

ELECTROWEAK STRUCTURE OF THE NUCLEON

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June 7, 2023



20th International Conference on Hadron Spectroscopy and Structure



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ELECTROWEAK STRUCTURE OF THE NUCLEON

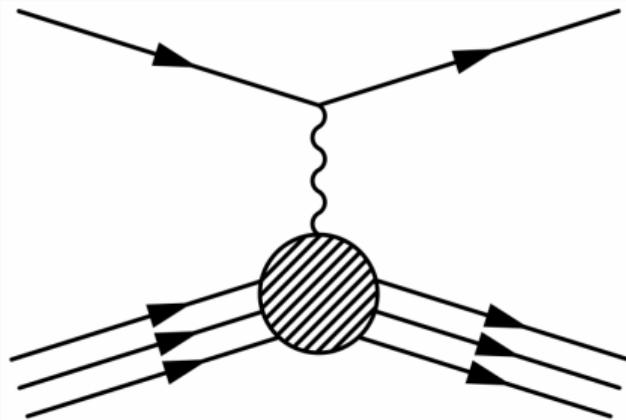
- $F_{V(\text{EM})}$ → charge distribution and magnetism

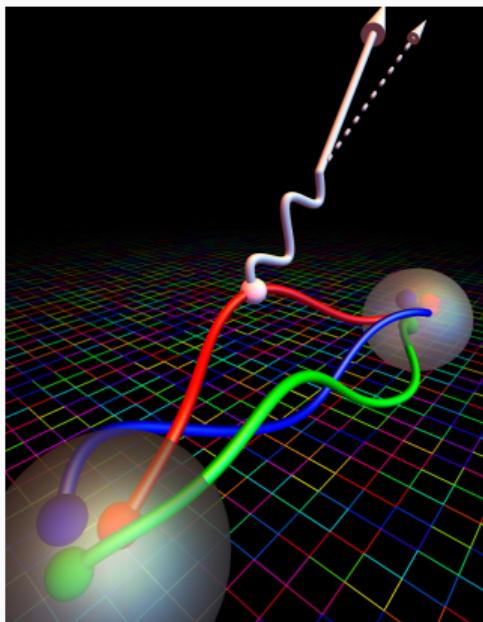
F_A → spin distribution

Weak interaction is $V - A$

- $V^\mu = \bar{q} Q \gamma^\mu q, \langle N(p') | V^\mu(0) | N(p) \rangle = \bar{u}' \left[\gamma^\mu F_1^N(q^2) + \frac{i \sigma^{\mu\nu} q_\nu}{2m_N} F_2^N(q^2) \right] u,$
- $A^\mu = \bar{q} \gamma^\mu \gamma_5 q, \langle N(p') | A^\mu(0) | N(p) \rangle = \bar{u}' \left\{ \gamma_\mu F_A(q^2) + \frac{q_\mu}{2m_N} G_P(q^2) \right\} \gamma_5 u$

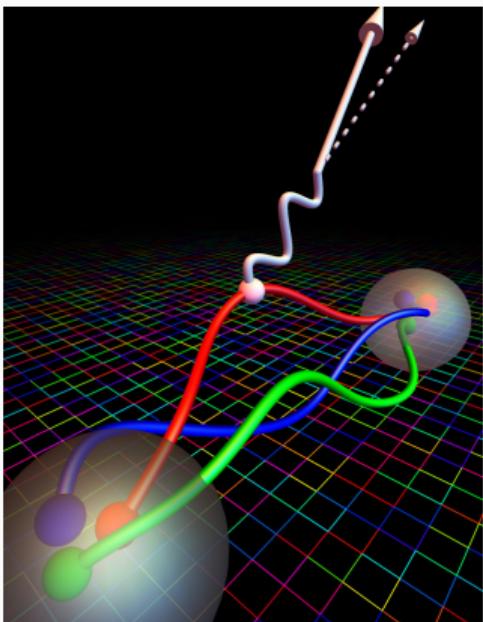
$$q = p' - p \quad N = p, n$$





- Experimental determinations
 - Proton radius puzzle [Pohl et al., Nature 466, 213 (2010)]
- Lattice QCD
 - Nucleon f.f. is a benchmark for LQCD
 - Uncertainties reduced for unphysical large M_π
 - Technical difficulties → recent progress (cite)
 - Experimental and lattice q^2 parametrisation:
 - dipole ansatz
 - z-expansion
 - ...

} \Rightarrow different $\langle r_1^2 \rangle$, and F_{EM} in general



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} ⇒ different $\langle r_i^2 \rangle$, and F_{EM} in general

- Chiral Perturbation Theory (χ PT)

- EFT for QCD at low energy
- QCD based parametrization of q^2 and M_π dependencies ⇒ extrapolate lattice results to the phys. point and extract $\langle r_i^2 \rangle$ and κ from the lattice simulations
- Account for finite volume, lattice spacing and excited states
- Determining χ PT LECs from the lattice ⇒ predicting other observables

- Dispersion theory

- Enlarge the q^2 range of plain χ PT

- Goal: Disp+ χ PT = good q^2 and M_π description

DISPERSION THEORY

- Enlarge the q^2 range of χ PT (ρ dynamics)

1. Disp. rel. (Cauchy)

$$F(q^2) = \int_{4M_\pi^2}^{\infty} \frac{ds}{\pi} \frac{\text{Im } F(s)}{s - q^2 - i\varepsilon}$$

2. Unitarity $\Rightarrow \text{Im } F = \frac{1}{2} \sum_n T_{\gamma^* n} T_{n \bar{N} N}^\dagger, n = \pi^+ \pi^-, \dots$

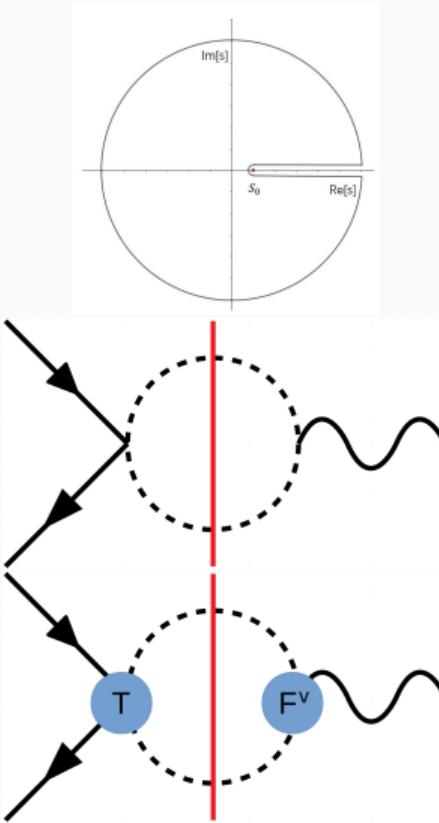
- $\ell = 1, \pi\pi$ must be iso-vector

$$\Rightarrow F_i = \frac{1}{2} F_i^{(s)} + F_i^{(v)} \frac{\sigma^3}{2}$$

3. Using full $N\bar{N}\pi\pi$ and $\gamma^*\pi\pi$ vertices with M_π dep.

$$F(q^2) = \frac{1}{12\pi} \int_{4M_\pi^2}^{\infty} \frac{ds}{\pi} \frac{T p_{\text{cm}}^3 F_\pi^{V*}}{s^{1/2}(s - q^2 - i\varepsilon)},$$

