

Revisiting the CMB Accretion bound on PBH abundance

Daniele Gaggero

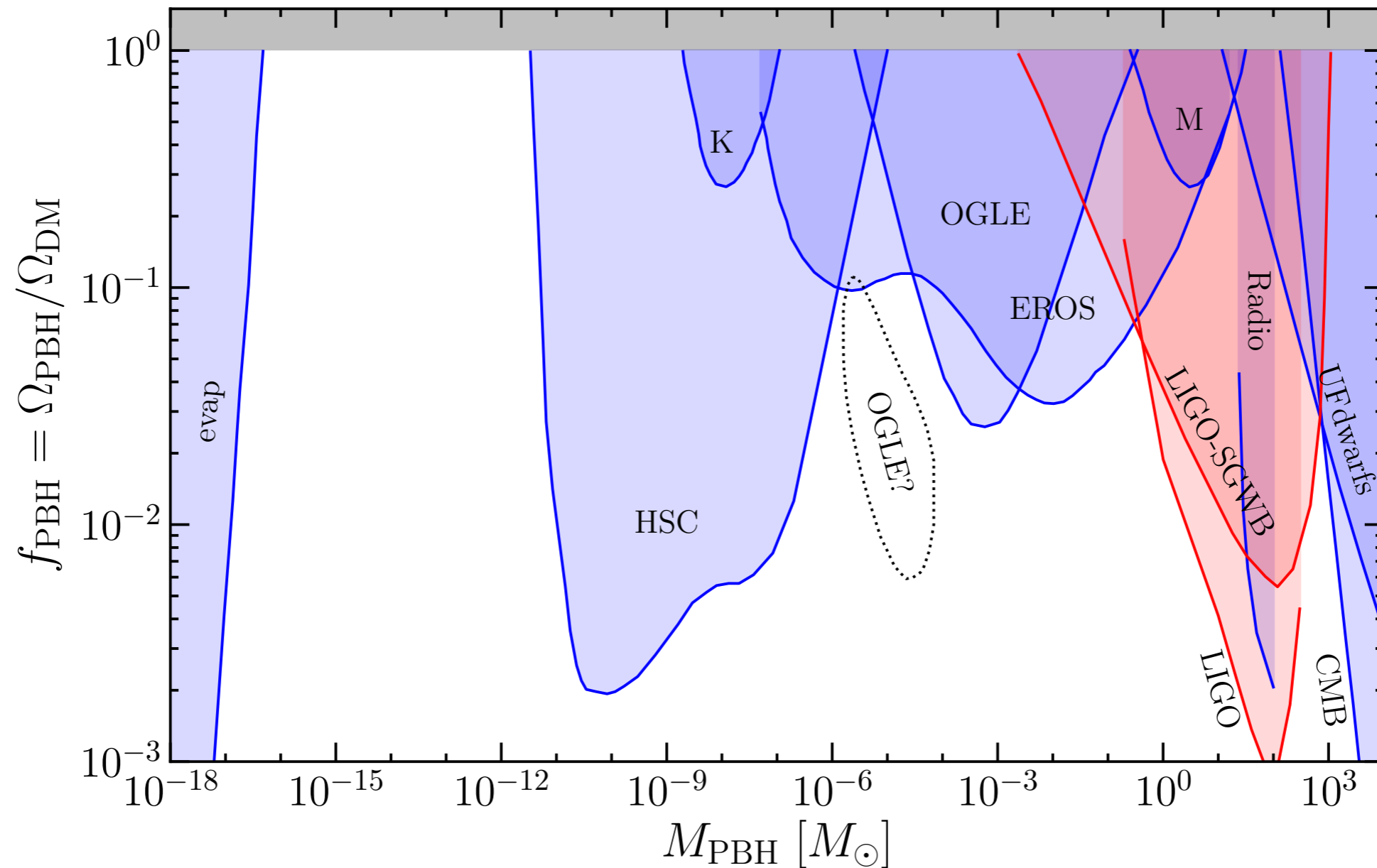
TEMPA
2023

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Sezione di Pisa

Black Holes and Dark Matter - Phenomenology

Constraints on PBHs as DM



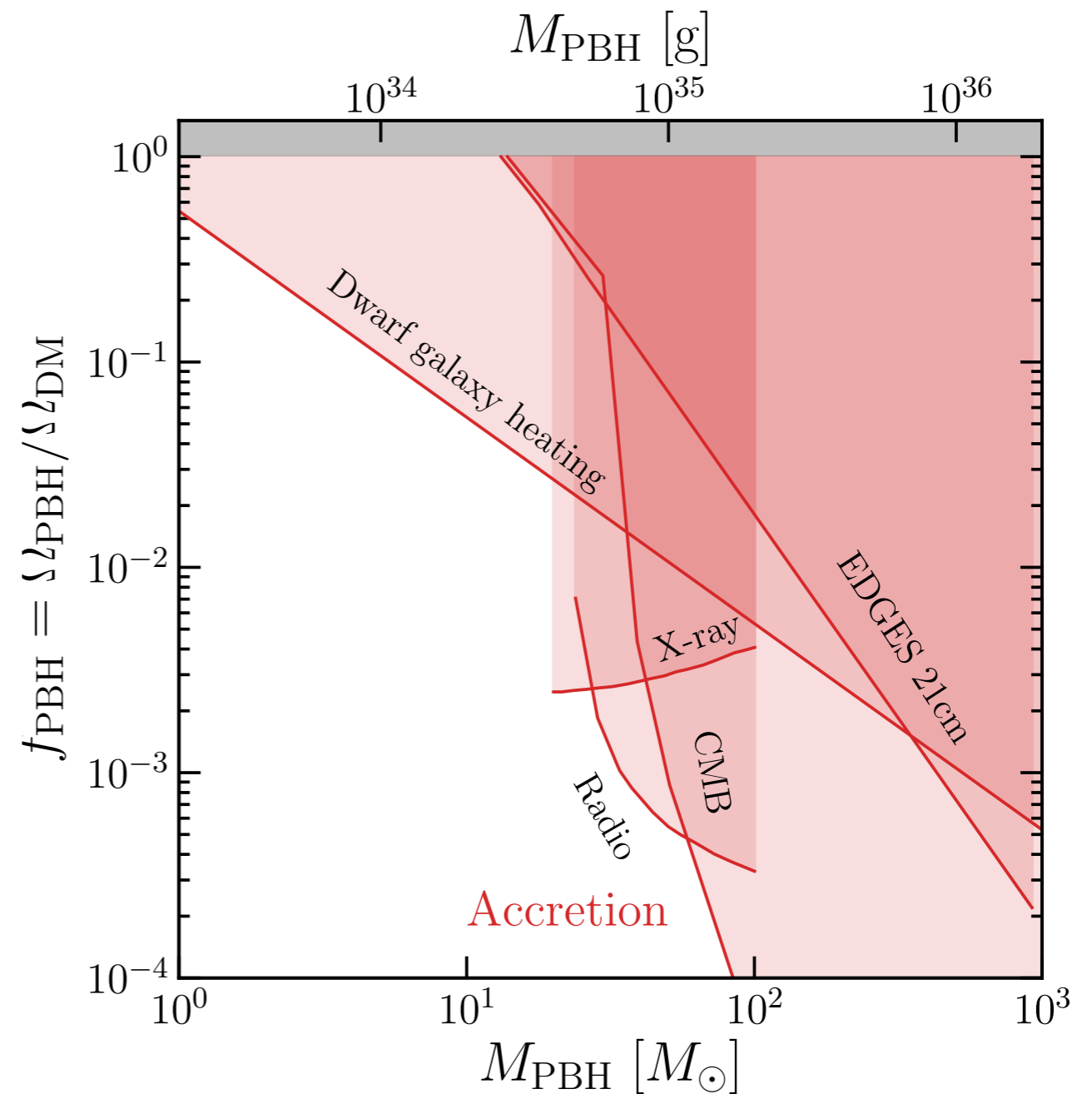
Credit: Bradley Kavanagh, <https://github.com/bradkav/PBHbounds>

Accretion bounds

Constraints from Accretion of baryonic matter



- *Very strong* especially for high masses
- **Astronomical environments:** X-ray/radio bounds (focus on Galactic center)
- **Cosmological bound** from CMB and 21 cm
- They rely on complicated accretion physics
- Comprehensive assessment of the uncertainties is very much needed!



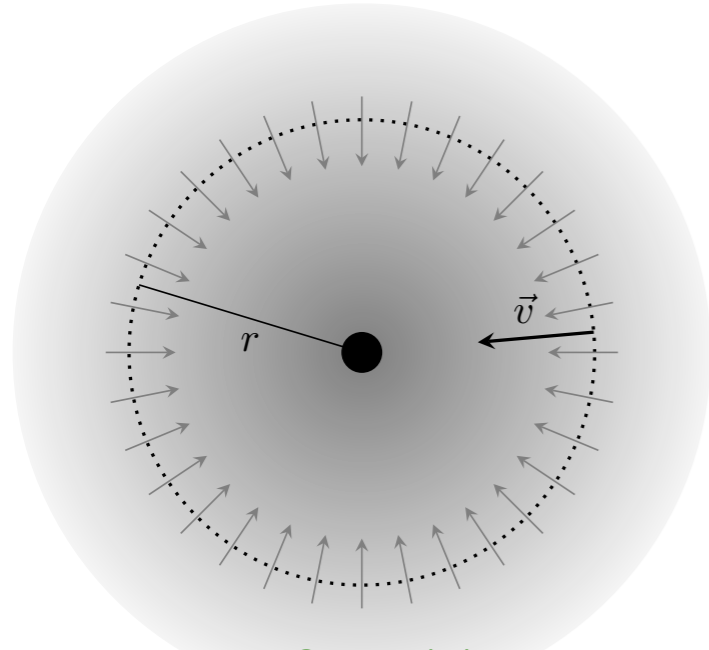
Accretion physics under the spotlight: BHL formalism.

