

The landscape of QCD axion models

Ne Ψ 2023 - Pisa - 17.02.2023

Luca Di Luzio



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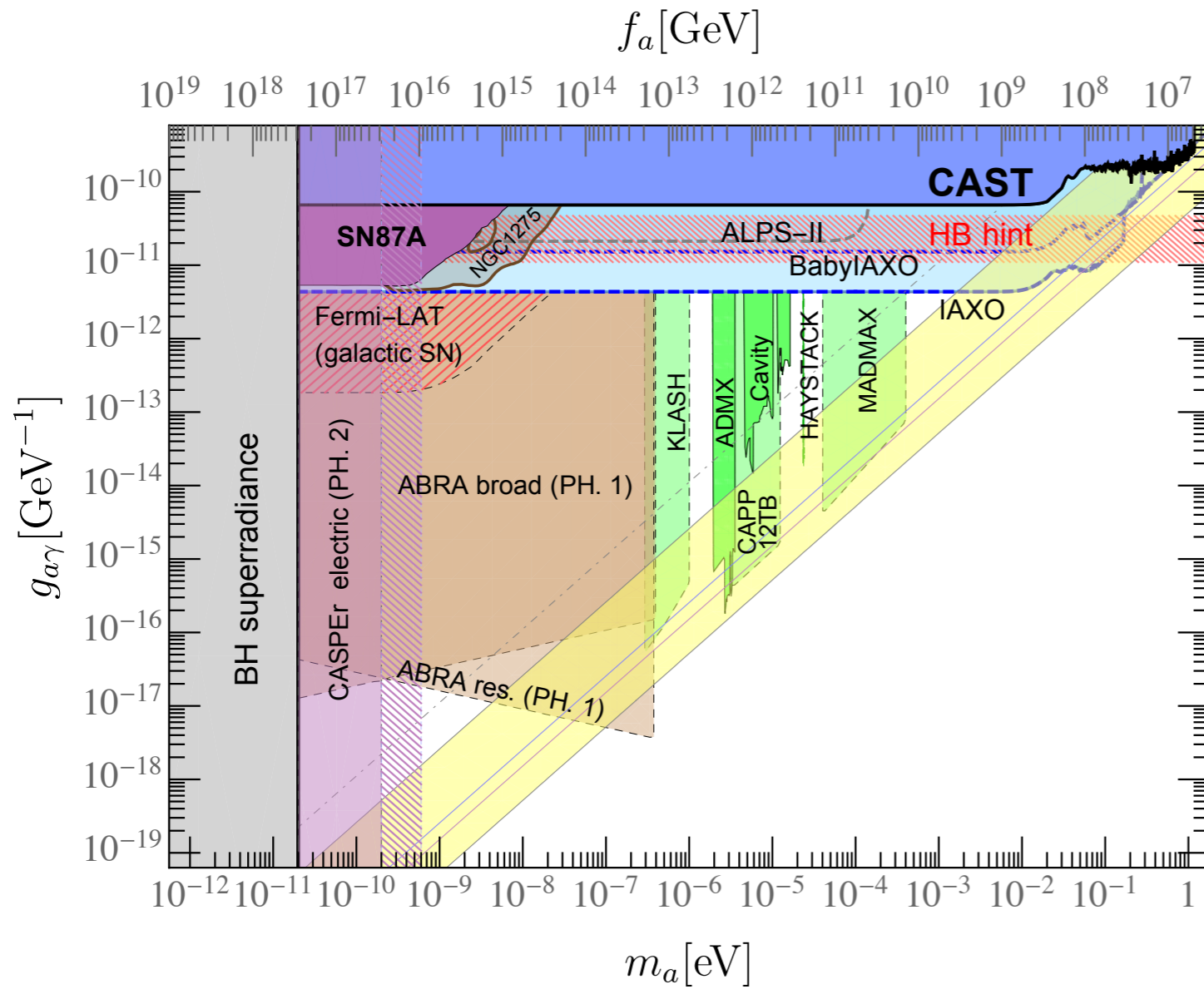


UNIVERSITÀ
DEGLI STUDI
DI PADOVA



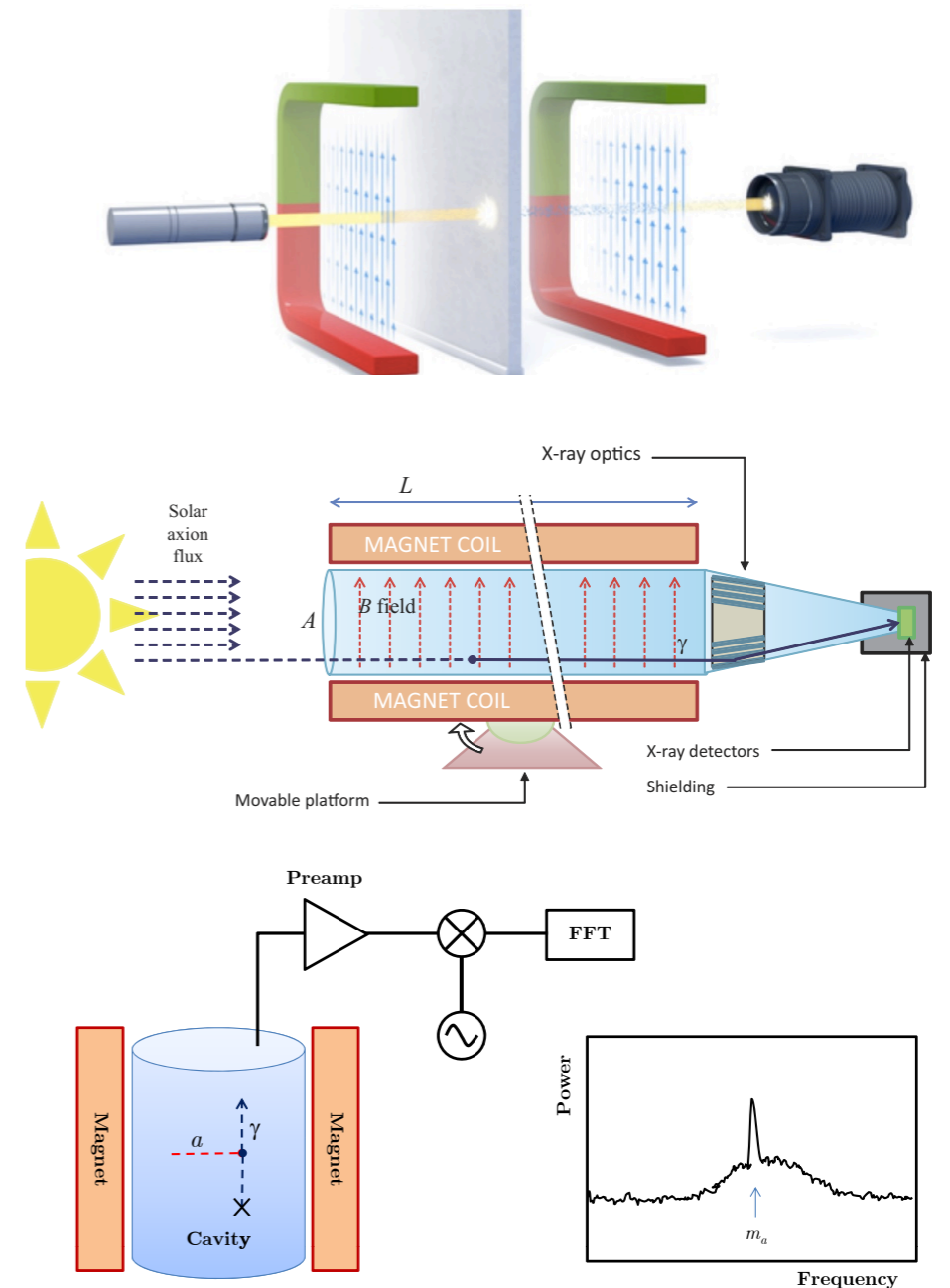
Dipartimento di Fisica e
Astronomia
"Galileo Galilei"

In 10 years from now ?

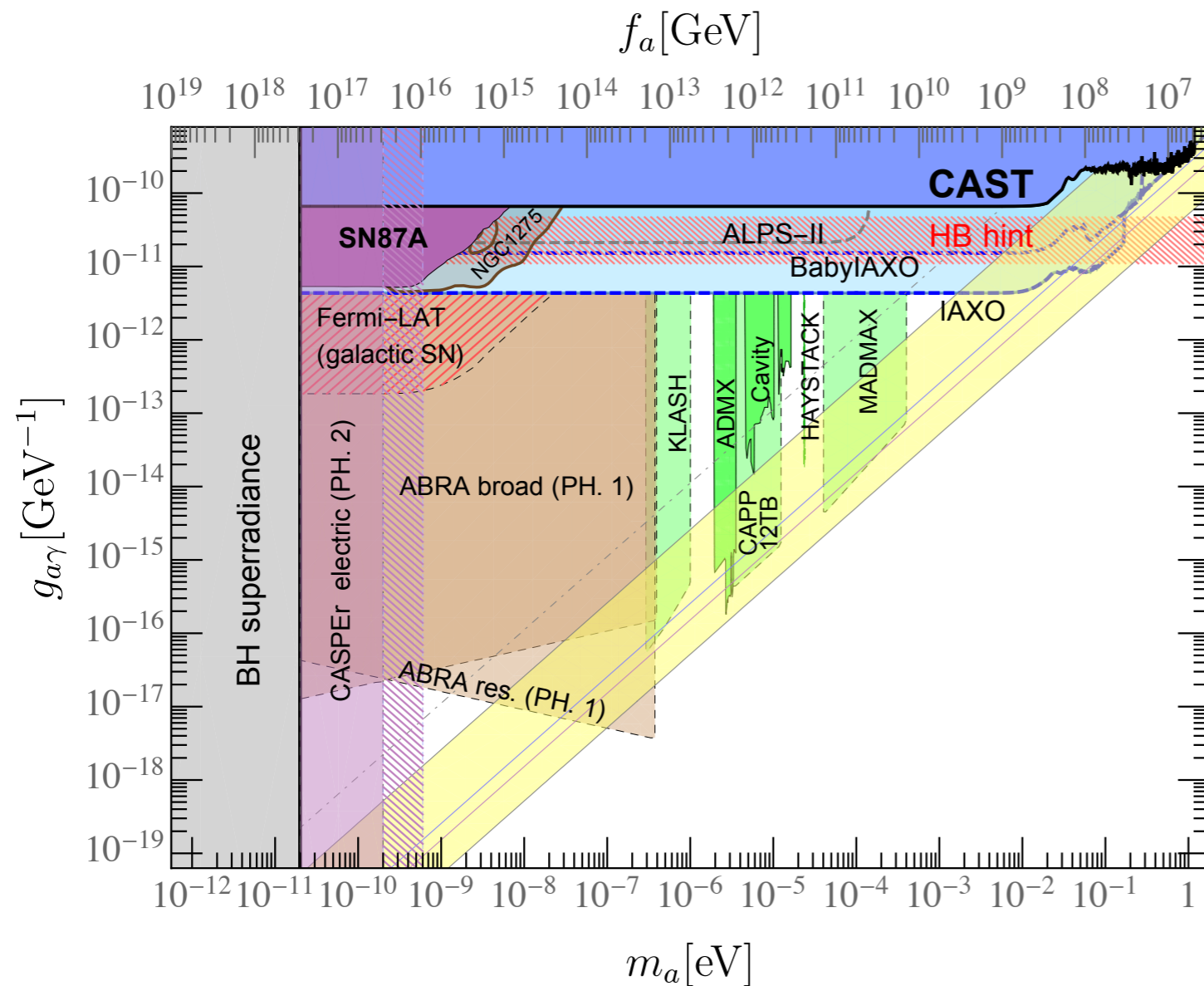


[LDL, Giannotti, Nardi, Visinelli 2003.01100 (Phys. Rept.)]

❖ An experimental opportunity



In 10 years from now ?



[LDL, Giannotti, Nardi, Visinelli 2003.01100 (Phys. Rept.)]

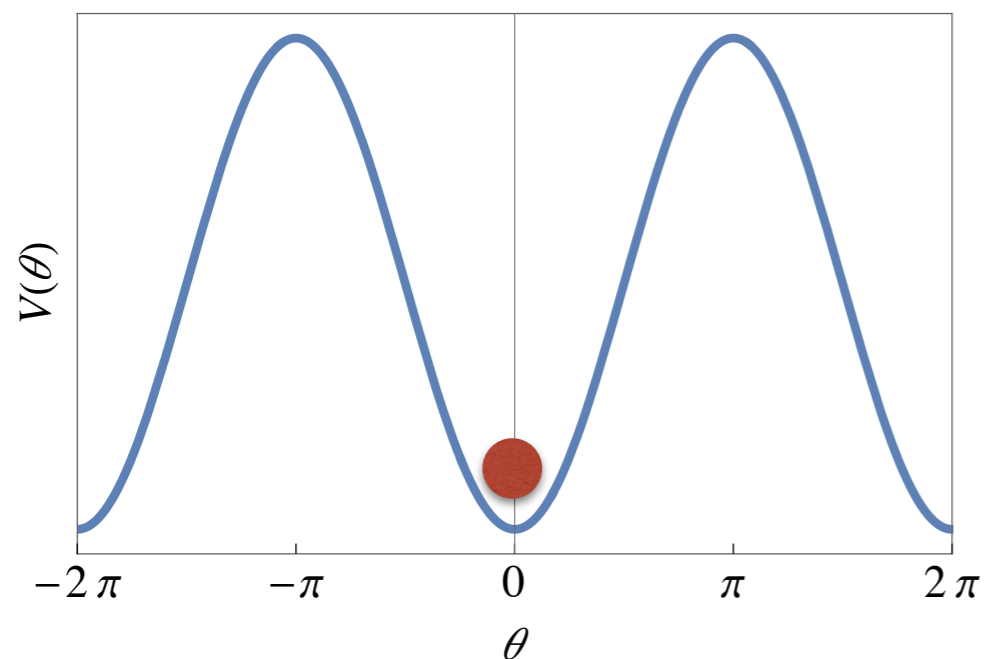
- ♣ An experimental opportunity
 - ★ Time now to rethink the QCD axion
1. PQ mechanism
 2. Axion couplings
[from EFTs to UV models]
 3. QCD axions beyond standard benchmarks

QCD axion

Strong CP problem

$$\delta\mathcal{L}_{\text{QCD}} = \theta \frac{g_s^2}{32\pi^2} G\tilde{G} \quad |\theta| \lesssim 10^{-10}$$

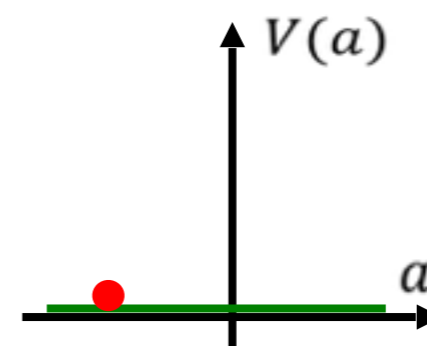
promote θ to a dynamical field,
which *washes-out* CP violation in QCD



$$\theta \rightarrow \frac{a}{f_a} \quad \text{with} \quad \langle a \rangle = 0$$

