

Nonresonant Searches for Axion-Like Particles in Vector Boson Scattering Processes at the LHC

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Based on: J. Bonilla, I. Brivio, J. Machado-Rodríguez, J. F. de Trocóniz

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Axion-Like Particles

- Axion-Like Particles (or ALPs) are **neutral pseudo scalar** pseudo Goldston Bosons
- **Effective Field Theory** (EFT) consistent with SM gauge and CP symmetries

- Either **shift-invariant** and/or **anomalous couplings** interactions

$$\left\{ \begin{array}{l} \frac{\partial_\mu a}{f_a} \bar{\psi} \gamma^\mu \psi \\ \frac{a}{f_a} X_{\mu\nu} \tilde{X}^{\mu\nu} \end{array} \right.$$

- ALP interactions with SM particles have a **derivative character**: they grow with momentum

Axion-Like Particles

$$\mathcal{L}_{ALP} \supset -c_{\tilde{B}} \frac{a}{f_a} B_{\mu\nu} \tilde{B}^{\mu\nu} - c_{\tilde{W}} \frac{a}{f_a} W_{\mu\nu}^i \tilde{W}^{i\mu\nu} - c_{\tilde{G}} \frac{a}{f_a} G_{\mu\nu}^A \tilde{G}^{A\mu\nu}$$

- Classical searches: ALP couplings to gluons and photons
- ALP couplings to **EWK bosons**: WW , ZZ , and $Z\gamma$
- **Depend on two parameters**
- ALP-gauge interactions at ATLAS and CMS:
 - Mono-X
 - Resonant
 - **New idea: nonresonant ALP searches**

$$\left\{ \begin{array}{l} g_{a\gamma\gamma} = \frac{4}{f_a} (s_\theta^2 c_{\tilde{W}} + c_\theta^2 c_{\tilde{B}}) \\ g_{a\gamma Z} = \frac{4}{f_a} s_{2\theta} (c_{\tilde{W}} - c_{\tilde{B}}) \\ g_{aZZ} = \frac{4}{f_a} (c_\theta^2 c_{\tilde{W}} + s_\theta^2 c_{\tilde{B}}) \\ g_{aWW} = \frac{4}{f_a} c_{\tilde{W}} \end{array} \right. \quad \theta: \text{Weinberg angle}$$

↑
Imposed by gauge invariance

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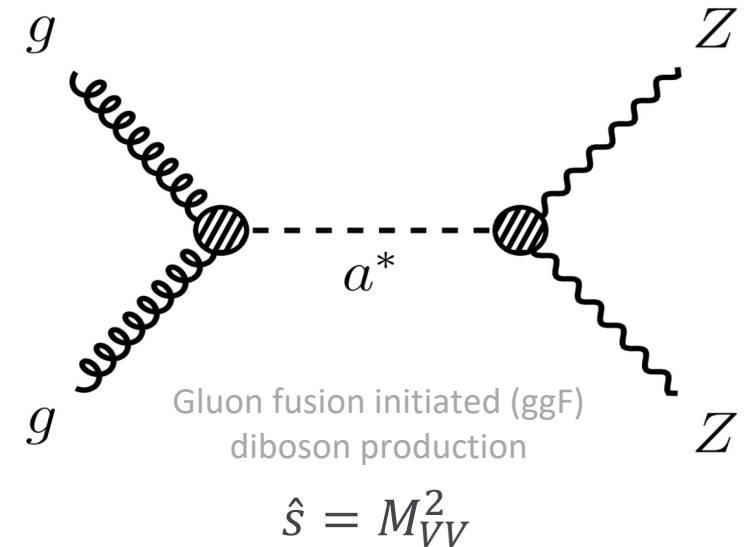
↑
Imposed by gauge invariance

A Novel Approach: Nonresonant ALP-mediated diboson production

- M.B. Gavela, J.M. No, V. Sanz and J.F. de Trocóniz [1905.12953]
- ALP acts as a **very off-shell mediator** $\longrightarrow m_a^2 \ll \hat{s}$
- Signals **independent of ALP mass m_a and its decay width Γ_a** up to $m_a \lesssim 100$ GeV: allows to explore large areas in the parameter space
- Suppression from \hat{s} **compensated by derivative character** of ALP interactions

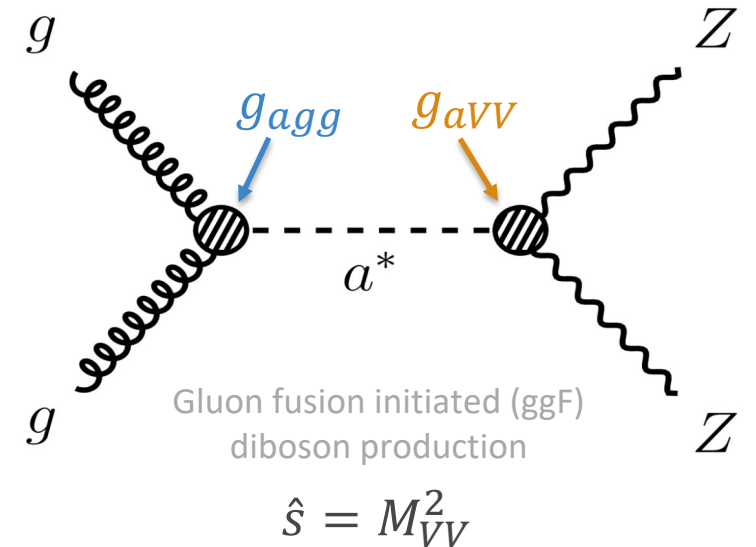
$$\hat{\sigma} \propto \hat{s}/f_a^4$$

Harder scaling



A Novel Approach: Nonresonant ALP-mediated diboson production

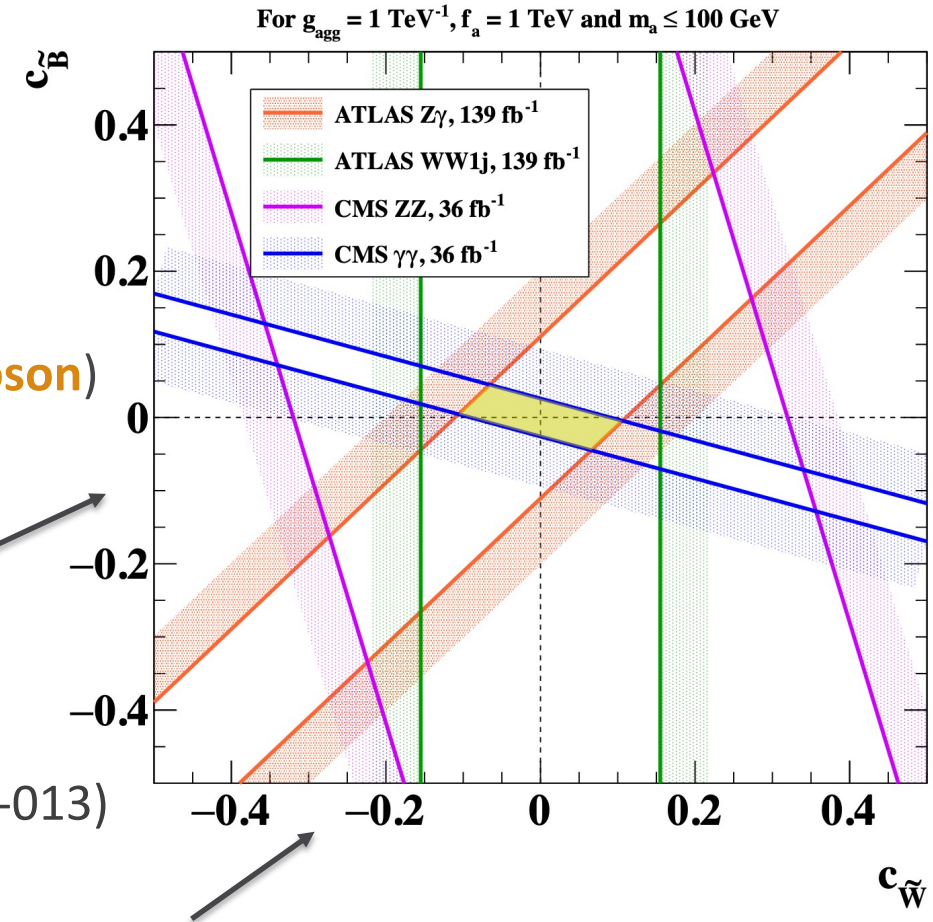
- Reinterpretation of CMS analyses:
 - $gg \rightarrow ZZ$ (CMS-B2G-17-013)
 - $gg \rightarrow \gamma\gamma$ (CMS-EXO-17-017)
- Sensitive to (ALP coupling to gluons x ALP coupling to EWK diboson)
$$g_{agg} \times g_{aVV}$$
- Cross-sections large enough to constrain significantly the theoretical models using Run 2 data.
- **Dedicated ALP search at CMS:** $gg \rightarrow a^* \rightarrow ZZ/ZH$ (CMS-B2G-20-013)



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S. Carrá, V. Goumarre, R. Gupta, S. Heim, B. Heinemann, J. Kuchler, F. Meloni, P. Quilez and Y.C. Yap
 reinterpretation of ATLAS analyses, [2106.10085]: $Z\gamma$ (CERN-EP-2019-228), WWj (CERN-EP-2021-030)

