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XVI Workshop on  
Particle Correlations and Femtoscopy  
& IV Resonance Workshop

**WPCF 2023**

Catania | 6-10 Nov 2023

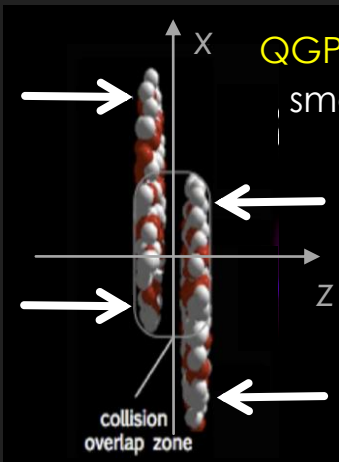
# Non-equilibrium effects and spherocity in relativistic proton-nucleus collisions



Lucia Oliva

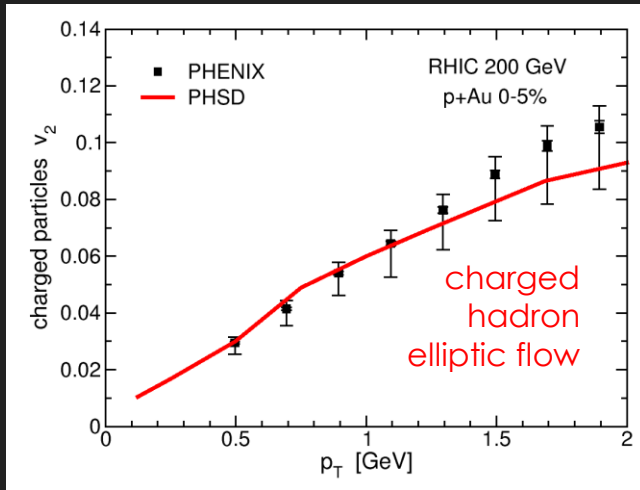
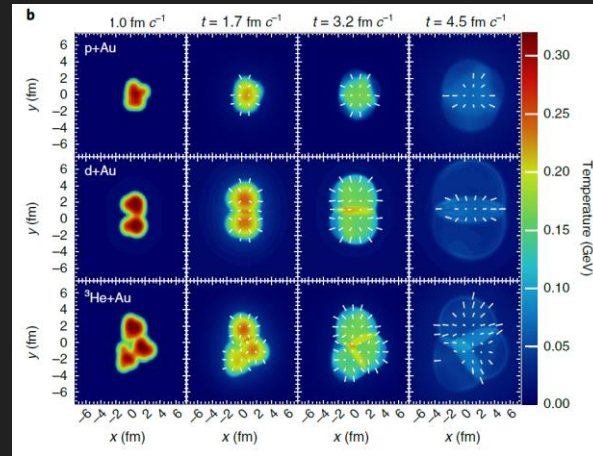
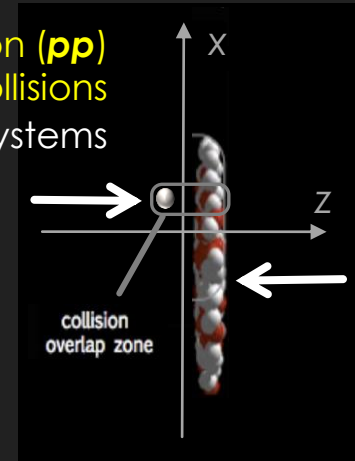
in collaboration with  
Wenkai Fan, Pierre Moreau, Steffen Bass,  
Elena Bratkovskaya, Vadim Voronyuk





QGP initially expected only in high energy collisions of two heavy ions (**AA**)  
 small colliding systems regarded as control measurements

Signatures of collective flow found in proton-proton (**pp**)  
 and proton-nucleus (**pA**) collisions  
 in high-multiplicity events of small systems



PHENIX Coll.,  
 Nature Phys. 15, 214 (2019)

creation of short-lived QGP  
 droplets in small systems

LO, P. Moreau, V. Voronyuk and E. Bratkovskaya,  
 Phys. Rev. C 101, 014917 (2020)

# Simulating large and small colliding systems

Two main approaches to describe the dynamics of hot QCD medium in relativistic nuclear collisions

## HYDRODYNAMIC MODELS

macroscopic description  
evolution based on conservation laws  
unreasonable effectiveness of hydro

## TRANSPORT MODELS

microscopic description  
evolution of particle distribution functions  
inherent inclusion of nonequilibrium

energy-momentum tensor

viscous corrections

$$T^{\mu\nu} = e u^\mu u^\nu - \Delta^{\mu\nu} (P + \Pi) + \pi^{\mu\nu}$$

$$\partial_\mu T^{\mu\nu} = 0$$

fluid 4-velocity

collision integral

$$(p^\mu \partial_\mu + g Q F^{\mu\nu} p_\nu \partial_\mu^p) f(x, p) = \mathcal{C}[f]$$

field interaction

one-particle distribution function

# Simulating large and small colliding systems

Two main approaches to describe the dynamics of hot QCD medium in relativistic nuclear collisions

Initial conditions +

## HYDRODYNAMIC MODELS

+ hadronic afterburner (transport)

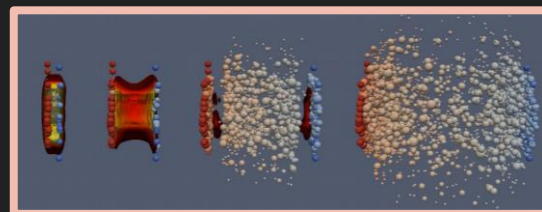
macroscopic description

evolution based on conservation laws  
unreasonable effectiveness of hydro

## TRANSPORT MODELS

microscopic description

evolution of particle distribution functions  
inherent inclusion of nonequilibrium  
suitable for the pre-equilibrium stage,  
for partonic and hadronic phases



Picture credit: MADAI Collaboration [🔗](#)



Picture credit: PHSD project [🔗](#)

# Simulating large and small colliding systems

Two main approaches to describe the dynamics of hot QCD medium in relativistic nuclear collisions

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## TRANSPORT MODELS

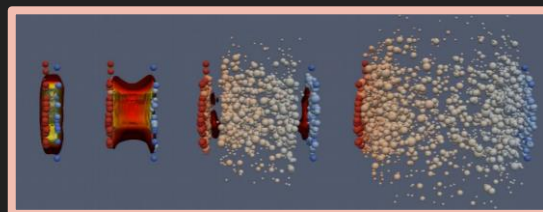
microscopic description

evolution of particle distribution functions

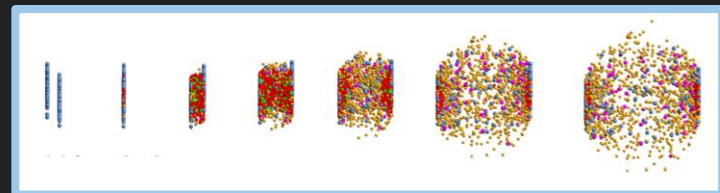
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Picture credit: MADAI Collaboration 



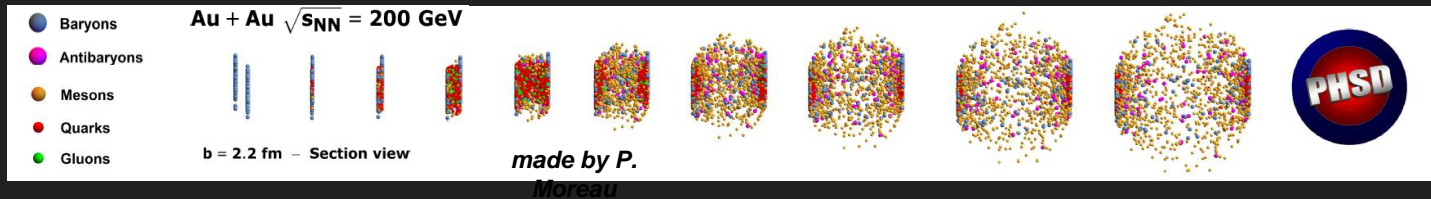
Picture credit: PHSD project 

Both **hybrid** and **transport** approaches are successful in describing AA and pA  
Different way of treating nonequilibrium effects

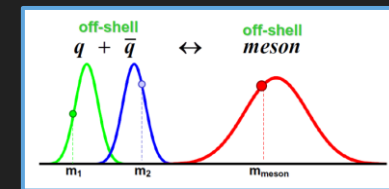
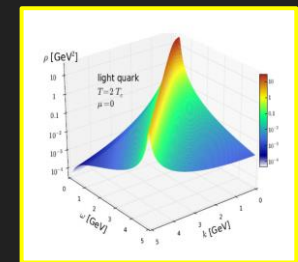
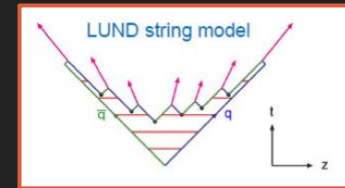
# Parton-Hadron-String Dynamics – PHSD

non-equilibrium off-shell transport approach

to study the phase transition from hadronic to partonic matter and the QGP properties from a microscopic origin



