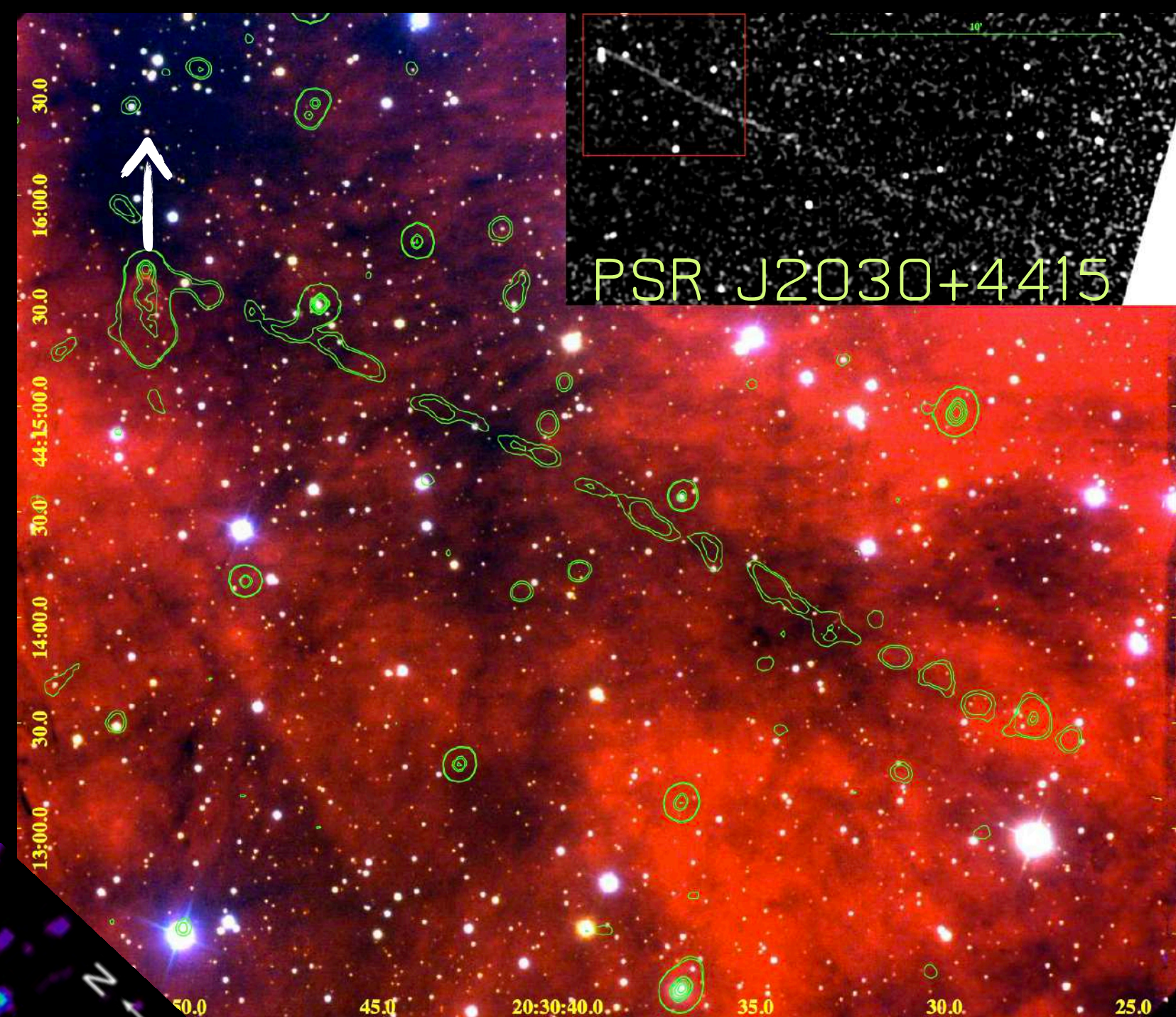
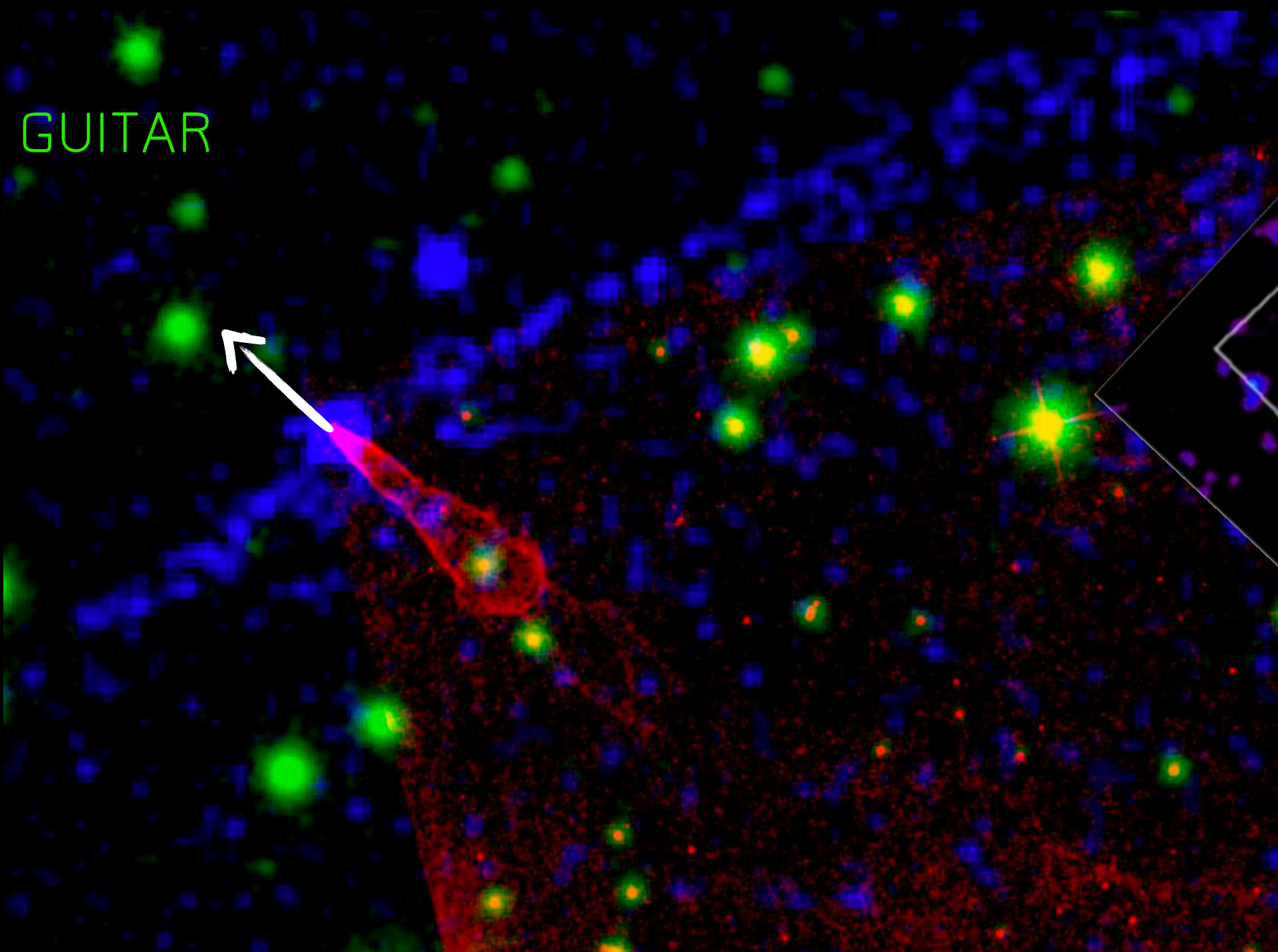


# EVIDENCE OF PARTICLE ESCAPE: MISALIGNED X-RAY TAILS

[Pavan et al. 2016]

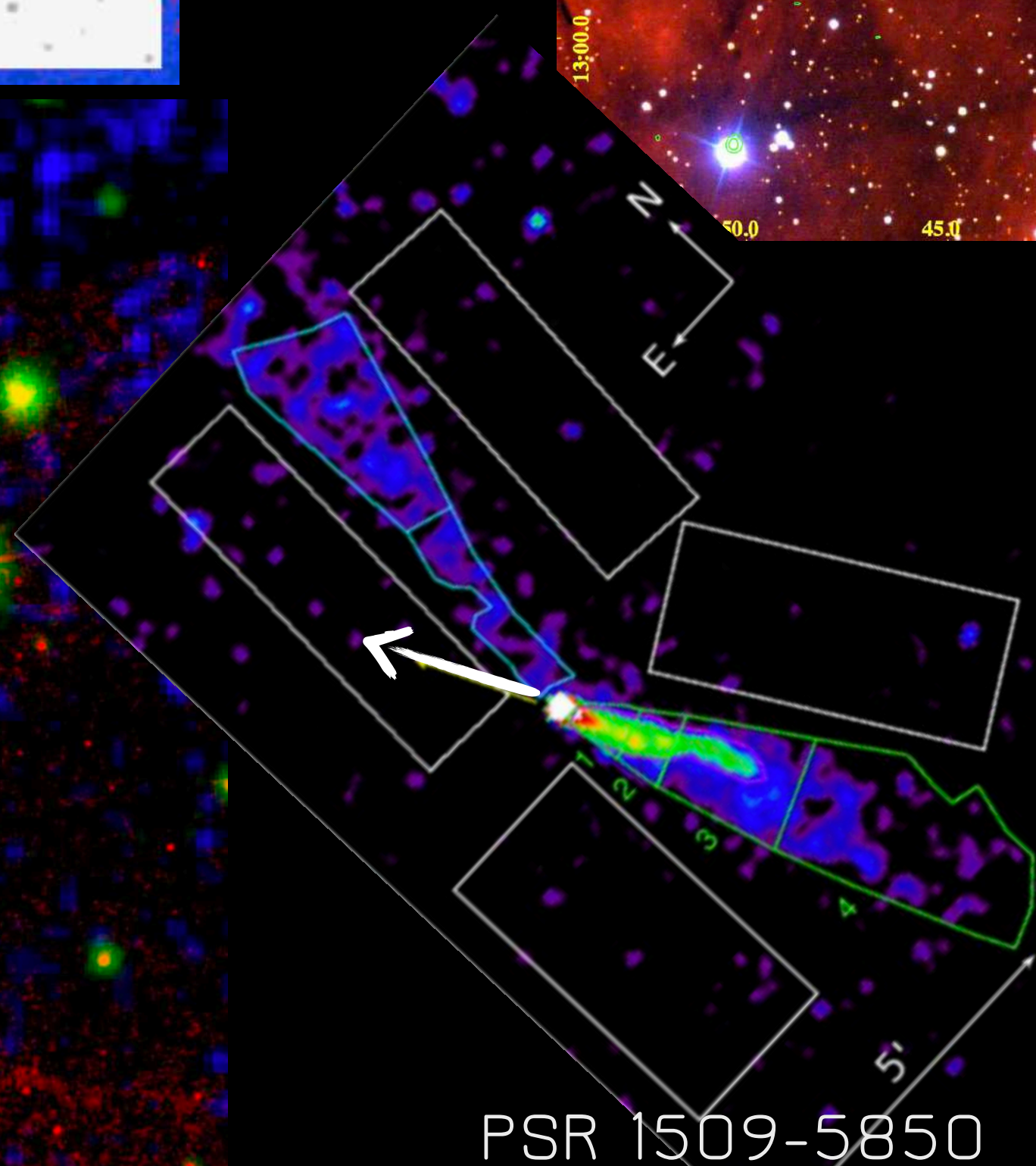


[Cordes et al. 1993,  
De Vries et al. 2022]

