

Earth tomography with supernova neutrinos at future neutrino detectors

based on RH, O. Mena and S. Palomares-Ruiz, Phys.Rev.D 108 (2023) 8, 083011

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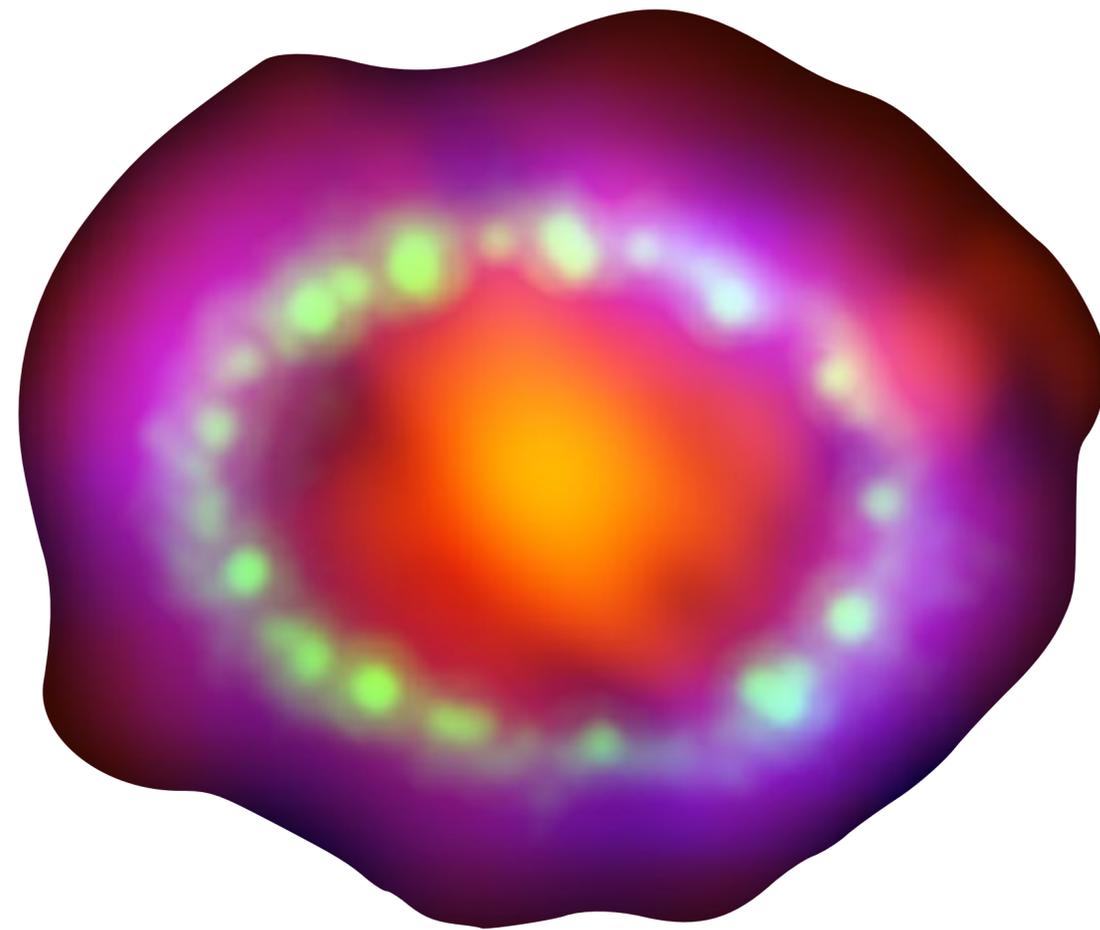
Asymmetry
Essential Asymmetries of Nature

IFIC
INSTITUT DE FÍSICA
CORPUSCULAR

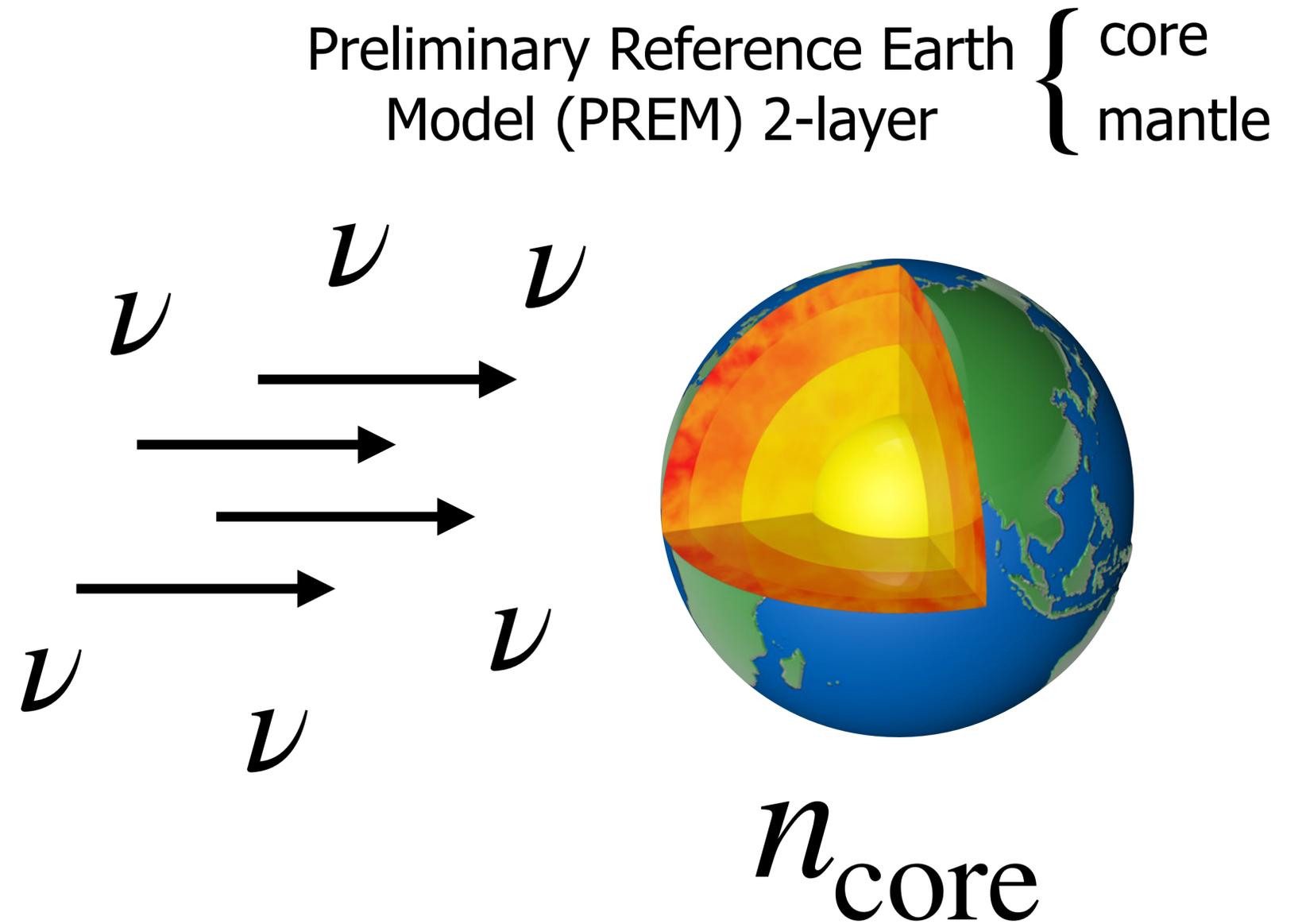
CSIC
CONSEJO SUPERIOR DE INVESTIGACIONES CIENTÍFICAS

SSM
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Main goal of this work

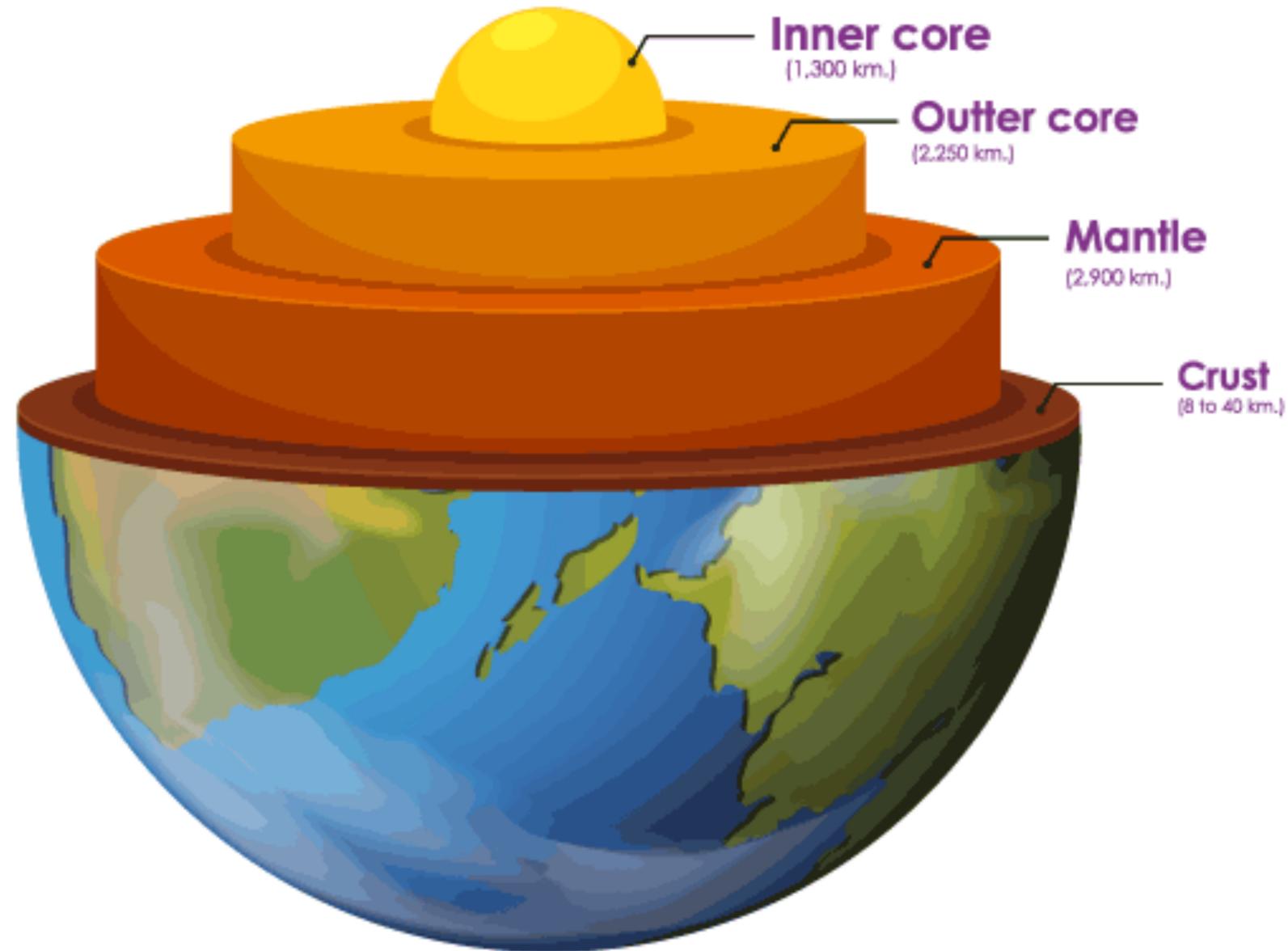


galactic SN $d_{\text{SN}} \sim 10 \text{ kpc}$



Earth inner structure

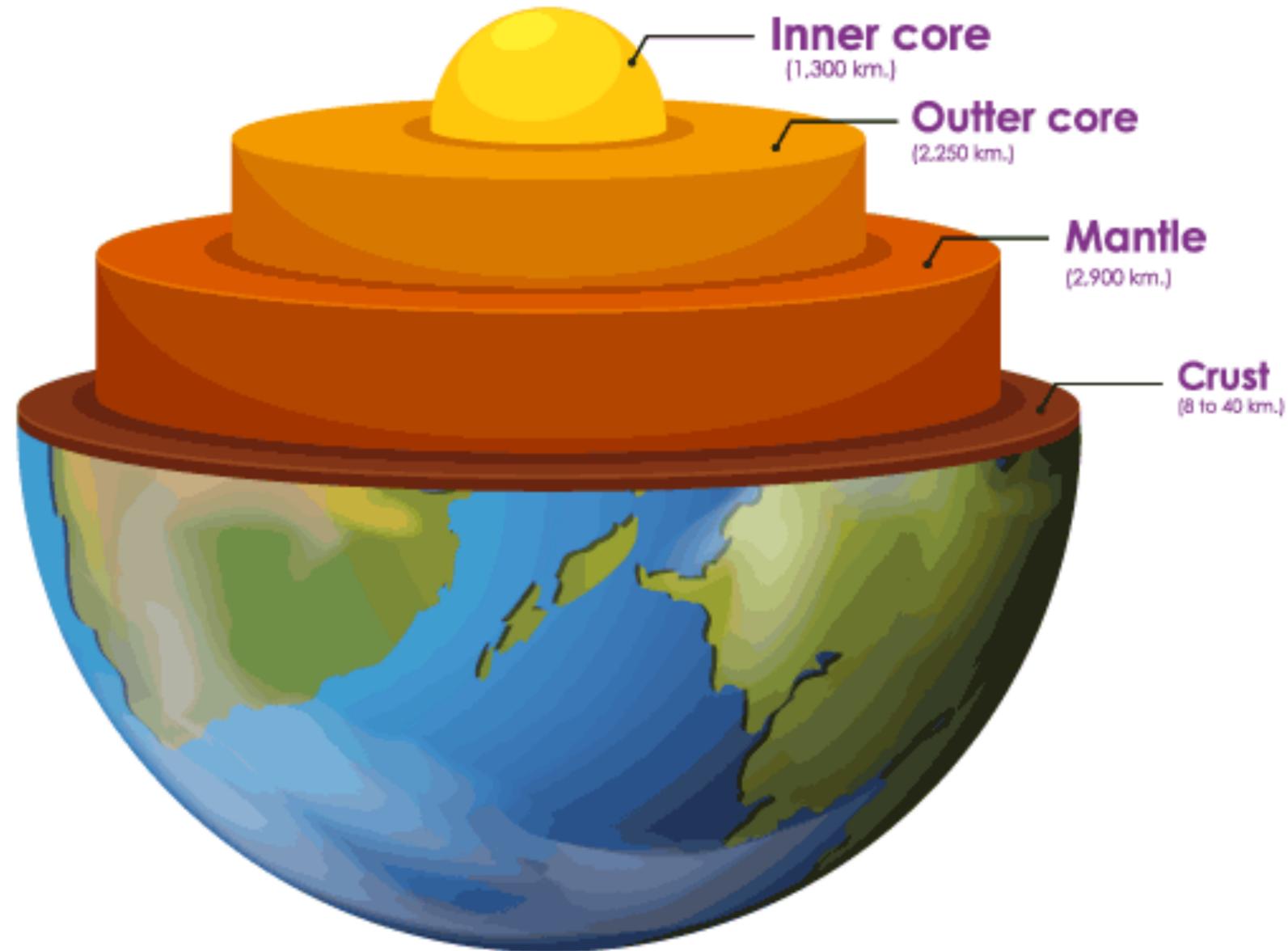
- Modern Earth model:



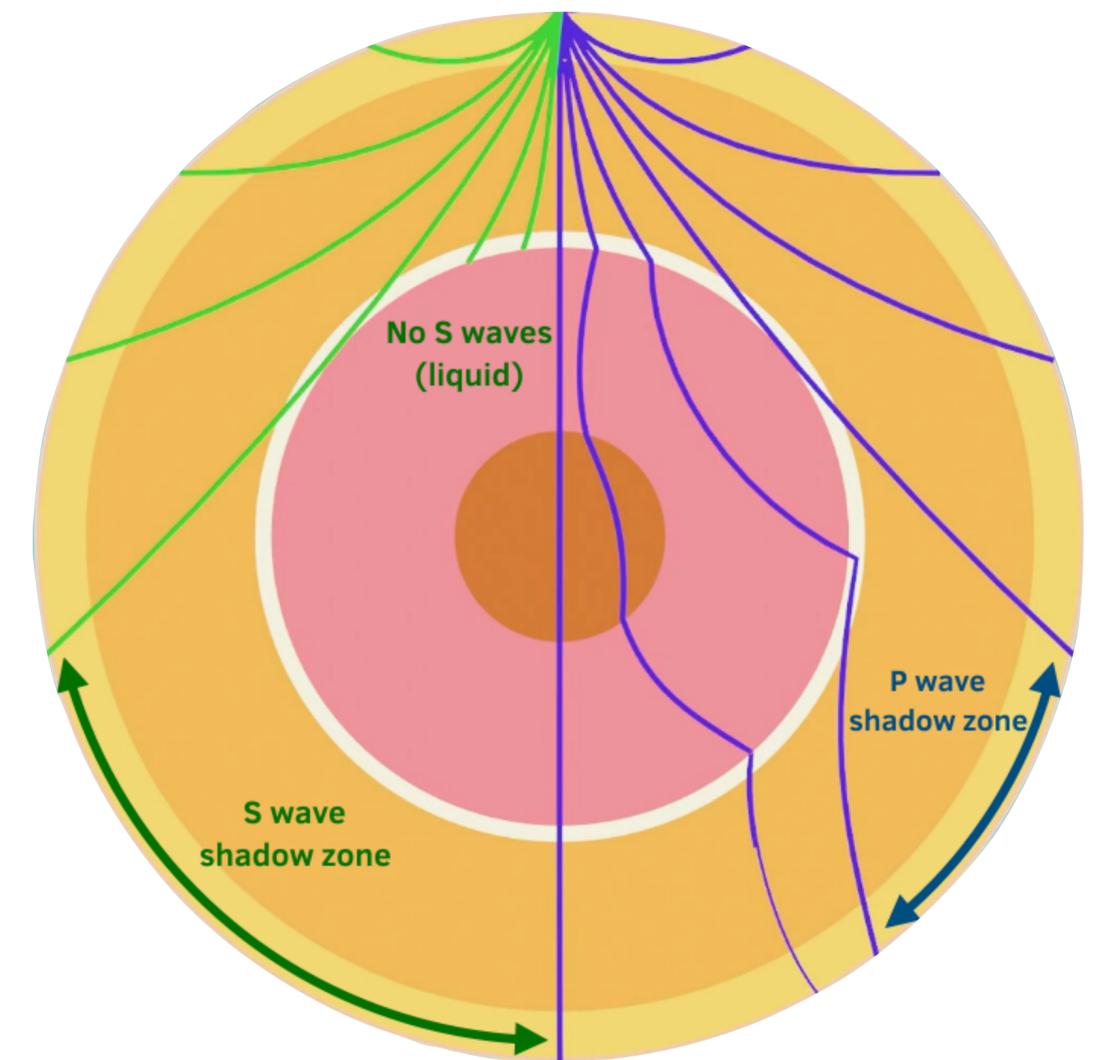
- How do we know this?

