

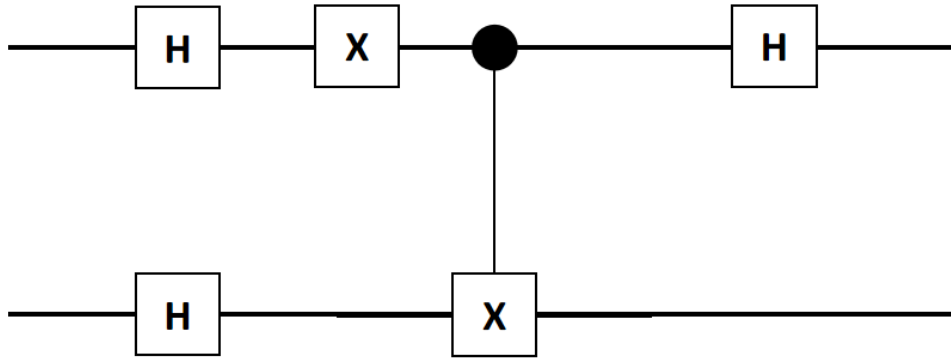
# Noisy Gates for Quantum Computing

Quantum Computing @INFN

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# Real Quantum Computers



In an ideal (= isolated) world, quantum computers run beautifully

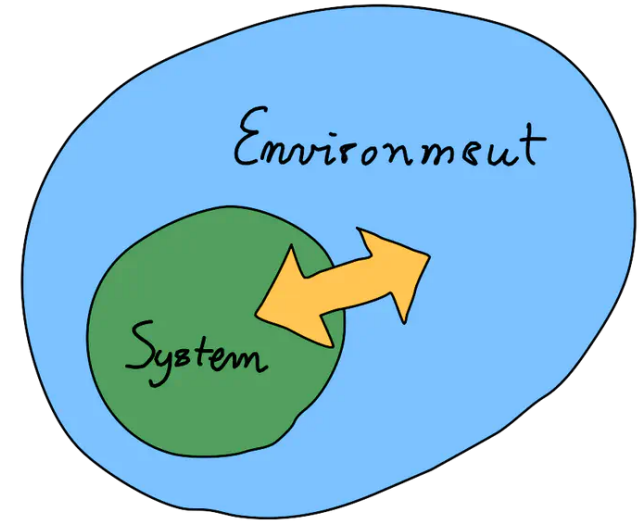
In real life, they are subject to **noise**

- Quantum Error Correction, but even more qubits are needed
- NISQ (Noise Intermediate-Scale Quantum) devices

# Study the noise

A **proper theoretical modelling** of the effect of the environment on a quantum systems allows to:

- Have a **physical understanding** of the sources of noise
- Suggest strategies to **mitigate errors**
- Perform **accurate simulations** to predict how the performances scale with the number of qubits/gates.



Georgopoulos, K., Emary, C., & Zuliani, P. (2021). Modeling and simulating the noisy behavior of near-term quantum computers. *Physical Review A*, 104(6), 062432.

Sun, J., Yuan, X., Tsunoda, T., Vedral, V., Benjamin, S. C., & Endo, S. (2021). Mitigating realistic noise in practical noisy intermediate-scale quantum devices. *Physical Review Applied*, 15(3), 034026.

Guerreschi, G. G., & Matsuura, A. Y. (2019). QAOA for Max-Cut requires hundreds of qubits for quantum speed-up. *Scientific reports*, 9(1), 1-7.

Xue, C., Chen, Z. Y., Wu, Y. C., & Guo, G. P. (2021). Effects of quantum noise on quantum approximate optimization algorithm. *Chinese Physics Letters*, 38(3), 030302.

Resch, S., & Karpuzcu, U. R. (2021). Benchmarking quantum computers and the impact of quantum noise. *ACM Computing Surveys (CSUR)*, 54(7), 1-35.

