

The Hubble Constant Troubled by Dark Matter

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Outline

- Introduction
- Dark Matter particles as the source of dark radiation
- Constraining the Hubble constant from non-thermal production of Dark Matter
 - LCDM
 - w CDM
- Discussion and perspectives

The concordance model - Λ CDM:

a set of assumptions

General Relativity + Cosmological Principle

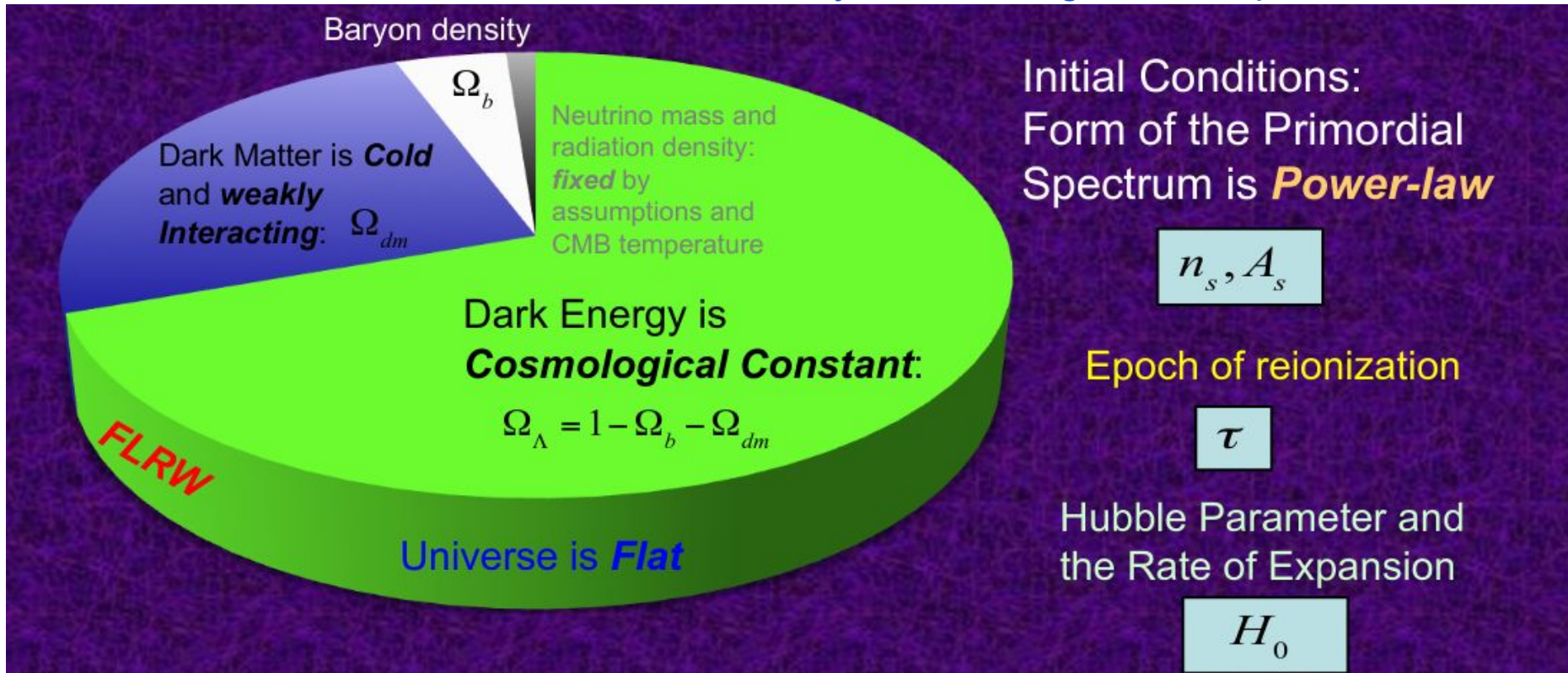


Figure from [A. Shafieloo](#)

The concordance model - Λ CDM: a set of assumptions confirmed by observational data

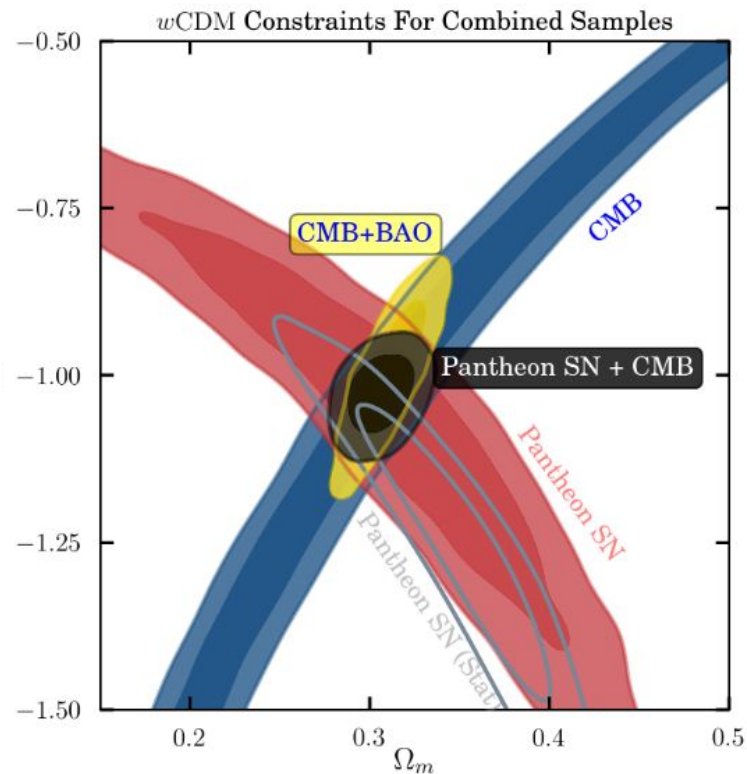
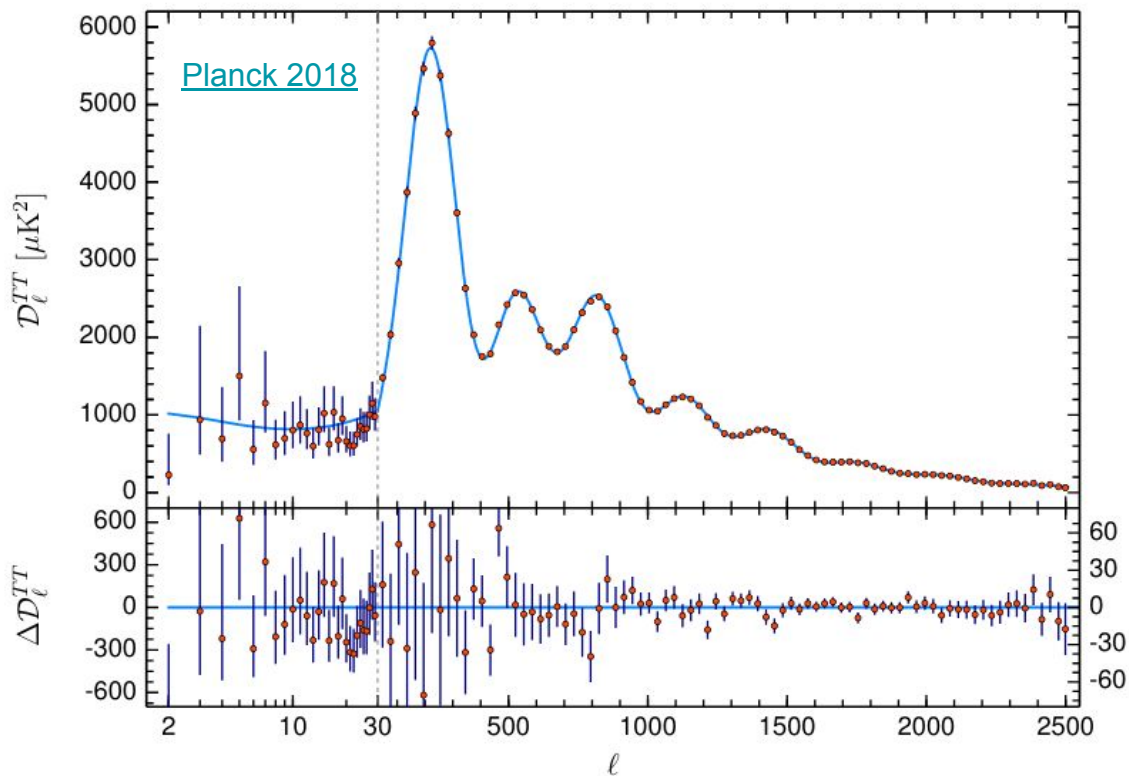


