## Entanglement-Enabled Spin Inter erence

By James Daniel Brandenburg

1. Mystery: Nuclear Radius Measurements
2. Solved?: Imaging with Polared Tight 3. The (Inverse) Cotler-Wilcrek Process

2023 Workshop on Particle Correlations and Femtoseopy Catania, Italy


## Shiming light on Gluons

- Photo-nuclear measurements have been used to study QCD matter already for decades[1-3]
[1] H1 Collaboration. J. High Energ. Phys. 2010, 32 (2010). [2] ZEUS Collaboration. Eur. Phys. J. C 2, 247-267 (1998).


Well known process for probing the hadronic structure of the photon and nucleon (nuclear) target

UPCs: The Strongest Electromagnetic Fields

$\triangleright$ In heavy-ion collisions:
$E_{\text {max }}=\frac{Z e \gamma}{b^{2}} \approx 5 \times 10^{16}-10^{18} \mathrm{~V} / \mathrm{cm}$ $B_{\max } \sim 10^{14}-10^{16} \mathrm{~T}$
$\triangleright$ Strongest EM fields in the Universe
$\triangleright$ But very short lifetime - not constant
Must be treated in terms of photon quanta

$$
\begin{array}{ll}
E_{\gamma, \text { max }} \approx \gamma \hbar c / R & 80 \mathrm{GeV} @ \mathrm{LHC} \\
3 \mathrm{GeV} @ \mathrm{RHIC}
\end{array}
$$

High energy (small wavelength) photons
can be used to 'image' the nucleus

## Past Photo-Nuclear Measurements

- Many studies of $\gamma \mathbb{P} \rightarrow \rho^{0} \rightarrow \pi^{+} \pi^{-}$in the past



## Coherent Diffractive Interactions:

- Photon interacts with the entire nucleus
- Diffractive structure in $p_{T}^{2} \approx-t$
- Transverse momentum related to Fourier transform of nuclear density distribution

$$
\sigma(\gamma p \rightarrow V p)=\left.\frac{\mathrm{d} \sigma}{\mathrm{~d} t}\right|_{t=0} \int_{t_{\min }}^{\infty}|F(t)|^{2} \mathrm{~d} t
$$

## Past Photo-Nuclear $\mathbb{M}$ Measurements



Other measurements at RHIC \& LHC include:

Photoproduction of J $/ \psi$ in Au+Au UPC at $\sqrt{S_{N N}}=200 \mathrm{GeV}$
PHENIX Phys.Lett.B679:321-329,2009
$\rho^{0}$ vector mesons in $\mathrm{Pb}-\mathrm{Pb}$ UPC at $\sqrt{S_{N N}}=$ 5.02 TeV

ALICE, JHEPO6 (2020) 35
$\mathrm{J} / \Psi$ in $\mathrm{Pb}+\mathrm{Pb}$ UPC at $\sqrt{S_{N N}}=2.76 \mathrm{TeV}$ CMS, Phys. Lett. B 772 (2017) 489
... and many more

## So what's the problem?



Photo-nuclear measurements have historically produced a |t| slope that corresponds to a mysteriously large source!

STAR (2017): $|t|$ slope $=407.8 \pm 3(\mathrm{GeV} / \mathrm{c})^{-2}$
$\rightarrow$ Effective radius of 8 fm
$\left(R_{\text {Au }}^{\text {charged }} \approx 6.38 \mathrm{fm}\right)$
ALICE (Pb) : $|\mathrm{t}|$ slope $=426 \pm 6 \pm 15(\mathrm{GeV} / \mathrm{c})^{-2}$
$\rightarrow$ Effective radius of 8.1 fm
$\left(R_{P b}^{\text {charged }} \approx 6.62 \mathrm{fm}\right)$

## Extracted nuclear radii are way too large to be explainable

STAR Collaboration, L. Adamczyk, et al., Phys. Rev. C 96, 054904 (2017).
J. Adam et al. (ALICE Collaboration), J. High Energy Phys. 1509 (2015) 095.