# Standard noise model

Breuer and Petruccione: *The Theory of Open Quantum Systems*, Oxford University Press (2002)

## Theory of **open quantum systems**

$$|\psi\rangle \rightarrow \rho = |\psi\rangle\langle\psi|$$

State vector

Density matrix

$$\frac{d}{dt}\rho_t = -\frac{i}{\hbar}[H_t, \rho_t] + \sum_k \gamma_k \left[ L_k \rho_t L_k^\dagger - \frac{1}{2} \{L_k^\dagger L_k, \rho_t\} \right]$$

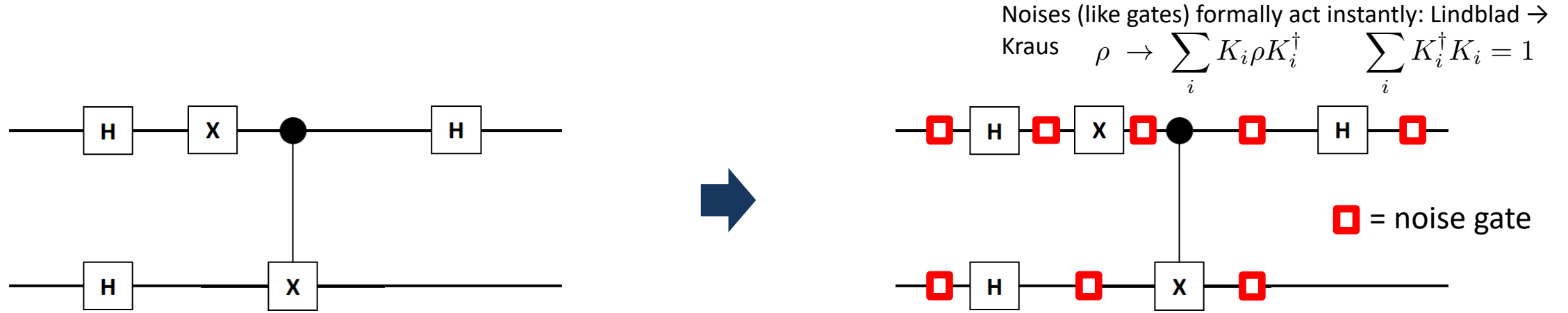
Internal evolution

Effect of the environment

Issues to deal with:

- More complicated dynamics; how to model the environment efficiently
- With the density matrix, the problem scales quadratically with the size of the system.

# How to describe noises

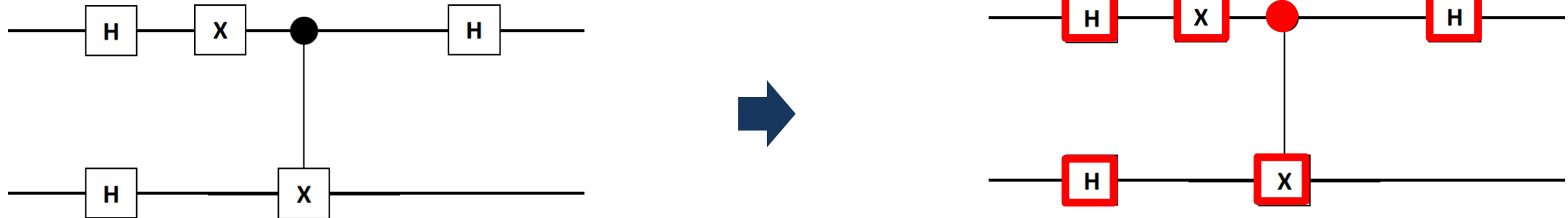


Standard noise simulation (e.g. in Qiskit)

- Gates and noise are formally **decoupled** (a sort of Trotterization), because time scales are small (IBM: gate time  $\sim 10^{-8}$  s, decoherence times  $\sim 10^{-4}$  s) ?
- Use the **quantum-jump-like approach** to replace the density matrix with (stochastic) state vector  $\rightarrow$  stochastic dynamics ✓

# Noisy Gates

Our approach: provide a more accurate description of the noisy behaviour of a quantum computer



- Noises are **embedded** in the gate → more realistic picture



- State vector (stochastic) description



# Noisy Gates

Bassi, A., & Deckert, D. A. (2008). Noise gates for decoherent quantum circuits. *Physical Review A*, 77(3), 032323.

