Introduction:
Science Case for the EIC

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Lecture 1 of 2

Introduction to EIC – two lectures

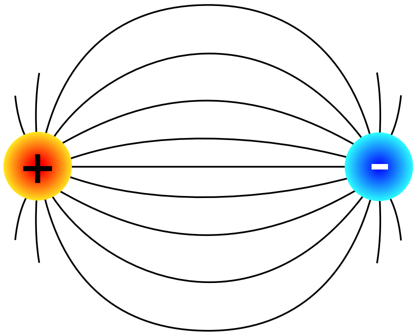
- Hour 1: Historical: Why EIC? –
  Limitations of past & current experiments

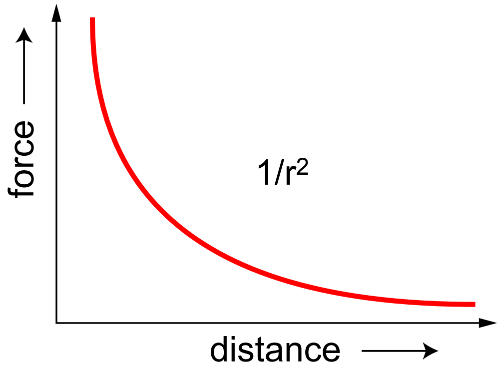
- Hour 2: Today: why EIC?
  What the EIC could deliver

Quantum Electrodynamics (QED)

Theory of electromagnetic interactions

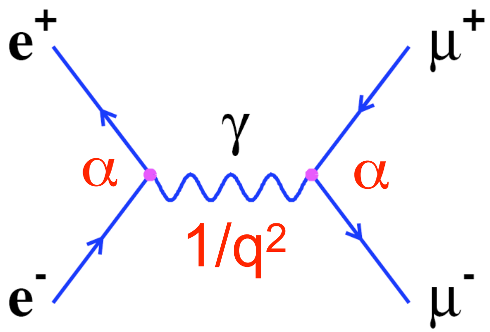
- Exchange particles (photons) do not carry electric charge
- Flux is not confined: $V(r) \sim 1/r$, $F(r) \sim 1/r^2$





$$V(r) = -\frac{q_1 \ q_2}{4\pi\varepsilon_0 \ r} = -\frac{\alpha_{em}}{r}$$

Example Feynman Diagram:
e^+e^- annihilation



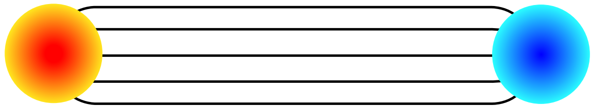
Coupling constant (α): Interaction Strength
In QED: $\alpha_{em} = 1/137$

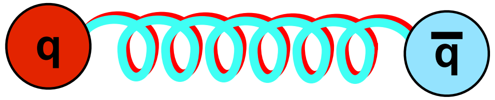
Quantum Chromodynamics (QCD)

Quantum Chromo Dynamics is the “nearly perfect” fundamental theory of the strong interactions

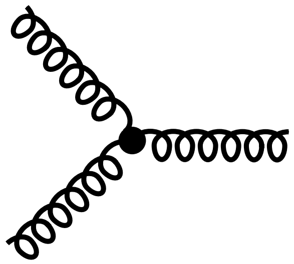
F. Wilczek, hep-ph/9907340

- Three color charges: red, green and blue





- Exchange particles (gluons) carry color charge and can self-interact



Self-interaction: QCD significantly harder to analyze than QED

- Flux is *confined*: $$V(r) = -\frac{4}{3}\frac{\alpha_s}{r} + kr$$

~1/r at short range

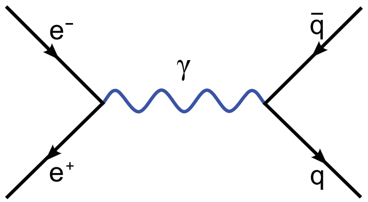
long range ~ r

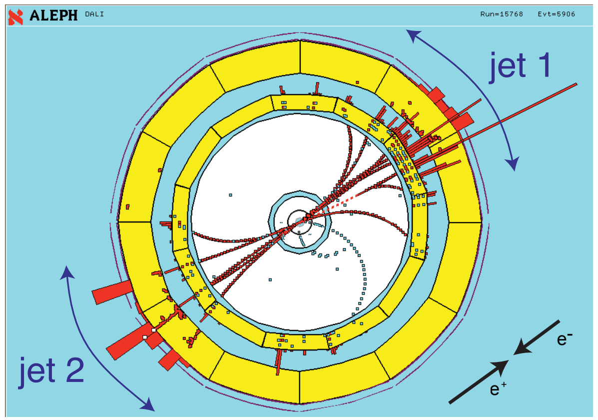
Long range aspect ⇒ quark confinement and existence of nucleons

Gluons!

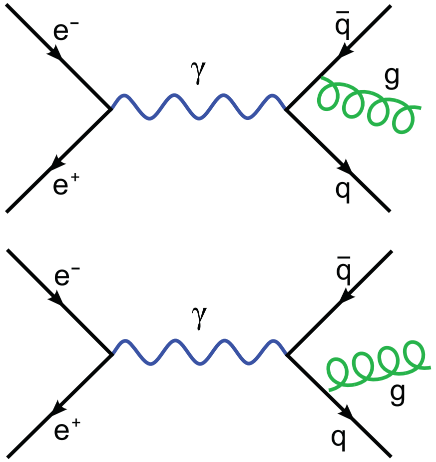
Discovery of gluons: Mark-J, Tasso, Pluto, Jade experiments at PETRA (e+e- collider) at DESY (CM energy 13-32 GeV)

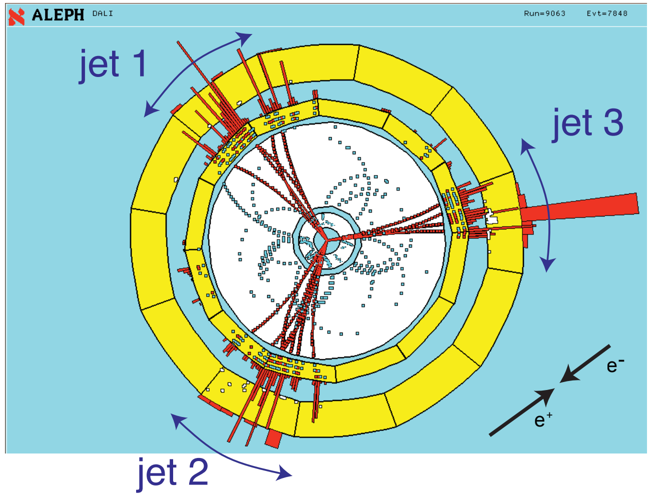
- $e^+ e^- \to q\ \bar{q} \to$ 2-jets





- $e^+ e^- \to q\ \bar{q}\ g \to$ 3-jets





Deep Inelastic Scattering (DIS)

Study of internal structure of a watermelon:

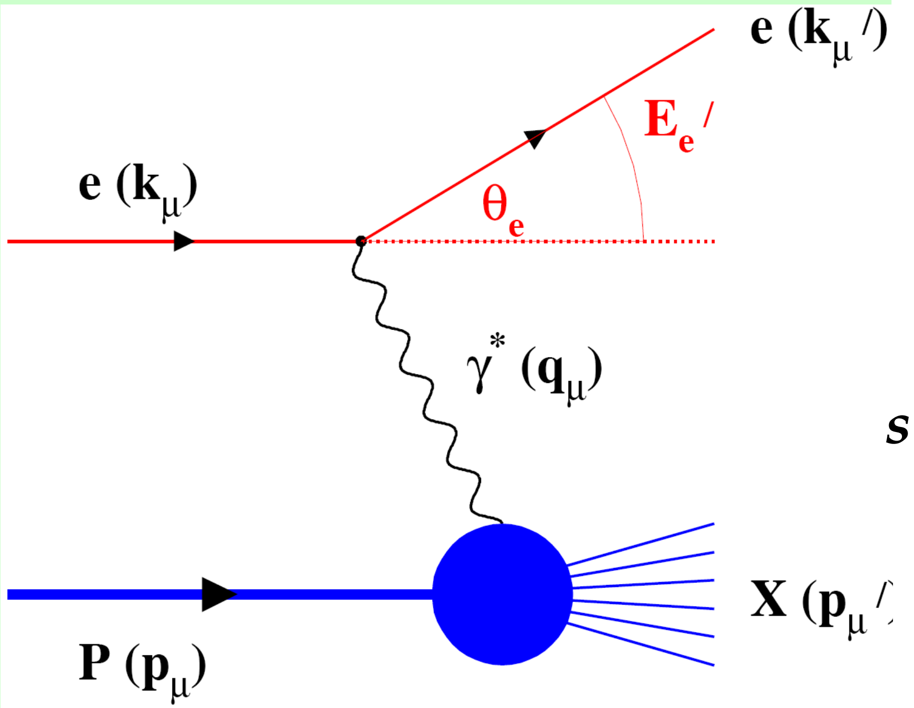
A-A (RHIC/LHC)

1) Violent collision of melons

2) Cutting the watermelon with a knife

Violent DIS e-A (Deep Inelastic Scattering -- DIS)

Deep Inelastic Scattering: Precision & Control



Kinematics:

$$s = 4\,E_e\,E_p$$

Inclusive events: e+p/A → e'+X

Semi-Inclusive events: e+p/A → e'+h(π,K,p,jet)+X

Exclusive events: e+p/A → e'+ p'/A'+ h(π,K,p,jet)

The only dimension considered comes in through "x".

Fraction of momentum carried by the quark/parton.

It is obviously moving in the direction of the proton. – Only one-dimensional information is explored & hence revealed….

All transverse motion was ignored. However, now we have more precision…

Unpolarized e-p/A DIS

DIS without Spin:

Inclusive Cross-Section:

$$\frac{d^2\sigma^{eA\to eX}}{dxdQ^2} = \frac{4\pi\alpha^2}{xQ^4}\left[\left(1-y+\frac{y^2}{2}\right)F_2(x,Q^2) - \frac{y^2}{2}F_L(x,Q^2)\right]$$

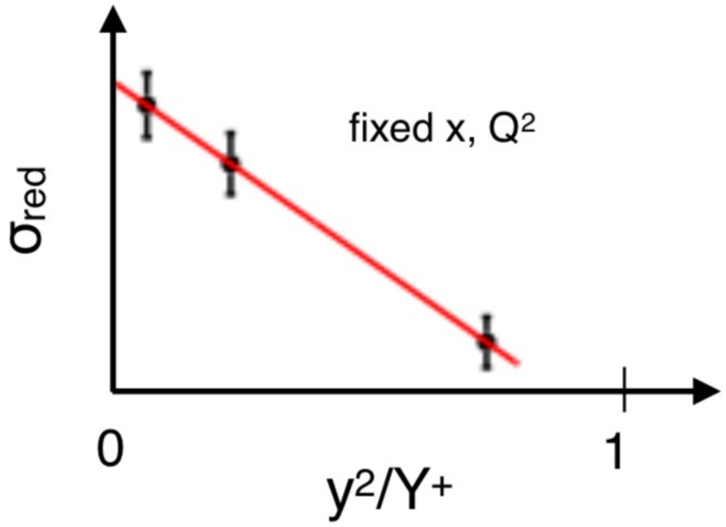
Reduced Cross-Section:

$$\sigma_r = \left(\frac{d^2\sigma}{dxdQ^2}\right)\frac{xQ^4}{2\pi\alpha^2[1+(1-y)^2]} = F_2(x,Q^2) - \frac{y^2}{1+(1-y)^2}F_L(x,Q^2)$$

$$\sigma_r(x,Q^2) = F_2^A(x,Q^2) - \frac{y^2}{Y+}F_L^A(x,Q^2)$$

Rosenbluth Separation:

- Recall $Q^2 = x\ y\ s$
- Measure at different $\sqrt{s}$
- Plot σ_{red} versus $y2/Y^+$ for fixed x, Q^2
- F_2 is σ_{red} at $y2/Y^+ = 0$
- F_L = Slope of $y2/Y^+$



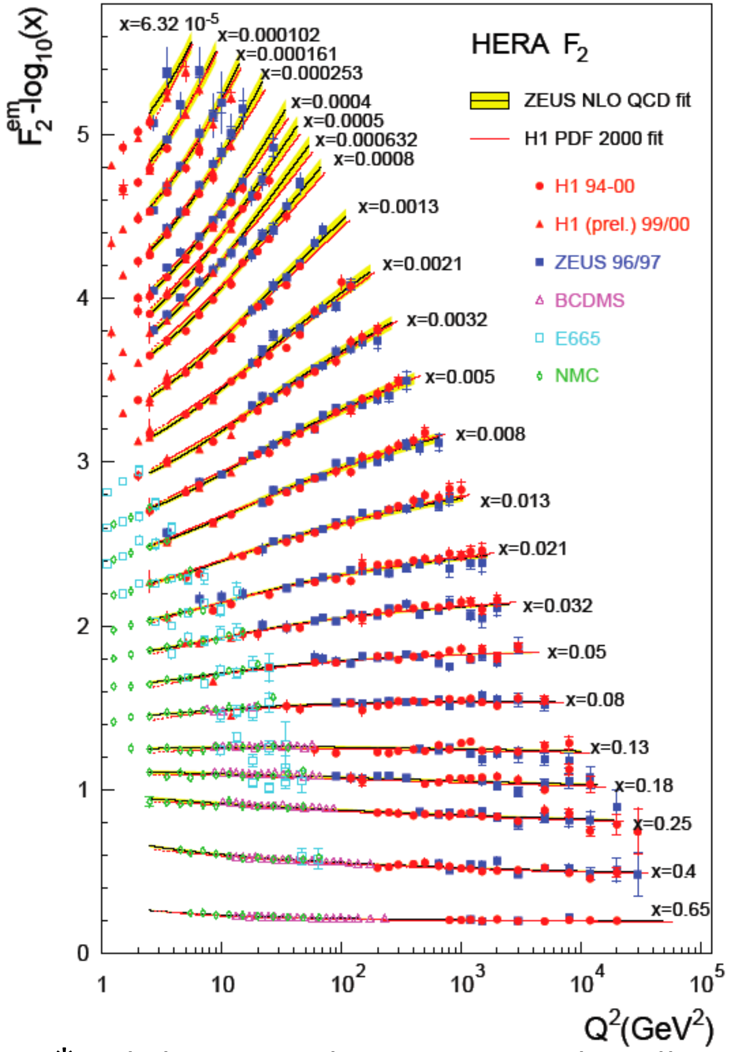
Recall from
Enrico Tassi's Lecture

Measurement of unpolarized glue at HERA

F_2 Structure Function

Vs.

Q^2 mom. exchanged

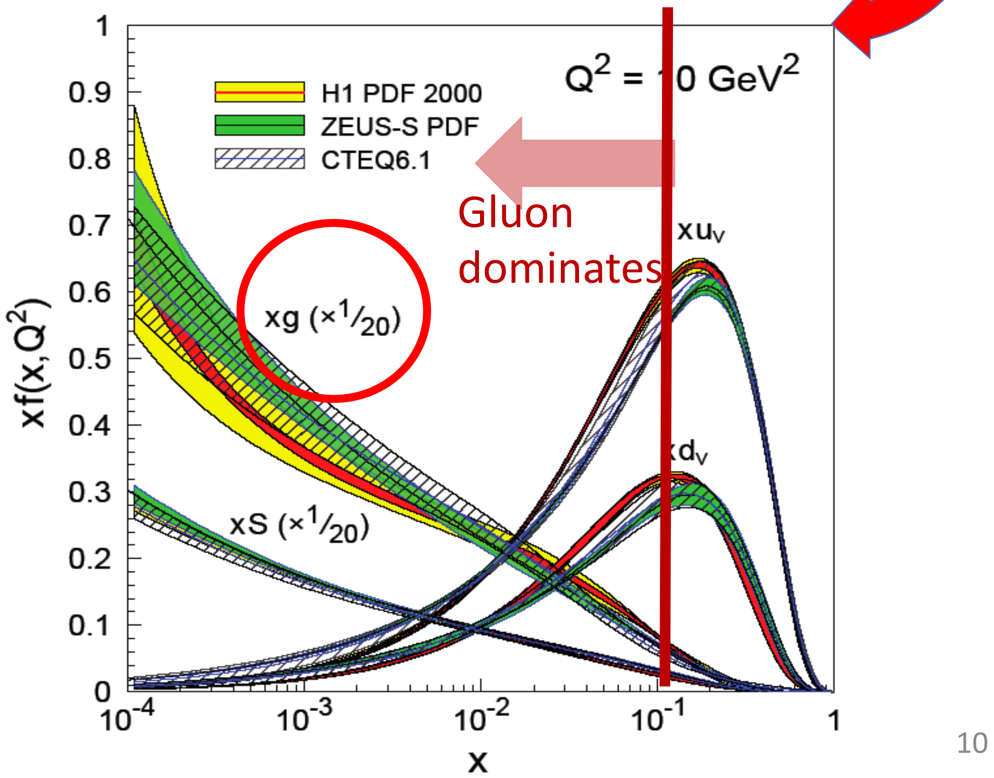


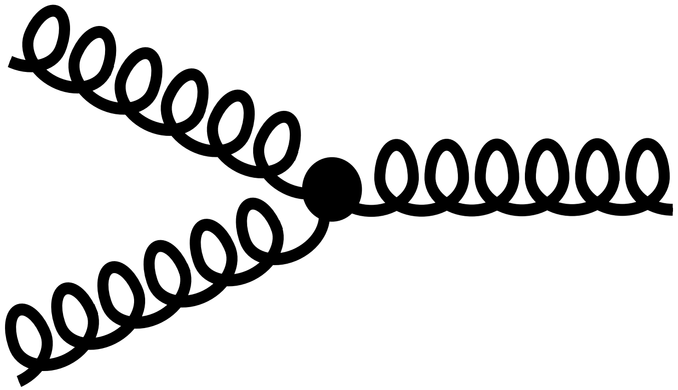
*Dokshitzer, Gribov, Lipatov, Altarelli, Parisi

- Scaling violations of $F_2(x,Q^2)$

$$\frac{\partial F_2(x,Q^2)}{\partial \ln Q^2} \propto G(x,Q^2)$$

- NLO pQCD analyses: fits with **linear** DGLAP* equations





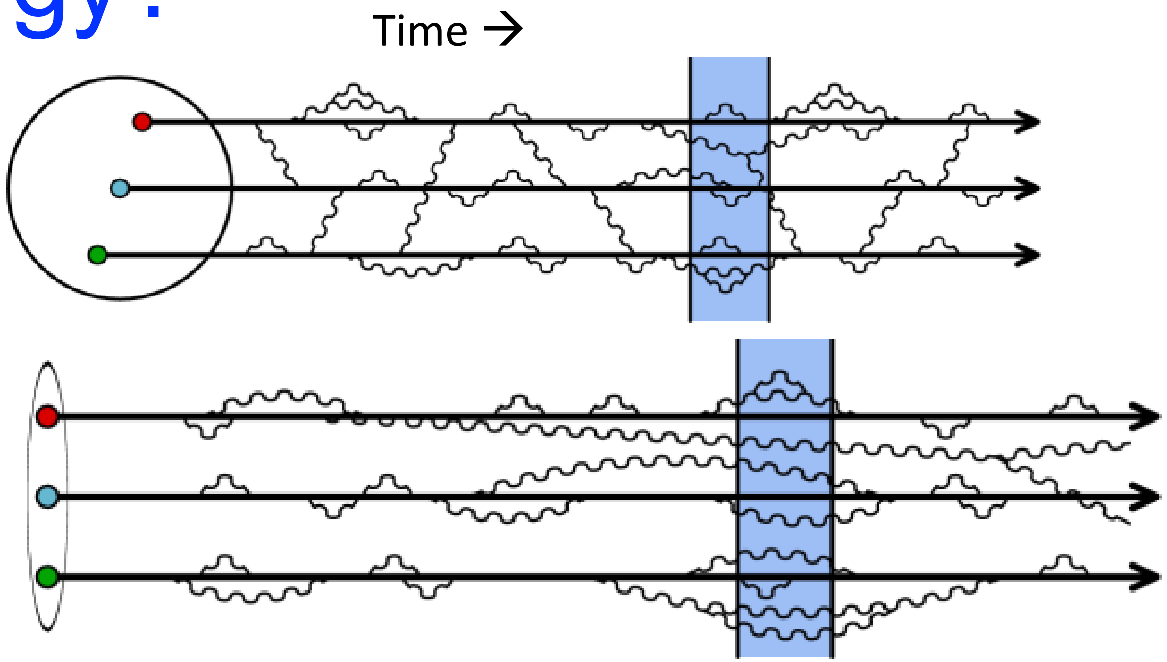
Consequence of gluon self interactions ➔ non-linear GDLAP evolution… ?

Particularly at high energy ➔ low-x

How does a Proton look at low and very high energy?

Low energy: High x
Regime of fixed target exp.

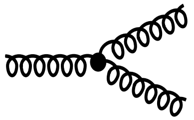
High energy: Low- x
Regime of a Collider



Cartoon of boosted proton

At high energy:

- Wee partons fluctuations are time dilated in strong interaction time scales
- Long lived gluons radiate smaller x gluons ➔ which in turn radiate more… a chain reaction leading to a **runaway growth?**

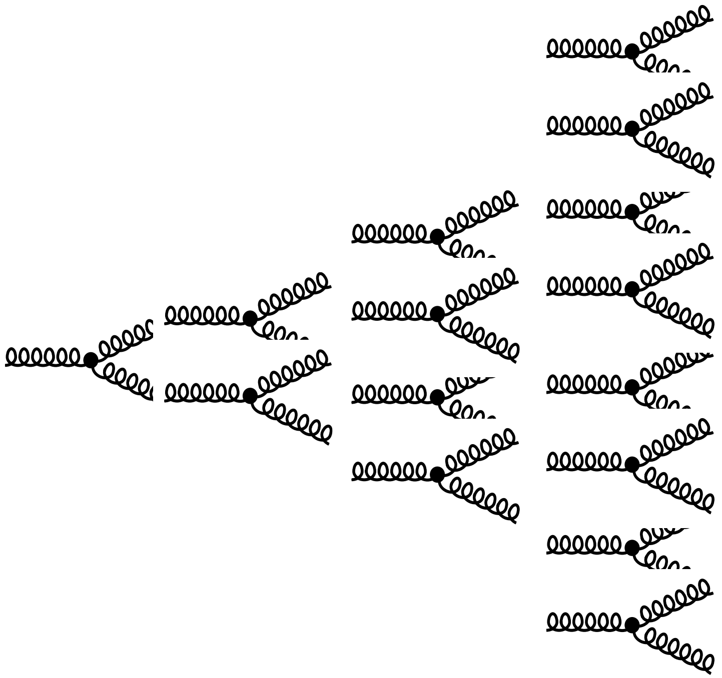

Gluon splitting

Gluon and the consequences of its interesting properties:

Gluons carry color charge ➔ Can interact with other gluons!

“…The result is a self catalyzing enhancement that leads to a runaway growth. A small color charge in isolation builds up a big color thundercloud….”

F. Wilczek, in “Origin of Mass”
Nobel Prize, 2004



? Infinity?
No!

We will come back to this a little later…

$$\frac{1}{2} = \left[\frac{1}{2}\Delta\Sigma + L_Q\right] + [\Delta g + L_G]$$

Quark Spin. Quark Ang. Mom. Gluon Spin. Gluon Ang. Mom

Spin Crisis

Life was easy in the Quark Parton Model until first spin experiments were performed!

“*spin* has killed more theories in physics than any other single observables”
-- Elliot Leader

“If theorists had their way, they would ban all experiments with Spin”
-- James D. Bjorken (jokingly)

Levitating top

Despite understanding gravity, and rotational motion individually, when combined it produces unexpected, unusual and interesting results.

In nature, we observe such things and try to understand the physics behind it.

Unpolarized e-p/A DIS

DIS without Spin:

Inclusive Cross-Section:

$$\frac{d^2\sigma^{eA\to eX}}{dxdQ^2} = \frac{4\pi\alpha^2}{xQ^4}\left[\left(1-y+\frac{y^2}{2}\right)F_2(x,Q^2) - \frac{y^2}{2}F_L(x,Q^2)\right]$$

Reduced Cross-Section:

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Rosenbluth Separation:

- Recall $Q^2 = x\,y\,s$
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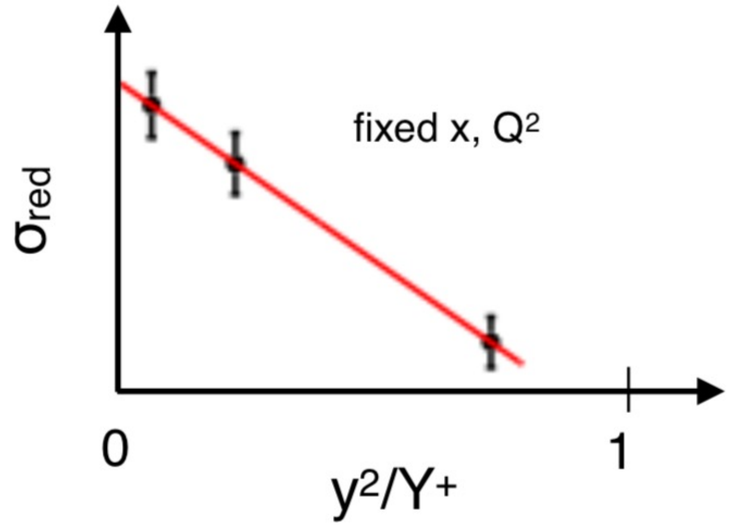
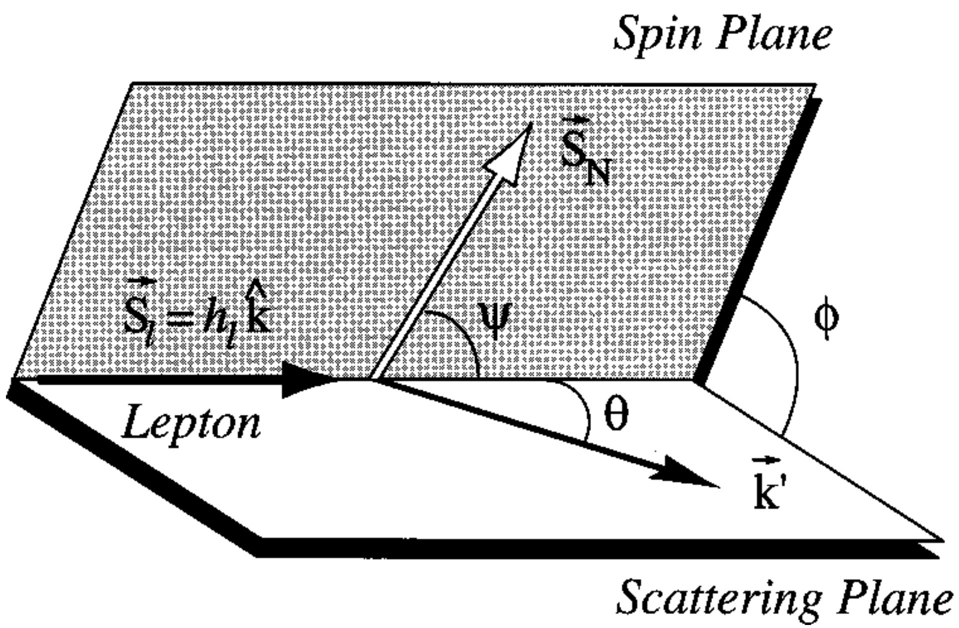
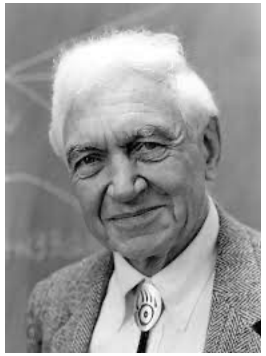


Lepton-nucleon cross section…*with spin*



$$\Delta\sigma = \cos\psi\,\Delta\sigma_{\parallel} + \sin\psi\,\cos\phi\,\Delta\sigma_{\perp}$$



V. W. Hughes
1922-2003

$$\gamma = \frac{2Mx}{\sqrt{Q^2}} = \frac{\sqrt{Q^2}}{\nu}.$$

For high energy scattering γ is small

$$\frac{d^2\Delta\sigma_{\parallel}}{dx\,dQ^2} = \frac{16\pi\alpha^2 y}{Q^4}\left[\left(1 - \frac{y}{2} - \frac{\gamma^2 y^2}{4}\right) g_1 - \frac{\gamma^2 y}{2}\, g_2\right]$$

$$\frac{d^3\Delta\sigma_T}{dx\,dQ^2\,d\phi} = -\cos\phi\,\frac{8\alpha^2 y}{Q^4}\,\gamma\sqrt{1 - y - \frac{\gamma^2 y^2}{4}}\left(\frac{y}{2}\, g_1 + g_2\right)$$

Cross section asymmetries….

- $\Delta\sigma_{||} =$ anti-parallel – parallel spin cross sections
- $\Delta\sigma_{perp}$= lepton-nucleon spins orthogonal
- Instead of measuring cross sections, it is prudent to measure the differences: Asymmetries in which many measurement imperfections might cancel:

$$A_{||} = \frac{\Delta\sigma_{||}}{2\bar{\sigma}}, \quad A_{\perp} = \frac{\Delta\sigma_{\perp}}{2\bar{\sigma}},$$

which are related to virtual photon-proton asymmetries A_1, A_2:

$$A_{||} = D(A_1 + \eta A_2), \quad A_{\perp} = d(A_2 - \xi A_1)$$

$$A_1 = \frac{\sigma_{1/2} - \sigma_{3/2}}{\sigma_{1/2} + \sigma_{3/2}} = \frac{g_1 - \gamma^2 g_2}{F_1}$$

$$A_2 = \frac{2\sigma^{TL}}{\sigma_{1/2} + \sigma_{3/2}} = \gamma\,\frac{g_1 + g_2}{F_1}$$

First Moments of SPIN SFs

$$\Delta q = \int_0^1 \Delta q(x)dx \qquad g_1(x) = \frac{1}{2}\Sigma_f e_f^2\{q_f^+(x) - q_f^-(x)\} = \frac{1}{2}\Sigma_f e_f^2 \Delta q_f(x)$$

$$\Gamma_1^p = \frac{1}{2}\left[\frac{4}{9}\Delta u + \frac{1}{9}\Delta d + \frac{1}{9}\Delta s\right] = \frac{1}{12}\underbrace{(\Delta u - \Delta d)}_{a_3=g_a} + \frac{1}{36}\underbrace{(\Delta u + \Delta d - 2\Delta s)}_{a_8} + \frac{1}{9}\underbrace{(\Delta u + \Delta d + \Delta s)}_{a_0}$$

$a_3=g_a$: Neutron decay

a_8: (3F-D)/3 Hyperon Decay

a_0: $\Delta\Sigma$

$$\Gamma_1^{p,n} = \frac{1}{12}\left[\pm a_3 + \frac{1}{\sqrt{3}}a_8\right] + \frac{1}{9}a_0$$

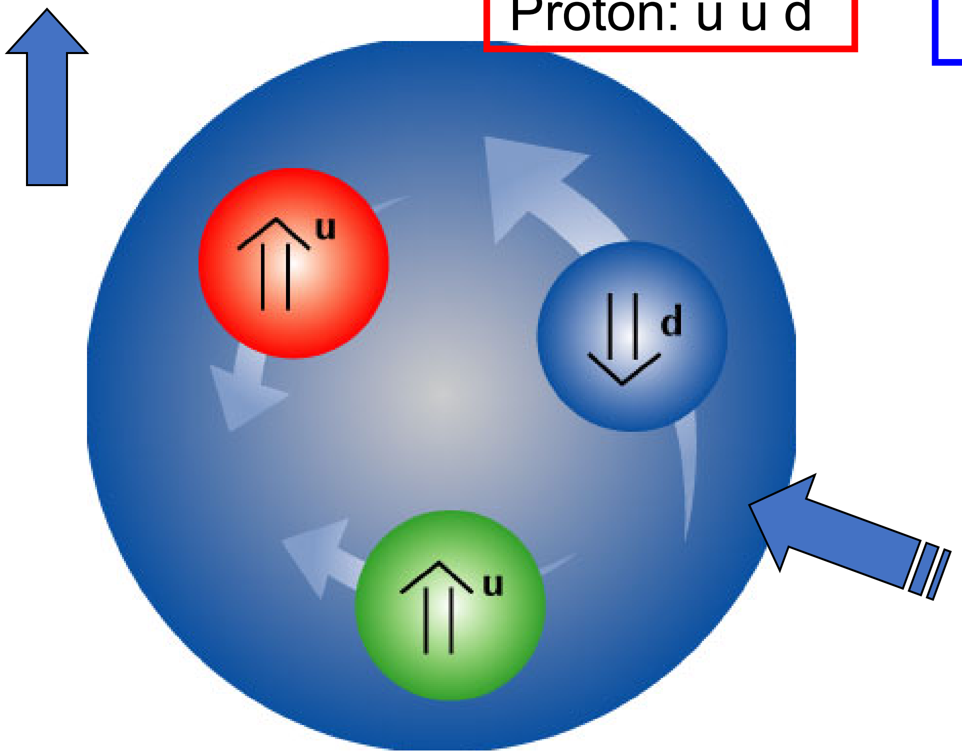
The measurements of spin

Nucleon's Spin: Naïve Quark Parton Model (ignoring relativistic effects… now, illustration only, but historically taken seriously)

- Protons and Neutrons are spin 1/2 particles
- Quarks that constitute them are also spin 1/2 particles
- And there are three of them in the

Proton: u u d

Neutron: u d d



S proton = Sum of all quark spins!

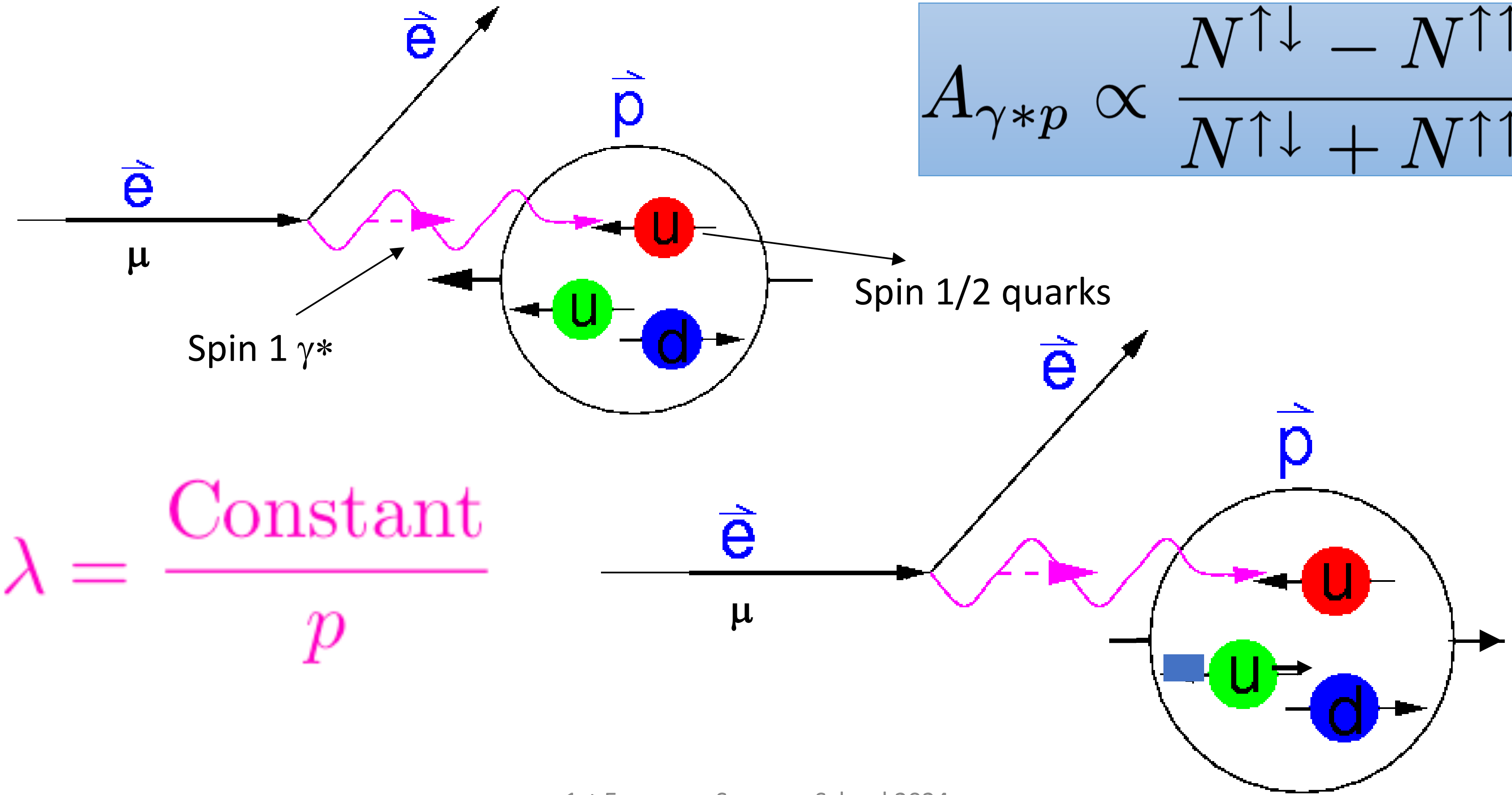
$$\frac{1}{2} \stackrel{?}{=} \frac{1}{2} + \frac{1}{2} + \frac{1}{2}$$

$$\frac{1}{2} = \frac{1}{2} - \frac{1}{2} + \frac{1}{2}$$

How was the Quark Spin measured?

