

Timelike pion form factor from lattice QCD

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June - 9 - 2023

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2023



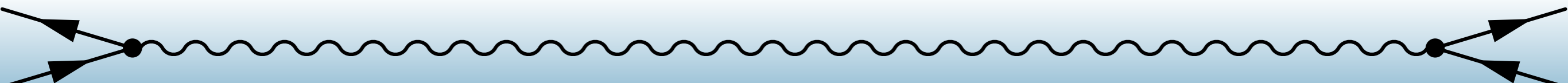
In collaboration with J. Dudek



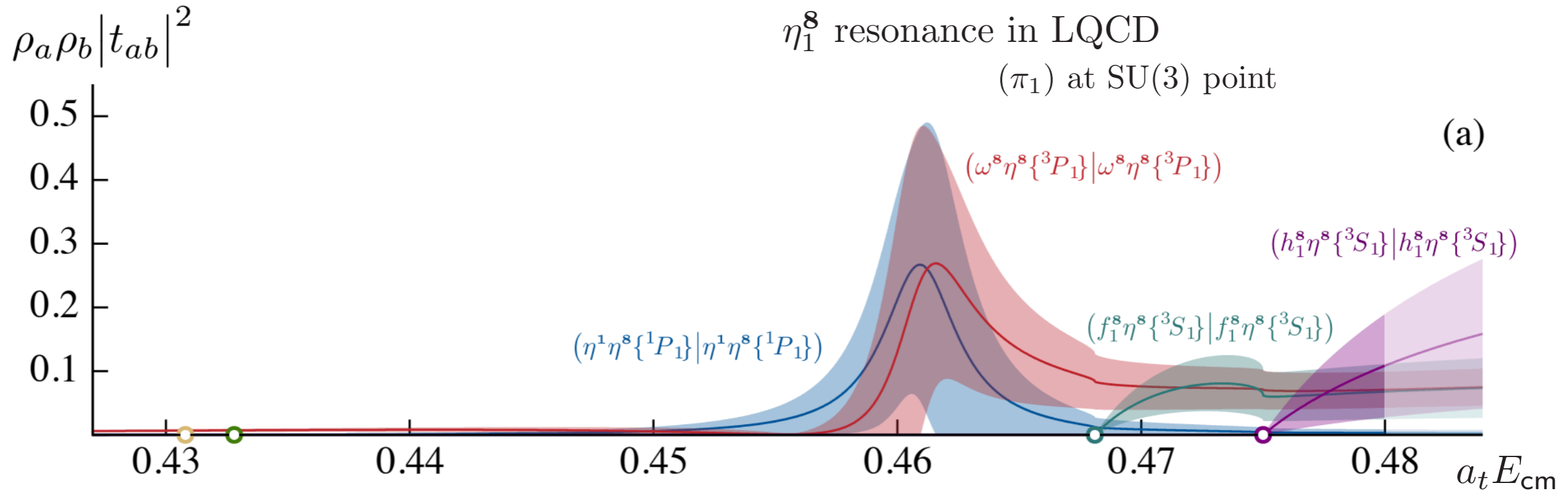
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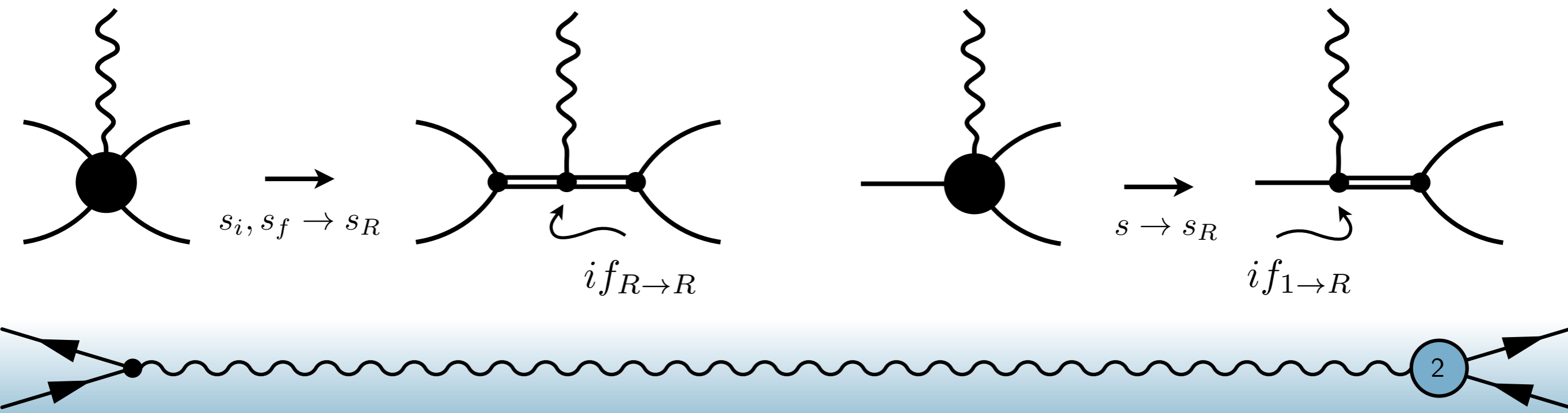
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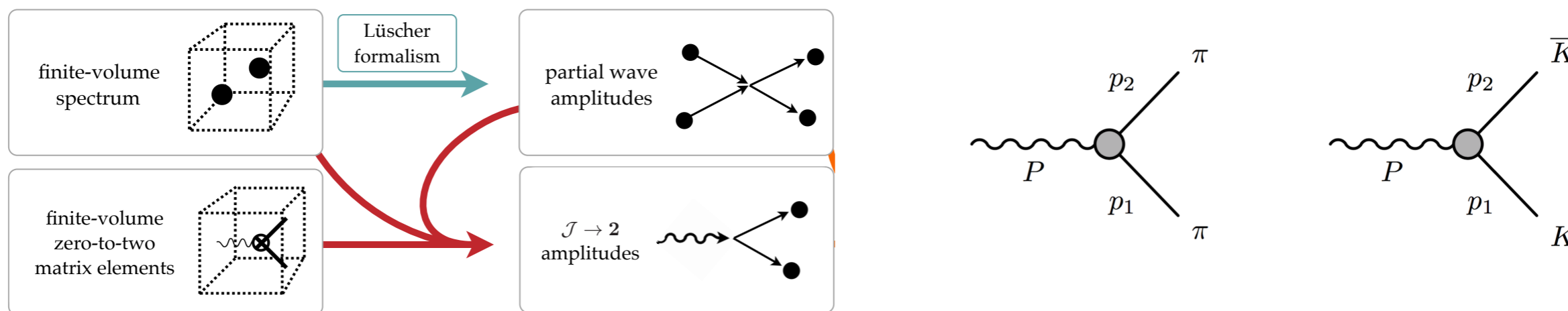
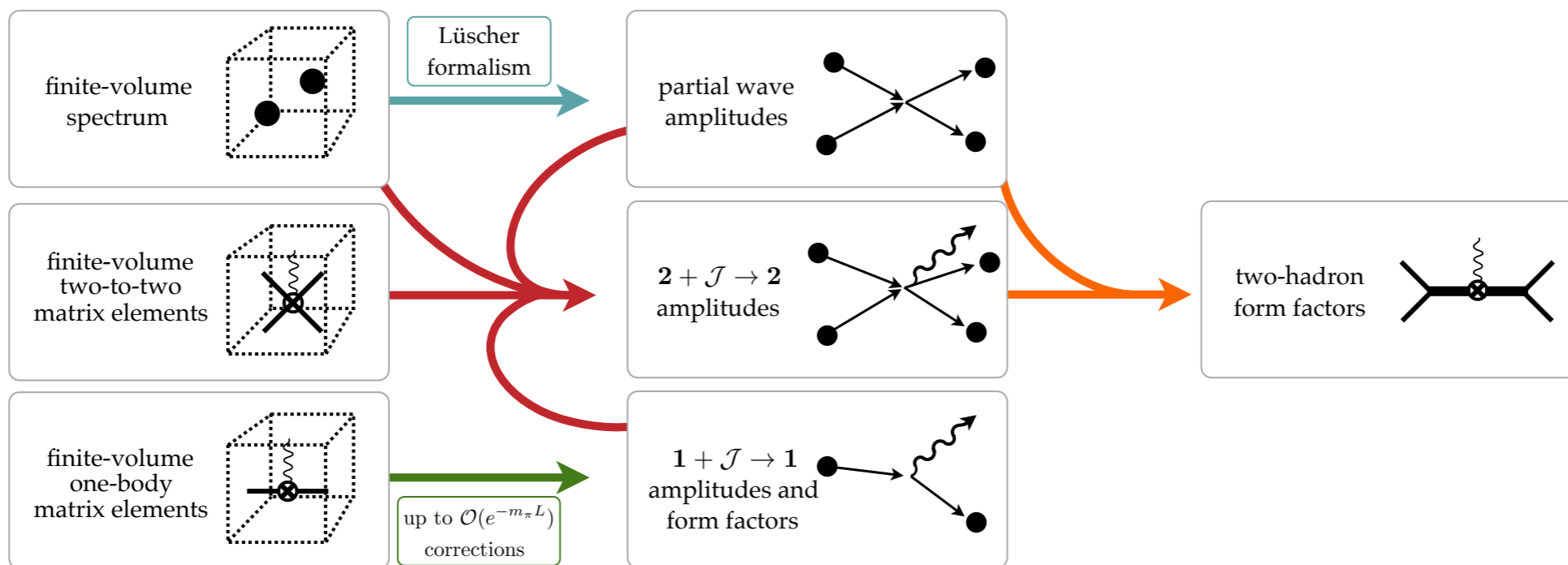
Resonance structure and production



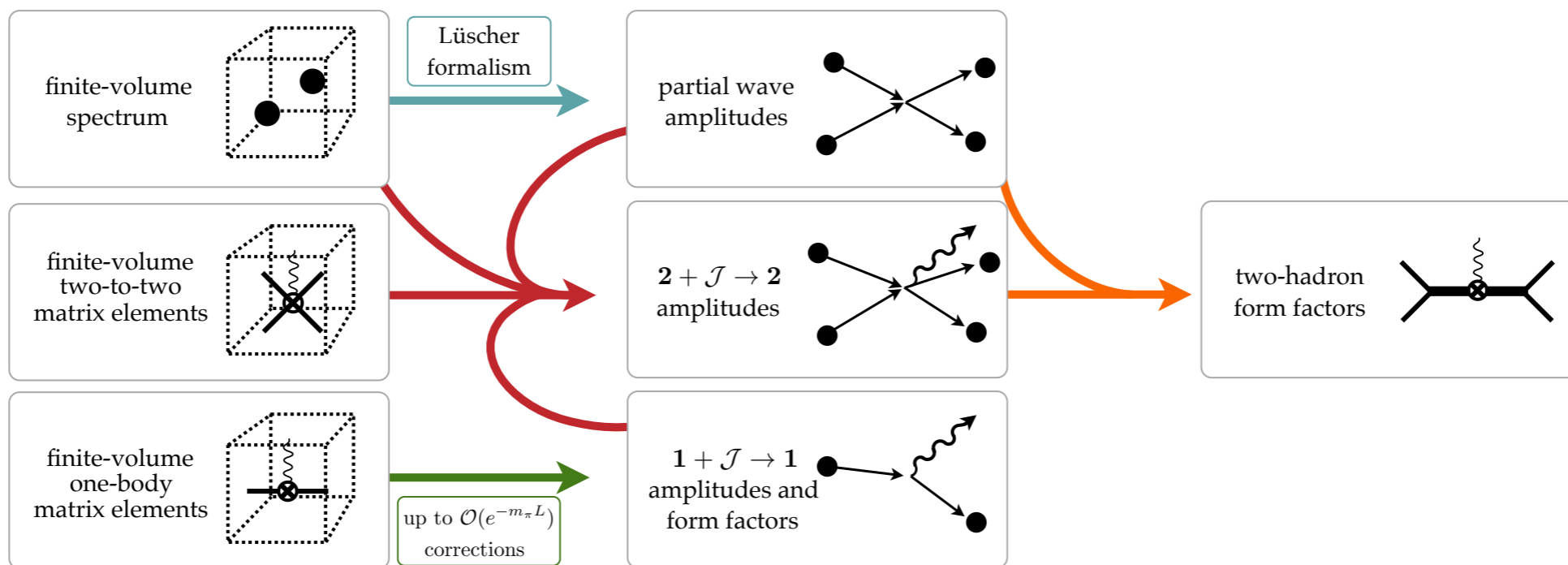
Diagonal elements of the scattering amplitude from [2009.10034 Woss et al.]