$$\frac{d}{dt}\rho_t = \underbrace{-\frac{i}{\hbar}[H_t, \rho_t]}_{\text{Gate}} + \underbrace{\sum_k \gamma_k \left[ L_k \rho_t L_k^\dagger - \frac{1}{2} \{L_k^\dagger L_k, \rho_t\} \right]}_{\text{Noise}}$$

$$d|\psi_s\rangle = \left[ -\frac{i}{\hbar} H_s ds + \sum_{k=1}^{N^2-1} \left[ i\epsilon dW_{k,s} - \frac{\epsilon^2}{2} ds L_k^\dagger \right] L_k \right] |\psi_s\rangle$$

\* - reparametrized time in gate time units  
- diagonalized Lindblad in canonical form

Stochastic evolution for the state vector (stochastic unravelling)

1) Formal equivalence:  $\rho_t = \mathbb{E}[|\psi_t\rangle\langle\psi_t|]$

2) The dynamics is **linear**, therefore it can be represented as a gate (noisy gate)



$$|\psi_{s=1}(\xi)\rangle = \bar{N}(\xi) |\psi_0\rangle$$

\* Due to the noises  $\xi$ , the gate is not unitary and norm preserving. But at the statistical level the trace is preserved, and one recovers the standard (Lindblad) behaviour.

# Solution of the SDE

Gardiner, C. W. (1985). *Handbook of stochastic methods* (Vol. 3, pp. 2-20). Berlin: Springer.

Arnold, L. (1974). *Stochastic differential equations*. New York: John Wiley & Sons

$$\bar{N}(\xi) = U_g e^{\Lambda} e^{\Xi}(\xi)$$

$U_g$  = noiseless gate  $s \in [0, 1]$

$$\Lambda = -\frac{\epsilon^2}{2} \int_0^1 ds \sum_{k=1}^{N^2-1} [L_{k,s}^\dagger L_{k,s} - L_{k,s}^2]$$

Deterministic contribution of the noise, to order  $O(\epsilon^2)$

$$\Xi(\xi) := i\epsilon \sum_{k=1}^{N^2-1} \int_0^1 dW_{k,s} L_{k,s}$$

Stochastic contribution of the noise, to order  $O(\epsilon^2)$

Operators in the interaction picture

- Find  $H(U_g)$  and  $L_k$  for a given device
- Compute  $L_{k,s}$  and  $\Lambda$
- Compute stochastic properties of Itô processes in  $\Xi(\xi)$
- Sample to get a realization of  $\bar{N}(\xi)$



# IBMQ devices

Krantz, P., Kjaergaard, M., Yan, F., Orlando, T. P., Gustavsson, S., & Oliver, W. D. (2019). A quantum engineer's guide to superconducting qubits. *Applied Physics Reviews*, 6(2), 021318.  
Georgopoulos, K., Emary, C., & Zuliani, P. (2021). Modeling and simulating the noisy behavior of near-term quantum computers. *Physical Review A*, 104(6), 062432.

## Gates

Native gate set  $\{RZ(\phi), X, SX, \text{CNOT}\}$

Cross resonance (CR) gate

$\theta$  : rotation angle

$\phi$  : phase, realizes  
virtual Z gates

$$H(\theta, \phi) = \frac{\theta \hbar}{2} R_{xy}(\phi)$$
$$R_{xy}(\phi) = \cos(\phi)X + \sin(\phi)Y$$



$$H^{(1,2)}(\theta, \phi) = \frac{\hbar \theta}{2} Z^{(1)} \otimes R_{xy}^{(2)}$$
$$R_{xy}(\phi) = \cos(\phi)X + \sin(\phi)Y$$

Note: how to implement the pulse

## Noises

-Single qubit depolarization:  $\gamma_d$

-Single qubit amplitude and phase damping:  $\gamma_1, \gamma_z$

$$L_1 = \sqrt{\frac{\lambda_1}{\lambda}} \sigma^-, \quad L_2 = \sqrt{\frac{\lambda_2}{\lambda}} \sigma^+, \quad L_3 = \sqrt{\frac{\lambda_3}{\lambda}} Z$$
$$\lambda_1 = 2\gamma_d, \quad \lambda_2 = 2\gamma_d + \gamma_1, \quad \lambda_3 = \gamma_d + \gamma_z$$
$$\lambda = \lambda_1 + \lambda_2 + \lambda_3$$

