

PAUL SCHERRER INSTITUT



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Search for a muon EDM

Muon4Future, Venice, 05/30/23

Project funded by



Schweizerische Eidgenossenschaft
Confédération suisse
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Swiss Confederation

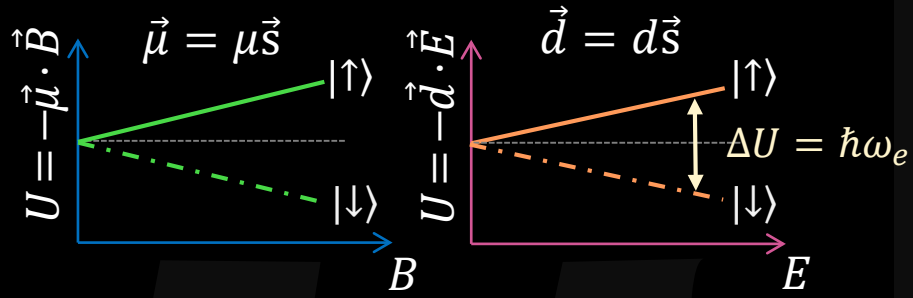
Federal Department of Economic Affairs,
Education and Research EAER
State Secretariat for Education,
Research and Innovation SERI

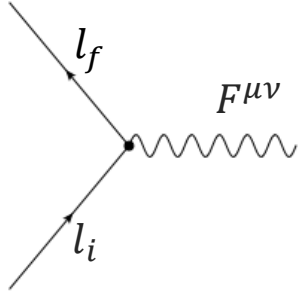


SWISS NATIONAL SCIENCE FOUNDATION



CP violation & edm





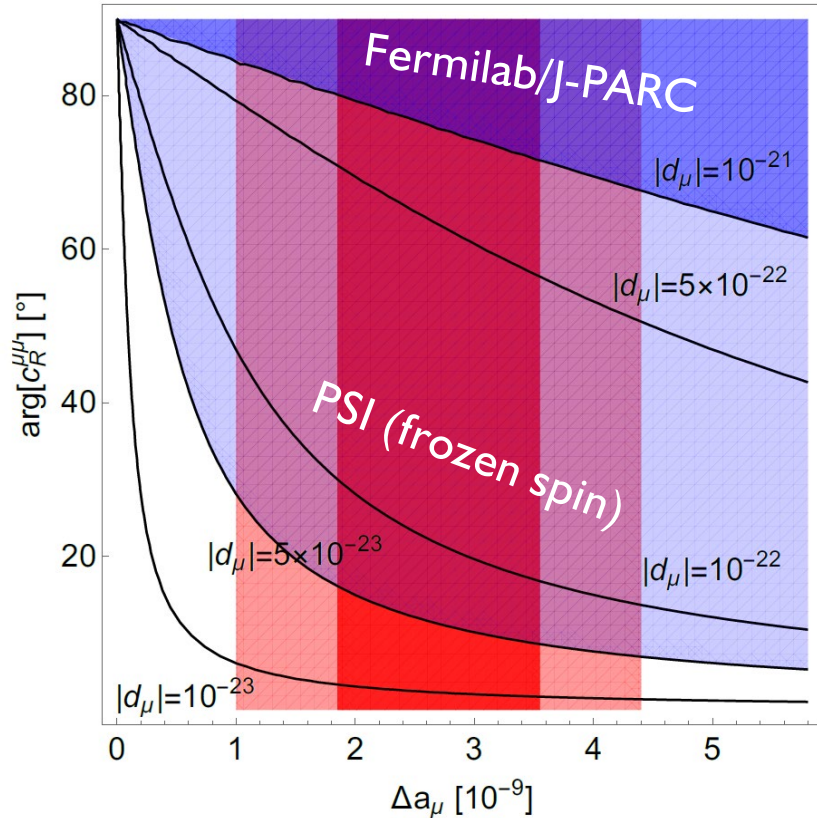
Effective Hamiltonian: $\mathcal{H}_{\text{eff}} = c_R^{l_f l_i} \bar{l}_f \sigma_{\mu\nu} P_R l_i F^{\mu\nu} + \text{h. c.}$

$$\langle p' | J_{\mu}^{\text{EM}} | p \rangle = \bar{\Psi}(p') \left[F_1 \gamma_{\mu} + \underbrace{\frac{iF_2}{2M} \sigma_{\mu\nu} q^{\nu}}_{\text{magnetic-dipole}} + \underbrace{\frac{iF_3}{2M} \sigma_{\mu\nu} \gamma_5 q^{\nu}}_{\text{electric-dipole}} + \frac{F_4}{M^2} (q^2 \gamma_{\mu} - \gamma^{\mu} q_{\mu} q_{\nu}) \right] \Psi(p)$$

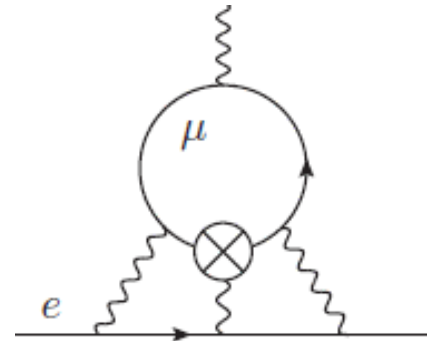
charge
Anapole - moment

$$\delta F_2 = a_{l_i} = -\frac{2m_{l_i}}{e} \left(c_R^{l_i l_i} + c_R^{l_i l_i^*} \right) = -\frac{4m_{l_i}}{e} \text{Re } c_R^{l_i l_i}$$

$$F_3 = d_{l_i} = i \left(c_R^{l_i l_i} - c_R^{l_i l_i^*} \right) = -2 \text{Im } c_R^{l_i l_i}$$



- MFV: $|d_{\mu \leftarrow e}^{\text{MFV}}| < 8.5 \times 10^{-28} \text{ ecm}$
- Contribution only starts at the 3-loop level*
 $|d_{\mu \leftarrow e}| < 4 \times 10^{-20} \text{ ecm}$
- Y. Ema et al., PRL **128**, 131801 (2022)
 $|d_\mu(^{199}\text{Hg})| < 6 \times 10^{-20} \text{ ecm}$
 $|d_\mu(\text{ThO})| < 2 \times 10^{-20} \text{ ecm}$
- Bennett et al., PRD **80**, 052008 (2009)
 $|d_\mu| < 1.5 \times 10^{-19} \text{ ecm}$



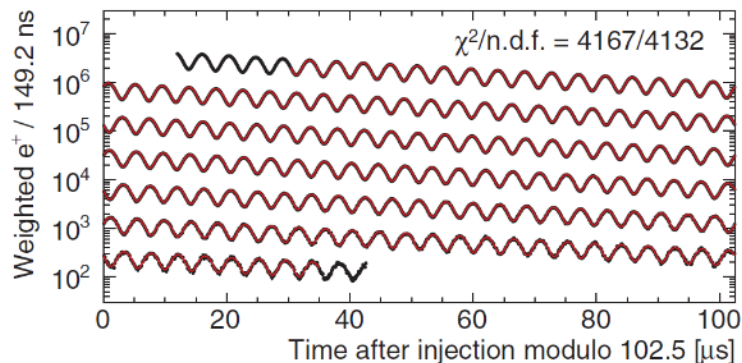
*A.Crivellin, M. Hoferichter, PSW PRD 98, 113002 (2018)

Muon dipole moments and frequencies

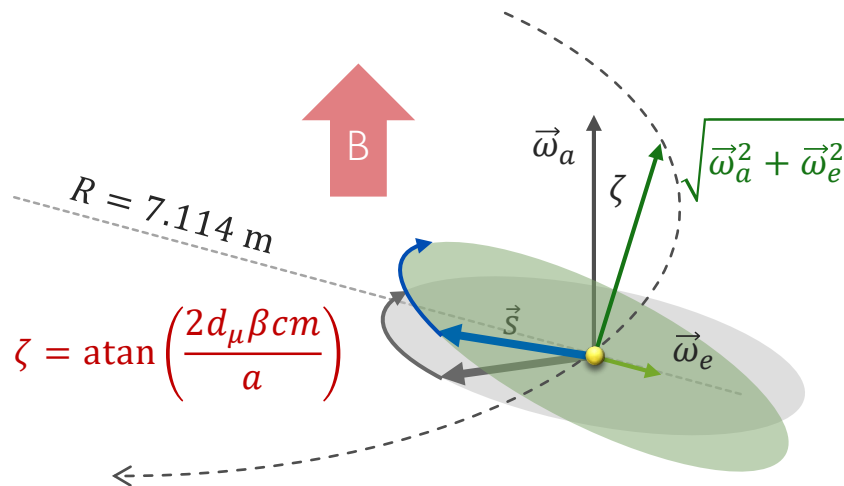


$$\vec{\omega} = \vec{\omega}_L - \vec{\omega}_c = -\frac{q}{m} \left[a\vec{B} + \left(\frac{1}{\gamma^2 - 1} - a \right) \frac{\vec{\beta} \times \vec{E}}{c} \right]$$

g-2 term ω_a



FNAL* & JPARC**: $\sigma(d_\mu) \approx 10^{-21} ecm$



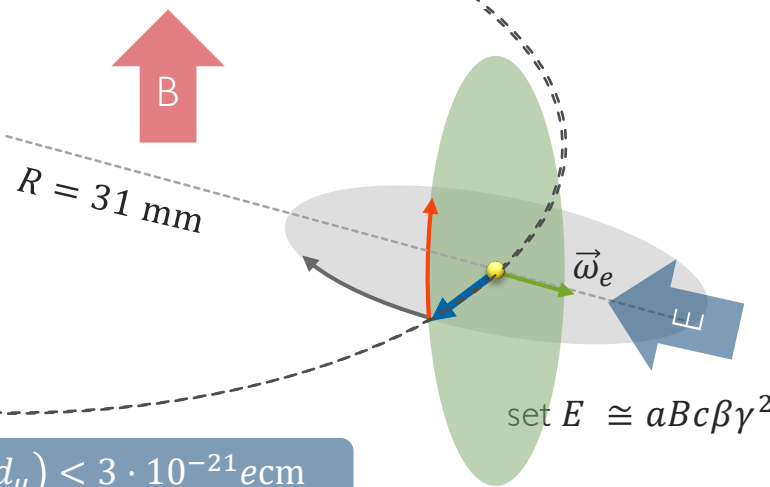
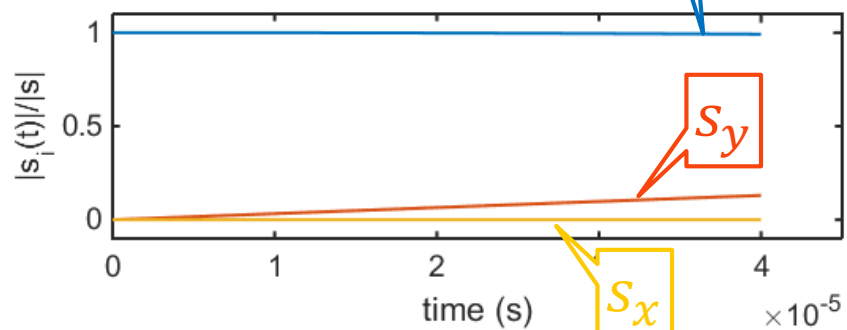
$$\zeta = \text{atan} \left(\frac{2d_\mu \beta c m}{a} \right)$$

Muon dipole moments –freezing the spin at PSI



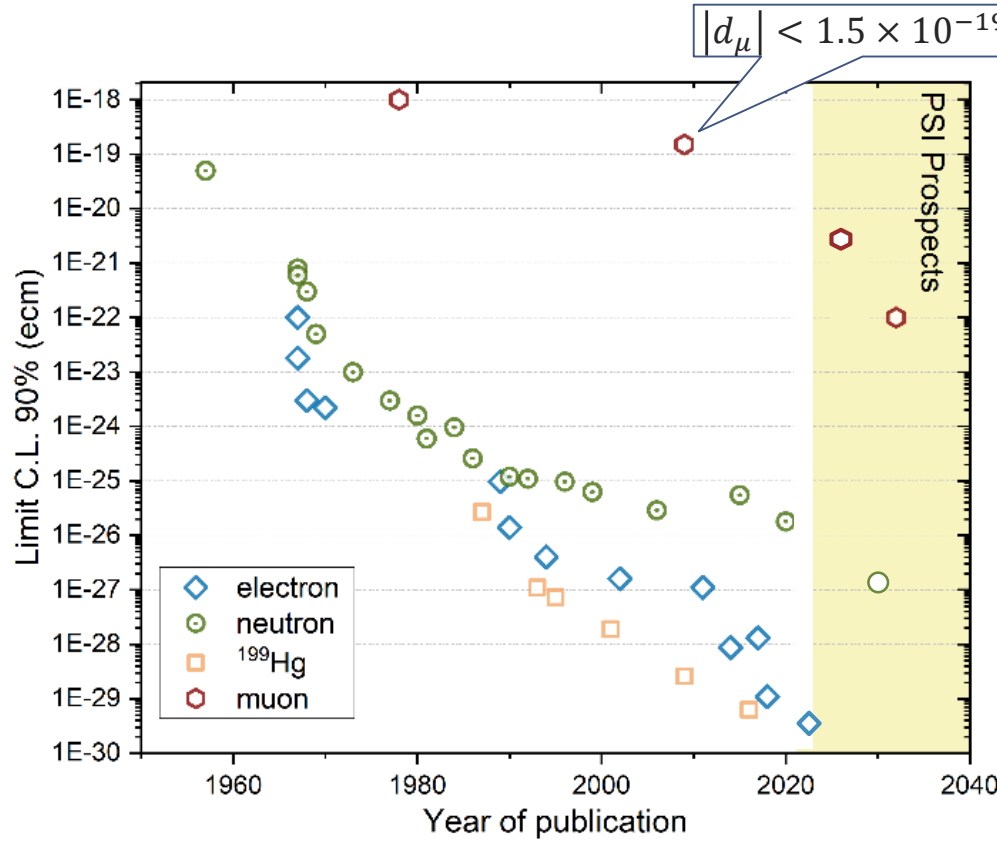
$$\vec{\omega} = -\frac{q}{m} \left[\underbrace{a\vec{B} + \left(\frac{1}{1-\gamma^2} - a \right) \frac{\vec{\beta} \times \vec{E}}{c}}_{\text{g-2 term } \omega_a} + \underbrace{\frac{2d_\mu mc}{q\hbar} \left(\frac{\vec{E}}{c} + \vec{\beta} \times \vec{B} \right)}_{\text{EDM term } \omega_e} \right]$$