- **INITIAL NUCLEI COLLISION:** nucleon-nucleon collisions lead to the formation of strings that decay to pre-hadrons
- **FORMATION OF QGP:** if  $\epsilon > \epsilon_c$  pre-hadrons dissolve in massive off-shell quarks and gluons + mean-field potential
- **PARTONIC STAGE:** off-shell transport equations with Dynamical Quasi-Particle Model (DQPM)
- **HADRONIZATION:** partons with broad spectral functions hadronize to off-shell baryons and mesons
- **HADRONIC PHASE:** off-shell transport equations with hadron-hadron interactions



W. Cassing and E. Bratkovskaya, Phys. Rev. C 78, 034919 (2008) [\[1\]](#); Nucl. Phys. A 831, 215 (2009) [\[2\]](#)  
 P. Moreau, O. Soloveva, LO, T. Song et al., Phys. Rev. C 100, 014911 (2019) [\[3\]](#); PHSD website [\[4\]](#)

# VISHNew (+ hadronic afterburner)

## 2+1D viscous hydrodynamic model

to study the QGP medium and its properties from a macroscopic point of view

$$\partial_\mu T^{\mu\nu} = 0$$

space-time evolution of the QGP  
via conservation equations  
of the energy-momentum tensor

$$T^{\mu\nu} = e u^\mu u^\nu - \Delta^{\mu\nu} (P + \Pi) + \pi^{\mu\nu}$$

hydro equations closed by  
an equation of state  $P=P(e)$   
(lattice QCD + HRG)

time evolution of the viscous corrections  
via 2<sup>nd</sup> order Israel-Stewart equations

$$\begin{aligned} \tau_\Pi \dot{\Pi} + \Pi &= -\zeta\theta - \delta_{\Pi\Pi}\Pi\theta + \phi_1\Pi^2 \\ &\quad + \lambda_{\Pi\pi}\pi^{\mu\nu}\sigma_{\mu\nu} + \phi_3\pi^{\mu\nu}\pi_{\mu\nu} \\ \tau_\pi \dot{\pi}^{\langle\mu\nu\rangle} + \pi^{\mu\nu} &= 2\eta\sigma^{\mu\nu} + 2\pi_\alpha^{\langle\mu}w^{\nu\rangle\alpha} - \delta_{\pi\pi}\pi^{\mu\nu}\theta \\ &\quad + \phi_7\pi_\alpha^{\langle\mu}\pi^{\nu\rangle\alpha} - \tau_{\pi\pi}\pi_\alpha^{\langle\mu}\sigma^{\nu\rangle\alpha} \\ &\quad + \lambda_{\pi\Pi}\Pi\sigma^{\mu\nu} + \phi_6\Pi\pi^{\mu\nu} \end{aligned}$$

$e$ : local energy density  
 $P$ : local isotropic pressure  
 $\Pi$ : bulk viscous pressure  
 $\pi^{\mu\nu}$ : shear stress tensor

- **PARTONIC STAGE**: VISHNew
- **HADRONIZATION**: Cooper-Frye
- **HADRONIC PHASE**: UrQMD

H. Song and U. W. Heinz, Phys. Rev. C 77, 064901 (2008) [\[1\]](#); Phys. Rev. C 78, 024902 (2008) [\[2\]](#)  
Version of the Ohio State Uni code slightly modified by Duke Uni group [\[3\]](#)

# Coarse-graining PHSD medium

Study the hot medium evolution with transport and hydro

- same initial conditions to reduce the early stage impact
- similar equation of state and specific shear viscosity

Y. Xu et al., Phys. Rev. C 96, 024902 (2017) 

T. Song et al., Phys. Rev. C 101, 044903 (2020) 

LO et al., Phys. Rev. C 106, 044910 (2022) 

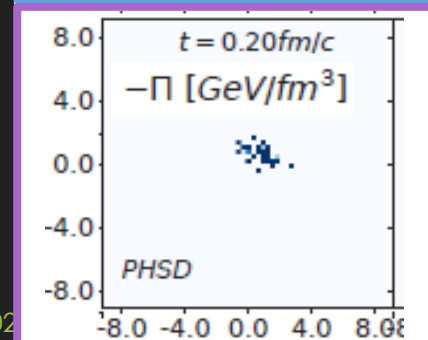
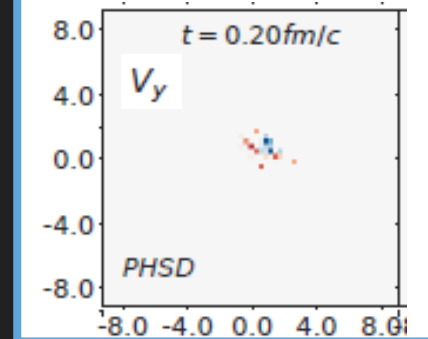
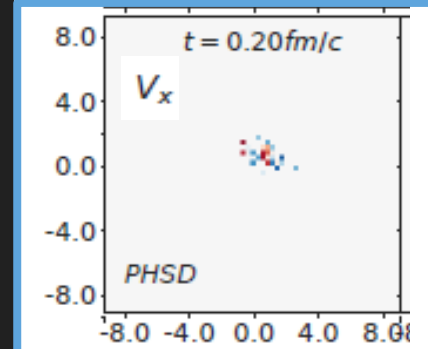
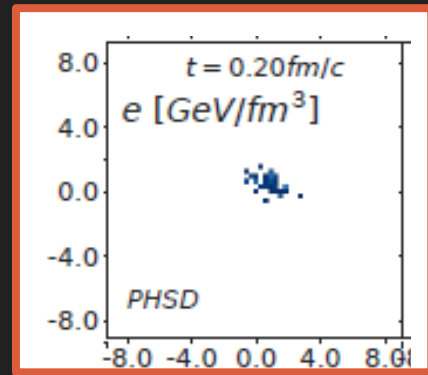
$$T^{\mu\nu} = e u^\mu u^\nu - \Delta^{\mu\nu} (P + \Pi) + \pi^{\mu\nu}$$

VISHNew  
vs  
PHSD

initialized  
with

$$T^{\mu\nu}(x) = \sum_i \int_0^\infty \frac{d^3 p_i}{(2\pi)^3} f_i(E_i) \frac{p_i^\mu p_i^\nu}{E_i} = \frac{1}{V} \sum_i \frac{p_i^\mu p_i^\nu}{E_i}$$

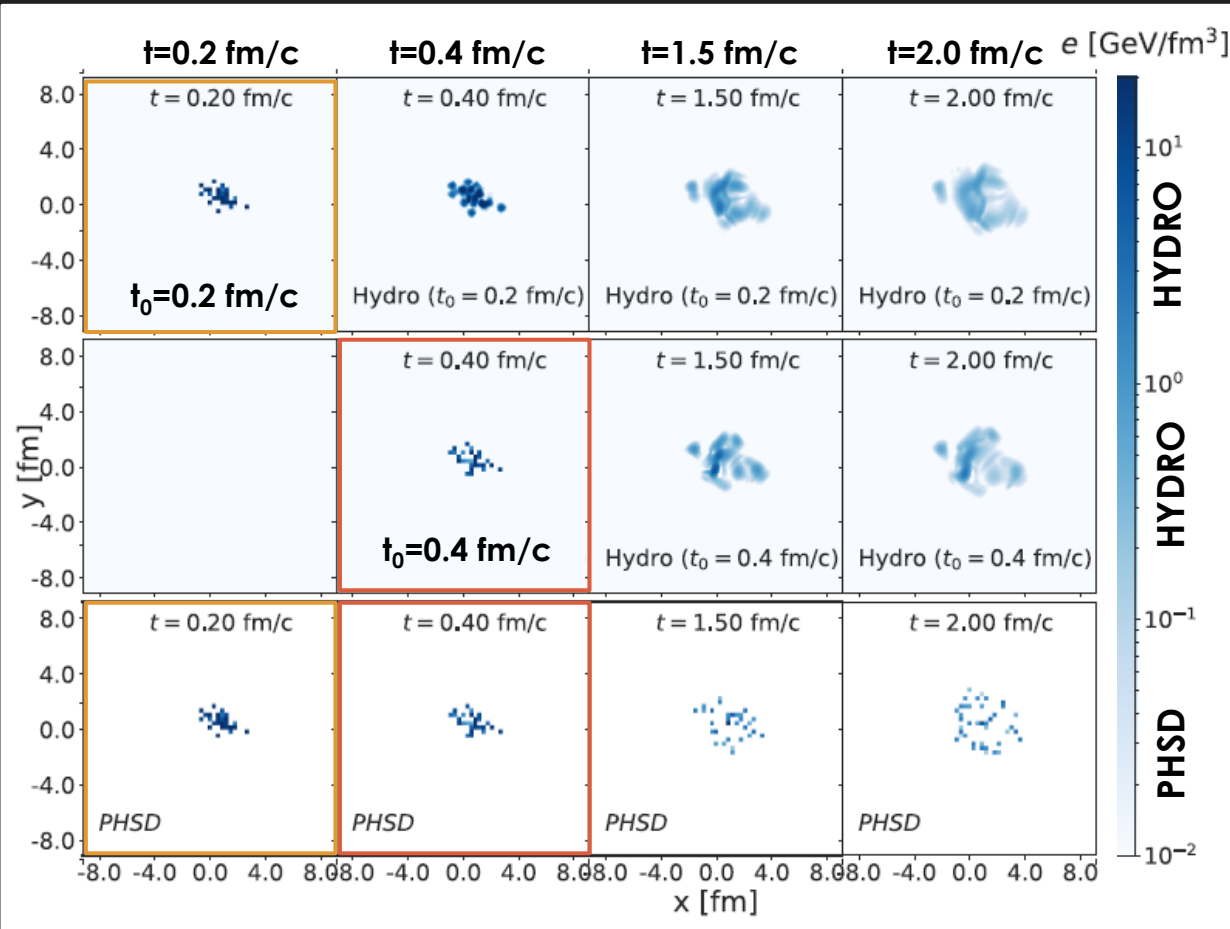
Hydro initialization time? 0 0.2 0.4 0.6 0.8 t [fm/c] 





