

Istituto Nazionale di Fisica Nucleare



HAS QCD



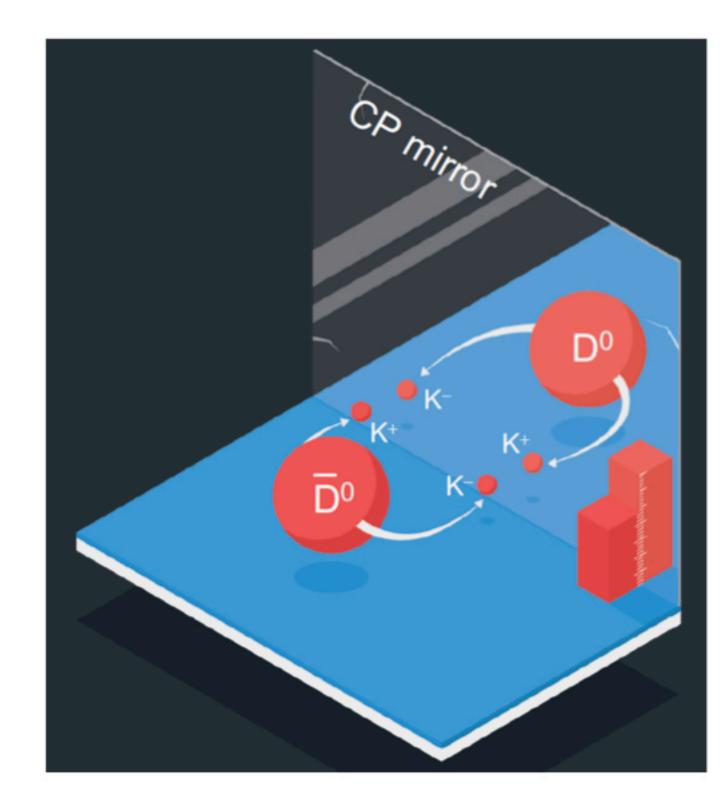
Searching strong parity violation in the proton structure

Matteo Cerutti

in collaboration with A. Bacchetta, L. Manna, M. Radici and X. Zheng

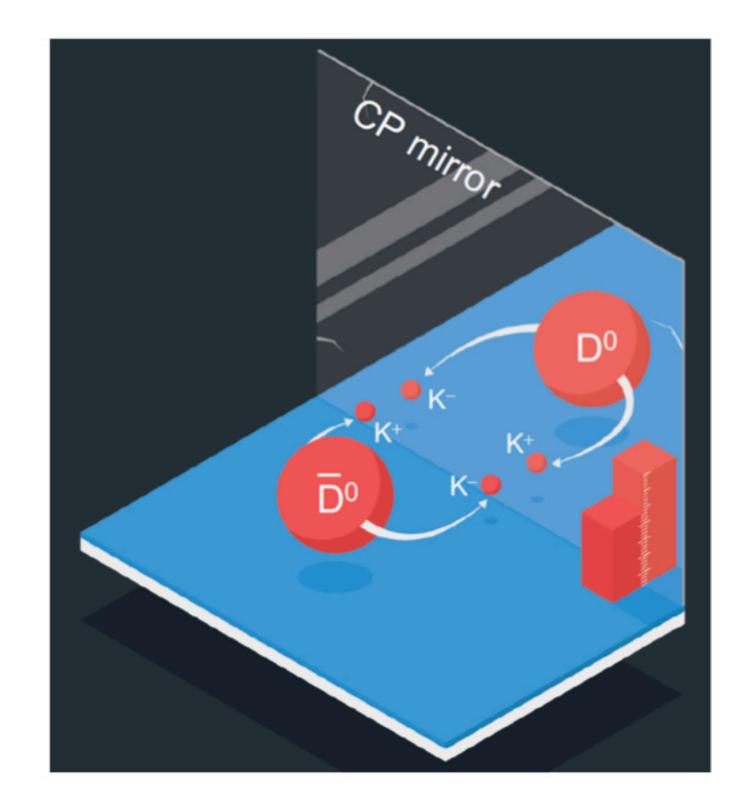
Hadron 2023 — 08/06/2023

Investigation of the "Strong CP problem"



Investigation of the "Strong CP problem"

Matter-Antimatter imbalance





EW sector

CP violation is included



EW sector

CP violation is included

Weak CP



EW sector

CP violation is included

Weak CP

too small...





EW sector

CP violation is included

Weak CP

too small...



QCD sector

EW sector

CP violation is included

Weak CP

too small...



QCD sector

Strong CP



EW sector

CP violation is included

Weak CP

too small...



QCD sector

$$\mathcal{L}_{\rm QCD}' = \mathcal{L}_{\rm QCD} + \mathcal{L}^{\rm CP}$$

Strong CP



EW sector

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Weak CP

too small...



QCD sector

Strong CP

$$\mathcal{L}_{\rm QCD}' = \mathcal{L}_{\rm QCD} + \mathcal{L}^{\rm CP}$$

 θ -term SMEFT operators



EW sector

CP violation is included

Weak CP

too small...



QCD sector

Strong CP

$$\mathcal{L}'_{\rm QCD} = \mathcal{L}_{\rm QCD} + \mathcal{L}^{\rm CP}$$

 θ -term SMEFT operators



Nucleon electric dipole moment

EW sector

CP violation is included

Weak CP

too small...



QCD sector

Strong CP

$$\mathcal{L}'_{\rm QCD} = \mathcal{L}_{\rm QCD} + \mathcal{L}^{\rm CP}$$

 θ -term SMEFT operators



Nucleon electric dipole moment

never measured...



P-symmetry

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QCD sector

QCD Lagrangian is assumed to be invariant under parity transformations

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Are there any effects of QCD P-violation on the internal structure of nucleons?

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Terms from EW sector

Wealk P-violation

Terms from QCD sector

P-symmetry

QCD sector

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Terms from EW sector

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Terms from QCD sector

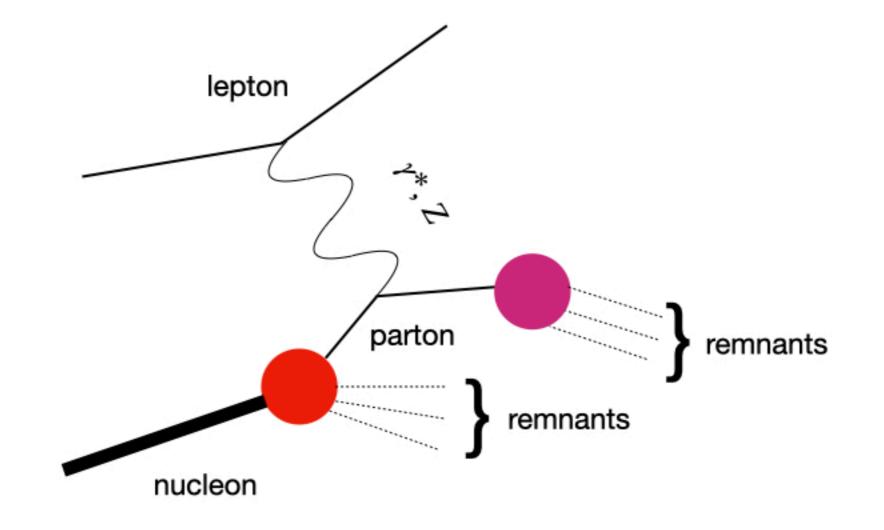
Strong P-violation



Which implications could the presence of strong P-violation cause to inclusive DIS?



$l(\ell) + N(P) \to \gamma^*(q) \to l(\ell') + X$



Cross Section

 $\frac{d^3\sigma}{dx_B dy d\phi_S} = \frac{\alpha^2 y}{2 Q^4} L_{\mu\nu}(l,l',\lambda_e) 2M W^{\mu\nu}(q,P,S)$

In general

Cross Section

$$\frac{d^3\sigma}{dx_B dy d\phi_S} = \frac{\alpha^2 y}{2 Q^4} \underbrace{L_{\mu\nu}(l,l',\lambda_e)}_{2MW^{\mu\nu}(q,P,S)} 2MW^{\mu\nu}(q,P,S)$$

In general

$$\frac{d^3\sigma}{dx_B dy d\phi_S} = \frac{\alpha^2 y}{2Q^4} \sum_{j=\gamma,\gamma Z,Z} \eta^j L^{(j)}_{\mu\nu}(l,l';\lambda_e) 2MW^{\mu\nu}(q,P,S)$$

Cross Section

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$$\eta^{\gamma} = 1 \qquad \qquad \eta^{\gamma Z} = \left(\frac{G_F M_Z^2}{2\sqrt{2}\pi\alpha}\right) \frac{Q^2}{Q^2 + M_Z^2} \qquad \qquad \eta^Z = (\eta^{\gamma Z})^2$$

$$2MW_{\mu\nu}(q,P) = \sum_{X} \int \frac{d^3 P_X}{2E_X} \delta^4(P+q-P_X) \langle P|J^{\dagger}_{\mu}(0)|P_X\rangle \langle P_X|J_{\nu}(0)|P\rangle$$