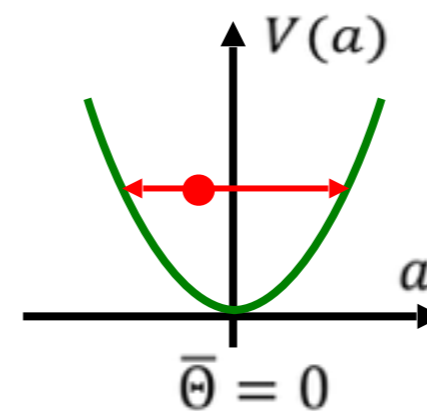
Dark Matter

misalignment + topological defects



$T \gg 1 \text{ GeV}$

[Raffelt]



$T \sim 1 \text{ GeV}$

$$w_a = p_a / \rho_a \simeq 0$$

$$\ddot{a} + 3H\dot{a} + m_a^2(T) f_a \sin\left(\frac{a}{f_a}\right) = 0$$

PQ mechanism

- Assume a new spin-0 boson with a pseudo-shift symmetry $a \rightarrow a + \alpha f_a$

broken by $\frac{a}{f_a} \frac{g_s^2}{32\pi^2} G\tilde{G}$  $E(0) \leq E(\langle a \rangle)$ [Vafa-Witten, PRL 53 (1984)]

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$$\theta_{\text{eff}} = \frac{\langle a \rangle}{f_a}$$

$$\begin{aligned} e^{-V_4 E(\theta_{\text{eff}})} &= \int \mathcal{D}\varphi e^{-S_0 + i\theta_{\text{eff}} \int G\tilde{G}} \\ &= \left| \int \mathcal{D}\varphi e^{-S_0 + i\theta_{\text{eff}} \int G\tilde{G}} \right| \\ &\leq \int \mathcal{D}\varphi \left| e^{-S_0 + i\theta_{\text{eff}} \int G\tilde{G}} \right| = e^{-V_4 E(0)} \end{aligned}$$

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- Assume a new spin-0 boson with a pseudo-shift symmetry $a \rightarrow a + \alpha f_a$

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$$\theta_{\text{eff}} = \frac{\langle a \rangle}{f_a}$$

$$e^{-V_4 E(\theta_{\text{eff}})} = \int \mathcal{D}\varphi e^{-S_0 + i\theta_{\text{eff}} \int G\tilde{G}}$$

$$= \left| \int \mathcal{D}\varphi e^{-S_0 + i\theta_{\text{eff}} \int G\tilde{G}} \right|$$

$$\leq \int \mathcal{D}\varphi \left| e^{-S_0 + i\theta_{\text{eff}} \int G\tilde{G}} \right| = e^{-V_4 E(0)}$$

- Does the axion really relax to zero ?

$$\mathcal{D}\varphi \equiv dA_\mu^a \det(\not{D} + M)$$



path-integral measure positive definite only for a vector-like theory (e.g. QCD)
does not apply to the SM !

PQ mechanism

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broken by $\frac{a}{f_a} \frac{g_s^2}{32\pi^2} G\tilde{G}$  $E(0) \leq E(\langle a \rangle)$ [Vafa-Witten, PRL 53 (1984)]


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$$\leq \int \mathcal{D}\varphi \left| e^{-S_0 + i\theta_{\text{eff}} \int G\tilde{G}} \right| = e^{-V_4 E(0)}$$

- Does the axion really relax to zero ?

 $\theta_{\text{eff}} \sim G_F^2 f_\pi^4 j_{\text{CKM}} \approx 10^{-18}$ [Georgi Randall, NPB276 (1986)
Okawa, Pospelov, Ritz, 2111.08040]

PQ mechanism works accidentally in the SM !

$$j_{\text{CKM}} = \text{Im} V_{ud} V_{cd}^* V_{cs} V_{us}^* \approx 10^{-5}$$

PQ mechanism

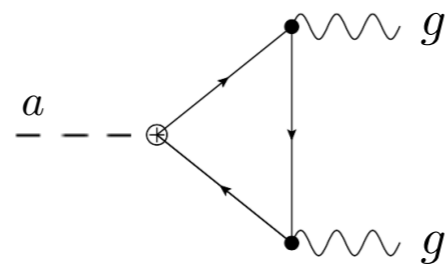
- Assume a new spin-0 boson with a pseudo-shift symmetry $a \rightarrow a + \alpha f_a$

broken by $\frac{a}{f_a} \frac{g_s^2}{32\pi^2} G\tilde{G}$  $E(0) \leq E(\langle a \rangle)$

- its origin can be traced back to a global $U(1)_{PQ}$ [Peccei, Quinn '77, Weinberg '78, Wilczek '78]

1. *spontaneously broken (axion is the associated pNGB)*

2. *QCD anomalous*

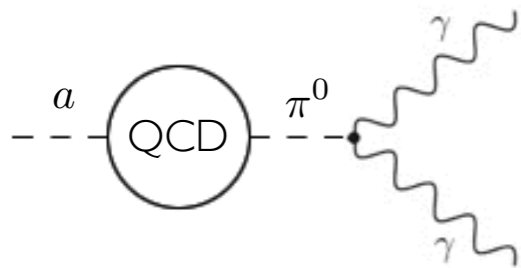


$$\partial^\mu J_\mu^{PQ} = \frac{N\alpha_s}{4\pi} G \cdot \tilde{G}$$

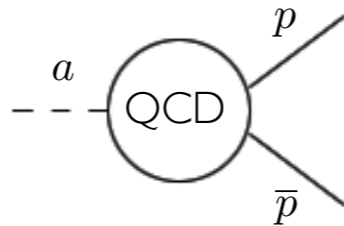
Axion properties [model-indep.]

- Consequences of $\frac{a}{f_a} \frac{g_s^2}{32\pi^2} G\tilde{G}$

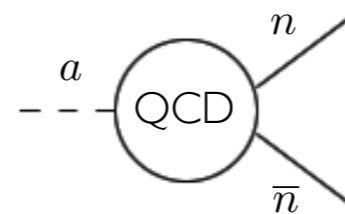
2. 'model-independent' axion couplings to photons, nucleons, electrons, ...



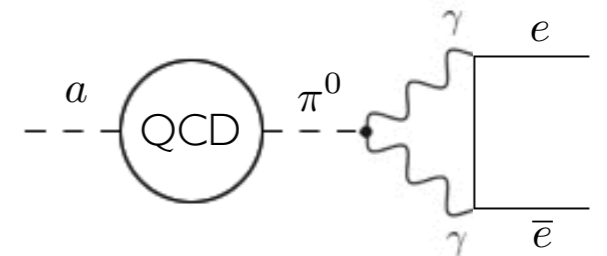
$$C_\gamma = -1.92(4)$$



$$C_p = -0.47(3)$$



$$C_n = -0.02(3)$$



$$C_e = -7.8(2) \times 10^{-6} \log\left(\frac{f_a}{m_e}\right)$$

$$\mathcal{L}_a \supset \frac{\alpha}{8\pi} \frac{C_\gamma}{f_a} a F_{\mu\nu} \tilde{F}^{\mu\nu} + \frac{C_f}{2f_a} \partial_\mu a \bar{f} \gamma^\mu \gamma_5 f \quad (f = p, n, e)$$

[Grilli di Cortona, Hardy, Vega, Villadoro, 1511.02867
 Lu, Du, Guo, Meißner, Vonk, 2003.01625
 Choi, Im, Kim, Seong, 2106.05816]

[see talk by G. Piazza for axion-pion coupling]

Axion properties [model-indep.]

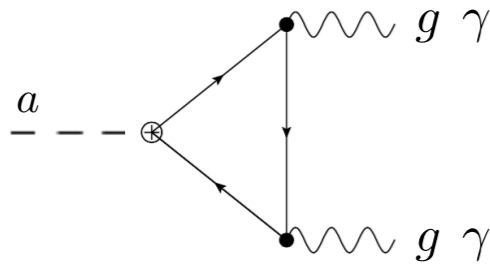
- Consequences of $\frac{a}{f_a} \frac{g_s^2}{32\pi^2} G\tilde{G}$

3. EFT breaks down at energies of order f_a

 *UV completion can drastically affect low-energy axion properties!*

Axion properties [model-dep.]

I. Axion-photon



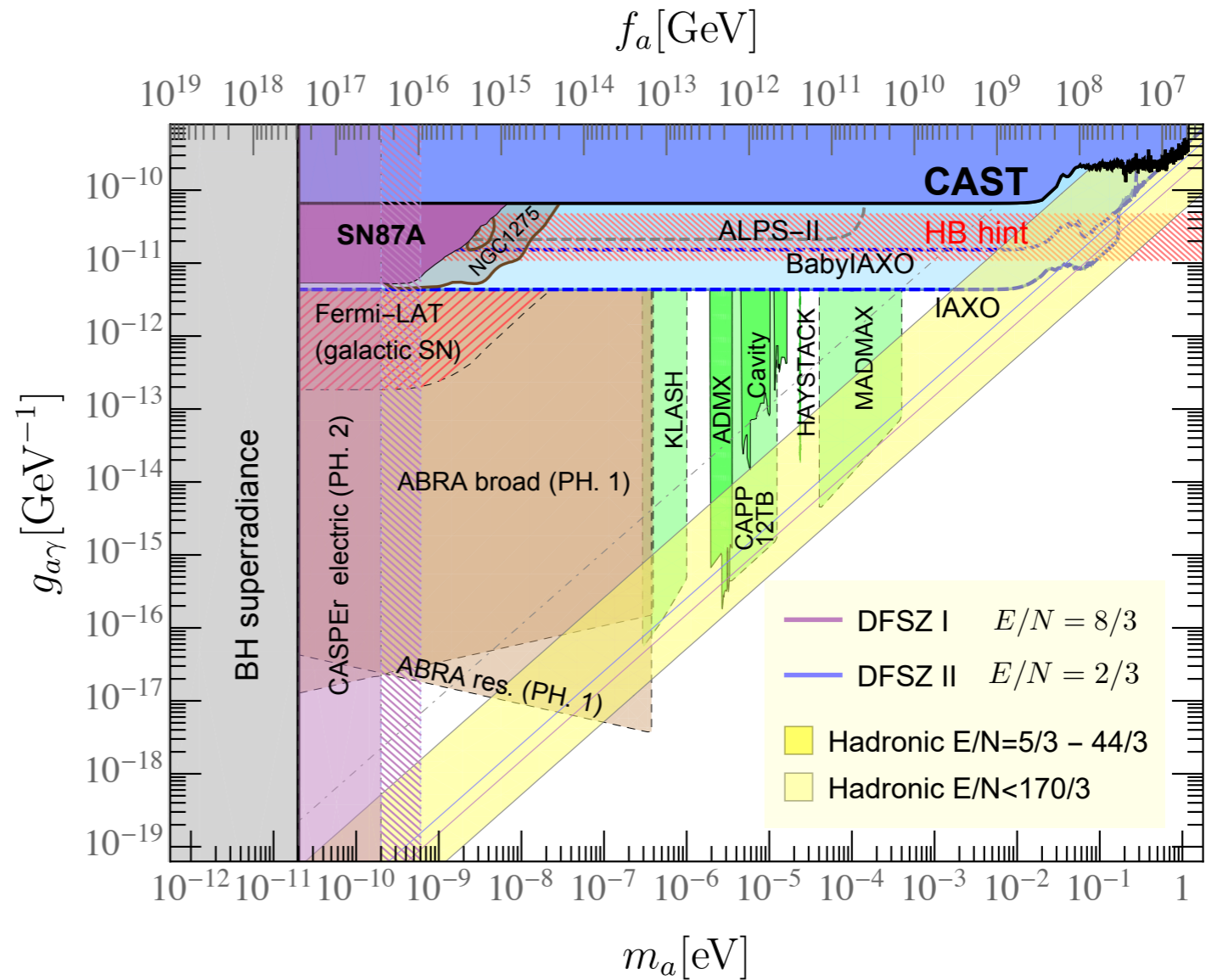
$$\partial^\mu J_\mu^{PQ} = \frac{N\alpha_s}{4\pi} G \cdot \tilde{G} + \frac{E\alpha}{4\pi} F \cdot \tilde{F}$$

$$C_\gamma = E/N - 1.92(4)$$

model-independent

depends on UV completion

enhance/suppress C_γ

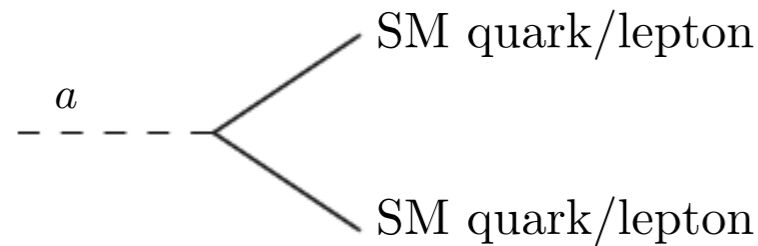


[LDL, Giannotti, Nardi, Visinelli 2003.01100 (Phys. Rept.)]

Axion properties [model-dep.]

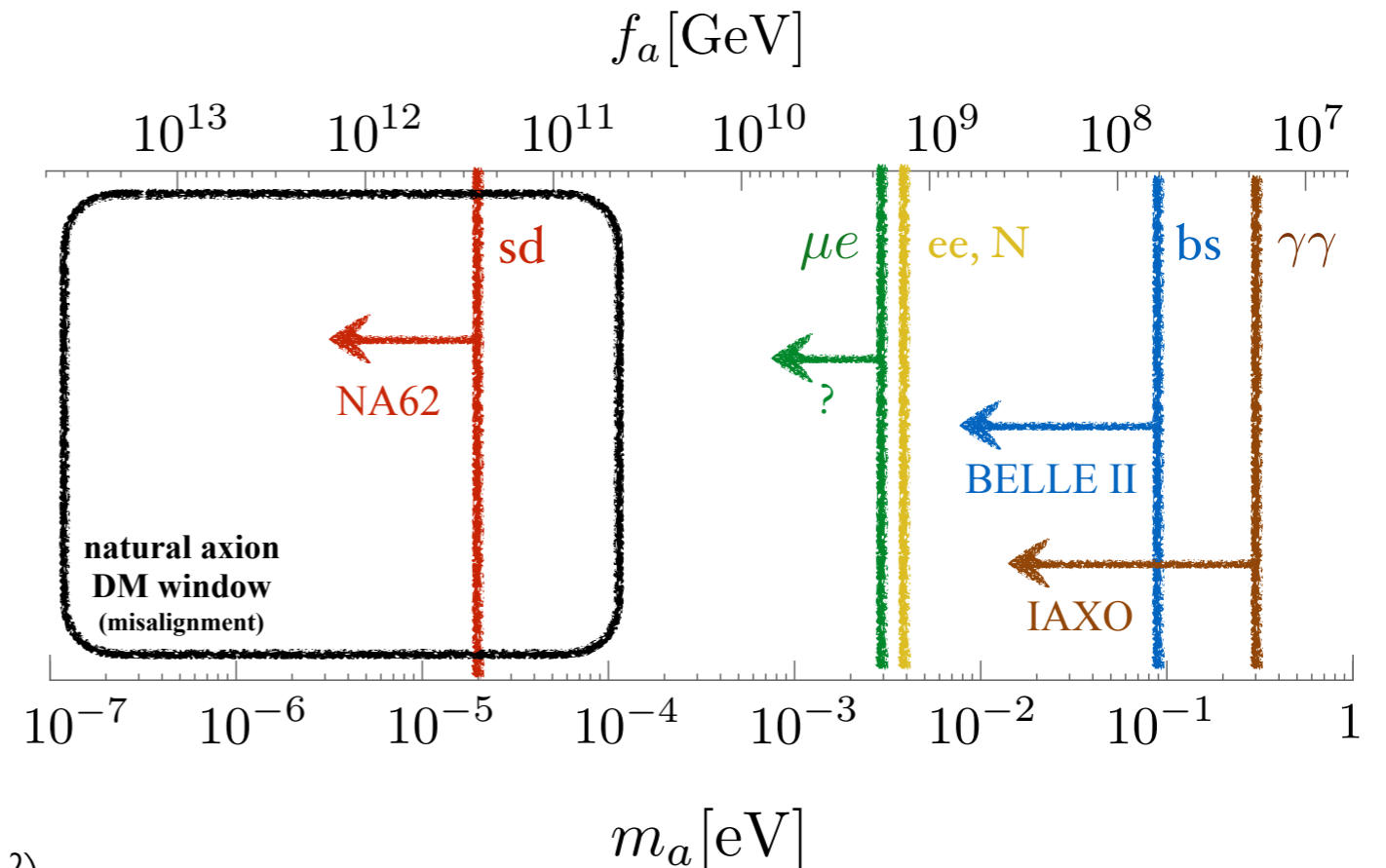
2. Axion-SM fermions

[Robert Ziegler, La Thuile'19]



