

"hep-ph" contributions

Q



TALKS WITH GENERAL SCOPE

MALTONI

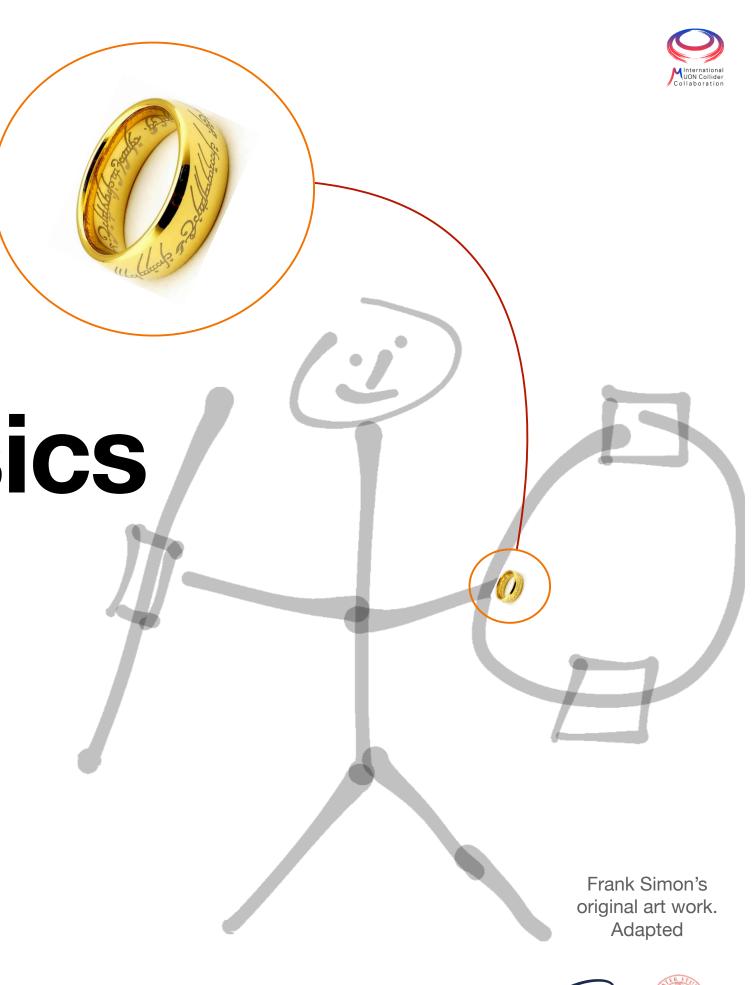
"One ring to rule them all, one ring to find them" ... - J.R.R.Tolkien

Muon collider: physics

Fabio Maltoni Université catholique de Louvain Università di Bologna



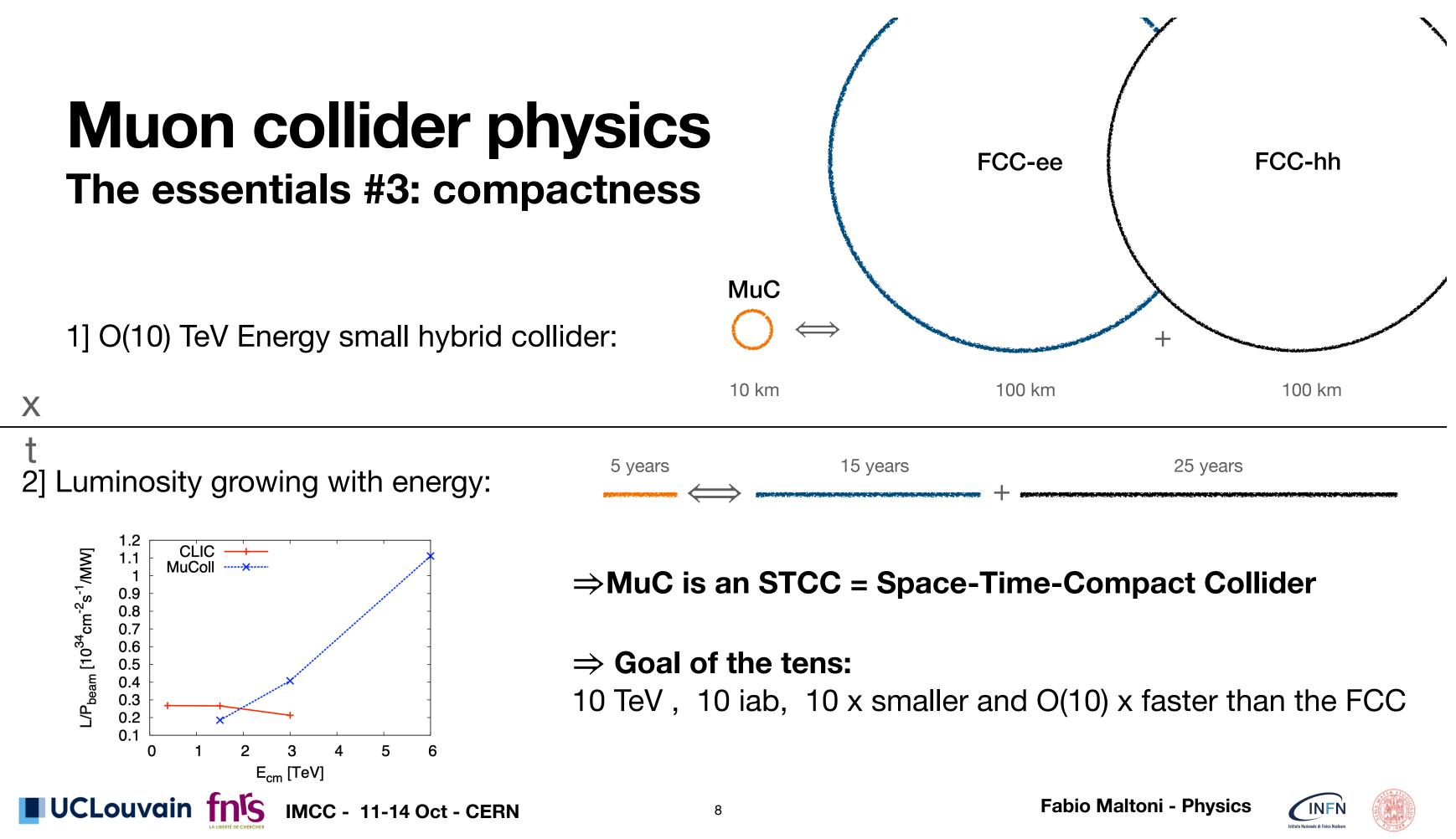
Muon collider collaboration Annual Meeting 11-14 Oct - CERN





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MALTONI



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Muon collider physics The essentials

- •A O(10 TeV) muC is in the range of what could be technically achievable. R&D is necessary.
- •It would radically change the way we do collider physics, opening the exploration of EW phenomena at higher scales through an hybrid direct/indirect approach in a clean environment.
- •Given what we know now from the LHC + what will learn from HL-LHC what are the muC physics drivers?















Reports of the working groups: Physics

30/7-018 - Kjell Johnsen Auditorium, CERN

Physics Summary

Andrea Wulzer





First Collaboration Meeting (Oct. 11-14 2022)

Andrea Wulzer

08:50 - 09:05

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WULZER

Reports of the working groups: Physics

30/7-018 - Kjell Johnsen Auditorium, CERN

Dark Matter [Xiaoran Zhao]

- mono-X,di-X and DT for low mass region $M_{\chi} < \frac{\sqrt{s}}{2}$
- Indirect probes are good at thresholds $M_{\chi} \sim \frac{\sqrt{s}}{2}$
- and can probe high mass $M_{\chi} > rac{\sqrt{s}}{2}$
- Soft/collinear radiations shift NC to NC+CC
- Hard radiations affect the dynamics and sensitivities
- Statistic uncertainties in $\mathcal{O}(0.1 \sim 1\%)$ level: need further improvements on theoretical predictions(NLO+NLL or higher?)

	$3 { m TeV}$			$10 { m TeV}$			30 TeV		
	DL	$e^{\mathrm{DL}}-1$	$\operatorname{SL}(\frac{\pi}{2})$	DL	$e^{\mathrm{DL}}-1$	$\operatorname{SL}(\frac{\pi}{2})$	DL	$e^{\mathrm{DL}}-1$	$\operatorname{SL}(\frac{\pi}{2})$
$\ell_L \to \ell'_L$	-0.46	-0.37	0.25	-0.82	-0.56	0.33	-1.23	-0.71	0.41
$\ell_L \to q_L$	-0.44	-0.36	0.25	-0.78	-0.54	0.34	-1.18	-0.69	0.42
$\ell_L \to e_R$	-0.32	-0.27	0.13	-0.56	-0.43	0.17	-0.85	-0.57	0.21
$\ell_L \to u_R$	-0.27	-0.24	0.11	-0.48	-0.38	0.15	-0.72	-0.51	0.18
$\ell_L \to d_R$	-0.24	-0.21	0.10	-0.43	-0.35	0.13	-0.64	-0.47	0.16
$\ell_R \to \ell'_L$	-0.32	-0.27	0.13	-0.56	-0.43	0.17	-0.85	-0.57	0.21
$\ell_R \to q_L$	-0.30	-0.26	0.12	-0.53	-0.41	0.16	-0.79	-0.55	0.21
$\ell_R \to \ell'_R$	-0.17	-0.16	0.07	-0.30	-0.26	0.09	-0.46	-0.37	0.12
$\ell_R \to u_R$	-0.12	-0.12	0.05	-0.22	-0.20	0.07	-0.33	-0.28	0.08
$\ell_R \to d_R$	-0.09	-0.09	0.04	-0.17	-0.16	0.05	-0.25	-0.22	0.06

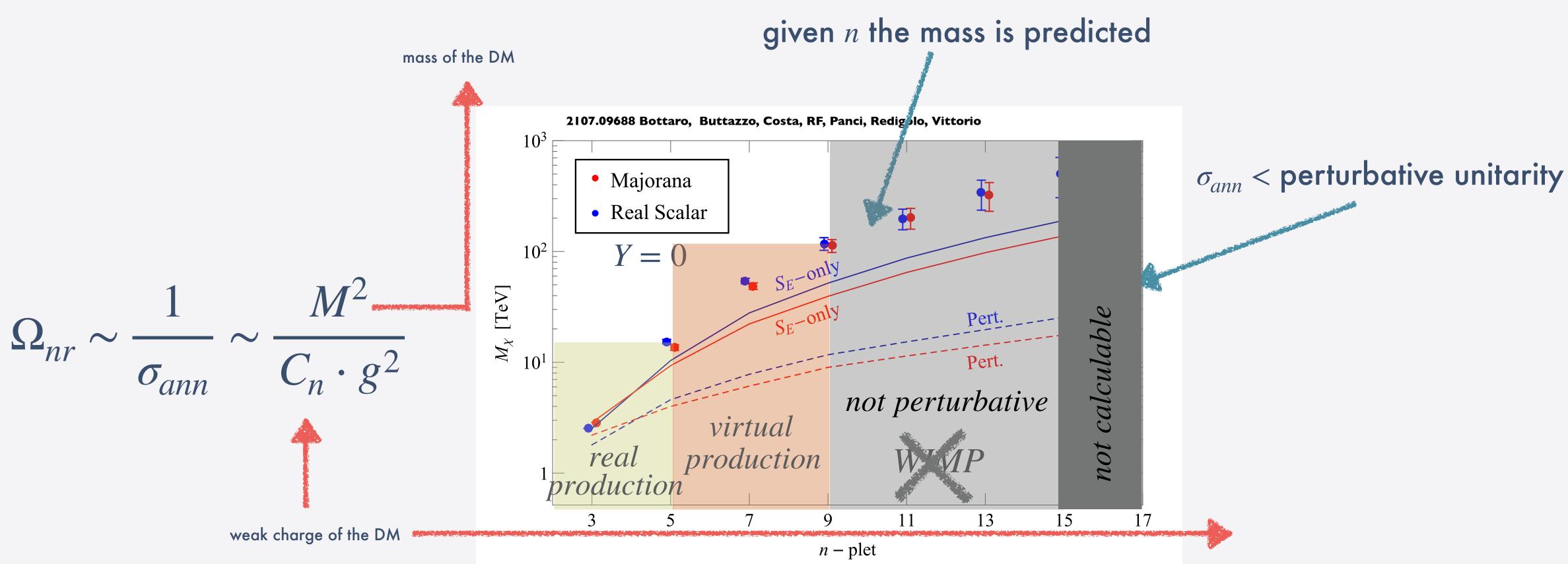
Andrea Wulzer

08:50 - 09:05

Note: we are currently not able to make sufficiently accurate predictions for this study (and many others)



AN "INTERPOLATOR" MODEL

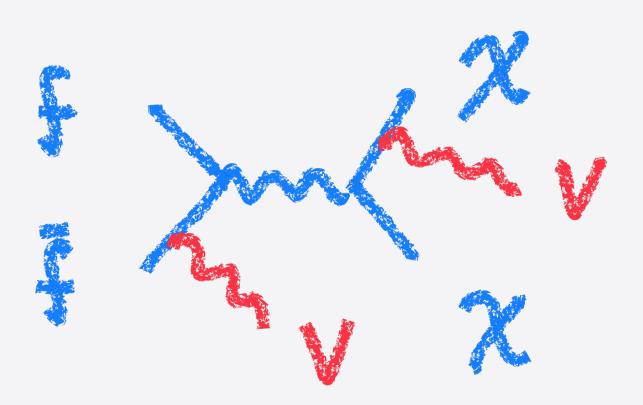


If Dark Matter feels SM weak interactions we can use the general *n*-plet WIMP to measure how well we are able to test this hypothesis and possibly discover or exclude one or several or the whole category of DM candidates.

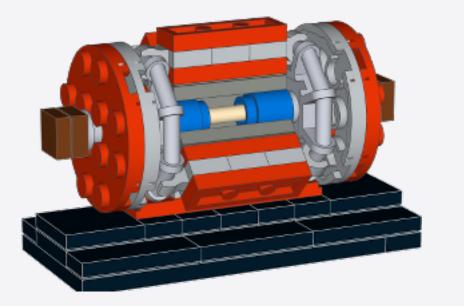




DIRECT PRODUCTION

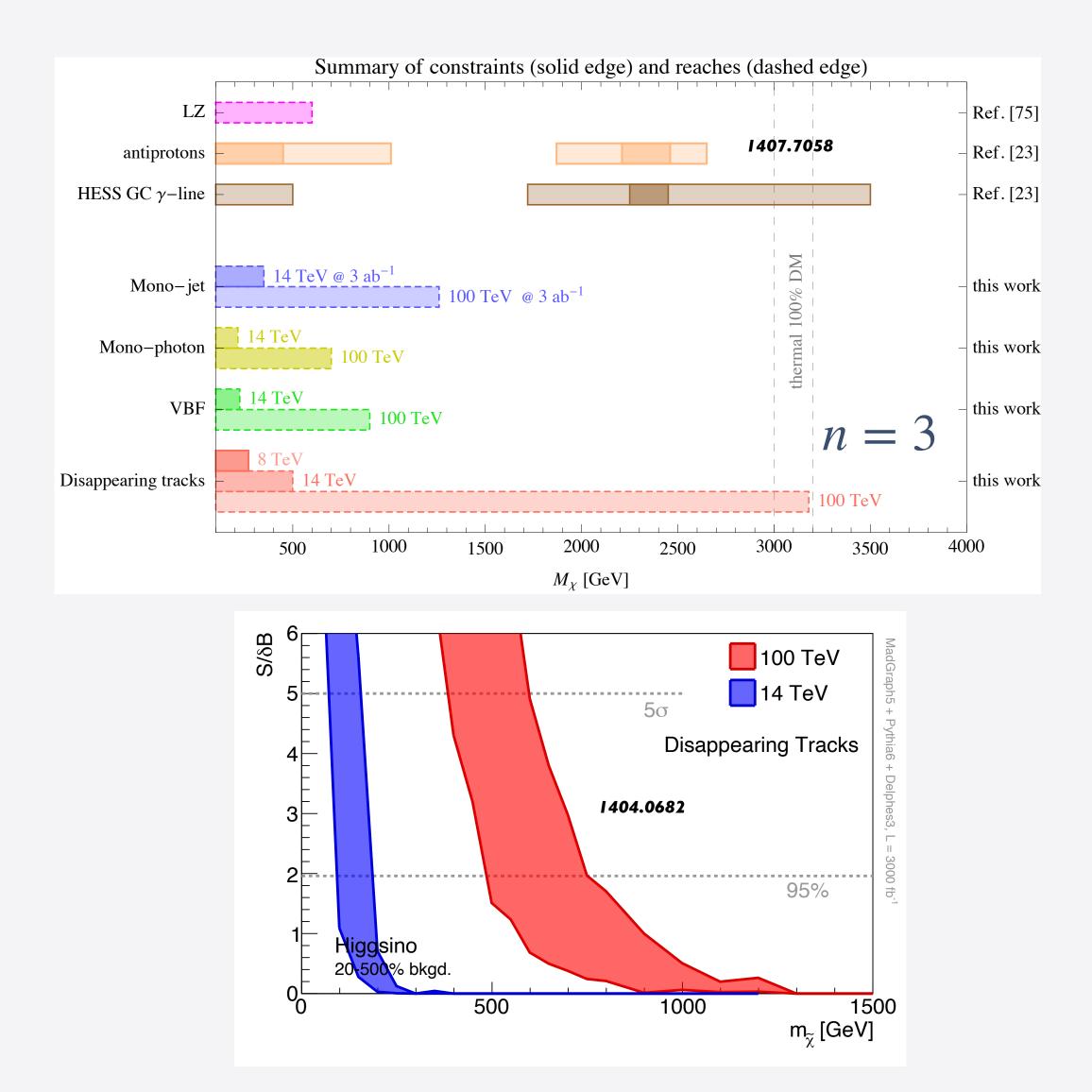


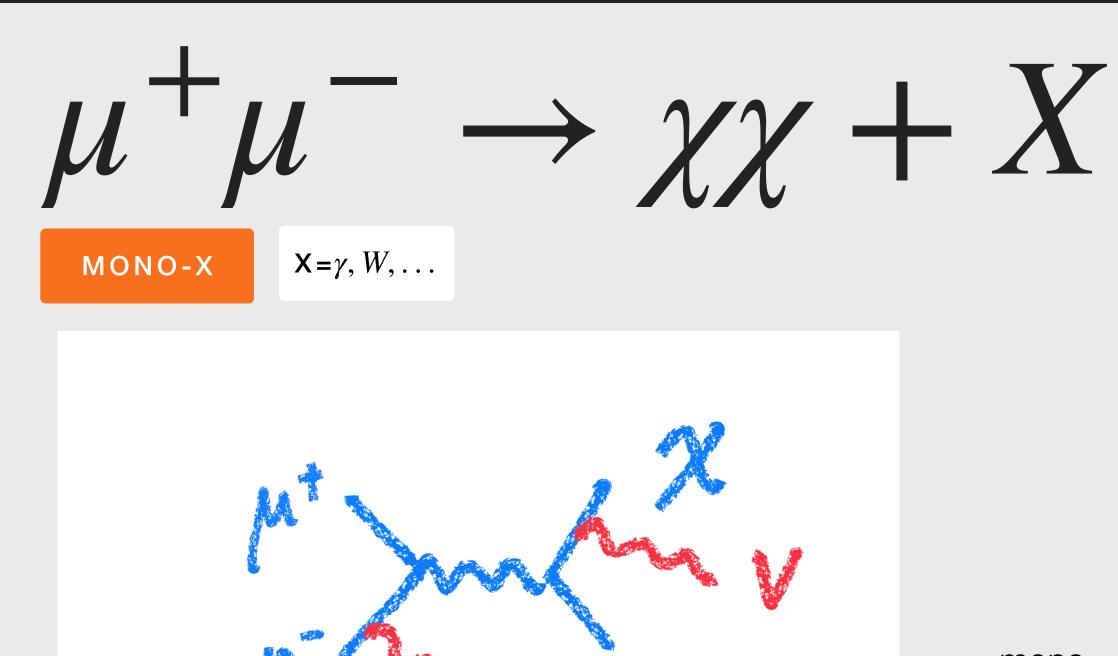
Production of Dark Matter weak multiplet states and observation of the decay products or associated productions









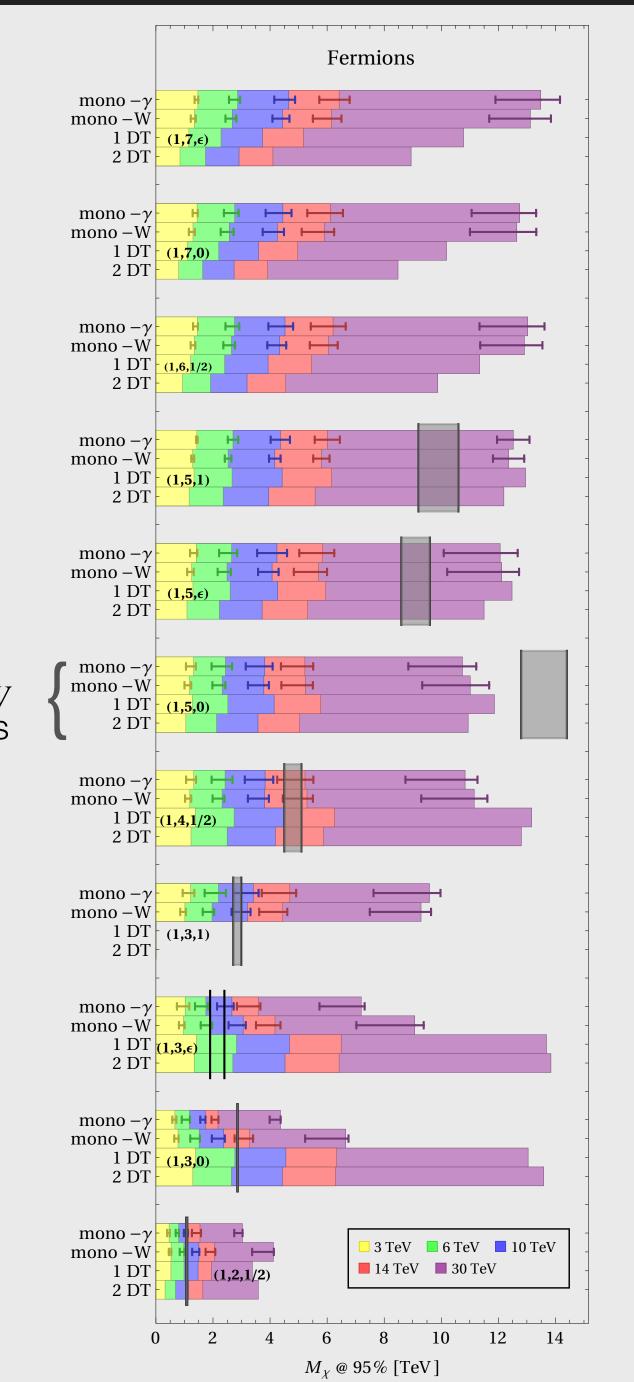


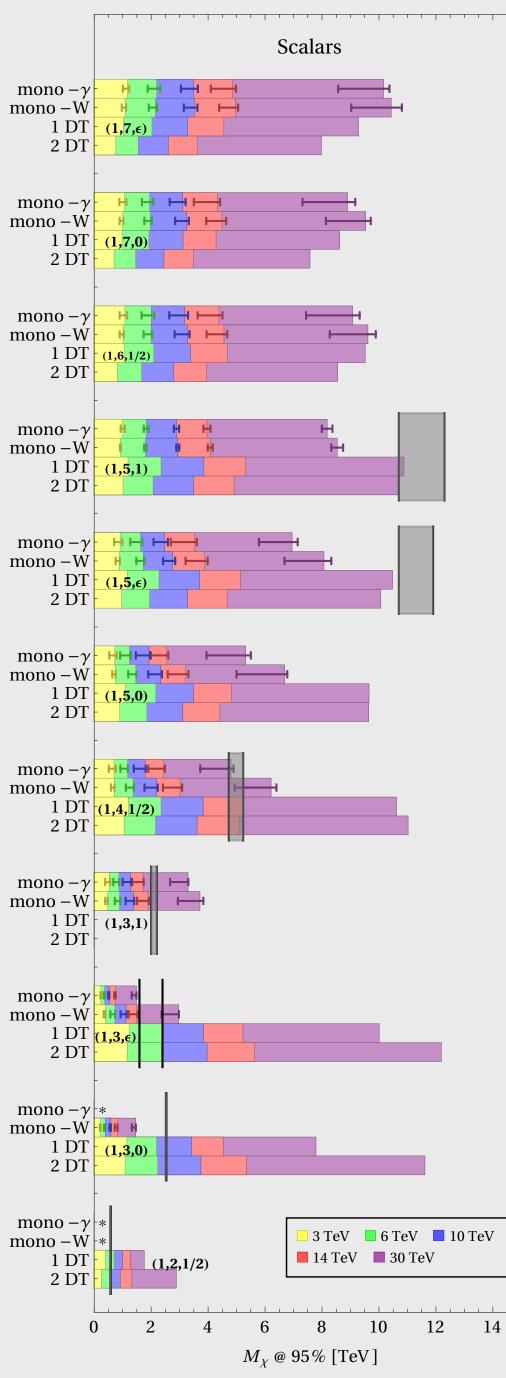
Large χ mass requires CoM energy!

Weak radiation yield the most constraining channel "mono-W"

mono-γ mono-W tracklets

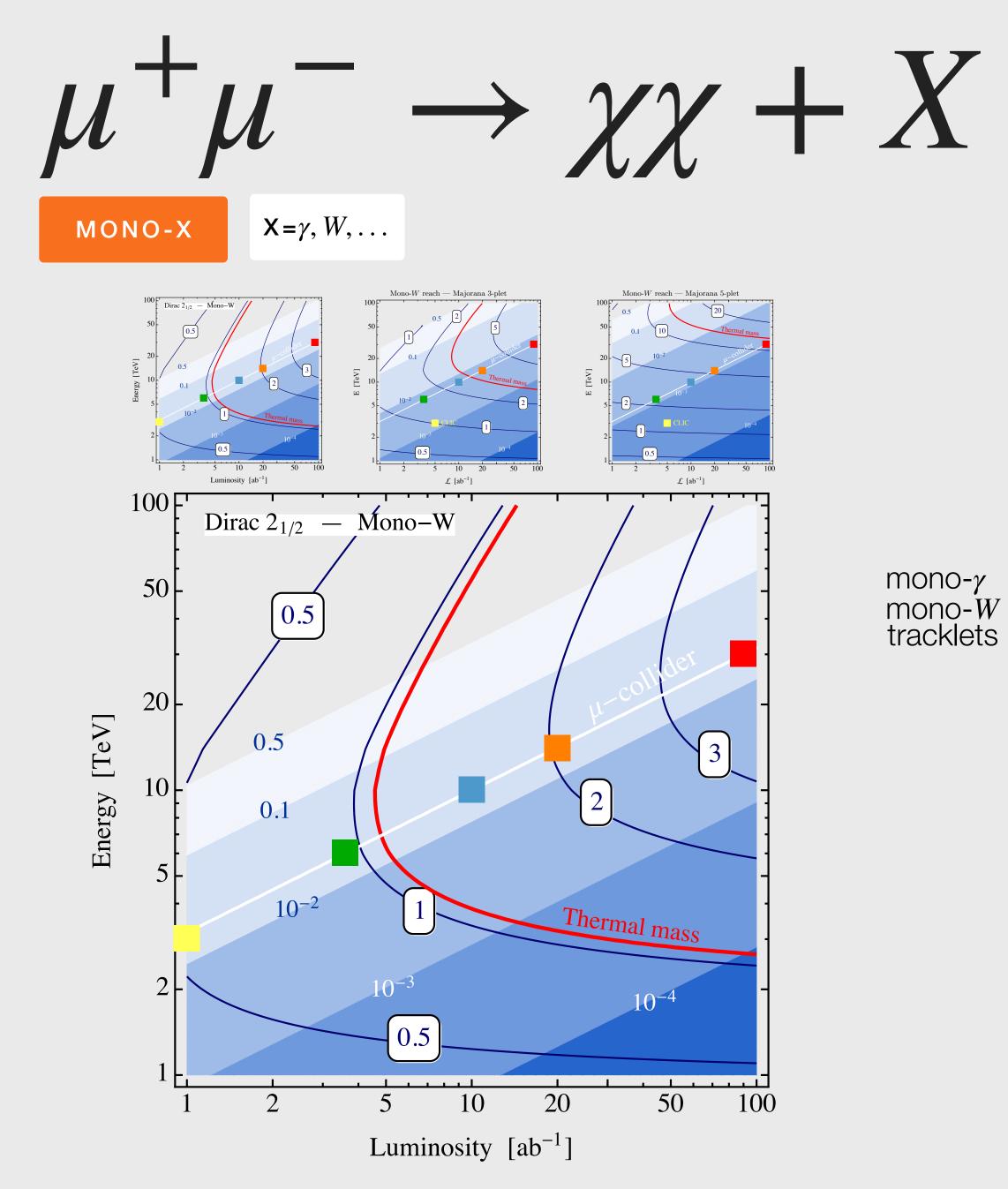
Roberto Franceschini - Physics at the high-energy muon collider



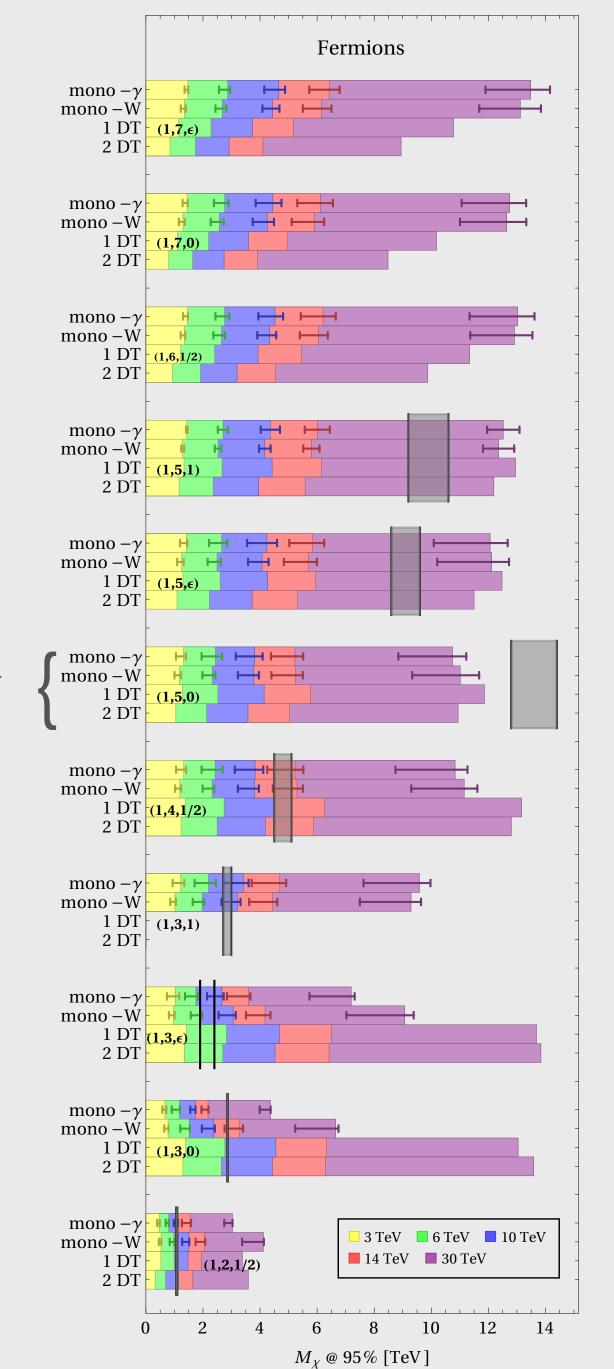


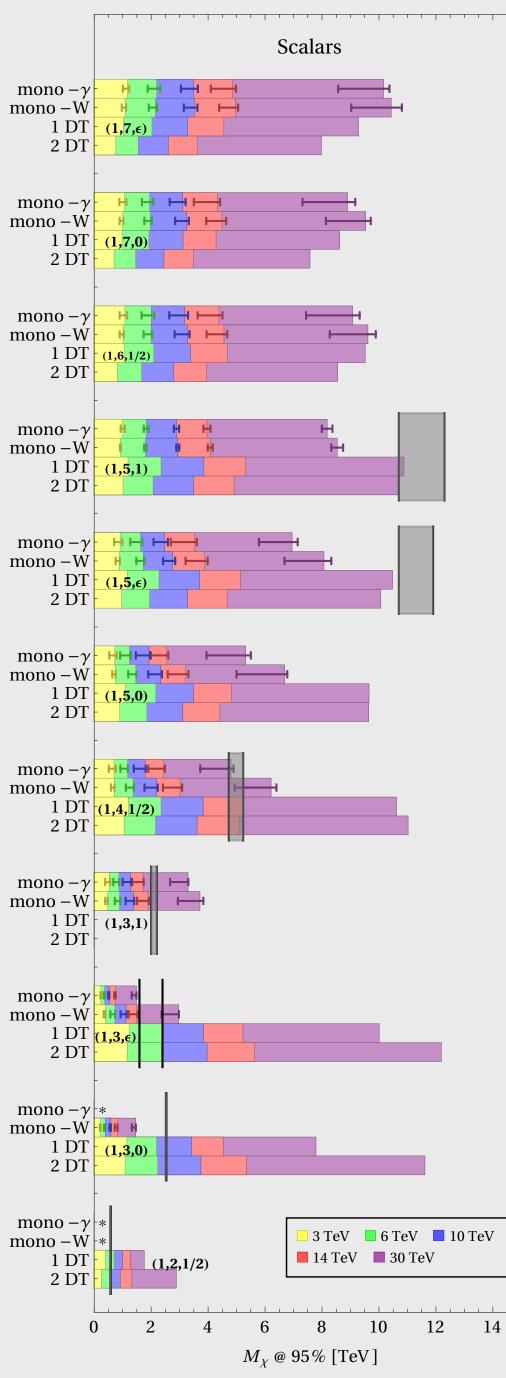


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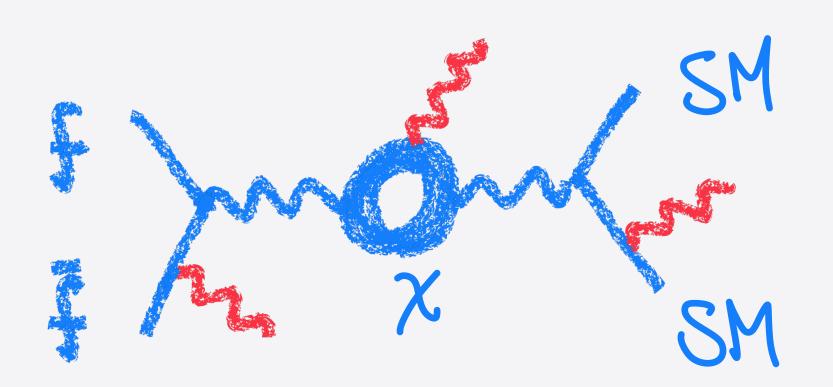
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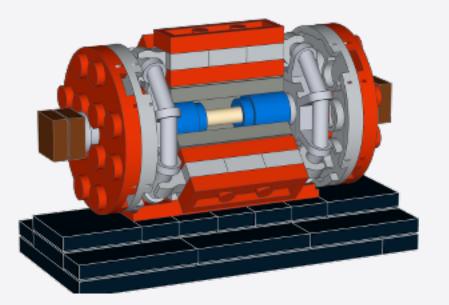