Continuity equation for steady-state flow

$$\frac{1}{r^2} \frac{\partial}{\partial r} (r^2 \rho v) = 0$$

Euler equation

$$\rho v \frac{dv}{dr} = -\frac{dP}{dr} - \frac{GM\rho}{r^2}$$



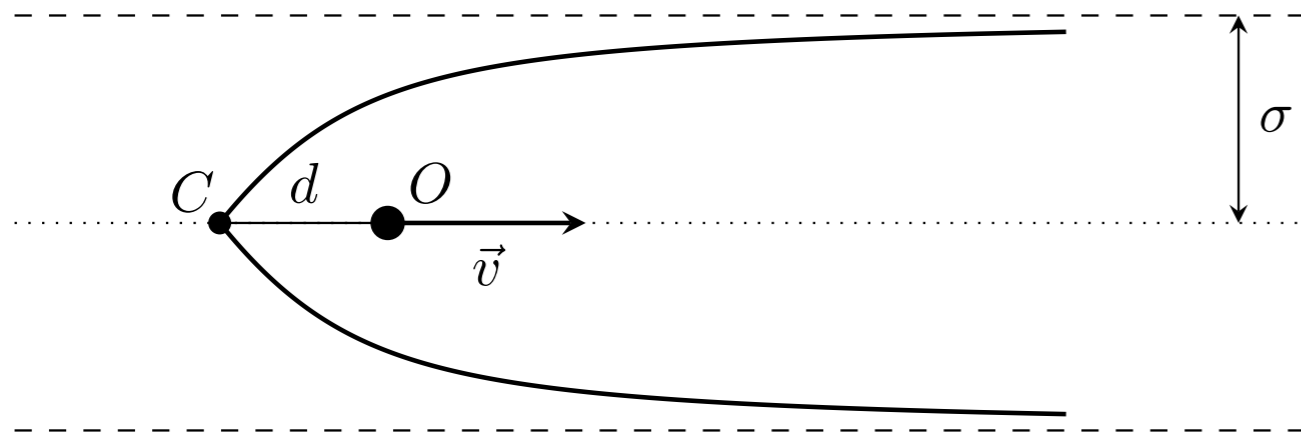
BH at rest: *Bondi* accretion rate

$$\dot{M} = 4\pi r_s^2 \rho(r_s) c_s(r_s) = \pi \frac{(GM)^2 \rho(\infty)}{c_s^3(\infty)} \left(\frac{2}{5-3\gamma} \right)^{\frac{5-3\gamma}{2(\gamma-1)}}$$

H. Bondi, MNRAS 112(2):195–204, 1952

H. Bondi and F. Hoyle, MNRAS 104(5):273–282, 1944

Moving BH: *Bondi-Hoyle-Littleton* accretion rate



$$\dot{M}_{\text{BHL}} = 4\pi \frac{(GM)^2 \rho_\infty}{(v^2 + c_\infty^2)^{3/2}}$$

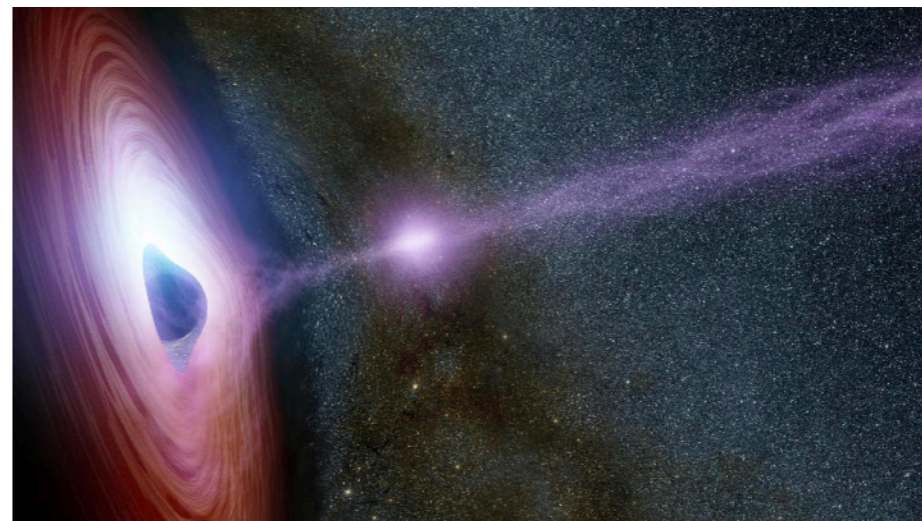
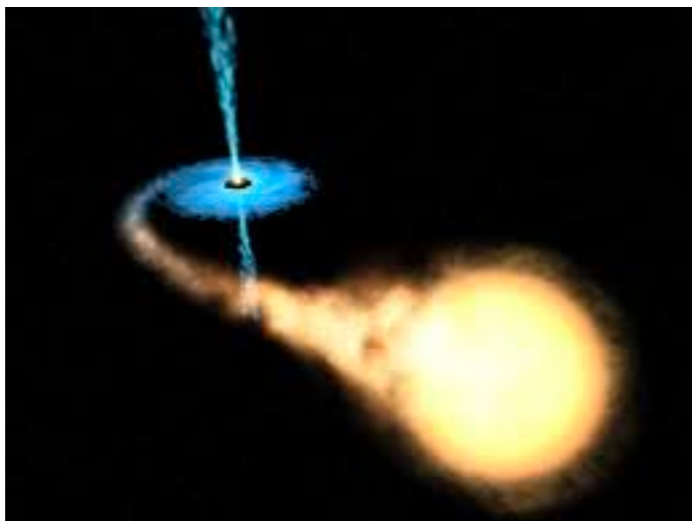
Accretion physics under the spotlight: BHL formalism.

Bondi-Hoyle-Littleton formula needs to be “fudged” because of observational constraints related to local neutron stars, the SMBH at the center of the Galaxy, and AGNs.

