

# Investigating collective effects in small collision systems using PYTHIA 8 and EPOS4 simulations

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C. Pruneau





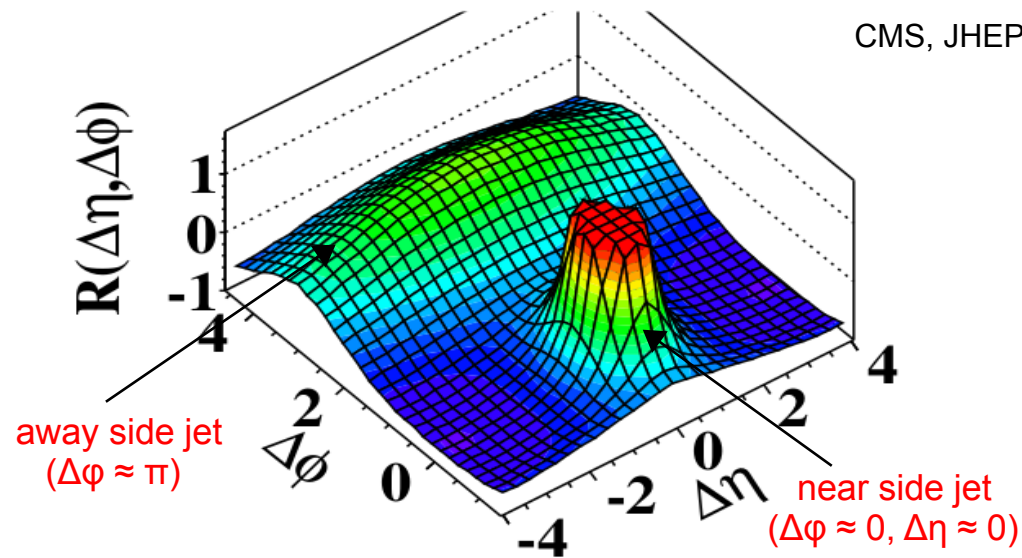
# Outline

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- Introducing experimental methods of flow analysis using cumulant and scalar product
- Determination of  $v_2$  and  $c_2\{2\}$ ,  $c_2\{4\}$  for EPOS4 and PYTHIA 8 and comparing results
- Comparison of balance function analysis of both event generators

(b) CMS MinBias,  $1.0\text{GeV}/c < p_T < 3.0\text{GeV}/c$

CMS, JHEP 1009 (2010) 091

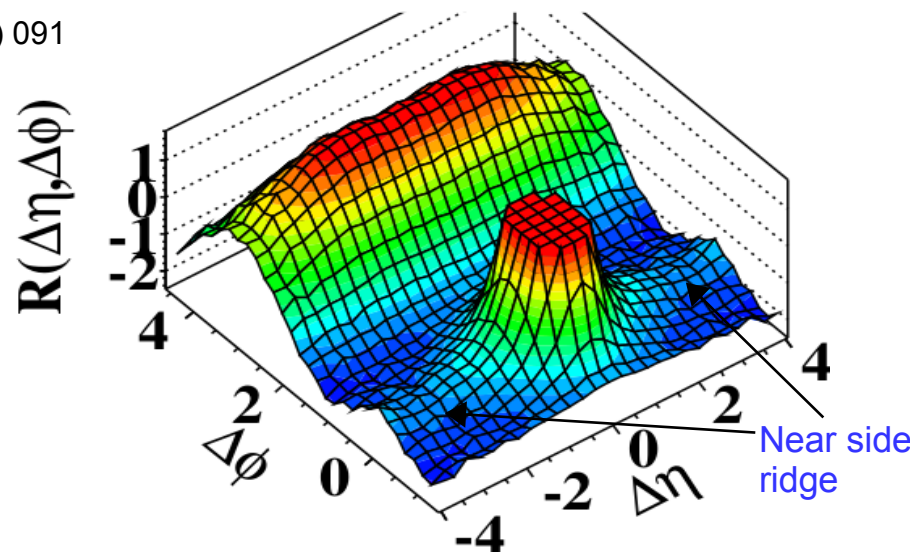
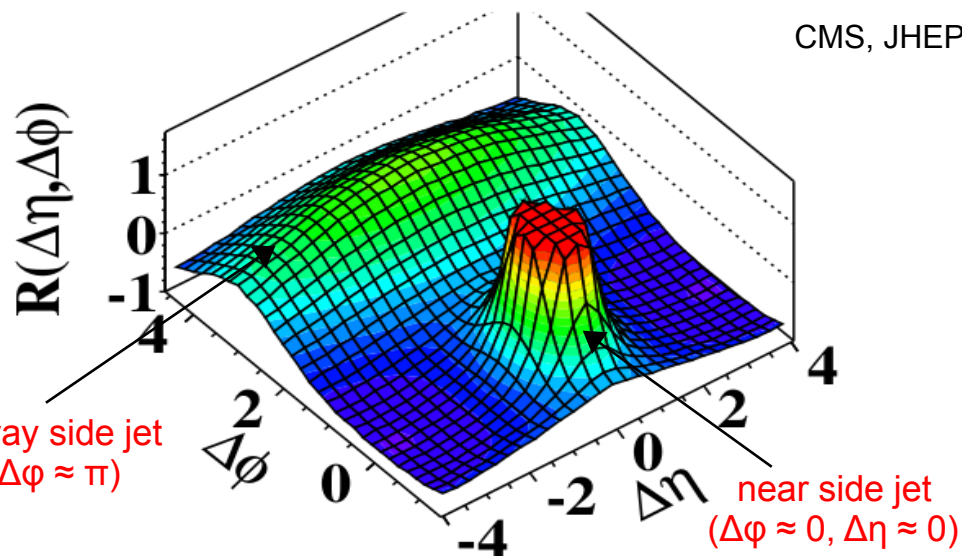


- Minimum bias pp
  - Nonflow contributions
    - Near-side jet peak (+resonances, HBT effects)
    - Recoil jet in away side

(b) CMS MinBias,  $1.0\text{GeV}/c < p_T < 3.0\text{GeV}/c$

(d) CMS  $N \geq 110$ ,  $1.0\text{GeV}/c < p_T < 3.0\text{GeV}/c$

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  - Nonflow contributions
    - Near-side jet peak (+resonances, HBT effects)
    - Recoil jet in away side