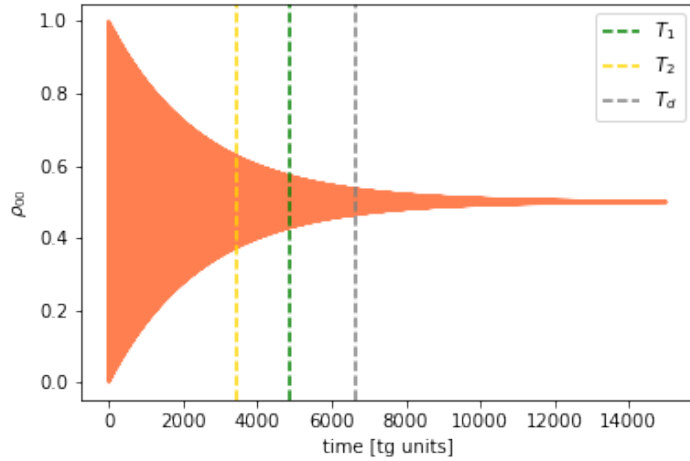
$$\lambda_k \sim 10^4 \text{ Hz}$$

$$t_g \sim 10^{-8} \text{ s}$$

$$\epsilon = \sqrt{\lambda t_g} \ll 1$$

# Simulation of the noisy X gate

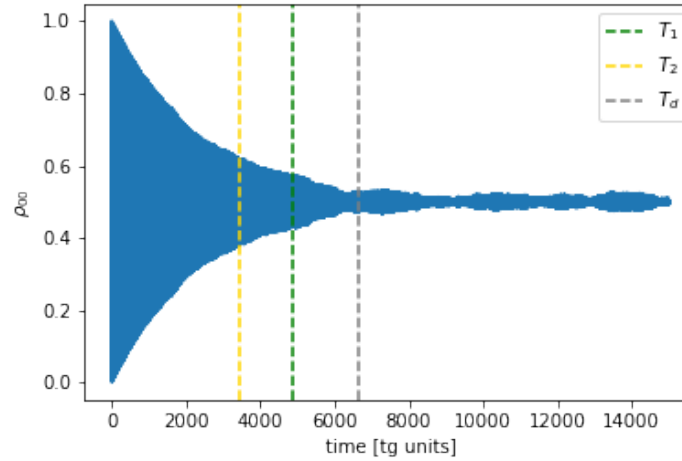
Lindblad equation