$$\dot{M} = 4\pi\lambda(GM_{BH})^2\rho(v_{BH}^2 + c_s^2)^{-3/2}$$

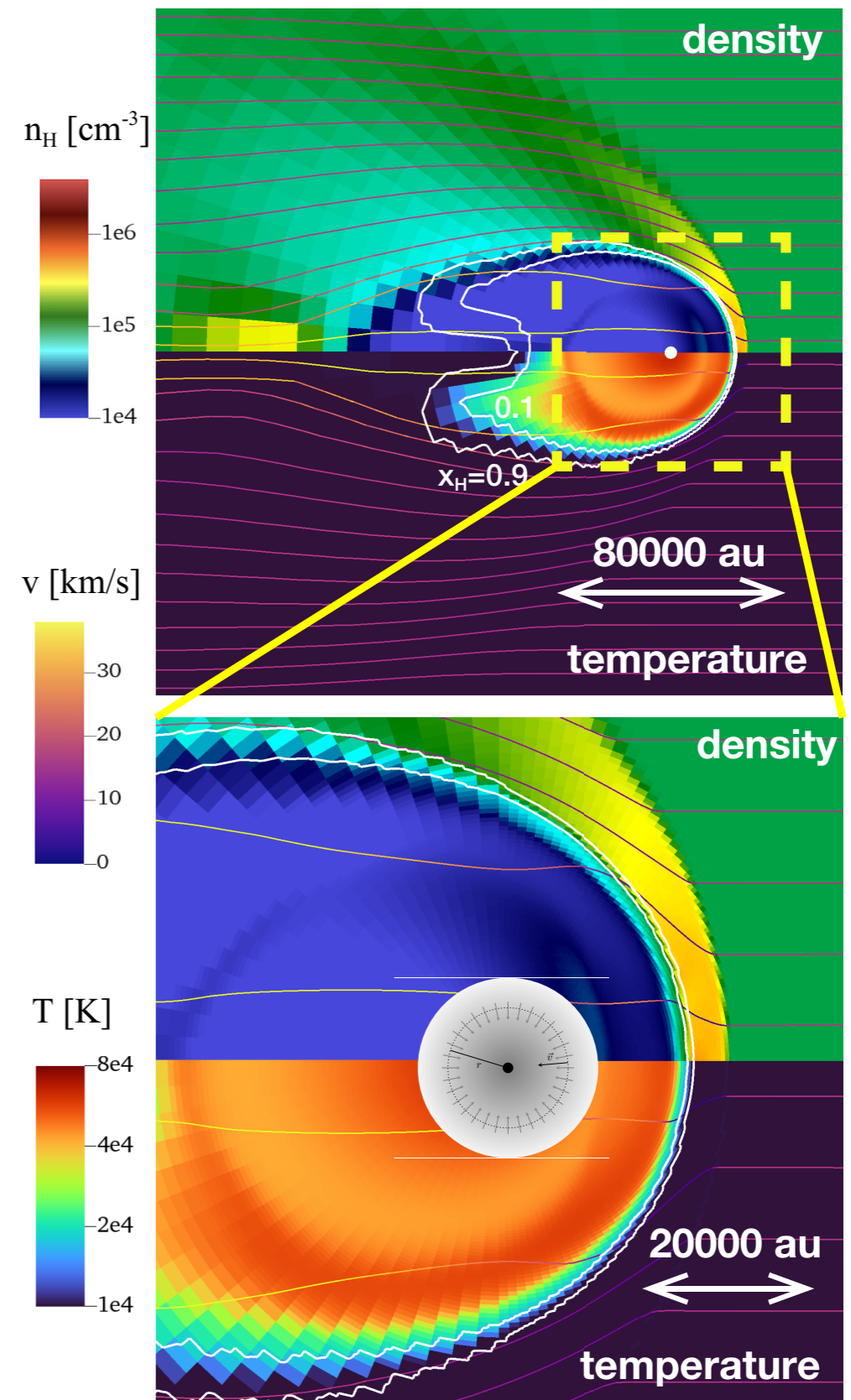
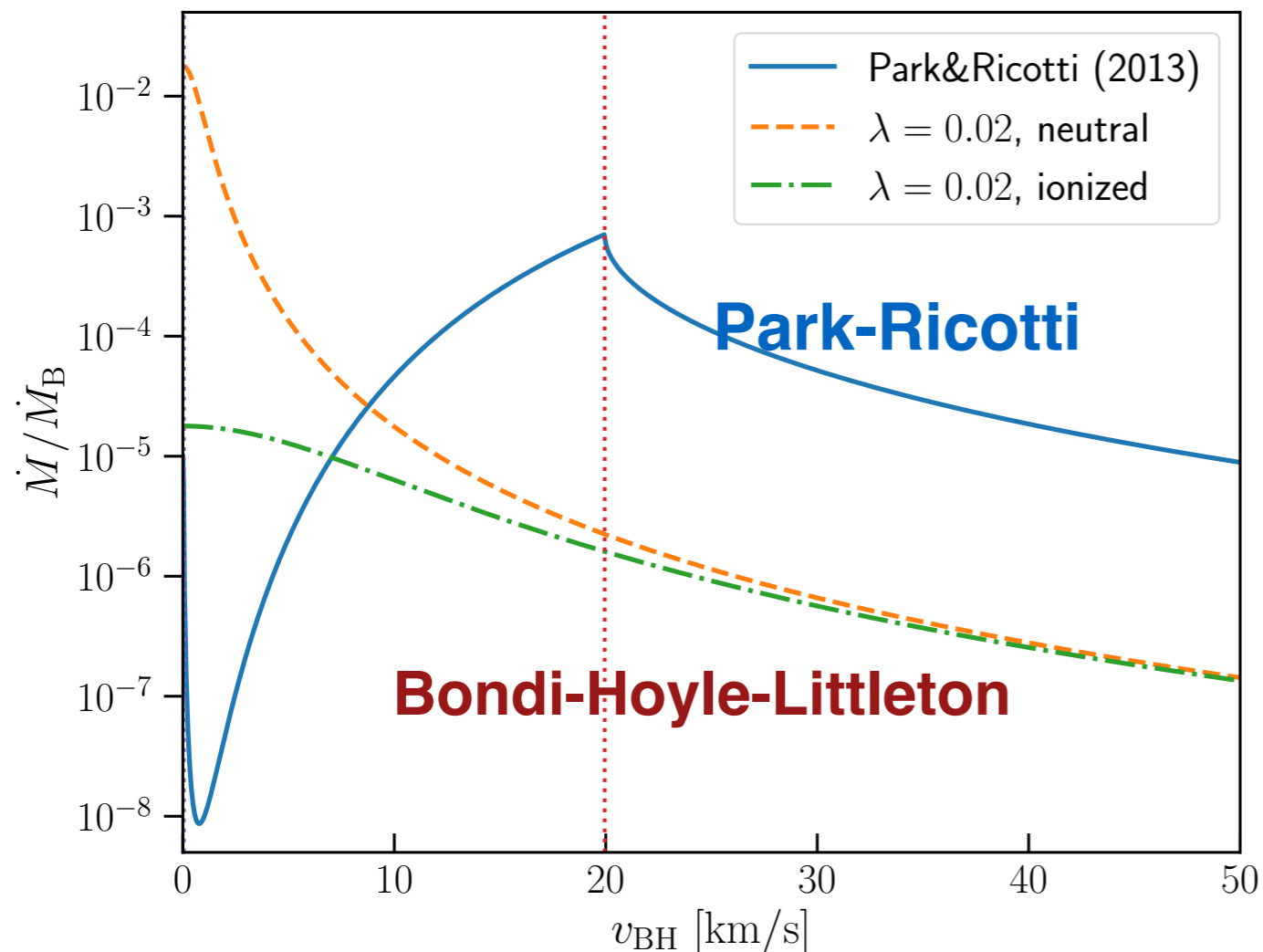
- [Perna et al. 2003](#), “*Bondi accretion and the problem of missing isolated neutron stars*”
- [S. Pellegrini 2005](#), “*Nuclear Accretion in Galaxies of the Local Universe: Clues from Chandra Observations*” (explanation for the radiative quiescence of supermassive black holes in the local Universe)
- [Wang et al. 2013](#), “*Dissecting X-ray-emitting Gas around the Center of our Galaxy*”

The fudge factor takes into account several effects, including the role of outflows



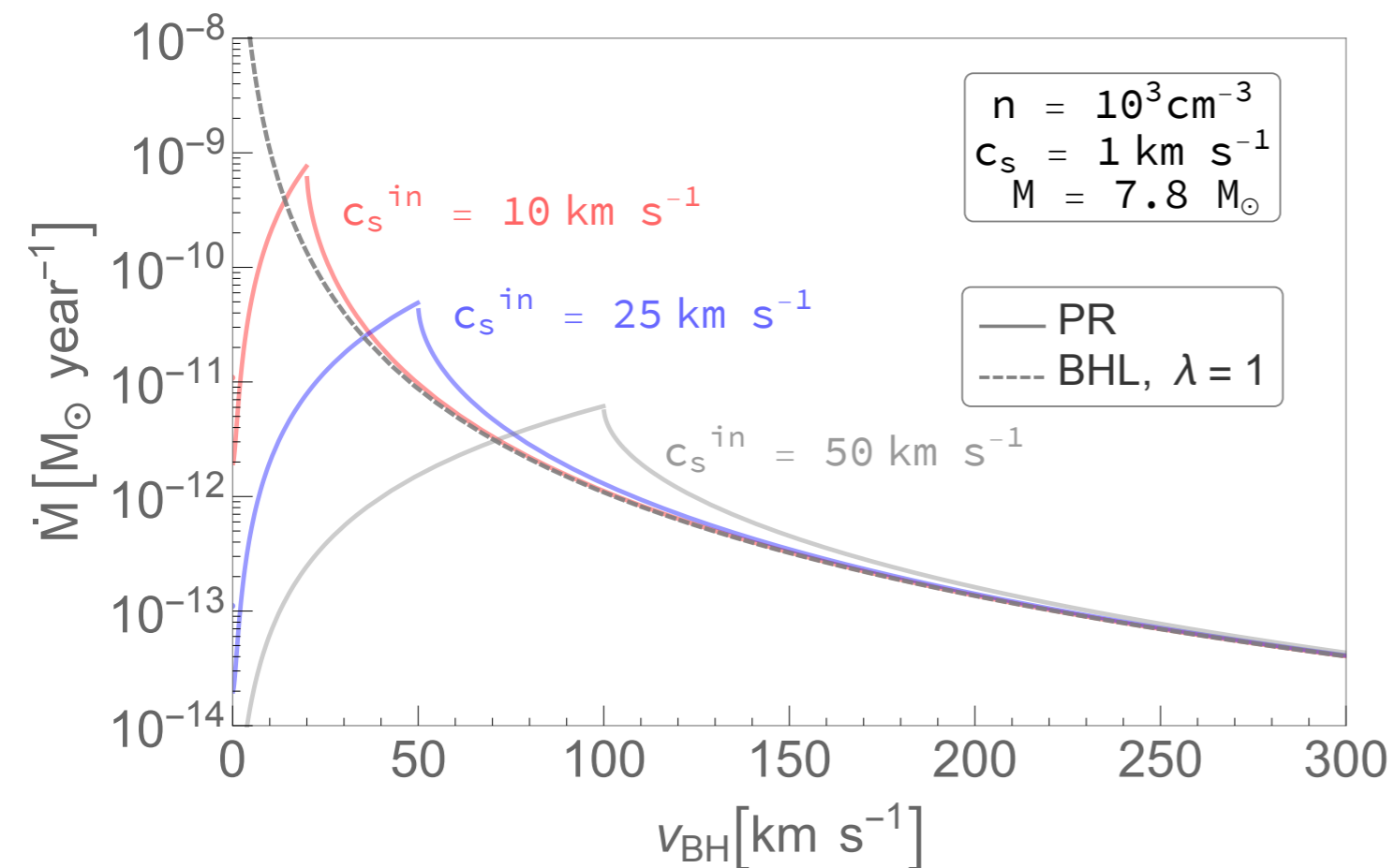
The Park-Ricotti model

- **Park-Ricotti model:** numerical simulations + semi-analytical parametrization in presence of radiative feedback.
- Suppression of the accretion rate at low velocity, due to the formation of an ionized bubble

