

CORRELATIONS BETWEEN ANISOTROPY FLOW AND MEAN TRANSVERSE MOMENTUM USING SUBEVENT CUMULANTS IN SMALL SYSTEMS AT CMS

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In heavy ion collisions:

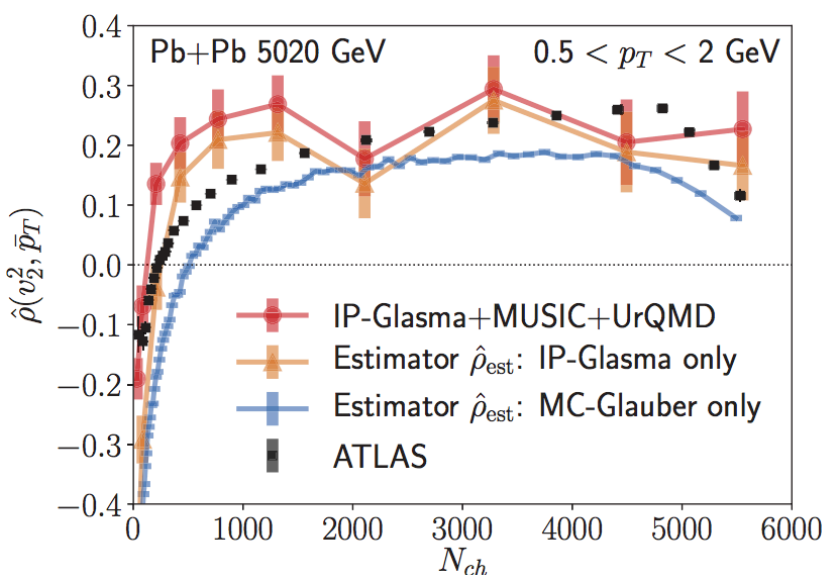
- Anisotropic flow in the final state arises from the hydrodynamic response to the initial geometric anisotropy
- Mean p_T ($[p_T]$) reflects the strength of radial flow push, which is related to the initial energy density of the fireball
- The correlations between v_n and $[p_T]$ probe the fluctuations present in the initial density profile

$$\rho(v_n^2, [p_T]) = \frac{\text{cov}(v_n^2, [p_T])}{\sqrt{\text{var}(v_n^2)}\sqrt{\text{var}([p_T])}}$$

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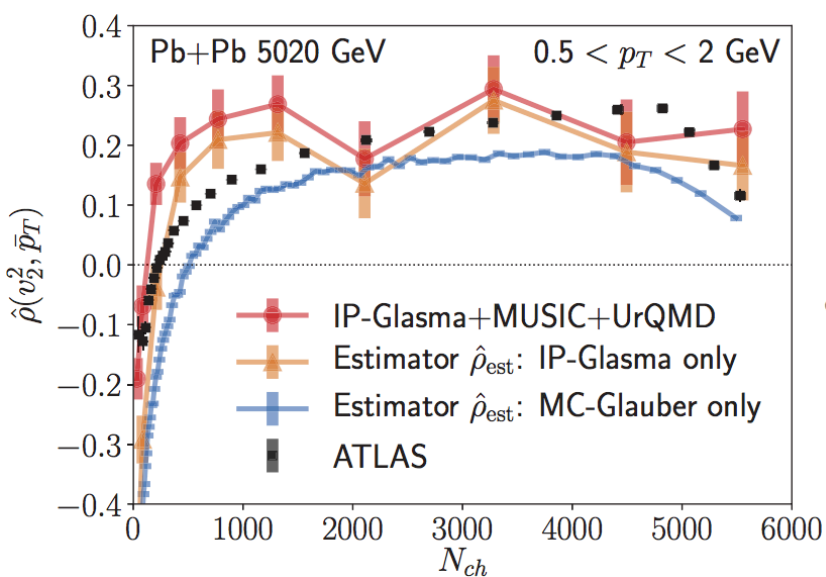
Phys. Rev. C 102, 034905 (2020)

Sensitive to the degree of sub-nucleon fluctuations

Introduction

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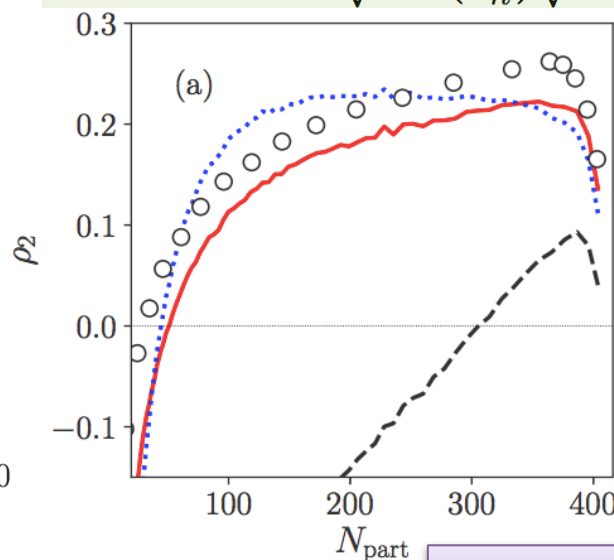
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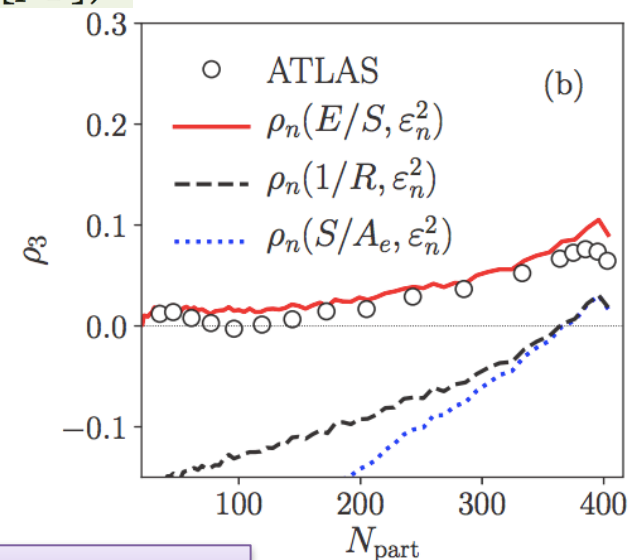
Phys. Rev. C 102, 034905 (2020)

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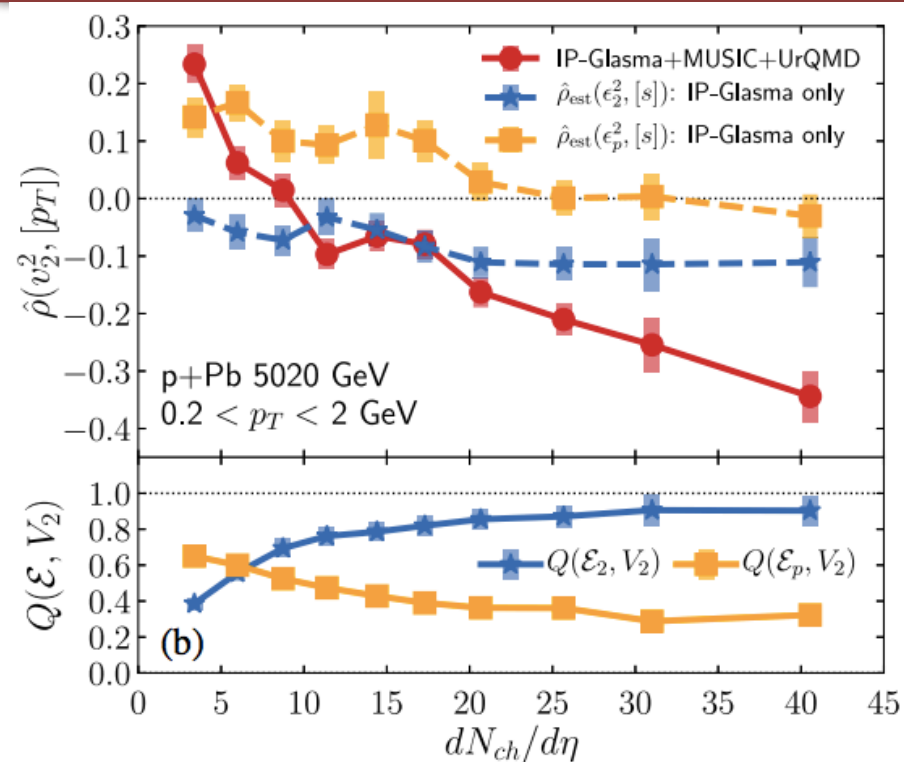


Phys. Rev. C 103, 024909 (2021)

