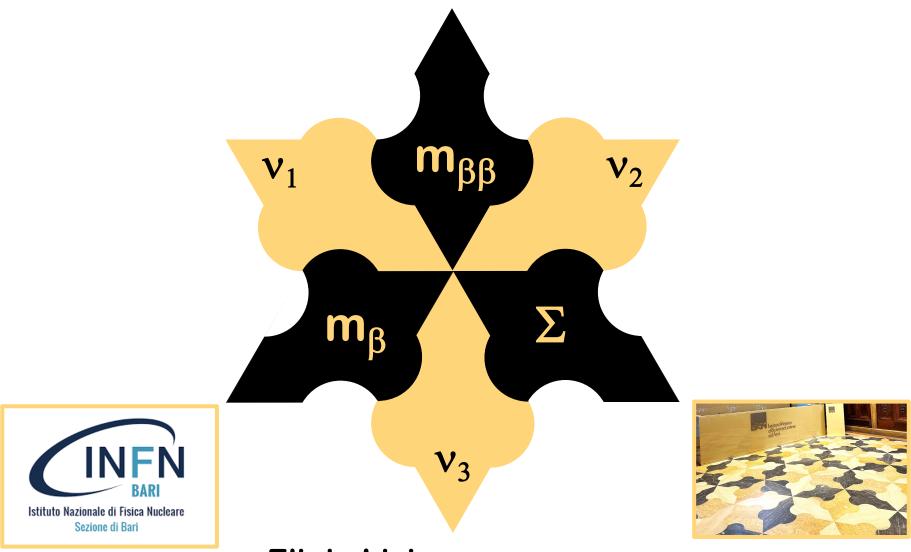
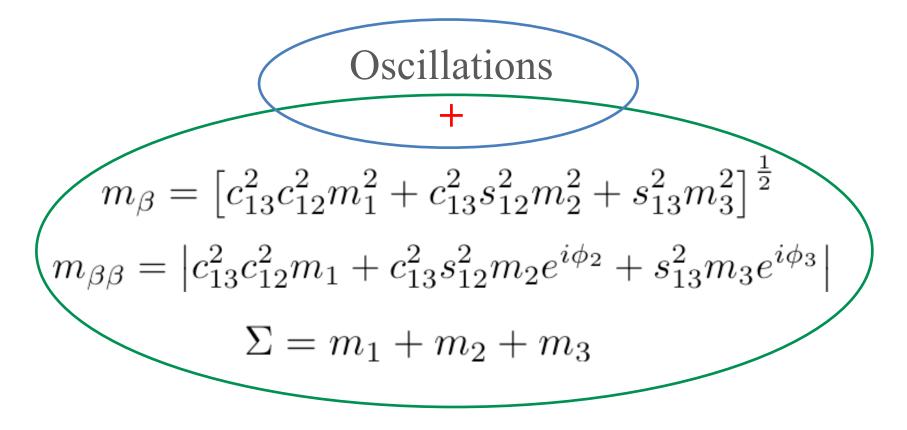
# Towards a global analysis of absolute $\nu$ masses



Eligio Lisi (INFN, Bari, Italy)

XX Int. Workshop on Neutrino Telescopes – Palazzo Franchetti, Venezia, 27 Oct 2023

In perspective, the global analysis of absolute 3v masses from:

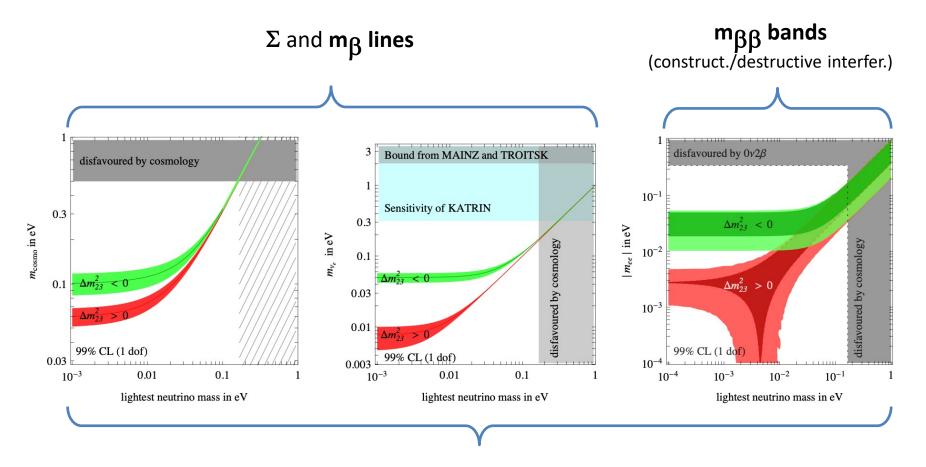


involves several issues that are worth discussing, in the light of (far) future  $m_{\beta}$ ,  $m_{\beta\beta}$ ,  $\Sigma$  signals ...and of possible new physics

# **Outline:**

Graphs of 3v osc. bounds Towards a  $\Sigma$  signal Towards a  $m_{\beta}$  signal Towards  $m_{\beta\beta}$  (& beyond 3v) Epilogue

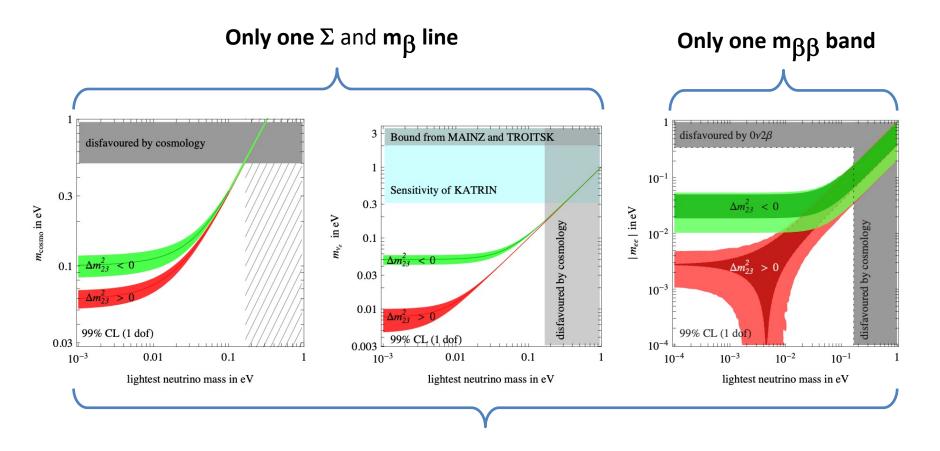
# Graphs of 3v osc. bounds (1): $(m_{\beta}, m_{\beta\beta}, \Sigma)$ vs $m_{\text{lightest}}$ in NO/IO



Lines and bands somewhat smeared by oscillation parameter uncertainties

*Figure from Strumia & Vissani, hep-ph/0606054* 

### Precise oscillometry in next decade → Negligible smearing & NO/IO selection



ightarrow Progress in these planes will be driven only by absolute mass observables

(within the standard *3v* framework)

# **Comment on mass ordering through oscillations**

No bump/dip/kink... but small/smooth differences in spectral templates S(E) → requires statistical comparison of template shapes vs data

Probes:	Templates:	Oscill. physics:
(1) MBL reactors:	S(E)	$\pm \Delta m^2$ vs δm <sup>2</sup>
(2) LBL acceler.:	S(E, flavor)	$\pm \Delta m^2$ vs MSW, δm <sup>2</sup>
(3) Atmospheric:	S(E, flavor, zenith)	$\pm \Delta m^2$ vs MSW, δm <sup>2</sup>

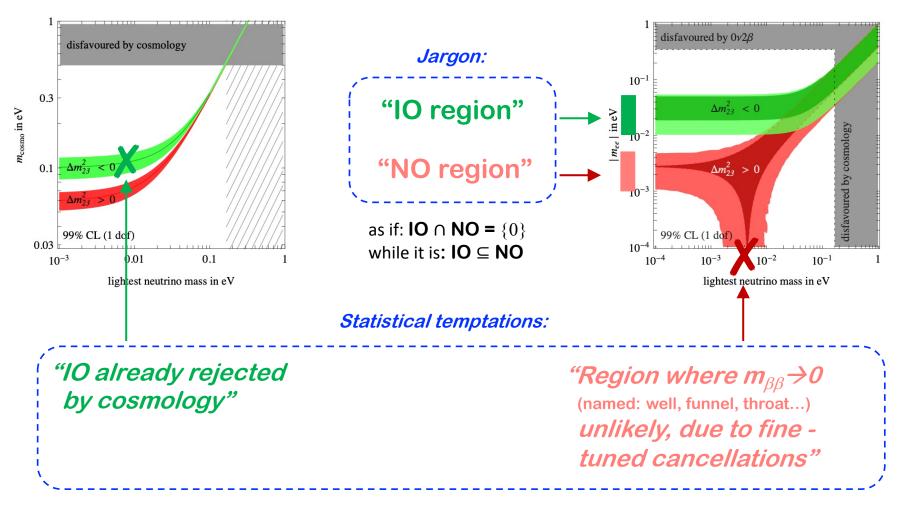
In addition, "synergy" or "complementarity" of different probes:

(4)  $\geq$  2 probes: Spread of  $\{+\Delta m_i^2\}$  vs  $\{-\Delta m_i^2\}$  smaller for true ordering

Currently: Some hints from (2-4), sum up to  $\sim 2.5\sigma$  in favor of  $+\Delta m^2$ Future: from hints to discovery, as lines of evidence (1 - 4) grow & converge

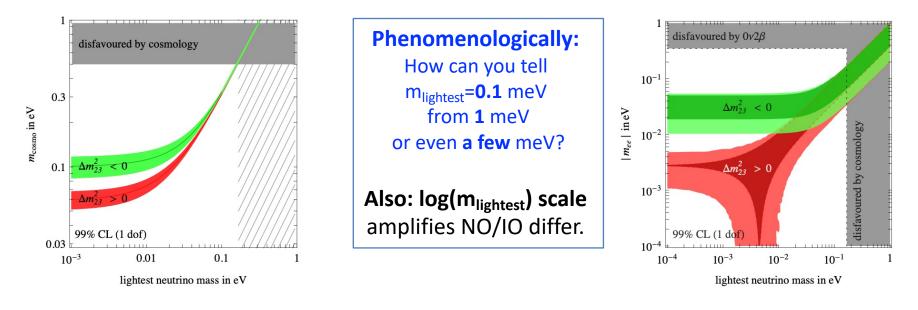
See also talk by S. Parke

### In the meantime... avoid "jargon" and "statistical temptations" ...e.g.:

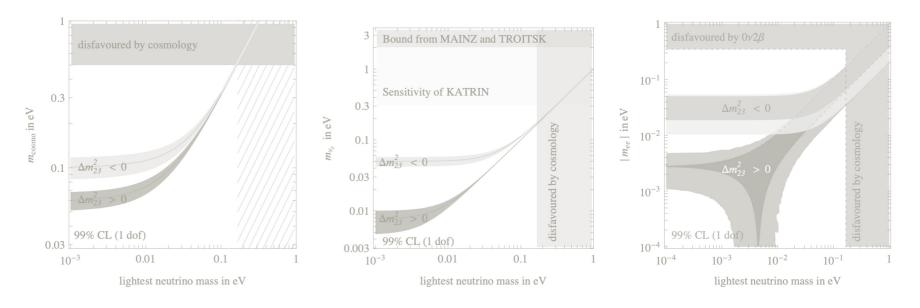


Nature does not care about our "naturalness" criteria or phase-space (under)sampling!

### ... also notice that m<sub>lightest</sub> is not really measured



It makes sense to project away  $m_{lightest} \rightarrow$ 



It makes sense to project away  $m_{lightest} \rightarrow$ 

**But** ... keep in mind that the case  $m_{lightest} = 0$  guarantees futuristic implications:

- a relativistic CvB component
- a 0vββ lower bound in NO
- a v component with v=c

up to redshift z=0 ( $m_{lightest} < T_v \sim 0.1 \text{ meV suffices}$ )  $m_{\beta\beta} > 1 \text{ meV} (m_{lightest} < 1 \text{ meV suffices})$ from multimessenger astrophysical sources

# Graphs of 3v osc. bounds (2): $(m_{\beta}, m_{\beta\beta}, \Sigma)$ without $m_{\text{lightest}}$

Only measurable quantities; graphically amplified structures are squeezed away

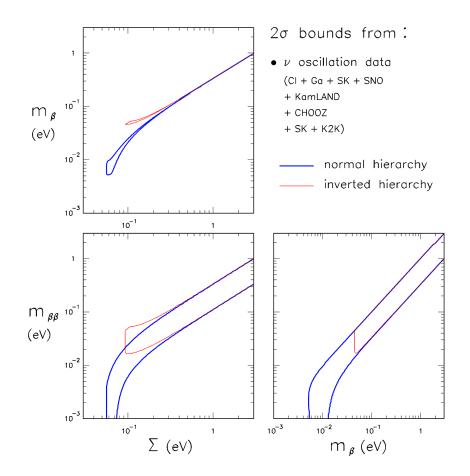
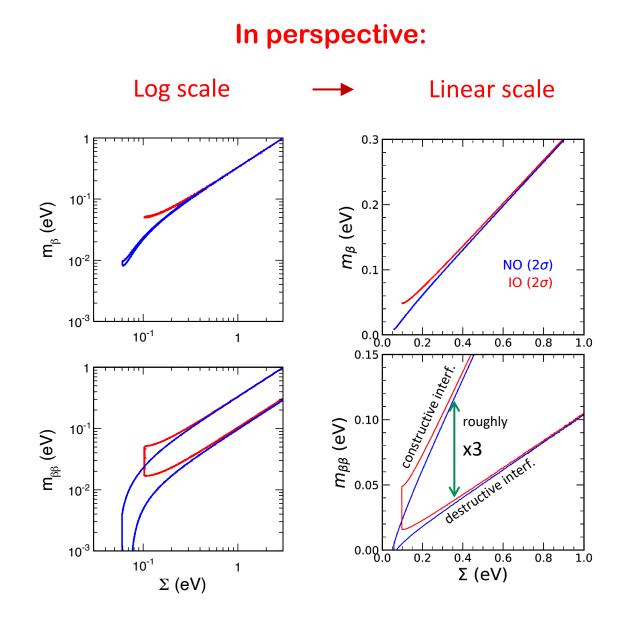
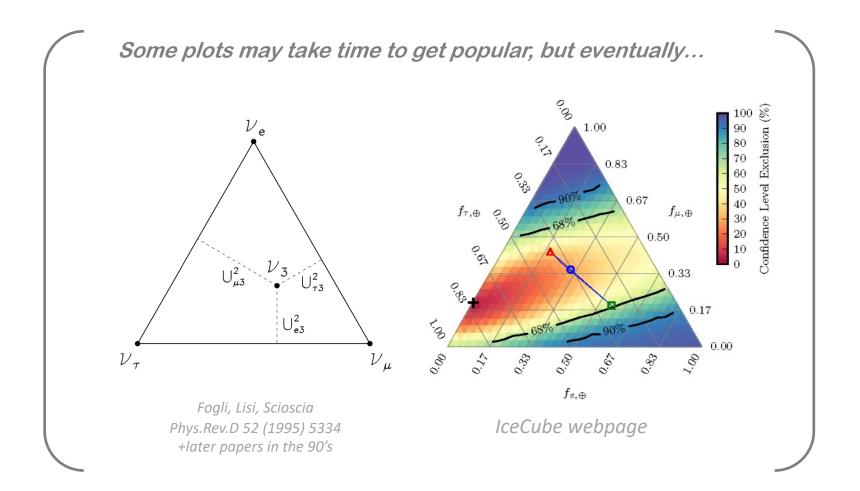
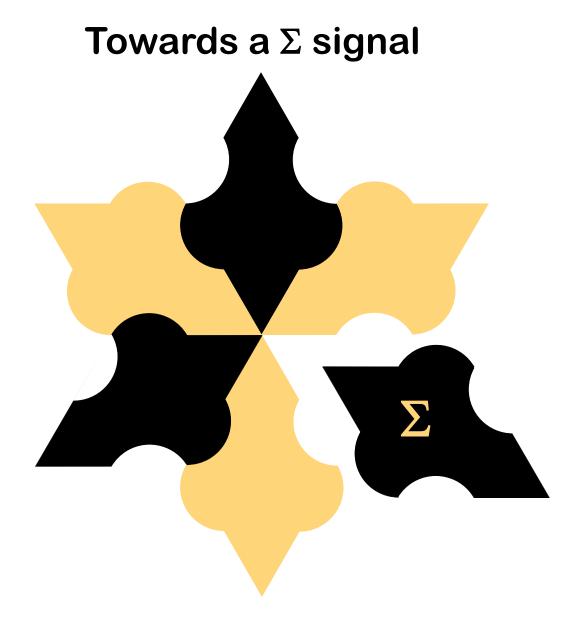


Figure from Fogli, Lisi, Marrone, Melchiorri, Palazzo, Serra and Silk, hep-ph/0408045





#### 



*Current bounds: freely adapted from PDG quoted values and from 2107.00532 Forecasts: mainly adapted from M. Lattanzi, talks at NOW 2022 and TAUP 2023 See also talks by O. Mena and R. Laureijs at this Workshop*   $\Sigma$  signal is guaranteed: min  $\Sigma \simeq \begin{cases} 60 \text{ meV} (\text{NO}) \\ 100 \text{ meV} (\text{IO}) \end{cases}$ 

