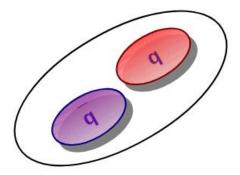




Non-Strange Light-Meson Spectroscopy at COMPASS

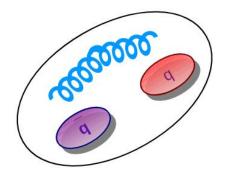
Philipp Haas for the COMPASS Collaboration 06.06.2023 – HADRON 2023

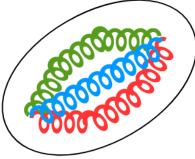
Motivation

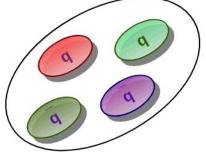


- The Constituent Quark Model predicts mesons as $|q\bar{q}\rangle$ states
- QCD allows meson configurations beyond $|q\bar{q}\rangle$ so-called exotics:
 - Hybrids $|q\overline{q}g\rangle$, Glueballs $|gg\rangle$, Multiquarks $|qq\overline{qq}\rangle$

Exotic mesons at COMPASS (talk by B. Ketzer, Tue. 15:00)

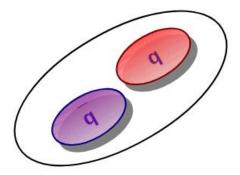






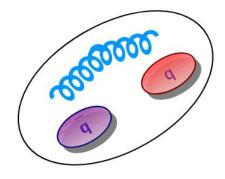
https://arxiv.org/pdf/1405.4195.pdf

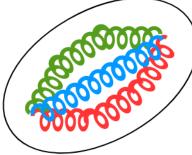
Motivation

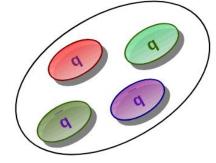


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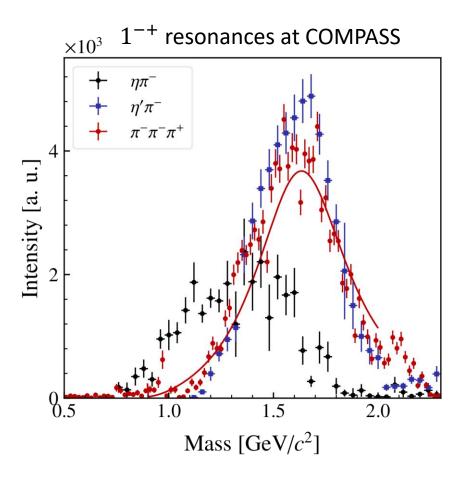


https://arxiv.org/pdf/1405.4195.pdf

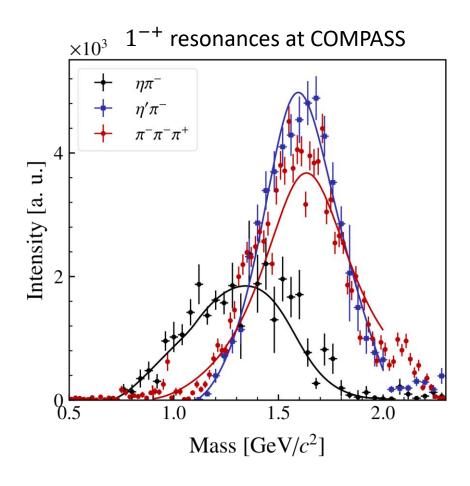
- Light non-strange $|q\bar{q}\rangle$ states cannot make up states with spin quantum numbers $J^{PC}=0^{--}$, even⁺⁻, odd⁻⁺
 - "Spin-exotic" mesons
 - Direct access to find states beyond $|q\bar{q}\rangle$ states

- Lattice QCD predicts the lightest exotic in 1⁻⁺
 - Single pole around 1.6 GeV/ c^2
 - Dominant decay to $b_1\pi$

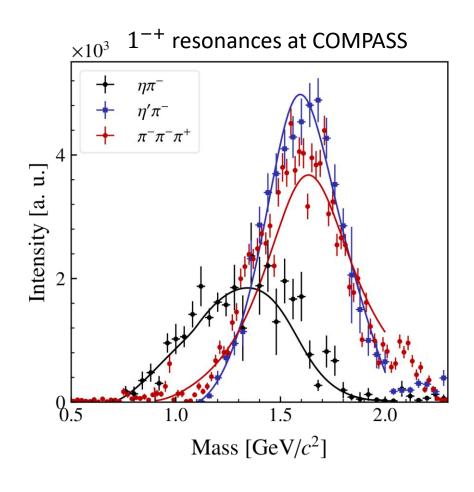
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- 1^{-+} signals at 1.4 GeV/ c^2 and 1.6 GeV/ c^2 seen at COMPASS and other experiments
- JPAC found single pole $\pi_1(1600)$ sufficient for $\eta^{(\prime)}\pi$ COMPASS data
- BNL claimed $\pi_1(2015)$ in $\omega\pi^-\pi^0$ and $f_1\pi$



Experimental Setup

- Located at CERN SPS
- 190 GeV/c negative hadron beam
- Various targets:
 - <u>Liquid-hydrogen</u>
 - Heavy solid-state targets
 Pb, Ni → Primakoff reactions
 - Pb, Ni → Primakoff reactions (talk by D. Ecker, Thu. 17:20)

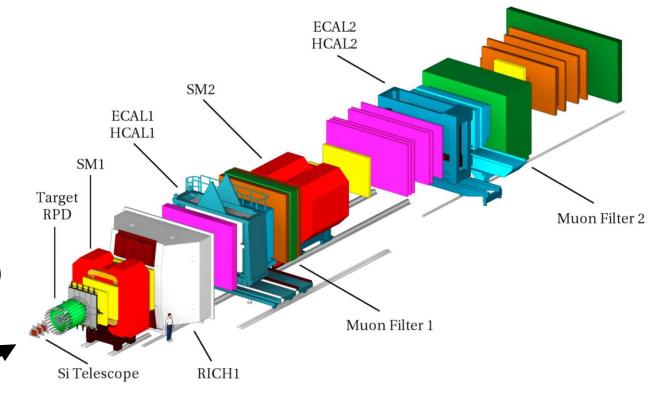
Beam:

 $97\% \pi^{-}$

 $2\% K^{-}$

 $1\% \, \bar{p}$

- Inelastic high-energy $\pi^- p$ scattering
 - Isovector light mesons $X^ (a_I \text{ and } \pi_I)$



Light-Meson Spectroscopy at COMPASS

Analyzed channels:

- $\pi^-\pi^-\pi^+/\pi^-\pi^0\pi^0$
- $\eta \pi^-/\eta' \pi^-$
- $K^-\pi^-\pi^+ \longrightarrow \frac{\text{Strange-meson spectroscopy}}{\text{(talk by S. Wallner, Thu. 14:00)}}$
- $\omega \pi^- \pi^0$

Upcoming channels under study:

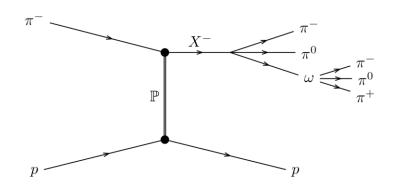
K_SK^-	Search for $a_6(2450)$
$K_S K_S \pi$	Investigate nature of $a_1(1420)$
$f_1\pi^-$	Search for π_1 states
$K_{\scriptscriptstyle S}\pi^-$	Ctranga masans spactrascony
$\Lambda \bar{p}$	Strange mesons spectroscopy

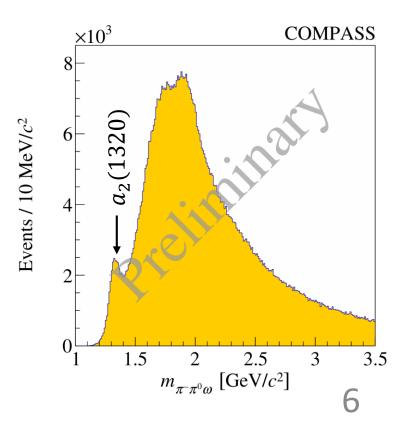
Analysis of $\omega(782)\pi^-\pi^0$

- Overlapping and interfering X^- states
 - m_X spectrum shows no clear peaks above 1.5 ${\rm GeV}/c^2$
- Disentangling the different contributions requires partial-wave analysis