Figure from [Scolnic et al. \(2018\)](#)

The concordance model - Λ CDM

General Relativity + Cosmological Principle

Perfect fluid

$$g_{\mu\nu}dx^\mu dx^\nu = -dt^2 + a^2 d\chi^2$$

Scale factor

$$\frac{\ddot{a}}{a} = -\frac{4\pi G}{3} \left(\rho + \frac{3p}{c^2} \right) + \frac{\Lambda c^2}{3}$$

**FRIEDMANN
EQUATIONS**

$$\frac{\dot{a}^2 + kc^2}{a^2} = \frac{8\pi G\rho + \Lambda c^2}{3}$$



$$H^2 = H_0^2 \left(\frac{\Omega_m}{a^3} + \frac{\Omega_r}{a^4} + \frac{\Omega_k}{a^2} + \Omega_\Lambda \right)$$

$$\rho_c = \frac{3H^2}{8\pi G}$$

$$\Omega_M = \frac{\rho}{\rho_{crit}}$$

The concordance model - Λ CDM

How can we determine H_0 ?

$$g_{\mu\nu}dx^\mu dx^\nu = -dt^2 + a^2 d\chi^2$$



$$ds^2 = 0$$

For a photon

$$dt = a(t) dr$$

$$a = \frac{1}{1+z} \quad H(t) = \frac{1}{a(t)} \frac{da}{dt}$$

$$r = \int_{t_e}^{t_0} \frac{dt}{a(t)} = \int_0^z \frac{dz}{H(z)}$$

Distance measurements

- Luminosity distances from Cepheids and Supernovae
- Angular diameter distances from BAO, galaxy clustering

Using CMB observations

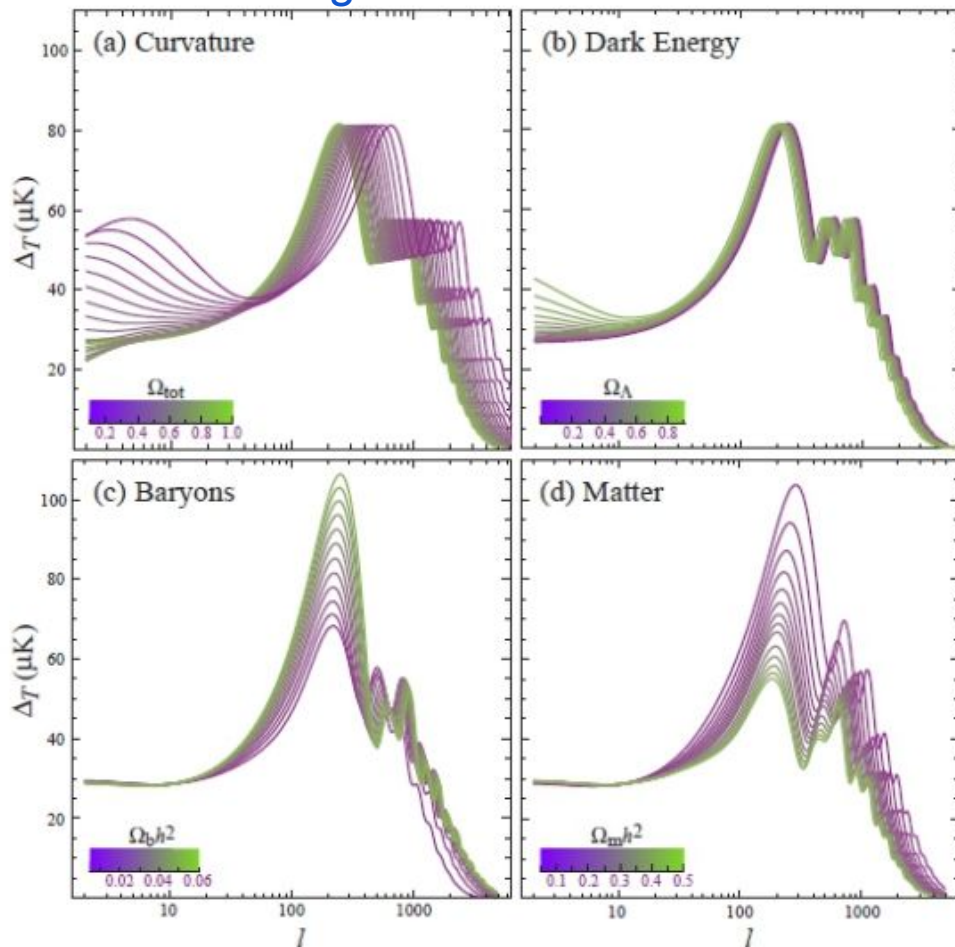
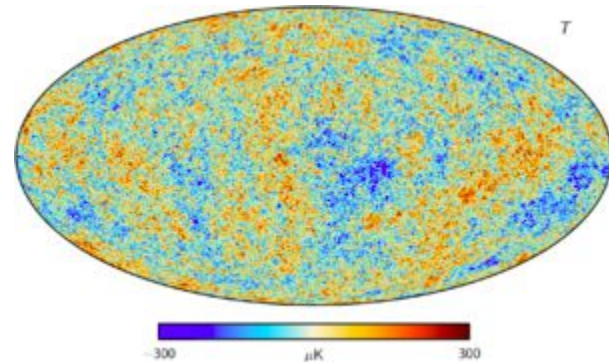
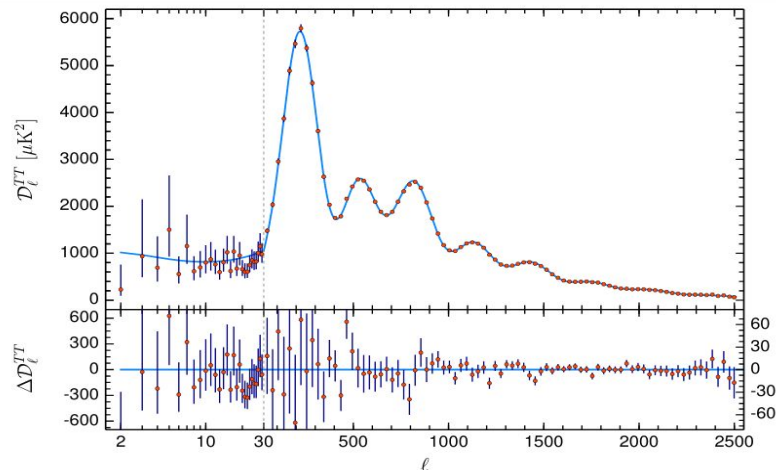