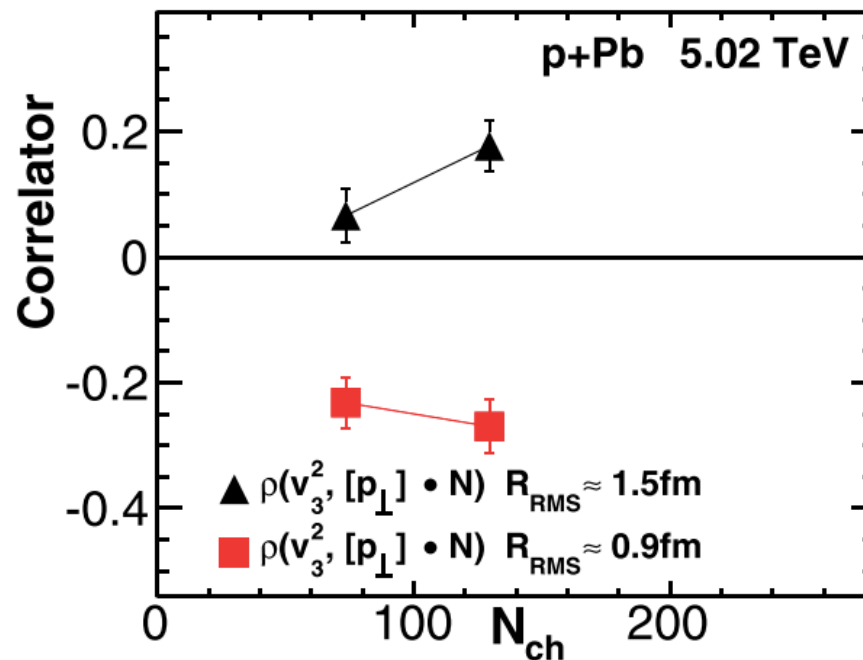


The mechanism driving the correlations can be traced back to the initial density profile

Motivations for small systems (1)



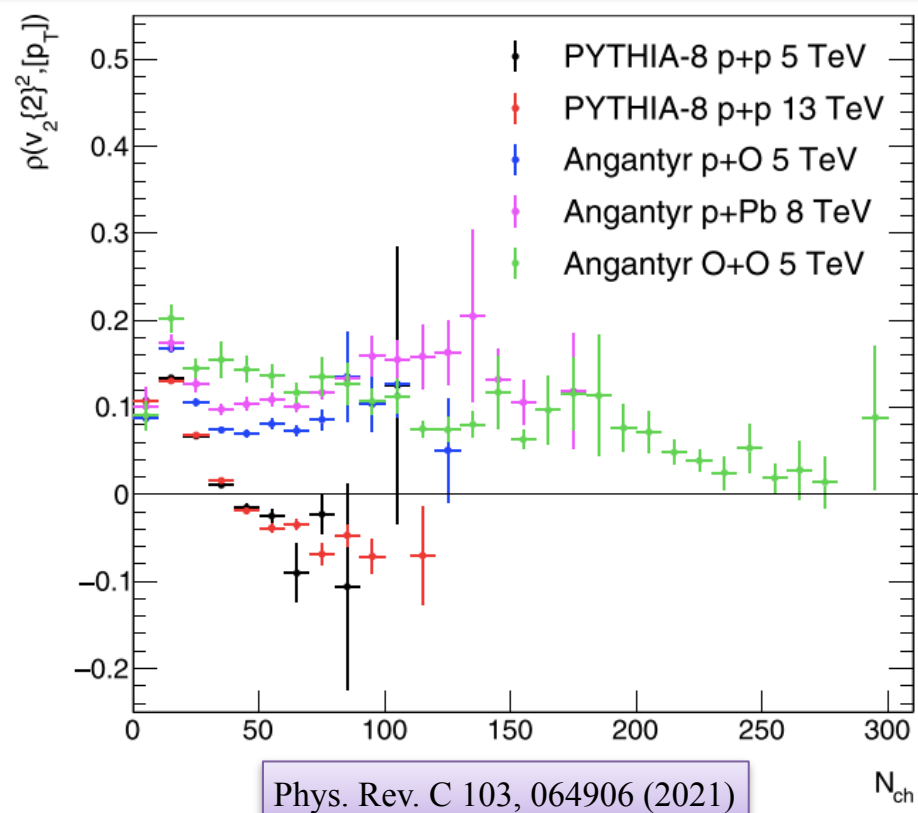
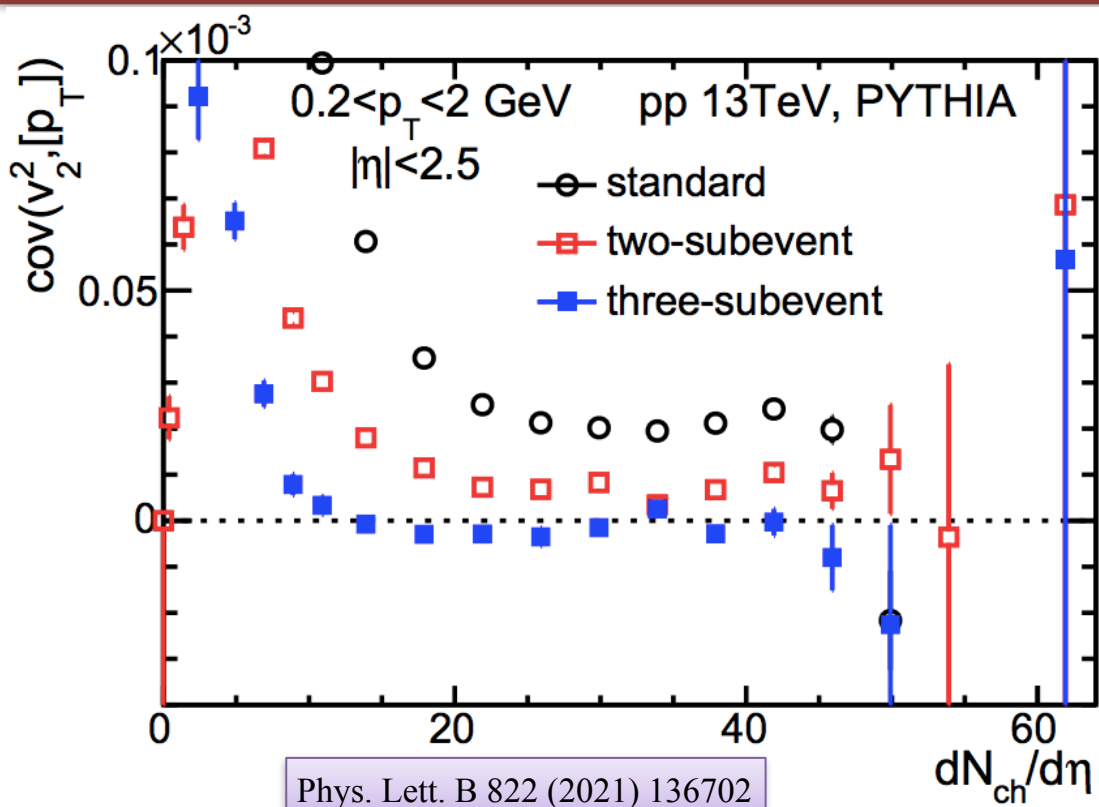
Phys. Rev. Lett. 125, 192301 (2020)



Phys. Rev. C 101, 064902 (2020)

- $\hat{\rho}$ exhibit a sign change with increasing multiplicity when an initial momentum anisotropy is present
- Carries information about the origin of the observed momentum anisotropy
- Sensitive to the transverse size of the initial fireball

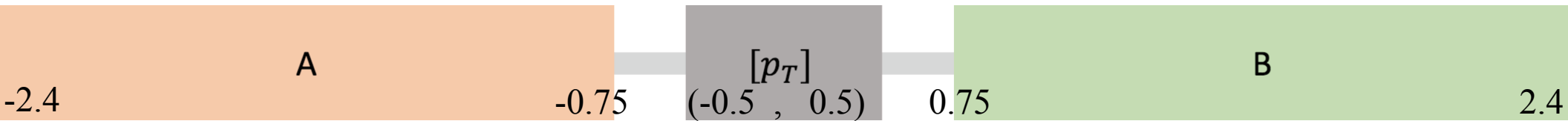
Motivations for small systems (2)



- However, the correlations from PYTHIA alone have a sign change
- Nonflow has to be carefully dealt
- The goal of this analysis:
 - Introduce a new variable to remove more nonflow
 - Search for sign change at the lowest possible multiplicity in pp/pPb/PbPb collisions

$$\rho(v_2^2, [p_T]) = \frac{\text{cov}(v_2^2, [p_T])}{\sqrt{\text{Var}(v_2^2)_{dyn}} \sqrt{\text{Var}([p_T])_{dyn}}} \quad (1)$$

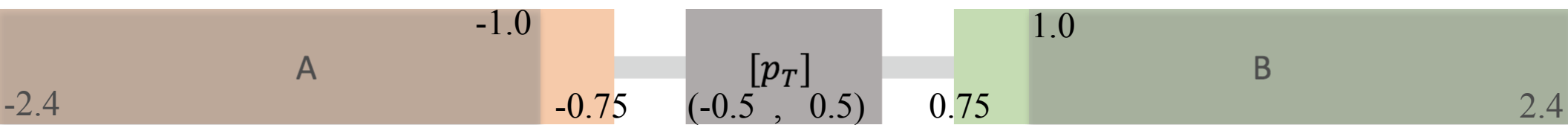
$\rho(c_2\{2\}, [p_T])$



Subevents A and B are used for calculation of $c_2\{2\}$; $|\eta| < 0.5$ for $[p_T]$

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$\rho(c_2\{2\}, [p_T])$



Subevents A and B are used for calculation of $c_2\{2\}$; $|\eta| < 0.5$ for $[p_T]$

Covariance between $c_2\{2\}$ and $[p_T]$:

$$\text{cov}(c_2\{2\}, [p_T]) = \Re \left\langle \sum_{a,b} \exp^{2i(\phi_a - \phi_b)} ([p_T] - \langle [p_T] \rangle) \right\rangle \quad (2)$$

