

The case for ions: the physics of nuclear PDFs and hadronization studies

Lecture 1

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Outline

- Why use a nucleus as a target/beam?
- How to deal with nuclei.
- Data used in nPDF fits.
- Sets (until 2015) and comparison.
- 2015-present: improving.
- Issues in nPDF extractions.
- What can the EIC do for nPDFs?
- Summary.

Why use a nucleus a a target/beam?

$$\frac{d^2\sigma}{dxdQ^2} \propto F_2(x,Q^2) - \frac{y^2}{1 + (1-y)^2} F_L(x,Q^2)$$

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$$f_{proton}^{S}(x, Q^2) = \sum_{i=1}^{N_f} f_i(x, Q^2)$$

$$f_{proton}^{NS}(x,Q^2) = \sum_{i=1}^{N_f} e_i^2 f_i(x,Q^2)$$

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We have a minimum of seven PDFs to determine: (anti-)up, (anti-)down, (anti-)strange, and gluon.

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We can have I.i. combinations of PDFs if we use a neutron instead. Assuming isospin symmetry (i.e. $u_{proton} = d_{neutron}$) is valid

$$f_{proton}^{S}(x,Q^{2}) = f_{neutron}^{S}(x,Q^{2})$$

$$f_{proton}^{NS}(x,Q^2) \neq f_{neutron}^{NS}(x,Q^2)$$

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But we could use **deuterium**. A (very light) nucleus.

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In that case we have one more structure function, xF_3 , which enters the cross-section with different signs (depending on the sign of the W boson). We can have I.i. combinations of PDFs if we use CC DIS (exchange of a W boson).

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But we could use **larger nuclei** as targets.

 The expectation was that using a nucleus would be similar to using a proton.

 At most one would need to account for the non-isoscalar nuclei
and the internal motion
of nucleons in the nucleus (Fermi motion).





actual nuclear modification

(only statistical uncertainty shown)



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"The result is **in complete disagreement with the calculations** illustrated in Fig. 1." "The result is **in complete disagreement with the calculations** illustrated in Fig. 1."

"We are not aware of any published detailed prediction presently available which can explain the behaviour of these data. However, there are several effects known and discussed which can change the quarks distributions in a high A nucleus compared to the free nucleon case and can contribute to the observed effect."

PLB 202, 603



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The effects follow a very particular pattern:





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Moreover, we want to study nuclei also because:

- It is interesting and we want to know how and why this occurs.
- Tightly linked to neutrino physics.
- HI and the QGP benchmarking.
- Saturation studies (e.g. CGC).

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DONUT Kamioka Observatory NA61 Super-Kamiokano	ACE	Enriched Xenon	Large Volume Detector	NESTOR Project
	AMANDA	Observatory	LAND	NEVOD
	ANTARES	EMC	LHCb	Kolar Gold Fields
	ArgoNeuT	FASER	MINOS	PHENIX
	ATLAS	Fermilab E-906/SeaQuest	Modular Neutron Array	PUMA
	Bevatron	Gargamelle	Monopole, Astrophysics	Rutherford gold foil
	Borexino	Germanium Detector Array	and Cosmic Ray	experiment
	Bubble Chamber	HARP	Observatory	SAGE
	CDHS	HERA-B	Mu to E Gamma	SciBooNE
	CLAS detector	HERMES	Mu2e	SNO+
	CMS	IceCube	Mu3e	Soudan 1
	COMPASS (NA58)	Irvine-Michigan-	NA32	Soudan 2
	Cowan-Reines experiment	Brookhaven	NA35	STAR
	CUORE	Kamioka Liquid Scintillator	NA49	Sudbury Neutrino
	DAPHNE	Antineutrino Detector	NA60	Observatory
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	DONUT	Kamioka Observatory	NA61	Super-Kamiokande

How to deal with nuclei

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• Do nothing.



- Build theoretical models:
 - shadowing ~ 400
 - anti-shadowing ~ 40
 - EMC effect ~ 370
 - Fermi motion ~ 90



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 A purely phenomenological approach, assuming all we do for the proton will remain valid: nuclear PDFs.



Build theoretical models:

- shadowing ~ 400
- anti-shadowing ~ 40
- EMC effect ~ 370
 - Fermi motion ~ 90

- A purely phenomenological approach, assuming all we do for the proton will remain valid: nuclear PDFs.
- We must test this idea.

Just like proton PDFs, these nPDFs should:

- be **universal**.
- **evolve** with the scale with DGLAP.
- be non-perturbative (we need data to determine them).
- contain all the information about the behaviour of partons in a nucleus.

Remember that we're talking about the collinear framework.

We can always study more involved observables in the nuclear environment.

