

# Femtoscscopy with Lévy sources at NA61/SHINE

XVI Workshop on Particle Correlations and Femtoscopy  
IV Resonance Workshop  
Catania, Italy

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9 November, 2023




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# Two-pion femtoscopic correlations in Be+Be collisions at $\sqrt{s_{NN}} = 16.84$ GeV measured by the NA61/SHINE at CERN

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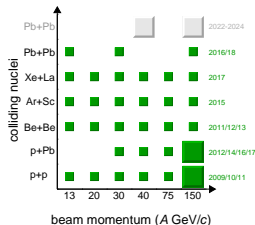
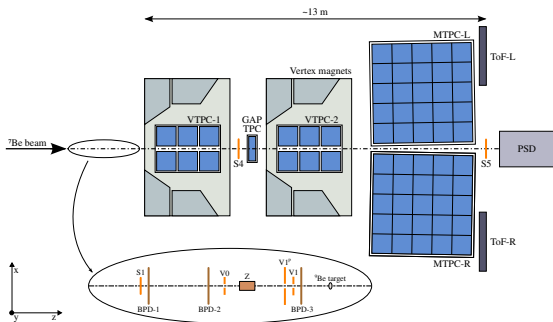
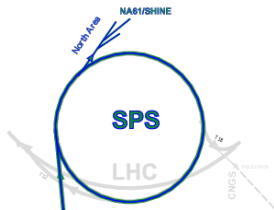
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B.Porfy for NA61/SHINE Collab. WPCF2023

- 1 Experiment
- 2 Femtoscopy analysis details
- 3 Lévy HBT results
- 4 Intermittency analysis
- 5 Conclusion

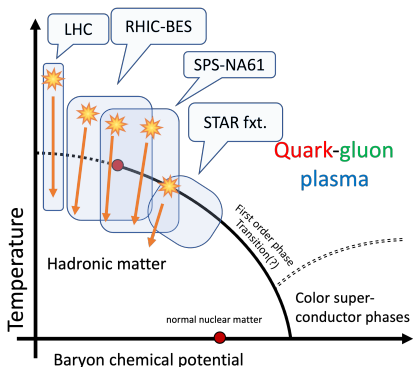
# The NA61/SHINE Detector

- Located at CERN SPS, North Area
- Fixed target experiment; upgrade during LS2
- Large acceptance hadron spectrometer (TPC)
  - ▶ Covering the full forward hemisphere
  - ▶ Outstanding tracking, down to  $p_T = 0$  GeV/ $c$
- Different systems scanned in beam energy
- Strong interactions programme:
  - ▶ Search for Critical Point: **femtoscopy, intermittency, fluctuations, ...**



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# Critical Point Search Using Femtoscopy



- At Critical Point - fluctuations at all scales
- Power-law in spatial correlations
- Critical exponent  $\eta$
- QCD universality class  $\leftrightarrow$  3D Ising:
  - Halasz et al., Phys.Rev.D58 (1998) 096007
  - Stephanov et al., Phys.Rev.Lett.81 (1998) 4816
  - ▶ 3D Ising:  $\eta = 0.03631(3)$   
El-Showk et al., J.Stat.Phys.157 (2014) 869
  - ▶ Random field 3D Ising  $\eta = 0.50(5)$   
Rieger, Phys.Rev.B52 (1995) 6659

- Spatial correlation exponent near **Critical Point?**

- Lévy sources applicable in HBT

Csörgő, Hegyi, Zajc, Eur.Phys.J.C36 (2004) 67, nucl-th/0310042

- System size scan progress at 150A GeV/c: **Be+Be, Ar+Sc**

Be+Be: NA61/SHINE, Eur.Phys.J.C 83 (2023) 10, 919; Ar+Sc: Universe 2023, Volume 9, Issue 7, 298  
next Xe+La or Pb+Pb

- Energy scan ongoing in Ar+Sc

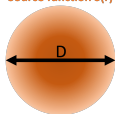
# Bose-Einstein Correlations in Heavy-Ion Physics