Dynamic variance of $c_2\{2\}$:

$$\text{Var}(c_2\{2\})_{dyn} = \langle \langle 4 \rangle \rangle - \langle \langle 2 \rangle \rangle^2 \quad (3)$$

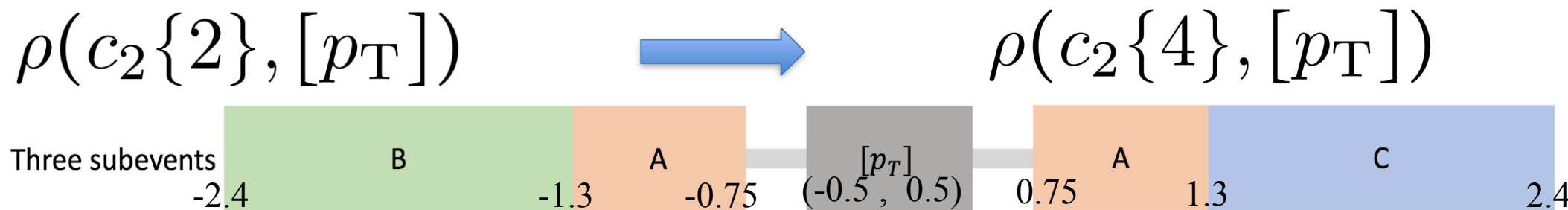
Variance of $[p_T]$ from dynamic $[p_T]$ fluctuation c_k :

$$c_k = \left\langle \left[(p_{Ti} - \langle [p_T] \rangle) (p_{Tj} - \langle [p_T] \rangle) \right] \right\rangle \quad (4)$$

$$\rho(v_2^2, [p_T]) = \frac{\text{COV}(v_2^2, [p_T])}{\sqrt{\text{Var}(v_2^2)_{dyn}} \sqrt{\text{Var}([p_T])_{dyn}}} \quad (1)$$



Extend and study the new variable to remove more nonflow



$c_2\{4\}$ is analyzed with three subevent method

$$c_2\{4\}_{3\text{-sub}} = \langle 4 \rangle_{a,a|b,c} - 2\langle 2 \rangle_{a|b} \langle 2 \rangle_{a|c}$$

Phys. Rev. C 96, 034906 (2017)

- This analysis focuses on small systems
- It is the first paper to :
 - use multiparticle correlations for flow when correlating with $[p_T]$
 - explore the correlator with different η gaps to study nonflow effects
 - measure $v_3 - [p_T]$ correlations in small systems
 - include measurements in pp collisions

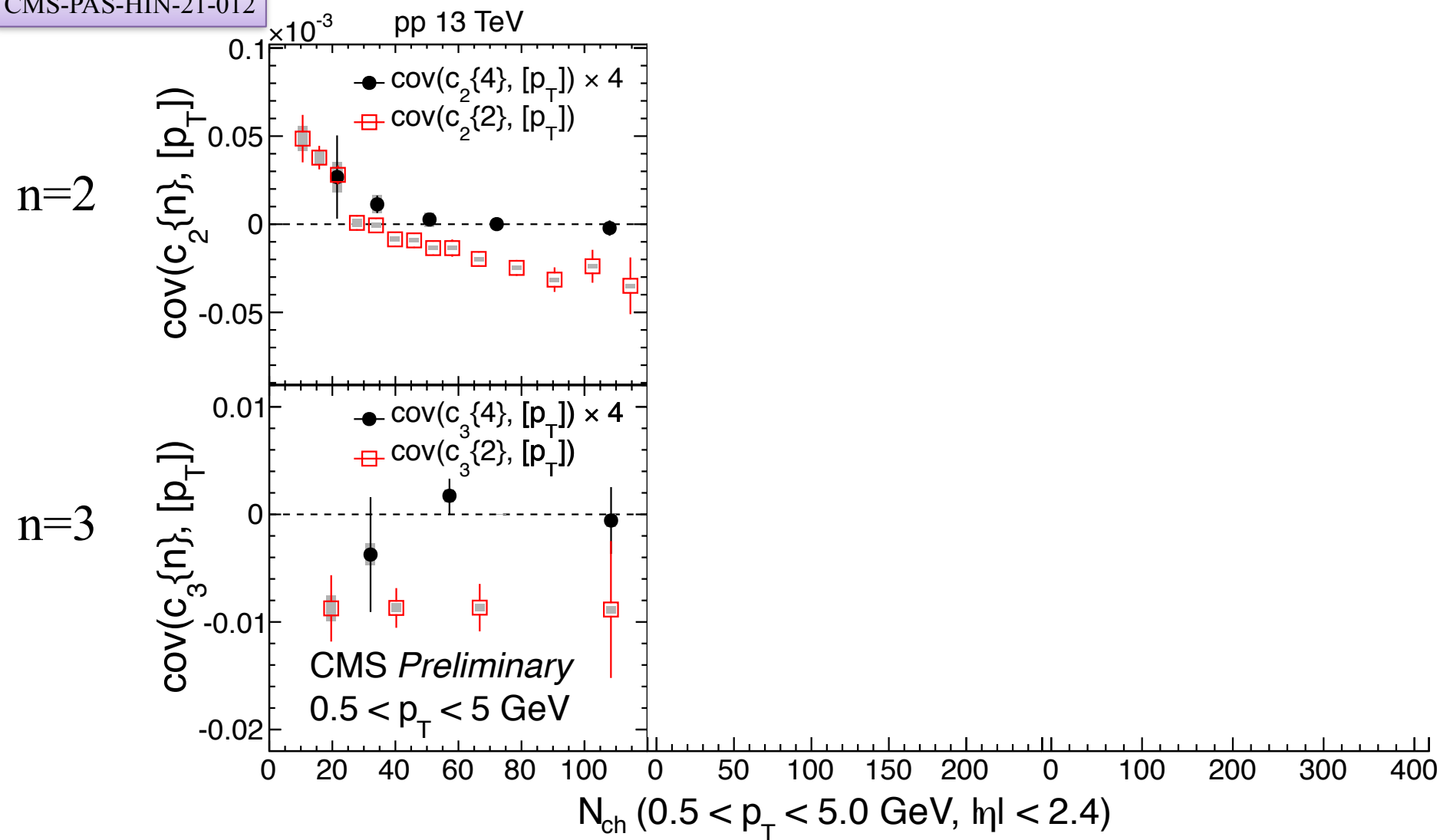
Results

Results are presented as a function of N_{ch} , which is defined in $0.5 < p_{\text{T}} < 5.0$ GeV, $|\eta| < 2.4$, and corrected for tracking efficiency

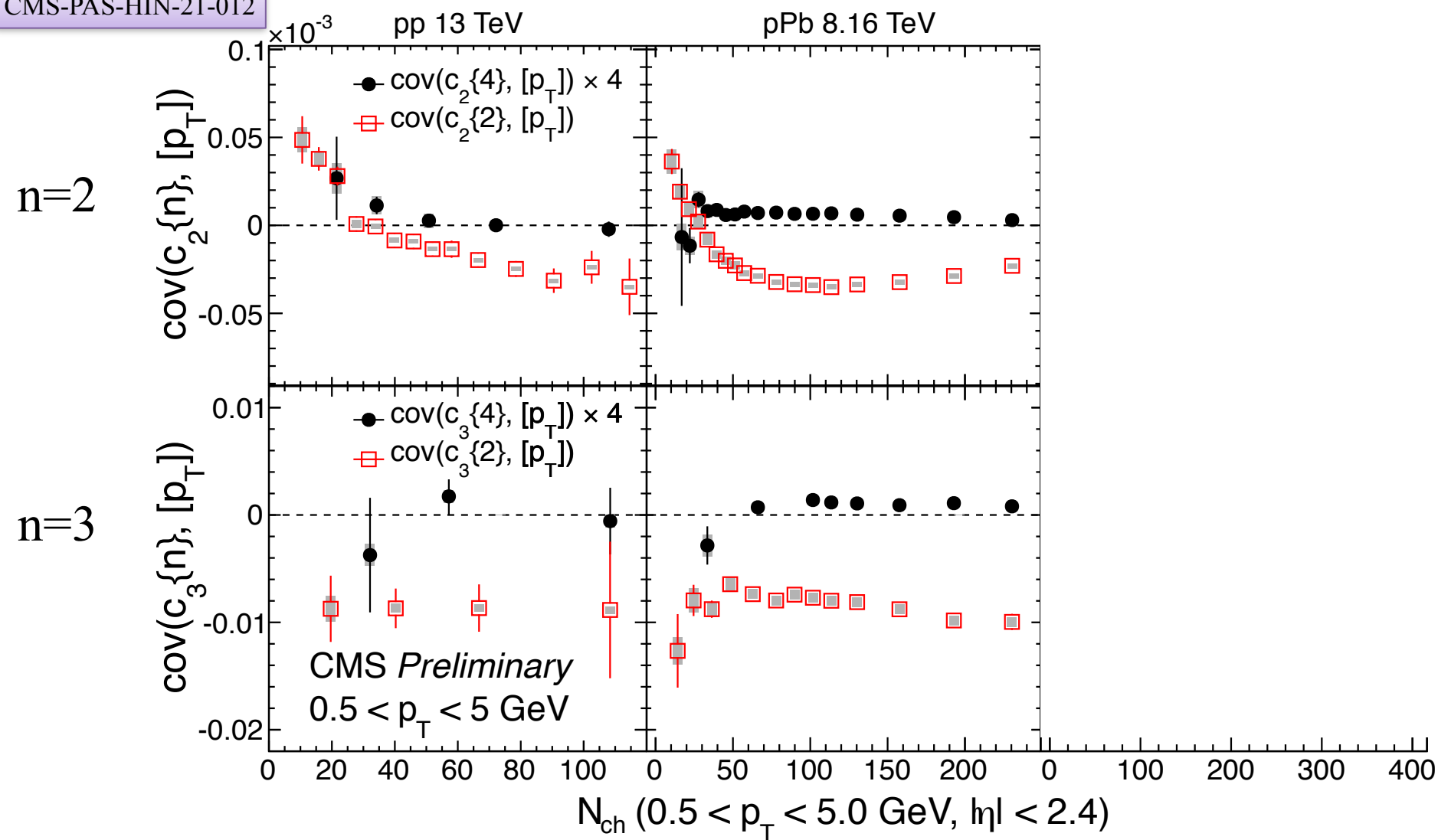
CMS-PAS-HIN-21-012 ([link](#))

