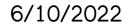




# TO THE THERMAL LIMIT AND BELOW

#### G. Dho, E. Baracchini, S. Piacentini, D. Marques

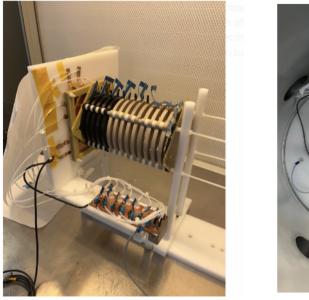




#### Previously

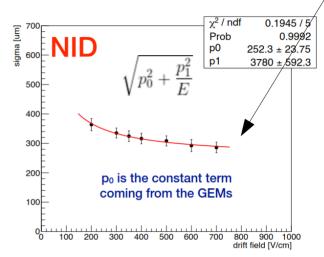
G.Dho

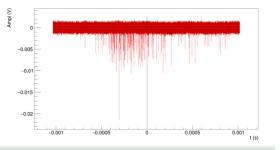
• MANGO was put in the keg to measure NID diffusion with more drift length to improve this

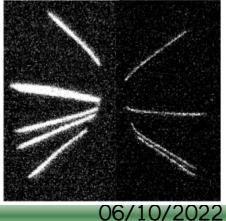




Pressure was reduced at 650 mbar for more stable operations

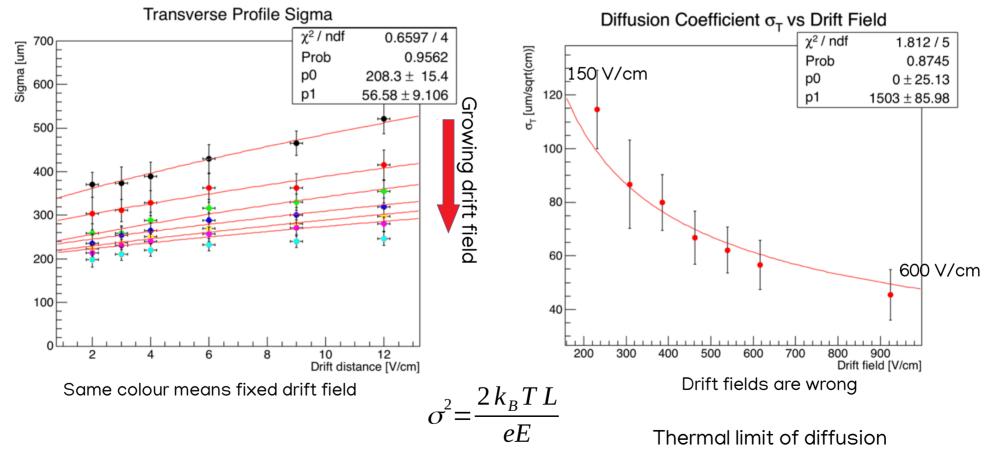






#### Previously

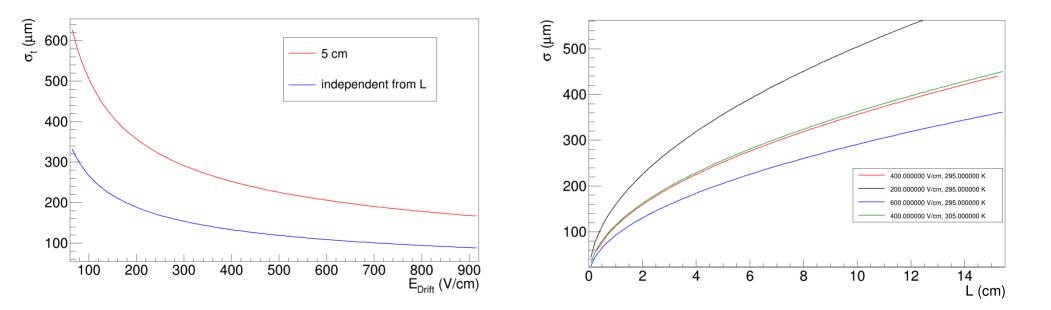
• This new configuration latest measurement presented on 16/6/2022



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#### Too Good to Be True?

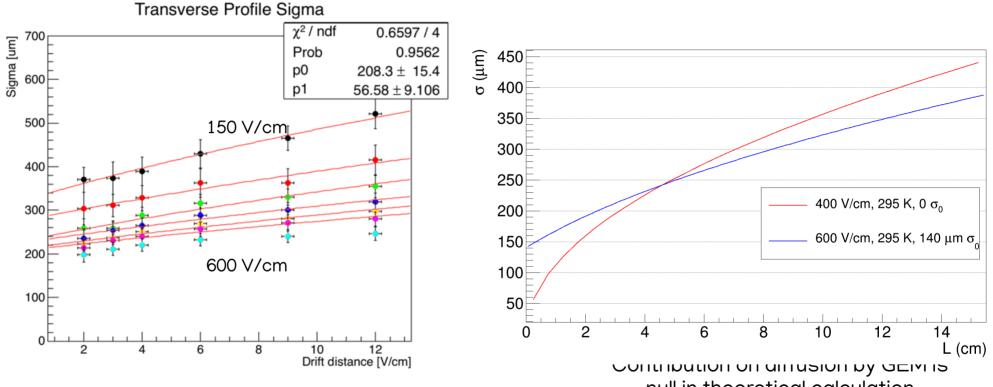
• Our data suggest a diffusion coefficient smaller than what thermally possible