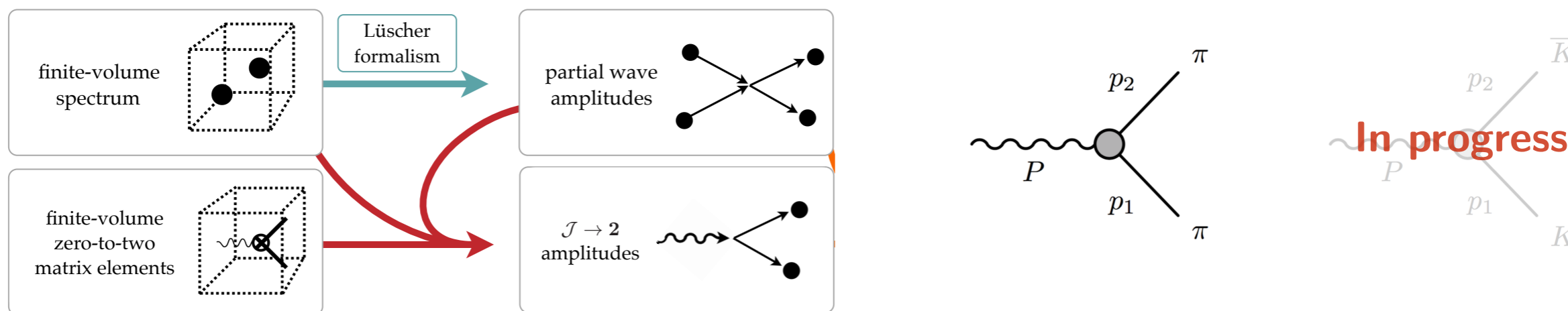
Finite-volume roadmap



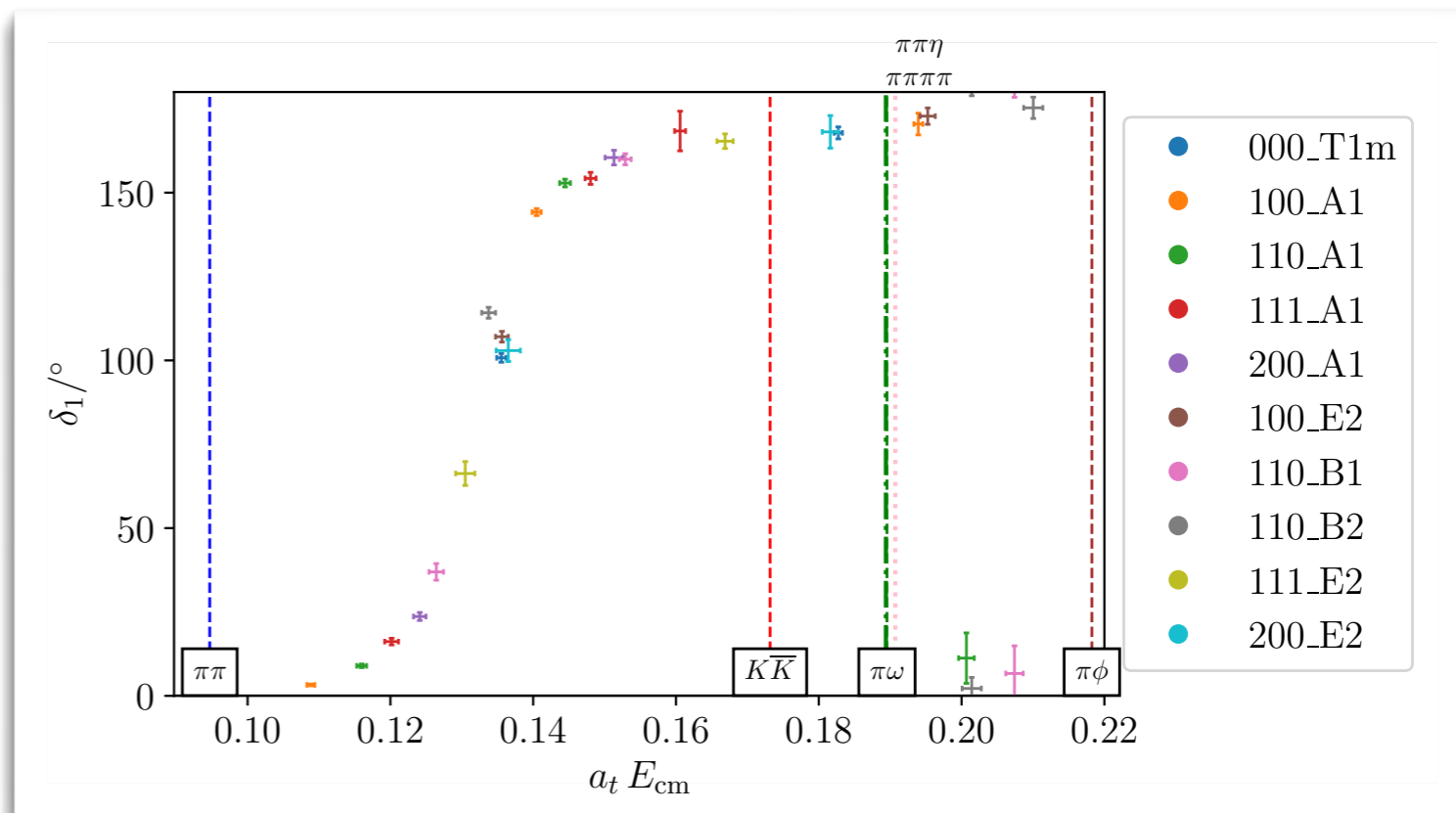
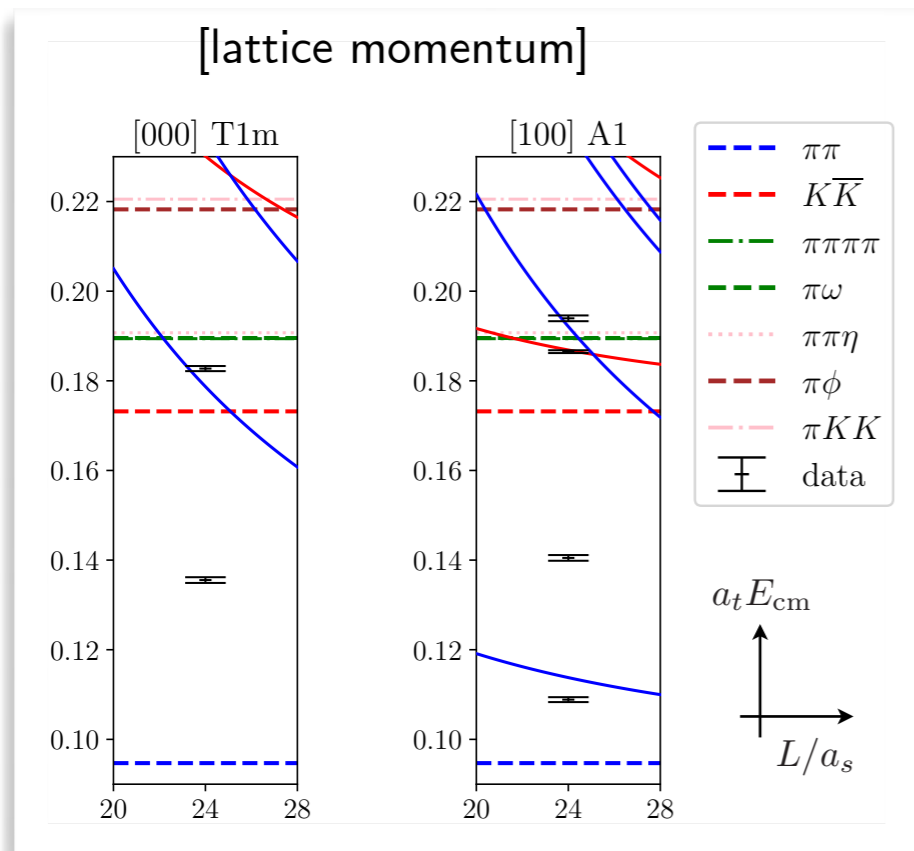
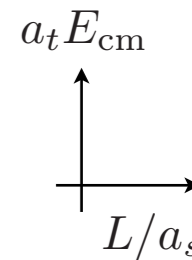
Finite-volume roadmap



This talk:



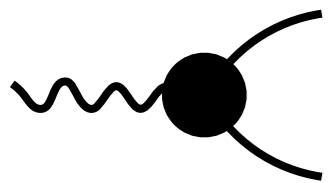
Scattering in the finite volume



$$\det \left(F^{-1}(E_n, L) + \mathcal{M}(E_n) \right) = 0$$

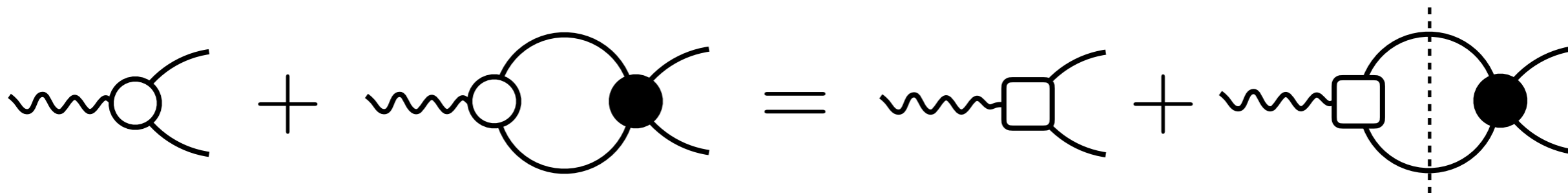
Known geometric function

Finite-volume
correlation function poles
in momentum space

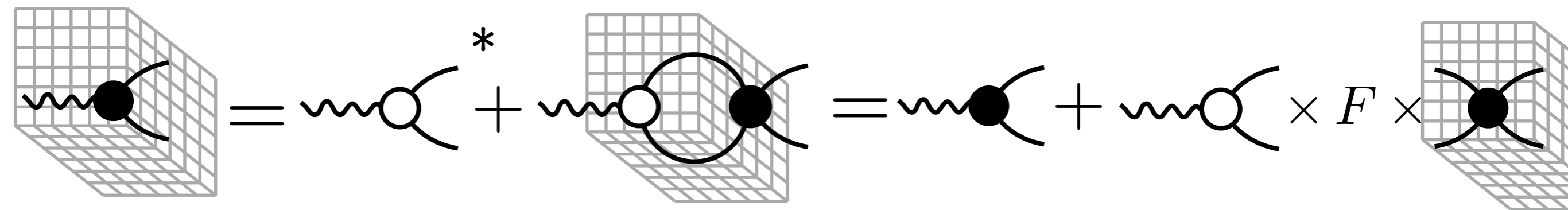


Finite volume corrections

Infinite volume: Watson's theorem. $\mathcal{H}(s) = \mathcal{A}_{02}(s) \mathcal{M}(s)$



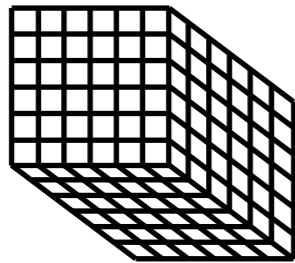
Finite volume: Lellouch-Lüscher factor.



$$\text{wavy line} \times \lim_{E \rightarrow E_n} \frac{E - E_n}{F^{-1} + \mathcal{M}} = \text{wavy line in finite volume}$$

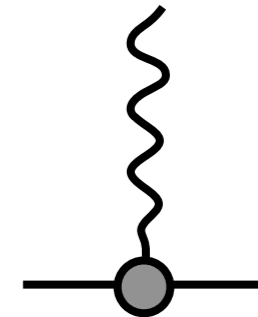
*Up to exponentially small corrections

Scattering and pair-production $J^P(I^G)=1^-(1^+)$

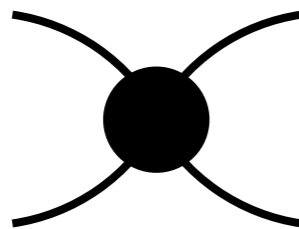


$$L = 2.7 \text{ fm}$$

	$a_t m$	m/MeV
π	0.0474	284
K	0.0866	519



For renormalization of the current

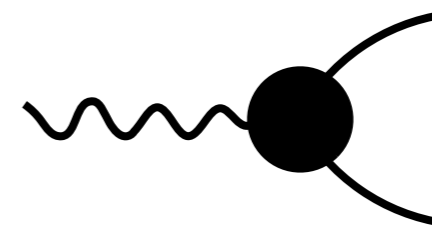


$$C_{ab}(t - t_s) = \langle O_a(t) O_b^\dagger(t_s) \rangle$$

◆ Two-meson operators

◆ $q\bar{q}$ operators

1. GEVP $C_{ab}(t)v_b^n = C_{ab}(t_0)v_b^n \lambda_n(t - t_0)$
2. Fit eigenvalues $\lambda_n(t - t_0) = e^{-E_n(t-t_0)}$

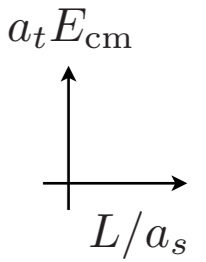


$$\langle \mathcal{J}(t) \Omega_n^\dagger(0) \rangle \propto \langle 0 | \mathcal{J}(0) | n \rangle$$

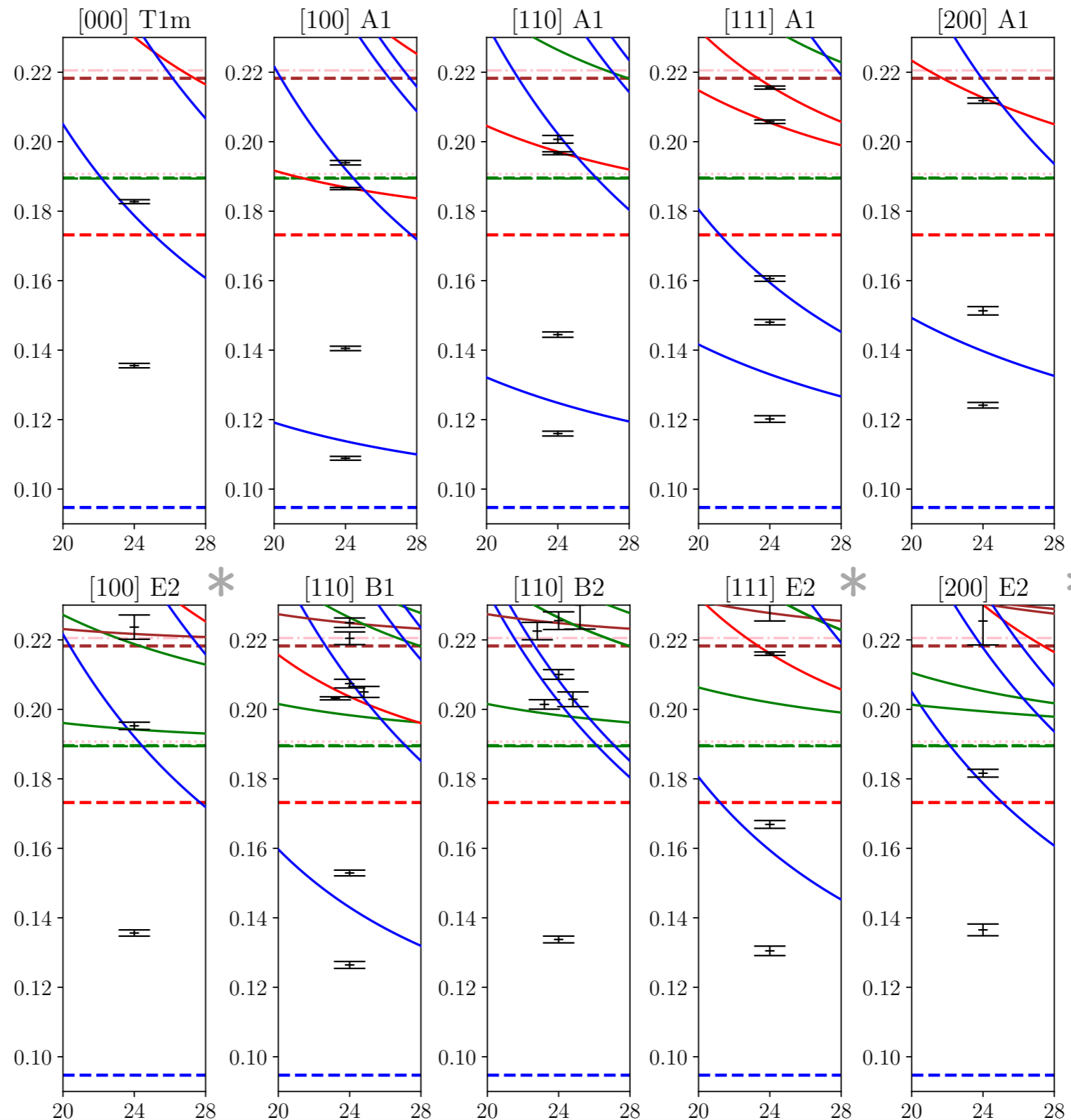
◆ Optimized operators

1. Extract invariant $\mathcal{H}_L^\mu = K^\mu \mathcal{F}_L$
2. FV correction $\mu_0'^* = \frac{\partial}{\partial E^*} (\mathcal{M}^{-1} + F)$

Finite volume spectrum



- 17 elastic levels
- 25 $\pi\pi$ -like levels
- 7 KK -like levels



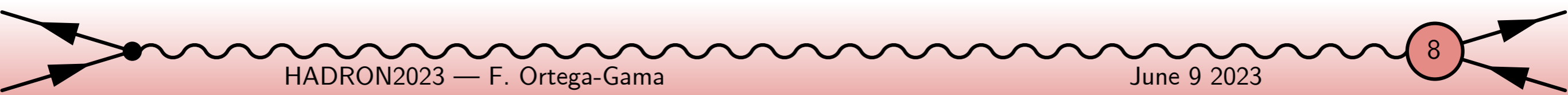
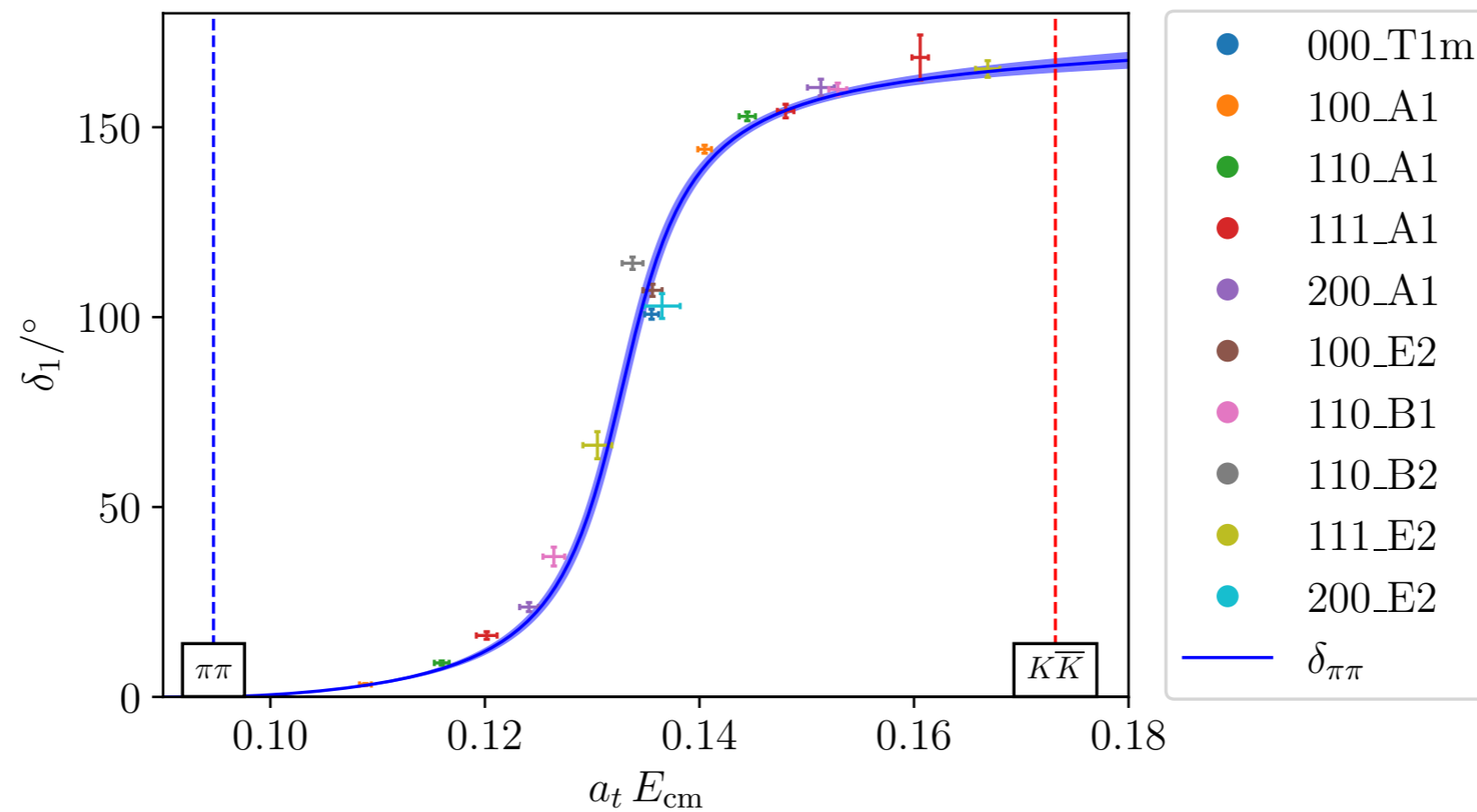
*without $\omega\pi$ operators