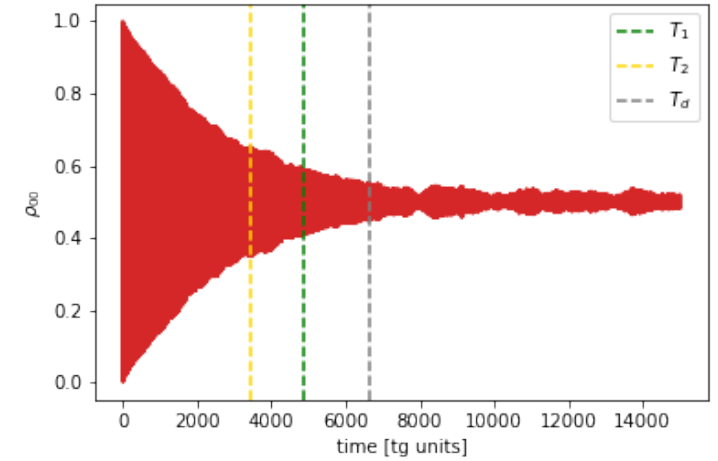
Initial state =  $|0\rangle$

Shown:  $\langle 0|\rho|0\rangle$

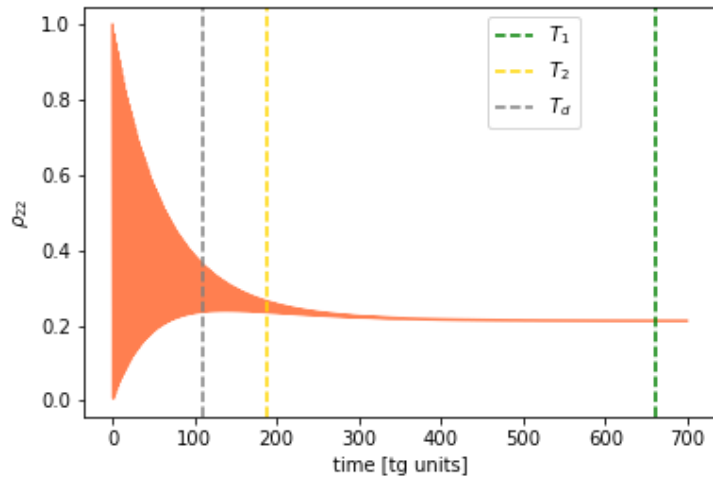
Noisy gates



Qiskit simulator



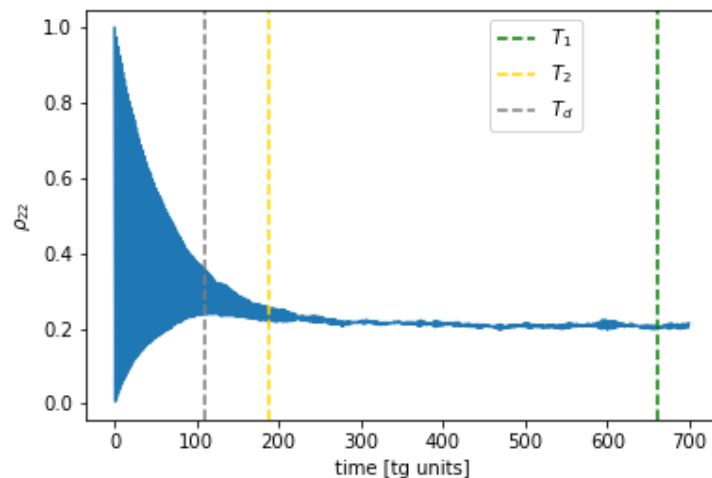
