

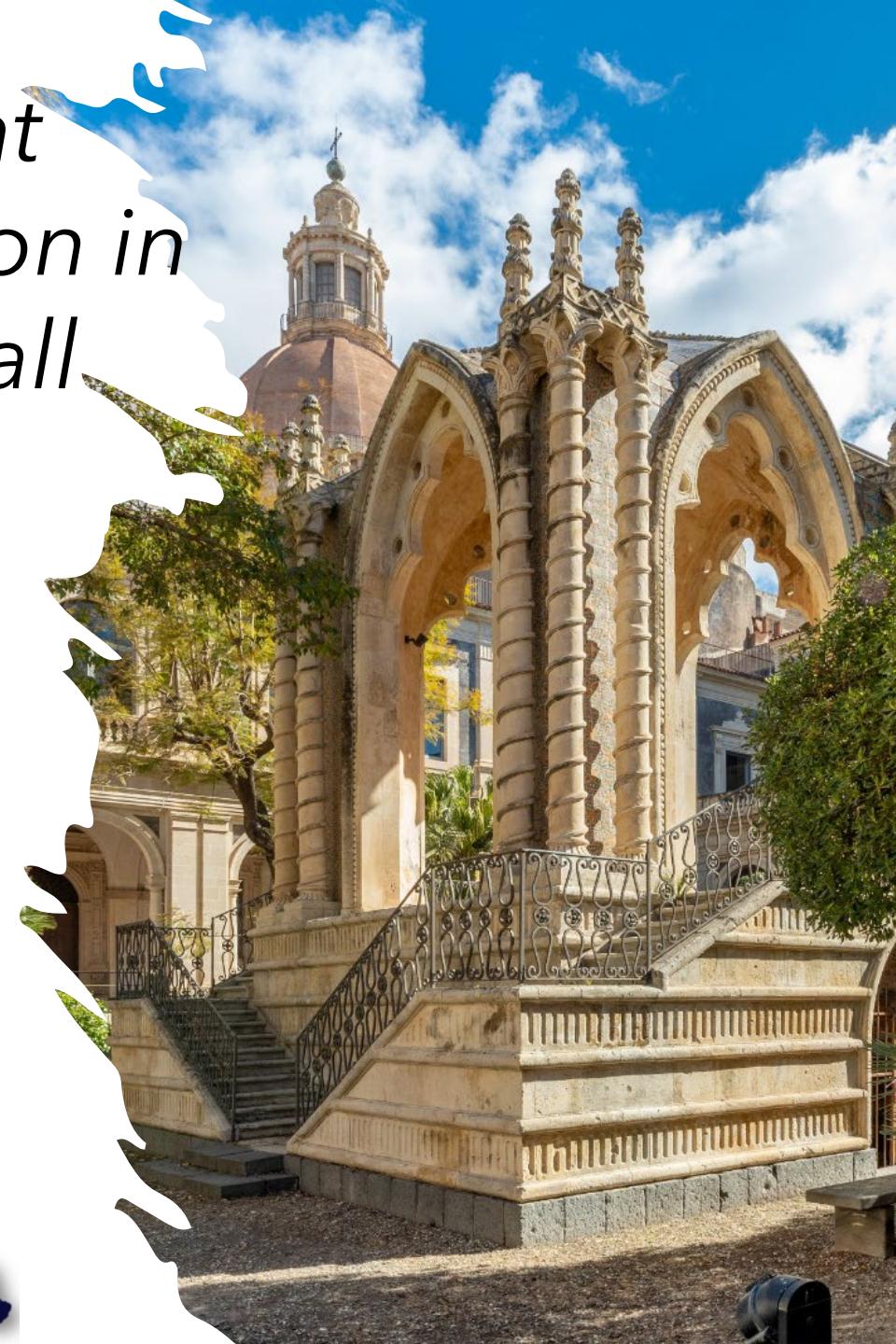
# *Constraining the light (anti)nuclei production in and out of jets in small systems with ALICE*

Marika Rasà<sup>1</sup> on behalf of the  
ALICE Collaboration

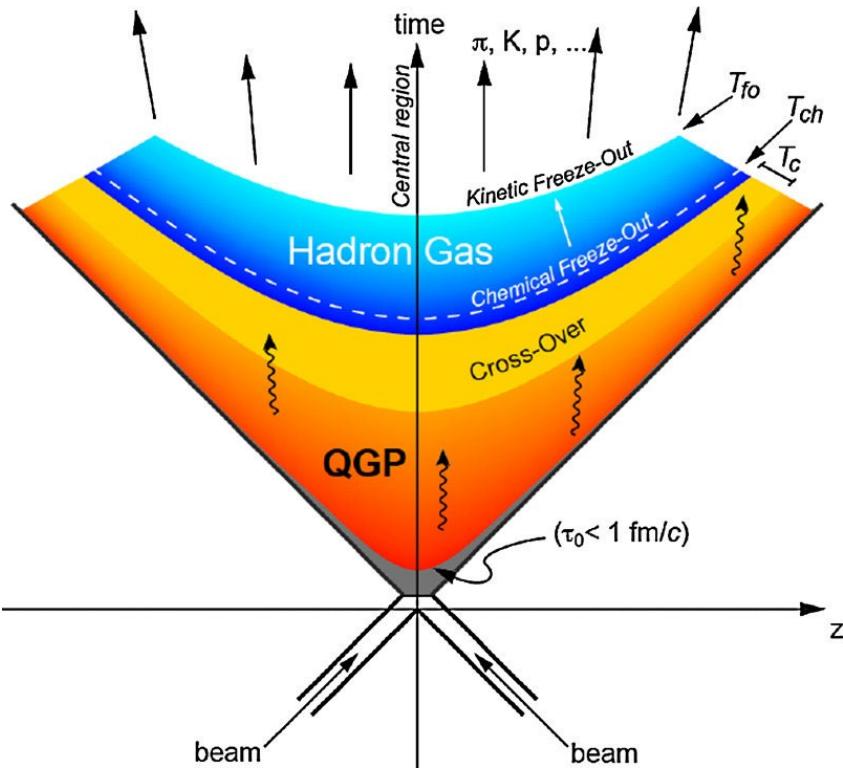
1. University and INFN, Catania



Università  
di Catania

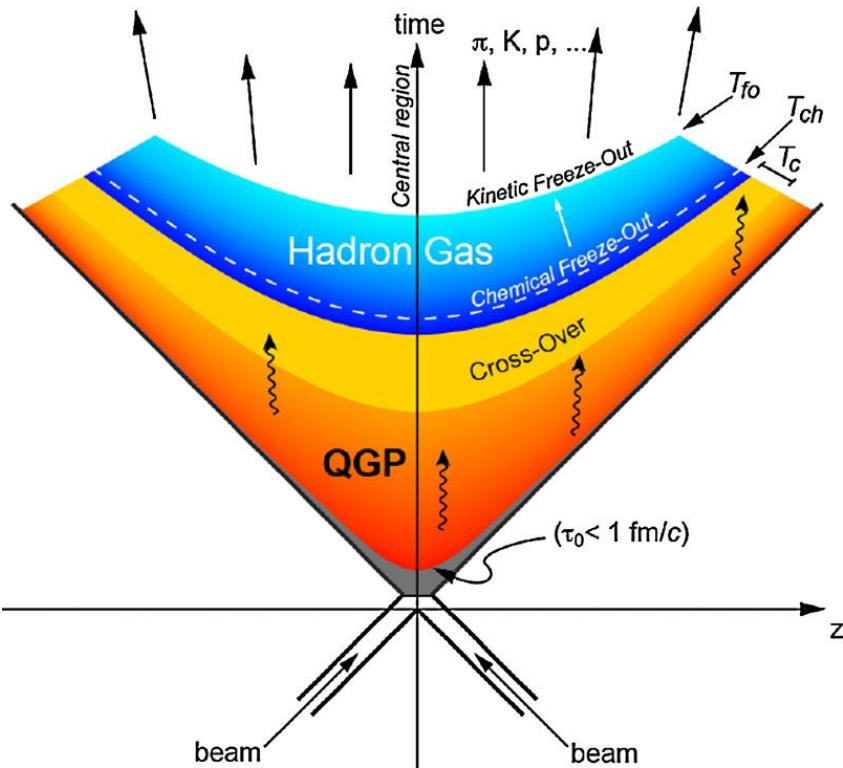


# Why studying light (anti)nuclei?



- Light (anti)nuclei are produced in high-energy hadronic collisions at the LHC
- Their production mechanism is still not understood
- Two phenomenological models:
  - Statistical hadronization
  - Coalescence

# Why studying light (anti)nuclei?



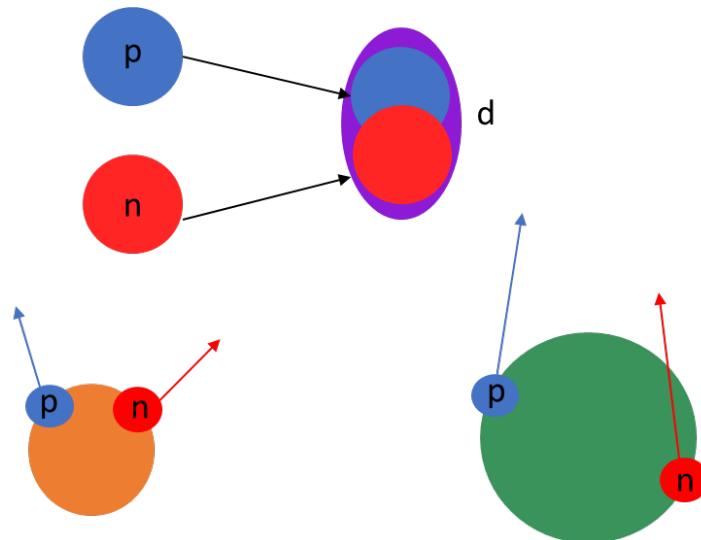
- Light (anti)nuclei are produced in high-energy hadronic collisions at the LHC
- Their production mechanism is still not understood
- Two phenomenological models:
  - Statistical hadronization
  - Coalescence

*Focus on it*

# Simple coalescence model

S. T. Butler et al., *Phys. Rev.* 129 (1963) 836

- If (anti)nucleons are close in phase space and match the spin states, they can form an (anti)nucleus
- Coalescence parameter  $B_A$  is the key observable:



$$E_A \frac{d^3 N_A}{dp_A^3} = B_A \cdot \left( E_p \frac{d^3 N_p}{dp_p^3} \right)^A$$

Invariant yield  
of nucleus

Coalescence  
parameter

Invariant yield  
of protons

# State of the art coalescence model

- Coalescence parameter depends on both the source size and radial extension of the nucleus wave function
- Wigner function formalism

$$N_A = g_a \cdot \int d^3x_1 \dots d^3x_A \cdot d^3k_1 \dots d^3k_A \cdot f_1(x_1, k_1) \dots f_A(x_A, k_A) \cdot W_A(x_1, \dots, x_A, k_1, \dots, k_A)$$

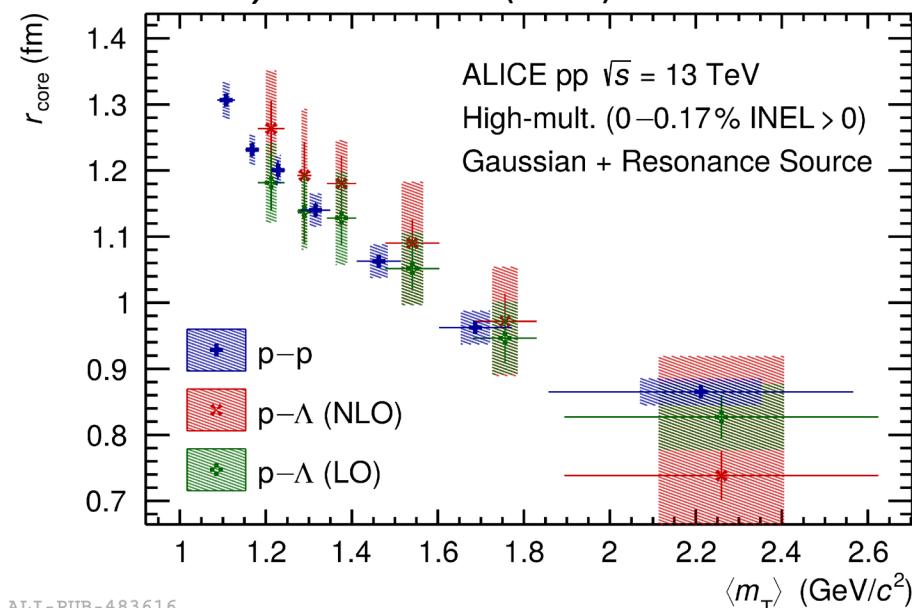
# State of the art coalescence model

- Coalescence parameter depends on both the source size and radial extension of the nucleus wave function
- Wigner function formalism

$$N_A = g_a \cdot \int d^3x_1 \dots d^3x_A \cdot d^3k_1 \dots d^3k_A \cdot f_1(x_1, k_1) \dots f_A(x_A, k_A) \cdot W_A(x_1, \dots, x_A, k_1, \dots, k_A)$$

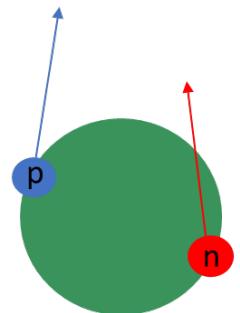
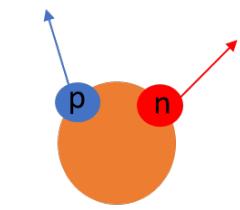
phase space distributions of nucleons → dependence on the source size

*Phys. Lett. B* 811 (2020) 135849



ALI-PUB-483616

- Femtoscopy measurement to the source size available
- Dependence on the source size:
- **Small source size → Large  $B_A$**   
(pp ~ 1 fm, p–Pb ~ 1.5 fm)
- **Large source size → Small  $B_A$**   
(Pb–Pb ~ 3-6 fm)



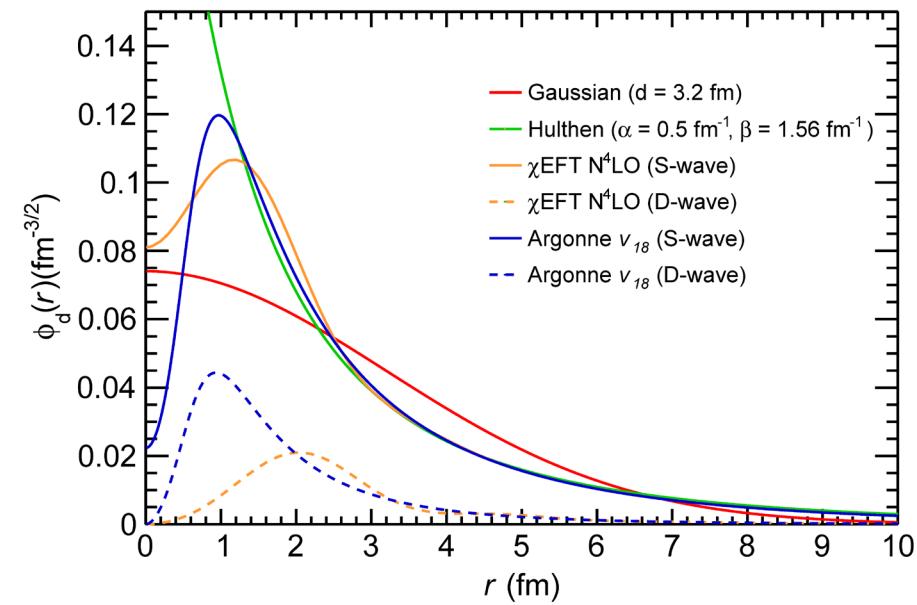
# State of the art coalescence model

- Coalescence parameter depends on both the source size and radial extension of the nucleus wave function
- Wigner function formalism

$$N_A = g_a \cdot \int d^3x_1 \dots d^3x_A \cdot d^3k_1 \dots d^3k_A \cdot f_1(x_1, k_1) \dots f_A(x_A, k_A) \cdot W_A(x_1, \dots, x_A, k_1, \dots, k_A)$$