Imaging the $\mathbb{N} u c l e u s$ with Polarized Photons

## What is NEW with transversely polarized photons?



What is NEW with transversely polarized photons?
C. Li, J. Zhou, Y. Zhou, Phys. Lett. B 795, 576 (2019)
C. Li, J. Zhou \& Y. Zhou Phys. Rev. D 101, 034015 (2020)


Gluons from nucleus


Recently realized that asymmetries in angle $\phi$ related to polarization

## Access to initial photon polarization

## Imaging the Nucleus with Polarized Photons

What is NEW with transversely polarized photons?


Gluons from nucleus
C. Li, J. Zhou, Y. Zhou, Phys. Lett. B 795, 576 (2019)
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- Intrinsic photon spin transferred to $\rho^{0}$
- $\rho^{0}$ spin converted into orbital angular momentum between pions
o Observable as anisotropy in $\pi^{ \pm}$ momentum


## Access to initial photon polarization

# Observation of Strong Asymmetry in $\rho^{0} \rightarrow \pi^{+} \pi^{-}$ 




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## Observation of Strong Asymmetry in $\rho^{0} \rightarrow \pi^{+} \pi^{-}$

STAR: Signal Trtr pairs with $P_{T}<60 \mathrm{MeV}$


- Intrinsic photon spin transferred to $\rho^{0}$
- $\rho^{0}$ spin converted into orbital angular momentum between pions
- Observable as anisotropy in $\pi^{ \pm}$

STAR Collaboration, Sci. Adv. 9, eabq3903 (2023). momentum

## Trivial Spin-Momentum Alignment?

## For a single diagram (pA)



Gluons from nucleus


VM inherits the spin from photon (no helicity flip)
Diffractive -> VM momentum dominantly from the Pomeron
$\rightarrow$ VM has no alignment between spin and momentum

Imaging the $\mathbb{N} u c l e u s$ with Polarized Photons

## What is NEW with transversely polarized photons?



Imaging the $\mathbb{N} u c l e u s$ with Polarized Photons

## What is NEW with transversely polarized photons?



What is NEW with transversely polarized photons?


Both possibilities occur simultaneously

## Interference of two amplitudes



2

## Interference of two amplitudes

interference - So What?


Interference of Amplitudes, so what?


## Robust Theoretical Description

- First theoretical prediction for deformed Uranium
- Sensitivity to nuclear geometry!


## $\boldsymbol{\beta}_{2}$



3. The CotlerWilczek
Process


- Two photon measurement from incoherent source
- "image" encoded in transverse correlations
- Requires photons be indistinguishable


## Intensity Interferometry

- Incoherent Source
- Interference results from second-order coherence
- Quantum statistics determines bunching vs. anti-bunching $\mathbf{g}^{(2)}(\mathbf{t})$ second-order correlation


Photon detections as function of time for a) antibunched, b) random, and c) bunched light

## Intensity Interferometry

- Results from higher order coherence

$$
\begin{aligned}
|\phi\rangle= & \left(A_{1 \alpha} A_{2 \beta}+A_{2 \alpha} A_{1 \beta}\right)|\omega, \omega\rangle \\
\langle\phi \mid \phi\rangle= & \left|A_{1 \alpha}\right|^{2}\left|A_{2 \beta}\right|^{2}+\left|A_{2 \alpha}\right|^{2}\left|A_{1 \beta}\right|^{2} \\
& +A_{1 \alpha} A_{2 \beta} A_{2 \alpha}^{*} A_{1 \beta}^{*}+A_{1 \alpha}^{*} A_{2 \beta}^{*} A_{2 \alpha} A_{1 \beta} \\
& \left\langle A_{1 \alpha} A_{1 \beta}^{*}\right\rangle_{E} \neq 0
\end{aligned}
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$$

Requires indistinguishable states!