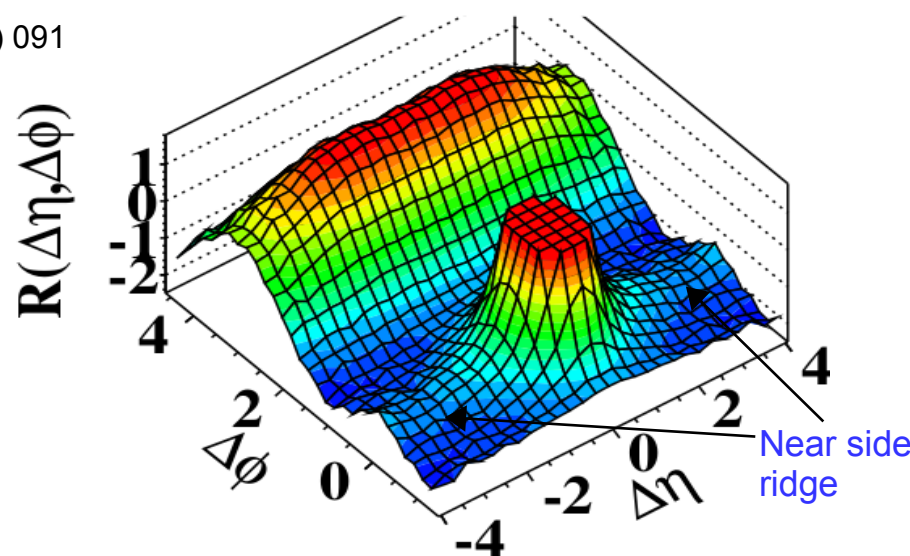
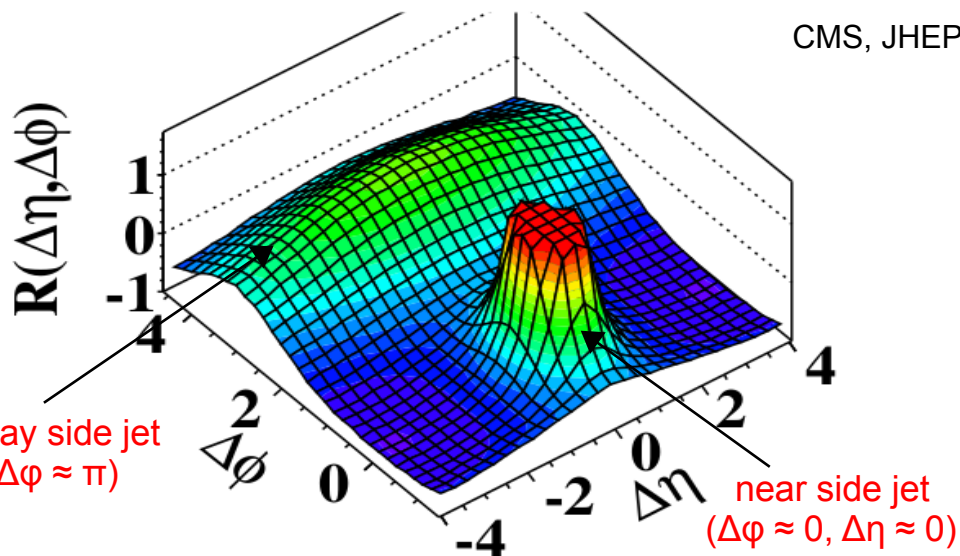
- High multiplicity pp
  - Near side ridge, typical of collective systems

- Decomposed into Fourier harmonics
 
$$1 + \sum_{n=1}^{\infty} 2 v_n \cos(n(\varphi - \Psi_n))$$

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- Minimum bias pp

- Nonflow contributions

- Near-side jet peak (+resonances, HBT effects)
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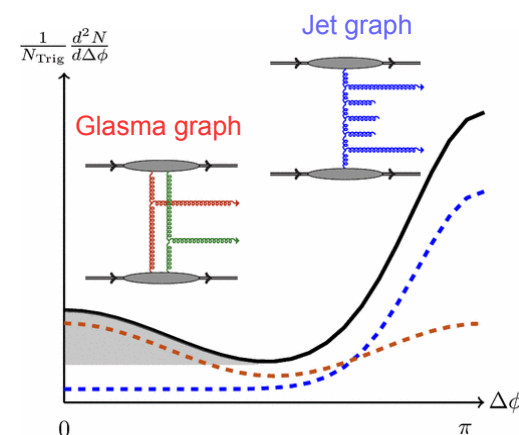
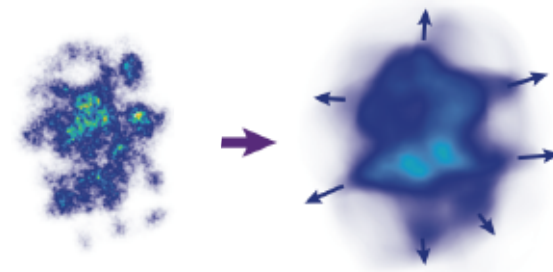
- High multiplicity pp

- Near side ridge, typical of collective systems

- Decomposed into Fourier harmonics
- $$1 + \sum_{n=1}^{\infty} 2 v_n \cos(n(\varphi - \Psi_n))$$

What is the origin of these collective effects?

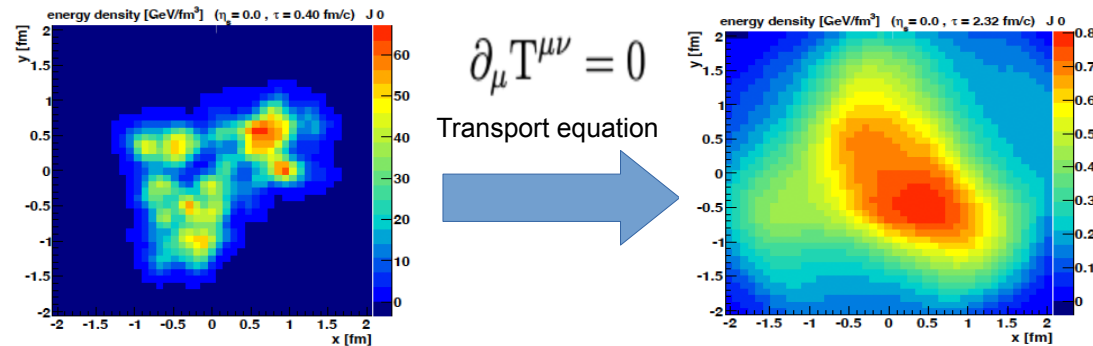
- Final state effects
  - Initial spatial eccentricities converted into momentum anisotropies via final state interactions
    - Hydrodynamics
    - Parton transport
    - Parton escape
  
- Initial state effects
  - Initial momentum anisotropies from initial interactions
    - Color Glass Condensate (CGC) Glasma
    - Color-field domains
    - Numerical solutions



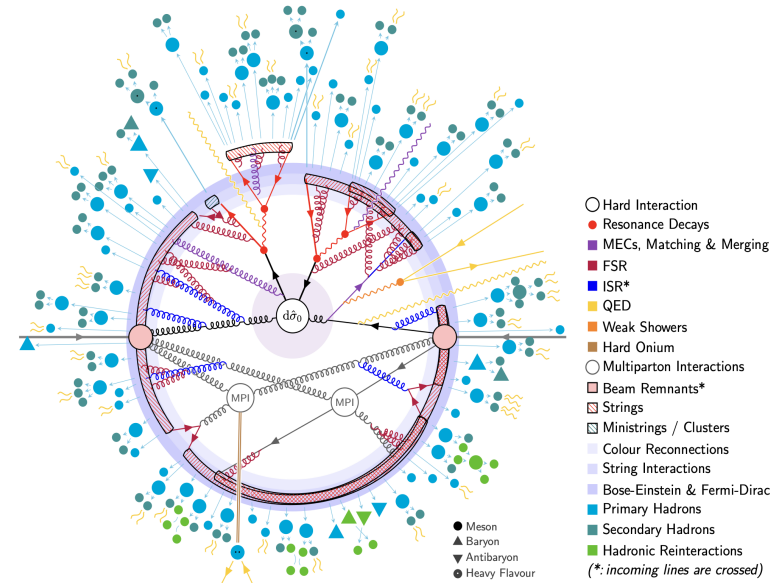
How to disentangle different regimes?

# Our approach: macroscopic vs microscopic models

K. Werner, arXiv: 2306.10277



C. Bierlich et al., arXiv: 2203.11601



- Macroscopic model: EPOS4

- Core-corona model with statistical hadronization
- Collective effects from hydrodynamical evolution of the medium

- Microscopic model: PYTHIA 8

- QCD strings with LUND fragmentation
- Collective effects from new processes
  - Color reconnection, rope hadronization, ...

