



# Measurements of p- $\Lambda$ and d- $\Lambda$ correlations in 3 GeV Au+Au collisions at STAR

Zhi Qin (秦 智)

for the STAR collaboration  
*Tsinghua University*



2023.11.10



Supported in part by  
U.S. DEPARTMENT OF  
**ENERGY**

# QCD Dense Matter & Nucleon-Nucleon/Hyperon Interactions

## INSIDE A NEUTRON STAR

A NASA mission will use X-ray spectroscopy to gather clues about the interior of neutron stars — the Universe's densest forms of matter.

### Outer crust

Atomic nuclei, free electrons

### Inner crust

Heavier atomic nuclei, free neutrons and electrons

### Outer core

Quantum liquid where neutrons, protons and electrons exist in a soup

### Inner core

Unknown ultra-dense matter. Neutrons and protons may remain as particles, break down into their constituent quarks, or even become 'hyperons'.

### Atmosphere

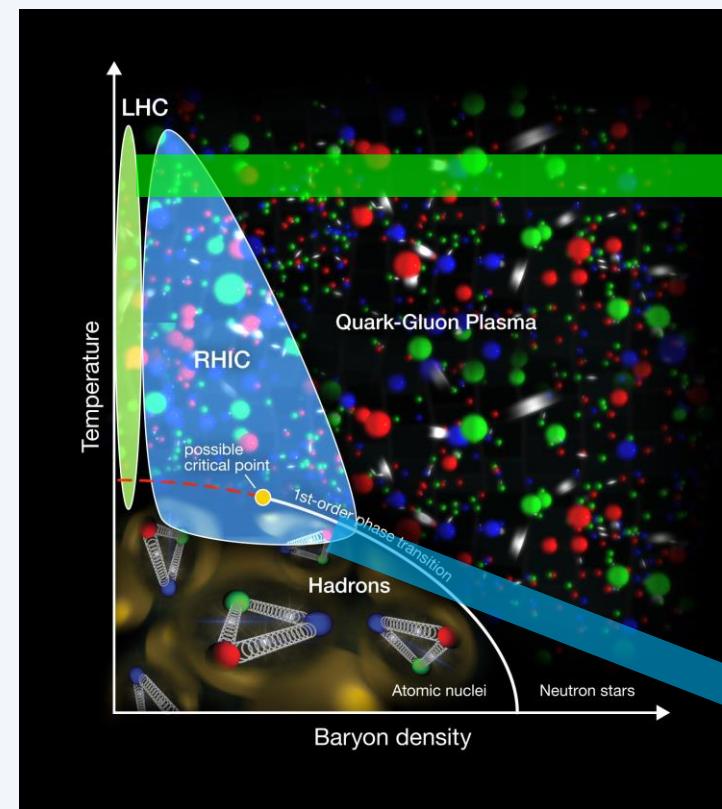
Hydrogen, helium, carbon

Beam of X-rays coming from the neutron star's poles, which sweeps around as the star rotates.

©nature

Credit: Source: Adapted from NASA Goddard SVS  
Nature volume 546, page18 (2017)

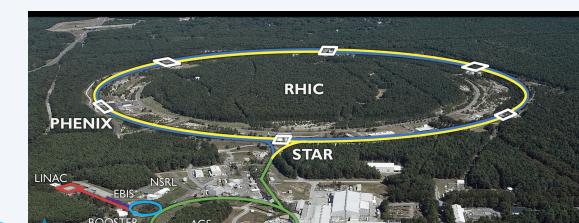
- ❖ Heavy ion collisions – laboratory for dense QCD matter
- ❖ Role of Nucleon-Nucleon (N-N) and Hyperon-Nucleon (Y-N) interactions in the Equation-of-State



<https://www.bnl.gov/newsroom/news.php?a=219079>

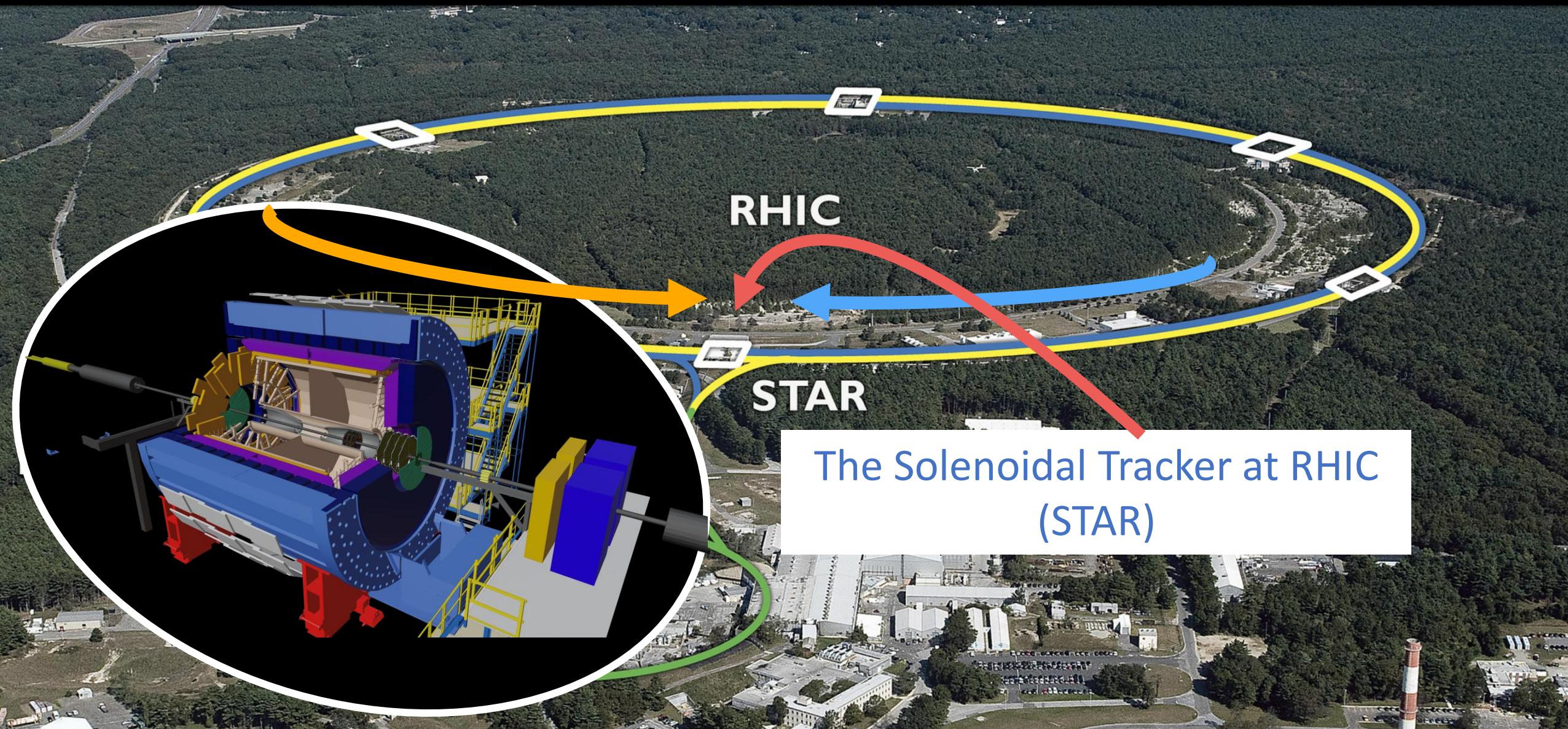


The Large Hadron Collider (LHC)