[Granados et al, EPJ A 53 (2017)]



DISPERSION THEORY

- Our two vertices, T and F_π^V , include nonperturbatively the $\pi\pi$ scattering amplitude, t
- $\mathcal{A}(s) = \frac{8}{\pi} \sum_{\ell} (2\ell+1) P_{\ell}(\cos \theta) t_{\ell}(s)$
- $\ell = 1, t = \frac{\sqrt{s}}{2p_{\text{cm}}} \sin \delta e^{i\delta}$
- $\Omega(s) = \exp \left\{ s \int_{4M_\pi^2}^{\infty} \frac{ds'}{\pi} \frac{\delta(s')}{s'(s'-s-i\epsilon)} \right\}$
- We fit t_{IAM} to δ from [Garcia-Martin PRD 83(2011)]
- We check that the M_π dependence is realistic

$$F_\pi^V(s) = [1 + \alpha_v s] \Omega(s)$$

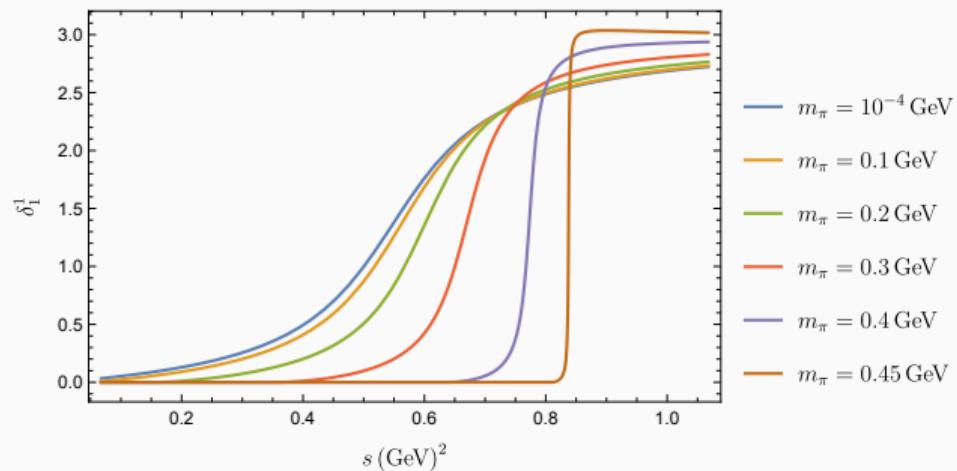
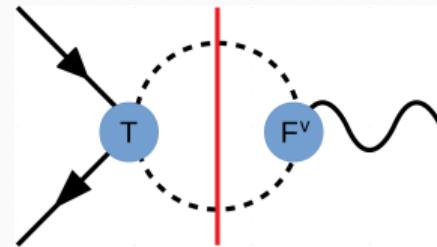
$$T_i(s) = K_i(s) + \Omega(s) P_i + I_i(s)$$

$$I_i(s) = \Omega(s) s \int_{4M_\pi^2}^{\infty} \frac{ds'}{\pi} \frac{K_i(s') \sin \delta(s')}{|\Omega(s')|(s'-s-i\epsilon)}$$

K and P from $NN \rightarrow \pi\pi$ in χPT

$$\mathcal{M} = A \bar{v} u - \frac{1}{2} B \bar{v} \not{u} \longleftrightarrow K, P$$

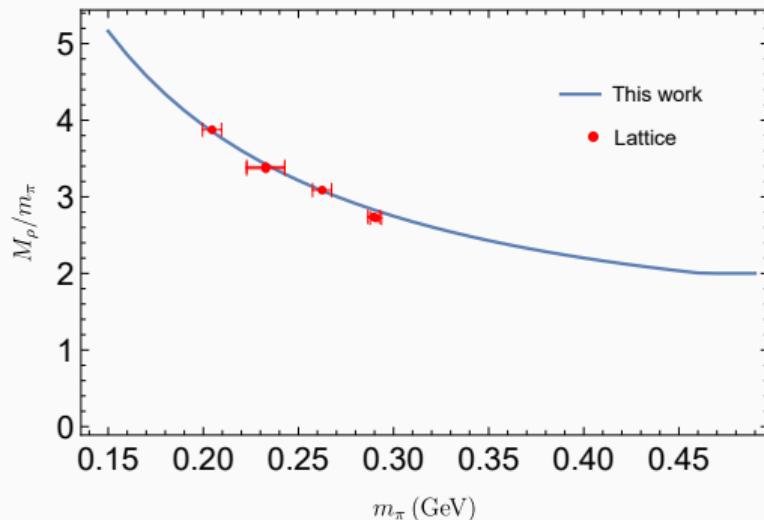
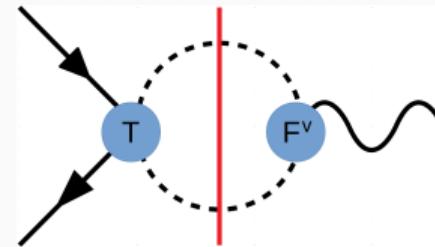
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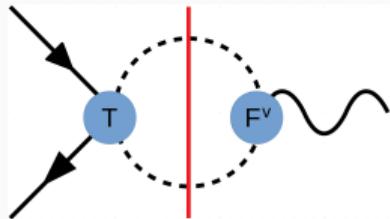
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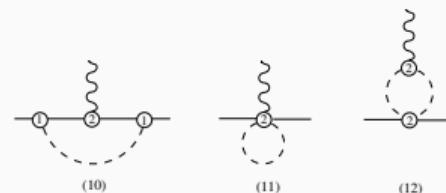
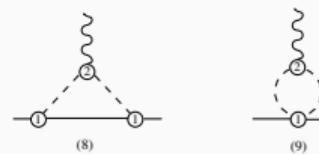
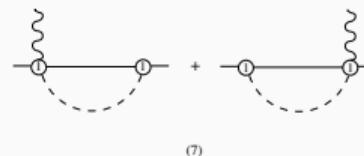
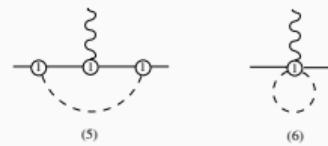
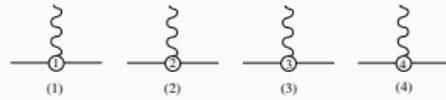


- $\text{Disp} + \chi\text{PT} = F_{\text{EM}}^{\text{disp}} + \text{diagrams without } 2\pi \text{ cut from } \chi\text{PT}$