# Simulation of the noisy CR gate



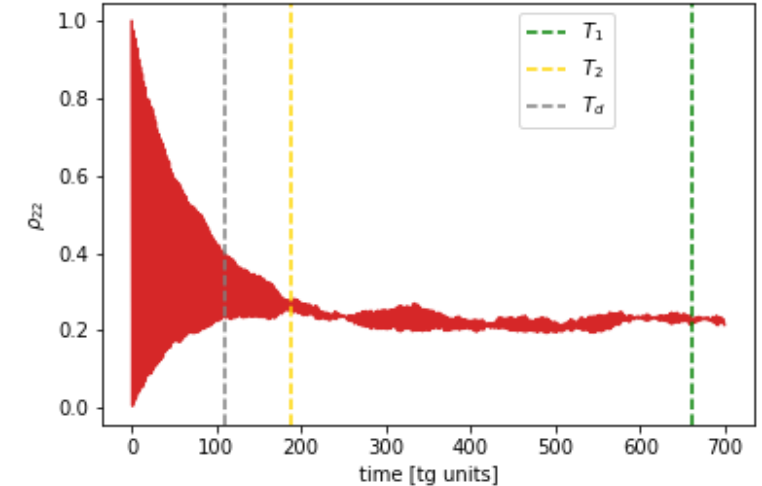
Initial state =  $|10\rangle$

Shown:  $\langle 10|\rho|10\rangle$

Numerical solution

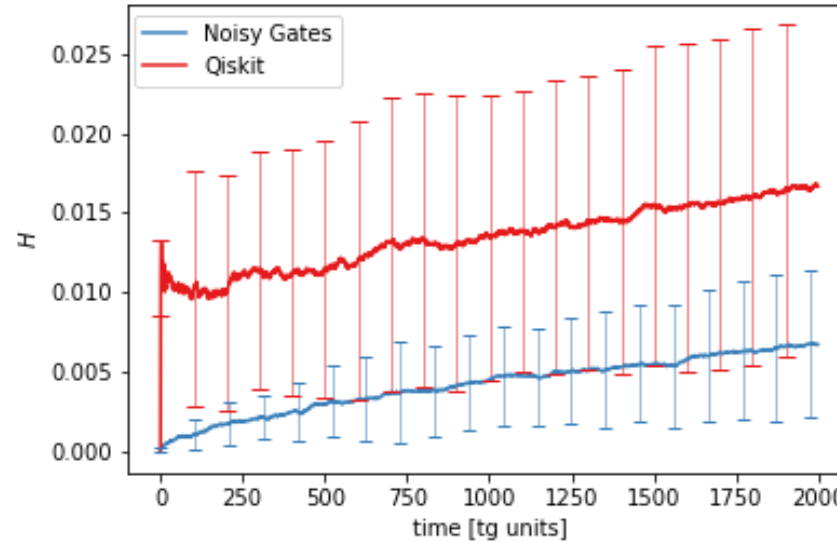
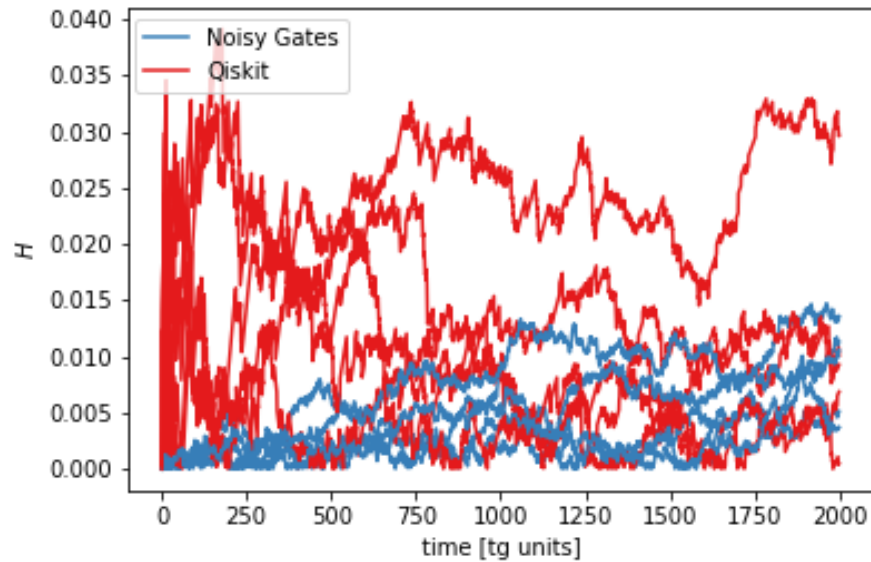