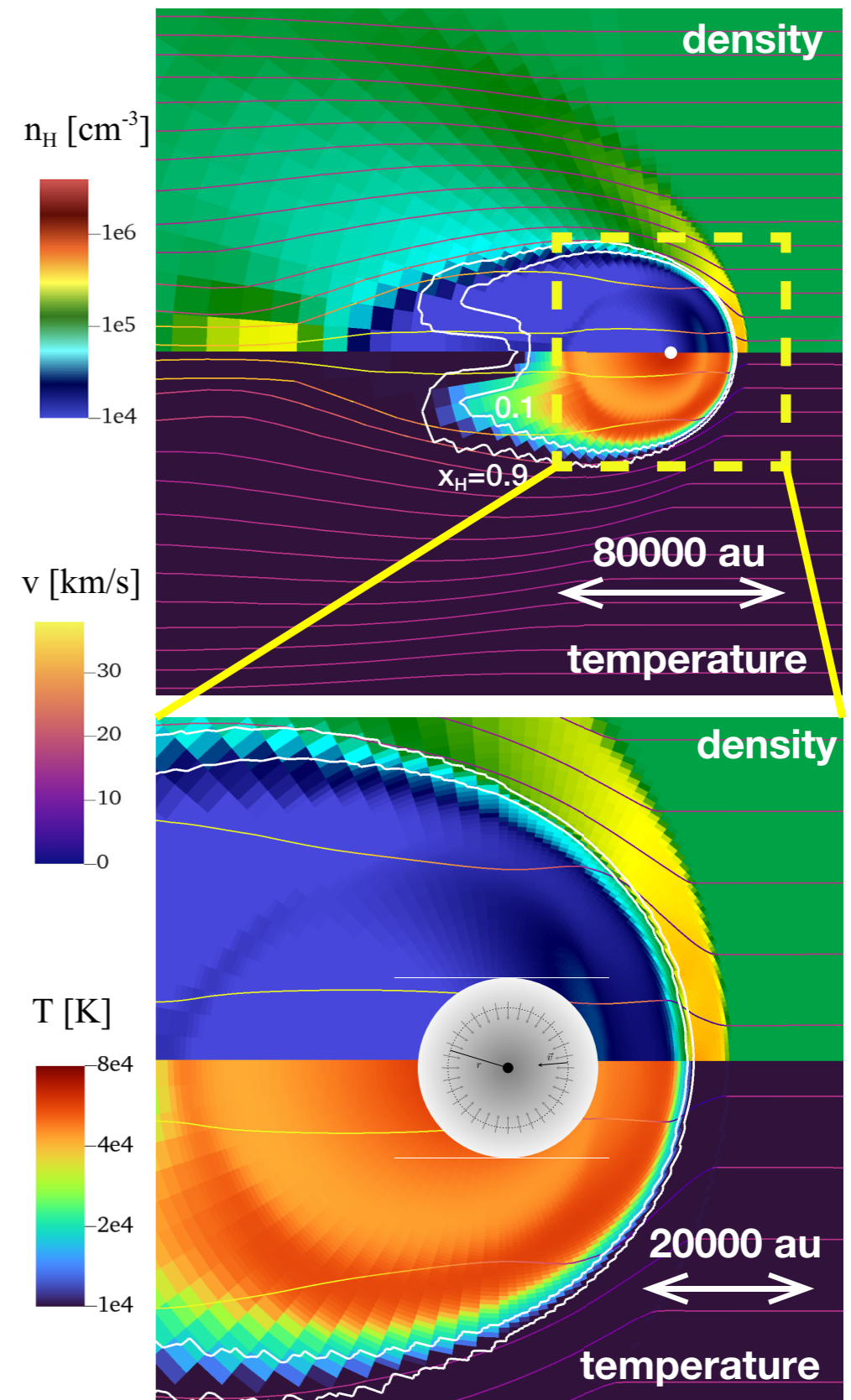


The Park-Ricotti model

- **Park-Ricotti model:** numerical simulations + semi-analytical parametrization in presence of radiative feedback.
- Peaks of accretion rate depends on ionized sound speed



Scarcella+ 2012.10421



The physics behind the bound

- PBHs accrete baryonic matter.
 - The accretion rate \dot{M} depends on ambient density and PBH - baryon relative speed. **BHL** and **PR** model.
 - Ambient density dilutes with decreasing redshift

$$\rho_{\infty} = m_p n_{\infty} \approx m_p 200 \text{ cm}^{-3} \left(\frac{1+z}{1000} \right)^3 \quad \text{Poulin+ 1707.04206}$$

- PBH speed relative to baryons also decreases according to linear theory:

$$\sqrt{\langle v_L^2 \rangle} \simeq \min \left[1, \frac{1+z}{1000} \right] \times 30 \text{ km/s}.$$

Revisiting the Cosmological constraint

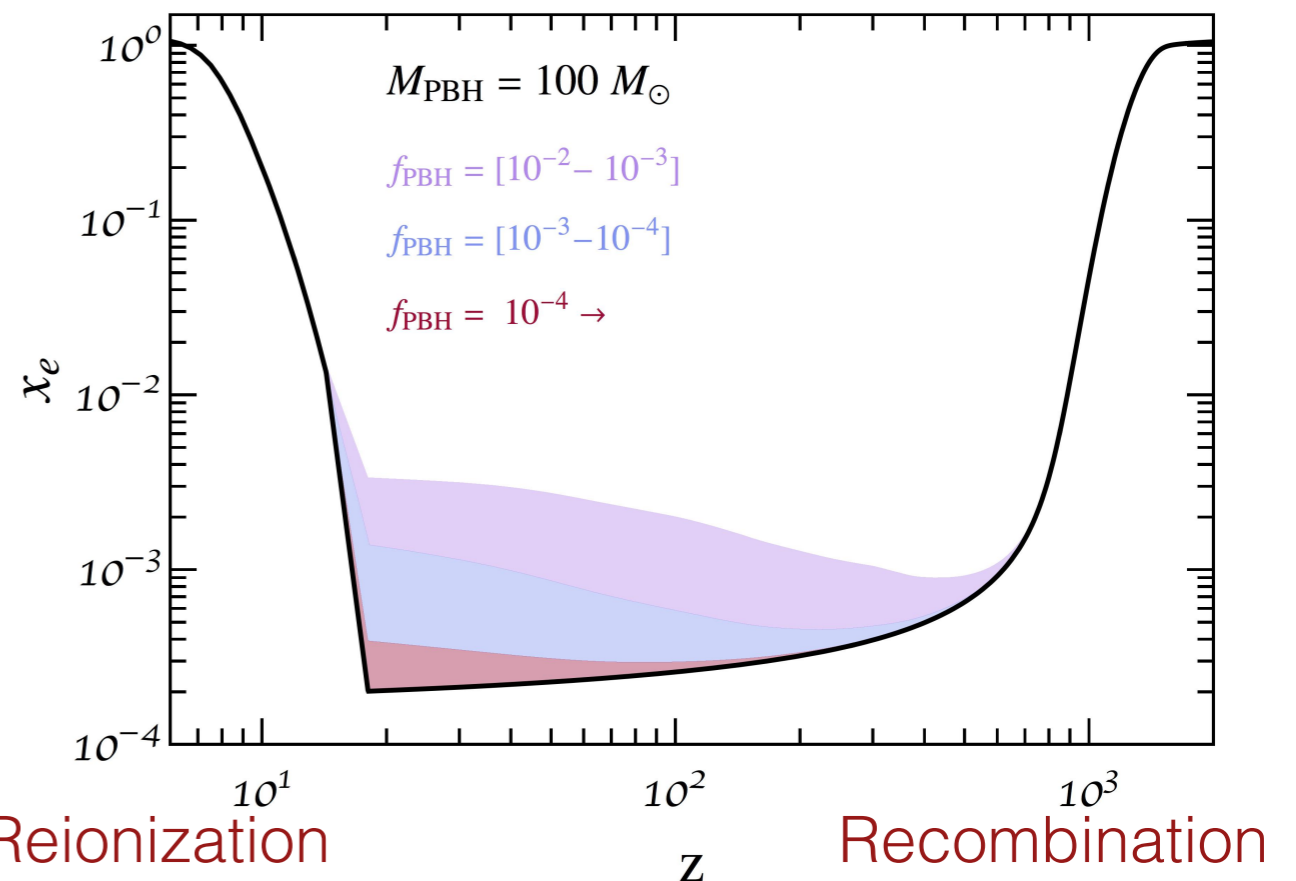
The physics behind the bound

- Accretion disks emits *ionizing radiation* during the Dark Ages (between Recombination and Reionization):
 - IGM is **heated up** (alteration of T_M)
 - IGM is also **partially ionized** (alteration of the *free electron fraction* X_e)

$$\frac{dx_e(z)}{dz} = \frac{1}{(1+z)H(z)} (R(z) - I(z) - I_X(z)) ,$$
$$\frac{dT_M}{dz} = \frac{1}{1+z} \left[2T_M + \gamma(T_M - T_{\text{CMB}}) \right] + K_h .$$

