

Theory overview of muon $g-2$ and EDM

Paride Paradisi

University of Padova and INFN

Muon4Future workshop
May 29th-31st 2023, Venice - IT

1 The high-intensity frontier

- ▶ Experimental status/prospects

2 “Old muon $g-2$ puzzle” (pre BMW 2021)

- ▶ Possible new physics interpretations

3 Leptonic $g-2$, EDMs & LFV interrelationship

- ▶ Model independent considerations

4 “New muon $g-2$ puzzle” (post BMW 2021)

- ▶ Possible new physics interpretations

5 Conclusions and future prospects

Where to look for New Physics at low-energy?

- Processes very **suppressed** or even **forbidden** in the SM

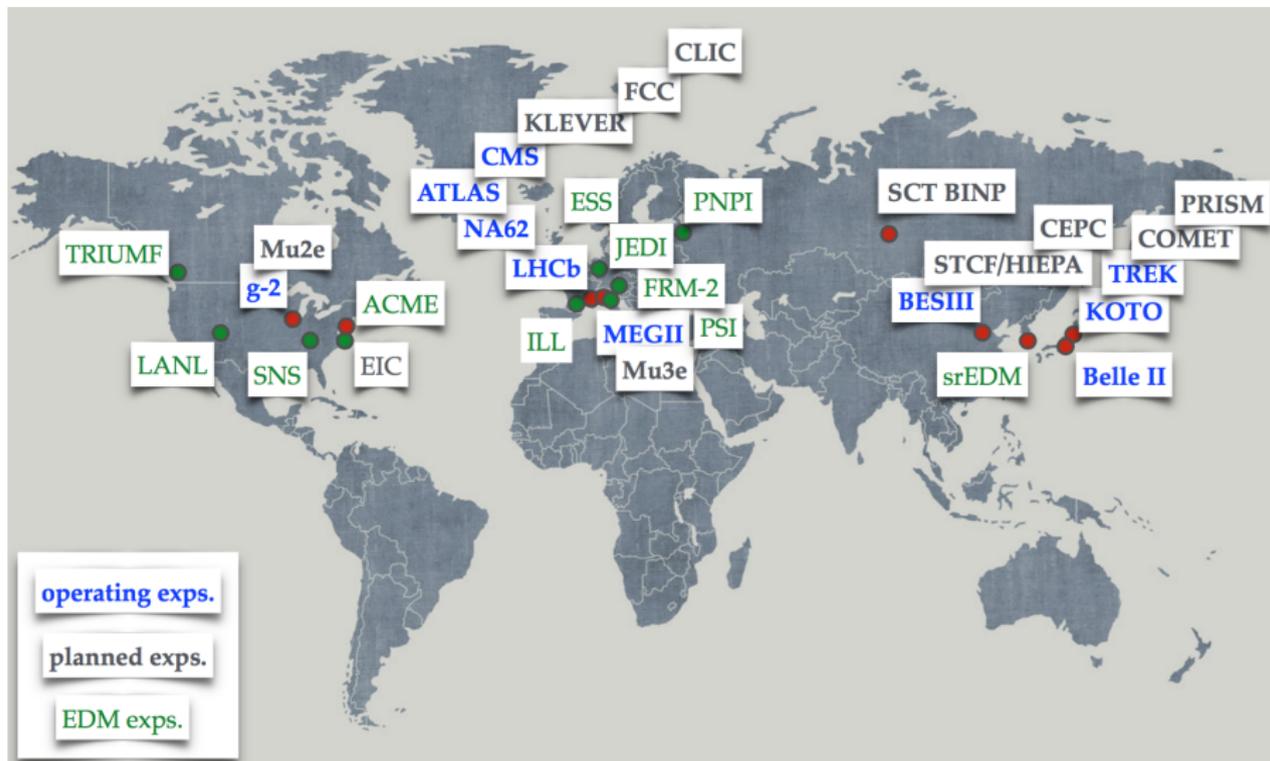
- ▶ LFV processes ($\mu \rightarrow e\gamma$, $\mu \rightarrow e$ in N, $\tau \rightarrow \mu\gamma$, $\tau \rightarrow 3\mu, \dots$)
- ▶ CPV effects in the leptonic (e, μ) and neutron EDMs
- ▶ FCNC & CPV in $B_{s,d}$ & D decay/mixing amplitudes

- Processes predicted with **high precision** in the SM

- ▶ EWPO as $(g-2)_\mu$: $\Delta a_\mu = a_\mu^{\text{exp}} - a_\mu^{\text{SM}} = (2.51 \pm 0.59) \times 10^{-9}$ (4.2σ discrepancy!)
- ▶ LFUV in $M \rightarrow \ell\nu$ (with $M = \pi, K, B$), $B \rightarrow D^{(*)}\ell\nu$, $B \rightarrow K\ell\ell'$, τ and Z decays

- **High-intensity frontier**: A collective effort to determine the NP symmetries