VIRTUAL* PRODUCTION

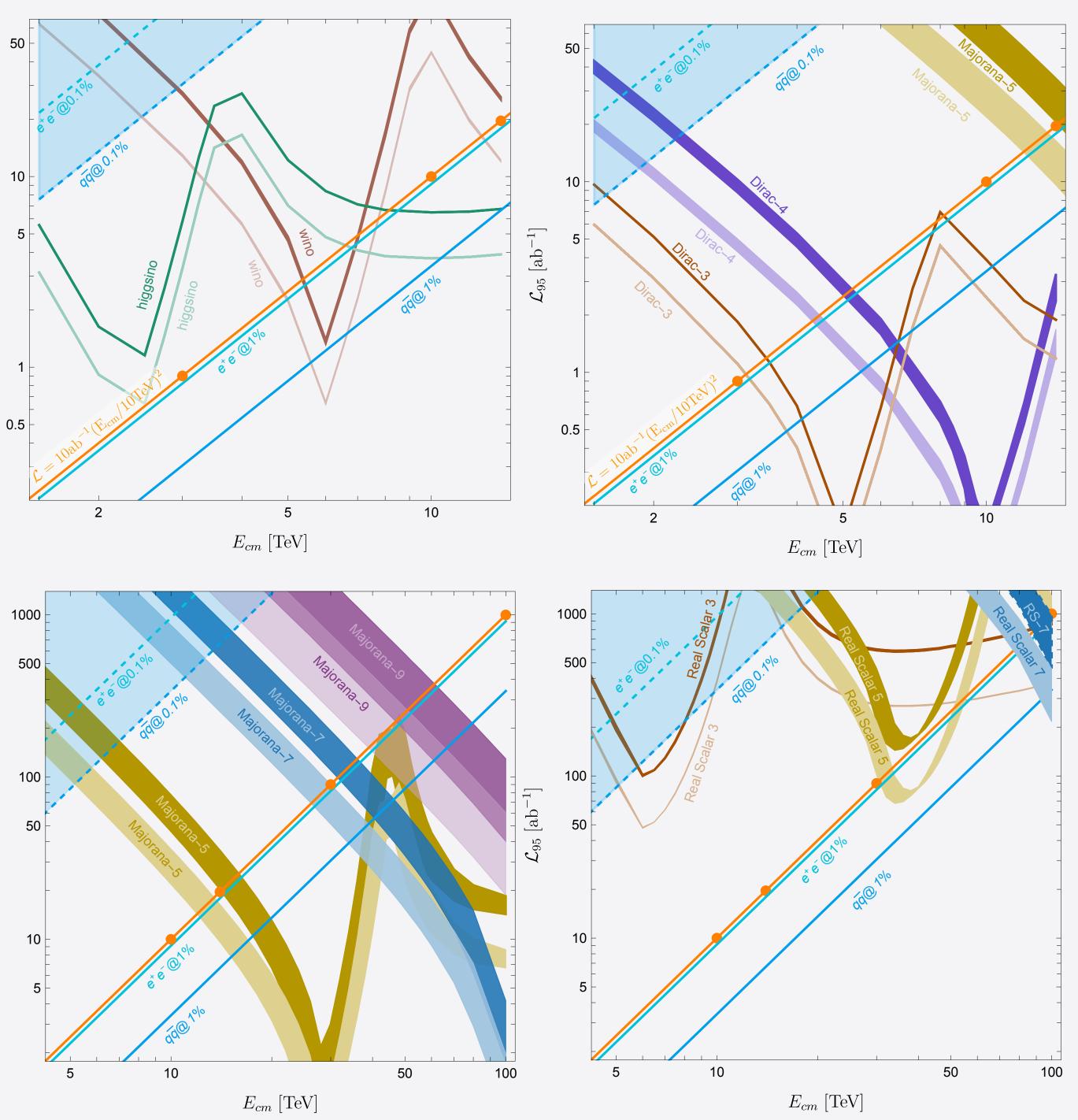


Virtual or propagating DM affects SM production rates

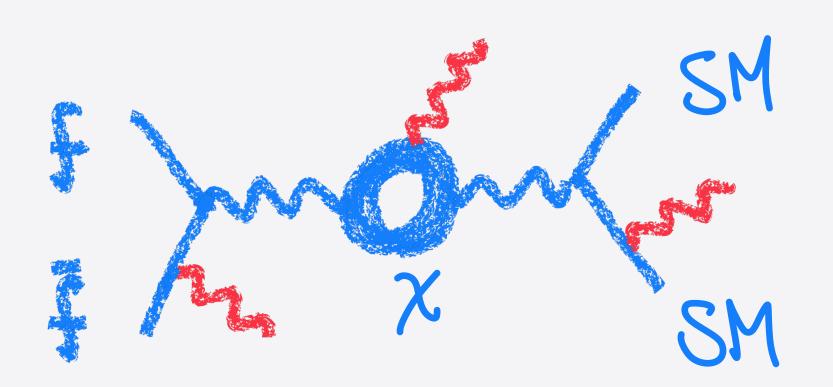


2040sup to 10+ TeV

 $\mathcal{L}_{95} ~[\mathrm{ab}^{-1}]$



VIRTUAL* PRODUCTION

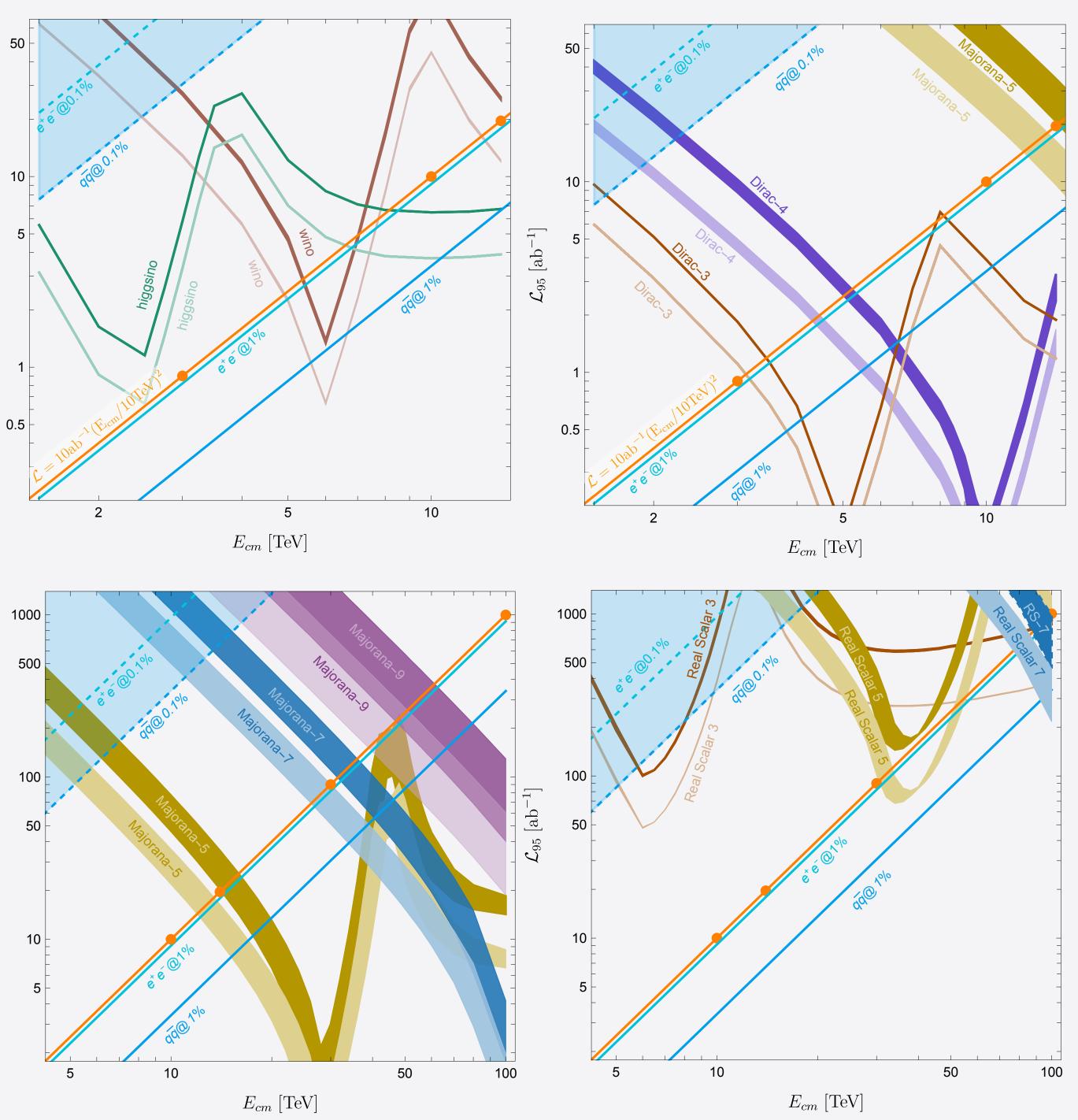


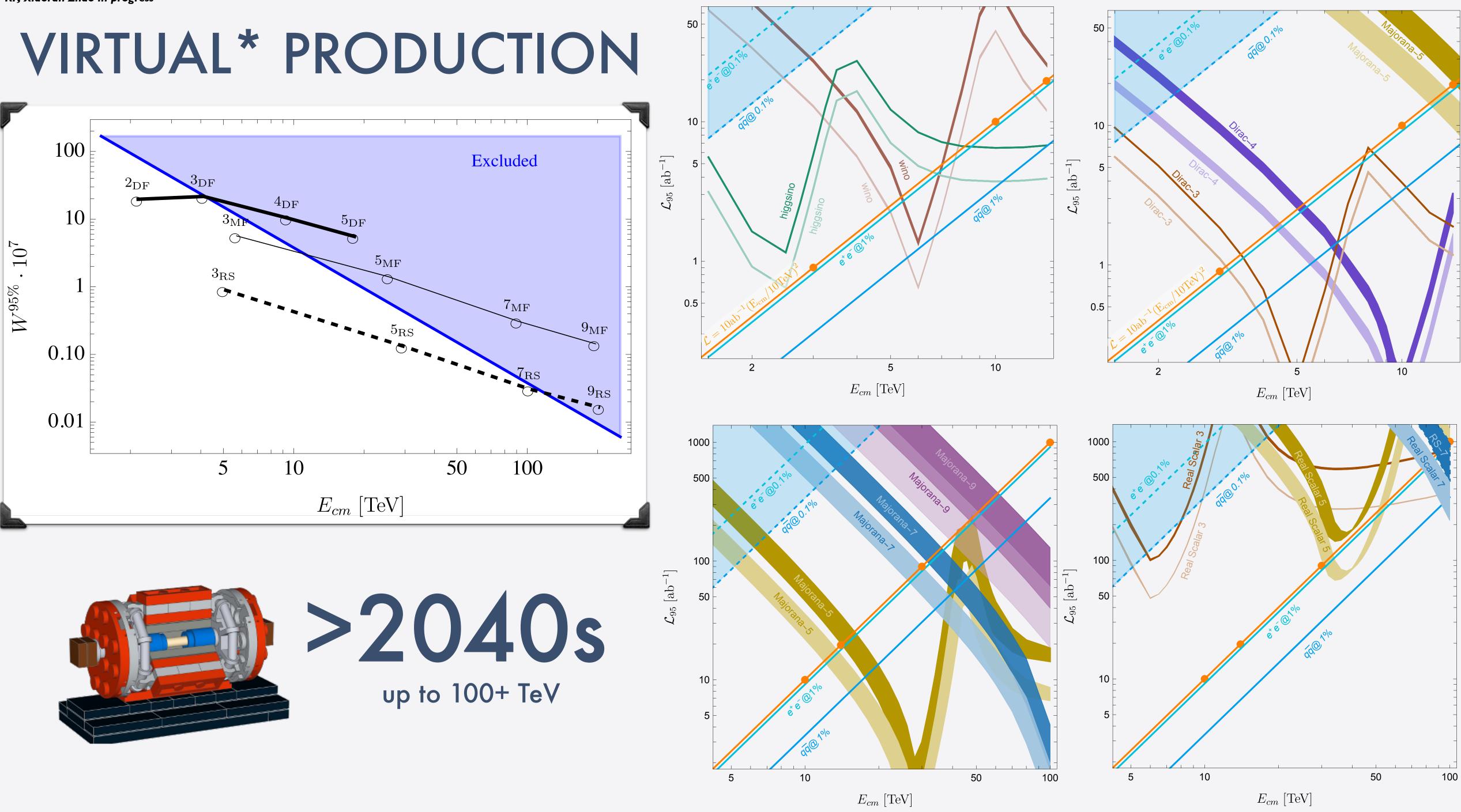
Virtual or propagating DM affects SM production rates



 $\mathcal{L}_{95} ~[\mathrm{ab}^{-1}]$

 $\mathcal{L}_{95} \left[\mathrm{ab}^{-1}
ight]$







DM spin	EW n-plet	M_{χ} (TeV)	$\Lambda_{\rm Landau}/M_{\chi}$	$(\sigma v)_{\rm tot}^{J=0}/(\sigma v)_{\rm max}^{J=0}$	$\Delta M_0 [{ m MeV}]$	$\Lambda_{\rm UV}^{\rm max}(\Delta M_0^{\rm min})/M_{\chi}$	ΔM_{-} [MeV]
Complex scalar	2	0.58 ± 0.01	$> M_{\rm Pl}$	_	$0.22 - 4.6 \times 10^4$	_	4.2 - 9600
	4	4.98 ± 0.05	$> M_{\rm Pl}$	0.004	0.22 - 10^4	_	3.2 - 2000
	6	34.9 ± 0.5	$\simeq 6 \times 10^{13}$	0.016	0.54 - 2300	_	280 - 660
	8	88 ± 2	2×10^4	0.12	$0.89 - 1.2 \times 10^3$	_	324 - 507
	10	167 ± 4	20	0.45	1.27 - 800	_	340 - 450
Dirac fermion	2	1.08 ± 0.01	$> M_{\rm Pl}$	_	0.22 - 5000	2×10^{5}	4.8 - 7800
	4	4.8 ± 0.1	$\simeq M_{\rm Pl}$	0.013	0.21 - 2200	$ imes 10^5$	3.6 - 2600
	6	31.7 ± 0.5	2×10^4	0.057	0.51 - 510	$\times 10^4$	185 - 780
	8	82 ± 2	14	0.37	0.86 - 800	3000	290 - 550

DM spin	EW n-plet	M_{χ} (TeV)	$(\sigma v)_{\rm tot}^{J=0}/(\sigma v)_{\rm max}^{J=0}$	$\Lambda_{ m Landau}/M_{ m DM}$	$\Lambda_{ m UV}/M_{ m DM}$
	3	2.53 ± 0.01		3×10^{37}	$4 \times 10^{24} *$
	5	15.4 ± 0.7	0.002	5×10^{36}	2×10^{24}
Real scalar	7	54.2 ± 3.1	0.022	2×10^{19}	2×10^{24}
	9	117.8 ± 15.4	0.088	3×10^3	2×10^{24}
	11	199 ± 42	0.25	20	3×10^{24}
	13	338 ± 102	0.6	3.5	3×10^{24}
	3	2.86 ± 0.01		3×10^{37}	$8 \times 10^{12*}$
Majorana fermion	5	13.6 ± 0.8	0.003	3×10^{17}	5×10^{12}
	7	48.8 ± 3.3	0.019	1×10^4	4×10^7
	9	113 ± 15	0.07	30	3×10^7
	11	202 ± 43	0.2	6	3×10^7
	13	324.6 ± 94	0.5	2.6	3×10^7

PRELIMINARY

PRELIMINARY

WULZER

Reports of the working groups: Physics

30/7-018 - Kjell Johnsen Auditorium, CERN

Monte Carlo Tools [Mario Chiesa]

Conclusions

- only few MC event generators for high-energy μ -colls available, generally only at LO accuracy
- in principle it is possible to transpose all the technology developed for high-precision generators for hadron colliders to the leptonic environment
- in practice, some new challenges arise, mainly connected to the large \sqrt{s} and the large final-state multiplicities, Sudakov corrections

Andrea Wulzer

08:50 - 09:05



WULZER

Reports of the working groups: Physics

30/7-018 - Kjell Johnsen Auditorium, CERN

EW and QCD physics at the muon collider [Yang Ma]

Improving understanding of partons in the proton MuC is **also** a Vector Bosons Collider (but **also** a muon collider: this is its **dual** nature)

Compare the "EW LHC" with LHC

pp VS $\mu\mu$

$$\mathcal{L}_{W_{\lambda_1}^+ W_{\lambda_2}^-} = \int_{\tau}^1 \frac{\mathrm{d}\xi}{\xi} f_{W_{\lambda_1}}(\xi, \mu_f) f_{W_{\lambda_1}}(\frac{\tau}{\xi}, \mu_f)$$

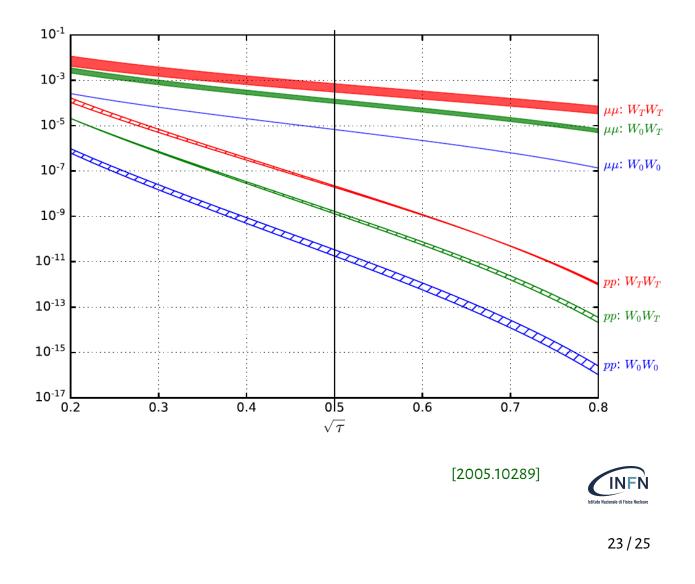
Consider the two colliders in the same ring

$$\sqrt{s}_{\mu\mu} = \sqrt{s}_{pp}$$

For $2 \rightarrow 1$ processes, take a benchmark

$$\sqrt{\tau} = \frac{M}{\sqrt{s}} = \frac{1}{2}$$

The ratio $\mu\mu/pp$ is larger than 10^4 !



Andrea Wulzer

08:50 - 09:05



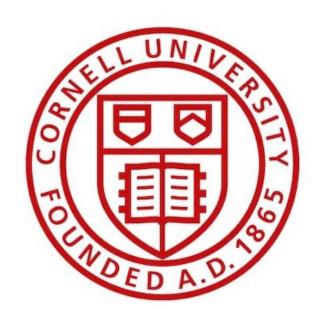
PARALLEL TALKS

The physics case of a very forward muons detector

40/S2-D01 - Salle Dirac, CERN

The physics case of a very forward muon detector

Maximilian Ruhdorfer Cornell University





Muon Collider Collaboration Meeting October 12, 2022

Work in progress with R. Masarotti, E. Salvioni and A. Wulzer



09:00 - 09:30



The physics case of a very forward muons detector

40/S2-D01 - Salle Dirac, CERN

This Talk

 ϕ can be Dark Matter

 ϕ can be part of the model that stabilizes the Higgs mass

 ϕ can modify the Higgs potential

1. Physics case for very forward muon detector (idealized)

• Focus on scalar Higgs portal to invisible new physics

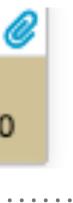
$$-rac{\lambda}{2}\phi^2|H|^2$$

marginal portal

• Assume perfect resolution of MIM,...

2. Realistic case study: invisible Higgs decays

- lacksquare
- New BGs become important



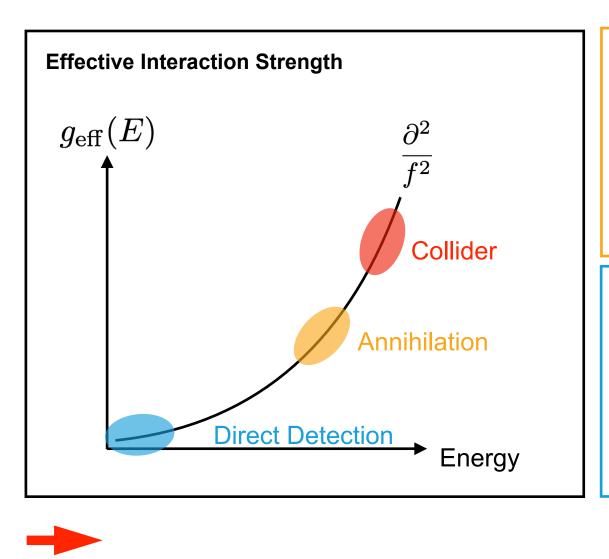
09:00 - 09:30

$$\frac{c_d}{2f^2}\partial_\mu\phi^2\partial^\mu|H|^2$$

derivative portal

Include accelerator and detector effects (beam energy spread,...)





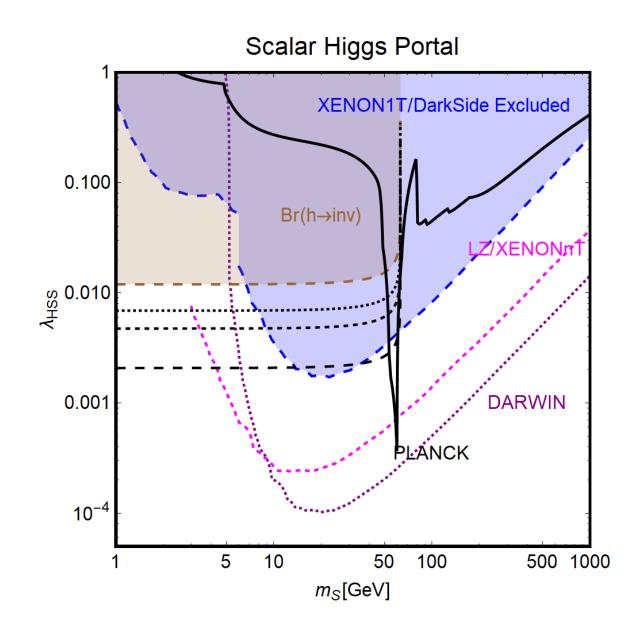
The physics case of a very forward muons detector

40/S2-D01 - Salle Dirac, CERN

 ϕ can be Dark Matter

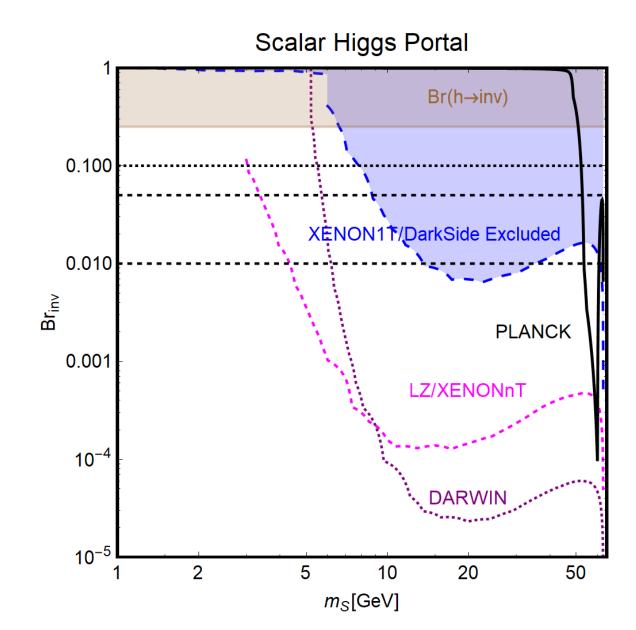
 ϕ can be part of the model that stabilizes the Higgs mass

 ϕ can modify the Higgs potential

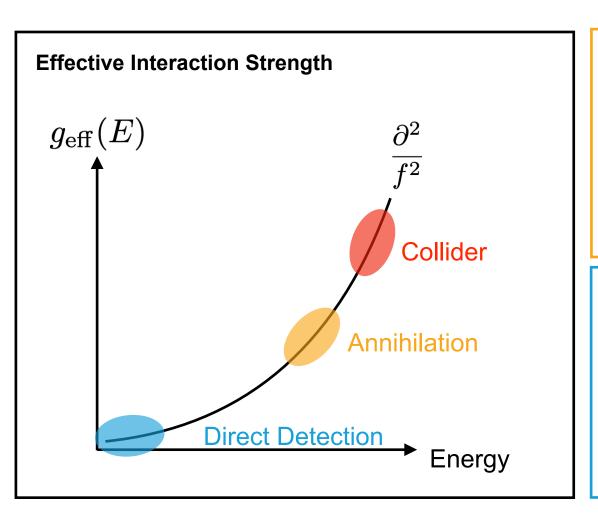




09:00 - 09:30



ϕ can be Dark Matter



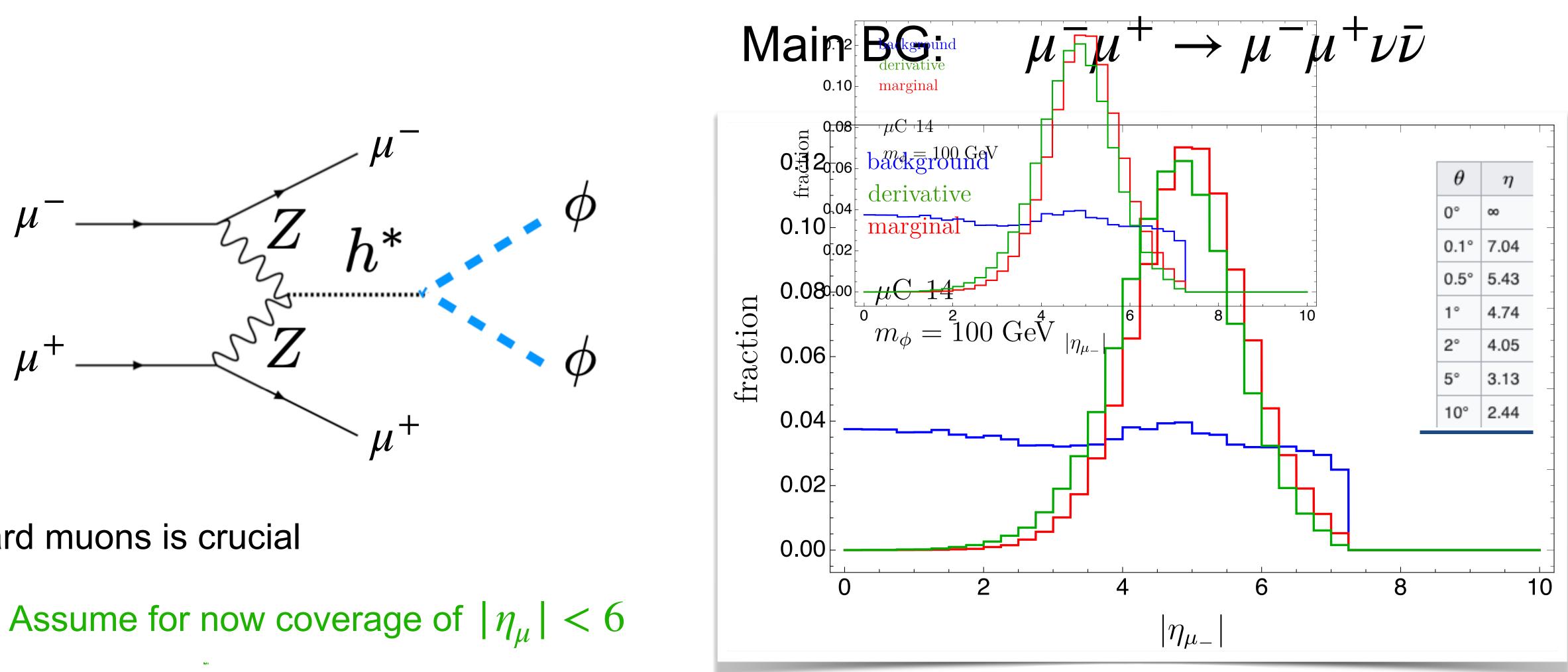
Goldstone bosons tend to not be stable, e.g. $\pi^0 \to \gamma\gamma$ Stability of ϕ is unwarranted unless a symmetry is imposed





The physics case of a very forward muons detector

40/S2-D01 - Salle Dirac, CERN



prward muons is crucial

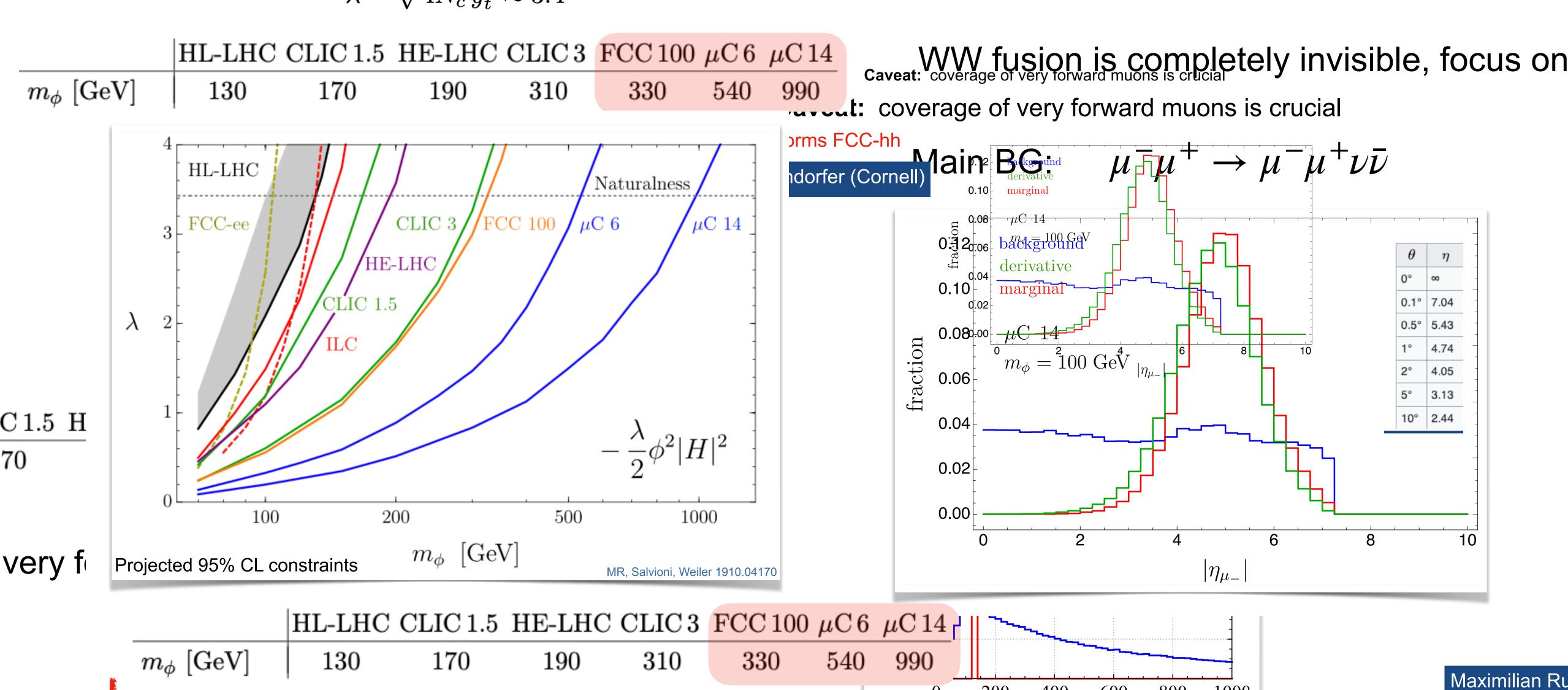


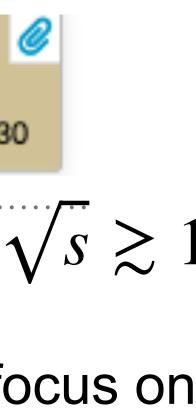
09:00 - 09:30 WW fusion is completely invisible, focus Caveat: coverage of very forward muons is crucial

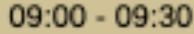
The physics case of a very forward muons detector

40/S2-D01 - Salle Dirac, CERN

$$\lambda = \sqrt{4N_c} \, y_t^2 \approx 3.4$$



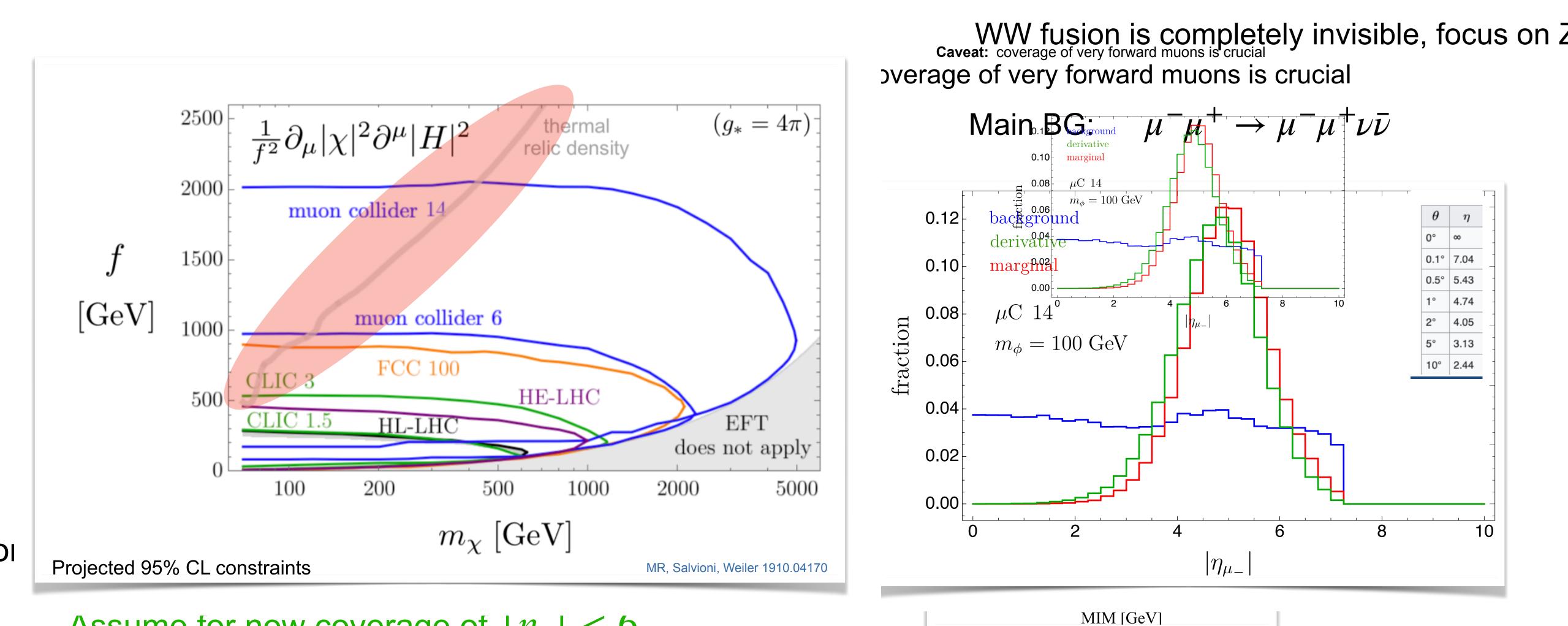




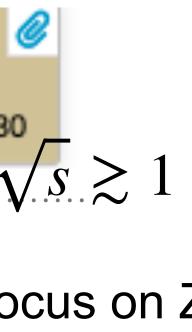
Maximilian Ru

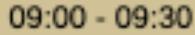
The physics case of a very forward muons detector