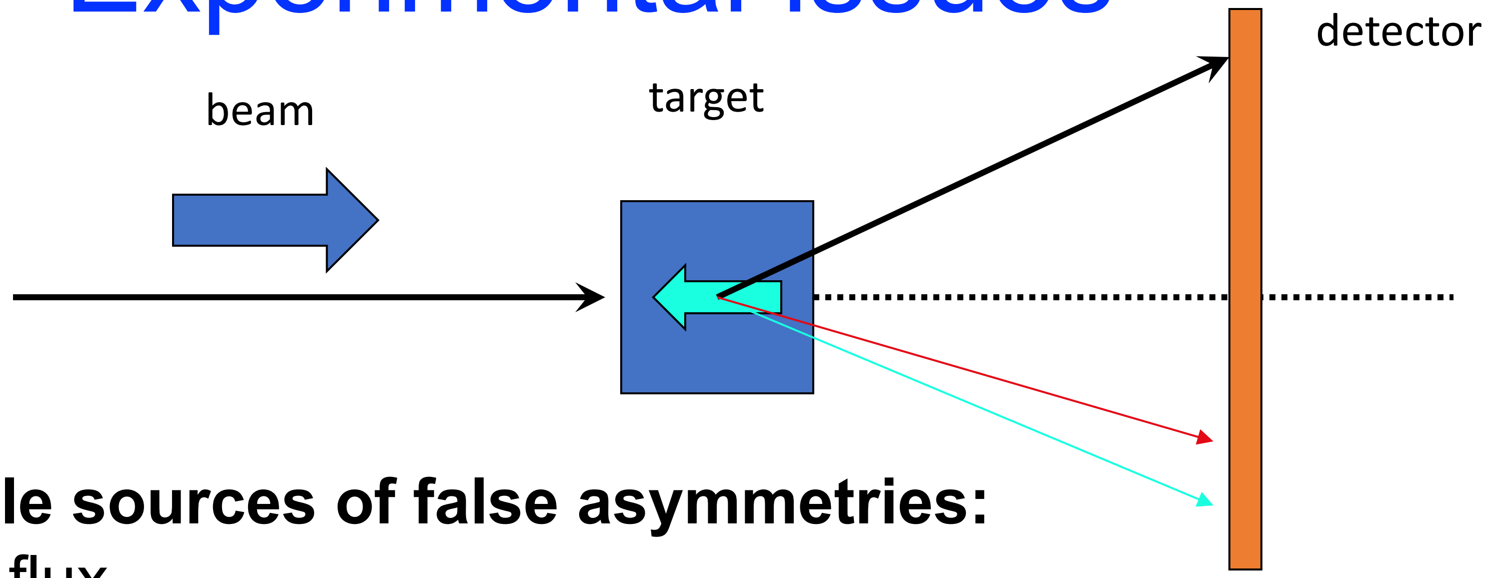
- Deep Inelastic polarized electron or muon scattering

$$A_{\gamma * p} \propto \frac{N^{\uparrow\downarrow} - N^{\uparrow\uparrow}}{N^{\uparrow\downarrow} + N^{\uparrow\uparrow}}$$



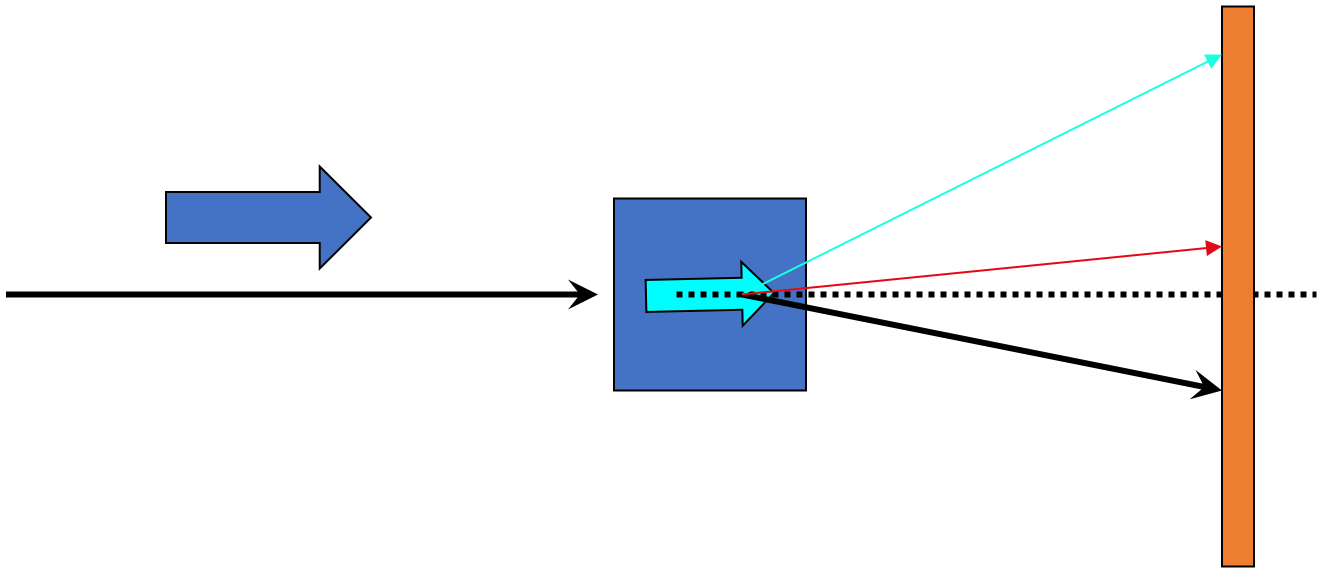
$$\lambda = \frac{\text{Constant}}{p}$$

Experimental issues



Possible sources of false asymmetries:

- beam flux
- target size
- detector size
- detector efficiency



$A_{measured} = A_{LL}$
Double Longitudinal Spin asymmetry

$$A_{measured} = \frac{N^{\rightarrow\leftarrow} - N^{\rightarrow\rightarrow}}{N^{\rightarrow\leftarrow} + N^{\rightarrow\rightarrow}}$$

$$N^{\leftarrow\rightarrow} = N_b \cdot N_t \cdot \sigma^{\leftarrow\rightarrow} \cdot D_{acc} \cdot D_{eff}$$

$$N^{\rightarrow\rightarrow} = N_b \cdot N_t \cdot \sigma^{\rightarrow\rightarrow} \cdot D_{acc} \cdot D_{eff}$$

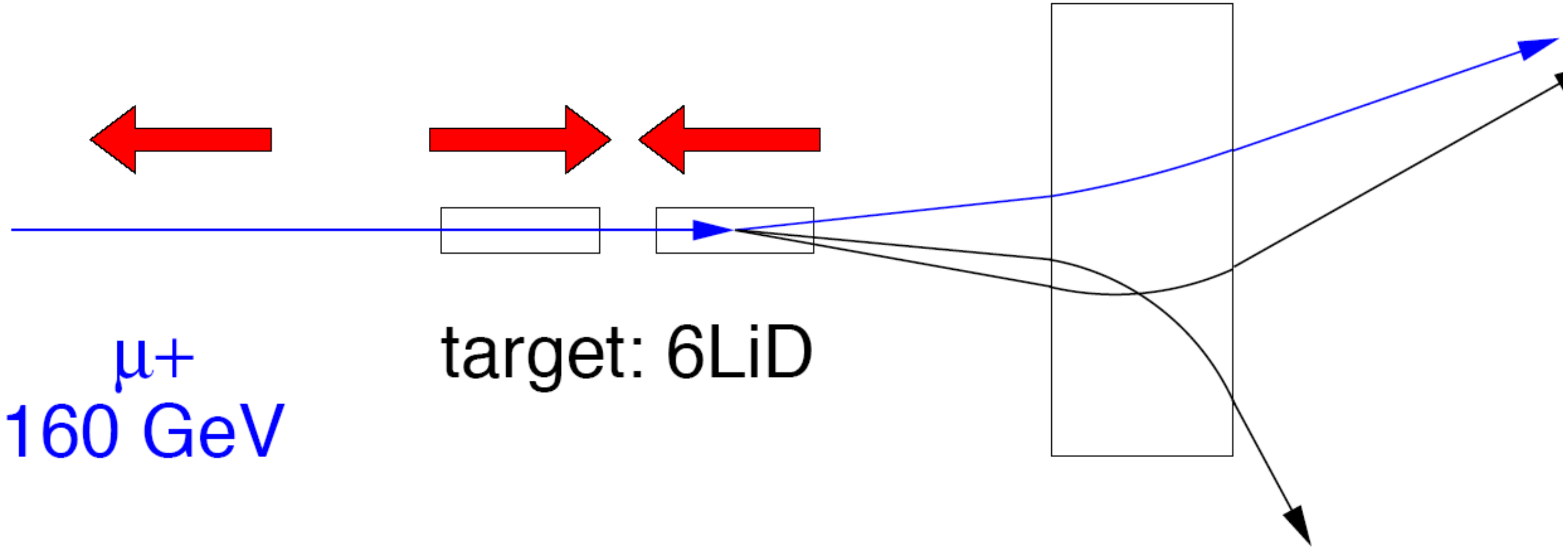
If all other things are equal,

they cancel in the ratio

$$A_{measured} = \frac{\sigma^{\rightarrow\leftarrow} - \sigma^{\rightarrow\rightarrow}}{\sigma^{\rightarrow\leftarrow} + \sigma^{\rightarrow\rightarrow}}$$

A Typical Setup

- Experiment setup (EMC, SMC, COMPASS@CERN)



- Target polarization direction reversed every 6-8 hrs
- Typically experiments try to limit false asymmetries to be about 10 times smaller than the physics asymmetry of interest

Experimental Needs in DIS

Polarized target, polarized beam

- Polarized targets: hydrogen (p), deuteron (pn), helium (^{3}He: 2p+n)
- Polarized beams: electron,muon used in DIS experiments

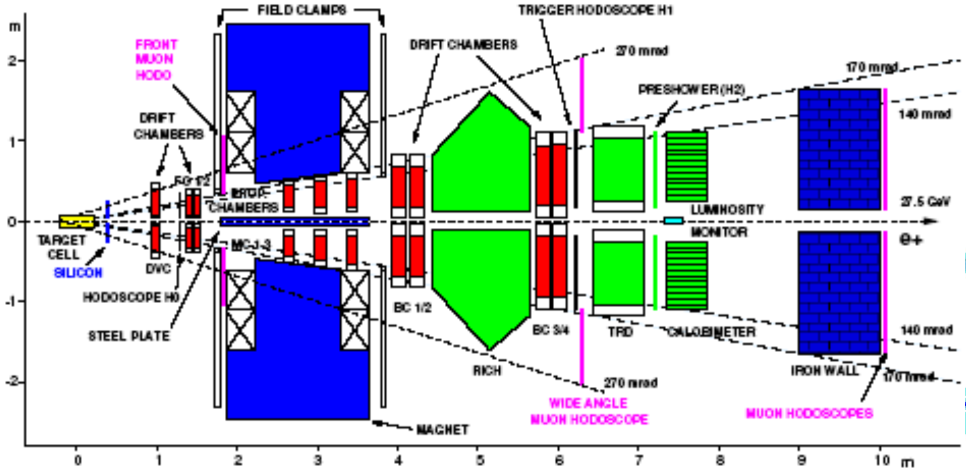
Determine the kinematics: measure with high accuracy:

- Energy of **incoming lepton**
- Energy, direction of **scattered lepton**: energy, direction
- Good identification of **scattered lepton**

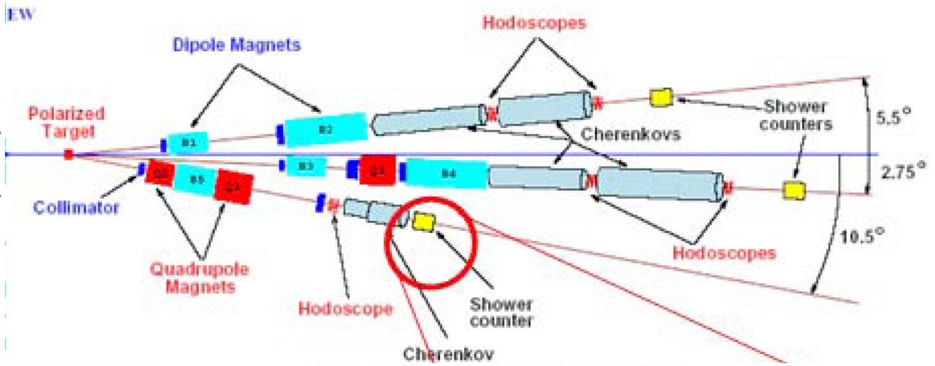
Control of false asymmetries:

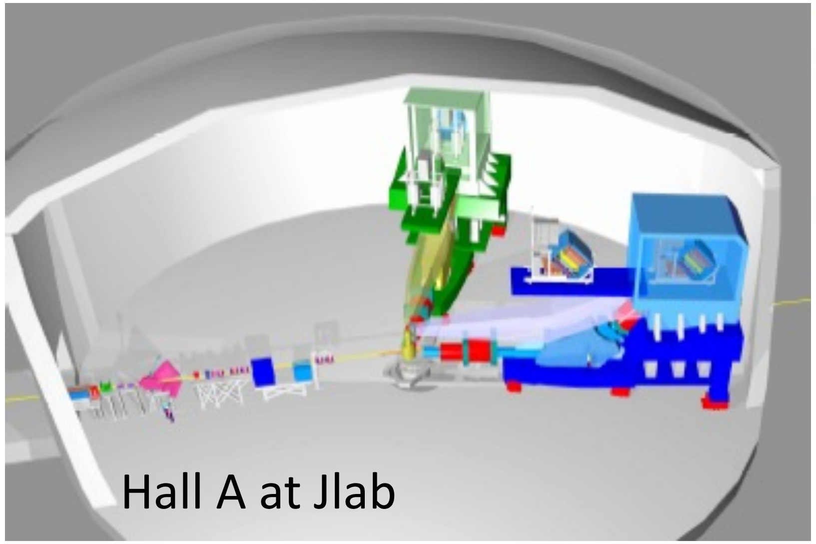
- Need excellent understanding and control of **false asymmetries** (time variation of the detector efficiency etc.)

Experiments

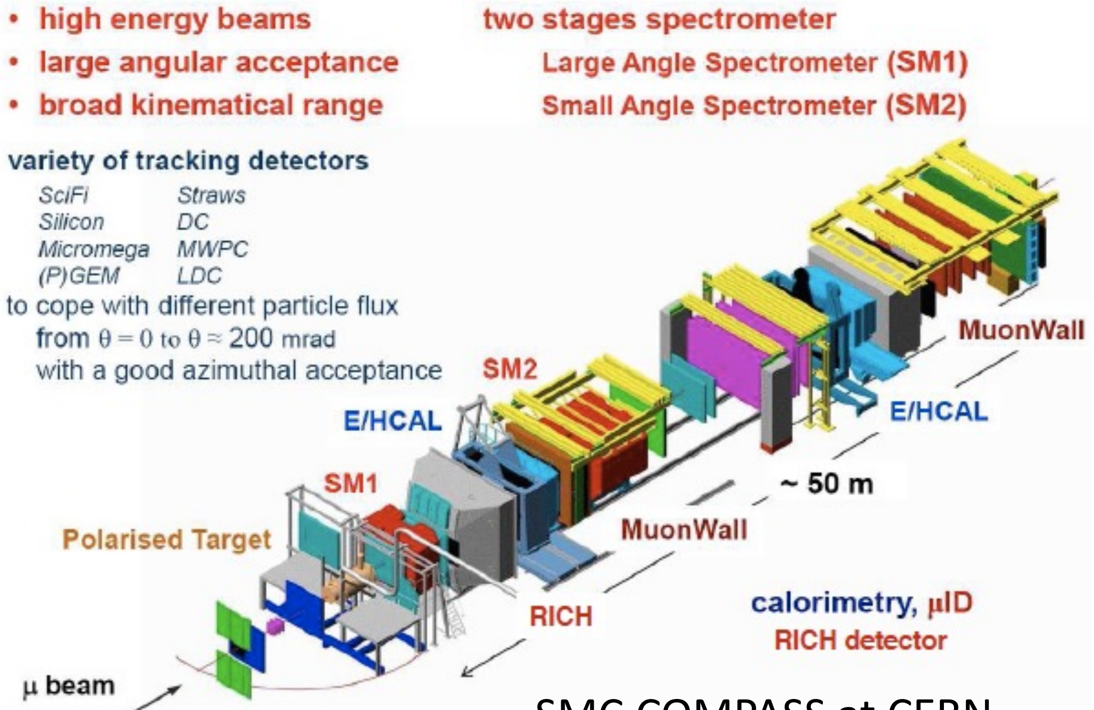


HERMES at DESY





Hall A at Jlab



SMC,COMPASS at CERN

First moment of $g_1^p(x)$: Ellis-Jaffe SR

$$\Gamma_1^{p,n} = \frac{1}{12}\left[\pm a_3 + \frac{1}{\sqrt{3}} a_8\right] + \frac{1}{9} a_0$$

$$a_3 = \frac{g_A}{g_V} = F + D = 1.2601 \pm 0.0025 \qquad a_8 = 3F - D \Longrightarrow F/D = 0.575 \pm 0.016$$

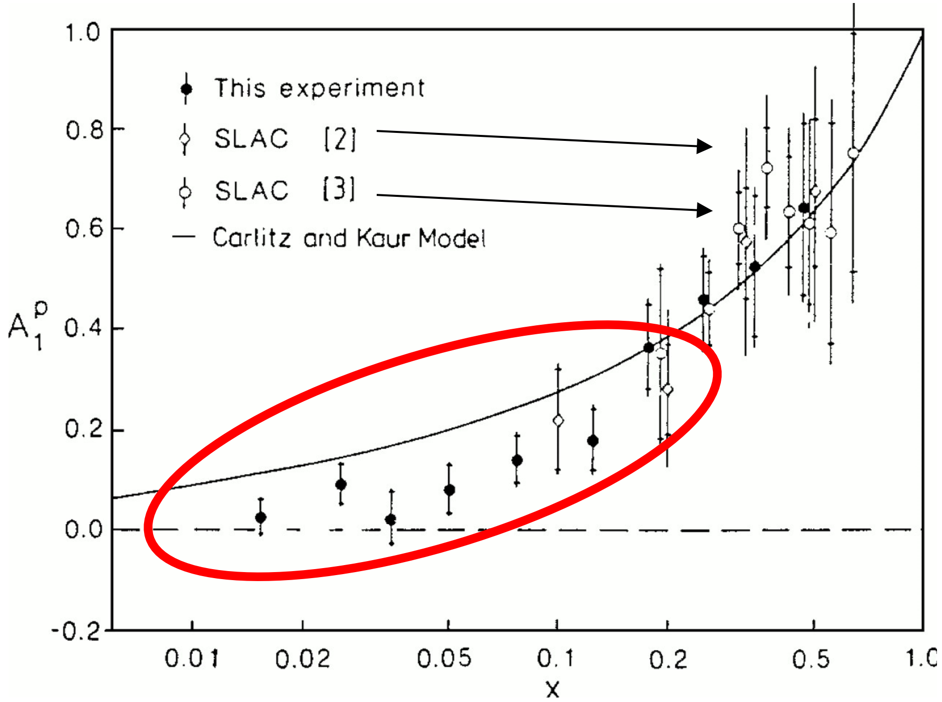
Assuming $SU(3)_f$ & $\Delta s = 0$, Ellis & Jaffe: $\Gamma_1^p = 0.170 \pm 0.004$

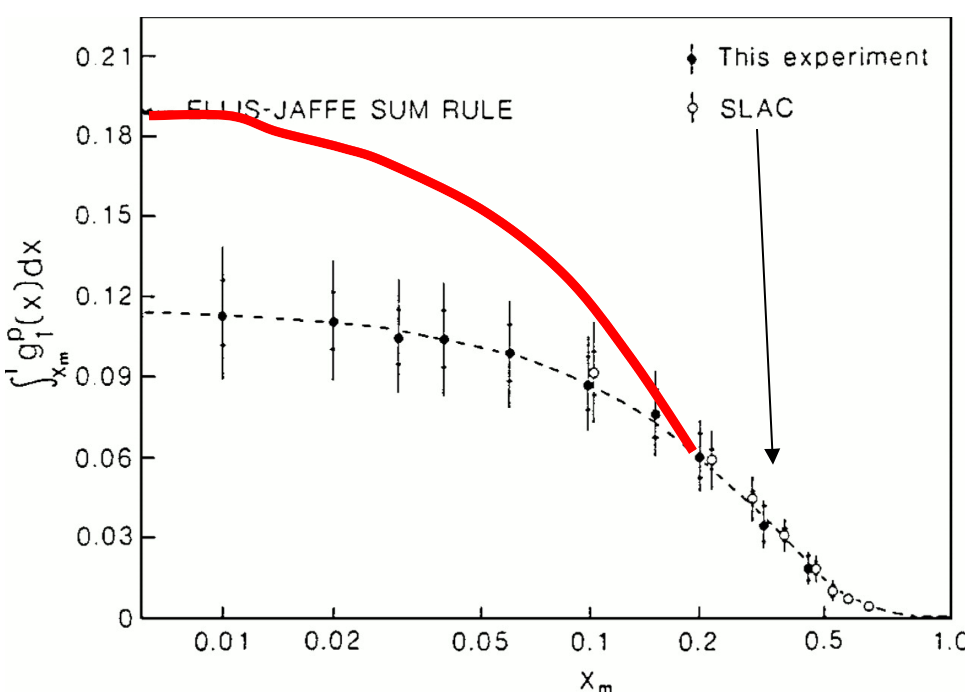
Measurements were done at SLAC (E80, E130) Experiments:

- Low 8-20 GeV electron beam on fixed target
- Did not reach low enough x → $x_{min} \sim 10^{-2}$
- Found consistency of data and E-J sum rule above

Proton Spin Crisis (1989)!

EMC experiment at CERN: high energy muon beam – reached lower x





$\Delta\Sigma/2 = (0.12)$ +/- (0.17) (EMC, 1989)

$\Delta\Sigma/2 = 0.58$ expected from E-J sum rule....

If the quarks did not carry the nucleon's spin, what did? → Gluons?

Consequence:

- Quark (+anti-quark) contribution to nucleon spin is definitely small:

$$\text{Quark's contribution to nucleon spin} \rightarrow \frac{1}{2}\Delta\Sigma = 0.15 \pm 0.03$$

  - Is this smallness due to some cancellation between quark & anti-quark polarization?

- Or does glue makes a very large contribution? $\Delta G = 1 \pm 1.5$

- Most NLO analyses by consistent with HIGH gluon contribution
  - Direct measurement of gluon spin with other probes warranted.
  - Seeded the RHIC Spin program

Measurement of unpolarized glue at HERA

- Scaling violations of $F_2(x,Q^2)$

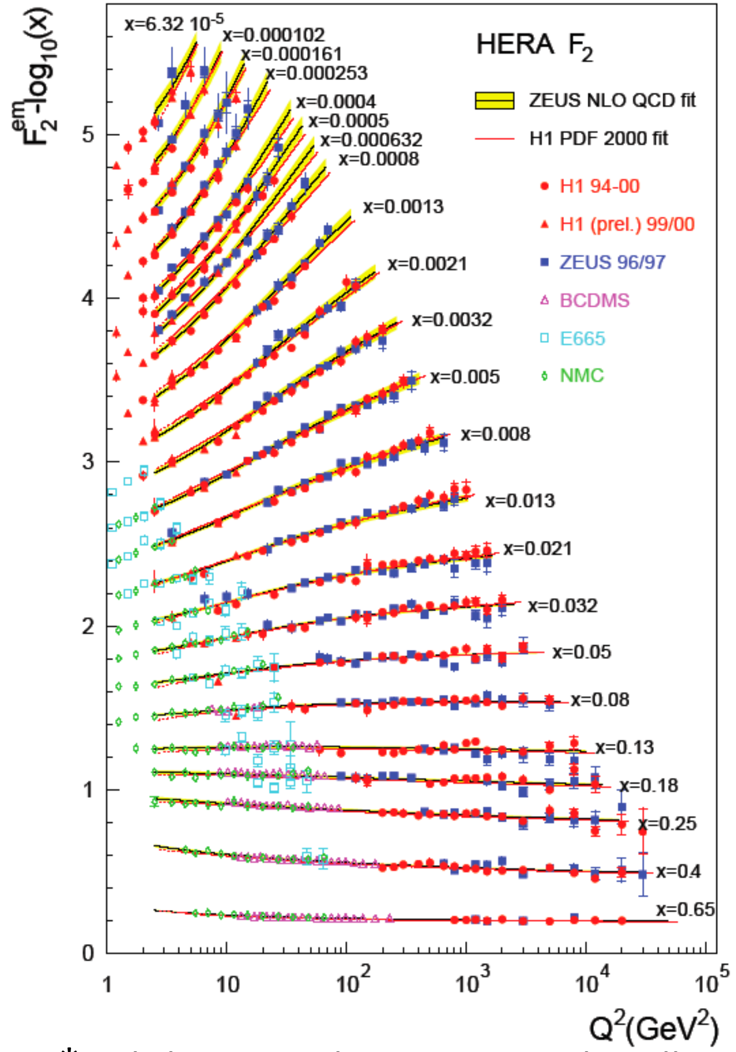
$$\frac{\partial F_2(x,Q^2)}{\partial \ln Q^2} \propto G(x,Q^2)$$

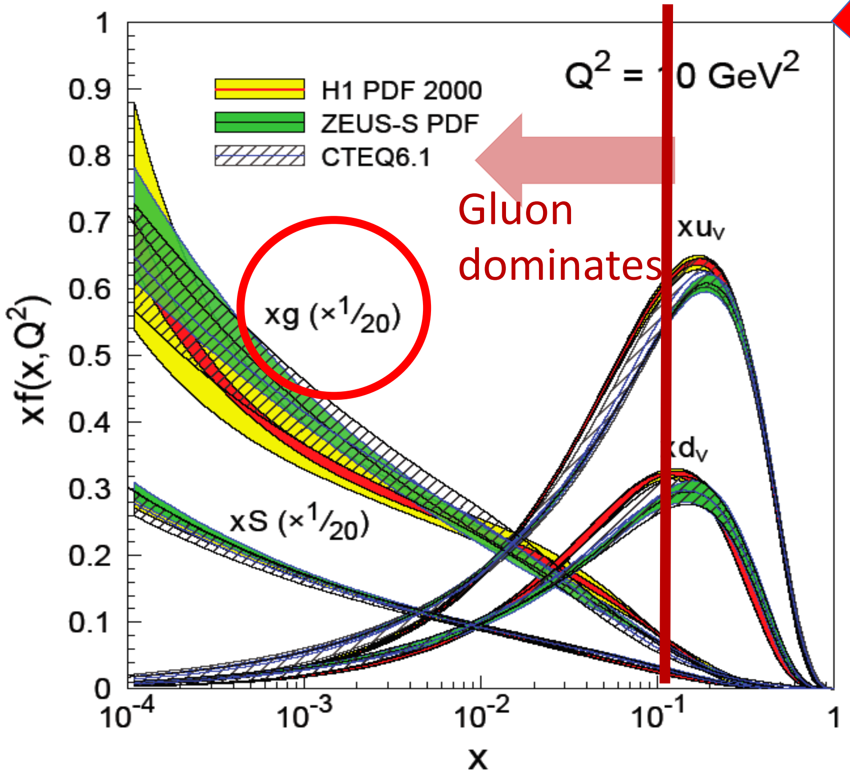
- NLO pQCD analyses: fits with **linear** DGLAP* equations

F_2 Structure Function

Vs.

Q^2 mom. exchanged



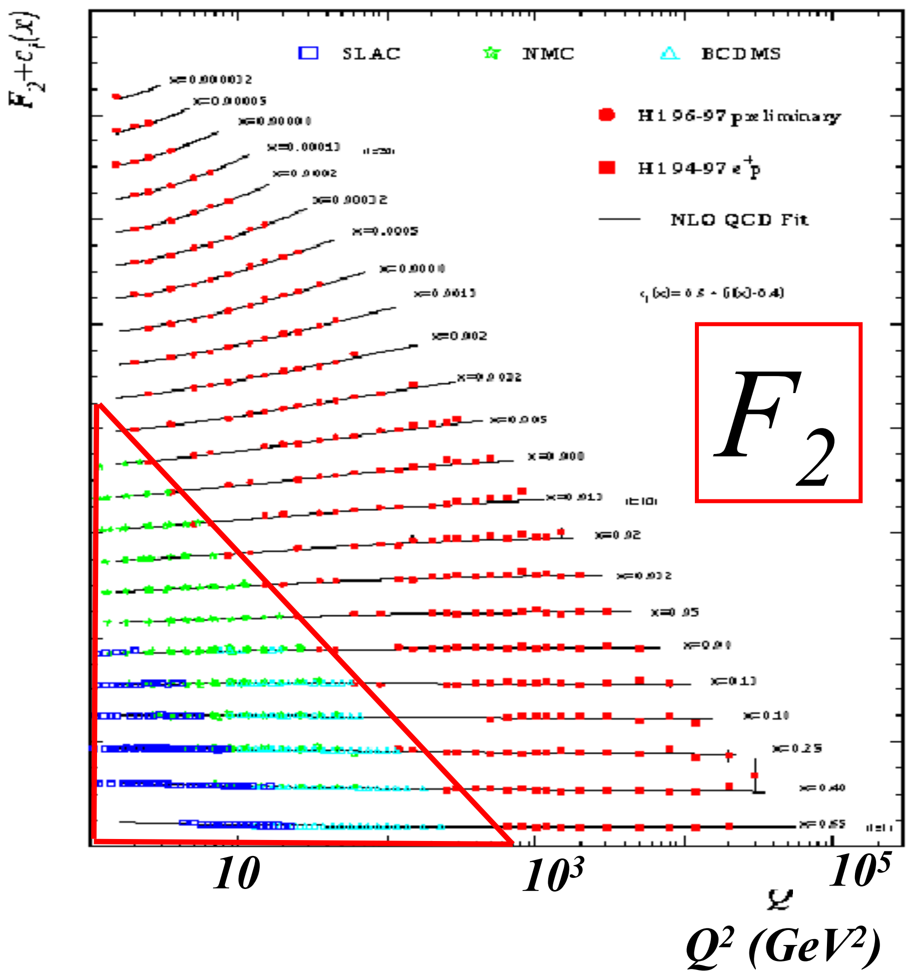


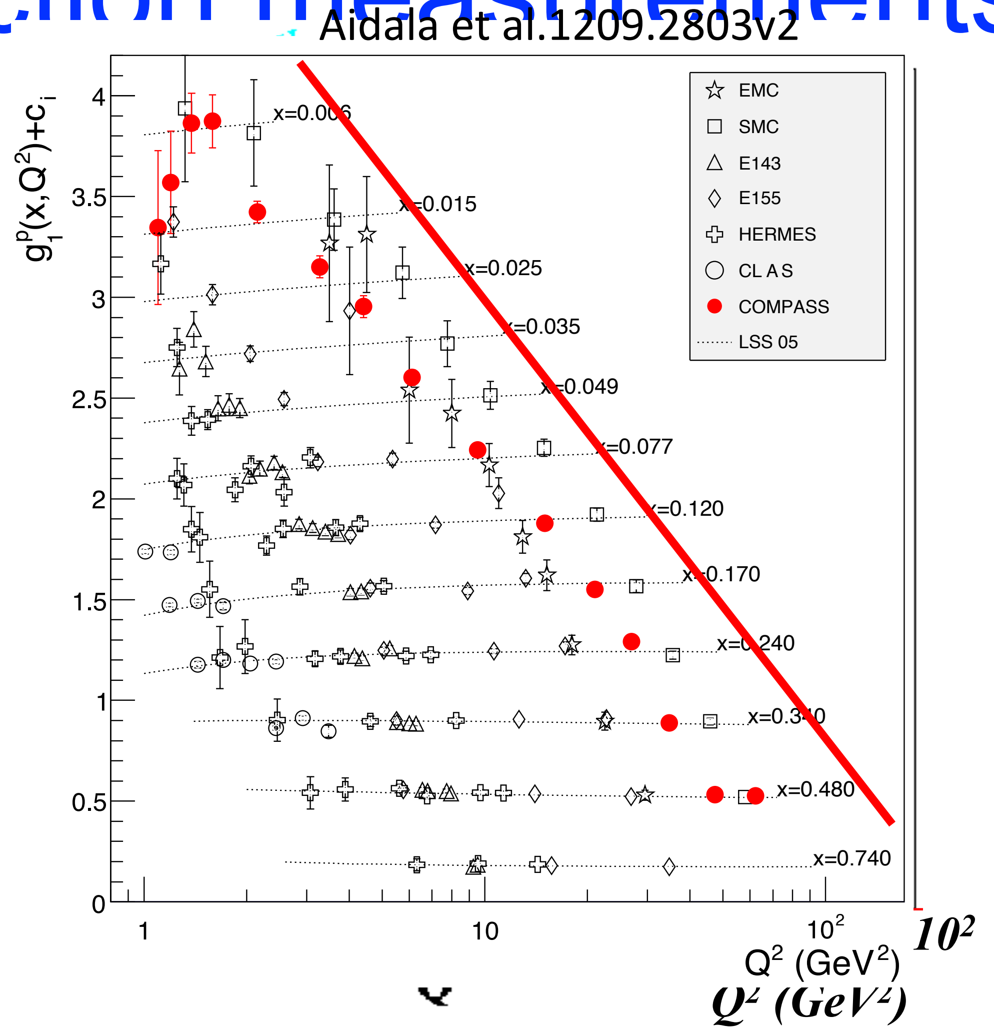
*Dokshitzer, Gribov, Lipatov, Altarelli, Parisi

Can one do the same thing for spin structure function g_1?

Spin contribution of the gluon to the proton from scaling violation g_1 spin structure function?

F_2 vs. g_1 structure function measurements

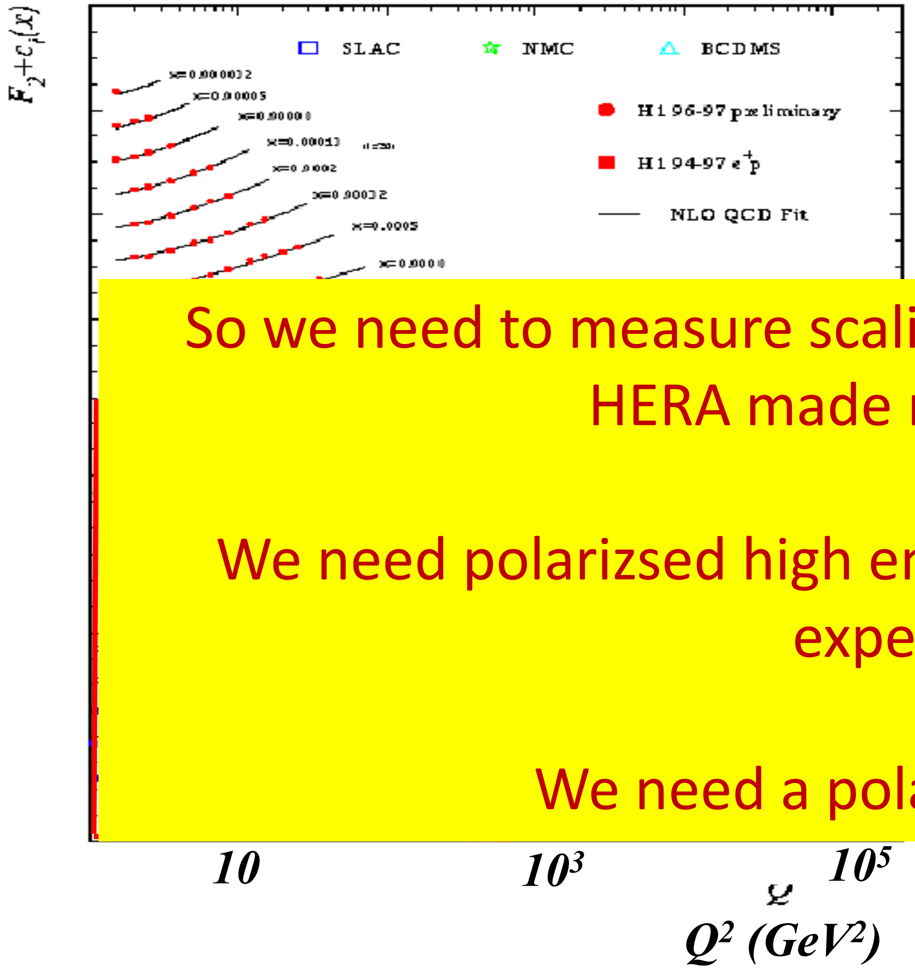


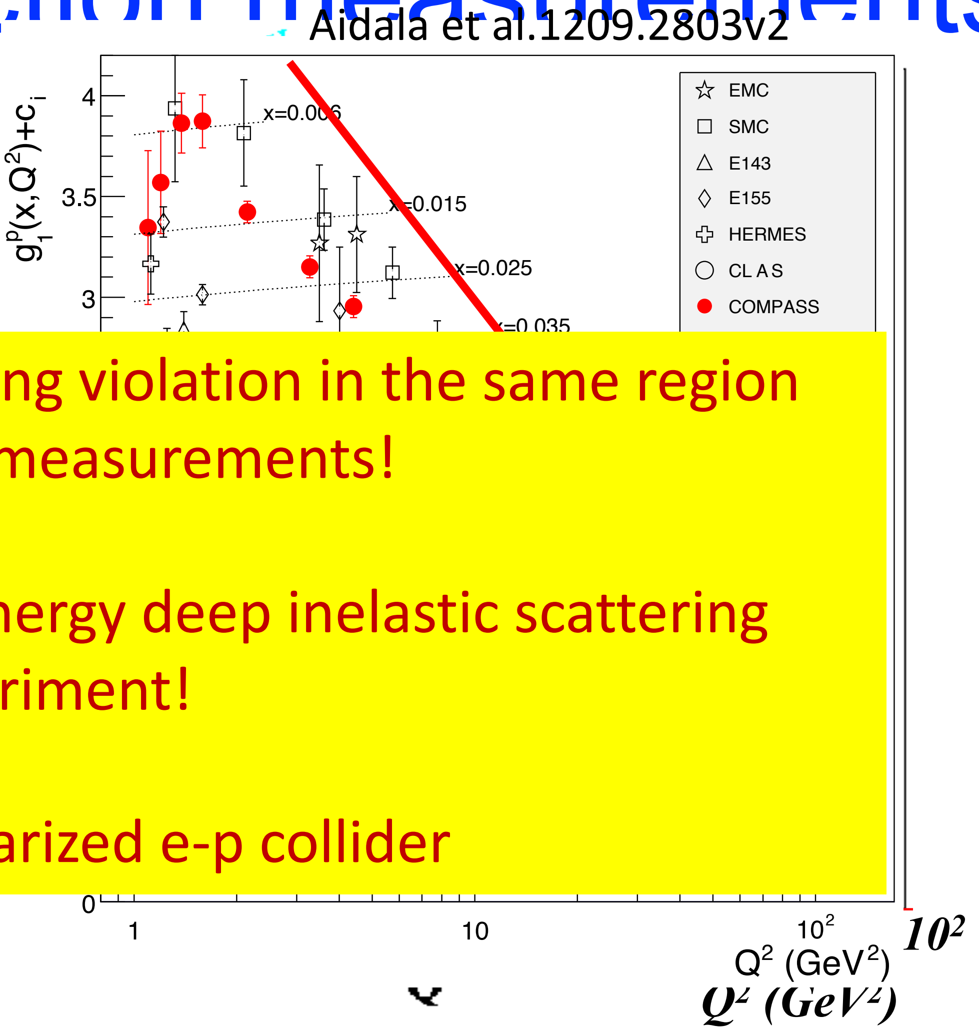


Large amount of polarized data since 1998… but not in NEW kinematic region!

Large uncertainty in gluon polarization (+/-1.5) results from lack of wide Q^2 arm

F_2 vs. g_1 structure function measurements





So we need to measure scaling violation in the same region HERA made measurements!

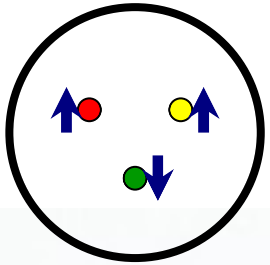
We need polarizsed high energy deep inelastic scattering experiment!

We need a polarized e-p collider

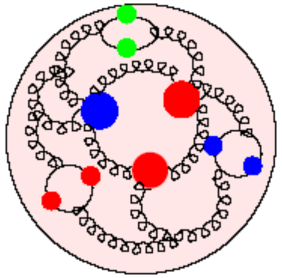
Large amount of polarized data since 1998… but not in NEW kinematic region!

Large uncertainty in gluon polarization (+/-1.5) results from lack of wide Q^2 arm

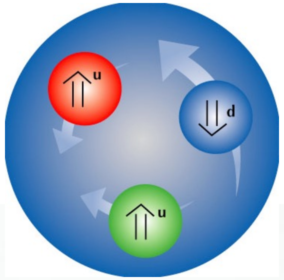
Our Understanding of Nucleon Spin Puzzle



1980s



+



1990/2000s

$$\frac{1}{2} = \left[\frac{1}{2}\Delta\Sigma + L_Q\right] + [\Delta g + L_G]$$

$\Delta\Sigma /2$ = 0.12 +/- 0.17

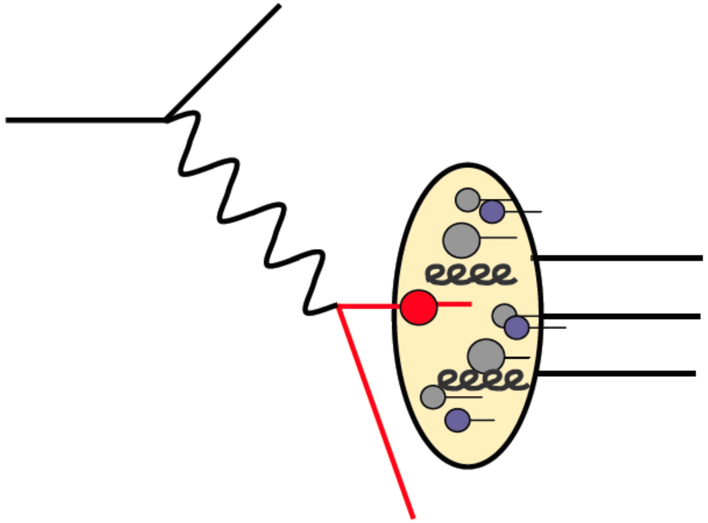
Need information about transverse dimensions of the proton

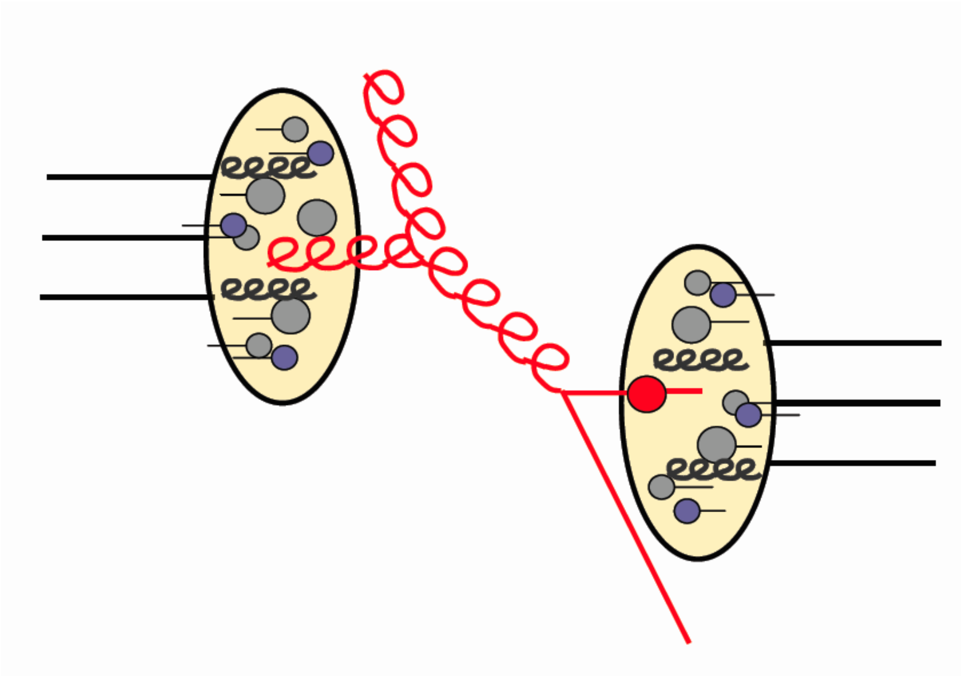
Spin discovered a problem…. What now? Need precision and investigations of gluons….