Earth inner structure

- Modern Earth model:



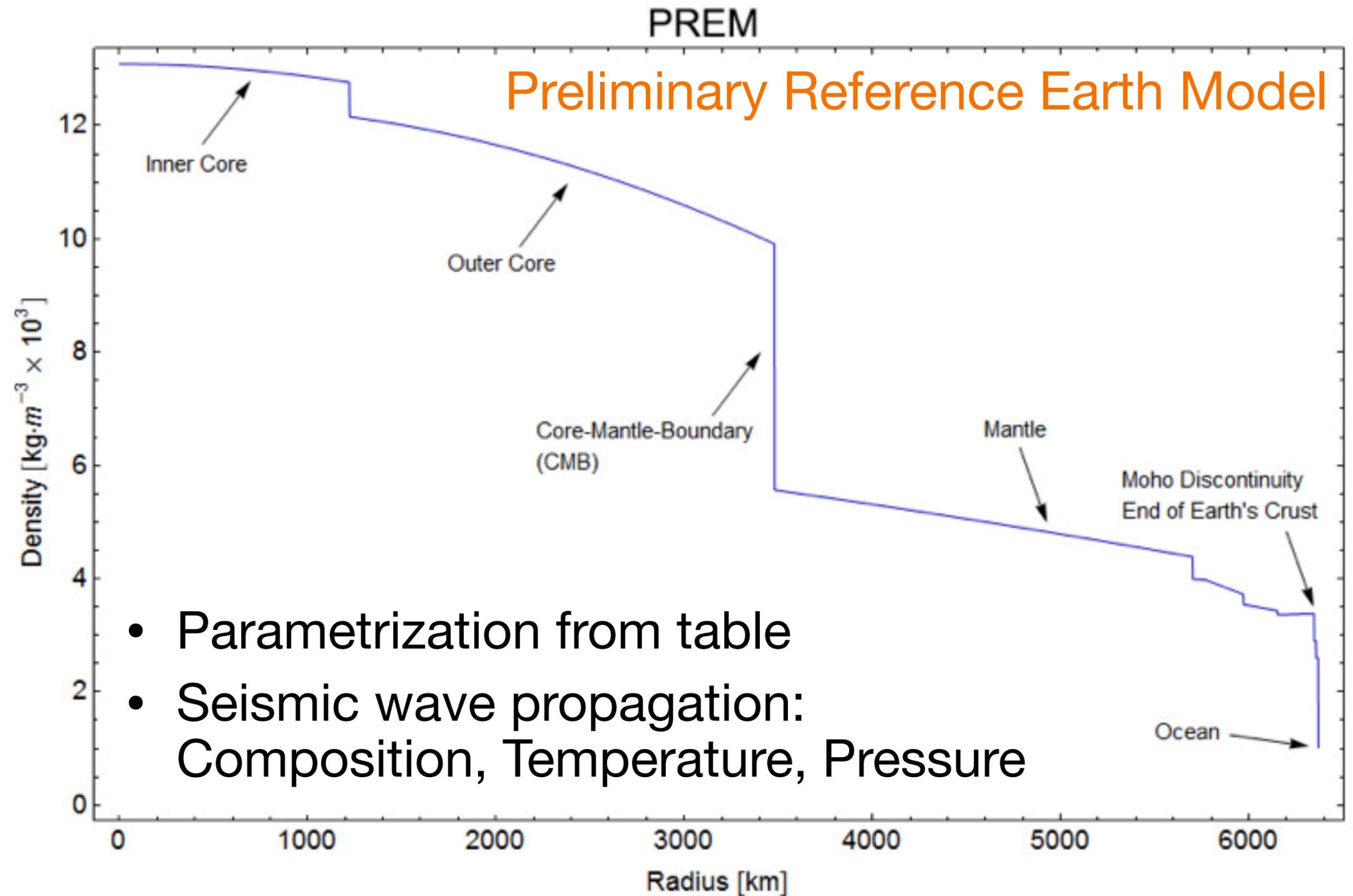
- How do we know this?
Seismic waves!



Earth inner structure

A.M. Dziewonski, D.L. Anderson,
Phys.Earth Planet.Interiors 25 (1981), 297-356

LEVEL	RADIUS KM	DEPTH KM	DENSITY G/CCH
1	0.0	6371.0	13.08848
2	100.0	6271.0	13.08630
3	200.0	6171.0	13.07977
4	300.0	6071.0	13.06888
5	400.0	5971.0	13.05364
6	500.0	5871.0	13.03404
7	600.0	5771.0	13.01009
8	700.0	5671.0	12.98178
9	800.0	5571.0	12.94912
10	900.0	5471.0	12.91211
11	1000.0	5371.0	12.87073
12	1100.0	5271.0	12.82501
13	1200.0	5171.0	12.77493
14	1221.5	5149.5	12.76360
15	1221.5	5149.5	12.16634
16	1300.0	5071.0	12.12500
17	1400.0	4971.0	12.06924
18	1500.0	4871.0	12.00989
19	1600.0	4771.0	11.94682
20	1700.0	4671.0	11.87990



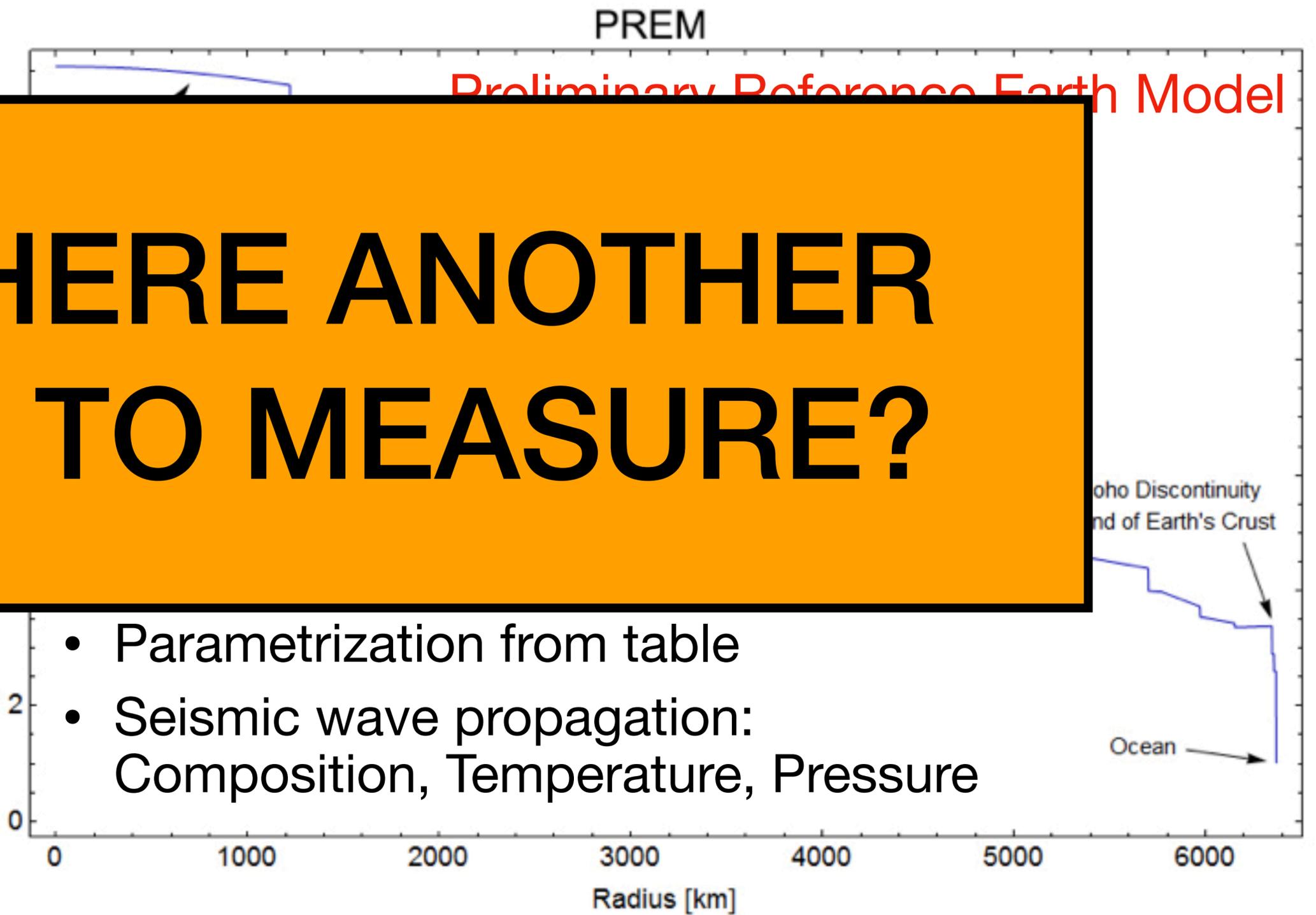
Earth inner structure

A.M. Dziewonski, D.L. Anderson,
Phys.Earth Planet.Interiors 25 (1981), 297-356

LEVEL	RADIUS KM
1	0.0
2	100.0
3	200.0
4	300.0
5	400.0
6	500.0
7	600.0
8	700.0
9	800.0
10	900.0
11	1000.0
12	1100.0
13	1200.0
14	1221.5
15	1221.5
16	1300.0
17	1400.0
18	1500.0
19	1600.0
20	1700.0

**IS THERE ANOTHER
WAY TO MEASURE?**

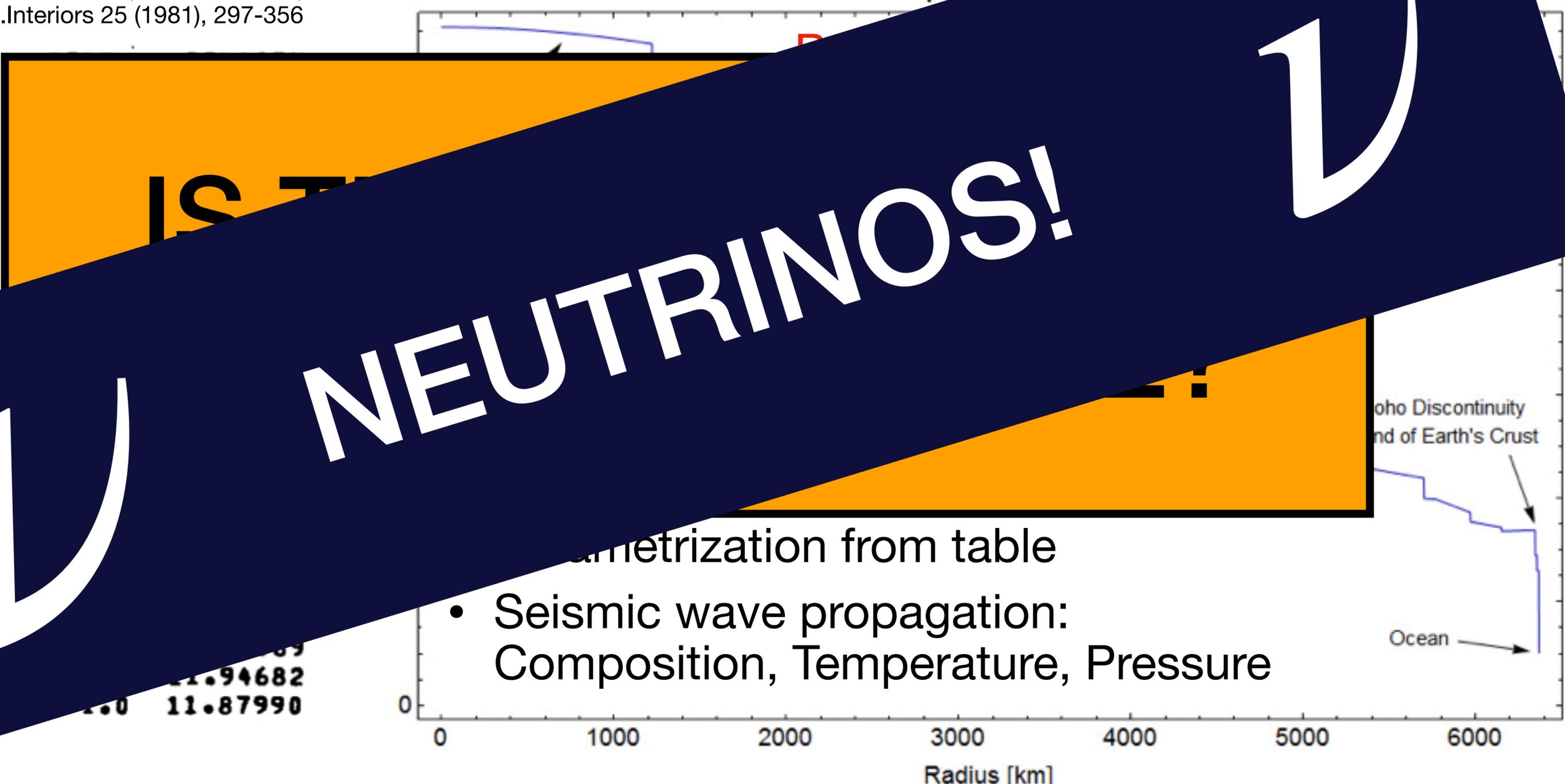
- Parametrization from table
- Seismic wave propagation:
Composition, Temperature, Pressure



Earth inner structure

A.M. Dziewonski, D.L. Anderson,
Phys.Earth Planet.Interiors 25 (1981), 297-356

LEVEL	RADIUS KM
1	0.0
2	100.0
3	200.0
4	300.0
5	400.0
6	500.0
7	



IS T

NEUTRINOS!

Moho Discontinuity
End of Earth's Crust

Ocean

- Seismic wave propagation:
Composition, Temperature, Pressure

Neutrino tomography types

OSCILLATION TOMOGRAPHY

$$\frac{d\phi(E, x)}{dx} = -i \mathcal{H}_{\text{flavor}} \phi(E, x)$$

$$\mathcal{H}_{\text{flavor}} = \frac{1}{2E} U M^2 U^\dagger + \mathbb{V}$$

Vacuum

Matter

$$P_{2\nu}(\nu_\alpha \rightarrow \nu_\beta) = \sin^2 2\theta_m \sin^2 \left(\frac{\Delta_m L}{4E} \right)$$

$$\Delta_m = \sqrt{(\Delta m^2 \sin 2\theta \mp 2E \mathbb{V})^2 + (\Delta m^2 \cos 2\theta)^2}$$

- Man-made beams
- Solar neutrinos

- Supernova neutrinos
- Atmospheric Neutrinos

$$\sin^2 2\theta_m = \sin^2 2\theta \left(\frac{\Delta m^2}{\Delta_m} \right)^2$$

Neutrino tomography types

OSCILLATION TOMOGRAPHY

$$\frac{d\phi(E, x)}{dx} = -i \mathcal{H}_{\text{flavor}} \phi(E, x)$$

$$\mathcal{H}_{\text{flavor}} = \frac{1}{2E} U M^2 U^\dagger + \mathbf{V}$$

$$P_{2\nu}(\nu_\alpha \rightarrow \nu_\beta) = \sin^2 2\theta_m \sin^2 \left(\frac{\Delta m^2 L}{4E} \right)$$

Vacuum

Matter

$$\Delta m^2 = \sqrt{(\Delta m^2 \sin 2\theta \mp 2E \mathbf{V})^2 + (\Delta m^2 \cos 2\theta)^2}$$

- Man-made beams
- Solar neutrinos

- Supernova neutrinos
- Atmospheric Neutrinos

$$\sin^2 2\theta_m = \sin^2 2\theta \left(\frac{\Delta m^2}{\Delta m^2} \right)^2$$

ABSORPTION TOMOGRAPHY

$$\frac{d\phi(E, x)}{dx} = -n(x) \sigma \phi(E, x)$$

- Man-made beams

- Cosmic neutrinos

- Atmospheric neutrinos

MeV
↑
GeV

Neutrino tomography results!

Slide taken from S. Palomares-Ruiz talk

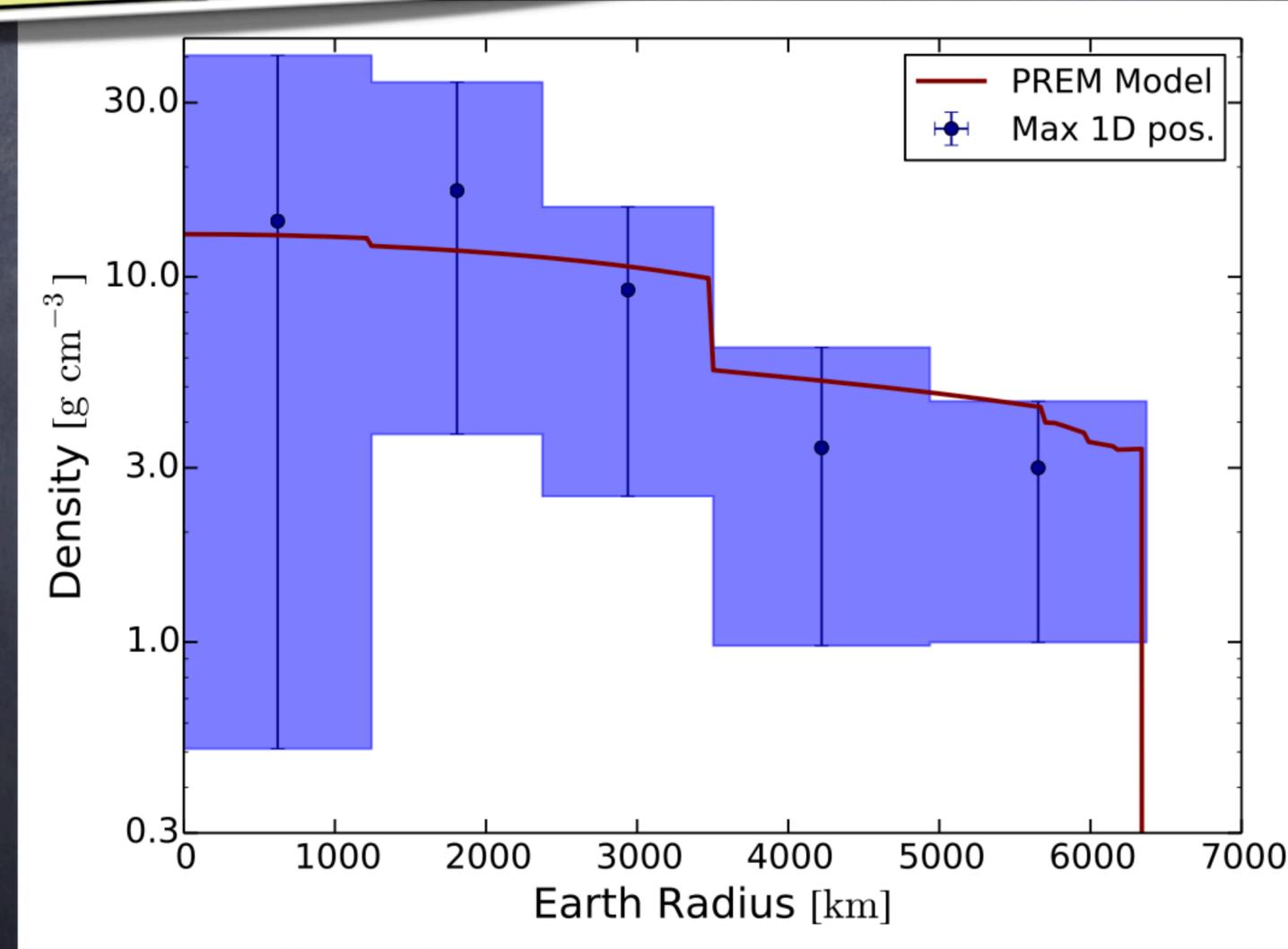
https://drive.google.com/file/d/1yaF_Wo9SiDd92Ub34ji00s5YyIpc9JHU/view

- 1 year data of atmospheric neutrinos at Ice-Cube neutrino detector

First Earth tomography with neutrinos!

