



Ne Ψ 23 | **NePSi 23**



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Università di Pisa – Dipartimento di Fisica
Europe/Rome timezone

QCD Axion vs Lattice QCD: **Past successes and future challenges**

Giovanni Villadoro



Definition of QCD axion:

pseudo-Goldstone boson whose shift symmetry is only/mostly broken by the QCD anomaly

Peccei Quinn '77
Weinberg, Wilczek '78
(KSVZ, DSFZ...)

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$$\mathcal{L}_{SM}^{\theta=0} + \frac{1}{2}(\partial_\mu a)^2 + \frac{a}{f_a} \frac{\alpha_s}{8\pi} G\tilde{G} + \dots$$

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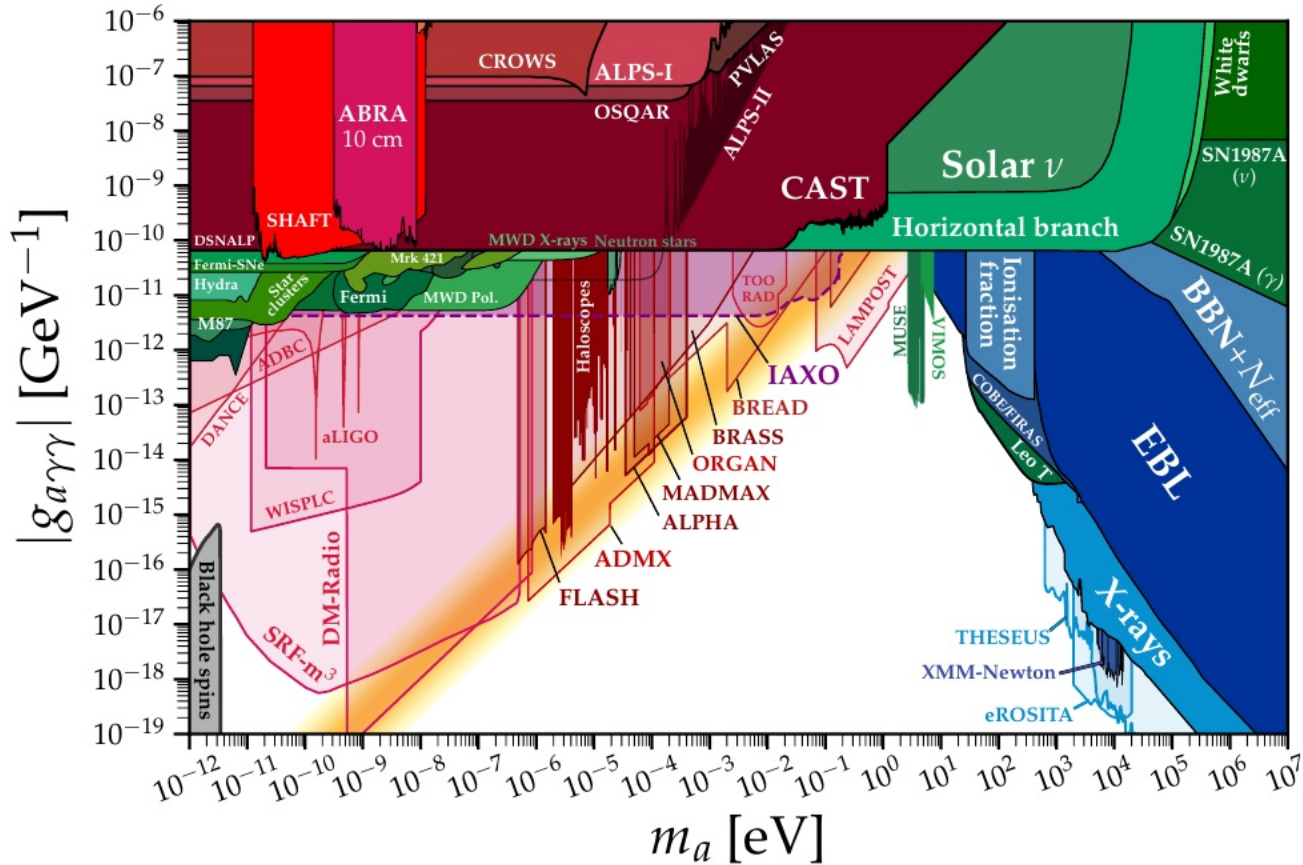


QCD dynamics relaxes θ_{QCD} to zero: no neutron EDM (up to CKM)

QCD axion can naturally explain observed DM abundance

Predictive/Testable Theory: $m_a \leftrightarrow g_{a\text{-SM}} \propto 1/f_a$

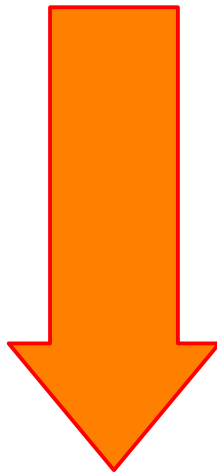
Axion Present and Future Searches



2203.14923

Axions Vs Lattice QCD:

- The Mass
- The Couplings to Nucleons
- The Couplings to Photons
- The Thermal Mass
- The Thermal Width



The Mass:



$$m_a^2 = \frac{\chi_{\text{top}}}{f_a^2}$$

$$\chi_{\text{top}} = \int d^4x \left\langle \frac{G\tilde{G}(x)}{32\pi^2} \frac{G\tilde{G}(0)}{32\pi^2} \right\rangle$$

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From Chiral Perturbation Theory

Weinberg '78

$$m_a^2 = \frac{m_u m_d}{(m_u + m_d)^2} \frac{m_\pi^2 f_\pi^2}{f_a^2}$$

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Grilli, Hardy, Pardo, GV - 1511.02867

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lattice average:

$$z \equiv \frac{m_u^{\overline{\text{MS}}}(2 \text{ GeV})}{m_d^{\overline{\text{MS}}}(2 \text{ GeV})} = 0.48(3)$$

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Grilli, Hardy, Pardo, GV - 1511.02867

$$7(4) \cdot 10^{-3} \Delta m_\pi^2|_{(\text{strong})}$$

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$$m_a = 5.70(6)(4) \mu\text{eV} \left(\frac{10^{12} \text{ GeV}}{f_a} \right)$$

The Mass:

Gorghetto, GV - 1812.01008

$$m_a = \left[\underbrace{5.815(22)_z(04)_{f_\pi}}_{\text{LO}} \underbrace{-0.121(38)_{\ell_i^r}}_{\text{NLO}} \underbrace{-0.022(07)_{\ell_i^r}(05)_{c_i^r}}_{\text{NNLO}} \underbrace{+0.019(06)_{k_i^r}}_{\text{EM}} \right] \mu\text{eV} \frac{10^{12} \text{ GeV}}{f_a}$$

The Mass:

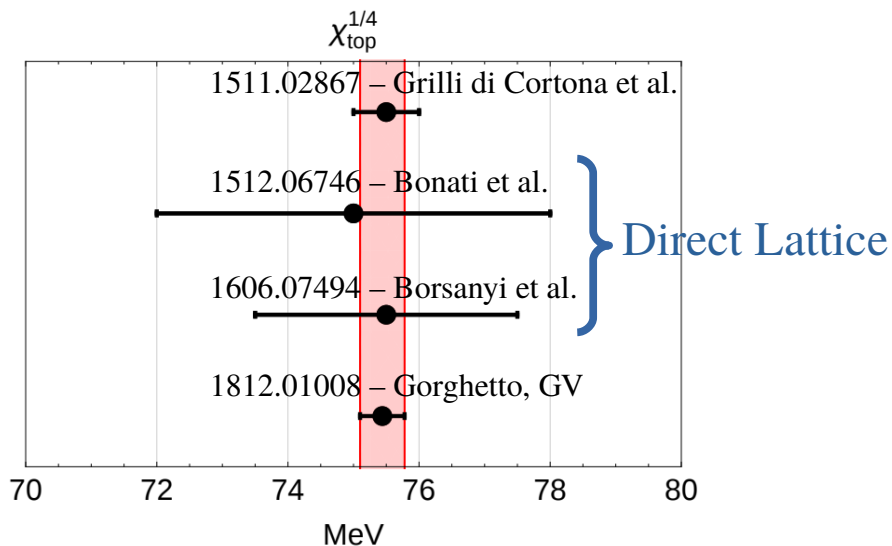
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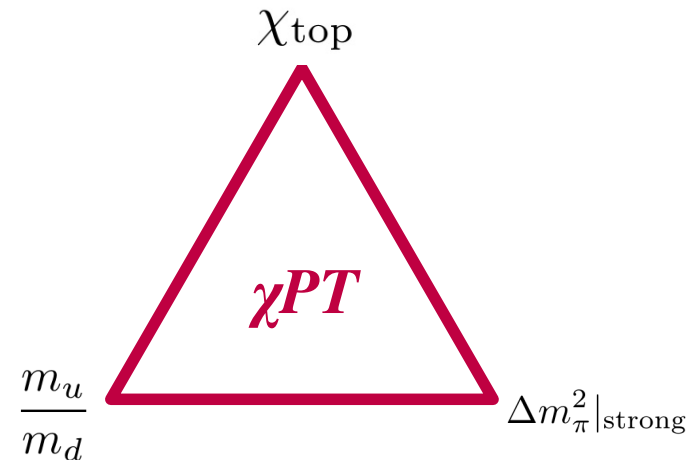
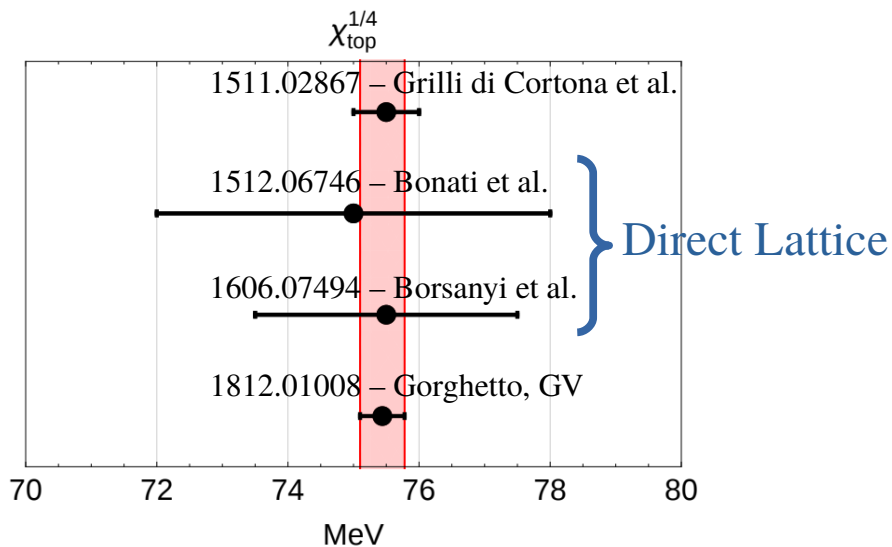
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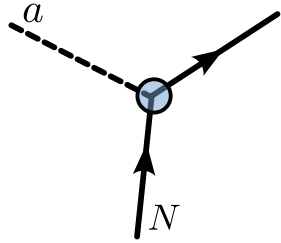
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The Coupling to Nucleons:

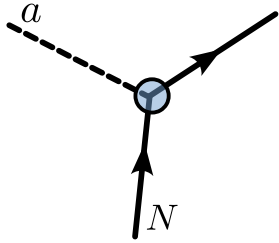
Grilli, Hardy, Pardo, GV - 1511.02867



$$\frac{\partial_\mu a}{2f_a} \sum_q c_q \bar{q} \gamma^\mu \gamma_5 q$$

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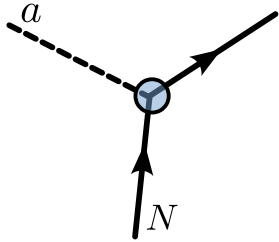
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$$\frac{\partial_\mu a}{2f_a} c_N \bar{N} \gamma^\mu \gamma_5 N$$

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Grilli, Hardy, Pardo, GV - 1511.02867



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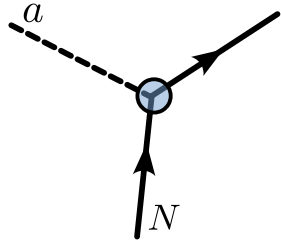


$$c_N = \sum_q c_q \Delta q$$

$$\langle N | \bar{q} \gamma^\mu \gamma_5 q | N \rangle = s^\mu \Delta q$$

The Coupling to Nucleons:

Grilli, Hardy, Pardo, GV - 1511.02867



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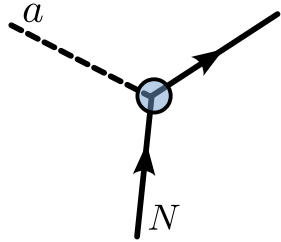
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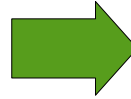
$$\Delta u - \Delta d = g_A = 1.2723(23)$$

The Coupling to Nucleons:

Grilli, Hardy, Pardo, GV - 1511.02867



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$$g_0^{ud} = \Delta u + \Delta d = 0.521(53)$$

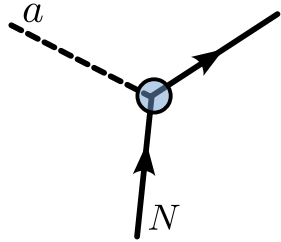
$$\Delta s = -0.026(4)$$

$$\Delta c = \pm 0.004$$

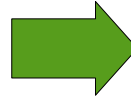
from Lattice @ 2015...

