

The RBC/UKQCD $g-2$ program

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February 15, 2023 – Pisa, NePSi 23

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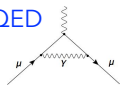
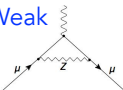
[Stony Brook University](#)

Jun-Sik Yoo

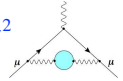
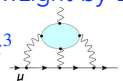
Sergey Syritsyn (RBRC)

Contributions from known particles: The Standard Model

$$a_{\mu}(\text{SM}) = a_{\mu}(\text{QED}) + a_{\mu}(\text{Weak}) + a_{\mu}(\text{Hadronic})$$

<p>QED</p>  <p>+ ...</p>	$116\,584\,718.9(1) \times 10^{-11}$	0.001 ppm
<p>Weak</p>  <p>+ ...</p>	$153.6(1.0) \times 10^{-11}$	0.01 ppm

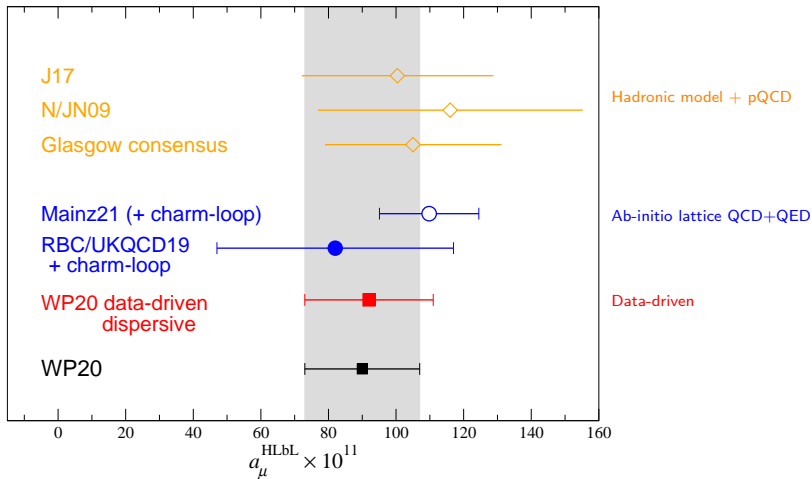
Hadronic...

<p>...Vacuum Polarization (HVP)</p> <p>α^2</p>  <p>+ ...</p>	$6845(40) \times 10^{-11}$ [0.6%]	0.37 ppm
<p>...Light-by-Light (HLbL)</p> <p>α^3</p>  <p>+ ...</p>	$92(18) \times 10^{-11}$ [20%]	0.15 ppm

Numbers from Theory Initiative Whitepaper

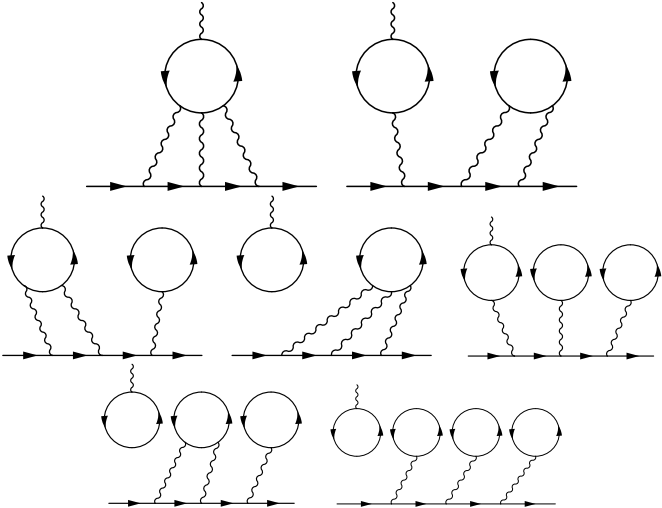
Uncertainty dominated by hadronic contributions

Status of hadronic light-by-light contribution



Systematically improvable methods are maturing; uncertainty to a_{μ} controlled at 0.15ppm; **cross-checks detailed in Theory Initiative whitepaper**

Diagrams to calculate



The first complete ab-initio LQCD result for HLbL

Hadronic Light-by-Light Scattering Contribution to the Muon Anomalous Magnetic Moment from Lattice QCD

Thomas Blum,^{1,2} Norman Christ,³ Masashi Hayakawa,^{4,5} Taku Izubuchi,^{6,2}
Luchang Jin^{1,2,*}, Chulwoo Jung,⁶ and Christoph Lehner^{7,6}