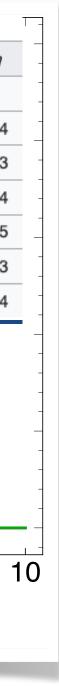
40/S2-D01 - Salle Dirac, CERN



Assume for now coverage of $|\eta_{\mu}| < 6$

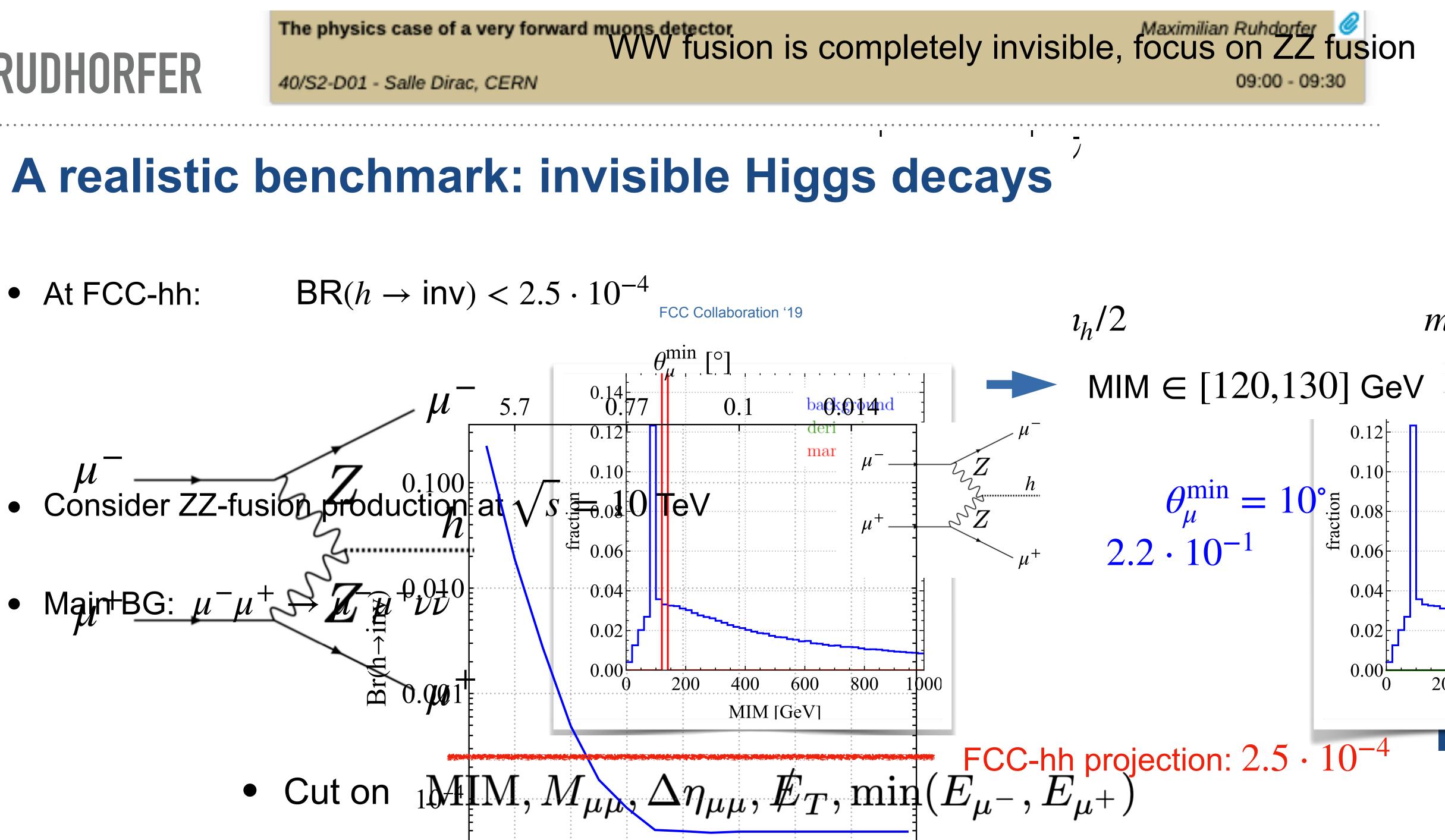






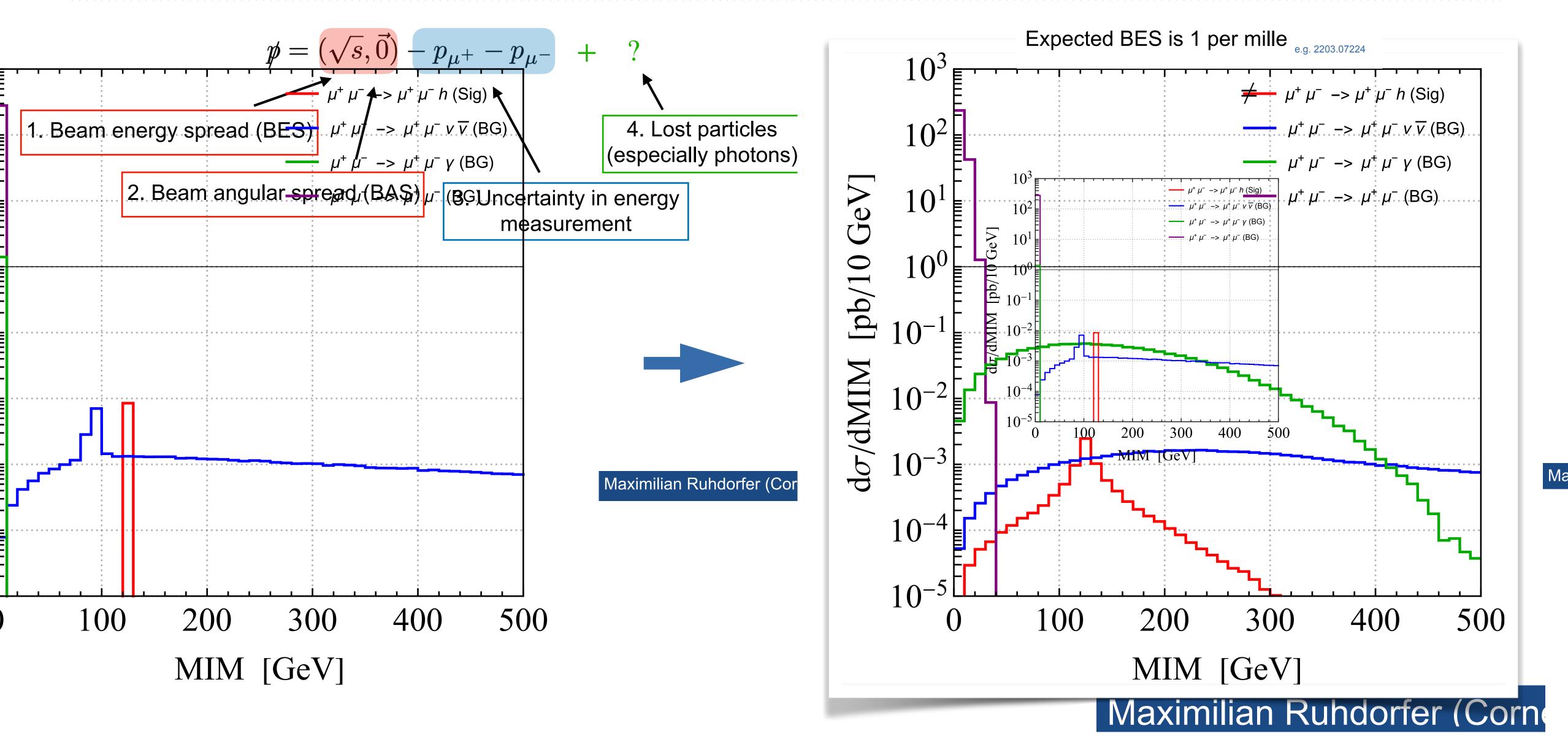


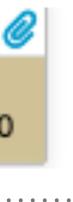
• At FCC-hh:



The physics case of a very forward muons detector

40/S2-D01 - Salle Dirac, CERN

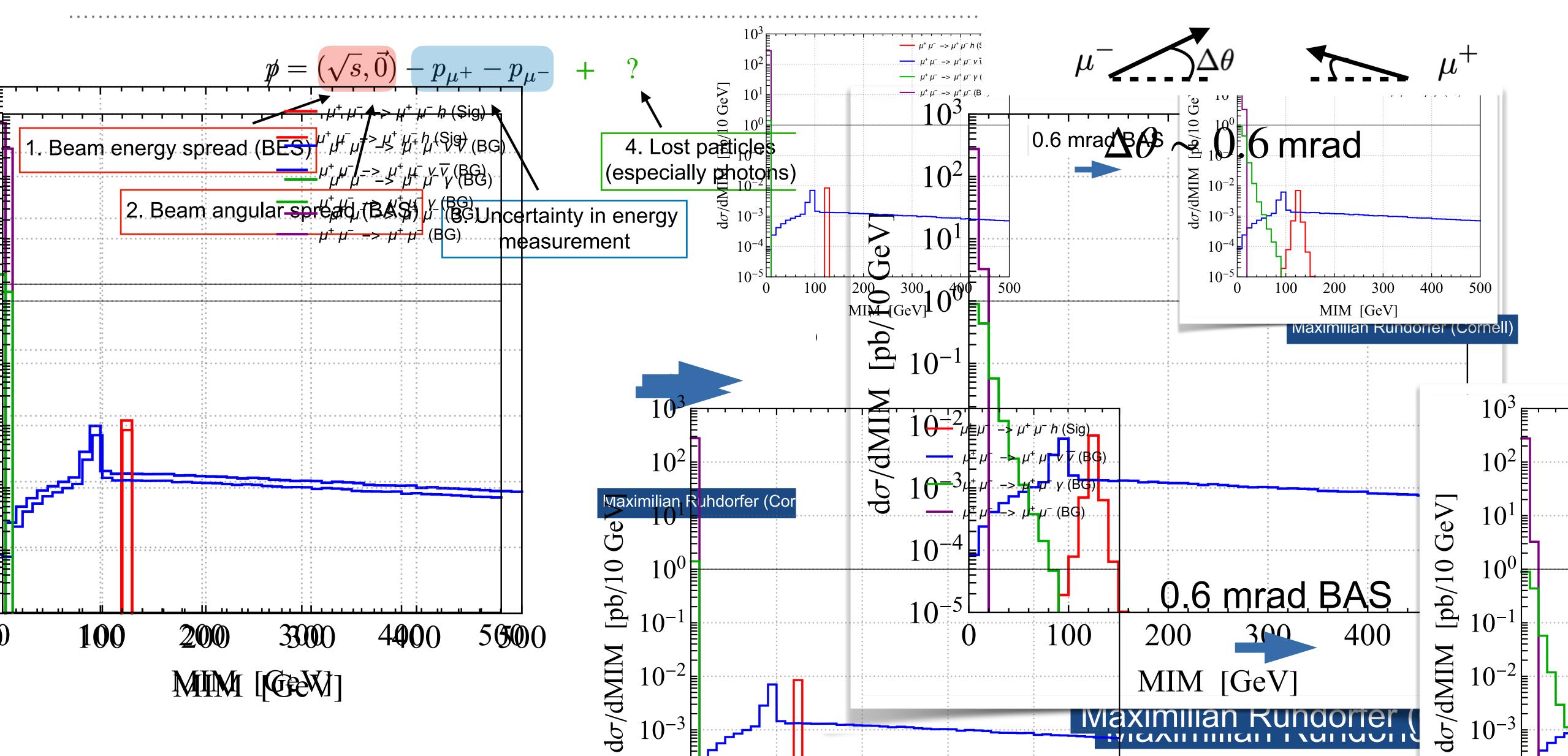




09:00 - 09:30

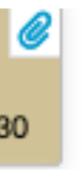
The physics case of a very forward muons detector

40/S2-D01 - Salle Dirac, CERN

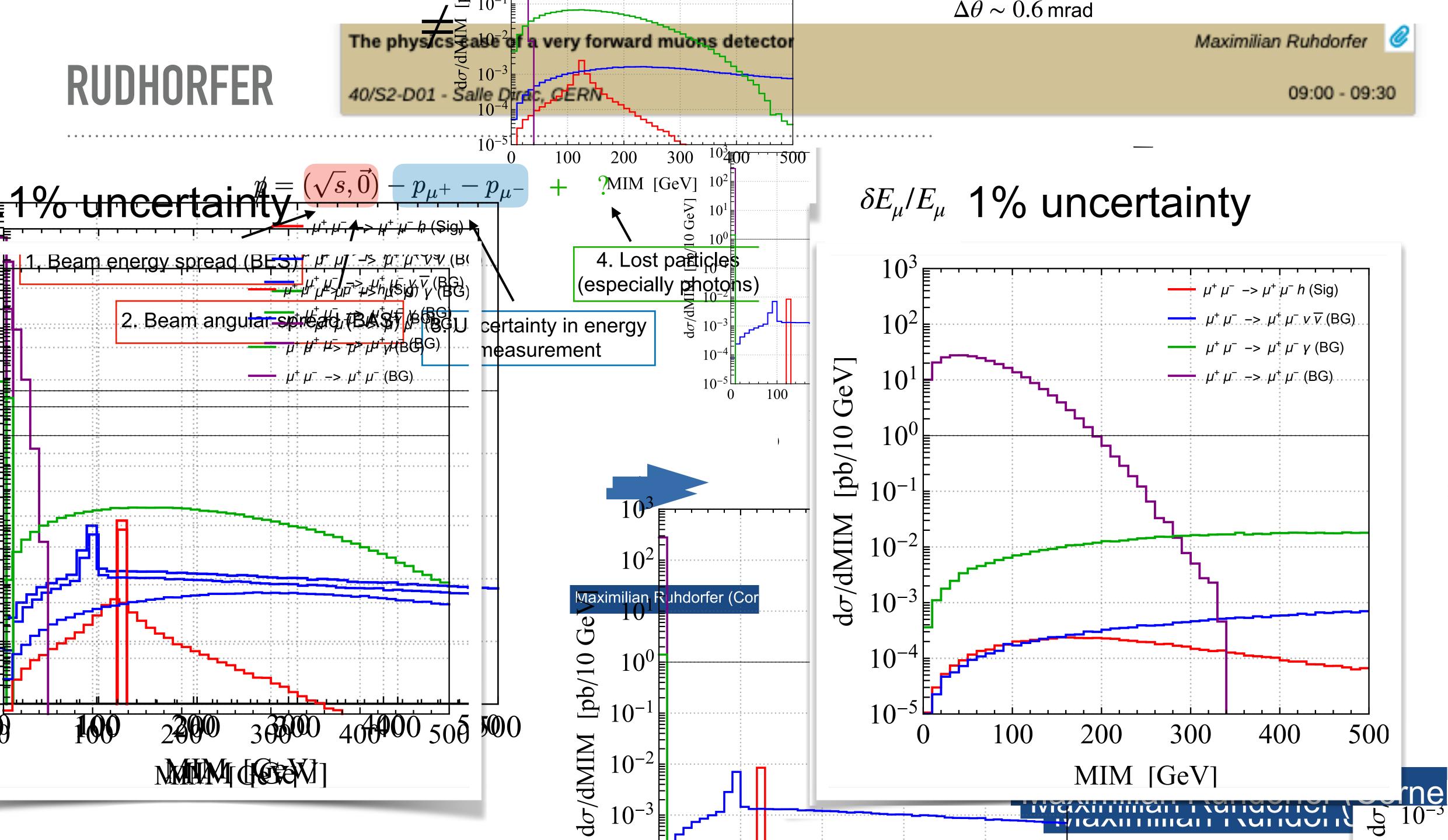


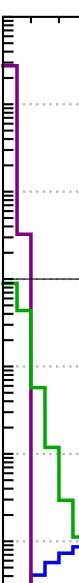
$\Delta\theta \sim 0.6 \,\mathrm{mrad}$

Maximilian Ruhdorfer



09:00 - 09:30

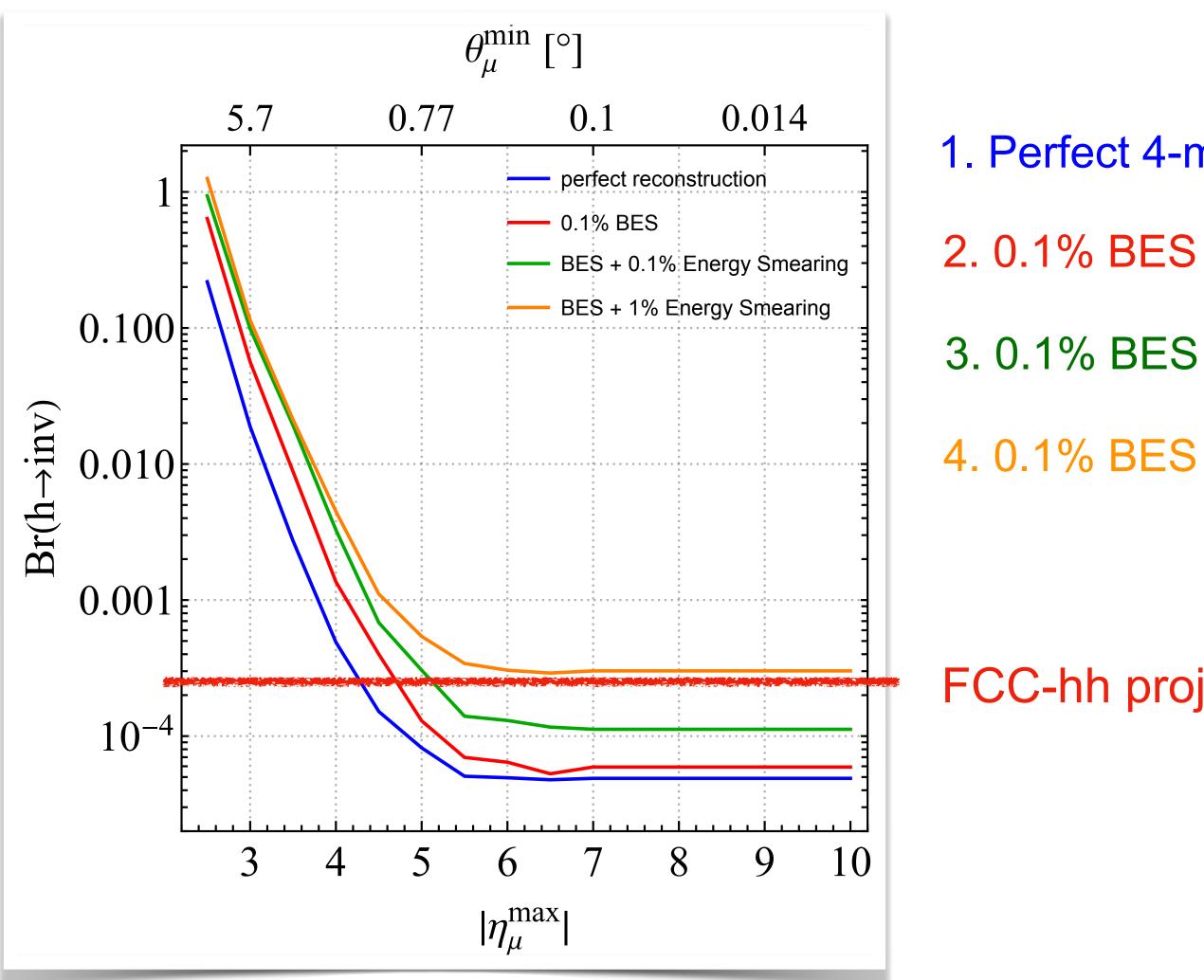




The physics case of a very forward muons detector

40/S2-D01 - Salle Dirac, CERN

• Sensitivity to $BR(h \rightarrow inv)$ with all effects combined





09:00 - 09:30

- 1. Perfect 4-momentum reconstruction
- 3. 0.1% BES + 0.1% energy uncertainty
- 4. 0.1% BES + 1% energy uncertainty

FCC-hh projection: $2.5 \cdot 10^{-4}$

EW and QCD physics at the muon collider

MA

40/S2-D01 - Salle Dirac, CERN

EW and QCD physics at the muon collider

Yang Ma

INFN Bologna

The 1st Collaboration Meeting of the Muon Collider Study October 12, 2022

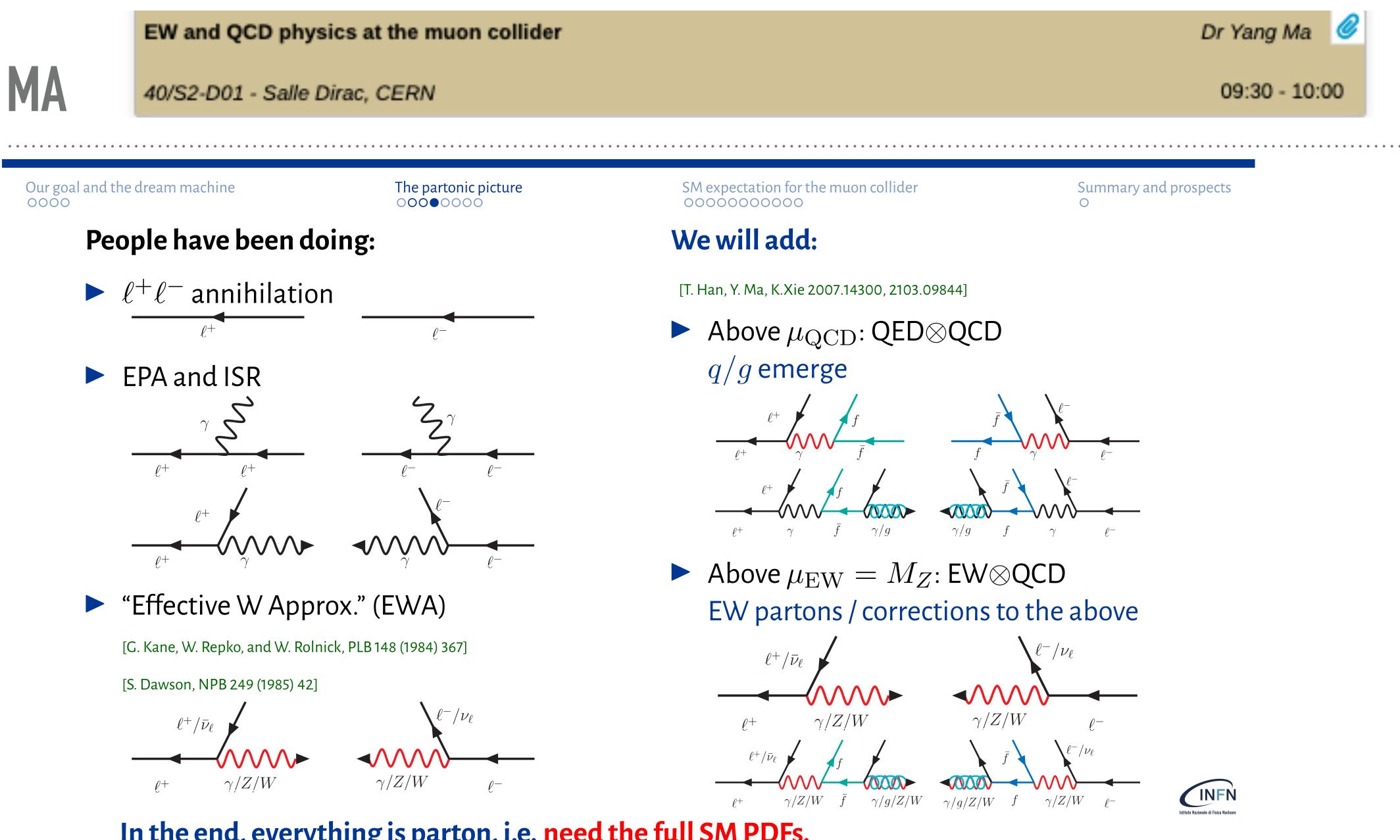




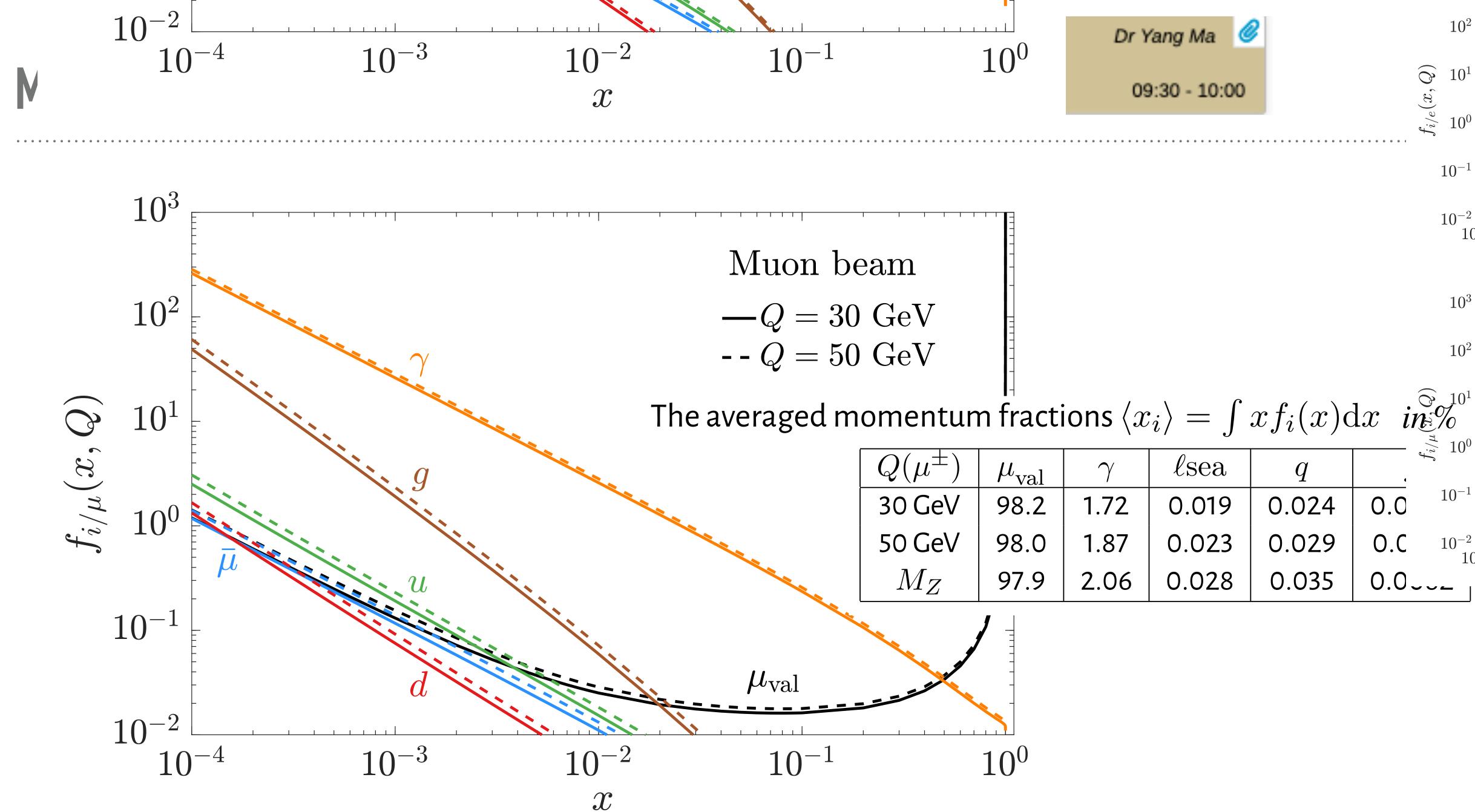
Istituto Nazionale di Fisica Nucl

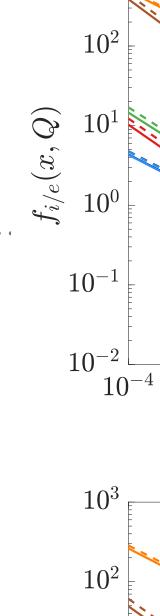


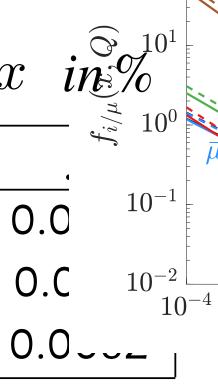
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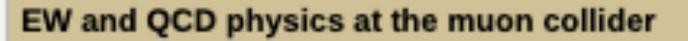


In the end, everything is parton, i.e. need the full SM PDFs.









40/S2-D01 - Salle Dirac, CERN

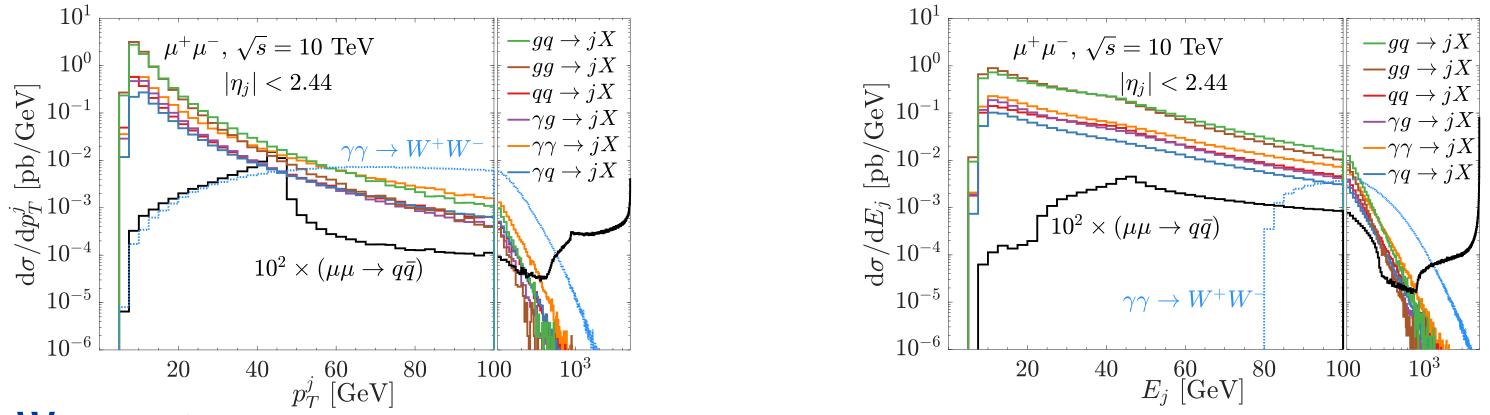
Our goal and the dream machine 0000

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The partonic picture 0000000

SM expectation for the muon collider 00000000000

Inclusive jet distributions at a 10 TeV muon collider **Important guidelines for future analysis**



We expect

- \blacktriangleright Jet production dominates over WW production until $p_T > 60$ GeV;
- WW production takes over around energy ~ 200 GeV.

The SM EW sector, as well as any possible BSM, can only be seen in the high p_T (E_j) range.



Summary and prospects

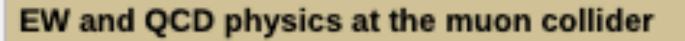
Copious q/g-jets within acceptance.

Concrete impact in analyses, i.e. h or resonance searches, yet to be evaluated.



