

Model selection in kaon photoproduction

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Motivation for the work on kaon photo/electroproduction

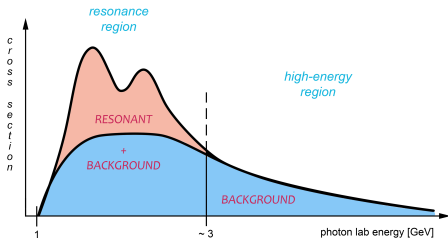
- We aim at understanding the baryon spectrum and production dynamics of particles with strangeness at low energies.
- Constituent Quark Model predicts a lot more N^* states than was observed in pion production experiments → “missing” resonance problem.
- Models for the description of elementary hyperon electroproduction are a suitable tool for hypernuclear physics calculations.
- New good-quality photoproduction data from LEPS, GRAAL, MAMI and (particularly) CLAS collaborations allow us to tune free parameters of the models.
- As the α_s increases with decreasing energy, we cannot use perturbative QCD at low energies → need for introducing effective theories and models.

Introduction

Photoproduction process



- Threshold: $E_\gamma^{lab} = 0.911 \text{ GeV}$, $W = 1.609 \text{ GeV}$
- In the lowest order, the reaction is described by the exchange of hadrons.
 - *The 3rd nucleon-resonance region:*
many resonant states and no dominant one in the $K^+\Lambda$ production
→ need to assume a large number of nucleon resonances with mass $< 2.5 \text{ GeV}$



- **Resonance region:**
resonance contributions dominate (N^*)
- **Background:**
a plenty of nonresonant contributions (p , K , Λ ; K^* and Y^*)

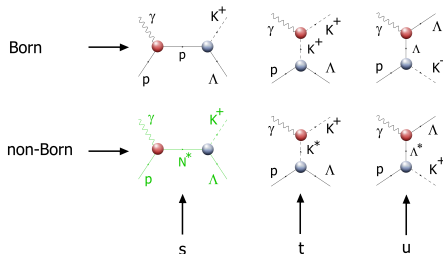
Isobar model

Single-channel approximation

- higher-order contributions (rescattering, FSI) included, to some extent, by means of effective values of the coupling constants

Use of effective hadron Lagrangian

- hadrons either in their ground or excited states
- amplitude constructed as a sum of tree-level Feynman diagrams
 - background and resonant part



- hadronic form factors account for the inner structure of hadrons

Free parameters (couplings, hff's cutoffs) adjusted to experimental data.

Satisfactory agreement with the data in the energy range $W = 1.6 - 2.5 \text{ GeV}$.

Isobar model

Novel features of our isobar model

Exchanges of high-spin resonant states

- non physical lower-spin components removed by appropriate choice of \mathcal{L}_{int}

$$V_S^\mu \mathcal{P}_{ij,\mu\nu}^{(1/2)} V_{EM}^\nu = 0$$

Energy-dependent decay widths of nucleon resonances \rightarrow restoration of unitarity

$$\Gamma(\vec{q}) = \Gamma_{N^*} \frac{\sqrt{s}}{m_{N^*}} \sum_i x_i \left(\frac{|\vec{q}_i|}{|\vec{q}_i^{N^*}|} \right)^{2l+1} \frac{D(|\vec{q}_i|)}{D(|\vec{q}_i^{N^*}|)},$$

Extension from photoproduction to electroproduction

- Phenomenological form factors in the electromagnetic vertex
- Longitudinal couplings of N^* 's to γ^* (crucial at small Q^2)

$$V_{\mu}^{EM}(N_{1/2}^* p \gamma) = -i \frac{g_3^{EM}}{(m_R + m_p)^2} \Gamma_{\mp} \gamma_{\beta} \mathcal{F}^{\beta},$$

$$V_{\mu}^{EM}(N_{3/2}^* p \gamma) = -i \frac{g_3^{EM}}{m_R(m_R + m_p)^2} \gamma_5 \Gamma_{\mp} (\not{q} g_{\mu\beta} - q_{\beta} \gamma_{\mu}) \mathcal{F}^{\beta},$$

$$V_{\mu\nu}^{EM}(N_{5/2}^* p \gamma) = -i \frac{g_3^{EM}}{(2m_p)^5} \Gamma_{\mp} (q_{\alpha} q_{\beta} g_{\mu\nu} + q^2 g_{\alpha\mu} g_{\beta\nu} - q_{\alpha} q_{\nu} g_{\beta\mu} - q_{\beta} q_{\nu} g_{\alpha\mu}) p^{\alpha} \mathcal{F}^{\beta}.$$