**But:**  $\Sigma$  = **output** of a multi-parameter **fit** to cosmological data within  $\Lambda$ CDM,

## $\mathcal{L}(\Sigma) \rightarrow \text{(roughly)} \Sigma \simeq \Sigma_0 \pm \sigma$

**Currently: Variants** in #parameters, datasets, model...  $\rightarrow$  various 95% CL limits:

 $\Sigma_i < \Sigma_{0i} + 2\sigma_i$ ,  $i \in \{variants\}$ 

	-					
Cos	osmological inputs for nonoscillation data analysis Results: Cosmo only		$Cosmo + m_{\beta} + m_{\beta\beta}$			
#	Model	Data set	$\Sigma$ (2 $\sigma$ )	$\Delta\chi^2_{\rm IO-NO}$	$\Sigma$ (2 $\sigma$ )	$\Delta\chi^2_{ m IO-NO}$
0	$\Lambda {\rm CDM} + \Sigma$	Planck TT, TE, EE	$< 0.34~{\rm eV}$	0.9	$< 0.32 \ \mathrm{eV}$	1.0
1	$\Lambda \mathrm{CDM} + \Sigma$	Planck TT, TE, EE + lensing	$< 0.30~{\rm eV}$	0.8	$< 0.28~{\rm eV}$	0.9
<b>2</b>	$\Lambda {\rm CDM} + \Sigma$	Planck TT, TE, EE + BAO	$< 0.17~{\rm eV}$	1.6	$<0.17~{\rm eV}$	1.8
3	$\Lambda {\rm CDM} + \Sigma$	Planck TT, TE, EE + BAO + lensing	$< 0.15~{\rm eV}$	2.0	$<0.15~{\rm eV}$	2.2
4	$\Lambda \mathrm{CDM} + \Sigma$	Planck TT, TE, EE + lensing + $H_0(R19)$	$< 0.13 \ {\rm eV}$	3.9	$<0.13~{\rm eV}$	4.0
5	$\Lambda \mathrm{CDM} + \Sigma$	Planck TT, TE, EE + BAO + $H_0$ (R19)	$< 0.13~{\rm eV}$	3.1	$<0.13~{\rm eV}$	3.2
6	$\Lambda {\rm CDM} + \Sigma$	Planck TT, TE, EE + BAO + lensing + $H_0$ (R19)	$< 0.12~{\rm eV}$	3.7	$<0.12~{\rm eV}$	3.8
7	$\Lambda {\rm CDM} + \Sigma + A_{\rm lens}$	Planck TT, TE, EE + lensing	$< 0.77~{\rm eV}$	0.1	$< 0.66~{\rm eV}$	0.1
8	$\Lambda {\rm CDM} + \Sigma + A_{\rm lens}$	Planck TT, TE, EE + BAO	$< 0.31~{\rm eV}$	0.2	$< 0.30~{\rm eV}$	0.3
9	$\Lambda {\rm CDM} + \Sigma + A_{\rm lens}$	Planck TT, TE, EE + BAO + lensing	$< 0.31~{\rm eV}$	0.1	$< 0.30~{\rm eV}$	0.2
10	$\Lambda {\rm CDM} + \Sigma$	$ACT + WMAP + \tau_{prior}$	$< 1.21~{\rm eV}$	-0.1	$< 1.00~{\rm eV}$	0.1
11	$\Lambda {\rm CDM} + \Sigma$	ACT + WMAP + Planck lowE	$< 1.12~{\rm eV}$	-0.1	$< 0.87~{\rm eV}$	0.1
12	$\Lambda \mathrm{CDM} + \Sigma$	ACT + WMAP + Planck lowE + lensing	$< 0.96~{\rm eV}$	0.0	$< 0.85~{\rm eV}$	0.1

	Model	95% CL (eV)	Ref.				
CMB alone							
Pl18[TT+lowE]	$\Lambda \text{CDM} + \sum m_{\nu}$	< 0.54	[22]				
Pl18[TT,TE,EE+lowE]	$\Lambda CDM + \sum m_{\nu}$	< 0.26	[22]				
CMB + probes of background evolution							
Pl18[TT+lowE] + BAO	$\Lambda \text{CDM} + \sum m_{\nu}$	< 0.13	[43]				
Pl18[TT,TE,EE+lowE]+BAO	$\Lambda CDM + \sum m_{\nu} + 5$ params.	< 0.515	[23]				
$\overline{\text{CMB} + \text{LSS}}$							
Pl18[TT+lowE+lensing]	$\Lambda CDM + \sum m_{\nu}$	< 0.44	[22]				
Pl18[TT,TE,EE+lowE+lensing]	$\Lambda CDM + \sum m_{\nu}$	< 0.24	[22]				
CMB + probes of background evolution + LSS							
P118[TT, TE, EE + lowE] + BAO + RSD	$\Lambda \text{CDM} + \sum m_{\nu}$	< 0.10	[43]				
$Pl18[TT+lowE+lensing] + BAO + Lyman-\alpha$	$\Lambda \text{CDM} + \sum m_{\nu}$	< 0.087	[44]				
Pl18[TT,TE,EE+lowE] + BAO + RSD + Pantheon	+ DES $\Lambda CDM + \sum m_{\nu}$	< 0.13	[45]				

*E.g., Capozzi+ 2107.00532* 

Lesgourgues, Verde PDG 2022

+talk by Olga Mena

# $\Sigma$ signal is guaranteed: min $\Sigma \simeq \begin{cases} 60 \text{ meV} (\text{NO}) \\ 100 \text{ meV} (\text{IO}) \end{cases}$

**But:**  $\Sigma =$ **output** of a multi-parameter **fit** to cosmological data within  $\Lambda$ CDM,

$$\mathcal{L}(\Sigma) \rightarrow$$
 (roughly)  $\Sigma \simeq \Sigma_0 \pm \sigma$ 

Currently: Variants in #parameters, datasets, model... → various 95% CL limits:

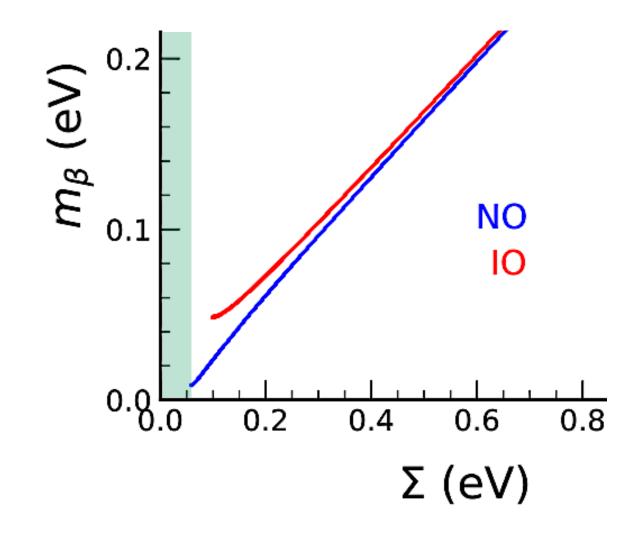
$$\Sigma_i < \Sigma_{0i} + 2\sigma_i$$
,  $i \in \{variants\}$ 

**Strongest** current limits [PDG,  $\Sigma_i < 90-130$  meV at  $2\sigma$ ] roughly correspond to:

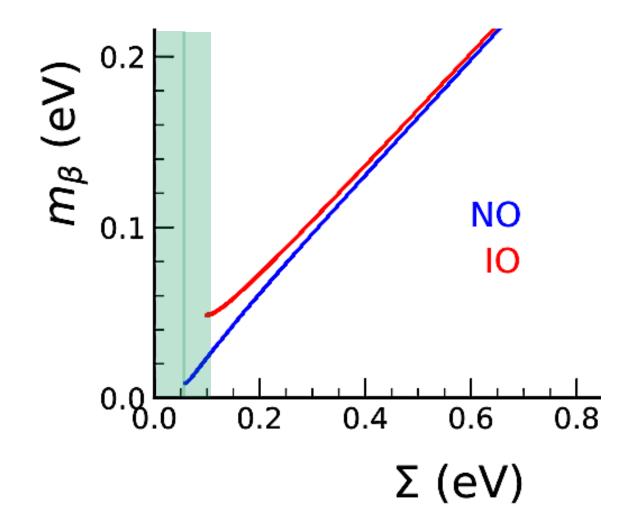
$$\Sigma_{0i} \sim 0 \text{ meV}$$
  
 $\sigma_i \sim 45-65 \text{ meV}$ 

Weaker limits involve larger uncertainties  $\sigma_i$  and/or nonzero best fits  $\Sigma_{0i} \sim O(\sigma_i)$ 