A tool to measure spatial correlations:

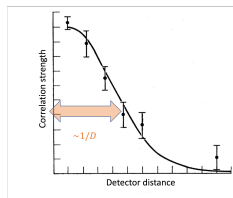
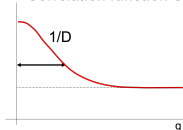
Bose-Einstein relative momentum correlations

- R. Hanbury Brown, R.Q.Twiss observed Sirius with two optical telescopes  
R. Hanbury Brown and R. Q. Twiss Nature 178 (1956)
  - ▶ Intensity correlations as a function of detector distance
  - ▶ Measuring angular size of point-like sources
- Goldhaber, Goldhaber, Lee and Pais: applicable in high energy physics: (for identical pions)  
Goldhaber, Goldhaber, Lee and Pais Phys.Rev.Lett.3 (1959) 181
  - ▶ Momentum correlation  $C(q)$ ,  $q = |p_1 - p_2|$ , is related to the source  $S(x)$   
 $C(q) \cong 1 + |\tilde{S}(q)|^2$  where  $\tilde{S}(q)$  Fourier transform of  $S(q)$

Source function  $S(r)$



Correlation function  $C(q)$



- $S(r)$  frequently assumed to be Gaussian  $\rightarrow$  Gaussian  $C(q)$

# Lévy Distribution in Heavy-Ion Physics

- Measurements not fully supporting Gaussian  $\rightarrow$  Generalized CLT

**Lévy-stable distribution:**  $\mathcal{L}(\alpha, R, r) = \frac{1}{(2\pi)^3} \int d^3q e^{iqr} e^{-\frac{1}{2}|qR|^\alpha}$

- $\triangleright$  From generalization of Gaussian, power-law tail:  $\sim r^{-(d-2+\alpha)}$ 
  - $\triangleright \alpha = 1$  Cauchy,  $\alpha = 2$  Gaussian
- The shape of the correlation function with Lévy source:  $C(q) = 1 + \lambda \cdot e^{-|qR|^\alpha}$ 
  - $\triangleright \alpha = 1$ : Exponential,  $\alpha = 2$ : Gaussian Csörgő, Hegyi, Zajc, Eur.Phys.J.C36 (2004) 67-78
- Reasons for Lévy source:

- $\triangleright$  QCD jets; Anomalous diffusion; Critical phenomena; ...

Csörgő, Hegyi, Novák, Zajc, Acta Phys.Polon.B36 (2005) 329-337

Csanád, Csörgő, Nagy, Braz.J.Phys.37 (2007) 1002

Csörgő, Hegyi, Novák, Zajc, AIP Conf.Proc.828 (2006) 525-532

Metzler, Klafter, Physics Reports 339 (2000) 1-77

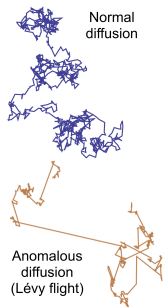
Kincses, Stefaniak, Csanád, Entropy 24 3 (2022) 308

Kórodi, Kincses, Csanád, arXiv:2212.02980

- Lévy distributions lead to power-law spatial correlations

- Spatial correlation at the Critical Point:  $\sim r^{-(d-2+\eta)}$

- $\triangleright$  Lévy-exponent  $\alpha$  identical to correlation exponent  $\eta$





# Lévy Distribution in Heavy-Ion Physics

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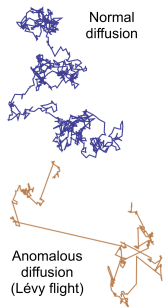
Kincses, Stefaniak, Csanád, Entropy 24 3 (2022) 308

Kórodi, Kincses, Csanád, arXiv:2212.02980

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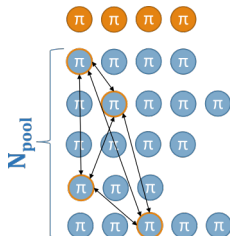
- $\triangleright$  Lévy-exponent  $\alpha$  identical to correlation exponent  $\eta$