21/25





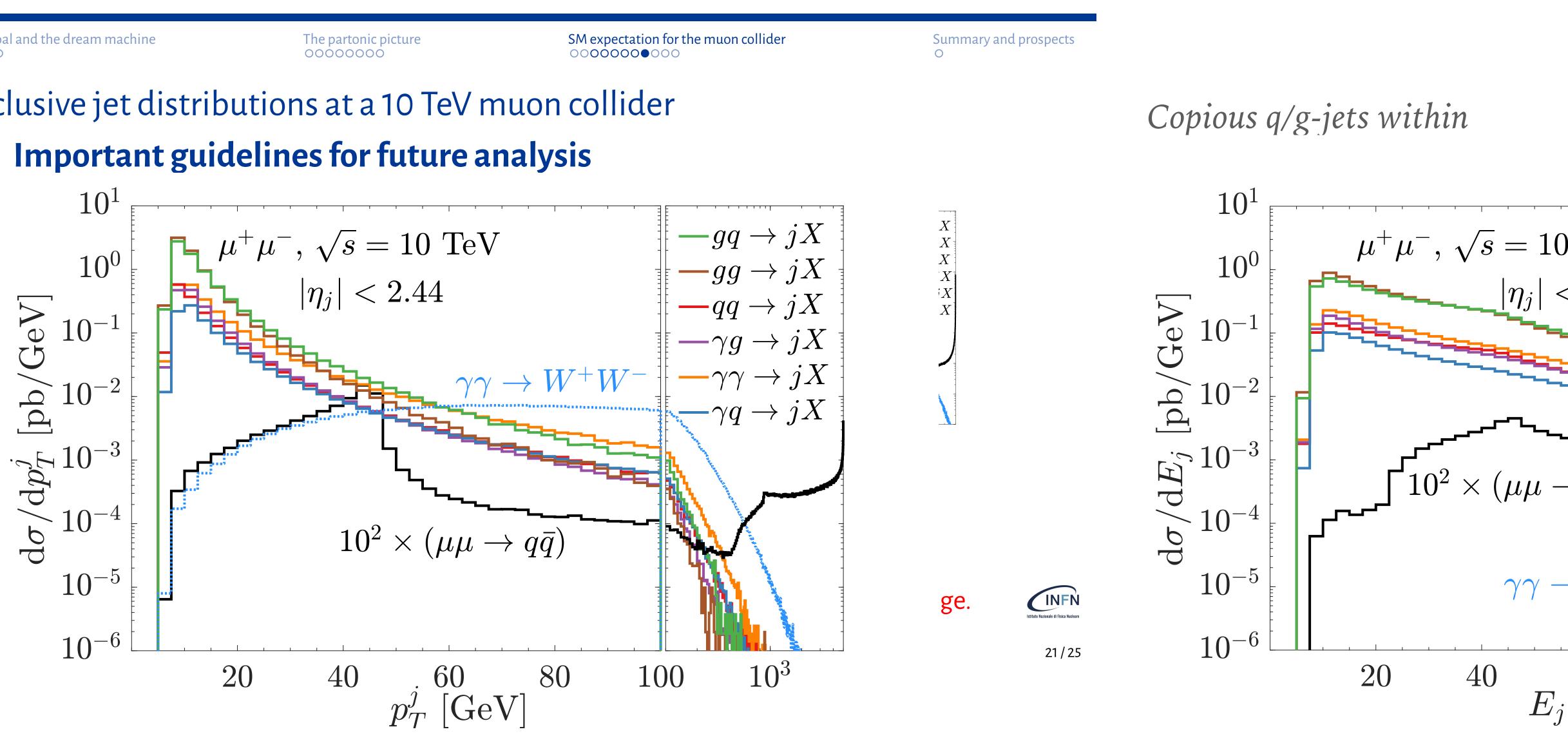
40/S2-D01 - Salle Dirac, CERN

Our goal and the dream machine 0000

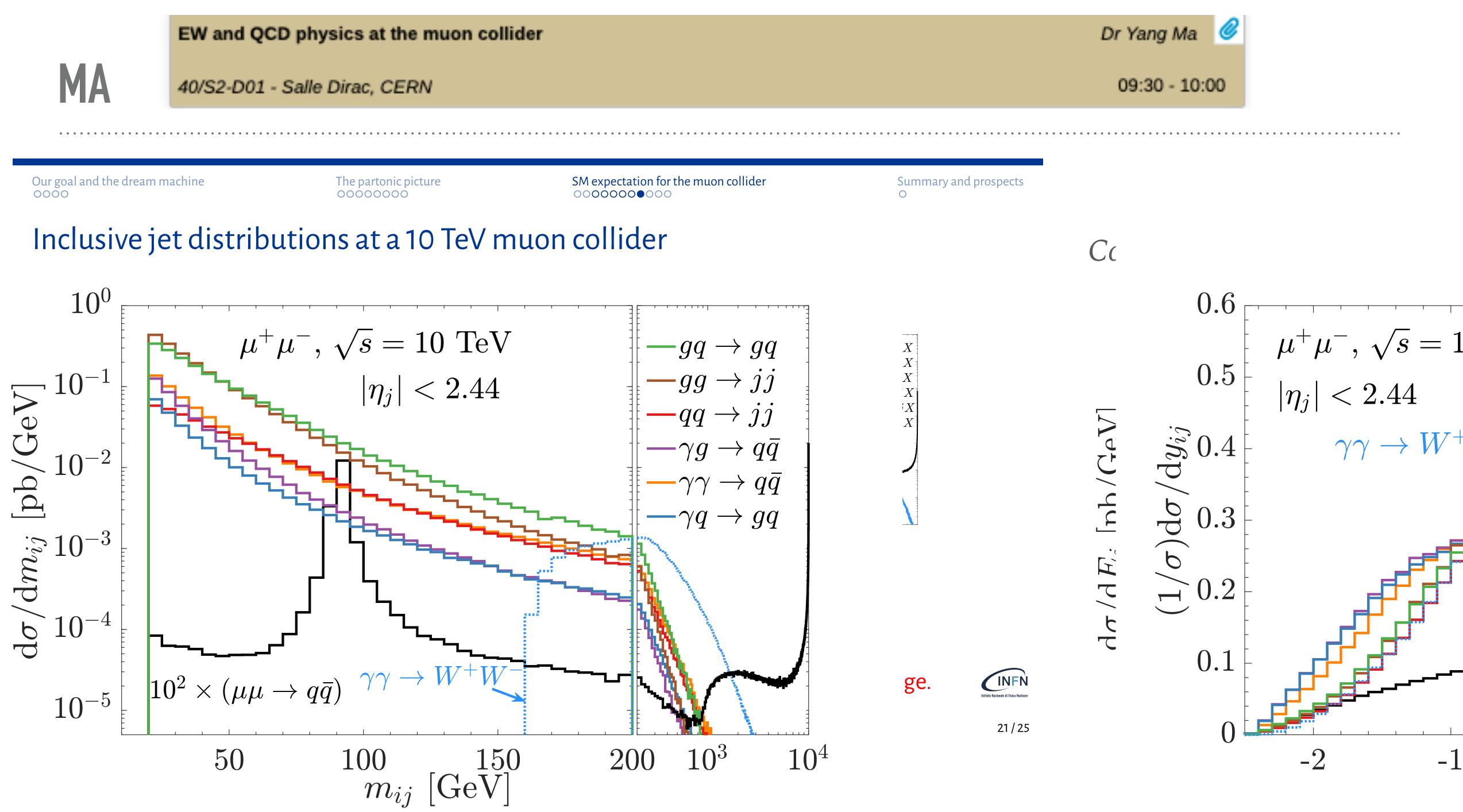
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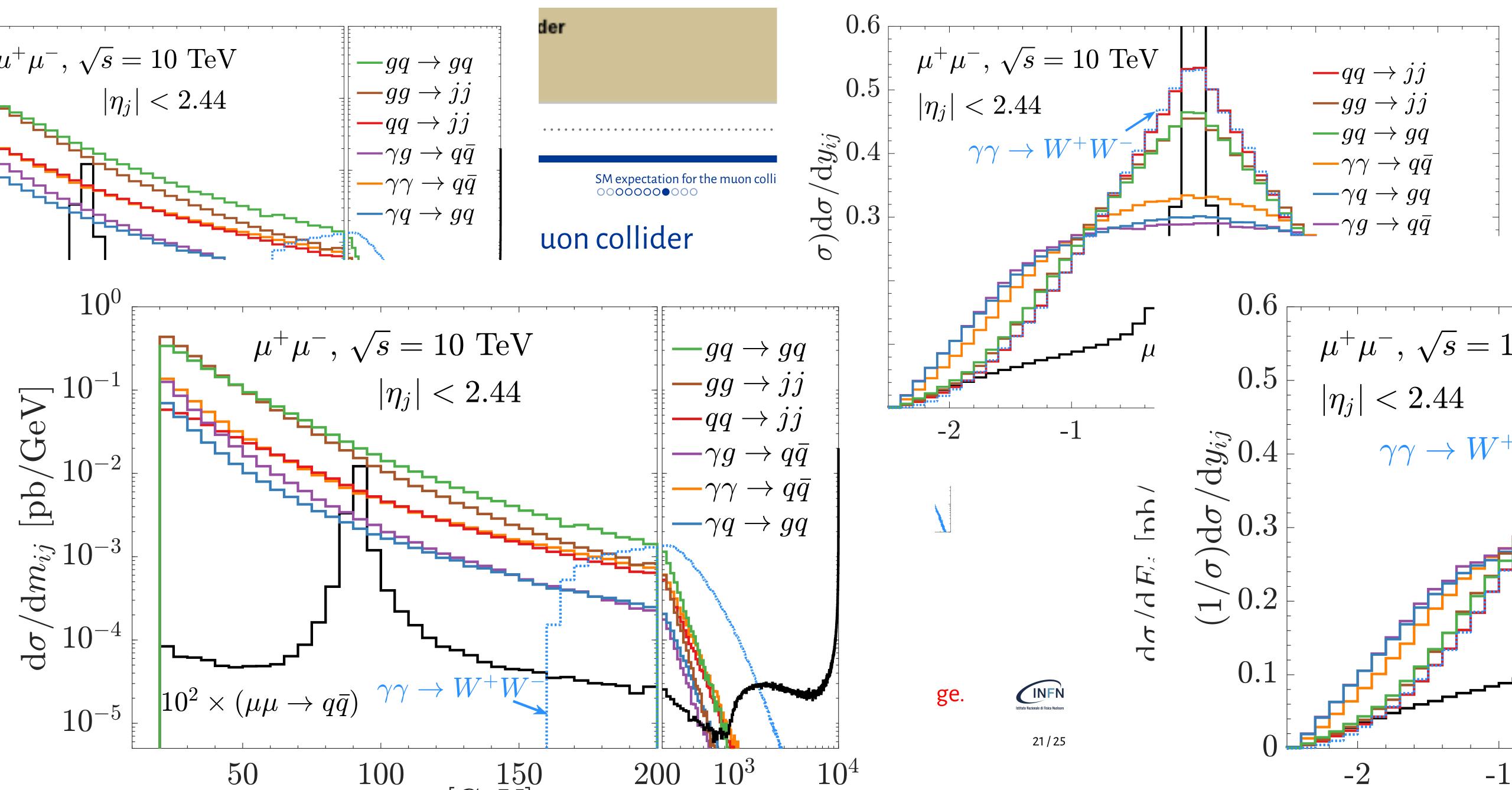
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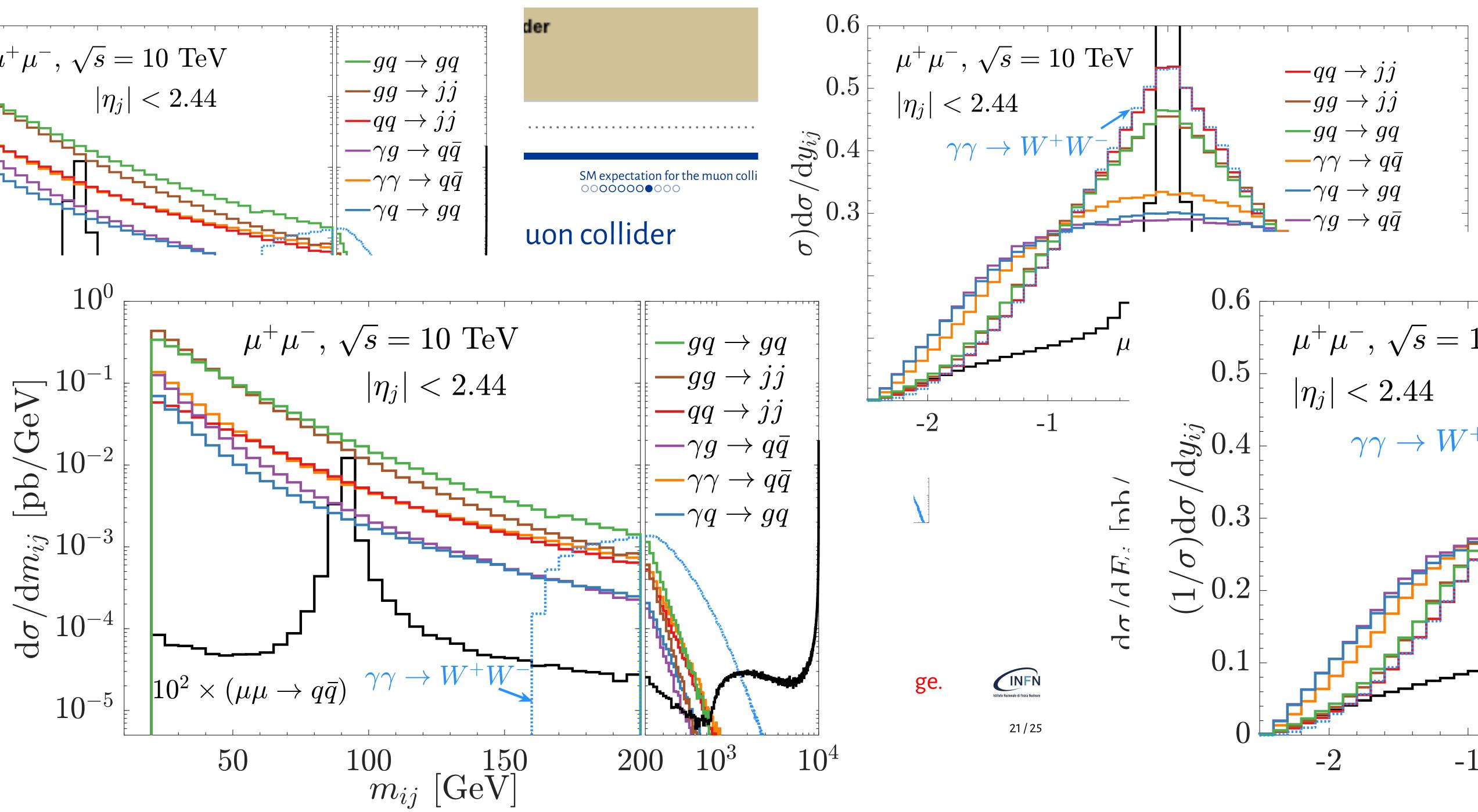
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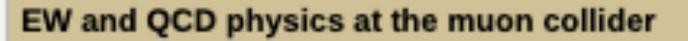












40/S2-D01 - Salle Dirac, CERN

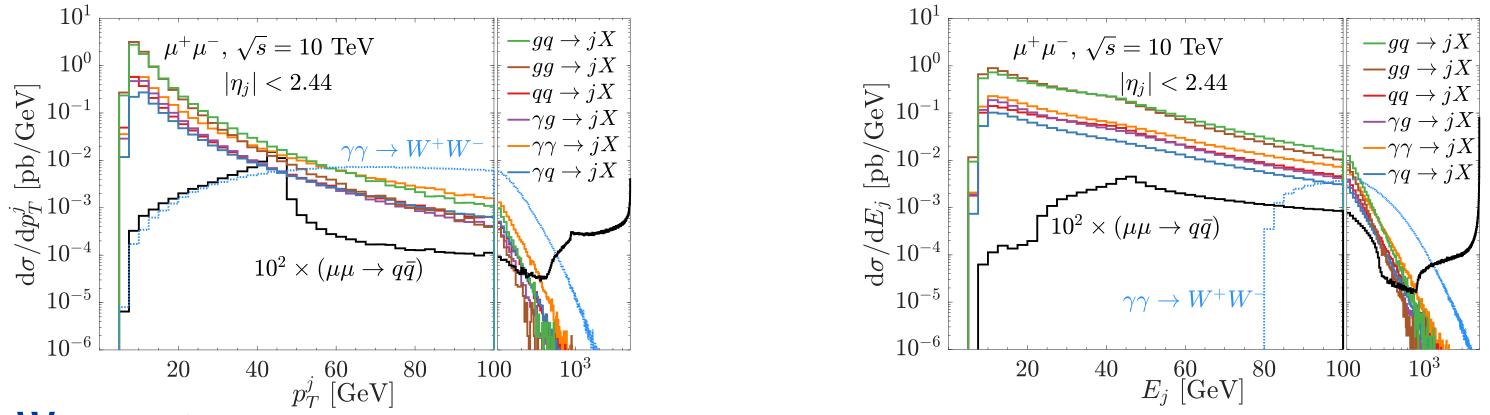
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SM expectation for the muon collider 00000000000

Inclusive jet distributions at a 10 TeV muon collider **Important guidelines for future analysis**



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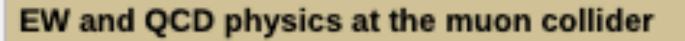
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21/25





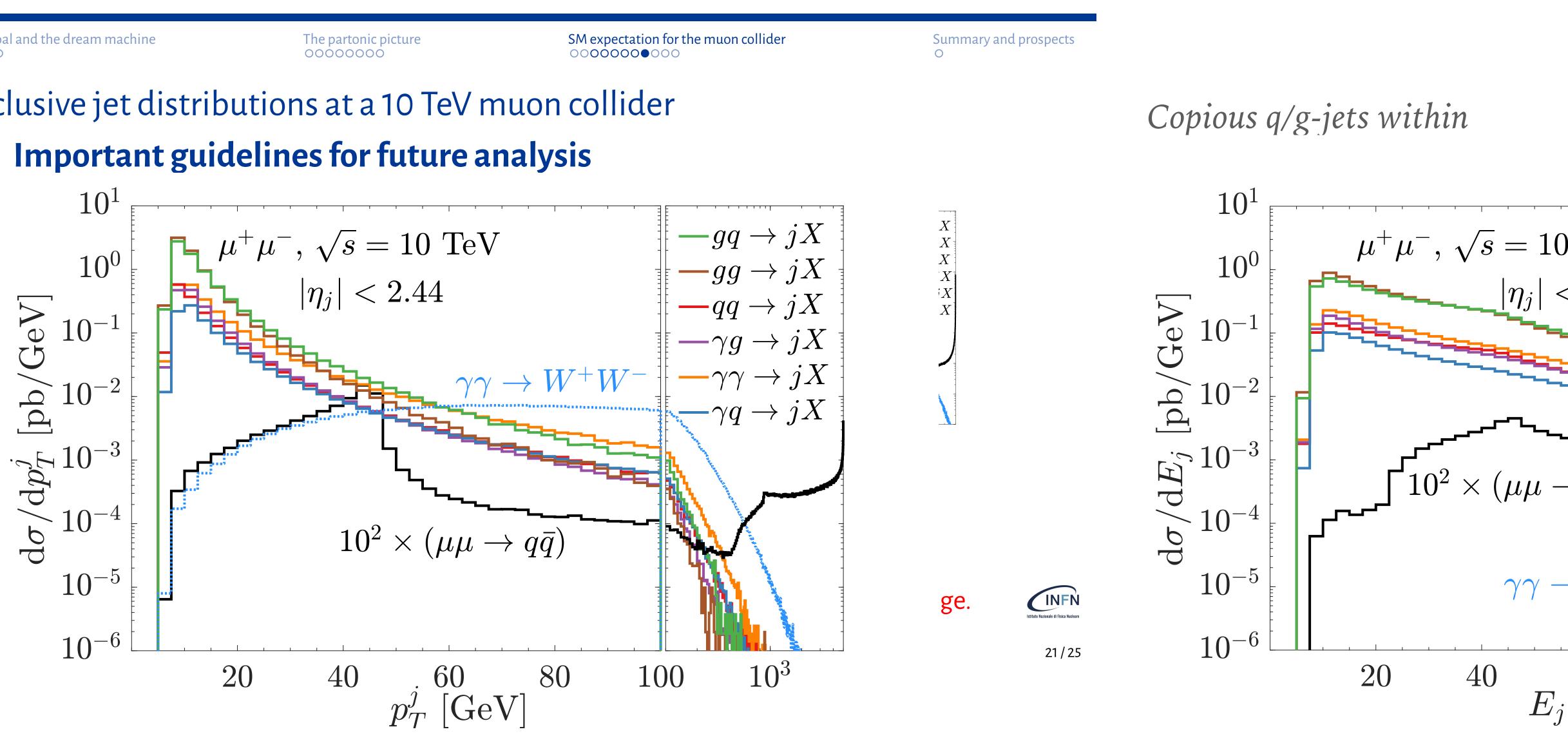
40/S2-D01 - Salle Dirac, CERN

Our goal and the dream machine 0000

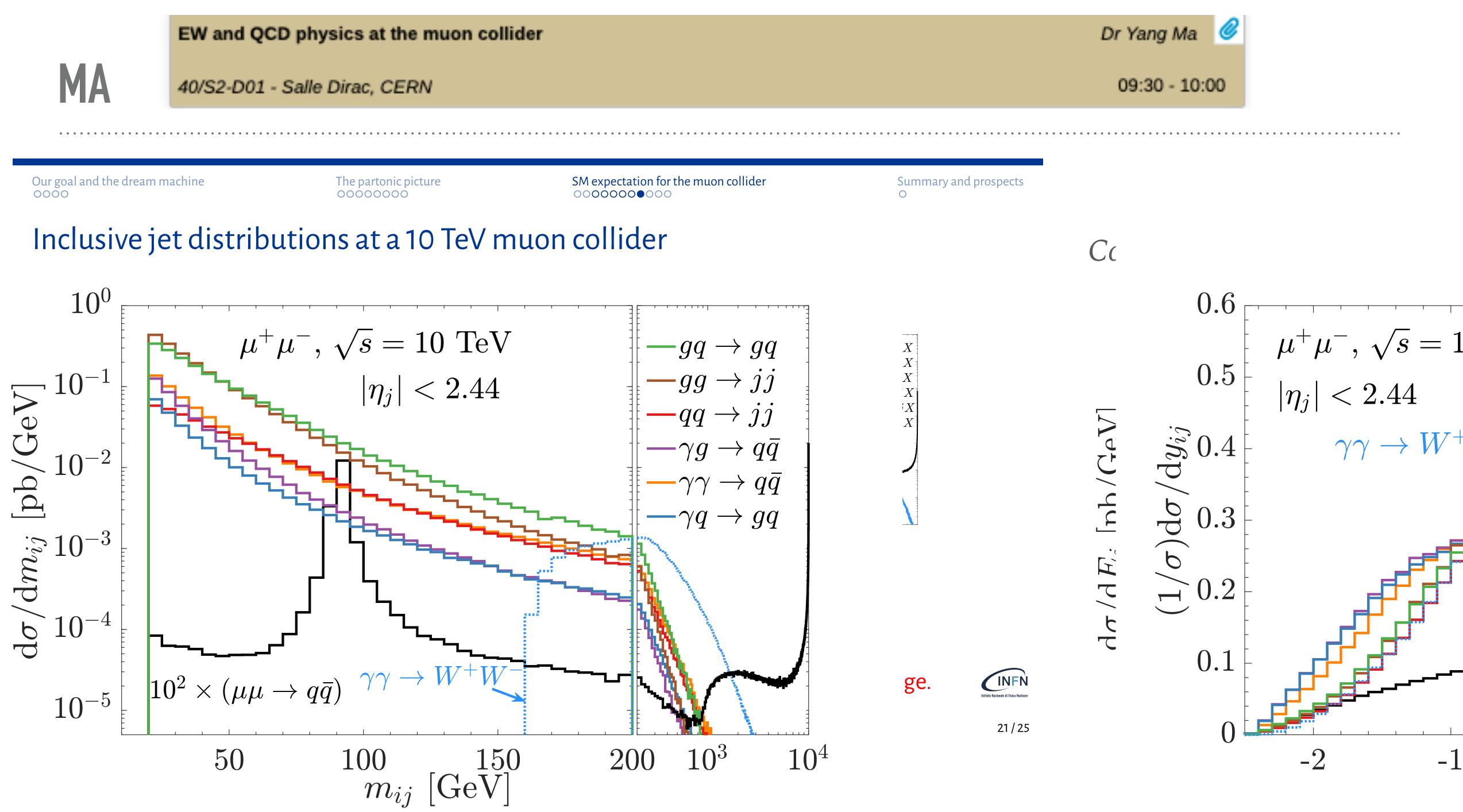
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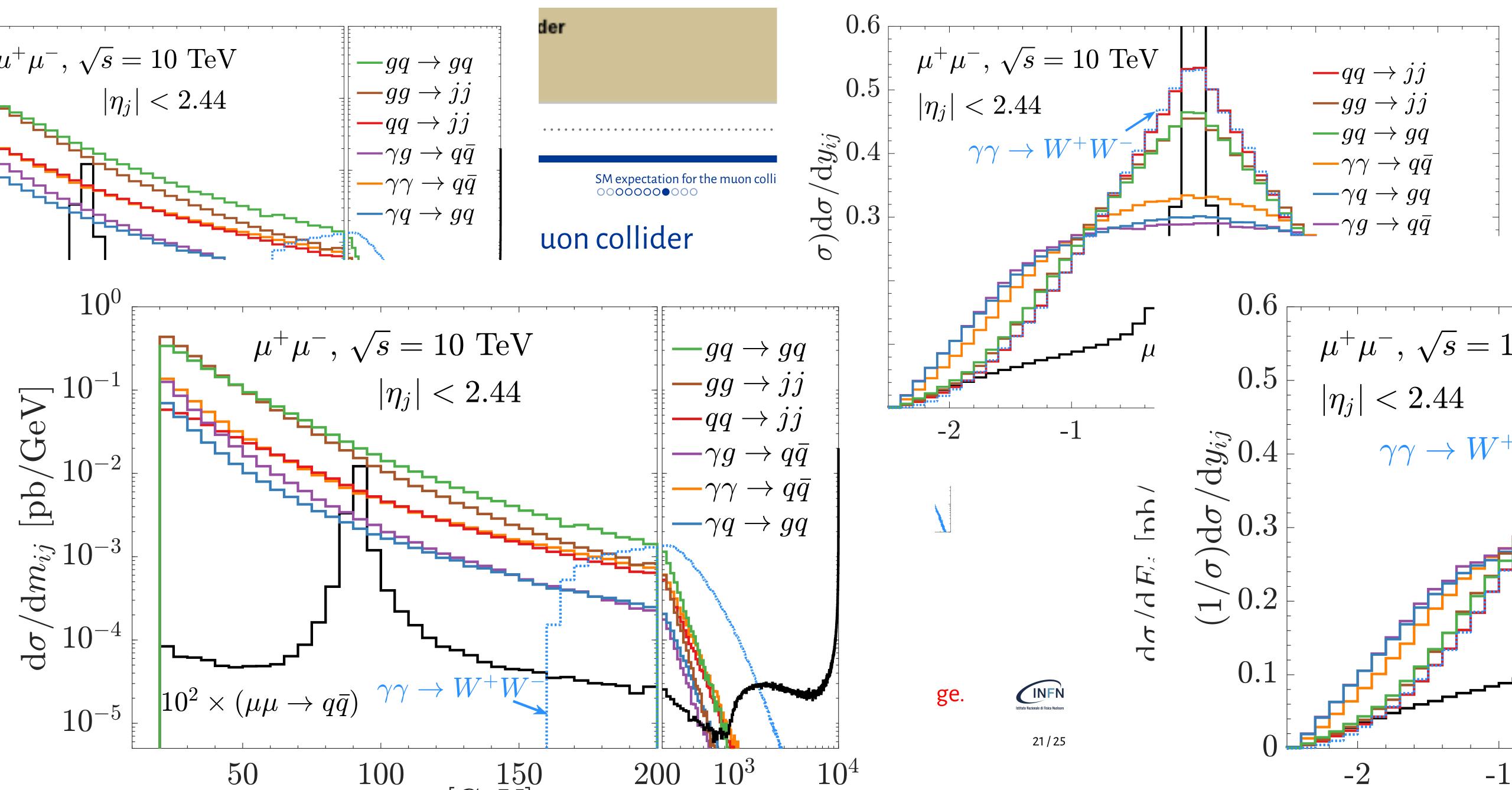
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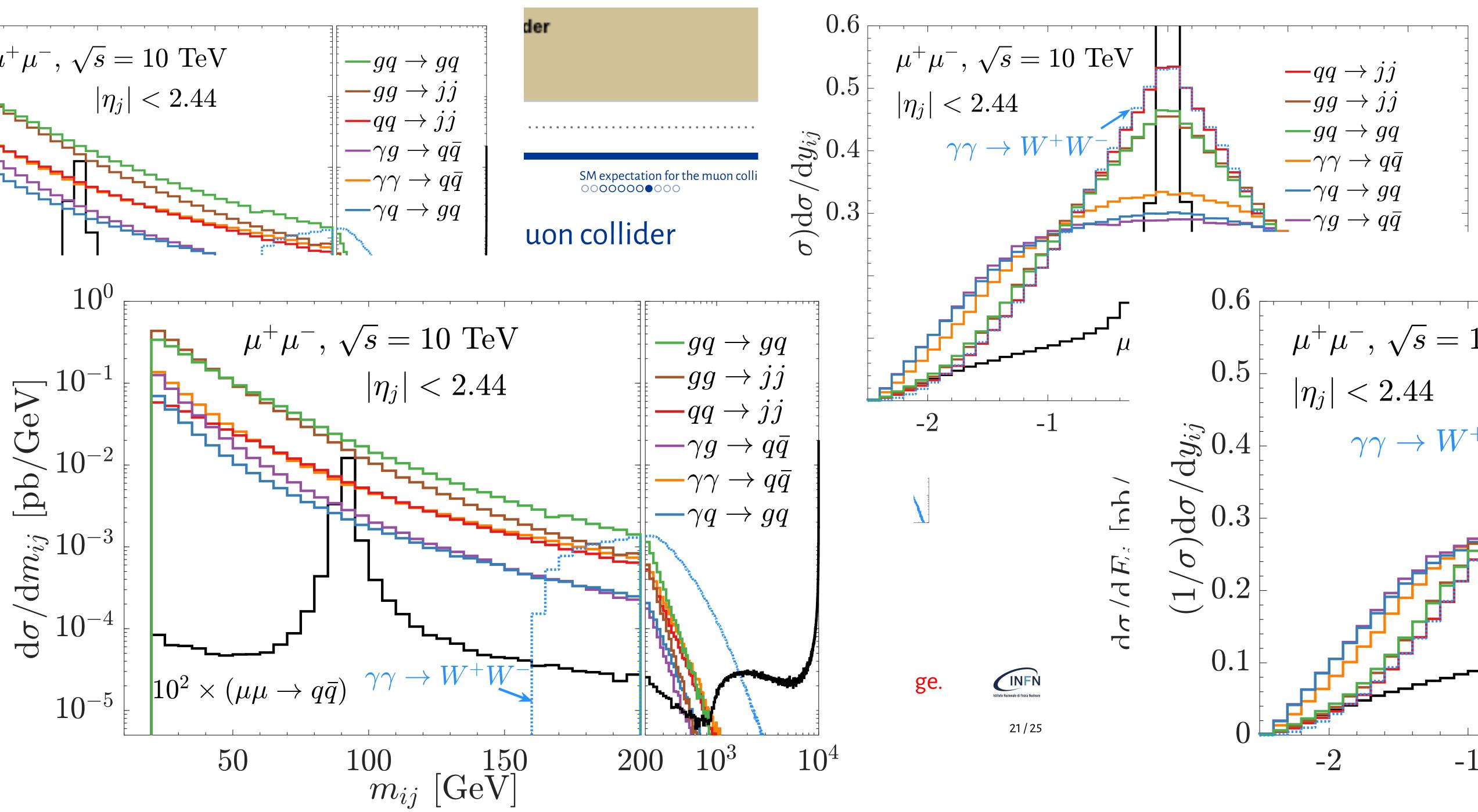
Inclusive jet distributions at a 10 TeV muon collider





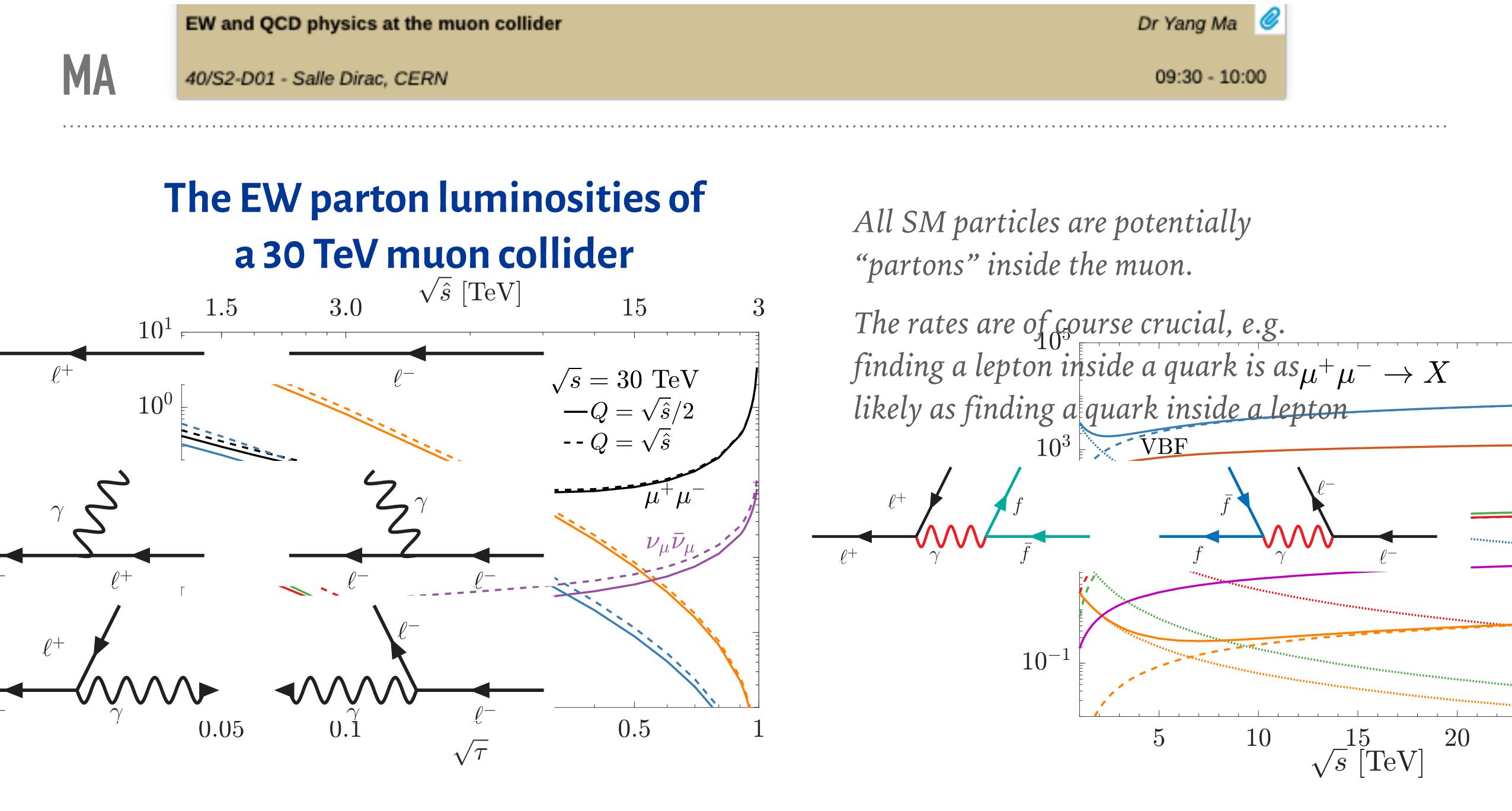








a 30 TeV muon collider





LFUV at muon collider

6/R-012 - conference room, CERN



Admir Greljo





13.10.2022, Muon Collider Collaboration Annual Meeting, CERN





GRELJO

LFUV at muon collider

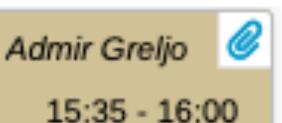
6/R-012 - conference room, CERN

Admir Greljo | LFUV at MuC

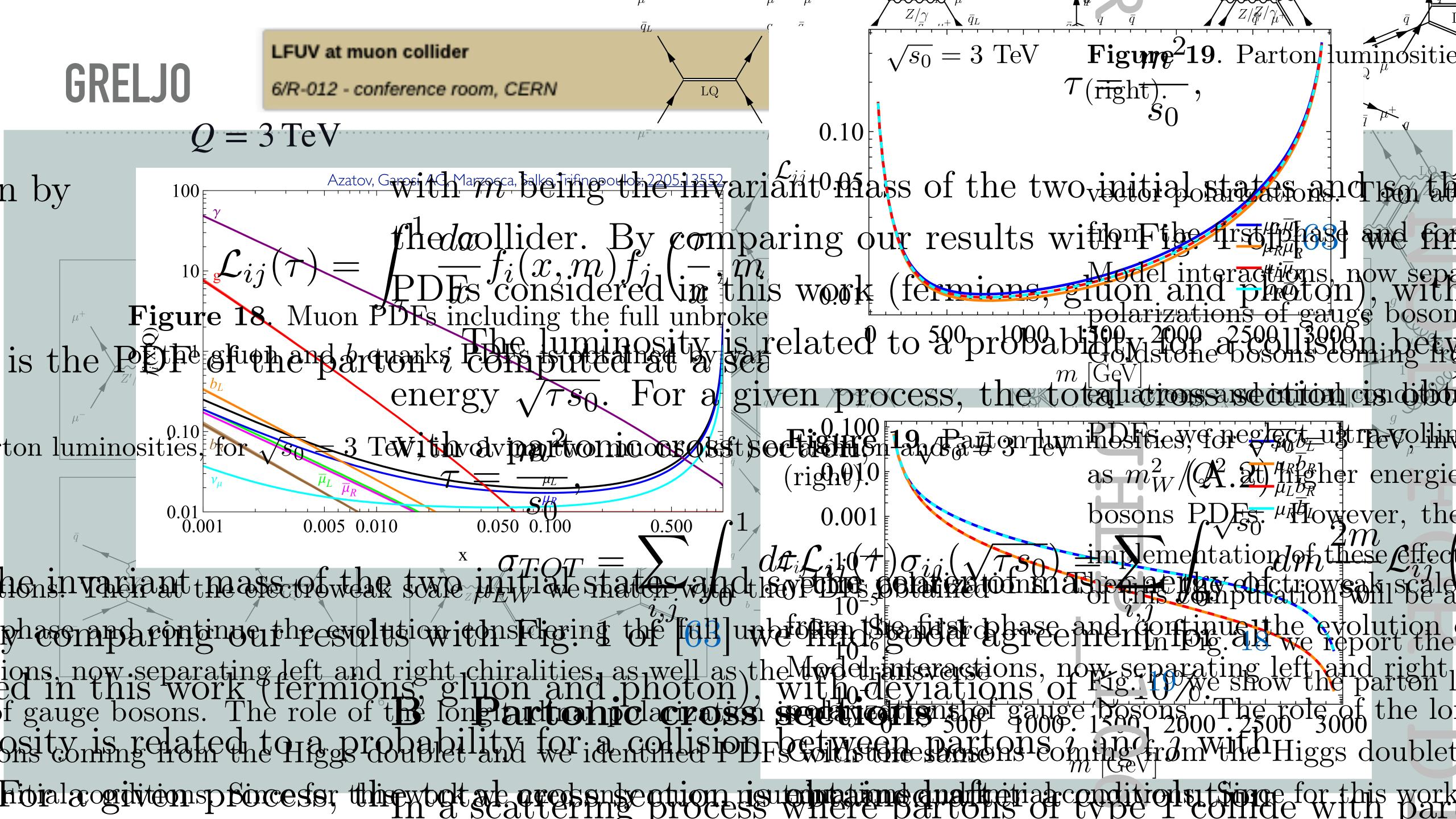
Motivation

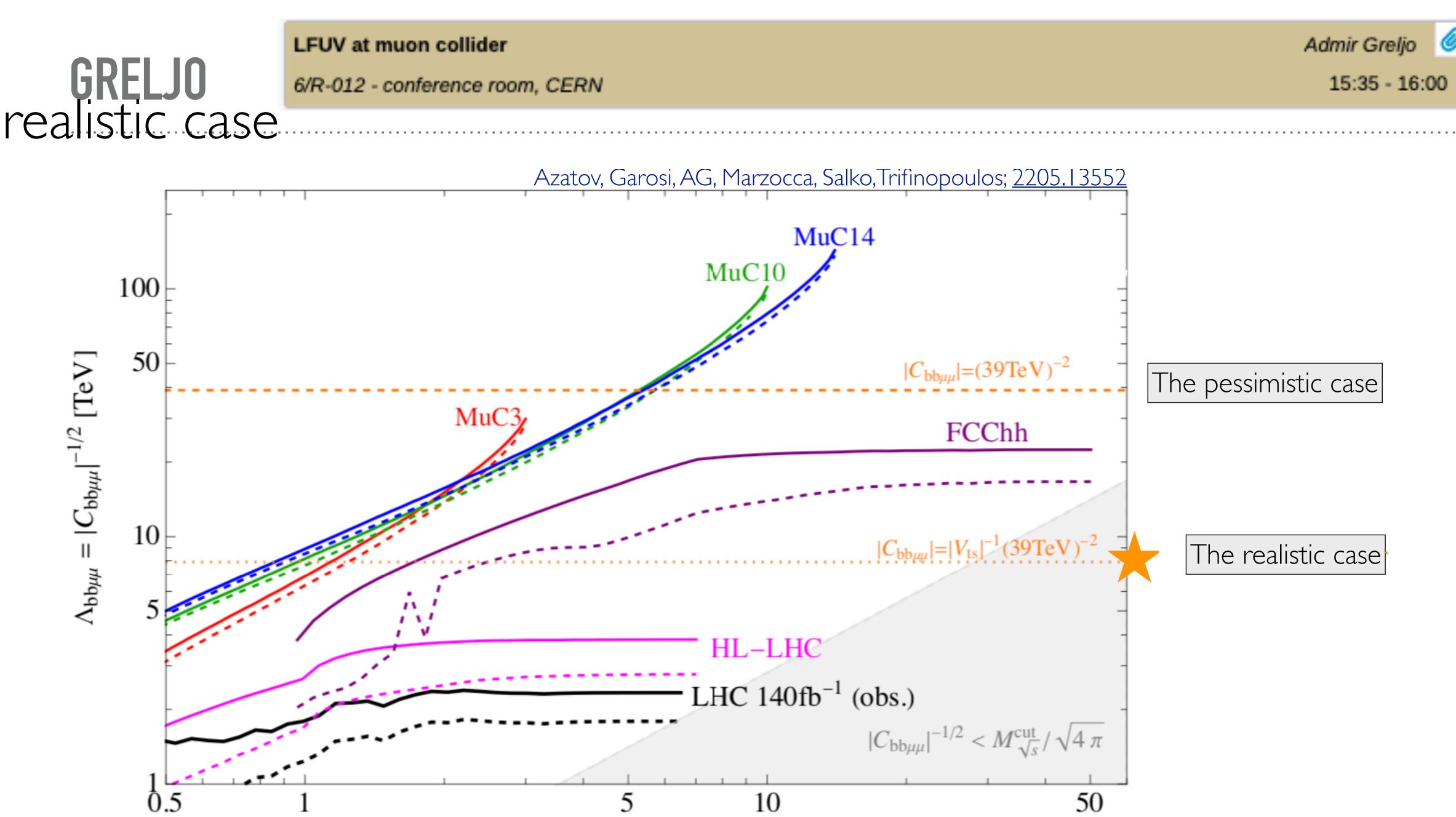
- The LHCb anomalies in $b \rightarrow s\mu^+\mu^-$ decays remind us that — New Physics might take an exotic form an option we should embrace given the present status of the field.
- Several anomalous observables: BRs, angular distributions, LFUV ratios. see LHCb Implications next week
- Coherent explanation by a short-distance $bs\mu\mu$ contact interaction $-\mathcal{O}(10^{-5})G_F$ — the violation of perturbative unitarity $\leq 100 \,\mathrm{TeV}$
- New mass threshold in the vicinity of colliders?
- Today: Azatov, Garosi, AG, Marzocca, Salko, Trifinopoulos; <u>2205.13552</u>