MAIN RESULT: 1-D DENSITY PROFILE

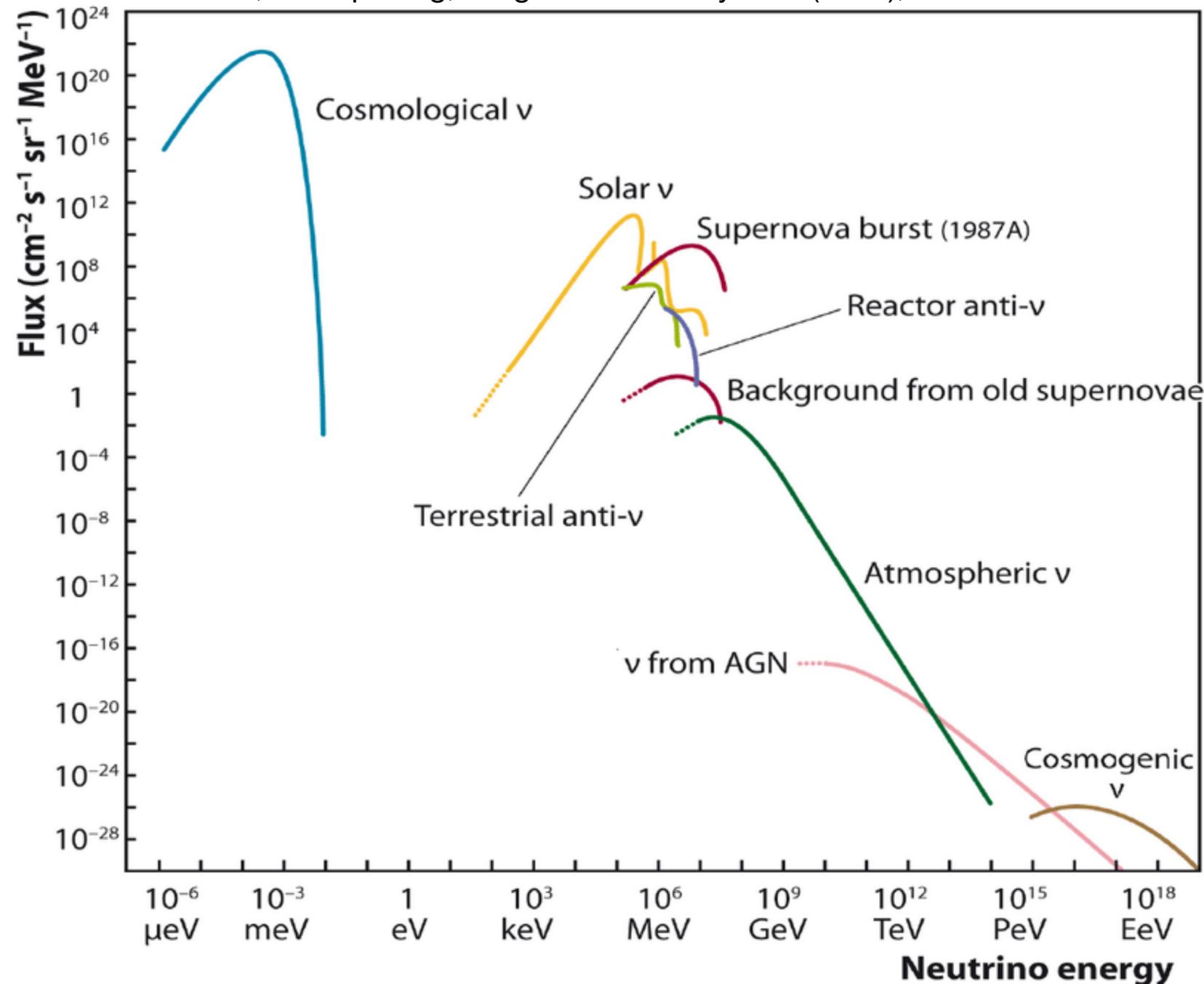
unlike reconstructions with seismic data, no constraint on the Earth mass or moment of inertia



A. Donini, SPR and J. Salvado, Nature Physics 15:37, 2019

Supernova neutrinos

U.F. Katz, Ch. Spiering, Prog.Part.Nucl.Phys. 67 (2012), 651-704



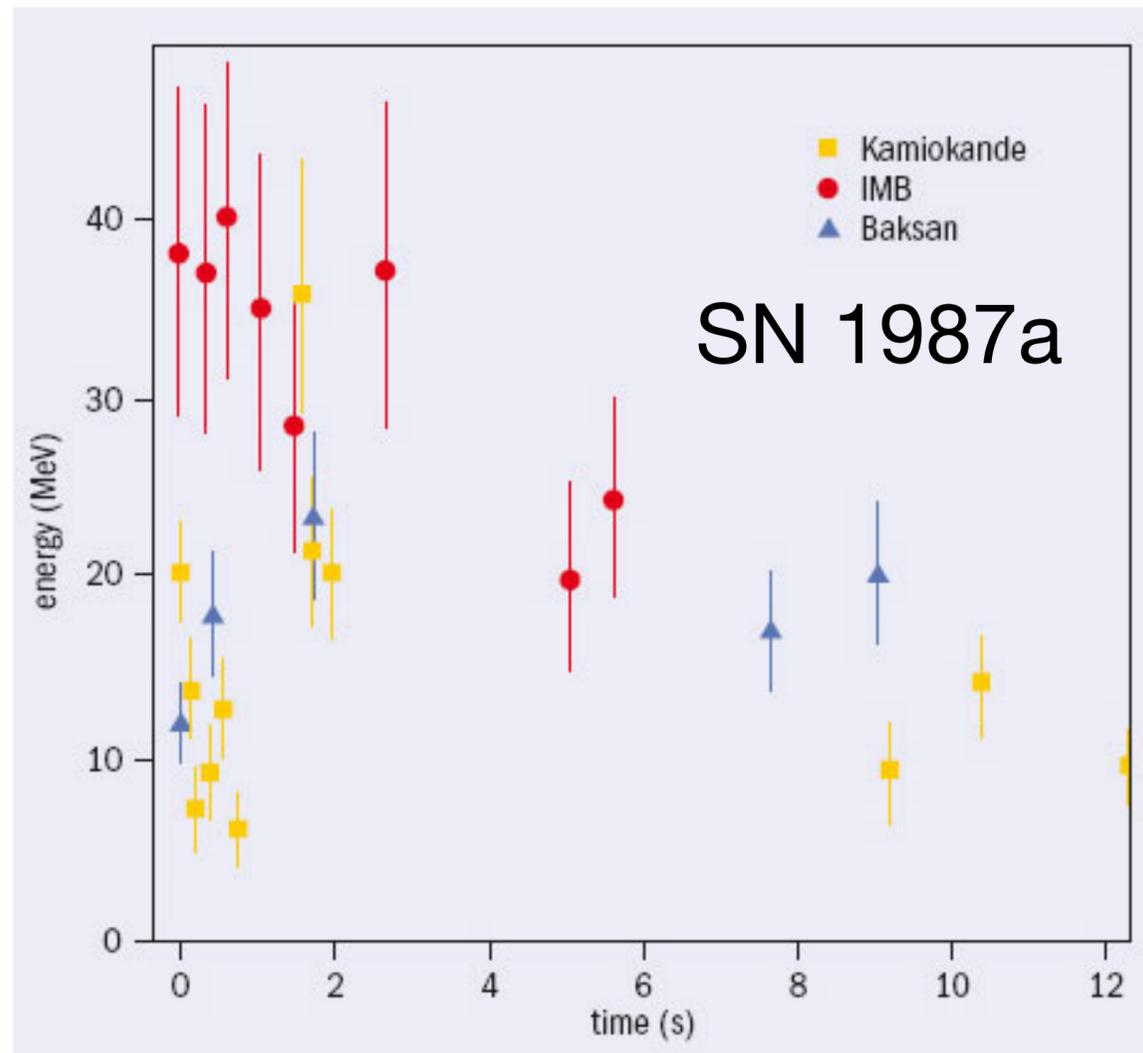
- Core-collapse SN is the violent explosion during death of massive stars.
- 99% energy of star ($\sim 10^{53}$ erg) is released in the form of neutrinos.

Excellent source due to high flux and low background (applying temporal cut ~ 10 s).

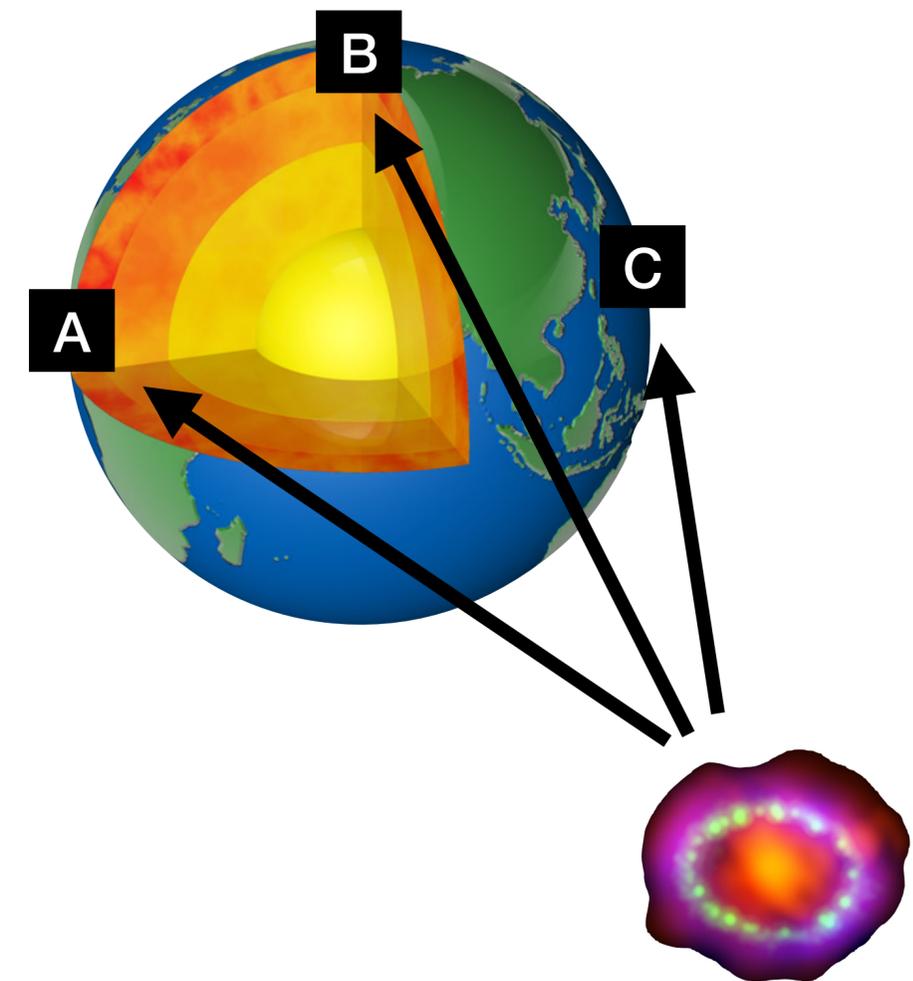
Supernova neutrinos

Main drawbacks

Uncertainty on fluxes



One direction per detector

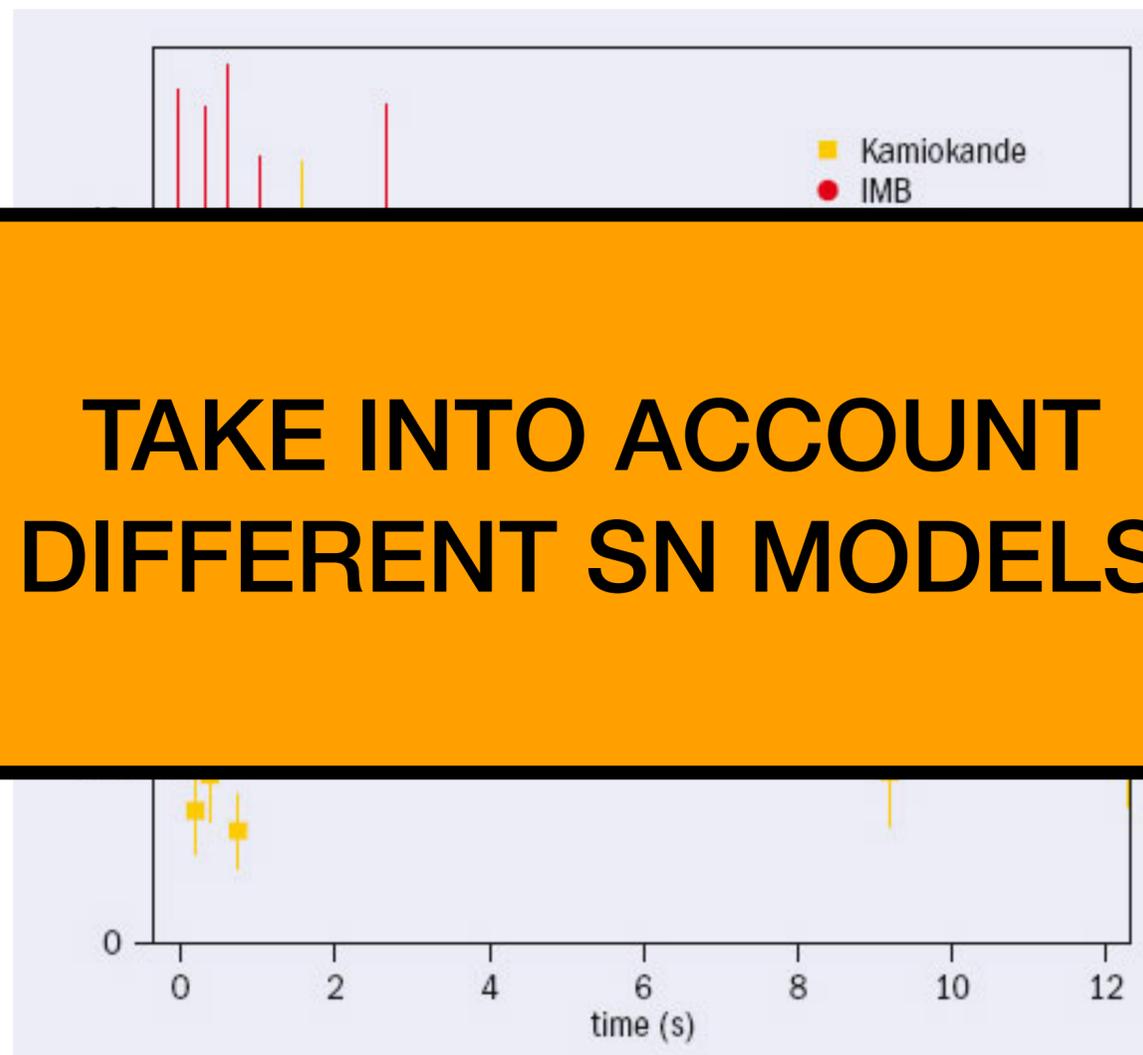


Supernova neutrinos

Main drawbacks

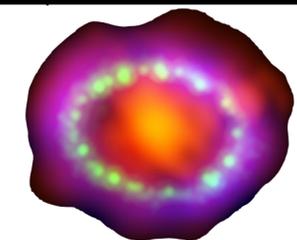
Uncertainty on fluxes

One direction per detector



**TAKE INTO ACCOUNT
DIFFERENT SN MODELS**

**TAKE INTO ACCOUNT
DIFFERENT INCIDENT
DIRECTIONS**



Supernova neutrinos

Main drawbacks

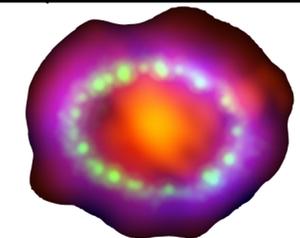
Uncertainty on fluxes

**GALACTIC SN RATE:
~1 PER CENTURY**

**ACCOUNT
DIFFERENT INCIDENT
DIRECTIONS**

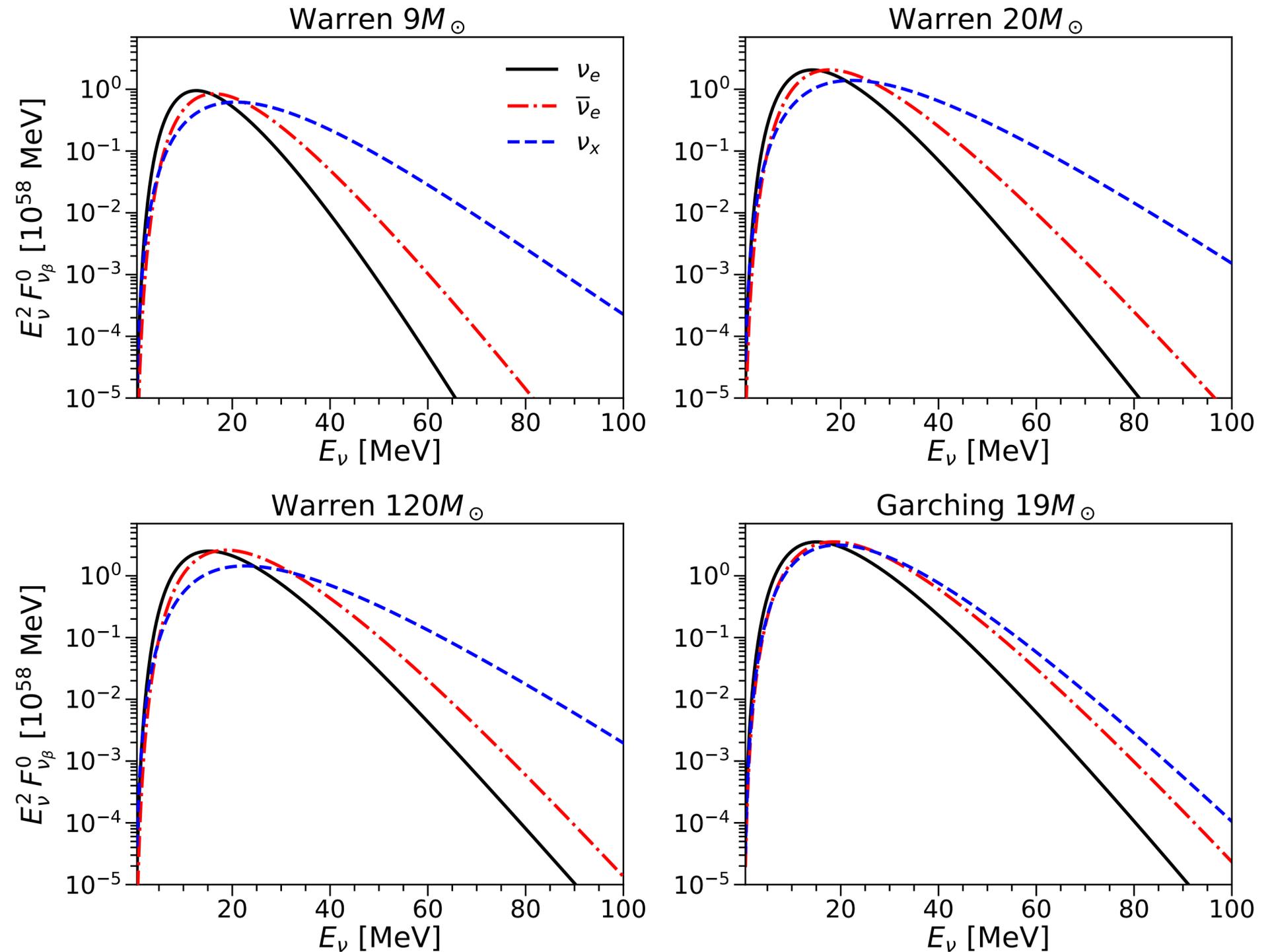


2 4 6 8 10 12
time (s)



Supernova neutrinos

- Initial fluxes from different simulations and masses:
 - **Garching:** R.Bollig et. al. *Astrophys.J.* 915 (2021) 1, 28
 - **Warren:** M.L. Warren et. al. *Astrophys.J.* 898 (2020) 2, 139



Supernova neutrinos

- Adiabatic transitions make neutrinos go out from the SN as mass eigenstates.