Results for covariance

CMS-PAS-HIN-21-012



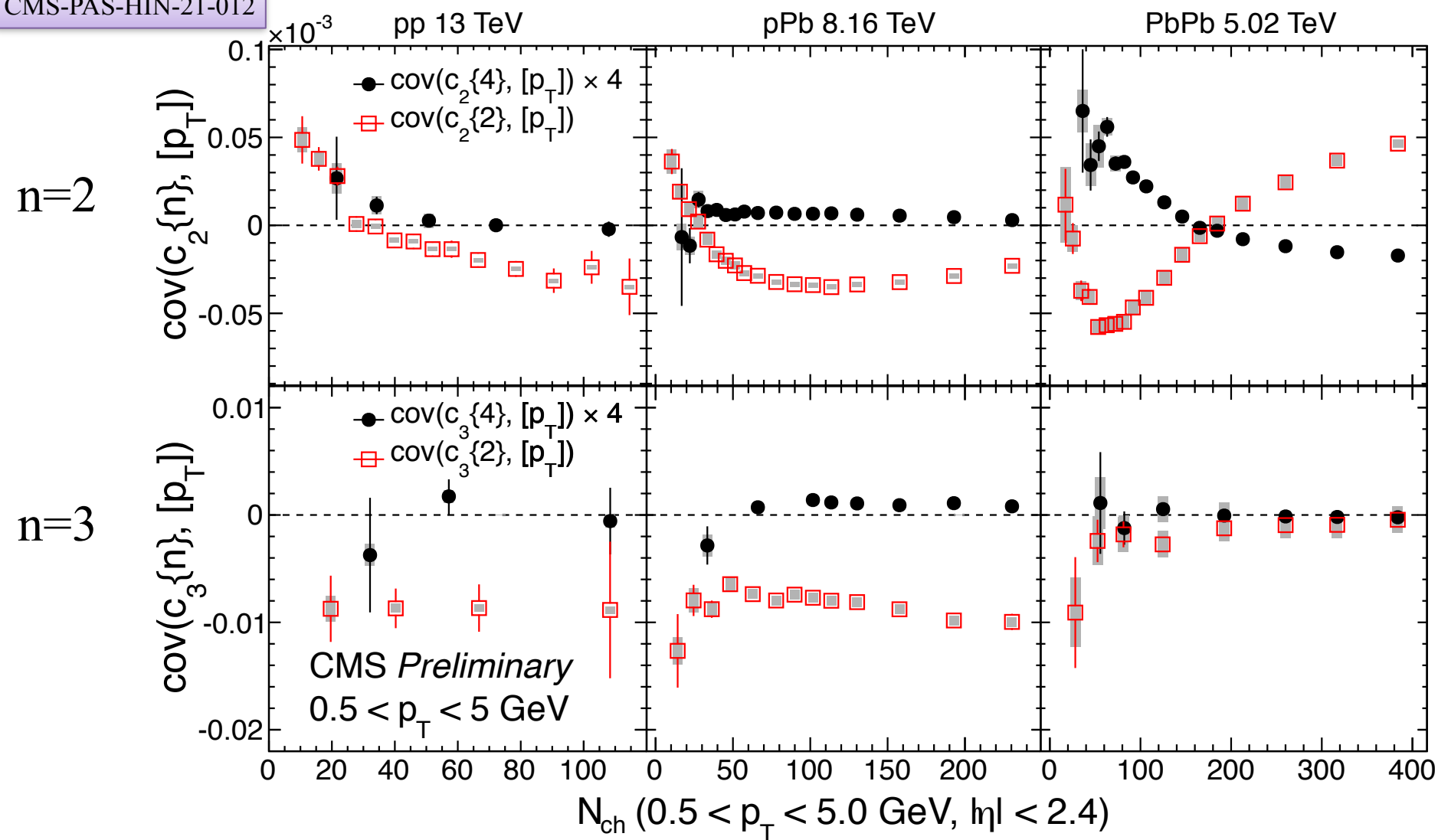
- Clear sign change for pp collisions with $c_2\{2\}$
- No sign change at low N_{ch} using multiparticle correlations with current statistics
- The sign of the normalized correlator is determined by the covariance



- Clear sign change for pp and pPb collisions with $c_2\{2\}$
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Results for covariance

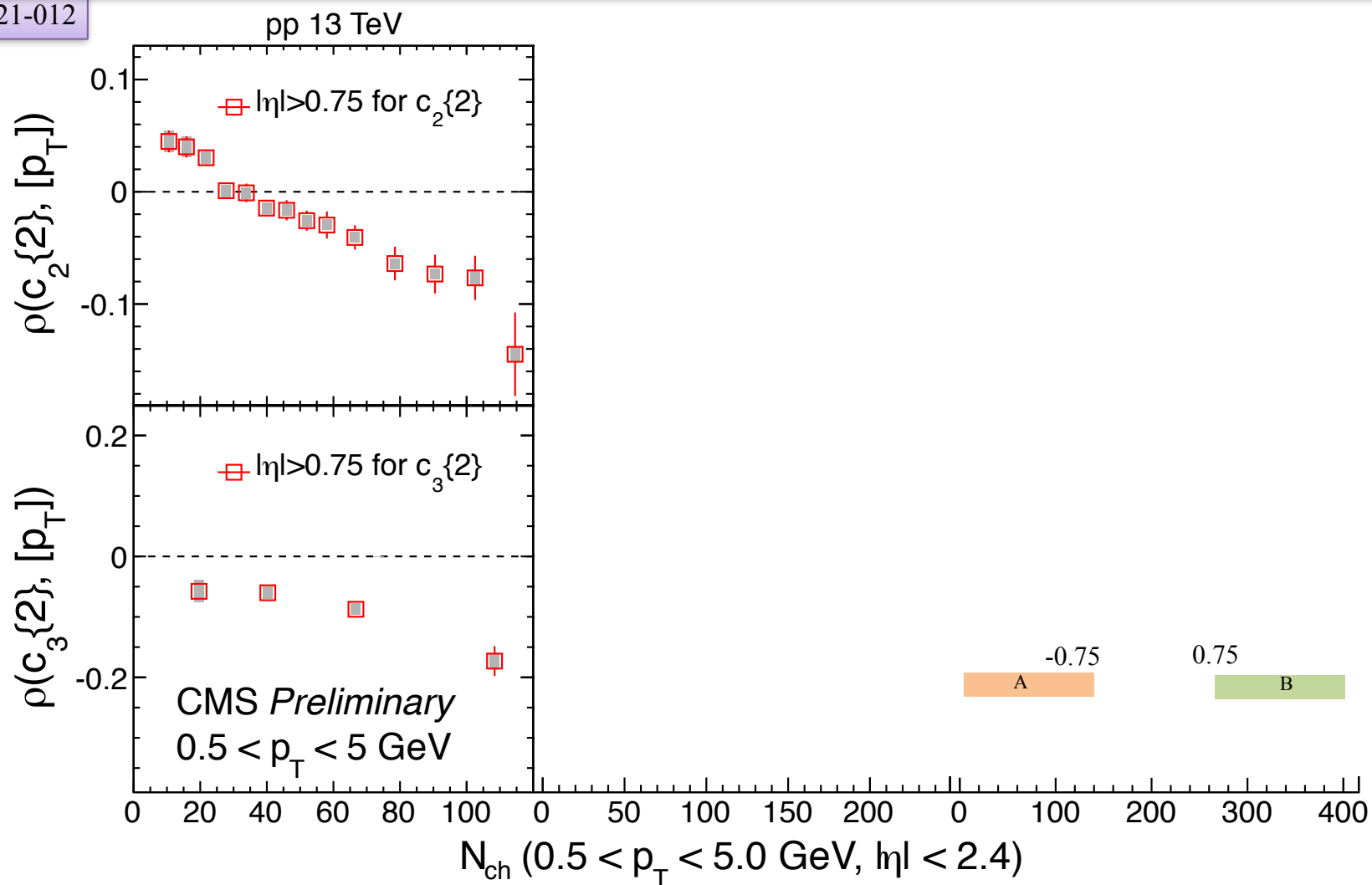
CMS-PAS-HIN-21-012



- Clear sign change for pp and pPb collisions with $c_2\{2\}$
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n=2

