

Back-to-back hadrons with unpol. and longitudinal pol. targets

Harut Avakian (JLab)

Workshop on kaons and TFR, LNF, Dec 15, 2022

Introduction

- Locating the longitudinally polarized quarks in unpolarized protons
- Separating the kinematics of current and target fragmentation

Test results with Longitudinally polarized protons (CLAS12 RGC)

– $e p \rightarrow e' p X$

– $e p \rightarrow e' p \pi + X$

Summary

Hadron production in hard scattering

X. Artru & Z. Belghobsi

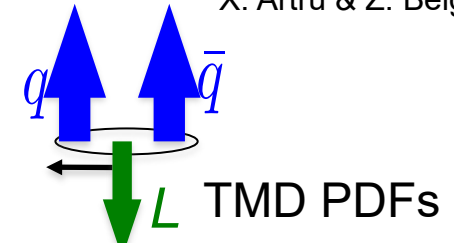
x_F – fractional momentum in the CM frame

polarized quark

$x_F > 0$ (current fragmentation)

$x_F < 0$ (target fragmentation)

FSI



TMD PDFs

Fracture Functions

		Quark polarization		
		U	L	T
Nucleon Polarization	U	\hat{u}_1^{\uparrow}	$\frac{\mathbf{k}_T \times \mathbf{P}_T}{m_N m_b} \hat{u}_1^{\uparrow b}$	$\frac{e_f^2 P_T^{\uparrow}}{m_b} \hat{u}_1^{\uparrow b} + \frac{e_f^2 k_T^{\uparrow}}{m_N} \hat{u}_1^{\uparrow}$
	L	$\frac{S_L(\mathbf{k}_T \times \mathbf{P}_T)}{m_N m_b} \hat{u}_1^{\uparrow b}$	$S_L \hat{u}_1^{\uparrow b}$	$\frac{S_L P_T}{m_b} \hat{u}_1^{\uparrow b} + \frac{S_L \mathbf{k}_T}{m_N} \hat{u}_1^{\uparrow}$
	T	$\frac{P_T \times S_T}{m_b} \hat{u}_1^{\uparrow b} + \frac{\mathbf{k}_T \times S_T}{m_N} \hat{u}_1^{\uparrow}$	$\frac{P_T S_T}{m_b} \hat{u}_1^{\uparrow b} + S_T \hat{u}_1^{\uparrow}$	$\frac{P_T(P_T S_T)}{m_b^2} \hat{u}_1^{\uparrow b} + \frac{\mathbf{k}_T(\mathbf{k}_T S_T)}{m_N^2} \hat{u}_1^{\uparrow} + \frac{P_T(\mathbf{k}_T S_T) - \mathbf{k}_T(P_T S_T)}{m_N m_b} \hat{u}_1^{\uparrow b}$

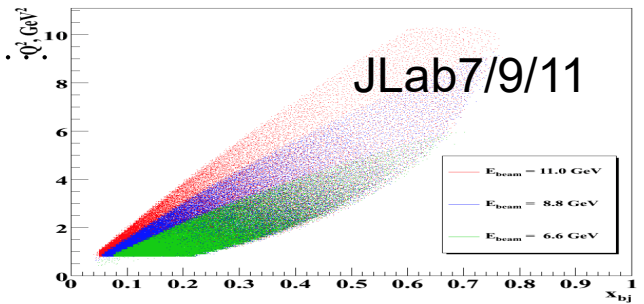
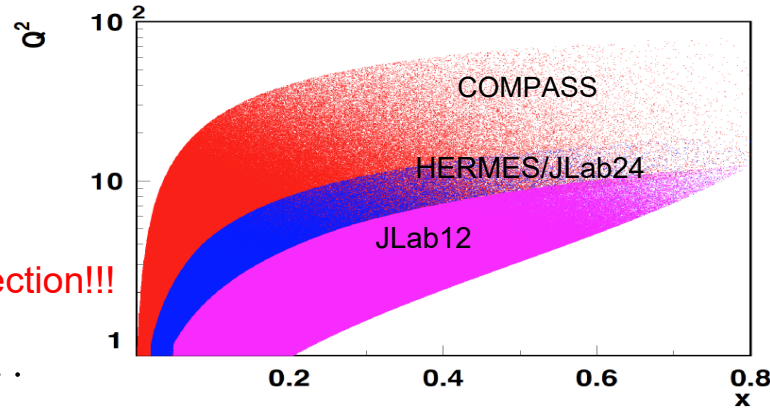
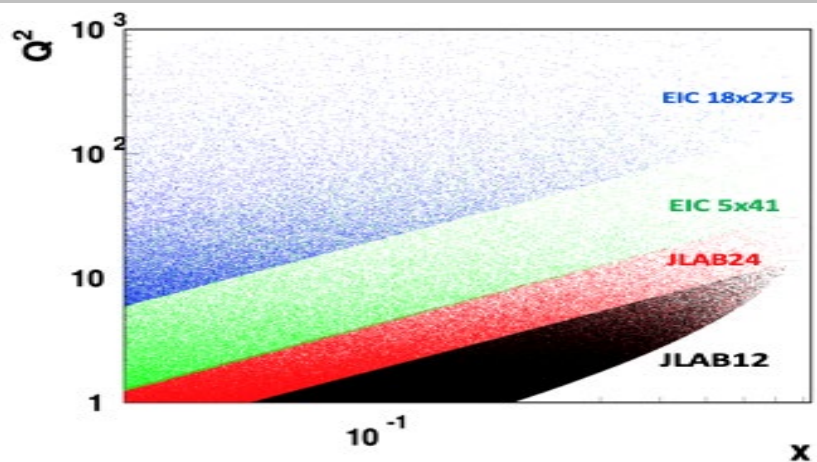
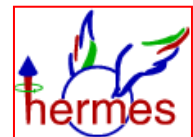
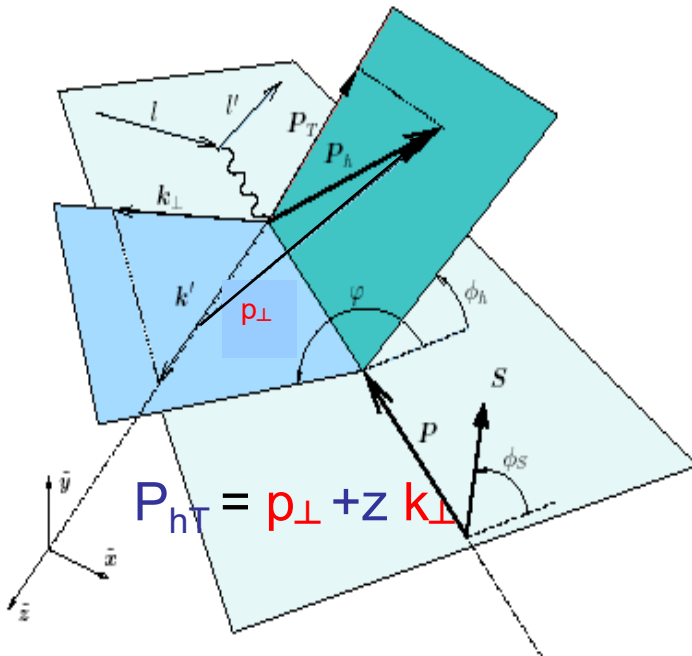
Anselmino, Barone, Kotzinian

		Quark polarization		
		U	L	T
Nucleon Polarization	U	$f_1^q(x, k_T^2)$	X	$\frac{e_f^2 k_T^J}{M} h_1^{qJ}(x, k_T^2)$
	L	X	$S_L g_{1L}^q(x, k_T^2)$	$S_L \frac{\mathbf{k}_T}{M} h_L^q(x, k_T^2)$
	T	$\frac{[\mathbf{k}_T \times S_T]}{M} f_T^q(x, k_T^2)$	$\frac{\mathbf{k}_T \cdot S_T}{M} g_T^q(x, k_T^2)$	$S_T h_T^q(x, k_T^2) + \frac{\mathbf{k}_T (\mathbf{k}_T \cdot S_T)}{M} h_T^q(x, k_T^2)$

CFR/Current Fragmentation Region
 → hadrons produced from struck quark
 TFR/Target Fragmentation Region
 → hadrons produced from remnant

Correlations of the spin of the target or/and the momentum and the spin of quarks, combined with final state interactions define the azimuthal distributions of produced particles (different in CFR and TFR)

SIDIS kinematical coverage and observables



SIDIS experiments measure azimuthal dependence of the cross section!!!

$$\sigma \propto F_{UU} + P_b \sqrt{2\epsilon(1-\epsilon)} F_{LU}^{\sin \phi} \sin \phi + P_t \epsilon F_{UL}^{\sin 2\phi} \sin 2\phi + \dots$$

$$+ \epsilon F_{UU,L} + |S_{\perp}| [F_{UT}^{\sin \phi - \phi_S} \sin(\phi - \phi_S) + \sqrt{2\epsilon(1+\epsilon)} F_{UT}^{\sin \phi_S} \sin \phi_S] + \dots$$

- Studies of azimuthal modulations give access to underlying 3D partonic distributions
- QCD predicts only the Q^2 -dependence of 3D PDFs

WG3 hadronization and TMDs plan

3 hours of discussion: (15+5) = 8 talks + 20mins flash

1.Role of the JLab upgrade to address critical elements for interpretation of world SIDIS data

1. Role of $F_{UU,L}$
 1. theory perspective (Alessandro/Alexey)
 2. experimental perspective (Ed Kinney)
2. Role of VM decays and p_T distributions (Albi)
3. Radiative Corrections in SIDIS (Jianwei/Nobuo)
4. Role of SIDIS regions (correlations between CFR/TFR/SFR..),
 1. Theory perspective (Prokudin/Sato?)

$$F_{UU,L} \sim P_T^2 Z^2 / Q^2$$

2.Opportunities of High luminosity JLab upgrade for precision TMD studies in valence region, the role of multidimensional binning

1. Experimental perspective (Bakur Parsamyan)
2. Theory perspective (Alexey Vladimirov)

3.Understanding hadronization [2 talks]

1. Independent fragmentation and role of charge symmetry (Whit Armstrong)
2. Correlations in Di-hadron production (Dilks)

4.Flash talks [5 talks] (3 mins)

1. **Semi-exclusive processes: VM production** (Andrei Afanasev)
2. **Understanding the scale dependence of spin and azimuthal asymmetries**
 1. Cahn effect (Andrea Moretti)
 2. SSAs (Stefan Diehl)
 3. Evolution studies (Harut)
3. **Opportunities with Tagged SIDIS**
 1. Theory perspective
 2. Theory perspective (Wim Cosyn?)
4. **Role of SIDIS regions (correlations between CFR/TFR/SFR..), experimental perspective**(Timothy Hayward)

JLab22

Search JLab22

Threads

Direct messages

Mentions & reactions

Drafts & sent

Slack Connect

More

Channels

general

hadron-structure-plans

wg1-hadron-spectra-and-structure-as-probes-of-qcd

wg2-partonic-structure-and-spin

wg3-hadronization-and-transverse-momentum

wg4-quark-gluon-form-factors-and-generalized-parton-distributions

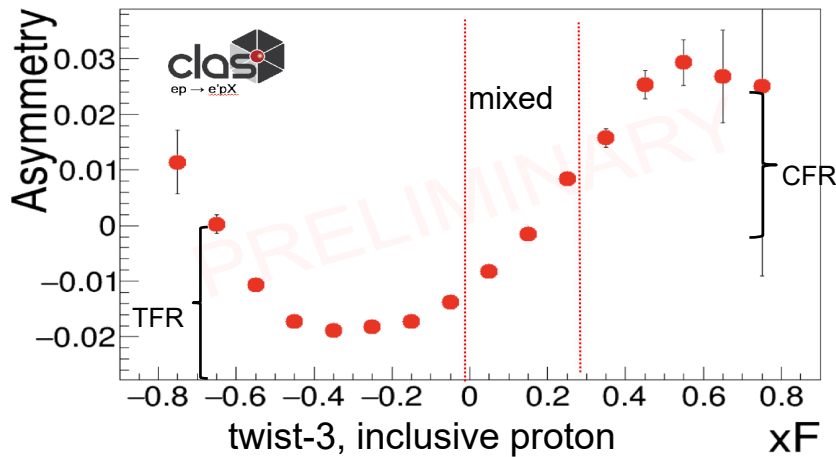
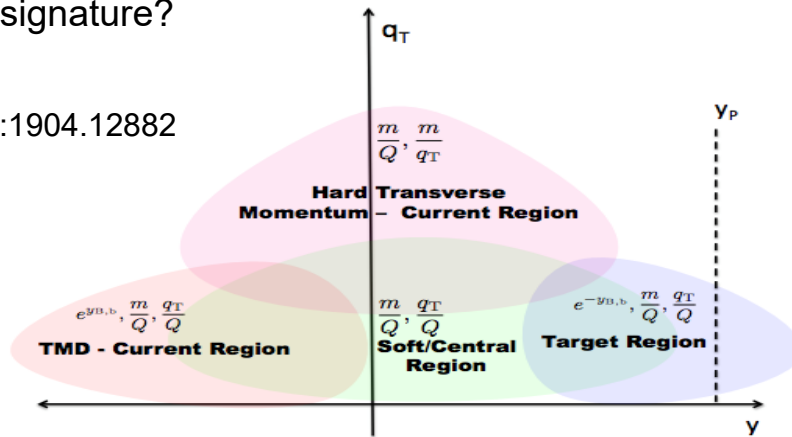
wg5-hadron-quark-transition-and-nuclear-dynamics-at-extreme-conditions

wg6-low-energy-tests-of-standard-model-and-fundamental-symmetries

Beam SSAs: Where is the struck quark?

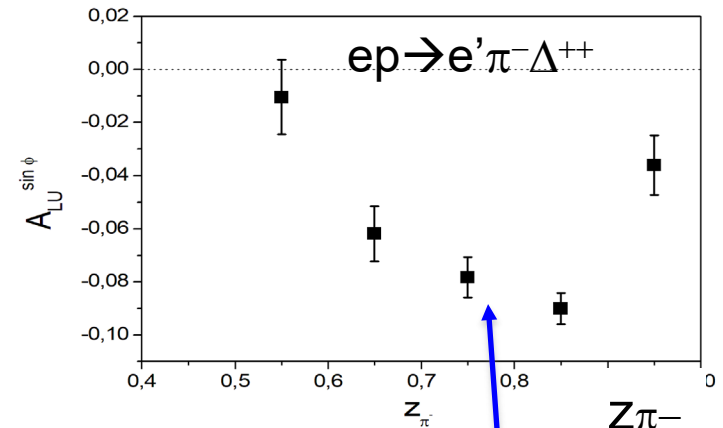
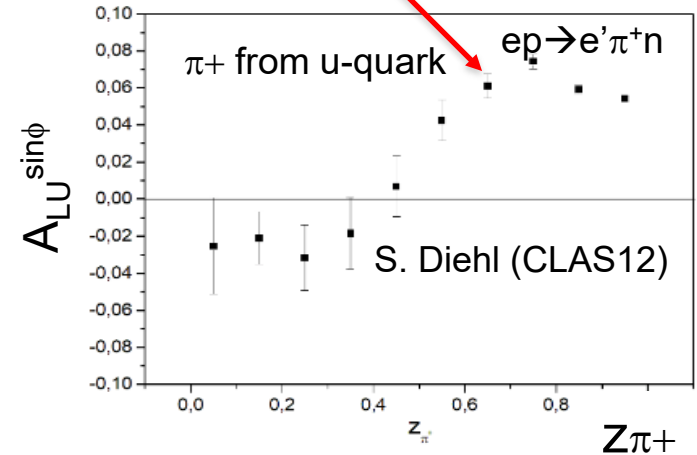
- Can we separate the CFR from TFR using the SSA as signature?

arXiv:1904.12882



Negative sign of the SSA (plateau) defines the TFR dominance

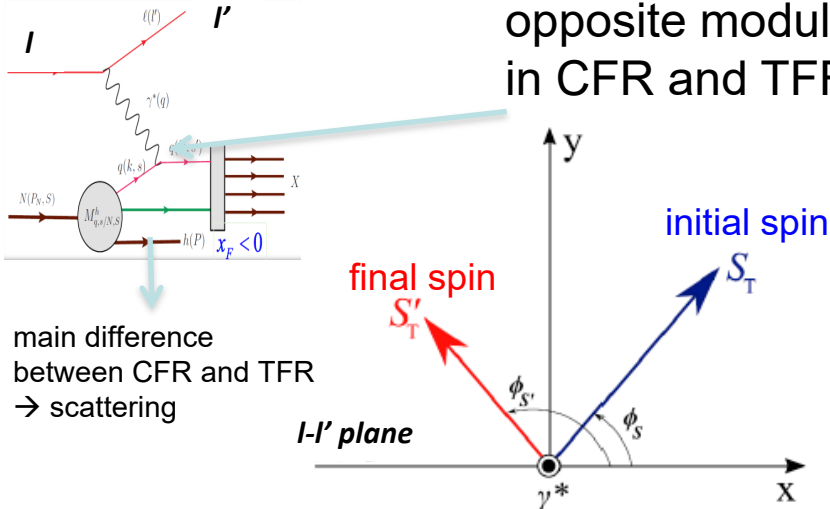
Polarized u-quark, dominates
→ SSA positive



Polarized d-quark, is hard to locate, and one obvious process where we can guarantee it was hit, is the production of Δ^{++} (negative SSA)

Hadron production in TFR

X. Tong (ECT* 2022)



$$\phi_{s'} = \pi - \phi_s$$

Collins SSAs, $\sin(\phi + \phi_s), \sin 2\phi$

spin-flip causing Collins contributions in CFR

Target fragmentation provides complementary information on proton structure and is simpler in interpretation (no Collins effect, both in leading and subleading contributions)

$$F_{UU}^{\cos(\phi_h)} = -2 \sum_a e_a^2 x_B^2 \frac{|k_\perp|}{Q} u(x_B, \xi, k_\perp),$$

$$F_{LU}^{\sin(\phi_h)} = 2 \sum_a e_a^2 x_B^2 \frac{|k_\perp|}{Q} l(x_B, \xi, k_\perp)$$

$$F_{LL}^{\cos(\phi_h)} = -2 \sum_a e_a^2 x_B^2 \frac{|k_\perp|}{Q} l_L(x_B, \xi, k_\perp)$$

$$F_{UL}^{\sin(\phi_h)} = -2 \sum_a e_a^2 x_B^2 \frac{|k_\perp|}{Q} u_L(x_B, \xi, k_\perp)$$

4 contributions for each in CFR!!!

• Two with unpolarized target; Two with longitudinally polarized target;

Proton polarization: x-sections

Semi-Inclusive:

Lepton helicity → “+1” along the beam

$$\frac{d\sigma}{dx dy d\psi dz d\phi_h dP_{h\perp}^2} = \frac{\alpha^2}{xyQ^2} \frac{y^2}{2(1-\varepsilon)} \left(1 + \frac{\gamma^2}{2x}\right) \left\{ F_{UU,T} + \varepsilon F_{UU,L} + \lambda_e \sqrt{2\varepsilon(1-\varepsilon)} \sin\phi_h F_{LU}^{\sin\phi_h} \right. \\ \left. + S_{\parallel} \left[\sqrt{2\varepsilon(1+\varepsilon)} \sin\phi_h F_{UL}^{\sin\phi_h} + \varepsilon \sin(2\phi_h) F_{UL}^{\sin 2\phi_h} \right] + S_{\parallel} \lambda_e \sqrt{1-\varepsilon^2} F_{LL} \right\}$$

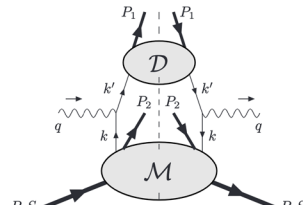
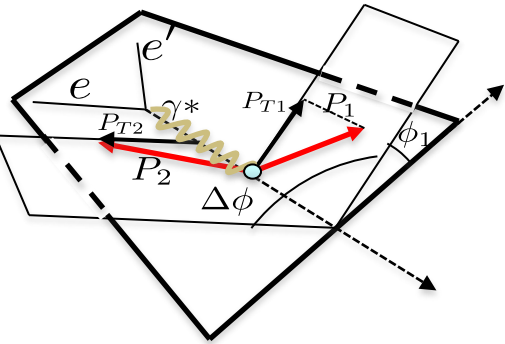
Proton helicity → “+1” opposite to the beam

NH3 – 14 runs from latest RGC run group with beam energy 10.55 GeV used for tests, first preliminary results, to be released using pass0 RGC processing in January

Double polarized experiments allow studies of single beam-spin, single target-spin and double spin asymmetries