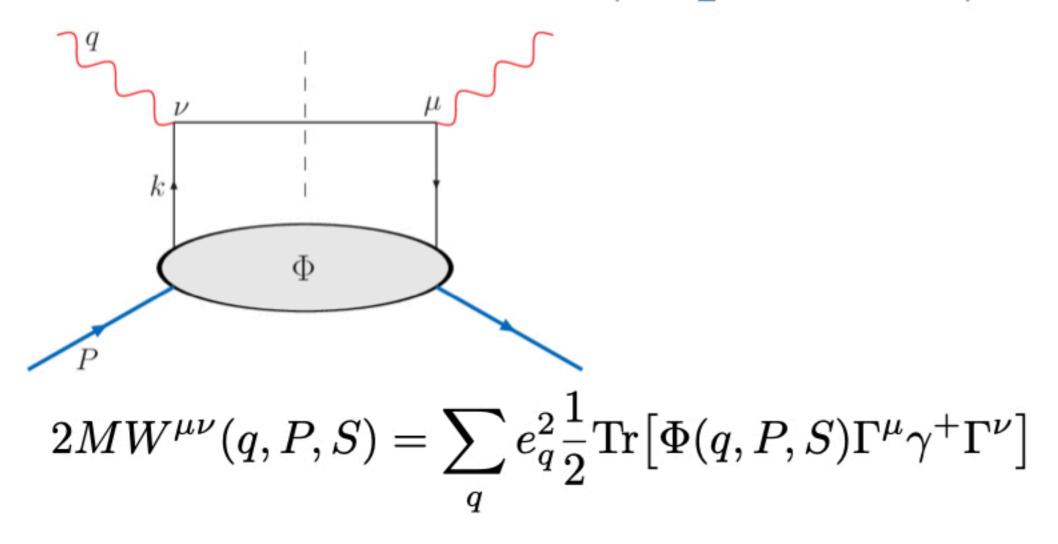
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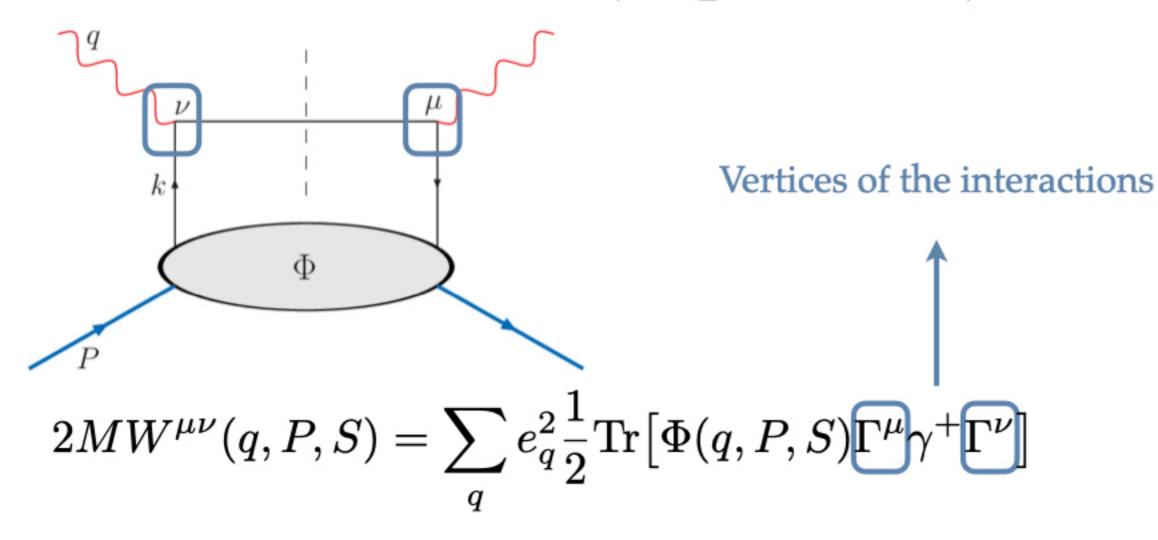
Dominant contribution on the Light-Cone

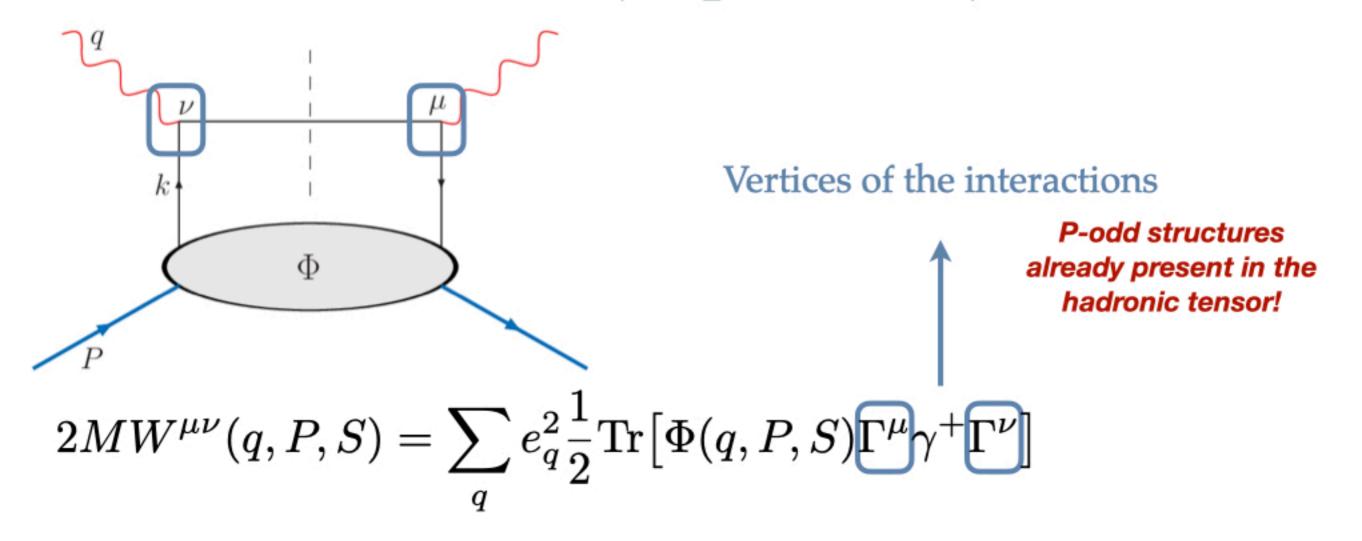
$$2MW_{\mu\nu}(q,P) = \sum_{X} \int \frac{d^{3}P_{X}}{2E_{X}} \delta^{4}(P+q-P_{X}) \langle P|J_{\mu}^{\dagger}(0)|P_{X}\rangle \langle P_{X}|J_{\nu}(0)|P\rangle$$

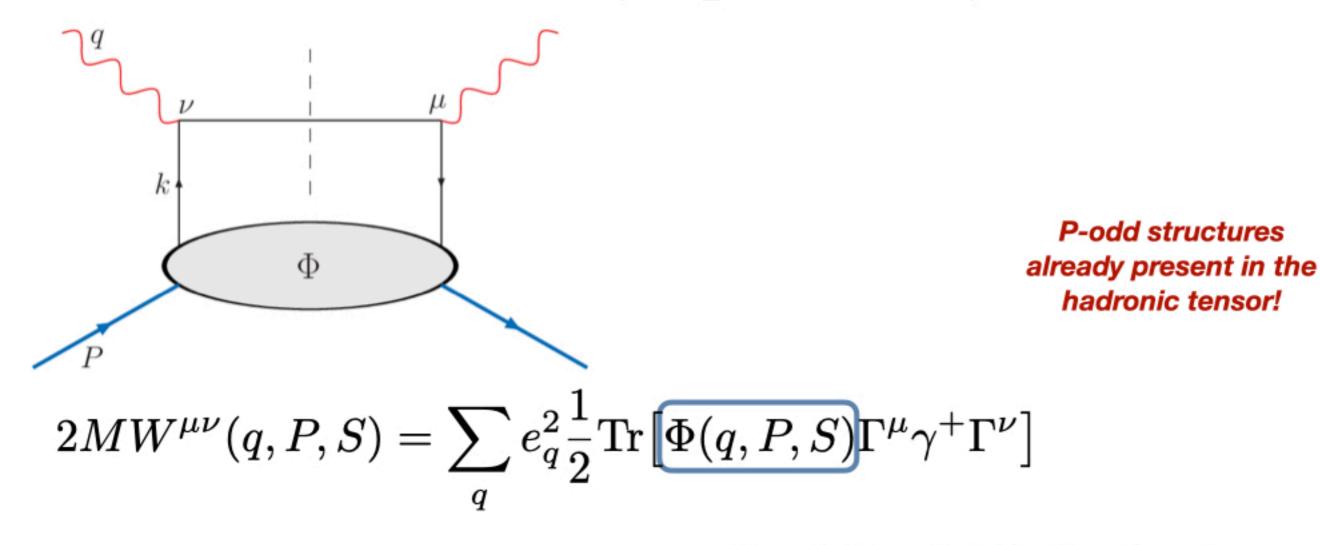
Dominant contribution on the Light-Cone

$$2MW^{\mu\nu}(q,P,S) = \sum_{q} e_q^2 \frac{1}{2} \operatorname{Tr} \left[\Phi(q,P,S) \Gamma^{\mu} \gamma^{+} \Gamma^{\nu} \right]$$