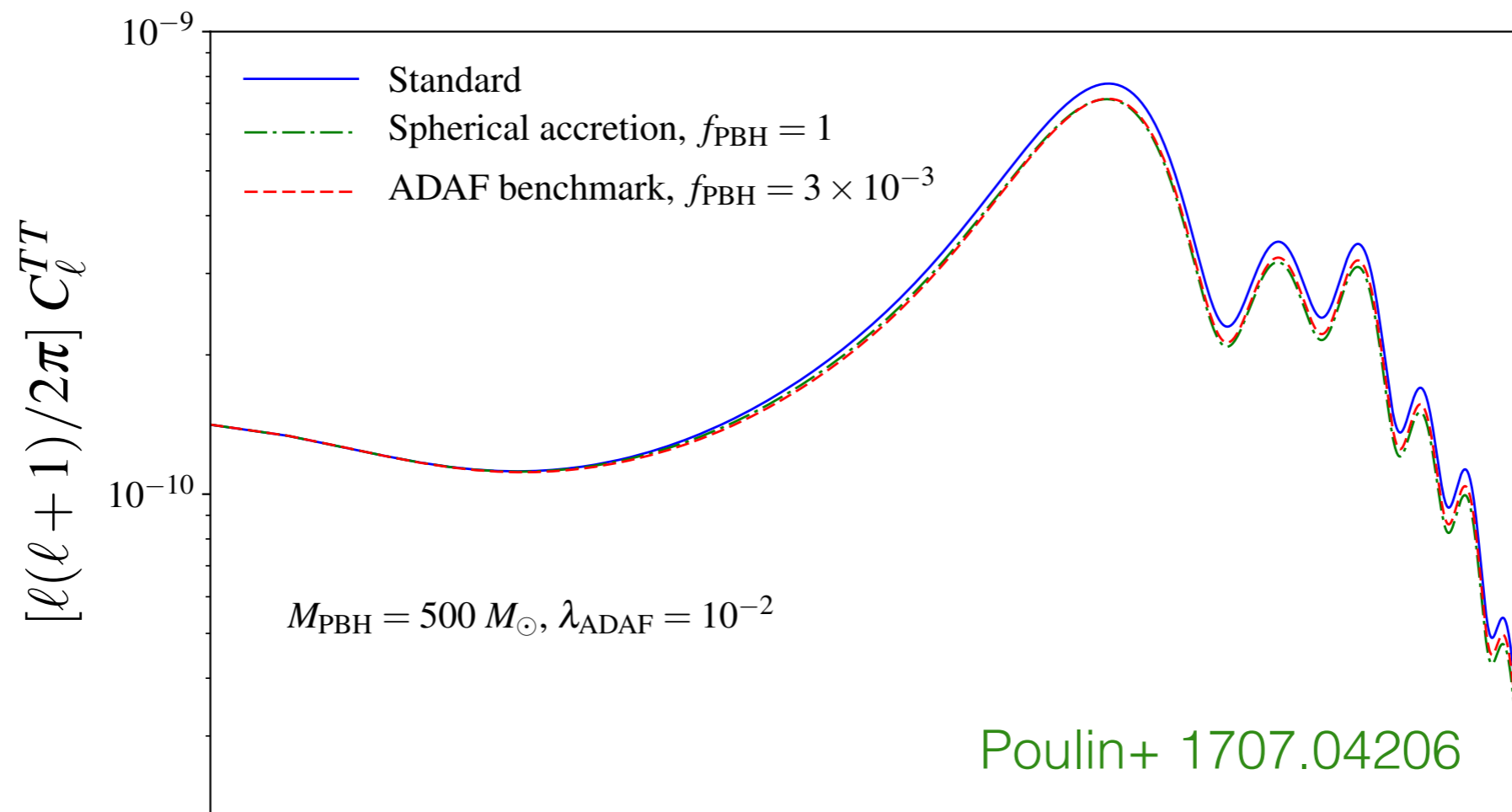
Stocker+ 1801.01871

Mena+ 1906.07735



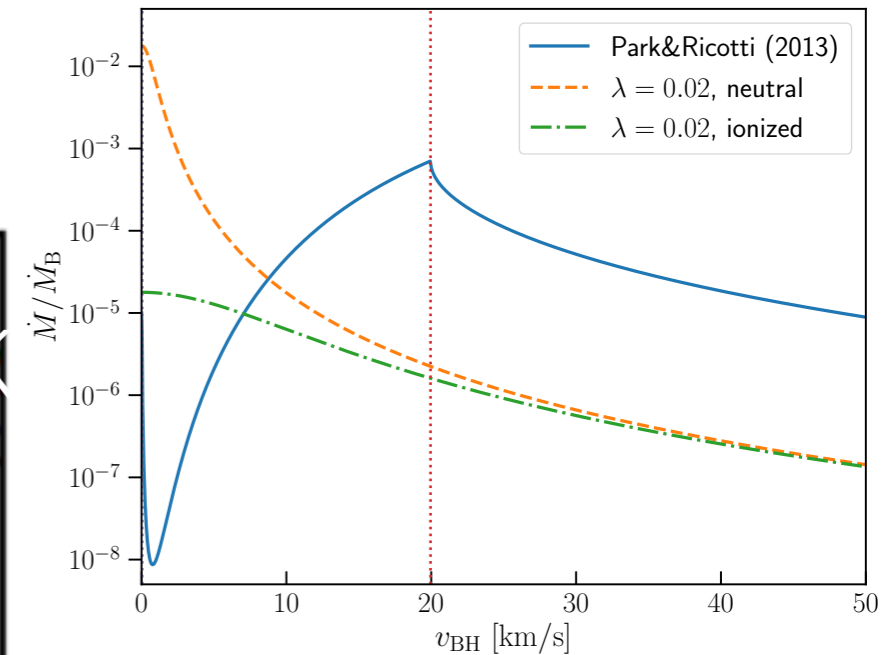
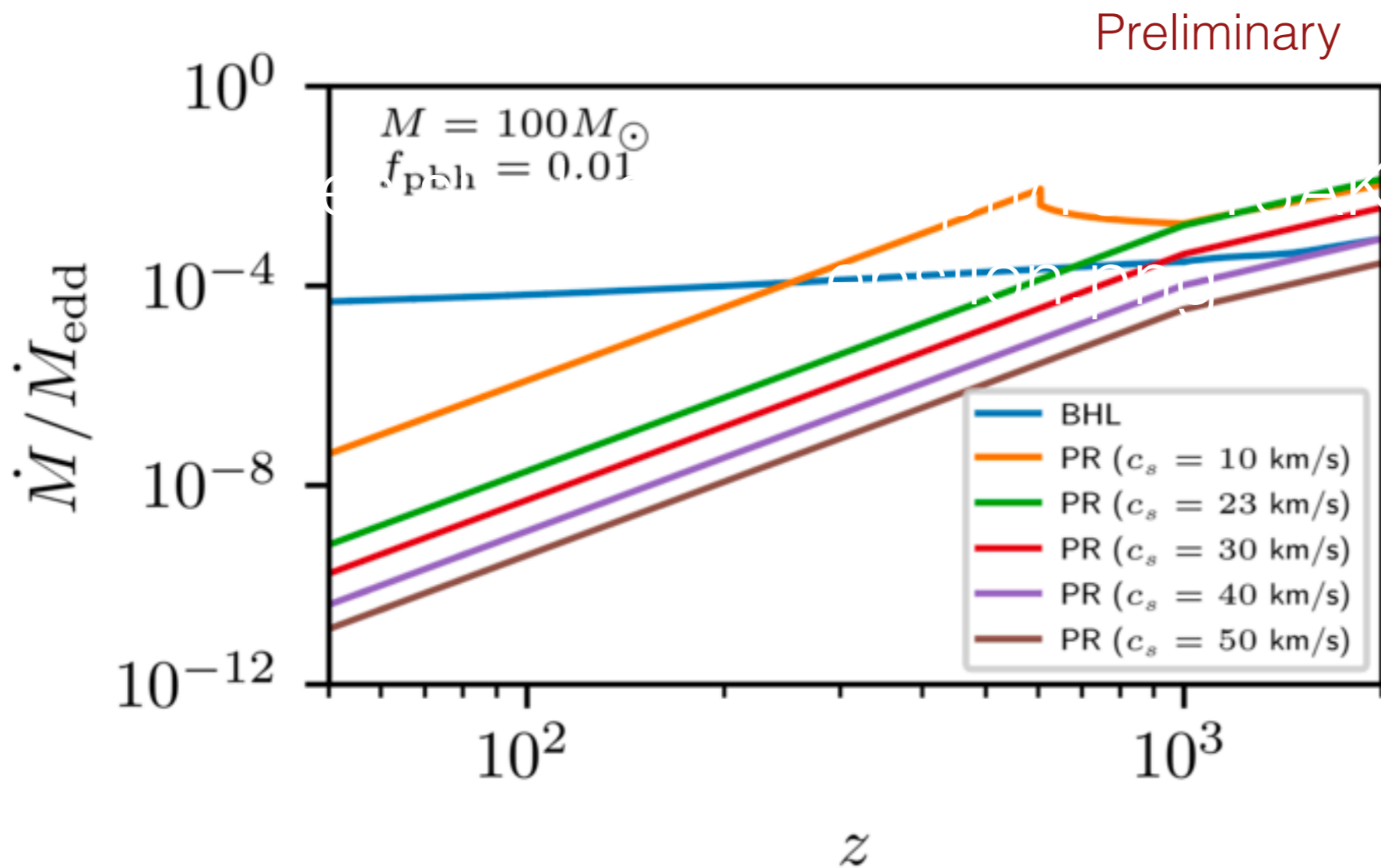
The physics behind the bound