Figure from Hu and Dodelson (2001)

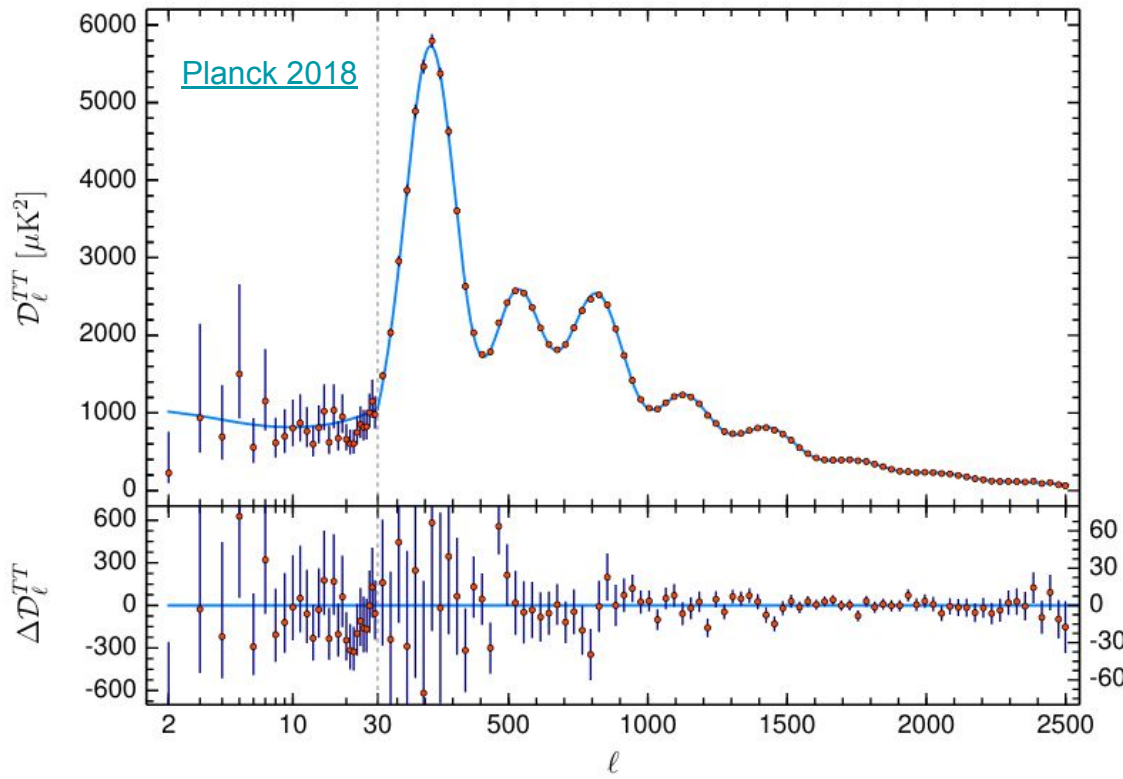


confront data and theory



The concordance model - Λ CDM:

a set of assumptions confirmed by observational data



Parameter	TT,TE,EE+lowE+lensing+BAO 68% limits
$\Omega_b h^2$	0.02242 ± 0.00014
$\Omega_c h^2$	0.11933 ± 0.00091
$100\theta_{\text{MC}}$	1.04101 ± 0.00029
τ	0.0561 ± 0.0071
$\ln(10^{10} A_s)$	3.047 ± 0.014
n_s	0.9665 ± 0.0038
H_0 [km s ⁻¹ Mpc ⁻¹]	67.66 ± 0.42
Ω_Λ	0.6889 ± 0.0056
Ω_m	0.3111 ± 0.0056

The concordance (??) model - Λ CDM: combination of reasonable assumptions

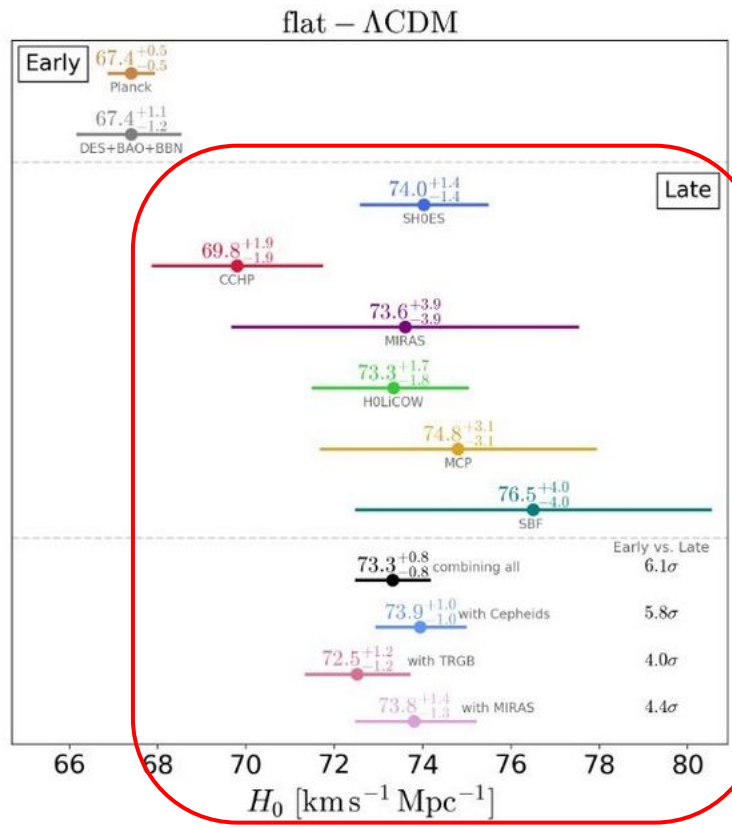
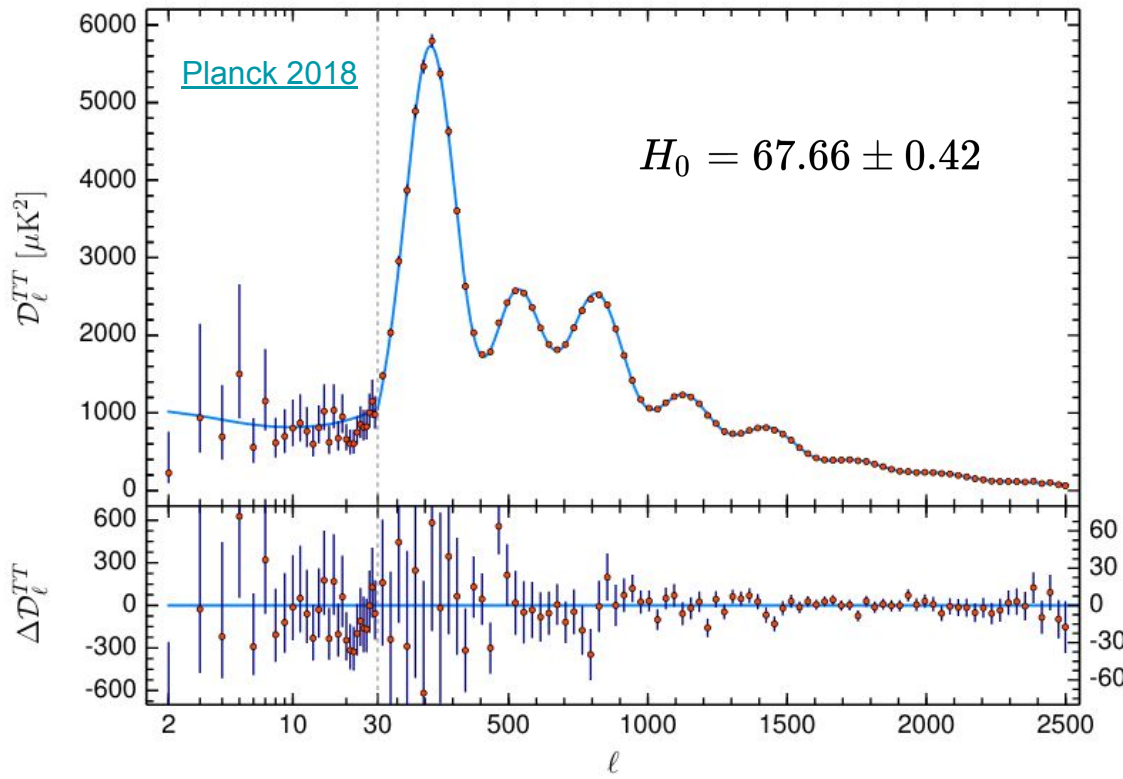