Elastic scattering

$$\mathcal{M}_\ell = \frac{E_{\text{cm}}}{2q^*} e^{i\delta_\ell} \sin \delta_\ell$$

$$\mathcal{M}_\ell^{-1} = \frac{1}{(2q^*)^\ell} K_\ell^{-1} \frac{1}{(2q^*)^\ell} - i\rho_{\text{CM}}$$

$$K_1 = \frac{g^2}{-s + m^2} + \gamma$$

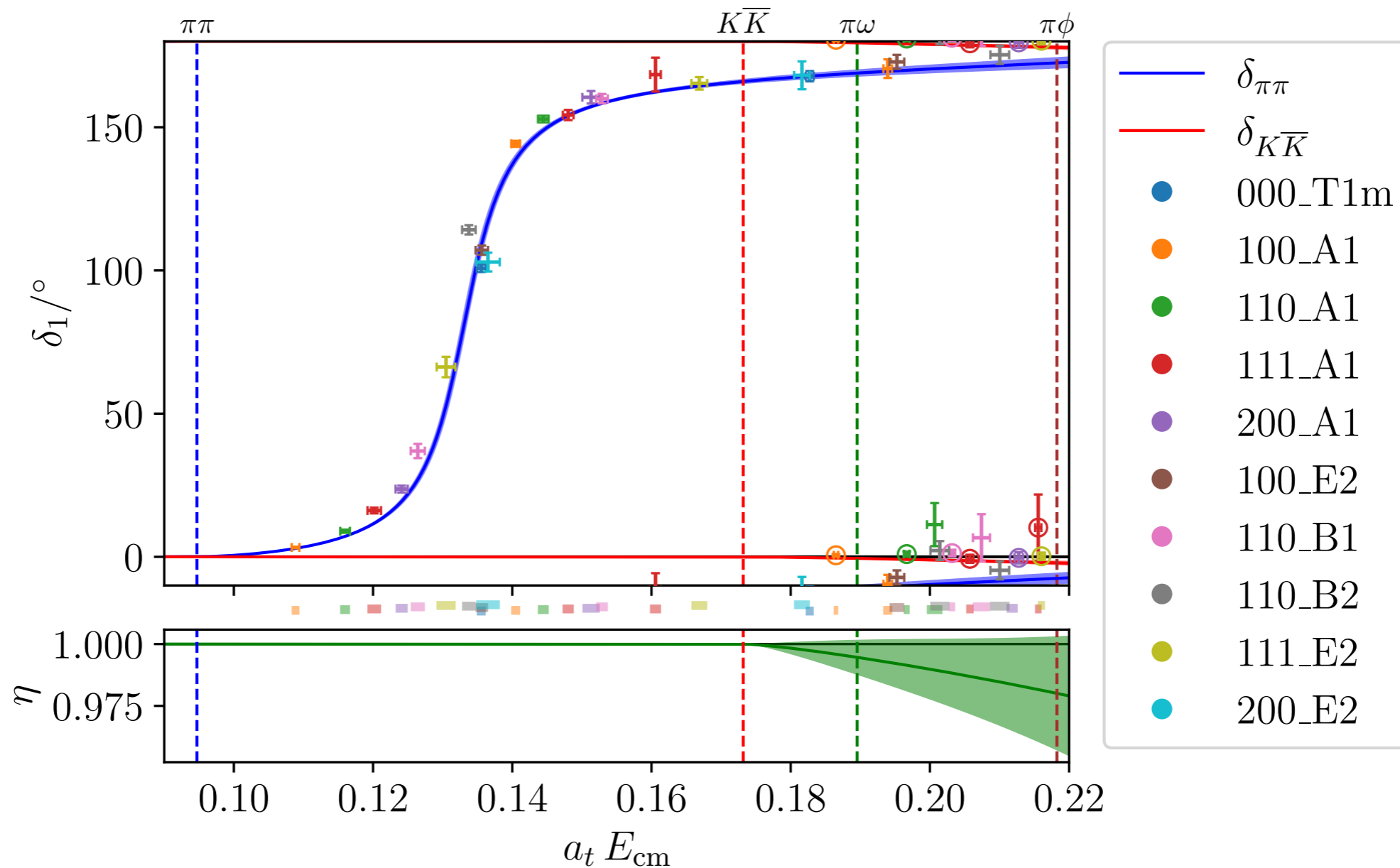


Coupled channel fit

$$\mathcal{M}_{\ell,aa} = \frac{\eta e^{2i\delta_{\ell,a}} - 1}{2i\rho_a}$$

$$K_{ab} = \frac{g_a g_b}{-s + m_r^2} + \gamma_{ab}$$

$$\mathcal{M}_{\ell,ab} = \frac{\sqrt{1 - \eta^2} e^{i(\delta_{\ell,a} + \delta_{\ell,b})}}{2\sqrt{\rho_a \rho_b}}$$



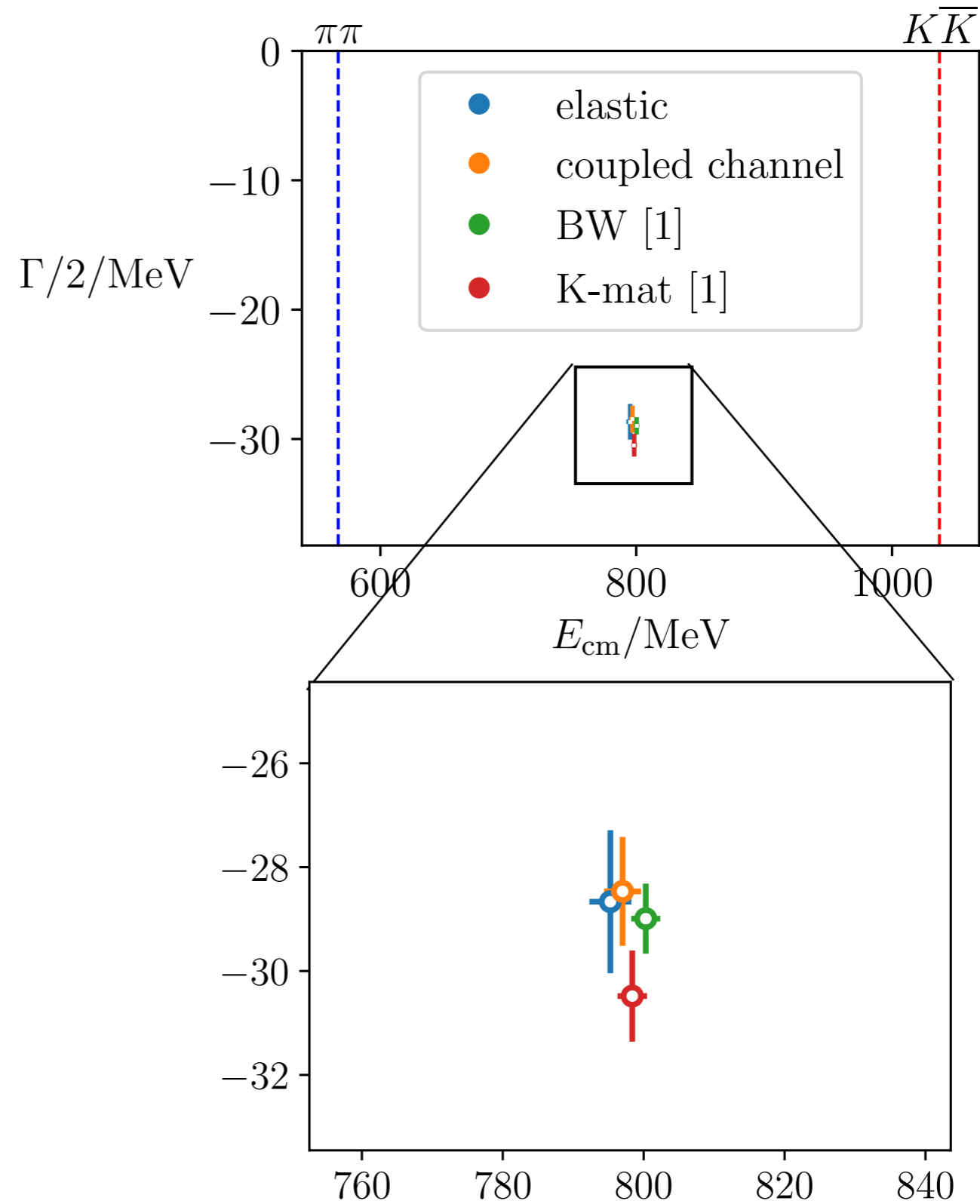
Rho resonance

$$\mathcal{M}(s) \sim \frac{g_R^2}{(m_\rho - i\Gamma_\rho/2)^2 - s}$$

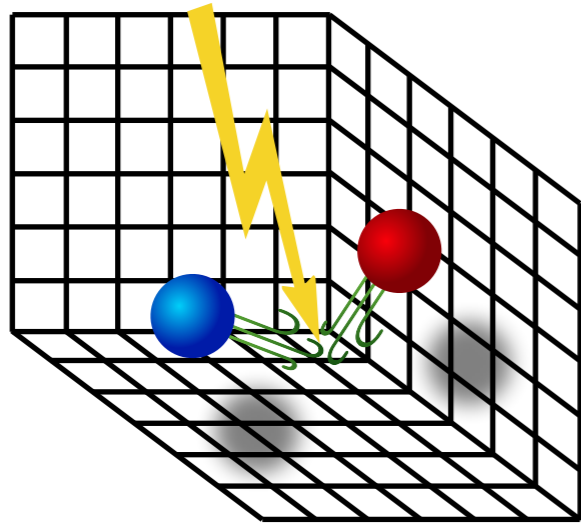
$$g_R = 254 \pm 5 \text{ MeV}$$

$$\text{Re}(m_\rho) = 797 \pm 2.6 \text{ MeV}$$

$$\Gamma_\rho/2 = 28.5 \pm 1.0 \text{ MeV}$$

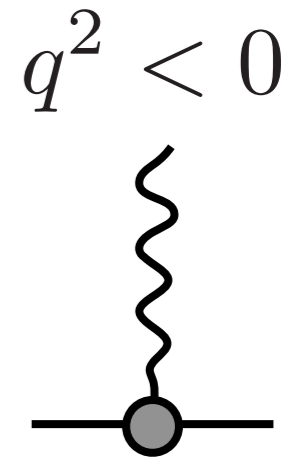
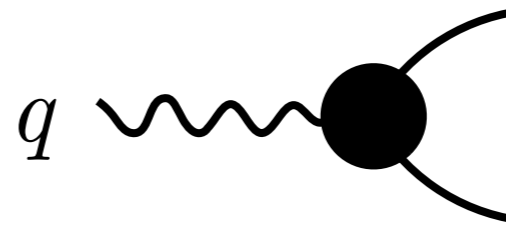


Pair production



$$\mathcal{H}_L^\mu = K^\mu \mathcal{F}_L$$

Finite-volume matrix elements



$$\mathcal{H}^\mu(q) = K^\mu f(q^2)$$

$$f = \frac{\mathcal{M}}{q^{*2}} \frac{1}{\tilde{r}_n} \mathcal{F}_L$$

Lellouch-Lüscher factor

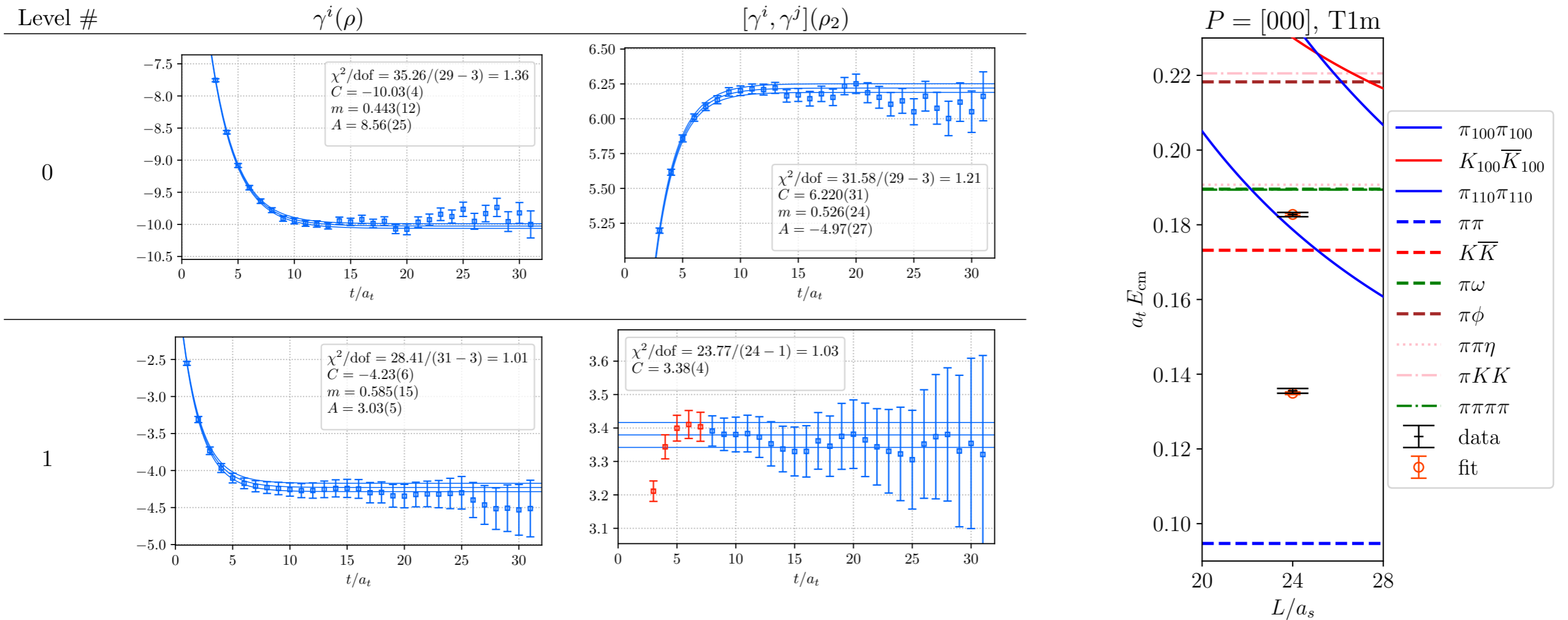
q^*

Relative momentum
in cm frame

Matrix elements

Optimized operators $\rightarrow \frac{\langle \mathcal{J}(t)\Omega_n^\dagger(0) \rangle}{\langle \Omega_n(t)\Omega_n^\dagger(0) \rangle} = \langle 0|\mathcal{J}(0)|n\rangle Z_n^{-1/2} + \mathcal{O}(e^{-(E_N-E_n)t})$