Experimental status



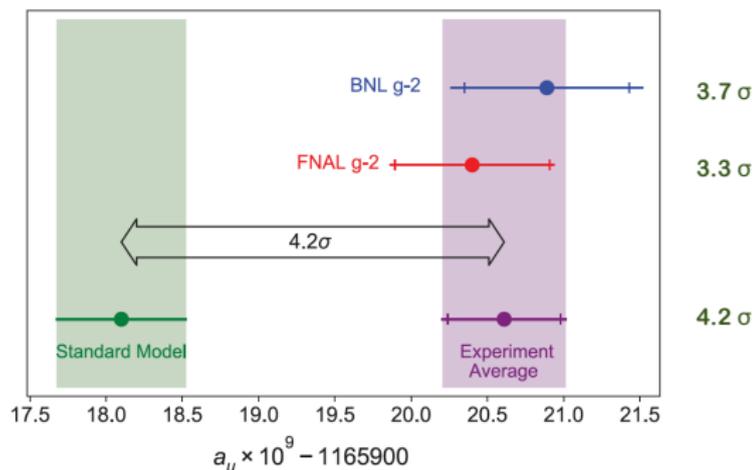
Process	Present	Experiment	Future	Experiment
$\mu \rightarrow e\gamma$	4.2×10^{-13}	MEG	$\approx 6 \times 10^{-14}$	MEG II
$\mu \rightarrow 3e$	1.0×10^{-12}	SINDRUM	$\approx 10^{-16}$	Mu3e
$\mu^- \text{ Au} \rightarrow e^- \text{ Au}$	7.0×10^{-13}	SINDRUM II	?	
$\mu^- \text{ Ti} \rightarrow e^- \text{ Ti}$	4.3×10^{-12}	SINDRUM II	?	
$\mu^- \text{ Al} \rightarrow e^- \text{ Al}$	—		$\approx 10^{-16}$	COMET, MU2e
$\tau \rightarrow e\gamma$	3.3×10^{-8}	Belle & BaBar	$\sim 10^{-9}$	Belle II
$\tau \rightarrow \mu\gamma$	4.4×10^{-8}	Belle & BaBar	$\sim 10^{-9}$	Belle II
$\tau \rightarrow 3e$	2.7×10^{-8}	Belle & BaBar	$\sim 10^{-10}$	Belle II
$\tau \rightarrow 3\mu$	2.1×10^{-8}	Belle & BaBar	$\sim 10^{-10}$	Belle II
$d_e(\text{e cm})$	1.1×10^{-29}	ACME	$\sim 3 \times 10^{-31}$	ACME III
$d_\mu(\text{e cm})$	1.8×10^{-19}	Muon (g-2)	$\sim 10^{-22}$	PSI

Table: Present and future experimental sensitivities for relevant low-energy observables.

- So far, only upper bounds. Still excellent prospects for exp. improvements.
- We can expect a NP signal in all above observables below the current bounds.

Experimental status of the muon $g - 2$

- **April 7th 2021: Muon $g - 2$ experiment at FNAL confirms BNL!**



$$a_\mu^{\text{EXP}} = (116592089 \pm 63) \times 10^{-11} [0.54\text{ppm}] \text{ BNL E821}$$

$$a_\mu^{\text{EXP}} = (116592040 \pm 54) \times 10^{-11} [0.46\text{ppm}] \text{ FNAL E989 Run 1}$$

$$a_\mu^{\text{EXP}} = (116592061 \pm 41) \times 10^{-11} [0.35\text{ppm}] \text{ WA}$$

- **FNAL aims at 16×10^{-11} . A new FNAL release expected by this summer!** [see Foster's talk]
- **Muon $g - 2$ proposal at J-PARC: Phase-1 with similar BNL precision.**

New Physics for the muon $g - 2$: at which scale?

- Δa_μ discrepancy at $\sim 4.2 \sigma$ level:

$$\Delta a_\mu = a_\mu^{\text{EXP}} - a_\mu^{\text{SM}} \equiv a_\mu^{\text{NP}} = (2.51 \pm 0.59) \times 10^{-9}$$

$$\Delta a_\mu \equiv a_\mu^{\text{NP}} \approx (a_\mu^{\text{SM}})_{\text{weak}} \approx \frac{m_\mu^2}{16\pi^2 v^2} \approx 2 \times 10^{-9}$$

- ▶ NP is at the weak scale ($\Lambda \approx v$) and weakly coupled to SM particles.*
- ▶ NP is very light ($\Lambda \lesssim 1 \text{ GeV}$) and feebly coupled to SM particles.
- ▶ NP is very heavy ($\Lambda \gg v$) and strongly coupled to SM particles.

*Favoured by the *hierarchy problem* and by a WIMP DM candidate but disfavoured by the LEP and LHC bounds (supersymmetry being the most prominent example).

[For a through compilation of models, see Athron, Balazs, Jacob, Kotlarski, Stockinger, Stockinger-Kim, '21.]

$\Lambda \approx \nu$: SUSY and the muon ($g - 2$)

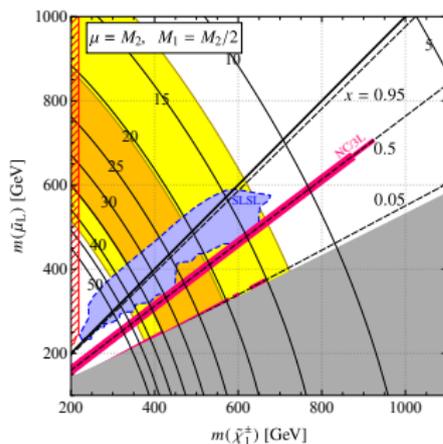
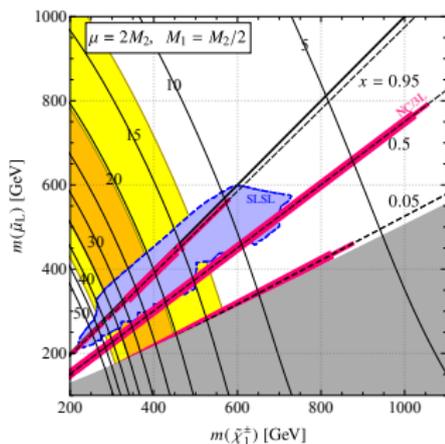
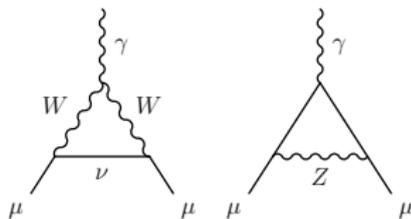
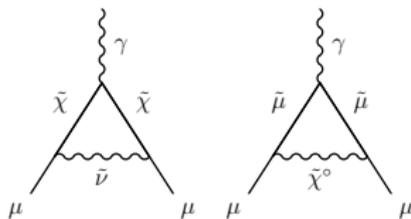