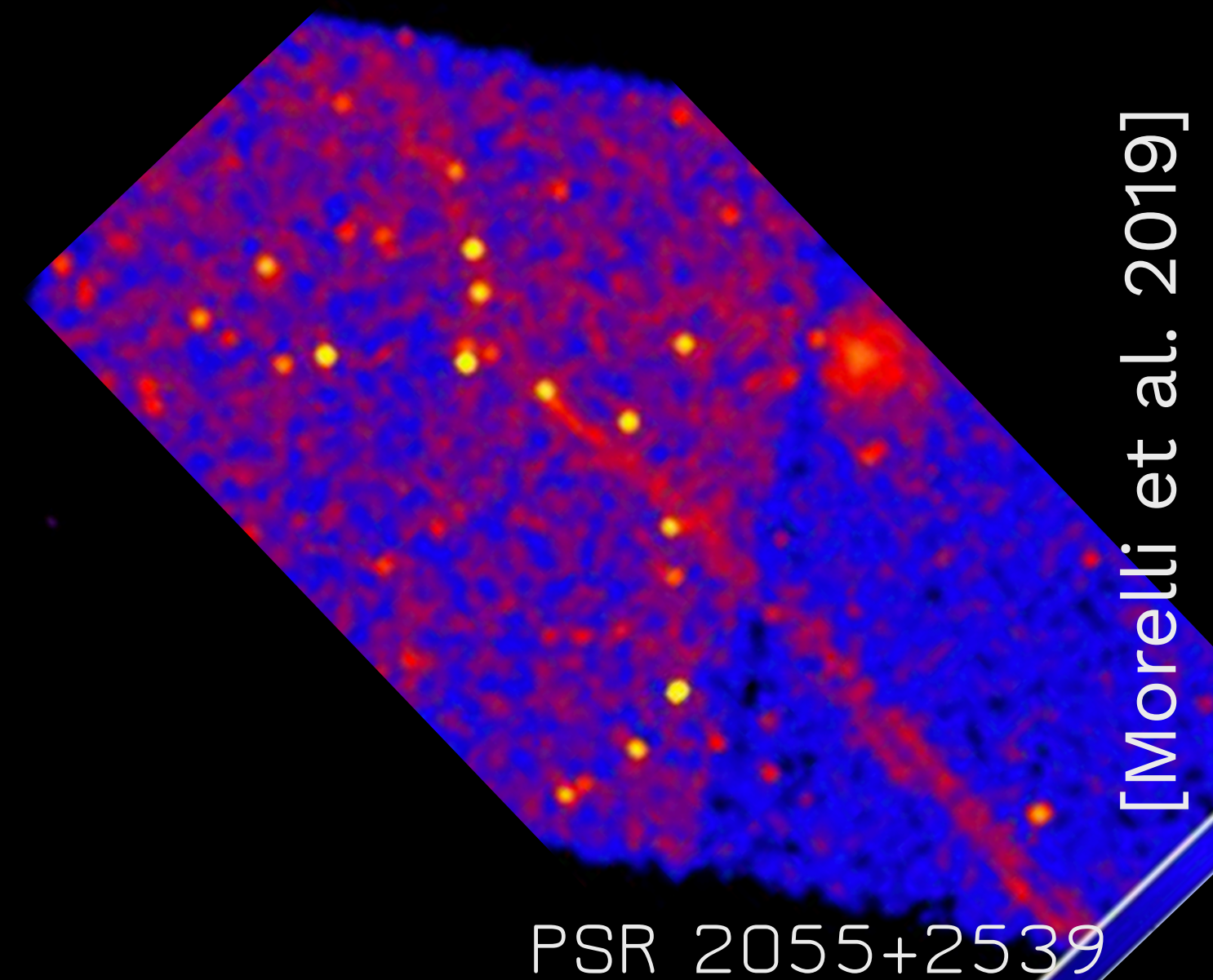


[De Vrie & Romani 2020,2022]

[Klingler et al. 2016]

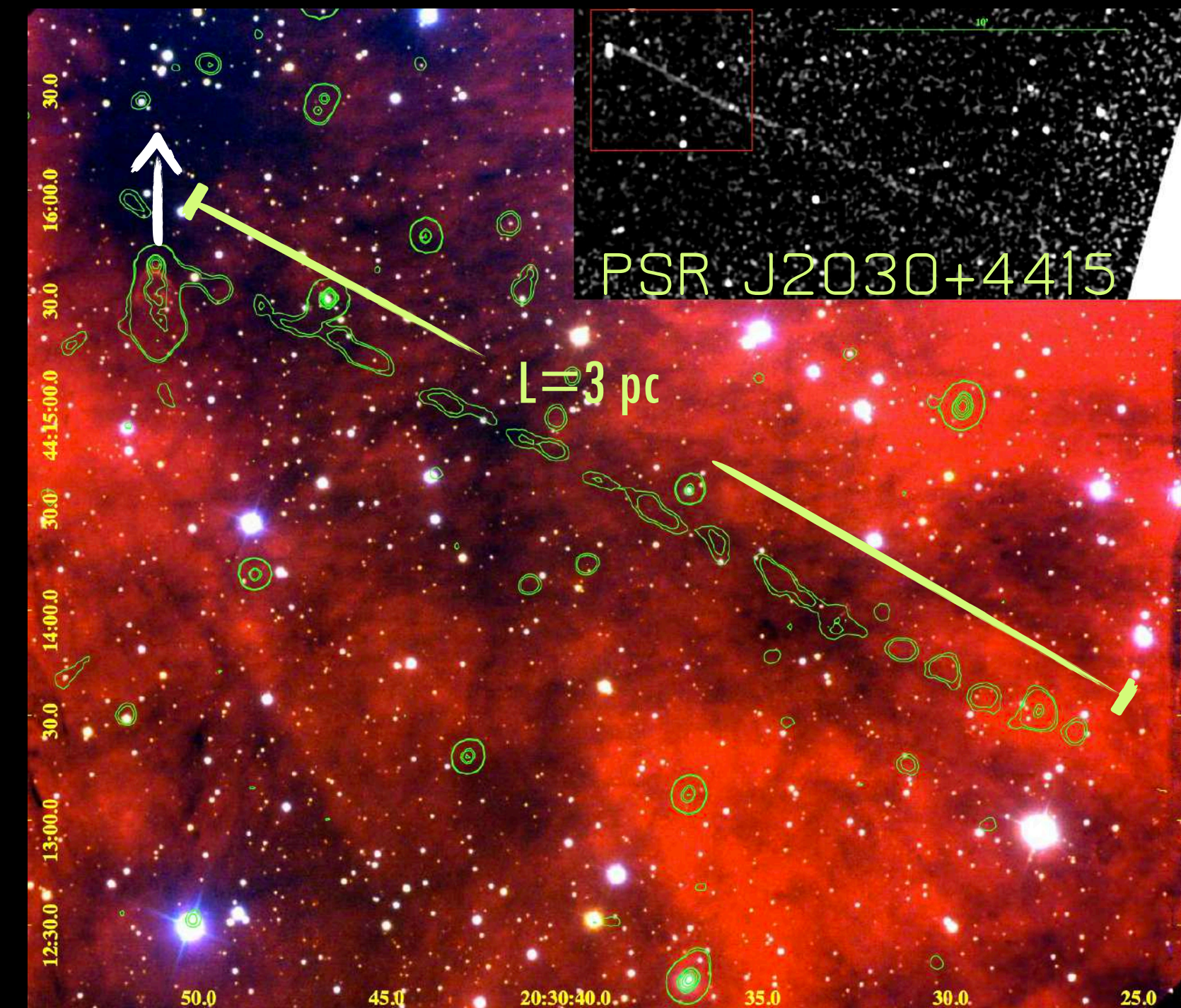
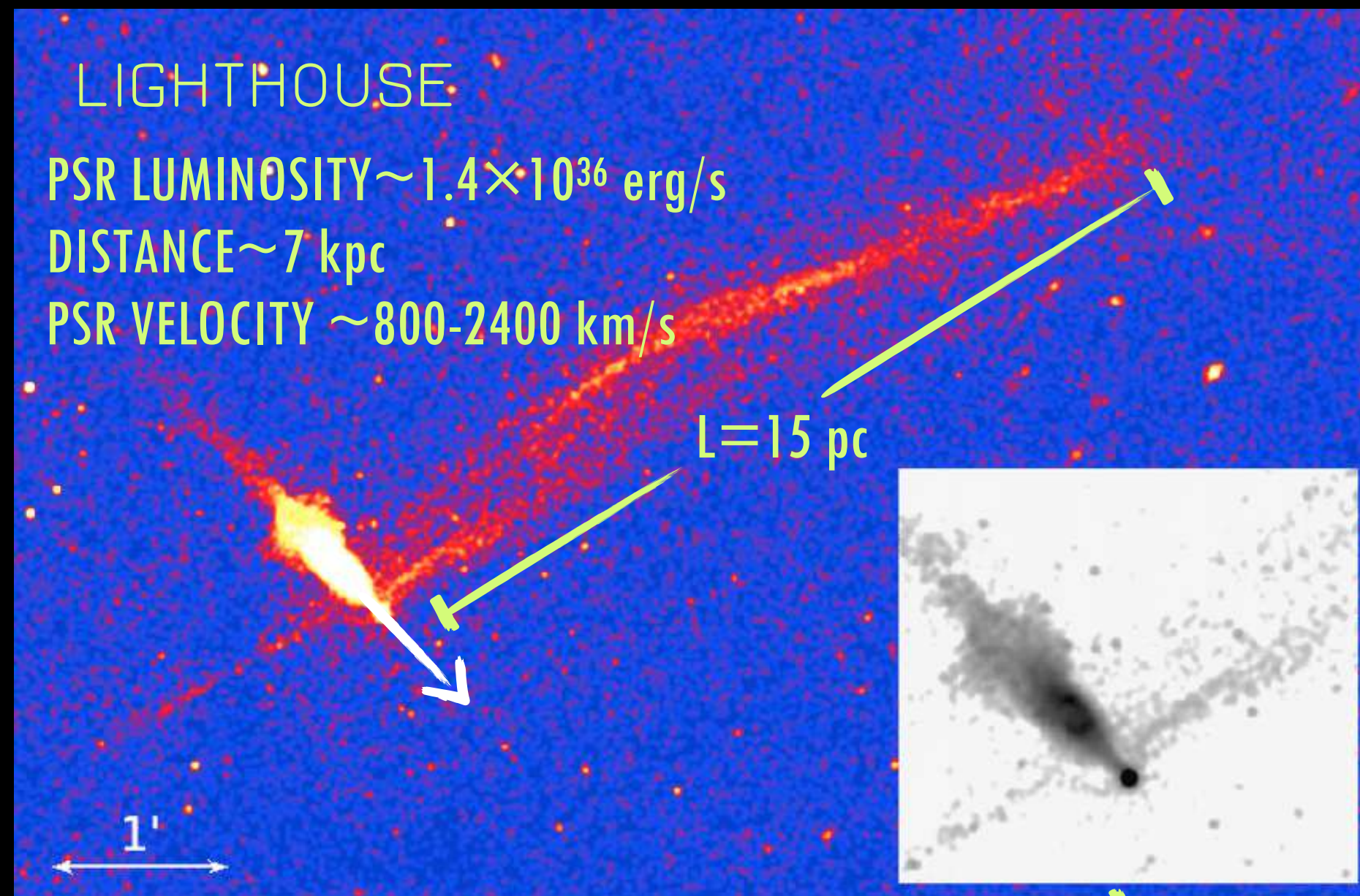