The Coupling to Nucleons:

Grilli, Hardy, Pardo, GV - 1511.02867



$$\frac{\partial_\mu a}{2f_a} \sum_q c_q \bar{q} \gamma^\mu \gamma_5 q$$



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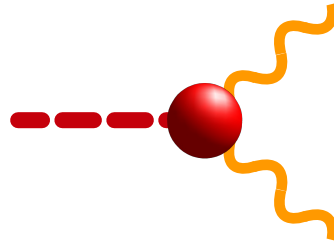
$$\Delta c = \pm 0.004$$

$$\left\{ \begin{array}{l} c_p = -0.47(3) + 0.88(3)c_u^0 - 0.39(2)c_d^0 - 0.038(5)c_s^0 \\ \quad \quad \quad - 0.012(5)c_c^0 - 0.009(2)c_b^0 - 0.0035(4)c_t^0, \\ c_n = -0.02(3) + 0.88(3)c_d^0 - 0.39(2)c_u^0 - 0.038(5)c_s^0 \\ \quad \quad \quad - 0.012(5)c_c^0 - 0.009(2)c_b^0 - 0.0035(4)c_t^0, \end{array} \right.$$

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The Coupling to Photons:

$$g_{a\gamma\gamma} = \frac{\alpha_{em}}{2\pi f_a} \left\{ \frac{E}{N} - \frac{2}{3} \frac{4m_d + m_u}{m_d + m_u} \right\}$$



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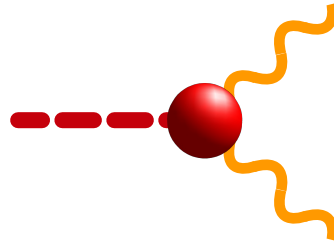
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$E/N =$

0 (KSVZ,...)

8/3 (DFSZ, GUT-KSVZ,...)

2 (Unification,...)



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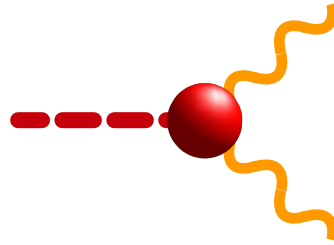
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tree ~ -2

$a \rightarrow \pi \rightarrow \gamma\gamma$



The Coupling to Photons:

Grilli, Hardy, Pardo, GV - 1511.02867

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NLO

$= 0.033(6)$

from $\pi \rightarrow \gamma\gamma$ $\eta \rightarrow \gamma\gamma$

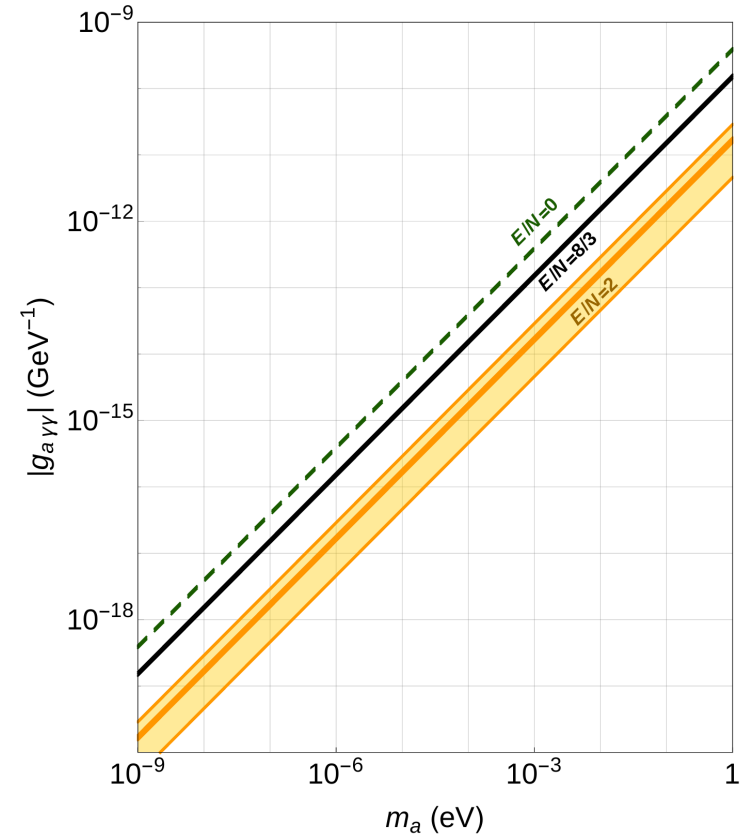
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$$g_{a\gamma\gamma} = \frac{\alpha_{em}}{2\pi f_a} \left[\frac{E}{N} - 1.92(4) \right]$$

$$g_{a\gamma\gamma} = \begin{cases} -2.227(44) \cdot 10^{-3} / f_a & E/N = 0 \\ 0.870(44) \cdot 10^{-3} / f_a & E/N = 8/3 \\ 0.095(44) \cdot 10^{-3} / f_a & E/N = 2 \end{cases}$$



The Coupling to Photons:

Grilli, Hardy, Pardo, GV - 1511.02867

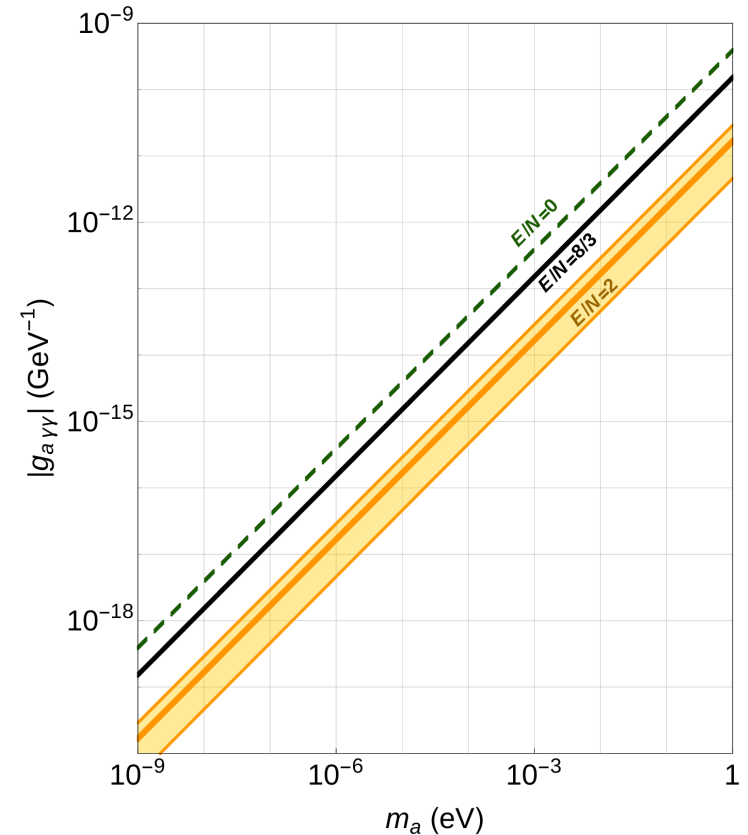
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$a\gamma\gamma$ from Lattice?

Any hope for $\pi\gamma\gamma$ from Lattice?

$$\langle j_A^\mu j_{em}^\nu j_{em}^\rho \rangle$$



The Thermal Mass:

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$T < T_c \sim m_\pi$	$\chi_{\text{top}}(T) \simeq \chi_{\text{top}}(0) \left[1 - \frac{3}{2(2\pi)^{3/2}} \frac{m_\pi^2}{f_\pi^2} \left(\frac{m_\pi}{T} \right)^{\frac{3}{2}} e^{-m_\pi/T} \right]$	Grilli, Hardy, Pardo, GV - 1511.02867
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$T \gg \Lambda_{QCD}$	$\chi_{\text{top}}(T) \sim \kappa \left(\frac{T_c}{T} \right)^\alpha \quad \alpha(N_f = 3) \simeq 8$	Gross, Pisarski, Yaffe - '81
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dilute instanton gas approx. (NLO corrections?)