Nonresonant ALPs in VBS

- Vector Boson Scattering (VBS): $q_1 q_2 \rightarrow V'_1 V'_2 q'_1 q'_2$

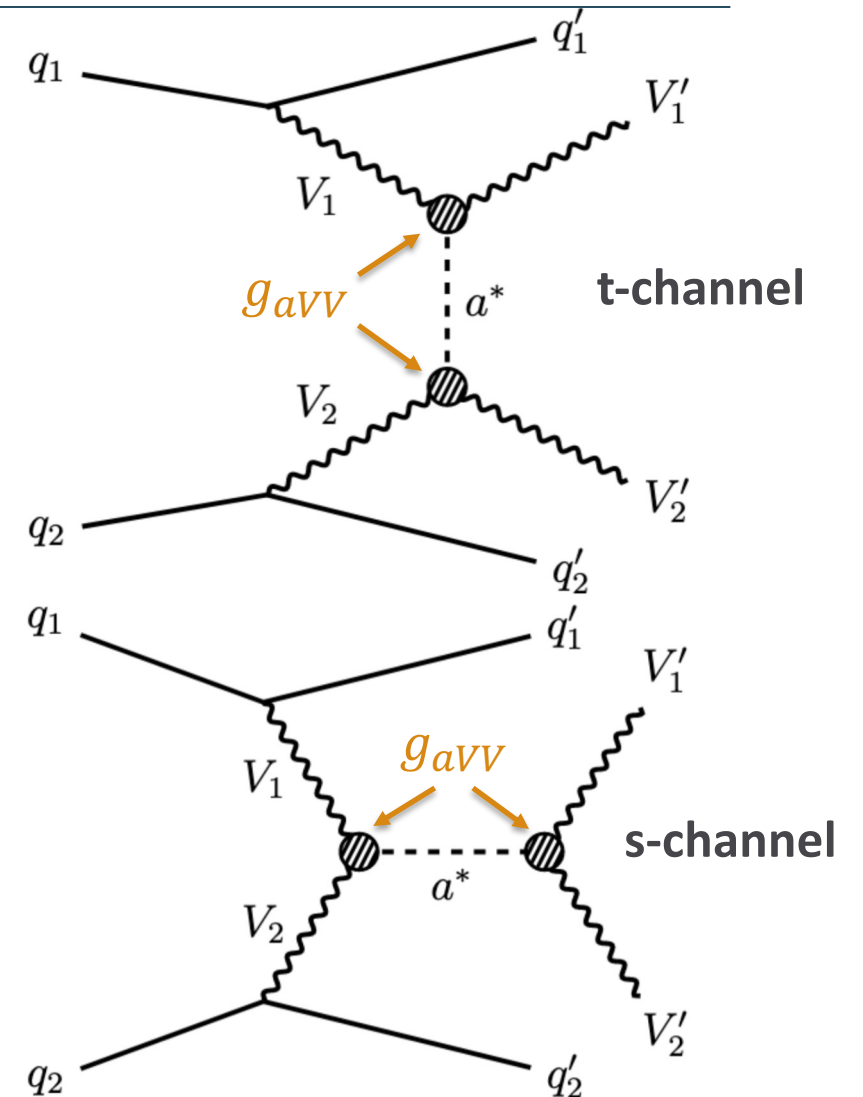
$$V'_1 V'_2 = ZZ, Z\gamma, W^\pm \gamma, WZ, W^\pm W^\pm$$

- VBS limits on ALP couplings to vector boson independently of the gluon coupling

$$\longrightarrow c_{\tilde{B}} \quad c_{\tilde{W}} \quad \cancel{c_{\tilde{G}}} \longrightarrow 2 \text{ parameters}$$

- Nonresonant ALP: independent of ALP mass m_a and its decay width up to $m_a \lesssim 100 \text{ GeV}$

Why VBS?



Nonresonant ALPs in VBS

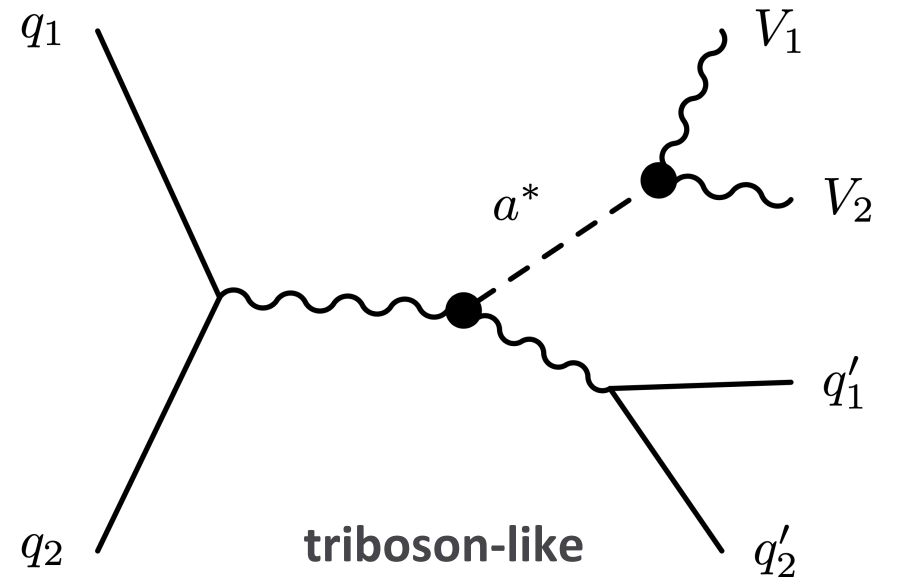
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Efficiently suppressed by VBS cut:
 $M_{q_1 q_1 q_2} > 120 \text{ GeV}$

Why VBS?