Thermal limit of diffusion

### Too Good to Be True?

• Our data suggest a diffusion coefficient smaller than what thermally possible



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null in theoretical calculation



### MATHEMATICAL PROVE: MOBILITY

• Starting from the Blum-Rolandi book calculation the w drift velocity is

$$w = v + c_d = \frac{eE}{m}\tau + c_d$$

- E electric field
- *m* mass of the target gas
- $\tau$  average time between two collisions
- $c_d$  velocity in the direction of the drift as a result of the collision with molecules
- The authors evaluate  $c_d$  and get the mobility

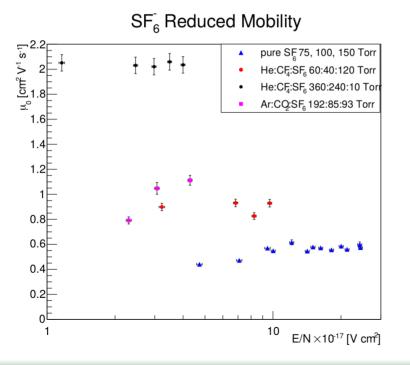
$$\mu = e\tau \left(\frac{1}{m} + \frac{1}{M}\right) \quad \longrightarrow \quad \mu = \sqrt{\frac{1}{m} + \frac{1}{M}} \sqrt{\frac{1}{3k_B T}} \frac{e}{N\sigma}$$

M mass of the negative ion drifting



## MATHEMATICAL PROVE: MOBILITY

 This allows to correctly predict the relative difference of the mobilities for the gas mixtures of the NITEC paper, by averaging the gas components weighted on the average time between two collision of each molecule



$$\mu = \sqrt{\frac{1}{m} + \frac{1}{M}} \sqrt{\frac{1}{3k_B T}} \frac{e}{N\sigma}$$

# When the gas is He dominated the NID are faster

06/10



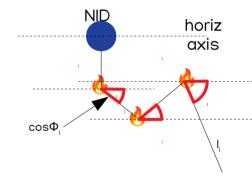
• For the diffusion the Blum-Rolandi writes the gaussian diffusion

$$\sigma_{tot}^2 = 2Dt = 2\frac{D}{w}L = 2\frac{D}{\mu}\frac{L}{E} = \sigma_{diff}^2L$$

• But it can evaluate with the average quadratic dispersion

$$\sum_{i}^{n} \int (l_{i} \cos \phi_{i})^{2} \frac{1}{\lambda} e^{-\frac{l_{i}}{\lambda}} f(\cos \phi_{i}) d\lambda_{i} d\cos \phi_{i}$$

- $l_i$  distance l traveled in the between the ith and ith-1 collision
- $\cos \phi_i$  angle between the ion direction and the chosen direction after the ith collision
- $\lambda$  mean free path
- $f(\cos \varphi)$  probability distribution of the  $\cos \varphi$  due to ion to ion collision
- n number of collision after a time t (large number of collisions)



 $\begin{array}{c} \text{Cos} \Phi_{_i} \text{ is not the} \\ \text{scattering angle of the} \\ \text{ion drifting} \end{array}$ 

06/10

 Solving the integral, defining y the ratio of the masses M/m and F(y) the angular part of the integral

$$\sigma_{tot}^2 = \sum_{i}^{n} 2\lambda^2 \int \cos^2 \phi_i f(\cos \phi_i) d\cos \phi_i = 2\frac{2\epsilon}{m}\tau F(y)t$$
$$\sigma_{tot}^2 = \frac{6k_B T \tau}{m\mu} F(y) \frac{L}{E}$$

• Inserting the mobility calculated before and dividing by L

$$\sigma_{diff}^2 = \frac{6k_BT}{e} \cdot \frac{F(y)}{\left(1 + \frac{1}{y}\right)} \frac{1}{E}$$



$$\sigma_{diff}^2 = \frac{6k_BT}{e} \cdot \frac{F(y)}{\left(1 + \frac{1}{y}\right)} \frac{1}{E}$$

If an isotropic distribution is chosen

• F(y) = 2/3

If also y=1 (ions on ions)

• The thermal limit formula is obtained



06/10/20

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$$\sigma_{diff}^2 = \frac{6k_BT}{e} \cdot \frac{F(y)}{\left(1 + \frac{1}{y}\right)} \frac{1}{E}$$

If an isotropic distribution is chosen

• F(y) = 2/3

If also y=1 (ions on ions)