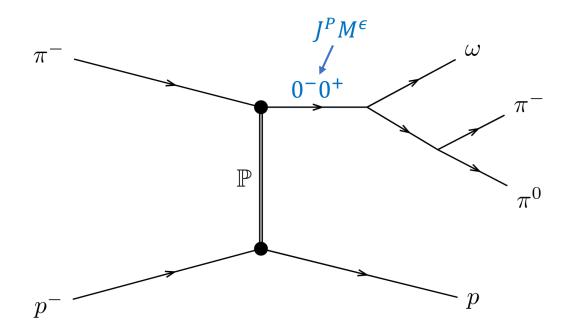
Talk by J. Beckers on Thu. 14:00: Progress in the Partial-Wave Analysis Methods at COMPASS

 Partial-wave decomposition splits the total amplitude in the different contributions

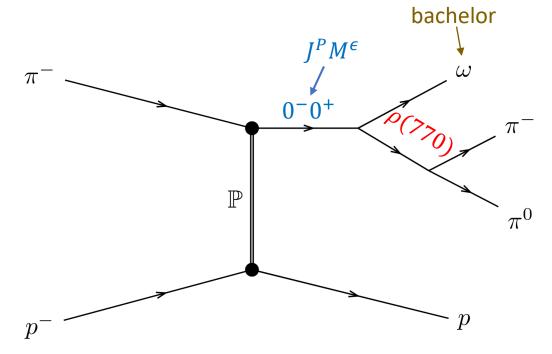




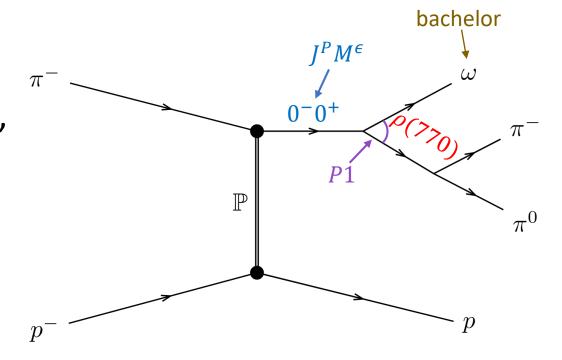
• Exited meson X^- with quantum numbers 0^-0^+ is produced



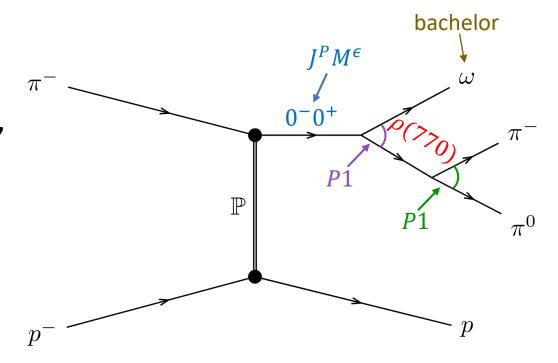
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- Isobar model: X^- decays to $\omega \rho$ (770), where ρ (770) an unstable intermediate state the isobar



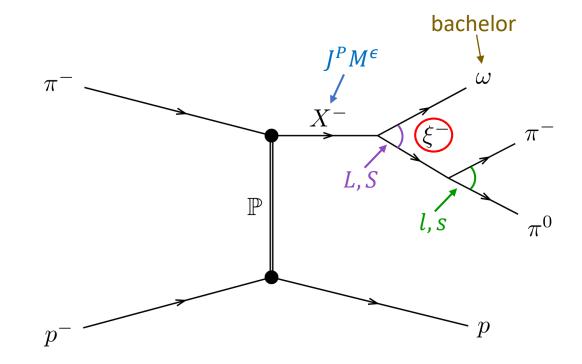
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 - P1 coupling between ω and $\rho(770)$



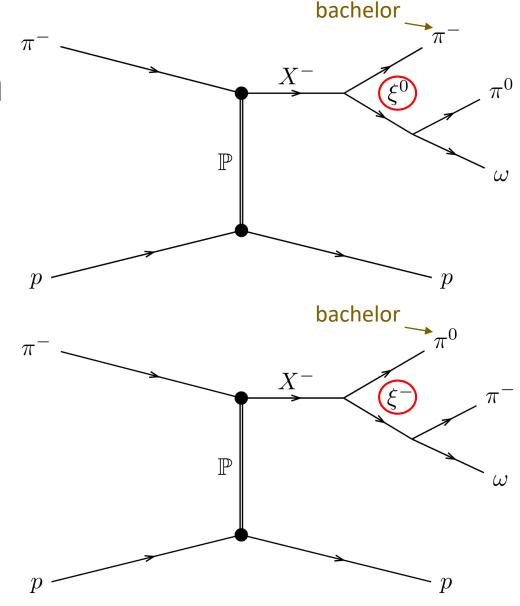
- Exited meson X^- with quantum numbers 0^-0^+ is produced
- Isobar model: X^- decays to $\omega \rho$ (770), where ρ (770) an unstable intermediate state the isobar
 - P1 coupling between ω and $\rho(770)$
- $\rho(770)$ decays to $\pi^-\pi^0$
 - second P1 coupling
- $i = 0^-0^+[\rho(770)P] \omega P1$



- Exited meson X^- with quantum numbers $J^P M^{\epsilon}$ is produced
- Isobar model: X^- decays to $\omega \xi^-$, where ξ^- is an unstable intermediate state the isobar
 - L, S coupling between ω and ξ^-
- ξ^- decays to $\pi^-\pi^0$
 - second *l*, *s* coupling
- $i = J^P M^{\epsilon} [\xi l] \omega LS$



- Further decay channels of X^- :
 - $\pi^0 \xi^-, \pi^- \xi^0$
- Both decays have the same amplitude
 - \Rightarrow Coherently sum over both isospin configurations $\pi^0 \xi^-$, $\pi^- \xi^0$
- $i = J^P M^{\epsilon} [\xi l]$ bachelor LS
 - ξ either decays to $\omega\pi$ or $\pi\pi$



- Coherent superposition of partial-waves:
 - $i = J^P M^{\epsilon} [\xi l]$ bachelor LS

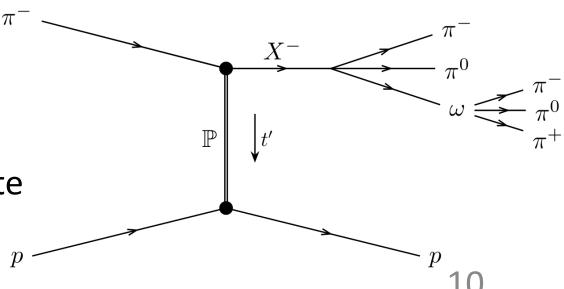
$$I(m_X, t', \tau) = \left| \sum_i \mathcal{T}_i(m_X, t') \psi_i(m_X, \tau) \right|^2$$

with:

 m_X : mass of the $\omega(782)\pi^-\pi^0$ system

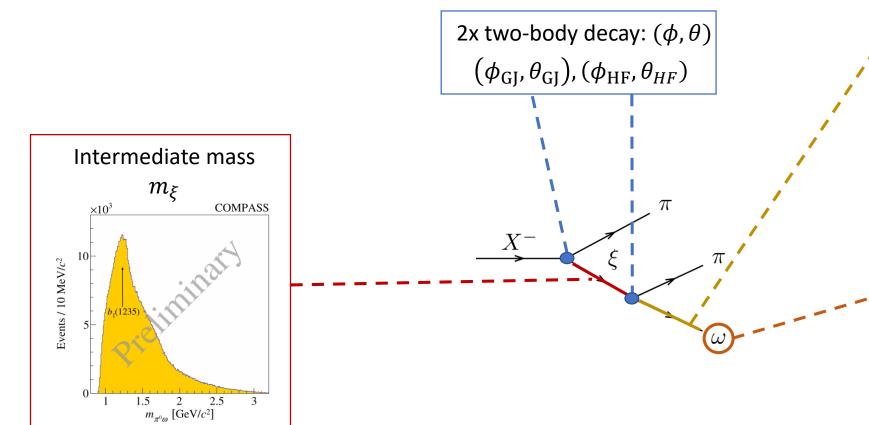
t': squared four-momentum transfer

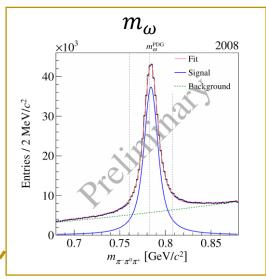
 τ : phase-space variables of the final state

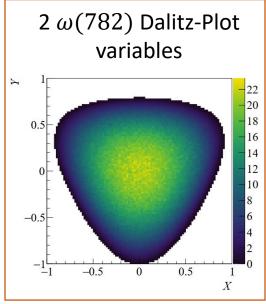


Phase-Space Variables

• τ : Total of 8 phase-space variables







- Coherent superposition of partial-waves:
 - $i = J^P M^{\epsilon} [\xi l]$ bachelor LS

$$I(m_X, t', \tau) = \left| \sum_i \mathcal{T}_i(m_X, t') \psi_i(m_X, \tau) \right|^2$$