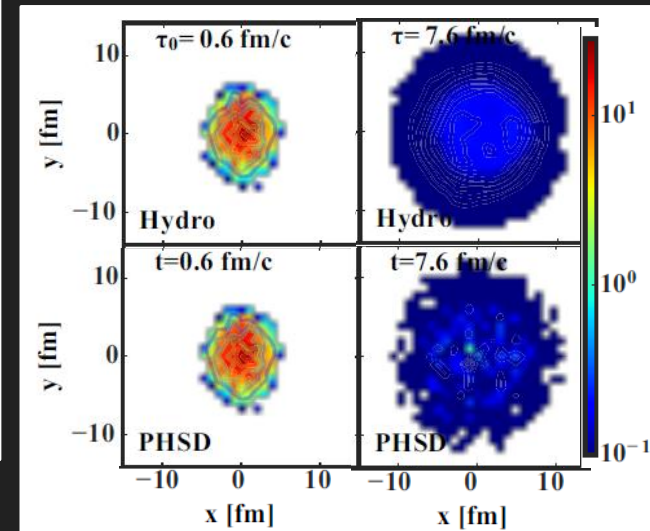
# Medium evolution: hydrodynamics vs PHSD

*p+Pb @ LHC 5.02 TeV –  $b = 2$  fm*



- ❖ PHSD evolution more chaotic all times
- ❖ VISHNew smooths the initial PHSD profile
- ❖ Visible impact of hydro initialization time

*Au+Au @ RHIC 200 GeV  
 $b = 6$  fm*



Higher degree of inhomogeneity in pA than in AA

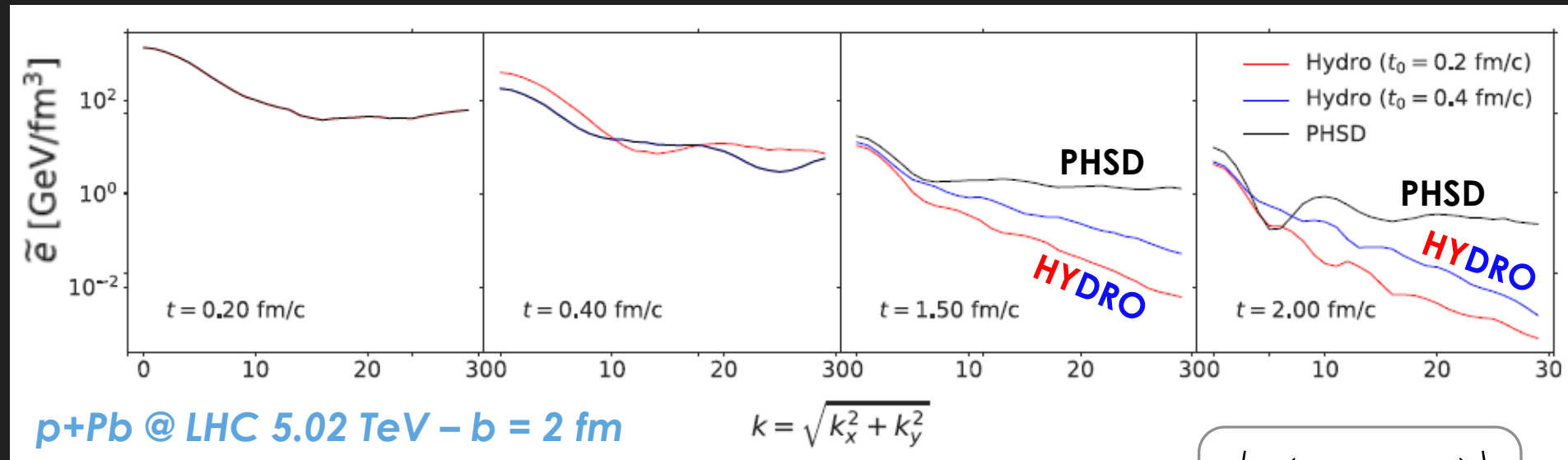
LO, W. Fan, P. Moreau, S.A. Bass and E. Bratkovskaya,  
Phys. Rev. C 106, 044910 (2022) [\[4\]](#)

Y. Xu et al.,  
Phys. Rev. C 96, 024902 (2017) [\[4\]](#)

# Quantifying medium inhomogeneity

Fourier transform of the energy density profile

$$\tilde{e}(k_x, k_y) = \frac{1}{m} \frac{1}{n} \sum_{x=0}^{m-1} \sum_{y=0}^{n-1} e(x, y) e^{2\pi i \left( \frac{xk_x}{L_x m} + \frac{yk_y}{L_y n} \right)}$$



**radial distribution of the Fourier modes of the energy density**

$$\left\langle \tilde{e} \left( \sqrt{k_x^2 + k_y^2} \right) \right\rangle$$

**Shorter wavelength modes survive only in PHSD**

- constant inhomogeneity of QGP medium in the microscopic description
- nonequilibrium dynamics able to preserve the medium irregularities

LO, W. Fan, P. Moreau, S.A. Bass and E. Bratkovskaya, Phys. Rev. C 106, 044910 (2022) [\[1\]](#)

# Viscous corrections

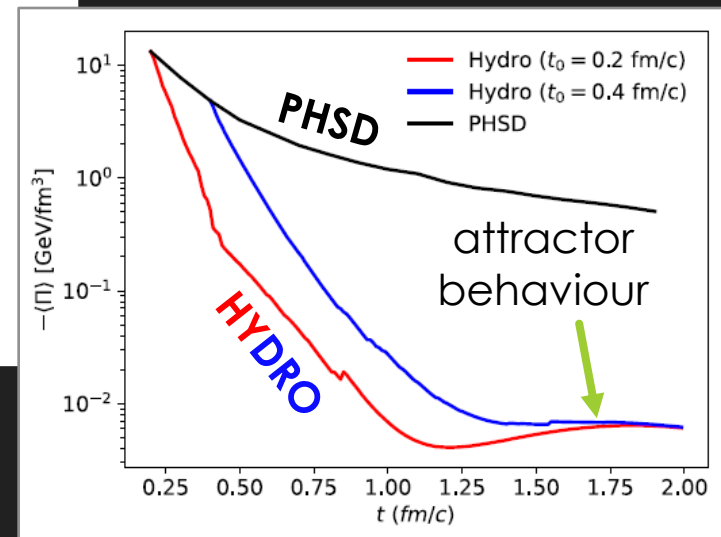
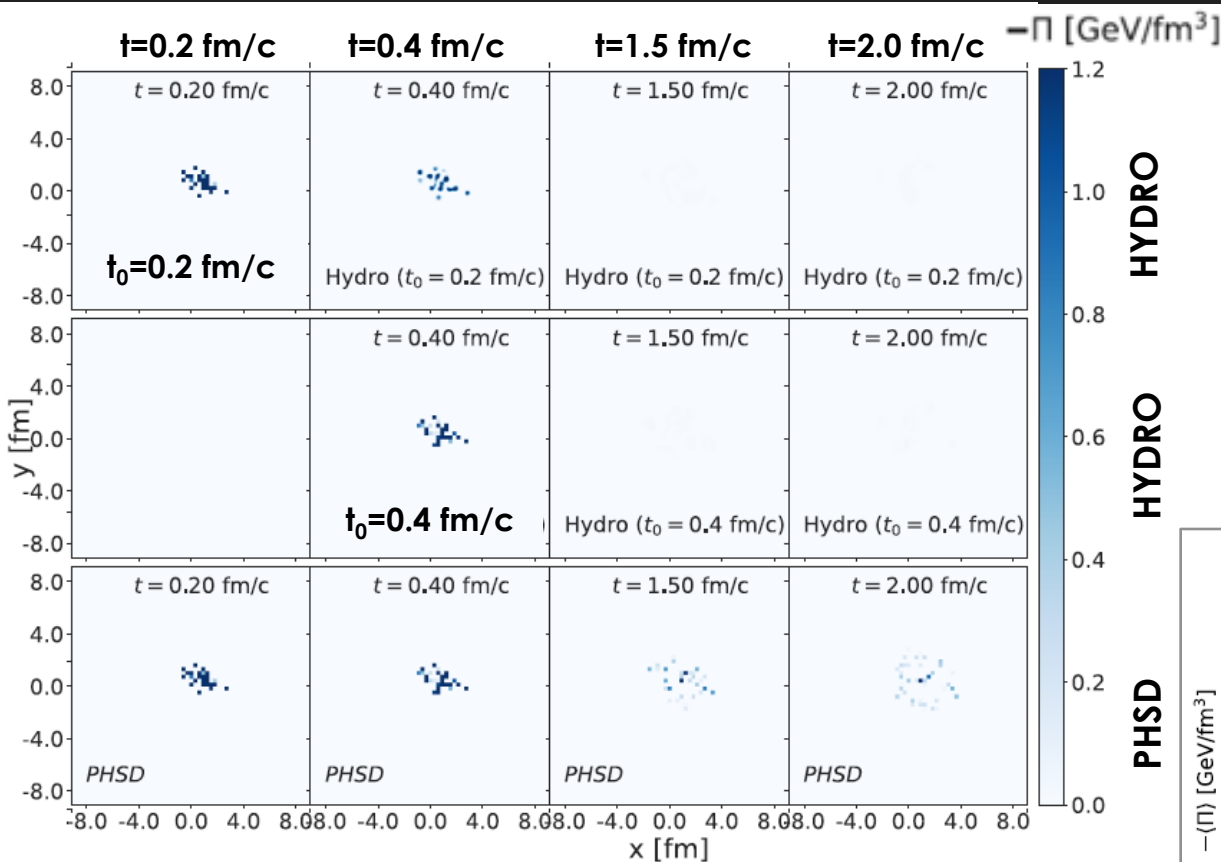
$p+Pb$  @ LHC 5.02 TeV –  $b = 2$  fm

