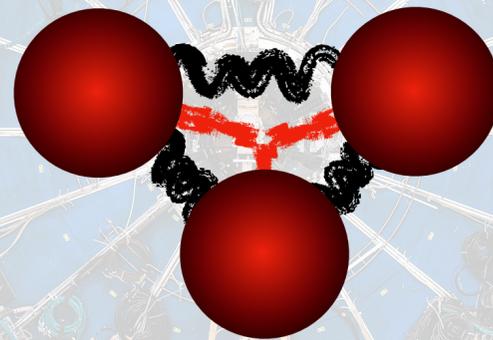




Novel technique to access the three-body interactions with ALICE at the LHC

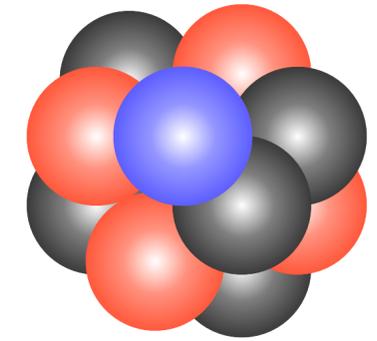
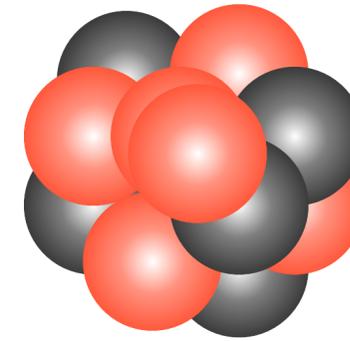


Laura Šerkšnytė
Technical University of Munich
On behalf of the ALICE Collaboration
Genova 07.06.23

Many-body systems

- Properties of nuclei and hypernuclei cannot be described satisfactorily with two-body forces only

L. Girlanda et al., PRC 102, 064003 (2020)



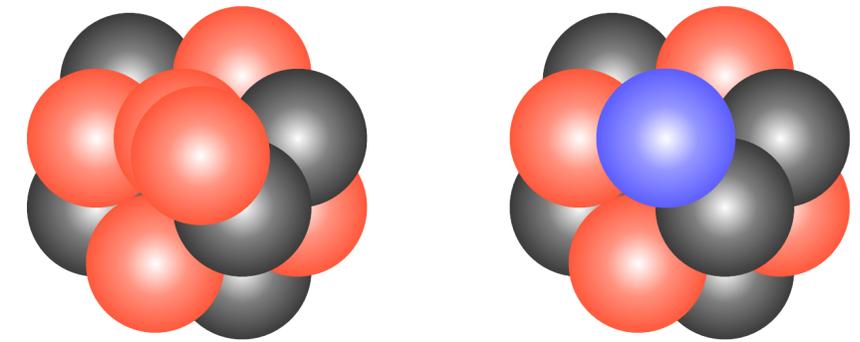
ρ



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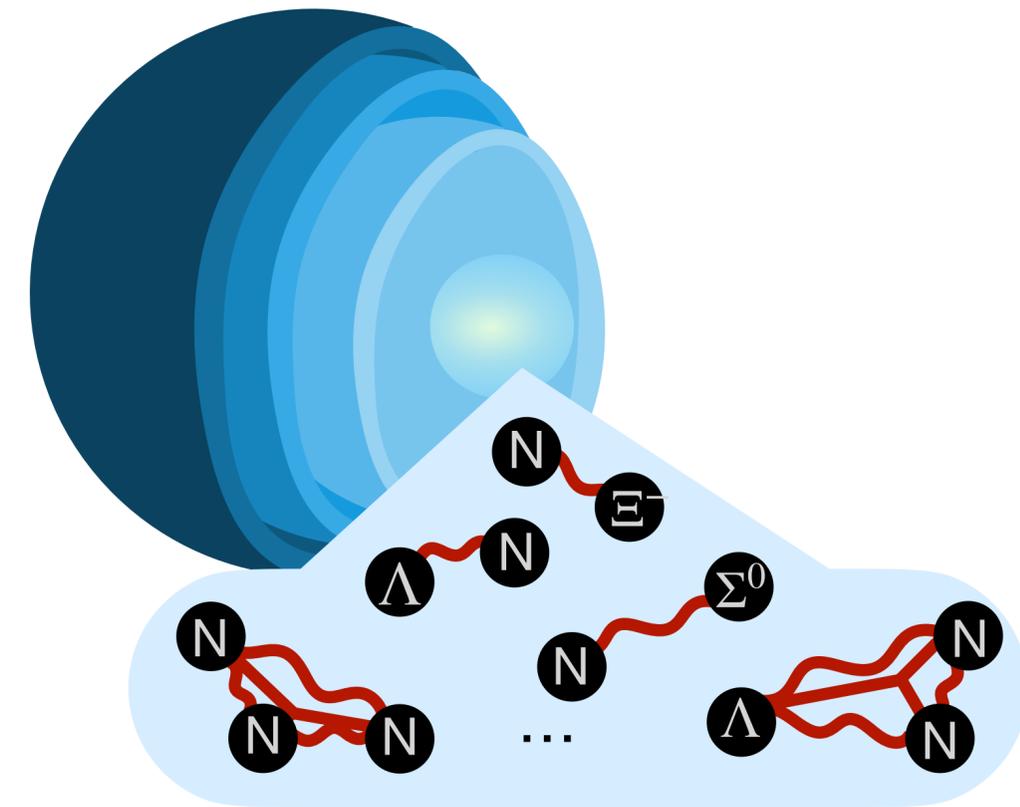
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L. Girlanda et al., PRC 102, 064003 (2020)



- ρ
- The equation of state of dense matter is increasingly sensitive to the three-body forces with increasing density
 - Description of such dense systems as neutron stars also requires three-body interactions

D. Lonardonì et al., PRL 114, 092301 (2015)



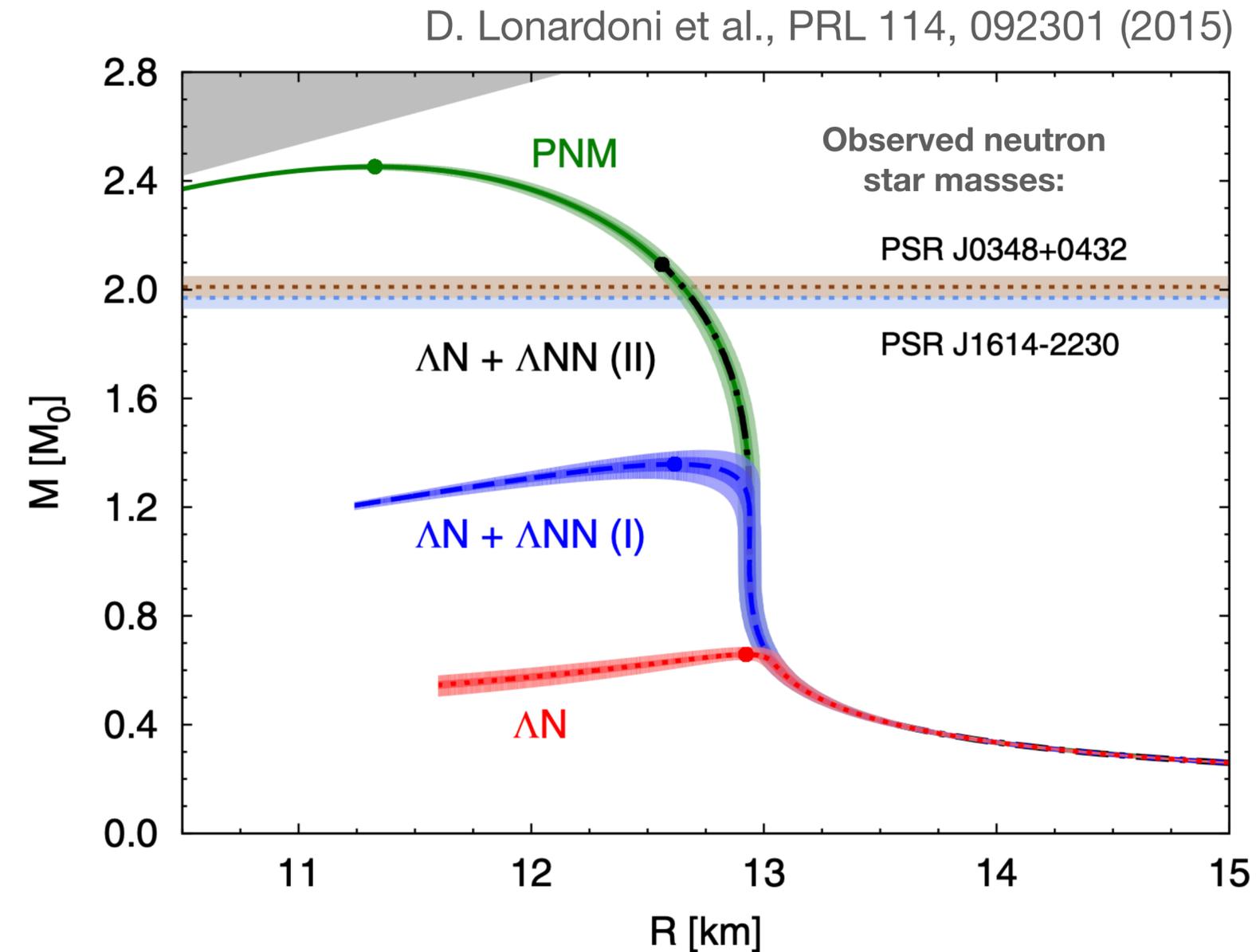
Marcel Lesch 8 Jun 2023, 15:12

Neutron stars and three-body forces

- Three-body interaction models are fitted to reproduce measured (hyper)nuclei properties

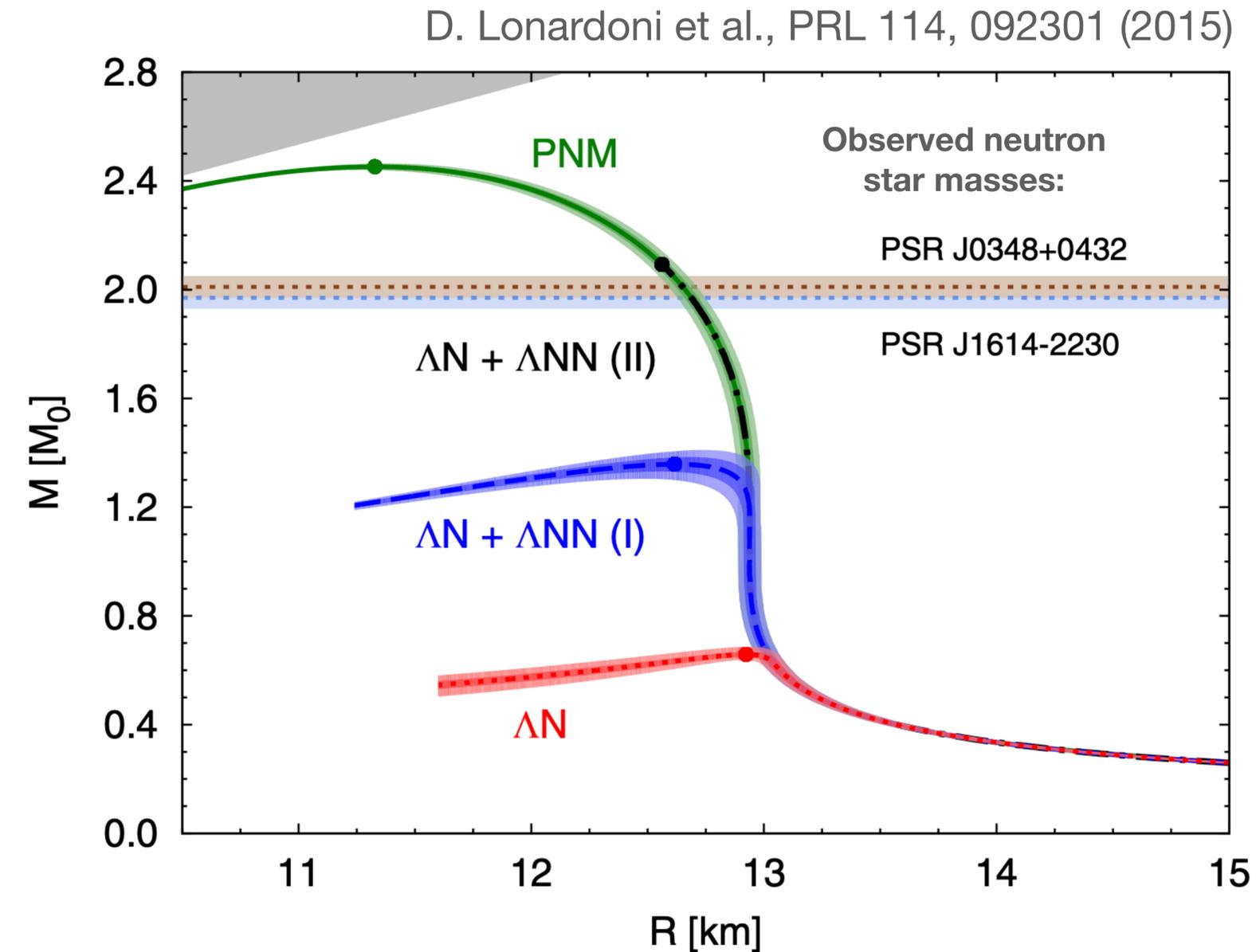
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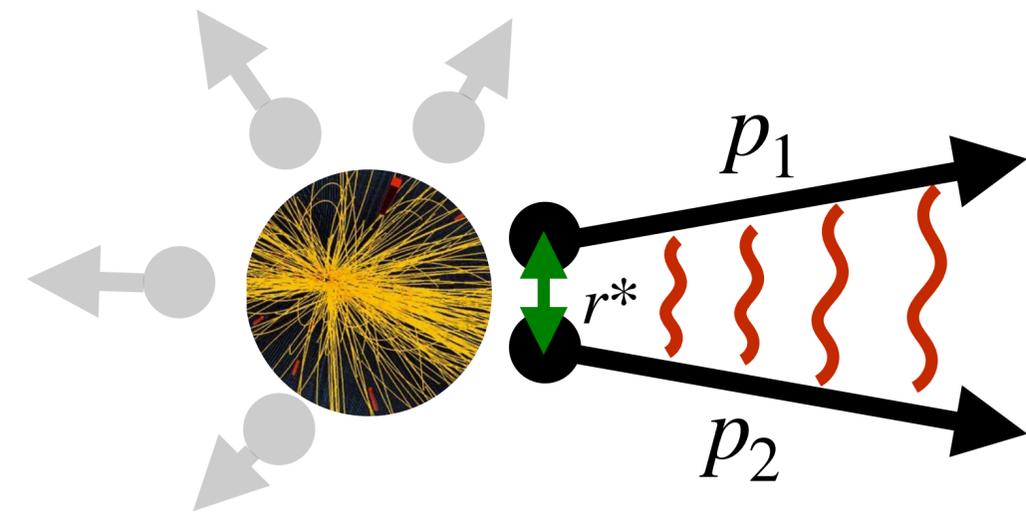
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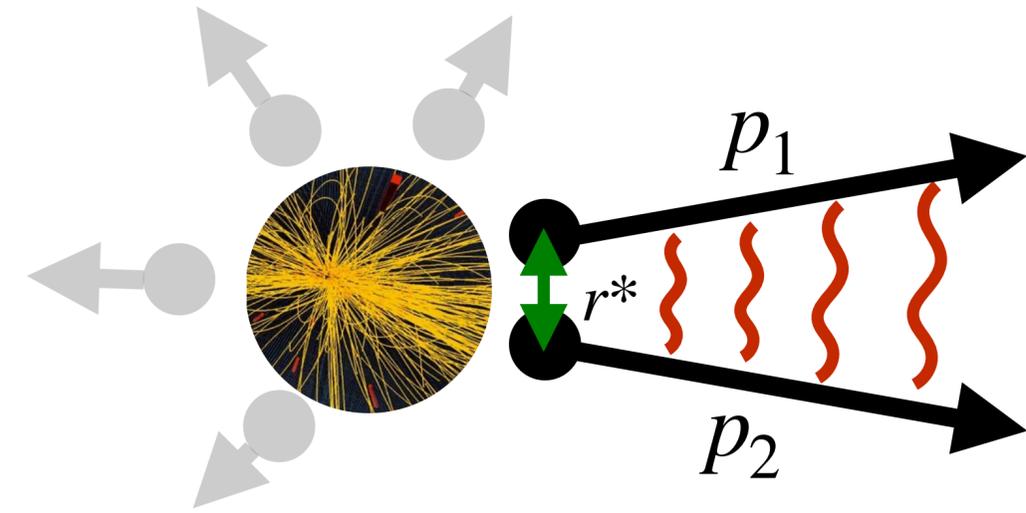


New observables are required to solve the three-body problem!

Femtoscscopy

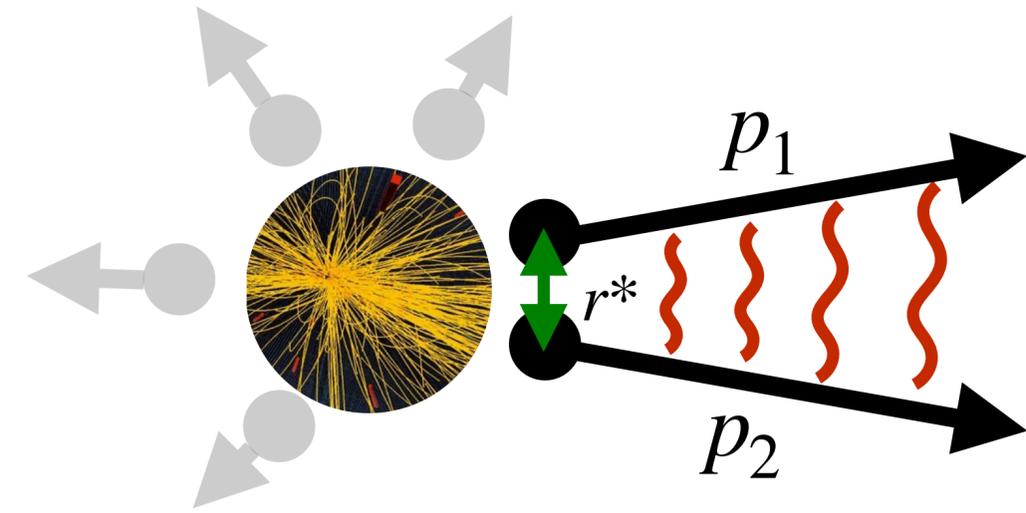


Emission source $S(r^*)$

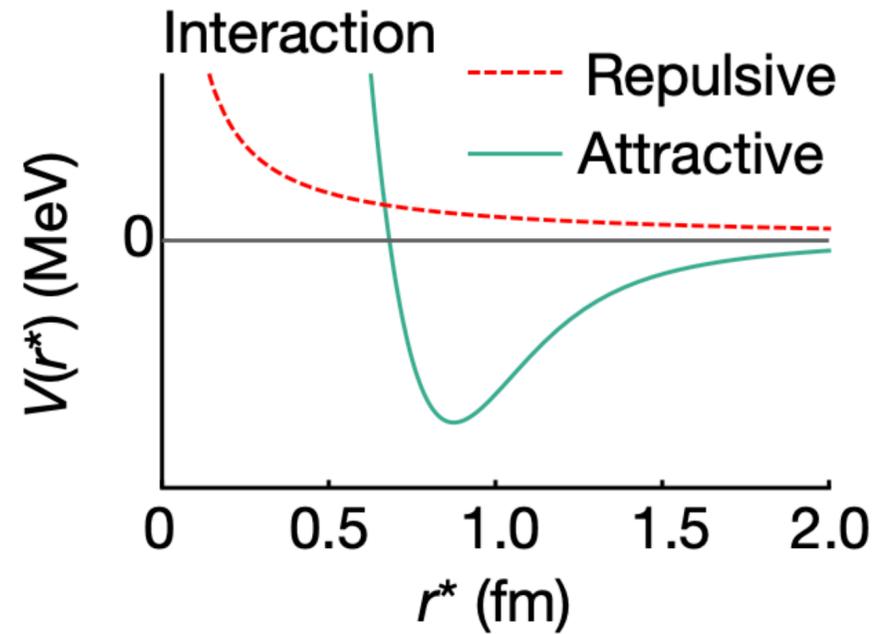


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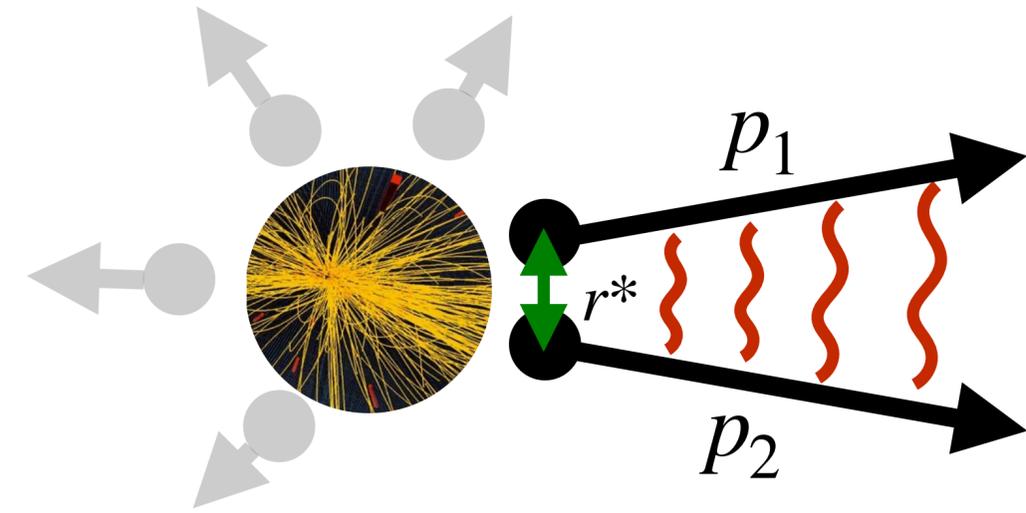


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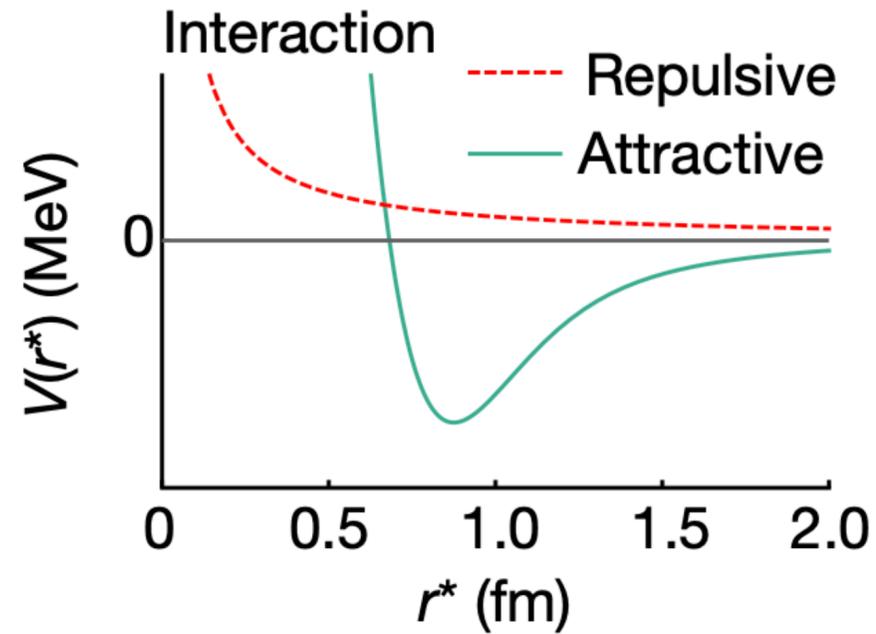


Schrödinger equation
Two-particle wave function
 $|\psi(\mathbf{k}^*, \mathbf{r}^*)|$

