



Nuclear physics for the precise extraction of V_{ud}

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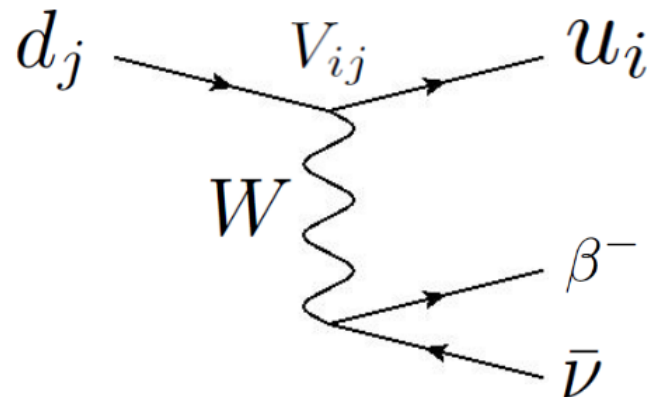
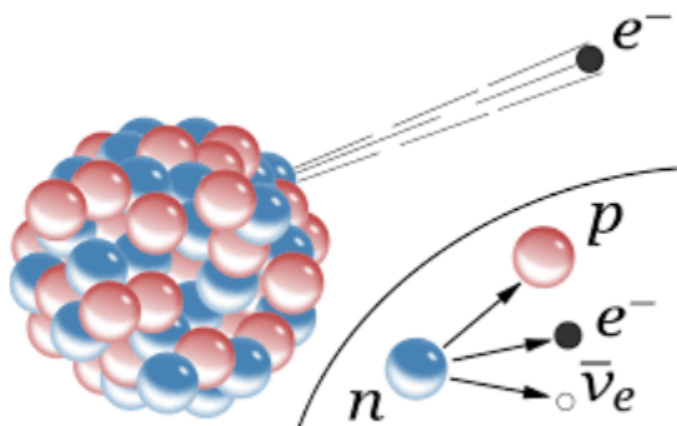
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Marciana 2023 – Lepton Interactions with Nucleons and Nuclei

5 September, 2023

Nuclear beta decay



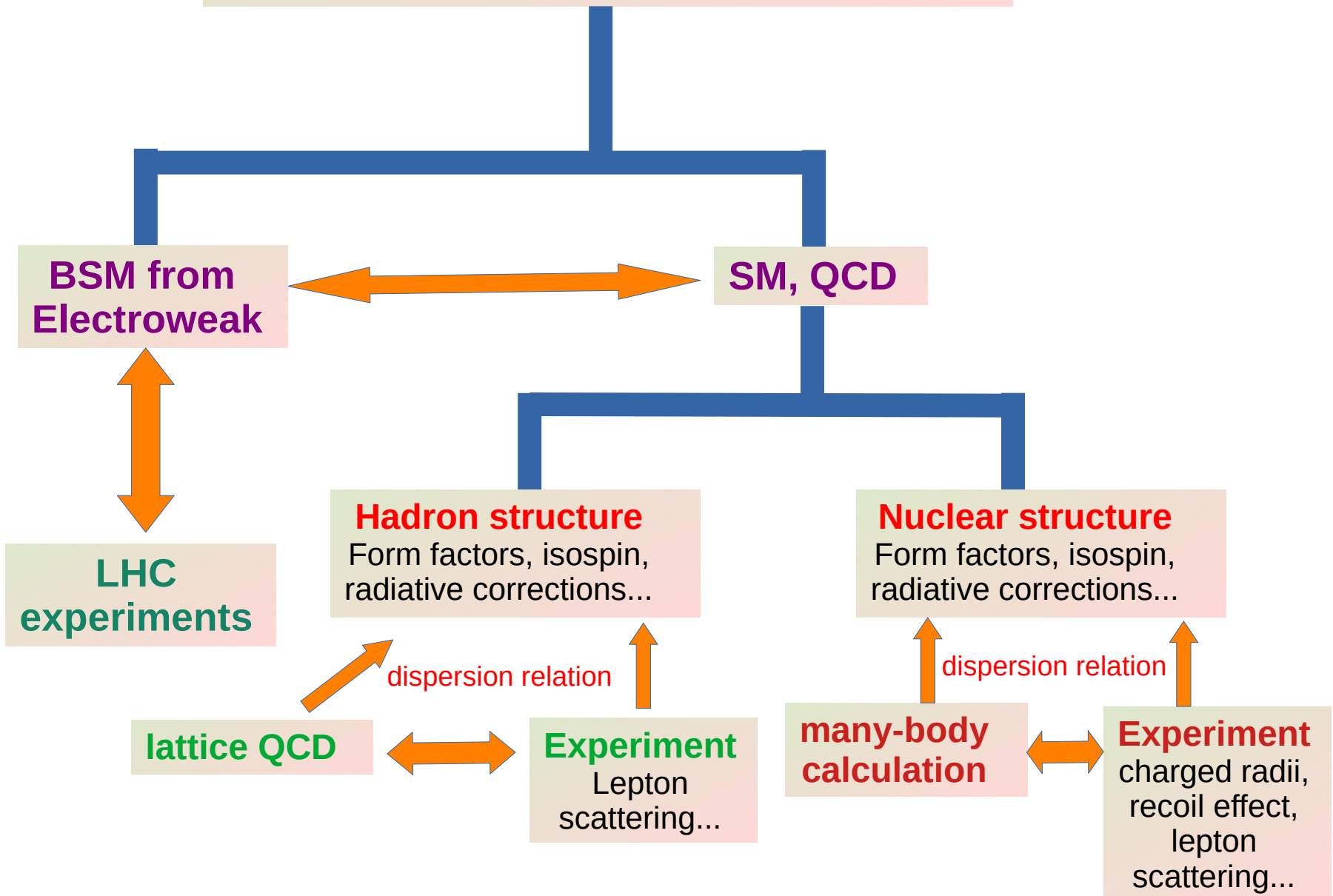
Observables measured at **0.01%** level probes new physics at the scale:

$$\left(\frac{v_H}{\Lambda_{\text{BSM}}} \right)^2 \sim 0.01\% \implies \Lambda_{\text{BSM}} \sim 20 \text{ TeV}$$

