

Dark Matter Capture, Thermalisation and Annihilation in Neutron Stars

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in collaboration with

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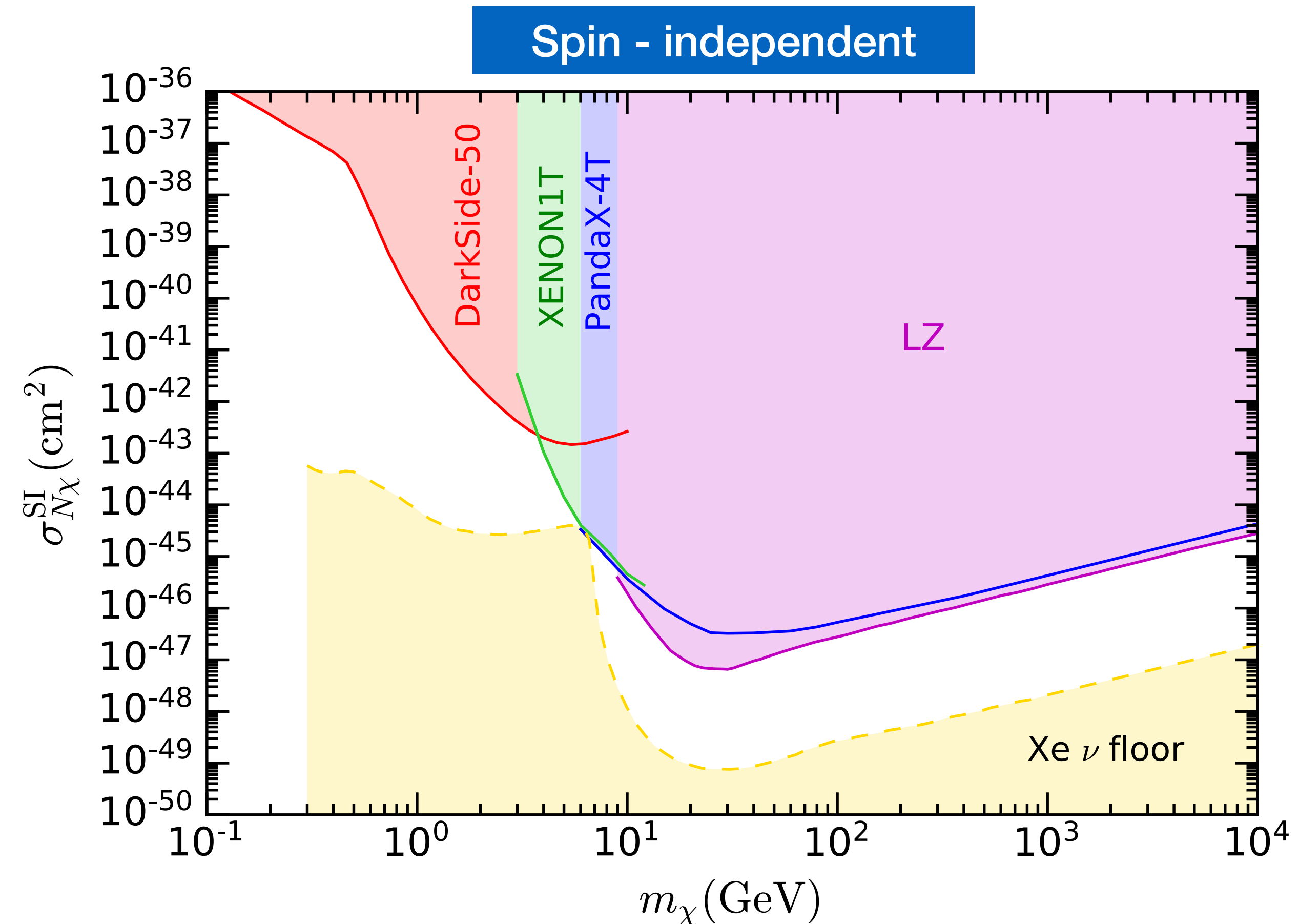
[arXiv:2004.14888](#) (JCAP), [arXiv:2010.13257](#) (JCAP), [arXiv:2012.08918](#) (PRL), [arXiv: 2108.02525](#) (JCAP), [arXiv:23XX.XXXX](#)



Introduction

Direct Detection

- Stringent constraints on spin-independent (SI) interactions.
- Restricted by
 - ➔ Nuclear mass of the target
 - ➔ Recoil threshold
- Less sensitivity to interactions with **momentum or velocity suppressed cross sections**.

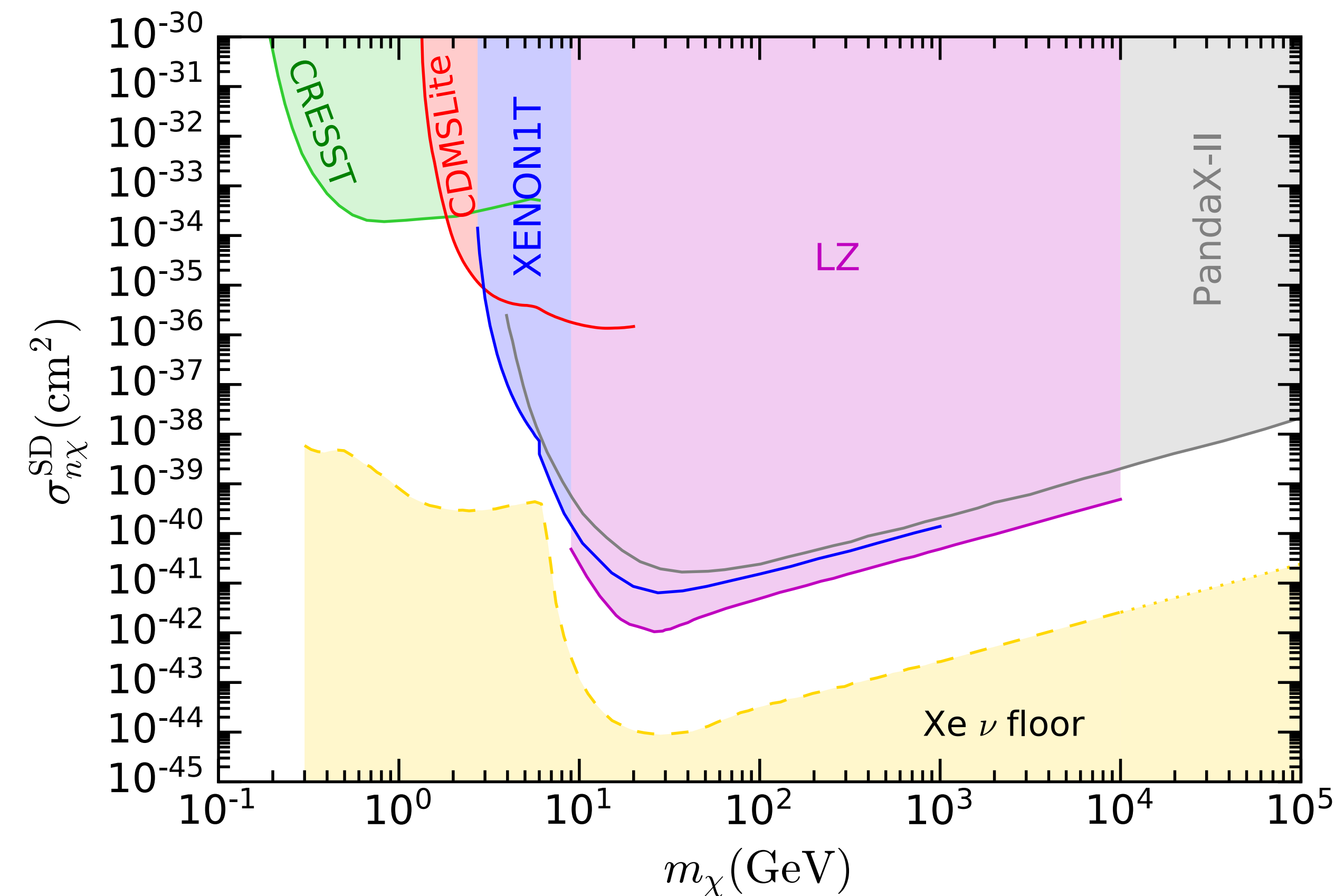


Introduction

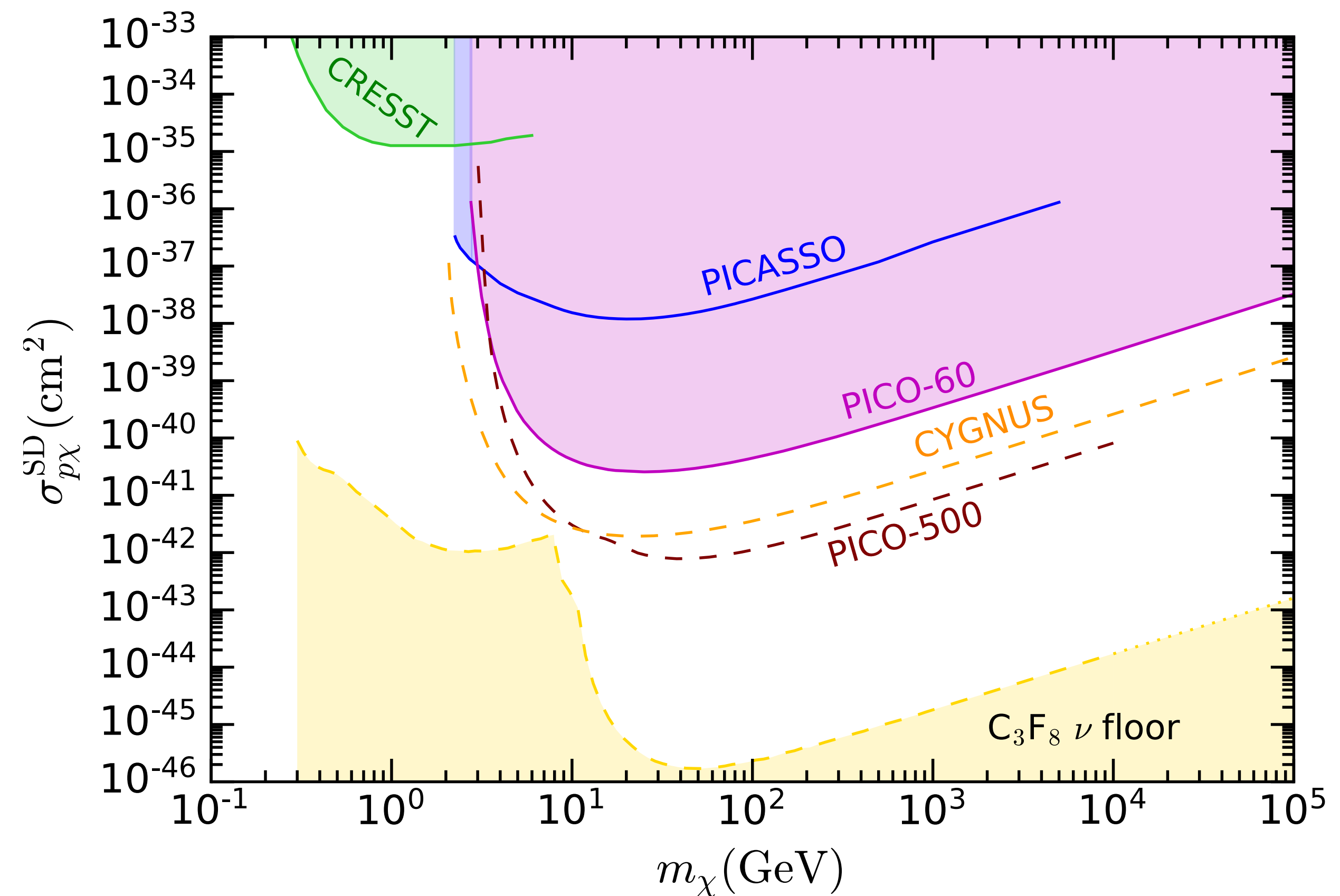
Direct Detection

- Much weaker sensitivity to spin-dependent interactions.

SD neutron



SD proton





DM Capture in the Sun

- DM scatters, loses energy, becomes gravitationally bound to the Sun. [Gould 1987](#)
- Accumulates and annihilates in the centre of the Sun.
- In equilibrium, annihilation rate proportional to the **DM-nucleon scattering cross section**.
- Neutrinos from DM annihilation can be detected in the Earth (Super-Kamiokande, Antares, IceCube).