- Scalar product (SP) method

$$v_n\{\text{SP}\} = \frac{\langle\langle \mathbf{u}_{n,k} \mathbf{Q}_n^* / M \rangle\rangle}{\sqrt{\langle\langle \mathbf{Q}_n^{*a} \mathbf{Q}_n^{*b} / (M^a M^b) \rangle\rangle}}$$

Particles of Interest  
(POI)

$$u_{n,x} = \cos(n\varphi)$$

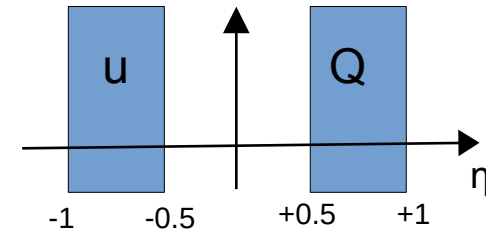
$$u_{n,y} = \sin(n\varphi)$$

Reference Particles  
(RPs)

$$Q_{n,x} = \sum_i \cos(n\varphi_i)$$

$$Q_{n,y} = \sum_i \sin(n\varphi_i)$$

S. Voloshin et al., arXiv:0809.2949





- Scalar product (SP) method

$$v_n\{SP\} = \frac{\langle\langle \mathbf{u}_{n,k} Q_n^* / M \rangle\rangle}{\sqrt{\langle\langle Q_n^{*a} Q_n^{*b} / (M^a M^b) \rangle\rangle}}$$

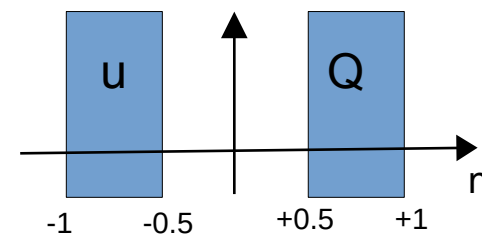
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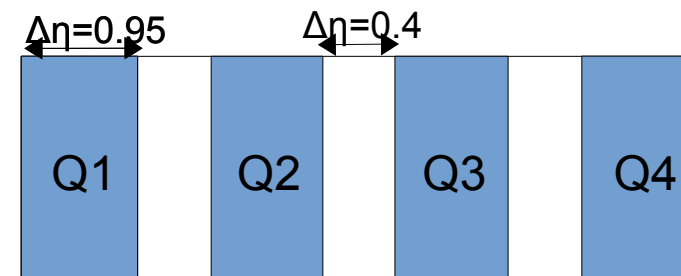
S. Voloshin et al., arXiv:0809.2949



- Cumulant method

- 2- and 4-particle azimuthal correlations for an event
- Averaging over all events  $\rightarrow$  2<sup>nd</sup> and 4<sup>th</sup> order cumulants

$$\begin{aligned} c_n\{2\} &= \langle\langle 2 \rangle\rangle = v_n^2 \\ c_n\{4\} &= \langle\langle 4 \rangle\rangle - 2\langle\langle 2 \rangle\rangle^2 = -v_n^4 \end{aligned}$$



A. Bilandzic et al., PRC 83, 044913 (2011)  
J. Jia et al., PRC 96, 034906 (2017)

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$$v_n\{\text{SP}\} = \frac{\langle\langle \mathbf{u}_{n,k} \mathbf{Q}_n^* / M \rangle\rangle}{\sqrt{\langle\langle \mathbf{Q}_n^{*a} \mathbf{Q}_n^{*b} / (M^a M^b) \rangle\rangle}}$$

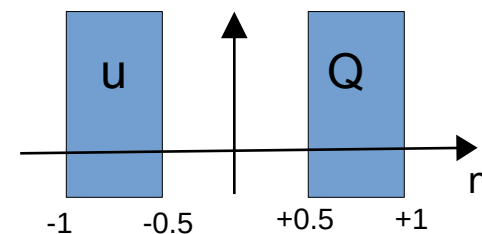
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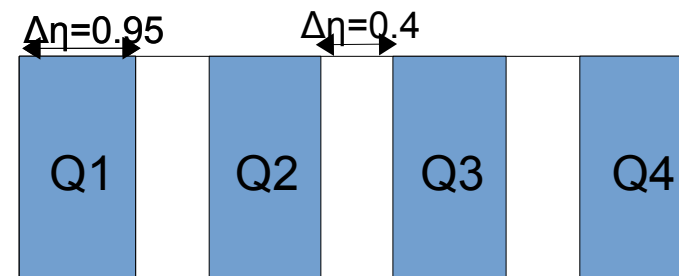
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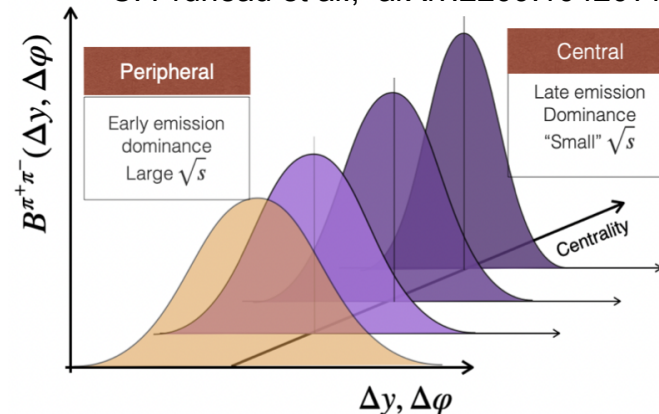
- Balance function

$$A_2^{\alpha\beta}(y_1, y_2) = \frac{\rho_2^{\alpha\beta}(y_1, y_2)}{\rho_1^\beta(y_2)} - \rho_1^\alpha(y_1)$$

$$B^{\alpha|\bar{\beta}}(y_1|y_2) = A_2^{\alpha|\bar{\beta}}(y_1|y_2) - A_2^{\bar{\alpha}|\bar{\beta}}(y_1|y_2)$$

$$B^{\bar{\alpha}|\beta}(y_1|y_2) = A_2^{\bar{\alpha}|\beta}(y_1|y_2) - A_2^{\alpha|\beta}(y_1|y_2)$$

C. Pruneau et al., arXiv:2209.10420v1

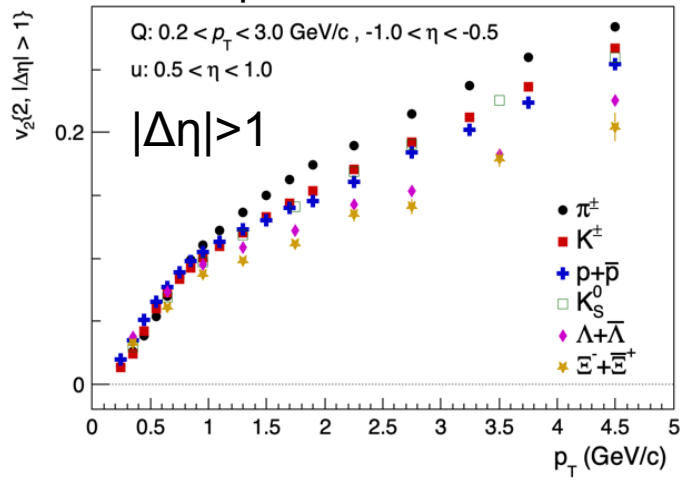


# $v_n$ in pp collisions @ 13.6 TeV

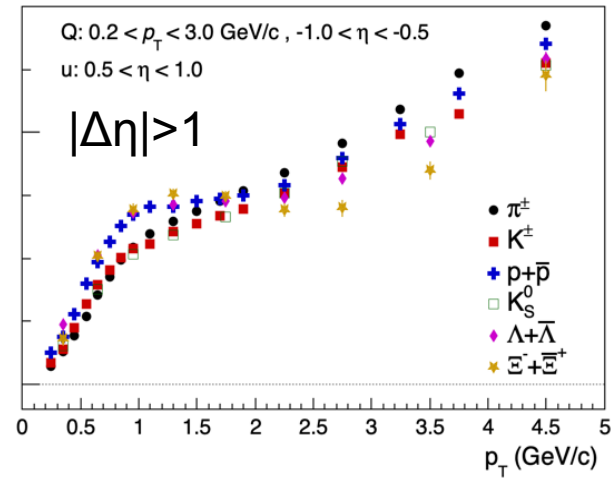
- PYTHIA 8
  - Rope hadronization <https://gitlab.com/Pythia8/releases/-/issues/80>
  - Monash tune
- EPOS4
  - Full simulation (core+corona+hadronic afterburner)

