

DarkCosmoGrav

January 23–25, 2023



New Frontiers in Particle Physics, Gravity,
and Cosmology



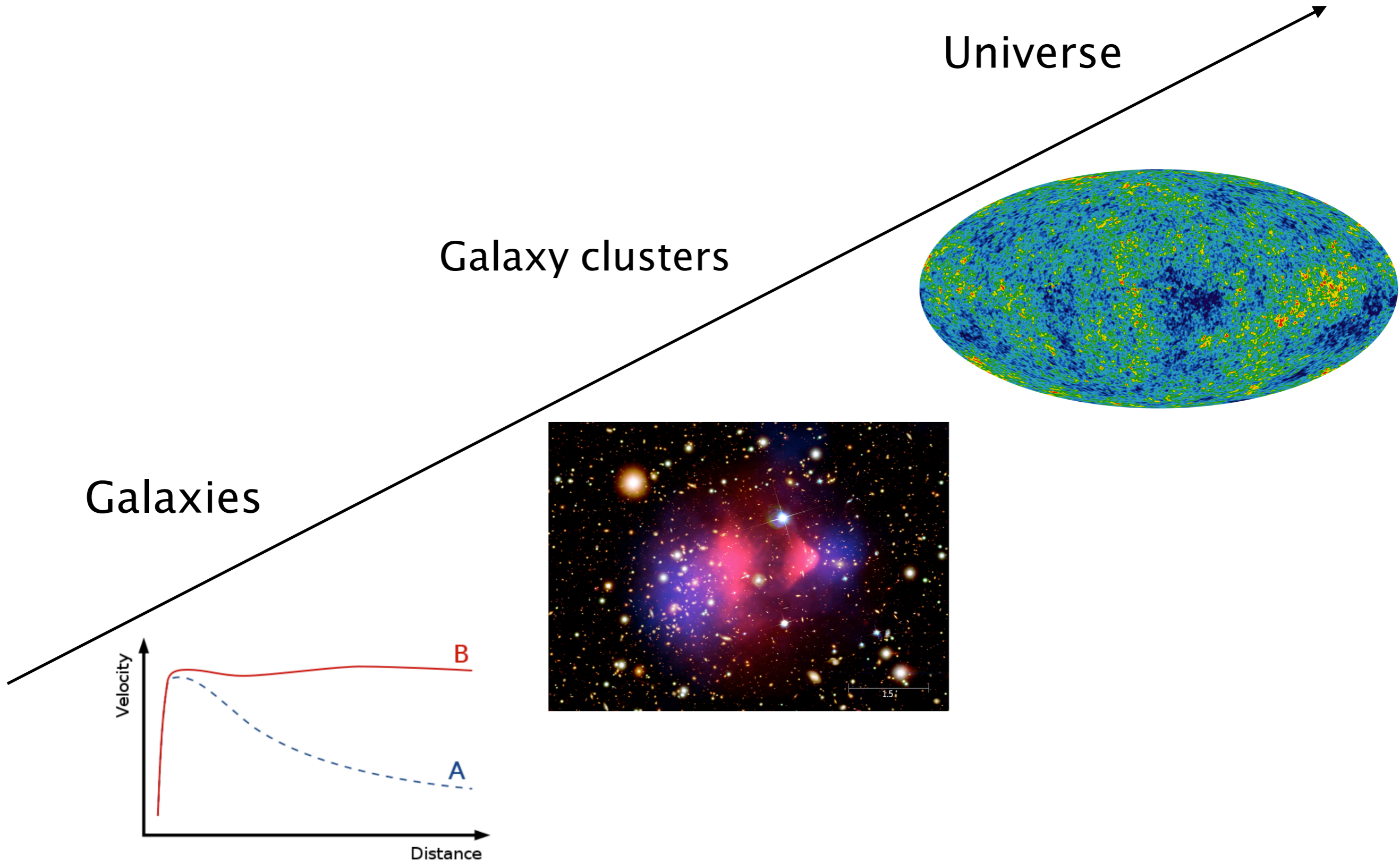
Axion Free-kick Misalignment Mechanism

Ling-Xiao Xu

University of Padova & INFN, Padova

with Seokhoon Yun, [hep-ph/2211.13074](https://arxiv.org/abs/hep-ph/2211.13074)