$$\frac{\partial_\mu a}{2f_a} \bar{\psi}_i \gamma^\mu (C_{ij}^V + C_{ij}^A \gamma_5) \psi_j$$

$$J_{PQ}^\mu$$



$$C_{ij}^{V,A} \propto (V_\psi^\dagger PQ_\psi V_\psi)_{ij} \quad (\text{PQ as a flavour symmetry ?})$$

flavour-violating axion coupling

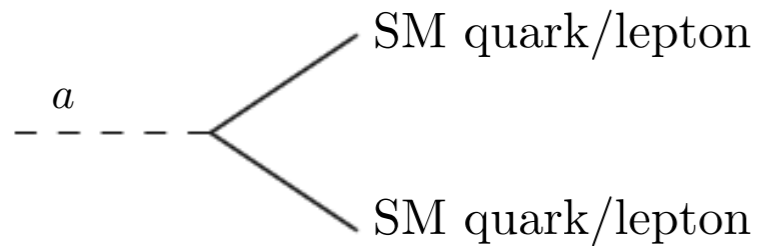
enhance/suppress $C_{p,n,e}$

for $C_i = \{C_\gamma, C_e, C_N, C_{sd}^V, C_{bs}^V\} = 1$
flavour beats astrophysics!

Axion properties [model-dep.]

2. Axion-SM fermions

[LDL, Mescia, Nardi, Panci, Ziegler, 1712.04940 + 1907.06575
"Astrophobic Axions"]

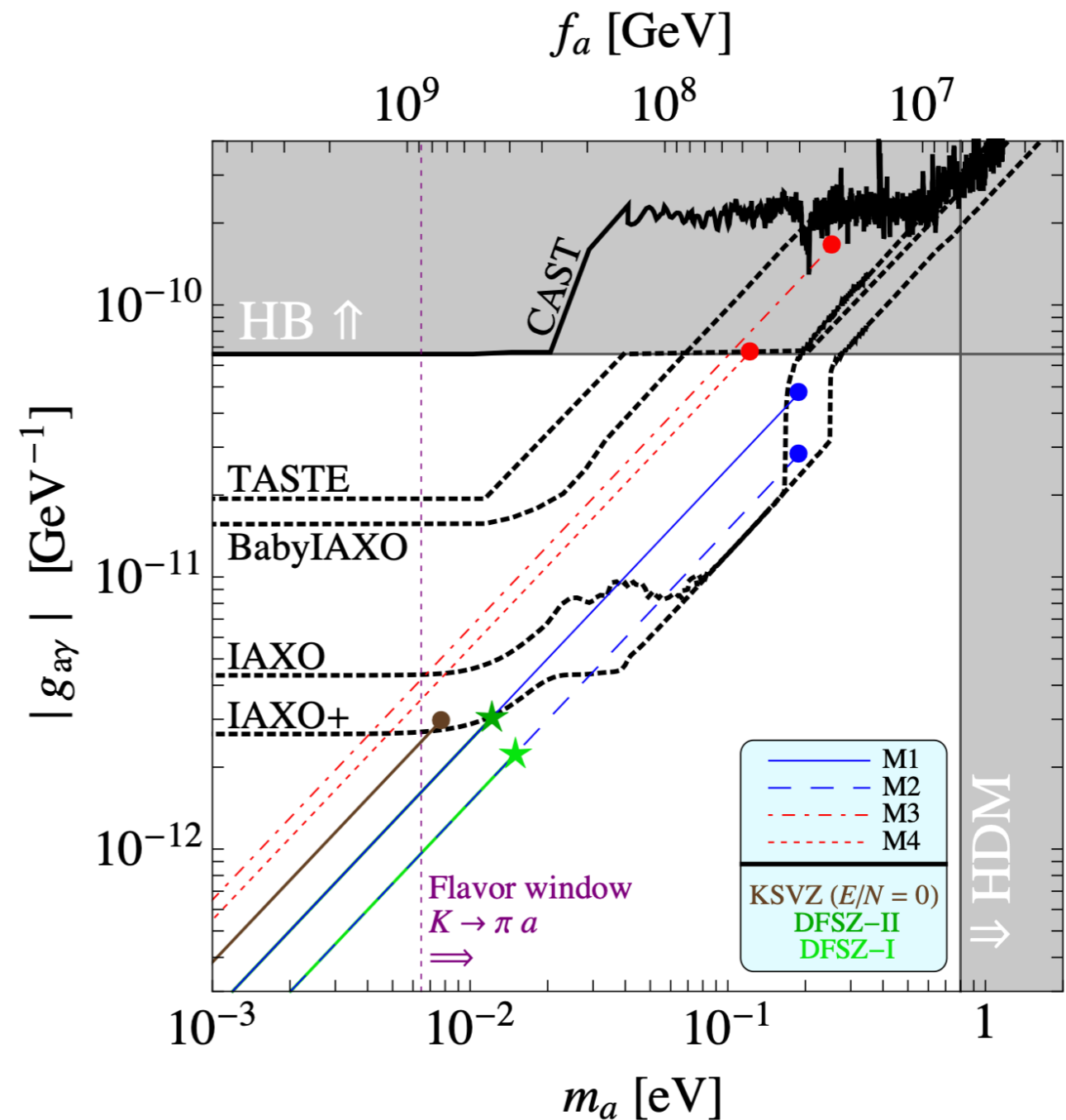


$$\frac{\partial_\mu a}{2f_a} \bar{\psi}_i \gamma^\mu (C_{ij}^V + C_{ij}^A \gamma_5) \psi_j$$

$$J_{PQ}^\mu$$

flavour-violating axion coupling

enhance/suppress $C_{p,n,e}$



Axion properties [model-dep.]

3. CP-violating axions

[Moody, Wilczek PRD30 (1984)]

$$\mathcal{L} \supset g_{aN}^S a \bar{N} N + g_{af}^P a \bar{f} i \gamma_5 f$$

$$g_{aN}^S \sim \frac{f_\pi}{f_a} \theta_{\text{eff}}$$



from UV sources of CP-violation
e.g. $\mathcal{O}_{\text{CPV}} = (\bar{u}u)(\bar{d}i\gamma_5 d)$

[Barbieri, Romanino, Strumia hep-ph/9605368
Pospelov hep-ph/9707431
Bigazzi, Cotrone, Jarvinen, Kiritsis 1906.12132
Bertolini, LDL, Nesti 2006.12508
Okawa, Pospelov, Ritz, 2111.08040
Dekens, de Vries, Shain, 2203.11230]

$$V(a) \simeq \frac{1}{2} \frac{a^2}{f_a^2} \underbrace{\langle G\tilde{G}, G\tilde{G} \rangle}_\chi + \frac{a}{f_a} \underbrace{\langle G\tilde{G}, \mathcal{O}_{\text{CPV}} \rangle}_{\chi'}$$



$$\theta_{\text{eff}} \equiv \frac{\langle a \rangle}{f_a} = -\frac{\chi'}{\chi}$$



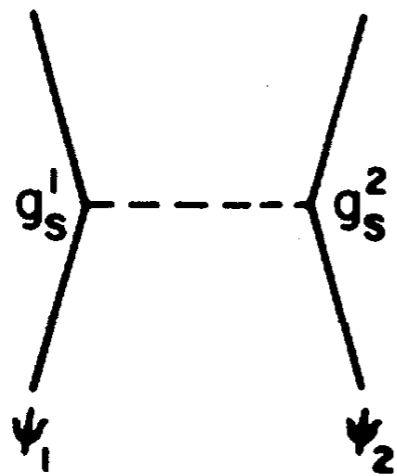
Axion properties [model-dep.]

3. CP-violating axions

[Moody, Wilczek PRD30 (1984)]

$$\mathcal{L} \supset g_{aN}^S a \bar{N} N + g_{af}^P a \bar{f} i \gamma_5 f \quad g_{aN}^S \sim \frac{f_\pi}{f_a} \theta_{\text{eff}} \quad \leftarrow \text{from UV sources of CP-violation}$$

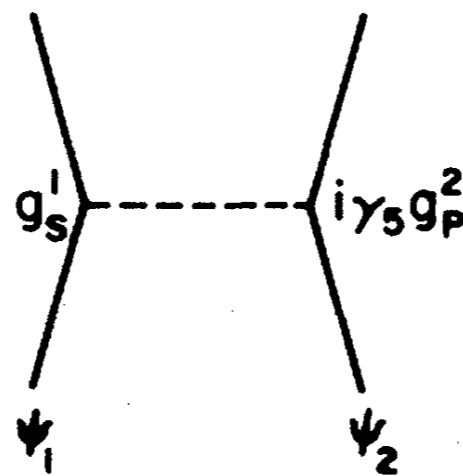
New macroscopic forces from non-relativistic potentials



(a)

monopole-monopole

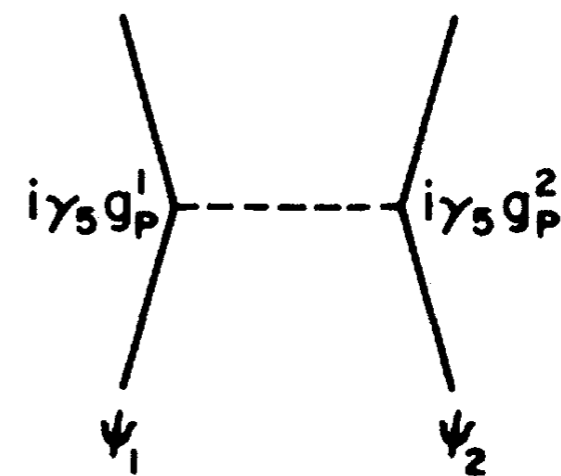
double θ_{eff} suppression



(b)

monopole-dipole

ARIADNE, QUAX-gpgs, ...
NMR enhancement



(c)

dipole-dipole

spin suppression + bkgd
from ordinary magnetic forces

Axion properties [model-dep.]

3. CP-violating axions

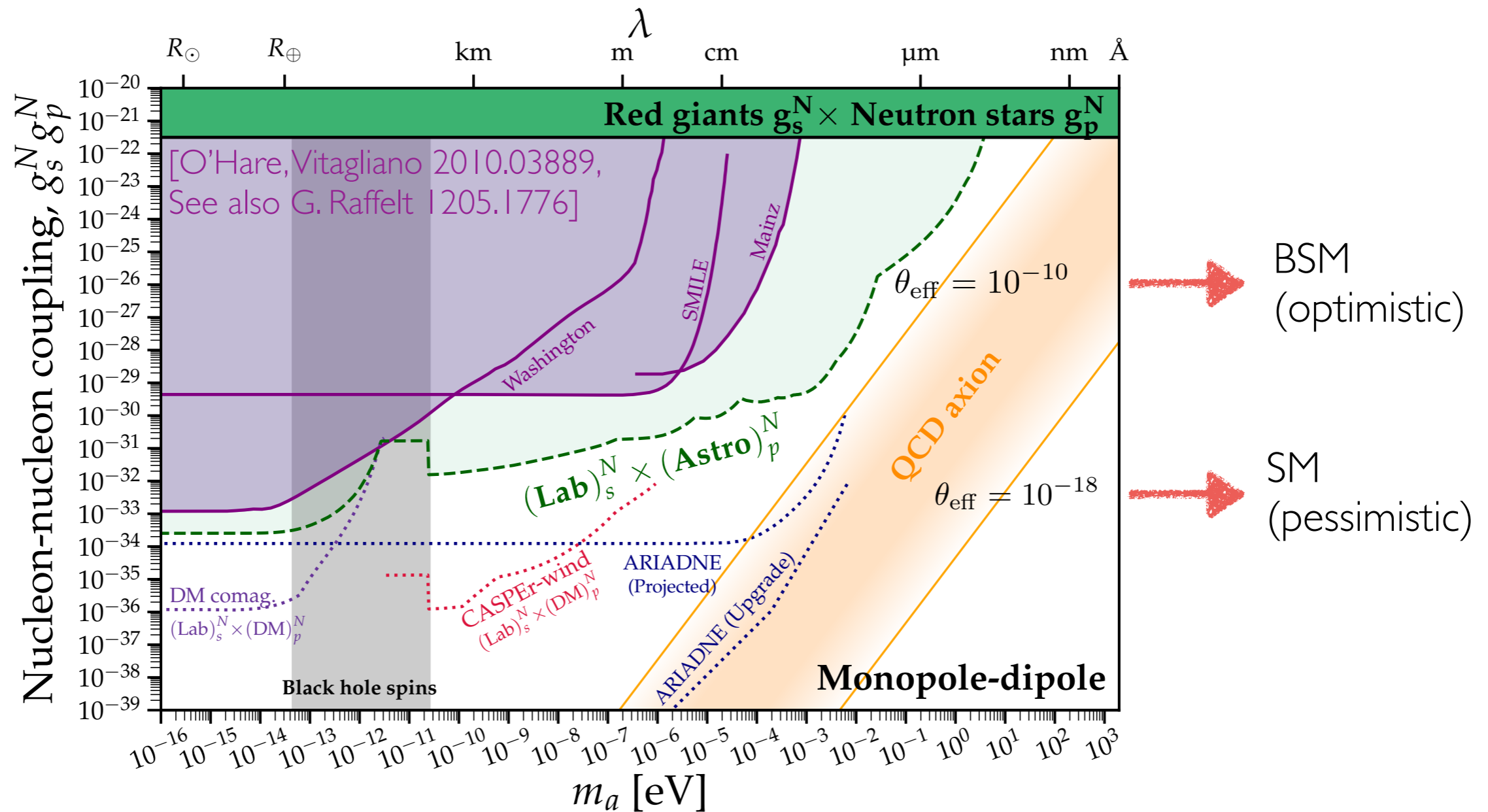
[Moody, Wilczek PRD30 (1984)]

$$\mathcal{L} \supset g_{aN}^S a \bar{N} N + g_{af}^P a \bar{f} i \gamma_5 f$$

$$g_{aN}^S \sim \frac{f_\pi}{f_a} \theta_{\text{eff}}$$



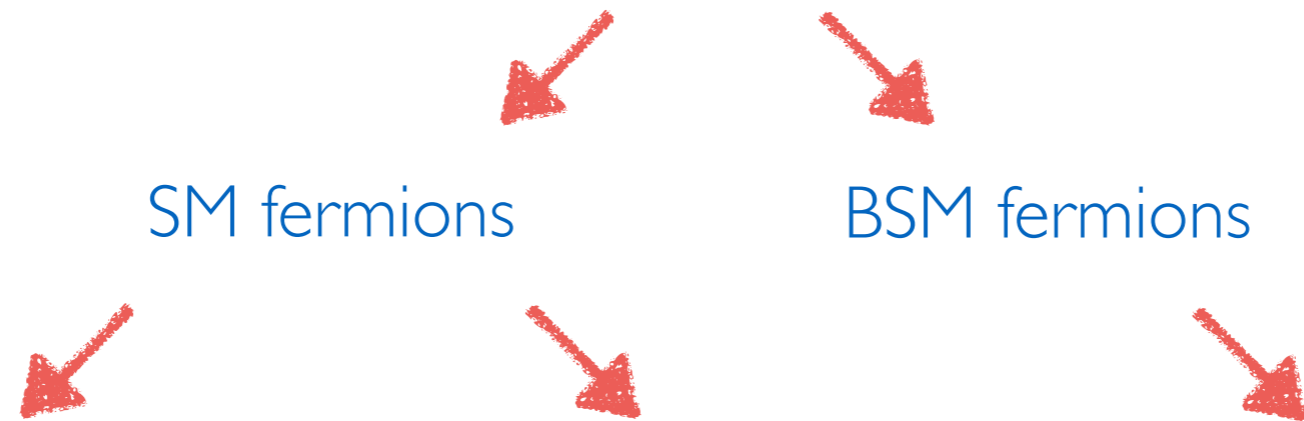
from UV sources of CP-violation



Benchmark axion models

- global $U(1)_{PQ}$ (*QCD anomalous* + *spontaneously broken*)

$$U(1)_{PQ} \times SU(3)_c^2$$



2Higgs

PQWW

[Peccei, Quinn '77,
Weinberg '78, Wilczek '78]

