

HADRON2023

June 8th, 2023
Genova

Modeling spin effects in electron-positron annihilation to hadrons

Albi Kerbizi

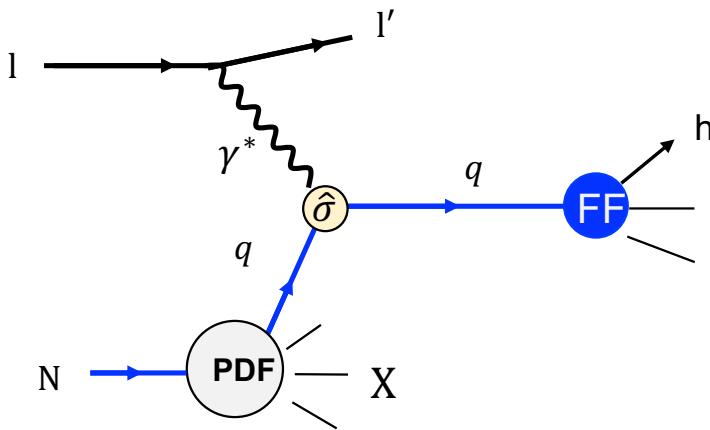
University of Trieste and INFN Trieste

work done in the context of the POLFRAG project



Istituto Nazionale di Fisica Nucleare

Nucleon structure and hadronization



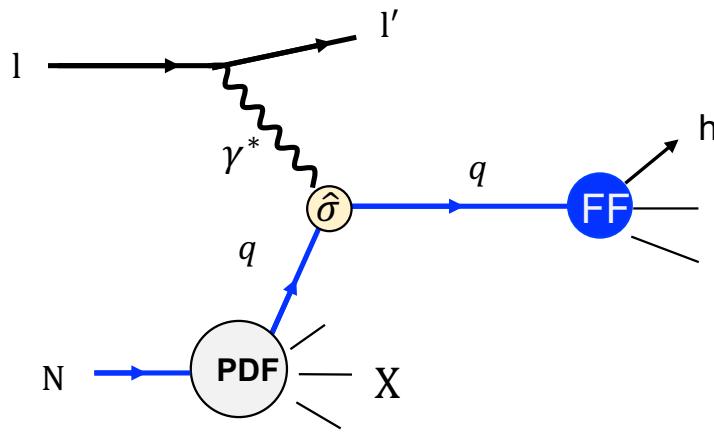
Semi-inclusive deep inelastic scattering (SIDIS)
powerful tool used to study the partonic structure of
nucleons

Couples PDFs and fragmentation functions (FFs)

transverse spin structure → involves the
fragmentation of transversely polarized quarks
described by the **Collins FF $H_{1q}^{\perp h}$**

Collins, NPB 396, 161 (1993).

Nucleon structure and hadronization



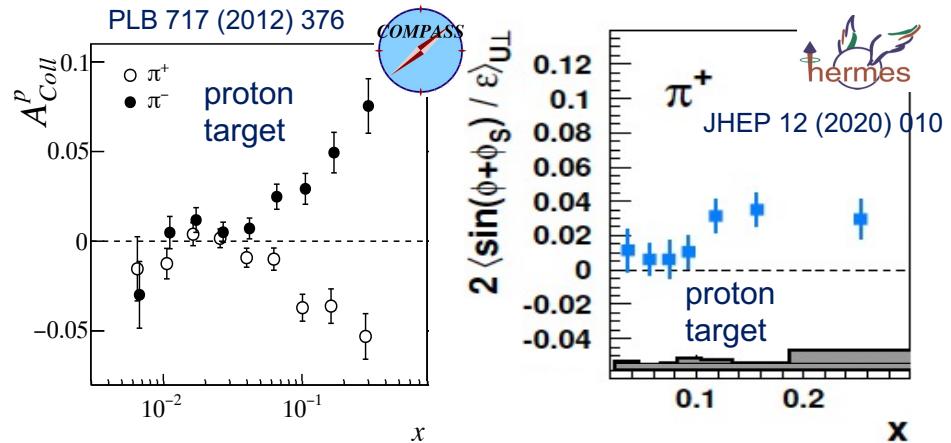
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Example: Collins asymmetry

$$A_{UT}^{\sin \phi_h + \phi_S - \pi} = \frac{\sum_q e_q^2 h_1^q \otimes H_{1q}^{\perp h}}{\sum_q e_q^2 f_1^q \otimes D_{1q}^h}$$

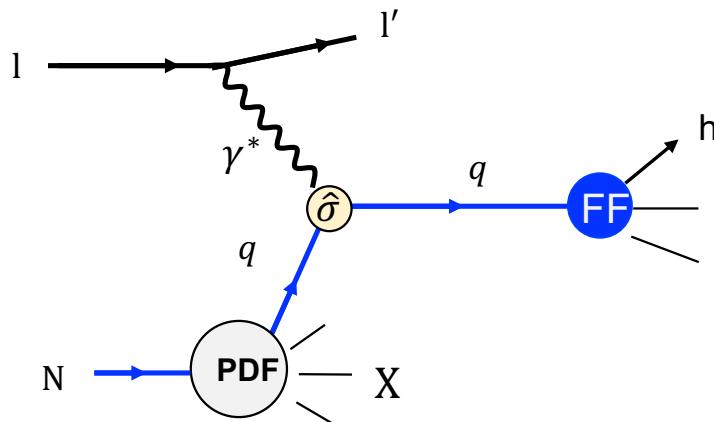
$h_1^q \rightarrow$ transversity PDF

transverse polarization of quarks in a transversely polarized nucleon

Measured by HERMES (p), COMPASS (p,d), Jlab (n)

To extract transversity the Collins FF is needed

Nucleon structure and hadronization



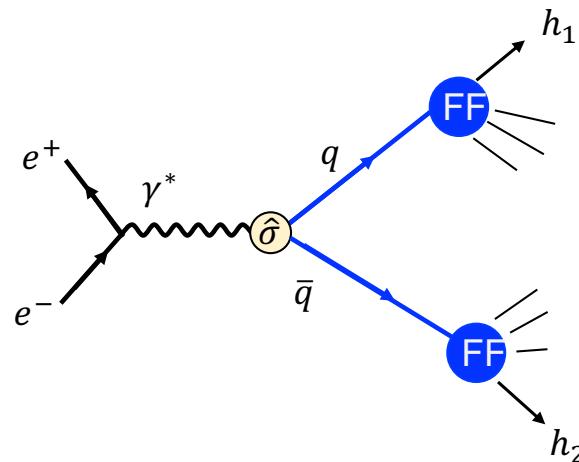
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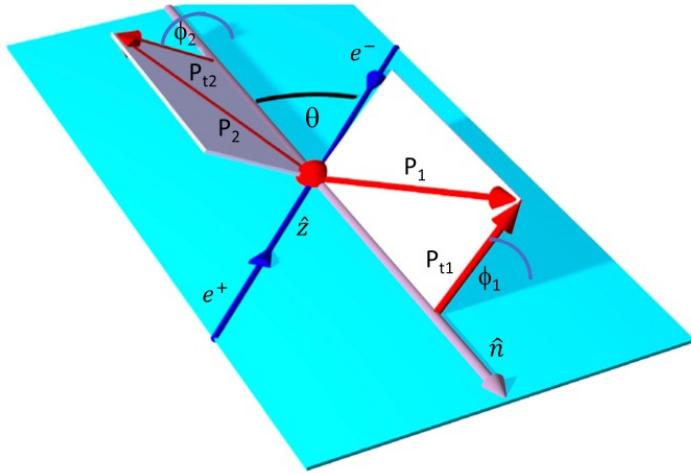
e^+e^- annihilation to hadrons
 an important tool to study hadronization