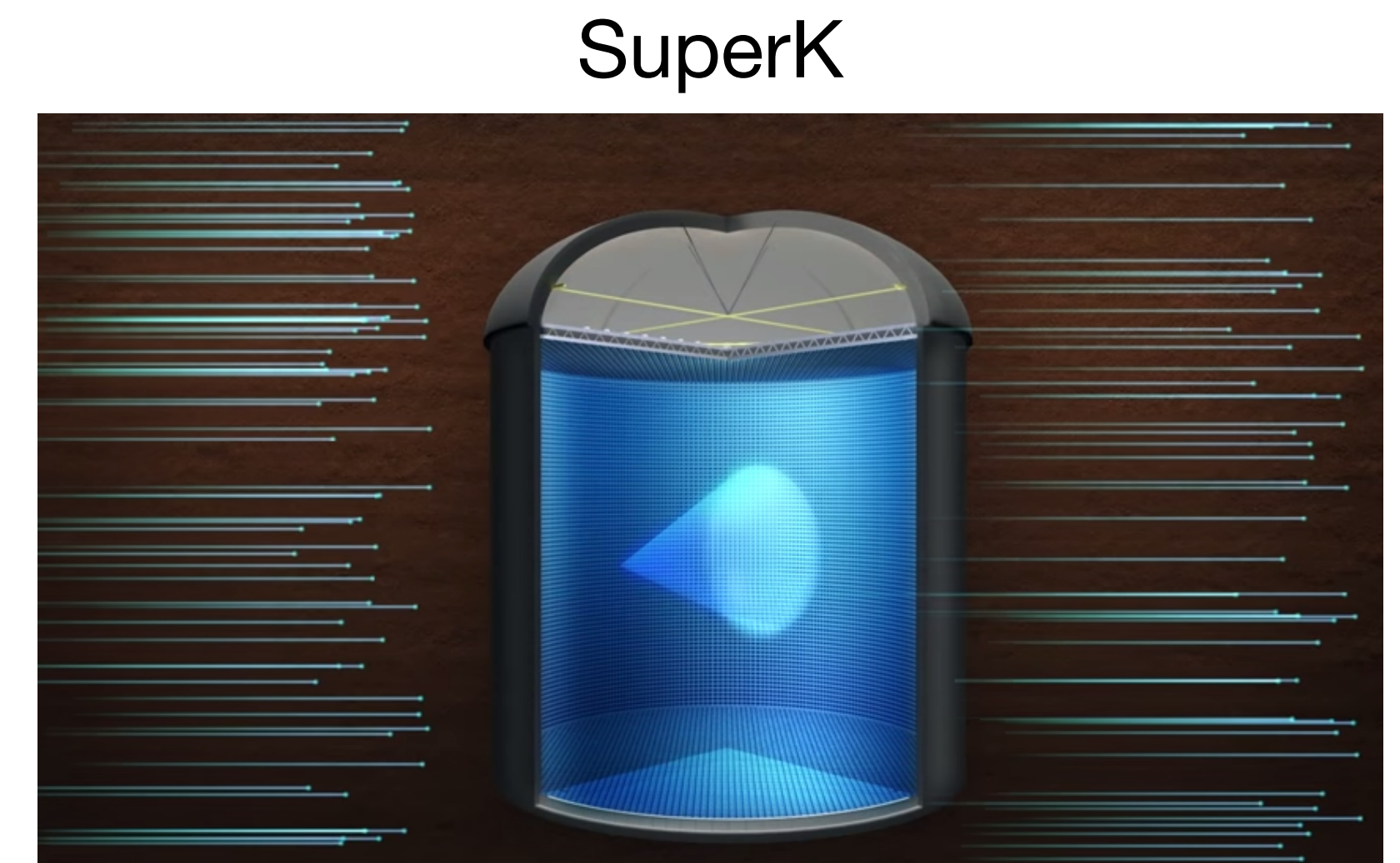
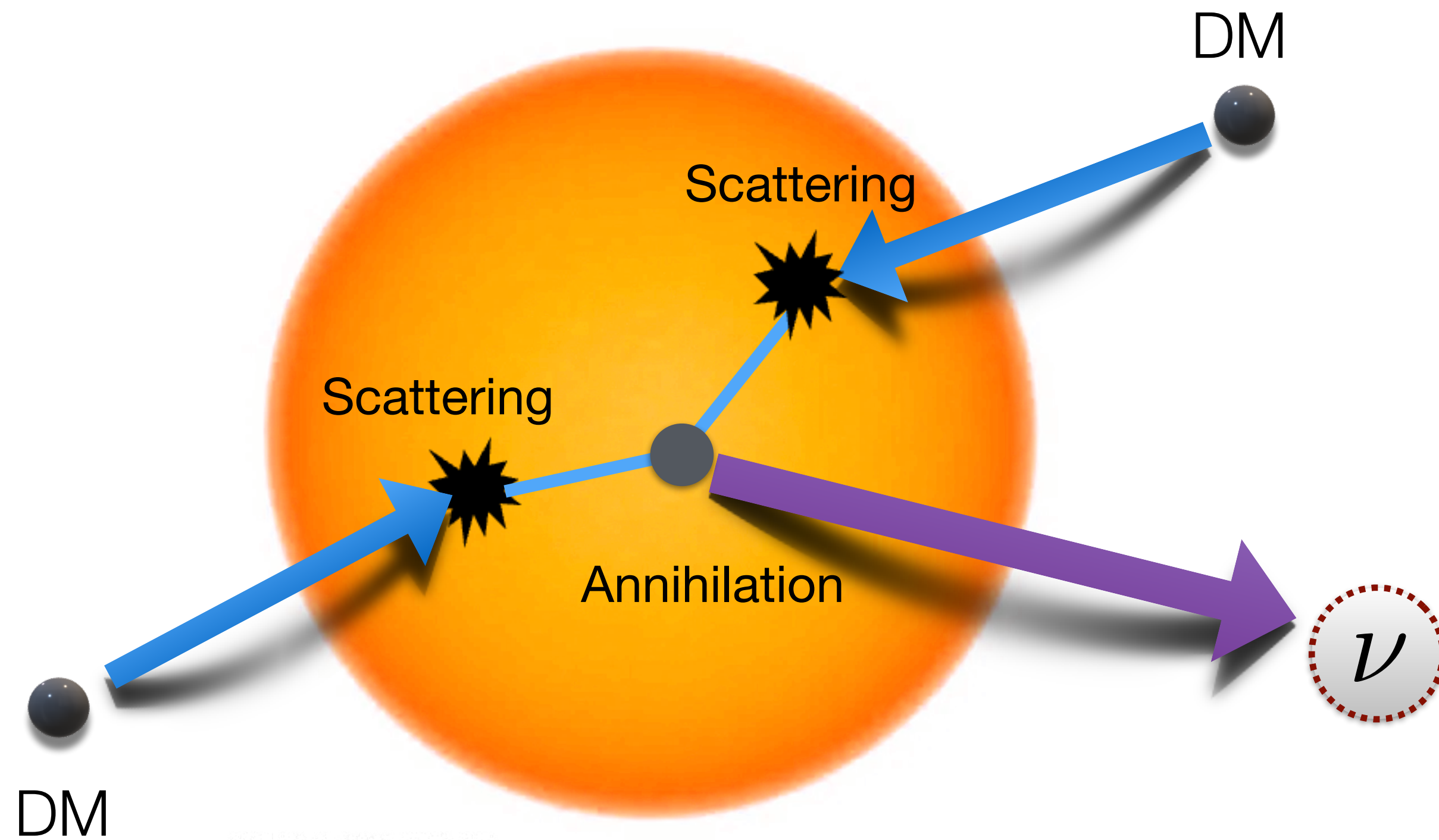
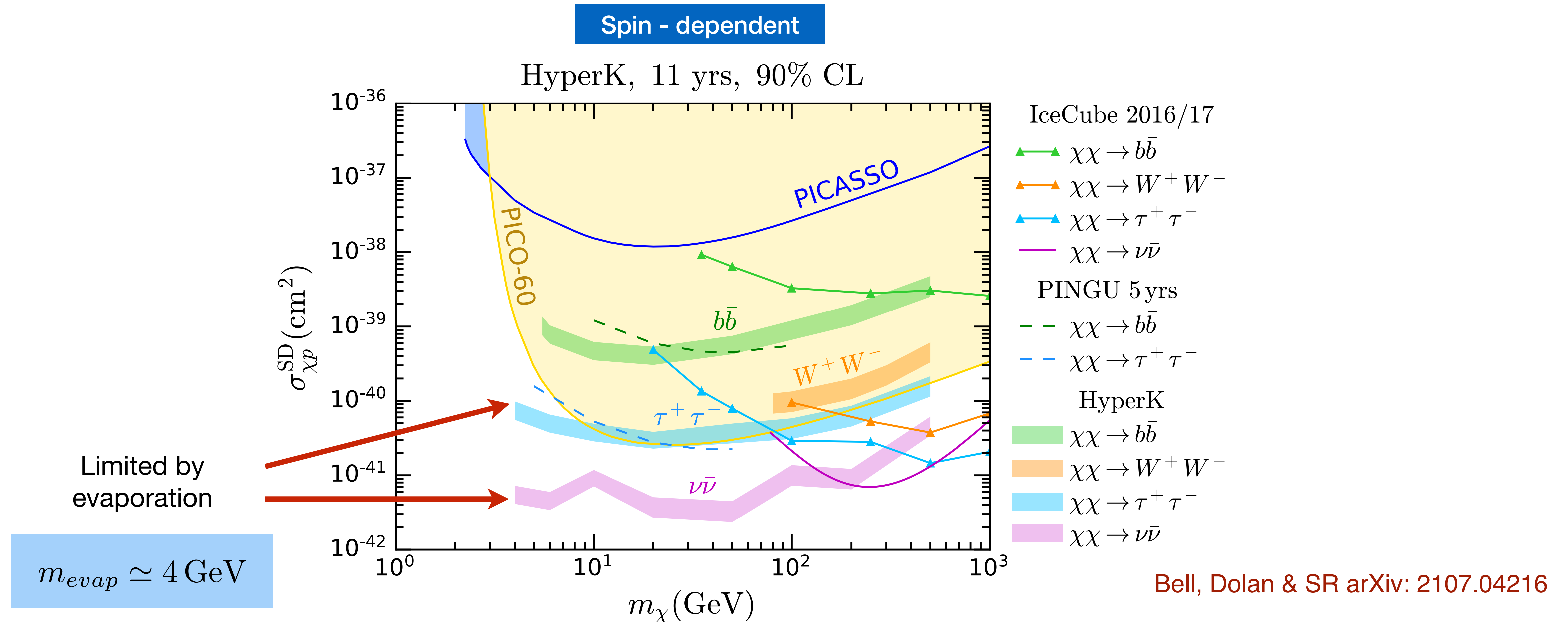


Image credit: Institute for Cosmic Ray Research, The University of Tokyo

Captured DM annihilating in the Sun

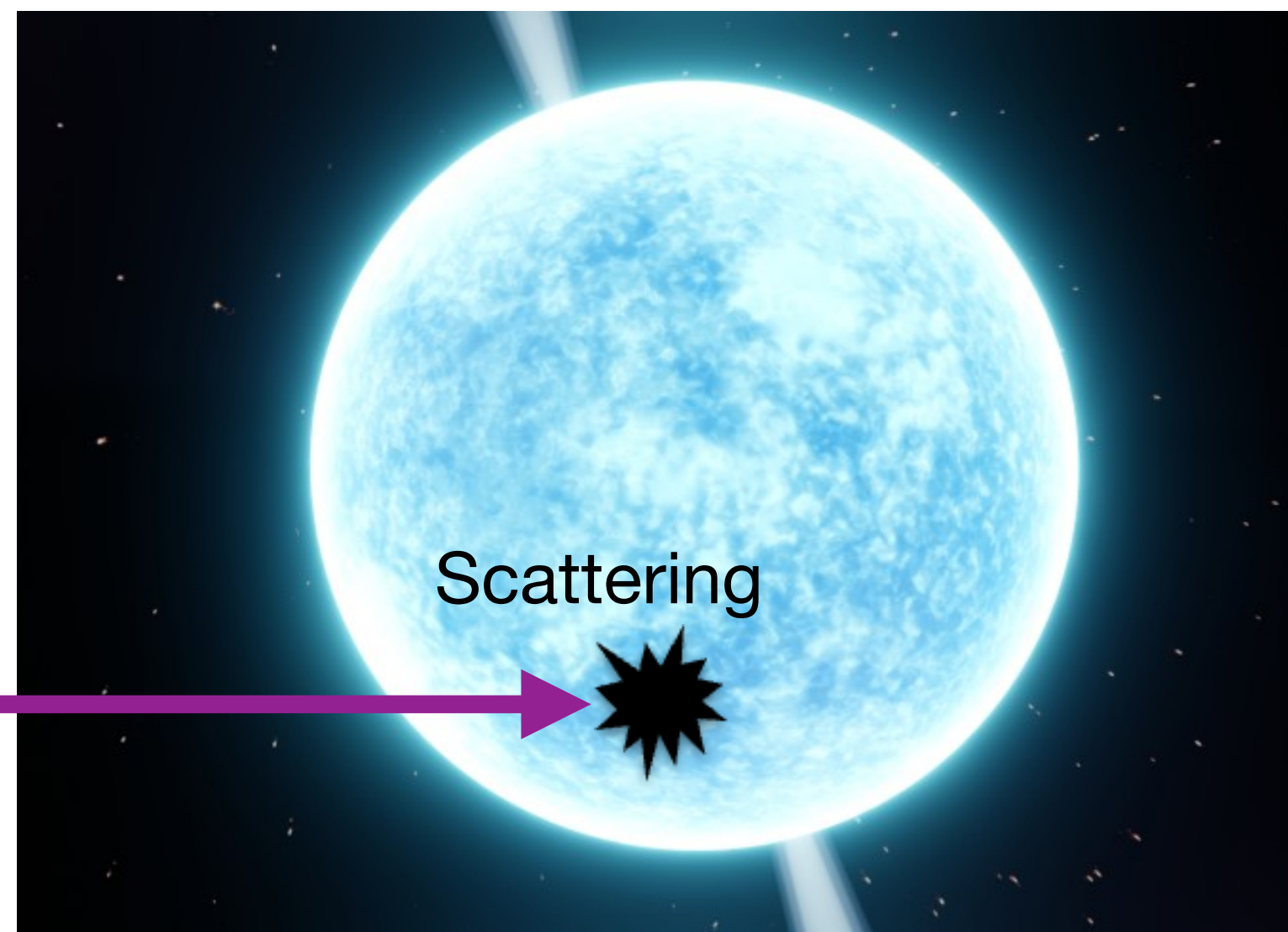
- Limits on the SI cross section from DM annihilation to neutrinos **much weaker than DD**.