$f_a \sim v$ ruled out

2Higgs+Singlet

DFSZ

[Zhitnitsky '80,
Dine, Fischler, Srednicki '81]

$f_a \gg v$ "Invisible" axion (phase of singlet field)

BSM fermions

Higgs+Singlet

KSVZ

[Kim '79,
Shifman, Vainshtein, Zakharov '80]

Benchmark axion models

- global $U(1)_{\text{PQ}}$ (*QCD anomalous* + *spontaneously broken*)

$$U(1)_{\text{PQ}} \times SU(3)_c^2$$



SM fermions

BSM fermions



2Higgs+Singlet

Higgs+Singlet

DFSZ

KSVZ

$$C_\gamma = E/N - 1.92(4)$$

$$E/N = 8/3$$

$$E/N = 0$$

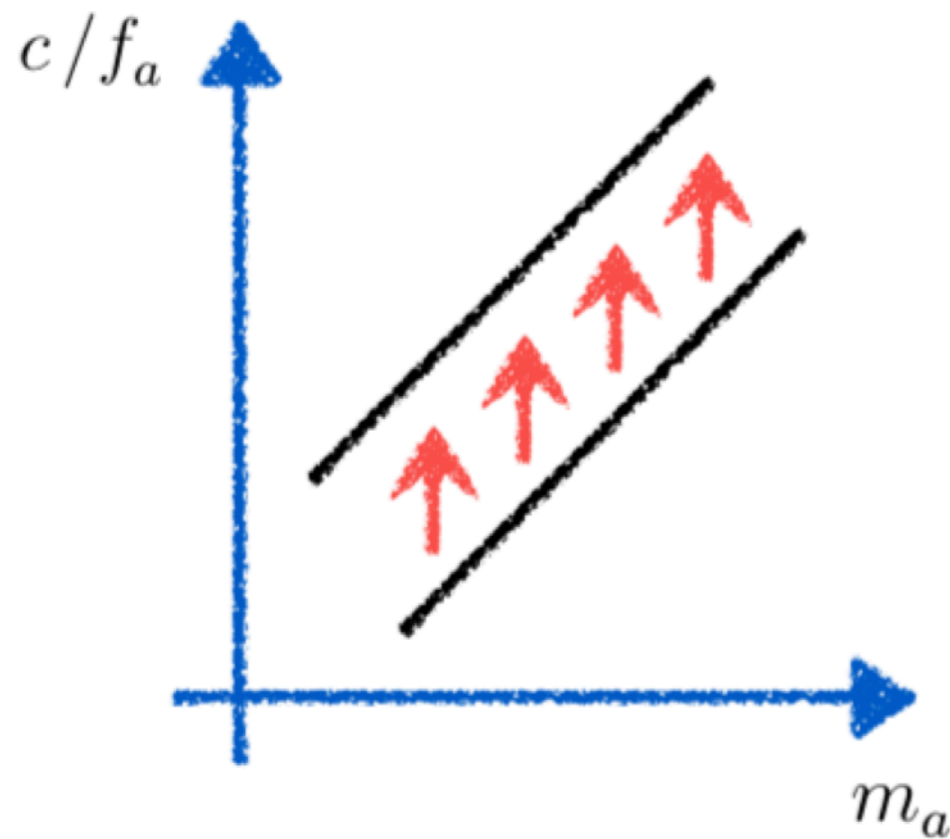
$$C_{p,n,e}(\beta) \sim \mathcal{O}(1)$$

$$C_p \simeq -0.5$$

$$C_{n,e} \simeq 0$$

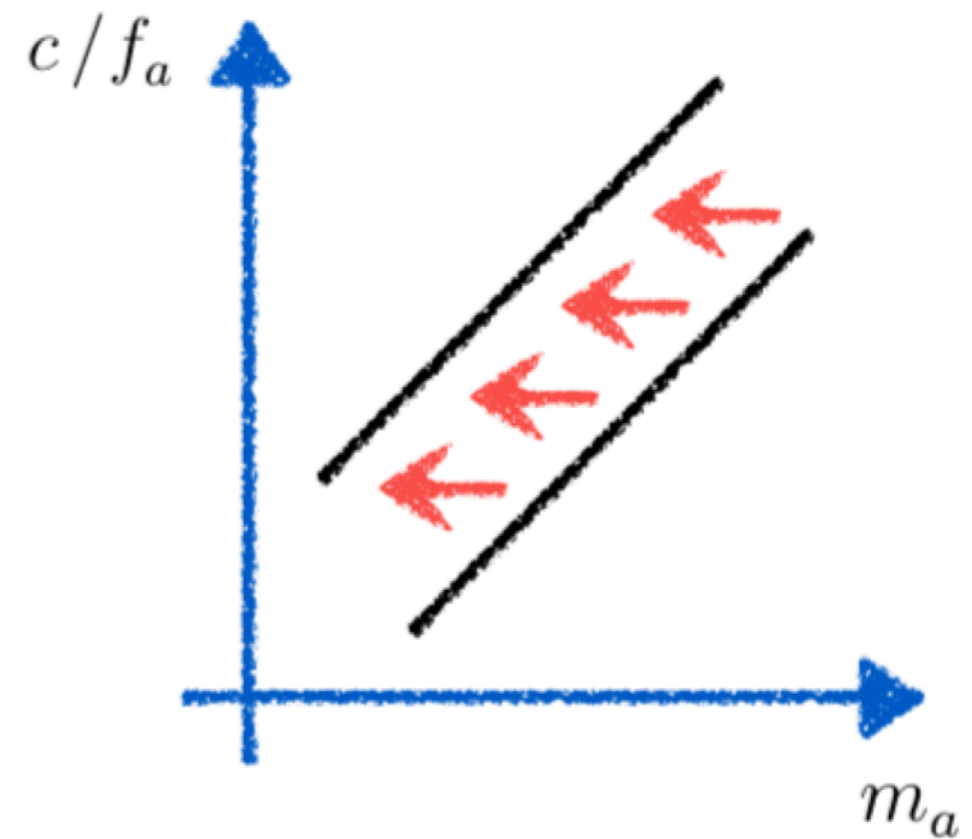
(also no flavour and CP-violating effects)

Axions beyond benchmarks



enhance Wilson coefficient for fixed m_a

[LDL, Mescia, Nardi 1610.07593 + 1705.05370
Farina, Pappadopulo, Rompineve, Tesi 1611.09855
Agrawal, Fan, Reece, Wang 1709.06085
Darne', LDL, Giannotti, Nardi 2010.15846
Ringwald, Sokolov 2104.02574]



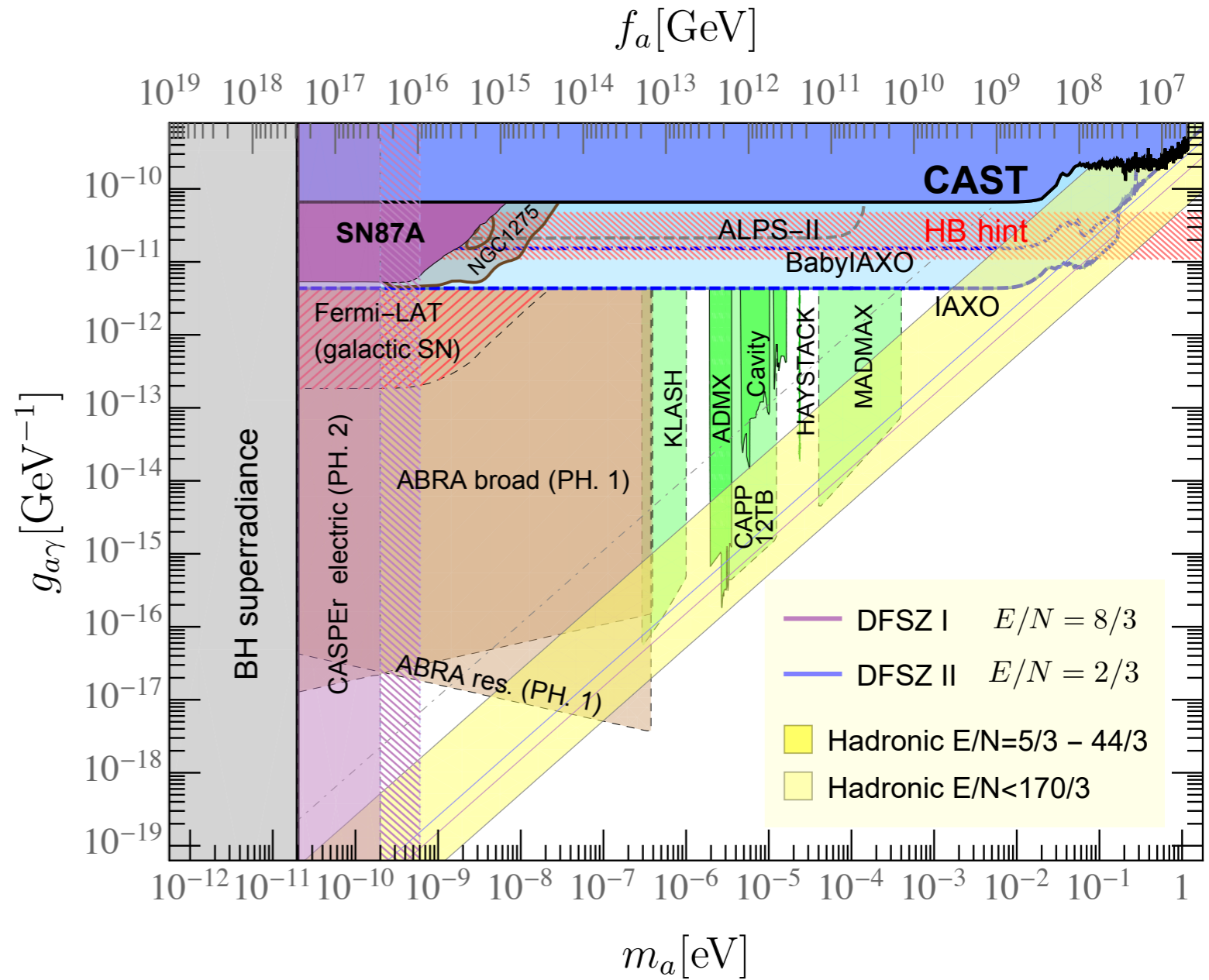
suppress axion mass for fixed f_a

[Hook 1802.10093,
LDL, Gavela, Quilez, Ringwald 2102.00012
+ 2102.01082]

→ QCD axion parameter space much larger than what traditionally thought

Axion-Photon

EXP	STATUS
CAST (CERN)	finished
ADMX (Seattle)	running
HAYSTAC (New Haven)	running
ALPs-II (DESY)	construction
CAPP (South Korea)	construction
ORGAN (Perth)	prototype
ABRACADABRA (MIT)	prototype
(Baby)IAXO (DESY)	preparation
MADMAX (DESY)	preparation
ACTION (South Korea)	proposed
FLASH (Frascati)	CDR
QUAX- $a\gamma$ (Legnaro)	running
...	...



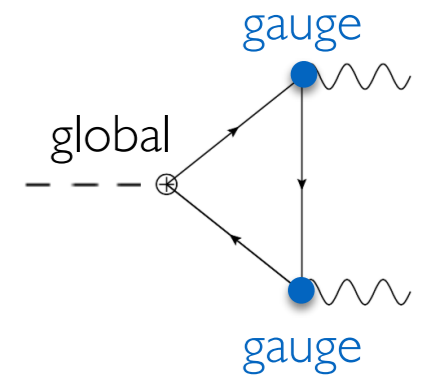
[LDL, Giannotti, Nardi, Visinelli 2003.01100 (Phys. Rept.)]

