



IFIC (CSIC-UV), València

**Resonance Chiral Theory and its predictions,
case of 2π and 3π final states in TAUOLA MC:
status report**

O. Shekhovtsova

*work done together with
Z. Was and P. Roig*

- Theoretical calculation. Resonance Chiral Theory.
- TAUOLA, implementation RCHT results for $3\pi/2\pi$ modes.

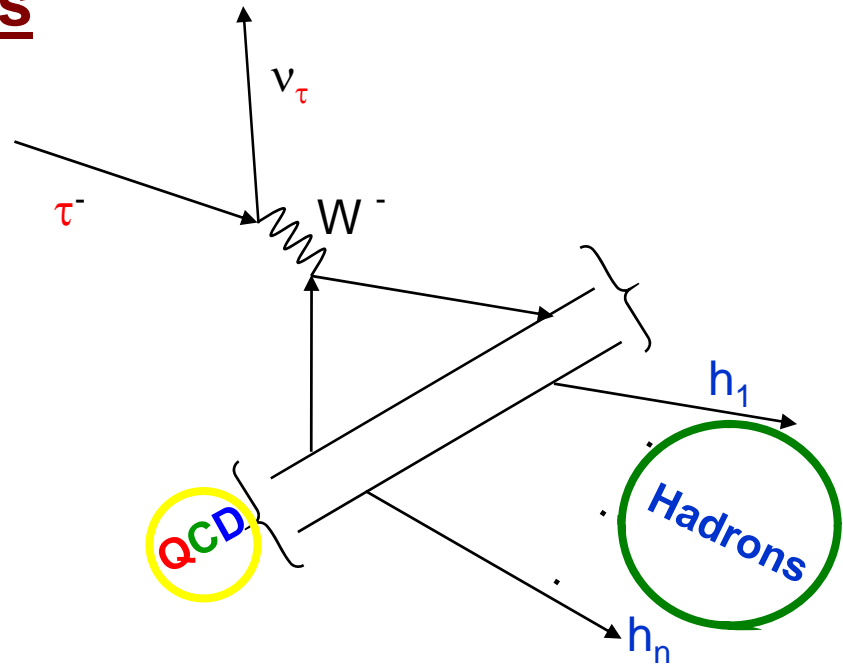
Tests numerical stability

- Comparison with BABAR for 3π mode
- Conclusion (inspiration for future work)

Basics of hadronic tau decays

$$M = \frac{G_F}{\sqrt{2}} V_{CKM} \bar{u}(v_\tau) \gamma^\mu (1 - \gamma_5) u(\tau) T_\mu$$

$$d\Gamma = \frac{G_F^2}{4M_\tau^2} |V_{CKM}|^2 d\Phi^{(n+1)} L_{\mu\nu} T^\mu T^{\nu*}$$



$$T_\mu = \langle \text{Hadrons} | (\mathbf{V} - \mathbf{A})_\mu e^{iS_{\text{QCD}}} | 0 \rangle = \sum_i (\text{Lorentz Structure})^i F_i(Q^2, s_j)$$

3 pions: 3 Lorentz independent structure

2 pions: 1 Lorentz independent structure (vector)

$$F_i(Q^2, s_j) \longrightarrow R_{\chi T}$$

Resonance Chiral Theory (Chiral Theory with the explicit inclusion of ρ and a_1 mesons)

*G.Ecker et al., Nucl. Phys B321,
311 (1989)*

1. To enlarge the energy region for ChPT (≤ 1.5 GeV) the resonance fields ($V_{\mu\nu}, A_{\mu\nu}$) should be added in explicit way
2. $E \ll m_\rho$ the singularity associated with the pole of a resonance propagator is replaced by momentum expansion
3. The main contribution to the coupling constants (at least up to p^4) comes from the meson resonance exchange

The antisymmetric tensor field for the vector (axial-vector) meson

About the parameter values: experimental data or theoretical relations?

Parameters: f_π, F_V, G_V, F_A

for example

$$\Gamma(\rho \rightarrow e^+e^-) = \frac{4\pi\alpha^2 F_V^2}{3m_\rho}$$

Theory: $F_V G_V = f_\pi^2, F_V^2 - F_A^2 = f_\pi^2, F_V^2 M_V^2 = F_A^2 M_A^2$

$2F_V G_V \neq F_V^2$ M_V, M_A – mass of vector (axial-) nonet

chiral and large N_c limit:

$M_\rho \sim 770$ MeV, $M_A \sim 998$ MeV

$$\Gamma(\rho \rightarrow \pi\pi) = \frac{G_V^2 m_\rho^3}{48\pi f_\pi^4} \left(1 - \frac{4m_\pi^2}{m_\rho^2}\right)^{3/2}$$

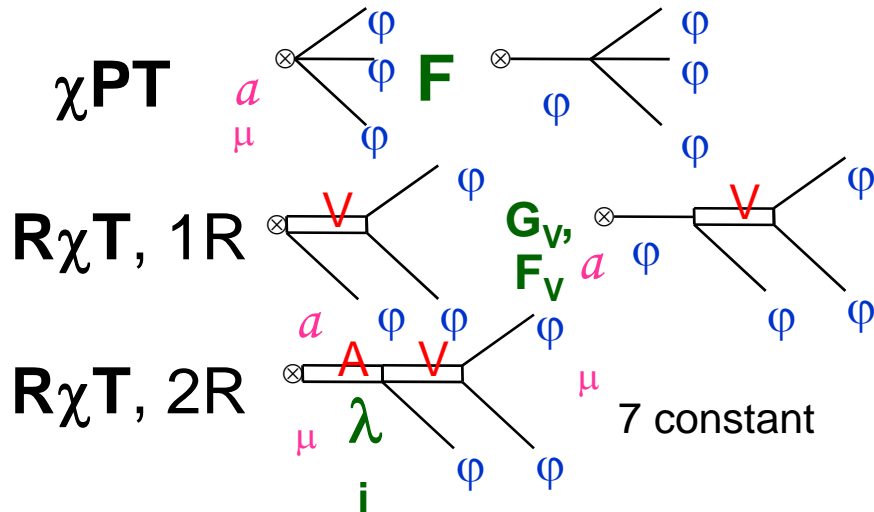
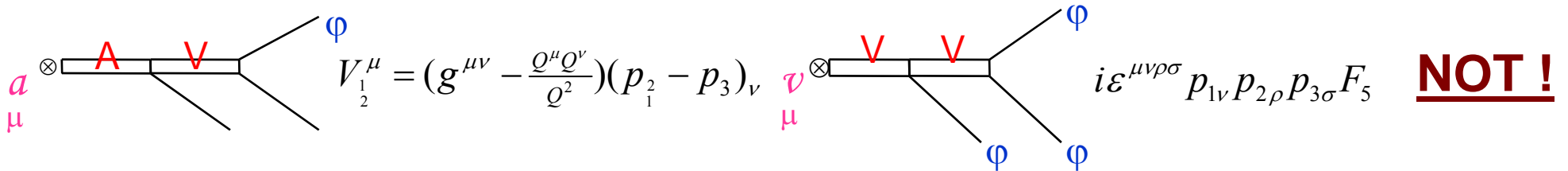
$$\Gamma(a_1 \rightarrow \pi\gamma) = \frac{\alpha F_a^2 m_a}{24 f_\pi^2} \left(1 - \frac{m_\pi^2}{m_a^2}\right)^3$$