Wigner density of the bound state → dependence on the wave function

M. Mahlein et al., *Eur. Phys. J. C* 83 (2023) 804



- Different wave functions available

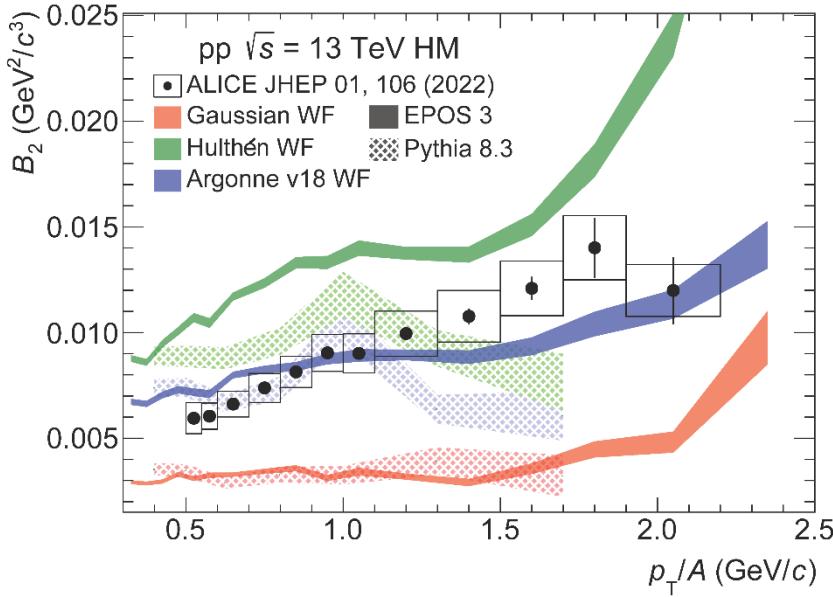
# State of the art coalescence model

- Coalescence parameter depends on both the source size and radial extension of the nucleus wave function
- Wigner function formalism

$$N_A = g_a \cdot \int d^3x_1 \dots d^3x_A \cdot d^3k_1 \dots d^3k_A \cdot f_1(x_1, k_1) \dots f_A(x_A, k_A) \cdot W_A(x_1, \dots, x_A, k_1, \dots, k_A)$$

Wigner density of the bound state → dependence on the wave function

M. Mahlein et al., *Eur. Phys. J. C* 83 (2023) 804

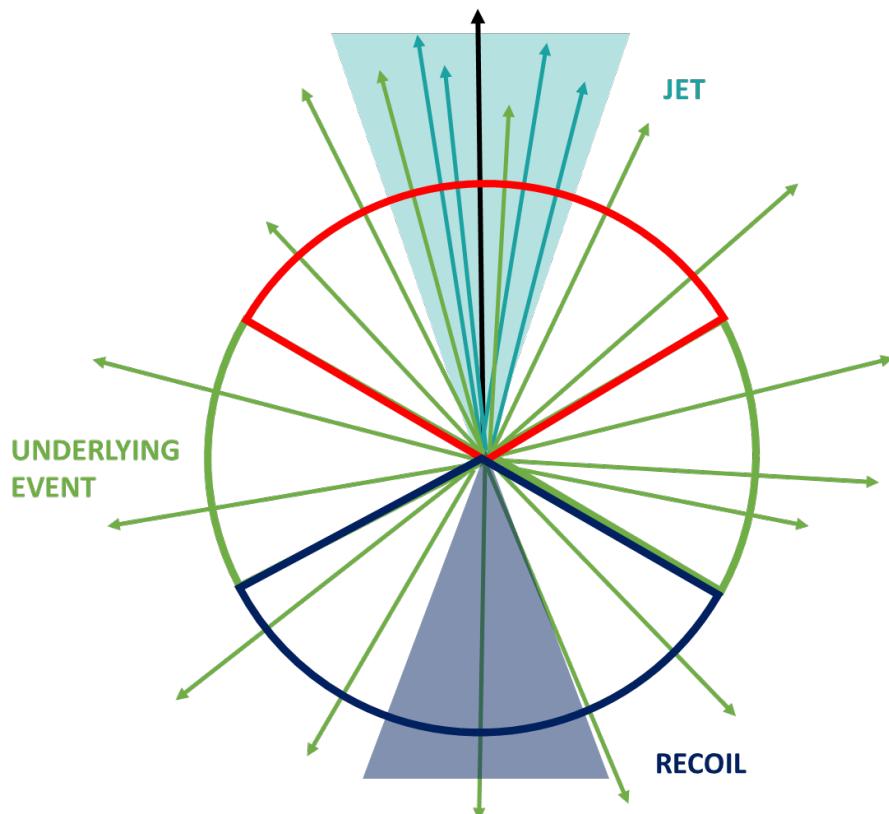


- Different wave functions available
- Best results for deuteron with Argonne  $v_{18}$

# In-jet and underlying event

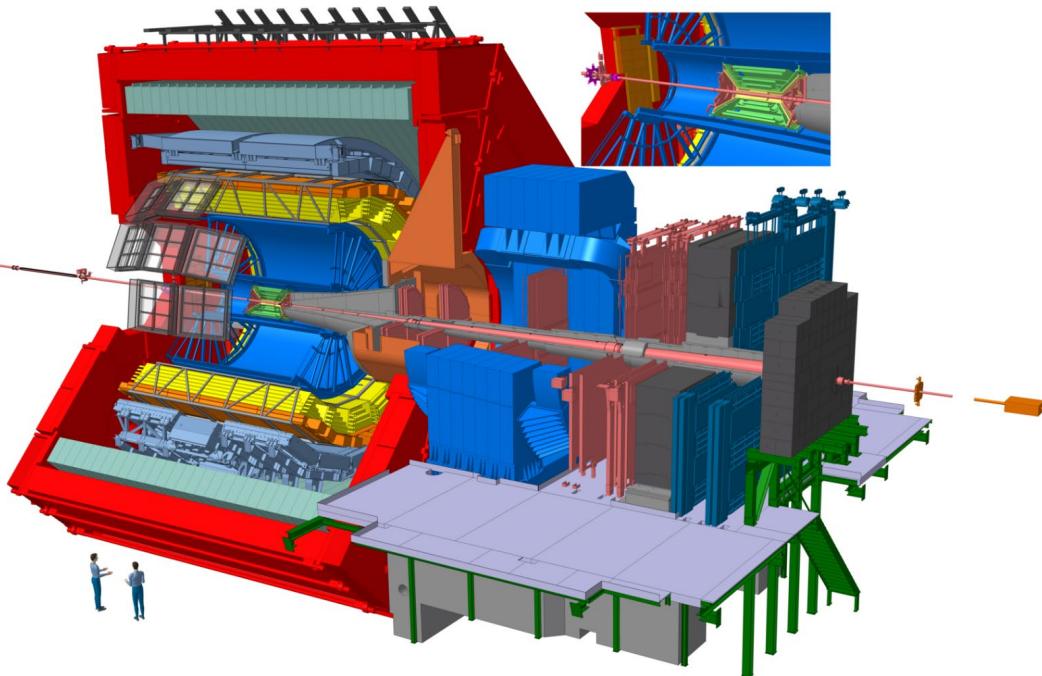
- Constraint of the coalescence model: study production of nuclei in and out of jets
  - nucleons in jets are closer in phase space → larger coalescence probability expected wrt UE
- CDF technique:
  - **Toward** ( $|\Delta\phi| < 60^\circ$ ) : contains JET and UE
  - **Transverse** ( $60^\circ < |\Delta\phi| < 120^\circ$ ) : dominated by the Underlying Event (UE)
  - **Away** ( $|\Delta\phi| > 120^\circ$ ): contains RECOIL JET and UE
- **Jet: Toward – Transverse**

**Leading particle (highest  $p_T$  and  $p_T > 5 \text{ GeV}/c$ )**



T. Martin et al, *Eur. Phys. J. C* **76**, 299 (2016)

# The ALICE detector in Run1&2



- Most suited LHC experiment to study light (anti)nuclei production
- Excellent PID capabilities

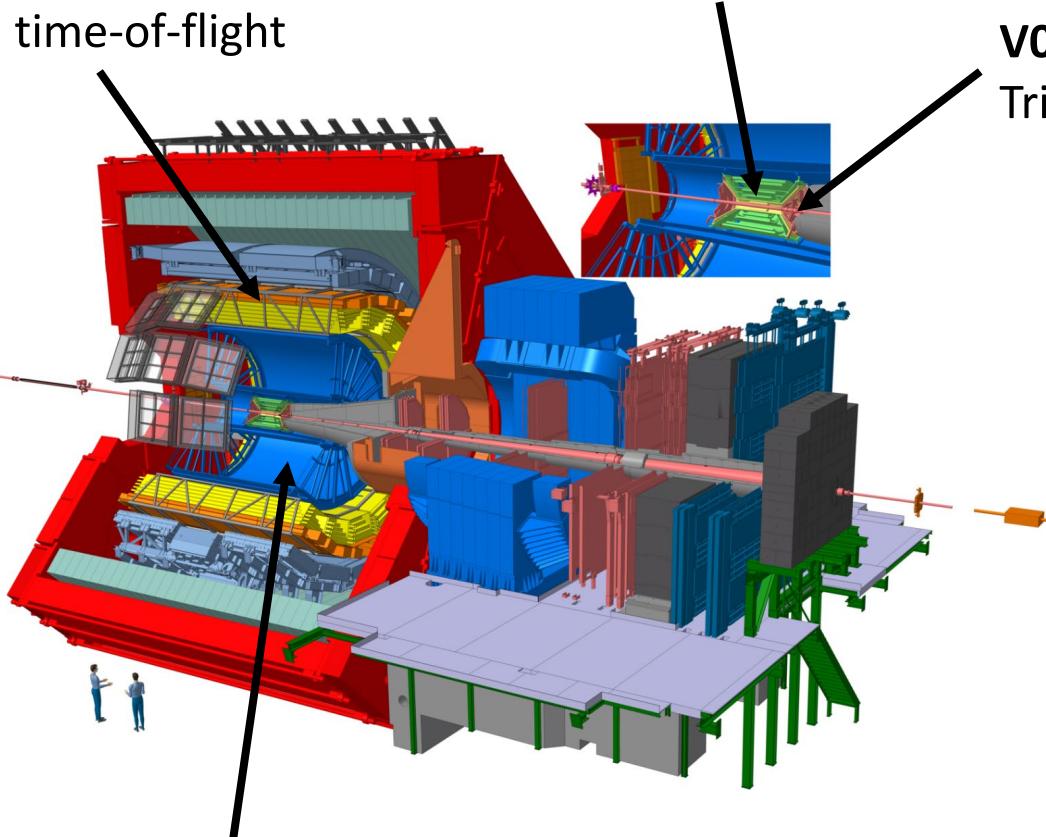
*JINST* 3 (2008) S08002

*Int. J. Mod. Phys. A* 29 (2014) 1430044

# The ALICE detector in Run1&2

## Time Of Flight (TOF)

PID via time-of-flight



## Inner Tracking System (ITS)

Tracking, vertex, PID

V0

Trigger, multiplicity

- Most suited LHC experiment to study light (anti)nuclei production
- Excellent PID capabilities

## Time Projection Chamber (TPC)

Tracking, PID via  $dE/dx$

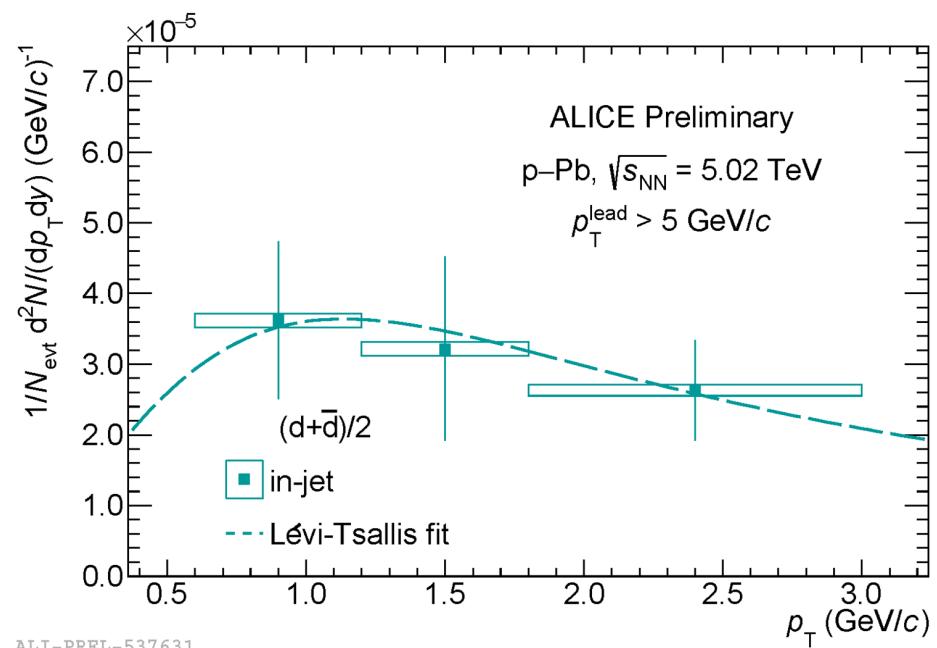
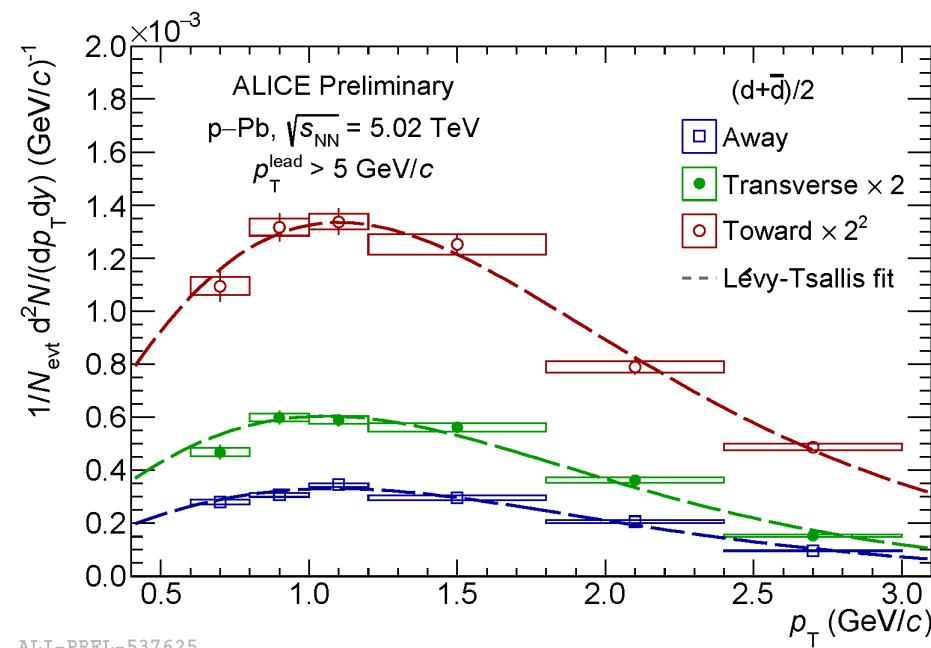
JINST 3 (2008) S08002