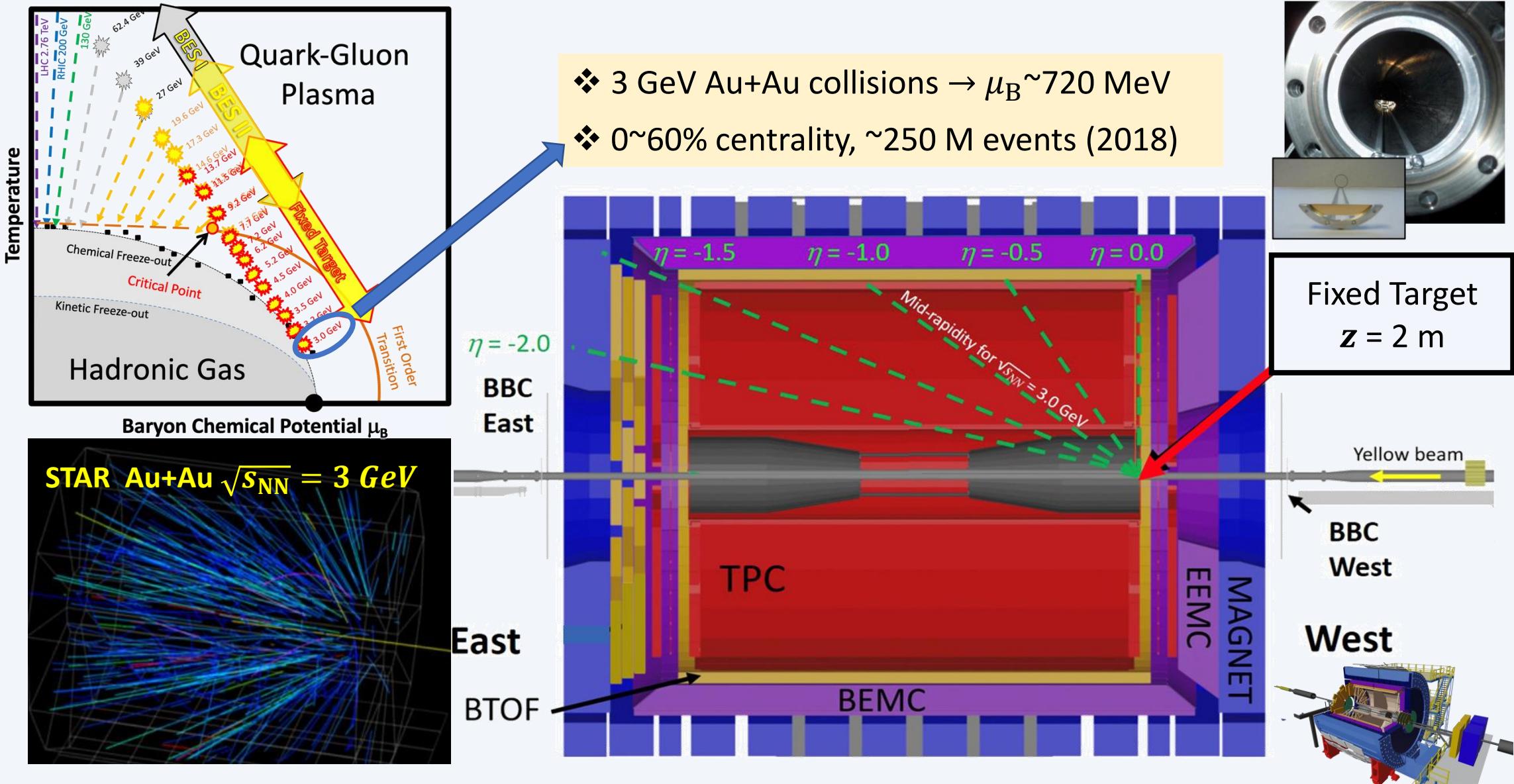


Relativistic Heavy Ion Collider (RHIC)

# STAR Detector



# Beam Energy Scan – II & Fixed Target Setup



# Baryon Correlation Function (CF)

Momentum correlation function:

$$C(\mathbf{p}_1, \mathbf{p}_2) \equiv \frac{P(\mathbf{p}_1, \mathbf{p}_2)}{P(\mathbf{p}_1) \cdot P(\mathbf{p}_2)}$$

Single-particle momentum

*Statistical*

Approximating the emission process and the momenta of the particles:

$$C(\mathbf{k}^*) = \int d^3r^* S(r^*) |\Psi(r^*, \mathbf{k}^*)|^2$$

*Modeling*

Distribution of the relative distance of particle pair

Relative wave function of the particle pair

*Experimental*

$$C(k^*) = \mathcal{N} \frac{A(k^*)}{B(k^*)}$$

Signal

Background

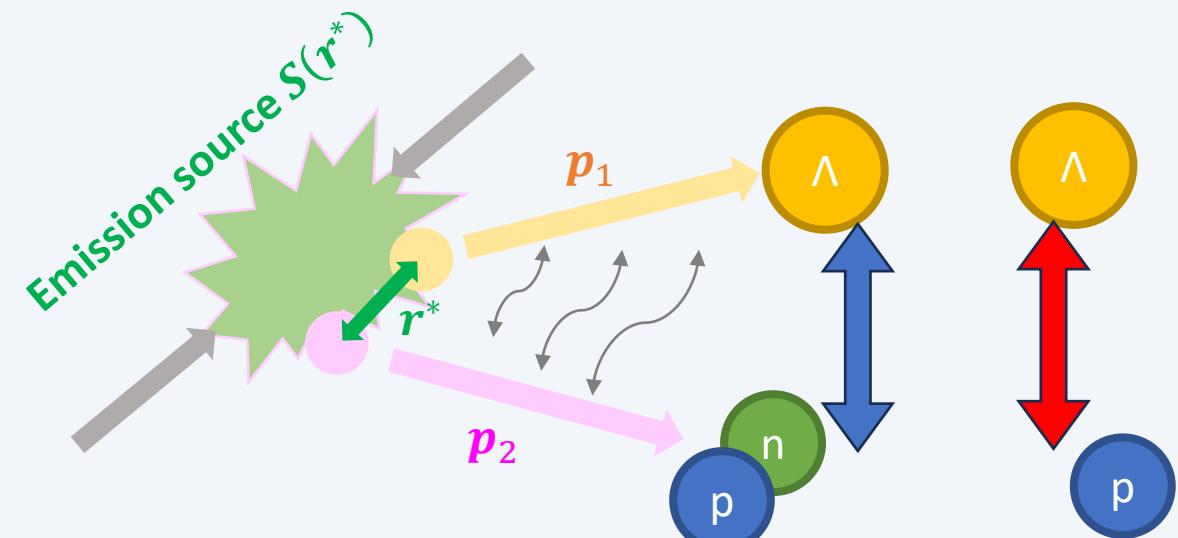
Normalization factor

$k^*$ : particle momentum in the pair rest frame

R. Lednický, et al. Sov.J.Nucl.Phys. 35 (1982) 770

L. Michael, et al. Ann.Rev.Nucl.Part.Sci. 55 (2005) 357-402

J. Haidenbauer, Phys.Rev.C 102 (2020) 3, 034001



- ❖ Space and time evolution of particle-emitting source
  - ❖ Final state interactions
- p- $\Lambda$  & d- $\Lambda$  correlations:  
N(-N)-Y interactions / hypernuclei structure