NO

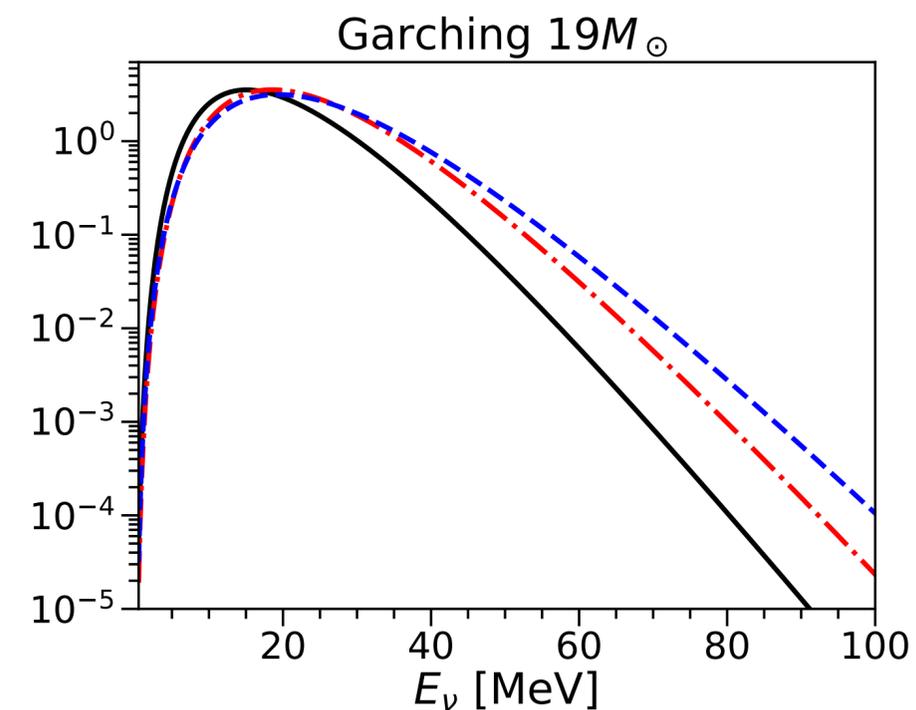
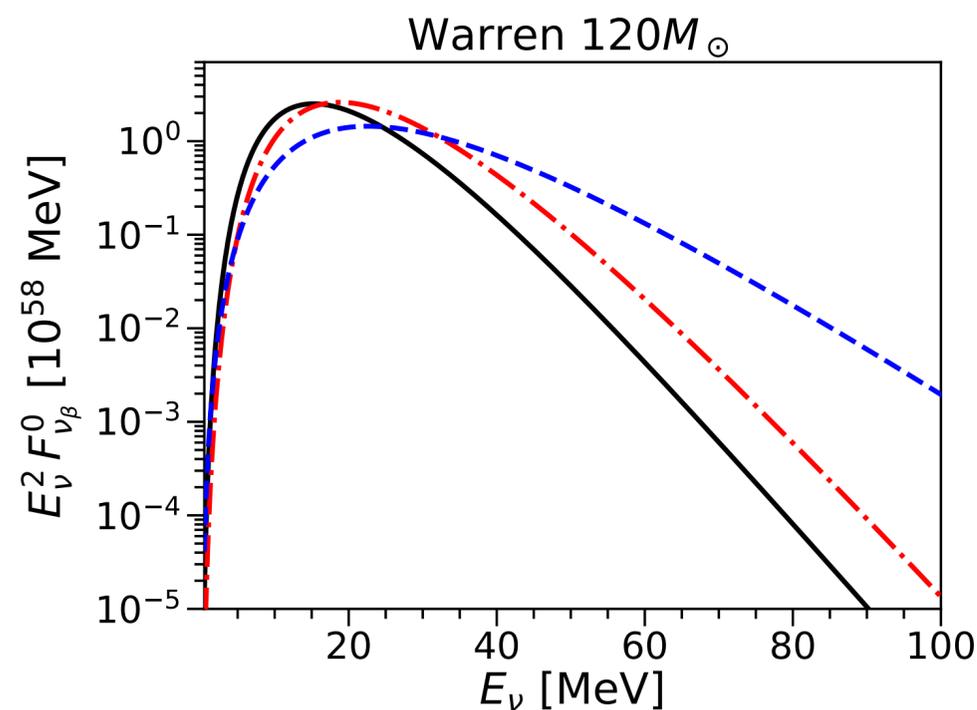
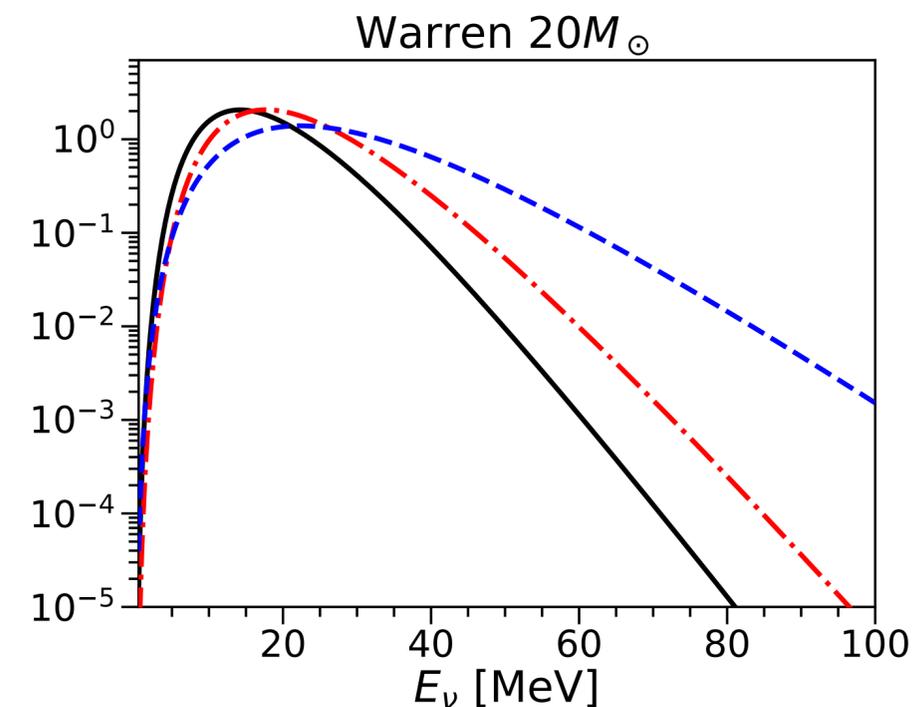
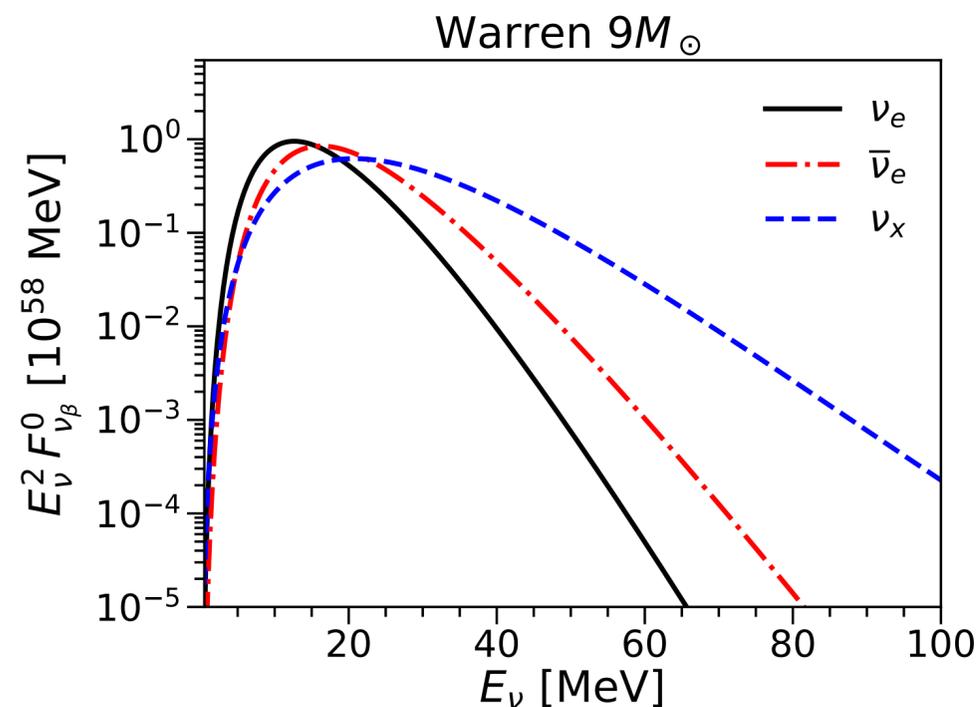
IO

$$F_{\nu_3}^0 = F_{\nu_e}^0$$

$$F_{\bar{\nu}_1}^0 = F_{\bar{\nu}_e}^0$$

$$F_{\nu_2}^0 = F_{\nu_e}^0$$

$$F_{\bar{\nu}_3}^0 = F_{\bar{\nu}_e}^0$$

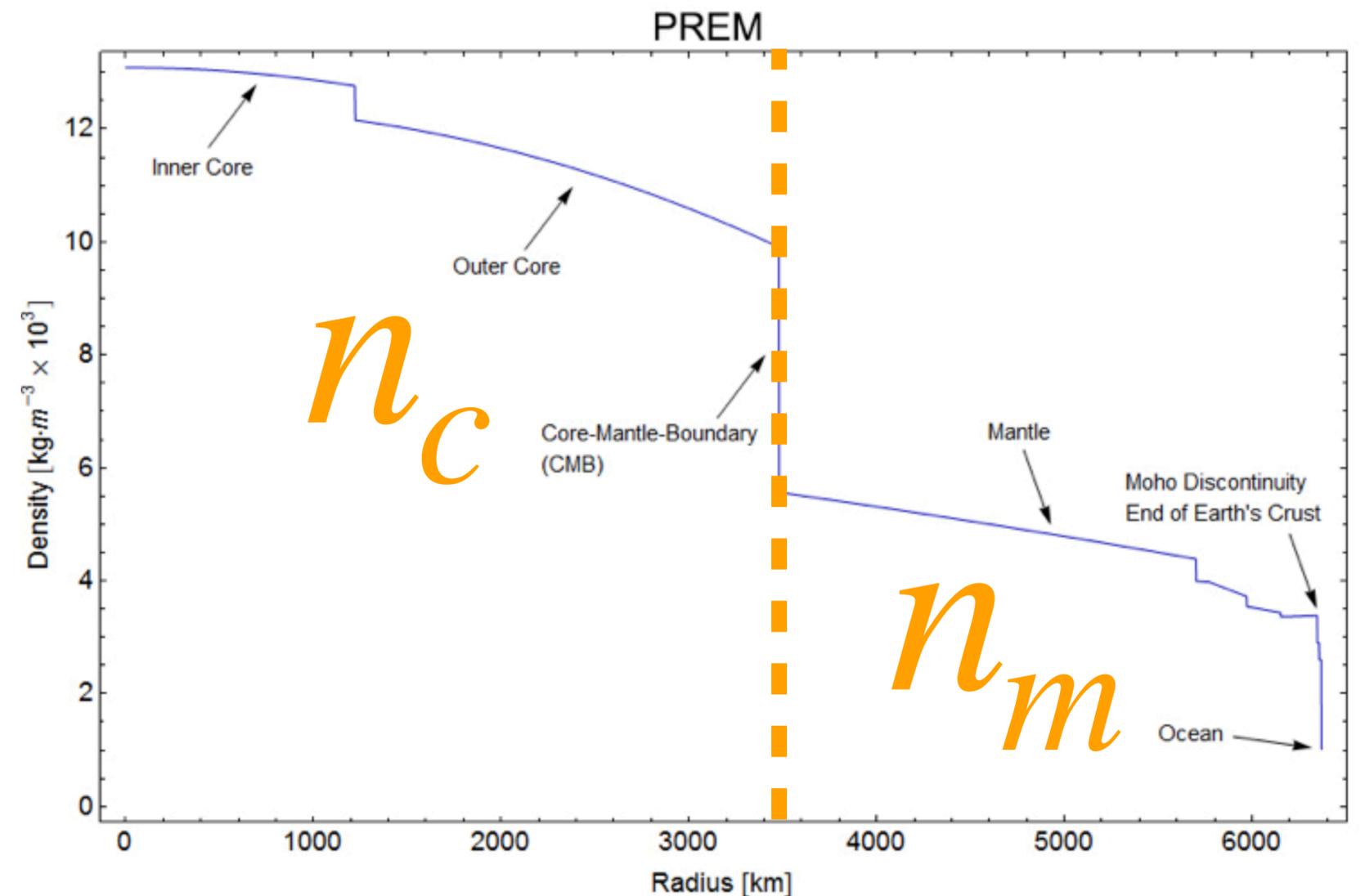


Earth matter effects

$$\rho_{\oplus}(n_c) = \begin{cases} n_c \rho^{\text{PREM}}(r) & , \quad 0 \leq r \leq R_c , \\ n_m \rho^{\text{PREM}}(r) & , \quad R_c < r \leq R_{\oplus} , \end{cases}$$

2-layer profile:
PREM normalized

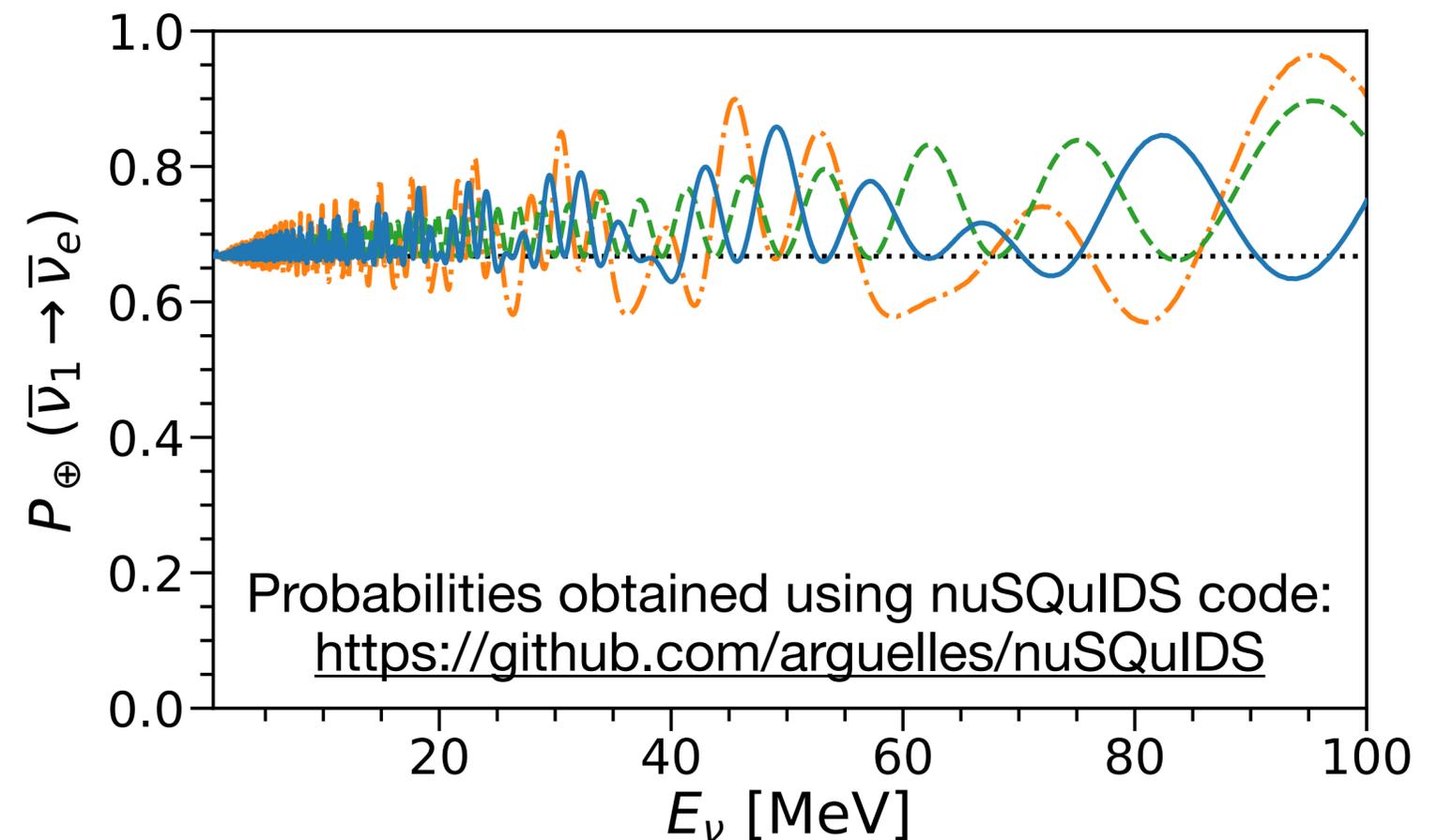
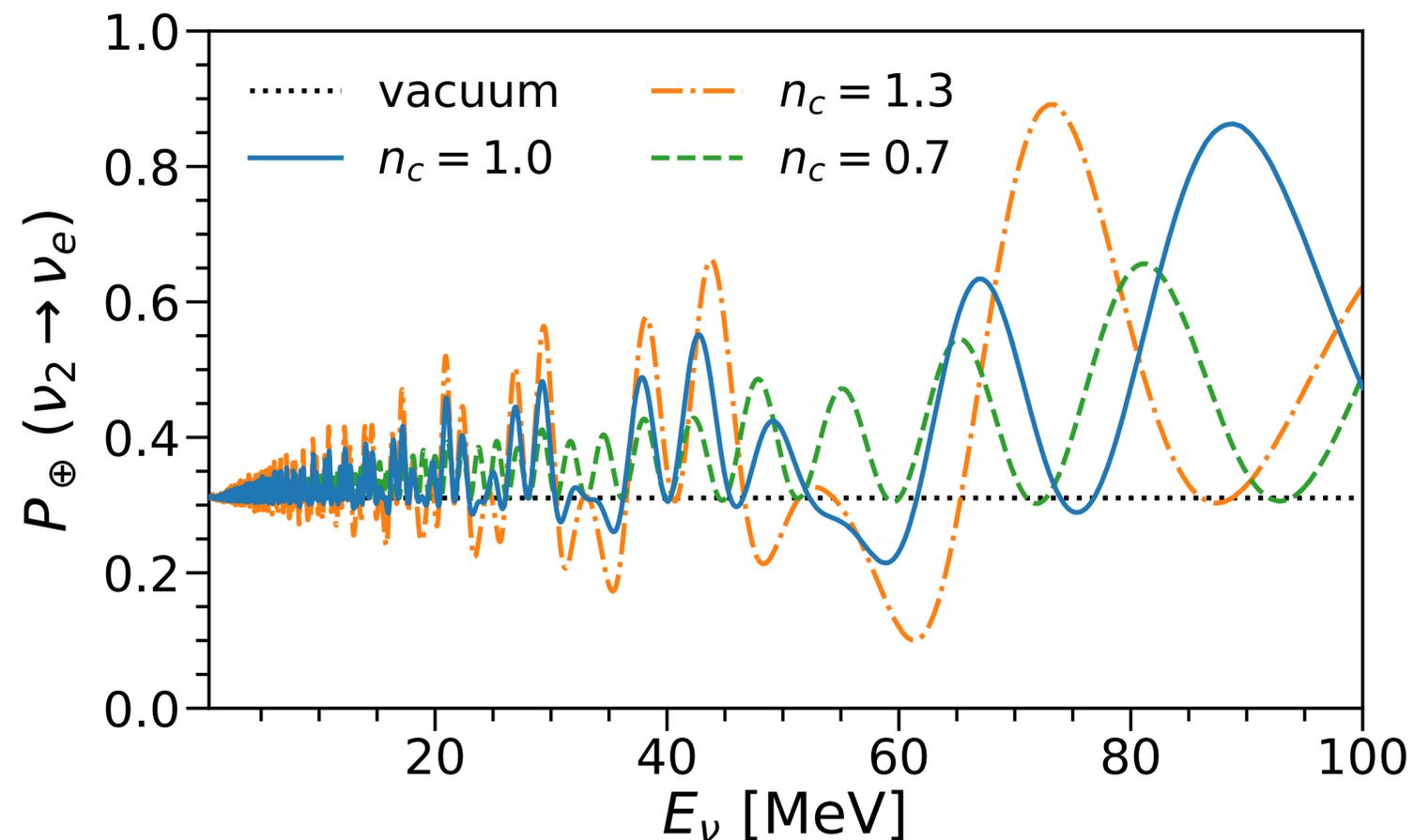
- PREM model with only one free parameter: core norm (n_c)
- We fix mantle norm (n_m) imposing we know the mass of the Earth
- Density also depends on the incident direction ($\cos \theta_z \equiv c_z$)



Earth matter effects

$$F_{\nu_e}^D = p F_{\nu_e}^0 + (1 - p) F_{\nu_x}^0$$

- We aim to be sensitive to these changes on the probabilities.
- Neutrinos have maximal matter effects at $E \sim (40-100)$ MeV driven by the solar mass squared difference: $\Delta m_{21}^2 = 7.5 \cdot 10^{-5} \text{ eV}^2$ (direction dependent).