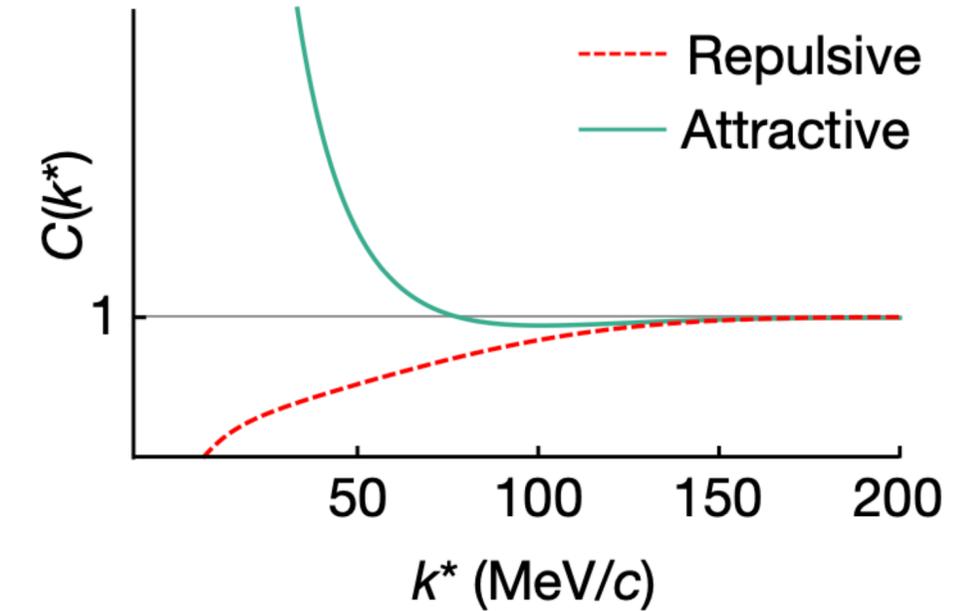
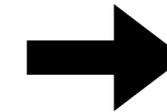
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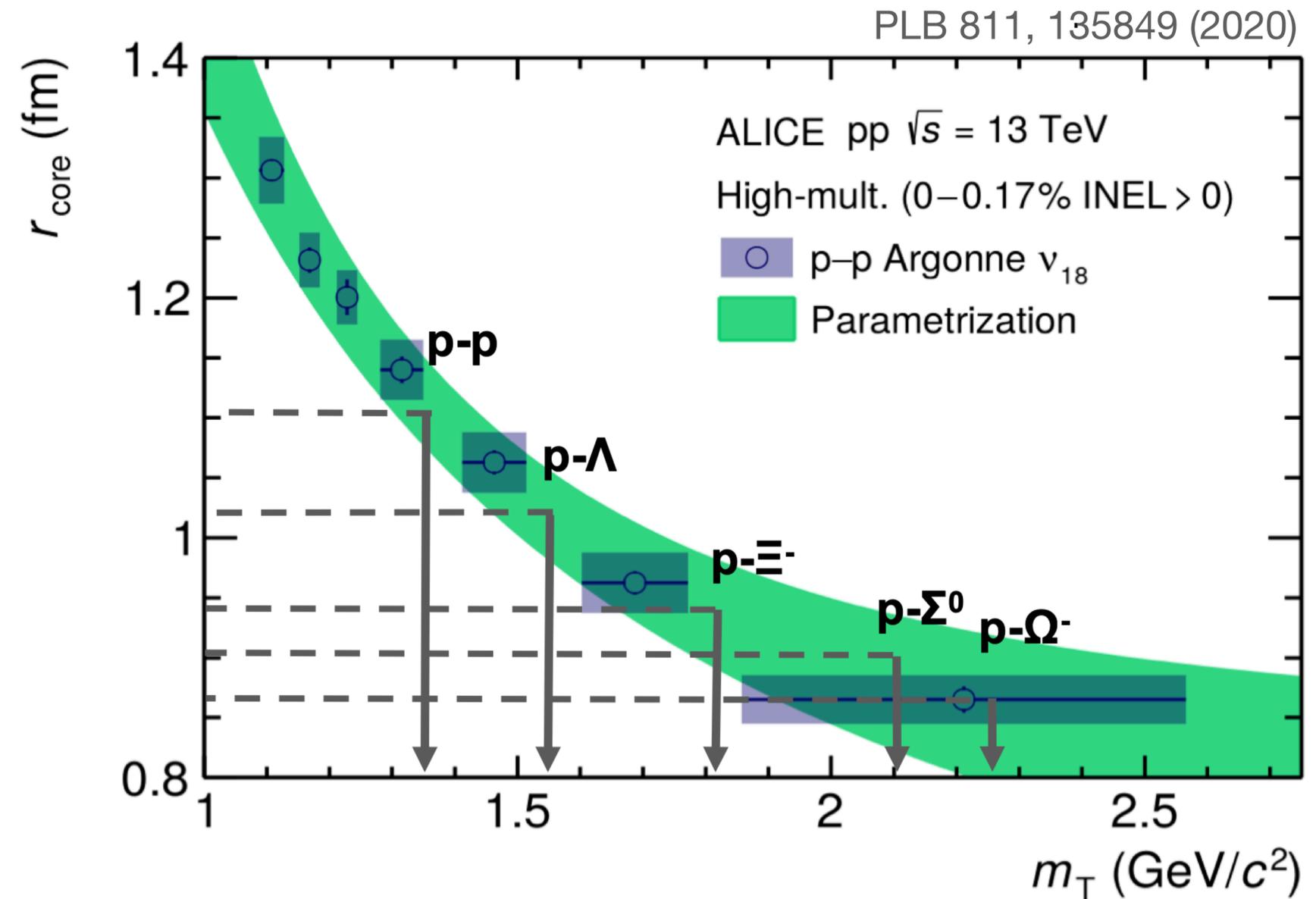


Correlation function $C(k^*)$

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Emission source

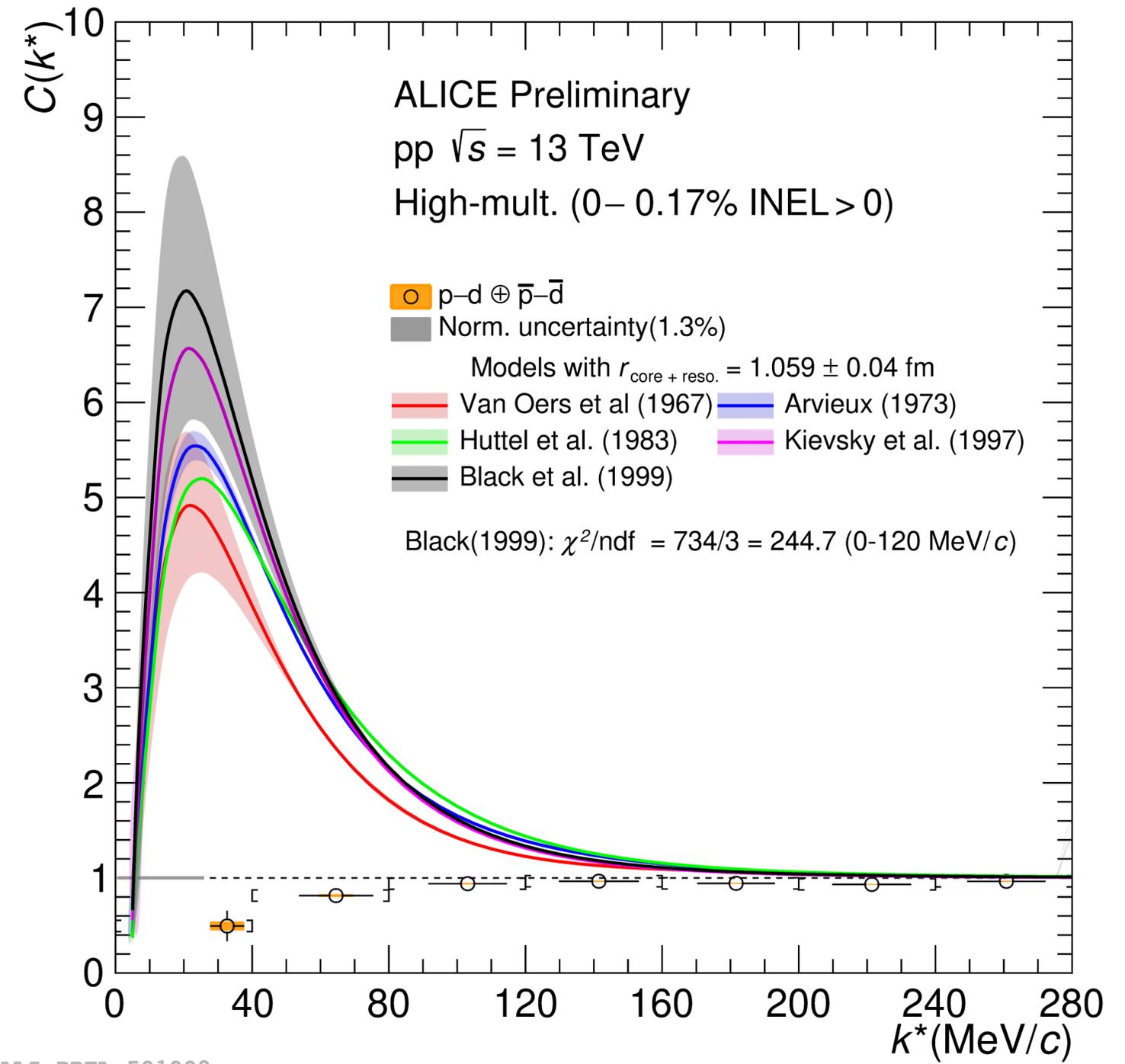
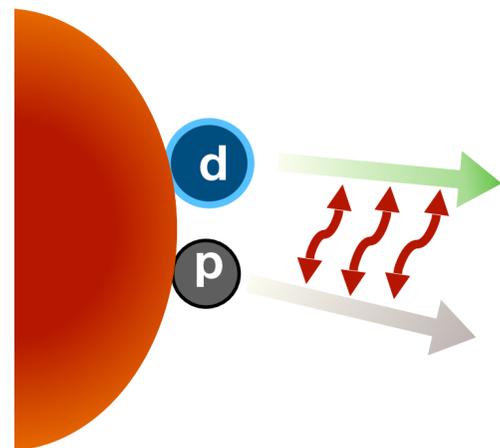
- Two main contributions:
 - general: Collective effects result in Gaussian core
 - specific: Decaying resonances require source correction
- Access to very small distances in pp collisions at the LHC



Proton-deuteron correlation

- Effective two-body system
 - Coulomb + Strong interactions via Lednický model; only s-wave
R. Lednický, Phys. Part. Nuclei 40, 307–352 (2009)
- Anchored to scattering experiments
- Emission source: from m_T scaling

$$r_{\text{eff}}^{\text{pd}} = 1.059^{+0.04}_{-0.04} \text{ fm}$$



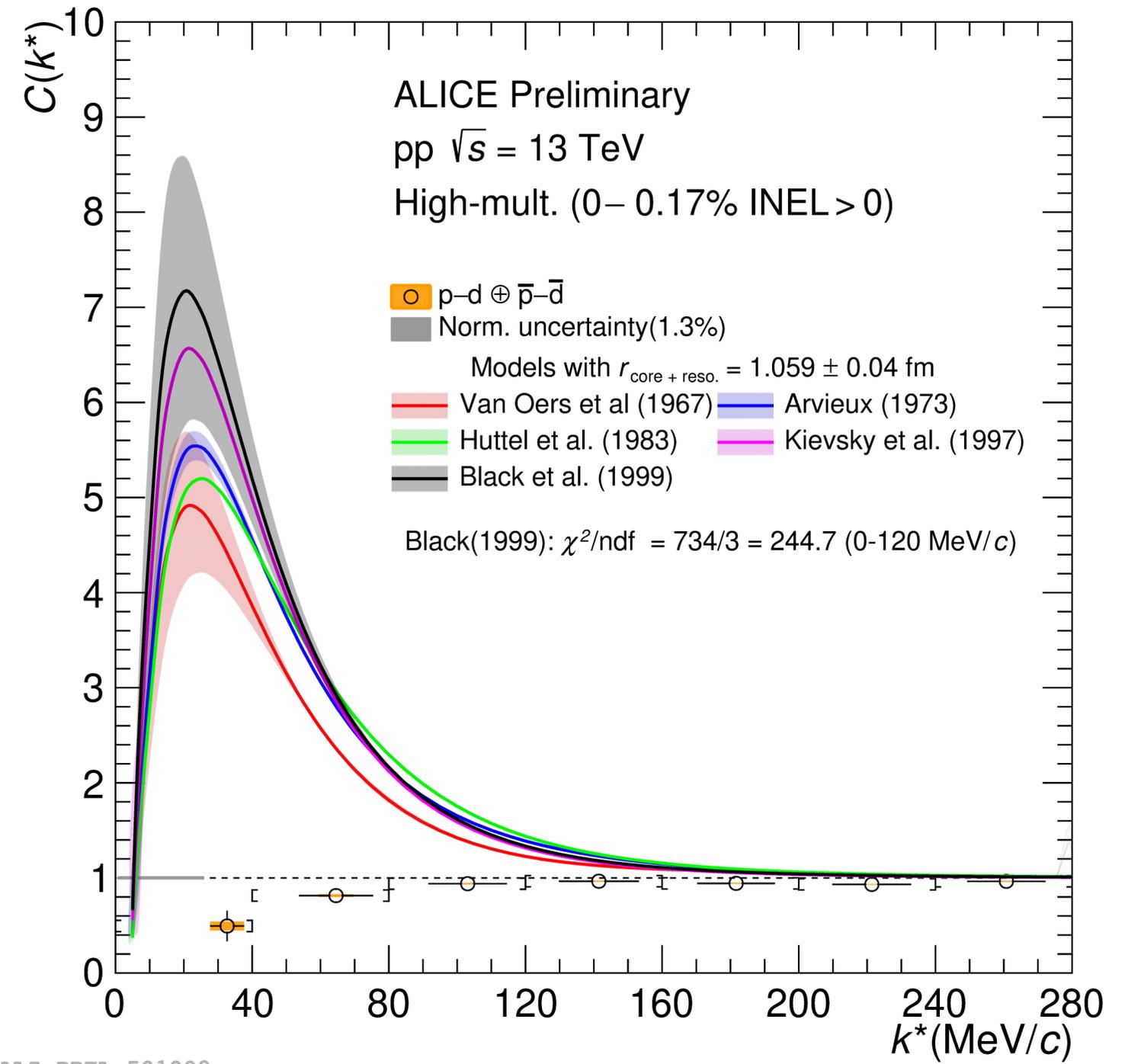
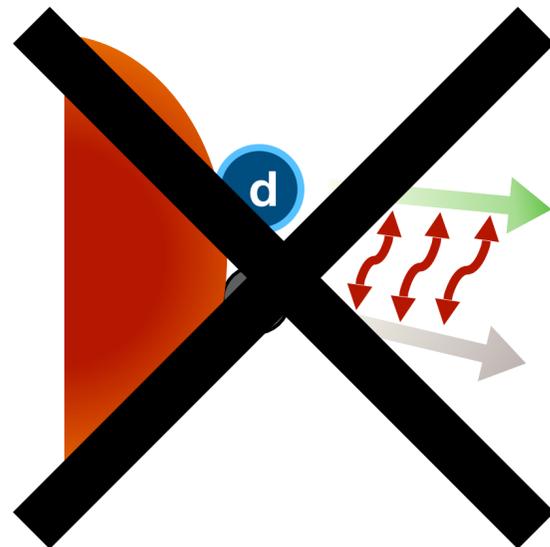
ALI-PREL-501009

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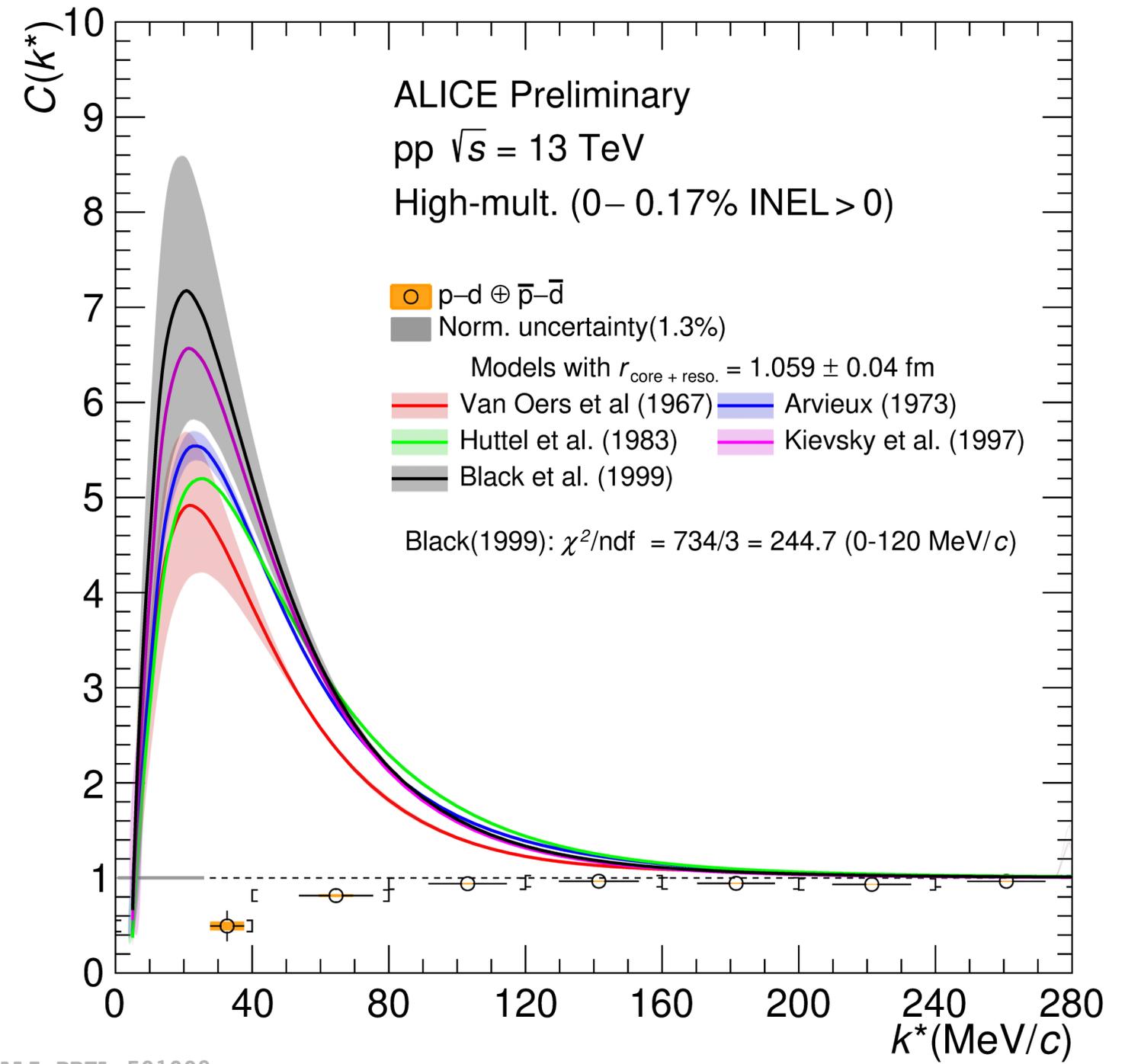
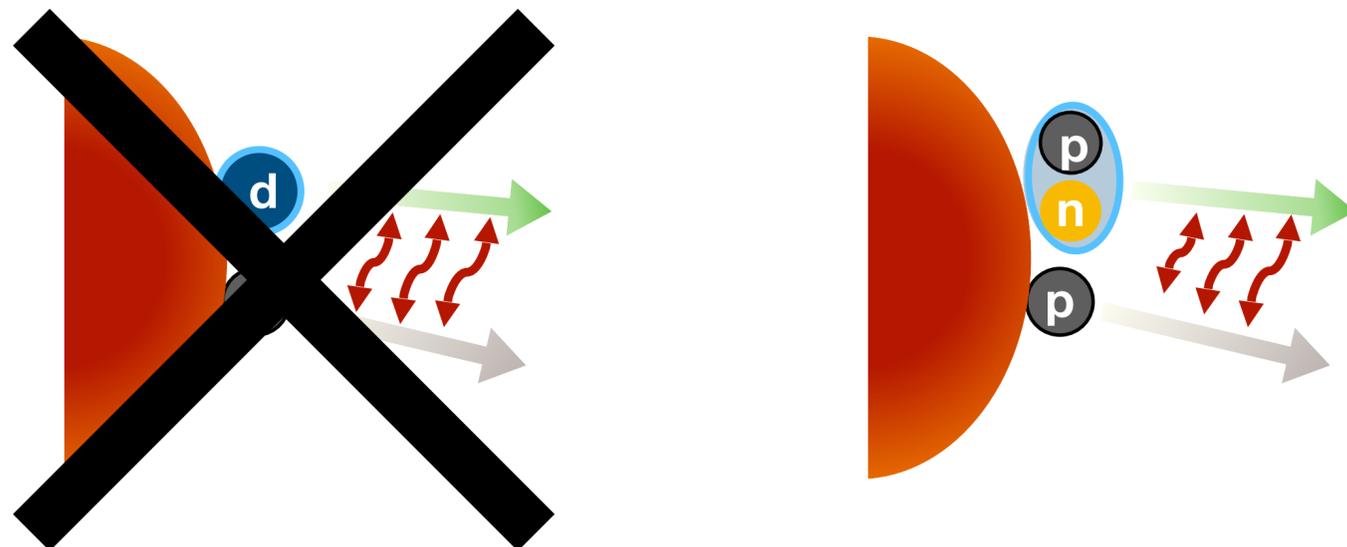
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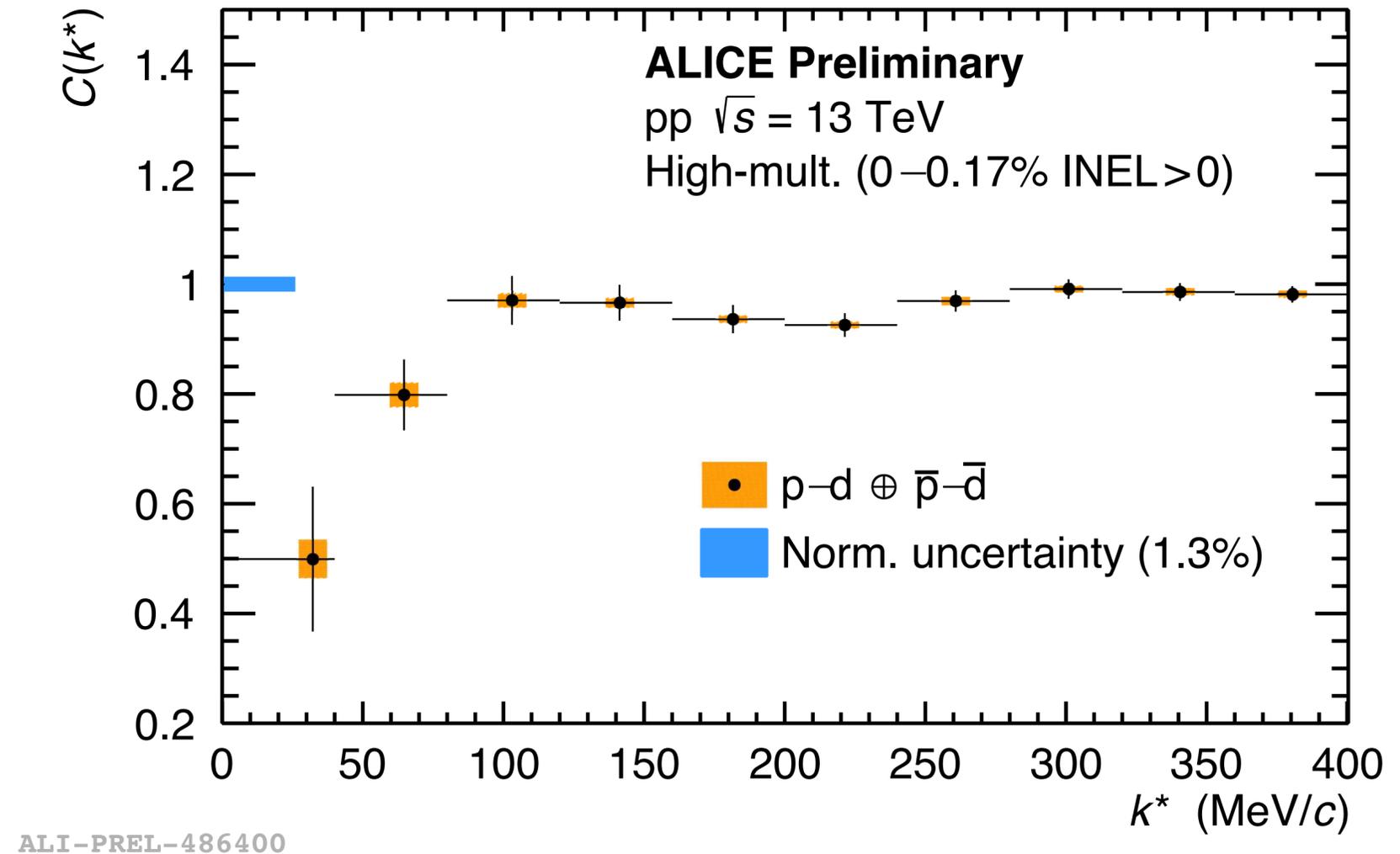
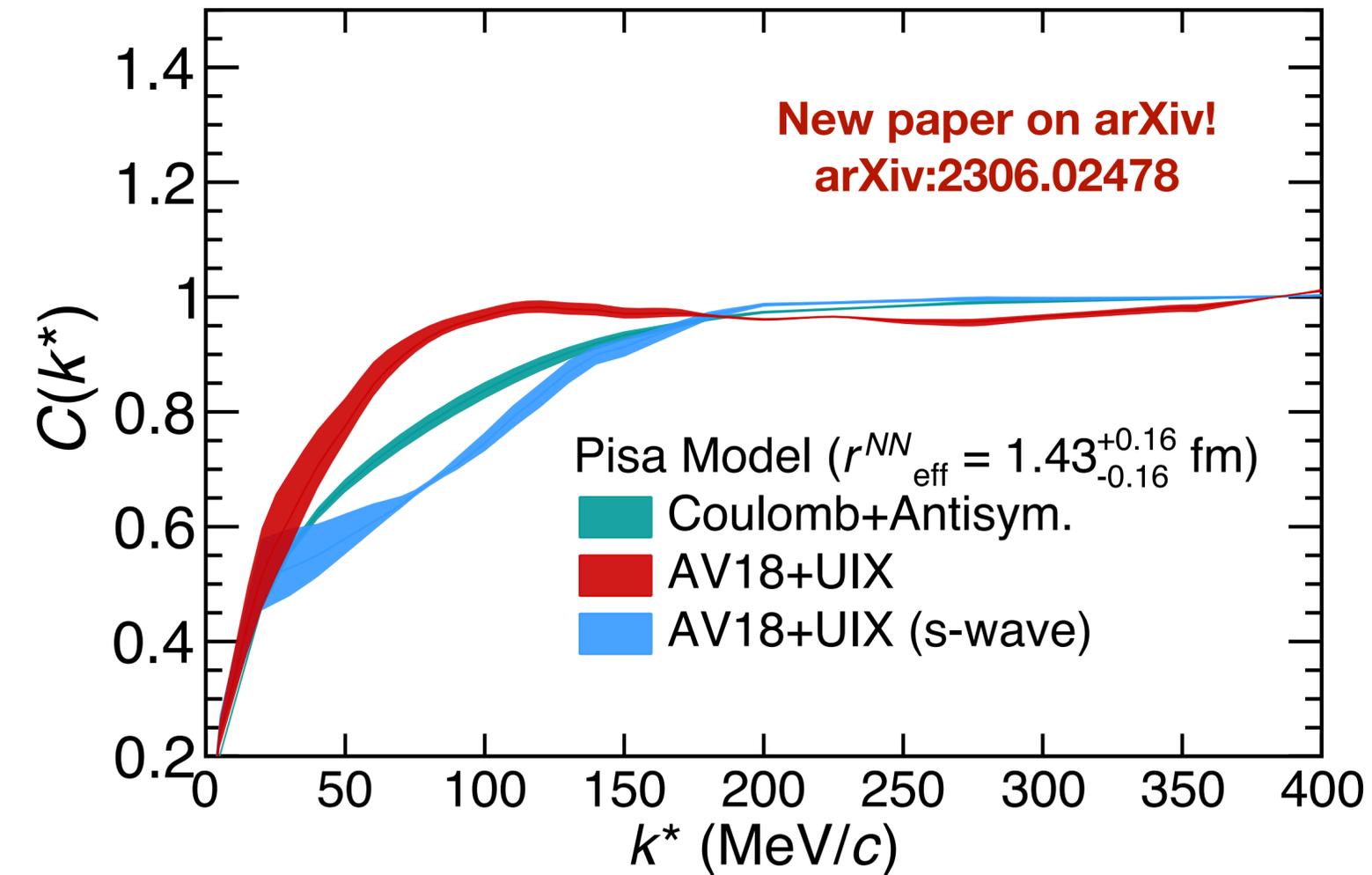
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ALI-PREL-501009