## The Cotler-Wilczek Process

Sources
$|\psi\rangle=A_{1 \alpha} A_{2 \beta}\left|\omega_{1}, \omega_{2}\right\rangle+A_{2 \alpha} A_{1 \beta}\left|\omega_{2}, \omega_{1}\right\rangle$

$$
\langle\psi \mid \psi\rangle=\left|A_{1 \alpha} A_{2 \beta}\right|^{2}+\left|A_{2 \alpha} A_{1 \beta}\right|^{2}
$$

## Distinguishable states = NO Interference!


arXiv:1502.02477


## The Cotler-Wilczek Process

$|\psi\rangle=A_{1 \alpha} A_{2 \beta}\left|\omega_{1}, \omega_{2}\right\rangle+A_{2 \alpha} A_{1 \beta}\left|\omega_{2}, \omega_{1}\right\rangle$
Sources


Interference Recovered! $\left\langle A_{1 \alpha} A_{1 \beta}^{*}\right\rangle_{E} \neq 0$

1. Entangler performs unitary transformation:

$$
\begin{aligned}
U\left|\omega_{1}\right\rangle & =\cos (\theta)\left|\omega_{1}\right\rangle+\sin (\theta) e^{i \omega_{0}}\left|\omega_{2}\right\rangle \\
U\left|\omega_{2}\right\rangle & =\sin (\theta) e^{-i \omega_{0}}\left|\omega_{1}\right\rangle+\cos (\theta)\left|\omega_{2}\right\rangle
\end{aligned}
$$

2. Filter projects common state:

$$
\begin{aligned}
& \left|\omega_{1} \omega_{2}\right\rangle \rightarrow \cos (\theta) \sin (\theta) e^{-i \omega_{0}}\left|\omega_{1}, \omega_{1}\right\rangle \\
& \left|\omega_{2} \omega_{1}\right\rangle \rightarrow \cos (\theta) \sin (\theta) e^{-i \omega_{0}}\left|\omega_{1}, \omega_{1}\right\rangle
\end{aligned}
$$

## The Cotler-Willczelk Process

$$
|\psi\rangle=A_{1 \alpha} A_{2 \beta}\left|\omega_{1}, \omega_{2}\right\rangle+A_{2 \alpha} A_{1 \beta}\left|\omega_{2}, \omega_{1}\right\rangle \quad \text { Sources }_{2}
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2. Filler projects common state:

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Interference Recovered! $\left\langle A_{1 \alpha} A_{1 \beta}^{*}\right\rangle_{E} \neq 0$


Entanglement Enabled Intensity Interference


> Entanglement enabled Intensity Interferometry from exclusive $\pi^{+} \pi^{-}$measurements in UPC's as an inverse Cotler-Wilczek process
> Haowu Duan, Raju Venugopalan, Zhoudunming Tu, Zhangbu Xu, James Daniel Brandenburg, In preparation

Inverse Cotler-Wilczek Process: 'Filter' $\rho^{0}$ state comes first.
Entanglement of daughter pions enables interference

$$
\begin{aligned}
<N_{A} N_{B} \mid \pi^{+} \pi^{-}> & =<N_{A} N_{B}\left|\rho_{A}><\rho_{A}\right| \pi^{+} \pi^{-}, A>\square \text { Filler }^{\circ} \\
& \times<\pi^{+} \pi^{-}, A \mid\left(\left|\pi^{+}, 1>\left|\pi^{-}, 2>+\left|\pi^{+}, 2>\right| \pi^{-}, 1>\right)\right.\right. \\
& +<N_{A} N_{B}\left|\rho_{B}><\rho_{B}\right| \pi^{+} \pi^{-}, B>\square \text { Filter } \\
& \times<\pi^{+} \pi^{-}, B \mid\left(\left|\pi^{+}, 1>\left|\pi^{-}, 2>+\left|\pi^{+}, 2>\right| \pi^{-}, 1>\right)\right.\right.
\end{aligned}
$$

Entangler

Entangler
(16)

## Interference only occurs if final state particles are entangled!

"What's so wonderful," Cotler says, "is that these contemporary experiments are still pushing the boundaries of our understanding of both quantum mechanics and measurement and opening up new horizons for both theory and experiment."