Frozen-spin potential at PSI: $\sigma(d_\mu) < 6 \cdot 10^{-23} \text{ ecm}$

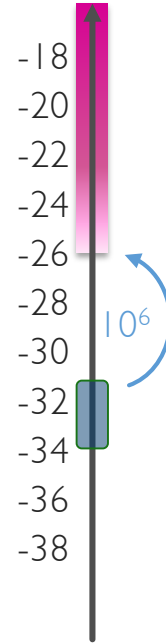


[*Farley et al., PRL93 042001 (2004)] ,

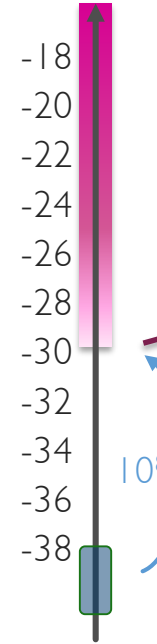
A not so brief history of EDM searches



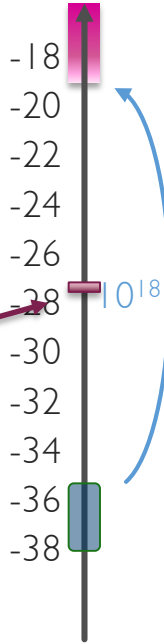
Neutron
log(d) /ecm



Electron
log(d) /ecm



Muon
log(d) /ecm

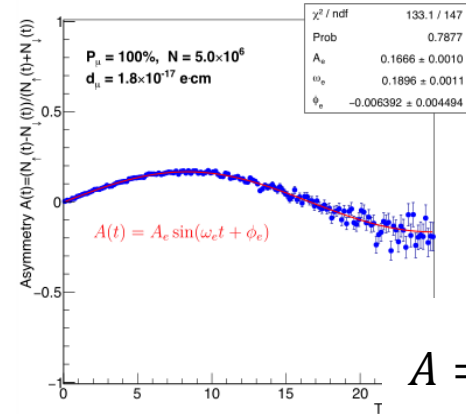
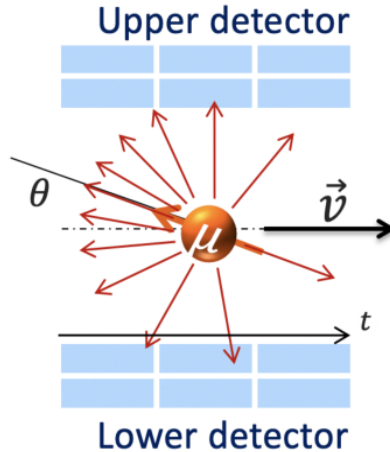
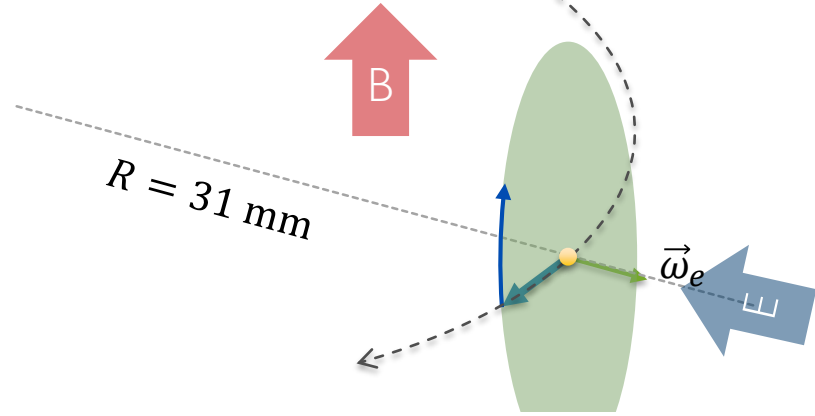


*Bennett et al., PRD80(2009)052008
** Abel et al., PRL124(2020)081803

The general experimental idea

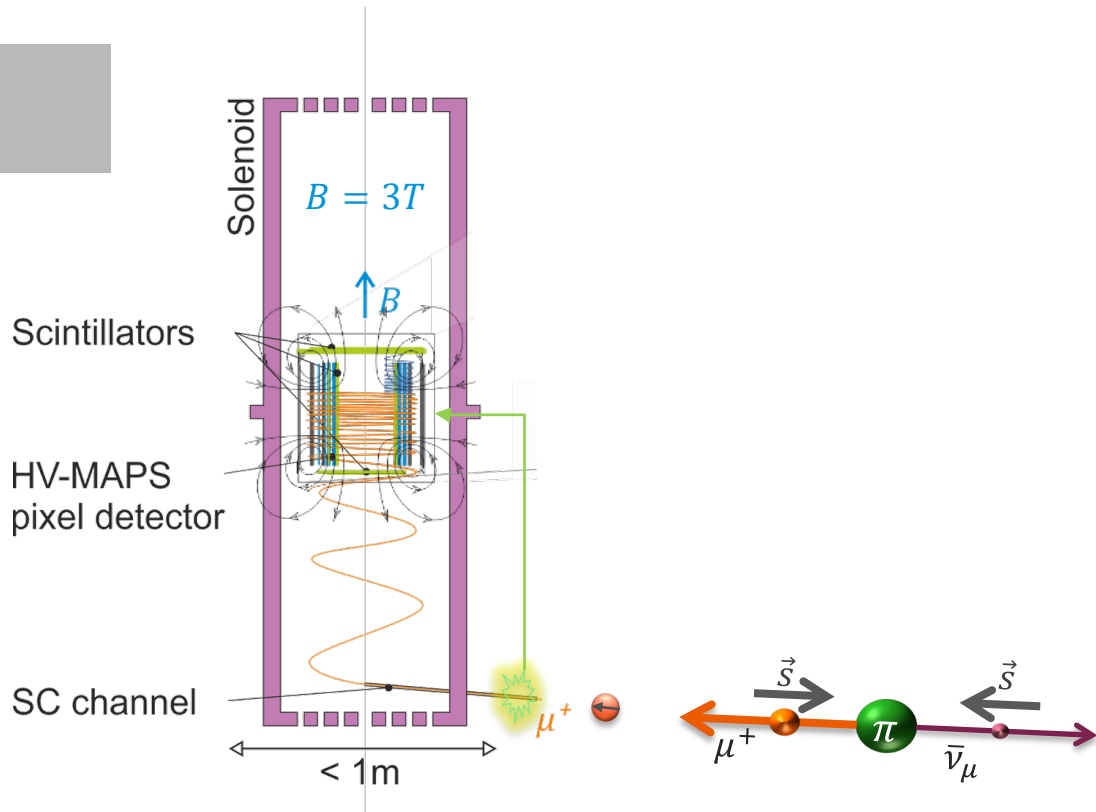


- If the EDM $\neq 0$, then there will be a vertical precession out of the plane of the orbit
- An asymmetry increasing with time will be observed recording decay positrons
- If the EDM = 0, then the spin should always be parallel to the momentum asymmetry should be zero
- Some asymmetry could still be observed due to systematic effects



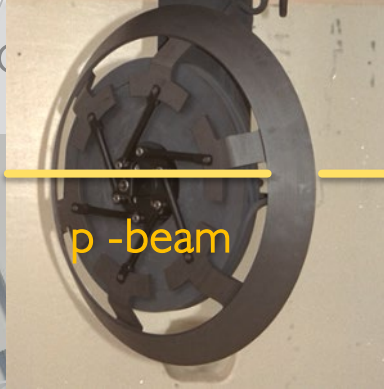
$$A = \frac{N_u - N_d}{N_u + N_d}$$

Search for a muon EDM using the frozen-spin technique with longitudinal injection

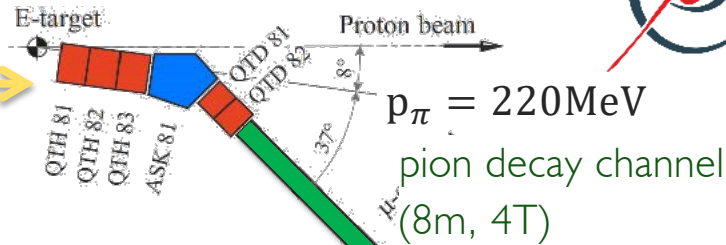


- μ^+ from **Pion-decay** → high polarization $p \approx 95\%$
- Injection through **superconducting channel**
- **Fast** scintillator **triggers** pulse
- Magnetic **pulse stops** longitudinal motion of μ^+
- Weakly focusing field for **storage**
- **Thin electrodes** provide electric field for **frozen spin**
- Pixelated detectors for **e^+ - tracking**

μ^+ @ 125MeV/c or 28MeV/c



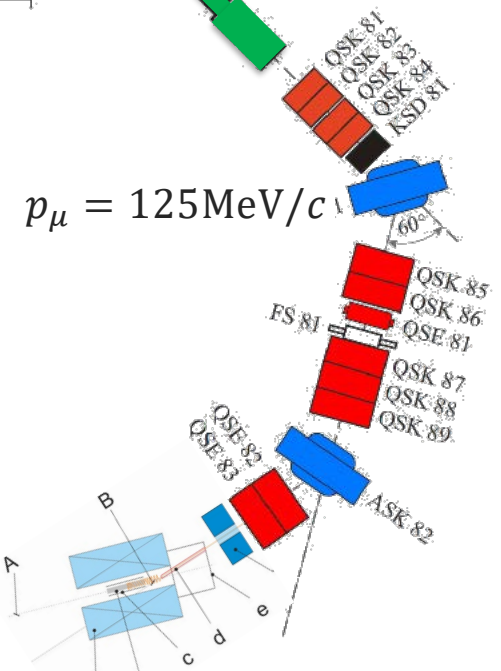
p-beam



1 m

μ -beam:

- $10^8 \mu^+ / \text{s}$
- Large phase space



A phased approach



Phase I (small solenoid, surface muons)



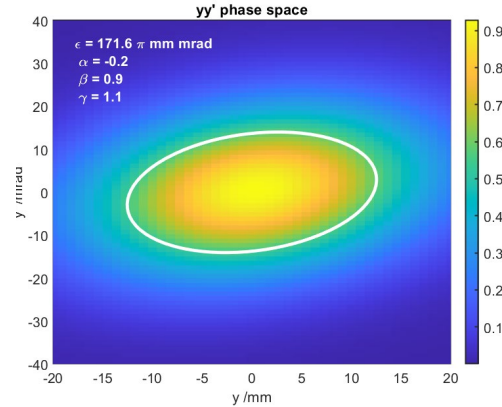
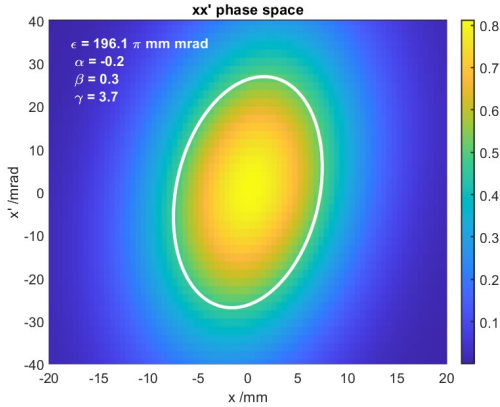
- Existing solenoid at PSI, max 5T
- Bore diameters 200mm
- Field was measured in 2022 (found suitable for injection)

Phase 2 (dedicated magnet muon momentum $\geq 125\text{MeV}/c$)