Correlations in back-to-back 2 hadron production

M. Anselmino, V. Barone and A. Kotzinian,
Physics Letters B 713 (2012)



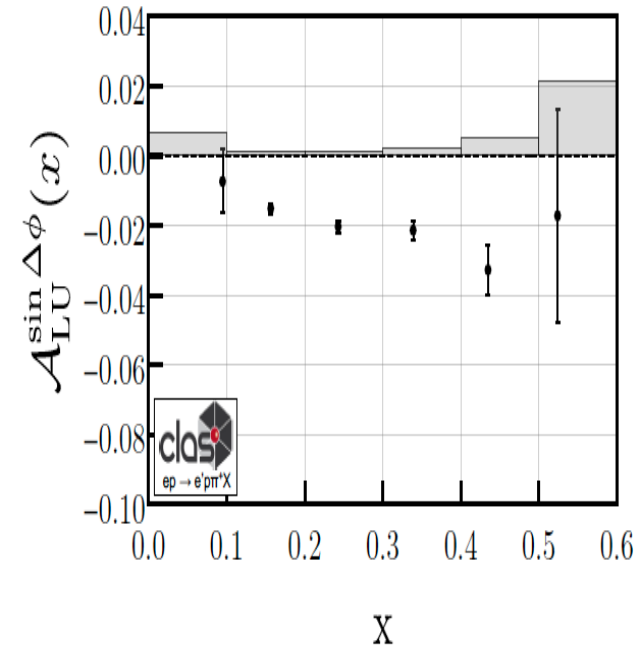
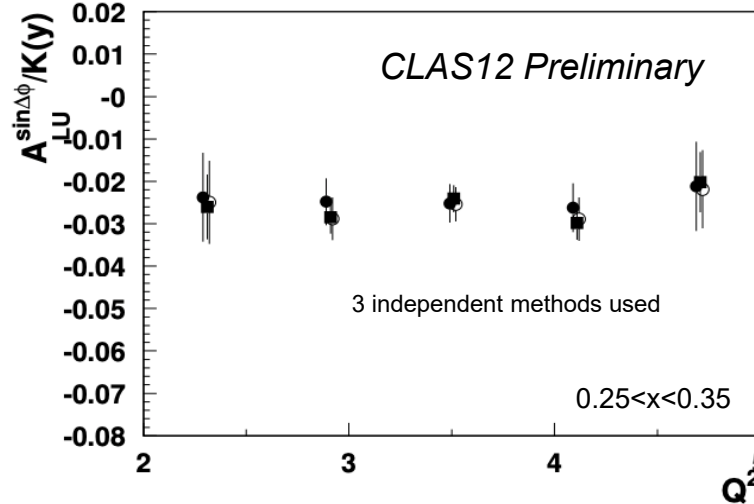
$$A_{LU} \propto \frac{C[w_5 \hat{l}_1^{\perp h} D_1]}{C[\hat{u}_1 D_1]} \sin \Delta\phi$$

arXiv: [2208.05086](https://arxiv.org/abs/2208.05086)

PRL

Twist-2 table
(Fracture Functions)

N/q	U	L	T
U	\hat{u}_1	\hat{l}_1^h	$\hat{t}_1^h, \hat{t}_1^\perp$
L	\hat{u}_{1L}^\perp	\hat{l}_{1L}	$\hat{t}_{1L}^h, \hat{t}_{1L}^\perp$
T	$\hat{u}_{1T}^h, \hat{u}_{1T}^\perp$	$\hat{l}_T^h, \hat{l}_T^\perp$	$\hat{t}_{1T}, \hat{t}_{1T}^h, \hat{t}_{1T}^\perp, \hat{t}_{1T}^{hh}$

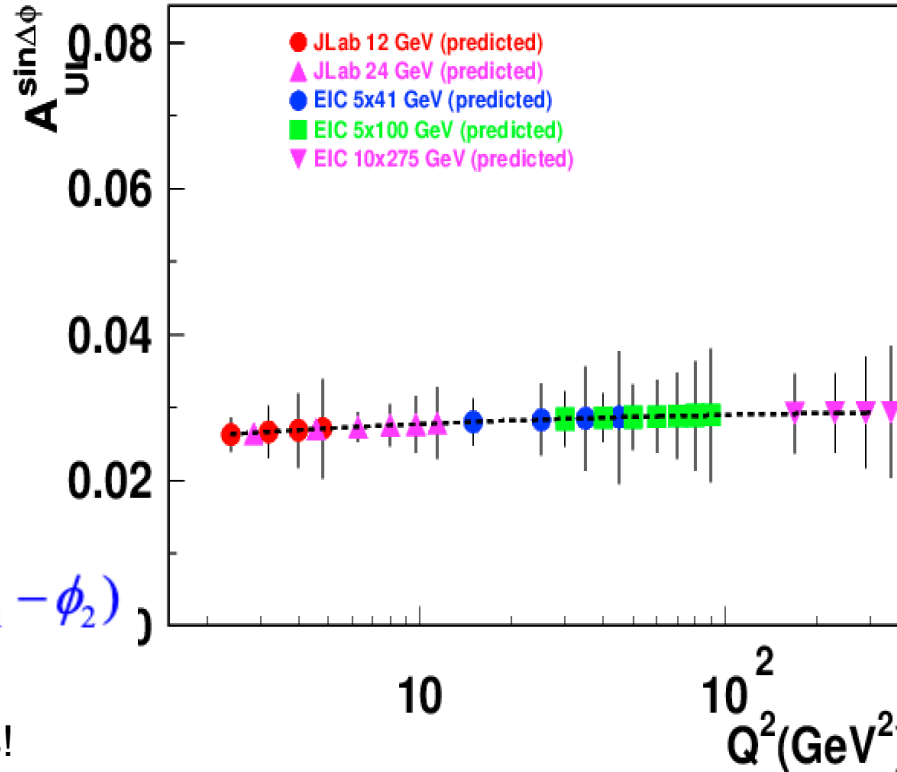


- SSA significant at large x , where the valence quarks (non-perturbative sea) dominate?
- Correlation asymmetry is linked to Leading Twist(LT) distributions of **longitudinally polarized quarks**
- First indication in large x SIDIS of a LT observable
- **Correlation between the struck quark and the remnant produces correlation between hadrons (entanglement)**
- Multidimensional measurements crucial for evolution studies

B2B correlations with long. Pol. Target

N/q	U	L	T
U	\hat{u}_1	$\hat{l}_1^{\perp h}$	$\hat{t}_1^h, \hat{t}_1^{\perp}$
L	$\hat{u}_{1L}^{\perp h}$	\hat{l}_{1L}	$\hat{t}_{1L}^h, \hat{t}_{1L}^{\perp}$
T	$\hat{u}_{1T}^h, \hat{u}_{1T}^{\perp}$	$\hat{l}_{1T}^h, \hat{l}_{1T}^{\perp}$	$\hat{t}_{1T}^h, \hat{t}_{1T}^{hh}, \hat{t}_{1T}^{\perp}, \hat{t}_{1T}^{\perp h}$

Lumi: JLab 10^{35} , EIC4x51/5x100/10x275 0.044,0.6,1x10³⁴)
 $y > 0.05, 100$ days



CLAS12 proposals

NH3/ND3
[E12-09-009](#)

[E12-07-107](#)
[E12-09-007A](#)

³He
[C12-20-002](#)

⁷LiD
[E12-14-001](#)

A. Kotzinian, arXiv:1107.2292

$$\sigma_{UU} = F_0^{\hat{u} \cdot D_1}$$

$$\sigma_{UL} = -\frac{P_{T1} P_{T2}}{m_2 m_N} F_{k1}^{\hat{u}_{1L}^{\perp h} \cdot D_1} \sin(\phi_1 - \phi_2)$$

No depolarization, like Sivers!

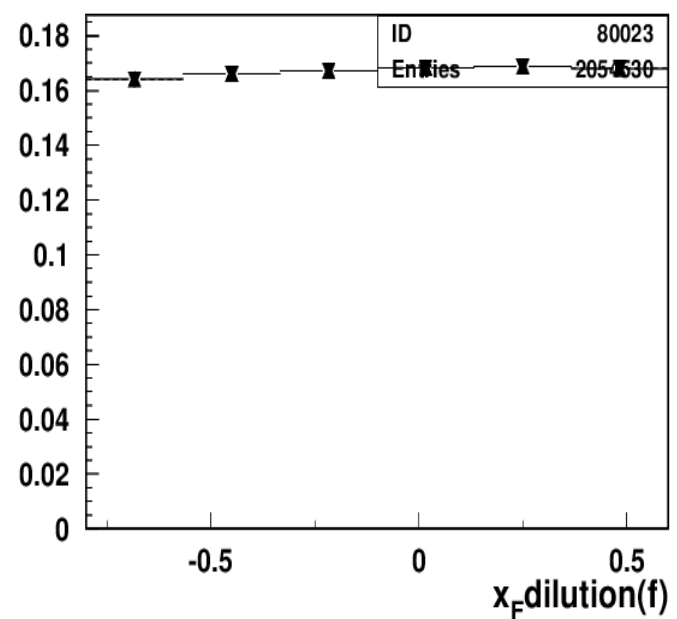
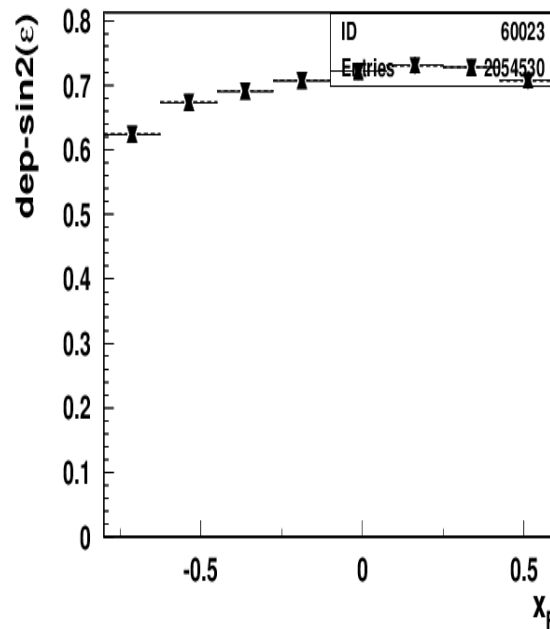
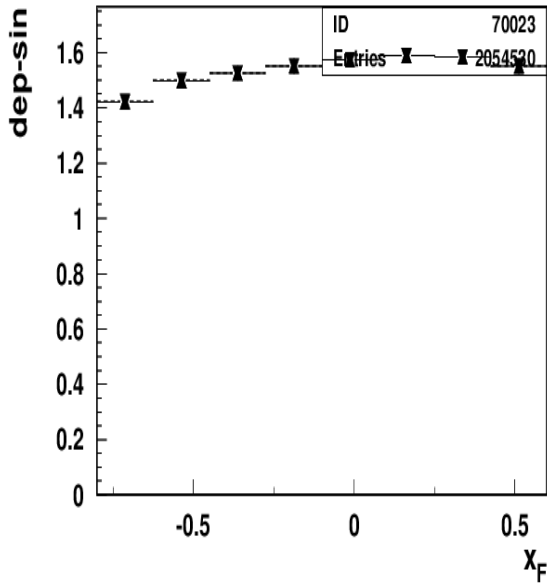
- Target SSA can be measured in the full Q^2 range, combining different facilities
- Advantages: Higher Lumi for JLab, less suppression at high Q^2 for EIC
- JLab24 will be crucial to bridge the studies of FFs between JLab12 and EIC in the valence region

Proton polarization: Double spin asymmetries

Semi-Inclusive:

$$\frac{d\sigma}{dx dy d\psi dz d\phi_h dP_{h\perp}^2} = \frac{\alpha^2}{xyQ^2} \frac{y^2}{2(1-\varepsilon)} \left(1 + \frac{\gamma^2}{2x}\right) \left\{ F_{UU,T} + \varepsilon F_{UU,L} + \lambda_e \sqrt{2\varepsilon(1-\varepsilon)} \sin\phi_h F_{LU}^{\sin\phi_h} \right.$$

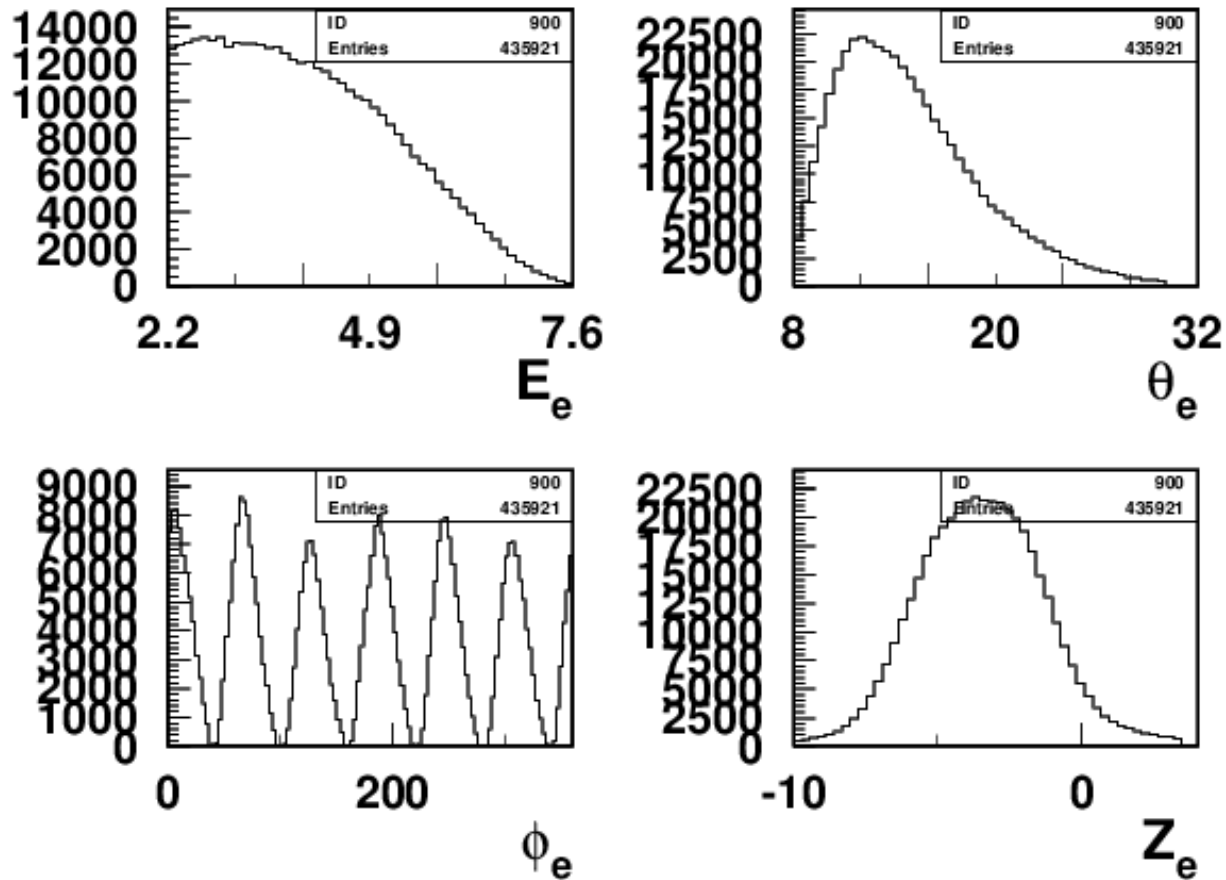
$$\left. + S_{\parallel} \left[\sqrt{2\varepsilon(1+\varepsilon)} \sin\phi_h F_{UL}^{\sin\phi_h} + \varepsilon \sin(2\phi_h) F_{UL}^{\sin 2\phi_h} \right] + S_{\parallel} \lambda_e \sqrt{1-\varepsilon^2} F_{LL} \right\}$$



The $\sin\phi$ modulation is enhanced compared to all other by a factor of ~ 2

RGC NH3: $ep \rightarrow e'pX$

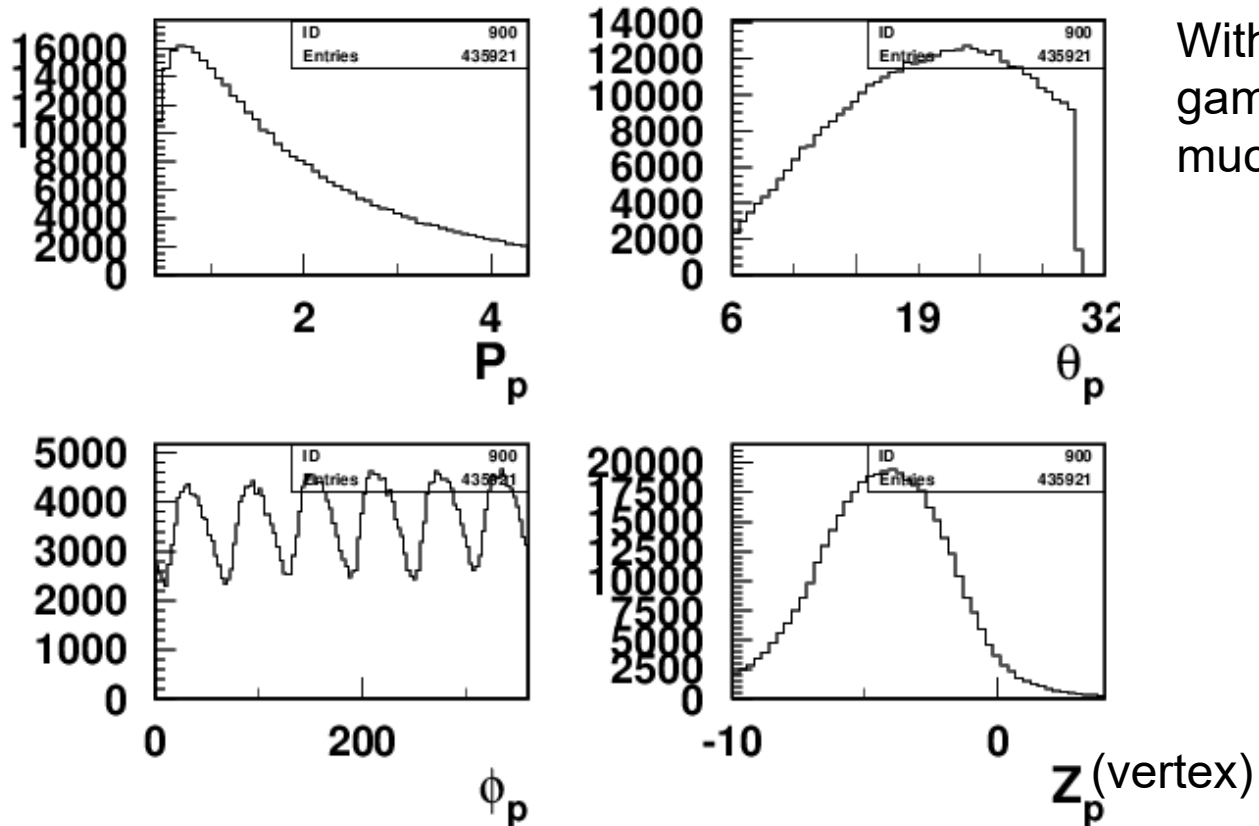
Kinematical distributions in epX: electrons



Cuts on electrons: $E > 2.6$, $W > 2$, $Q^2 > 1.0$, remove edges in phi

RGC NH3: $ep \rightarrow e'pX$

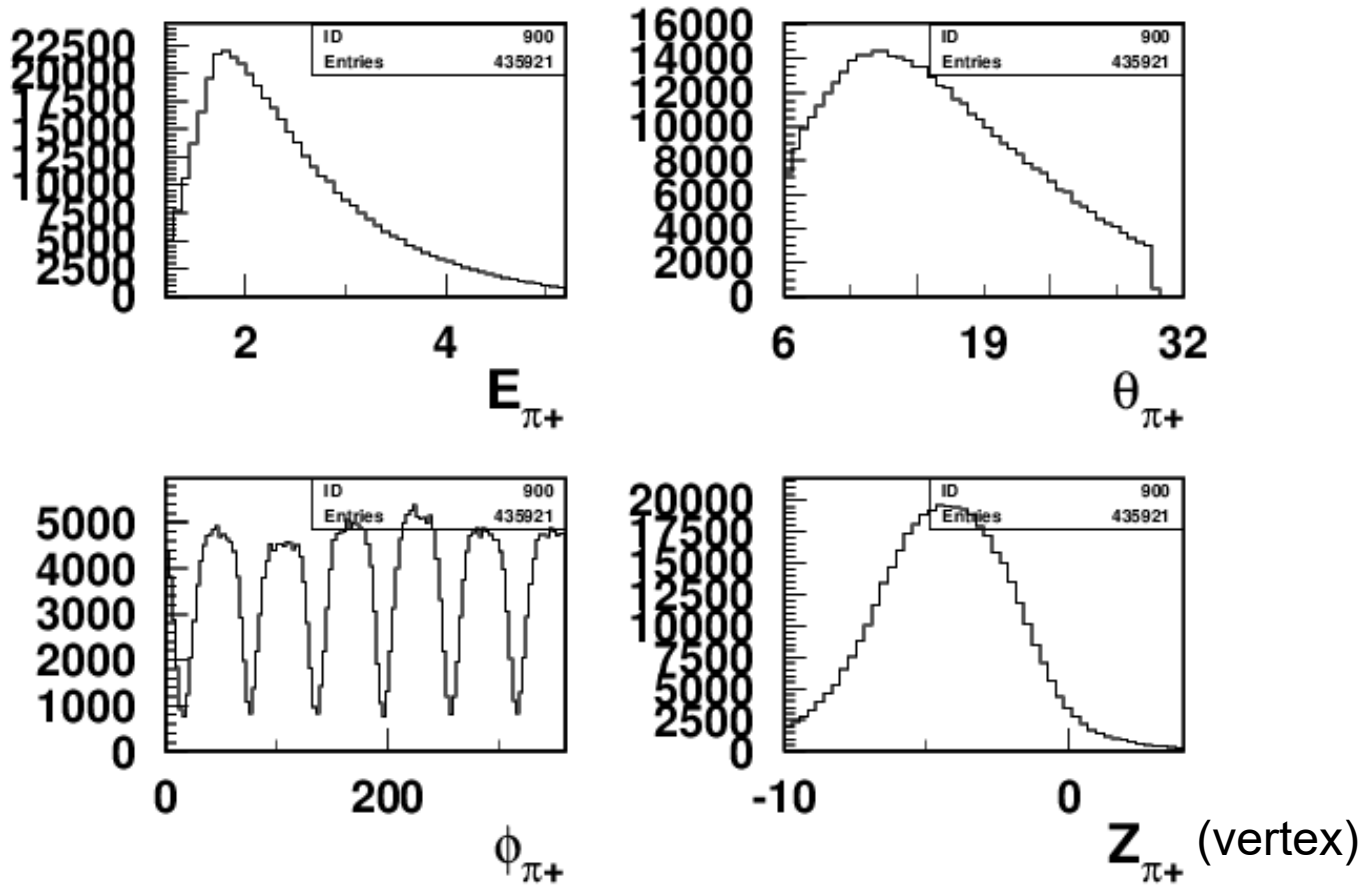
Kinematical distributions in epX: protons