Three pion mode: $\tau^- \rightarrow (3\pi)^- \nu_\tau$

Hadronic tensor: $\tau \rightarrow h_1(p_1)h_2(p_2)h_3(p_3)\nu_\tau \quad T^\mu = V_1^\mu F_1 + V_2^\mu F_2 + Q F_4 + i\varepsilon^{\mu\nu\rho\sigma} p_{1\nu} p_{2\rho} p_{3\sigma} F_5$

$$V_1^\mu = (g^{\mu\nu} - \frac{Q^\mu Q^\nu}{Q^2})(p_2 - p_3)_\nu, \quad V_2^\mu = (g^{\mu\nu} - \frac{Q^\mu Q^\nu}{Q^2})(p_1 - p_3)_\nu,$$

$$Q^\mu = p_1^\mu + p_2^\mu + p_3^\mu$$



$$F_{1,2} = F_{1,2}^\chi + F_{1,2}^R + F_{1,2}^{RR},$$

$$F_4 = F_p \sim m_\pi^2 / Q^2$$

D. Gomez-Dumm, A. Pich, J. Portoles, P. Roig
(arXiv:0911.4436)

$$V = \rho, \rho' \quad A = a_1$$

How to include ρ' ?

$$\rho : F_V, G_V \longrightarrow \rho' : F_{V1}, G_{V1}$$

Too many constants

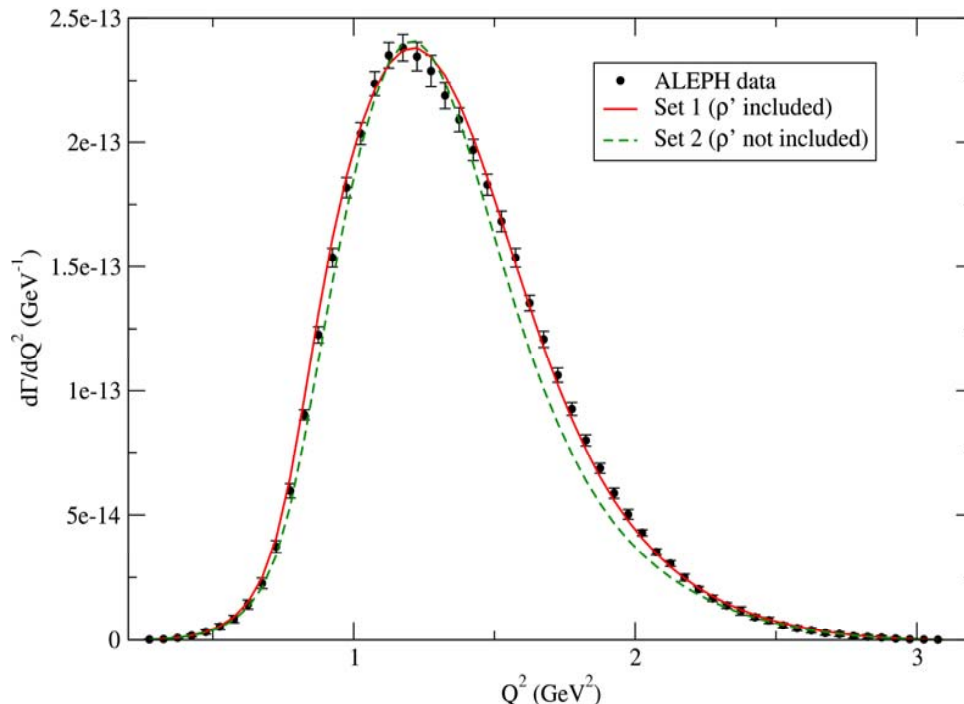
Following K&S:
$$\frac{1}{M_\rho^2 - q^2 - iM_\rho\Gamma_\rho(q^2)} \longrightarrow \frac{1}{1 + \beta_{\rho'}} \left[\frac{1}{M_\rho^2 - q^2 - iM_\rho\Gamma_\rho(q^2)} + \frac{\beta_{\rho'}}{M_{\rho'}^2 - q^2 - iM_{\rho'}\Gamma_{\rho'}(q^2)} \right]$$

$$\Gamma_\rho(s) = \frac{M_\rho s}{96 \pi F^2} \left[\sigma_\pi^3 \theta(s - 4m_\pi^2) + \frac{1}{2} \sigma_K^3 \theta(s - 4m_K^2) \right]$$

$$\sigma_P = \sqrt{1 - 4m_P^2/s}$$

$$\Gamma_{\rho'}(q^2) = \Gamma_{\rho'}(M_{\rho'}^2) \frac{M_{\rho'}}{\sqrt{q^2}} \left(\frac{p(q^2)}{p(M_{\rho'}^2)} \right)^3 \theta(q^2 - 4m_\pi^2),$$

$$p(x) = \frac{1}{2} \sqrt{x - 4m_\pi^2}.$$



a1 width:

$$\Gamma_{a_1}^{\pi, K}(Q^2) = \frac{-1/n!}{192(2\pi)^3 F_A^2 M_{a_1}} \left(\frac{M_{a_1}^2}{Q^2} - 1 \right)^2 \int ds dt T_{1, \nu}^{\pi, K \mu} T_{1, \mu}^{\pi, K \nu}$$

$$\Gamma_{a_1}(Q^2) = \Gamma_{a_1}^\pi(Q^2) \theta(Q^2 - 9m_\pi^2) + \Gamma_{a_1}^K(Q^2) \theta(Q^2 - (2m_K + m_\pi)^2)$$

$$T_{1+}^\mu = V_1^\mu F_1 + V_2^\mu F_2.$$

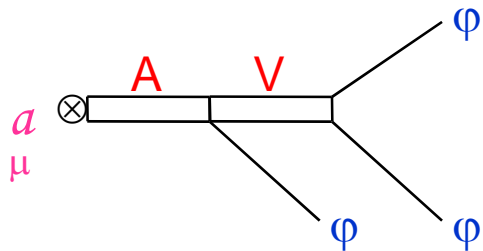
$$F_V = 0.180 \text{ GeV} \quad , \quad F_A = 0.149 \text{ GeV} \quad , \quad \beta_{\rho'} = -0.25 \quad , \\ M_V = 0.775 \text{ GeV} \quad , \quad M_{K^*} = 0.8953 \text{ GeV} \quad , \quad M_{a_1} = 1.120 \text{ GeV}$$

and

$$F_V = 0.206 \text{ GeV} \quad , \quad F_A = 0.145 \text{ GeV} \quad , \quad \beta_{\rho'} = 0 \quad , \\ M_V = 0.775 \text{ GeV} \quad , \quad M_{K^*} = 0.8953 \text{ GeV} \quad , \quad M_{a_1} = 1.115 \text{ GeV}$$

TAUOLA

Published version: J.H.Kuhn, A. Santamaría, Z. Phys. C48 (1990) 445



$$\underline{V = \rho, \rho'} \quad \underline{A = a_1}$$

3 pion:

- Vector Meson Dominance $BW_R(x^2) = \frac{M_R^2}{M_R^2 - x^2 - i\sqrt{x^2} \Gamma_R(x^2)}$
- Asymptotic behaviour ruled by **QCD**
- χ PT $\mathcal{O}(p^2)$ **YES** χ PT $\mathcal{O}(p^4)$ **NO**

$$T_{\pm\mu}^{\chi PT} \Big|_{1+} = \mp \frac{2\sqrt{2}}{3F} \left[\left(1 + \frac{3s}{2M_V^2}\right) V_{1\mu} + \left(1 + \frac{3t}{2M_V^2}\right) V_{2\mu} \right] + \text{chiral loops} + \mathcal{O}(p^6)$$

$$T_{\pm\mu}^{(KS)} \xrightarrow{s,t \ll M_V^2} \mp \frac{2\sqrt{2}}{3F} \left[\left(1 + \frac{s}{M_V^2}\right) V_{1\mu} + \left(1 + \frac{t}{M_V^2}\right) V_{2\mu} \right]$$