Average over 1000 realizations

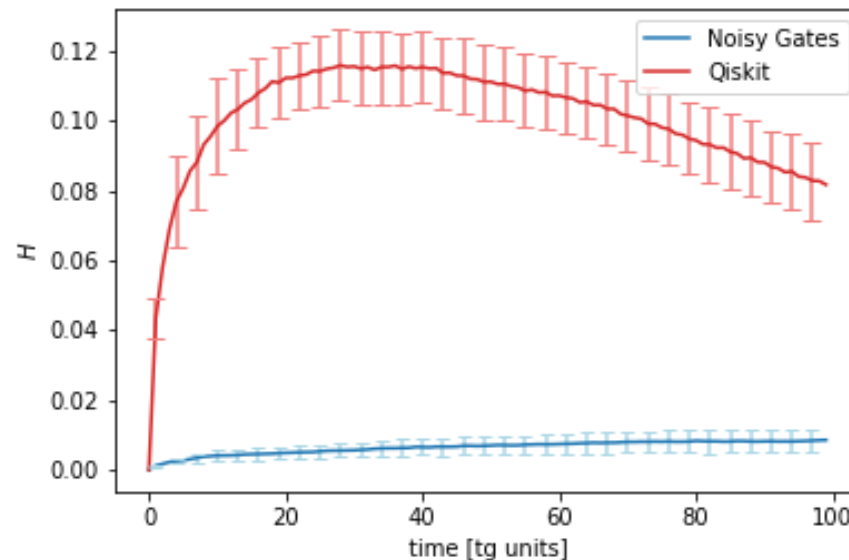
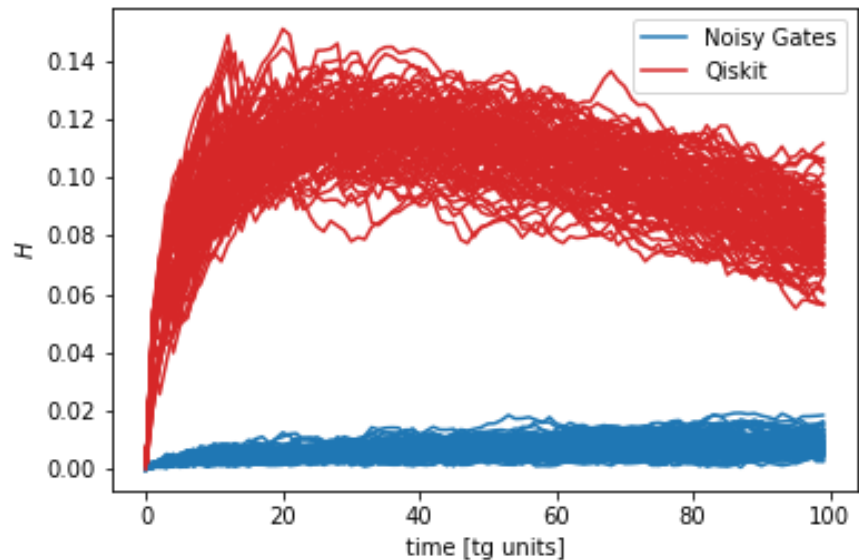


Average over 1000 realizations

# Simulation of the noisy X gate



# Simulation of the noisy CR gate



**Hellinger Distance**

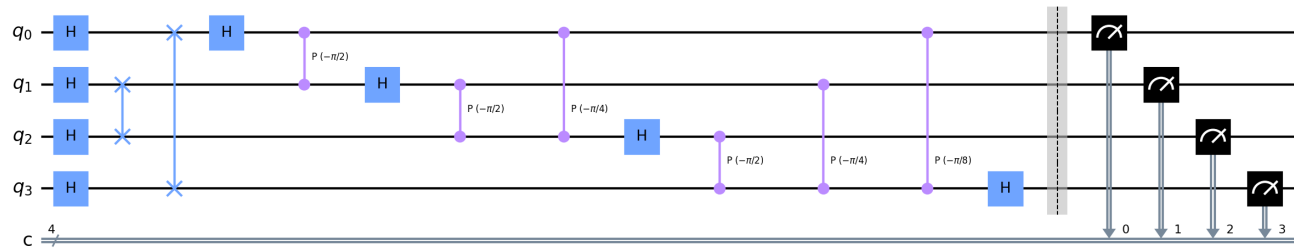
Lindblad and noisy  
gates

Lindblad and Qiskit  
simulator

$$H(P, Q) = \frac{1}{\sqrt{2}} \sqrt{\sum_{i=1}^n (\sqrt{p_i} - \sqrt{q_i})^2}$$

100 independent simulations each including 1000 runs

# The inverse QFT algorithm



$$|0\rangle^{\otimes N} \xrightarrow{H^{\otimes N}} |+\rangle^{\otimes N} \xrightarrow{QFT^{-1}} |0\rangle^{\otimes N}$$

**Note:** same noise model (depolarization + relaxation) **during** (after) gates + relaxation on idle qubits and bit-flip error before measurements