Cuts on electrons: $P > 0.5$, $ID = 2212$, only FD, $|\chi^2| < 5$, $|\Delta Vz| < 7$

RGC NH3: $ep \rightarrow e' p \pi + X$

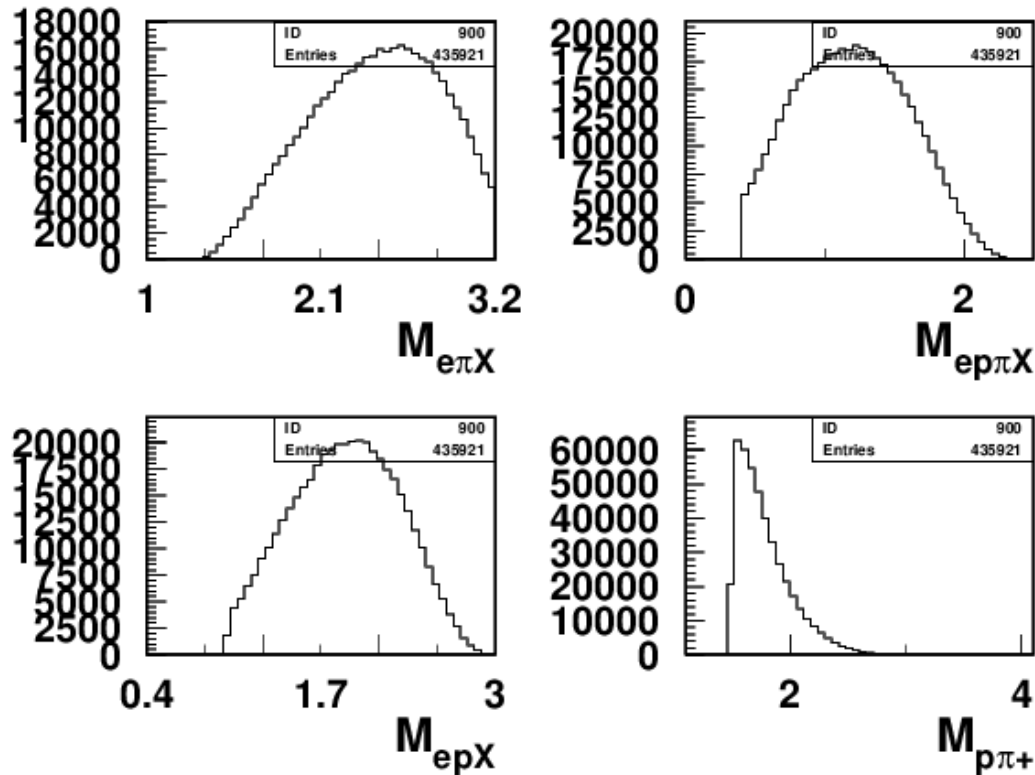
Kinematical distributions in $ep\pi X$: pions



Cuts on electrons: $P > 1.2$, ID=211, only FD, $|\chi^2| < 5$, $|\Delta Vz| < 7$

RGC NH3: $ep \rightarrow e' p \pi + X$

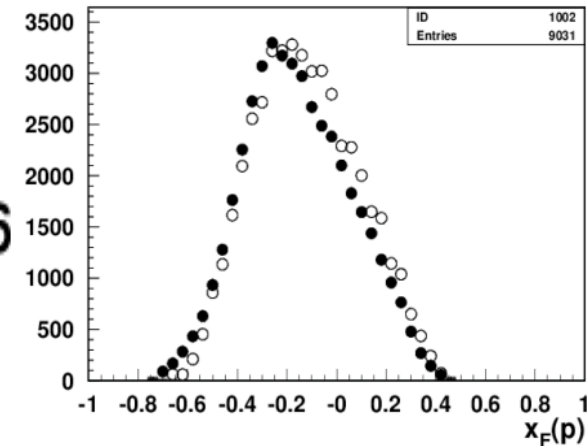
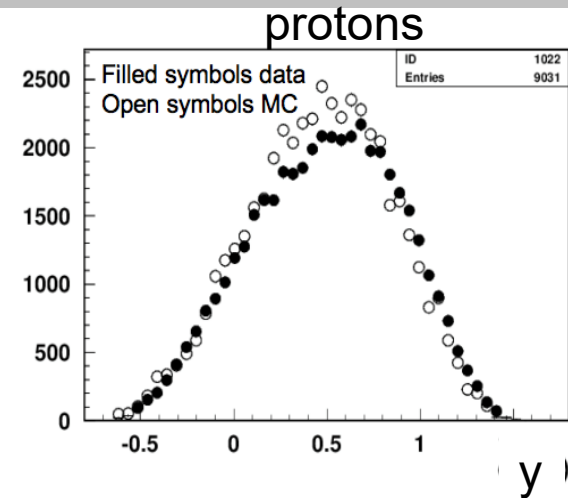
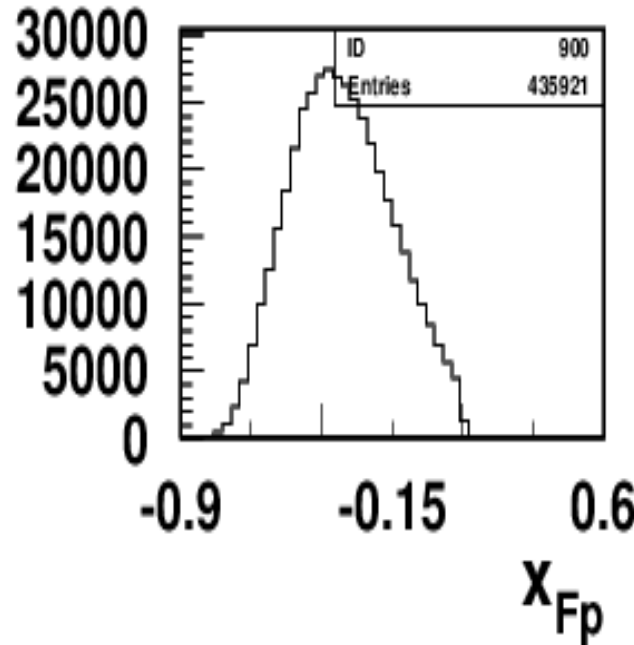
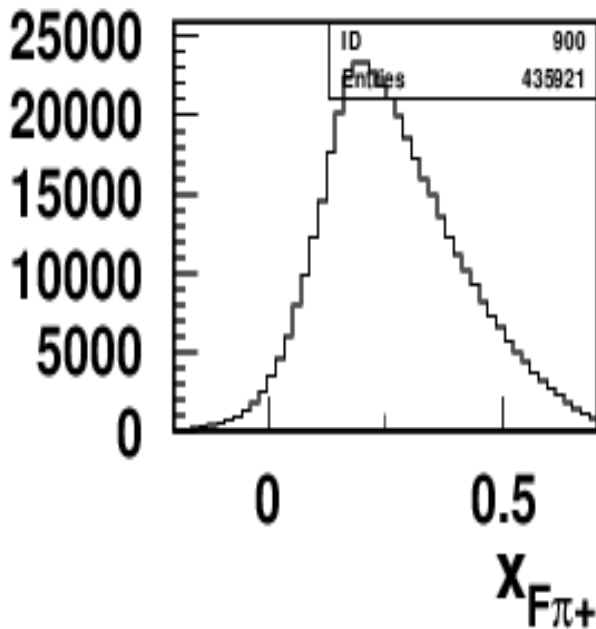
Kinematical distributions in $ep\pi X$: missing masses



So far using cuts on missing masses to exclude exclusive states
Missing mass dependences will be extracted to understand the origin

RGC NH3: $ep \rightarrow e' p \pi + X$

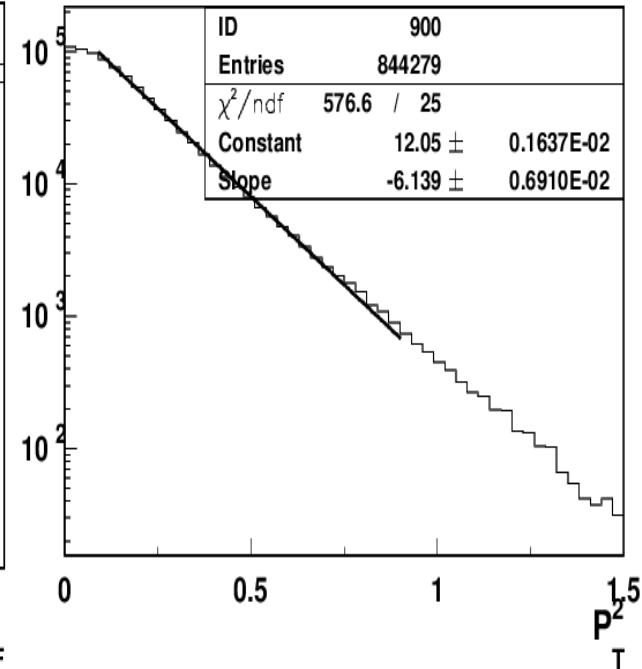
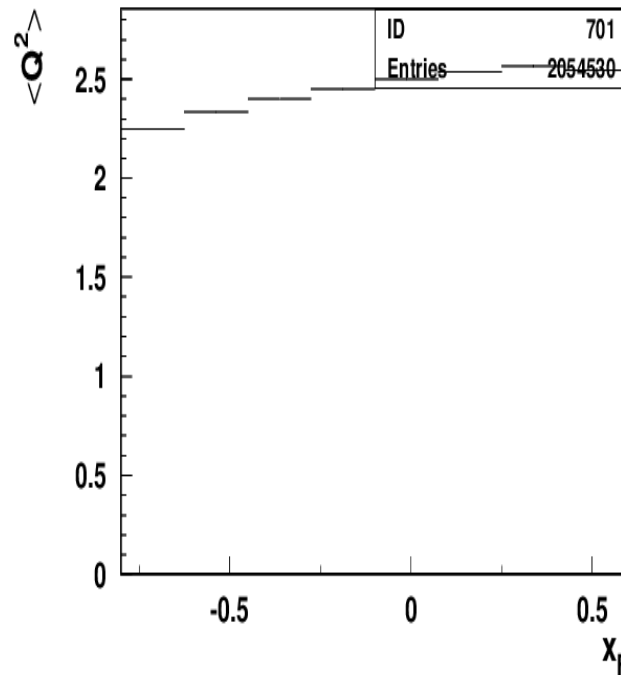
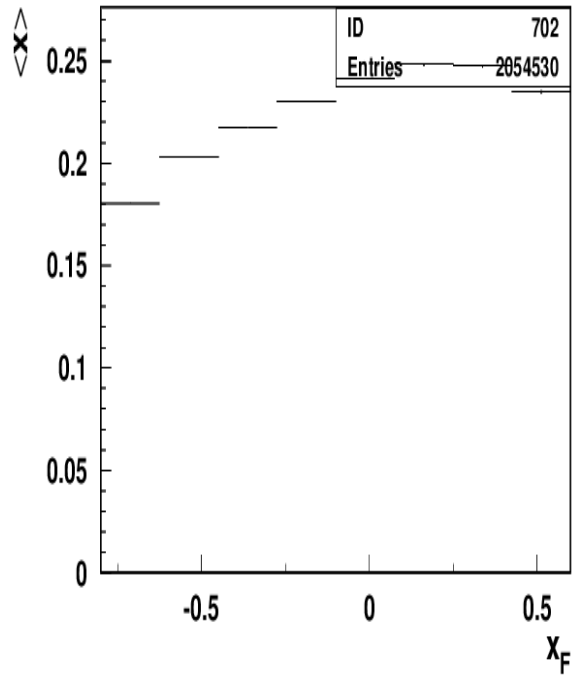
Kinematical distributions in $ep\pi X$: x_F -distributions



Protons in the negative hemisphere (TFR) pions in the positive (CFR)
 Distributions in x_F and rapidity (y) in good agreement with MC

RGC NH3: $ep \rightarrow e'pX$

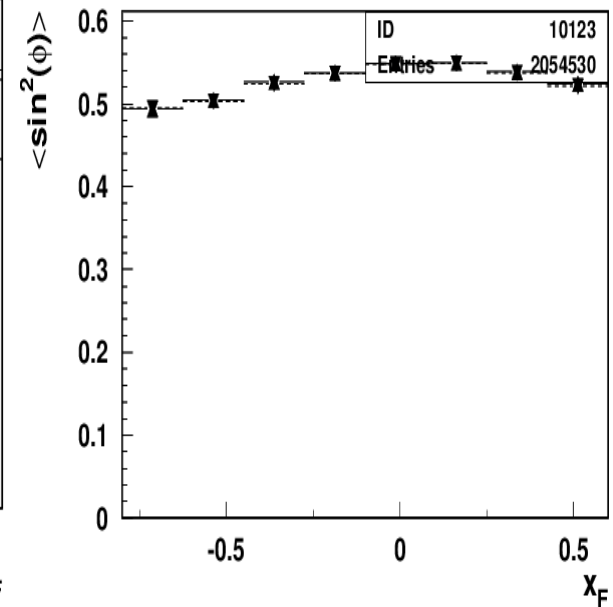
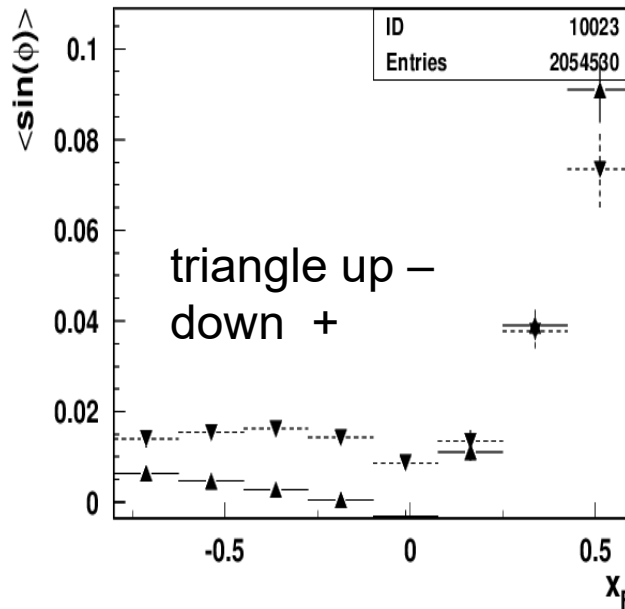
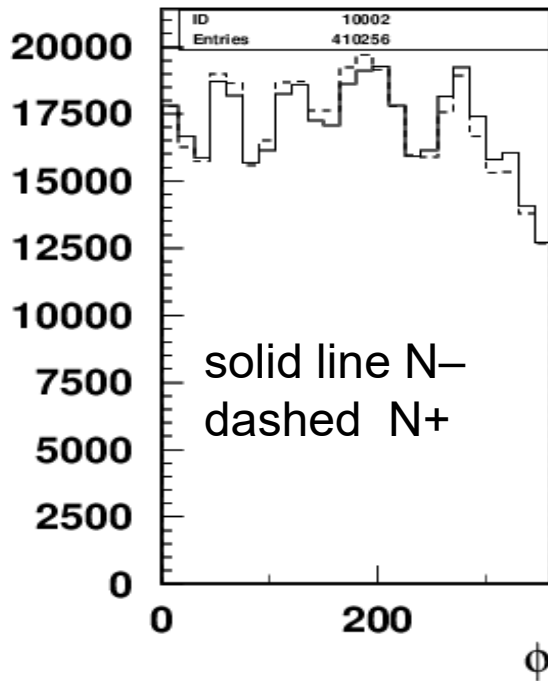
Kinematical correlations of variables for epX



No significant variations of the electron kinematics for different bins in x_F
Proton distributions in transverse momentum close to gaussian

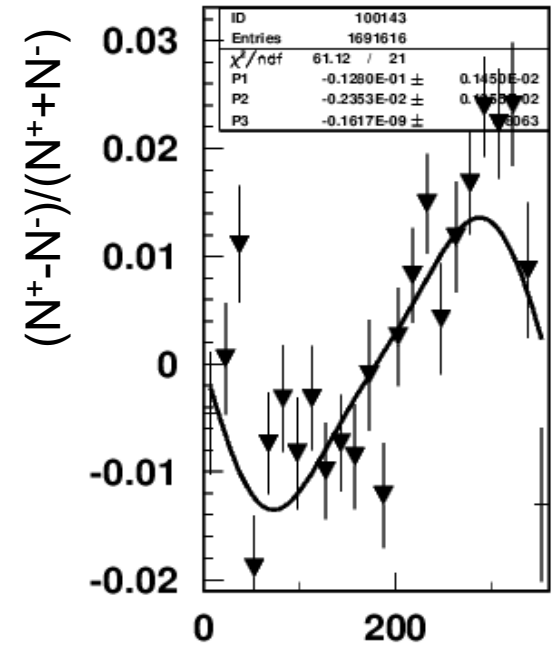
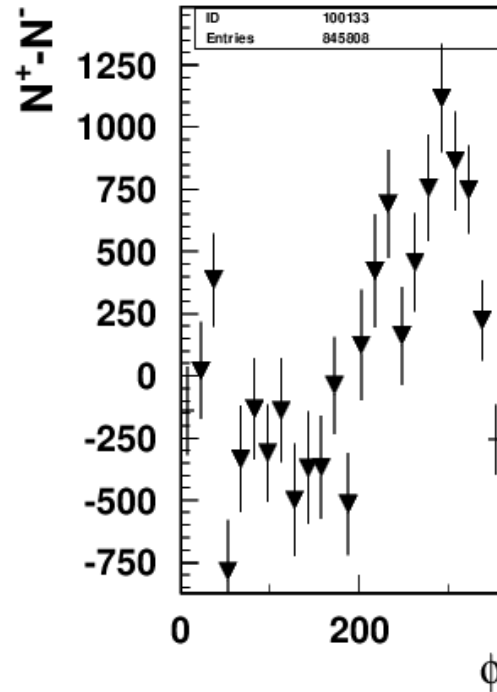
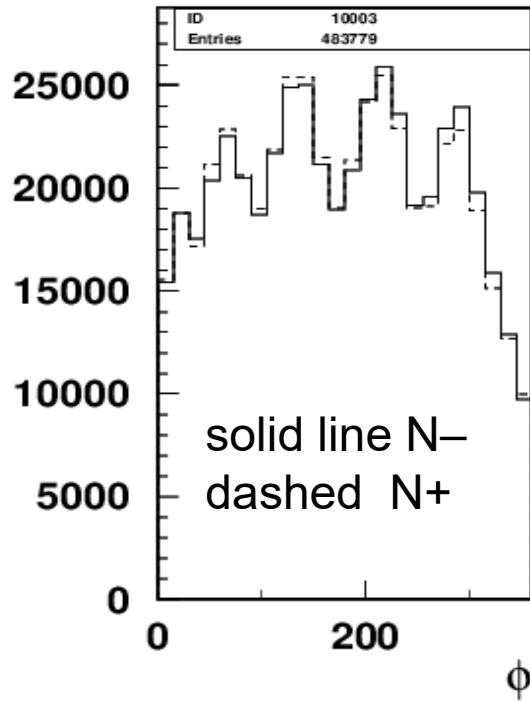
RGC NH3: azimuthal modulations $ep \rightarrow e'pX$

$$\frac{F_{UL}}{F_{UU}} = 1/f/D(y)/P_T \times \frac{\langle \sin \phi \rangle}{\langle \sin^2 \phi \rangle}$$



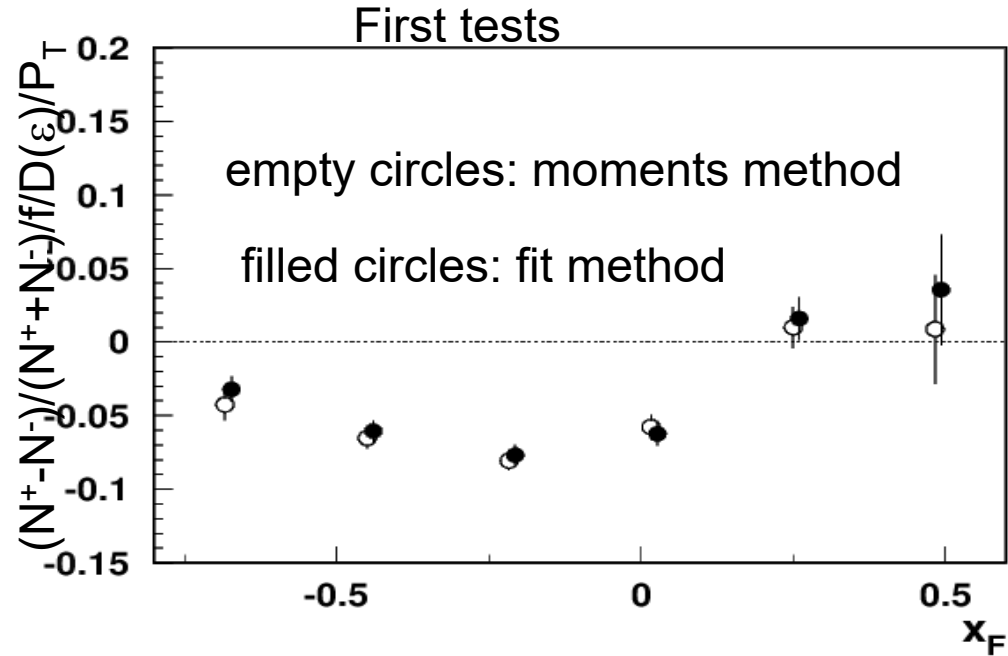
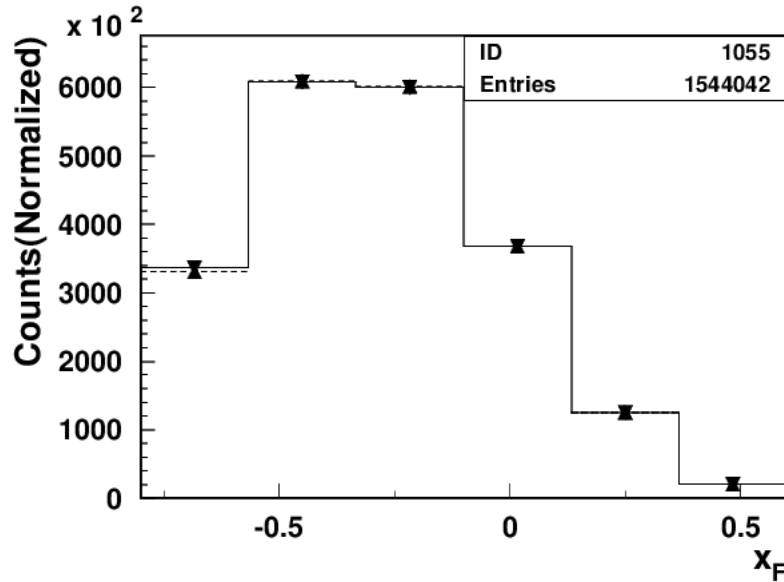
While azimuthal distributions complex, need both polarizations to define the A_{UL}