Argonne 4T solenoid

- Large bore (up to 900 mm diameter)
- High Temporal field stability (10ppb/h)
- Excellent spatial field uniformity (< 1 ppb/mm)



xx' : $\epsilon = 196\pi \text{ mm mrad}$

yy' : $\epsilon = 171 \pi \text{ mm mrad}$

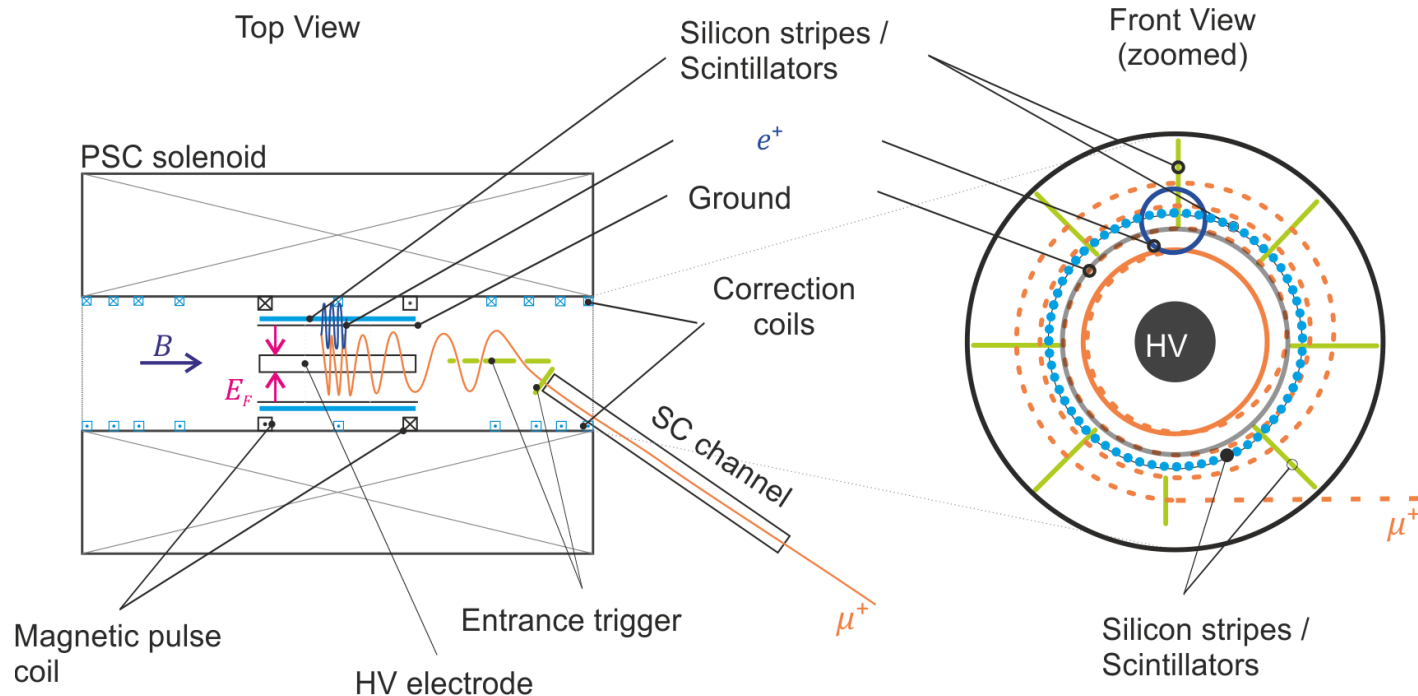
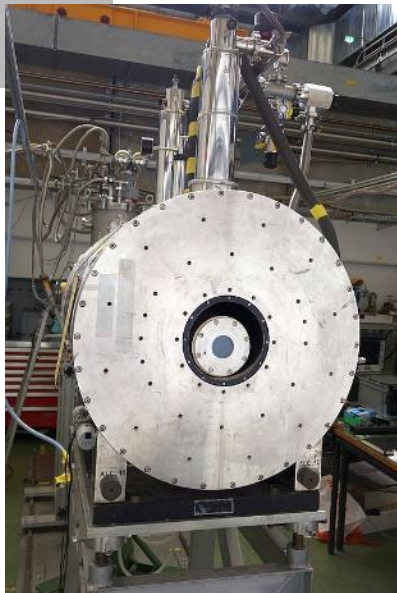
- Large phase space at exit of beam collimated by passage through a collimation channel
- Due to adiabatic magnetic collimation large part of transmitted μ^+ are reflected.
- Storage efficiency $\sim 5 \times 10^{-4}$

$$\sigma(d_\mu) = \frac{\hbar \gamma a_\mu}{2pE_f \sqrt{N} \tau_\mu \alpha}$$

	$\pi E1$	$\mu E1$
Muon flux (μ^+/s)	4×10^6	1.2×10^8
Channel transmission	0.03	0.005
Injection efficiency	0.017	0.60
Muon storage rate (1/s)	2×10^3	360×10^3
Gamma factor γ	1.04	1.56
e^+ detection rate (1/s)	500	90×10^3
Detections per 200 days	8.64×10^9	1.5×10^{12}
Mean decay asymmetry A	0.3	0.3
Initial polarization P_0	0.95	0.95
Sensitivity in one year (e-cm)	$< 3 \times 10^{-21}$	$< 6 \times 10^{-23}$

The muEDM phase I on piE1

Test bed and frozen spin demonstrator



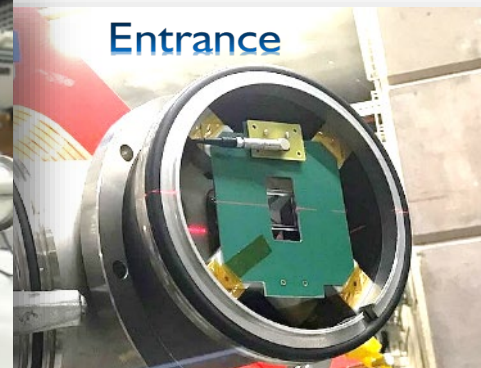
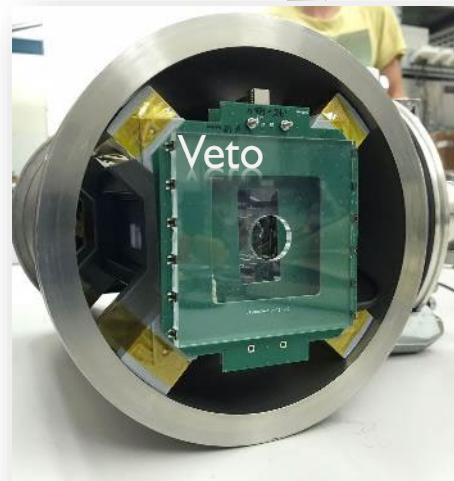
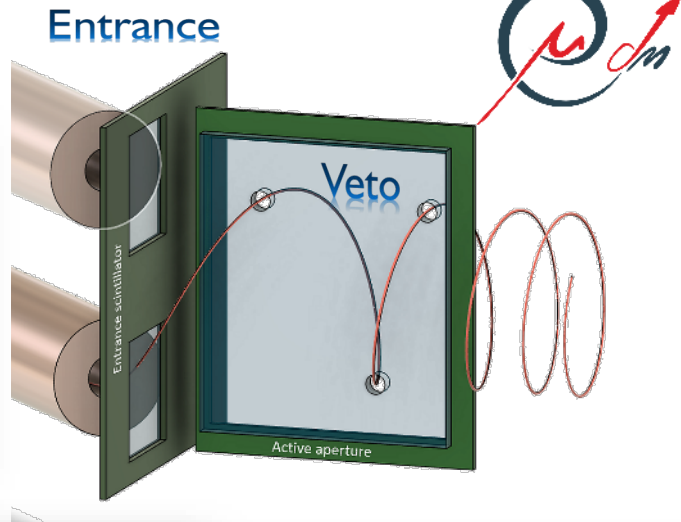
$$\text{muEDM measurement} < 3 \cdot 10^{-21} \text{ ecm}$$

Muon entrance trigger

- Magnetic pulse needs to be triggered by incident muon
- Only about 1% of muons passing through the collimation channel are within the acceptance phase space
- Scattering in scintillators increase beam divergence

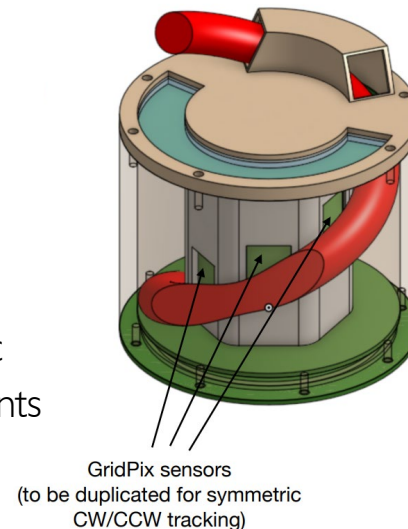
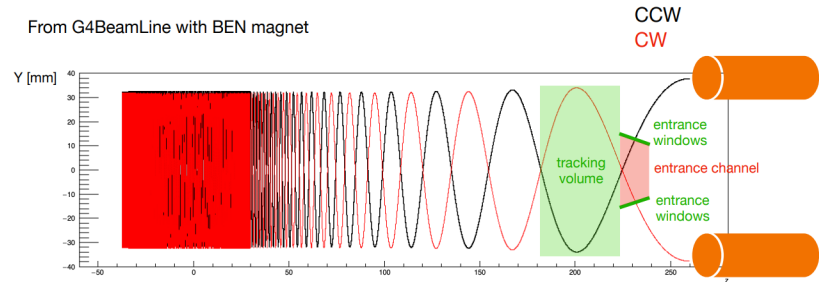


- Combine thin ($\leq 100\mu\text{m}$) entrance scintillator with an active aperture as veto

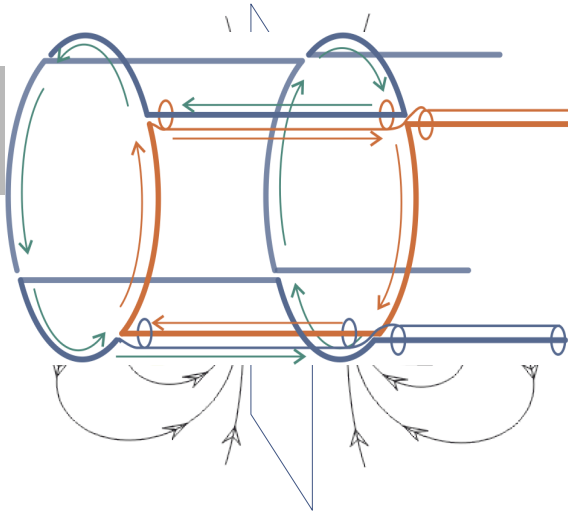




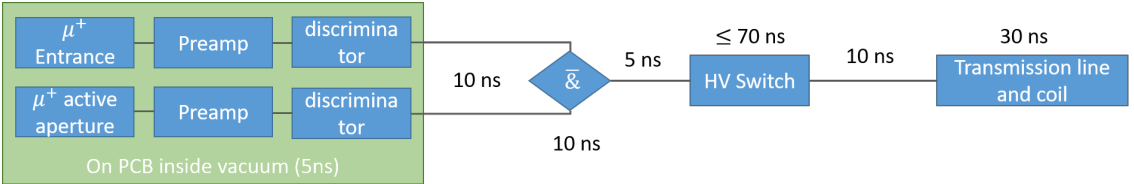
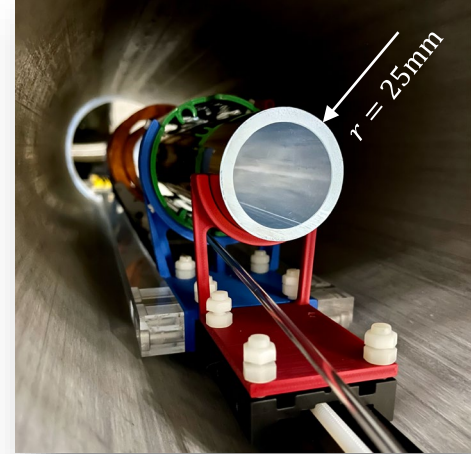
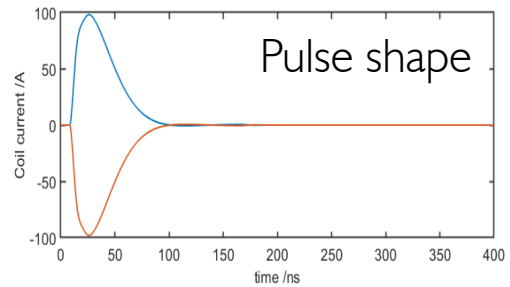
- Need to characterize muon trajectory before EDM measurement
- Requires \sim mrad angular and $\sim 0.1\%$ momentum resolution
- Gaseous TPC with 2 geometries possible: longitudinal drift (for momentum), radial drift (for angle).
- Design satisfying constraints with sufficient phase-space acceptance possible with current trajectory parameter
- Resolution of the phase space reconstruction, with realistic ionization and drift properties from beam test measurements and investigation of low pressure gas options



Radial magnetic field pulse to kick muons



Pulse generator



- PS delay input \rightarrow output $\delta t < 60\text{ ns}$
- Pulse FWHM $\sim 40\text{ ns}$
- Peak current per coil $\sim 170\text{ A}$
- Includes damping effect (factor 2) by Eddy currents
- Suppression of oscillation in tail $< 1\text{ A}$ (corresponds to $B_{osc} < 5\mu\text{T}$)