Int. J. Mod. Phys. A 29 (2014) 1430044

# In-jet and UE spectra

- What do we need for the coalescence parameter?
- **First ingredient:** (anti)deuteron spectra

$$B_2 = \frac{\frac{1}{2\pi p_T^d} \left( \frac{d^2 N}{dy dp_T} \right)_d}{\left( \frac{1}{2\pi p_T^p} \left( \frac{d^2 N}{dy dp_T} \right)_p \right)^2}$$

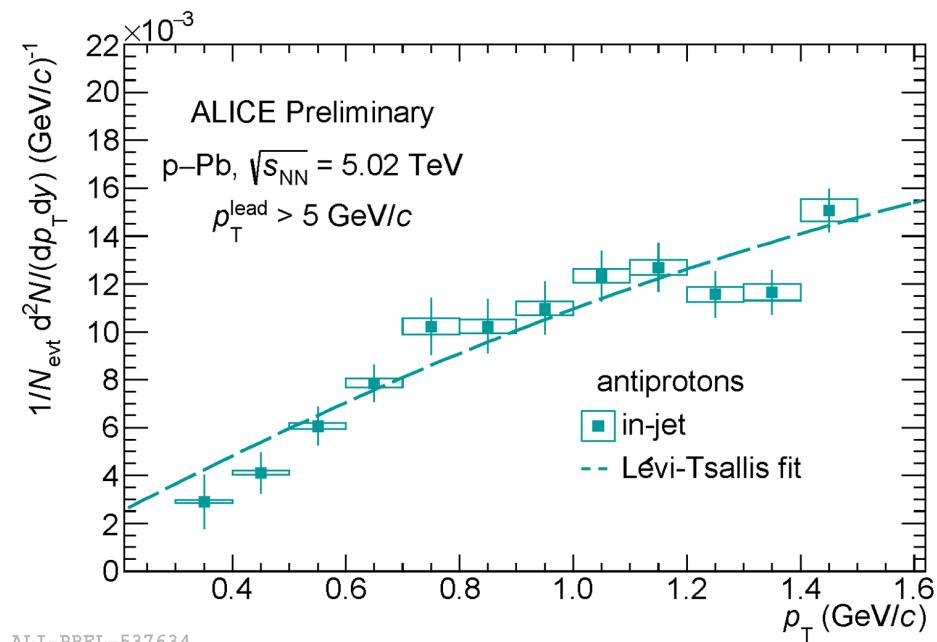
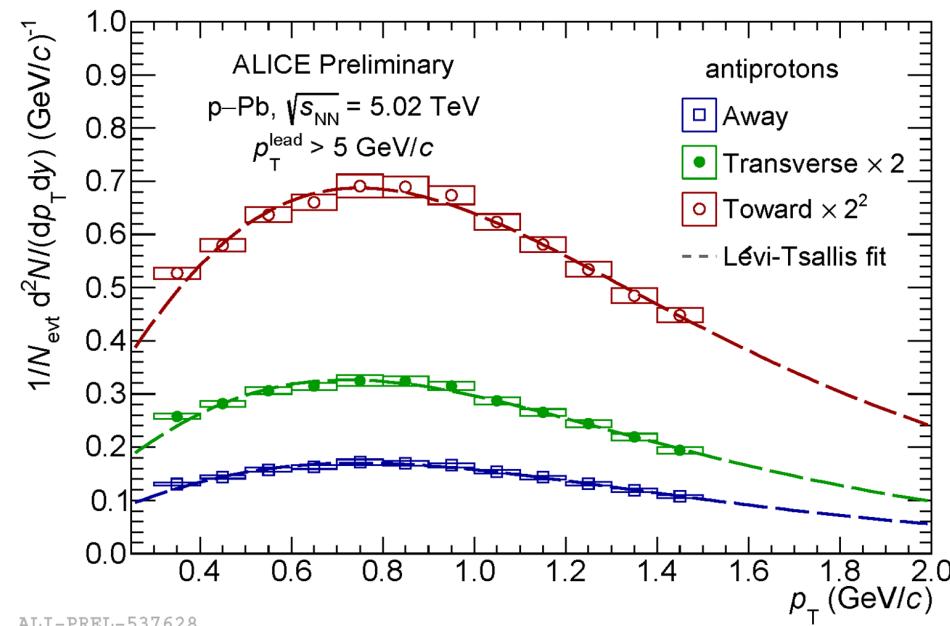


Jet = Toward - Transverse

# In-jet and UE spectra

- What do we need for the coalescence parameter?
- **Second ingredient:** (anti)proton spectra

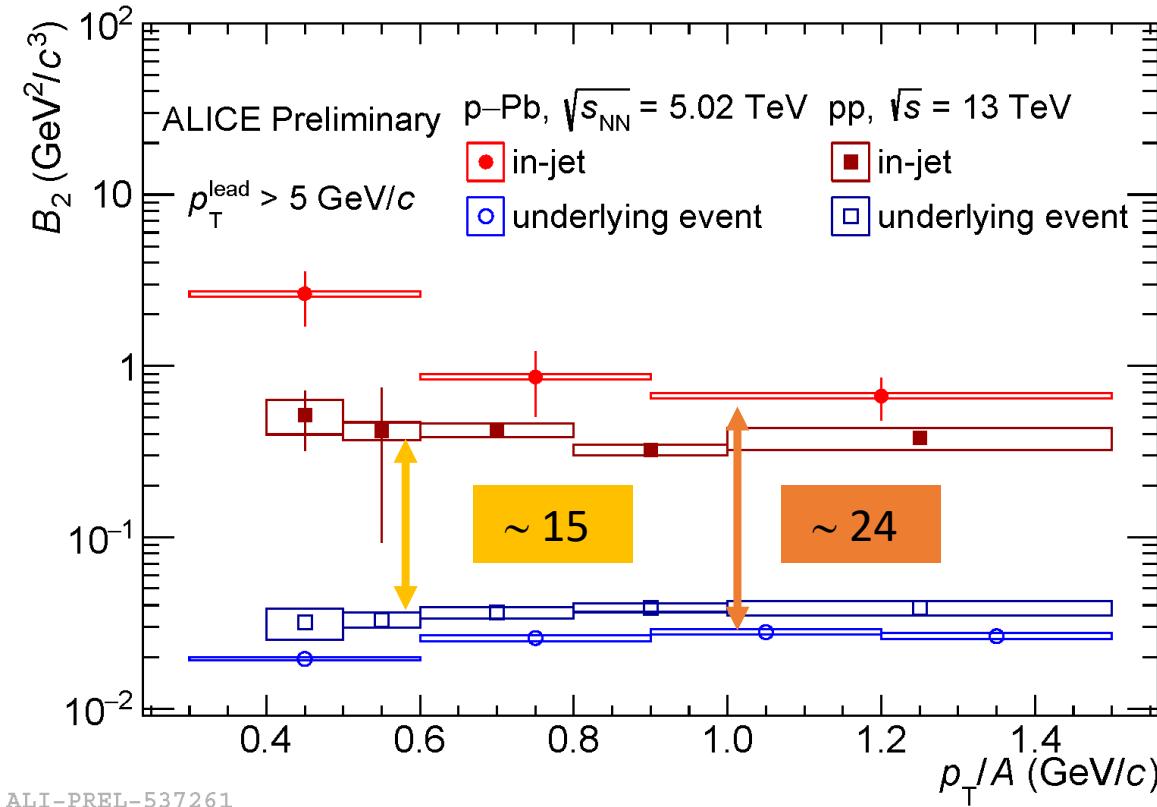
$$B_2 = \frac{\frac{1}{2\pi p_T^d} \left( \frac{d^2N}{dydp_T} \right)_d}{\left( \frac{1}{2\pi p_T^p} \left( \frac{d^2N}{dydp_T} \right)_p \right)^2}$$



Jet = Toward - Transverse

# $B_2$ in jet and UE

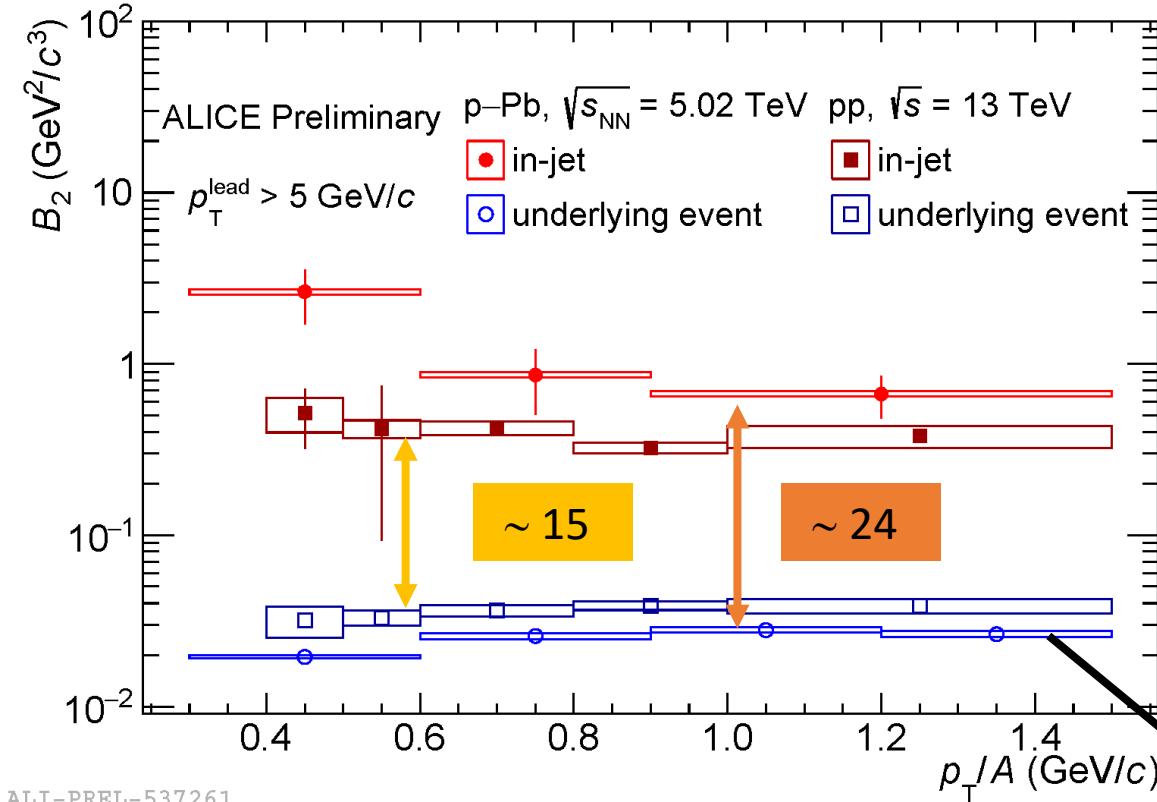
pp: *Phys. Rev. Lett.* 131 (2023) 042301



- **Striking gap** between  $B_2^{\text{jet}}$  and  $B_2^{\text{UE}} \rightarrow$  compatible with the coalescence picture
- **Larger gap** in p-Pb collision wrt pp collisions

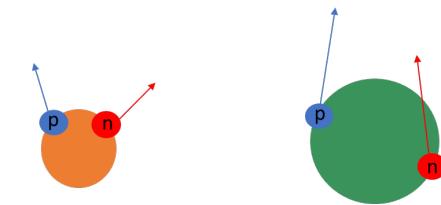
# $B_2$ in jet and UE

pp: *Phys. Rev. Lett.* 131 (2023) 042301



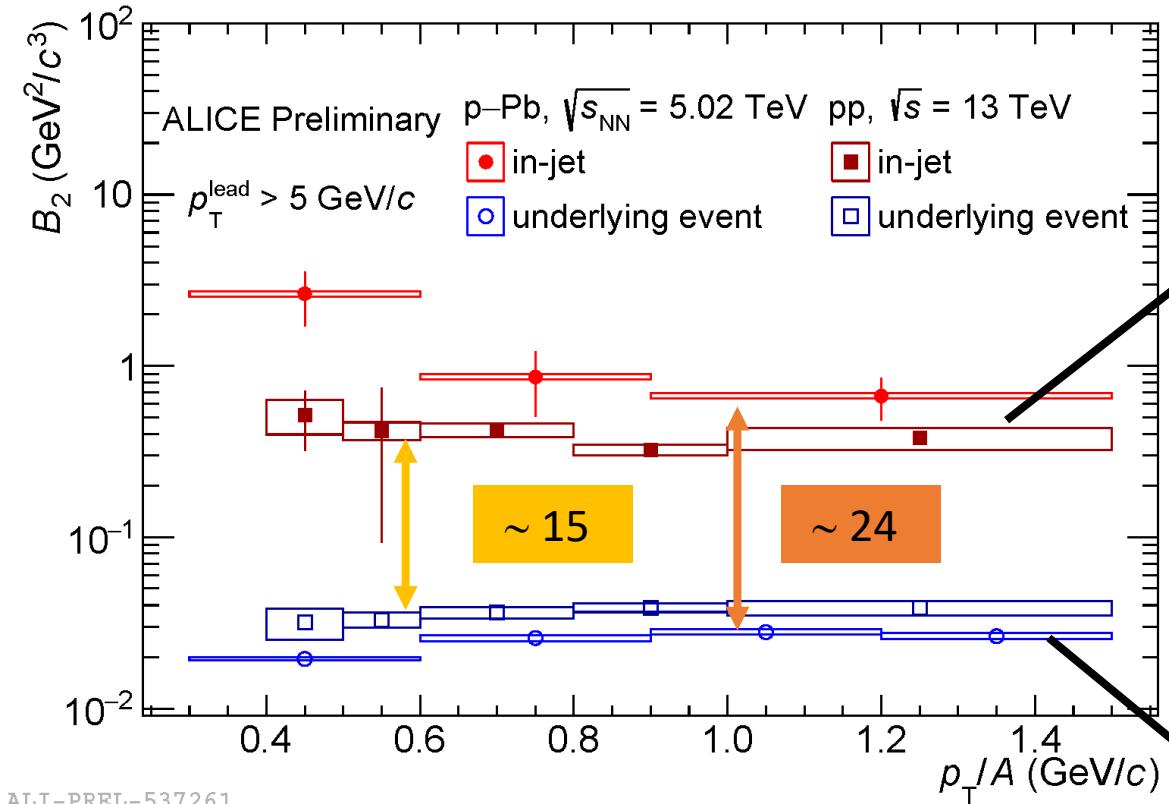
- **Striking gap** between  $B_2^{\text{jet}}$  and  $B_2^{\text{UE}}$  → compatible with the coalescence picture
- **Larger gap in p-Pb collision** wrt pp collisions

$B_2^{\text{UE}}(\text{p-Pb}) < B_2^{\text{UE}}(\text{pp})$  since p-Pb source size is larger than pp source size