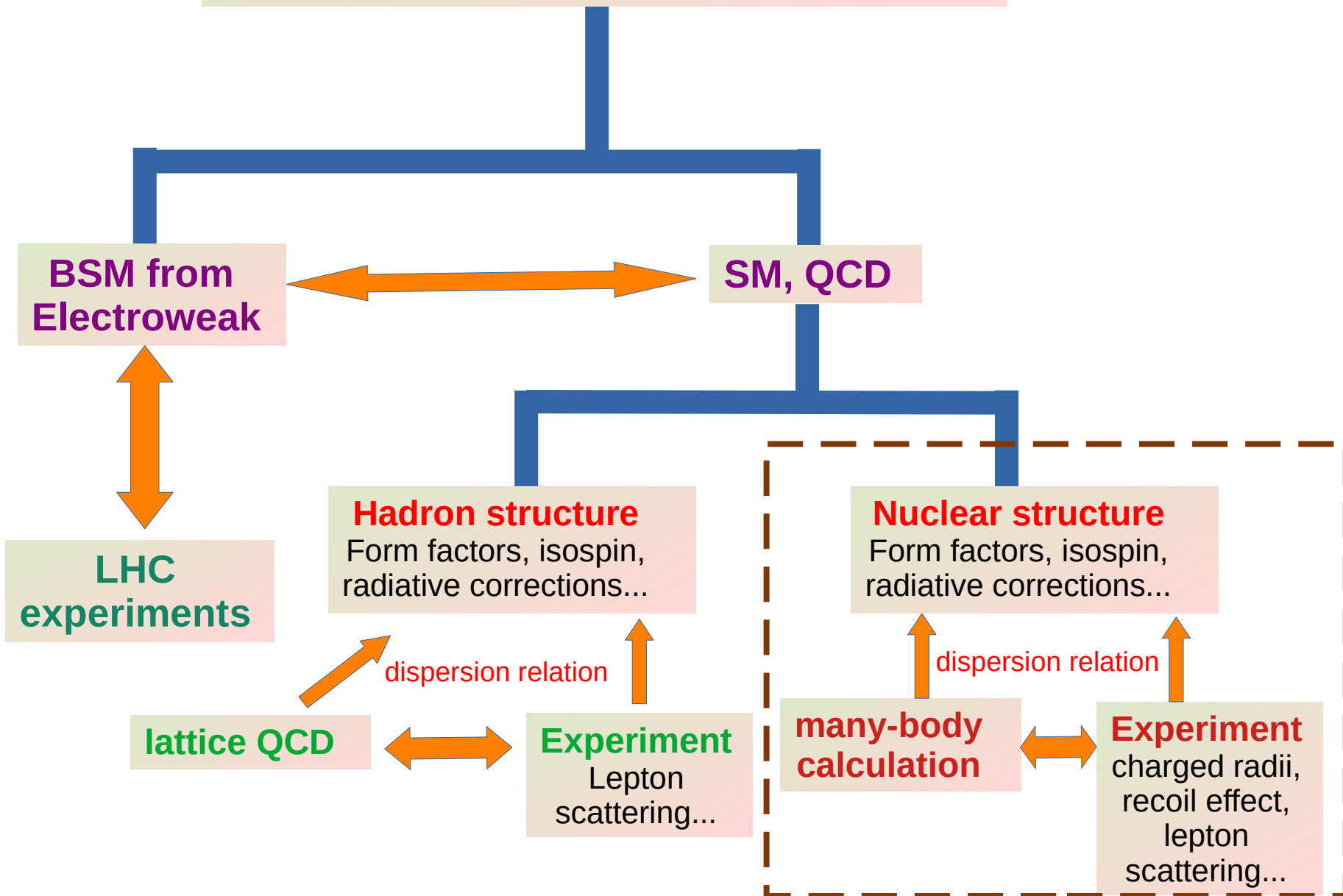
Competitive to high-energy experiments at LHC!

See Vincenzo's talk!

Nuclear Beta Decays as Precision Laboratories

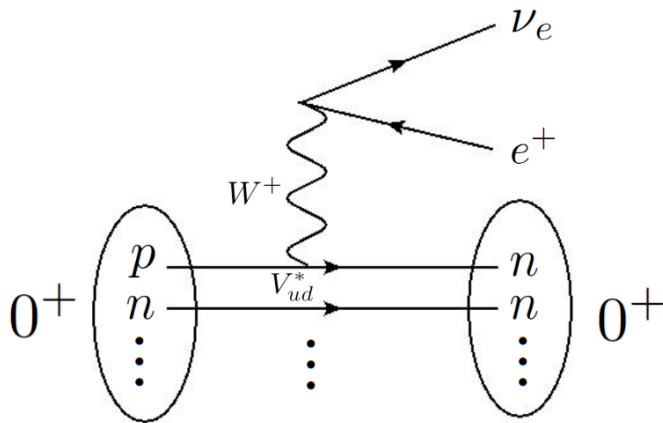


Nuclear Beta Decays as Precision Laboratories



“Superallowed” beta decays of $T=1, J^p=0^+$ nuclei

$$i(0^+) \rightarrow f(0^+) + e^+ + \nu_e$$



Provides the **best measurement of V_{ud}** :

- **23** measured transitions
- **15** with lifetime precision better than **0.23%**

Hardy and Towner, 2020 PRC

$T_z = -1$
${}^6_6\text{C} \rightarrow {}^6_5\text{B}$
${}^{14}_8\text{O} \rightarrow {}^{14}_7\text{N}$
${}^{18}_{10}\text{Ne} \rightarrow {}^{18}_9\text{F}$
${}^{22}_{12}\text{Mg} \rightarrow {}^{22}_{11}\text{Na}$
${}^{26}_{14}\text{Si} \rightarrow {}^{26}_{13}\text{Al}$
${}^{30}_{16}\text{S} \rightarrow {}^{30}_{15}\text{P}$
${}^{34}_{18}\text{Ar} \rightarrow {}^{34}_{17}\text{Cl}$
${}^{38}_{20}\text{Ca} \rightarrow {}^{38}_{19}\text{K}$
${}^{42}_{22}\text{Ti} \rightarrow {}^{42}_{21}\text{Sc}$
${}^{46}_{24}\text{Cr} \rightarrow {}^{46}_{23}\text{V}$
${}^{50}_{26}\text{Fe} \rightarrow {}^{50}_{25}\text{Mn}$
${}^{54}_{28}\text{Ni} \rightarrow {}^{54}_{27}\text{Co}$