Figure: LHC Run 2 bounds on SUSY scenario for the muon $g - 2$ anomaly for $\tan \beta = 40$. Orange (yellow) regions satisfy the muon $g - 2$ anomaly at the 1σ (2σ) level [Endo et al., '20].



$$(a_\mu^{\text{SM}})_{\text{weak}} \approx \frac{g^2 m_\mu^2}{32\pi^2 M_W^2} \approx 2 \times 10^{-9}$$



$$a_\mu^{\text{SUSY}} \approx \frac{g^2 m_\mu^2 \tan \beta}{32\pi^2 \tilde{m}^2} \approx 2 \times 10^{-9}$$

$\tilde{m} = 500 \text{ GeV}$ & $\tan \beta = 40$

Axion-like Particle effective Lagrangian

$$\mathcal{L} = e^2 C_{\gamma\gamma} \frac{a}{\Lambda} F_{\mu\nu} \tilde{F}^{\mu\nu} + \frac{c_{\mu\mu}}{2} \frac{\partial^\nu a}{\Lambda} \bar{\mu} \gamma_\nu \gamma_5 \mu$$

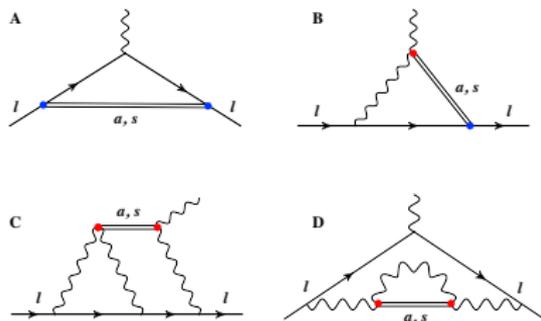


Figure: Contributions of a scalar ‘s’ and a pseudoscalar ‘a’ ALP to the $(g - 2)_\ell$.

[Marciano, Masiero, PP, Passera '16]

[Cornella, P.P., Sumensari '19]

$$\Delta a_\mu = \frac{m_\mu^2}{\Lambda^2} \left[\frac{12\alpha^3}{\pi} C_{\gamma\gamma}^2 \ln^2 \frac{\Lambda^2}{m_\mu^2} - \frac{(c_{\mu\mu})^2}{16\pi^2} h_1 \left(\frac{m_a^2}{m_\mu^2} \right) - \frac{2\alpha}{\pi} c_{\mu\mu} C_{\gamma\gamma} \ln \frac{\Lambda^2}{m_\mu^2} \right]$$

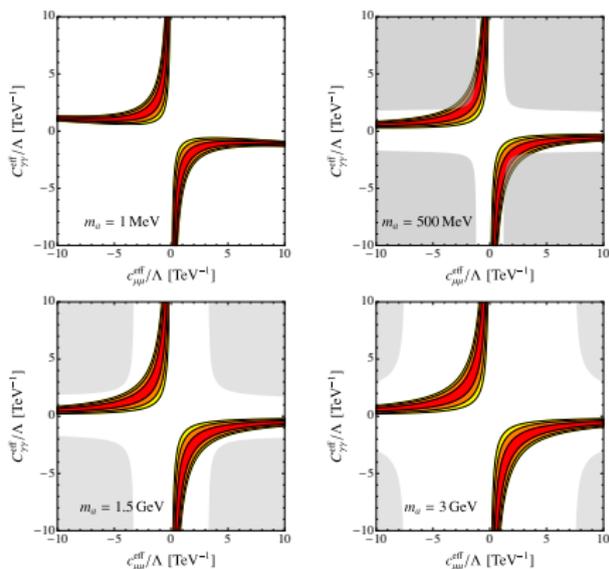


Figure: Δa_μ regions favoured at 68% (red), 95% (orange) and 99% (yellow) CL. Gray regions are excluded by the BaBar search $e^+e^- \rightarrow \mu^+\mu^- + \mu^+\mu^-$ [Bauer, Neubert, Thamm, '17]

- SMEFT Lagrangian relevant for Δa_μ** [Buttazzo and P.P., '20]

$$\mathcal{L} = \sum_{V=B,W} \frac{C_{eV}^\ell}{\Lambda^2} (\bar{\ell}_L \sigma^{\mu\nu} e_R) H V_{\mu\nu} + \sum_{q=c,t} \frac{C_T^{\ell q}}{\Lambda^2} (\bar{\ell}_L \sigma_{\mu\nu} e_R) (\bar{Q}_L \sigma^{\mu\nu} q_R) + h.c.$$