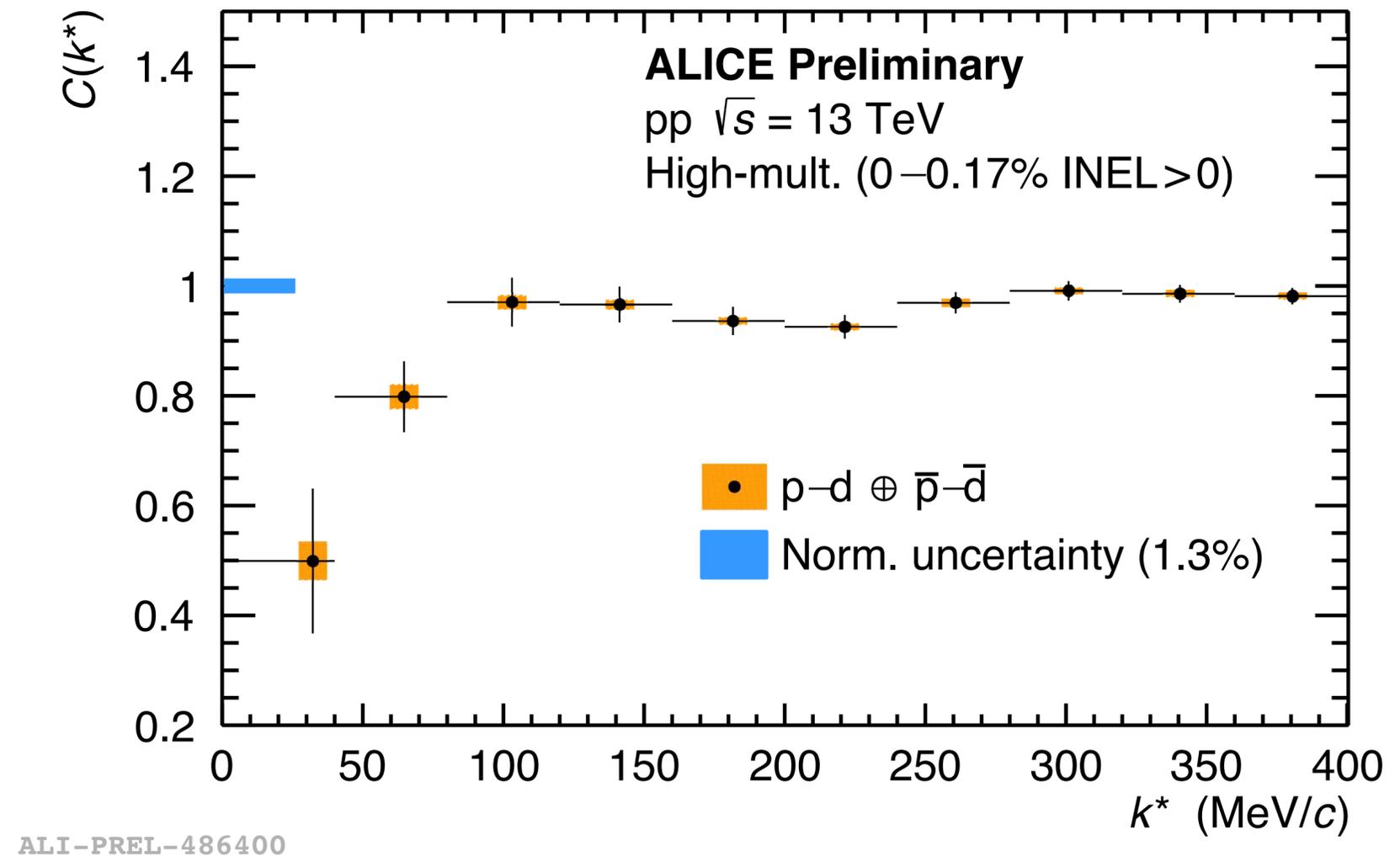
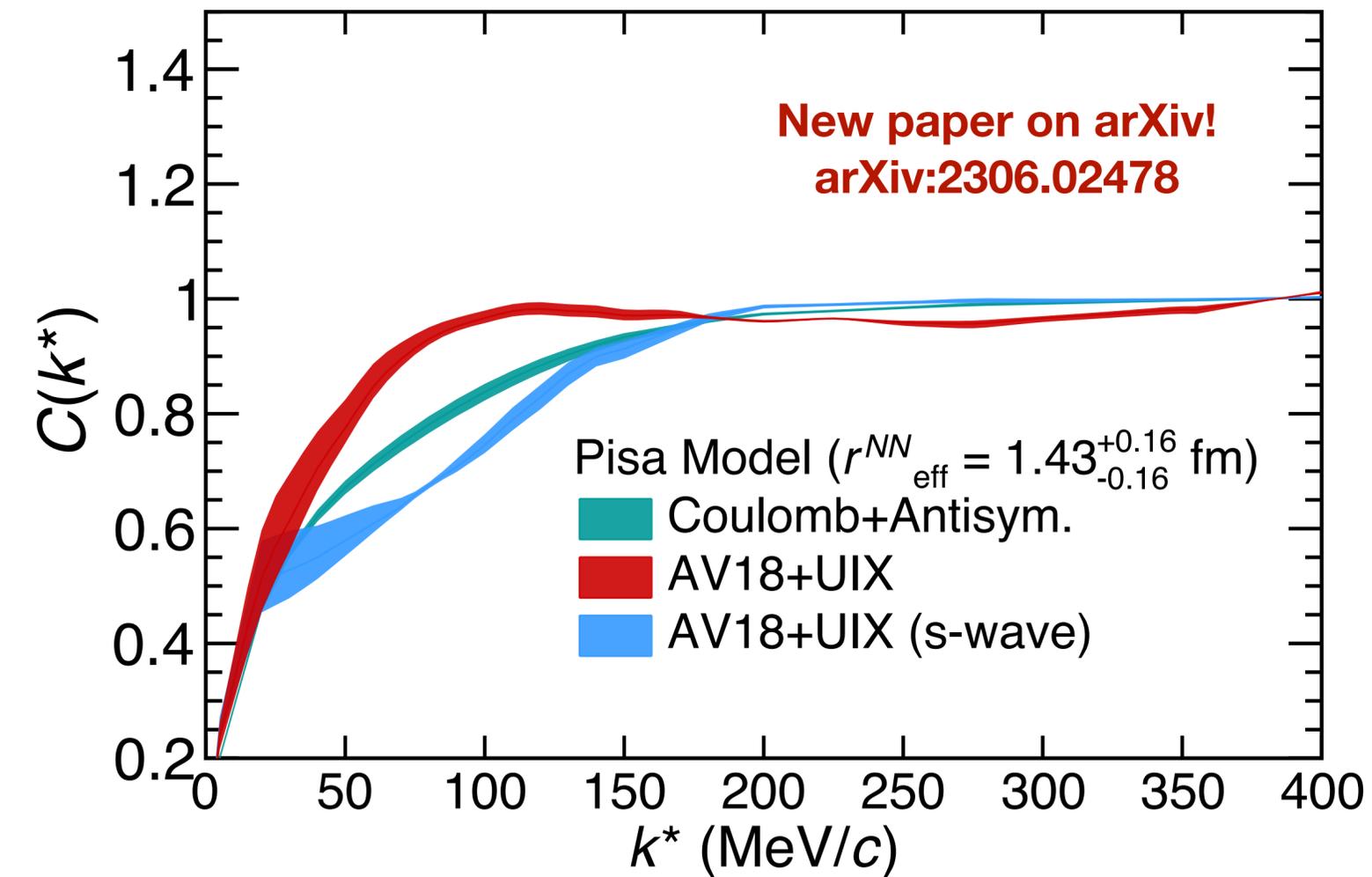
Three-body dynamics in p-d system

- Modelling of the p-d correlation function as a three-body system can describe the data well
→ Includes two-body (AV18) and three-body (UIX) interactions



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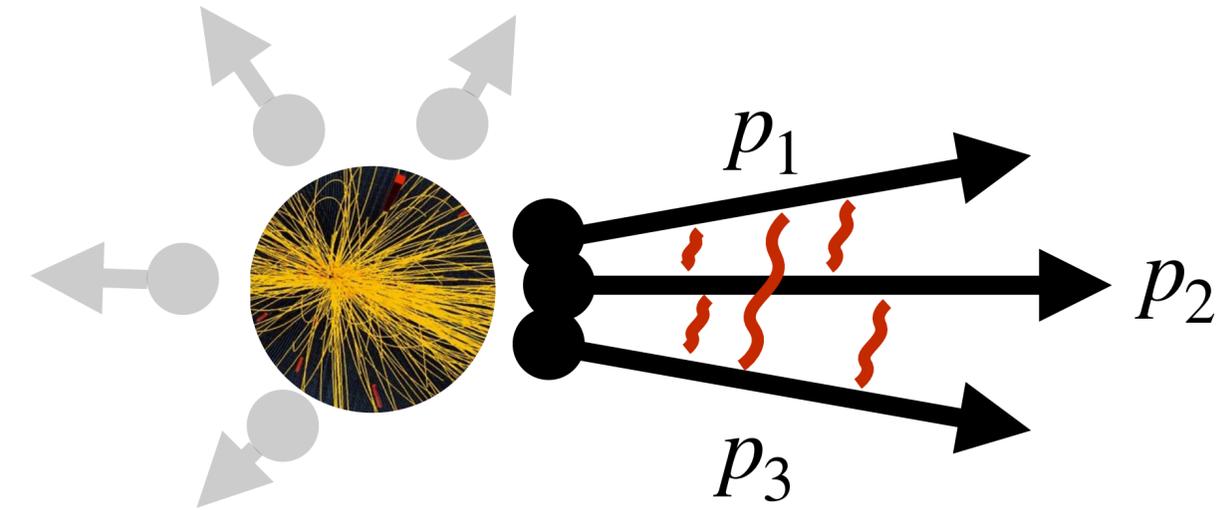
System sensitive to the three-body dynamics and also the three-body interaction!

Three-particle femtoscopy

- Experimental three-particle correlation function

$$C(Q_3) = \mathcal{N} \frac{N_{\text{same}}(Q_3)}{N_{\text{mixed}}(Q_3)}$$

$$Q_3 = \sqrt{-q_{ij}^2 - q_{jk}^2 - q_{ki}^2}$$

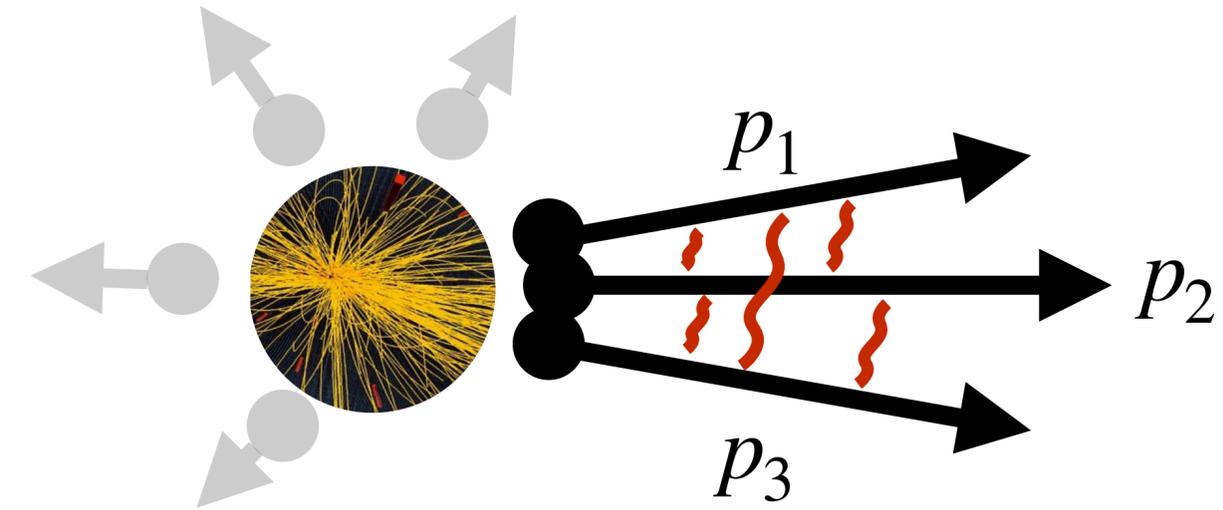


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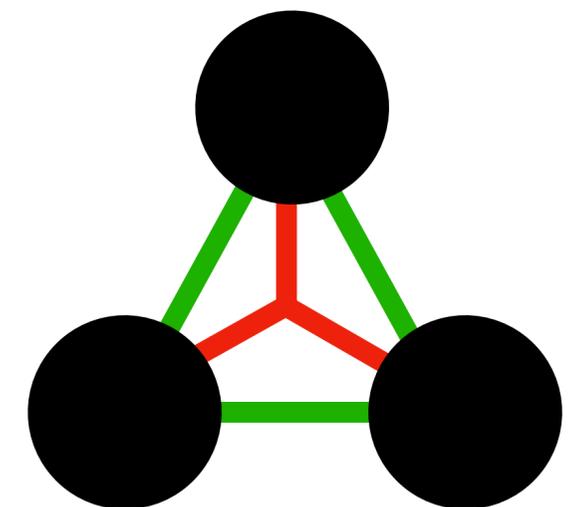
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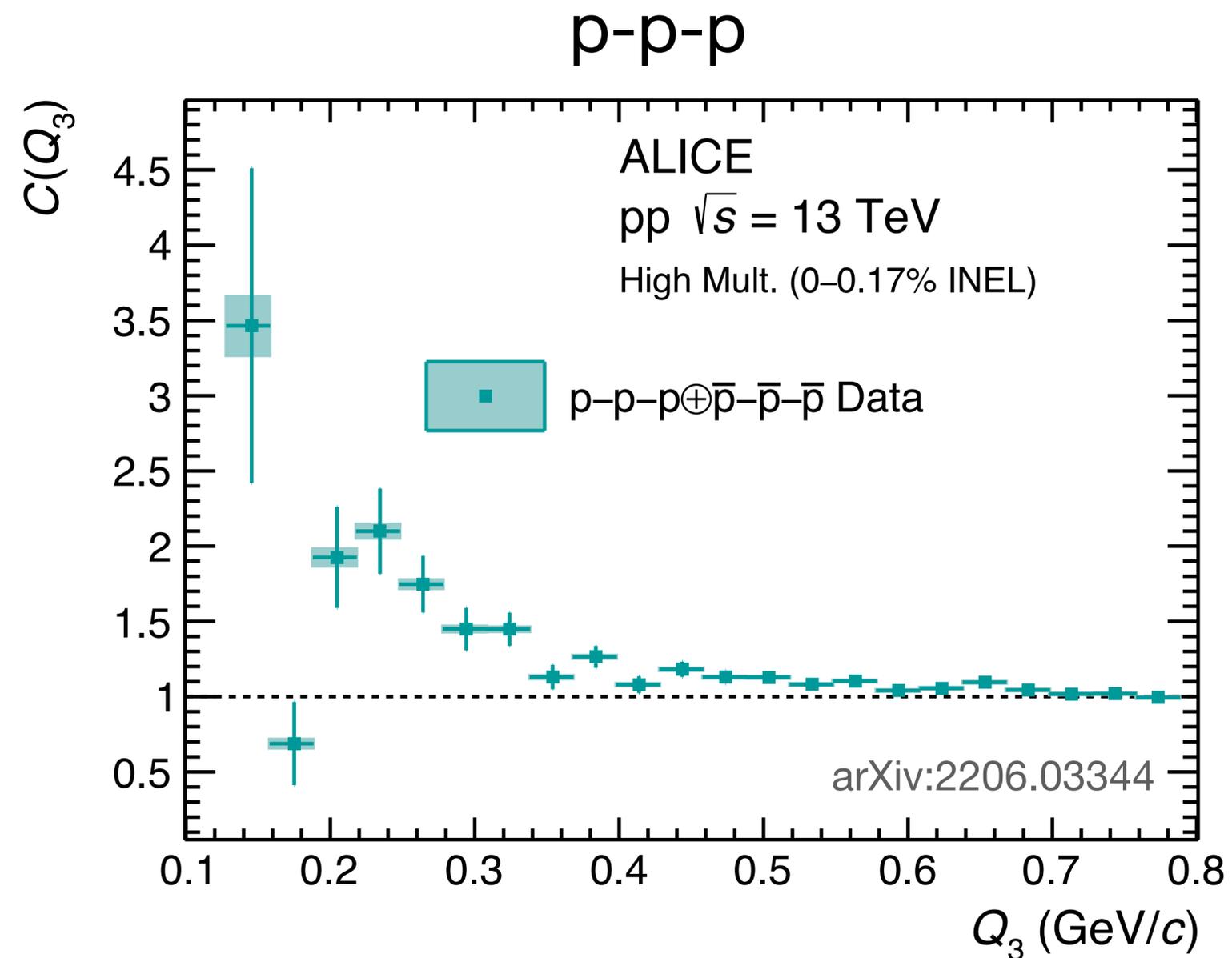
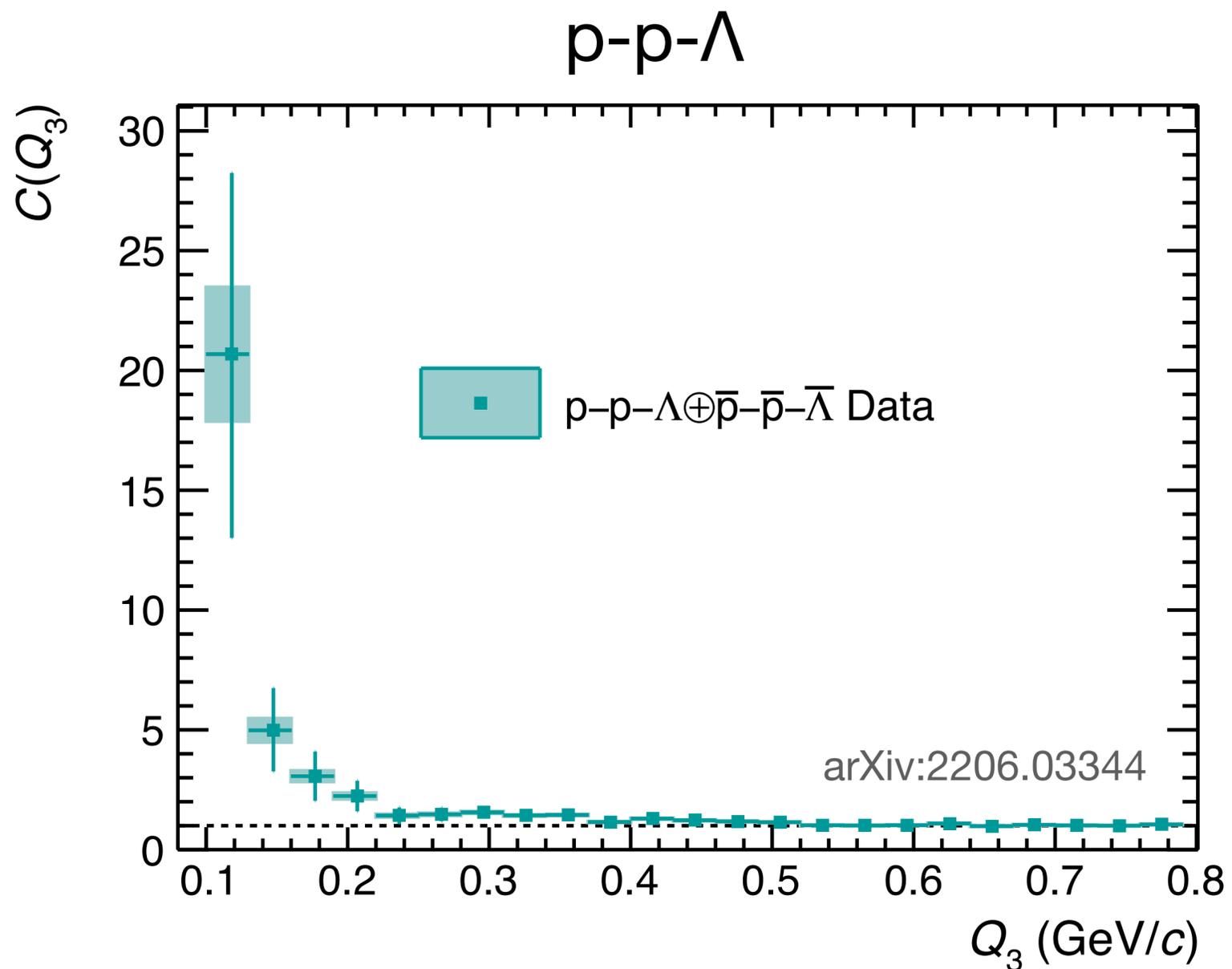
- two-body interactions
- three-body interaction

$$C_3(\mathbf{p}_1, \mathbf{p}_2, \mathbf{p}_3) = \iiint S_3(\mathbf{r}_1, \mathbf{r}_2, \mathbf{r}_3) \left| \Psi(\mathbf{r}_1, \mathbf{r}_2, \mathbf{r}_3, \mathbf{p}_1, \mathbf{p}_2, \mathbf{p}_3) \right|^2 d^3\mathbf{r}_1 d^3\mathbf{r}_2 d^3\mathbf{r}_3$$



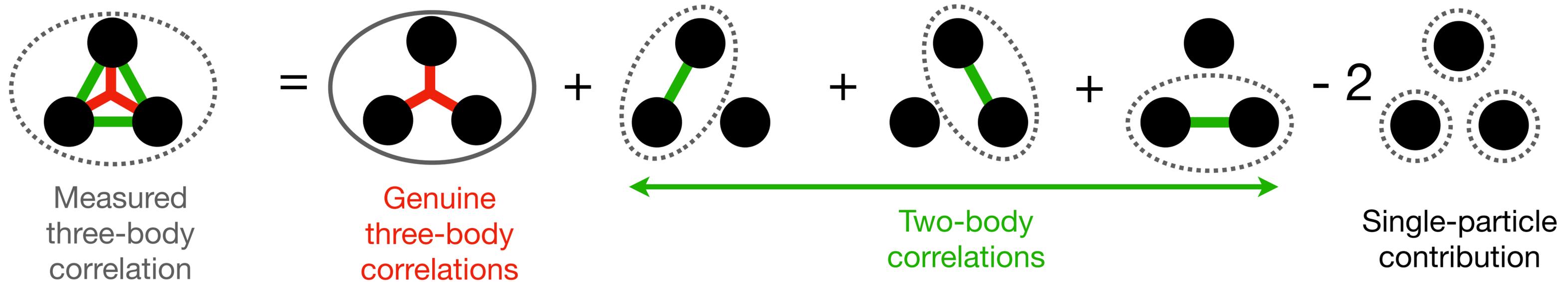
Measured correlation functions

- Measured correlation functions are not equal to unity
 - Are two- or/and three-body interactions responsible?