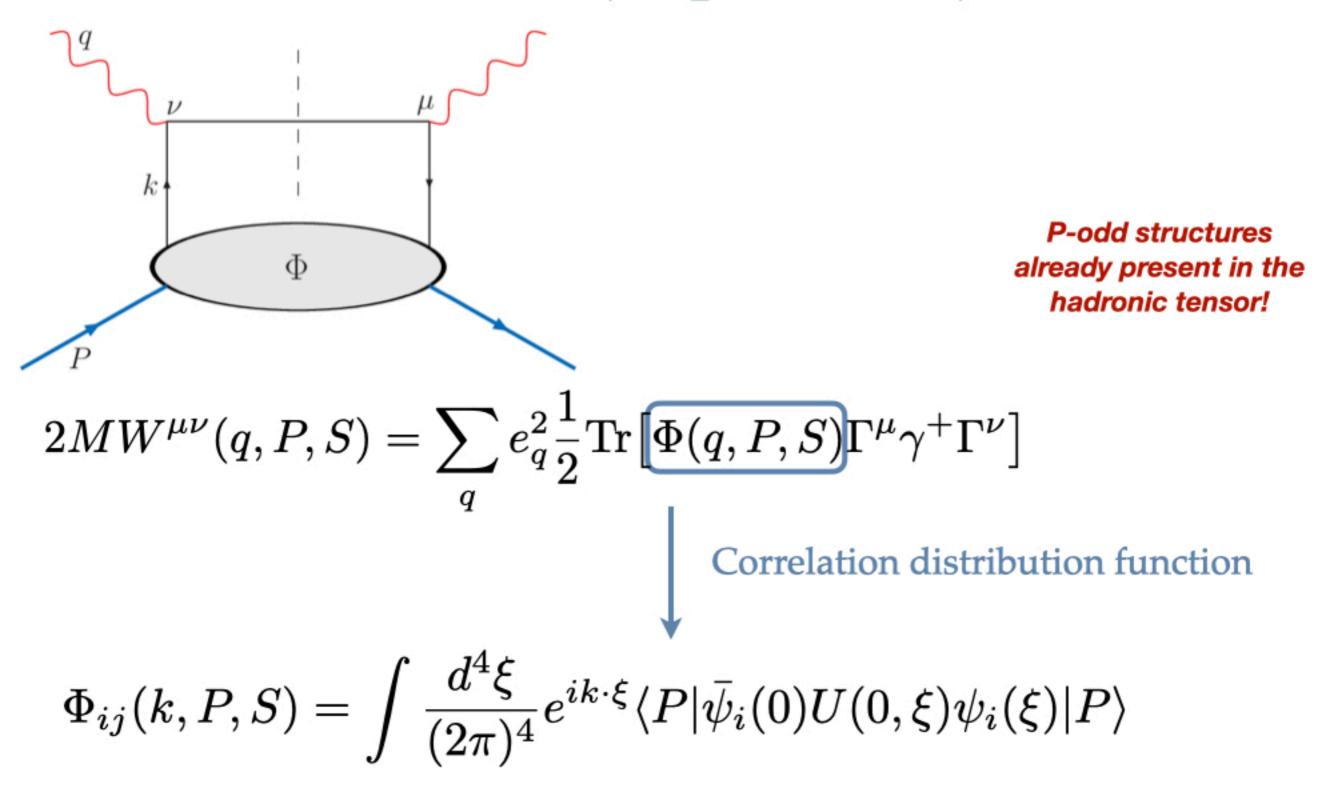


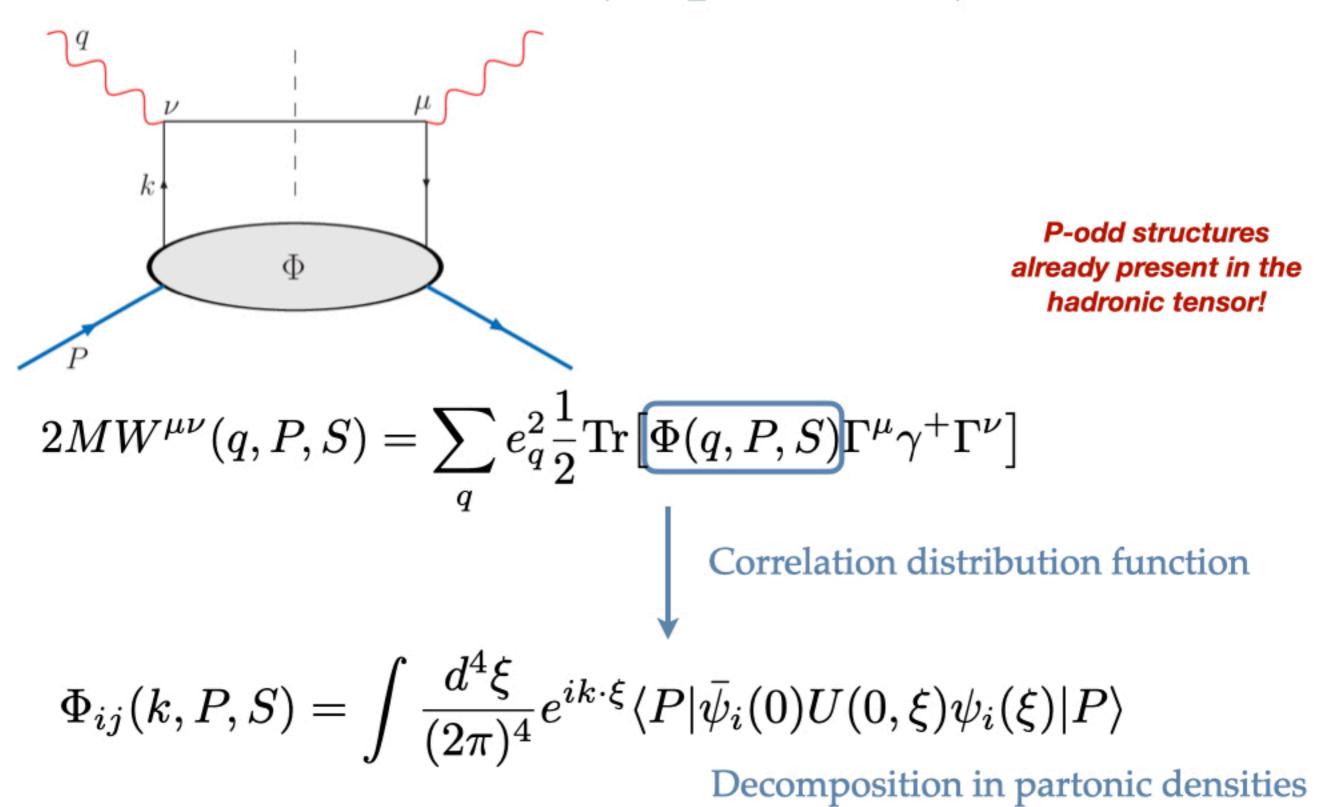






Correlation distribution function





Partonic correlator (unpolarized)

Integrated correlator

$$\Phi_{ij}(x_B) = \int \frac{d\xi^-}{2\pi} e^{ik\cdot\xi} \langle P|\bar{\psi}_j(0)\psi_i(\xi)|P\rangle_{\xi^+=\xi_T=0}$$

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Lorenz scalar

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Lorenz scalar Hermiticity Lorenz scalar Hermiticity

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Lorenz scalar Hermiticity Parity invariance Lorenz scalar Hermiticity

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Lorenz scalar Hermiticity Parity invariance Lorenz scalar Hermiticity

 $\mathbb{1},~\gamma^{\mu},~\sigma^{\mu\nu}$

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Lorenz scalar Hermiticity Parity invariance

 $\mathbb{1},~\gamma^{\mu},~\sigma^{\mu\nu}$

Lorenz scalar Hermiticity Parity invariance

$$i\gamma^5,~\gamma^\mu\gamma^5,~i\gamma^5\sigma^{\mu
u}$$

Integrated correlator

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 $\mathbb{1}, \ \gamma^{\mu}, \ \sigma^{\mu\nu}$

 $i\gamma^5, \ \gamma^\mu\gamma^5, \ i\gamma^5\sigma^{\mu
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Leading twist contributions

Integrated correlator

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Lorenz scalar Hermiticity Parity invariance Lorenz scalar Hermiticity Parity invariance

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$$i\gamma^5,\;\gamma^\mu\gamma^5,\;i\gamma^5\sigma^{\mu
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Leading twist contributions

 $\Phi_{\rm PE}(x) \simeq rac{1}{2} f_1(x) \gamma^-$

Integrated correlator

$$\Phi_{ij}(x_B) = \int \frac{d\xi^-}{2\pi} e^{ik\cdot\xi} \langle P|\bar{\psi}_j(0)\psi_i(\xi)|P\rangle_{\xi^+=\xi_T=0}$$

Lorenz scalar Hermiticity Parity invariance Lorenz scalar Hermiticity Parity invariance

1	$\gamma^{\mu},$	$\sigma^{\mu u}$
1,	γ,	0.

$$i\gamma^5,\;\gamma^\mu\gamma^5,\;i\gamma^5\sigma^{\mu
u}$$

Leading twist contributions

$$\Phi_{
m PE}(x)\simeq rac{1}{2}f_1(x)\gamma^- \qquad \qquad \Phi_{
m PV}(x)\simeq rac{1}{2}g_1^{
m PV}(x)\gamma^5\gamma^-$$

Integrated correlator

$$\Phi_{ij}(x_B) = \int \frac{d\xi^-}{2\pi} e^{ik\cdot\xi} \langle P|\bar{\psi}_j(0)\psi_i(\xi)|P\rangle_{\xi^+=\xi_T=0}$$

Lorenz scalar Hermiticity Parity invariance Lorenz scalar Hermiticity Parity invariance

1,
$$\gamma^{\mu}$$
, $\sigma^{\mu\nu}$ $i\gamma^5$, $\gamma^{\mu}\gamma^5$, $i\gamma^5\sigma^{\mu\nu}$

Leading twist contributions

$$\Phi_{\rm PE}(x) \simeq rac{1}{2} f_1(x) \gamma^ \Phi_{\rm PV}(x) \simeq rac{1}{2} g_1^{\rm PV}(x) \gamma^5 \gamma^-$$

$$\Phi(x) = \Phi_{\rm PE}(x) + \Phi_{\rm PV}(x)$$

$$\begin{aligned} \frac{d\sigma^{\pm}}{dxdy} &= \frac{2\pi\alpha^2}{xyQ^2} \Biggl[\left(Y_+ + \gamma^2 y^2/2 \right) \left(F_{2UU} + \lambda F_{2LU}^{\pm} \right) \\ &- y^2 \left(F_{L,UU} + \lambda F_{L,LU}^{\pm} \right) \\ &- \frac{Y_-}{\sqrt{1+\gamma^2}} \left(x F_{3UU}^{\pm} + \lambda x F_{3LU} \right) \Biggr] \end{aligned}$$

$$\frac{d\sigma^{\pm}}{dxdy} = \frac{2\pi\alpha^2}{xyQ^2} \Big[Y_+ F_2^{\pm} - y^2 F_L^{\pm} \mp Y_- x F_3^{\pm} \Big]$$
 PDG 2023