\author{

- Jordan Cotler
}


## SCIENTIFIC

## AMERICAN Scientists See Quantum Interference between Different Kinds of Particles for First Time

A newly discovered interaction related to quantum entanglement between dissimilar particles opens a new window into the nuclei of atoms

## Thank you for your attention! I hope you can at least say:

Before I came here I was confused about this subject. Having listened to your lecture I am still confused. But on a higher level.



## $\mathbb{N e u t r o n ~ S k i n s ~ a t ~ H i g h - E n e r g y ~}$



$$
\begin{gathered}
S_{U}=0.44 \pm 0.05 \text { (stat.) } \\
\pm 0.08 \text { (syst.) fm }
\end{gathered}
$$

- Uranium neutron skin appears surprisingly large?
- Above trend and lowenergy measurements?


## Which Radius is "correct'?

Now instead of $p_{x}$ and $p_{y}$ lets look at $|t|$ with a 2D approach


STAR: Au+Au $\sqrt{\mathrm{s}_{\mathrm{NN}}}=200 \mathrm{GeV}$


- Drastically different radius depending on $\phi$, still way too big
- Notice how much better the Woods-Saxon dip is resolved for $\phi=\pi / 2$-> experimentally able to remove photon momentum, which blurs diffraction pattern
- Can we extract the 'true' nuclear radius from $|\mathrm{t}|$ vs. $\phi$ information?


## Analogy to

Interferometry in
Astro-Physics

## Quantum Interference provides subdiffraction limited imaging

M87 Supermassive Black hole

## Analogy to

Interferometry in Astro-Physics

Quantum Interference provides subdiffraction limited imaging



Access to details of gluon distribution and neutron skin at high energy

Nuclear Gluon distribution

## $\mathbb{N e u t r o n ~ S k i n s ~ a t ~ H i g h - E n e r g y ~}$

## Uranium



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## Robust <br> Theoretical <br> Description

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- Sensitivity to nuclear geometry!


## $\boldsymbol{\beta}_{2}$




STAR: Photonuclear $\rho^{0} \rightarrow \pi^{+} \pi^{-}$


Interference pattern used for diffraction tomography of gluon distribution $\rightarrow$ analog to $x$-ray diffraction tomography

First high-energy measurements of gluon distribution with sub-femtometer resolution

Technique provides quantitative access to gluon saturation effects
BUT measurements via other vector mesons are needed for to validate QCD theoretical predictions/interpretations
Future measurements with $\phi$ meson and J/ $\psi$ are important
$\mathbb{N u c l e a r}$ Radius Comparison

| $\mathrm{Au}+\mathrm{Au}(\mathrm{fm})$ | U+U (fm) |
| :---: | :---: |
| 6.38 (long: 6.58, short: 6.05 ) | 6.81 (long: 8.01, short: 6.23) |
| $7.95 \pm 0.03$ | -- |
| $7.47 \pm 0.03$ | $7.98 \pm 0.03$ |
| $6.53 \pm 0.03$ (stat.) $\pm 0.05$ (syst.) | $7.29 \pm 0.06$ (stat.) $\pm 0.05$ (syst.) |
| $6.45 \pm 0.27$ | $6.90 \pm 0.14$ |
| $6.74 \pm 0.06$ | -- |
| $\begin{gathered} 0.17 \pm 0.03 \text { (stat.) } \pm 0.08 \text { (syst.) } \\ \sim 2 \sigma \end{gathered}$ | $\begin{aligned} & 0.44 \pm 0.05 \text { (stat.) } \pm 0.08 \text { (syst.) } \\ & \sim 4.7 \sigma \quad \text { (Note: for } \mathrm{Pb} \approx 0.3 \text { ) } \end{aligned}$ |