RGC NH3: azimuthal modulations $ep \rightarrow e'pX$



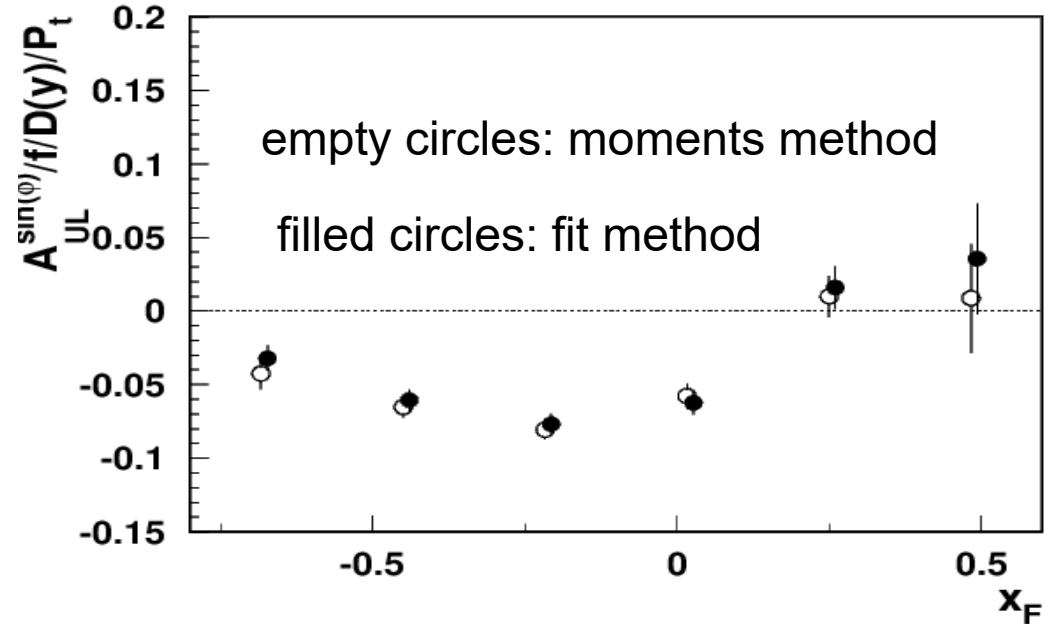
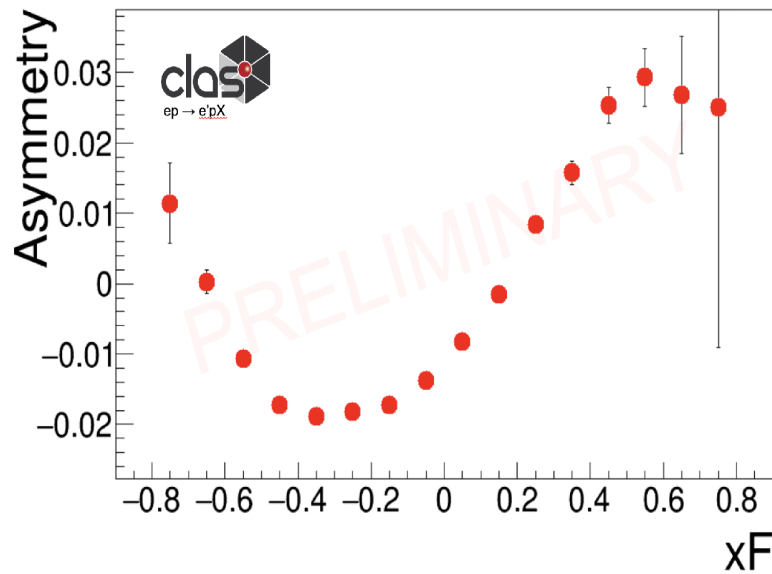
While azimuthal distributions complex, differences look like sinusoids

RGC NH3: azimuthal modulations $ep \rightarrow e'pX$



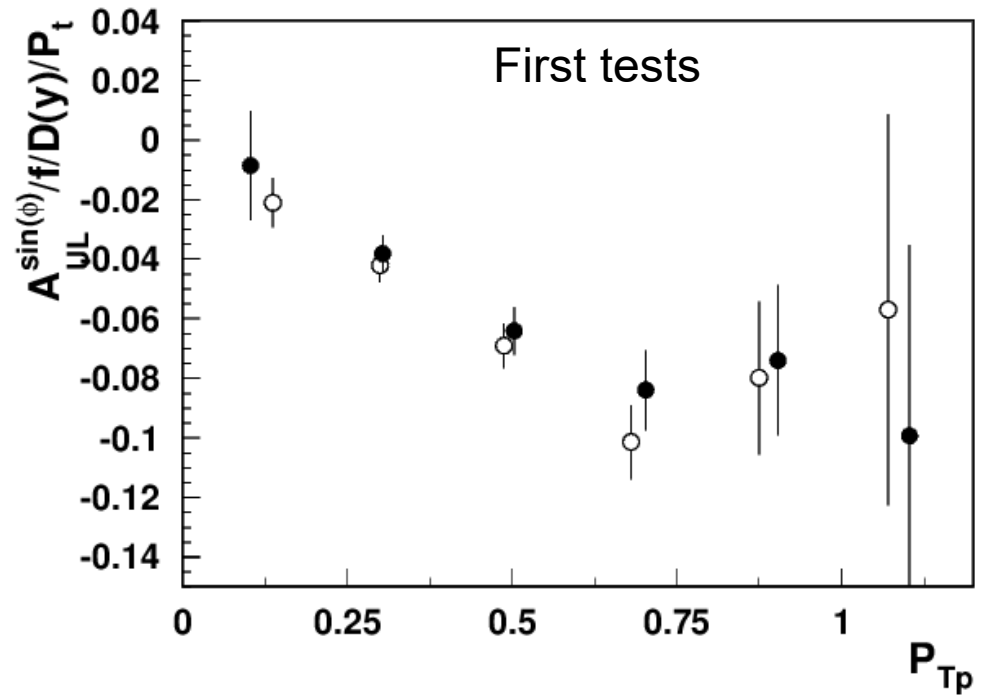
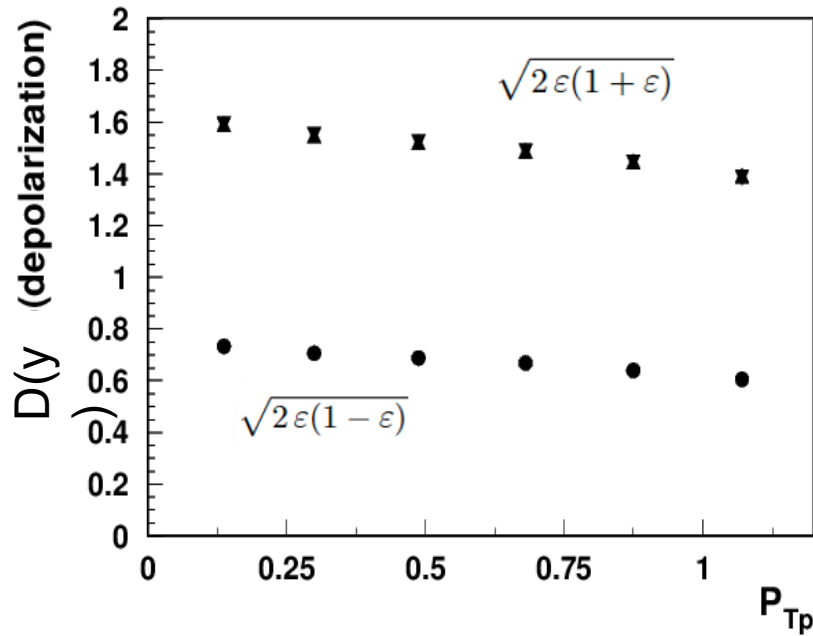
Test results: Both extraction methods agree indicating a huge SSA

RGC NH3: azimuthal modulations $ep \rightarrow e'pX$



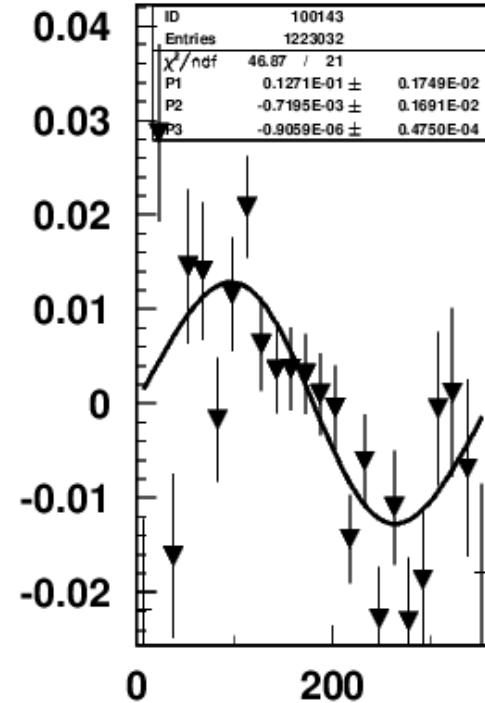
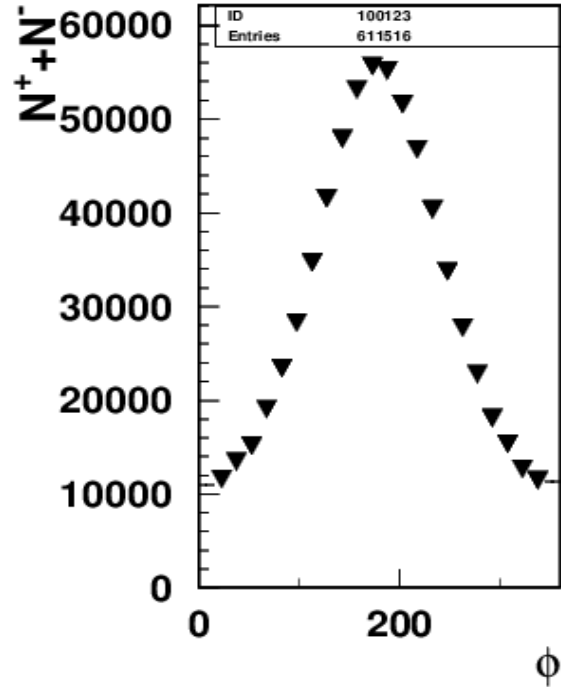
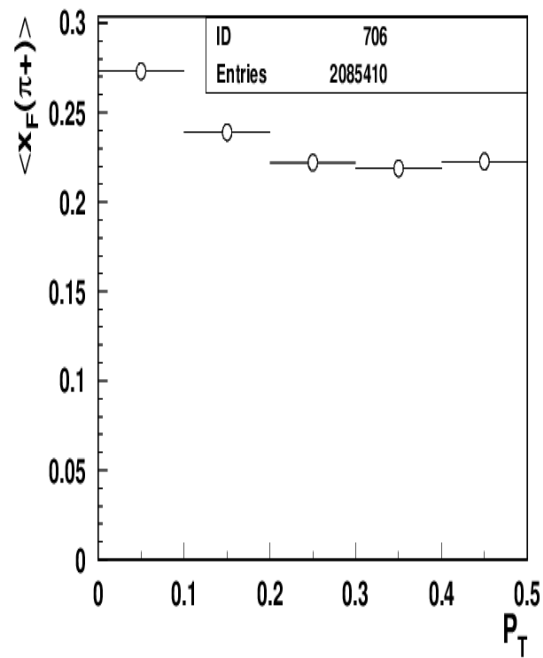
- Beam and target sinusoidal modulations most likely related to longitudinally polarized quarks
- SSAs indicate that protons up to $x_F \sim 0.3$ come mostly from target fragmentation

RGC NH3: azimuthal modulations $ep \rightarrow e'pX$



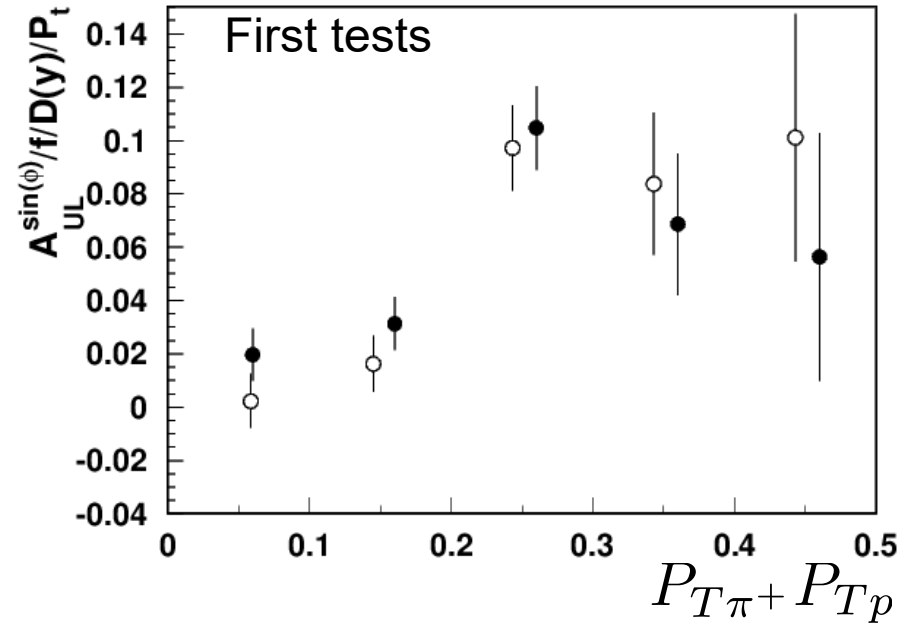
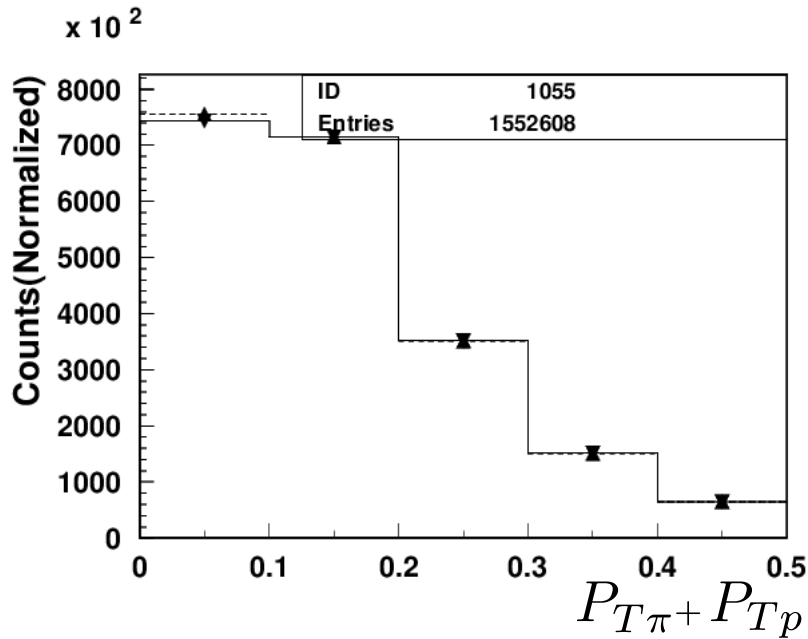
- Beam and target sinusoidal modulations most likely related to longitudinally polarized quarks
- SSAs indicate that protons up to $x_F \sim 0.3$ come mostly from target fragmentation

RGC NH3: azimuthal modulations $ep \rightarrow e' p \pi^+ X$



- Beam and target sinusoidal modulations most likely related to longitudinally polarized quarks
- SSAs indicate that protons up to $x_F \sim 0.3$ come mostly from target fragmentation

RGC NH3: azimuthal modulations $ep \rightarrow e' p \pi + X$



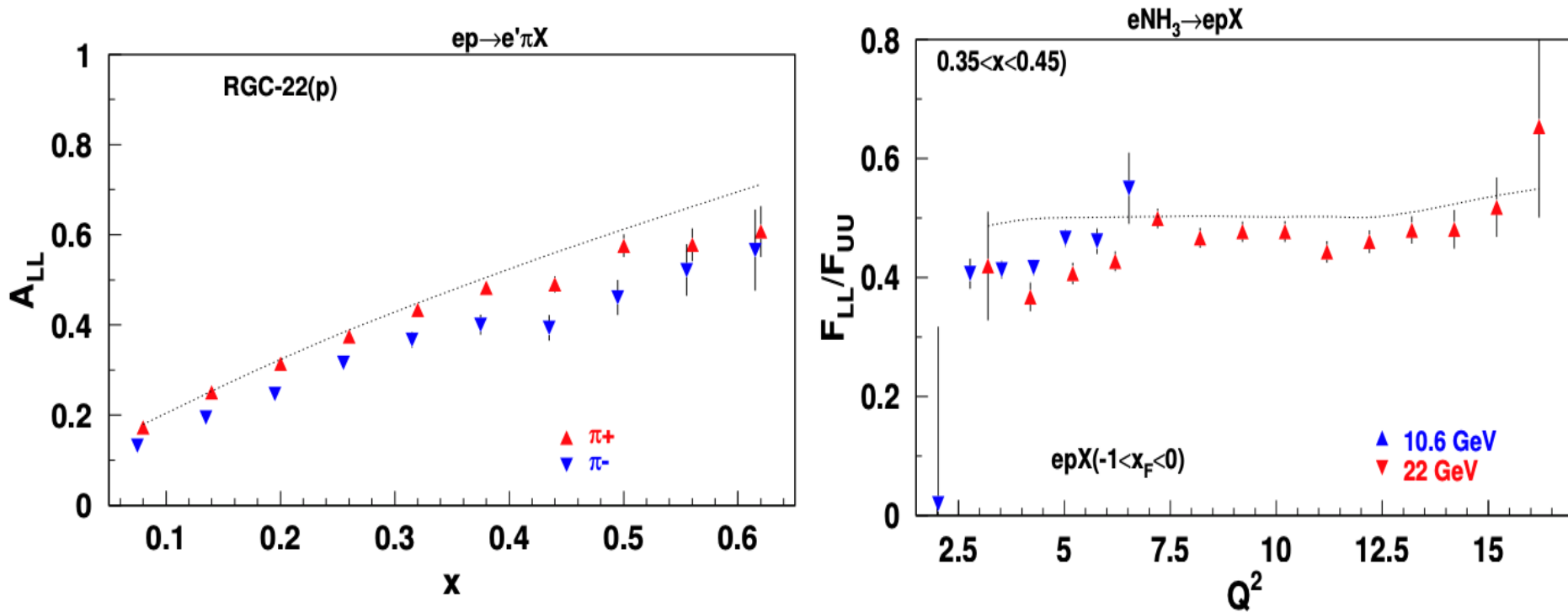
$$\sigma_{UU} = F_0^{\hat{u} \cdot D_1}$$

$$\sigma_{UL} = -\frac{P_{T1} P_{T2}}{m_2 m_N} F_{k1}^{\hat{u}_{1L}^{lh} \cdot D_1} \sin(\phi_1 - \phi_2)$$

- Significant correlations observed in b2b pion production with longitudinally polarized NH3
- Order of magnitude more statistics already available
- Q^2 -dependence will be studied (like for beam SSA)

CLAS12 at 22 GeV with longitudinally polarized target

Full simulations using LUND-based generator and full CLAS12 reconstruction chain



- Studies of evolution of observed double spin asymmetries will be a critical task in validating the QCD predictions $g_1(x, k_T)$ -studies CLAS12
- Asymmetries measured with input polarized and unpolarized PDFs, can be used to test the flavor decomposition capabilities
- Kinematical correlations, even for small bins relevant (multidimensional bins critical)

summary

- Sinusoidal modulations of the hadronic cross sections observed with unpolarized and longitudinally polarized targets, provide direct access to correlation of orbital momentum and longitudinal polarization of quarks
- Test results indicate significant single target spin asymmetries in epX and $b2b\ ep\pi X$ SSAs from RGC longitudinally polarized target data
- SSAs can be used in separation of kinematical regions TFR/Soft/CFR
- Kinematical dependence of SSA of TFR protons with longitudinally polarized target support the origin of beam/target SSAs related to longitudinally polarized quarks.
- Full RGC data set would allow studies P_T and Q^2 dependence of single and double spin asymmetries