$T_z = 0$
${}^{26m}_{13}\text{Al} \rightarrow {}^{26}_{12}\text{Mg}$
${}^{34}_{17}\text{Cl} \rightarrow {}^{34}_{16}\text{S}$
${}^{38m}_{19}\text{K} \rightarrow {}^{38}_{18}\text{Ar}$
${}^{42}_{21}\text{Sc} \rightarrow {}^{42}_{20}\text{Ca}$
${}^{46}_{23}\text{V} \rightarrow {}^{46}_{22}\text{Ti}$
${}^{50}_{25}\text{Mn} \rightarrow {}^{50}_{24}\text{Cr}$
${}^{54}_{27}\text{Co} \rightarrow {}^{54}_{26}\text{Fe}$
${}^{62}_{31}\text{Ga} \rightarrow {}^{62}_{30}\text{Zn}$
${}^{66}_{33}\text{As} \rightarrow {}^{66}_{32}\text{Ge}$
${}^{70}_{35}\text{Br} \rightarrow {}^{70}_{34}\text{Se}$
${}^{74}_{37}\text{Rb} \rightarrow {}^{74}_{36}\text{Kr}$

“Trajectory of V_{ud} ” in the past few years

CKM unitarity
good

Early 2018: 0.97417(21)

Single-nucleon RC
by Marciano and Sirlin, 2006



~4 σ unitarity
violation

Late 2018: 0.97370(14)

Single-nucleon RC revisited
with dispersion relation



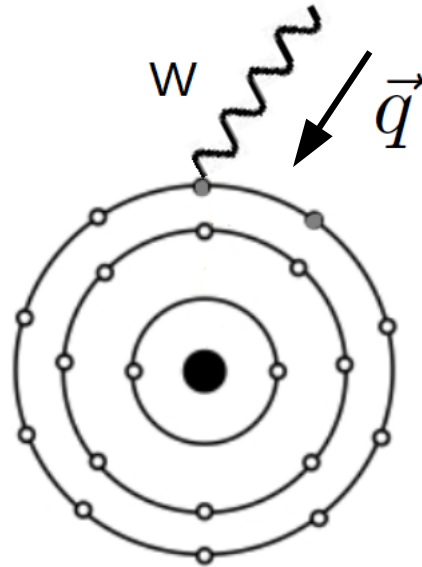
~3 σ unitarity
violation

Late 2020: 0.97373(31)

New nuclear effects discovered,
final review by Hardy/Towner

$$|V_{ud}|^2 + |V_{us}|^2 + |V_{ub}|^2 \stackrel{?}{=} 1$$

Superallowed decays at tree level



Charged weak decay form factor:

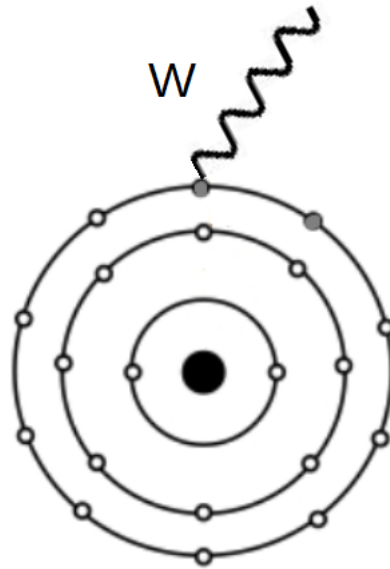
$$f(\vec{q}^2) = M_F (1 + \bar{f}(\vec{q}^2))$$

Fermi matrix element

q^2 -dependence

Nuclear (charged) weak distribution

$\rho_{CW}(r)$: Distribution of “active” nucleons eligible for weak transitions in a nucleus



Old : Impulse approximation in **non-relativistic (NR) nuclear models**.

$$\langle f|O|i\rangle = \sum_{\alpha\beta} \langle \alpha|O|\beta\rangle \langle f|a_{\alpha}^{\dagger}a_{\beta}|i\rangle$$

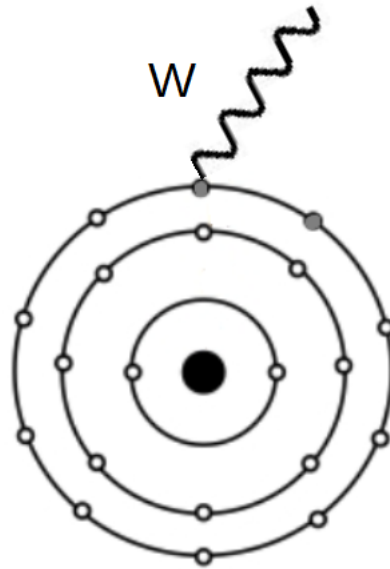
α, β : single-nucleon states

Hardy and Towner, 2005 PRC

How reliable is such approximation at **0.01% level**?

Nuclear (charged) weak distribution

$\rho_{\text{CW}}(r)$: Distribution of “active” nucleons eligible for weak transitions in a nucleus



New: Utilize **CVC** to relate (charged) weak distribution to nuclear **charge distributions**

Hostein, 1974 RMP; CYS, 2023 PRL

$$\begin{aligned}\rho_{\text{CW}}(r) &= \rho_{\text{ch},1}(r) + Z_0 (\rho_{\text{ch},0}(r) - \rho_{\text{ch},1}(r)) \\ &= \rho_{\text{ch},1}(r) + \frac{Z_{-1}}{2} (\rho_{\text{ch},-1}(r) - \rho_{\text{ch},1}(r))\end{aligned}$$

Form factor expansion:

$$F(q^2) = 1 + \frac{q^2}{6} \langle r^2 \rangle + \dots$$

M.S. radius

Data + isospin predicts **unexpectedly large (charged) weak radii!**

$$\begin{aligned} \text{E.g. } \langle r_{\text{Ch}}^2 \rangle_{^{38}\text{Ar}}^{1/2} &= 3.4028(19) \text{ fm} \\ \langle r_{\text{Ch}}^2 \rangle_{^{38}\text{Ca}}^{1/2} &= 3.467(1) \text{ fm} \end{aligned} \quad \longrightarrow \quad \langle r_{\text{CW}}^2 \rangle^{1/2} = 4.00(4) \text{ fm}$$

Simplified shell model prediction: $\langle r_{\text{CW}}^2 \rangle^{1/2} \approx 3.62 \text{ fm}$

$$|V_{ud}|^2 \propto \frac{1}{ft}$$

(Preliminary) data-driven re-analysis of the “**statistical rate function**” f :