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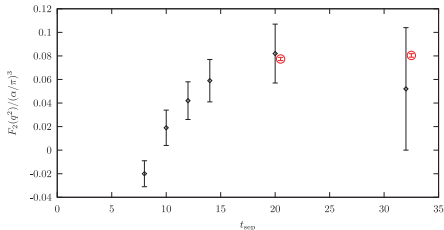
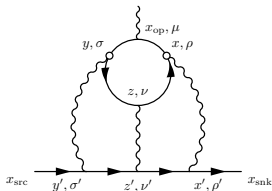
(Received 18 December 2019; accepted 27 February 2020; published 1 April 2020)

We report the first result for the hadronic light-by-light scattering contribution to the muon anomalous magnetic moment with all errors systematically controlled. Several ensembles using $2 + 1$ flavors of physical mass Möbius domain-wall fermions, generated by the RBC and UKQCD collaborations, are employed to take the continuum and infinite volume limits of finite volume lattice QED + QCD. We find $a_{\mu}^{\text{HLbL}} = 7.87(3.06)_{\text{stat}}(1.77)_{\text{sys}} \times 10^{-10}$. Our value is consistent with previous model results and leaves little room for this notoriously difficult hadronic contribution to explain the difference between the standard model and the BNL experiment.

It took many years and improvements to get there . . .

PRD93(2015)014503:

We introduce a new sampling strategy with 10x reduced noise for same cost (red versus black):

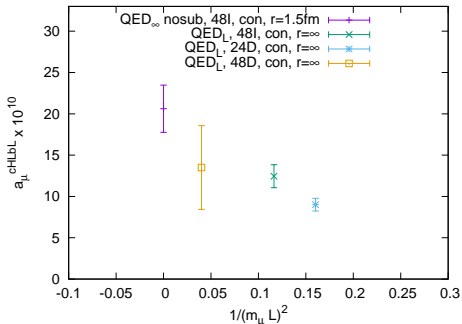
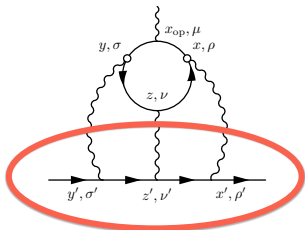


Stochastically evaluate the sum over vertices x and y :

- ▶ Pick random point x on lattice
- ▶ Sample all points y up to a specific distance $r = |x - y|$, see vertical red line
- ▶ Pick y following a distribution $P(|x - y|)$ that is peaked at short distances

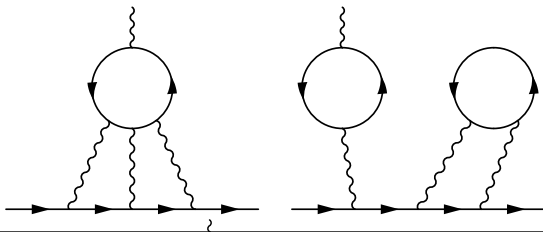
PRD96(2017)034515:

We remove power-law like finite-volume errors by computing the muon-photon part of the diagram in infinite volume (similar idea proposed by Mainz group)



Infinite-volume kernel projects to larger four-point function \Rightarrow large noise at physical pion mass

Paper in preparation using explicit pion-pole cancellation



$$\lim_{R \rightarrow \infty} \frac{a_{\mu}^{\text{discon}}(R_{\text{max}} > R)}{a_{\mu}^{\text{con}}(R_{\text{max}} > R)} = -\frac{1}{2} \cdot \frac{(e_u^2 + e_d^2)^2}{e_u^4 + e_d^4} = -\frac{25}{34}.$$

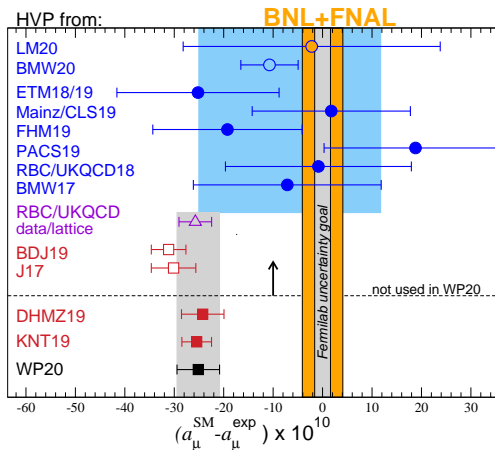
$$a_{\mu}^{\text{no-pion}} = a_{\mu}^{\text{discon}} + \frac{25}{34} a_{\mu}^{\text{con}},$$

$$a_{\mu}^{\text{discon}} = a_{\mu}^{\text{no-pion}} - \frac{25}{34} a_{\mu}^{\text{con}},$$

$$a_{\mu}^{\text{total}} = a_{\mu}^{\text{no-pion}} + \frac{9}{34} a_{\mu}^{\text{con}}.$$

Long-distance pion-pole cancels in $a_{\mu}^{\text{no-pion}}$ such that it can be evaluated at shorter distances. Only less noisy connected diagram remains.