Detection of g-2 precession ω_a

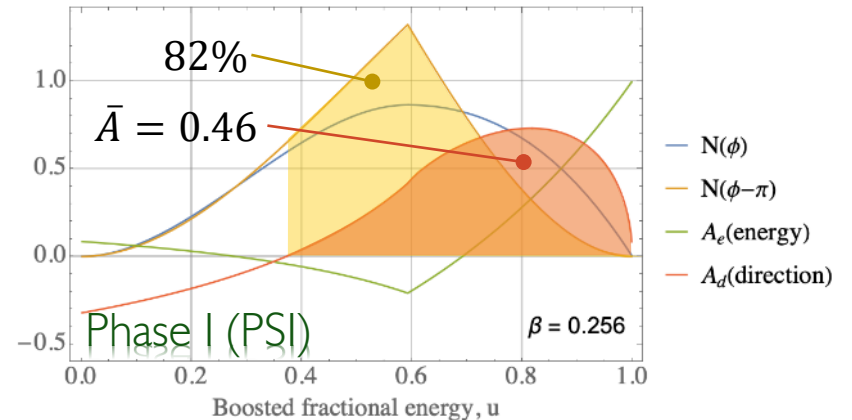
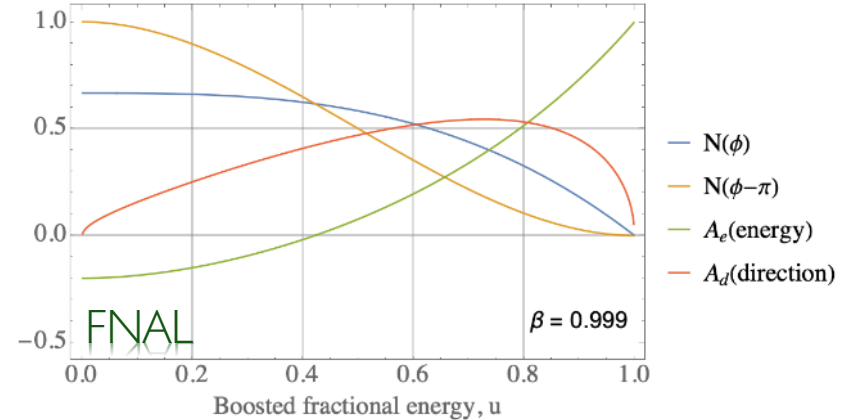
- Measurement of mean magnetic field $\langle B \rangle$
- Measure $\omega_a(E)$ to tune electric field to frozen-spin condition

Requires momentum resolution

Detection of EDM polarization

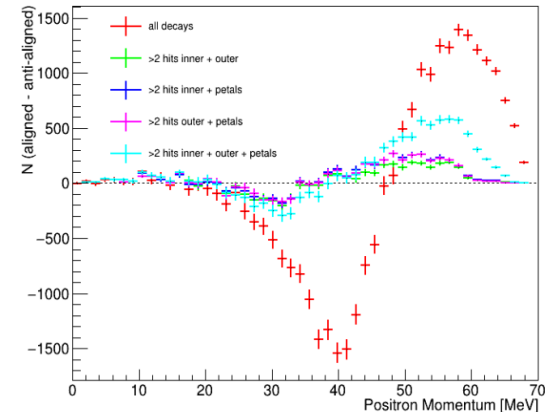
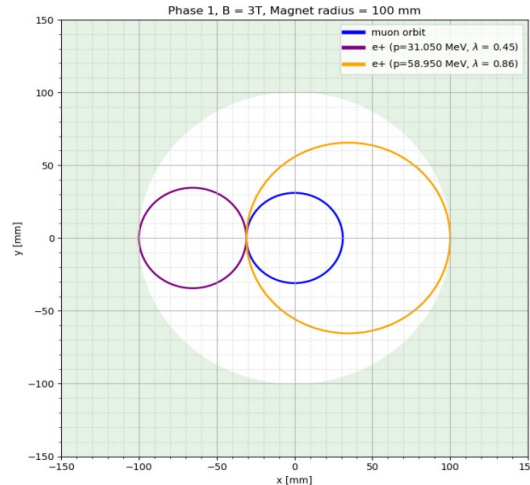
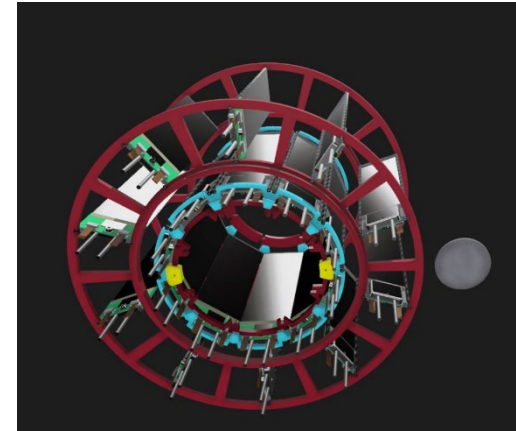
- Measurement of Asymmetry as function of time $A(t)$

Requires spatial resolution along cylinder





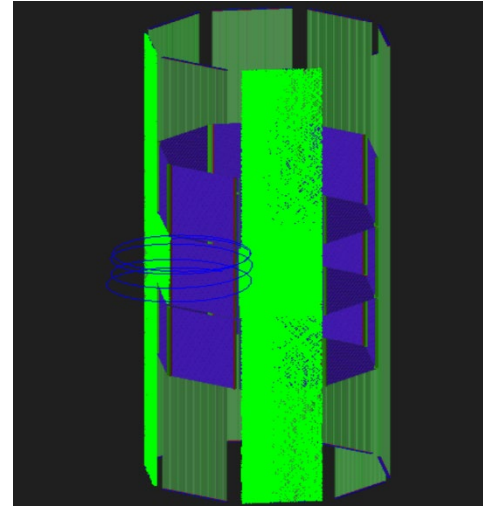
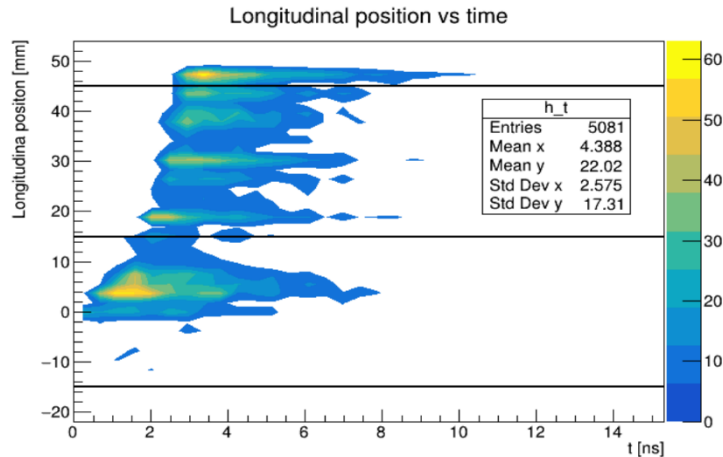
- Measurement of g-2 to measure B-field and determine voltage for frozen spin condition
- For Phase I ($p = 28\text{MeV}$) maximum energy is $\sim 68.65\text{ MeV}$
- Positrons with momentum greater than 59 MeV will hit the magnet regardless of decay direction.
- For ($31 < p < 59$) MeV positrons it depends on decay direction
- Cylindrical (silicon strip) tracking detectors at $r=35\text{ mm}, 47.5\text{ mm}$



(courtesy J. Price)



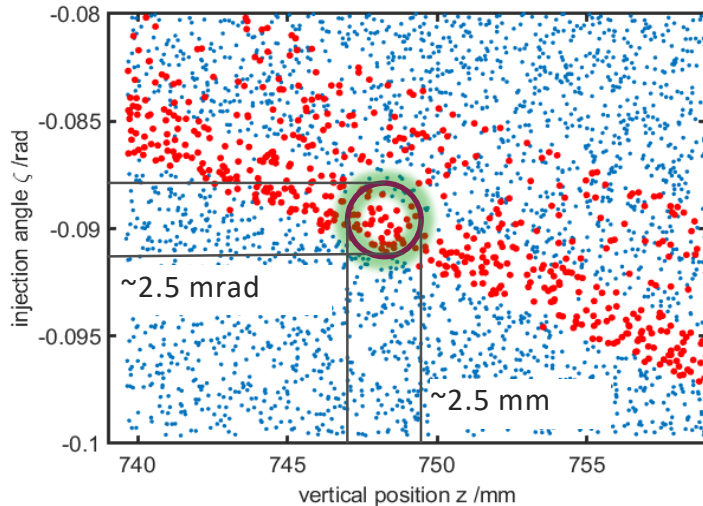
- For EDM signal, detect up-down asymmetry in photons
- Double barrel Scifi tracker, radius of the inner detector currently equal to 50 mm
- Bundles of fibers with good resolution
 - transverse and longitudinal fibers
 - transverse fibers with longitudinal straw/pix



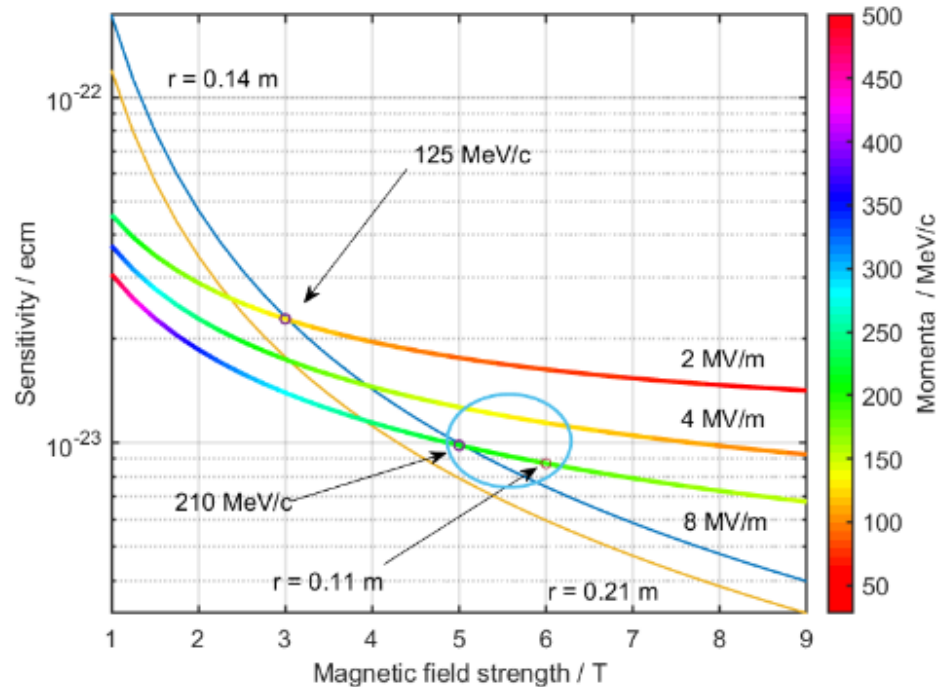
- Photon time and position (longitudinal info on internal barrel)
- Large number of readout channel a challenge
- Considering other possible geometries

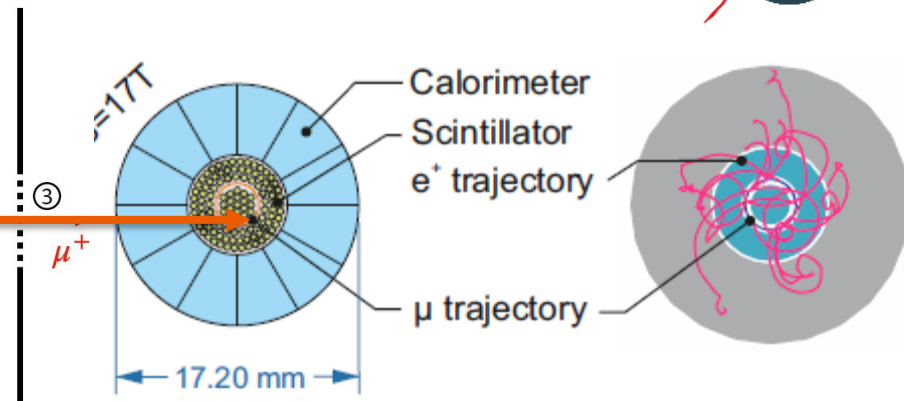
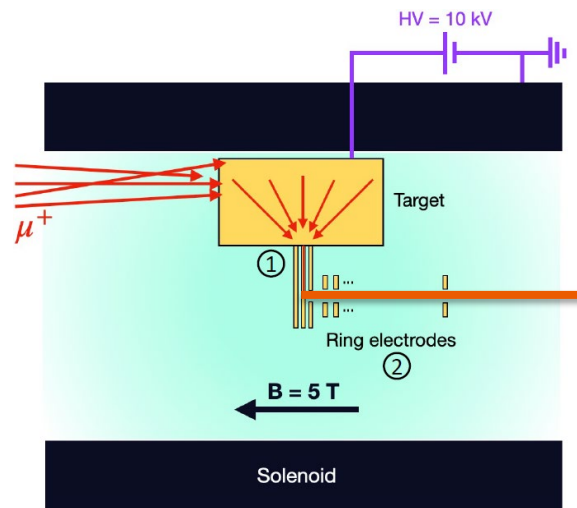


- Low emittance & high flux to inject all delivered muons
- “Muon on request” or pulsed beam
- The more muons the merrier