Transition	f_{new}	f_{HT}	$\frac{f_{\text{new}} - f_{\text{HT}}}{f_{\text{new}}} (\%)$
$^{18}\text{Ne} \rightarrow ^{18}\text{F}$	134.62(0) _{rad} (0) _{shape} (17) _{Q_{EC}}	134.64(17) _{Q_{EC}}	-0.01(0) _{rad} (0) _{shape}
$^{22}\text{Mg} \rightarrow ^{22}\text{Na}$	418.27(1) _{rad} (1) _{shape} (13) _{Q_{EC}}	418.35(13) _{Q_{EC}}	-0.02(0) _{rad} (0) _{shape}
$^{34}\text{Ar} \rightarrow ^{34}\text{Cl}$	3409.89(16) _{rad} (18) _{shape} (25) _{Q_{EC}}	3410.85(25) _{Q_{EC}}	-0.03(0) _{rad} (1) _{shape}
$^{38}\text{Ca} \rightarrow ^{38m}\text{K}$	5327.49(14) _{rad} (36) _{shape} (31) _{Q_{EC}}	5328.88(31) _{Q_{EC}}	-0.03(0) _{rad} (1) _{shape}
$^{42}\text{Ti} \rightarrow ^{42}\text{Sc}$	7124.3(5.7) _{rad} (0.8) _{shape} (1.4) _{Q_{EC}}	7130.1(1.4) _{Q_{EC}}	-0.08(8) _{rad} (1) _{shape}
$^{50}\text{Fe} \rightarrow ^{50}\text{Mn}$	15053(18) _{rad} (3) _{shape} (60) _{Q_{EC}}	15060(60) _{Q_{EC}}	-0.04(12) _{rad} (2) _{shape}
$^{54}\text{Ni} \rightarrow ^{54}\text{Co}$	21137(3) _{rad} (1) _{shape} (52) _{Q_{EC}}	21137(57) _{Q_{EC}}	+0.00(2) _{rad} (0) _{shape}
$^{34}\text{Cl} \rightarrow ^{34}\text{S}$	1995.076(81) _{rad} (103) _{shape} (94) _{Q_{EC}}	1996.003(96) _{Q_{EC}}	-0.05(0) _{rad} (1) _{shape}
$^{38m}\text{K} \rightarrow ^{38}\text{Ar}$	3296.32(8) _{rad} (21) _{shape} (15) _{Q_{EC}}	3297.39(15) _{Q_{EC}}	-0.03(0) _{rad} (1) _{shape}
$^{42}\text{Sc} \rightarrow ^{42}\text{Ca}$	4468.53(3.36) _{rad} (0.52) _{shape} (0.46) _{Q_{EC}}	4472.46(46) _{Q_{EC}}	-0.09(8) _{rad} (1) _{shape}
$^{50}\text{Mn} \rightarrow ^{50}\text{Cr}$	10737.93(11.50) _{rad} (2.02) _{shape} (0.50) _{Q_{EC}}	10745.99(49) _{Q_{EC}}	-0.08(11) _{rad} (2) _{shape}
$^{54}\text{Co} \rightarrow ^{54}\text{Fe}$	15769.4(2.3) _{rad} (0.7) _{shape} (2.7) _{Q_{EC}}	15766.8(2.7) _{Q_{EC}}	+0.02(1) _{rad} (0) _{shape}
$^{74}\text{Rb} \rightarrow ^{74}\text{Kr}$	47326(127) _{rad} (18) _{shape} (94) _{Q_{EC}}	47281(93) _{Q_{EC}}	+0.10(27) _{rad} (4) _{shape}

CYS, PRELIMINARY RESULTS

Shifts at 0.01% level are observed!

Upward shift of V_{ud} ?

Fermi matrix element and isospin-symmetry breaking (ISB)

$$|M_F|^2 = |\langle f | \hat{\tau}_+ | i \rangle|^2 = 2(1 - \delta_C)$$

$$|V_{ud}|^2 \propto \frac{1}{1 - \delta_C}$$

Starts at **second order** in ISB interaction:

$$H = H_0 + \textcircled{V} \text{ ISB}$$

*Behrends and Sirlin, 1960 PRL;
Ademollo and Gatto, 1964 PRL*

$$\delta_C = \mathcal{O}(V^2) \quad \mathbf{0.1\% \sim 1\%}$$

- Computing δ_C : Classic problem over **6 decades!** *MacDonald, 1958 Phys.Rev*

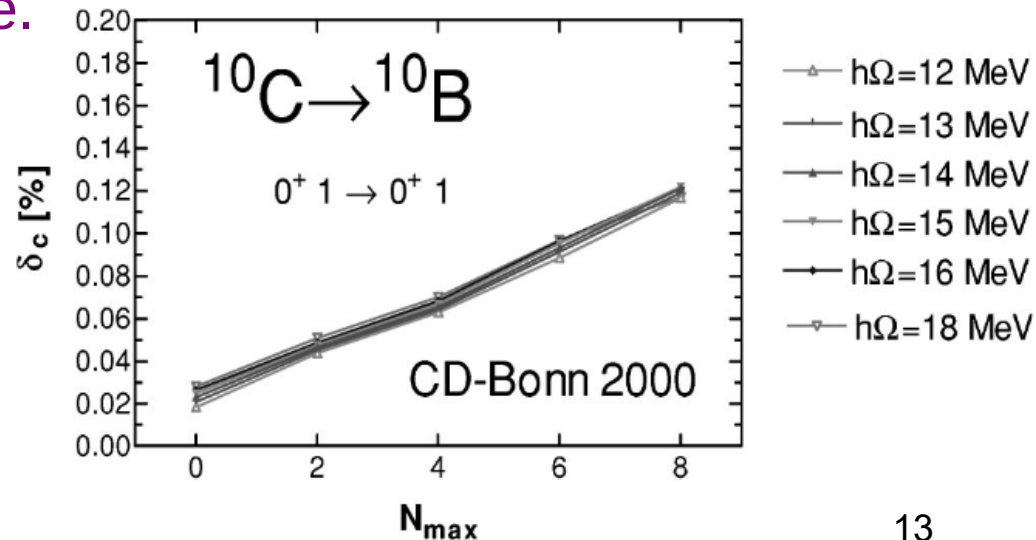
Transitions	δ_C (%)				
	WS	DFT	HF	RPA	Micro
$^{26m}\text{Al} \rightarrow ^{26}\text{Mg}$	0.310	0.329	0.30	0.139	0.08
$^{34}\text{Cl} \rightarrow ^{34}\text{S}$	0.613	0.75	0.57	0.234	0.13
$^{38m}\text{K} \rightarrow ^{38}\text{Ar}$	0.628	1.7	0.59	0.278	0.15
$^{42}\text{Sc} \rightarrow ^{42}\text{Ca}$	0.690	0.77	0.42	0.333	0.18
$^{46}\text{V} \rightarrow ^{46}\text{Ti}$	0.620	0.563	0.38	/	0.21
$^{50}\text{Mn} \rightarrow ^{50}\text{Cr}$	0.660	0.476	0.35	/	0.24
$^{54}\text{Co} \rightarrow ^{54}\text{Fe}$	0.770	0.586	0.44	0.319	0.28