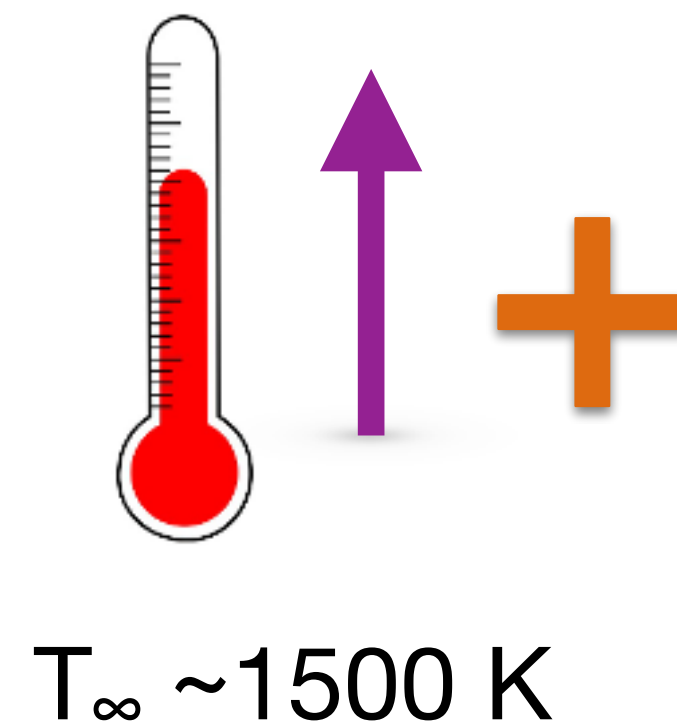
DM capture in Neutron Stars

- Extremely efficient at capturing DM, capture probability order 1 for $\sigma_{n\chi} \sim \mathcal{O}(10^{-45} - 10^{-44} \text{cm}^2)$
 - Capture plus subsequent annihilation can heat up local NSs (10 pc) [Baryakhtar et al. arXiv:1704.01577 \(PRL\)](#)
- ➡ Possible detection with JWST NIRCam (SNR = 5) [Chatterjee et al. arXiv: 2205.05048](#)

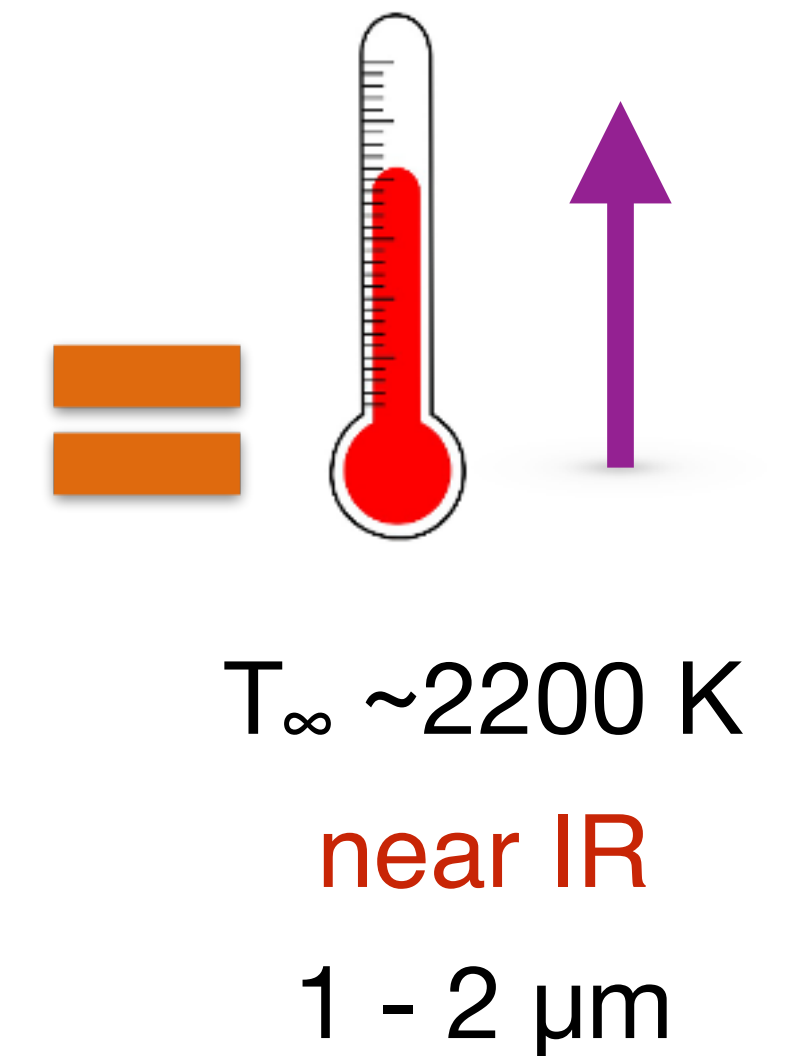
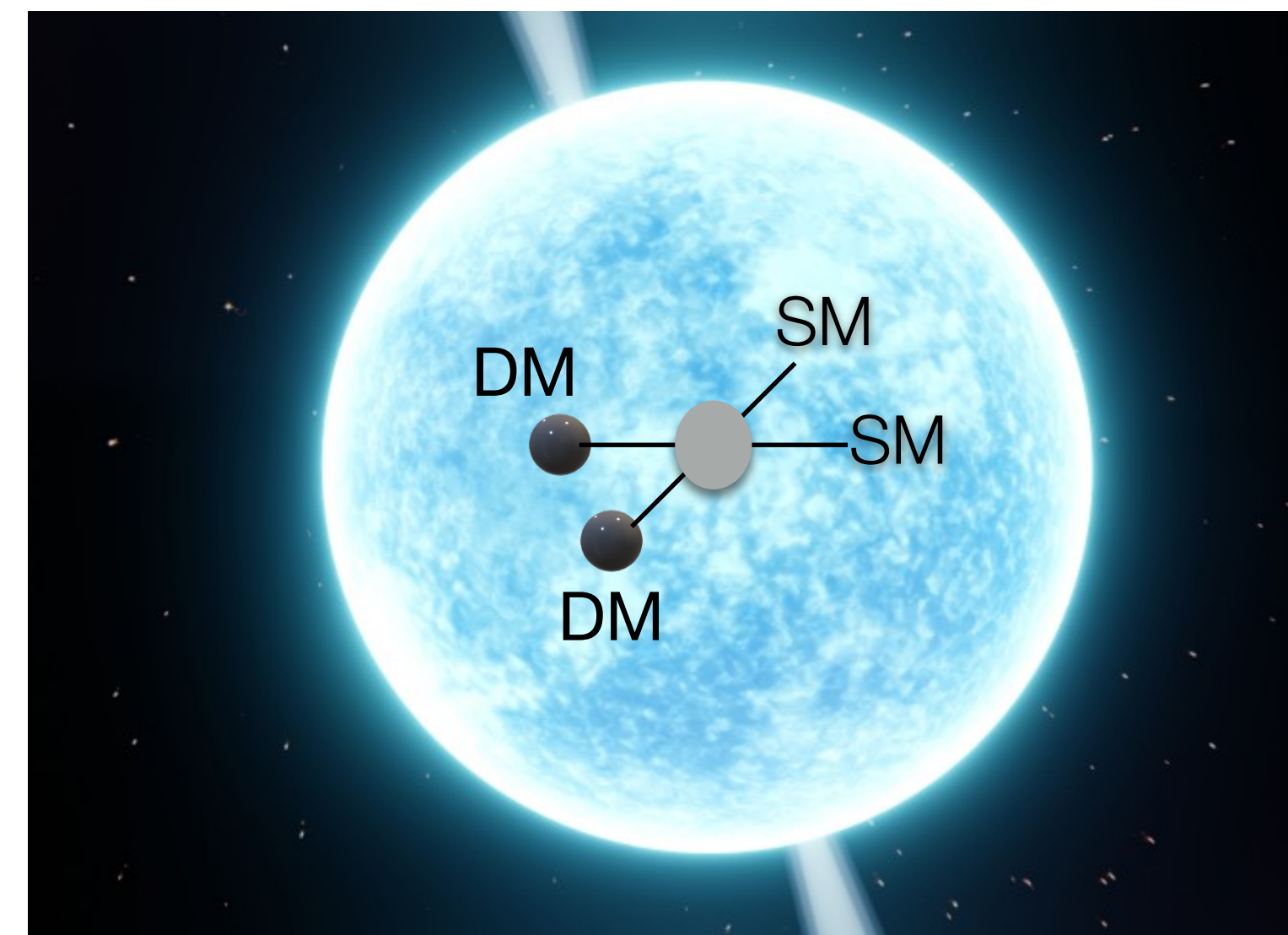
DM capture



Nearby,
isolated NS



DM annihilation



Neutrons Stars

- The densest stars known.
- Supported against collapse by **neutron** degeneracy pressure.

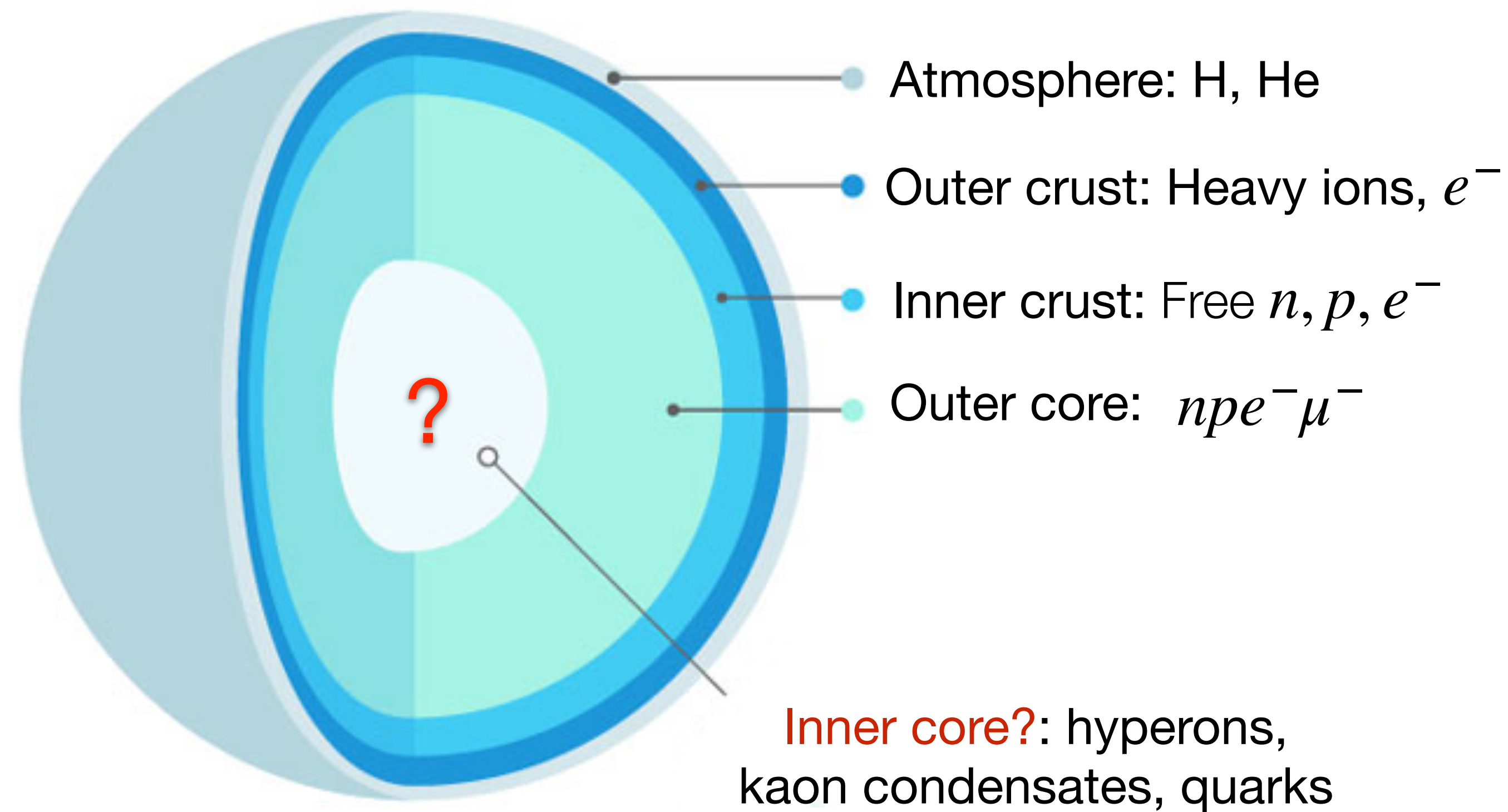


Image Credit: Feryal Özel

Equation of State (EoS)