# Methodology

## Experimental

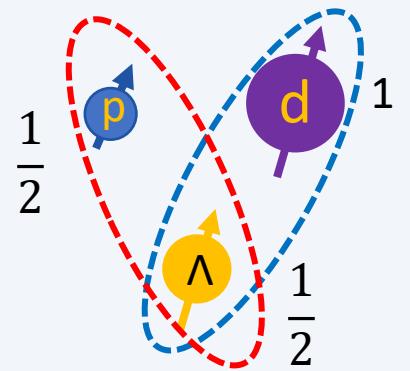
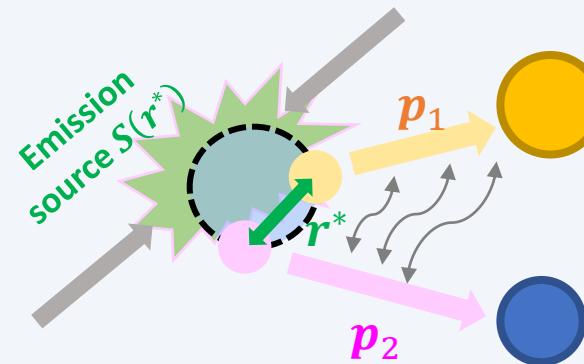
$$C(k^*) = \mathcal{N} \frac{A(k^*)}{B(k^*)}$$

Same events
Mixed events

Normalization factor

## Corrections on:

- ❖ Purity
- ❖ Feed-down effect
- ❖ Track splitting & merging
- ❖ Momentum resolution



Singlet State	$^1S_0$	(S)
Triplet State	$^3S_1$	(T)
Doublet State	$^2S_{1/2}$	(D)
Quartet State	$^4S_{3/2}$	(Q)

## Physics Image Extraction

Formalism with  
Lednicky-Lyuboshitz  
(L-L) approach

$R_G$  : spherical Gaussian source of pairs  
 $f_0$  : scattering length  
 $d_0$  : effective range

## Major assumptions:

- ❖ Ignore multiplicity correlations and charge of source
- ❖ Smoothness approximation
- ❖ Effective range expansion for  $\Psi(\mathbf{r}^*, \mathbf{k}^*)$
- ❖ Static and spherical Gaussian source

- Single particle source:  $S_i(x_i, p_i^*)$
- Pair source (radius  $R_G$ ):  $S(x, p^*) \propto e^{-x^2/4R_G^2} \delta(t - t_0)$

$$C(k^*) \approx 1 + \frac{|f(k)|^2}{2R_G^2} F(d_0) + \frac{2\text{Re}f(k)}{\sqrt{\pi}R_G} F_1(2kR) - \frac{\text{Im}f(k)}{R_G} F_2(2kR_G)$$

$$\frac{1}{f(k)} \approx \frac{1}{f_0} + \frac{d_0 k^2}{2} - ik$$

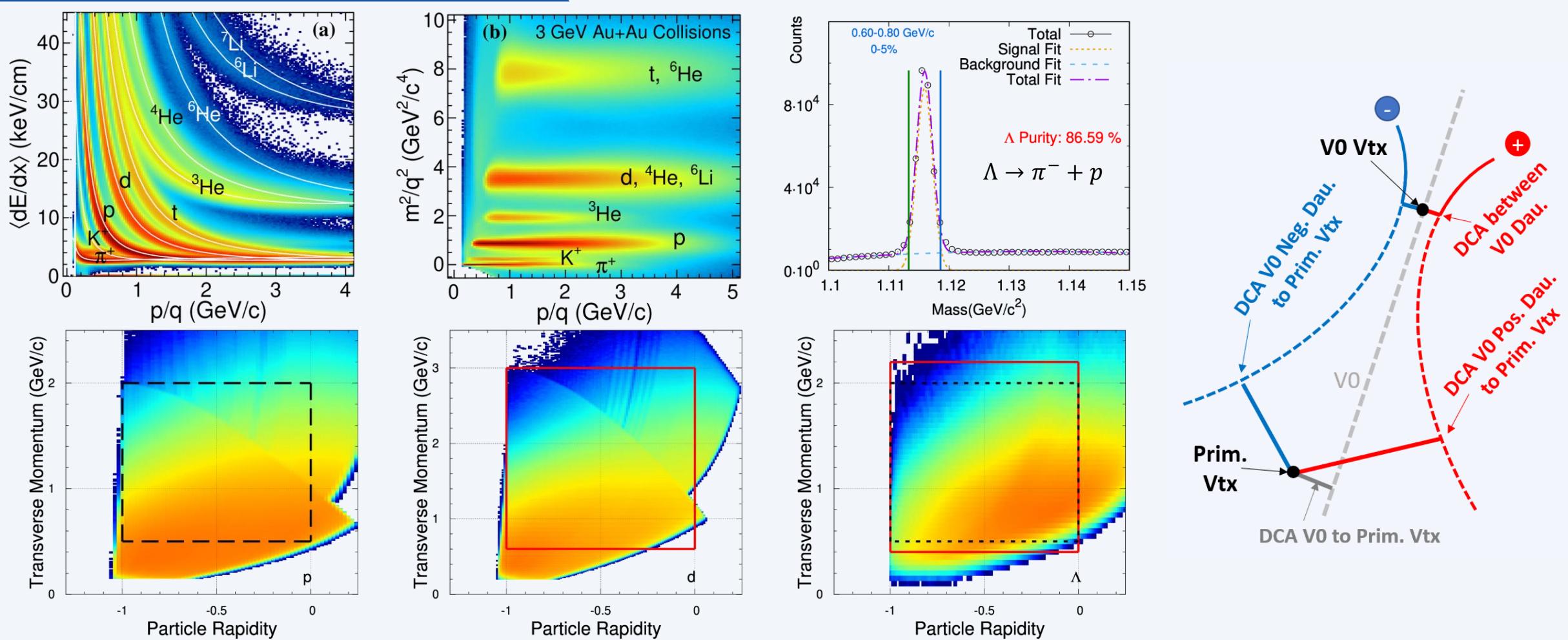
Different  $f_0$  and  $d_0$  for different spin states

R. Lednicky, et al. Sov.J.Nucl.Phys. 35 (1982) 770

L. Michael, et al. Ann.Rev.Nucl.Part.Sci. 55 (2005) 357-402

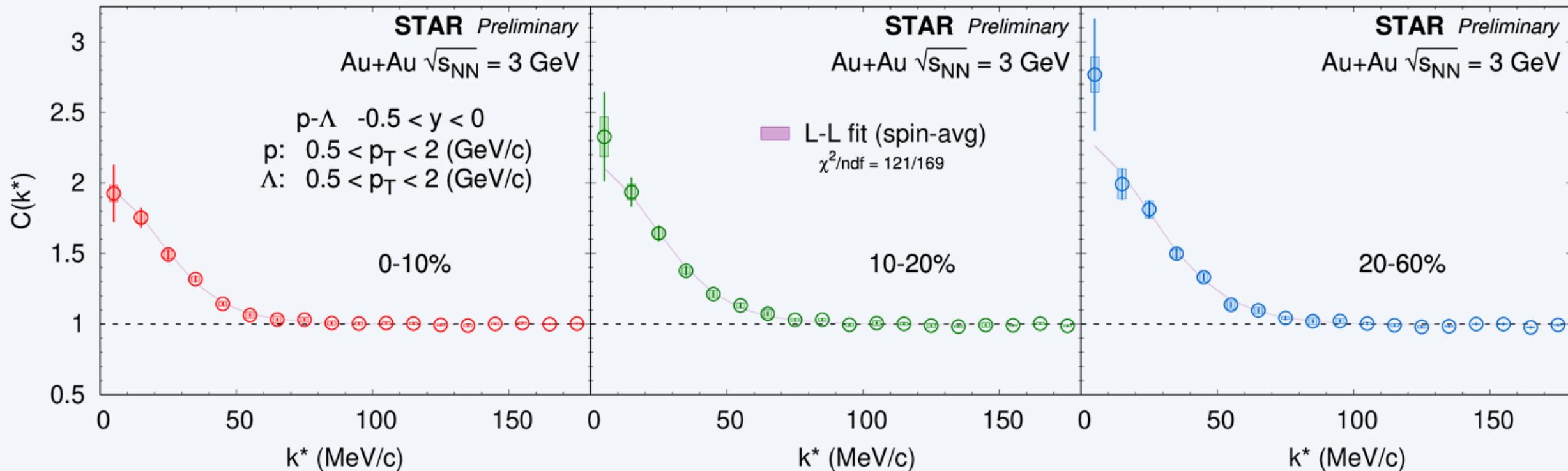
J. Haidenbauer, Phys.Rev.C 102 (2020) 3, 034001