q and \bar{q} have correlated transverse polarizations

→ access to **Collins FF $H_{1q}^{\perp h}$** via the Collins asymmetry in e^+e^-

other FFs as well: unpolarized, interference..

Nucleon structure and hadronization



Cross section for the production of two back-to-back hadrons h_1 and h_2

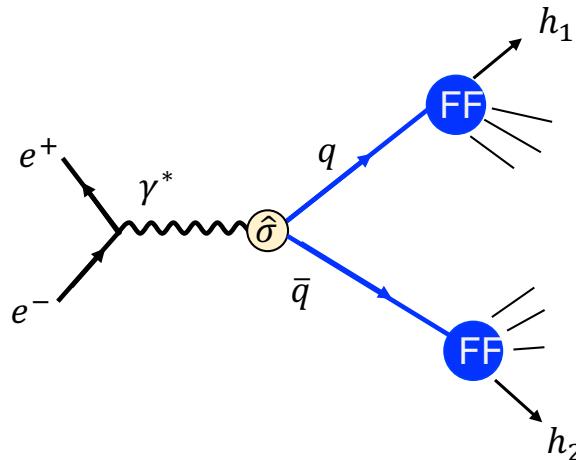
Boer, NPB, 806:23–67, 2009
D'Alesio et al., JHEP 10 (2021) 078

$$d\sigma^{e^+e^- \rightarrow h_1 h_2 X} \propto 1 + \frac{\sin^2 \theta}{1 + \cos^2 \theta} A_{12} \cos(\phi_1 + \phi_2)$$

Collins asymmetry

$$A_{12}(z_1, z_2, P_{T1}, P_{T2}) = \frac{\sum_q e_q^2 H_{1q}^{\perp h_1}(z_1, P_{T1}) \times H_{1\bar{q}}^{\perp h_2}(z_2, P_{T2})}{\sum_q e_q^2 D_{1q}^{h_1}(z_1, P_{T1}) \times D_{1\bar{q}}^{h_2}(z_2, P_{T2})}$$

$$z_i = 2E_{h_i}/\sqrt{s}$$



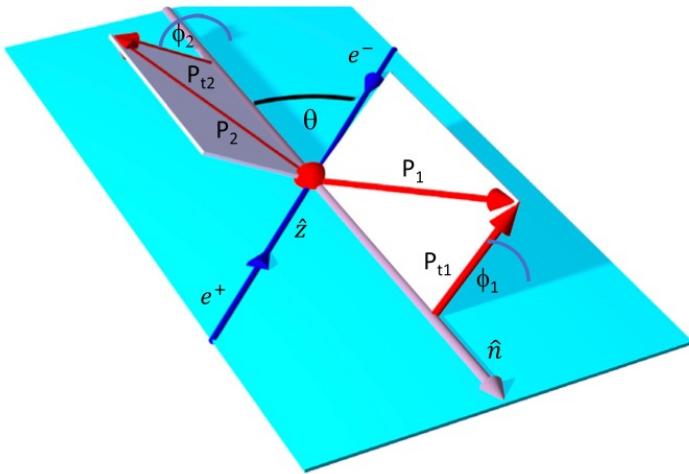
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Nucleon structure and hadronization



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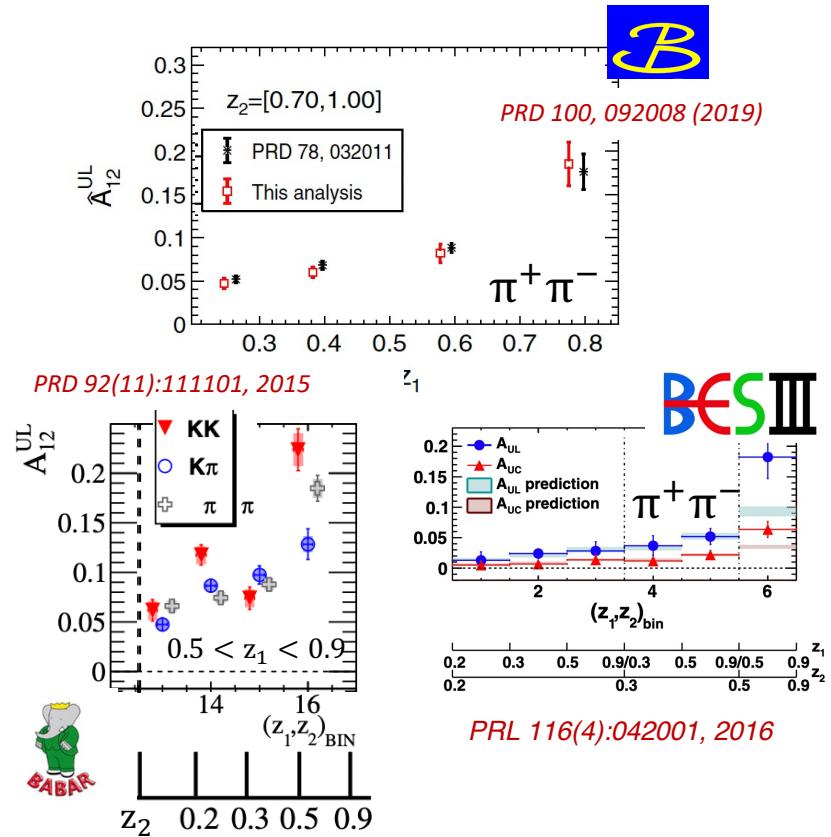
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$$z_i = 2E_{h_i}/\sqrt{s}$$



many measurements by *BELLE*, *BABAR*, *BESIII*

Used for the extractions of transversity PDFs

Anselmino et al, PRD 92 (11) (2015) 114023
Martin et al., PRD 91(1):014034, 2015
Kang et al., PRD 93 (1) (2016) 014009
...

Benchmark for hadronization models!