Dark matter exists



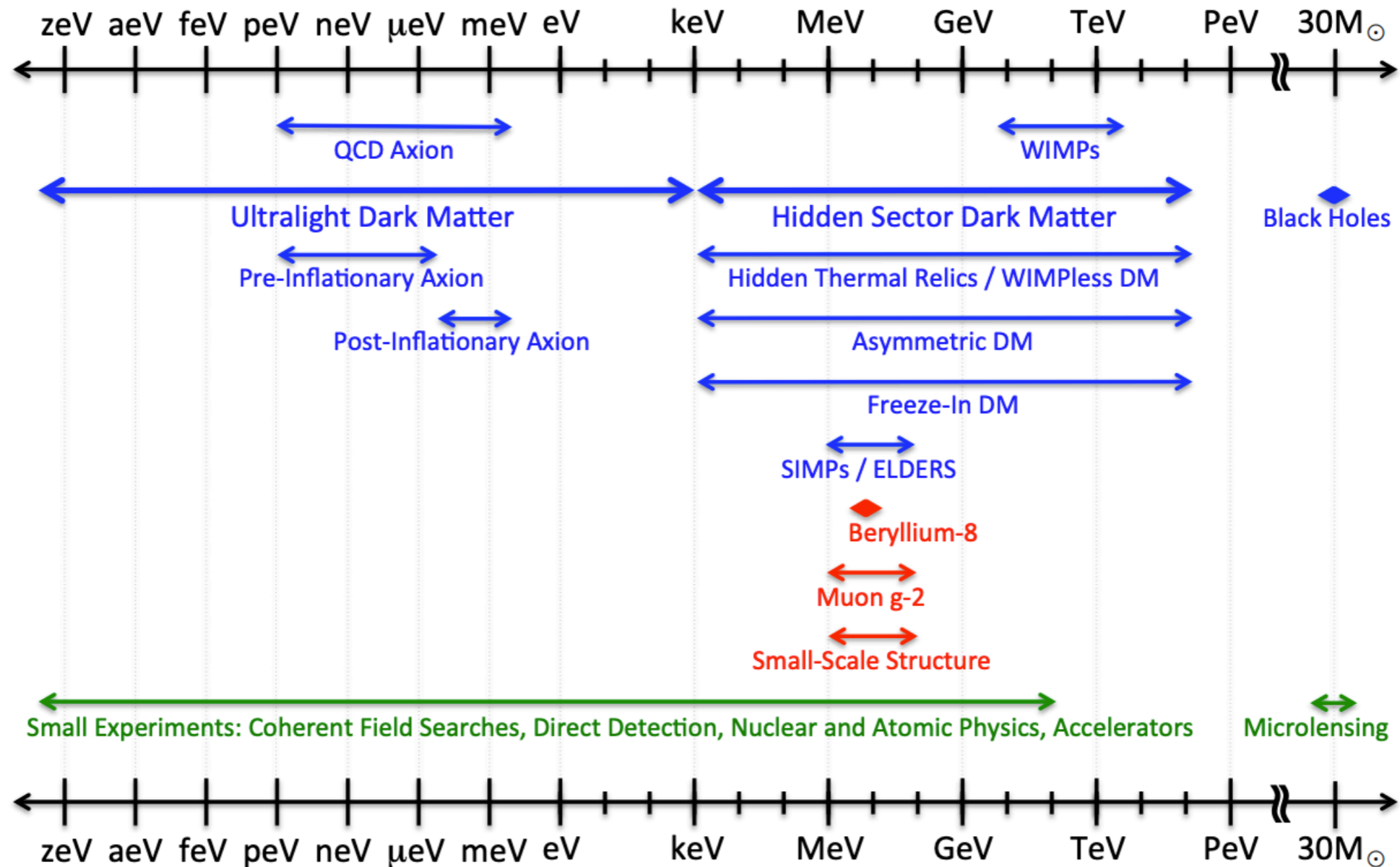
Microscopic nature of dark matter

very few constraints:

- Upper bound $\sim 10^{48}$ GeV, due to the observed absence of gravitational lensing.
- Lower bound $\sim 10^{-22}$ eV for boson, otherwise the de Broglie wave length would be too large for it to fit within galaxy.
- Lower bound ~ 10 eV for fermion, otherwise the number density would be too large for it to fit in galaxy while still obey Pauli exclusion principle.
- Self interactions of DM should be small enough, there is less than $\mathcal{O}(1)$ scattering at the time scale when galaxy clusters merge.
- DM interacts weakly (if at all) with SM besides gravity, and needs to be stable/long-lived.