# Particle Identification & Reconstruction @ 3 GeV



- ❖  $\pi^-$ , p, and d particles are identified by TPC and TOF
- ❖ A larger acceptance is used in d- $\Lambda$  correlation measurement (red) due to statistics

# p- $\Lambda$ Correlation Measurement @ STAR

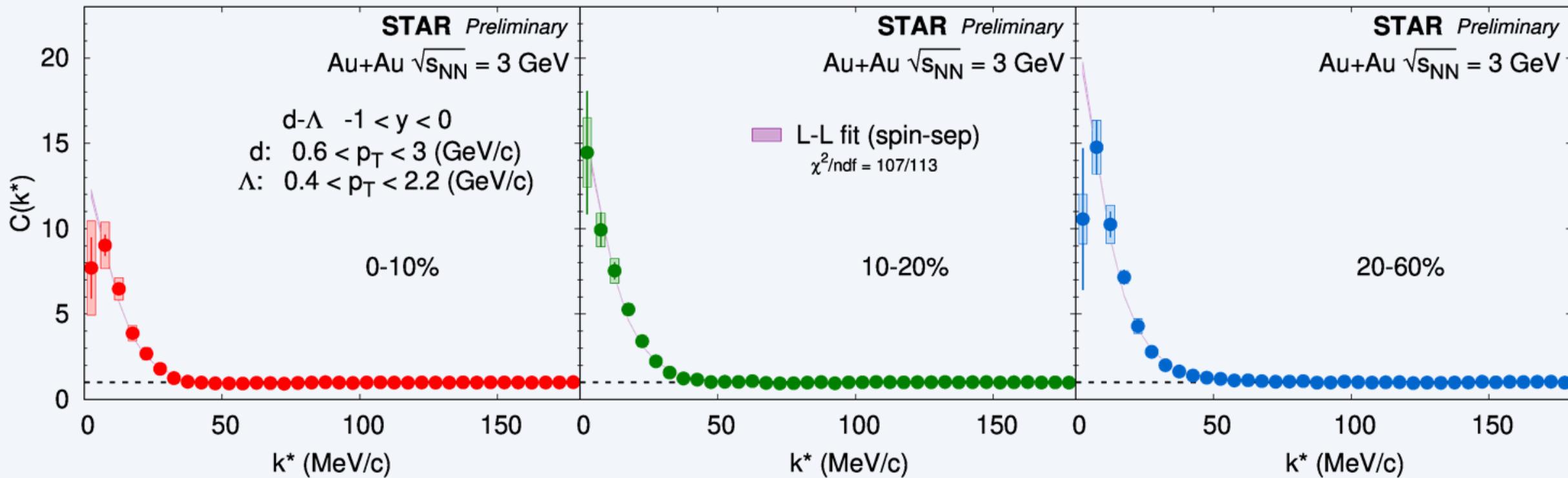


## Corrections

1. Purity correction
2.  $\Lambda$  feed-down correction
3. Track splitting & merging
4. Momentum smearing effect

- ❖ Simultaneous fit to data in different centralities/rapidity
    - ❖  $R_G^i$ , spin-avg  $f_0$  and  $d_0$  with Lednicky-Lyuboshitz approach
  - ❖ Spin-avg scattering length ( $f_0$ ) and effective range ( $d_0$ ):
- $f_0 = 2.32^{+0.12}_{-0.11} \text{ fm}$        $d_0 = 3.5^{+2.7}_{-1.3} \text{ fm}$

# d- $\Lambda$ Correlation Measurement @ STAR



## Corrections

1. Purity correction
2. Track splitting & merging

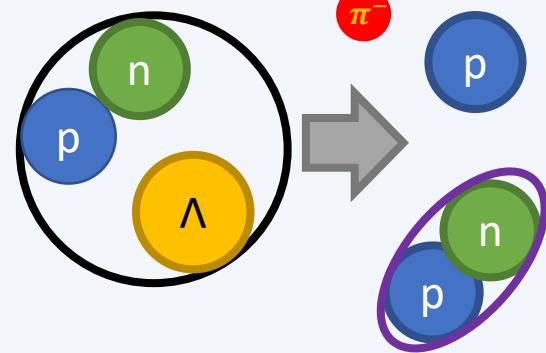
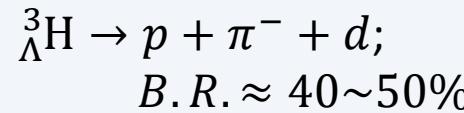
**3. Contamination from**  
 ${}^3\text{H} \rightarrow \pi^- + p + d$  decay