Status and impact of hadronic vacuum polarization contribution



Ab-initio lattice QCD(+QED) calculations are maturing

Difficult problem: scales from $2m_{\pi}$ to several GeV enter; cross-checks needed at high precision

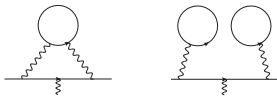
Hybrid window method restricts scales that enter from lattice/dispersive data

Dispersive, $e^{+}e^{-} \rightarrow \text{hadrons}$ (20+ years of experiments)

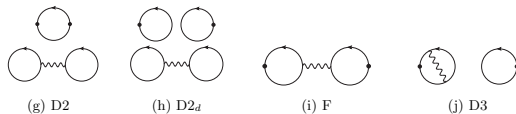
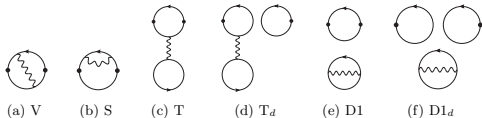
Now first published lattice result with sub-percent precision available (BMW20), cross-checks are crucial to establish or refute high-precision lattice methodology (same situation as for HLbL)

Diagrams to calculate

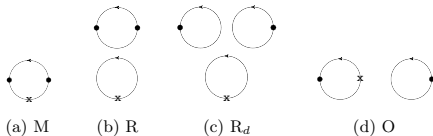
Isospin
limit



QED
corrections



Strong
isospin
breaking



Calculation of the Hadronic Vacuum Polarization Contribution to the Muon Anomalous Magnetic Moment

T. Blum,¹ P. A. Boyle,² V. Gülpers,³ T. Izubuchi,^{4,5} L. Jin,^{1,5} C. Jung,⁴ A. Jüttner,³ C. Lehner,^{4,*} A. Portelli,² and J. T. Tsang²

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(Received 25 January 2018; published 12 July 2018)

We present a first-principles lattice QCD + QED calculation at physical pion mass of the leading-order hadronic vacuum polarization contribution to the muon anomalous magnetic moment. The total contribution of up, down, strange, and charm quarks including QED and strong isospin breaking effects is $a_{\mu}^{\text{HVP LO}} = 715.4(18.7) \times 10^{-10}$. By supplementing lattice data for very short and long distances with R -ratio data, we significantly improve the precision to $a_{\mu}^{\text{HVP LO}} = 692.5(2.7) \times 10^{-10}$. This is the currently most precise determination of $a_{\mu}^{\text{HVP LO}}$.

Pure lattice result and dispersive result with reduced $\pi\pi$ dependence (window method)

Lattice QCD – Time-Moment Representation

Starting from the vector current $J_\mu(x) = i \sum_f Q_f \bar{\Psi}_f(x) \gamma_\mu \Psi_f(x)$ we may write

$$a_\mu^{\text{HVP LO}} = \sum_{t=0}^{\infty} w_t C(t)$$

with

$$C(t) = \frac{1}{3} \sum_{\vec{x}} \sum_{j=0,1,2} \langle J_j(\vec{x}, t) J_j(0) \rangle$$

and w_t capturing the photon and muon part of the HVP diagrams ([Bernecker-Meyer 2011](#)).

The correlator $C(t)$ is computed in lattice **QCD+QED** at **physical pion mass** with **non-degenerate** up and down quark masses including up, down, strange, and charm quark contributions. The missing bottom quark contributions are computed in pQCD.