Cumulants in femtoscopy

The total three-particle correlations can be expressed as a sum of genuine three-body correlation and the lower-order contributions employing Kubo's cumulants [1]:



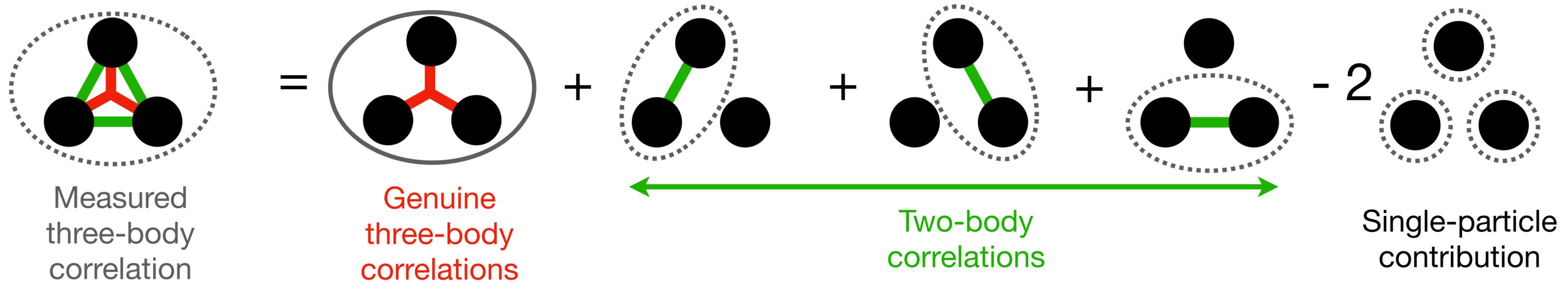
In terms of correlation functions:

$$c_3(Q_3) = C(Q_3) - C_{12}(Q_3) - C_{23}(Q_3) - C_{31}(Q_3) + 2$$

[1] R. Kubo, J. Phys. Soc. Jpn. 17, 1100-1120 (1962)

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Lower-order contributions

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Lower-order contributions

- Two methods:
 - Data-driven method: event mixing
 - Projector method: project two-body correlation function on the three-particle phase space

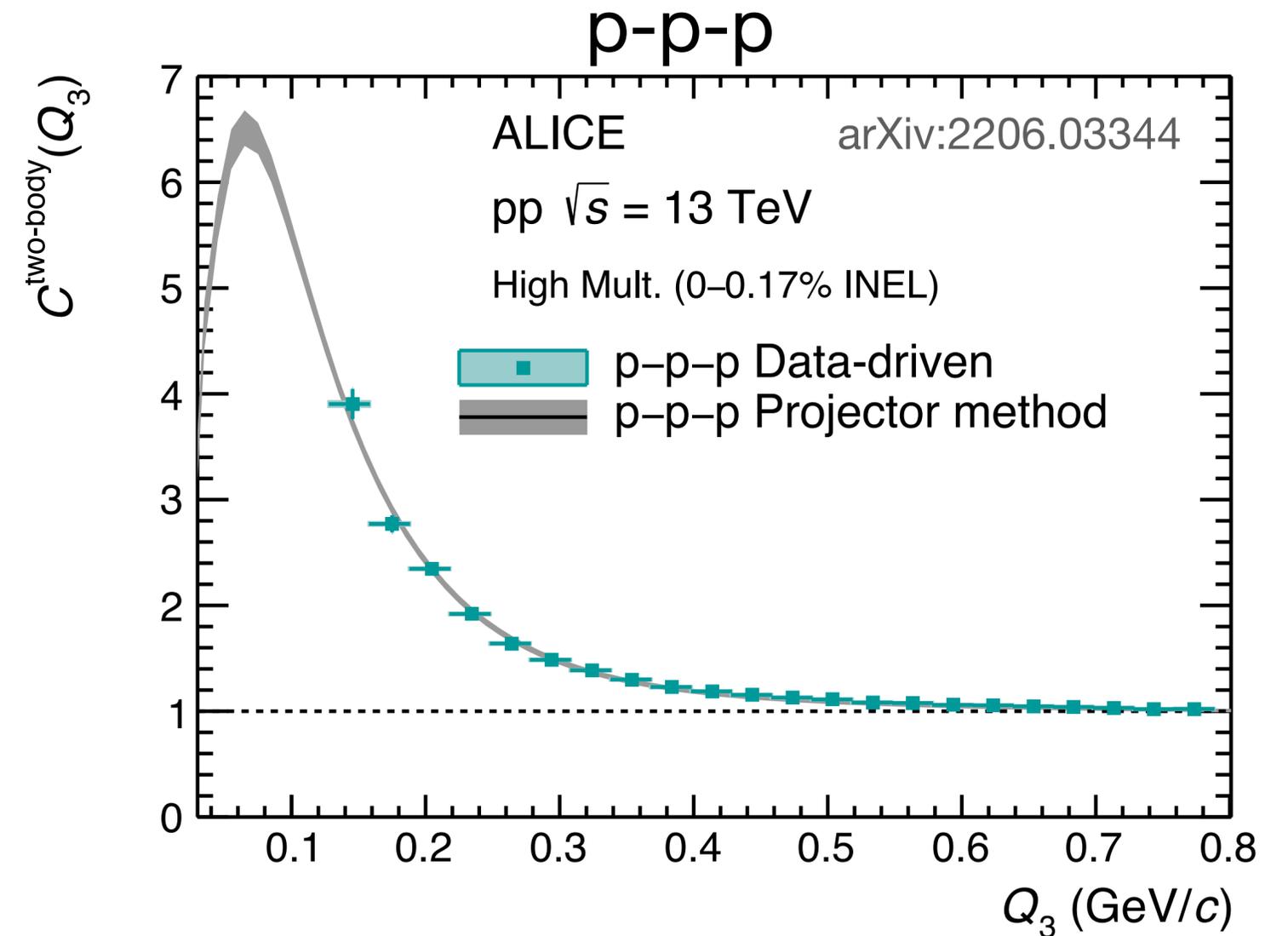
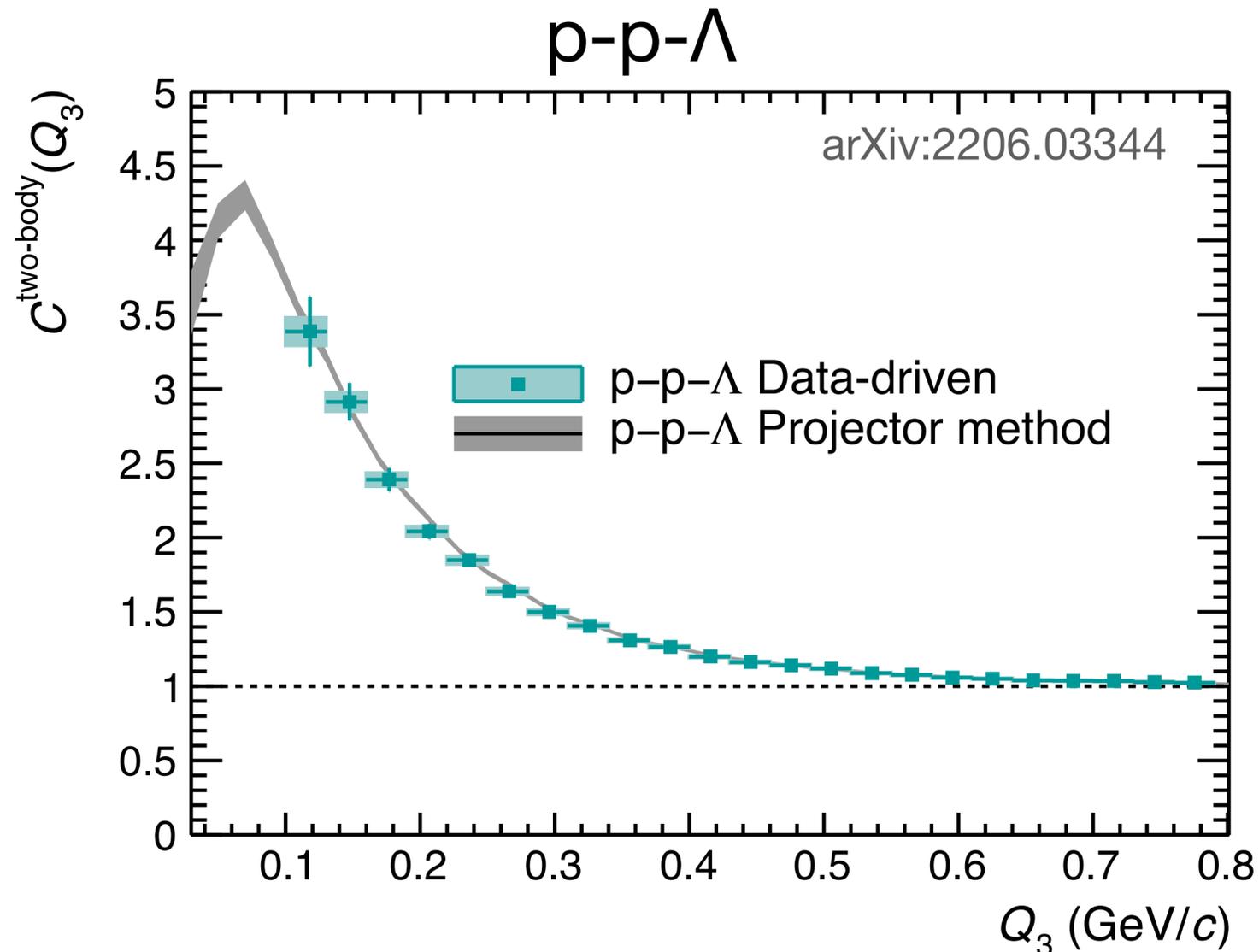
Del Grande, Šerkšnytė et al. EPJC 82 (2022) 244

Lower-order contributions

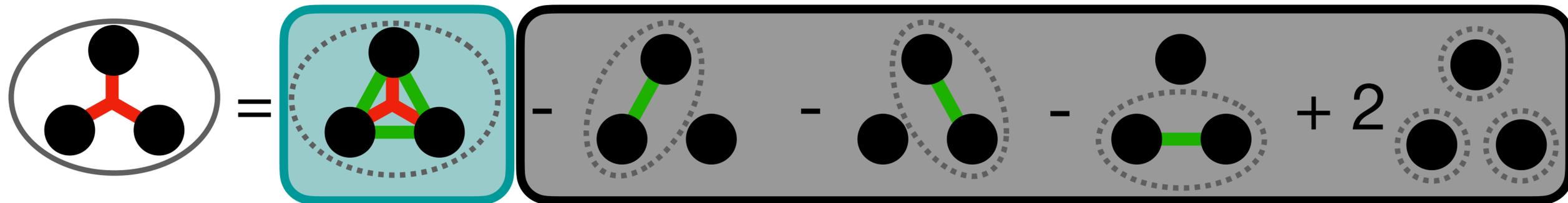
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Lower-order contributions under control!

Del Grande, Šerkšnytė et al. EPJC 82 (2022) 244

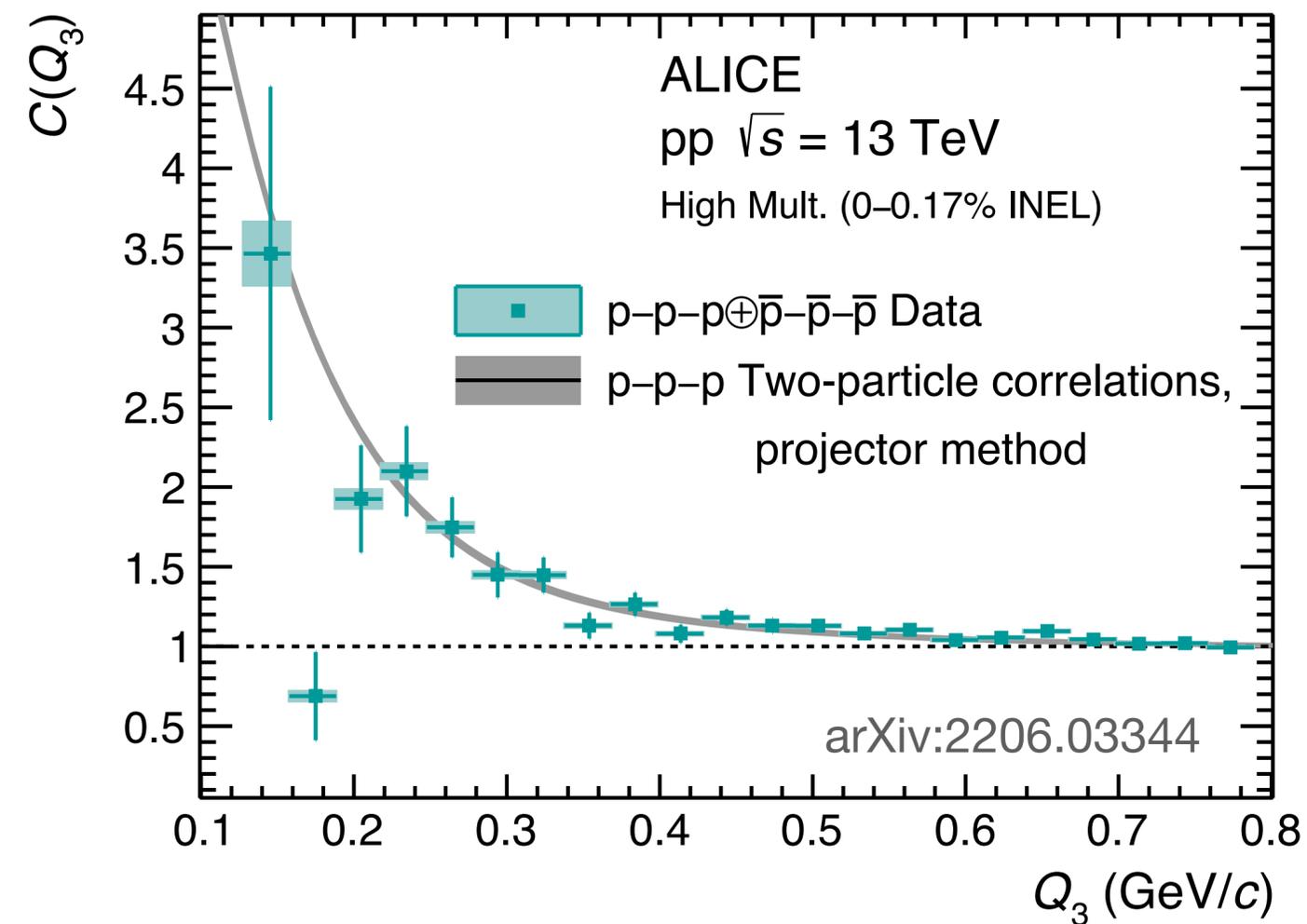
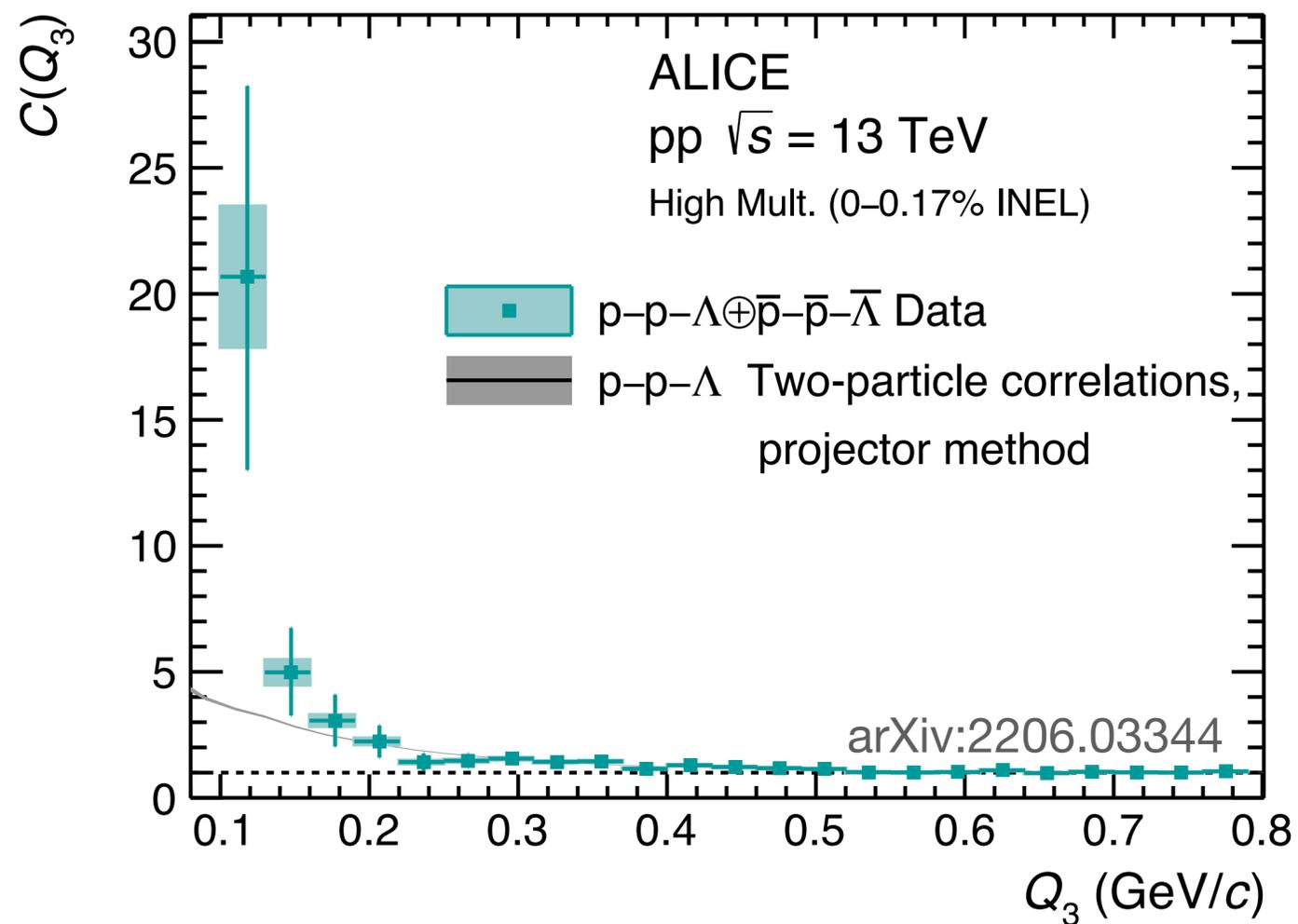


p-p- Λ and p-p-p correlation functions

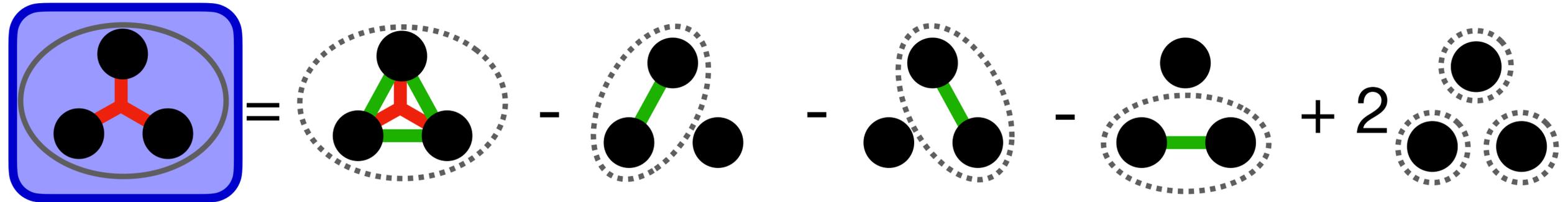


p-p- Λ

p-p-p



p-p- Λ cumulant



Hint of a positive cumulant for p-p- Λ

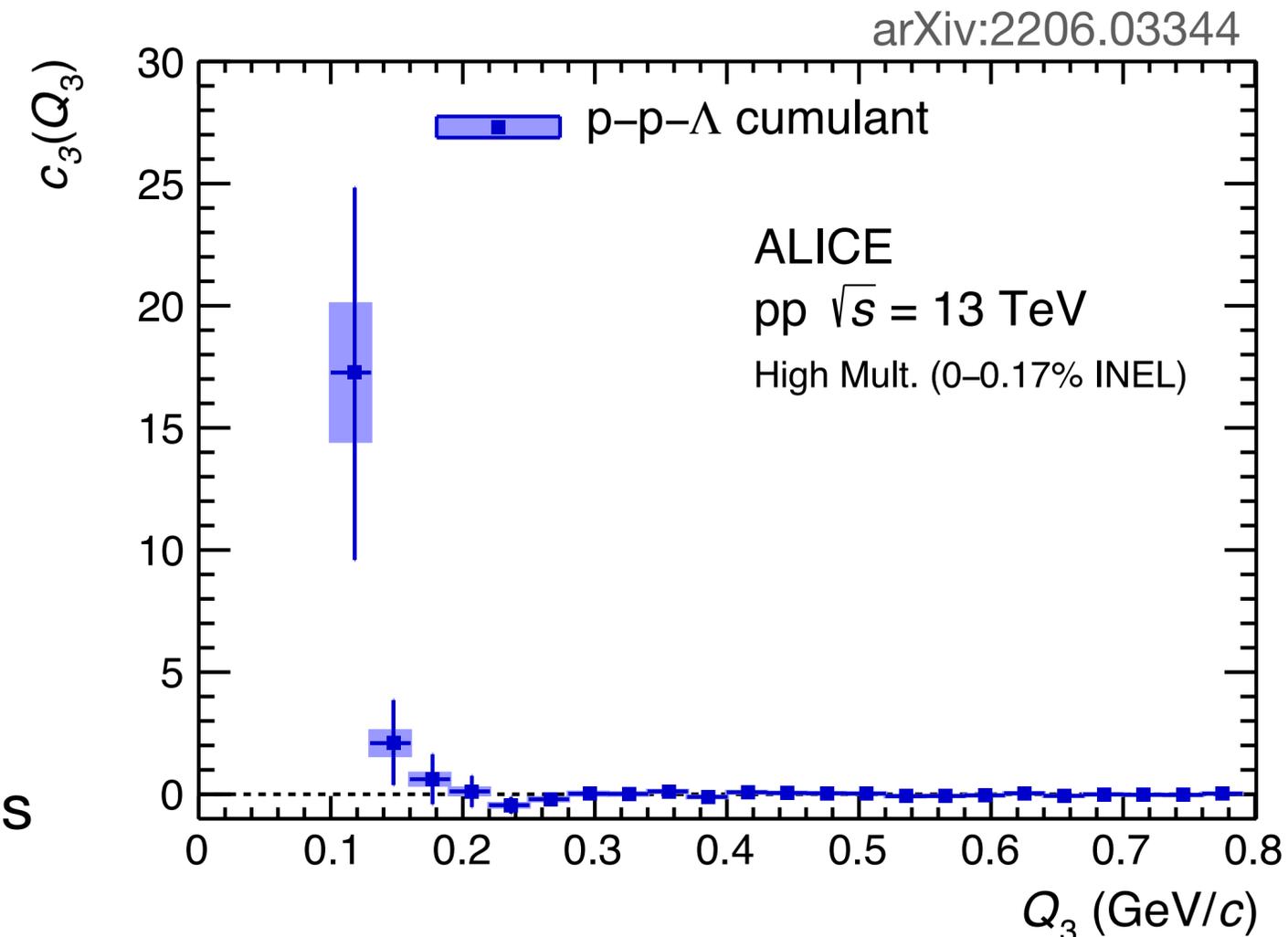
- Only two identical and charged particles
 - ✓ Main expected contribution from three-body strong interaction
- Relevant measurement for equation of state of neutron stars

Statistical significance:

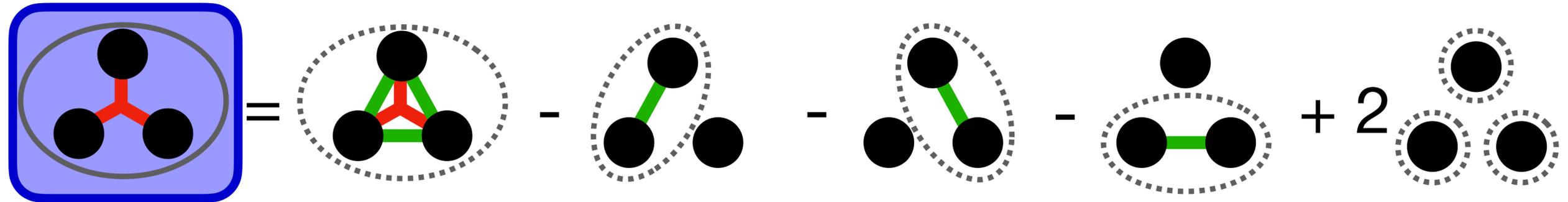
$$n_\sigma = 0.8 \text{ for } Q_3 < 0.4 \text{ GeV}/c$$

Conclusion: no significant deviation from null hypothesis

In upcoming Run 3, two orders of magnitude gain in statistics expected!