# Correlation Function Measurement Details

- Be+Be at 150A GeV/c beam momentum, 0 - 20% centrality
- Ar+Sc at 150A GeV/c beam momentum, 0 - 10% centrality
- Track and pair cuts
- Particle Identification via  $dE/dx$  method  $\pi^-$ ,  $\pi^+$

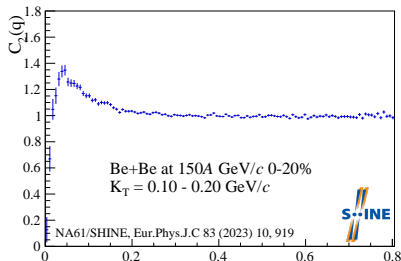
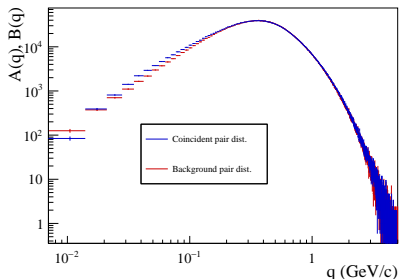
- $A(q)$  - Pairs from same event
- $B(q)$  - Pairs from mixed events
- $C(q)$  - Correlation function,  $C(q) = A(q)/B(q)$



- Correlation function  $q_{LCMS}$  1D variable
- LCMS: Longitudinally CoMoving System
- $m_T \equiv \sqrt{m^2 + (K_T/c)^2}$ ;  $K = (p_1 + p_2)/2$ ;  $K_T = \frac{1}{2}\sqrt{K_x^2 + K_y^2}$
- 4  $m_T$  bins Be+Be, 8  $m_T$  bins Ar+Sc

# Bose–Einstein Correlation Function

- $C_2(q)$ : B–E peak and Coulomb hole, at low  $q$  values:



- Like charged pairs: Coulomb interaction  $\rightarrow$  Coulomb correction (CC)
  - ▶ Calc: complicated numerical integral
  - ▶ Numerically possible: look-up table  $\rightarrow$  physical parameter parametrization  
Nagy, M., Purzsa, Á., Csanád, M. and Kincses, Dániel, arXiv:2308.10745
- Meas.: LCMS, CC.: PCMS (pair center of mass) negligible, BUT
- 1D spher. symm. source LCMS not spherical PCMS

$$R \rightarrow R_{\text{PCMS}} = \sqrt{\frac{1 - \frac{2}{3}\beta_T^2}{1 - \beta_T^2}} \cdot R_{\text{LCMS}}, \quad q_{\text{inv}} \approx \sqrt{1 - \beta_T^2/3} \cdot q_{\text{LCMS}}, \quad \beta_T = \frac{K_T}{m_T}$$

B. Kurgys, D. Kincses, M. Nagy, and M. Csanád, Universe 2023, 9(7), 328

# Parameters of Lévy-source

- Fit function: Bowler-Sinyukov

$$C(q) = 1 - \lambda + (1 + e^{-|qR|^\alpha}) \cdot \lambda \cdot K(q)$$

Yu. Sinyukov et al., Phys.Lett.B432 (1998) 248,  
M.G. Bowler, Phys.Lett.B270 (1991) 69

- $R$  Lévy-scale parameter:

- ▶ Length of homogeneity

- ▶ From simple hydro calc.:

$$R_{\text{HBT}} = R / \sqrt{1 + (m_T / T_0) \cdot u_T^2}$$

- $\lambda$  correlation strength:

- ▶ Core-halo ratio:

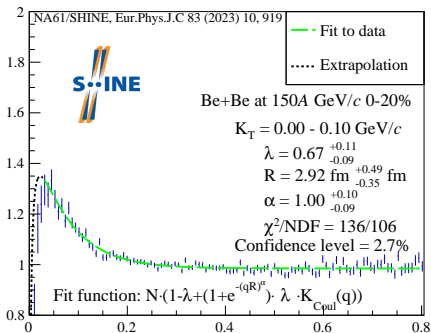
$$\lambda = \left( \frac{N_{\text{core}}}{N_{\text{core}} + N_{\text{halo}}} \right)^2$$