# PID $v_2$

“Rope hadronization”



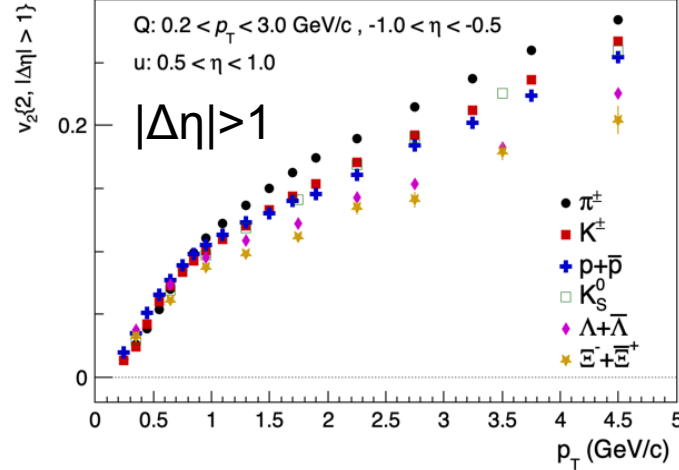
“Monash tune”



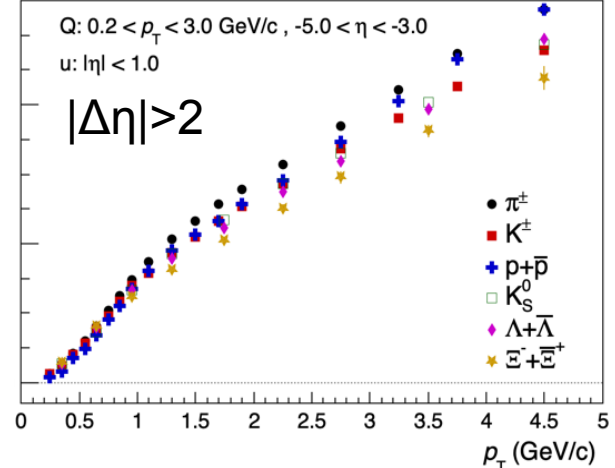
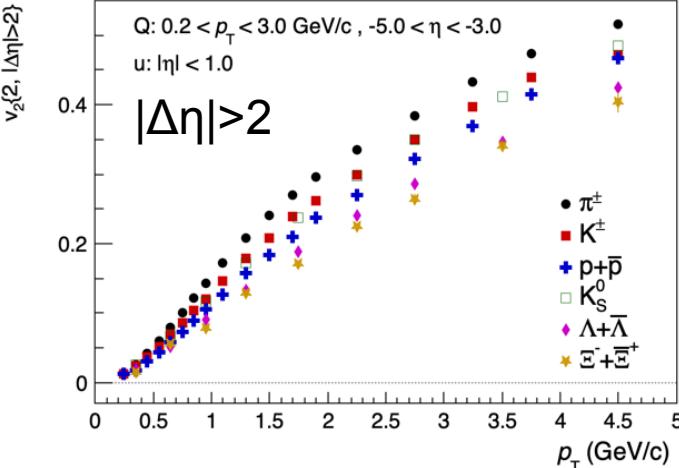
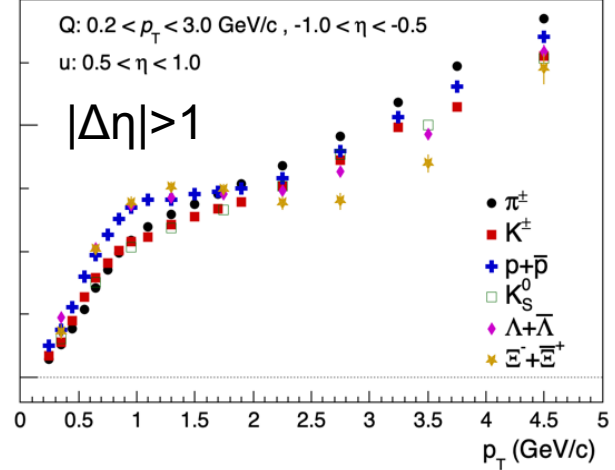
- PYTHIA 8
  - Mass ordering broken for  $|\Delta\eta| > 1$

# PID $v_2$

“Rope hadronization”

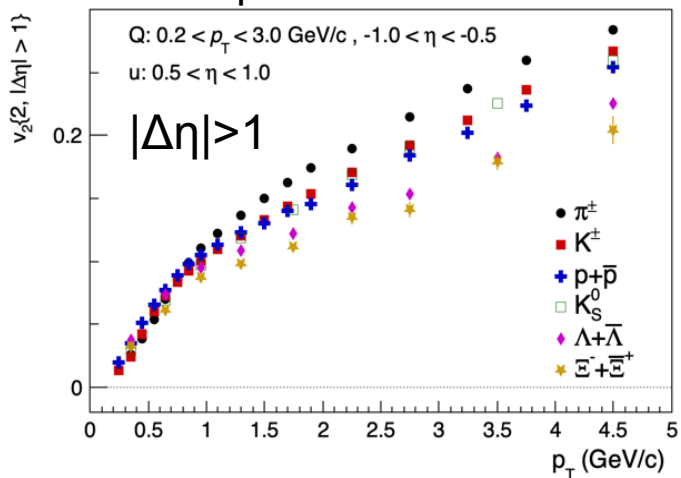


“Monash tune”

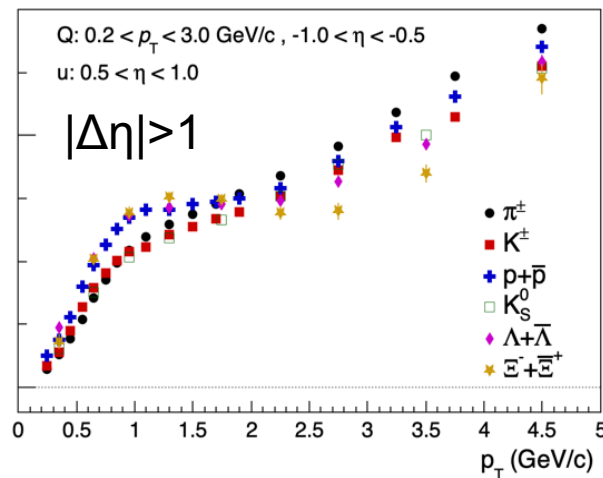


- PYTHIA 8
  - Mass ordering broken for  $|\Delta\eta|>1$
  - Small mass ordering for  $|\Delta\eta|>2$ 
    - More pronounced for rope hadronization
  - No particle type grouping