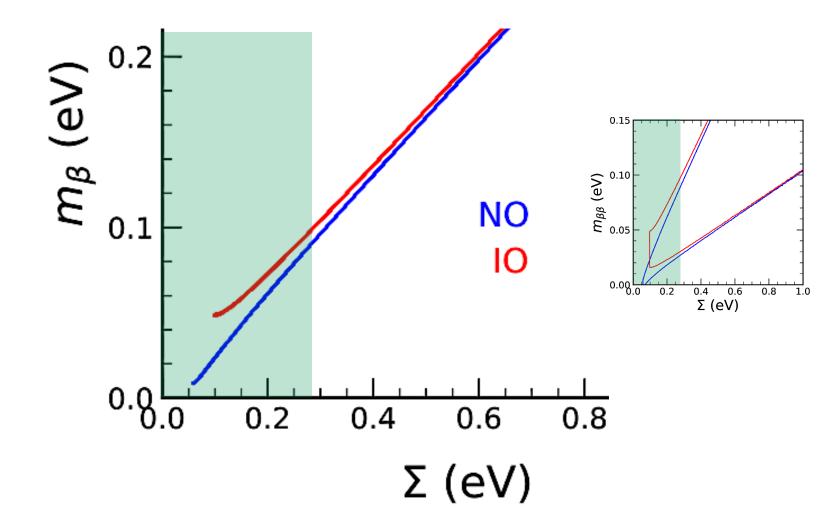
Implications of a current strong limit, e.g.,  $\Sigma = 0 \pm 60$  meV:



Unphysical best fit, but ... compatible with min(NO) at  $\sim 1\sigma$  and min(IO) at  $< 2\sigma$ To some extent, best fit may be an artifact of degenerate mass approximation  $\rightarrow$  For nondegenerate v masses get, e.g.,  $\Sigma = 60 \pm 60$  meV:



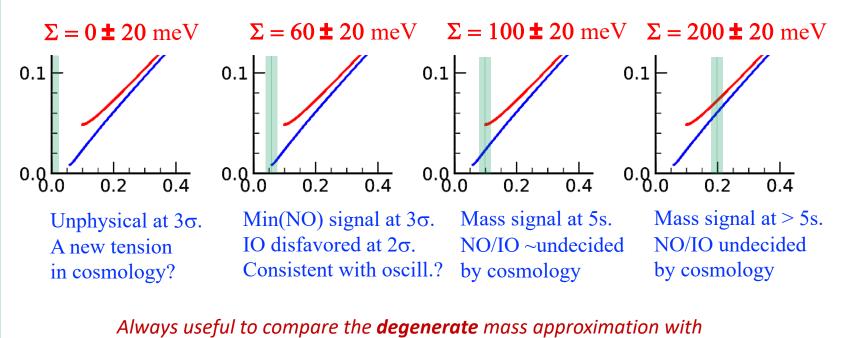
Physical best fit sitting at min(NO), compatible with min(IO) at  $<1\sigma$ Note: small but nonzero fit difference by taking  $\Sigma$ =60=**0**+**9**+**51** rather than **20**+**20**+**20**  More variants can cover up to, say,  $\Sigma < 270$  meV at  $1\sigma$  (akin to weakest PDG limits)



 $\begin{array}{l} \mbox{Rather conservative $\Sigma$ bound, implying $m_{\beta}$ and $m_{\beta\beta}$ (much) below 100 meV} \\ \mbox{\it Mass ordering undecided by cosmology} \end{array}$ 

 $\sigma \sim 45-65 \text{ meV} (\text{now}) \rightarrow \sigma \sim 30 \text{ meV} (\text{baseline}) \rightarrow \sigma \sim 20 \text{ meV} (\text{goal})$ 

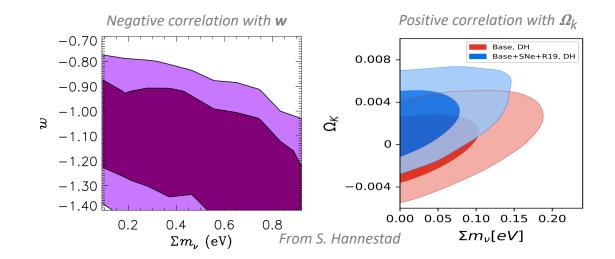
#### Different (and very interesting!) implications, depending on central value of $\Sigma$ , e.g.:



the full-fledged **nondegenerate** case including oscillation  $\Delta m^2_{ij}$ .

Any such result/implication will emerge gradually, and not without debate. Saga of multi-parameter fit variants is likely to continue (focus: from limits to signals):

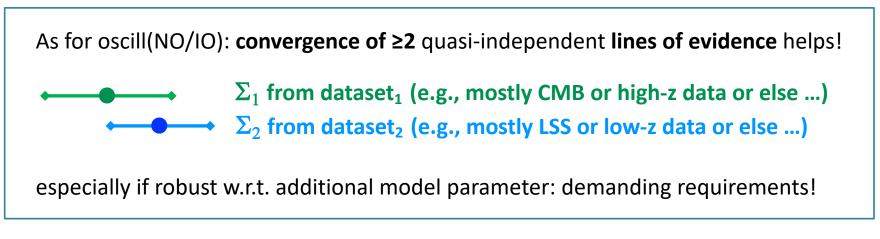
- Old tensions (e.g., H<sub>0</sub>) might not be solved by new data; new tensions may appear
- The  $\Lambda$ CDM model might evolve into a richer model as DE and DM get "understood"
- New model parameters (e.g.,  $w \neq -1$ , curvature...) may be correlated with  $\Sigma$  (see below)
- "Statistical temptations" might enhance claims about  $\Sigma$  signal significance

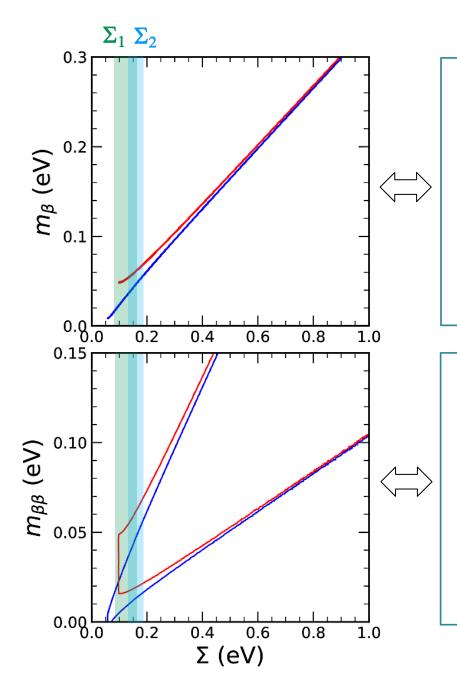


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- Old tensions (e.g., H<sub>0</sub>) might not be solved by new data; new tensions may appear
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- New model parameters (e.g., w  $\neq$  -1, curvature...) may be correlated with  $\Sigma$
- "Statistical temptations" might enhance claims about  $\Sigma$  signal significance

### What will it take to get a convincing signal $\Sigma \simeq \Sigma_0 \pm \sigma$ ?



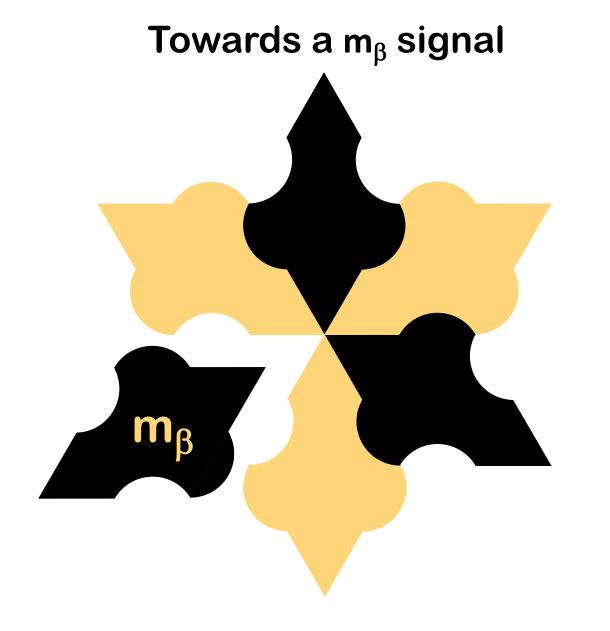


In any case: for settled NO/IO, any estimate for  $\Sigma$  will be in one-to-one correspondence with a m<sub> $\beta$ </sub> estimate

Viceversa, a  $m_{\beta}$  measurement can (dis)confirm  $\Sigma$  and (de)stabilize this corner of cosmology.

Weaker correspondence of  $\Sigma$  with m<sub> $\beta\beta$ </sub>, due to x3 variation from interference of unknown Majorana phases.

Viceversa:  $m_{\beta\beta} > 0$  signal with less than x3 error may constrain cases of max constructive vs destructive interfer.



# $\mathbf{m}_{\beta}$ signal is guaranteed: min $\mathbf{m}_{\beta} \simeq \begin{cases} 9 \text{ meV} (\text{NO}) \\ 50 \text{ meV} (\text{IO}) \end{cases}$

While  $\Sigma$  requires to model the whole universe,  $\mathbf{m}_{\beta}$  requires to model source + detector  $\rightarrow$  Instrinsically robust and pivotal role of  $\beta$  decay.

One must find the  $\mathbf{m}_{\beta}$  signal at any cost!