Figure from Verde et al. (2019)

The H_0 tension

A huge discrepancy between the Hubble constant inferred from the CMB, fitting the standard Λ CDM model, and the one obtained from local measurements.

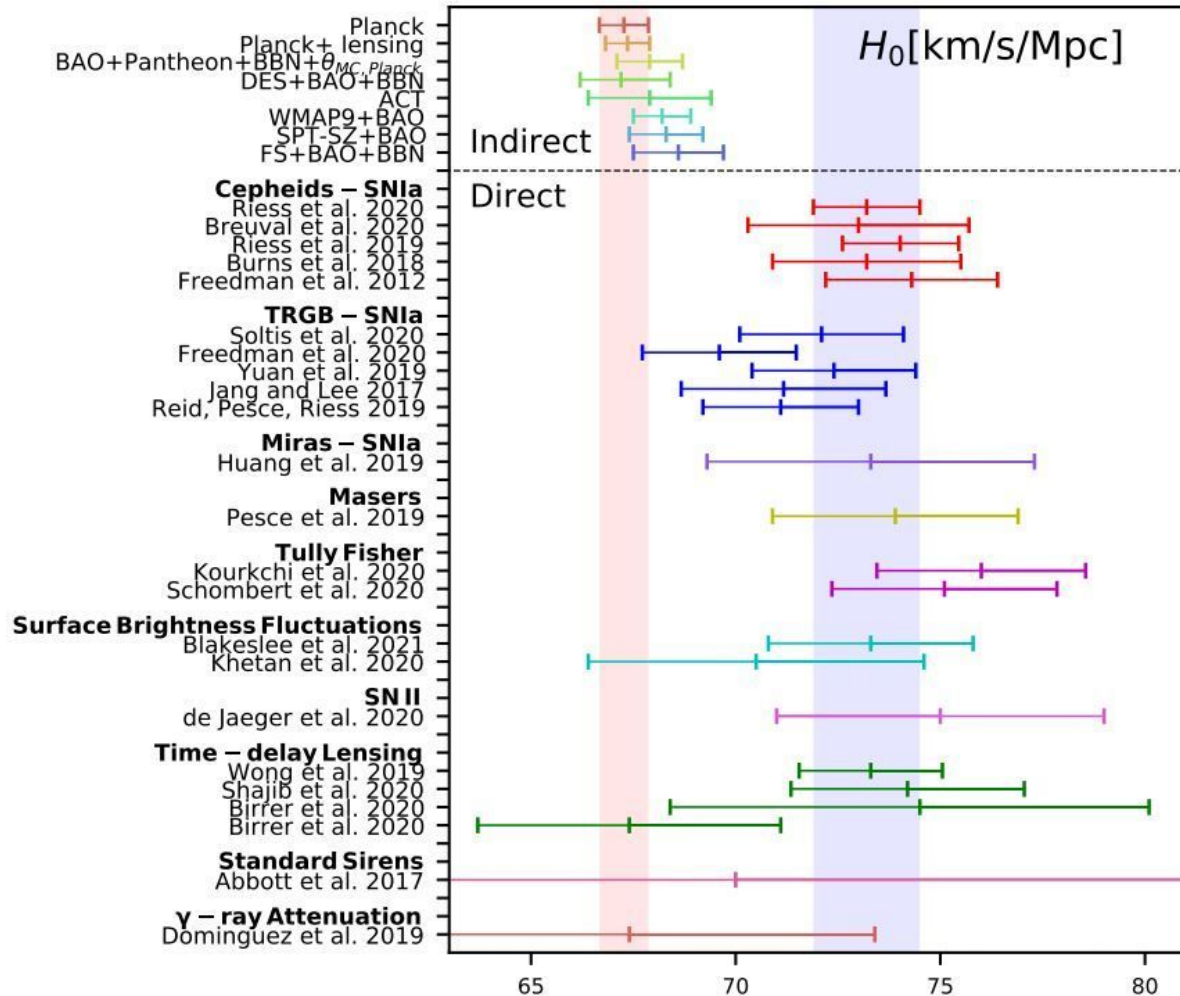


Figure from Di Valentino [\[2011.00246\]](#)

A lot of proposals to address the Hubble constant tension...

- The most famous attempts involve an **Early Dark Energy** component and **extra relativistic species at recombination**, but they can not increase the hubble constant enough to solve the tension with local measurements below 3σ .
- **Late time** modifications, such as **Phantom Dark Energy** and **Phenomenologically Emergent Dark Energy**, can completely solve the hubble tension within 1σ but leave the sound horizon unaltered introducing a tension with the BAO data.
- Last, another promising possibility is an **interaction between the dark matter and the dark energy models**, where the flux of energy between them allows for a lower matter density value and a larger H_0 value.

See [Di Valentino \(2020\)](#) and references therein.

Dark Matter particles as the source of dark radiation

In the realm of particle physics, the majority of models that try to solve the H_0 tension rely on new interactions involving the standard model neutrinos or decaying dark matter models [\[Abdalla, E. et al. \(2022\)\]](#).

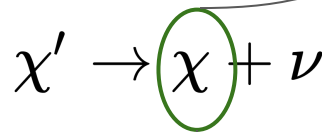
Here, instead, we introduce a non-thermal production mechanism of dark matter to increase the relativistic degrees of freedom and consequently raise H_0 , via the decay

$$\chi' \rightarrow \chi + \nu \qquad m_{\chi'} \gg m_{\chi}$$

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$$m_{\chi'} \gg m_{\chi}$$

mimic the effect of extra dark radiation, i.e. relativistic degrees of freedom N_{eff}

just a fraction of the dark matter abundance coming from this mechanism

Dark Matter particles as the source of dark radiation

Reminding that the radiation density (ρ_{rad}) is determined by the photon's temperature (T) and the relativistic degrees of freedom (g_*), i.e.,

$$\rho_{rad} = \frac{\pi^2}{30} g_* T^4. \quad g_* = 2 + \frac{7}{4} \left(\frac{4}{11} \right)^{4/3} N_{eff}.$$

where N_{eff} is the effective number of relativistic neutrino species, where in the Λ CDM is $N_{eff} = 3$.