Huge theory landscape

Dark Sector Candidates, Anomalies, and Search Techniques



hep-ph/1707.04591

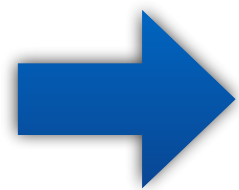
US Cosmic Visions: New ideas in Dark Matter

The conventional misalignment mechanism

Preskill, Wise, Wilczek; Abbott, Sikivie; Dine, Fischler, 83

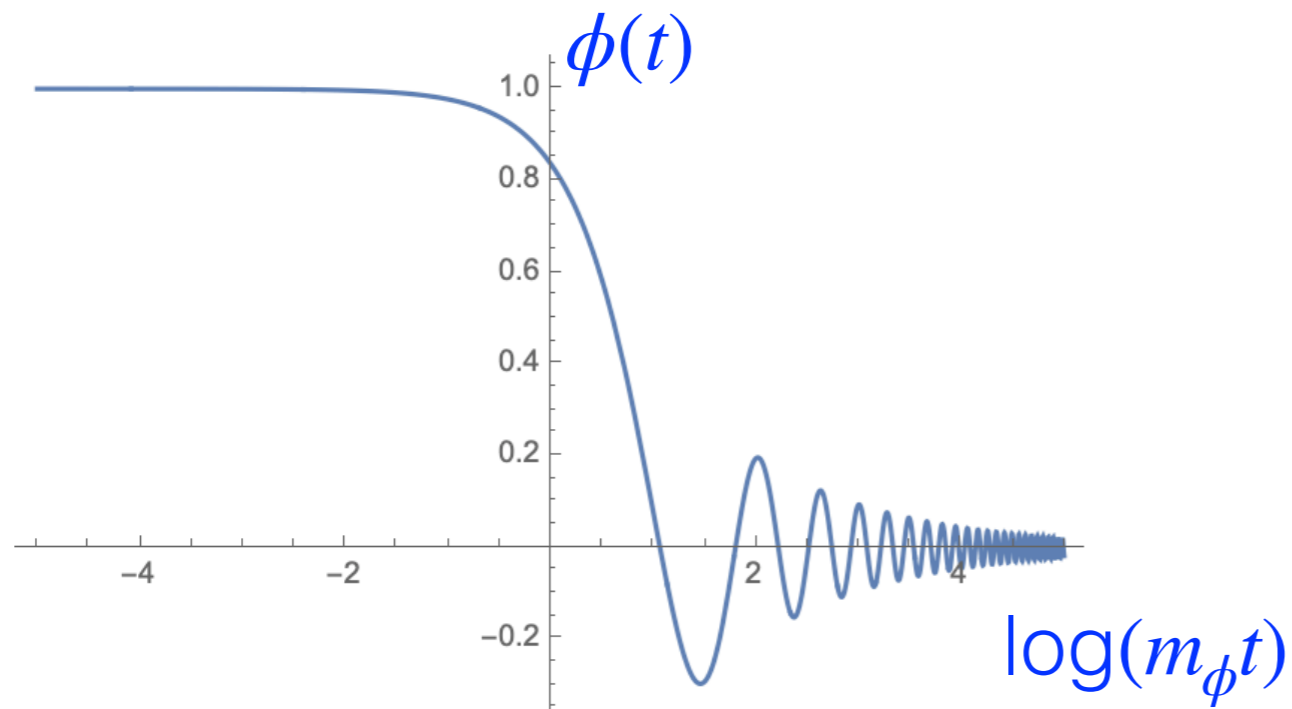
Consider a spatially-homogeneous real scalar field in an expanding universe:

$$\frac{1}{\sqrt{|g|}} \mathcal{L} = \frac{1}{2} (\partial^\mu \phi^*) (\partial_\mu \phi) - \frac{m_\phi^2}{2} \phi^* \phi$$