n=3



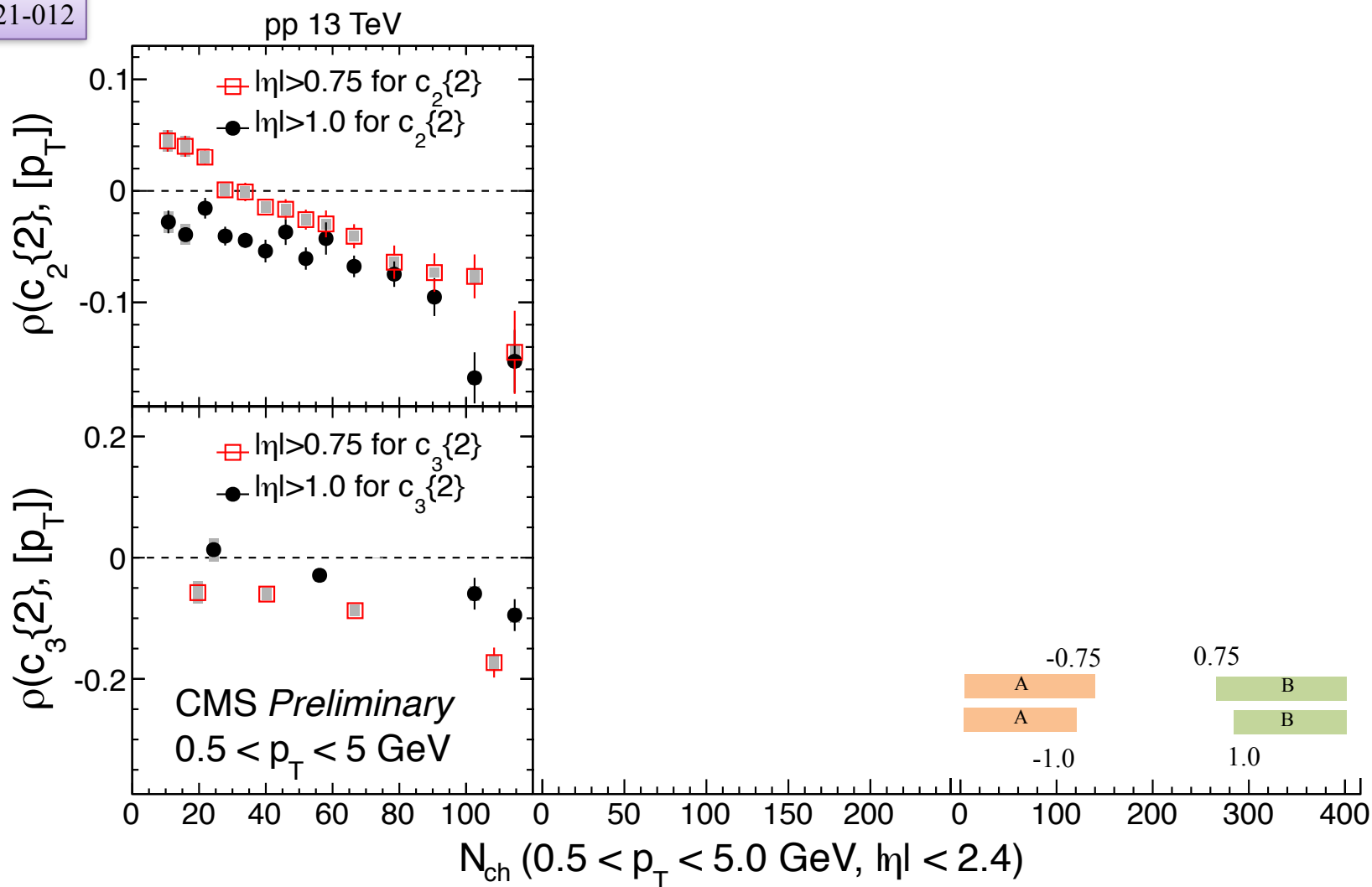
- Apparent sign change for $\rho(c_2\{2\}, [p_T])$ in pp collisions

Results for the correlator

CMS-PAS-HIN-21-012

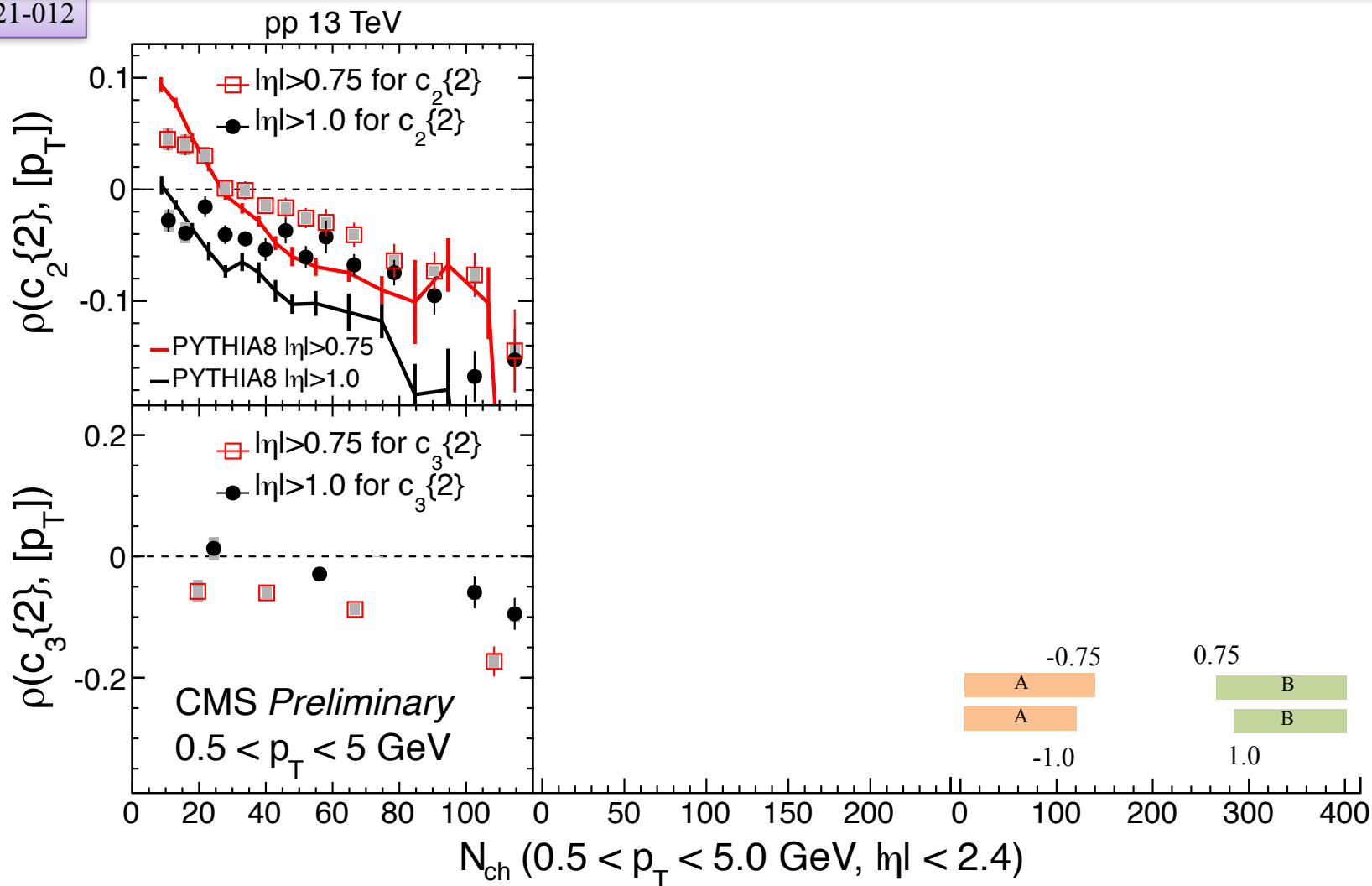
n=2

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- Apparent sign change for $\rho(c_2\{2\}, [p_T])$ in pp collisions
- However, no sign change is observed when using $|\eta| > 1.0$ for $c_2\{2\}$