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$T \gtrsim T_c$?

see M.P. Lombardo's talk

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dilute instanton gas approx. (NLO corrections?)

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$$m_a^2(T) = \frac{\chi_{\text{top}}(T)}{f_a^2} \quad \chi_{\text{top}}(T) = \int d^4x \left\langle \frac{G\tilde{G}(x)}{32\pi^2} \frac{G\tilde{G}(0)}{32\pi^2} \right\rangle_T$$

$T < T_c \sim m_\pi$	$\chi_{\text{top}}(T) \simeq \chi_{\text{top}}(0) \left[1 - \frac{3}{2(2\pi)^{3/2}} \frac{m_\pi^2}{f_\pi^2} \left(\frac{m_\pi}{T} \right)^{\frac{3}{2}} e^{-m_\pi/T} \right]$	Grilli, Hardy, Pardo, GV - 1511.02867
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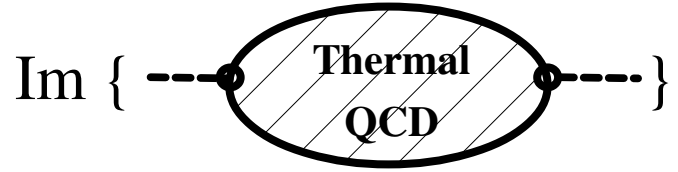
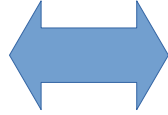
$T \gtrsim T_c$?  Controls Axion Abundance $\Omega_a = 0.1 k_\alpha \left[\frac{\theta_0}{2.15} \right]^2 \left[\frac{f_a}{10^{12} \text{ GeV}} \right]^{1+\frac{1}{2+\alpha/2}} \left[\frac{60}{g_\star} \right]^{\frac{1}{2}-\frac{1}{4+\alpha}}$

see M.P. Lombardo's talk

$T \gg \Lambda_{QCD}$	$\chi_{\text{top}}(T) \sim \kappa \left(\frac{T_c}{T} \right)^\alpha \quad \alpha(N_f = 3) \simeq 8$	Gross, Pisarski, Yaffe - '81
	dilute instanton gas approx. (NLO corrections?)	

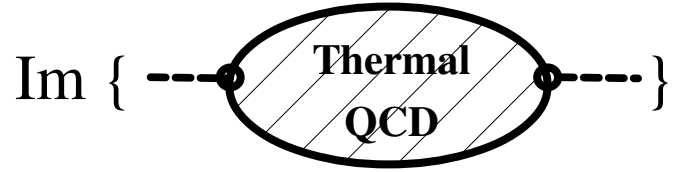
The Thermal Width:

$$\Gamma_{\text{top}}^{\gt} \equiv \int d^4x e^{ik^\mu x_\mu} \left\langle \frac{\alpha_s}{8\pi} G\tilde{G}(x^\mu) \frac{\alpha_s}{8\pi} G\tilde{G}(0) \right\rangle$$



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Thermal Width \rightarrow Relativistic Axions from SM bath:

Abundance depends on decoupling temperature
(which depends on coupling f_a)

Boltzmann Eq.

$$\frac{df_{\mathbf{p}}}{dt} = (1 + f_{\mathbf{p}}) \Gamma^{\lt} - f_{\mathbf{p}} \Gamma^{\gt}$$

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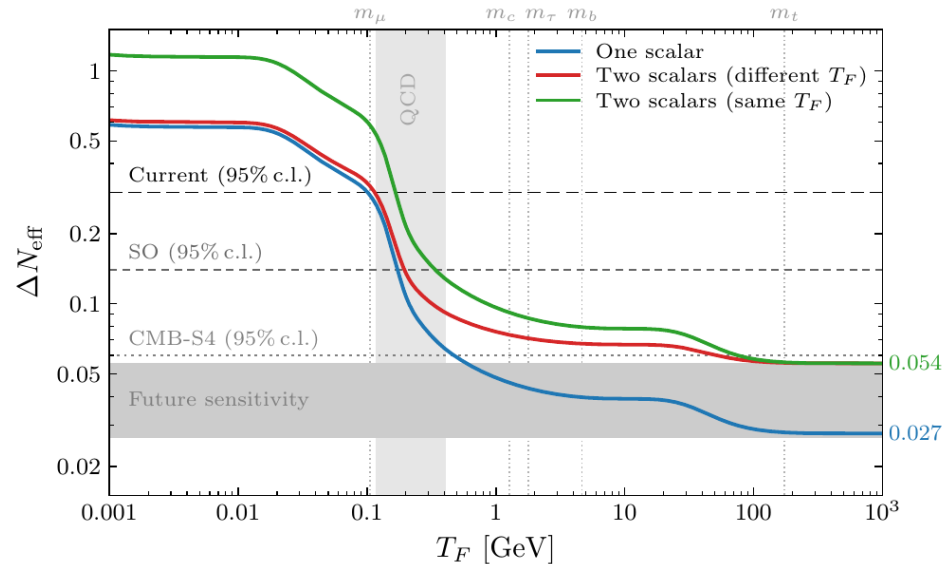
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from 2109.12088