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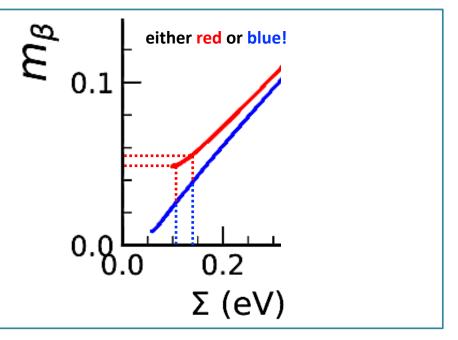
One must find the  $\mathbf{m}_{\beta}$  signal at any cost!

There is realistic path to go from  $\sim 200$  meV (KATRIN) to  $\sim 50$  meV (PROJECT 8)

Timescale: ~10 yrs. Other projects explored, in R&D phase (J. Formaggio's talk)

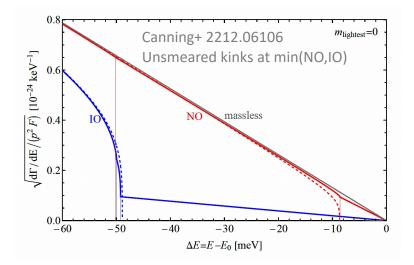
If lucky, in **203X** we might see **up to two absolute mass signals** and analyze them in fine details: **a new frontier of global fits** 

If not: **path**  $m_{\beta} \sim 50 \rightarrow \sim 9 \text{ meV}$ needs to be envisaged. Hard but absolutely necessary!



## Fine details in future global analyses...

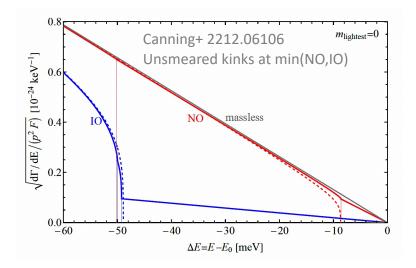
Improvements in  $\mathbf{m}_{\beta}$  sensitivity might come with improvements in resol.  $\Delta E_{\beta}$  from current  $\Delta E_{\beta} \sim 1 \text{ eV}$  (KATRIN) to, hopefully,  $\Delta E_{\beta} \sim O(\sqrt{\Delta m^2}) \sim 50 \text{ meV}$  or less  $\rightarrow$  possible sensitivity to kink(s) info rather than just overall smeared distortion



Concerning  $\Sigma$ : as noted, it will be worthwhile to check small differences between the **degenerate** mass approximation and **nondegenerate** masses

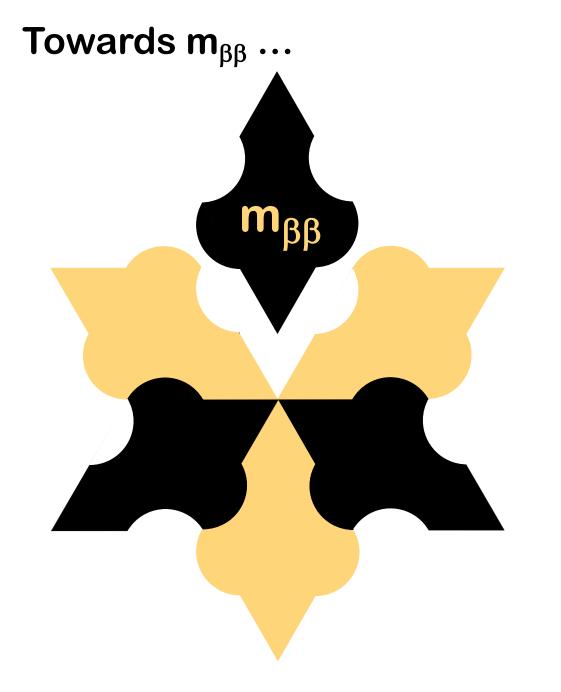
# Fine details in future global analyses...

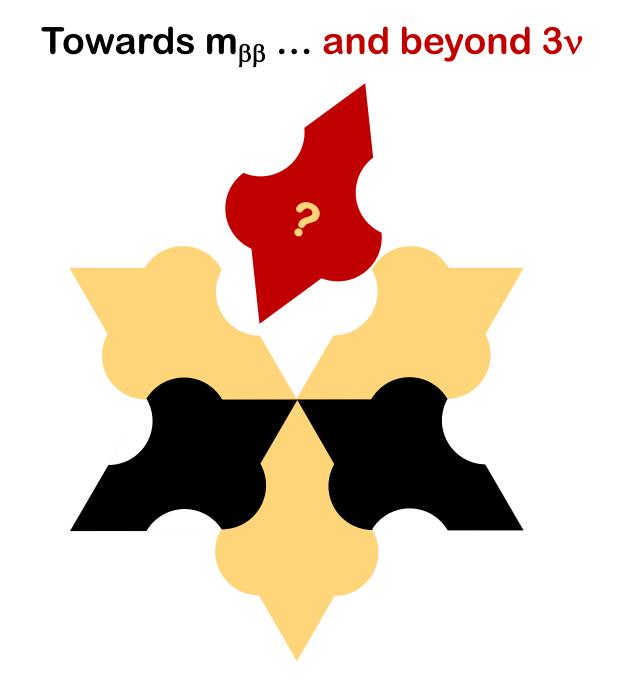
Improvements in  $\mathbf{m}_{\beta}$  sensitivity might come with improvements in resol.  $\Delta E_{\beta}$  from current  $\Delta E_{\beta} \sim 1 \text{ eV}$  (KATRIN) to, hopefully,  $\Delta E_{\beta} \sim O(v\Delta m^2) \sim 50 \text{ meV}$  or less  $\rightarrow$  possible sensitivity to **kink(s) info** rather than just overall smeared distortion



Concerning  $\Sigma$ : as noted, it will be worthwhile to check small differences between the **degenerate** mass approximation and **nondegenerate** masses

There may be a little bit more information than just 2 param. ( $m_{\beta}$  and  $\Sigma$ )! Possible slight sensitivity to the  $v_i$  mass distribution, hopefully consistent with the one dictated by the true mass ordering + oscillation splittings.

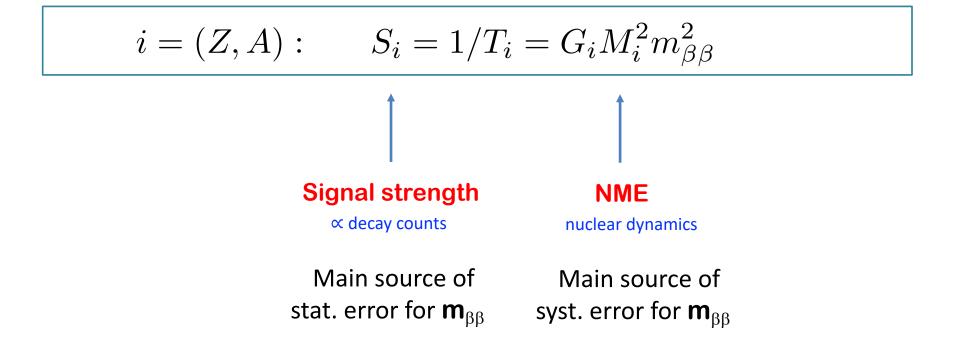




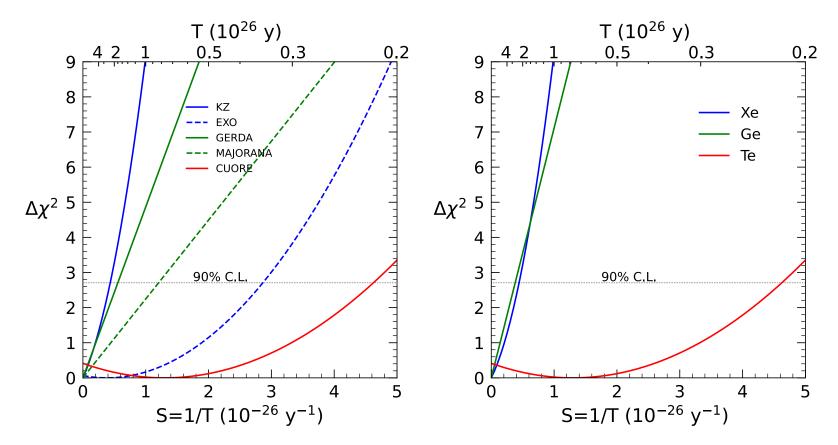
# $\mathbf{m}_{\beta\beta}$ signal is not guaranteed: $\min \mathbf{m}_{\beta\beta} \simeq \begin{cases} 0 \text{ meV} & (NO) \\ 18 \text{ meV} & (IO) \end{cases}$

But Majorana/Dirac discrimination is of fundamental importance! (talks: M. Agostini, S. Petcov)

Signal estimates depend on nuclear model of (Z,A) + model of source/detector



#### Signal strength likelihood for latest results



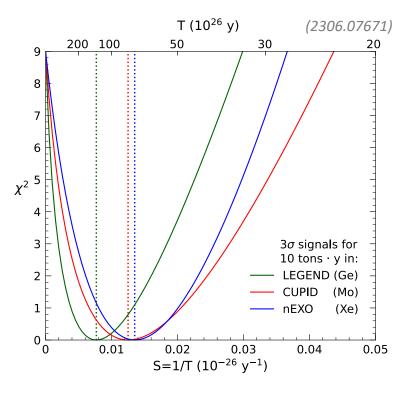
Best fit at (or close to) null signal  $\rightarrow$  NME-dependent upper limits on  $m_{\beta\beta}$ 