- $F_{\text{EM}}^{\text{disp}}$

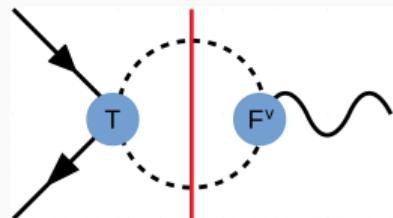


- $+\chi\text{PT}$ diagrams



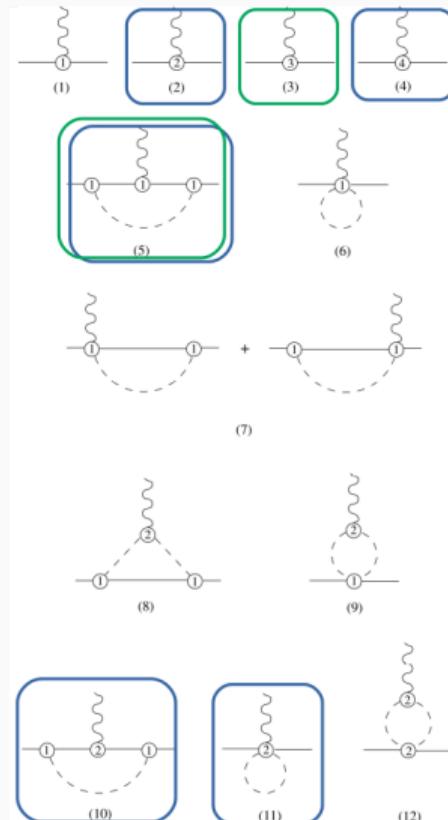
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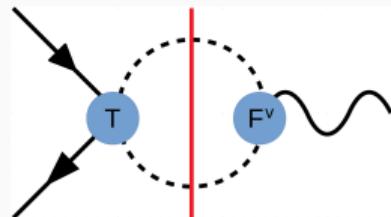
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- Relativistic and with explicit $\Delta(1232)$ [Bauer et al., PRC 86 (2012)]
- green: $F_1 - F_1(0)$ (the charge is trivial)
- blue: $F_2 \rightarrow$ we add the $\mathcal{O}(p^4) \Delta$ terms



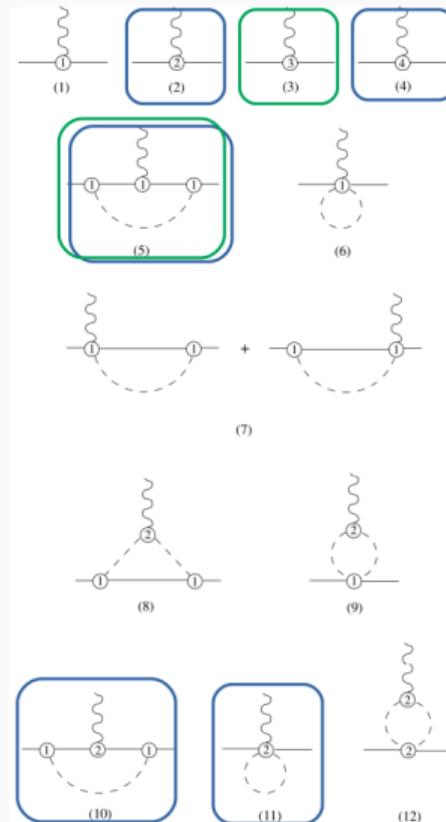
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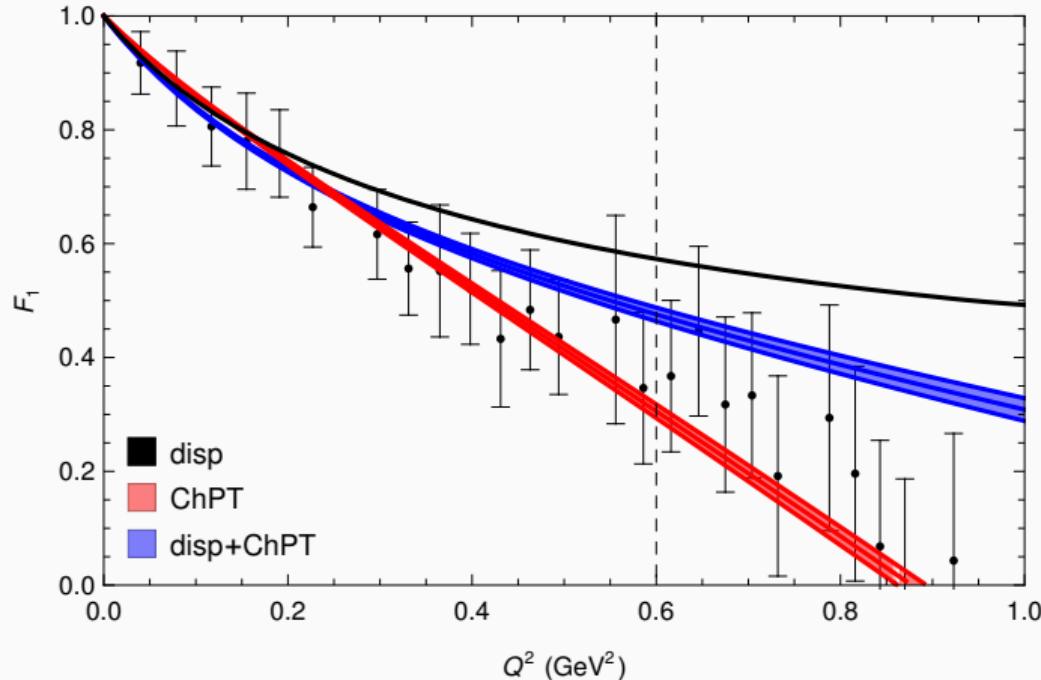


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- green: $F_1 - F_1(0)$ (the charge is trivial)
- blue: $F_2 \rightarrow$ we add the $\mathcal{O}(p^4) \Delta$ terms
- Disp and χPT differ in the renormalization (UV)
 - At $\mathcal{O}(p^3)$ disp and χPT agree on the M_π nonanalyticities
 - Example: the dispersive contribution from $T^{\text{point}} \sim \frac{1}{F^2}$ agrees with χPT
 - $F_1^{\text{point}} \sim F_1^{(9)} \sim q^2 \log M_\pi$
 - differences absorbed in LECs



- Dirac f.f., F_1
- $\langle N | V^\mu | N \rangle = \bar{u}' \left[\gamma^\mu F_1^N(q^2) + \frac{i\sigma^{\mu\nu} q_\nu}{2m_N} F_2^N(q^2) \right] u, F_1 = F_E + \frac{Q^2}{4m^2} \frac{(F_M - F_E)}{(1+Q^2/(4m^2))}, Q^2 = -q^2.$
- $F_1 = 1 + \frac{q^2}{6} \left[-12d_6 + \langle r_1^{2(\log M_\pi)} \rangle \log M_\pi \right] + \mathcal{O}(p^4)$
- Comparison with LQCD data [Djukanovic PRD 103(2021)] ← controlled FV and discret. effects
- In the χ PT and disp+ χ PT F_1 , d_6 is fitted to LQCD
- Disp → q^2 curvature