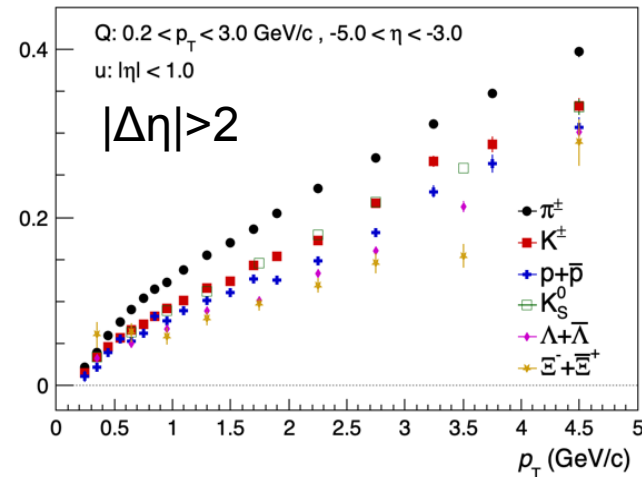
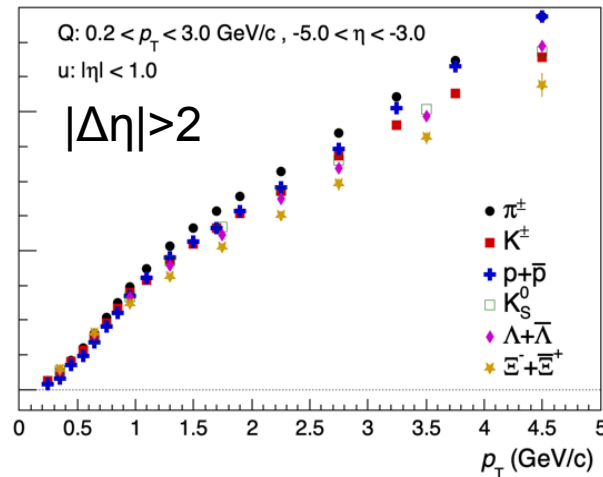
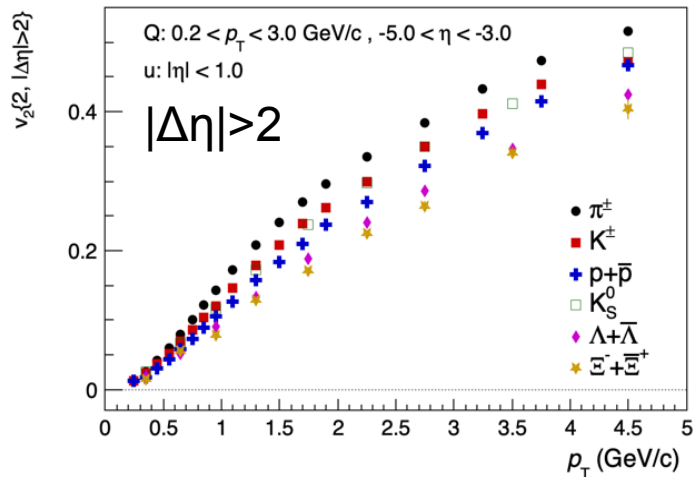
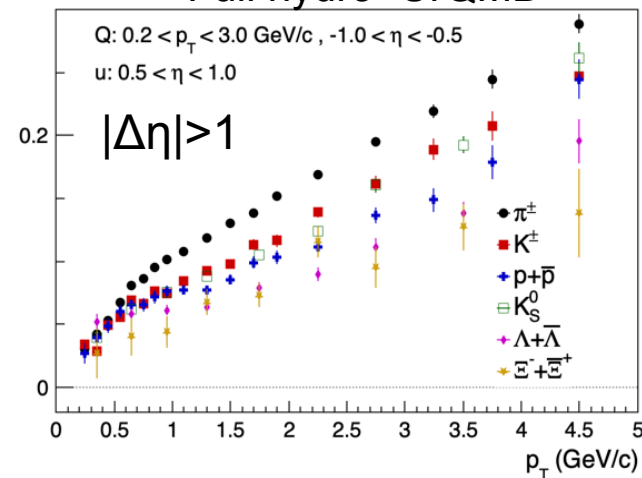
“Rope hadronization”



“Monash tune”



“Full hydro+UrQMD”



- PYTHIA 8

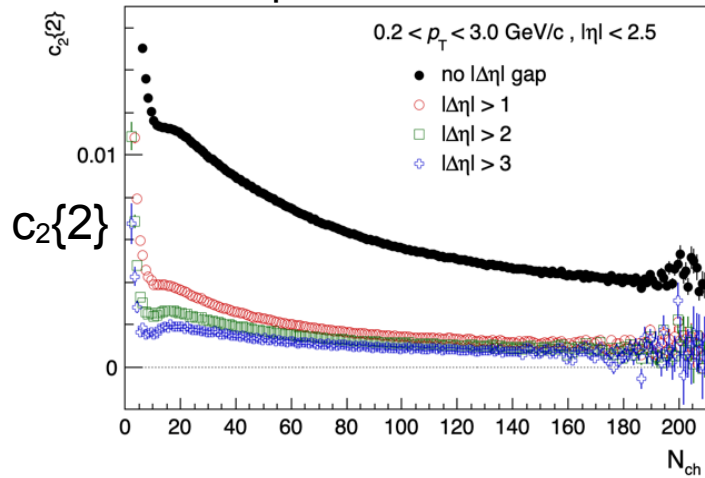
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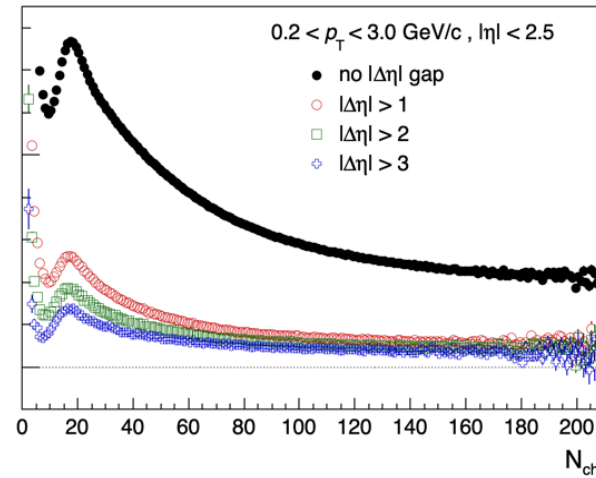
- Different trends than in PYTHIA
  - Influenced by core and hadronic afterburner (see backup)

# $c_2\{2\}$ and $c_2\{4\}$

“Rope hadronization”



“Monash tune”