[Morelli et al. 2019]

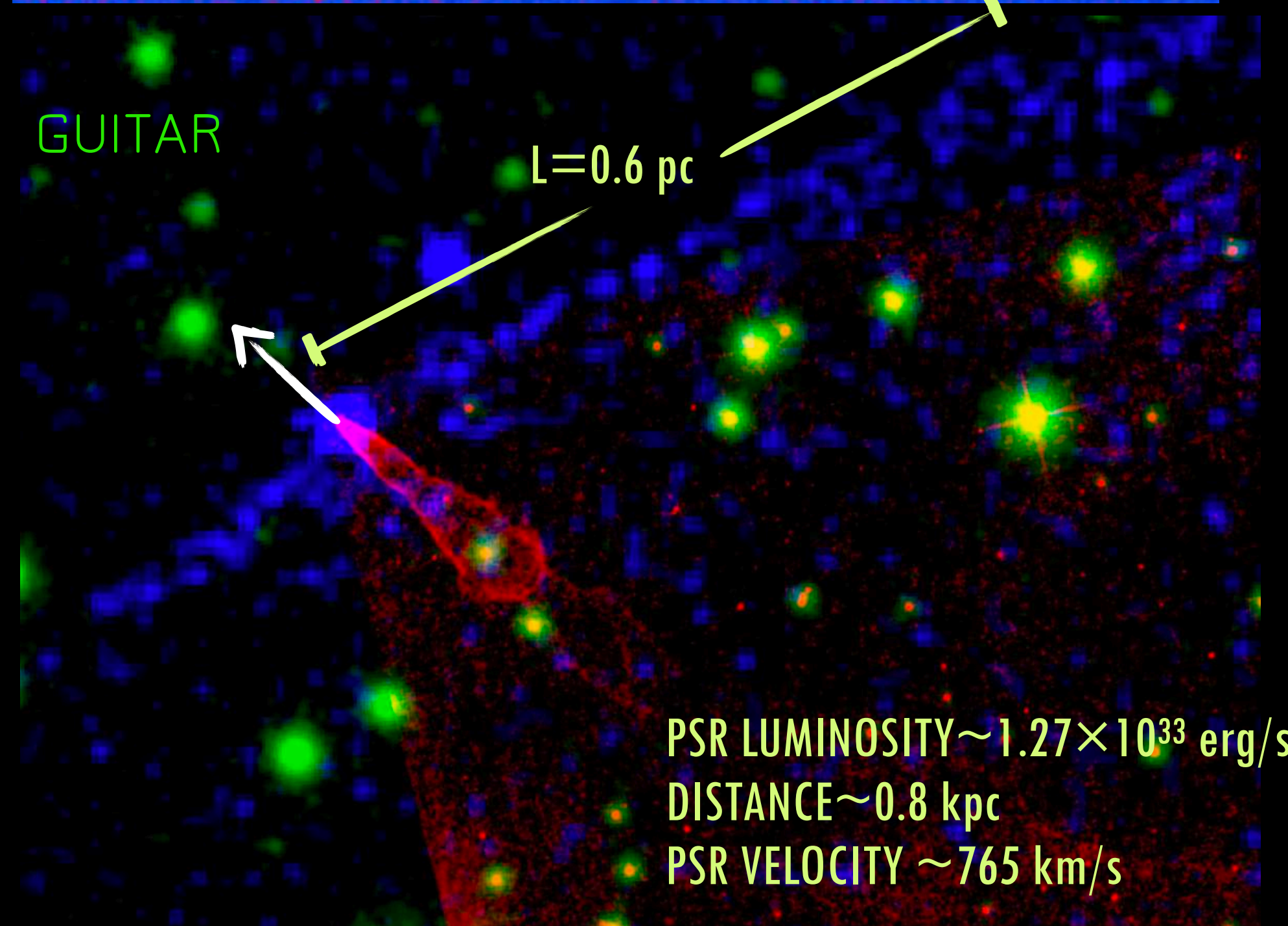




# EVIDENCE OF PARTICLE ESCAPE: MISALIGNED X-RAY TAILS



PSR LUMINOSITY  $\sim 2.2 \times 10^{34}$  erg/s  
DISTANCE  $\sim 720$  pc  
PSR VELOCITY  $\sim 290$  km/s





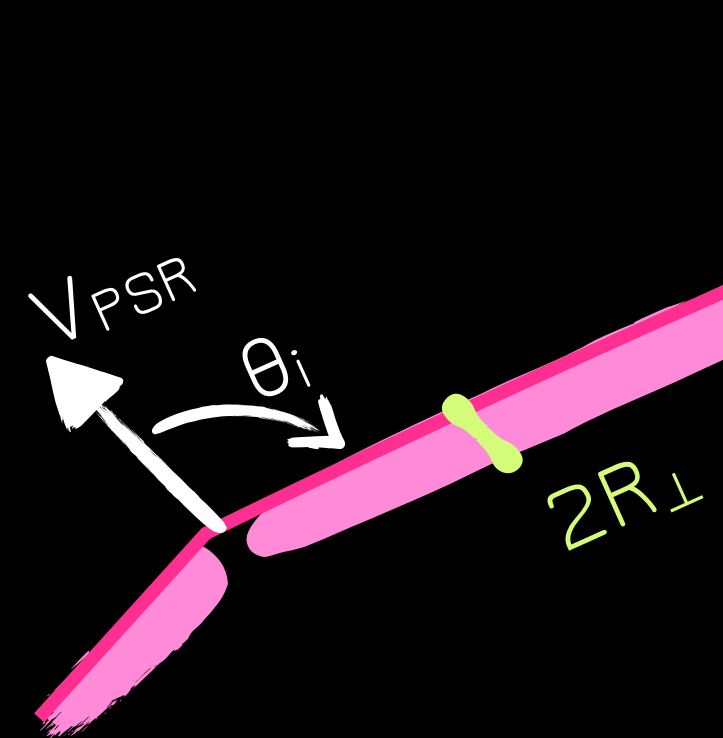
# INTERPRETING THE X-RAY FEATURES

Can be explained as synchrotron radiation emitted by particles escaping the nebula, for which you need:

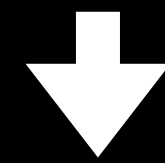
- 1 Massive particle escape
- 2 Particles of energy close to the maximum theoretically available one, connected with the pulsar potential drop:

$$\gamma_{\text{MPD}} = \frac{e}{mc^2} \sqrt{\frac{\dot{E}}{c}} \simeq 3 \times 10^8 \left( \frac{\dot{E}}{10^{34} \text{ erg/s}} \right)$$

- 3 A mechanism to amplify the ambient magnetic field, estimated from the transverse size of the feature:



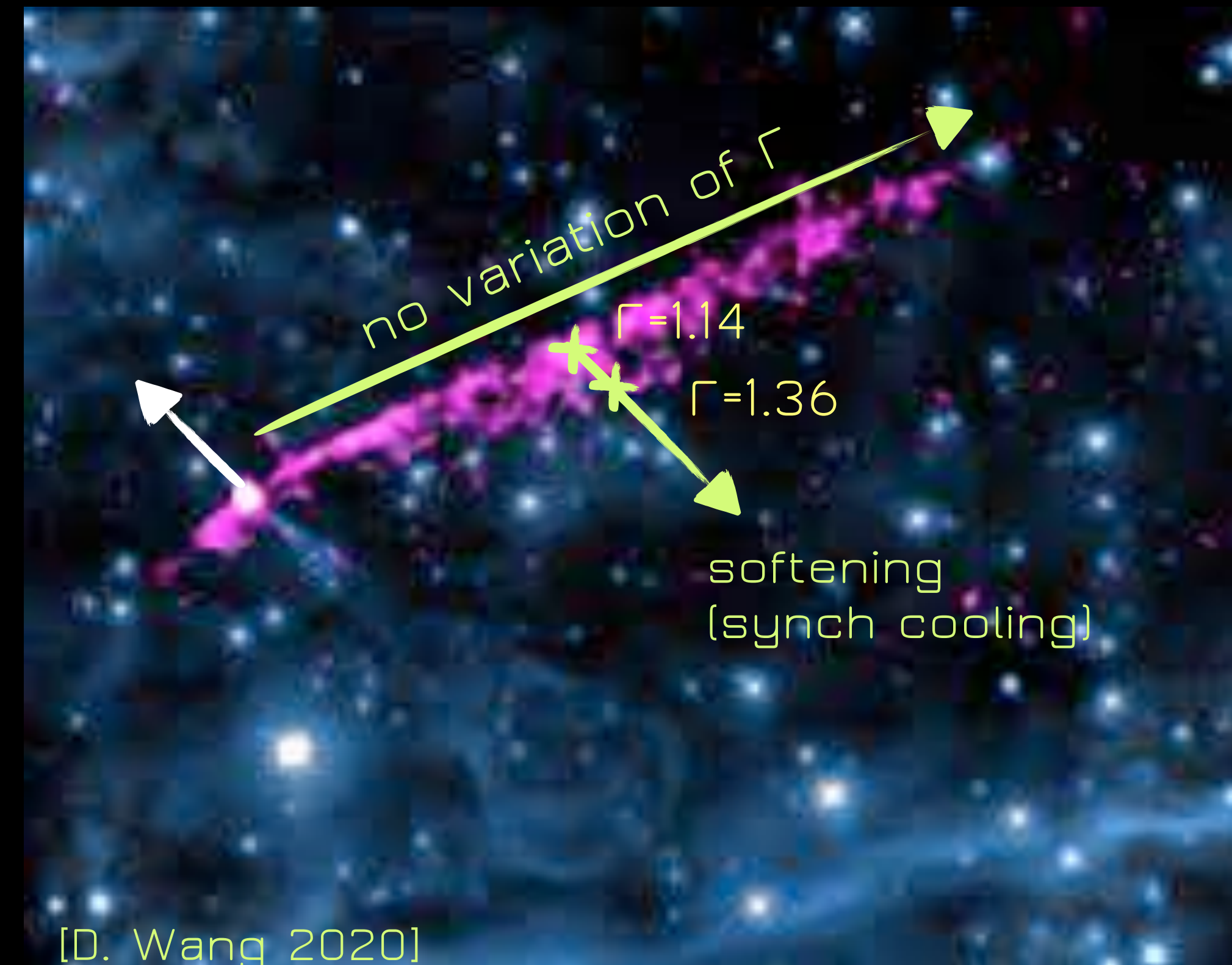
$$2R_{\perp} / (V_{\text{PSR}} \sin \theta_i) \simeq \tau_{\text{sync}}$$



$$\Delta B \simeq 40 \left( \frac{E_{\text{ph}}}{2 \text{ keV}} \right)^{-1/3} \mu\text{G}$$

[Bandiera 2008]

Guitar nebula [0.5-7 keV + H $\alpha$ ]





# HOW TO AMPLIFY THE FIELD?

An efficient way is to excite the non-resonant Bell instability, for which you need:

- ① An electric current  $J$
- ② The escaping particles to carry an energy density larger than the magnetic energy density of the ambient medium:

$$n_{\text{esc}} mc^2 (\gamma_{\text{esc}} - 1) \geq B_0^2 / (8\pi)$$

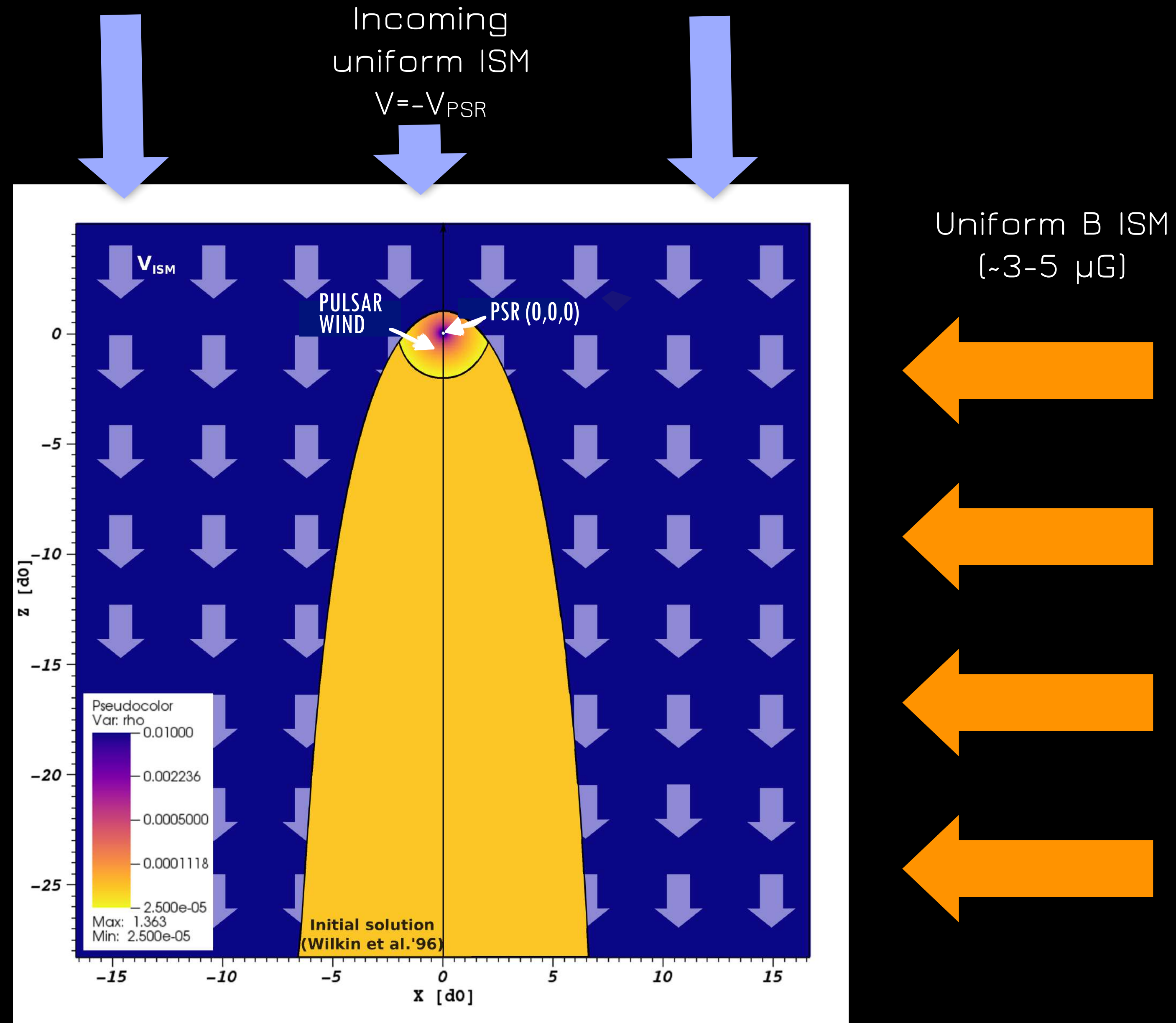


# HOW DO PARTICLES ESCAPE? INVESTIGATING THE PROCESS THROUGH NUMERICAL SIMULATIONS

Simulations with the PLUTO code [Mignone et al. 2007]