Precision measurement of nuclear interaction radius at high-energy Measured radius of Uranium shows evidence of significant neutron skin
[1] STAR Collaboration, L. Adamczyk, et al., Phys. Rev. C 96, 054904 (2017).
[2] H. Alvensleben, et al., Phys. Rev. Lett. 24, 786 (1970). [3] G. McClellan, et al., Phys. Rev. D 4, 2683 (1971).
$\mathbb{N u c l e a r}$ Radius Comparison
$\mathrm{Au}+\mathrm{Au}(\mathrm{fm}) \quad \mathrm{U}+\mathrm{U}(\mathrm{fm})$
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Inclusive |t| slope (STAR 2017) [1]
$7.95 \pm 0.03$
Inclusive |t| slope (WSFF fit)*
$7.47 \pm 0.03$
$7.98 \pm 0.03$
Tomographic technique*
$6.53 \pm 0.03$ (stat.) $\pm 0.05$ (syst.)
$7.29 \pm 0.06$ (stat.) $\pm 0.05$ (syst.)
DESY [2]
$6.45 \pm 0.27$
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Cornell [3]
$6.74 \pm 0.06$
$0.17 \pm 0.03$ (stat.) $\pm 0.08$ (syst.)
$0.44 \pm 0.05$ (stat.) $\pm 0.08$ (syst.)
(Tomographic Technique)
~ $2 \sigma$
$\sim 4.7 \sigma$
(Note: for $\mathrm{Pb} \approx 0.3$ )

## Precision measurement of nuclear interaction radius at high-energy Measured radius of Uranium shows evidence of significant neutron skin <br> [1] STAR Collaboration, L. Adamczyk, et al., Phys. Rev. C 96, 054904 (2017). <br> [2] H. Alvensleben, et al., Phys. Rev. Lett. 24, 786 (1970). [3] G. McClellan, et al., Phys. Rev. D 4, 2683 (1971).

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\begin{gathered}
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\pm \mathbf{0 . 0 8} \text { (syst.) fm } \\
\boldsymbol{S}_{\boldsymbol{A} \boldsymbol{u}}^{\boldsymbol{M R}-\boldsymbol{E D F}}=\mathbf{0 . 1 7} \mathbf{f m} \\
\text { Bally, B., Giacalone, G. \& Bender, M. } \\
\text { Eur. Phys. J. A 59, 58 (2023). }
\end{gathered}
$$

- Gold agrees well with state-of-the-art energy density functional calculations
- Consistent with trend from low energy measurements


## $\mathbb{N e u t r o n ~ S k i n s ~ a t ~ H i g h - E n e r g y ~}$

## Uranium



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## $\boldsymbol{\beta}_{2}$




## Confirmation from ALICE (New at QM Sept 2023)



Neutron emission categories test the impact parameter dependence


- New STAR measurementof $\Psi$ $J / \psi$ at QM in Sept 2023
- Consistent within error with Diffraction + Interference (Diff+Int) effect at low $p_{T}$
- Effect of Soft Photon radiation (Rad) visible at higher $p_{T}$



## Access to Hadronic Light-by-Light



Interference with the hadronic light-by-light diagram
Leads to a unique signature -> odd spin configurations

Contribution fromPladronic Vacuum Polarization and Hadronic Light-by-Light are the largest theoretical uncertainties for Standard Model muon g-2


## Elliptic Glluon Tomography (Tensor Pomeron)



Phys. Rev. D 104, 094021 (2021)
Elliptic gluon distribution: correlation
between impact parameter and momentum

- Clear signature of elliptic gluon distribution within nuclei.
- Complimentary measurements at RHIC and EIC




## Shinning lisht on Grluons

- Photo-nuclear measurements have been used to study QCD matter already for decades[1-3]
[1] H1 Collaboration. J. High Energ. Phys. 2010, 32 (2010) [2] ZEUS Collaboration. Eur. Phys. J. C 2, 247-267 (1998).
[3] See refs 1-25 in [2]


Photon energies $\gtrsim 10 \mathrm{GeV}$ : probe gluon distribution - Interaction through
Pomeron (two gluon state at lowest order)