Caveats:

- Significant **model dependence**.
- **No direct experimental constraint**
- Ab-initio calculations still in preliminary stages

Caurier et al., 2002 PRC

(Selected results)



Probing isospin mixing in experiment?

Nuclear mass splitting: $\delta E_n = {}_0\langle n|V|n\rangle_0$ **X**

$$|n\rangle = |n\rangle_0 + \underbrace{\sum_{k \neq n} |k\rangle_0 \frac{{}_0\langle k|V|n\rangle_0}{E_n^{(0)} - E_k^{(0)}}}_{\text{We need this}} + \dots$$

We need this

Second-class current: doubly-suppressed (kinematics*ISB), null results so far *Minamisono et al., 2011 PRC*

$$V_\mu = \bar{u}(p_2) \left(g_V \gamma_\mu - i \frac{g_M - g_V}{2M} \sigma_{\mu\nu} q^\nu + \frac{g_S}{2M} q_\mu \right) u(p_1),$$

$$A_\mu = -\bar{u}(p_2) \gamma_5 \left(g_A \gamma_\mu - i \frac{g_T}{2M} \sigma_{\mu\nu} q^\nu + \frac{g_P}{2M} q_\mu \right) u(p_1).$$

Isospin mixing can be probed through measurements of **electroweak nuclear radii**

The non-zeroneess of

$$\Delta M_A^{(1)} \equiv -\underbrace{\langle r_{\text{cw}}^2 \rangle}_{\substack{\text{cw radius} \\ \text{(directly} \\ \text{measured)}}} + \underbrace{(\langle r_{n,1}^2 \rangle - \langle r_{p,1}^2 \rangle)}_{\text{related to neutron skin}}$$

$$\Delta M_B^{(1)} \equiv \frac{1}{2} \underbrace{(Z_1 \langle r_{\text{ch},1}^2 \rangle + Z_{-1} \langle r_{\text{ch},-1}^2 \rangle)}_{\text{THREE charge radii}} - Z_0 \langle r_{\text{ch},0}^2 \rangle$$

signify ISB. No double suppression!

Seng and Gorchtein, 2023 PLB

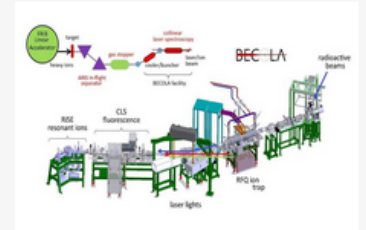
A	$R_{\text{Ch},-1}$ (fm)	$R_{\text{Ch},0}$ (fm)	$R_{\text{Ch},1}$ (fm)
10	^{10}C	$^{10}\text{B}(\text{ex})$	^{10}Be : 2.3550(170) ^a
14	^{14}O	$^{14}\text{N}(\text{ex})$	^{14}C : 2.50 25(87) ^a
18	^{18}Ne : 2.9714(76) ^a	$^{18}\text{F}(\text{ex})$	^{18}O : 2.77 26(56) ^a
22	^{22}Mg : 3.0691(89) ^b	$^{22}\text{Na}(\text{ex})$	^{22}Ne : 2.9525(40) ^a
26	^{26}Si	^{26m}Al	^{26}Mg : 3.0337(18) ^a
30	^{30}S	$^{30}\text{P}(\text{ex})$	^{30}Si : 3.1336(40) ^a
34	^{34}Ar : 3.3654(40) ^a	^{34}Cl	^{34}S : 3.2847(21) ^a
38	^{38}Ca : 3.467(1) ^c	^{38m}K : 3.437(4) ^d	^{38}Ar : 3.4028(19) ^a
42	^{42}Ti	^{42}Sc : 3.5702(238) ^a	^{42}Ca : 3.5081(21) ^a
46	^{46}Cr	^{46}V	^{46}Ti : 3.6070(22) ^a
50	^{50}Fe	^{50}Mn : 3.7120(196) ^a	^{50}Cr : 3.6588(65) ^a
54	^{54}Ni : 3.738(4) ^e	^{54}Co	^{54}Fe : 3.6933(19) ^a
62	^{62}Ge	^{62}Ga	^{62}Zn : 3.9031(69) ^b
66	^{66}Se	^{66}As	^{66}Ge
70	^{70}Kr	^{70}Br	^{70}Se
74	^{74}Sr	^{74}Rb : 4.1935(172) ^b	^{74}Kr : 4.1870(41) ^a

DOE highlight:

July 7, 2023

Nuclear Charge Distribution Measurements May Solve Outstanding Puzzle In Particle Physics

By reanalyzing the distribution of active protons in nuclei, researchers found a possible solution to a particle physics puzzle involving quarks.



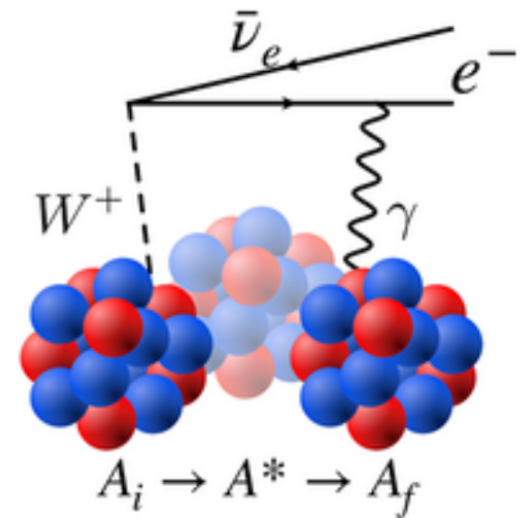
Experimental opportunities:

- Nuclear charge radii measurements: **FRIB, TRIUMF, PSI** (muX) ...
- Recoil effects in nuclear beta decay: **FRIB** (SALER)...
- Neutron skin: **JLab, Mainz** (P2)...