• Decay amplitude $\psi_i(m_X, \tau)$: calculated using the isobar model

- Coherent superposition of partial-waves:
 - $i = J^P M^{\epsilon} [\xi l]$ bachelor LS

$$I(m_X, t', \tau) = \left| \sum_i \mathcal{T}_i(m_X, t') \psi_i(m_X, \tau) \right|^2$$

- Decay amplitude $\psi_i(m_X, \tau)$: calculated using the isobar model
- Transition amplitude $\mathcal{T}_i(m_X, t')$:
 - $\Rightarrow \mathcal{T}_i(m_X,t')$ contains production, propagation, and coupling of i
 - No assumptions about the resonant content of X^-
 - \Rightarrow Extract $\mathcal{T}_i(m_X,t')$ by independent maximum-likelihood fits of $I(\tau)$ in bins of (m_X,t')

Partial-Wave Decomposition – Wave Set

ullet In principle: Infinite number of partialwaves i

$$I(m_X, t', \tau) = \left| \sum_i \mathcal{T}_i(m_X, t') \psi_i(m_X, \tau) \right|^2$$

Partial-Wave Decomposition – Wave Set

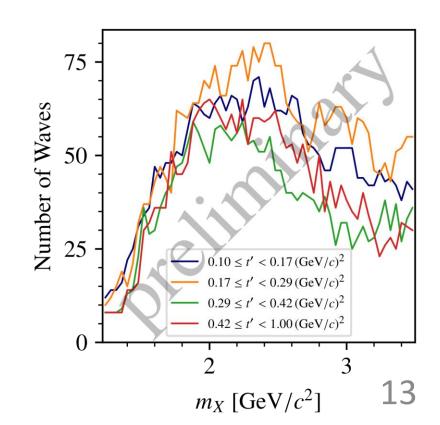
- In principle: Infinite number of partialwaves i
- Construct a wave pool of 893 allowed waves by systematic constraints
 - $\xi \to \pi \pi$: $\rho(770)$, $\rho(1450)$, $\rho_3(1690)$
 - $\xi \to \omega \pi$: $b_1(1235)$, $\rho(1450)$, $\rho_3(1690)$
 - $J \le 8, M \le 2, L \le 8$

$$I(m_X, t', \tau) = \left| \sum_i \mathcal{T}_i(m_X, t') \psi_i(m_X, \tau) \right|^2$$

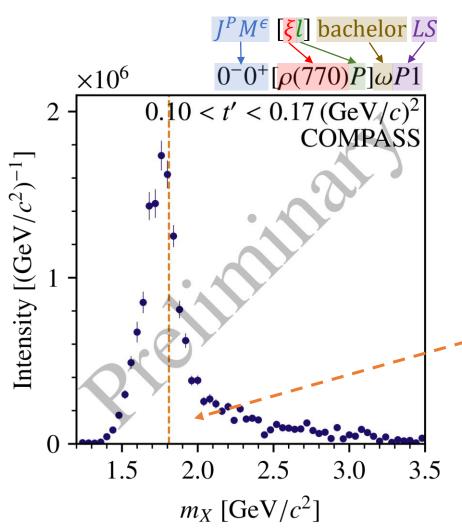
Partial-Wave Decomposition – Wave Set

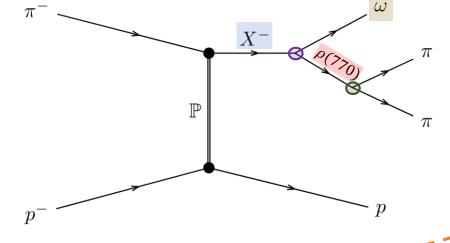
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 - $\xi \to \omega \pi$: $b_1(1235)$, $\rho(1450)$, $\rho_3(1690)$
 - $J \le 8, M \le 2, L \le 8$
- Wave set selected using regularizationbased model-selection
 - Unique wave set for each (m_X, t') cell

$$I(m_X, t', \tau) = \left| \sum_i \mathcal{T}_i(m_X, t') \psi_i(m_X, \tau) \right|^2$$



Results $J^{PC} = 0^{-+}$



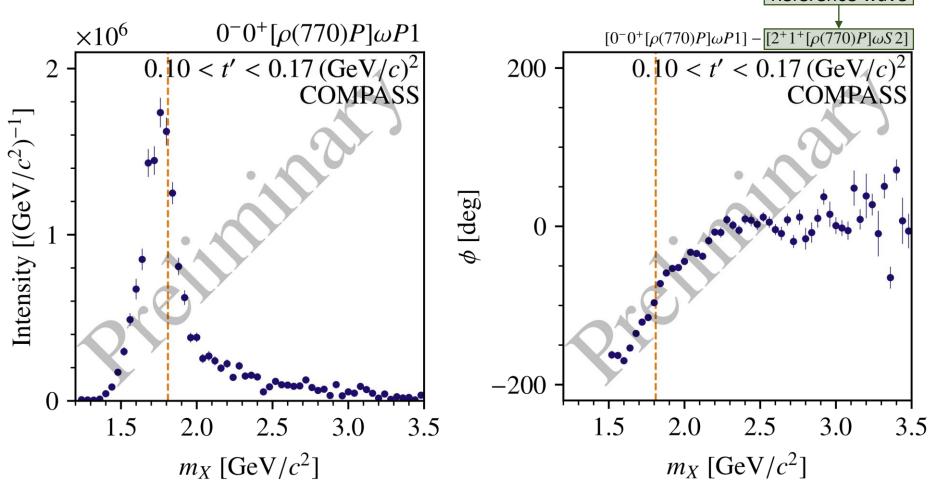


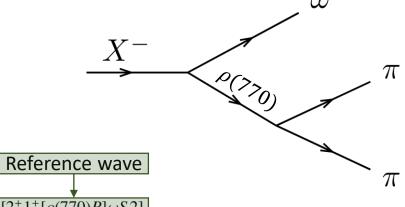
States listed in PDG

$$\pi(1800)$$
 $m = 1810^{+9}_{-11} \text{ MeV}$
 $\Gamma = 215^{+7}_{-8} \text{ MeV}$

 Dashed lines to indicate nominal PDG masses of resonances

Results $J^{PC} = 0^{-+}$

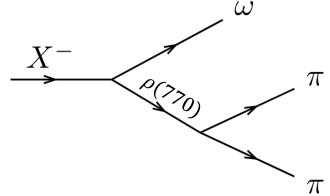


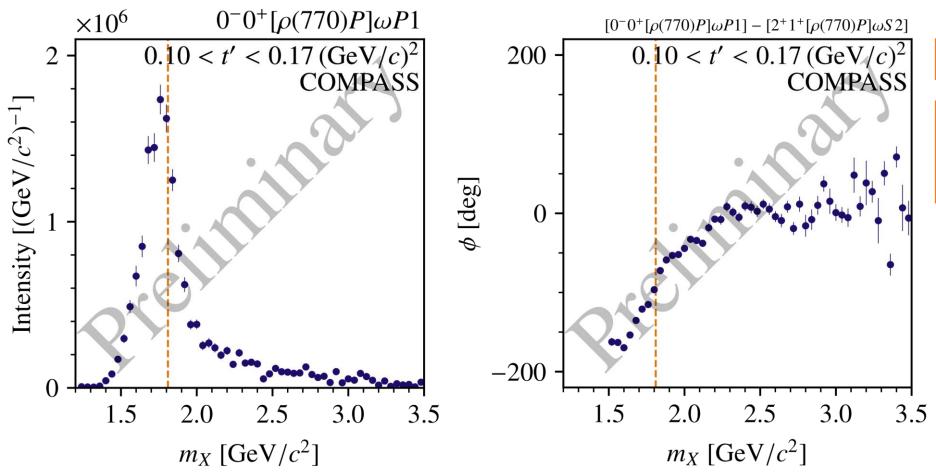


States listed in PDG

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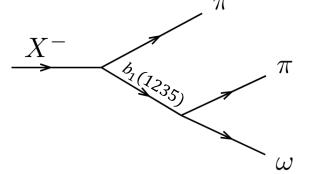