- Impact on CMB anisotropy is due to the alteration of the visibility function and the recombination optical depth



Revisiting the Cosmological constraint: Results

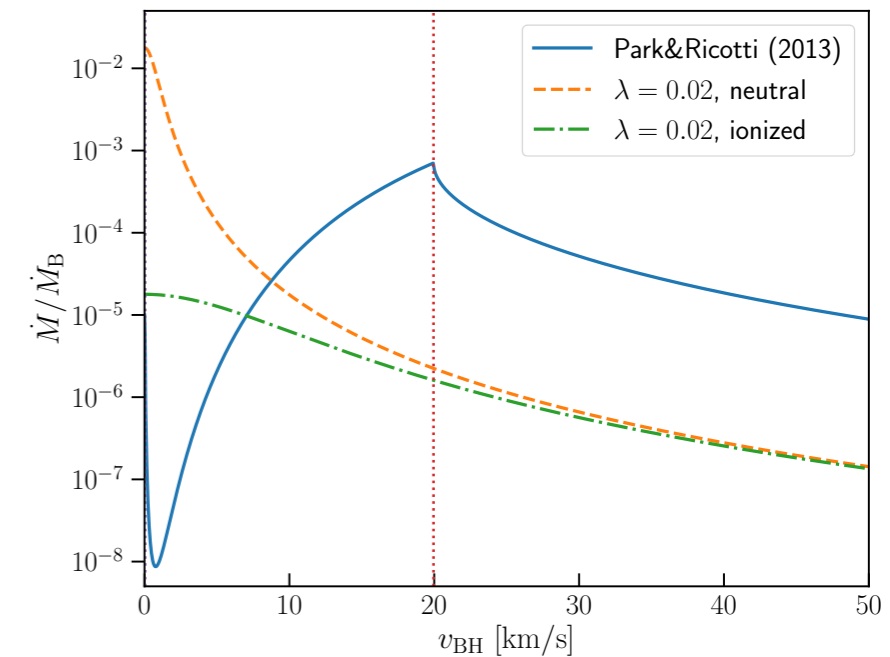
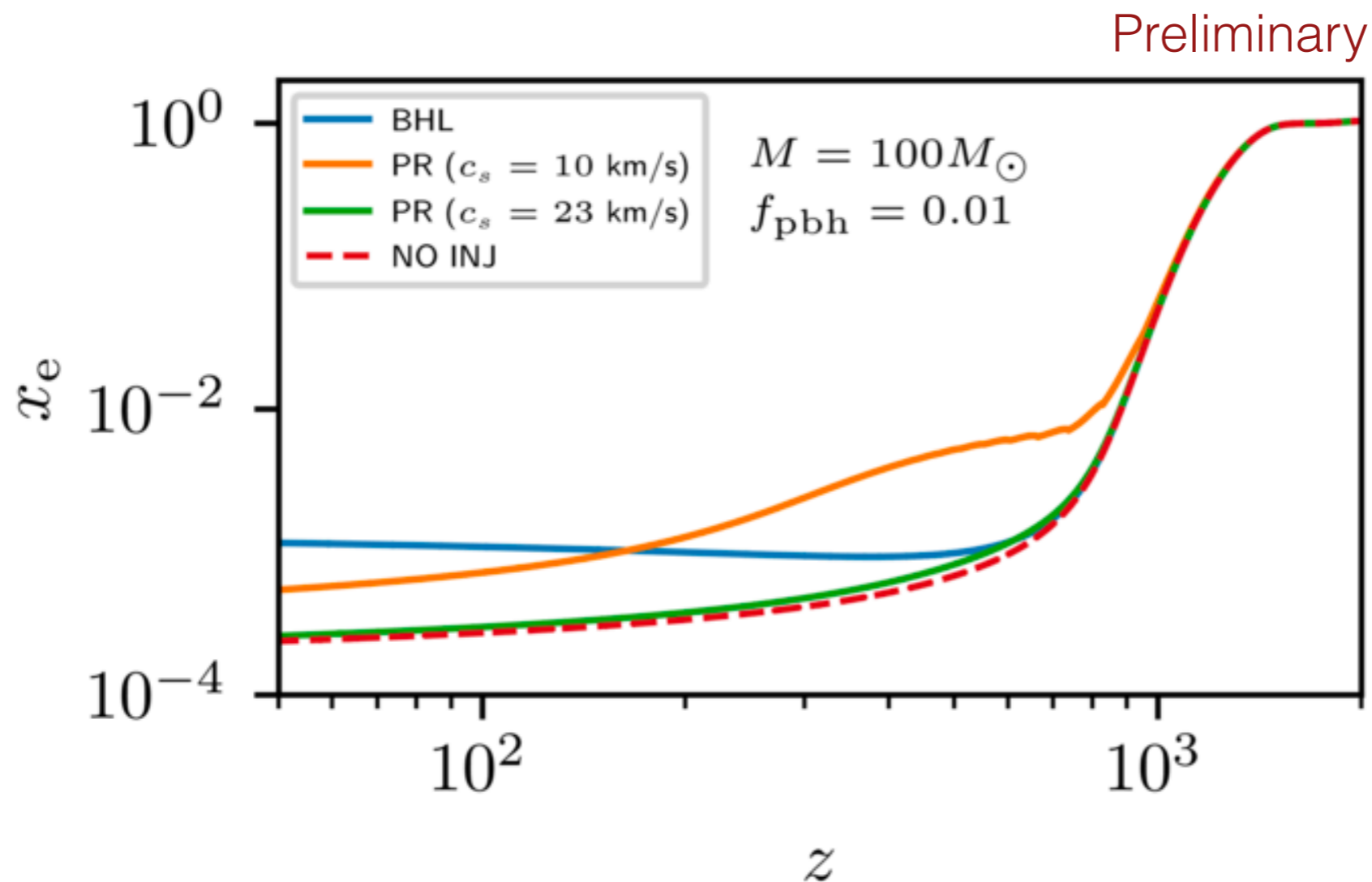
- Accretion rate **suppression** around PBHs is very relevant
- Dependence on the ionized sound speed
- May weaken the bound



Dominic Agius, DG, Rouven Essig, Francesca Scarcella, Gregory Suczewski, Mauro Valli, *in preparation*

Revisiting the Cosmological constraint: Results

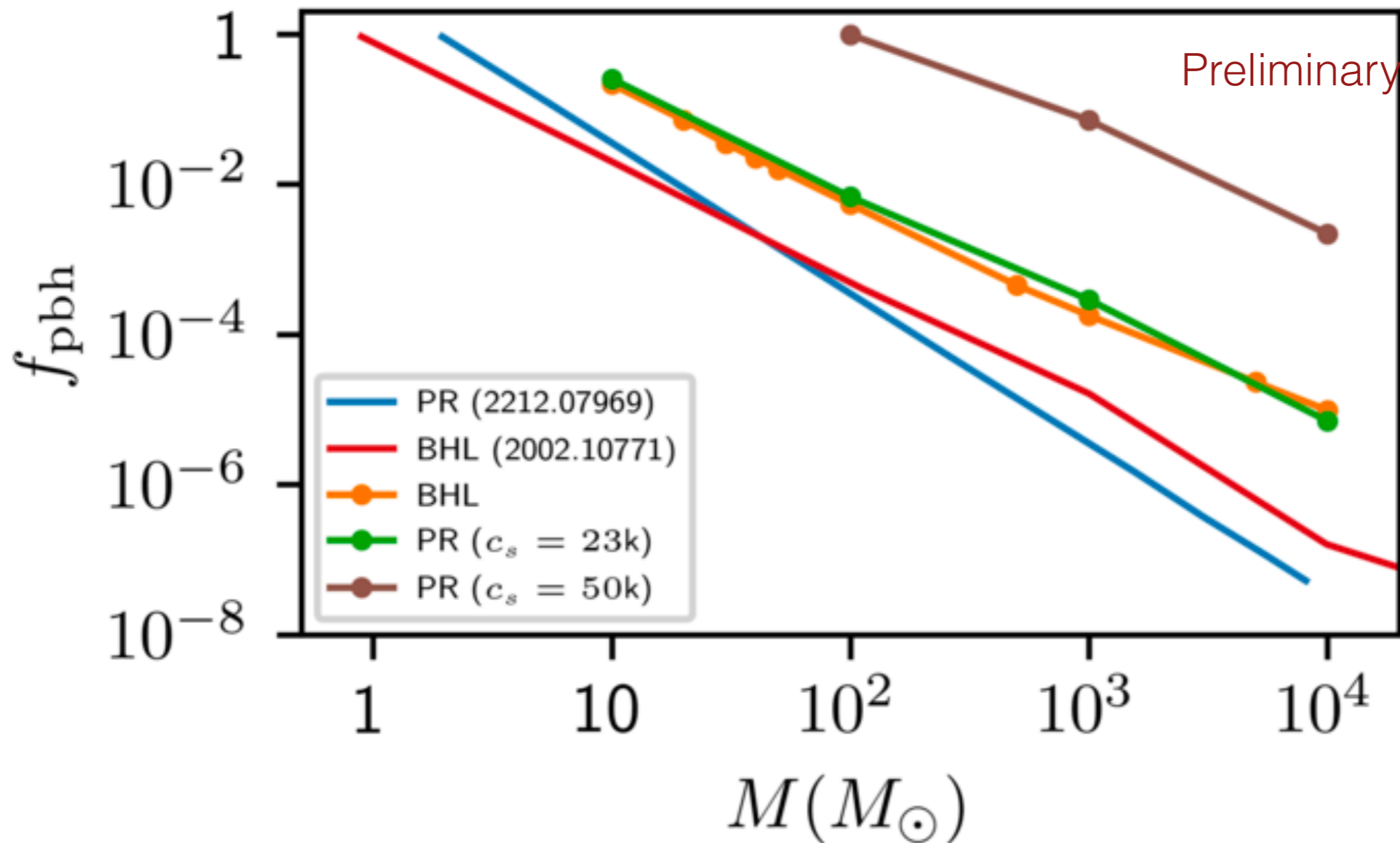
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Revisiting the Cosmological constraint