 $xF_{3LU}(x,Q^2) = xF_3^{(\gamma)} - g_V^e \eta_{\gamma Z} xF_3^{(\gamma Z)} + \left(g_V^{e^2} + g_A^{e^2}\right) \eta_Z xF_3^{(Z)}$

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$$egin{aligned} xF_3^{(\gamma)}(x,Q^2) &= 0 \ xF_3^{(\gamma Z)}(x,Q^2) &= \sum_q 2e_q g_A^q x f_1^{(q-ar q)} \ xF_3^{(Z)}(x,Q^2) &= \sum_q 2g_V^q g_A^q x f_1^{(q-ar q)} \end{aligned}$$

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 $xF_3^{(\gamma)}(x,Q^2) = 0$ $xF_3^{(\gamma Z)}(x,Q^2) = \sum 2e_q g_A^q x f_1^{(q-\bar{q})}$ q $xF_3^{(Z)}(x,Q^2) = \sum 2g_V^q g_A^q x f_1^{(q-\bar{q})}$ $x\Delta F_3^{(\gamma)}(x,Q^2) = -\sum_q e_q^2 x g_1^{\text{PV}(q+\bar{q})}$ $x\Delta F_3^{(\gamma Z)}(x,Q^2) = -\sum 2e_q g_V^q x g_1^{\mathrm{PV}(q+\bar{q})}$ Additional contributions due to the new PV parton $x\Delta F_3^{(Z)}(x,Q^2) = -\sum (g_V^{q2} + g_A^{q2}) x g_1^{\mathrm{PV}(q+\bar{q})}$ distribution

$$xF_{3LU}(x,Q^2) = xF_3^{(\gamma)} - g_V^e \eta_{\gamma Z} xF_3^{(\gamma Z)} + \left(g_V^{e^2} + g_A^{e^2}\right) \eta_Z xF_3^{(Z)}$$

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 $xF_3^{(\gamma Z)}(x,Q^2) = \sum_q 2e_q g_A^q x f_1^{(q-\bar{q})}$
 $xF_3^{(Z)}(x,Q^2) = \sum_q 2g_V^q g_A^q x f_1^{(q-\bar{q})}$
 $x\Delta t$
Additional contributions

MAIN INNOVATION OF PV-HYPOTESIS

$$egin{aligned} &x\Delta F_3^{(\gamma)}(x,Q^2)=-\sum_q e_q^2 x g_1^{ ext{PV}(q+ar q)} \ &x\Delta F_3^{(\gamma Z)}(x,Q^2)=-\sum_q 2 e_q g_V^q x g_1^{ ext{PV}(q+ar q)} \ &x\Delta F_3^{(Z)}(x,Q^2)=-\sum_q ig(g_V^{q2}+g_A^{q2}ig) x g_1^{ ext{PV}(q+ar q)} \end{aligned}$$

Additional contributions due to the new PV parton distribution

$$\begin{aligned} \frac{d\sigma^{\pm}}{dxdy} &= \frac{2\pi\alpha^2}{xyQ^2} \Biggl[\left(Y_+ + \gamma^2 y^2 / 2 \right) \left(F_{2UU} + \lambda F_{2LU}^{\pm} \right) \\ &- y^2 \left(F_{L,UU} + \lambda F_{L,LU}^{\pm} \right) \\ &- \frac{Y_-}{\sqrt{1+\gamma^2}} \left(x F_{3UU}^{\pm} + \lambda x F_{3LU} \right) \Biggr] \end{aligned}$$

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Standard DIS structure functions

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Standard DIS structure functions

$$\begin{split} F_{2UU}(x,Q^2) &= F_2^{(\gamma)} - g_V^e \eta_{\gamma Z} F_2^{(\gamma Z)} + \left(g_V^{e^2} + g_A^{e^2}\right) \eta_Z F_2^{(Z)}, \\ F_{2LU}^{\pm}(x,Q^2) &= \mp g_A^e \eta_{\gamma Z} F_2^{(\gamma Z)} \pm 2g_V^e g_A^e \eta_Z F_2^{(Z)}, \\ xF_{3UU}^{\pm}(x,Q^2) &= \mp g_A^e \eta_{\gamma Z} xF_3^{(\gamma Z)} \pm 2g_V^e g_A^e \eta_Z xF_3^{(Z)}, \\ xF_{3LU}(x,Q^2) &= xF_3^{(\gamma)} - g_V^e \eta_{\gamma Z} xF_3^{(\gamma Z)} + \left(g_V^{e^2} + g_A^{e^2}\right) \eta_Z xF_3^{(Z)}, \end{split}$$

Phenomenology

Experimental observable

PVDIS Asymmetry

$$A_{\rm PV} \equiv \frac{d\sigma(\lambda=1) - d\sigma(\lambda=-1)}{d\sigma(\lambda=1) + d\sigma(\lambda=-1)}$$

PVDIS Collaboration, Nature 506 (2014) D. Wang et al., Phys.Rev.C 91 (2015)

Experimental observable

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PVDIS Collaboration, Nature 506 (2014) D. Wang et al., Phys.Rev.C 91 (2015)

$$=\frac{Y_{+}F_{2LU}-y^{2}F_{L,LU}-Y_{-}xF_{3LU}}{Y_{+}F_{2UU}-y^{2}F_{L,UU}-Y_{-}xF_{3UU}}$$

 $Y_{\pm} = 1 \pm (1 - y)^2$

Experimental observable

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PVDIS Collaboration, Nature 506 (2014) D. Wang et al., Phys.Rev.C 91 (2015)

$$=rac{Y_+F_{2LU}-y^2F_{L,LU}-Y_-xF_{3LU}}{Y_+F_{2UU}-y^2F_{L,UU}-Y_-xF_{3UU}}$$

Contribution of g_1^{PV} in each of the structure functions due to γZ and Z channels

 $Y_{\pm} = 1 \pm (1 - y)^2$

HERA dataset (Run I + II combined)

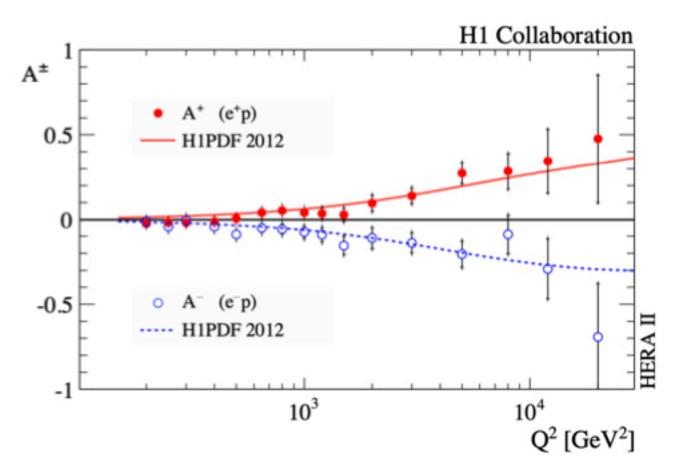
H1 Collaboration, Eur. Phys. J. C 78 (2018)

HERA dataset (Run I + II combined)

H1 Collaboration, Eur. Phys. J. C 78 (2018)



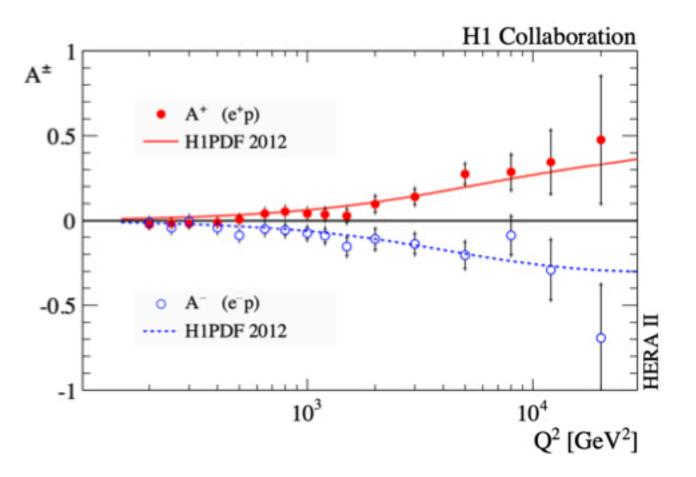
e⁻ asymmetry: 138 data



HERA dataset (Run I + II combined)

H1 Collaboration, Eur. Phys. J. C 78 (2018)

 e^+ asymmetry: 136 data e^- asymmetry: 138 data



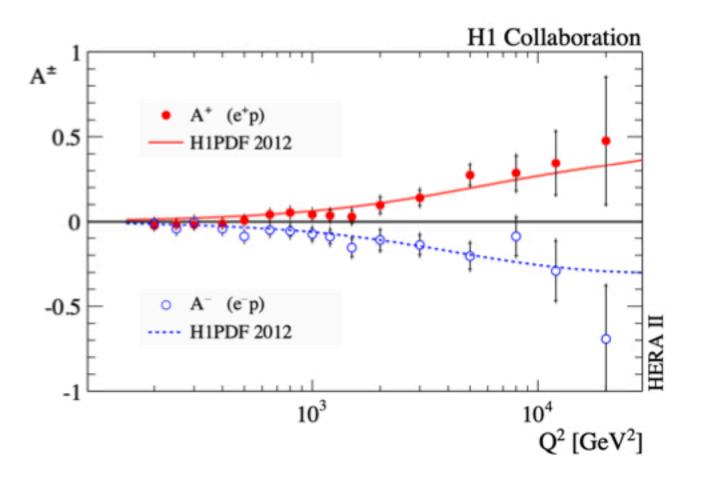
JLab6 PVDIS dataset

PVDIS Collaboration, *Nature* 506 (2014) D. Wang et al., Phys.Rev.C 91 (2015)