Modeling hadronization

We have developed a model for the simulation of the fragmentation of a quark with a given polarization → string+ 3P_0 model

AK, Artru, Belghobsi, Bradamante, Martin, PRD 97, 074010 (2018)	PS mesons
AK, Artru, Belghobsi, Martin, PRD 100, 014003 (2019)	PS mesons
AK, Artru, Martin, PRD 104, 114038 (2021)	PS + VM

Implemented in Pythia for SIDIS → StringSpinner

AK, L. Lönnblad, CPC 272 (2022) 108234	PS, Pythia 8.2
AK, L. Lönnblad, arXiv: 2305.05058	PS + VM, Pythia 8.3

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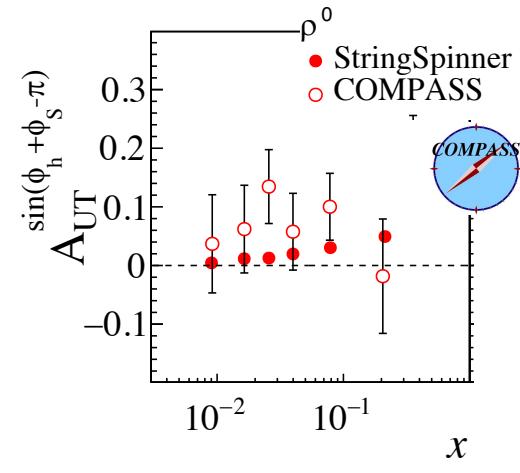
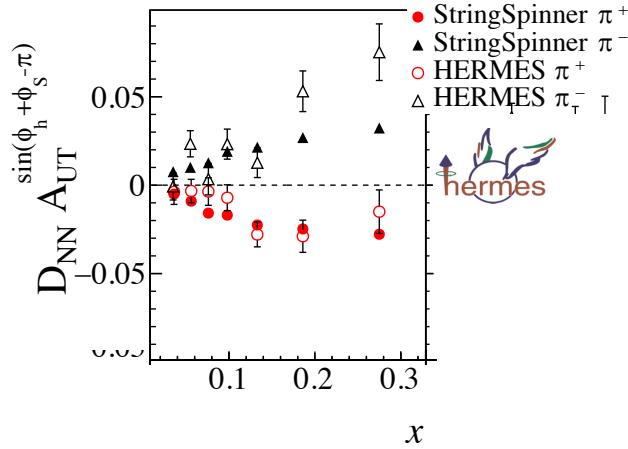
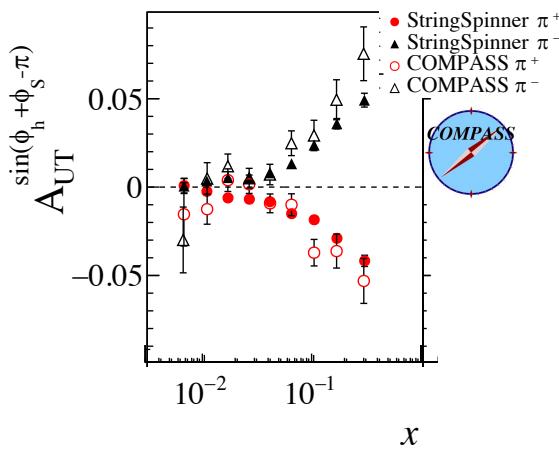
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PS + VM, Pythia 8.3



Promising results for SIDIS! (more results in arXiv: 2305.05058)

A similar work for e^+e^- annihilation does not exist

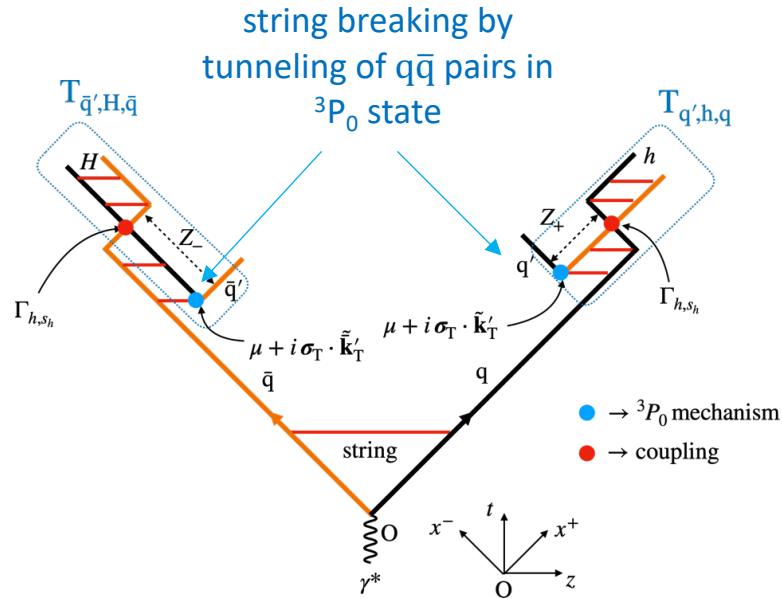
requires extension of string+ 3P_0 model → this talk

In the following slides

- i) recall of the string+ 3P_0 model
- ii) recipe for the simulation of $e^+e^- \rightarrow q\bar{q} \rightarrow \text{hadrons}$
application of the model to the hadronization of a quark-antiquark pair with correlated spin states

in collaboration with X. Artru

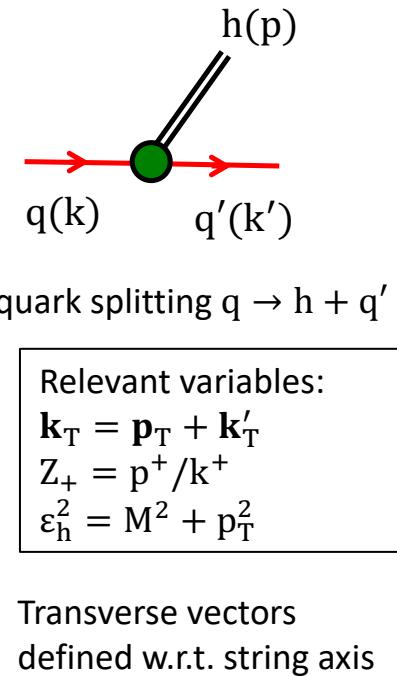
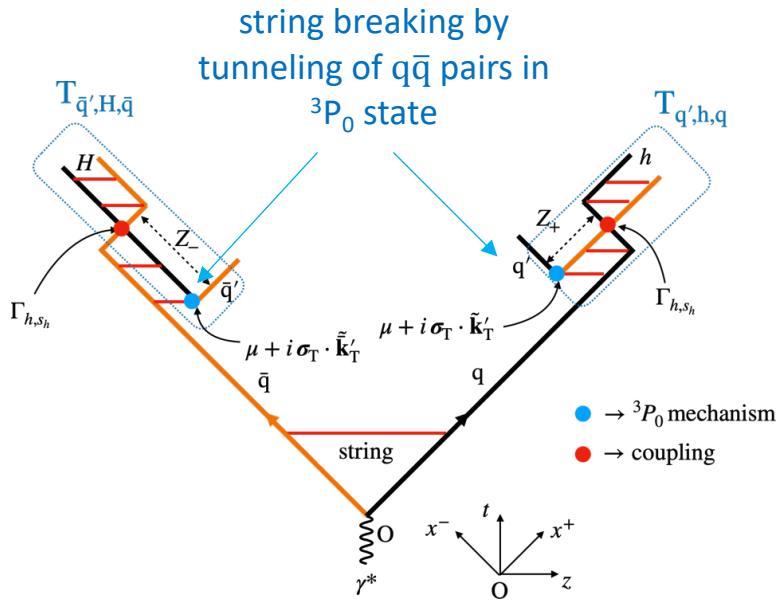
The hadronization model: string+ 3P_0