Nonresonant ALPs in VBS

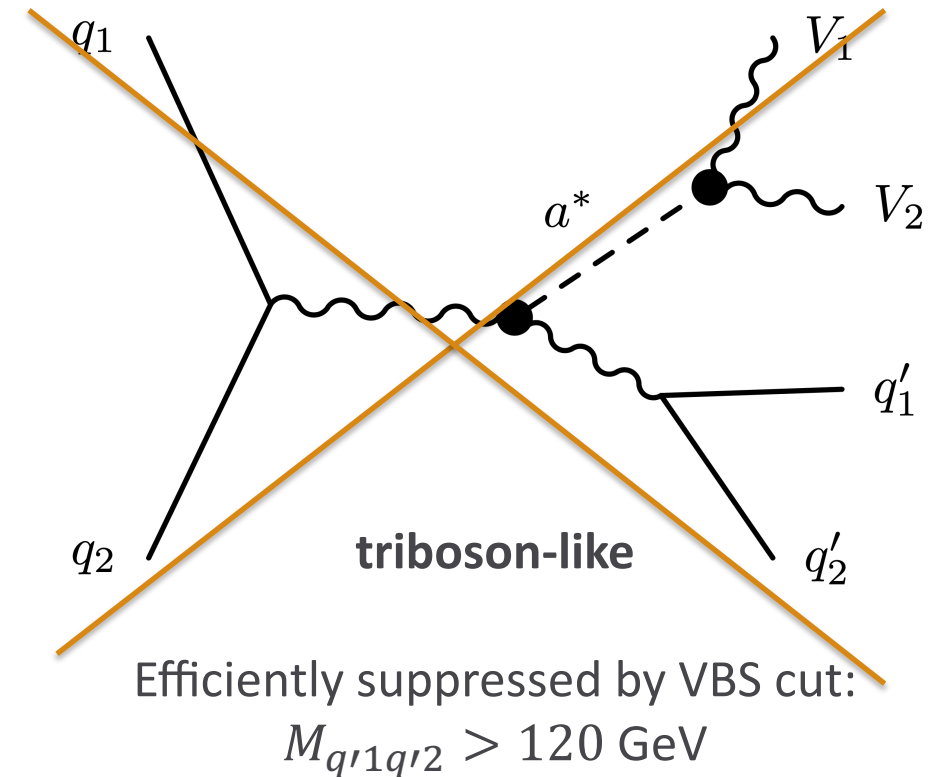
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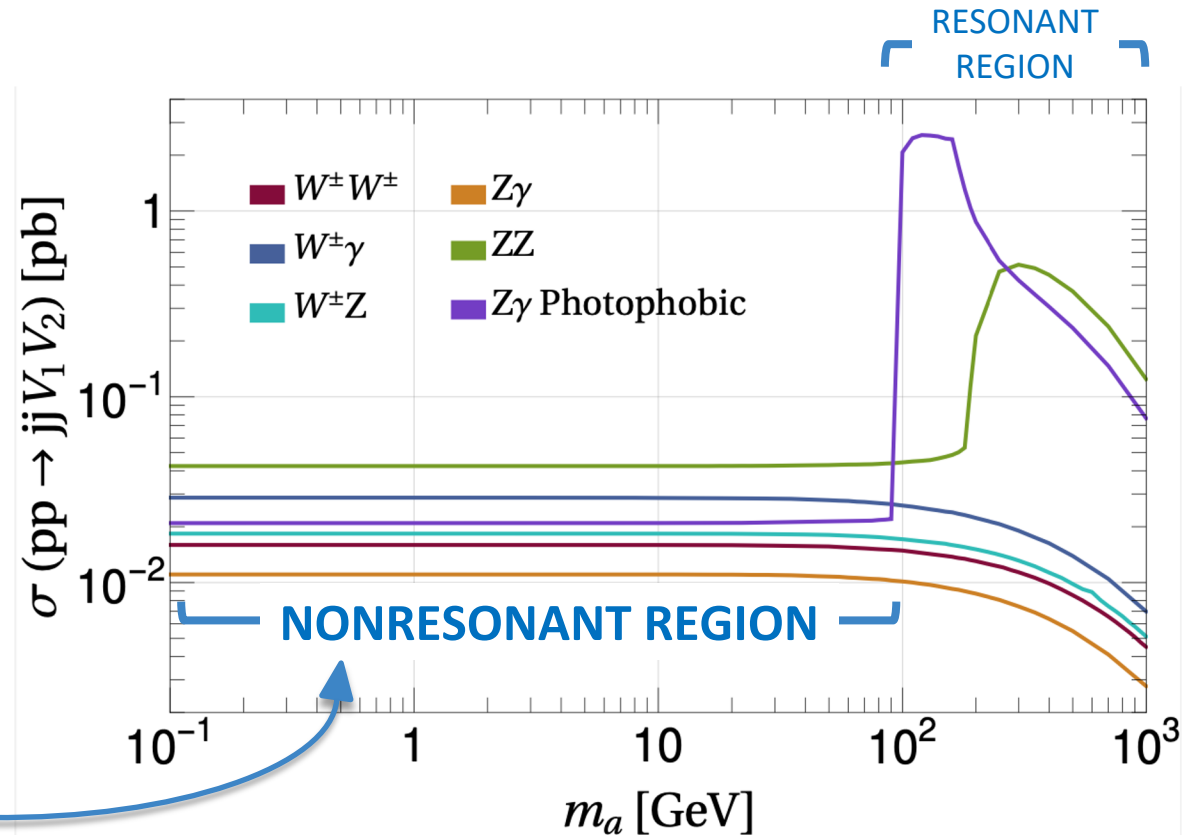
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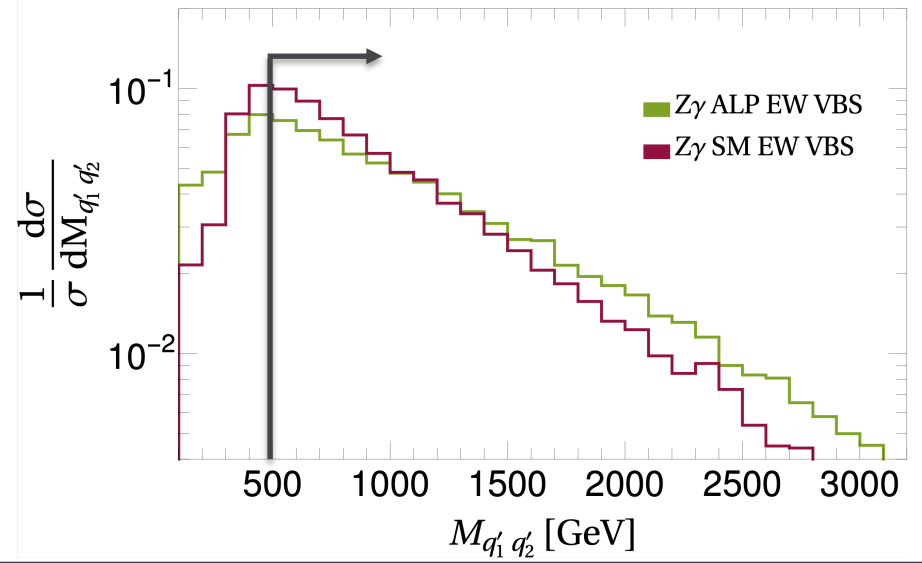
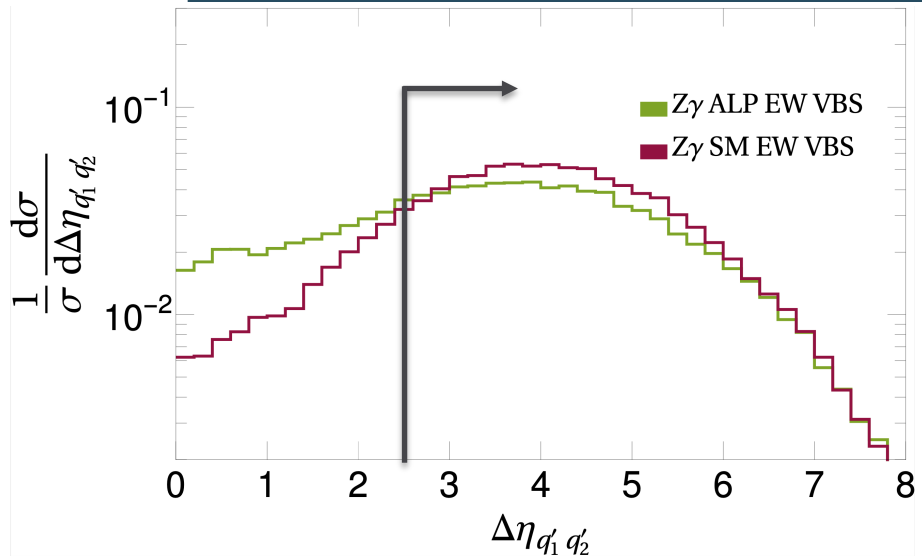
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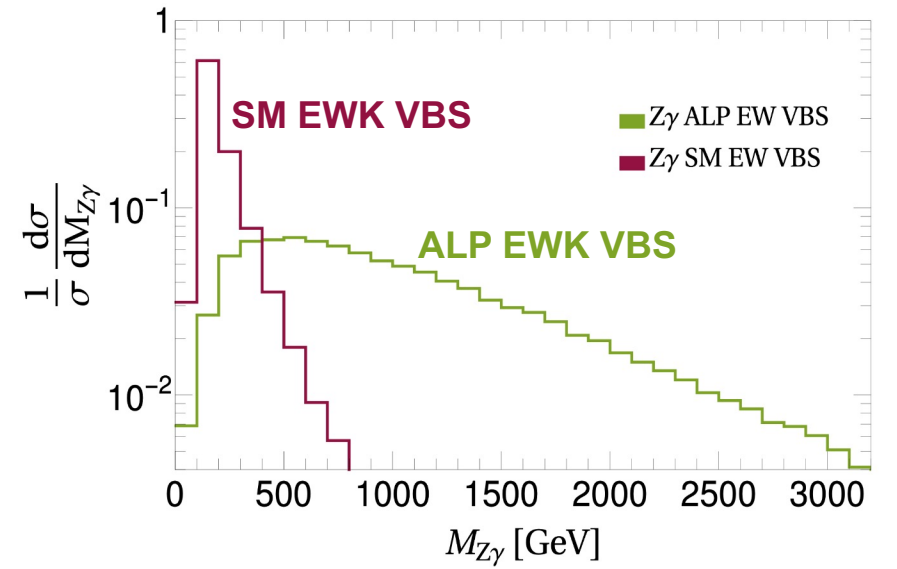


Wide rapidity separation

VBS CHARACTERISTIC OBSERVABLES

Large dijet mass

Diboson Mass



- Derivative nature of ALP interactions: **deviations in the tail** of the **diboson mass** with respect to the SM

ALP Diboson Mass in CMS Leptonic Analyses

- ATLAS/CMS Run 2 measurements: first comparison to data, calibration of simulation tools and calculation of educated predictions for higher luminosities.
- **Reinterpretation of five CMS VBS analyses** with lepton/photon final states:
 - ZZ: CMS-SMP-20-01
 - $W^\pm W^\pm$ and WZ: CMS-SMP-19-012
 - $Z\gamma$: CMS-SMP-20-016
 - $W\gamma$: CMS-SMP-19-008
- Look at **high energy deviations** in the tail of the transverse momentum/mass spectra
- Selections cuts, data and backgrounds in the CMS papers
- **Generation of ALP VBS**: MadGraph_aMC@NLO + Pythia8 + Delphes3

ALP Diboson Mass in CMS Leptonic Analyses

- **Calibrate** our **Delphes detector simulation** using the SM EWK channel \longrightarrow **EWK scale factor**

- Compare EWK SM VBS expected yields from the CMS simulation and ours \longrightarrow

$$\rho = \frac{N_{SM\ EWK}^{simulated}}{N_{SM\ EWK}^{CMS}}$$

Channel	Obs.	Lum. [fb^{-1}]	Selection Criteria	ρ
ZZ	M_{ZZ}	137	$M_{jj} > 100 \text{ GeV}$	0.8 ± 0.1
$Z\gamma$	$M_{Z\gamma}$	137	$M_{jj} > 500 \text{ GeV}, \Delta\eta_{jj} > 2.5, p_T^\gamma > 120 \text{ GeV}$	1.4 ± 0.2
$W^\pm\gamma$	$M_{W\gamma}$	35.9*	$M_{jj} > 800 \text{ GeV}, \Delta\eta_{jj} > 2.5, p_T^\gamma > 100 \text{ GeV}$	3.1 ± 0.5
$W^\pm Z$	M_{WZ}^T	137	$M_{jj} > 500 \text{ GeV}, \Delta\eta_{jj} > 2.5$	1.5 ± 0.4
$W^\pm W^\pm$	M_{WW}^T	137	$M_{jj} > 500 \text{ GeV}, \Delta\eta_{jj} > 2.5$	1.3 ± 0.2

Table 3. Summary of the CMS VBS analyses: the diboson mass observable, the integrated luminosity, the most important selection criteria and the normalization scale factor ρ .