*hep-ph/0213283, G. Gómez-Dumm,
A. Pich, J.Portolés*

RχT 3pion in TAUOLA

(only F_1, F_2 for the moment)

$$m_{\pi^+} = m_{\pi^0}$$

Check of precision integration

comparison with numerical integration of analytical formulae
(Gauss method integration)

- $F_1 = 1, F_2 = 0$ *no singularity phase space*
MC: $\sigma = (2.7402 \pm 0.0004) \text{ E-17 GeV}$ Analyt: $\sigma = (2.7408 \pm 0.00005) \text{ E-17 GeV}$
- $F_1 = \text{true}, F_2 = 0$
MC: $\sigma = (1.8720 \pm 0.0004) \text{ E-13 GeV}$ Analyt: $\sigma = (1.8720 \pm 0.0004) \text{ E-13 GeV}$
- $F_1 = \text{true}, F_2 = \text{true}$
MC: $\sigma = (4.2013 \pm 0.0006) \text{ E-13 GeV}$ Analyt: $\sigma = (4.2014 \pm 0.0008) \text{ E-13 GeV}$

$m_{\pi^+} \neq m_{\pi^0}$ for phase space $F_1 = \text{true}, F_2 = \text{true}$

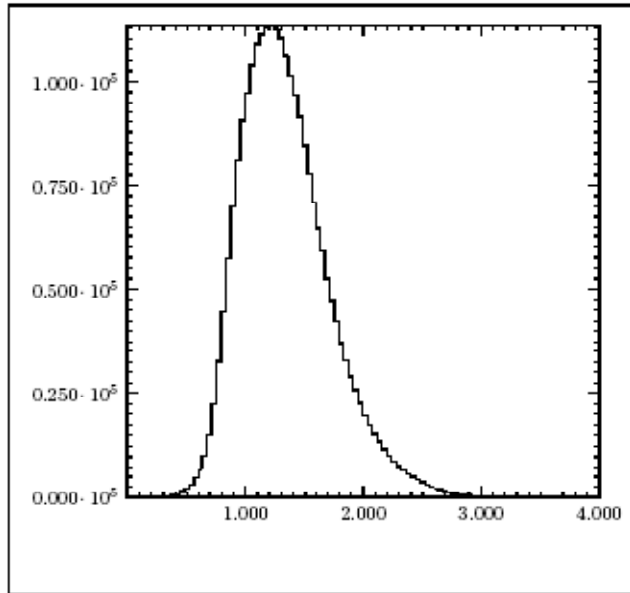
MC: $\sigma = (4.2046 \pm 0.0006) \text{ E-13 GeV}$

Elementary technical test:

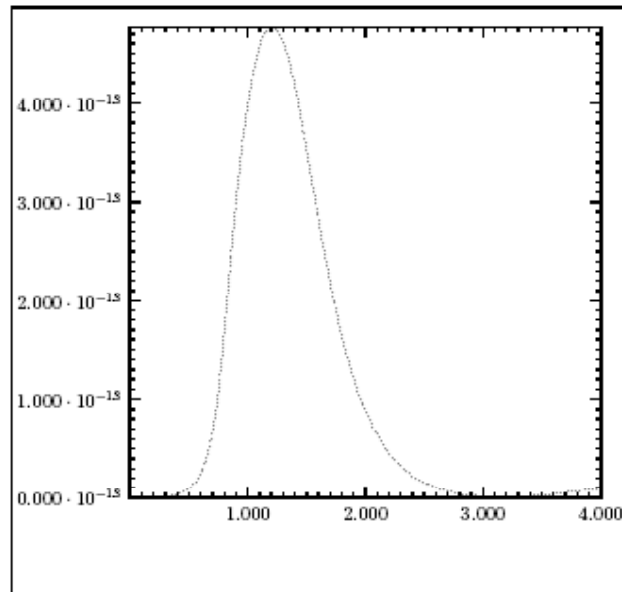
Comparison with linearly interpolated spectrum ($F_1 = true$, $F_2 = true$)

Difference is ~10% (first 5 points from 1000) at the beginning of the spectrum and less than 0.5%

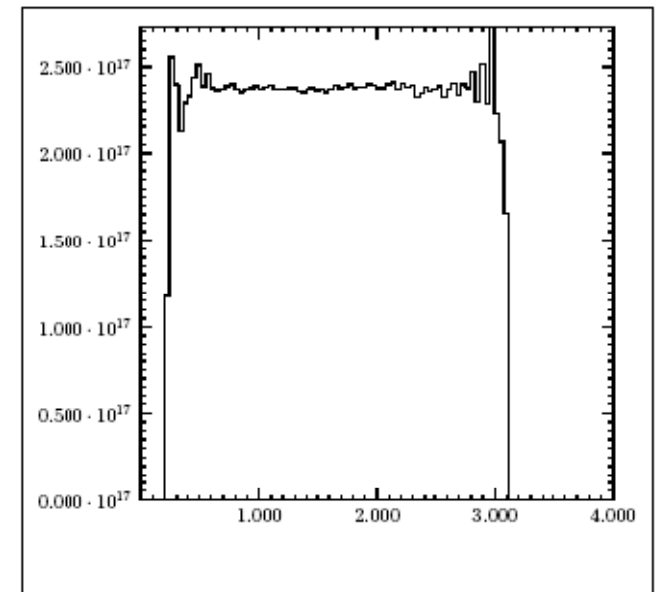
Monte CARLO for QQ



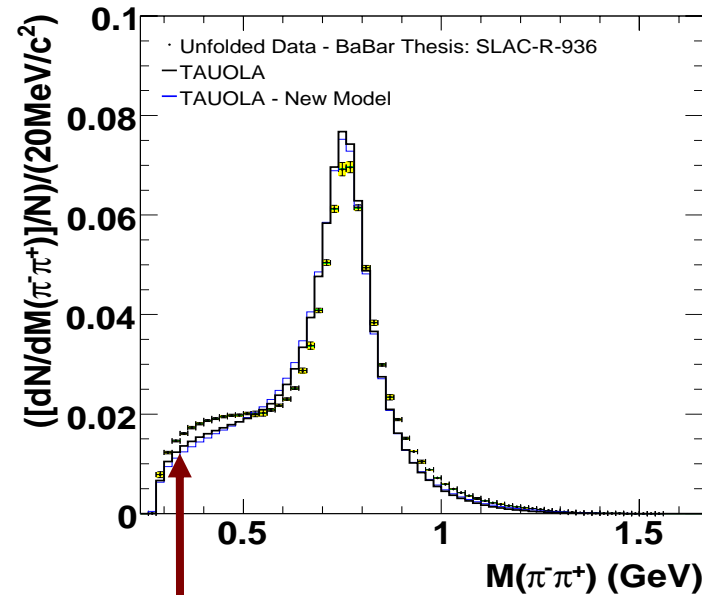
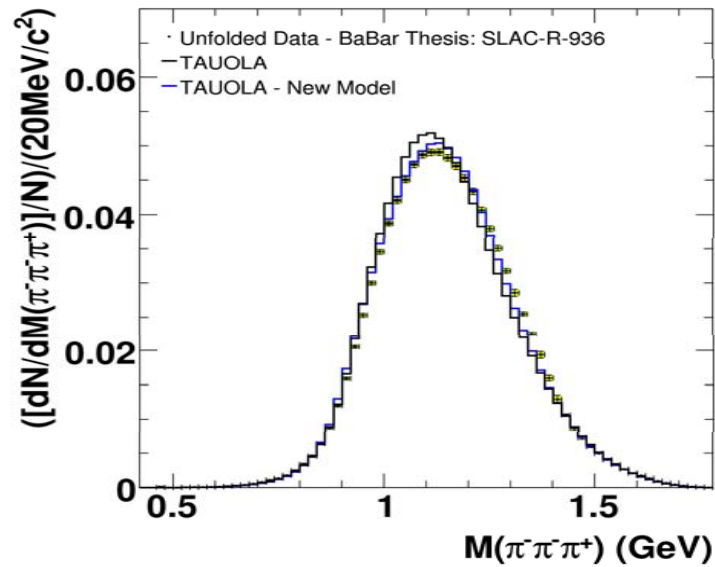
analytical function for QQ