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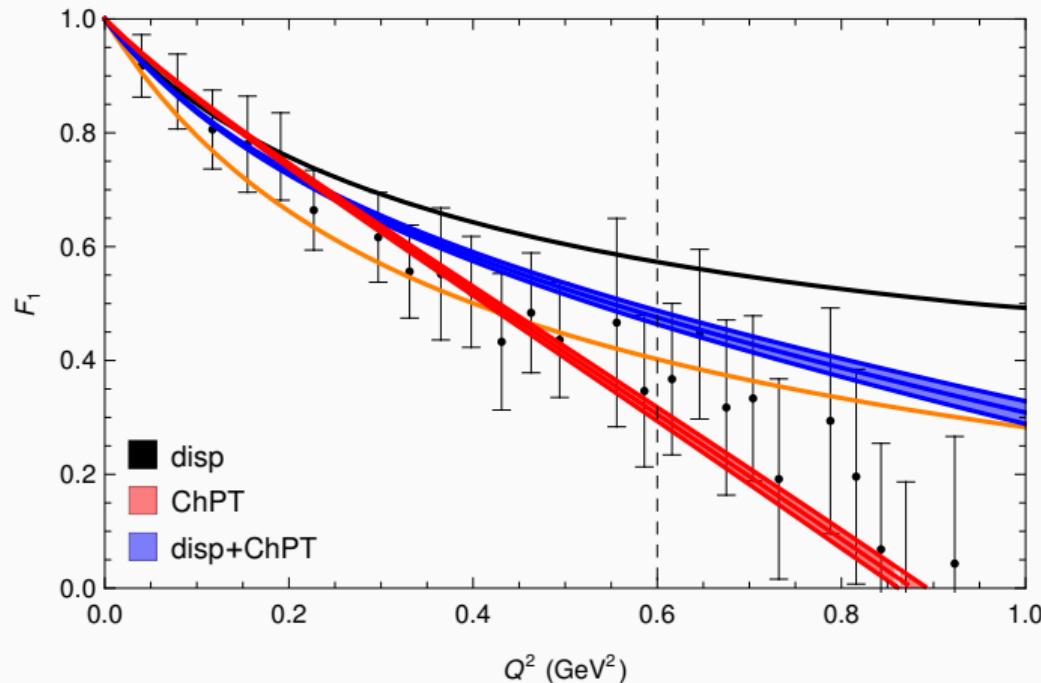
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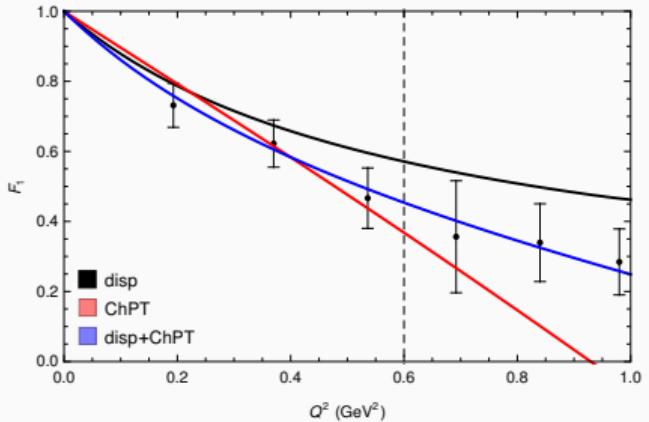
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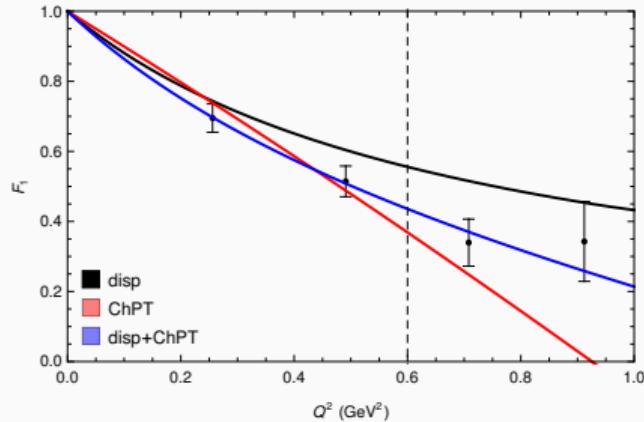
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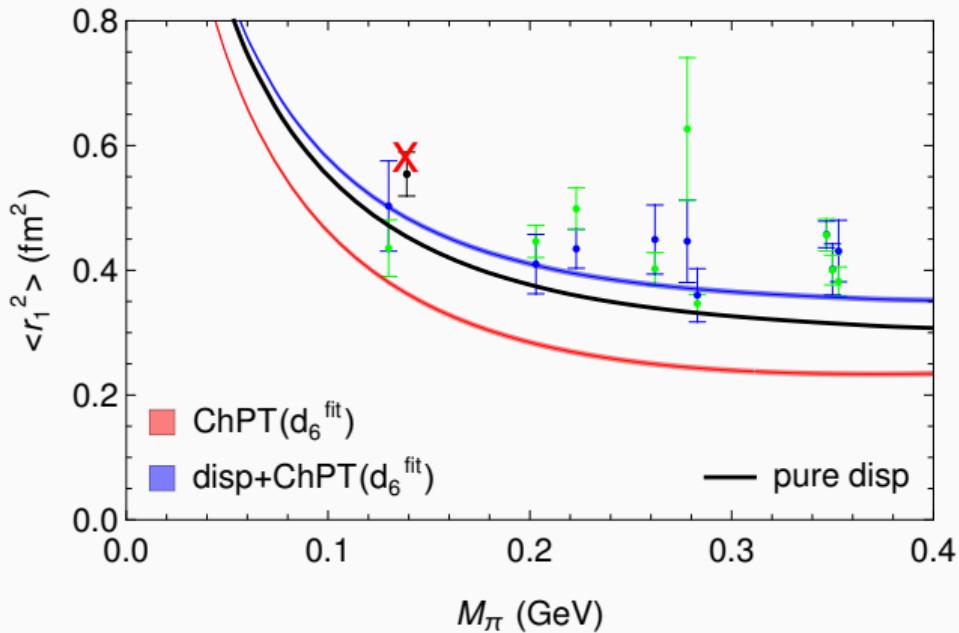
(h) H105 $M_\pi = 0.278$ GeV



(i) N302 $M_\pi = 0.353$ GeV

- Disp+ χ PT describes well the M_π dep.
 - good d_6 fit for $Q^2 < 0.6$ GeV 2 and $M_\pi \lesssim 350$ MeV
 - outperforms the pure dispersive and plain χ PT descriptions

	Disp (prediction)	χ PT	Disp+ χ PT
$d_6(\mu = m_\rho)$ (GeV $^{-2}$)	-	0.074 ± 0.010	0.416 ± 0.010
$d_6(\mu = m_N)$ (GeV $^{-2}$)	-	-0.422 ± 0.010	0.155 ± 0.010
χ^2/dof	$108.9/47 = 2.32$	$73.9/(47 - 1) = 1.61$	$24.6/(47 - 1) = 0.53$
$\langle r_1^2 \rangle_{\text{phys}}$ (fm 2)	0.4541	0.3626 ± 0.0047	0.4838 ± 0.0047