As we are trying to raise H_0 by increasing N_{eff} , N_{eff} tell us how much extra radiation we are adding to the universe via our mechanism

$$\Delta N_{eff} = \frac{\rho_{extra}}{\rho_{1\nu}}.$$

Dark Matter particles as the source of dark radiation

In principle, we may reproduce the effect of an extra neutrino species by adding any other kind of radiation source. Calculating the ratio between one neutrino species density and cold dark matter density at the matter-radiation equality ($t = t_{eq}$) we get,

$$\frac{\rho_{1\nu}}{\rho_{DM}} \Big|_{t=t_{eq}} = \frac{\Omega_{\nu,0}\rho_c}{3a_{eq}^4} \times \left(\frac{\Omega_{DM,0}\rho_c}{a_{eq}^3} \right)^{-1} = 0.16.$$

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The 4-momentum conservation implies,

$$E_\chi(\tau) = m_\chi \left(\frac{m_{\chi'}}{2m_\chi} + \frac{m_\chi}{2m_{\chi'}} \right) \equiv m_\chi \gamma_\chi(\tau), \quad \Rightarrow \quad \frac{E_\chi(t)}{m_\chi} = \left[1 + \left(\frac{a(\tau)}{a(t)} \right)^2 (\gamma_\chi^2(\tau) - 1) \right]^{1/2} \equiv \gamma_\chi(t).$$

$$\gamma_\chi(t) = \sqrt{\frac{(m_\chi^2 - m_{\chi'}^2)^2}{4m_\chi^2 m_{\chi'}^2} \left(\frac{\tau}{t} \right) + 1}.$$

Dark Matter particles as the source of dark radiation

Let the dark matter energy be written as $E_\chi = m_\chi (\gamma_\chi - 1) + m_\chi$

In the ultra relativistic regime the first term dominates, such that we can rewrite the above equation as

$$E_{DM} = N_{HDM} m_\chi (\gamma_\chi - 1) + N_{CDM} m_\chi \quad N_{HDM} \ll N_{CDM}$$

the total number of relativistic dark matter particles (hot particles) the total number of nonrelativistic DM (cold particles)

The ratio between relativistic and nonrelativistic dark matter density energy is,

$$\frac{\rho_{HDM}}{\rho_{CDM}} = \frac{N_{HDM} m_\chi (\gamma_\chi - 1)}{N_{CDM} m_\chi} \equiv f (\gamma_\chi - 1)$$

Dark Matter particles as the source of dark radiation

We find that the extra radiation produced via this mechanism is

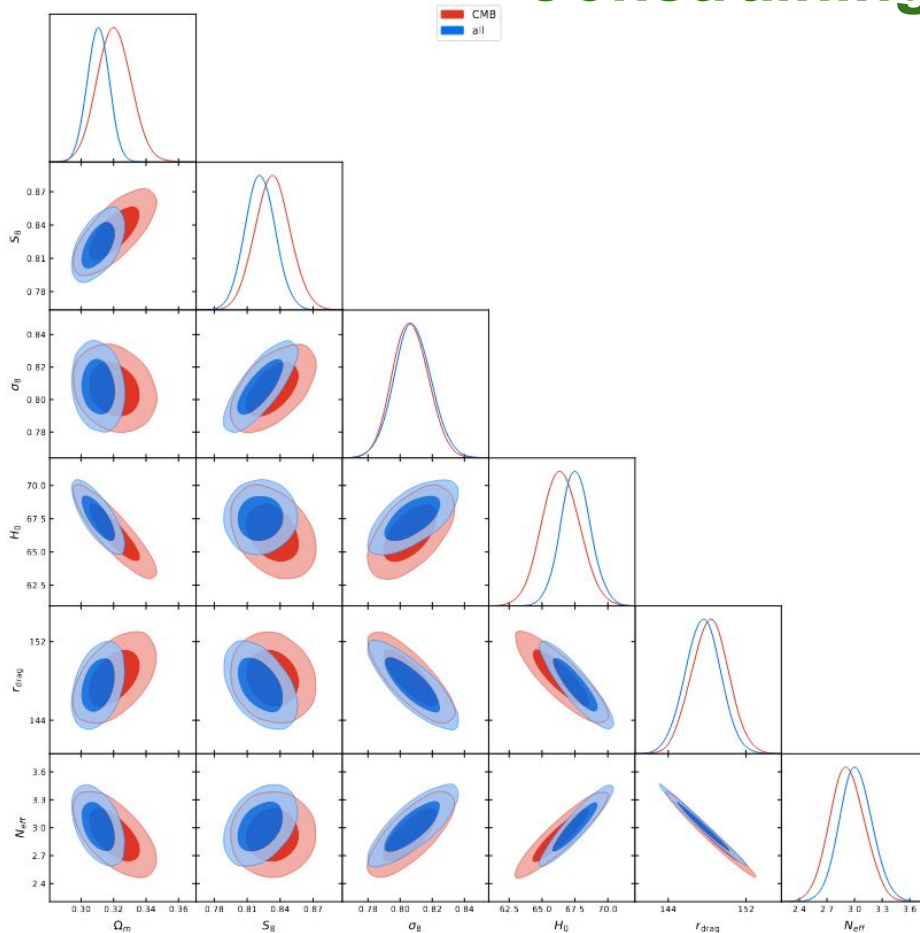
$$\Delta N_{eff} = \lim_{t \rightarrow t_{eq}} \frac{f(\gamma_\chi - 1)}{0.16},$$

In the regime $m_{\chi'} \gg m_\chi$

$$\Delta N_{eff} \approx 2.5 \times 10^{-3} \sqrt{\frac{\tau}{10^6 s}} \times f \frac{m_{\chi'}}{m_\chi}.$$

for $t_{eq} \approx 50,000$ years $\approx 1.6 * 10^{12}$ s

Constraining the Hubble constant from varying N_{eff}



The standard Λ CDM model with N_{eff} free to vary already allow for higher values of H_0 .

Now, we are able to explore the origin of this increase on H_0 through our expression for ΔN_{eff}

How to constrain the Hubble constant from this non-thermal production of Dark Matter?

Embed ΔN_{eff} on a Boltzmann code: [CAMB](#)