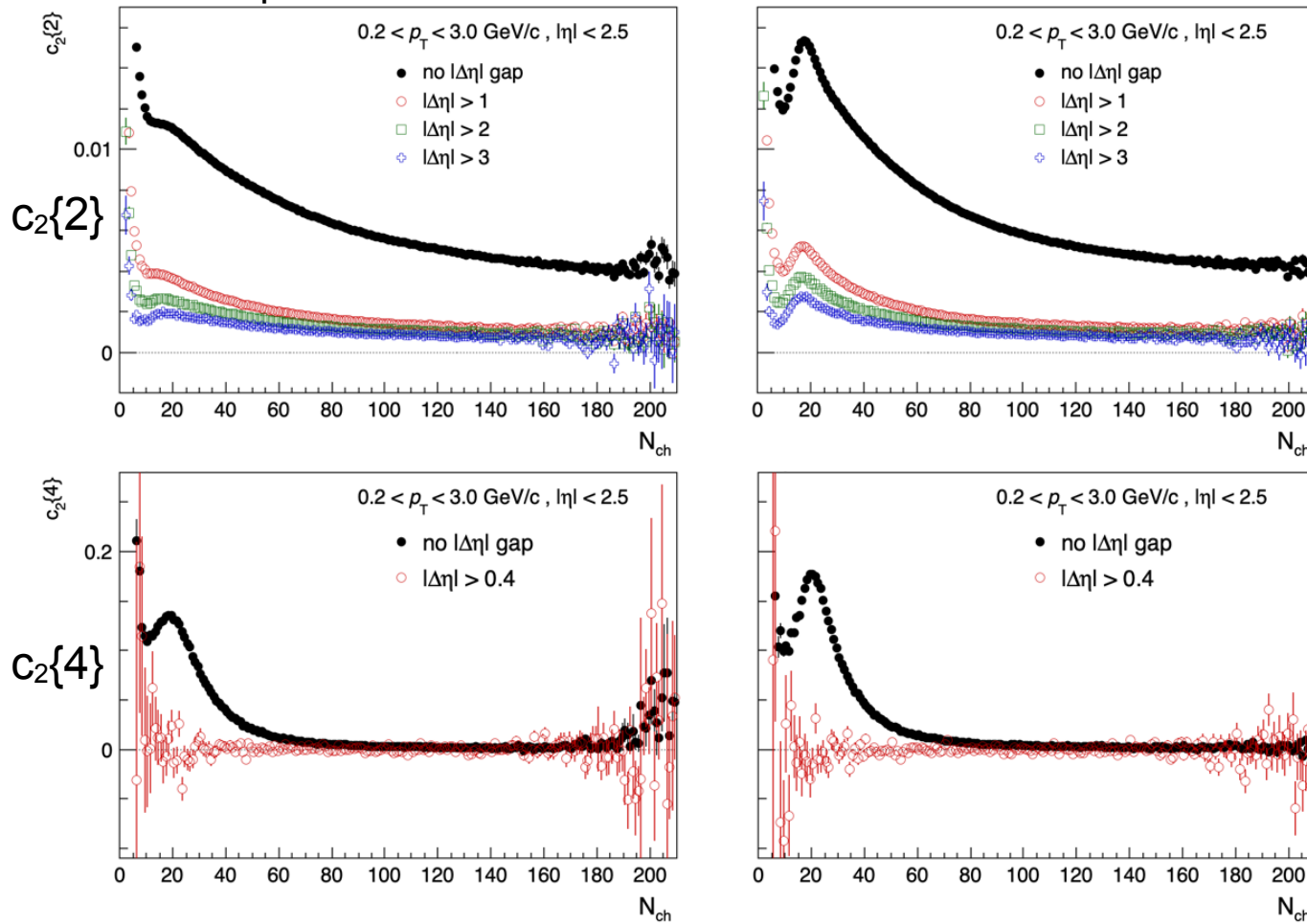


- $c_2\{2\} > 0$  at high multiplicities
  - Small dependence on  $|\Delta\eta|$  gap

# $c_2\{2\}$ and $c_2\{4\}$

“Rope hadronization”

“Monash tune”



- $c_2\{2\} > 0$  at high multiplicities
  - Small dependence on  $|\Delta\eta|$  gap
- $c_2\{4\} \sim 0 \rightarrow$  expected from Gaussian fluctuations

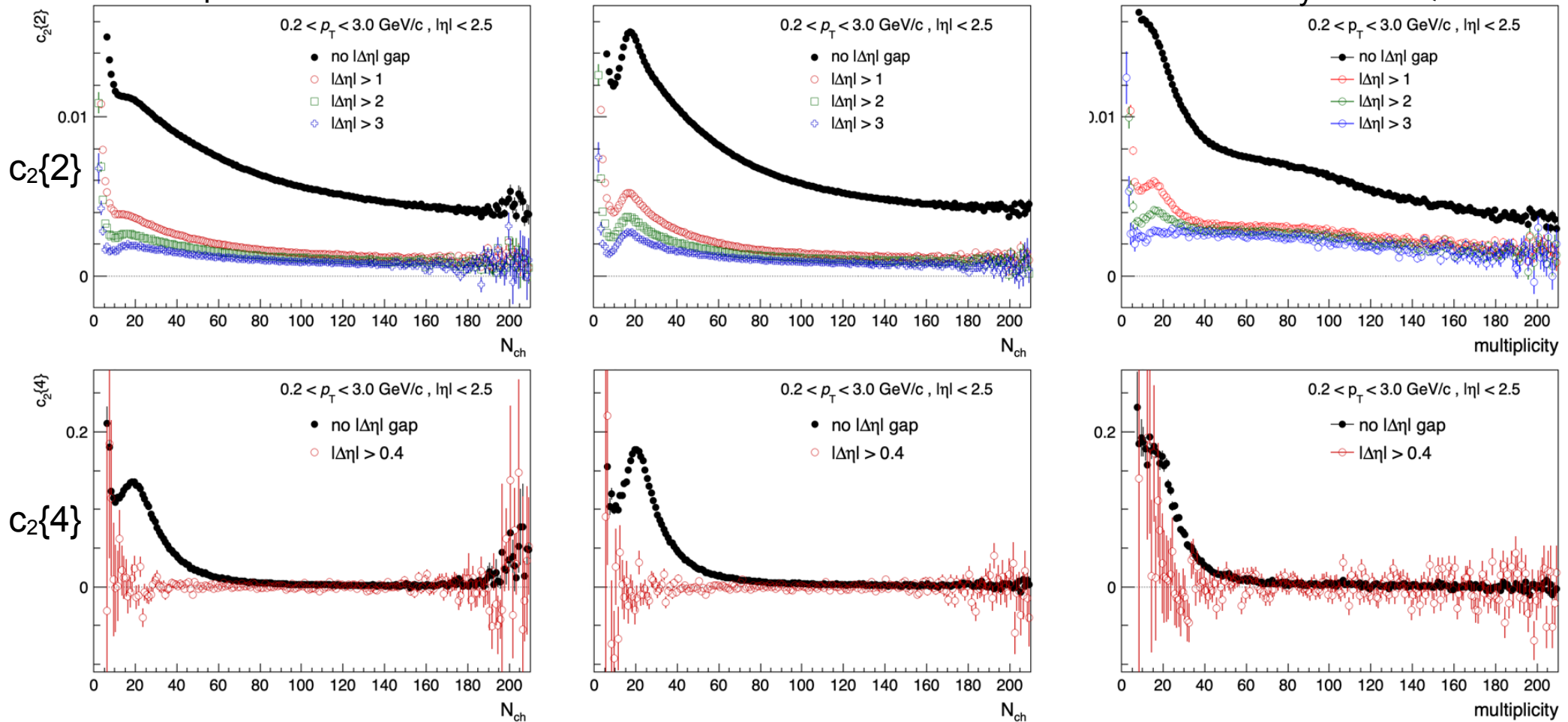


# $c_2\{2\}$ and $c_2\{4\}$

“Rope hadronization”

“Monash tune”

“Full hydro+UrQMD”



- $c_2\{2\} > 0$  at high multiplicities
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- $c_2\{4\} \sim 0 \rightarrow$  expected from Gaussian fluctuations
- Different trends in EPOS than in PYTHIA
  - More pronounced at low multiplicities

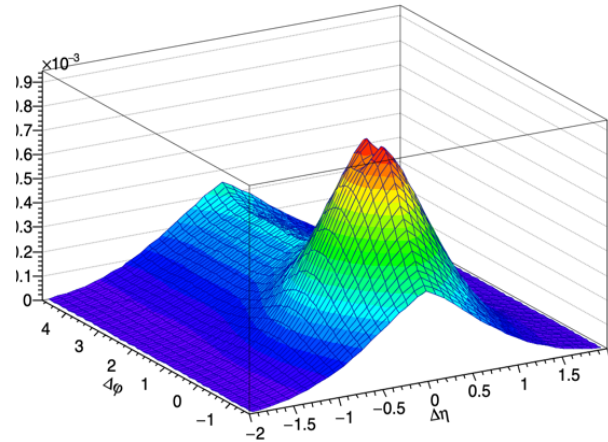
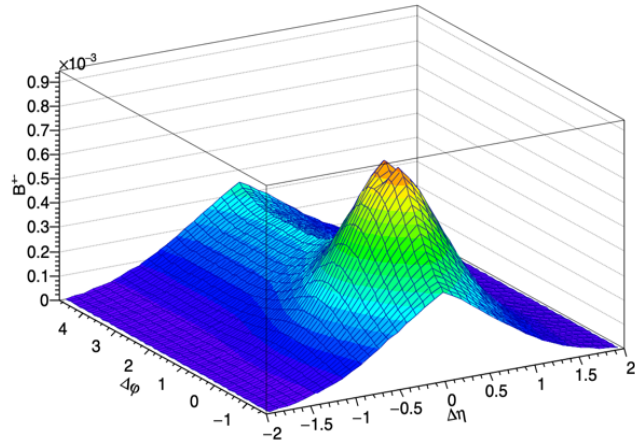
# Balance Function in pp collisions @ 13.6 TeV