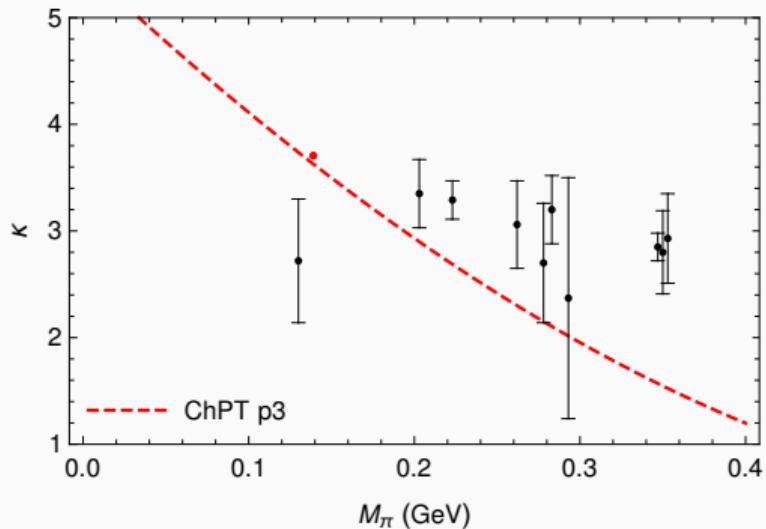


- $F_1^{(v)} = 1 + \frac{1}{6} \langle r_1^2 \rangle^{(v)} q^2 + \mathcal{O}(q^4)$,
 $\langle r_1^2 \rangle^{\text{PDG}} = 0.577 \text{ fm}^2$
- Heavy baryon fit to LQCD from [Djukanovic PRD 103(2021)]: $\langle r_1^2 \rangle^{\text{HB}} = 0.554 \pm 0.035 \text{ fm}^2$

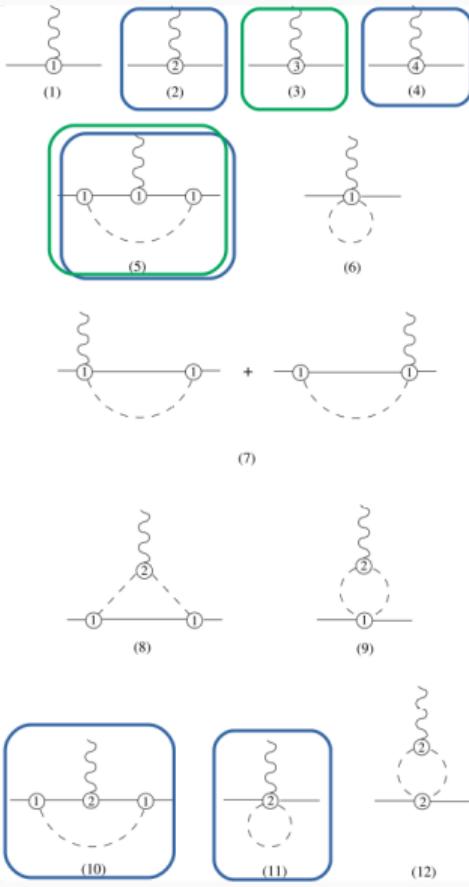
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• Pauli f.f., F_2

- $\langle N | V^\mu | N \rangle = \bar{u}' \left[\gamma^\mu F_1^N(q^2) + \frac{i\sigma^{\mu\nu} q_\nu}{2m_N} F_2^N(q^2) \right] u, F_2 = \frac{F_M - F_E}{(1+Q^2/(4m^2))}$
 - $F_2^{(v)} = \kappa^{(v)} [1 + \frac{1}{6} \langle r_2^2 \rangle^{(v)} q^2 + \mathcal{O}(q^4)]$
 - $\mathcal{O}(p^3) \chi\text{PT}$ is not enough (and Δ subtleties)
- $$F_2(0)^{\text{o}(3)} = c_6 - \left(\frac{\pi m_N g_A^2}{4\pi^2 F^2} \right) M_\pi$$

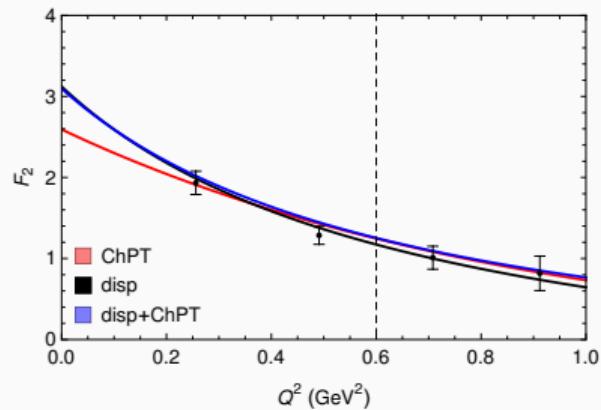
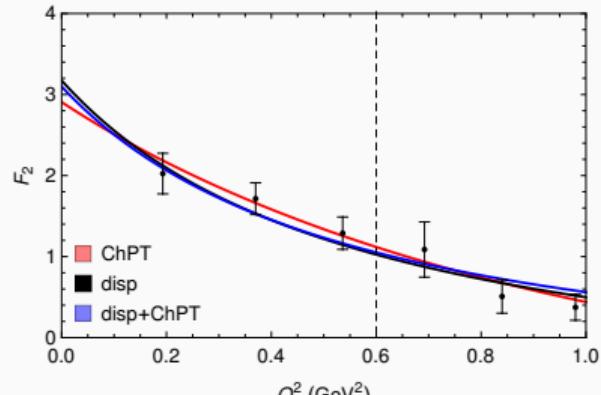
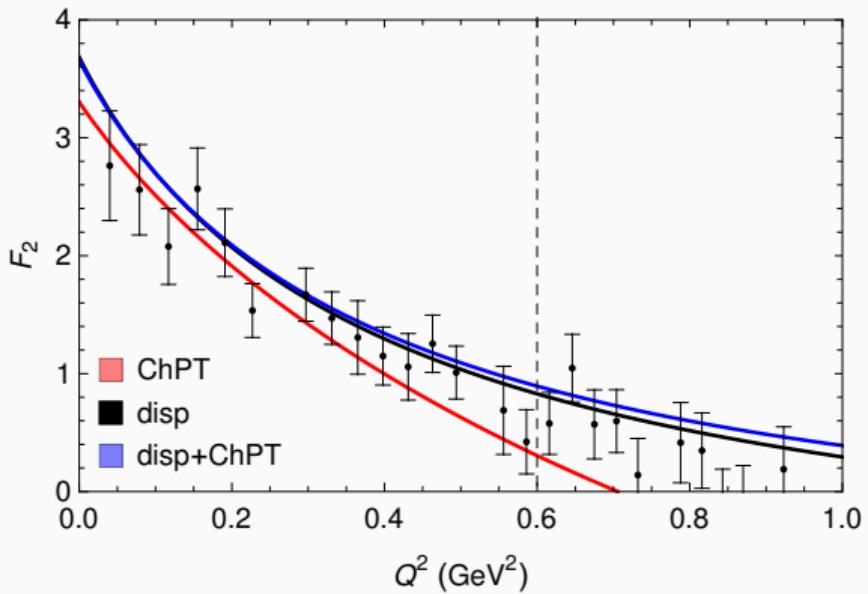


- ⇒ include $\Delta \mathcal{O}(p^4)$
- disp+ χPT $\mathcal{O}(p^4)$: $F_2 = F_2^{\text{disp}} + F_2^{\text{tree}} + F_2^{\chi\text{PTloop}},$
 $F_2^{\text{tree}} = c_6 - 16e_{106}m_NM_\pi^2 + 2q^2(d_6 + 2e_{74}m_N)$