Complementary high- p_T searches at future colliders: FCC-hh versus MuC



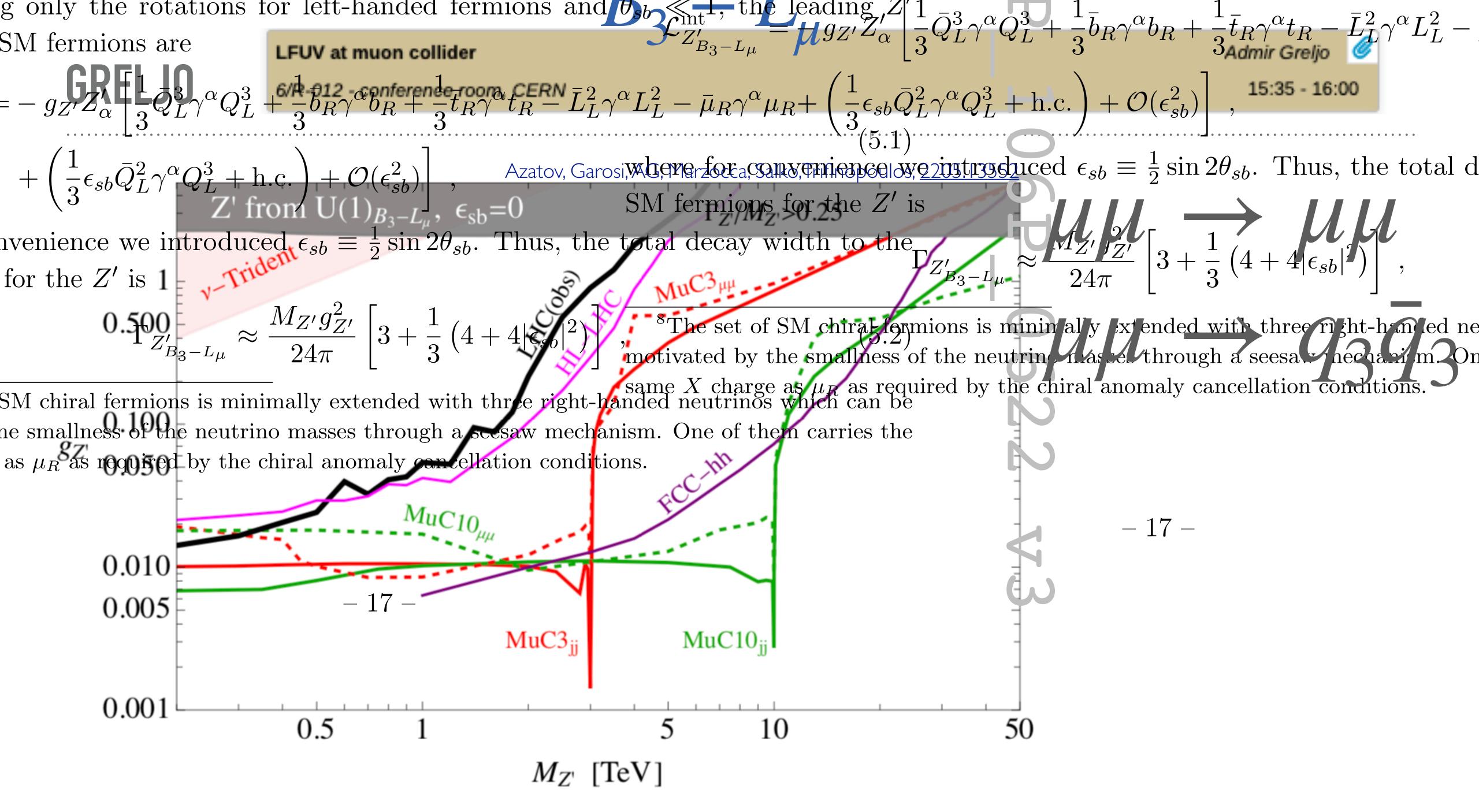


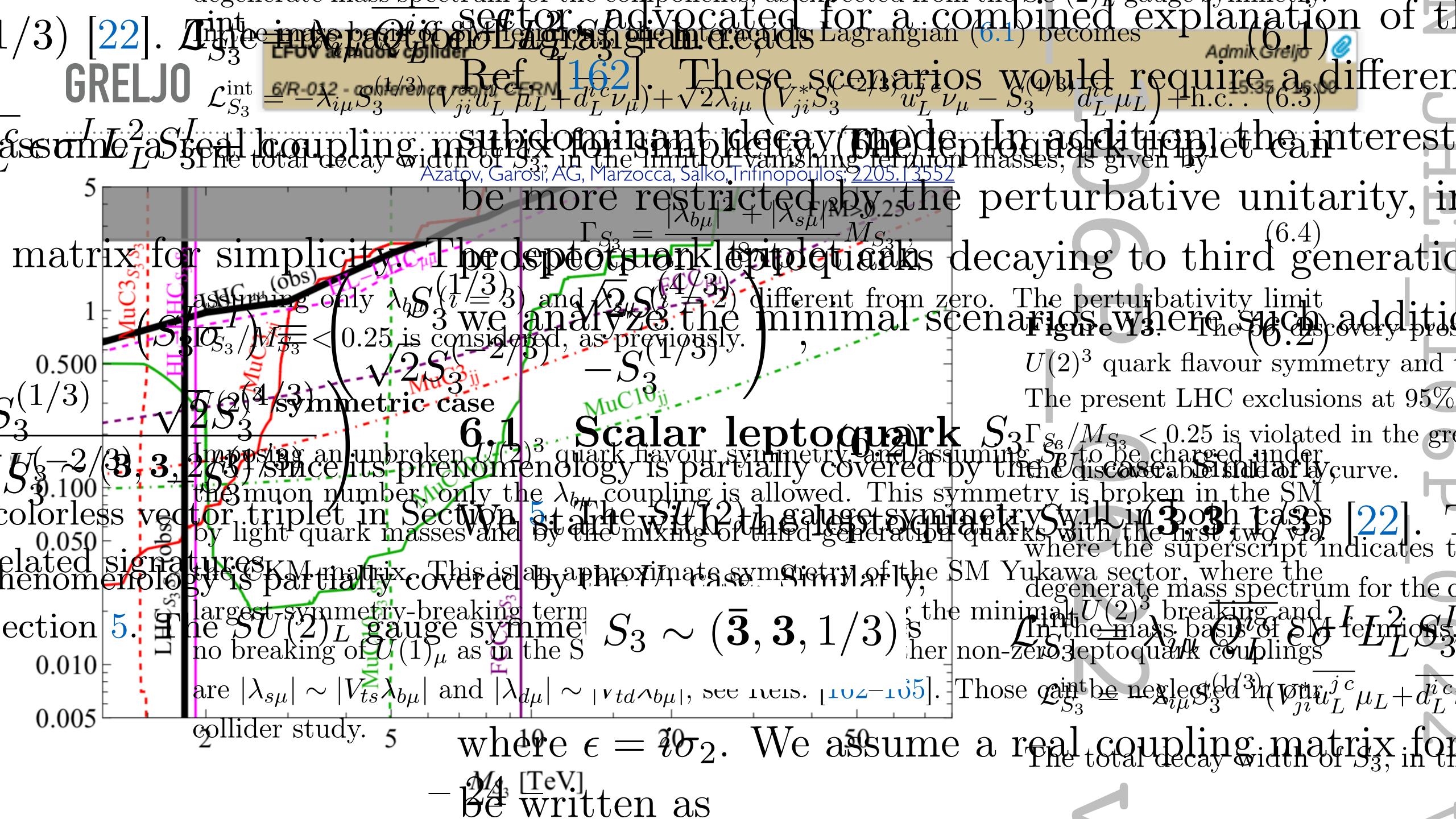




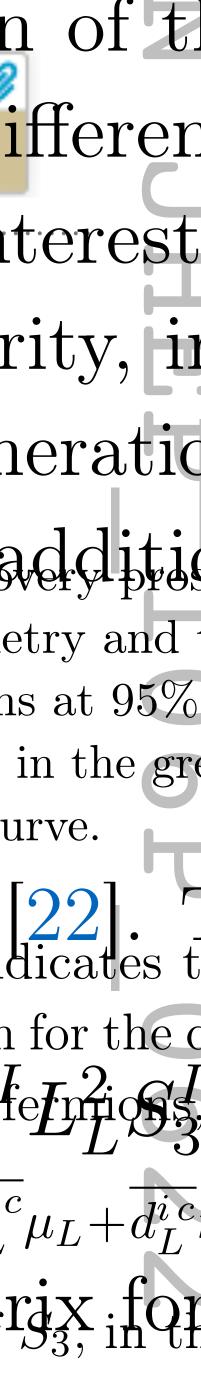
 $M_{\sqrt{s}}^{\text{cut}}$ [TeV]

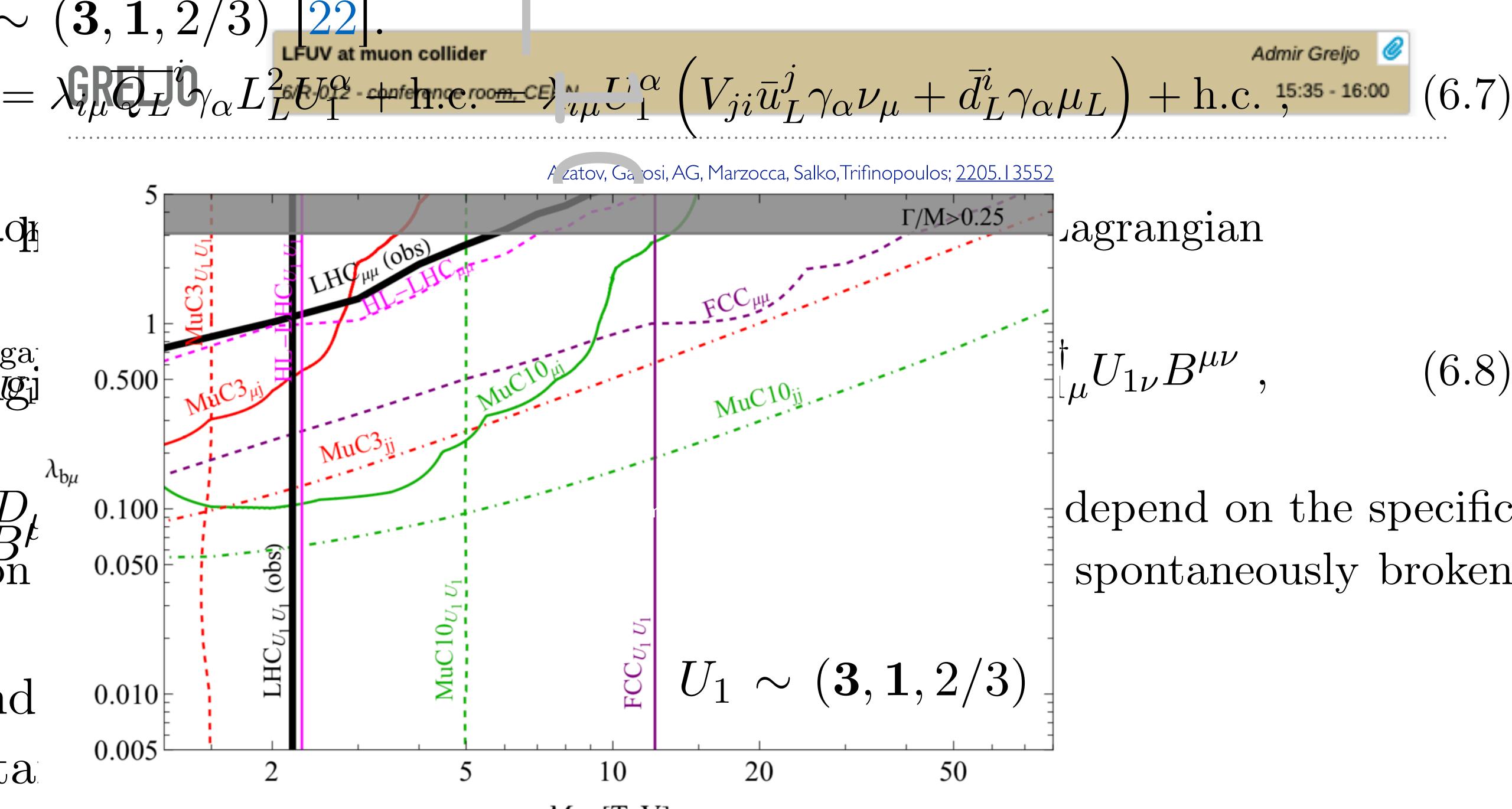
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GRELJO $\mathcal{L}_{S_3}^{\text{int}} = -\chi_{i\mu} S_3^{(1/3)} (V_{ji}^* u_L^i \mu_L + d_L^{ic} \nu_{\mu}) + \sqrt{2} \lambda_{i\mu} (V_{ji}^* S_3^{ic} \mu_L - S_3^{ic} \nu_{\mu}) + \sqrt{2} \lambda_{i\mu} (V_{ji}^* S_3^{ic} \mu_L - S_3^{ic} \nu_{\mu}) + \sqrt{2} \lambda_{i\mu} (V_{ji}^* S_3^{ic} \mu_L - S_3^{ic} \nu_{\mu}) + \sqrt{2} \lambda_{i\mu} (V_{ji}^* S_3^{ic} \mu_L - S_3^{ic} \nu_{\mu}) + \sqrt{2} \lambda_{i\mu} (V_{ji}^* S_3^{ic} \mu_L - S_3^{ic} \mu_L - S_3^{ic} \mu_L) + \sqrt{2} \lambda_{i\mu} (V_{ji}^* S_3^{ic} \mu_L - S_3^{ic} \mu_L - S_3^{ic} \mu_L) + \sqrt{2} \lambda_{i\mu} (V_{ji}^* S_3^{ic} \mu_L - S_3^{ic} \mu_L) + \sqrt{2} \lambda_{i\mu} (V_{ji}^* S_3^{ic} \mu_L - S_3^{ic} \mu_L) + \sqrt{2} \lambda_{i\mu} (V_{ji}^* S_3^{ic} \mu_L - S_3^{ic} \mu_L) + \sqrt{2} \lambda_{i\mu} (V_{ji}^* S_3^{ic} \mu_L - S_3^{ic} \mu_L) + \sqrt{2} \lambda_{i\mu} (V_{ji}^* S_3^{ic} \mu_L - S_3^{ic} \mu_L) + \sqrt{2} \lambda_{i\mu} (V_{ji}^* S_3^{ic} \mu_L - S_3^{ic} \mu_L) + \sqrt{2} \lambda_{i\mu} (V_{ji}^* S_3^{ic} \mu_L - S_3^{ic} \mu_L) + \sqrt{2} \lambda_{i\mu} (V_{ji}^* S_3^{ic} \mu_L - S_3^{ic} \mu_L) + \sqrt{2} \lambda_{i\mu} (V_{ji}^* S_3^{ic} \mu_L - S_3^{ic} \mu_L) + \sqrt{2} \lambda_{i\mu} (V_{ji}^* S_3^{ic} \mu_L - S_3^{ic} \mu_L) + \sqrt{2} \lambda_{i\mu} (V_{ji}^* S_3^{ic} \mu_L - S_3^{ic} \mu_L) + \sqrt{2} \lambda_{i\mu} (V_{ji}^* S_3^{ic} \mu_L - S_3^{ic} \mu_L) + \sqrt{2} \lambda_{i\mu} (V_{ji}^* S_3^{ic} \mu_L - S_3^{ic} \mu_L) + \sqrt{2} \lambda_{i\mu} (V_{ji}^* S_3^{ic} \mu_L - S_3^{ic} \mu_L) + \sqrt{2} \lambda_{i\mu} (V_{ji}^* S_3^{ic} \mu_L - S_3^{ic} \mu_L) + \sqrt{2} \lambda_{i\mu} (V_{ji}^* S_3^{ic} \mu_L - S_3^{ic} \mu_L) + \sqrt{2} \lambda_{i\mu} (V_{ji}^* S_3^{ic} \mu_L - S_3^{ic} \mu_L) + \sqrt{2} \lambda_{i\mu} (V_{ji}^* S_3^{ic} \mu_L - S_3^{ic} \mu_L) + \sqrt{2} \lambda_{i\mu} (V_{ji}^* S_3^{ic} \mu_L - S_3^{ic} \mu_L) + \sqrt{2} \lambda_{i\mu} (V_{ji}^* S_3^{ic} \mu_L - S_3^{ic} \mu_L) + \sqrt{2} \lambda_{i\mu} (V_{ji}^* S_3^{ic} \mu_L - S_3^{ic} \mu_L) + \sqrt{2} \lambda_{i\mu} (V_{ji}^* S_3^{ic} \mu_L - S_3^{ic} \mu_L) + \sqrt{2} \lambda_{i\mu} (V_{ji}^* S_3^{ic} \mu_L - S_3^{ic} \mu_L) + \sqrt{2} \lambda_{i\mu} (V_{ji}^* \mu_L)$ be more restricted by the perturbative unitarity, is matrix for simplicity The dependent on the same decaying to third generation assume only $\lambda_b S(\frac{1}{3})$ and $\lambda_b S(\frac{4}{3})$ and $\lambda_b S(\frac{4}{3})$ different from zero. The perturbativity limit 3 We and $\sqrt{2}S(\frac{4}{3})$ different from zero. The perturbativity limit 0.25, 3 We and $\sqrt{2}S(\frac{4}{3})$ different from zero. The perturbativity limit 0.25, 3 We and $\sqrt{2}S(\frac{4}{3})$ different from zero. The perturbativity limit 0.25, 3 We and $\sqrt{2}S(\frac{4}{3})$ different from zero. The perturbativity limit 0.25, 3 We and $\sqrt{2}S(\frac{4}{3})$ different from zero. The perturbativity limit 0.25, 3 We and $\sqrt{2}S(\frac{4}{3})$ different from zero. The perturbativity limit 0.25, 3 We and $\sqrt{2}S(\frac{4}{3})$ different from zero. The perturbativity limit 0.25, 3 We and $\sqrt{2}S(\frac{4}{3})$ different from zero. The perturbativity limit 0.25 difference of the provided of the provide $U(2)^3$ quark flavour symmetry and The present LHC exclusions at 95% $\frac{3}{5} - \frac{3}{3}$, $\frac{3}{3}$, $\frac{3}{5}$, $\frac{3}{5}$, $\frac{6.1}{3}$, $\frac{6.1}{3}$, $\frac{3}{5}$, $\frac{3}{5}$, $\frac{10}{5}$, $\frac{3}{5}$, $\frac{3}{$ the muon number, only the $\lambda_{b\mu}$ coupling is allowed. This symmetry is broken in the SN tor, triplet in Section of the with 2 th sates symmetry will in Soch task by light quark masses and by the mixing of third september of quarks with the first two variables of the section of the first two variables are by the mixing of third september of quarks with the first two variables are by the mixing of third september of quarks with the first two variables are by the mixing of third september of the first two variables are by the mixing of third september of the first two variables are by the mixing of third september of the first two variables are by the first two variables are by the mixing of third september of the first two variables are by the first two variables her non-zero dept are $|\lambda_{s\mu}| \sim |V_{ts}\lambda_{b\mu}|$ and $|\lambda_{d\mu}| \sim |v_{td}\lambda_{b\mu}|$, see news. [102 - 105]. Those \mathfrak{B}_{S_3} is new level $\mathfrak{B}_{s\mu}$ and $|\lambda_{d\mu}| \sim |v_{td}\lambda_{b\mu}|$, see news. [102 - 105]. Those \mathfrak{B}_{S_3} is new level $\mathfrak{B}_{s\mu}$ and $|\lambda_{d\mu}| \sim |v_{td}\lambda_{b\mu}|$, see news. [102 - 105]. where $\epsilon = i \sigma_2$. We assume a real coupling matrix if q_3 , q_4





 M_{U_1} [TeV]



Dark Matter at muon colliders

40/S2-D01 - Salle Dirac, CERN

Dark matter at the muon colliders

Xiaoran Zhao¹

In collaboration with R. Franceschini

Based on arXiv:22xx.xxxx

¹Dipartimento di Matematica e Fisica, Università di Roma Tre and INFN, sezione di Roma Tre, I-00146 Rome, Italy

October 12, 2022

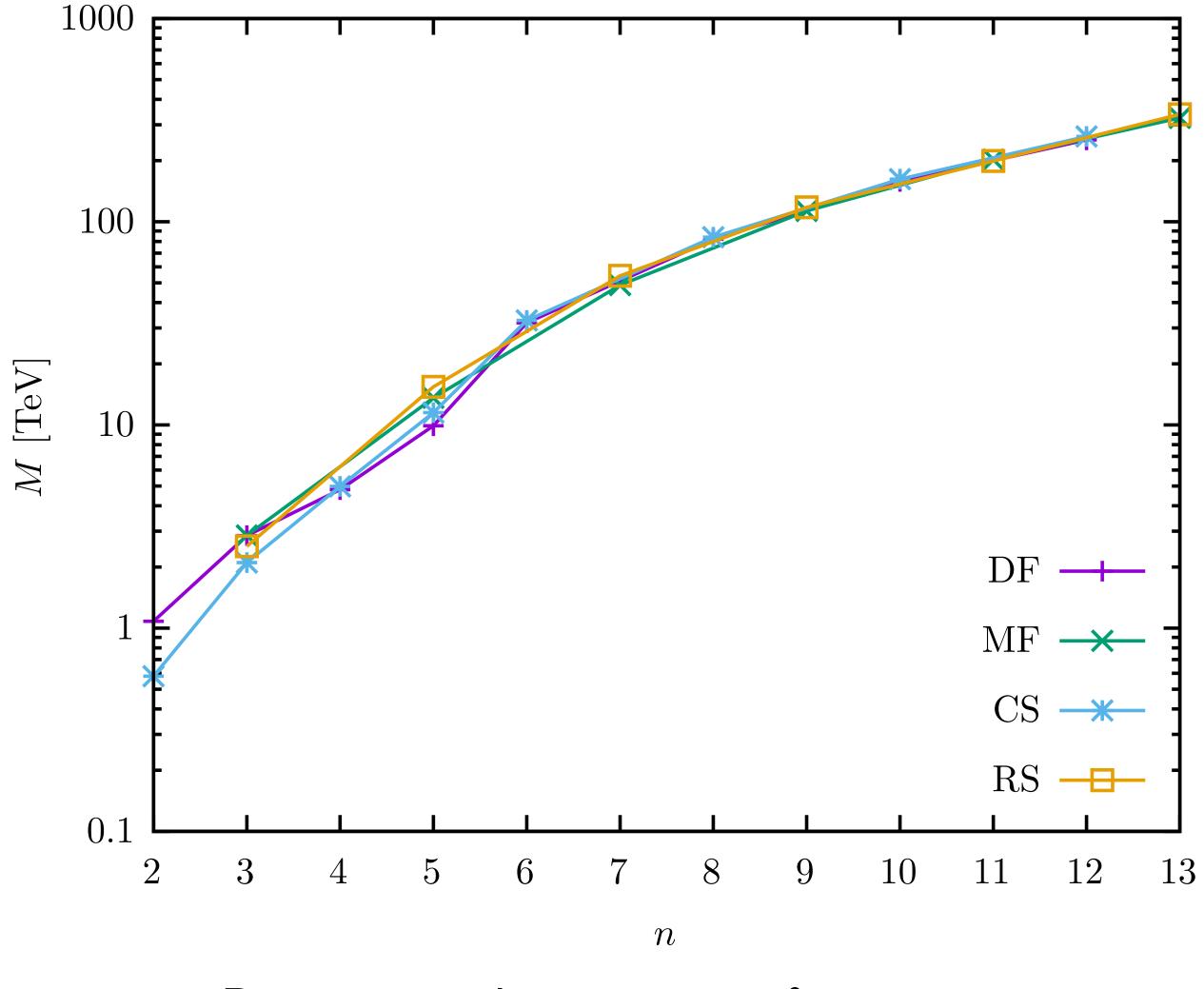




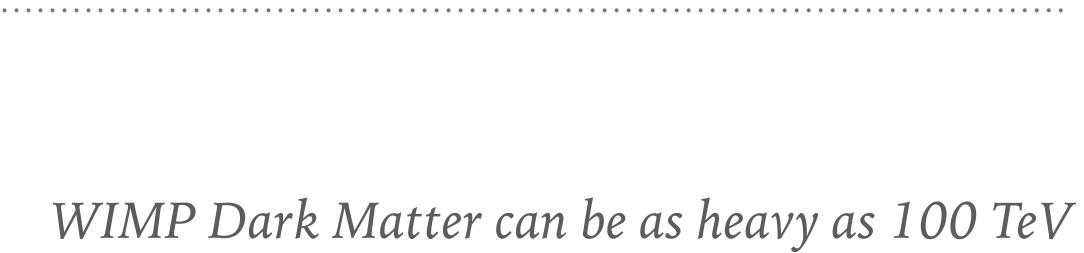


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Bottaro et.al.2107.09688&2205.04486

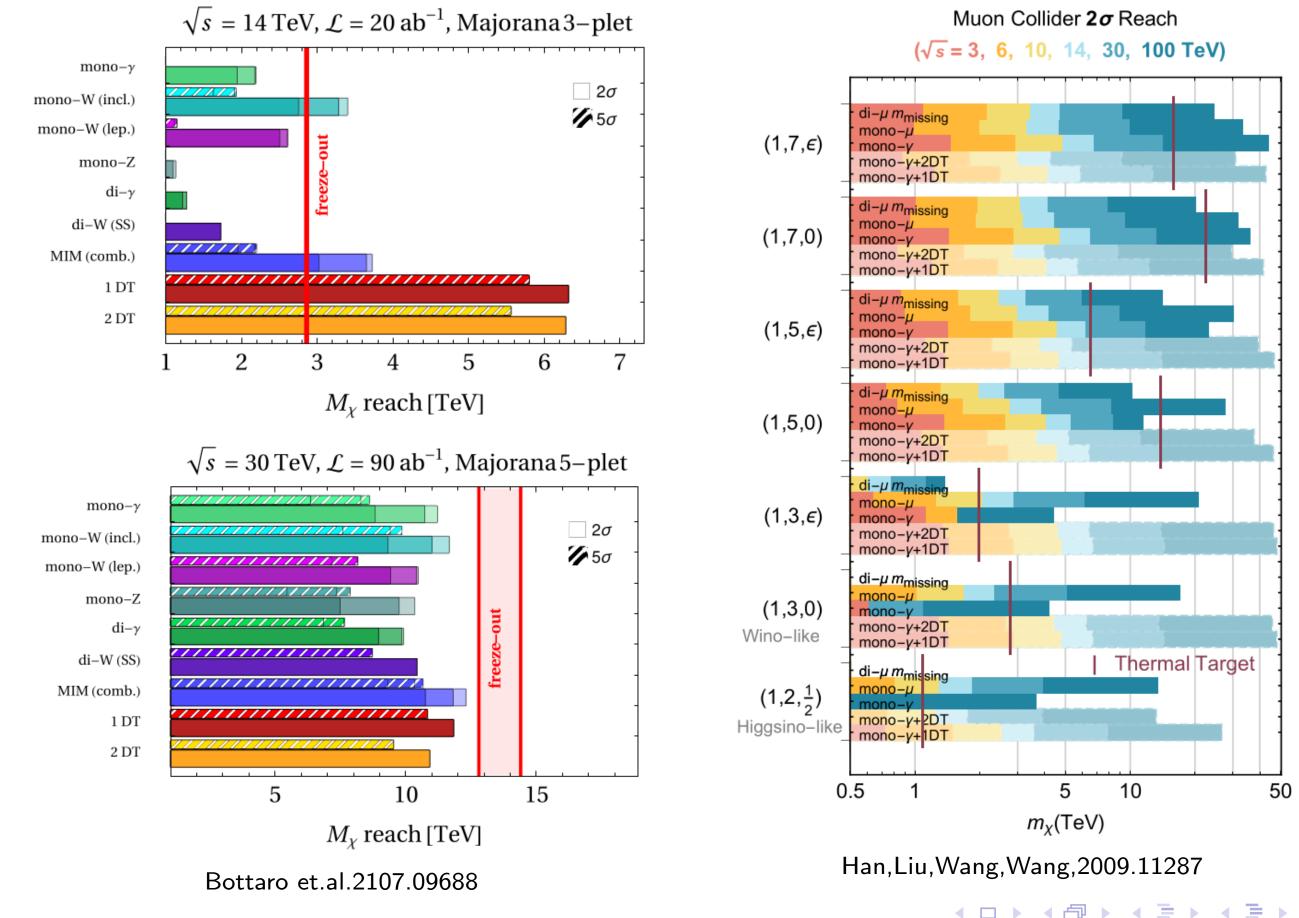


ZHAO

Dark Matter at muon colliders

40/S2-D01 - Salle Dirac, CERN

direct production at muon colliders







ZHAO

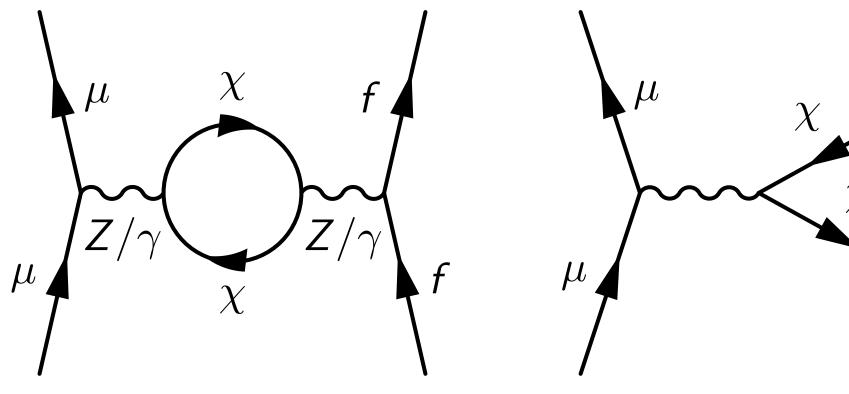
Dark Matter at muon colliders

40/S2-D01 - Salle Dirac, CERN

Indirect probes

• Direct production is limited by $M < \frac{\sqrt{s}}{2}$.

Indirect probes through loop corrections: no such limit!





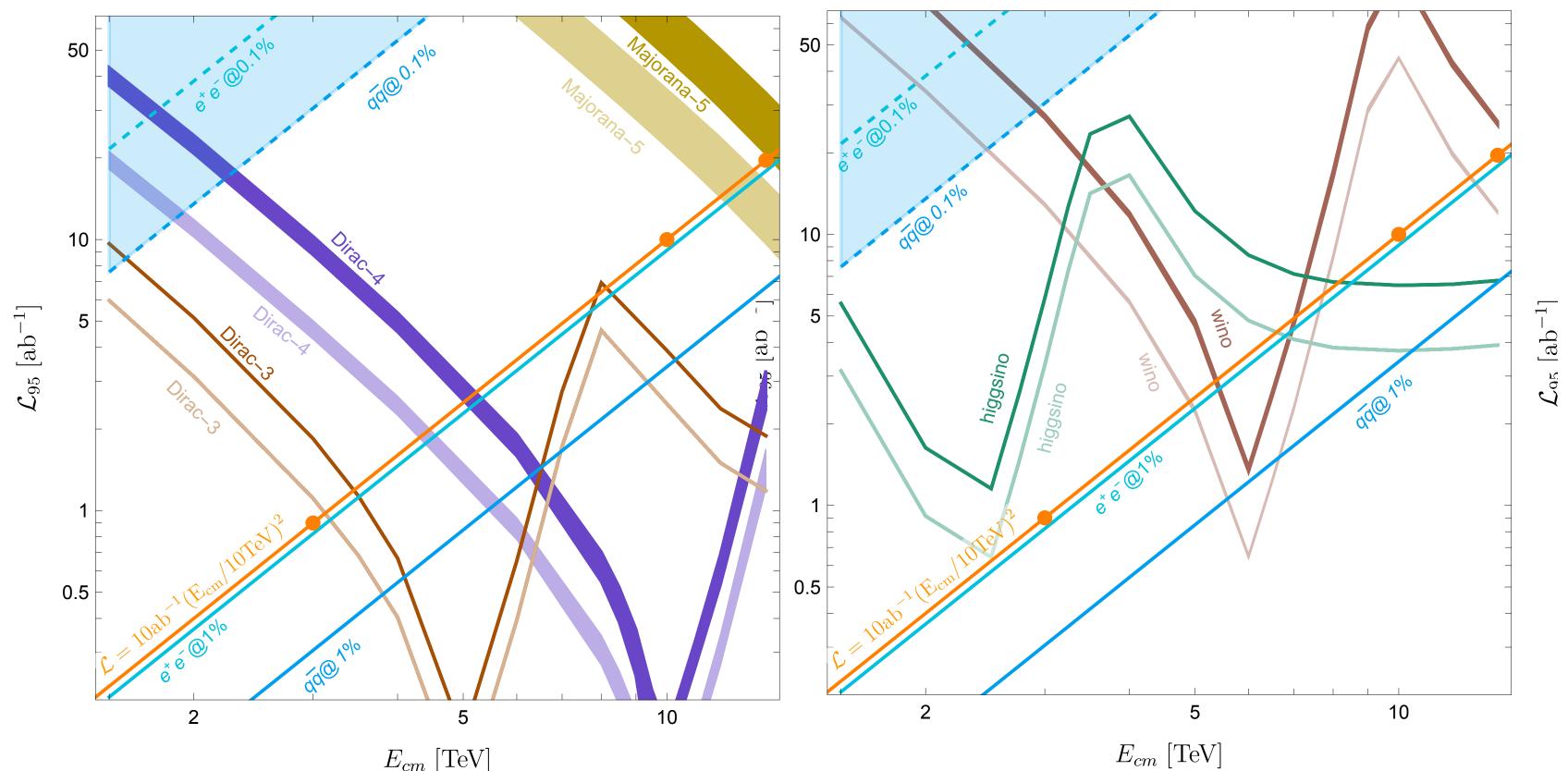
 \mathcal{N} $\land \chi$ χ \mathcal{W}



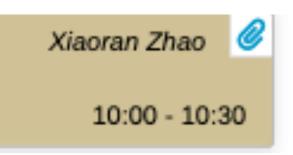
ZHAO

Dark Matter at muon colliders

40/S2-D01 - Salle Dirac, CERN

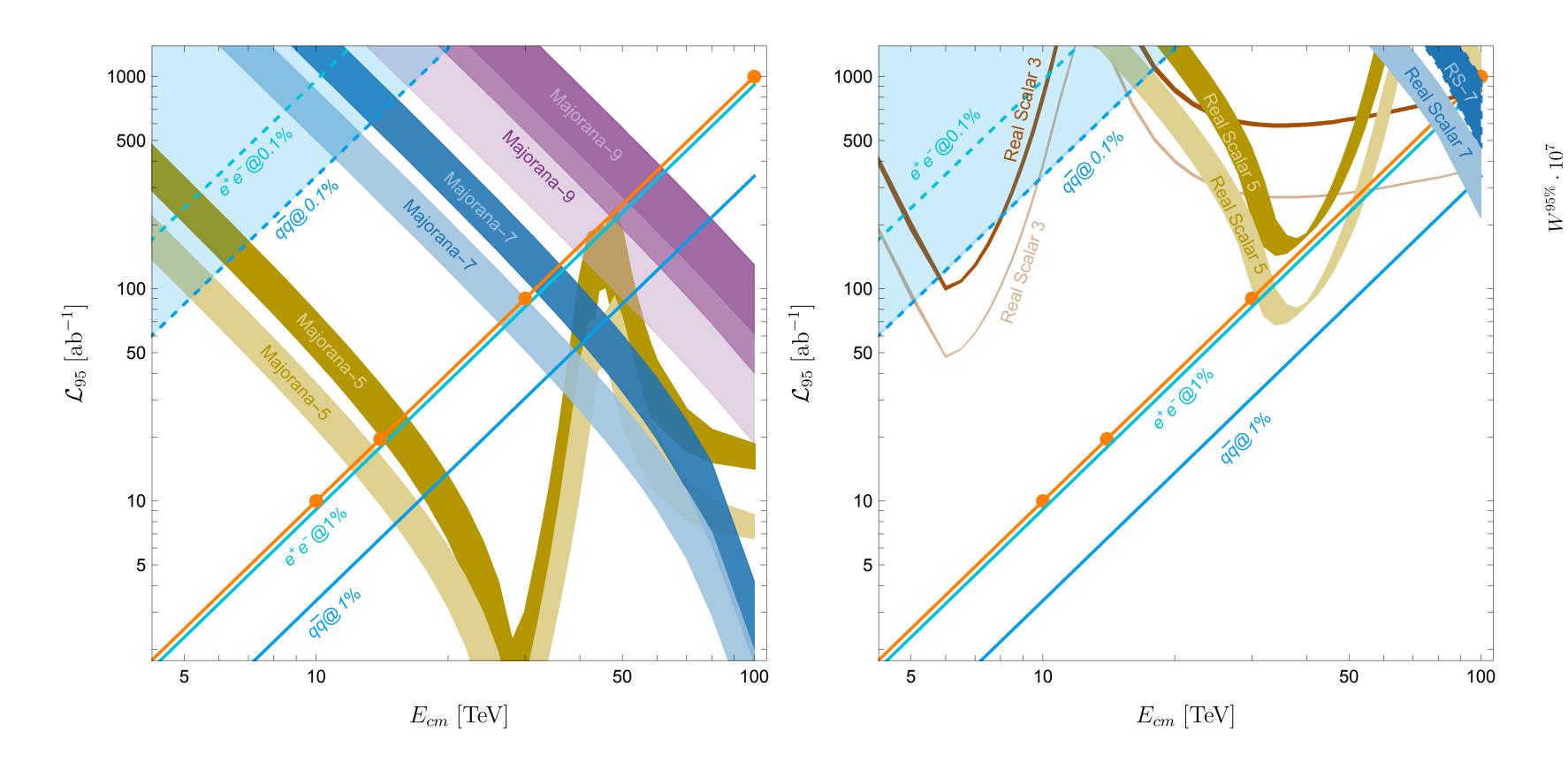


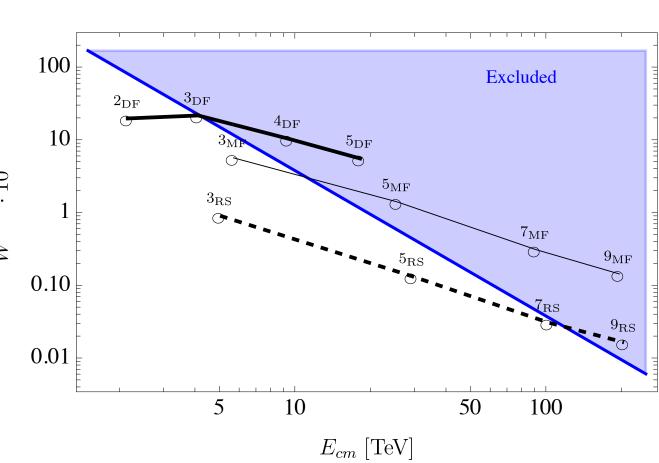
 E_{cm} [TeV]