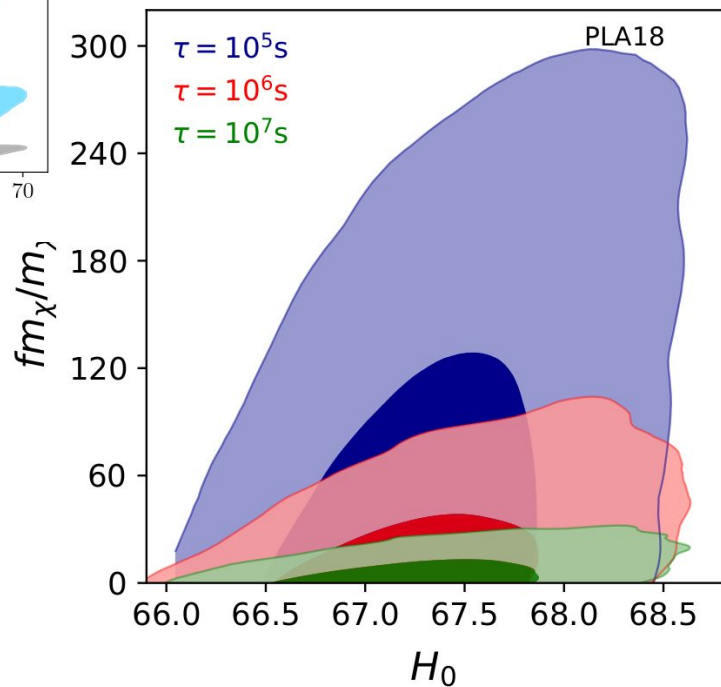
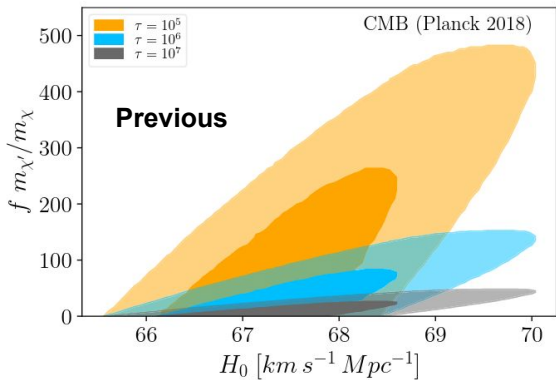
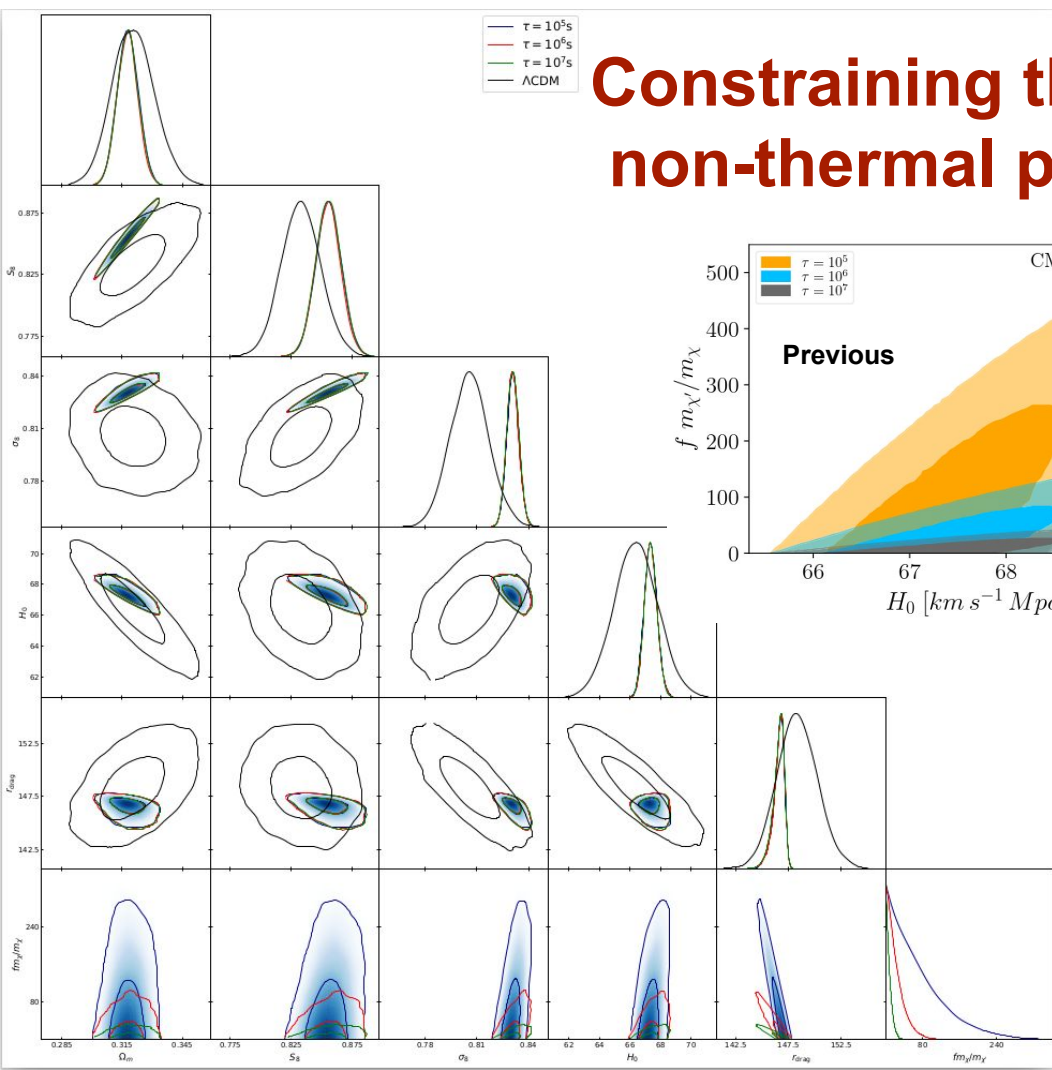
Establish the priors on \mathcal{T} and $f \frac{m_{\chi'}}{m_{\chi}}$

Perform an MCMC analysis: [CosmoMC](#)

\mathcal{T}	$f \frac{m_{\chi'}}{m_{\chi}}$	ΔN_{eff}
$10^5 s$	0 – 800	0 – 0.6
$10^6 s$	0 – 300	0 – 0.6
$10^7 s$	0 – 100	0 – 0.6

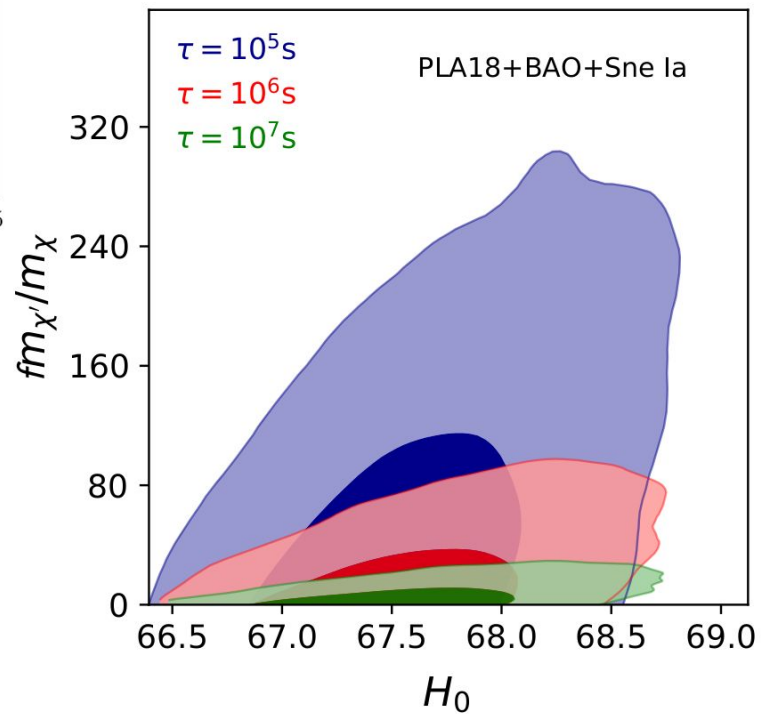
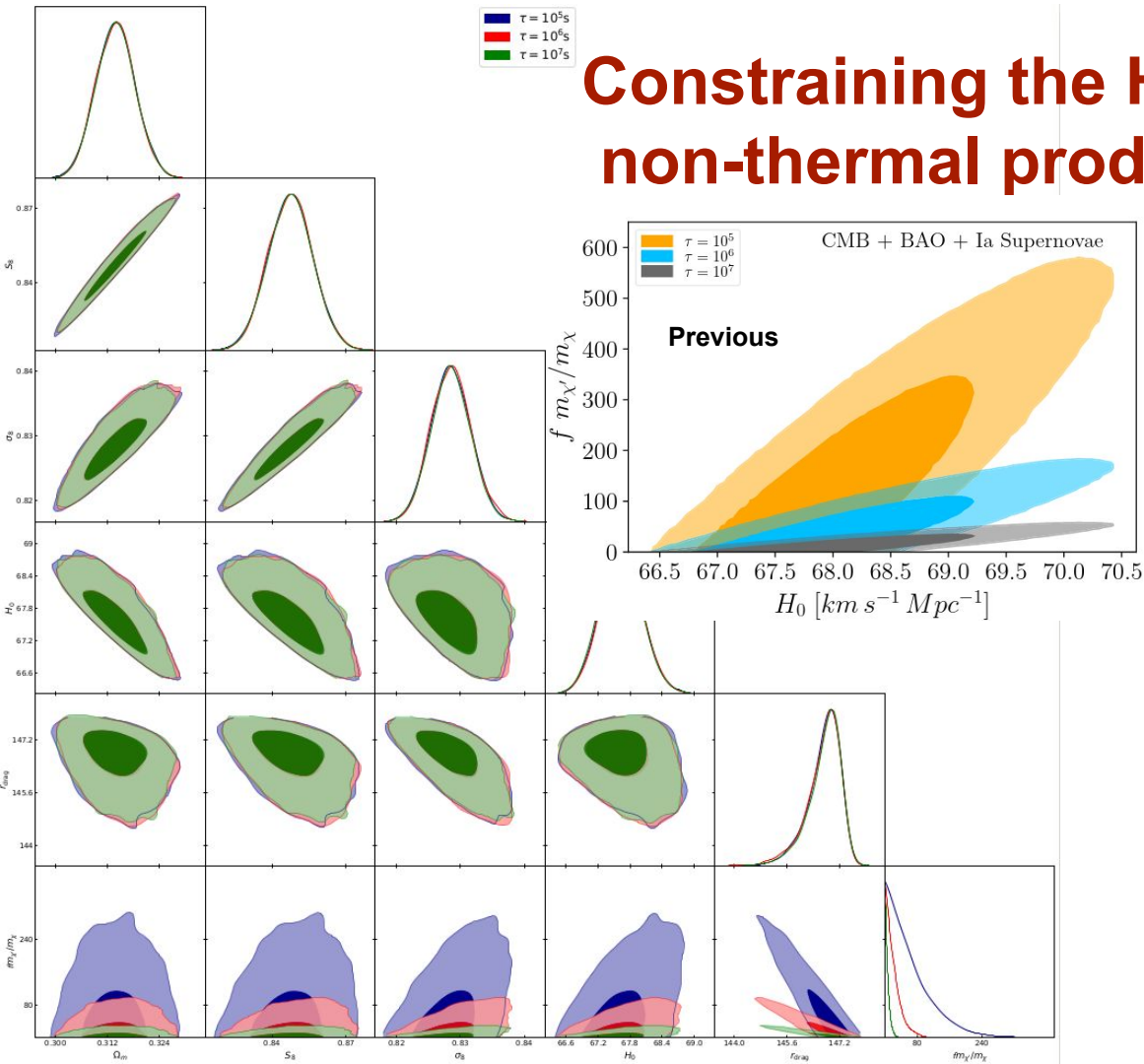
Constraining the Hubble constant from non-thermal production of Dark Matter