n=2



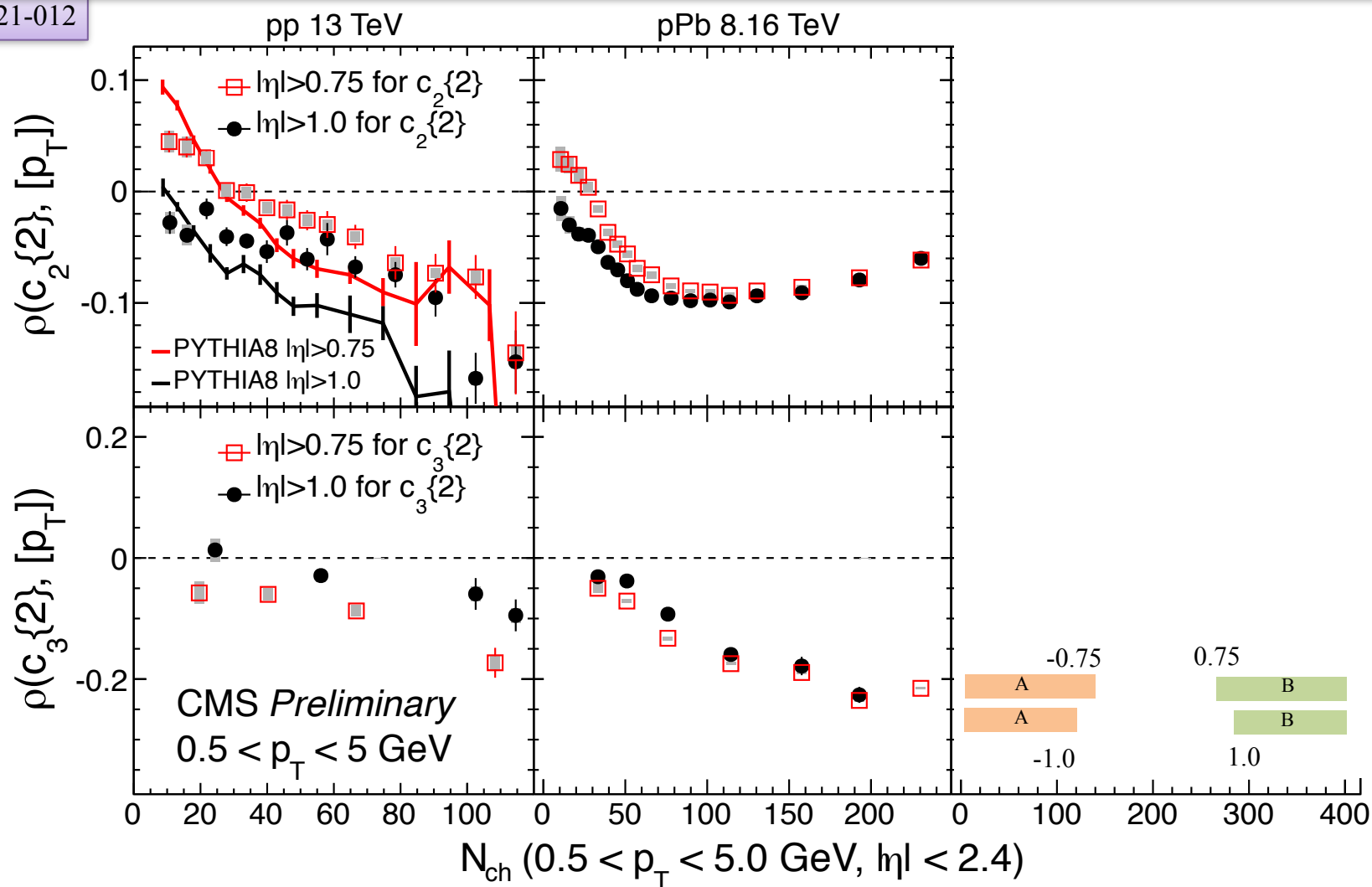
- Apparent sign change for $\rho(c_2^{2}, [p_T])$ in pp collisions
- However, no sign change is observed when using $|\eta| > 1.0$ for c_2^{2}
- Also true for PYTHIA8 events

Results for the correlator

CMS-PAS-HIN-21-012

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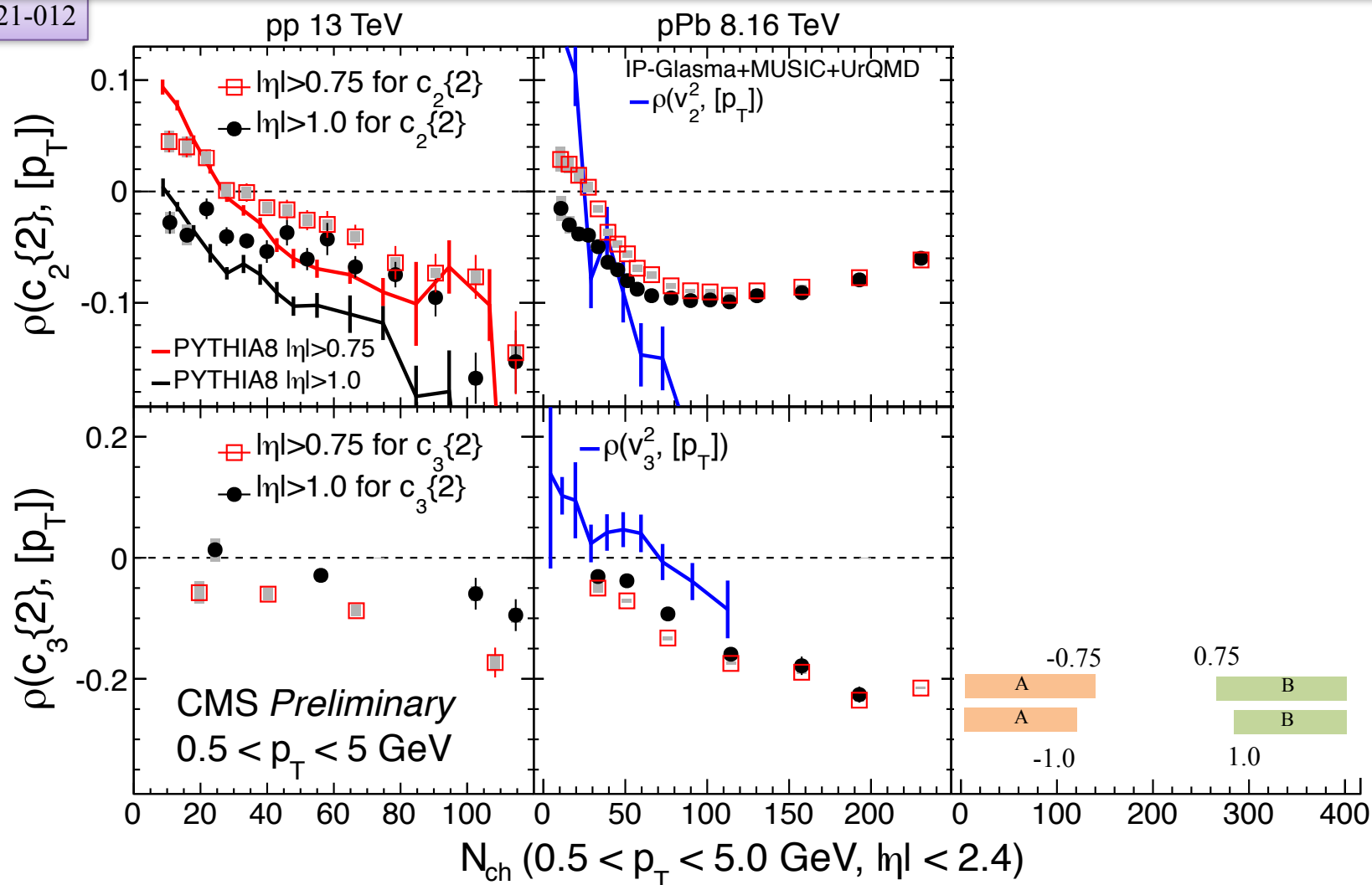
- Apparent sign change for $\rho(c_2\{2\}, [p_T])$ in pPb -> agree with IP-Glasma+hydro
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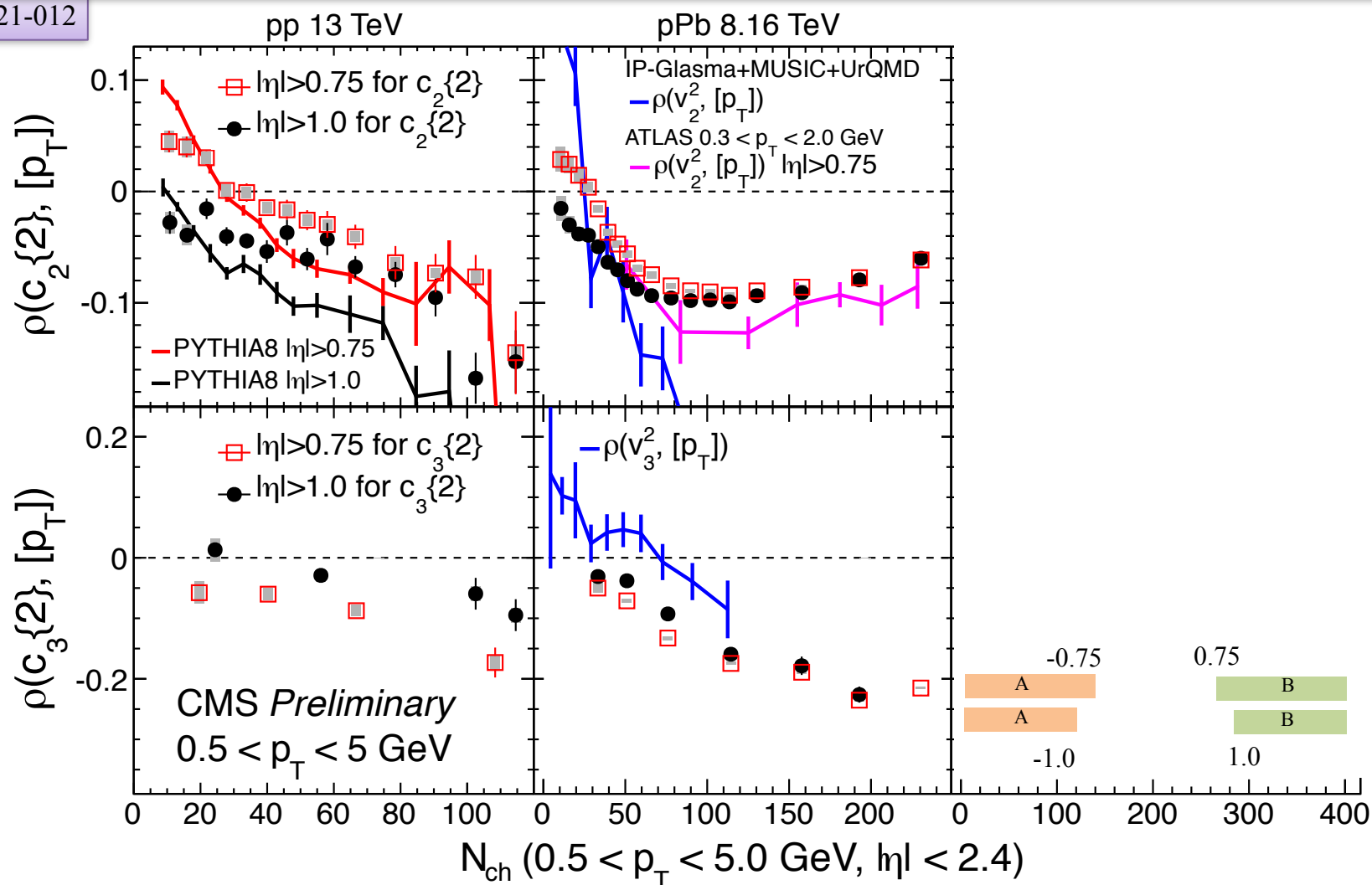
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- Apparent sign change for $\rho(c_2\{2\}, [p_T])$ in pPb \rightarrow agree with IP-Glasma+hydro
- However, no sign change is observed when using $|\eta| > 1.0$ for $c_2\{2\}$
- After removing nonflow with larger η gap, no evidence of CGC in data