RHIC Spin program and the Transverse Spin puzzle

Pre-cursor to a polarized e-p collider

Complementary techniques





Photons colorless: forced to interact at NLO with gluons

Can't distinguish between quarks and anti-quarks either

Why not use polarized quarks and gluons abundantly available in protons as probes ?

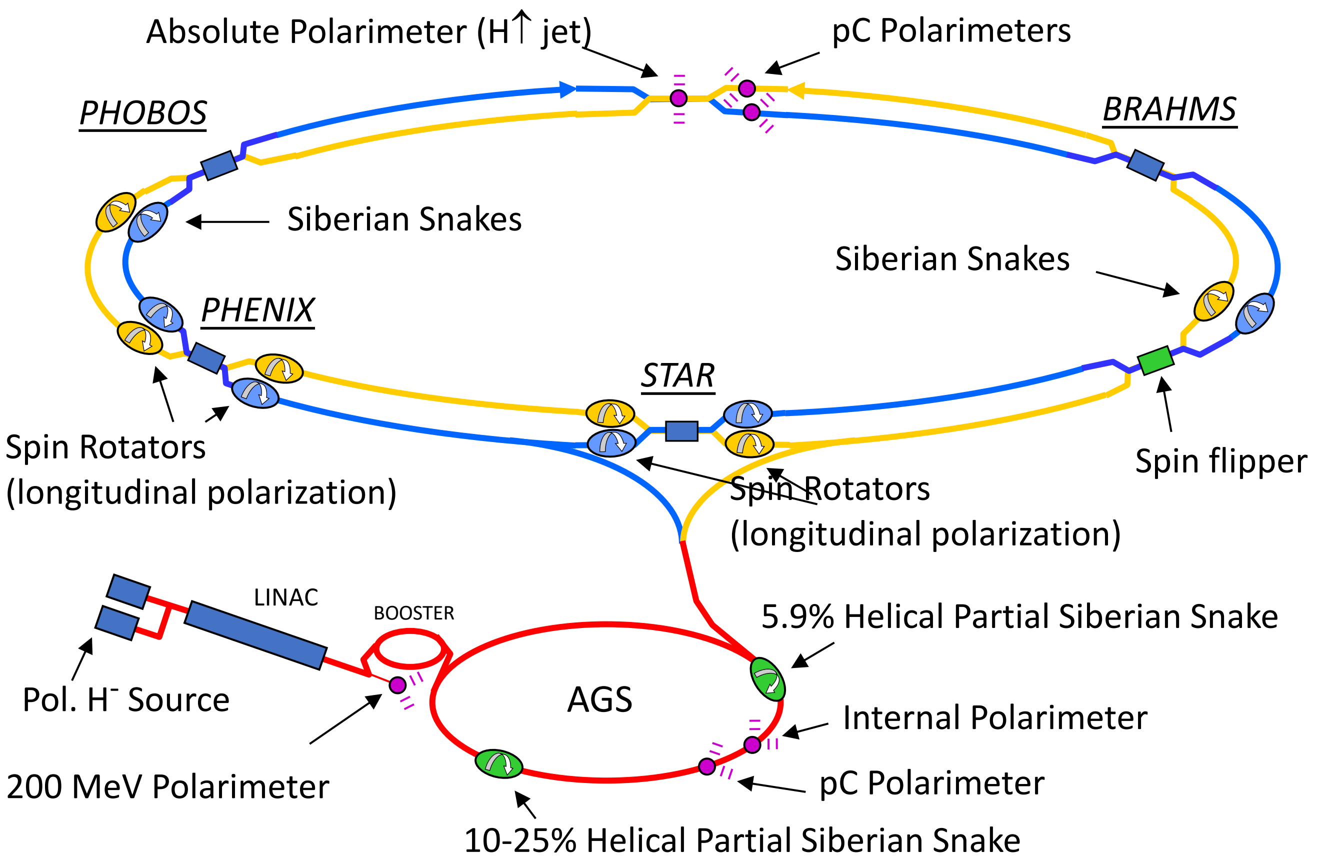
Seeds for RHIC Spin program:

If one wants to study gluon's spin contribution to proton's spin, ***why not directly explore the gluon spin with polarized proton p-p collisions?***

Curious and bothersome transverse spin asymmetries in p-p scattering persistent in every experiment performed…. US physicists heavily involved… decided to investigate further at high energy

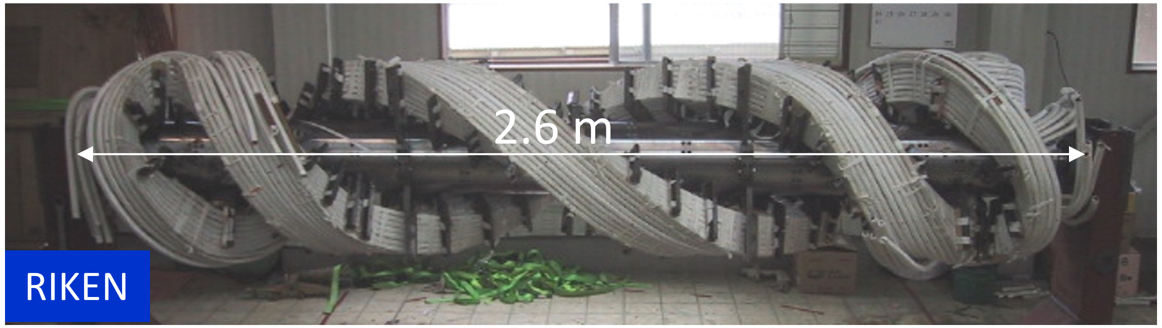
Technical know-how of polarizing proton beams at high energy became available in the mid-late 1990's

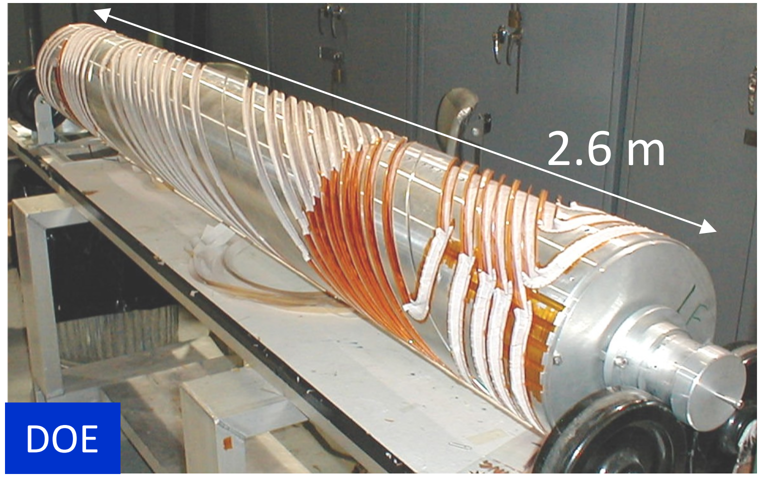
RHIC as a Polarized Proton Collider



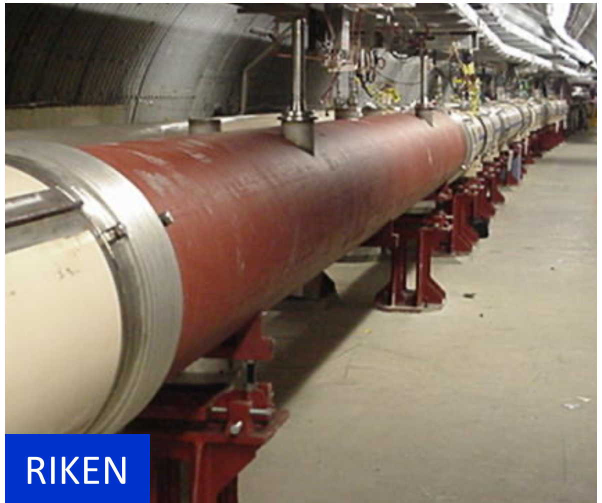
Without Siberian snakes: $\nu_{sp} = G\gamma = 1.79\ E/m \rightarrow$ ~1000 depolarizing resonances
With Siberian snakes (local 180° spin rotators): $\nu_{sp} = ½ \rightarrow$ no first order resonances
Two partial Siberian snakes (11° and 27° spin rotators) in AGS

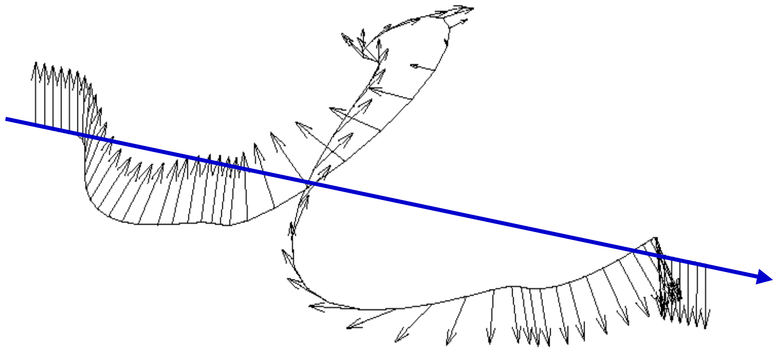
Siberian Snakes



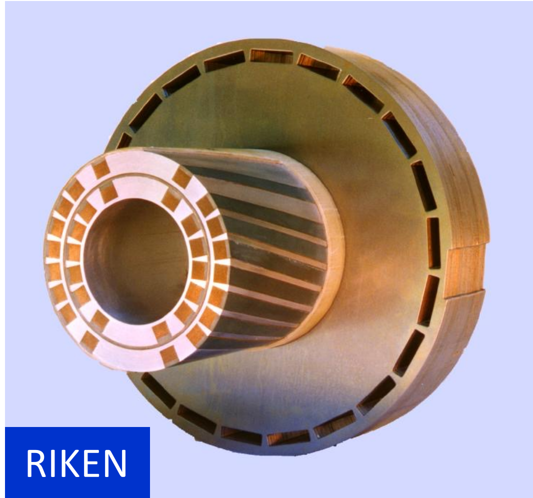


- AGS Siberian Snakes: variable twist helical dipoles, 1.5 T (RT) and 3 T (SC), 2.6 m long
- RHIC Siberian Snakes: 4 SC helical dipoles, 4 T, each 2.4 m long and full 360° twist





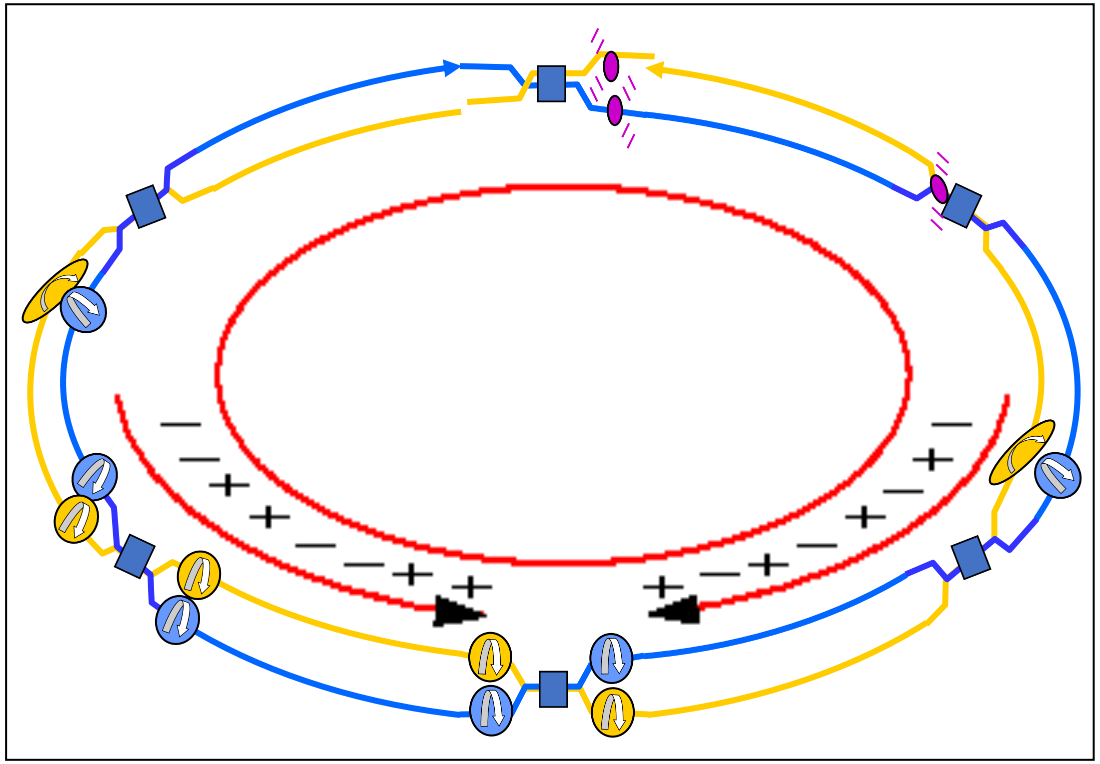
Courtesy of A. Luccio



Measuring A_{LL}

Longitudinal Spin Asymmetry using polarized proton bunches in the RHIC ring

$$A_{LL} = \frac{d\sigma_{++} - d\sigma_{+-}}{d\sigma_{++} + d\sigma_{+-}} = \frac{1}{|P_1 P_2|}\frac{N_{++} - RN_{+-}}{N_{++} - RN_{+-}}; \qquad R = \frac{L_{++}}{L_{+-}}$$

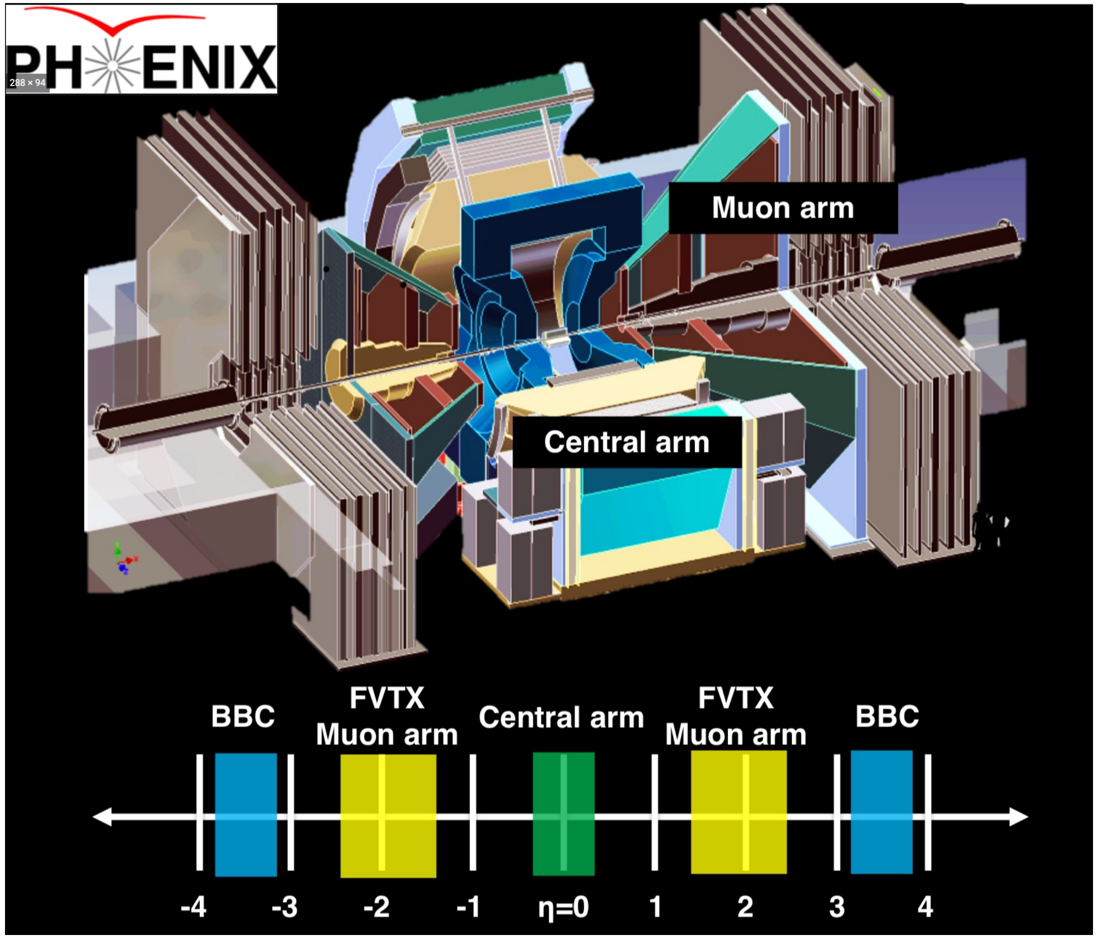


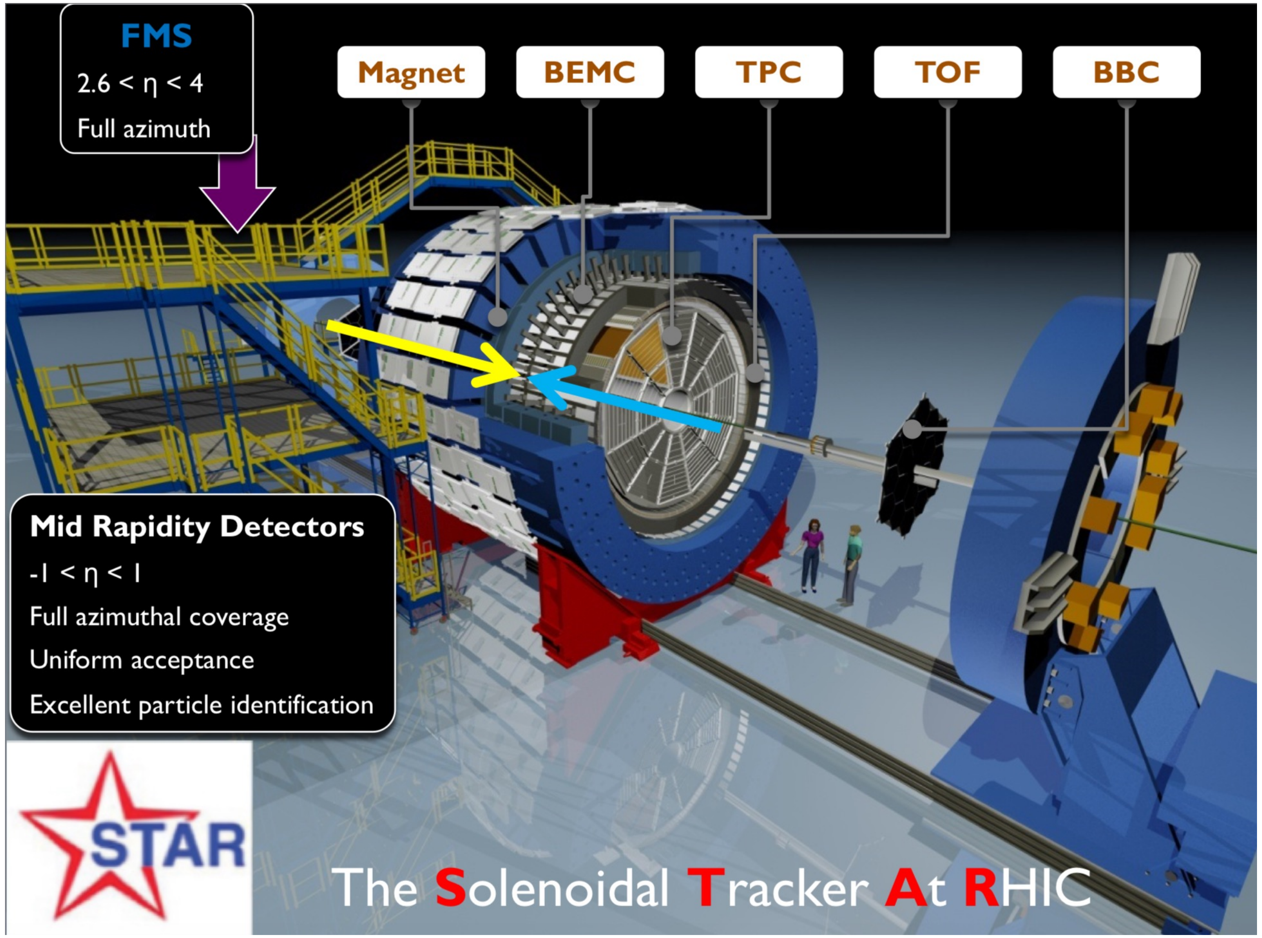
(N) Yield
(R) Relative Luminosity
(P) Polarization

Exquisite control over false asymmetries due to ultra fast rotations of the target and probe spin.

✓ Bunch spin configuration alternates every 106 ns
✓ Data for all bunch spin configurations are collected at the same time
⇒ Possibility for false asymmetries are greatly reduced

Two main detectors for spin studies



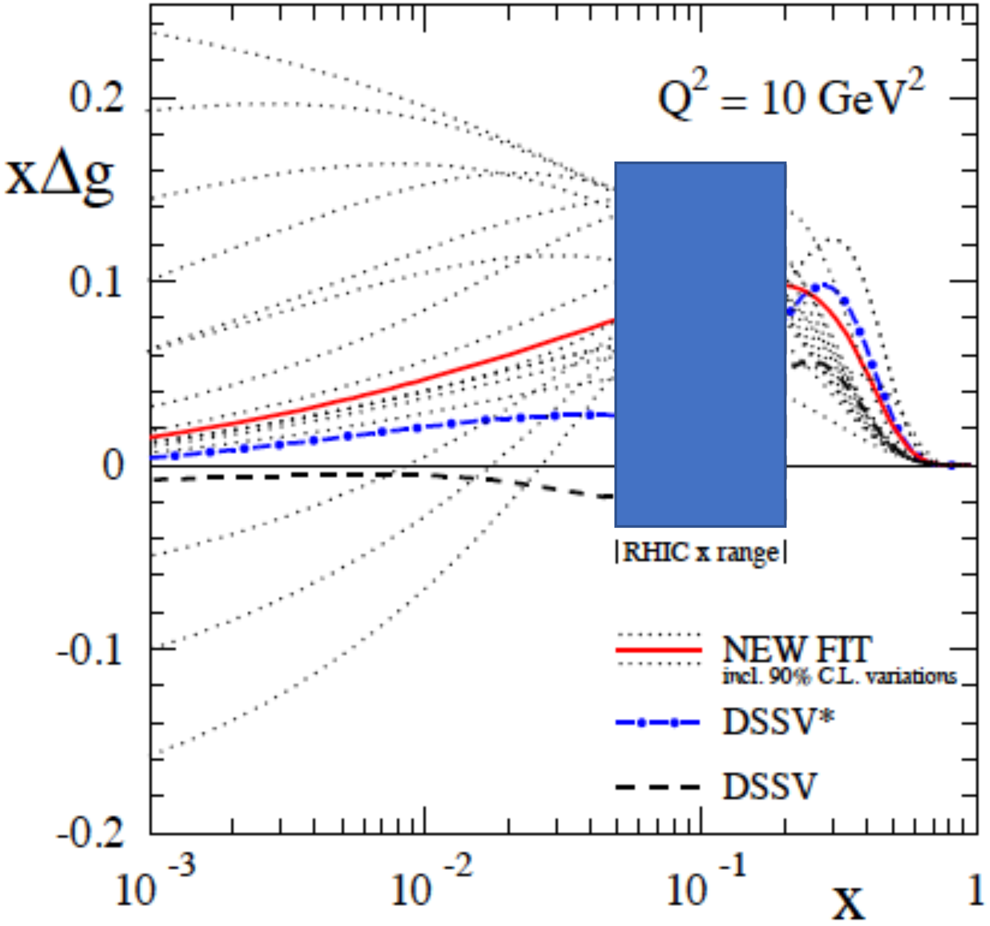


Recent global analysis: DSSV

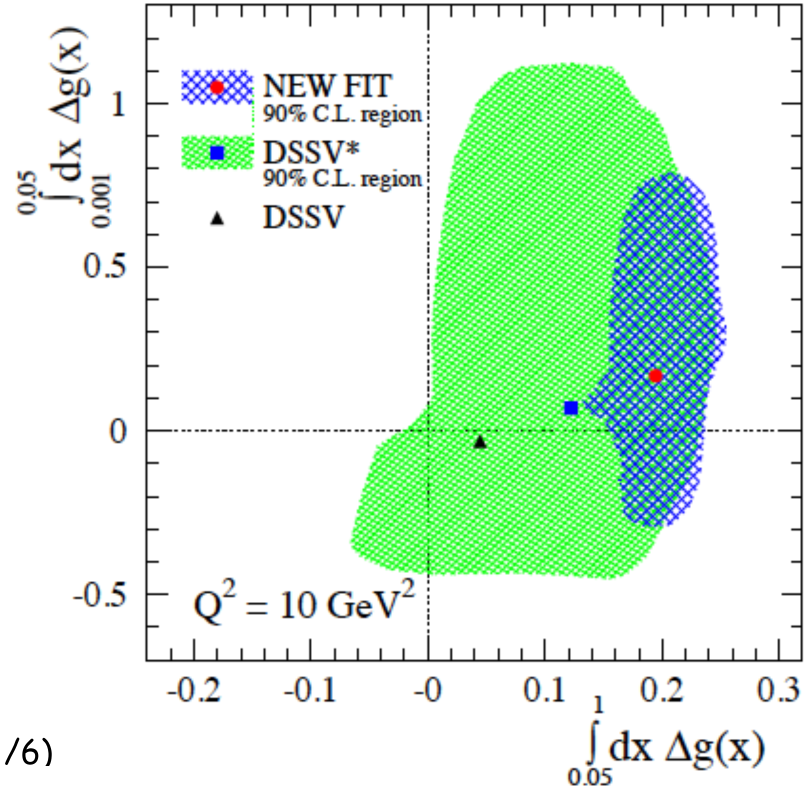
D. deFlorian et al., arXiv:1404.4293

$\Delta G = 0.2$ +/- 0.02 +/- 0.5

Wide spread at low x (x<0.05) of alternative fits consistent within 90% of C.L.

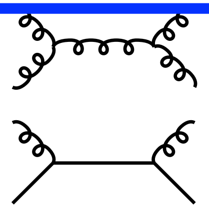


/6)



While RHIC made a huge impact on ΔG

large uncertainties to remain in the low-x unmeasured region!

Reaction	Dom. partonic process	probes	LO Feynman diagram
$\vec{p}\vec{p} \to \pi + X$ [61, 62]	$\vec{g}\vec{g} \to gg$ $\vec{q}\vec{g} \to qg$	Δg	
$\vec{p}\vec{p} \to \text{jet(s)} + X$ [71, 72]	$\vec{g}\vec{g} \to gg$ $\vec{q}\vec{g} \to qg$	Δg	(as above)

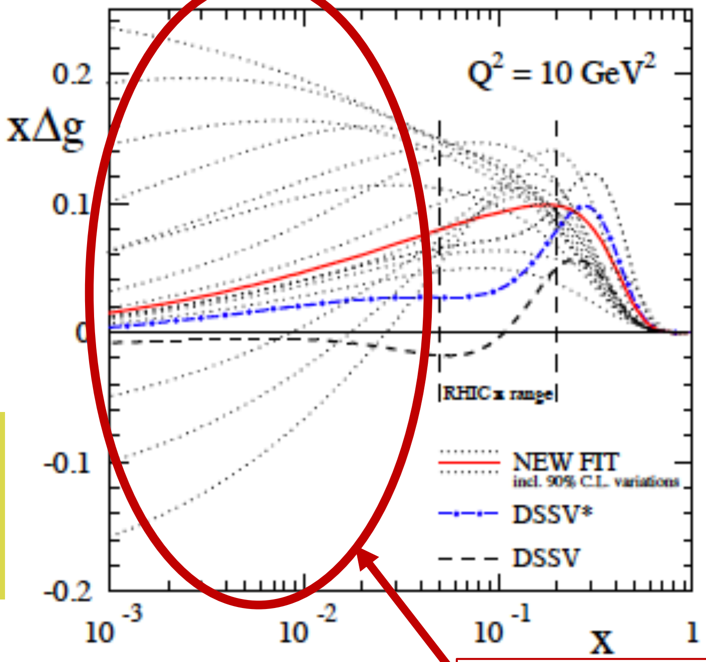
2009 RHIC data established non-zero ΔG

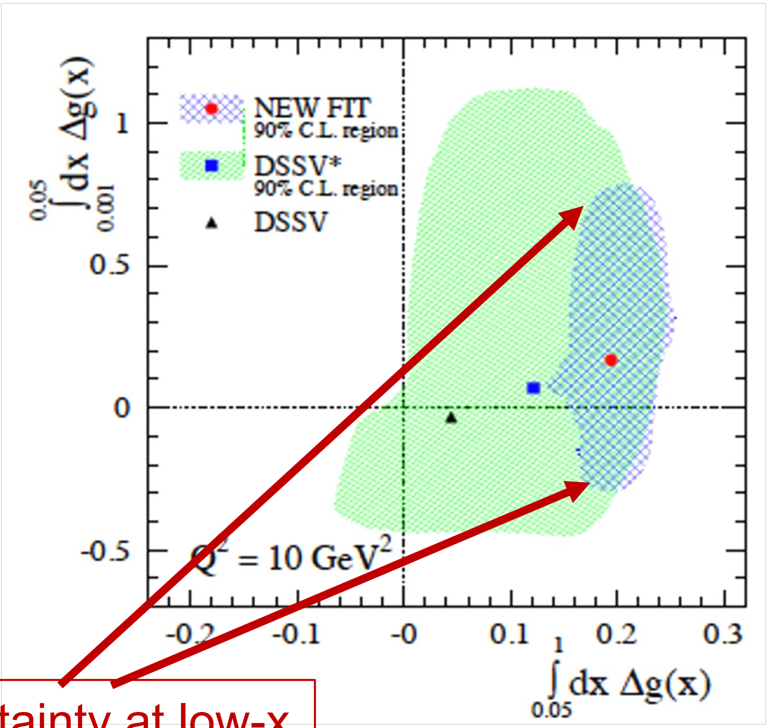
-- PHENIX 2005-9, PRD 90, 12007 (2014)

-- STAR 2009, PRL 115 (2015) 92002

-- DSSV PRL (113) 12001 (2014)

$$\int_{0.05}^{1.0} dx\Delta g \sim 0.2 \pm^{0.06}_{0.07} @ \ 10\ \text{GeV}^2$$



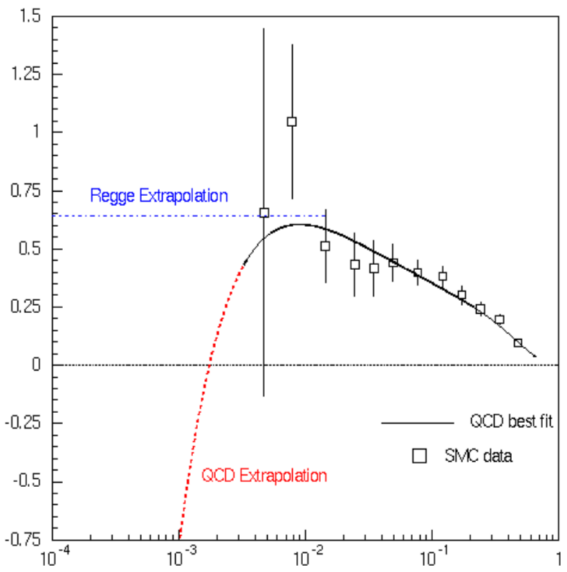


Large uncertainty at low-x

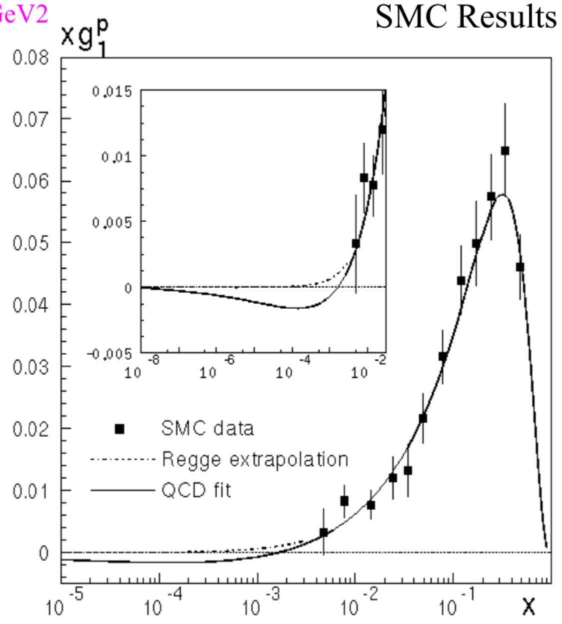
Recall from “EIC A Dream to Reality” on Day 1

Lack of low x data… consequences

Q2 = 10 GeV2



SMC Results



- Regge extrapolation:

$$\int_0^{0.003} g_1^p(x, Q_0^2)dx = 0.002 \pm 0.002$$

- QCD fit extrapolation:

$$\int_0^{0.003} g_1^p(x, Q_0^2)dx = -0.011 \pm 0.011$$

$$g_1(x \to 0) \propto x^{\alpha} \quad \text{as } 0 < \alpha < 0.5$$

Regge/QCD

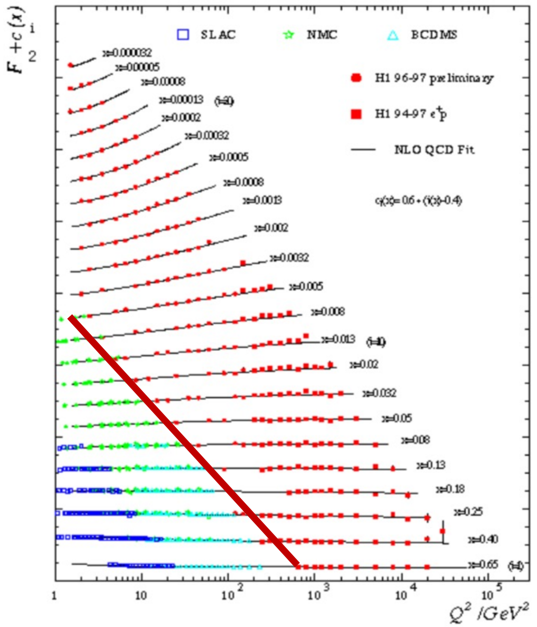
In these discussions, while many focused on the low-x Extrapolations.

SMC PRD98 (112002) 1998

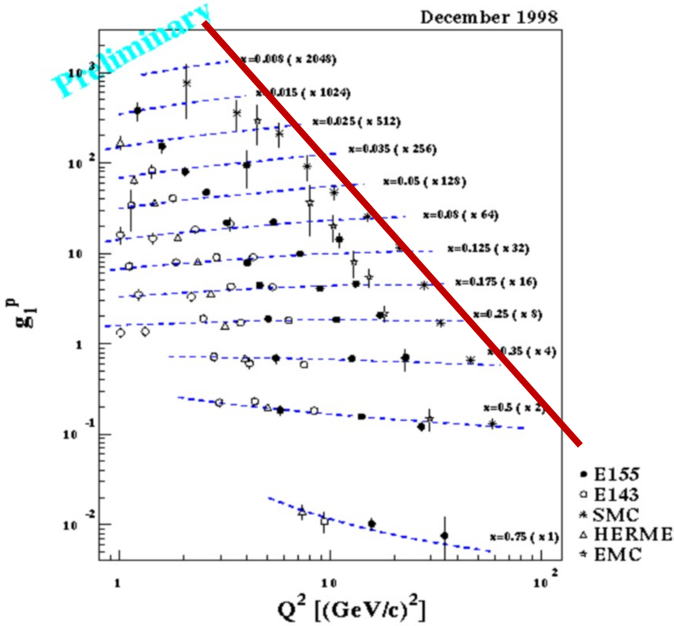
Seeds for a **polarized collider**

How far does polarized DIS have to go!

World data on F_1^p



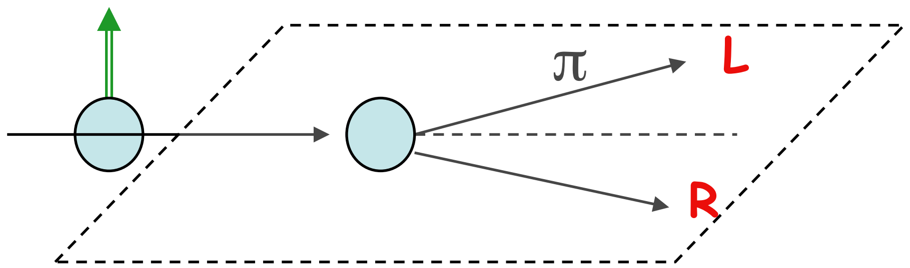
World data on g_1^p



Transverse Spin effects in p-p observed but ignored for 40+ years

Recent developments and state of the art in Alessandro Bacchetta's and Silvia Dalla Torre's lectures

Transverse spin introduction



$$A_N = \frac{N_L - N_R}{N_L + N_R}$$

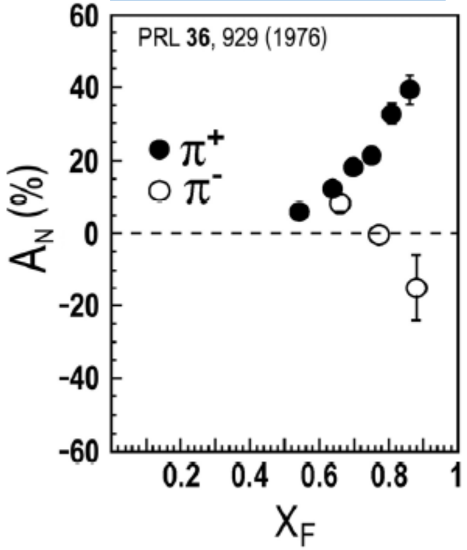
$$A_N \sim \frac{m_q}{p_T} \cdot \alpha_S \sim 0.001$$

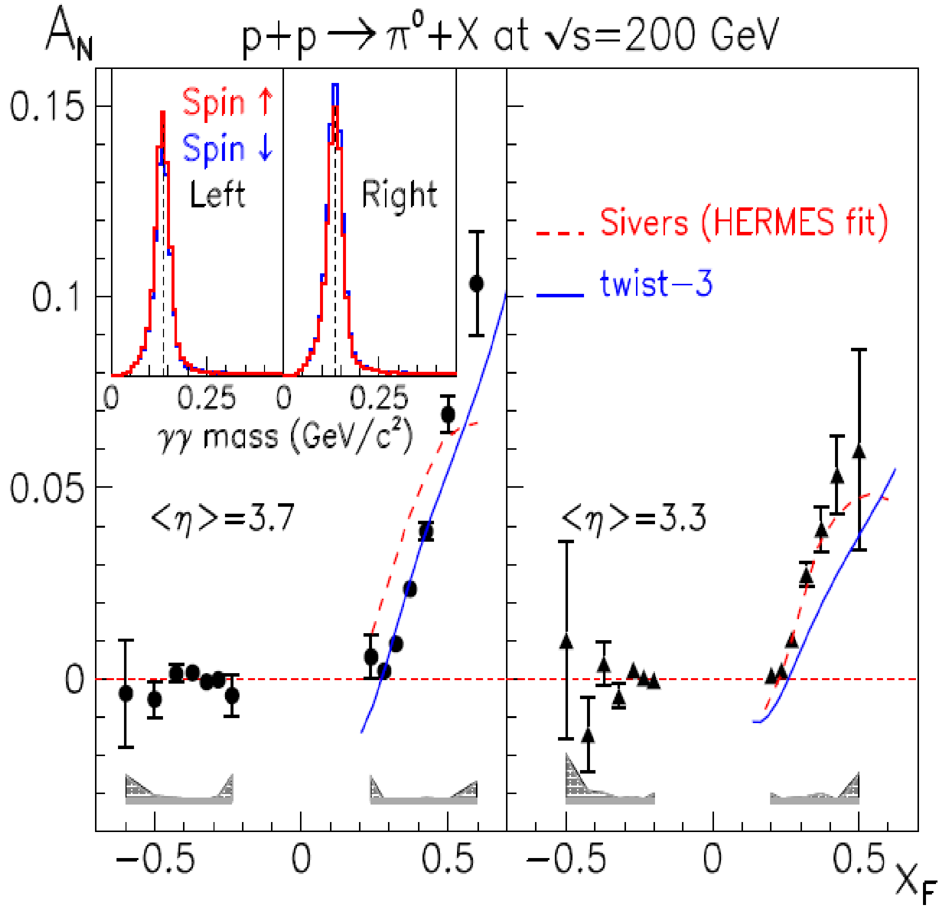
Kane, Pumplin and Repko
PRL 41 1689 (1978)

- Since people focused at high p_T to interpret them in pQCD frameworks, this (expected small effect) was “neglected However….
- Pion production in single transverse spin collisions showed us something different….