Fit function: $C + Ae^{-mt}$



$$\langle 0|\mathcal{J}_{\text{impro}}^\rho|n\rangle = \langle 0|\rho|n\rangle - \frac{1}{4}(1 - \xi)a_t E_n \langle 0|\rho_2|n\rangle$$

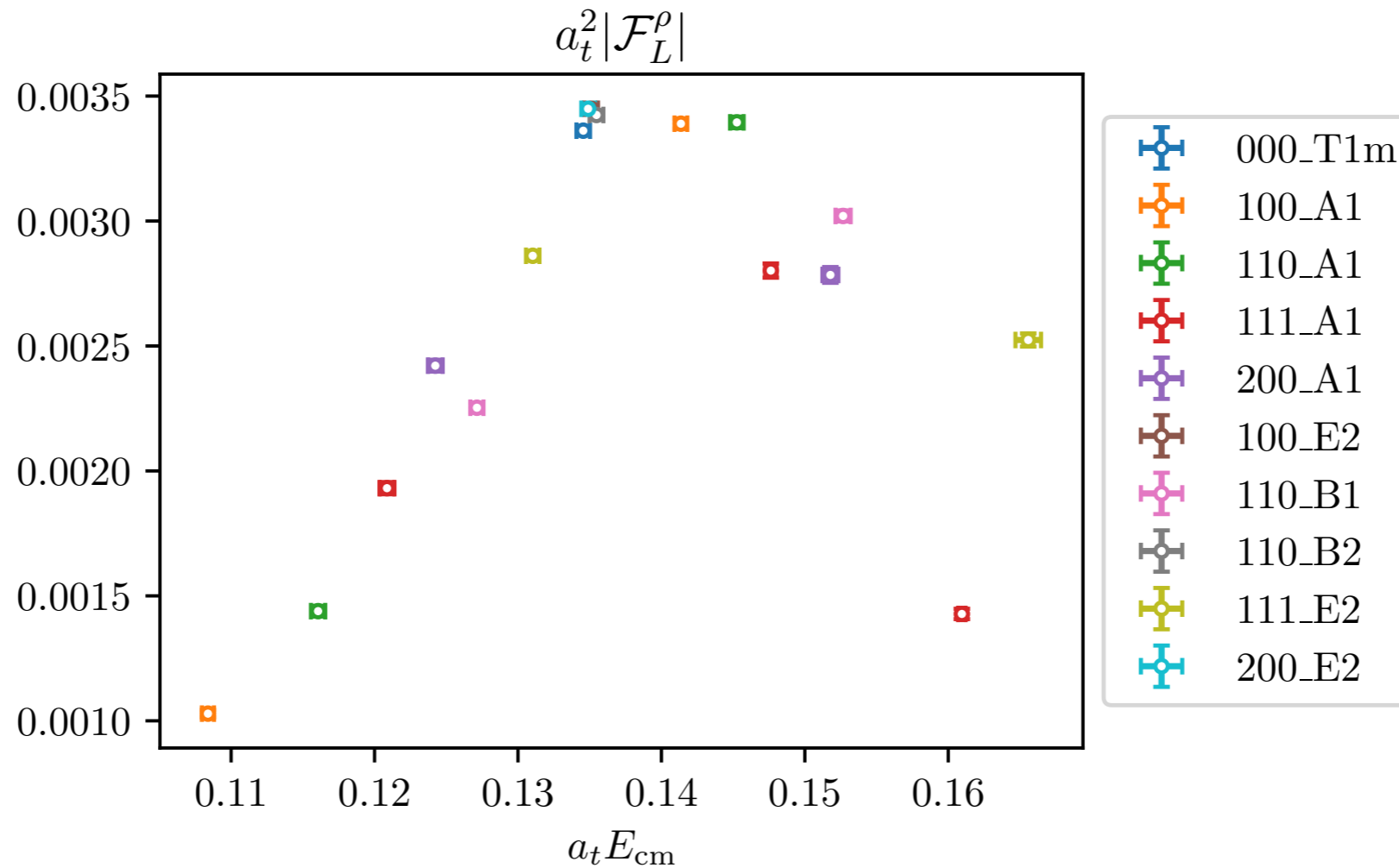
anisotropy

Invariant matrix element

$$\mathcal{H}_L^\mu = K^\mu \mathcal{F}_L$$

$$\langle 0 | \mathcal{J}^\mu(0) | n, P, 1\lambda \rangle = \frac{2}{\sqrt{3}} \epsilon^\mu(P, \lambda) \mathcal{F}_L(P^2)$$

$$\mathcal{F}_L(P^2) = \sqrt{\frac{2E_n}{L^3}} \frac{Z_n^{1/2} Z_V C_n^{\mathbf{d}, \Lambda, \rho}}{K_n(\vec{d}, \Lambda, J)}$$

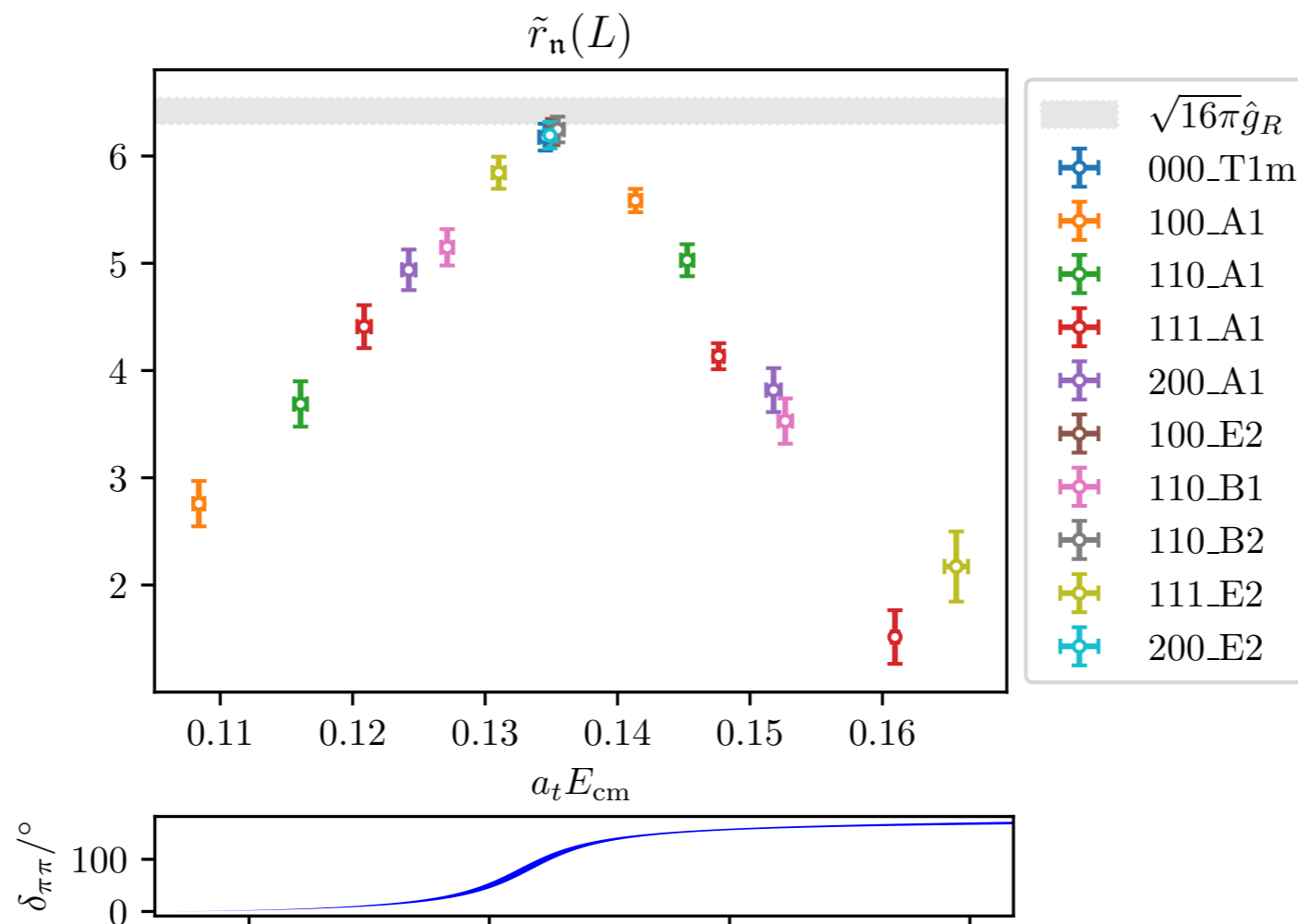


Finite volume correction

- Lellouch-Lüscher factor

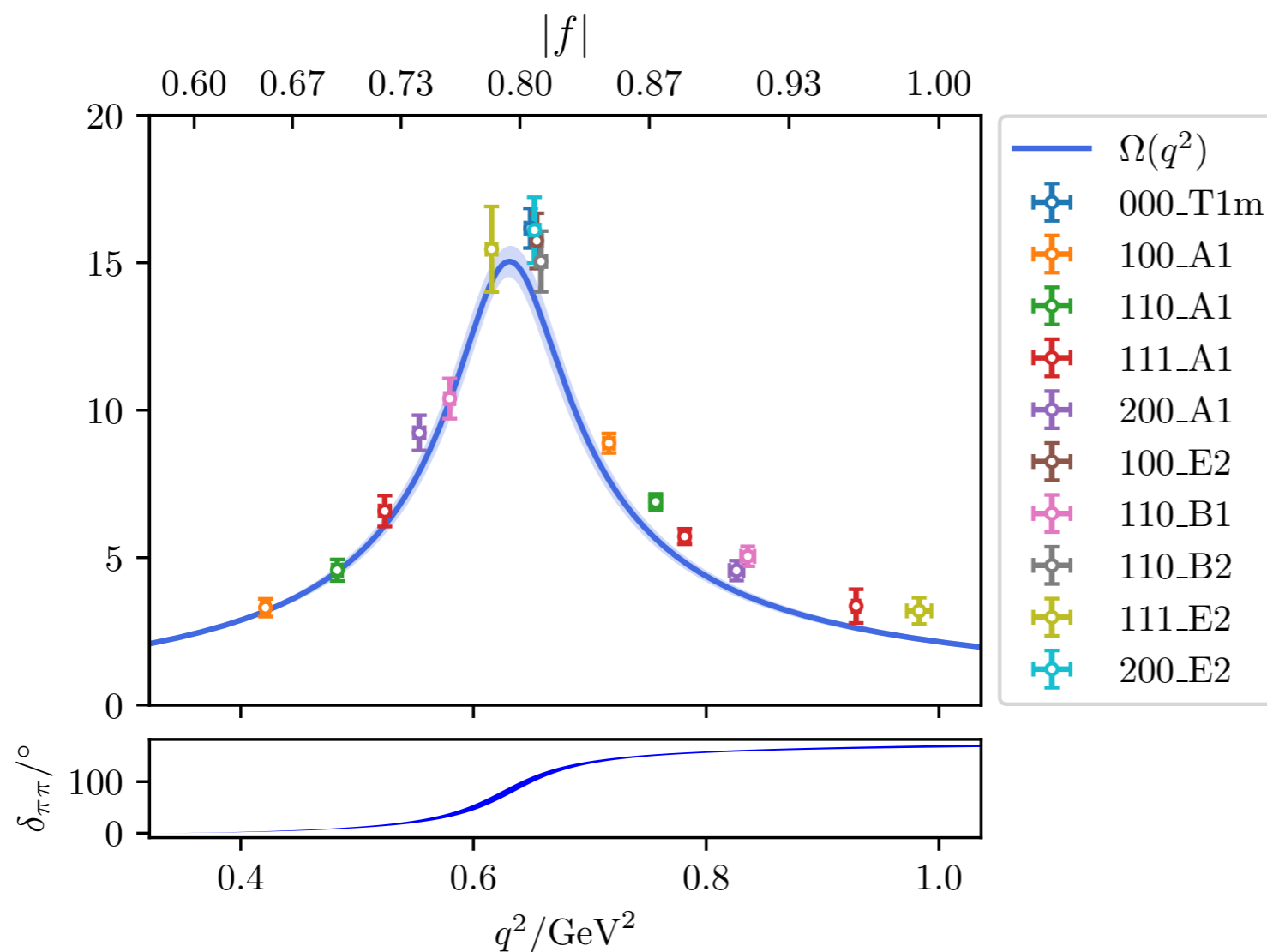
$$\mu_0'^* = \frac{\partial}{\partial E^*} (\mathcal{M}^{-1} + F) \quad \tilde{r}_n = \frac{1}{q^*} \sqrt{\frac{-2E_n^*}{\mu_0'^*}} \sim \sqrt{16\pi g_R/q^*}$$

In the narrow width limit



Timelike form factor

$$f = \frac{\mathcal{M}}{q^{*2}} \frac{1}{\tilde{r}_n} \mathcal{F}_L$$



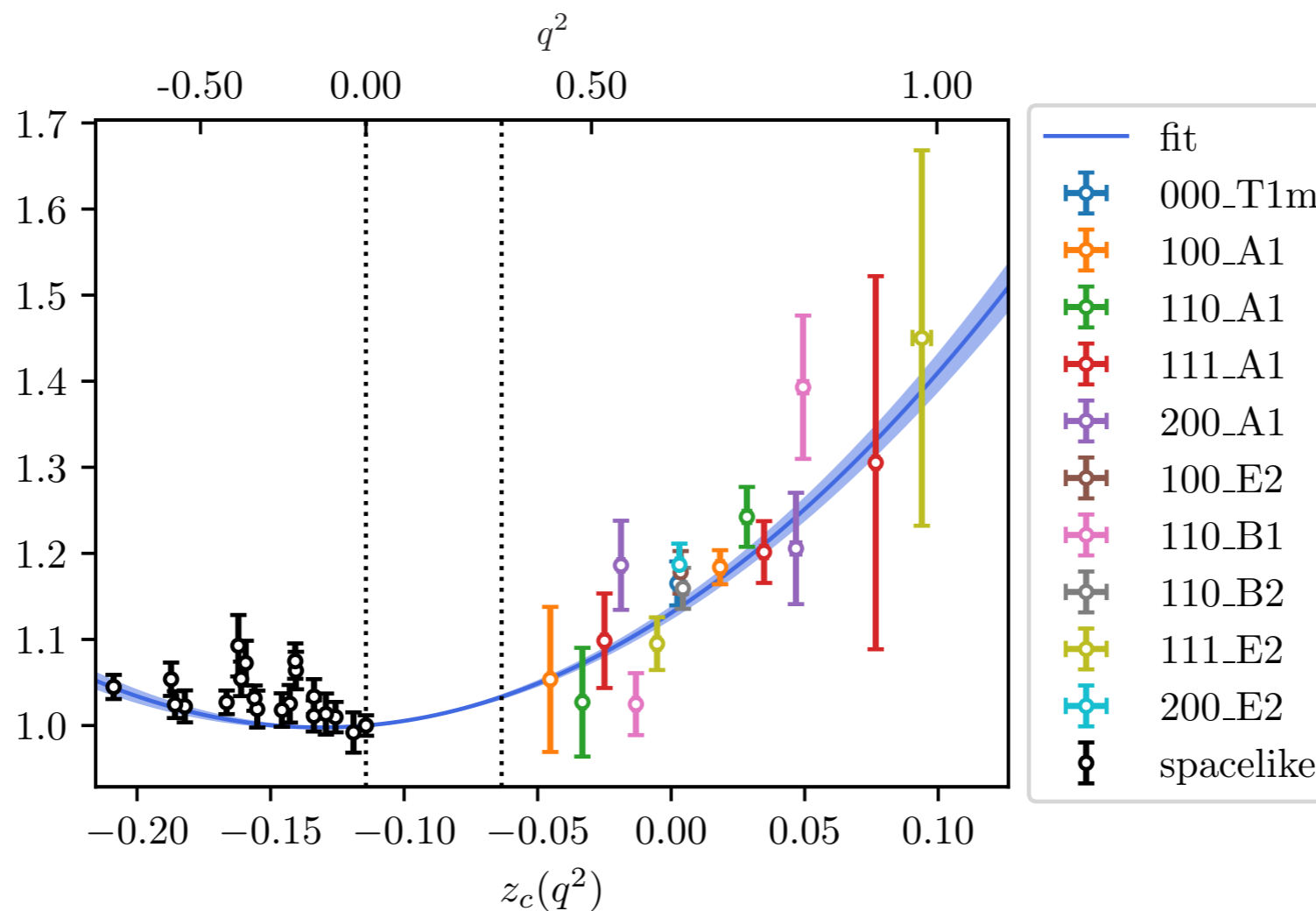
Watson's theorem:

$$f = \Omega \times \mathcal{F}$$

$$\Omega(s) = \exp\left(\frac{s}{\pi} \int_{s_{\text{thr}}}^{\infty} ds' \frac{\delta(s')}{s'(s' - s)}\right)$$