Superaligned decays at loop level

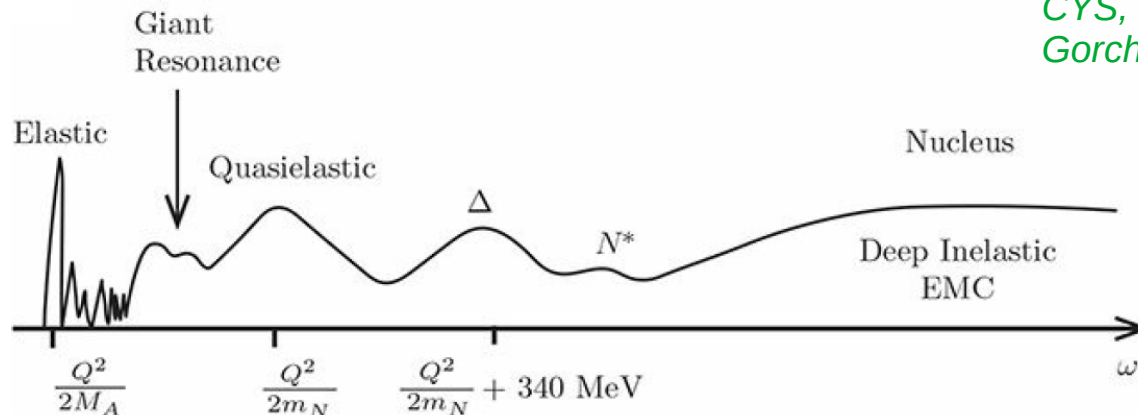
One-loop **radiative corrections** probe QCD at all scales:

- Parton scale ($\gg 1\text{GeV}$) : Theoretically calculable
- Hadron scale ($\sim 1\text{GeV}$) : Data from lepton-nucleon scattering + lattice QCD
- Nuclear scale ($\sim 100\text{ MeV}$) : Most uncertain, requires many-body calculations**



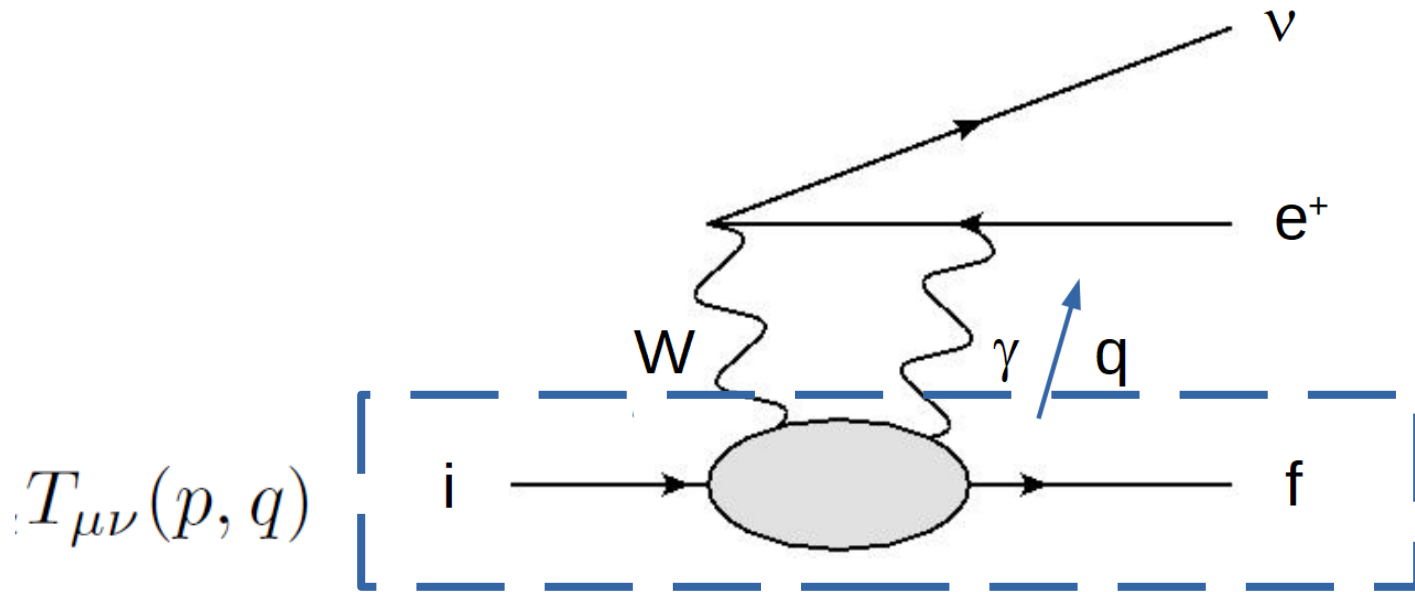
Previous shell-model-based calculations had missed important nuclear effects