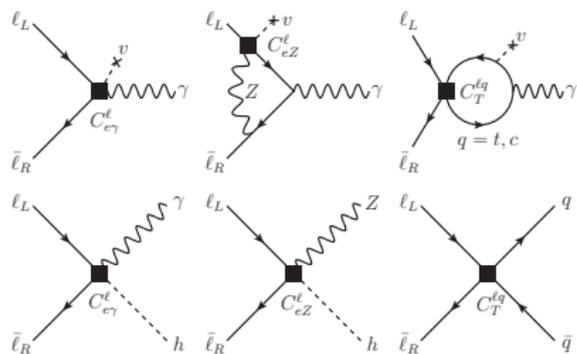


Figure: SMEFT Feynman diagrams for the $g-2$ (upper row) and scattering processes (lower row): $H = v + h/\sqrt{2}$

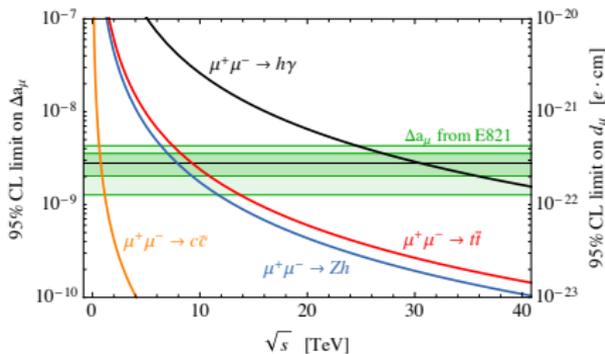


Figure: 95% C.L. reach on Δa_μ vs \sqrt{s} from various processes assuming $\mathcal{L} = (\sqrt{s}/10 \text{ TeV})^2 \times 10 \text{ ab}^{-1}$.

$$\Delta a_\mu \sim \frac{m_\mu v}{\Lambda^2} C_{eV,T} \iff \sigma_{\mu\mu \rightarrow f} \sim \frac{s}{\Lambda^4} |C_{eV,T}|^2 \quad (f = e\gamma, eZ, q\bar{q})$$

- At high energy $\sigma_{\mu\mu \rightarrow f}$ can compete with Δa_μ to test the very same NP!**

- **NP effects are encoded in the effective Lagrangian**

$$\mathcal{L} = e \frac{m_\ell}{2} (\bar{\ell}_R \sigma_{\mu\nu} A_{\ell\ell'} \ell'_L + \bar{\ell}'_L \sigma_{\mu\nu} A_{\ell\ell'}^* \ell_R) F^{\mu\nu} \quad \ell, \ell' = e, \mu, \tau,$$

- ▶ **Branching ratios of $\ell \rightarrow \ell' \gamma$** [see Teixeira's and Renner's talks]

$$\frac{\text{BR}(\ell \rightarrow \ell' \gamma)}{\text{BR}(\ell \rightarrow \ell' \nu_\ell \bar{\nu}_{\ell'})} = \frac{48\pi^3 \alpha}{G_F^2} (|A_{\ell\ell'}|^2 + |A_{\ell'\ell}|^2).$$

- ▶ **Δa_ℓ and leptonic EDMs**

$$\Delta a_\ell = 2m_\ell^2 \text{Re}(A_{\ell\ell}), \quad \frac{d_\ell}{e} = m_\ell \text{Im}(A_{\ell\ell}).$$

- ▶ **“Naive scaling”**: a broad class of NP theories contributes to Δa_ℓ and d_ℓ as

$$\frac{\Delta a_\ell}{\Delta a_{\ell'}} = \frac{m_\ell^2}{m_{\ell'}^2}, \quad \frac{d_\ell}{d_{\ell'}} = \frac{m_\ell}{m_{\ell'}}.$$

- $\text{BR}(\ell_i \rightarrow \ell_j \gamma)$ vs. $(g - 2)_\mu$

$$\text{BR}(\mu \rightarrow e \gamma) \approx 3 \times 10^{-13} \left(\frac{\Delta a_\mu}{3 \times 10^{-9}} \right)^2 \left(\frac{\theta_{e\mu}}{10^{-5}} \right)^2$$

$$\text{BR}(\tau \rightarrow \mu \gamma) \approx 4 \times 10^{-8} \left(\frac{\Delta a_\mu}{3 \times 10^{-9}} \right)^2 \left(\frac{\theta_{\mu\tau}}{10^{-2}} \right)^2$$

- EDMs vs. $(g - 2)_\mu$

$$d_e \approx \left(\frac{\Delta a_\mu}{3 \times 10^{-9}} \right) 10^{-29} \left(\frac{\phi_e^{CPV}}{10^{-5}} \right) e \text{ cm},$$

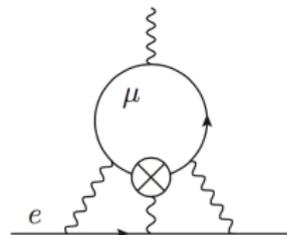
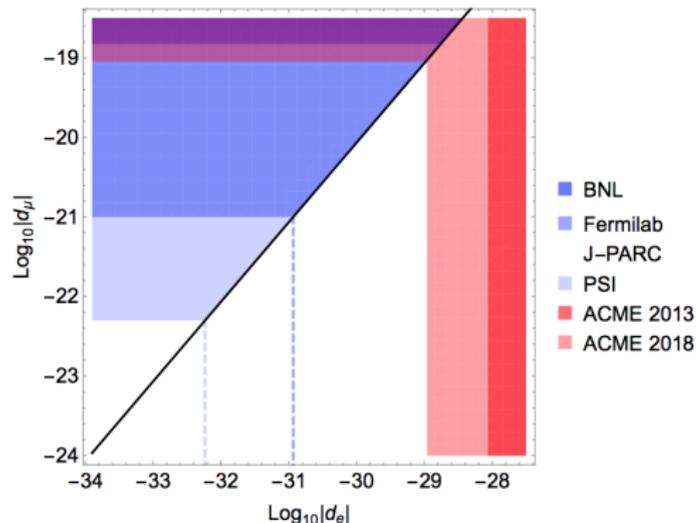
$$d_\mu \approx \left(\frac{\Delta a_\mu}{3 \times 10^{-9}} \right) 2 \times 10^{-22} \phi_\mu^{CPV} e \text{ cm},$$