p-p-p cumulant



arXiv:2206.03344

Negative cumulant for p-p-p

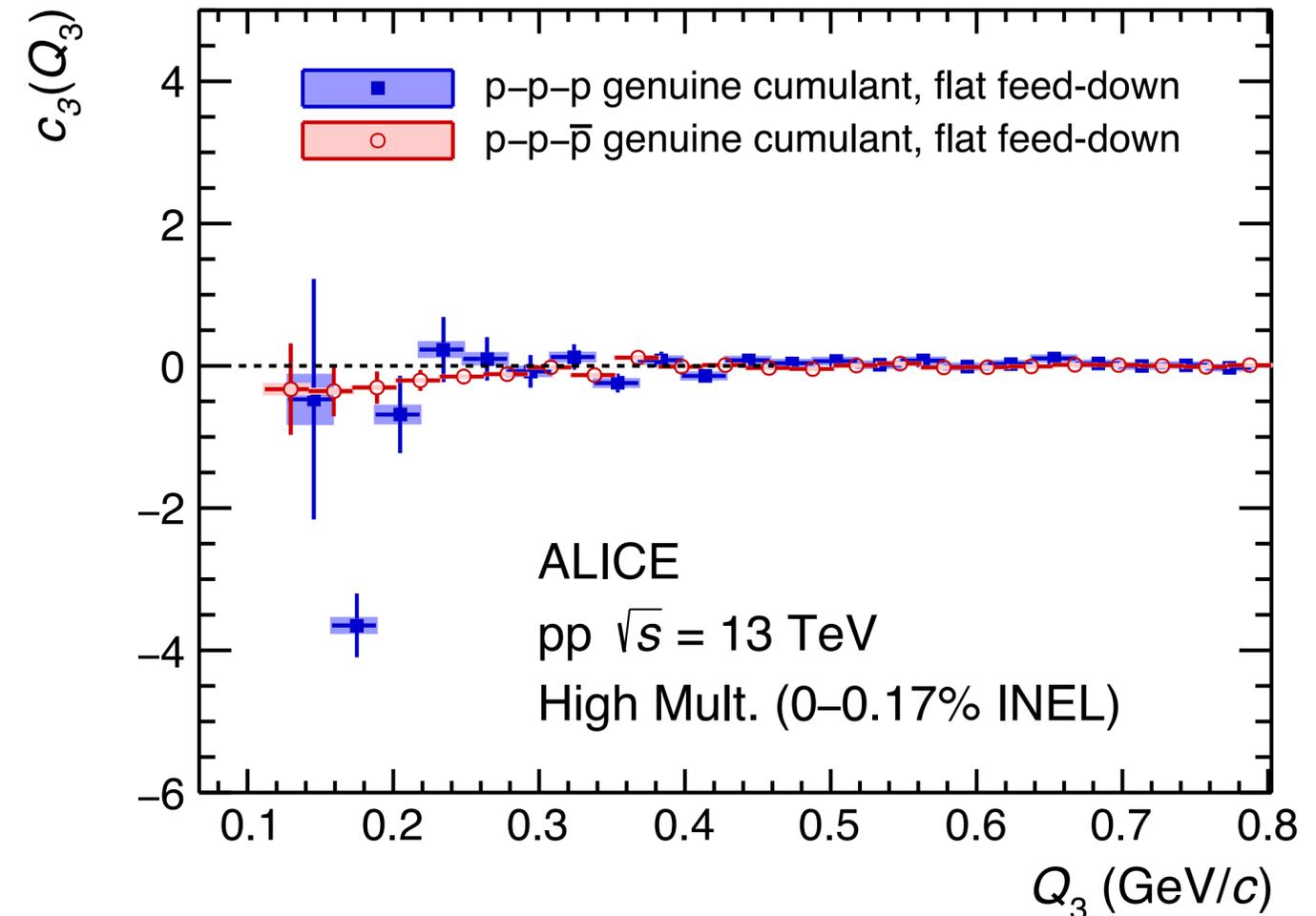
Possible forces at play:

- *Pauli blocking at the three-particle level*
- three-body strong interaction
- long-range Coulomb

Statistical significance:

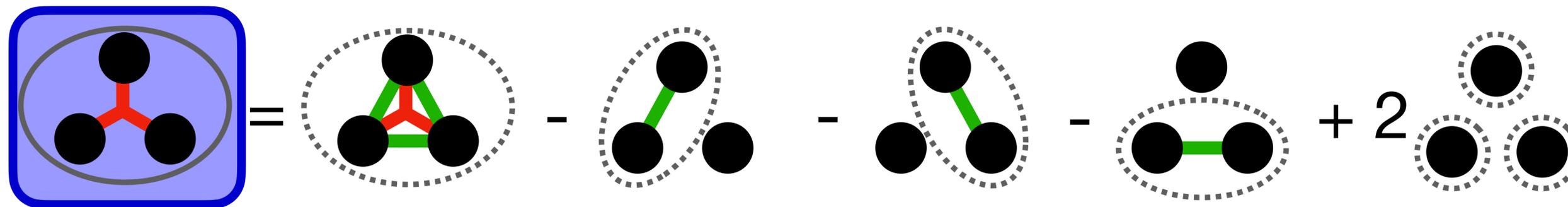
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Conclusion: significant deviation from null hypothesis; ongoing collaboration with A. Kivsky, L. Marcucci and M. Viviani (Pisa University - INFN) for the theoretical interpretation



Test with mixed-charge particles, cumulant negligible.

p-p-p cumulant



arXiv:2206.03344

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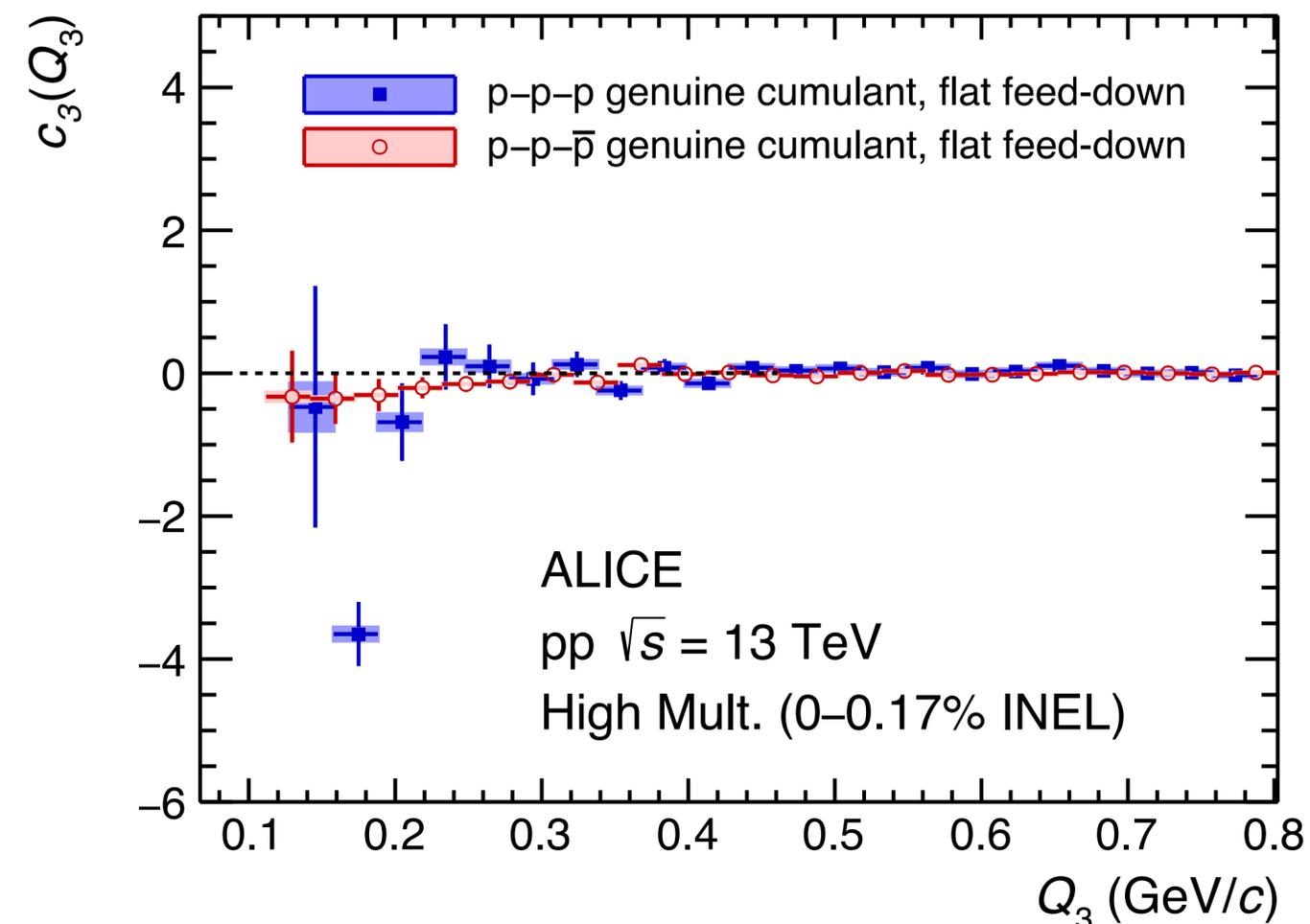
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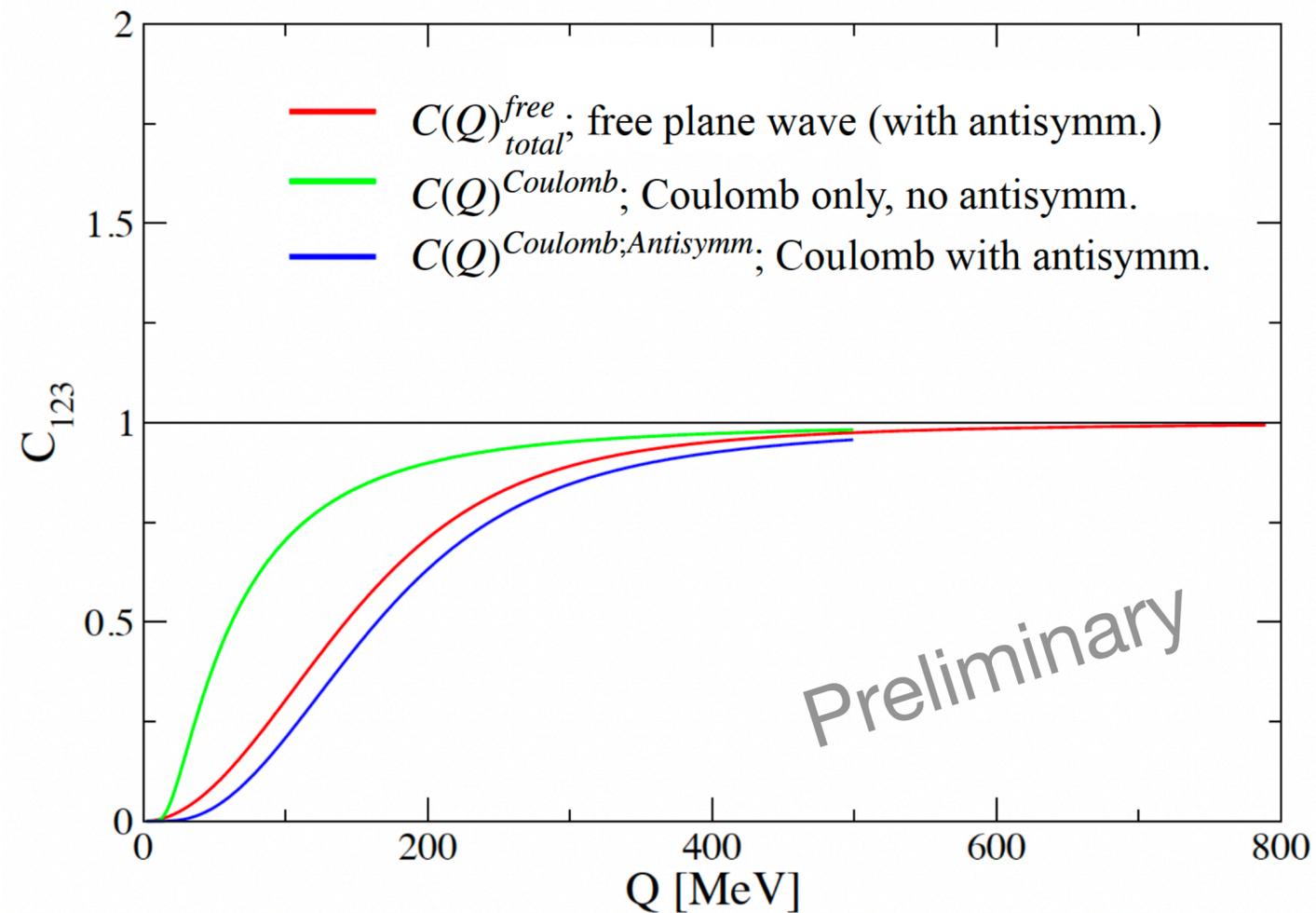
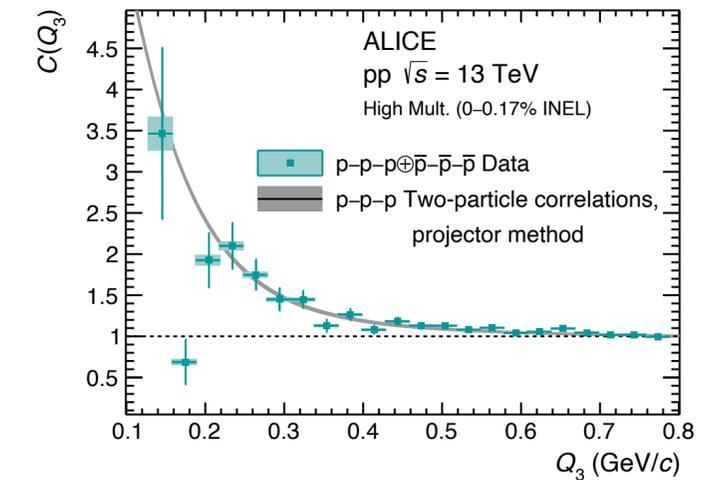
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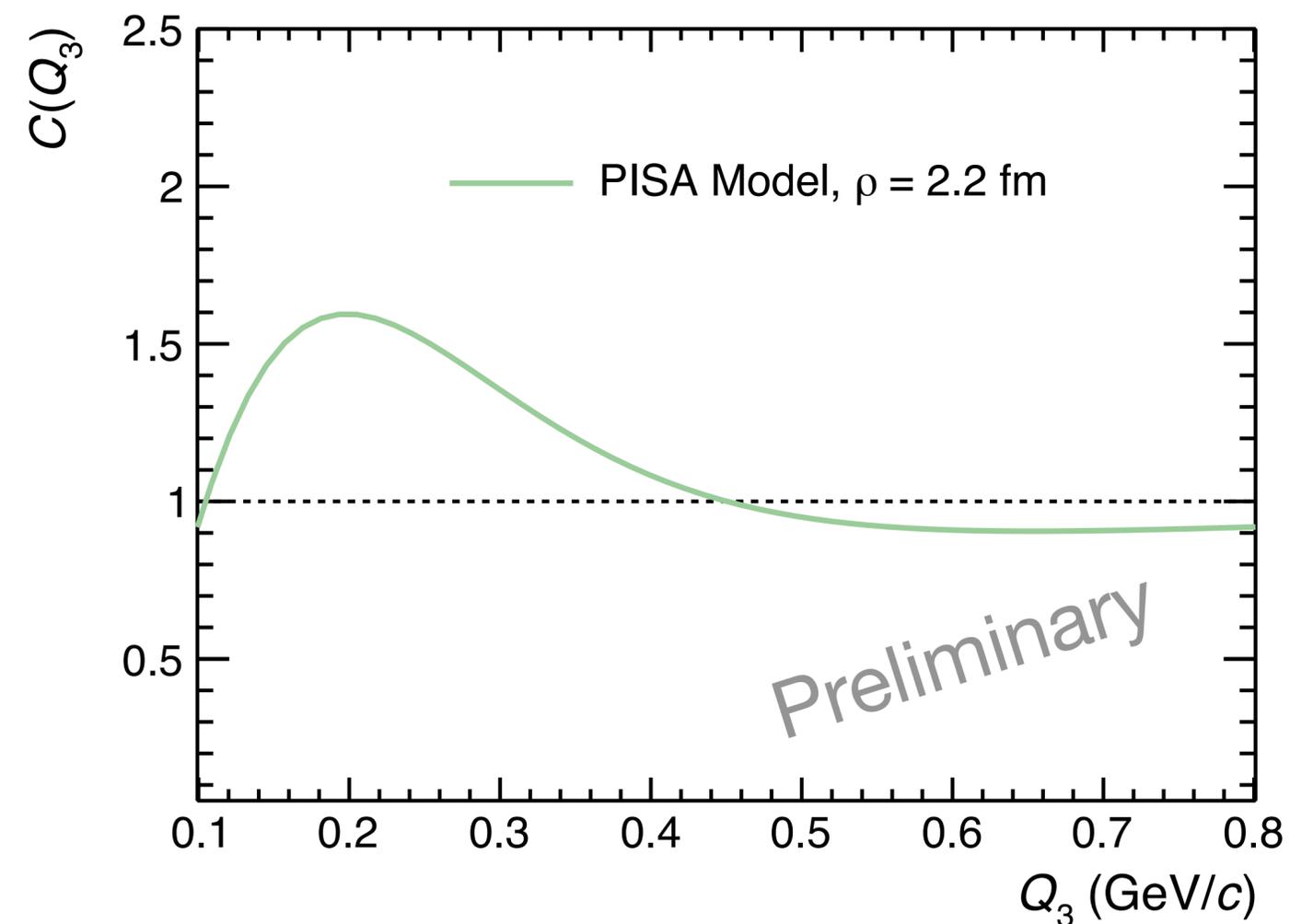
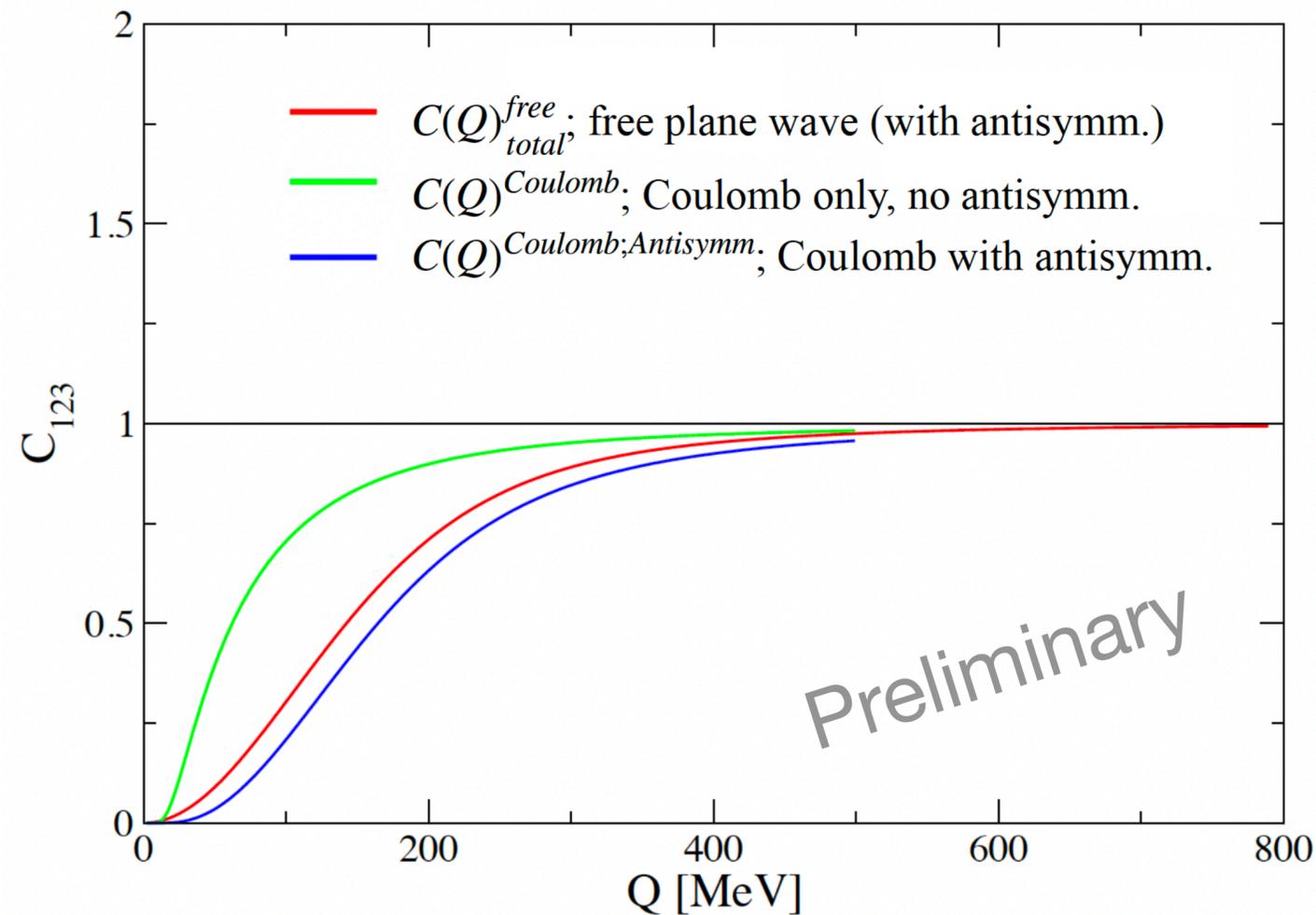
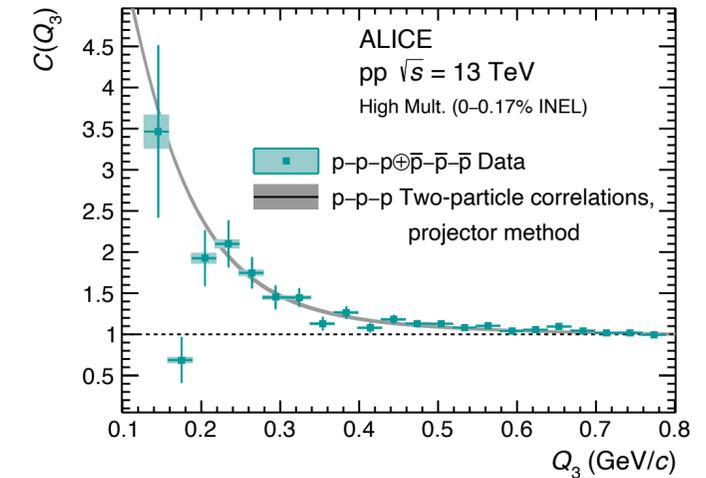
p-p-p calculations (work in progress)

- Calculations performed by Alejandro Kievsky (PISA group)
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 - only two-body strong interaction included
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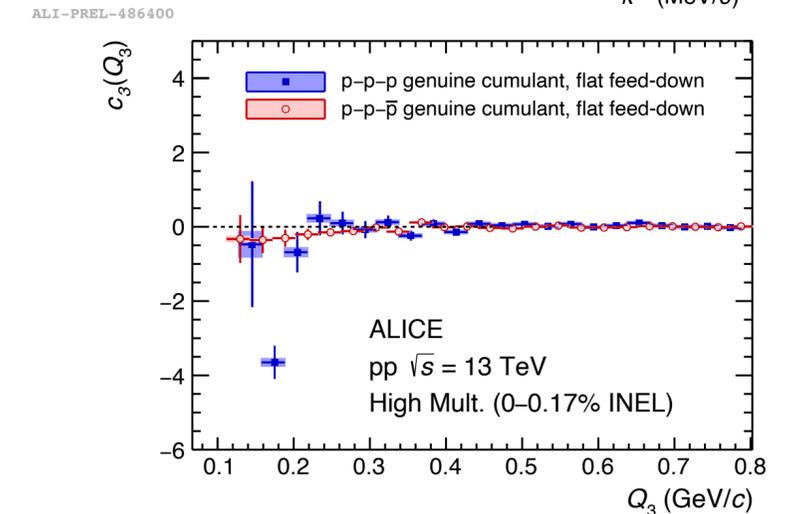
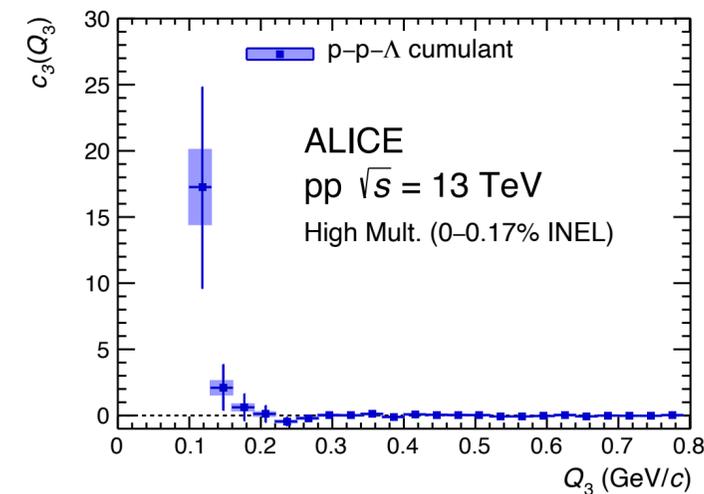
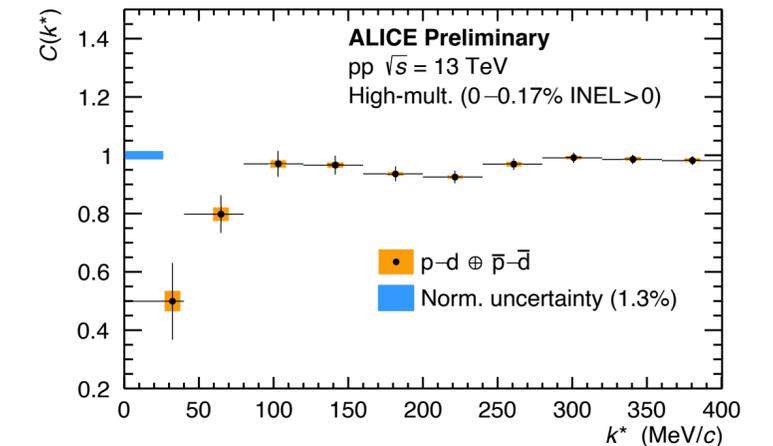
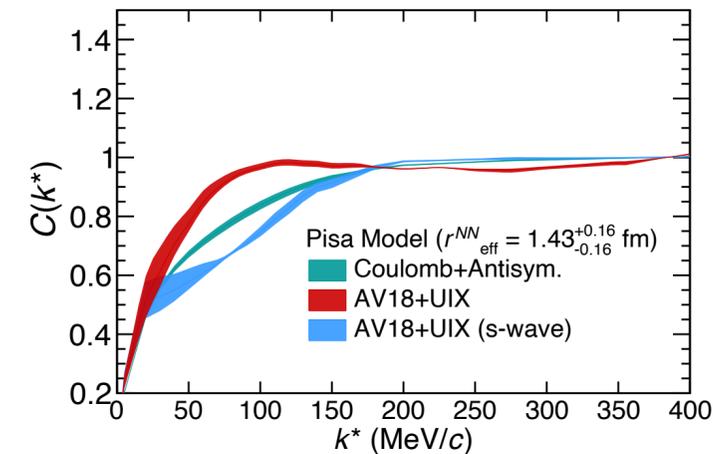
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Conclusions

First measurements tackling the problem of genuine three-body interactions using femtoscopy!