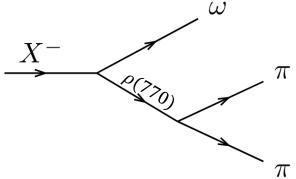


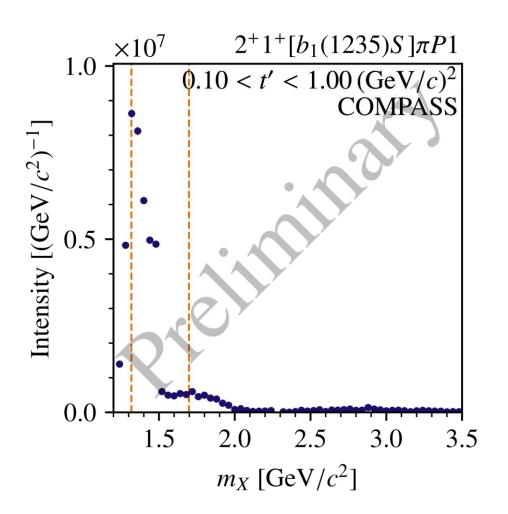
States listed in PDG

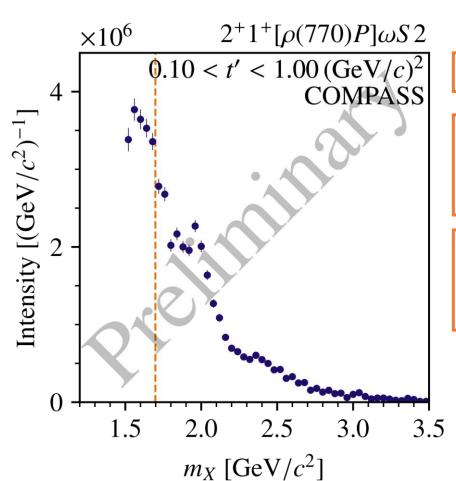
$$\pi(1800)$$
 $m = 1810^{+9}_{-11} \text{ MeV}$
 $\Gamma = 215^{+7}_{-8} \text{ MeV}$

Results $J^{PC} = 2^{++}$







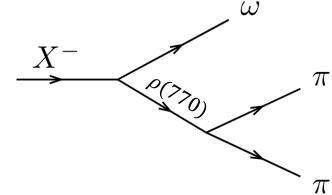


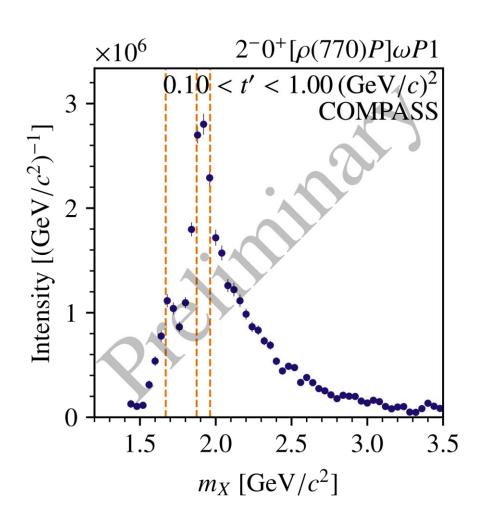
States listed in PDG

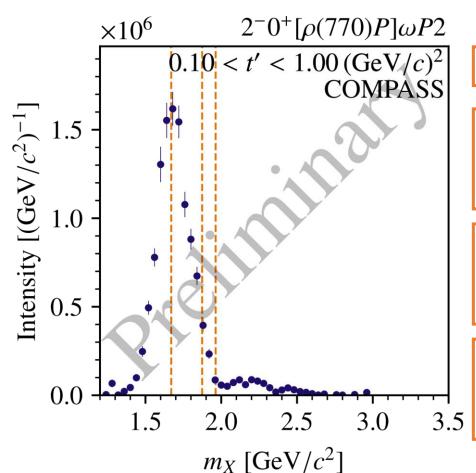
 $a_2(1320)$ $m = 1318.2 \pm 0.6 \text{ MeV}$ $\Gamma = 105^{+1.7}_{-1.9} \text{ MeV}$

 $a_2(1700)$ $m = 1698 \pm 40 \text{ MeV}$ $\Gamma = 265 \pm 60 \text{ MeV}$

Results $J^{PC} = 2^{-+}$







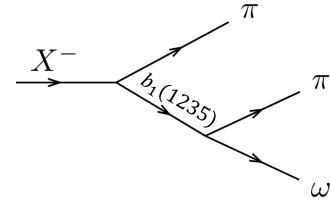
States listed in PDG

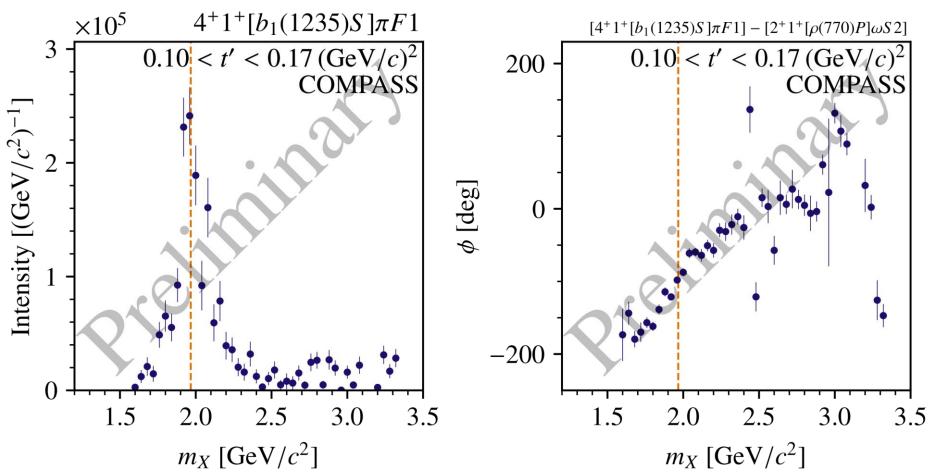
$$\pi_2(1670)$$
 $m = 1670^{+2.9}_{-1.2} \text{ MeV}$
 $\Gamma = 258^{+8}_{-9} \text{ MeV}$

$$\pi_2(1880)$$
 $m = 1874^{+26}_{-5} \text{ MeV}$
 $\Gamma = 237^{+33}_{-30} \text{ MeV}$

$$\pi_2(2005)$$
 $m = 1963^{+17}_{-27} \text{ MeV}$
 $\Gamma = 370^{+16}_{-90} \text{ MeV}$