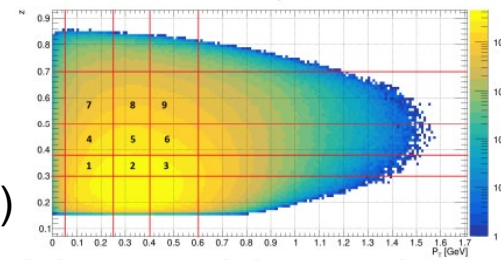
TODO list

- Precision measurements for dilution factors for all relevant processes will be required for more precise measurements of all observables (cross checked by MC), systematics from target polarization
- Define the role of correlations in CFR SSAs

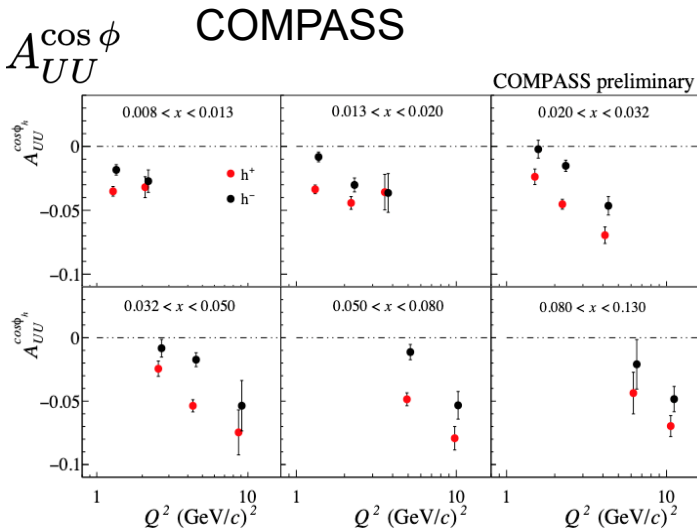
support

Attempts to understand Q^2 -dependence of HT

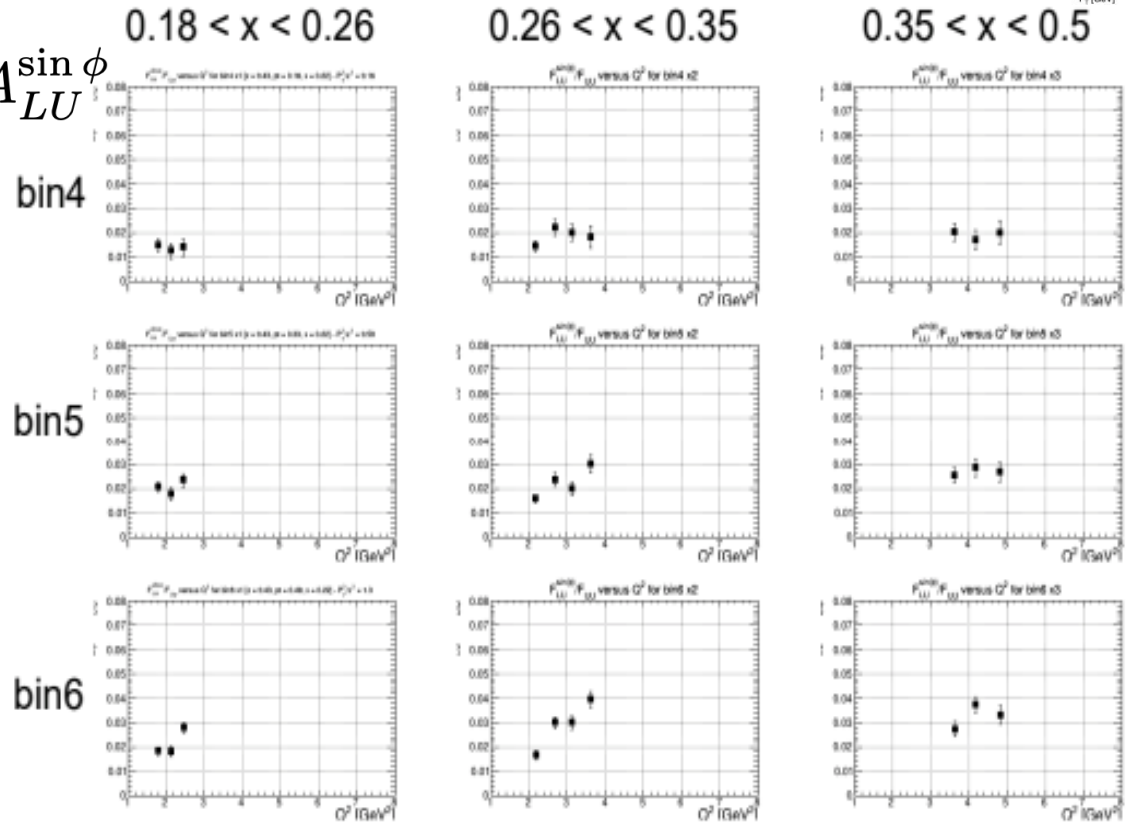
z vs P_T



CLAS12(preliminary)



$A_{LU}^{\sin \phi}$



- We always measure ratio to

$$F_{UU,T} + \epsilon F_{UU,L}$$

- The moments defined as a ratio to ϕ -independent x-section(to $F_{UU,T}$), are not decreasing with Q !!!
- The HT observables, don't look much like HT observables, something missing in understanding
- **Understanding of these behavior can be a key to understanding of other inconsistencies**
- Checking the Q^2 and P_T -dependences of the $F_{UU,L}$ may provide crucial input for validation

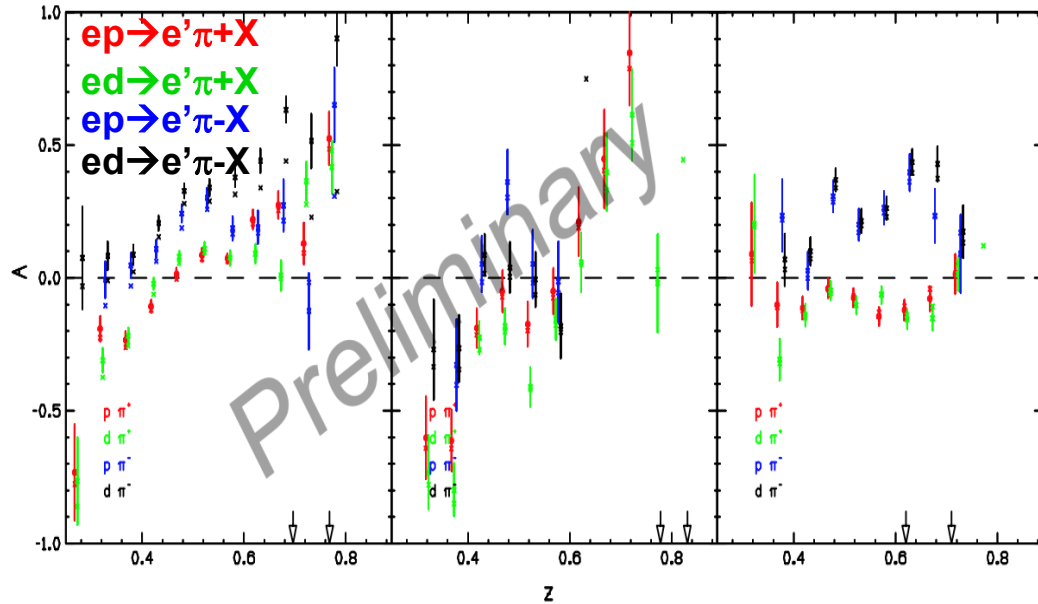
sin and cos azimuthal modulations from JLab

P. Bosted et al → Hall-C (ECT*-2022)

$x=0.31$ $Q^2=3.1$ GeV^2
VERY PRELIMINARY

$x=0.30$ $Q^2=4.1$ GeV^2
VERY PRELIMINARY

$x=0.45$ $Q^2=4.5$ GeV^2
VERY PRELIMINARY



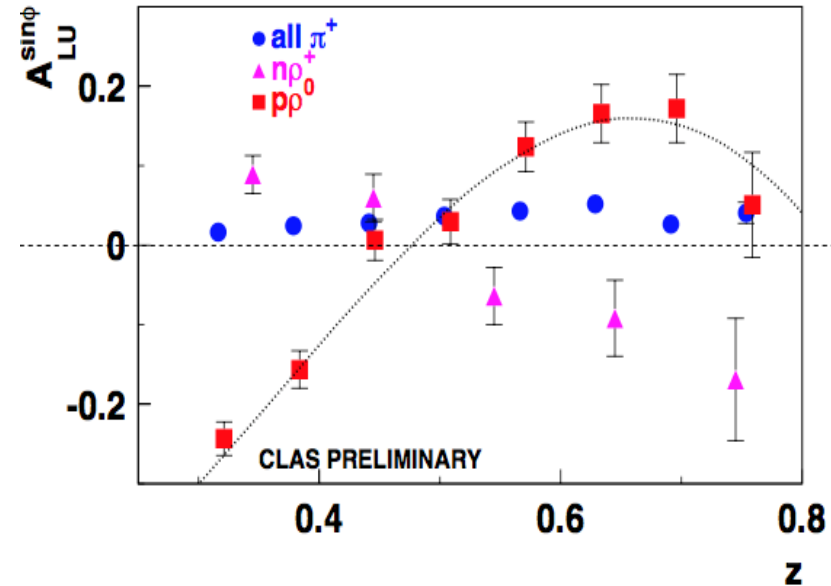
$$y = M_0 b e^{-b p_T^2} (1 + A p_T \cos \phi)$$

- "A" (cosine modulation) generally close to zero or positive.
- Cahn effect should give negative cosine

- Indication of dominant rho contributions, in particular for pions at low P_T , specially for π^- ,
- Understanding of the production mechanism is critical in understanding of QCD dynamics

H.A. → Hall-B (GDH-2004)

$e p \rightarrow e' \pi^+ X$ ($n\rho^+, p\rho^0$)



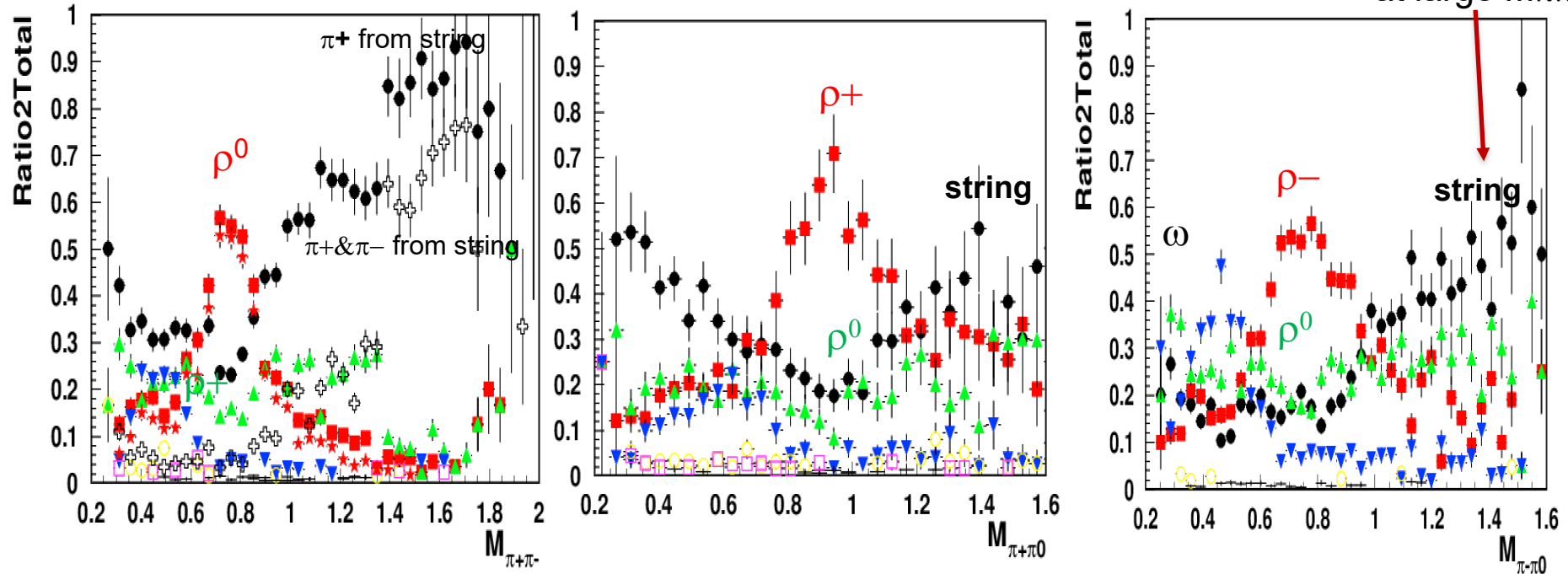
rho decay pions produce pions with SSA flipping the sign $\sim z=0.5$)

Sources of inclusive pions: CLAS12 MC

Detection of π^0 s allows also studies of

ρ^\pm

High P_T pions
at large $M_{\pi\pi}$



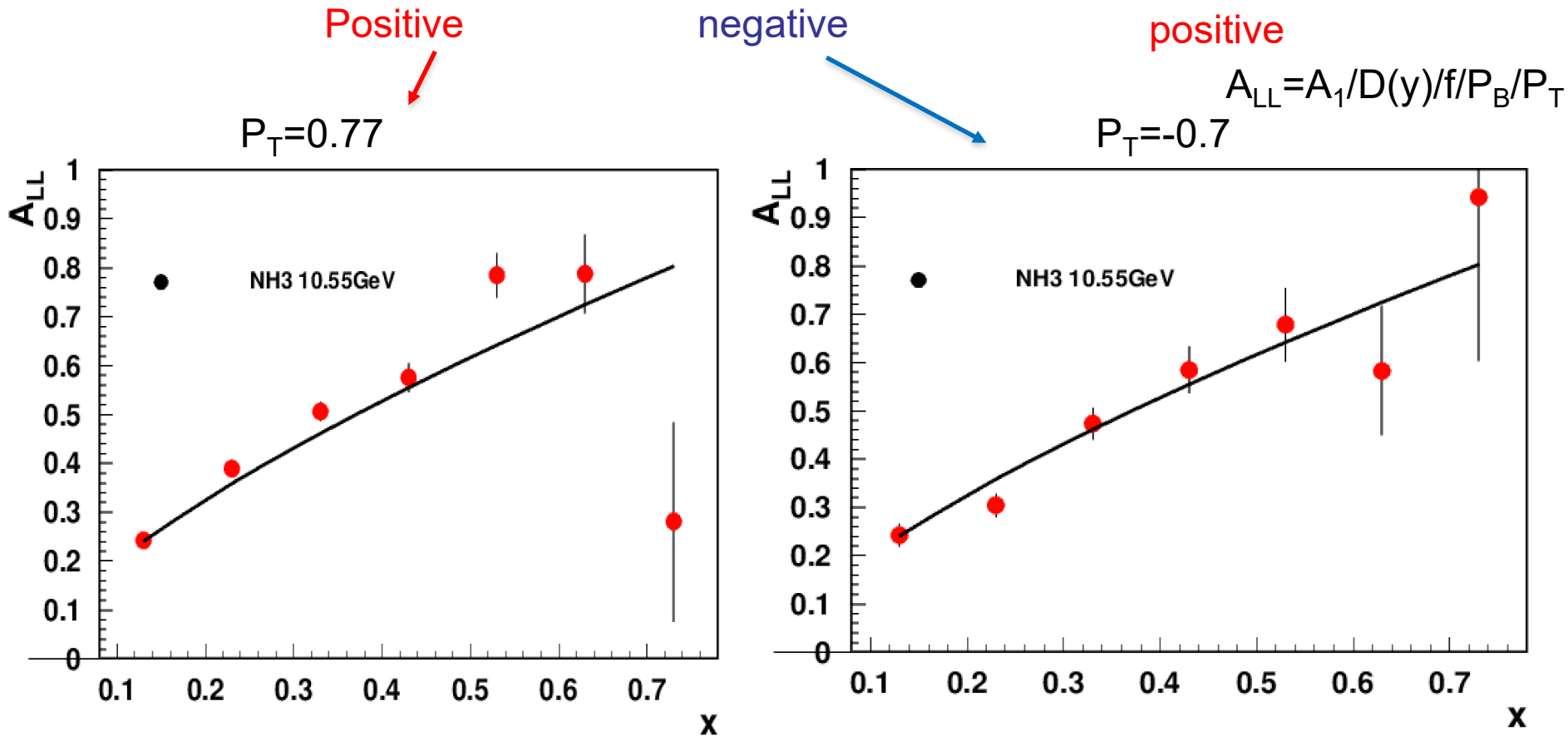
Dominant fraction of 2 pion combinations come from VM decays

- ρ
- string
- ω

All measured 2 pion combinations are dominated by VM decays, indicate that all inclusive pions are dominated by VM decays at small P_T s, and in particular at lower z !!!

Target A_{LL} in $ep \rightarrow e'X$ (RGC)

NH3 - FTout: 16983, 16990, 16992, 17001, 17015, 17019, 17076, 17088, 17092



$P_b=0.83$, $f=0.15+0.075*x$ (from MC), the line $x^{0.7}$

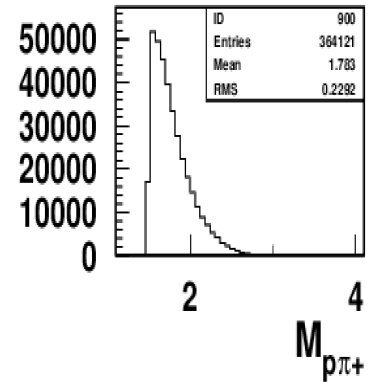
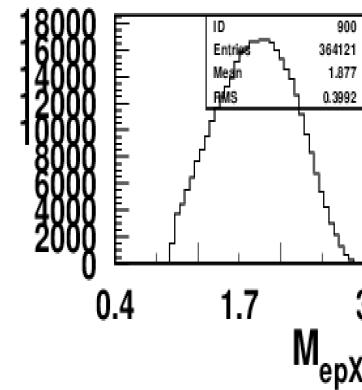
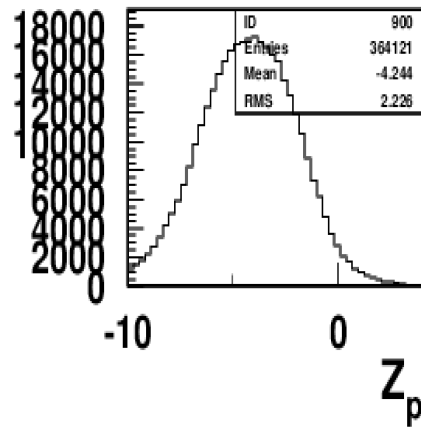
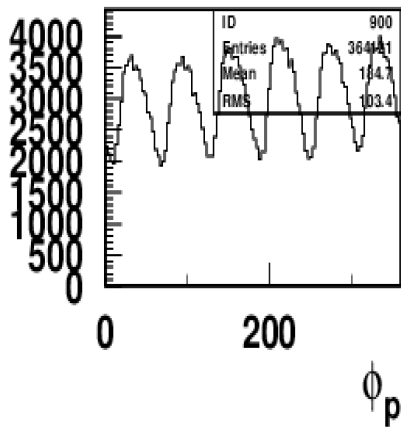
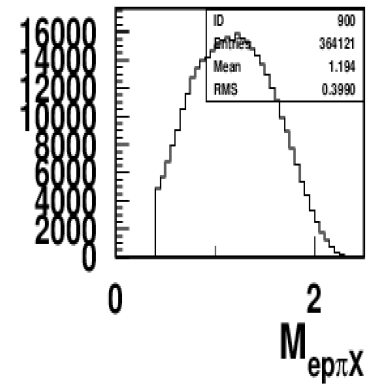
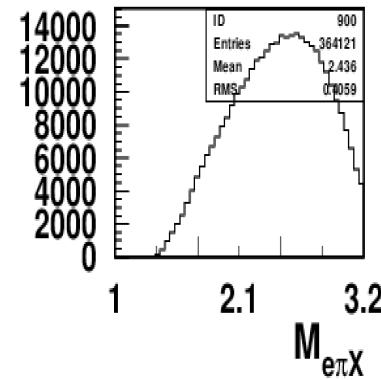
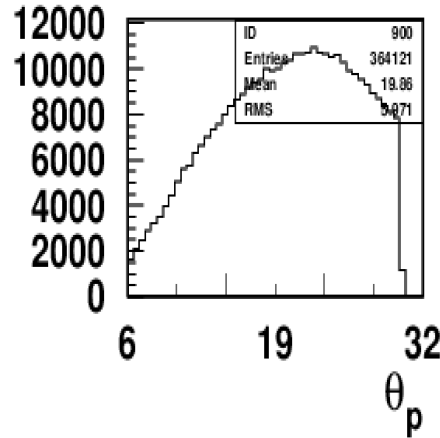
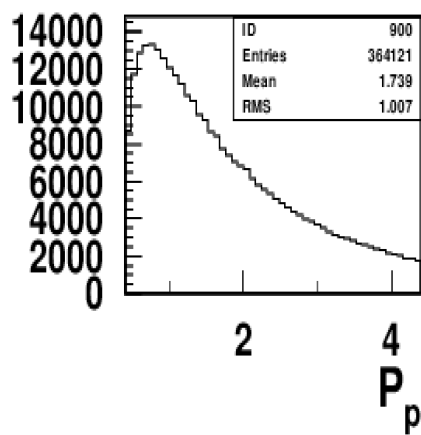
RGC NH3, Ftout: $e p \rightarrow e' p \pi + X$ (RGC)