ratio with finest WIDOLGA3PI



Comparison with BABAR data: *Ian M. Nugent, (Victoria U.) . SLAC-R-936, Dec16, 2009. Ph.D. Thesis (Advisor: Dr. J. Michael Roney).*



F_4 will help (scalar contribution)?

Other th. or exp. problems?

The status: *without real fit, better result for q^2 spectrum should not be expected*

MORE WORK NEEDED

angular distributions should help to separate scalar from vector

Two pion case: $\tau \rightarrow \pi\pi V$

Hadronic tensor $T^\mu = (p_1 - p_2)^\mu F_\pi(q^2)$

R χ T model for pion form factor

$$F(q^2) = \left(\frac{m_v^2 + (\gamma e^{-i\phi_1} + \delta e^{-i\phi_2})s}{D_\rho(q^2)} - \frac{\gamma s e^{-i\phi_1}}{D_{\rho'}(q^2)} - \frac{\delta s e^{-i\phi_2}}{D_{\rho''}(q^2)} \right) e^{-((\pi\pi + KK) \text{ loops})}$$

Features:

correct up to χ PT $\mathcal{O}(p^4)$ behavior

correct high energy behaviour

Parameters from BELLE spectrum

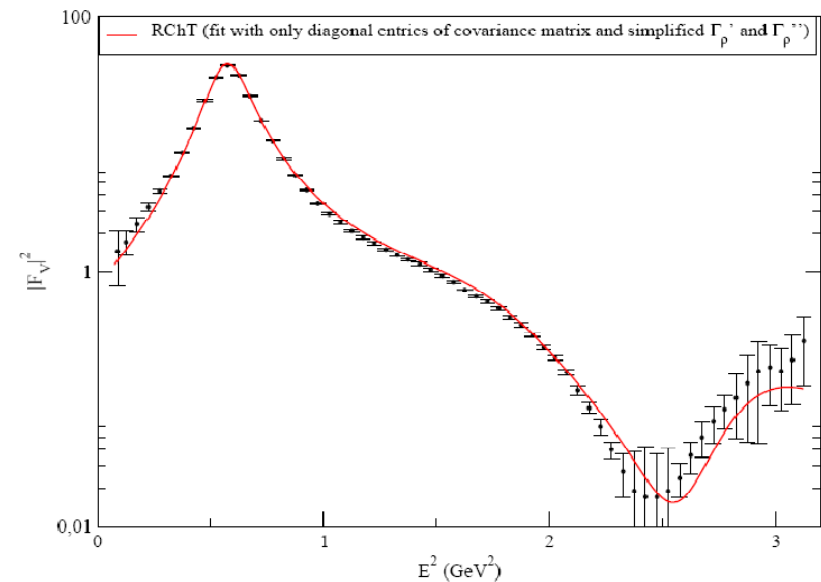
Phys. Rev. D 78 (2008) 072006

Checks:

comparison with linearly interpolated spectrum

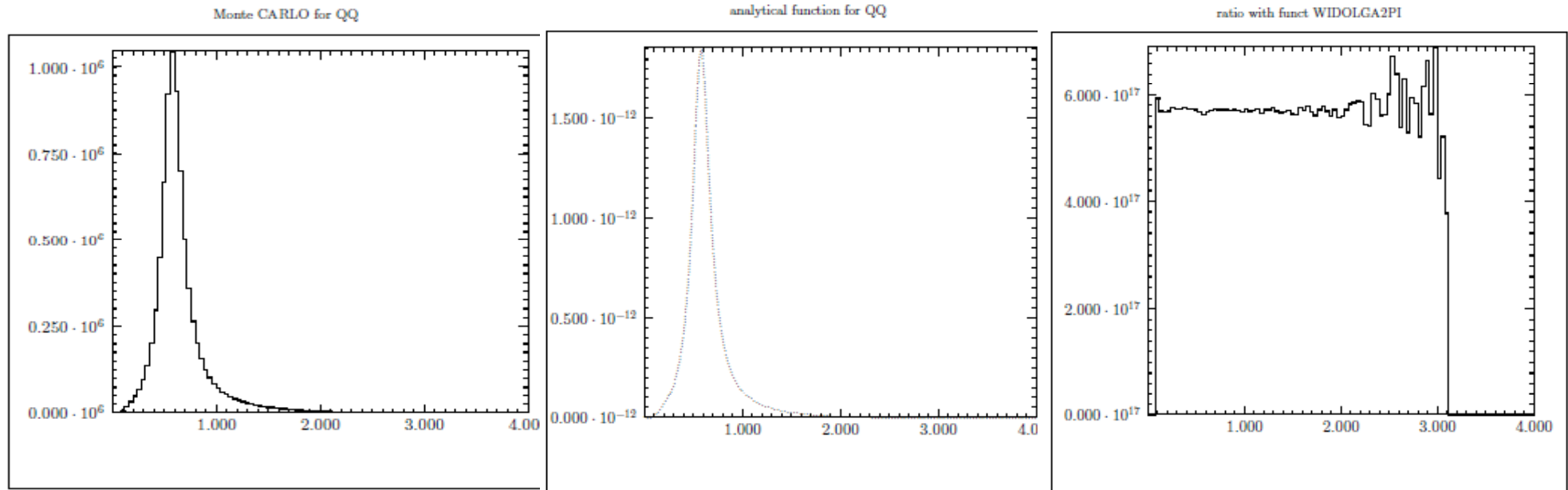
comparison MC and analytical results

after Gauss integration



Technical test:

Monte Carlo results (8e6 events)



MC: $\sigma = (5.617 \pm 0.0005) \text{ E-13 GeV}$ Analyt: $\sigma = (5.616 \pm 0.0005) \text{ E-13 GeV}$

CONCLUSIONS

- Theoretical results in **$R\chi T$** for 2 and 3 pion modes are put in TAUOLA.
- All work is performed with support of code manager SVN.
- A starting version for code development TAUOLA is Belle/BaBar version.
That guarantees that once work is completed can be installed into collaboration software.
- Once BaBar and/or Belle will confirm the upgrades to be valuable and fits to the data will be improved the SVN reference will make installation of such upgrade to C++ and FORTRAN applications of TAUOLA at LHC rather straightforward too.