BHL vs PR: the “**Unexpected robustness**” of the bound



See also
Facchinetti+
2212.07969

Made with Cobaya+CLASS (modified to account for energy injection)
2018 low- l Planck TT.EE, high- l Planck TT.TE.EE, lensing, ACT, BAO

Other uncertainties to be assessed

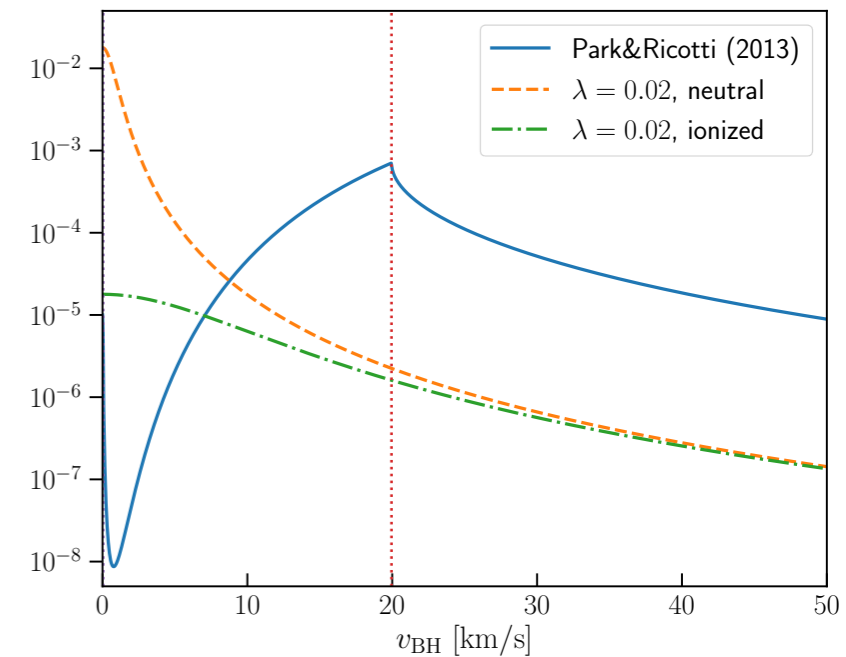
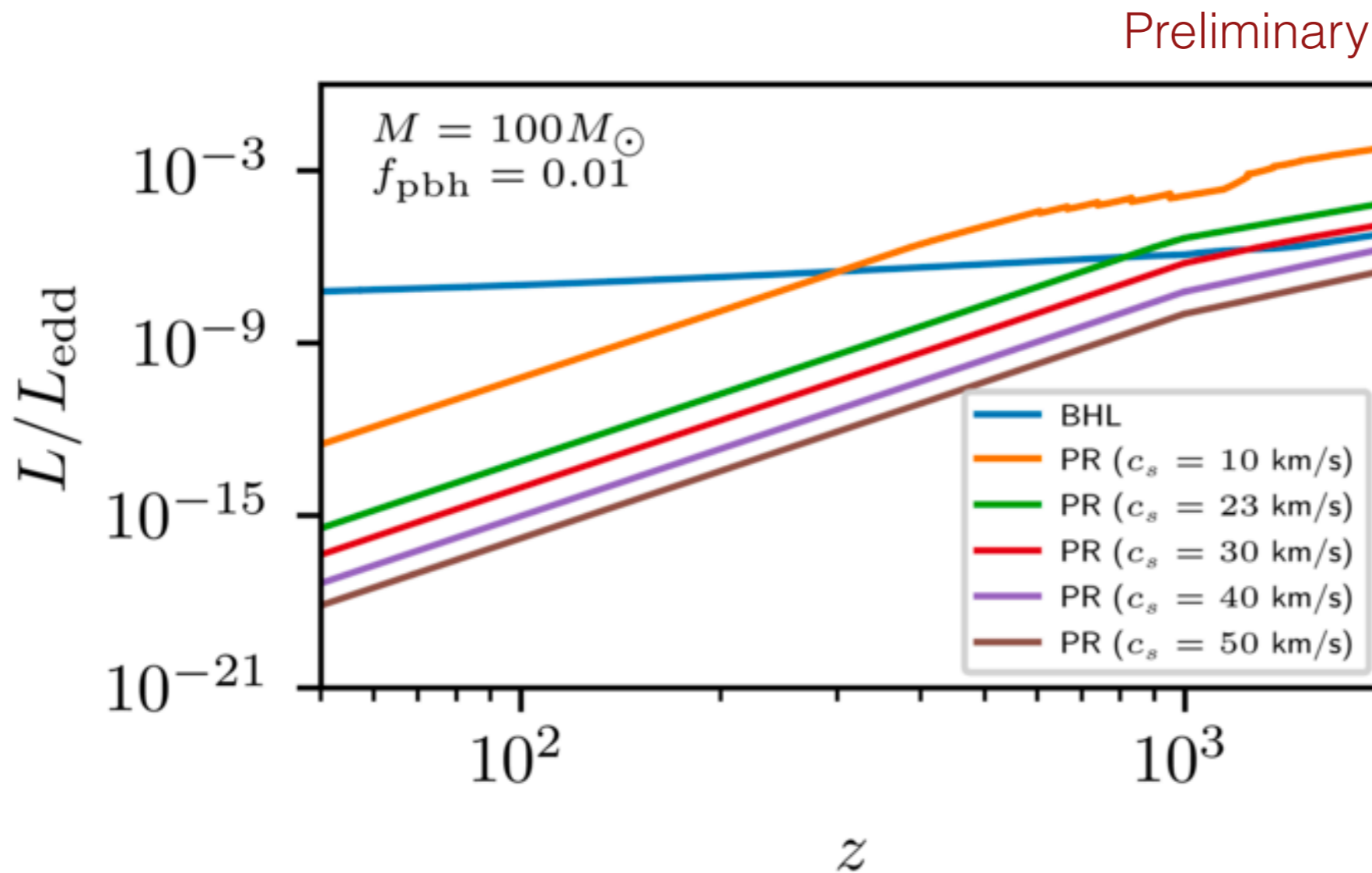
- Evolution of the DM-baryon velocity
(see for instance the discussion in [Poulin+ 1707.04206](#))
- Luminosity of the accretion disks around PBHs
- Energy deposition functions
([Galli+ 1306.0563](#); [Slatyer 1506.03812](#), [Liu+ 2303.07366](#))

Conclusions

- Revisitation of the **Cosmological bound on the PBH abundance**
 - Very strong at high masses.
 - Assessment of uncertainties is needed.
 - Different models of accretion predict very different time evolution of X_e .
 - The bound is not heavily affected.
 - More investigation on DM-baryon relative speed and radiative efficiency is needed.