Enhancing $g_{a\gamma}$

$$g_{a\gamma} = \frac{\alpha}{2\pi} \frac{C_{a\gamma}}{f_a}$$

$$C_{a\gamma} = E/N - 1.92(4)$$

$$\partial^\mu J_\mu^{PQ} = \frac{N\alpha_s}{4\pi} G \cdot \tilde{G} + \frac{E\alpha}{4\pi} F \cdot \tilde{F}$$

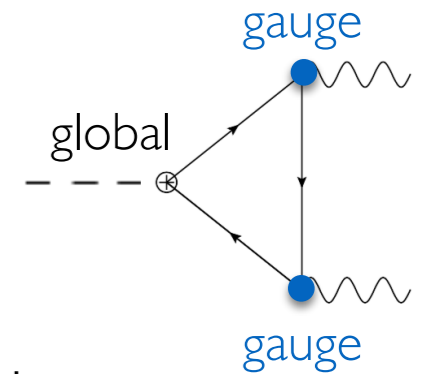


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	R_Q	\mathcal{O}_{Qq}	$\Lambda_{\text{Landau}}^{2\text{-loop}} [\text{GeV}]$	E/N
R_Q^w	(3, 1, -1/3)	$\bar{Q}_L d_R$	$9.3 \cdot 10^{38} (g_1)$	2/3
	(3, 1, 2/3)	$\bar{Q}_L u_R$	$5.4 \cdot 10^{34} (g_1)$	8/3
	(3, 2, 1/6)	$\bar{Q}_R q_L$	$6.5 \cdot 10^{39} (g_1)$	5/3
	(3, 2, -5/6)	$\bar{Q}_L d_R H^\dagger$	$4.3 \cdot 10^{27} (g_1)$	17/3
	(3, 2, 7/6)	$\bar{Q}_L u_R H$	$5.6 \cdot 10^{22} (g_1)$	29/3
	(3, 3, -1/3)	$\bar{Q}_R q_L H^\dagger$	$5.1 \cdot 10^{30} (g_2)$	14/3
R_Q^s	(3, 3, 2/3)	$\bar{Q}_R q_L H$	$6.6 \cdot 10^{27} (g_2)$	20/3
	(3, 3, -4/3)	$\bar{Q}_L d_R H^{\dagger 2}$	$3.5 \cdot 10^{18} (g_1)$	44/3
	($\bar{6}$, 1, -1/3)	$\bar{Q}_L \sigma_{\mu\nu} d_R G^{\mu\nu}$	$2.3 \cdot 10^{37} (g_1)$	4/15
	($\bar{6}$, 1, 2/3)	$\bar{Q}_L \sigma_{\mu\nu} u_R G^{\mu\nu}$	$5.1 \cdot 10^{30} (g_1)$	16/15
	($\bar{6}$, 2, 1/6)	$\bar{Q}_R \sigma_{\mu\nu} q_L G^{\mu\nu}$	$7.3 \cdot 10^{38} (g_1)$	2/3
	(8, 1, -1)	$\bar{Q}_L \sigma_{\mu\nu} e_R G^{\mu\nu}$	$7.6 \cdot 10^{22} (g_1)$	8/3
	(8, 2, -1/2)	$\bar{Q}_R \sigma_{\mu\nu} \ell_L G^{\mu\nu}$	$6.7 \cdot 10^{27} (g_1)$	4/3
	(15, 1, -1/3)	$\bar{Q}_L \sigma_{\mu\nu} d_R G^{\mu\nu}$	$8.3 \cdot 10^{21} (g_3)$	1/6
	(15, 1, 2/3)	$\bar{Q}_L \sigma_{\mu\nu} u_R G^{\mu\nu}$	$7.6 \cdot 10^{21} (g_3)$	2/3

- Pheno preferred hadronic axions

1. Q-fermions short lived (no coloured relics)
2. No Landau poles below Planck



$$E/N \in [5/3, 44/3]$$

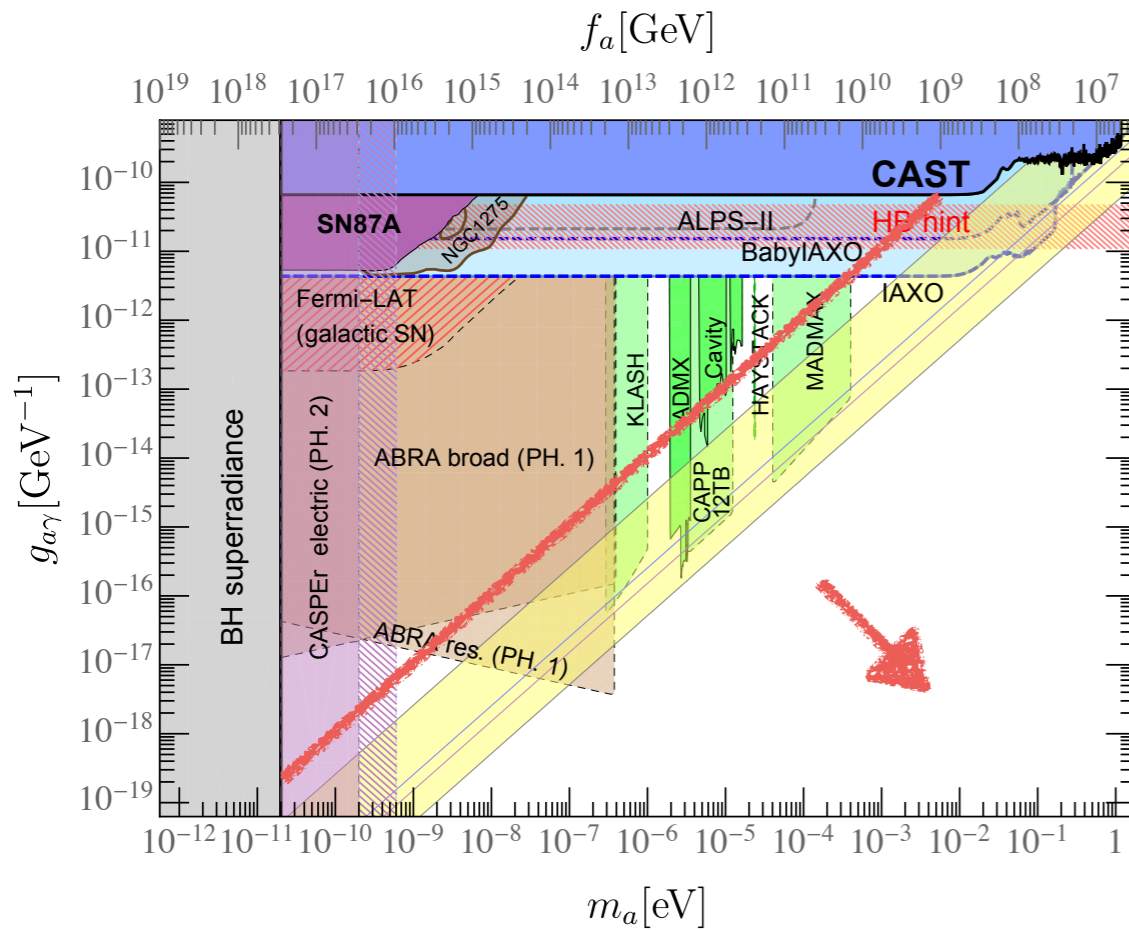
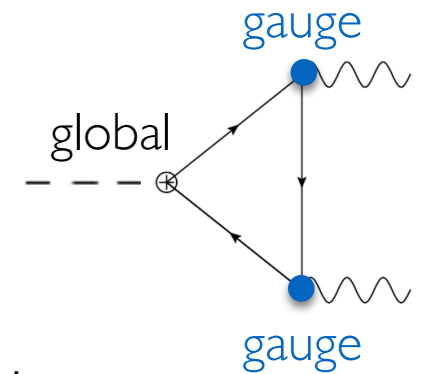
[LDL, Mescia, Nardi | 6 | 10.07593]

Enhancing $g_{a\gamma}$

$$g_{a\gamma} = \frac{\alpha}{2\pi} \frac{C_{a\gamma}}{f_a}$$

$$C_{a\gamma} = E/N - 1.92(4)$$

$$\partial^\mu J_\mu^{PQ} = \frac{N\alpha_s}{4\pi} G \cdot \tilde{G} + \frac{E\alpha}{4\pi} F \cdot \tilde{F}$$



- Pheno preferred hadronic axions

- More Q's? [LDL, Mescia, Nardi 1705.05370]

$$E/N < 170/3 \quad (\text{perturbativity})$$

$$g_{a\gamma} \rightarrow 0$$

[“such a cancellation is immoral, but not unnatural”,
D. B. Kaplan, NPB260 (1985)]

- Going above $E/N = 170/3$?

- boost global charge (clockwork) → backup slides

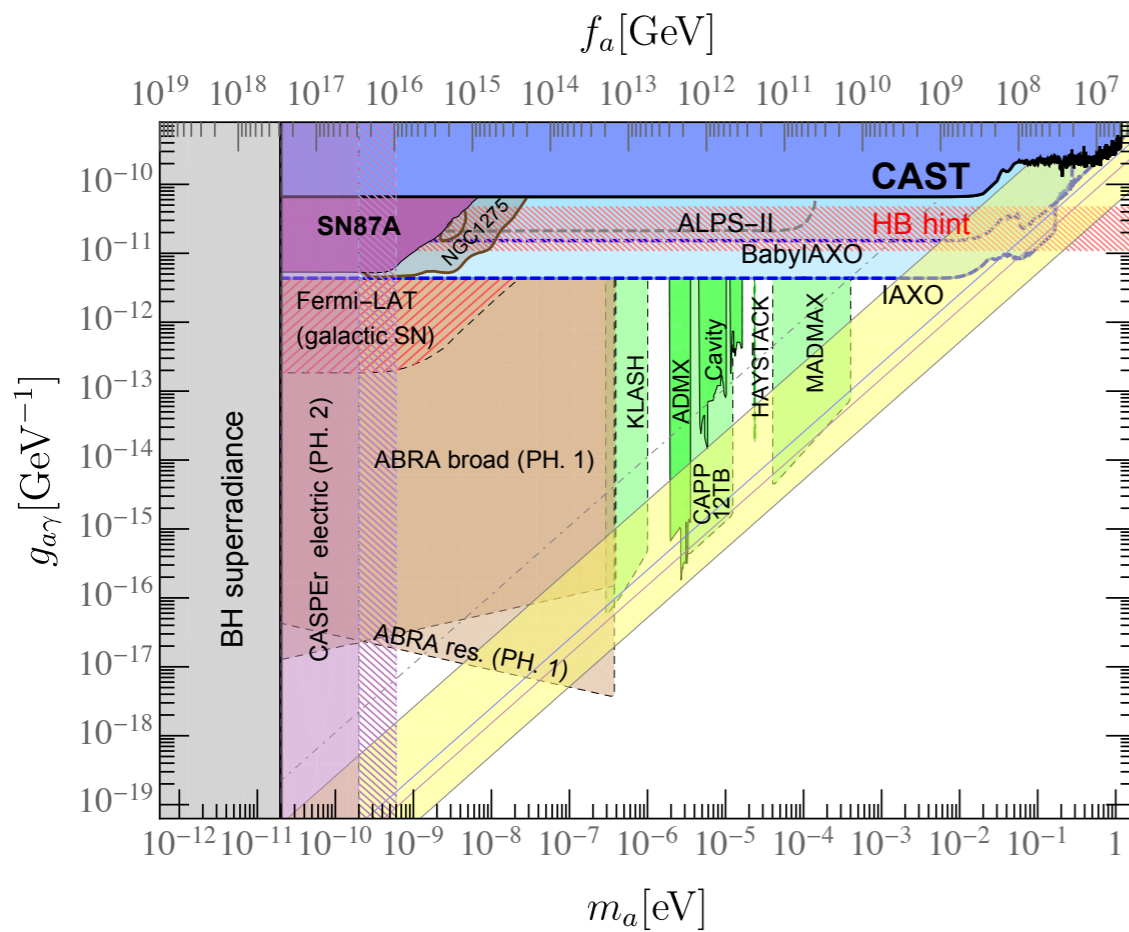
- be agnostic, E/N is a free parameter

Enhancing $g_{a\gamma}$

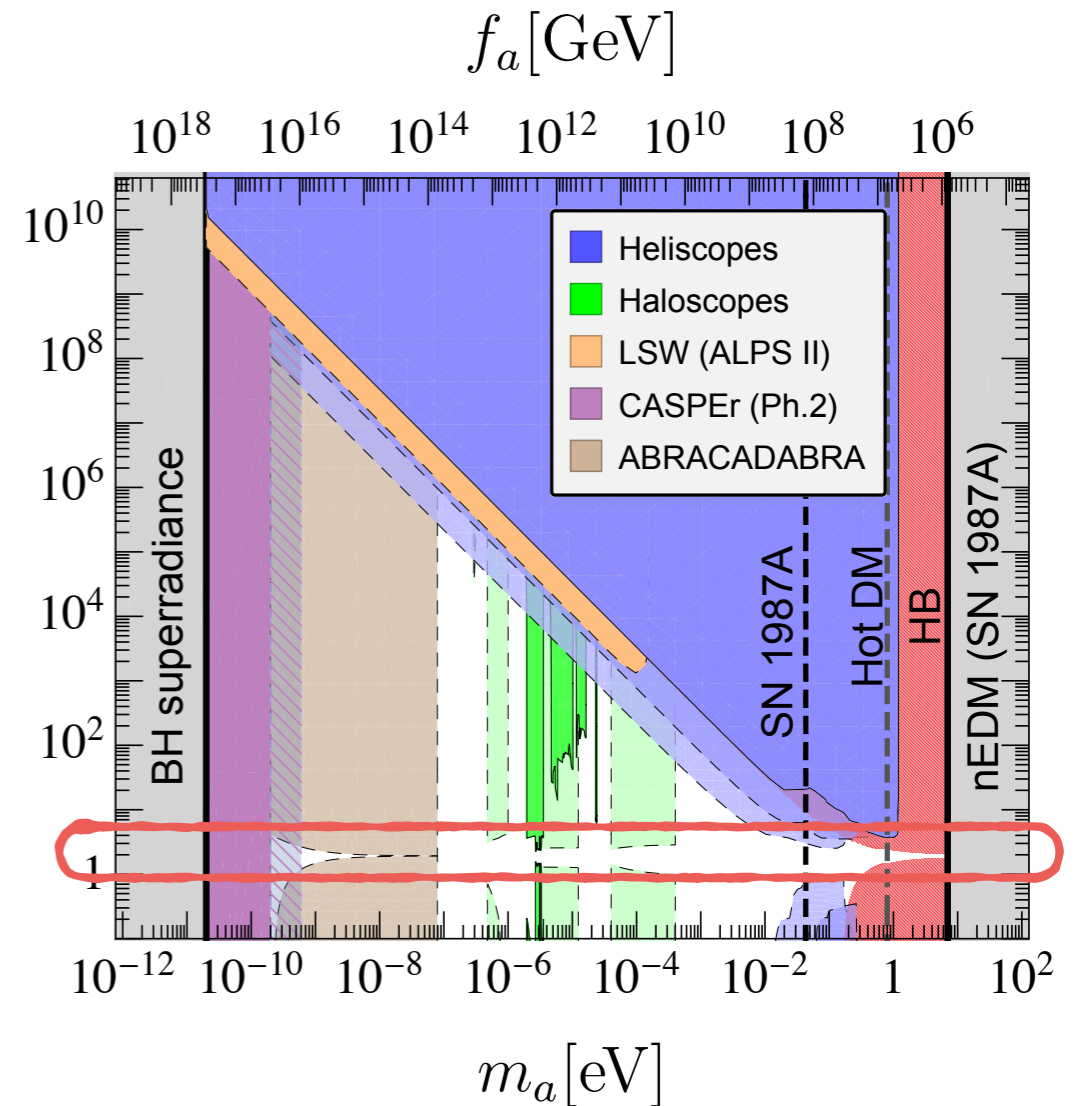
$$g_{a\gamma} = \frac{\alpha}{2\pi} \frac{C_{a\gamma}}{f_a}$$

$$C_{a\gamma} = E/N - 1.92(4)$$

[LDL, Giannotti, Nardi, Visinelli 2003.01100 (Phys. Rept.)]