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H1 Collaboration, Eur. Phys. J. C 78 (2018)

 e^+ asymmetry: 136 data e^- asymmetry: 138 data



JLab6 PVDIS dataset

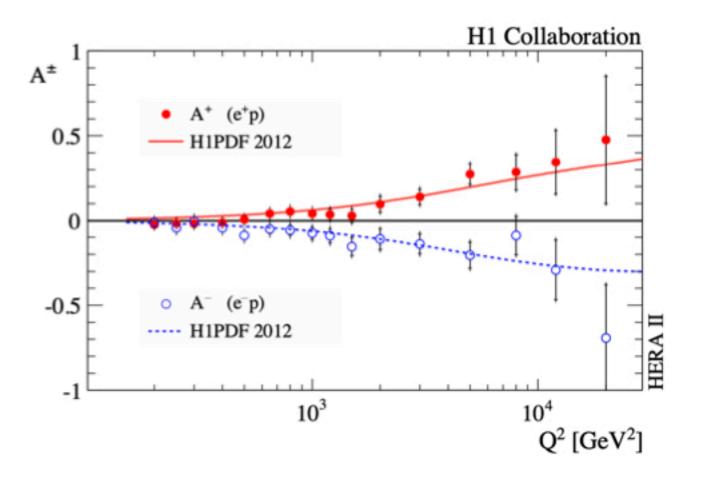
PVDIS Collaboration, *Nature* 506 (2014) D. Wang et al., Phys.Rev.C 91 (2015)

e^- asymmetry: 2 data

HERA dataset (Run I + II combined)

H1 Collaboration, Eur. Phys. J. C 78 (2018)

 e^+ asymmetry: 136 data e^- asymmetry: 138 data



JLab6 PVDIS dataset

PVDIS Collaboration, *Nature* 506 (2014) D. Wang et al., Phys.Rev.C 91 (2015)

SLAC-E122 dataset

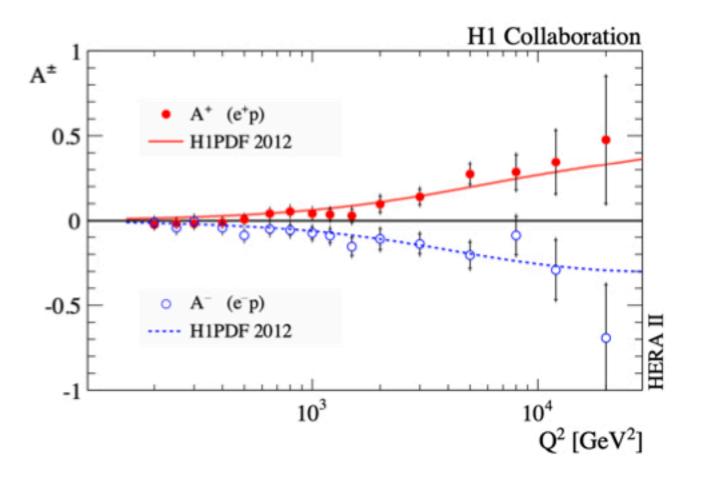
C.Y. Prescott et al., Phys. Lett. B (1979)

e⁻ asymmetry: 2 data

HERA dataset (Run I + II combined)

H1 Collaboration, Eur. Phys. J. C 78 (2018)

e⁺ asymmetry: 136 data e⁻ asymmetry: 138 data



JLab6 PVDIS dataset

PVDIS Collaboration, *Nature* 506 (2014) D. Wang et al., Phys.Rev.C 91 (2015)

SLAC-E122 dataset

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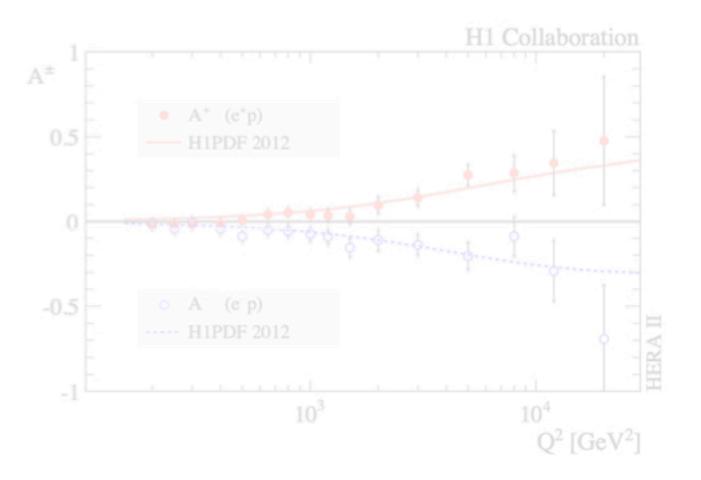
e⁻ asymmetry: 2 data

e⁻ asymmetry: 11 data

HERA dataset (Run I + II combined)

H1 Collaboration, Eur. Phys. J. C 78 (2018)

e⁺ asymmetry: 136 data e⁻ asymmetry: 138 data



JLab6 PVDIS dataset

PVDIS Collaboration, *Nature* 506 (2014) D. Wang et al., Phys.Rev.C 91 (2015)

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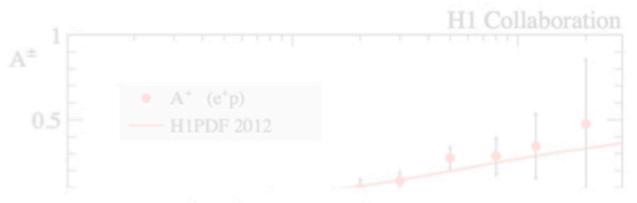
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HERA dataset (Run I + II combined)

H1 Collaboration, Eur. Phys. J. C 78 (2018)

 e^+ asymmetry: 136 data e^- asymmetry: 138 data



Imbalance between information from electron and positron beams

JLab6 PVDIS dataset

PVDIS Collaboration, *Nature* 506 (2014) D. Wang et al., Phys.Rev.C 91 (2015)

SLAC-E122 dataset

C.Y. Prescott et al., Phys. Lett. B (1979)

e⁻ asymmetry: 2 data

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 O^2 [GeV²]

HERA dataset

HERA dataset

 $Q^2 \in (200, 30000) \ {
m GeV}^2$

HERA dataset

 $Q^2 \in (200, 30000) \text{ GeV}^2$

high-energy $Q^2 \gg M_N^2$ no need of modification of the theory

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JLab6 + SLAC-E122 datasets

HERA dataset

 $Q^2 \in (200, 30000) \text{ GeV}^2$

high-energy $Q^2 \gg M_N^2$ no need of modification of the theory

JLab6 + SLAC-E122 datasets

$$Q^2 \in (0.9, 1.9) \text{ GeV}^2$$

low-energy $Q^2 \simeq M_N^2$

HERA dataset

 $Q^2 \in (200, 30000) \text{ GeV}^2$

high-energy $Q^2 \gg M_N^2$ no need of modification of the theory

JLab6 + SLAC-E122 datasets

 $Q^2 \in (0.9, 1.9) \text{ GeV}^2$

low-energy $Q^2 \simeq M_N^2$ applicability of the theory?

HERA dataset

high-energy $Q^2 \gg M_N^2$ no need of modification of the theory

JLab6 + SLAC-E122 datasets

low-energy $Q^2 \simeq M_N^2$ applicability of the theory? $Q^2 \in (200, 30000) \text{ GeV}^2$

$$Q^2 \in (0.9, 1.9) \text{ GeV}^2$$

Target-Mass Corrections

e.g., A. Bacchetta et al., JHEP 02 (2007)

Experimental data: energy range

HERA dataset

high-energy $Q^2 \gg M_N^2$ no need of modification of the theory

JLab6 + SLAC-E122 datasets

low-energy $Q^2 \simeq M_N^2$ applicability of the theory? $Q^2 \in (200, 30000) \text{ GeV}^2$

$$Q^2 \in (0.9, 1.9) \text{ GeV}^2$$

Target-Mass Corrections

e.g., A. Bacchetta et al., JHEP 02 (2007)

EW radiative corrections

J. Erler, S. Su, Prog.Part.Nucl.Phys. 71 (2013)

Experimental data: energy range

HERA dataset

high-energy $Q^2 \gg M_N^2$ no need of modification of the theory

JLab6 + SLAC-E122 datasets

low-energy

 $Q^2 \simeq M_N^2$

applicability of the theory?

$$\begin{aligned} C_{1u} &= 2g_A^e g_V^u = 2\left(-\frac{1}{2}\right)\left(\frac{1}{2} - \frac{4}{3}\sin^2\theta_W\right) \\ C_{2u} &= 2g_V^e g_A^u = 2\left(-\frac{1}{2} + 2\sin^2\theta_W\right)\left(\frac{1}{2}\right) \\ C_{1d} &= 2g_A^e g_V^d = 2\left(-\frac{1}{2}\right)\left(-\frac{1}{2} + \frac{2}{3}\sin^2\theta_W\right) \\ C_{2d} &= 2g_V^e g_A^d = 2\left(-\frac{1}{2} + 2\sin^2\theta_W\right)\left(-\frac{1}{2}\right) \end{aligned}$$

 $Q^2 \in (200, 30000) \text{ GeV}^2$

$$Q^2 \in (0.9, 1.9) \text{ GeV}^2$$

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e.g., A. Bacchetta et al., JHEP 02 (2007)