$$P = P(\rho)$$

➔ APR

➔ BSk

➔ QMC

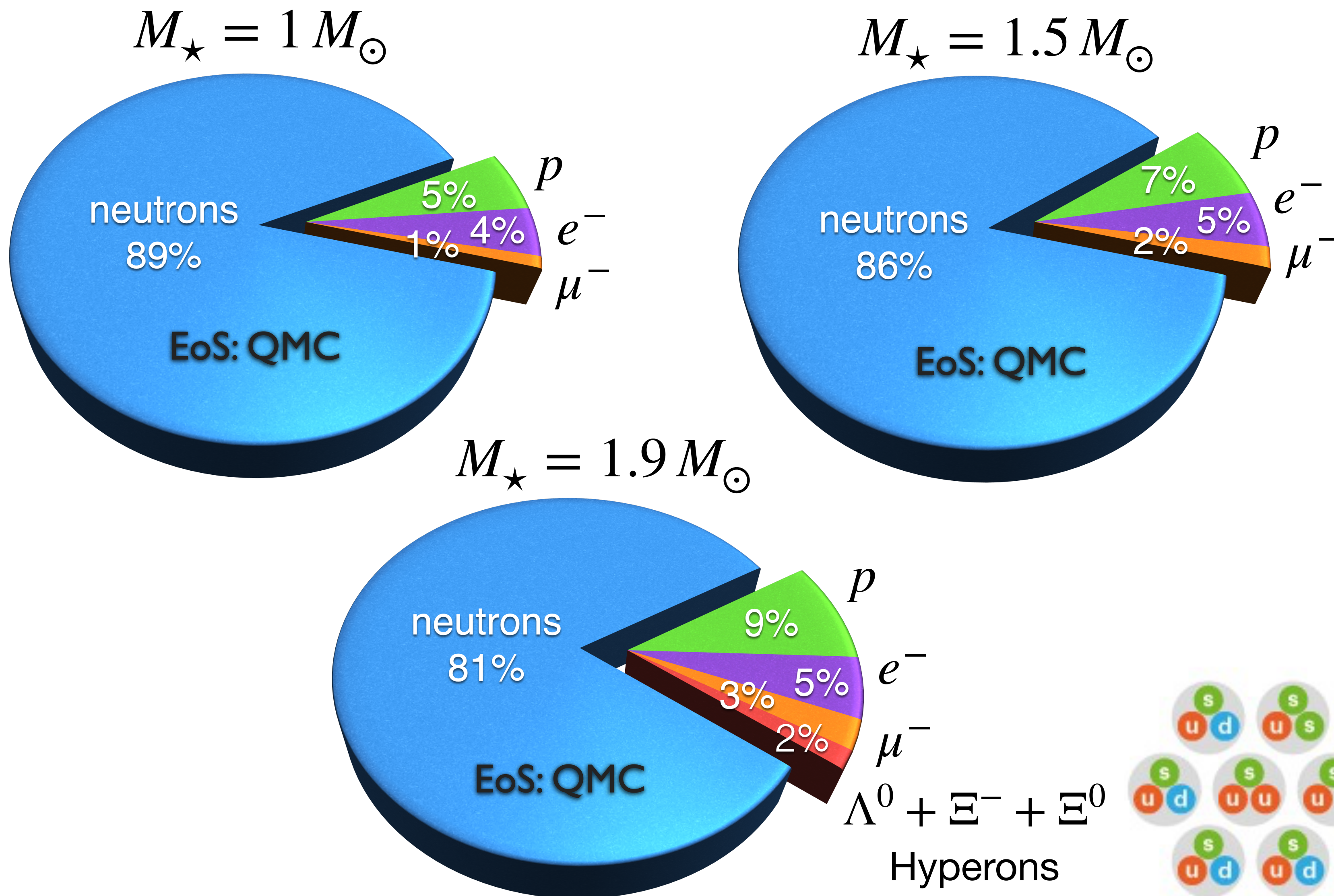
➔ SLy, ...

+

Hydrostatic Equilibrium

➔ Tolman-Oppenheimer-Volkoff (TOV) equations

DM capture in Neutron Stars



Targets

- Baryons
 - ➔ Strongly interacting
 - ➔ Pauli blocking (interacting Fermi gas)
- Leptons
 - ➔ Relativistic
 - ➔ Pauli blocking (free Fermi gas)

DM capture in Neutron Stars

Scattering off a Fermi gas of interacting baryons

- Different kinematic regime from DM capture in the Sun.

➡ DM accelerated to quasi-relativistic speeds

➡ TOV equations and Schwarzschild metric

$$B(r) \sim 1 - v_{esc}^2(r)$$

Capture rate

$$C = \frac{\rho_\chi}{m_\chi} \int_0^\infty \frac{f_{MB}(u_\chi) du_\chi}{u_\chi} \int_0^{R_\star} 4\pi r^2 \frac{\sqrt{1 - B(r)}}{B(r)} \Omega^-(r)$$

Interaction rate

$$\Omega^-(r) = \frac{1}{2\pi^2} \int dt dE_i ds \frac{E_i}{m_\chi} \sqrt{\frac{B(r)}{1 - B(r)}} \frac{s}{\beta(s, m_i^{\text{eff}}) \gamma(s, m_i^{\text{eff}})} \frac{d\sigma_{i\chi}}{d \cos \theta_{cm}} f_{\text{FD}}(E_i, r) (1 - f_{\text{FD}}(E'_i, r))$$

DM flux

Prob. to scatter to $v \leq v_{esc}$

Relativistic kinematics

Pauli Blocking target initial and final states

Bell, Busoni, SR & Virgato, arXiv: 2004.14888

DM capture in Neutron Stars

Scattering off a Fermi gas of interacting baryons

- Two important effects missing in all previous calculations:
 - ➔ Momentum transfer $\mathcal{O}(10 \text{ GeV})$
 - ➔ Momentum dependence of the hadronic matrix elements

Nucleon couplings

$$Q_0 \sim 1 \text{ GeV}$$

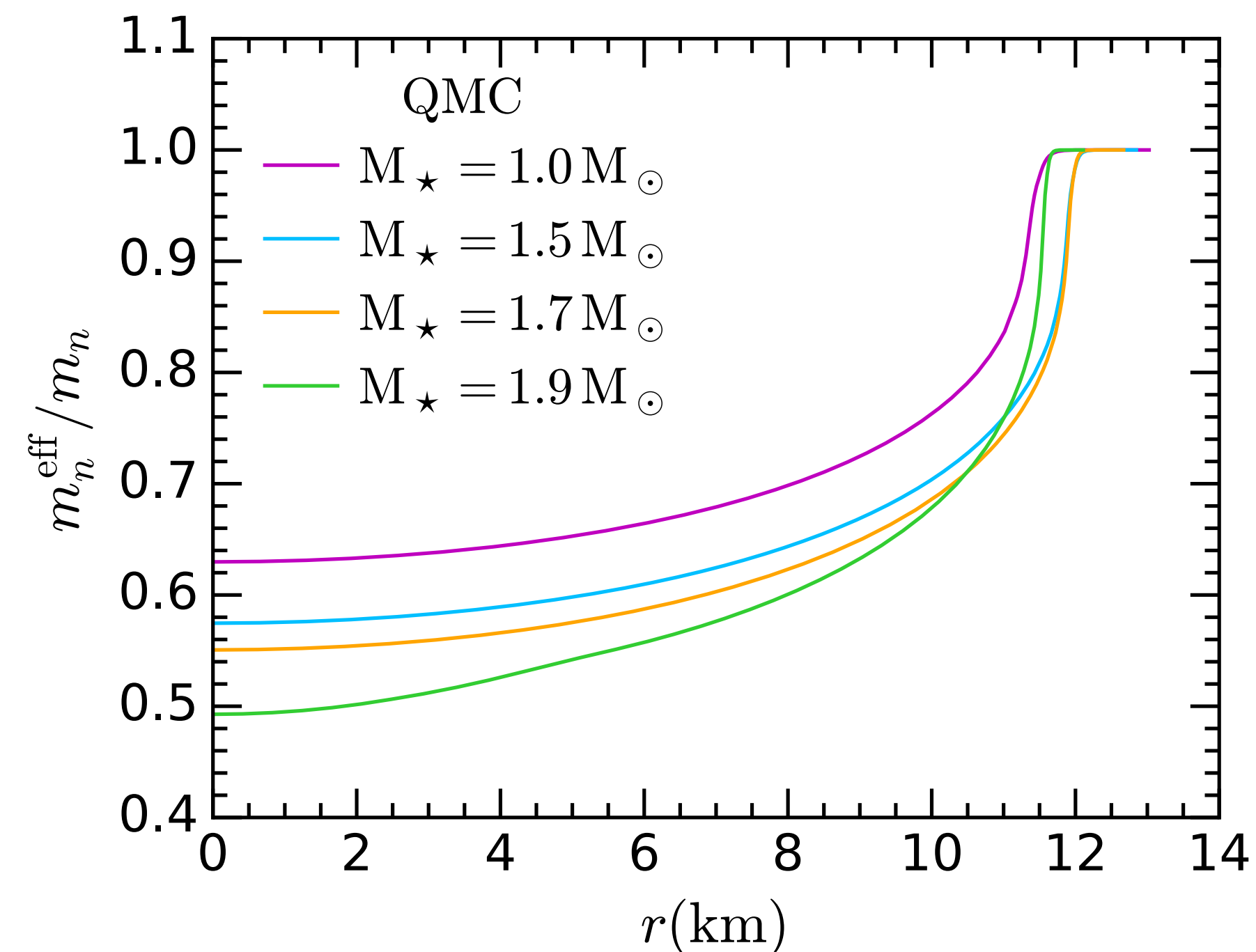
$$c_n(q) = \frac{c_n(0)}{(1 - q^2/Q_0^2)^2}$$

$$\frac{d\sigma_{n\chi}}{d \cos \theta_{cm}}(m_n^{\text{eff}}(r), c_n(q), s, t)$$

- ➔ Nucleons undergo strong interactions, **free Fermi gas is not a good approximation.**