E/N



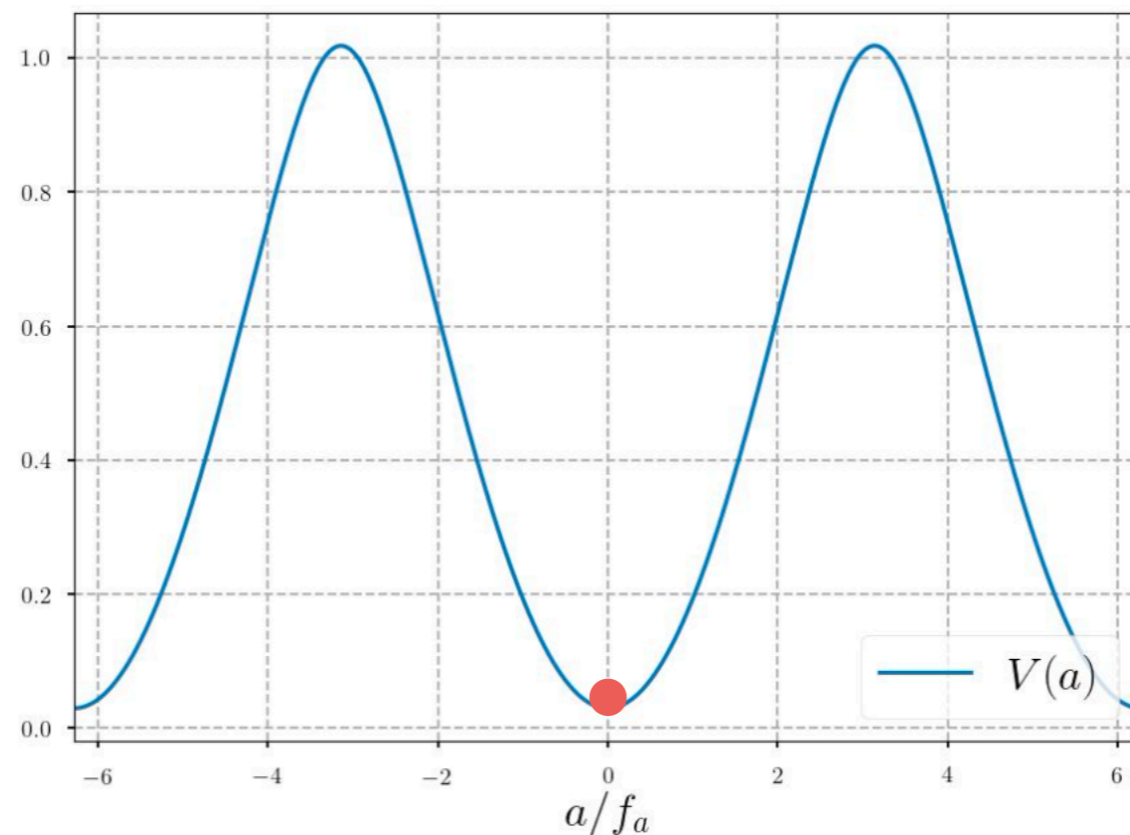
1. exp.s have just started to constrain E/N from above
2. $E/N \sim 1.92$ appears as a tuned region in theory space

Suppressing m_a

- Standard QCD axion

[Di Vecchia, Veneziano, NPB171 (1980)
Leutwyler, Smilga, PRD46 (1992)
Grilli di Cortona, Hardy, Vega, Villadoro, I511.02867]

$$\frac{a}{f_a} \frac{\alpha_s}{8\pi} G\tilde{G} \quad \longrightarrow \quad V(a) = -m_\pi^2 f_\pi^2 \sqrt{1 - \frac{4m_u m_d}{(m_u + m_d)^2} \sin^2\left(\frac{a}{2f_a}\right)}$$

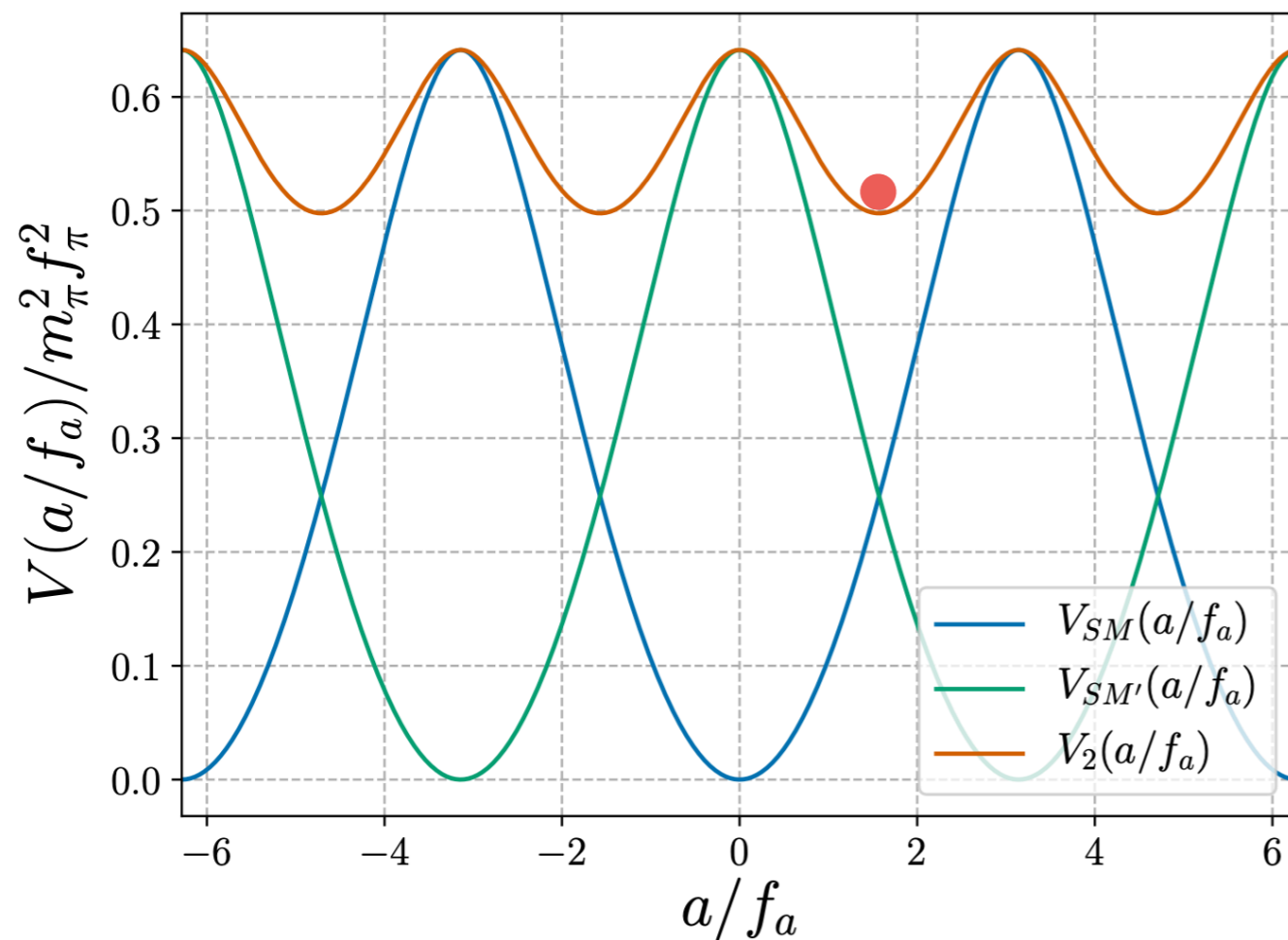


Suppressing m_a

- Z_2 axion: mirror world

$$\begin{aligned} \text{SM} &\longleftrightarrow \text{SM}' \\ a &\longrightarrow a + \pi f_a \end{aligned}$$

$$\mathcal{L} = \mathcal{L}_{\text{SM}} + \mathcal{L}_{\text{SM}'} + \frac{\alpha_s}{8\pi} \left(\frac{a}{f_a} - \theta \right) G\tilde{G} + \frac{\alpha_s}{8\pi} \left(\frac{a}{f_a} - \theta + \pi \right) G'\tilde{G}'$$

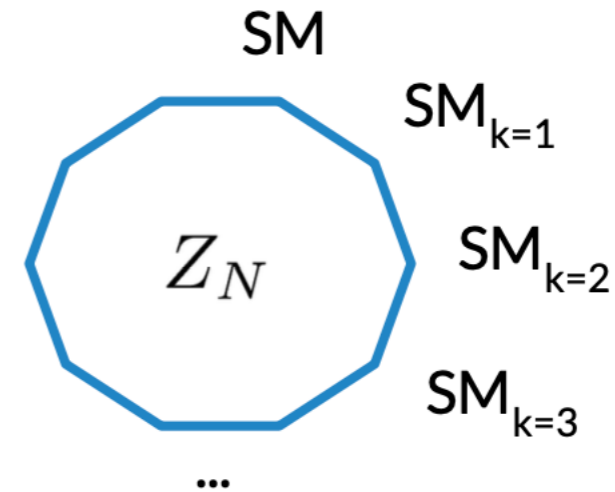


axion mass is suppressed
but minimum in $\pi/2$

Suppressing m_a

- Z_N axion: N mirror worlds [Hook 1802.10093]

$$\begin{aligned} \text{SM}_k &\longrightarrow \text{SM}_{k+1 \pmod{\mathcal{N}}} \\ a &\longrightarrow a + \frac{2\pi k}{\mathcal{N}} f_a, \end{aligned}$$



the axion ($\theta_a \equiv a/f_a$) realizes the Z_N symmetry non-linearly

$$\mathcal{L} = \sum_{k=0}^{\mathcal{N}-1} \left[\mathcal{L}_{\text{SM}_k} + \frac{\alpha_s}{8\pi} \left(\theta_a + \frac{2\pi k}{\mathcal{N}} \right) G_k \tilde{G}_k \right]$$

[LDL, Gavela, Quilez, Ringwald 2102.00012]

$$\longrightarrow V_{\mathcal{N}}(\theta_a) = -m_{\pi}^2 f_{\pi}^2 \sum_{k=0}^{\mathcal{N}-1} \sqrt{1 - \frac{4z}{(1+z)^2} \sin^2 \left(\frac{\theta_a}{2} + \frac{\pi k}{\mathcal{N}} \right)}$$

$$z \equiv \frac{m_u}{m_d} \sim 1/2$$

$$\simeq \frac{m_{\pi}^2 f_{\pi}^2}{\sqrt{\pi}} \sqrt{\frac{1-z}{1+z}} \mathcal{N}^{-1/2} (-1)^{\mathcal{N}} z^{\mathcal{N}} \cos(\mathcal{N}\theta_a)$$

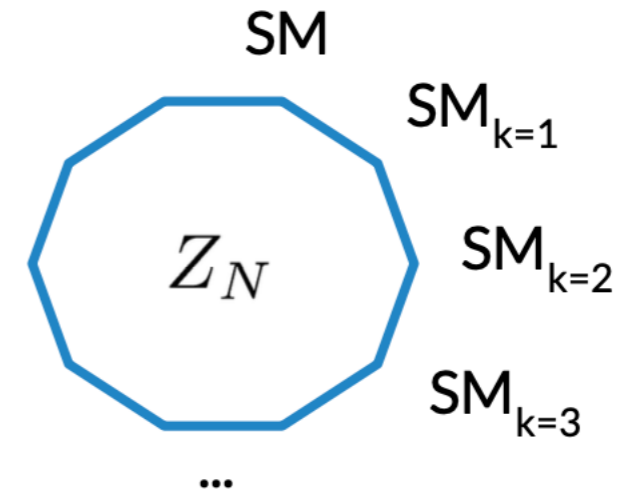
axion potential exponentially suppressed at large N

Suppressing m_a

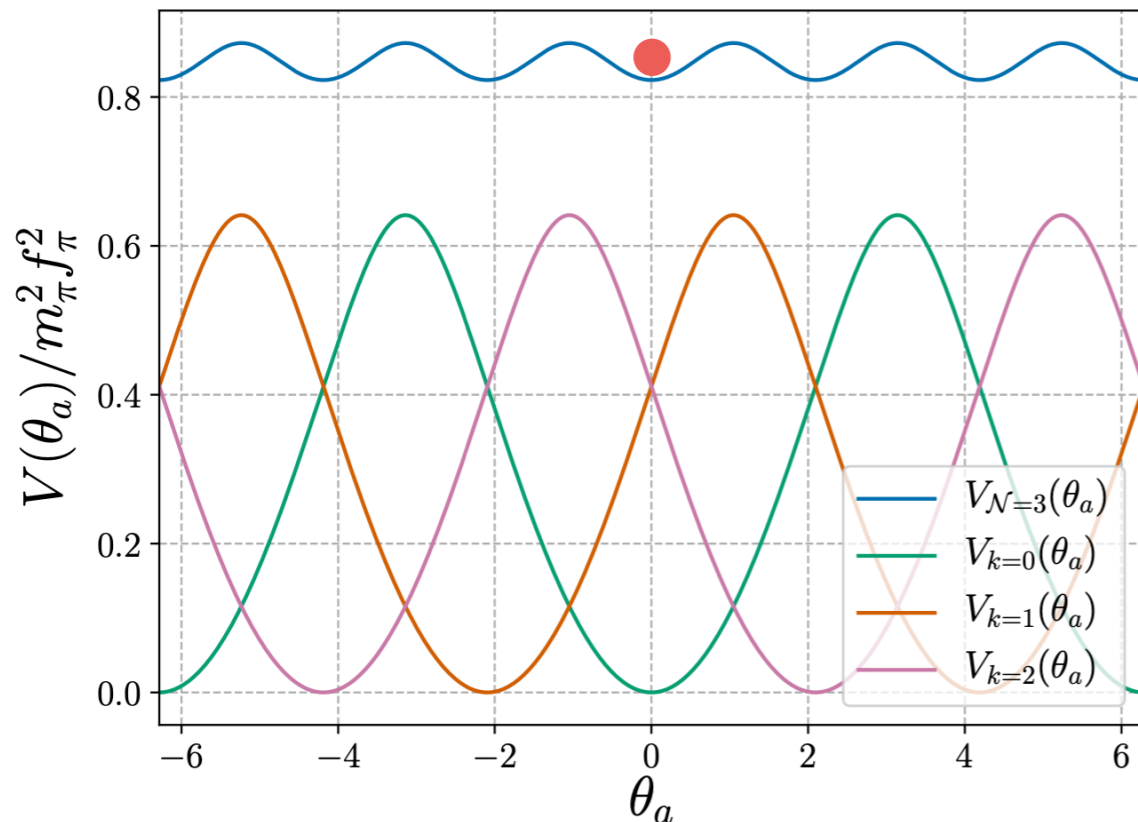
- Z_N axion: N mirror worlds [Hook 1802.10093]

$$SM_k \longrightarrow SM_{k+1 \pmod{N}}$$

$$a \longrightarrow a + \frac{2\pi k}{N} f_a,$$



e.g. Z_3 axion



[LDL, Gavela, Quilez, Ringwald 2102.00012]