- **p-d**: can be described with full three-body calculations
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- **p-p-p**: negative cumulant with a significance of 6.7σ



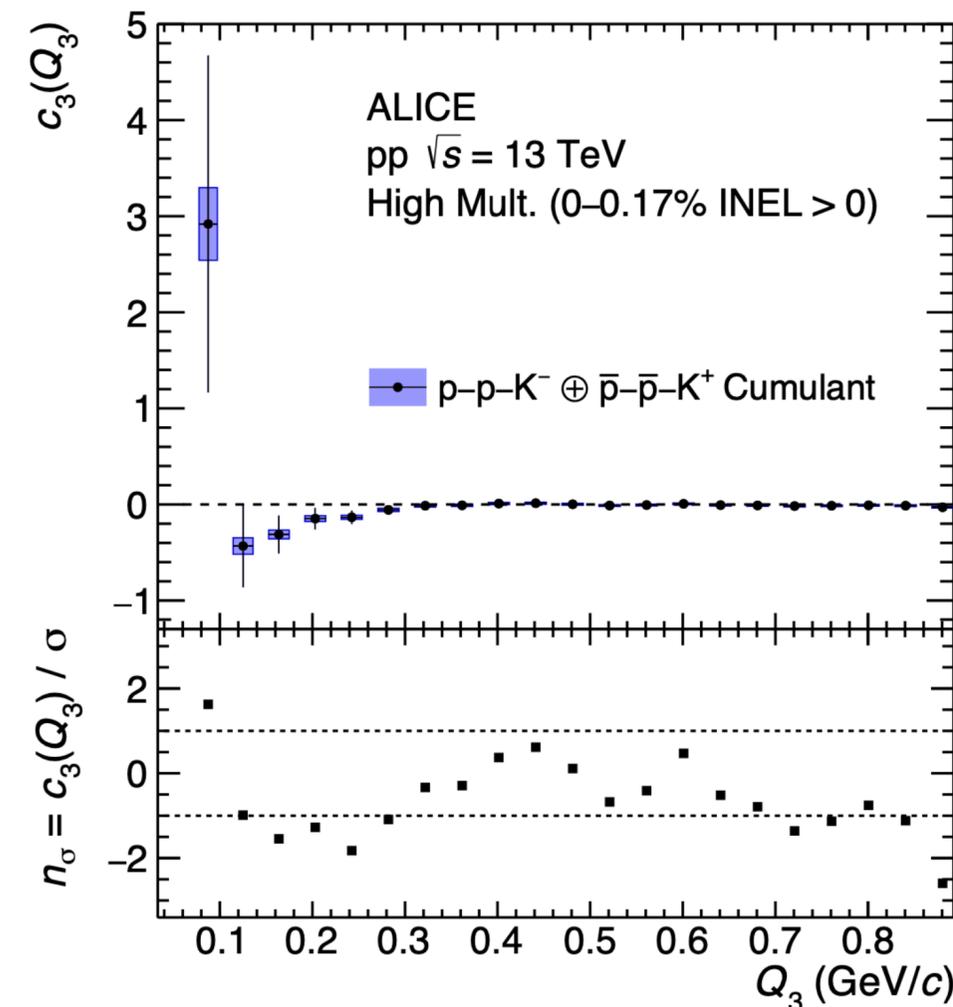
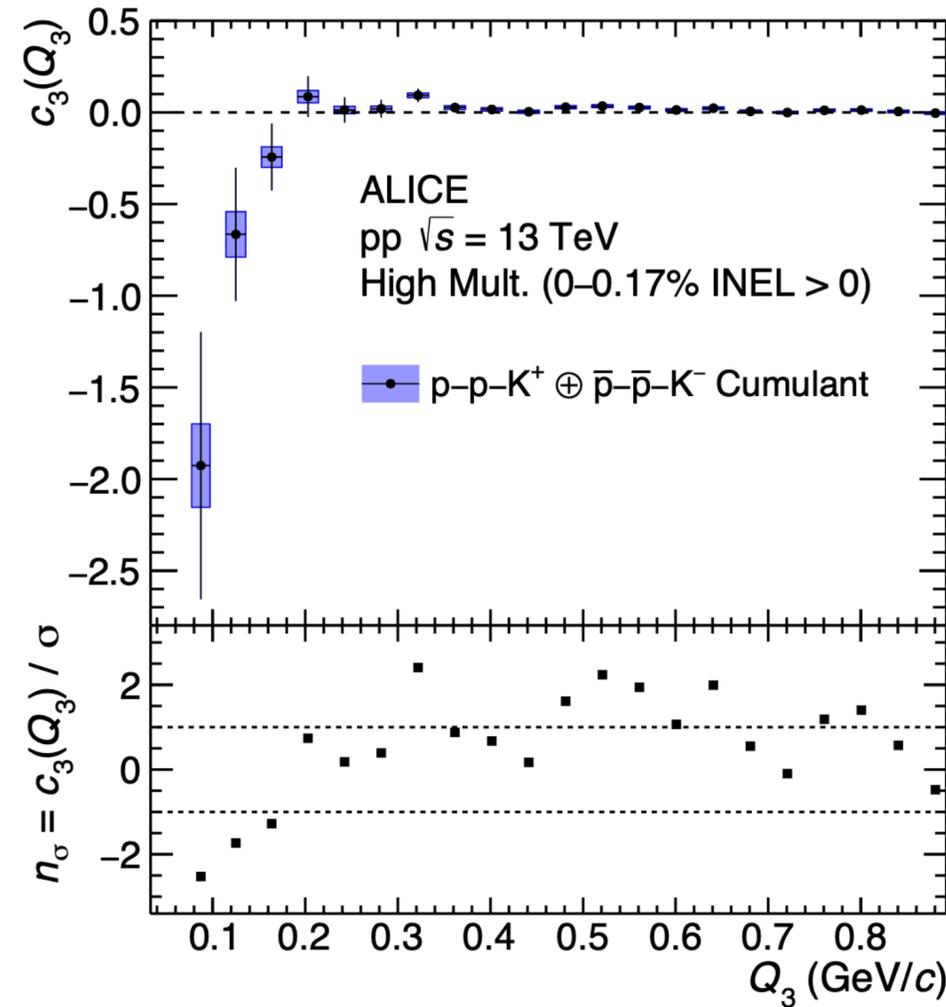
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New paper: [arXiv:2303.13448](https://arxiv.org/abs/2303.13448)

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Valentina Mantovani Sarti 5 Jun 2023, 14:30

Dimitar Mihaylov 5 Jun 2023, 17:40

Wioleta Rzeża 7 Jun 2023, 14:24

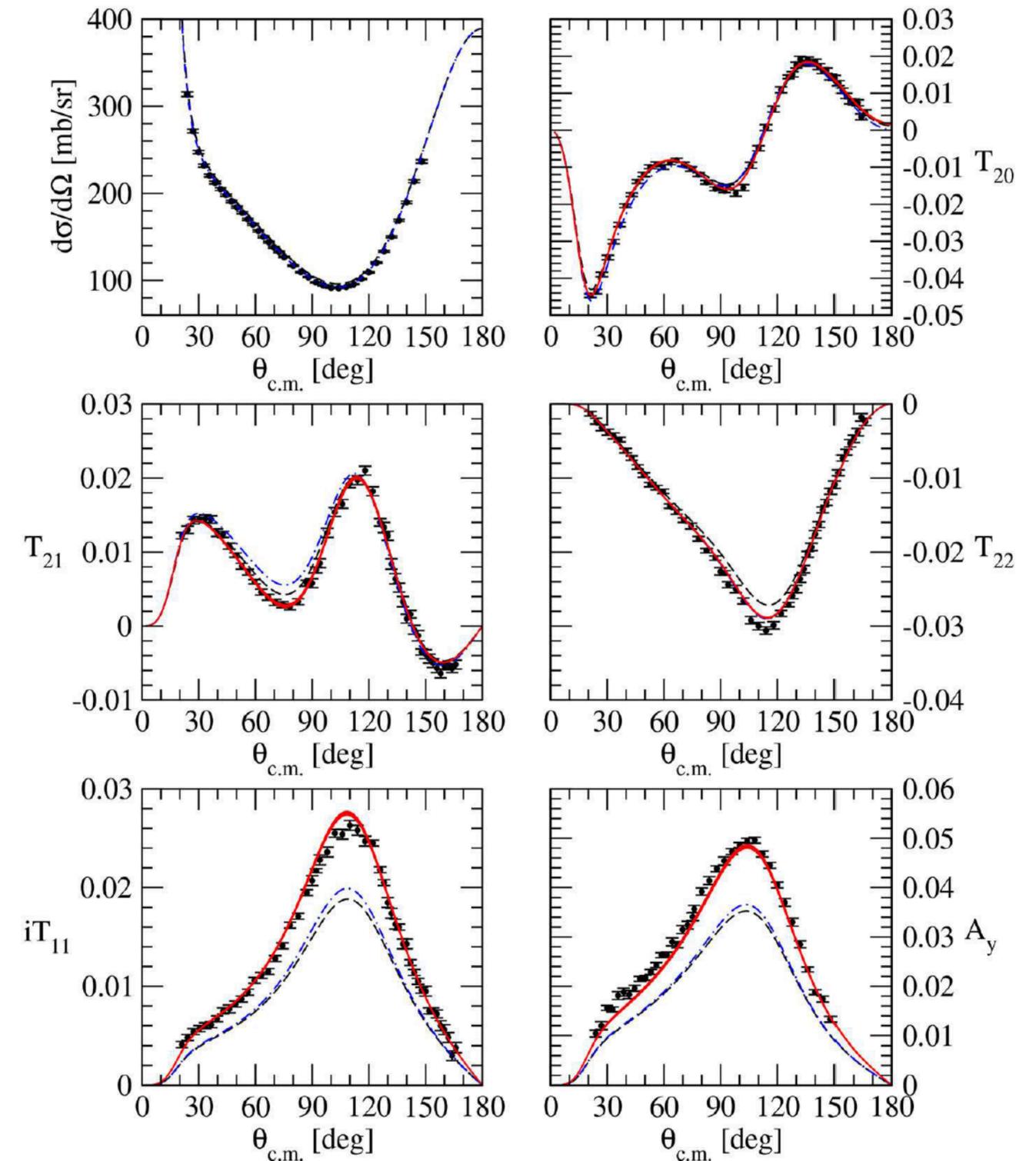
Marcel Lesch 8 Jun 2023, 15:12

Ramona Lea 8 Jun 2023, 15:42

Back-up

p-d scattering

- Three body interactions are required to reproduce scattering data

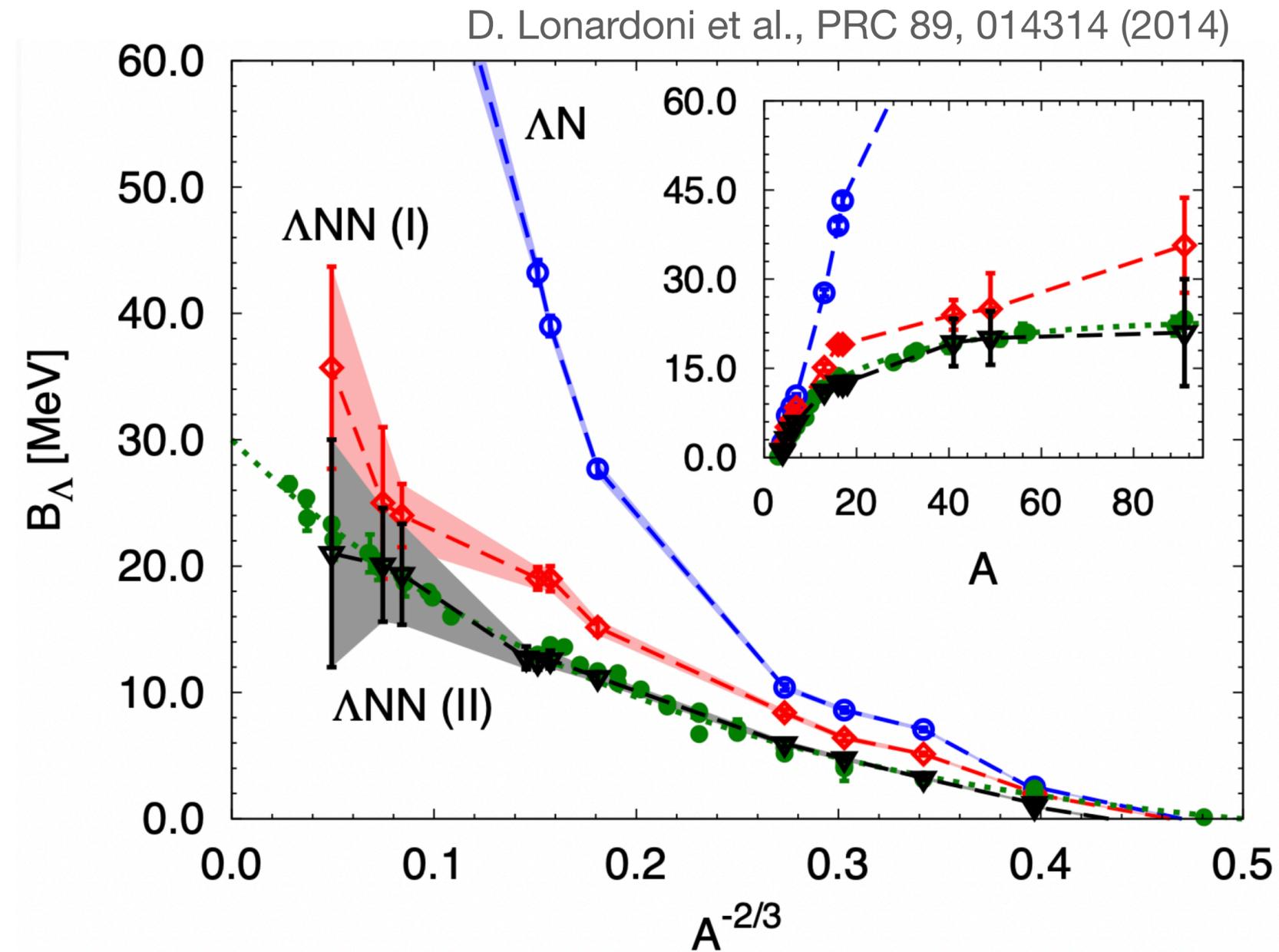


L.E. Marcucci et al., Front. Phys. 8, 69 (2020)

How to constrain three-body forces?

- Models are fitted to reproduce measured (hyper)nuclei properties
 - Access only to nuclear densities
 - Strongly dependent on the assumed two-body and many-body interactions
 - Different parametrisations of three-body forces describe better different nuclei

Parameters	System	B_{Λ}^{CSB}
Set (I)	${}^4_{\Lambda}\text{H}$	1.89(9)
	${}^4_{\Lambda}\text{He}$	2.13(8)
Set (II)	${}^4_{\Lambda}\text{H}$	0.95(9)
	${}^4_{\Lambda}\text{He}$	1.22(9)
Expt. [12]	${}^4_{\Lambda}\text{H}$	2.04(4)
	${}^4_{\Lambda}\text{He}$	2.39(3)



ALICE detector

- Excellent tracking and particle identification (PID) capabilities
- Most suitable detector at the LHC to study (anti-)nuclei production and annihilation

Inner Tracking System

Tracking, vertex, PID (dE/dx)

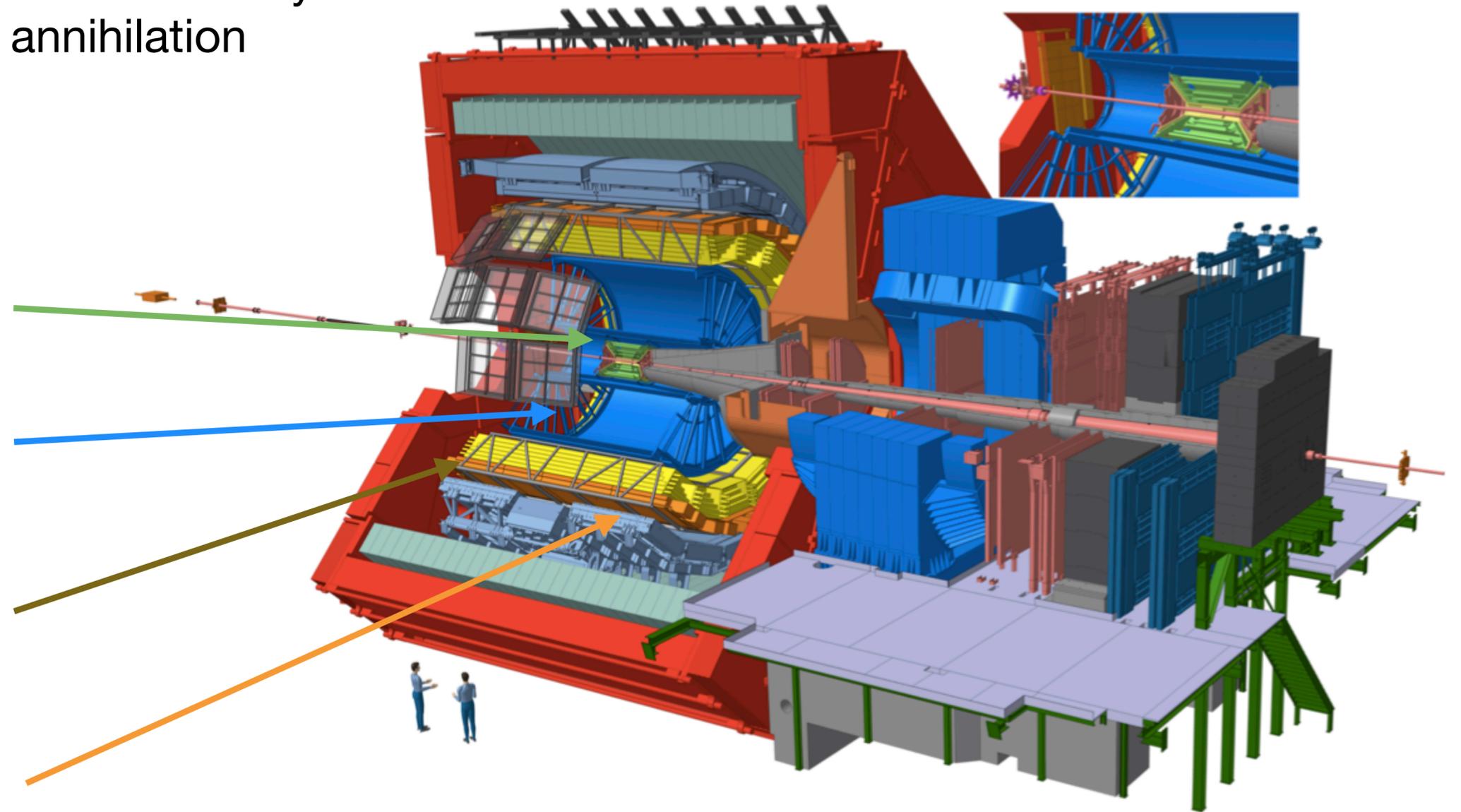
Time Projection Chamber

Tracking, PID (dE/dx)

Transition Radiation Detector

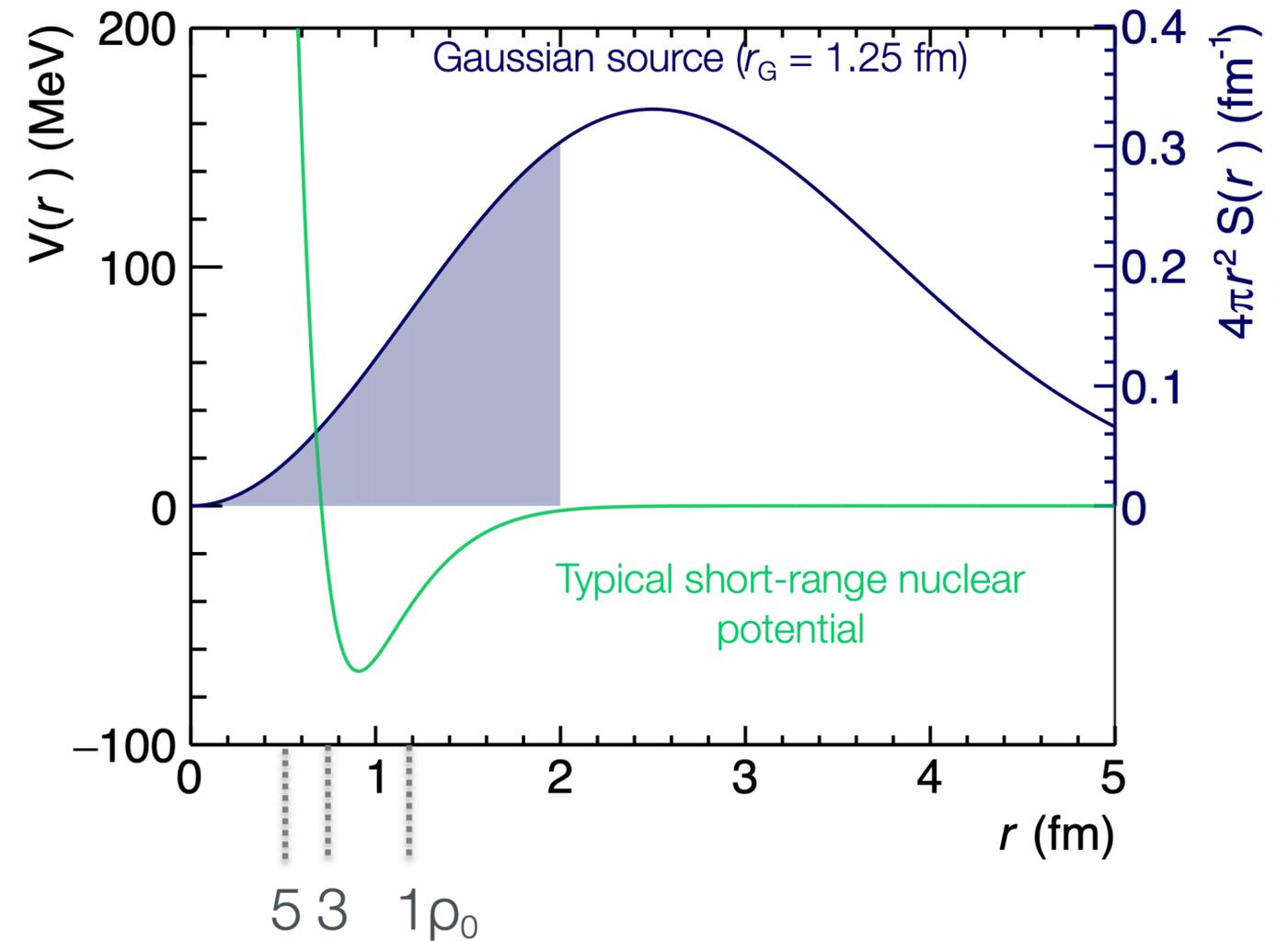
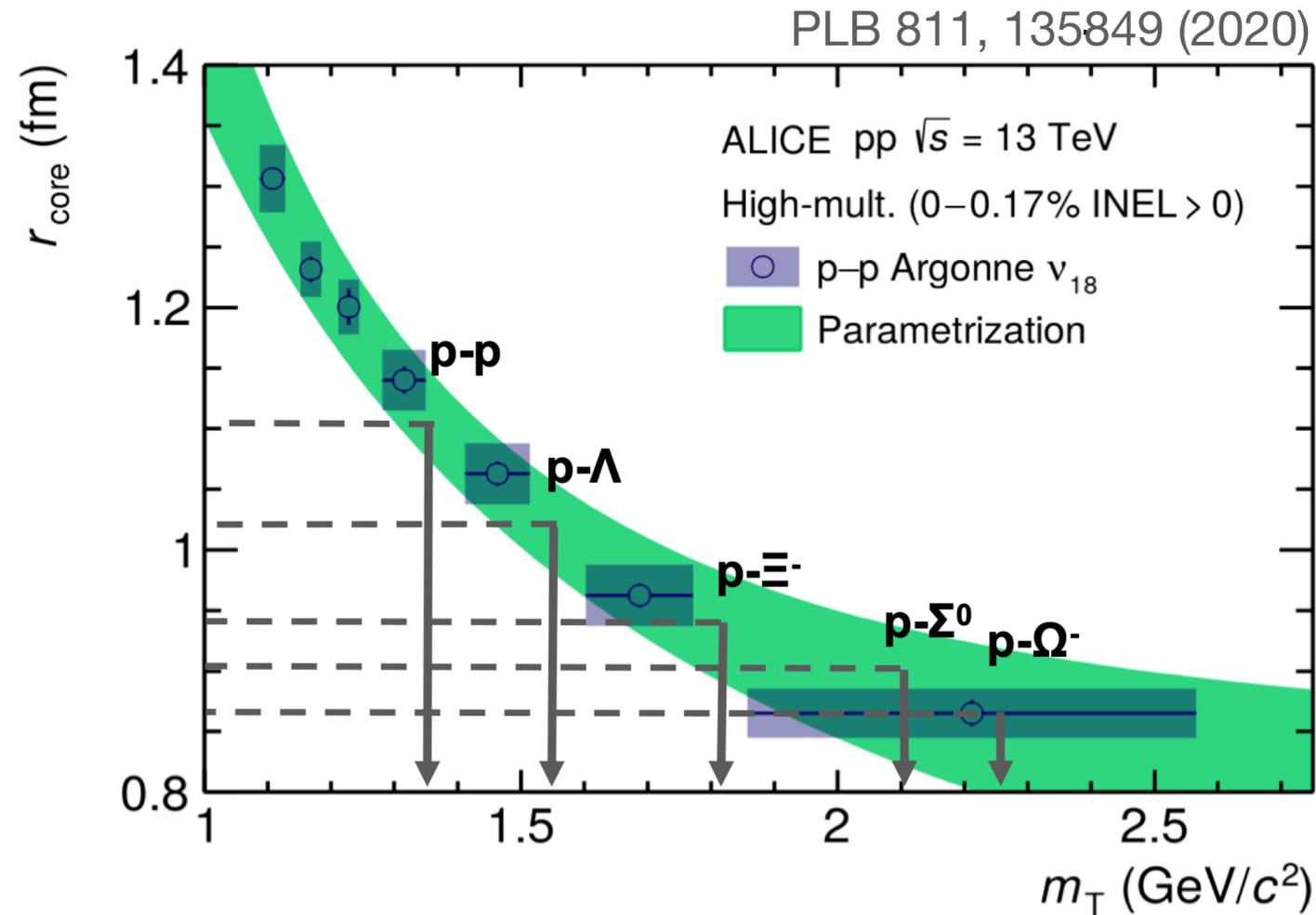
Time Of Flight detector