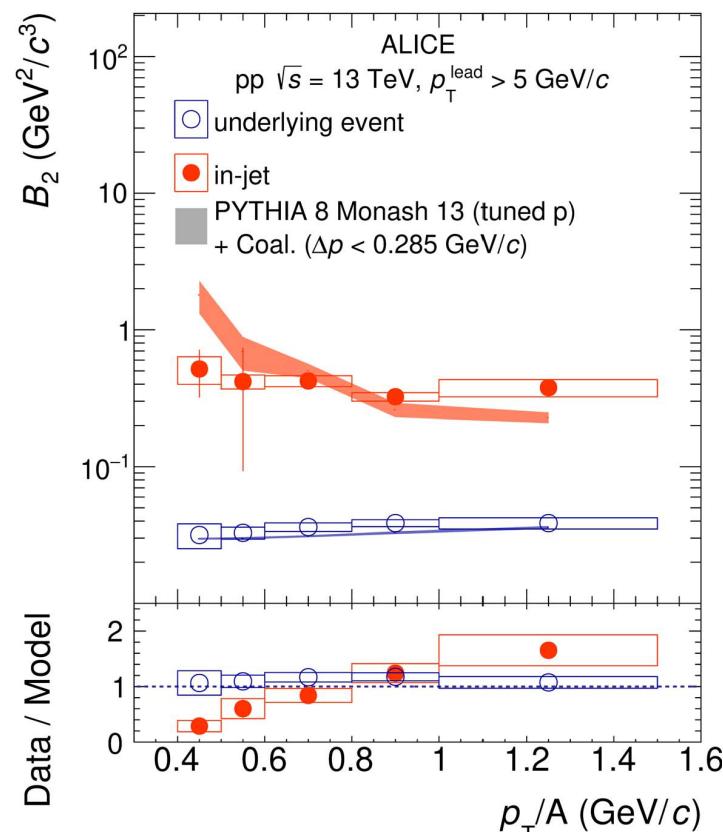
# $B_2$ in jet and UE

pp: *Phys. Rev. Lett.* 131 (2023) 042301



- **Striking gap** between  $B_2^{\text{jet}}$  and  $B_2^{\text{UE}} \rightarrow$  compatible with the coalescence picture
  - **Larger gap** in p-Pb collision wrt pp collisions
    - $B_2^{\text{jet}} (\text{p-Pb}) > B_2^{\text{jet}} (\text{pp})$
    - Different particle composition of jets in pp and p-Pb?
    - Stronger momentum correlations between nucleons?
- $B_2^{\text{UE}} (\text{p-Pb}) < B_2^{\text{UE}} (\text{pp})$  since p-Pb source size is larger than pp source size

# $B_2$ in jet and UE – model comparison

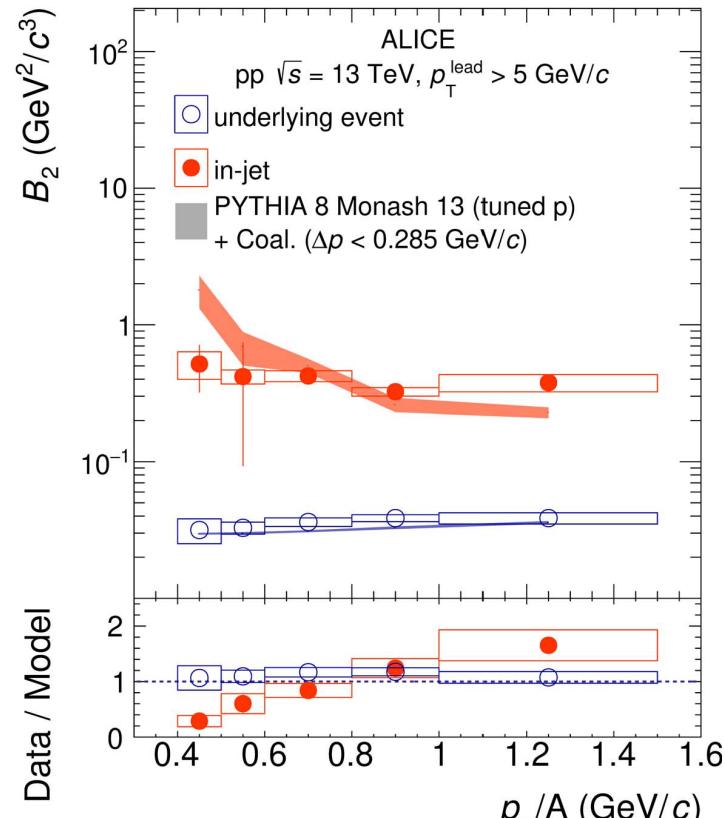


ALI-PUB-533071

PYTHIA 8 Monash 13 +  
simple coalescence

*Phys. Rev. Lett.* 131 (2023) 042301

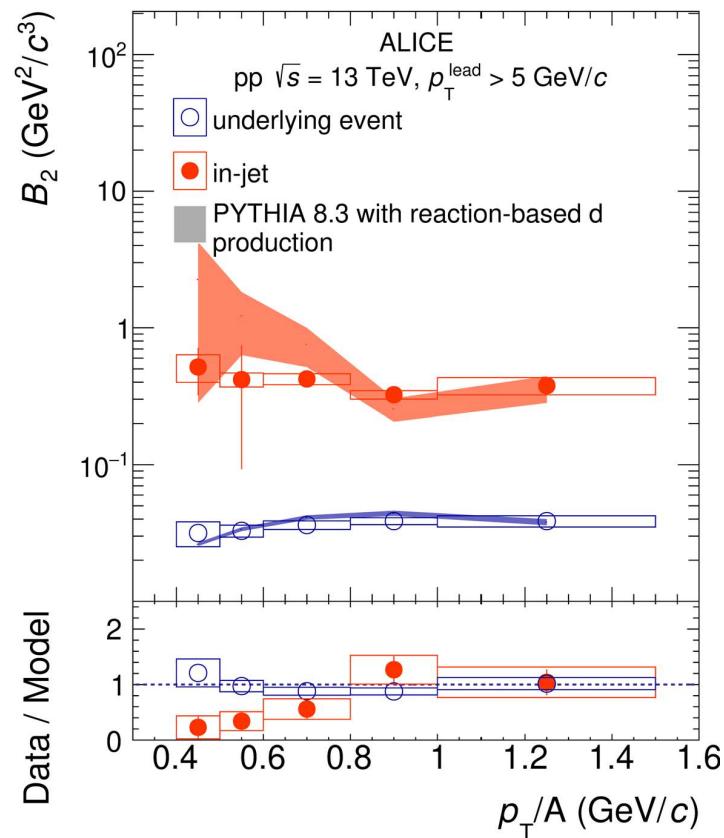
# $B_2$ in jet and UE – model comparison



PYTHIA 8 Monash 13 +  
simple coalescence

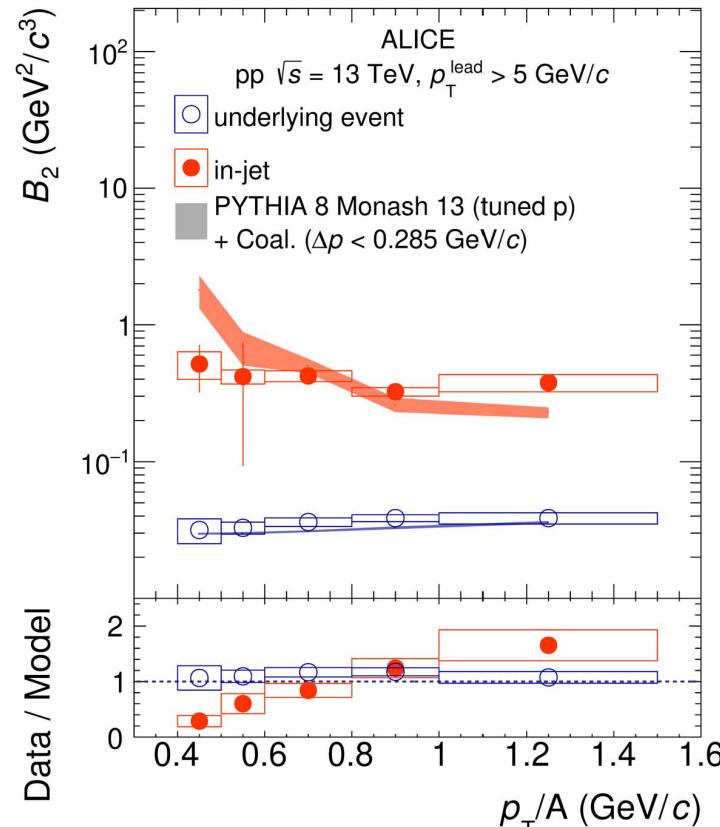
Phys. Rev. Lett. 131 (2023) 042301

PYTHIA 8.3 with reaction-based deuteron production (Bierlich et al., arXiv:2203.11601)



ALI-PUB-533075

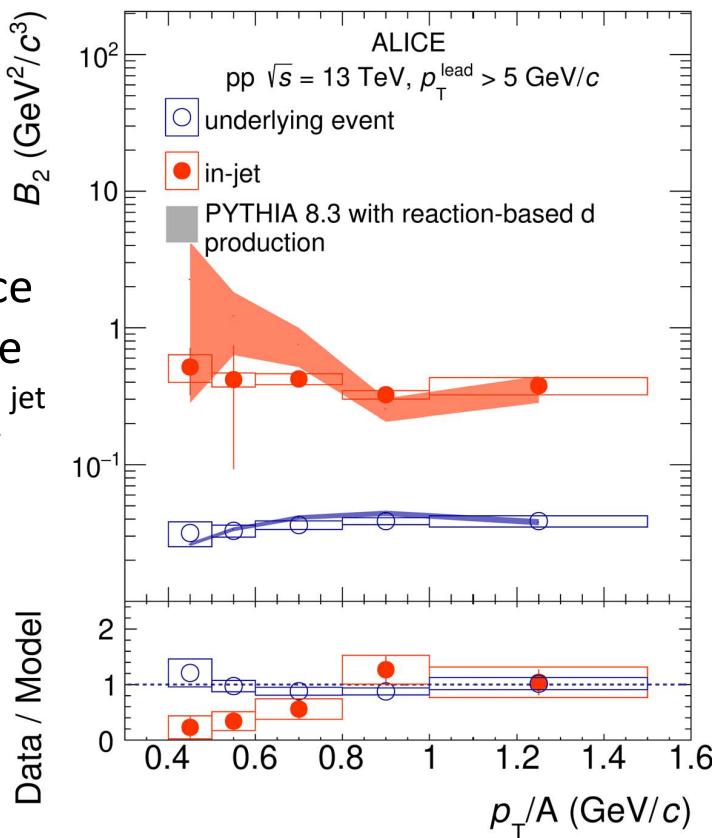
# $B_2$ in jet and UE – model comparison



PYTHIA 8 Monash 13 +  
simple coalescence

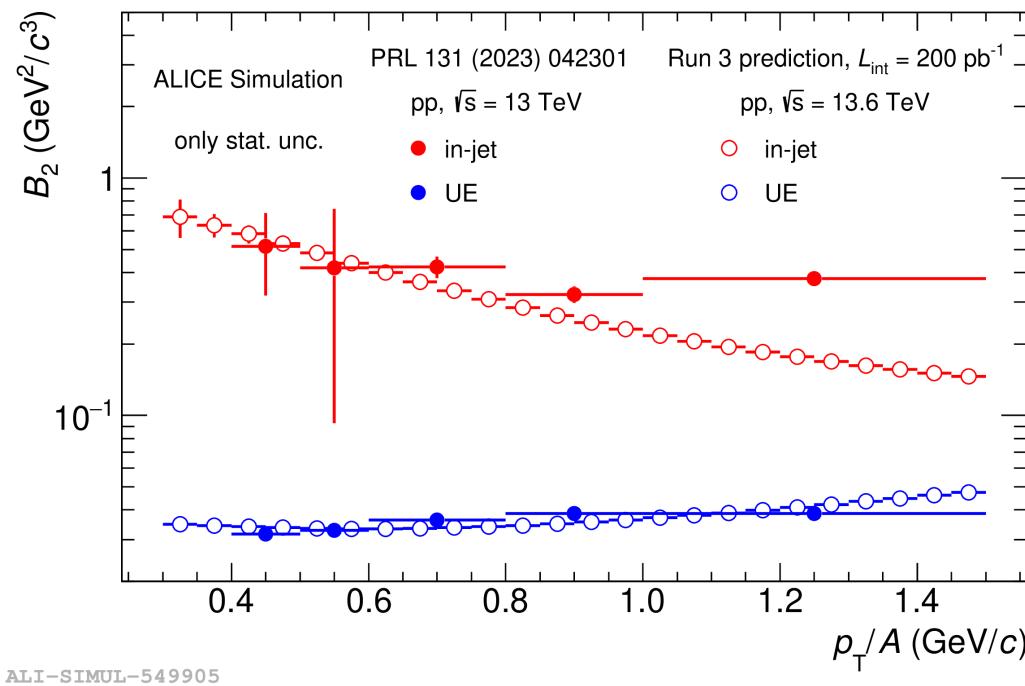
Both models  
qualitatively reproduce  
the data and the large  
difference between  $B_2^{\text{jet}}$   
and  $B_2^{\text{UE}}$

PYTHIA 8.3 with reaction-based deuteron production (Bierlich et al., arXiv:2203.11601)



Phys. Rev. Lett. 131 (2023) 042301

# Prospect for Run 3 measurements



- Simulation studies for Run 3 prospects
- The measured  $p_T$  spectra [1][2] are parametrized and used as inputs for the simulation
- Assumed same efficiency and  $\sigma_{\text{inel}}$ <sup>[3]</sup> of Run 2
- **Improvement of the statistical uncertainties:**
  - factor 4 for  $B_2^{\text{jet}}$
  - factor 3 for  $B_2^{\text{UE}}$
- **Promising results**, multi-differential measurements (e.g. vs multiplicity in the transverse region) could be performed

[1] *Phys. Rev. Lett.* 131 (2023) 042301

[2] *JHEP* 06 (2023) 027

[3] LHCb Collaboration, *JHEP* 06 (2018) 100

# Conclusions

---

- Light (anti)nuclei production have been studied in depth by the ALICE experiment, in order to constraint their production mechanism
- State of the art coalescence model based on femtoscopy measurements and Wigner formalism
- The coalescence model can be tested looking at the light (anti)nuclei production in and out of jets
- **Striking gap between  $B_2^{\text{jet}}$  and  $B_2^{\text{UE}}$**   $\rightarrow$  compatible with the coalescence picture
- Good agreement with model comparison in pp collisions
- **New studies could be performed** thanks to the high integrated luminosity that will be collected at the end of Run 3



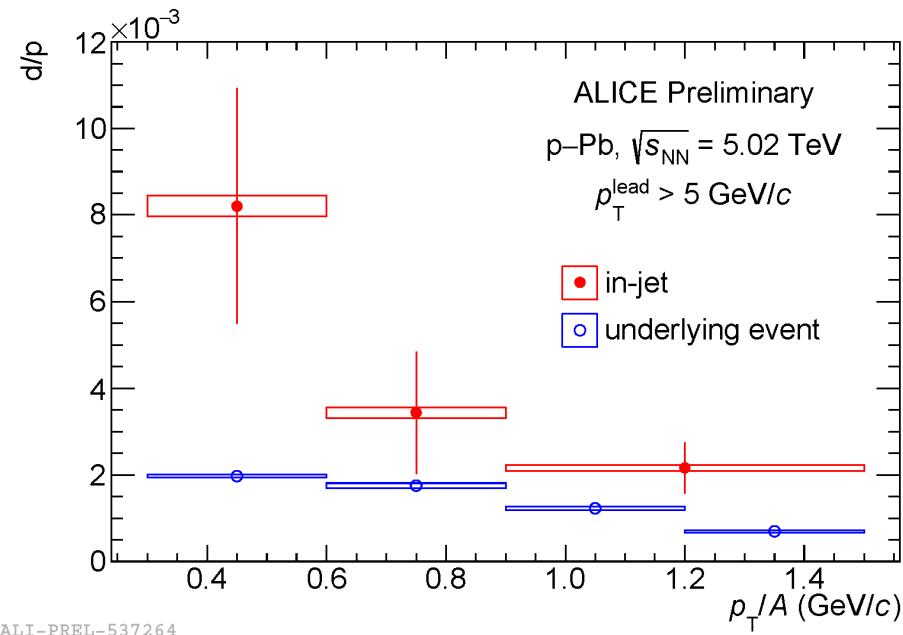
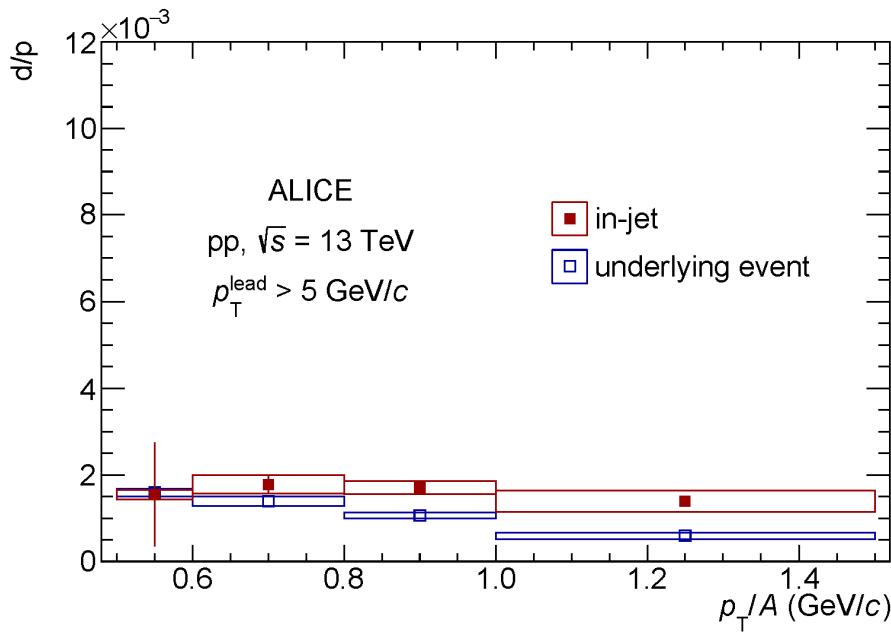
*Thank  
you for  
your kind  
attention!*

# Backup

---

# d/p injet and UE

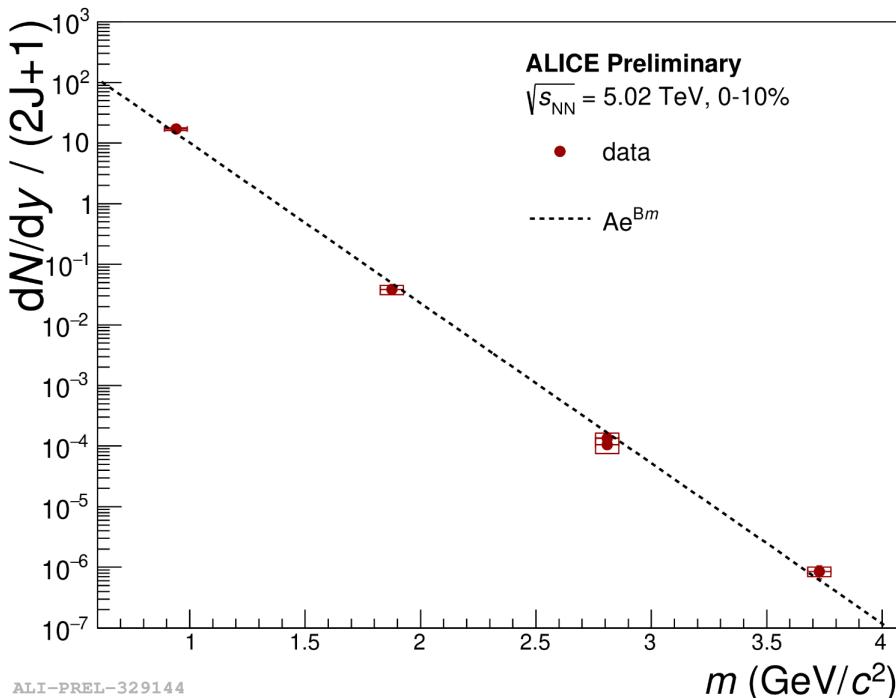
d: *Phys. Rev. Lett.* 131 (2023) 042301 p: *JHEP* 06 (2023) 027



- $d/p^{\text{jet}}$  is higher than  $d/p^{\text{UE}}$
- **Higher  $d/p^{\text{jet}}$  in p-Pb collisions wrt pp collisions**
- Different particle composition → could affect the coalescence probability

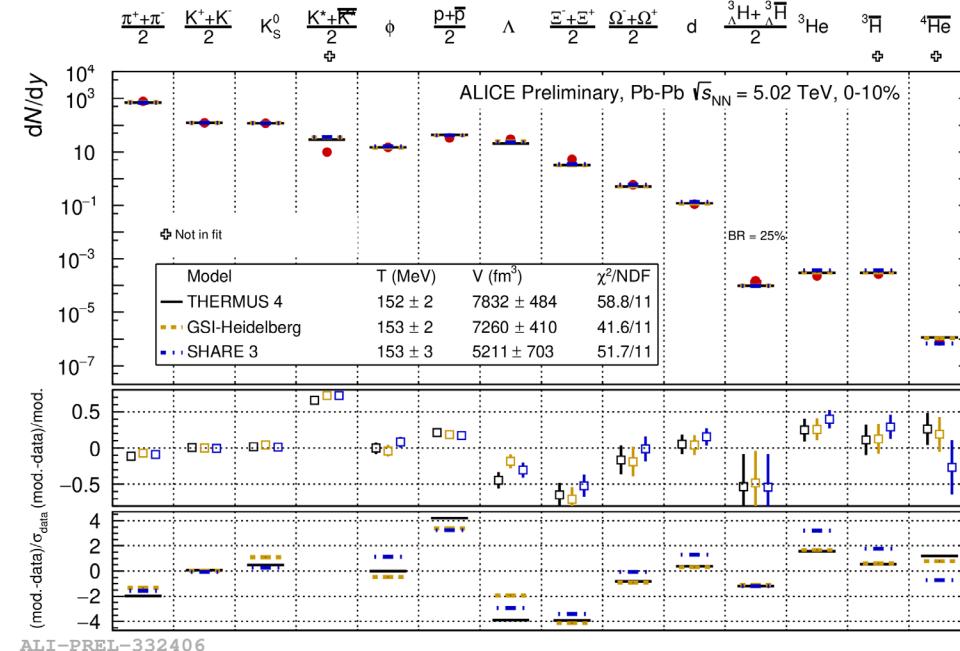
# Statistical Hadronization Model

- Hadrons emitted from a system in statistical and chemical equilibrium
- $T_{\text{chem}}$  is the key parameter
- $dN/dy \propto \exp(-m/T_{\text{chem}}) \rightarrow$  nuclei are sensitive to  $T_{\text{chem}}$  due their large mass
- Particle yield well described with a common  $T_{\text{chem}}$  of  $\sim 156$  MeV



Andronic et al, *Nature* vol. 561 (2018) 321-330

# Statistical Hadronization Model



- Hadrons emitted from a system in statistical and chemical equilibrium
- $T_{\text{chem}}$  is the key parameter
- $dN/dy \propto \exp(-m/T_{\text{chem}})$  → nuclei are sensitive to  $T_{\text{chem}}$  due their large mass
- Particle yield well described with a common  $T_{\text{chem}}$  of  $\sim 156$  MeV
- Comparison between measured and expected yield, evaluated with different SHM implementations
- Nuclei binding energy  $\sim$  few MeV → how can they survive?

Andronic et al, *Nature* vol. 561 (2018) 321-330

THERMUS 4: *Comput.Phys.Commun.* 180 (2009) 84-

106

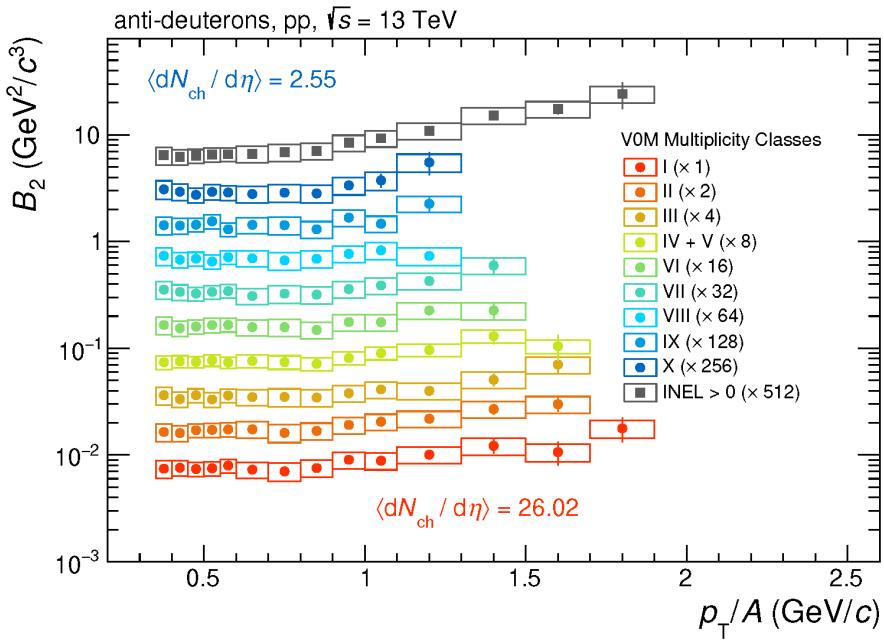
GSI-Heidelberg: *Phys.Lett. B* 673 (2009) 142

SHARE 3: *Comput.Phys.Commun.* 167 (2005) 229-

251

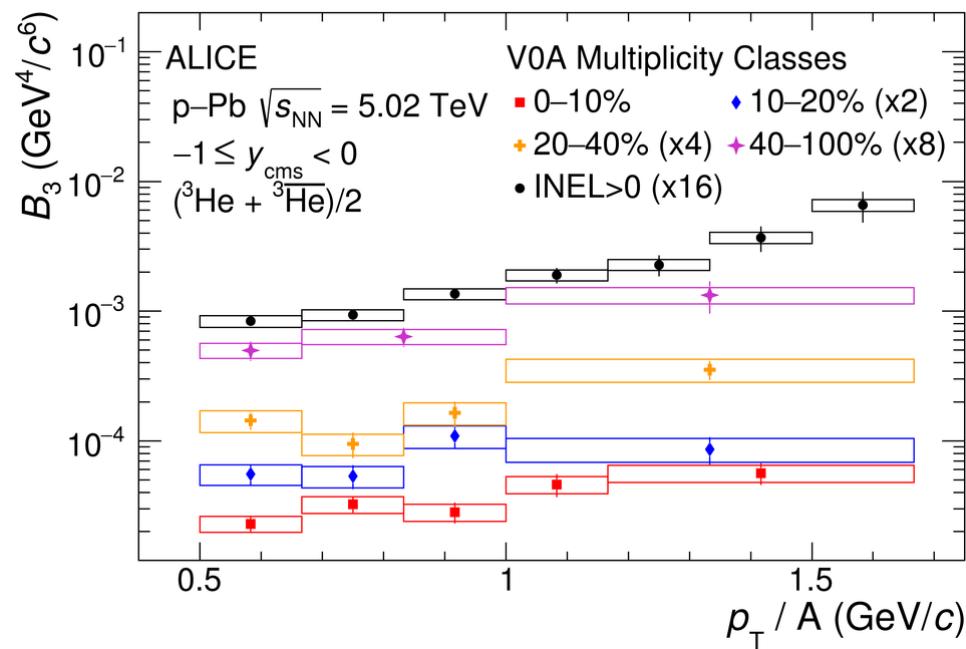
# Coalescence parameter

Eur. Phys. J. C 80 (2020) 889



ALI-PUB-483726

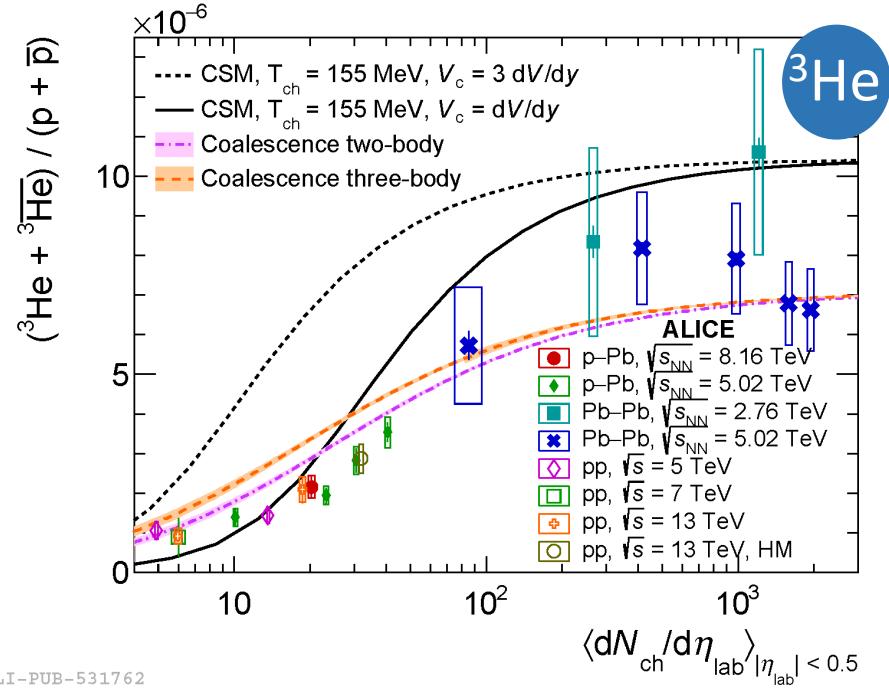
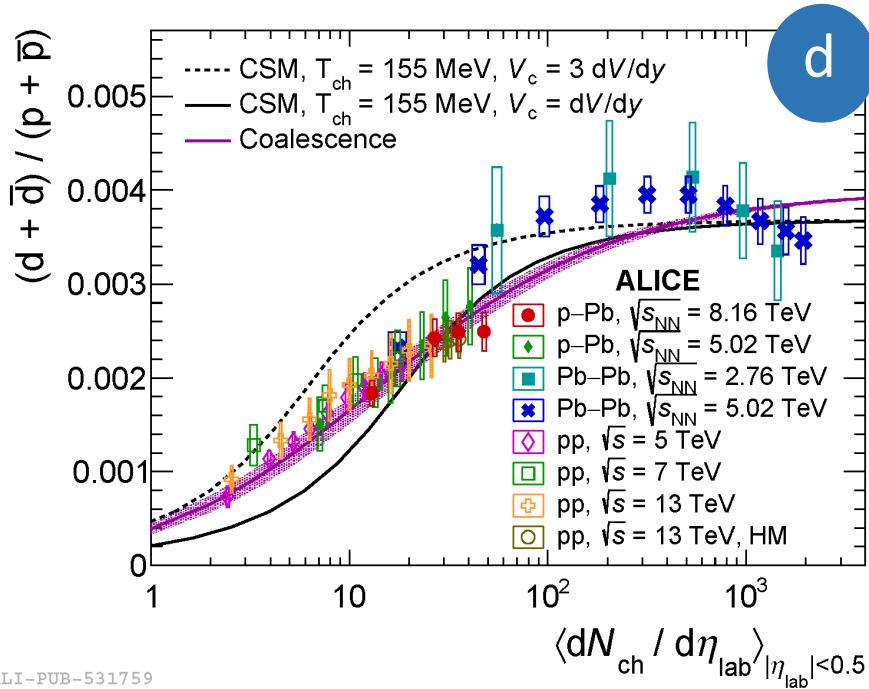
Phys. Rev. C 101 (2020) 044906



- $B_A$  is rather flat in all multiplicity classes, but increase at high  $p_T/A$  in the MB class

# Comparison with models – ratio to p

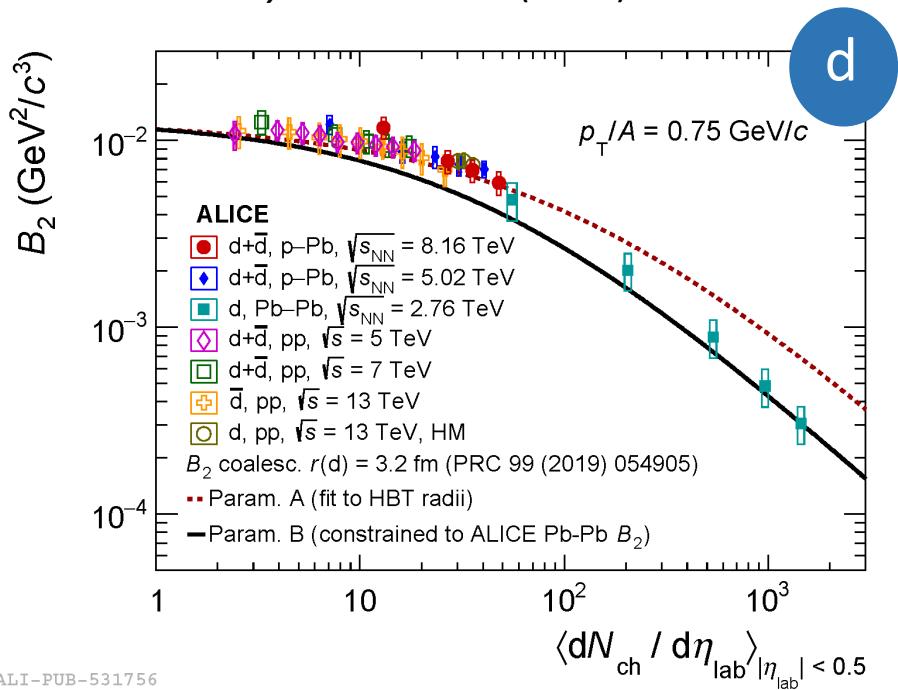
*Phys. Lett. B* 846 (2023) 137795



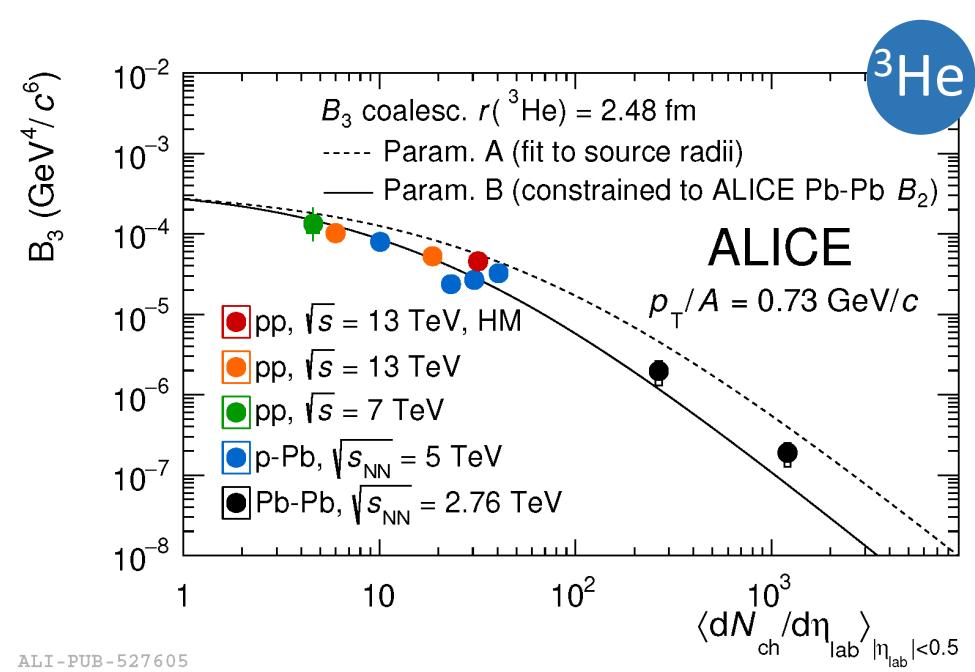
- d/p and He/p ratio evolves smoothly as a function of multiplicity → dependence on the system size
- Observed saturation at multiplicity that corresponds to Pb–Pb collisions
- Ratio compared to predictions from Thermal-FIST CSM and coalescence model
- SHM and coalescence give similar prediction for d, while they diverge for  $^3\text{He}$  → **need new observables!**

# Comparison with models – $B_A$

*Phys. Lett. B* 846 (2023) 137795

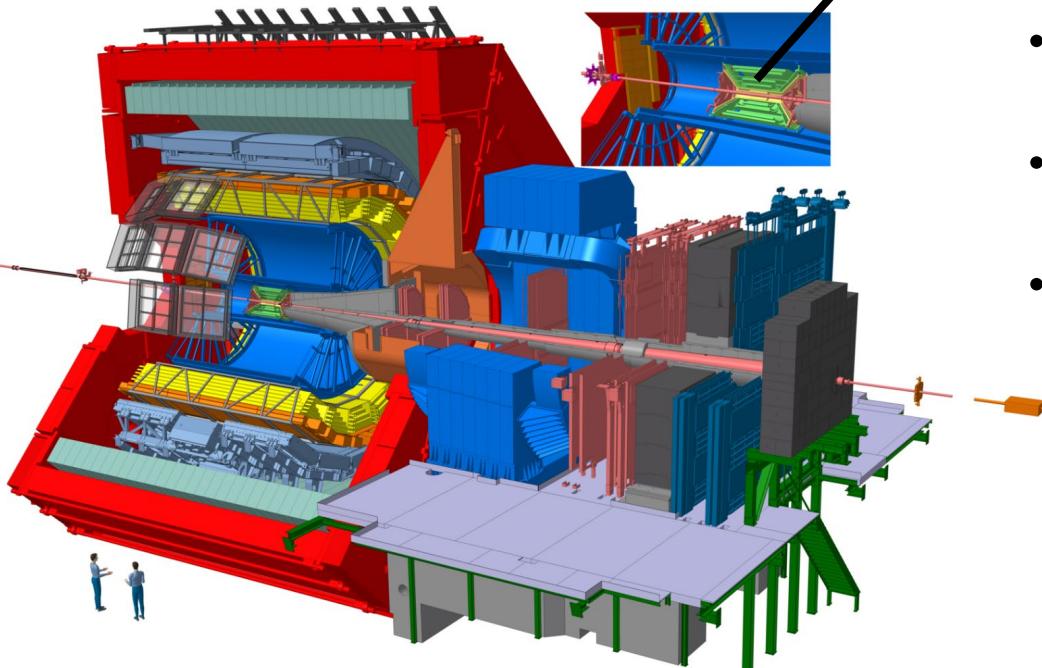


*JHEP* 01 (2022) 106



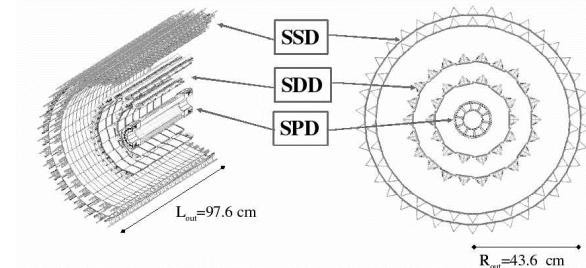
- Coalescence parameter evolves smoothly with multiplicity and decreases with source size
- Different parametrization of source size as a function of  $dN_{\text{ch}}/d\eta$  available
- The parametrizations diverge at high multiplicity →  $B_A$  is a **good observable!**

# The ALICE detector in Run 1&2



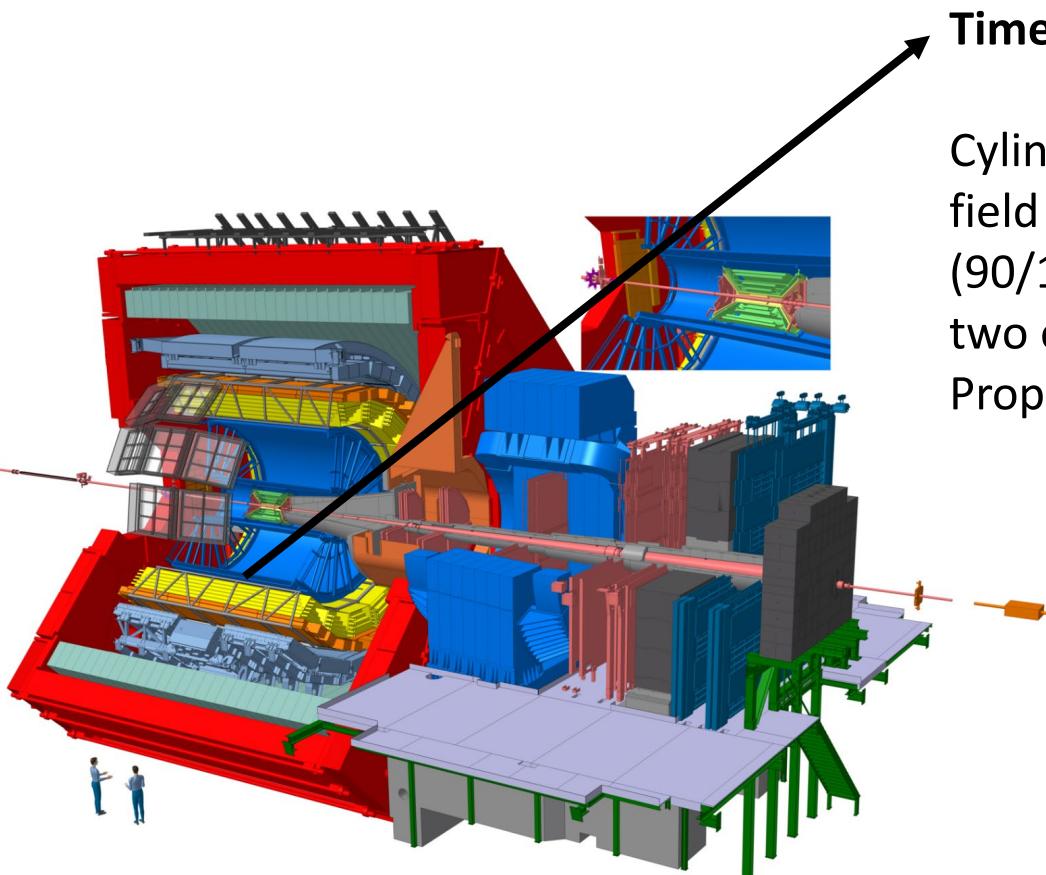
**Inner Tracking System (ITS)**  
Six concentrical layer of silicon sensors:

- 2 layers of Silicon Pixel Detectors (SPD);
- 2 layers of Silicon Drift Detectors (SDD);
- 2 layers of Silicon micro-Strip Detectors (SSD).



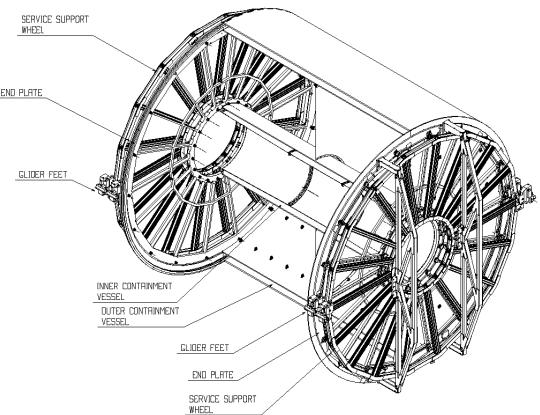
*JINST 3 (2008) S08002*  
*Int. J. Mod. Phys. A 29 (2014) 1430044*

# The ALICE detector in Run 1&2



**Time Projection Chamber (TPC)**

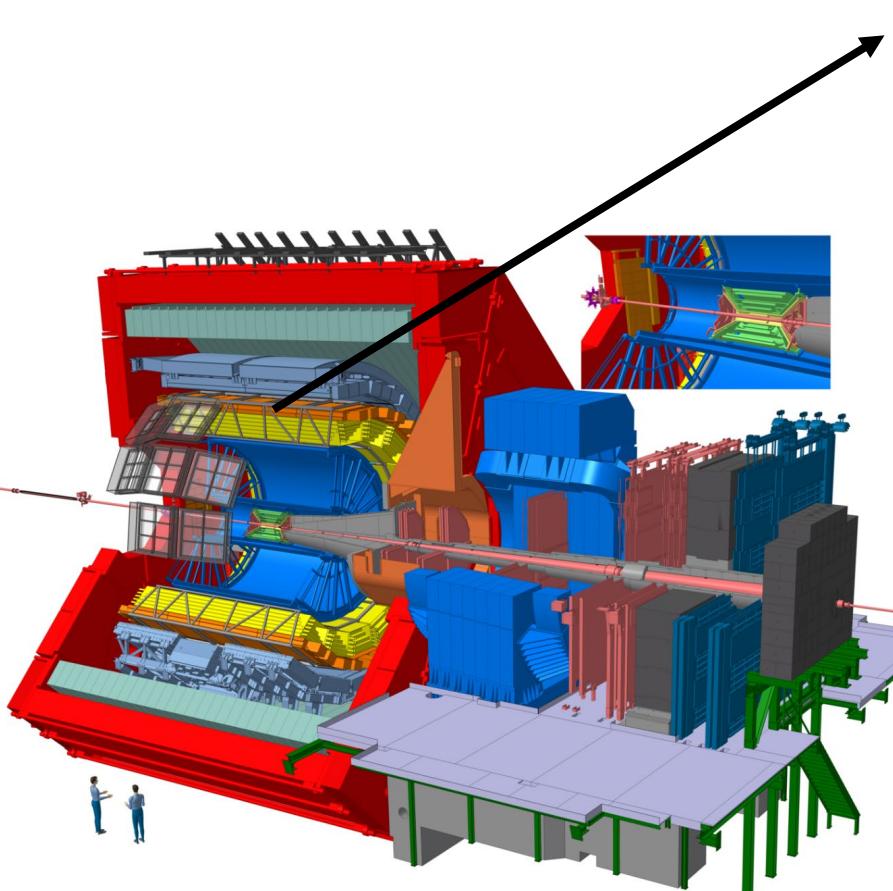
Cylindrical gas detector, made by a field cage filled with Ne/CO<sub>2</sub>/N<sub>2</sub> (90/10/5). The cage is closed with two endcaps made of Multi-Wire Proportional Chambers (MWPC).



JINST 3 (2008) S08002

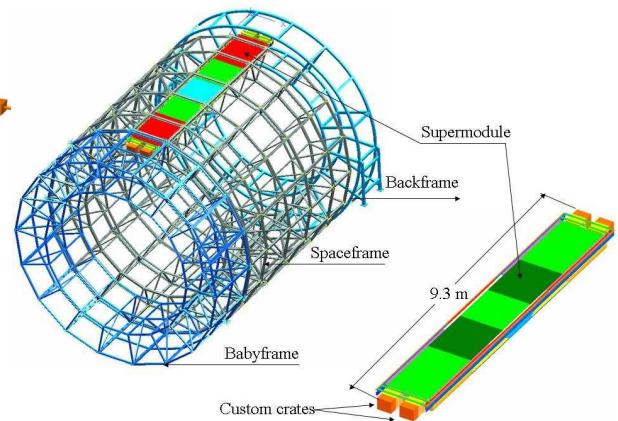
*Int. J. Mod. Phys. A* 29 (2014) 1430044

# The ALICE detector in Run 1&2



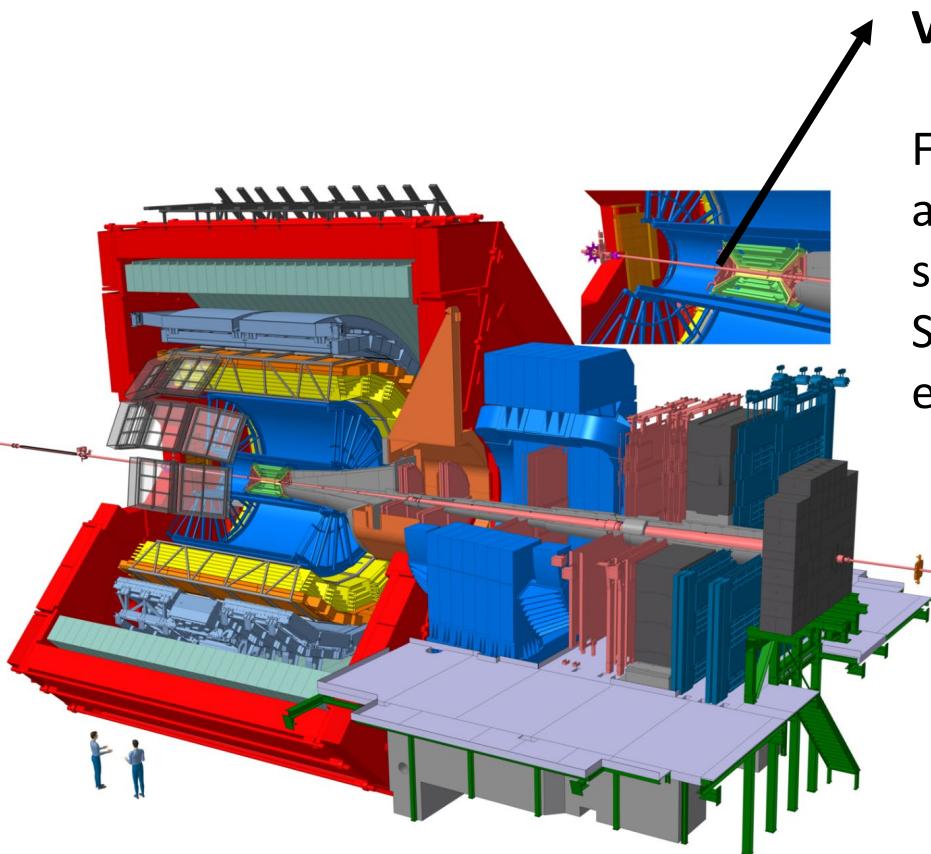
## Time Of Flight (TOF)

90 modules formed by a system of 10 gaps double stack Multigap Resistive Plate Chambers (MRPC). The resistive plates are made with commercially available soda-lime glass sheets with a gap of 250 µm.



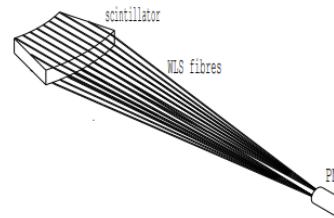
*JINST 3 (2008) S08002  
Int. J. Mod. Phys. A 29 (2014) 1430044*

# The ALICE detector in Run 1&2

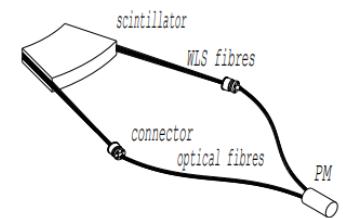


V0

Formed by two different modules, V0A and V0C, consisting of two arrays of scintillator counters and Wave-Length Shifting (WLS) fibres installed on either sides of the interaction point.



V0A



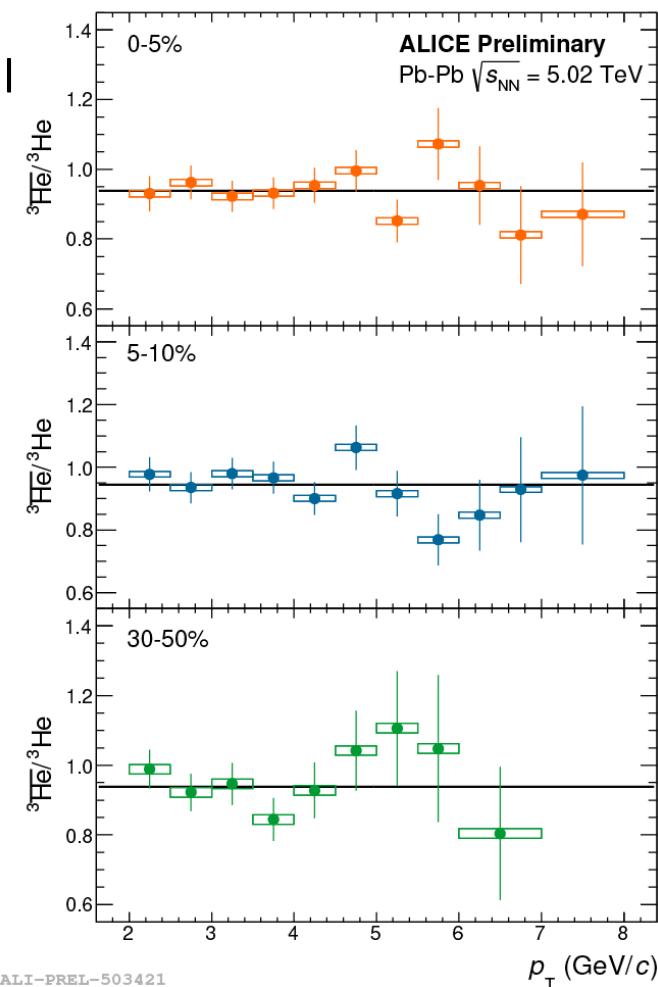
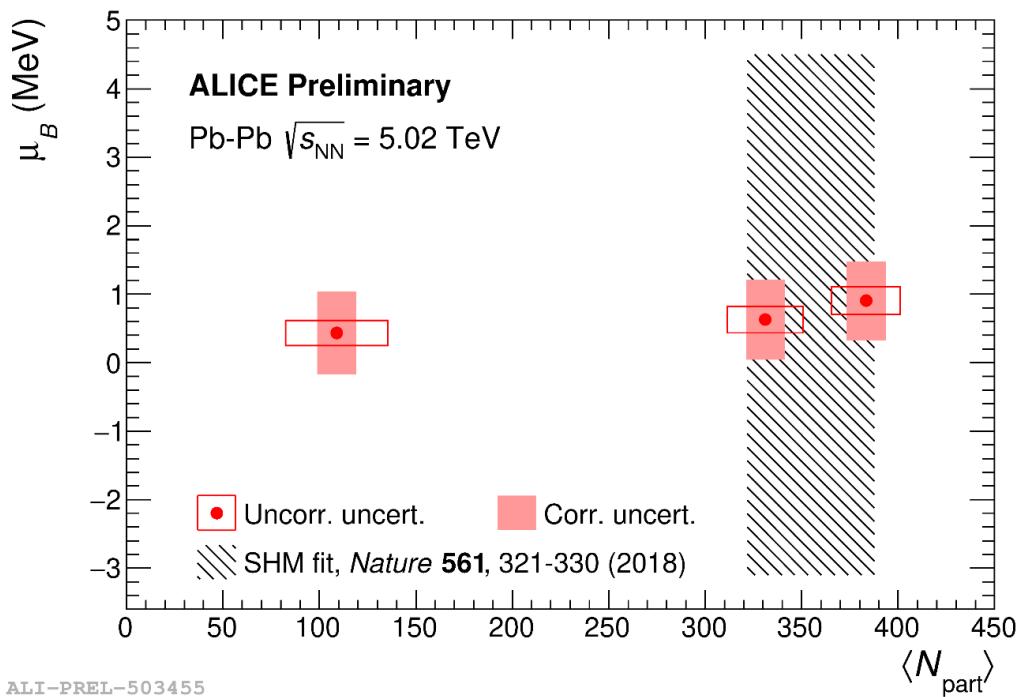
V0C

JINST 3 (2008) S08002

Int. J. Mod. Phys. A 29 (2014) 1430044

# LHC: an (anti)nuclei factory

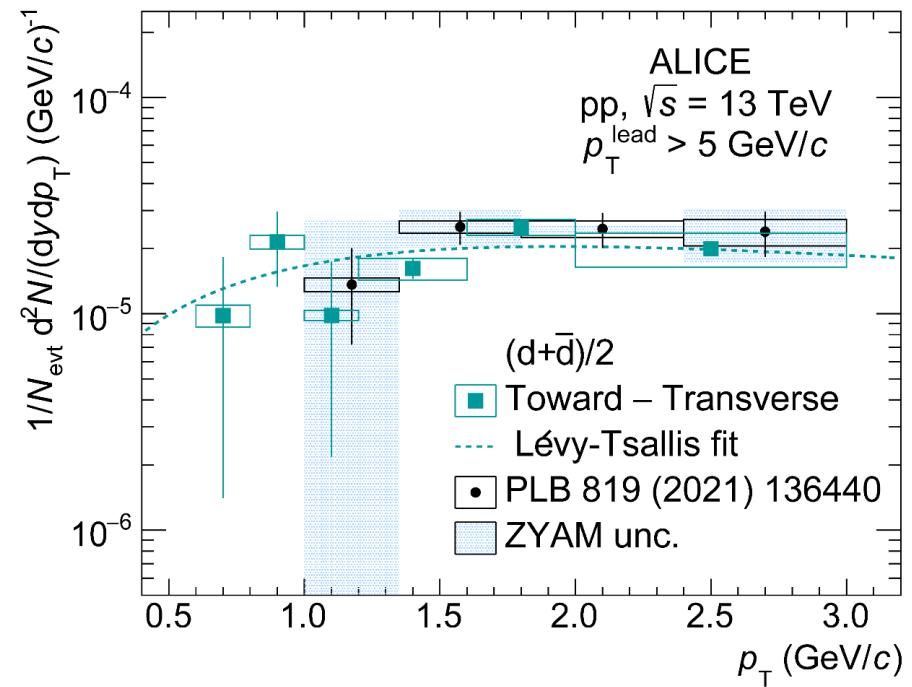
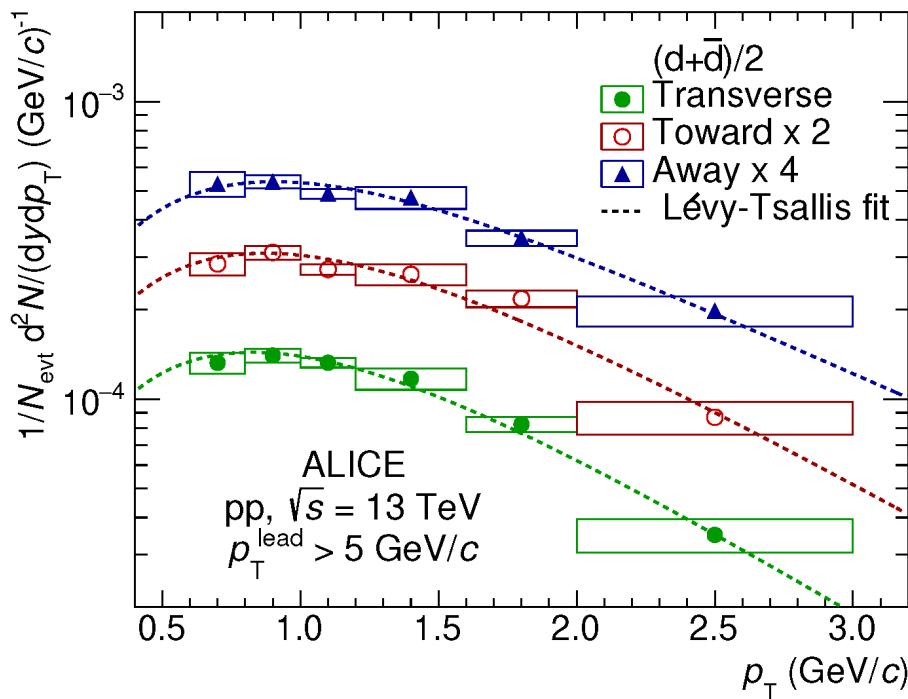
- At the LHC energies the same quantity of matter and antimatter are produced at midrapidity → baryochemical potential  $\mu_B \approx 0$
- Antimatter-to-matter ratio consistent with unity



# In-jet and UE spectra

*Phys. Rev. Lett.* 131 (2023) 042301

pp @ 13 TeV



ALI-PUB-533063

Deuteron production in events with  
 $p_T^{\text{lead}} > 5 \text{ GeV}/c$

Jet = Toward – Transverse

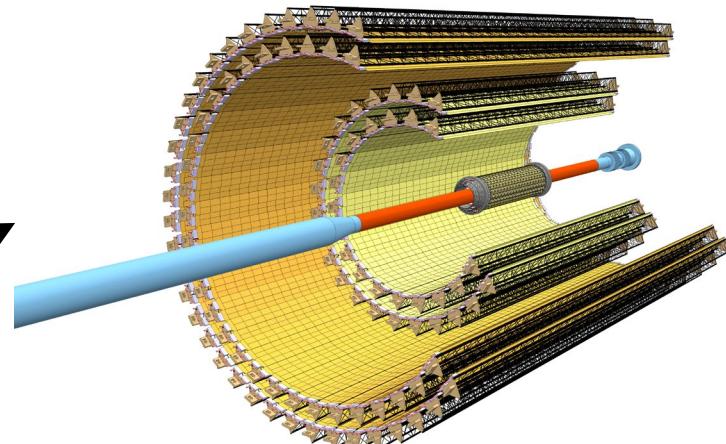
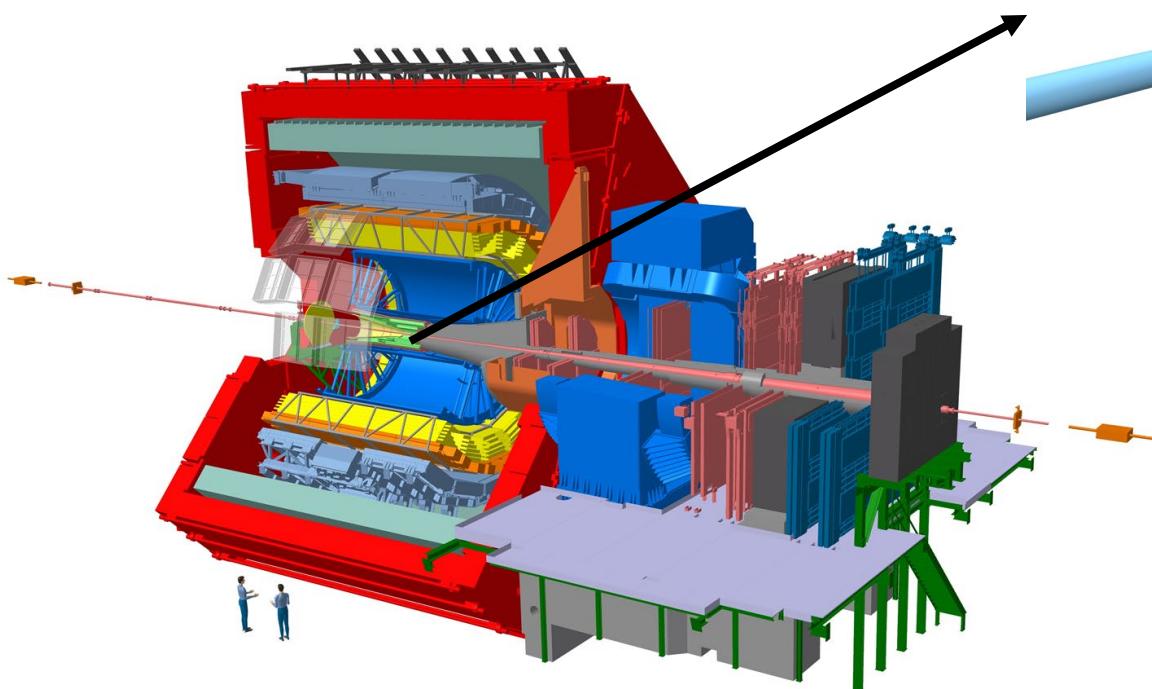
The results are consistent with those obtained  
using the two-particle correlation method

# PYTHIA simulation details

- PYTHIA 8.3:
  - d production via ordinary reactions
  - Energy dependent cross sections parametrized based on data
  - Reactions:

$p + n \rightarrow \gamma + d$	$p + p \rightarrow \pi^+ + d$
$p + n \rightarrow \pi^0 + d$	$p + p \rightarrow \pi^+ + \pi^0 + d$
$p + n \rightarrow \pi^0 + \pi^0 + d$	$n + n \rightarrow \pi^- + d$
$p + n \rightarrow \pi^+ + \pi^- + d$	$n + n \rightarrow \pi^- + \pi^0 + d$
- PYTHIA 8 Monash:
  - Simple coalescence
  - d is formed if  $\Delta p < p_0$ , with  $p_0 = 285 \text{ MeV}/c$

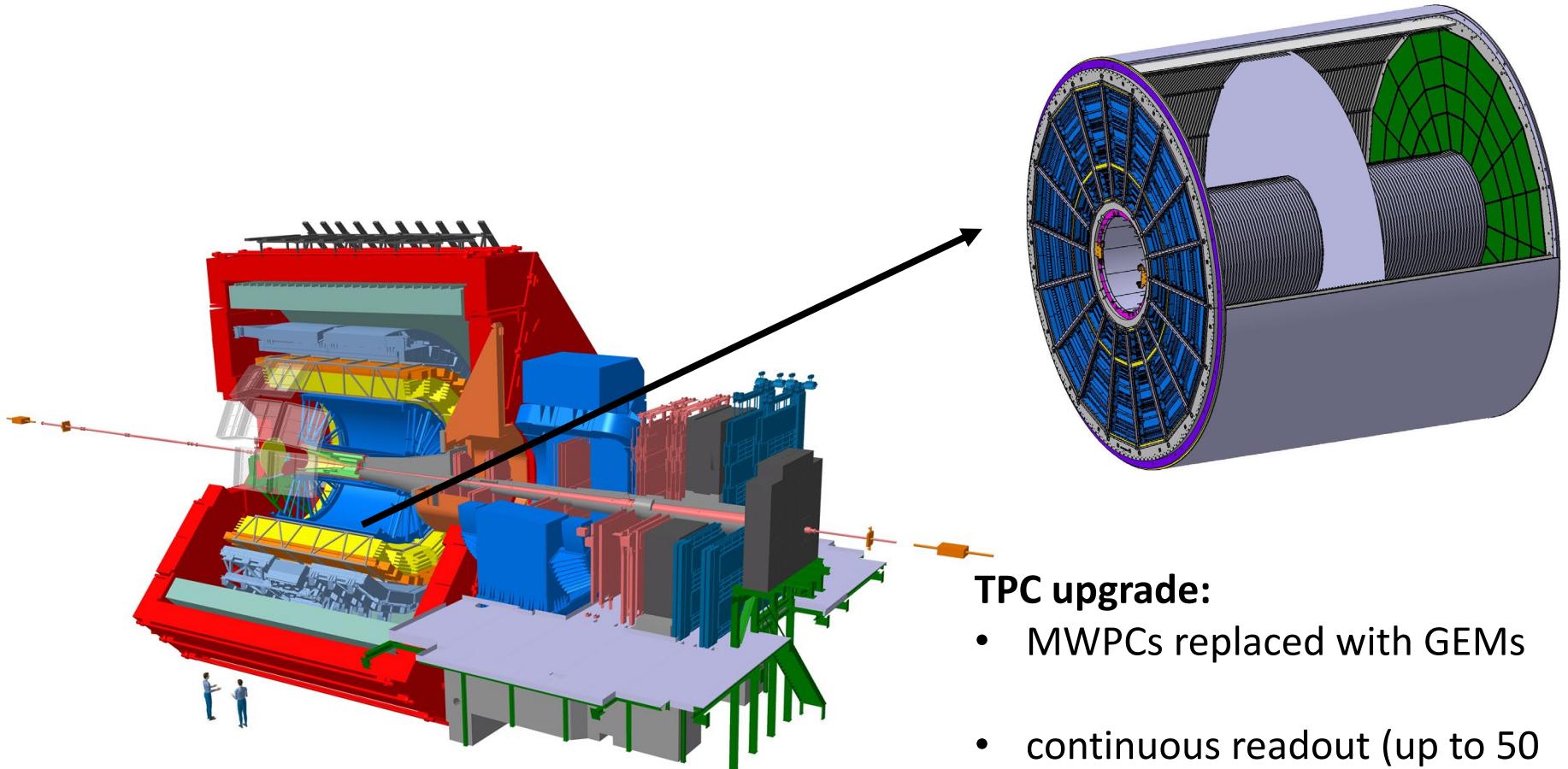
# The new ALICE detector



## ITS upgrade:

- 7 layers of silicon pixel detectors with low material budget ( $0.35\% X_0$  per layer for inner barrel)
- 1st layer closer to interaction point (22 mm) → improved pointing resolution ( $\times 3$ )

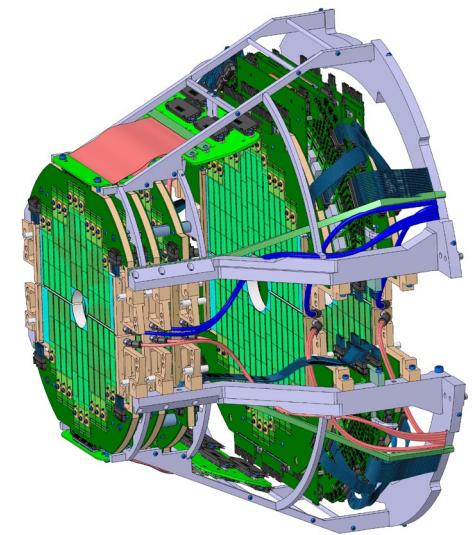
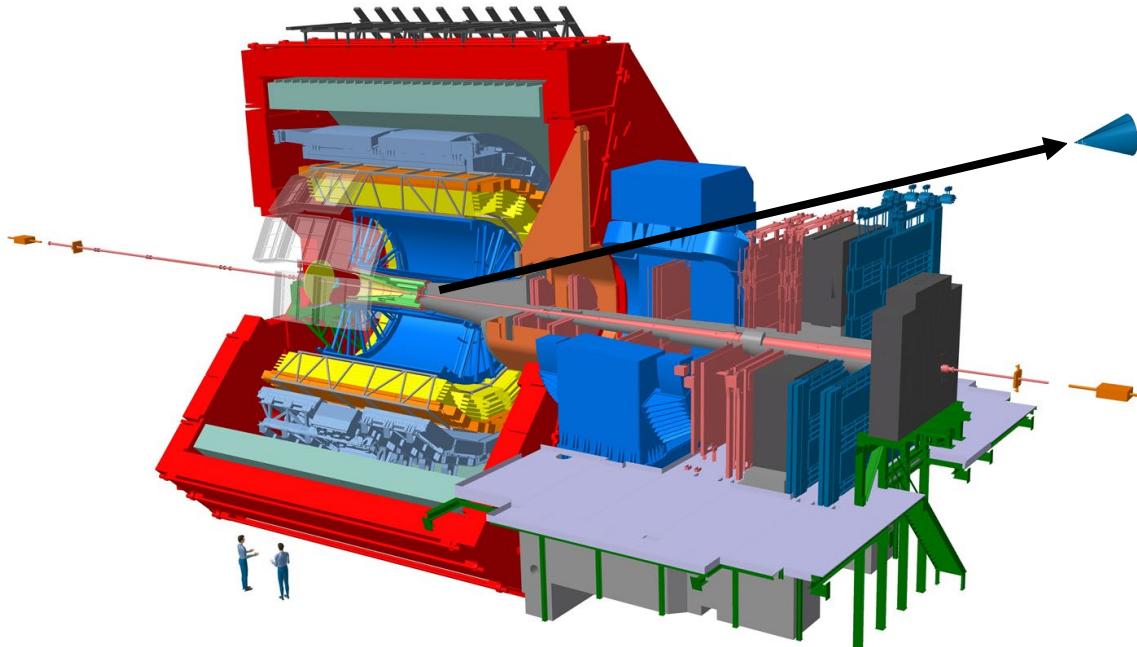
# The new ALICE detector



## TPC upgrade:

- MWPCs replaced with GEMs
- continuous readout (up to 50 kHz in Pb-Pb collisions) → about x50 more events wrt Run 2 for Pb-Pb collisions

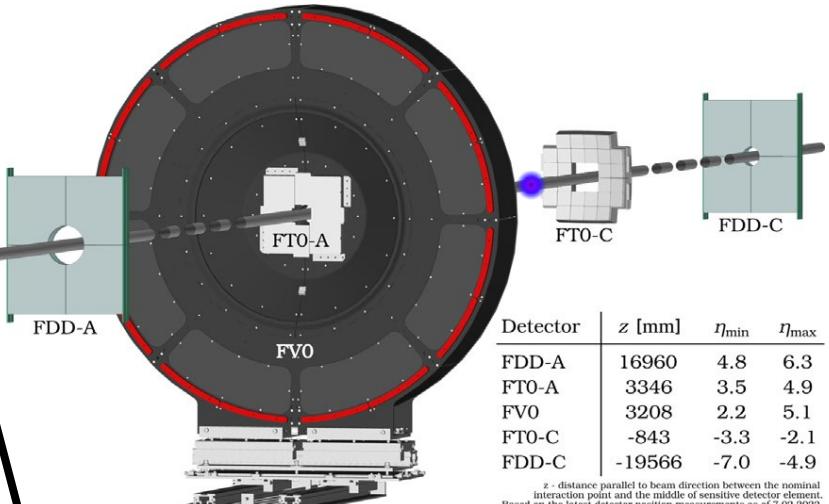
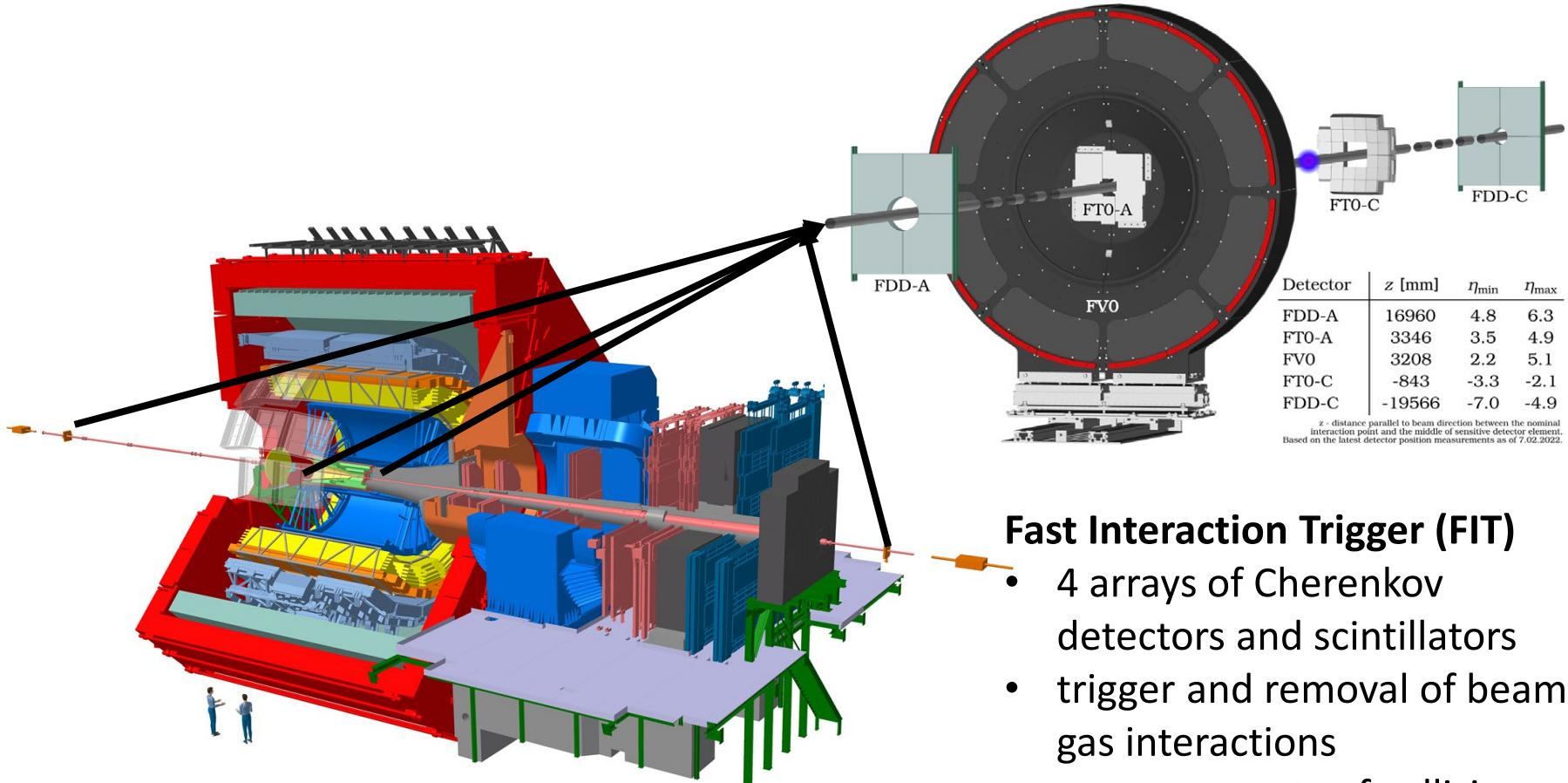
# The new ALICE detector



## MFT (Muon Forward Tracker)

- High-resolution silicon tracker installed before the forward absorber
- Improve muon pointing and separation of prompt and non-prompt muons

# The new ALICE detector



## Fast Interaction Trigger (FIT)

- 4 arrays of Cherenkov detectors and scintillators
- trigger and removal of beam-gas interactions
- measurements of collision time and geometry (centrality and event plane)