• The thermal limit formula is obtained

If an isotropic distribution is chosen

• F(y) = 2/3

#### Depending on the target y varies

 y<1 the diffusion goes below thermal limit

• Averaging on the different components of the gas (k, percentage of gas in mixture)

$$\bar{\sigma}_{diff}^2 = \frac{6k_BT}{eE} \frac{\sum_{i}^{gas} \delta_i F_i(y_i) \frac{\rho_i}{A_i} k_i (R_i + R_{SF_6})^2}{\sum_{i}^{gas} \frac{\rho_i}{A_i} k_i (R_i + R_{SF_6})^2} \xrightarrow{\text{$k$, percentage of gas in mixture}}{\delta_i \text{ average momentum loss}} = y/(y+1)}$$

• Calculating the sigma with a field of 600 V/cm and using F(y)=2/3, I get

 $92 \frac{\mu m}{\sqrt{cm}}$ 

**Thermal limit** 

$$65 \frac{\mu m}{\sqrt{cm}}$$

# $47 \frac{\mu m}{\sqrt{cm}}$

#### Calculated value

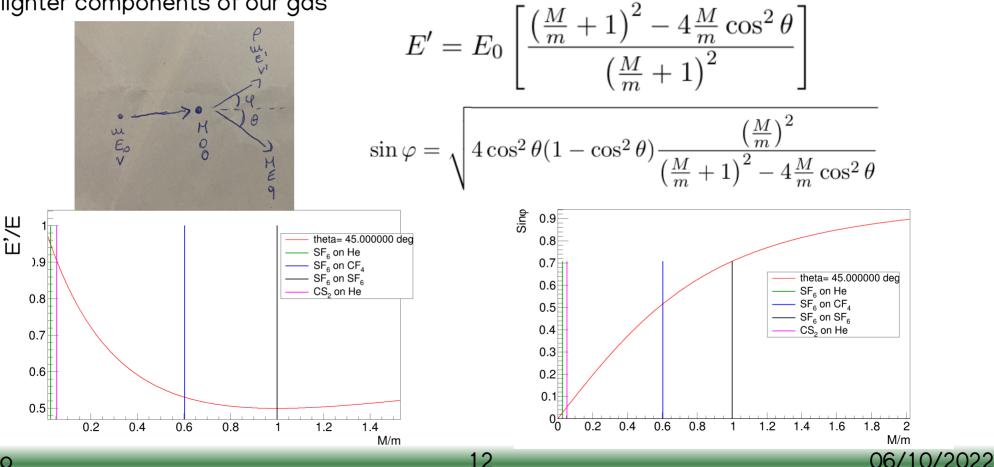
This is only an upper limit as the distribution of the angle was considered isotropic

#### Measured value

Error under evaluation... Could be consistent with the calculated one

#### **KINEMATICALLY**

• The concept can be understood looking at the kinematics of the very heavy SF, on the lighter components of our gas



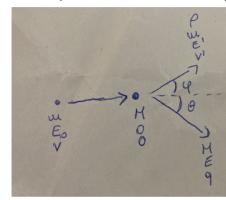
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#### KINEMATICALLY

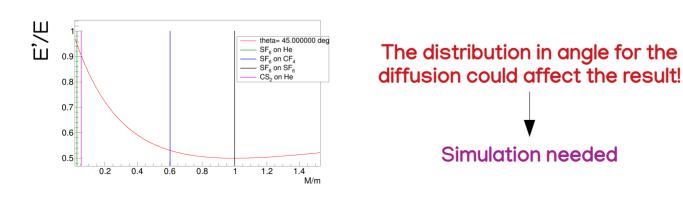
G.Dho

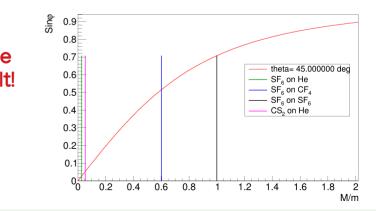
• The concept can be understood looking at the kinematics of the very heavy SF<sub>6</sub> on the lighter components of our gas  $\Gamma(M) = 2$ 

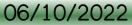
si



$$E' = E_0 \left[ \frac{\left(\frac{M}{m} + 1\right)^2 - 4\frac{M}{m}\cos^2\theta}{\left(\frac{M}{m} + 1\right)^2} \right]$$
$$\ln\varphi = \sqrt{4\cos^2\theta(1 - \cos^2\theta)\frac{\left(\frac{M}{m}\right)^2}{\left(\frac{M}{m} + 1\right)^2 - 4\frac{M}{m}\cos^2\theta}}$$







#### CONCLUSION

G.Dho

- The data on NID from the MANGO in a KEG campaign are being analyzed
- The measured diffusion seems to go below the best expectations of the thermal limt
- From prime principles it is possible to see that the large mass of SF<sub>6</sub> on the light He may be the reason for this behaviuor (Supergaseous helium-helped Anion drift??, Cold Anion Drift??...)
- A simulation of the collisions may be needed to obtain the actual values measured
- Some hard NID data taken with 50:50  $\rm He: CF_4$  and 70:30 planned, to see a difference in the diffusion coefficient changing the helium content

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Credit: Samuele Torelli