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  - Monash tune
- EPOS4
  - Full simulation (core+corona+hadronic afterburner)

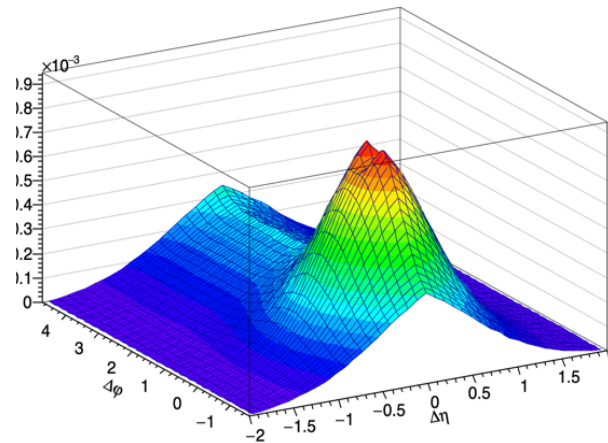
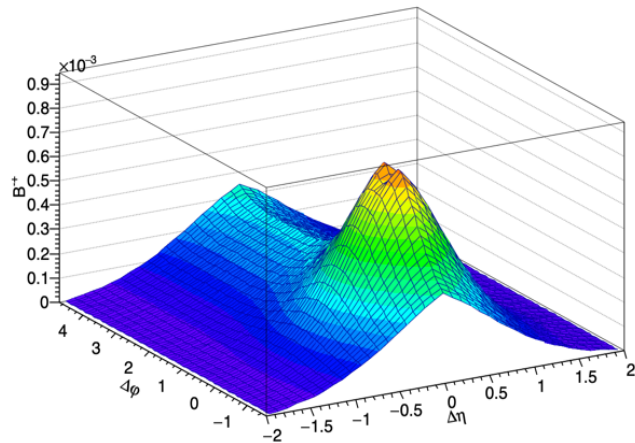
“Rope hadronization”

“Monash tune”

$B^{+-}$



$B^{-+}$



Integral definition:

$$I^{\alpha\bar{\beta}} = \frac{\langle N_2^{\alpha\bar{\beta}} \rangle}{\langle N_1^{\bar{\beta}} \rangle} - \frac{\langle N_2^{\bar{\alpha}\bar{\beta}} \rangle}{\langle N_1^{\bar{\beta}} \rangle}$$

$$I^{\bar{\alpha}\beta} = \frac{\langle N_2^{\bar{\alpha}\beta} \rangle}{\langle N_1^{\beta} \rangle} - \frac{\langle N_2^{\alpha\beta} \rangle}{\langle N_1^{\beta} \rangle}$$

Integral value  $B^{+-}$  : 0.469

0.490

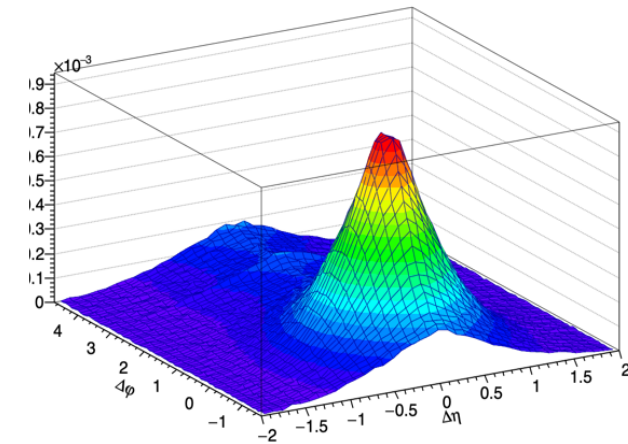
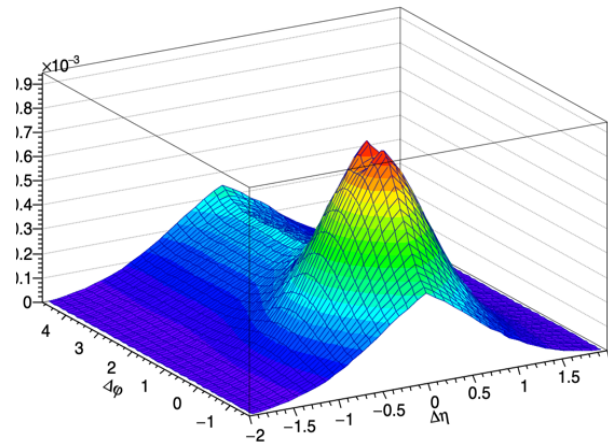
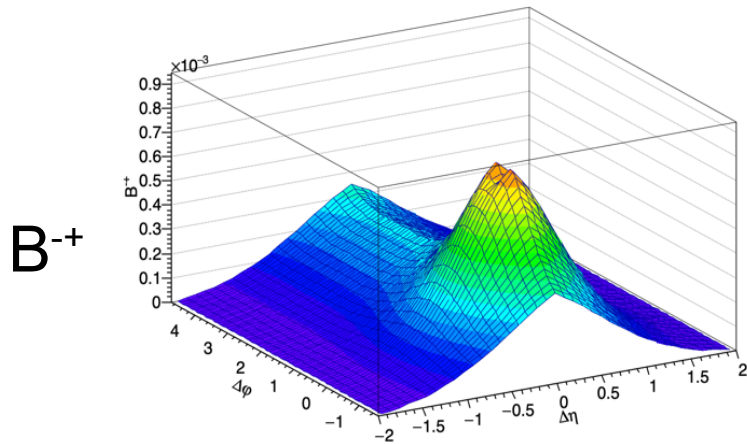
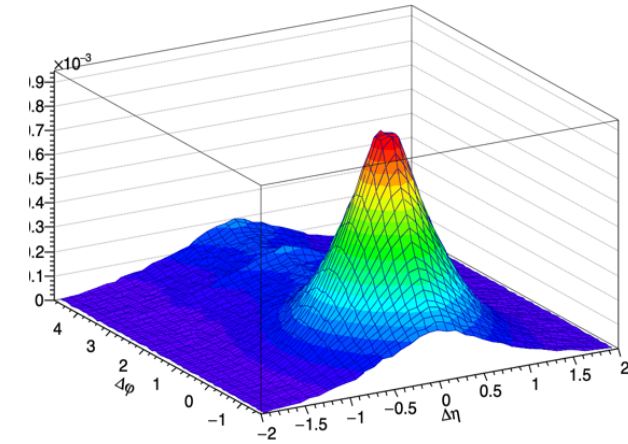
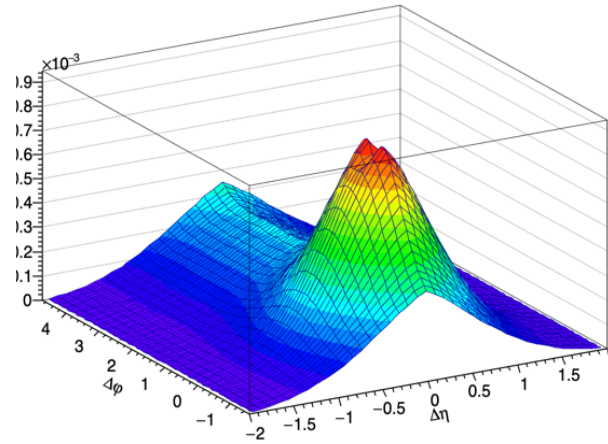
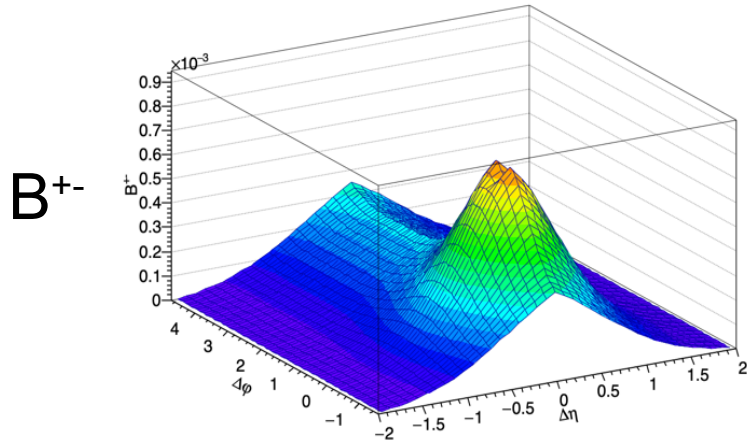
Integral value  $B^{-+}$  : 0.474