Smooth function fit

$$\mathcal{F} = f/\Omega$$

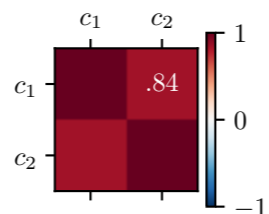


$$\chi^2/n_{\text{dof}} = 87.7/(37 - 2) = 2.51$$

$$Q \text{ [fixed]} = 1(0)$$

$$c_1 = 2.03(11)$$

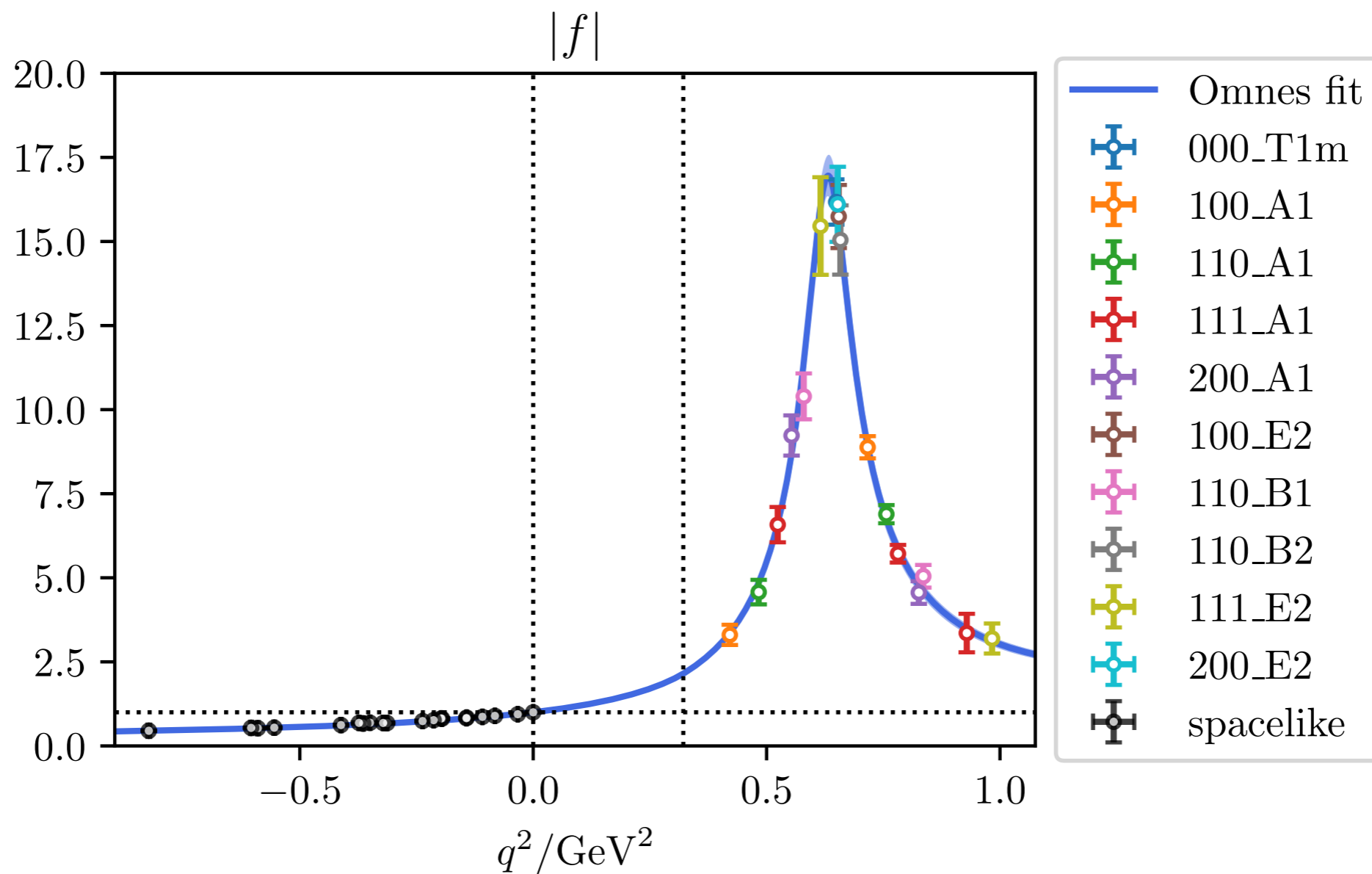
$$c_2 = 7.7(5)$$



$$\mathcal{F}(q^2) = Q + \sum_{n=1}^N c_n (z_c(q^2) - z_c(0))^n$$

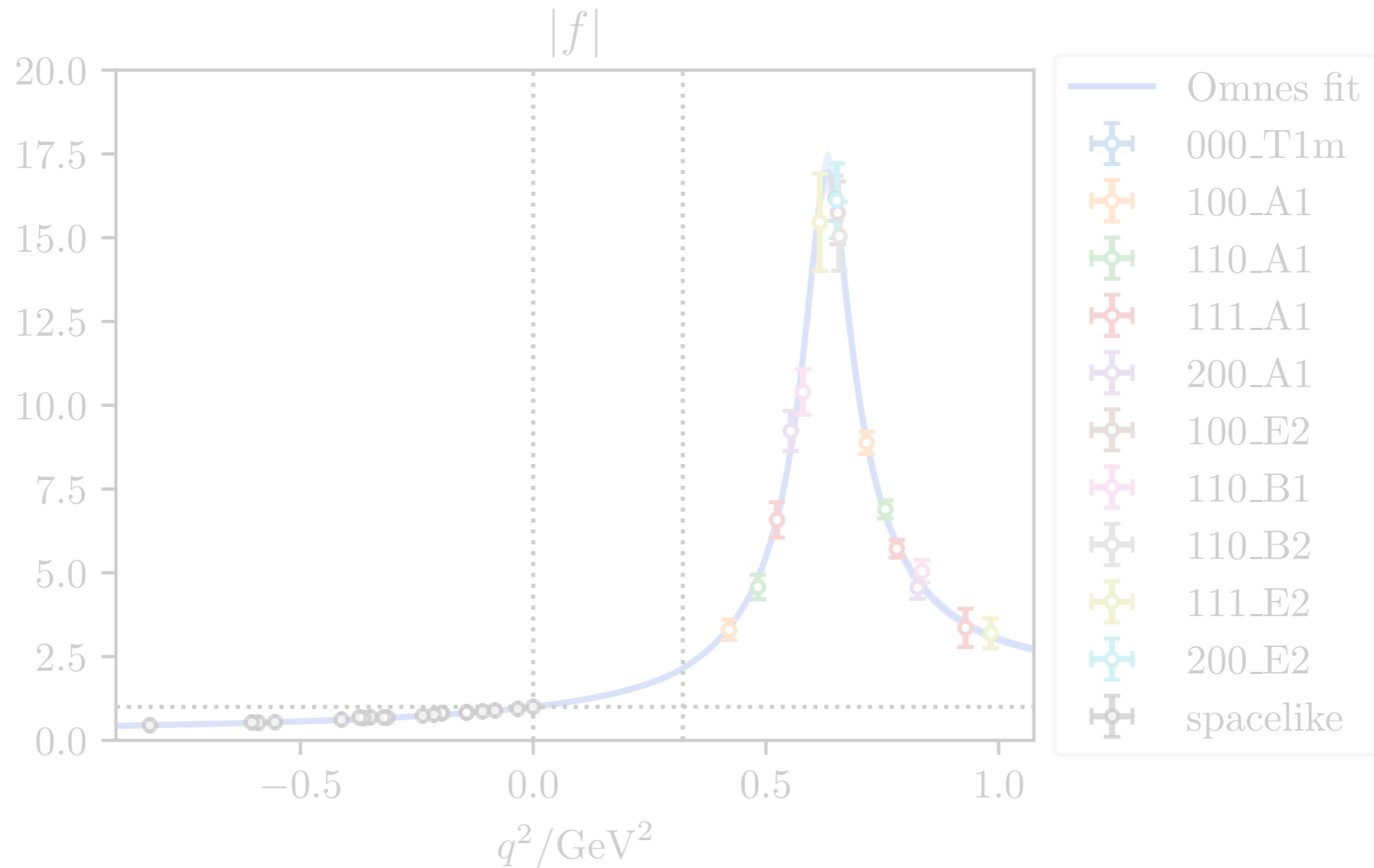
$$z_c(q^2) = \frac{\sqrt{s_c - s_{in}} - \sqrt{s_c - q^2}}{\sqrt{s_c - s_{in}} + \sqrt{s_c - q^2}}$$

Form factor fit: spacelike and timelike region



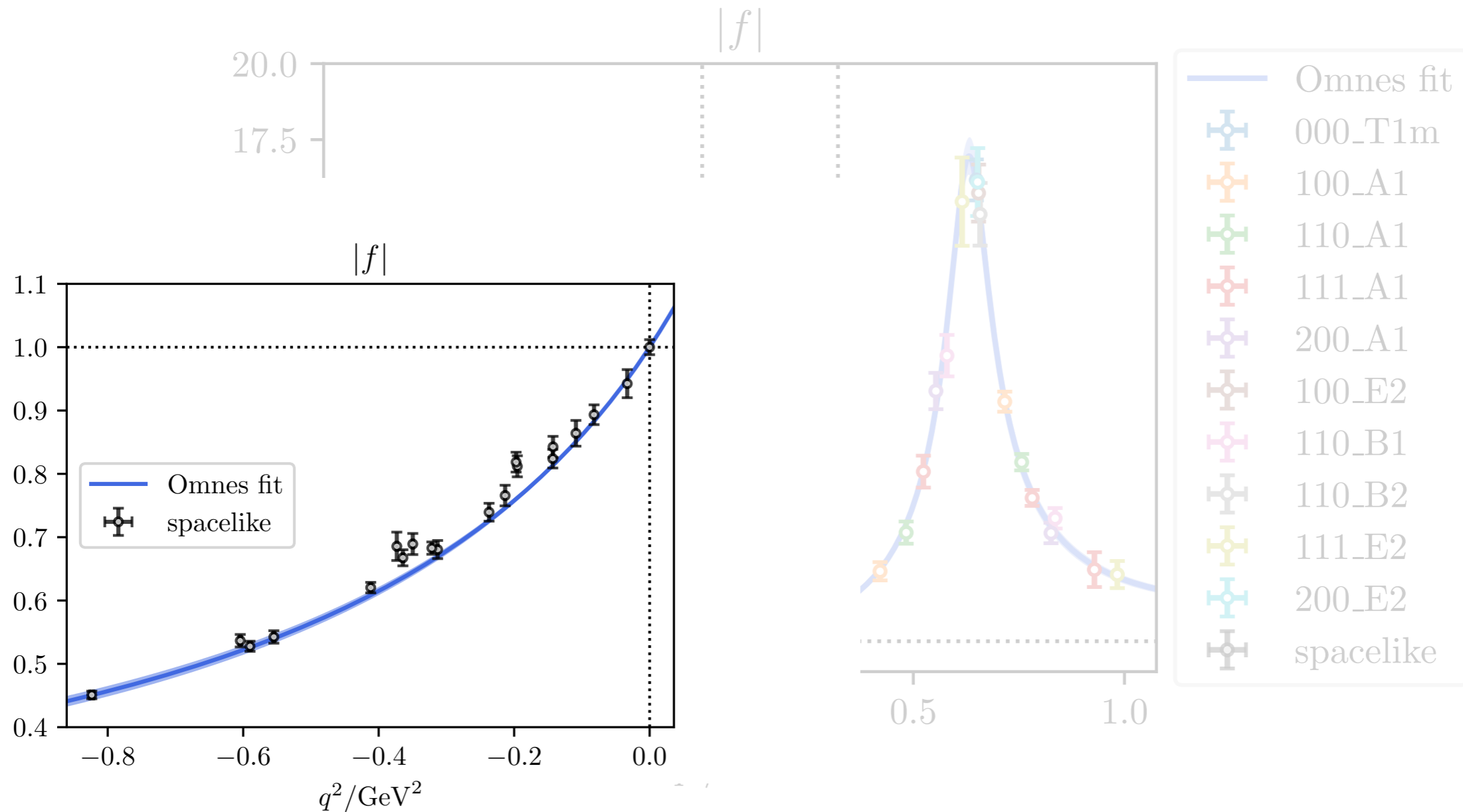
$$\sqrt{\langle r_\pi^2 \rangle} = 0.616(7) \text{ fm}$$

Form factor fit: spacelike and timelike region



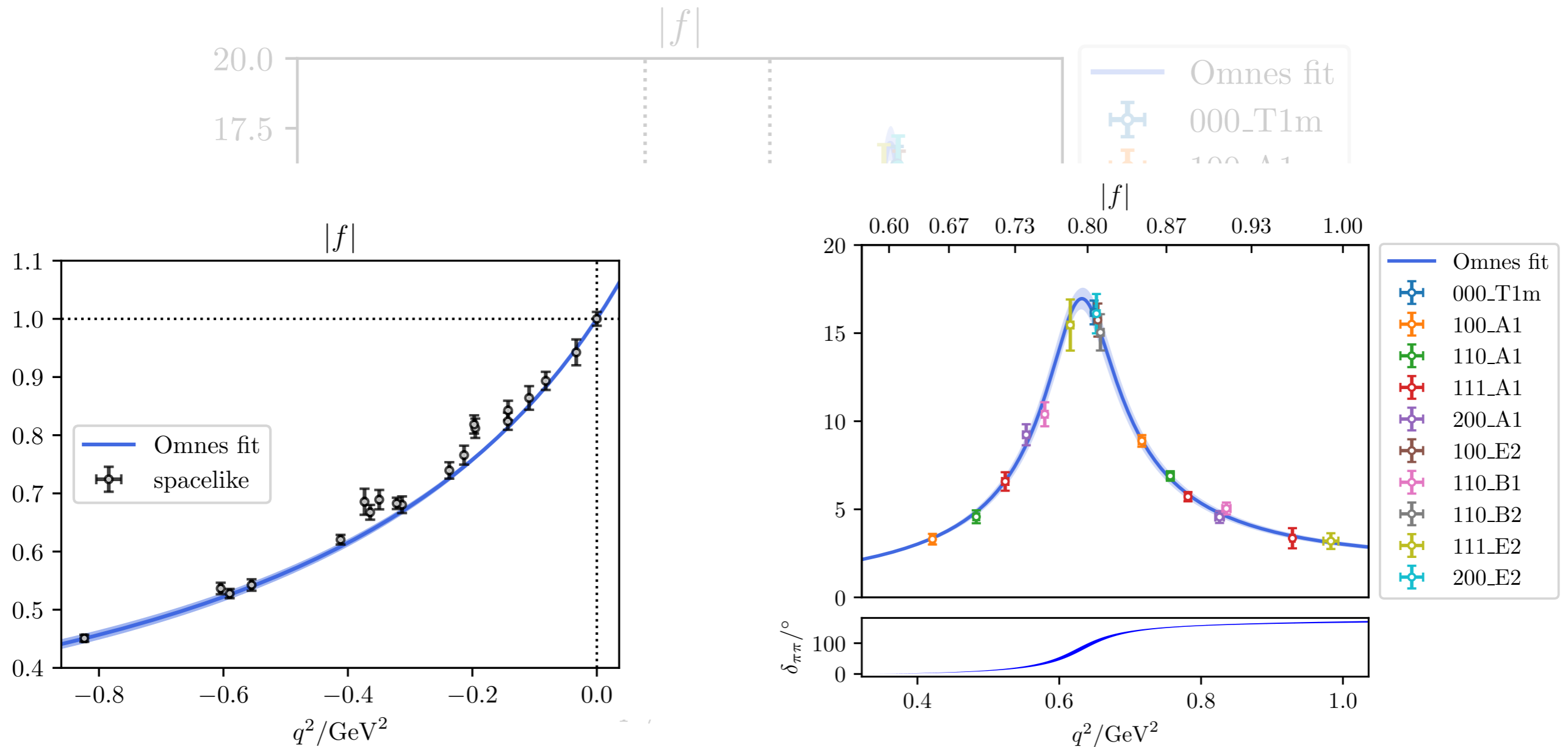
$$\sqrt{\langle r_\pi^2 \rangle} = 0.616(7) \text{ fm}$$