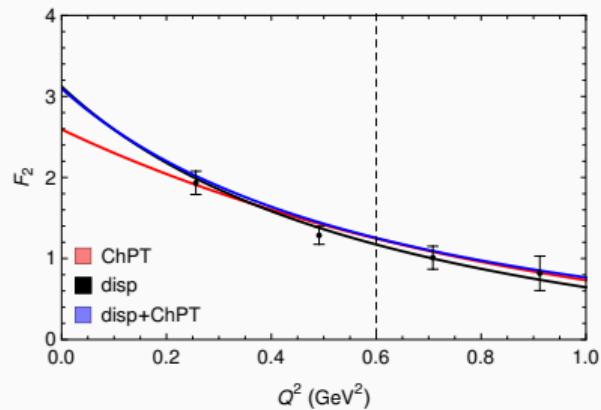
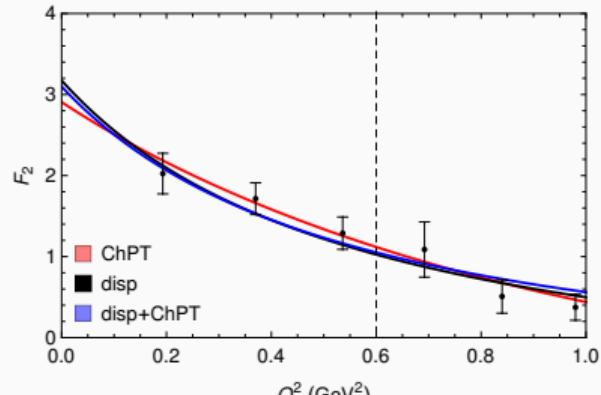
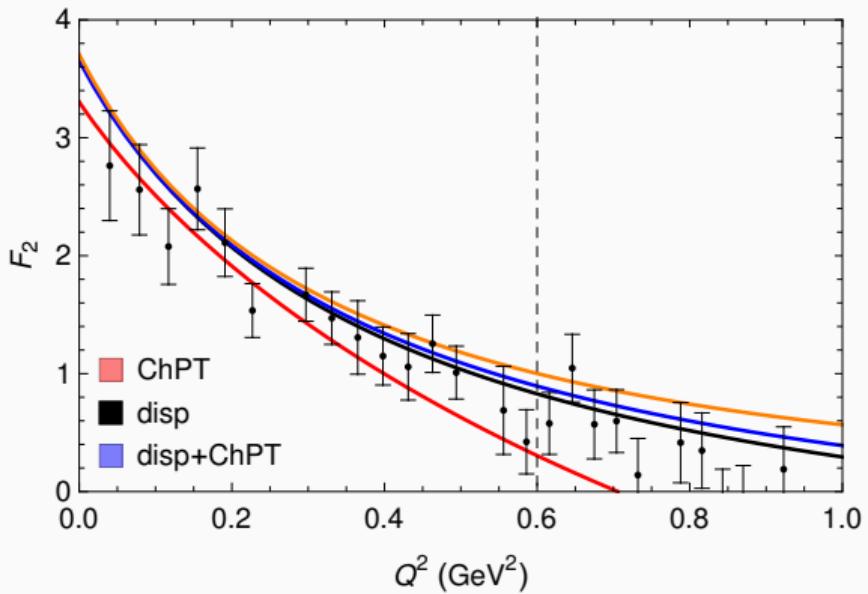
- F_2 fit to LQCD

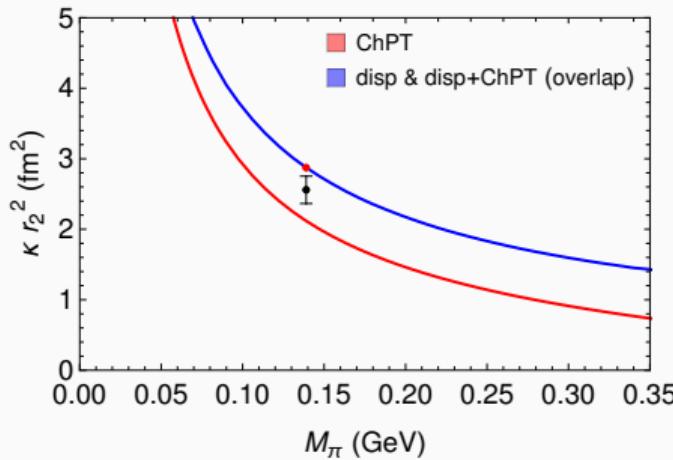
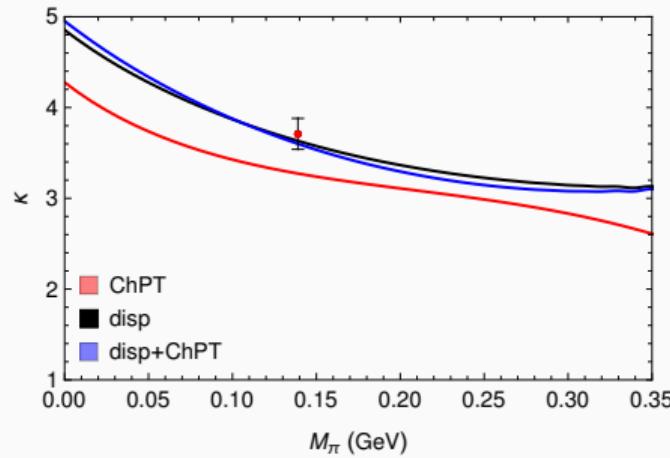
- In F_2^{disp} free c_6
- In $F_2^{\chi\text{PT}}$ and $F_2^{\text{disp}+\chi\text{PT}}$, free c_6, e_{106}, e_{74}
- $\chi\text{PT } \mathcal{O}(p^4)$ and disp separately are good enough to describe the data



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- $F_2^{(v)} = \kappa^{(v)} \left[1 + \frac{1}{6} \langle r_2^2 \rangle^{(v)} q^2 + \mathcal{O}(q^4) \right]$
- $\kappa_{\text{PDG}} = 3.706,$
- $\kappa_{\text{HB}} = 3.71 \pm 0.17,$
- $\langle r_2^2 \rangle_{\text{PDG}} = 0.7754 \text{ fm}^2,$
- $\langle r_2^2 \rangle_{\text{HB}} = 0.690 \pm 0.042 \text{ fm}^2.$

	Disp+c ₆	χ^{PT}	Disp+ χ^{PT}
χ^2/dof	$\frac{49.95}{47-2} = 1.11$	$\frac{44.18}{47-4} = 1.027$	$\frac{56.08}{47-4} = 1.304$
χ_0^2/dof	1.09	1.027	1.283
κ_{phys}	3.64	3.42	3.61
$\langle r_2^2 \rangle_{\text{phys}}$ (fm ²)	0.673	0.619	0.668

AXIAL FORM FACTOR

- $F_A \rightarrow$ spin distribution

Weak interaction is $V-A$

$$A^{i\mu} = \bar{q} \gamma^\mu \gamma_5 \frac{\tau^i}{2} q, \langle N | A^\mu | N \rangle = \bar{u}' \left[\gamma_\mu F_A(q^2) + \frac{q_\mu}{2m_N} G_P(q^2) \right] \gamma_5 \frac{\tau^i}{2} u$$