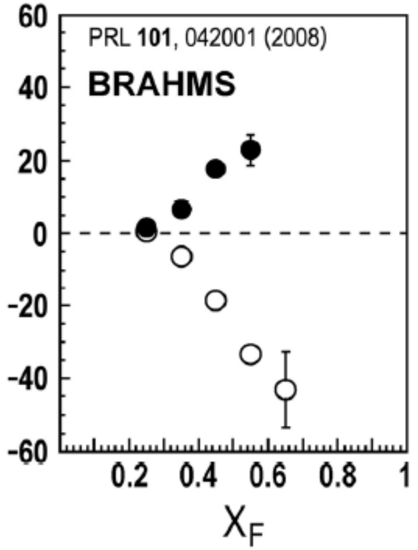
Pion asymmetries: at broad range in CM energies!

ZGS/ANL
√s=4.9 GeV





RHIC
√s=62.4 GeV

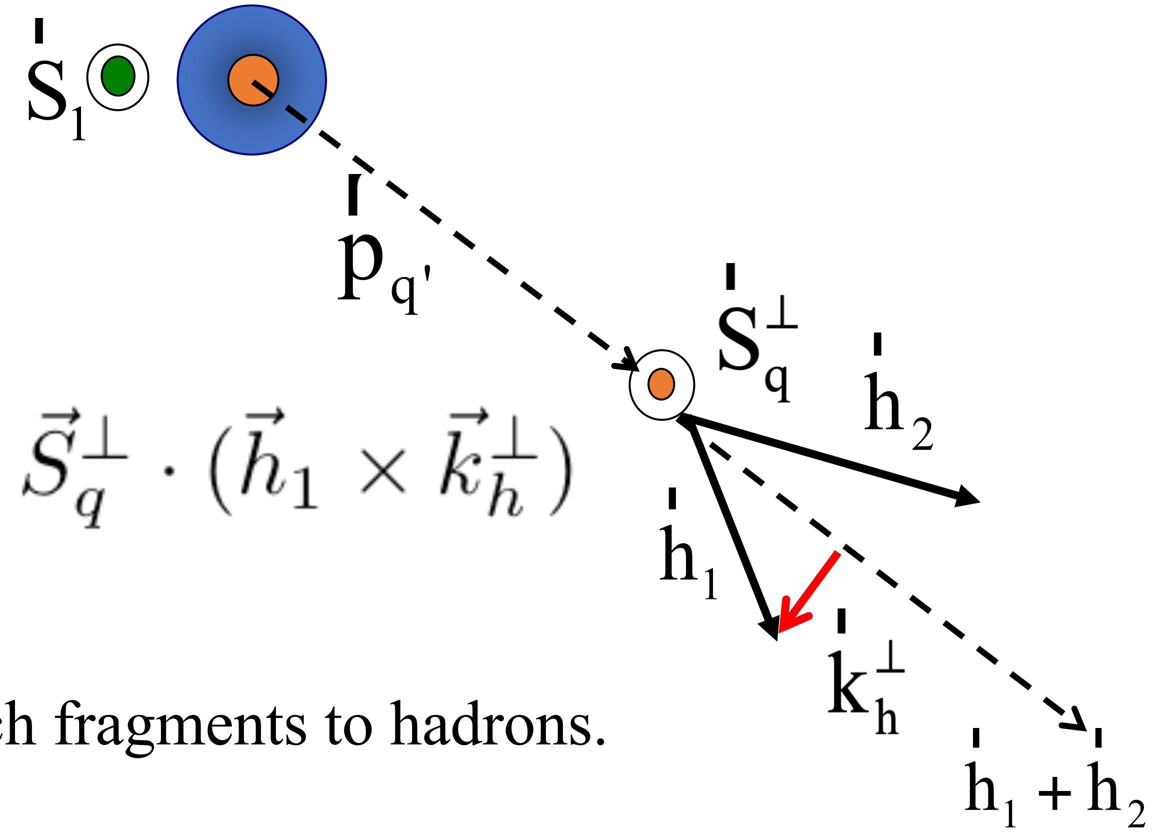


Suspect soft QCD effects at low scales, but they seem to remain relevant to perturbative regimes as well ➔ **0.001 expected 0.2-0.6 observed at all Center of Mass Energies**

What could be the origin of such effect?
Collins (Heppelmann) effect: Asymmetry in the fragmentation hadrons

Example: $p^{\uparrow} + p \rightarrow h_1 + h_2 + X$

Nucl Phys B396 (1993) 161,
Nucl Phys B420 (1994) 565

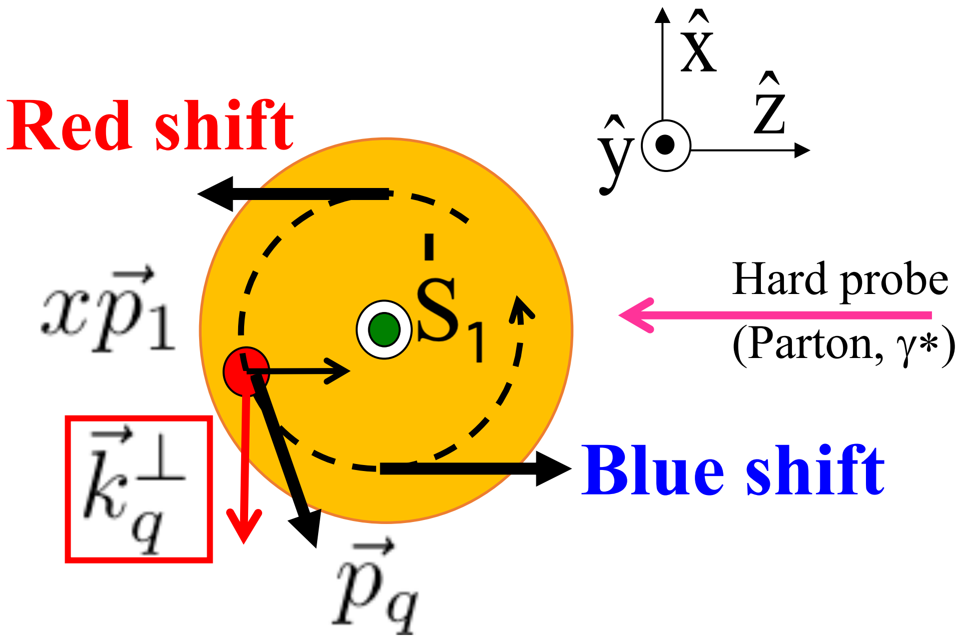


$$SSA_{Collins} \propto \vec{S}_q^{\perp} \cdot (\vec{h}_1 \times \vec{k}_h^{\perp})$$

Polarization of struck quark which fragments to hadrons.

Other possibility: What does “Sivers effect” probe?

Top view, Breit frame



Quarks orbital motion adds/ subtracts longitudinal momentum for negative/positive $\hat{x}$.

PRD66 (2002) 114005

Parton **D**istribution **F**unctions rapidly fall in longitudinal momentum fraction x.

Final **S**tate **I**nteraction between outgoing quark and target spectator.

Sivers function

$$f_{1T}^{\perp}(x, \vec{k}_q^{\perp})$$

hep-ph/ 0703176

Quark Orbital angular momentum

Generalized **P**arton **D**istribution Functions

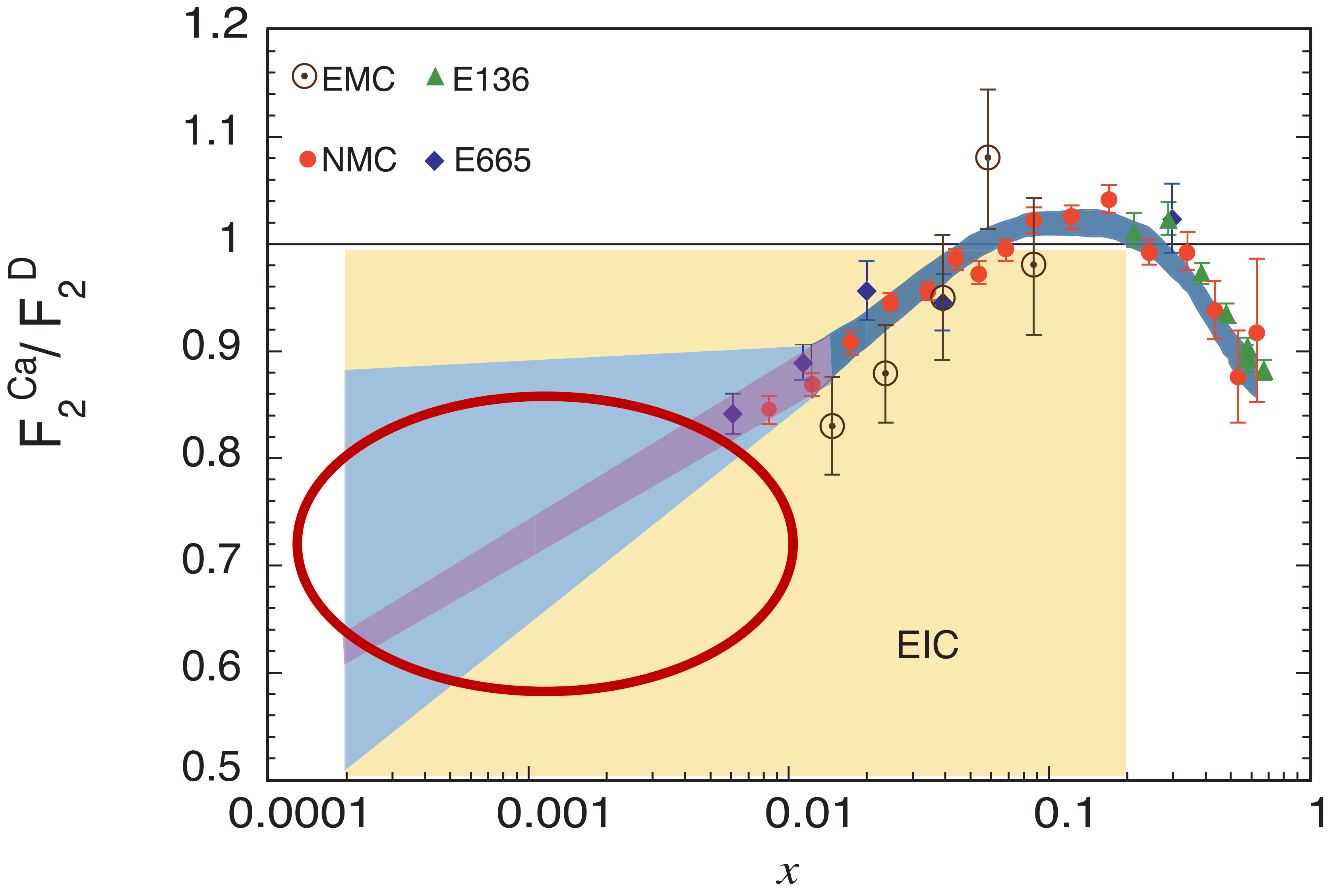
PRD59 (1999) 014013

Lepton nucleus scattering for understanding the nuclear structure and dynamics:

Nuclear structure a known unknown….

Much more in Pia Zurita's lecture

PDFs in nuclei are different than in protons!



Since 1980's we know the ratio of F_2's of nuclei to that of Deuteron (or proton) are different.

Nuclear medium modifies the PDF's.

Fair understanding of what goes on, in the x > 0.01.

However, what happens at low x?

Does this ratio saturate? Or keep on going? – Physics would be very different depending on what is observed.

Data needed at low-x

Lessons learned:

- Proton and neutrons spin not just alignment of quarks and gluons….
  - Proton's spin is complex: alignment of quarks, **gluons and orbital motion**
- To fully understand proton structure (including the partonic dynamics) one needs to explore over a **broader x-Q2 range (not in fixed target but in collider experiment)**
- **Low-x behavior of gluons in proton:** Precise measurements of gluons critical.
- Low x behavior of partons in Nuclei essential to complete our understanding of structure of matter…
- To understand the nuclear fragments – target fragment – one needs to measure e-A in a collider geometry

We need a new high-luminosity polarized **e-p/A** collider….

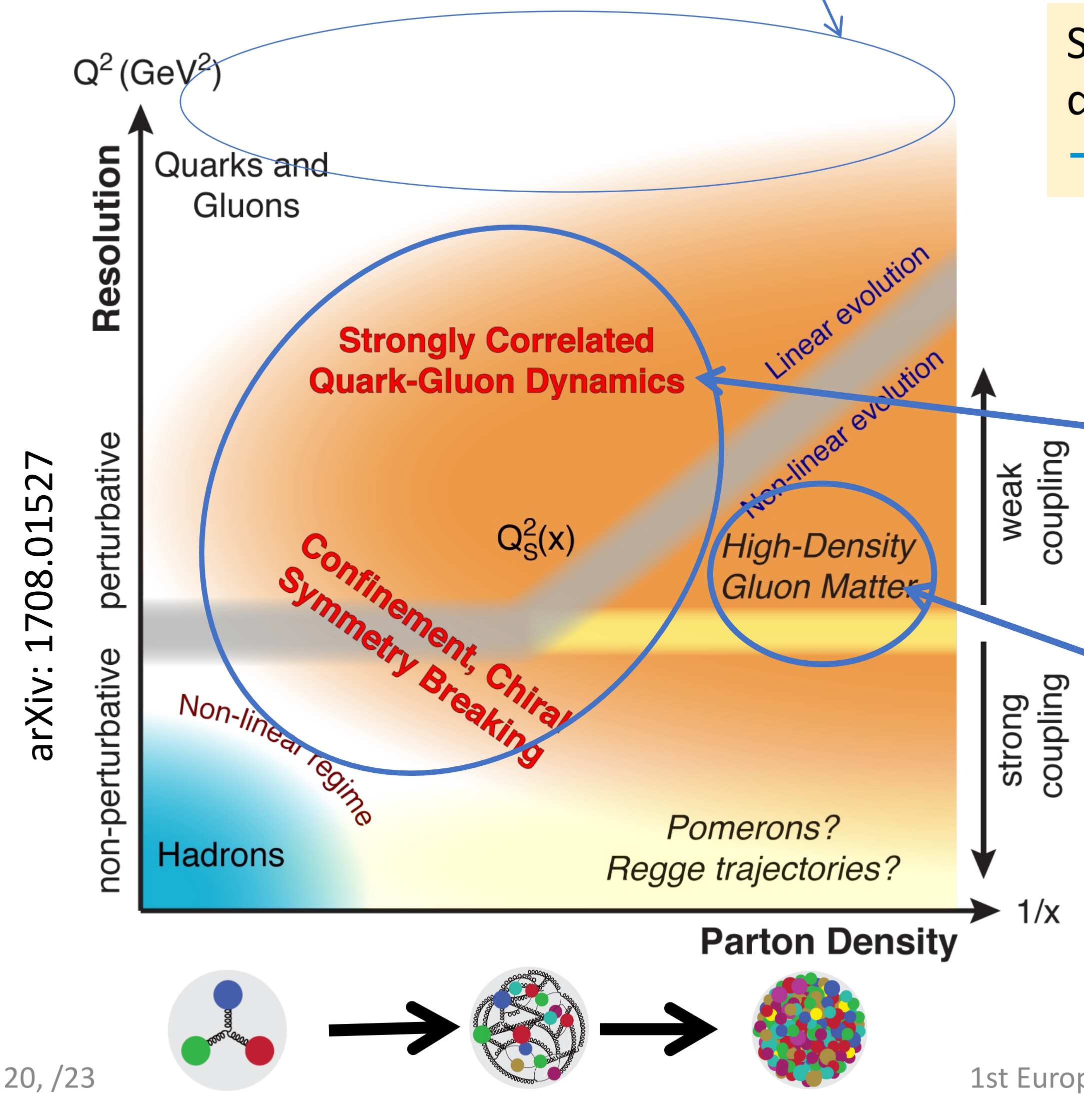
Center for Frontiers in Nuclear Science

Lecture Part 2: Why EIC now?

QCD Landscape to be explored by a future facility

QCD at high resolution (Q^2) — weakly correlated quarks and gluons are well-described

Strong QCD dynamics creates many-body correlations between quarks and gluons
→ hadron structure emerges

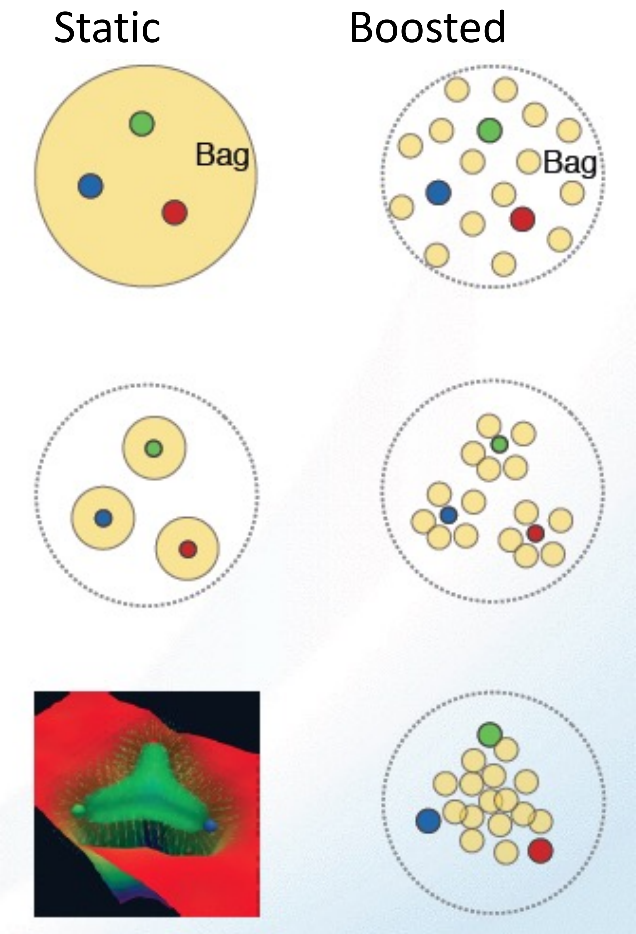


Systematically explore correlations in this region.

An exciting opportunity: Observation of a new regime in QCD of weakly coupled high-density matter

Need Precision and Control

What does a proton look like in transverse dimension?



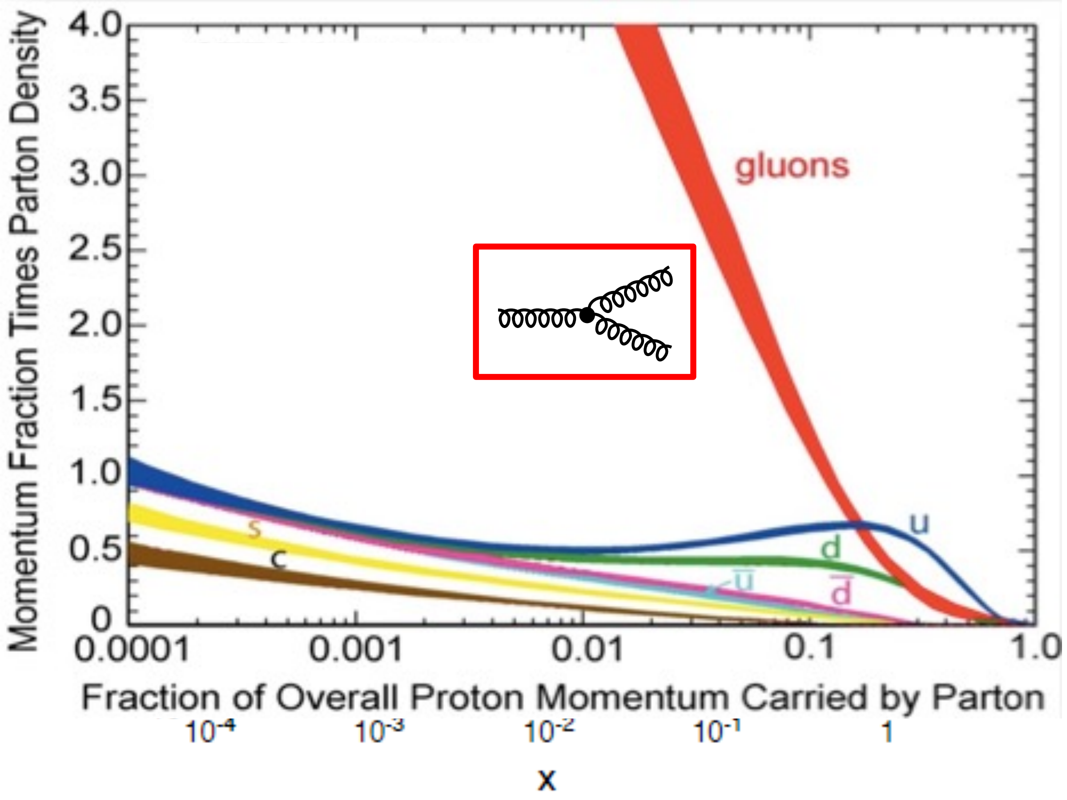
Bag Model: Gluon field distribution is wider than the fast moving quarks. Color (Gluon) radius > Charge (quark) Radius

Constituent Quark Model: Gluons and sea quarks hide inside massive quarks. Color (Gluon) radius ~ Charge (quark) Radius

Lattice Gauge theory (with slow moving quarks), gluons more concentrated inside the quarks: Color (Gluon) radius < Charge (quark) Radius

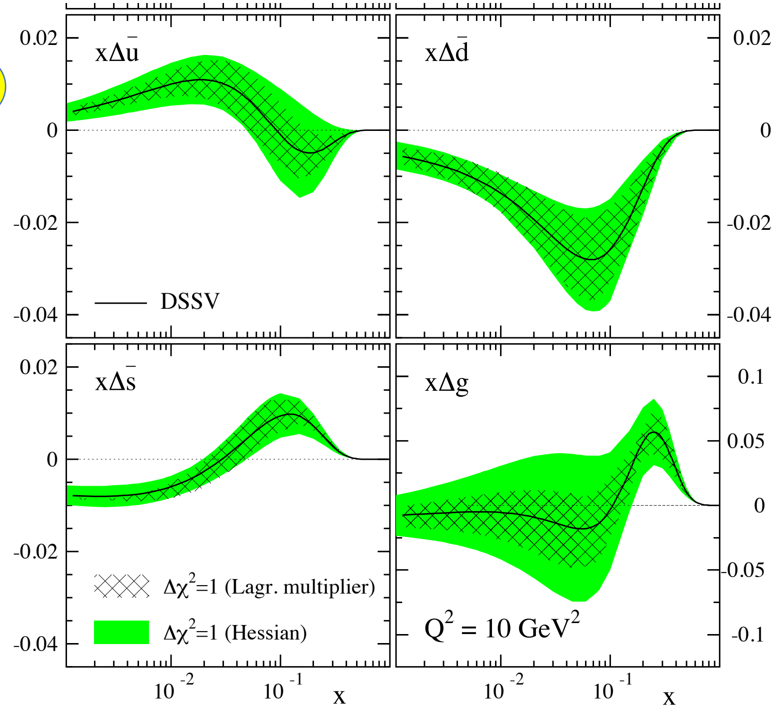
Need <u>*transverse*</u> images of the quarks <u>*and gluons*</u> in protons

What do ***gluons*** in protons look like?
Unpolarized & polarized parton distribution functions



QCD
Terra-incognita!

High Potential for Discovery



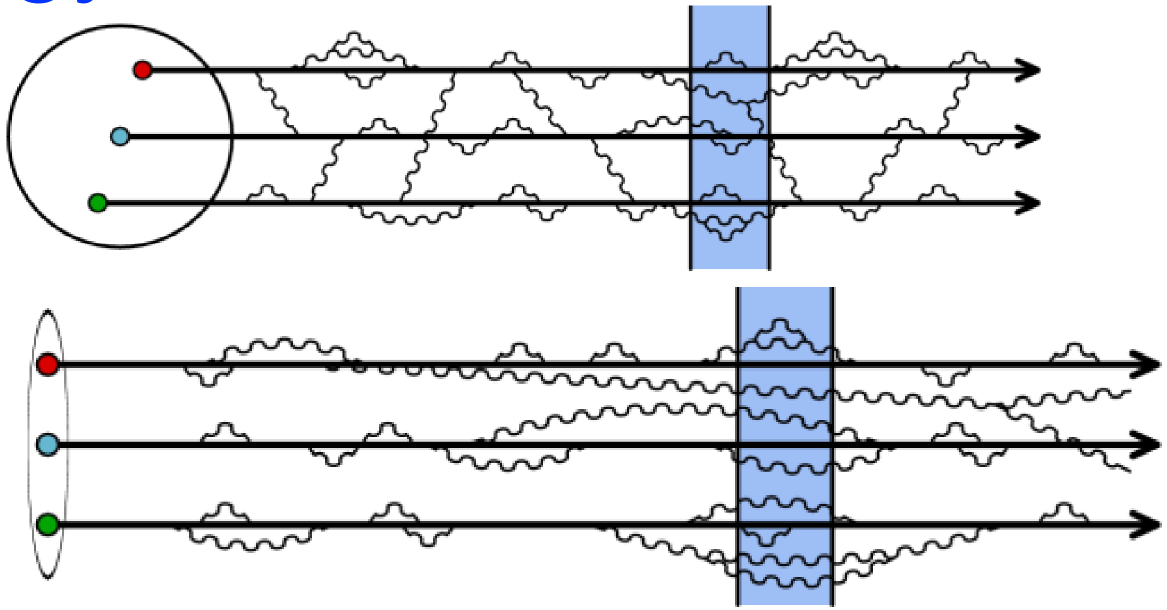
Need to go beyond 1-dimension!

Need (2+1)D image of gluons in a nucleon in position & momentum space

How does a Proton look at low and very high energy?

Low energy: High x
Regime of fixed target exp.

High energy: Low- x
Regime of a Collider



Cartoon of boosted proton

At high energy:

- Wee partons fluctuations are time dilated in strong interaction time scales
- Long lived gluons radiate further smaller x gluons ➔ which intern radiate more……. Leading to a **runaway growth?**

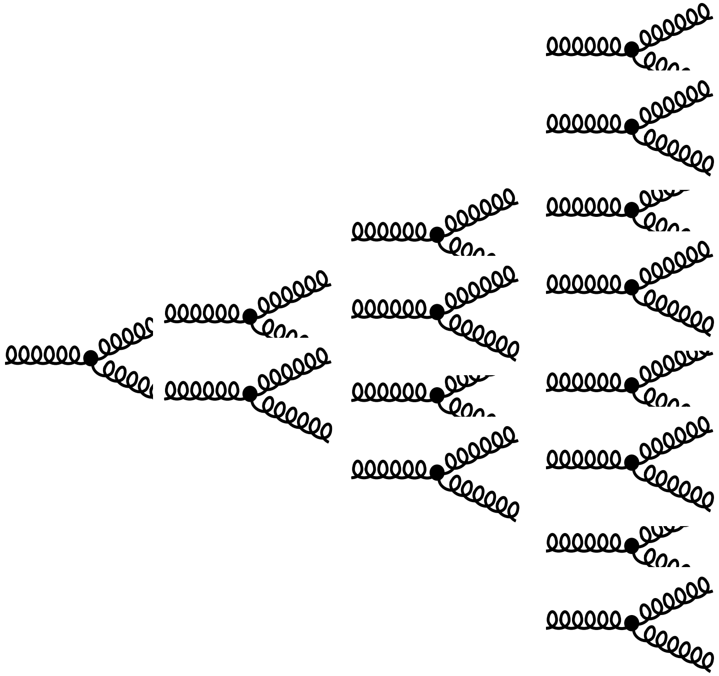
Recall Marco Radici's comment

Gluon and the consequences of its interesting properties:

Gluons carry color charge ➔ Can interact with other gluons!

"…The result is a self catalyzing enhancement that leads to a runaway growth. A small color charge in isolation builds up a big color thundercloud…."

F. Wilczek, in "Origin of Mass"
Nobel Prize, 2004



? Infinity?

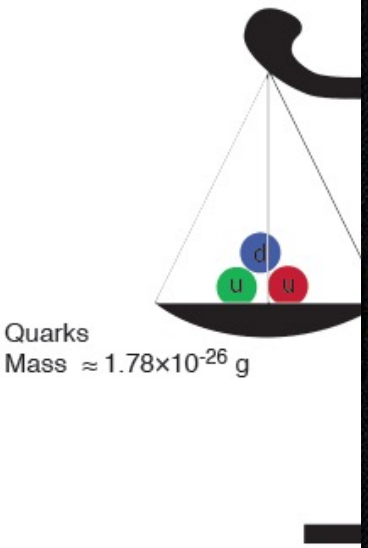
No!

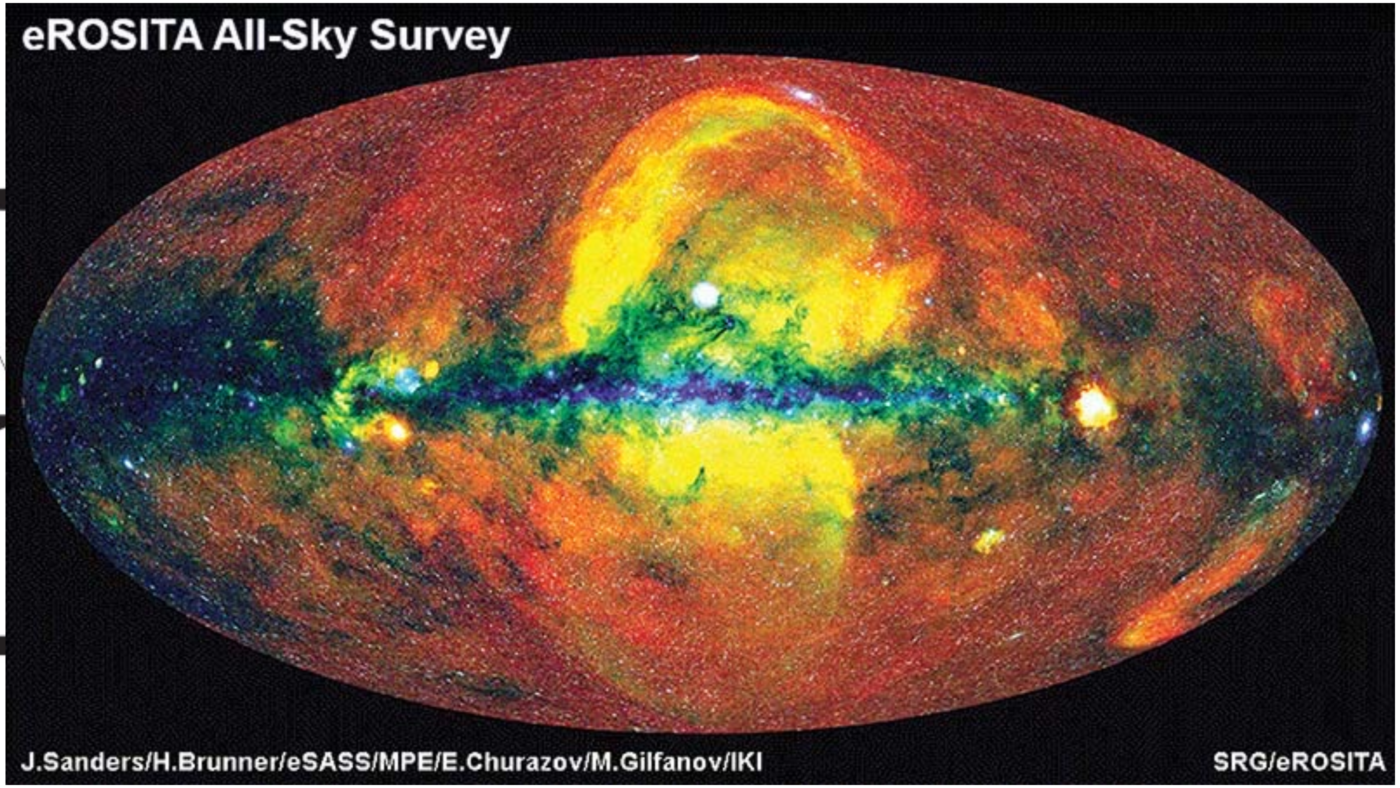
Proton mass puzzle

François Englert

Peter W. Higgs

Nobel 2013 With
Francois Englert
"Higgs Boson" that gives mass to quarks, electrons,….



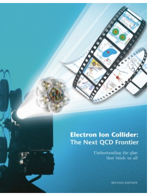


Add the masses of the quarks (HIGGS mechanism) together 1.78 x 10^{-26} grams

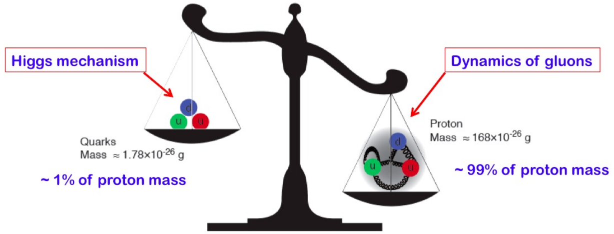
But the proton's mass is 168 x 10^{-26} grams

→ **only 1% of the mass of the protons (neutrons)** → Hence the Universe

→Where does the rest of the mass come from?

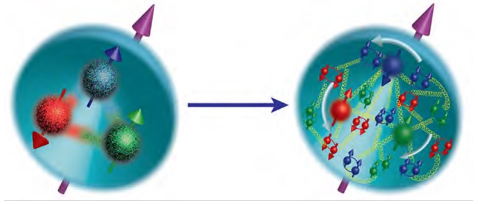


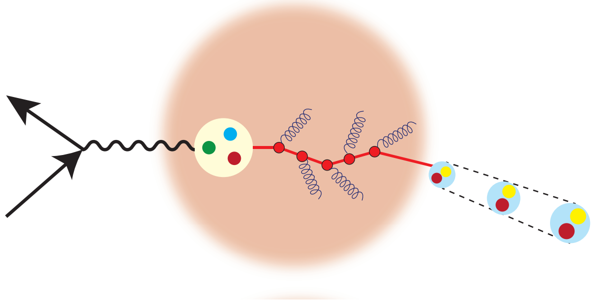
EIC Physics at-a-Glance



How are the sea quarks and gluons, and their spins, distributed in space and momentum inside the nucleon?

How do the nucleon properties (mass & spin) emerge from their interactions?

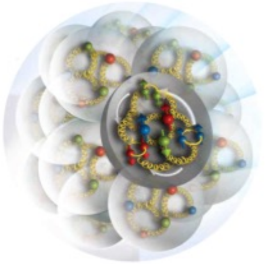




How do color-charged quarks and gluons, and colorless jets, interact with a nuclear medium?

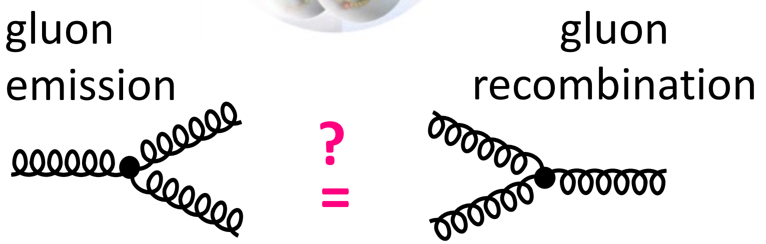
How do the confined hadronic states emerge from these quarks and gluons?

How do the quark-gluon interactions create nuclear binding?

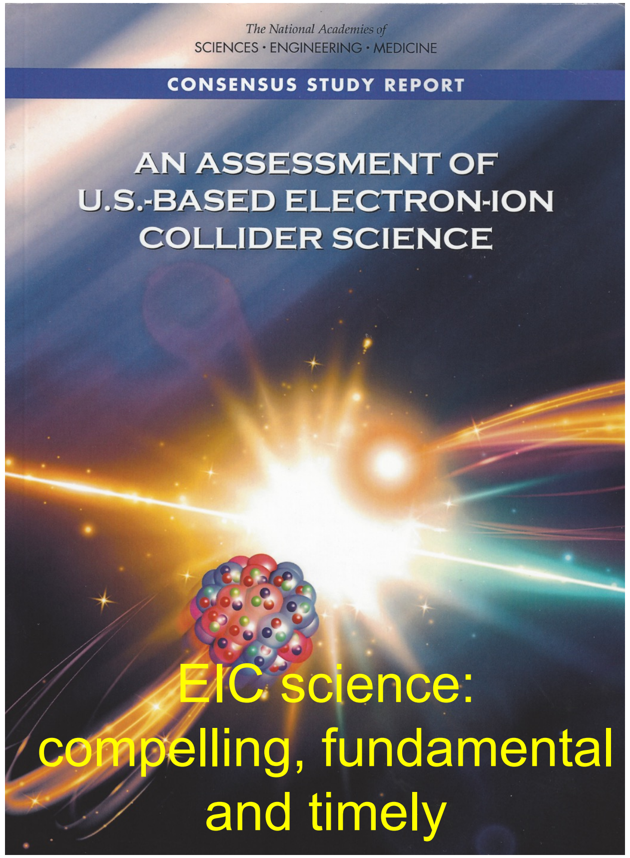


How does a dense nuclear environment affect the quarks and gluons, their correlations, and their interactions?

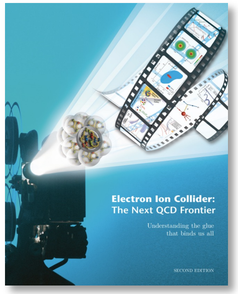
What happens to the gluon density in nuclei? Does it saturate at high energy, giving rise to a gluonic matter with universal properties in all nuclei, even the proton?



National Academy's Assessment



EIC science: compelling, fundamental and timely

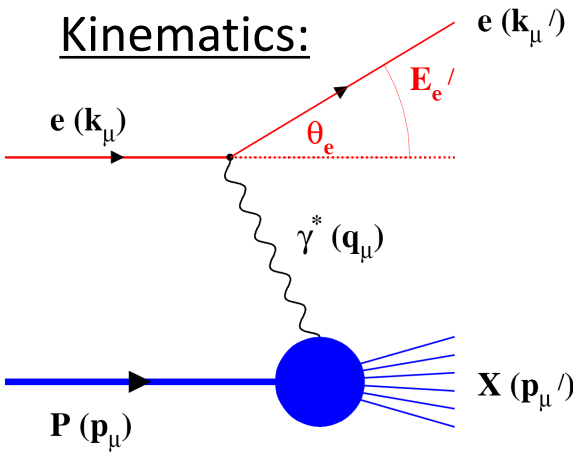


Machine Design Parameters:

- High luminosity: up to 10^{33}-10^{34} $cm^{-2}sec^{-1}$
  - a factor ~100-1000 times HERA
- Broad range in center-of-mass energy: ~20-100 GeV upgradable to 140 GeV
- Polarized beams e-, p, and light ion beams with flexible spin patterns/orientation
- Broad range in hadron species: protons…. Uranium
- Up to two detectors well-integrated detector(s) into the machine lattice

Deep Inelastic Scattering: Precision and control

Kinematics:



$$Q^2 = -q^2 = -(k_\mu - k'_\mu)^2$$

$$Q^2 = 2E_e E'_e(1-\cos\Theta_{e'})$$

Measure of resolution power

$$y = \frac{pq}{pk} = 1 - \frac{E'_e}{E_e}\cos^2\left(\frac{\theta'_e}{2}\right)$$

Measure of inelasticity

$$x = \frac{Q^2}{2pq} = \frac{Q^2}{sy}$$

Measure of momentum fraction of struck quark

Hadron :

$$z = \frac{E_h}{\nu}; p_t$$ with respect to γ^*

$$s = 4E_h E_e$$

High lumi & acceptance

Low lumi & acceptance

Exclusive DIS

detect & identify *everything* e+p/A → e'+h(π,K,p,jet)+…

Semi-inclusive events:

e+p/A → e'+h(π,K,p,jet)+X

detect the scattered lepton in coincidence with identified hadrons/jets

Inclusive events:

e+p/A → e'+X

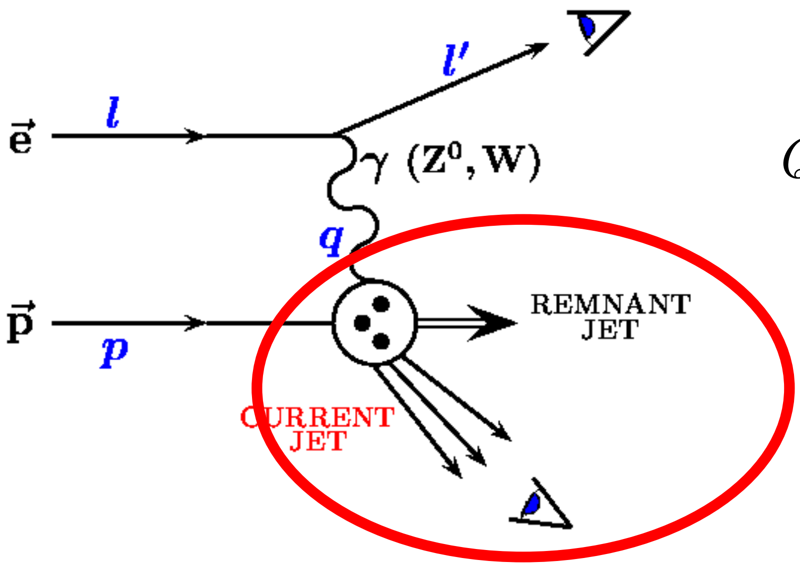
detect *only the scattered lepton* in the detector

Some times scattered electron can't be measured….

