Spin 1 resonances as probes of spin-vorticity dynamics





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Hydrodynamics with spin





A remarkable experimental discovery, which opened a fascinating field of theoretical investigation.



What does a relativistic system look like where

- Close to local equilibrium (strongly coupled, many DoFs)
- Dofs have spin (not "colliding balls" but "shapes")



Unresolved very non-trivial statistical mechanics problem. Spin not the same as "small vortex".

Spin <u>quantum</u> microstate, how angular momentum is shared determines macroscopic entropy

Vorticity A <u>classical</u> collective excitation carrying angular momentum

So **spin hydrodynamics** means <u>backreaction</u> of microscopic DoFs on macroscopic perturbations!

Phenomenology so far:Becattini et al, 1303.3431

GC ensemble with angular momentum as a conserved quantity, fermions (1 species)

$$\exp\left(-\frac{p_{\alpha}u^{\alpha}}{T}\right) \to \exp\left(-\frac{p_{\alpha}u^{\alpha}}{T}\right)(\bar{u},\bar{v})\exp\left[\frac{\Sigma_{\mu\nu}\omega^{\mu\nu}}{T}\right]\left(\begin{array}{c}u\\v\end{array}\right)$$

And Fermi-Dirac statistics. Here

- $\omega^{\mu\nu}$ vorticity tensor
- $\Sigma^{\mu\nu}$ spin projection tensor $\sim \begin{pmatrix} 0 & \vec{\sigma} \\ \vec{\sigma} & 0 \end{pmatrix}$

local equilibrium isentropic particle production with spin

Fits Λ global polarization

Doesen't fit local polarization (wrong phase).

- use T-vorticity (why ?) only vorticity that makes theoretical sense (Angular momentum, conserved circulation) is $\nabla \times [(e+p)\vec{u}])$
- Symmetric shear and isothermal freeze-out (but ...)

Underestimates some spin-1 particles, ϕ , J/Ψ

But... and this is the main point! Theoretically This can not be the whole story! Perfect equilibrium between vorticity and spin <u>acausal!</u>

GT+Montene gro, 1807.02796, GT, Montene gro, Tinti, 1701.08263



Vorticity+spin must be <u>aligned</u> at <u>equilibrium</u> for entropy to have a well-defined minimum

But this cannot occur instantaneusly in a dynamical system

 $u_{\alpha}\partial^{\alpha}s_{\beta\gamma} + s_{\beta\gamma} = \omega_{\beta\gamma}$

GT, Tinti, Montenegro dispersion relation <u>quartic</u>, always non-causal (Ostrogradski's theorem)

Nora Weickgenannt et al transport theory with spin has non-local collision term

E.Speranza et al non-linear stability analysis

Qualitative picture :GT,D Montenegro, 1807.02796 lower limit to η/s in polarizeable fluids

$$\tau_Y^2 \ge \frac{8c\chi^2(b_o, 0)}{(1 - c_s^2)b_o F'(b_o)} \quad , \quad \frac{\eta}{s} \ge T\tau_Y$$



Many attemps at quantitative picture Stephanov, Hongo, Florkowski, Ryblewski, Torrier Phenomenology needed!

Spin and vorticity: classical-quantum interaction

$$\rho_{spin} = \text{Tr}_{bath} \left[\rho_{spin} \times \underbrace{\rho_{bath}}_{vorticity} \right]$$

- ρ_{spin} Mixed, evolves under Lindblad equation
- ρ_{bath} Maximally mixed, evolves under classical equation of motion
- H_{int} spin-orbit coupling
- **Cooper-Frye** limit of maximal mixing
- Need model-building for this



Global vorticity formed <u>earlier</u>, local vorticity <u>later</u>. Thus the former should be more equilibrated with vorticity than the latter! important consequences for Λ , vector meson. But we need more quantitative observables!