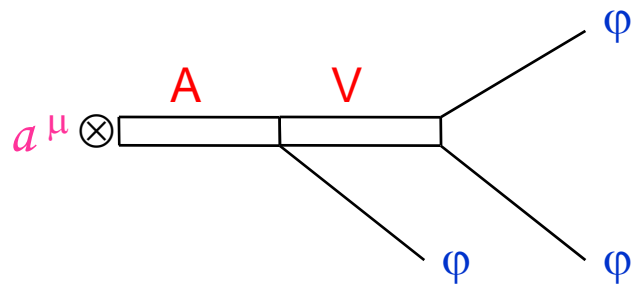
WORK IN PROGRESS

BACKUP

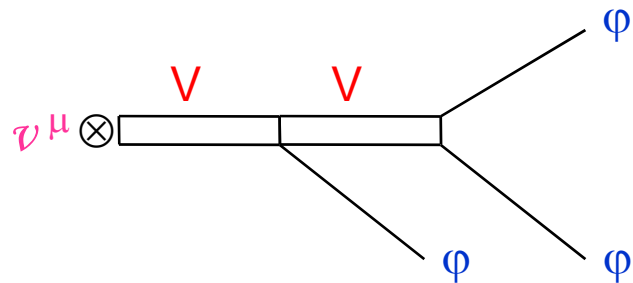
SLIDES

K-S-like works & HADRONIZATION IN TAUOLA

(Finkemeier, Mirkes '95,'96)
 (Finkemeier, Kühn, Mirkes '96)



$$\rightarrow V_{1\mu} = \left(g_{\mu\nu} - \frac{Q_\mu Q_\nu}{Q^2} \right) (p_2 - p_1)^\nu$$



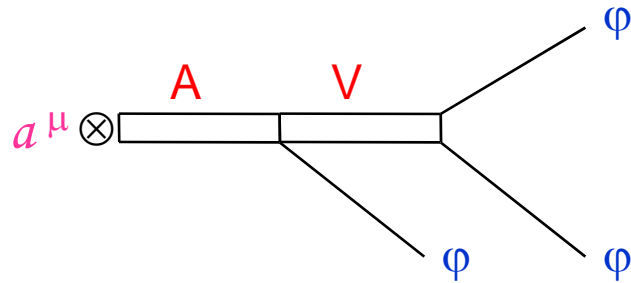
$$\rightarrow i\varepsilon_{\mu\nu\rho\sigma} p_1^\nu p_2^\rho p_3^\sigma$$

$\chi, 1R \text{ \& } 2R$
 obtained from:

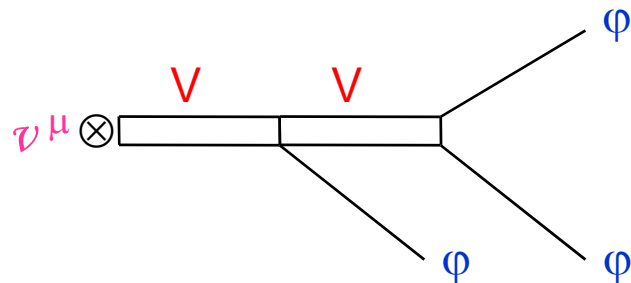
$$\frac{M_{R1}^2}{M_{R1}^2 - x^2 - i\sqrt{x^2} \Gamma_{R1}(x^2)} \frac{M_{R2}^2}{M_{R2}^2 - y^2 - i\sqrt{y^2} \Gamma_{R2}(y^2)}$$

K-S-like works & HADRONIZATION IN TAUOLA

(Finkemeier, Mirkes '95,'96)
 (Finkemeier, Kühn, Mirkes '96)



Some allowed $\rho^{0,-}$, K^{*0} contributions are lacking in the modes: $K^+K^-\pi^-$, $K^-K^0\pi^0$, $K^-\pi^-\pi^+$, $K^0\pi^0\pi^-$



$$\left\{ \begin{array}{l} M_{\rho'}^{V_\mu} \neq M_{\rho'}^{A_\mu} \\ \Gamma_{\rho'}^{V_\mu} \neq \Gamma_{\rho'}^{A_\mu} \end{array} \right.$$



3 Multiplets^{V_μ} ≠ 2 Multiplets^{A_μ}

χ PT: The low-energy

EFT of QCD

(Gasser & Leutwyler '84, '85)

$$\phi(x) = \begin{pmatrix} \frac{\pi^0}{\sqrt{2}} + \frac{\eta_8}{\sqrt{6}} & \pi^+ & K^+ \\ \pi^- & -\frac{\pi^0}{\sqrt{2}} + \frac{\eta_8}{\sqrt{6}} & K^0 \\ K^- & \bar{K}^0 & -\frac{2\eta_8}{\sqrt{6}} \end{pmatrix}$$

**Goldstone
Bosons**

$$SU(3)_L \otimes SU(3)_R \rightarrow SU(3)_V$$

$$u(x) = \exp\left(\frac{i\phi(x)}{\sqrt{2}F}\right), \quad u_\mu = i\left[u^\dagger(\partial_\mu - ir_\mu)u - u(\partial_\mu - il_\mu)u^\dagger\right]$$

$$\chi = 2B_0(s + ip), \quad \chi_\pm = u^\dagger \chi u^\dagger \pm u \chi u$$

$$f_\pm^{\mu\nu} = u F_L^{\mu\nu} u^\dagger \pm u^\dagger F_R^{\mu\nu} u$$

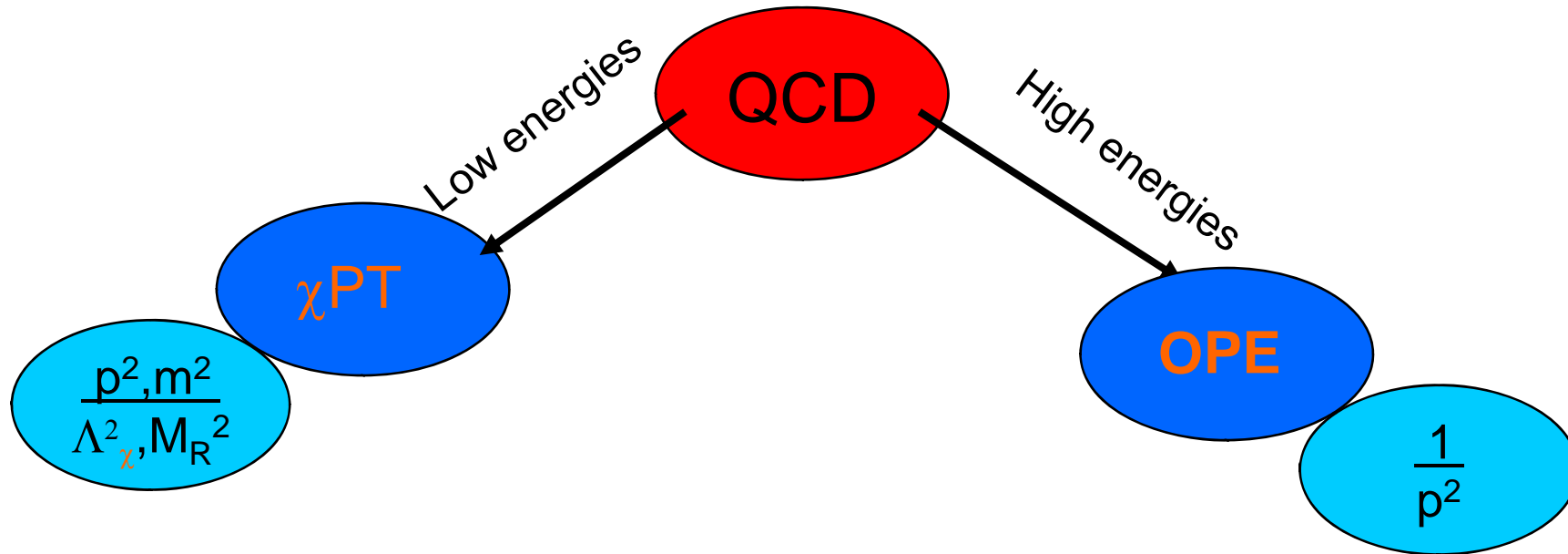
$$\mathcal{L}_\chi^{(2)} = \frac{F^2}{4} \langle u_\mu u^\mu + \chi_+ \rangle$$