PID (TOF measurement)



Emission source

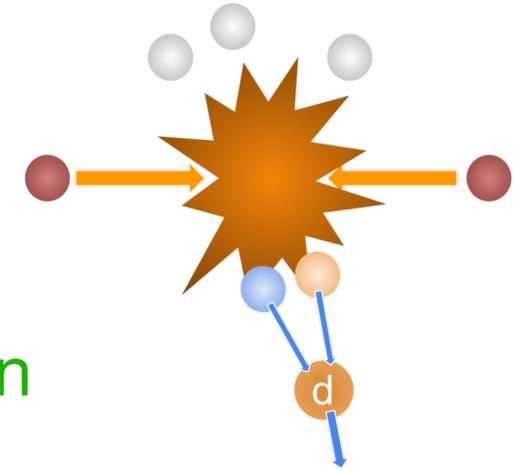
- Two main contributions:
 - general: Collective effects result in Gaussian core
 - specific: Decaying resonances require source correction



How to access three-body systems?

Three-body dynamics

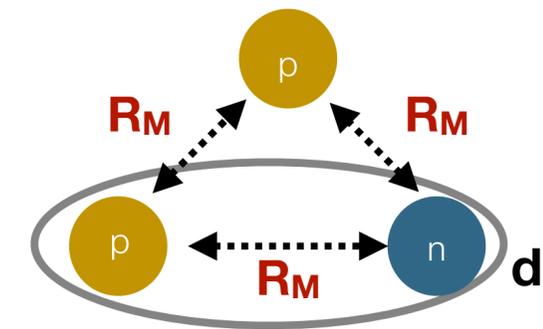
- Start with p-p-n state:
 - single-particle Gaussian emission source
 - three-nucleon wave function asymptotically behaves as p-d state
 - account for the probability to form deuteron employing deuteron wave function



$$C_{pd}(k) = \frac{1}{A_d} \frac{1}{6} \sum_{m_2, m_1} \int d^3 r_1 d^3 r_2 d^3 r_3 S_1(r_1) S_1(r_2) S_1(r_3) \left| \Psi_{m_2, m_1} \right|^2$$

- Rewritten as a function of the known source size R_M constrained by p-p

$$C_{pd}(k) = \frac{1}{A_d} \frac{1}{6} \sum_{m_2, m_1} \int \rho^5 d\rho d\Omega \frac{e^{-\rho^2/4R_M^2}}{(4\pi R_M^2)^3} \left| \Psi_{m_2, m_1} \right|^2$$



Proton-deuteron wave function

The three body wave function with proper treatment of 2N and 3N interaction at very short distances goes to a p-d state.

- Three-body wavefunction for p-d: $\Psi_{m_2, m_1}(x, y)$ describing three-body dynamics, anchored to p-d scattering observables.
 - x = distance of p-n system within the deuteron
 - y = p-d distance
 - m_2 and m_1 deuteron and proton spin
- $\Psi_{m_2, m_1}(x, y)$ three-nucleon wave function asymptotically behaves as p-d state:

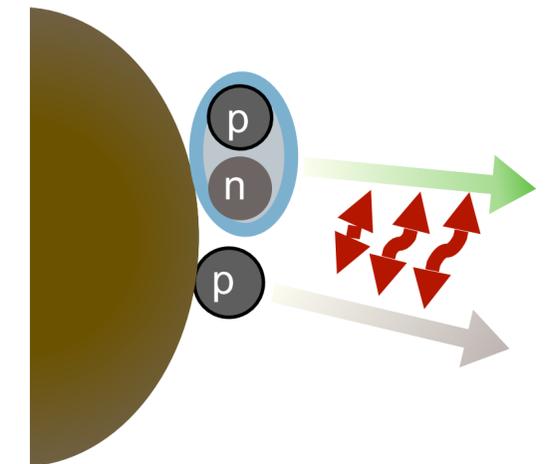
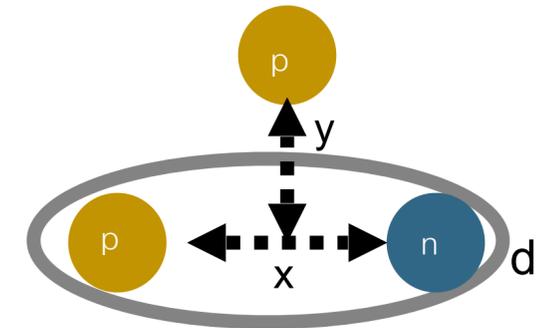
$$\Psi_{m_2, m_1}(x, y) = \underbrace{\Psi_{m_2, m_1}^{(\text{free})}}_{\text{Asymptotic form}} + \sum_{LSJ}^{J \leq \bar{J}} \underbrace{\sqrt{4\pi i^L} \sqrt{2L+1} e^{i\sigma_L} (1m_2 \frac{1}{2} m_1 |SJ_z)(LOSJ_z | JJ_z)}_{\text{Strong three-body interaction}} \tilde{\Psi}_{LSJJ_z}.$$

Asymptotic form

Strong three-body interaction

→ $\tilde{\Psi}_{LSJJ_z}$ describe the configurations where the three particles are close to each other

→ $\Psi_{m_1, m_2}^{(\text{free})}$ an asymptotic form of p-d wave function



Kievsky et al, Phys. Rev. C 64 (2001) 024002
 Kievsky et al, Phys. Rev. C 69 (2004) 014002
 Deltuva et al, Phys. Rev. C 71 (2005) 064003

Proton-deuteron correlations

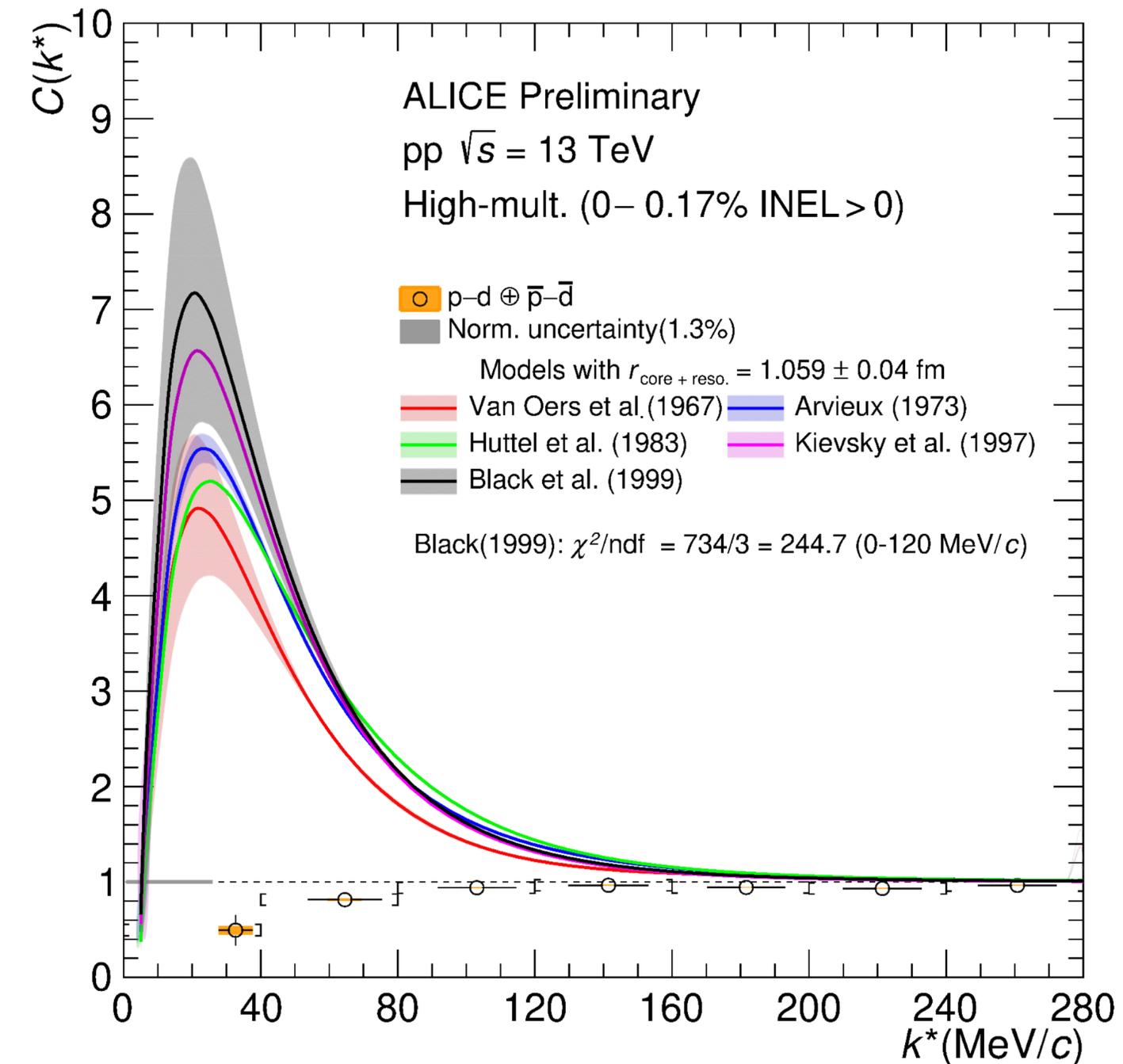
Point-like particle models anchored to scattering experiments

	$S = 1/2$		$S = 3/2$	
	$f_0(\text{fm})$	$d_0(\text{fm})$	$f_0(\text{fm})$	$d_0(\text{fm})$
Van Oers et al. (1967)	$-1.30^{+0.20}_{-0.20}$	—	$-11.40^{+1.20}_{-1.80}$	$2.05^{+0.25}_{-0.25}$
Arvieux (1973)	$-2.73^{+0.10}_{-0.10}$	$2.27^{+0.12}_{-0.12}$	$-11.88^{+0.10}_{-0.40}$	$2.63^{+0.01}_{-0.02}$
Huttel et al. (1983)	-4.0	—	-11.1	—
Kievsky et al. (1997)	-0.024	—	-13.7	—
Black et al. (1999)	$0.13^{+0.04}_{-0.04}$	—	$-14.70^{+2.30}_{-2.30}$	—

W. T. H. Van Oers, & K. W. Brockman Jr, NPA 561 (1967);
 J. Arvieux et al., NPA 221 (1973); E. Huttel et al., NPA 406 (1983);
 A. Kievsky et al., PLB 406 (1997); T. C. Black et al., PLB 471 (1999);

- Coulomb + strong interaction using the Lednický model
Lednický, R. Phys. Part. Nuclei 40, 307–352 (2009)
- Only s-wave interaction
- Source radius evaluated using the hadron-hadron universal m_T scaling

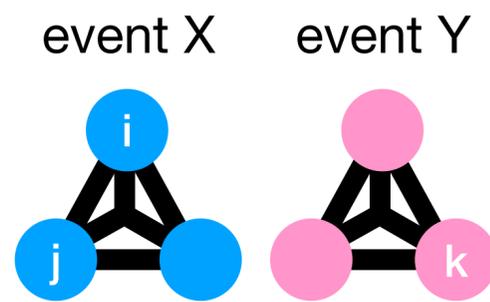
Point-like particle description doesn't work for p-d



ALI-PREL-501009

Data-driven method

- Use event mixing
- Two particles from the same event and one particle from another:



$$C_{ij} \left(\left[\mathbf{p}_i, \mathbf{p}_j \right], \mathbf{p}_k \right) = \frac{N_2 \left(\mathbf{p}_i, \mathbf{p}_j \right) N_1 \left(\mathbf{p}_k \right)}{N_1 \left(\mathbf{p}_i \right) N_1 \left(\mathbf{p}_j \right) N_1 \left(\mathbf{p}_k \right)}$$

- Calculate Lorentz-invariant scalar Q_3 for every triplet $\mathbf{p}_i, \mathbf{p}_j, \mathbf{p}_k$ to obtain $C_{ij}(Q_3)$

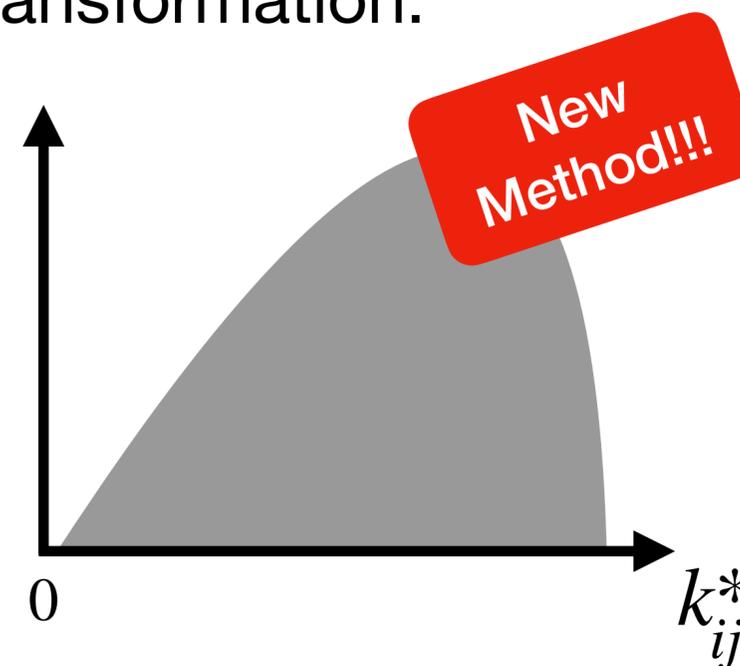
Projector method

- Use two-particle measured or theoretical correlation function $C([\mathbf{p}_i, \mathbf{p}_j])$
- Perform kinematic transformation:

$$C_2 \left(k_{ij}^* \right) \rightarrow C_{ij} \left(Q_3 \right)$$

$$k_{ij}^* \text{ (pair)} \rightarrow Q_3 \text{ (triplet)}$$

For one Q_3 value \rightarrow

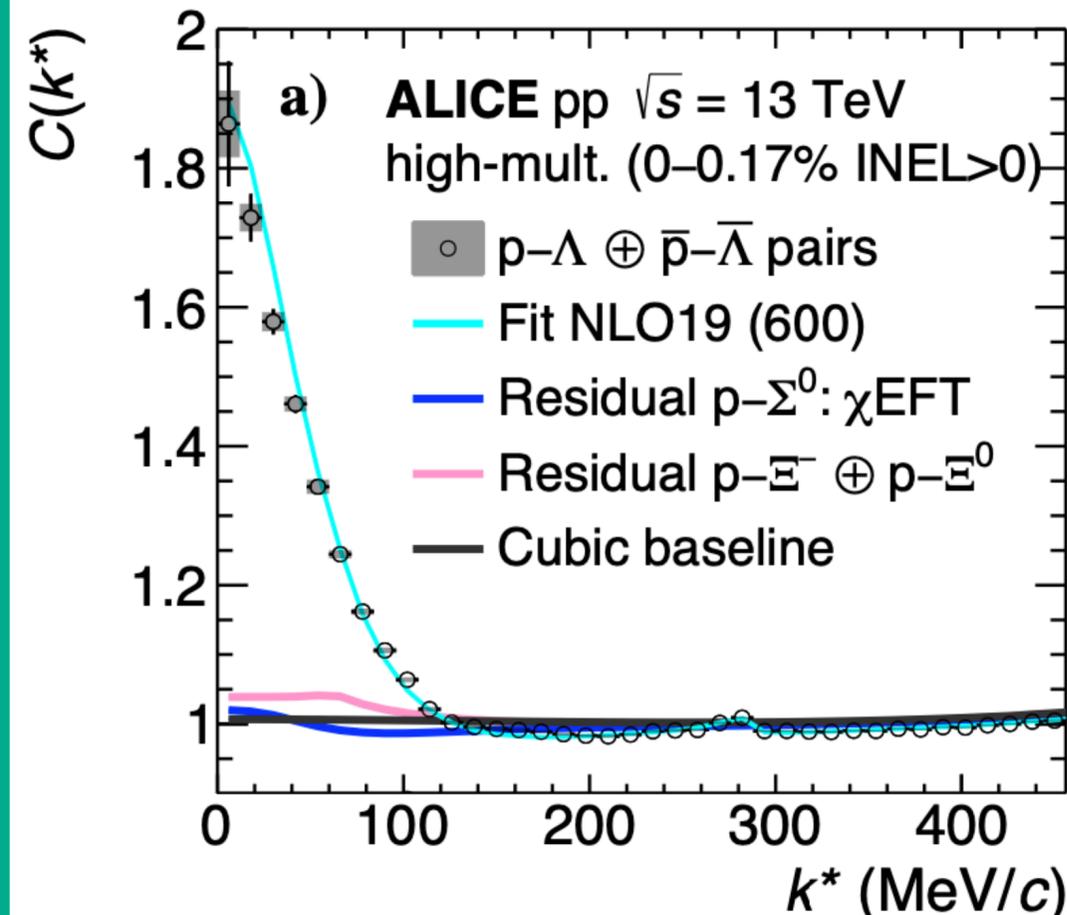
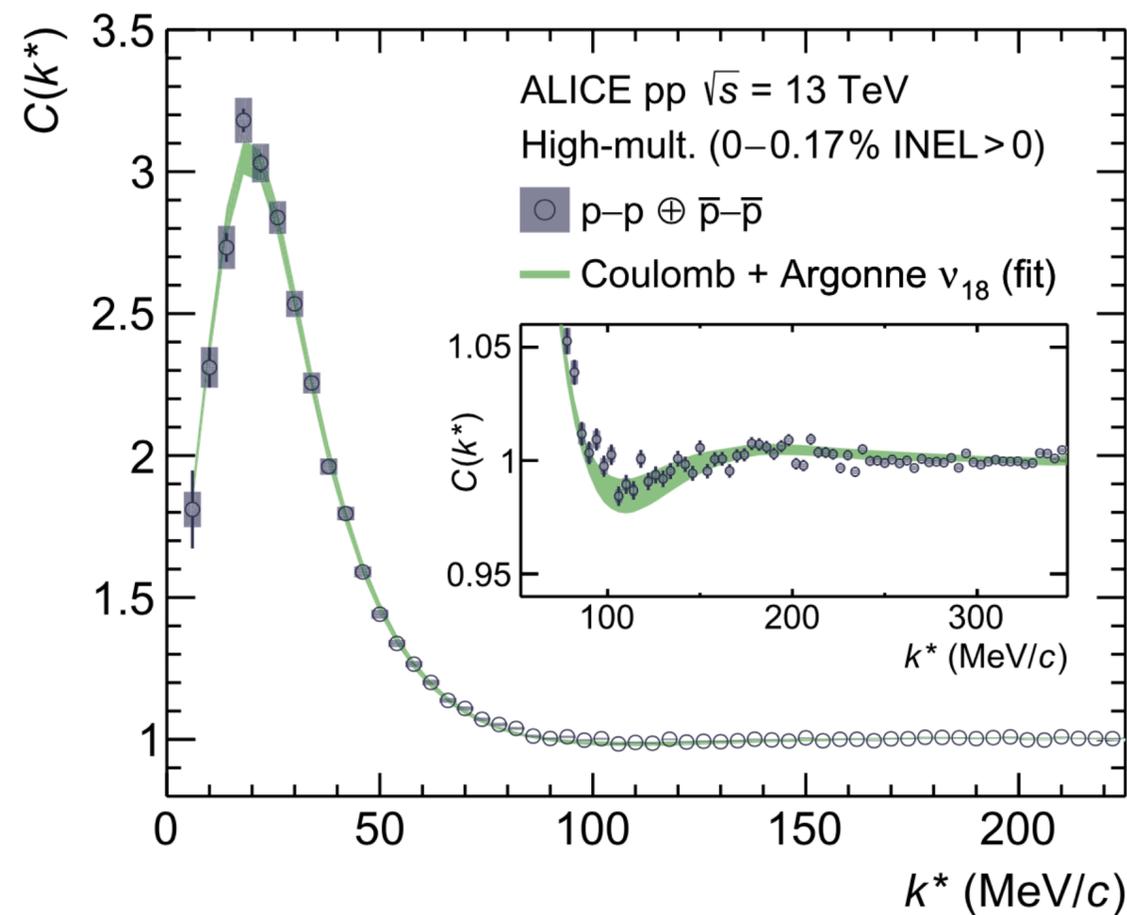
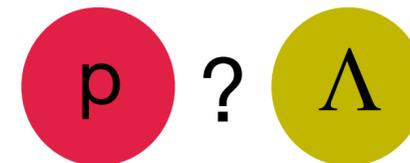
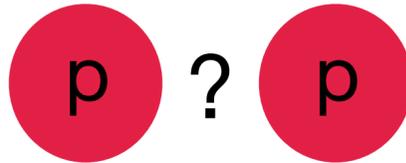