In the LCDM scenario



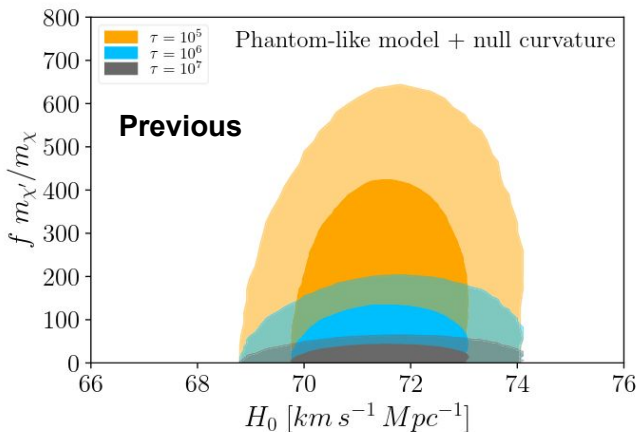
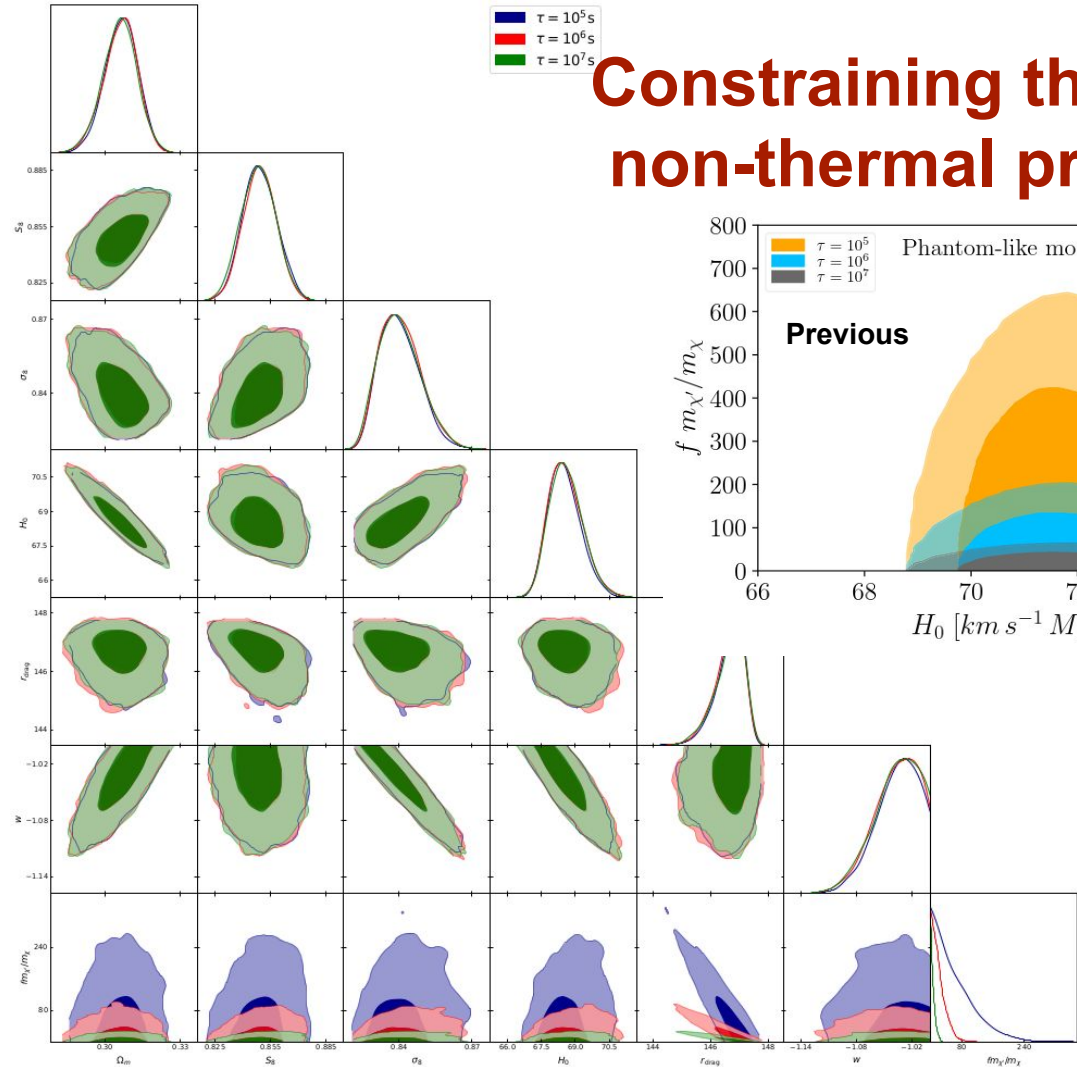
Constraining the Hubble constant from non-thermal production of Dark Matter