Future detection of Supernova neutrinos

$$\frac{dR(E_{\text{rec}})}{dE_{\text{rec}}} = N_t \int dE_{\text{true}} dE_\nu \varepsilon(E_{\text{true}}) \mathcal{R}(E_{\text{true}}, E_{\text{rec}}) \frac{d\Phi_\nu^D(E_\nu)}{dE_\nu} \frac{d\sigma(E_\nu, E_e)}{dE_e}$$
$$\mathcal{R}(E_{\text{true}}, E_{\text{rec}}) = \frac{1}{\sqrt{2\pi} \sigma_{\text{det}}} \exp\left(-\frac{(E_{\text{true}} - E_{\text{rec}})^2}{2\sigma_{\text{det}}^2}\right)$$

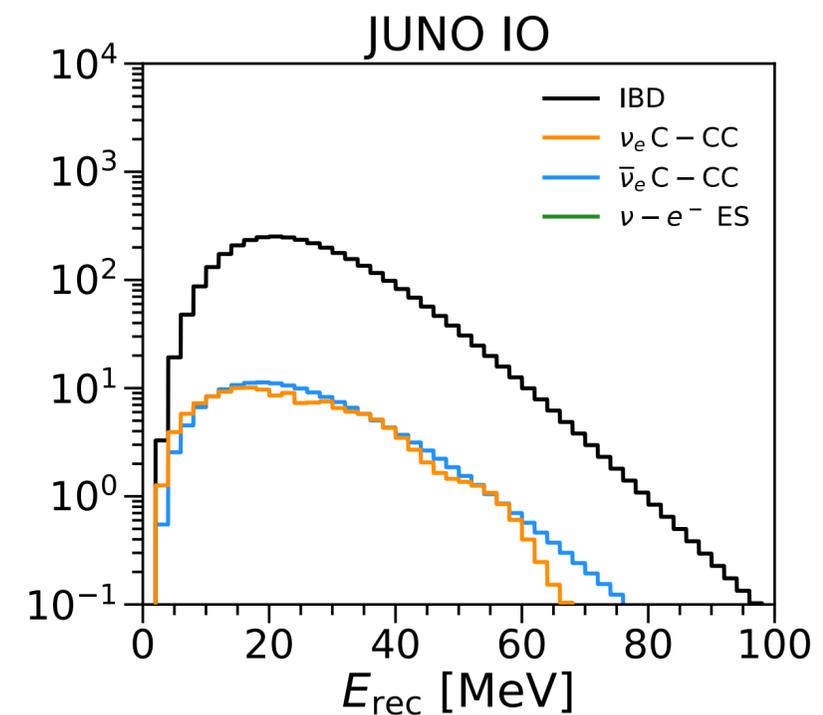
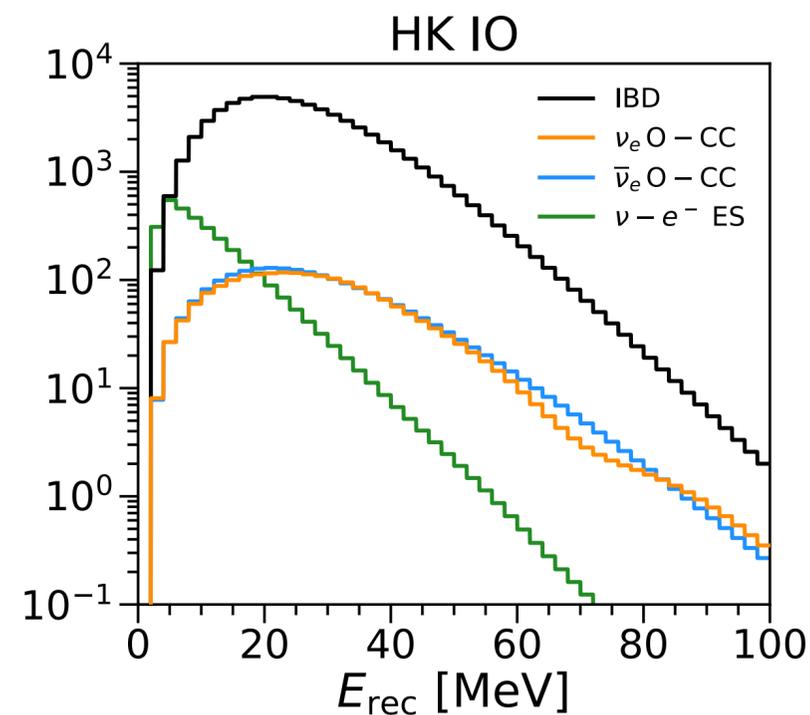
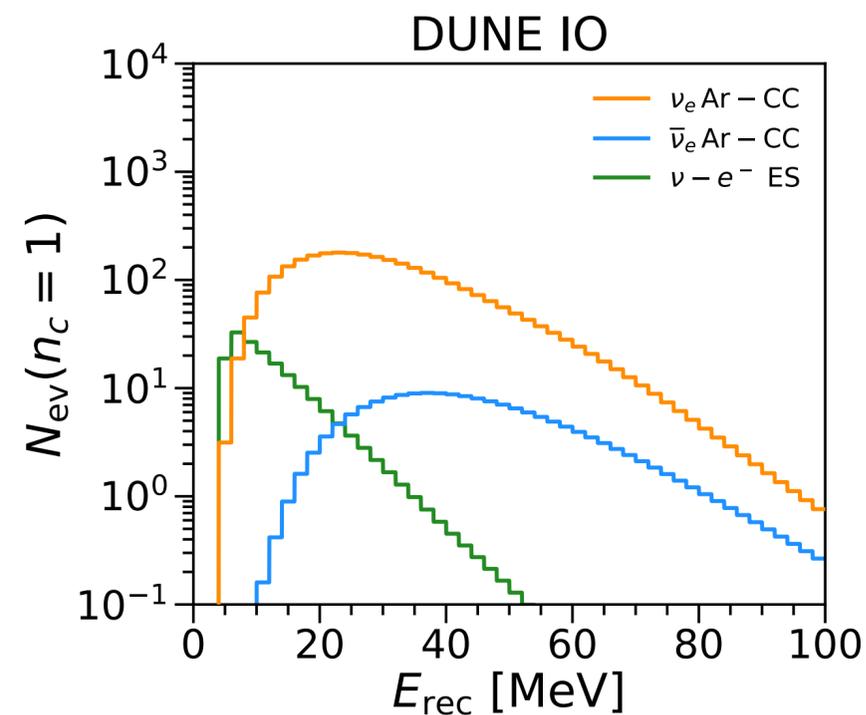
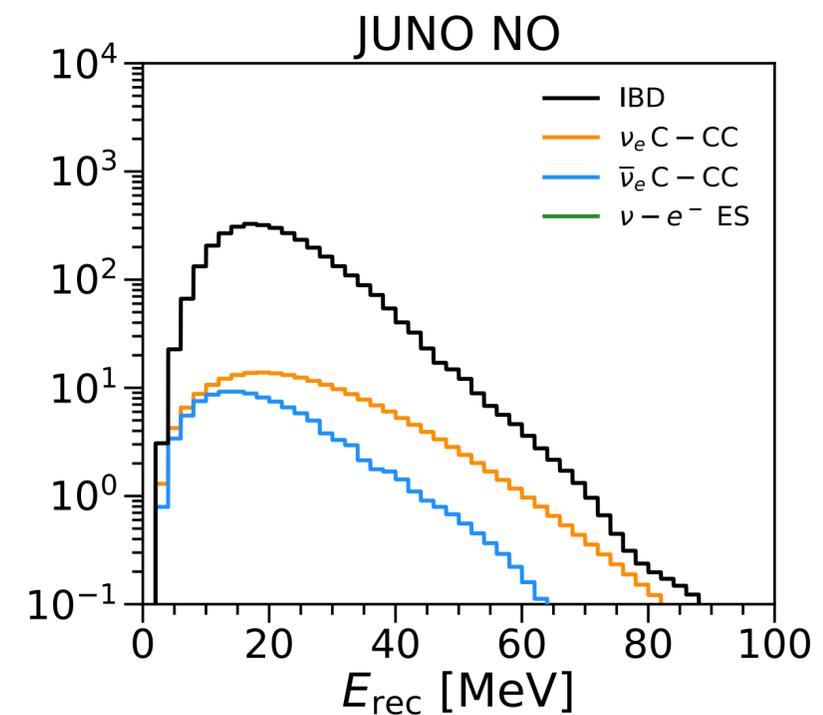
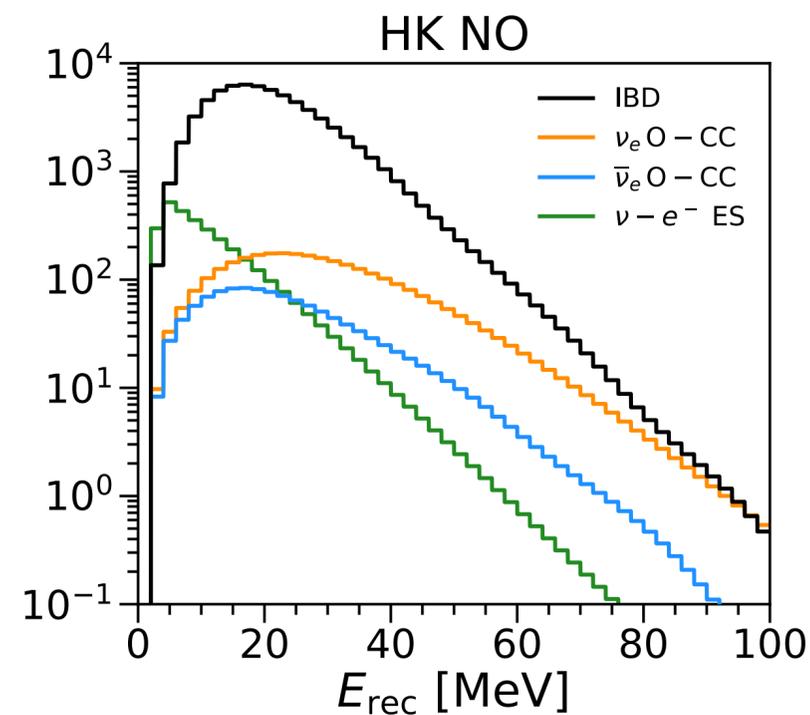
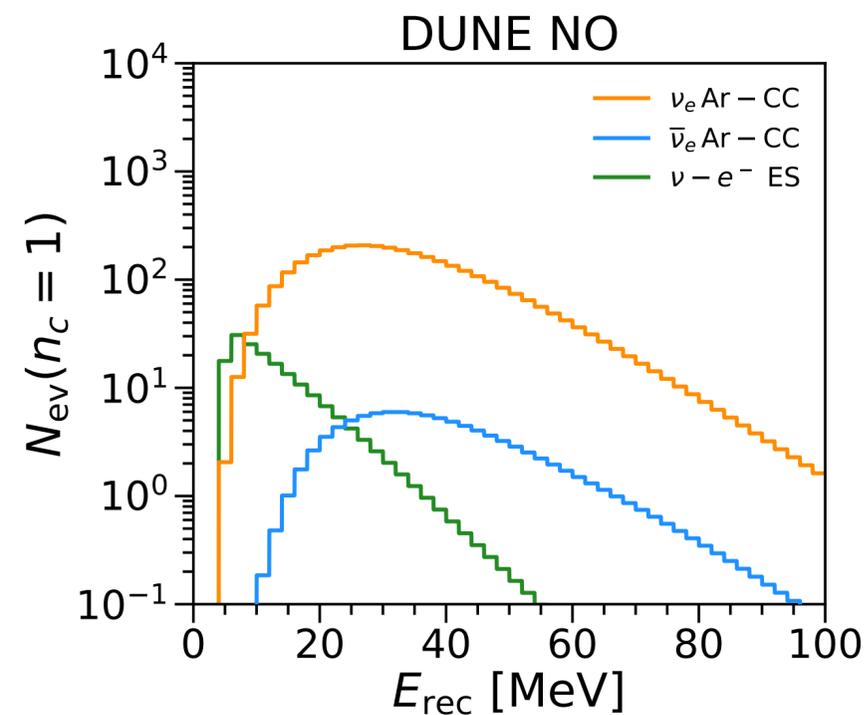
DUNE (LIQUID ARGON)	HK (WATER CHERENKOV)	JUNO (LIQUID SCINTILLATOR)
$\nu_e \text{Ar} - \text{CC} : \nu_e + {}^{40}\text{Ar} \rightarrow e^- + \text{X} ,$ $\bar{\nu}_e \text{Ar} - \text{CC} : \bar{\nu}_e + {}^{40}\text{Ar} \rightarrow e^+ + \text{X} ,$ $\nu - e^- \text{ES} : \nu + e^- \rightarrow \nu + e^- .$	$\text{IBD} : \bar{\nu}_e + p \rightarrow e^+ + n ,$ $\nu_e \text{O} - \text{CC} : \nu_e + {}^{16}\text{O} \rightarrow e^- + \text{X} ,$ $\bar{\nu}_e \text{O} - \text{CC} : \bar{\nu}_e + {}^{16}\text{O} \rightarrow e^+ + \text{X} ,$ $\nu - e^- \text{ES} : \nu + e^- \rightarrow \nu + e^- .$	$\text{IBD} : \bar{\nu}_e + p \rightarrow e^+ + n ,$ $\nu_e \text{C} - \text{CC} : \nu_e + {}^{12}\text{C} \rightarrow e^- + \text{X} ,$ $\bar{\nu}_e \text{C} - \text{CC} : \bar{\nu}_e + {}^{12}\text{C} \rightarrow e^+ + \text{X} ,$ $\nu - e^- \text{ES} : \nu + e^- \rightarrow \nu + e^- .$
$N_t^{\text{Ar}} = 6.03 \cdot 10^{32}$ 20% ENERGY RESOLUTION	$N_t^{\text{P}} = 2.94 \cdot 10^{34}$ MEDIUM ENERGY RESOLUTION	$N_t^{\text{P}} = 1.47 \cdot 10^{33}$ GOOD ENERGY RESOLUTION

Future detection of Supernova neutrinos

- Warren20,
 $c_z = -1$,
 $d_{\text{SN}} = 10 \text{ kpc}$

**NO:
effect in
antineutrinos**

**IO:
effect in
neutrinos**



Sensitivity to Earth density profile

- Binned poissonian χ^2 distribution with 1 degree of freedom (n_c):

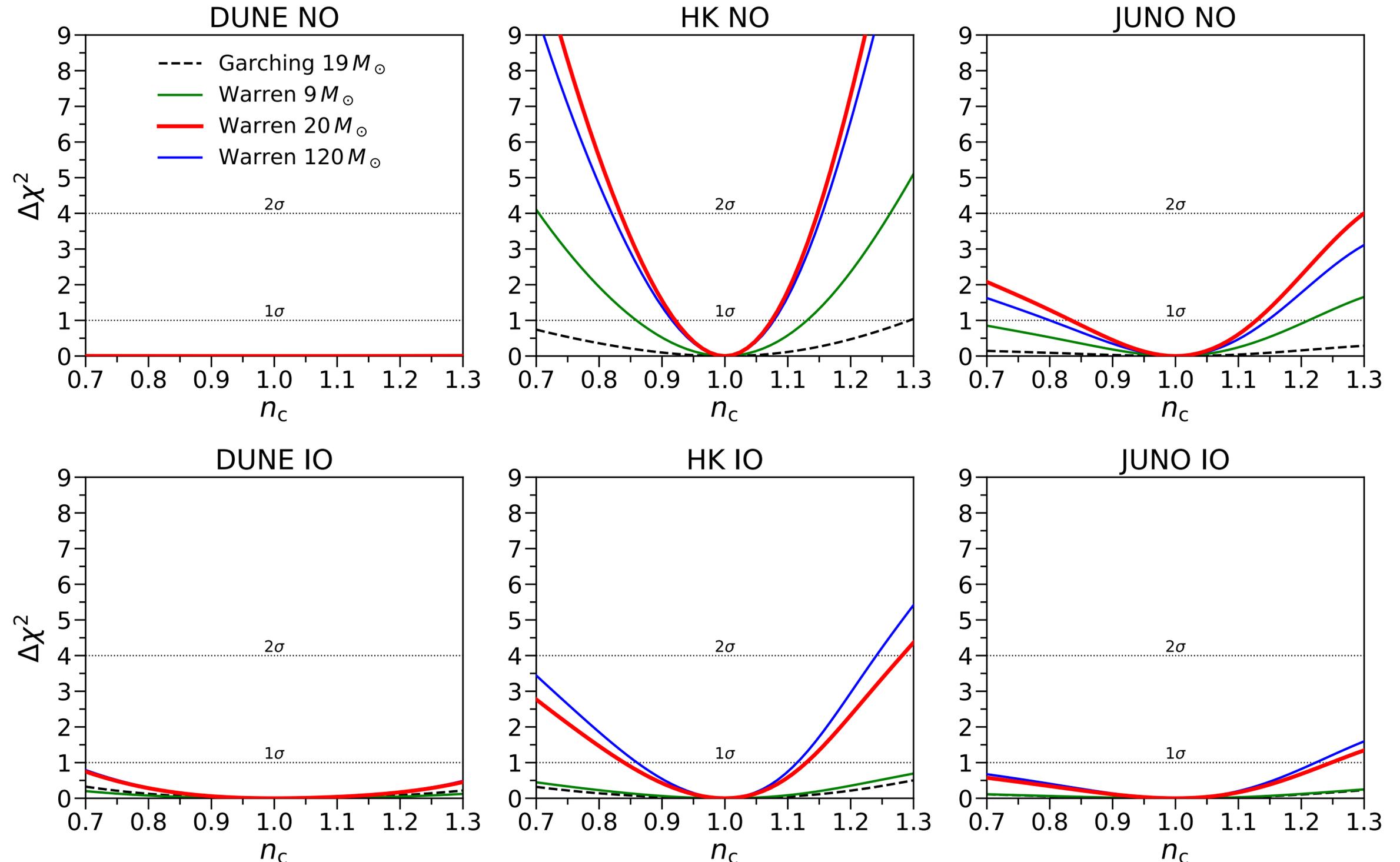
$$\Delta\chi^2(n_c; c_z) = 2 \sum_{i,s} \left[N_{i,s}(n_c; c_z) - N_{i,s}(n_c = 1; c_z) + N_{i,s}(n_c = 1; c_z) \ln \left(\frac{N_{i,s}(n_c = 1; c_z)}{N_{i,s}(n_c; c_z)} \right) \right]$$

- Detection channels of each detector (different topologies):

DUNE (LIQUID ARGON)	HK (WATER CHERENKOV)	JUNO (LIQUID SCINTILLATOR)
$\nu_e \text{Ar} - \text{CC} + \bar{\nu}_e \text{Ar} - \text{CC}$ $\nu - e^- \text{ES}$	0.9 IBD 0.1 IBD + $\nu_e \text{O} - \text{CC} +$ $+ \bar{\nu}_e \text{O} - \text{CC} + \nu - e^- \text{ES}$	0.95 IBD 0.05 IBD + $\nu_e \text{O} - \text{CC} +$ $+ \bar{\nu}_e \text{O} - \text{CC} + \nu - e^- \text{ES}$