All fermionic dark matter candidates with $n \leq 5$ ($M \lesssim 15$ TeV) can be excluded for some center-ofmass energy at or below 14 TeV







All dark matter candidates up to the maximal WIMP mass $M \simeq 100$ TeV can be excluded for some center-ofmass energy

CESAROTTI

Physics opportunities with a muon "beam dump"

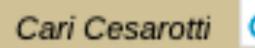
6/R-012 - conference room, CERN

Searching for New Physics at **High-Energy Future Muon Colliders**

Cari Cesarotti, MIT CTP Fellow **IMCC CERN Meeting** Oct 13, 2022







10:00 - 10:30

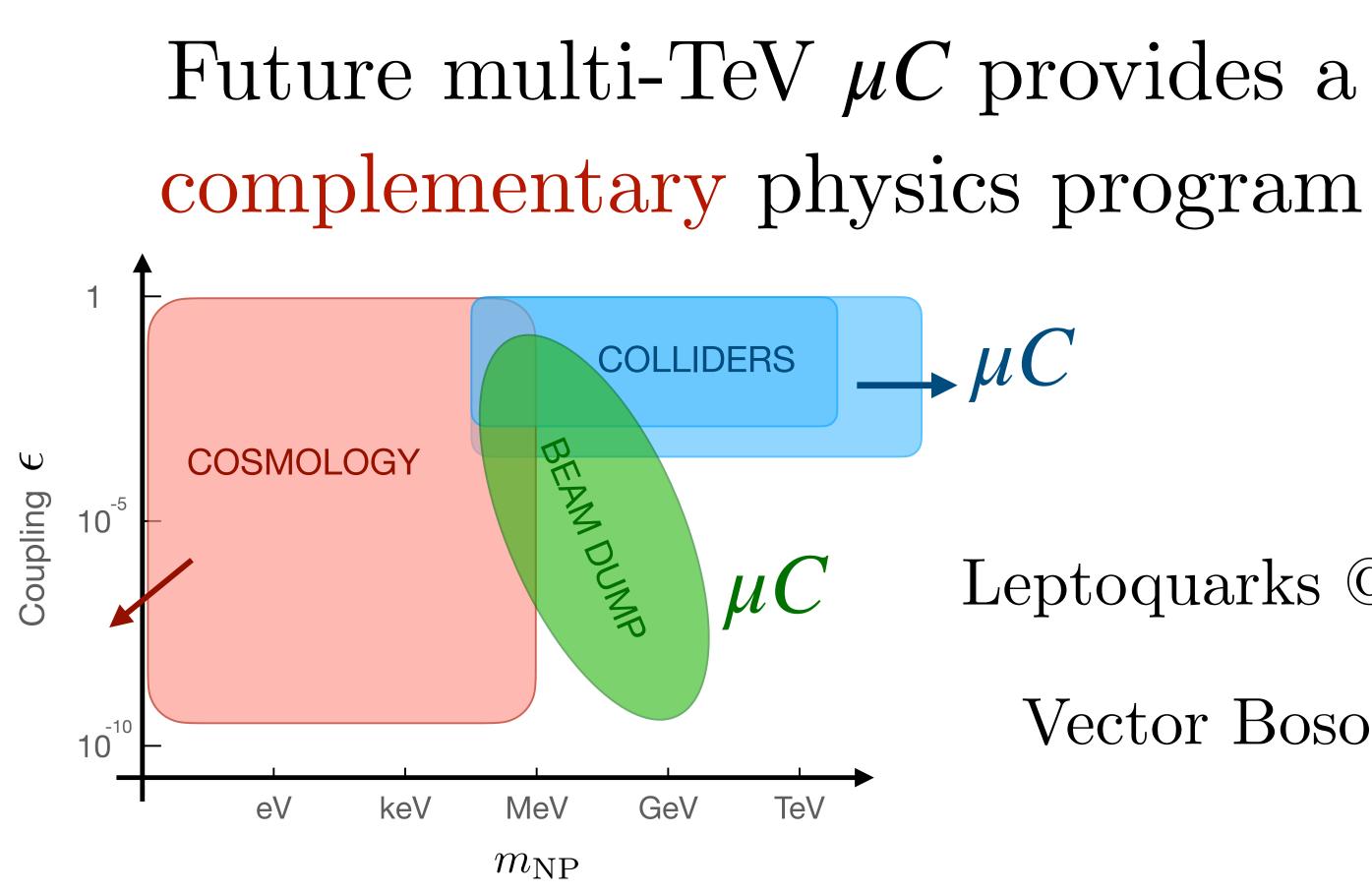
(2104.05720) w/ P. Asadi, R. Capdevilla, S. Homiller (2202.12302) w/S. Homiller, R. Mishra, M. Reece



CESAROTTI

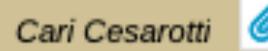
Physics opportunities with a muon "beam dump"

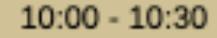
6/R-012 - conference room, CERN



2104.05720 CC, P. Asadi, R Capdevilla, S. Homiller







μĊ

Leptoquarks @ Collider Vector Bosons @ BD

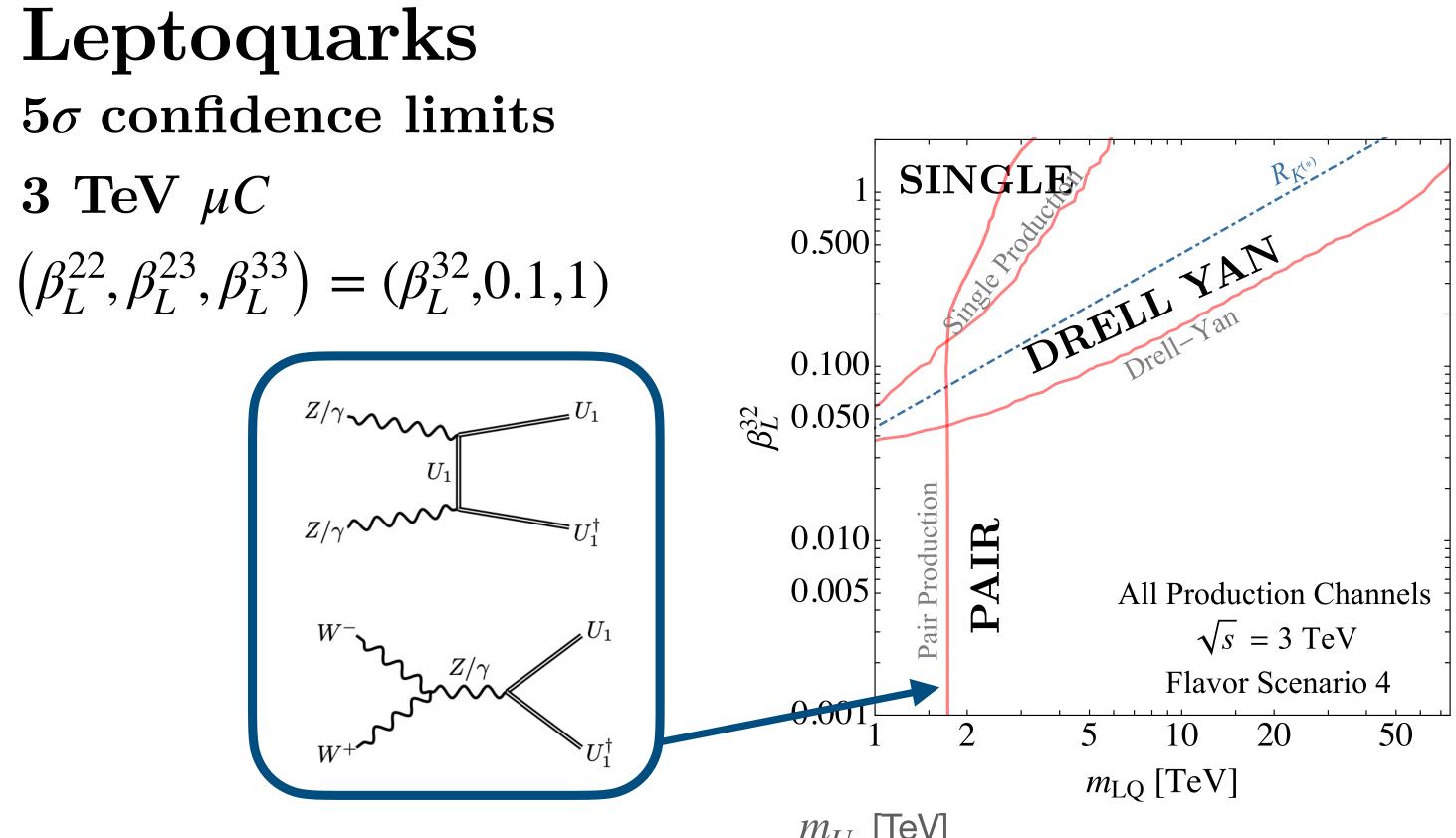
2202.12302 CC, S. Homiller, R. Mishra, M. Reece



Physics opportunities with a muon "beam dump"

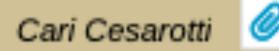
CESAROTTI

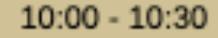
6/R-012 - conference room, CERN



2104.05720 CC, P. Asadi, R
 Capdevilla, S. Homiller $m_{U_1}\,[\mbox{TeV}]$







Similar results to Greljo



CESAROTTI

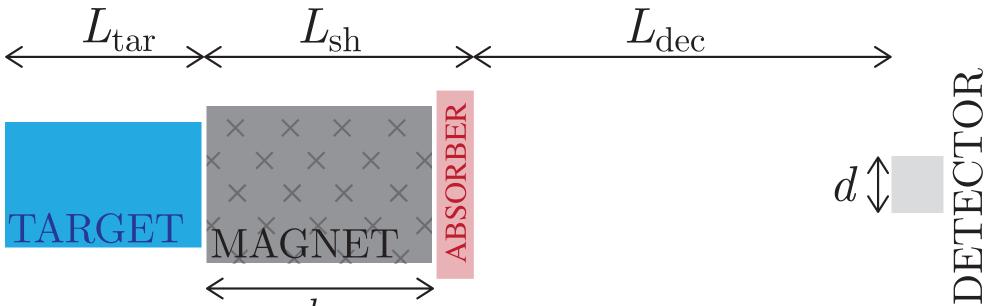
Physics opportunities with a muon "beam dump"

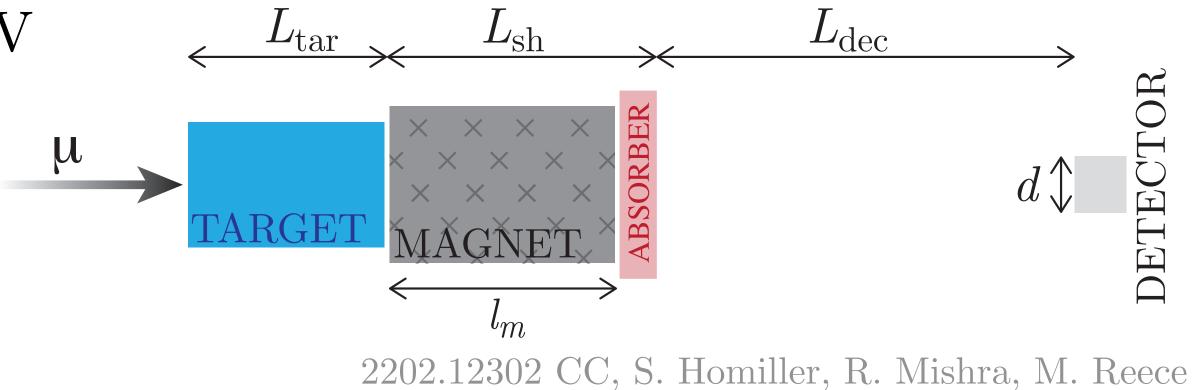
6/R-012 - conference room, CERN

A beam dump experiment at the μ C allows us to push into both the energy and the intensity frontier

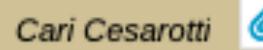
Can probe NP scenarios with:

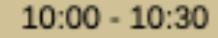
- Very weak couplings
- Couplings to 2nd gen. leptons
- Masses $\leq 100 \text{ GeV}$











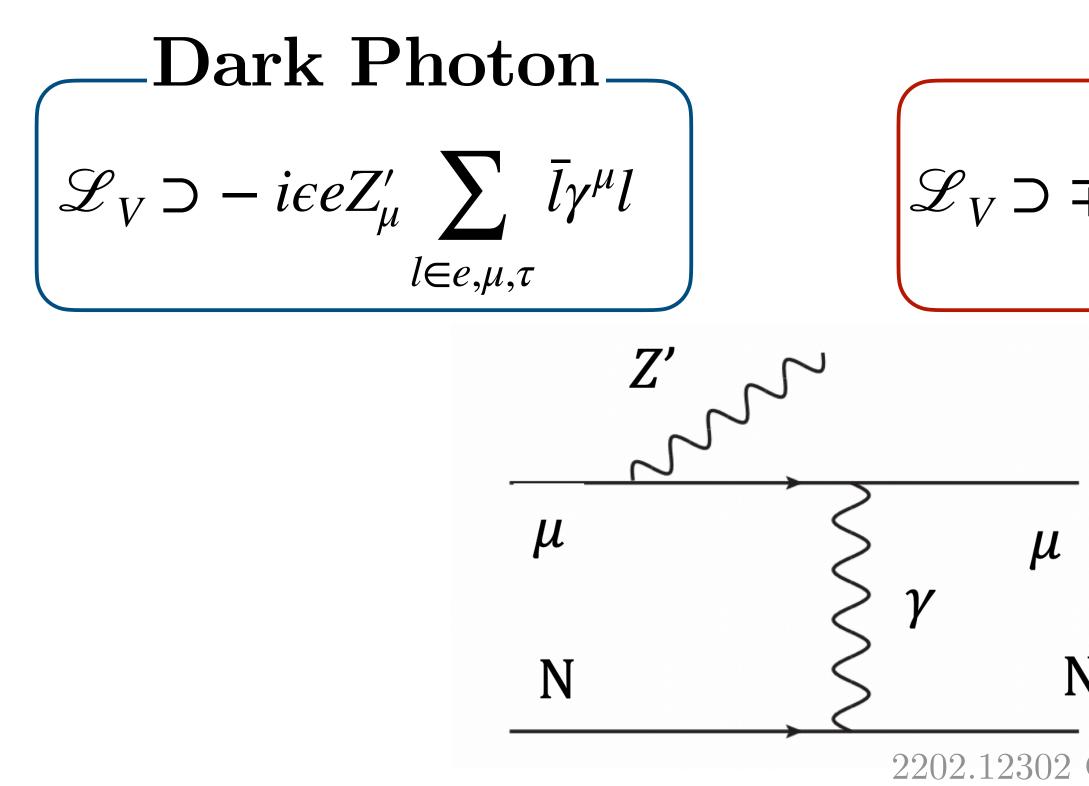


Physics opportunities with a muon "beam dump"

CESAROTTI

6/R-012 - conference room, CERN

New physics Z' Scenarios









 $\mathcal{L}_V \supset \mp i g Z'_{\mu} \sum \left(\bar{l} \gamma^{\mu} l + \bar{\nu}_l \sigma^{\mu} \nu_l \right)$ $l \in \mu, \tau$

2202.12302 CC, S. Homiller, R. Mishra, M. Reece

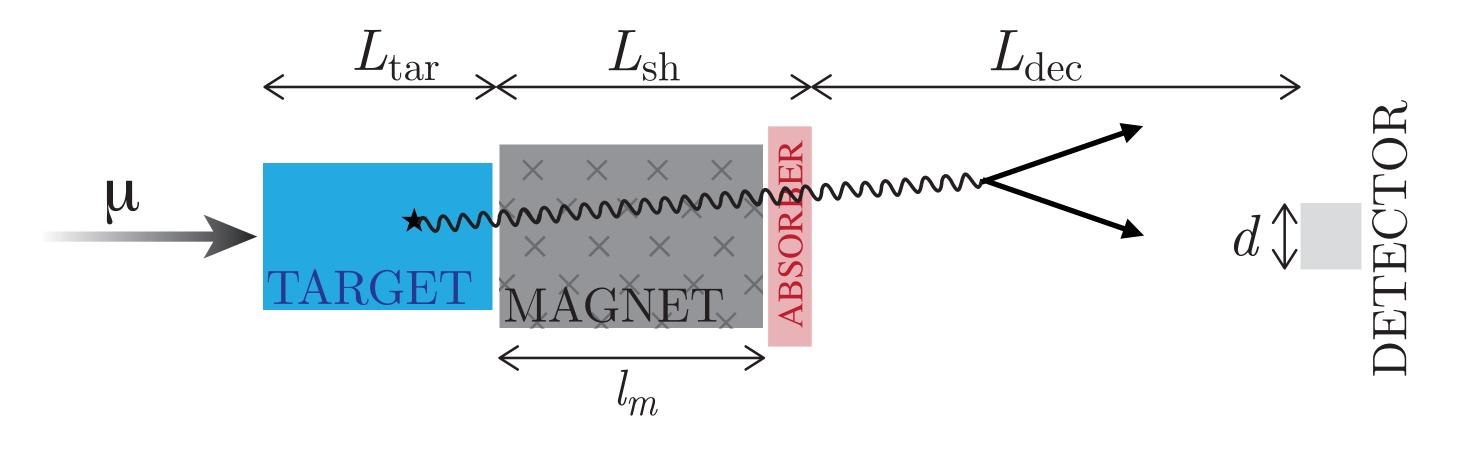


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Physics opportunities with a muon "beam dump"

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Beam Dump Setup

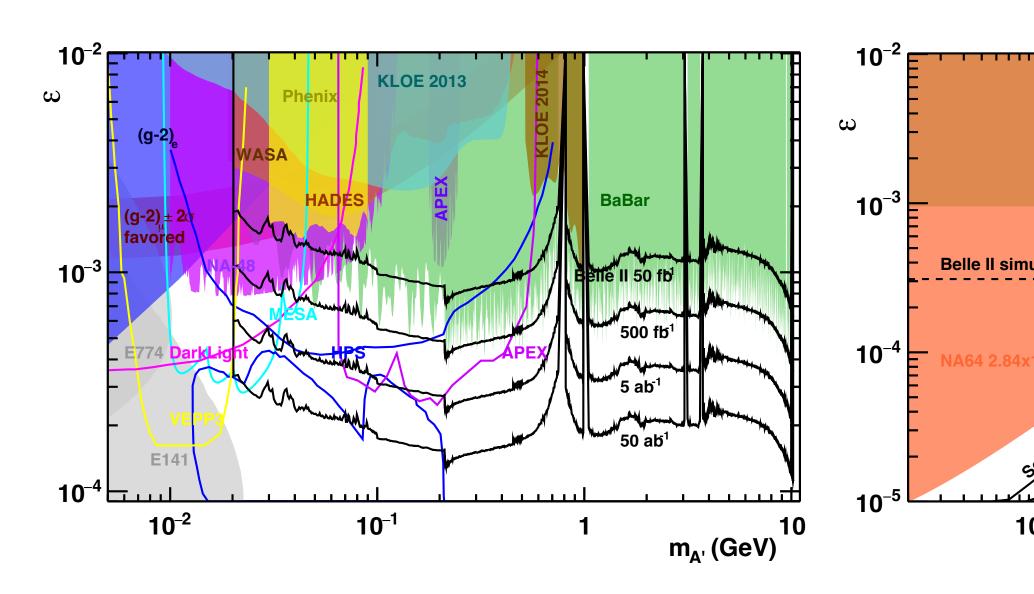


2202.12302 CC, S. Homiller, R. Mishra, M. Reece





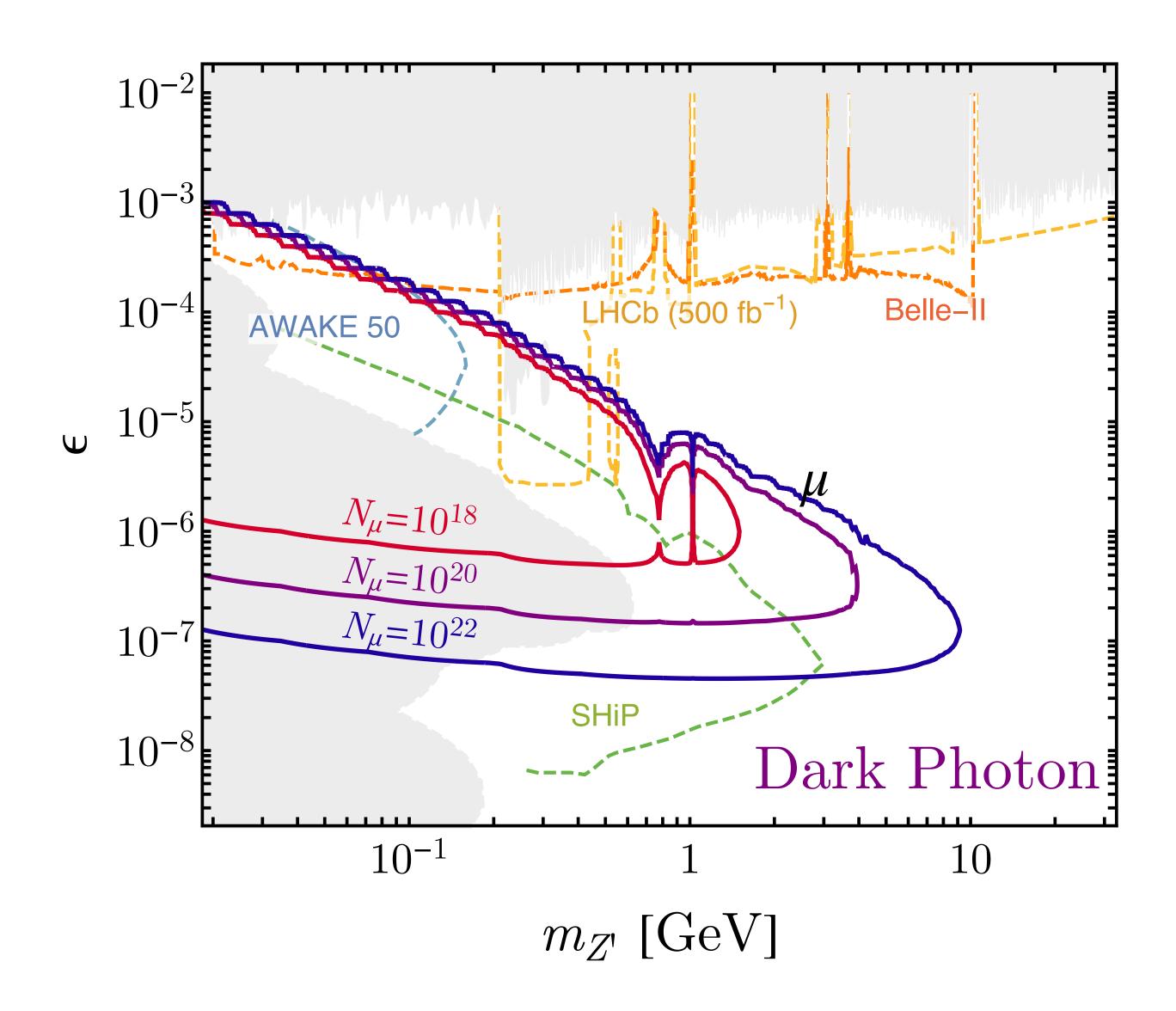
"dark photon" search here means: "shine through wall" experiment, also known as "thick target" beam dump visible search



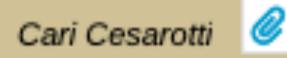
CESAROTTI

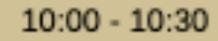
Physics opportunities with a muon "beam dump"

6/R-012 - conference room, CERN



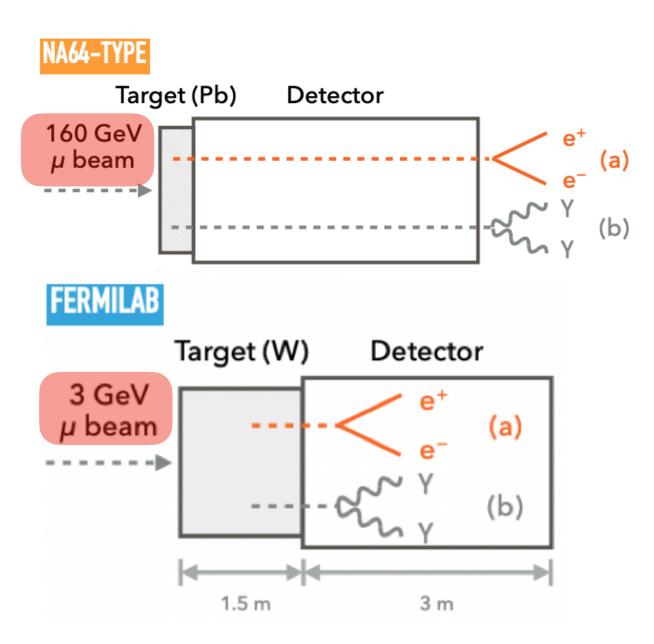






"dark photon" search here means: "shine through wall" experiment, also known as "thick target" beam dump visible search

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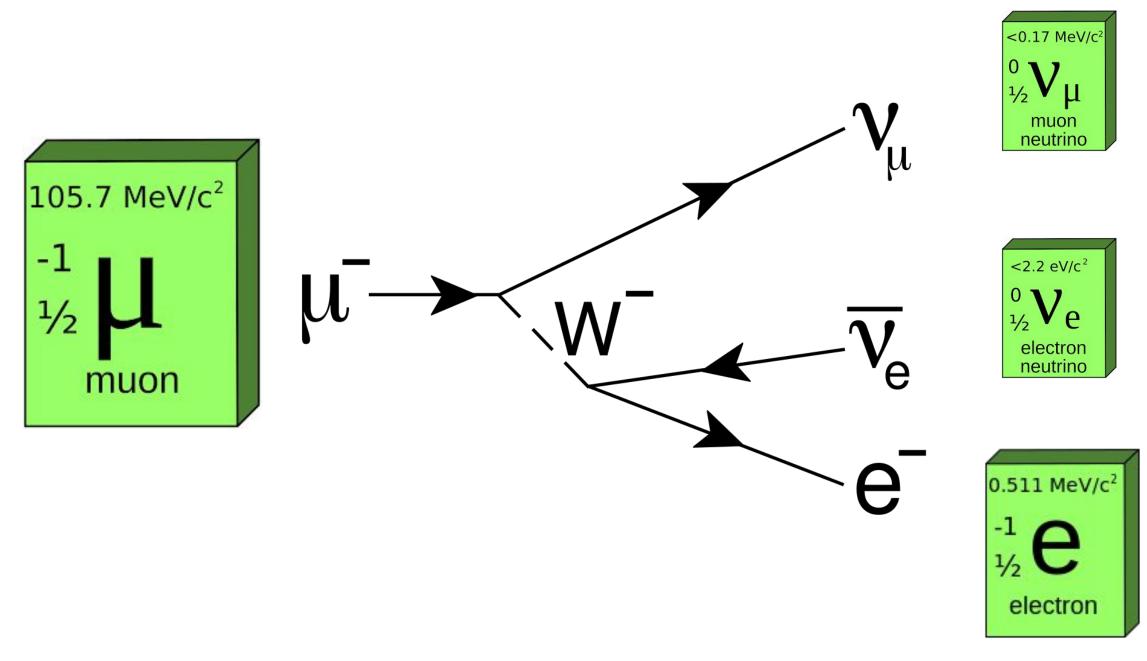
Possible intermediate steps towards a Muon collider

40/S2-D01 - Salle Dirac, CERN

Possible intermediate steps towards a Muon collider

from neutrino neutrino/lepton, electron muon to muon muon collisions?

"Exotica" Particle >> Construction "Exotica" Collider





Qiang Li, Peking University 2022/10/12

arXiv:2205.15350, arXiv:2204.11871, arXiv:2201.07808, arXiv:2109.01265, arXiv:2107.13581, arXiv:2010.15144







Possible intermediate steps towards a Muon collider

40/S2-D01 - Salle Dirac, CERN

Muon Collider: intermediate steps?

link

Matt Strassler | June 10, 2022 at 6:22 PM | Reply

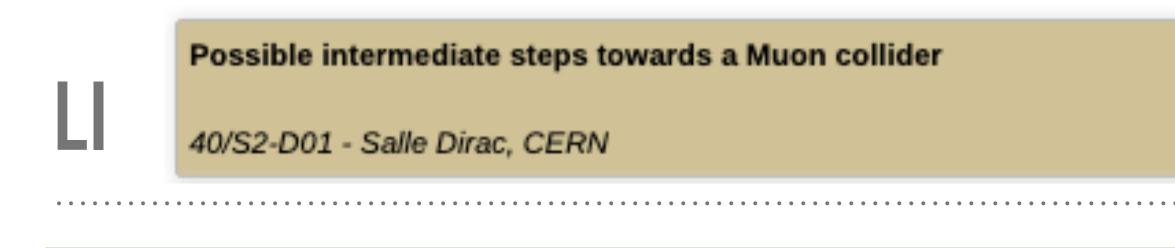
Andrew, these are very serious concerns too. But one cannot move before one has funding, and one cannot get funding without a clear argument as to why funding should be provided. At higher energy, the only clear arguments, right now, are for a Higgs/top factory. That will be an electron positron machine of some type, unless the ambitious muon collider project can demonstrate enough likelihood of success and enough intermediate physics goals (e.g. neutrino beams) that it can be justified as well. (Meanwhile other colliders at lower energy but very high luminosity might be pursued.)

19. Alain Blondel: alain.blondel@cern.ch question general but triggered by Steve Ritz. Establish the list of questions that are of great importance and should be answered across frontiers/experiments/facilities. Here is a question that I think if of key importance and is addressed in many 'frontiers' without being sufficiently asked as a unique question for which the various groups would gain to reflect in common: - given that neutrinos have masses, the question of existence and masses of right handed neutrinos (or their alternatives) should have a common discussion, formalism, expectations, visible consequences and what other problems they might solve, while understanding the possibilities, from the minimal one to those more complicated. This is certainly the most likely new physics there is, and it seems to naturally result from the present discoveries. It was evident from the presentations today that this question appears in the neutrino frontier, rare processes, cosmic and energy frontier as well as instrumentation, and in Hitoshi's presentation, and yet there is not a uniform language or momentum to look for it in all possible ways - so it remains somewhat confidential.

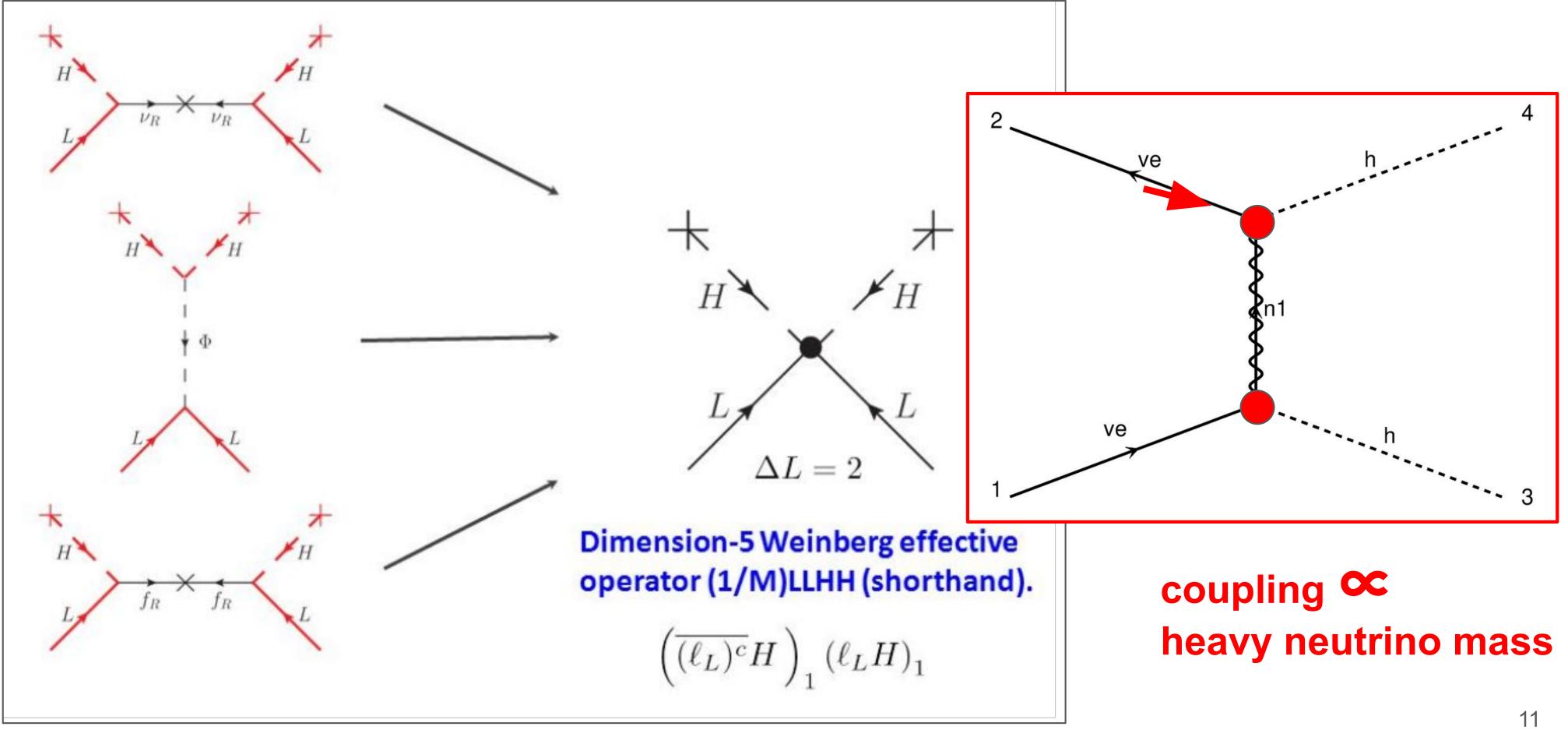
...enough intermediate physics goals (e.g. neutrino beams)" neutrino mass ... "This is certainly the most likely new physics there is..."