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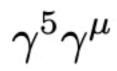
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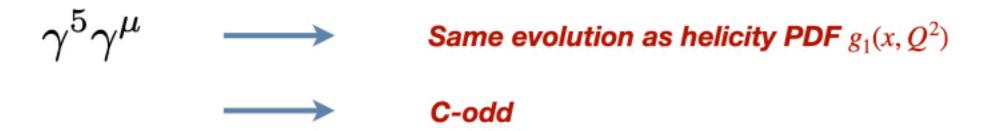
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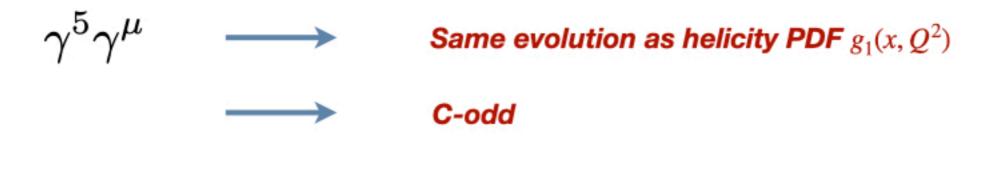
J. Erler, S. Su, Prog.Part.Nucl.Phys. 71 (2013)

$$\begin{array}{lll} C_{1u}^{\rm SM} &=& -0.1887 - 0.0011 \times \frac{2}{3} \ln(\langle Q^2 \rangle / 0.14 {\rm GeV}^2) \\ C_{1d}^{\rm SM} &=& 0.3419 - 0.0011 \times \frac{-1}{3} \ln(\langle Q^2 \rangle / 0.14 {\rm GeV}^2) \\ C_{2u}^{\rm SM} &=& -0.0351 - 0.0009 \ln(\langle Q^2 \rangle / 0.078 \; {\rm GeV}^2) \\ C_{2d}^{\rm SM} &=& 0.0248 + 0.0007 \ln(\langle Q^2 \rangle / 0.021 \; {\rm GeV}^2) \end{array}$$









$$xF_{3}^{j}(x,Q^{2}) = \sum_{q} C_{q}^{j} x f_{1}^{(q-\bar{q})}$$

$$\gamma^5 \gamma^\mu \longrightarrow Same evolution as helicity PDF $g_1(x, Q^2)$
 $\longrightarrow C-odd$$$

$$xF_{3}^{j}(x,Q^{2}) = \sum_{q} C_{q}^{j} x f_{1}^{(q-\bar{q})} \qquad \Delta xF_{3}^{j}(x,Q^{2}) = -\sum_{q} C_{q}^{'j} x \alpha g_{1}$$

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$$F_2^j(x,Q^2) = \sum_q \hat{C}_q^j x f_1^{(q+\bar{q})}$$

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PV parton density comes from the structure

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1 parameter to be fitted

PDF set for

PDF set for

 $f_1(x,Q^2)$

NNPDF4.0 Ball et al. (NNPDF), EPJ C 82 (2022)

PDF set for

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NNPDFpol1.1

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100 MC replicas of unpolarized PDF

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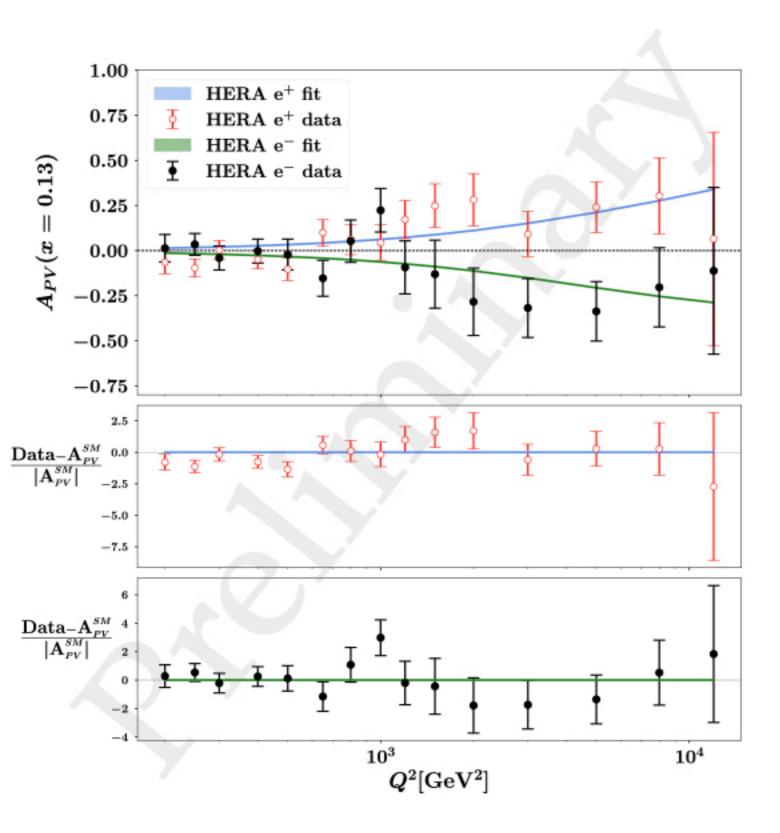
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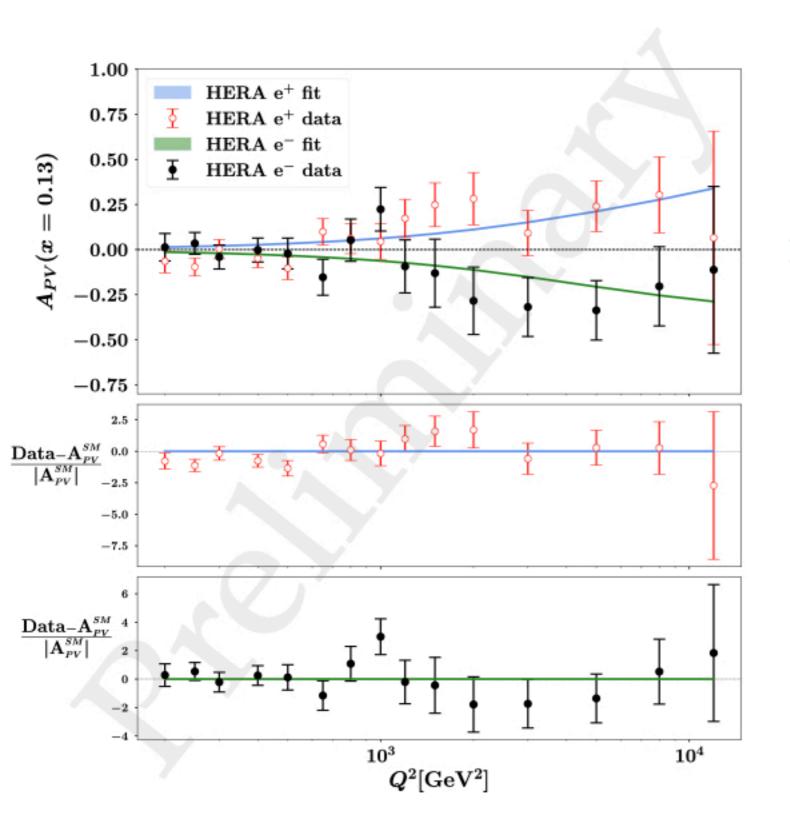
Statistical distribution of 100 values of parameter α

Results of the fit: χ^2 **values**

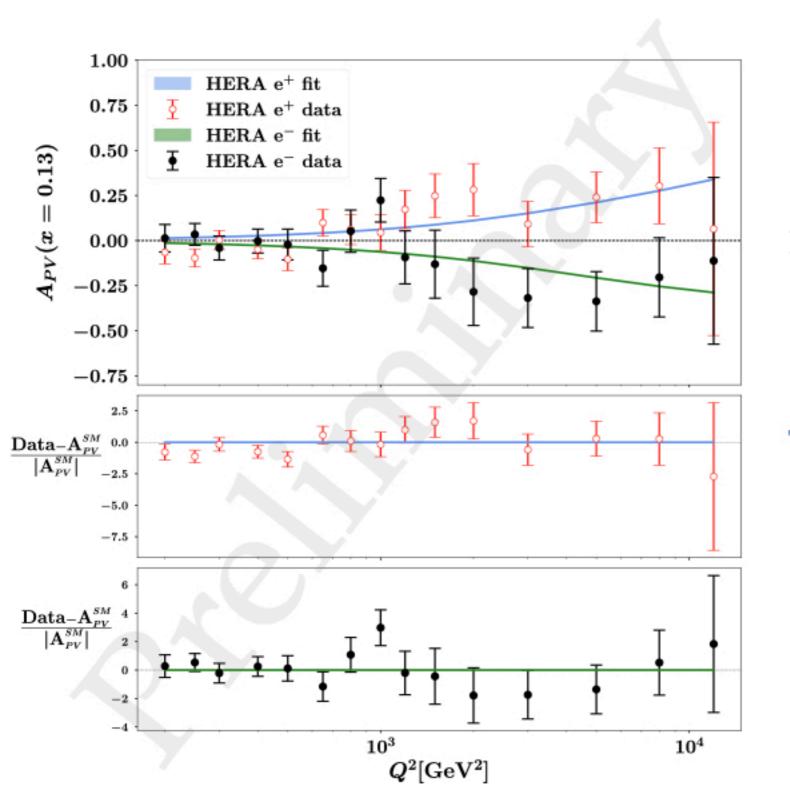
Fit WITH EW radiative corrections

	N of points	χ²/N _{data} (SM)	χ²/N _{data} (Fit)
HERA A^+	136	1.12	1.12
HERA A^-	138	0.98	0.98
JLab6 A^-	2	0.67	0.42
SLAC-E122 A	11	0.97	0.94
TOTAL	287	1.042	1.037