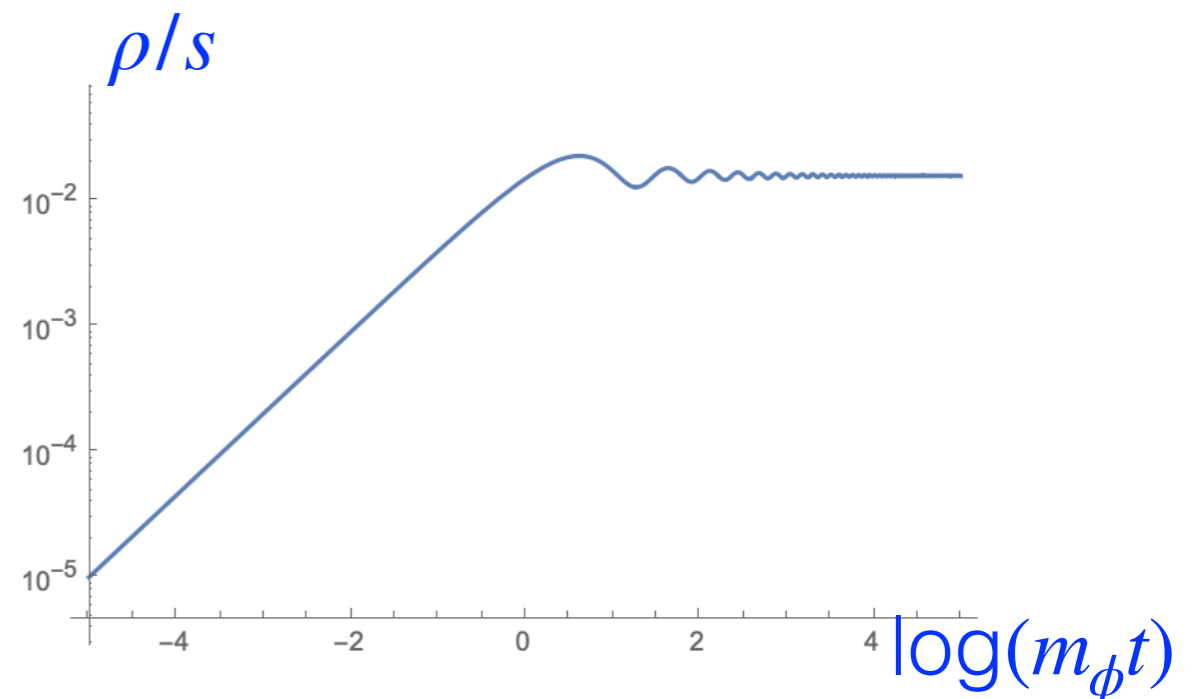


equation of motion:

$$\ddot{\phi}(t) + 3H\dot{\phi}(t) + m_\phi^2\phi(t) = 0$$



$$H \sim 1/t$$

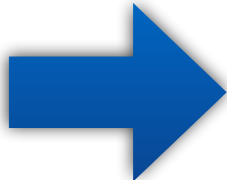


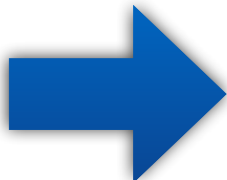
Ultralight scalars are ubiquitous

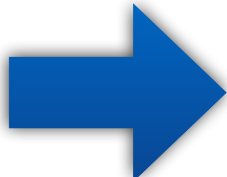
- Ultralight scalars can arise from many theoretical frameworks.
- Ultralight scalars are (will be) under phenomenological/experimental scrutiny.

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 { dark matter (Preskill, Wise, Wilczek; Abbott, Sikivie; Dine, Fischler, 83)
baryogenesis (Affleck, Dine, 85)
Inflation (Guth, 81; Linde, 82)

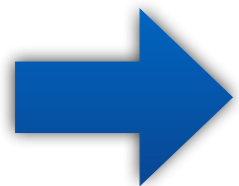
 { Solving strong CP problem – axion (Peccei, Quinn, 77; Weinberg, 78; Wilczek, 78)
Selecting cosmological constant (Abbott, 85)
Solving the hierarchy problem – relaxion (Graham, Kaplan, Rajendran, 15)

 String axiverse (Arvanitaki, Dimopoulos, Dubovsky, Kaloper, March-Russell, 09)

a lot more recent activities, e.g. many Snowmass 2022 papers

Ultralight scalars are ubiquitous

- Ultralight scalars can arise from many theoretical frameworks.
- Ultralight scalars are (will be) under phenomenological/experimental scrutiny.
- Theories of scalars are not 'rigid', as reflected in the hierarchy problem or the quality problem.