A plea to experimentalists: please always publish  $\mathcal{L}(S)$ , not just S at 90% CL! Otherwise: impossible to combine independent results, even in same (Z,A)

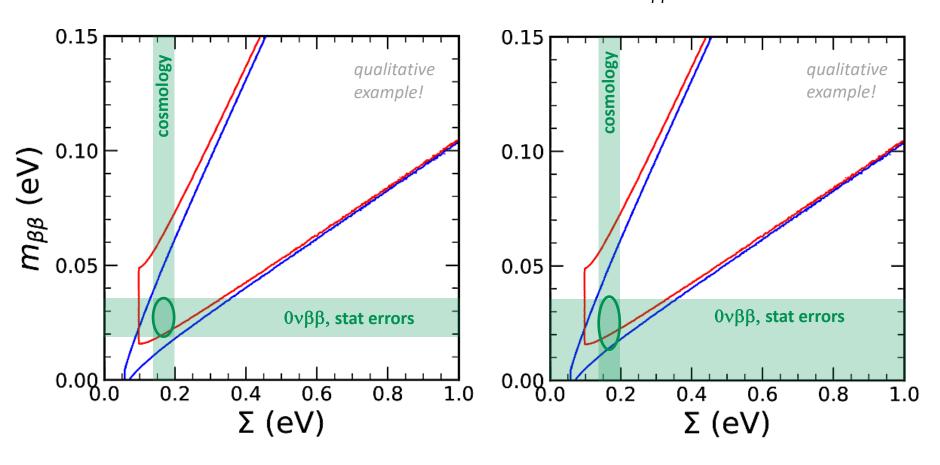
*Lisi, Marrone, Nath, 2306.07671* 

### Realistic path to reach $\geq 3\sigma$ evidence down to $m_{\beta\beta} \sim 18$ meV, even for lowest known NME: Ton-scale masses, 10-year time scale $\rightarrow 10$ ton yr exposure (talk by M. Agostini)

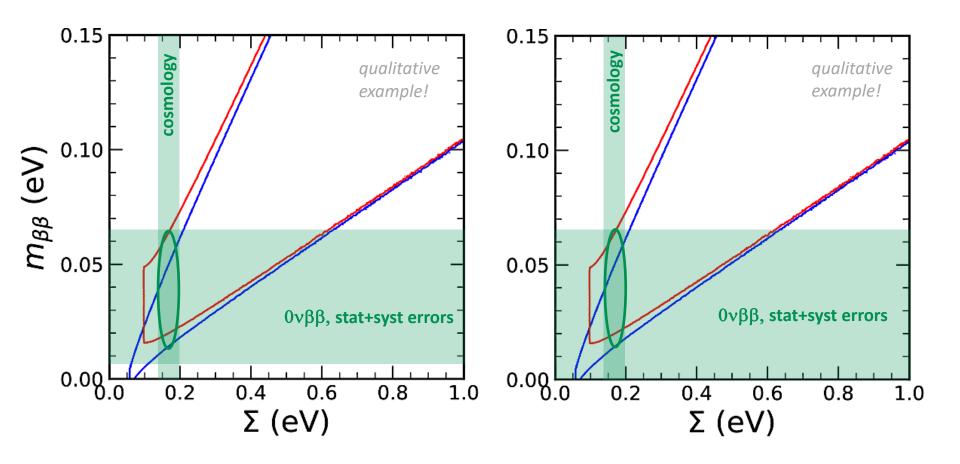
### Signal strength likelihood for prospective $3\sigma$ evidence:



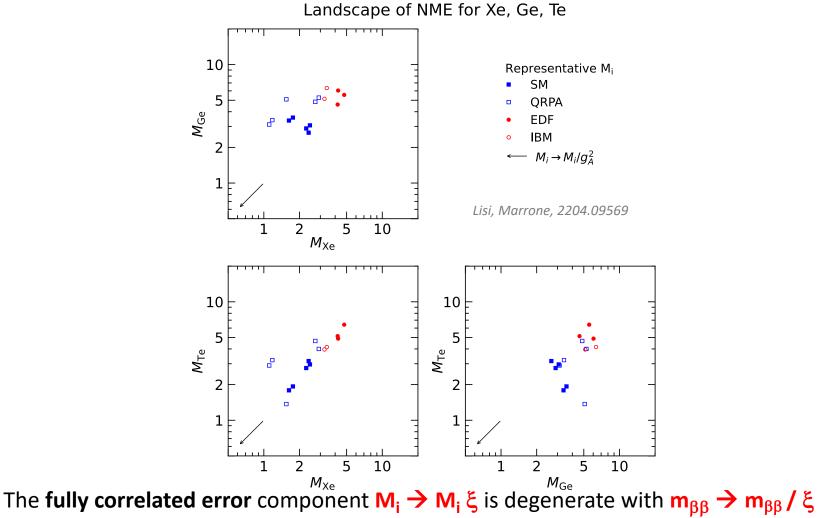
In each expt.,  $\pm 1\sigma$  stat. spread of  $m_{\beta\beta} \propto \sqrt{S}$  smaller than "x3 variation" (even better for >3 $\sigma$  evidence, or by combining  $\geq 2$  experiments) In combination with a signal for  $\Sigma$  (of for m<sub> $\beta$ </sub>, or both) some constraints on Majorana phases may emerge (even for *upper limits only* on m<sub> $\beta\beta$ </sub>)



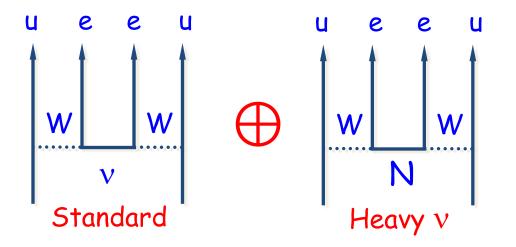
### Unfortunately... washed out by <u>current</u> x3 spread of NMEs



### M<sub>i</sub> spread dangerous because it's: (1) large; (2) correlated among i=(Z,A)



and is not reduced by combining multi-isotope signals (Faessler+, 1103.2504)

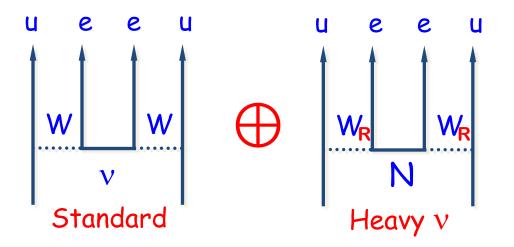


### New physics beyond 3v?

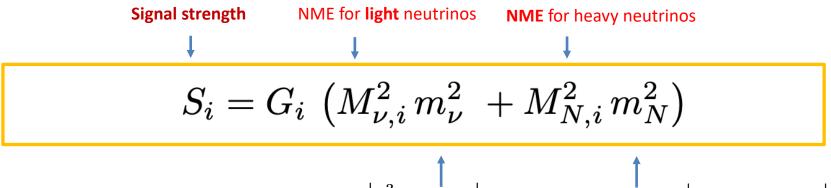
E.g., possible to have both light and heavy v in many theo. models, e.g. see-saw



Large and correlated NME spread may also prevent discrimination of **new physics contributions** (if any)



Light and heavy v exchange may be ~non-interfering\*, e.g. in LR-symmetric models: (\*simplest case, no extra phases)



$$m_{
u} = \left| \sum_{k=1}^{3} U_{ek}^2 m_k \right|$$

$$m_N = \frac{m_W^4}{m_{W_R}^4} \left| \sum_h V_{eh}^2 \frac{m_p m_e}{M_h} \right|$$

Effective Majorana mass (light) Effective Majorana mass (heavy) Need two equations (two isotopes i,j) for two mass unknowns:

$$\begin{bmatrix} S_i G_i^{-1} \\ S_j G_j^{-1} \end{bmatrix} = \begin{bmatrix} M_{\nu,i}^2 & M_{N,i}^2 \\ M_{\nu,j}^2 & M_{N,j}^2 \end{bmatrix} \begin{bmatrix} m_{\nu}^2 \\ m_{\nu}^2 \end{bmatrix}$$

$$\begin{array}{c} \mathbf{DATA} \\ \mathbf{NME} \\ \mathbf{hinematics} \end{array} \quad \begin{array}{c} \mathbf{NME} \\ \mathbf{Majorana\ masses} \\ (\text{nuclear physics}) \end{array} \quad \begin{array}{c} \text{Majorana\ masses} \\ (\text{particle physics}) \end{array}$$