$$X \rightarrow h(g, \Phi) X h(g, \Phi)^\dagger$$

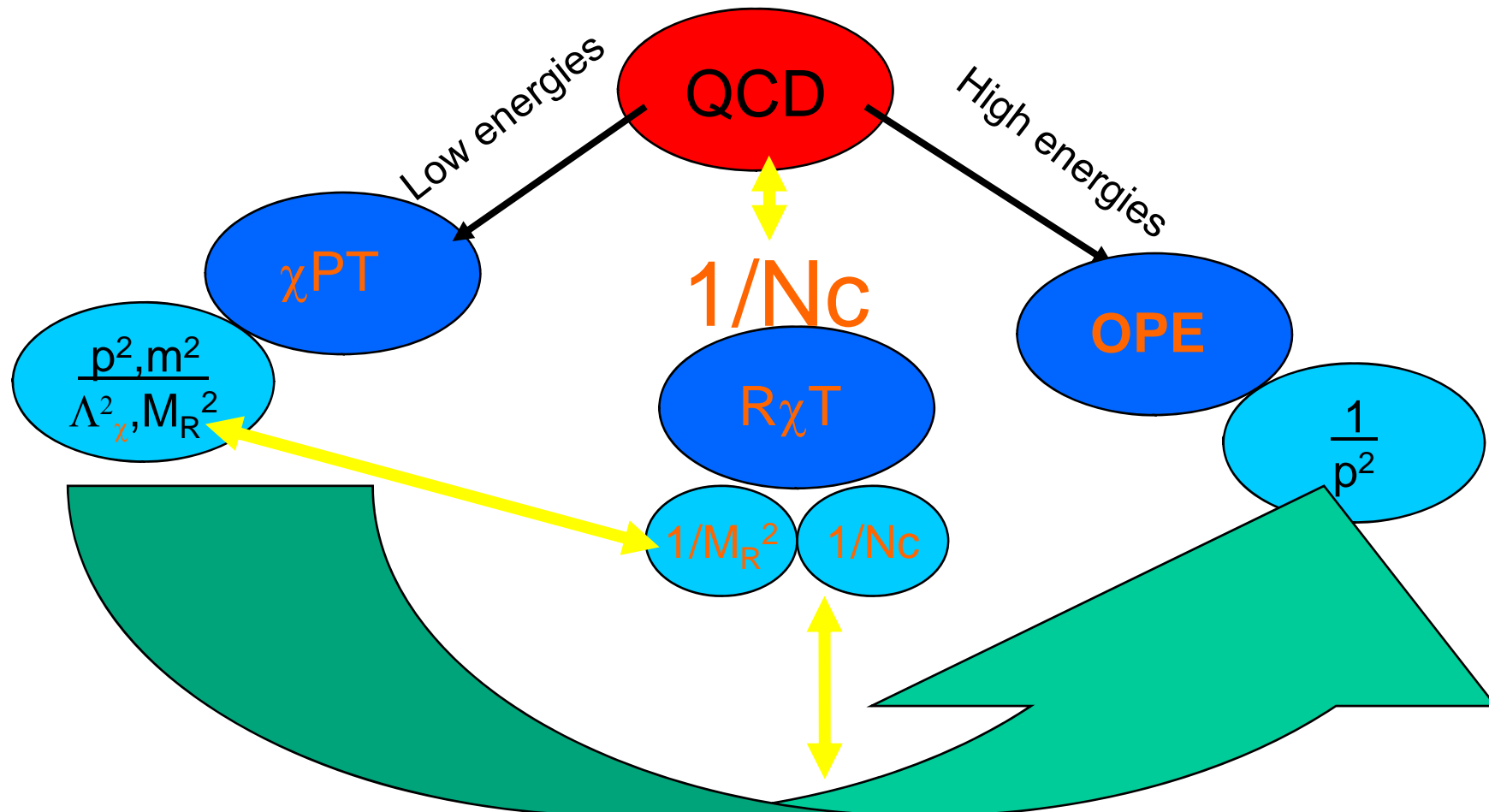
$$\mathcal{L}_\chi^{(4)} = L_1 \langle u_\mu u^\mu \rangle^2 + \dots + L_4 \langle u_\mu u^\mu \rangle \langle \chi_+ \rangle + \dots + L_7 \langle \chi_- \rangle^2 + \dots - iL_9 \langle f_+^{\mu\nu} u_\mu u_\nu \rangle + \dots$$

$$\mathcal{L}_\chi^{(4)} \text{ in the odd-intrinsic parity sector}$$

$R_\chi T$ matching to the OPE allows it to reproduce QCD high-energy behaviour:



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Resonances+ Goldstone Bosons

TOOLS : R_χT

(Ecker, Gasser, Pich, De Rafael '89)

(Ecker, Gasser, Leutwyler, Pich, De Rafael '89)

$$\mathcal{L}_{R\chi T}^{(P_I=+)} = \mathcal{L}_{\chi}^{(2)} + \mathcal{L}_{V,A}^{kin} + \mathcal{L}_V + \mathcal{L}_A + \mathcal{L}_{VAP};$$

$$\mathcal{L}_V = \frac{F_V}{2\sqrt{2}} \langle V_{\mu\nu} f_+^{\mu\nu} \rangle + \frac{iG_V}{\sqrt{2}} \langle V_{\mu\nu} u^\mu u^\nu \rangle$$

$$\mathcal{L}_A = \frac{F_A}{2\sqrt{2}} \langle A_{\mu\nu} f_-^{\mu\nu} \rangle$$

Antisymmetric tensor formalism

$$V_{\mu\nu}(x) = \begin{pmatrix} \frac{\rho^0}{\sqrt{2}} + \frac{\omega_8}{\sqrt{6}} & \rho^+ & K^{*+} \\ \rho^- & -\frac{\rho^0}{\sqrt{2}} + \frac{\omega_8}{\sqrt{6}} & K^{*0} \\ K^{*-} & \bar{K}^{*0} & -\frac{2\omega_8}{\sqrt{6}} \end{pmatrix}_{\mu\nu}$$

$$\mathcal{L}_{VAP} = \sum_{i=1}^5 \lambda_i O^i(V_{\mu\nu}, A^{\mu\nu}, \phi) = \lambda_1 \langle [V_{\mu\nu}, A^{\mu\nu}] \chi_- \rangle + \dots$$

(Gómez Dumm, Pich, Portolés '04)

$$\mathcal{L}_{R\chi T}^{(P_I=-)} = \mathcal{L}_{\chi(WZW)}^{(4)} + \mathcal{L}_{VJP} + \mathcal{L}_{VVP} + \mathcal{L}_{VPPP};$$