- **Main messages:**

- ▶ $\Delta a_\mu \approx (3 \pm 1) \times 10^{-9}$ requires a nearly flavor and CP conserving NP
- ▶ Large effects in the muon EDM $d_\mu \sim 10^{-22} e \text{ cm}$ are still allowed!

[Giudice, P.P., & Passera, '12]

Experimental status of the muon EDM



$$d_\mu \leq 10^{-21} \text{ e cm} \left(\frac{d_e}{10^{-31} \text{ e cm}} \right)$$

[Crivellin, Hoferichter & Schmidt-Wellenburg, '18, see Schmidt-Wellenburg's talk]

$$d_\mu \simeq \left(\frac{\Delta a_\mu}{3 \times 10^{-9}} \right) 2 \times 10^{-22} \phi_\mu^{CPV} \text{ e cm},$$

[Giudice, PP & Passera, '12]

- Longstanding muon $g - 2$ anomaly

$$\Delta a_\mu = a_\mu^{\text{EXP}} - a_\mu^{\text{SM}} \equiv a_\mu^{\text{NP}} = (2.51 \pm 0.59) \times 10^{-9}$$

$$\Delta a_\mu \equiv a_\mu^{\text{NP}} \approx (a_\mu^{\text{SM}})_{\text{weak}} \approx \frac{m_\mu^2}{16\pi^2 v^2} \approx 2 \times 10^{-9}$$

- Testing the muon $g - 2$ anomaly through the electron $g - 2$

$$\frac{\Delta a_e}{\Delta a_\mu} = \frac{m_e^2}{m_\mu^2} \iff \Delta a_e = \left(\frac{\Delta a_\mu}{3 \times 10^{-9}} \right) 0.7 \times 10^{-13}$$

- ▶ a_e has never played a role in testing NP effects. From $a_e^{\text{SM}}(\alpha) = a_e^{\text{EXP}}$, we extract α which was the most precise value of α up to 2018!
 - ▶ The situation has now changed thanks to th. and exp. progresses.
 - ▶ α can be extracted from atomic physics and a_e used to perform NP tests!
- $\Delta a_e \lesssim 10^{-13}$ will bring a_e to play a pivotal role in probing new physics!

[Giudice, P.P. & Passera, '12]

- **Status of Δa_e as of 2012**

$$\Delta a_e = a_e^{\text{EXP}} - a_e^{\text{SM}} = -9.2(8.1) \times 10^{-13},$$
$$\delta a_e \times 10^{13} : (0.6)_{\text{QED4}}, (0.4)_{\text{QED5}}, (0.2)_{\text{HAD}}, (7.6)_{\delta\alpha}, (2.8)_{\delta a_e^{\text{EXP}}}.$$

- ▶ The errors from QED4 and QED5 will be reduced soon to 0.1×10^{-13} [Kinoshita]
- ▶ We expect a reduction of δa_e^{EXP} to a part in 10^{-13} (or better). [Gabrielse]
- ▶ Work is also in progress for a significant reduction of $\delta\alpha$. [Nez]

- **Status of Δa_e as of 2018: 2.4σ discrepancy** [Parker et al., Science, '18]

$$\Delta a_e = a_e^{\text{EXP}} - a_e^{\text{SM}}(\alpha_{\text{Berkeley}}) = -8.8(3.6) \times 10^{-13}$$
$$\delta a_e \times 10^{13} : (0.1)_{\text{QED5}}, (0.1)_{\text{HAD}}, (2.3)_{\delta\alpha}, (2.8)_{\delta a_e^{\text{EXP}}}.$$

- **Status of Δa_e as of 2020: 1.6σ discrepancy** [Morel et al., Nature, '20]

$$\Delta a_e = a_e^{\text{EXP}} - a_e^{\text{SM}}(\alpha_{\text{LKB2020}}) = 4.8(3.0) \times 10^{-13}$$
$$\delta a_e \times 10^{13} : (0.1)_{\text{QED5}}, (0.1)_{\text{HAD}}, (0.9)_{\delta\alpha}, (2.8)_{\delta a_e^{\text{EXP}}}.$$