quicker decay than in AA

H. Song and U. W. Heinz,  
Phys. Rev. C 81, 024905 (2010) [\[4\]](#)

- $\Pi$  drops much faster in hydro than in PHSD
- Hydro lines lose memory of initial conditions at  $\sim 1.8$  fm/c

**averaged  $\Pi$**



**bulk viscous pressure**

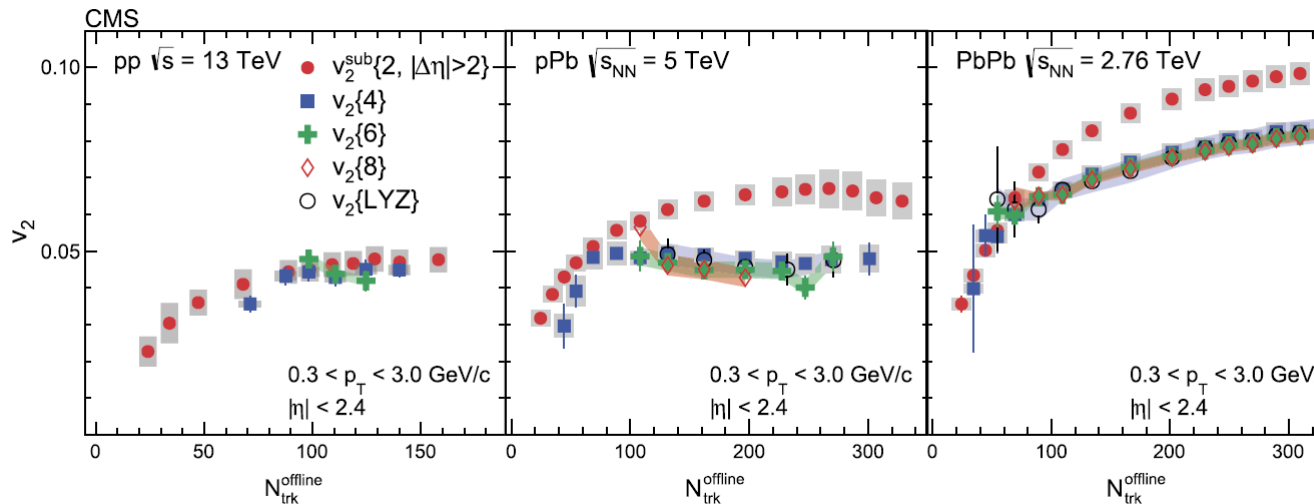
$$\Pi = -\frac{1}{3} \Delta_{\mu\nu} T^{\mu\nu} - P$$

LO, W. Fan, P. Moreau, S.A. Bass and E. Bratkovskaya, Phys. Rev. C 106, 044910 (2022) [\[4\]](#)

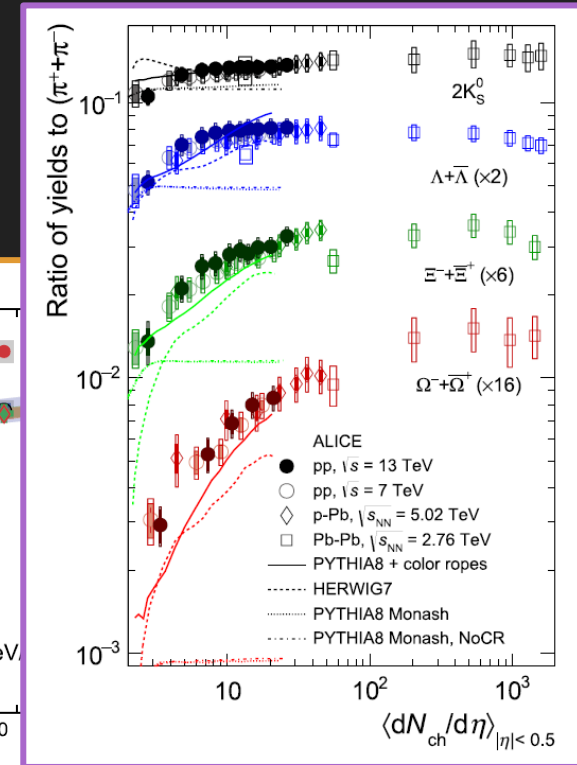
# QGP signals in small systems

Experimental evidence of collective-like behaviour in high-multiplicity  $pp$  and  $pA$  collisions

elliptic flow



CMS Collaboration, Phys. Lett. B 765, 193 (2017) [\[4\]](#)



ALICE Collaboration, Eur. Phys. J. C 80, 693 (2020) [\[4\]](#)

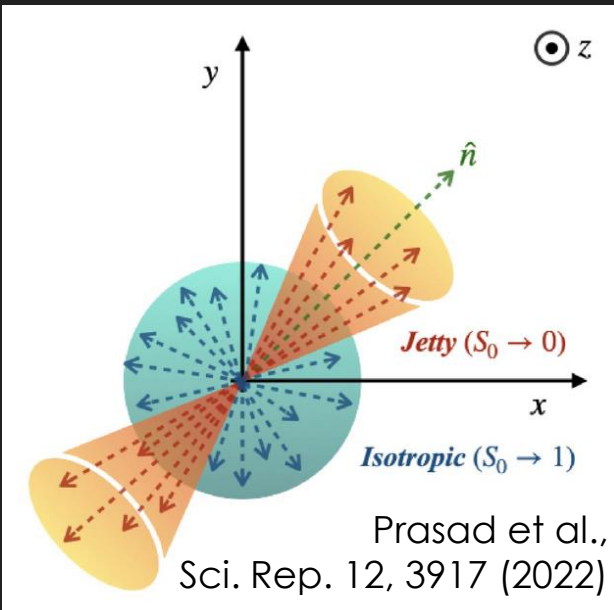
strangeness enhancement

difficulty to well identify QGP signals in small systems

→ attempts to study observables through novel **multi-differential methods**

→ **event-shape engineering**

# Event topology: transverse spherocity



$$S_0 \equiv \frac{\pi^2}{4} \min_{\hat{n}_s} \left( \frac{\sum_i |\mathbf{p}_{Ti} \times \hat{n}_s|}{\sum_i p_{Ti}} \right)^2$$

TRANSVERSE  
SPHEROCITY

event-shape  
observable

A. Banfi, G. Salam and G. Zanderighi,  
JHEP 06, 038 (2010) [\[1\]](#)

❖  $S_0 \rightarrow 0$ : **JETTY events**

all transverse momenta (anti)parallel  
or sum dominated by a single track  
→ dominated by hard physics in  $pp$

❖  $S_0 \rightarrow 1$ : **ISOTROPIC events**

transverse momentua isotropically distributed  
→ dominated by soft physics in  $pp$