NH3 - Ftout: 16983, 16990, 16992, 17001, 17015, 17019, 17076, 17088, 17092

Positive

negative

positive



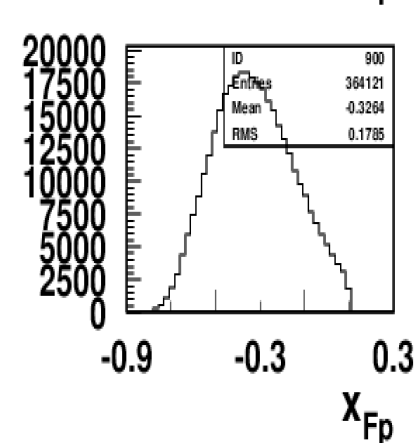
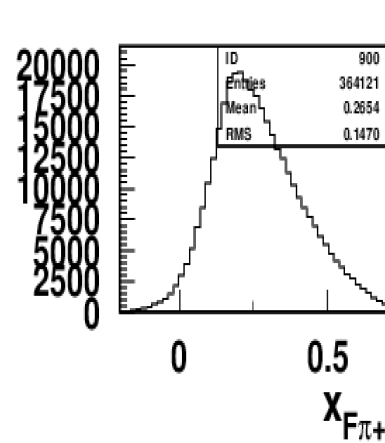
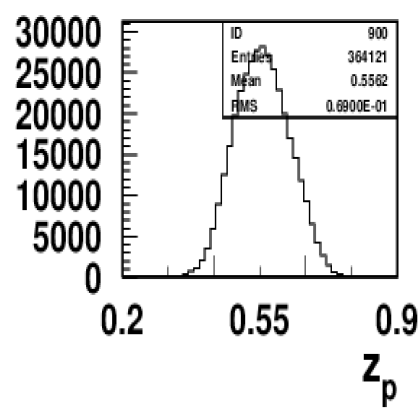
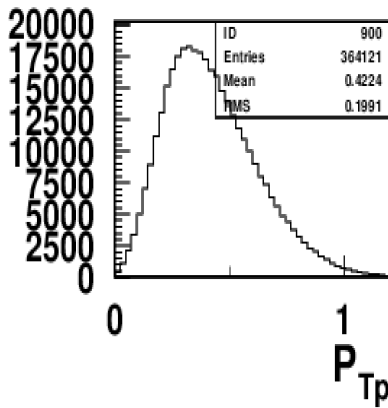
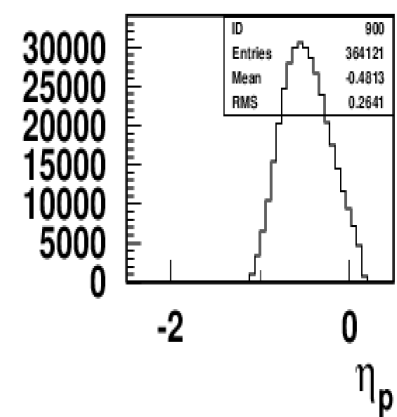
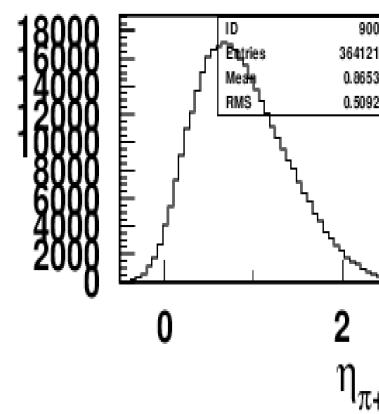
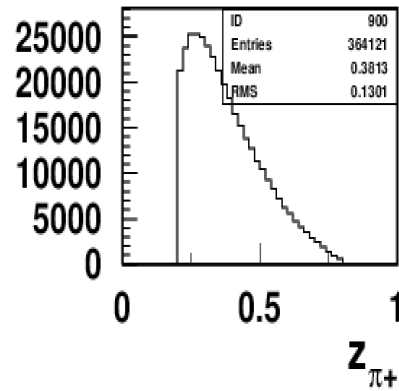
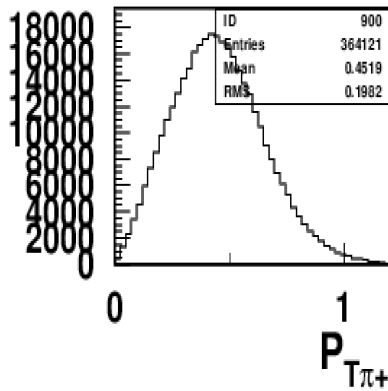
RGC NH3, Ftout: $e p \rightarrow e' p \pi + X$ (RGC)

NH3 - Ftout: 16983, 16990, 16992, 17001, 17015, 17019, 17076, 17088, 17092

Positive

negative

positive



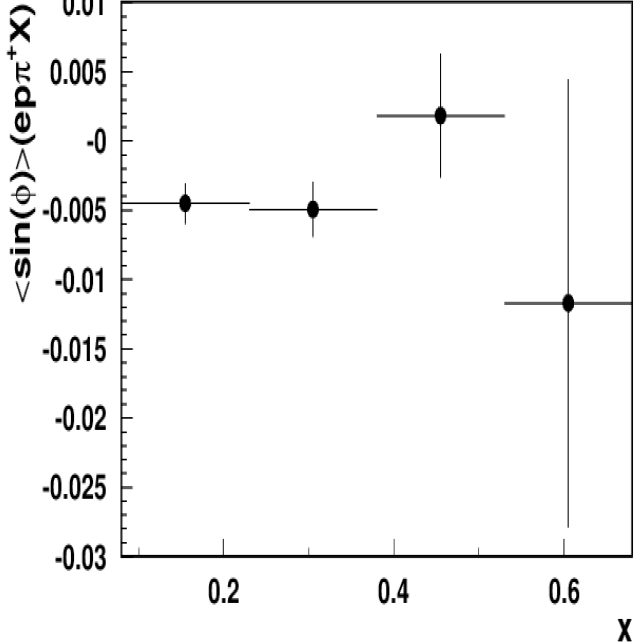
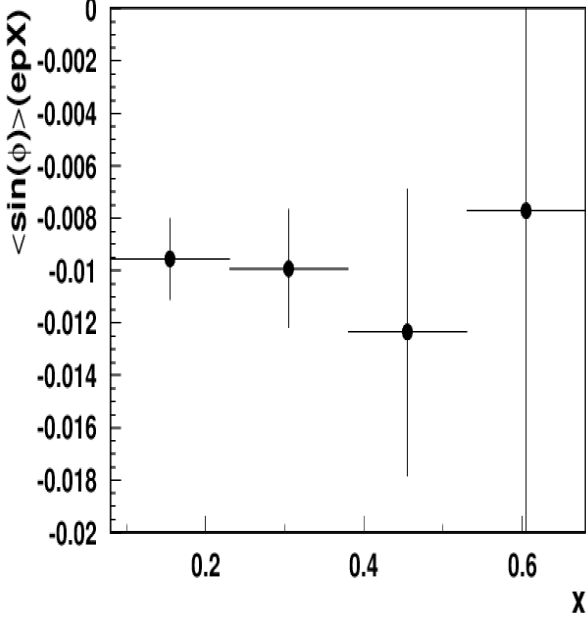
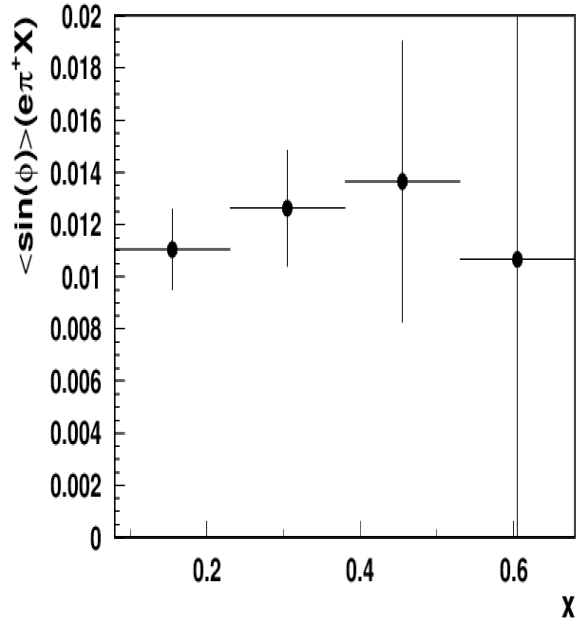
Beam SSAs in $ep \rightarrow e' p \pi + X$ (RGC)

NH3 - FTout: 16983, 16990, 16992, 17001, 17015, 17019, 17076, 17088, 17092

Positive

negative

positive



Helicity of the lepton is correct, reproducing signs of all observed beam SSAs

Target A_{LL} in $ep \rightarrow e'X$ (RGC)

NH3 - FTout: 16983, 16990, 16992, 17001, 17015, 17019, 17076, 17088, 17092

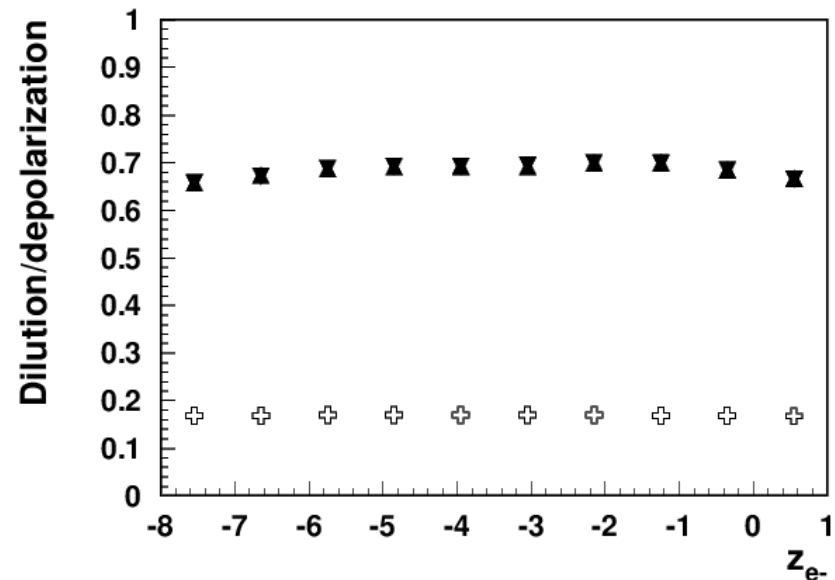
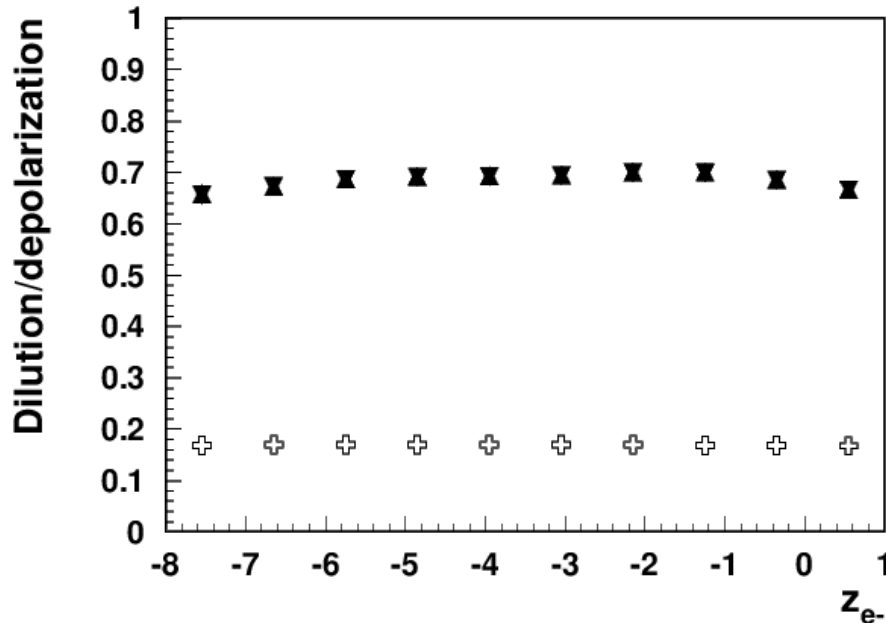
Positive

negative

positive

$P_T=0.77$

$A_{LL}=A_1/D(y)/f/P_B/P_T$
 $P_T=-0.7$



$P_b=0.83$, $f=0.15+0.075*x$ (from MC), the line $x^{0.7}$

Target A_{LL} in $ep \rightarrow e'X$ (RGC)

NH3 - FTout: 16983, 16990, 16992, 17001, 17015, 17019, 17076, 17088, 17092

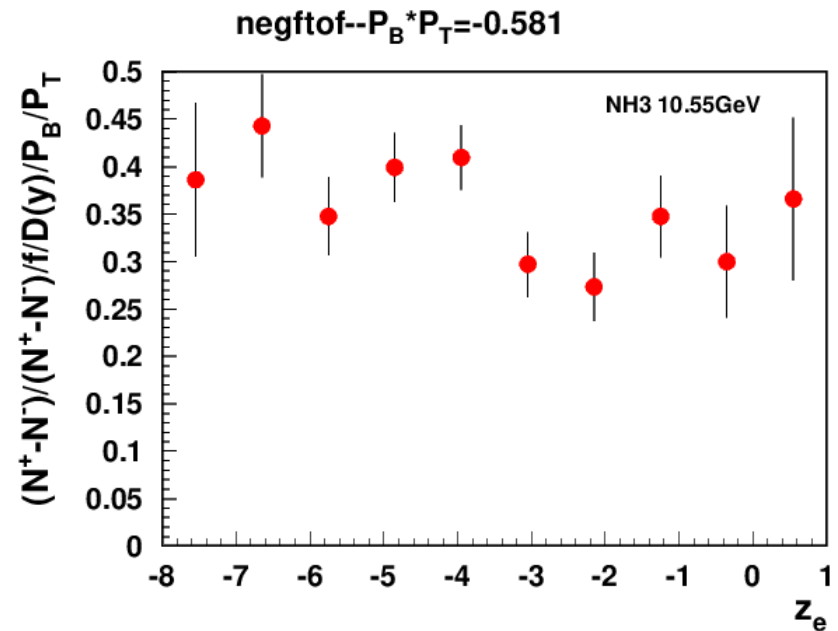
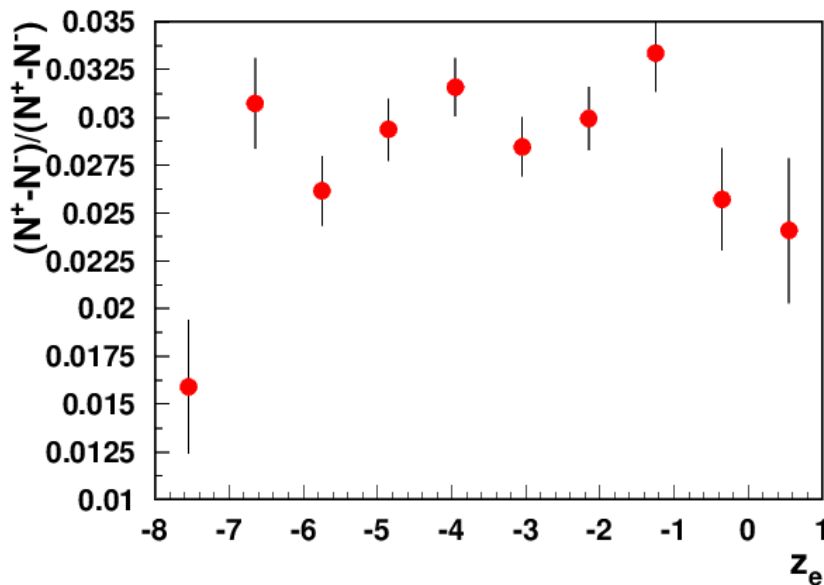
Positive

negative

positive

$P_T=0.77$

$A_{LL}=A_1/D(y)/f/P_B/P_T$
 $P_T=-0.7$



$P_b=0.83$, $f=0.15+0.075*x$ (from MC), the line $x^{0.7}$

Target A_{LL} in $ep \rightarrow e'X$ (RGC)

NH3 - FTin: 16681 16683 16717 16720 16768 16770 16772 16726 16728 16741 16746 16750 16759 16762

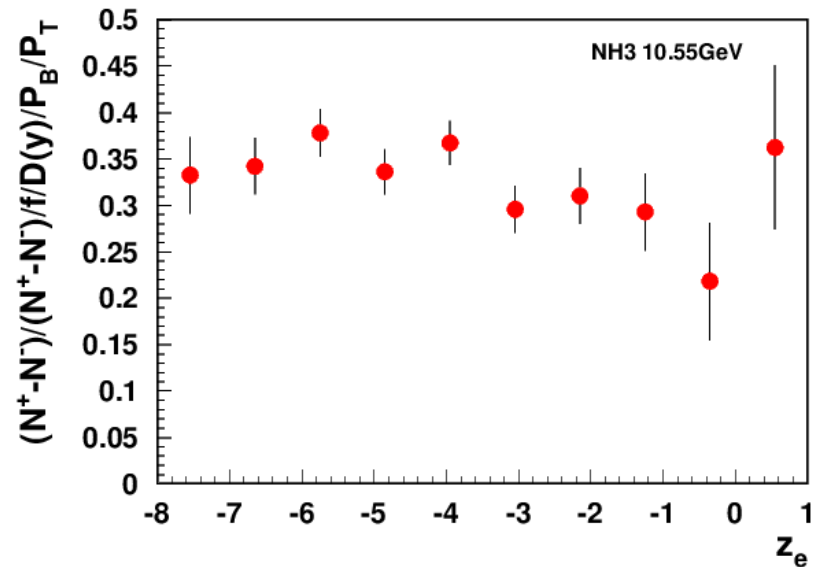
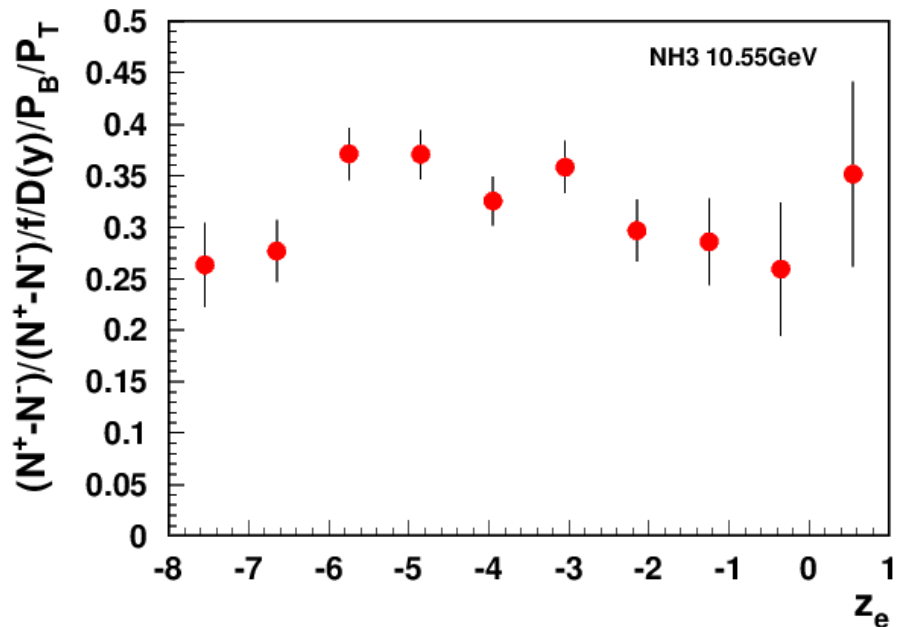
Positive

negative

$P_T=0.77$

$P_T=-0.7$

$$A_{LL} = A_1 / D(y) / f / P_B / P_T$$



$P_b=0.83$, $f=0.15+0.075*x$ (from MC), the line $x^{0.7}$

Kinematics

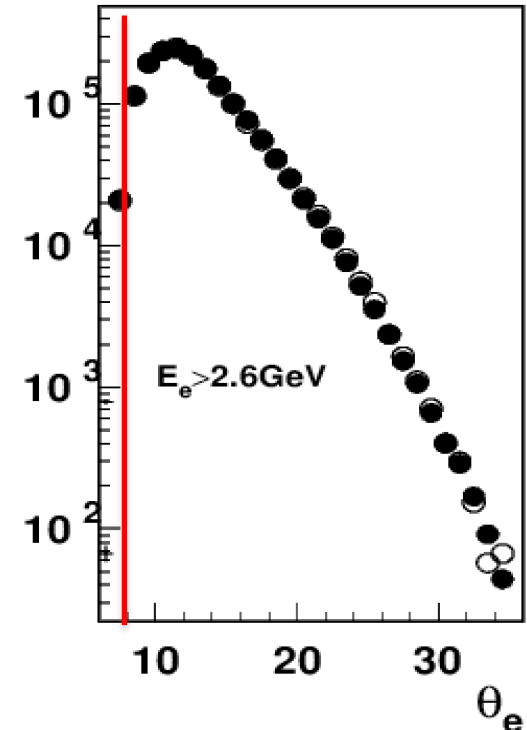
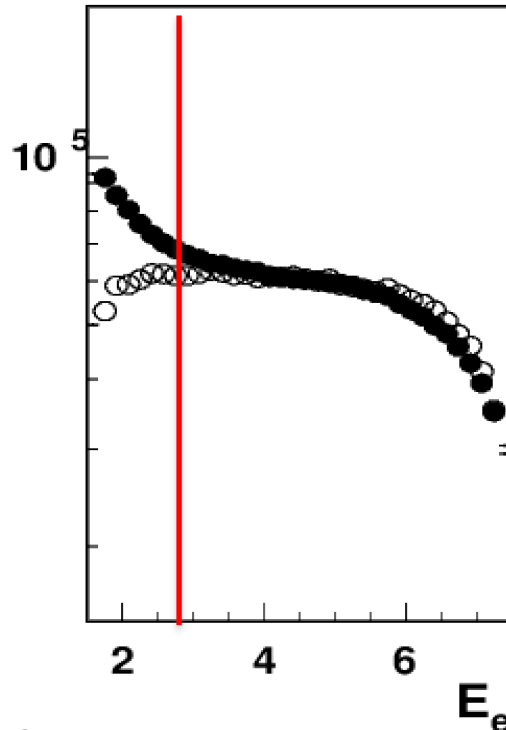
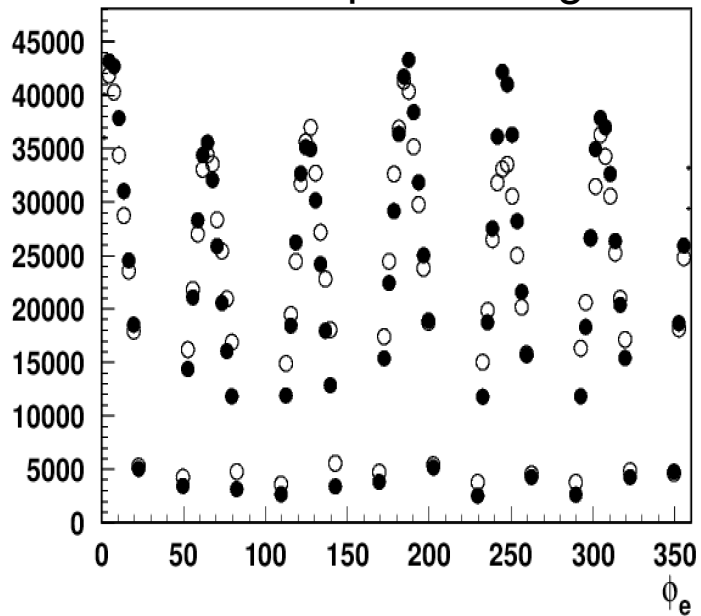
MC files with SIDIS events

