Sub-GeV dark matter search at beam dump experiments

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[attractive features]





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- Freeze-out mechanism can yield DM abundance \Rightarrow DM-SM reaction cross section can be large





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- Viable DM mass range is limited \Rightarrow Thermal mass window (~1 MeV to ~100 TeV)

• One of the DM candidates is thermal DM, i.e., DM is thermalized with SM particles in the early universe

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Sub-GeV dark matter $\sim 1 \text{ MeV} \sim 1 \text{ GeV}$

Sub-GeV

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For a fixed $\langle \sigma v \rangle_{ann}$, y becomes smaller (feebly coupled) when m_{χ} gets smaller





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Hight intensity experiments are needed to search for Sub-GeV DM

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• Beam dump experiments are high-intensity experiments and are sensitive to Sub-GeV DM

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Beam particle

e.g., electron, proton, and muon

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- three components of beam dump experiment:



Detector

Detect signature of DM

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[Goal of beam dump experiment]

To detect DM signatures produced by beam-target collision

[Center of mass energy]

$$\sqrt{s} = \sqrt{m_{\text{beam}}^2}$$

where m_{beam} is mass of beam particle, m_{target} is mass of target particle, and E_{beam} is beam energy

 $+ m_{\text{target}}^2 + 2E_{\text{beam}}m_{\text{target}}$



[Center of mass energy]

$$\sqrt{s} = \sqrt{m_{\text{beam}}^2}$$

where m_{beam} is mass of beam particle, m_{target} is mass of target particle, and E_{beam} is beam energy

kinematically allowed

$$+ m_{\text{target}}^2 + 2E_{\text{beam}}m_{\text{target}}$$

• Center of mass energy is smaller than collider energy scale, but Sub-GeV DM productions are



[Center of mass energy]

$$\sqrt{s} = \sqrt{m_{\text{beam}}^2}$$

kinematically allowed

Ex. Beam = electron (
$$m_{\text{beam}}$$
 = 0.5 MeV),

$$\sqrt{s} = \sqrt{m_{\text{beam}}^2 + }$$

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where m_{beam} is mass of beam particle, m_{target} is mass of target particle, and E_{beam} is beam energy

Center of mass energy is smaller than collider energy scale, but Sub-GeV DM productions are

target = nucleon (m_{target} = 1 GeV), E_{beam} = 10 GeV

 $+ m_{\text{target}}^2 + 2E_{\text{beam}}m_{\text{target}} \simeq 5 \text{ GeV}$



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Ex. Beam = electron (
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 = 0.5 MeV), target = nucleon (m_{target} = 1 GeV), E_{beam} = 10 GeV

$$\sqrt{s} = \sqrt{m_{\text{beam}}^2 + m_{\text{target}}^2 + 2E_{\text{beam}}m_{\text{target}}} \simeq 5 \text{ GeV}$$

Boosted Sub-GeV DM can be produced in beam dump

$$+ m_{\text{target}}^2 + 2E_{\text{beam}}m_{\text{target}}$$

where m_{beam} is mass of beam particle, m_{target} is mass of target particle, and E_{beam} is beam energy

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Beam dump experiment is high luminosity frontier



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(# of produced DM) = (DM production cross section $[L^2]$)

X



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(# of produced DM) = (DM production cross section $[L^2]$)

 \times (Beam flux [T⁻¹]) \times



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(# of produced DM) = (DM production cross section $[L^2]$)

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Beam dump experiment is high luminosity frontier

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Beam dump experiment is high luminosity frontier



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Ex. Target = Iron, # of injected proton beam = 10^{20}

Beam dump experiment is high luminosity frontier



Beam dump experiment is high luminosity frontier



Beam dump Beam particle

• Center of mass energy and luminosity are determined by beam and target \Rightarrow How about detector?

Detector





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e.g., E137, LDND, MiniBooNE, CCM120,

BDX, PIP2-BD, DUNE, CCM200, SHiP, ILC-BDX(proposed), etc







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Beam dump experiments are divided into three classes of experiments







Beam dump experiments are divided into three classes of experiments







Typical setup:

Beam

Beam dump

Detector



Beam dump

Typical setup:

Beam

Benchmark model: $\mathscr{L} \supset \epsilon \cdot eA'_{\mu}J^{\mu}_{\text{EM}} - g_DA'_{\mu}(i\bar{\chi}_2\gamma^{\mu}\chi_1 + \text{H.c.})A'$: Dark photon, $\chi_{1,2}$: Pseudo-Dirac DM, J^{μ}_{EM} : SM EM current

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of detected DM (signal events):



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~ (# of produced DM) ×



of detected DM (signal events):

 \sim (# of produced DM) \times (Probability DM reaches detector) \times



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cto	r



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of detected DM (signal events):