A. Khuntia et al., J. Phys. G 48, 035102 (2021) [\[1\]](#)

A. Ortiz et al., Nucl. Phys. A 941, 78 (2015) [\[1\]](#)

ALICE Coll., Eur. Phys. J. C 79, 857 (2019) [\[1\]](#)

A. Nassirpour, J. Phys. Conf. Ser. 1602, 012007 (2020) [\[1\]](#)

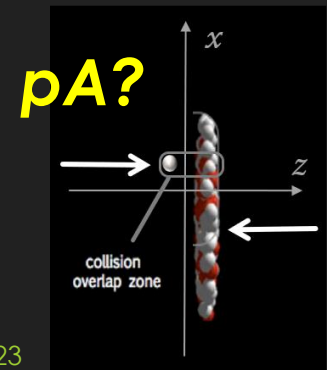
S. Prasad et al., Sci. Rep. 12, 3917 (2022) [\[1\]](#)

N. Mallick et al., J. Phys. G 48, 045104 (2021) [\[1\]](#)

S. Prasad et al., Phys. Rev. D 107, 074011 (2023) [\[1\]](#)

$pp$

$AA$



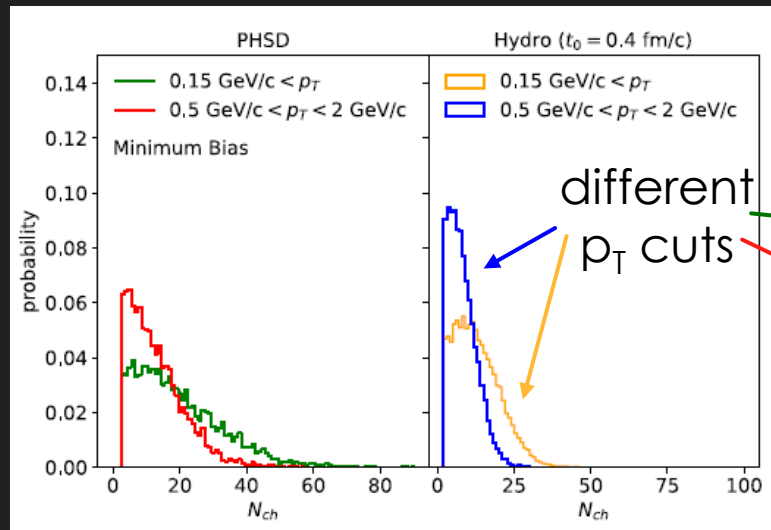
# Multi-differential event categorization

PHSD

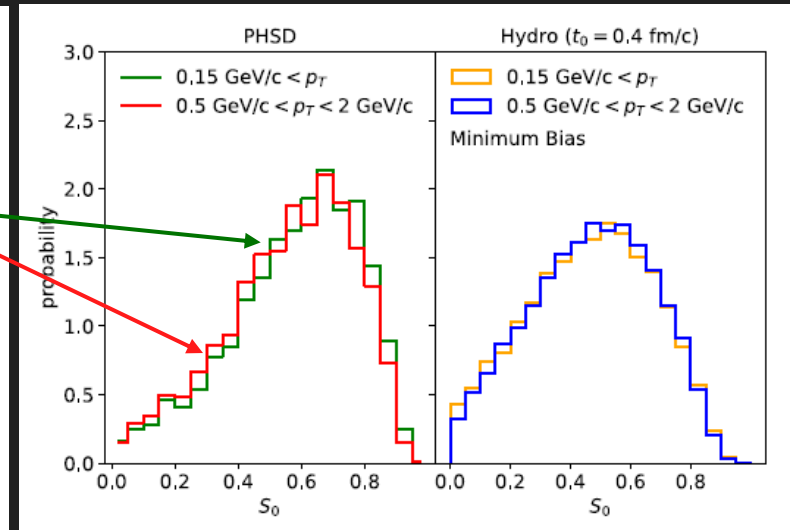
HYDRO

PHSD

HYDRO

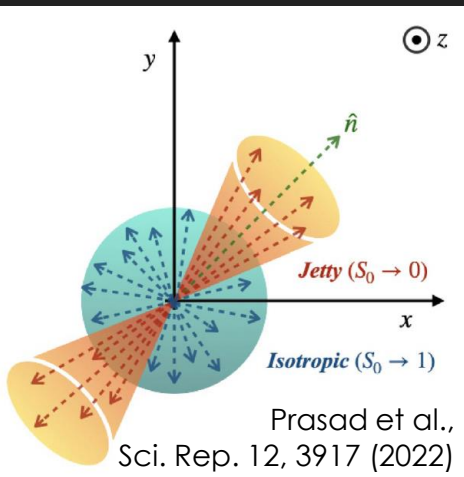


**charged particle distribution**



**transverse sphericity distribution**

p+Pb @ LHC 5.02 TeV



More isotropic event configurations in PHSD compared to hydro

- event topology connected to the different description of the medium in transport and hydro approaches
- multi-differential measurements promising tools to study medium properties in pA

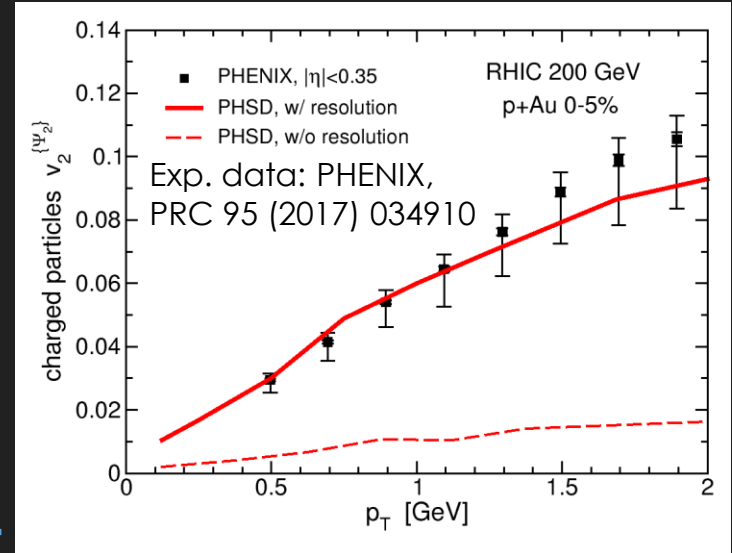
# Elliptic flow in pA

$$v_2(p_T) = \frac{\langle \cos[2(\varphi(p_T) - \Psi_2)] \rangle}{Res(\Psi_2)}$$

In high-multiplicity pA events  $v_2$  comparable to that found in AA

LO, P. Moreau, V. Voronyuk and E. Bratkovskaya,  
Phys. Rev. C 101, 014917 (2020) [\[1\]](#)

p+Au @ RHIC 200 GeV



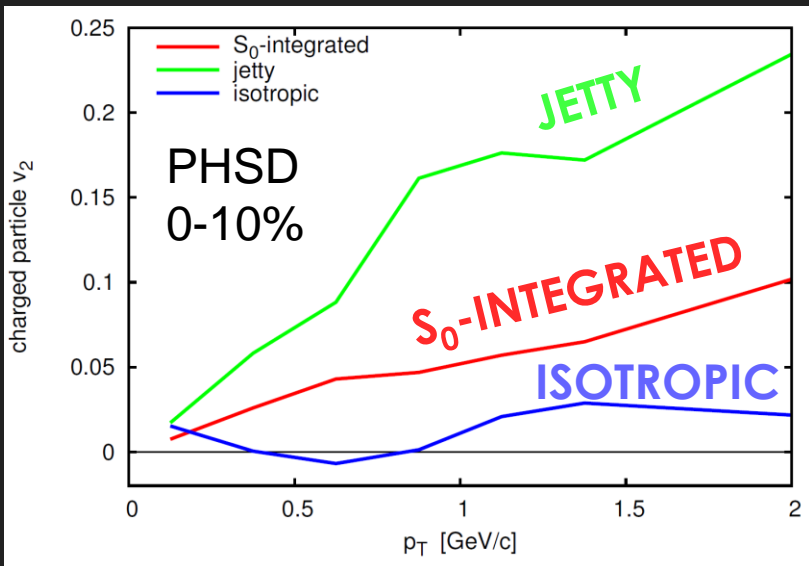
charged particle elliptic flow

# Elliptic flow and trasverse spherocity in pA

$$v_2(p_T) = \frac{\langle \cos[2(\varphi(p_T) - \Psi_2)] \rangle}{Res(\Psi_2)}$$

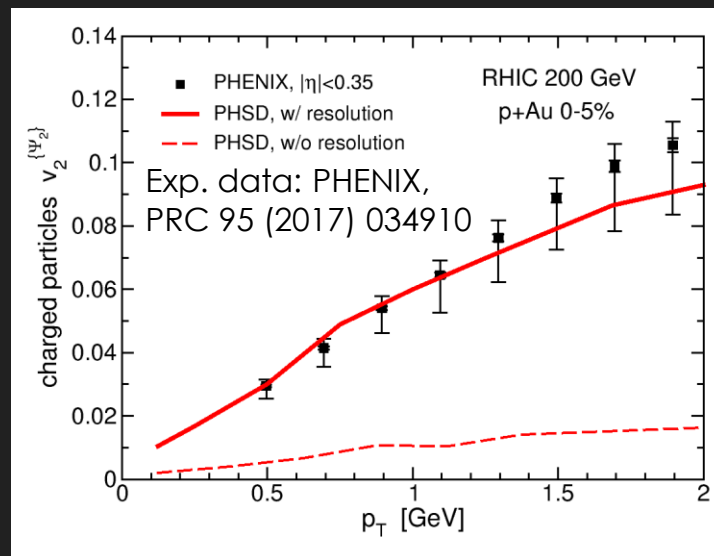
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p+Pb @ LHC 5.02 TeV



LO, P. Moreau, V. Voronyuk and E. Bratkovskaya, Phys. Rev. C 101, 014917 (2020) [\[1\]](#)

p+Au @ RHIC 200 GeV



## charged particle elliptic flow

- isotropic events (high 20%  $S_0$ )  $\rightarrow v_2 \approx 0$
- jetty events (low 20%  $S_0$ )  $\rightarrow$  predominant contribution to the  $v_2$  of spherocity-integrated events


in agreement with AMPT results for Pb+Pb

N. Mallick et al., J. Phys. G 48, 045104 (2021)

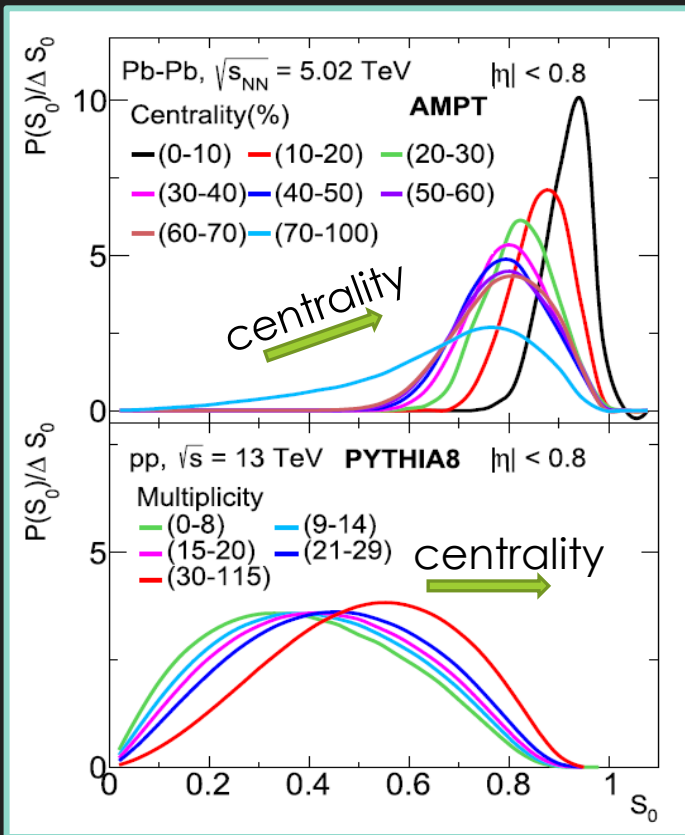
LO, W. Fan, P. Moreau, S.A. Bass and E. Bratkovskaya, Acta Phys. Polon. Supp. 16, 68 (2023) [\[1\]](#)



# Non-trivial relation between event classifiers

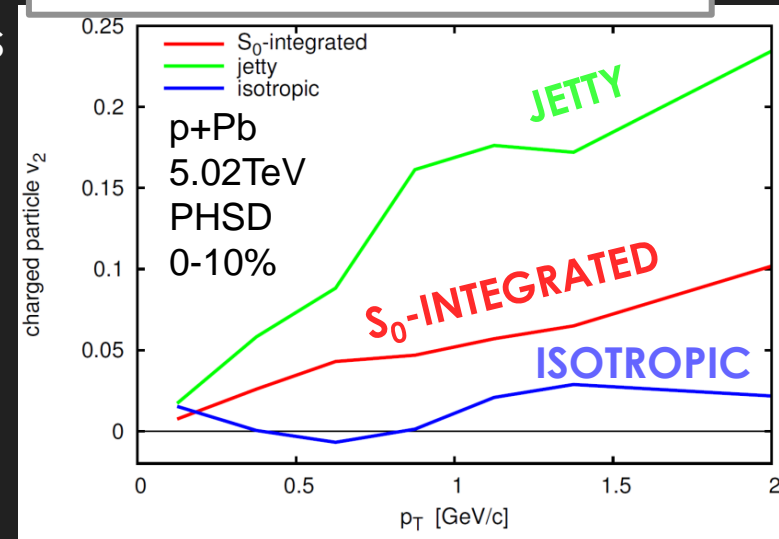
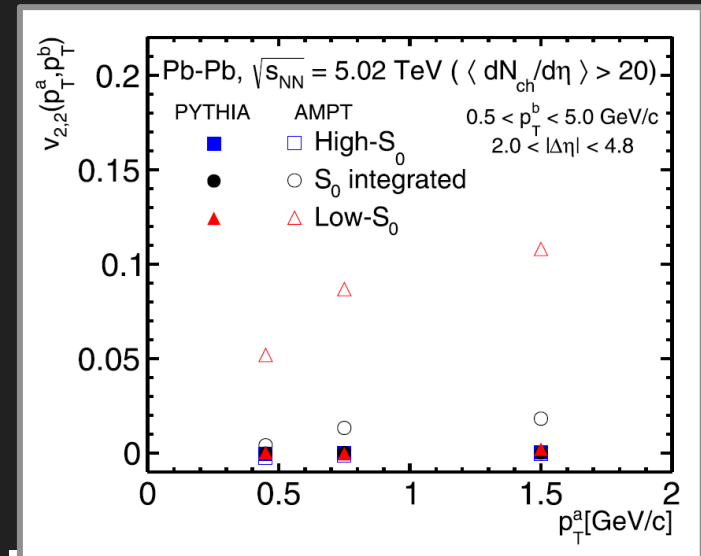
N. Mallick et al., J. Phys. G 48, 045104 (2021) 

In AA, pA, pp  
more central events  $\equiv$  high-multiplicity collisions in pA  
↓  
more isotropic topologies



In AA and pA  
higher  $v_2$  for  
jetty topologies

In pA  
higher  $v_2$  for  
high-multiplicity  
events



S. Prasad et al., Sci. Rep. 12, 3917 (2022) 

LO et al., Acta Phys. Polon. Supp. 16, 68 (2023) 

# Conclusions

Study of **nonequilibrium effects** and **transverse spherocity** in  $p+Pb$  collisions at LHC energy by comparing the PHSD **transport** approach and the VISHNew **hydro** model

- hydro dissolves efficiently initial hot spots in energy density  $e$ ; PHSD preserves medium irregularities
- The bulk viscous pressure  $\Pi$  quickly vanishes in hydro; in PHSD it remains nonzero during the QGP lifetime
- In  $pA$  collisions  $e$  and  $\Pi$  keep in both approaches a high degree of inhomogeneity than in  $AA$
- Transverse spherocity is an event-shape observable that separates jetty and isotropic topologies
- PHSD favors more isotropic event configurations compared to hydro
- Analysis on the elliptic flow  $v_2$  by applying multiplicity + spherocity event selection: in high-multiplicity class, jetty events contribute predominantly to  $v_2$  while isotropic events have  $v_2 \approx 0$

## Thank you for your attention



# Coarse-graining the PHSD medium

Study the hot medium evolution with transport and hydro

- same initial conditions to reduce the early stage impact
- similar equation of state and specific shear viscosity