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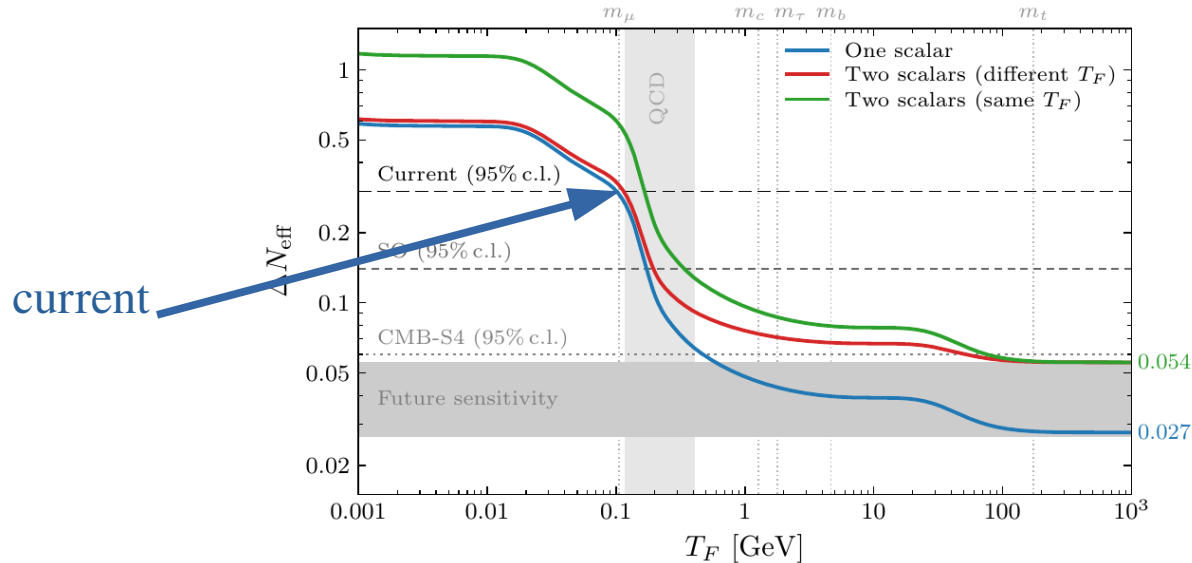
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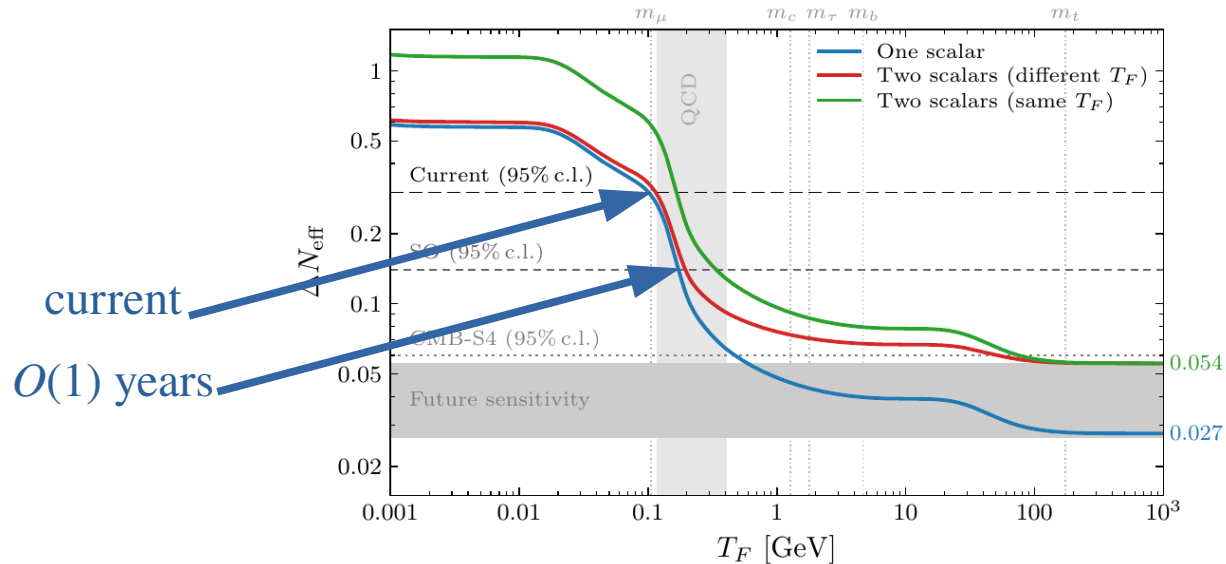
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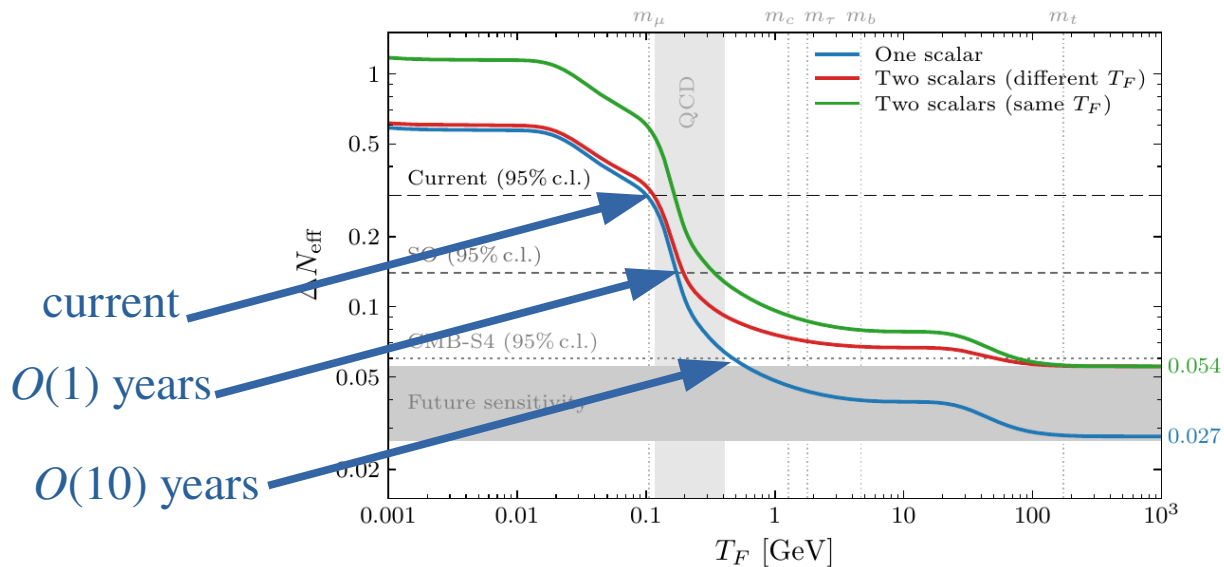
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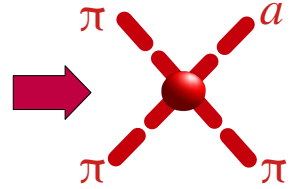


from 2109.12088

The Thermal Width:

$$T < T_c$$

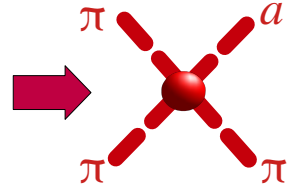
$$\Gamma^< = \frac{1}{2E} \int \left(\prod_{i=1}^3 \frac{d^3 \mathbf{k}_i}{(2\pi)^3 2E_i} \right) f_1^{\text{eq}} f_2^{\text{eq}} (1 + f_3^{\text{eq}}) (2\pi)^4 \delta^{(4)}(k_1^\mu + k_2^\mu - k_3^\mu - k^\mu) |\mathcal{M}|^2$$



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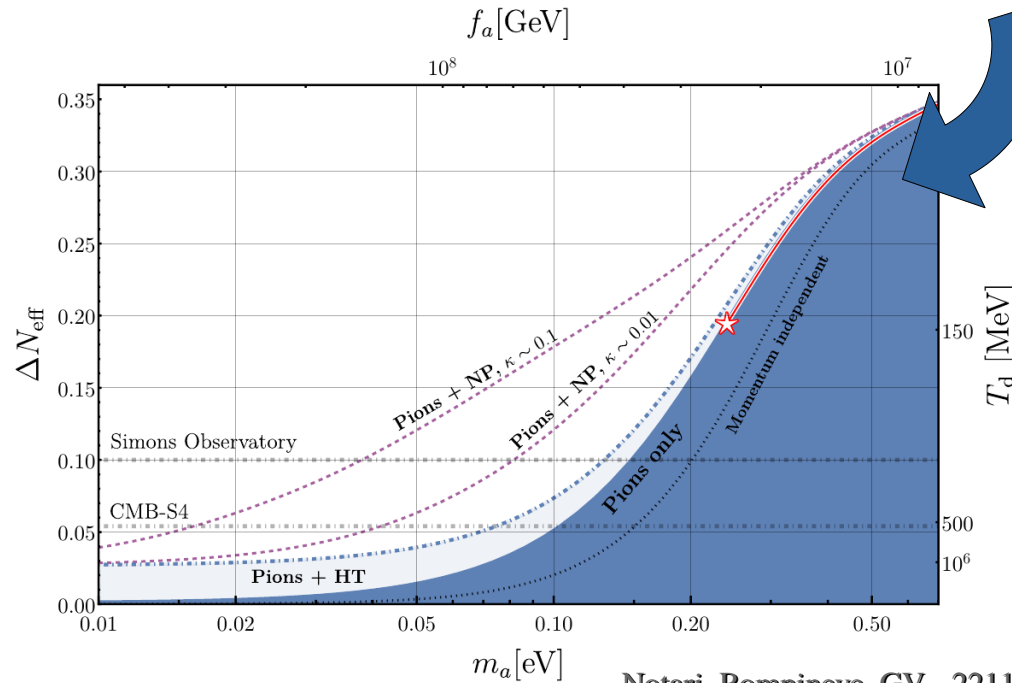
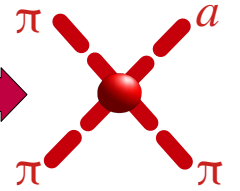
χ PT +
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see G.Piazza's talk