The standard misalignment mechanism can be over-simplified and be easily modified.

Nontrivial interplay are expected if multiple scalars are present.

A new misalignment mechanism

based on nontrivial interplay between axion (ALP) and a light dilaton (DLP)

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$$\mathcal{L} = \sqrt{|g|} \left\{ \frac{f_\theta^2 \hat{\chi}^2}{2} (\partial_\mu \theta)^2 + \frac{1}{2} (\partial_\mu \chi)^2 - V(\sigma, \theta) \right\}$$

$$\chi = F \hat{\chi} = F e^{\frac{\sigma(x)}{F}}$$

Weyl compensator

$$\left\{ \begin{array}{l} V(\sigma, \theta) = V(\sigma) + \hat{\chi}^4 V(\theta) \\ \quad = V(\sigma) + m_\theta^2 f_\theta^2 \hat{\chi}^4 (1 - \cos[\theta]) \\ \\ V(\sigma) = \lambda F^4 \left(-\frac{4}{4 - \epsilon} \hat{\chi}^{4 - \epsilon} + \hat{\chi}^4 \right) + \lambda F^4 \frac{\epsilon}{4 - \epsilon} \end{array} \right.$$

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$$\text{Global minima:} \quad \langle \sigma \rangle = 0 \quad \langle \theta \rangle = 0 \text{ mod } 2\pi$$

Equations of motion and initial conditions

$$\left\{ \begin{array}{l} \ddot{\theta}(t) + \left(3H + \frac{2}{F} \dot{\sigma}(t) \right) \dot{\theta}(t) + m_{\theta}^2 e^{2\frac{\sigma(t)}{F}} \sin[\theta(t)] = 0 \\ \frac{\ddot{\sigma}(t)}{F} + \left(3H + \frac{\dot{\sigma}(t)}{F} \right) \frac{\dot{\sigma}(t)}{F} + \frac{m_{\sigma}^2}{\epsilon} e^{2\frac{\sigma(t)}{F}} \left(1 - e^{-\epsilon\frac{\sigma(t)}{F}} \right) = -\mathcal{K}(\theta) \end{array} \right.$$

Kick term: $\mathcal{K}(\theta) \sim \frac{4m_{\theta}^2 f_{\theta}^2 (1 - \cos[\theta]) e^{2\frac{\sigma(t)}{F}} - f_{\theta}^2 \dot{\theta}^2}{F^2}$

- Dilaton field-dependent Hubble friction, masses
- Various possible histories for field evolution, depending on initial conditions