Results $J^{PC} = 4^{++}$



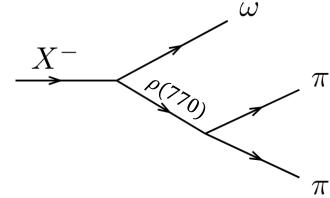


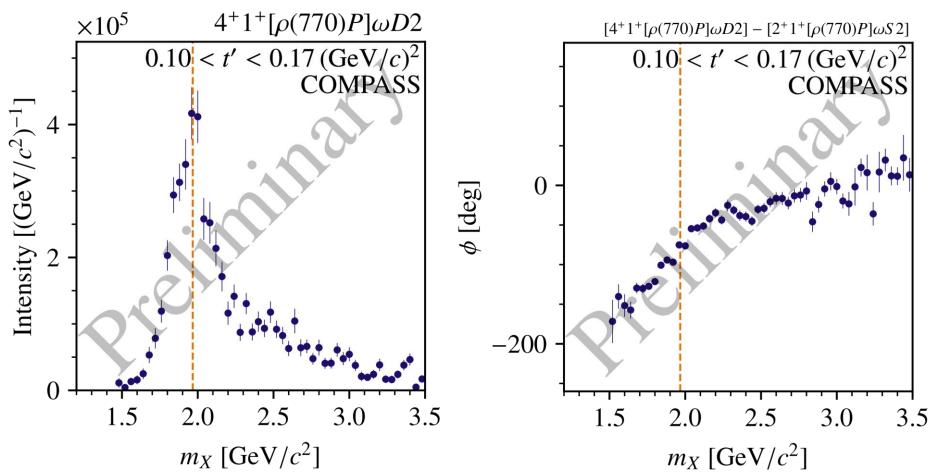
States listed in PDG

$$a_4(1970)$$

 $m = 1967 \pm 16 \text{ MeV}$
 $\Gamma = 324^{+15}_{-18} \text{ MeV}$

Results $J^{PC} = 4^{++}$

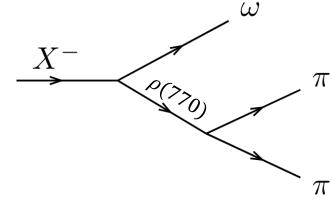


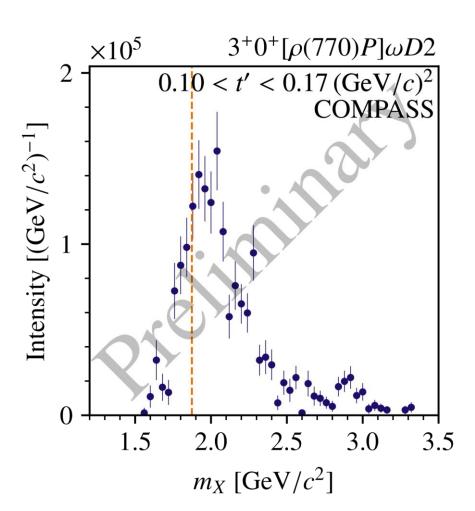


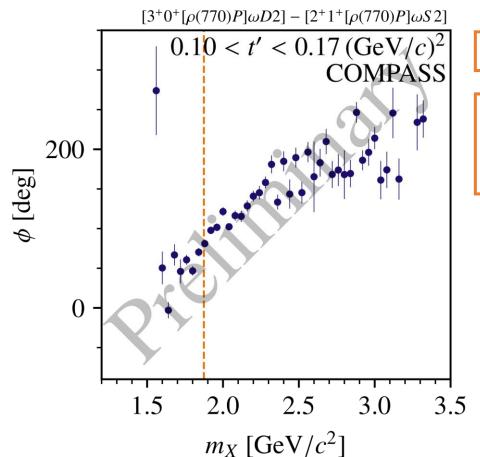
States listed in PDG

 $a_4(1970)$ $m = 1967 \pm 16 \text{ MeV}$ $\Gamma = 324^{+15}_{-18} \text{ MeV}$

Results $J^{PC} = 3^{++}$







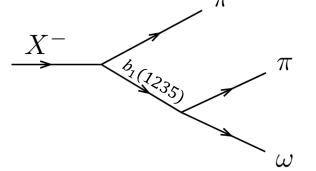
States listed in PDG

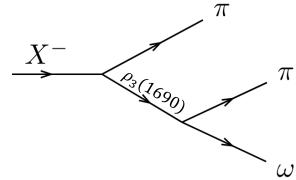
 $a_3(1875)$ $m = 1874 \pm 105 \text{ MeV}$ $\Gamma = 385 \pm 166 \text{ MeV}$

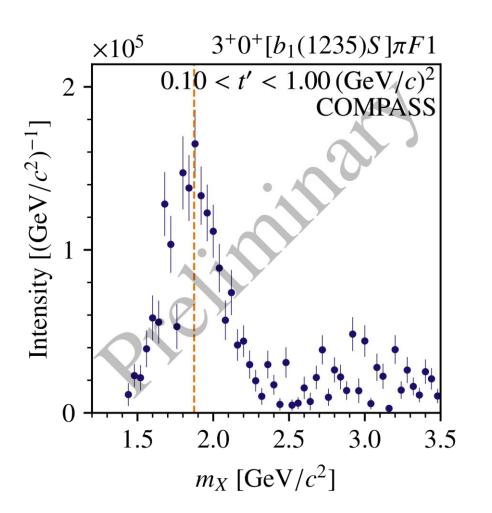
This only has been seen in $\pi^-\pi^-\pi^+$ at BNL E852

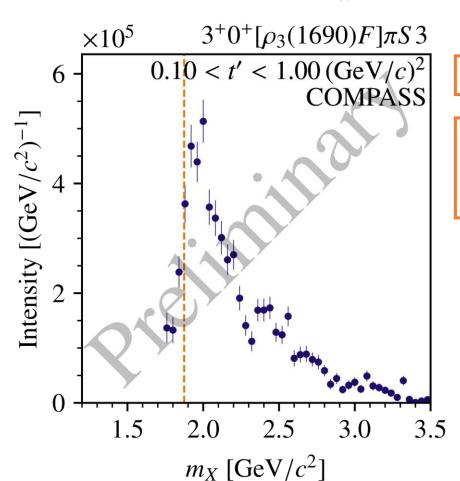
The PDG further lists a $a_3(2030)$

Results $J^{PC} = 3^{++}$









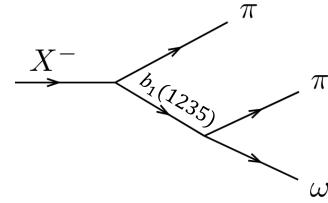
States listed in PDG

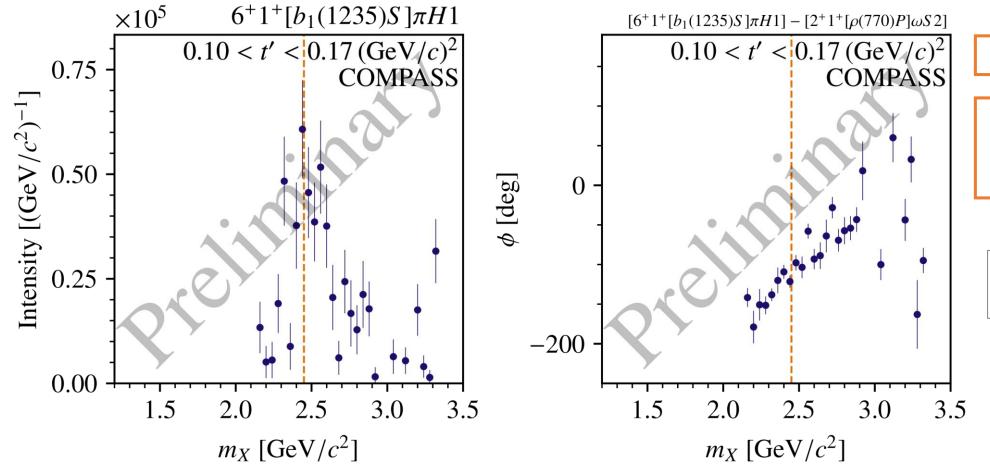
 $a_3(1875)$ $m = 1874 \pm 105 \text{ MeV}$ $\Gamma = 385 \pm 166 \text{ MeV}$

This only has been seen once in $\pi^-\pi^-\pi^+$

The PDG also lists a $a_3(2030)$

Results $J^{PC} = 6^{++}$





States listed in PDG

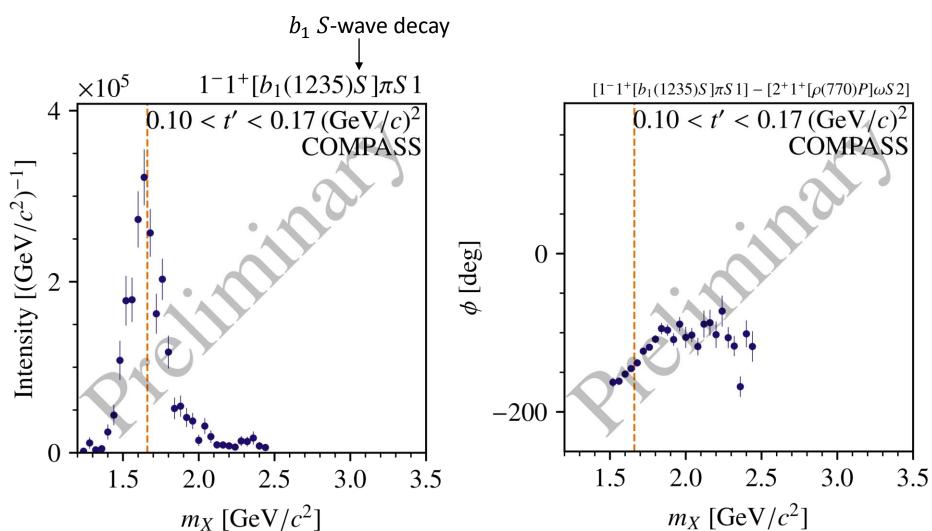
 $a_6(2450)$

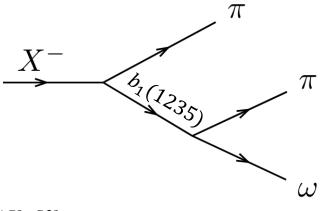
 $m = 2450 \pm 130 \text{ MeV}$

 $\Gamma = 400 \pm 250 \text{ MeV}$

This only has been seen once in K_SK

Results $J^{PC} = 1^{-+}$

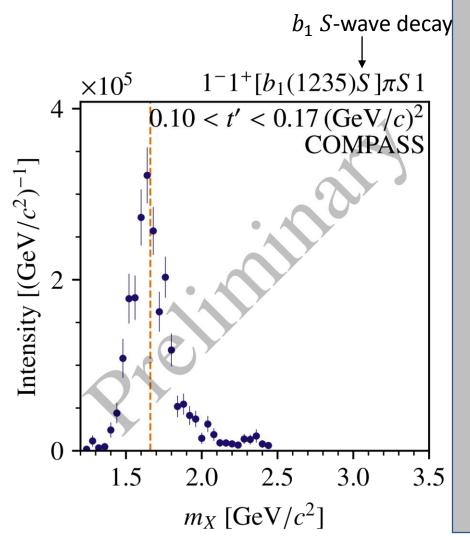




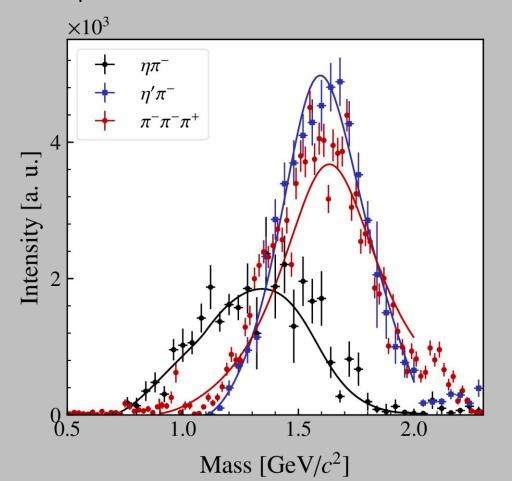
States listed in PDG

 $\pi_1(1600)$ $m = 1661^{+15}_{-11} \text{ MeV}$ $\Gamma = 240 \pm 50 \text{ MeV}$