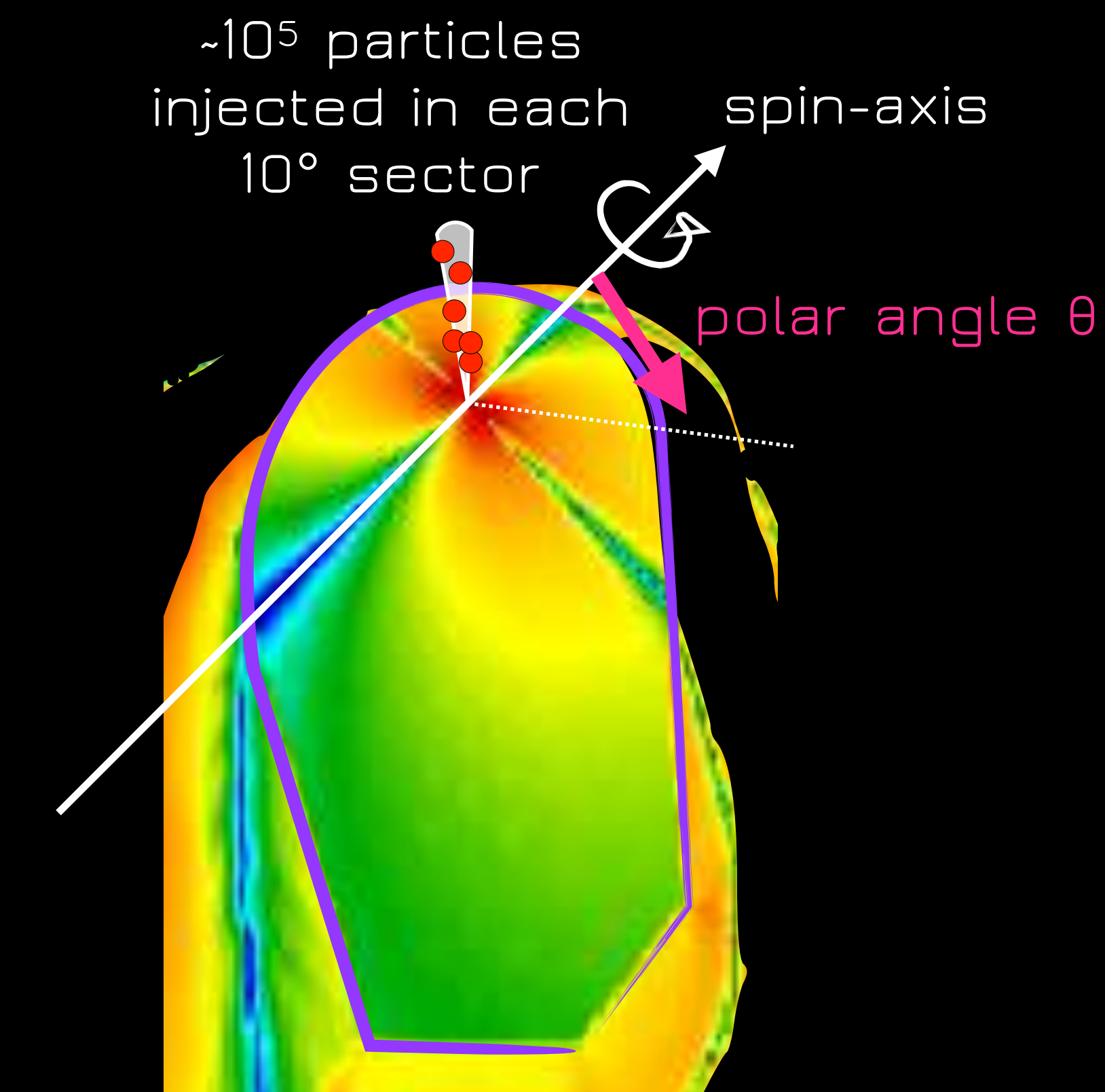
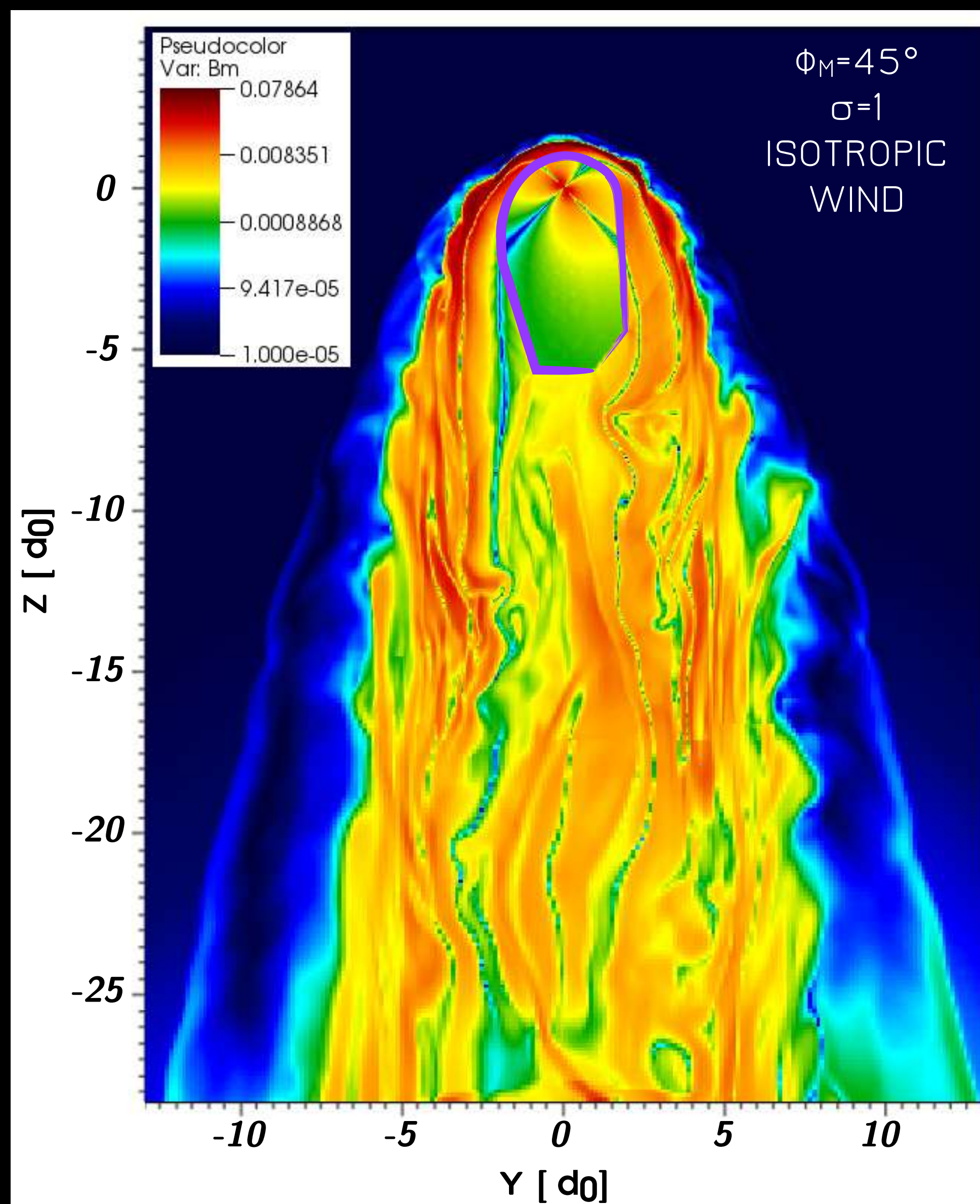
Class A project @



[Olmi & Bucciantini 2019-I/II]



# INJECTING PARTICLES IN THE PULSAR WIND



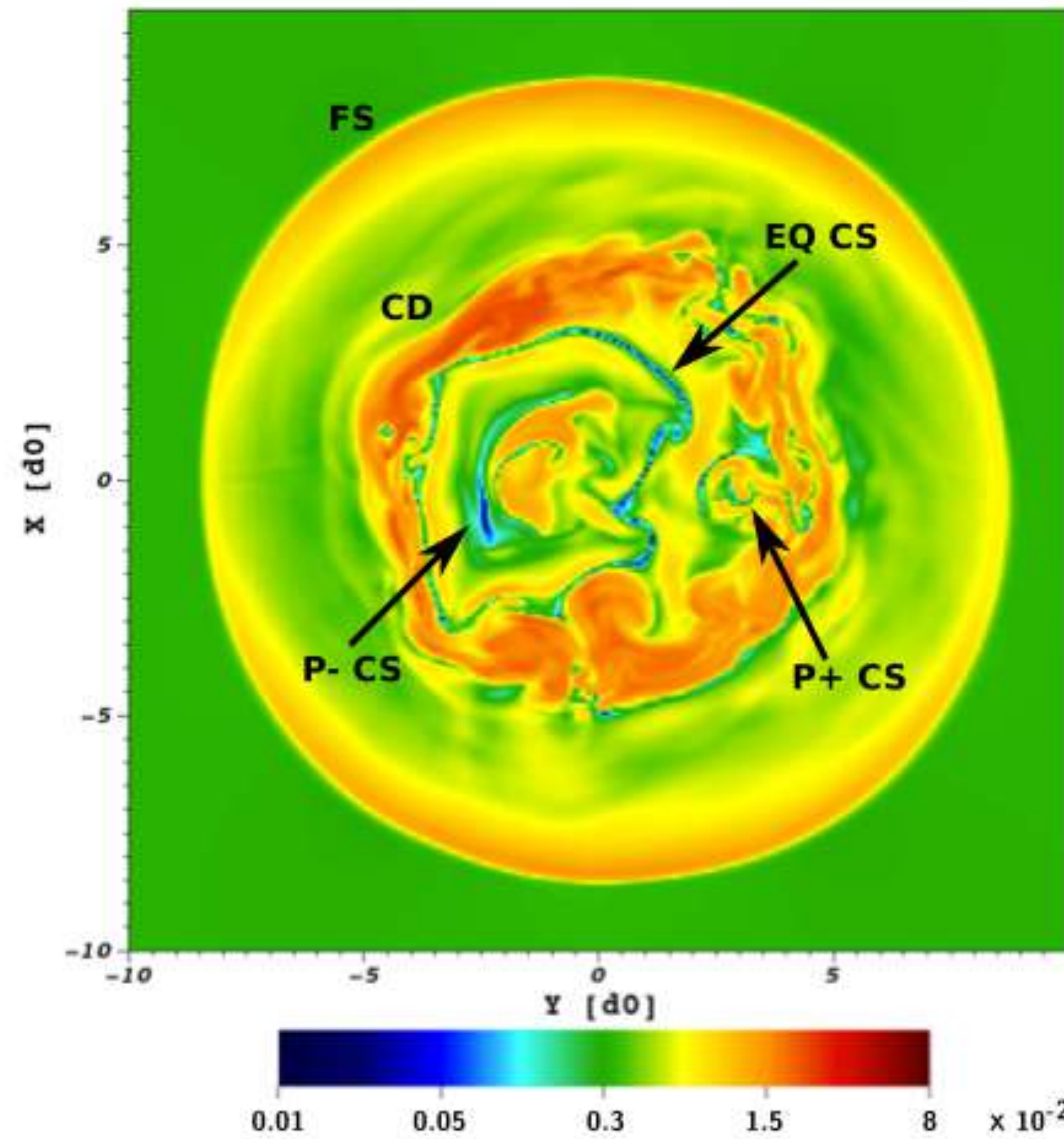
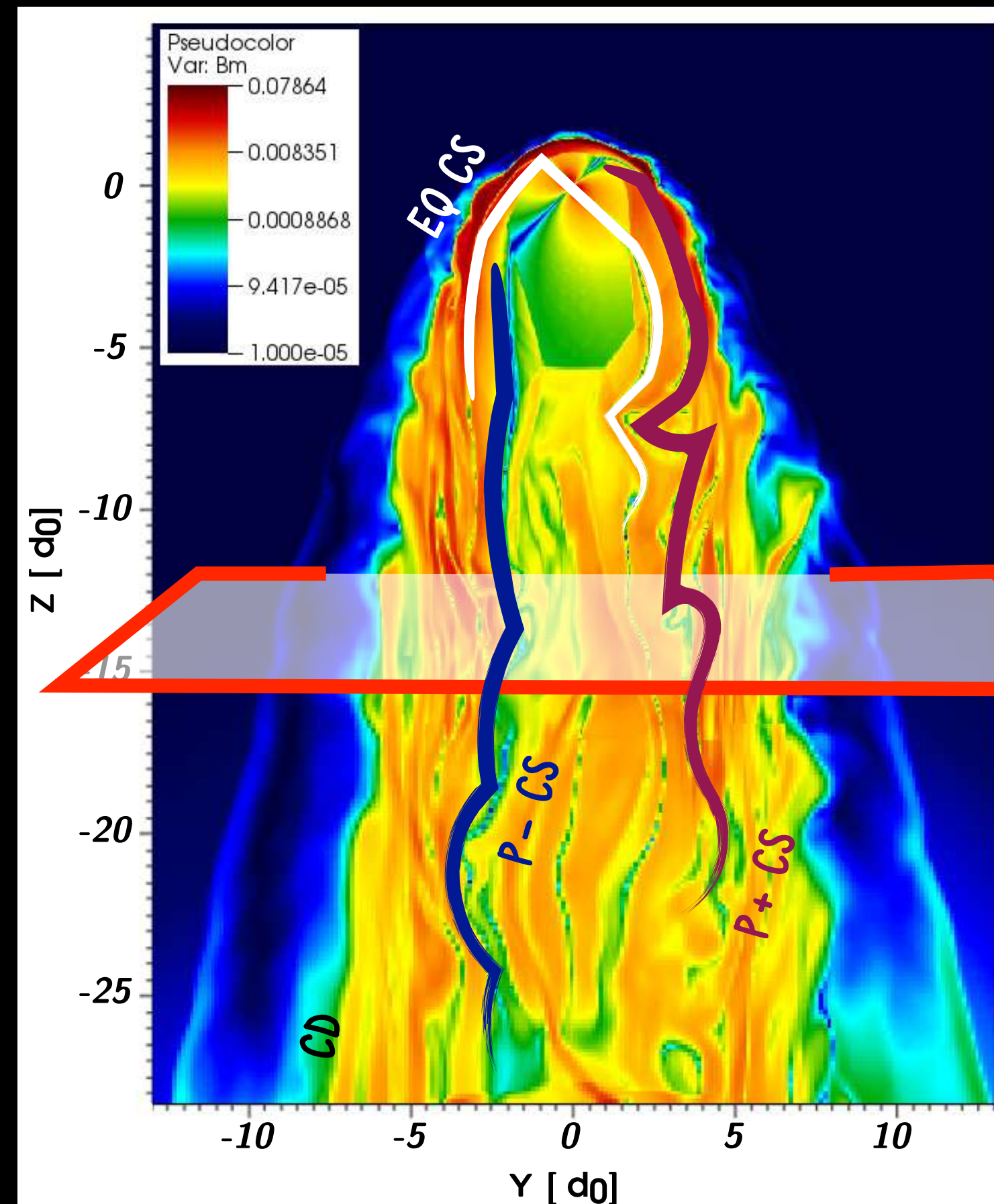
Electrons and positrons of  
different energies  
(1~50 TeV):