Nucleon effective mass

$$m_n \rightarrow m_n^{\text{eff}}(r)$$



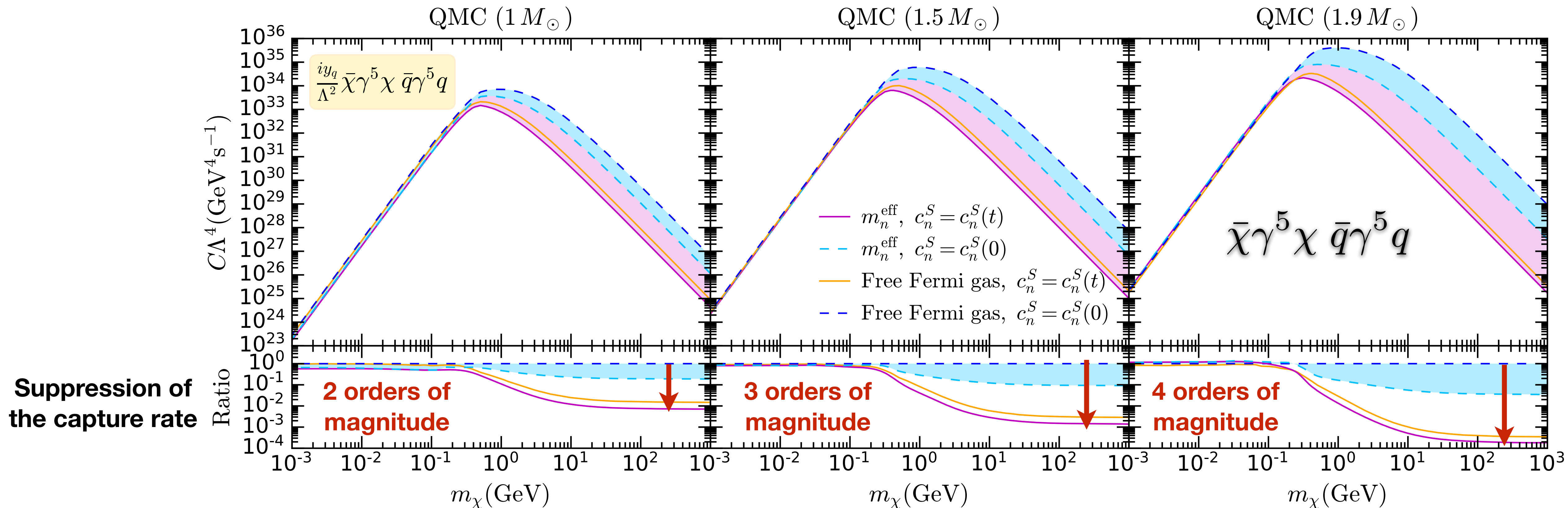
Bell, Busoni, Motta, SR, Thomas & Virgato, arXiv: 2012.08918, arXiv: 2108.02525

Scattering Operators for Fermionic DM

Operator	Coupling	Interaction	Momentum suppressed
$\bar{\chi}\chi \bar{q}q$	y_q/Λ^2	SI	✗
$\bar{\chi}\gamma^5\chi \bar{q}q$	iy_q/Λ^2	SI	✓
$\bar{\chi}\chi \bar{q}\gamma^5q$	iy_q/Λ^2	SD	✓
$\bar{\chi}\gamma^5\chi \bar{q}\gamma^5q$	y_q/Λ^2	SD	✓
$\bar{\chi}\gamma_\mu\chi \bar{q}\gamma^\mu q$	$1/\Lambda^2$	SI	✗
$\bar{\chi}\gamma_\mu\gamma^5\chi \bar{q}\gamma^\mu q$	$1/\Lambda^2$	SI, SD	✓
$\bar{\chi}\gamma_\mu\chi \bar{q}\gamma^\mu\gamma^5q$	$1/\Lambda^2$	SD	✓
$\bar{\chi}\gamma_\mu\gamma^5\chi \bar{q}\gamma^\mu\gamma^5q$	$1/\Lambda^2$	SD	✗
$\bar{\chi}\sigma_{\mu\nu}\chi \bar{q}\sigma^{\mu\nu}q$	$1/\Lambda^2$	SD	✗
$\bar{\chi}\sigma_{\mu\nu}\gamma^5\chi \bar{q}\sigma^{\mu\nu}q$	i/Λ^2	SI	✓

DM-neutron capture rate in NSs

- Accounting for nucleon structure and strong interactions **suppresses** the capture rate



Bell, Busoni, Motta, SR, Thomas & Virgato, arXiv: 2012.08918, arXiv: 2108.02525

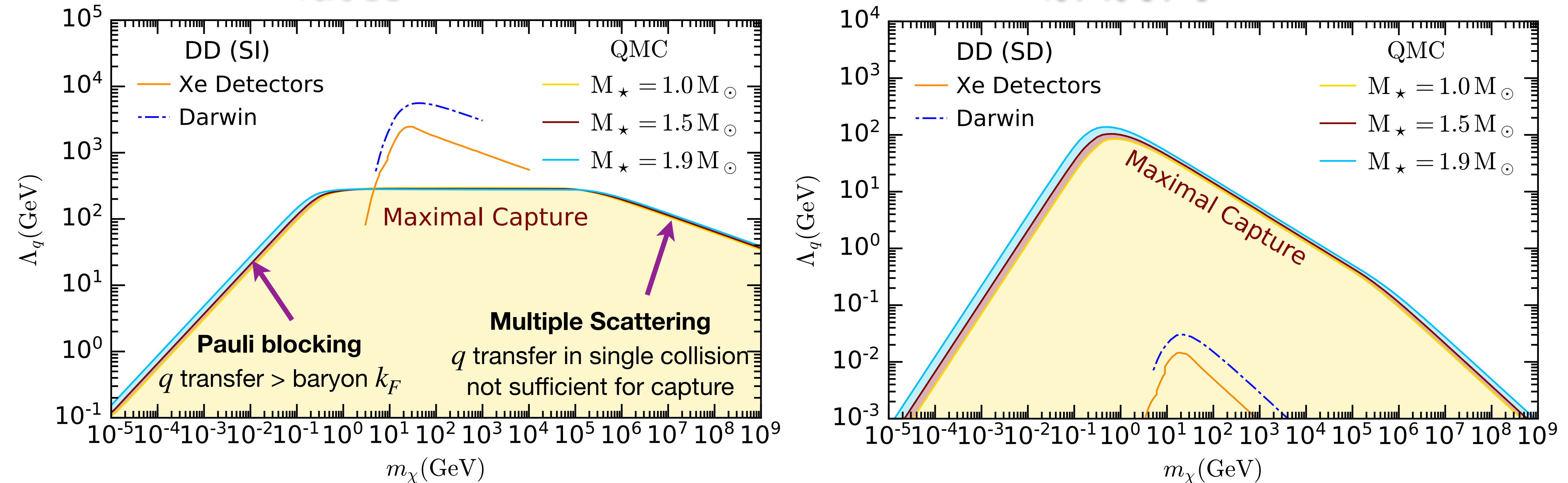
NS sensitivity to DM-nucleon interactions

EFT operators

$$\sigma_{i\chi} \propto \Lambda_q^{-4}$$

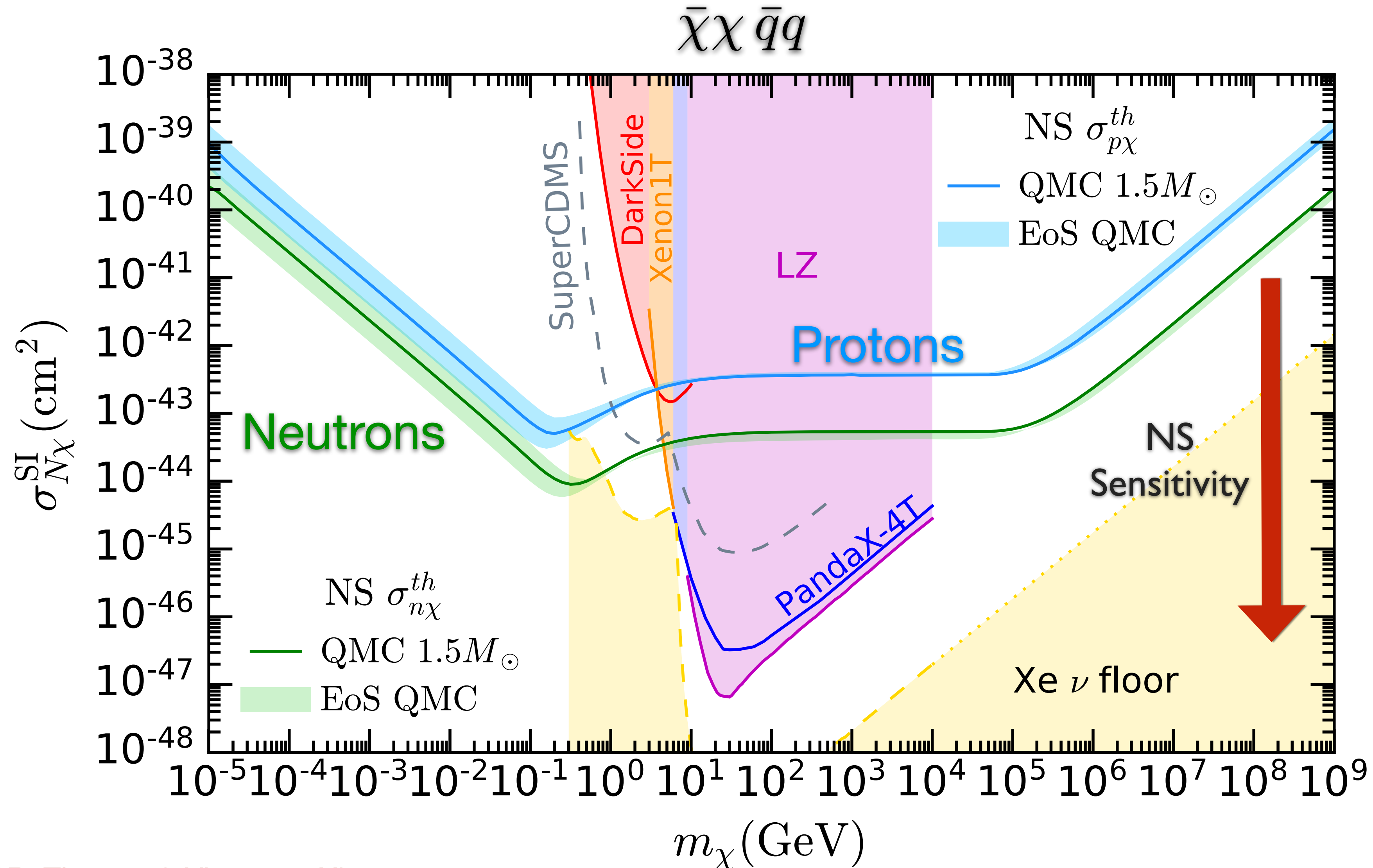
$\bar{\chi}\chi \bar{q}q$ unsuppressed

$\bar{\chi}\gamma^5\chi \bar{q}\gamma^5q$ q^4 suppressed



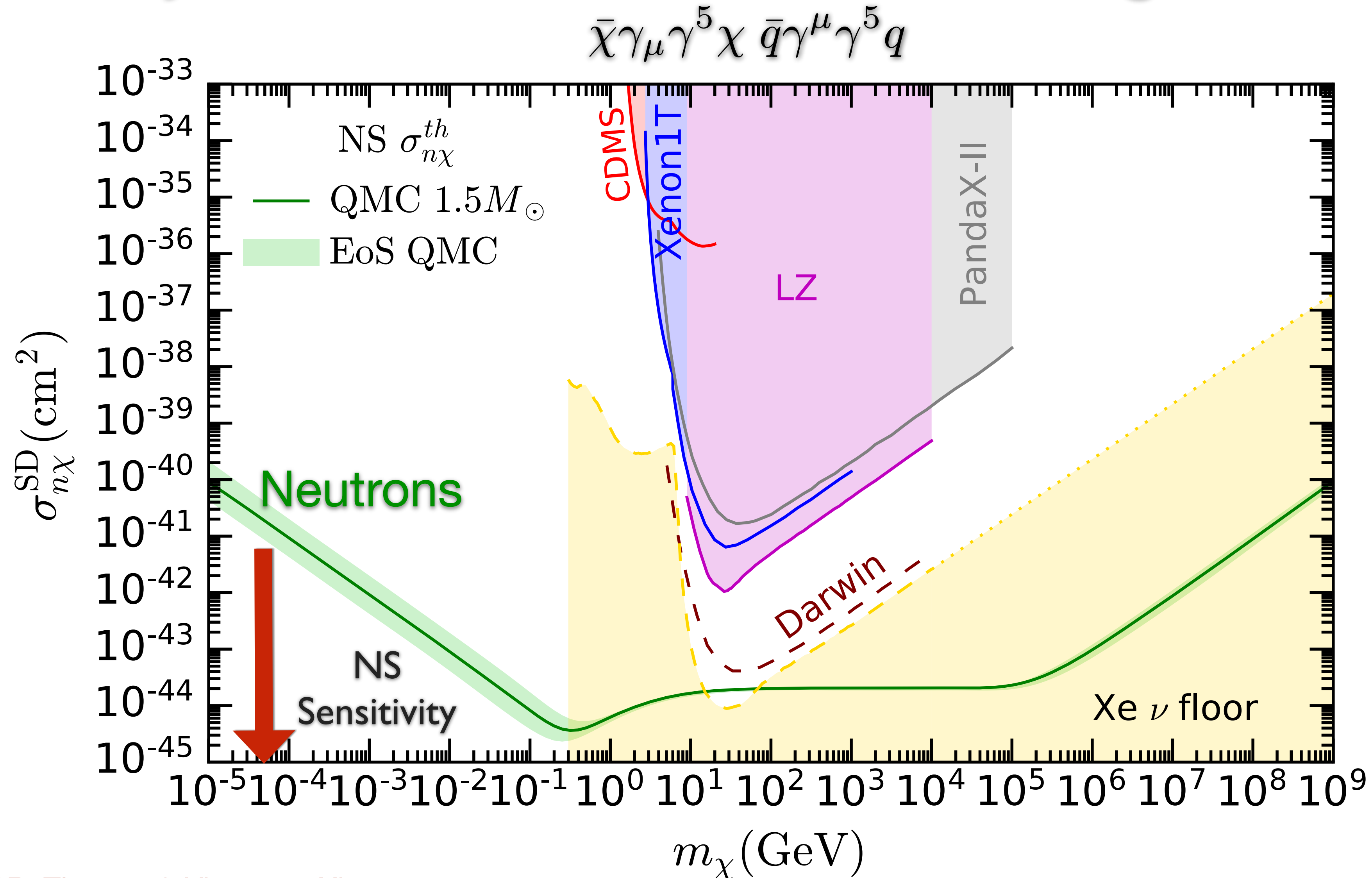
Bell, Busoni, Motta, SR, Thomas & Virgato, arXiv: 2108.02525