* Recent update at 138 fb^{-1} : CMS PAS SMP-21-011

ALP Diboson Mass in CMS Leptonic Analyses

- **~20 % signal systematics**: PDFs + renormalization and factorization scales + MadGraph@aMC
- **Background uncertainties** from CMS analyses
- Consistency of the ALP EFT and estimation of the impact of the highest-energy bins
 - ➔ **upper cut on diboson mass M_{VV}**
- Two benchmarks:
 - $M_{VV} < 2$ TeV: ~85 % efficiency
 - $M_{VV} < 4$ TeV: >99 % efficiency

ALP Diboson Mass in CMS Leptonic Analyses

	$c_{\tilde{W}} = c_{\tilde{B}}$ signal / interf. [fb]	Photophobic signal / interf. [fb]	Expected Lepton Events	Int. lum. [fb ⁻¹]
ZZ	42.4 / -13.5	18.5 / -9.3	9.3 / -3.2	137
WZ	18.4 / 1.7	23.9 / -0.14	4.2 / 0.05	137
$W^\pm W^\pm$	16.0 / -4.0	16.0 / -4.0	18 / -5.5	137
$W\gamma$	28.7 / 4.3	5.4 / 1.7	3.6 / -0.04	35.9
$Z\gamma$	11.1 / 0.3	20.9 / -9.1	15.1 / 0.07	137

ALP Diboson Mass in CMS Leptonic Analyses

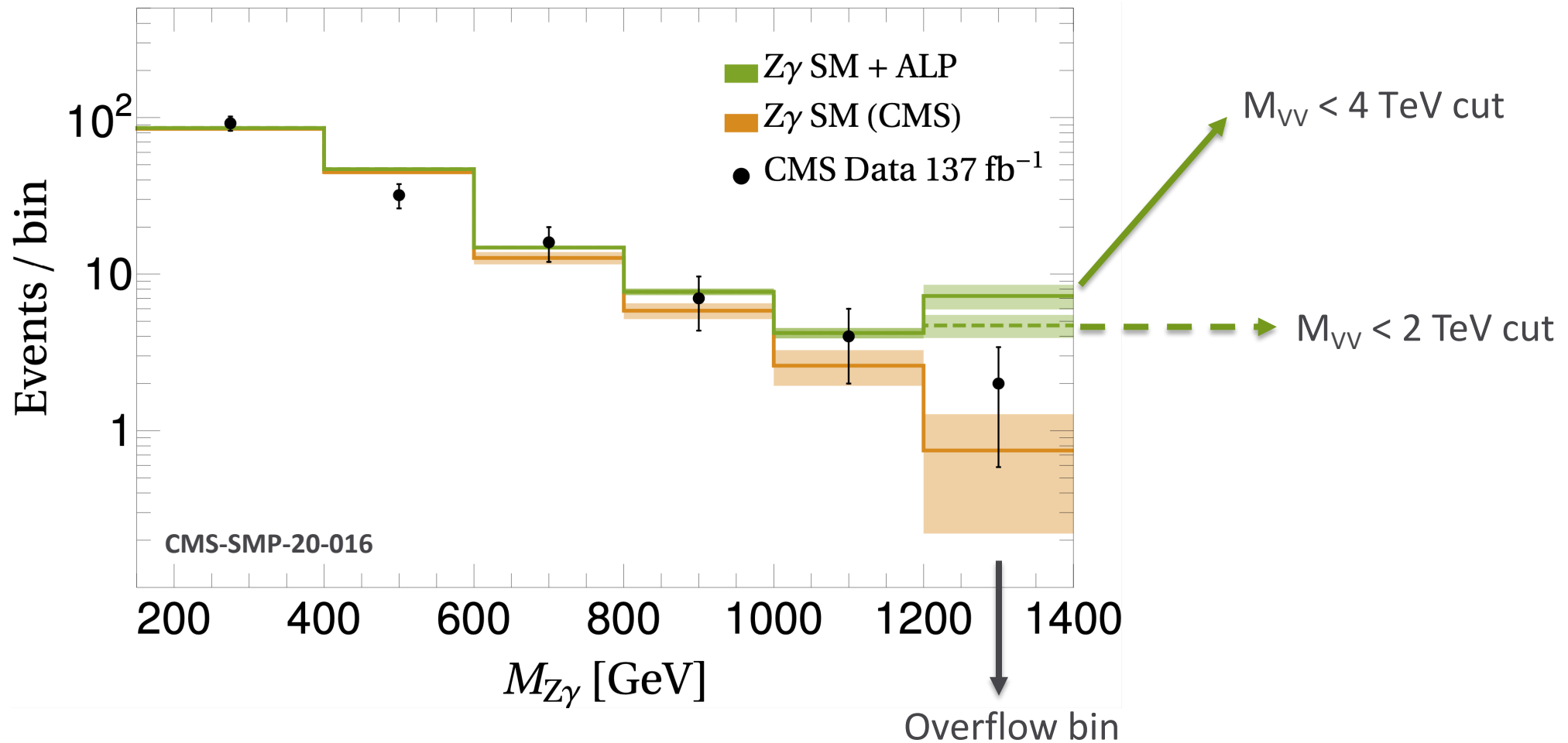
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$g_{\alpha\gamma Z} = 0$

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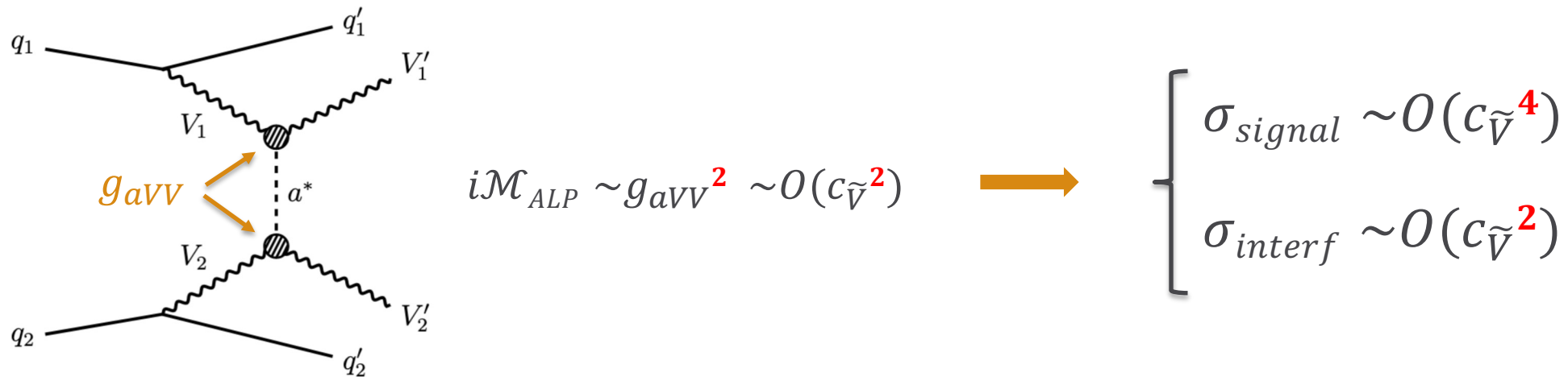


Results

- **Maximum likelihood fit** of signal and background to the **diboson invariant/transverse masses**
- No excess found with respect to SM expectations
- **Current limits** with CMS Run 2 data and **projected limits** at Run 3 and HL-LHC in the ALP ($c_{\tilde{W}}, c_{\tilde{B}}$) parameter space

Results

- **Diff. cross-sections are parameterized** in the $(c_{\tilde{W}}, c_{\tilde{B}})$ plane with quartic / quadratic **polynomials** for pure signal / interference ALP components.