Y. Xu et al., Phys. Rev. C 96, 024902 (2017) 

T. Song et al., Phys. Rev. C 101, 044903 (2020) 

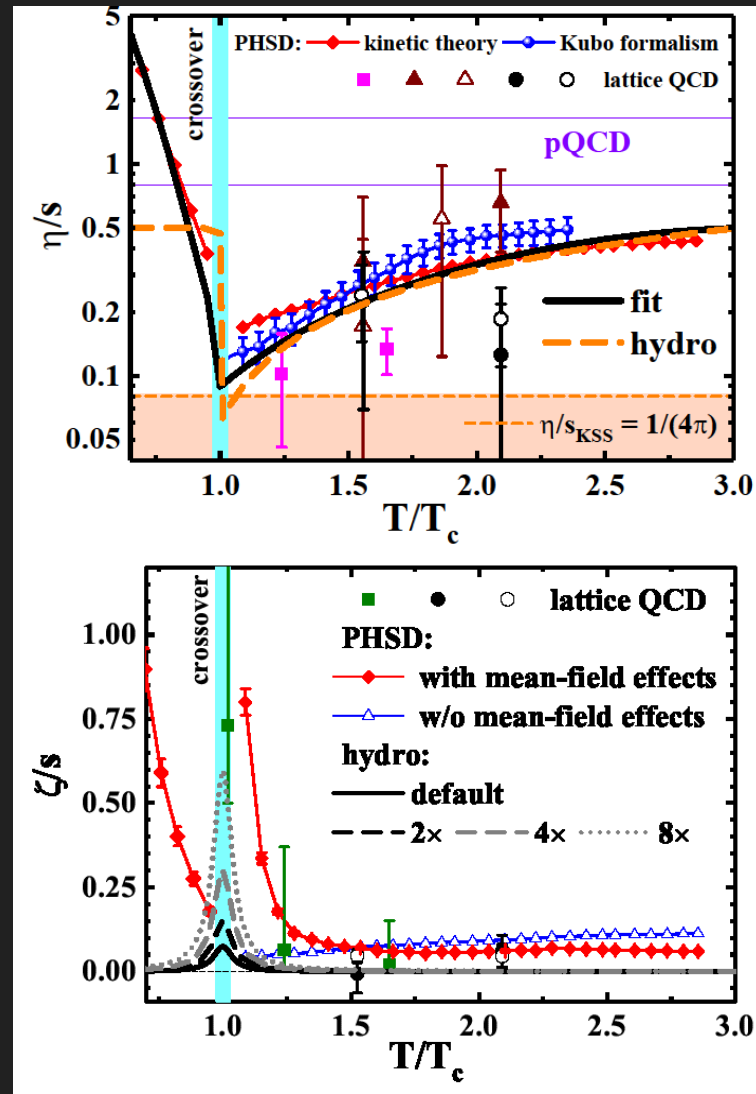
LO et al., Phys. Rev. C 106, 044910 (2022) 


0 0.2 0.4 0.6 0.8

$\tau$  [fm/c]

Initialization time for hydro?

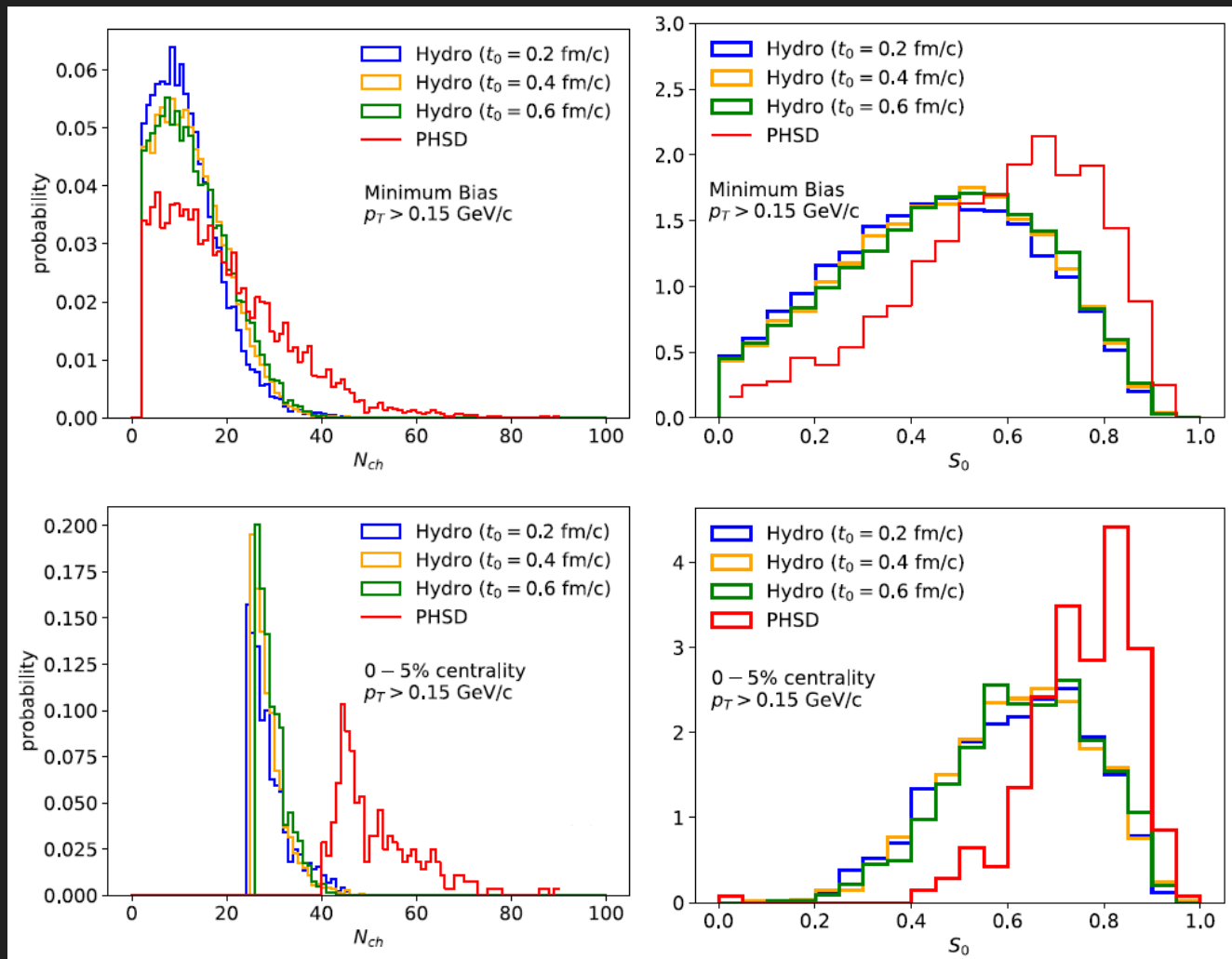
Specific viscosities in the two models  
 $\eta/s(T)$ : the hydro code uses a parametrization obtained from PHSD  
 $\zeta/s(T)$ : much smaller in the hydro code than in PHSD



LO, W. Fan, P. Moreau, S.A. Bass and E. Bratkovskaya, Phys. Rev. C 106, 044910 (2022) 

# Multi-differential event categorization

p+Pb @ LHC 5.02 TeV



charged particles  
 $|\eta| < 0.5$   
 $p_T > 0.15$  GeV/c  
 $N_{\text{trk}} \geq 3$  for  $S_0$

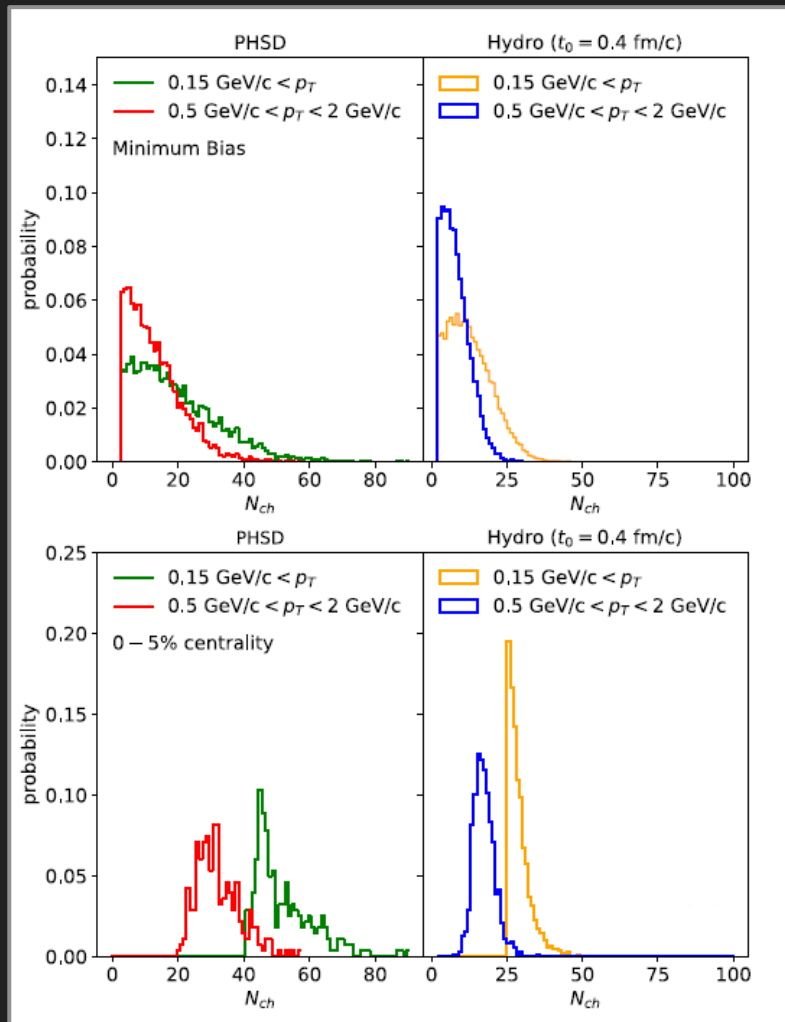
More isotropic configurations in PHSD than in hydro

