## HAS QCD

# Searching strong parity violation in 



## Motivations

Investigation of the "Strong CP problem"


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Matter-Antimatter imbalance


## Motivations

EW sector
CP violation is included

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EW sector
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Weak CP

## Motivations

EW sector
Weak CIP

CP violation is included too small...


## Motivations

EW sector

Weak Cl
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## QCD sector

## Motivations

EW sector

Weak Cl

## QCD sector

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## Motivations

EW sector

Feak Cl

QCD sector
Strong CP

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$\mathcal{L}_{\mathrm{QCD}}^{\prime}=\mathcal{L}_{\mathrm{QCD}}+\mathcal{L}^{\mathrm{CP}}$

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Weak Cl

QCD sector
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$\mathcal{L}_{\mathrm{QCD}}^{\prime}=\mathcal{L}_{\mathrm{QCD}}+\mathcal{L}^{\mathrm{CP}}$
$\theta$-term
SMEFT operators

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SMEFT operators


Nucleon electric dipole moment

## Motivations

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$\mathcal{L}_{\mathrm{QCD}}^{\prime}=\mathcal{L}_{\mathrm{QCD}}+\mathcal{L}^{\mathrm{CP}}$
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Nucleon electric dipole moment never measured...

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P-symmetry

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QCD sector

QCD Lagrangian is assumed to be invariant under parity transformations

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$P$-violation on the internal structure of nucleons?

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Terms from EW sector

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Wealk Paiclasiom
$\checkmark$

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QCD sector
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Terms from EW sector


Terms from QCD sector

## Motivations

P-symmetry
QCD sector
QCD Lagrangian is assumed to be invariant under parity transformations

## Are there any effects of QCD

$P$-violation on the internal structure of nucleons?

Terms from EW sector
Weak Pryiolatiom

0

## Which implications could the

presence of strong P-violation cause
to inclusive DIS?

## DIS process

$$
l(\ell)+N(P) \rightarrow \gamma^{*}(q) \rightarrow l\left(\ell^{\prime}\right)+X
$$



## Cross Section

$$
\frac{d^{3} \sigma}{d x_{B} d y d \phi_{S}}=\frac{\alpha^{2} y}{2 Q^{4}} L_{\mu \nu}\left(l, l^{\prime}, \lambda_{e}\right) 2 M W^{\mu \nu}(q, P, S)
$$

In general

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In general

$$
\frac{d^{3} \sigma}{d x_{B} d y d \phi_{S}}=\frac{\alpha^{2} y}{2 Q^{4}} \sum_{j=\gamma, \gamma Z, Z} \eta^{j} L_{\mu \nu}^{(j)}\left(l, l^{\prime} ; \lambda_{e}\right) 2 M W^{\mu \nu}(q, P, S)
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& \eta^{\gamma}=1 \quad \eta^{\gamma Z}=\left(\frac{G_{F} M_{Z}^{2}}{2 \sqrt{2} \pi \alpha}\right) \frac{Q^{2}}{Q^{2}+M_{Z}^{2}} \quad \eta^{Z}=\left(\eta^{\gamma Z}\right)^{2}
\end{aligned}
$$

## Hadronic Tensor (unpolarized)

$2 M W_{\mu \nu}(q, P)=\sum_{X} \int \frac{d^{3} P_{X}}{2 E_{X}} \delta^{4}\left(P+q-P_{X}\right)\langle P| J_{\mu}^{\dagger}(0)\left|P_{X}\right\rangle\left\langle P_{X}\right| J_{\nu}(0)|P\rangle$

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Dominant contribution on the Light-Cone

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Dominant contribution on the Light-Cone

$$
2 M W^{\mu \nu}(q, P, S)=\sum_{q} e_{q}^{2} \frac{1}{2} \operatorname{Tr}\left[\Phi(q, P, S) \Gamma^{\mu} \gamma^{+} \Gamma^{\nu}\right]
$$