Can we describe the behaviour with the average of protons and neutrons?

$$f_{i/A}(x,Q^2) = \frac{Z}{A} f_{i/p}(x,Q^2) + \frac{(A-Z)}{A} f_{i/n}(x,Q^2)$$

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We see is a **genuine modification of the initial state** due to the medium. The naive attempt fails. Miserably.

We really have to do something else.

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Options used so far:

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$$f_{i/p/A}(x, Q_0^2, A) = f_{i/p}(x, Q_0^2) R_i(x, A)$$



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 but extended

$$xf_{i/p}(x, Q_0^2) = c_0 x^{c_1}(1-x)^{c_2} e^{c_3 x}(1+e^{c_4 x})^{c_5}$$

$$c_k \to c_{k,0} + c_{k,1}(1 - A^{-c_{k,2}})$$
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$$f_{i/p/A}(x, Q_0^2, A) \propto x^{\alpha}(1-x)^{\beta} NN$$

to control the edges of the kin. space

very flexible, no need to assume the A dependence

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Now, do a **global** fit.

Data used in nPDF fits

- NC DIS: since ever.
- Drell-Yan with fixed target: since ever.
- Single inclusive hadron production: since 2009.
- CC DIS: since 2012.
- W and Z production at the LHC: since 2016.
- di-jets at the LHC: since 2016.
- D meson production at the LHC: since 2021.
- Prompt photon production at the LHC: since 2021.

Sets (until 2015) and comparison



When doing a fit one has to make choices, and these impact the outcome:

- initial scale for the evolution
- order of the perturbation theory
- heavy-flavour scheme
- data fitted
- parametrisation
- proton baseline
- ...



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- ...

It is customary to show the nPDFs as ratios. We distinguish valence ($f_{i,valence} = f_i - f_{\bar{i}}$), sea and gluon modifications.

The old days: LO

- nDS used a convolutional approach, the rest a multiplicative factor.
- No flavour separation: only valence, sea and gluon.



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JHEP 0904 (2009) 065













- This is a very involved observable (in terms of parton distributions).
- RHIC data are not very precise.
- EPS09 and DSSZ used a totally different
 treatment.

PRD 85, 074028

In general, up to 2012:

 The valence distributions are *well* constrained in the region where we have data.

 Idem for the non-strange sea densities, but with ad-hoc constraints. They don't look that bad.

• The strange quark lives up to its name.

The gluon density is a BLOOPY MESS.

2015-present: improving

To improve (n)PDFs we can:

- use the same observables with higher precision and/or broader kinematic coverage: NC DIS, Drell-Yan, single inclusive hadron production.
- use other observables: CC DIS, W/Z production, dijets, D meson production.
- improve the theoretical computation.

• CC DIS

• $R_{u_v} \neq R_{d_v}$ requieres data that can better distinguish up from down. 1.61.4



Clearly the green band looks very different. They didn't use CC
 DIS data.

• CC DIS

 Basically 4 experiments: CCFR, NuTeV, CDHSW, and Chorus.

- There are tensions among the different neutrino experiments if we try to fit the cross-sections.
- No problem to accommodate the structure functions in a global fit, nor with Chorus crosssections.

Could NOMAD solve this?





Orell-Yan

Scarce amount of data in older fits: only 92 points, given as ratios to p + d and p + Be.



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Since 2016, "new" Drell-Yan data (1981, 1987, 1989) has been to be considered.

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Since 2016, "new" Drell-Yan data (1981, 1987, 1989) has been to be considered. 28 points from $\pi^{\pm}+W$, and $\pi^{-}+Pt$. Requires π PDFs.

Single inclusive hadron production

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Single inclusive hadron production



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contribution of each initial state parton in p+Pb collisions at the LHC



W and Z boson production

Is there a problem with the normalisation?





Norm 0: no normalisation parameter.

Norm 2: normalisation parameters for CMS and ATLAS Run I.

Norm 3: normalisation parameters for CMS and ATLAS Run I, and CMS Run II.









better theoretical computation



- From NLO to NNLO (for some observables).
- Consistent treatment of heavy quarks, with correct mass schemes.
- Inclusion of nuclear effects in deuterium and other light nuclei.
- Deep learning to reduce the parametrisation bias.
Available sets*

- and March RD 93, 085037. nCTEQ15WZ: EPJC 80, 968.

 nCTEQ15HiX: PRD 103, 114015. nCTEQ15WZ+SIH: PRD 104, 094005.
- and [™]-[™]: nDS: prd 69, 074028. DSSZ: prd 85, 074028.
- **F**-**F**: **nTuJu19**: **prd** 100, 096015. **nTuJu21**: **prd** 105, 094031.
- EKS: EPJC 9, 61. EPSO9: JHEP 0904, 065. EPPS16: EPJC 77, 163.
 EPPS21: EPJC 82, 413.
- **IFENTIAL STATE OF STATE OF**
- Image: Image:
- NN: nNNPDF1.0: EPJC 79, 471. nNNPDF2.0: JHEP 09, 183.

 nNNPDF3.0: EPJC 82, 507.