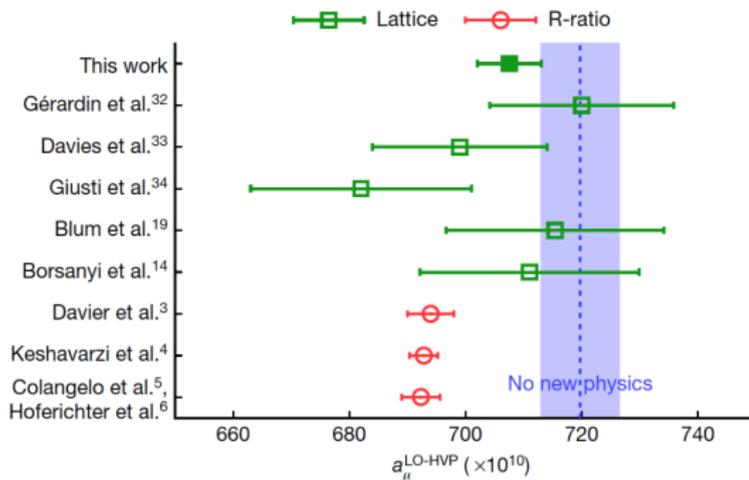
- **Status of Δa_e as of 2022: $(2.8)_{\delta a_e^{\text{EXP}}}$ reduced to $(1.3)_{\delta a_e^{\text{EXP}}}$** [Gabrielse et al., '22]

HLO contribution from lattice QCD

- Great progress also in lattice QCD, where spacetime is modeled as a discrete grid of points. The BMW collaboration reached a 0.8% precision!

$$a_\mu^{\text{HLO}} = 7075(23)_{\text{stat}}(50)_{\text{syst}} [55]_{\text{tot}} \times 10^{-11}$$

- 2–2.5 σ tension with the “data-driven” evaluations.



Borsanyi et al (BMWc), Nature 2021

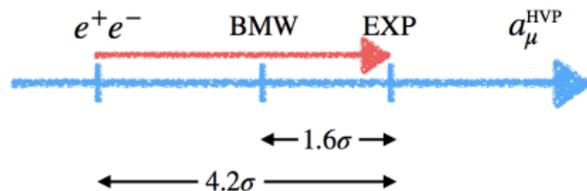
[see Colangelo's and Fodor's talks]

“New muon g-2 puzzle”

$$(a_\mu^{\text{HVP}})_{\text{EXP}} = a_\mu^{\text{EXP}} - a_\mu^{\text{SM, rest}}$$

$$(a_\mu^{\text{HVP}})_{e^+e^-}^{\text{WP20}} = 6931(40) \times 10^{-11}$$

$$(a_\mu^{\text{HVP}})_{\text{BMW}} = 7075(55) \times 10^{-11}$$



“new puzzle”: if BMW is correct, the “old” g-2 discrepancy (4.2σ) would be basically gone

→ however, this brings in a new tension with e^+e^- data (2.2σ)

Here, NP in $\sigma_{\text{had}}(e^+e^- \rightarrow \text{hadrons})$ such that [LDL, Masiero, Paradisi, Passera 2112.08312]

1. $(a_\mu^{\text{HVP}})_{e^+e^-}^{\text{WP20}} \approx (a_\mu^{\text{HVP}})_{\text{EXP}}$
2. the approximate agreement between BMW and EXP is not spoiled
3. w/o a direct contribution a_μ^{NP} (i.e. NP not in muons)

Muon $g-2 \Leftrightarrow \Delta\alpha$ connection

- Can Δa_μ be due to a **missing contribution** in σ_{had} ?

[Marciano, Passera, Sirlin 2008 & 2010;
Keshavarzi, Marciano, Passera, Sirlin 2020.
See also Crivellin, Hoferichter, Manzari, Montull 2020;
Malaescu, Schott 2020;
Colangelo, Hoferichter, Stoffer 2020]

→ a upward shift of σ_{had} induces an increase of $\Delta\alpha_{\text{had}}^{(5)}(M_Z)$

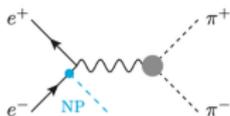
$$\alpha(M_Z) = \frac{\alpha}{1 - \Delta\alpha_{\text{lep}}(M_Z) - \Delta\alpha_{\text{had}}^{(5)}(M_Z) - \Delta\alpha_{\text{top}}(M_Z)}$$

- disfavoured by the EW fit (at about 2σ), if the shift happens at $\sqrt{s} \gtrsim 1 \text{ GeV}$

[Keshavarzi, Marciano, Passera, Sirlin 2020]

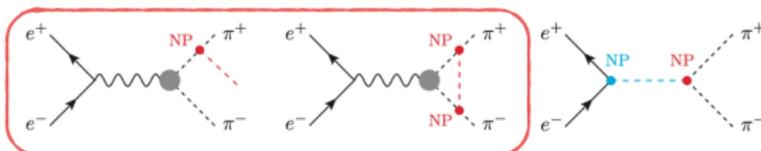
→ selects light NP inducing a sub-GeV modification of σ_{had}

- Light new physics inducing a sub-GeV modification of σ_{had} is the only possibility



- NP coupled only to **electrons** \rightarrow severe bounds

[See however
Darmé, Grilli di Cortona, Nardi 21/2.09/39
NP in Bhabha scattering? \rightarrow backup slides]

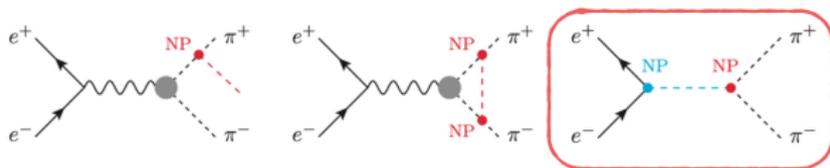


- NP coupled only to **hadrons**

FSR effects due to NP should be included into $\sigma_{\text{had}}(s)$, not easy to be accounted for...
(depend on exp. cuts and mass of NP)