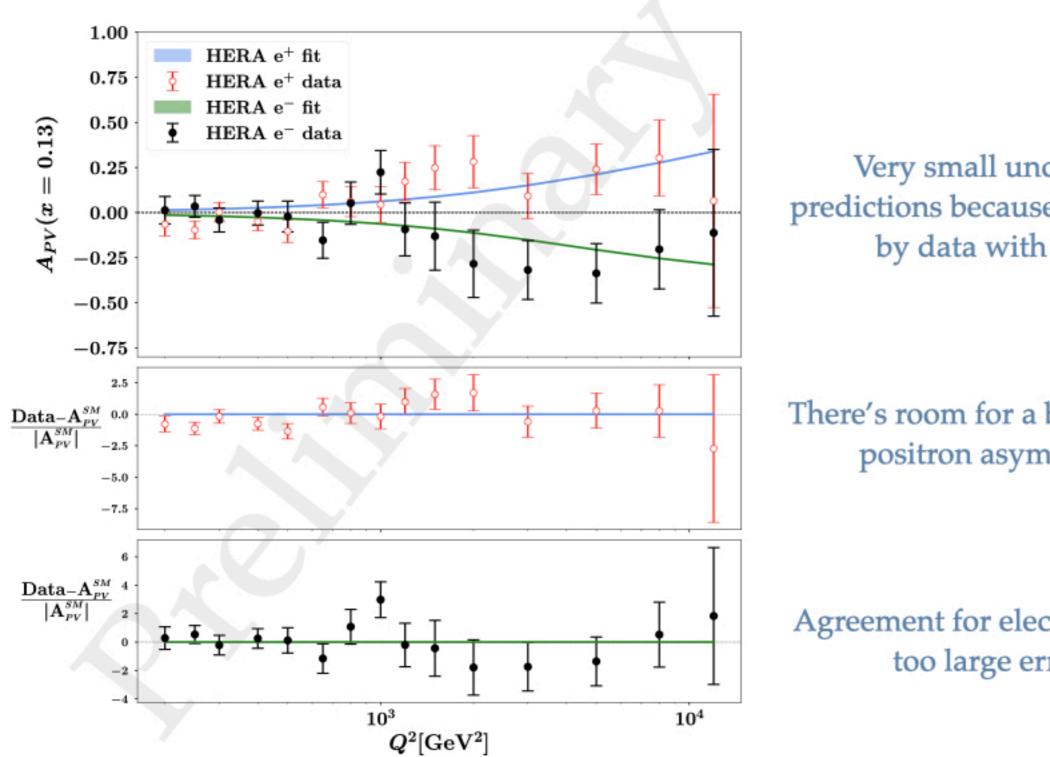


Very small uncertainties in the predictions because the fit is dominated by data with smaller errors



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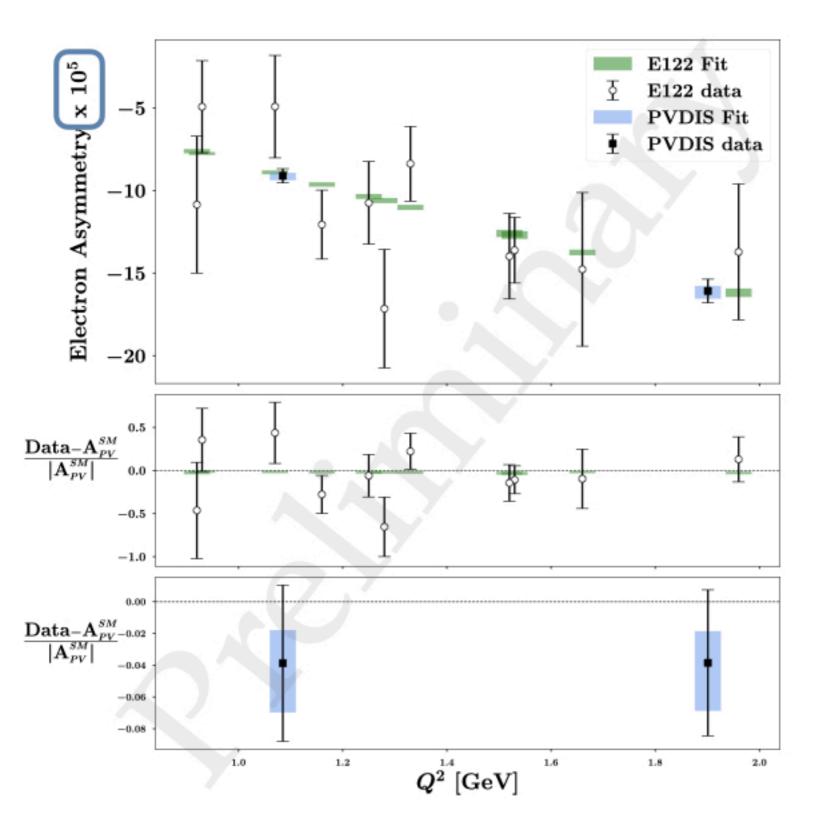
There's room for a better description for positron asymmetry at low-Q

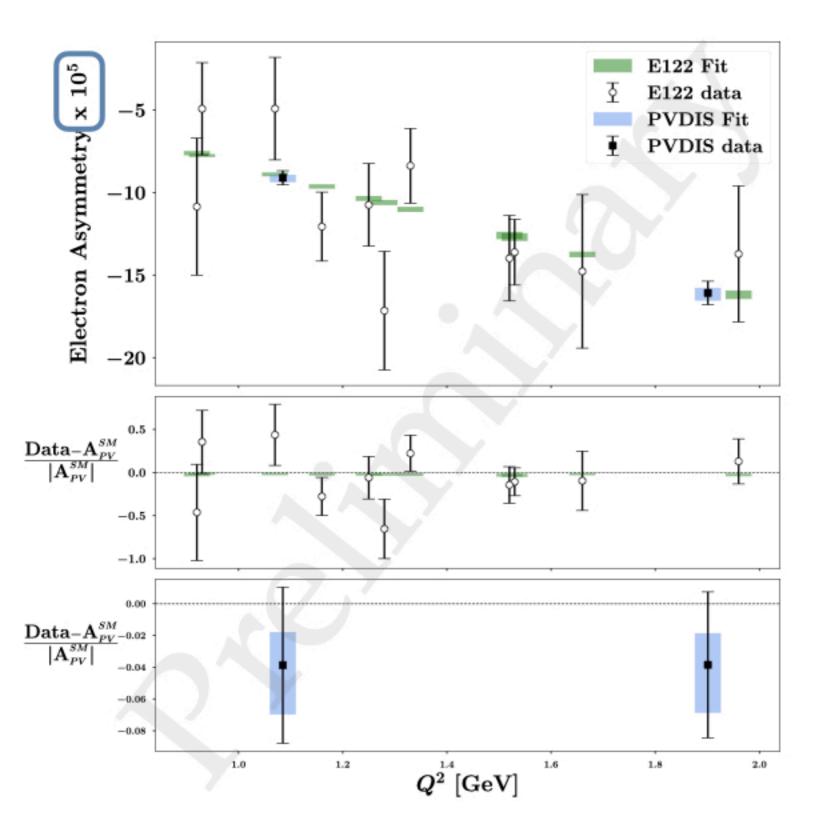


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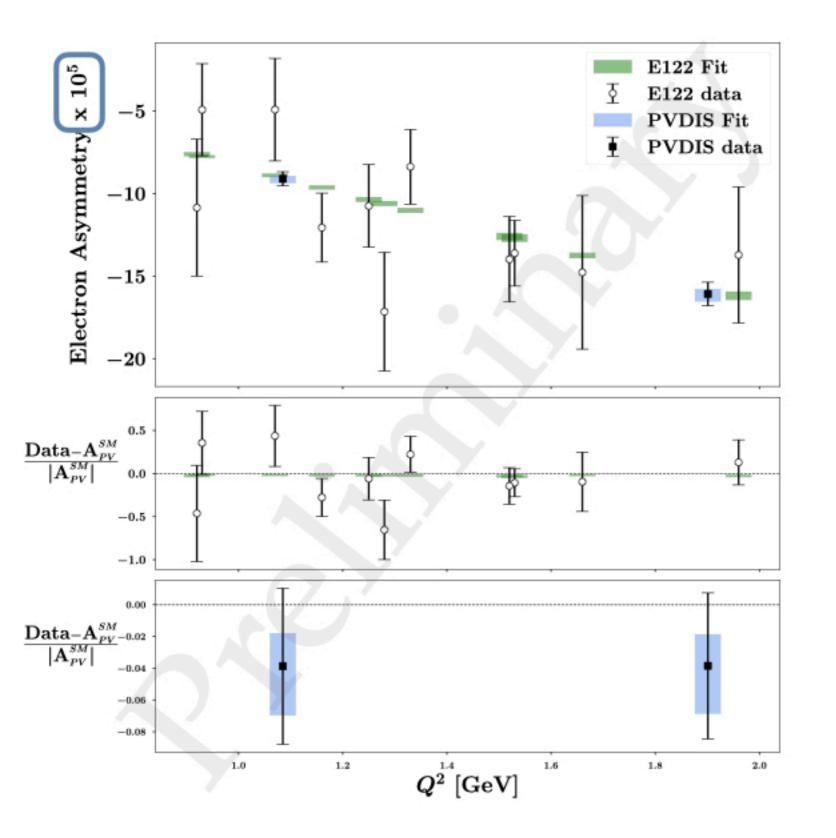
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Agreement for electron asymmetry, but too large errors at low-Q



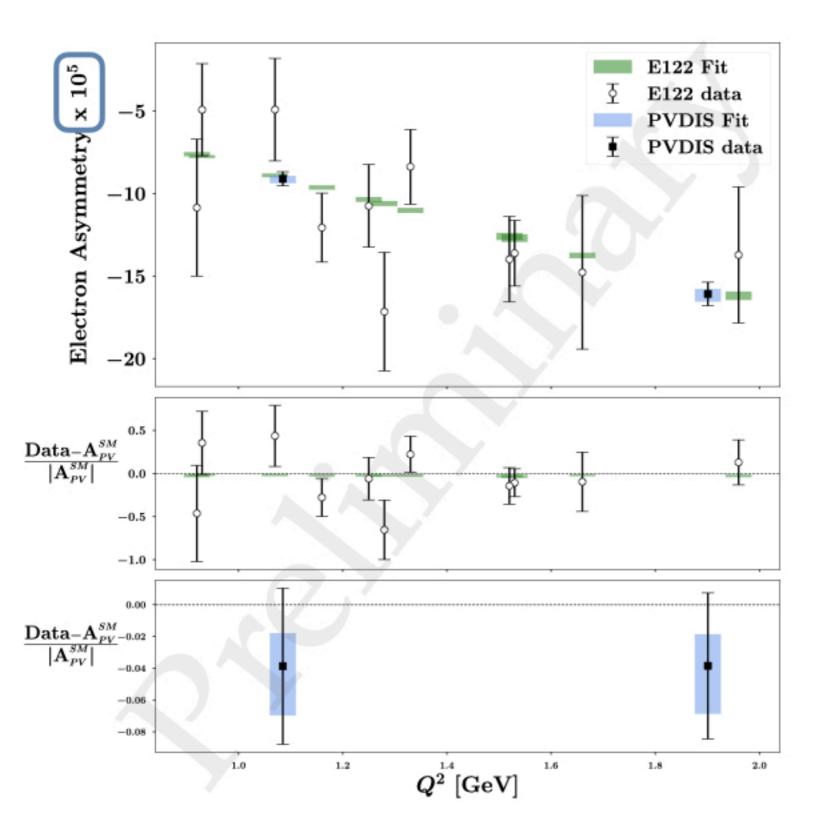


Sizeable improvement of the fit w.r.t. SM predictions



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Old dataset with still quite large experimental errors (> 20 %)