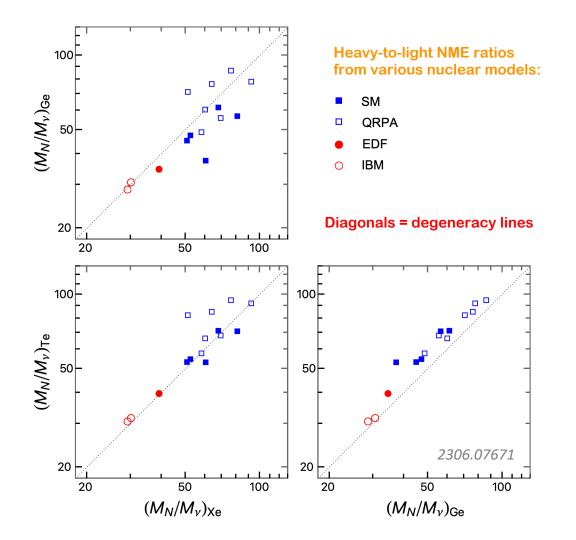
With three (or more) isotopes: can make further checks.  $\rightarrow$  Need multi-isotope  $0\nu\beta\beta$  decay searches

Non-degenerate solution iff matrix determinant is non-zero:

$$\frac{M_{N,i}}{M_{\nu,i}} \neq \frac{M_{N,j}}{M_{\nu,j}}$$

NME heavy/light <u>ratio</u> uncertainties  $\rightarrow$ 

#### Large spread of heavy/light ratios of NME around the degeneracy lines:



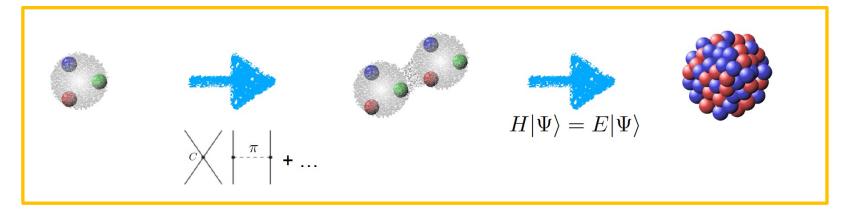
→ Difficult to separate heavy v contribution - and new physics in general [Taming degeneracy by error control will be easier for largely off-diagonal central values]

#### But...there is a realistic path towards improved NME estimates in the wider context of ab-initio approaches in nuclear physics

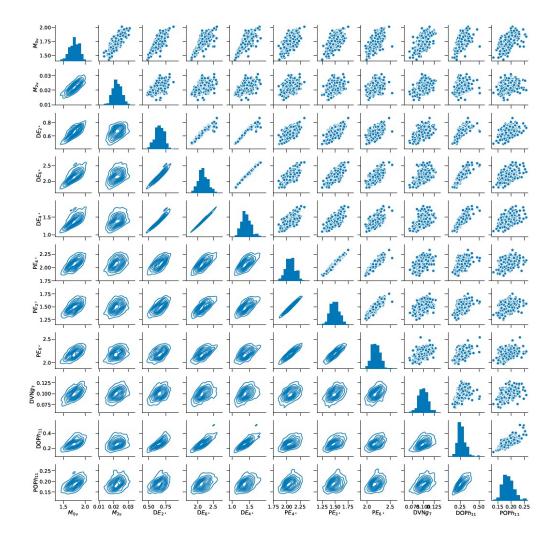
2203.12169

Neutrinoless Double-Beta Decay: A Roadmap for Matching Theory to Experiment

Ab-initio approaches: start from well-motivated NN and NNN forces and solve multi-N Schroedinger equation with systematically improvable methods



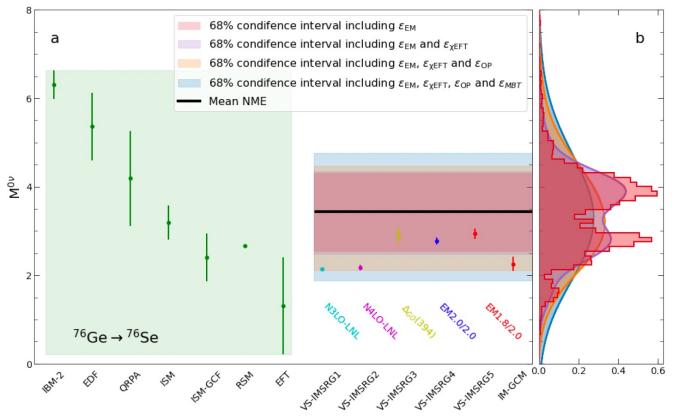
See talks by T. Miyagi at ISPUN 2023 J. Menendez at HADRON 2023 A. Ekstrom at HIRSCHEGG 2023 Benchmark method(s) with a variety of nuclear data and processes (including  $2\nu\beta\beta$ )



E.g., Horoi+ 2302.03664

#### Obtain probability distribution for calculated NME (not yet correlations etc.)

E.g., Belley+ 2308.1564 for <sup>76</sup>Ge



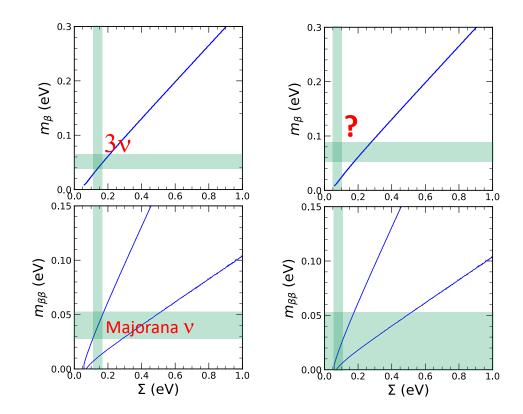
Already improvements w.r.t. usual x3 spread. Room for significant progress. We may hope in NME (co)variances commensurate to ton-scale requirements.

# Epilogue

### Conceivable to dream about scenarios like these at **NEUTEL 203X**:

We may experience some nightmares, as well as **surprises**...

... but we will learn a lot new from nature at very different scales



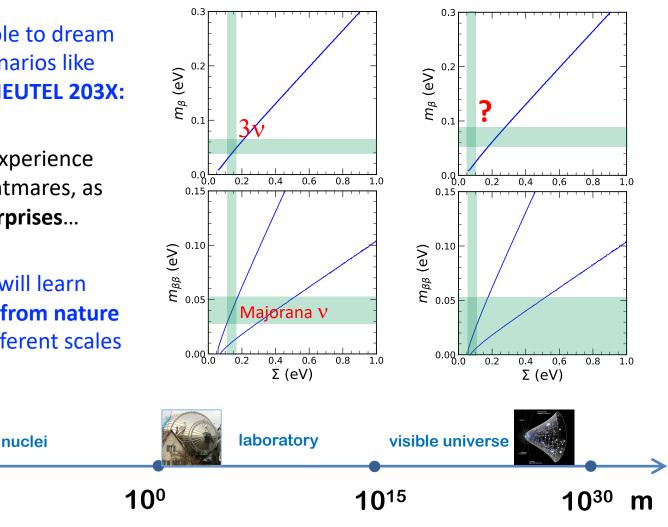
# **Epilogue**

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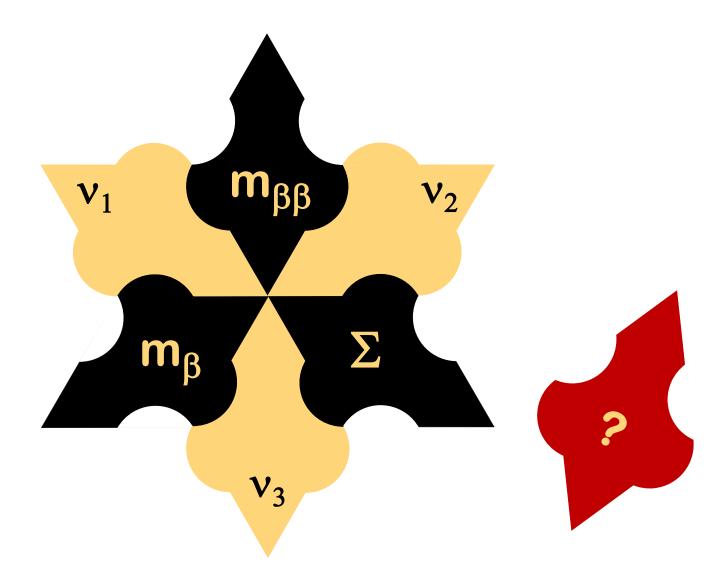
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**10**<sup>-15</sup>