[/work/ceba24gev/sidis/reconstructed/polarized-plus-10.5GeV-neutron/hipo/](#)
[/work/ceba24gev/sidis/reconstructed/polarized-minus-10.5GeV-neutron/hipo/](#)
[/work/ceba24gev/sidis/reconstructed/polarized-plus-10.5GeV-proton/hipo/](#)
[/work/ceba24gev/sidis/reconstructed/polarized-minus-10.5GeV-proton/hipo/](#)

Data from [/volatile/clas12/rg-c/production/ana_data/dst/train/sidisdvcs/](#)

Event selection The same cuts applied to MC and data $|z+3.5| < 4$, $E_e > 2.6$, $35 > \theta > 8$

Remove acceptance edges



Within fiducial region MC and data are consistent

Target A_{LL} in $eNH3 \rightarrow e'X$ (RGC-sol1)

NH3 : 17215

17191 17195

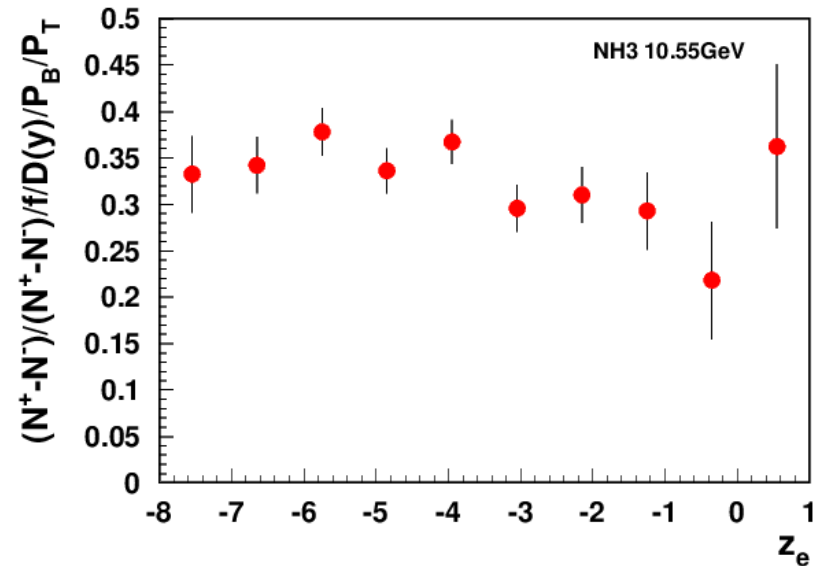
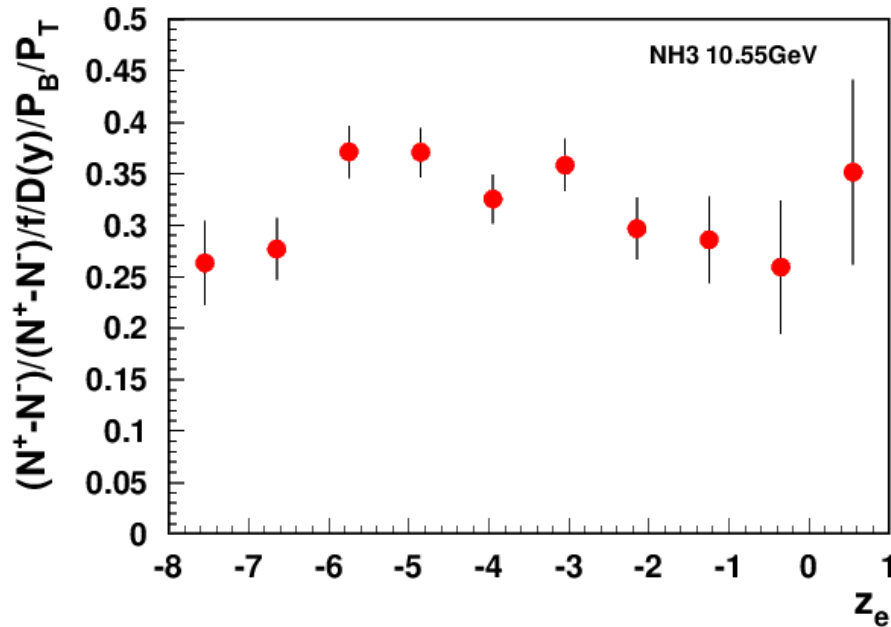
Positive

negative

$P_T=0.77$

$P_T=-0.7$

$$A_{LL} = A_1/D(y)/f/P_B/P_T$$



$P_b=0.83, f=0.15+0.075*x$ (from MC), the line $x^{0.7}$

Target A_{LL} in $eND3 \rightarrow e'X$ (RGC)

ND3: 17265 17269, 17273

17227 17237 17262

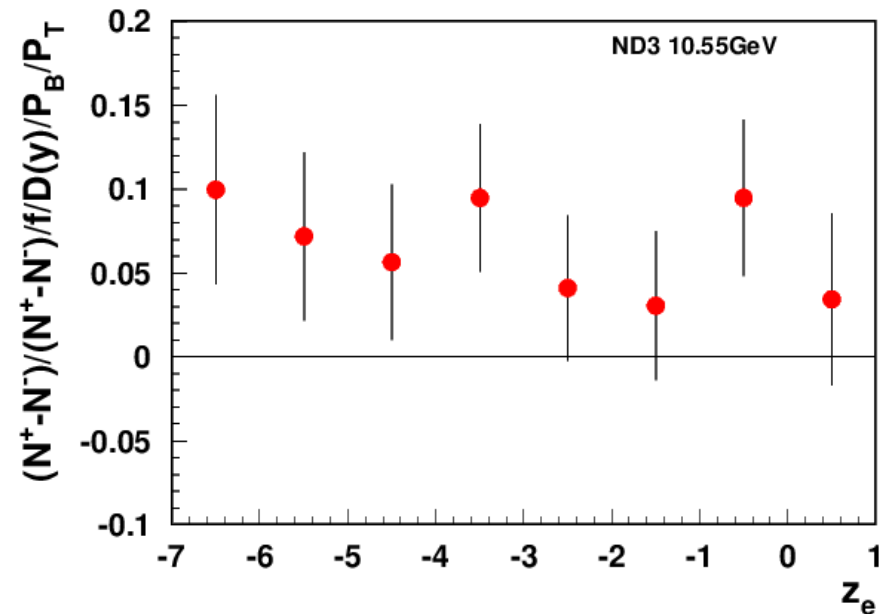
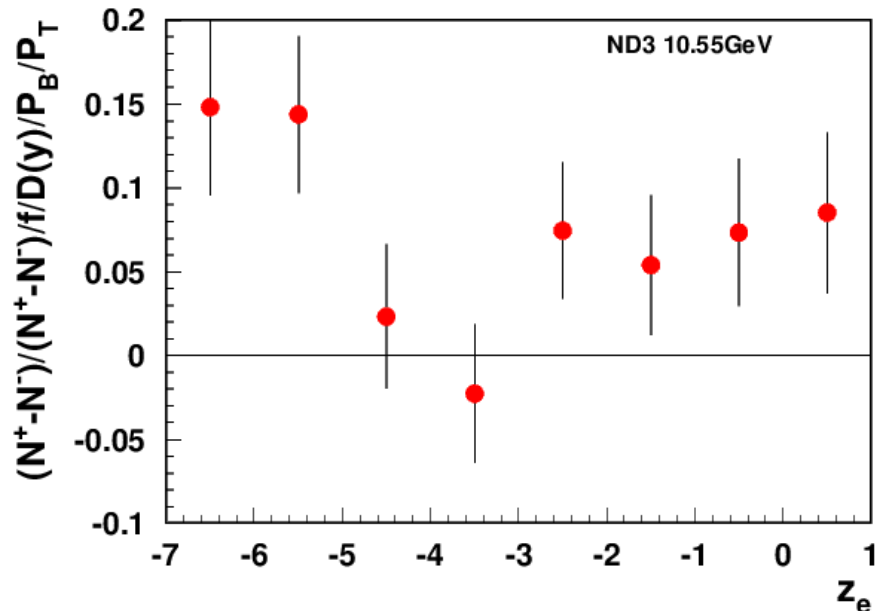
Positive

negative

$P_T=0.3$

$P_T=-0.3$

$$A_{LL} = A_1/D(y)/f/P_B/P_T$$

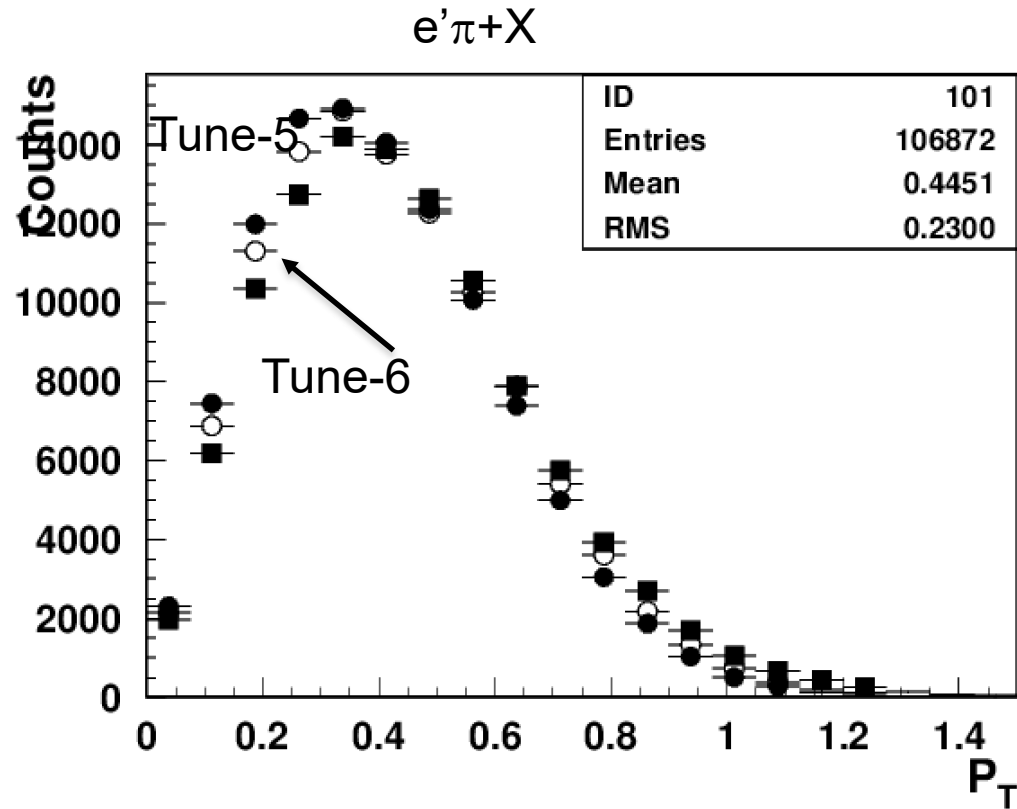
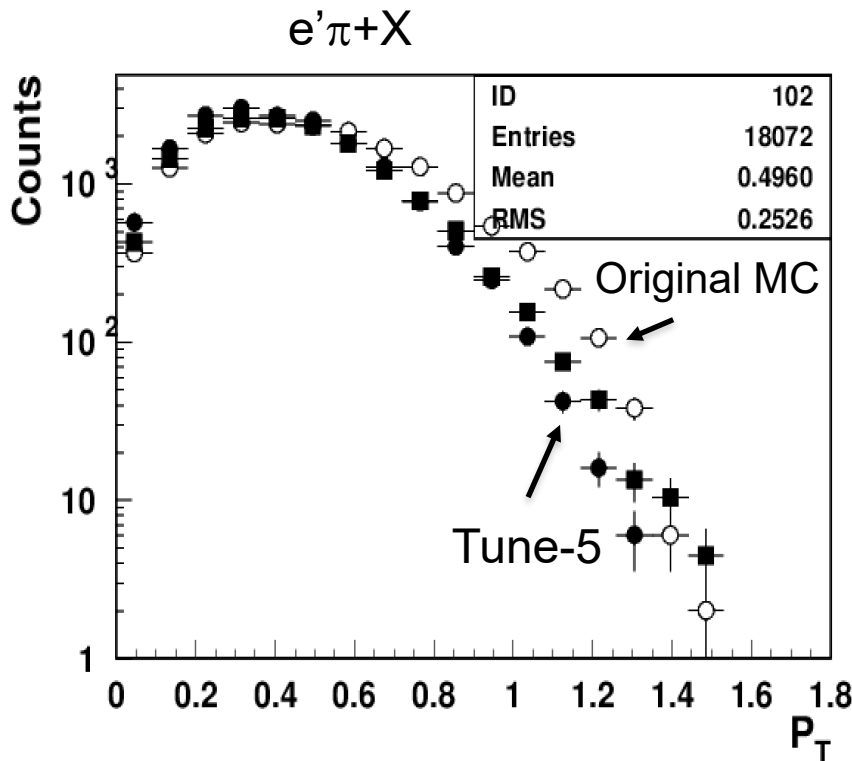


$P_b=0.83, f=0.84*(0.305-0.01*x)$ (from MC)

Adding nuclear background

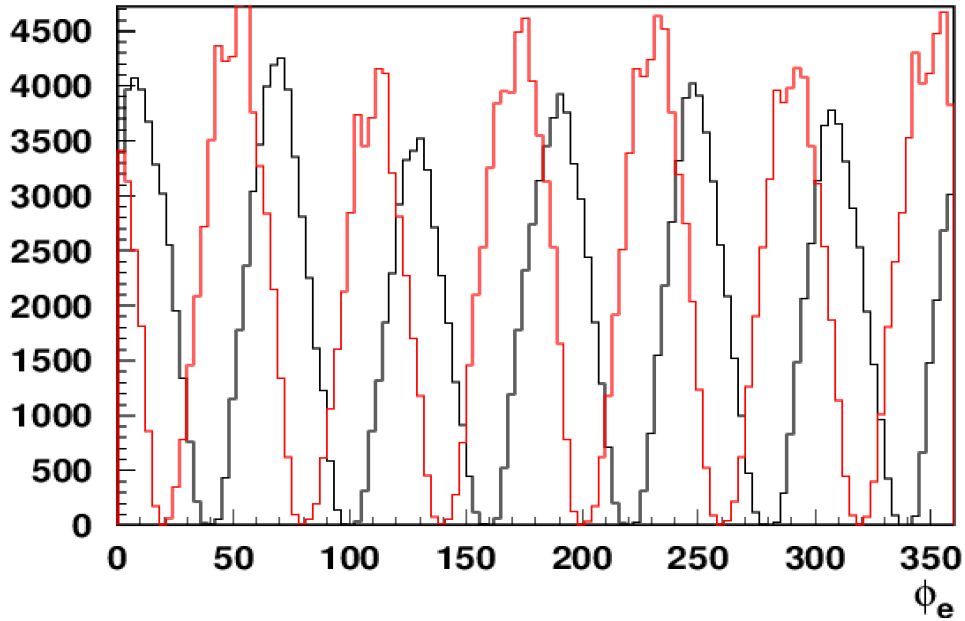
- Comparing nuclear MC (A. Alaoui/L. El Fassi) with RGC carbon

Squares RGC carbon run#16128

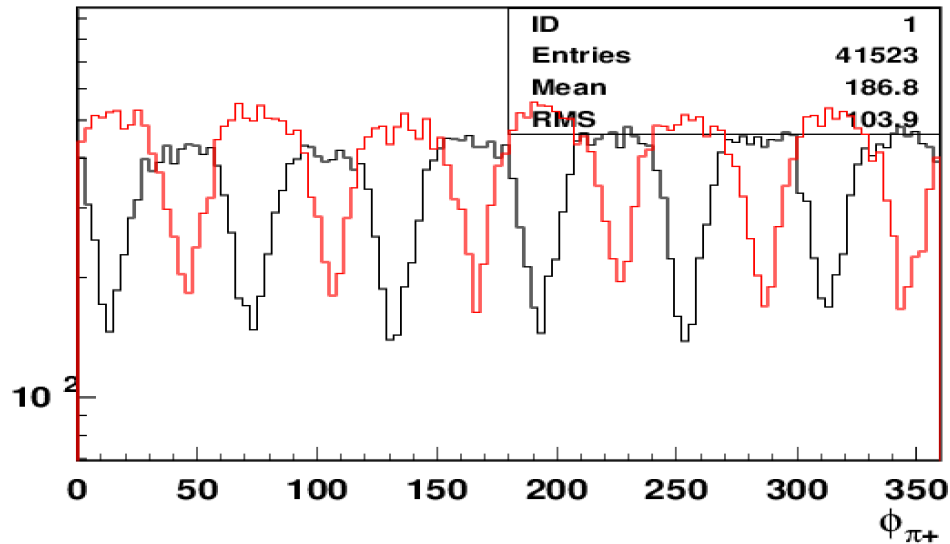


Nuclear MC (PYTHIA based) has been tuned in several iterations to get closer to RGC carbon data (main changes: use the same input as for clasdis with enhanced VM fractions)
Will need accounting of radiative corrections for fine tuning

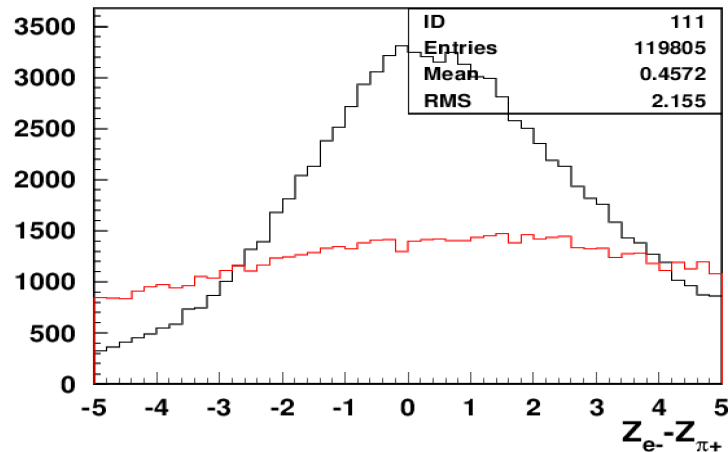
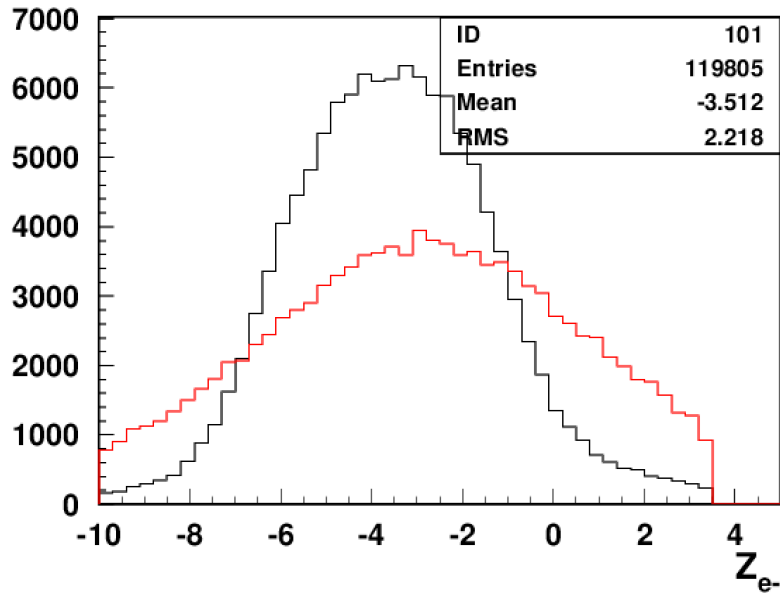
Azimuthal kicks