Form factor fit: spacelike and timelike region



$$\sqrt{\langle r_\pi^2 \rangle} = 0.616(7) \text{ fm}$$

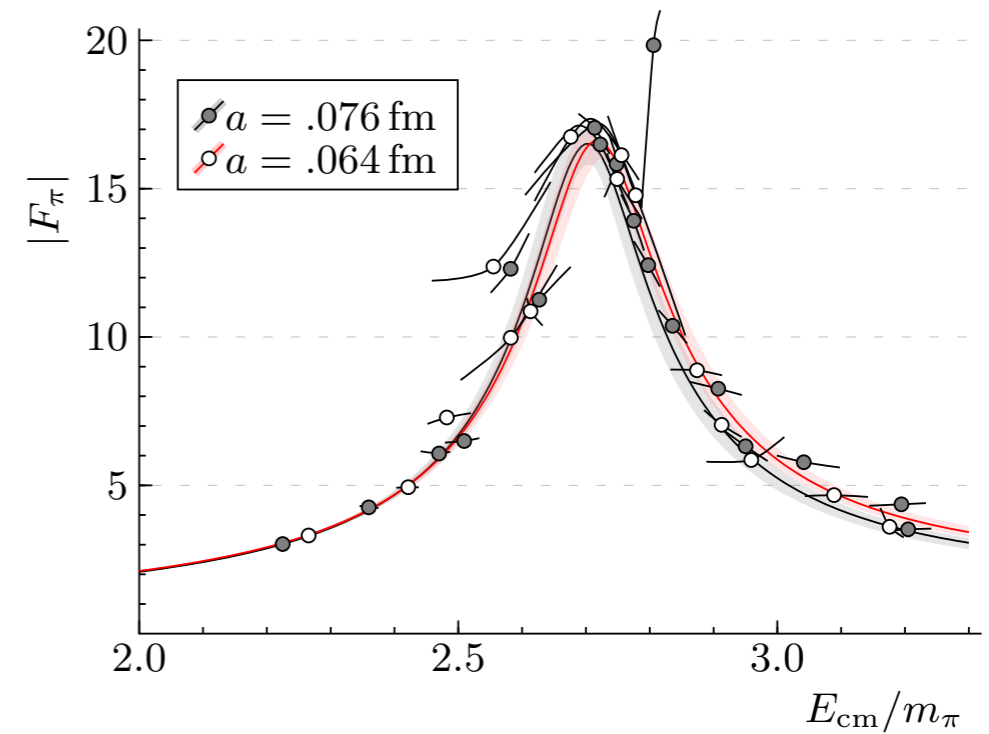
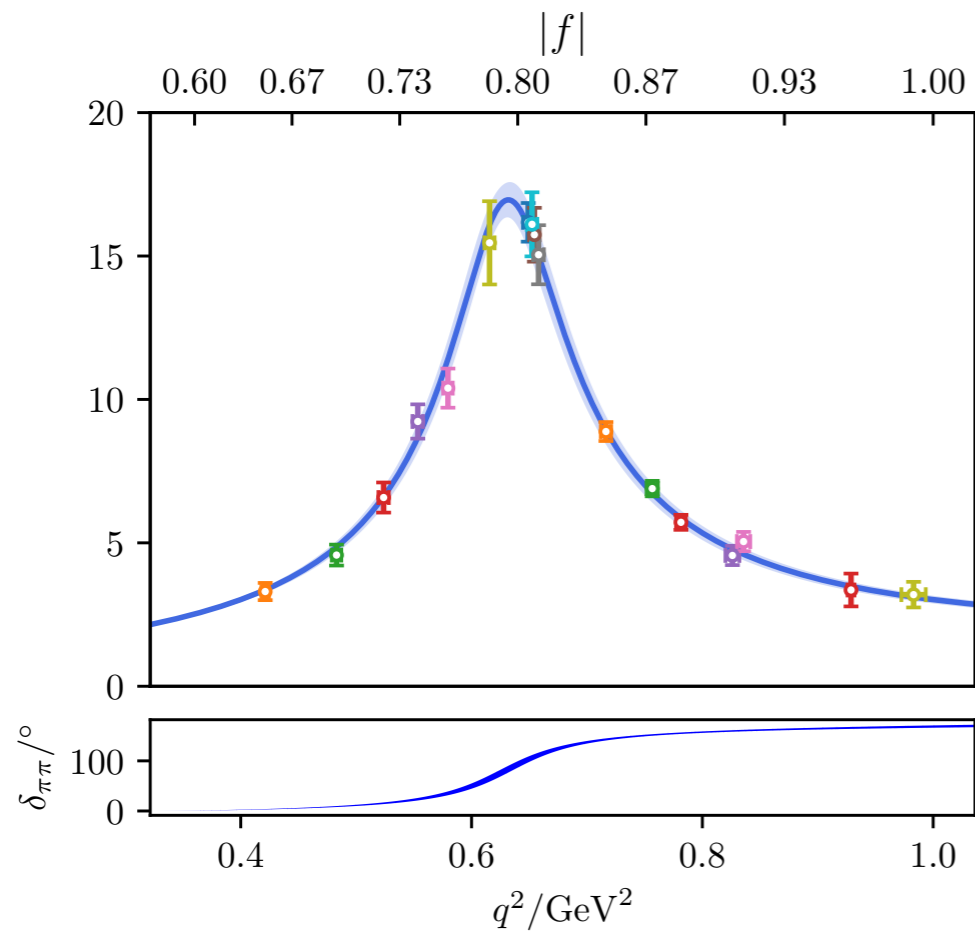
Form factor fit: spacelike and timelike region



$$\sqrt{\langle r_\pi^2 \rangle} = 0.616(7) \text{ fm}$$

Consistency with existing results

$$m_\pi \approx 280 \text{ MeV}$$

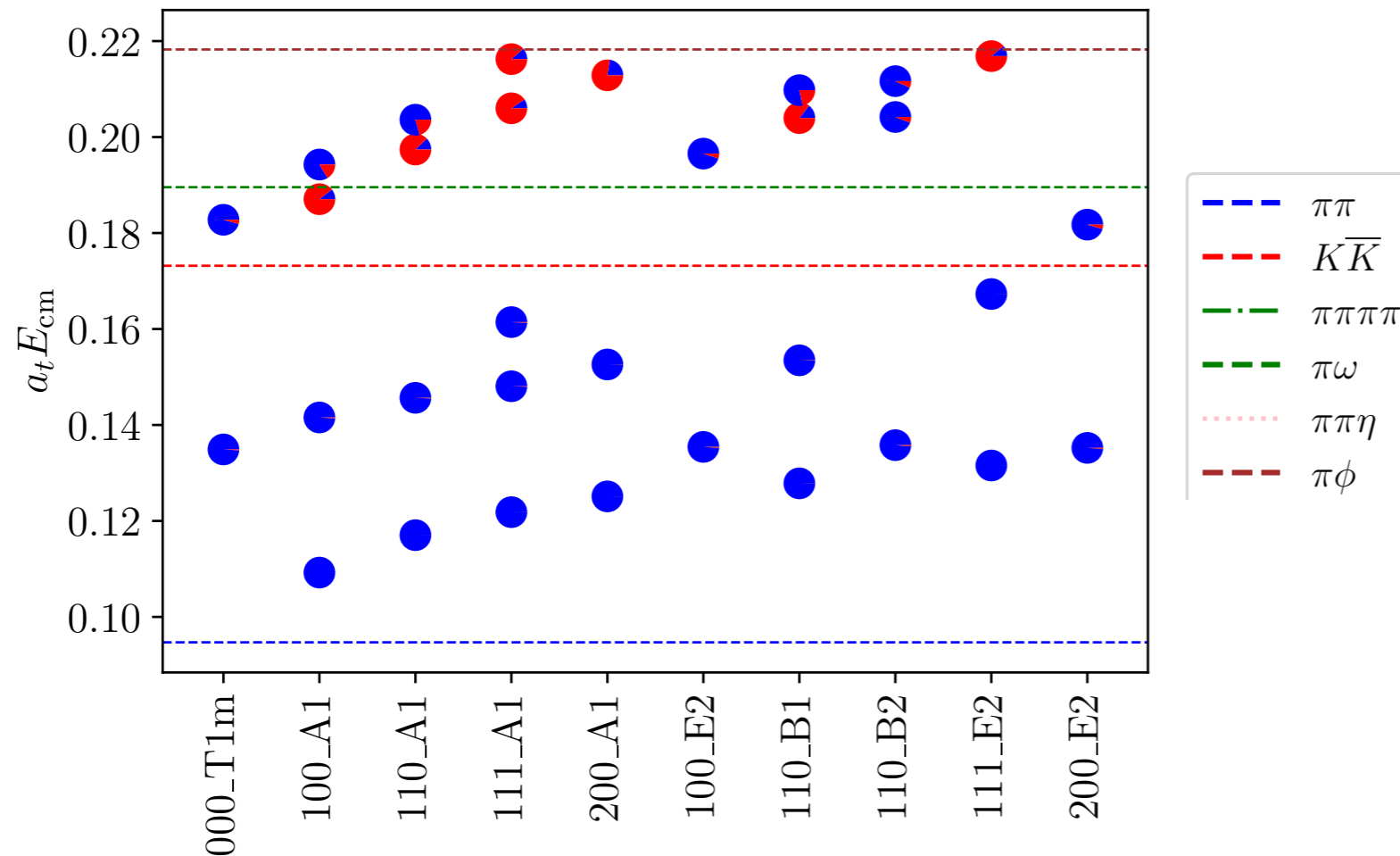


[arXiv:1808.05007] Andersen, et al.

Coupled channel Finite Volume correction

$$\lambda_0'^{\star} w_0 w_0^{\top} = \frac{\partial}{\partial E^{\star}} (\mathcal{M} + F^{-1})$$

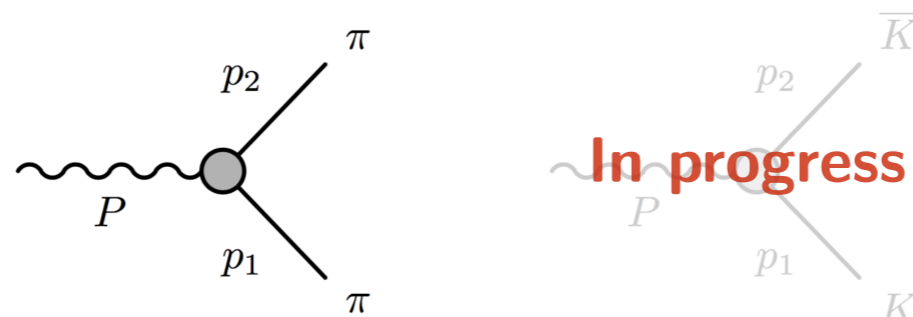
$$\sum_a w_{0,a}^2 = 1$$



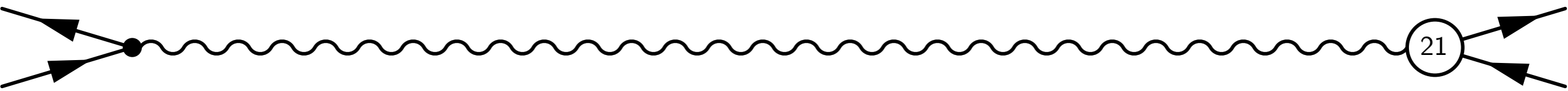
$$\mathcal{F}_L \propto w_{0,a} f_a$$

Summary and outlook

- Isovector p -wave in elastic and coupled channel regions
 - Finite-volume spectrum
 - Scattering amplitude
- Pair production amplitude
 - Zero-to-two finite volume matrix elements
 - Lellouch-Lüscher factor
 - Form factor fit across spacelike and timelike region
- Future work to extend the analysis to the coupled channel region.



Back up slides



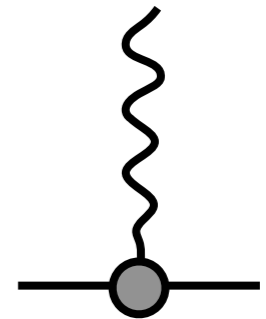
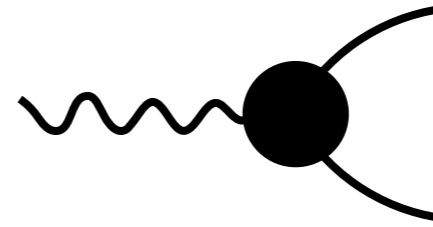
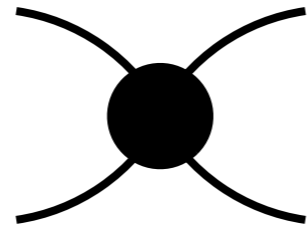
Operator basis

	[000] T_1^-	[100] A_1	[110] A_1	[111] A_1	[200] A_1
	$11 \times \bar{\psi}\mathbf{\Gamma}\psi$	$8 \times \bar{\psi}\mathbf{\Gamma}\psi$	$9 \times \bar{\psi}\mathbf{\Gamma}\psi$	$10 \times \bar{\psi}\mathbf{\Gamma}\psi$	$11 \times \bar{\psi}\mathbf{\Gamma}\psi$
	$\pi_{[100]}\pi_{[100]}$	$\pi_{[100]}\pi_{[000]}$	$\pi_{[110]}\pi_{[000]}$	$\pi_{[111]}\pi_{[000]}$	$\pi_{[200]}\pi_{[000]}$
		$K_{[100]}\bar{K}_{[000]}$	$K_{[110]}\bar{K}_{[000]}$	$\pi_{[110]}\pi_{[100]}$	$K_{[200]}\bar{K}_{[000]}$
		$\pi_{[110]}\pi_{[100]}$	$\pi_{[111]}\pi_{[100]}$	$K_{[111]}\bar{K}_{[000]}$	
				$K_{[110]}\bar{K}_{[100]}$	

	[100] E_2	[110] B_1	[110] B_2	[111] E_2	[200] E_2
	$17 \times \bar{\psi}\mathbf{\Gamma}\psi$	$12 \times \bar{\psi}\mathbf{\Gamma}\psi$	$15 \times \bar{\psi}\mathbf{\Gamma}\psi$	$12 \times \bar{\psi}\mathbf{\Gamma}\psi$	$15 \times \bar{\psi}\mathbf{\Gamma}\psi$
	$\pi_{[110]}\pi_{[100]}$	$\pi_{[100]}\pi_{[100]}$	$\omega_{[110]}\pi_{[000]}$	$\pi_{[110]}\pi_{[100]}$	$\pi_{[110]}\pi_{[110]}$
		$\omega_{[110]}\pi_{[000]}$	$\pi_{[111]}\pi_{[100]}$	$K_{[110]}\bar{K}_{[100]}$	
		$K_{[100]}\bar{K}_{[100]}$	$\pi_{[110]}\pi_{[110]}$		
		$\pi_{[110]}\pi_{[110]}$	$\phi_{[110]}\pi_{[000]}$		
		$\phi_{[110]}\pi_{[000]}$			
		$\omega_{[100]}\pi_{[100]}$			

	$a_t E_{\text{th}}$	E_{th}/MeV
$\pi\pi$	0.0947	567
$K\bar{K}$	0.1732	1037
$\pi\pi\pi\pi$	0.1894	1134
$\pi\omega$	0.1896	1135
$\pi\pi\eta$	0.1908	1142
$\pi\phi$	0.2183	1307
$\pi K\bar{K}$	0.2206	1320
$\pi\eta\omega$	0.2856	1710