- Nucleon axial isovector form factor

- $F_A(q^2) = g_A [1 + \frac{1}{6} \langle r_A^2 \rangle q^2 + \mathcal{O}(q^4)]$
- g_A and F_A dependence in q^2 are necessary in ν oscillations experiments
- μ capture, β -decay
- χ PT calculation of F_A
 \Rightarrow extract $\langle r_A^2 \rangle$ from lattice QCD without ad-hoc parametrization
- $g_A(M_\pi)$: test of πN scatt. in χ PT [Alvarado & Alvarez-Ruso PRD 105 \(2021\)](#)
- $\mathcal{O}(p^4) F_A$ in relativistic χ PT
- $F_A = \hat{g}_A + 4d_{16}M_\pi^2 + d_{22}q^2 + \text{loops}(M_\pi, q^2)$
- Difference between orders \simeq theoretical uncertainty

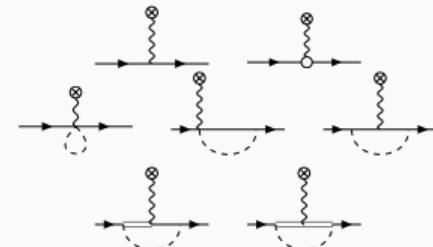


Figure: $\mathcal{O}(p)$ and $\mathcal{O}(p^3)$ (w. f. renormalisation not shown)

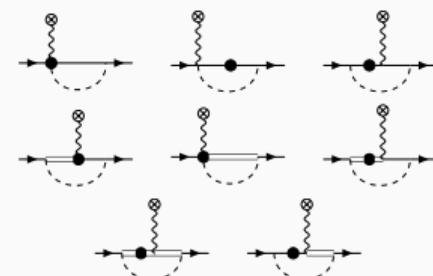
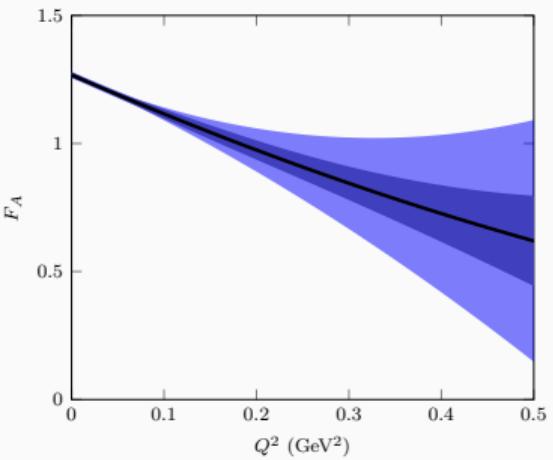
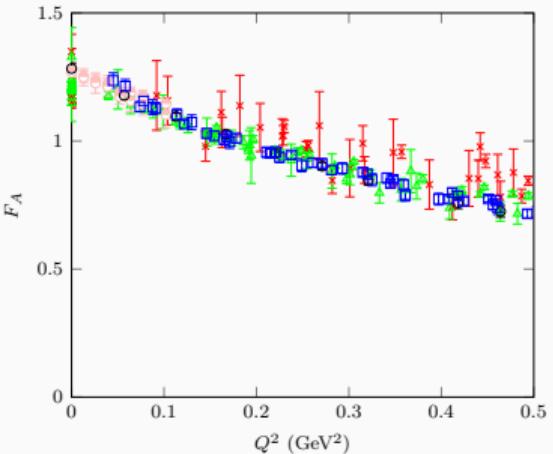


Figure: $\mathcal{O}(p^4)$

- F_A : Meta-analysis of large set of recent LQCD results

- Many recent works \Rightarrow substantial improvements
- RQCD^[1] + PNDME^[2] + "Mainz"^[3] + PACS^[4] + ETMC^[5]
- Reasonably good fit: χ^2 plateau
 $\implies M_\pi^{\text{cut}} \simeq 400 \text{ MeV}, Q_{\text{cut}}^2 = 0.36 \text{ GeV}^2$



[1] Bali et al. JHEP 05 (2020)

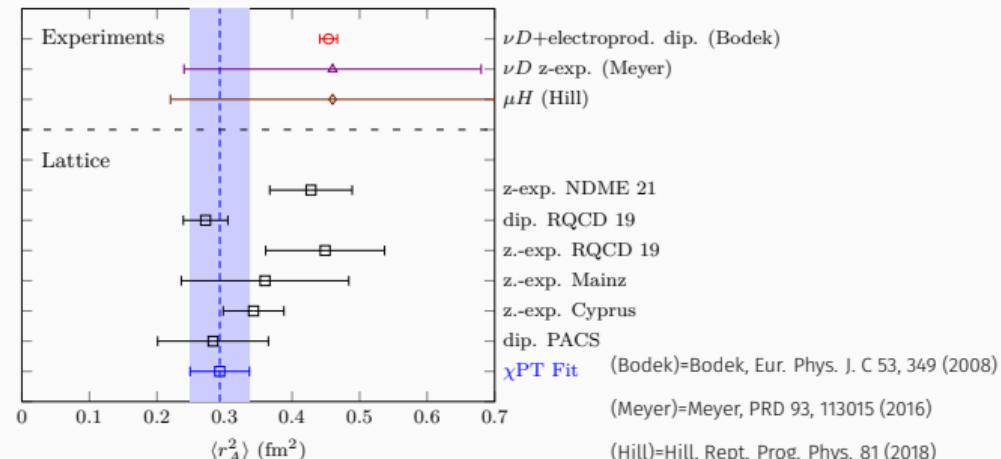
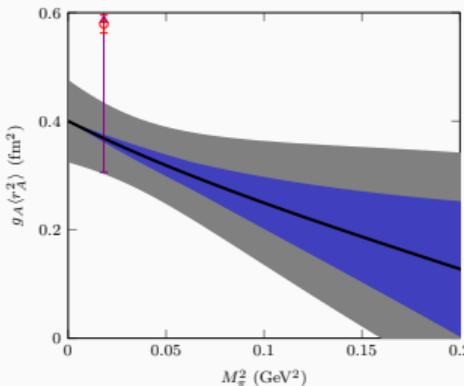
[2] Park et al. 2103.05599

[3] Meyer et al. Modern Phys. A 34 (2019)

[4] Shintani et al. PRD 102 (2020)

[5] Alexandrou et al. PRD 103 (2021)

$$F_A = g_A \left(1 + \frac{1}{6} \left[\langle r_A^2 \rangle \right] q^2 \right) \text{ AXIAL RADIUS}$$



- Our $\mathcal{O}(p^4)$ χ PT extraction:
 - M_π slope driven by loops with Δ
 - $d_{22} = 0.29 \pm 1.69 \text{ GeV}^{-2}$ (no assumptions on $\Delta\Delta\pi$ coupling enlarges error)
 - d_{22} compatible with $\mathcal{O}(p^3)$ π electroprod. [Guerrero et al. PRD, 102 \(2020\)](#)
 - $\boxed{\langle r_A^2 \rangle(M_{\text{phys}}) = 0.293 \pm 0.044 \text{ fm}^2}$
- Empirical determinations (model dependent) are in **tension** with ours and with most of LQCD extractions
- Typically the extracted $\langle r_A^2 \rangle^{\text{phys}}$ value varies depending on the parametrisation
- Our QCD based parametrisation leads to a value in line with most of the individual LQCD extractions