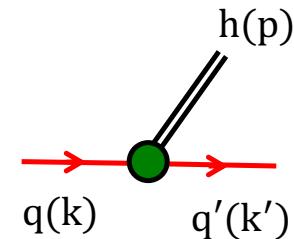
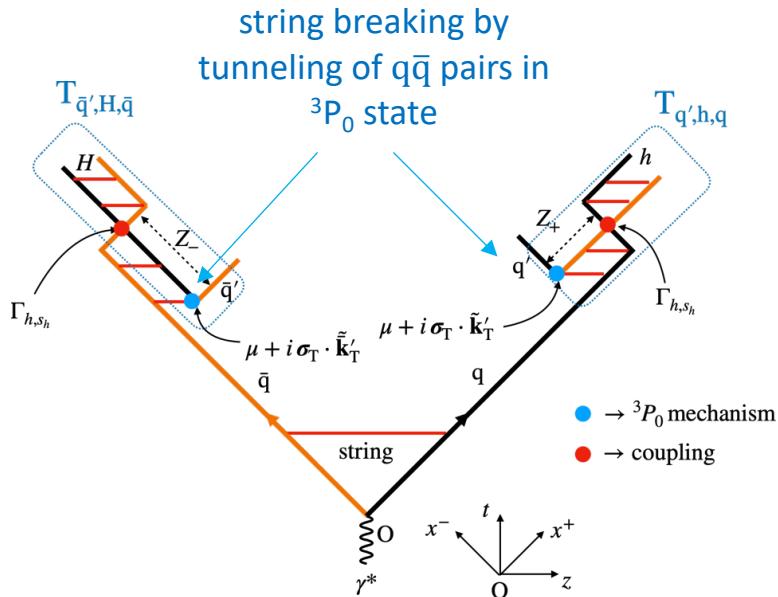
the hadronization of the $q\bar{q}$ pair is described in the string fragmentation framework supplemented with the 3P_0 model of quark tunneling

extension of the Lund string Model (Pythia)

The hadronization model: string+ 3P_0



The hadronization model: string+ 3P_0



quark splitting $q \rightarrow h + q'$

Relevant variables:
 $\mathbf{k}_T = \mathbf{p}_T + \mathbf{k}'_T$
 $Z_+ = p^+/k^+$
 $\varepsilon_h^2 = M^2 + p_T^2$

Transverse vectors
defined w.r.t. string axis

Quark splitting amplitude in the string+ 3P_0 model

$$T_{q',h,q} \propto C_{q',h,q} D_h(M^2) \left(\frac{1 - Z_+}{\varepsilon_h^2} \right)^{\frac{a}{2}} \underbrace{\exp \left[-\frac{b_L \varepsilon_h^2}{2Z_+} \right]}_{\text{longitudinal momentum}} N_a^{-\frac{1}{2}}(\varepsilon_h^2) e^{-\frac{b_T k'^2_T}{2}} \underbrace{\text{transverse momentum}}_{\text{(w.r.t string axis)}}$$

flavor mass

Free param. Lund

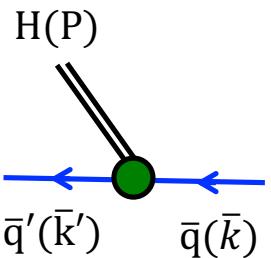
Free param. string+ 3P_0

$[\mu + \sigma_z \sigma_T \cdot \mathbf{k}'_T]$
 3P_0 mechanism
 $[\mu \text{ complex mass parameter}]$

Γ_{h,s_h}

Coupling
e.g.
 $\Gamma_{h=PS} = \sigma_z$

AK, Artru, Martin, PRD 104, 114038 (2021)



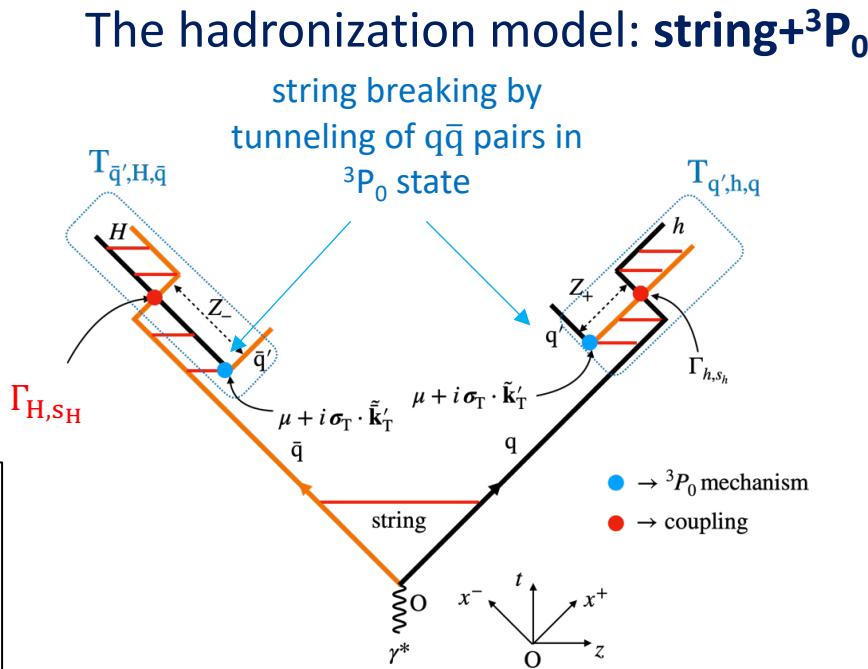
antiquark splitting
 $\bar{q} \rightarrow H + \bar{q}'$