- ▶ Core: primordial pions

- ▶ Halo: pions from long-lived resonances

- $\alpha$  Lévy-stability index

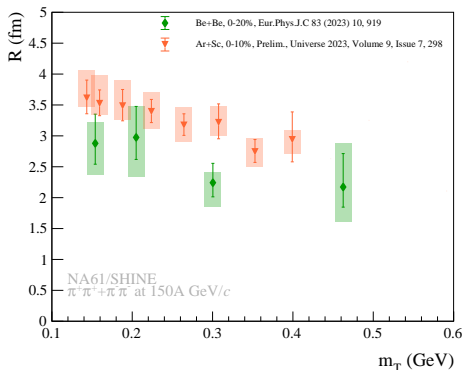
- ▶  $\alpha = 2$ : Gauss shape, simple hydro
- ▶  $\alpha < 2$ : Generalized central limit theorem
- ▶  $\alpha = 0.5$ : Conjectured value at CP



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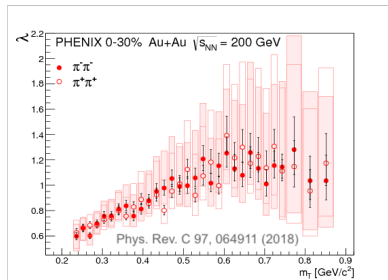
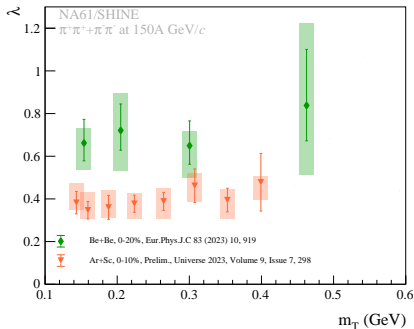
# Lévy-scale parameter $R$ vs. $m_T$

- Describes length of homogeneity
- From hydro:  $R \sim 1/\sqrt{m_T}$  (For Gaussian source)  
Csörgő, Lörstad, Phys.Rev.C54 (1996) 1390-1403  
S. V. Akkelin and Yu. M. Sinyukov, Phys.Lett.B356 (1995) 525-530  
S. Chapman, P. Scotto and U. W. Heinz, Phys.Rev.Lett.74 (1995) 4400-4403
- Visible  $m_T$  dependence - sign of transverse flow
  - ▶  $\alpha$  anticorrelates with  $R$ ,  $\lambda$ ; increase in  $\alpha \rightarrow$  decrease in  $R$



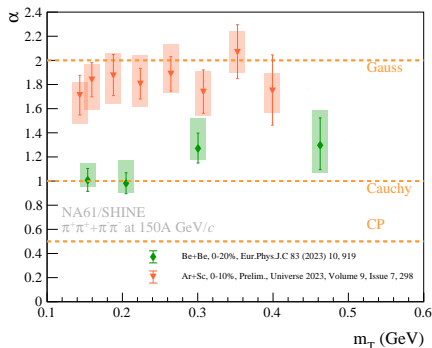
# Correlation Strength $\lambda$ vs. $m_T$

- Describes core-halo ratio Csörgő, Lörstad, Zimányi, Z.Phys.C71 (1996)  
Bolz et al, Phys.Rev.D47 (1993) 3860-3870
- Comparing with SPS and RHIC results: - also see talk of **D. Kincses**
  - ▶ Low  $m_T$  values show no decrease (sim. to my previous and other SPS results)
  - ▶ Halo component increases at RHIC (e.g. In-medium mass mod.)  
S. E. Vance et al, Phys.Rev.Lett.81 (1998) 2205-2208  
T. Csörgő et al, Phys.Rev.Lett.105 (2010) 182301  
A. Adare for PHENIX Collaboration, Phys.Rev.C97 (2018) no.6, 064911
- $\lambda$  value shows no  $m_T$  dependence