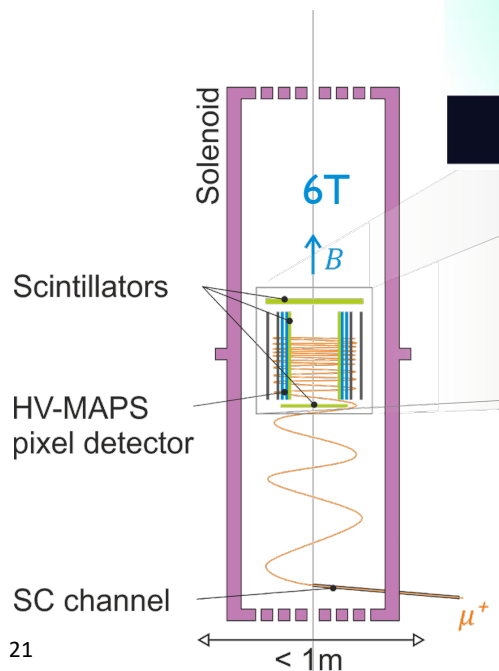


$$E \approx aBc\beta\gamma^2$$





c.f. D. Taqqu



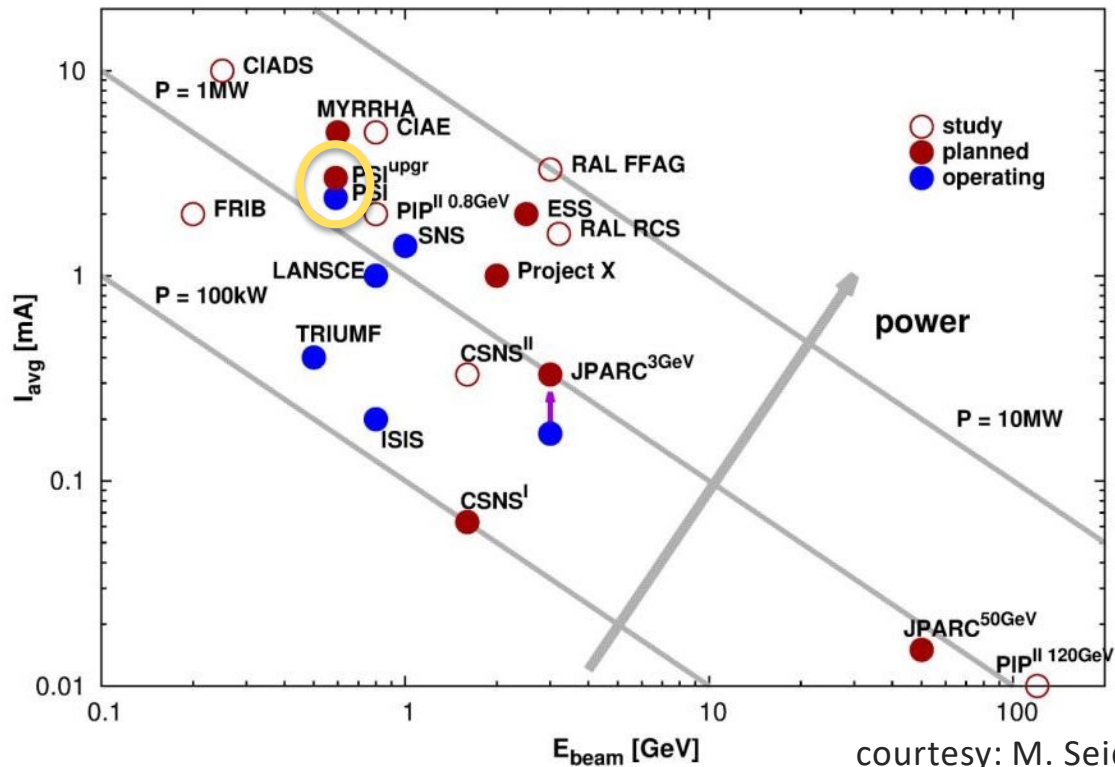
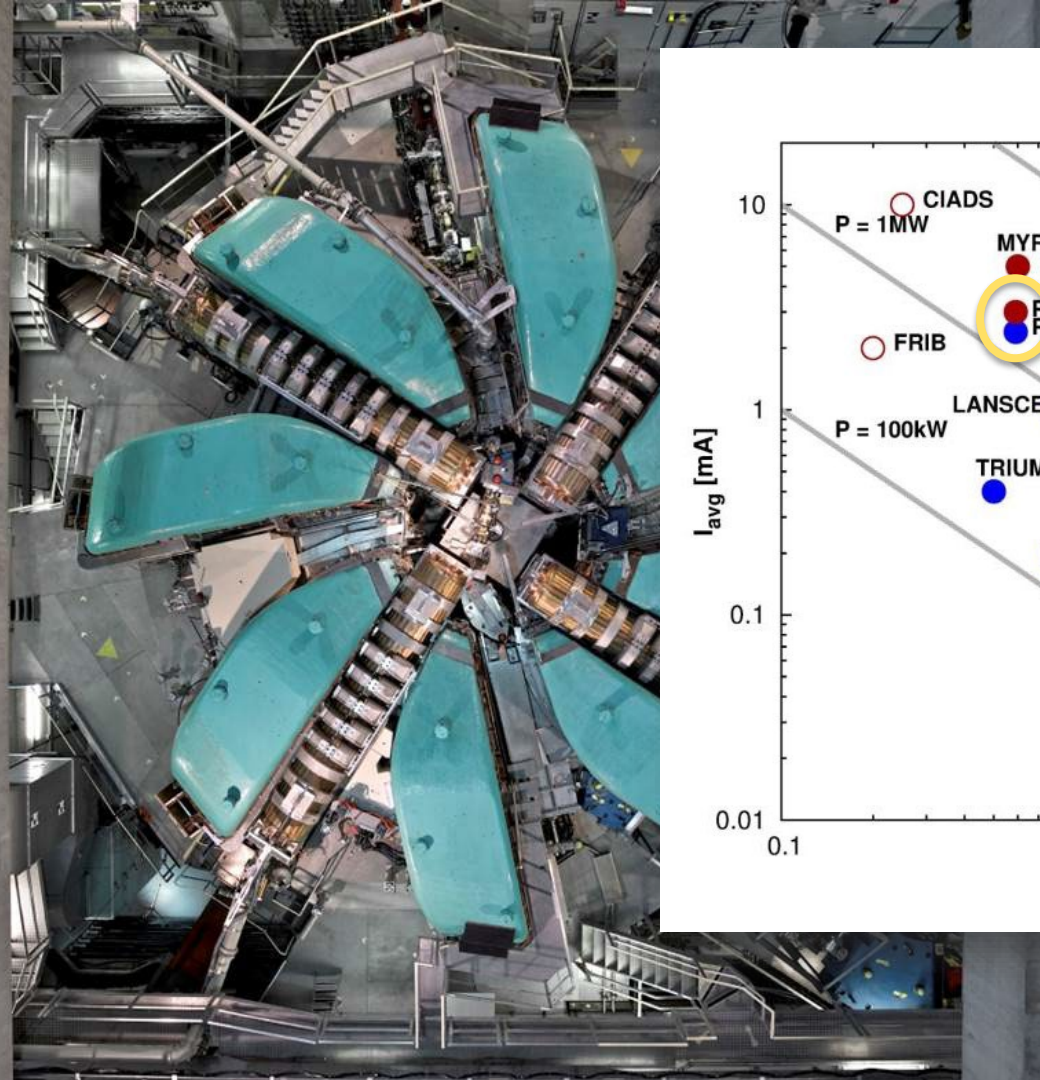
	Fermilab E989	J-PARC	AMMiED	TiAMMo
Muon momentum	30.9 GeV/c	300 MeV/c	125 MeV/c	14.6 MeV/c
Lorentz γ	29.3	3	1.56	1
Polarization	100%	50%	95%	95%
Magnetic field B	1.45 T	3 T	6 T	17 T
Focusing field	E-quadrupoles	weakly focusing (wf)	wf	wf
Cyclotron period	149 ns	7.4 ns	3.8 ns	0.4 ns
AMM precession period	4.37 μ s	2.11 μ s	1.05 μ s	0.37 μ s
Number of detected e^+	1.6×10^{11}	5.7×10^{11}	7×10^{12}	1.7×10^{12}
AMM precision (stat.)	0.1 ppm	0.45 ppm	0.1 ppm	0.08 ppm
AMM precision (sys.)	0.1 ppm	0.07 ppm	0.07 ppm	0.07 ppm



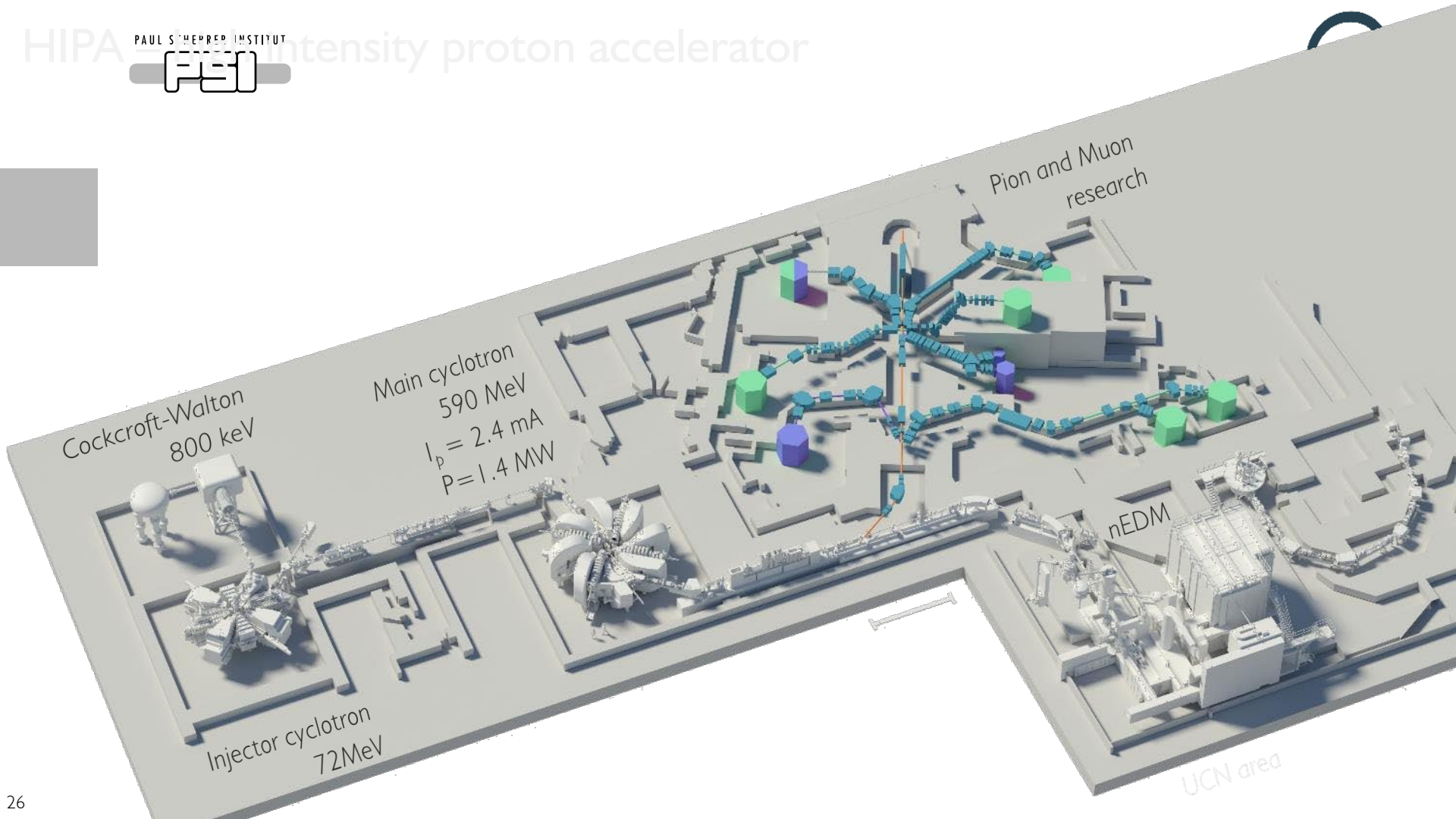
- A dedicated experiment to experimentally search for muon EDM is being set-up at PSI
- In case of null result, it will improve the current experimental upper limit by 3 orders of magnitude, **further improvement with pulsed beams possible (e.g. Fermilab FFA?)**
- The experiment will take place in two phases, where we will demonstrate the frozen-spin technique in phase I
- Optimization of the design for Phase I underway using simulation studies:
prototype building of important components
Test beams demonstrating feasibility
- A tinny g-2 could profit from a re-accelerated low emittance cold muon source
- Muon multiplicity would require much higher efficiency for muCool







courtesy: M. Seidel



Cockcroft-Walton
800 keV

Main cyclotron
590 MeV
 $I_p = 2.4 \text{ mA}$
 $P = 1.4 \text{ MW}$

Pion and Muon
research

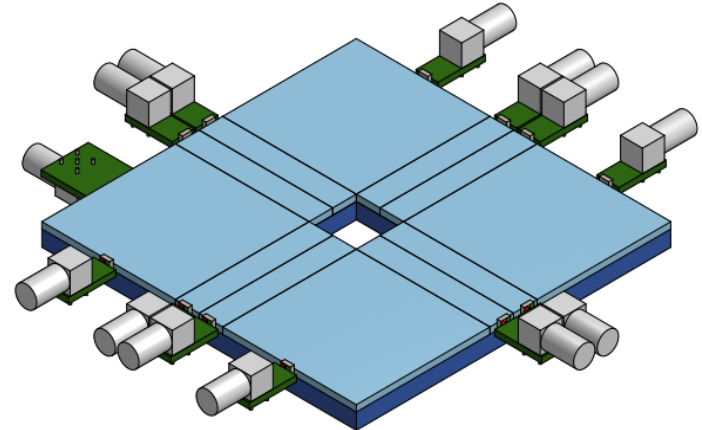
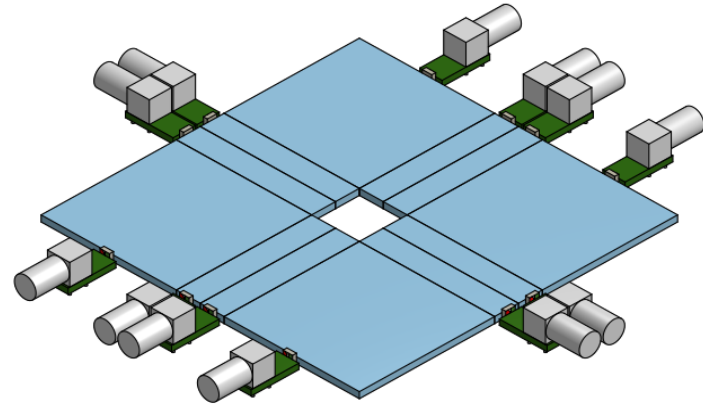
nEDM