Acceptance

 \sim (# of produced DM) \times (Probability DM reaches detector) \times (Probability DM is detected)

ct	0	r	
			/





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 \propto (Beam flux)×(ϵe)²

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of detected DM (signal events):

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 \propto (height of detector)² \times (length b/w beam dump and detector)⁻²

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* Detector cannot be too near because large beam dump or shield is needed to reduce beam-induced BG

Acceptance

\sim (# of produced DM) \times (Probability DM reaches detector) \times (Probability DM is detected)

 \propto (Length of detector)×(ee)²

High flux beam, near* and large detectors are suited for recoil and visible decay search





Background events in recoil and visible decay search

Background events in recoil and visible decay search

Beam dump

Beam

Detector


Detector



Detector

• Large beam dump or shield is needed to remove beam-induced particles, e.g., π^{\pm} , μ^{\pm} , K



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BG of recoil events

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- neutrinos are produced by decay of π^{\pm} , μ^{\pm} ,... in beam dump Ex. $\pi^+ \to \mu^+ + \nu_{\mu}$, $\mu^+ \to \bar{\nu}_{\mu} + e^+ + \nu_{e}$,...

Beam-induced neutrino can be main background events



- π^{\pm} , μ^{\pm} are absorbed or decay at rest by large beam dump \Rightarrow neutrino flux is reduced

Beam-induced neutrino can be main background events

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Missing energy search, e.g., NA64:

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EM and hadron calorimeter, i.e., active target



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Acceptance



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* not proportional to $(\epsilon e)^2$ in contrast to recoil and visible search

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• # of missing events:

~ (# of produced DM) × (Probability DM reaches detector) × (Probability DM is detected) \propto (Beam flux)×(ϵe)² * not proportional to $(\epsilon e)^2$ in contrast to recoil and visible search

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 \propto (Beam flux)×(ϵe)²

Acceptance is good, but the continuous beam(small flux) is needed to reconstruct SM

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Acceptance



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Tracker EM and hadron calorimeter, i.e., active target



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- Potential BG events:
 - Mistakenly tagged initial beam, e.g., incident lower energy beam

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EM and Hadron calorimeter



Missing energy search, e.g., NA64:

EM and hadron calorimeter, i.e., active target Tracker



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- Potential BG events:

 - Photo-nuclear reaction and muon production

Missing momentum search, e.g., LDMX:







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* Here, single electron events are assumed to be selected as signal events to remove BG events

Missing momentum search, e.g., LDMX:








Background events in missing energy/momentum search

Missing energy search, e.g., NA64:



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Missing momentum search, e.g., LDMX:



- Mistakenly tagged initial beam, e.g., incident lower energy beam \Rightarrow Reducible by tracker in front of target - Photo-nuclear reaction and muon production \Rightarrow Reducible by calorimeter and tracker behind target







Background events in missing energy/momentum search

Missing energy search, e.g., NA64:



- Potential BG events:

 - Neutrino background, e.g., neutrino trident process($eN \rightarrow e \nu \bar{\nu} N$)

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- Mistakenly tagged initial beam, e.g., incident lower energy beam \Rightarrow Reducible by tracker in front of target - Photo-nuclear reaction and muon production \Rightarrow Reducible by calorimeter and tracker behind target - Neutrino background, e.g., neutrino trident process($eN \rightarrow e\nu\bar{\nu}N$) \Rightarrow Irreducible but it's negligible Less than one predicted BG event in both NA64 and LDMX





Excluded regions by beam dump experiments



• Benchmark model: $\mathscr{L} \supset \epsilon \cdot eA'_{\mu}J^{\mu}_{EM} - g_DA'_{\mu}(i\bar{\chi}_2\gamma^{\mu}\chi_1 + H.c.)$ A': Dark photon, $\chi_{1,2}$: Pseudo-Dirac DM

Pseudo-Dirac DM saturates observed DM abundance

DM annihilation cross section







Excluded regions by beam dump experiments



* Boosted DM productions in the beam dump are unaffected much by the spin of the DM or the Lorentz structure of its interactions



Proton beam

- \Rightarrow parasitic running of neutrino experiment
 - * MiniBooNE is off-target running to reduce neutrino BG

Sensitivity of future beam dump experiments



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- NA64e is an improved NA64



Summary

- The beam dump(fixed target) experiment is high luminosity experiment sensitive to Sub-GeV scale
- Several factors determine the sensitivity of the beam dump experiment (fixed target), e.g., beam flux, beam energy, beam particle, acceptance, detection approach,...
- There are many experiments, but these are not simply competitions to achieve the same physics
 - Ex. Proton beam dump experiments are suited to leptophobic NP search
 - Muon beam dump experiments are suited to muon-philic NP search
 - LDMX is highly sensitive to the DM production process, but it is not sensitive to the DM detection process in contrast to recoil and visible decay searches. * LDMX is missing momentum search
 - \Rightarrow If a positive DM signal is observed, model distinction would be possible
- The beam dump(fixed target) experiment may play a significant role in searching for Sub-GeV DM models involving recoil, visible decay, and missing signatures