In the Λ CDM scenario

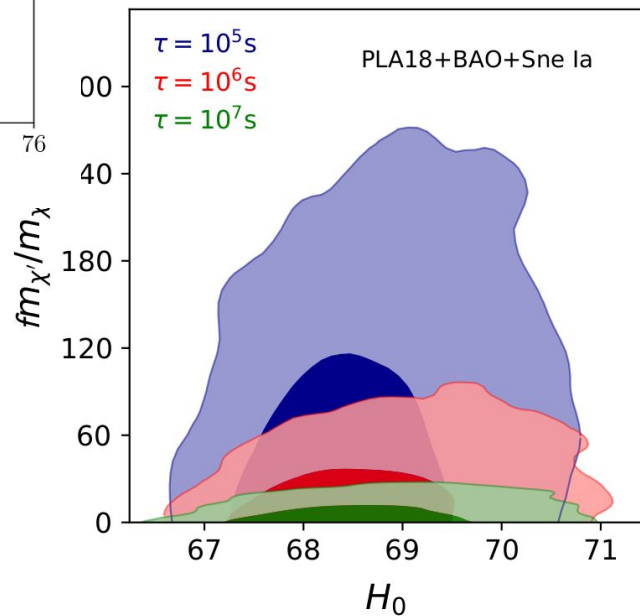


■ $\tau = 10^5$ s
■ $\tau = 10^6$ s
■ $\tau = 10^7$ s

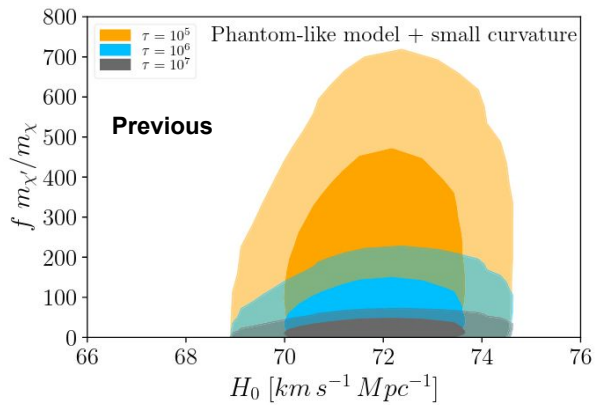
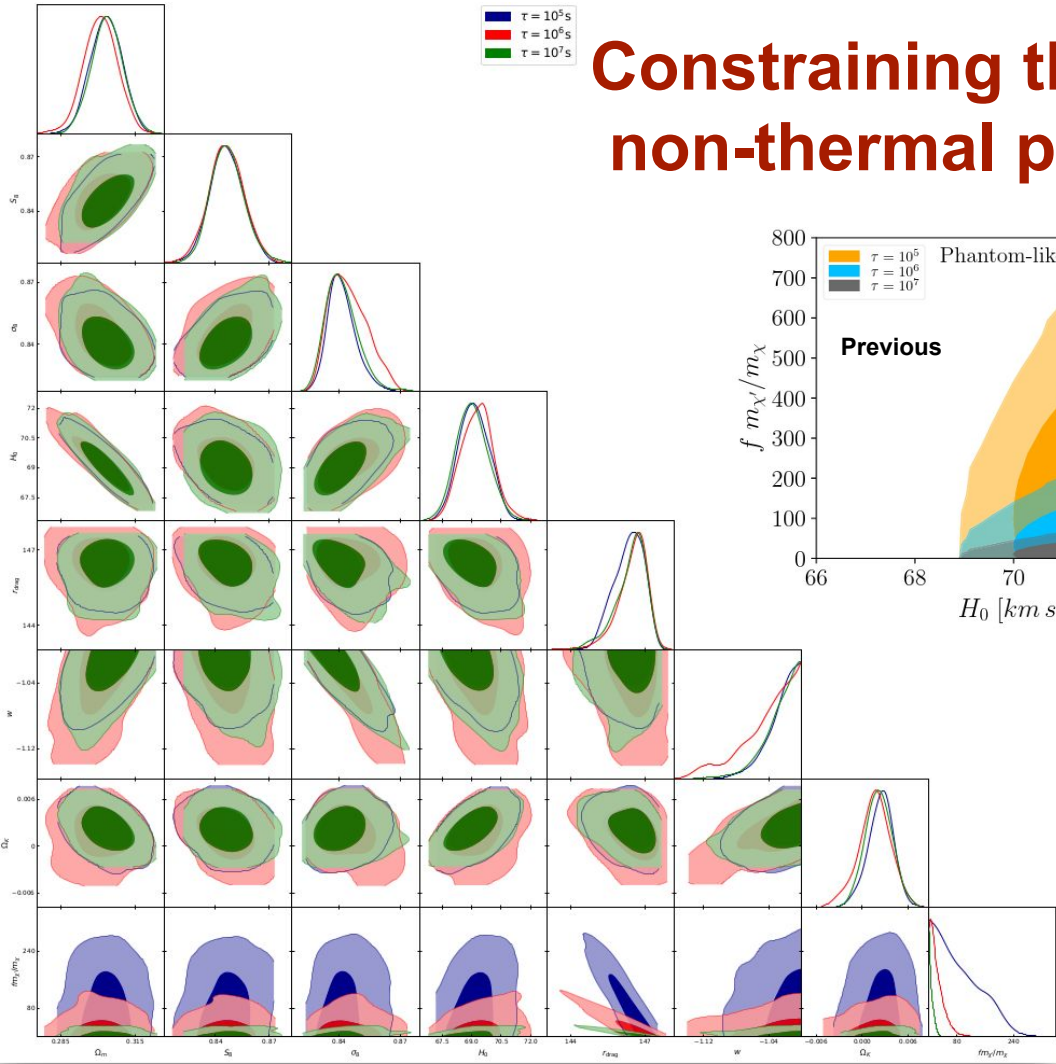
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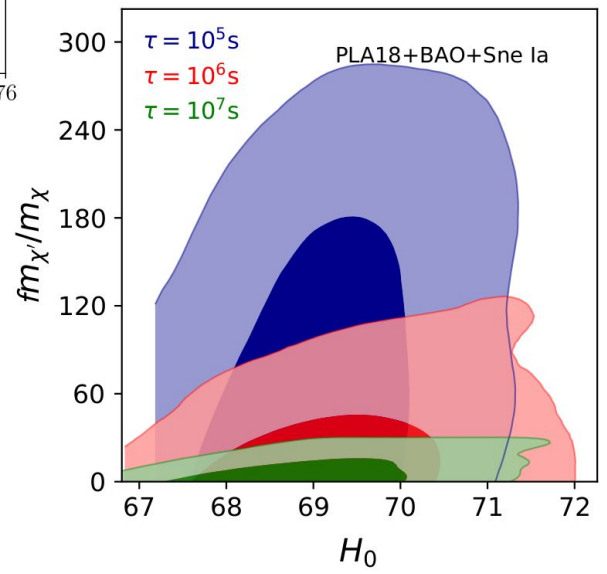
In the w CDM scenario



Constraining the Hubble constant from non-thermal production of Dark Matter



In the $w\text{CDM} + \text{curvature}$ scenario



Perspectives

BBN bounds constraints the decay process to happen between

$$10^2 s \leq \tau \leq 10^4 s$$

So, we need to **consider this new prior on our analysis!**

Structure formation needs $f \leq 0.01$ to this non-thermal dark matter production be consistent with clustering data.

We can use this information to put constraints on the fraction $\frac{m_{\chi'}}{m_{\chi}}$