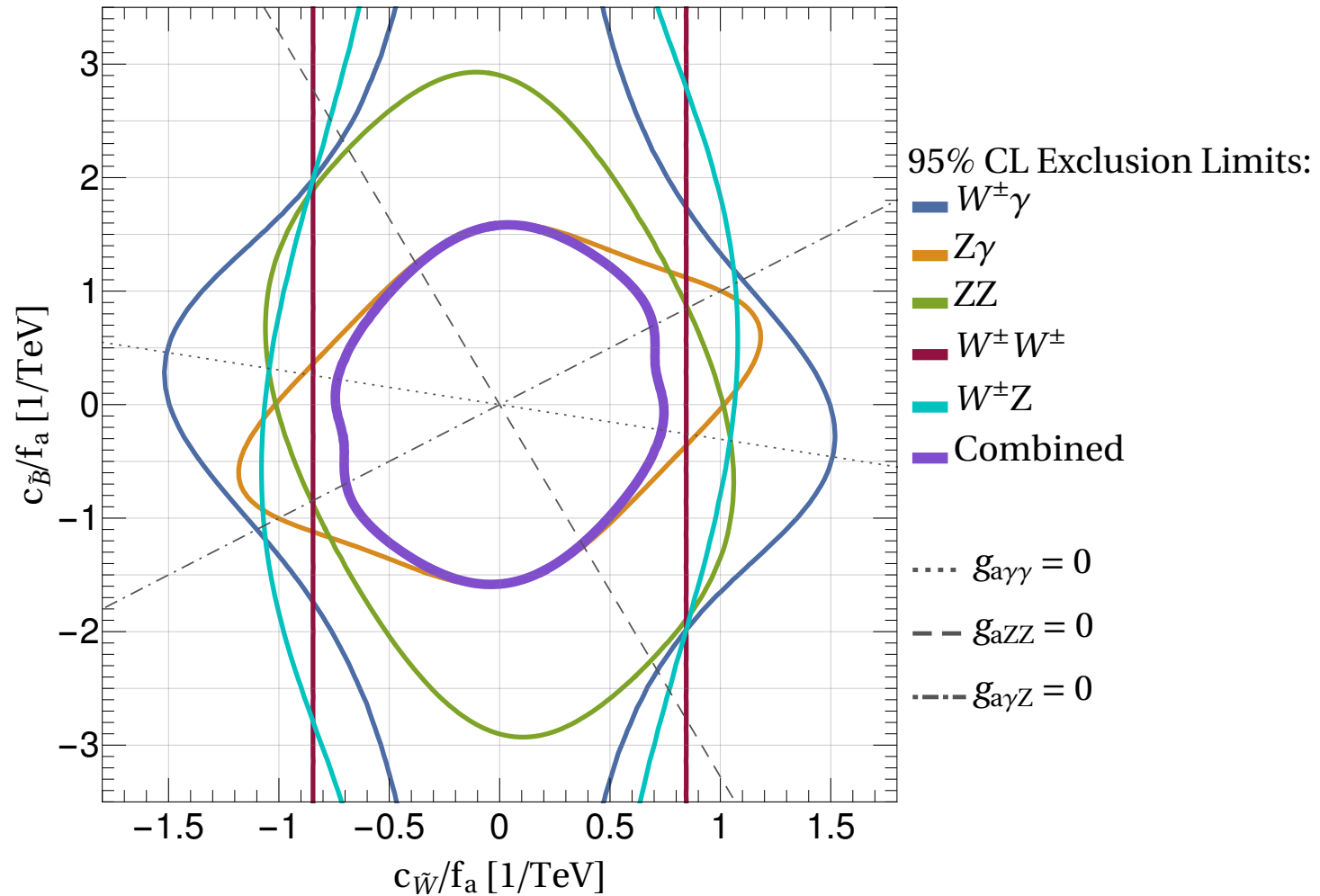
- MadGraph5_aMC@NLO reweighting tool for the generation at different points in the $(c_{\tilde{W}}, c_{\tilde{B}})$ plane:

$$\begin{array}{l}
 g_{a\gamma Z} = 0 \longrightarrow p_0 = (1, 1), \quad p_1 = (0, 2), \quad p_2 = (1, 0), \\
 p_3 = (1, -1), \quad p_4 = (1, -0.305), \quad p_5 = (1, -3.279) \longleftarrow g_{aZZ} = 0 \\
 \longleftarrow g_{a\gamma\gamma} = 0
 \end{array}$$

Results

CURRENT LIMITS

- Only ZZ and $Z\gamma$ can constrain $c_{\tilde{B}}/f_a$
- Most promising channels: $Z\gamma$ and $W^\pm W^\pm$

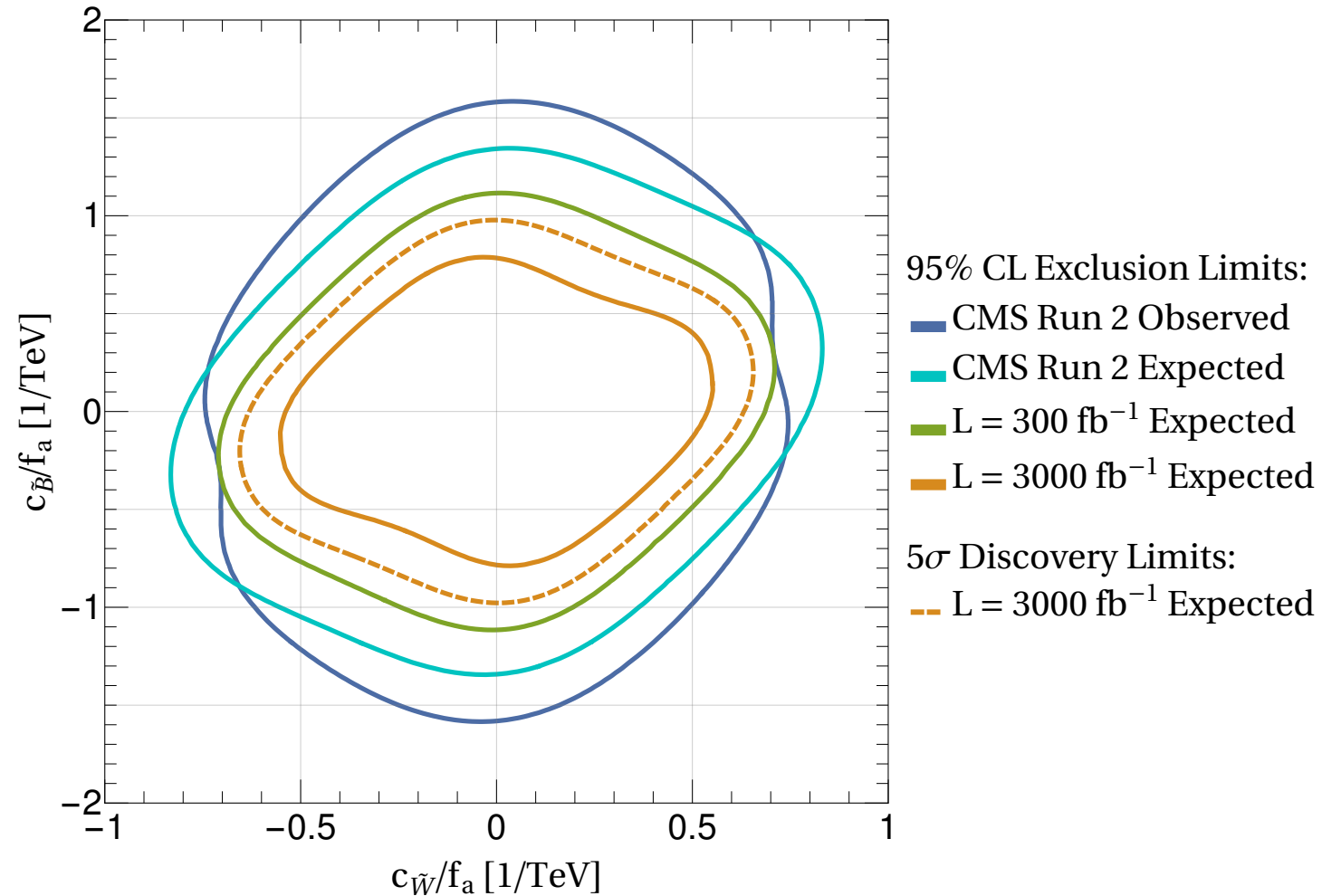


(for $f_a > 4$ TeV and $M_a < 100$ GeV)