Reason:

1) Scattering angle so small that it is too close to the beam pipe
2) Radiative correction too large, i.e. electron lost its energy due to Initial State Radiation or Brehmstrahlung through material -- So the kinematic reconstruction unreliable.

What to do? Then see if we can reconstruct the hadronic final state?



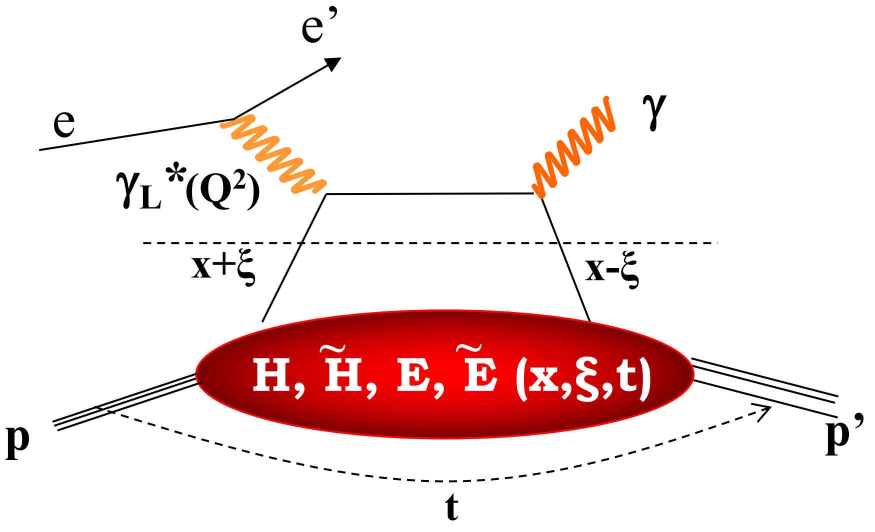
$$y = \frac{E_j}{2E_e}(1 - cos\theta_j)$$
$$Q^2 = E_j^2 sin^2\theta_j/(1 - y)$$

$$x = \frac{E_j}{2E_p}(1 + cos\theta_j)/(1 - y)$$
$$E_j = yE_e + x(1 - y)E_p$$
$$cos\theta_j = \frac{-yE_e + (1-y)xE_p}{y\,E_e + (1-y)xE_p}$$
$$E_j^2 sin^2\theta_j = 4\,xy(1-y)E_eE_p = Q^2\,(1-y)$$

$$y_{JB} = \frac{1}{2E_e}\sum_h (E_h - p_{Zh})$$
$$Q^2_{JB} = \frac{(\sum_h p_{Xh})^2 + (\sum_h p_{Yh})^2}{1 - y_{JB}}$$
$$x_{JB} = Q^2_{JB}/(y_{JB}s)$$

Deep Inelastic Scattering: Deeply Virtual Compton Scattering

Kinematics:



Exclusive measurement:
e + (p/A) → e'+ (p'/A')+ γ / J/ψ / ρ / ϕ
detect **all** event products in the detector

Special sub-event category rapidity gap events
e + (p/A) → e' + γ / J/ψ / ρ / ϕ / jet
Don't detect **(p'/A') in final state**

$$Q^2 = -q^2 = -(k_\mu - k'_\mu)^2$$

Measure of resolution power

$$Q^2 = 2E_e E'_e(1 - \cos\Theta_{e'})$$

$$y = \frac{pq}{pk} = 1 - \frac{E'_e}{E_e}\cos^2\left(\frac{\theta'_e}{2}\right)$$

Measure of inelasticity

$$x_B = \frac{Q^2}{2pq} = \frac{Q^2}{sy}$$

Measure of momentum fraction of struck quark

$$t = (p - p')^2, \xi = \frac{x_B}{2 - x_B}$$

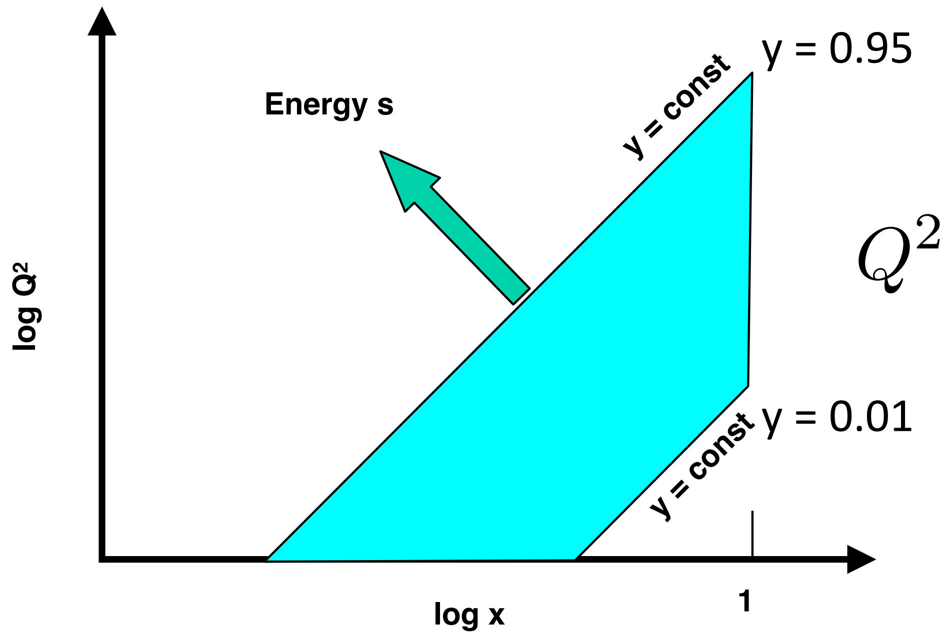
Complete set of variables for DIS e-p:

https://core.ac.uk/download/pdf/25211047.pdf

We will use some of these more often than others, you should know them all.

E_p	proton beam energy
E_e	electron beam energy
$p = (0, 0, E_p, E_p)$	four momentum of incoming proton with mass m_p
$e = (0, 0, -E_e, E_e)$	four momentum of incoming electron
$e' = (E'_e sin\theta'_e, 0, E'_e cos\theta'_e, E'_e)$	four momentum of scattered electron
$s = (e + p)^2 = 4E_pE_e$	square of total ep c.m. energy
$q^2 = (e - e')^2 = -Q^2$	mass squared of exchanged current J
	= square of four momentum transfer
$\nu = q \cdot p/m_p$	energy transfer by J in p rest system
$\nu_{max} = s/(2m_p)$	maximum energy transfer
$y = (q \cdot p)/(e \cdot p) = \nu/\nu_{max}$	fraction of energy transfer
$x = Q^2/(2q \cdot p) = Q^2/(ys)$	Bjorken scaling variable
$q_c = x \cdot p + (e - e')$	four momentum of current quark
$M^2 = (e' + q_c)^2 = x \cdot s$	mass squared of electron - current quark system.

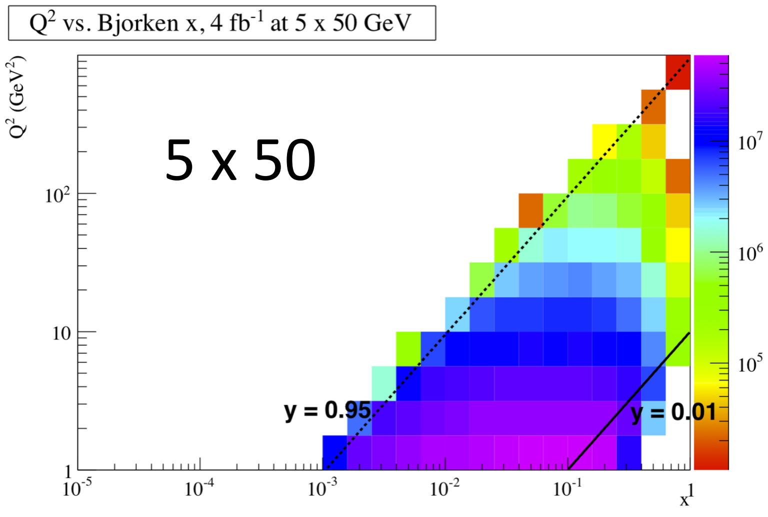
The x-Q^2 plane…

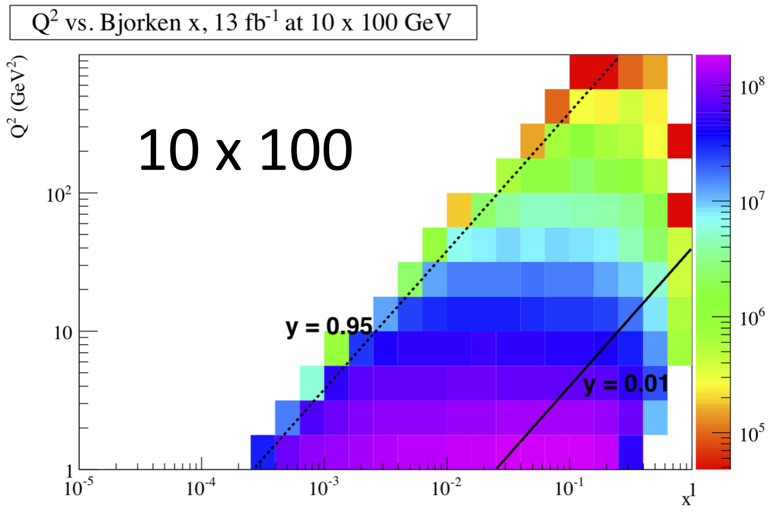


$$Q^2 \approx s \cdot x \cdot y$$

- Low-x reach requires large $\sqrt{s}$
- Large-Q^2 reach requires large $\sqrt{s}$
- *y* at colliders typically limited to 0.95 < y < 0.01

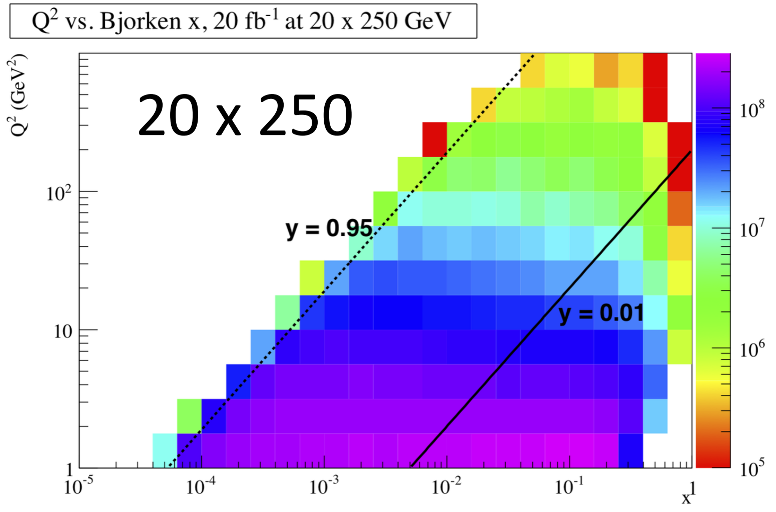
Kinematic coverage as a function of energy of collisions





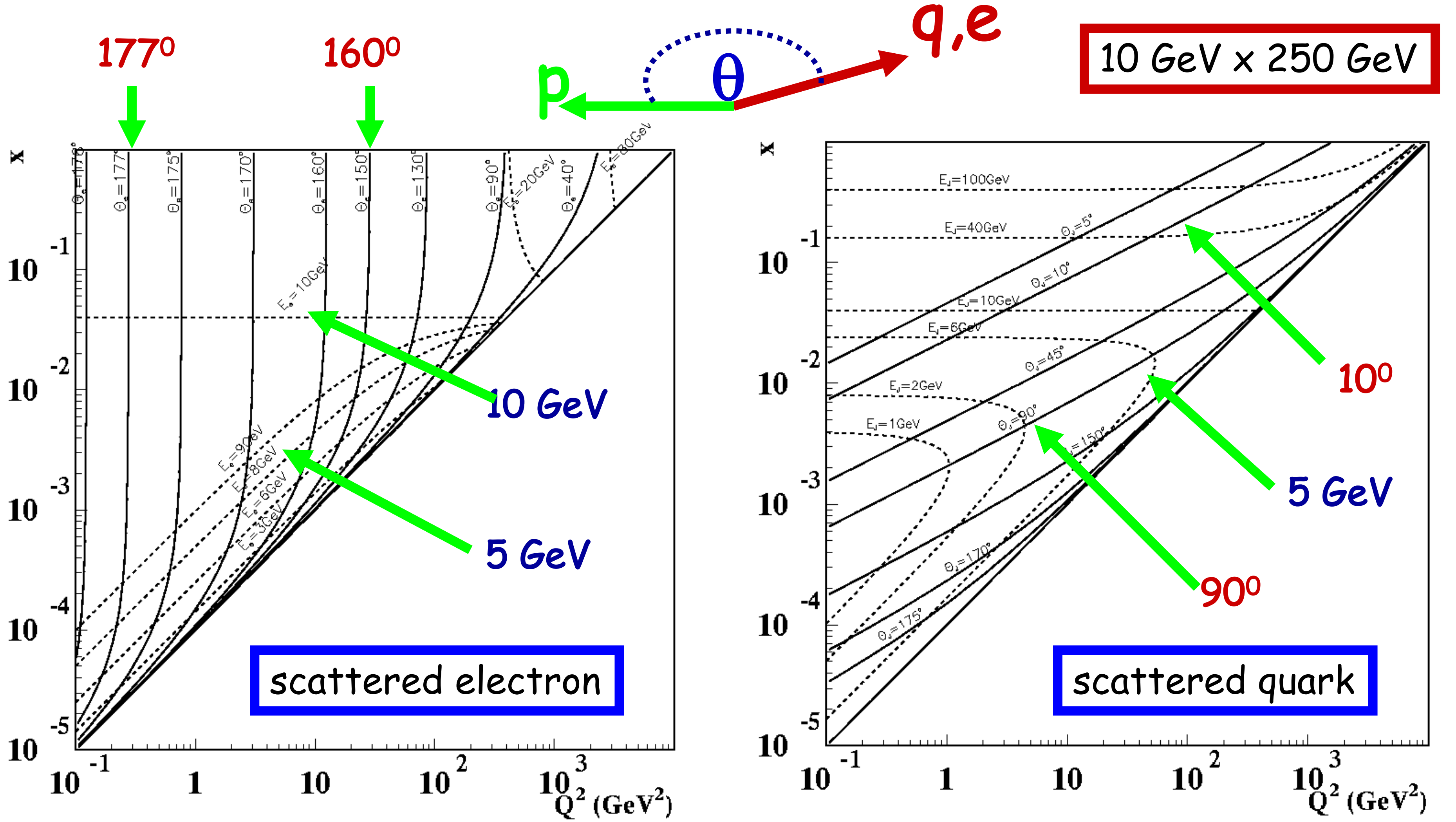
As beam energies increase, so does the x, Q^2 coverage of the collider: 5, 10 and 20 GeV electrons colliding with 50, 100 and 250 GeV protons

y = 0.95 and 0.01 are shown on all plots (they too shift as function of energy of collisions)

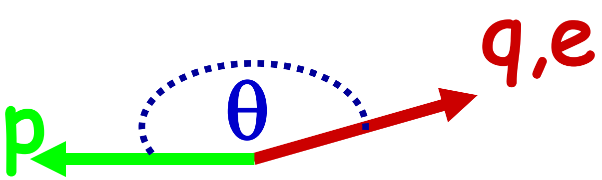


Home Work: Where do electrons and quarks go?

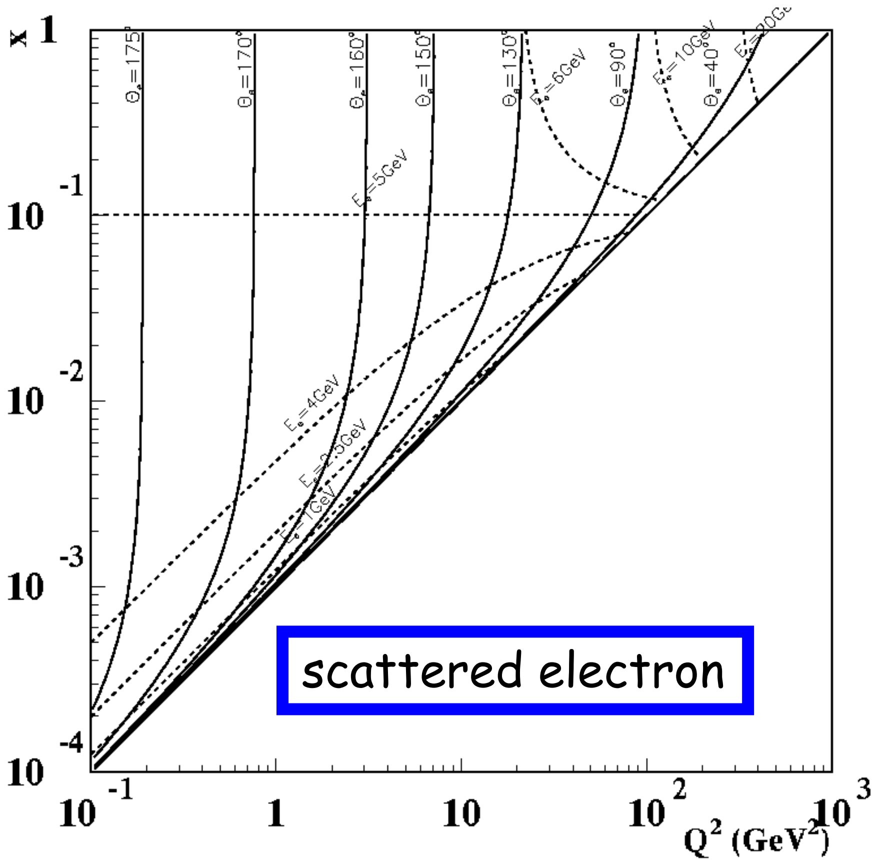
Angles measured w.r.t. proton direction

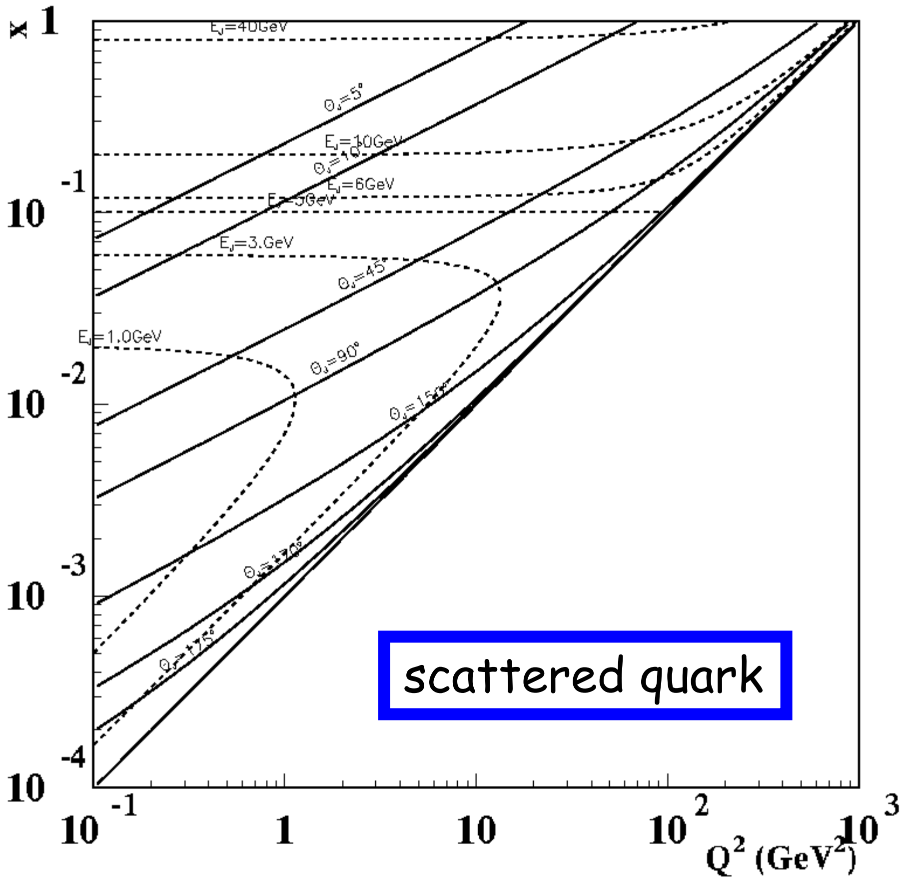


Electron, Quark Kinematics



5 GeV x 50 GeV

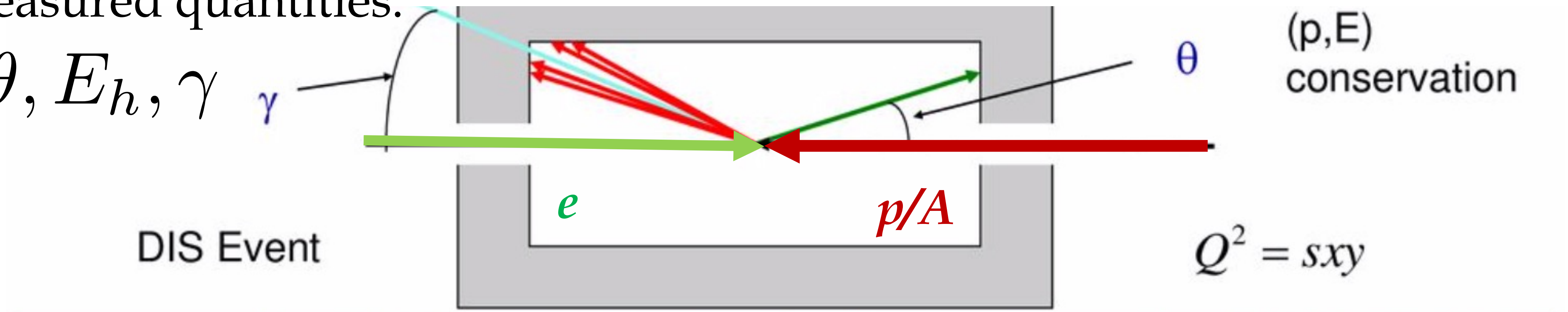




There are multiple ways to reconstruct events:

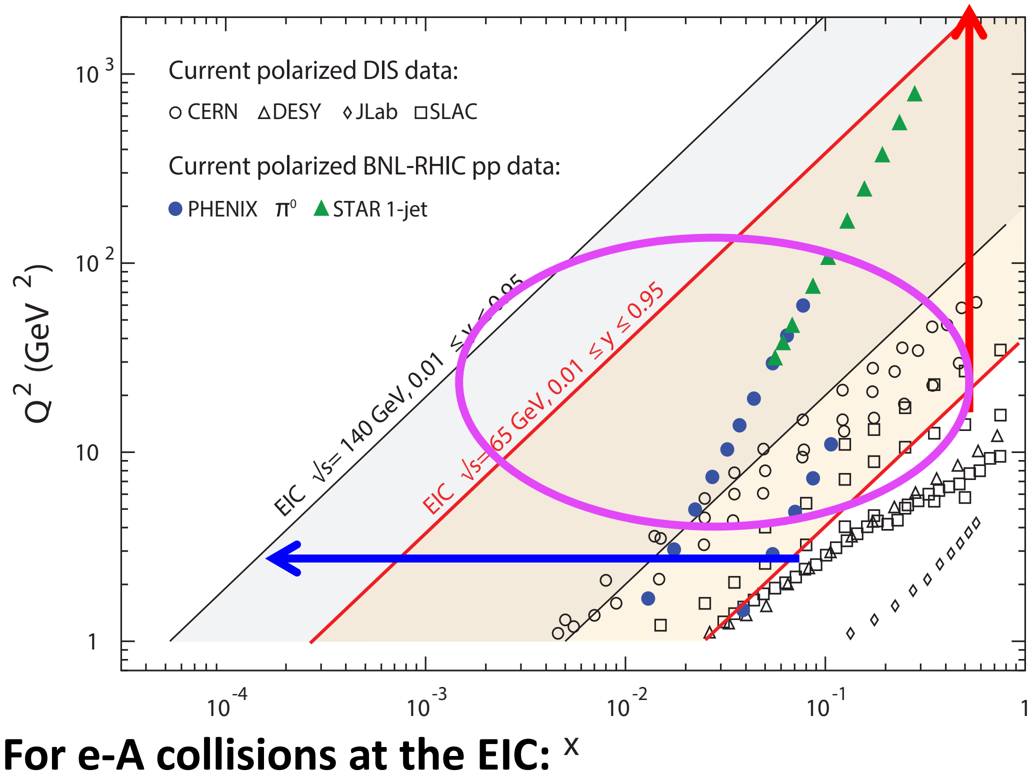
Four measured quantities:

$$E'_e, \theta, E_h, \gamma$$



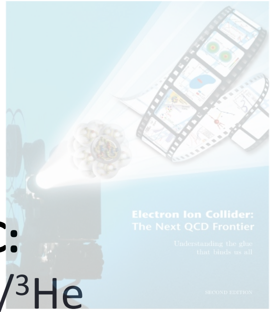
EIC Science ➔ what it could provide

EIC: Kinematic reach & properties



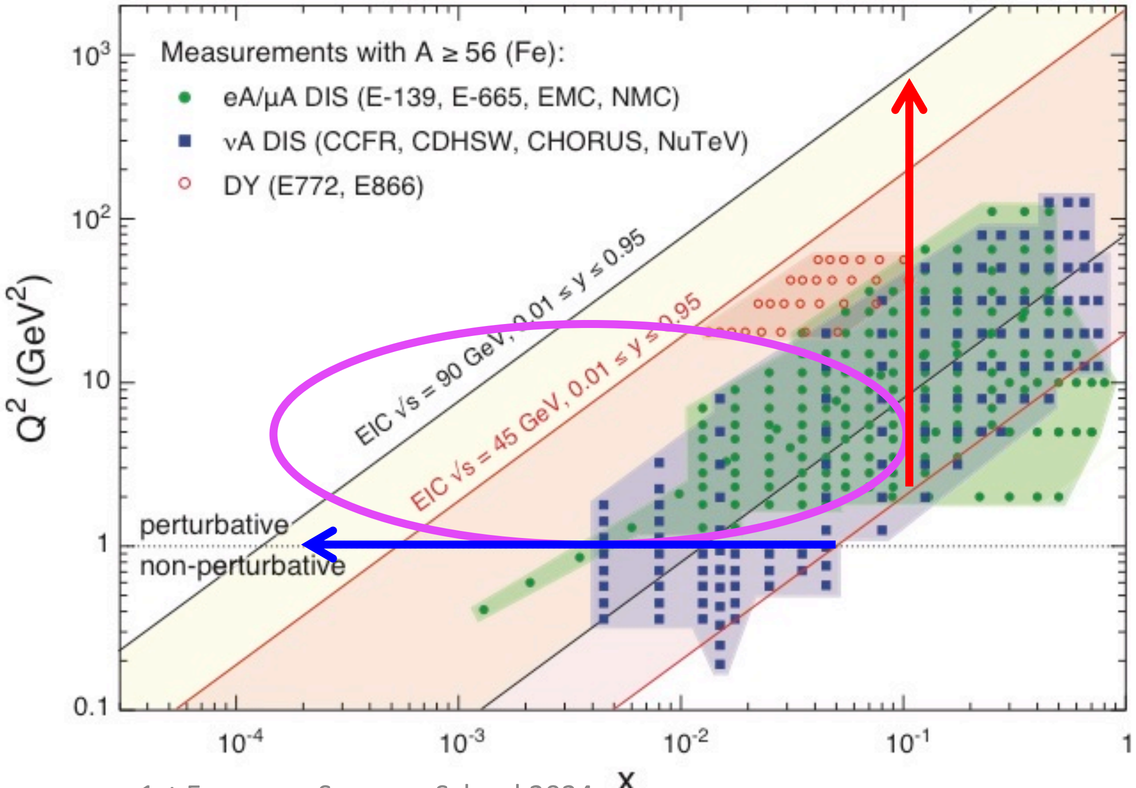
For e-A collisions at the EIC:

- ✓ **Wide range in nuclei**
- ✓ Lum. per nucleon same as e-p
- ✓ Variable center of mass energy
- ✓ Wide x range (evolution)
- ✓ Wide x region (reach high gluon densities)



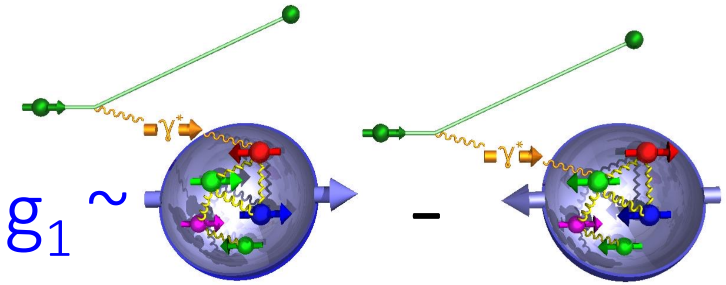
For e-N collisions at the EIC:

- ✓ **Polarized** beams: e, p, d/^{3}He
- ✓ Variable center of mass energy
- ✓ Wide Q^2 range → evolution
- ✓ Wide x range → spanning valence to low-x physics



Nucleon Spin: Precision with EIC

$$\frac{1}{2} = \left[\frac{1}{2}\Delta\Sigma + L_Q\right] + [\Delta g + L_G]$$

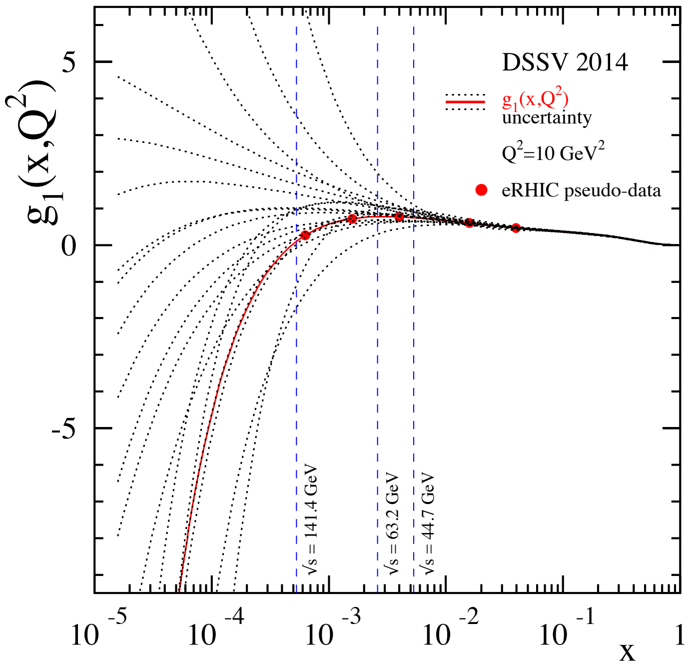


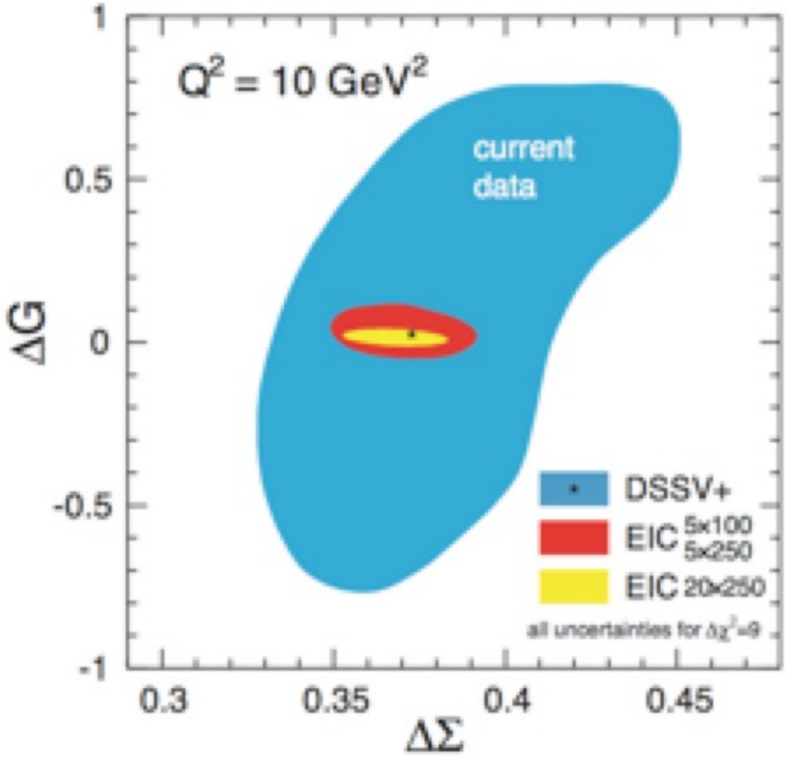
$\Delta\Sigma/2$ = Quark contribution to Proton Spin
Δg = Gluon contribution to Proton Spin
L_Q = Quark Orbital Ang. Mom
L_G = Gluon Orbital Ang. Mom

Spin structure function g_1 needs to be measured over a large range in x-Q^2

Precision in $\Delta\Sigma$ and Δg ➔ A clear idea Of the magnitude of L_Q+L_G = L

SIDIS: strange and charm quark spin contributions

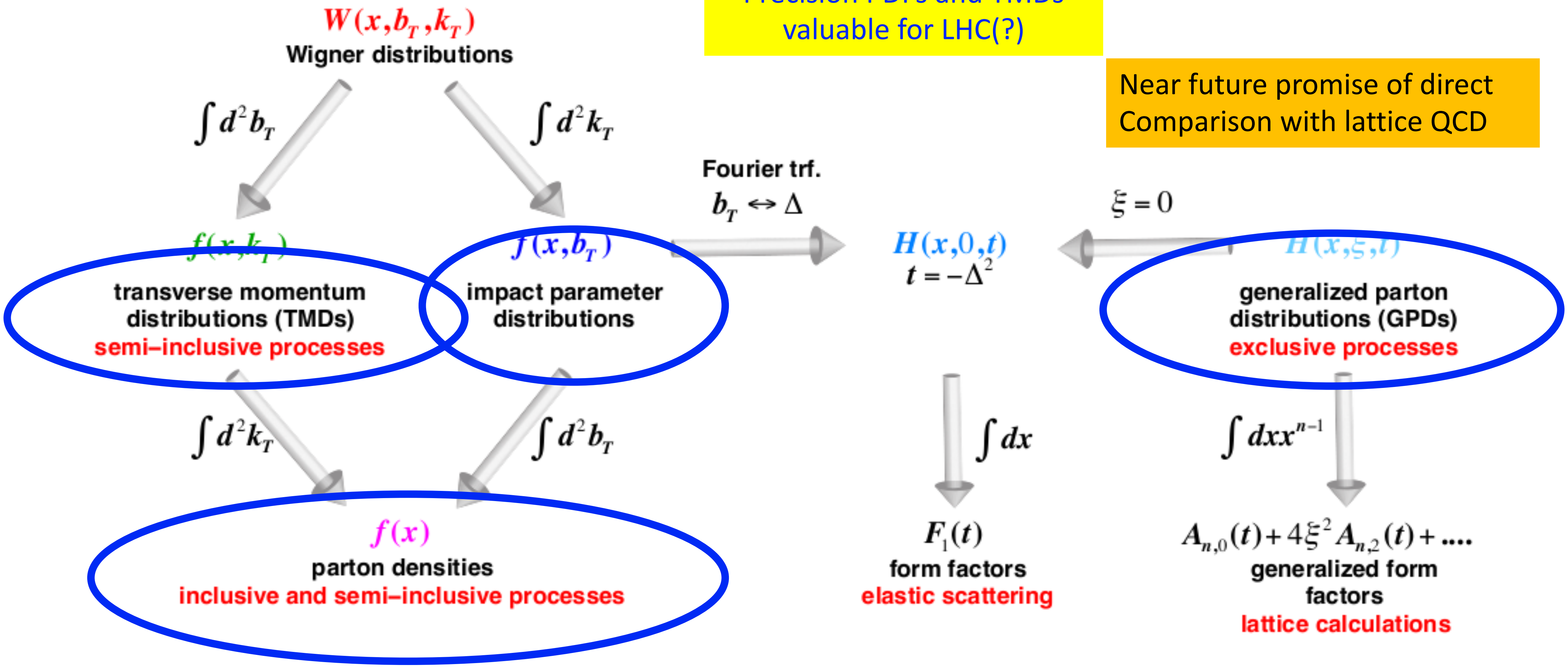




2+1D Imaging of hadrons: beyond precision PDFs

Precision PDFs and TMDs valuable for LHC(?)