Relevant variables:

$$\bar{\mathbf{k}}_T = \mathbf{P}_T + \bar{\mathbf{k}}'_T$$

$$Z_- = P^- / \bar{k}^-$$

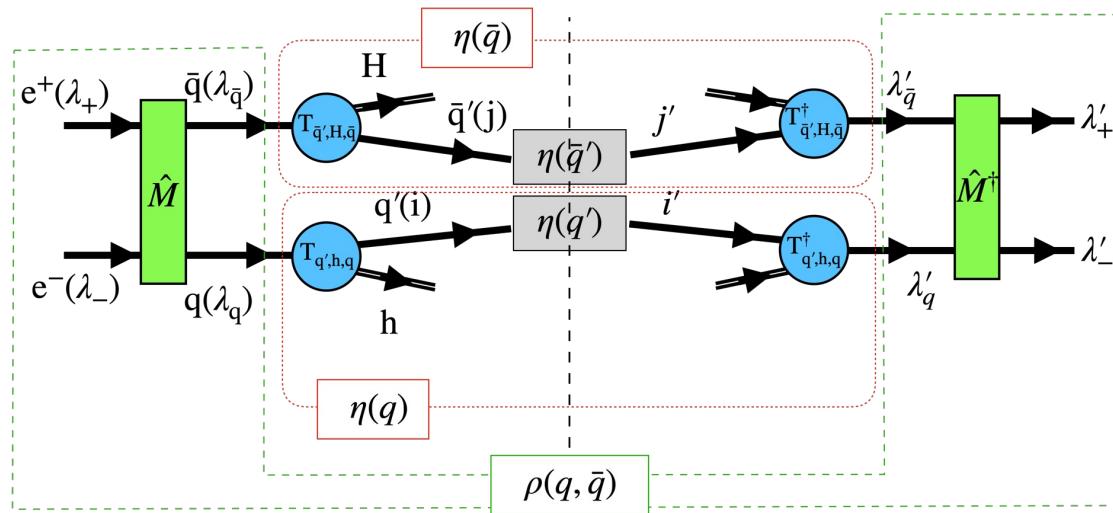
$$\varepsilon_H^2 = M^2 + P_T^2$$



Antiquark splitting amplitude in the string+ 3P_0 model obtained by the quark one by

$$\{q, h, q'\} \rightarrow \{\bar{q}, H, \bar{q}'\}, \quad Z_+ \rightarrow Z_-, \quad \{\mathbf{k}_T, \mathbf{p}_T, \mathbf{k}'_T\} \rightarrow \{\bar{\mathbf{k}}_T, \mathbf{P}_T, \bar{\mathbf{k}}'_T\}$$

The ingredients needed for e^+e^- annihilation



Squared amplitude associated to the unitarity diagram (factorized approach)

$$|A(e^+e^- \rightarrow h H X)|^2 = \langle |M|^2 \rangle \times \text{Tr}_{q\bar{q}} \rho(q, \bar{q}) \eta(q) \otimes \eta(\bar{q})$$

Spin-averaged squared matrix element for the hard scattering $e^+e^- \rightarrow q\bar{q}$

Joint spin density matrix of $q\bar{q}$

$$\rho(q, \bar{q}) = C_{\alpha\beta}^{q\bar{q}} \sigma_q^\alpha \otimes \sigma_{\bar{q}}^\beta$$

implements spin correlations

$$C_{\alpha\beta}^{q\bar{q}} \rightarrow \text{correlation coefficients}$$

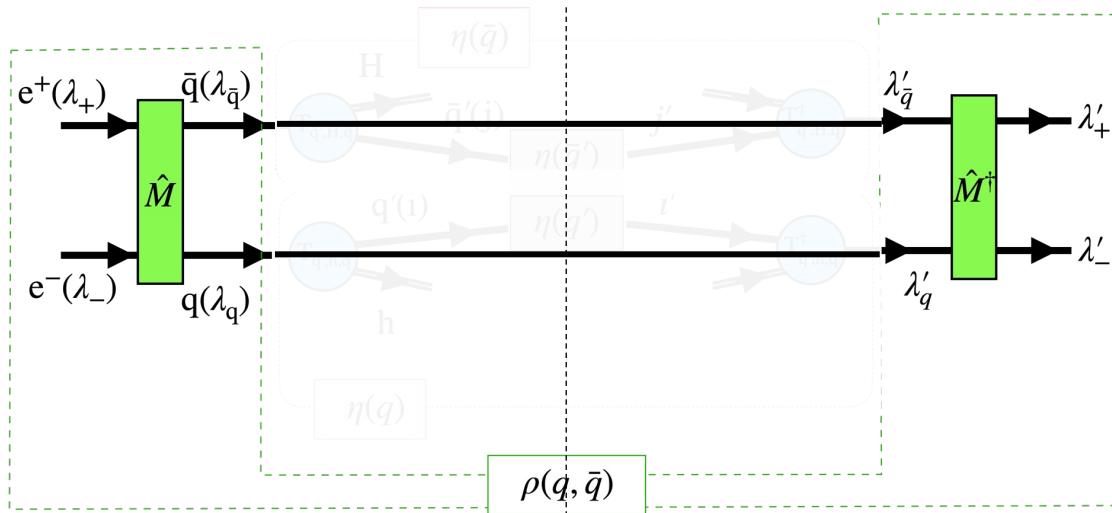
Artru et al., Phys. Rept., 470:1–92, 2009

Acceptance matrix of q
information coming to q from
“future” emissions

$$\eta(q) = \begin{cases} 1_q & (\text{no info}) \\ T_{q',h,q}^\dagger \eta(q') T_{q',h,q} & q \rightarrow h + q' \end{cases}$$

similarly for $\eta(\bar{q})$

The recursive recipe for simulating e^+e^- annihilation



Integrate over the emissions of q and \bar{q} and set up the pair

no info to q or $\bar{q} \rightarrow \eta(q) = 1_q, \eta(\bar{q}) = 1_{\bar{q}}$

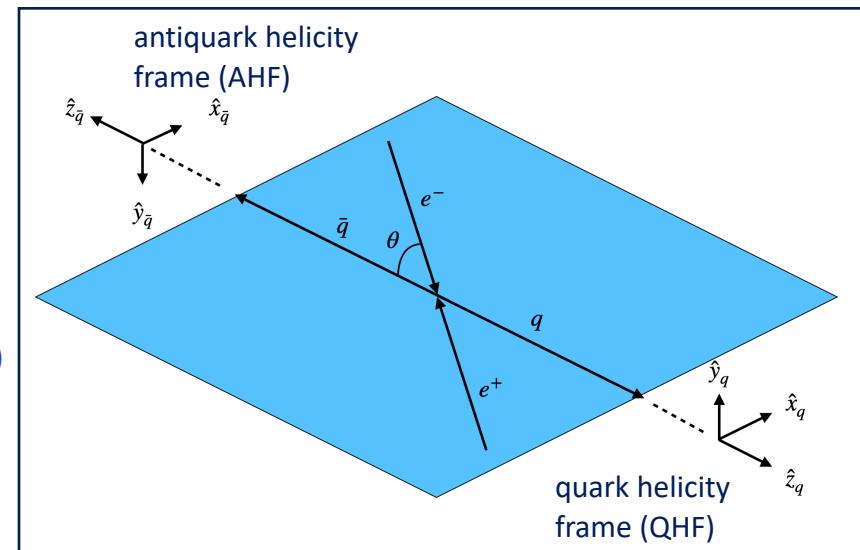
i. generate the quark flavors and kinematics using