# Lévy-stability index $\alpha$ vs. $m_T$

- Lévy-stability index  $\alpha$ : shape of spatial correlation
- Compatible with Lévy assumption, far from CP ( $\alpha = 0.5$ )
  - ▶ Be+Be: far from Gaussian ( $\alpha = 2$ ), close to Cauchy ( $\alpha = 1$ )
  - ▶ Ar+Sc: far from Cauchy, close to Gaussian
- NA61/SHINE:  $\alpha \approx 0.9 - 1.5$ ,  $\sqrt{s_{NN}} = 16.82$  GeV, Be+Be, Eur.Phys.J.C 83 (2023) 10, 919
- PHENIX:  $\alpha = 1.2$   $\sqrt{s_{NN}} = 200$  GeV, Au+Au, Phys.Rev.C97 (2018) no.6, 064911
- CMS:  $\alpha = 1.6-1.9$   $\sqrt{s_{NN}} = 5$  TeV, Pb+Pb arXiv:2306.11574





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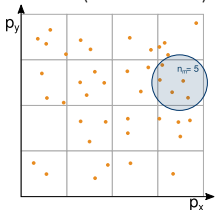
# Critical Point Search Using Intermittency

NA61/SHINE, arXiv:2305.07557 (accepted by EPJC), and this talk

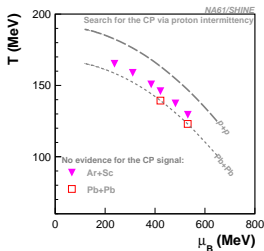
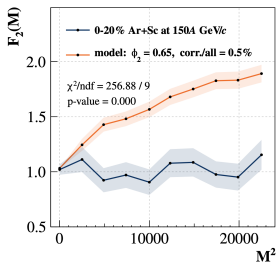
- Ar+Sc  $\sqrt{s_{NN}} \approx 5.1 - 16.8$  GeV, proton intermittency
- Pb+Pb  $\sqrt{s_{NN}} \approx 5.1, 7.6$  GeV, proton intermittency  
 $h^-$  intermittency

Scaled Factorial Moment:

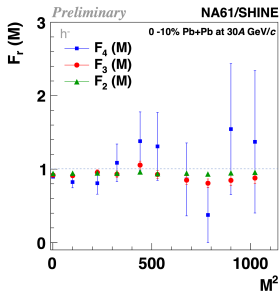
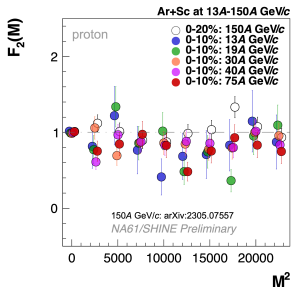
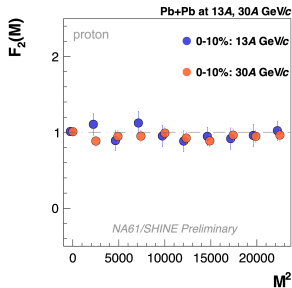
$$F_r(M) = \frac{\left\langle \frac{1}{M^2} \sum_{m=1}^{M^2} n_m(n_m-1)\dots(n_m-r+1) \right\rangle}{\left\langle \frac{1}{M^2} \sum_{m=1}^{M^2} n_m \right\rangle^r}$$



- Preserves power-law momentum correlation's power:  $F_r(M) \sim (M^2)^{\phi_r}$
- System freeze-out near Critical Point:  $F_2(M)$  power-law dependence,  $\phi_2 = 5/6$   
N. Antoniu et al., Phys.Rev.Lett.97 (2006) 032002



# Intermittency Results



Using cumulative  
transverse-momentum variable

A. Bialas and M. Gazdzicki Phys.Lett.B252 (1990)  
483–486

and statistically independent points

NA61/SHINE, arXiv:2305.07557 (accepted by EPJC)