Near future promise of direct Comparison with lattice QCD



3-Dimensional Imaging Quarks and Gluons

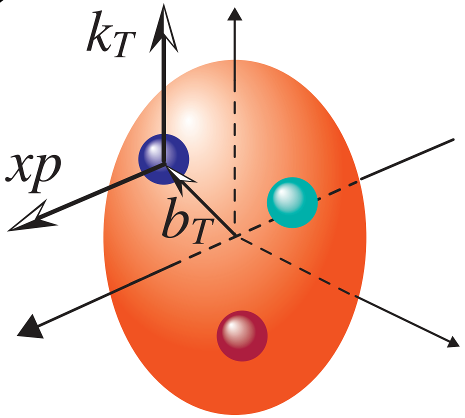
Wigner functions $W(x,b_T,k_T)$

offer unprecedented insight into confinement and chiral symmetry breaking.

$$W(x,b_T,k_T)$$

Momentum space $\int d^2b_T$ → $f(x,k_T)$

$\int d^2k_T$ ***Coordinate space*** → $f(x,b_T)$

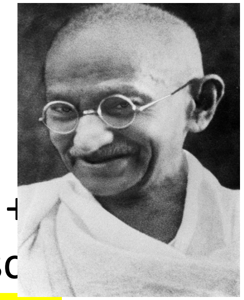


Spin-dependent 3D momentum space images from semi-inclusive scattering
→ **TMDs**

Spin-dependent 2D coordinate space (transverse) + 1D (longitudinal momentum) images from exclusive scattering
→ **GPDs**

Position and momentum → Orbital motion of quarks and gluons

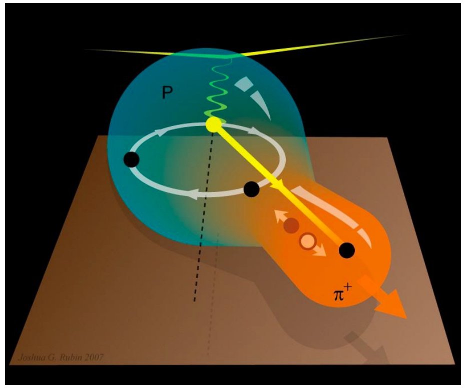
Possible direct access to gluon Wigner function through diffractive di-jet measurements at an EIC: Y. Hatta et al. PRL 16, 022301 (2016



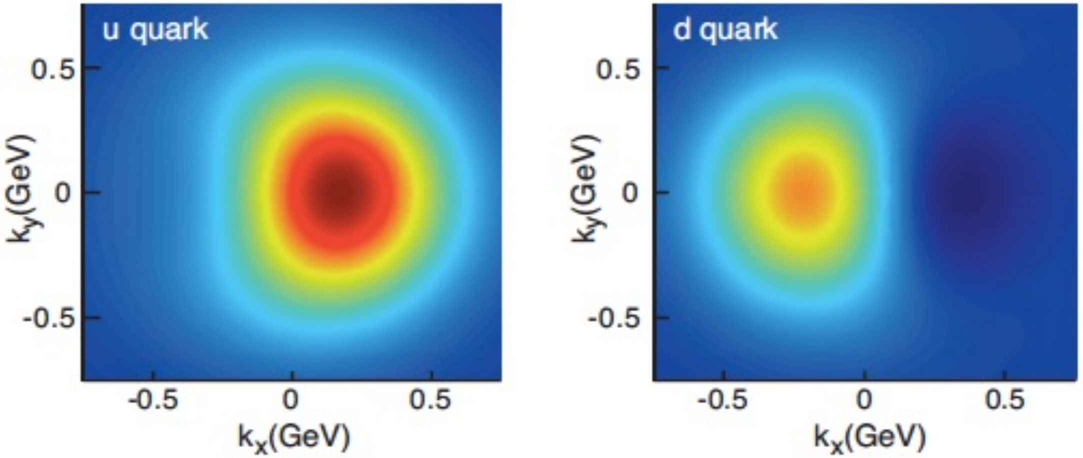
2+1 D partonic image of the proton with the EIC

Spin-dependent 3D momentum space images from semi-inclusive scattering (SIDS)

Transverse Momentum Distributions



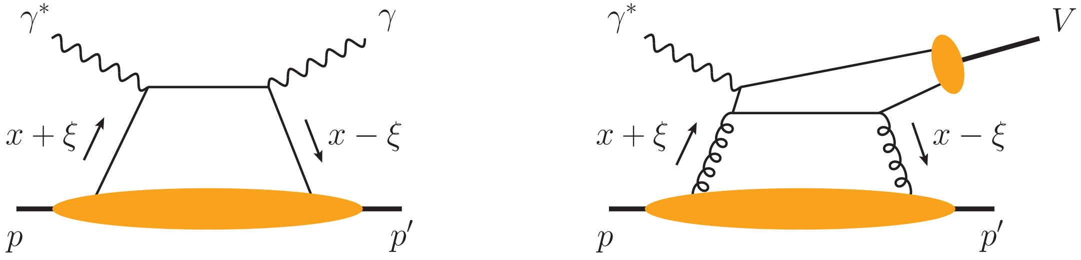
Possible measurements of K (s) and D (c)



Spin-dependent 2D coordinate space (transverse) + (longitudinal momentum) images from exclusive sc

Transverse Position Distributions

Quarks Motion



Gluons: Only @ Collider

Deeply Virtual Compton Scattering

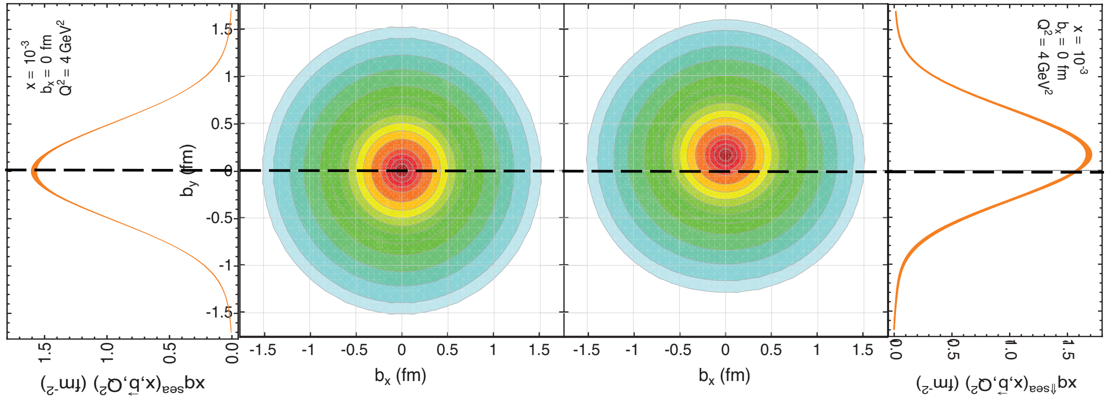
Measure all three final states

e + p → e'+ p'+ γ

Fourier transform of momentum transferred=(p-p') → Spatial distribution

2D position distribution for sea-quarks

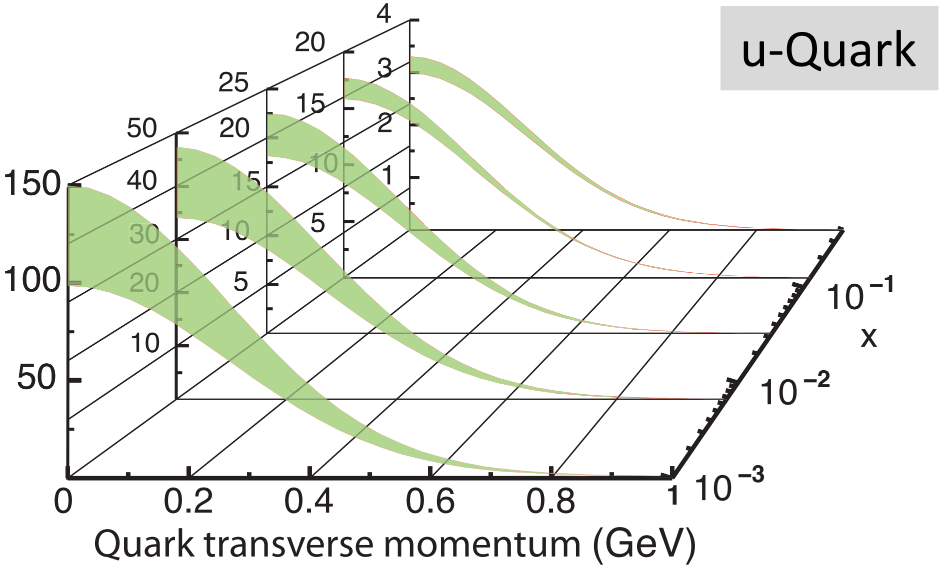
unpolarized polarized



2+1 D partonic image of the proton with the EIC

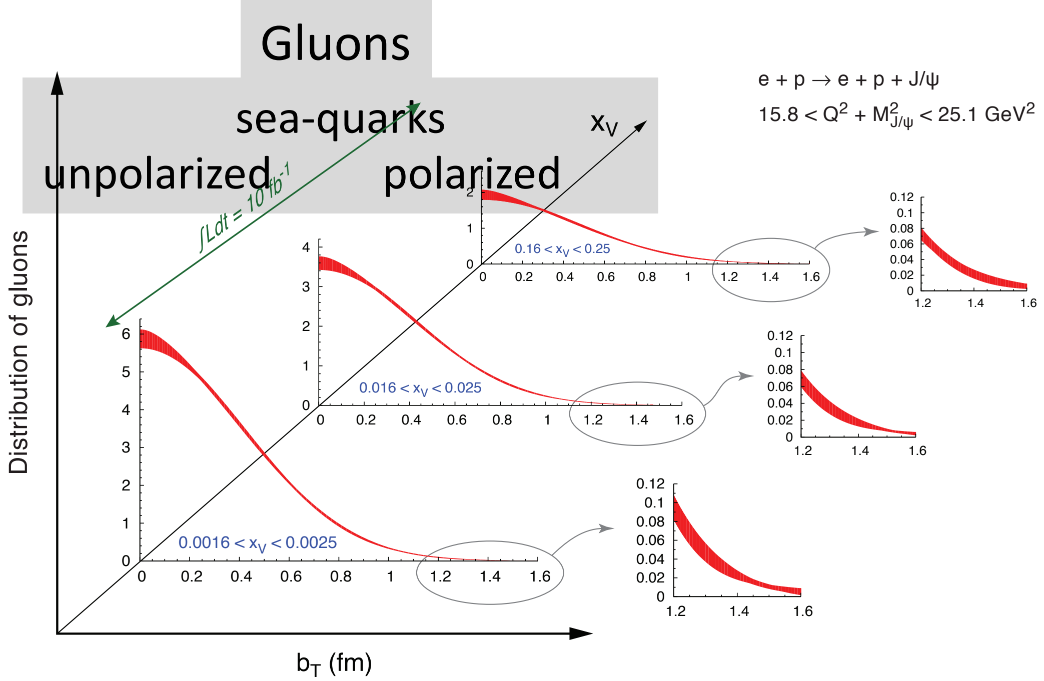
Spin-dependent 3D momentum space images from semi-inclusive scattering

Transverse Momentum Distributions



Spin-dependent 2D coordinate space (transverse) + 1D (longitudinal momentum) images from exclusive scattering

Transverse Position Distributions

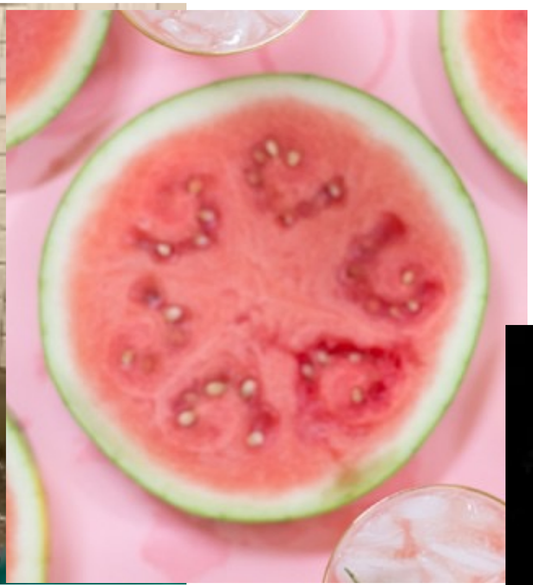


"Color form factor" of proton…

Study of internal structure of a watermelon:

A-A (RHIC)

1) Violent collision of melons

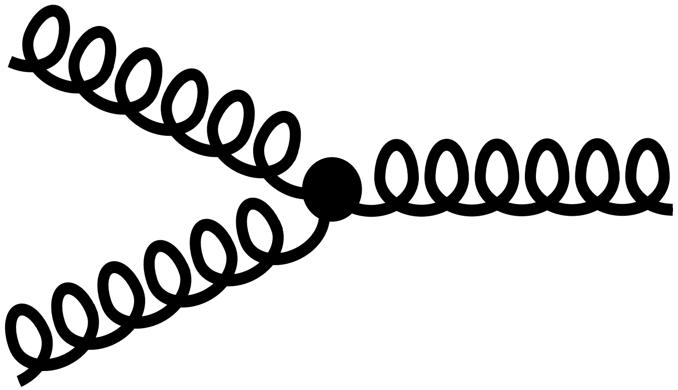


2) Cutting the watermelon with a knife

Violent DIS e-A (EIC)

3) MRI of a watermelon

Non-Violent e-A (EIC)



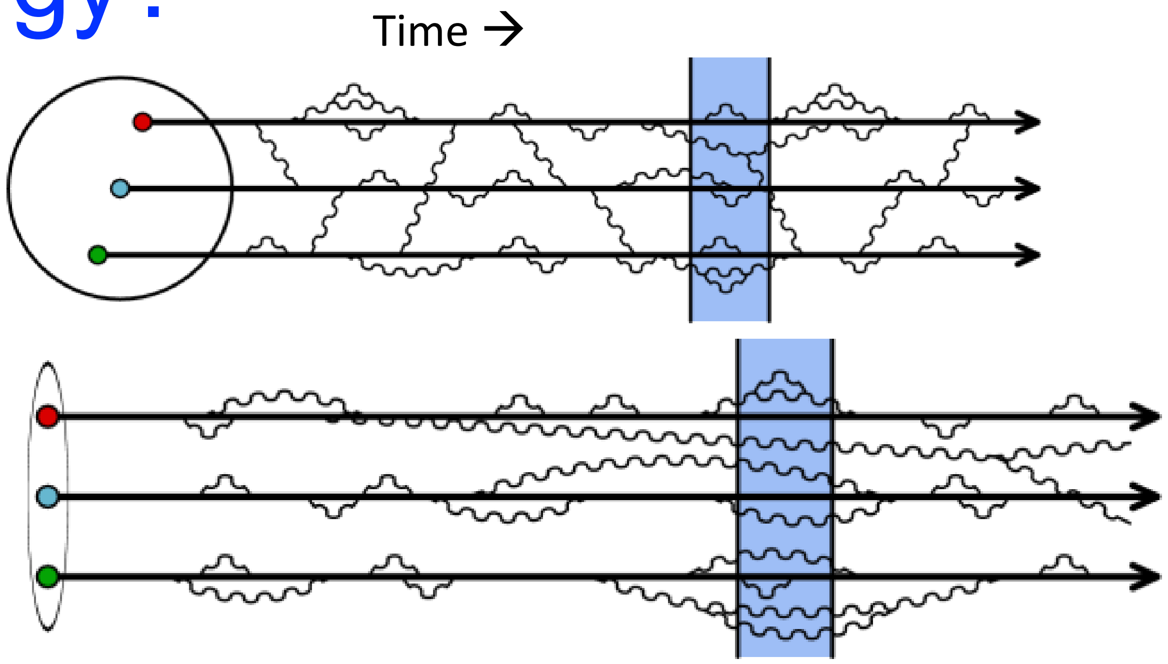
Consequence of gluon self interactions ➔ non-linear GDLAP evolution… ?

Particularly at high energy ➔ low-x

How does a Proton look at low and very high energy?

Low energy: High x
Regime of fixed target exp.

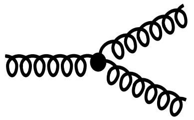
High energy: Low- x
Regime of a Collider



Cartoon of boosted proton

At high energy:

- Wee partons fluctuations are time dilated in strong interaction time scales
- Long lived gluons radiate smaller x gluons ➔ which in turn radiate more… a chain reaction leading to a **runaway growth?**



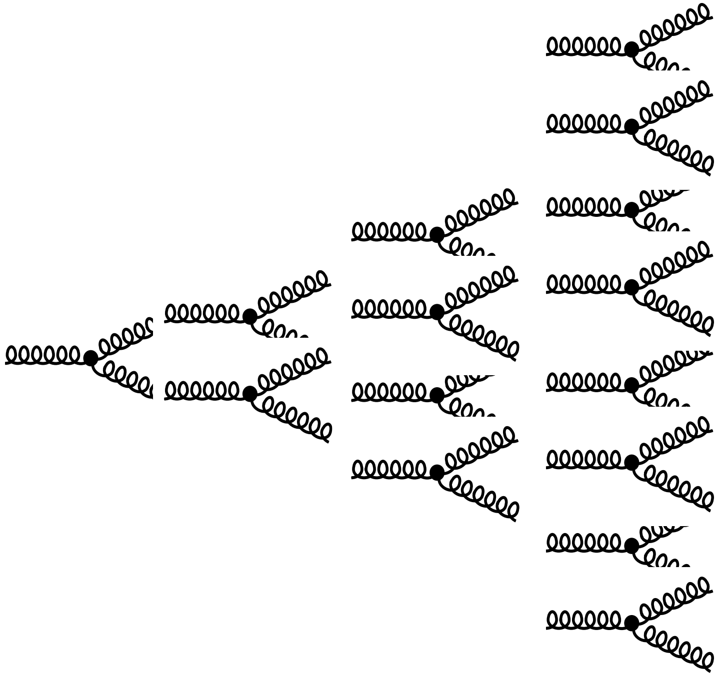
Gluon splitting

Gluon and the consequences of its interesting properties:

Gluons carry color charge ➔ Can interact with other gluons!

"…The result is a self catalyzing enhancement that leads to a runaway growth. A small color charge in isolation builds up a big color thundercloud…."

F. Wilczek, in "Origin of Mass"
Nobel Prize, 2004



? Infinity?

No!

Gluon and the consequences of its interesting properties:

Gluons carry color charge ➔ Can interact with other gluons!

“…The result is a self catalyzing enhancement that leads to a runaway growth.
A small color charge in isolation builds up a big color thundercloud….”

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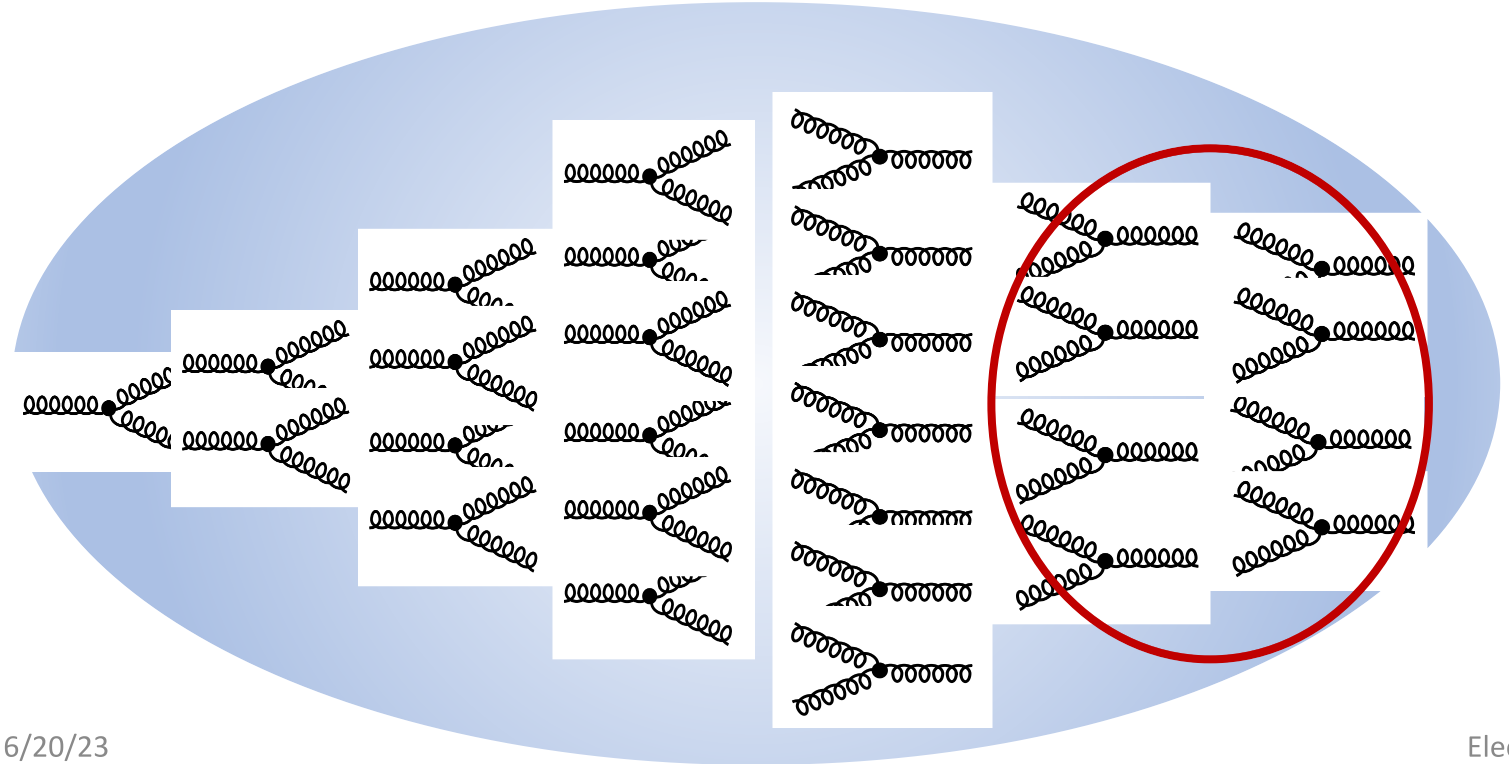
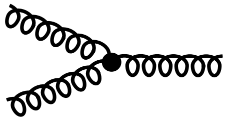
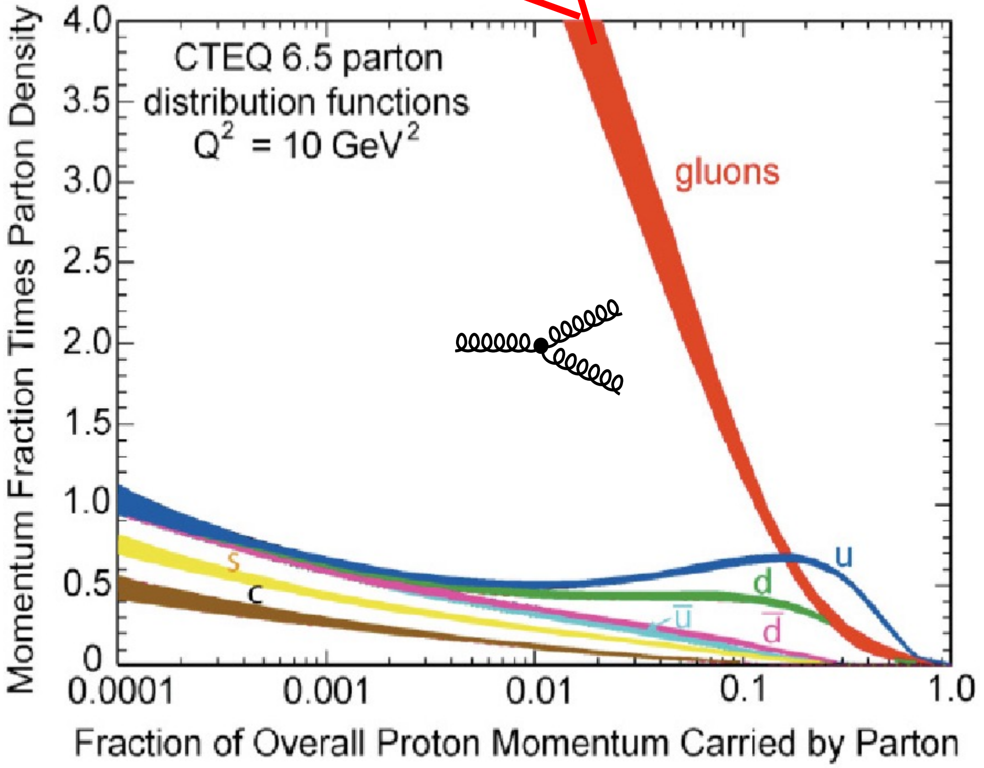


In search of a new state of matter!



?



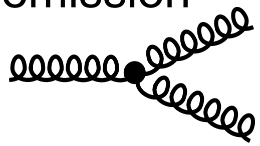
Experimental evidence needed

What could tame the low-x rise?

Can EIC access this region?

QCD inherently has the needed mechanism for this taming but we don't know when it gets triggered.

gluon emission



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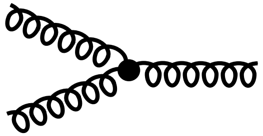
gluon recombination



At Q_S

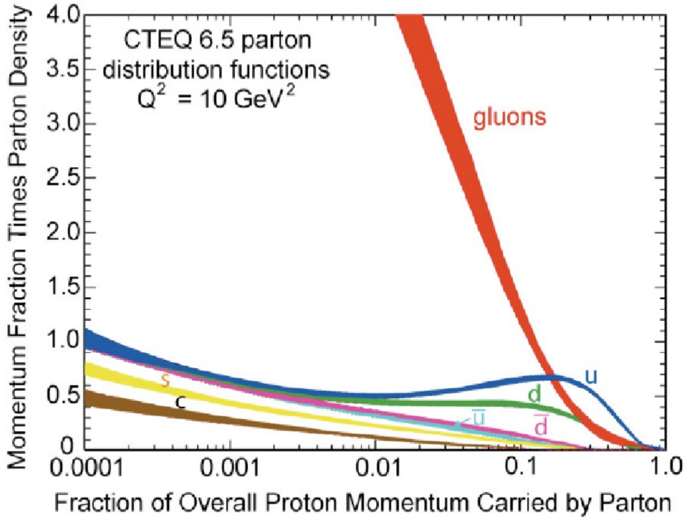
Observation of gluon recombination effects

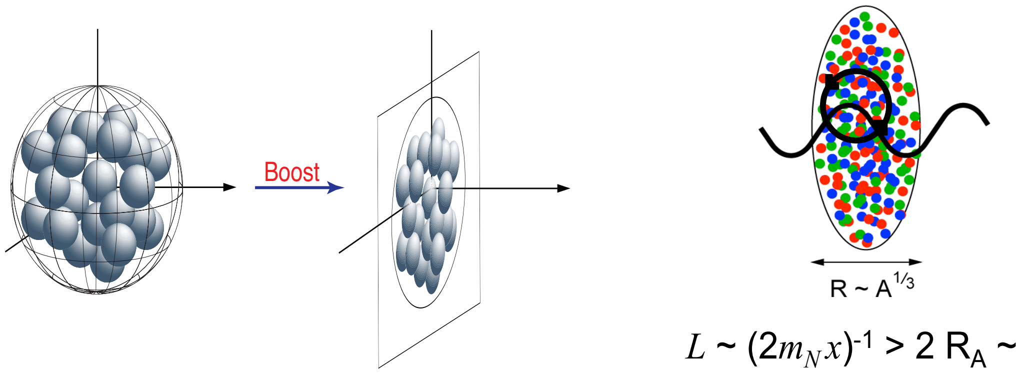
➔Is there such new state of matter?

→ "Color Glass Condensate"

→ 50-100 times higher energy density than the core of the neutron star

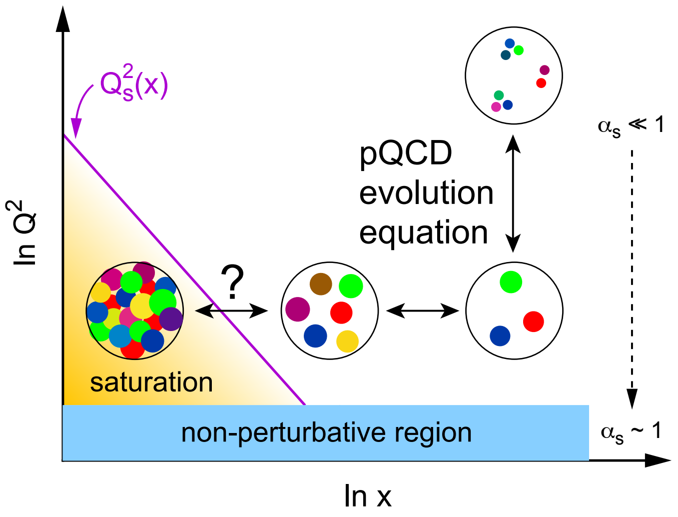
Low x physics with nuclei



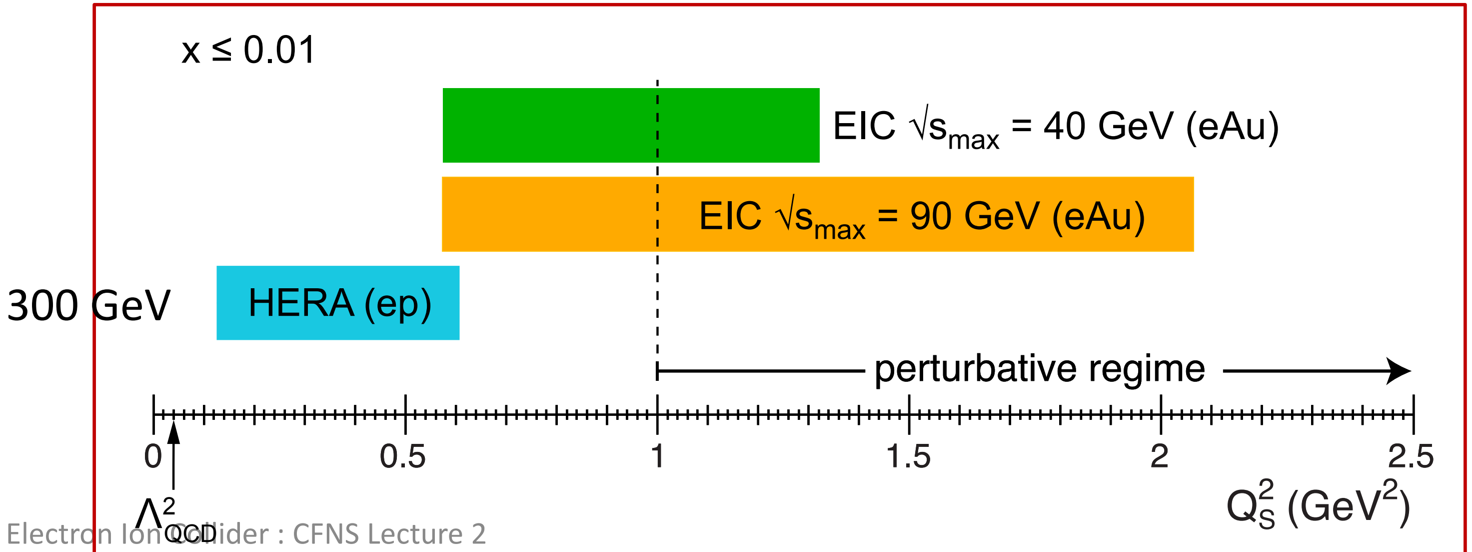


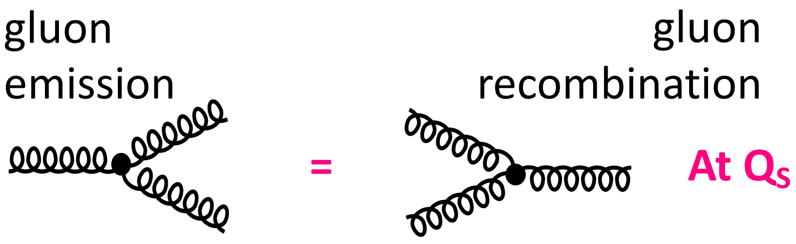
$$(Q_s^A)^2 \approx cQ_0^2 \left[\frac{A}{x}\right]^{1/3}$$

$L \sim (2m_N x)^{-1} > 2\, R_A \sim A^{1/3}$



Accessible range of saturation scale $Q_s^{\ 2}$ at the EIC with e+A collisions.
arXiv:1708.01527



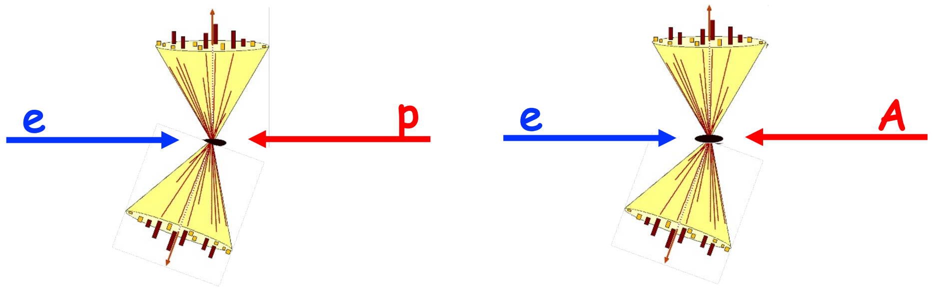


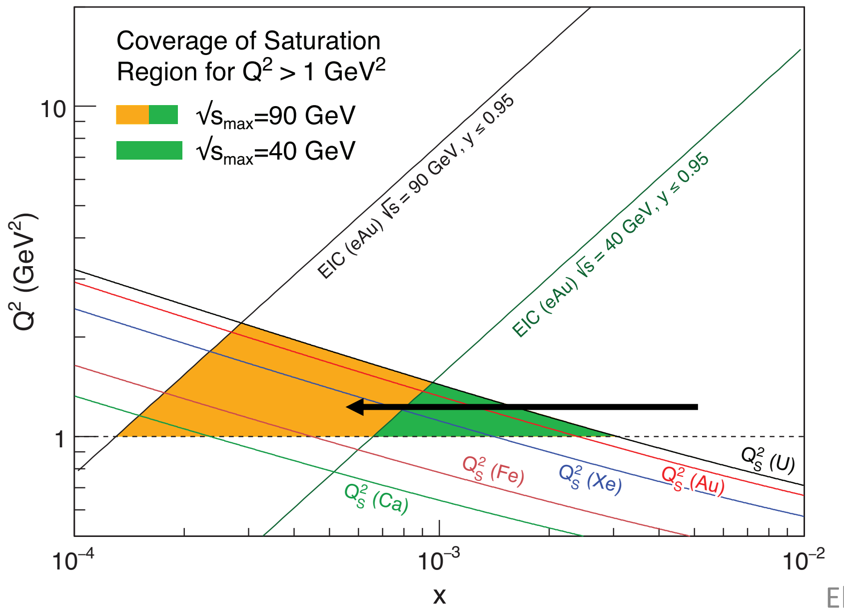
Can EIC discover a new state of matter?

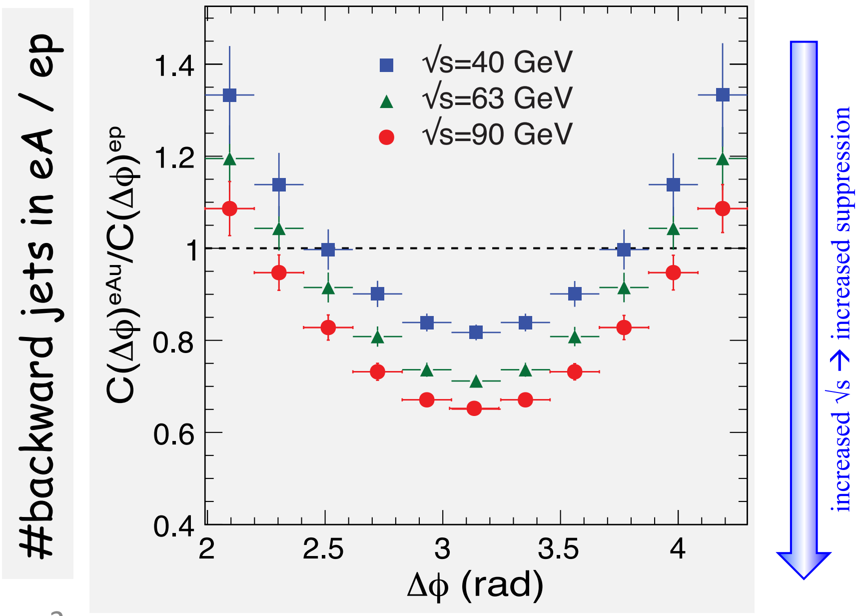
EIC provides an absolutely unique opportunity
to have very high gluon densities
→ electron – lead collisions
combined with an unambiguous observable

EIC will allow to unambiguously
map the transition from a non-saturated to
saturated regime

counting experiment of
Di-jets in ep and eA
Saturation:
Disappearance of backward jet in eA





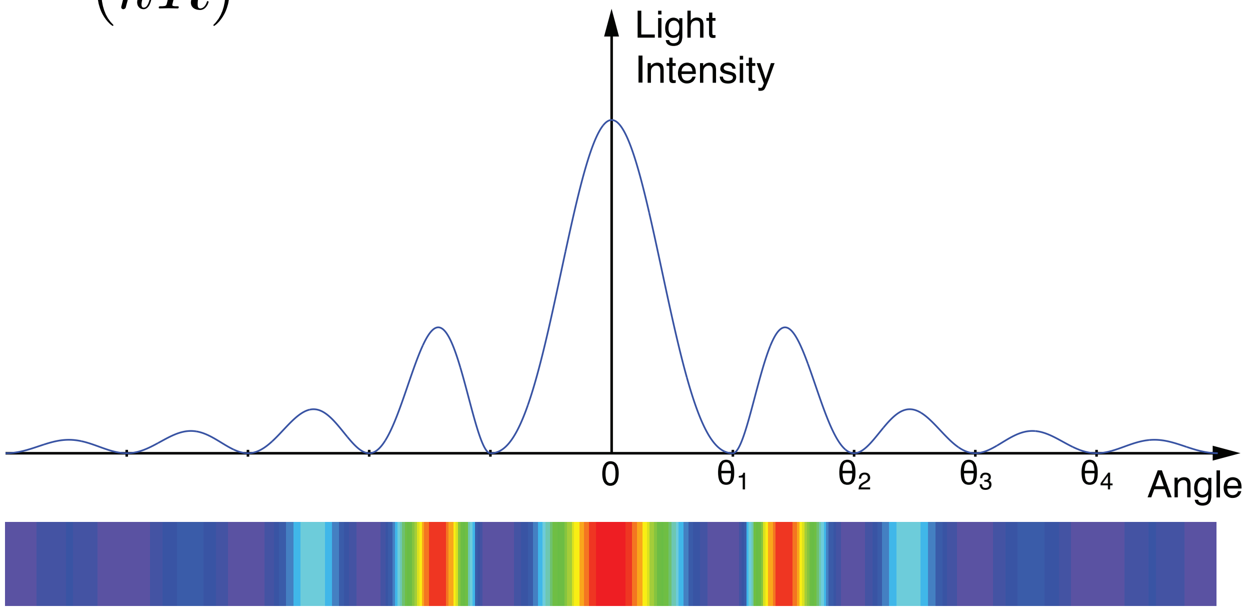


Diffraction in Optics and high energy scattering

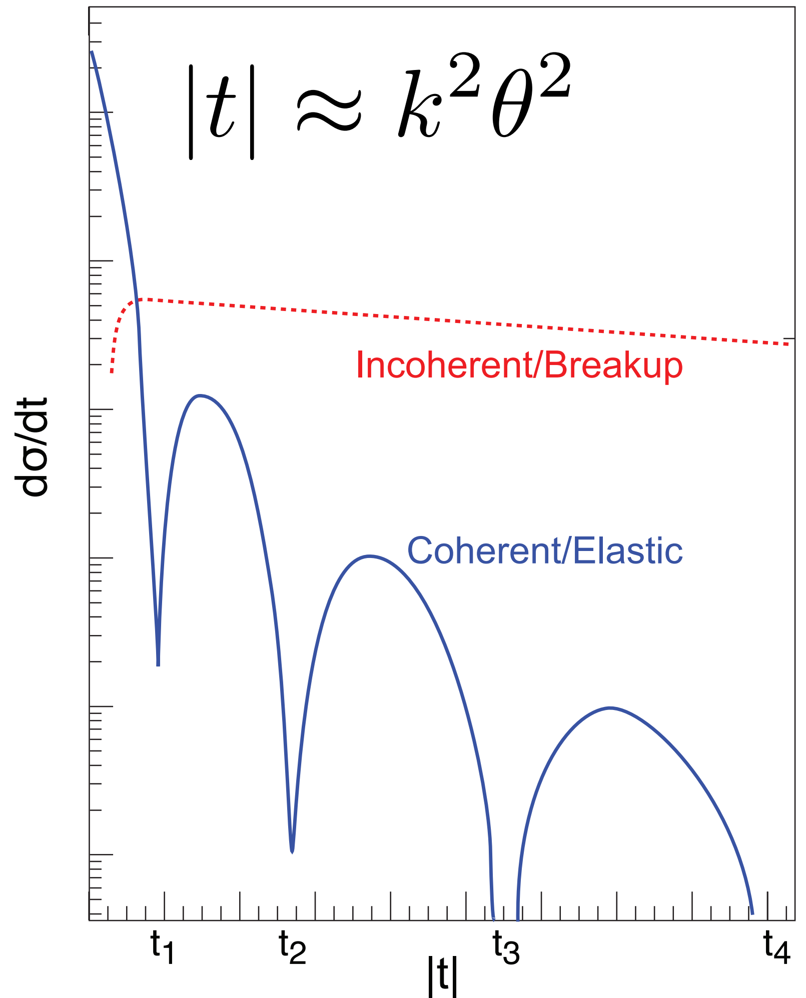
Light with wavelength λ obstructed by an opaque disk of radius R suffers diffraction:

k → wave number

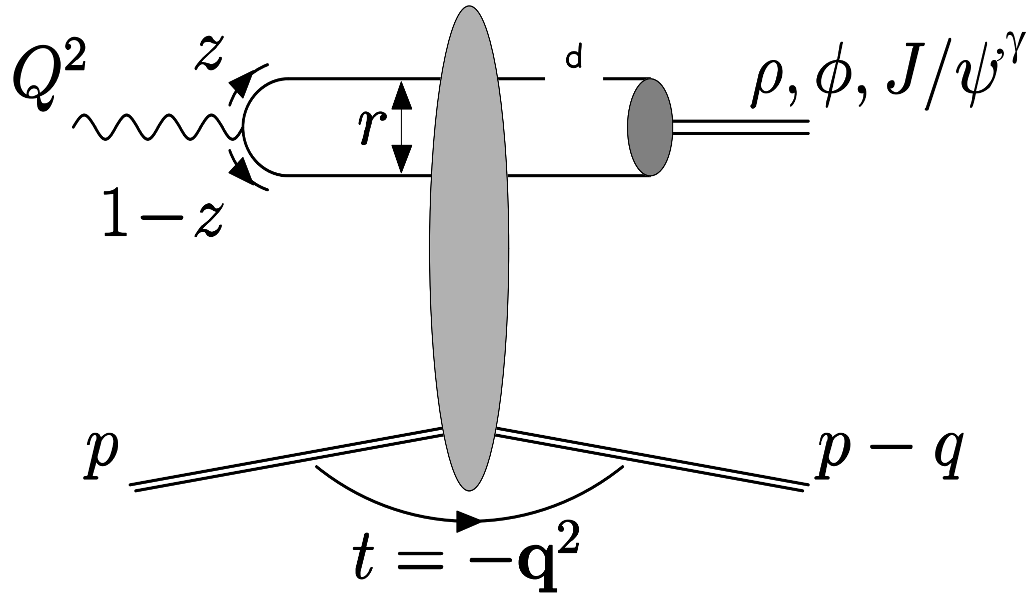
$$\theta_i \sim \frac{1}{(kR)}$$



Calculation of e-A diffraction



Transverse imaging of the gluons nuclei

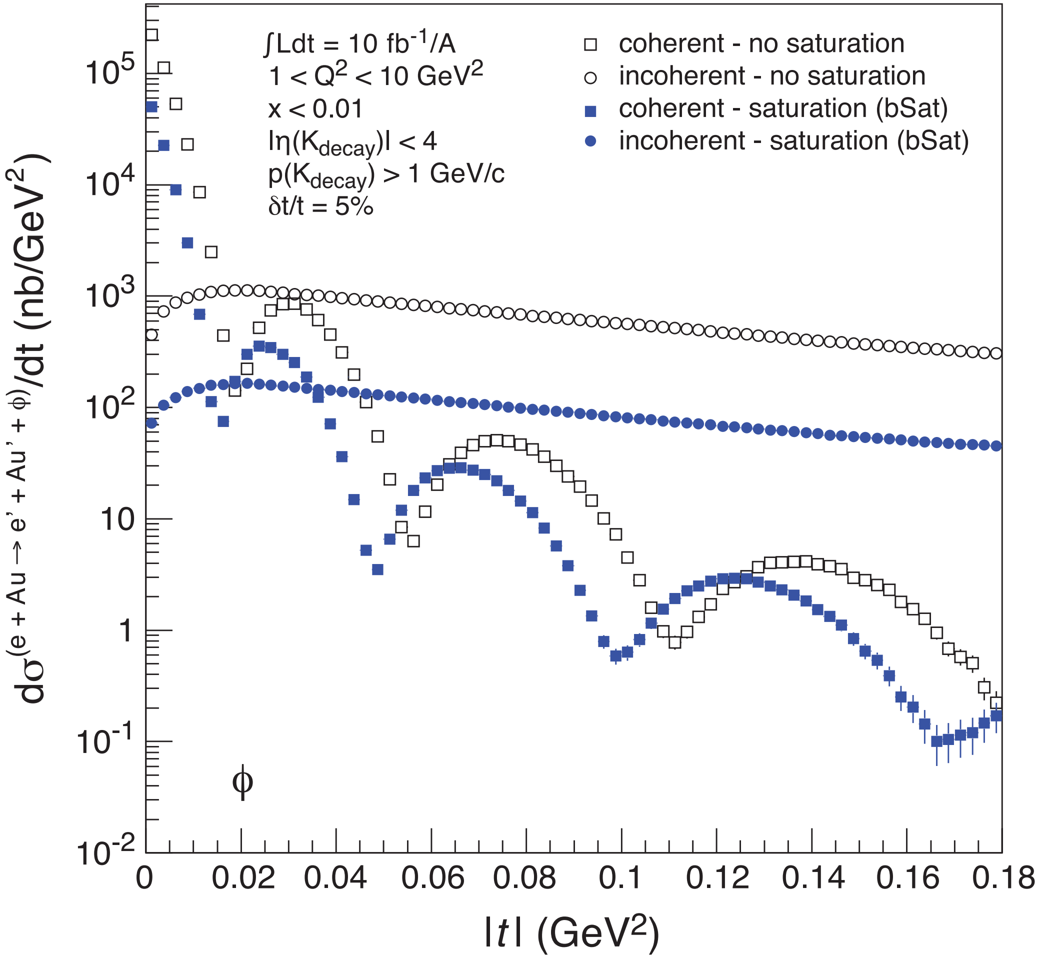


➔Does low x dynamics (Saturation) modify the transverse gluon distribution?