CONCLUSIONS

- F_{EM}
- Dirac f.f., F_1
 - The dispersive calculation supplemented with χ PT contributions outperforms the pure dispersive and plain χ PT descriptions
 - it fits well the LQCD F_1 at least for $Q^2 < 0.6 \text{ GeV}^2$ and $M_\pi \lesssim 350 \text{ MeV}$
 - $\langle r_1^2 \rangle_{\text{phys}} = 0.4838 \pm 0.0047 \text{ fm}^2$
 - value close to the LQCD HB [Djukanovic PRD 103(2021)] extraction and to the experimental one
- Pauli f.f., F_2
 - Disp, $\mathcal{O}(p^4)$ χ PT and disp+ χ PT describe the data well
 - $\kappa_{\text{phys}} = 3.61$ and $\langle r_2^2 \rangle_{\text{phys}} = 0.668 \text{ fm}^2$
 - values in line with the HB and the experimental ones
- F_A
 - Successful description of LQCD $F_A(q^2)$ using $\mathcal{O}(p^4)$ relativistic χ PT
 - Fit describes data in $M_\pi^{\text{cut}} \simeq 400 \text{ MeV}$, $Q_{\text{cut}}^2 = 0.36 \text{ GeV}^2$
 - There is tension between the experimental and lattice extraction of $\langle r_A^2 \rangle$
 - We extract $\langle r_A^2 \rangle_{\text{phys}} = 0.291 \pm 0.052 \text{ fm}^2$ without ad hoc parametrisations
- (Useful LEC values extracted from both calculations)

	Disp (prediction)	χ^{PT}	Disp+ χ^{PT}
$d_6(\mu = m_\rho) \text{ (GeV}^{-2})$	-	0.074 ± 0.010	0.416 ± 0.010
$d_6(\mu = m_N) \text{ (GeV}^{-2})$	-	-0.422 ± 0.010	0.155 ± 0.010
χ^2/dof	$108.9/47 = 2.32$	$73.9/(47 - 1) = 1.61$	$24.6/(47 - 1) = 0.53$
$\langle r_1^2 \rangle_{\text{phys}} \text{ (fm}^2)$	0.4541	0.3626 ± 0.0047	0.4838 ± 0.0047

	$\mathcal{O}(p^3) \Delta$	$\mathcal{O}(p^4) \Delta$	$\mathcal{O}(p^3) \Delta$	$\mathcal{O}(p^4) \Delta$
\tilde{g}_A (free)	1.1782 ± 0.0073		1.2041 ± 0.0074	1.274 ± 0.041
d_{16} (GeV^{-2}) (free)	-1.021 ± 0.048		0.983 ± 0.062	-1.46 ± 1.00
d_{22} (GeV^{-2}) (free)	1.275 ± 0.086		3.77 ± 1.96	0.29 ± 1.69 (free g_1)
h_A	-	-	1.35	1.35
g_1 (free)	-	-	-0.69 ± 0.69	0.66 ± 0.56
c_1 (GeV^{-1})	-	-0.89 ± 0.06	-	-1.15 ± 0.05
c_2 (GeV^{-1})	-	3.38 ± 0.15	-	1.57 ± 0.10
c_3 (GeV^{-1})	-	-4.59 ± 0.09	-	-2.54 ± 0.05
c_4 (GeV^{-1})	-	3.31 ± 0.13	-	2.61 ± 0.10
a_1 (GeV^{-1})	-	-	-	0.90
b_1 (GeV^{-2}) (free)	-	-	-	-0.27 ± 4.96
b_2 (GeV^{-2}) (free)	-	-	-	2.27 ± 2.28
\tilde{b}_4 (GeV^{-2}) (free)	-	-	-	-12.48 ± 1.28
x_1 (fm^{-2}) (free)	-8.4 ± 5.8	-	-5.6 ± 5.9	-0.25 ± 16.5 (consistent)
x_2 (fm^{-2}) (free)	-8.6 ± 2.6	-	-7.1 ± 2.6	-6.36 ± 4.20
x_3 (fm^{-1}) (free)	-0.25 ± 0.21	-	-0.08 ± 0.22	0.36 ± 0.47
y_1 ($\text{fm}^{-2} \text{ GeV}^{-2}$) (free)	-100 ± 40	-	-76 ± 44	-64 ± 121
y_2 ($\text{fm}^{-2} \text{ GeV}^{-2}$) (free)	-31 ± 21	-	-21 ± 22	-15 ± 46
y_3 ($\text{fm}^{-1} \text{ GeV}^{-2}$) (free)	-0.63 ± 1.49	-	0.36 ± 1.63	2.54 ± 3.98
\check{m} (GeV)	0.874	0.874	0.855	0.855
\check{m}_Δ (GeV)	-	-	1.166	1.166

EXTRA 2

	$\mathcal{O}(p^3) \Delta$	$\mathcal{O}(p^4) \Delta$	$\mathcal{O}(p^3) \Delta$	$\mathcal{O}(p^4) \Delta$
\dot{g}_A (free)	1.1782 ± 0.0073		1.2041 ± 0.0074	1.274 ± 0.041
d_{16} (GeV^{-2}) (free)	-1.021 ± 0.048		0.983 ± 0.062	-1.46 ± 1.00
d_{22} (GeV^{-2}) (free)	1.275 ± 0.086		3.77 ± 1.96	0.29 ± 1.69 (free g_1)
h_A	-	-	1.35	1.35
g_1 (free)	-	-	-0.69 ± 0.69	0.66 ± 0.56
c_1 (GeV^{-1})	-	-0.89 ± 0.06	-	-1.15 ± 0.05
c_2 (Gev^{-1})	-	3.38 ± 0.15	-	1.57 ± 0.10
c_3 (Gev^{-1})	-	-4.59 ± 0.09	-	-2.54 ± 0.05
c_4 (Gev^{-1})	-	3.31 ± 0.13	-	2.61 ± 0.10
a_1 (Gev^{-1})	-	-	-	0.90
b_1 (GeV^{-2}) (free)	-	-	-	-0.27 ± 4.96
b_2 (GeV^{-2}) (free)	-	-	-	2.27 ± 2.28
\tilde{b}_4 (GeV^{-2}) (free)	-	-	-	-12.48 ± 1.28
x_1 (fm^{-2}) (free)	-8.4 ± 5.8	-	-5.6 ± 5.9	-0.25 ± 16.5 (consistent)
x_2 (fm^{-2}) (free)	-8.6 ± 2.6	-	-7.1 ± 2.6	-6.36 ± 4.20
x_3 (fm^{-1}) (free)	-0.25 ± 0.21	-	-0.08 ± 0.22	0.36 ± 0.47
y_1 ($\text{fm}^{-2} \text{ GeV}^{-2}$) (free)	-100 ± 40	-	-76 ± 44	-64 ± 121
y_2 ($\text{fm}^{-2} \text{ Gev}^{-2}$) (free)	-31 ± 21	-	-21 ± 22	-15 ± 46
y_3 ($\text{fm}^{-1} \text{ Gev}^{-2}$) (free)	-0.63 ± 1.49	-	0.36 ± 1.63	2.54 ± 3.98
\dot{m} (GeV)	0.874	0.874	0.855	0.855
\dot{m}_Δ (GeV)	-	-	1.166	1.166
χ^2_0/dof	$857.31/(127 - 9) = 7.27$		$533.87/(127 - 10) = 4.45$	$196.58/(127 - 13) = 1.724$