- No structure indicating a power-law

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# Summary

- NA61/SHINE Lévy HBT analysis
  - ▶ 150A GeV/c beam momentum
  - ▶ Be+Be collisions, 0-20% centrality
  - ▶ Ar+Sc collisions, 0-10% centrality
- Measured momentum correlations of sum of like charged  $\pi$  pairs
- Fit done with correlation functions from symmetric Lévy source
- Parameter  $m_T$  dependence:
  - ▶  $\alpha(m_T)$ : Be+Be:  $\approx 0.9 - 1.5$ ; Ar+Sc:  $\approx 1.5 - 2.0 \rightarrow$  far from CP
  - ▶  $R(m_T)$ : visible  $m_T$  dependence - sign of transverse flow
  - ▶  $\lambda(m_T)$ : no dependence, no hole
- Symmetric Lévy source is a good assumption

Ongoing, Outlook:

- Energy scan with Ar+Sc; Larger system analysis

Intermittency analysis:

- No indication of Critical Point

Thank you for your attention!

# Backup

# Intermittency Summary

- Intermittency analysis measuring Scaled Factorial Moments in
  - ▶ Ar+Sc  $\sqrt{s_{NN}} \approx 5.1 - 16.8$  GeV, proton intermittency
  - ▶ Pb+Pb  $\sqrt{s_{NN}} \approx 5.1, 7.6$  GeV, proton intermittency, negatively charged hadrons ( $h^-$ ) intermittency
- Analysis was based on statistically independent data points
- Cumulative variables
  - ▶ remove dependence on the shape of the single-particle distribution
  - ▶ preserve critical behaviour
- No signal of a power-law increase was observed
  - ▶ in Second Scale Factorial Moment Ar+Sc, Pb+Pb proton intermittency,
  - ▶ in Second, Third and Fourth Scale Factorial Moment Pb+Pb negatively charged hadrons ( $h^-$ ) intermittency,on the cumulative momentum bin size



# Intermittency Methodology

In NA61/SHINE, intermittency analysis is performed at mid-rapidity, and particle fluctuations are studied in the transverse momentum plane.

$$F_r(M) = \frac{\left\langle \frac{1}{M^D} \sum_{m=1}^{M^D} n_m(n_m-1)\dots(n_m-r+1) \right\rangle}{\left\langle \frac{1}{M^D} \sum_{m=1}^{M^D} n_m \right\rangle^r}$$

- ▶  $M^D$ : number of equally sized cells in D-dimensional space
- ▶  $n_m$ : number of particles in  $m^{\text{th}}$  bin
- ▶  $\langle \dots \rangle$ : averaging over events
- At the second order phase transition, the system is a simple fractal, and the factorial moment exhibits a power law dependence:  $F_r(M) = F_r(\Delta) \cdot (M^D)^{\phi_q}$
- $\phi_r$  are predicted to follow the pattern:  $D \cdot \phi_r = (r-1) \cdot d_r$ ,  
where  $d_r$ : anomalous fractal dimension

Wosiek, Acta Phys.Polon.B19, 863 (1988)

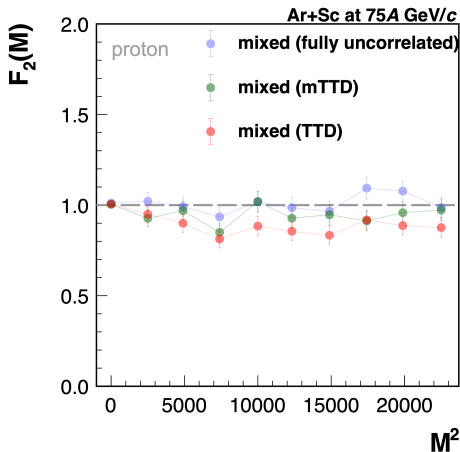
Bialas, Hwa, Phys.Lett.B253 (1991) 436

Bialas, Peschanski, NPB 273(1986) 703

# Intermittency - Effect of mTTD cut

Possible explanation:

anti-correlation by momentum based Two-Track Distance cut  
necessary to account for close-in-space tracks