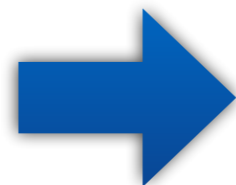
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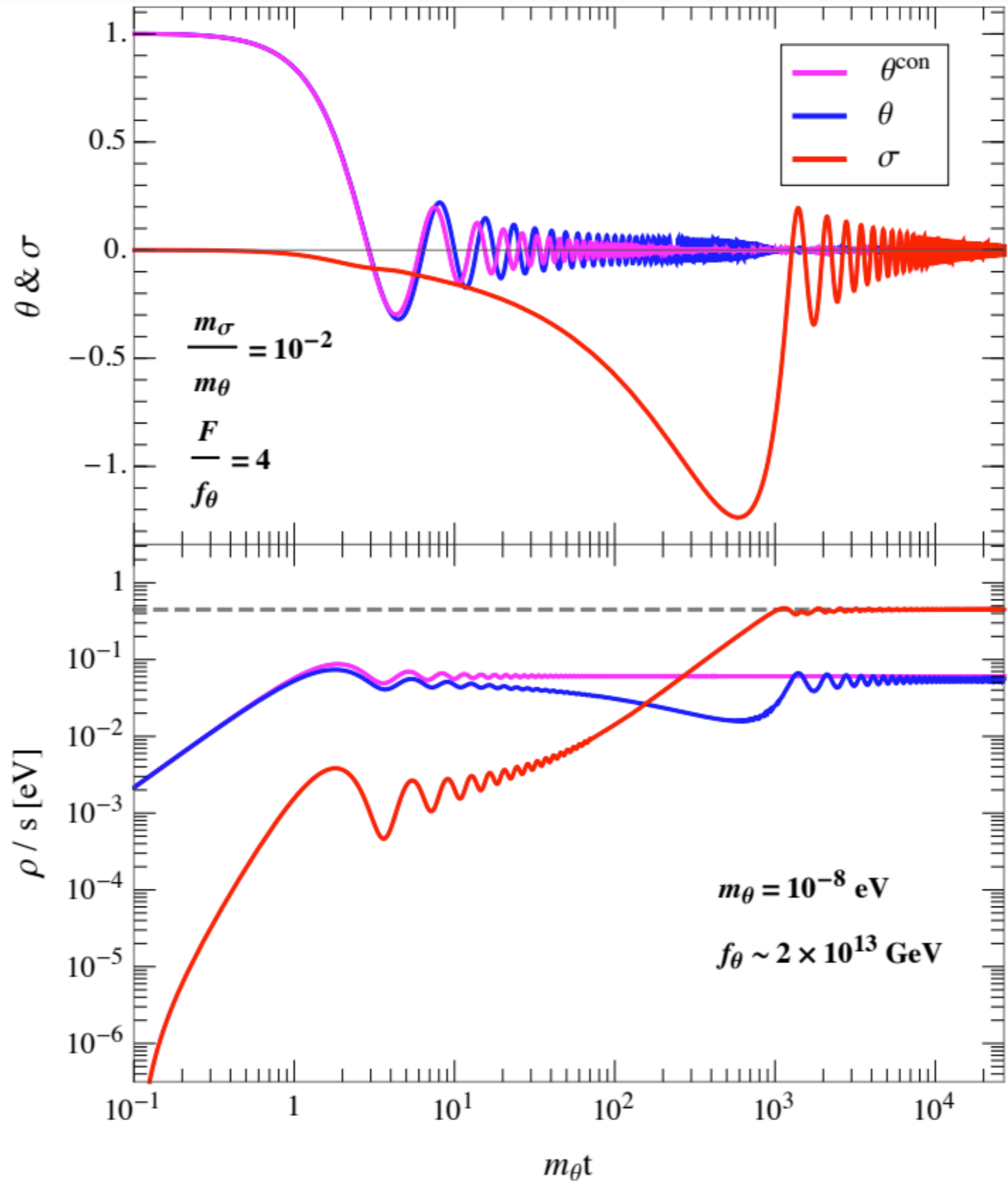
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- Dilaton field-dependent Hubble friction, masses
- Various possible histories for field evolution, depending on initial conditions
- Let us consider $m_{\theta} > m_{\sigma}$ $\dot{\theta}(t_i) = \dot{\sigma}(t_i) = 0$