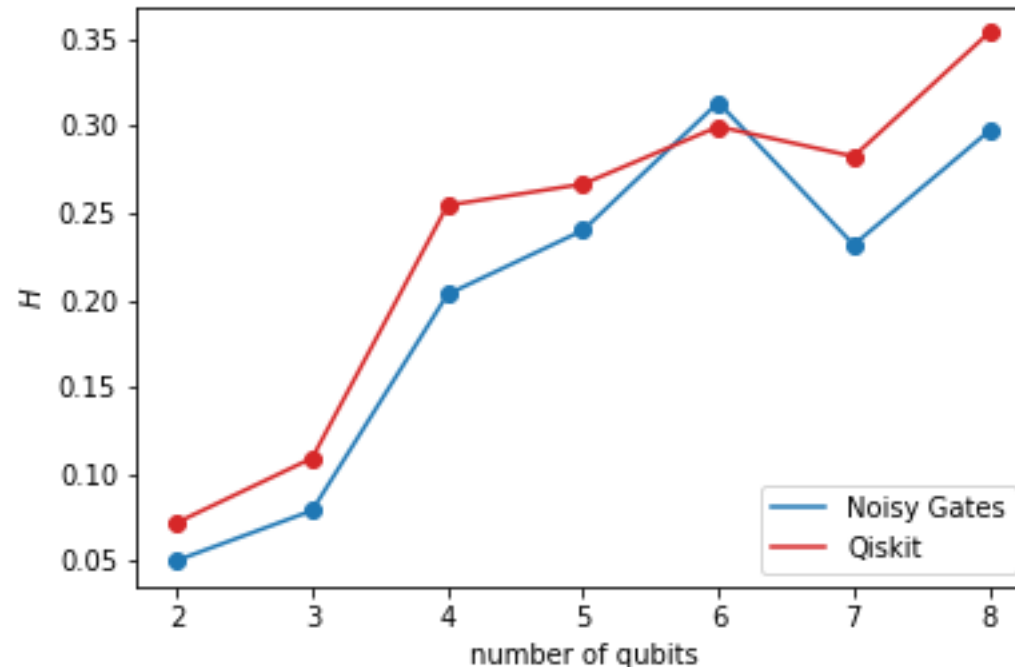
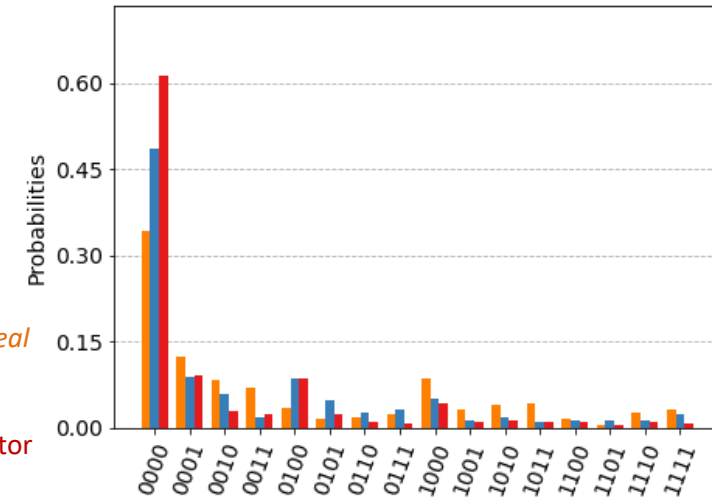
Transpile

Run

*ibmq\_montreal*

Noisy gates

Qiskit simulator



**Hellinger Distance**

*ibmq\_montreal* and  
noisy gates

*ibmq\_montreal* and  
Qiskit simulator

# Future work

Better analysis of noises, especially for two qubit gates

Application to algorithm of interest, to extract “long-time” behaviour

# The Team & Collaborators



Angelo Bassi



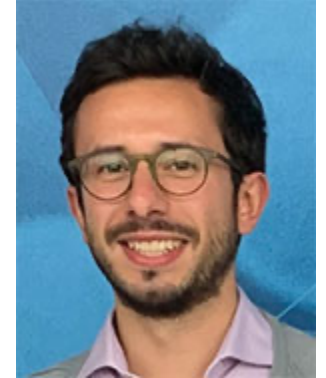
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Thank you for your attention