MC statistics



For renormalization
of the current

MC configurations

400

348

348

time sources

4

1

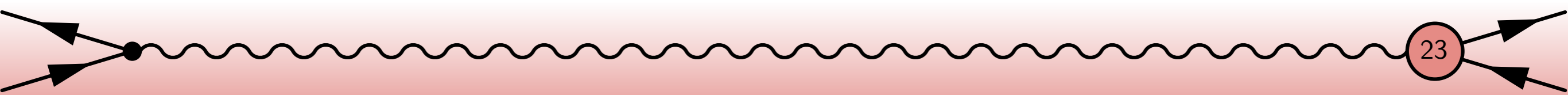
1

Correlation timeslice extent

40

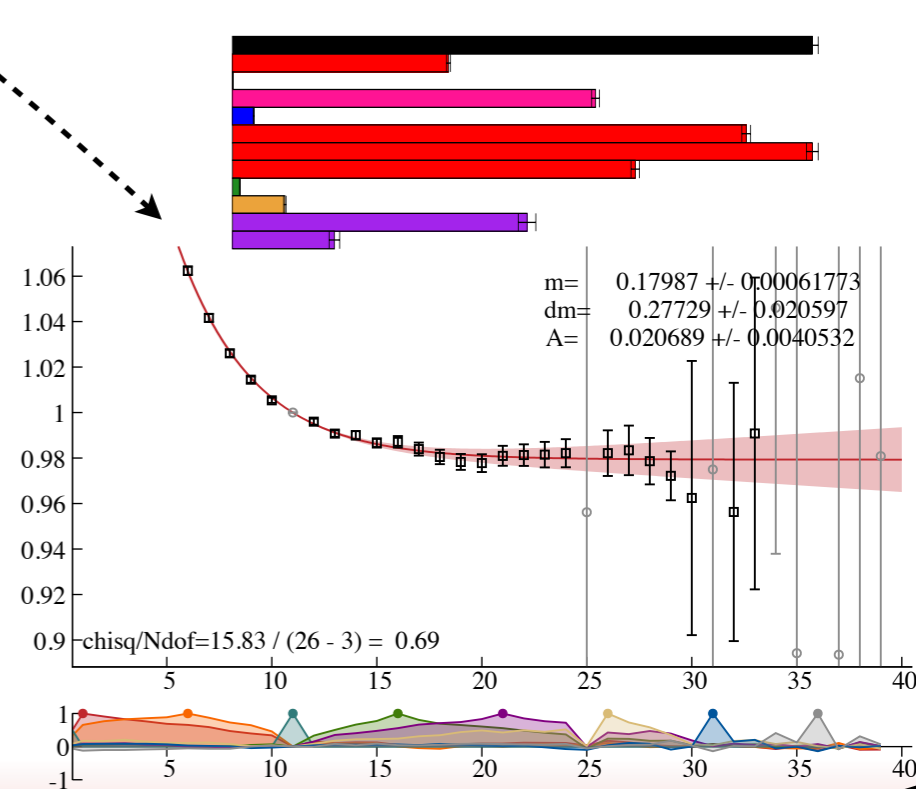
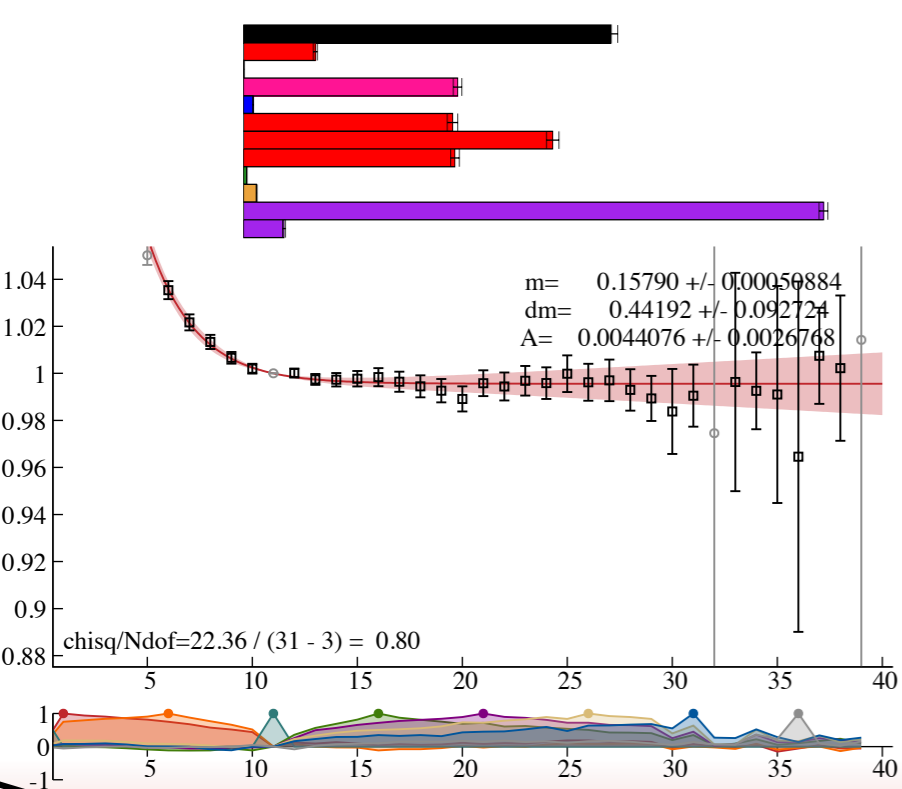
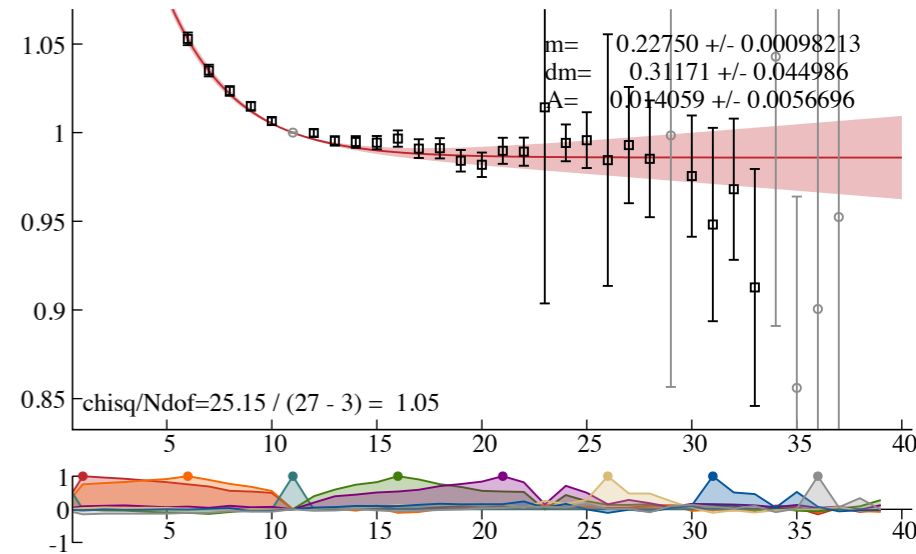
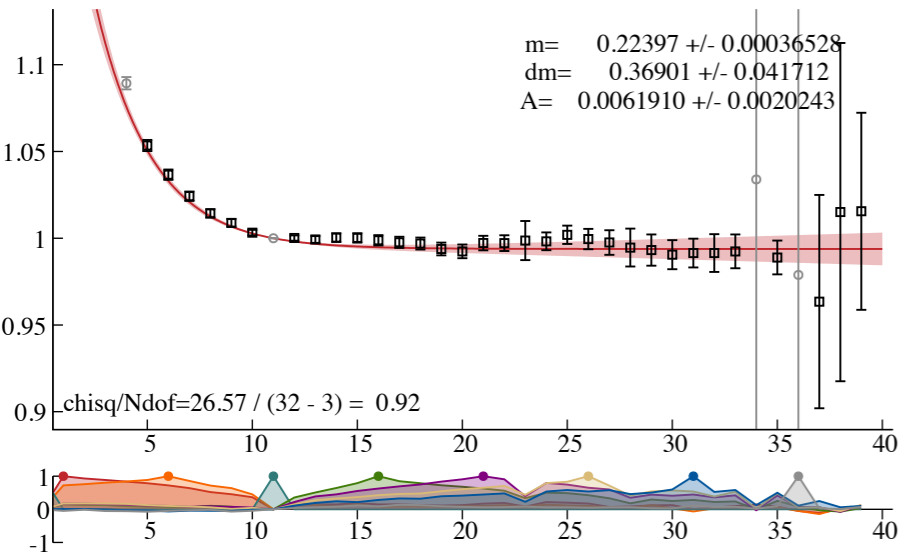
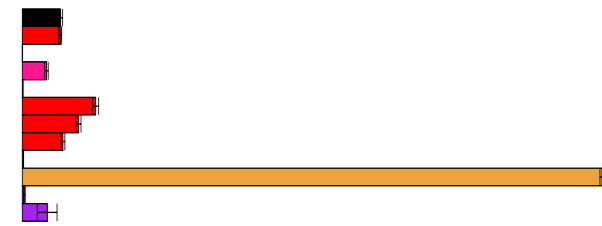
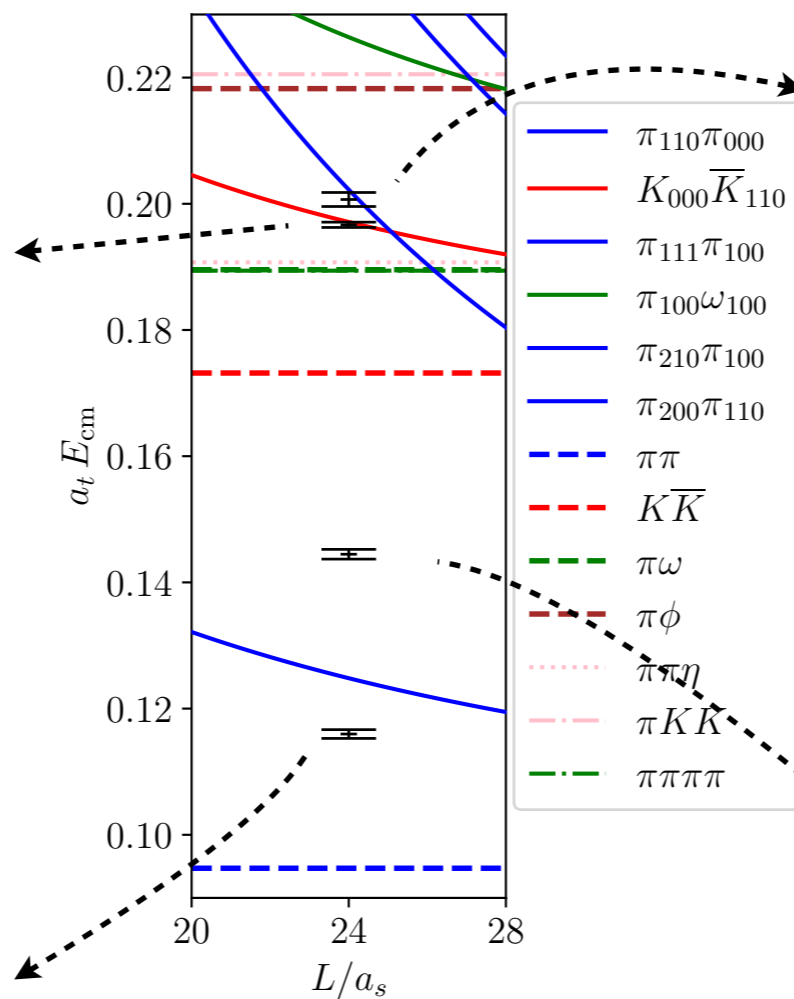
32

32



Typical spectrum

$P = [110], A1$
 $t_0 = 11$



Multimeson $\pi\pi$ $K\bar{K}$

qbarq ρ ρ_2 ρ_3 b_0

hybrids $(\gamma^0 \gamma^5 \times D_{J=1}^{[2]})_{J=1}$

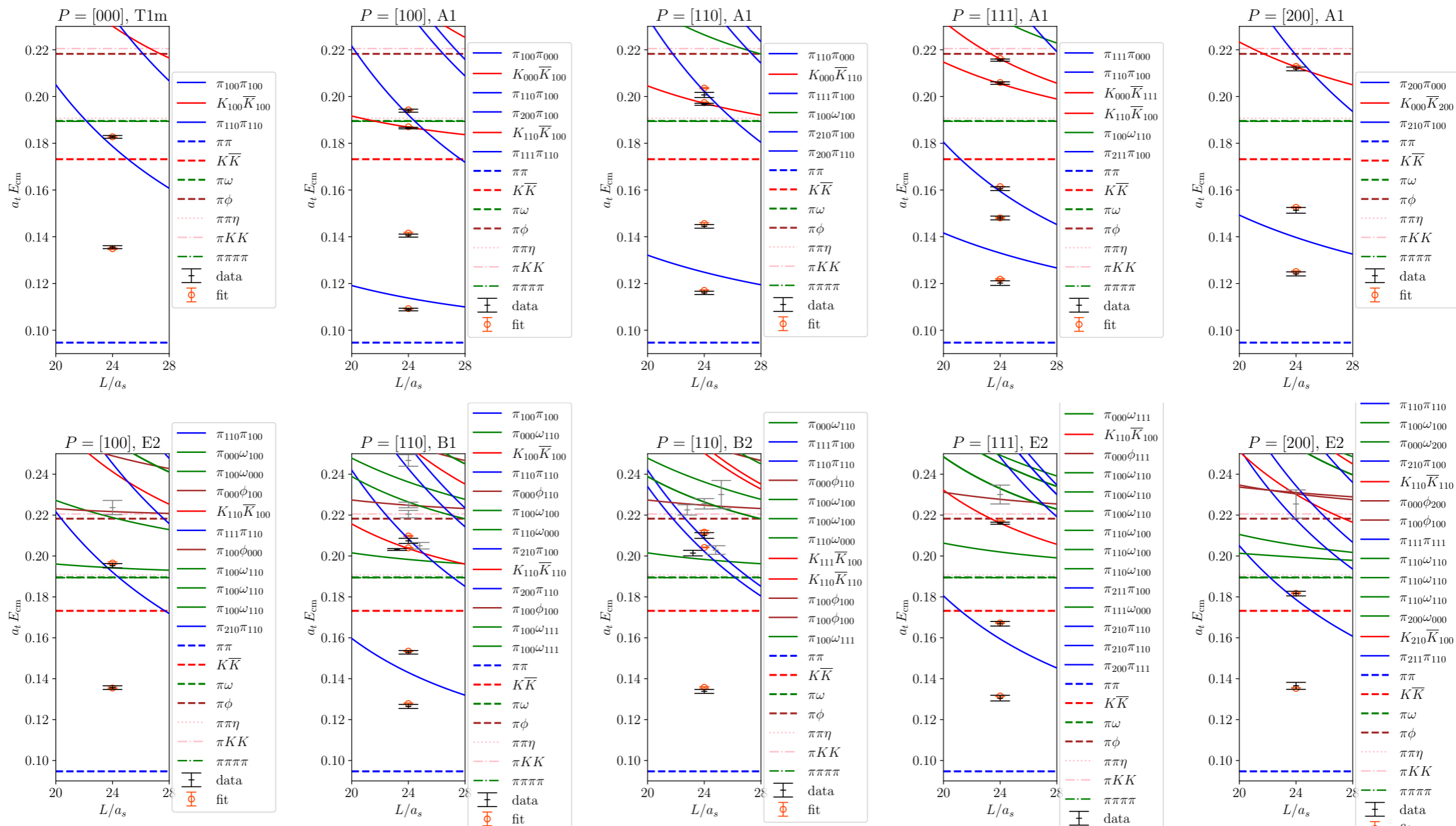
Coupled channels fit:

$\pi\pi$
 $K\bar{K}$

$$\mathcal{M}_{\ell,ab}^{-1} = \frac{1}{(2q_a^*)^\ell} K_{\ell,ab}^{-1} \frac{1}{(2q_b^*)^\ell} - i\rho_{\text{CM},ab}$$

$$K_{ab} = \frac{g_a g_b}{-s + m_r^2} + \gamma_{ab}$$

$$\chi^2/\text{dof} = 28.70/(32 - 6) = 1.10$$

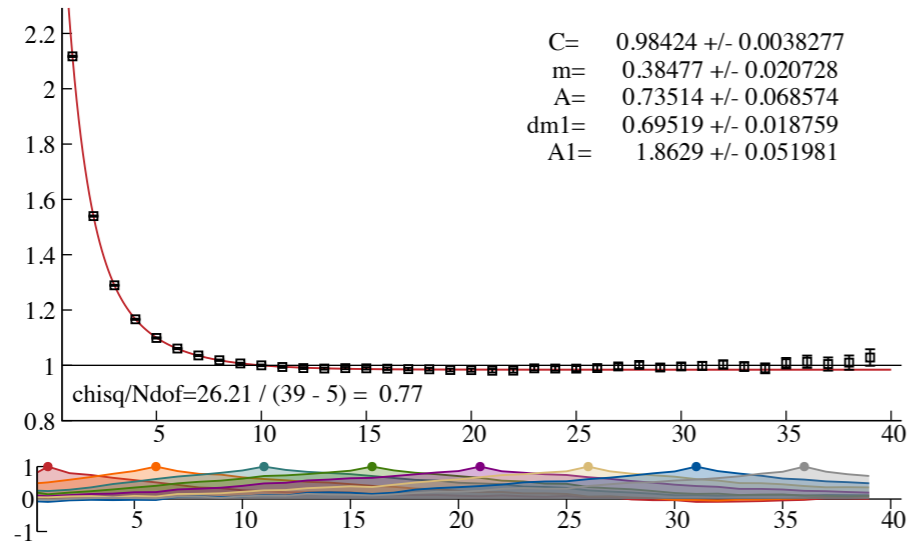


Optimized operators

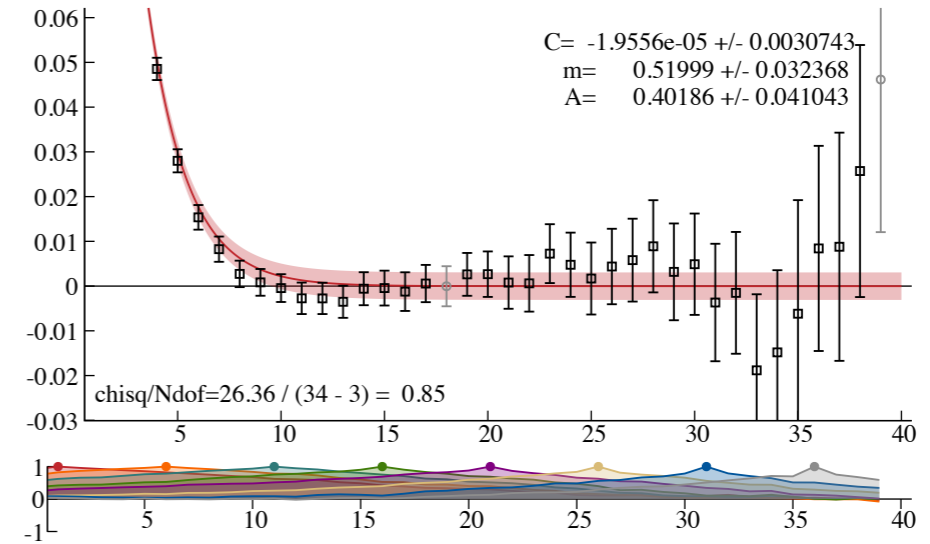
[000] T1m irrep

$$e^{\sqrt{E_n E_m} t} \langle 0 | \Omega_n(t) \Omega_m^\dagger(0) | 0 \rangle = \delta_{n,m} + \mathcal{O}(e^{-(E_N - \sqrt{E_n E_m}) t})$$

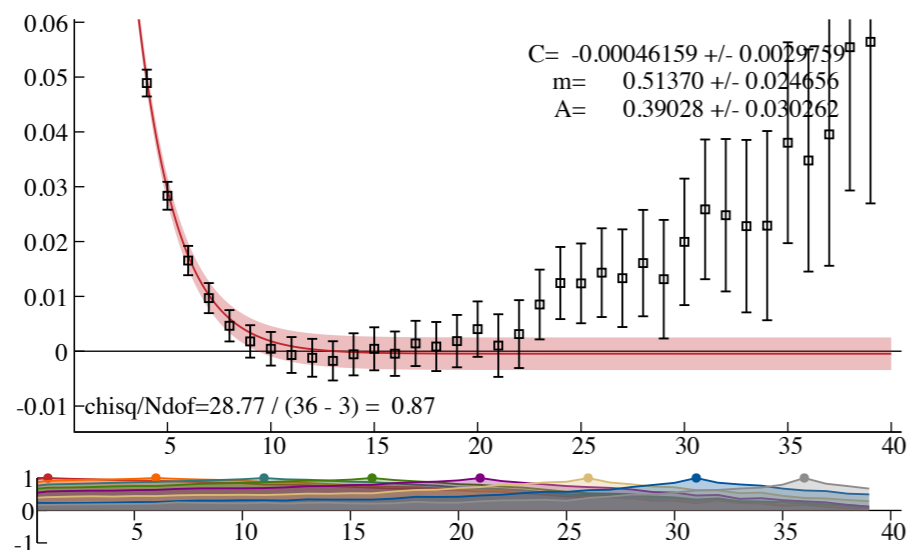
C00



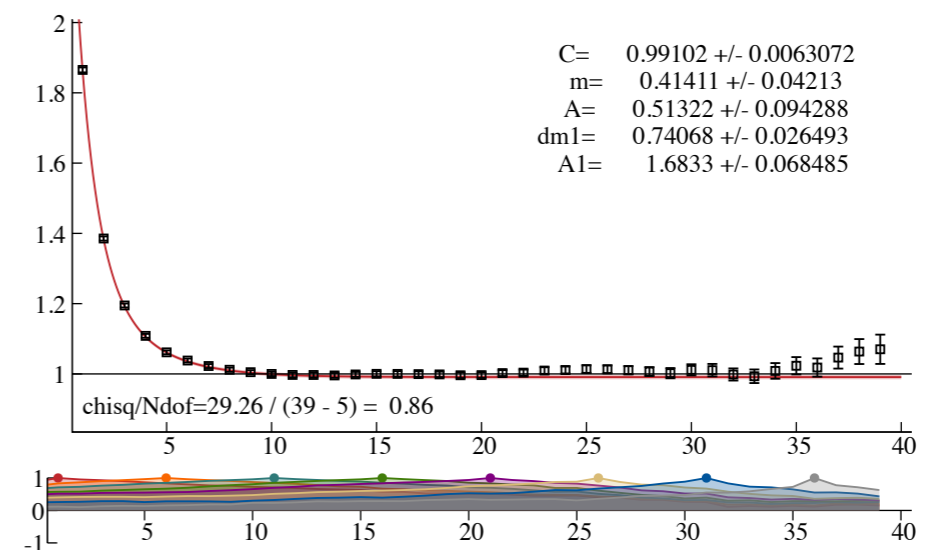
C01



C10



C11



Typical three point fit

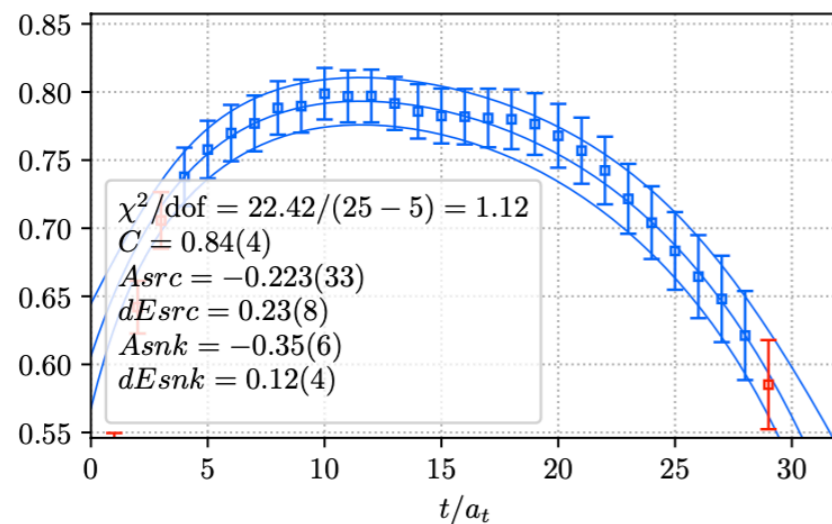
$$\langle \pi_{[1-10]} | \mathcal{J}_{\text{impro}}^\rho | \pi_{[200]} \rangle$$

$$\frac{\langle \pi(\tau, \vec{p}_f) \mathcal{J}(t) \pi^\dagger(0, \vec{p}_i) \rangle_{\text{rel}}}{\langle \pi(\tau - t, \vec{p}_f) \pi^\dagger(0, \vec{p}_f) \rangle \langle \pi(t, \vec{p}_i) \pi^\dagger(0, \vec{p}_i) \rangle} = Z_f^{-1/2} Z_i^{-1/2} \langle \pi(\vec{p}_f) | \mathcal{J}(0) | \pi(\vec{p}_i) \rangle + \dots$$

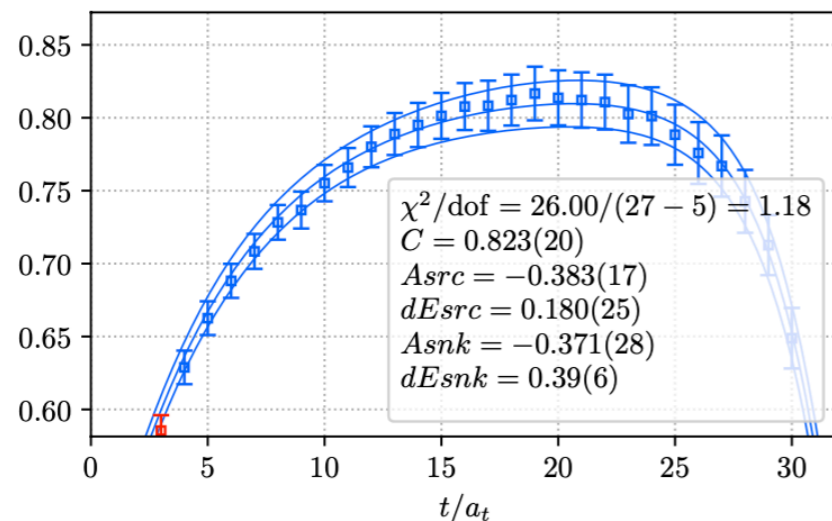
$$C + A_{\text{src}} e^{-dE_{\text{src}} t} + A_{\text{snk}} e^{-dE_{\text{snk}}(\tau - t)}$$

Current
irrep

A1



B1



Kinematic factor

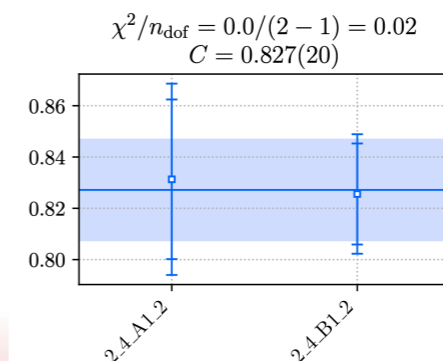
$$\langle \pi(p_1) | \mathcal{J}^i(\vec{q}) | \pi(p_2) \rangle = (p_1 + p_2)^i f(Q^2),$$

$$\mathcal{S}_m^{\Lambda, \mu} \varepsilon_i^m(\vec{q}) \langle \pi(p_1) | \mathcal{J}^i(\vec{q}) | \pi(p) \rangle = \langle \pi(p_1) | \mathcal{J}^{\Lambda, \mu}(\vec{q}) | \pi(p_2) \rangle,$$

Improvement

$$\langle \mathbf{m} | \mathcal{J}_{\text{impro}}^\rho | \mathbf{n} \rangle = \langle \mathbf{m} | \rho | \mathbf{n} \rangle + \frac{1}{4} (1 - \xi) a_t (E_m - E_n) \langle \mathbf{m} | \rho_2 | \mathbf{n} \rangle$$

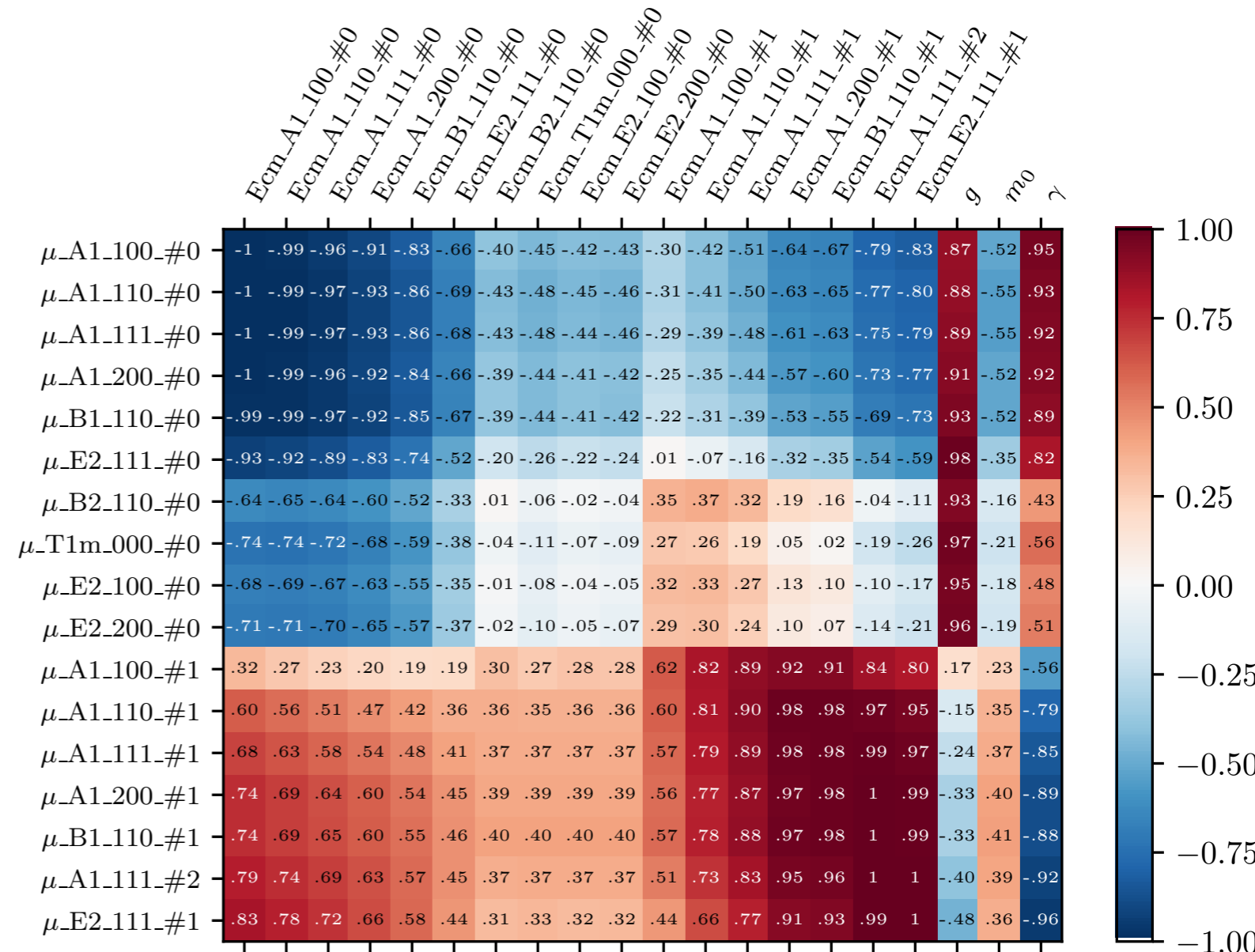
Average over irreps



B2 kinematic factor is zero.

Correlation between slope and energies

- Slopes away from the resonance are highly correlated to the Lüscher energy.
- The slopes close to the resonance have high correlation to g .



Correlation matrix between slopes, Lüscher energies and scattering parameters