## Hadronic Tensor (unpolarized)



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Correlation distribution function
J. Collins, "Foundation of Perturbative QCD"
M. Anselmino et al., Z. Phys. C 64, 267 (1997)

## Hadronic Tensor (unpolarized)



P-odd structures already present in the hadronic tensor!

$$
2 M W^{\mu \nu}(q, P, S)=\sum_{q} e_{q}^{2} \frac{1}{2} \operatorname{Tr}\left[\begin{array}{|}
\Phi(q, P, S) & \left.\Gamma^{\mu} \gamma^{+} \Gamma^{\nu}\right] \\
\text { Correlation distribution function }
\end{array}\right.
$$

$$
\Phi_{i j}(k, P, S)=\int \frac{d^{4} \xi}{(2 \pi)^{4}} e^{i k \cdot \xi}\langle P| \bar{\psi}_{i}(0) U(0, \xi) \psi_{i}(\xi)|P\rangle
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Decomposition in partonic densities
J. Collins, "Foundation of Perturbative QCD"
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## Partonic correlator (unpolarized)

Integrated correlator

$$
\Phi_{i j}\left(x_{B}\right)=\int \frac{d \xi^{-}}{2 \pi} e^{i k \cdot \xi}\langle P| \bar{\psi}_{j}(0) \psi_{i}(\xi)|P\rangle_{\xi^{+}=\xi_{T}=0}
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Leading twist contributions

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i \gamma^{5}, \gamma^{\mu} \gamma^{5}, i \gamma^{5} \sigma^{\mu \nu}
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Leading twist contributions

$$
\Phi_{\mathrm{PE}}(x) \simeq \frac{1}{2} f_{1}(x) \gamma^{-}
$$

$$
\Phi_{\mathrm{PV}}(x) \simeq \frac{1}{2} g_{1}^{\mathrm{PV}}(x) \gamma^{5} \gamma^{-}
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## Partonic correlator (unpolarized)

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$\Phi_{i j}\left(x_{B}\right)=\int \frac{d \xi^{-}}{2 \pi} e^{i k \cdot \xi}\langle P| \bar{\psi}_{j}(0) \psi_{i}(\xi)|P\rangle_{\xi^{+}=\xi_{T}=0}$

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$\mathbb{1}, \gamma^{\mu}, \sigma^{\mu \nu}$
$i \gamma^{5}, \gamma^{\mu} \gamma^{5}, i \gamma^{5} \sigma^{\mu \nu}$
Leading twist contributions

$$
\begin{gathered}
\Phi_{\mathrm{PE}}(x) \simeq \frac{1}{2} f_{1}(x) \gamma^{-} \quad \Phi_{\mathrm{PV}}(x) \simeq \frac{1}{2} g_{1}^{\mathrm{PV}}(x) \gamma^{5} \gamma^{-} \\
\Phi(x)=\Phi_{\mathrm{PE}}(x)+\Phi_{\mathrm{PV}}(x)
\end{gathered}
$$

## Neutral-current DIS

$$
\begin{array}{rlr}
\frac{d \sigma^{ \pm}}{d x d y}=\frac{2 \pi \alpha^{2}}{x y Q^{2}}[ & \left(Y_{+}+\gamma^{2} y^{2} / 2\right)\left(F_{2 U U}+\lambda F_{2 L U}^{ \pm}\right) & \\
& -y^{2}\left(F_{L, U U}+\lambda F_{L, L U}^{ \pm}\right) & \frac{d \sigma^{ \pm}}{d x d y}=\frac{2 \pi \alpha^{2}}{x y Q^{2}}\left[Y_{+} F_{2}^{ \pm}-y^{2} F_{L}^{ \pm} \mp Y_{-} x F_{3}^{ \pm}\right] \\
& \left.-\frac{Y_{-}}{\sqrt{1+\gamma^{2}}}\left(x F_{3 U U}^{ \pm}+\lambda x F_{3 L U}\right)\right] &
\end{array}
$$