UCN area

Injector cyclotron
72 MeV

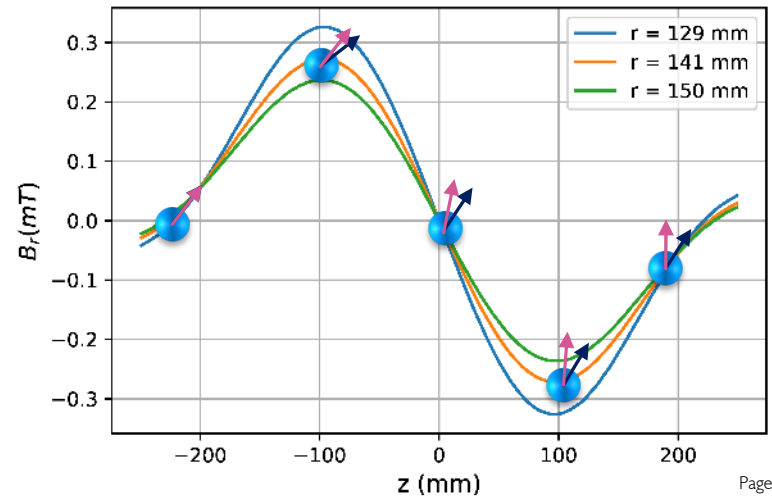
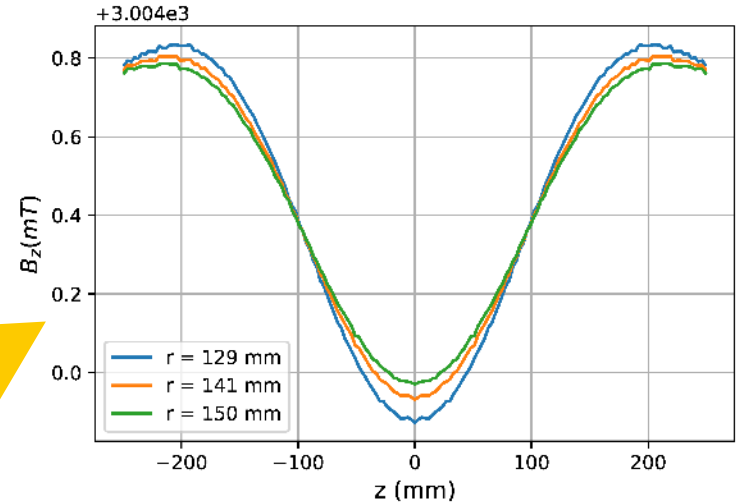
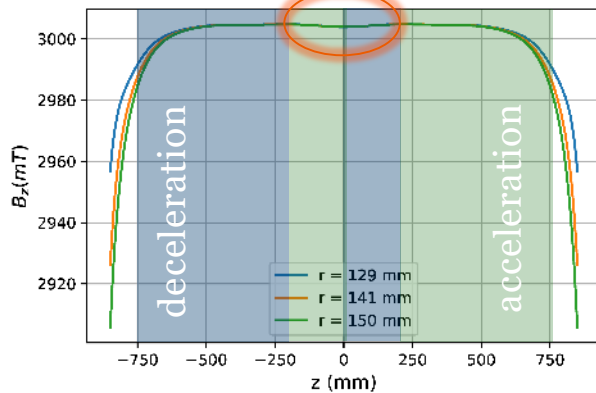


- Focus muon beam onto opening of injection channel
- Scintillator tiles coupled to SiPMs
- Hole in center to let muon beam pass
- Front tile thickness 1-2 mm to stop surface muons
- A thicker (up to ~5 mm) scintillator layer could be added to better discriminate muons and positrons
- Centering procedure optimized in simulation
- Next step, prototype building



Storage and injection

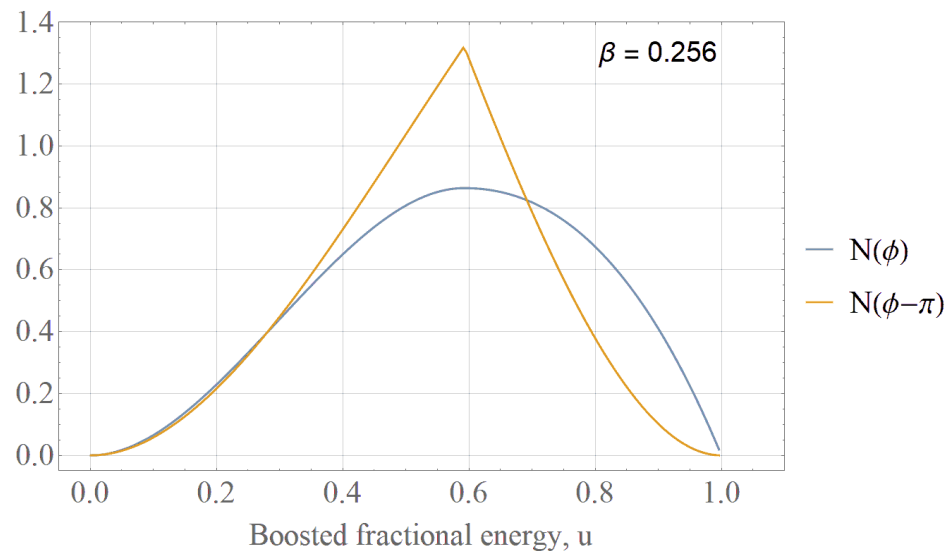
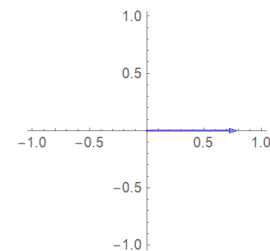
- Strength of weakly focusing field in the center region defines “depth” of storage
- The deeper/stronger the weakly-focusing field, the stronger the pulse needs to be



Tuning the electric field to the frozen-spin condition



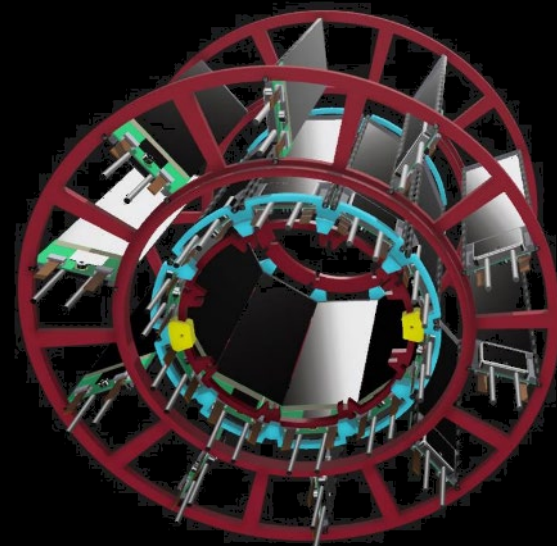
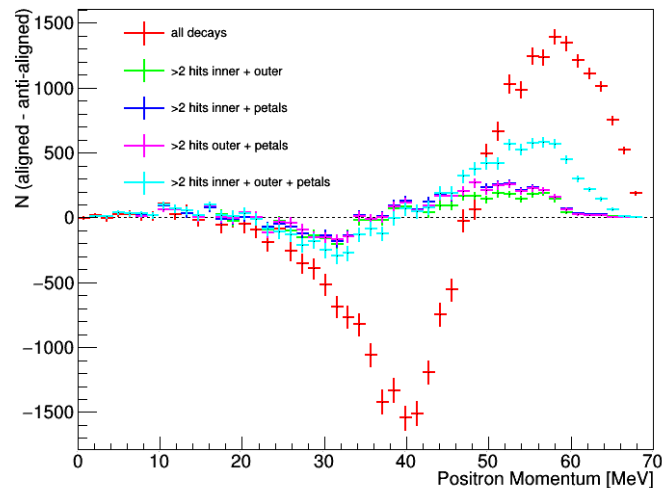
- Measure the g-2 frequency ω_a
- Two momentum bins
 $28 \text{ MeV}/c < p_1 < 50 \text{ MeV}/c$
 $50 \text{ MeV}/c < p_2$
- Change E field in the range
 $\pm E_{\text{frozen}} \approx \pm 3 \text{ kV}/\text{cm}$
- Extrapolate to E_{frozen} where $\omega_a = 0$





Silicon strip detector for g-2 detection

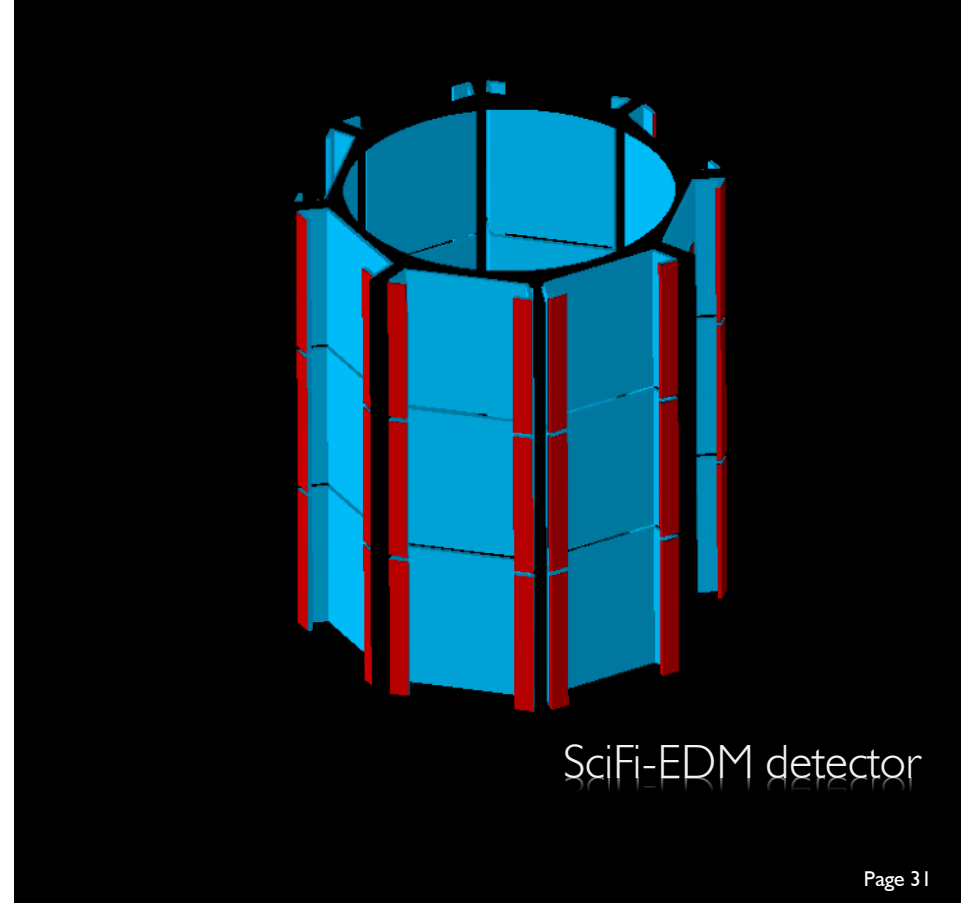
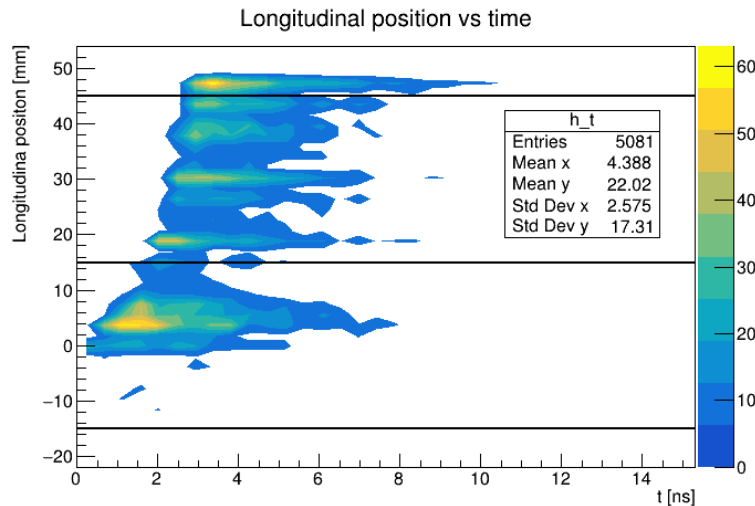
- Reconstruction of transverse positron momentum ($\Delta p \approx 5 \text{ MeV}/c$)
- Timing $\Delta t \approx 2 \text{ ns}$
- Spatial resolution $\approx 0.1 \text{ mm}$ (lateral)