$$\gamma = E/(m_e c^2) = [0.5, 1.0, 3.0, 10] \times 10^7$$



# THE ROLE OF CURRENT SHEETS IN THE TAIL

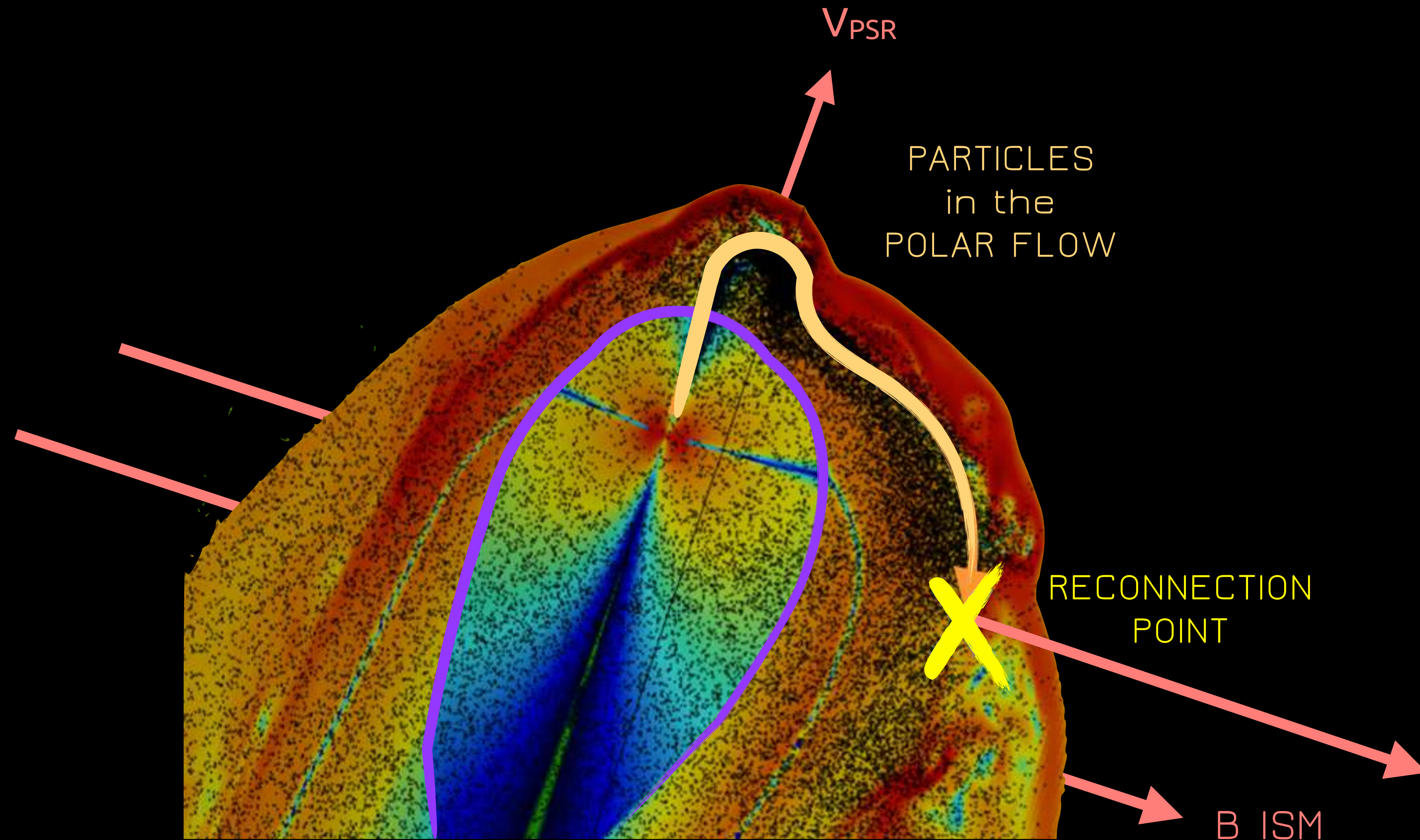
Particles remain confined in current sheets from injection





# ESCAPE OF PARTICLES: HOW AND FROM WHERE

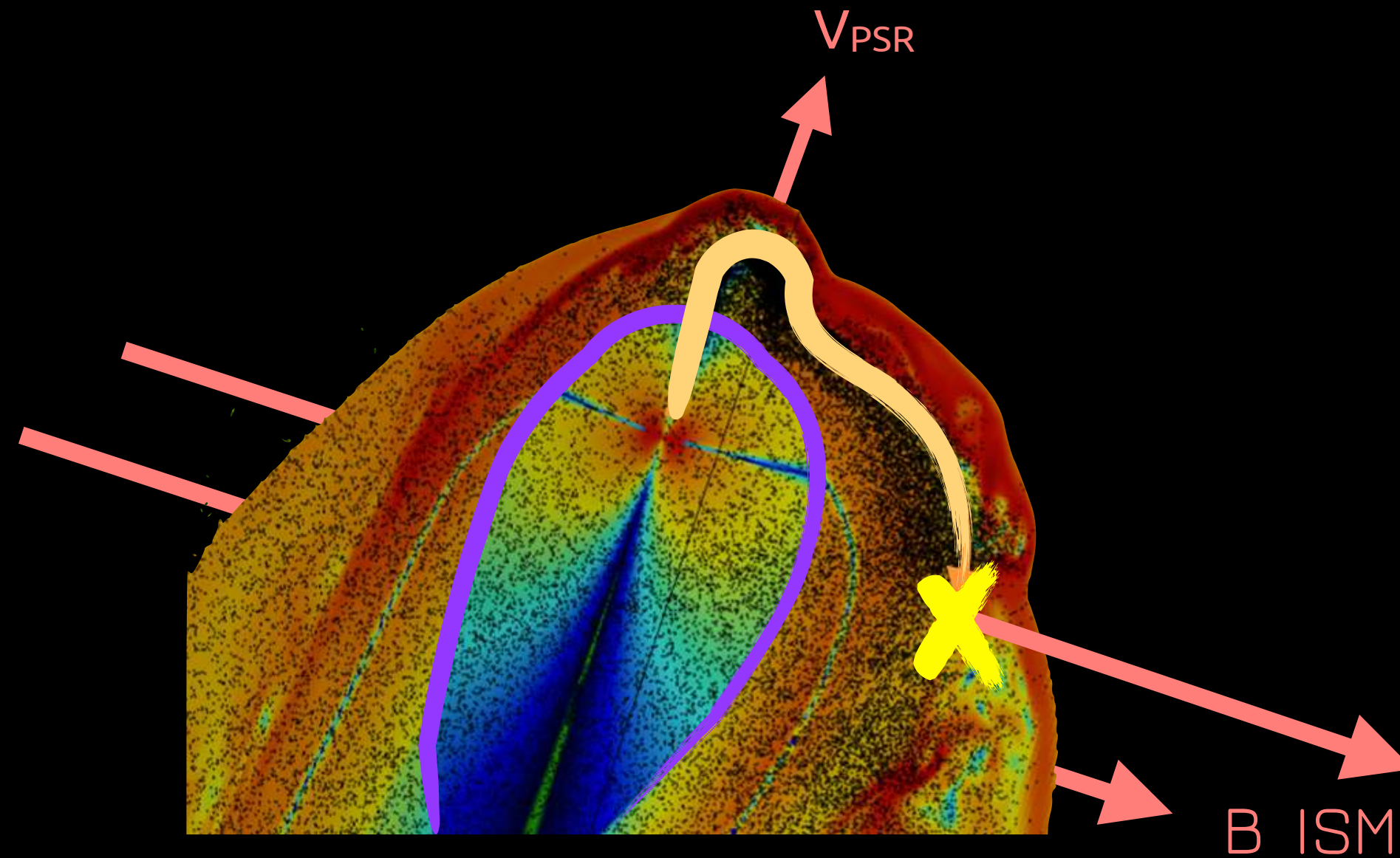
When the current sheet points towards the CD, particles can escape at reconnection points of the magnetopause





# ESCAPE OF PARTICLES: EFFICIENCY

How efficient the process is depends on the particle energy = its Larmor radius  $R_L$



Low energy particles have small  $R_L$ , are closely attached to the streamlines and can only escape directly at reconnection points.

High energy particles have large  $R_L$ , and efficiently escape in the vicinity of reconnection points.