Window method (introduced in RBC/UKQCD 2018)

We also consider a window method. Following Meyer-Bernecker 2011 and smearing over t to define the continuum limit we write

$$a_\mu = a_\mu^{\text{SD}} + a_\mu^{\text{W}} + a_\mu^{\text{LD}}$$

with

$$a_\mu^{\text{SD}} = \sum_t C(t) w_t [1 - \Theta(t, t_0, \Delta)],$$

$$a_\mu^{\text{W}} = \sum_t C(t) w_t [\Theta(t, t_0, \Delta) - \Theta(t, t_1, \Delta)],$$

$$a_\mu^{\text{LD}} = \sum_t C(t) w_t \Theta(t, t_1, \Delta),$$

$$\Theta(t, t', \Delta) = [1 + \tanh [(t - t')/\Delta]] / 2.$$

All contributions are well-defined individually and can be computed from lattice or R-ratio via $C(t) = \frac{1}{12\pi^2} \int_0^\infty d(\sqrt{s}) R(s) s e^{-\sqrt{s}t}$ with $R(s) = \frac{3s}{4\pi\alpha^2} \sigma(s, e^+ e^- \rightarrow \text{had})$.

a_μ^{W} has small statistical and systematic errors on lattice!

An update of Euclidean windows of the hadronic vacuum polarization

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 C. Lehner,^{6,13} N. Matsumoto,⁷ R. D. Mawhinney,⁹ A. S. Meyer,^{13,14} and J. T. Tsang^{10,15}
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¹⁴*Lawrence Berkeley National Laboratory, Berkeley, CA 94720, USA*

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University of Southern Denmark, Campusvej 55, 5230 Odense M, Denmark

(Dated: February 14, 2023)

We compute the standard Euclidean window of the hadronic vacuum polarization using multiple independent blinded analyses. We improve the continuum and infinite-volume extrapolations of the dominant quark-connected light-quark isospin-symmetric contribution and address additional sub-leading systematic effects from sea-charm quarks and residual chiral-symmetry breaking from first principles. We find $a_\mu^W = 235.56(65)(50) \times 10^{-10}$, which is in 3.8σ tension with the recently published dispersive result of $a_\mu^W = 229.4(1.4) \times 10^{-10}$ [1] and in agreement with other recent lattice determinations. We also provide a result for the standard short-distance window. The results reported here are unchanged compared to our presentation at the Edinburgh workshop of the g-2 Theory Initiative in 2022 [2].

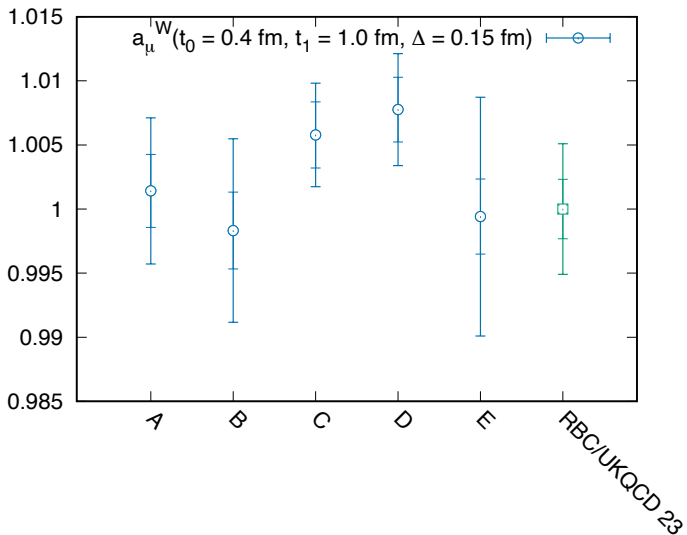
Blinding

- ▶ 2 analysis groups for ensemble parameters (not blinded)
- ▶ 5 analysis groups for vector-vector correlators (blinded, to avoid bias towards other lattice/R-ratio results)
- ▶ Blinded vector correlator $C_b(t)$ relates to true correlator $C_0(t)$ by

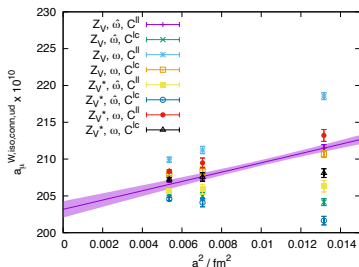
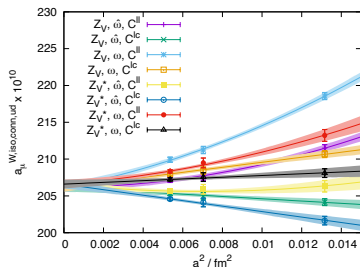
$$C_b(t) = (b_0 + b_1 a^2 + b_2 a^4) C_0(t) \quad (1)$$

with appropriate random b_0, b_1, b_2 , different for each analysis group. This prevents complete unblinding based on previously shared data on coarser ensembles.