Details in the next page

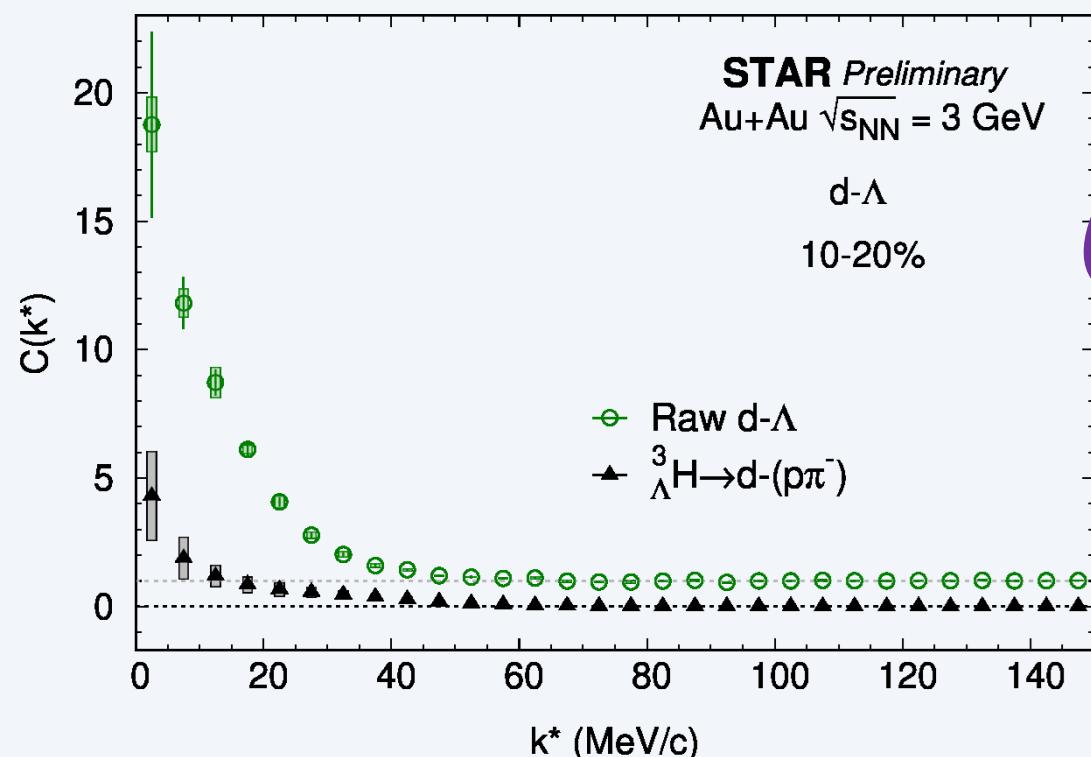
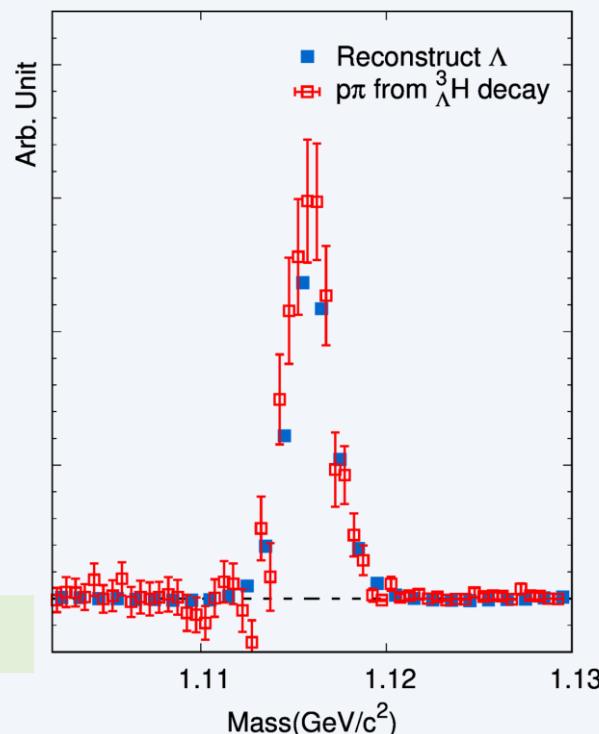
- ❖ First d- $\Lambda$  correlation measurements in the heavy-ion collision experiment
  - ❖ Simultaneous fit to data in different centralities
    - ❖  $R_G^i, f_0(D), d_0(D), f_0(Q)$ , and  $d_0(Q)$  with Lednicky-Lyuboshitz approach
- |                             |                           |
|-----------------------------|---------------------------|
| $f_0(D) = -20^{+3}_{-3}$ fm | $d_0(D) = 3^{+2}_{-1}$ fm |
| $f_0(Q) = 16^{+2}_{-1}$ fm  | $d_0(Q) = 2^{+1}_{-1}$ fm |

- ❖  $\Lambda$  feed-down correction not applied due to unknown d- $\Sigma/\Xi$  correlation
- ❖ Momentum smearing effect negligible

# Contamination Correction from ${}^3_{\Lambda}\text{H} \rightarrow p\pi^- + d$ Decay



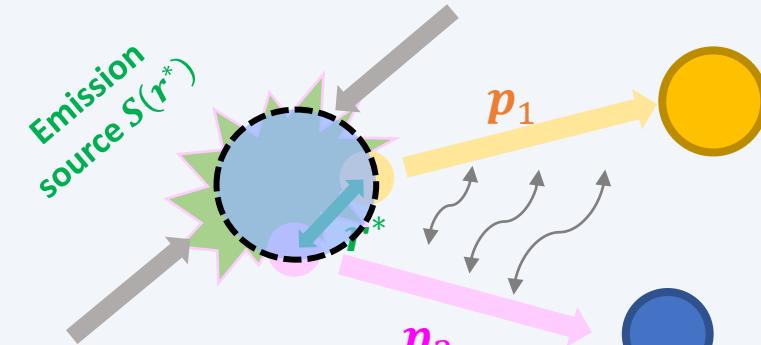
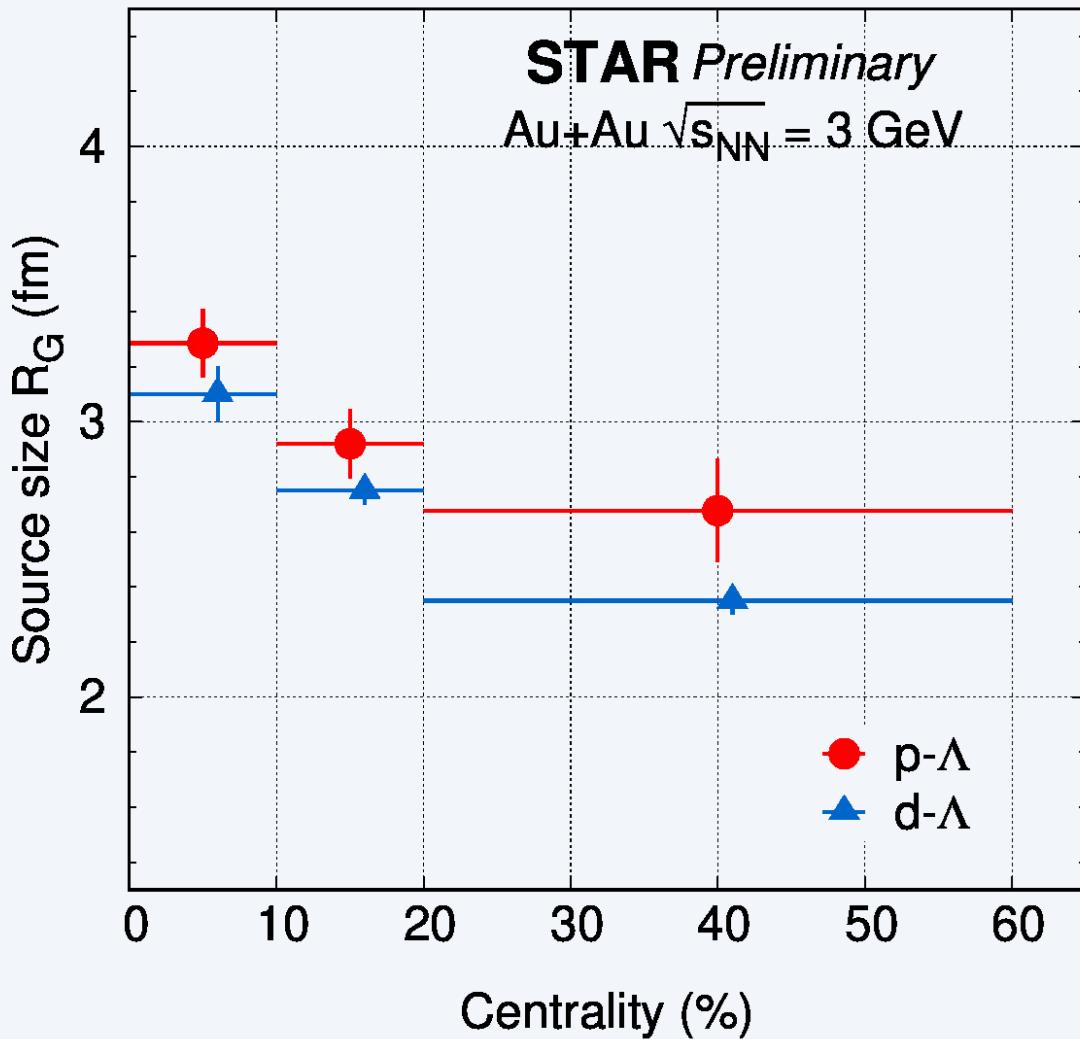
Violation of energy conservation