Si-Petal detector

Scintillating fiber detector for EDM
asymmetry measurement and timing

- Horizontal fiber ribbons with **250 μm** pitch and **100 μm** resolution
- Timing resolution **< 2ns**
- Reconstruction of longitudinal momentum

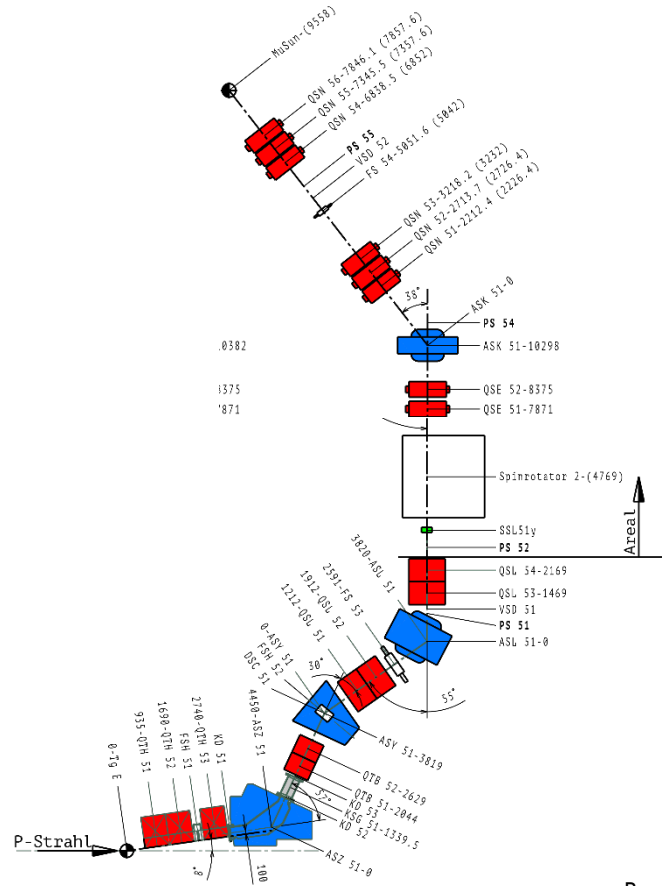




Horizontal Emittance: $200 \pi \text{ mm mrad}$

Vertical Emittance: $170 \pi \text{ mm mrad}$

- Beam rate about $4 \times 10^6 \text{ s}^{-1}$
- Acceptance phase space:
 - High transmission through channel 3%
 - Injection efficiency about 0.4%
 - Expected μ^+ storage rate 2kHz
 - Expected e^+ detection rate 0.5kHz
- Moderate E field 3kV/cm

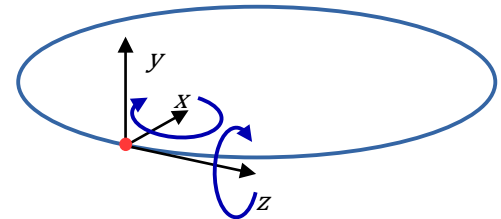




- Systematic effects: all effects that lead to a *real* or *apparent* precession of the spin around the radial axis that are not related to the EDM
- Major sources of systematic effects in the frozen spin technique:
 - Coupling of the magnetic moment with the EM fields of the experimental setup (*real*)
 - Early to late variation of detection efficiency of the EDM detectors (*apparent*)

$$\vec{\Omega}_{\text{MDM}} = -\frac{e}{m_0} \left[a\vec{B} - a\frac{\gamma-1}{\gamma} \frac{(\vec{\beta} \cdot \vec{B})\vec{\beta}}{\beta^2} + \left(\frac{1}{\gamma^2-1} - a \right) \frac{\vec{\beta} \times \vec{E}}{c} \right]$$

- Rotations that could mimic the EDM:
 - Radial around x
 - Azimutal around z

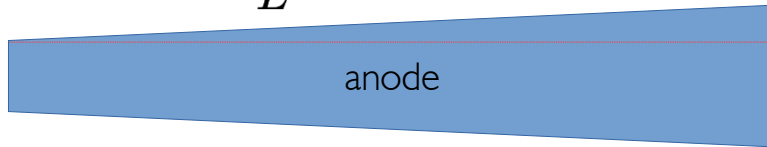


Sources of E_y field: conical central electrode

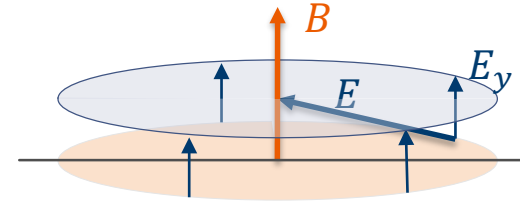


- None constant radius of cylindrical anode (cone)

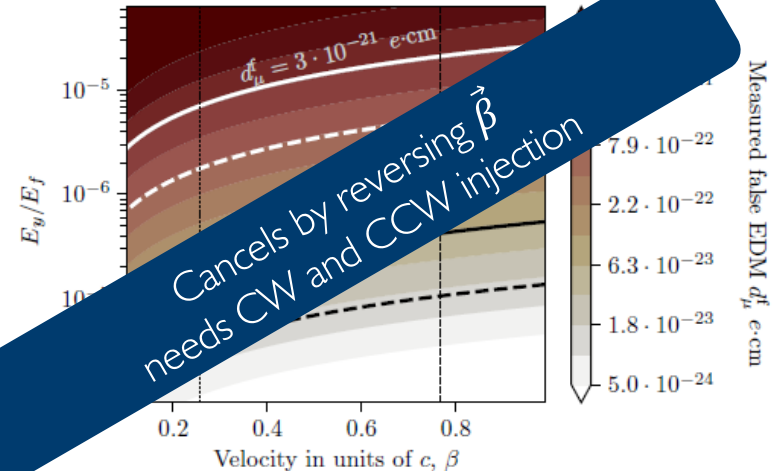
$$E_y \approx E_f \frac{\Delta R}{L} \approx E_f \alpha$$



- Cylindricity on the order of 50 nm is measurable even on large samples and possible to machine
- Ground electrode made of thin foil more difficult to keep deviations from cylindricity below $30\mu\text{m}$

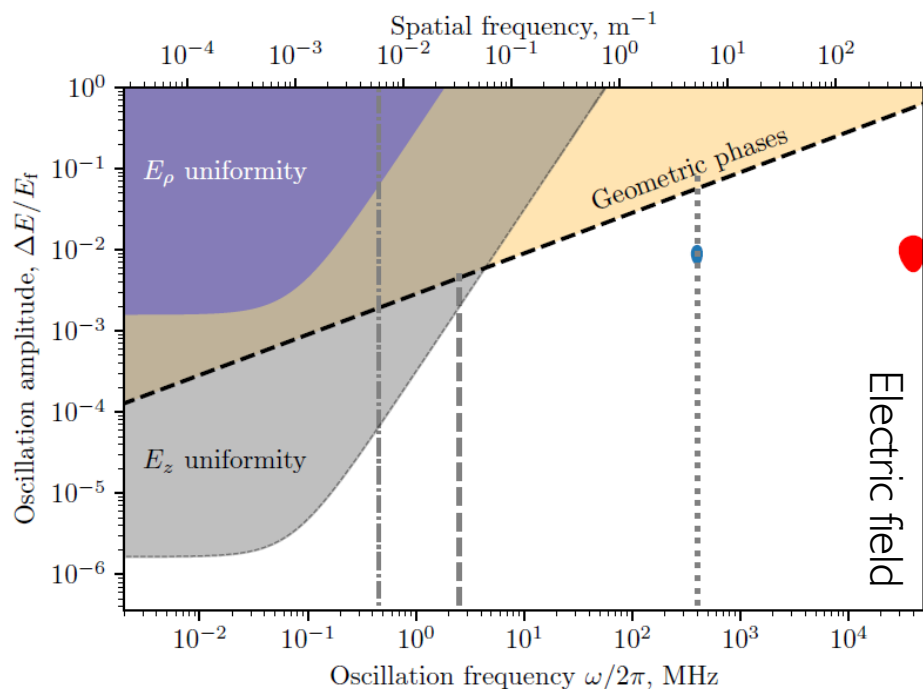
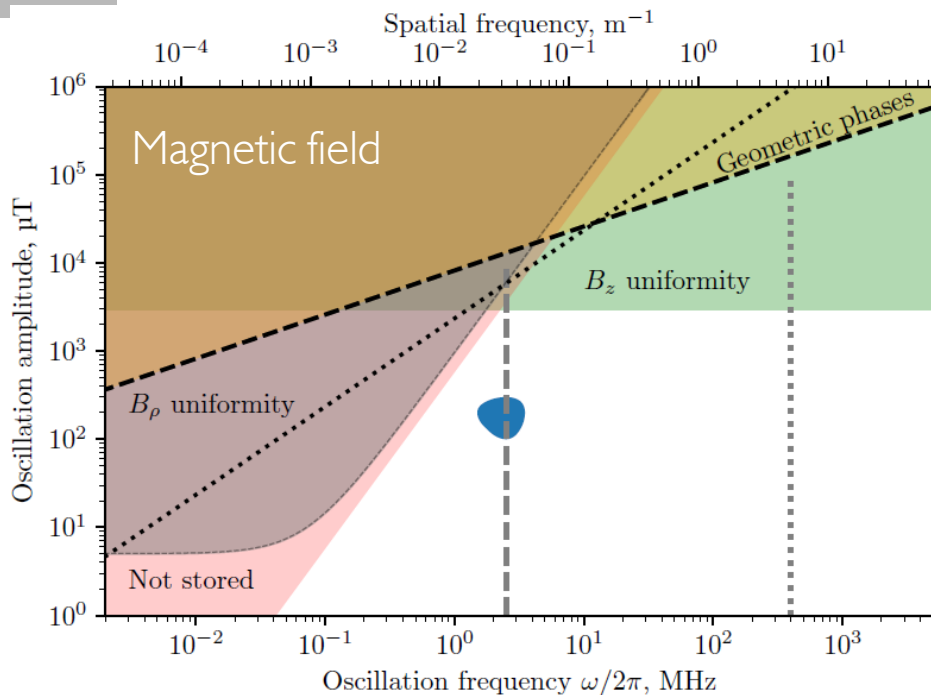


Will move orbit out of central plane until:
 $\langle B_r^* \rangle = -\langle E_v / \beta \gamma \rangle$



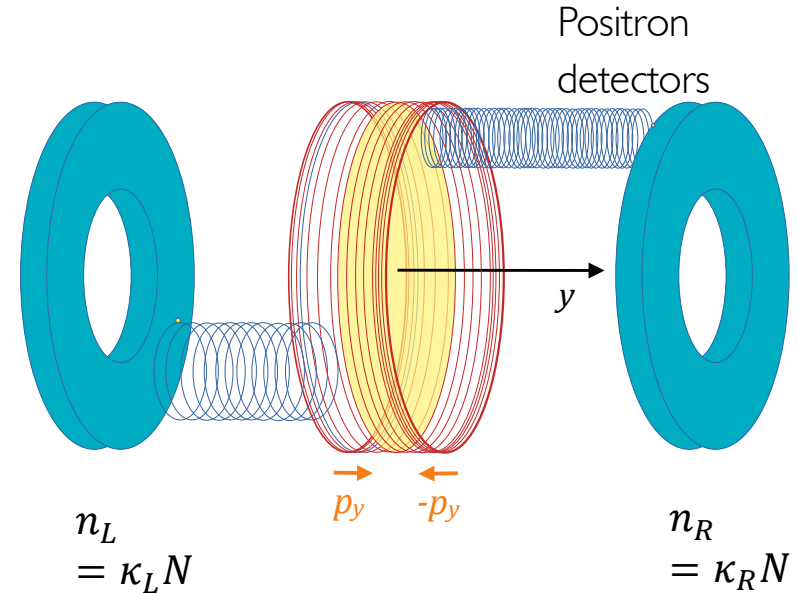


$$\frac{1}{4} \frac{\Omega_x \Omega_y}{\omega_x \omega_y} \left[\frac{\omega_x - \omega_y}{\omega_x + \omega_y} \cos((\omega_x + \omega_y)t + \beta_0) - \frac{\omega_x + \omega_y}{\omega_x - \omega_y} \cos((\omega_y - \omega_x)t + \beta_0) \right] \rightarrow -\frac{1}{2\omega} \Omega_x \Omega_y t \sin(\beta_0)$$





- The EDM will be deduced from the accumulation of asymmetry between the upstream and downstream detectors that increases with time
- Static differences in the detection efficiency of one detector compared to the other is not a problem
- Change of the detection efficiency with time is a problem as it will introduce time dependent asymmetry