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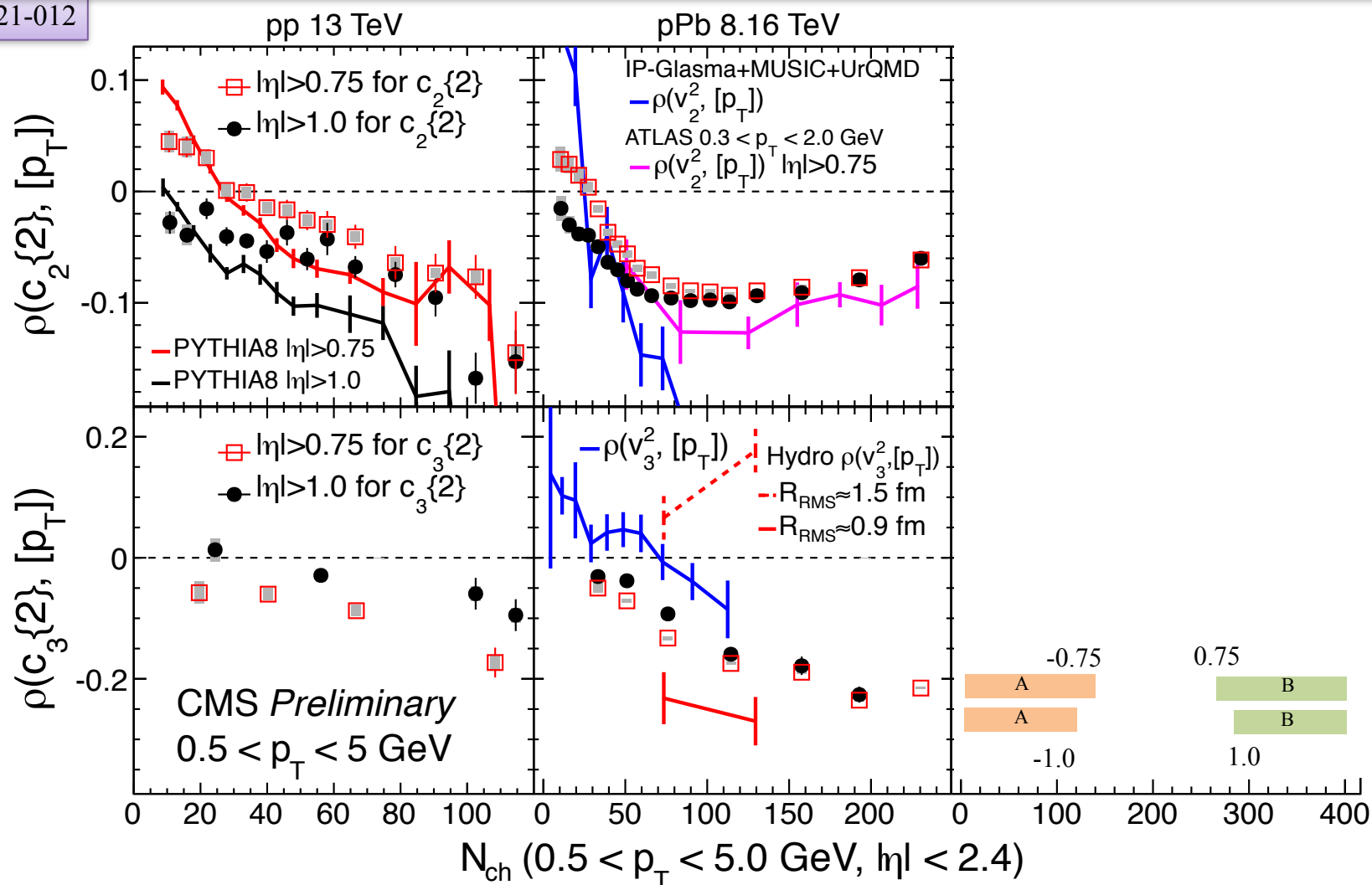
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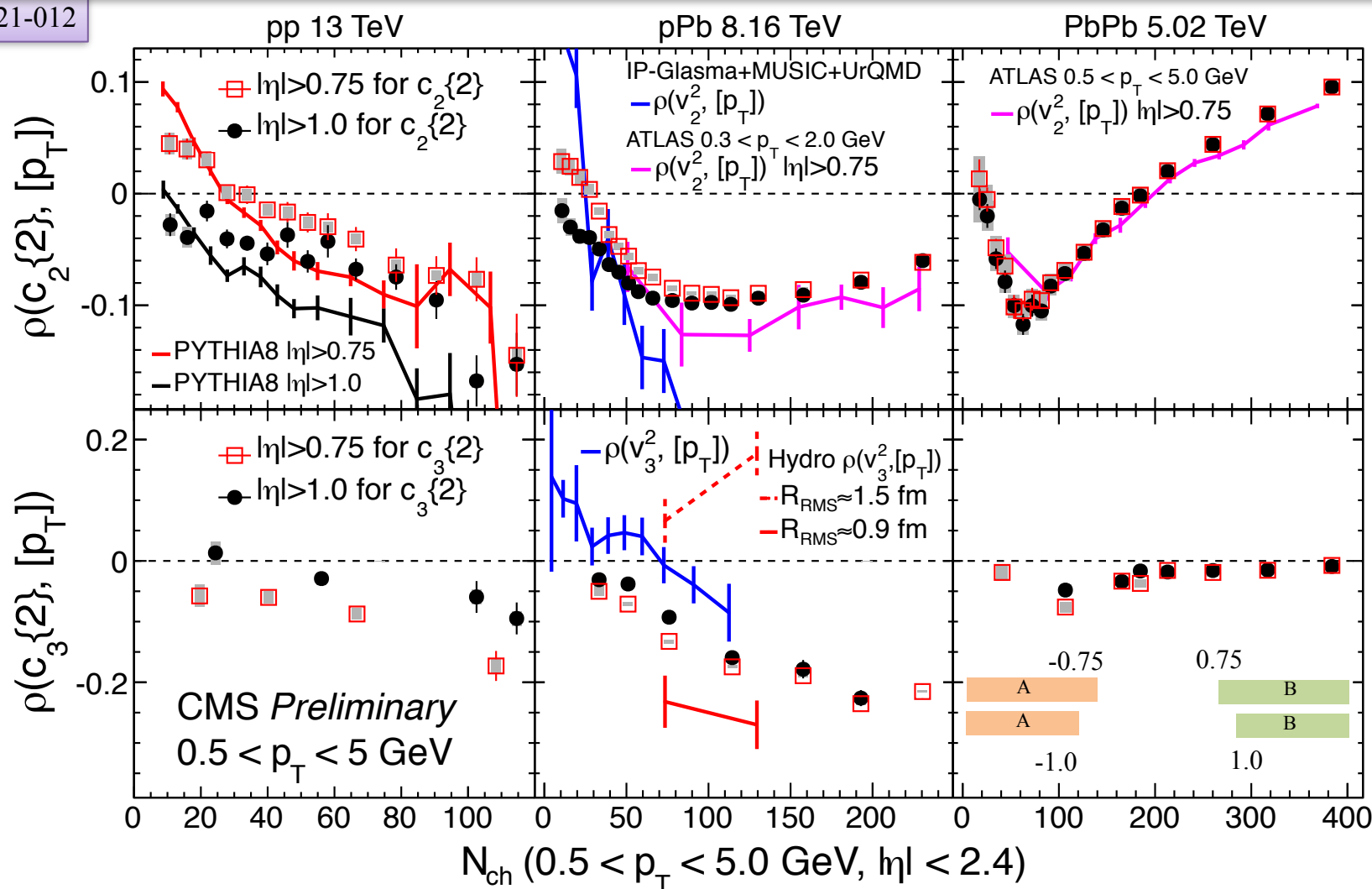


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- Data better described by the smaller initial fireball $R_{RMS} = 0.9 \text{ fm}$ in hydro