New fits for $K^+\Sigma^-$ channel

χ^2 minimization and overfitting

Fitting procedure with MINUIT library: **minimizing the χ^2**

$$\chi^2 = \sum_{i=1}^N \frac{[d_i - p_i(c_1, \dots, c_n)]^2}{\sigma_{d_i}^2},$$

(c_1, \dots, c_n) - set of free parameters, (d_1, \dots, d_N) - set of data points, p_i - theory, σ_{d_i} - error

Problem: χ^2 minimization cannot prevent **overfitting**

Example: polynomial curve fitting

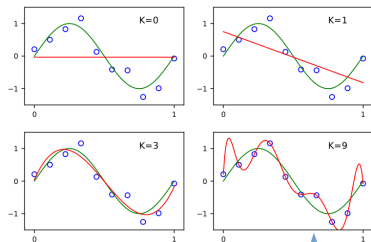
- $f(x, \mathbf{w}) = w_0 + w_1 x + w_2 x^2 + \dots + w_k x^k$

- increasing order of polynomial k fits the data well...

...but gives only poor description of the function which generated them...

...and may fail to generalize to new data

- **Occam's razor** (law of parsimony):
simpler models should be preferred



Model fits the noise in the sample

New fits of $K^+\Sigma^-$ channel

Least Absolute Shrinkage and Selection Operator (LASSO)

Remedy to the overfitting issue: regularization

- introduce a penalty term to the $\chi^2 \rightarrow$ penalization of large parameter values
- penalized χ^2_T : $\chi^2_T = \chi^2 + P(\lambda)$
- penalty term: $P(\lambda) = \lambda^4 \sum_{i=1}^{N_{res}} |g_i|$
 λ - regularization parameter, g_i - resonances' couplings
- LASSO forces some of the parameters to zero
 \rightarrow selection of a subset of the fit parameters
- λ controls the strength of the penalty and thus the complexity of the model
 \rightarrow higher powers of λ allow fine sampling of the region of small λ

New fits of $K^+\Sigma^-$ channel

Information criteria:

- Akaike information criterion

$$AIC = 2n_i + \chi_T^2$$

- Corrected Akaike information criterion

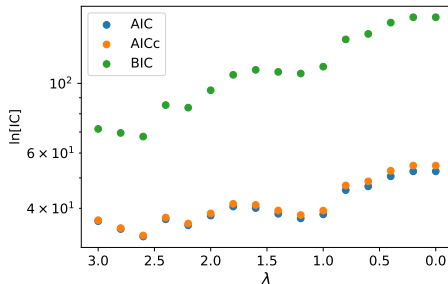
$$AICc = AIC + \frac{2n_i(n_i+1)}{N-n_i-1}$$

- Bayesian information criterion

$$BIC = n_i \ln(N) + \chi_T^2$$

n_i - no. of parameters corresponding to λ_i

N - number of data points



Applying the information criteria – forward selection

- 1 start with the full model: parameters initialized within $\langle -1; +1 \rangle$; use λ_{\max}
- 2 perform LASSO χ_T^2 minimization and compute IC
- 3 in each run reduce λ and run LASSO with the values of the previous run as starting values
- 4 repeat until λ_{\min} is reached