Results $J^{PC} = 1^{-+}$



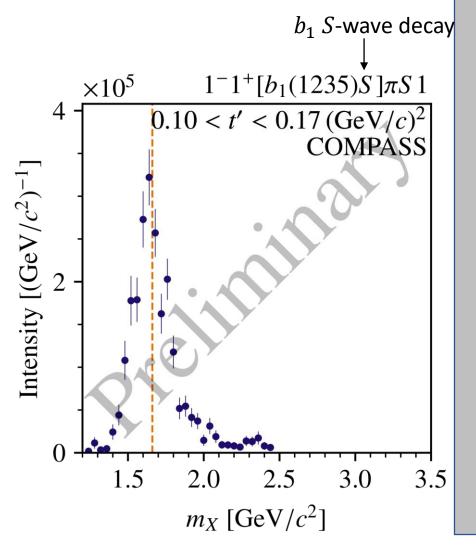


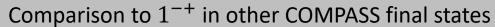


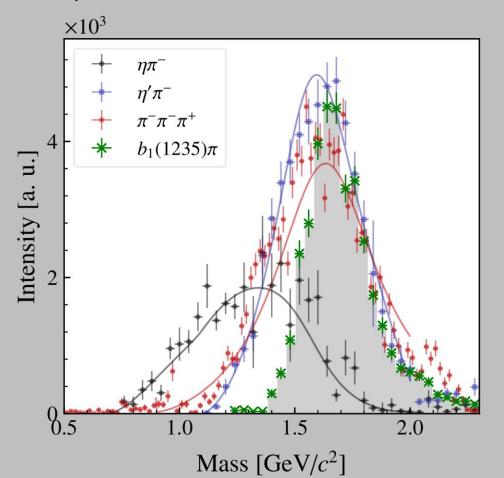
ites listed in PDG

 $\pi_1(1600)$ = 1661^{+15}_{-11} MeV
= 240 ± 50 MeV

Results $J^{PC} = 1^{-+}$





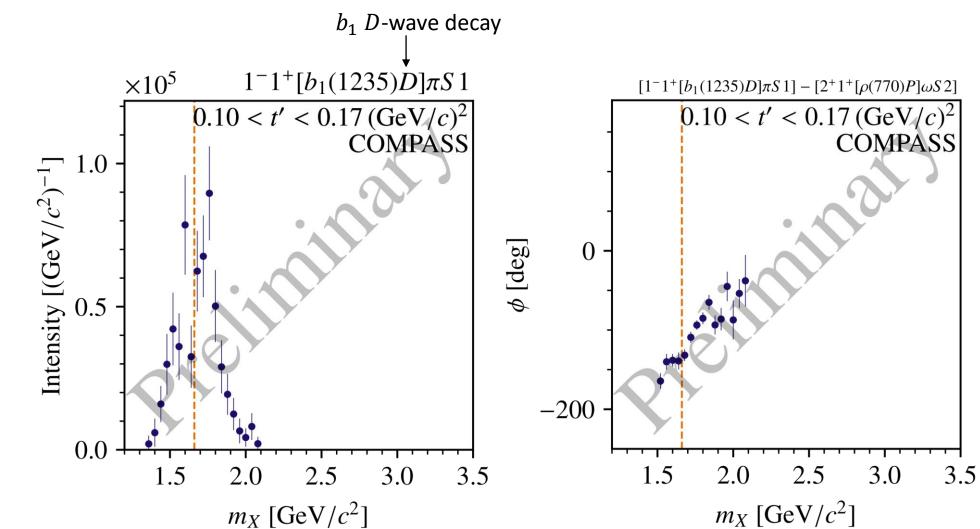


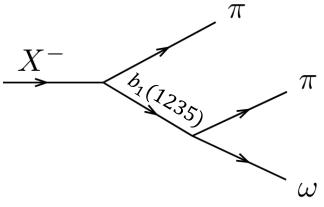
ω

ites listed in PDG

 $\pi_1(1600)$ = 1661^{+15}_{-11} MeV
= 240 ± 50 MeV

Results $J^{PC} = 1^{-+}$

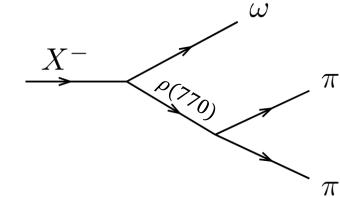


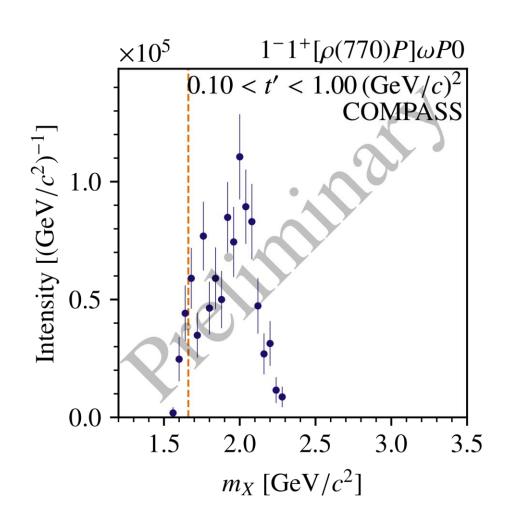


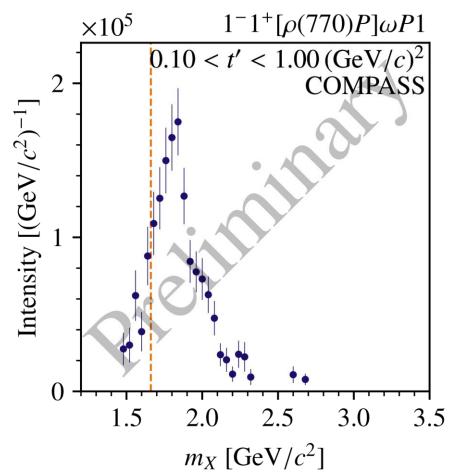
States listed in PDG

 $\pi_1(1600)$ $m = 1661^{+15}_{-11} \text{ MeV}$ $\Gamma = 240 \pm 50 \text{ MeV}$

Results $J^{PC} = 1^{-+}$







States listed in PDG

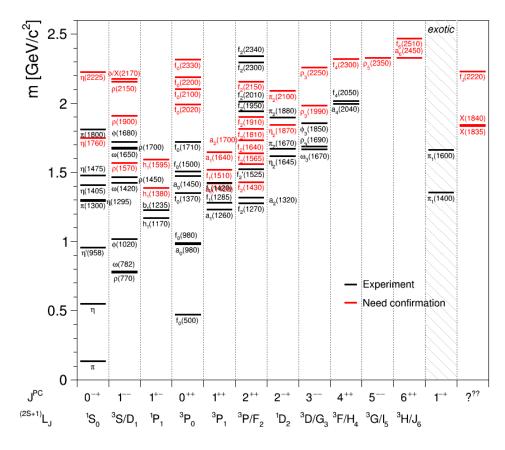
 $\pi_1(1600)$ $m = 1661^{+15}_{-11} \text{ MeV}$ $\Gamma = 240 \pm 50 \text{ MeV}$

Conclusion and Outlook

- Resonance-like signals for many well-established states visible
 - Clear peak for $\pi_1(1600) \to b_1(1235)\pi$
- Possible signals for further states: $a_3(1975), a_6(2450), \pi_1 \rightarrow \rho(770)\omega$
- Next step: Resonance-model fit to extract resonance parameters
 - First studies yield promising results