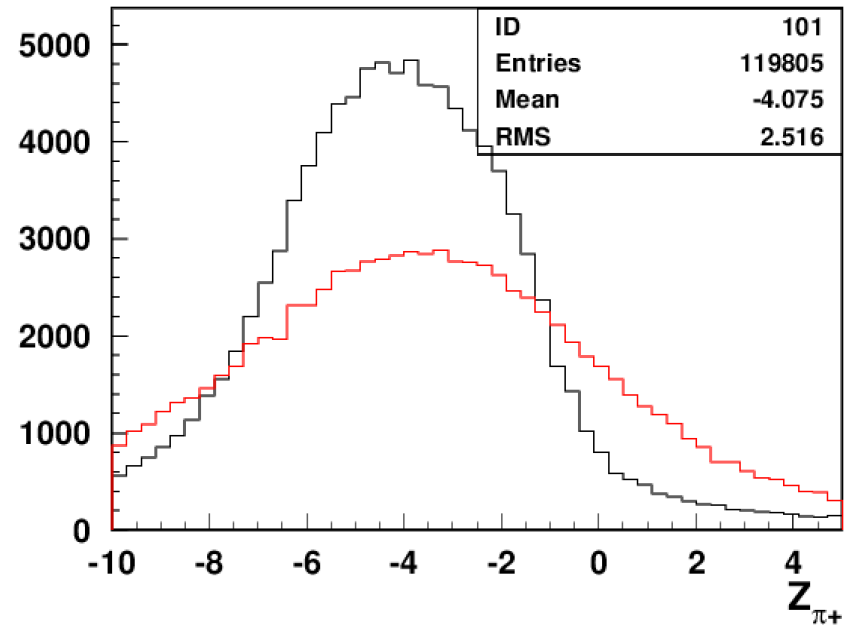
Red-inverted solenoid



Vertexes



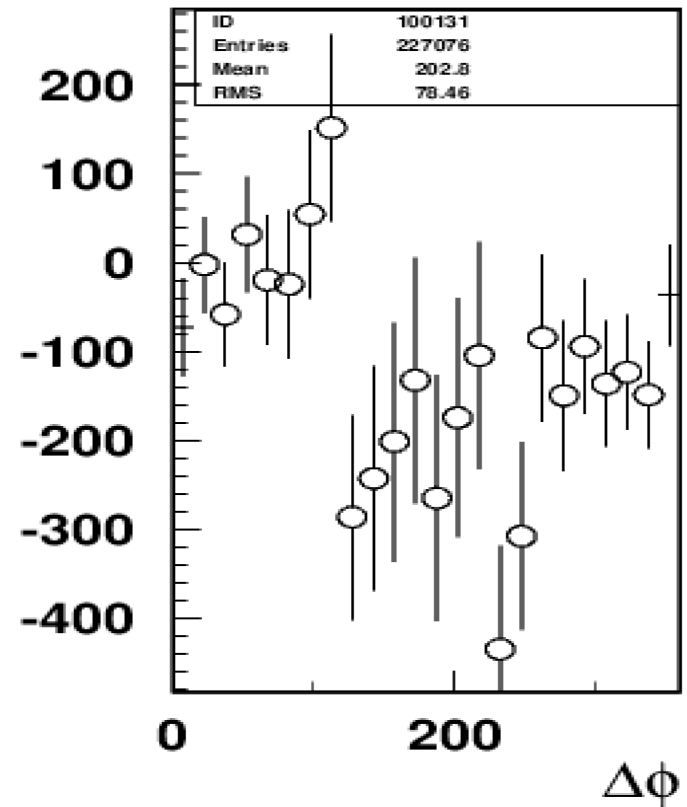
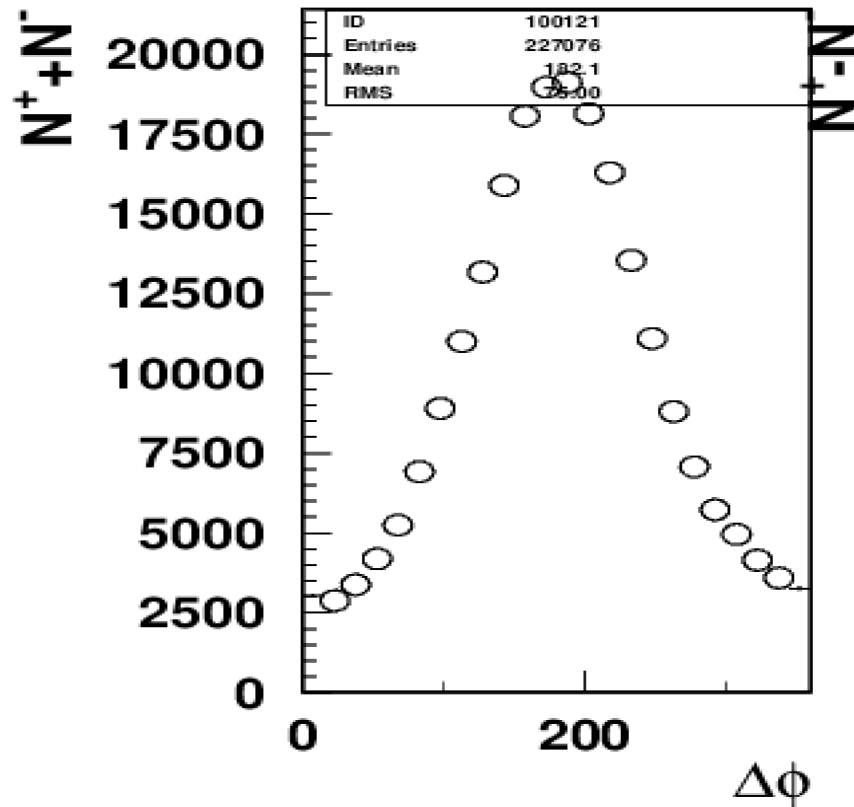
Red-inverted solenoid



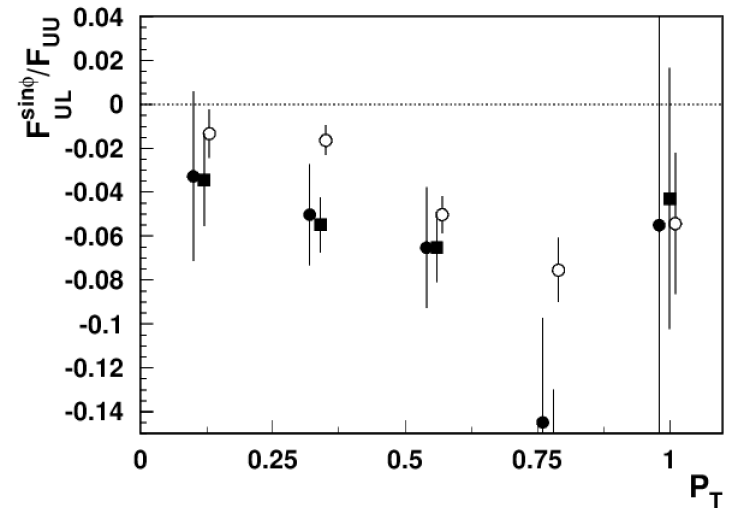
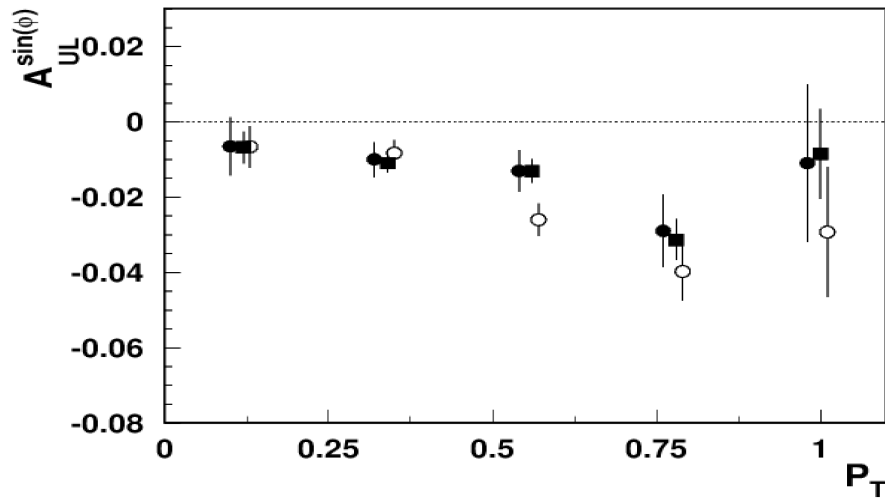
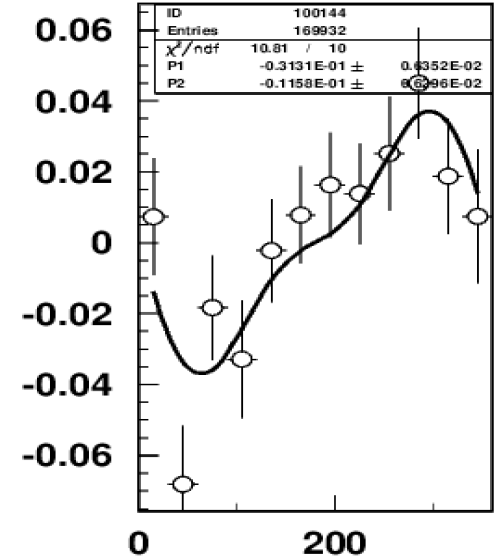
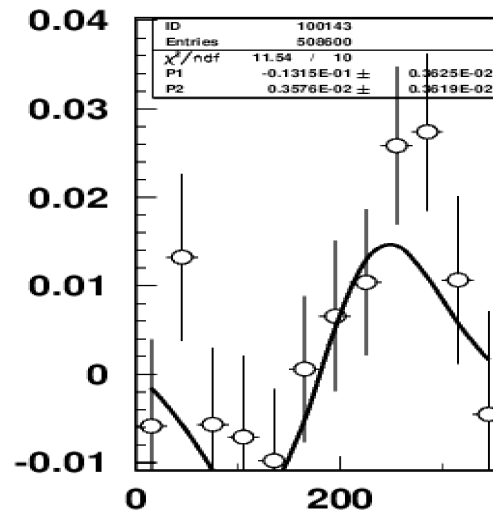
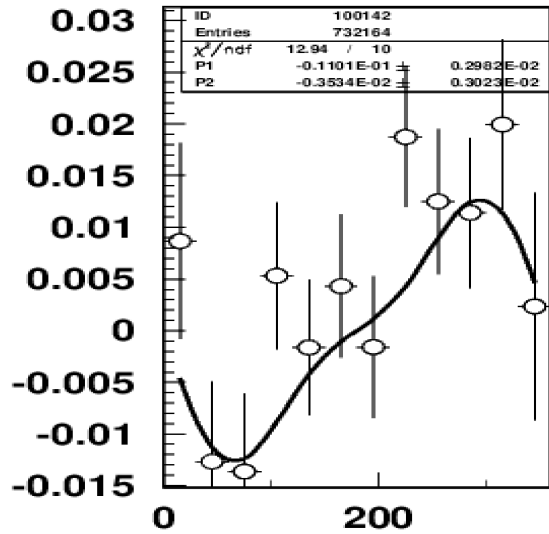
Bigger spread in z!
HBT vs TBT (will compare with full tracking)
May need calibration

Beam SSA in $ep \rightarrow e' p \pi + X$

Significant disbalance in helicities, not critical for SSAs, but will be important for DSAs



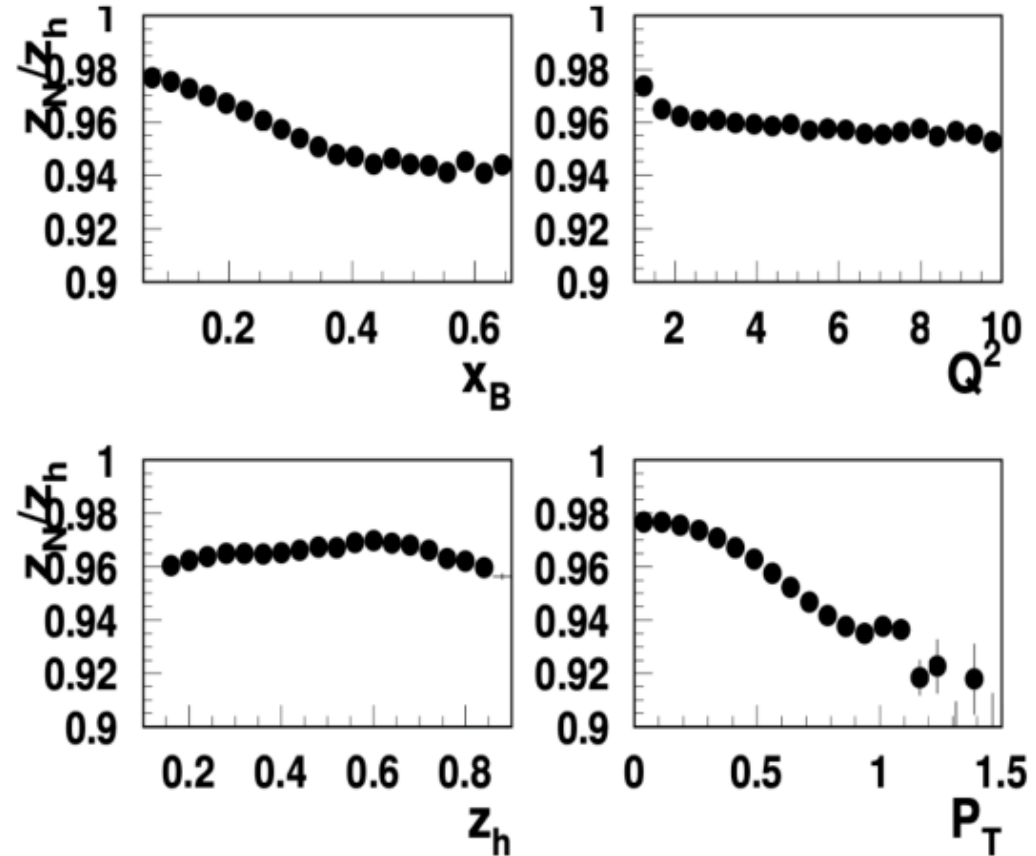
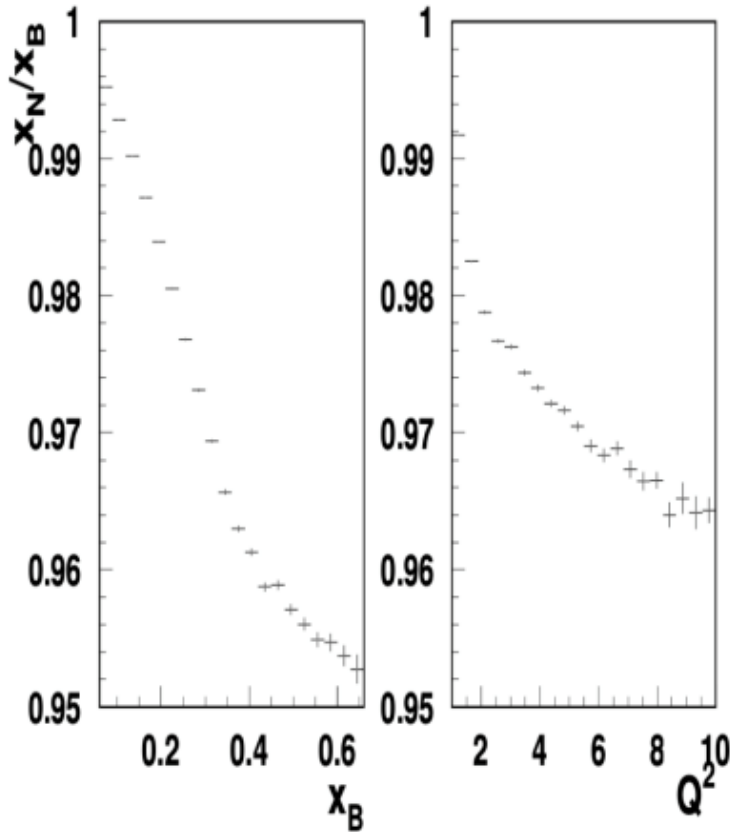
Target SSA in $ep \rightarrow e' p \pi + X$



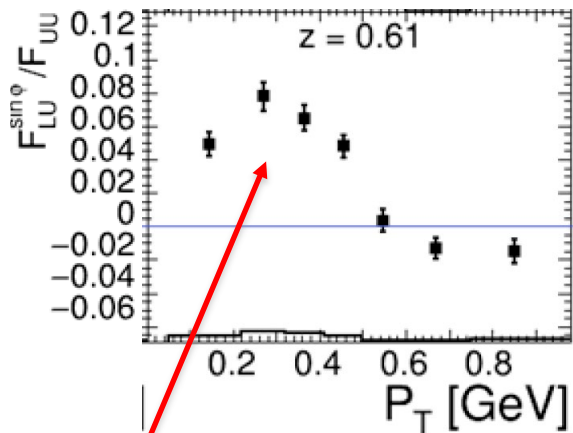
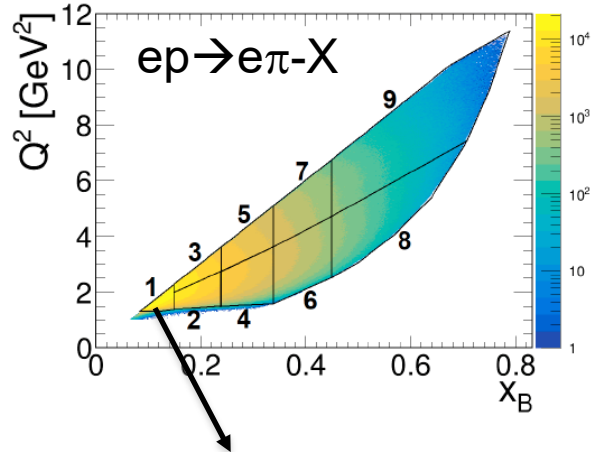
CLAS12: Mass corrections to x_B and z_h

$$x_N = -\frac{q^+}{P^+} = \frac{2x_{Bj}}{1 + \sqrt{1 + \frac{4x_{Bj}^2 M^2}{Q^2}}}$$

$$z_N = \frac{Q^4 x_N z_h \left(1 \pm \sqrt{1 - \frac{4M^2 M_B^2 x_{Bj}^2 (Q^4 + x_N^2 M^2 q_T^2)}{Q^8 z_h^2}} \right)}{2x_{Bj} (Q^4 + x_N^2 M^2 q_T^2)}$$

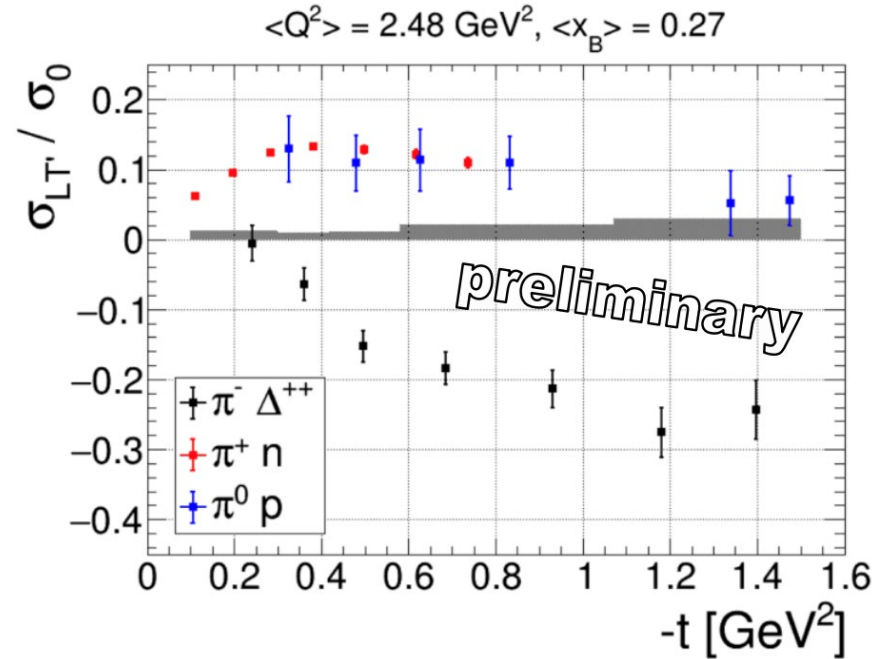


Multidimensional measurements



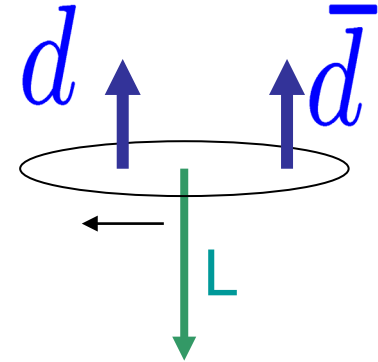
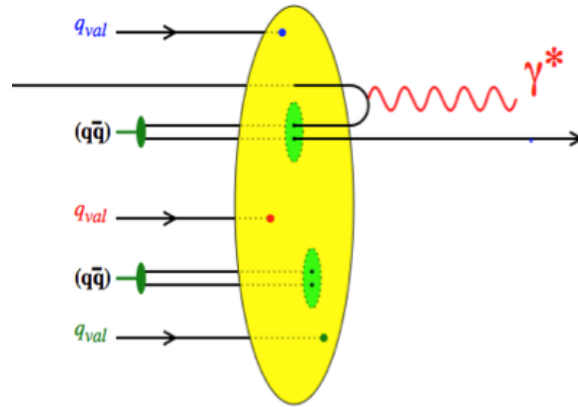
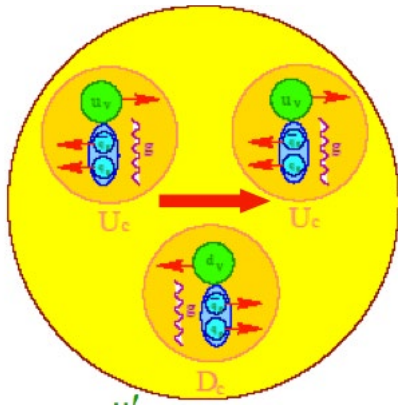
Low P_T dominated by VM decays from u-quarks

Beam SSAs as a tool to access the longitudinally polarized quarks



π^- has sensitivity to polarized d-quarks, but require multidimensional measurements

Correlations between target and current



- how the remnant system dresses itself up to become a full-fledged hadron
- correlation with the spin of the target or/and the produced particles