Optimal λ occurs at the minimum of the IC

New fits of $K^+\Sigma^-$ channel

Fitting procedure

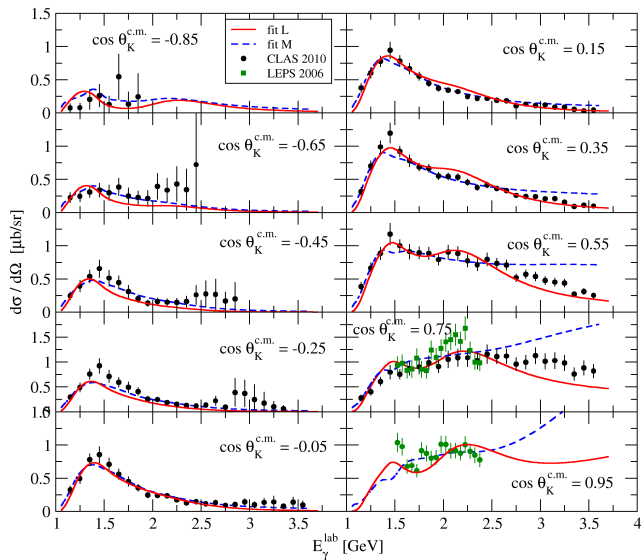
- non resonant part: Born terms and exchanges of K^* and K_1 and Σ^* 's
- resonant part: exchanges of N^* 's and Δ^* 's in the s channel
- around 600 data utilized to fit ≤ 25 parameters
- result with the smallest $\chi^2/\text{ndf} = 2.3 \rightarrow$ **fit M** (25 parameters, 14 resonances)
- **LASSO method** used: $\chi^2/\text{ndf} = 3.4 \rightarrow$ **fit L** (17 parameters, 9 resonances)

Characteristics of models

- only one Δ resonance introduced
- no hyperon resonances needed for reliable data description
- results in very good agreement with the cross-section and beam-asymmetry data
- fit L is very economical

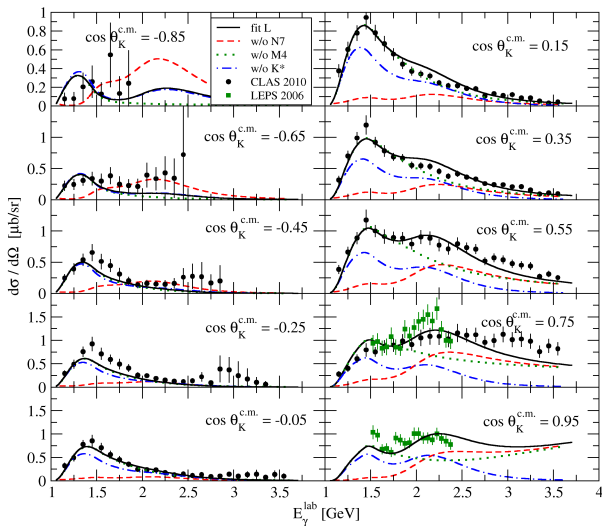
New fits of $K^+\Sigma^-$ channel

Differential cross section in dependence on the photon lab energy



New fits of $K^+\Sigma^-$ channel

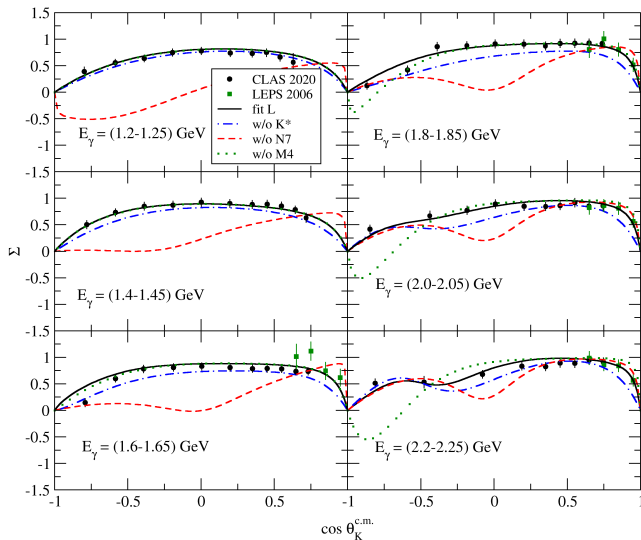
Differential cross section in dependence on the photon lab energy - fit L w/o individual resonances



N7: $N(1720)3/2^+$, M4: $N(2060)5/2^-$

New fits of $K^+\Sigma^-$ channel

Beam asymmetry in dependence on the kaon center-of-mass angle - fit L w/o individual resonances



Refitting the model's parameters in the $K^+\Lambda$ channel

Ridge regression and cross validation for suppressing hyperon couplings

Why refit?