* not all

Issues in nPDF extractions

• The kinematic coverage of the data.



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NC DIS data	Fixed target (FT)	FT deuterium	Collider	
proton PDF fit e.g. EPJC 81, 341	433	513	1264	

one can use these to do a proton PDF fit: HERAPDF

• The quantity of data.

NC DIS data	Fixed target (FT)	FT deuterium	Collider
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nuclear case	2309	812	0

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• How the data were/are published

• ~ 60% of the NC DIS data are $F_2, F_L, R = F_L/F_2$:

 $\frac{d^2 \sigma^{NC}}{dx dQ^2} \propto F_2 - \frac{y^2}{Y_+} F_L \equiv \sigma_{\rm f}^{NC}$

computable/what the measurement can be turned into • ~ 60% of the NC DIS data are $F_2, F_L, R = F_L/F_2$:

$$\frac{d^2 \sigma^{NC}}{dx dQ^2} \propto F_2 - \frac{y^2}{Y_+} F_L \equiv \sigma_{\Gamma}^{NC}$$

computable/must be extracted from the measurement

$$\frac{d^2 \sigma^{NC}}{dx dQ^2} \propto F_2 - \frac{y^2}{Y_+} F_L \equiv \sigma_{\Gamma}^{NC}$$

• To extract the structure functions one plots σ_r as a function of y^2/Y_+ , the slope is $-F_L$. To vary y^2 , we measure for different \sqrt{s} .



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 If you go through the literature, you will find that *F*₂ has been determined from either a parametrisation of R, R=0 or R=0.2. Take your pick. Some of the new data available has been published in a way that can't be computed in pQCD for proper comparisons. We must rely on theoretical models or nag the experimentalists.
 E.g.

 $\frac{d^2 N_{pPb}/dy dp_T}{\langle T_{pPb} \rangle d^2 \sigma_{nn}^{INEL}/dy dp_T}$ $R_{pPb} =$

published

 $\frac{1}{N_{ev}}\frac{d^2N}{dydp_T}$ $d^2\sigma$ σ^{INEL} dydp_T can be missing part calculated

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 E.g.

 $R_{pPb} = \frac{d^2 N_{pPb} / dy dp_T}{\langle T_{pPb} \rangle d^2 \sigma_{pp}^{INEL} / dv dp_T}$ $\frac{1}{N_{ev}}\frac{d^2N}{dydp_T}$ $d^2\sigma$ σ^{INEL} dydp₇ can be published missing part calculated

Some of the data that we're using might suffer from a medium modification of the final state: SIH, D0 production, di-jets in p+Pb collisions. • From the phenomenological side we have "contamination":





We should be extremely careful not to double, triple, quadruple count effects.



What can the EIC do for nPDFs?

• Improve the kinematic coverage of NC DIS: *nuclear HERA*.



"This broad kinematic coverage ... will revolutionize our current understanding of partonic distributions in nuclei."

Reduce the data uncertainty (at least according to simulations)



18x110 e-A N.C. Uncertainties

- Provide data for more nuclei: from deuterium to ²³⁸U.
- We will be able to use the **EIC data as basis for nPDF fits**.

 We will also be able to properly separate the longitudinal structure function that is sensitive to the gluon density.



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NPA 1026, 122447

 Using the PID needed for SIDIS, we can identify e.g. kaons coming from c quarks.

For the first time we will have jets and di-jets in e+A!

Phys.Rev.C 101 (2020) 6, 065204

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 The presence of a nuclear medium affects (non-trivially) the measured observables in high energy physics.

The differences proton/nuclei can be explained by interacting mechanisms between the partons in bound nucleons, resulting in the need for medium-modified or nuclear PDFs.

While one can propose theoretical models for the nPDFs, using the factorised framework of pQCD we can find model independent distributions, just as in the proton case.

- There are available several sets of nPDFs, and they all provide very good description of the data considered.
- Despite all the effort, nPDFs are very much behind proton PDFs. Mostly due to the amount and limited kinematic coverage of the data.

 While waiting for *clean* data, fitting groups have turned to more involved observables.

 These are sensitive to kinematic regions unreachable otherwise (for now) and/or to poorly constrained densities.

nPDF fitters waiting for the EIC (LHeC? FCC-eh?) data