CYS, Gorchtein and Ramsey-Musolf, 2019 PRD; Gorchtein, 2019 PRL



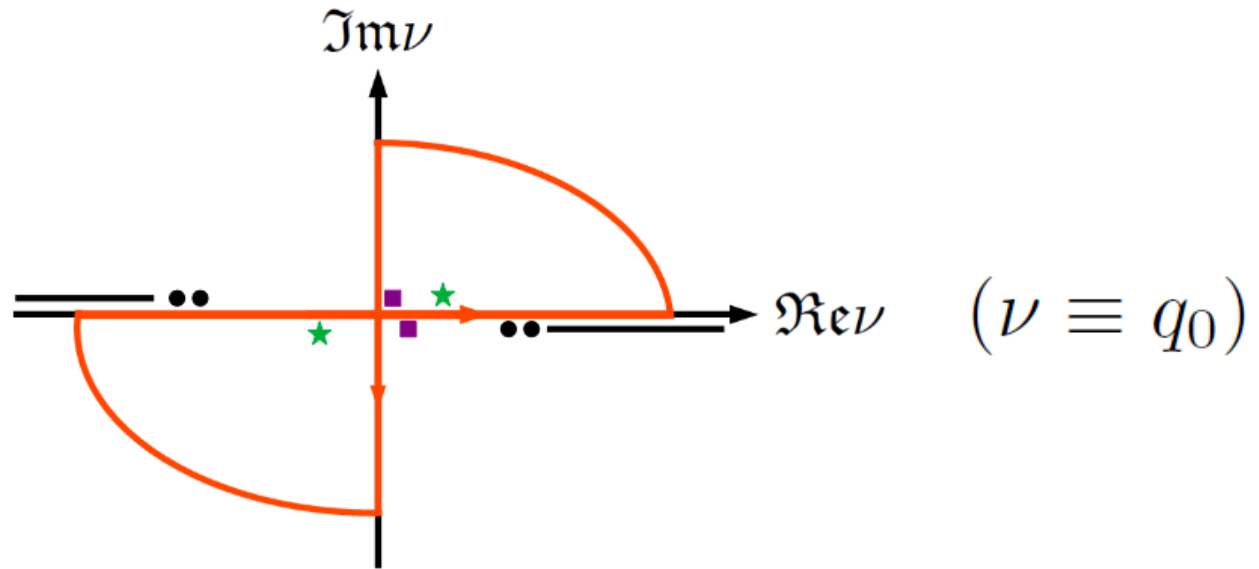
“Generalized forward Compton tensor”:

$$\begin{aligned}
 T^{\mu\nu}(p, q) &= \frac{1}{2} \int d^4x e^{iq \cdot x} \langle \phi_f(p) | T [J_{\text{em}}^\mu(x) \left(J_W^{\dagger\nu}(0) \right)_A] | \phi_i(p) \rangle \\
 &= \frac{i\epsilon^{\mu\nu\alpha\beta} p_\alpha q_\beta}{2M\nu} T_3(\nu, Q^2)
 \end{aligned}$$



The loop diagram depends on the **invariant function T_3**

Performing the q_0 -integral using **Wick rotation**



$$\square_{\gamma W}^b = \underbrace{(\square_{\gamma W}^b)_{\text{Wick}} + (\square_{\gamma W}^b)_{\text{res},e}}_{\text{Regular to } E_e} + \underbrace{(\square_{\gamma W}^b)_{\text{res},T_3}}_{\text{"Residue contribution", singular to } E_e}$$

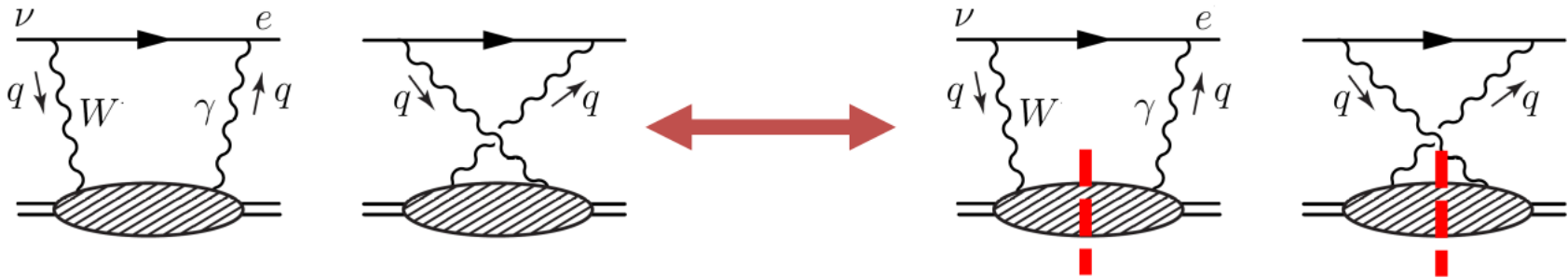
Regular to E_e

"Residue contribution",
singular to E_e

Expansion of the **“regular” terms** in powers of electron energy:

$$(\square_{\gamma W}^b)_{\text{Wick}} + (\square_{\gamma W}^b)_{\text{res},e} = \Xi_0 + \Xi_1 E_e + \mathcal{O}(E_e^2)$$

The coefficients can be related to **nuclear response functions** through **dispersion relation** *Seng and Gorchtein, 2023 PRC*



$$R^{xy}(\nu, Q^2) \sim \sum_X \delta(M + q_0 - E_X) \langle f | J_{\text{em}}^x | X \rangle \langle X | J_W^y | i \rangle$$

which serves as starting point for **ab-initio nuclear theory calculation**

- **Light nuclei:** NCSM, QMC...
- **Medium nuclei:** IMSRG, Coupled-Cluster, NLEFT...

Residue contribution

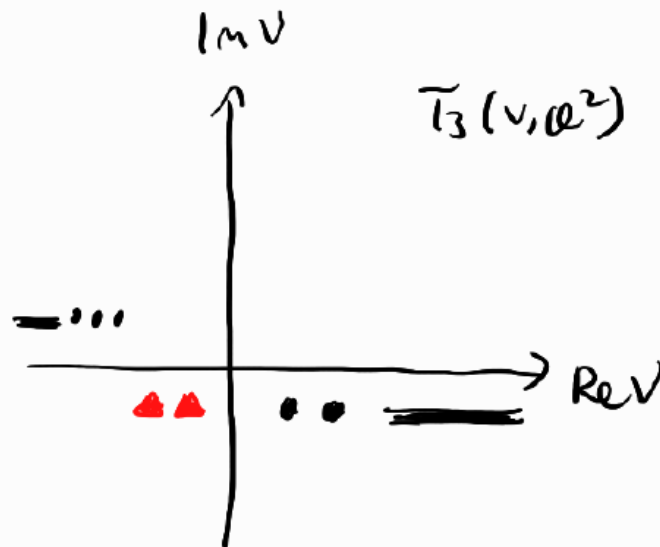
(First pointed out by Michael Gennari, TRIUMF)

For $A=10,14,18,22,26,30$ and 38 , the $T_z=0, J^p=0^+$ state is an excited state.