## ENERGY of ESCAPING PARTICLES

$\gamma \sim 10^6$ 1 TeV	$\gamma \sim 10^7$ 10 TeV	$\gamma \sim 10^8$ 10ns of TeV
-----------------------------	------------------------------	-----------------------------------

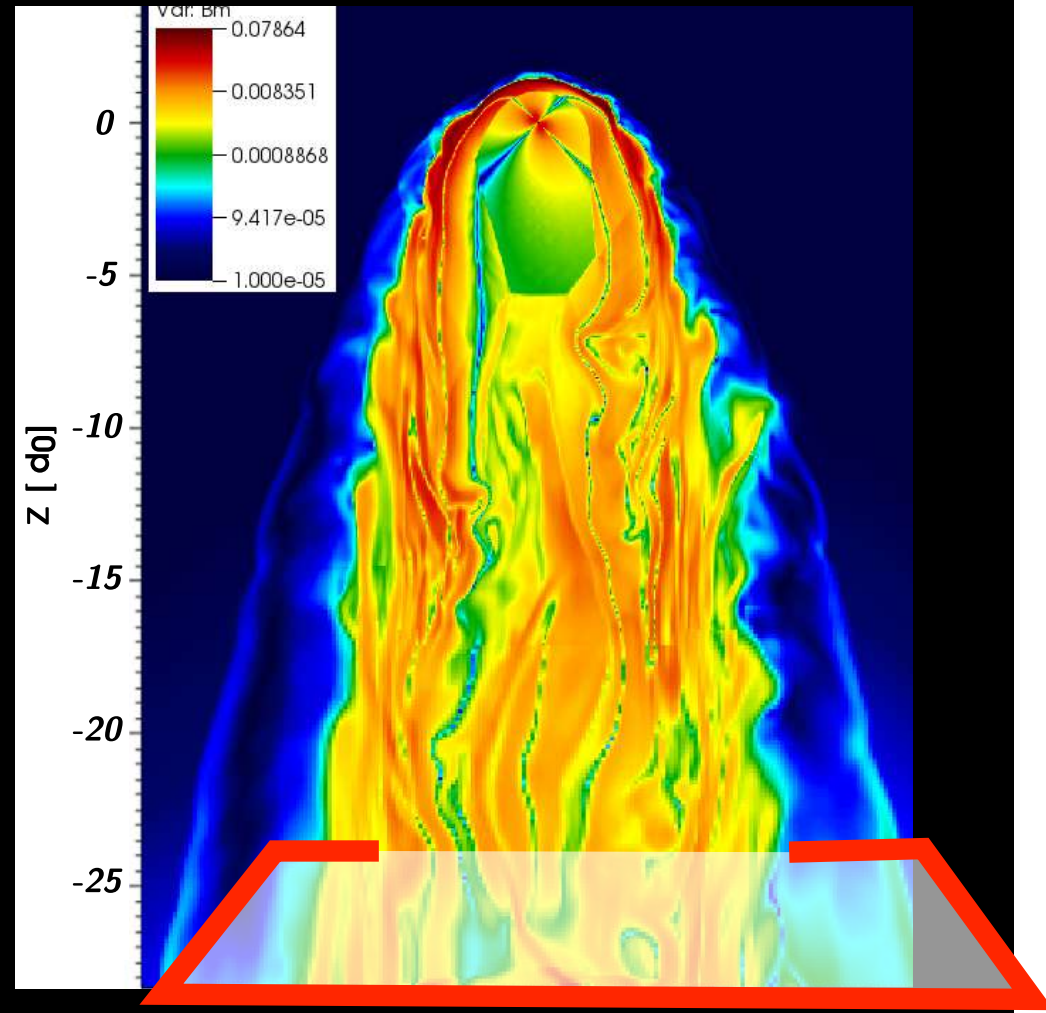


only particles carrying a large enough fraction of the maximum available energy escape efficiently:

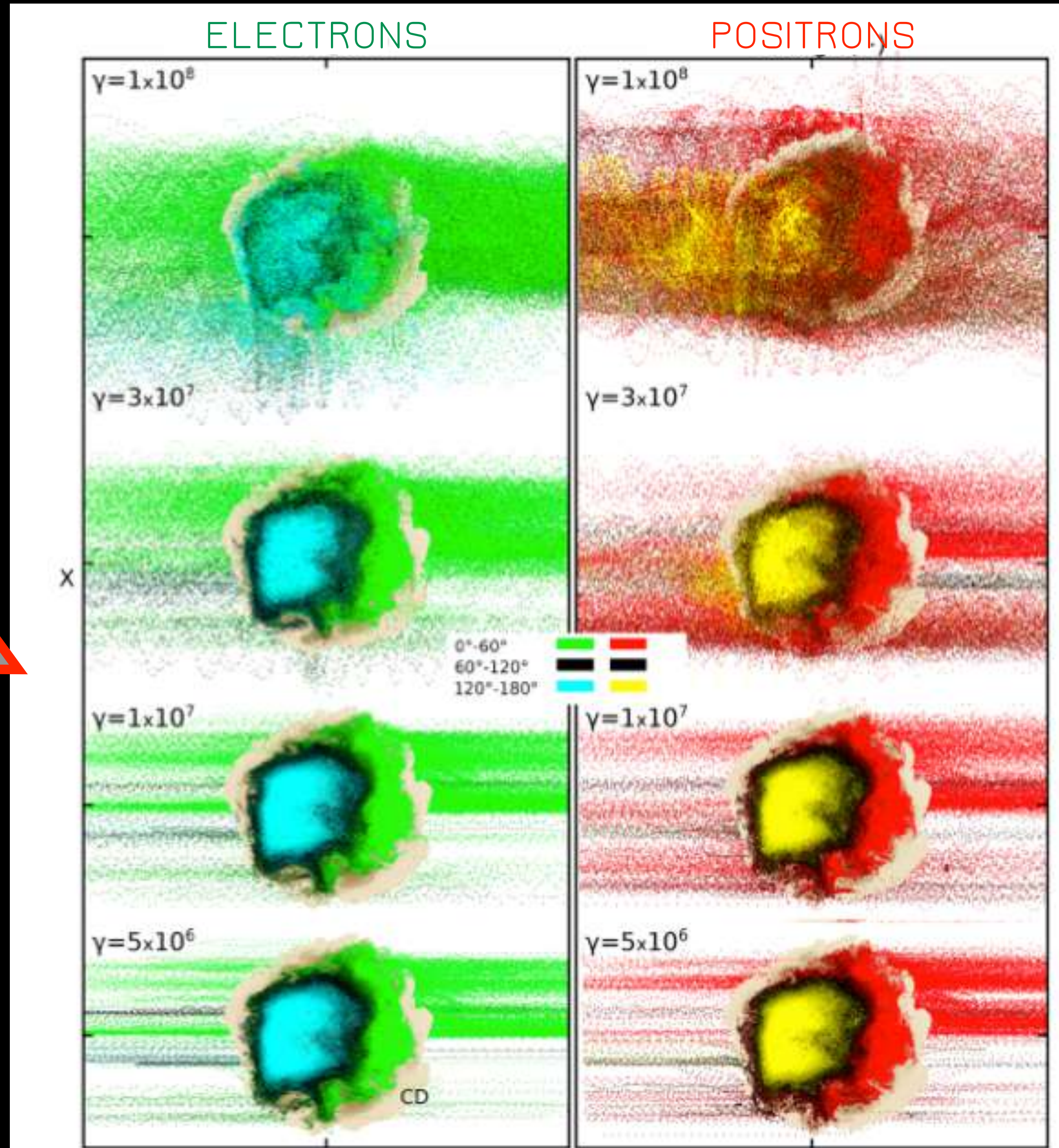
$$\gamma_{esc} \approx 10-50\% \gamma_{MPD}$$



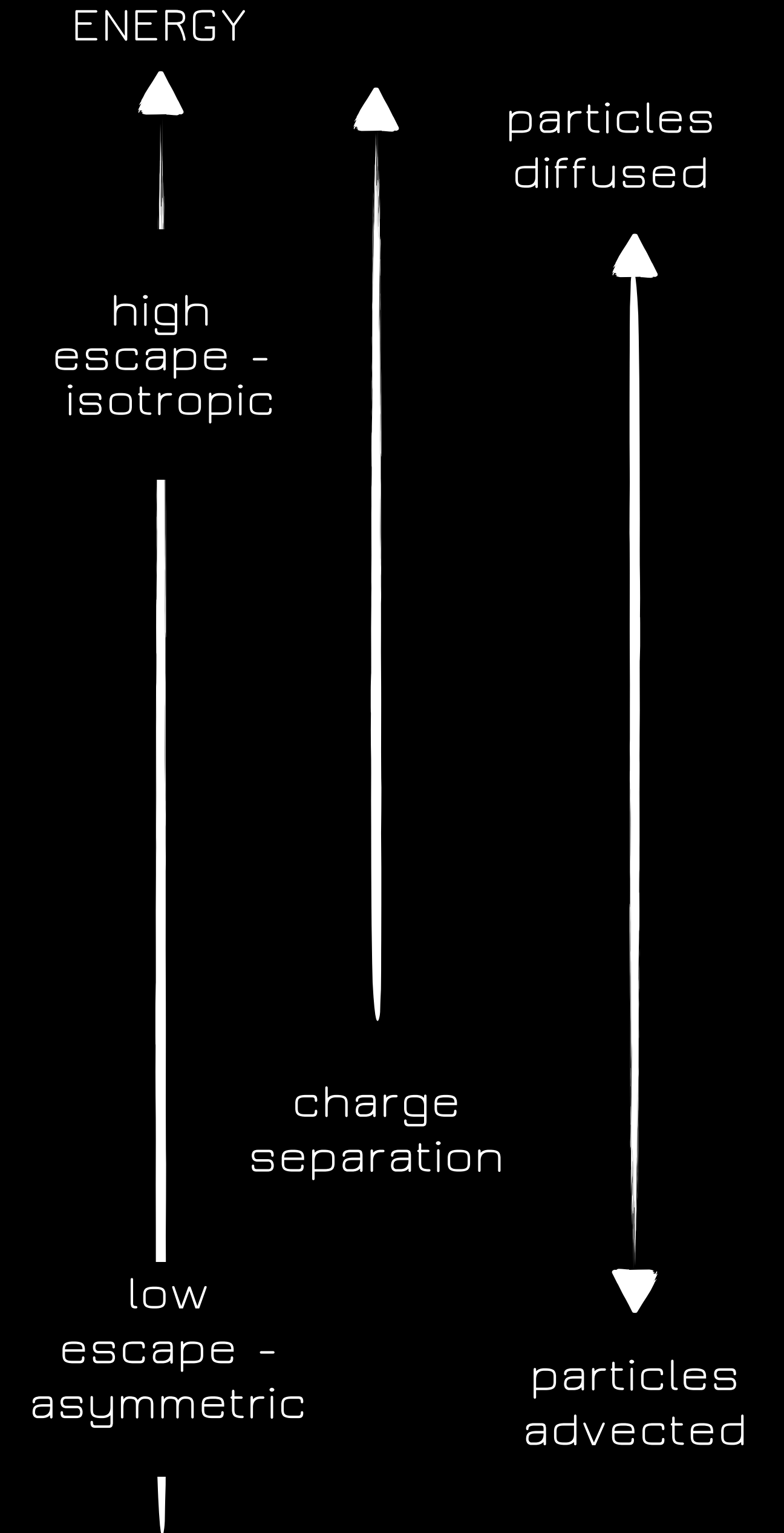
# ESCAPE OF PARTICLES: CHARGE DEPENDENCE AND ASYMMETRY



view from  
the back of  
the tail



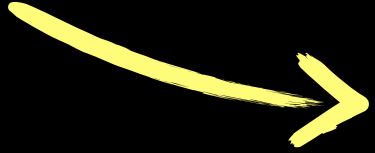
[Olmi & Bucciantini 2019-III]