0.486

“Rope hadronization”

“Monash tune”

“Full hydro+UrQMD”



Integral value  $B^{+-}$  : 0.469

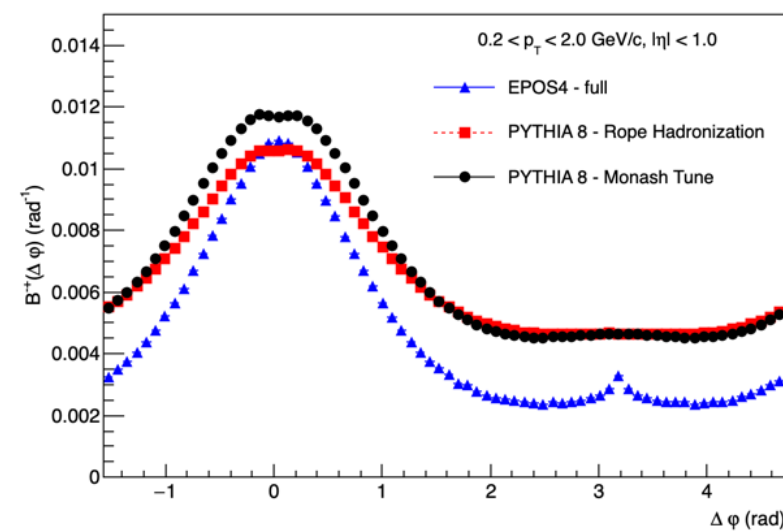
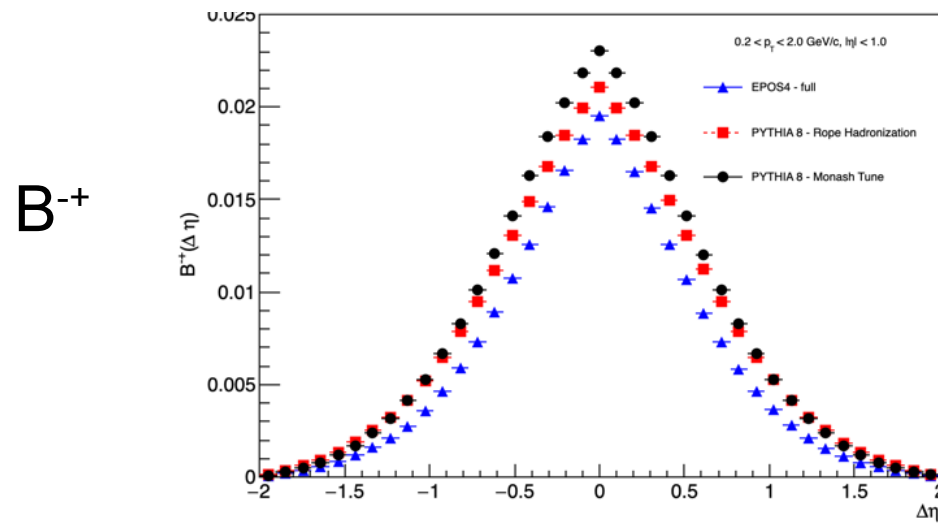
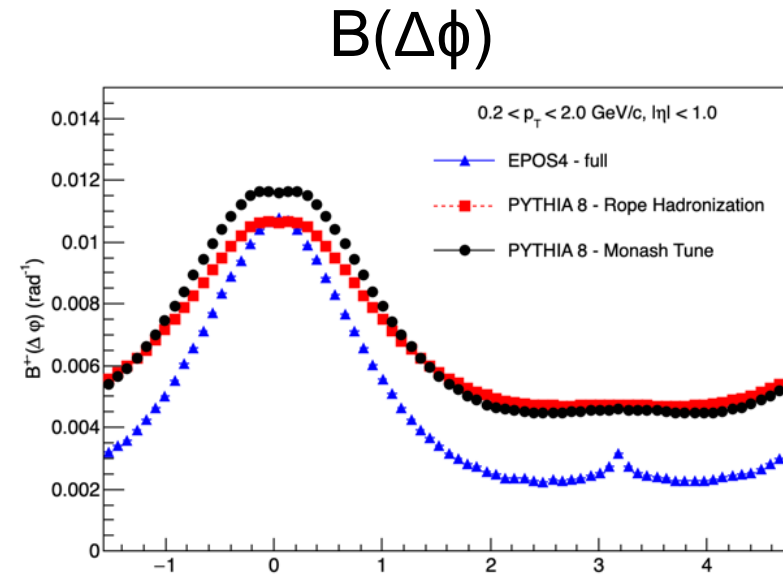
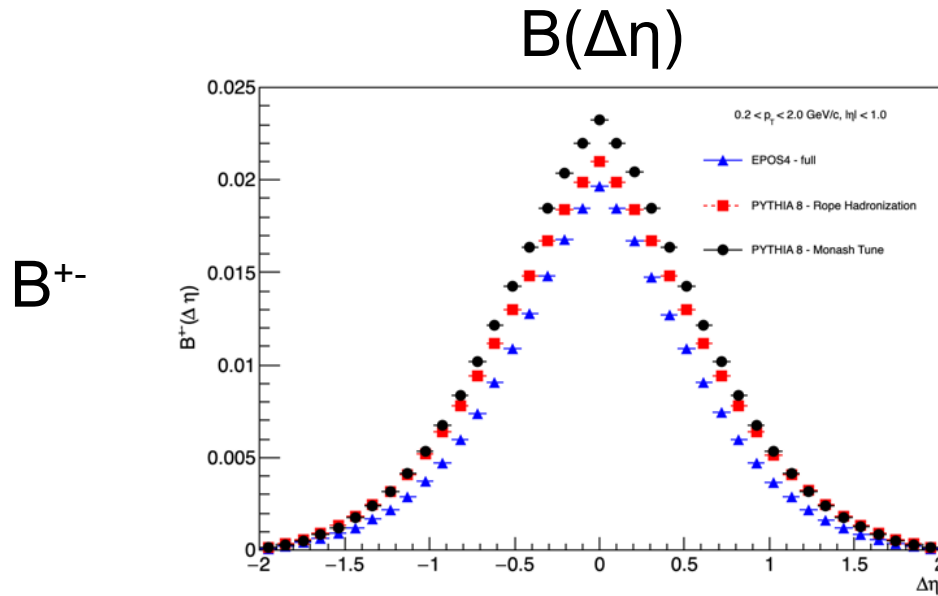
0.490

0.344

Integral value  $B^{-+}$  : 0.474

0.486

0.336



Different trends in EPOS than in PYTHIA

- Investigate collective effects in EPOS4 and PYTHIA 8
  - Different trends for various settings
- $c_2\{2\}$  decreasing with increasing multiplicity and  $|\Delta\eta|$  gap
  - Small dependence on  $|\Delta\eta|$  gap
- $c_2\{4\} \sim 0$  at high multiplicities
  - Expected for Gaussian fluctuations
- PID  $v_2$ : mass ordering for large  $|\Delta\eta|$  gap
  - No particle type grouping
- Balance function: different trends in away side

Thank you!



# Backup

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