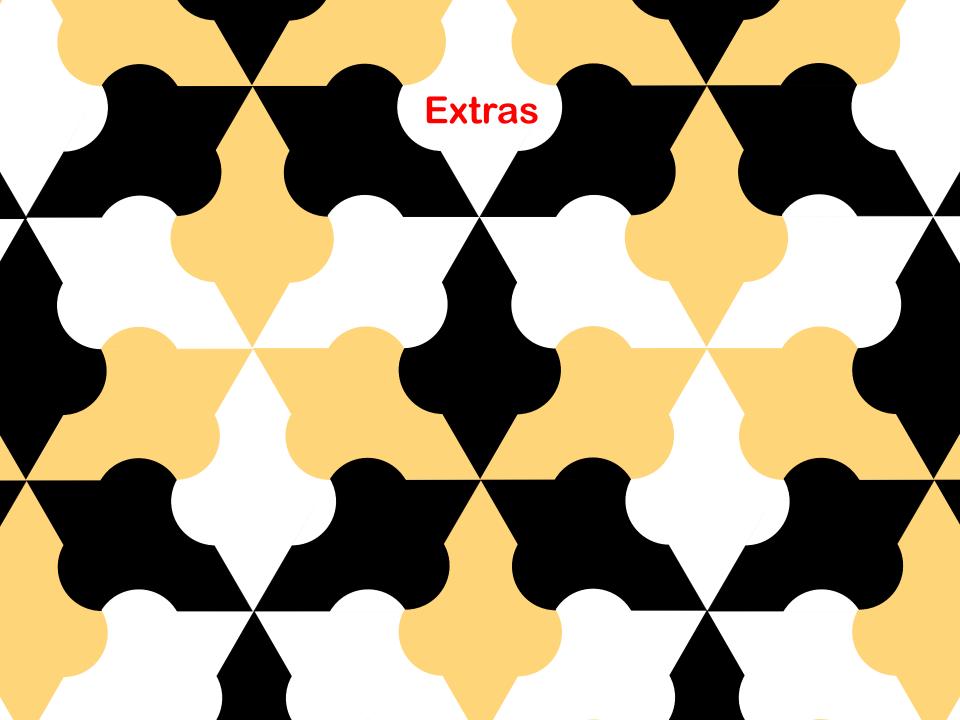


[... here, a log scale is appropriate!]

## Thank you for your attention!



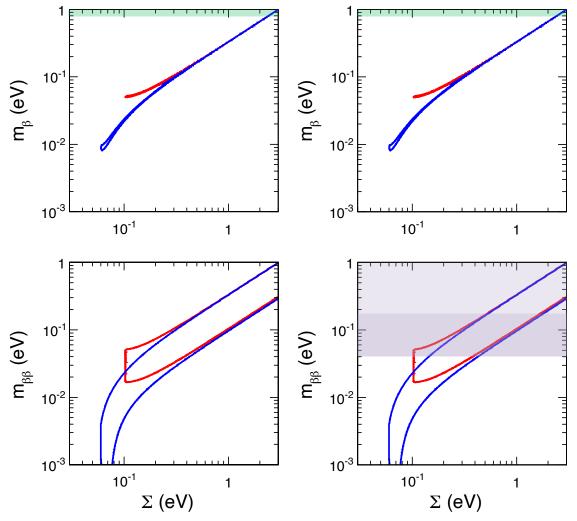
Work supported by PRIN 2022 "PANTHEON" (Italian MUR) & Network "TASP" (INFN)



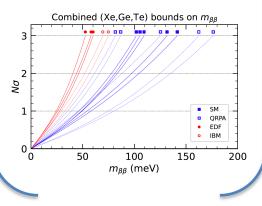


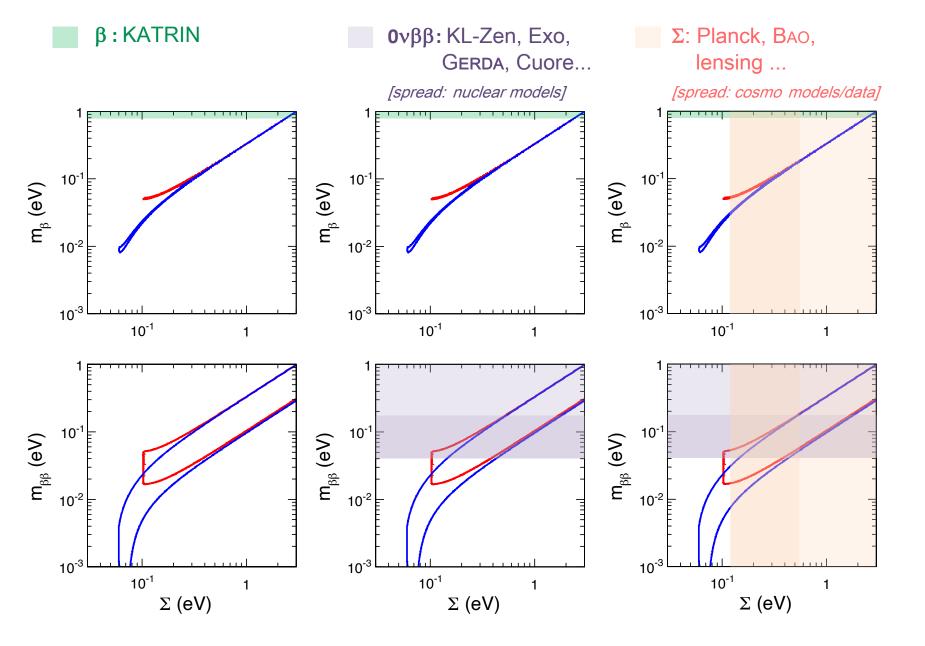
**0** $\nu$ ββ: KL-Zen, Exo, Gerda, Cuore...

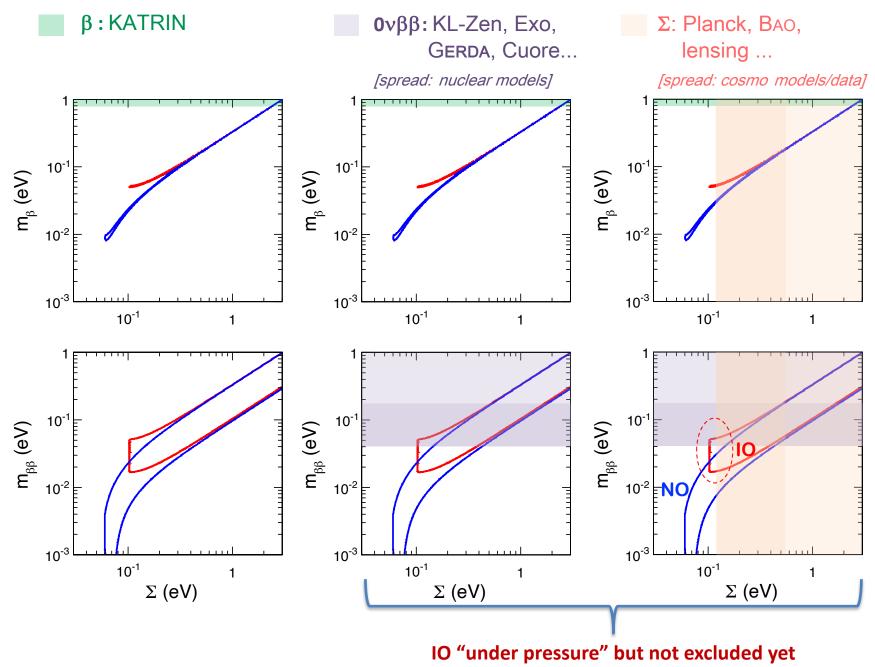
[spread: nuclear models]

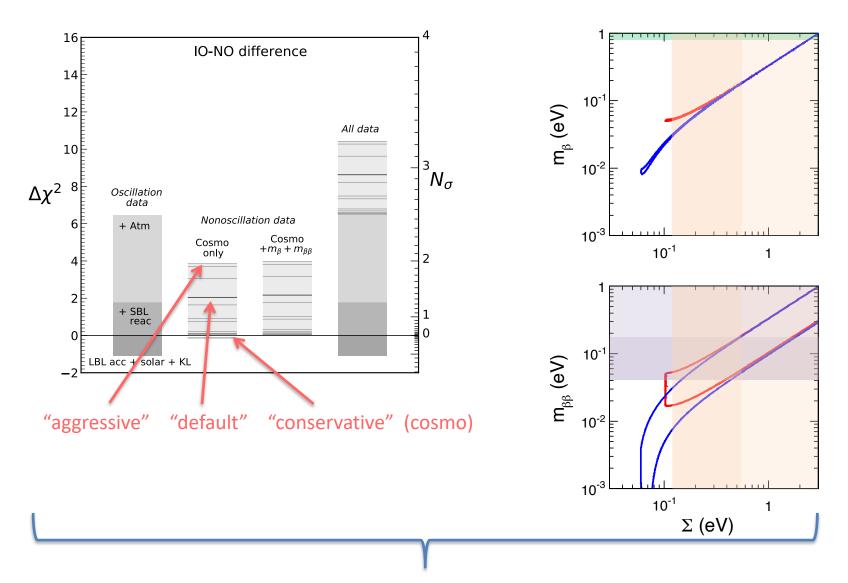


E.g., spread of upper bounds from Xe+Ge+Te data by using 15 nuclear matrix elements from 4 classes of nucl. models. e-print 2204.09569









IO currently disfavored at ~ $3\sigma$  by combining oscillation + nonoscillation data