NS sensitivity to SI DM-nucleon scattering cross section



Bell, Busoni, Motta, SR, Thomas & Virgato, arXiv: 2108.02525

NS sensitivity to SD DM-neutron scattering cross section



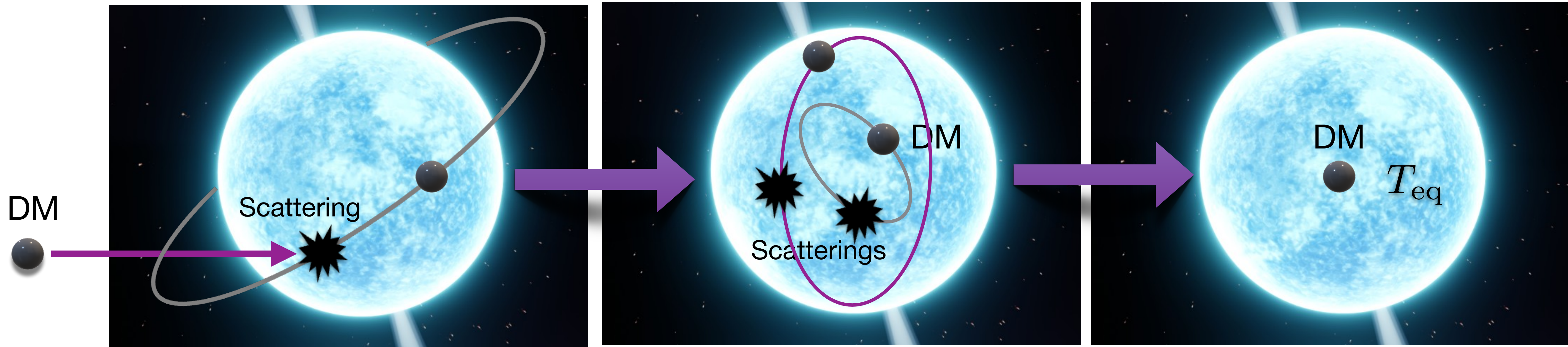
Bell, Busoni, Motta, SR, Thomas & Virgato, arXiv: 2108.02525

DM Thermalisation in NSs

- Capture

- Further scatterings

- Thermalised



1st stage:
 N_1 scatterings

2nd stage:
 N_2 scatterings

$$t_1^{\text{therm}}$$

$$\ll$$

$$t_2^{\text{therm}}$$

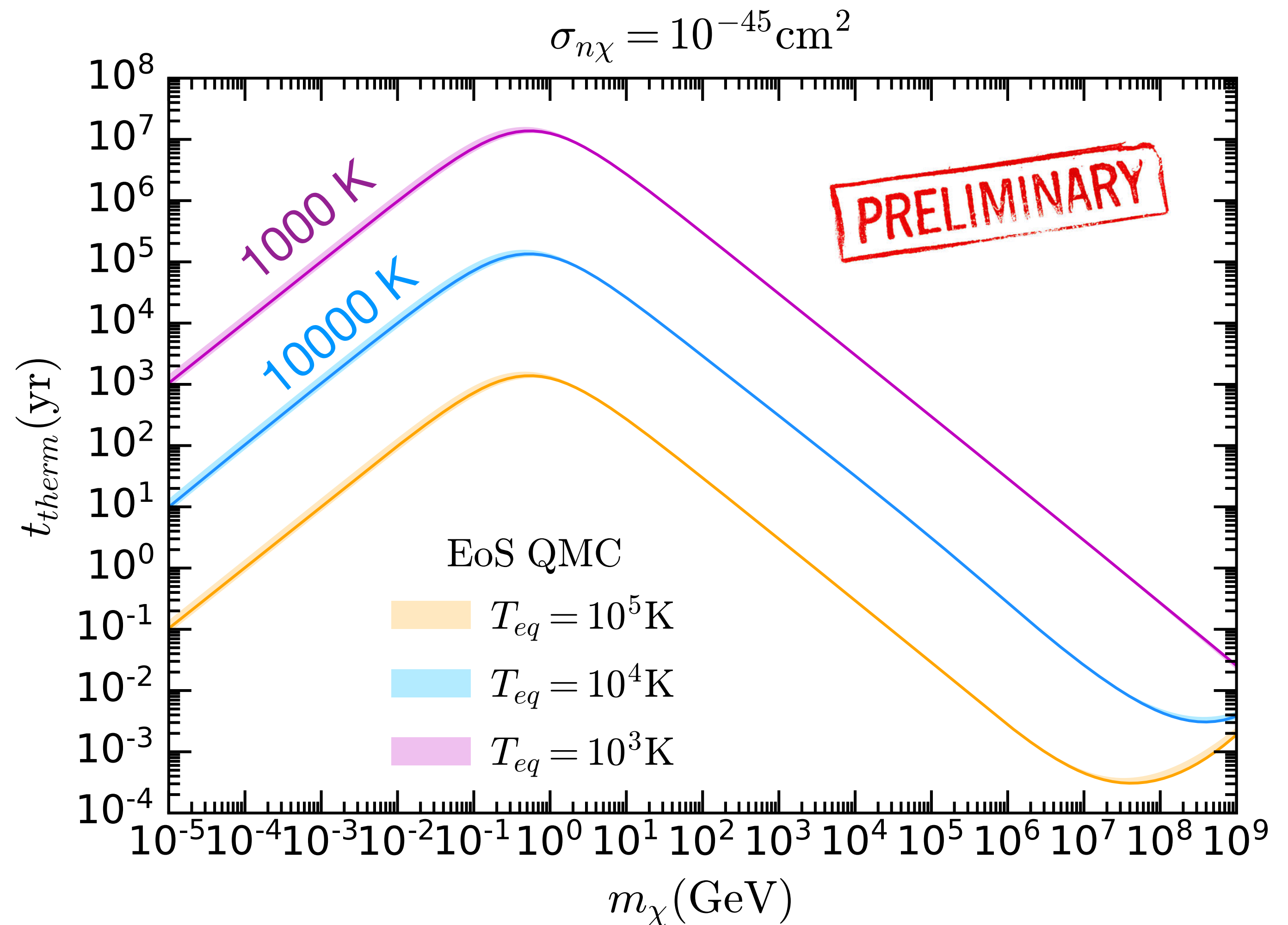
$$t_{\text{therm}} \simeq t_2^{\text{therm}}$$

DM Thermalisation in NSs

- After $N_1 + N_2$ scatterings DM reaches equilibrium temperature T_{eq}
- Thermalisation time (Pauli Blocking)
 - ➔ Sum of average time between collisions
 - ➔ Final energy transfer = T_{eq}

$$t_{\text{therm}} \simeq \sum_{n=N_1}^{N_2} \frac{1}{\Gamma^-(K_n)}$$

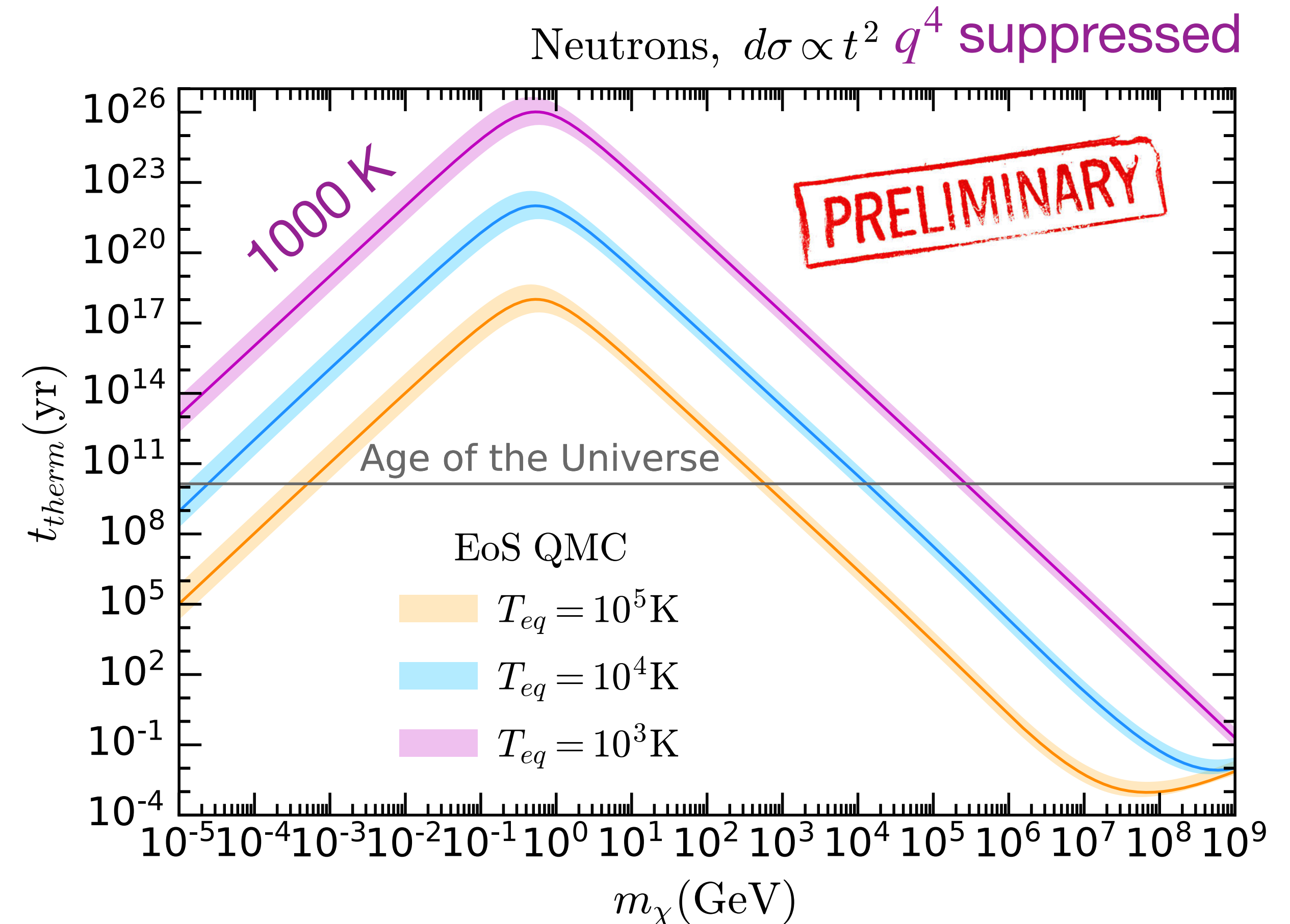
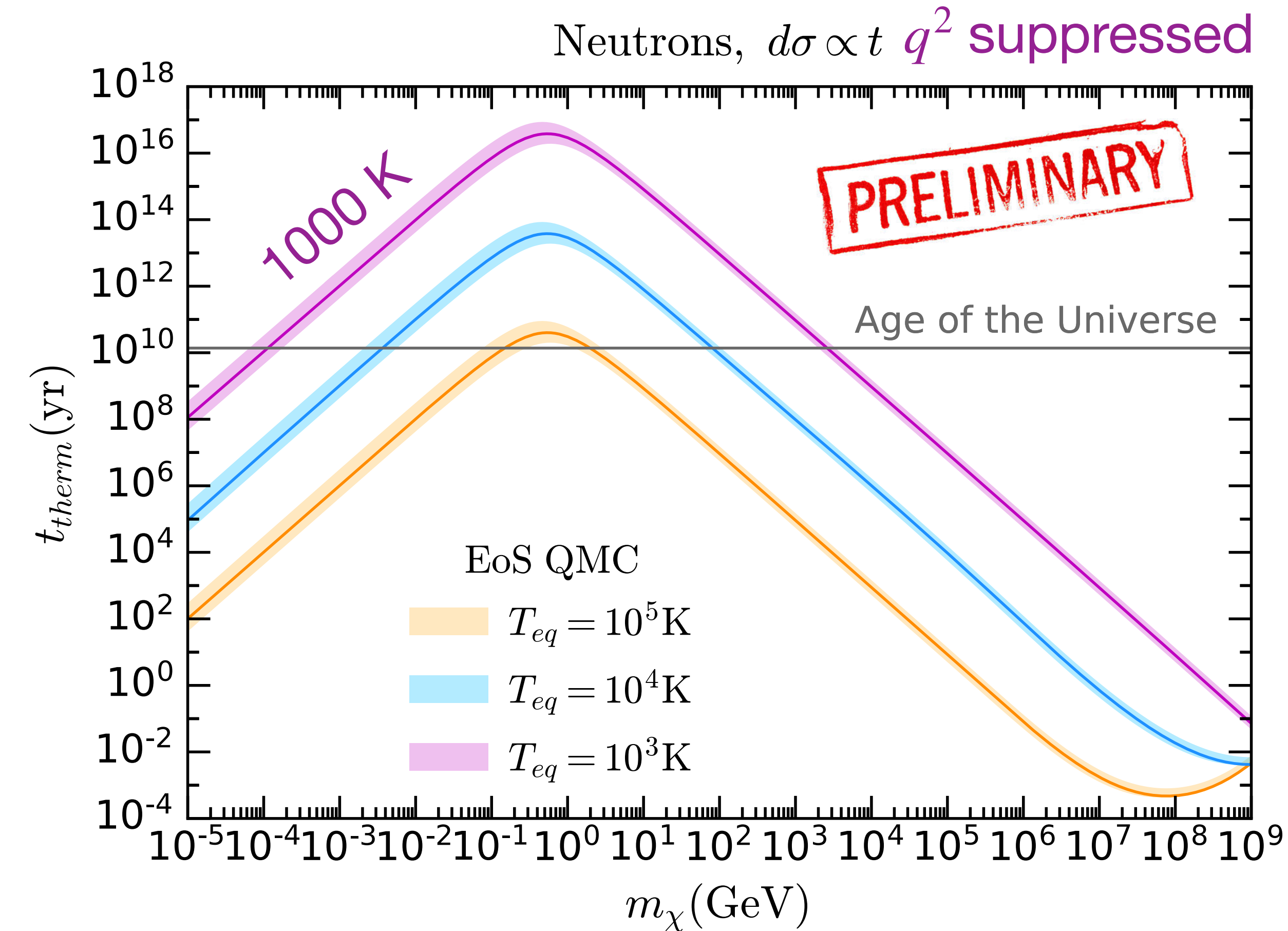
Bertoni, Nelson & Reddy, arXiv: 1309.1721