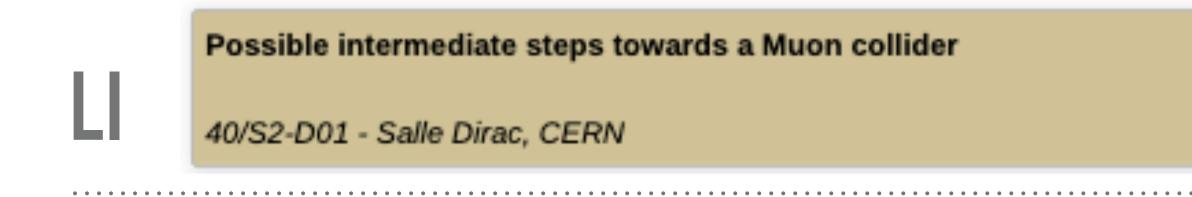
Seattle Snowmass Summer Meeting 2022



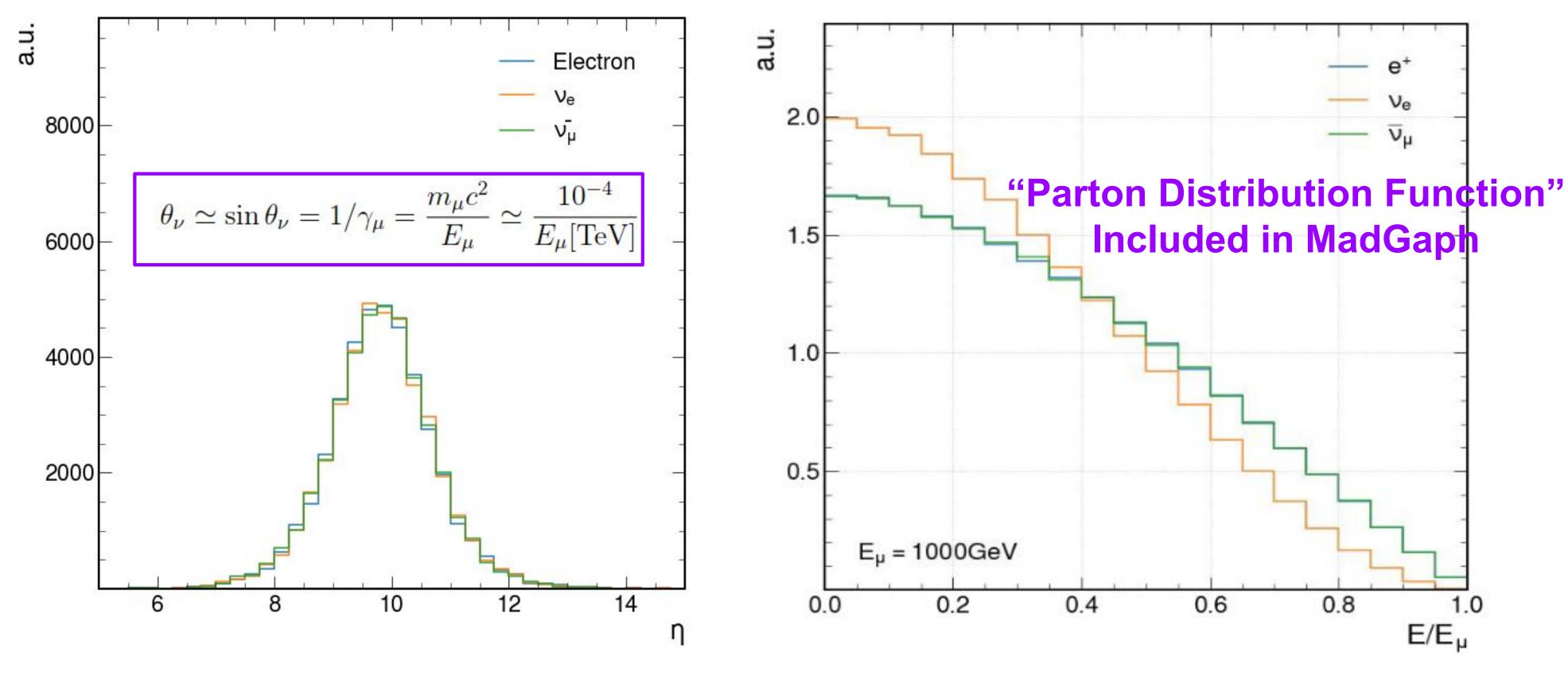
Neutrino Portal to BSM





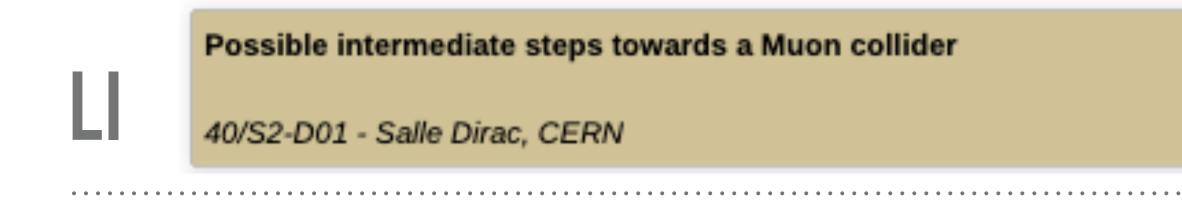


Neutrino Beam from 1TeV Muon beam



Highly collimated in angle, yet widely distributed in Energy





Neutrino Beam from 1TeV Muon beam

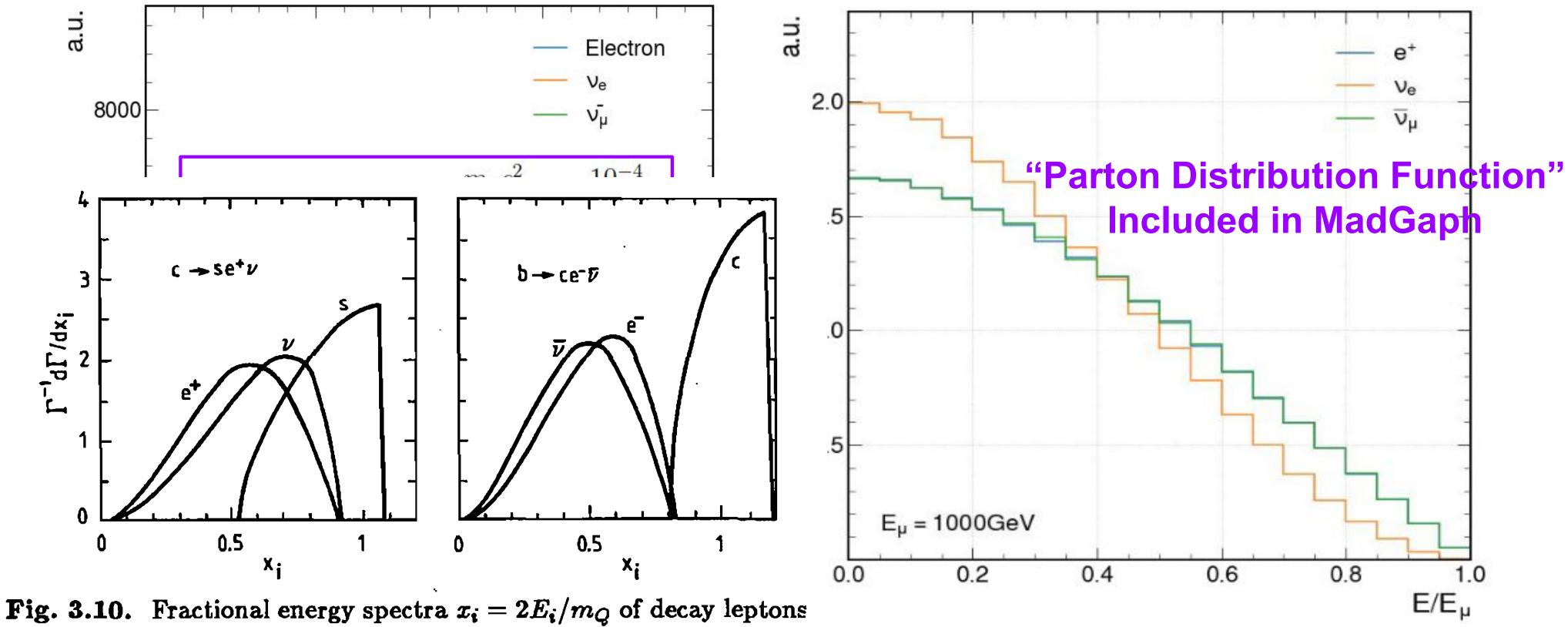


Fig. 3.10. Fractional energy spectra $x_i = 2E_i/m_Q$ of decay leptons and quarks from Q = c, b semileptonic decay.



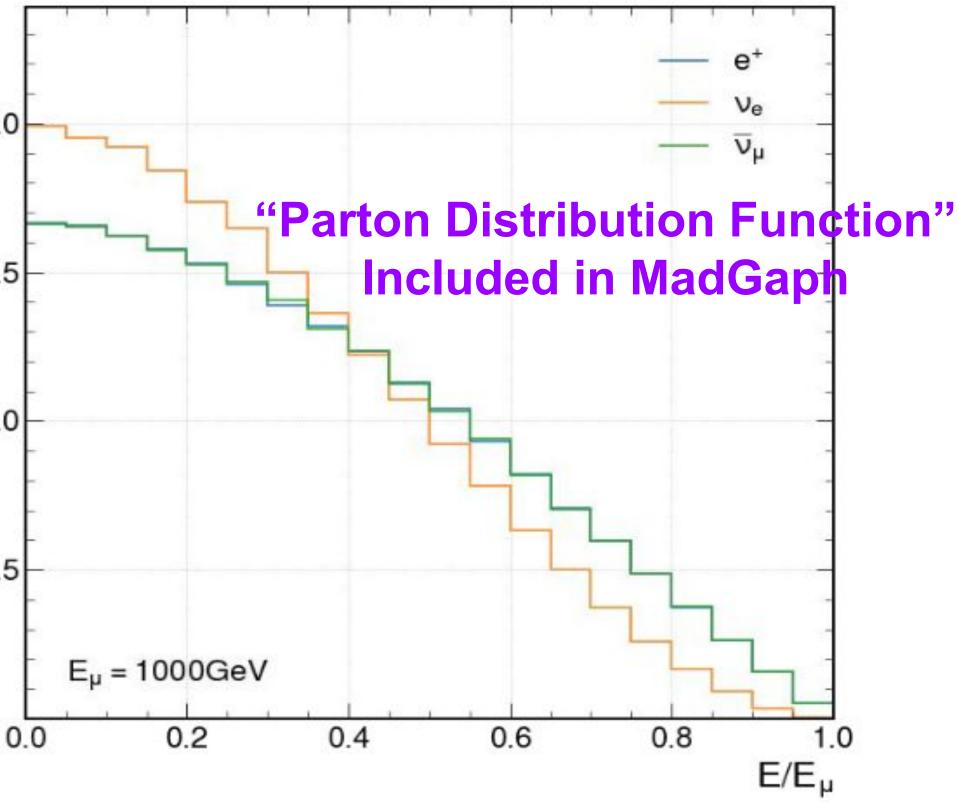
widely distributed in Energy

.

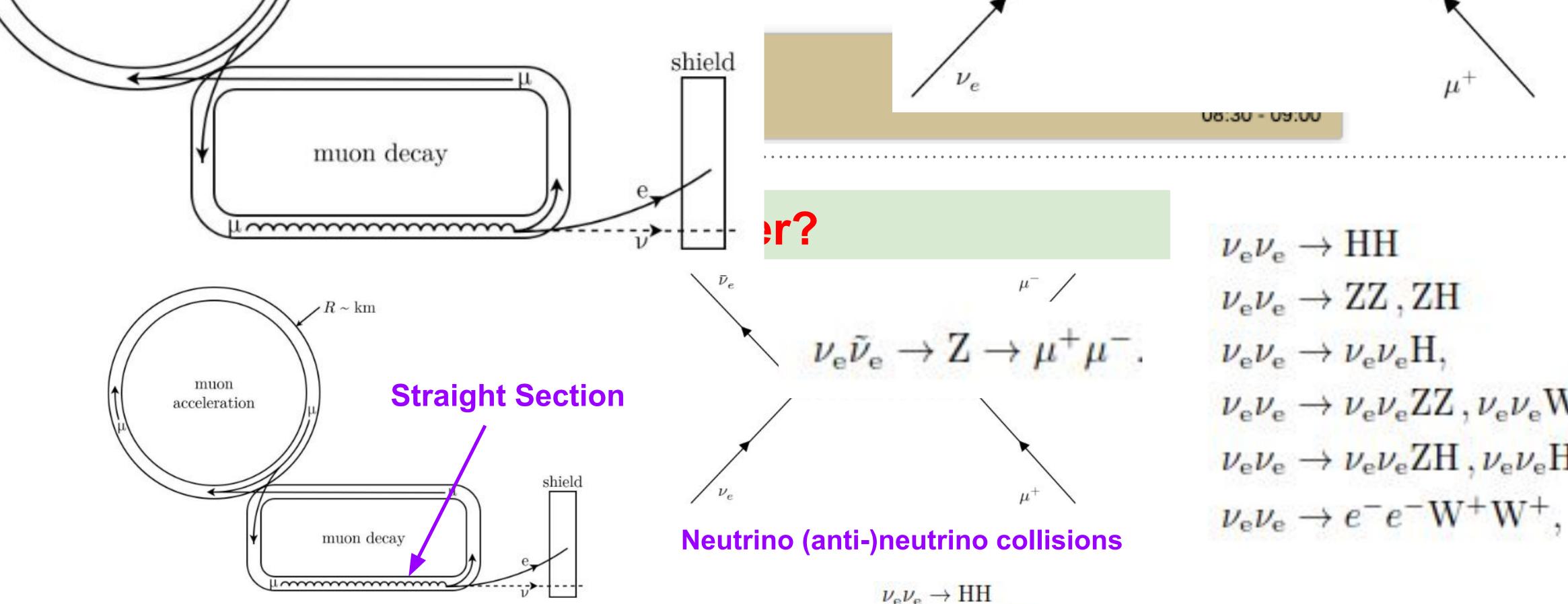
Possible intermediate steps towards a Muon collider 40/S2-D01 - Salle Dirac, CERN $\frac{1}{\Gamma}\frac{d\Gamma}{dx}$ Fig. 3.2. Positron decay spectrum from muon decay; DM 1TeV Muon beam 1.6 1.2 data from Phys. Rev. 119, 1400 (1960). The theoreti-0.8 a.u. cal curve includes radiative corrections and experimen-0.4 tal resolution; the latter ex-2.0 plains the tail above x = 1. 0.8 1.0 0.2 0.4 0.6 1.2 $c \rightarrow se^+ \nu$ b→ce⁻₽ Γ⁻¹dΓ/dx_i E_µ = 1000GeV 0.5 0.5 0 Xi

Fig. 3.10. Fractional energy spectra $x_i = 2E_i/m_Q$ of decay leptons and quarks from Q = c, b semileptonic decay.





widely distributed in Energy



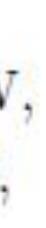
A small modulation of the muon decay angle through vertical bending, symbolized by the squiggly line, may be used to focus the neutrino beam.

 $\nu_{\rm e}\tilde{\nu}_{\rm e} \to {\rm Z} \to \mu^+\mu$

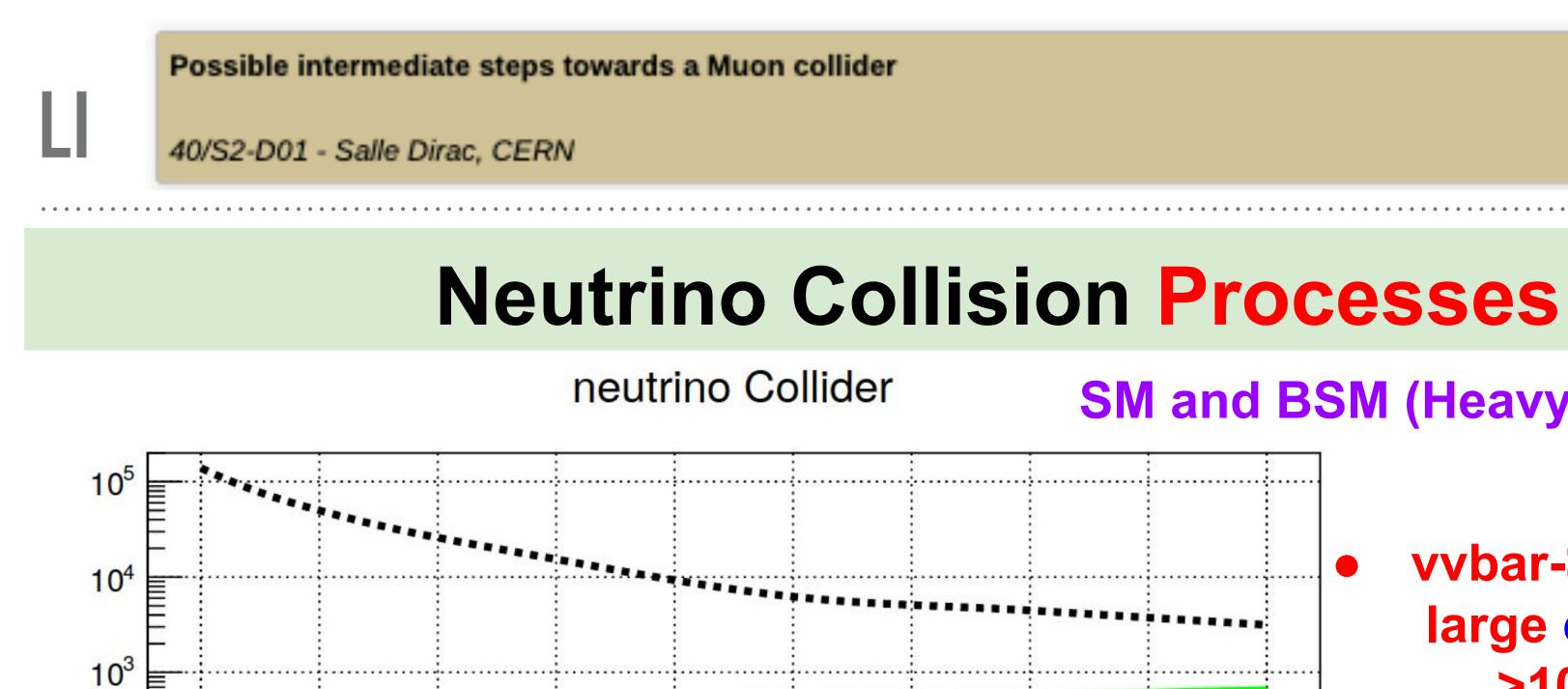
Question: ?/fb in 1-10 years

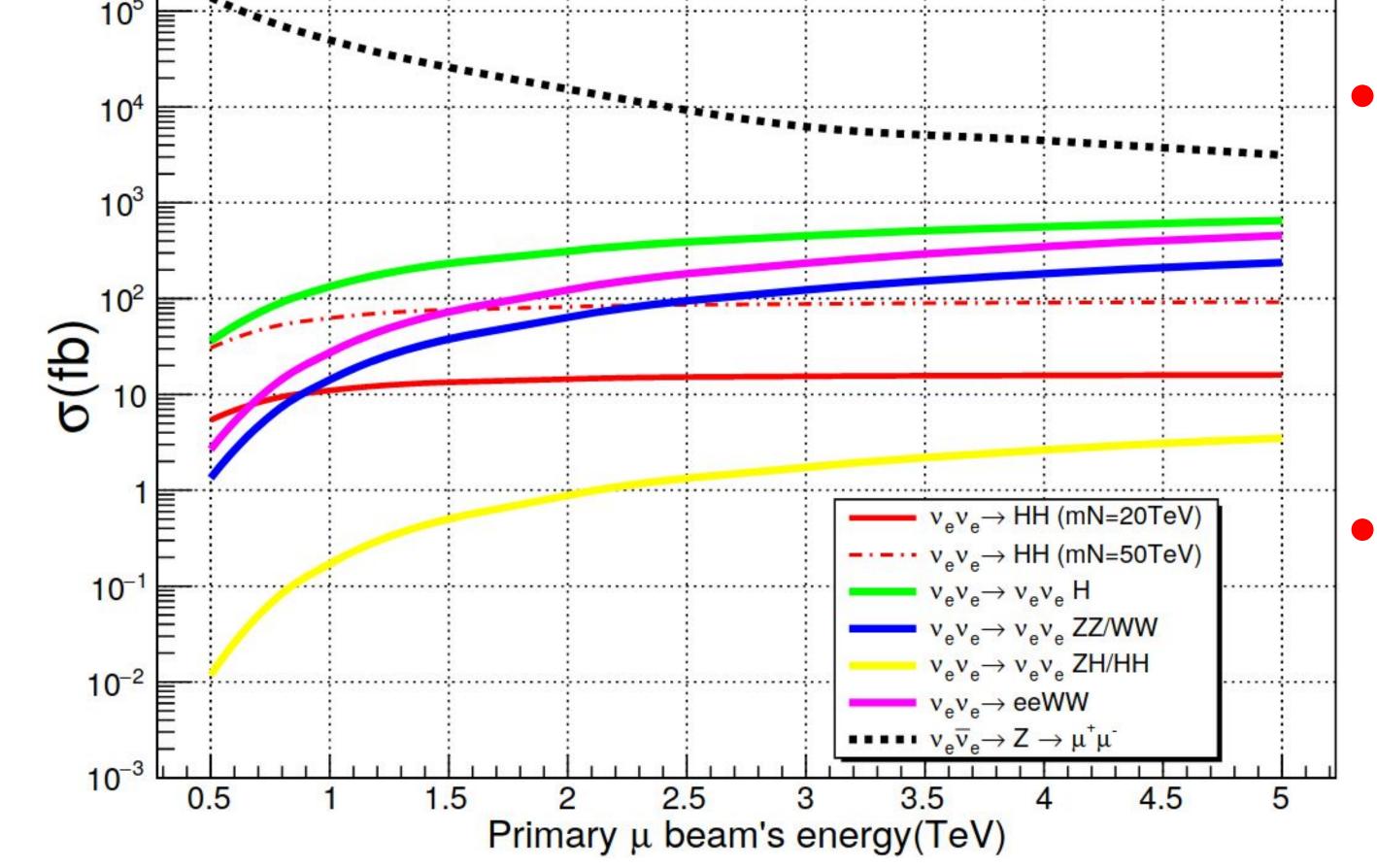
$$\begin{split}
\nu_{e}\nu_{e} \to HH \\
\nu_{e}\nu_{e} \to ZZ, ZH \\
\mu^{-}. \quad \nu_{e}\nu_{e} \to \nu_{e}\nu_{e}H, \\
\nu_{e}\nu_{e} \to \nu_{e}\nu_{e}ZZ, \nu_{e}\nu_{e}WW, \\
\nu_{e}\nu_{e} \to \nu_{e}\nu_{e}ZH, \nu_{e}\nu_{e}HH, \\
\nu_{e}\nu_{e} \to e^{-}e^{-}W^{+}W^{+},
\end{split}$$
14

 $\nu_{\rm e}\nu_{\rm e} \rightarrow \nu_{\rm e}\nu_{\rm e}ZZ, \nu_{\rm e}\nu_{\rm e}WW,$ $\nu_{\rm e}\nu_{\rm e} \rightarrow \nu_{\rm e}\nu_{\rm e} ZH$, $\nu_{\rm e}\nu_{\rm e} HH$,



14







SM and BSM (Heavy Majorana)

vvbar->Z: large cross section >100pb can be observed in short time! ~days to weeks

May loosen requirement on beam quality!



NUMBER BREAKING

L – violation

(1,1,0) (at least 2)

(1,1,0) (at least 2+1)

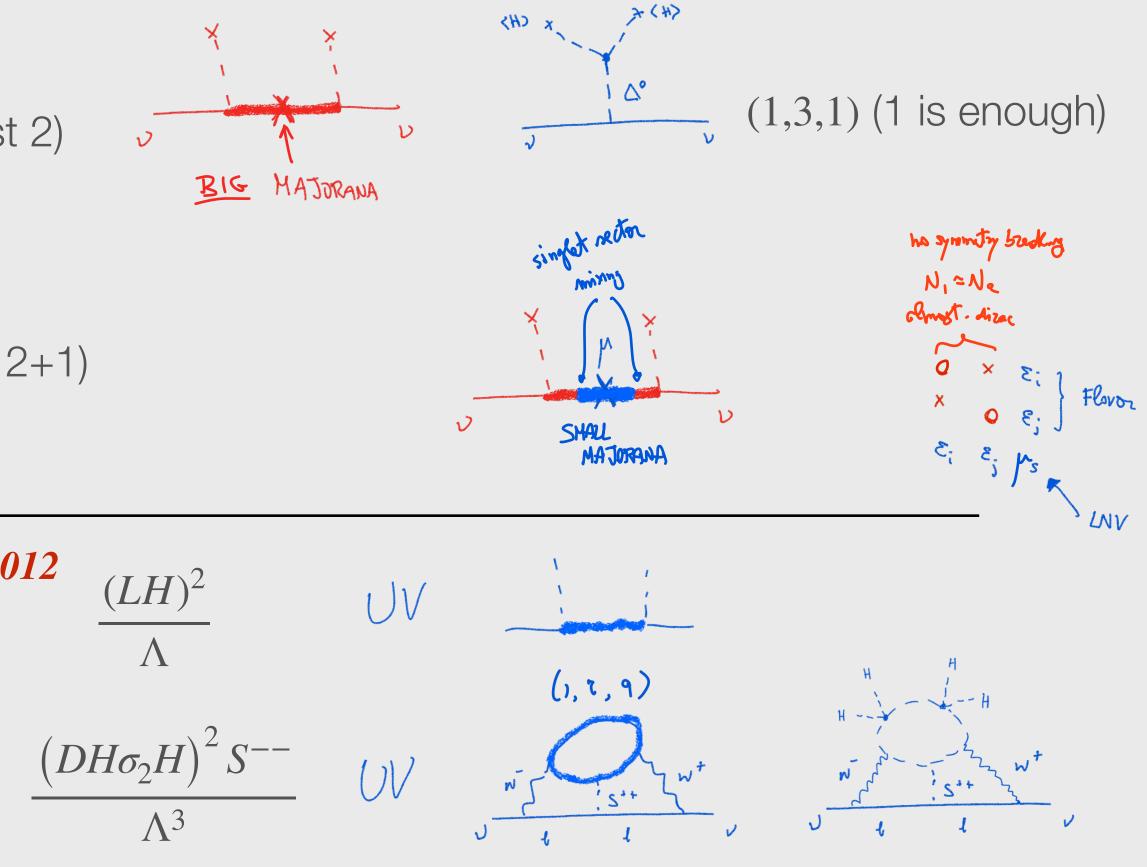
L – not accidental new physics before 2012 d = 5 (1,2,1/2)

d = 7 (1,1,2)

L – gauged, SSB

Roberto Franceschini - Moriond QCD 2022 - https://moriond.in2p3.fr/2022/QCD/

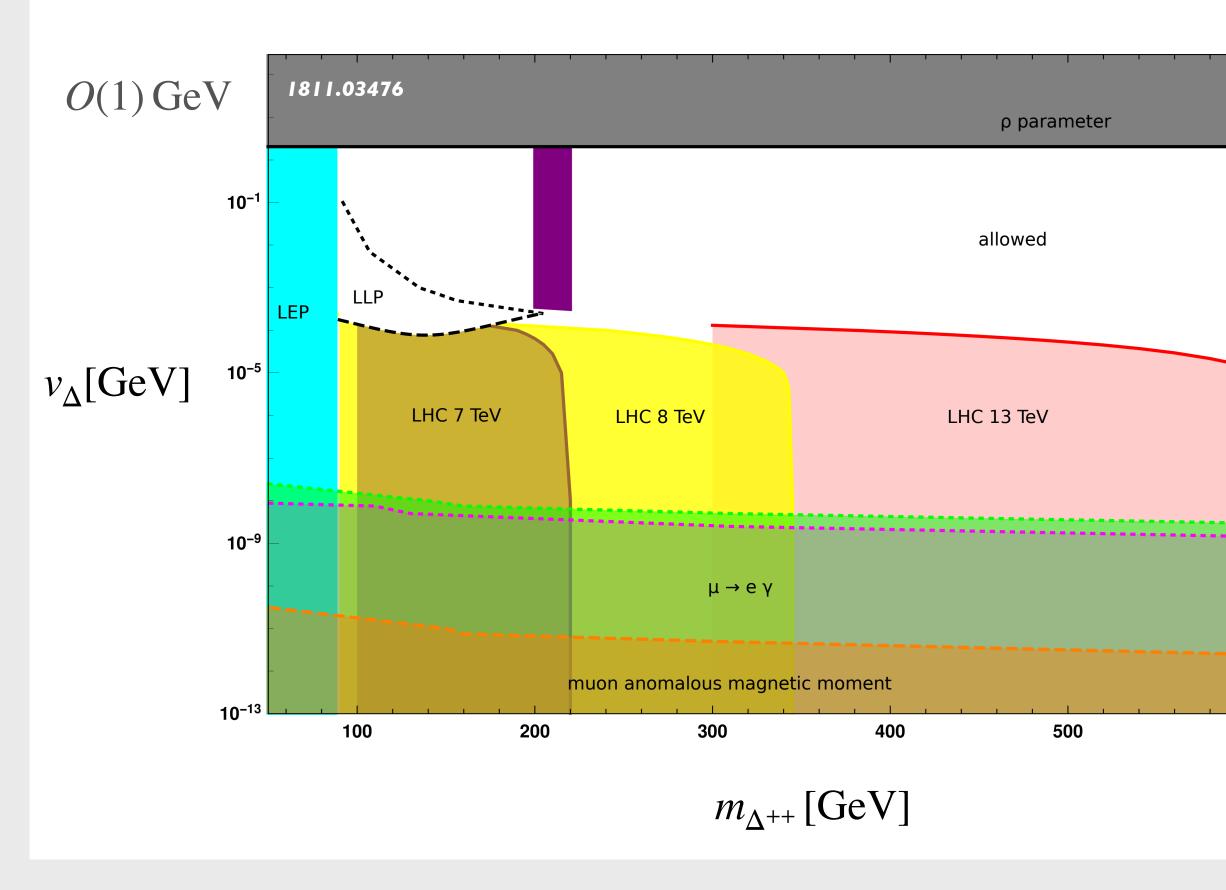
Neutrino mass mechanisms



 $SU(3) \otimes SU(2)_L \otimes SU(2)_L \otimes U(1)_{B-L}$ (1,2,1,1), (1,1,2,1), (1,2,2,1), (1,1,1,2),

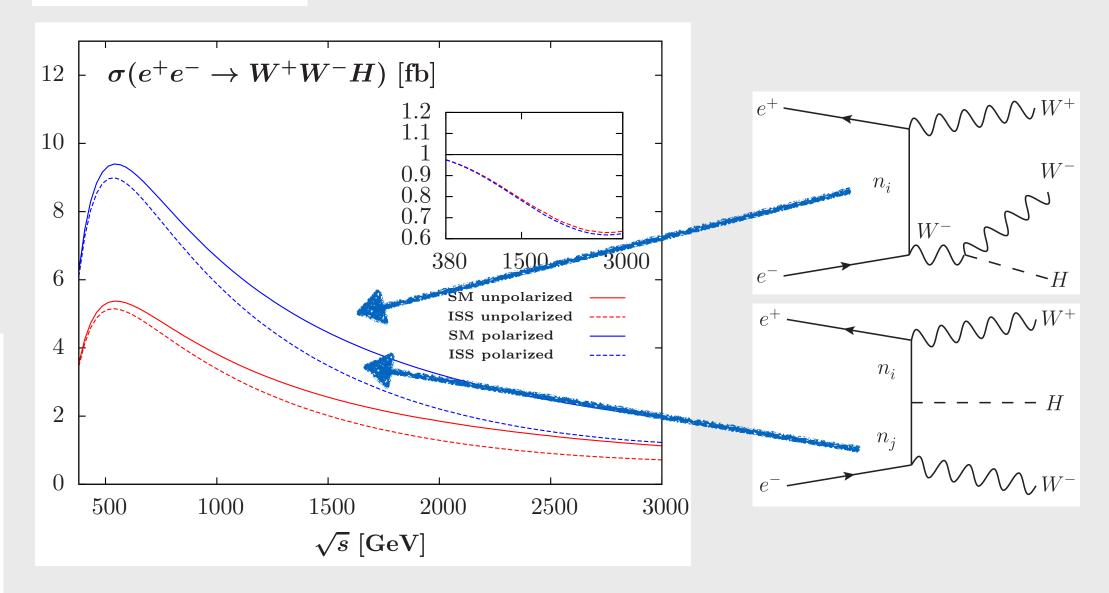
Plenty of neutrino mass models in reach

Type-2 See-Saw 1803.00677 - Agrawal, Mitra, Niyogi, Shil, Spannowsky



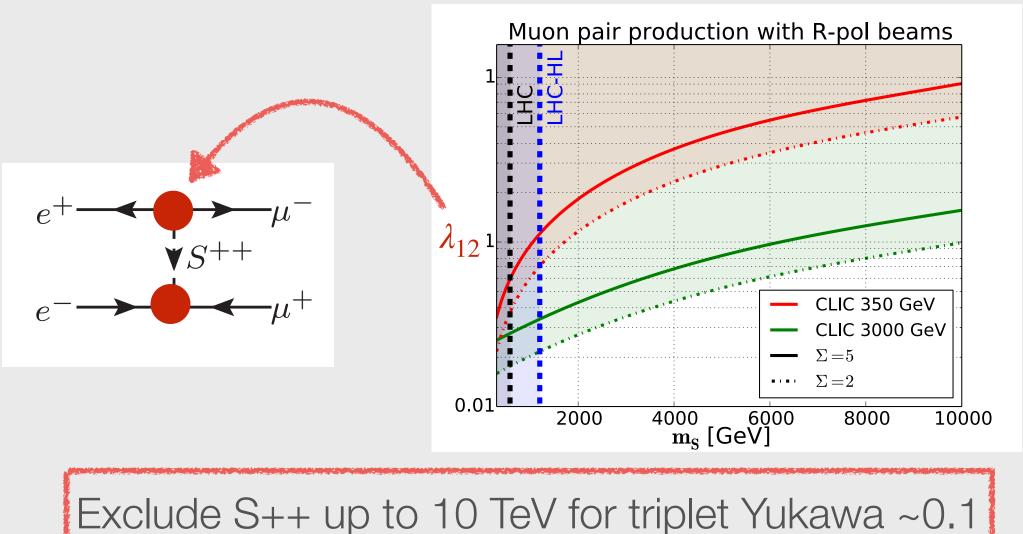
Roberto Franceschini - Moriond QCD 2022 - https://moriond.in2p3.fr/2022/QCD/