Sensitivity to Earth density profile

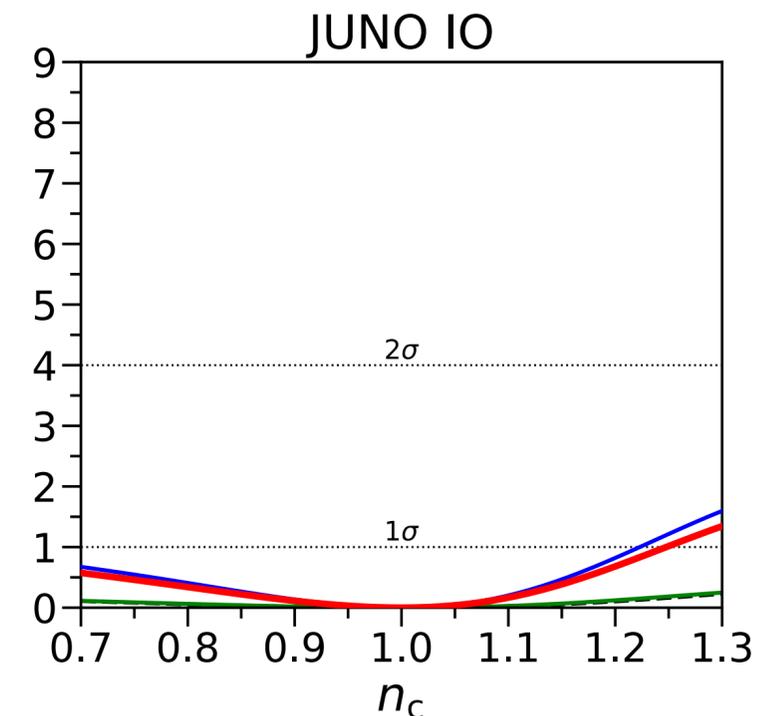
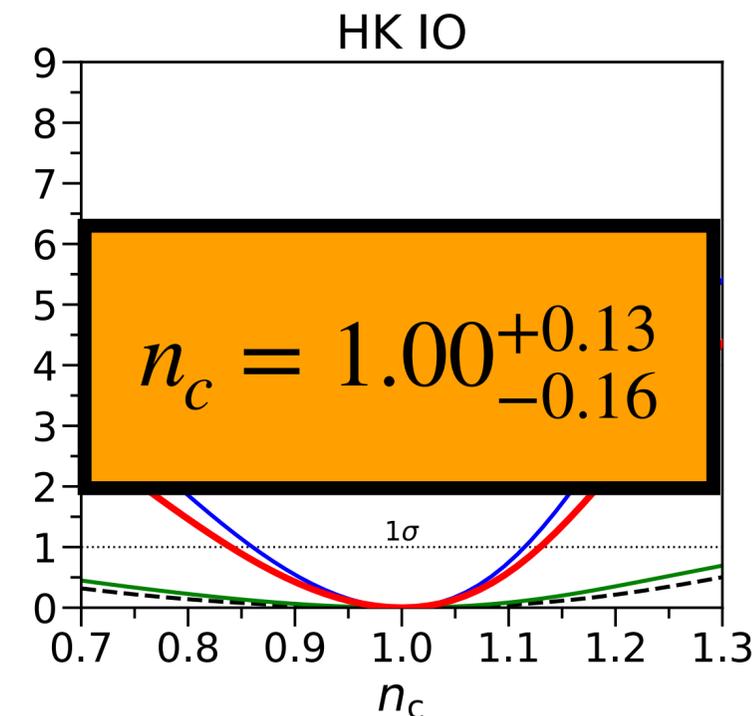
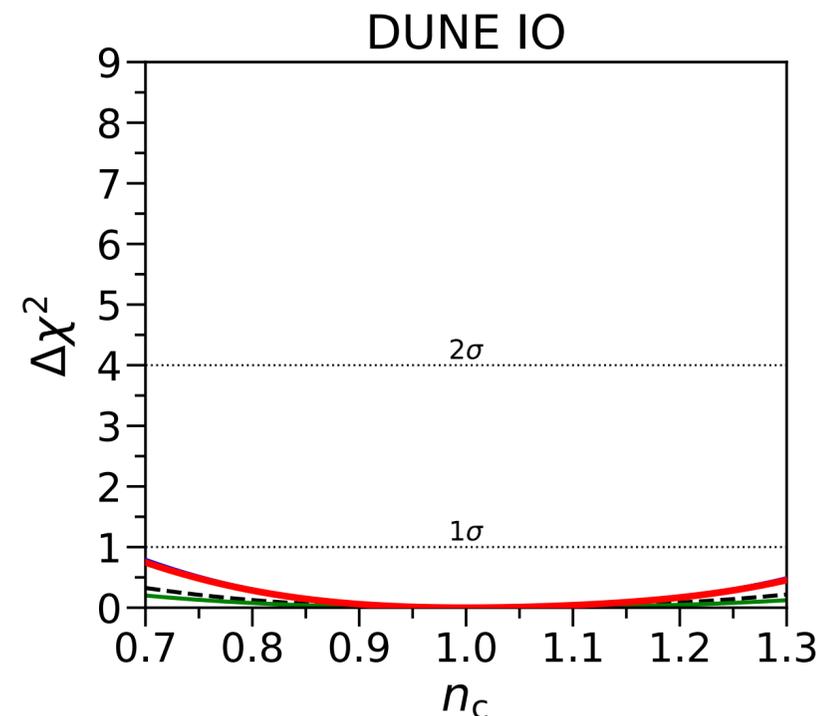
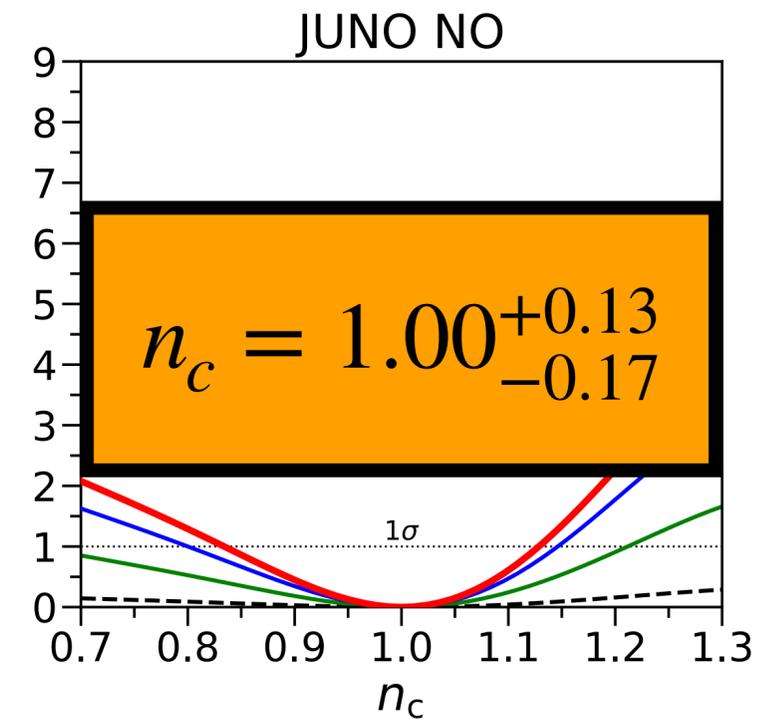
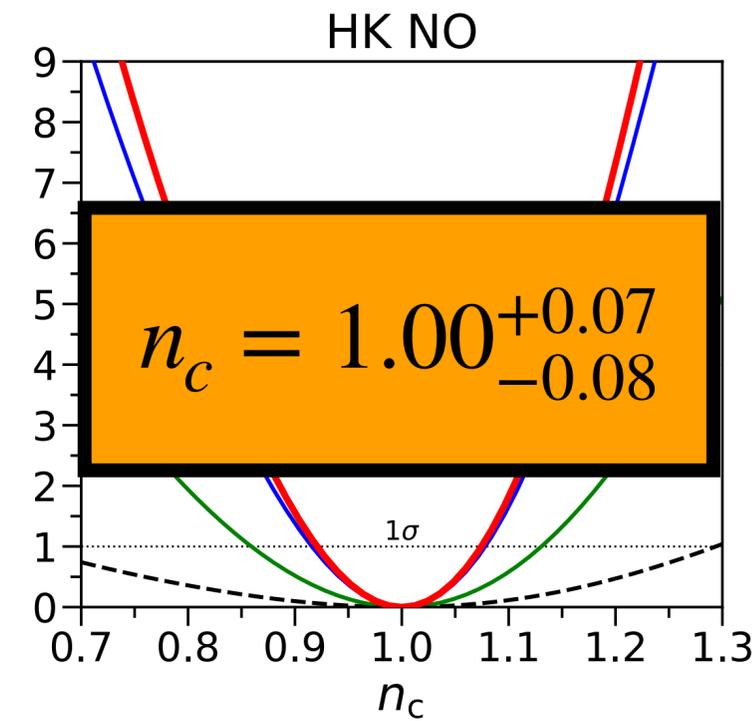
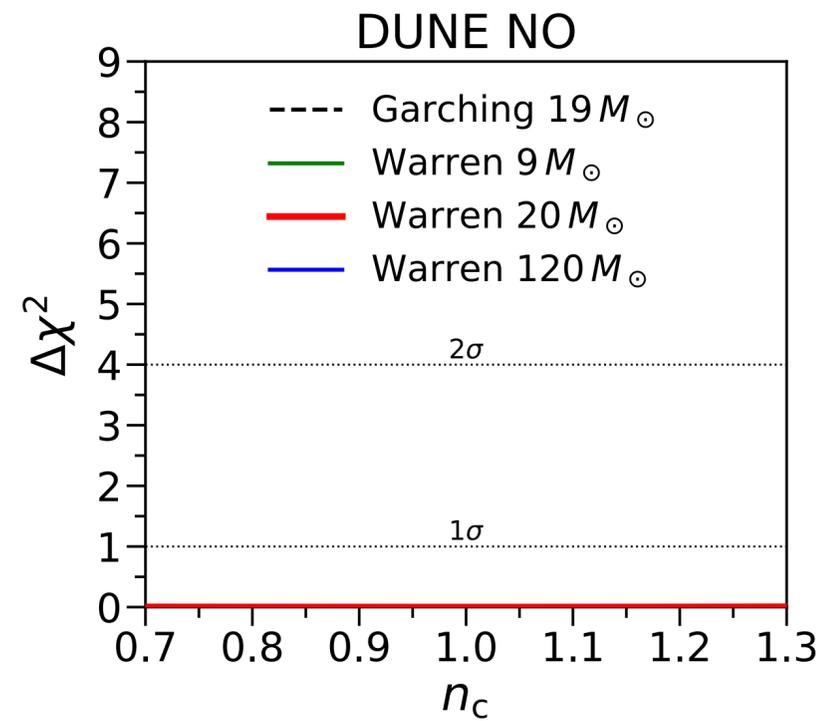
- Sensitivity for different models of SN neutrino burst.
- $c_z = -1$,
 $d_{\text{SN}} = 10$ kpc
- Very model dependent!
- HK will be the detector providing the best results.



Sensitivity to Earth density profile

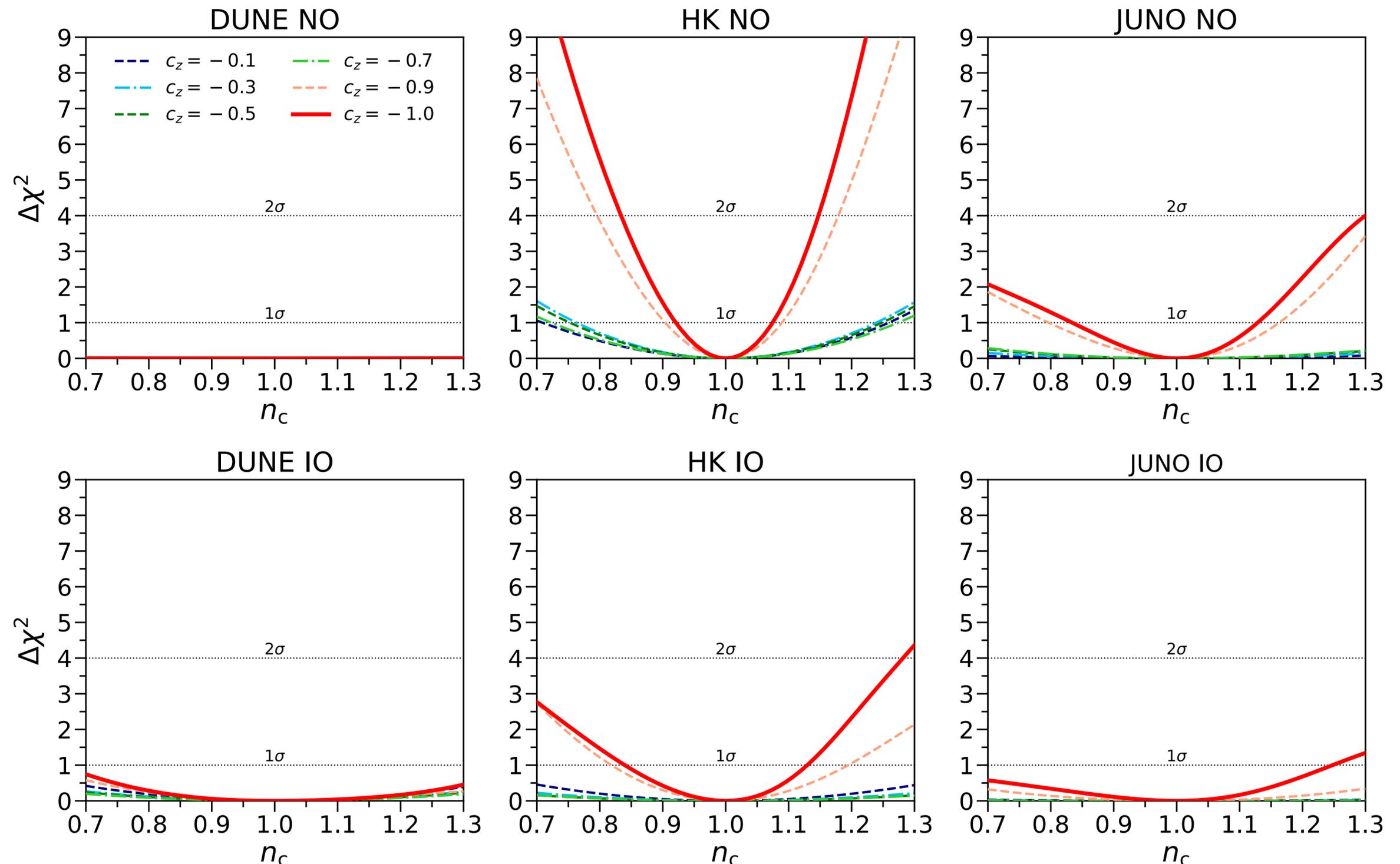
Warren $20M_{\odot}$

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Sensitivity to Earth density profile

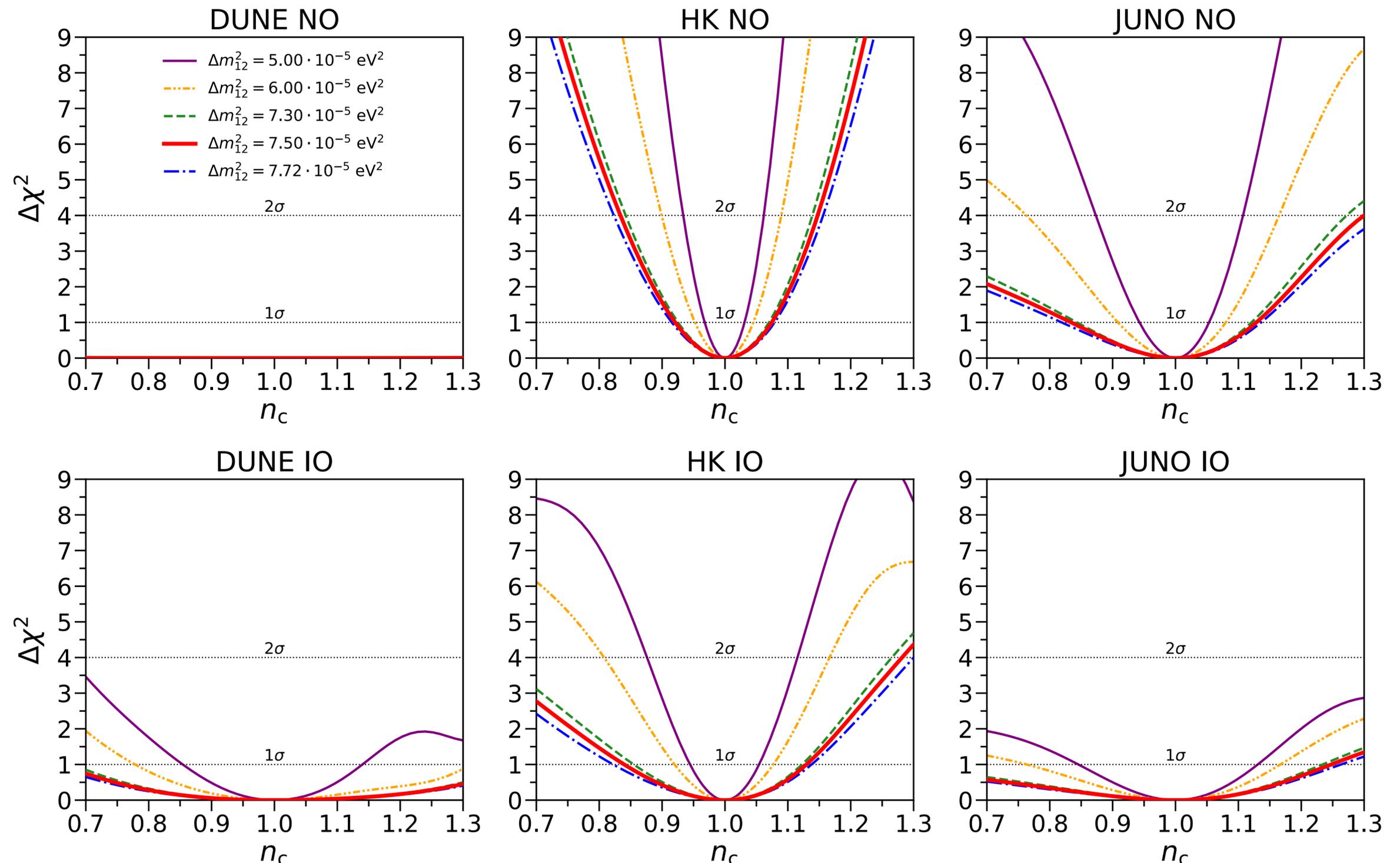
- Effect of incident direction of the SN neutrinos
- Warren20, $d_{\text{SN}} = 10$ kpc
- HK and JUNO NO: core trajectories.
- We need luck: SN neutrinos through core!



Sensitivity to Earth density profile

- Effect of solar mass splitting
- Warren20,
 $d_{\text{SN}} = 10$ kpc
- Optimistic result in Lindner et. al. *Astropart. Phys.* 19, 755 (2003) with

$$\Delta m_{21}^2 = 5 \cdot 10^{-5} \text{ eV}^2$$



Conclusions

- We forecast an oscillation tomography of the Earth with SN neutrinos at future neutrino detectors DUNE, HK and JUNO.
- Assuming adiabatic transitions in the SN, Earth matter effects happen mainly in antineutrinos for NO and neutrinos in IO.
- We studied a different set of initial fluxes and incident directions.
- Most optimistic case: HK and JUNO could determine the average Earth's core density within $\lesssim 10\%$ at 1σ CL with galactic SN neutrinos (at 10 kpc).

A future SN burst could aid in future neutrino Earth tomography studies, and be competitive with, and complementary to other analyses.