## Focus: structure function $x F_{3}\left(x, Q^{2}\right)$

$$
x F_{3 L U}\left(x, Q^{2}\right)=x F_{3}^{(\gamma)}-g_{V}^{e} \eta_{\gamma Z} x F_{3}^{(\gamma Z)}+\left(g_{V}^{e}{ }^{2}+g_{A}^{e}{ }^{2}\right) \eta_{Z} x F_{3}^{(Z)}
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$$

$$
\begin{aligned}
x F_{3}^{(\gamma)}\left(x, Q^{2}\right) & =0 \\
x F_{3}^{(\gamma Z)}\left(x, Q^{2}\right) & =\sum_{q} 2 e_{q} g_{A}^{q} x f_{1}^{(q-\bar{q})} \\
x F_{3}^{(Z)}\left(x, Q^{2}\right) & =\sum_{q} 2 g_{V}^{q} g_{A}^{q} x f_{1}^{(q-\bar{q})}
\end{aligned}
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x F_{3}^{(Z)}\left(x, Q^{2}\right) & =\sum_{q} 2 g_{V}^{q} g_{A}^{q} x f_{1}^{(q-\bar{q})}
\end{aligned}
$$

$$
\begin{aligned}
x \Delta F_{3}^{(\gamma)}\left(x, Q^{2}\right) & =-\sum_{q} e_{q}^{2} x g_{1}^{\mathrm{PV}(q+\bar{q})} \\
x \Delta F_{3}^{(\gamma Z)}\left(x, Q^{2}\right) & =-\sum_{q} 2 e_{q} g_{V}^{q} x g_{1}^{\mathrm{PV}(q+\bar{q})} \\
x \Delta F_{3}^{(Z)}\left(x, Q^{2}\right) & =-\sum_{q}\left(g_{V}^{q 2}+g_{A}^{q 2}\right) x g_{1}^{\mathrm{PV}(q+\bar{q})}
\end{aligned}
$$

Additional contributions due to the new PV parton distribution

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$$

$$
x F_{3}^{(\gamma)}\left(x, Q^{2}\right)=0
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$$



## MAIN INNOVATION OF PV-HYPOTESIS



$$
x \Delta F_{3}^{(\gamma)}\left(x, Q^{2}\right)=-\sum_{q} e_{q}^{2} x g_{1}^{\mathrm{PV}(q+\bar{q})}
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Additional contributions due to the new PV parton distribution

## Neutral-current DIS

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\begin{aligned}
\frac{d \sigma^{ \pm}}{d x d y}=\frac{2 \pi \alpha^{2}}{x y Q^{2}}[ & \left(Y_{+}+\gamma^{2} y^{2} / 2\right)\left(F_{2 U U}+\lambda F_{2 L U}^{ \pm}\right) \\
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## Standard DIS structure functions

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## Standard DIS structure functions

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\begin{aligned}
F_{2 U U}\left(x, Q^{2}\right) & =F_{2}^{(\gamma)}-g_{V}^{e} \eta_{\gamma Z} F_{2}^{(\gamma Z)}+\left(g_{V}^{e}{ }^{2}+g_{A}^{e}{ }^{2}\right) \eta_{Z} F_{2}^{(Z)} \\
F_{2 L U}^{ \pm}\left(x, Q^{2}\right) & =\mp g_{A}^{e} \eta_{\gamma Z} F_{2}^{(\gamma Z)} \pm 2 g_{V}^{e} g_{A}^{e} \eta_{Z} F_{2}^{(Z)} \\
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\end{aligned}
$$

## Phenomenology

## Experimental observable

## PVDIS Asymmetry

$$
A_{\mathrm{PV}} \equiv \frac{d \sigma(\lambda=1)-d \sigma(\lambda=-1)}{d \sigma(\lambda=1)+d \sigma(\lambda=-1)}
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$$