Experimental challenges being studied.

Diffractive vector meson production in **e-Au**

Diff. MC: “Sartre”



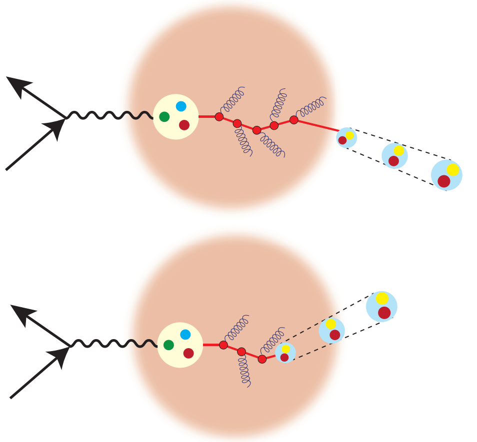
Simulation study by Toll & Ullrich

Emergence of Hadrons from Partons

Nucleus as a Femtometer sized filter

Unprecedented ν, the virtual photon energy range @ EIC : *precision & control*

$$\nu = \frac{Q^2}{2mx}$$



Control of ν by selecting kinematics;

Also under control the nuclear size.

(colored) Quark passing through cold QCD matter emerges as color-neutral hadron ➔

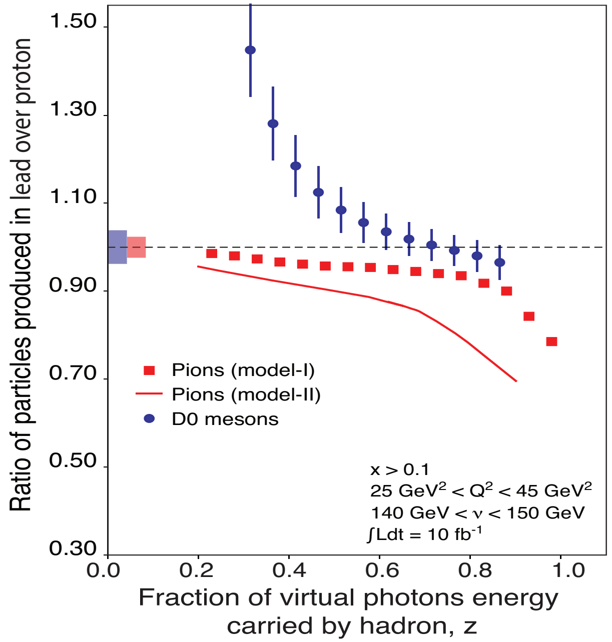
Clues to color-confinement?

Study in light quarks

vs.

heavy quarks

Energy loss by light vs. heavy quarks:



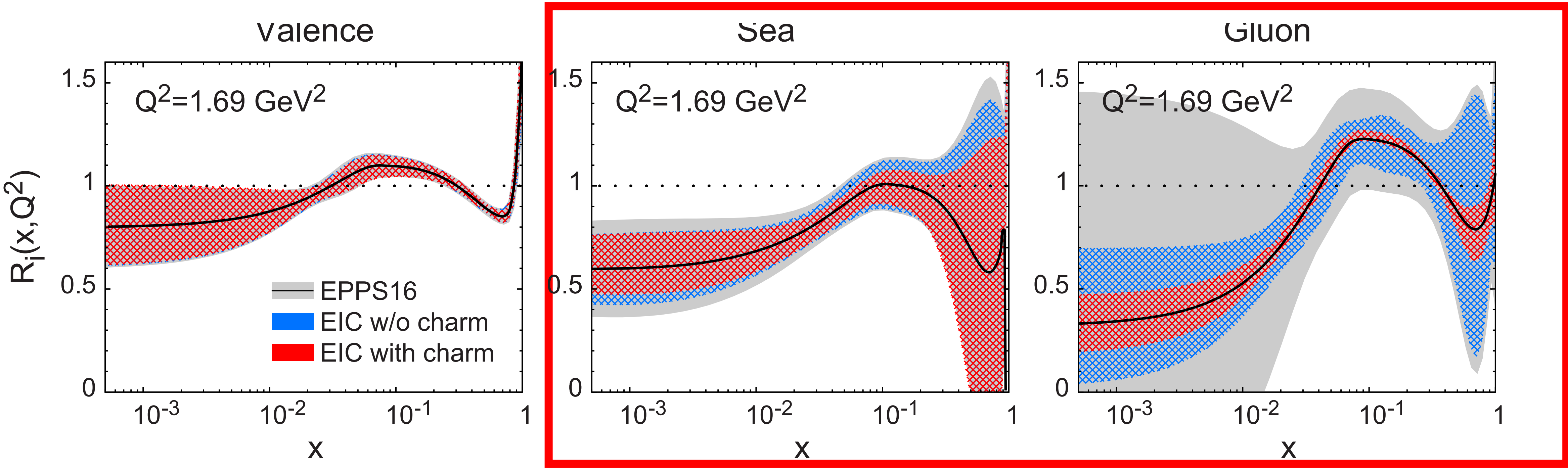
Identify π vs. D^0 **(charm)** mesons in e-A collisions:

Understand energy loss of light vs. heavy quarks traversing the cold nuclear matter:

Connect to energy loss in Hot QCD

Need the collider energy of EIC and its control on parton kinematics

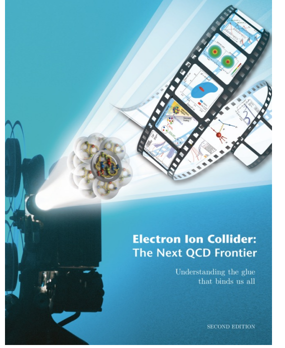
EIC: impact on the knowledge of 1D Nuclear PDFs



Ratio of Parton Distribution Functions of Pb over Proton:

- ❖ Without EIC, large uncertainties in **nuclear sea quarks and gluons** ➔ With EIC significantly reduced uncertainties
- ❖ Complementary to RHIC and LHC pA data. Provides information on initial state for heavy ion collisions.
- ❖ Does the nucleus behave like a proton at low-x? ➔ such color correlations relevant to the understanding of astronomical objects

Physics @ the US EIC beyond the EIC’s core science



New Studies with proton or neutron target:

- Impact of precision measurements of unpolarized PDFs at high x/Q^2, on LHC-Upgrade results(?)
- What role would TMDs in e-p play in W-Production at LHC? Gluon TMDs at low-x!
- Heavy quark and quarkonia (c, b quarks) studies with 100-1000 times lumi of HERA
- Does polarization of play a role (in all or many of these?)

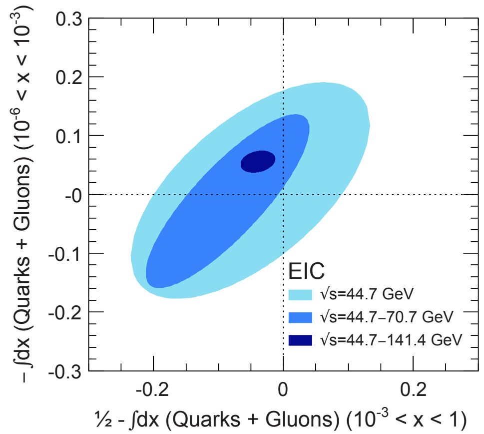
Physics with nucleons and nuclear targets:

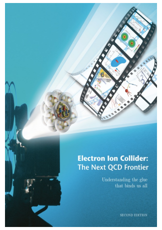
- Quark Exotica: 4,5,6 quark systems…? Much interest after recent LHCb led results.
- Physic of and with jets with EIC as a precision QCD machine:
  - Internal structure of jets : novel new observables, energy variability, polarization, beam species
  - Entanglement, entropy, connections to fragmentation, hadronization and confinement
  - Studies with jets: Jet propagation in nuclei… energy loss in cold QCD medium
- Connection to p-A, d-A, A-A at RHIC and LHC
- Polarized light nuclei in the EIC

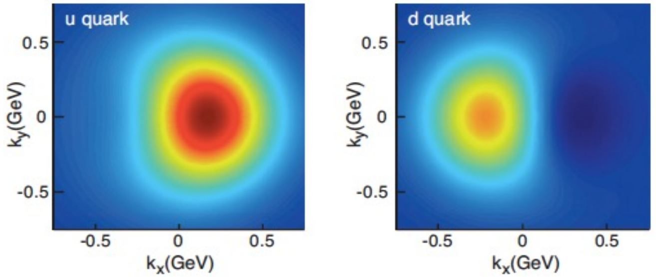
Precision electroweak and BSM physics:

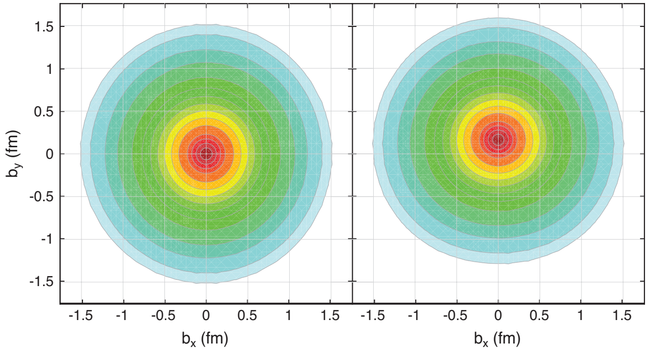
- Electroweak physics & searches beyond the SM: Parity, charge symmetry, lepton flavor violation

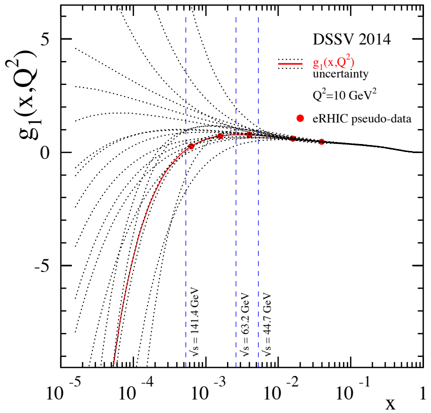
EIC science highlights

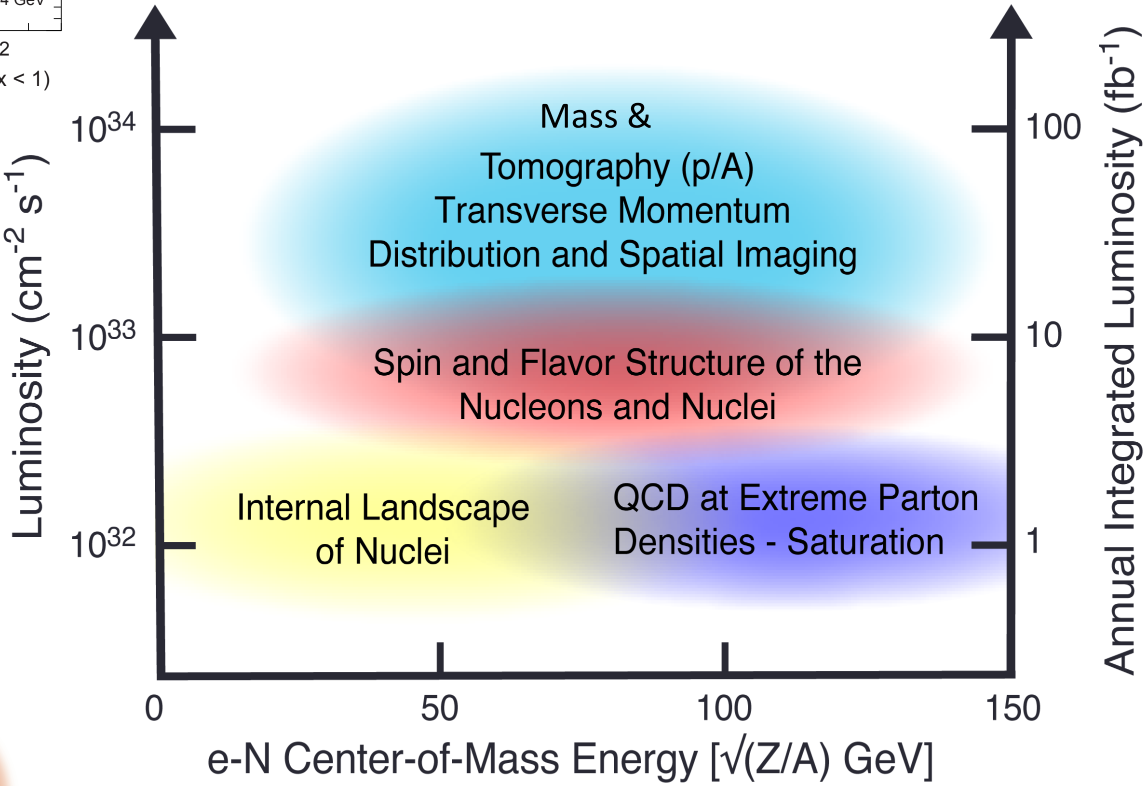


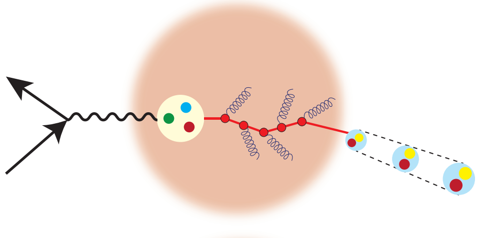


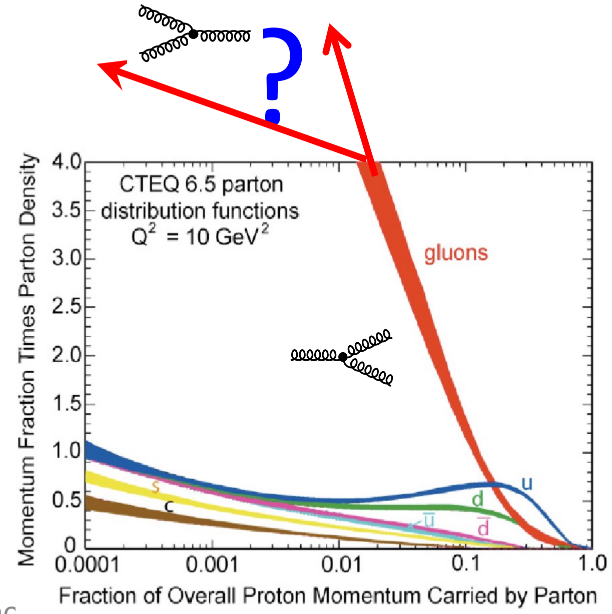


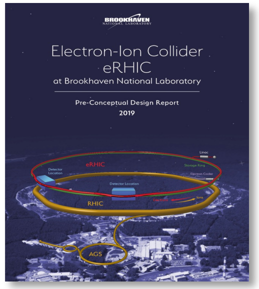




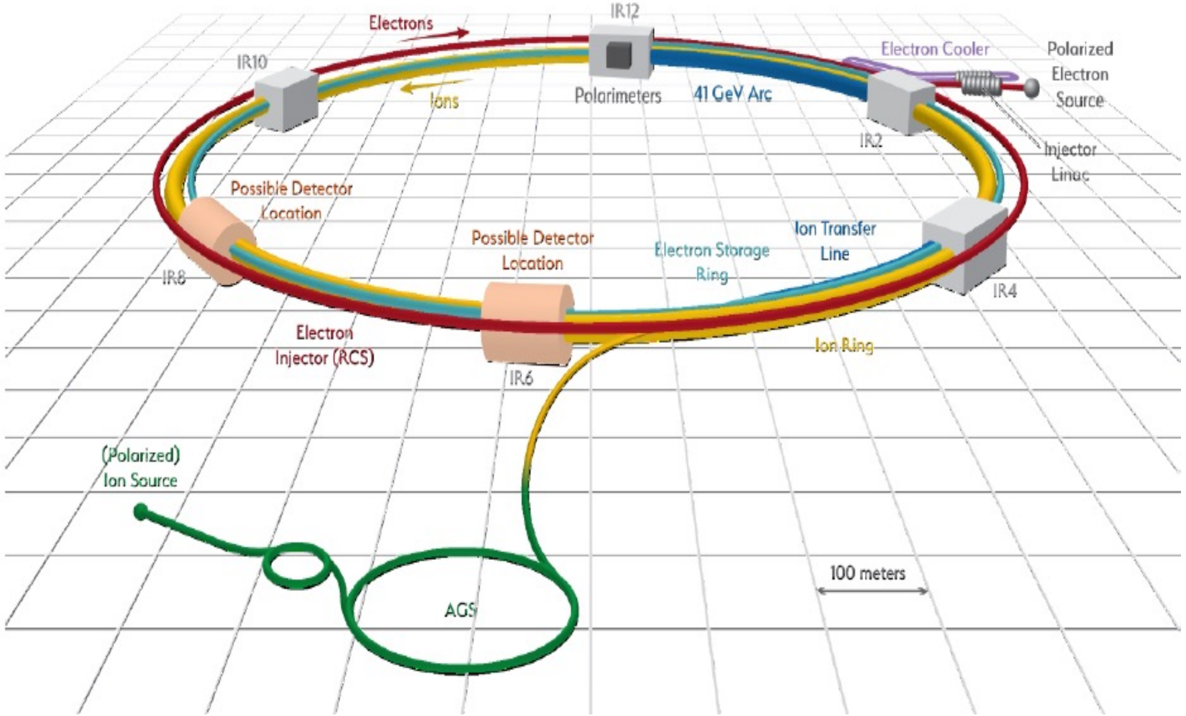








The US Electron Ion Collider



Hadrons up to 275 GeV

- Existing RHIC complex: Storage (Yellow), injectors (source, booster, AGS)
- Need few modifications
- RHIC beam parameters fairly close to those required for EIC@BNL

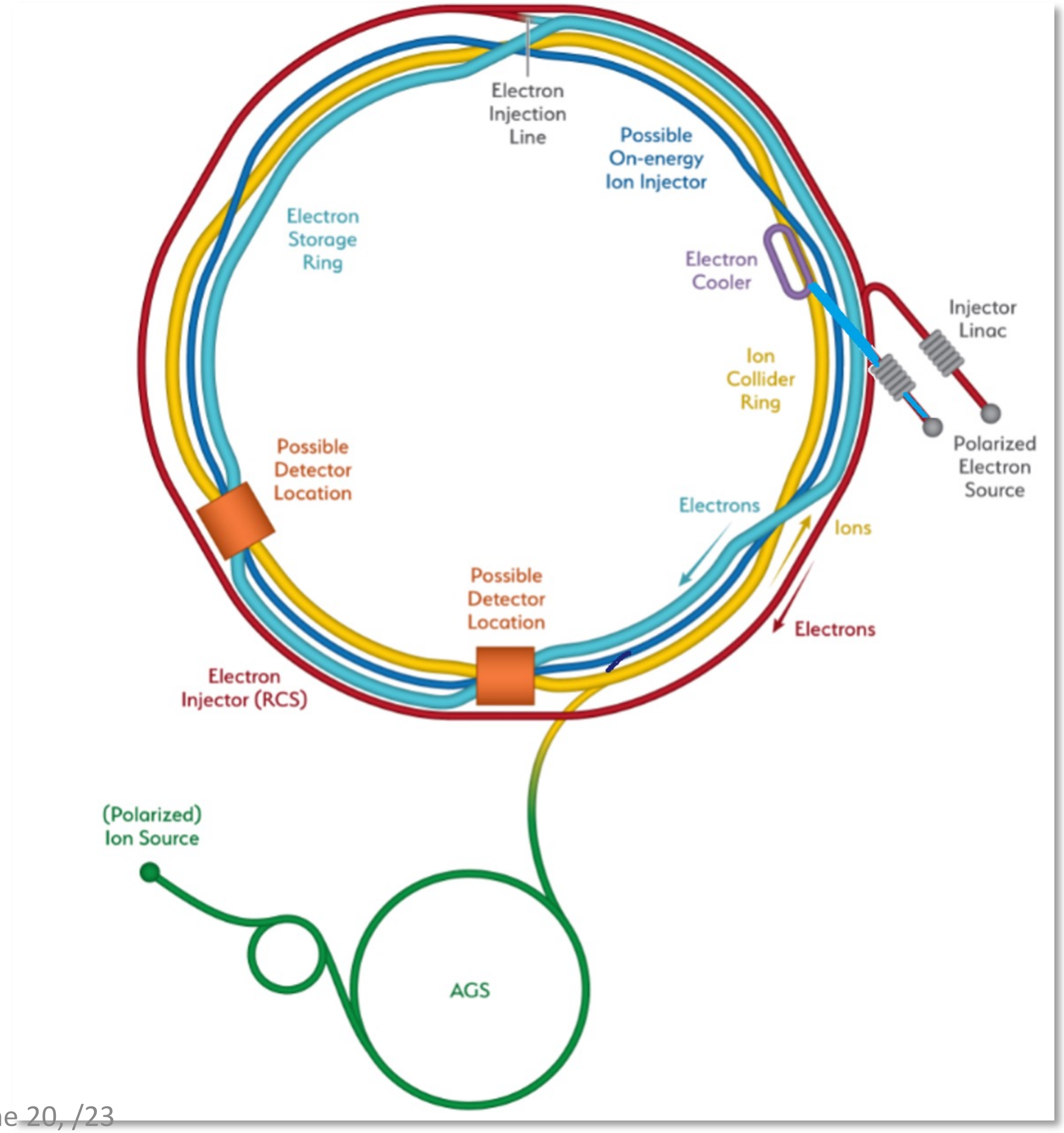
Electrons up to 18 GeV

- Storage ring, provides the range sqrt(s) = 20-140 GeV. Beam current limited by RF power of 10 MW
- Electron beam with variable spin pattern (s) accelerated in on-energy, spin transparent injector (Rapid-Cycling-Synchrotron) with 1-2 Hz cycle frequency
- Polarized e-source and a 400 MeV s-band injector LINAC in the existing tunnel

- Electron storage ring with frequent injection of fresh polarized electron bunches
- Hadron storage ring with strong cooling or frequent injection of hadron bunches

Design optimized to reach 10^{34} $cm^{-2}sec^{-1}$

Reference Detector – Location



Two possible locations – IP6 and IP8 – for detectors and Interaction Regions.

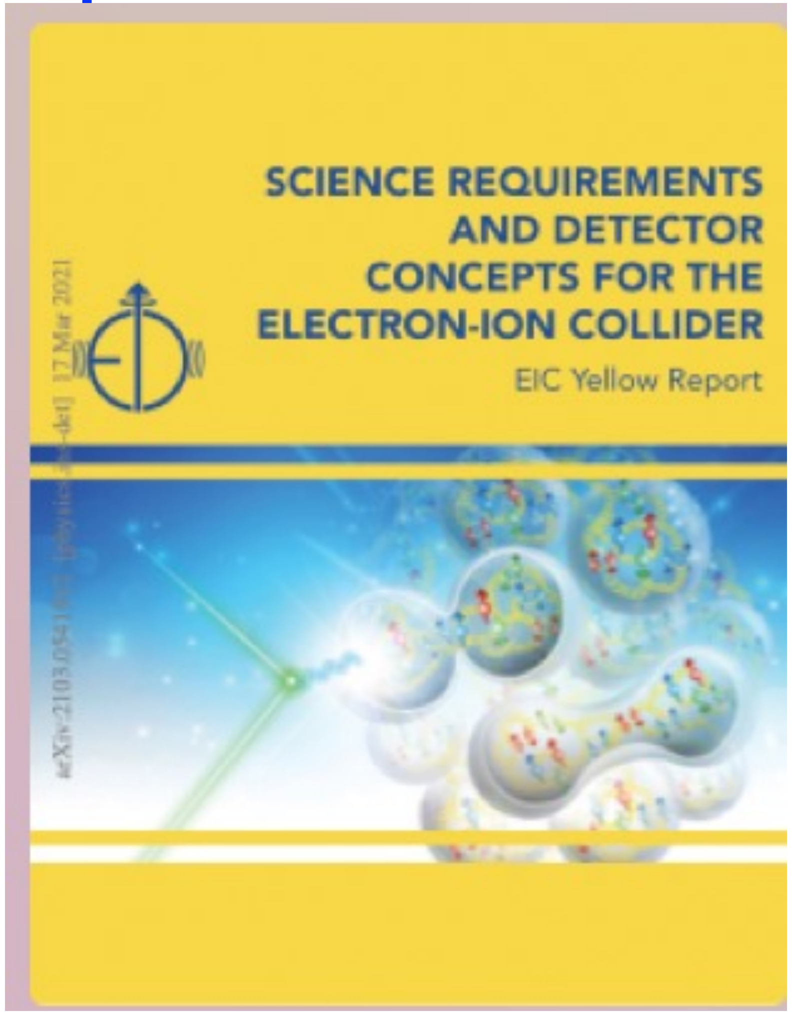
IP6 is the assumed detector location from project risk view (mainly schedule).

- **IP8 is also suitable.**

- Hadron Storage Ring
- Electron Storage Ring
- Electron Injector Synchrotron
- Possible on-energy Hadron injector ring
- Hadron injector complex

December 2019 – March 2021
EICUG Yellow Report

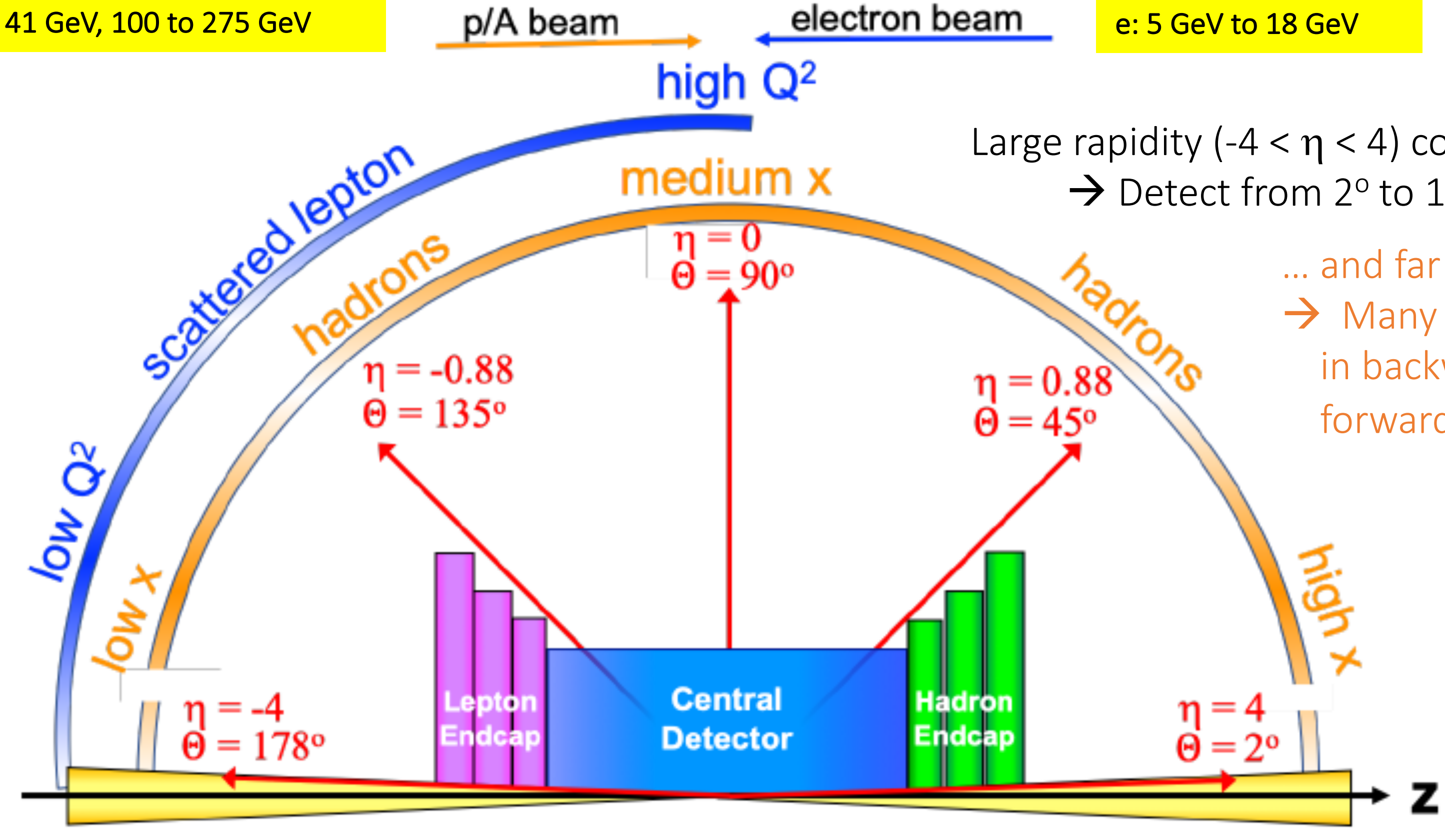
- Led by EICUG Steering Committee, with R. Ent & T. Ullrich as point people for the effort, initiated a UG-wide effort towards a detailed detector design effort with a detailed document.
- Kick off meeting at MIT in December 2019 followed by 4 more meetings in 2020 all remote: Philadelphia, Pavia, Miami, Washington DC, Berkeley



Design by Yulia Furetova and Shannon West

Details in Silvia Dalla Torre's lectures p: 41 GeV, 100 to 275 GeV

e: 5 GeV to 18 GeV



Large rapidity ($-4 < \eta < 4$) coverage
→ Detect from 2° to 178°…

… and far beyond
→ Many detectors in backward + forward region

low Q^2 scattered electrons
Bethe-Heitler photons
for luminosity

Backward

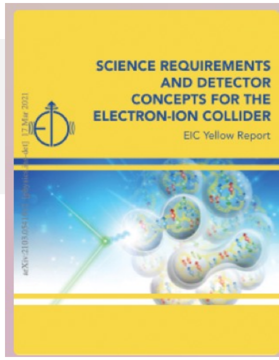
Main detector =
Barrel +
Lepton Endcap +
Hadron Endcap

particles from nuclear breakup, i.e. neutrons, scattered protons, and ions from diffractive reactions

Forward

Endcaps critical because of largely asymmetric beam energies (with different beams).
Similar large impact on IR/accelerator.

Resulting Experimental Requirements



More and more demanding moving from inclusive to fully exclusive scattering

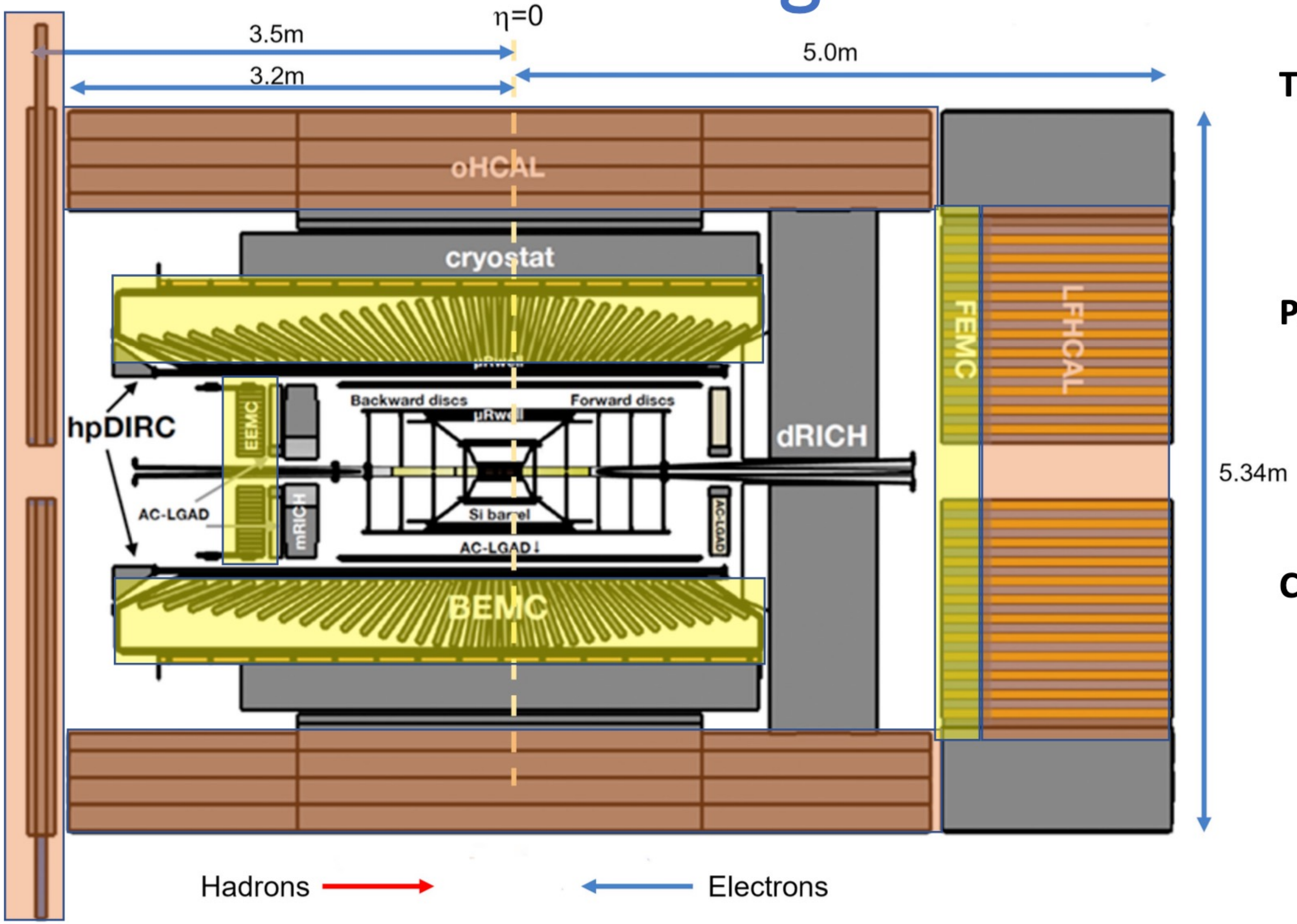
- **Inclusive measurements (DIS), required:**
  - Precise scattered electron identification (e.m. calorimetry, e/h PID) and extremely fine resolution in the measurement of its angle (tracking) and energy (calorimetry)
- **Semi-inclusive measurements (SI-DIS), also required:**
  - excellent hadron identification over a wide momentum and rapidity range (h-PID)
  - full 2π acceptance for tracking (tracking) and momentum analysis (central magnet)
  - excellent vertex resolution (low-mass vertex detector)
- **Exclusive measurements also required:**
  - Tracker with excellent space-point resolution (high resolution vertex) and momentum measurement (tracking),
  - Jet energy measurements (h calorimetry)
  - very forward detectors also to detect n and neutral decay products (Roman pots, large acceptance zero-degree calorimetry)
- **And luminosity control, e and A polarimeters, r-o electronics, DAQ, data handing**

EIC Detector Advisory Committee (DAC) Meeting, 28-29 September 2020

S. Dalla Torre (INFN) & T. Horn (CUA)

More in Silvia Dalla Torre's lesson

ePIC Detector Design



Tracking:

- New 1.7T solenoid
- Si MAPS Tracker
- MPGDs (μRWELL/μMegas)

PID:

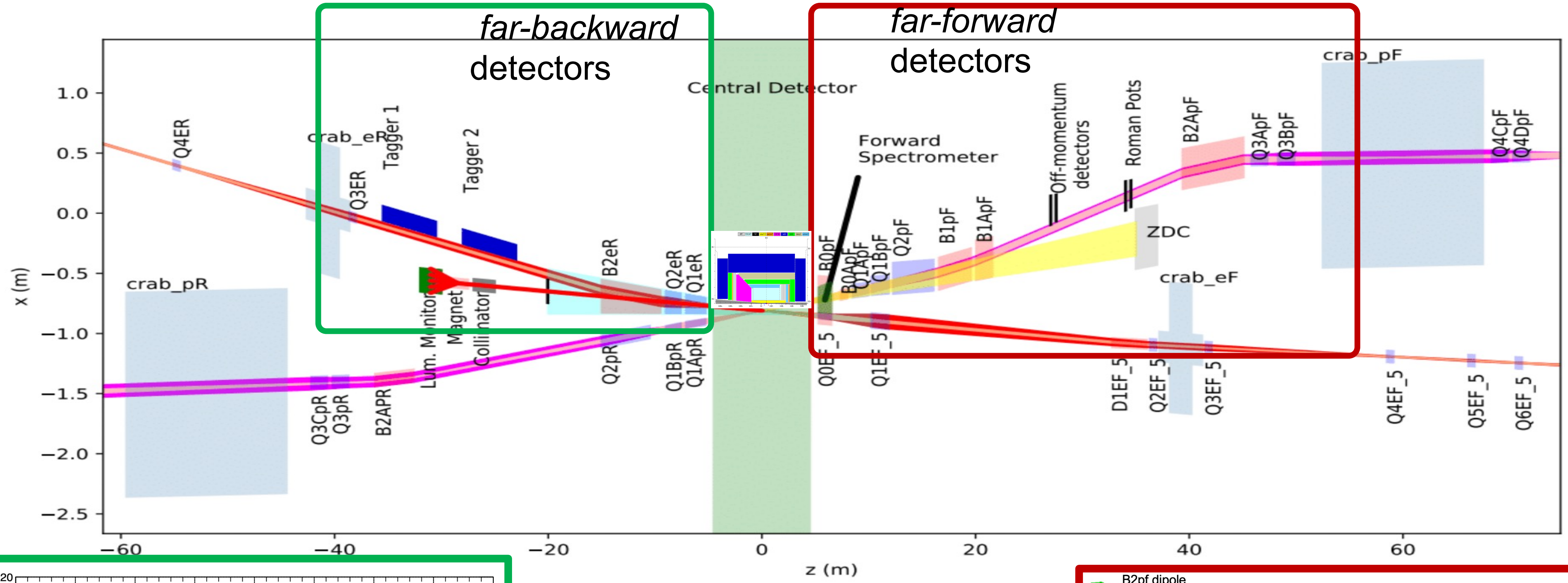
- hpDIRC
- mRICH/pfRICH
- dRICH
- AC-LGAD (~30ps TOF)