Backup

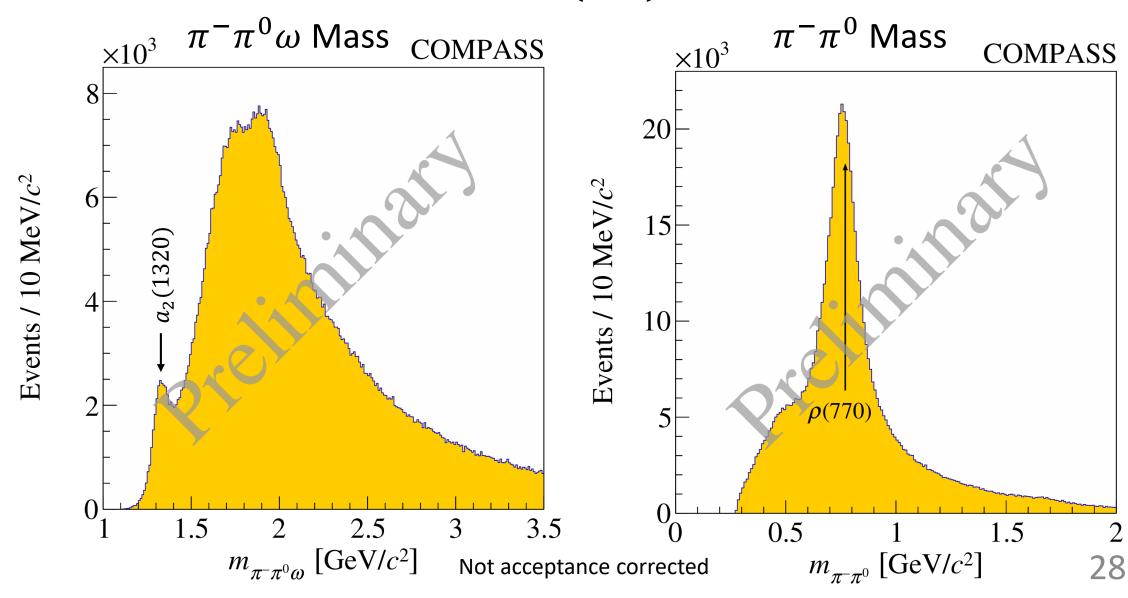
Mesons in QCD



- Many short-lived, exited states with similar masses
- ⇒ All possible intermediate states X for one final-state configuration interfere
- ⇒PWA necessary to determine contributions of certain X

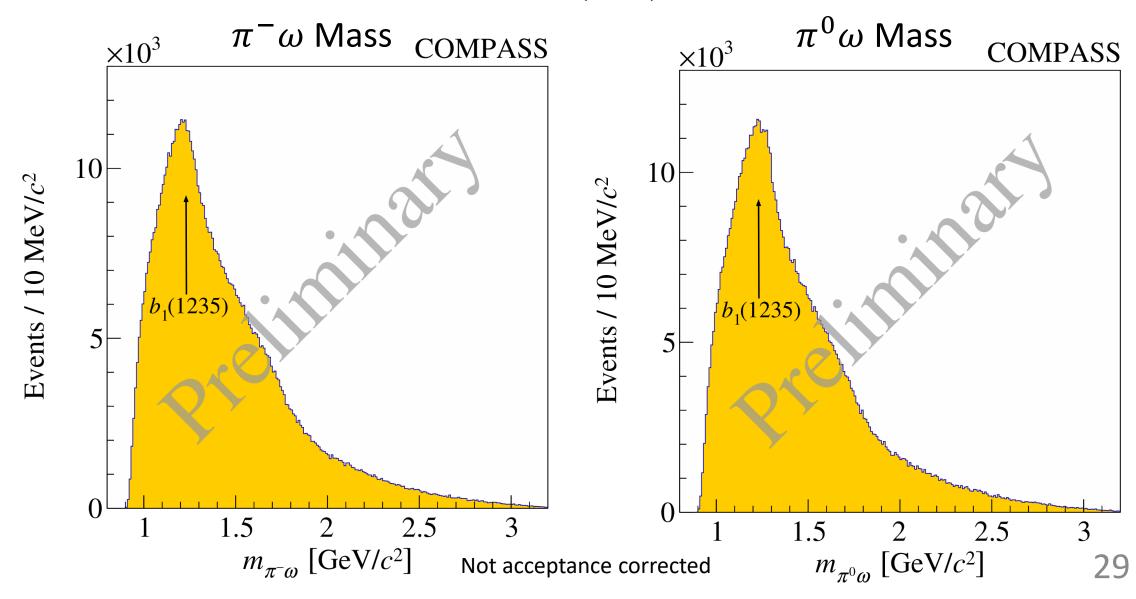
Kinematic Distributions - $\omega(782)\pi^-\pi^0$

• Total of 720,000 selected $\pi^-\pi^0\omega(782)$ events

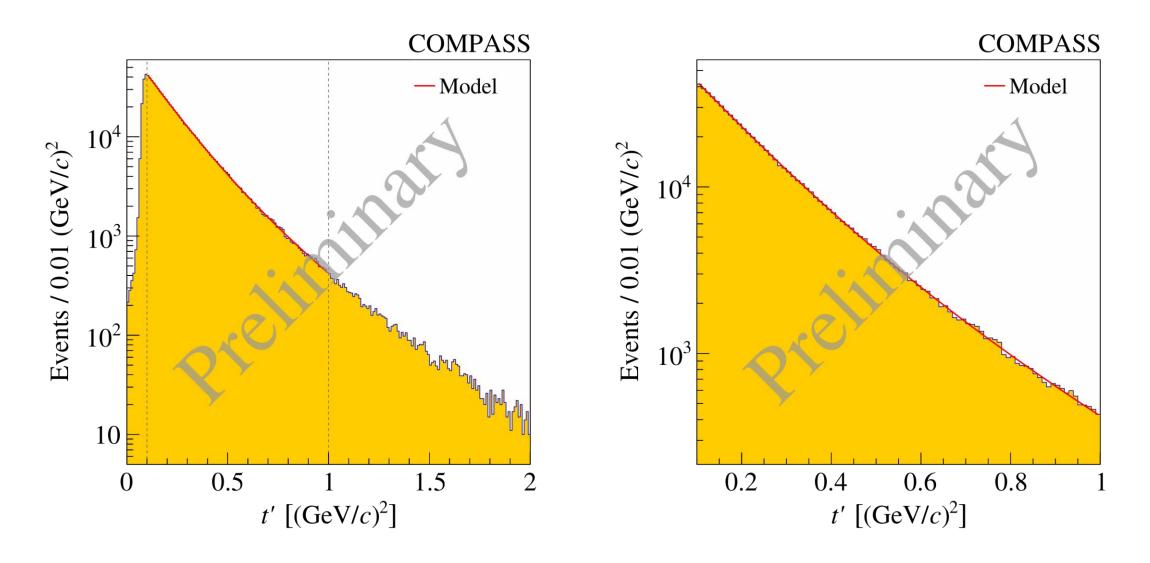


Kinematic Distributions - $\omega(782)\pi^-\pi^0$

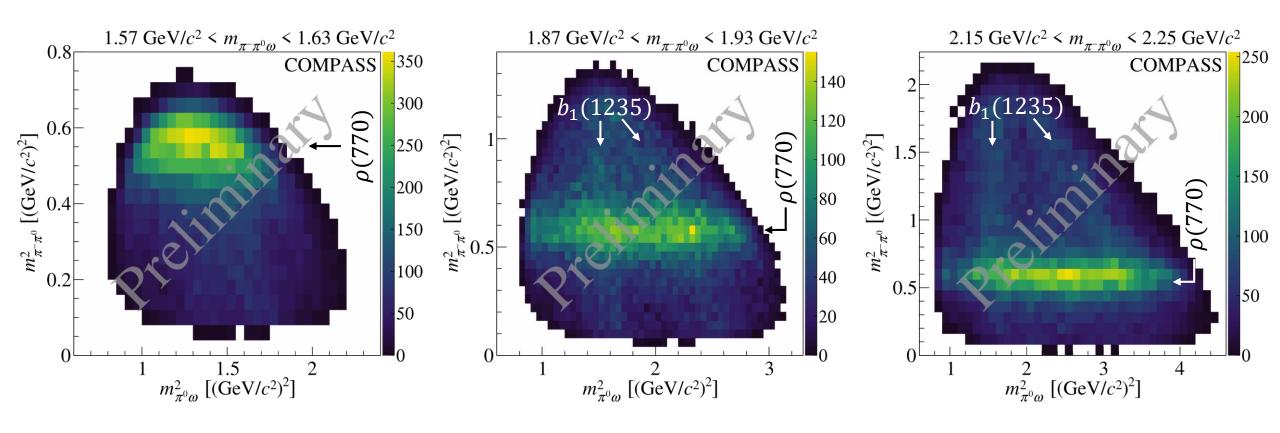
• Total of 720,000 selected $\pi^-\pi^0\omega(782)$ events