$T_z=0$ initial nucleus \rightarrow extra poles of T_3 in the first quadrant;

$T_z=0$ final nucleus \rightarrow extra poles of T_3 in the third quadrant

For $^{10}\text{C} \rightarrow ^{10}\text{B}$:



Excited Nuclear States for B-10 (Boron)

Energy levels

E^*	J^π	T	l_p
-------	---------	-----	-------

[keV]

0.0	3^+	0	1
718.35(4)	1^+	0	1
1740.2(2)	0^+	1	1
2154.3(5)	1^+	0	1
3587.1(5)	2^+	0	1
4774.0(5)	3^+	0	1

22

$$\nu_k = \sqrt{M_k^2 + \mathbf{q}^2} - M - i\varepsilon$$

In third quadrant if $\mathbf{q}^2 < M^2 - M_k^2$

Residue of T_3 due to a low-lying $J^P=1^+$ state:

$$\text{Res}T_3 \propto \underbrace{\langle f(0^+) | J_{\text{em}}^x(\vec{q}) | f(1^+) \rangle}_{\text{M1-transition}} \underbrace{\langle f(1^+) | J_W^y(-\vec{q}) | i(0^+) \rangle}_{\text{GT-transition}}$$

Matrix elements can be inferred from **M1** and **GT**-transition rates!

	GT, $\log_{10} ft$ (s)	M1, $t_{1/2}$
$^{10}\text{C} \rightarrow ^{10}\text{B}$	3.0426(7)	4.9(2.1) fs
$^{14}\text{O} \rightarrow ^{14}\text{N}$	7.279(8)	68(3) fs
$^{18}\text{Ne} \rightarrow ^{18}\text{F}$	3.091(4)	1.77(31) fs
$^{22}\text{Mg} \rightarrow ^{22}\text{Na}$	3.64	19.6(7) ps
$^{30}\text{S} \rightarrow ^{30}\text{P}$	4.322(11)	96(10) fs

From IAEA / NNDC website

Future experimental opportunities??

Ab-initio study of $^{10}\text{C} \rightarrow ^{10}\text{B}$ radiative corrections with no-core shell model

Collaboration: Michael Gennari, Petr Navratil (TRIUMF)

$$|V_{ud}|^2 \propto \frac{1}{1 + \delta_{NS}}$$

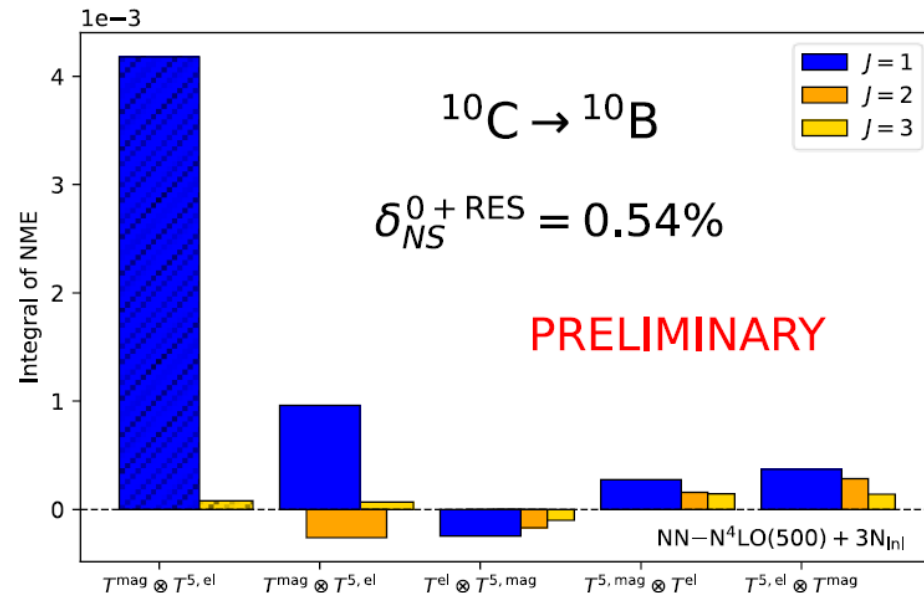
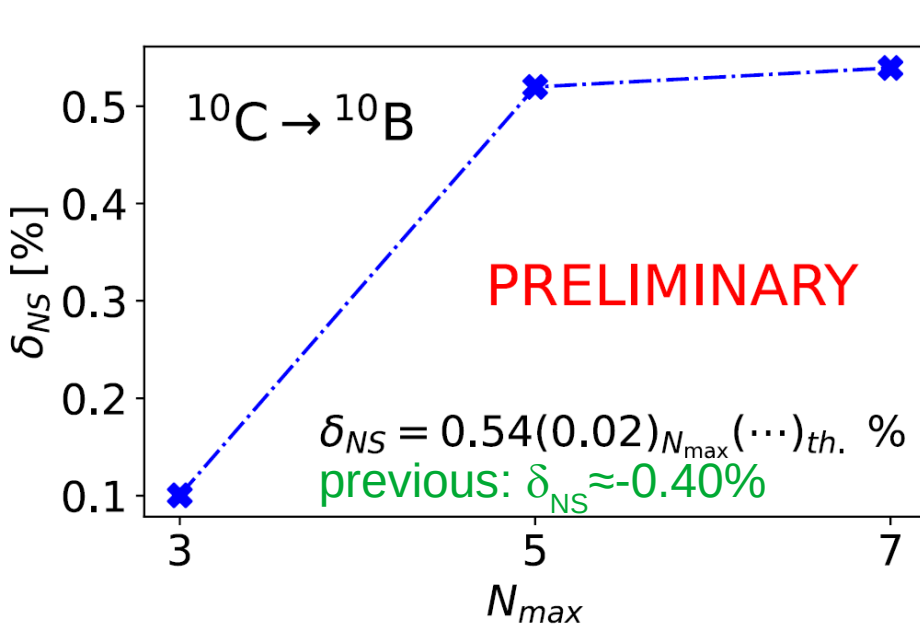


Image credit: Michael Gennari

Dominated by the “**residue contribution**”, which agrees with experiment, but was **completely absent in earlier studies!**

Summary

- **Nuclear beta decays** offer precision tests of SM
 - **Various anomalies** hint at the existence of **BSM physics**
 - **High-precision calculations** test our understanding of **QCD**
- Theory foundations are developed to relate the nuclear decay form factor (Fermi matrix element and q^2 -dependence) to measurable **electroweak nuclear radii**
- Rigorous dispersive formalism is constructed to enable **ab-initio studies** of **nucleus-dependent radiative corrections**; **M1** and **GT** rates for light nuclei are useful experimental inputs
- Research may foster new theory-experiment interplay, and could significantly alter the landscape of precision physics in the charged weak sector

Thanks for your attention!