- To obtain the correlation function:

$$C_{ij}(Q_3) = \int C(k_{ij}^*) W_{ij}(k_{ij}^*, Q_3) dk_{ij}^*$$

Two-body measurements

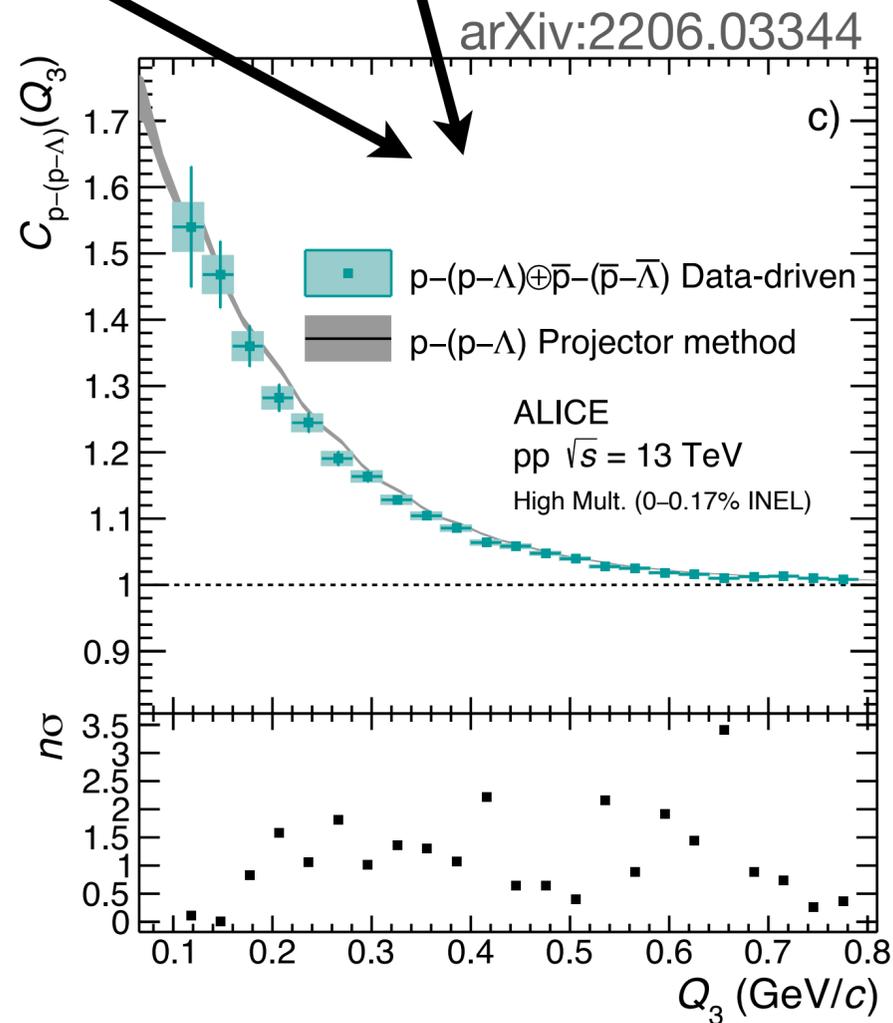
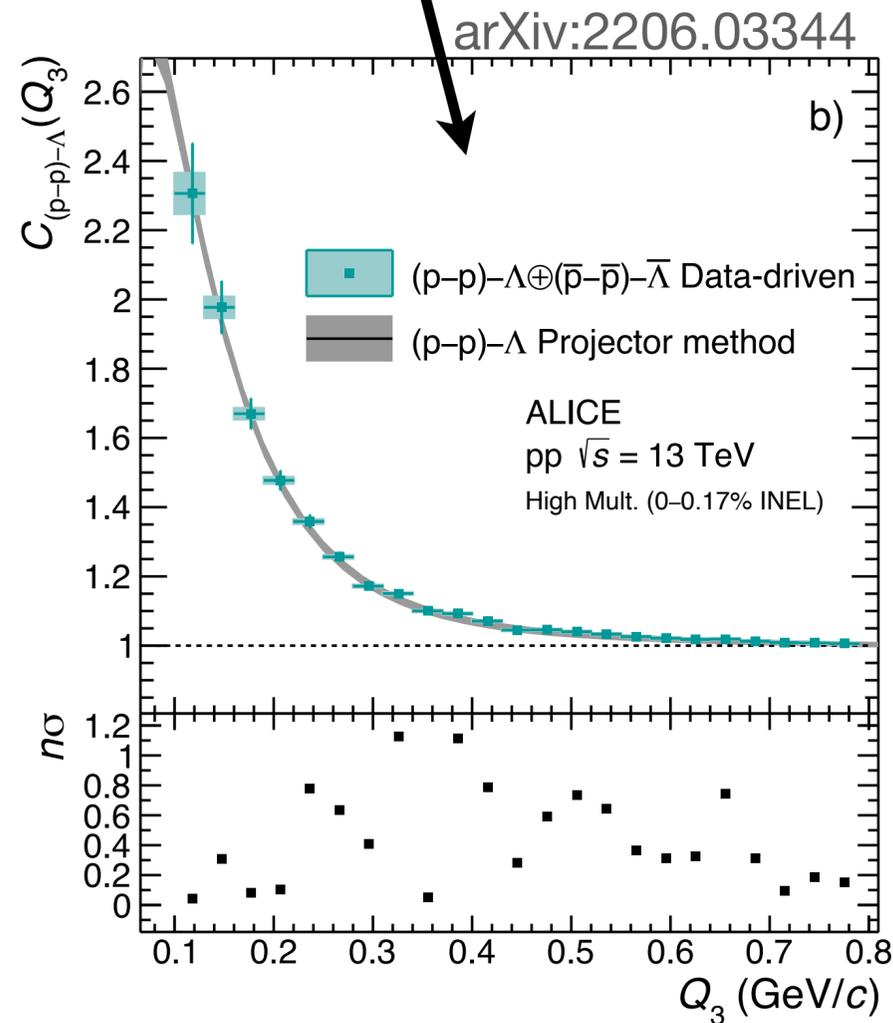
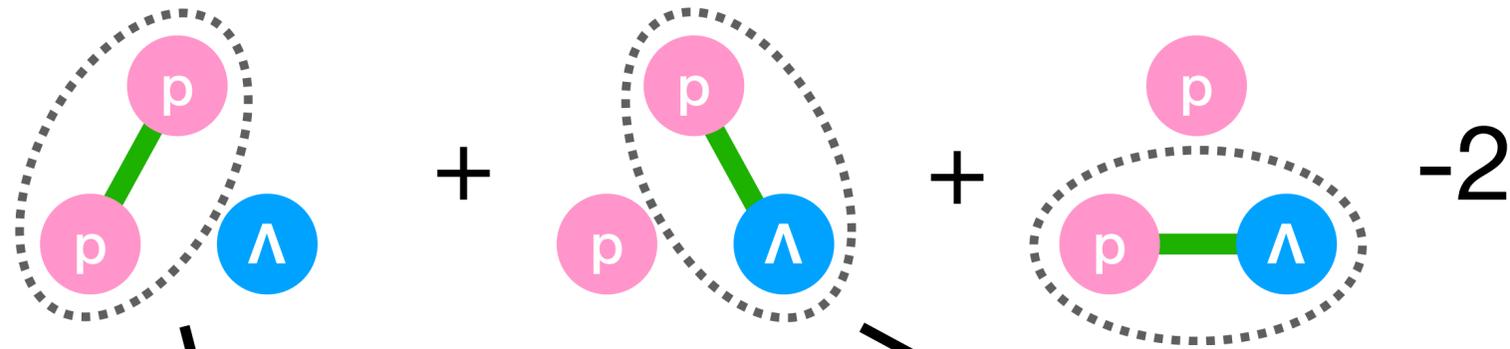
- Many different two-body interactions measured successfully!



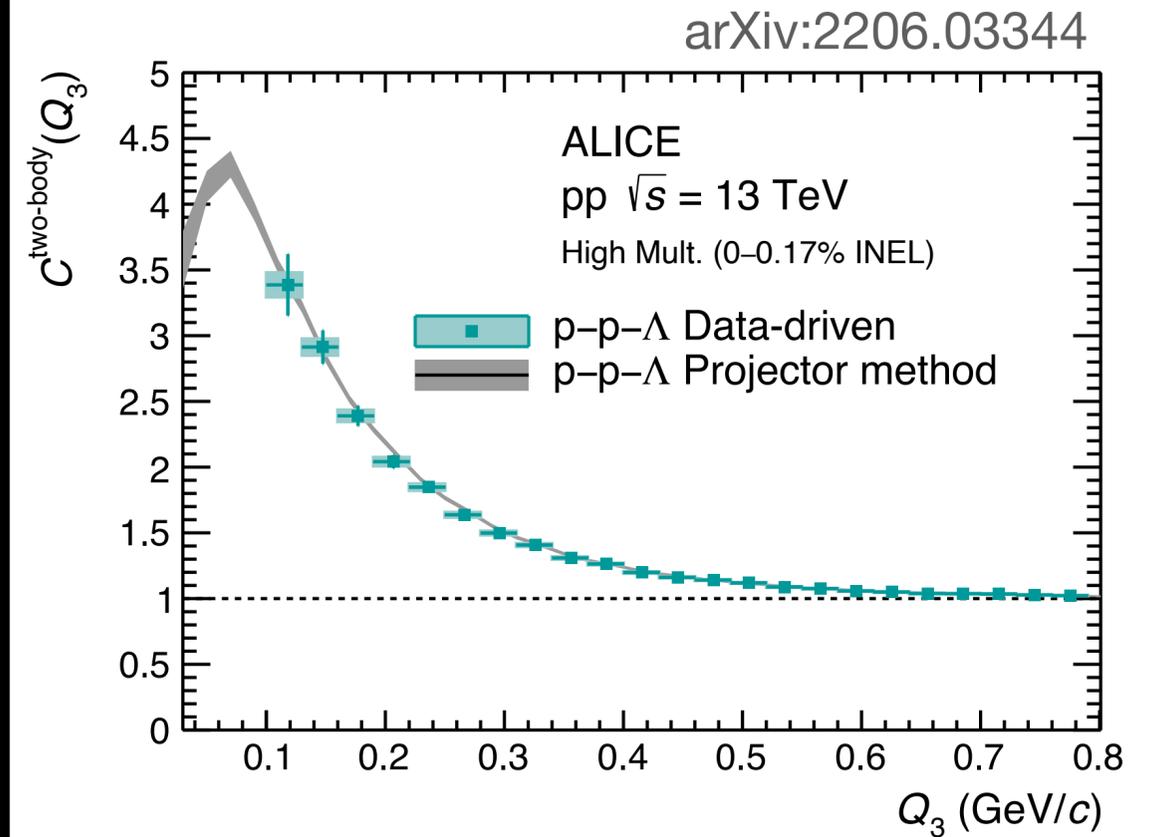
TUM Group:
EPJC 78 (2018) 394
arXiv:2107.10227

ALICE:
PRC 99 (2019) 024001
PLB 797 (2019) 134822
PRL 123 (2019) 112002
PRL 124 (2020) 09230
PLB 805 (2020) 135419
PLB 811 (2020) 135849
Nature 588 (2020) 232-238
[arXiv:2104.04427](https://arxiv.org/abs/2104.04427)
[arXiv:2105.05578](https://arxiv.org/abs/2105.05578)
[arXiv:2105.05683](https://arxiv.org/abs/2105.05683)
[arXiv:2105.05190](https://arxiv.org/abs/2105.05190)

Lower-order contributions: p-p- Λ



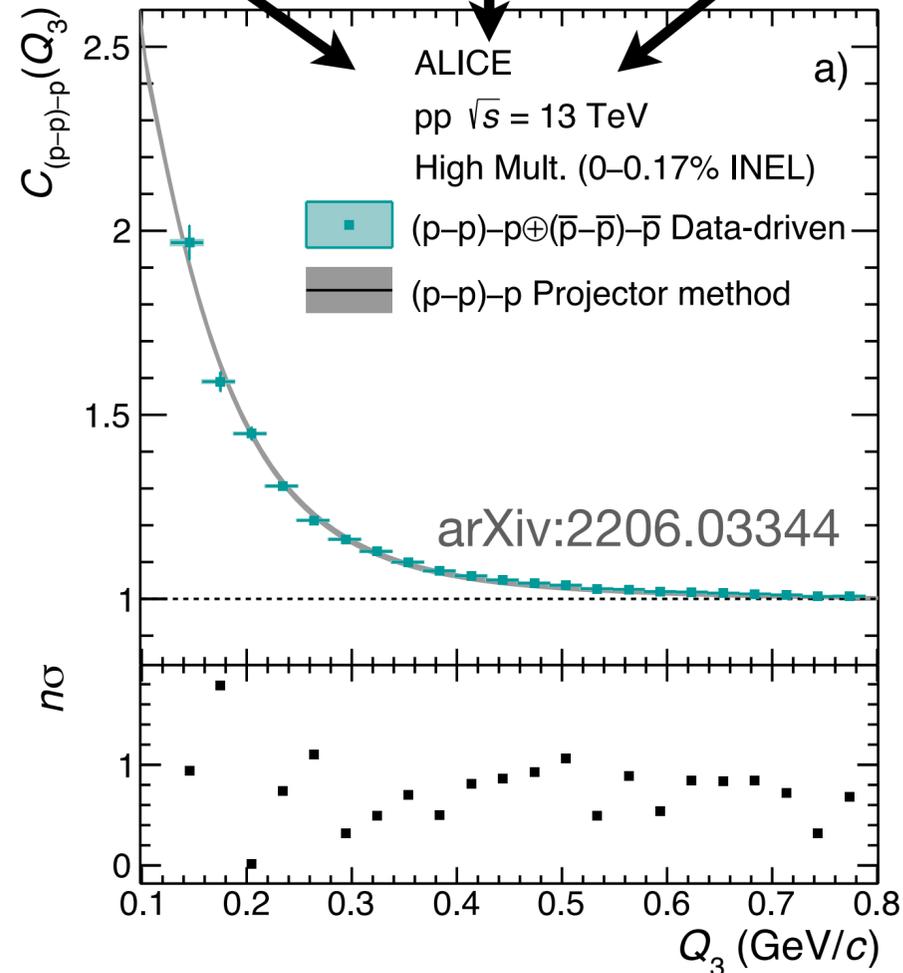
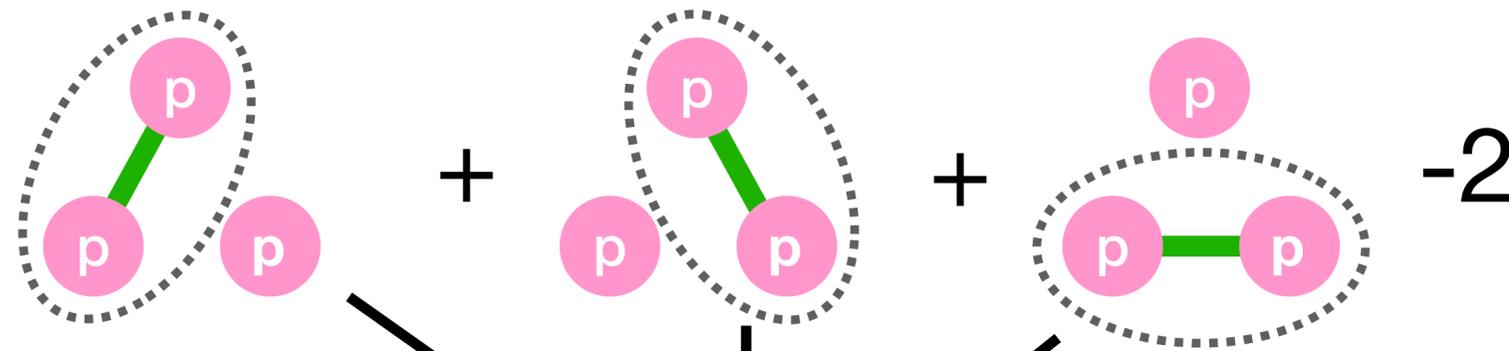
Total lower-order contributions



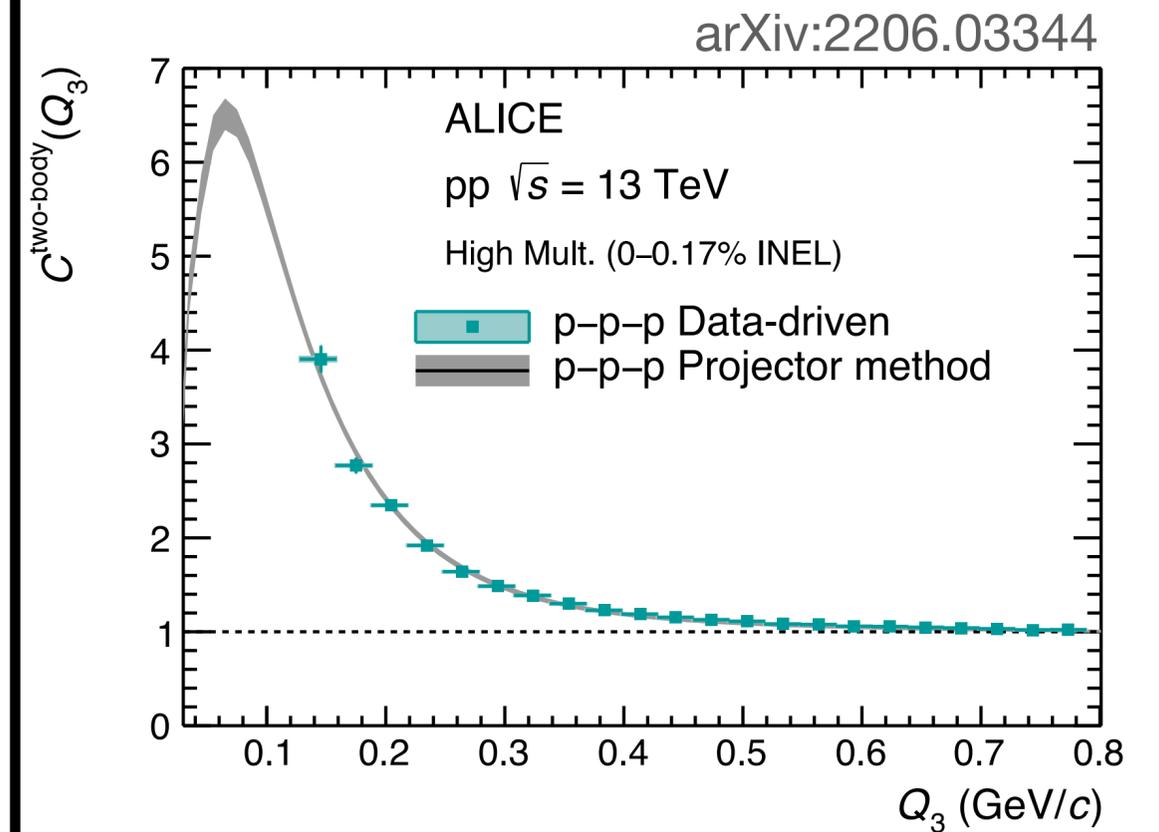
Already measured p-p [1] and p- Λ [2] correlation functions used for projection

[1] PLB 805 (2020) 135419; [2] arXiv:2104.04427

Lower-order contributions: p-p-p



Total lower-order contributions



Already measured p-p [1] correlation function used for projection.

[1] PLB 805 (2020) 135419

Projector

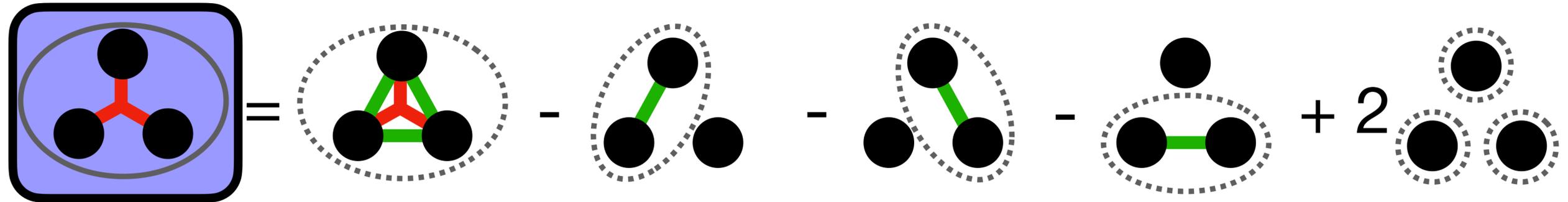
- Looking at 2-body correlation function in 3-body space requires to account for the phase-space of the particles.
- The projection onto Q_3 is performed by integrating the correlation function over all the configurations in the momentum phase space having the same value of Q_3

$$C(Q_3) = \iiint_{Q_3=\text{constant}} C([\mathbf{p}_i, \mathbf{p}_j], \mathbf{p}_k) d^3\mathbf{p}_i d^3\mathbf{p}_j d^3\mathbf{p}_k = \int C_2(k_{ij}^*) W_{ij}(k_{ij}^*, Q_3) dk_{ij}^*$$

$$W_{ij}(k_{ij}^*, Q_3) = \frac{16(\alpha\gamma - \beta^2)^{3/2} k_{ij}^{*2}}{\pi\gamma^2 Q_3^4} \sqrt{\gamma Q_3^2 - (\alpha\gamma - \beta^2) k_{ij}^{*2}}$$

- The α, β, γ depend only on the masses of the three particles.

p-p-K⁺ cumulant

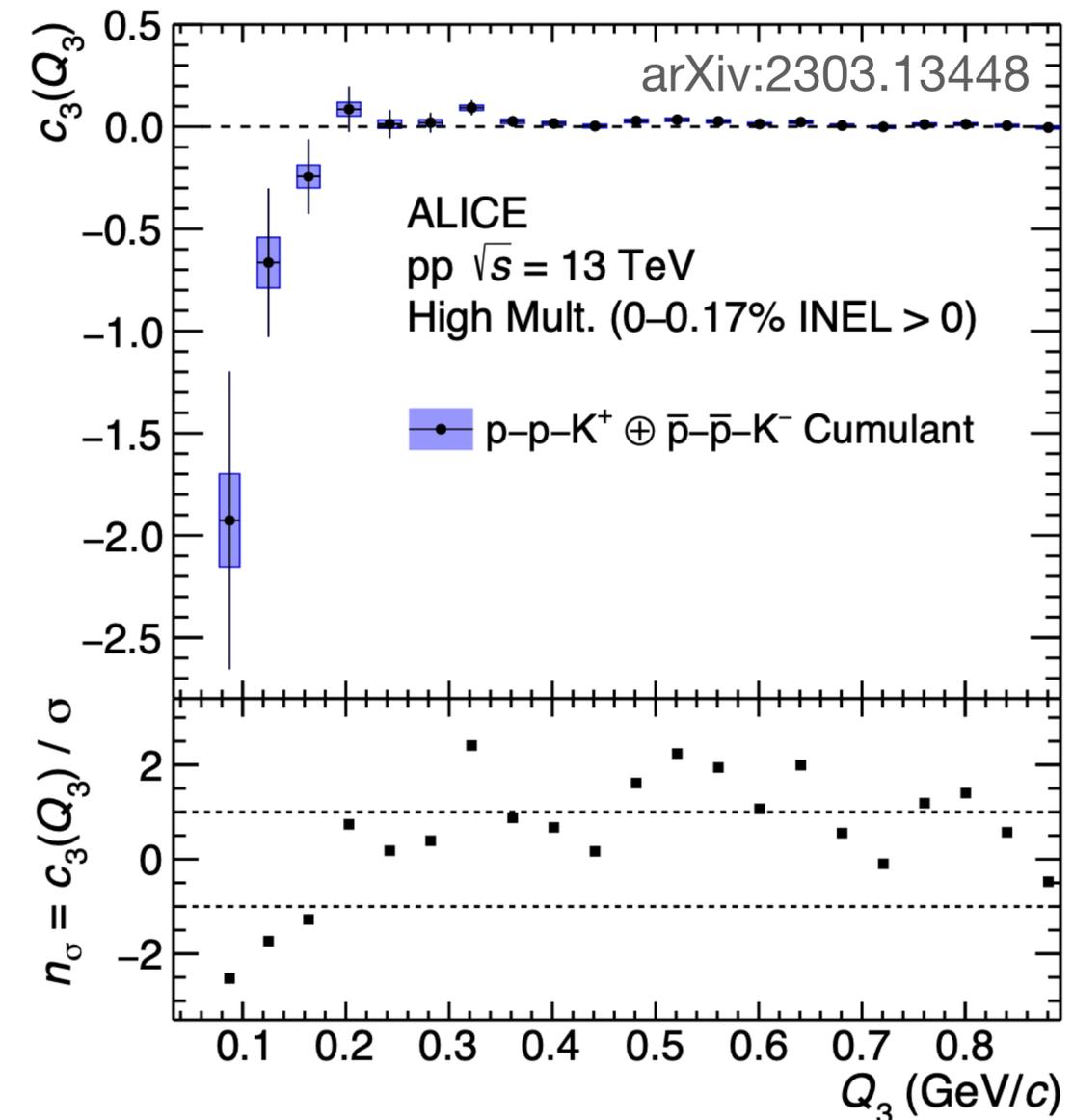


Hint of a negative cumulant for p-p-K⁺

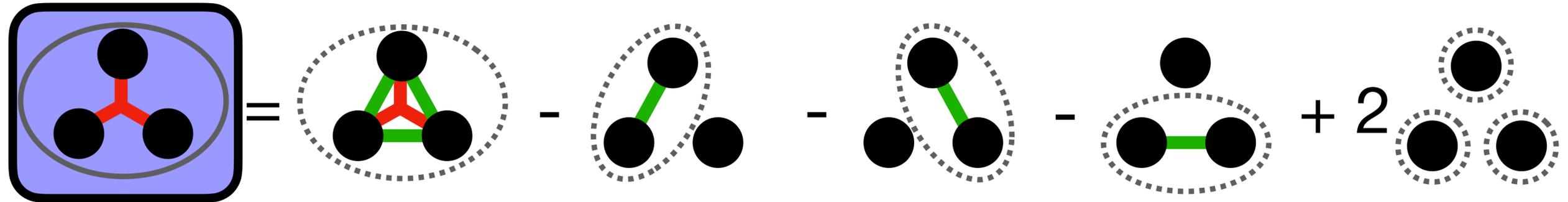
Statistical significance:

$$n_\sigma = 2.3 \text{ for } Q_3 < 0.4 \text{ GeV}/c$$

Conclusion: the measured cumulant is compatible with zero within the uncertainties



p-p-K⁻ cumulant



Zero cumulant for p-p-K⁻

Statistical significance:

$$n_\sigma = 0.5 \text{ for } Q_3 < 0.4 \text{ GeV}/c$$

Conclusion: the measured cumulant is compatible with zero within the uncertainties

p-p-K⁻ system shows only two-body interactions.

- ✓ The measurement confirms that three-body strong interaction should not be relevant in the formation of exotic kaonic bound states!