Sensitivity of beam dump experiments



• Benchmark model: $\mathscr{L} \supset \epsilon \cdot eA'_{\mu}J^{\mu}_{EM} - g_DA'_{\mu}(i\bar{\chi}_2\gamma^{\mu}\chi_1 + H.c.)$ A': Dark photon, $\chi_{1,2}$: Pseudo-Dirac DM



Continuous beam

The continuous beam doesn't have beam bunches, which makes the trigger easier





Continuous beam with high flux is needed for the missing signature search * The missing signature search experiments are not conducted much because of this

severe beam condition.

e.g. NA64



Beam unrelated background

- Beam-unrelated background: cosmic muons
 - deep underground location of detector
 - time window based on pulsed beam



- directionality of signal events
- muon veto

Beam-unrelated background events are reducible

Pseudo-Dirac DM mass eigenstates

$$-\mathscr{L} \supset m_D \eta \xi + \frac{1}{2} m_M (\eta^2 + \xi^2) + \text{h.c.}$$

$$= m_D(\xi^c)^{\dagger} \eta + \frac{m_M}{2} \left((\eta^c)^{\dagger} \eta + (\xi^c)^{\dagger} (\xi) \right) + 1$$

$$\eta = \frac{1}{\sqrt{2}}(\chi_2 - i\chi_1), \quad \xi = i\frac{1}{\sqrt{2}}(\chi_1 - i\chi_2)$$

$$=\frac{1}{2}(m_D - m_M)(\chi_1^c)^{\dagger}\chi_1 + \frac{1}{2}(m_D + m_M)(\chi_2^c)^{\dagger}\chi_1$$

 $\psi = (\eta, \xi^c)^T, \ \xi^c = i\sigma^2 \xi^*$

h.c.

 $(\chi_2^c)^{\dagger}\chi_2 + h.c.$

 $(\chi_1^c)^{\dagger}\chi_2 = (\chi_2^c)^{\dagger}\chi_1$

Pseudo-Dirac DM kinetic term

$$\mathscr{L} = \bar{\psi} i \gamma^{\mu} D_{\mu} \psi = i \left(\psi_L^{\dagger} \bar{\sigma}^{\mu} D_{\mu} \psi_L + \psi_R^{\dagger} \sigma^{\mu} D_{\mu} \psi_R \right)$$

$$\supset g_D \cdot \left(\eta^{\dagger} \bar{\sigma}^{\mu} A'_{\mu} \eta + (\xi^c)^{\dagger} \sigma^{\mu} A'_{\mu} \xi^c \right)$$
$$\eta^{\dagger} \bar{\sigma}^{\mu} \eta = \frac{1}{2} \left(\chi_2^{\dagger} \bar{\sigma}^{\mu} \chi_2 + \chi_1^{\dagger} \bar{\sigma}^{\mu} \chi_1 + i \chi_1^{\dagger} \bar{\sigma}^{\mu} \chi_2 - i \chi_2^{\dagger} \bar{\sigma}^{\mu} \chi_1 \right)$$
$$(\xi^c)^{\dagger} \sigma^{\mu} \xi^c = \frac{1}{2} \left(-\chi_1^{\dagger} \bar{\sigma}^{\mu} \chi_1 + i \chi_1^{\dagger} \bar{\sigma}^{\mu} \chi_2 - i \chi_2^{\dagger} \bar{\sigma}^{\mu} \chi_1 - \chi_2^{\dagger} \bar{\sigma}^{\mu} \chi_2 \right)$$

$$= ig_D \cdot A'_{\mu} \left(\chi_1^{\dagger} \bar{\sigma}^{\mu} \chi_2 - \chi_2^{\dagger} \bar{\sigma}^{\mu} \chi_1 \right)$$

$$= ig_D \cdot A'_{\mu} \left(\chi_1^{\dagger} \bar{\sigma}^{\mu} \chi_2 + (\chi_1^c)^{\dagger} \sigma^{\mu} \chi_2^c \right)$$

 $= ig_D \cdot A'_{\mu} \bar{\chi}_1 \gamma^{\mu} \chi_2 \quad \checkmark$

 $\psi = (\psi_L, \psi_R) = (\eta, \xi^c)^T, \ \xi^c = i\sigma^2 \xi^*$ $\bar{\sigma}^{\mu} = (1, -\vec{\sigma}), \quad \sigma^{\mu} = (1, \vec{\sigma})$ $D_{\mu} = \partial_{\mu} - ig_D A'_{\mu}$ $\sigma_2 \sigma^\mu \sigma_2 = (\bar{\sigma}^\mu)^T$