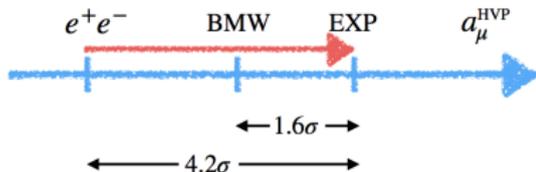
\rightarrow however, we know that in the QED case

$$(a_{\mu}^{\text{HVP}})_{e^+e^-}^{\text{FSR}} \approx 50 \times 10^{-11} \quad \longleftrightarrow \quad |(a_{\mu}^{\text{HVP}})_{\text{BMW}} - (a_{\mu}^{\text{HVP}})_{e^+e^-}^{\text{WP20}}| \approx 150 \times 10^{-11}$$



3. NP coupled both to hadrons and electrons

$$(a_\mu^{\text{HVP}})_{e^+e^-} = \frac{\alpha}{\pi^2} \int_{m_{\pi^0}^2}^{\infty} \frac{ds}{s} K(s) \text{Im} \Pi_{\text{had}}(s) = \frac{1}{4\pi^3} \int_{m_{\pi^0}^2}^{\infty} ds K(s) \sigma_{\text{had}}(s) \quad \sigma_{\text{had}} = \sigma_{\text{had}}^{\text{SM}} + \Delta\sigma_{\text{had}}^{\text{NP}}$$



$$\sigma_{\text{had}} - \Delta\sigma_{\text{had}}^{\text{NP}}$$

should be "subtracted" by NP,
since NP does not contribute to HVP at the LO
(but it contributes at the LO to the x-section)

\rightarrow a positive sift on $(a_\mu^{\text{HVP}})_{e^+e^-}$ requires $\Delta\sigma_{\text{had}}^{\text{NP}} < 0$ (negative interference)

A new light Z' vector boson

- Requirements:

1. σ_{had} modified at $\sqrt{s} \lesssim 1 \text{ GeV}$ (mostly $\pi^+\pi^-$ channel)

sub-GeV mediator

2. A sizeable negative interference with the SM

couples to u, d quarks (and electrons)

tree-level mediator

couples via a vector current

→ a light spin-1 mediator with vector couplings to first generation SM fermions

$$\mathcal{L}_{Z'} \supset (g_V^e \bar{e}\gamma^\mu e + g_V^q \bar{q}\gamma^\mu q) Z'_\mu \quad q = u, d \quad m_{Z'} \lesssim 1 \text{ GeV}$$

- It can be shown that (neglecting iso-spin breaking corrections due to NP)

$$\frac{\sigma_{\pi\pi}^{\text{SM+NP}}}{\sigma_{\pi\pi}^{\text{SM}}} = \left| 1 + \frac{g_V^e (g_V^u - g_V^d)}{e^2} \frac{s}{s - m_{Z'}^2 + im_{Z'}\Gamma_{Z'}} \right|^2$$

A new light Z' vector boson

1. Semi-leptonic processes

$e^+e^- \rightarrow q\bar{q}$ has been measured with per-cent accuracy at LEP-II

$$\frac{\sigma_{qq}^{\text{SM+NP}}}{\sigma_{qq}^{\text{SM}}} \approx 1 + 2 \frac{g_V^e g_V^q}{e^2 Q_q} \quad \longrightarrow \quad |g_V^e g_V^q| \lesssim 4.6 \cdot 10^{-4} |Q_q| \quad (\epsilon \lesssim 3.3 \cdot 10^{-3})$$

2. Leptonic processes

- for $m_{Z'} \lesssim 0.3$ GeV ($Z' \rightarrow e^+e^-$ is the main decay mode)

$$e^+e^- \rightarrow \gamma Z' \text{ @ BaBar} \quad \longrightarrow \quad g_V^e \lesssim 2 \cdot 10^{-4}$$

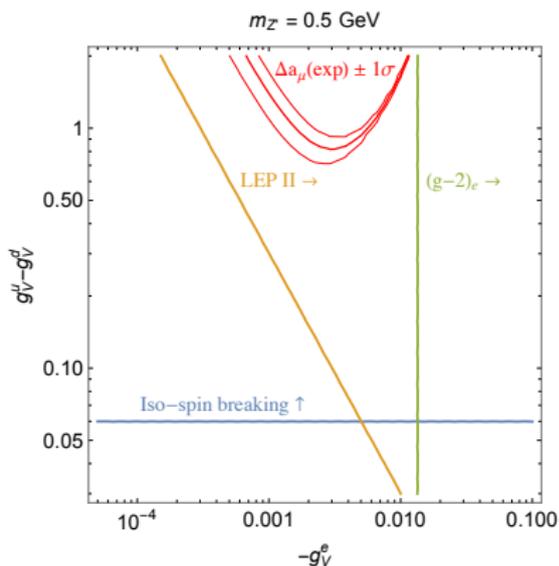
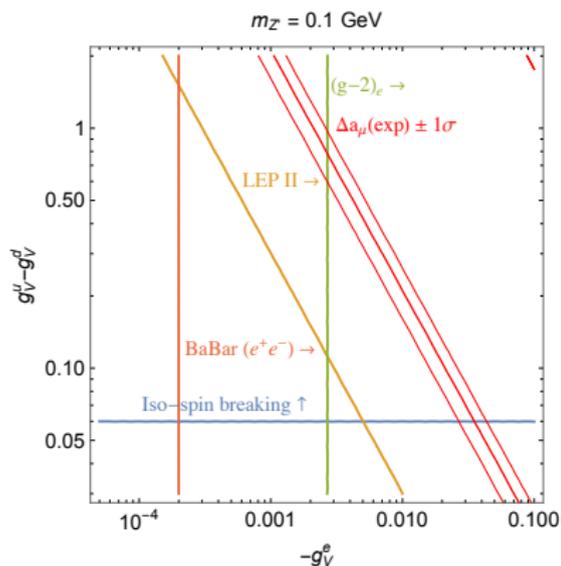
3. Iso-spin breaking observables