- include recent measurements of polarization observables
- need to investigate more the role of hyperon resonances in KY photoproduction
- large values of hyperon couplings: ridge regression to suppress them during the fitting procedure

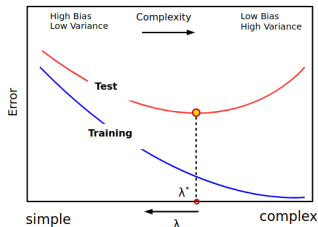
Ridge regularization

- penalized χ^2_T : $\chi^2_T = \chi^2 + \lambda^4 \sum_{i=1}^{n_\Lambda} g_i^2$, (n_Λ = no. of Y couplings)
- parameter values reduced but they are *not* reduced to zero

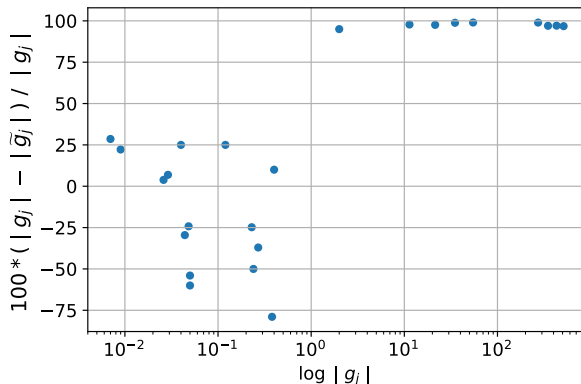
Cross validation



Bias-Variance trade-off



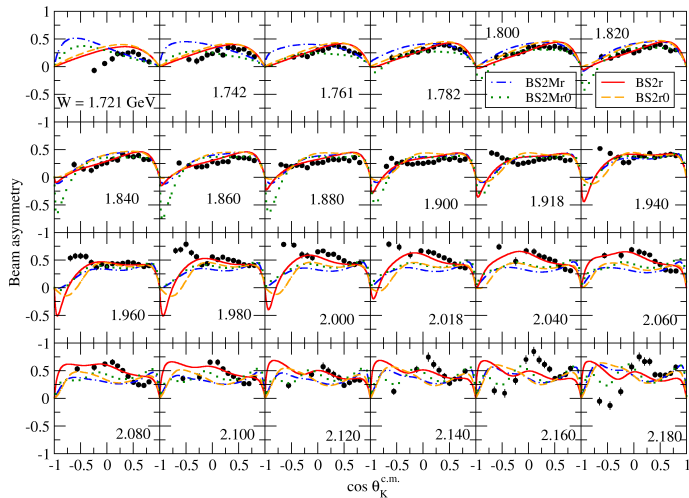
Relative percentage reduction of the resonance couplings



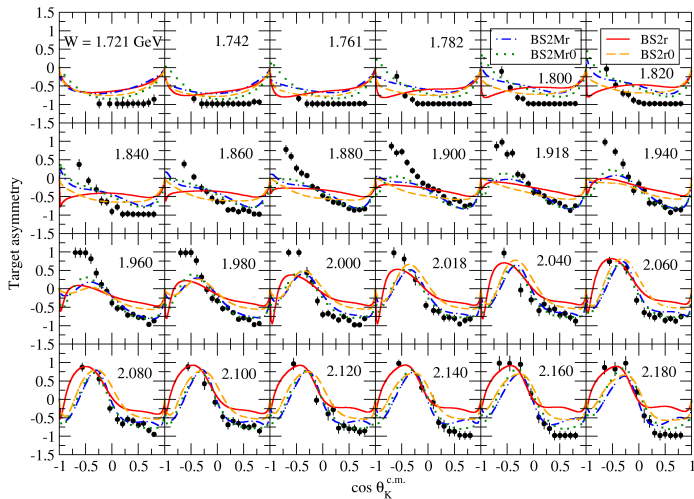
g_j - values from the unregularized fitting

\tilde{g}_j - values after performing Ridge regularization

$K^+\Lambda$ channel: beam asymmetry Σ



$K^+\Lambda$ channel: target asymmetry T



Summary

New version of **isobar model** for the $K^+\Lambda$ channel

- available for calculations **online** at:

[http://www.ujf.cas.cz/en/departments/
department-of-theoretical-physics/isobar-model.html](http://www.ujf.cas.cz/en/departments/department-of-theoretical-physics/isobar-model.html)

Description extended from the $K^+\Lambda$ channel to the $K^+\Sigma^-$ channel.

Regularization methods (LASSO, ridge) introduced as a remedy for overfitting.

Outlook

- testing the models in the DWIA calculations for hypernucleus production
- performing a multi-channel analysis of all Σ photoproduction channels
- extending the analysis of electroproduction beyond $Q^2 = 1 \text{ GeV}^2$
- studying the production of Ξ hypernuclei

Thank you for your attention!