PVDIS Collaboration, Nature 506 (2014)
D. Wang et al., Phys.Rev.C 91 (2015)

$$
Y_{ \pm}=1 \pm(1-y)^{2}
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## Experimental observable

PVDIS Asymmetry

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A_{\mathrm{PV}} \equiv \frac{d \sigma(\lambda=1)-d \sigma(\lambda=-1)}{d \sigma(\lambda=1)+d \sigma(\lambda=-1)}
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$$
=\frac{Y_{+} F_{2 L U}-y^{2} F_{L, L U}-Y_{-} x F_{3 L U}}{Y_{+}\left(F_{2 U U}\right)-y^{2}\left(F_{L, U U}-Y-x F_{3 U U}\right.}
$$

$$
Y_{ \pm}=1 \pm(1-y)^{2}
$$

Contribution of $g_{1}^{P V}$ in each of the structure functions due to
$\gamma Z$ and $Z$ channels

## Available experimental data

HERA dataset<br>(Run I + II combined)

H1 Collaboration, Eur. Phys. J. C 78 (2018)

## Available experimental data

HERA dataset
(Run I + II combined)
H1 Collaboration, Eur. Phys. J. C 78 (2018)
$e^{+}$asymmetry: 136 data
$e^{-}$asymmetry: 138 data


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## JLab6 PVDIS dataset

PVDIS Collaboration, Nature 506 (2014)
D. Wang et al., Phys.Rev.C 91 (2015)

## Available experimental data

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$e^{+}$asymmetry: 136 data
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$e^{-}$asymmetry: 2 data

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## SLAC-E122 dataset

C.Y. Prescott et al., Phys. Lett. B (1979)

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$e^{+}$asymmetry: 136 data
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$e^{-}$asymmetry: 11 data

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Imbalance between information from electron and positron beams
$\mathrm{Q}^{2}\left[\mathrm{GeV}^{2}\right]$
$e^{-}$asymmetry: 2 data
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## Experimental data: energy range

HERA dataset

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HERA dataset
$Q^{2} \in(200,30000) \mathrm{GeV}^{2}$

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HERA dataset

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Q^{2} \in(200,30000) \mathrm{GeV}^{2}
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high-energy
$Q^{2} \gg M_{N}^{2}$
no need of modification of the theory

## Experimental data: energy range

HERA dataset

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high-energy
$Q^{2} \gg M_{N}^{2}$
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JLab6 + SLAC-E122 datasets

## Experimental data: energy range

HERA dataset
high-energy
$Q^{2} \gg M_{N}^{2}$
no need of modification of the theory
JLab6 + SLAC-E122 datasets
low-energy
$Q^{2} \simeq M_{N}^{2}$

$$
Q^{2} \in(200,30000) \mathrm{GeV}^{2}
$$

$$
Q^{2} \in(0.9,1.9) \mathrm{GeV}^{2}
$$

## Experimental data: energy range

HERA dataset
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\begin{aligned}
& C_{1 u}=2 g_{A}^{e} g_{V}^{u}=2\left(-\frac{1}{2}\right)\left(\frac{1}{2}-\frac{4}{3} \sin ^{2} \theta_{W}\right) \\
& C_{2 u}=2 g_{V}^{e} g_{A}^{u}=2\left(-\frac{1}{2}+2 \sin ^{2} \theta_{W}\right)\left(\frac{1}{2}\right) \\
& C_{1 d}=2 g_{A}^{e} g_{V}^{d}=2\left(-\frac{1}{2}\right)\left(-\frac{1}{2}+\frac{2}{3} \sin ^{2} \theta_{W}\right) \\
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\begin{aligned}
& C_{1 u}^{\mathrm{SM}}=-0.1887-0.0011 \times \frac{2}{3} \ln \left(\left\langle Q^{2}\right\rangle / 0.14 \mathrm{GeV}^{2}\right) \\
& C_{1 d}^{\mathrm{SM}}=0.3419-0.0011 \times \frac{-1}{3} \ln \left(\left\langle Q^{2}\right\rangle / 0.14 \mathrm{GeV}^{2}\right) \\
& C_{2 u}^{\mathrm{SM}}=-0.0351-0.0009 \ln \left(\left\langle Q^{2}\right\rangle / 0.078 \mathrm{GeV}^{2}\right) \\
& C_{2 d}^{\mathrm{SM}}=0.0248+0.0007 \ln \left(\left\langle Q^{2}\right\rangle / 0.021 \mathrm{GeV}^{2}\right)
\end{aligned}
$$