VMD

$$\mathcal{L}_{VJP} = \sum_{i=1}^7 \frac{c_i}{M_V} O^i(V_{\mu\nu}, j^\nu, \partial^\mu \phi) = \frac{c_5}{M_V} \varepsilon_{\mu\nu\rho\sigma} \langle \{ \nabla_\alpha V^{\mu\nu}, f_+^{\rho\alpha} \} u^\sigma \rangle + \dots$$

(Ruiz-Femenía, Pich, Portolés '03)

$$\mathcal{L}_{VVP} = \sum_{i=1}^5 d_i O^i(V_{\mu\nu}, V_{\rho\sigma}, \phi) = d_1 \varepsilon_{\mu\nu\rho\sigma} \langle \{ V^{\mu\nu}, V^{\rho\alpha} \} \nabla_\alpha u^\sigma \rangle + \dots$$

$$\mathcal{L}_{VPPP} = \sum_{i=1}^5 \frac{g_i}{M_V} O^i(V_{\mu\nu}, \phi) = \frac{g_4}{M_V} \varepsilon_{\mu\nu\alpha\beta} \langle \{ V^{\mu\nu}, u^\alpha u^\beta \} \chi_- \rangle + \dots$$

(Gómez Dumm, Pich, Portolés, R. to appear)

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HADRONIZATION IN 3 MESON CHANNELS
Pablo Roig (IFIC & INFN & TUM)

**Resonances+
Goldstone
Bosons**

TOOLS : R_χT

(Ecker, Gasser, Pich, De Rafael '89)

(Ecker, Gasser, Leutwyler, Pich, De Rafael '89) ,...

$$\mathcal{L}_{R\chi T}^{(P_I=+)} = \mathcal{L}_{\chi}^{(2)} + \mathcal{L}_{V,A}^{kin} + \mathcal{L}_V + \mathcal{L}_A + \mathcal{L}_{VAP};$$

$$\mathcal{L}_V = \frac{F_V}{2\sqrt{2}} \langle V_{\mu\nu} f_+^{\mu\nu} \rangle + \frac{iG_V}{\sqrt{2}} \langle V_{\mu\nu} u^\mu u^\nu \rangle$$

$$\mathcal{L}_A = \frac{F_A}{2\sqrt{2}} \langle A_{\mu\nu} f_-^{\mu\nu} \rangle$$

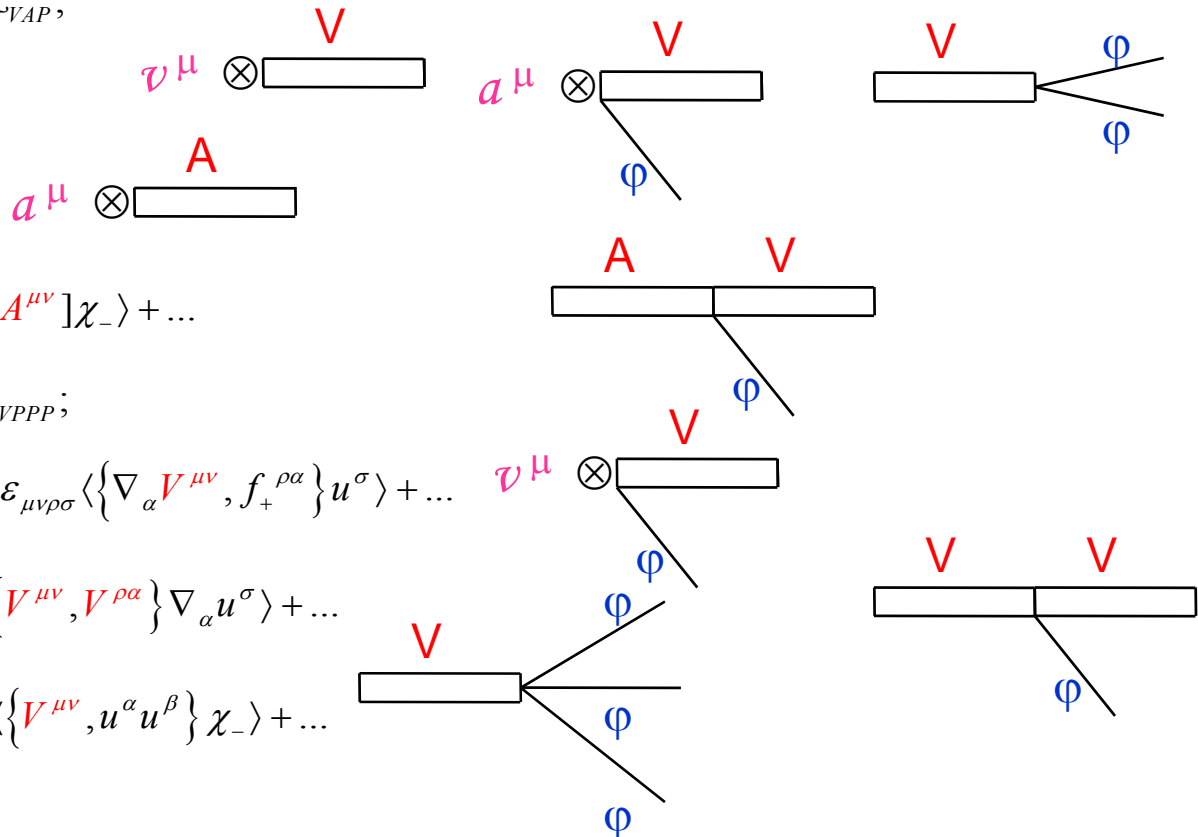
$$\mathcal{L}_{VAP} = \sum_{i=1}^5 \lambda_i O^i(V_{\mu\nu}, A^{\mu\nu}, \phi) = \lambda_1 \langle [V_{\mu\nu}, A^{\mu\nu}] \chi_- \rangle + \dots$$

$$\mathcal{L}_{R\chi T}^{(P_I=-)} = \mathcal{L}_{\chi(WZW)}^{(4)} + \mathcal{L}_{VJP} + \mathcal{L}_{VVP} + \mathcal{L}_{VPPP};$$

$$\mathcal{L}_{VJP} = \sum_{i=1}^7 \frac{c_i}{M_V} O^i(V_{\mu\nu}, j^\nu, \partial^\mu \phi) = \frac{c_5}{M_V} \epsilon_{\mu\nu\rho\sigma} \langle \{ \nabla_\alpha V^{\mu\nu}, f_+^{\rho\alpha} \} u^\sigma \rangle + \dots$$

$$\mathcal{L}_{VVP} = \sum_{i=1}^5 d_i O^i(V_{\mu\nu}, V_{\rho\sigma}, \phi) = d_1 \epsilon_{\mu\nu\rho\sigma} \langle \{ V^{\mu\nu}, V^{\rho\alpha} \} \nabla_\alpha u^\sigma \rangle + \dots$$

$$\mathcal{L}_{VPPP} = \sum_{i=1}^5 \frac{g_i}{M_V} O^i(V_{\mu\nu}, \phi) = \frac{g_4}{M_V} \epsilon_{\mu\nu\alpha\beta} \langle \{ V^{\mu\nu}, u^\alpha u^\beta \} \chi_- \rangle + \dots$$



CLEO analyses & ~~KS~~-like models

(Liu '03) (CLEO-III '04)