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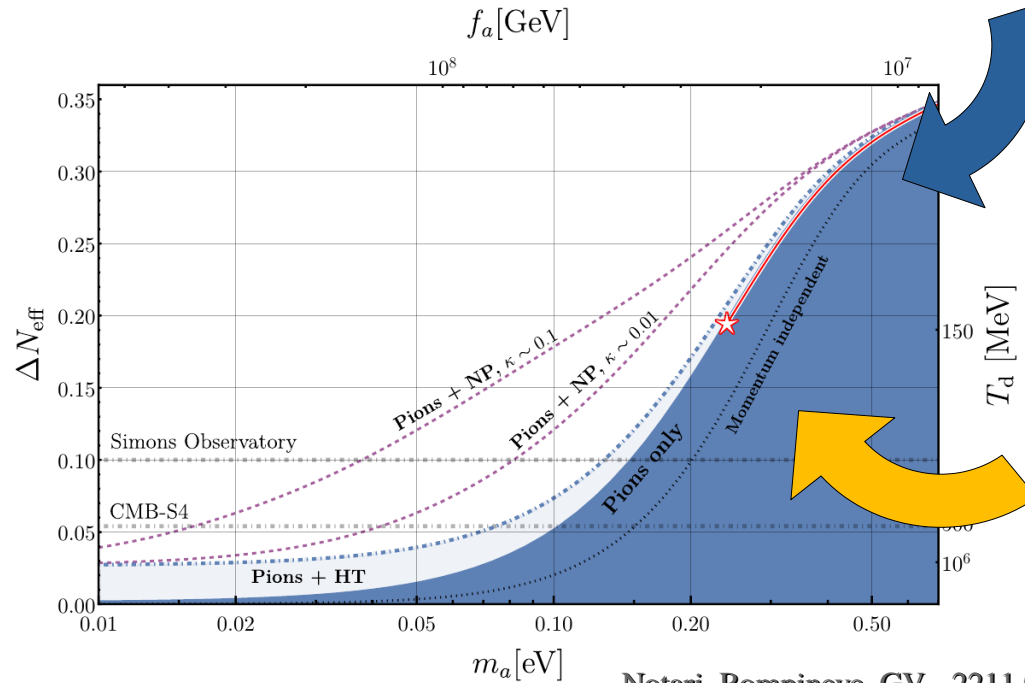
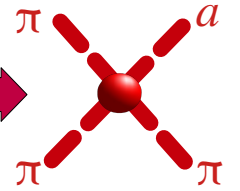
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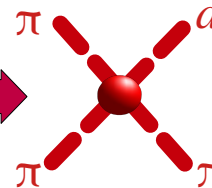
see G.Piazza's talk

should be considered
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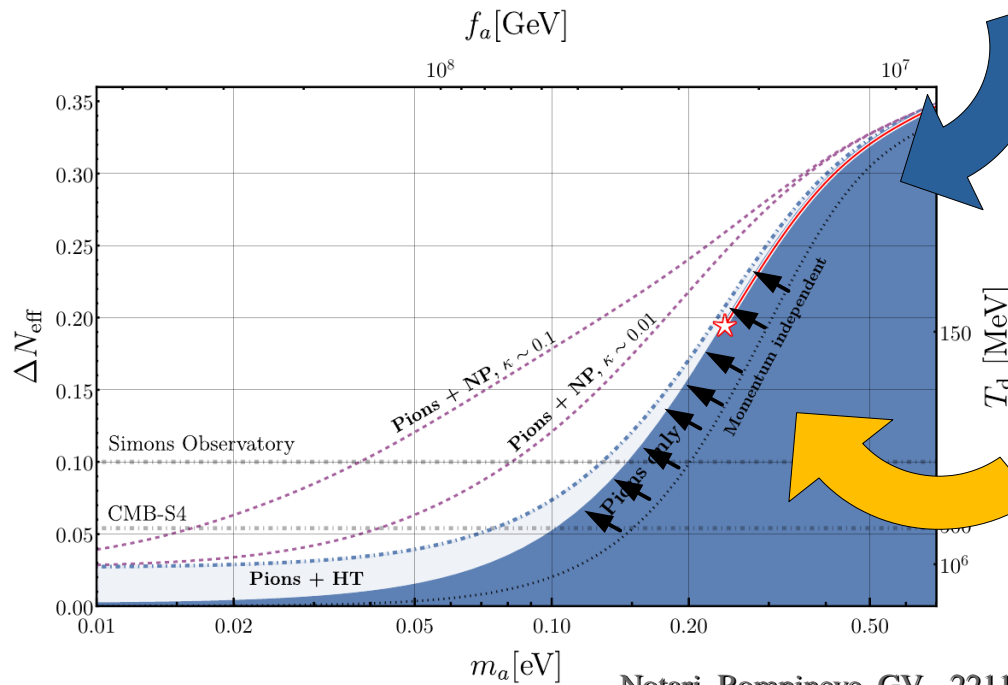
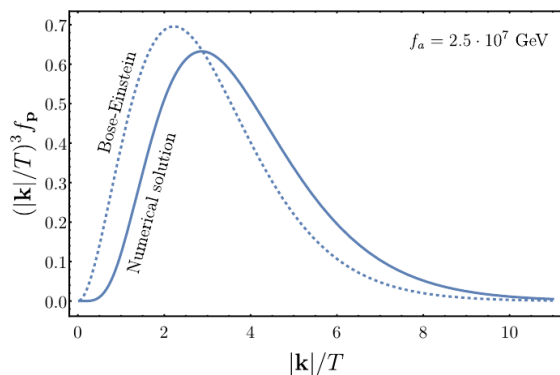
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Important Effects:

1. momentum dependence
~ 40 %



χ^2 PT +
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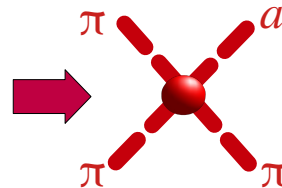
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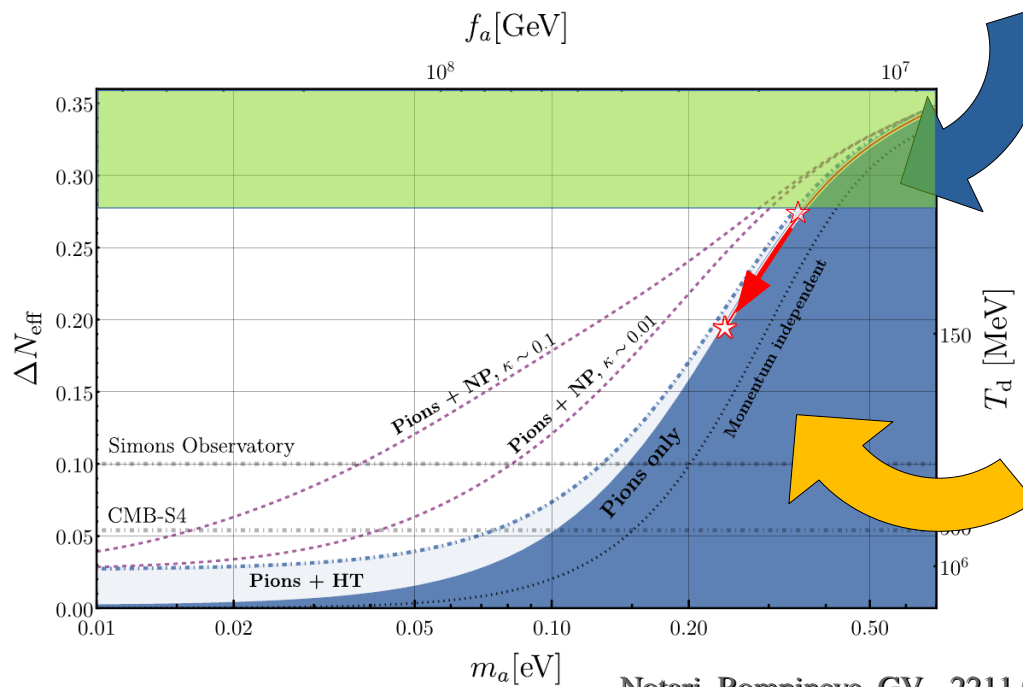
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Important Effects:

1. momentum dependence
~ 40 %
2. finite axion mass effects
~ 30 %



χ PT +
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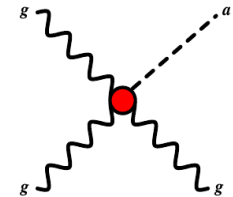
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Upcoming data requires understanding Thermal Width @ $T > T_c$

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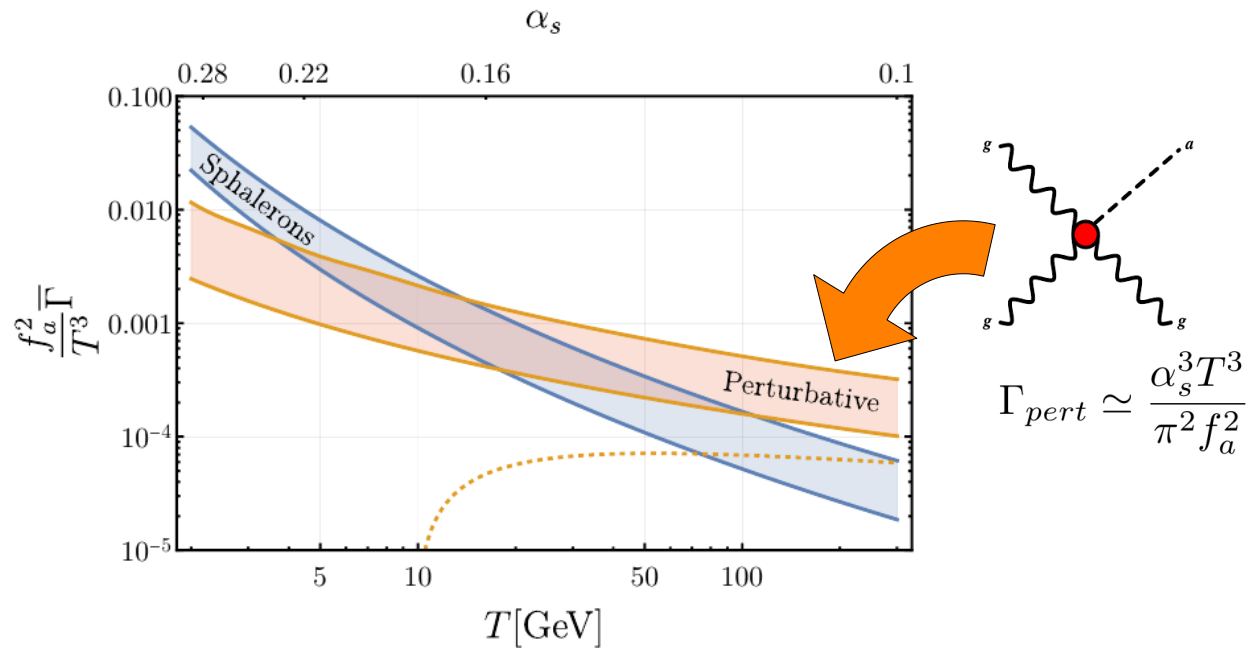
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$$\Gamma_{pert} \simeq \frac{\alpha_s^3 T^3}{\pi^2 f_a^2}$$