$$\theta(t_i) \sim \mathcal{O}(1), \sigma(t_i) \simeq 0.$$

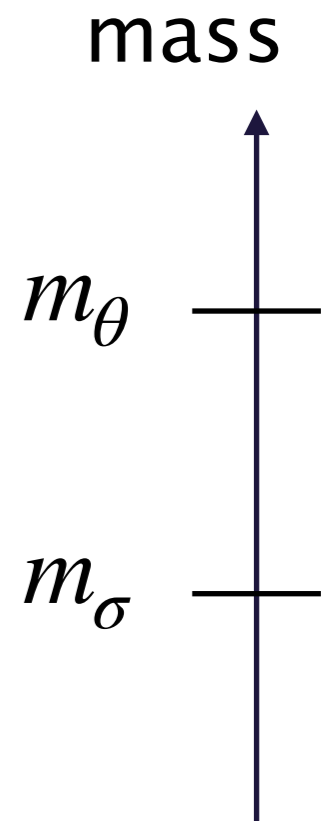
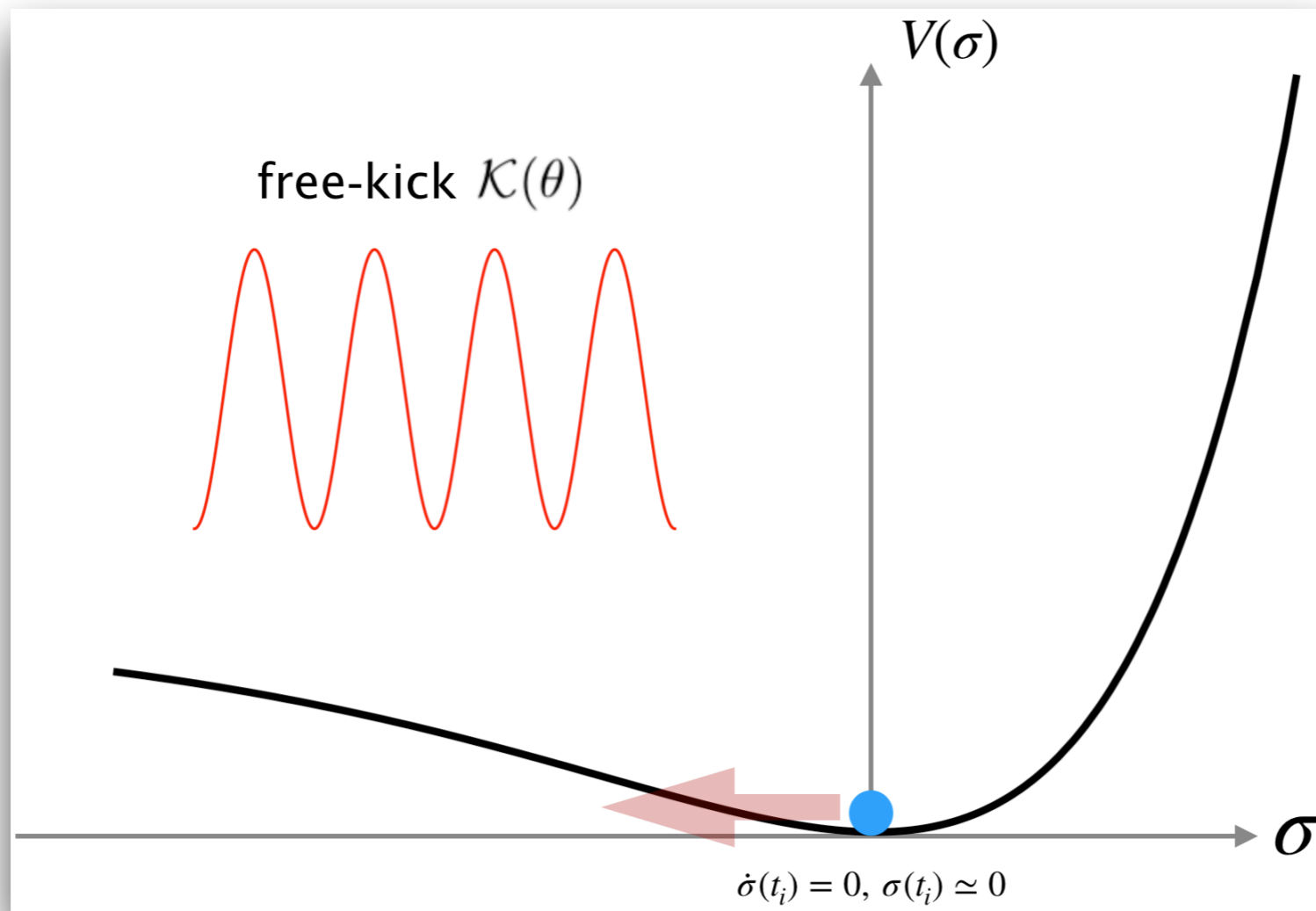


axion misalignment



For fixed values of m_θ and f_θ , DM abundance is enhanced!