- The  ${}^3_{\Lambda}\text{H}$  decayed  $p + \pi^-$  are **not experimentally distinguishable** with the reconstructed  $\Lambda$
- $(p\pi^-) - d$  from  ${}^3_{\Lambda}\text{H}$  will affect small  $k^*$  region

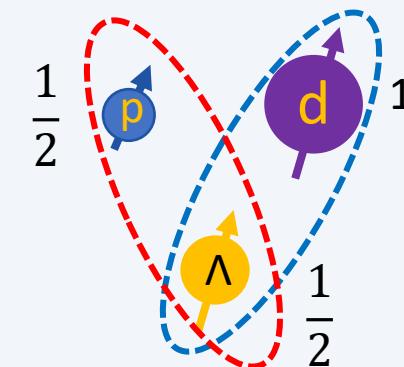
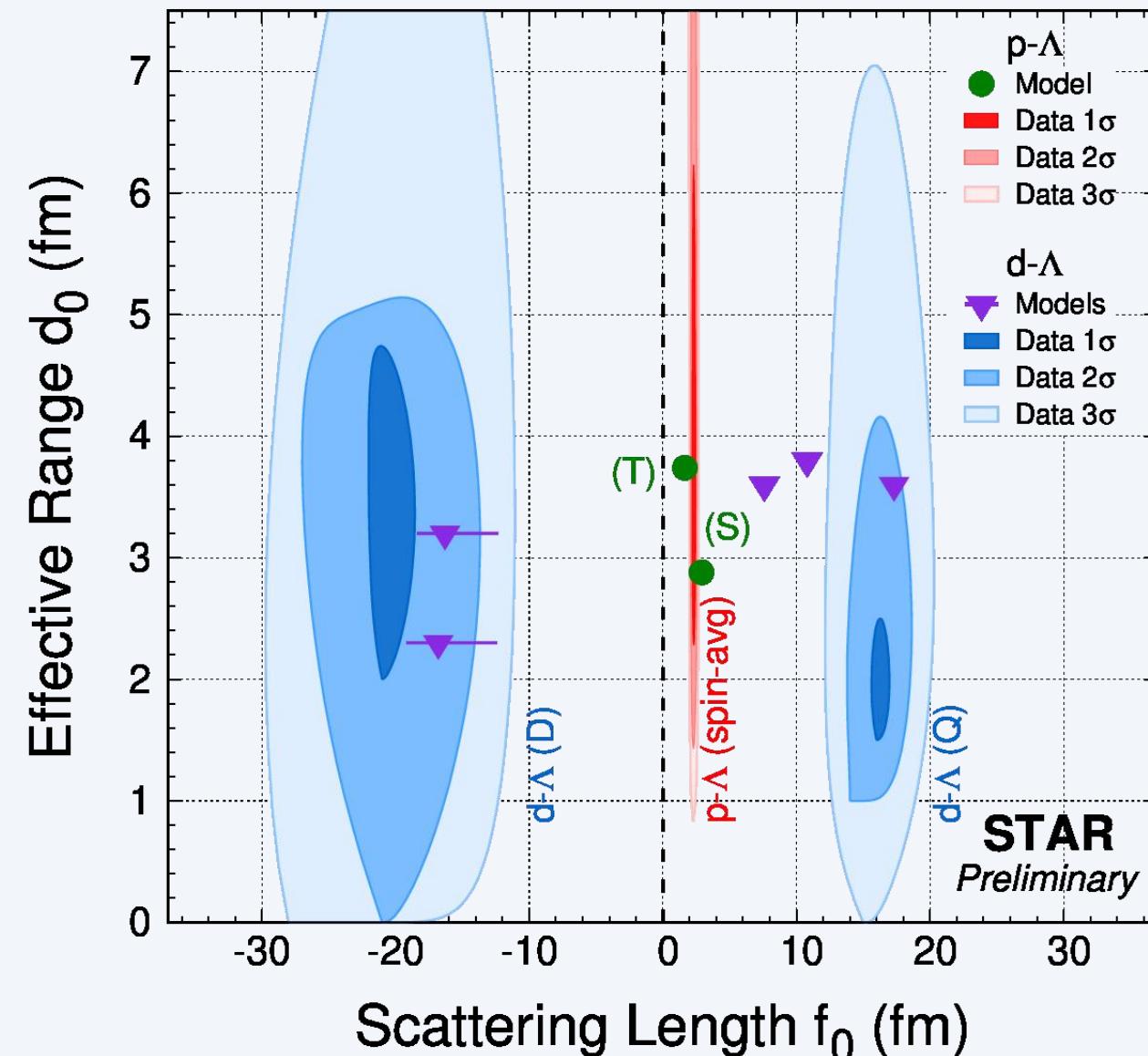
- Simulation based on STAR  ${}^3_{\Lambda}\text{H}$  yield measurement:  
4~8% of d- $\Lambda$  entries from  ${}^3_{\Lambda}\text{H}$  decay at  $k^* < 100$  MeV/c  
in 10~20% centrality
- Contamination subtracted from inclusive d- $\Lambda$  correlation

# Source Size with L-L approach



- ❖  $R_G$ : **spherical Gaussian source of pairs** by Lednicky-Lyuboshits approach
- ❖ Separation of emission source from final state interaction
- ❖ Collision dynamics as expected:
  - ❖  $R_G^{\text{central}} > R_G^{\text{peripheral}}$
  - ❖  $R_G(p - \Lambda) > R_G(d - \Lambda)$

# Scatterings Length ( $f_0$ ) and Effective Range ( $d_0$ )



$$\frac{1}{f(k)} \approx \frac{1}{f_0} + \frac{d_0 k^2}{2} - ik$$

❖ The constraint of the effective range ( $d_0$ ) is weaker

- ❖ The measurement is done at freeze-out
- ❖ Spin-avg for  $f_0$  &  $d_0$  p- $\Lambda$  system