## Shining lisht on Gluons

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The amplitude has three components:

$$
\begin{aligned}
& T^{\gamma^{\star} p \rightarrow V p}(x ; t)=\int_{0}^{1} \mathrm{~d} z \int \mathrm{~d}^{2} \mathbf{r} \Psi^{\gamma}(z, \mathbf{r}) \cdot \sigma^{q \bar{q}-p}(x, \mathbf{r} ; t) \cdot \\
& \text { Photon } \begin{array}{l}
\text { Diffractive } \\
\\
\\
\text { Dipole }
\end{array} \\
& \begin{array}{l}
\text { Vector } \\
V \\
\text { Deson }
\end{array} \\
& \hline
\end{aligned}
$$

Photon quantum numbers $J^{P C}=1^{--}$: Can transform into a 'heavy photon' i.e. a vector meson $\left(\rho^{0}, \phi, J / \psi\right)$ with $J^{P}=1^{-}$

Hanbury Brown and Twiss effect is a two (identical) particle interference due to quantum statistics


States must be identical to interfere, otherwise incoherent sum:
$\left.\left|D_{1 A} D_{2 B}\right| \mathrm{RB}\right\rangle+\left.D_{2 A} D_{1 B}|\mathrm{BR}\rangle\right|^{2}=\left|D_{1 A} D_{2 B}\right|^{2}+\left|D_{2 A} D_{1 B}\right|^{2}$


After entangling interference is restored:

$$
\left|D_{1 A}\right|^{2}\left|D_{2 B}\right|^{2}+\left|D_{2 A}\right|^{2}\left|D_{1 B}\right|^{2}+2 \operatorname{Re} D_{1 A} D_{2 B} D_{2 A}^{*} D_{1 B}^{*}
$$

## The Breit-Wheeler Process


$-\vec{E}--\vec{B} \quad \otimes z:$ Beam Direction


For decades it was believed the polarization info was lost due to random event-by-event orientation!


## Experimental access to photon polarization demonstrated

(STAR Collaboration)
Phys. Rev. Lett. 127, 052302 (2021)


- Polarization vector $\xi$ : aligned radially with the "emitting" source
- Intrinsic photon spin converted into orbital angular momentum
- Observable as anisotropy in $e^{ \pm}$ momentum
S. Bragin, et. al., Phys. Rev. Lett. 119 (2017), 250403
R. P. Mignani, et al., Mon. Not. Roy. Astron. Soc. 465 (2017), 492


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VM inherits the spin from photon (no helicity flip)
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Measurements from H1, ZEUS etc. explored proton via diffractive $\rho^{0}$ and $\phi$ production

## Past Photo-Nuclear $\mathbb{M}$ easurements

- STAR has studied $\gamma \mathbb{P} \rightarrow \rho^{0} \rightarrow \pi^{+} \pi^{-}$(and direct $\pi^{+} \pi^{-}$production) in the


Line shape results from
amplitude level interference:
$\rho^{0} \rightarrow \pi^{+} \pi^{-}+$Drell Söding
(direct $\pi^{+} \pi^{-}$) $+\omega \rightarrow \pi^{+} \pi^{-}$
$\propto\left|\frac{\sqrt{m_{\pi \pi} m_{\rho} \Gamma\left(m_{\pi \pi}\right)}}{m_{\rho}^{2}-m_{\pi \pi}^{2}+i m_{\rho} \Gamma\left(m_{\pi \pi}\right)}+\frac{f_{I}}{2}\right|^{2}$,

STAR Collaboration et al. Phys. Rev. Lett. 89, 272302 (2002). STAR Collaboration et al. Phys. Rev. Lett. 102, 112301 (2009). STAR Collaboration et al. Phys. Rev. C 96, 054904 (2017).

I will take just this one experiment, which has been designed to contain all of the mystery of quantum mechanics, ... Any other situation in quantum mechanics, it turns out, can allways be explained by saying, "Your remember the case of the experiment with the two holes? It's the same thing .
-Richard Feynman


Whem:



