



**Faculty
of Physics**

WARSAW UNIVERSITY OF TECHNOLOGY



Summary of recent developments in di-hadron correlations of identified particles (experiment & theory)

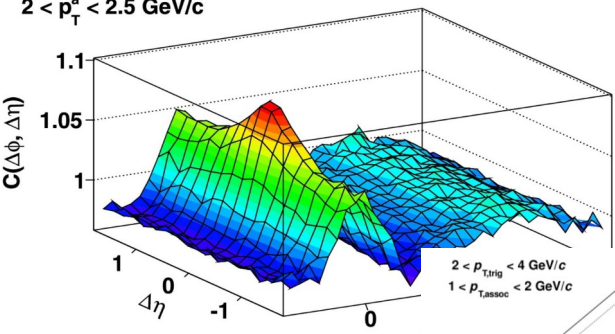


Łukasz Graczykowski
in collaboration with
Małgorzata Janik



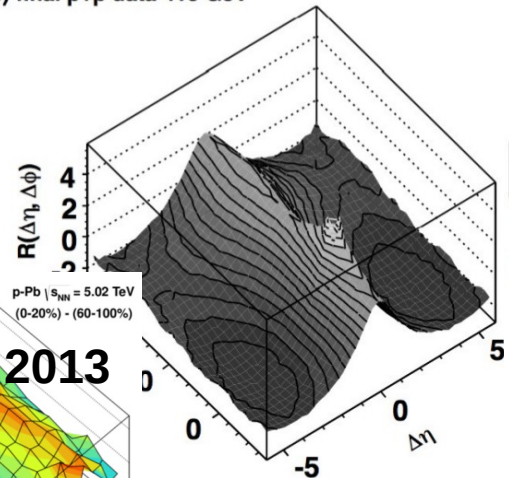
Catania (Italy), November 6-10, 2023

$3 < p_T^1 < 4 \text{ GeV}/c$
 $2 < p_T^a < 2.5 \text{ GeV}/c$



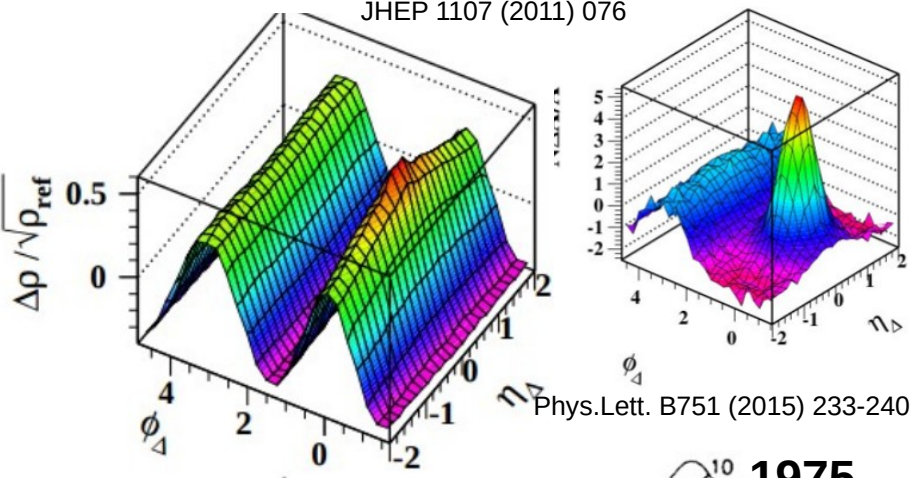
Pb-Pb 2.76 TeV
 0-10%

b) final p+p data 410 GeV



2007

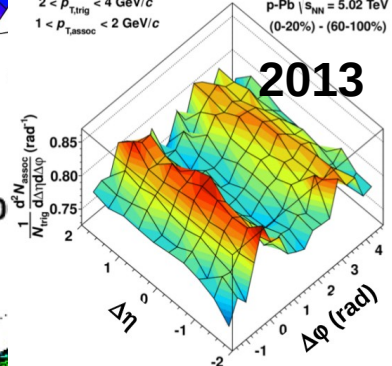
JHEP 1107 (2011) 076



Phys.Lett. B751 (2015) 233-240

CERN-PH-EP-2015-308
 Phys. Lett. B746 (2015) 1

(b) MinBias, $1.0 \text{ GeV}/c < p_T < 3.0$

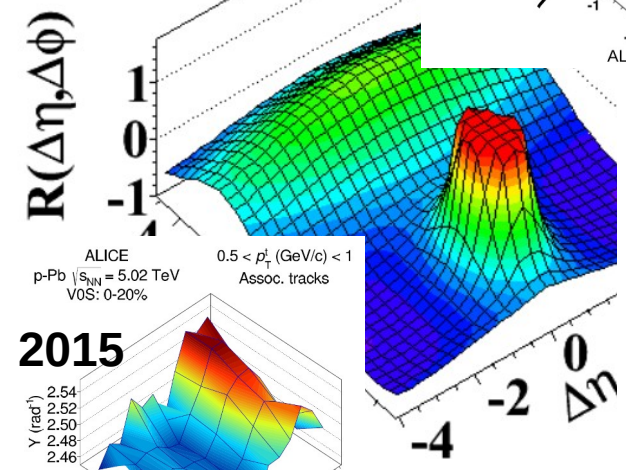


2013

$1 \text{ GeV}/c < p_T < 3.0 \text{ GeV}/c$

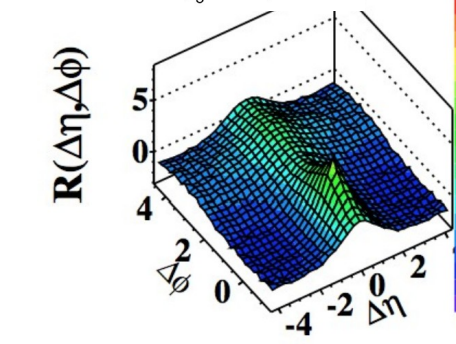
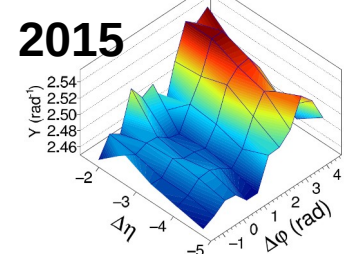
JHEP 1205 (2012) 157

ALICE, PLB719 (2013) 29



2015

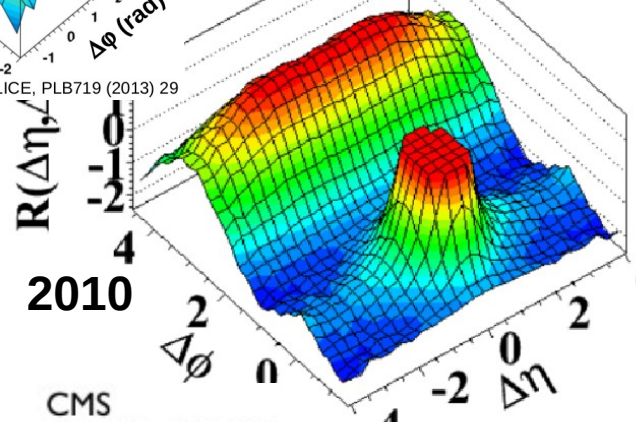
ALICE
 p-Pb $\sqrt{s_{NN}} = 5.02 \text{ TeV}$
 V0S: 0-20%



Phys.Rev.Lett. 117 (2016) 182301
 (b) CMS $\sqrt{s} = 2.36 \text{ TeV}$

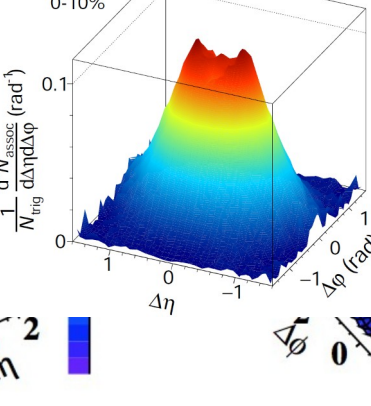
2010

CMS
 JHEP 1009 (2010) 091



ALICE, Pb-Pb
 $\sqrt{s_{NN}} = 2.76 \text{ TeV}$
 0-10%

$1 < p_{T, \text{trig}} < 2 \text{ GeV}/c$
 $1 < p_{T, \text{assoc}} < 2 \text{ GeV}/c$

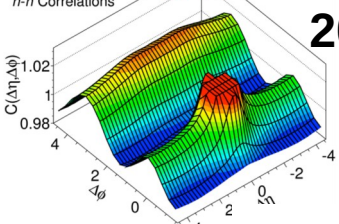


ATLAS Preliminary p+Pb $0.5 < p_T^{a,b} < 5 \text{ GeV}$
 $\sqrt{s_{NN}} = 8.16 \text{ TeV}, 171 \text{ nb}^{-1}$
 h-h Correlations

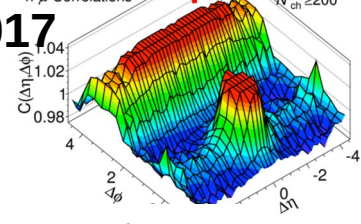
h-h

ATLAS Preliminary p+Pb $0.5 < p_T^1 < 5 \text{ GeV}$
 $\sqrt{s_{NN}} = 8.16 \text{ TeV}, 171 \text{ nb}^{-1}$
 $200 \leq N_{ch}^{rec} < 220$
 h-μ Correlations

μ-h

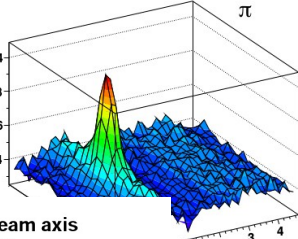


2017



Phys. Lett. B 753 (2016) 126-139

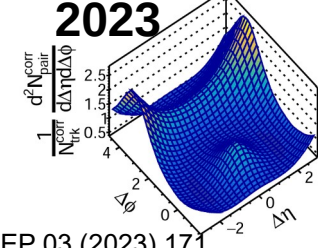
2010



2015

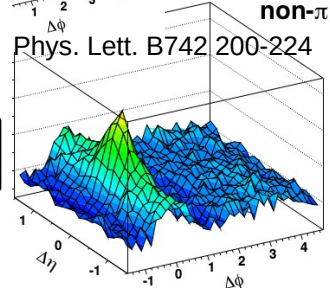
Belle e⁺e⁻, beam axis
 $\sqrt{s} = 10.52 \text{ GeV}$

2023

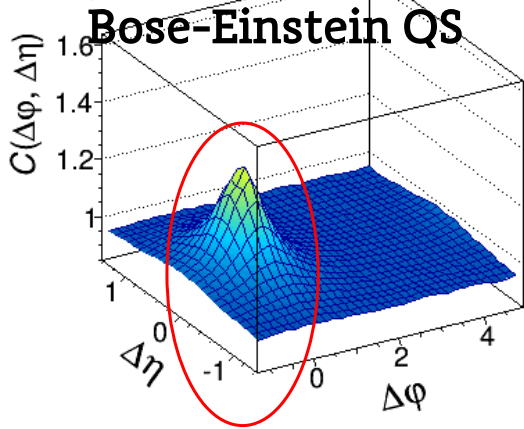


JHEP 03 (2023) 171

Phys. Lett. B742 200-224

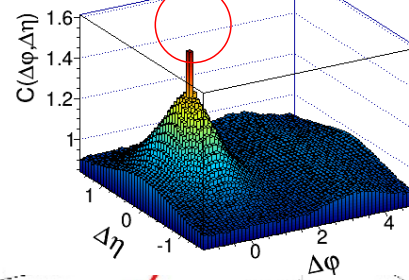


$0.4 < p_{T\text{-sum}} < 0.8 \text{ GeV}/c$

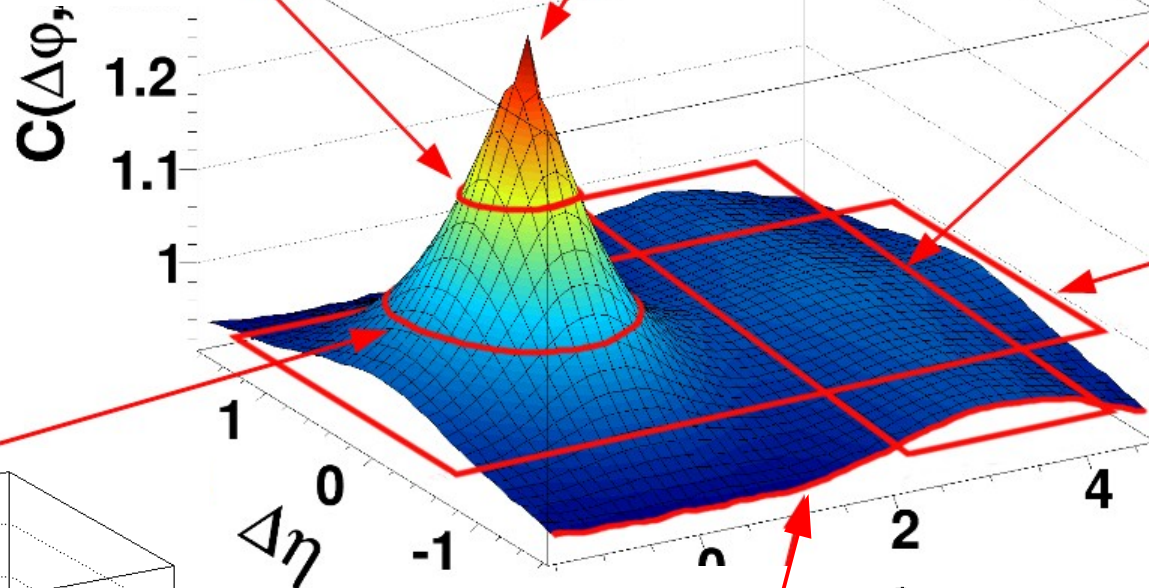
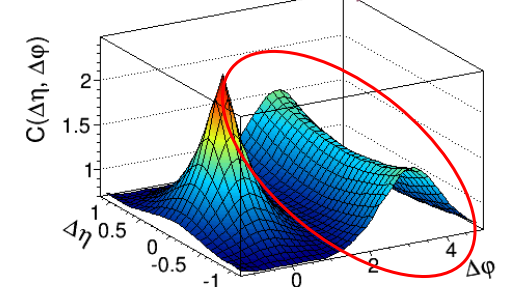


ALICE pp @ 7 TeV
(b) all unlike-sign pairs

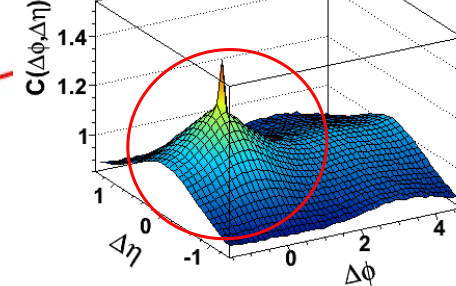
Photon conversion



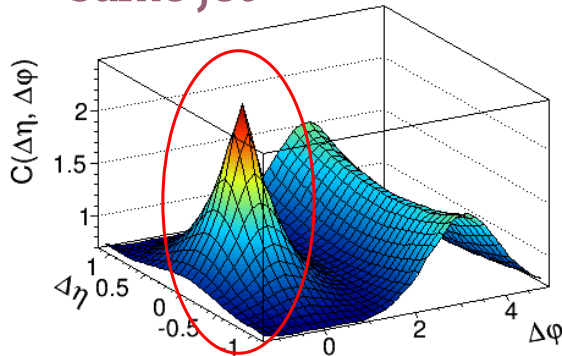
Back-to-back jets



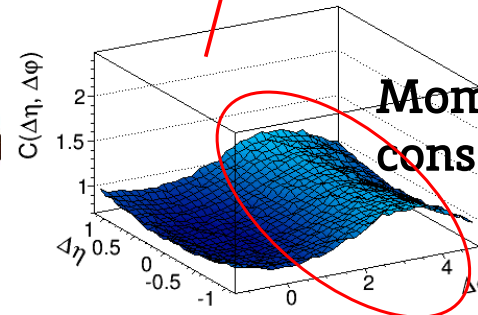
Resonances



Same jet

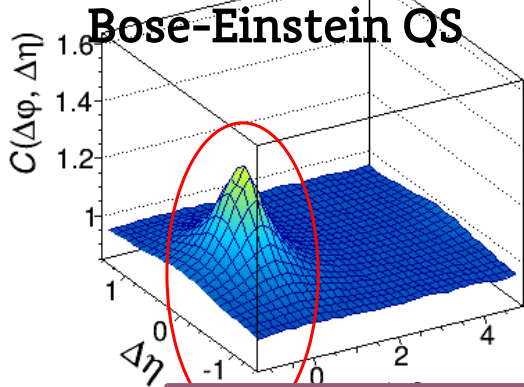


Momentum conservation



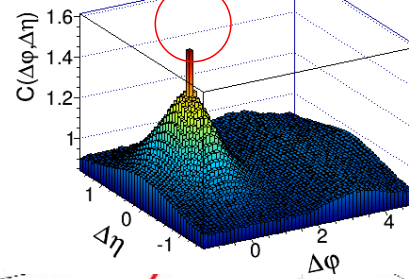
$$\Delta \eta = \eta_1 - \eta_2$$
$$\Delta \varphi = \varphi_1 - \varphi_2$$

$0.4 < p_{T-sum} < 0.8 \text{ GeV}/c$

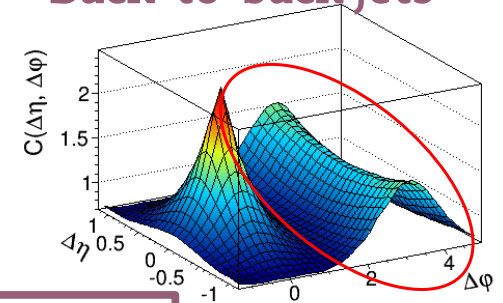


ALICE pp @ 7 TeV
(b) all unlike-sign pairs

Photon conversion

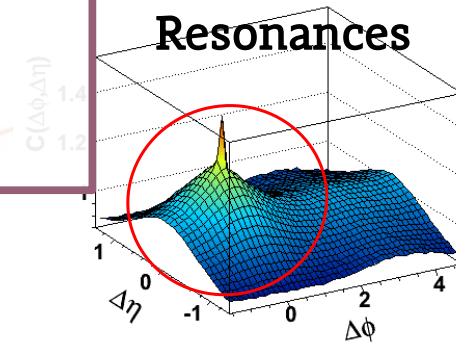


Back-to-back jets

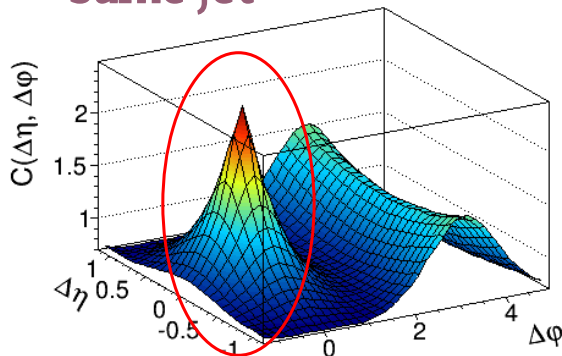


Does this picture always hold?

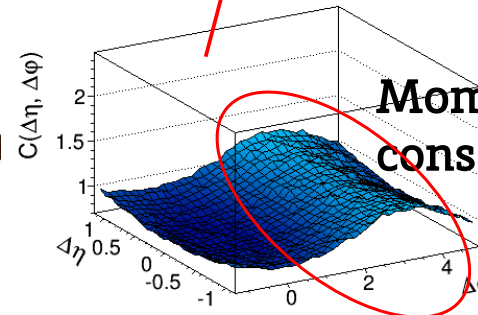
Resonances



Same jet



Momentum conservation



$$\Delta \eta = \eta_1 - \eta_2$$

$$\Delta \varphi = \varphi_1 - \varphi_2$$

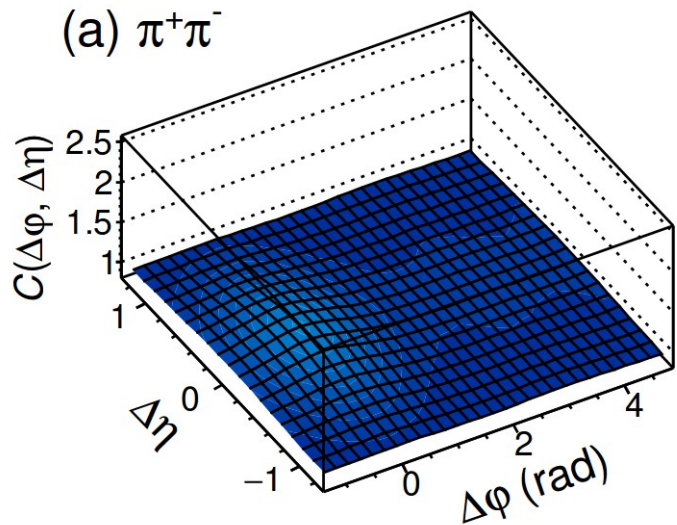


ALICE 7 TeV pp data – identified particles

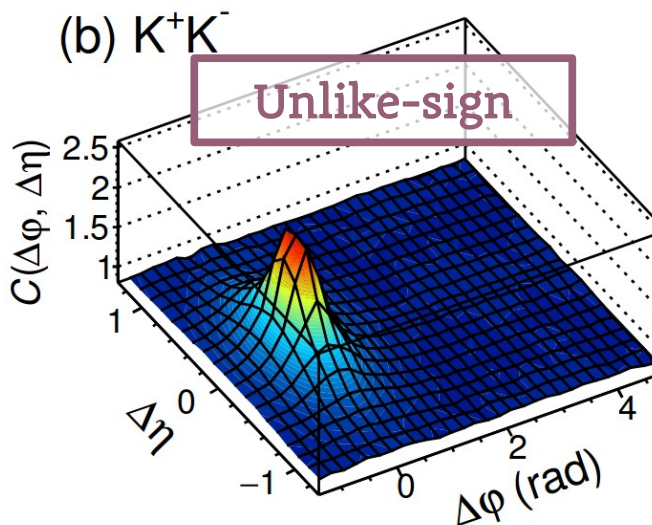


ALICE, Eur. Phys. J. C 77 (2017) 569, Ph.D. thesis of M. Janik <https://cds.cern.ch/record/2093543>

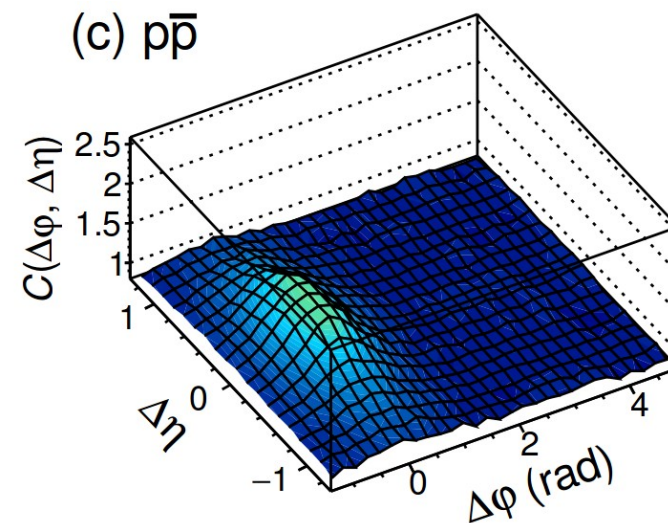
(a) $\pi^+\pi^-$



(b) K^+K^-



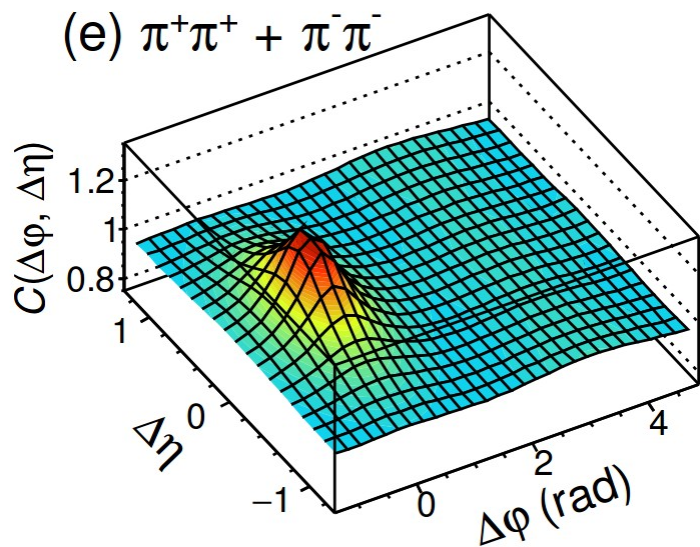
(c) $p\bar{p}$



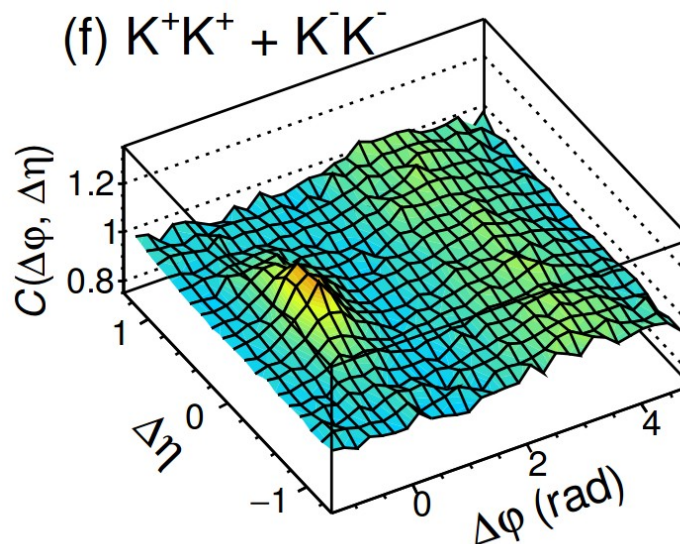
ALICE pp @ 7TeV

Like-sign

(e) $\pi^+\pi^+ + \pi^-\pi^-$

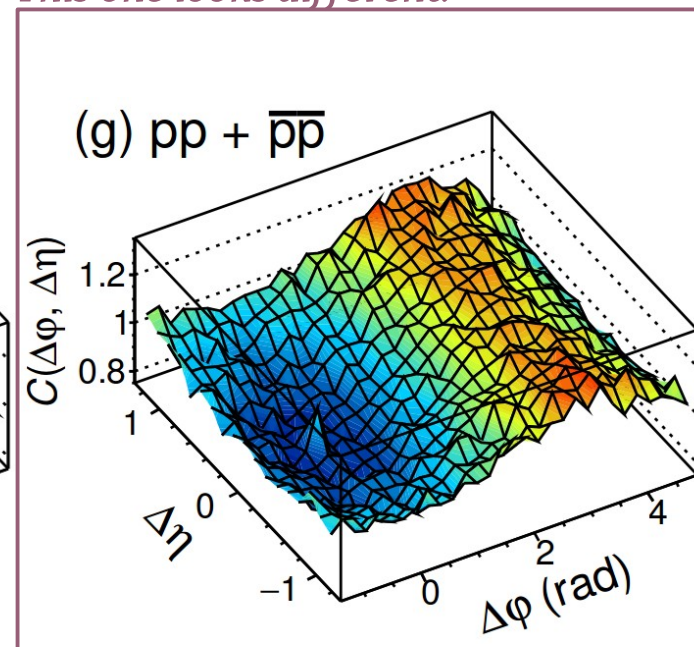


(f) $K^+K^+ + K^-K^-$

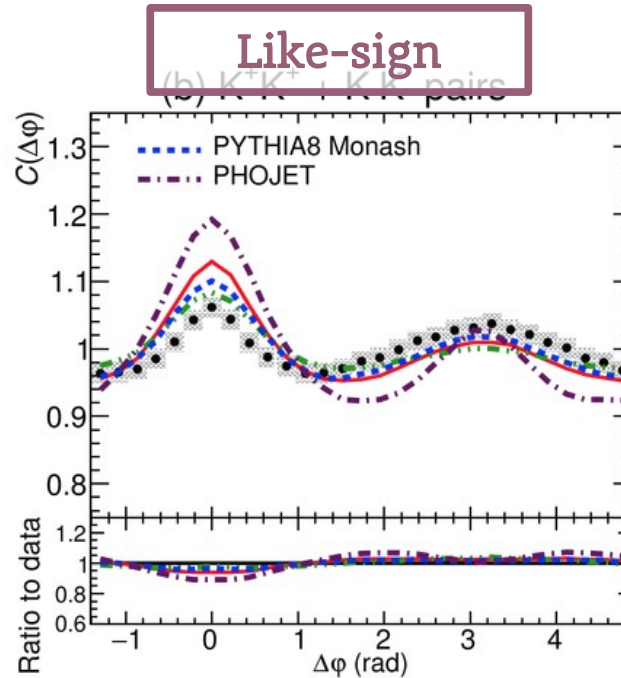
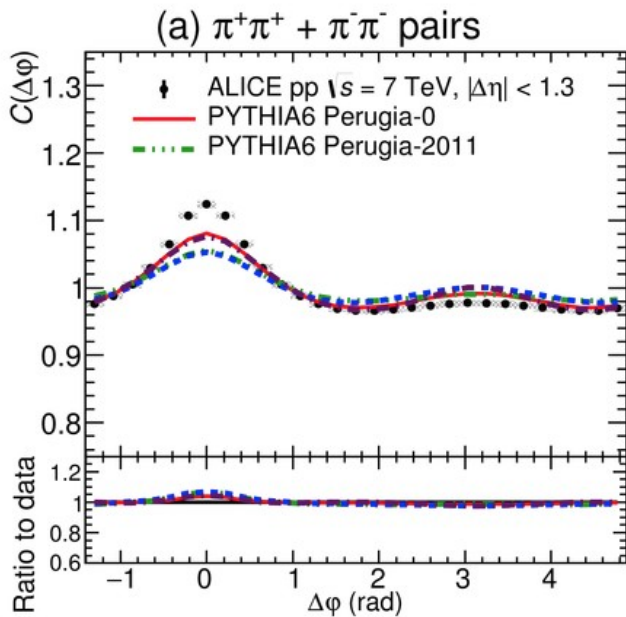
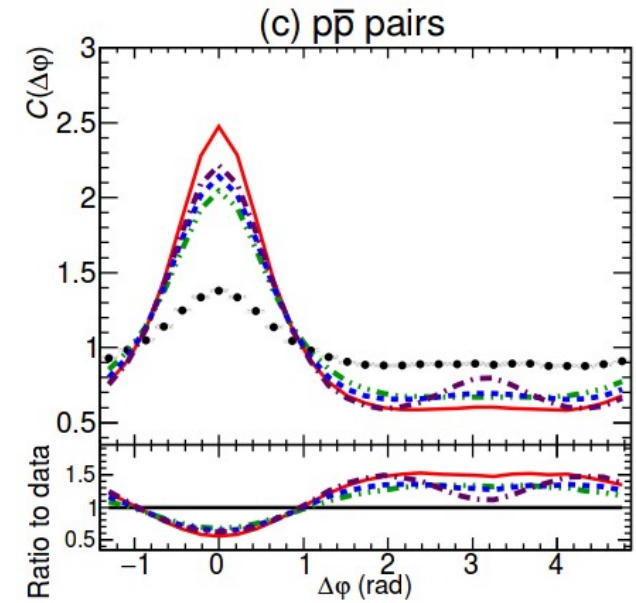
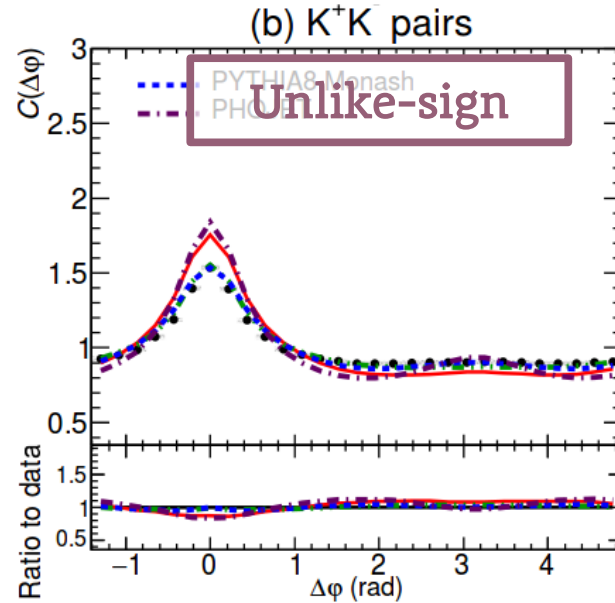
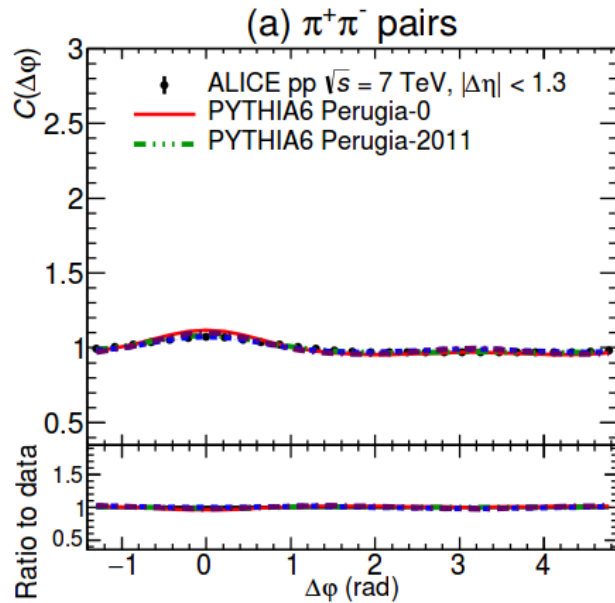


This one looks different!

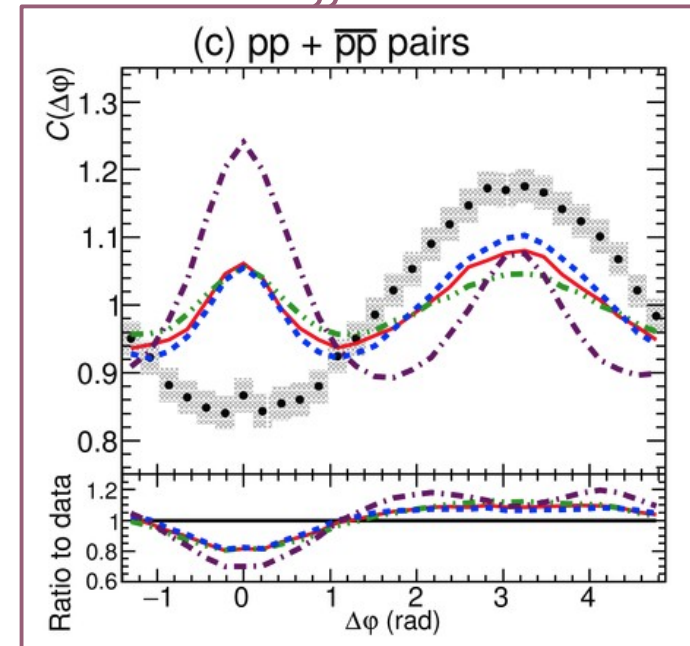
(g) $pp + \bar{p}\bar{p}$



ALICE, Eur. Phys. J. C 77 (2017) 569



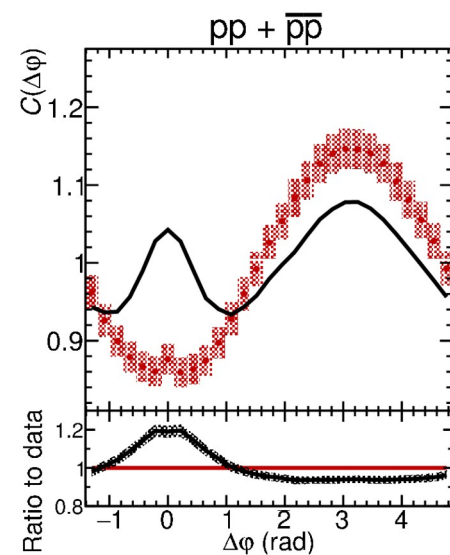
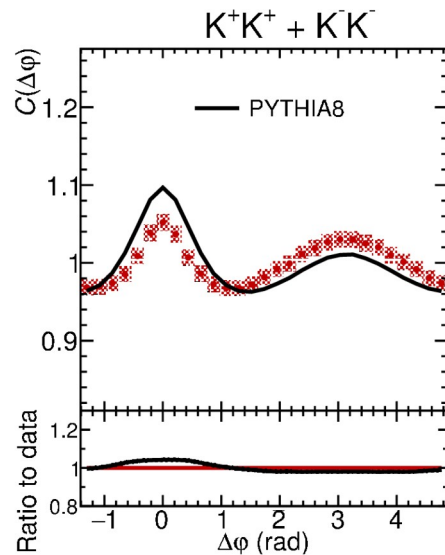
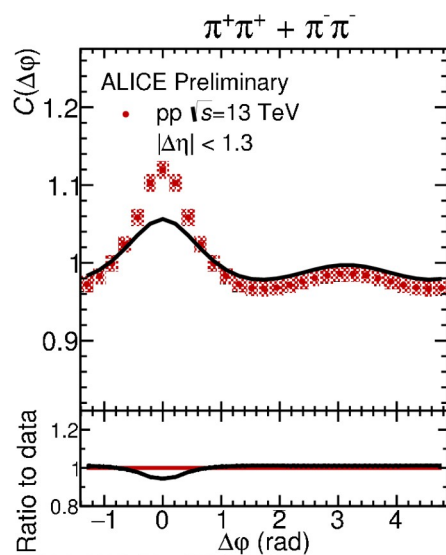
This one looks different!



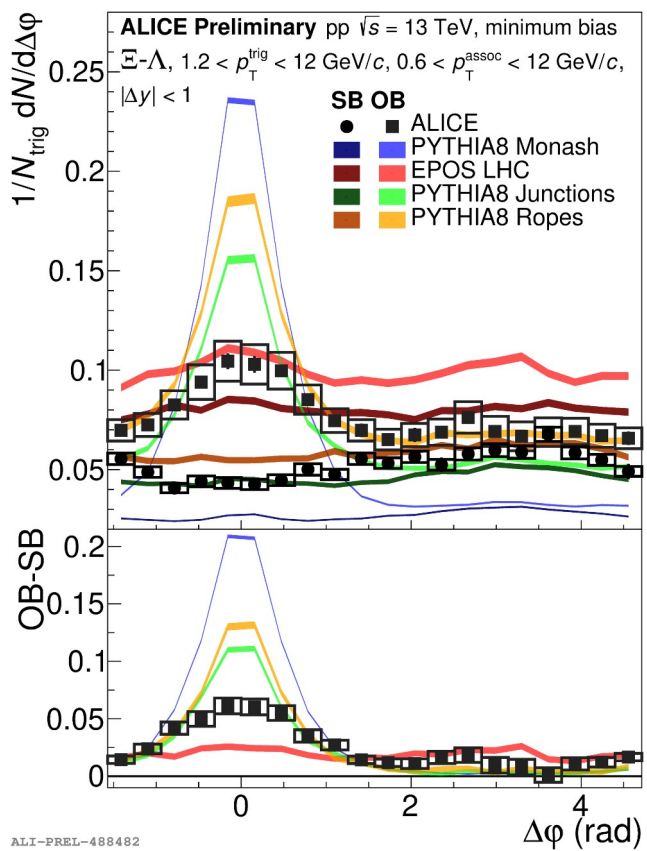


**Are there any advances since the 2017
ALICE paper?**

- The anticorrelation persists at 13 TeV collision energy
- It also persists for higher mass multi-strange baryons

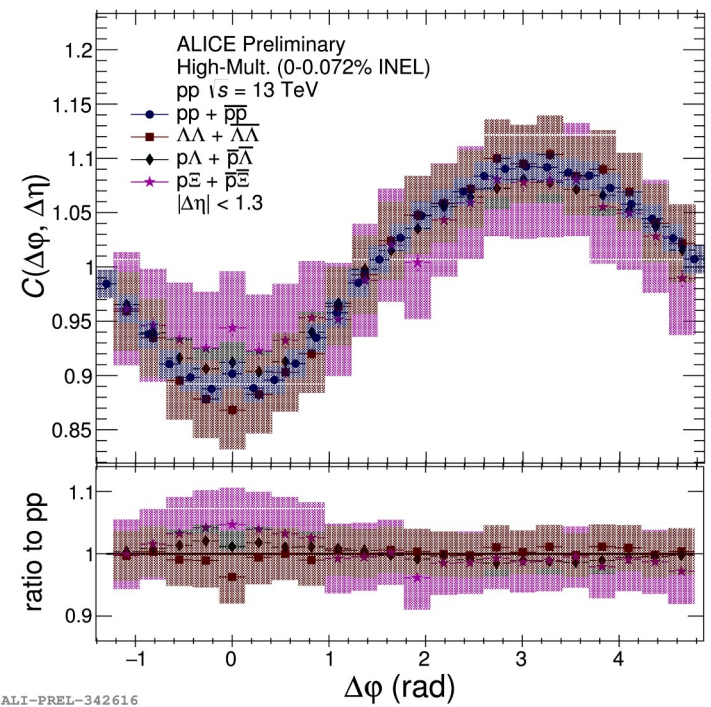
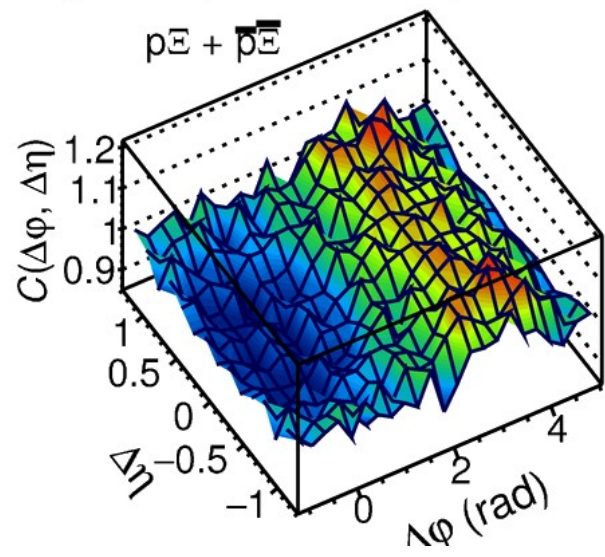


ALI-PREL-338139



ALI-PREL-488482

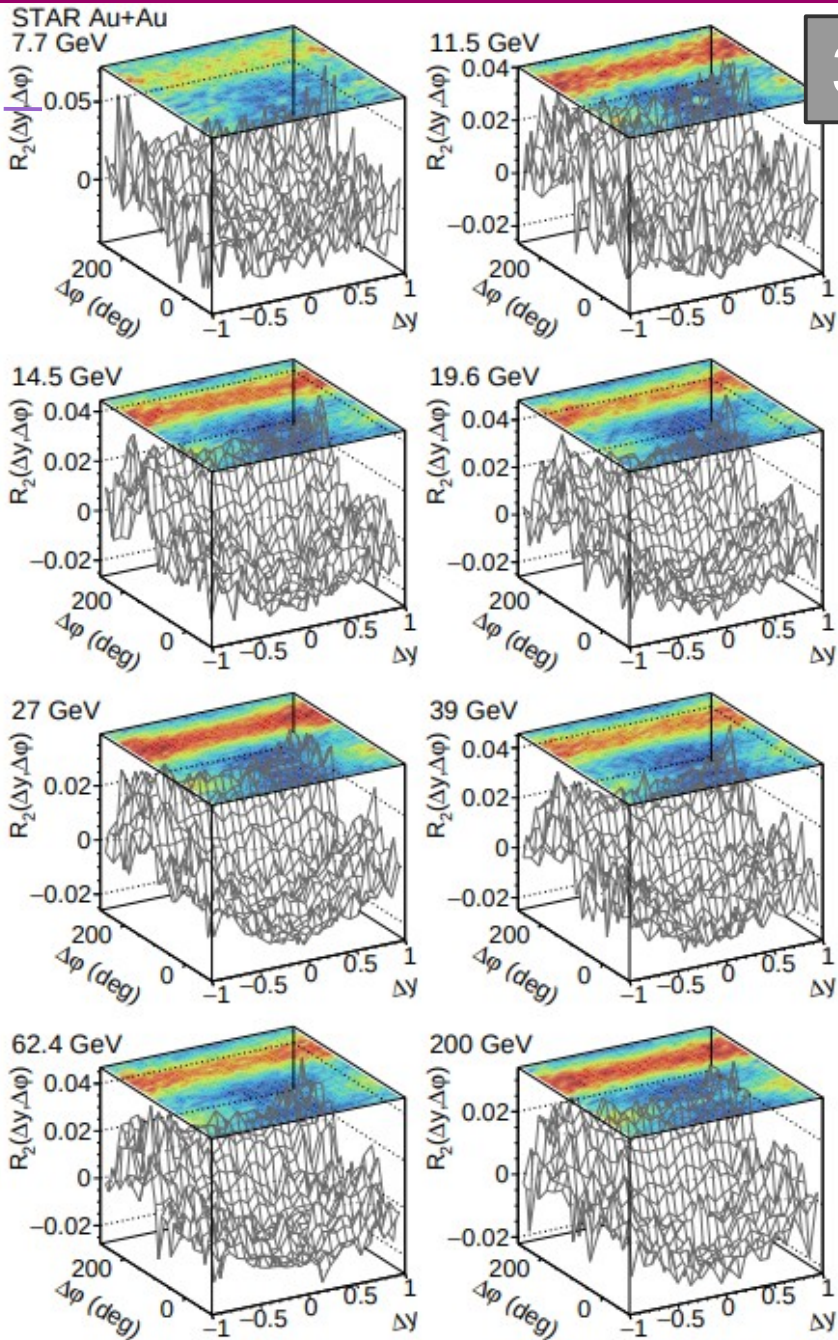
ALICE Preliminary, pp $\sqrt{s} = 13$ TeV
High-Mult. (0-0.072% INEL)



ALI-PREL-342616

30-40%

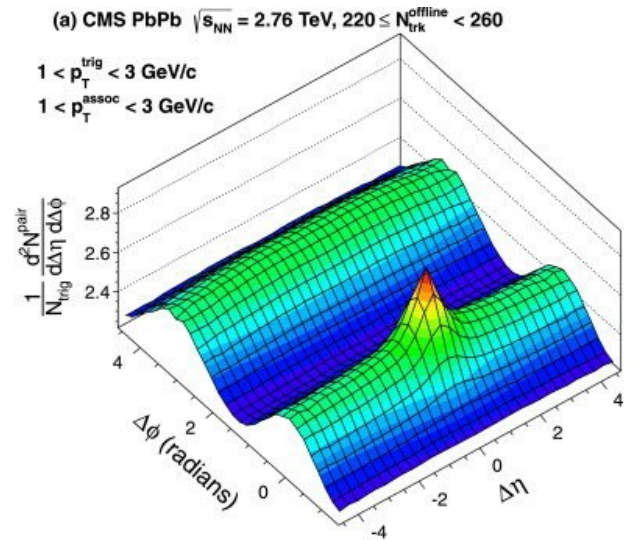
Au+Au collision energy



(a) Like-sign protons

STAR, Phys. Rev. C 101, 014916 (2020)

- The anticorrelation effect is present for Au-Au results
- It is convoluted with the flow double-ridge structure



CMS, Phys. Lett. B 724 (2013) 213



What about the theory side?



A Parametrization of the Properties of Quark Jets
 R.D. Field, R.P. Feynman
 Nucl. Phys. B 136 (1978) 131

From mechanism of jet production:
 Two primary hadrons with the same
baryon number
 are separated by at least
 two steps in "rank"

R. Feynman
 "Quark Jets"
 8th ISMD 1977

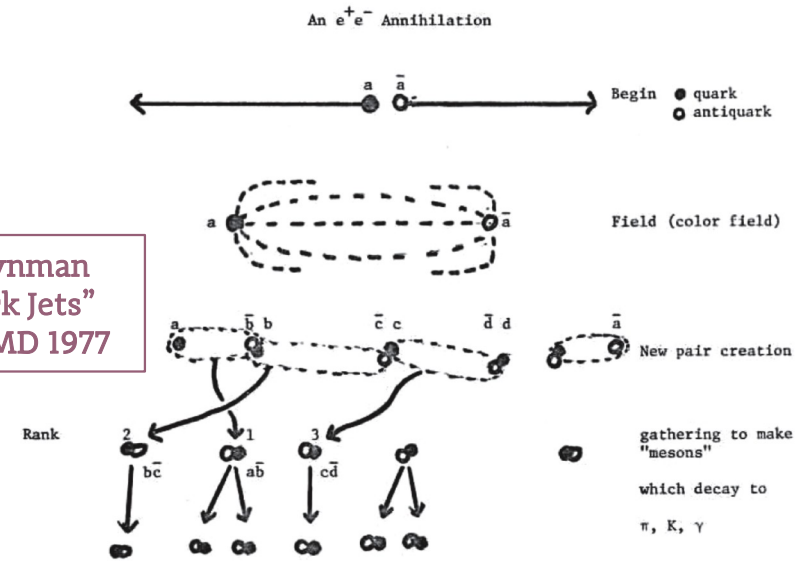
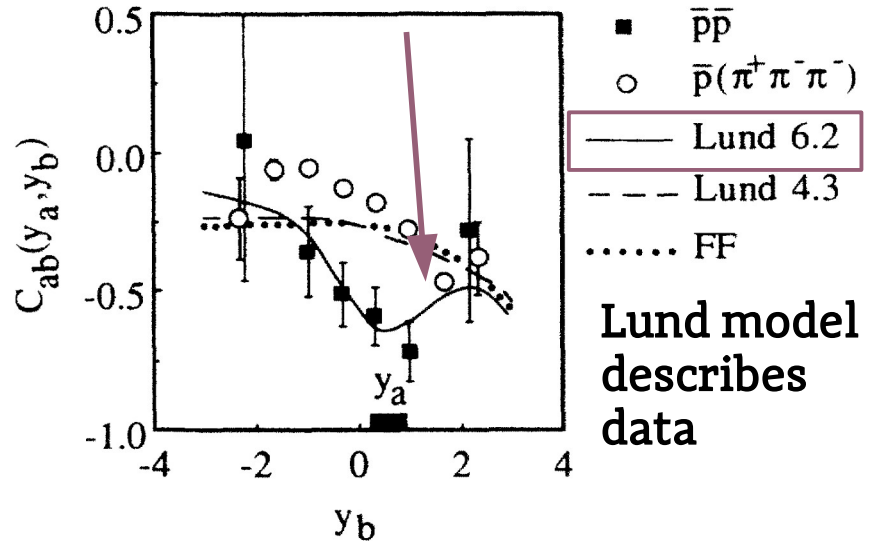
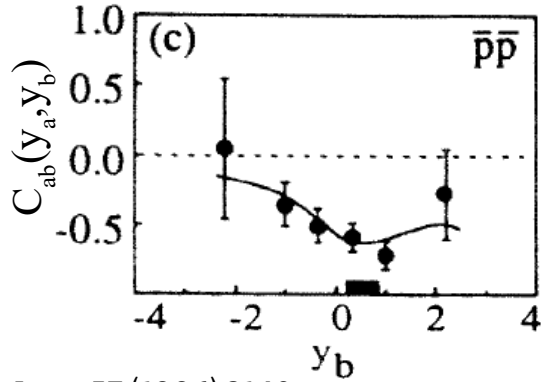
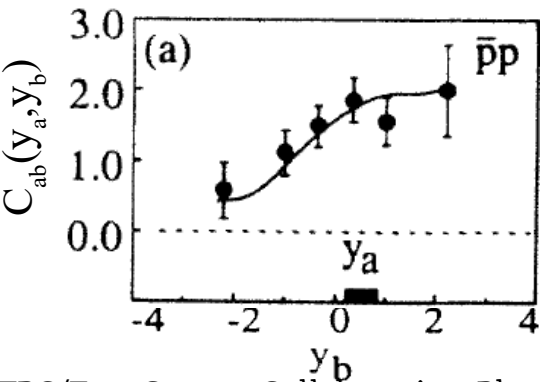


Fig. 10. Transparency from a talk Feynman gave on our model for how quarks fragment into hadrons at the International Symposium on Multiparticle Dynamics (ISMD), Kaysersberg, France, June 12, 1977.

We are not likely to find two baryons or two antibaryons very close to each other

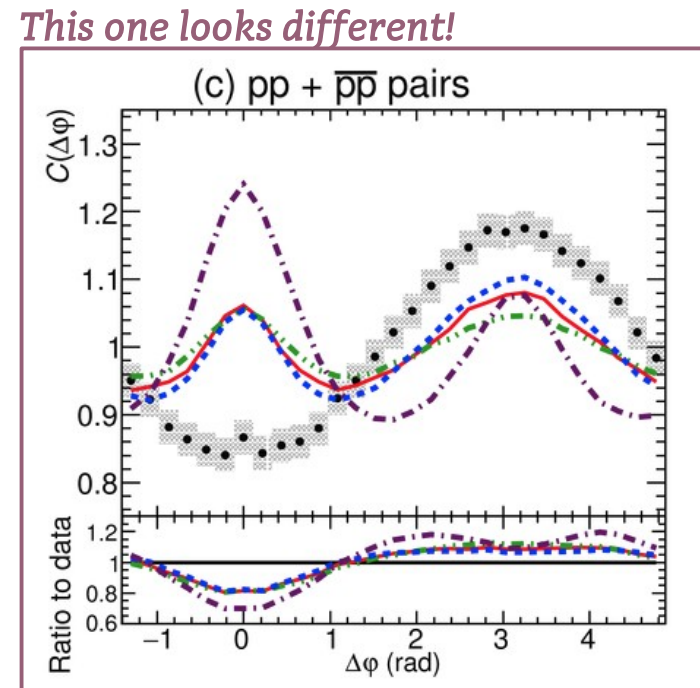
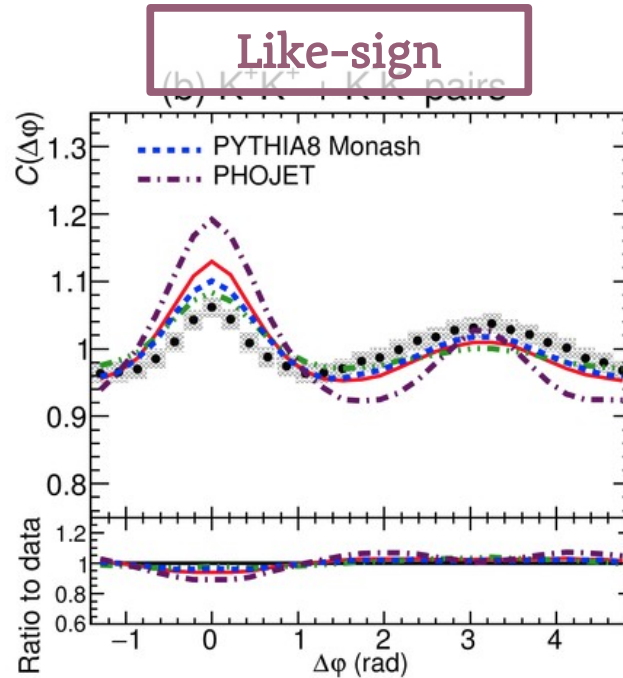
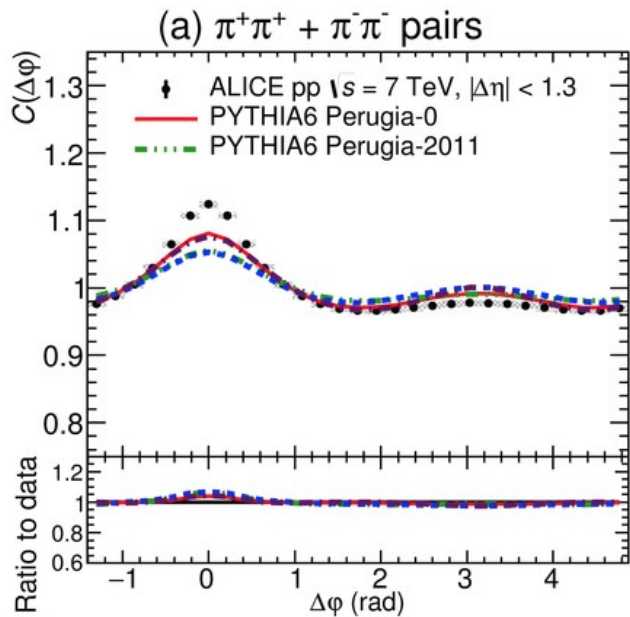
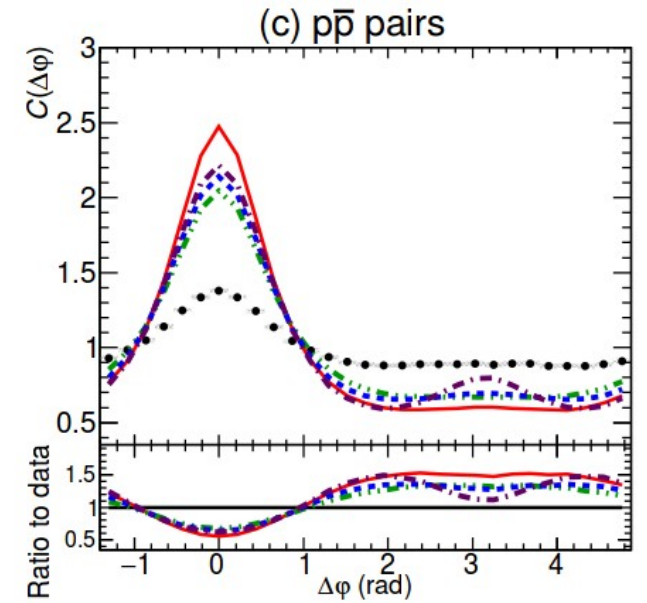
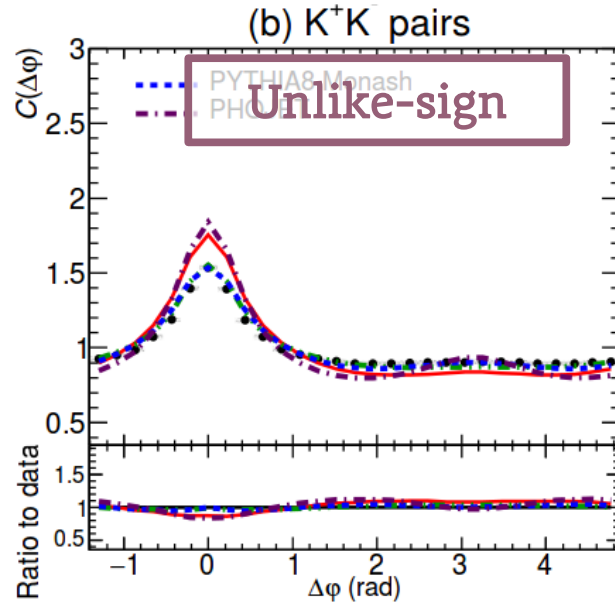
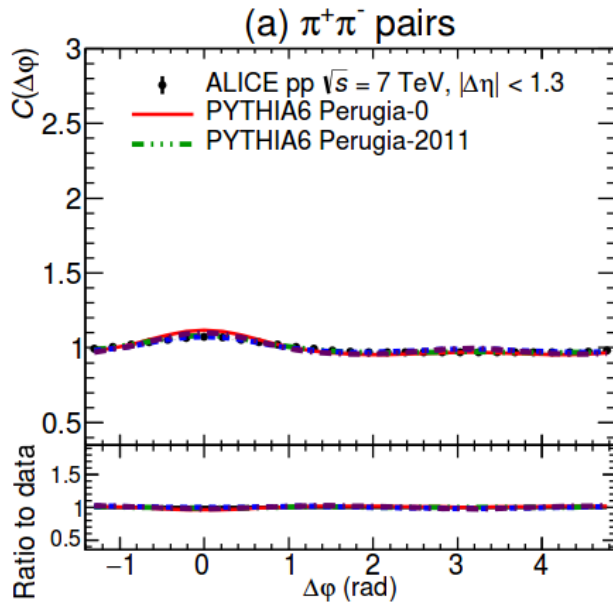
Local baryon number conservation partially responsible for anticorrelation at 29 GeV collision energy

Models at lower energies agree with observations seen in data

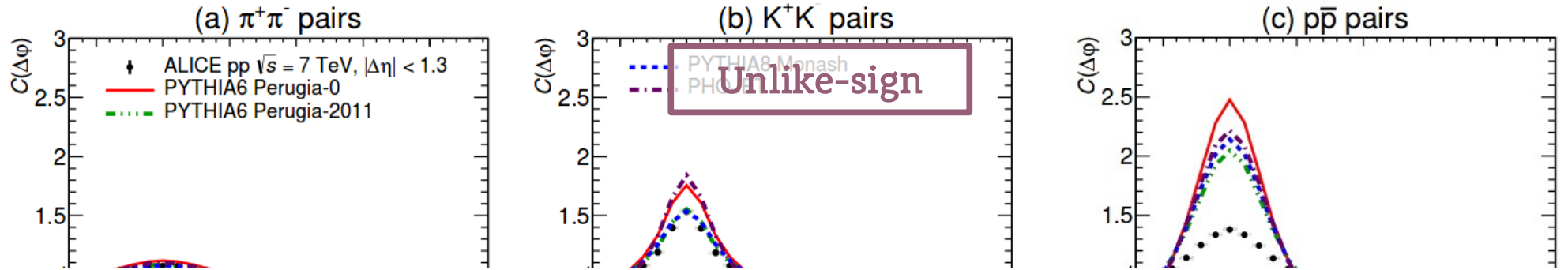


TPC/Two Gamma Collaboration, Phys.Rev.Lett. 57 (1986) 3140

ALICE, Eur. Phys. J. C 77 (2017) 569



ALICE, Eur. Phys. J. C 77 (2017) 569



T. Sjostrand, QM 2018, plenary talk
<https://indico.cern.ch/event/656452/contributions/2899749/>



Nucl. Phys. A 982 (2019) 43-49

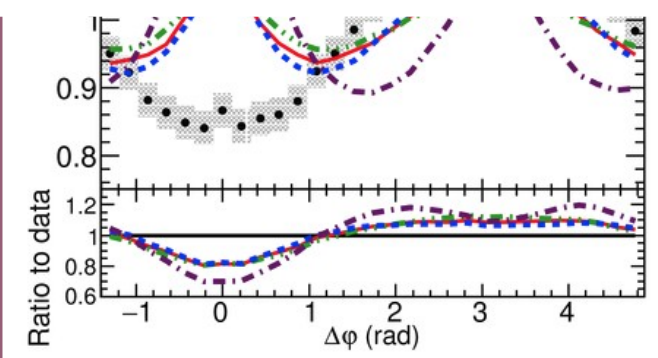
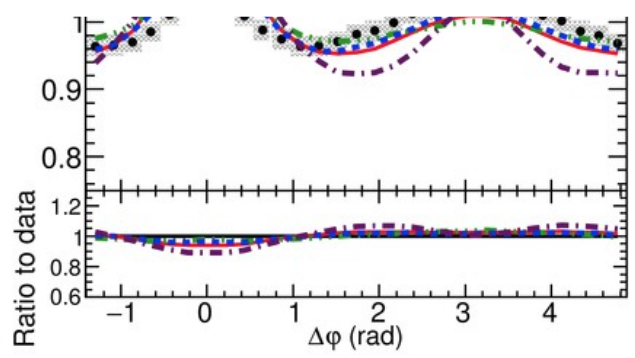
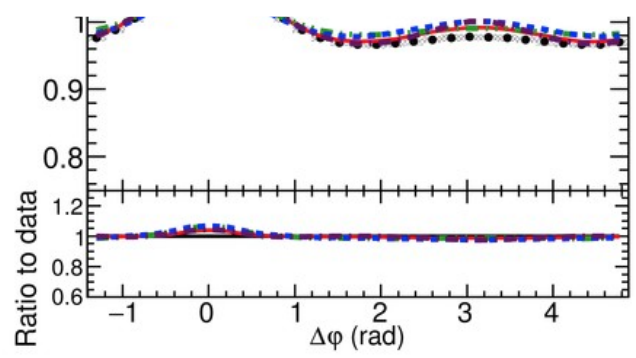
“The real problem is baryon production. [...] so it is clear we still lack some fundamental insight on baryon production, at least in the string context.”

Collective Effects:
 the viewpoint of HEP MC codes

Torbjörn Sjöstrand

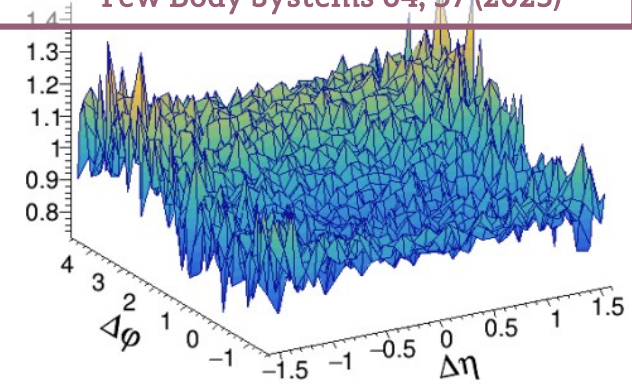
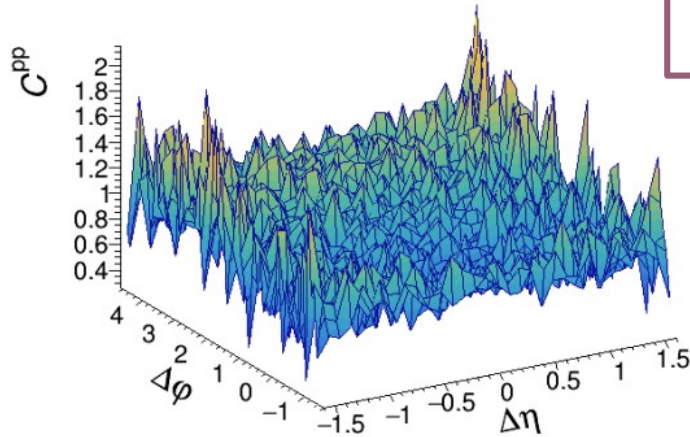
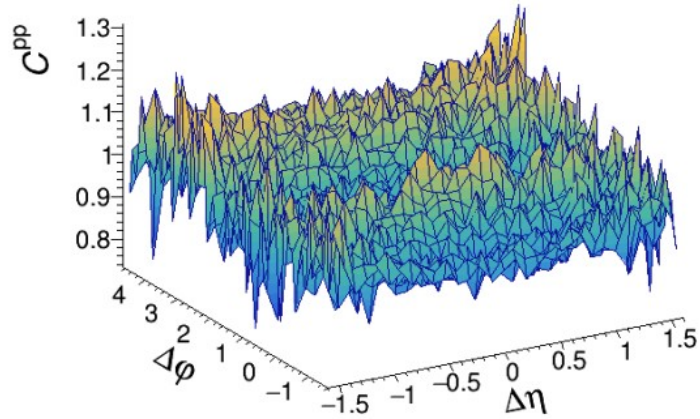
Department of Astronomy and Theoretical Physics
 Lund University
 Sölvegatan 14A, SE-223 62 Lund, Sweden

Quark Matter 2018, Venice, 13–19 May 2018

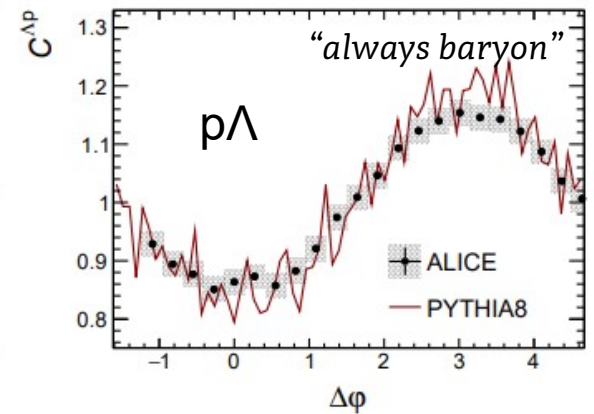
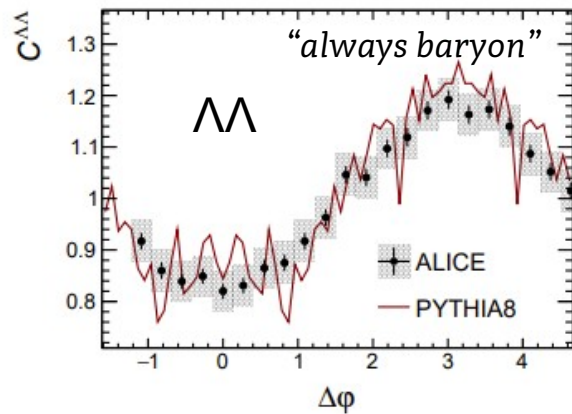
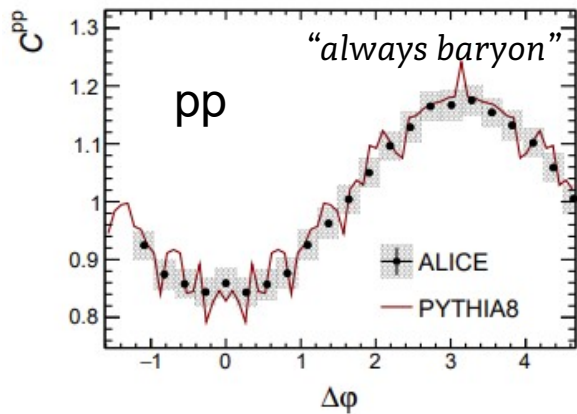




N. Demazure, V. Gonzalez, F. Llanes-Estrada
Few Body Systems 64, 57 (2023)



(a) pp C correlation with unmodified PYTHIA (b) pp C correlation, one-baryon per string policy (c) pp C correlation, always-baryon policy

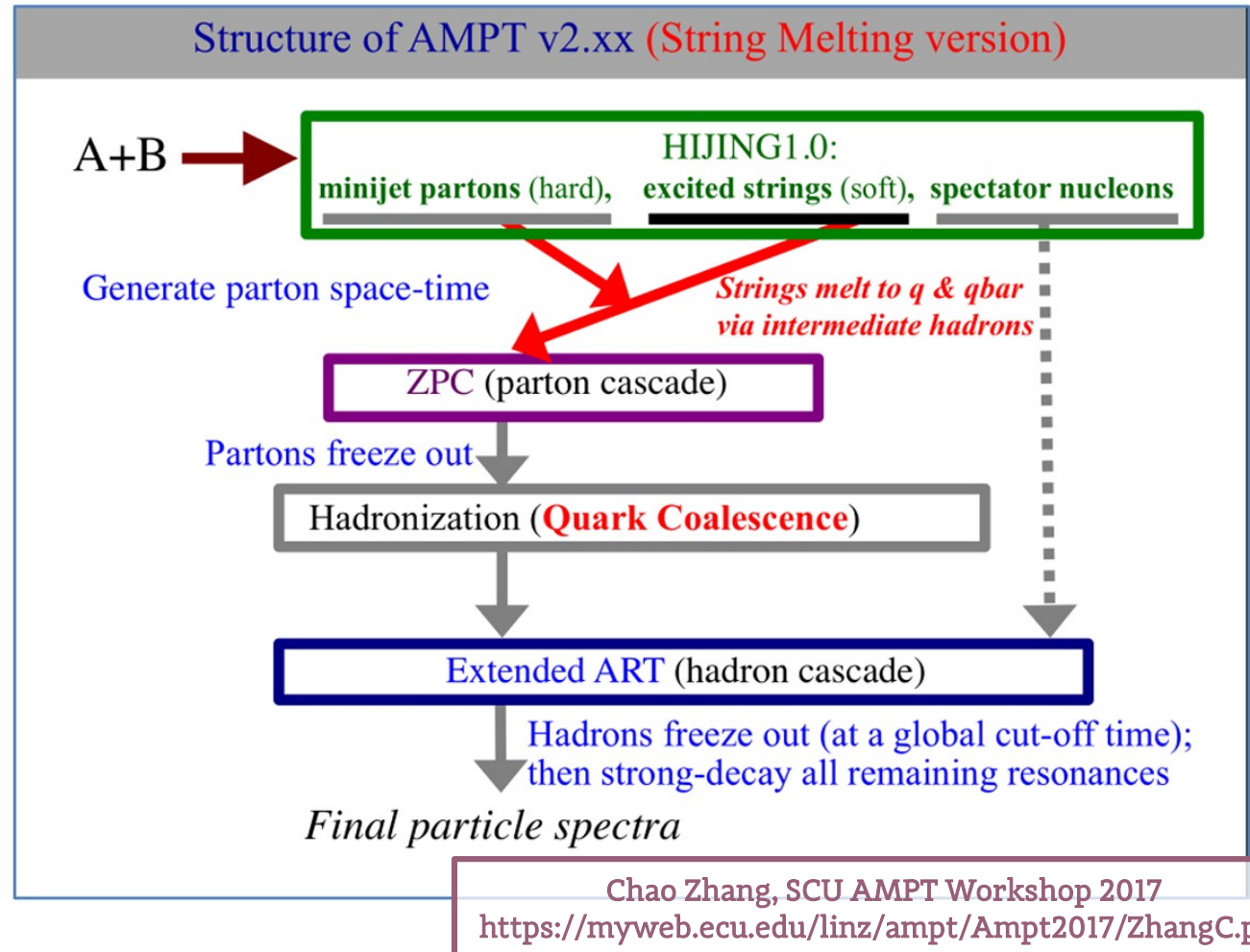
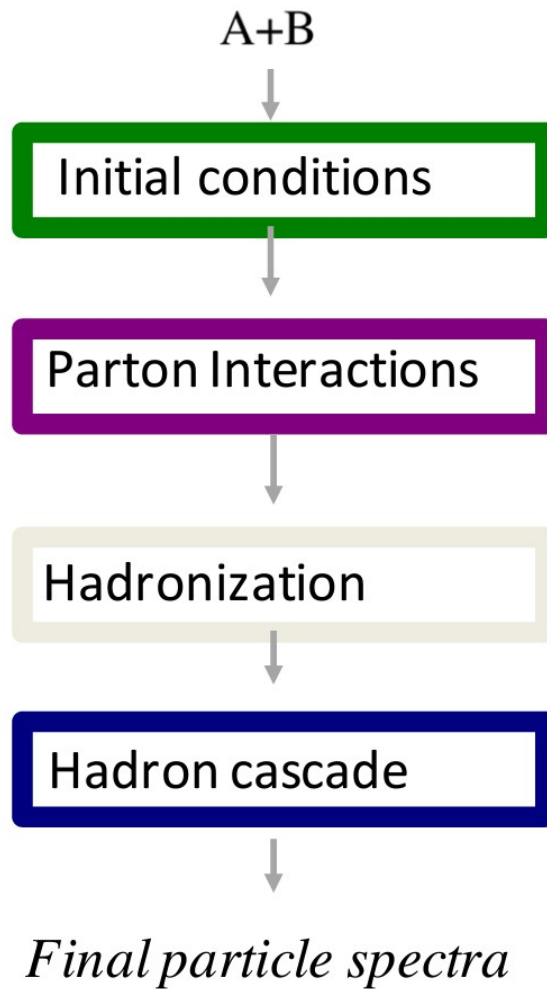


- Two modifications to PYTHIA string fragmentation allow the model to describe the data:
 - *one baryon* – each string must *produce at most one baryon* (a way to impose Pauli principle to baryons, but lowers the baryon-to-meson ratio)
 - *always baryon* – each string must *always produce one baryon* (no physical meaning, but produces very good agreement with data)

AMPT model



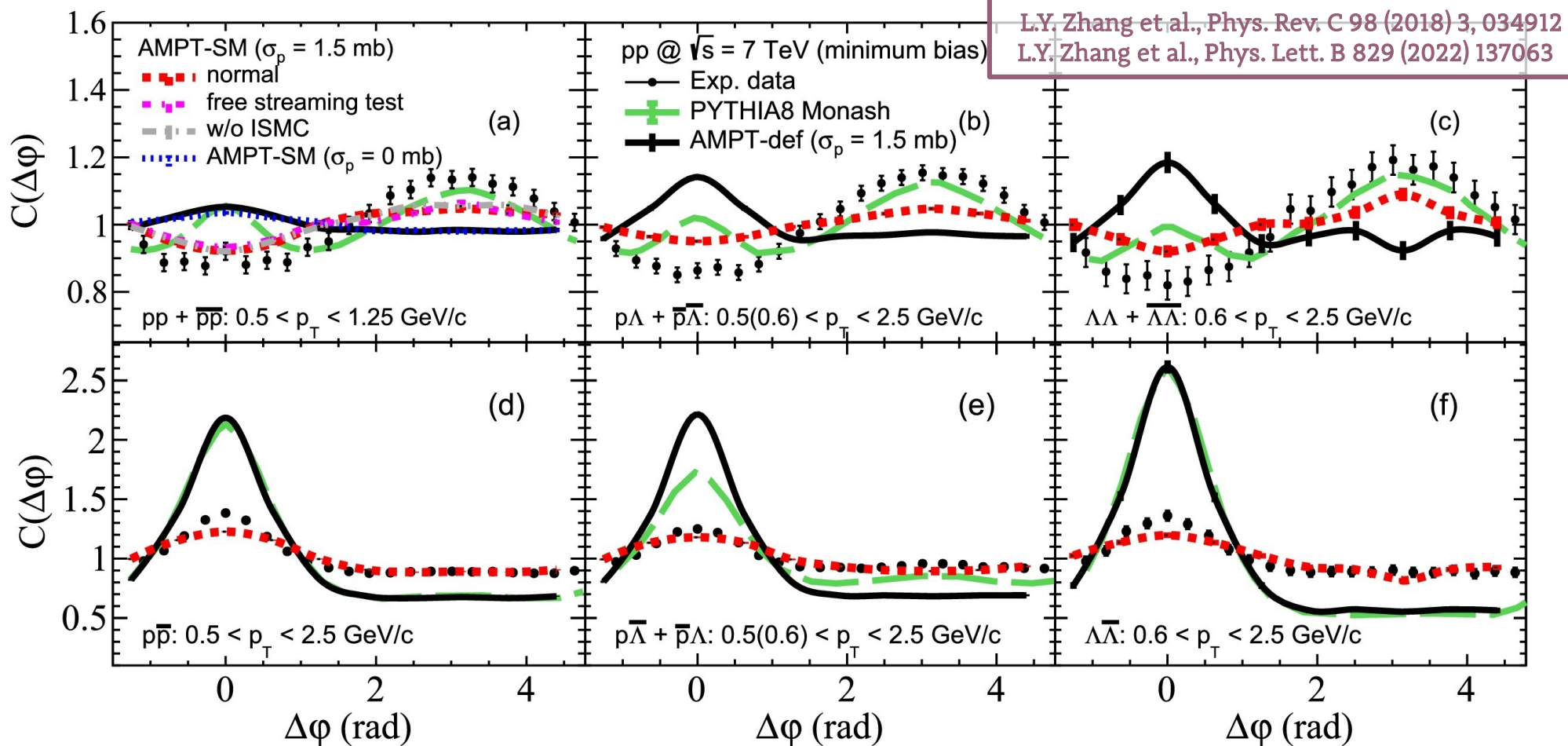
- Contains *4 main components* to describe the whole phase space of heavy-ion collisions
- *String melting*: convert hadrons from string fragmentation into quarks and antiquarks
- *Coalescence*: when partons stop interacting, combine nearest quark and antiquark to meson, and nearest three quarks to baryon



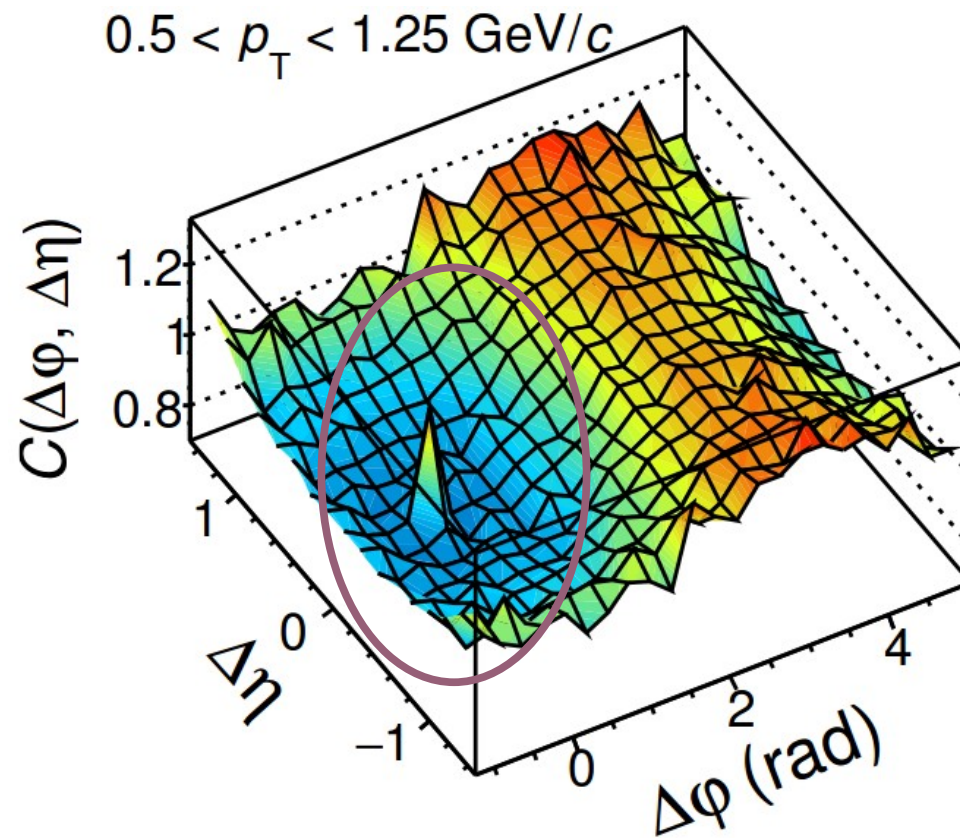
Modified AMPT



- *Improved* coalescence (removed separate conservation for mesons and baryons)
- String melting (SM) → parton degrees of freedom are expected in the initial state
 - **AMPT-SM** with non-zero parton cross section describes the data
 - test of **SM with parton cross section set to 0 mb** does not describe the data
- If initial state momentum correlation (ISMC) are removed → the result is similar to standard AMPT-SM version → describes anticorrelation

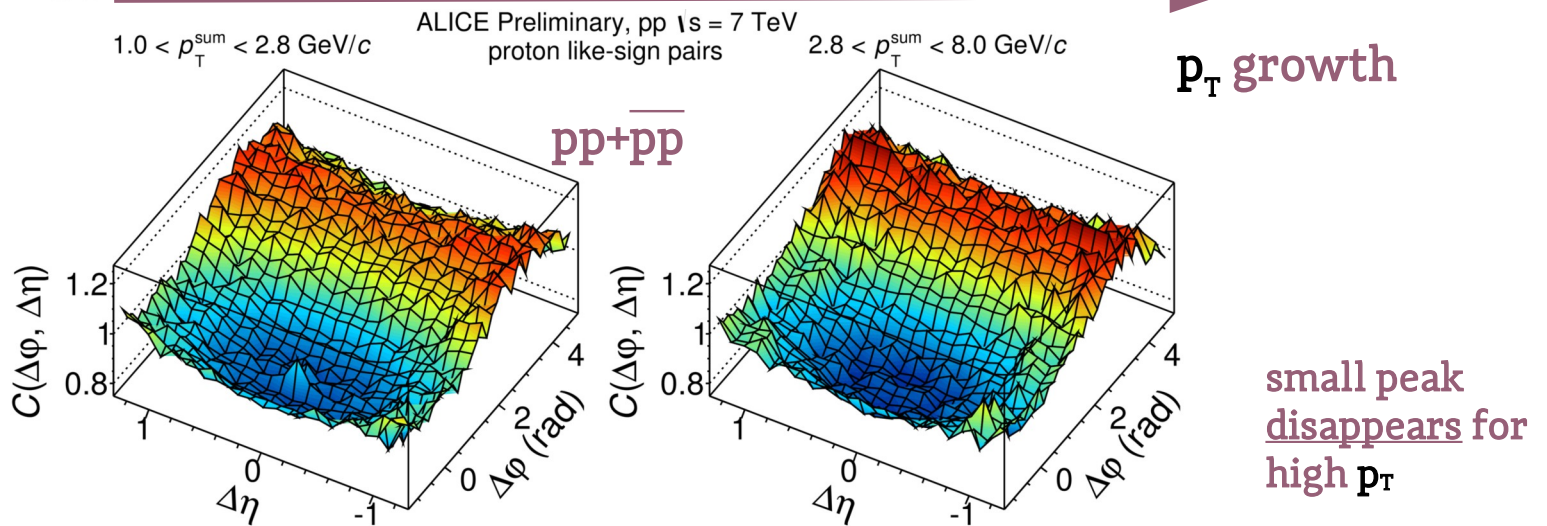
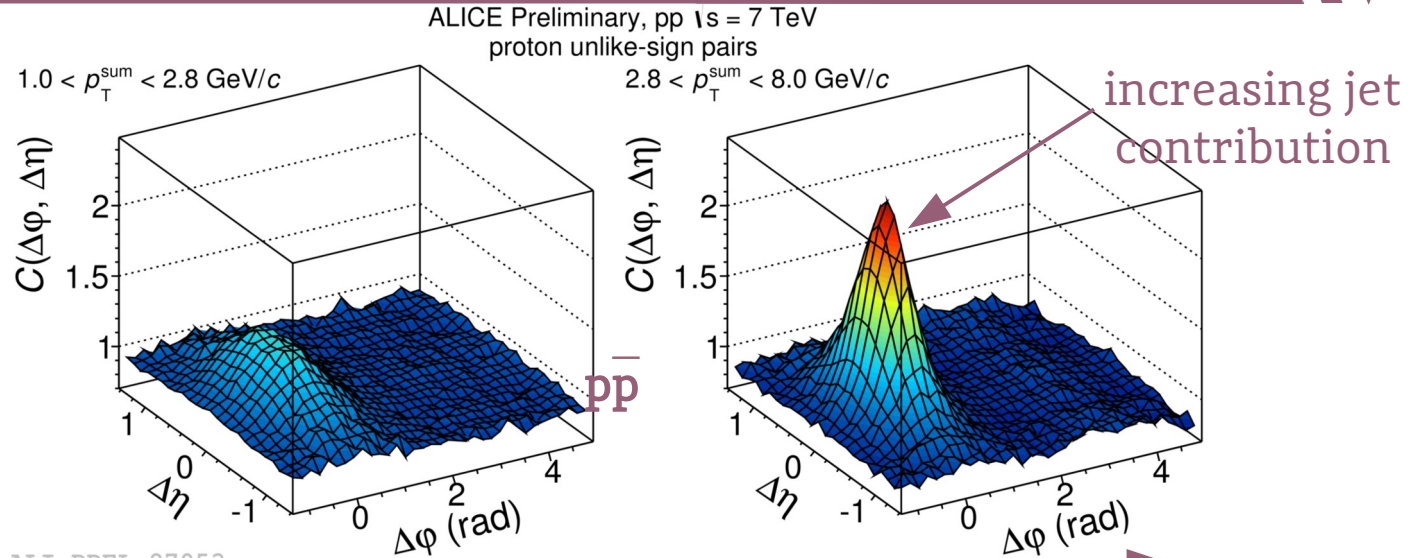


What is the origin of the “small peak” in pp correlations?





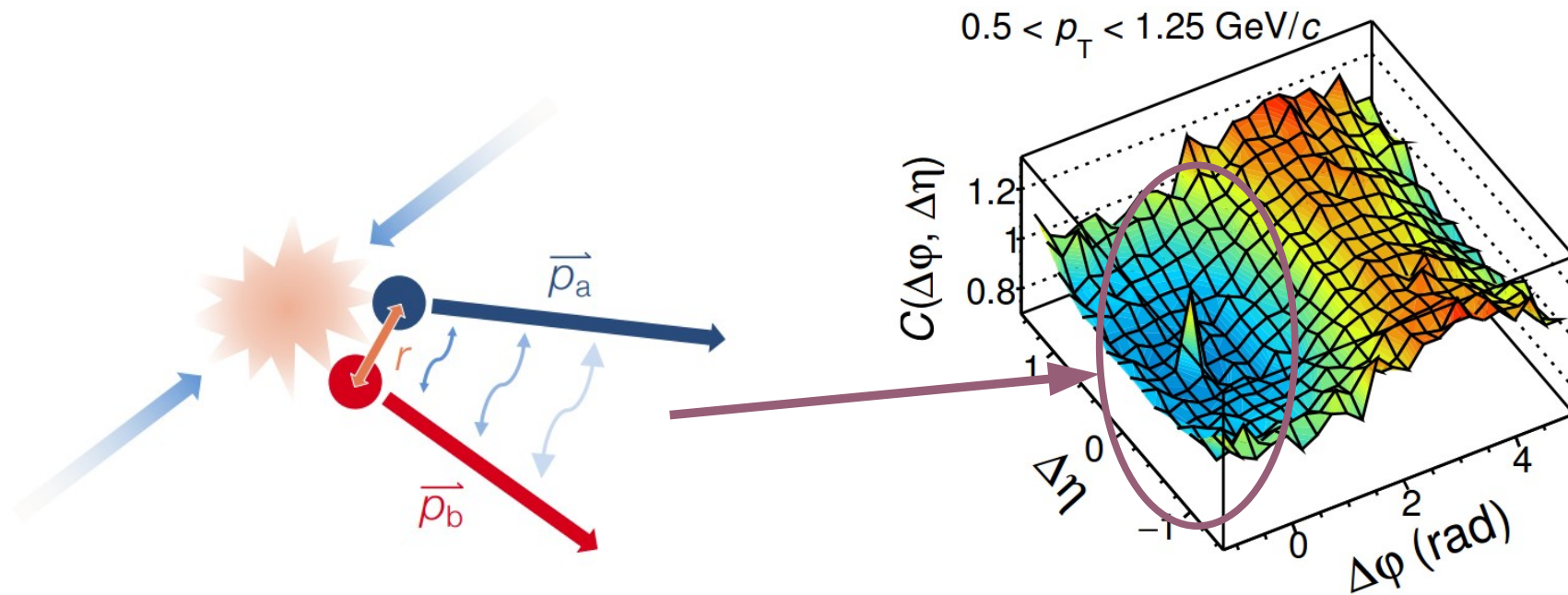
Near-side peak grows with p_T (more contribution from jets)



p_T growth

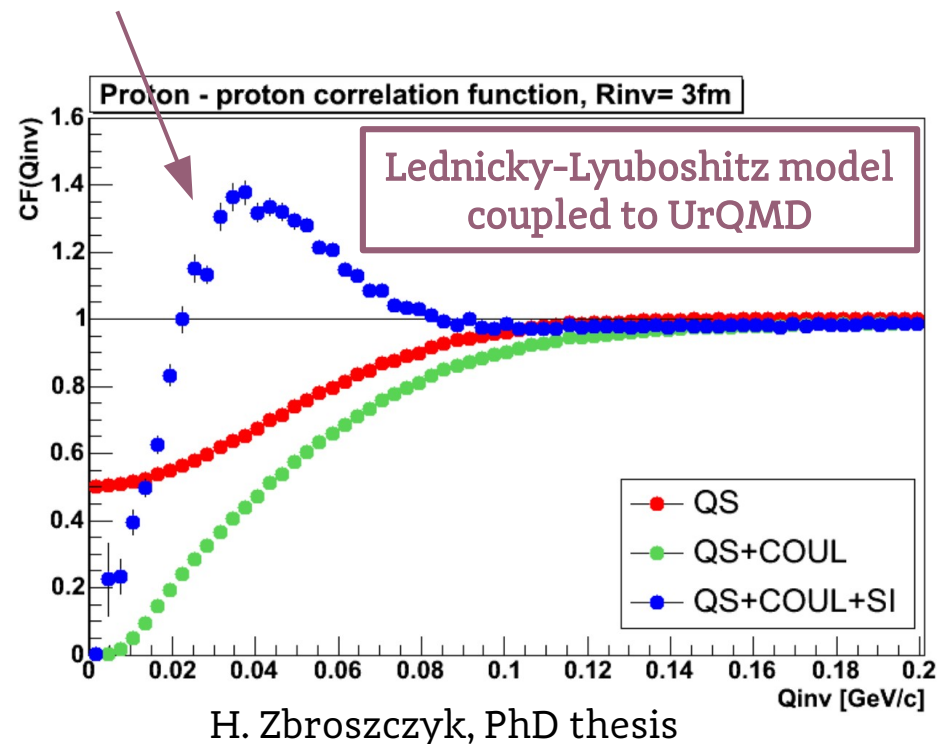
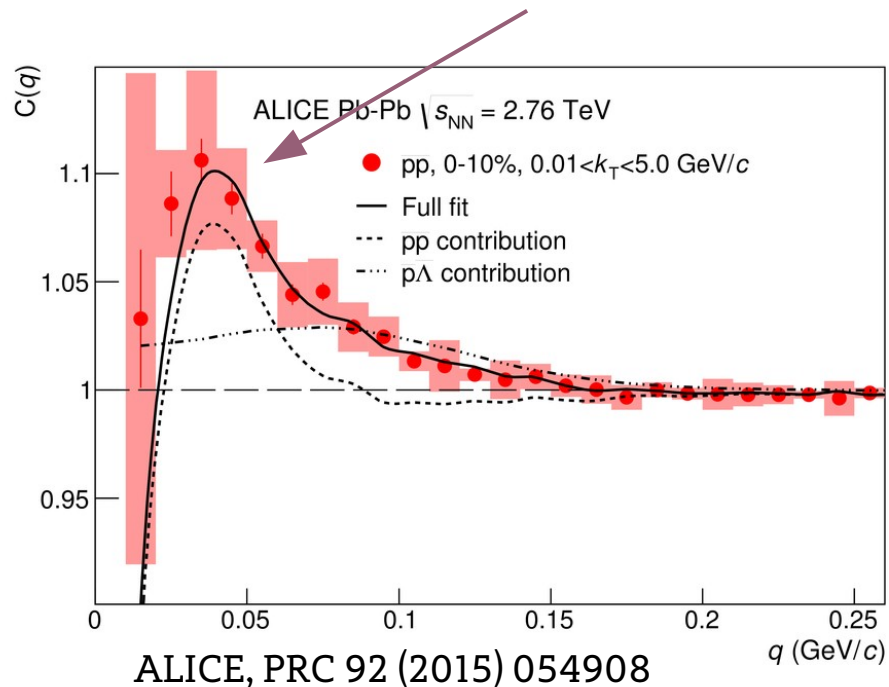
- The small peak seems to behave **strangely** → decreases with increasing p_T
- Is it an unnoticed and not removed **detector effect** OR is there some **physics** behind it?

- In the ALICE paper we *hypothesized* the small peak could be of the strong final-state interaction (FSI) origin:



Ingredients of the proton-proton femtosopic correlation function:

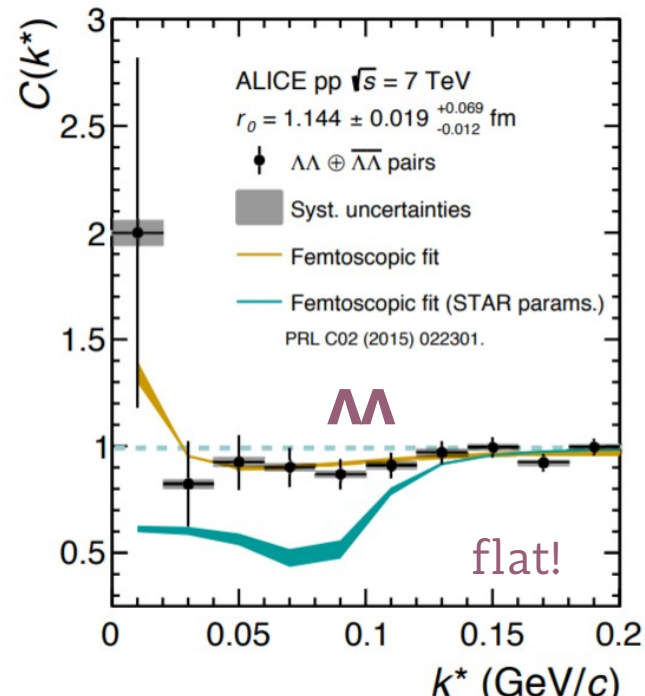
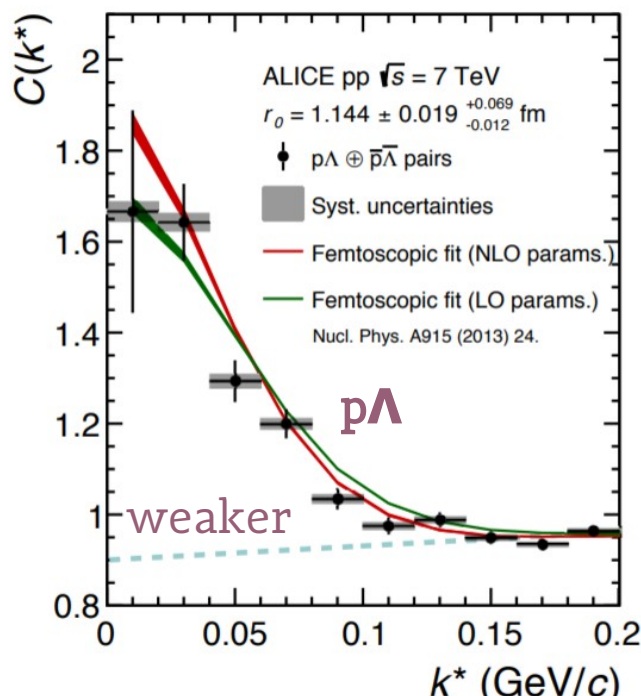
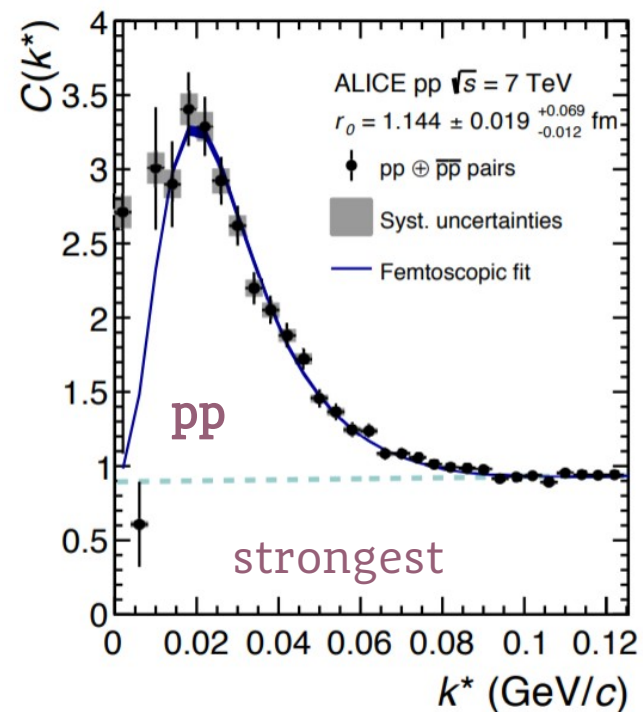
- strong FSI is significant in pp femtosopic correlation function
- dominant effect around $q = 0.04$ GeV/c
- **strong interaction** is the only source of positive correlation for baryons



Strong FSI for other baryon pairs



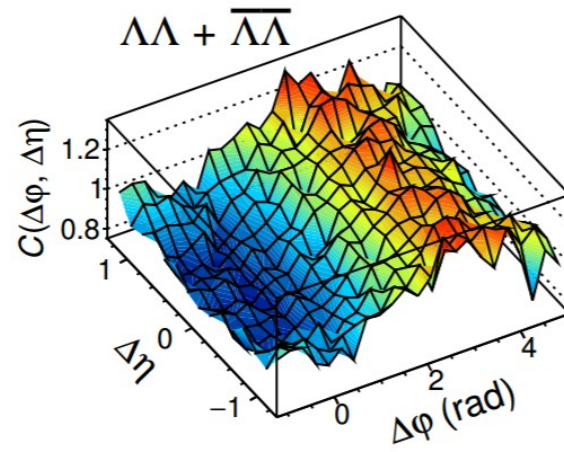
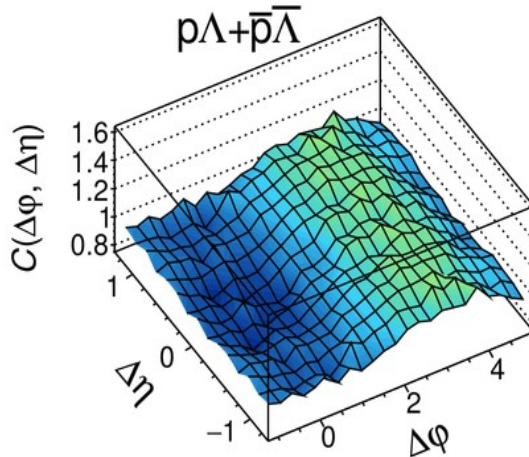
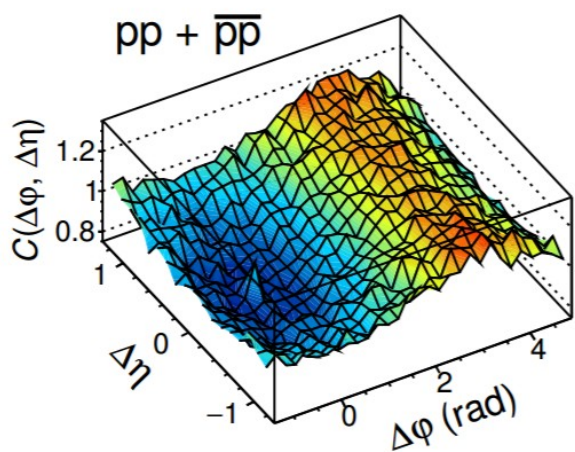
ALICE, PRC 99, 024001 (2019)



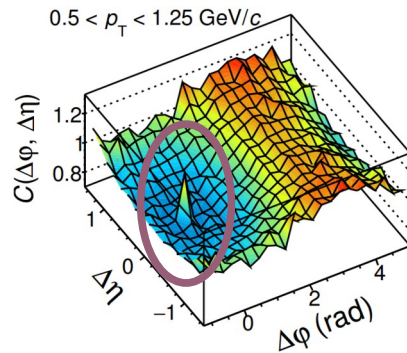
correlation strength



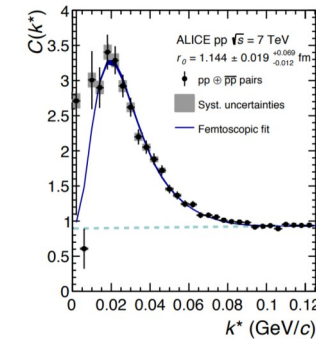
→ correlation **weakens** from pp to $\Lambda\Lambda$ pairs, same as the small peak in angular correlations



- **Direct transformation from $C(k^*)$ to $C(\Delta\eta, \Delta\varphi)$ is not possible**



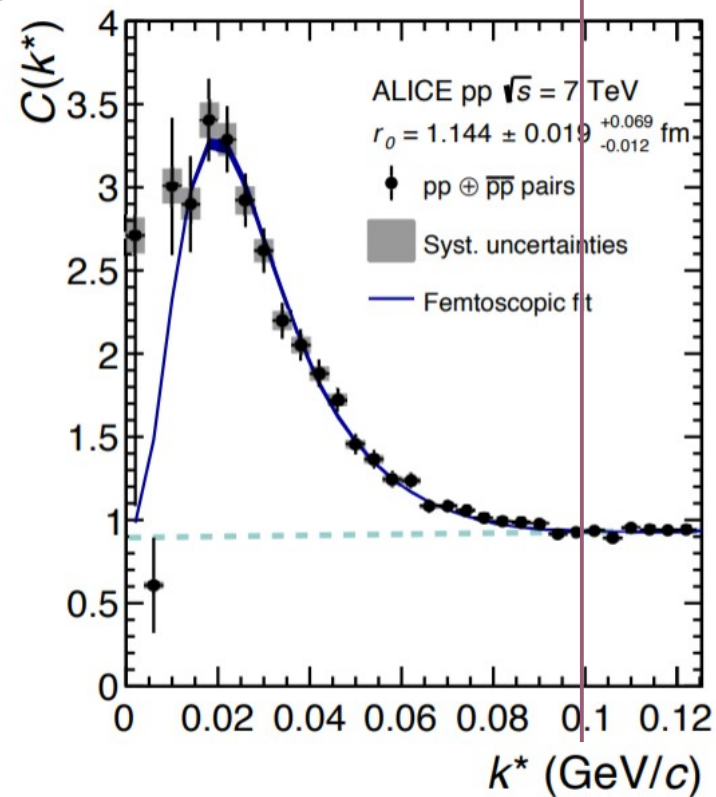
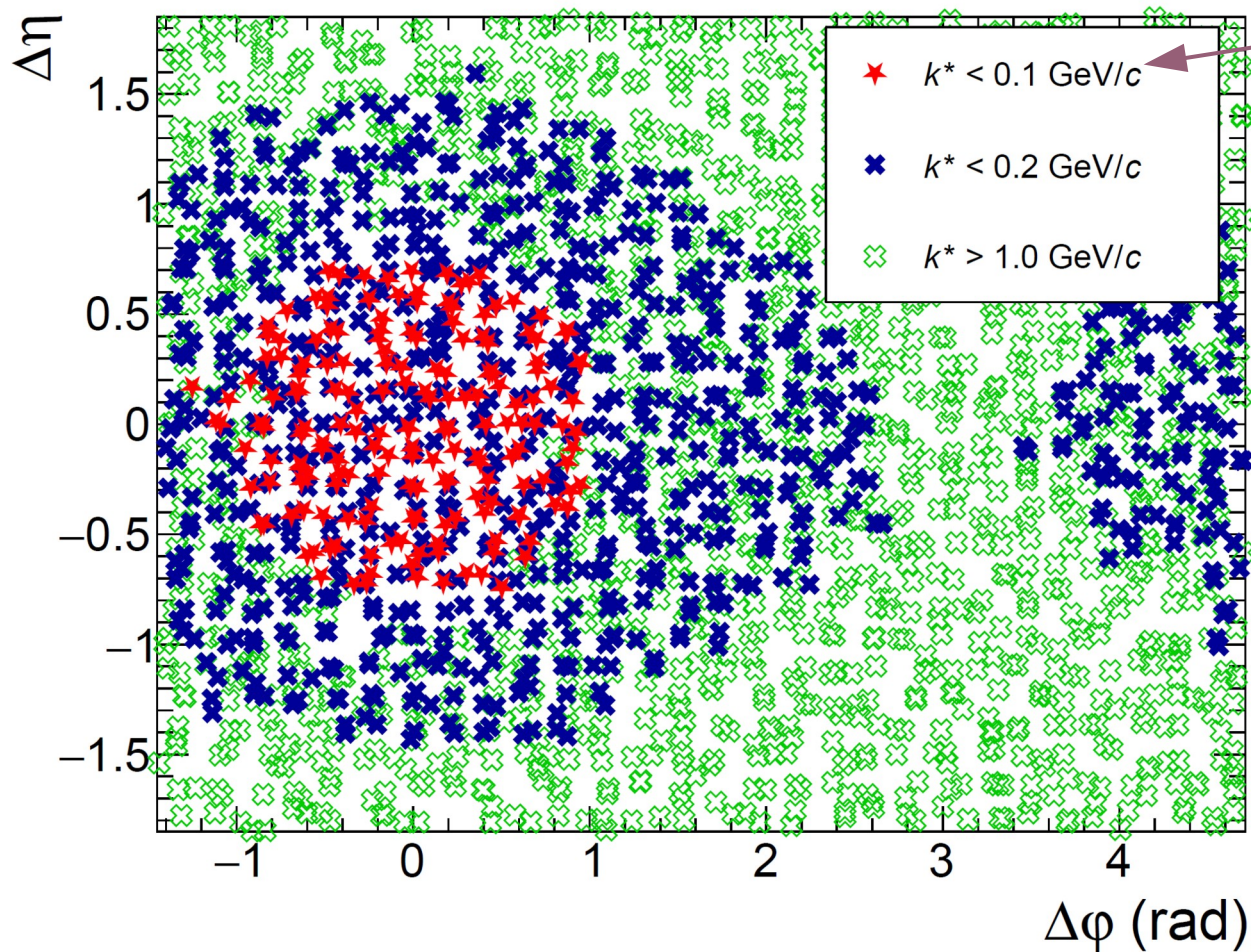
unfolding?



- **We propose a very simple Monte Carlo algorithm to unfold the angular correlation from the femtoscopic one**
 - we tested the method with PYTHIA 8 simulations coupled to Lednicky-Lyuboshitz (L-L) formalism for QS and FSI effects
 - we show how the effects of FSI and QS manifest in angular correlations

- **Femtoscopic region** (small k^*) translates directly to the near-side region (0,0) in the angular correlation

→ QS+FSI effects should be possible to be quite precisely unfolded from the femtoscopic correlation function



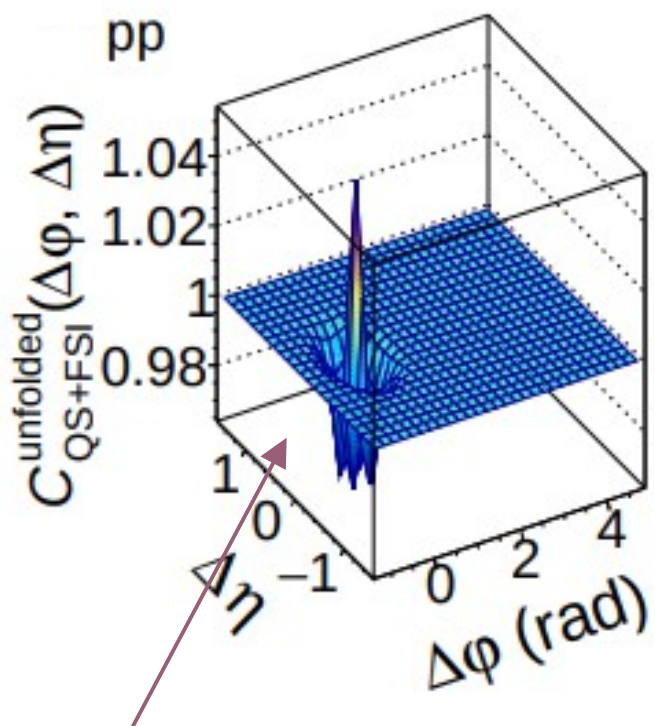


The unfolding of the QS+FSI works very well

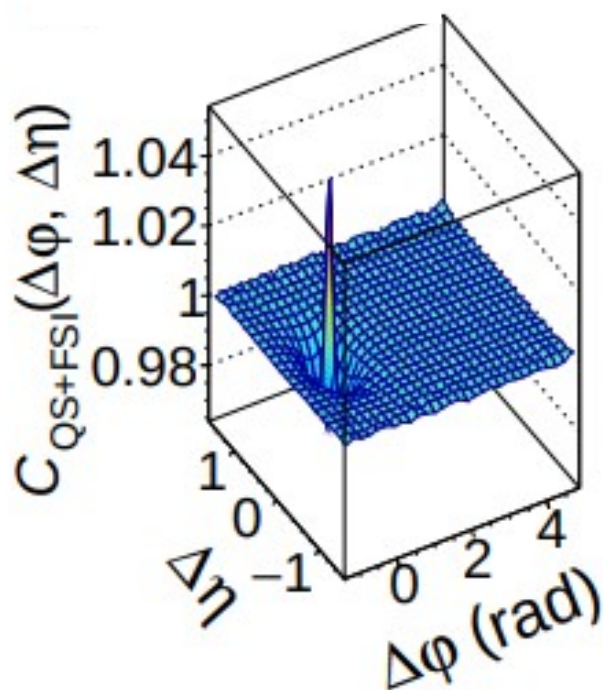
Ł.G. & M.J., PRC 104, 054909 (2021)

Protons

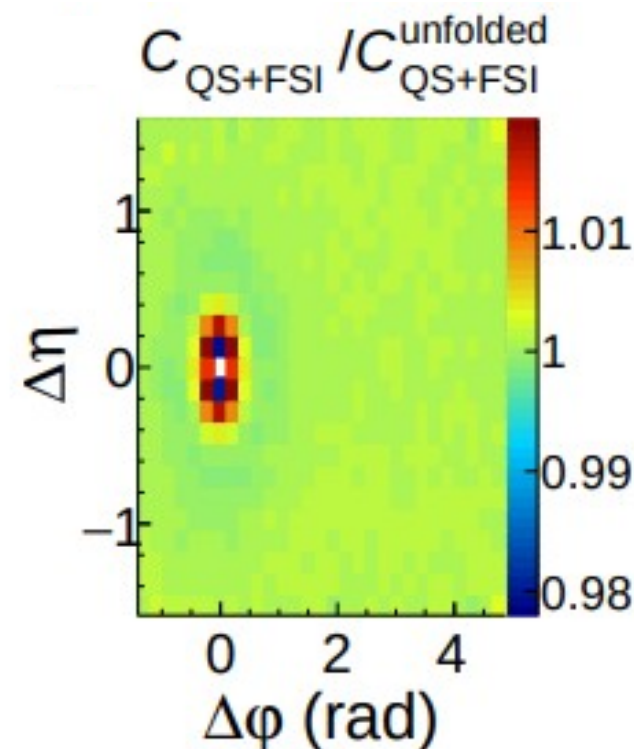
UNFOLDED FROM FEMTO CF
QS+FSI



QS+FSI ONLY
SIMULATED



RATIO



narrow dip + spike in the middle



Correlations of baryons reveal an interesting anticorrelation effect:

- Present also in 13 TeV pp ALICE data and in Au-Au collisions at various energies from STAR BES
- Interesting theoretical developments for AMPT and PYTHIA → are we on a good path to solving the puzzle? Is it a signature of a partonic matter in small systems?

Clear connection between femtoscopic and angular correlations:

- The small peak in pp correlations and the dip in pp proved to come from the strong FSI
- Femtoscopic correlations can be used to unfold the effects of QS+FSI in angular correlations, especially for pairs where MC models do not work (baryons)



THANK YOU!



BACKUP

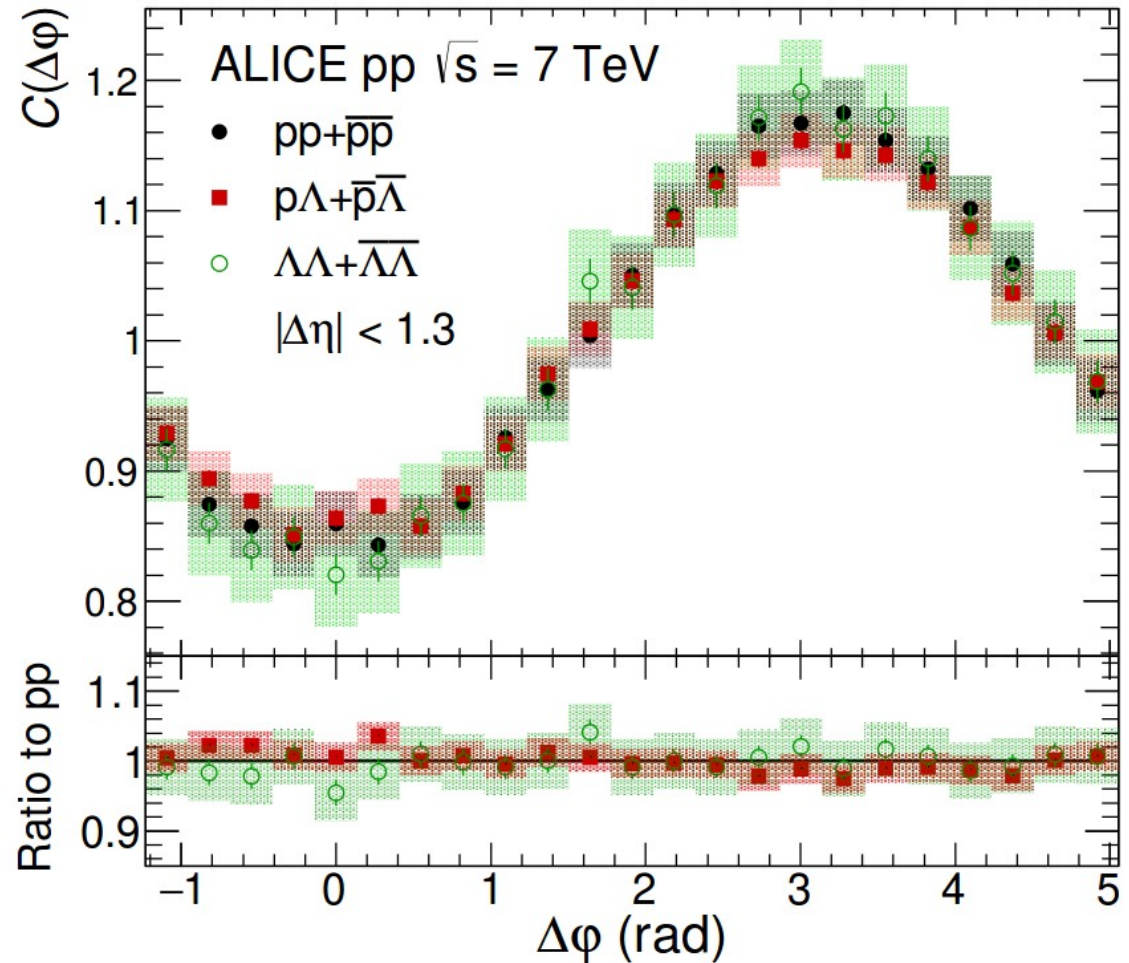


$\Lambda\Lambda$ and $p\Lambda$ correlation functions



ALICE, Eur. Phys. J. C 77 (2017) 569

- Useful to check if effect persists for other baryons than protons – is this a common effect for all baryons?
- Correlation functions were calculated for $\Lambda\Lambda$ and $p\Lambda$ pairs
- Λ baryons are neutral \rightarrow no Coulomb repulsion
- p and Λ are not identical \rightarrow no effect from Fermi-Dirac statistics
- All observations from pp can be extended to $\Lambda\Lambda$ and $p\Lambda$





A multiphase transport (AMPT) model

Default: Lin, Pal, Zhang, Li & Ko, PRC 61, 067901 (00); 64, 041901 (01);
72, 064901 (05); <http://www-cunuke.phys.columbia.edu/OSCAR>

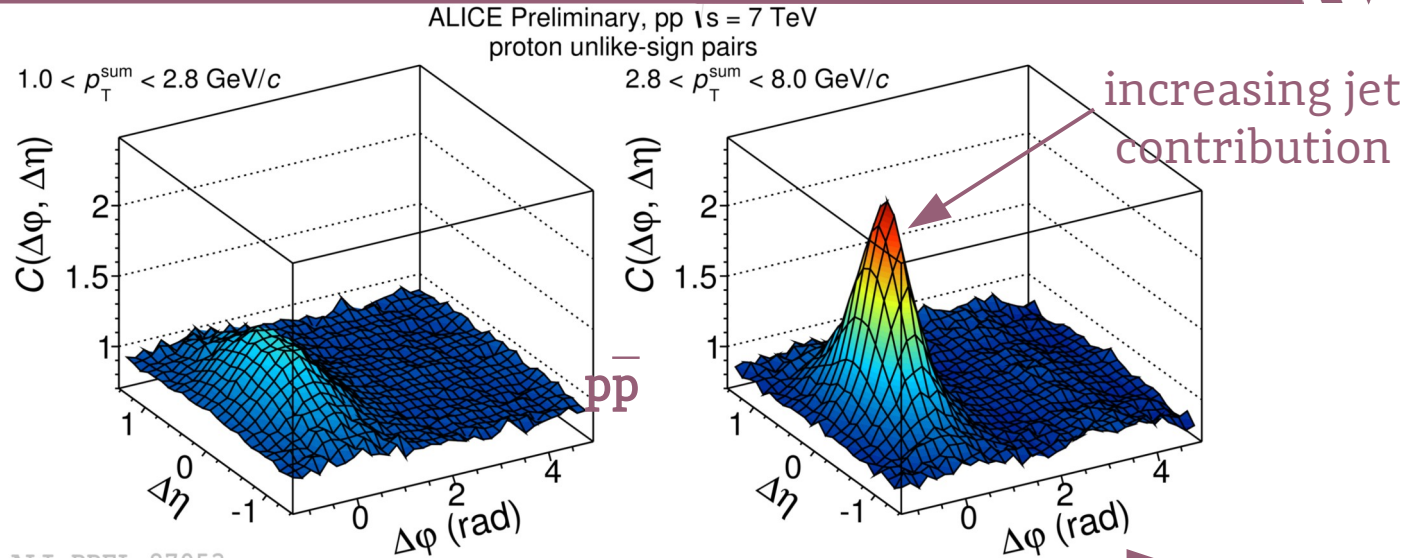
- Initial conditions: HIJING (soft strings and hard minijets)
- Parton evolution: ZPC
- Hadronization: Lund string model for default AMPT
- Hadronic scattering: ART

String melting: PRC 65, 034904 (02); PRL 89, 152301 (02)

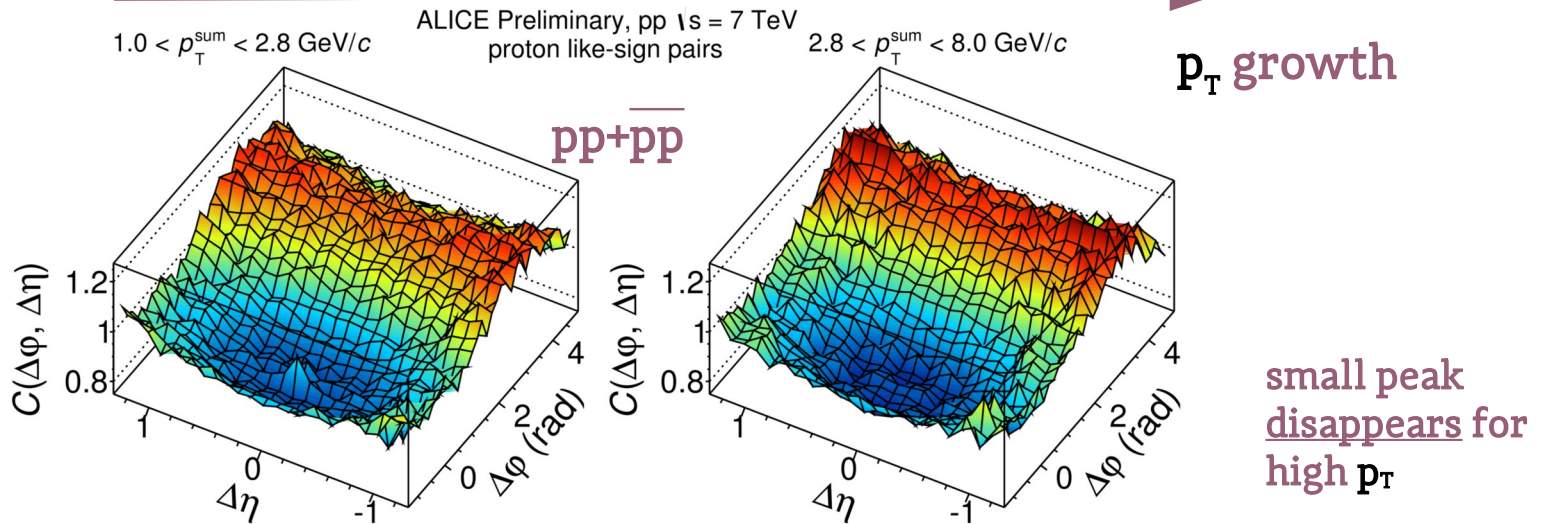
- Convert hadrons from string fragmentation into quarks and antiquarks
- Evolve quarks and antiquarks with ZPC
- When partons stop interacting, combine nearest quark and antiquark to meson, and nearest three quarks to baryon (coordinate-space coalescence)
- Hadron flavors are determined by the invariant mass of quarks



Near-side peak grows with p_T (more contribution from jets)



ALI-PREL-87053



ALI-PREL-87049

- The small peak seems to behave **strangely** → decreases with increasing p_T
- Is it an unnoticed and not removed **detector effect** OR is there some **physics** behind it?



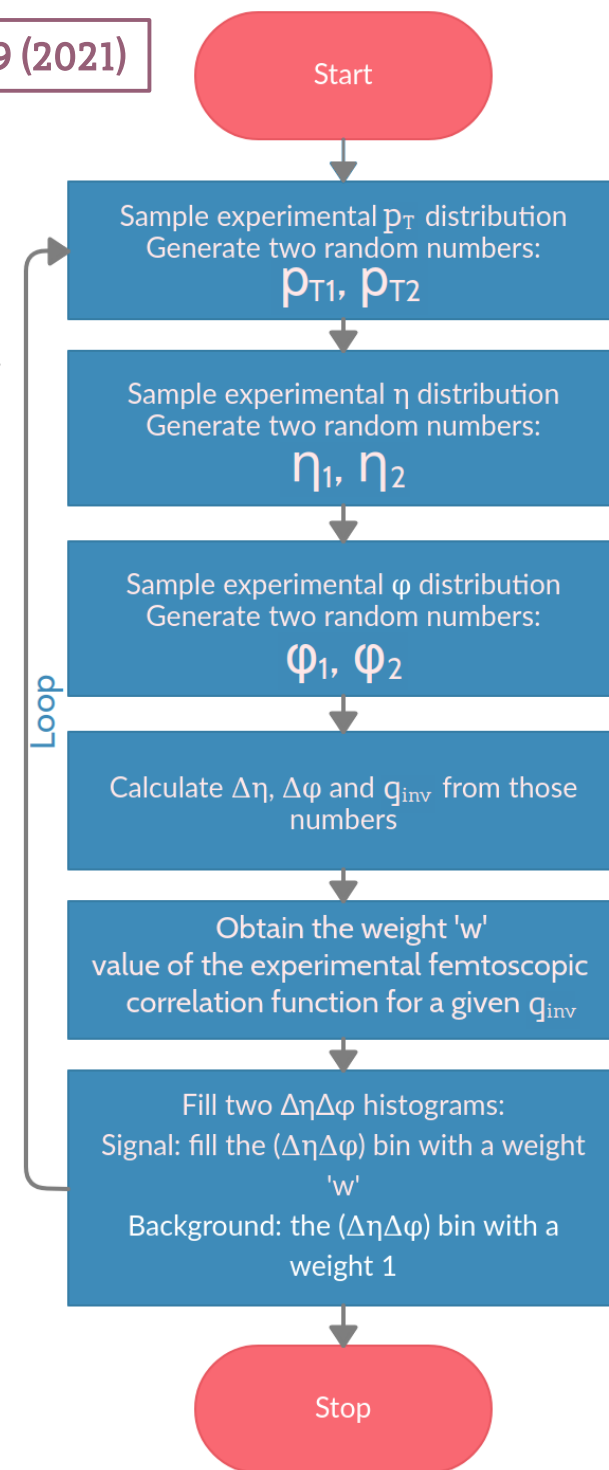
Ł.G. & M.J., PRC 104, 054909 (2021)

How does the unfolding work?

- we sample (twice) single-particle kinematic distributions (p_T, η, φ)
- for each iteration we calculate q_{inv} (or k^*) from those randomly sampled quantities
- we obtain the weight 'w' for a given q_{inv} (or k^*)
 - value of the femtosopic correlation
- then, we calculate $\Delta\eta$ and $\Delta\varphi$ and fill two histograms
 - signal with the weight 'w'
 - background, with weight = 1

By definition, such simple procedure will work **ONLY** for those effects to which the femtosopic CF is sensitive the most

It will **NOT** work for long-range effects (i.e. jets, momentum conservation)



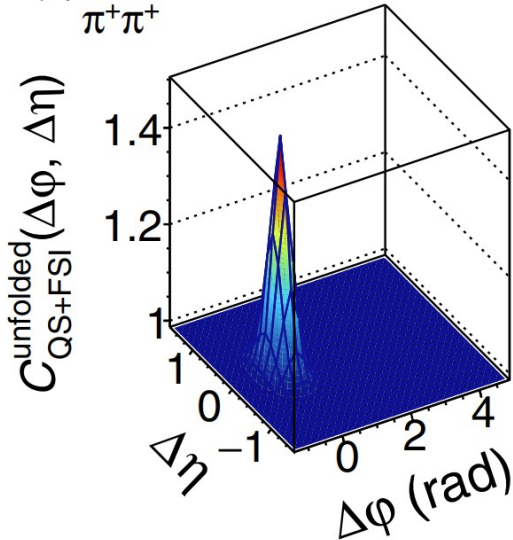
The unfolding of the QS+FSI works very well

Ł.G. & M.J., PRC 104, 054909 (2021)

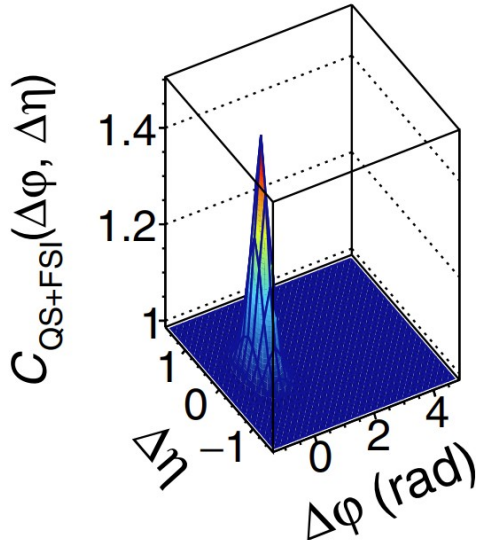
Lednický-Lyuboshitz model coupled to PYTHIA 8

Pions

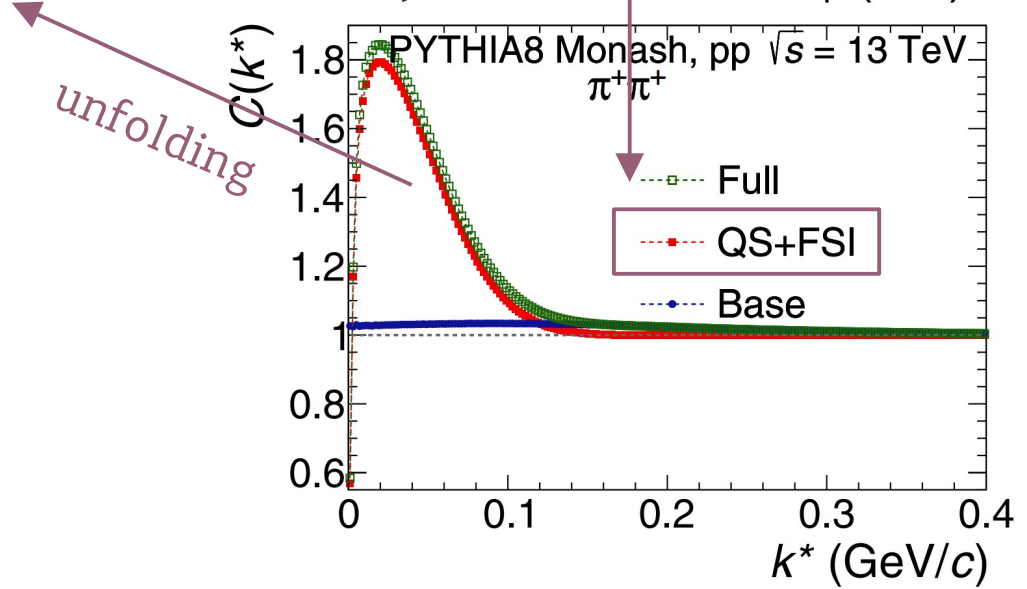
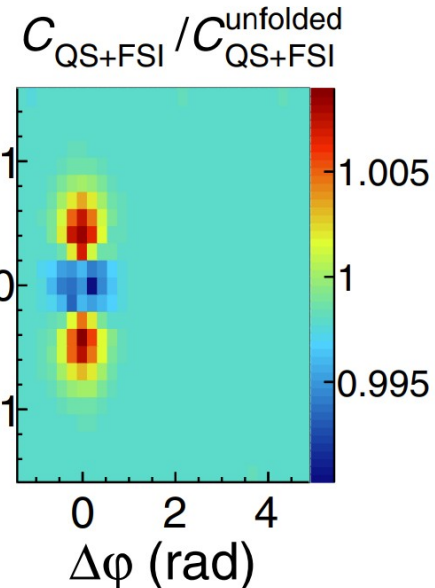
UNFOLDED FROM FEMTO CF
QS+FSI



QS+FSI ONLY
SIMULATED

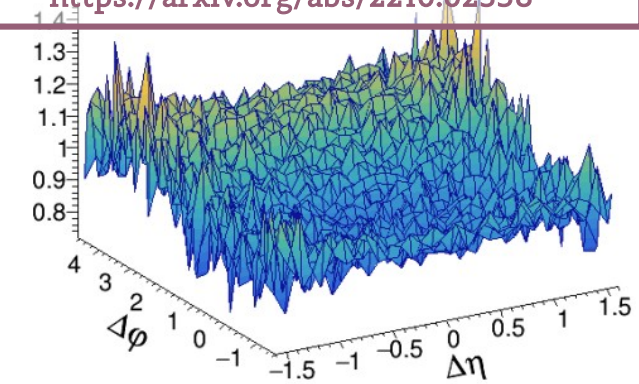
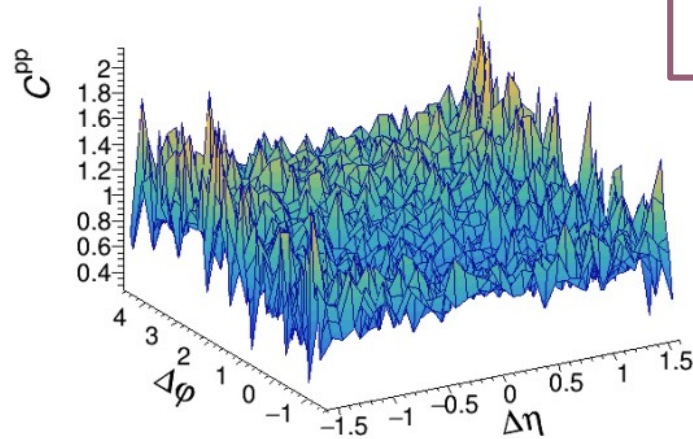
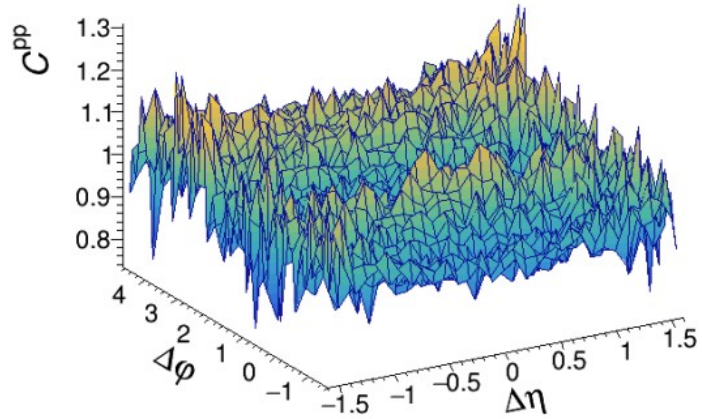


RATIO

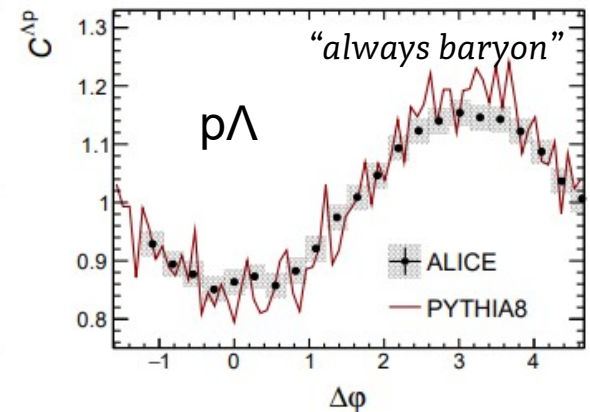
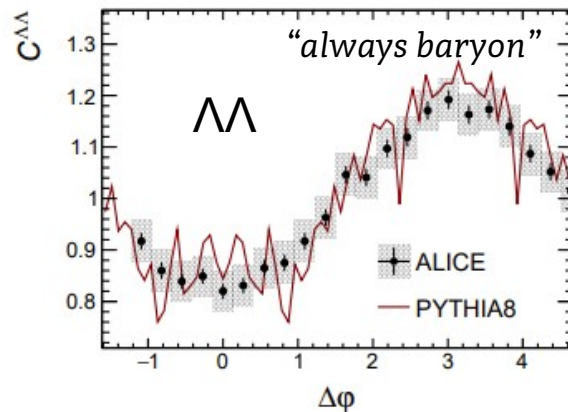
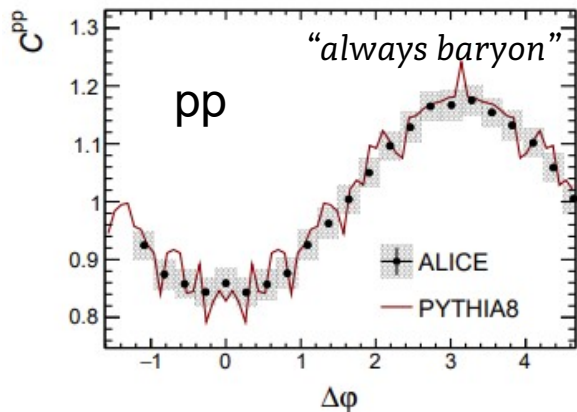


unfolding

N. Demazure, V. Gonzalez, F. Llanes-Estrada
<https://arxiv.org/abs/2210.02358>



(a) pp C correlation with unmodified PYTHIA (b) pp C correlation, one-baryon per string policy (c) pp C correlation, always-baryon policy



The LEP baryon correlation data could be reasonably fit by PYTHIA as is, given that the color string did form linking a back-to-back primary quark-antiquark pair; this means that baryons from the same string did not form positive correlations near $\Delta\eta \simeq 0 \simeq \Delta\phi$ in OPAL data, as they were somewhat randomized, with the string frame not too far from the laboratory frame.

At the LHC strings are however formed at various rapidities and azimuths, with a natal Lorentz boost. Because of that string boost, two baryons formed from the same string will create that positive correlation in the laboratory frame. Therefore, to avoid it and bring about the anticorrelation seen in the data, two-baryon production from the same string should be suppressed: our way of achieving it is the very rough pair of policies (one-baryon and all-baryon) that certainly need to be improved in future work.

- **Improved coalescence (removed separate conservation for mesons and baryons)**

• **String** A physics picture for the near-side depression feature emerges out of these

→ **AM** results: a parton matter is created and then expands to a finite volume, then

→ **tes** the hadronization process (quark coalescence in the AMPT model for this

• **If init** produces multiple hadrons relatively close in the coordinate space including

to the spatial azimuth. The finite expansion of the parton matter creates a finite

1.6 space-momentum correlation (often called the radial flow), and the

1.4 transverse flow velocity of a local parton volume tends to change its hadrons

from being close in spatial azimuth to being close in momentum azimuth.

Results from the AMPT-SM model with a finite parton cross section are

0.8 qualitatively consistent with the experimental data. In addition, the PYTHIA

2.5 model with string fragmentation alone, the AMPT-SM model with no parton

2 expansion, and the AMPT-def model that is dominated by the hadron cascade

1.5 all fail to produce the depression feature. Therefore, our study implies a finite

parton expansion before hadronization, where a prerequisite is the existence

0.5 of parton degrees of freedom, in pp collisions $\sqrt{s} = 7$ TeV.

$C(\Delta\phi)$

$C(\Delta\phi)$

0

2

4

0

2

4

0

2

4

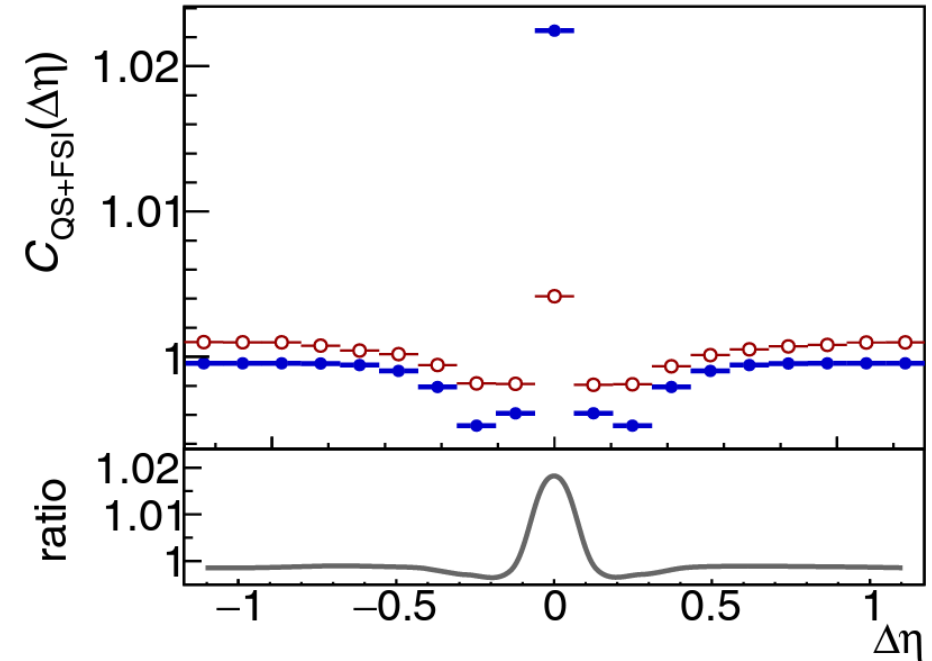
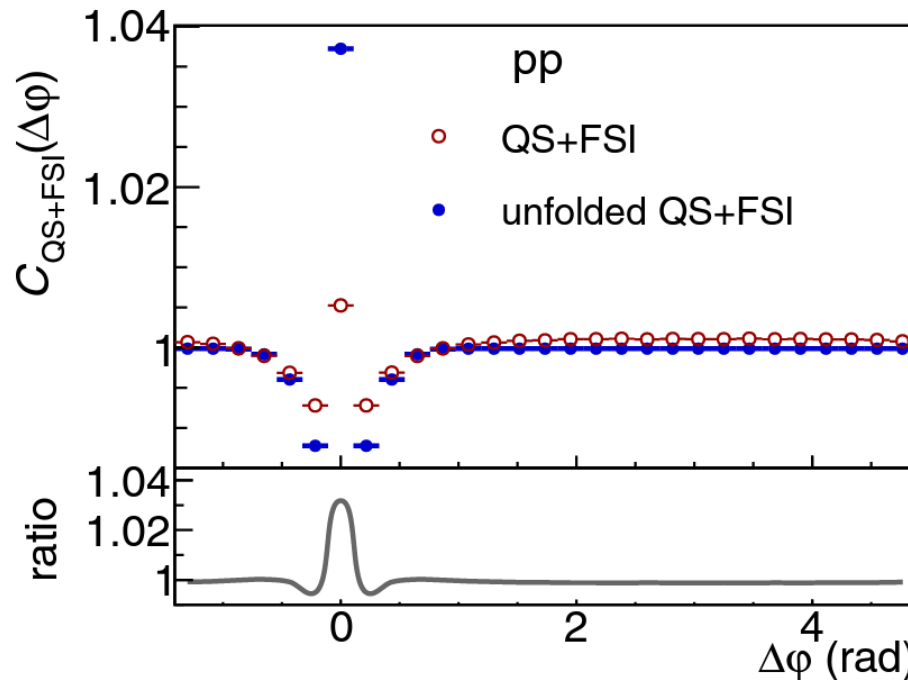
$\Delta\phi$ (rad)

$\Delta\phi$ (rad)

$\Delta\phi$ (rad)

34912
37063

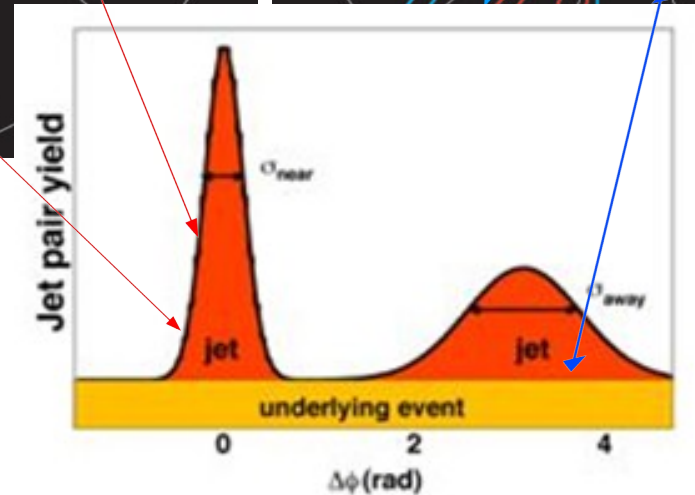
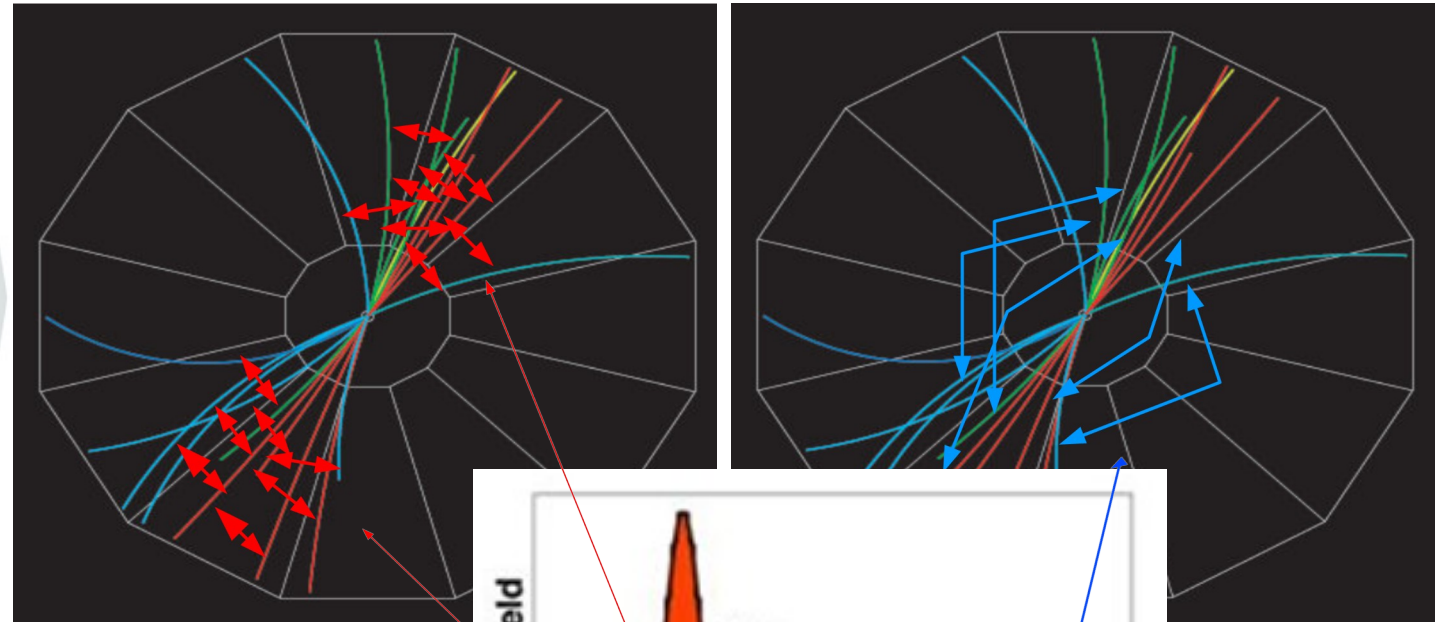
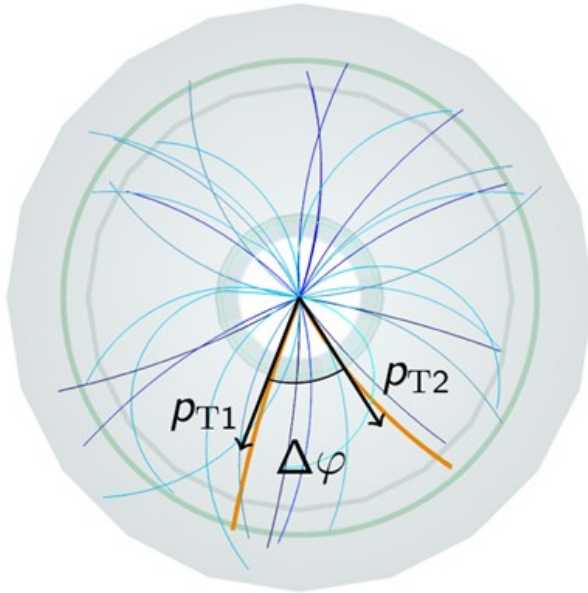
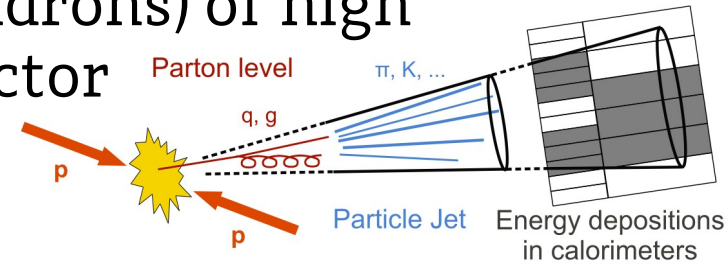




a depletion in the correlation function around $(\Delta\eta, \Delta\phi) \approx (0, 0)$ is visible with an additional peak structure directly at $(\Delta\eta, \Delta\phi) = (0, 0)$. Although the magnitude of the peak is substantially smaller in the unfolded correlation, the procedure is able to describe the shape qualitatively. A qualitatively similar peak structure, located at $(\Delta\eta, \Delta\phi) = (0, 0)$, in the middle of the depletion, was observed experimentally by the ALICE Collaboration and postulated to result from the strong two-proton interaction. This paper validates this ansatz.



- “Jet” is a collimated stream of particles (hadrons) of high momentum (energy) which reach the detector
- How to experimentally measure jets?
- We can look at the collision in the transverse plane and calculate azimuthal angle difference distribution:

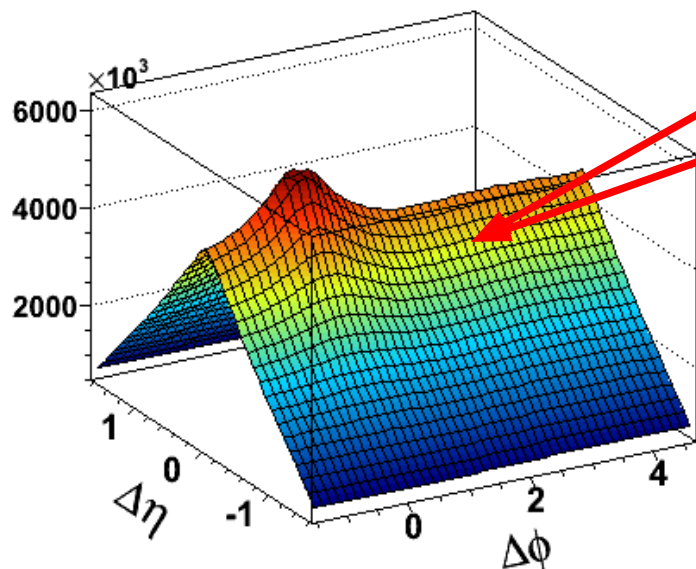


p_T - transverse momentum;
 φ - azimuthal angle;

Signal distribution

$$S(\Delta\eta, \Delta\phi) = \frac{d^2 N^{signal}}{d\Delta\eta d\Delta\phi}$$

Same event pairs



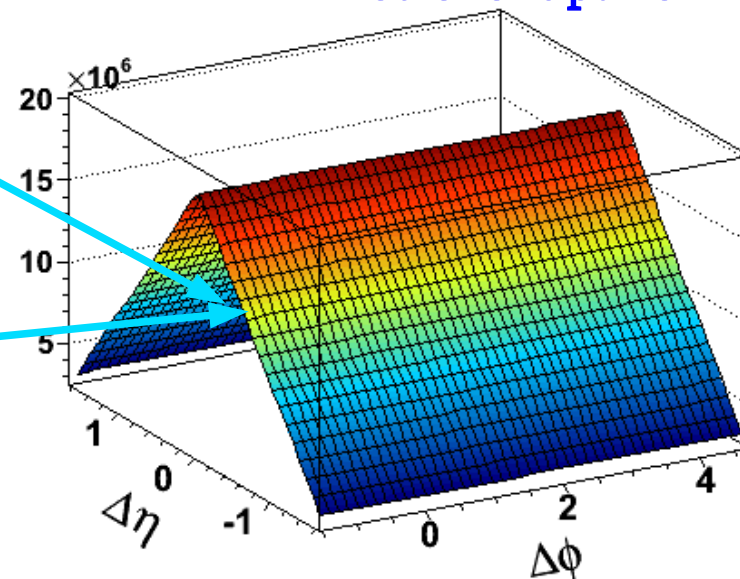
$$\Delta\eta = \eta_1 - \eta_2$$

$$\Delta\phi = \phi_1 - \phi_2$$

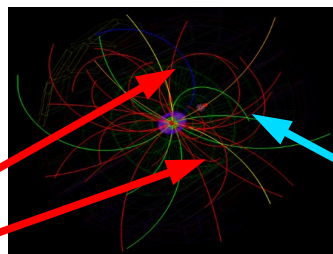
Uncorrelated reference

$$B(\Delta\eta, \Delta\phi) = \frac{d^2 N^{mixed}}{d\Delta\eta d\Delta\phi}$$

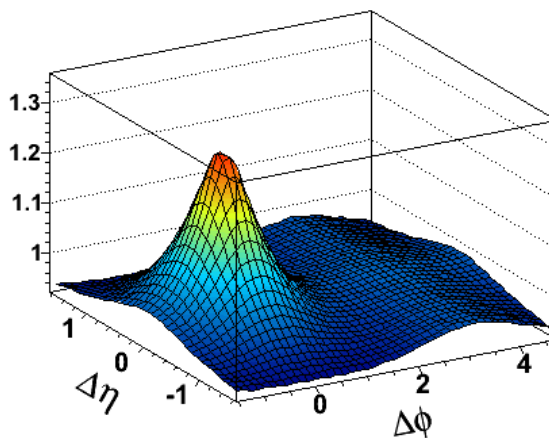
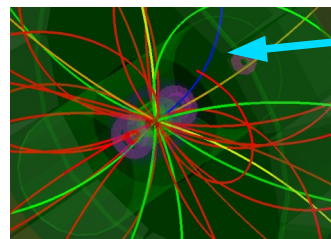
Mixed event pairs



Event 1



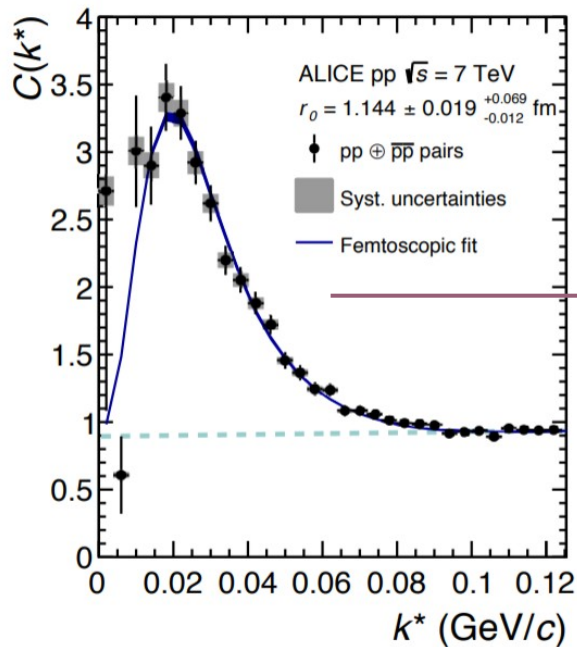
Event 2



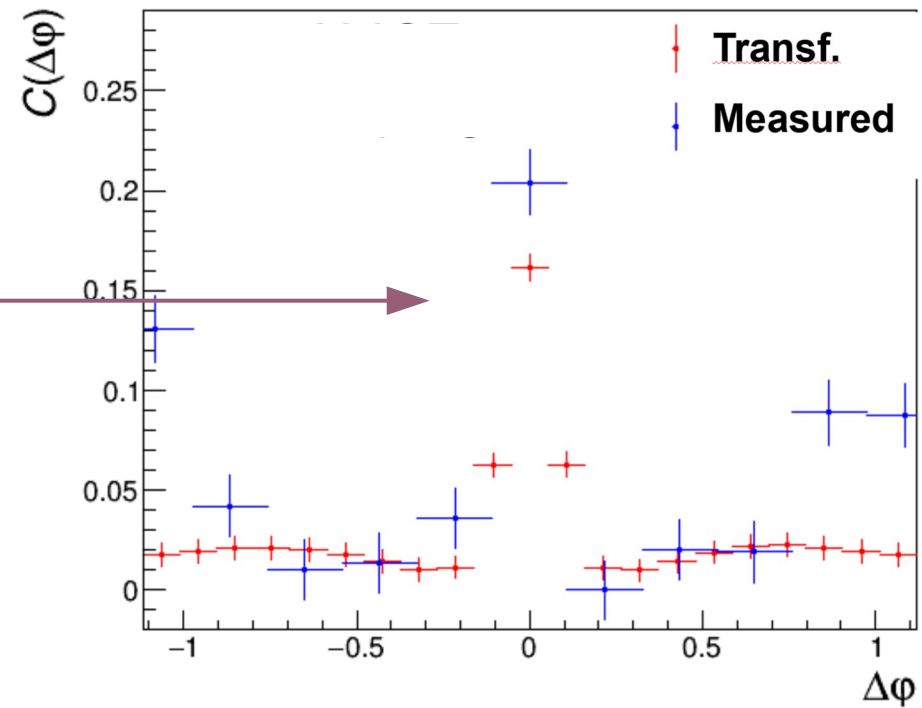
$$C(\Delta\eta, \Delta\phi) = \frac{N_{pairs}^{mixed}}{N_{pairs}^{signal}} \frac{S(\Delta\eta, \Delta\phi)}{B(\Delta\eta, \Delta\phi)}$$

Probability ratio

$$C(\mathbf{p}_1, \mathbf{p}_2) = \frac{P_{1,2}(\mathbf{p}_1, \mathbf{p}_2)}{P_1(\mathbf{p}_1) \cdot P_2(\mathbf{p}_2)}$$



unfolding



- Femto correlation produces spike at $(\Delta\eta, \Delta\phi) = (0, 0)$
- Comparison of two peaks: 1-bin wide projection on $\Delta\phi$ (subtract minimum)
- **Both the height and the width of two peaks are comparable!**



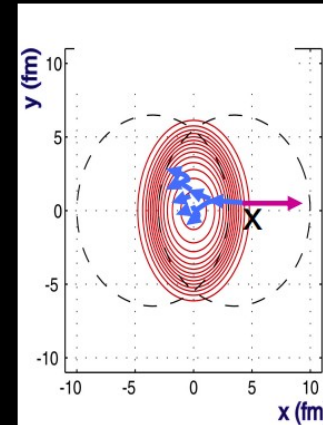
History – High Momentum Particle & Jet Correlations

FERMILAB-Pub-82/59-THY
August, 1982

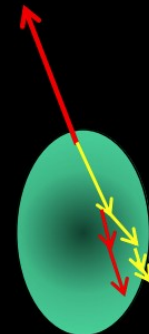
Energy Loss of Energetic Partons in Quark-Gluon Plasma:
Possible Extinction of High p_T Jets in Hadron-Hadron Collisions.

J. D. BJORKEN
Fermi National Accelerator Laboratory
P.O. Box 500, Batavia, Illinois 60510

this effect. An interesting signature may be events in which the hard collision occurs near the edge of the overlap region, with one jet escaping without absorption and the other fully absorbed.



Trigger particle



Away-side particles

Back-to-back Jets Away-side jets NOT quenched in pp collisions

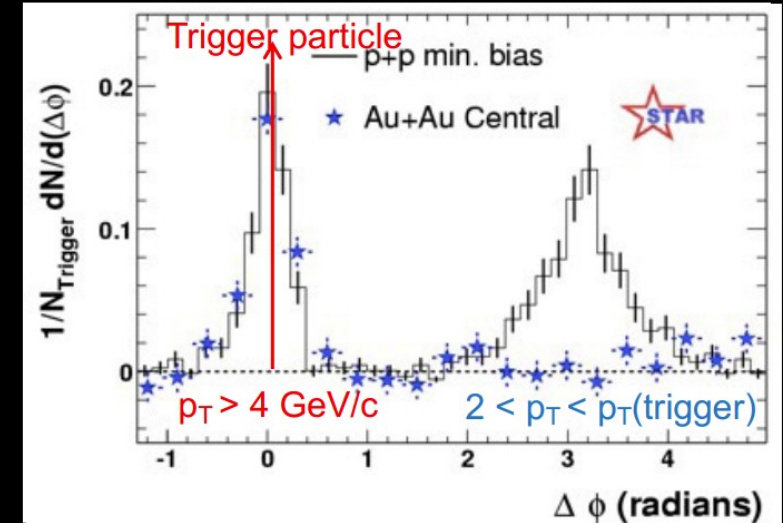
Back-to-back Jets Away-side jets observed as quenched in central Au + Au

Not quenched in Hi Mult d+Au

→ trigger particle origin near surface

→ strongly interacting medium

STAR, Phys.Rev.Lett. 91 (2003) 072304





History -

Energy Loss of Energetic Parton
Possible Extinction of High p_T Jets

J. D. BJORK
Fermi National Accelerator
P.O. Box 500, Batavia, IL

this effect. An interesting signature
collision occurs near the edge of
escaping without absorption and the c

Back-to-back Jets **A**
in pp collisions

Back-to-back Jets **A**
central Au + Au

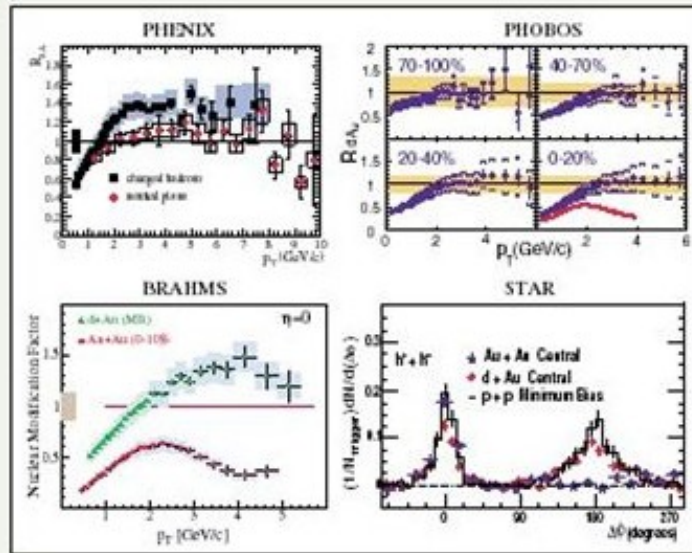
→ trigger particle
→ strongly interact

STAR, Phys.R

John Harris (Yale)

PHYSICAL REVIEW LETTERS

Articles published week ending
15 AUGUST 2003
Volume 91, Number 7



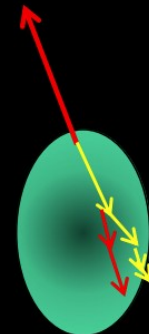
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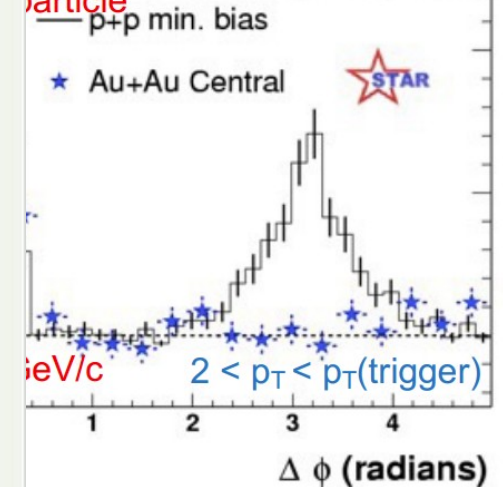
Correlations

trigger
particle



Away-side
particles

particle



Puerto Vallarta, Mexico, January 2003

J. Harris, WWND 2003
<https://indico.cern.ch/event/119>

6/11/2023, WPCF 2023

Łukasz Graczykowski (WUT)

40/25

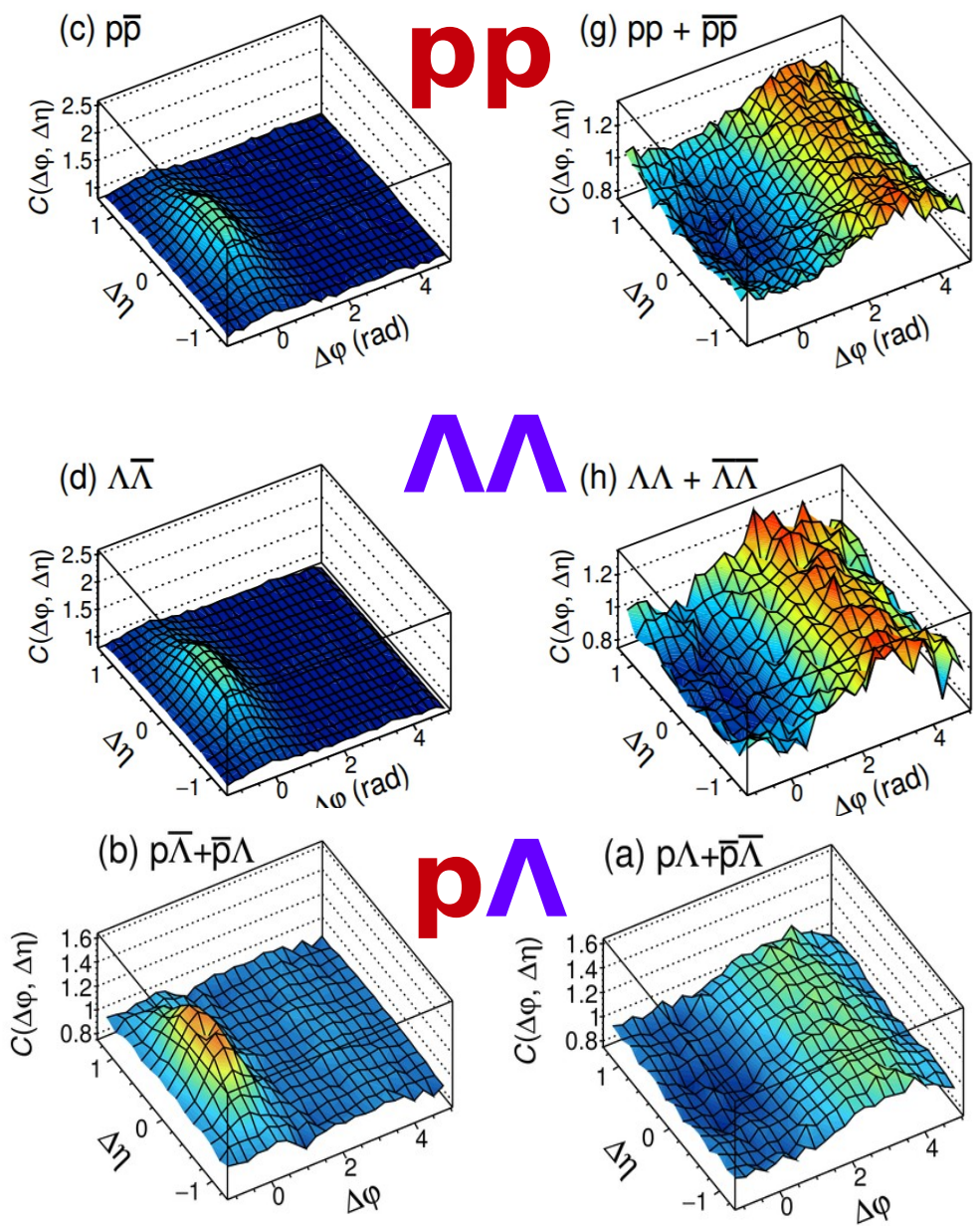
Λ and $p\Lambda$ correlation functions

ALICE, Eur. Phys. J. C 77 (2017) 569

Baryon-Antibaryon

Baryon-Baryon

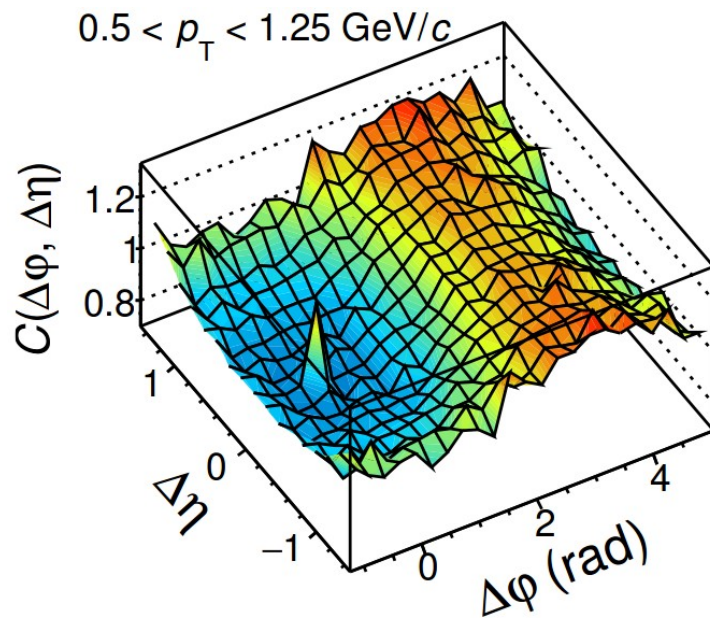
- Useful to check if effect persists for other baryons than protons – is this a common effect for all baryons?
- Correlation functions were calculated for $\Lambda\Lambda$ and $p\Lambda$ pairs
- Λ baryons are neutral \rightarrow no Coulomb repulsion
- p and Λ are not identical \rightarrow no effect from Fermi-Dirac statistics
- All observations from pp can be extended to $\Lambda\Lambda$ and $p\Lambda$



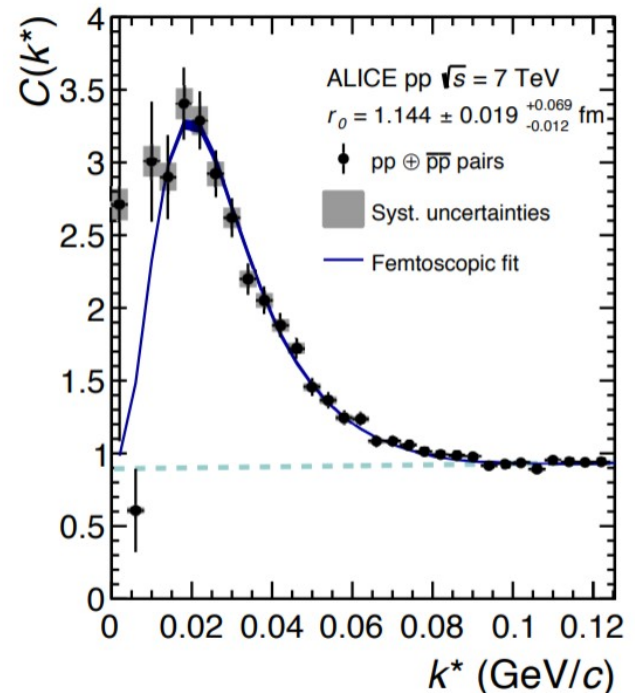


In our new paper we propose a **simple algorithm** to **unfold** the angular correlation from measured femtoscopic one

- we test the method with PYTHIA 8 simulations coupled to Lednicky and Lyuboshitz formalizm
- we show **how** the effects of **strong FSI and QS** manifest in angular correlations



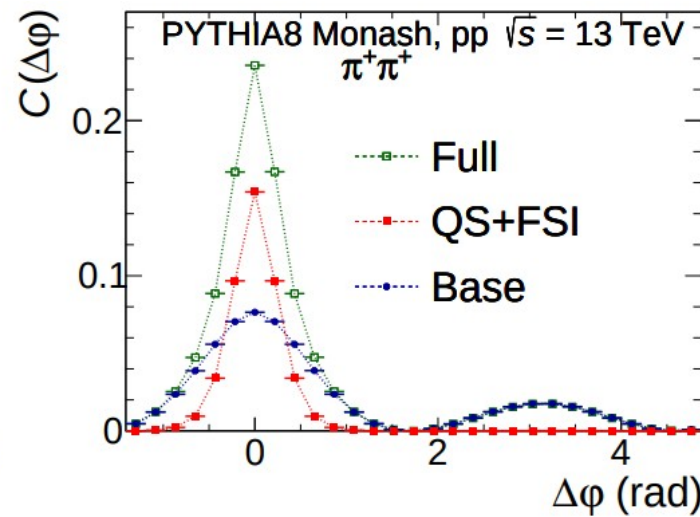
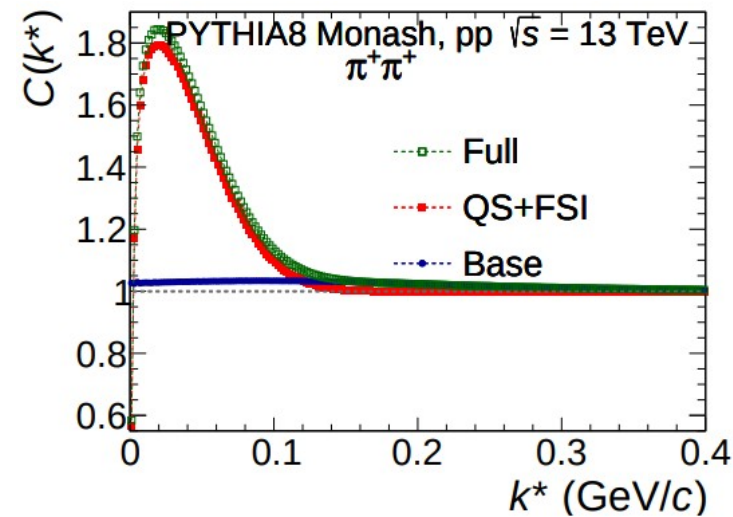
← unfolding



Ł.G. & M.J., PRC 104, 054909 (2021)



Ł.G. & M.J., PRC 104, 054909 (2021)

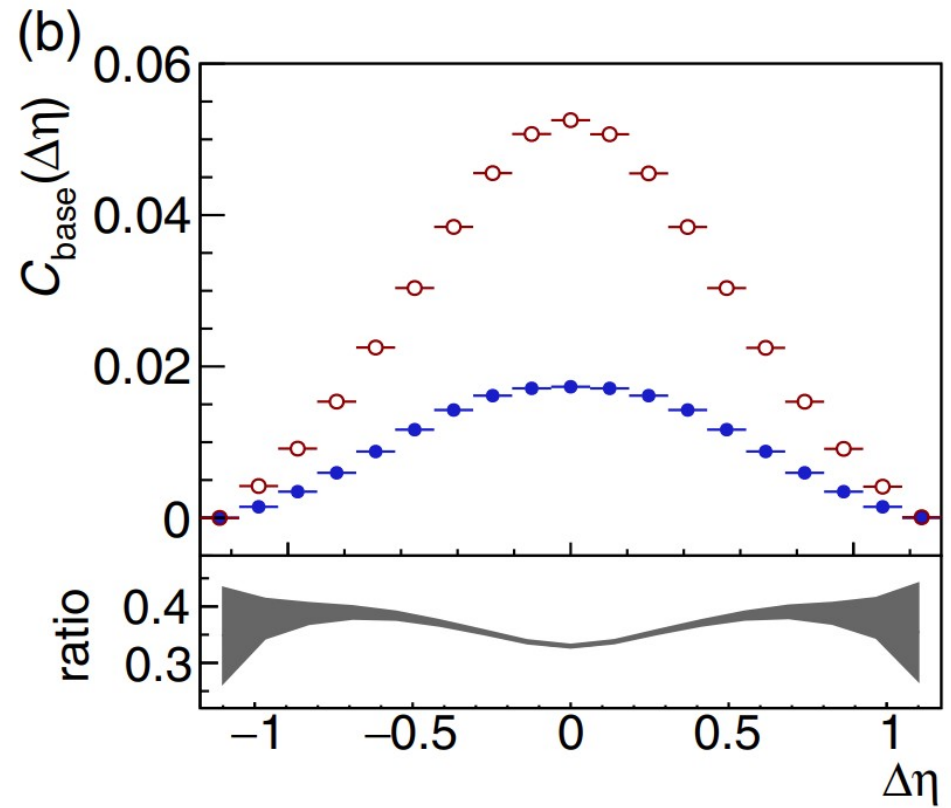
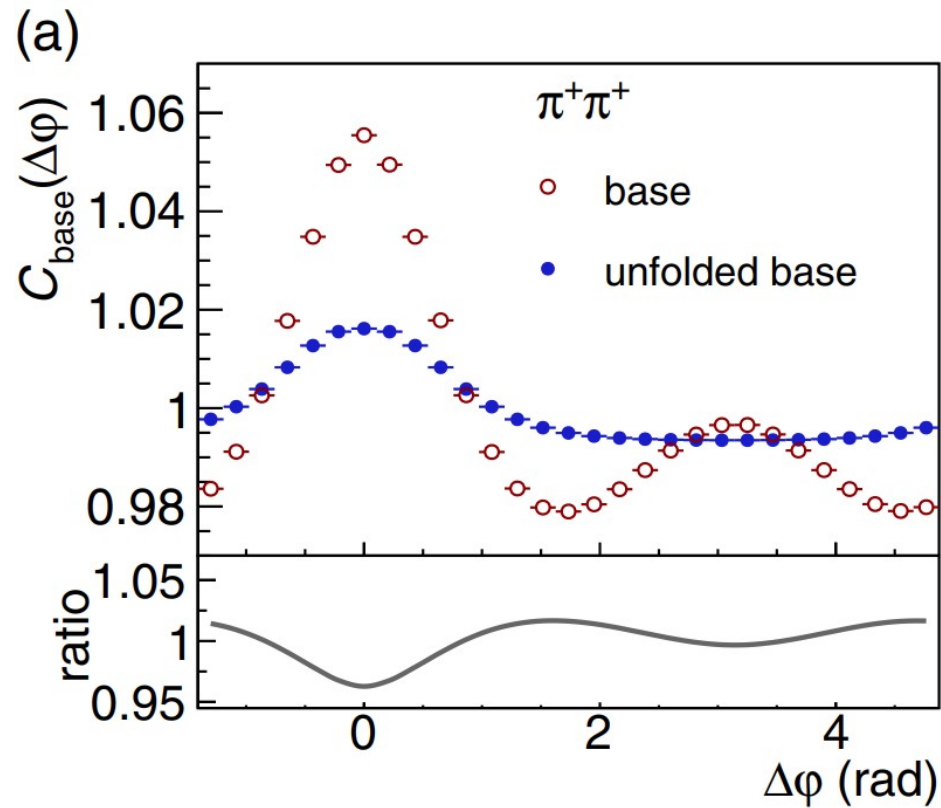


1. $C_{\text{base}} = S/M$, where M is the mixed-event distribution, contains only the event-wide correlations, without the QS and FSI effects added by the afterburner;
2. $C_{\text{full}} = S_w/M$ contains the full information, that is the event-wide correlations with additional effects of QS and FSI added by the afterburner;
3. $C_{\text{QS+FSI}} = M_w/M$ contains only the effects related to QS and FSI and is an equivalent to numerical integration of Eq. (2).

S – same event distribution
 M – mixed event distribution
 w – weight from Lednicky model

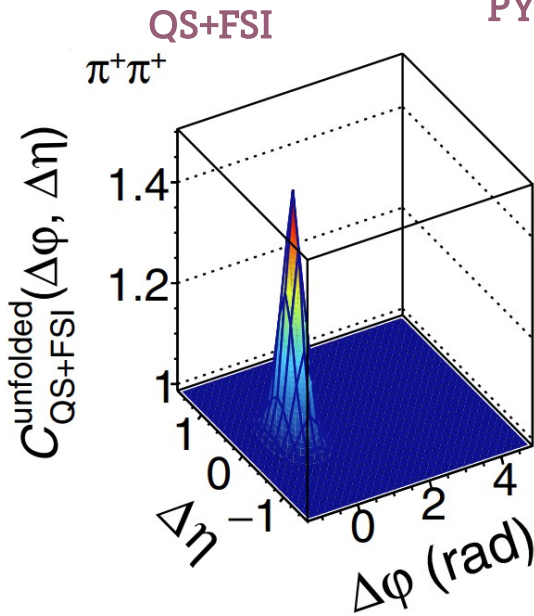


→ the global energy-momentum conservation shape is, obviously, not preserved in unfolded angular CF

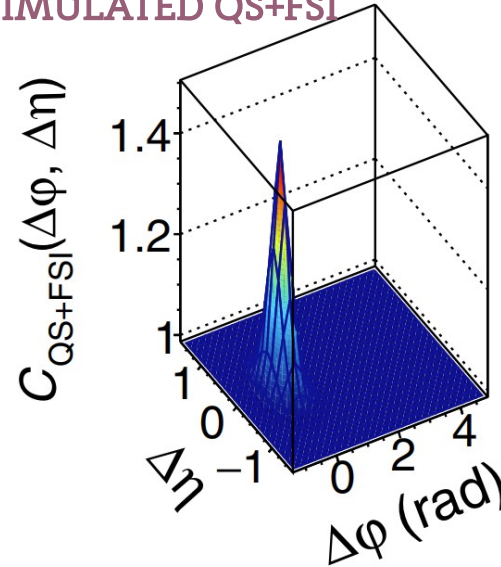


→ unfolding of the QS+FSI correlation, which is limited in k^* , works very well, here example for pions

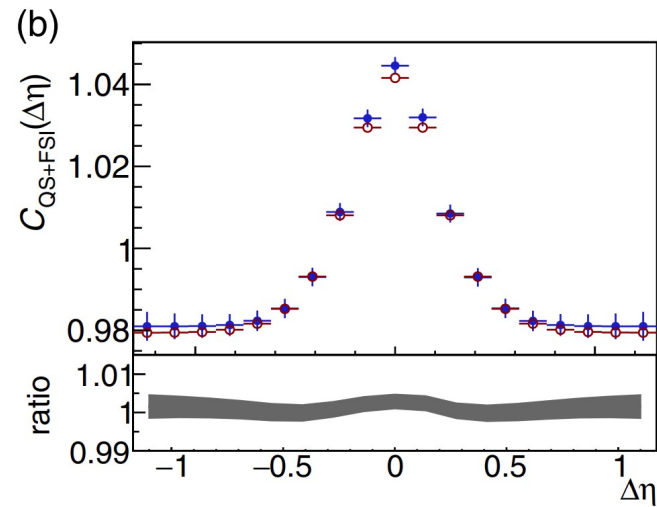
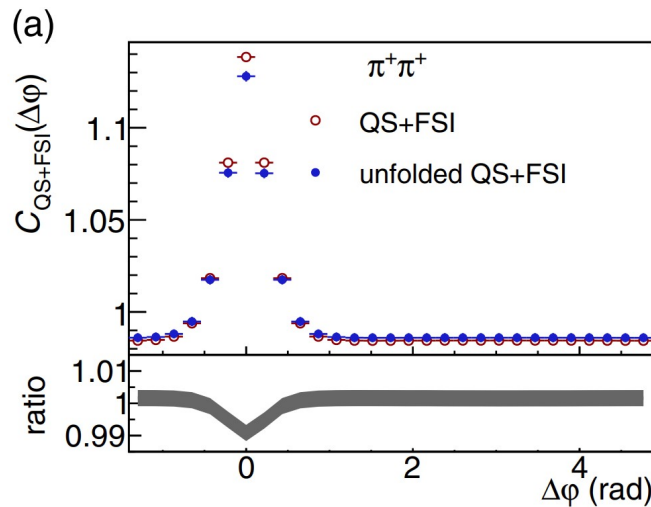
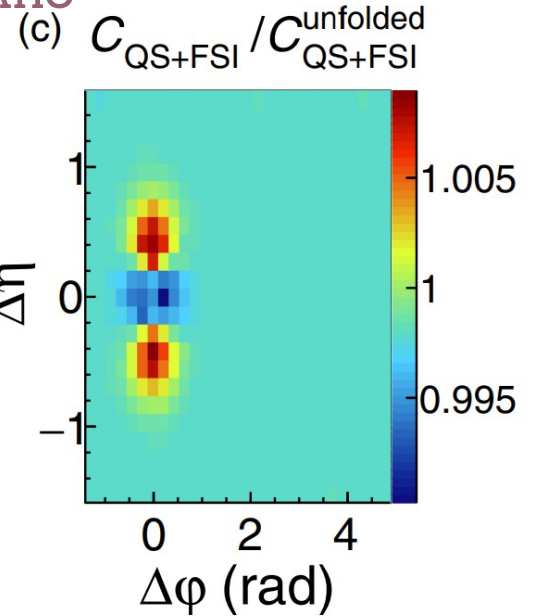
UNFOLDED FROM FEMTO CF



PYTHIA WITH LEDNICKY
SIMULATED QS+FSI



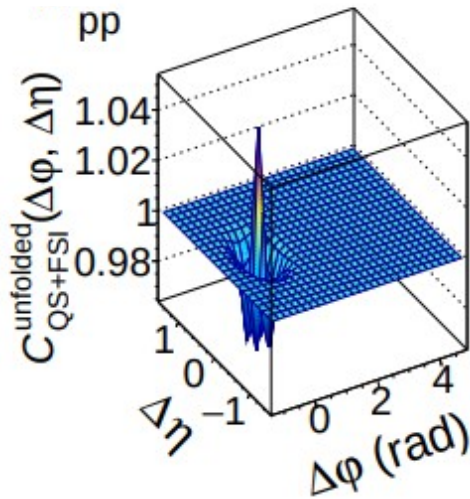
RATIO



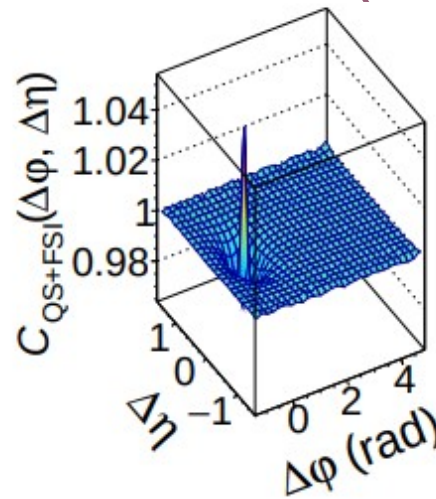


→ in the case of protons, the QS strong FSI is well-preserved and clearly seen as a sharp, narrow peak at (0,0), which proves the ALICE hypothesis

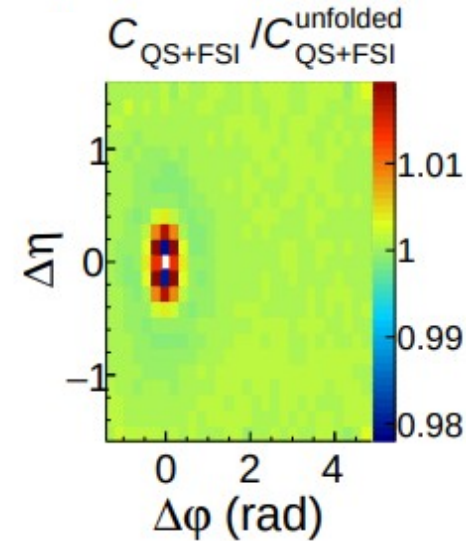
UNFOLDED FROM FEMTO CF
QS+FSI



PYTHIA WITH LEDNICKY
SIMULATED QS+FSI



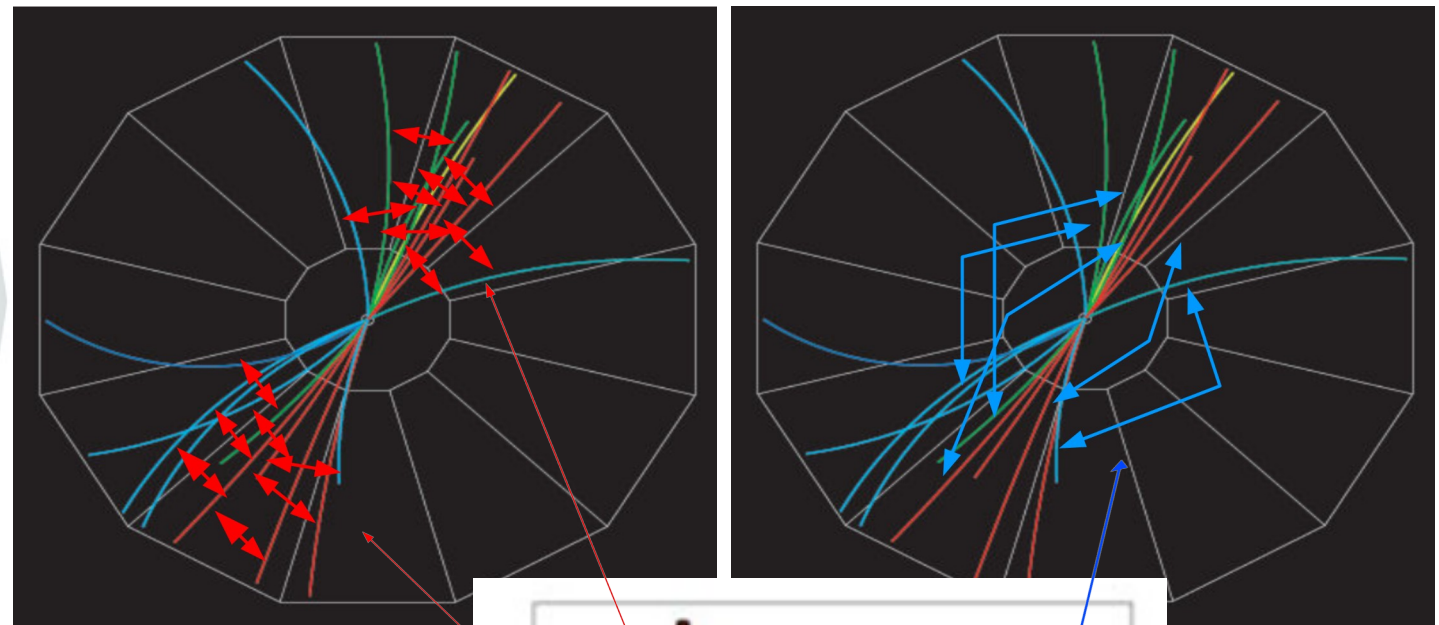
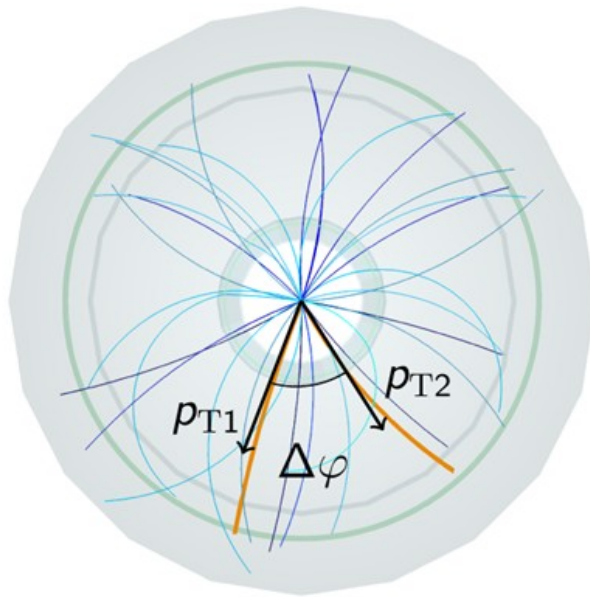
RATIO



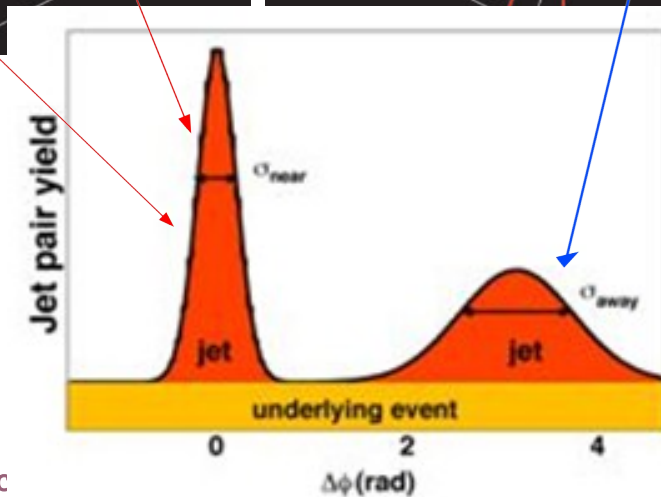
→ weaker femto CF for $p\Lambda$ and $\Lambda\Lambda$ pairs (weaker contribution from strong FSI) → less prominent “small peaks” in angular CF



- How to experimentally measure jets?
- We can look at the collision in the transverse plane and calculate azimuthal angle difference distribution:



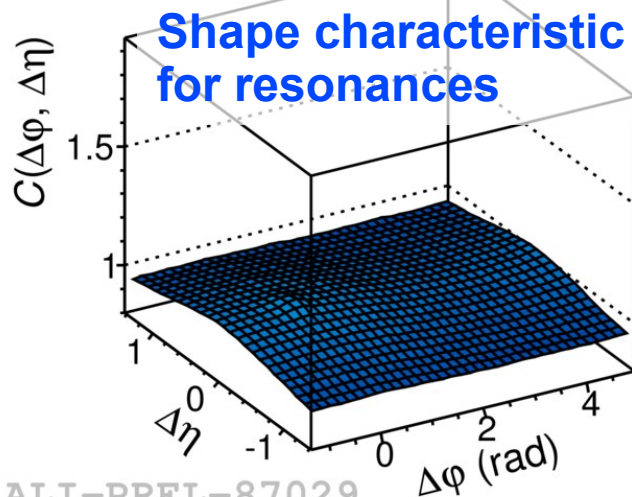
p_T - transverse momentum;
 φ - azimuthal angle;



ALICE 7 TeV pp data - pions



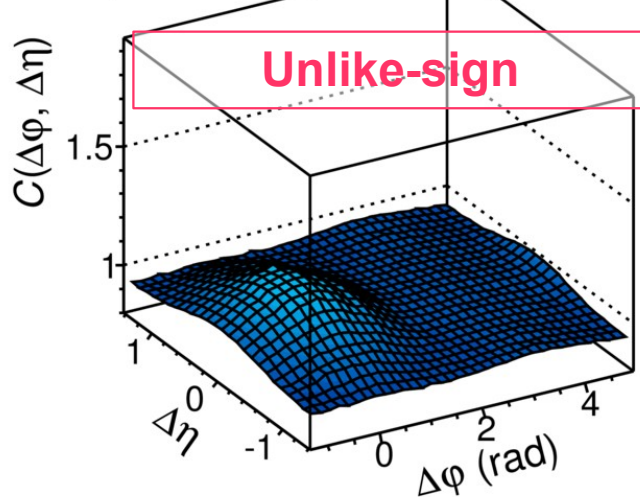
$0.4 < p_T^{\text{sum}} < 0.8 \text{ GeV}/c$



ALI-PREL-87029

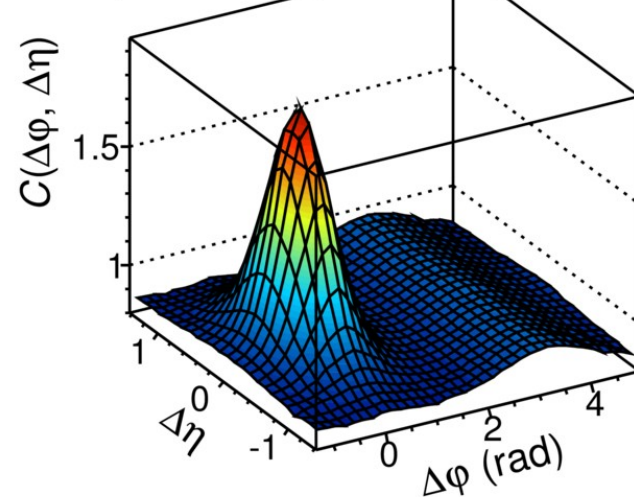
ALICE Preliminary, pp $\sqrt{s} = 7 \text{ TeV}$
pion unlike-sign pairs

$1.2 < p_T^{\text{sum}} < 1.4 \text{ GeV}/c$



$$p_{T\text{sum}} = |\vec{p}_{T1}| + |\vec{p}_{T2}|$$

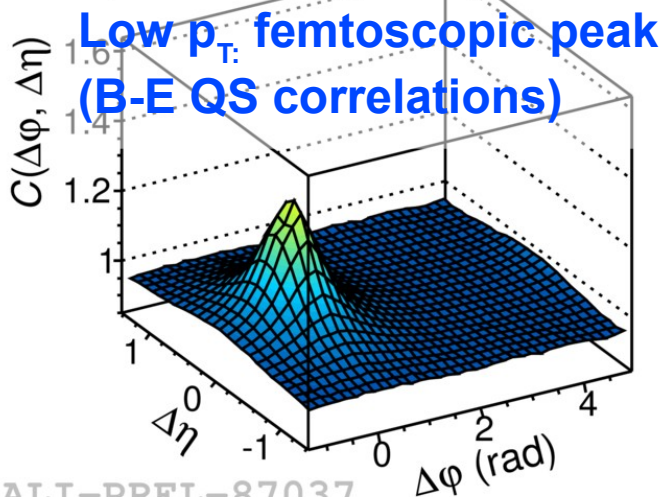
$2.5 < p_T^{\text{sum}} < 8.0 \text{ GeV}/c$



High p_T : jets

p_T growth

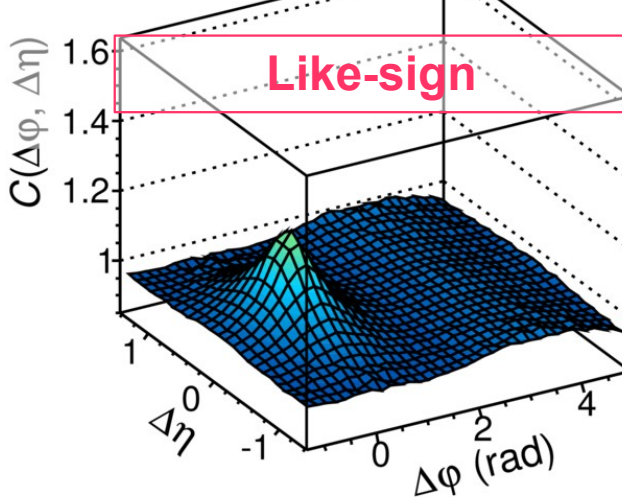
$0.4 < p_T^{\text{sum}} < 0.8 \text{ GeV}/c$



ALI-PREL-87037
07/11/2025, WU CF 2025

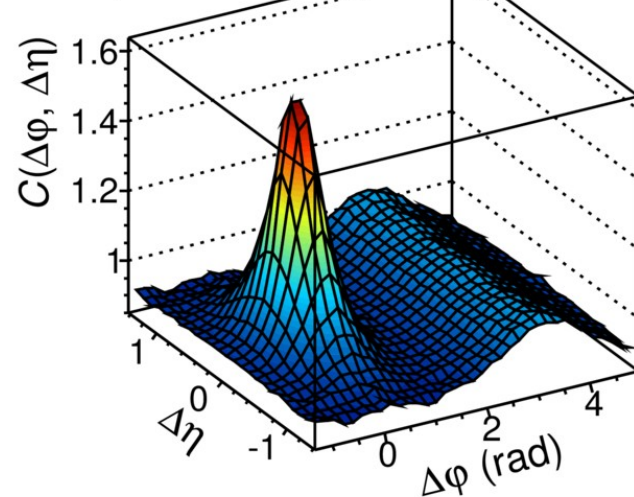
ALICE Preliminary, pp $\sqrt{s} = 7 \text{ TeV}$
pion like-sign pairs

$1.2 < p_T^{\text{sum}} < 1.4 \text{ GeV}/c$



LUKASZ GLACZYKOWSKI (WU)

$2.5 < p_T^{\text{sum}} < 8.0 \text{ GeV}/c$



40/43

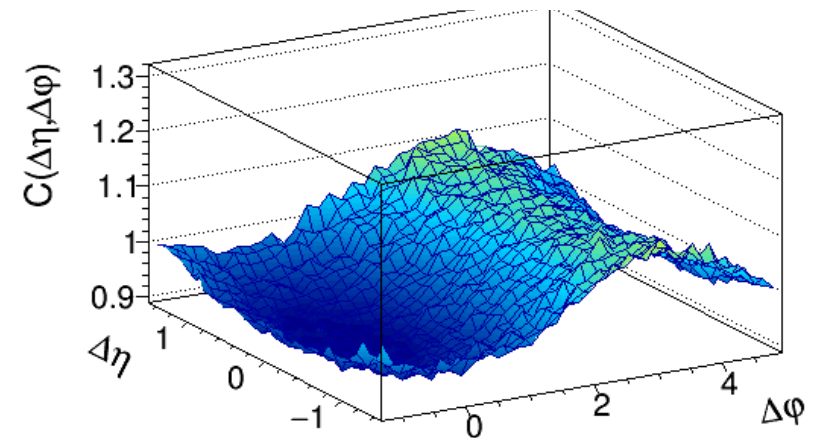
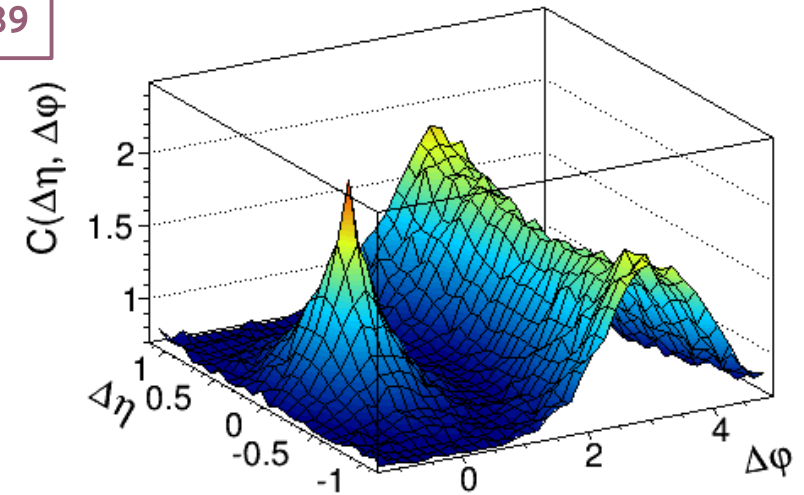
M. Janik, A. Kisiel, Ł. Graczykowski Nucl. Phys. A 956 (2016) 886-889

“Toy” Monte Carlo:

- Inclusion of conservation of energy, momentum and all quantum numbers local to the emission
- Our toy MC reproduces the standard “jet” correlation shape with near-side peak and away-side ridge

BUT

- Two-particle baryon-baryon correlation in data shows only global energy-momentum conservation features
- Yet, baryons **are** produced in jets (see e.g. proton-antiproton correlations), just no more than one



The puzzle remains unsolved!

Collective Effects:
the viewpoint of HEP MC codes

Torbjörn Sjöstrand
Department of Astronomy and Theoretical Physics
Lund University
Sölvegatan 14A, SE-223 62 Lund, Sweden

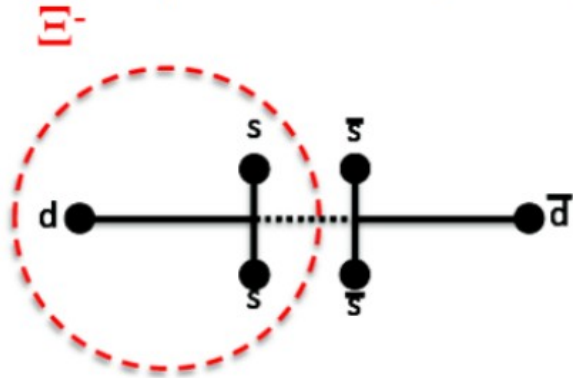
Quark Matter 2018, Venice, 13–19 May 2018

Nucl. Phys. A 982 (2019) 43-49

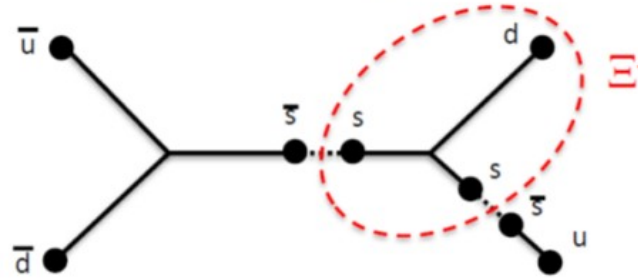
“The real problem is baryon production. [...] so it is clear we still lack some fundamental insight on baryon production, at least in the string context.”

Further studies

PYTHIA (standard configuration):



PYTHIA with junctions:

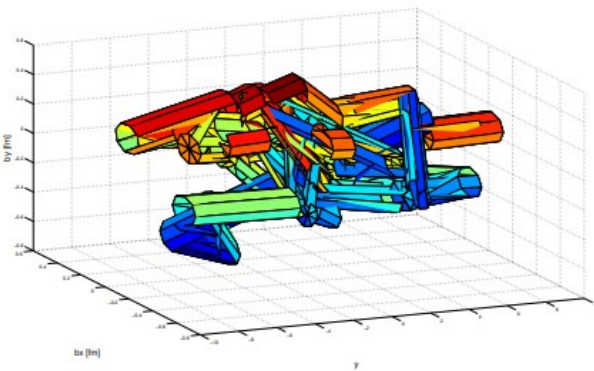


J. Adolfsson et al. *Eur. Phys. J. A* 56, 288 (2020)
(figures created by David Chinellato).

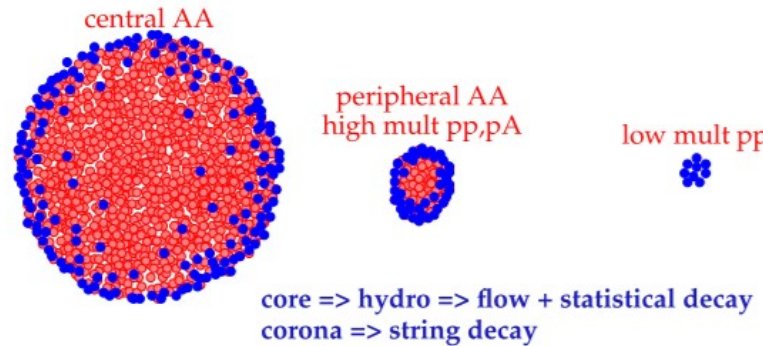
Predictions:

- PYTHIA: most quarks are produced at hadronisation \Rightarrow short-ranged correlations
- EPOS: quarks are produced in the core and diffuse before hadronisation \Rightarrow more long-ranged correlations

PYTHIA with ropes:



EPOS model:



C. Bierlich et al. *J. High Energ. Phys.* 2015, 148
Bottom right: K. Werner. hal-02434245 (2019)



First, let's define three variants of the model correlation function:

1. $C_{\text{base}} = S/M$, where M is the mixed-event distribution, contains only the event-wide correlations, without the QS and FSI effects added by the afterburner;
2. $C_{\text{full}} = S_w/M$ contains the full information, that is the event-wide correlations with additional effects of QS and FSI added by the afterburner;
3. $C_{\text{QS+FSI}} = M_w/M$ contains only the effects related to QS and FSI and is an equivalent to numerical integration of $C(\mathbf{k}^*) = \int S(\mathbf{k}^*, \mathbf{r}^*) |\Psi(\mathbf{k}^*, \mathbf{r}^*)|^2 d^4\mathbf{r}^*$.

S – same event distribution

M – mixed event distribution

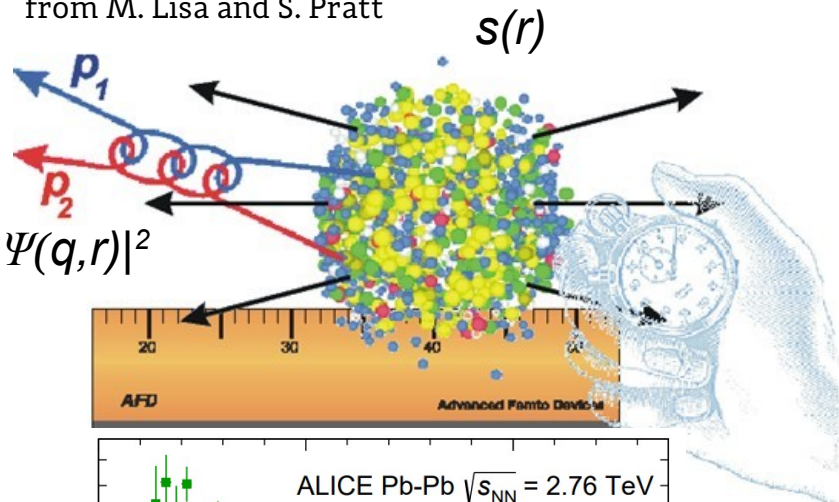
w – weight from Lednicky model

This can be done for both femtoscopic and angular CFs

Femtoscscopy - "traditional"



from M. Lisa and S. Pratt



$$C(q) = \int S(r) |\Psi(q,r)|^2 d^3r$$

measured correlation

Obtained by experiment

emission function (source size/shape)

Unknown

Two-particle wave function

Interaction known

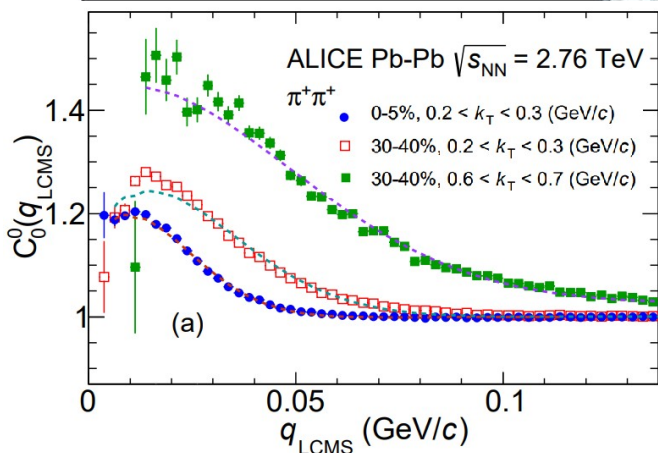
$$C(q) = N \frac{A(q)}{B(q)}$$

Probability ratio

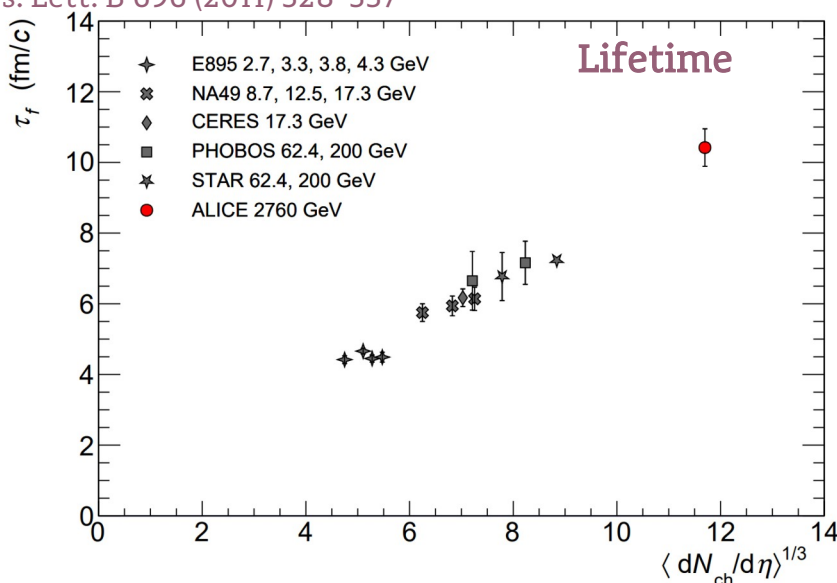
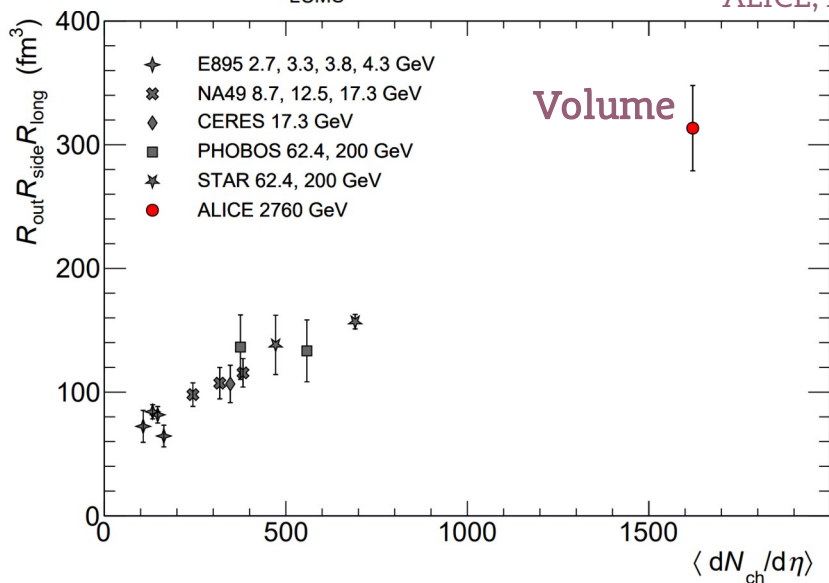
$$C(\mathbf{p}_1, \mathbf{p}_2) = \frac{P_{1,2}(\mathbf{p}_1, \mathbf{p}_2)}{P_1(\mathbf{p}_1) \cdot P_2(\mathbf{p}_2)}$$

$$q = 2 \cdot k^* = p_1 - p_2$$

$$r = x_1 - x_2$$

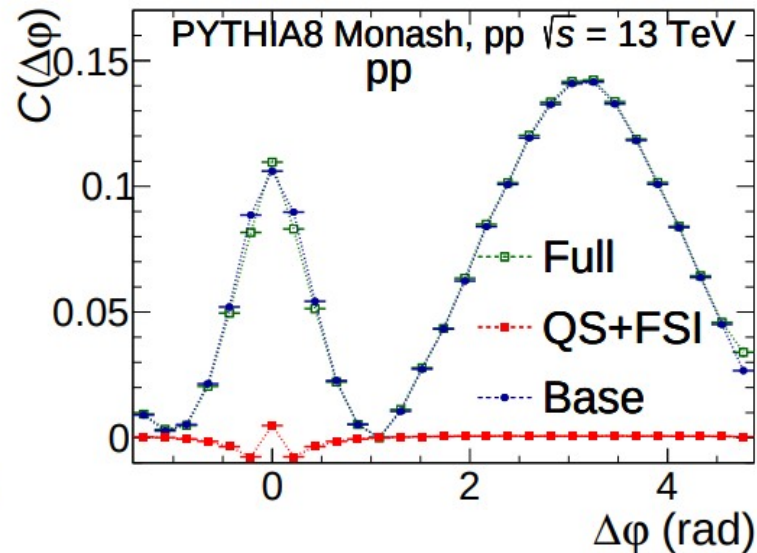
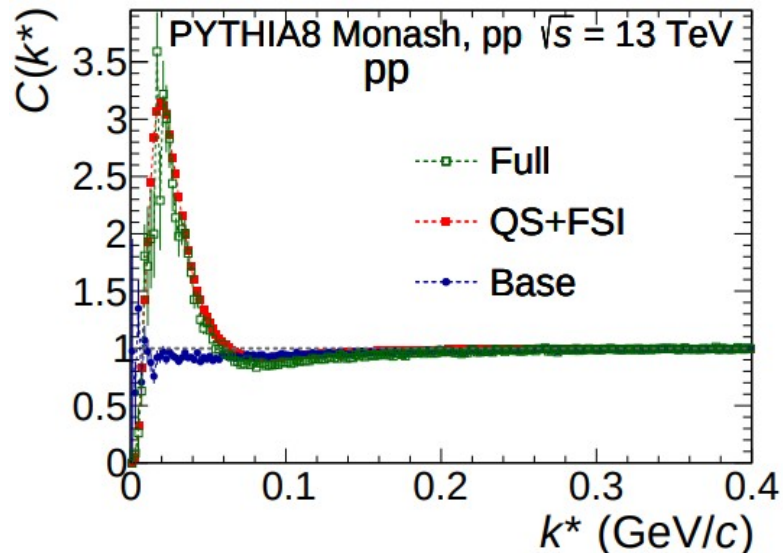
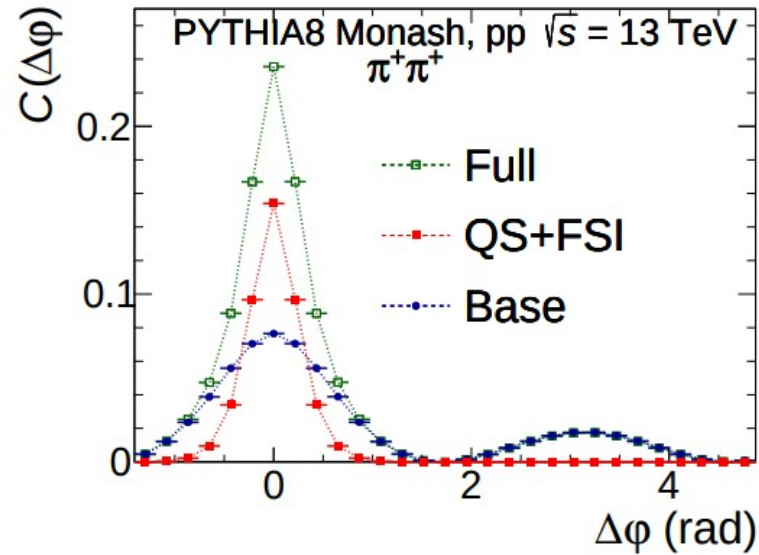
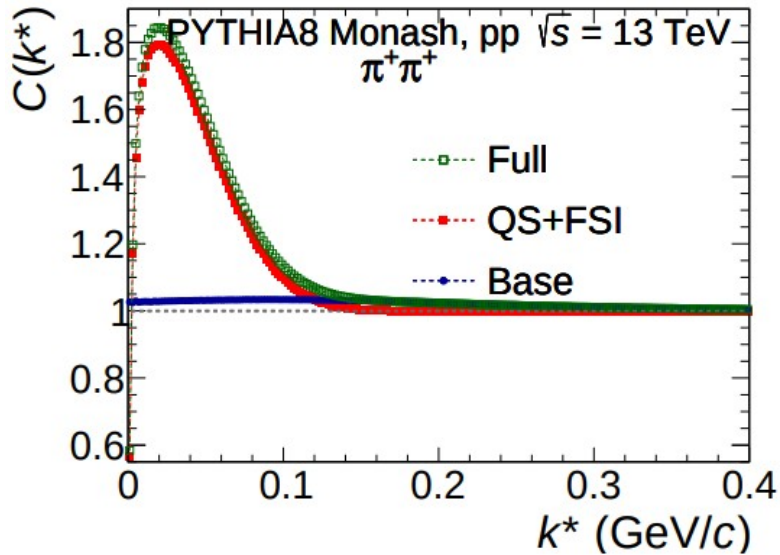


ALICE, Phys. Lett. B 696 (2011) 328-337



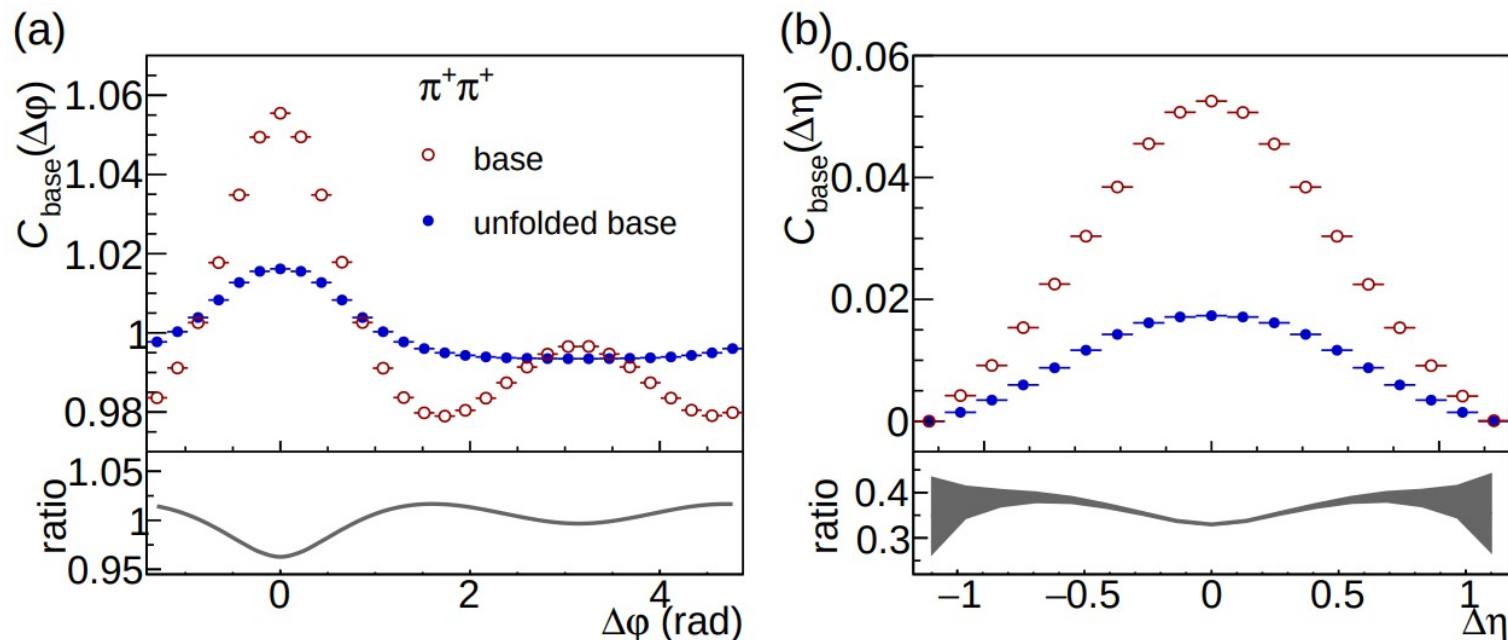


Calculated variants of femtoscopic and angular CFs using PYTHIA simulated events coupled to the L-L code





- The proposed unfolding procedure will work **ONLY** for short-range correlations, which include FSI and QS
 - for long-range (large k^*) correlations, i.e. jets, our algorithm is **too simple**
 - i.e. no energy-momentum conservation with such simple sampling



- Nevertheless, the algorithm works well for our use case and explains the origin of the small peak

$$C(q) = \int S(r) |\Psi(q, r)|^2 d^3 r$$

$q = 2 \cdot k^* = p_1 - p_2$

measured correlation emission function (source size/shape) pair wave function (includes cross section)

pair wave function $\longrightarrow \Psi = \exp(-ik^* r) + f \frac{\exp(ik^* r)}{r}$ s-wave scattering approximation

scattering amplitude $\longrightarrow f^{-1}(k^*) = \frac{1}{f_0} + \frac{1}{2} d_0 k^{*2} - ik^*$ effective range approximation

If **only** Strong FSI is present:

Lednický-Lyuboshitz equation

$$C(k^*) = 1 + \sum_s \rho_s \left[\frac{1}{2} \left| \frac{f^s(k^*)}{R} \right|^2 \left(1 - \frac{d_0^s}{2\sqrt{\pi}R} \right) + \frac{2\Re f^s(k^*)}{\sqrt{\pi}R} F_1(2k^*R) - \frac{\Im f^s(k^*)}{R} F_2(2k^*R) \right]$$

Sov. J. Nucl. Phys., 35, 770 (1982)

where ρ_s are the spin fractions

The correlation function is characterized by **three parameters**:

– **radius** R , **scattering length** f_0 , and **effective radius** d_0

– **cross section** σ (at low k^*) is simply: $\sigma = 4\pi |f|^2$