$$d\hat{\sigma}(q\bar{q})/d\cos\theta \propto \langle |\hat{M}|^2 \rangle$$

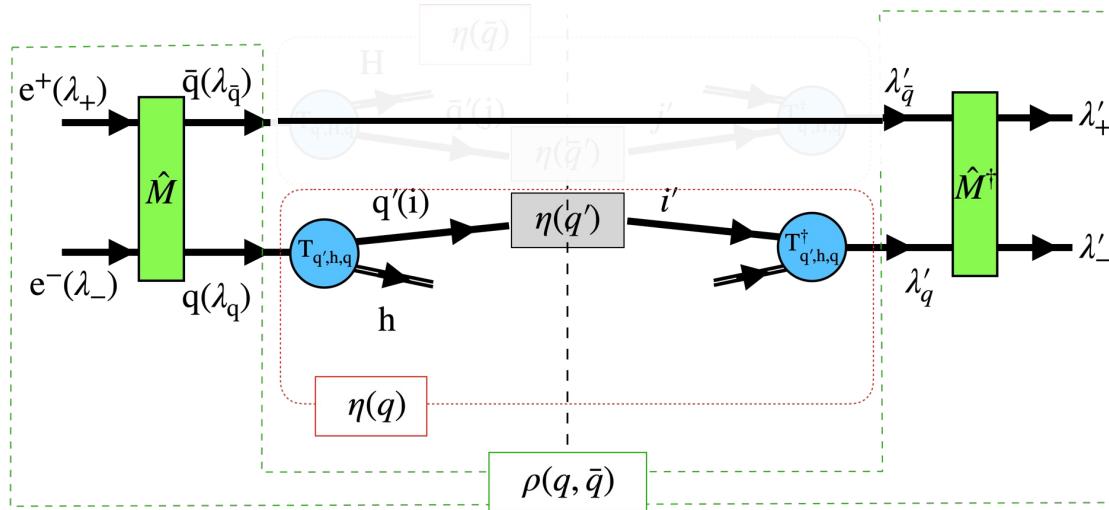
ii. set up the joint spin density matrix

$$\begin{aligned} \rho(q, \bar{q}) &\propto 1_q \otimes 1_{\bar{q}} - \sigma_q^z \otimes \sigma_{\bar{q}}^z \\ &+ \sin^2\theta [\sigma_q^x \otimes \sigma_{\bar{q}}^x + \sigma_q^y \otimes \sigma_{\bar{q}}^y]/(1 + \cos^2\theta) \end{aligned}$$

$\sigma_a^\alpha \rightarrow$ Pauli matrices along $\alpha = 0, x, y, z$
in the helicity frame of $a = q, \bar{q}$



The recursive recipe for simulating e^+e^- annihilation



Emit the first hadron the q side
assuming $h=PS$ from

$$\rightarrow \eta(q) = T_{q',h,q}^\dagger \eta(q') T_{q',h,q} \quad (\text{no info to } q' \rightarrow \eta(q') = 1_{q'})$$

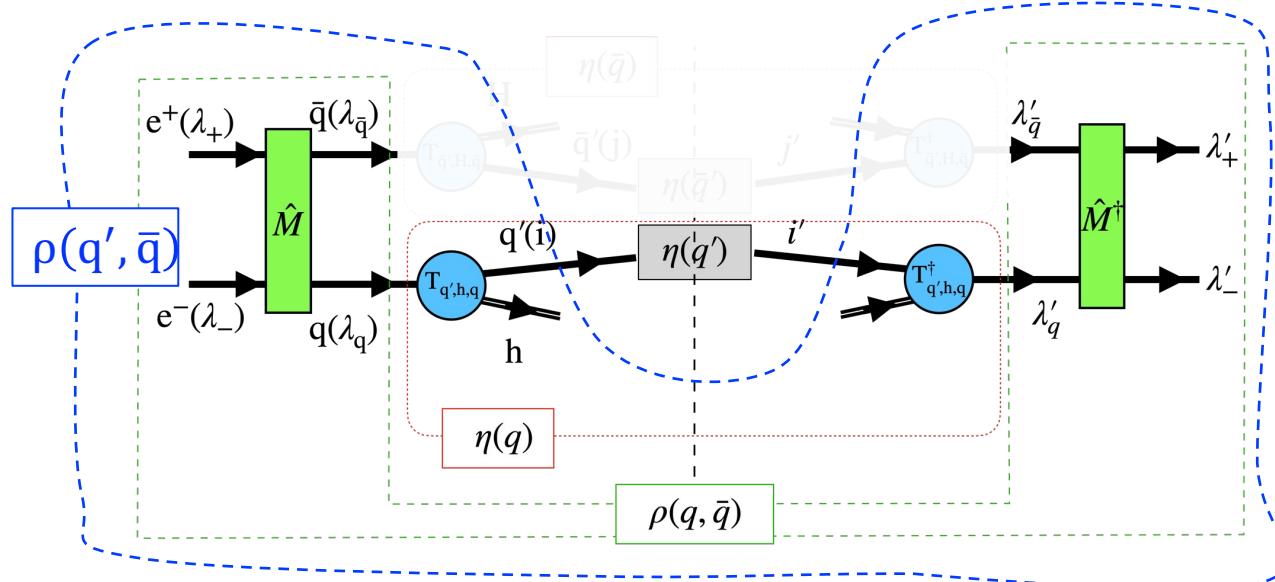
Emission probability density (splitting function)

$$\frac{dP(q \rightarrow h + q'; q\bar{q})}{dZ_+ Z_+^{-1} d^2 p_T} = \text{Tr}_{q'\bar{q}} T_{q',h,q} \rho(q, \bar{q}) T_{q',h,q}^\dagger = F_{q',h,q}(Z_+, p_T; k_T, C^{q\bar{q}})$$