Why spin 1 resonances can help: qbits vs qtrits In general spin 1/2 is a qbit of information

$$\rho = \frac{1}{2} \begin{pmatrix} \rho_{1,1} & \rho_{1,-1} \\ \rho_{-1,1} & \rho_{-1,-1} \end{pmatrix} , \quad \text{Tr}\rho = 1 \quad , \quad U^{-1}\rho U = \text{Diag}\left(\alpha, 1 - \alpha\right)$$

Whatever mechanism combines spin and angular momentum will result in some support of +1,-1 states.

$$\frac{dN_{decay}}{d\theta} = 1 + \alpha \cos\theta$$

Due to parity violation can get polarization vector. Thus in practice any state indistinguishable from equilibrium $\rho_{equilibrium} = \begin{pmatrix} e^{w/T} & 0 \\ 0 & e^{-w/T} \end{pmatrix}$ Can study Dependence of α on y, p_T but need to fit many parameters. Need qualitative signature! Spin 1 contains more information even if parity conserved!

Where

$$\frac{dN}{d\theta d\phi} \propto \cos^2 \theta \rho_{00} + \sin^2 \theta \left(\frac{1-\rho_{00}}{2} + r_{1,-1}\cos(2\phi) + \alpha_{1,-1}\sin(2\phi)\right) +$$

 $+\sin(2\theta)\left(r_{10}\cos\phi + \alpha_{10}\sin\phi\right)$

Variable	Element	coefficient $\times \frac{3}{4\pi}$
$ ho_{00}$	$ ho_{00}$	$\cos^2 \theta$
$\frac{1- ho_{00}}{2}$	$\frac{\rho_{11}+\rho_{-1-1}}{2}$	$\sin^2 heta$
r_{10}^2	$Re[\rho_{-10} - \rho_{10}]$	$\sin(2 heta)\cos(\phi)$
$lpha_{10}$	$Im[-\rho_{-10} + \rho_{10}]$	$\sin(2 heta)\sin(\phi)$
$r_{1,-1}$	$Re[ho_{1,-1}]$	$\sin^2 heta \cos(2\phi)$
$\alpha_{1,-1}$	$Im[ho_{1,-1}]$	$\sin^2 heta\sin(2\phi)$

Spin-1 is a qutrit! Two vectors! Can be mixed $\rho^2 \neq \rho$ in all frames!



$$\rho_8(n_3, n_8) = \frac{1}{3} U^{-1}(\theta, \phi) \begin{pmatrix} 1 + \sqrt{3} n_3 + n_8 & 0 & 0\\ 0 & 1 - \sqrt{3} n_3 + n_8 & 0\\ 0 & 0 & 1 - 2n_8 \end{pmatrix} U(\theta, \phi)$$

generic state neither pure nor thermal! But what is angle ϕ ? Need "objective" ebye definition! beam axis

From angular distributions to purity

$$\frac{dN}{d\theta d\phi} \propto \cos^2 \theta \rho_{00} + \sin^2 \theta \left(\frac{1-\rho_{00}}{2} + r_{1,-1}\cos(2\phi) + \alpha_{1,-1}\sin(2\phi)\right) +$$

 $+\sin(2\theta)\left(r_{10}\cos\phi + \alpha_{10}\sin\phi\right)$

Matrix elements and purity directly related

$$\frac{1}{12} \left(3 \left(n_8 - \sqrt{3} \, n_3 \right) \cos \left(2\theta_r \right) - \sqrt{3} \, n_3 + n_8 + 4 \right) = \rho_{00}$$
$$\frac{\left(n_8 - \sqrt{3} \, n_3 \right) \sin \left(\theta_r \right) \cos \left(\theta_r \right) \cos \left(\phi_r \right)}{\sqrt{2}} = r_{10}$$
$$\frac{\left(\sqrt{3} \, n_3 + 3n_8 \right) \sin \left(\theta_r \right) \sin \left(\phi_r \right)}{3\sqrt{2}} = \alpha_{10} \quad , \quad \phi_r = -\frac{1}{2} \tan^{-1} \left(\frac{\alpha_{1,-1}}{r_{1,-1}} \right)$$