Bell, Busoni, SR & Virgato, in preparation

DM Thermalisation in NSs

- Momentum suppressed interactions



Bell, Busoni, SR & Virgato, in preparation

DM Thermalisation in NSs

EFT operators

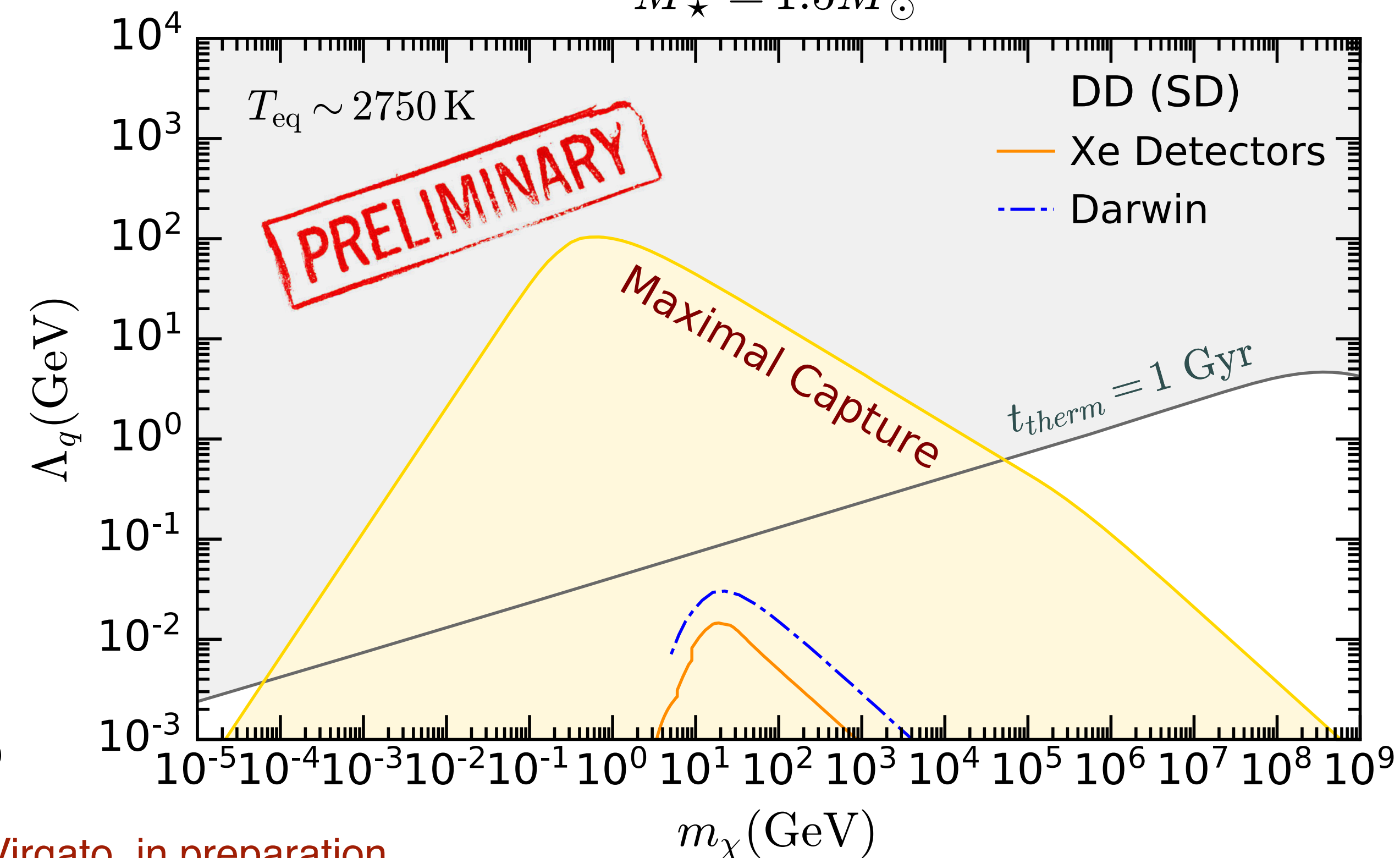
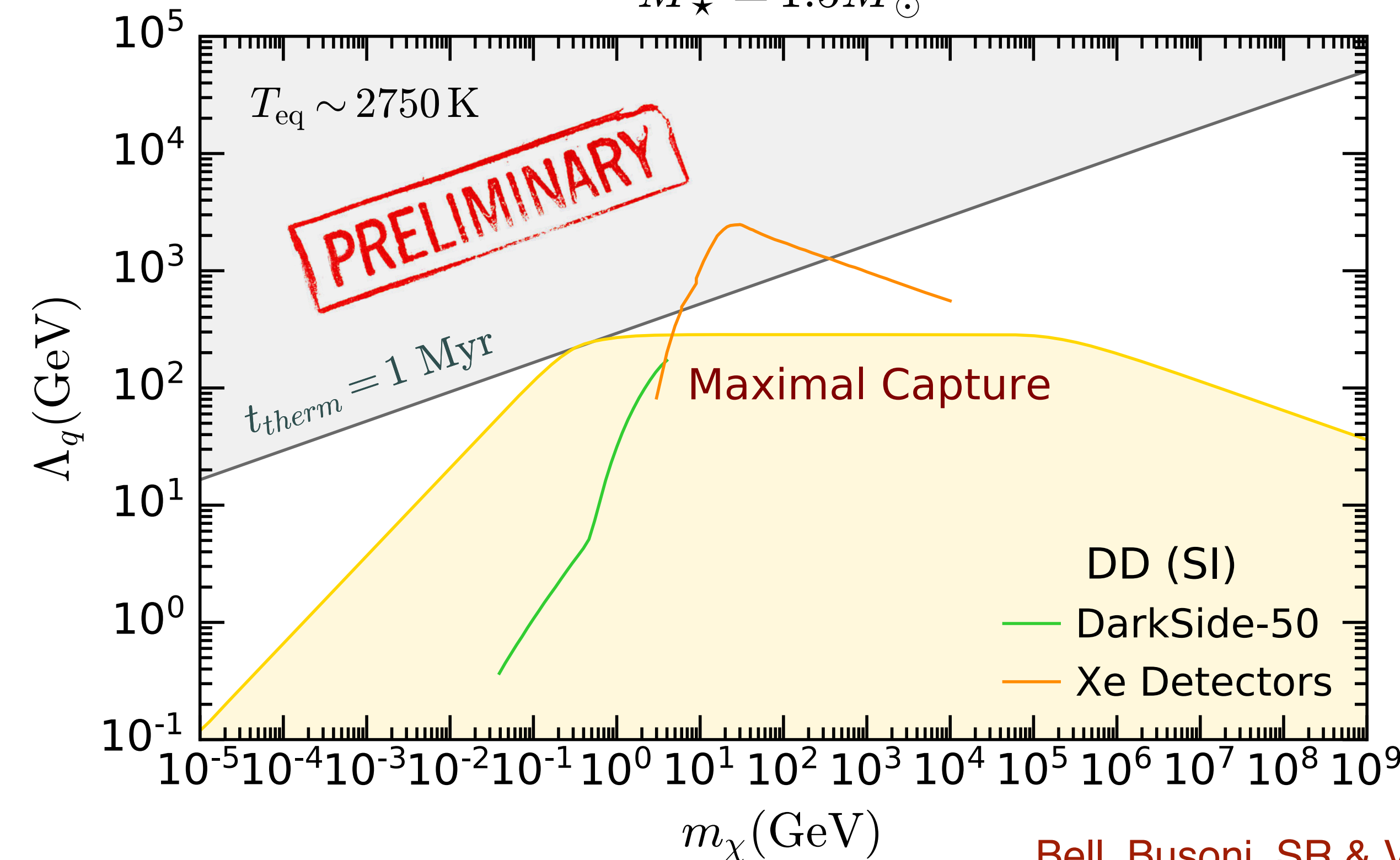
- Captured DM thermalises in ~ 1 Myr (unsuppressed interactions)

$\bar{\chi}\chi \bar{q}q$ **unsuppressed**

$\bar{\chi}\gamma^5\chi \bar{q}\gamma^5q$ q^4 **suppressed**

$M_\star = 1.5M_\odot$

$M_\star = 1.5M_\odot$



Bell, Busoni, SR & Virgato, in preparation

Capture and Annihilation Equilibrium

- Number of accumulated DM particles depends on the capture, evaporation and annihilation rates

$$\frac{dN_\chi}{dt} = C - EN_\chi - AN_\chi^2$$

Annihilation rate: $\Gamma_{ann} = \frac{1}{2} AN_\chi^2$

- When evaporation is negligible $m_\chi \gtrsim m_{evap}$

$$m_{evap} \sim \mathcal{O}(10\text{eV})$$

Bell, Busoni, SR & Virgato,
arXiv: 2010.13257

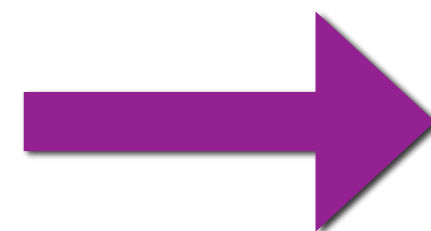
$$N_\chi(t) = \sqrt{\frac{C}{A}} \tanh\left(\frac{t}{t_{eq}}\right)$$

where

$$t_{eq} = \frac{1}{\sqrt{CA}}$$

$$A \simeq \frac{\langle \sigma_{ann} v_\chi \rangle}{(2\pi)^{3/2} r_\chi^3}$$

- If $t \gg t_{eq}$



$$\Gamma_{ann} = \frac{1}{2} C(\sigma)$$

capture - annihilation equilibrium

DM Annihilation in NSs

- If DM has not yet thermalized

$$t_{\text{eq}} = \frac{1}{\sqrt{CA}} \left(\frac{t_{\text{therm}} + t}{t} \right)^{\frac{\alpha}{2(2+n)}}$$