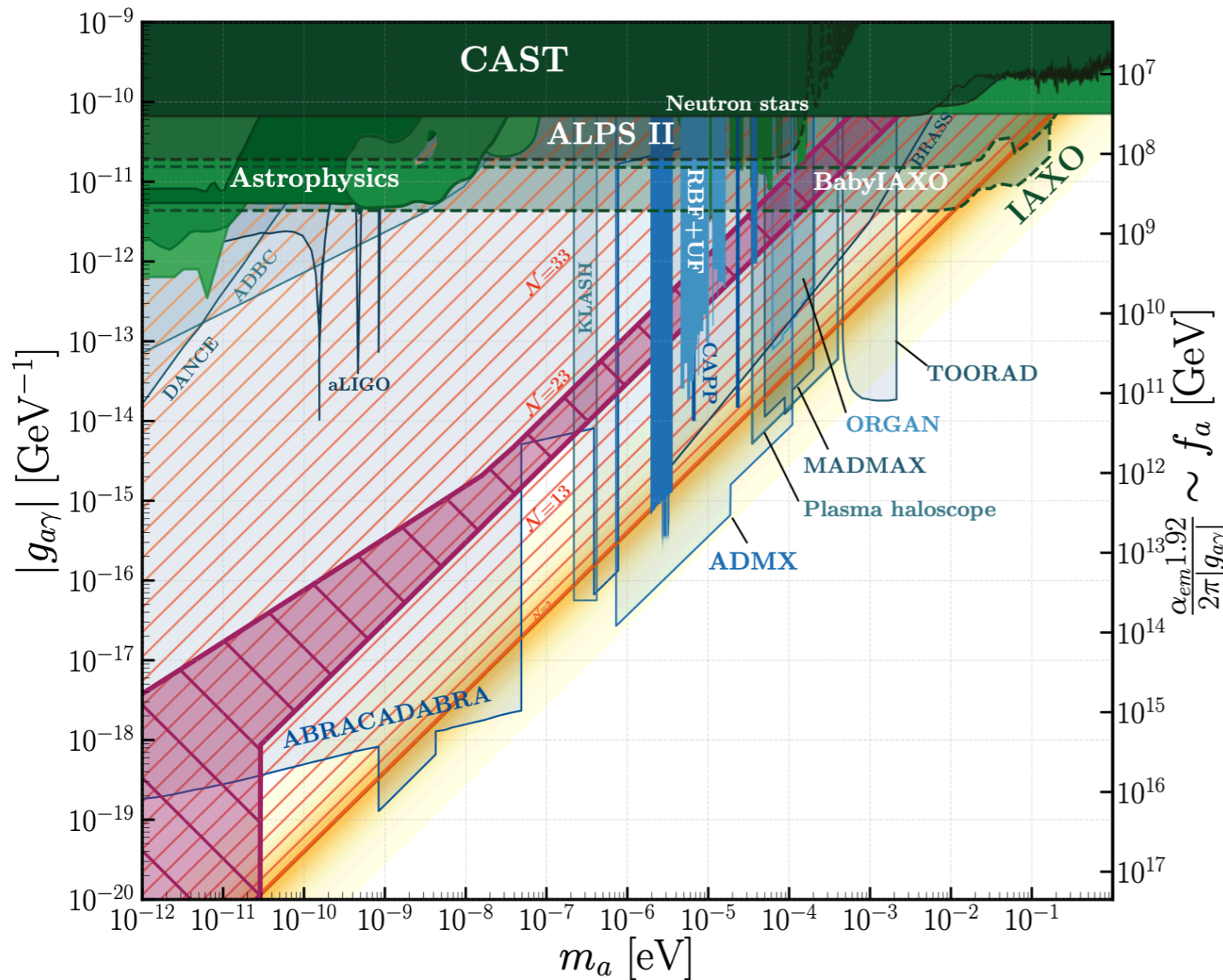
N needs to be odd in order to have a minimum in zero

(strong CP problem is solved with $1/N$ probability)

Suppressing m_a

- Z_N axion: N mirror worlds

[LDL, Gavela, Quilez, Ringwald 2102.00012 + 2102.01082]



$$m_a^2 \simeq \frac{m_\pi^2 f_\pi^2}{f_a^2} \frac{1}{\sqrt{\pi}} \sqrt{\frac{1-z}{1+z}} \mathcal{N}^{3/2} z^{\mathcal{N}}$$

universal enhancement of all axion couplings w.r.t. standard QCD axion

CASPER-Electric could disentangle enhanced coupling vs. suppressed mass mechanism → backup slides

Conclusions

- Take home message

axion properties are UV dependent

1. enhanced/suppressed axion couplings
2. modified $m_a - f_a$ relation
3. flavour violating axions
4. CP-violating axions

if an “axion-like particle” will be ever discovered away from the canonical QCD window, it might still have something to do with strong CP violation

Backup slides

A photo- and electro-philic Axion ?

- Consider a DFSZ-like construction with $2 + n$ Higgs doublets + a SM singlet Φ

$$\mathcal{L}_Y = Y_u \bar{Q}_L u_R H_u + Y_d \bar{Q}_L d_R H_d + Y_e \bar{L}_L e_R H_e$$

$$\frac{E}{N} = \frac{\frac{4}{3}\mathcal{X}(H_u) + \frac{1}{3}\mathcal{X}(H_d) + \mathcal{X}(H_e)}{\frac{1}{2}\mathcal{X}(H_u) + \frac{1}{2}\mathcal{X}(H_d)} \quad g_{ae} = \frac{\mathcal{X}(H_e) m_e}{2N f_a}$$

naively, a large PQ charge for H_e would make the job... but, enhanced global symmetry

$$U(1)^{n+3} \rightarrow U(1)_{\text{PQ}} \times U(1)_Y$$

must be explicitly broken in the scalar potential via non-trivial invariants (e.g. $H_u H_d \Phi^2$)



non-trivial constraints on PQ charges

A photo- and electro-philic Axion ?

- Consider a DFSZ-like construction with $2 + n$ Higgs doublets + a SM singlet Φ

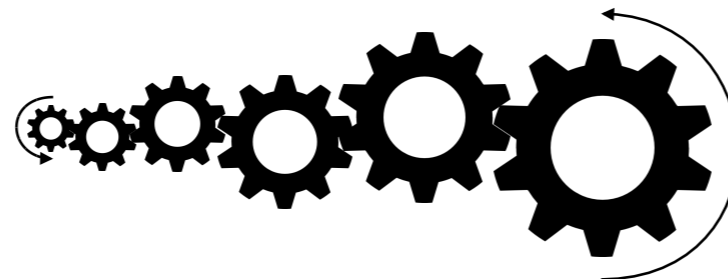
clockwork-like scenarios allow to consistently **boost** E/N [LDL, Mescia, Nardi 1705.05370]

$$\frac{E}{N} = \frac{\frac{4}{3}\mathcal{X}(H_u) + \frac{1}{3}\mathcal{X}(H_d) + \mathcal{X}(H_e)}{\frac{1}{2}\mathcal{X}(H_u) + \frac{1}{2}\mathcal{X}(H_d)} \quad g_{ae} = \frac{\mathcal{X}(H_e)m_e}{2N f_a}$$

$$(H_u H_d \Phi^2)$$

$$(H_k H_{k-1}^*)(H_{k-1}^* H_d^*)$$

$$(H_e H_n)(H_n H_d)$$



[Giudice, McCullough]

$$E/N \sim 2^{n+1}$$

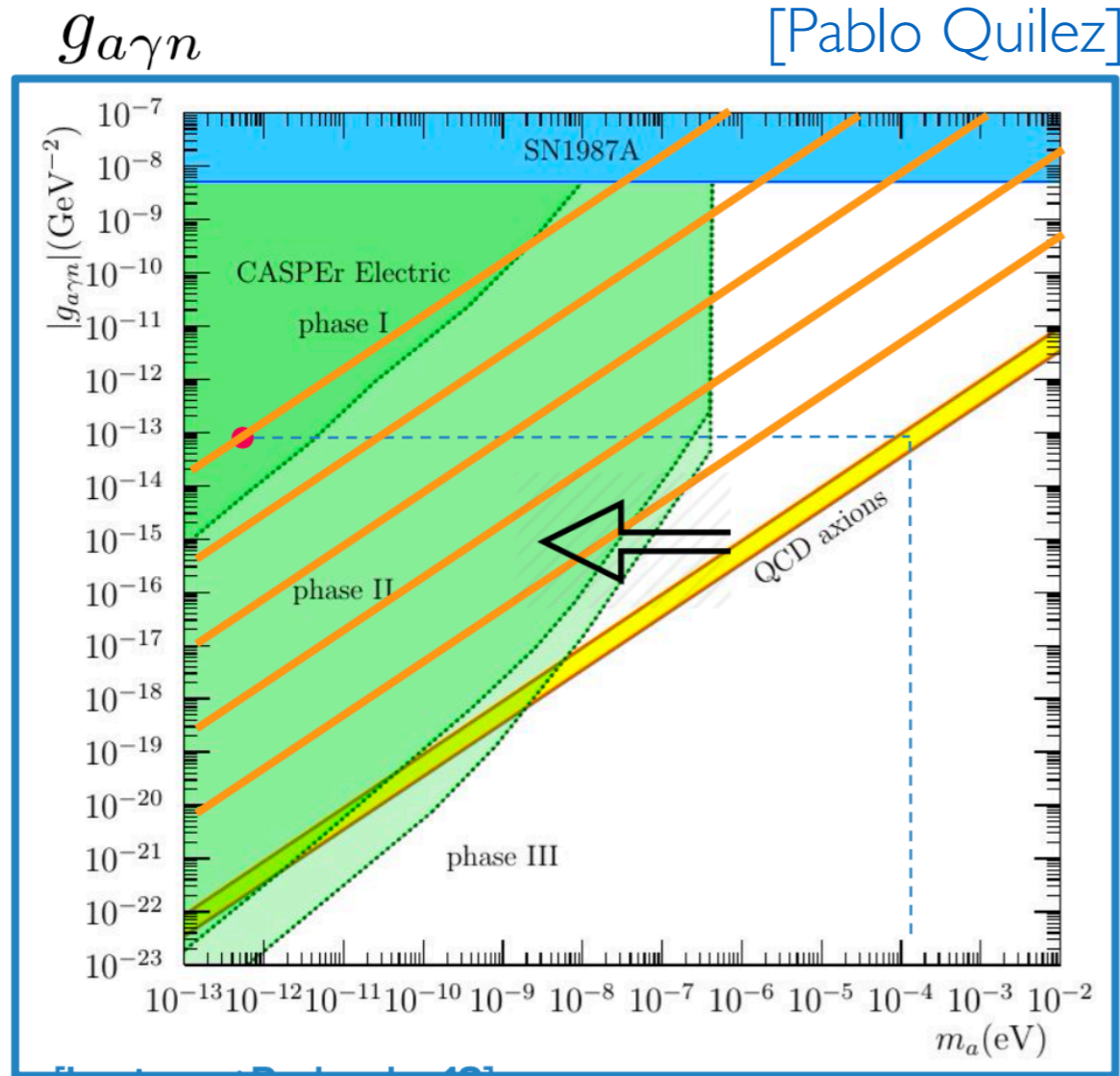
[See also Farina et al. 1611.09855, for KSVZ clockwork]

$$\mathcal{X}(H_e) = 2^{n+1} \left(1 - \frac{v_e^2}{v^2} \right) - \sum_{k=2}^n 2^k \frac{v_k^2}{v^2}$$

How to tell which mechanism ?



CASPEr-Electric could disentangle enhanced coupling vs. suppressed mass



[Irastorza+Redondo, 18]

Based on **2102.00012** and **2102.01082**

$$\mathcal{L} \supset \frac{a}{f_a} \frac{\alpha_s}{8\pi} G\tilde{G} + \dots$$

$$\delta\mathcal{L} \equiv -\frac{i}{2} \frac{0.011 e}{m_n} \frac{a}{f_a} \bar{n} \sigma_{\mu\nu} \gamma_5 n F^{\mu\nu} \equiv g_{a\gamma n}$$
~~$$m_a^2 f_a^2 \simeq m_\pi^2 f_\pi^2 \frac{m_u m_d}{(m_u + m_d)^2}$$~~

Coupling to the nEDM

$$m_a^2 f_a^2 \simeq \frac{m_\pi^2 f_\pi^2}{\sqrt{\pi}} \sqrt{\frac{1-z}{1+z}} \mathcal{N}^{3/2} z^{\mathcal{N}}$$

Axion mass

$m_a (\text{eV})$

CASPER-Electric

- Cosmic Axion Spin Precession Experiment

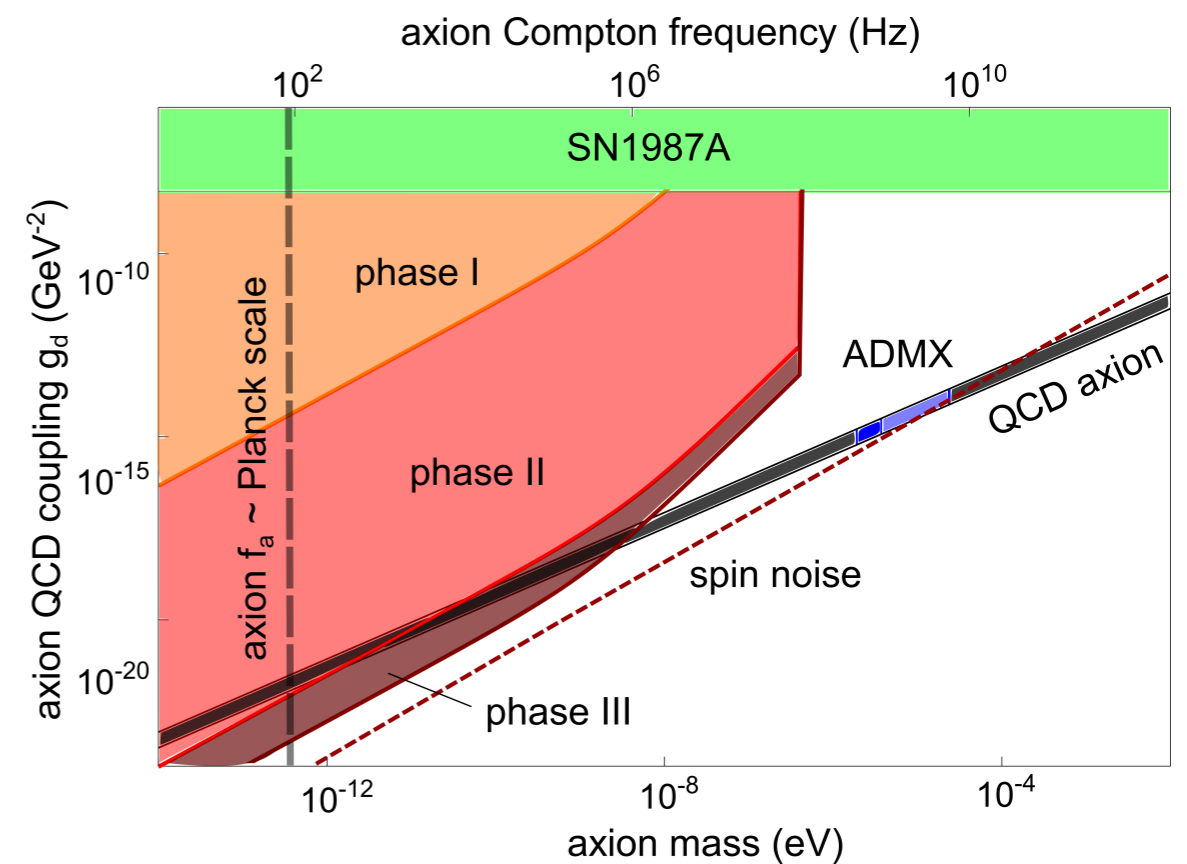
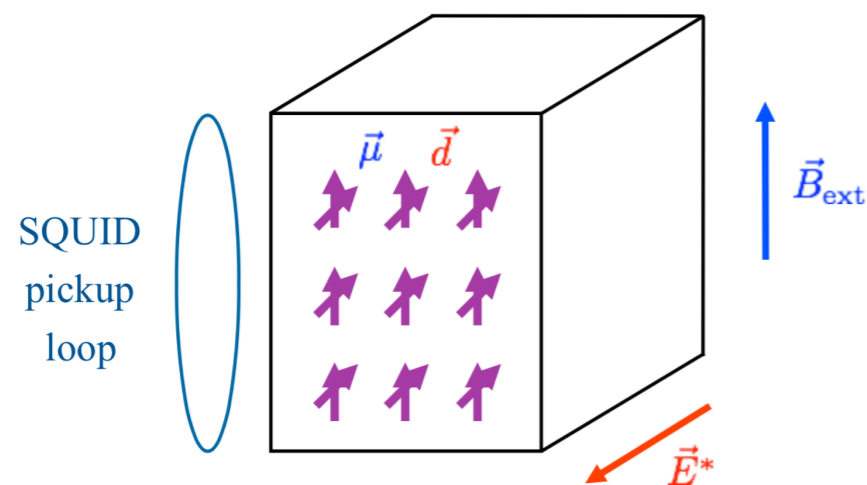
[Graham, Rajendran [1306.6088](#), Budker+ [1306.6089](#), Jackson Kimball+ [1711.08999](#)]