$$T_{q',h,q} \equiv T_{q',h,q} \otimes 1_{\bar{q}}$$

emission of a vector meson $h=VM$ more involved (but similar steps)
→ see backup

The recursive recipe for simulating e^+e^- annihilation

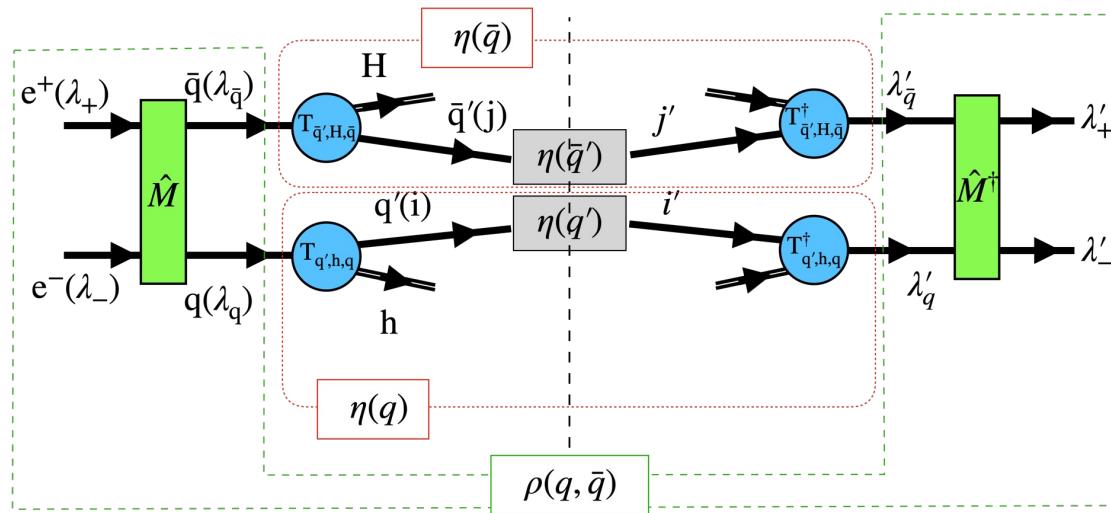


Evaluate the spin density matrix $\rho(q', \bar{q})$

$$\rho(q', \bar{q}) = T_{q',h,q} \rho(q, \bar{q}) T_{q',h,q}^\dagger$$

includes the information on the emission of h

The recursive recipe for simulating e^+e^- annihilation



Evaluate the spin density matrix $\rho(q'\bar{q})$

$$\rho(q', \bar{q}) = \mathbf{T}_{q', h, q} \rho(q, \bar{q}) \mathbf{T}_{q', h, q}^\dagger$$

includes the information on the emission of h

Emit a hadron H , e.g. $H = PS$, from the antiquark

$$\rightarrow \eta(\bar{q}) = \mathbf{T}_{\bar{q}', h, \bar{q}}^\dagger \eta(\bar{q}') \mathbf{T}_{\bar{q}', h, \bar{q}} \quad (\text{no info to } \bar{q}' \rightarrow \eta(\bar{q}') = 1_{\bar{q}'})$$

which gives the emission probability density

$$\frac{dP(\bar{q} \rightarrow H + \bar{q}'; q'\bar{q})}{dZ_- Z_-^{-1} d^2 P_T} = \text{Tr}_{q'\bar{q}'} \mathbf{T}_{\bar{q}', h, \bar{q}} \rho(q', \bar{q}) \mathbf{T}_{\bar{q}', h, \bar{q}}^\dagger = F_{\bar{q}', h, \bar{q}}(Z_-, P_T; \bar{k}_T, C^{q'\bar{q}})$$

conditional probability of emitting H , having emitted h

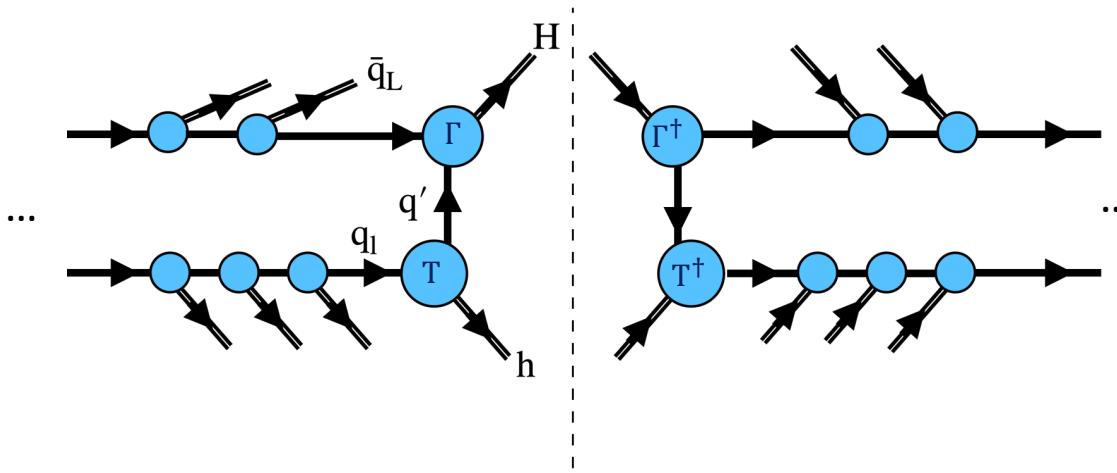
→ correlations between their transverse momenta

[Collins NPB, 304:794–804, 1988, Knowles NPB, 310:571–588, 1988]

Depend on the azimuthal angle h

Expressed in the AHF

The recursive recipe for simulating e^+e^- annihilation: exit condition



After several emissions hadronize the last pair $q_l\bar{q}_L$
joint spin-density matrix $\rho(q_l, \bar{q}_L)$

Emit the hadron $h = q_l\bar{q}'$ from q_l and project $\bar{q}_L q'$ to the state H

$$dP(q_l \rightarrow h + q'; q_l \bar{q}_L) = \text{Tr}_{q' \bar{q}_L} [T_{q', h, q_l} \otimes \Gamma_{H, s_H}] \quad \rho(q_l, \bar{q}_L) \quad [T_{q', h, q_l}^\dagger \otimes \Gamma_{H, s_H}^\dagger]$$

or emit the hadron $H = q'\bar{q}_L$ from \bar{q}_L and project $q_l\bar{q}'$ to the state h

Application of the recipe to the first two hadrons produced

Application of the recipe to $e^+e^- \rightarrow h H X$

$h = PS$ and $H = PS$ being the first two hadrons produced

$$dP(e^+e^- \rightarrow h H X) = \hat{\sigma}^{-1} \frac{d\hat{\sigma}}{d\cos\theta} \times F_{q',h,q}(Z_+, \mathbf{p}_T; \mathbf{k}_T, C^{q\bar{q}}) \times F_{\bar{q}',H,\bar{q}}(Z_-, \bar{\mathbf{p}}_T; \bar{\mathbf{k}}_T, C^{q'\bar{q}})$$
$$\text{Prob}(e^+e^- \rightarrow q\bar{q}) \quad \text{Prob}(q \rightarrow h + q') \quad \text{Prob}(\bar{q} \rightarrow H + \bar{q}'; q \rightarrow h + q')$$
$$\propto (1 + \cos^2\theta) \times (\dots) \times [1 + \frac{\sin^2\theta}{1 + \cos^2\theta} \frac{2\text{Im}(\mu)p_T}{|\mu|^2 + p_T^2} \frac{2\text{Im}(\mu)P_T}{|\mu|^2 + P_T^2} \cos(\phi_h + \phi_H)]$$

expected form for the azimuthal distribution of back-to-back hadrons!