**charged particle distr.**      **transverse sphericity distr.**

LO, W. Fan, P. Moreau, S.A. Bass and E. Bratkovskaya, Phys. Rev. C 106, 044910 (2022) [arXiv](#)

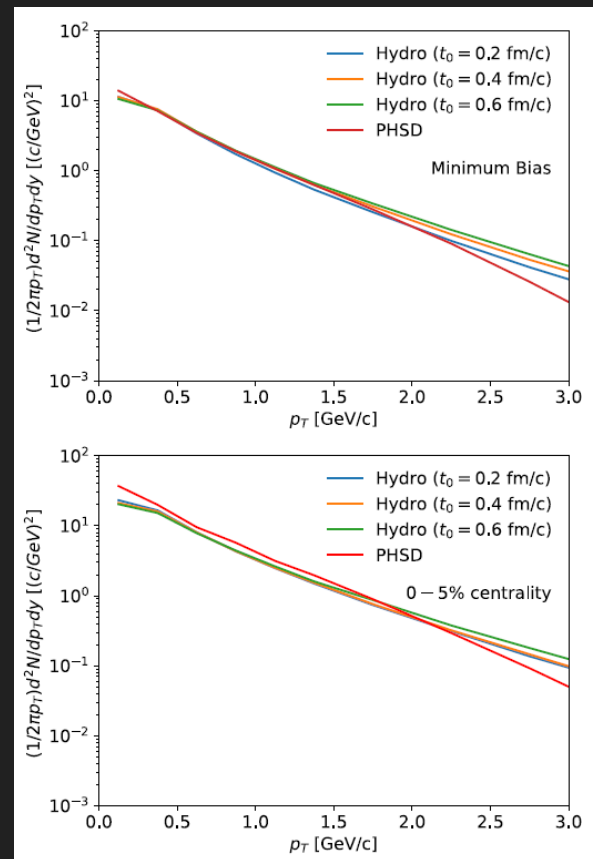
# Multiplicity: PHSD vs hydrodynamics

## charged particle distribution



event distribution  
in multiplicity  
CHANGES applying  
different  $p_T$  cuts

$p+Pb$  @ LHC 5.02 TeV



## charged particle $p_T$ -spectrum

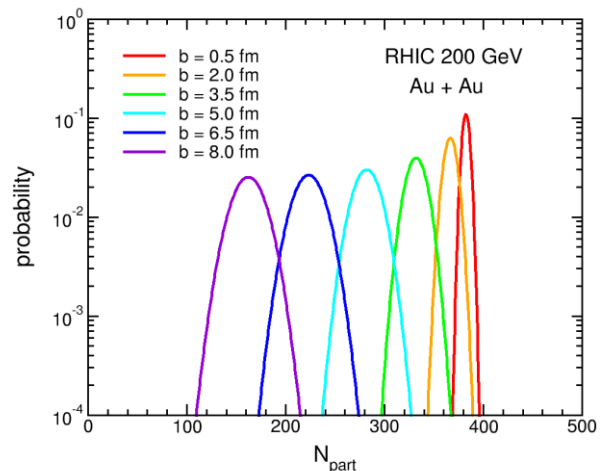
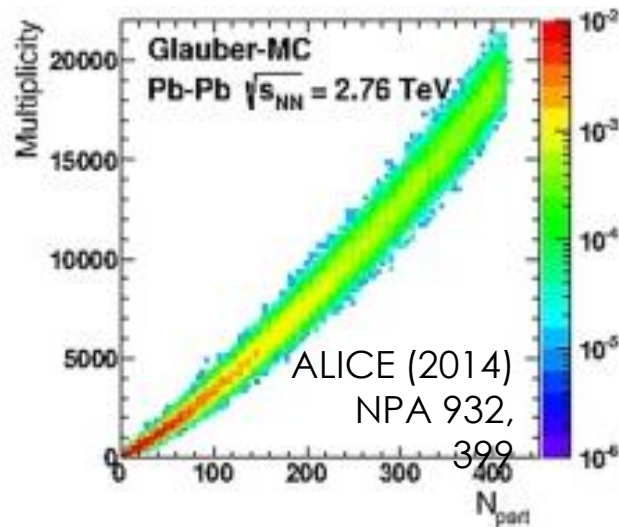
LO, W. Fan, P. Moreau, S.A. Bass and E. Bratkovskaya, Phys. Rev. C 106, 044910 (2022) [arXiv](#)



# Centrality determination : A+A vs p+A

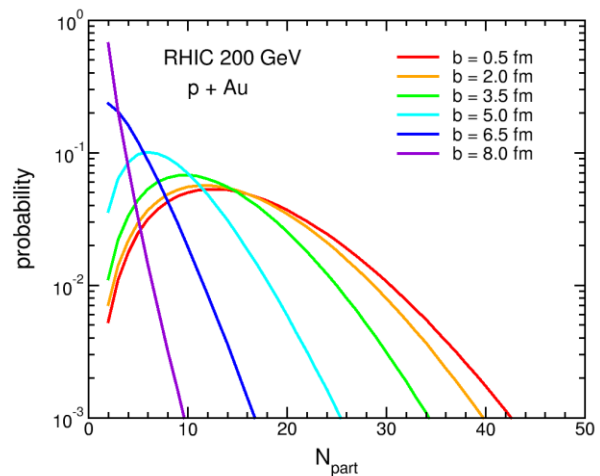
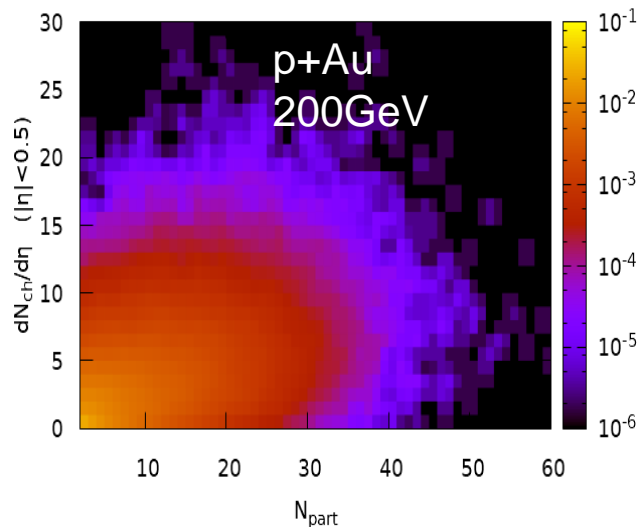
## A+A

centrality  
characterizes  
the amount of  
overlap  
in the interaction  
area



## p+A

multiplicity  
fluctuation mixes  
events  
from different  
impact parameters



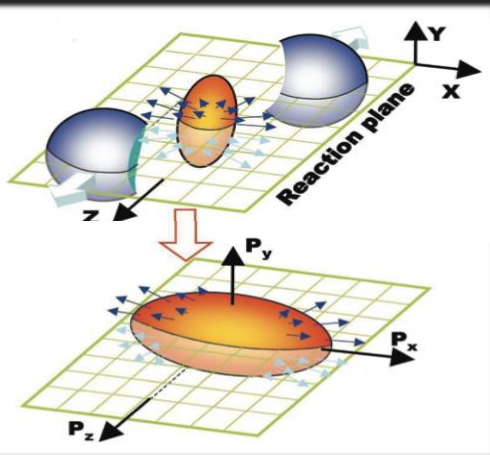
LO, P. Moreau, V. Voronyuk and E. Bratkovskaya, Phys. Rev. C 101, 014917 (2020) 



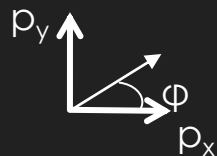
# Anisotropic radial flow $v_n$

azimuthal particle distributions  
w.r.t. the reaction plane

**QGP:** hydrodynamic behavior with very low  $\eta/s$  and formation of collective flows



heavy-ion collisions:  
not a simple **almond shape**  
but a **“lumpy” profile**



$$\frac{dN}{d\phi} \propto 1 + \sum_n 2 v_n(p_T) \cos[n(\phi - \Psi_n)]$$

flow coefficients

event-plane angle

$$v_n = \frac{\langle \cos[n(\phi - \Psi_n)] \rangle}{Res(\Psi_n)}$$

event-plane  
angle resolution

$$\Psi_n = \frac{1}{n} \text{atan2}(Q_n^y, Q_n^x)$$

$$Q_n^x = \sum_i \cos[n\phi_i]$$

$$Q_n^y = \sum_i \sin[n\phi_i]$$

important especially for  
small system, e.g. p+A

finite number of particles produces limited  
resolution in  $\Psi_n$  determination

$\Rightarrow v_n$  must be corrected up to what they  
would be relative to real reaction plane

A. Poskanzer and S. Voloshin, Phys. Rev. C 58, 1671 (1998)