Axion DM field induces an oscillating nEDM

$$\mathcal{L} \supset -\frac{i}{2} g_d a \bar{n} \sigma_{\mu\nu} \gamma_5 n F^{\mu\nu}$$

$$d_n(t) = g_d \frac{\sqrt{2\rho_{\text{DM}}}}{m_a} \cos(m_a t)$$

...which is detected via NMR techniques



CPV axion & long-range forces

- New CP violation in the UV can source a *scalar* axion-nucleon coupling

$$\frac{f_\pi}{2} \frac{a^2}{f_a^2} \bar{N}N \longrightarrow \bar{g}_{aN} a \bar{N}N \quad \bar{g}_{aN} \sim \frac{f_\pi}{f_a} \theta_{\text{eff}} \quad \theta_{\text{eff}} = \frac{\langle a \rangle}{f_a} \neq 0$$

[Moody, Wilczek PRD30 (1984)]

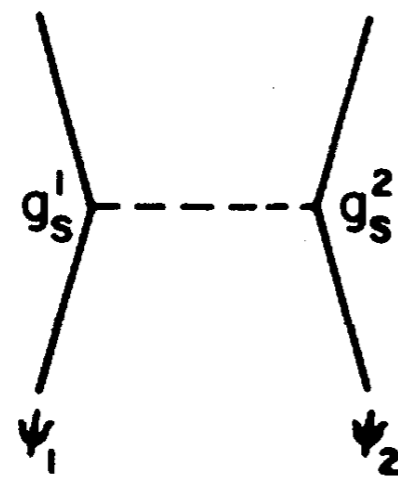
$$\frac{1}{2} \frac{a^2}{f_a^2} \underbrace{\langle G\tilde{G}, G\tilde{G} \rangle}_\chi + \frac{a}{f_a} \underbrace{\langle G\tilde{G}, \mathcal{O}_{\text{CPV}} \rangle}_{\chi'} \longrightarrow \frac{\langle a \rangle}{f_a} = -\frac{\chi'}{\chi}$$

CPV axion & long-range forces

- New CP violation in the UV can source a *scalar* axion-nucleon coupling

$$\frac{f_\pi}{2} \frac{a^2}{f_a^2} \bar{N}N \longrightarrow \bar{g}_{aN} a \bar{N}N \quad \bar{g}_{aN} \sim \frac{f_\pi}{f_a} \theta_{\text{eff}} \quad \theta_{\text{eff}} = \frac{\langle a \rangle}{f_a} \neq 0$$

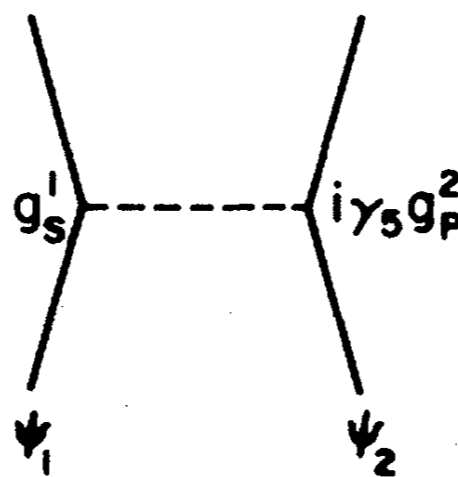
[Moody, Wilczek PRD30 (1984)]



(a)

monopole-monopole

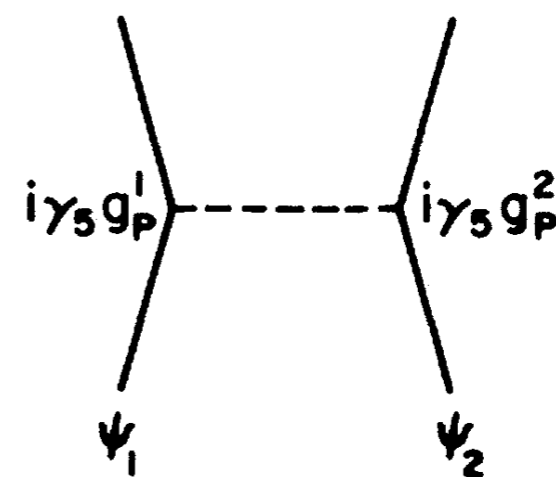
$$V(r) = \frac{-g_s^1 g_s^2 e^{-m_\phi r}}{4\pi r}$$



(b)

monopole-dipole

$$V(r) = (g_s^1 g_P^2) \frac{\hat{\sigma}_2 \cdot \hat{r}}{8\pi M_2} \left[\frac{m_\phi}{r} + \frac{1}{r^2} \right] e^{-m_\phi r}$$



(c)

dipole-dipole

A new master formula

- Moody-Wilczek formula

[Moody, Wilczek PRD30 (1984)]

$$\bar{g}_{aN} = \frac{1}{2} \frac{\bar{\theta}_{\text{eff}}}{f_a} \frac{m_u m_d}{m_u + m_d} \langle N | \bar{u}u + \bar{d}d | N \rangle \simeq \frac{1}{2} \bar{\theta}_{\text{eff}} \left(\frac{17 \text{ MeV}}{f_a} \right)$$

- From bary-meson chiral Lagrangian

[Bertolini, LDL, Nesti [2006.12508](#)]

$$\bar{g}_{an,p} \simeq \frac{4B_0 m_u m_d}{f_a(m_u+m_d)} \left[\pm (b_D + b_F) \frac{\langle \pi^0 \rangle}{F_\pi} + \frac{b_D - 3b_F}{\sqrt{3}} \frac{\langle \eta_8 \rangle}{F_\pi} - \sqrt{\frac{2}{3}} (3b_0 + 2b_D) \frac{\langle \eta_0 \rangle}{F_\pi} - \left(b_0 + (b_D + b_F) \frac{m_{u,d}}{m_d+m_u} \right) \bar{\theta}_{\text{eff}} \right]$$

meson tadpoles

iso-spin breaking

MW missed a factor 1/2

An application: Left-Right

- Low-scale (PQ)Left-Right with P-parity

[Bertolini, LDL, Nesti [2006.12508](#)
PRL 126 (2021)]

4-quark op. from W_R exchange

chiral representation \rightarrow

$$\mathcal{O}_1^{ud} = (\bar{u}u)(\bar{d}i\gamma_5 d) \quad \rightarrow \quad c_3(U_{11}^\dagger U_{22} - U_{11}U_{22}^\dagger)$$

$$U = \exp \left[\frac{2i}{\sqrt{6}F_0} \eta_0 I + \frac{2i}{F_\pi} \Pi \right]$$

$$\frac{\langle \pi^0 \rangle}{F_\pi} \simeq \frac{G_F}{\sqrt{2}} \mathcal{C}_1^{[ud]} \frac{c_3}{B_0 F_\pi^2} \frac{m_u + m_d + 4m_s}{m_u m_d + m_d m_s + m_s m_u}$$

$$\frac{\langle \eta_8 \rangle}{F_\pi} \simeq \frac{G_F}{\sqrt{2}} \mathcal{C}_1^{[ud]} \frac{\sqrt{3}c_3}{B_0 F_\pi^2} \frac{m_d - m_u}{m_u m_d + m_d m_s + m_s m_u}$$

$$\theta_{\text{eff}} \simeq \frac{G_F}{\sqrt{2}} \mathcal{C}_1^{[ud]} \frac{2c_3}{B_0 F_\pi^2} \frac{m_d - m_u}{m_u m_d}$$

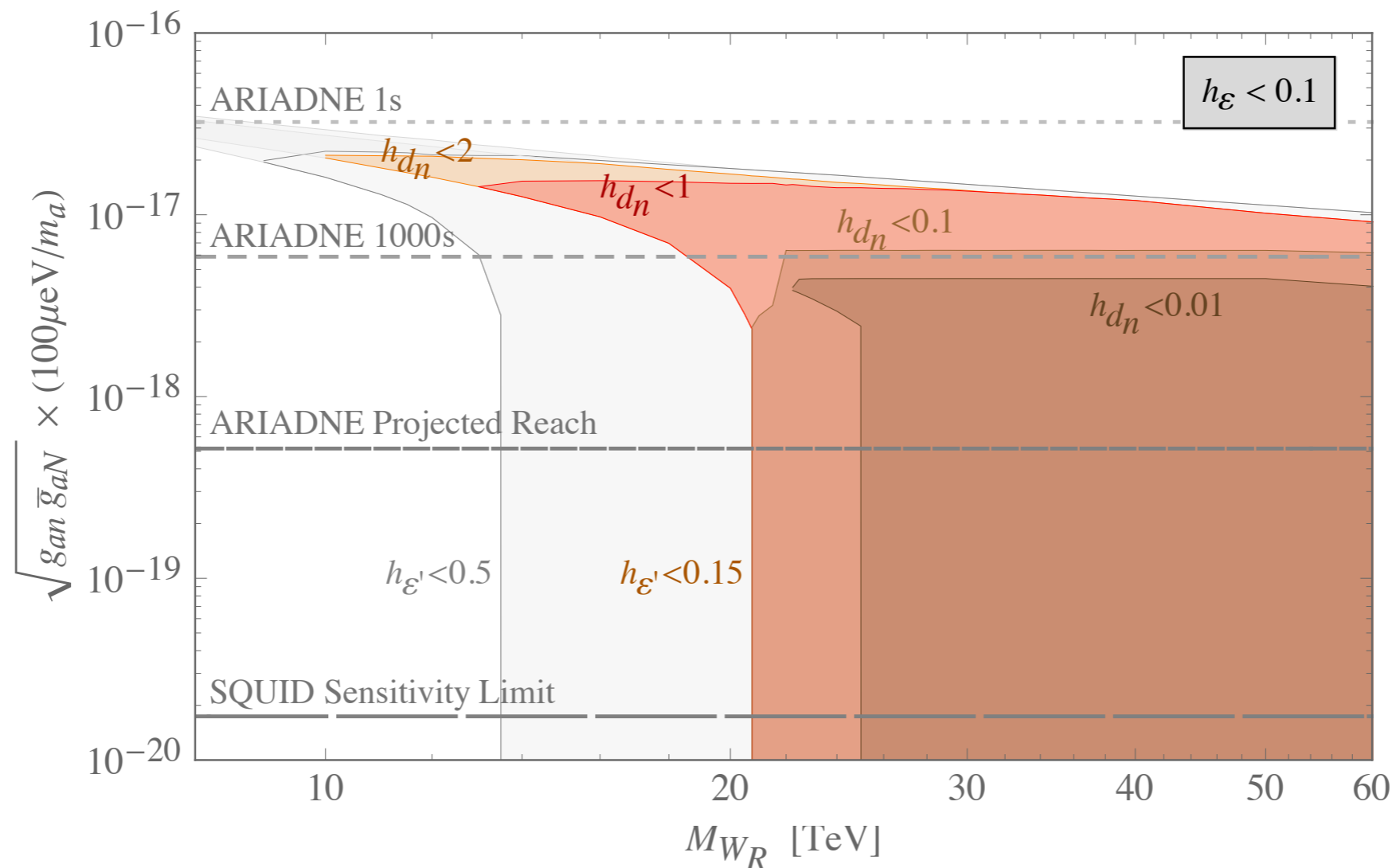
An application: Left-Right

- Low-scale (PQ)Left-Right with P-parity

[Bertolini, LDL, Nesti [2006.12508](#)
PRL 126 (2021)]

4 CPV observables (ε , ε' , d_n , \bar{g}_{aN}) function of a single phase α

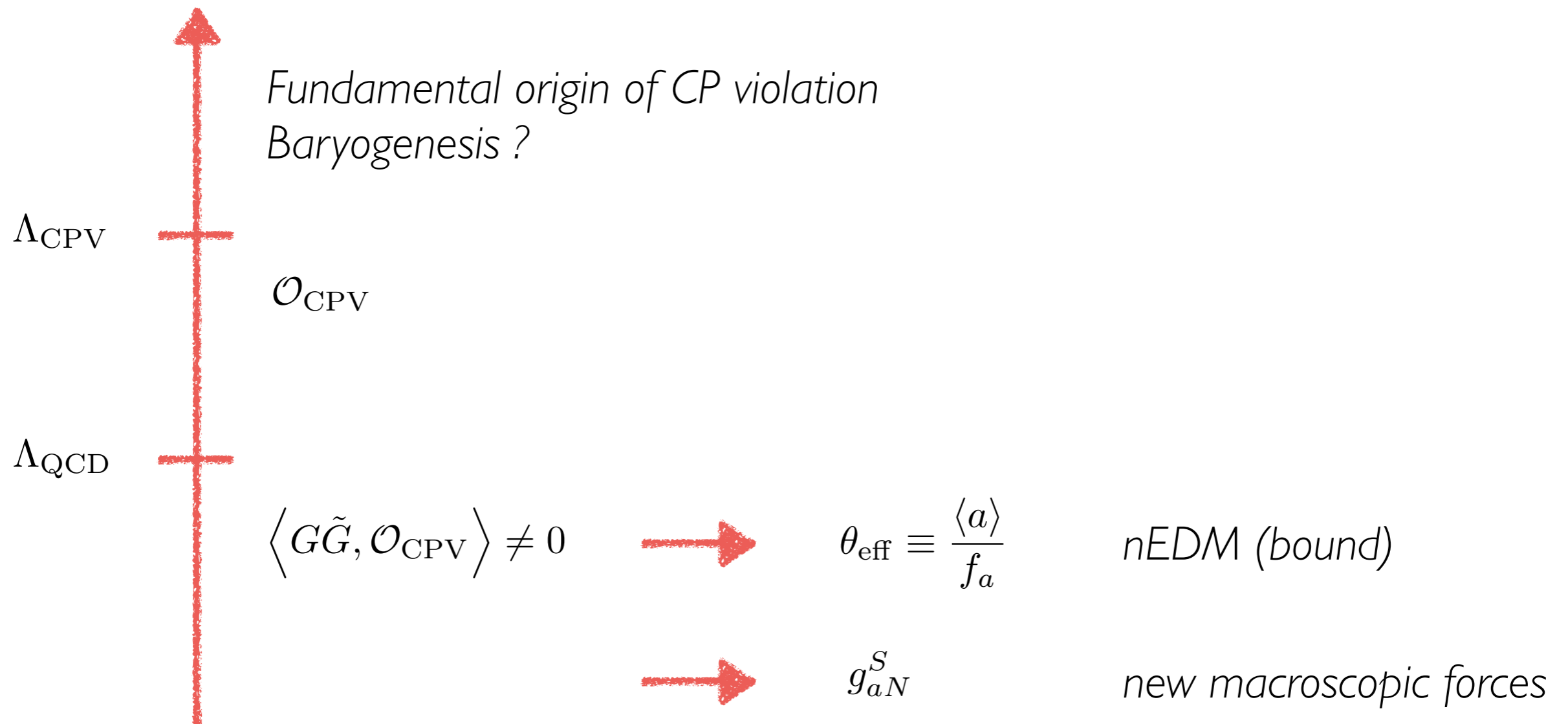
$$\langle \Phi \rangle = \text{diag} \{v_1, e^{i\alpha} v_2\}$$



$$h_{\mathcal{O}} \equiv \frac{\mathcal{O}^{\text{th}}}{\mathcal{O}^{\text{exp}}}$$

CP-violating axions

- Rethinking the axion as a portal to UV sources of CP-violation



strong CP problem or strong CP opportunity ?