Results

PROJECTED LIMITS

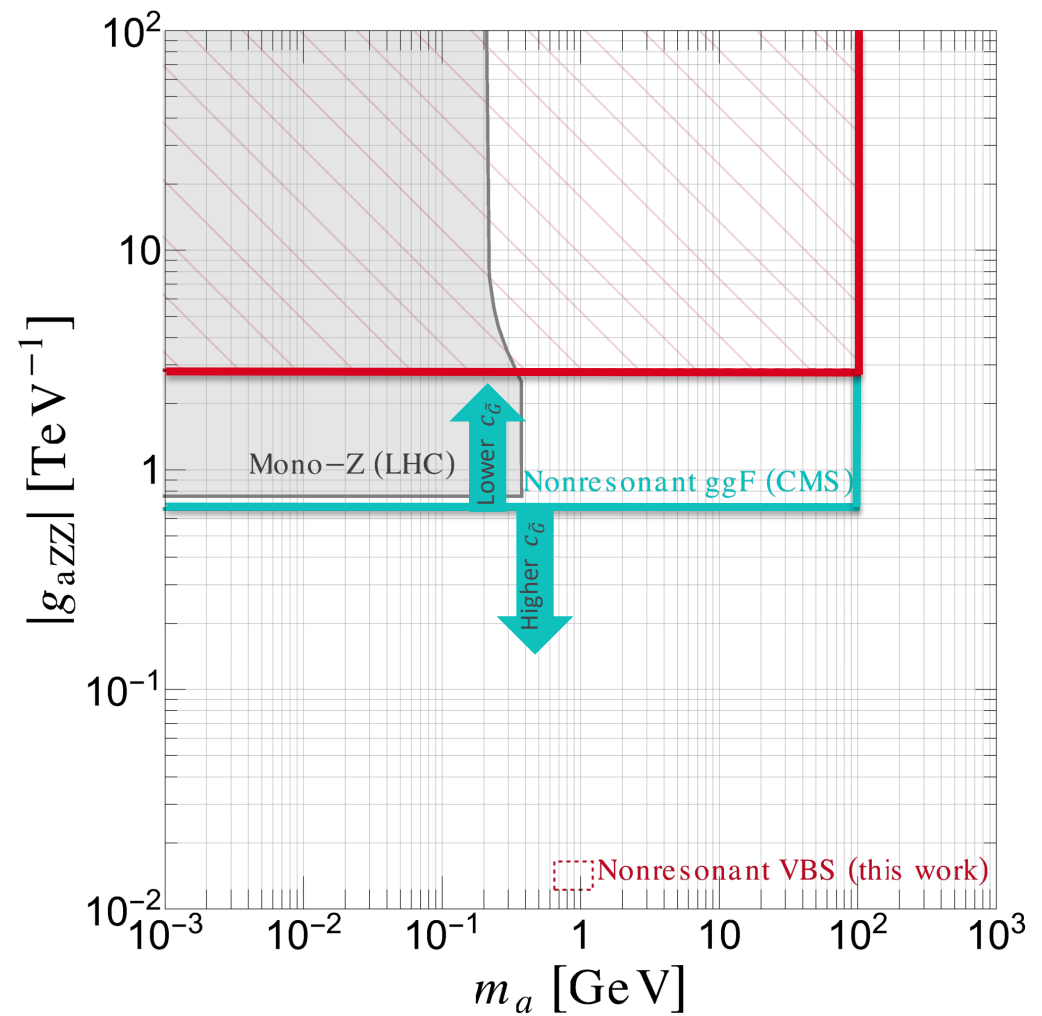
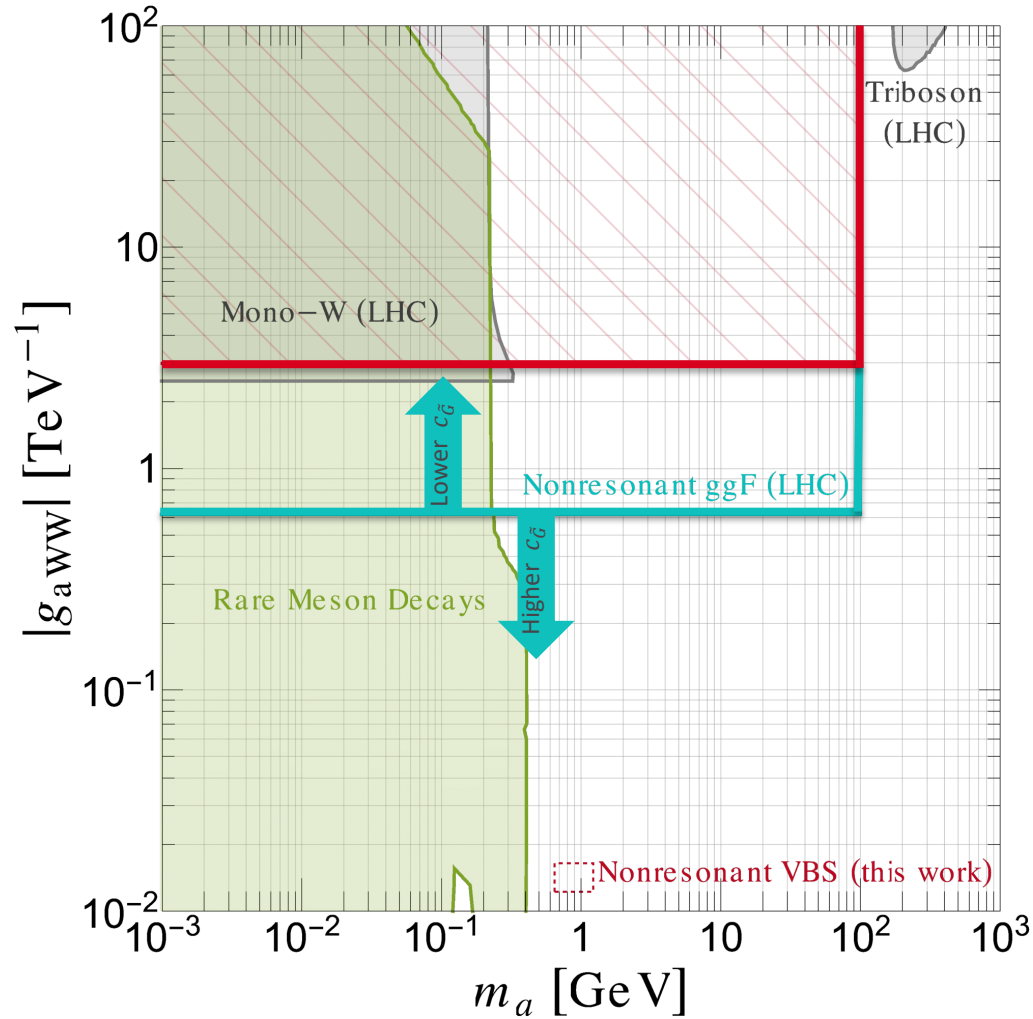
- HL-LHC limits on ALP couplings decrease by a factor $\sim 1.5 - 1.7$
- HL-LHC limits on ALP cross sections decrease by a factor $\sim 5 - 8$



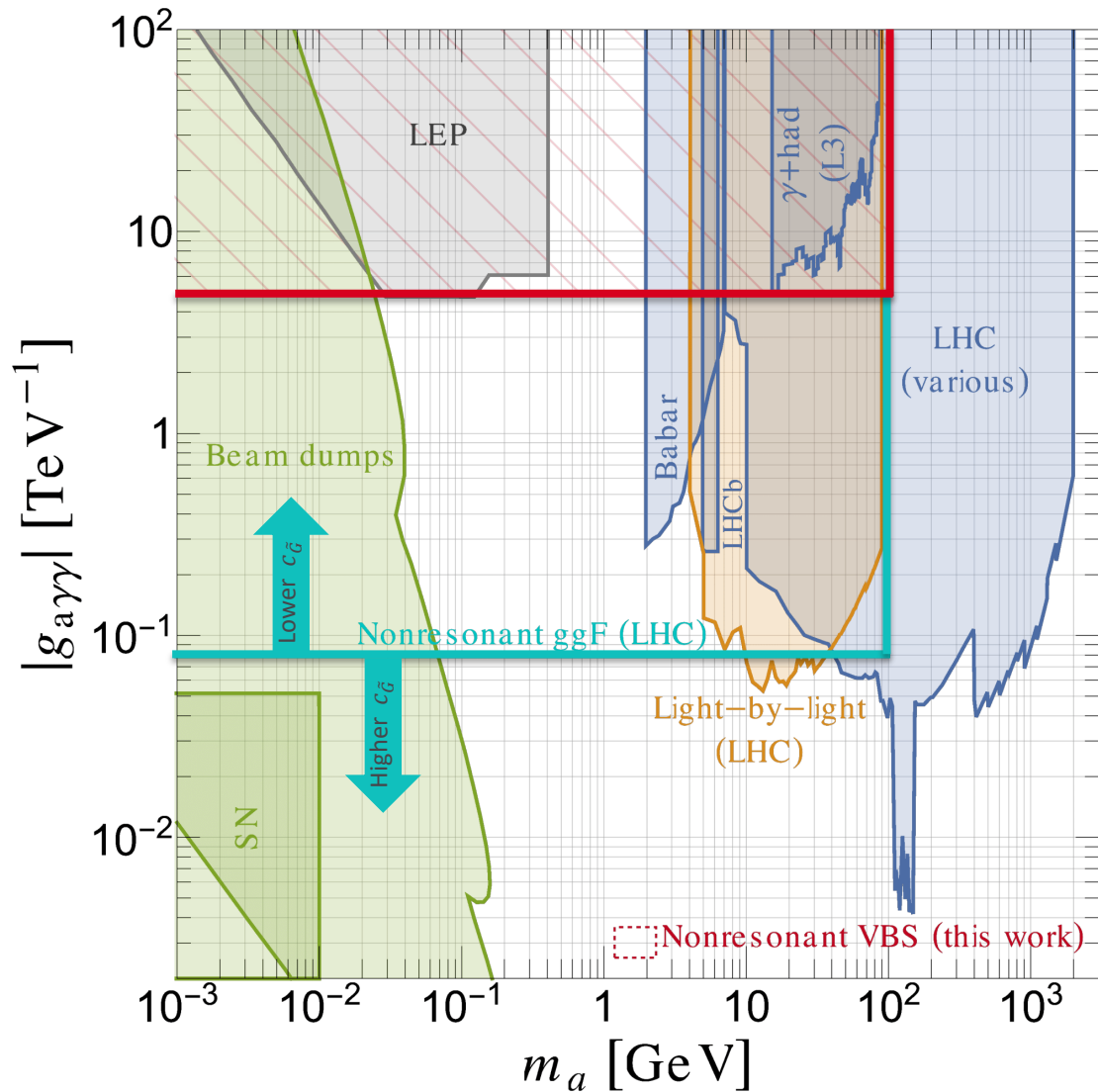
(for $f_a > 4 \text{ TeV}$ and $M_a < 100 \text{ GeV}$)

Results: comparison to existing bounds

- Limits are **very competitive** and probe **previously unexplored regions** of the param. space



Results: comparison to existing bounds



- **Red:** this work
- **Green:** no assumptions
- **Light blue:** nonresonant ggF. Depend on the coupling to gluons and assume $g_{agg} = 1 \text{ TeV}^{-1}$
- **Dark blue:** gluon dominance, i.e., $g_{agg} \gg g_{aV_1V_2}$
- **Orange:** $BR(a \rightarrow \gamma\gamma) = 1$
- **Grey:** more elaborate assumptions on the EWK sector

Conclusions

- Access to **EWK couplings independently of the gluons**
- Current limits (CMS Run 2 data) and projected limits (Run 3 and HL)
- Limits **independent** of the **ALP mass and decay width** ($m_a \lesssim 100$ GeV)
- **Limits** are **very competitive** and probe **previously unexplored regions** of the param. space
- Great opportunity for **dedicated ALP searches** at Run 3 and HL-LHC

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Thank you!