Revisiting the Cosmological constraint: Results

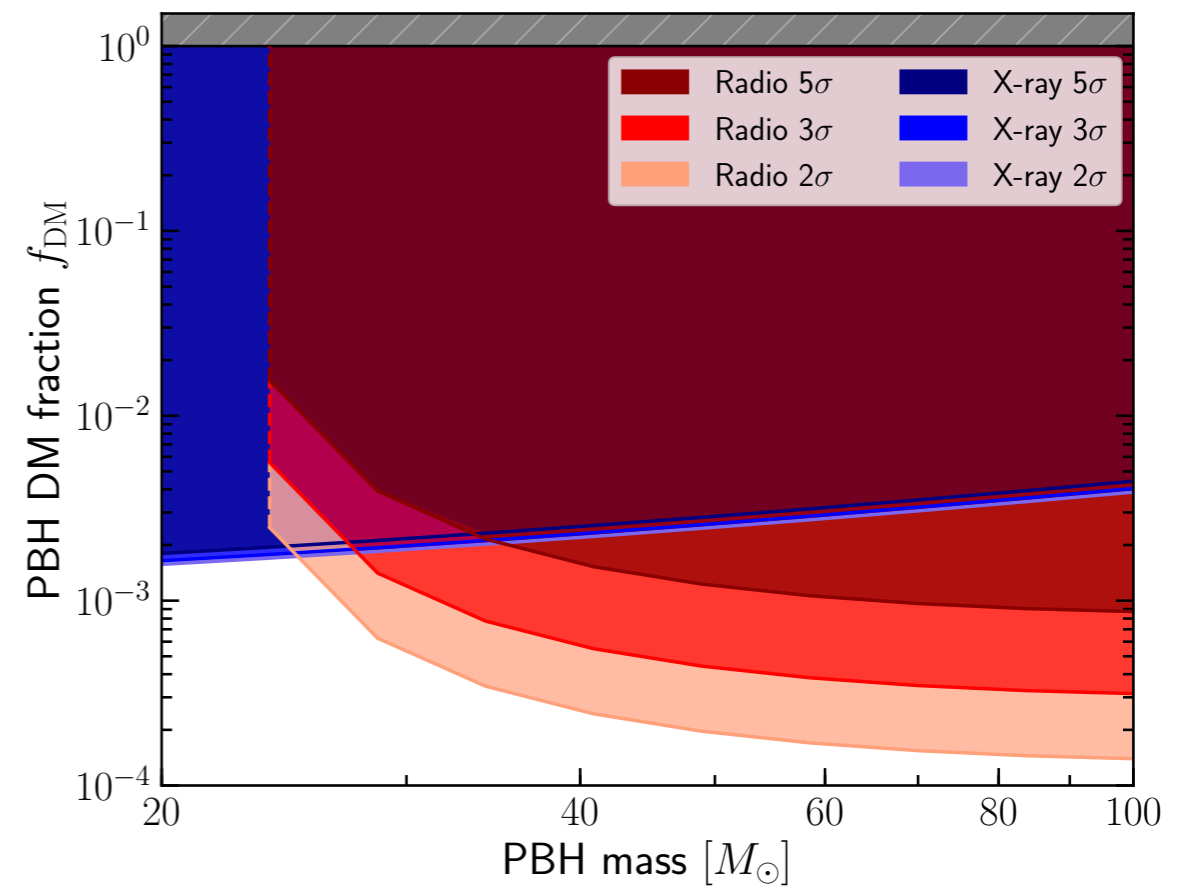
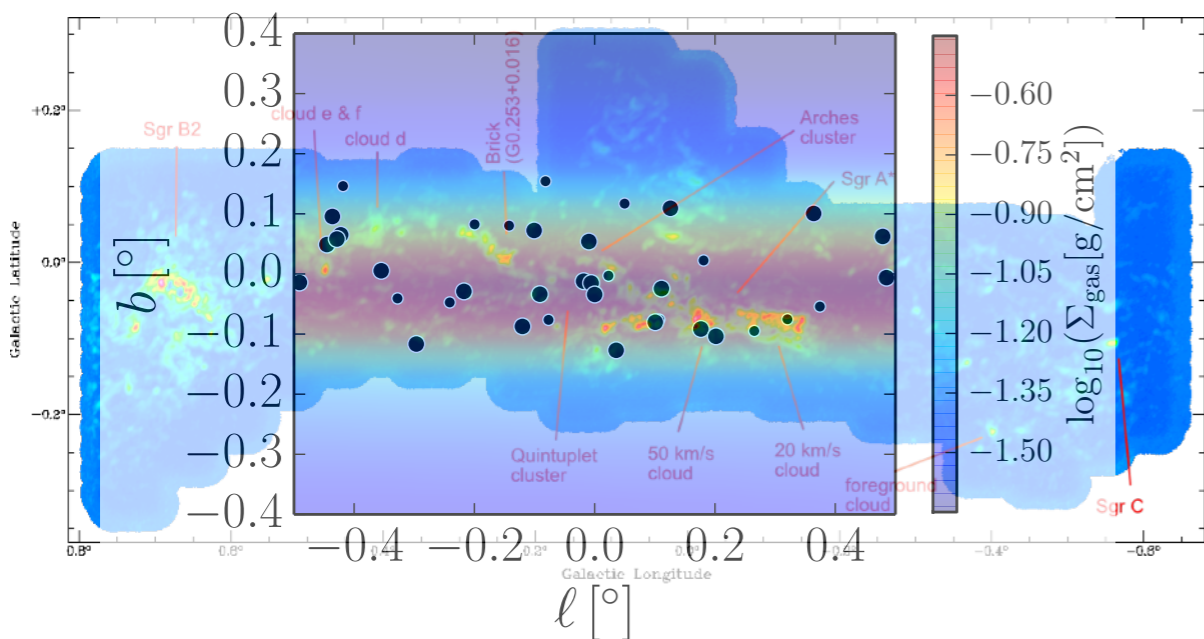
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Revisiting the Astronomical constraint

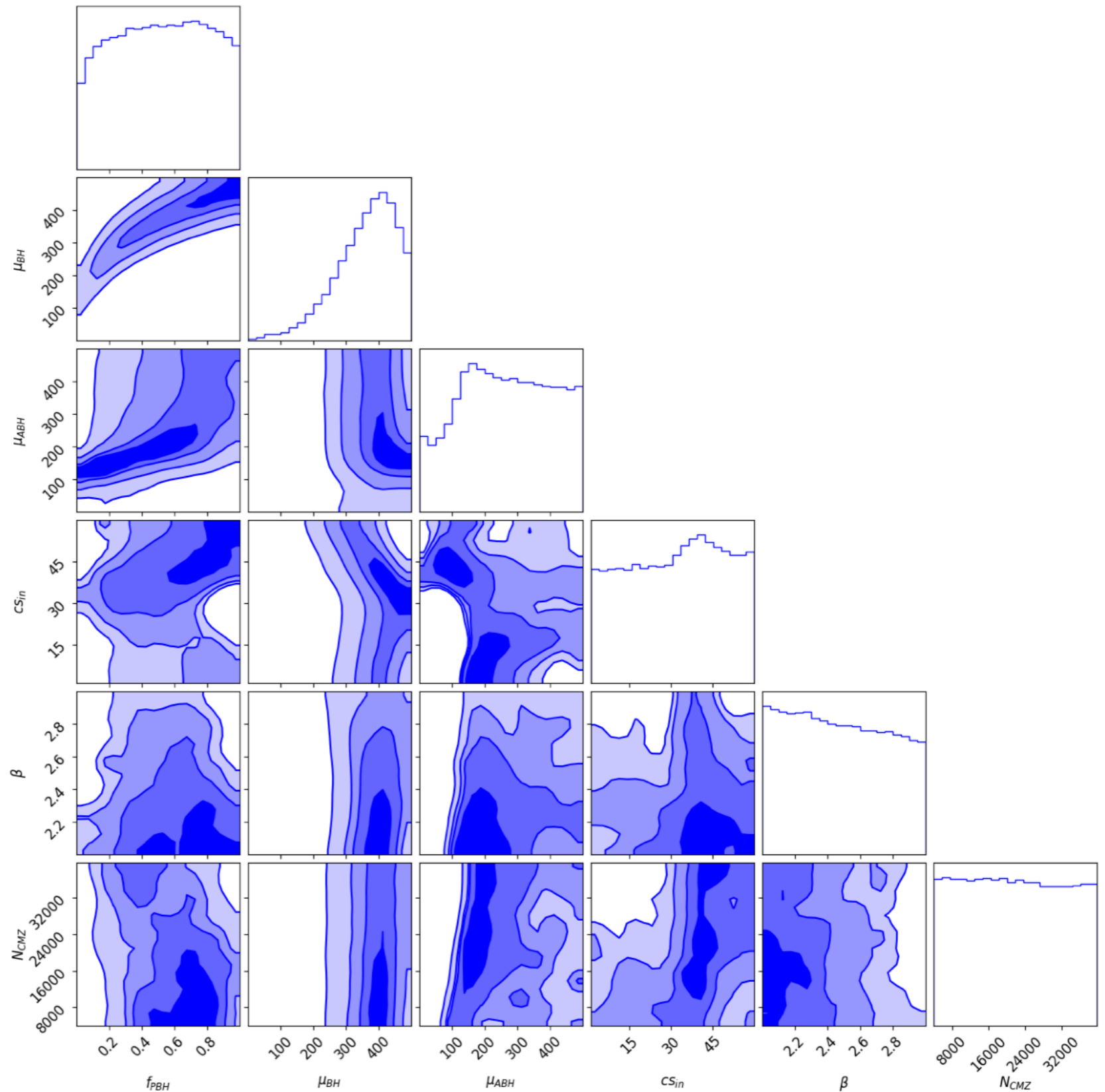
- Focus on Galactic Center region - rich of gas
- If PBHs exist, they accrete gas and emit
- We require not to overshoot data from X-ray and radio point source catalogues
- Both BHL and PR already considered



J. Manshanden, **DG**, R. Connors, G. Bertone, M. Ricotti, 1812.07967

Revisiting the Astronomical constraint

DG, Juan García-Bellido,
Francesca Scarcella, Tejas
Satheesh, *in preparation*



Revisiting the Astronomical constraint