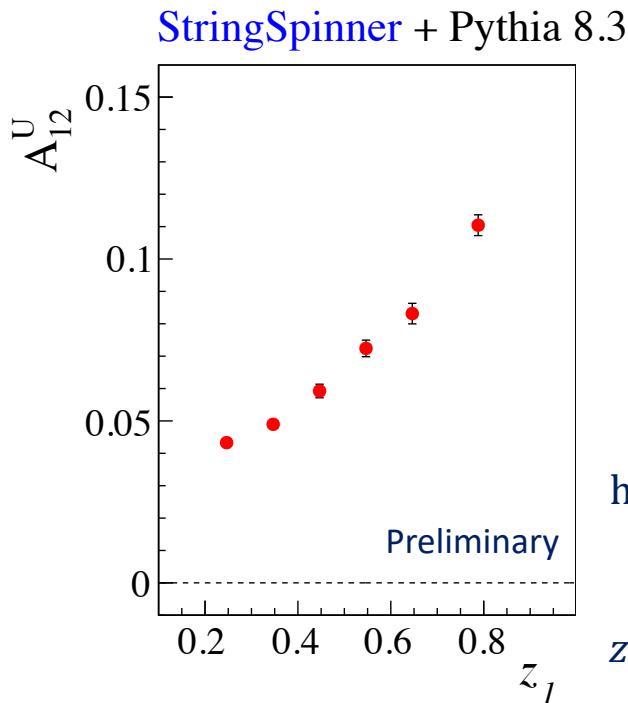
For quantitative results and phenomenology
→ implementation of the model in Pythia 8.3 for of $e^+e^- \rightarrow$ hadrons

we are extending the StringSpinner package (currently for SIDIS) to simulate also e^+e^-

ongoing work in collaboration with L. Lönnblad and A. Martin

Example of simulation results with Pythia 8.3

Collins asymmetry for back-to-back π^+ and π^- in $e^+e^- \rightarrow u\bar{u} \rightarrow \pi^+\pi^-X$ at $\sqrt{s} = 10$ GeV
only PS mesons in simulations



The asymmetry reproduces the main features of the data
positive sign
rising trend with fractional energy

more quantitative analysis will be performed

$$h_1 = \pi^\pm, \quad h_2 = \pi^\pm$$
$$z_i = \frac{2E_i}{\sqrt{s}}$$
$$z_1 > 0.2, \quad z_2 > 0.2$$

Conclusions

We generalized the string+ 3P_0 model of hadronization to $e^+e^- \rightarrow q\bar{q} \rightarrow \text{hadrons}$
recursive quantum mechanical recipe

The recipe is general, independent on the production mechanism of the $q\bar{q}$ pair

The implementation in Pythia 8.3 is ongoing
preliminary Collins asymmetry for back-to-back pions is as expected

(More) phenomenological studies ongoing
the goal is to publish the results in few months..

Backup

Relevant free parameters for string fragmentation used in simulations

(see AK, L. Lönnblad, arXiv: 2305.05058)

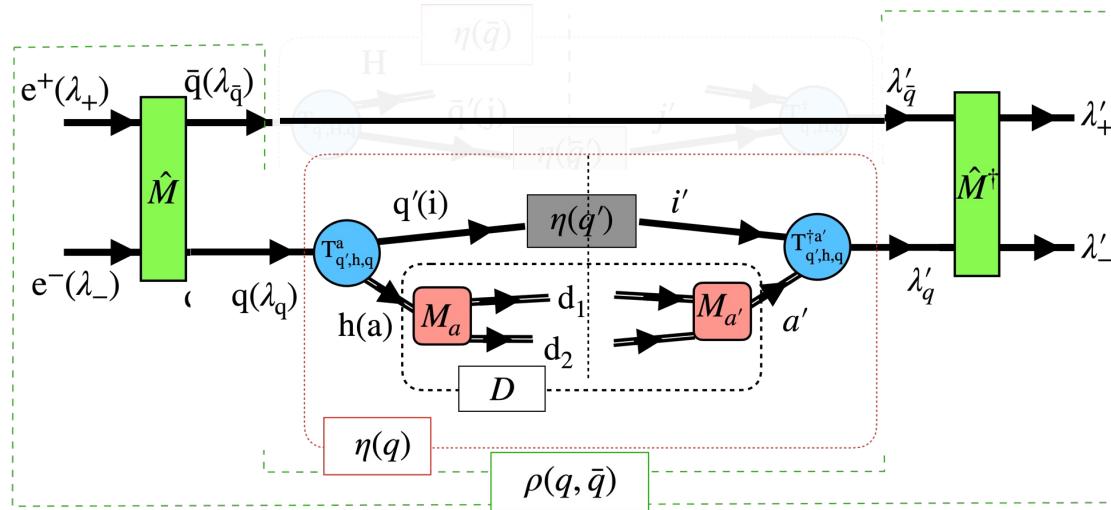
Pythia parameters

StringZ:aLund	default
StringZ:bLund	default
StringPT:sigma	default
StringPT:enhancedFraction	0.0
StringPT:enhancedWidth	0.0 GeV/c

String+ 3P_0 parameters

$\text{Re}(\mu)$	0.42 GeV/ c^2
$\text{Im}(\mu)$	0.76 GeV/ c^2
f_L	0.93
θ_{LT}	0

The recursive recipe for simulating e^+e^- annihilation: VM emission



For a vector meson $h=VM$

$$\rightarrow \eta(q) = T_{q',h=VM,q}^{a'\dagger} \eta(q') T_{q',h=VM,q}^a D_{a'a}, \quad \eta(q') = 1_{q'}, \text{ and } \eta(\bar{q}) = 1_{\bar{q}}$$

Steps:

i) Emission probability density (summing over decay information, i.e. $D_{a'a} = \delta_{a'a}$)

$$\frac{dP(q \rightarrow h = VM + q'; q\bar{q})}{dM^2 dZ_+ Z_+^{-1} d^2 p_T} = \text{Tr}_{q'\bar{q}} T_{q',h,q}^a \rho(q, \bar{q}) T_{q',h,q}^{a\dagger} = F_{q',h,q}(M^2, Z_+, p_T; k_T, C^{q\bar{q}})$$

ii) Calculate the spin density matrix of $h=VM$, and decay the meson

$$\rho_{aa'}(h) = \text{Tr}_{q'\bar{q}} T_{q',h,q}^a \rho(q, \bar{q}) T_{q',h,q}^{a\dagger}$$

iii) Decay the meson $p \rightarrow p_1 p_2 ..$

$$dN(p_1, p_2, \dots) / d\Omega \propto M_{\text{dec.}}^a(p \rightarrow p_1 p_2, \dots) \rho_{aa'}(h) M_{\text{dec.}}^{a\dagger a'}(p \rightarrow p_1 p_2, \dots)$$

iv) Build the decay matrix $D_{a'a}(p_1, p_2, \dots) = M_{\text{dec.}}^{a\dagger a'}(p \rightarrow p_1 p_2, \dots) M_{\text{dec.}}^a(p \rightarrow p_1 p_2, \dots)$