Inverse See-Saw 1712.07621 - Baglio, Pascoli, Weiland



Exclude ISS RH Neutrino up to 10 TeV for Yukawa ~1

1807.10224 - Crivellin, Ghezzi, Panizzi, Pruna, Signer

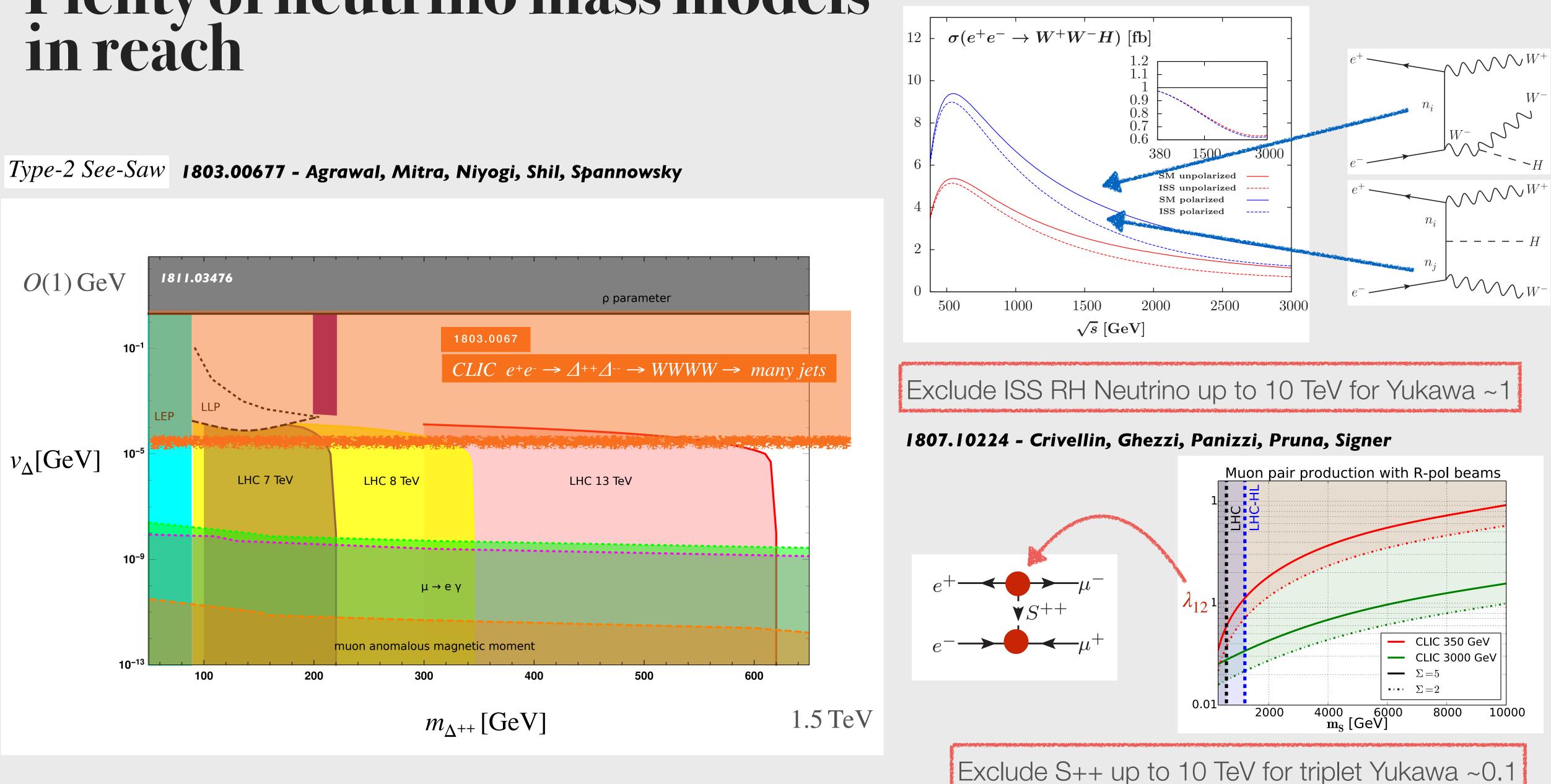


600

1.5 TeV



Plenty of neutrino mass models

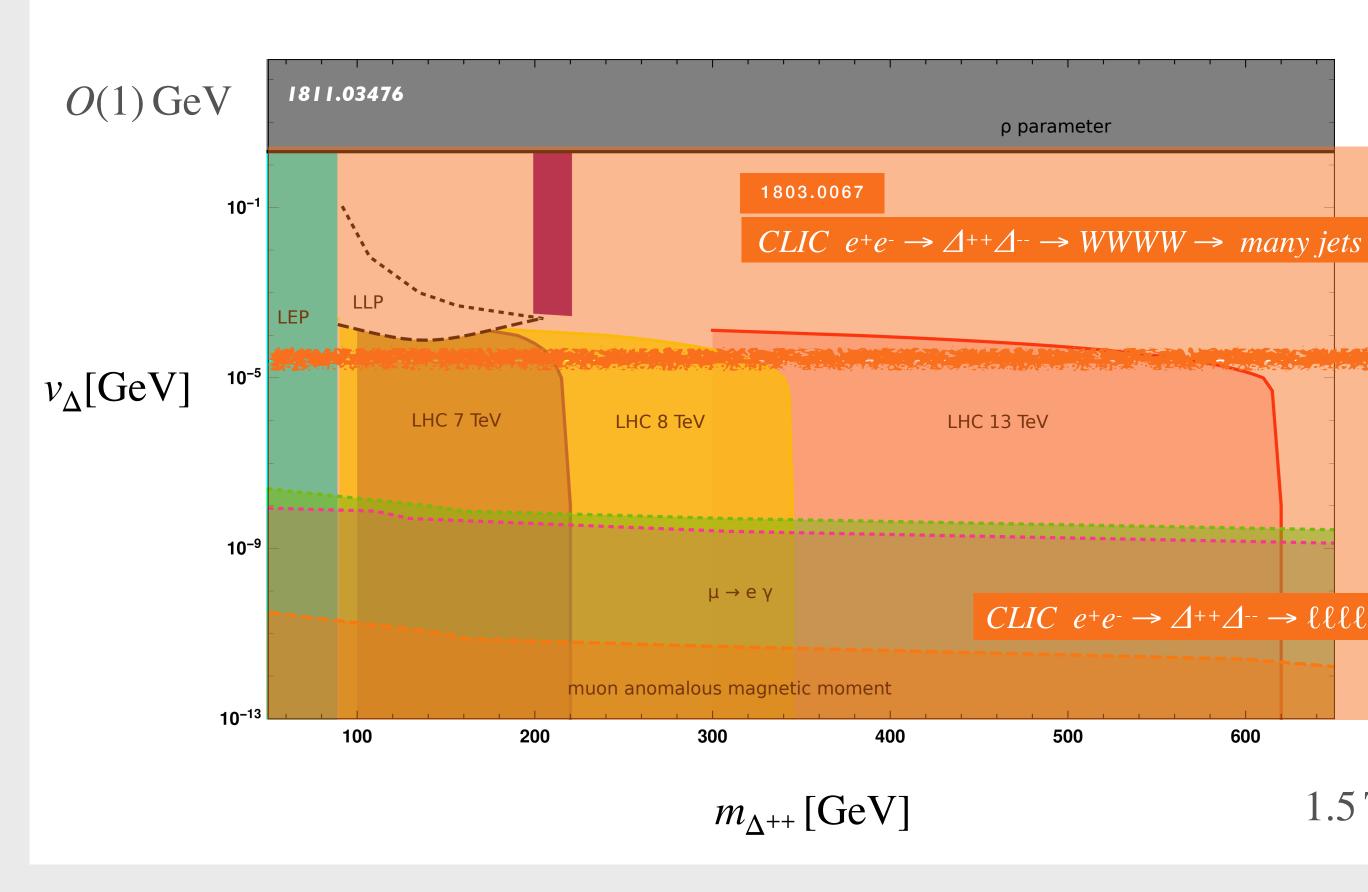


Roberto Franceschini - Moriond QCD 2022 - https://moriond.in2p3.fr/2022/QCD/

Inverse See-Saw 1712.07621 - Baglio, Pascoli, Weiland

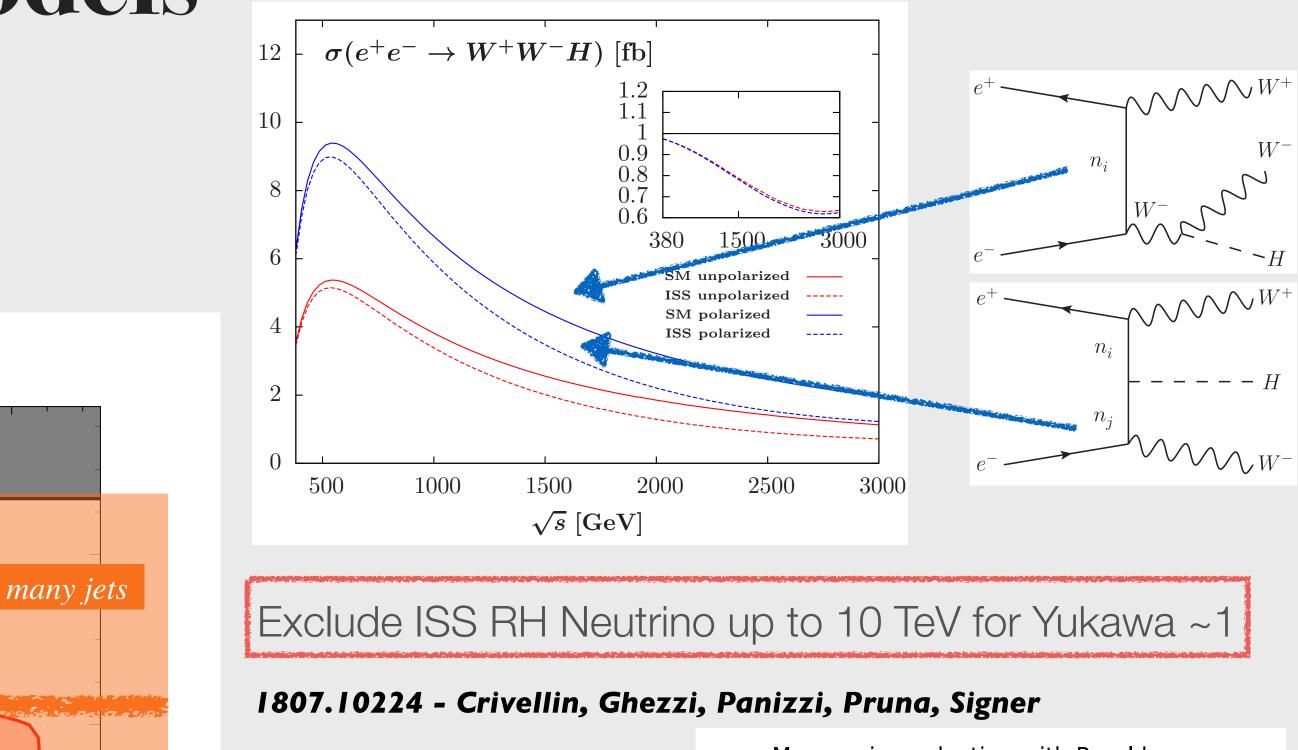
Plenty of neutrino mass models in reach

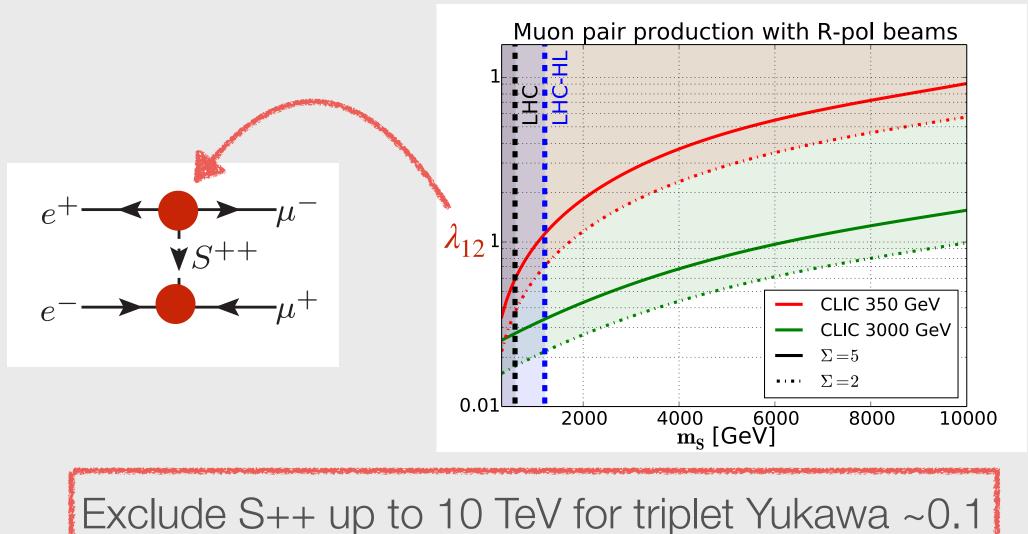
Type-2 See-Saw 1803.00677 - Agrawal, Mitra, Niyogi, Shil, Spannowsky



Roberto Franceschini - Moriond QCD 2022 - https://moriond.in2p3.fr/2022/QCD/

Inverse See-Saw 1712.07621 - Baglio, Pascoli, Weiland

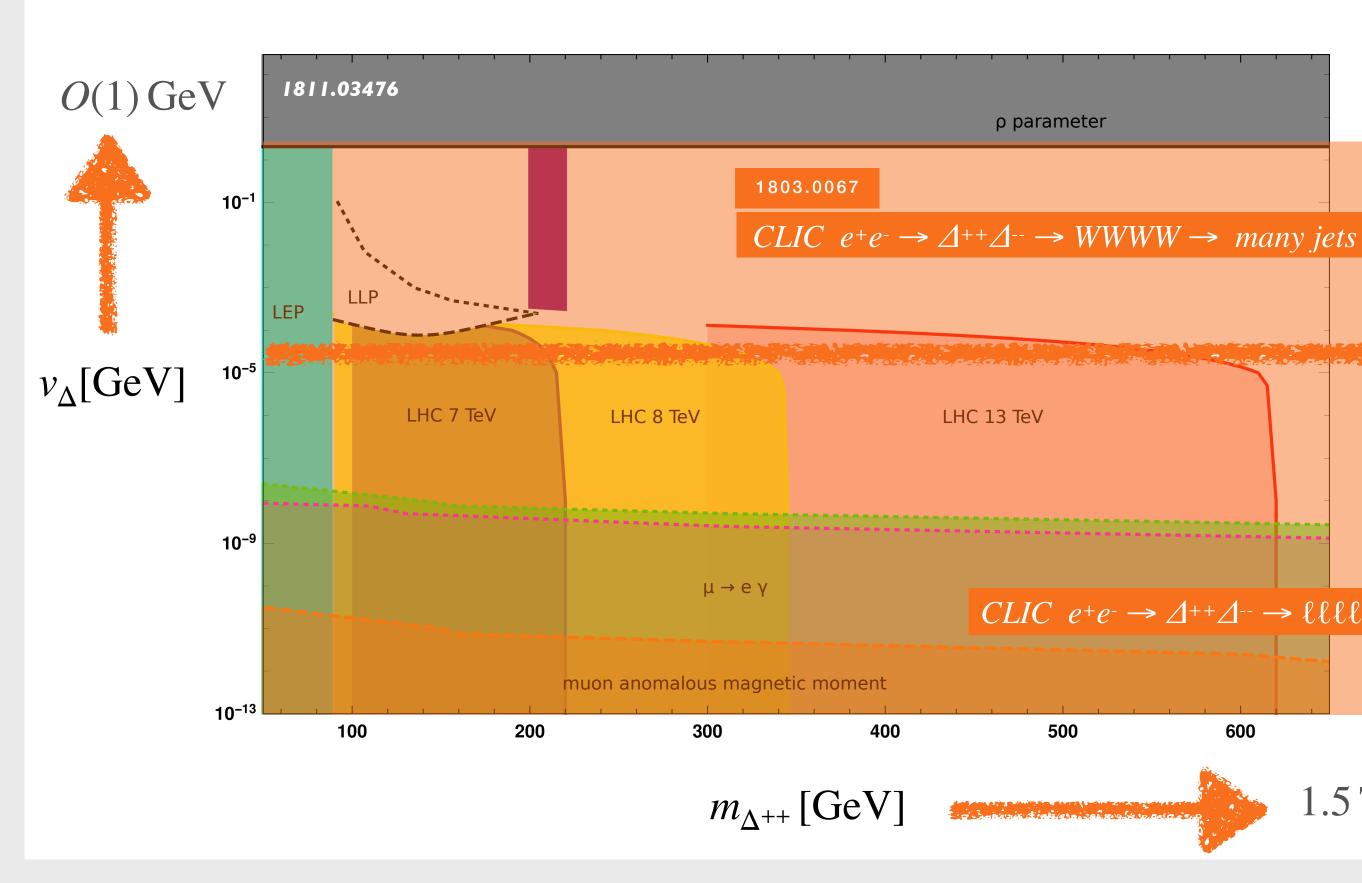




1.5 TeV

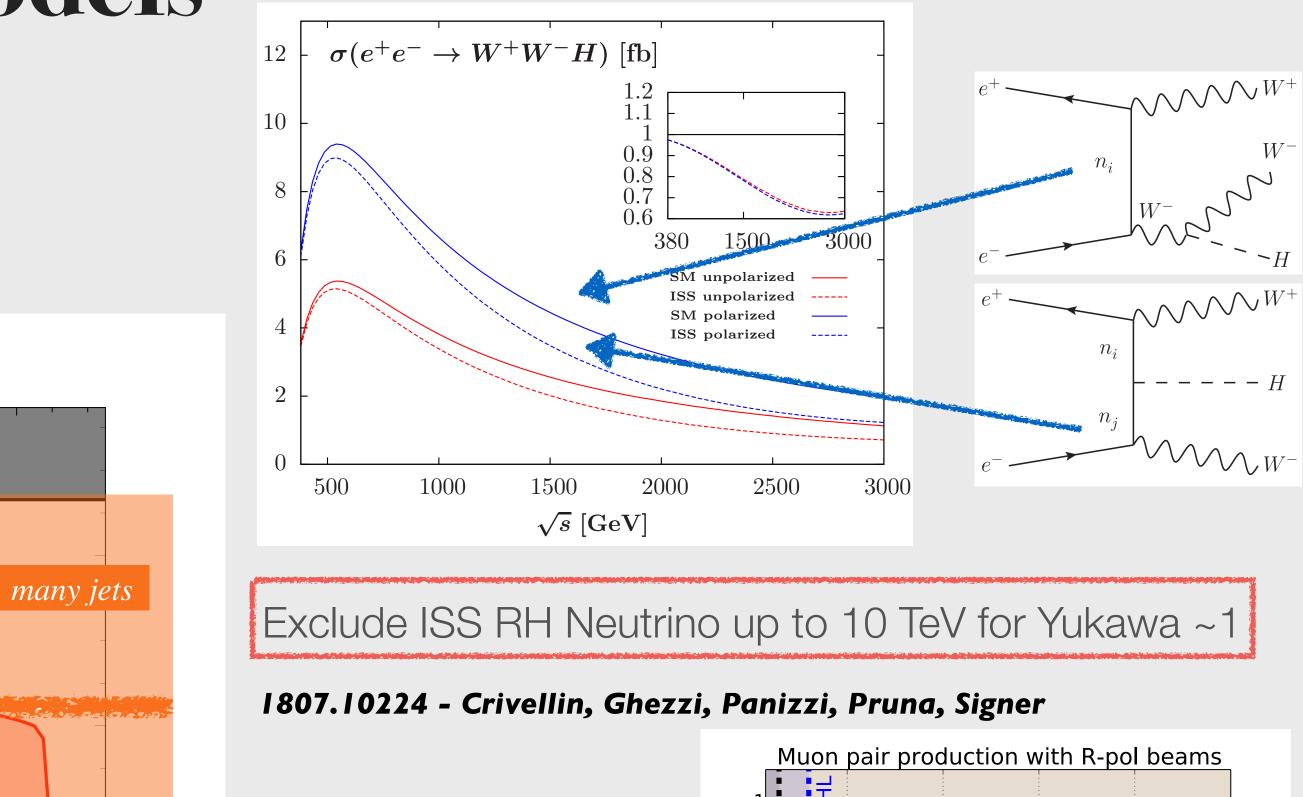
Plenty of neutrino mass models in reach

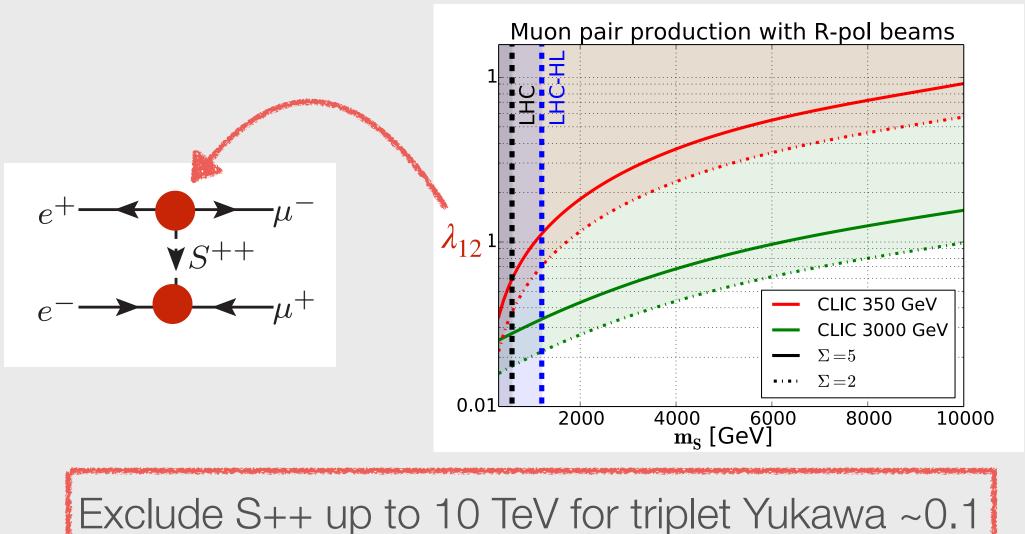
Type-2 See-Saw 1803.00677 - Agrawal, Mitra, Niyogi, Shil, Spannowsky



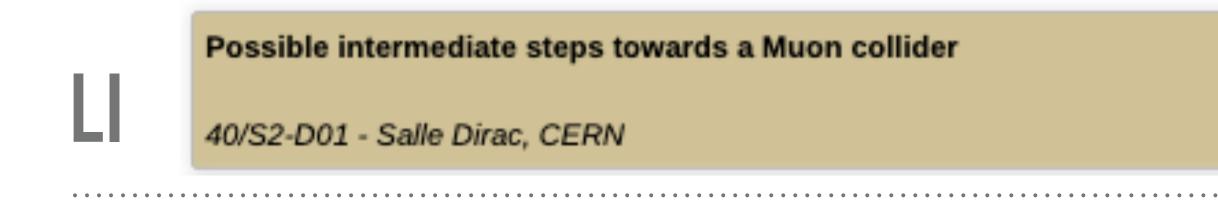
Roberto Franceschini - Moriond QCD 2022 - https://moriond.in2p3.fr/2022/QCD/

Inverse See-Saw 1712.07621 - Baglio, Pascoli, Weiland

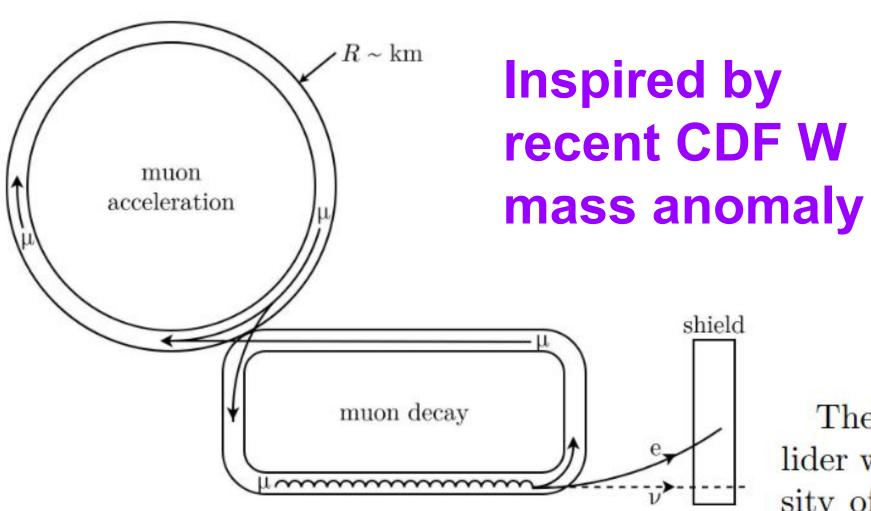




1.5 TeV



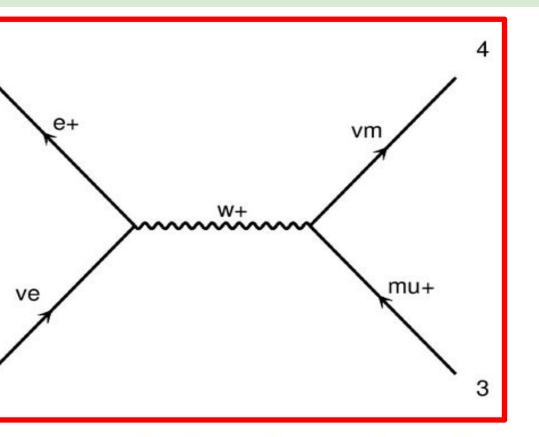
Neutrino lepton Collider



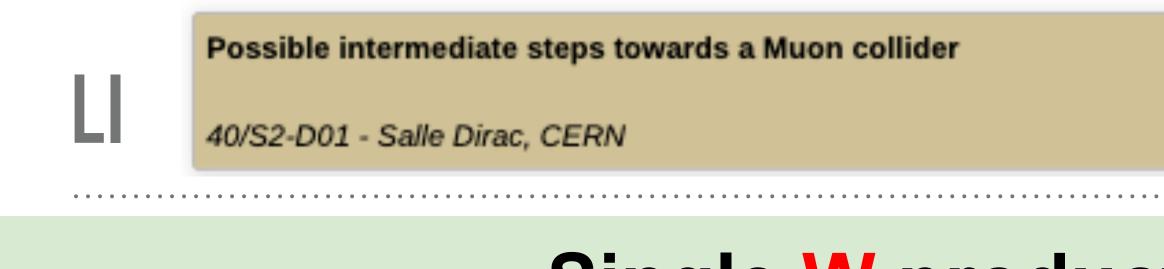
Similar design, but with only one sided neutrino beam, 0.1-1/fb in 10 years?

more focused neutrino beam 24.

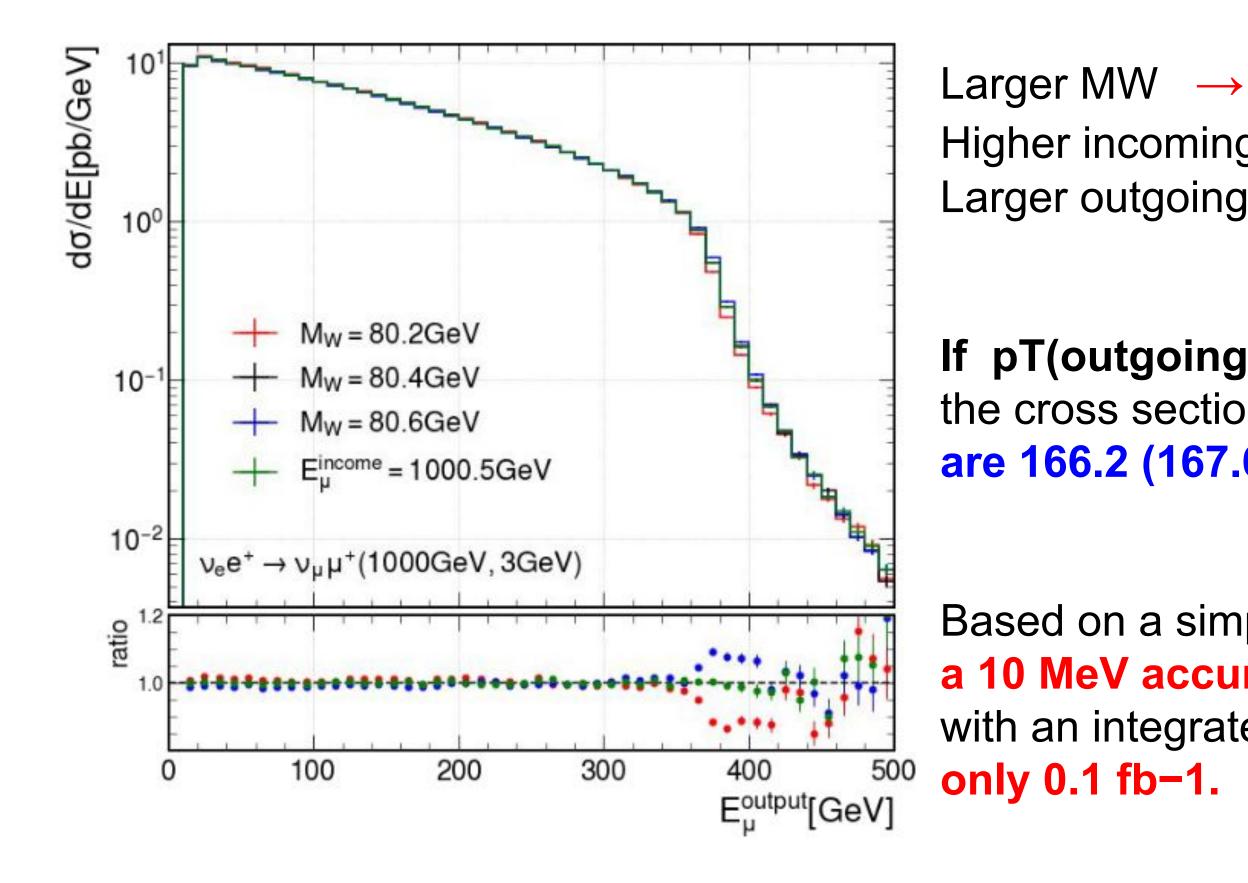


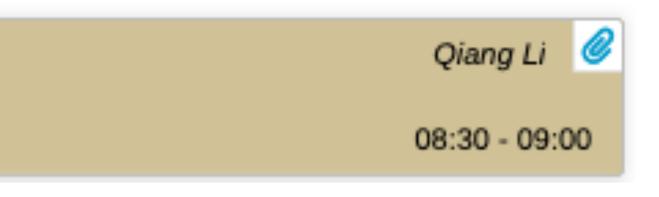


The instantaneous luminosity of a neutrino lepton collider would be limited by two main factors: 1) the intensity of the neutrino beam compared with the incoming muon beam is suppressed by roughly $L_l/L_c \sim 0.1$, i.e., the fraction of the collider ring circumference occupied by the production straight section [22], 2) the neutrino beam spread, which may still be kept at 10 to 100 microns at the interaction point, by applying a small modulation on muon decay angle through vertical bending to achieve 20



Single W production





Higher incoming neutrino Energy \rightarrow Larger outgoing Muon Energy (More boosted)

If pT(outgoing muon) > 40 GeV the cross sections with MW = 80.4 (80.41) are 166.2 (167.6) pb.

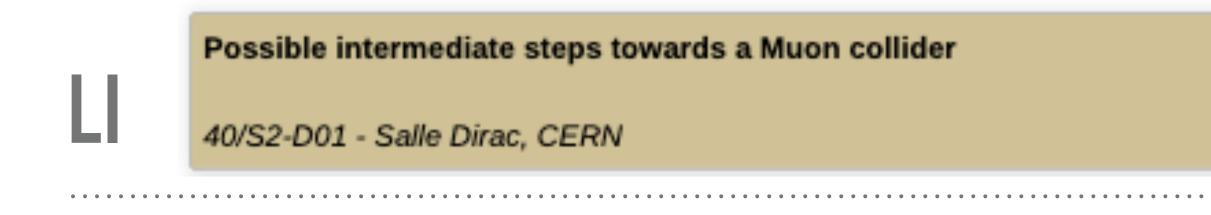
Based on a simple counting experiment, a 10 MeV accuracy on MW can be achieved with an integrated luminosity of

fiducial rate measurement $\delta\sigma/\sigma = 70 \cdot 10^{-4}$ yields $\delta m_W / m_W = 1.2 \cdot 10^{-4}$

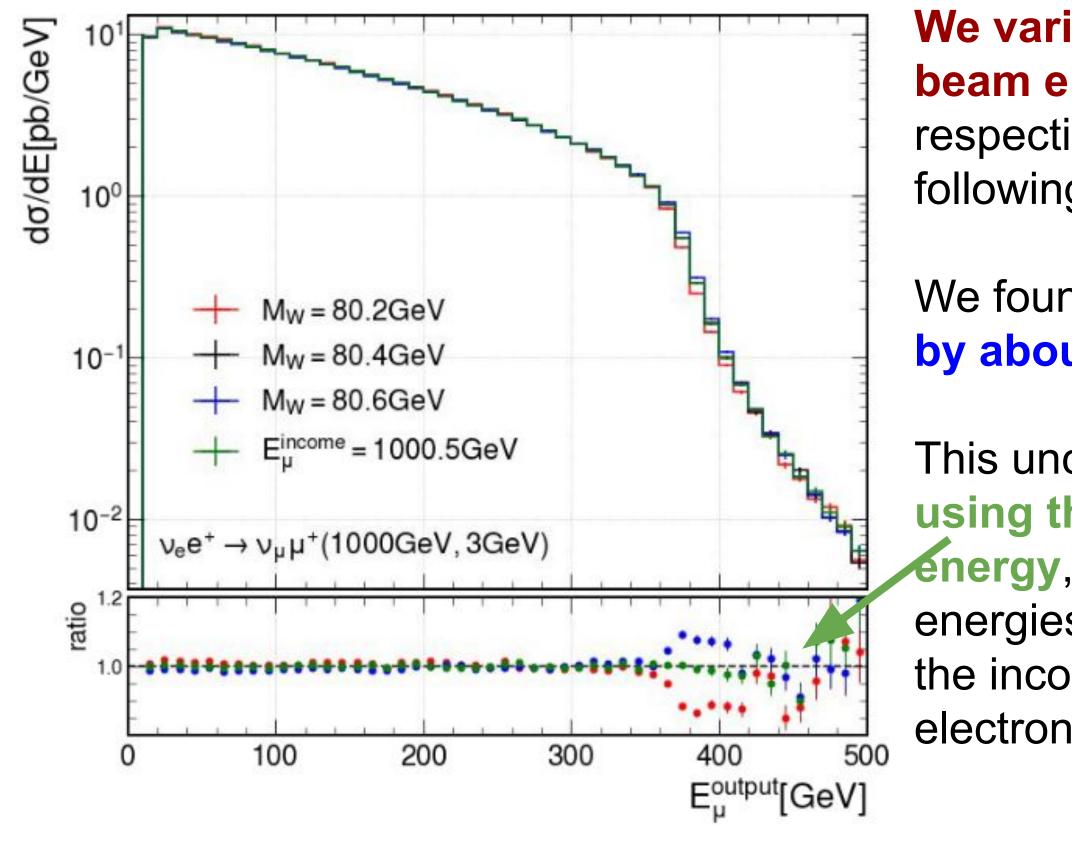
evidently the fiducial region of phase-space has a sensitivity to m_W







Robustness on W mass precision





We varied the incoming muon and electron beam energy by 0.5 GeV and 10 MeV,

respectively, which are quite conservative following previous refs.

We found that the cross sections changed by about 0.6 pb for both variations.

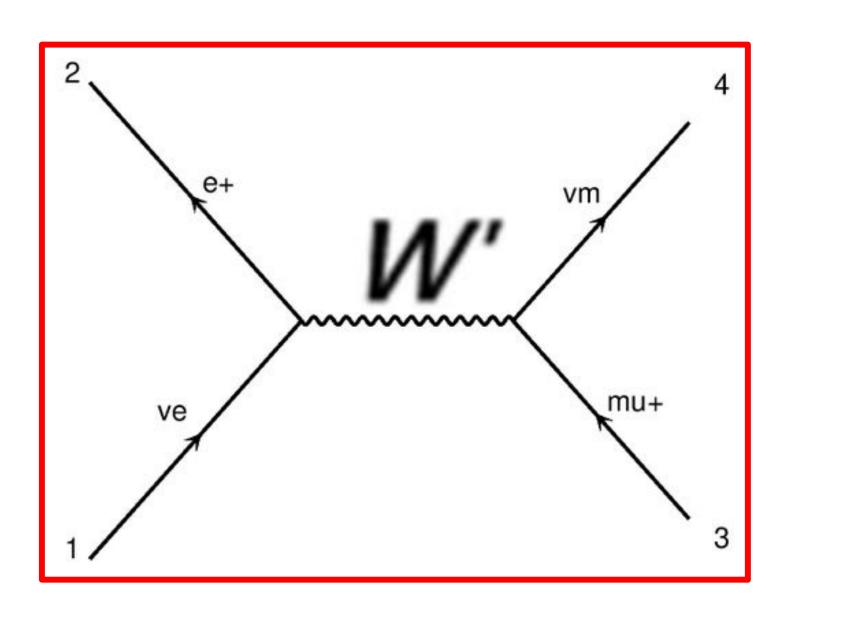
This uncertainty could be **mitigated by using the shape of the outgoing muon energy**, by scanning different incoming beam energies, or by calibrating the incoming muon beam energy with the electron decay products.

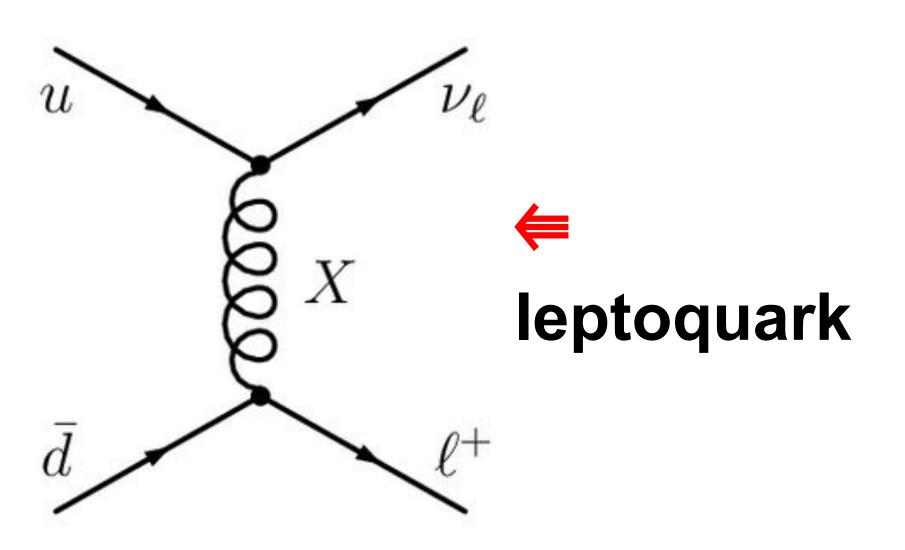
Possible intermediate steps towards a Muon collider

40/S2-D01 - Salle Dirac, CERN

More Physics from neutrino-lepton collisions

$$\begin{split} \mathrm{e}^{+}\mathrm{e}^{-} &\to \mathrm{Z}^{0(*)}, \ \nu_{\mathrm{e}}\mathrm{e}^{-} \to \nu_{\mathrm{e}}\mathrm{e}^{-}, \ \tilde{\nu}_{\mu}\\ \nu_{\mathrm{e}}\mathrm{e}^{+} &\to \mathrm{W}^{+(*)}, \ \tilde{\nu}_{\mu}\mathrm{e}^{+} \to \tilde{\nu}_{\mu}\mathrm{e}^{+},\\ \tilde{\nu}_{\mu}\mu^{-} &\to \mathrm{W}^{-(*)}, \ \nu_{\mathrm{e}}\mu^{-} \to \nu_{\mathrm{e}}\mu^{-}, \end{split}$$

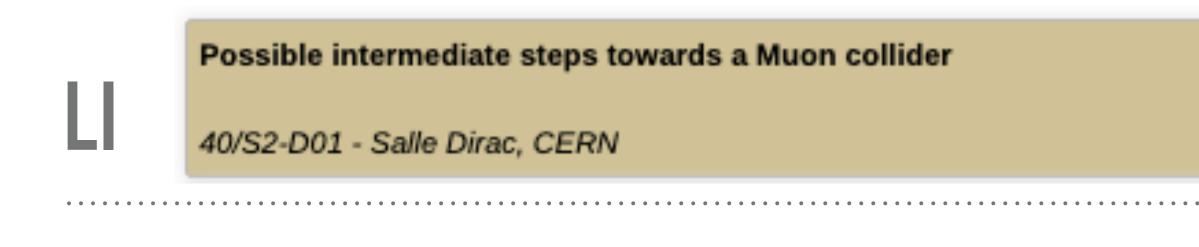






24

- $\tilde{\nu}_{\mu} e^{-} \rightarrow \tilde{\nu}_{\mu} e^{-},$ $\tilde{\nu}_{\mu} e^{+} \rightarrow \tilde{\nu}_{e} \mu^{+},$
- $\nu_{\rm e}\mu^- \to {\rm e}^-\nu_{\mu}.$
- **Anomalous Zvv couplings**



emu collider processes