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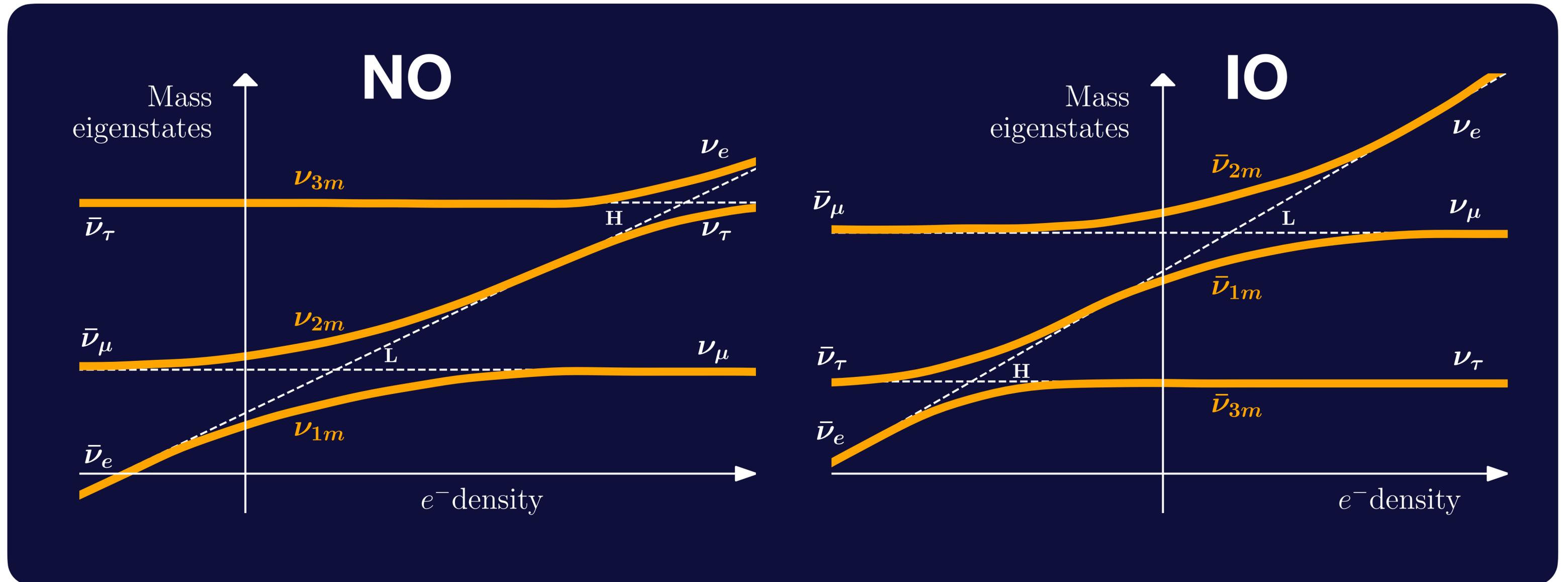
But we need luck!

Earth tomography with supernova neutrinos at future neutrino detectors

BACKUP SLIDES

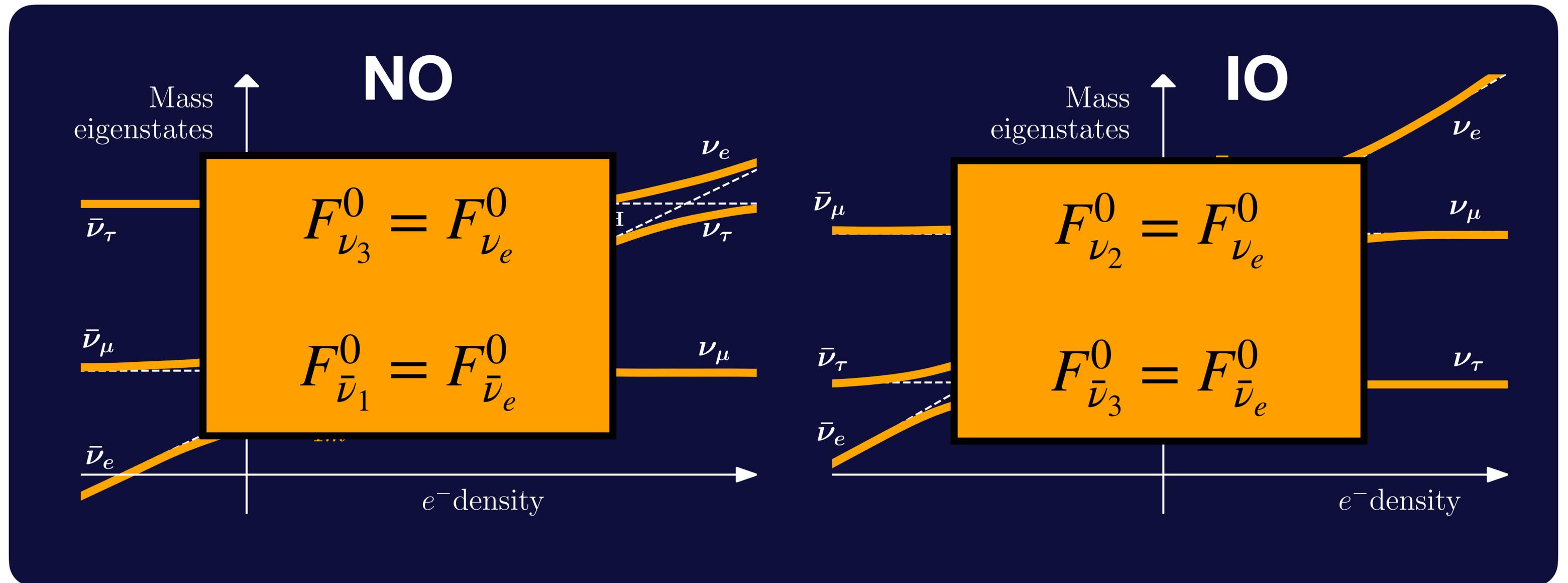
Supernova neutrinos

- Adiabatic transitions make neutrinos go out from the SN as mass eigenstates.



Supernova neutrinos

- Adiabatic transitions make neutrinos go out from the SN as mass eigenstates.



Earth matter effects

- Fluxes at detectors are a combination of fluxes at production:

$$F_{\nu_e}^D = p F_{\nu_e}^0 + (1 - p) F_{\nu_x}^0$$

$$F_{\nu_x}^D = \frac{1 - p}{2} F_{\nu_e}^0 + \frac{1 + p}{2} F_{\nu_x}^0$$

$$F_{\bar{\nu}_e}^D = \bar{p} F_{\bar{\nu}_e}^0 + (1 - \bar{p}) F_{\nu_x}^0$$

$$F_{\bar{\nu}_x}^D = \frac{1 - \bar{p}}{2} F_{\bar{\nu}_e}^0 + \frac{1 + \bar{p}}{2} F_{\nu_x}^0$$

If $F_{\nu_e}^0 = F_{\nu_x}^0$ we are not sensitive to the matter effects!

p and \bar{p} are the probability of transition from an initial mass state that depends on the neutrino mass ordering (SN emission) to a final flavor state that depends on the detection channel

Earth matter effects

- Fluxes at detectors are a combination of fluxes at production:

$$F_{\nu_e}^D = p F_{\nu_e}^0 + (1 - p) F_{\nu_x}^0$$

$$F_{\nu_x}^D = \frac{1 - p}{2} F_{\nu_e}^0 + \frac{1 + p}{2} F_{\nu_x}^0$$

$$F_{\bar{\nu}_e}^D = \bar{p} F_{\bar{\nu}_e}^0 + (1 - \bar{p}) F_{\nu_x}^0$$

$$F_{\bar{\nu}_x}^D = \frac{1 - \bar{p}}{2} F_{\bar{\nu}_e}^0 + \frac{1 + \bar{p}}{2} F_{\nu_x}^0$$

- In order to obtain p we need to know neutrino evolution:

$$\mathcal{H}_{\text{flavor}} = \underbrace{\frac{1}{2E} U M^2 U^\dagger}_{\text{Vacuum}} + \underbrace{V}_{\text{Matter}}$$

$$M^2 = \begin{pmatrix} 0 & 0 & 0 \\ 0 & \Delta m_{21}^2 & 0 \\ 0 & 0 & \Delta m_{31}^2 \end{pmatrix}$$

$$V = \begin{pmatrix} V(n_e) & 0 & 0 \\ 0 & 0 & 0 \\ 0 & 0 & 0 \end{pmatrix}$$

$$U = U_{23} \Gamma_\delta U_{13} U_{12}$$

PMNS matrix

Earth matter effects

- Fluxes at detectors are a combination of fluxes at production:

$$F_{\nu_e}^D = p F_{\nu_e}^0 + (1 - p) F_{\nu_x}^0$$

$$F_{\nu_x}^D = \frac{1 - p}{2} F_{\nu_e}^0 + \frac{1 + p}{2} F_{\nu_x}^0$$

$$F_{\bar{\nu}_e}^D = \bar{p} F_{\bar{\nu}_e}^0 + (1 - \bar{p}) F_{\nu_x}^0$$

$$F_{\bar{\nu}_x}^D = \frac{1 - \bar{p}}{2} F_{\bar{\nu}_e}^0 + \frac{1 + \bar{p}}{2} F_{\nu_x}^0$$

Vacuum probabilities

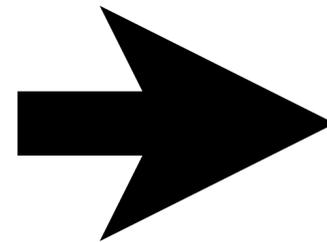
$$p_{\text{vac}}^{\text{NO}} \equiv P_{\text{vac}}(\nu_3 \rightarrow \nu_e) = |U_{e3}|^2 = \sin^2 \theta_{13}$$

$$\bar{p}_{\text{vac}}^{\text{NO}} \equiv P_{\text{vac}}(\bar{\nu}_1 \rightarrow \bar{\nu}_e) = |U_{e1}|^2 = \cos^2 \theta_{12} \cos^2 \theta_{13}$$

$$p_{\text{vac}}^{\text{IO}} \equiv P_{\text{vac}}(\nu_2 \rightarrow \nu_e) = |U_{e2}|^2 = \sin^2 \theta_{12} \cos^2 \theta_{13}$$

$$\bar{p}_{\text{vac}}^{\text{IO}} \equiv P_{\text{vac}}(\bar{\nu}_3 \rightarrow \bar{\nu}_e) = |U_{e3}|^2 = \sin^2 \theta_{13} \quad .$$

$V \neq 0$



Constant density probabilities

$$p_{\oplus}^{\text{NO}} \equiv P_{\oplus}(\nu_3 \rightarrow \nu_e) \simeq \sin^2 \theta_{13} \quad \times$$

$$\bar{p}_{\oplus}^{\text{NO}} \equiv P_{\oplus}(\bar{\nu}_1 \rightarrow \bar{\nu}_e) \simeq \cos^2 \theta_{13} (1 - \bar{P}_{\oplus}^{2\nu})$$

$$p_{\oplus}^{\text{IO}} \equiv P_{\oplus}(\nu_2 \rightarrow \nu_e) \simeq \cos^2 \theta_{13} P_{\oplus}^{2\nu}$$

$$\bar{p}_{\oplus}^{\text{IO}} \equiv P_{\oplus}(\bar{\nu}_3 \rightarrow \bar{\nu}_e) \simeq \sin^2 \theta_{13} \quad \times$$

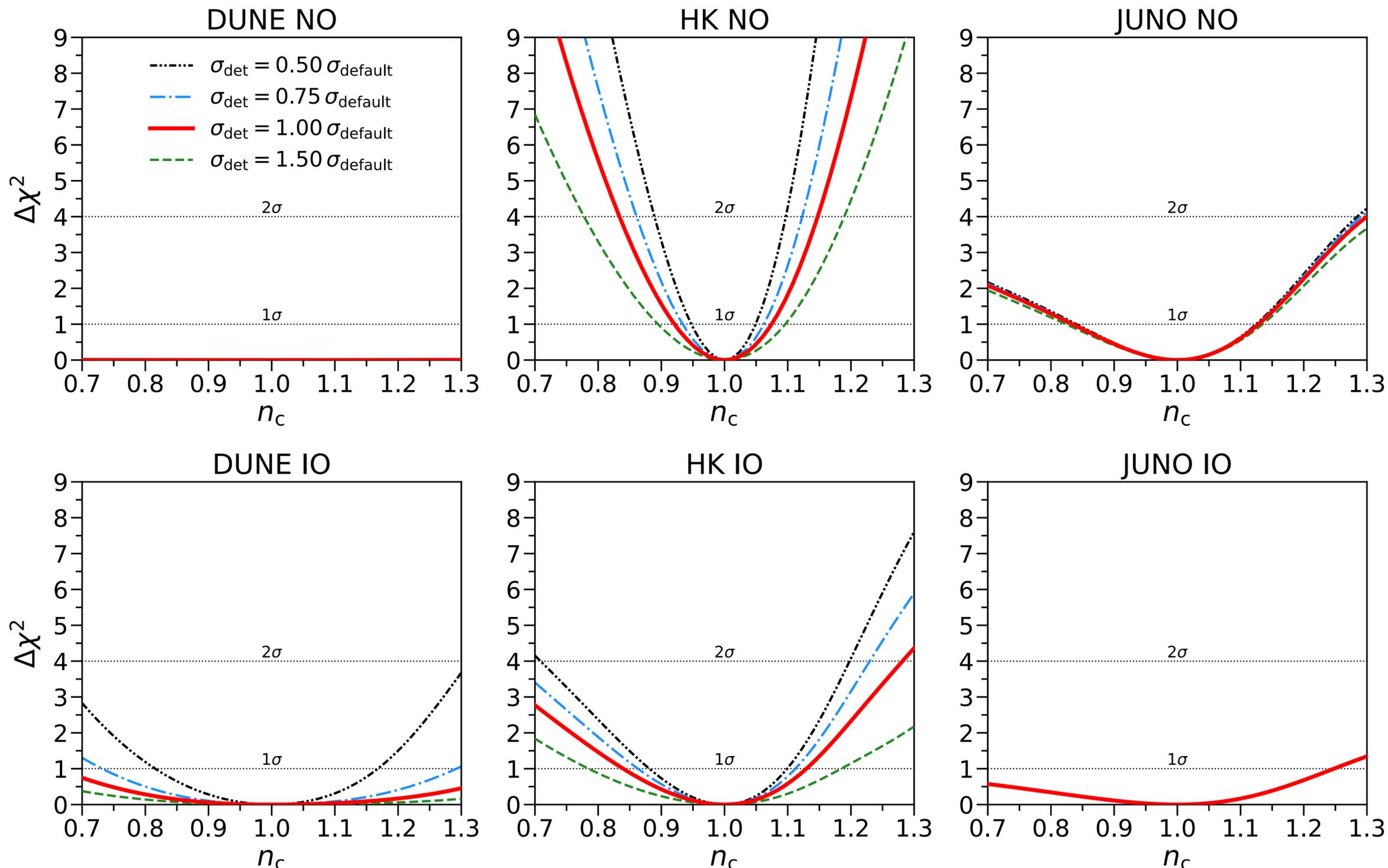
Sensitivity to Earth density profile

- Effect of energy resolution of detectors

- Warren20,
 $c_z = -1$,
 $d_{\text{SN}} = 10$ kpc

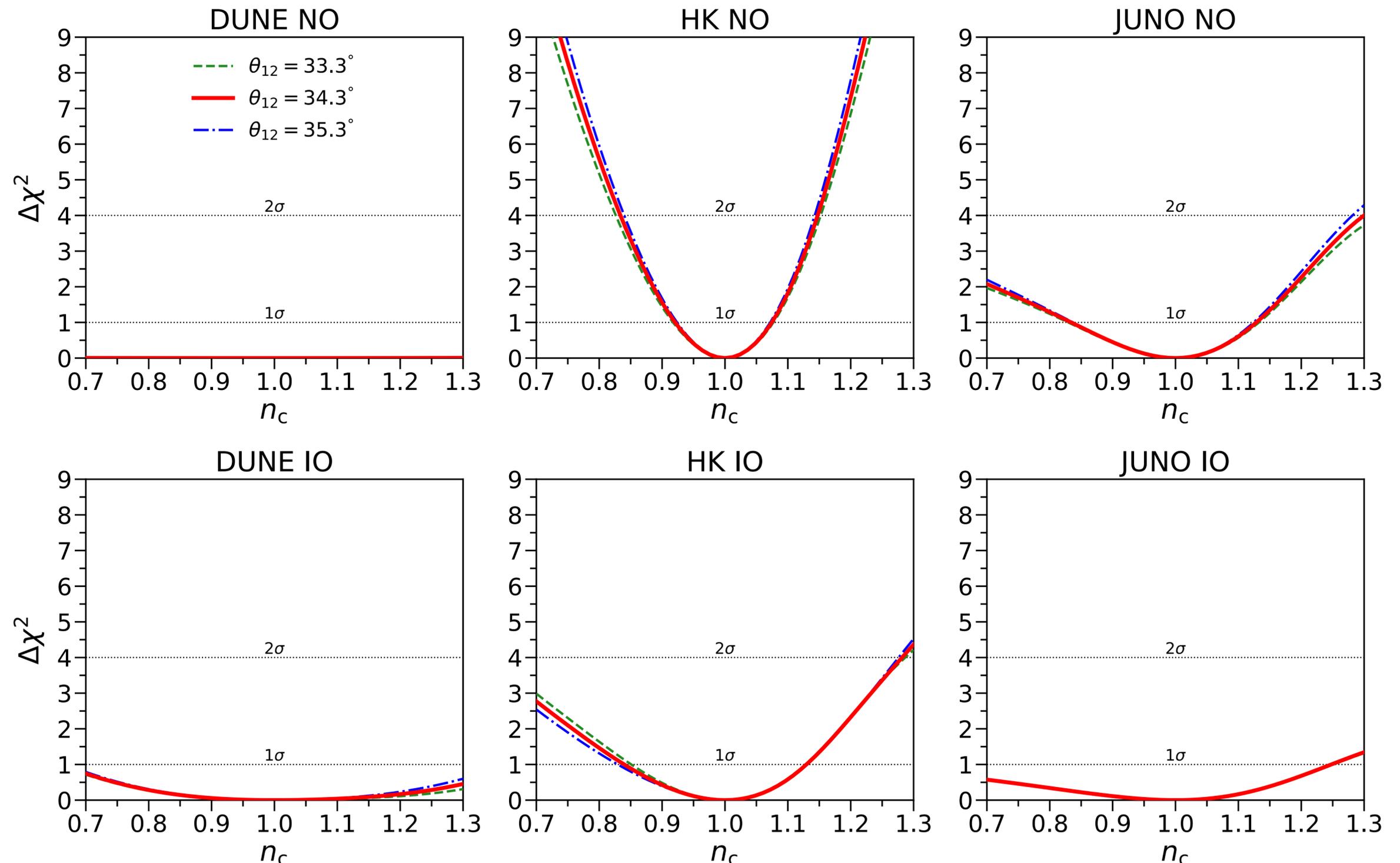
- JUNO has a superb resolution.

- DUNE and HK have huge improvements.