t' Distribution - $\omega(782)\pi^-\pi^0$

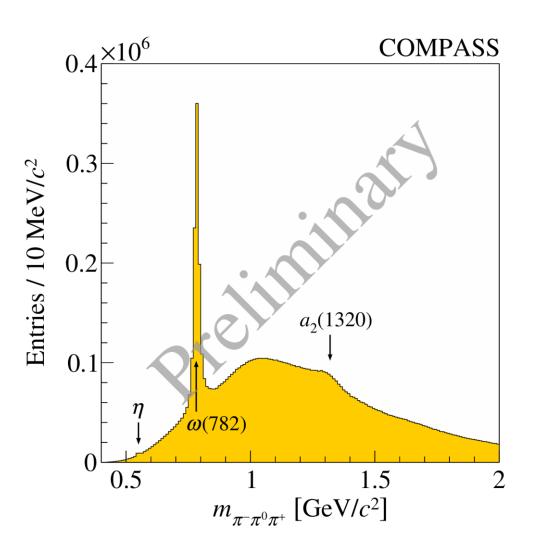


Dalitz Plots - $\omega(782)\pi^-\pi^0$



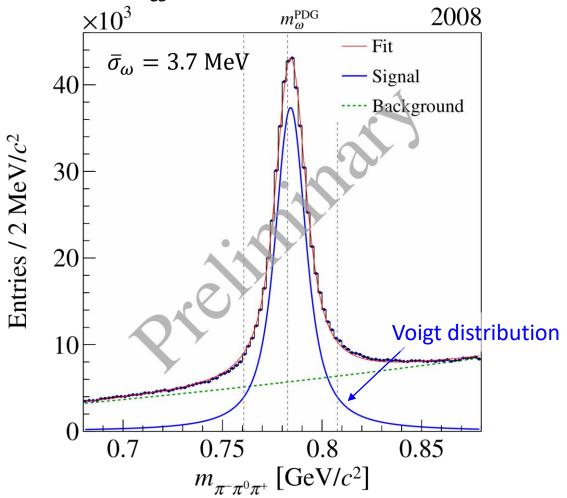
$\omega(782)$ Selection - $\omega(782)\pi^-\pi^0$

• Reconstruction of $\omega(782)$ from $\pi^-\pi^0\pi^+$ decay



$\omega(782)$ Selection - $\omega(782)\pi^-\pi^0$

- Reconstruction of $\omega(782)$ from $\pi^-\pi^0\pi^+$ decay
- Select events with exactly one $\pi^-\pi^0\pi^+$ combination within $\pm 3\sigma_\omega$ around the fitted m_ω



Partial-Wave Decomposition

$$I(m_X, t', \tau) = \left| \sum_i \mathcal{T}_i(m_X, t') \psi_i(m_X, \tau) \right|^2$$

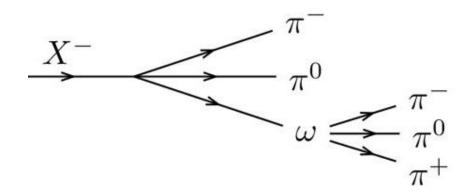
- Decay amplitude $\psi_i(m_X, \tau)$: calculated using the isobar model
- $\mathcal{T}_i(m_X,t')$ contains production, propagation, and coupling of
 - No assumptions about the resonant content of X^-
- Extract $\mathcal{T}_i(m_X,t')$ by independent maximum-likelihood fits of $I(\tau)$ in bins of (m_X,t')
 - Approximate \mathcal{T}_i by fitting step-wise constant functions in bins of (m_X, t')

$\omega(782)$ Decay in PWA Model

Factorisation of the decay amplitude

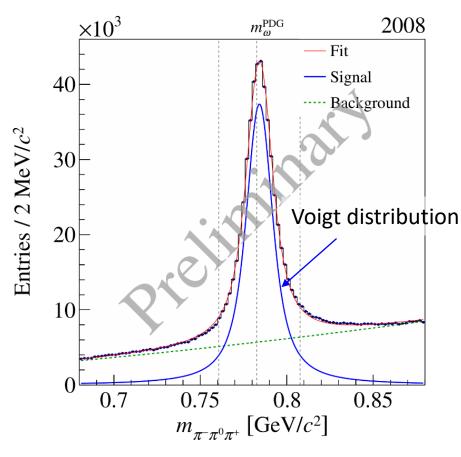
$$\psi_i = \Sigma_{\lambda_\omega} \psi_{i,X \to \omega \pi \pi}^{\lambda_\omega} \psi_{\omega \to 3\pi}^{\lambda_\omega}$$

- $\psi_{i,X o \omega\pi\pi}^{\lambda_\omega}$ calculated with isobar model
- $\psi_{\omega \to 3\pi}^{\lambda_{\omega}} = \mathcal{D}(m_{\omega})D_0^{\lambda_{\omega}}|p^+ \times p^-|$
 - $\mathcal{D}(m_{\omega})$ is the Breit-Wigner (BW) of ω
 - $D_0^{\lambda_\omega}$ and $|p^+ \times p^-|$ describe the orientation of ω and its P-wave Dalitz plot, respectively
 - Both are independent of m_{ω}



$\omega(782)$ Decay in PWA Model

- Problem: m_{ω} is only measured with limited resolution
 - \Rightarrow Intensity level: Convolution of BW with resolution function => m_{ω} follows Voigt distribution
 - ⇒ Convolution of the full intensity is not feasible
 - Solution: Neglect self-interference of ω as only one $\pi^-\pi^0\pi^+$ combination has a large amplitude
 - $\Rightarrow \mathcal{D}(m_{\omega})$ factorises out of the intensity: $I(m_X, t', \tau, m_{\omega}) = \tilde{I}(m_X, t', \tau) |\mathcal{D}(m_{\omega})|^2$
 - $\Rightarrow |\mathcal{D}(m_{\omega})|^2$ is modelled as Voigt distribution with parameters from fitted data

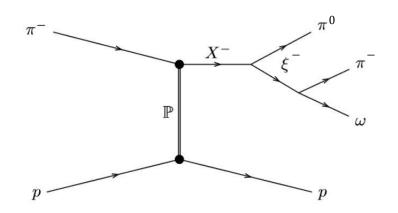


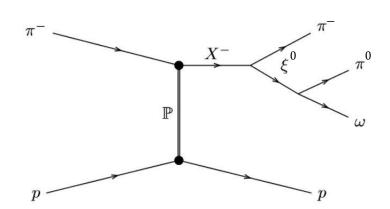
Isospin Symmetrization

• $X^- \to \xi^- \pi^0$ and $X^- \to \xi^0 \pi^-$ have the same amplitude (modulo a sign due to isospin Clebsch-Gordons)

 $\Rightarrow \mathcal{T}_i(m_X,t')$ is the same and we model the total decay amplitude as

$$\psi_i = +\frac{1}{2}\psi_{i,\xi^0\pi^-} - \frac{1}{2}\psi_{i,\xi^-\pi^0}$$

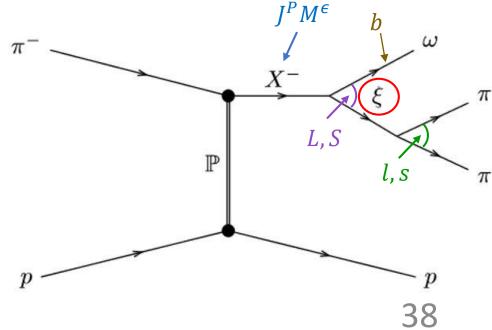




Wave Selection

- Method used for 3π , 5π and $K\pi\pi$
- Modified log-likelihood with penalties:
 - Cauchy regularization to suppress small waves
 - Connected bins over m_X to smoothen $\mathcal{T}_i(m_X)$
- Wave pool:
 - $J \le 8, M \le 2, \epsilon = +$
 - $\xi \to \pi \pi$: $\rho(770)$, $\rho(1450)$, $\rho_3(1690)$
 - $\xi \to \omega \pi$: $b_1(1235)$, $\rho(1450)$, $\rho_3(1690)$
 - *L* ≤ 8
 - 893 waves + flat wave

Notation: $i = I^P M^{\epsilon} [\xi l] b LS$



Flat Wave

- Isotropic in 5-body phase-space
- Used to describe background