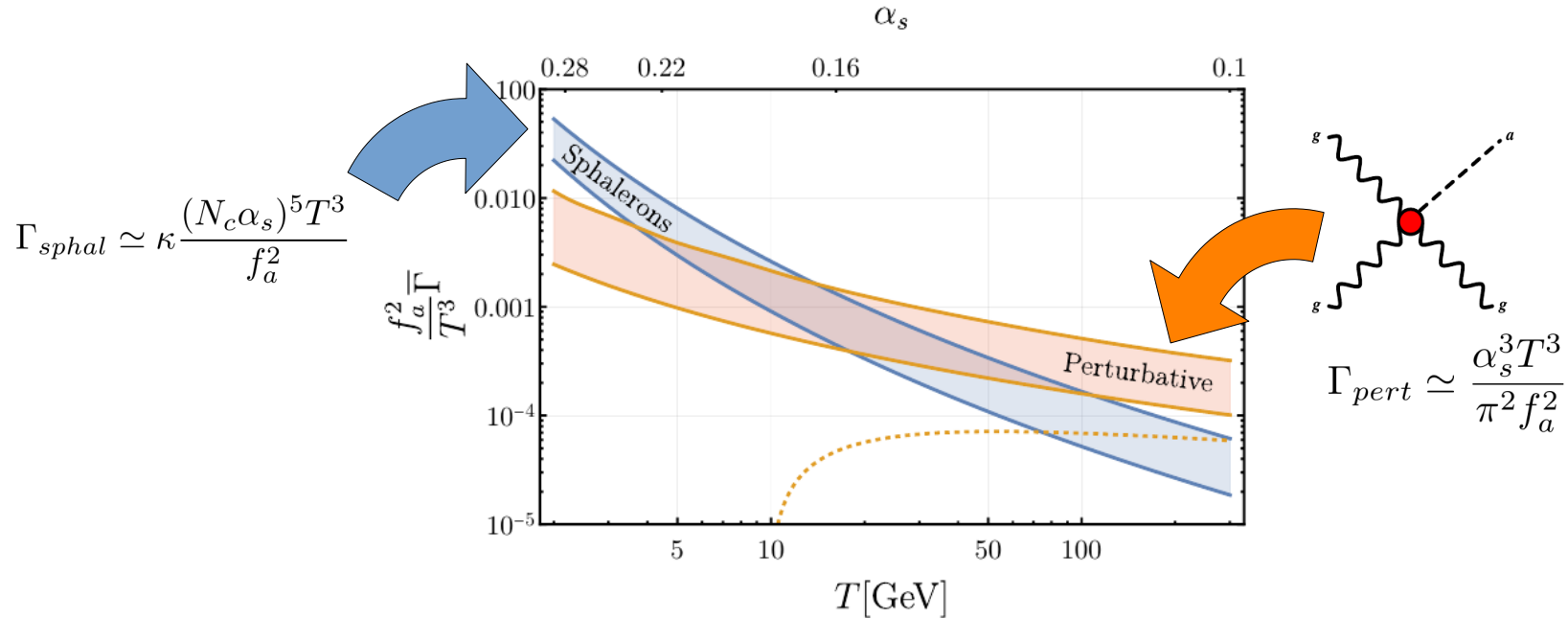
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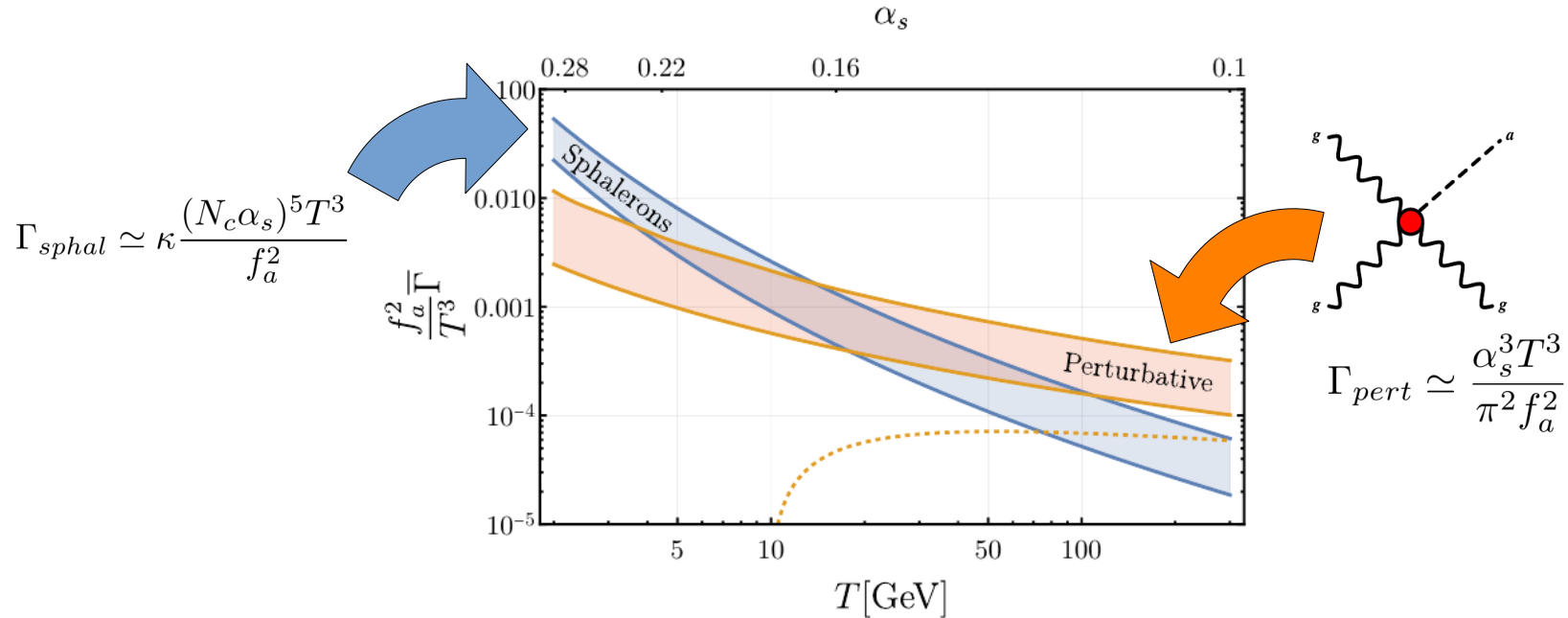
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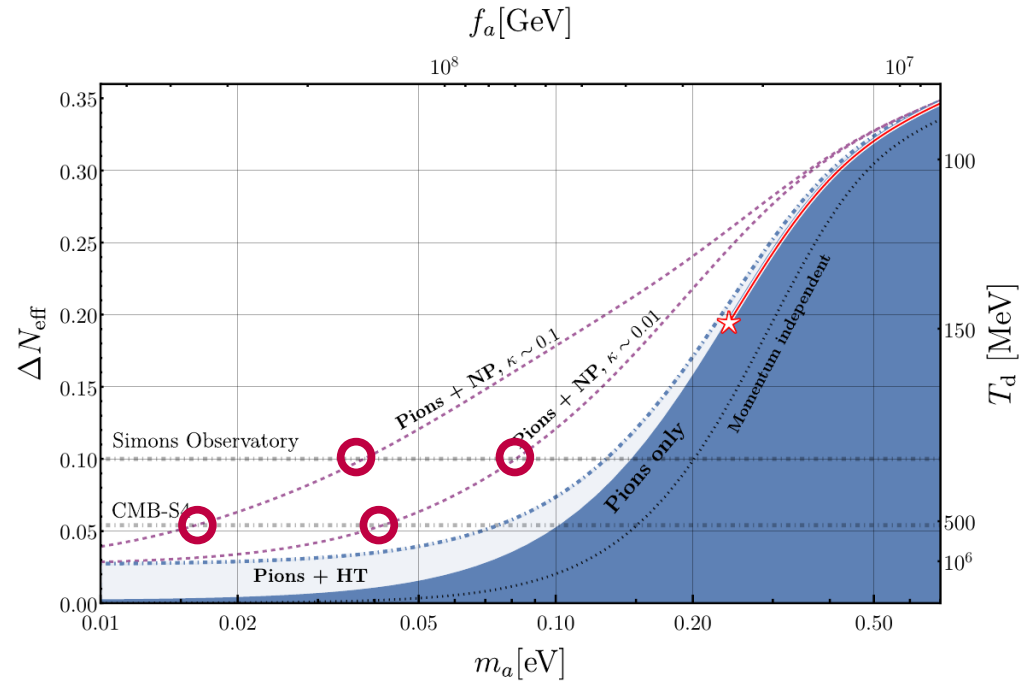


Plausible that both estimates are unreliable (both rely on perturbativity/semiclassics)
definitely below 1 GeV

The Thermal Width:

Upcoming data requires understanding Thermal Width @ $T > T_c$

$$\Gamma_{sphal} \simeq \kappa \frac{(N_c \alpha_s)^5 T^3}{f_a^2}$$



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Important to exploit upcoming experiments!

Thank you!