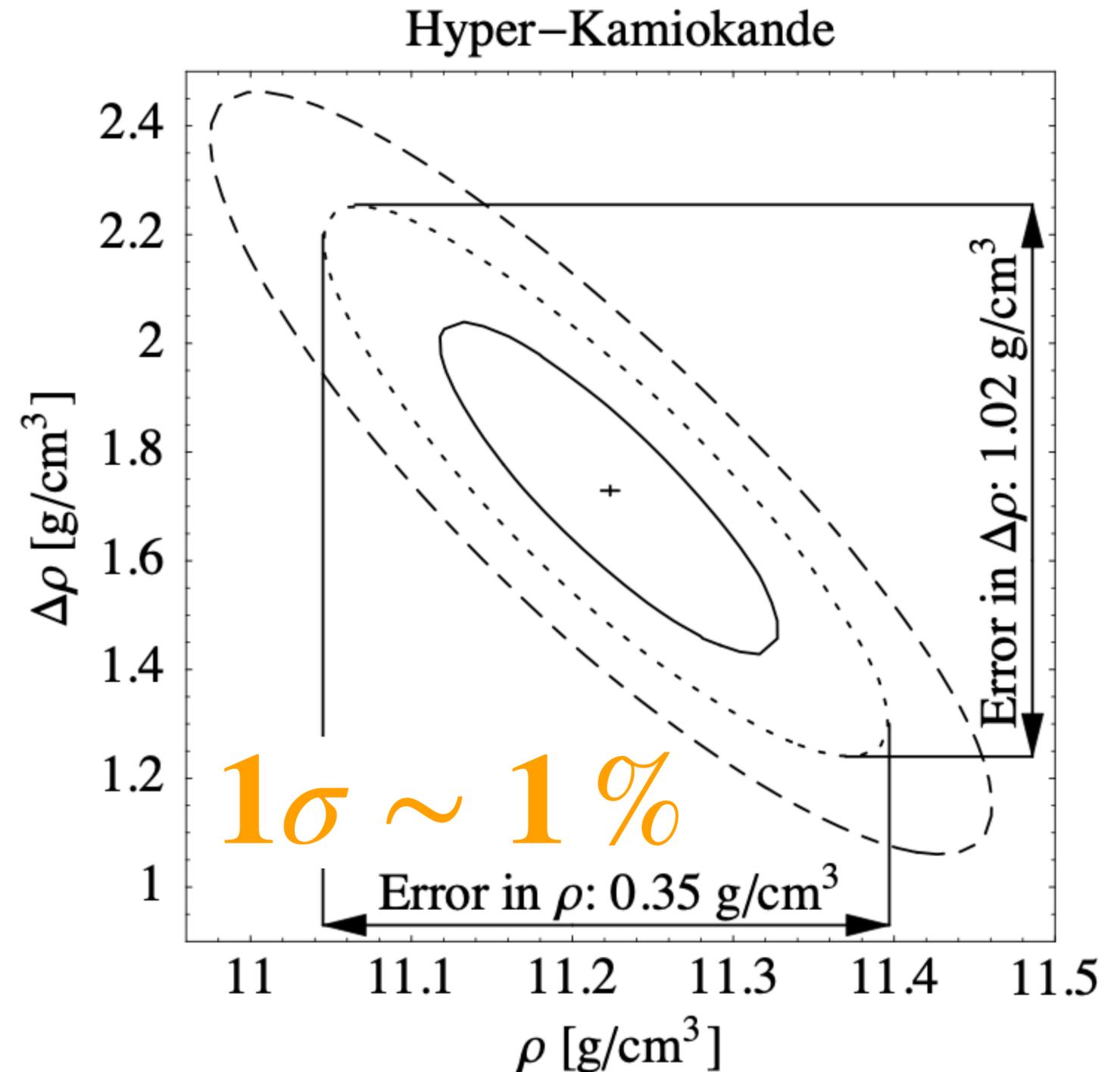
Sensitivity to Earth density profile

- Effect of solar mixing angle.
- Warren20,
 $d_{\text{SN}} = 10$ kpc
- $\pm 1\sigma$ errors in de Salas, Forero, Gariazzo, Martinez-Mirave, Mena, Ternes, Tortola, Valle *JHEP* 02 (2021) 071



Previous SN neutrino tomography

- M. Lindner, T. Ohlsson, R. Tomas, W. Winter, *Astropart.Phys.* 19 (2003) 755-770
- $\Delta m_{21}^2 = 5 \cdot 10^{-5} \text{ eV}^2$
- Fluxes shifted to higher Energy compared with current simulations.



Earth inner structure

Slide taken from
S. Palomares-
Ruiz talk

https://drive.google.com/file/d/1yaF_Wo9SiDd92Ub34ji0Os5YyIpc9JHU/view

- Estimated uncertainties of the different layers of the PREM model

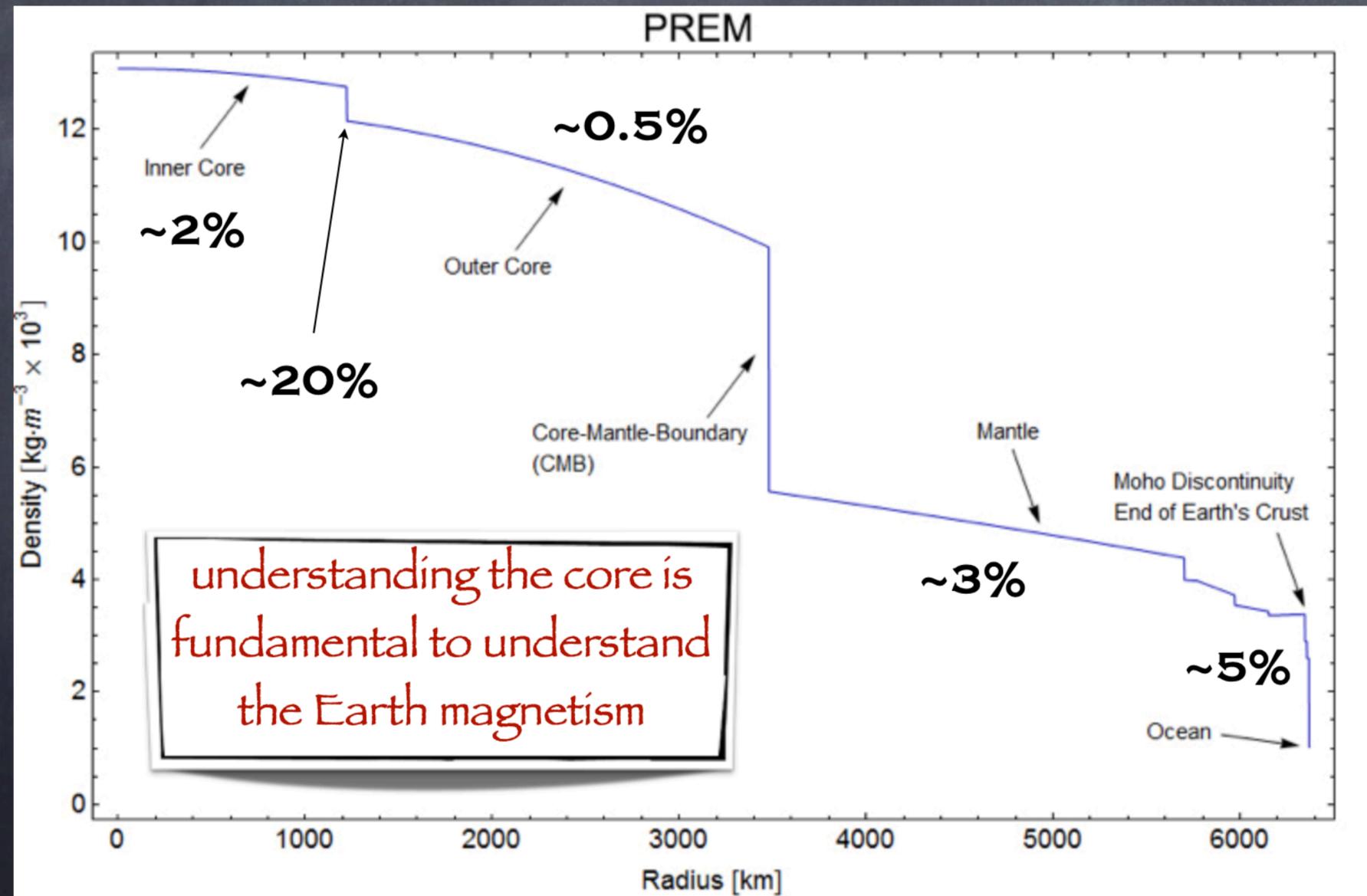


Sergio Palomares-Ruiz

PRELIMINARY REFERENCE EARTH MODEL (PREM)

A. M. Dziewonski and D. L. Anderson, Phys, Earth Planet. Inter. 25:297, 1981

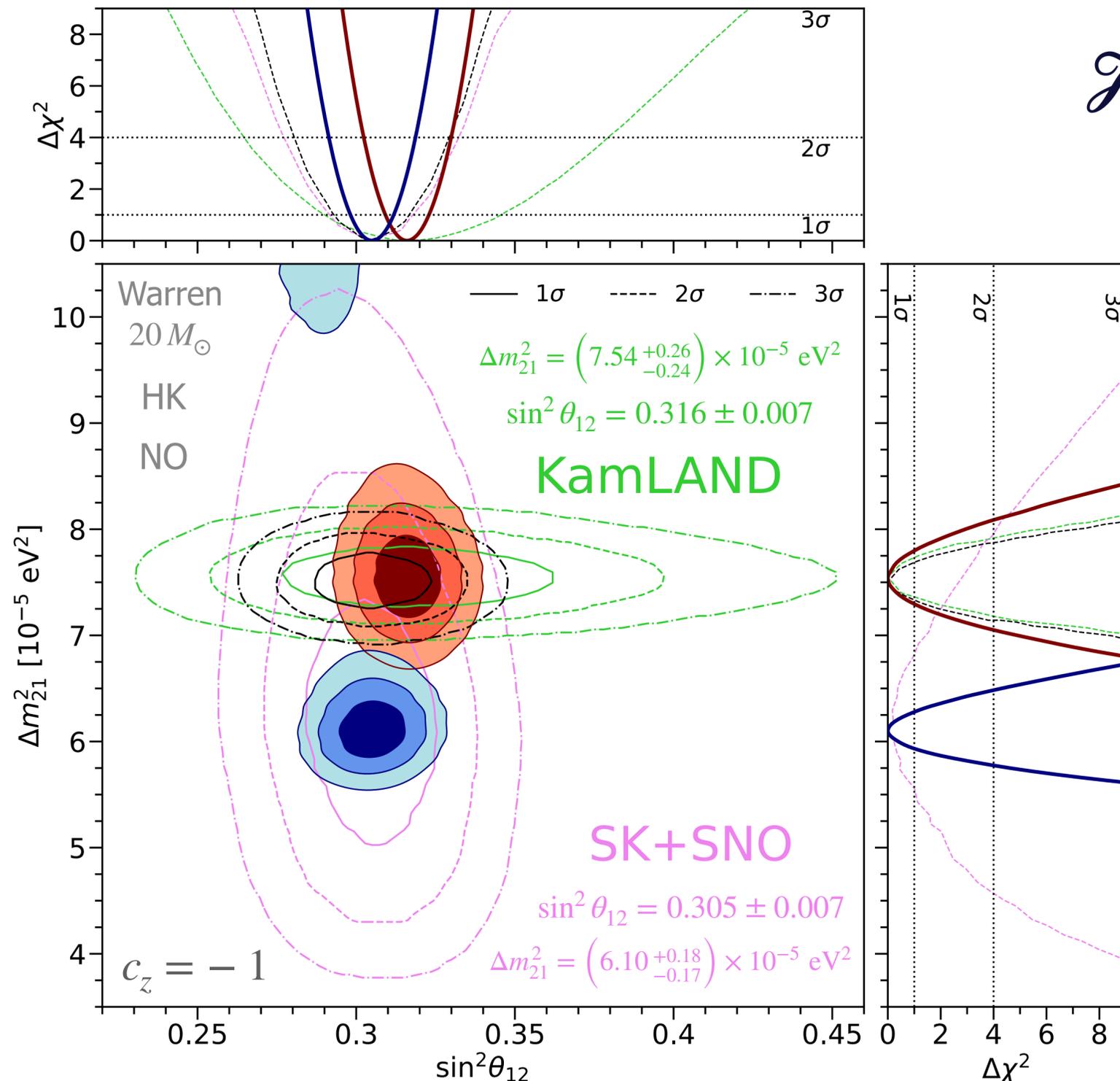
1-D density profile From seismic wave data and imposing the Earth's radius, mass and moment of inertia as additional constraints



Earth tomography with neutrinos

Sensitivity to Δm_{21}^2 with SN neutrinos

RH, O. Mena and
S. Palomares-Ruiz,
arXiv:2307.09509



$$\mathcal{H}_{\text{flavor}} = \frac{1}{2E} U M^2 U^\dagger + \mathbb{V}$$

Vacuum

Matter

$$P_{2\nu}(\nu_\alpha \rightarrow \nu_\beta) = \sin^2 2\theta_m \sin^2 \left(\frac{\Delta_m L}{4E} \right)$$

- Tomography:

$$\Delta_m = \sqrt{(\Delta m^2 \sin 2\theta \mp 2E \mathbb{V})^2 + (\Delta m^2 \cos 2\theta)^2}$$

↑ fixing solar mass squared

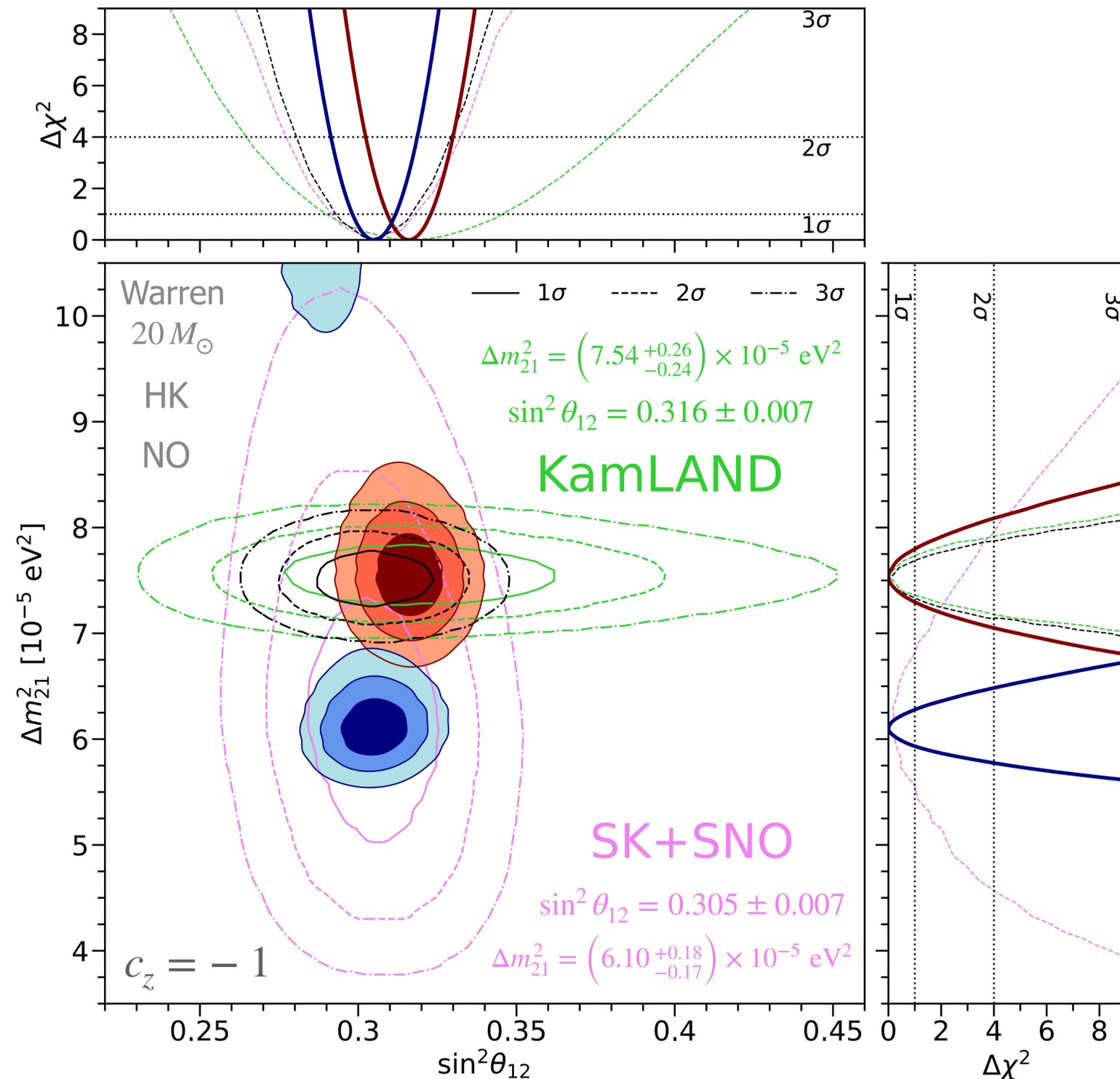
- Solar mass splitting:

$$\Delta_m = \sqrt{(\Delta m^2 \sin 2\theta \mp 2E \mathbb{V})^2 + (\Delta m^2 \cos 2\theta)^2}$$

↑ fixing potential

Sensitivity to Δm_{21}^2 with SN neutrinos

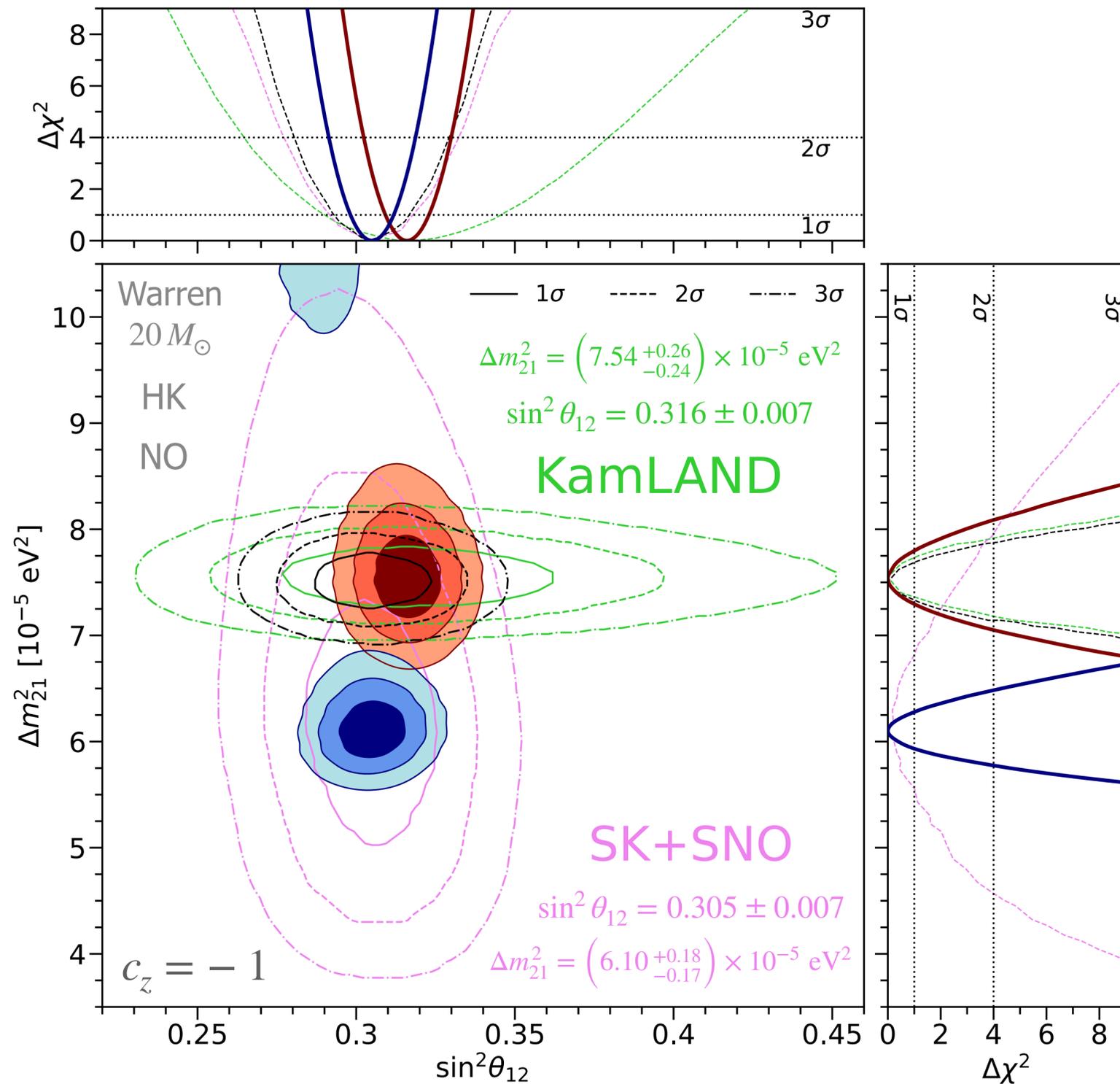
RH, O. Mena and
S. Palomares-Ruiz,
arXiv:2307.09509



- Forecasts for a SN burst at 10 kpc.
- Current KamLAND allowed regions
- Current SK+ SNO allowed regions
- Forecast assuming as “true=nature” value KamLAND best fit
- Alleviate tension between reactor and Earth matter effects.
- Forecast assuming as “true” value SK+SNO best fit
- Increase tension between reactor and Earth matter effects.

Sensitivity to Δm_{21}^2 with SN neutrinos

RH, O. Mena and S. Palomares-Ruiz, arXiv:2307.09509



- Forecasts for a SN burst at 10 kpc.

- Current KamLAND all-sky sensitivity

- Current SK+ SNO sensitivity

- Forecast of value K_{SN}

- Δm_{21}^2 from reactor

- Δm_{21}^2 from MS.

- Δm_{21}^2 as “true” value

- Δm_{21}^2 tension between reactor and earth matter effects.

> 5σ tension with KamLAND!