$$f_0 = 2.32^{+0.12}_{-0.11} \text{ fm}$$

$$d_0 = 3.5^{+2.7}_{-1.3} \text{ fm}$$

- ❖ Successfully separate two spin states in d- $\Lambda$

$$f_0(D) = -20^{+3}_{-3} \text{ fm}$$

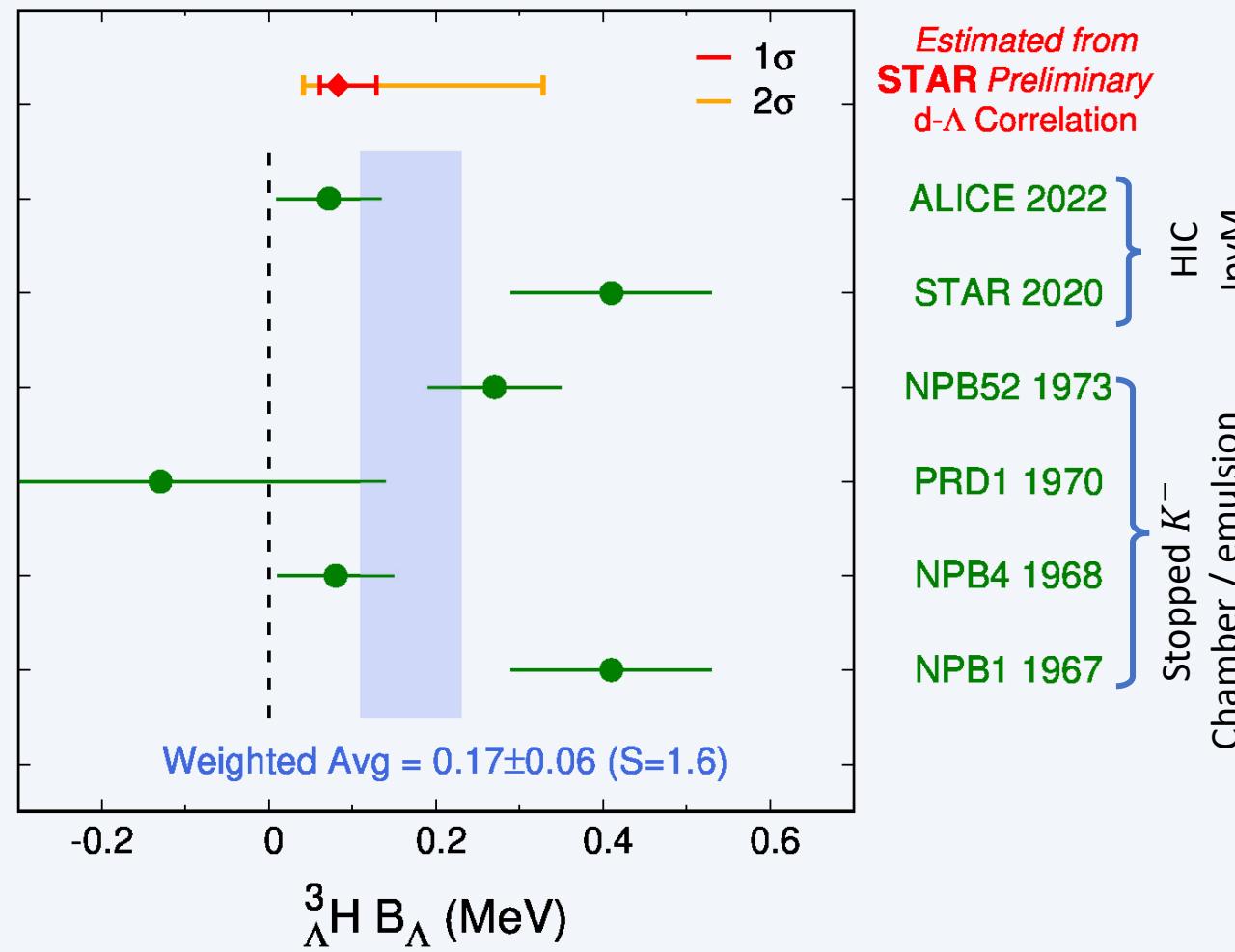
$$d_0(D) = 3^{+2}_{-1} \text{ fm}$$

$$f_0(Q) = 16^{+2}_{-1} \text{ fm}$$

$$d_0(Q) = 2^{+1}_{-1} \text{ fm}$$

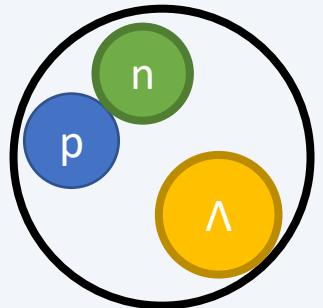
\*Edge of d- $\Lambda$  contours are shown with Bezier smooth to improve the visibility

# ${}^3_{\Lambda}\text{H}$ Binding Energy



${}^3_{\Lambda}\text{H}$  binding energy ( $B_{\Lambda}$ ):

- ❖ Bethe formula from Effective Range Expansion (ERE) parameters  $f_0(D)$  &  $d_0(D)$



$$\frac{1}{-f_0} = \gamma - \frac{1}{2} d_0 \gamma^2 \quad \begin{aligned} \diamond B_{\Lambda} &= \frac{\gamma^2}{2\mu_{d\Lambda}} \\ \diamond \mu_{d\Lambda} &\text{: reduced mass} \\ \diamond \gamma &\text{: binding momentum} \end{aligned}$$

- ❖  ${}^3_{\Lambda}\text{H} B_{\Lambda} = [0.04, 0.33]$  (MeV) @ 95% CL
  - Consistent with the world average
  - ❖ A new way to constrain the  ${}^3_{\Lambda}\text{H}$  structure

# Summary and outlook

- ❖ The first d- $\Lambda$  correlation function measurements in heavy-ion collisions
- ❖ New p- $\Lambda$  correlation function measurements with 3 GeV Au+Au collisions
- ❖ Successfully separated emission source size from final state interactions in p- $\Lambda$  & d- $\Lambda$  correlation functions

1.  $R_G^{\text{central}} > R_G^{\text{peripheral}}$  and  $R_G(p - \Lambda) > R_G(d - \Lambda)$
2. p- $\Lambda$  correlation spin-ave:

$$f_0 = 2.32^{+0.12}_{-0.11} \text{ fm}$$

$$d_0 = 3.5^{+2.7}_{-1.3} \text{ fm}$$

3. d- $\Lambda$  correlation spin-sep:

$$f_0(D) = -20^{+3}_{-3} \text{ fm}$$

$$d_0(D) = 3^{+2}_{-1} \text{ fm}$$

$$f_0(Q) = 16^{+2}_{-1} \text{ fm}$$

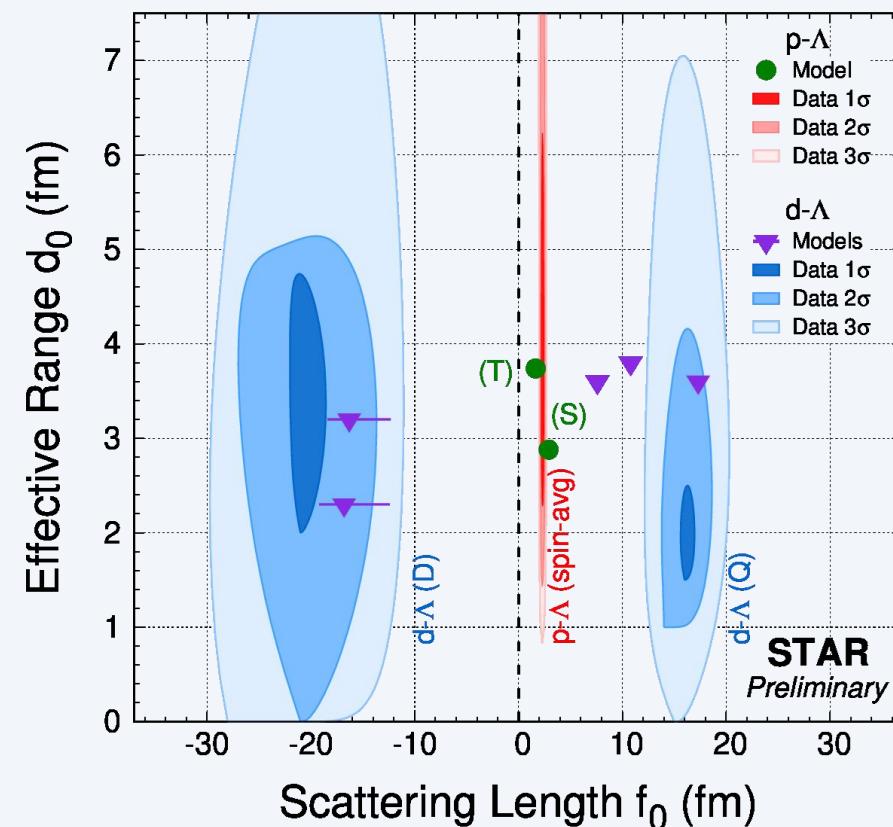
$$d_0(Q) = 2^{+1}_{-1} \text{ fm}$$

4.  ${}^3\text{H} B_\Lambda = [0.04, 0.33] \text{ (MeV)}$  @ 95% CL from d- $\Lambda$  correlation (D)