- Annihilation final states

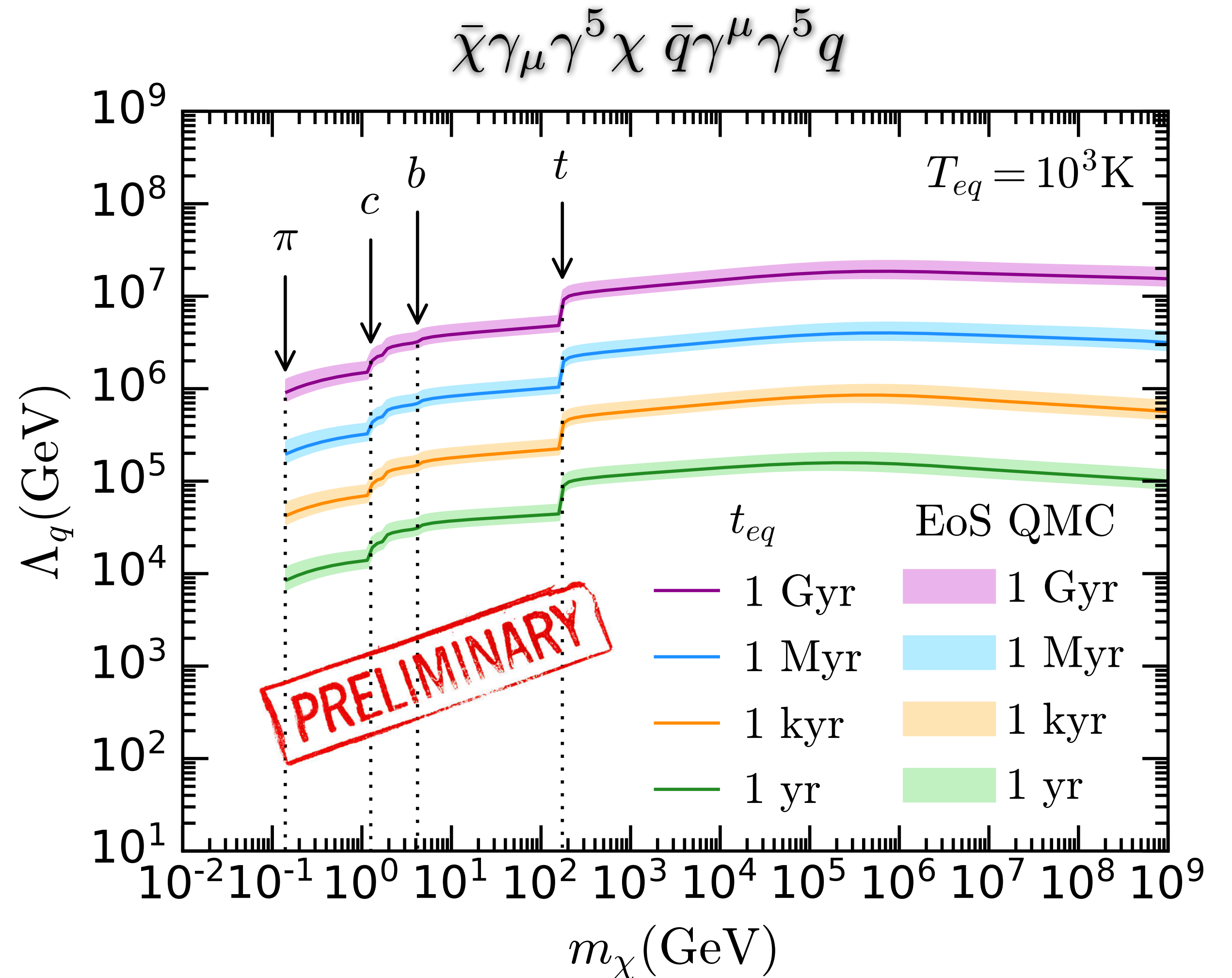
$$\chi\chi \rightarrow t\bar{t}, b\bar{b}, c\bar{c}, \pi^+\pi^-$$

- Annihilation to leptons

➔ Model dependent

➔ Pauli blocked

Bell, Busoni, SR & Virgato, in preparation



DM-induced Heating of NSs

EFT operators

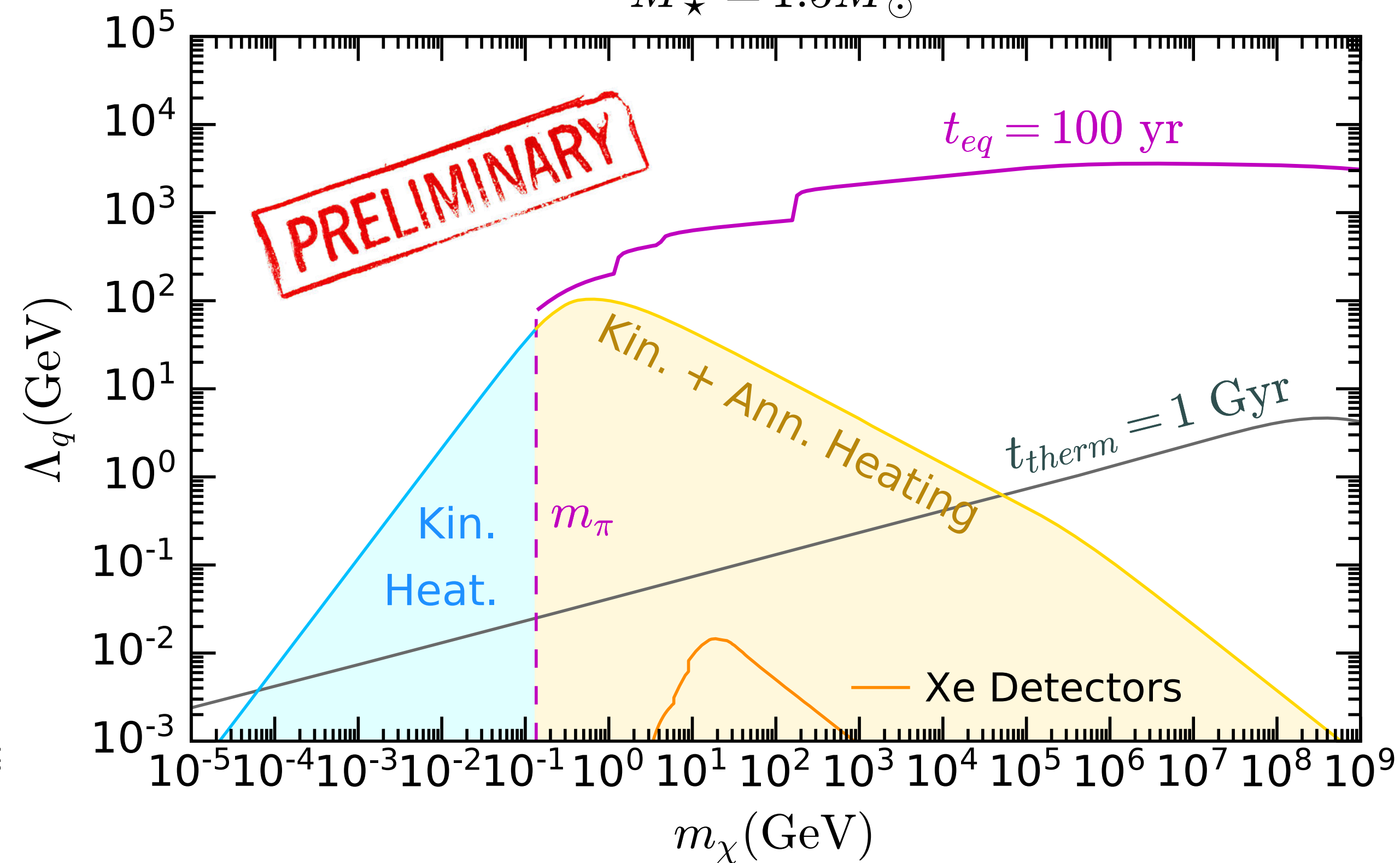
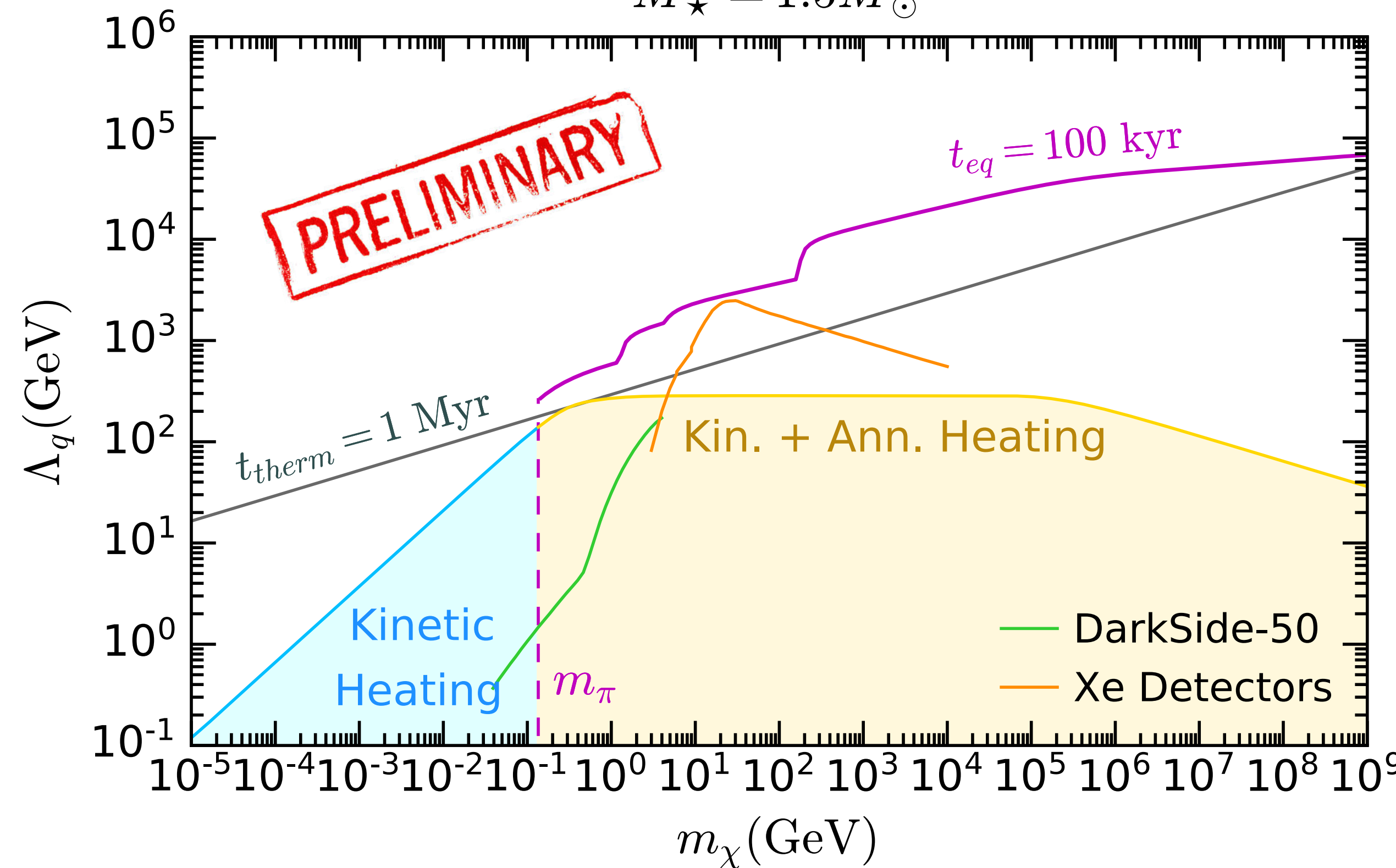
- Capture-annihilation equilibrium reached in ~ 1 yr (s-wave) up to 100 kyr (p-wave).

$$\bar{\chi}\chi \bar{q}q \quad \text{unsuppressed}$$

$$M_{\star} = 1.5M_{\odot}$$

$$\bar{\chi}\gamma^5\chi \bar{q}\gamma^5q \quad q^4 \text{ suppressed}$$

$$M_{\star} = 1.5M_{\odot}$$



Bell, Busoni, SR & Virgato, in preparation

Summary

- Improved calculation of the DM capture in neutron stars for (non-)relativistic, degenerate targets.
 - ➡ **Strong interactions** in NSs require treatment **beyond the free Fermi gas approximation**.
- Neutron stars could constrain different types of interactions, including those that are **velocity and momentum suppressed**.
- Captured DM would thermalise in ~ 1 Myr (unsuppressed interactions), momentum suppressed operators will need longer than the age of the Universe.
- Capture-annihilation equilibrium reached for all interactions in ~ 1 yr up to 100 kyr.
- Constraining DM interactions using DM-induced anomalous heating of neutron stars require
 - ➡ Observation of **old (cold)** neutron stars.
 - ➡ Better understanding of the **cooling process** in **neutron stars**.

Thank you for your
attention!