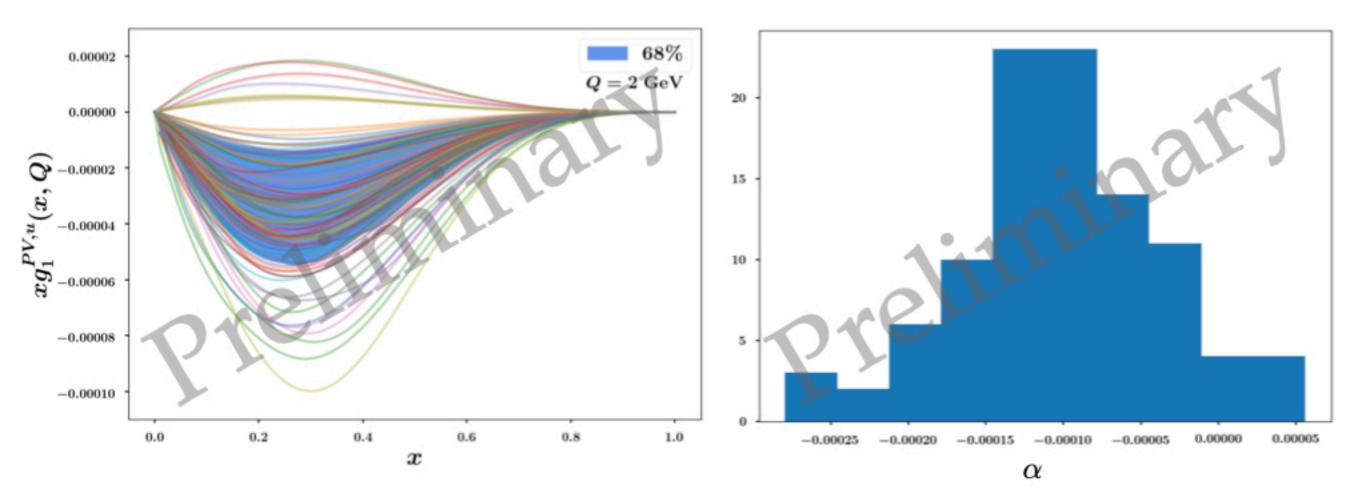
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Data points which actually drive the fit due to very small experimental errors (~%) **Results of the fit:** $g_1^{PV}(x, Q^2)$ **extraction**

$$g_1^{PV}(x) = \alpha \ g_1(x)$$

$$\alpha = (-1.01 \pm 0.66) \cdot 10^{-4}$$



 The strong P- violation can give origin to a new structure function in DIS cross section for one-photon exchange

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- A fit of present experimental data is compatible with a non-zero contribution from a new strong PV parton density at more than 1 sigma

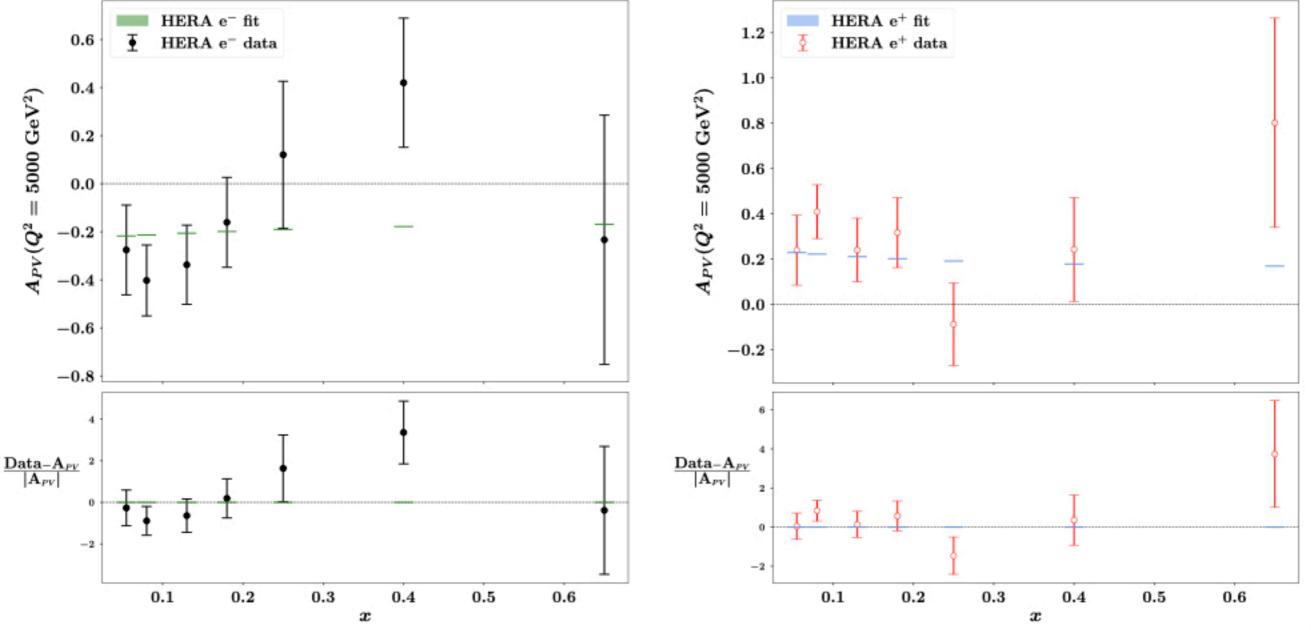
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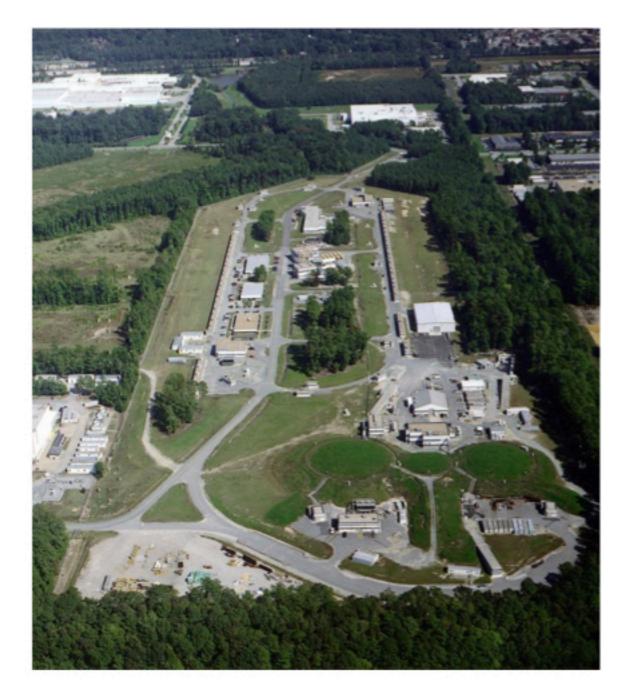
 Experimental data from positron beam are welcome to shed light on the complementarity with electron beam

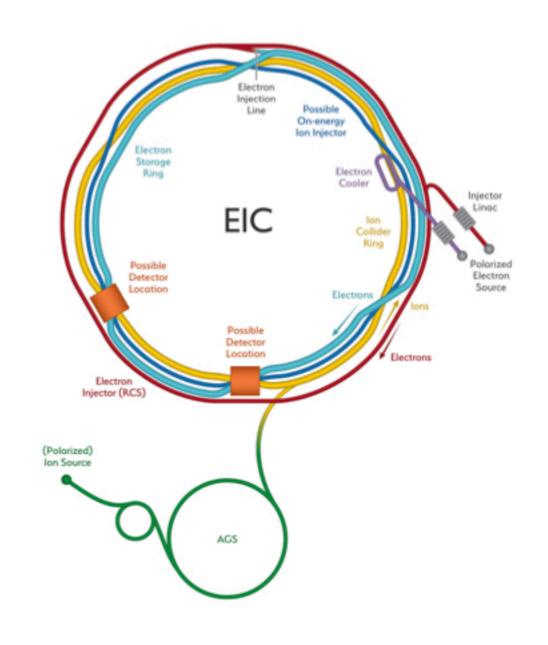
 A different behaviour of the PV parton distribution w.r.t. the variable x can be investigated

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 Predictions of the size of the PV distribution can be made in the kinematic domains of JLab12, JLab20+(?) and EIC





 Further investigations on a new P-odd, CP-odd distribution function arising when considering the polarisation of the target

$$\begin{split} \Phi^q(x,Q^2) &= \left\{ f_1^q(x,Q^2) + g_1^{\mathrm{PV}q}(x,Q^2)\gamma_5 \\ &+ S_L \Big(g_1^q(x,Q^2)\gamma_5 + f_{1L}^{\mathrm{PV}q}(x,Q^2) \Big) \\ &- \mathscr{G}_T \Big(h_1^q(x,Q^2)\gamma_5 - e_{1T}^{\mathrm{PV}q}(x,Q^2) \Big) \right\} \frac{\not{h}_+}{2} \end{split}$$

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$$\Delta x_B g_5(x_B, Q^2) \approx \Delta x_B g_5^{(\gamma)}(x_B, Q^2) = \frac{1}{2} \sum_q e_q^2 x_B f_{1L}^{\text{PV}(q-\bar{q})}$$