Results for the correlator

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- Apparent sign change for $\rho(c_2\{2\}, [p_T])$ in pPb \rightarrow agree with IP-Glasma+hydro
- However, no sign change is observed when using $|\eta| > 1.0$ for $c_2\{2\}$
- After removing nonflow with larger η gap, no evidence of CGC in data
- Data better described by the smaller initial fireball $R_{RMS} = 0.9 \text{ fm}$ in hydro

- The correlations between $[p_T]$ and cumulants from both two- and four-particle correlations in small systems are presented
- Apparent sign change is observed for $\rho(c_2\{2\}, [p_T])$ in pp and pPb
- However, no sign change is observed with larger η gap in $c_2\{2\}$
 - ATLAS default is $|\eta|>0.75$, with $|\Delta\eta|>1.5$
 - ALICE default is $|\eta|>0.4$, with $|\Delta\eta|>0.8$
 - CMS is studying both $|\eta|>0.75$ ($|\Delta\eta|>1.5$) and $|\eta|>1.0$ ($|\Delta\eta|>2.0$)
- After removing more nonflow with both two- and four-particle correlation cumulants, there is no evidence of CGC in data
- These high-precision data and the observables employing multiparticle correlations should provide new insight into the origin of azimuthal anisotropy in small collision systems

Keeping $\text{cov}(c_2\{4\}, [p_T])$ but drop $\rho(c_2\{4\}, [p_T])$ in this analysis

- The reason is we are not 100% sure if the variance $\text{Var}(c_2\{4\})_{\text{dyn}}$ in our new method is truly dynamic
- It may contain statistical fluctuations in our current method
- The measurement of v_n fluctuation in small systems is a task our community has not accomplished. The event-by-event v_n studies all stopped at 60-70% centrality in AA collisions

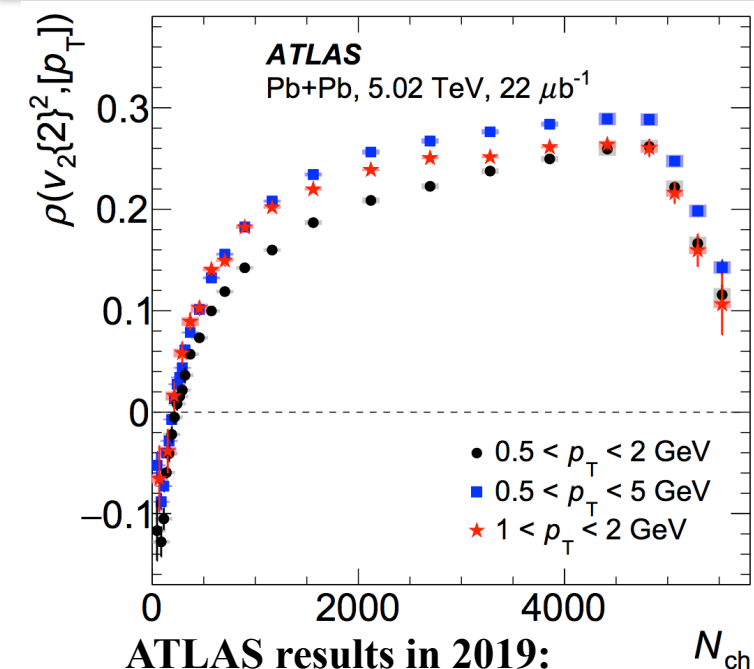
Table 1: Average multiplicity of reconstructed tracks per $N_{\text{ch}}^{\text{rec}}$ bin for N_{ch} and $N_{\text{trk}}^{\text{offline}}$ in pp, pPb, and PbPb collisions. Uncertainties for the tracking efficiency corrected N_{ch} are included.

$N_{\text{ch}}^{\text{rec}}$ range	pp		pPb		PbPb	
	$\langle N_{\text{ch}} \rangle$	$\langle N_{\text{trk}}^{\text{offline}} \rangle$	$\langle N_{\text{ch}} \rangle$	$\langle N_{\text{trk}}^{\text{offline}} \rangle$	$\langle N_{\text{ch}} \rangle$	$\langle N_{\text{trk}}^{\text{offline}} \rangle$
[0, 20)	8 ± 0.3	9	11 ± 0.4	12	16 ± 0.6	14
[20, 40)	34 ± 1	34	36 ± 1	36	57 ± 2	48
[40, 60)	58 ± 2	56	60 ± 2	60	96 ± 4	80
[60, 80)	82 ± 3	78	83 ± 3	82	135 ± 5	112
[80, 100)	106 ± 4	101	107 ± 4	105	175 ± 7	144
[100, 150)	132 ± 5	125	140 ± 6	137	240 ± 10	197
[150, 200)			198 ± 8	191	335 ± 13	276
[200, 250)			256 ± 10	246	434 ± 17	353
[250, 300)					535 ± 21	426

Phys. Lett. B 718 (2013) 795

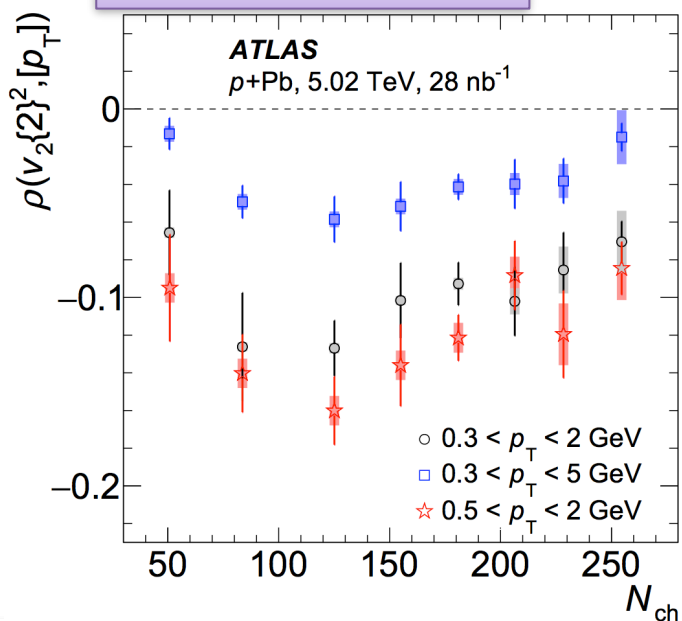
- The mapping table between N_{ch} and $N_{\text{trk}}^{\text{offline}}$

Existing measurements

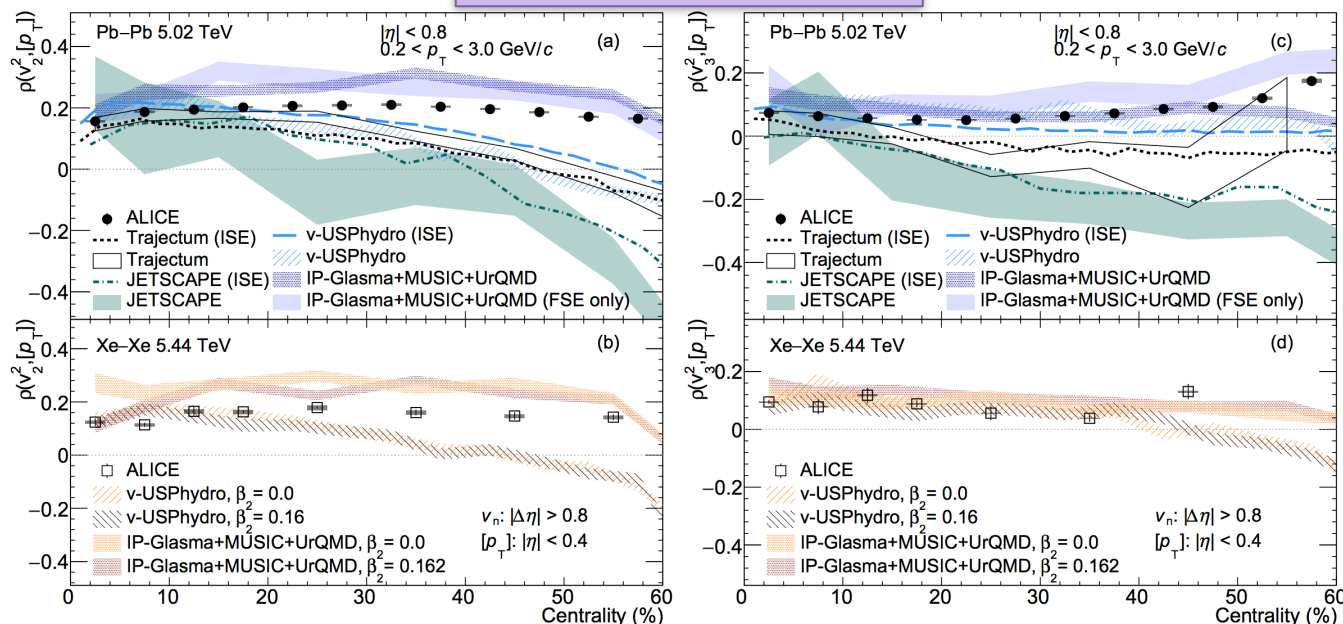


ATLAS results in 2019:

Eur. Phys. J. C 79 (2019) 985



ALICE: Phys.Lett.B 834 (2022) 137393



- ATLAS results in pPb, PbPb, and XeXe collisions
- ALICE results in PbPb and XeXe collisions
- Some recent studies from STAR