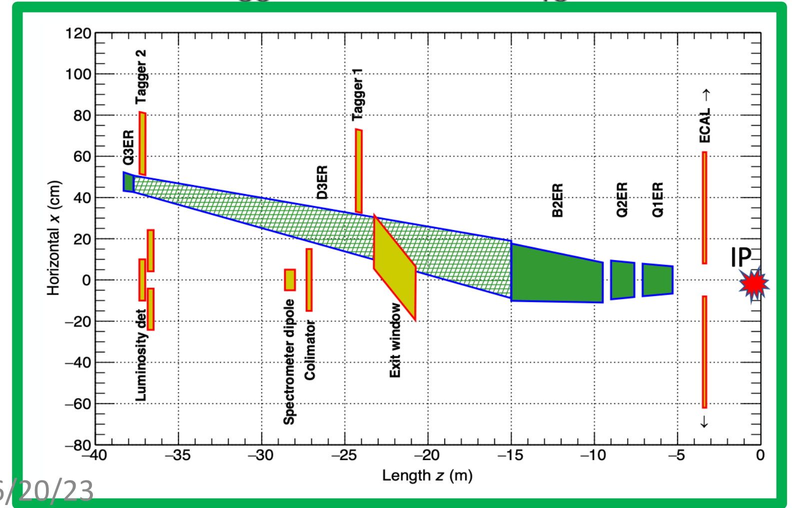
Calorimetry:

- SciGlass/Imaging Barrel EMCal
- PbWO4 EMCal in backward direction
- Finely segmented EMCal +HCal in forward direction
- Outer HCal (sPHENIX re-use)
- Backwards HCal (tail-catcher)

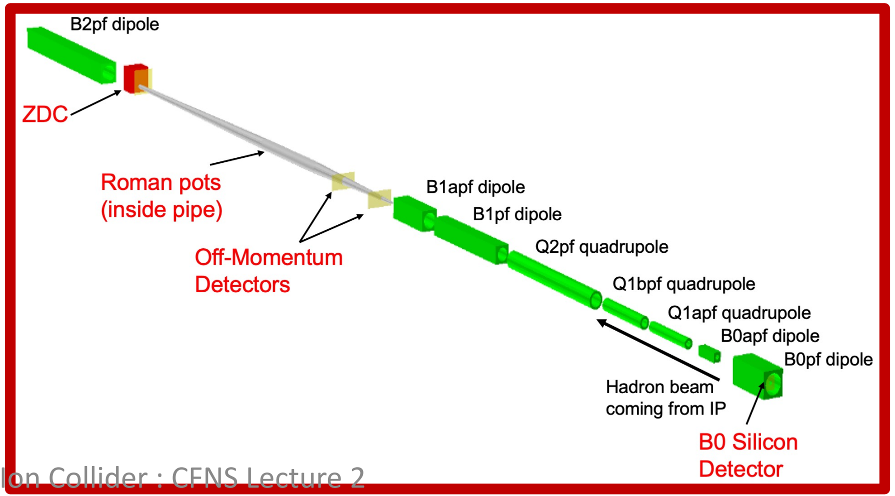
6

Reference Detector – Backward/Forward Detectors





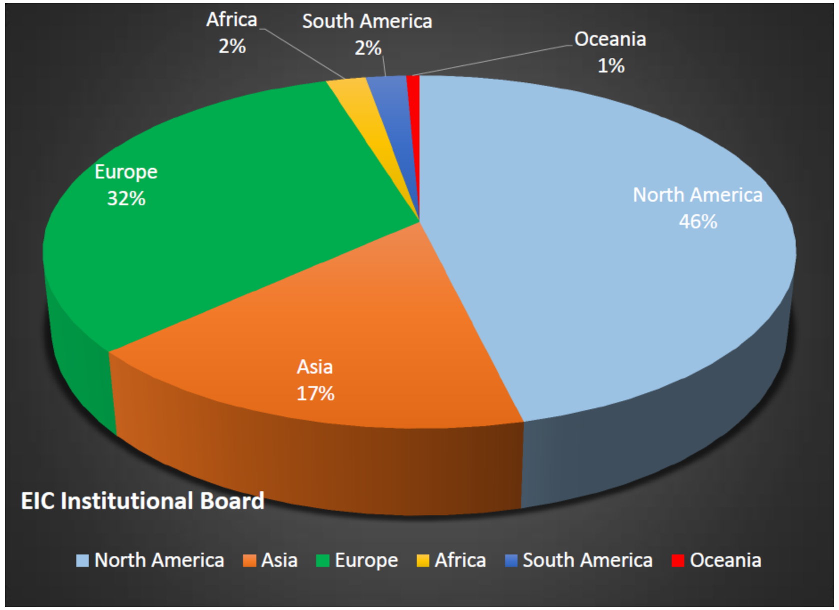
Extensive integration of forward and backward detector elements into the accelerator lattice



The EIC Users Group: EICUG.ORG

Formally established in 2016, now we have:
~1350 Ph.D. Members from 36 countries, 270 institutions
New members welcome

New:
Center for Frontiers in Nuclear Science (at Stony Brook/BNL)
[EIC²](EIC2) at Jefferson Laboratory



EICUG Structures in place and active:
EIC UG Steering Committee, Institutional Board, Speaker's Committee, Election & Nominations Committee

Year long workshops: Yellow Reports for detector design

Annual meetings: Stony Brook (2014), Berkeley (2015), **ANL (2016)**, **Trieste (2017)**, **CAU (2018)**, **Paris (2019)**, **FIU (2020)**, **Remote (2021)**, **Stony Brook (2022)**, **Warsaw (2023)**

Summary & Outlook

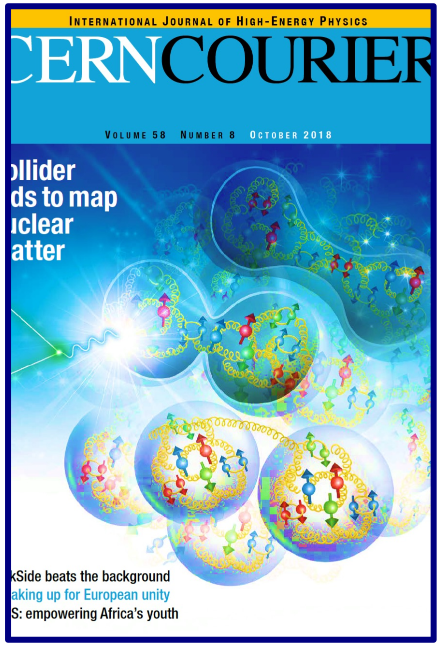
- Electron Ion Collider, a high-energy **high-luminosity polarized e-p, e-A collider**, funded by the DOE will be built in this decade and operate in 2030's.
  - Will address some of the most profound question yet unanswered in the Standard Model of Strong Interactions (and beyond)
- Up to two hermetic full acceptance detectors under consideration, currently **EIC project has funds for 1 detector**, cost of a second detector from non-DOE sources
  - Experimental collaboration formed: ePIC – See Sivlia Dalla Torre's talk(s)
  - EIC project assumes **an aggressive timeline : engineering collisions around 2031/2, physics collisions within 2-years of that.**
- **High interest in having international partners both on detector and accelerator**

- *For all early career scientists, graduate and undergraduate students: This machine is for you! Ample opportunity to contribute to machine, detector & physics of a new project.*

Welcome to the EIC family….

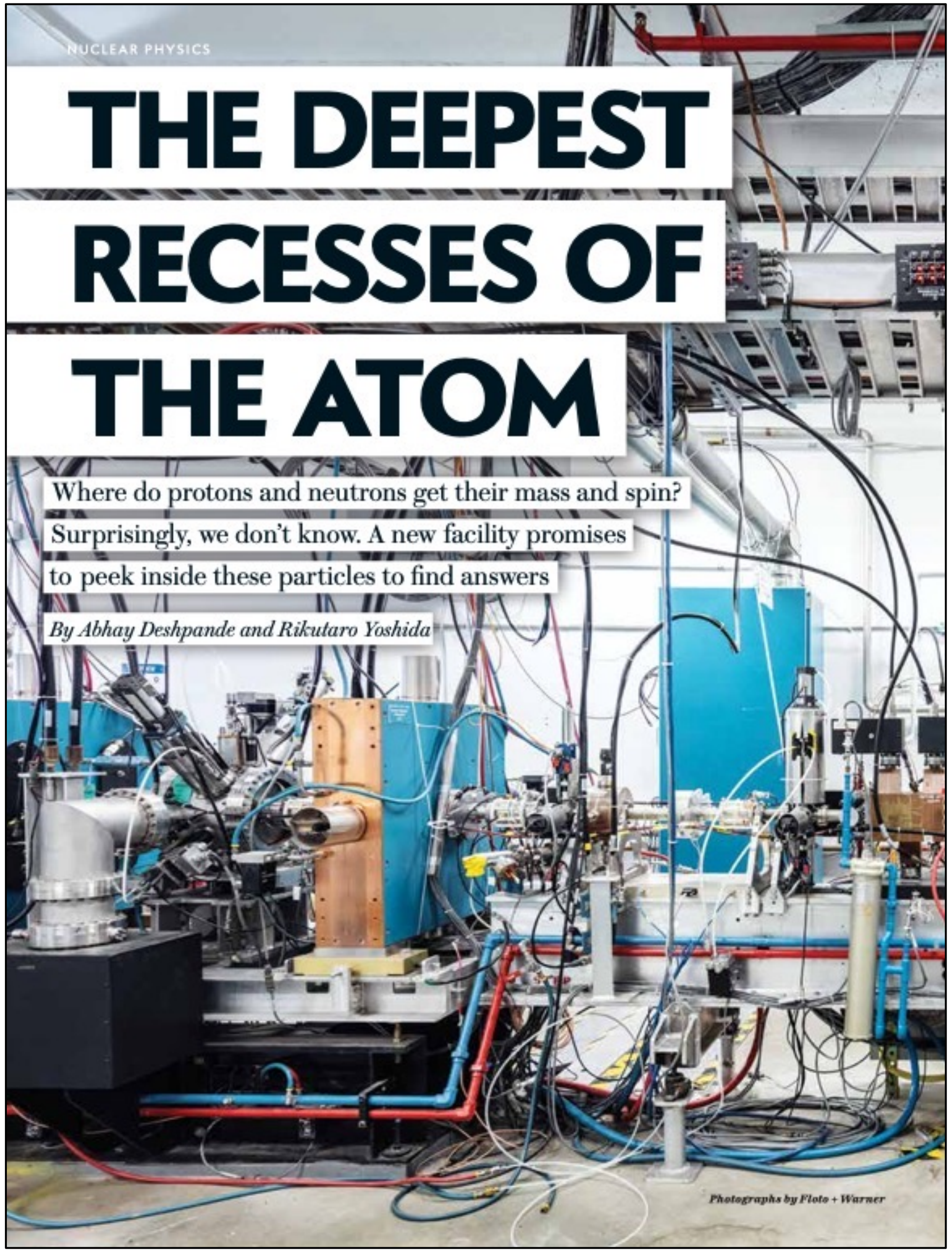

"New directions in science are launched by new tools much more often than by new concepts."
Freeman Dyson

R. Ent, T. Ullrich, R. Venugopalan
Scientific American (2015)
Translated into multiple languages



E. Aschenauer
R. Ent
October 2018

A. Deshpande
& R. Yoshida
June 2019
Translated in to multiple languages



Proton mass puzzle

© Nobel Media AB. Photo: A. Mahmoud

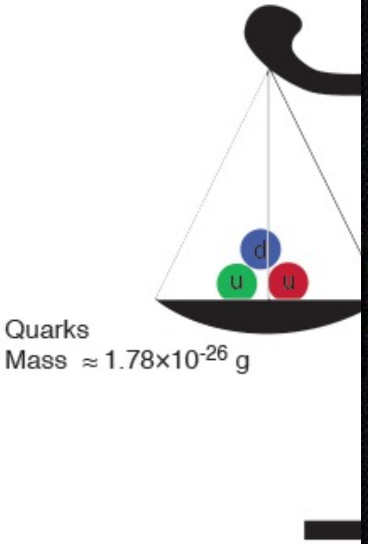
François Englert

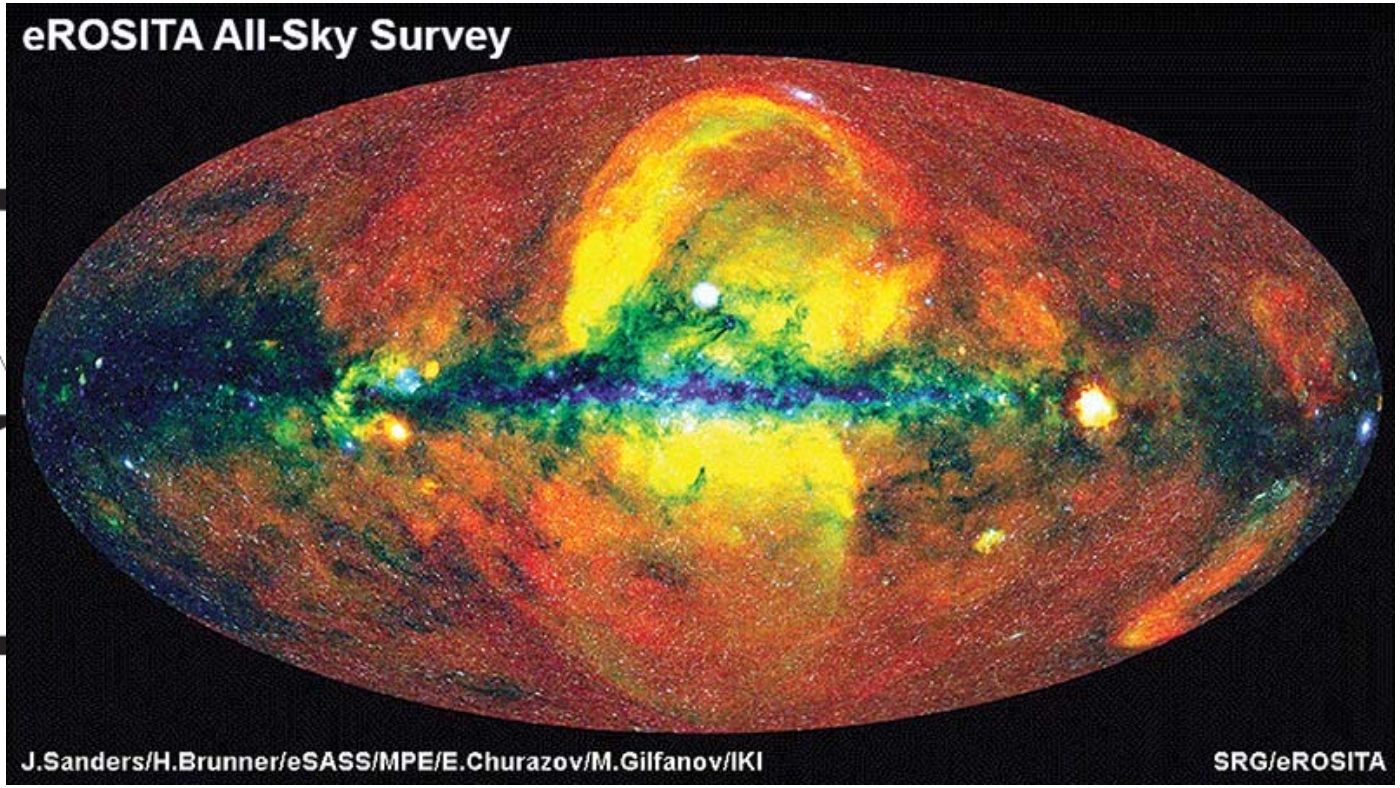
© Nobel Media AB. Photo: A. Mahmoud

Peter W. Higgs

Nobel 2013 With
Francois Englert

"Higgs Boson" that gives mass to quarks, electrons,….





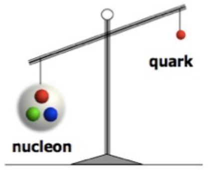
Add the masses of the quarks (HIGGS mechanism) together 1.78×10^{-26} grams

But the proton's mass is 168×10^{-26} grams

→ **only 1% of the mass of the protons (neutrons)** → Hence the Universe

→ Where does the rest of the mass come from?

Mass of the Nucleon (Pion & Kaon)



"The mass is the result of the equilibrium reached through dynamical processes." **X. Ji**

"… The vast majority of the nucleon's mass is due to quantum fluctuations of quark-antiquark pairs, the gluons, and the energy associated with quarks moving around at close to the speed of light. …"
-- The 2015 Long Range Plan for Nuclear Science

Relativistic Motion | ***Chiral Symmetry Breaking*** | ***Quantum Fluctuations***

X. Ji, PRL 74 1071 (1995)

$$M = E_q + E_g + \chi_{m_q} + T_g$$

Quark Energy | **Gluon Energy** | **Quark Mass** | **Trace Anomaly**

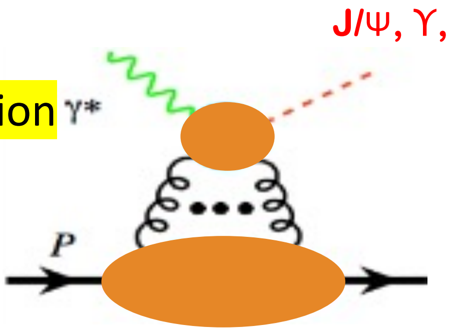
- Criticisms: not scale-invariant, decompositions: Lorentz invariant vs. rest frame
- Recent interest (workshops planned) to clarify how to determine the different contributions
- **Lattice QCD providing estimates**

E_q ~30% E_g ~40% χ_{m_q} ~10% T_g ~ 25%

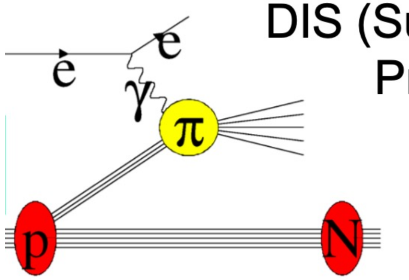
arXiv: 1710.09011

Trace anomaly:
J/Psi & Upsilon production near threshold:

SoLID@JLab & EIC

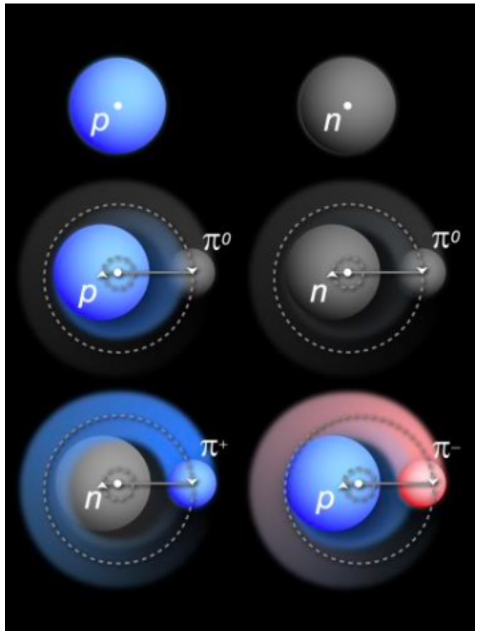


DIS (Sullivan Process)



(pion/Kaon) PDFs: P. C. Barry et al. PRL 127, 232001 (2021)

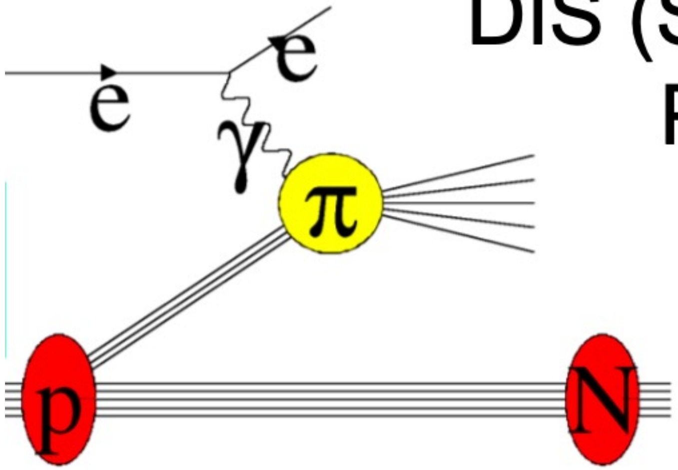
Pion/Kaon mass & PDFs



Relativistic Motion — *Chiral Symmetry Breaking* — *Quantum Fluctuations*

$$M = E_q + E_g + \chi_{m_q} + T_g$$

Quark Energy — Gluon Energy — Quark Mass — Trace Anomaly



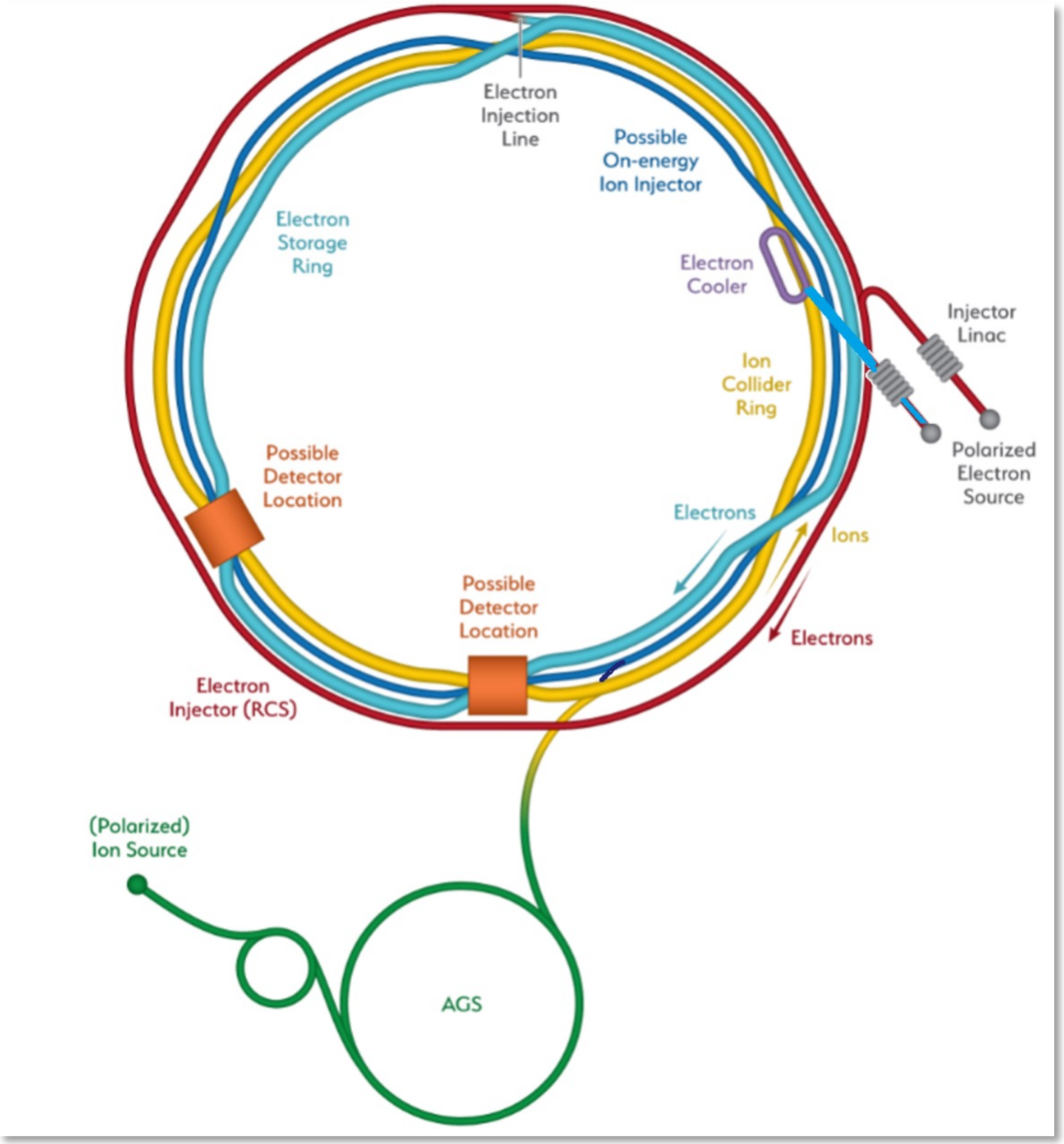
DIS (Sullivan Process)

For PDF studies

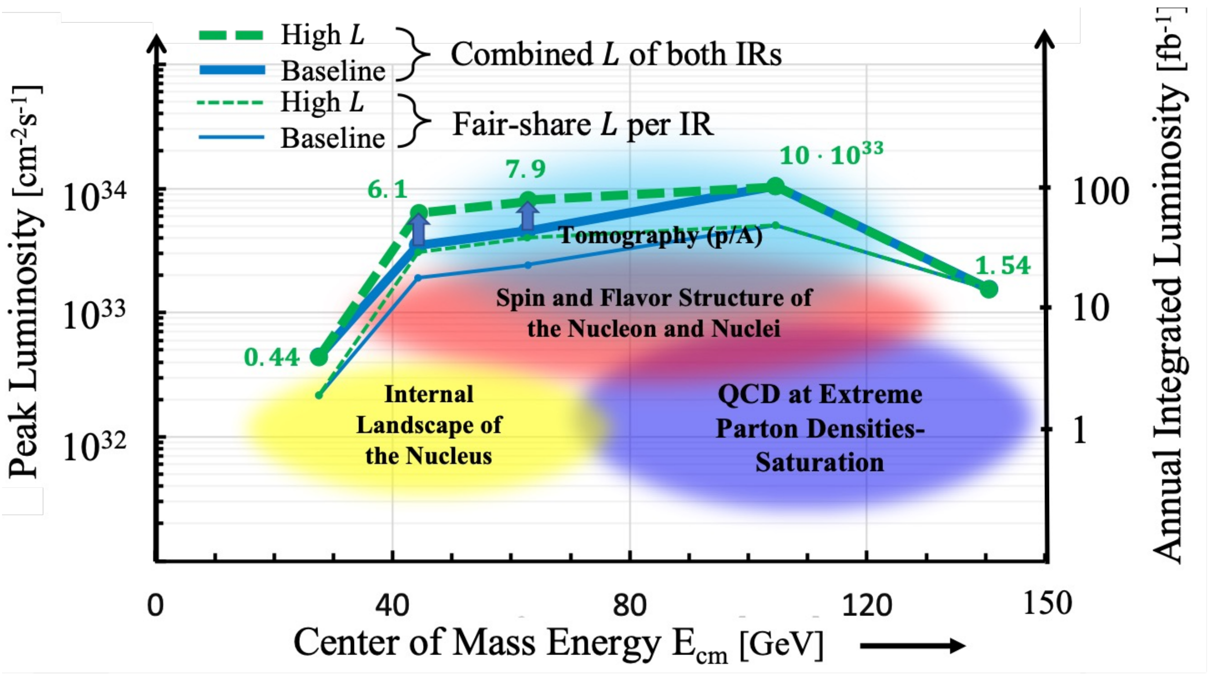
- How different are these terms in 2-quark systems? Light vs. heavy quarks?
- What can we learn from Sullivan Process about their structure?
- Hints for learning about origin of emergent mass?

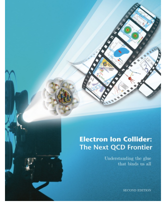
PDFs @ EIC : J> Arrington et al. J. of Physics G. 48 (2021) 075106
PDFs: P. C. Barry et al. PRL 127, 232001 (2021)

EIC Accelerator Design



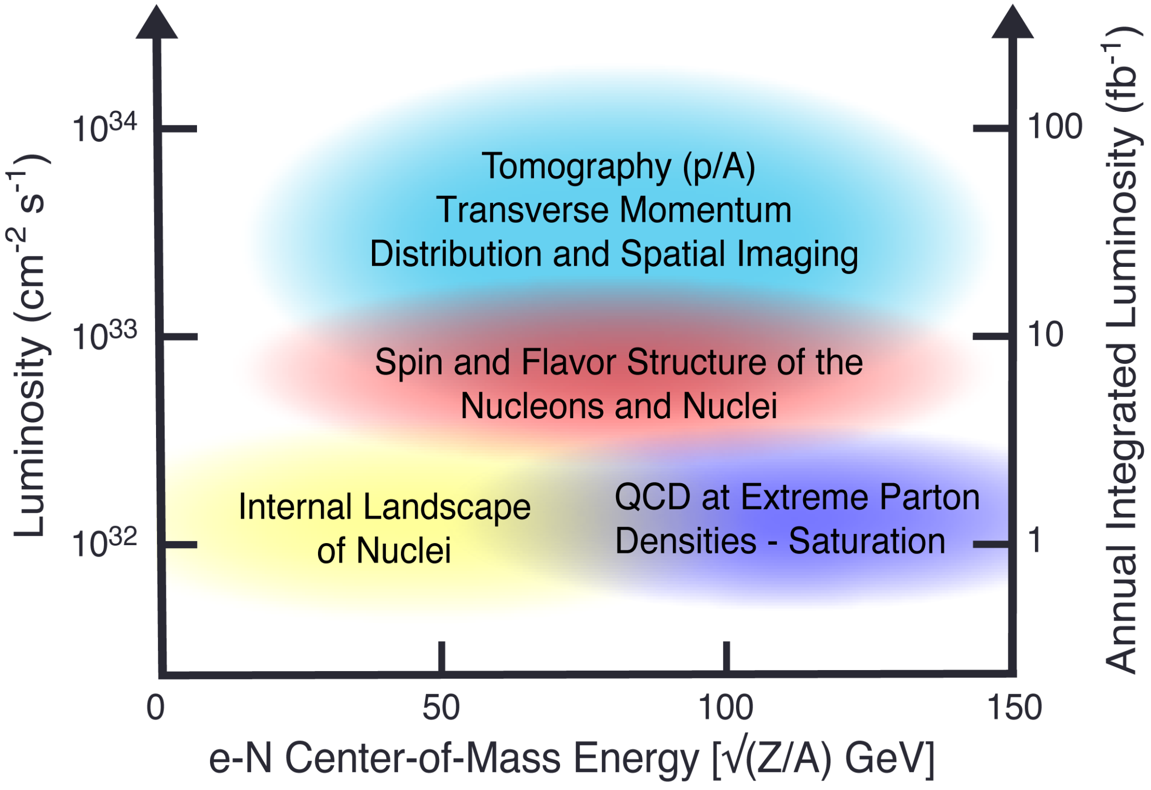
Center of Mass Energies:	20GeV - 140GeV
Luminosity:	10^{33} - 10^{34} $cm^{-2}s^{-1}$ / 10-100fb^{-1} / year
Highly Polarized Beams:	70%
Large Ion Species Range:	p to U
Number of Interaction Regions:	Up to 2!

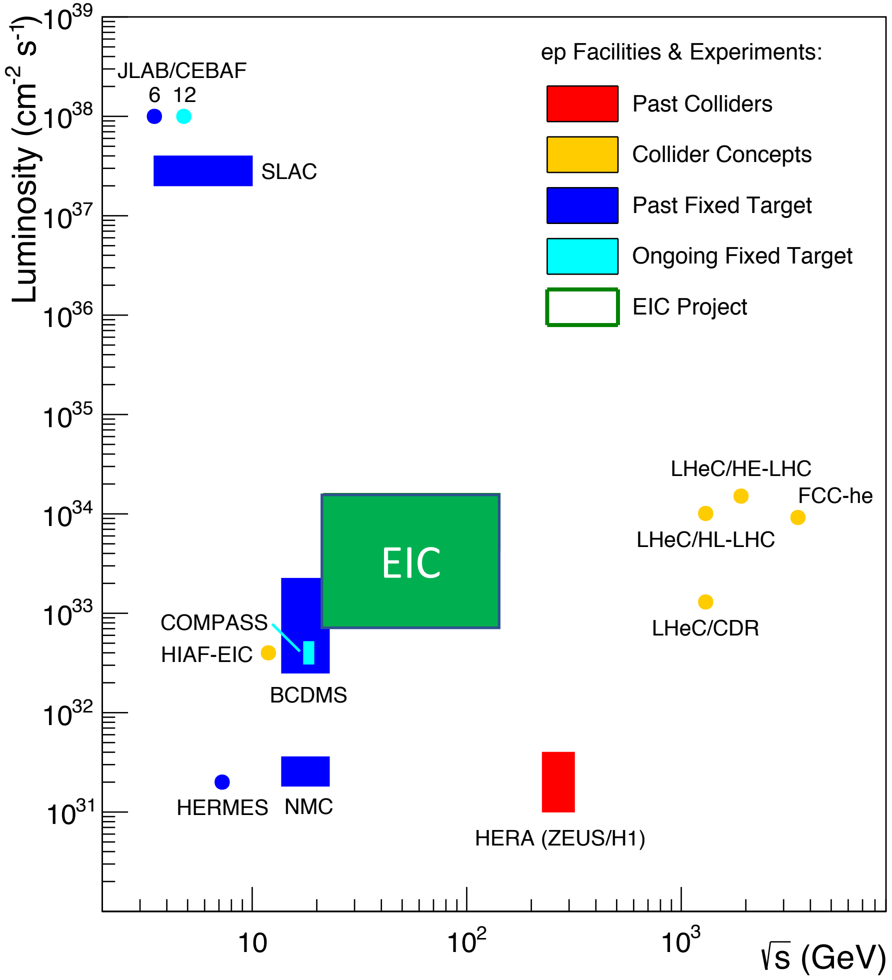




EIC Physics and the machine parameters

CM vs. Luminosity vs. Integrated luminosity



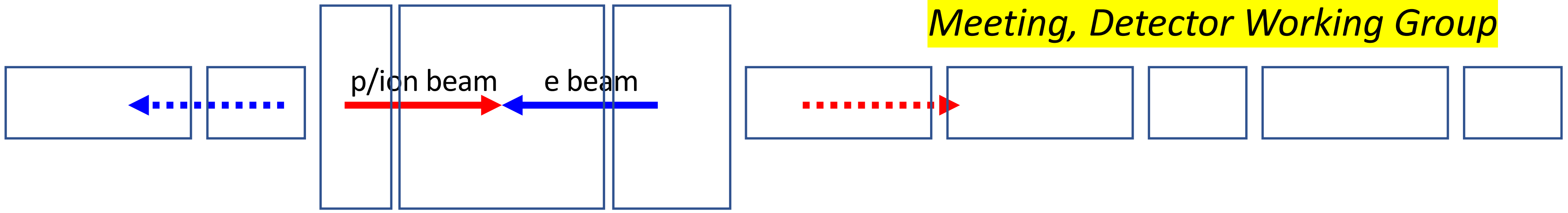


The US EIC with a wide range in $\sqrt{s}$, polarized electron, proton and light nuclear beams and luminosity makes it a unique machine in the world.

Cartoon/Model of the Extended Detector and IR

- ❑ EIC physics covers the entire region (backward, central, forward)
- ❑ Many EIC science processes rely on excellent and fully integrated forward detection scheme

Adapted from 2nd Yellow Report Meeting, Detector Working Group



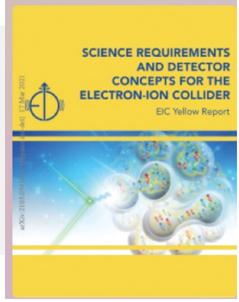
far-backward e-detection	e final-focus quads	"Central detector", includes e-endcap, central, and p/ion endcap detectors	forward dipole incl. h detection	Ion final-focus quads	forward dipole	far-forward h-detection	far-forward h-detection
Low-Q^2 spectroscopy		Inclusive Structure Functions, TMDs, heavy flavors and jets, electrons for GPDs	GPDs/DVCS, tagging, diffraction, high-medium t			Baryon decay π/K structure evaporated n	GPDs, tagging, diffraction, lowest-t
GEMs Diamond detectors?		Vertex and Tracking detectors, particle identification detectors, calorimetry detectors, muon detectors, etc.	Si/GEMs Roman pots, e/γ calorim.			GEMs Roman pots e/γ calorim.	Roman pots ZDCs

physics examples

detector examples

Concept DETECTOR

This detector concept was included in the EIC CDR prepared for the CD1 Review



Cherenkov

-1.00

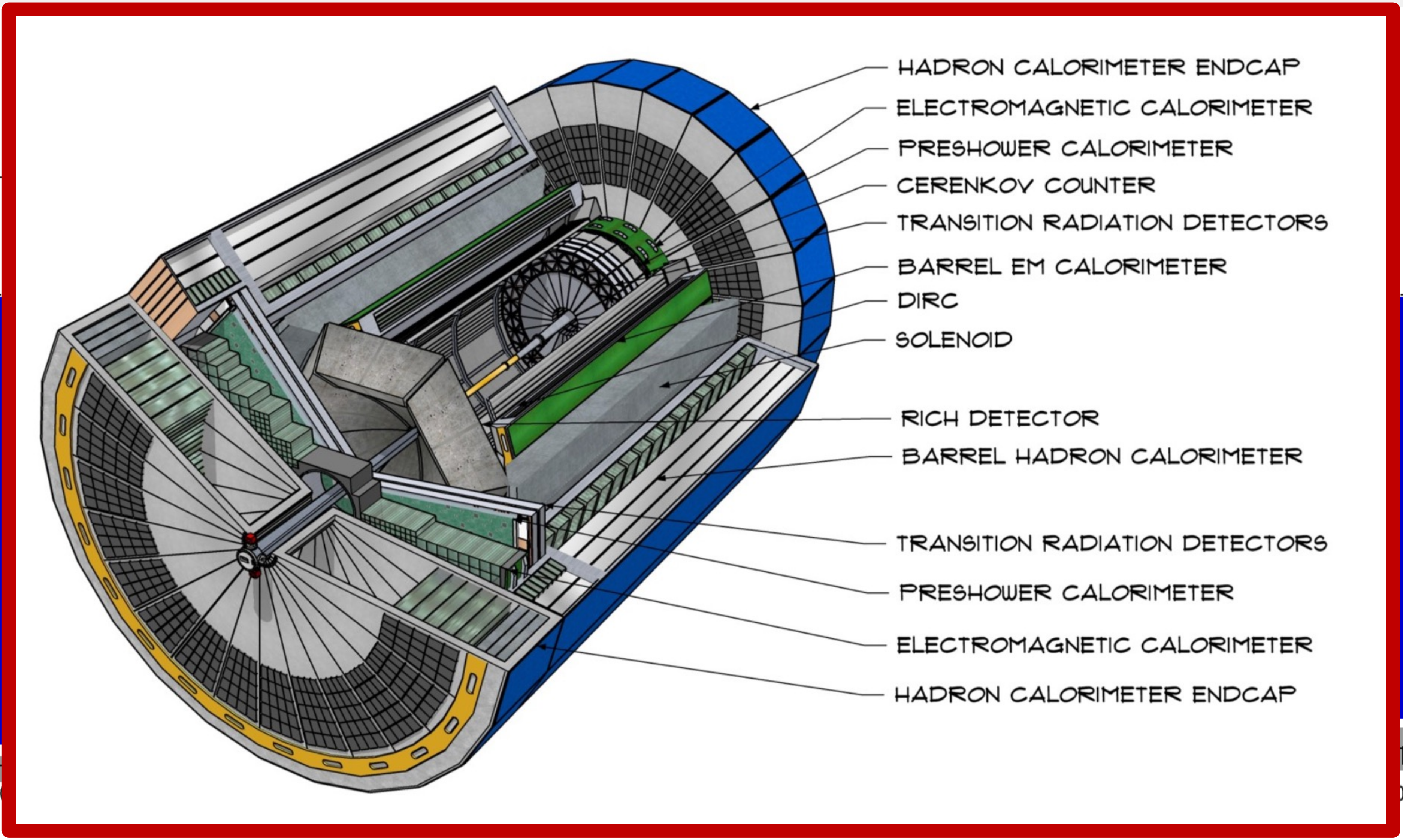
-4.20

-40

1.20

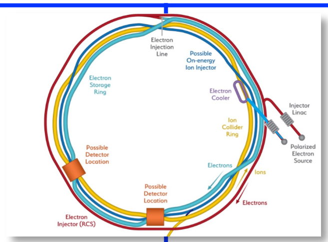
4.20

00.0



Complementarity for 1^{st}-IR & 2^{nd}-IR

	1^{st} IR (IP-6)	2^{nd} IR (IP-8)
Geometry:	ring inside to outside	ring outside to inside
	tunnel and assembly hall are larger Tunnel: ∅ 7m +/- 140m	tunnel and assembly hall are smaller Tunnel: ∅ 6.3m to 60m then 5.3m
Crossing Angle:	25 mrad	35 mrad secondary focus



different blind spots
different forward detectors and acceptances
different acceptance of central detector

Luminosity:

more luminosity at lower E_{CM}
optimize Doublet focusing FDD vs. FDF
→ impact of far forward p_T acceptance

Experiment:

1.5 Tesla or 3 Tesla
different subdetector technologies

ECA & RE