Axion 'free-kick' misalignment mechanism



Scalars

Roles

axion



a football player

dilaton



football

Conclusion and future directions

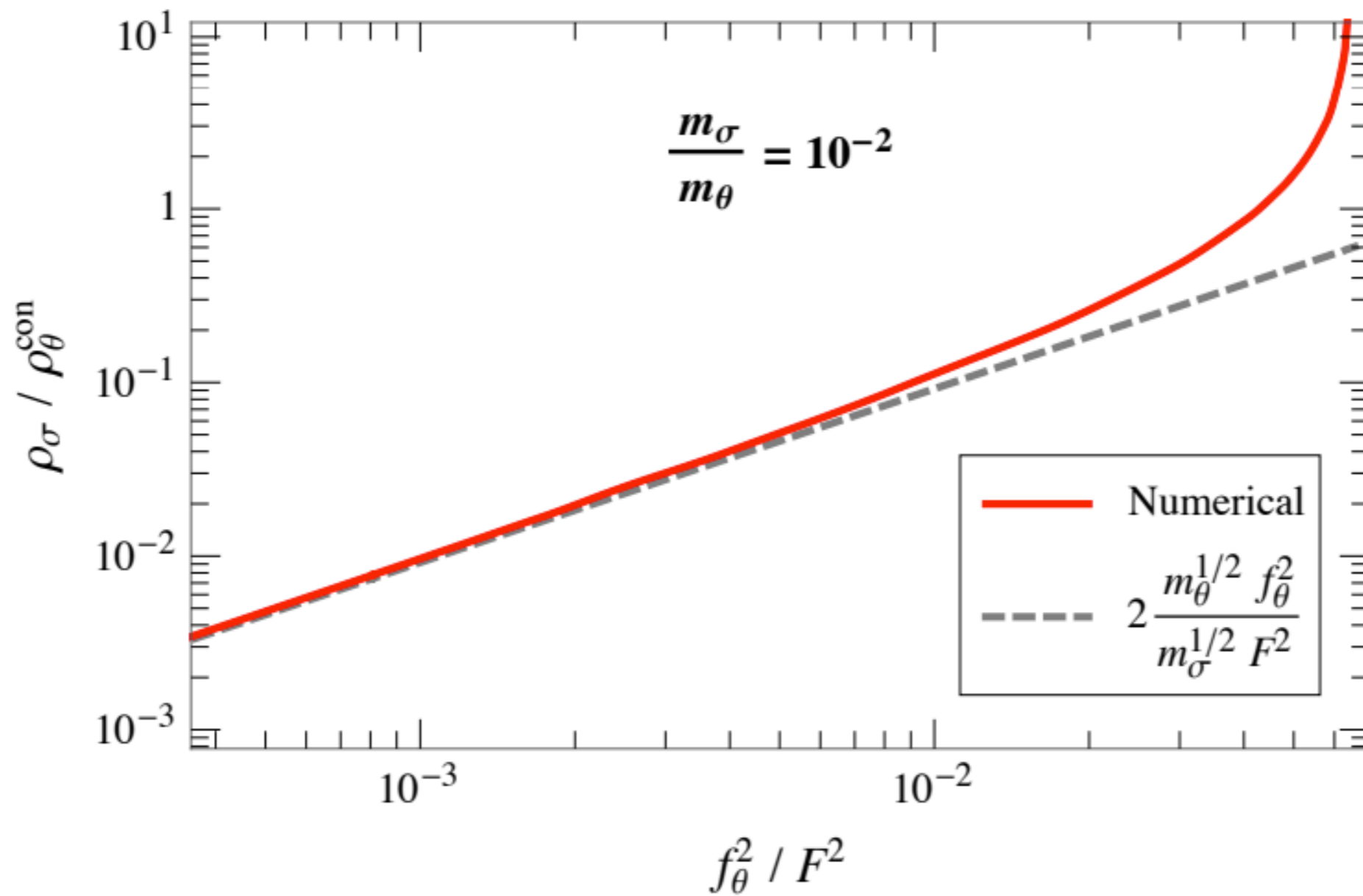
- We propose an alternative misalignment mechanism based on nontrivial interplay between axion and dilaton.
- The ‘kick’ effect is ubiquitous if multiple ultralight scalars are present in nature, as motivated by different problems and BSM scenarios.
- Work in progress to understand the implications to standard axion/relaxion dark matter scenarios.

Conclusion and future directions

- We propose an alternative misalignment mechanism based on nontrivial interplay between axion and dilaton.
- The ‘kick’ effect is ubiquitous if multiple ultralight scalars are present in nature, as motivated by different problems and BSM scenarios.
- Work in progress to understand the implications to standard axion/relaxion dark matter scenarios.
- What is the model-building price to pay to obtain a dilaton lighter than axion in a UV theory (i.e. hierarchy problem)?
- What is the optimal experimental strategy to measure the ‘kick’ effect between ultralight scalars?
- Implications to the “cosmological naturalness” paradigm?

Thank you!

backup



$$\frac{\rho_\sigma}{\rho_\theta^{\text{con}}} \sim \left(\frac{m_\theta}{m_\sigma} \right)^{1/2} \left(\frac{f_\theta}{F} \right)^2 (1 - \cos \theta_i)$$