A concrete example: Coalescence K^*,ϕ



Coalescence in a rotating medium

THe expression of the vector meson density matrix in terms of the quark density matrices and the vorticity is straight-forwardly

$$\left(\hat{\rho}^{M}\right)_{mn} = \sum_{ijkl} \left(P_{12}^{L}\right)_{ijklmn} U_{S}(\phi_{r},\theta_{r}) \left(U_{\omega}(\mu_{1},\nu_{1})\rho^{1}(\Omega)U_{\omega}^{-1}(\mu_{1},\nu_{1})\right)_{ij} \times$$

$\times \left(U_{\omega}(\mu_{2},\nu_{2})\rho^{2}(\Omega)U_{\omega}^{-1}(\mu_{2},\nu_{2}) \right)_{kl} U_{S}^{-1}(\phi_{r},\theta_{r})$

Where $P_L(w)$ is the (unknown) probability to aquire a spin quantum number from vorticity and the rest are 6-j and C-G coefficients!

Big approximation : non-relativistic. But with constituent quarks and moderate p_T, y not bad

Golden signature:off-diagonal elements



If ebye ϕ coefficients $\simeq 0$, Cooper-Frye, otherwise coalescence And then

we discovered it was all measured... for quarkonium!

$$\rho_{00} = \frac{1 + \lambda_{\theta}}{3 + \lambda_{\theta}} \quad , \quad r_{1,-1} = \frac{\lambda_{\phi}}{3 + \lambda_{\theta}} \quad , \quad r_{10} = \frac{\lambda_{\theta\phi}}{3 + \lambda_{\theta}}$$



Solving for $n_{1,8}$ gives close to maximally mixed state



Off-diagonal elements compatible with zero, Cooper-Frye! but...



Remember spin-orbit non-equilibrium is all about interplay between eaaly (global) and <u>late</u> (local) axes. Need modulation in angle from reaction plane of λ s!

Qualitatively one expects of diagonal elements to be "local", and ny harmonic function averages to zero. Is it zero in every azimuthal bin?

Getting more quantitative for Quarkonium Rotating Cornell potential with spin orbit interaction



Rotating Cornell potential with spin orbit interaction

$$\begin{cases} \mathbf{P} = \mathbf{p_1} + \mathbf{p_2}, & \mathbf{p} = \mu \left(\frac{\mathbf{p_1}}{m_1} - \frac{\mathbf{p_2}}{m_2} \right) \\ \mu = \frac{m_1 m_2}{m_1 + m_2}, & \mathbf{r_1} = \mathbf{R} + \frac{m_2}{m_1 + m_2} \mathbf{r}, & \mathbf{r_2} = \mathbf{R} - \frac{m_1}{m_1 + m_2} \mathbf{r} \end{cases}$$

$$\mathcal{H} = \frac{P^2}{2M} + \frac{p^2}{2\mu} - \mathbf{P} \cdot (\omega \times \mathbf{R}) - \mathbf{p} \cdot (\omega \times \mathbf{r}) - \omega \cdot (\mathbf{S}_1 + \mathbf{S}_2) + V(\mathbf{r})$$

With a Cornell-type potential. This way vorticity and spin interactions accounted for. Mass correction $\Delta E_{i,j}$ and Melting temperature from Debye formula

Mass becomes sensitive to relative direction of spin and vorticity



This is experimentally undetectable (widening), but...

One can get melting temperature from Debye formula!



Melting temperature depends non-monotonically on rotation <u>and</u> (anti)polarization. Could such "distillation" explain ϕ spin alignment? ϕ is quarkonium and only polarized survive! need abundance vs alignment scans with rapidity.

Can also have to do with quaronium being more suppressed at higher rapidity (Which is also higher vorticity)



"Distillation" is an important effect qualitatively different from Cooper-Frye Need: SUppression vs alignment in rapidity!

Correlation between invariant mass and off-equilibrium $\rho!$



Conclusions

Non-equilibrium between spin and vorticity theoretically well-established, need phenomenology

Spin and vorticity different objects

Cooper-Frye could be misleading

- Spin1 vector mesons can serve as such a link because of a rich structure in their density matrix. Coalescence in vector mesons, potential models in quarkonia
- **To do** baryon coalescence (can coalescence explain local polarization?), Ω (Spin 3/2), etc.

SPARE SLIDES