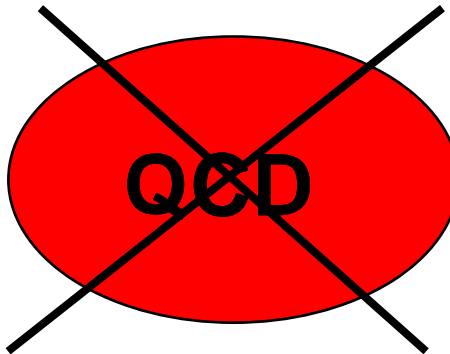
CLEO



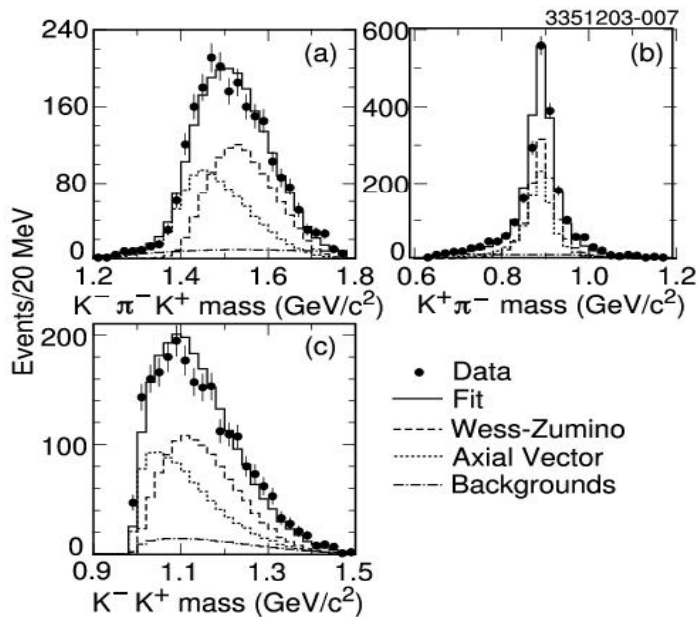
$$F_V = \underbrace{\frac{1}{2\sqrt{2}\pi^2 F^3}}_{L^{(4)}_{\chi,WZW}} \sqrt{R_B} \underbrace{\sum_i BW_i}_{\xrightarrow{Q^2 \rightarrow 0} 1}$$

$$BW_{V,A}(x = s_i, Q^2) = \frac{M_{V,A}^2}{M_{V,A}^2 - x - i\sqrt{x}\Gamma_{V,A}(x)}$$

1.80 ± 0.53

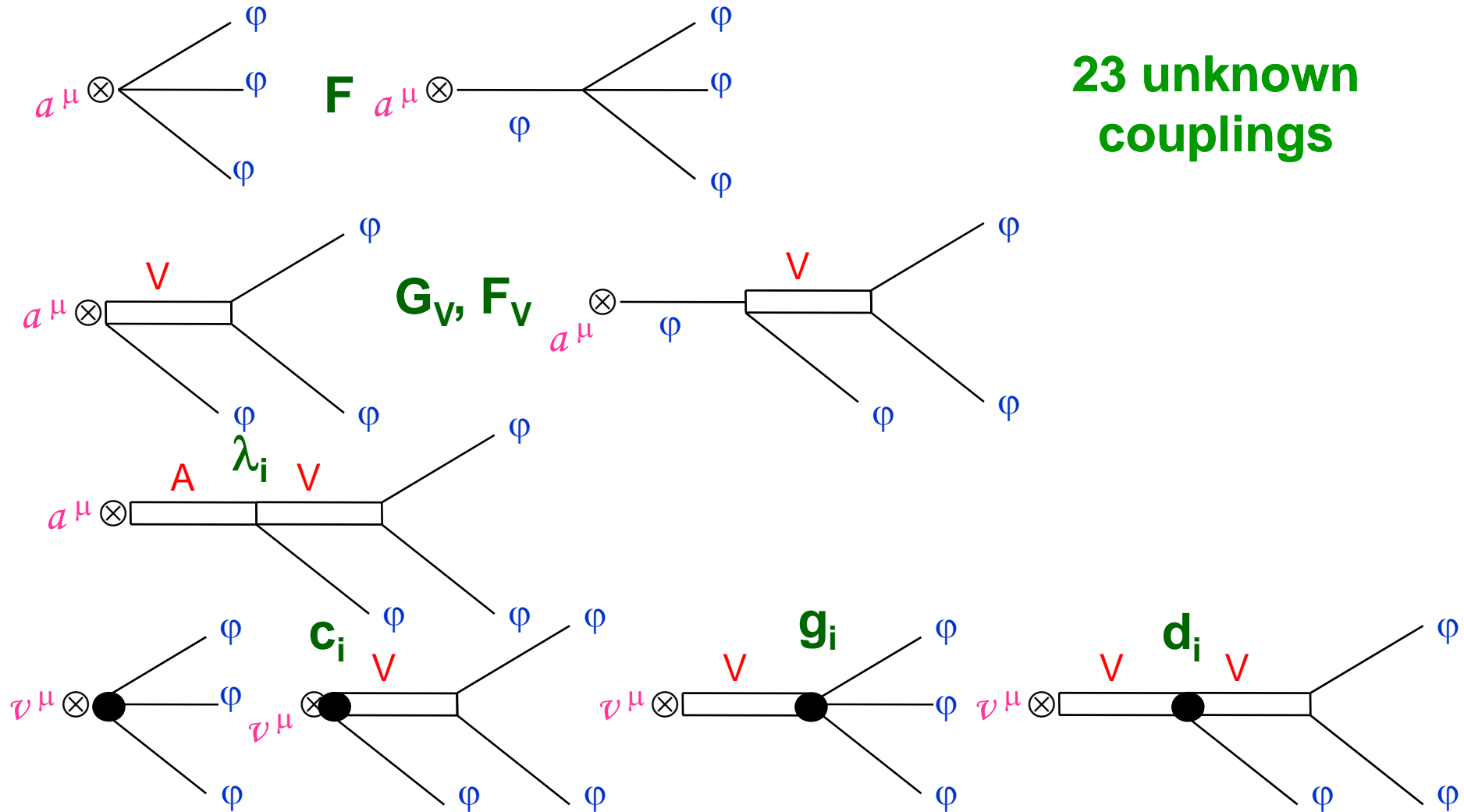


HADRONIZATION IN 3 MESON CHANNELS
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R_χT APPLIED



23 unknown couplings

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HADRONIZATION IN 3 MESON CHANNELS
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The axial-form factor and the $a_1: \tau^- \rightarrow (3\pi)^- \nu_\tau$

(Gómez-Dumm, Pich, Portolés '00)

$$\Gamma_\rho(s) = \frac{M_\rho s}{96\pi F^2} \left[\sigma_\pi^3 \Theta(s - 4m_\pi^2) + \frac{1}{2} \sigma_K^3 \Theta(s - 4m_K^2) \right]$$

(This work)

$$\Gamma_{a_1}(Q^2) = \Gamma_{a_1}^{3\pi}(Q^2) + \Gamma_{a_1}^{K\bar{K}\pi}(Q^2) + \Gamma_{a_1}^{(K\pi)^0 K^0}(Q^2),$$

$$\Gamma_{a_1}^{3\pi}(Q^2) = \frac{1}{48(2\pi)^3 M_{a_1}} \left(\frac{Q^2}{M_{a_1}^2} \right) \iint ds dt \left(F_1' V_{1\mu} + F_2' V_{2\mu} \right).$$

$$\left(F_1'^{\dagger} V_{1\mu} + F_2'^{\dagger} V_{2\mu} \right), \quad F_i' = F_i \frac{M_{a_1}^2 - Q^2}{\sqrt{2} F_A Q^2}$$