Relative unblinding



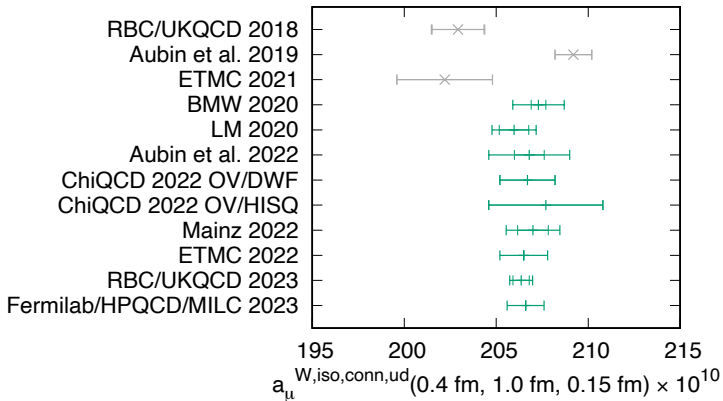
Full unblinding on August 31st, 2022



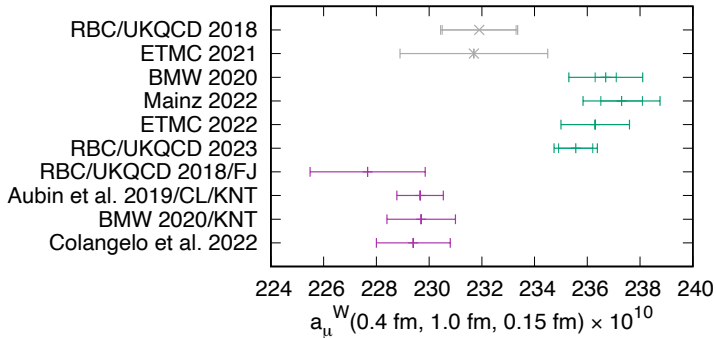
24 data points enter new continuum extrapolation, in 2018 had only 2 data points

Already generating data at two finer lattice spacings ($a^2 / \text{fm}^2 = 0.0017, 0.003$)

Isospin symmetric standard window in context



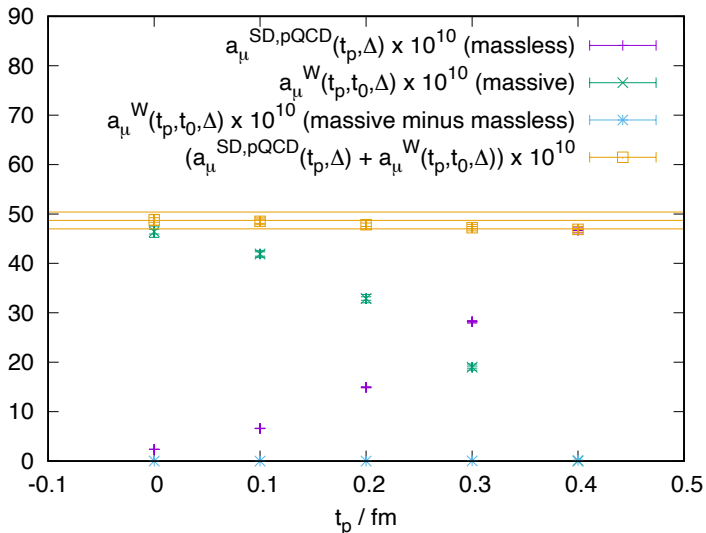
Adding back the QED+SIB parts calculated in RBC/UKQCD2018



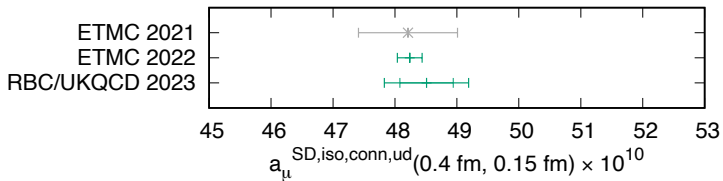
3.8 σ tension between lattice and $e^+e^- \rightarrow \text{hadrons}$

Short-distance window also computed

Stability tested against α_s^4 massless QCD calculation Chetyrkin & Maier, 2011:



Short-distance window in context



Summary:

- ▶ In isospin symmetric limit SD and standard windows are converging within lattice QCD.
- ▶ Tension for standard window with dispersive approach.
- ▶ Working on LD update. Possibly available this summer.
- ▶ Working on update to QED+SIB as well, possibly available this fall.
- ▶ Also: τ decays; on-going project for needed QED+SIB corrections
- ▶ Our final goal: match or surpass final FNAL experimental precision