Parameterization of $g_{1}^{P V}\left(x, Q^{2}\right)$

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PV parton density comes from the structure
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\end{array}\right] \begin{aligned}
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## Error propagation

PDF set for

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## Statistical distribution of

 100 values of parameter $\alpha$
## Results of the fit: $\chi^{2}$ values

Fit WITH EW radiative corrections

|  | $N$ of points | $x^{2} / N_{\text {data }}(\mathrm{SM})$ | $X^{2 /} \mathrm{N}_{\text {data }}$ (Fit) |
| :---: | :---: | :---: | :---: |
| HERA $A^{+}$ | 136 | 1.12 | 1.12 |
| HERA $A^{-}$ | 138 | 0.98 | 0.98 |
| JLab6 $A^{-}$ | 2 | 0.67 | 0.42 |
| SLAC-E122 $A^{-}$ | 11 | 0.97 | 0.94 |
| TOTAL | 287 | 1.042 | 1.037 |

## Results of the fit: data-theory comparison



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Very small uncertainties in the predictions because the fit is dominated by data with smaller errors

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There's room for a better description for positron asymmetry at low- Q


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There's room for a better description for positron asymmetry at low- Q

Agreement for electron asymmetry, but too large errors at low-Q

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Sizeable improvement of the fit w.r.t. SM predictions

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Old dataset with still quite large experimental errors ( $>20 \%$ )

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Data points which actually drive the fit due to very small experimental errors ( $\sim \%$ )

## Results of the fit: $g_{1}^{P V}\left(x, Q^{2}\right)$ extraction

$$
\begin{gathered}
g_{1}^{\mathrm{PV}}(x)=\alpha g_{1}(x) \\
\alpha=(-1.01 \pm 0.66) \cdot 10^{-4}
\end{gathered}
$$




## Conclusions and Outlook

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## Conclusions and Outlook

- Predictions of the size of the PV distribution can be made in the kinematic domains of JLab12, JLab20+(?) and EIC



## Conclusions and Outlook

- Further investigations on a new P-odd, CP-odd distribution function arising when considering the polarisation of the target

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\begin{aligned}
\Phi^{q}\left(x, Q^{2}\right)=\{ & f_{1}^{q}\left(x, Q^{2}\right)+g_{1}^{\mathrm{PV} q}\left(x, Q^{2}\right) \gamma_{5} \\
& +S_{L}\left(g_{1}^{q}\left(x, Q^{2}\right) \gamma_{5}+f_{1 L}^{\mathrm{PV} q}\left(x, Q^{2}\right)\right) \\
& \left.-\$_{T}\left(h_{1}^{q}\left(x, Q^{2}\right) \gamma_{5}-e_{1 T}^{\mathrm{PV} q}\left(x, Q^{2}\right)\right)\right\} \frac{\not x_{+}}{2}
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$\Delta x_{B} g_{5}\left(x_{B}, Q^{2}\right) \approx \Delta x_{B} g_{5}^{(\gamma)}\left(x_{B}, Q^{2}\right)=\frac{1}{2} \sum_{q} e_{q}^{2} x_{B} f_{1 L}^{\mathrm{PV}(q-\bar{q})}$