Lagrangian and physical couplings

$$\mathcal{L}_{ALP} = \frac{1}{4} \partial_\mu a \partial^\mu a - \frac{m_a^2}{2} a^2 - c_{\tilde{B}} \frac{a}{f_a} B_{\mu\nu} \tilde{B}^{\mu\nu} - c_{\tilde{W}} \frac{a}{f_a} W_{\mu\nu}^i \tilde{W}^{i\mu\nu} - c_{\tilde{G}} \frac{a}{f_a} G_{\mu\nu}^A \tilde{G}^{A\mu\nu}$$

EWSB



$$\mathcal{L}_{ALP,EW} = -\frac{g_{a\gamma\gamma}}{4} a F^{\mu\nu} \tilde{F}_{\mu\nu} - \frac{g_{a\gamma Z}}{4} a Z^{\mu\nu} \tilde{F}_{\mu\nu} - \frac{g_{aZZ}}{4} a Z^{\mu\nu} \tilde{Z}_{\mu\nu} - \frac{g_{aWW}}{2} a W^{+\mu\nu} \tilde{W}_{\mu\nu}^-$$

(Unitary gauge)

$$g_{a\gamma\gamma} = \frac{4}{f_a} (s_\theta^2 c_{\tilde{W}} + c_\theta^2 c_{\tilde{B}})$$

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$$g_{agg} = \frac{4}{f_a} c_{\tilde{G}}$$

$$g_{aZZ} = \frac{4}{f_a} (c_\theta^2 c_{\tilde{W}} + s_\theta^2 c_{\tilde{B}})$$

$$g_{aWW} = \frac{4}{f_a} c_{\tilde{W}}$$

Kinematic cuts

- Generation cuts: $p_T(q'_{1,2}) > 20 \text{ GeV}$, $\eta(q'_{1,2}) < 6$, $\Delta R(q'_1 q'_2) > 0.1$, $M_{q'_1 q'_2} > 120 \text{ GeV}$
 $p_T(\gamma) > 10 \text{ GeV}$, $\eta(\gamma) < 2.5$, $\Delta R(\gamma q'_{1,2}) > 0.4$,

- Selection cuts:

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Table 3. Summary of the CMS VBS analyses: the diboson mass observable, the integrated luminosity, the most important selection criteria and the normalization scale factor ρ .

Branching fractions and selection efficiencies

Analysis	ZZ	$Z\gamma$	$W^\pm\gamma$	$W^\pm Z$	$W^\pm W^\pm$
Branching fraction	0.45%	6.7%	22%	1.5%	4.8%
Efficiency	35.7%	14.0%	1.6%	11.3%	17.0%

Table 4. Summary of branching fractions and selection efficiencies for each VBS channel. The efficiencies are relative to the simulated events in which the W and Z bosons decay to electrons or muons.

Diboson mass upper cuts

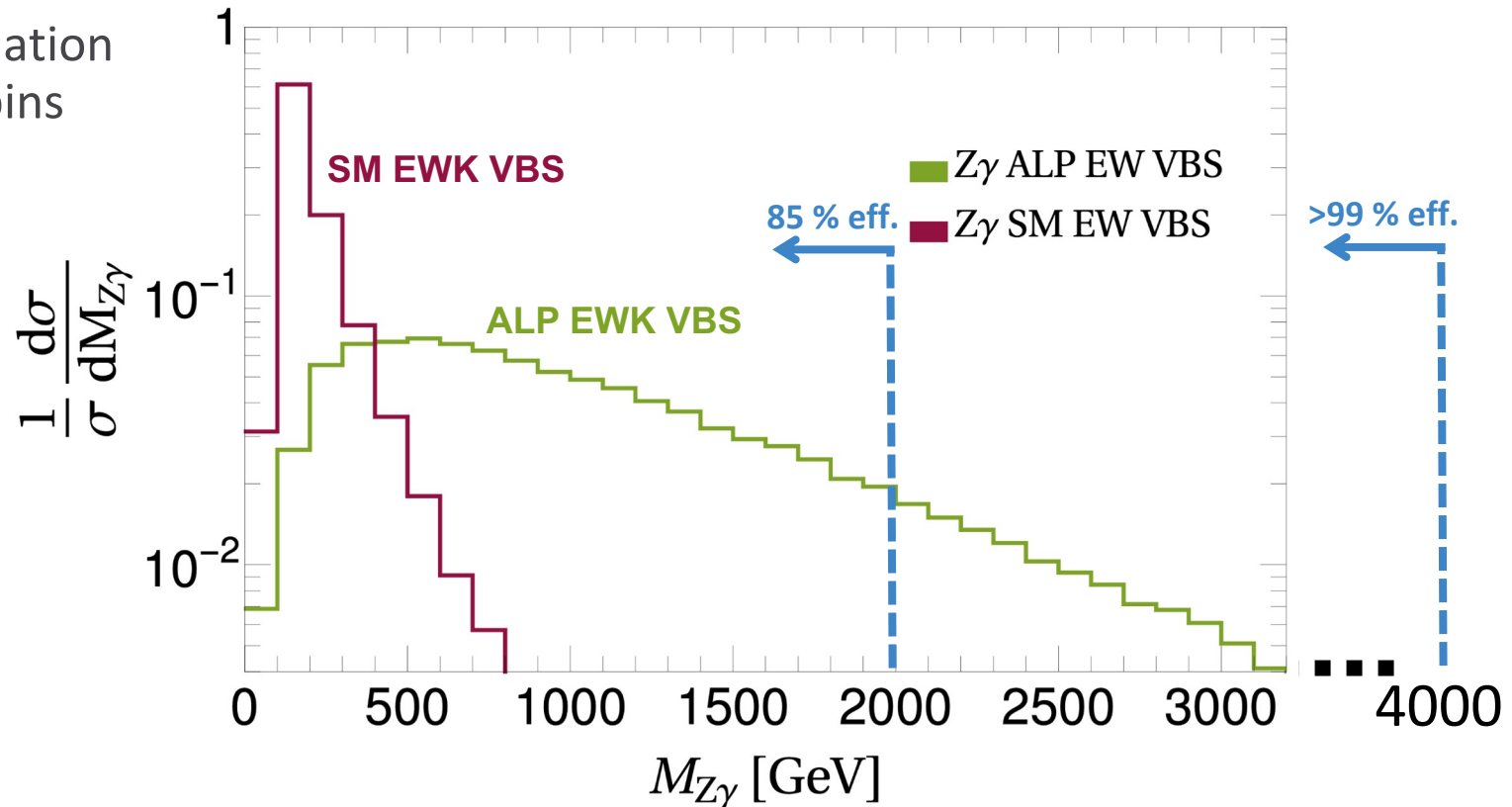
- Consistency of the ALP EFT and estimation of the impact of the highest-energy bins

→ upper cut on diboson mass

M_{VV}

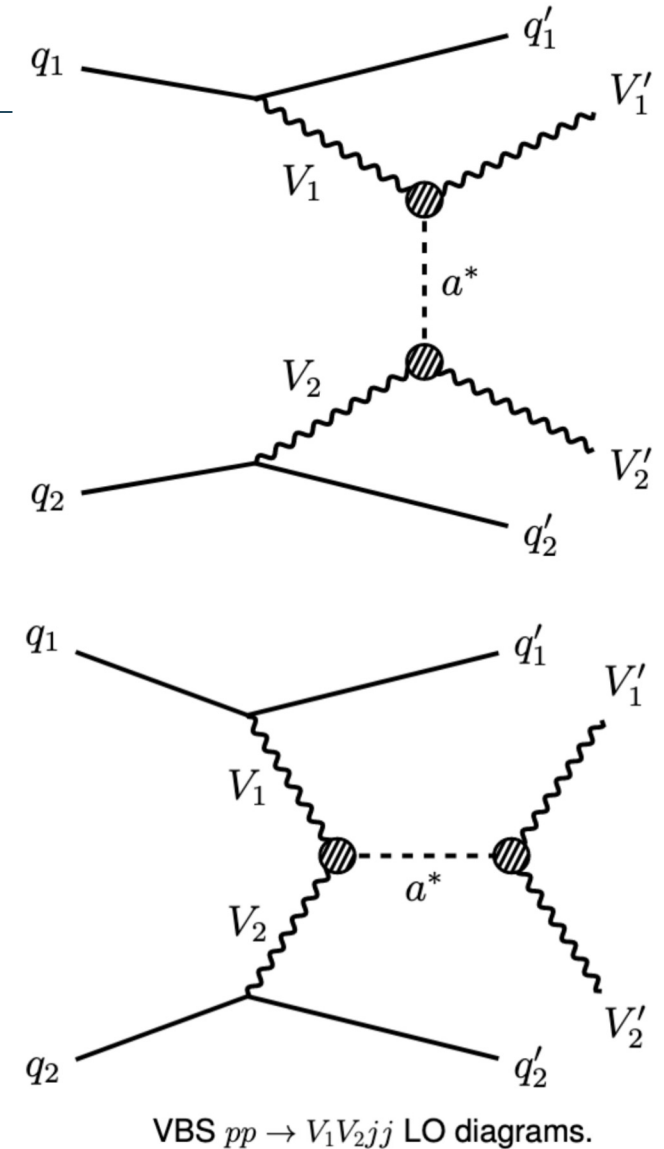
- Two benchmarks:

- $M_{VV} < 2 \text{ TeV}$: $\sim 85 \%$ efficiency
- $M_{VV} < 4 \text{ TeV}$: $> 99 \%$ efficiency

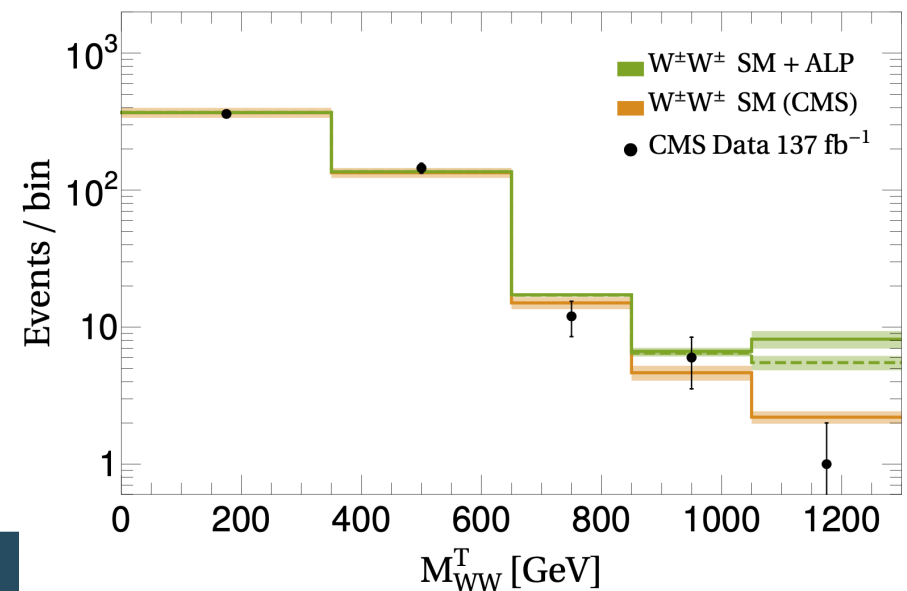
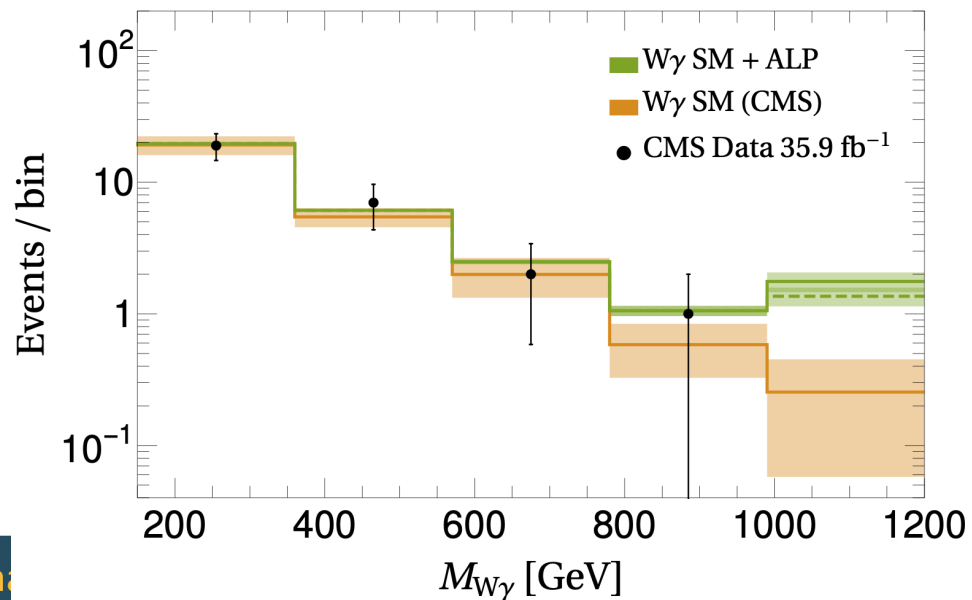
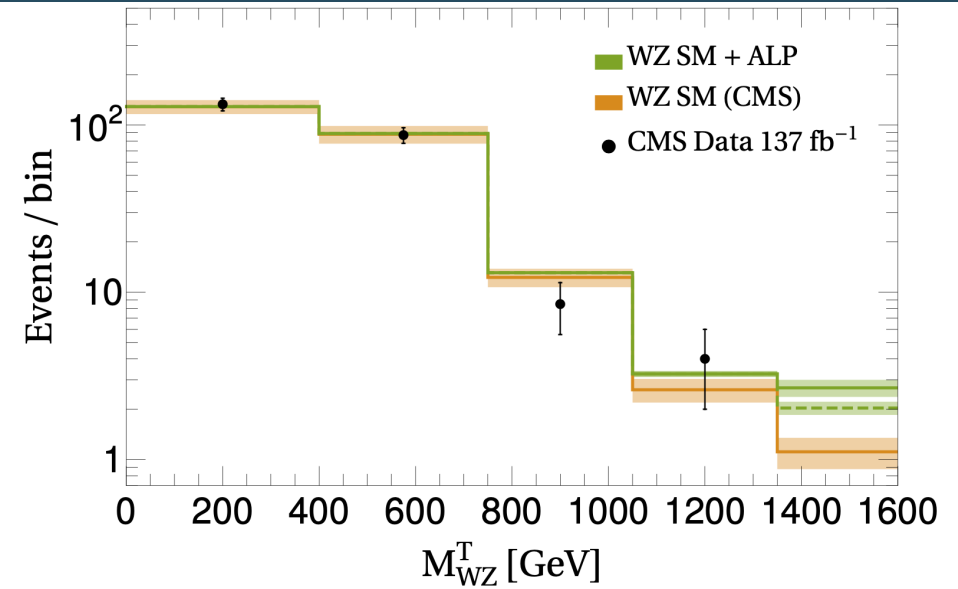
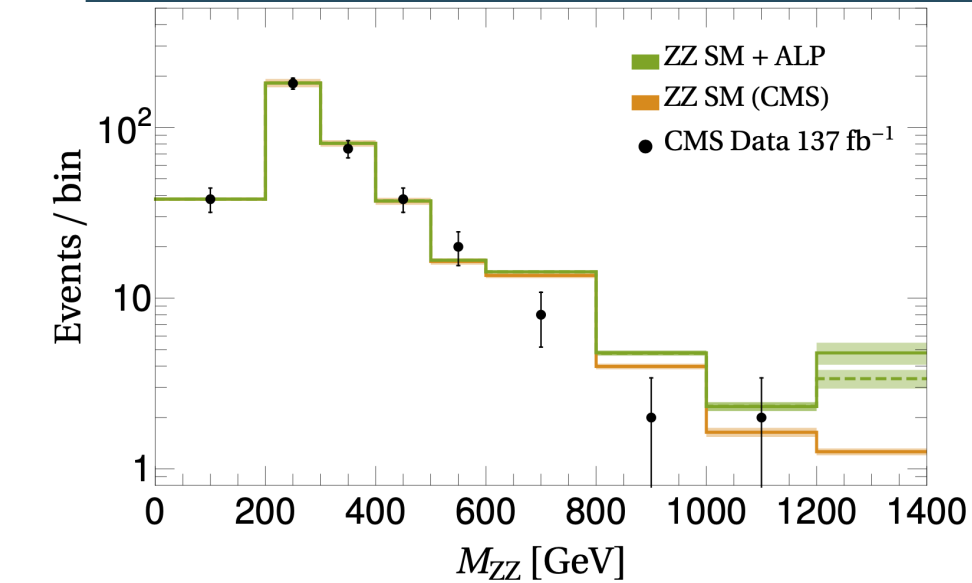


Contribution from gluons

- Same-sign WW, WZ, $W\gamma$: ALP QCD absent at tree level
- ZZ and $Z\gamma$: ALP QCD strongly reduced
 - Consistency with previous nonresonant limits [1905.12953], [2106.10085], [2111.13669]
 - VBS selection cuts
 - Large diboson masses
- For the tested region of the ALP parameter space, the theoretical prediction is dominated by ALP VBS
- Conservative: QCD ALP is positive with a subdominant contribution from its interference with EWK ALP



ALP Diboson Mass in CMS Leptonic Analyses



Results: comparison to existing bounds