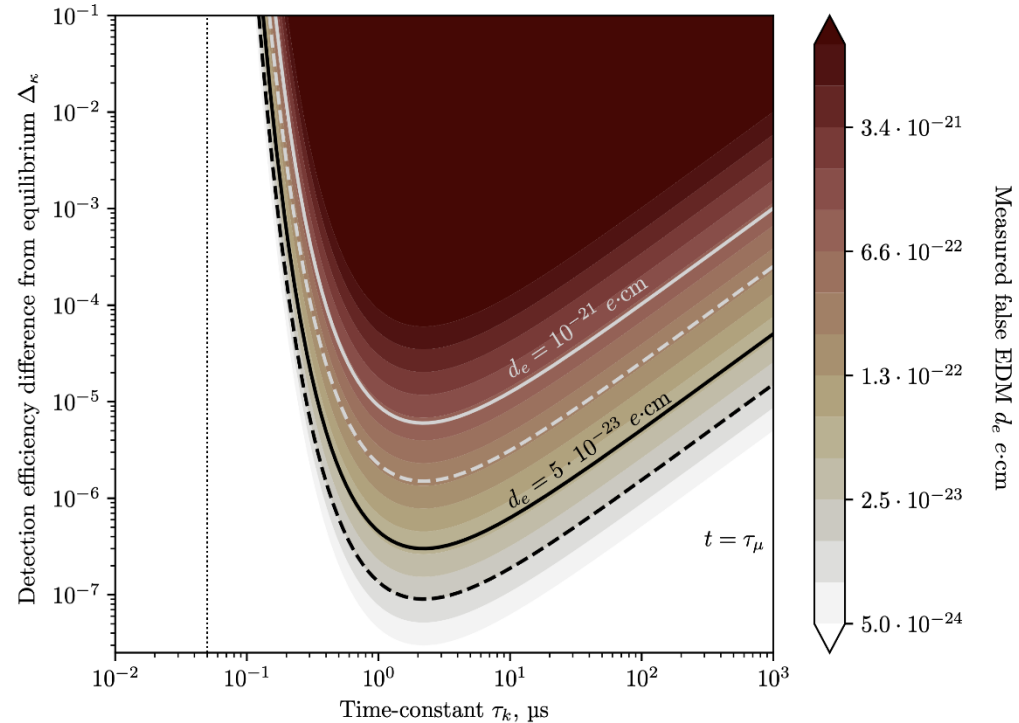
- Assumption: Change of detection efficiency triggered by pulse, exponential decay
- Detection efficiency of up and downstream detectors:

$$\kappa_u = \kappa_{u0} - \Delta_\kappa e^{-t/\tau_k},$$

$$\kappa_d = \kappa_{d0} + \Delta_\kappa e^{-t/\tau_k},$$

- Change in measured asymmetry with time:

$$\dot{A}_m = \frac{2}{\tau_k} \Delta_\kappa e^{-t/\tau_k}$$





- Systematic effects are studied using analytic expressions
- Comparison with GEANT4 spin tracking Monte Carlo for verification
- Deduce specifications for experiment

Next steps:

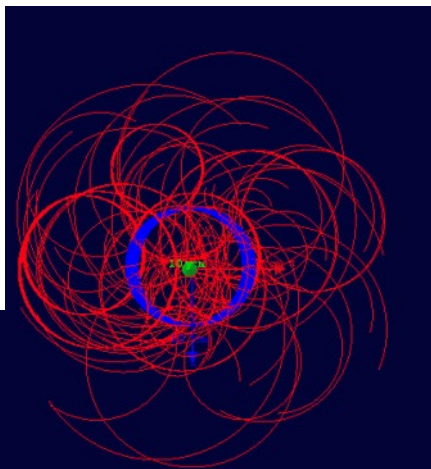
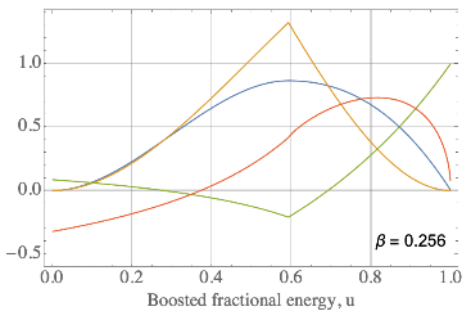
- Parametrization of magnetic-field non-uniformity
- Deduce magnetic-field requirements

Systematic effect	Constraints	Phase I	
		Expected value	Syst. ($\times 10^{-21} \text{ e}\cdot\text{cm}$)
Cone shaped electrodes (longitudinal E-field)	Up-down asymmetry in the electrode shape	$\Delta_R < 30 \mu\text{m}$	0.75
Residual B-field from kick	Decay time of kicker field	$< 50 \text{ ns}$	$< 10^{-2}$
Net current flowing muon orbit area	Wiring of electronics inside the orbit	$< 10 \text{ mA}$	$< 10^{-2}$
Longitudinal B-field uniformity	Solenoid alignment	$< 3 \text{ mT}$	-
Resonant geometrical phase accumulation	Misalignment of central axes	Pitch $< 1 \text{ mrad}$ Offset $< 2 \text{ mm}$	2×10^{-2}
TOTAL			1.1



Phase I

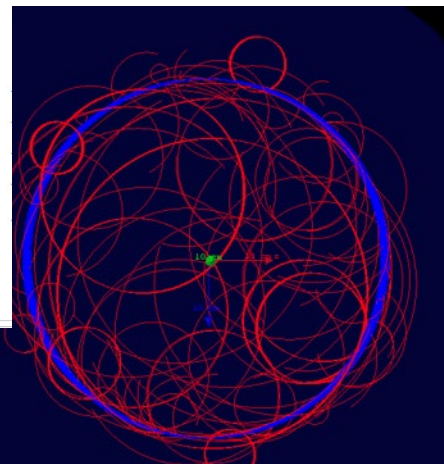
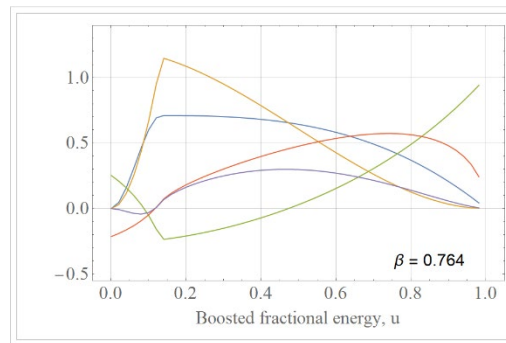
- B-Field 3T
- Momentum 28 MeV/c
- Muon radius 31 mm
- Most positrons outside



$$\leq 3 \times 10^{-21} \text{ ecm}$$

Phase II

- B-Field 3T
- Momentum 125 MeV/c
- Muon radius 141 mm
- Most positrons inside



$$\leq 6 \times 10^{-23} \text{ ecm}$$

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January 24, 2023



Conclusions and outlook

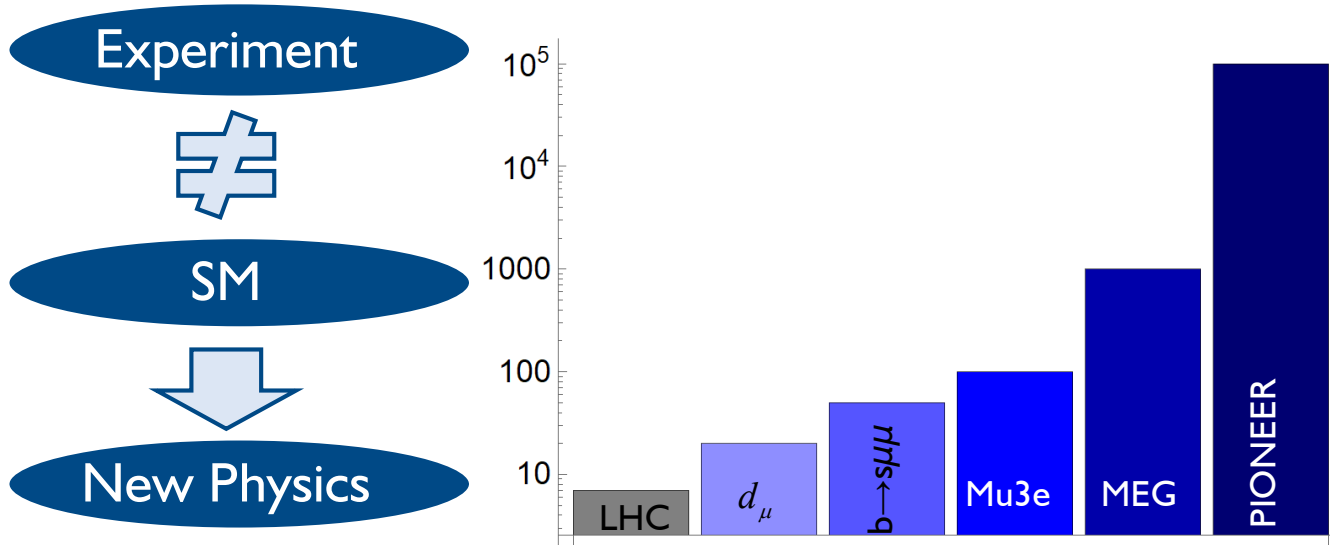
- At PSI we are setting up an experiment to demonstrate the frozen-spin technique with a sensitivity of a few 10^{-21} ecm to the EDM of the muon
- In the next four years we will demonstrate all necessary techniques, and design a dedicated hi-fidelity solenoid
- Using muons with a moment larger than 125 MeV/c and the new solenoid we will reach a sensitivity of 6×10^{-23} ecm



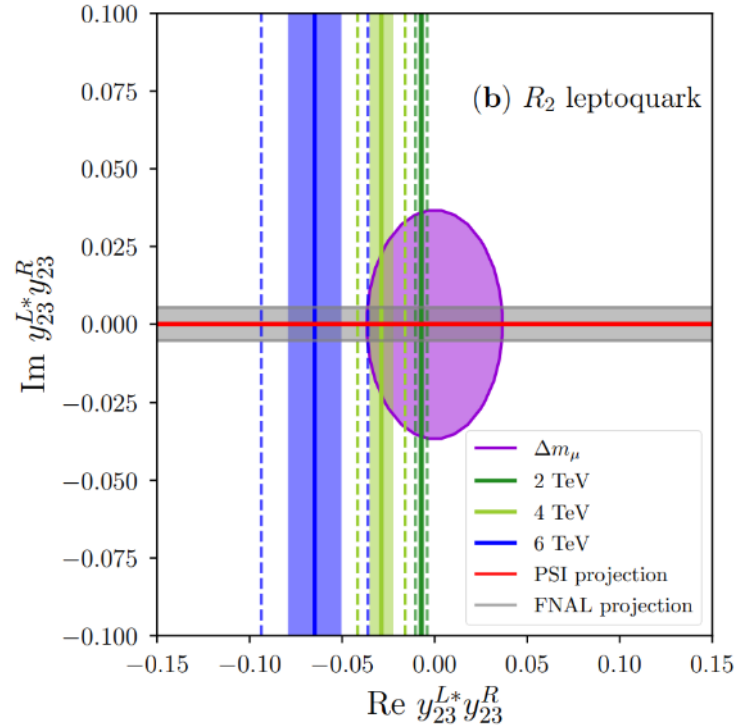
Backup



- At colliders one produces many (up to 10^{14}) heavy quarks or leptons and measures their decays into light flavors

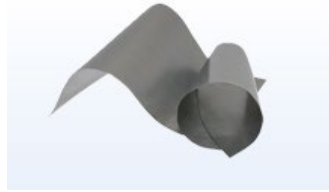


Flavor observables are sensitive to higher energy scales than collider searches

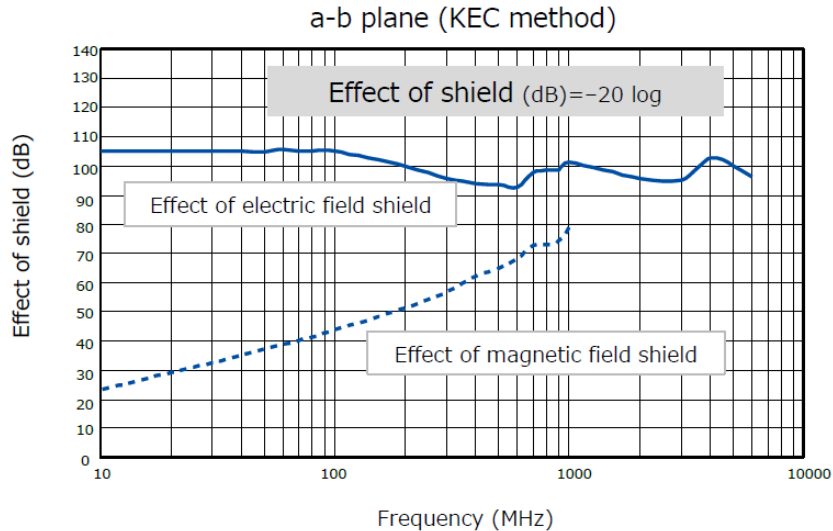


Bigaran, Volkas, 2110.03707

No significant tuning necessary



- Exist off the shelf without substrate down to $17\mu\text{m}$
- Still considerable damping of magnetic pulse possible
- Tests requires
- Alternative one dimensional wires (carbon fibers / tungsten)





- Characterization of potential electrode material with positrons and muons

$$50 \text{ MeV}/c < p < 145 \text{ MeV}/c$$

μ^+, e^+