## Outlook:

More than 10 times statistics from BES-II

- ❖ Emission source size vs. energy, rapidity...
- ❖ Baryon correlations with different species





# Thank you!



Catania (Italy), November 6-10, 2023

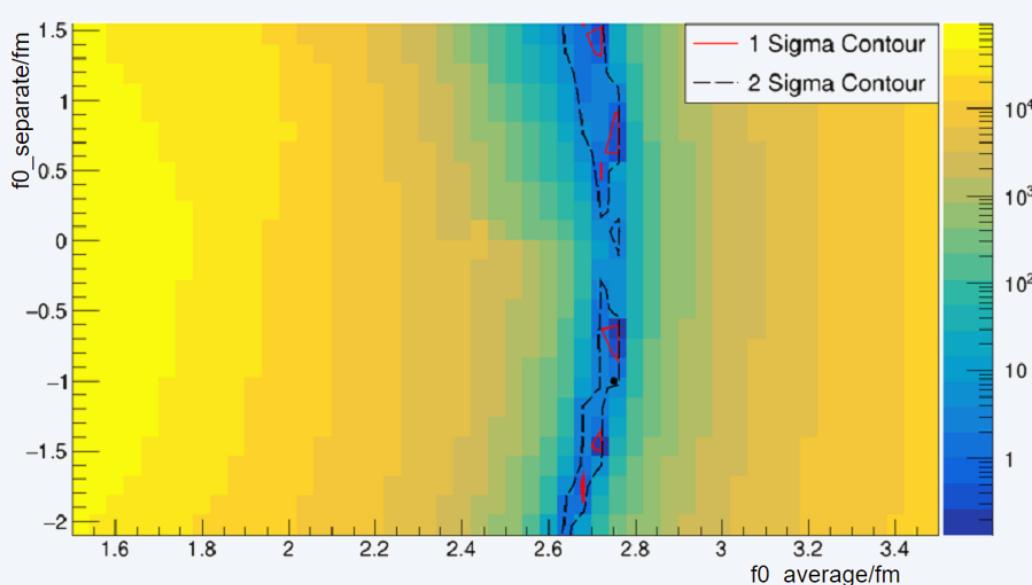
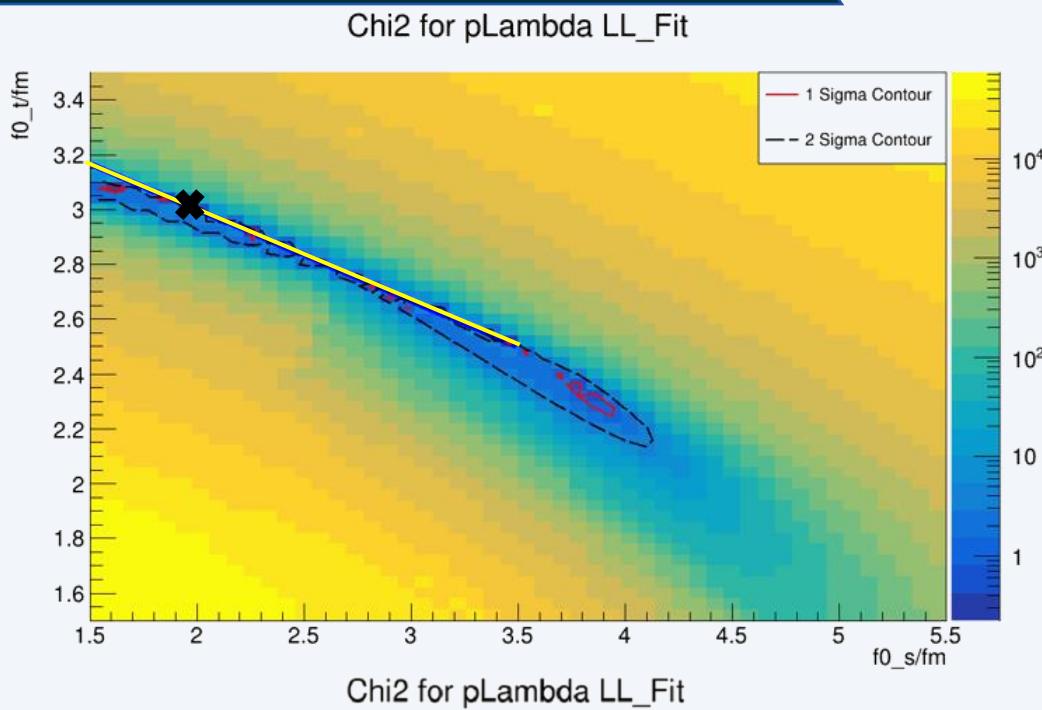
2023.11.10



Supported in part by  
U.S. DEPARTMENT OF  
**ENERGY**

**STAR** ★

# Backup—Spin averaged $f_0$ , $d_0$ for $p\Lambda$



Input parameters(theoretically calculation):

$$r_G = 2.2\text{fm}$$

$$f_{0,s} = 2\text{fm}, d_{0,s} = 2.46\text{fm}$$

$$f_{0,t} = 3\text{fm}, d_{0,t} = 3.99\text{fm}$$

Yellow Straight Line(in top plot):

$$\frac{1}{4}f_{0,s} + \frac{3}{4}f_{0,t} = \frac{1}{4} \times 2\text{fm} + \frac{3}{4} \times 3\text{fm} = 2.75\text{fm}$$

So, theoretically we can't distinguish  $f_0(d_0)$  of two spin from  $p\Lambda$  CF theoretically

We can define:

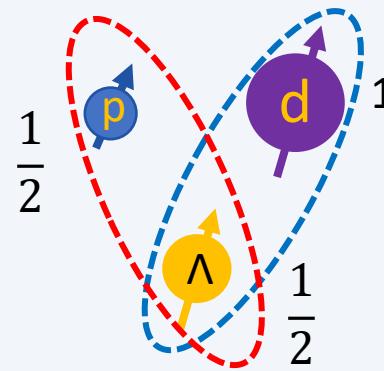
$$f_{0,\text{average}} = \frac{1}{4}f_{0,s} + \frac{3}{4}f_{0,t}$$

$$f_{0,\text{separate}} = f_{0,s} - f_{0,t}$$

and get left plot, it shows that  $f_{0,\text{average}}$  can be constrain well ( $\sim 2.75\text{fm}$ ), and no constrain for  $f_{0,\text{separate}}$

So we use this spin-average spin definition to our data, and fix  $f_{0,\text{separate}} = 0$

# Backup— Correlation Function & Spin States



Singlet State	$^1S_0$	(S)
Triplet State	$^3S_1$	(T)
Doublet State	$^2S_{1/2}$	(D)
Quartet State	$^4S_{3/2}$	(Q)

**p- $\Lambda$ :**  $|\psi(r, k)|^2 \rightarrow \frac{1}{4}|\psi_0(r, k)|^2 + \frac{3}{4}|\psi_1(r, k)|^2$

**d- $\Lambda$ :**  $|\psi(r, k)|^2 \rightarrow \frac{1}{3}|\psi_{1/2}(r, k)|^2 + \frac{2}{3}|\psi_{3/2}(r, k)|^2$

- Different spin states with different  $f_0$  and  $d_0$  parameters
- p- $\Lambda$  correlation:** current statistics is not enough to separate two spin states  $\rightarrow$  spin-averaged fit
- d- $\Lambda$  correlation:** very different  $f_0$  for (D) and (Q) are predicted  $\rightarrow$  **Spin-separated fit**