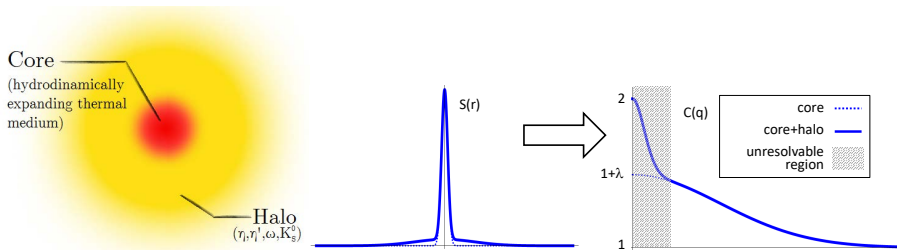


# Core-Halo Model

- Hydrodynamically increasing core  $\rightarrow$  pion emission
- Results in two component source:  $S(x) = S_M(x) + S_G(x)$
- Core  $\cong$  10 fm size, halo( $\omega, \eta \dots$ )  $>$  50 fm size
- Halo not seen due to detector resolution
- Real  $q \rightarrow 0$ , at  $C(q = 0) = 2$
- Results show  $C(q \rightarrow 0) = 1 + \lambda$ , where  $\lambda = \left( \frac{N_m}{N_g + N_m} \right)^2$

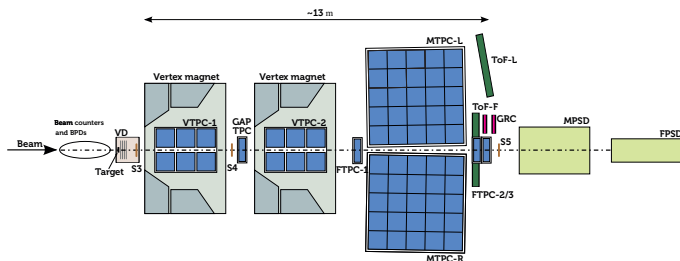
Bolz et al, Phys.Rev.D47 (1993) 3860-3870

Csörgő, Lörstad, Zimányi, Z.Phys.C71 (1996) 491-497



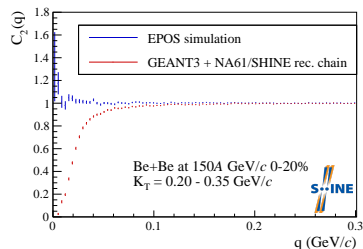
# The NA61/SHINE Detector Post LS2

- Upgrade of DAQ + new trigger system (TDAQ)
  - ▶ Detector readouts replaced → data taking rate up by 20x
  - ▶ TPCs - ALICE; other detectors - DRS4
- Construction of:
  - ▶ Vertex Detector - open-charm measurements
  - ▶ ToF-F wall
  - ▶ Multi-gap Resistive Plate Chamber based ToF-L (ToF-R under constr.)
  - ▶ Beam Position Detector
  - ▶ Geometry Reference Chamber - drift velocity measurements
- Upgrade of PSD to MPSD + FPSD



# Low-q Behavior

- B-E and Coul. effect not present in EPOS sim.
  - ▶  $C_2(q) \approx \text{const.}$
- Low-q range behavior in data:
  - ▶ Fits overestimate data
  - ▶ Theor. corr. func. cannot describe
  - ▶ Observed in Be+Be, Ar+Sc
- Strong cutoff observable
  - ▶ Several possibilities...
  - ▶ Might be experimental artefact?
- Visible deviation from generated (simulated)
  - ▶ Effects such as track merging present
- Low-q region (until reconstructed  $\approx 1$ ) can be excluded
  - ▶ Two Track Distance cut not needed



# Projectile Spectator Detector

- Centrality measurement with PSD
- Located on beam axis
- measures forward energy ( $E_F$ ) from spectators
- Intervals in  $E_F$  allows to select centrality classes