charged vs. neutral pion mass² difference $\Delta m^2 = m_{\pi^+}^2 - m_{\pi^0}^2$

$$(\Delta m^2)_{Z'} \sim \frac{(g_V^u - g_V^d)^2}{(4\pi)^2} \Lambda_\chi^2 \quad (\Lambda_\chi \approx 1 \text{ GeV})$$

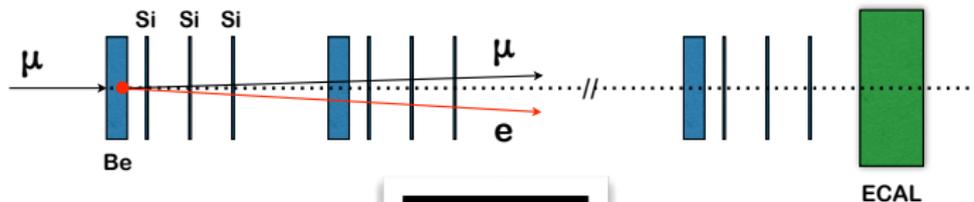
$$\longrightarrow \quad |g_V^u - g_V^d| \lesssim 0.06 \quad [\text{Rescaling lattice QCD calculation of Frezzotti et al 2112.01066}]$$

A new light Z' vector boson



Independent exp. bounds prevent to solve the “new muon $g-2$ puzzle”!

- $\Delta\alpha_{\text{had}}(t)$ can be measured via the **elastic scattering** $\mu e \rightarrow \mu e$.
- We propose to scatter a 150 GeV muon beam, available at CERN's North Area, on a fixed electron target (Beryllium). Modular apparatus: each station has one layer of Beryllium (target) followed by several thin Silicon strip detectors.

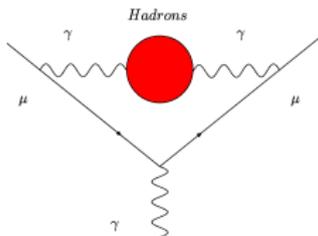


Abbiendi, Carloni Calame, Marconi, Matteuzzi, Montagna,
Nicosini, MP, Piccinini, Tenchini, Trentadue, Venanzoni
EPJC 2017 - arXiv:1609.08987

[Courtesy by M. Passera]

- Letter of Intent submitted to CERN SPSC in 2019 [see Pocanik's talk]

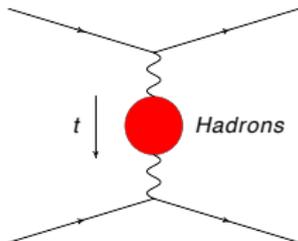
- The leading hadronic contribution a_{μ}^{HLO} computed via the **timelike** formula:



$$a_{\mu}^{\text{HLO}} = \frac{1}{4\pi^3} \int_{4m_{\pi}^2}^{\infty} ds K(s) \sigma_{\text{had}}^0(s)$$

$$K(s) = \int_0^1 dx \frac{x^2(1-x)}{x^2 + (1-x)(s/m_{\mu}^2)}$$

- Alternatively, simply exchanging the x and s integrations:



$$a_{\mu}^{\text{HLO}} = \frac{\alpha}{\pi} \int_0^1 dx (1-x) \Delta\alpha_{\text{had}}[t(x)]$$

$$t(x) = \frac{x^2 m_{\mu}^2}{x-1} < 0$$

Lautrup, Peterman, de Rafael, 1972

$\Delta\alpha_{\text{had}}(t)$ is the hadronic contribution to the running of α in the **spacelike** region: a_{μ}^{HLO} can be extracted from **scattering data!**

- The extraction of $\Delta\alpha_{\text{had}}$ is not contaminated by NP! [Masiero, PP, Passera, 2020]

- The muon $g - 2$ represents the most longstanding hint of New Physics now, thanks to the E989 experiment at FNAL, growing to 4.2σ .
- LQCD results by the BMWc weaken the muon $g - 2$ discrepancy to 1.6σ but they are in tension with the EW-fit and $e^+e^- \rightarrow \text{hadrons}$ experimental data:
 - ▶ The MUonE experiment can provide an independent measure of $\Delta\alpha_{\text{had}}$ which is not contaminated by new physics effects.
- Both heavy New Physics ($\nu \lesssim \Lambda \lesssim 100 \text{ TeV}$) and light New Physics ($\Lambda \lesssim 1 \text{ GeV}$) scenarios have the potential to account for the muon $g-2$ anomaly.
 - ▶ Different scenarios can be disentangled by dedicated searches at running or future experiments such as Belle II and a high-energy Muon Collider .
- If the muon $g - 2$ anomaly will survive, we expect relevant enhancements in leptonic EDMs (especially in the muon EDM) and LFV physics.
- Testing New Physics effects in the electron $g - 2$ at the 10^{-13} is not too far! This will bring a_e to play a pivotal role in probing New Physics in the leptonic sector.

Message: an exciting Physics program is in progress at the Intensity Frontier!