# HOW TO AMPLIFY THE FIELD?

An efficient way is to excite the non-resonant Bell instability, for which you need:

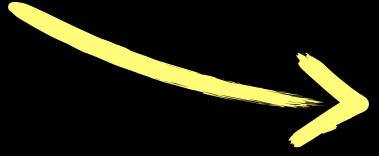
- ① An electric current  $J$   
 Charge separation in the escaping flow makes it possible!
- ② The escaping particles to carry an energy density larger than the magnetic energy density of the ambient medium:

$$n_{\text{esc}} mc^2 (\gamma_{\text{esc}} - 1) \geq B_0^2 / (8\pi)$$



# HOW TO AMPLIFY THE FIELD?

An efficient way is to excite the non-resonant Bell instability, for which you need:

- ① An electric current  $J$   
 Charge separation in the escaping flow makes it possible!
- ② The escaping particles to carry an energy density larger than the magnetic energy density of the ambient medium:

$n_{\text{esc}} mc^2 (\gamma_{\text{esc}} - 1) \geq B_0^2 / (8\pi)$   This condition is satisfied if a high enough fraction of the pulsar luminosity ( $\epsilon$ ) goes into escaping particles:

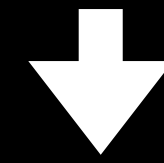
$$\epsilon \gtrsim 10^{-6} \dot{E}$$



# CAN WE REACH THE REQUIRED AMPLIFICATION?

The observed synchrotron emission in the feature required a factor  $\sim x5-10$  of amplification of the ambient field, that must be reached when the instability saturates:

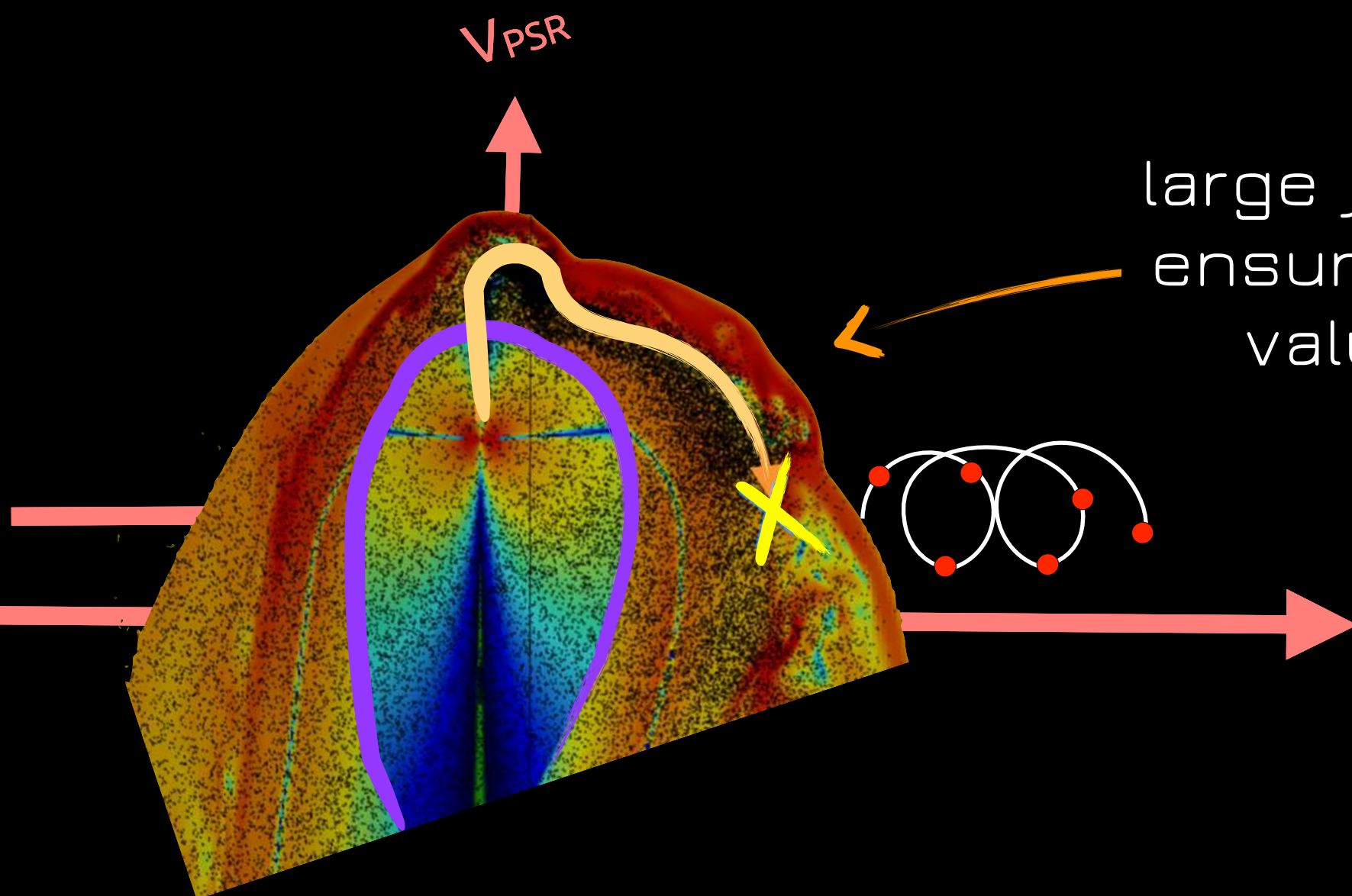
$$\text{saturation scale } \Delta x \approx R_{\text{gyro},\perp}(\gamma_{\text{esc}})$$



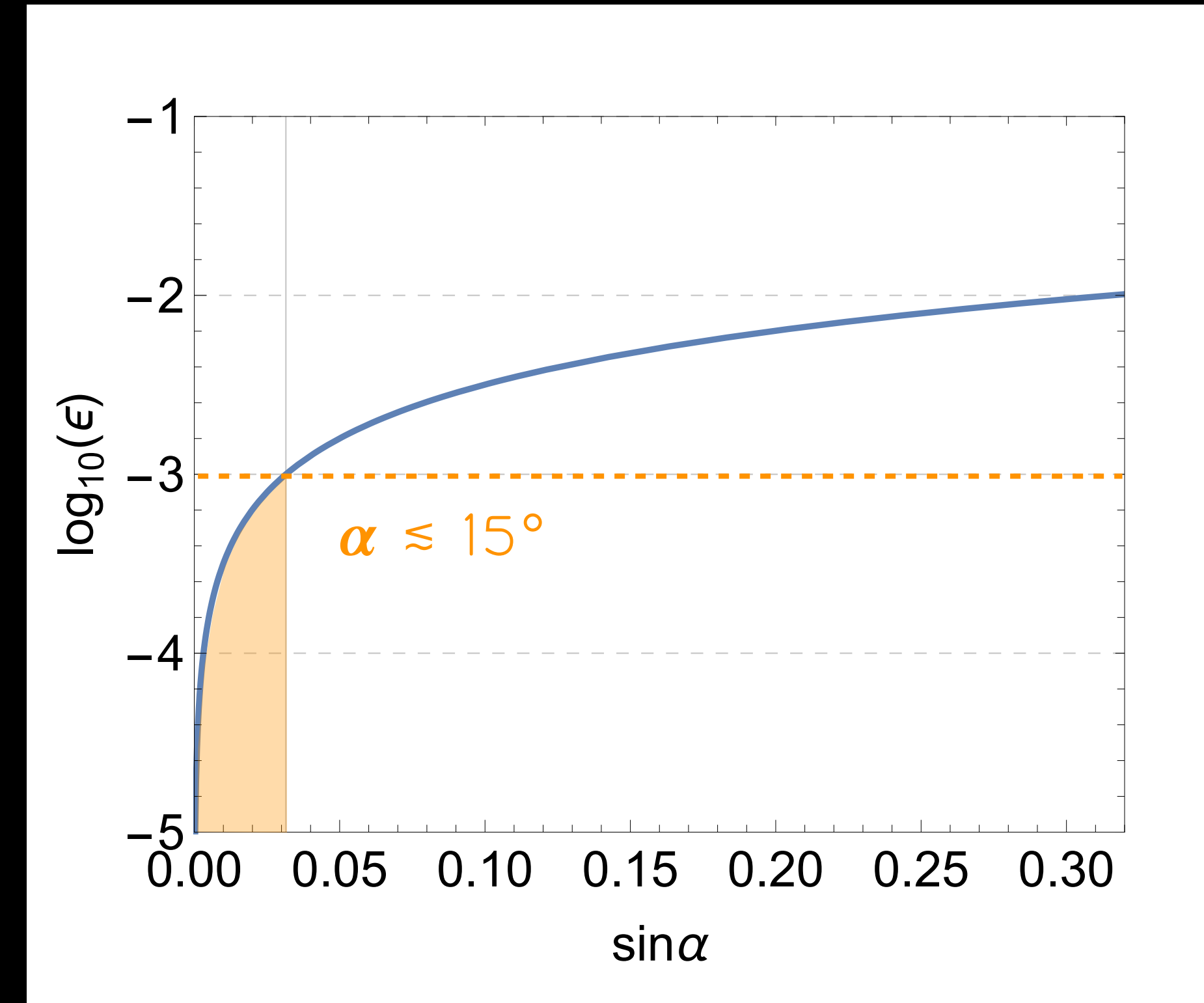
$$\left(\frac{\Delta B}{B_{\text{ISM}}}\right)_{\text{sat}} = \frac{2\sqrt{\epsilon}}{\sin \alpha} \left(\frac{\gamma_{\text{MPD}}}{\gamma_{\text{esc}}}\right)$$

$\alpha \rightarrow$  particle pitch angle

$\epsilon \rightarrow$  fraction of PSR luminosity going in the feature



large jump in B field ( $70 \rightarrow 3 \mu\text{G}$ ) ensures anisotropy and small values of escaping pitch angles



[Bandiera, Olmi, Blasi, Amato, in preparation]



# CONCLUSIONS

Evolved PWNe do show evident leakage of particles in the ambient medium.

In the last ~10 yrs the number of misaligned X-ray features detected increased → are them a universal feature of bow shock nebulae?

Understanding the nature of these features is fundamental to understand how PWNe load the Galaxy with particles (energy, number, charge, etc) and have important consequences on many open problems (CR physics, nature of TeV halos, PeVatrons)

